

中国科学技术大学
University of Science and Technology of China

Ying-Ying Li (李英英)

arXiv: 2302.10965, 2401.xxxxx, 2402.xxxxx
in collaboration with
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A. de Gouvêa, J. Lazar,
P. A. N. Machado, G. Li,
Y.-L. Zhou

Multi-messenger Opportunities for Heavy Neutral Leptons

Jan 09, 2024 @ ICISE, Quy Nhon

Heavy Neutral Lepton - Sterile Neutrino

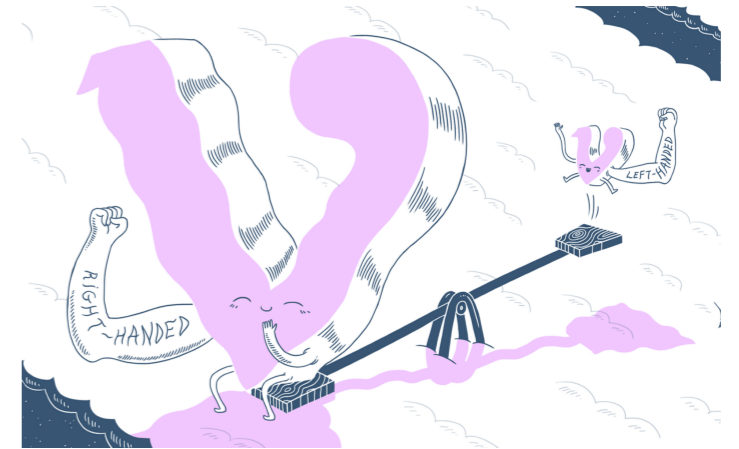
[Minkowski, Mohapatra, Senjanović, Gell-Mann, Ramond, Slansky, Yanagida]

Neutrino mass:
type-I seesaw
mechanism

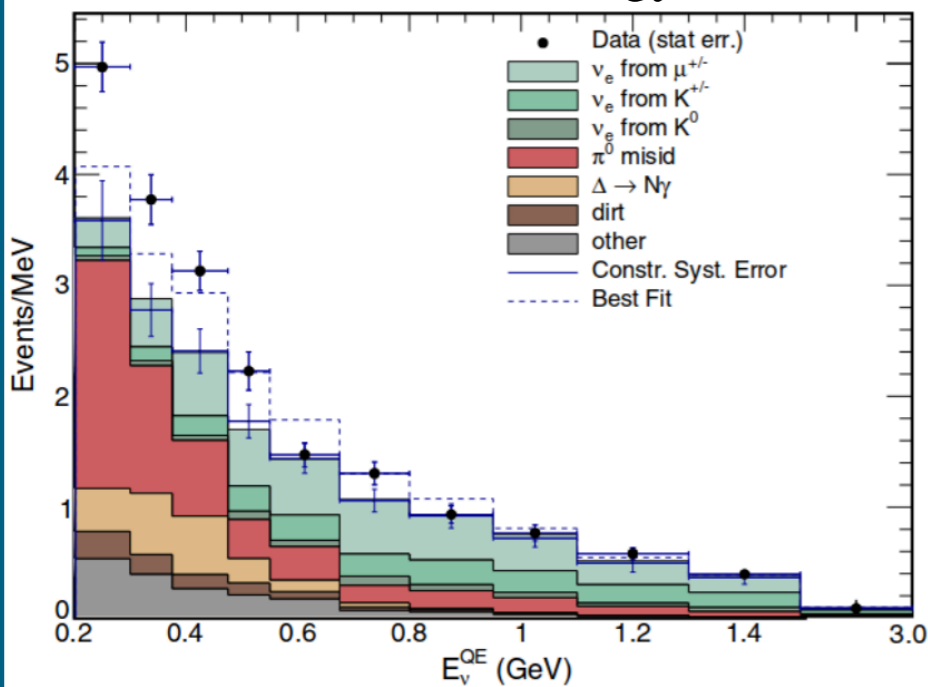
$$\mathcal{L} \supset \frac{1}{2} \bar{N}^c M_R N + \bar{L} Y_\nu \tilde{H} N + \text{h.c.}$$

$$\mathcal{M} = \begin{pmatrix} 0 & M_D \\ M_D & M_R \end{pmatrix}$$

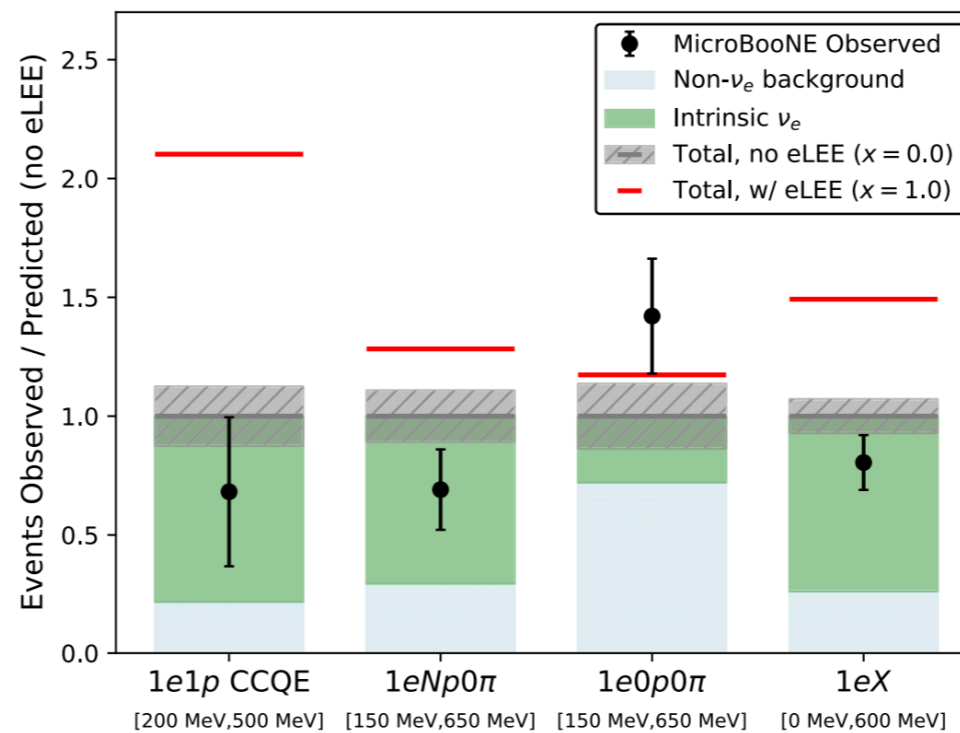
$$m_\nu = -M_D M_R^{-1} M_D^T$$



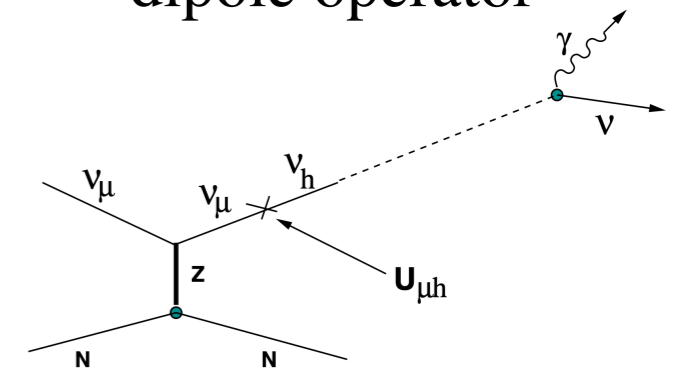
MinibooNE low energy excess



[PRL 128, 241801 (2022)]



Via mixing and
dipole operator



$$\theta^2 \sim 10^{-3}$$

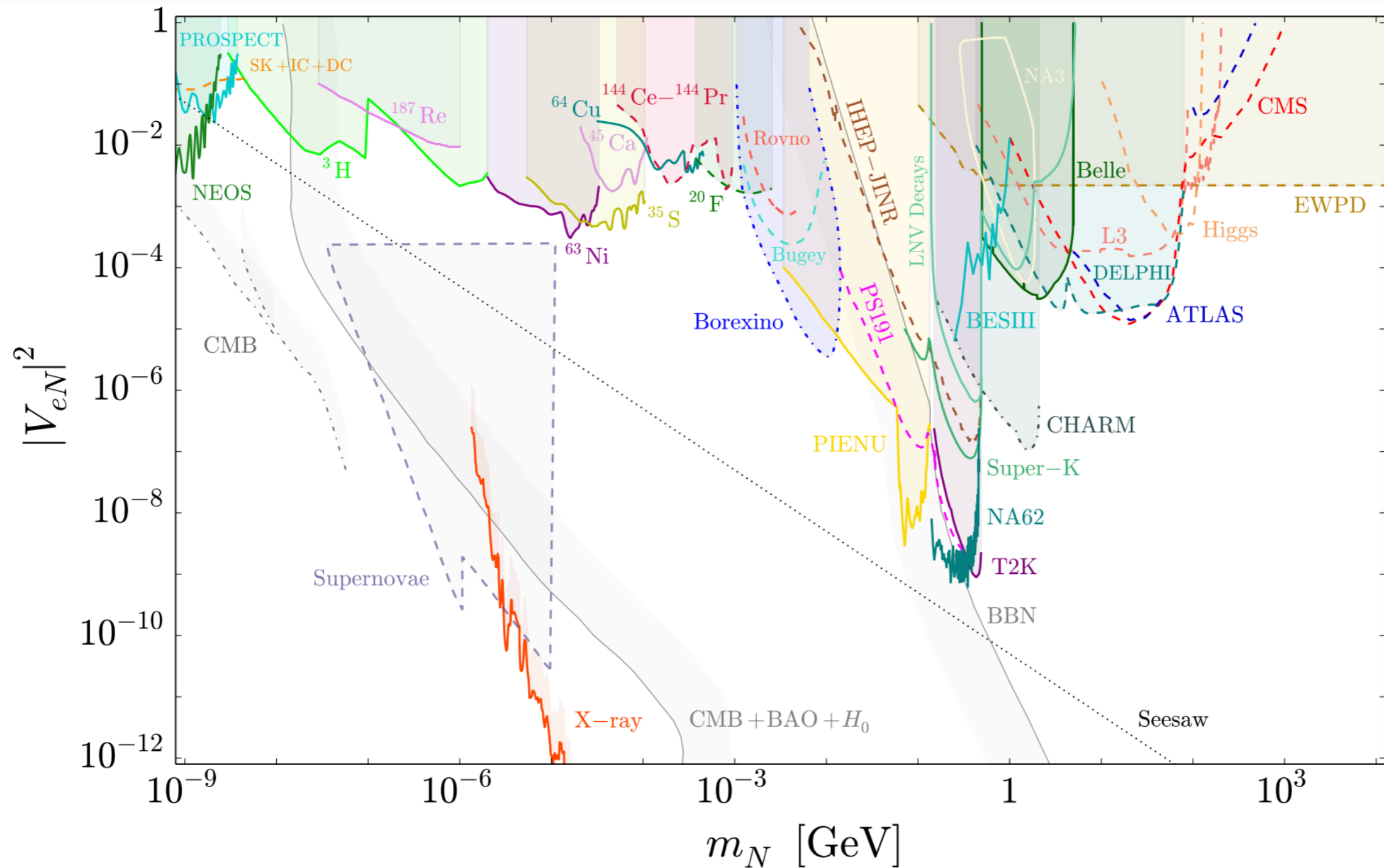
$$\mu \sim 10^{-9} \mu_B$$

[S. N. Gninenko, PRL 103, 241802 (2009)]

Heavy Neutral Lepton - Mixing Portal

$$\Gamma_{\nu_s \rightarrow \nu_a \gamma} = 1.38 \cdot 10^{-29} \text{ sec}^{-1} \left(\frac{\sin^2 2\theta}{10^{-7}} \right) \left(\frac{m_s}{1 \text{ keV}} \right)^5.$$

$$\theta^2 \sim \frac{m_\nu}{M_N}, \text{ keV} < M_N < 100 \text{ MeV}$$



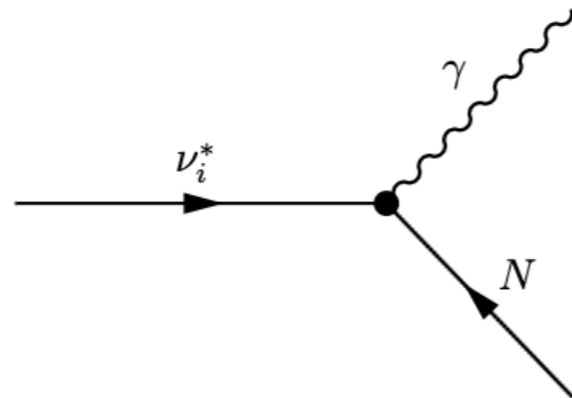
[P. D. Bolton, F. F. Deppisch, P. S. B. Dev, arXiv: 1912.03058]

Heavy Neutral Lepton - Dipole Portal

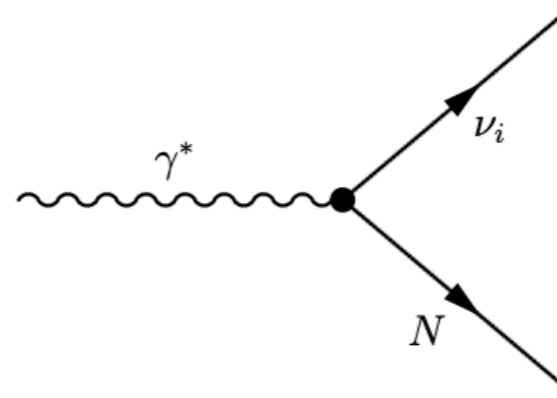
When the mixing effect is subdominant, ...

$$\mathcal{L} \supset \frac{1}{2} \mu_\nu \bar{\nu}_L^\alpha \sigma^{\mu\nu} N F_{\mu\nu}$$

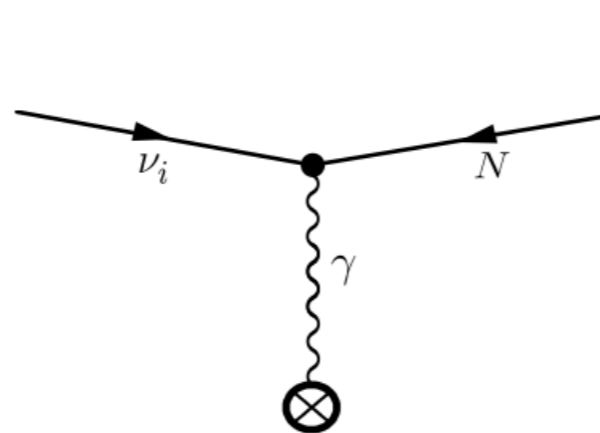
$$\pi^\pm, K^\pm \rightarrow \mu^\pm \left(\bar{\nu}_\mu^{(-)*} \rightarrow \gamma N^{(-)} \right) \quad \pi^0, \eta \rightarrow \gamma (\gamma^* \rightarrow \nu_a N)$$



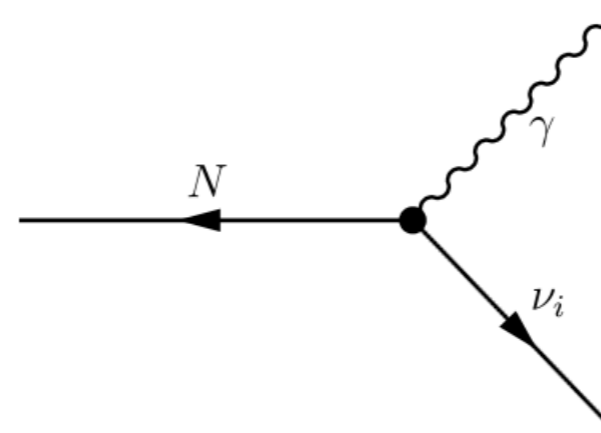
(a) Weak meson decays



(b) Dalitz-like decay



(c) Primakoff upscattering



(d) $N \rightarrow \gamma \nu$ (signal)

$$\text{Br}(M \rightarrow N) \propto \mu_\nu^2 m_M^2$$

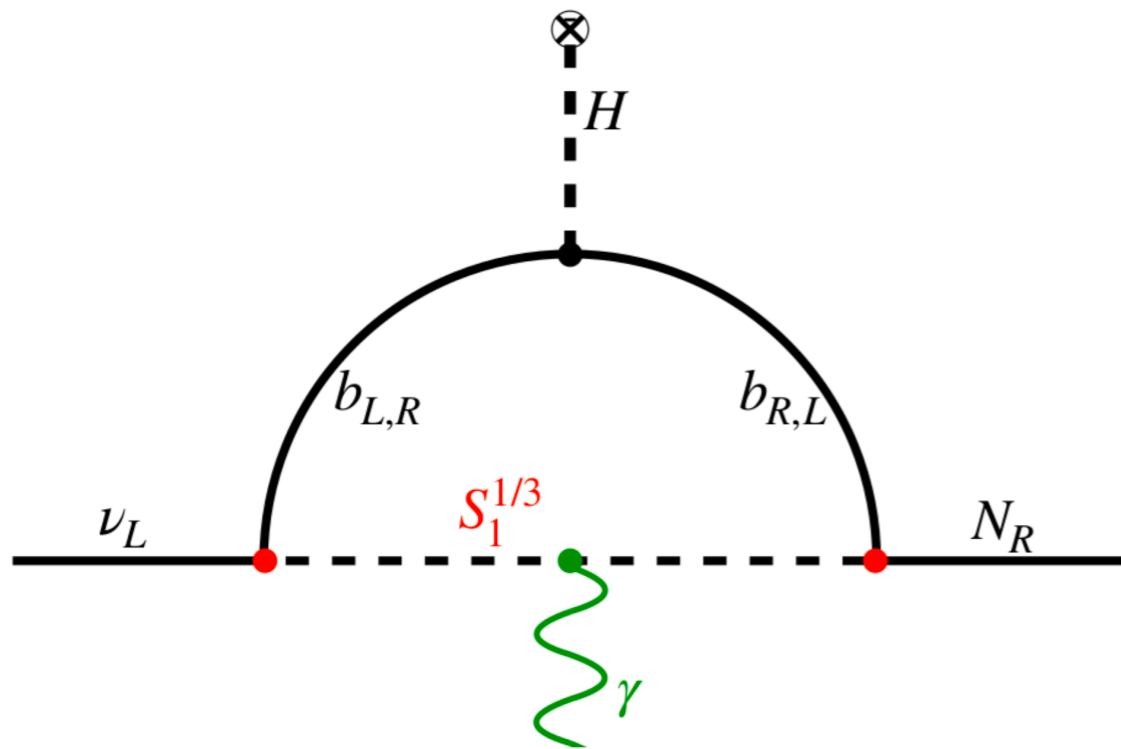
$$\Gamma_N = \frac{6}{4\pi} \mu_\nu^2 M_N^3 \quad (\text{flavor universal})$$

[Magill, Plestid, Pospelov, Tsai, arXiv:1803.03262]

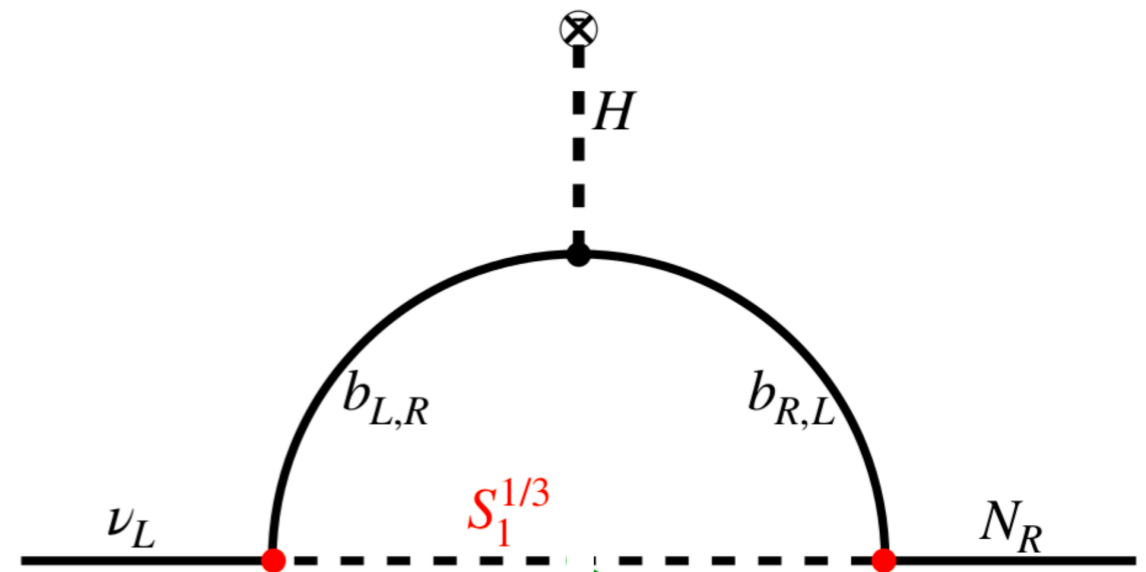
UV completion – Mixing vs Dipole

- Consider heavy scalar lepto-quark $S_1 \sim (\bar{3}, 1, 1/3)$

$$\mathcal{L} \supset y_1 \bar{b}_R^c N S_1 + y_2 \bar{Q}_L^3 L_L^i c S_1^\dagger + h.c.$$



$$\mu_\nu \approx \frac{e y_1 y_2}{8\pi^2 m_{LQ}^2} m_b \log \frac{m_b^2}{m_{LQ}^2}$$



$$m_{\nu N} \sim \frac{\mu_\nu}{\mu_B} \frac{m_{LQ}^2}{2m_e}$$

UV completion – Mixing vs Dipole

- Consider heavy scalar lepto-quark $S_1 \sim (\bar{3}, 1, 1/3)$

$$\mathcal{L} \supset y_1 \bar{b}_R^c N S_1 + y_2 \bar{Q}_L^3 L_L^i c S_1^\dagger + h.c.$$

$$m_{\nu N} \sim \frac{\mu_\nu m_{LQ}^2}{\mu_B 2m_e}$$

$$m_\nu \sim \frac{m_{\nu N}^2}{M_N}$$

$$\mu_\nu \sim 10^{-12} \mu_B \frac{\sqrt{m_\nu M_N}}{\text{MeV}}$$



for lepto-quark masses at the TeV scale

$$m_\nu \sim 0.1\text{eV}, \mu_\nu \sim 10^{-15} \mu_B$$

UV completion – Mixing vs Dipole

[Lindner et al., arXiv:1706.02555]

- Voloshin-type symmetry $SU(2)_\nu$

$$(N_R^C, \nu_L)^T \in \mathbf{2} \quad SU(2)_L \rightarrow SU(3)_L$$

$$\bar{\nu}_L N_R \rightarrow -\bar{\nu}_L N_R \quad \bar{\nu}_L \sigma_{\mu\nu} N_R F^{\mu\nu} \rightarrow \bar{\nu}_L \sigma_{\mu\nu} N_R F^{\mu\nu}$$

$$m_{\nu N} \sim \frac{\mu_\nu}{\mu_B} \frac{\alpha}{4\pi} \frac{m_V^2}{2m_e}$$

$$\mu_\nu \sim 10^{-8} \mu_B \frac{\sqrt{m_\nu M_N}}{\text{MeV}}$$

$SU(2)_\nu$ Symmetry-breaking scale m_V at the TeV scale

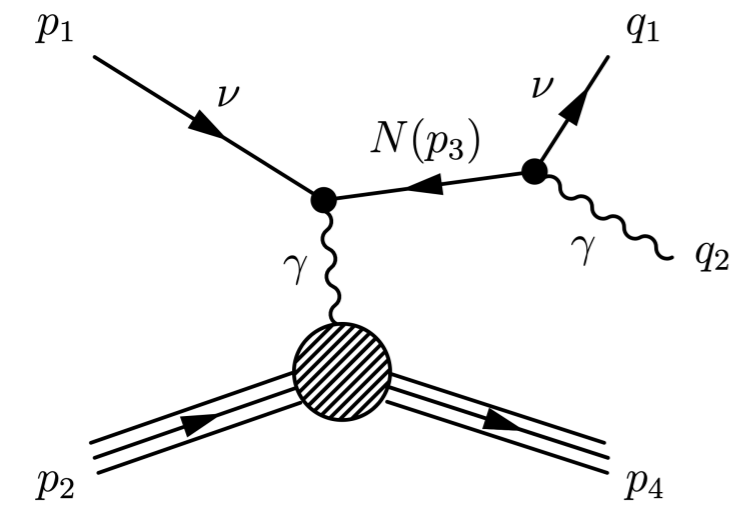
$$m_\nu \sim 0.1\text{eV}, \mu_\nu \sim 10^{-11} \mu_B$$

Current Probes: terrestrial experiments

Beam dump experiments:

MiniBoone, NONAD, CHARM

production via:

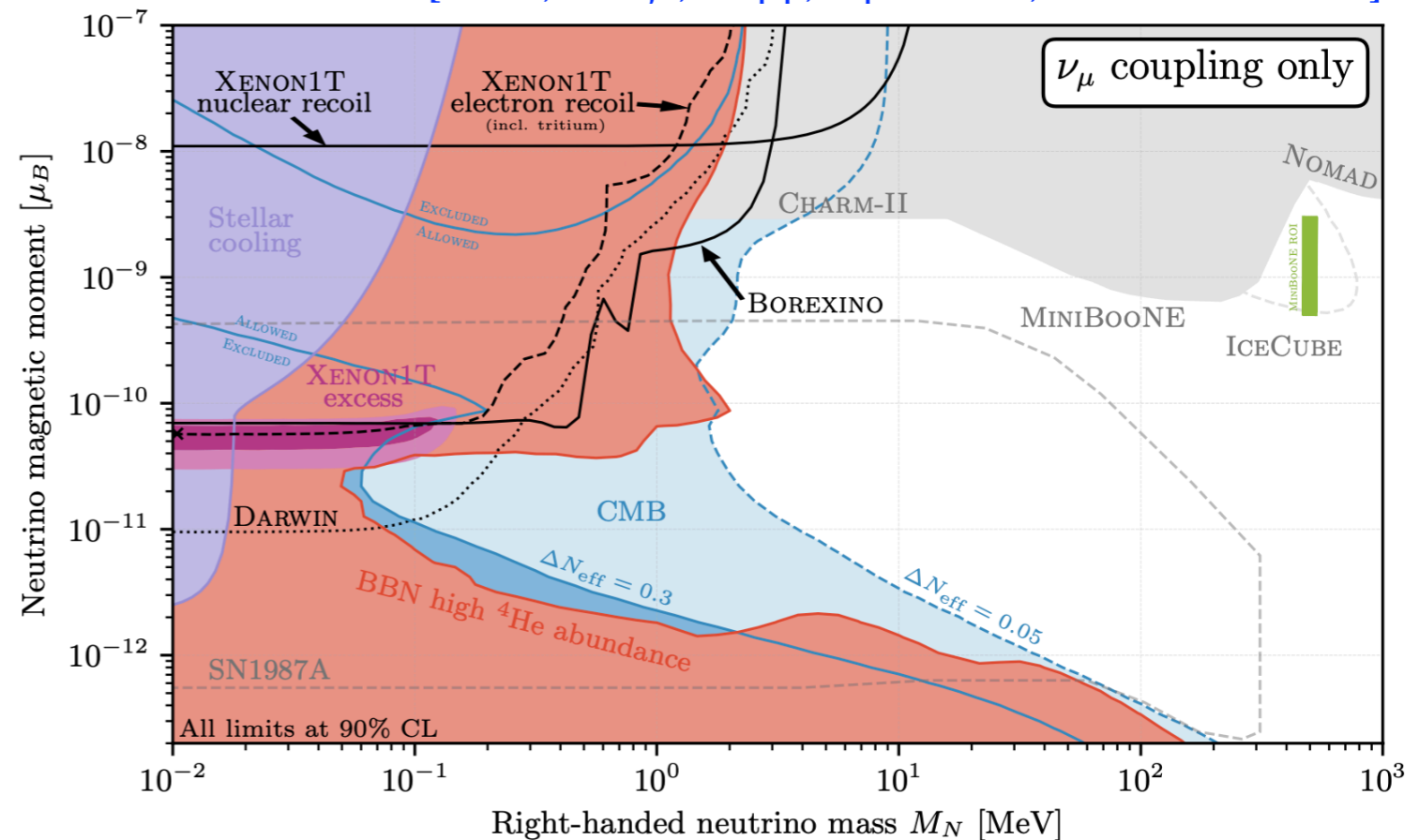


$$\pi^0, \eta \rightarrow \gamma(\gamma^* \rightarrow \nu_a N)$$

$$\pi^\pm, K^\pm \rightarrow \mu^\pm \left(\nu_\mu^{(-)*} \rightarrow \gamma N^{(-)} \right)$$

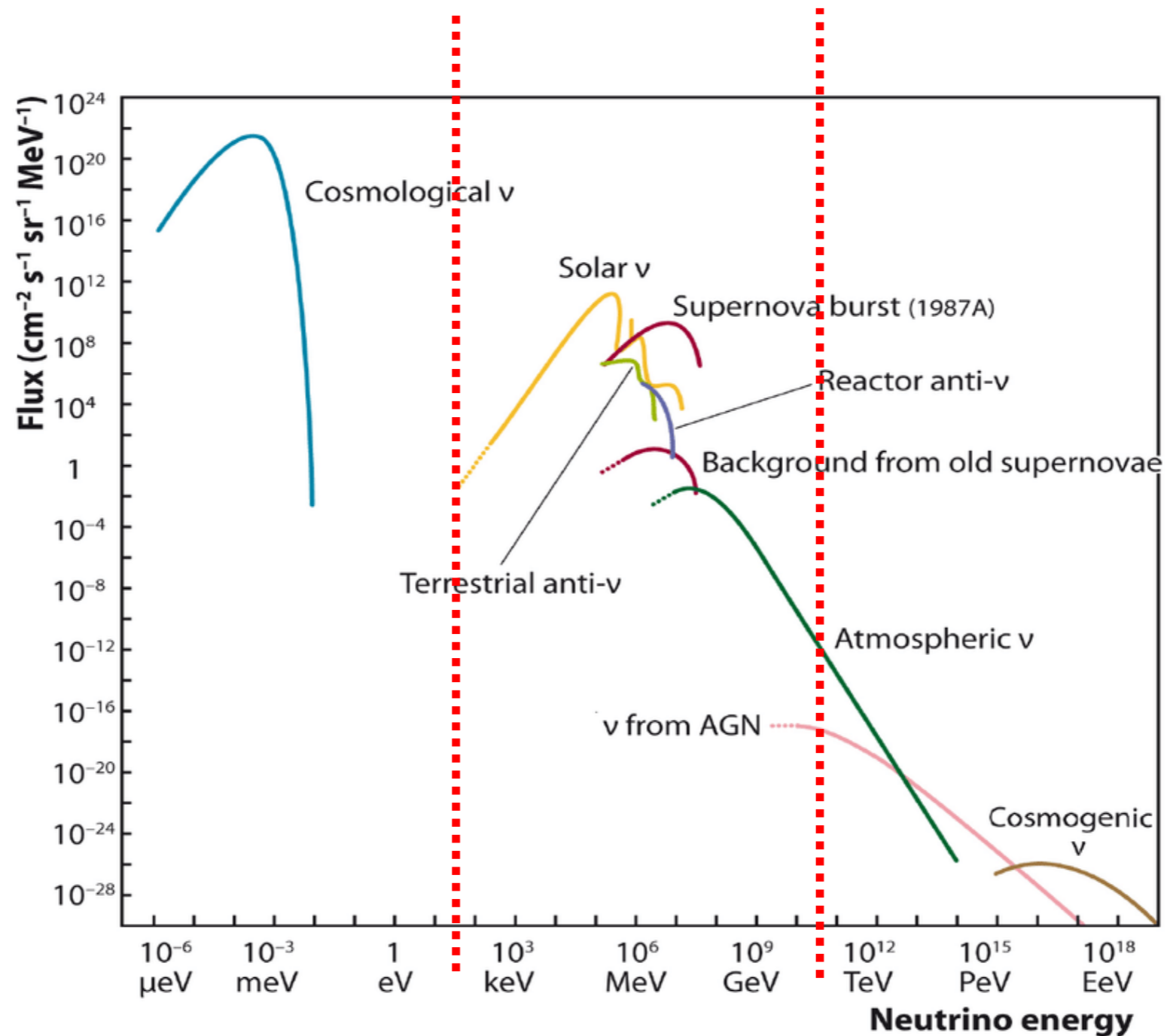
relevant for transition magnetic moments between ν_μ and N

[Magill, Plestid, Pospelov, Tsai, arXiv:1803.03262]
[Brdar, Greljo, Kopp, Opferkuch, arXiv:2007.15563]



Current Probes: terrestrial experiments

- looking to the sky



Current Probes: terrestrial experiments

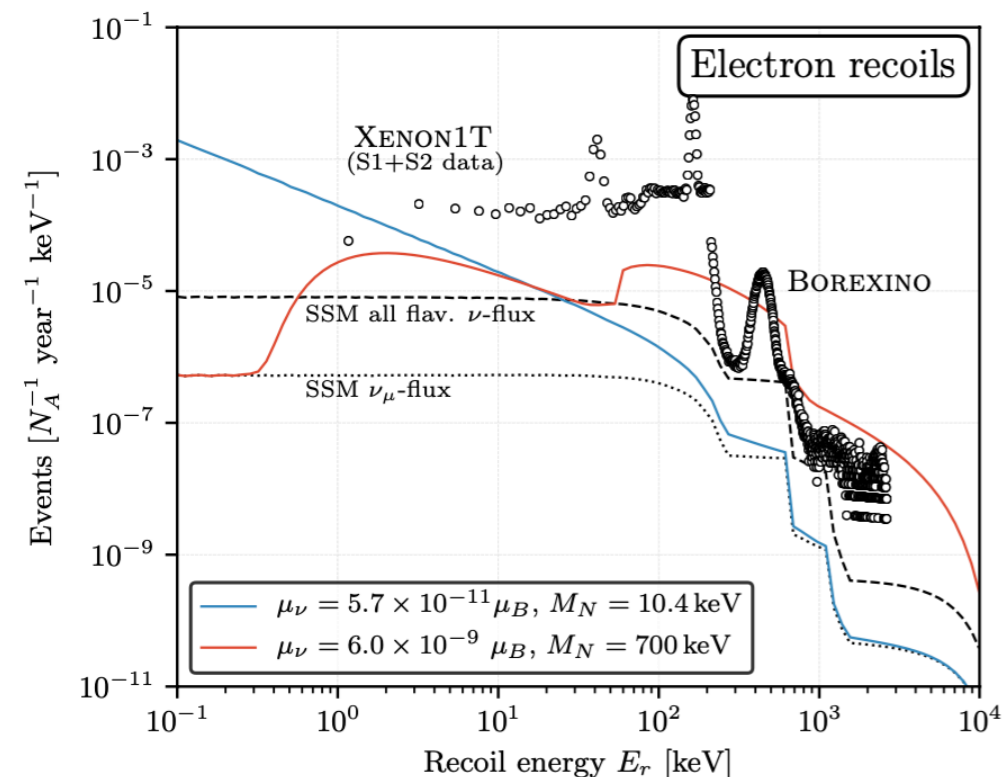
Solar neutrino spectrum:

Xenon1T, Borexino

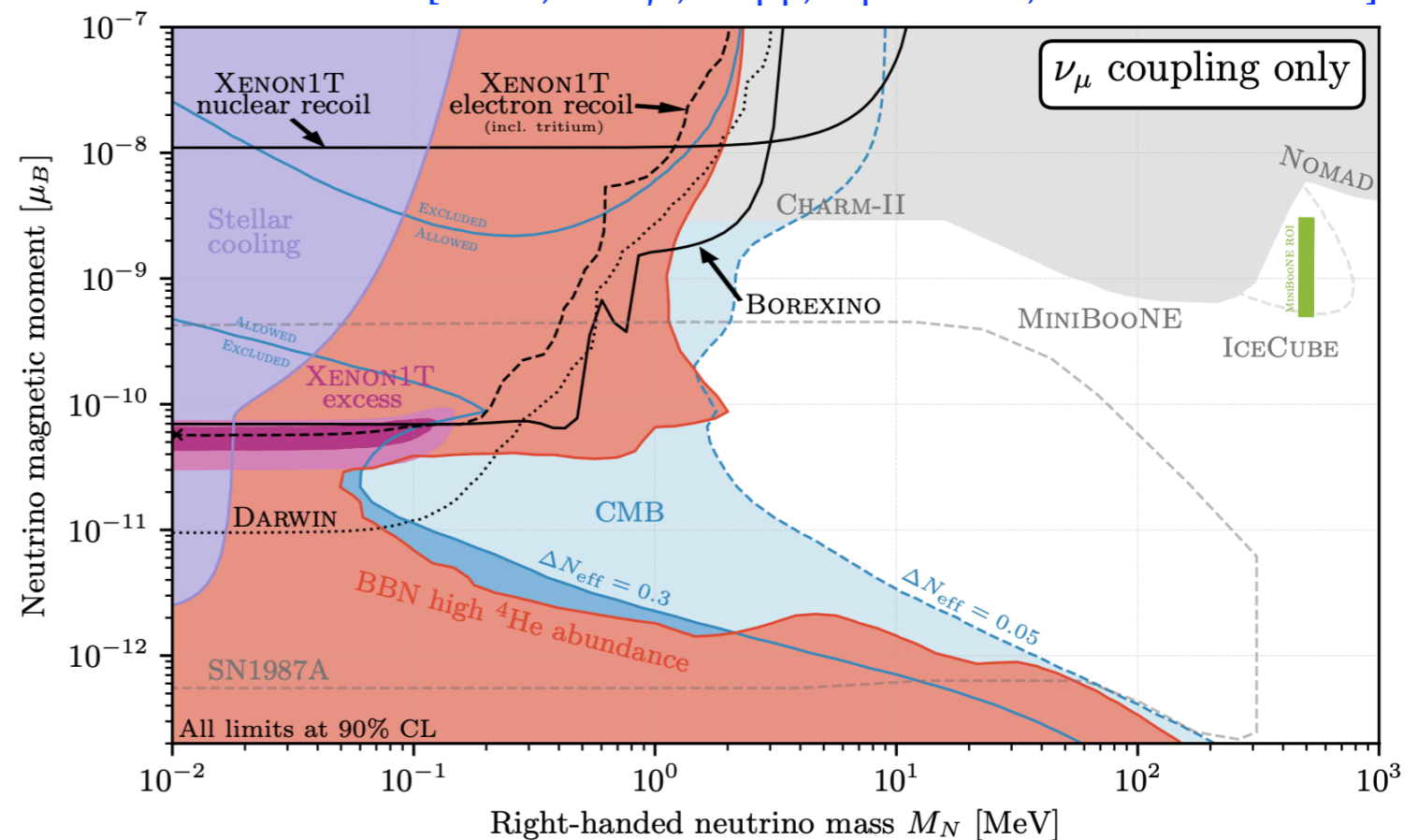
detection via:

$$\nu_L + e^- \rightarrow N + e^-$$

$$\nu_L + X_Z^A \rightarrow N + X_Z^A$$



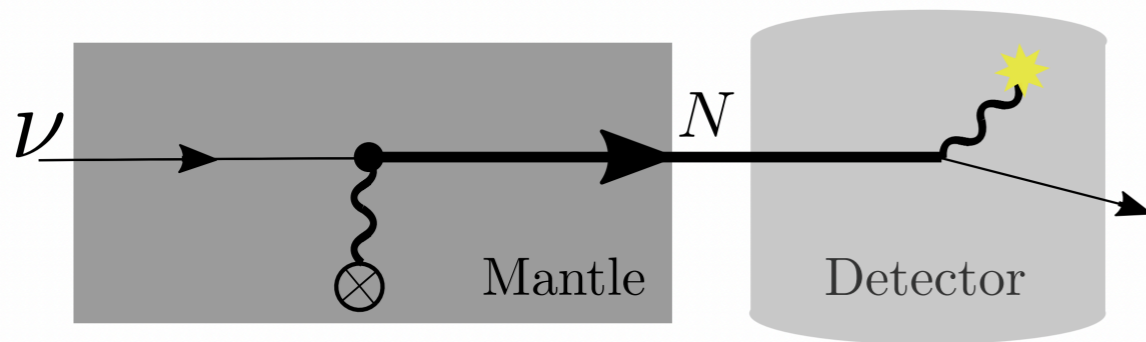
[Magill, Plestid, Pospelov, Tsai, arXiv:1803.03262]
 [Brdar, Greljo, Kopp, Opferkuch, arXiv:2007.15563]



Current Probes: terrestrial experiments

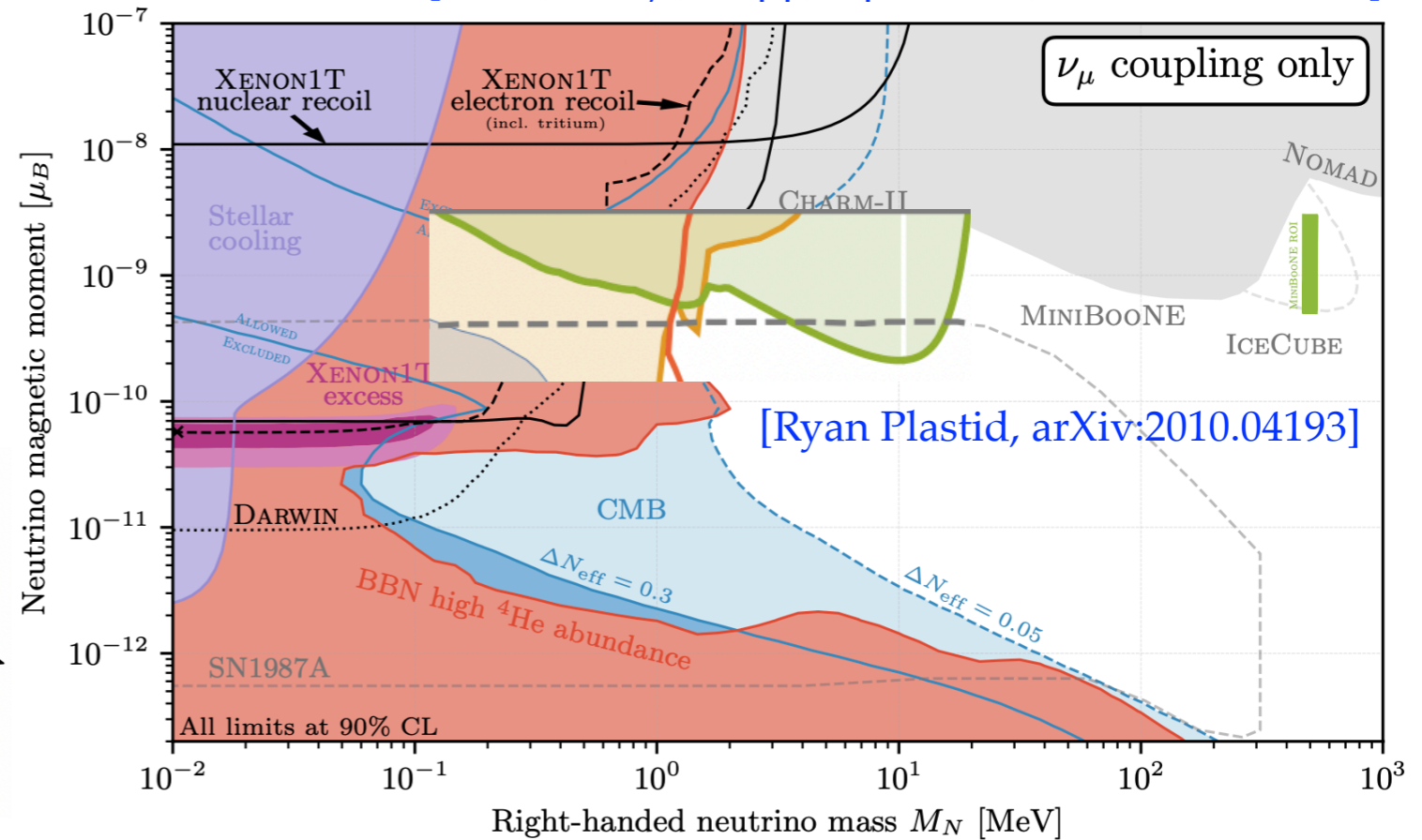
Solar neutrino
up-scattering and decays

production and decay



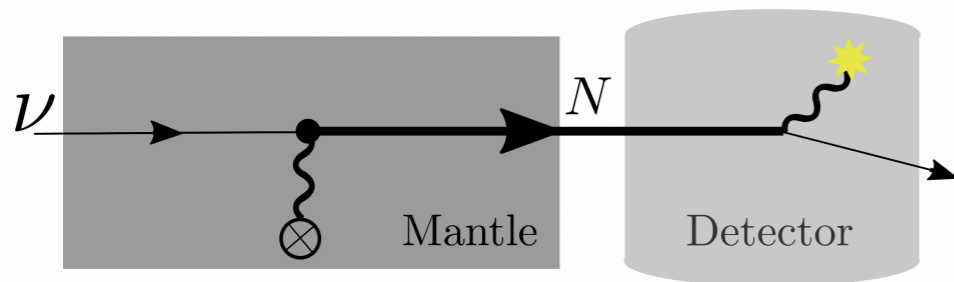
Borexino, SuperK

[Magill, Plestid, Pospelov, Tsai, arXiv:1803.03262]
[Brdar, Greljo, Kopp, Opferkuch, arXiv:2007.15563]

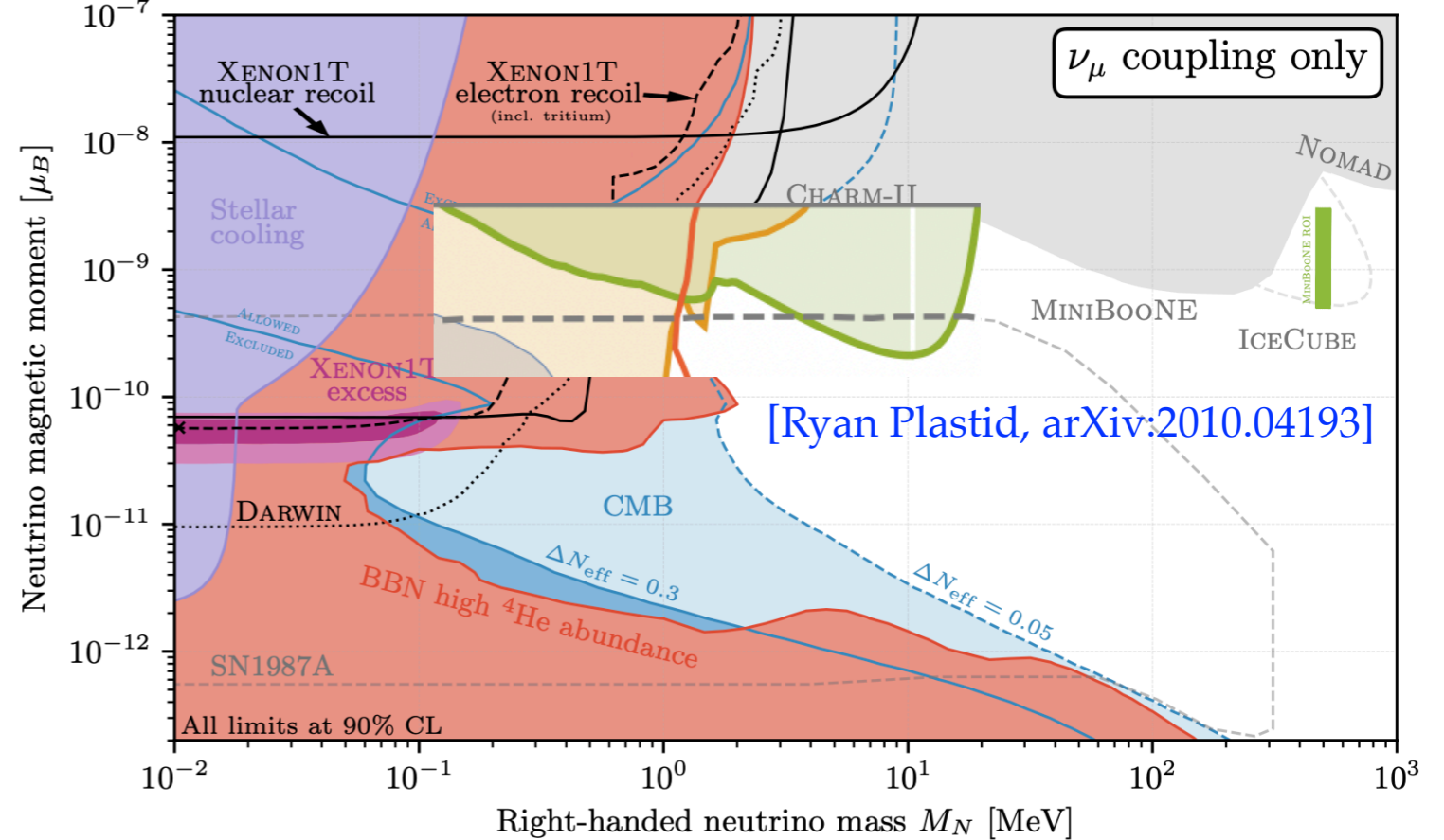
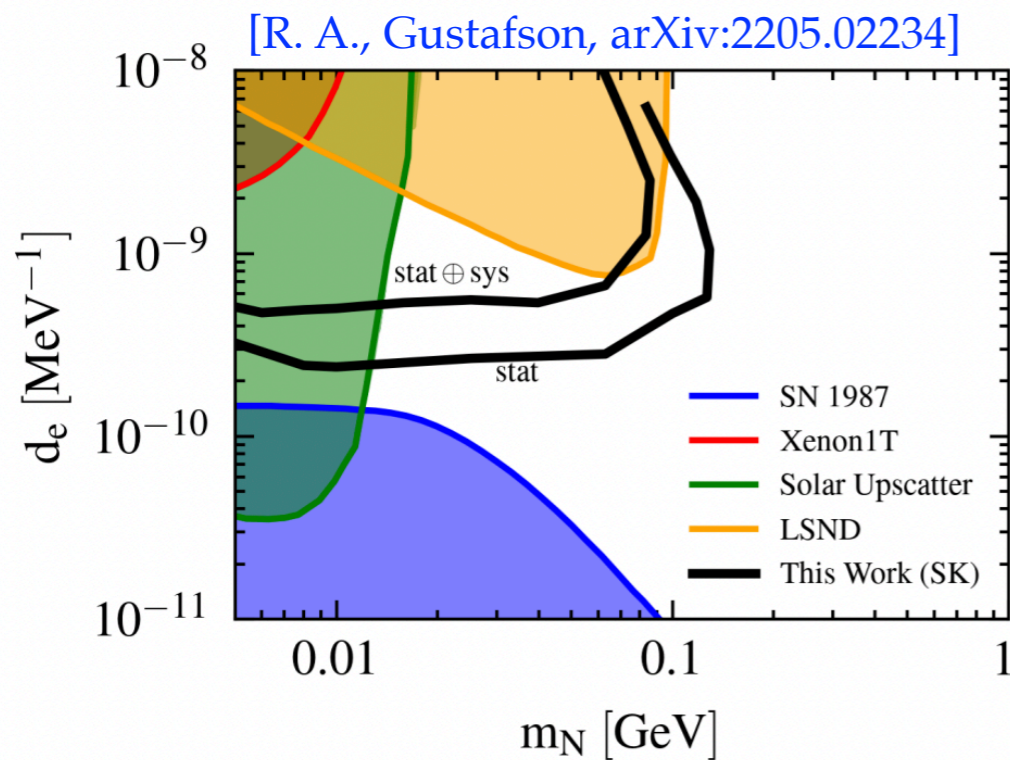


Current Probes: terrestrial experiments

Atmospheric neutrino DUNE, Super-K



[Magill, Plestid, Pospelov, Tsai, arXiv:1803.03262]
 [Brdar, Greljo, Kopp, Opferkuch, arXiv:2007.15563]



$$M_N \gtrsim 6\text{GeV (LHC, LEP): } e^+e^-(q\bar{q}) \rightarrow (N \rightarrow \gamma\nu)\bar{\nu} + h.c.$$

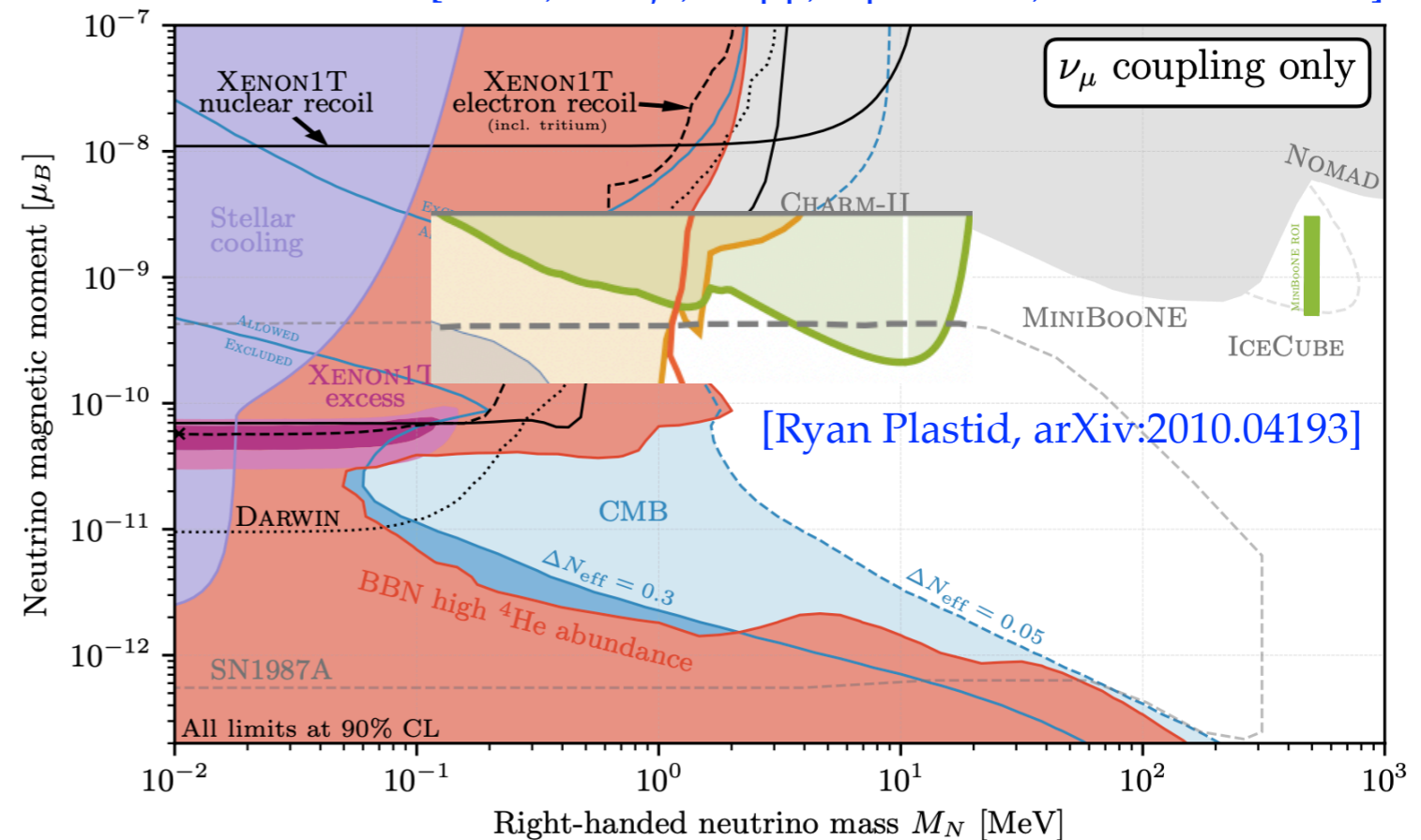
Current Probes: cosmology

CMB, BBN: N_{eff}

- Relativistic: N_{eff}
- Inject extra photons $N \rightarrow \nu\gamma$

$$\begin{aligned} \tau_N &= \frac{16\pi}{\mu_\nu^2 M_N^3} \\ &= 3760 \text{ sec} \times \left(\frac{1 \times 10^{-11} \mu_B}{\mu_\nu} \right)^2 \left(\frac{\text{MeV}}{M_N} \right)^3 \end{aligned}$$

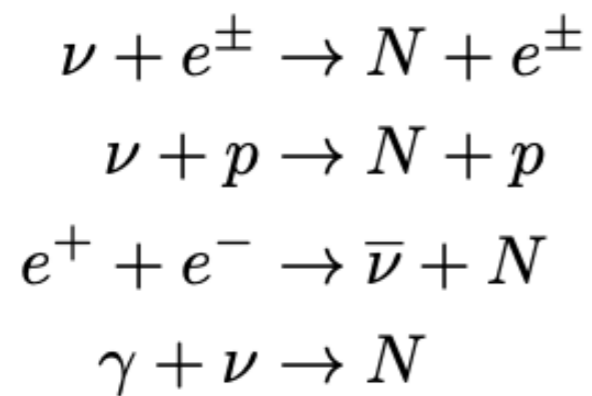
[Magill, Plestid, Pospelov, Tsai, arXiv:1803.03262]
[Brdar, Greljo, Kopp, Opferkuch, arXiv:2007.15563]



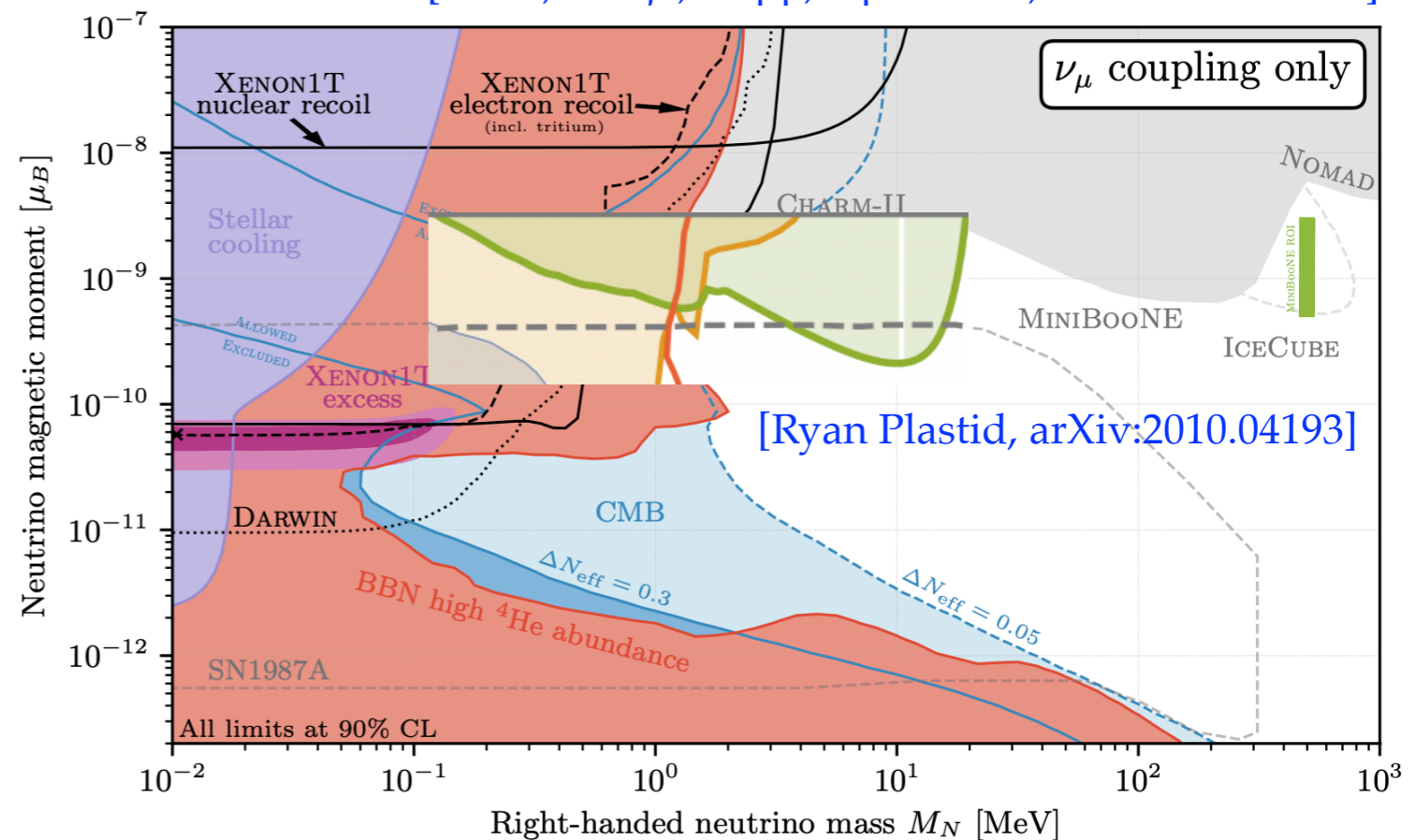
$$T_{\text{dec}} \simeq 1.28 \text{ GeV} \left(\frac{10^{-11} \mu_B}{\mu_\nu} \right)^2$$

Current Probes: Supernova

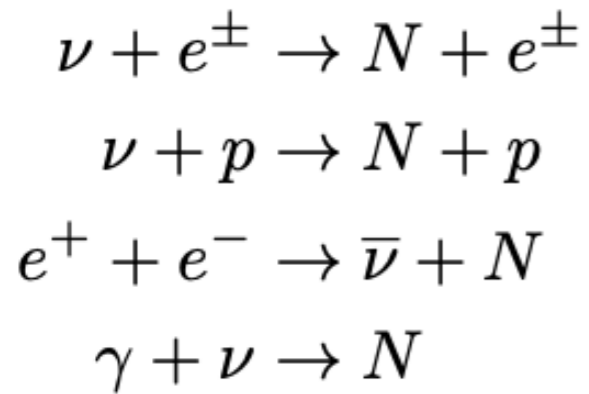
10% of energy loss
to sterile neutrino



[Magill, Plestid, Pospelov, Tsai, arXiv:1803.03262]
[Brdar, Greljo, Kopp, Opferkuch, arXiv:2007.15563]



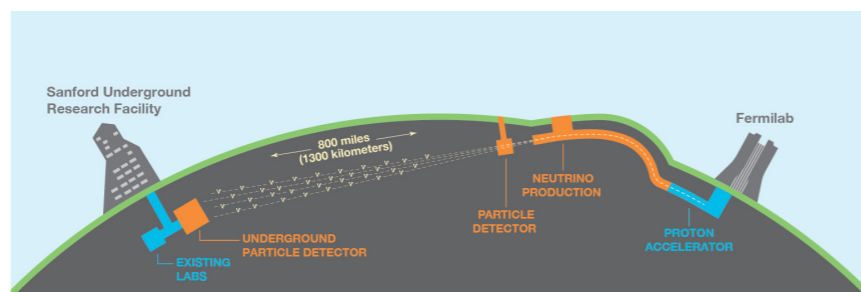
Multi-messenger Signals



$$\mathcal{L} \supset \frac{1}{2} \mu_\nu \bar{\nu}_L^\alpha \sigma^{\mu\nu} N F_{\mu\nu}$$



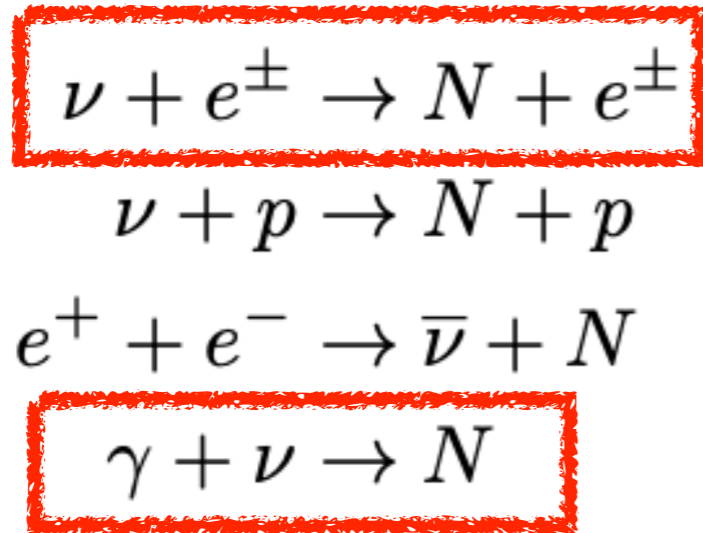
$$N \rightarrow \nu + \gamma$$



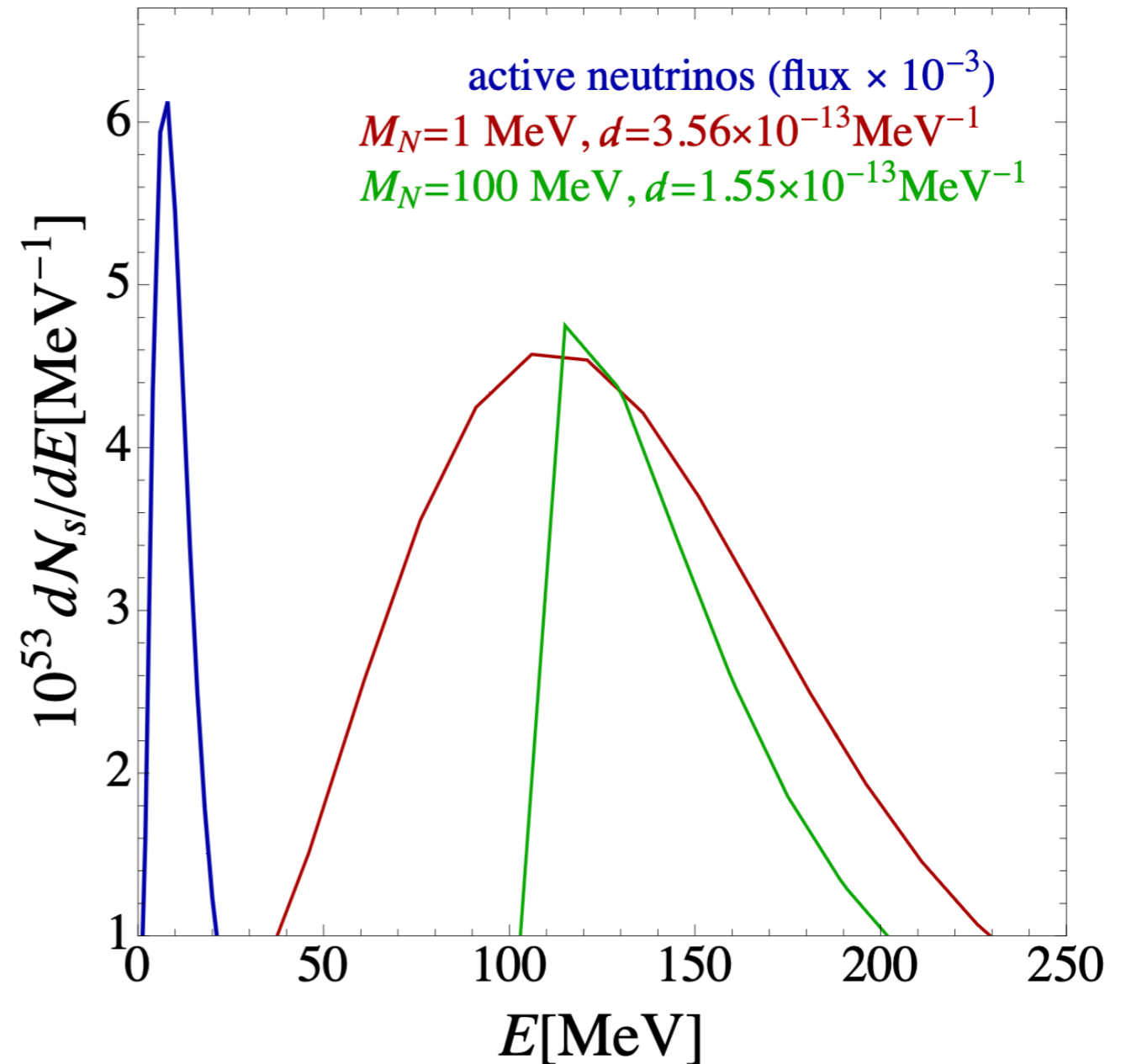
[V. Brdar, A. de Gouvêa, YYL, P. A. N. Machado, arXiv:2302.10965]

Multi-messenger Signals

[V. Brdar, A. de Gouvêa, YYL, P. A. N. Machado, arXiv:2302.10965]

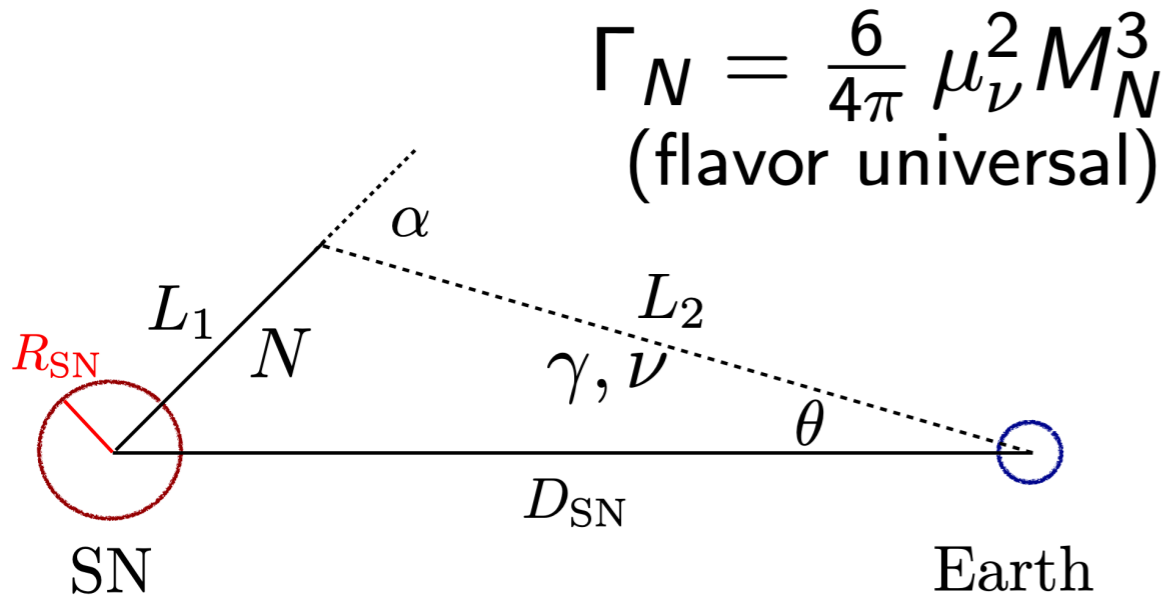


$$\frac{1}{4\pi r^2} \frac{\partial^2}{\partial r \partial t} \left(\frac{d\mathcal{N}_s}{dE_N} \right) = \sigma n_e \frac{dn_\nu}{dE}$$



HNL are produced in a dense $T \sim 100 \text{ MeV}$ core and leave subsequently the star without further interactions

Multi-messenger Signals



$$\cos \alpha = \frac{2E_N E_{\gamma/\nu} - M_N^2}{2E_{\gamma/\nu} \sqrt{E_N^2 - M_N^2}}$$

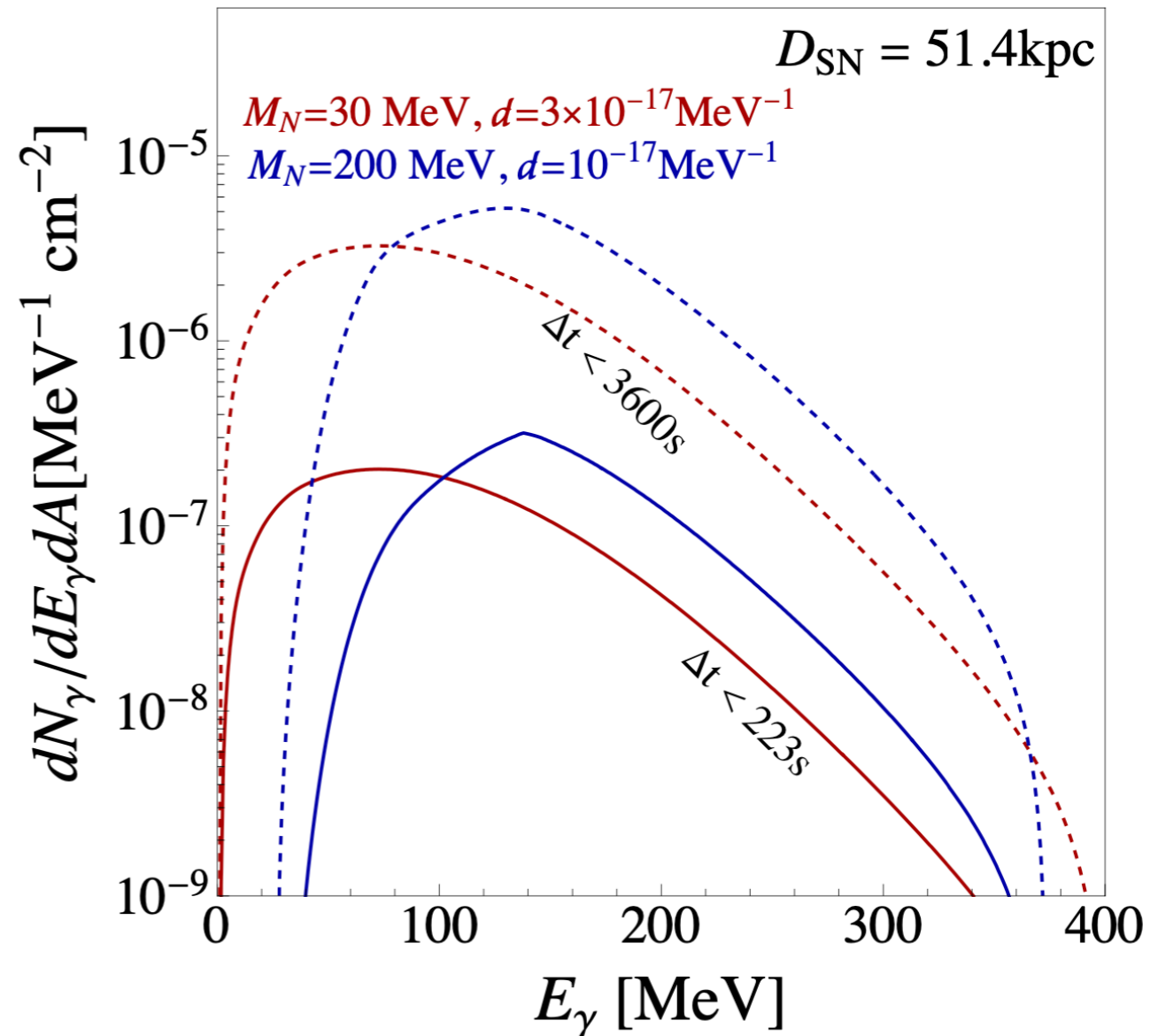
$$\Delta t = L_1/\beta + L_2 - D_{\text{SN}}$$

$$\beta = \sqrt{E_N^2 - M_N^2}/E_N$$

$$R_{\text{SN}}^{\gamma/\nu} \leq L_1 \leq L_1^{\text{max}}$$

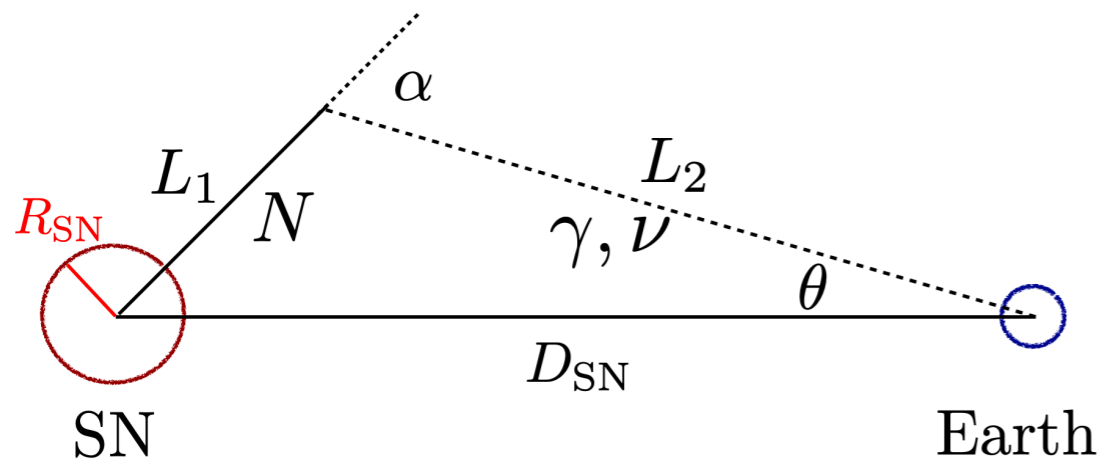
$$R_{\text{SN}}^\gamma = 3 \times 10^{10} \text{ m}$$

$$R_{\text{SN}}^\nu = 30 \text{ km}$$

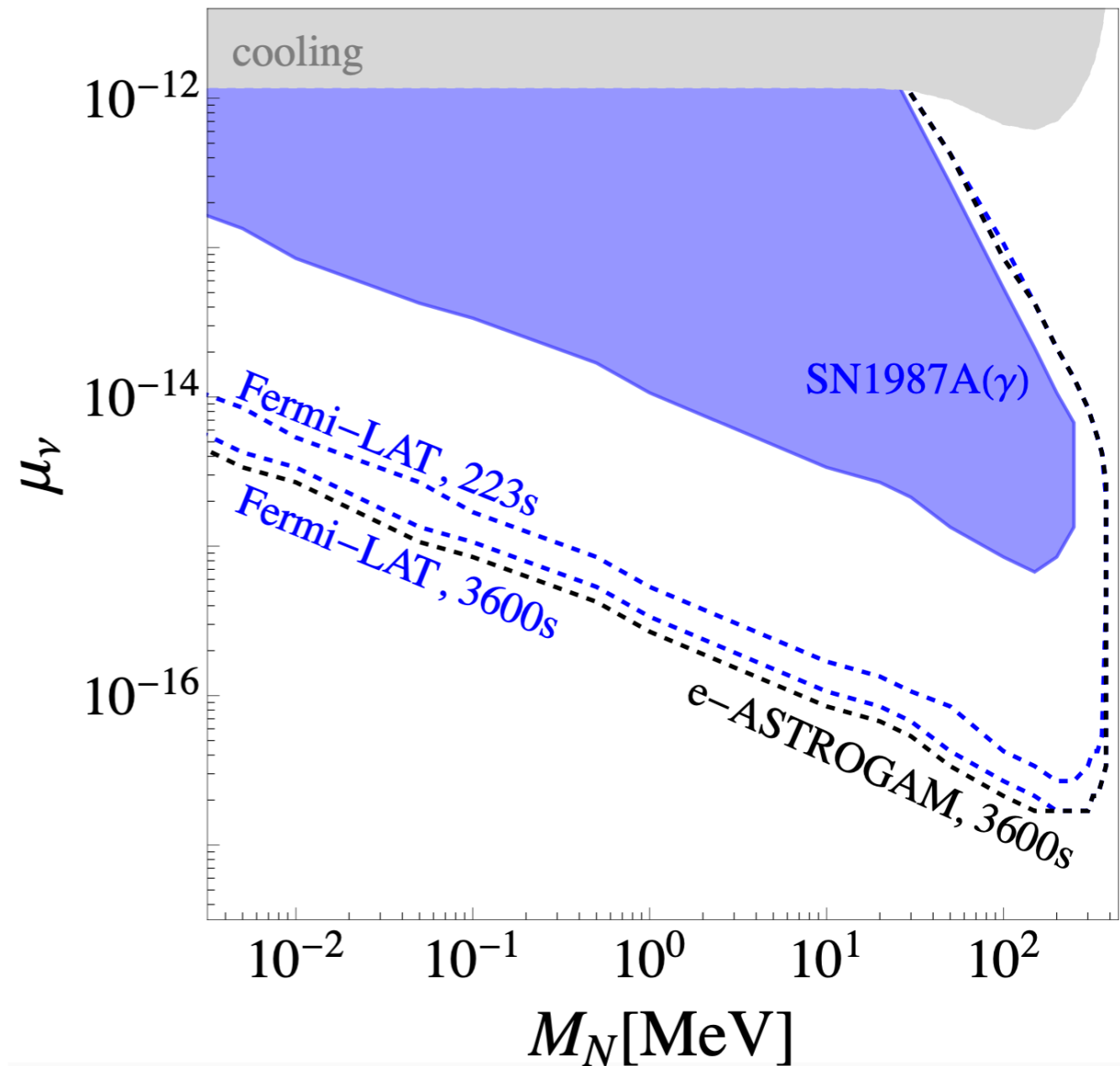


[V. Brdar, A. de Gouvêa, YYL, P. A. N. Machado, arXiv:2302.10965]

Multi-messenger Signals : γ - ray detection



- At the time of SN1987A, the **Gamma-Ray Spectrometer (GRS)** observed $N_{\text{obs}} = 1393$ photons with energy 25-100 MeV at $\Delta t < 223s$
- Assuming a SN event happens in the galaxy at a distance of $D_{\text{SN}} = 10\text{kpc}$,
Fermi-LAT: $E_{\gamma} > 100\text{MeV}, \theta < 5^{\circ}$
e-ASTROGAM: $E_{\gamma} > 1\text{MeV}, \theta < 1.25^{\circ}$

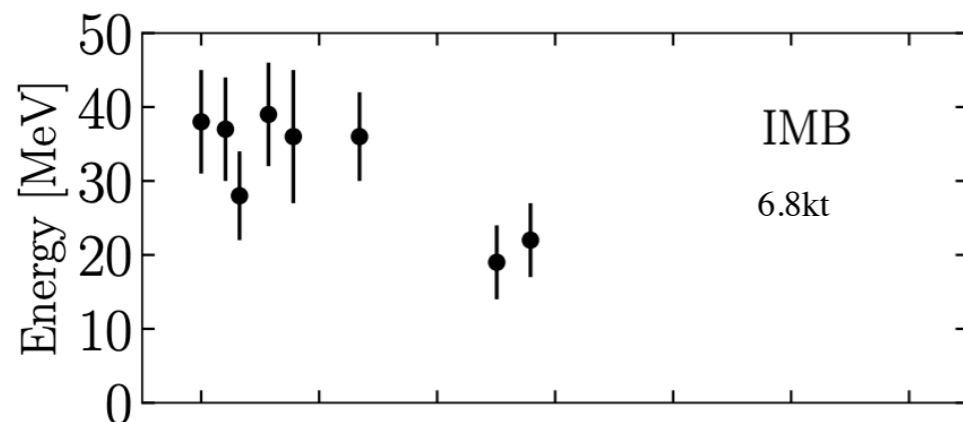
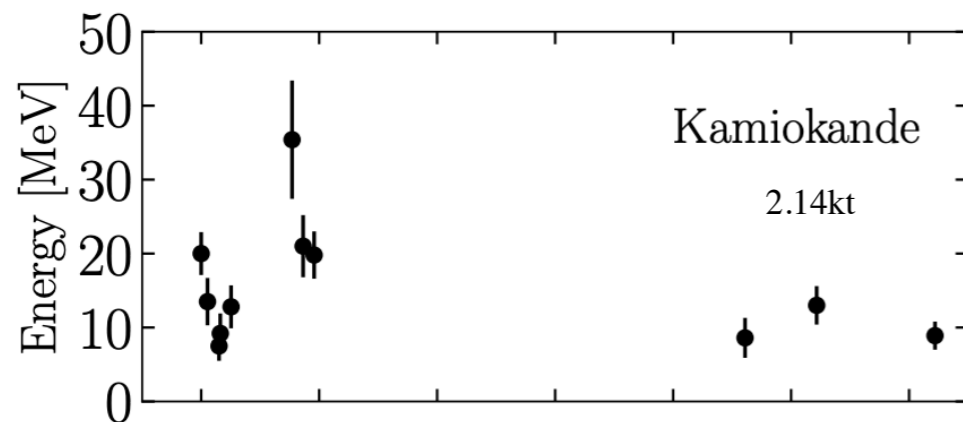


[V. Brdar, A. de Gouvêa, YYL, P. A. N. Machado, arXiv:2302.10965]

Multi-messenger Signals : neutrino detection

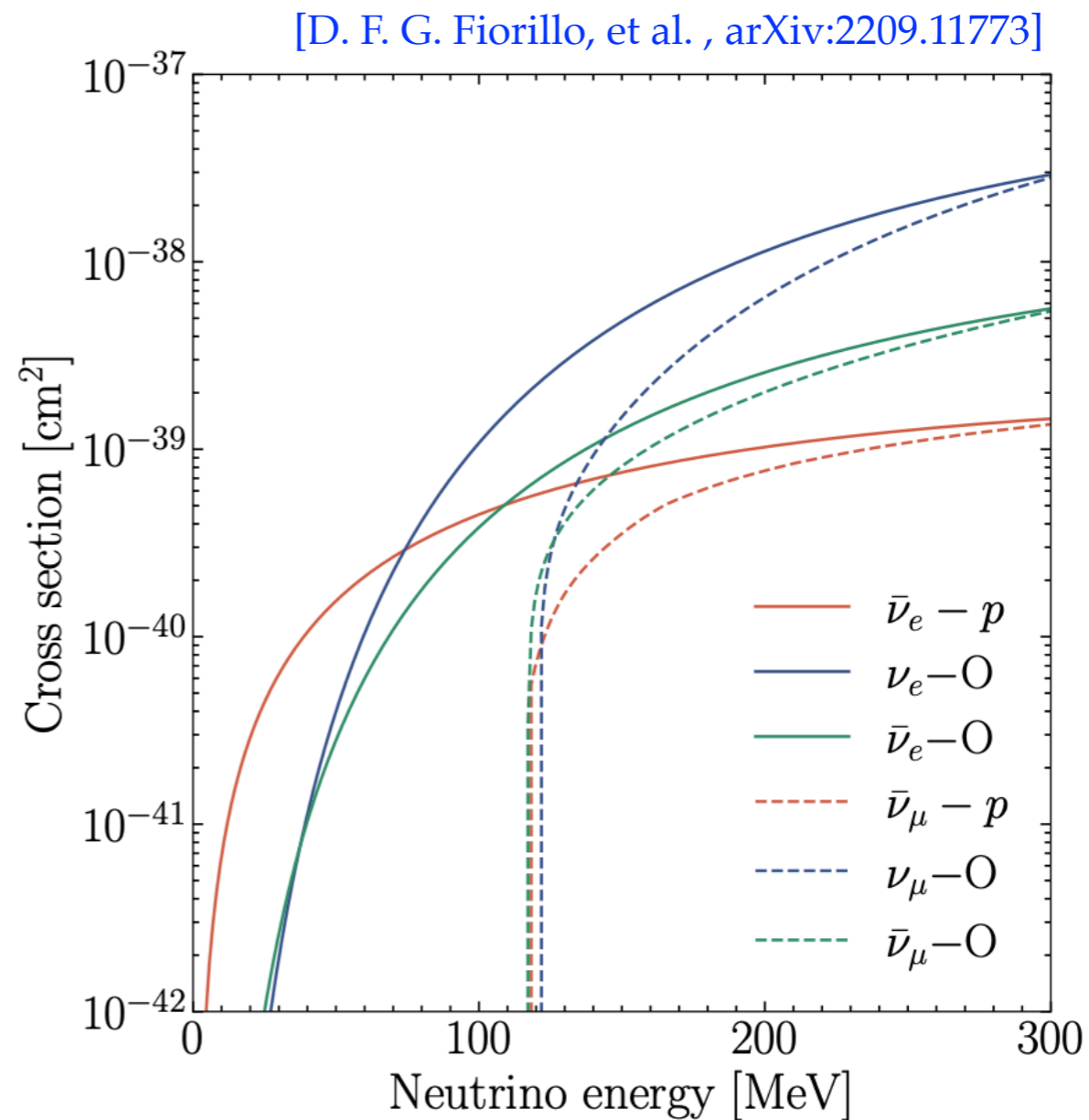
SN1987A, neutrino events

water-Cherenkov detectors



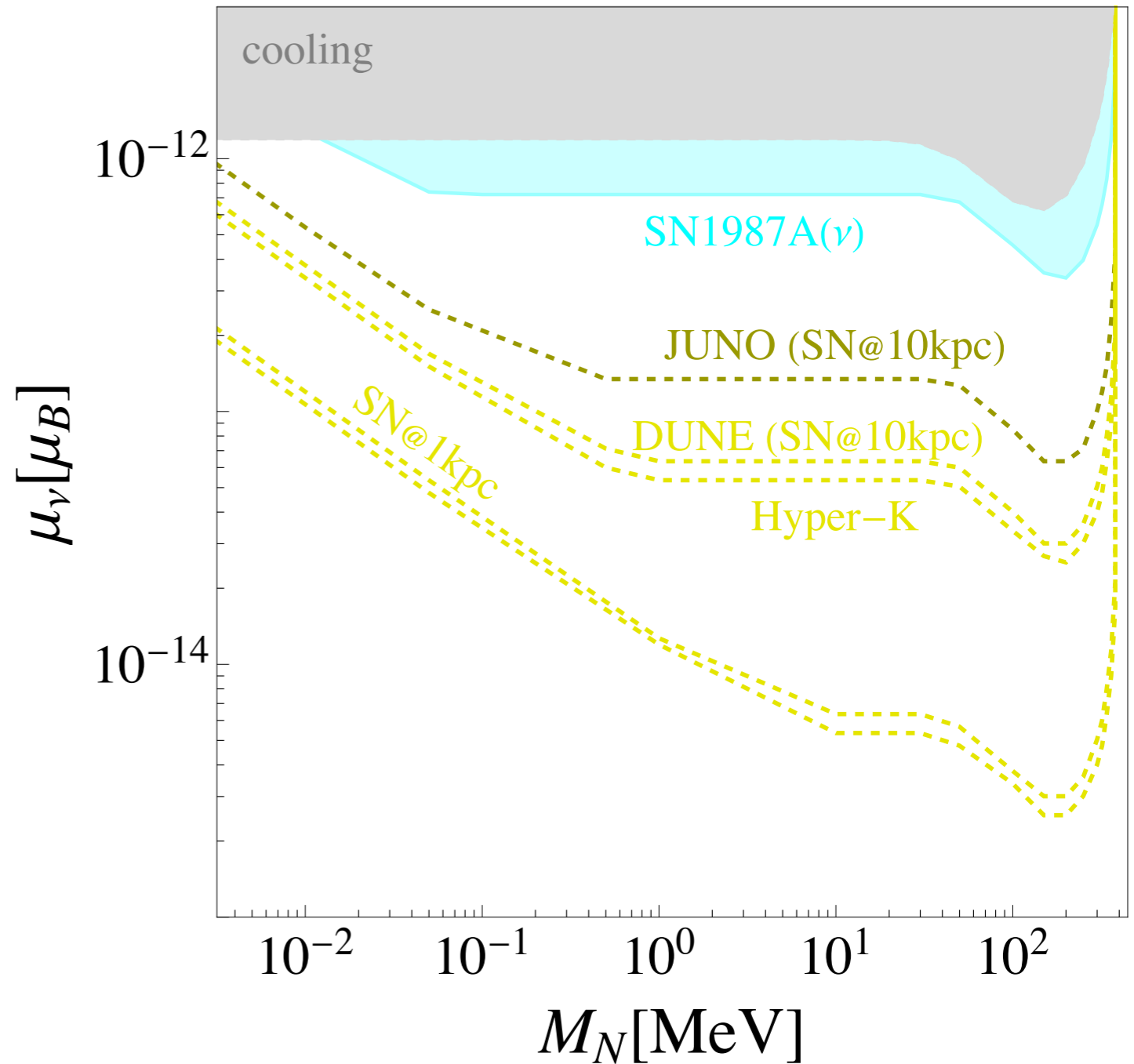
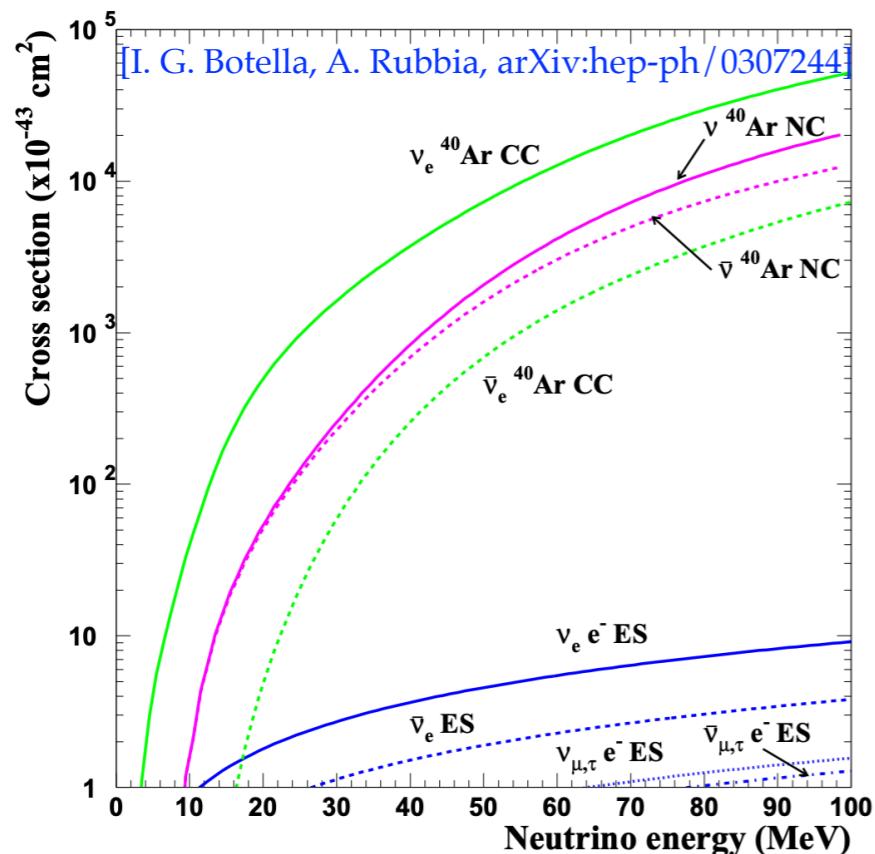
- No significant excess was observed by Kamiokande-II and IMB for $E_\nu > 50\text{MeV}, \Delta t < 2$ days

$$N_\nu^{\text{BSM}} = N_{\text{tgt}} \int dE_\nu \frac{dN_\nu}{dE_\nu dA} (E_\nu) \sigma_{\text{IBD}}(E_\nu) \varepsilon(E_\nu)$$



Multi-messenger Signals : neutrino detection

- Assuming a SN event happens in the galaxy at a distance of $D_{\text{SN}} = 10\text{kpc}$,
- JUNO**: 20kt fiducial volume, liquid scintillator detector
- DUNE**: 40kt, liquid argon
- Hyper-K**: 188kt fiducial volume water Cherenkov

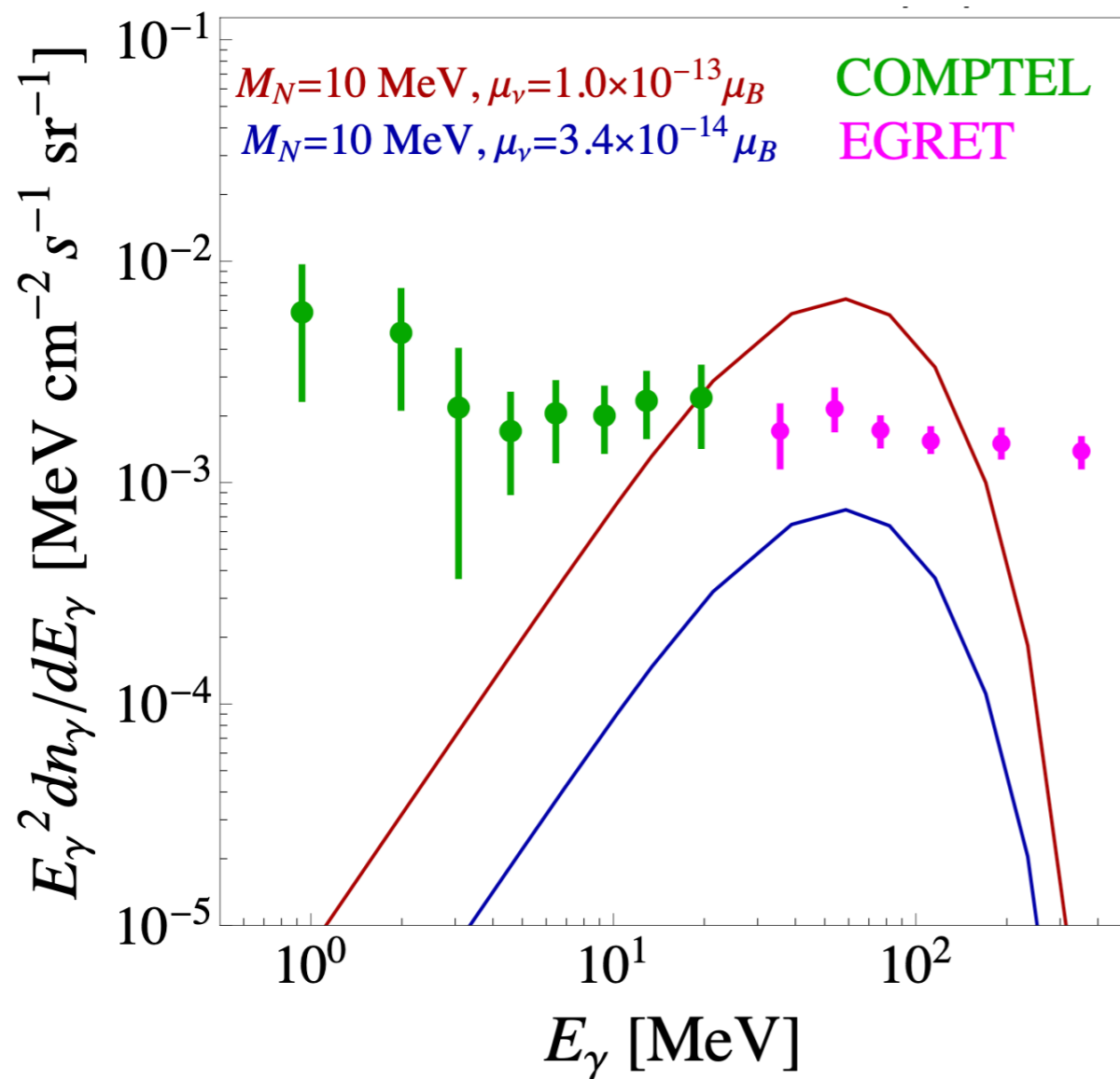


[V. Brdar, A. de Gouvêa, YYL, P. A. N. Machado, arXiv:2302.10965]

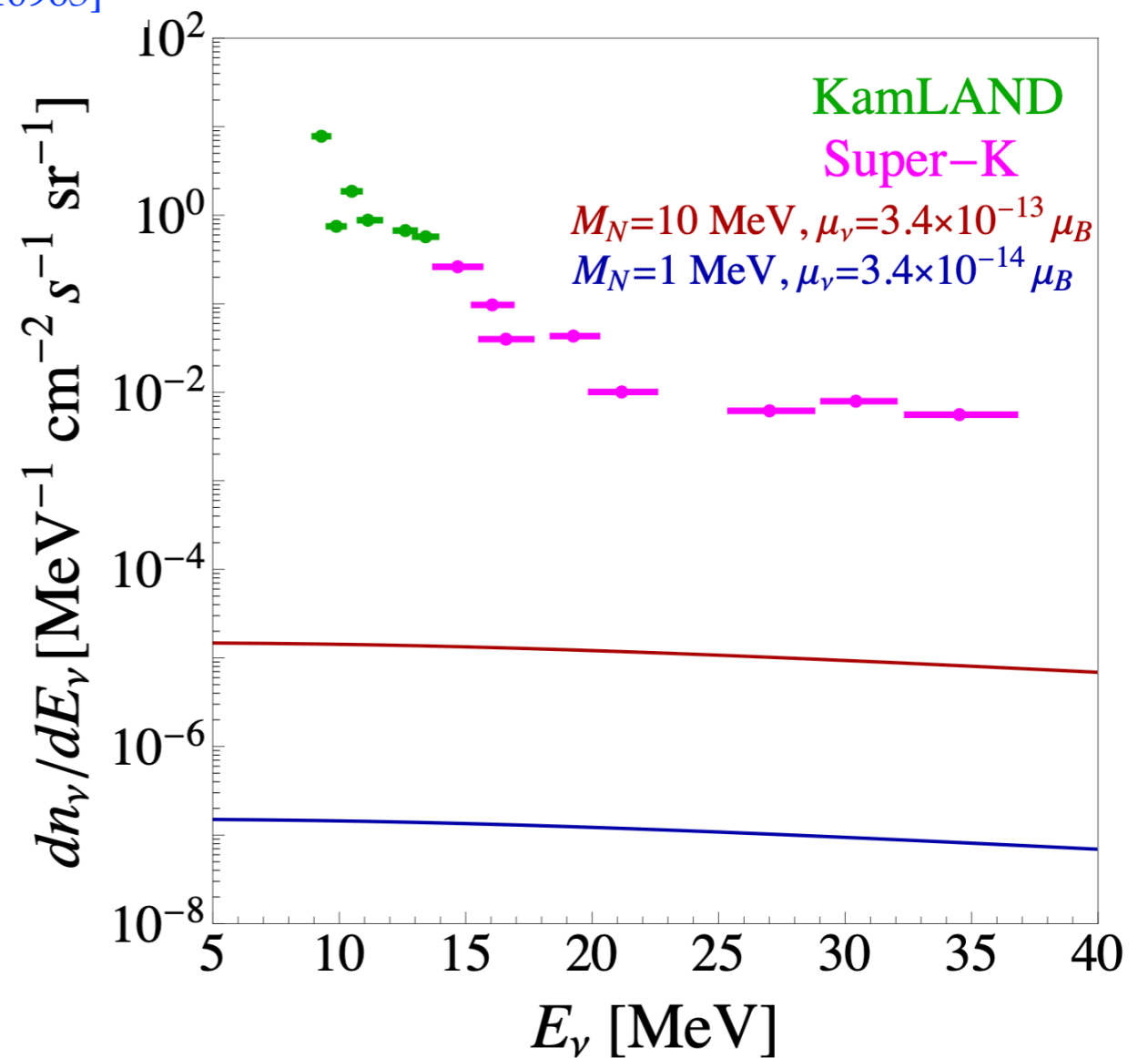
Multi-messenger Signals : diffused BSM Photon and neutrino background

$$\frac{dn_N}{dE} = \frac{c}{4\pi} \int_0^\infty dz (1+z) n'_{cc}(z) \frac{d\mathcal{N}_s}{dE}(E_z)$$

[V. Brdar, A. de Gouvêa, YYL, P. A. N. Machado, arXiv:2302.10965]



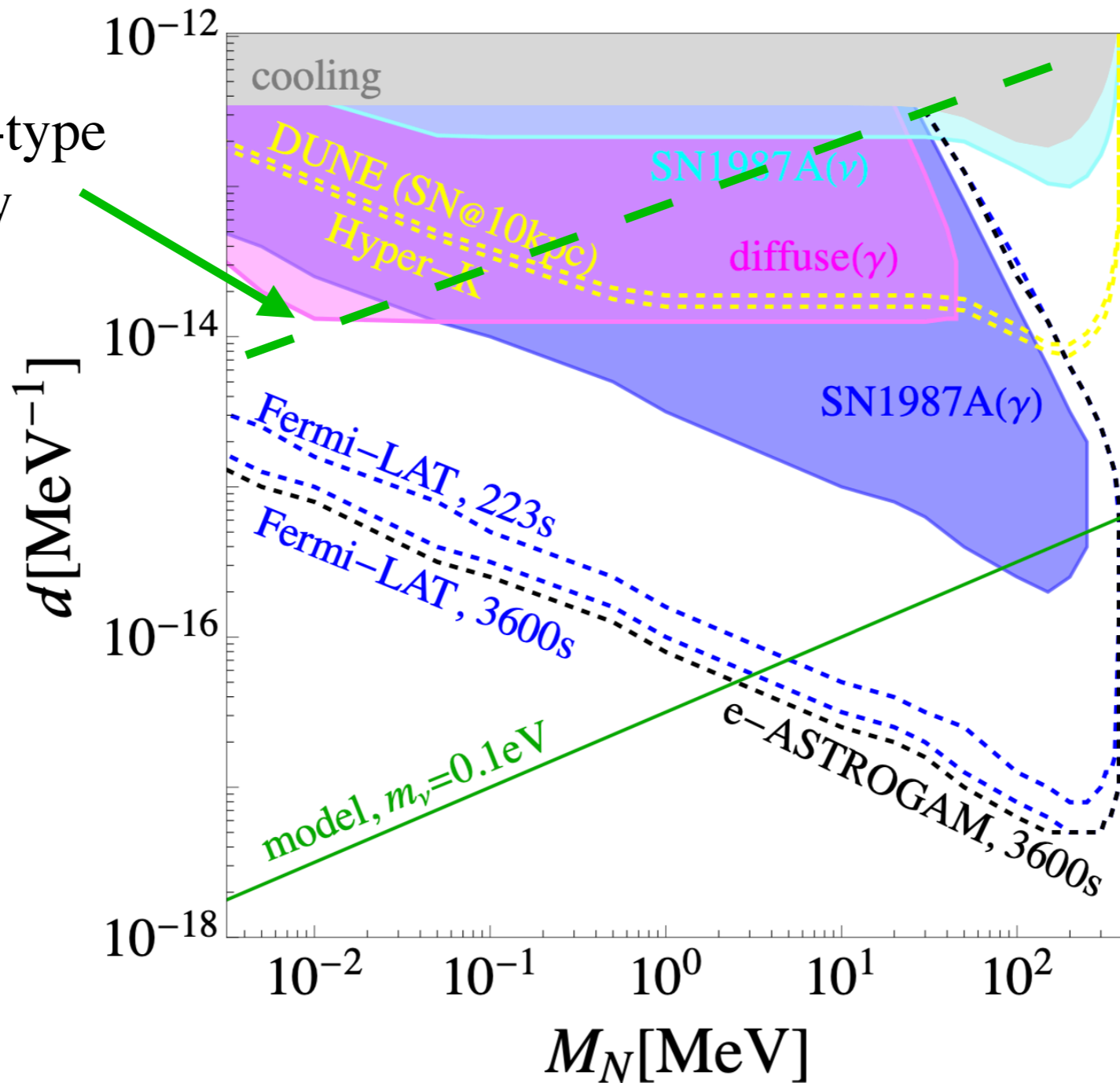
extragalactic background light



Multi-messenger Signals : diffused BSM Photon and neutrino background

$$\mathcal{L} \supset \frac{1}{2} \mu_\nu \bar{\nu}_L^\alpha \sigma^{\mu\nu} N F_{\mu\nu}$$

Voloshin-type
symmetry



[V. Brdar, A. de Gouvêa, YYL, P. A. N. Machado, arXiv:2302.10965]

Left - Right symmetric model

$$SU(2)_L \times SU(2)_R \times U(1)_{B-L}$$

$$\rightarrow SU(2)_L \times U(1)_Y \rightarrow U(1)_Q$$

$$\begin{pmatrix} W_L^\pm \\ W_R^\pm \end{pmatrix} = \begin{pmatrix} \cos \zeta & -\sin \zeta \\ \sin \zeta & \cos \zeta \end{pmatrix} \begin{pmatrix} W_1^\pm \\ W_2^\pm \end{pmatrix}$$

$$Q_L^i \left(\frac{1}{2}, 0, \frac{1}{3} \right): \begin{pmatrix} u_L \\ d_L \end{pmatrix}^\alpha, \quad \begin{pmatrix} c_L \\ s_L \end{pmatrix}^\alpha, \quad \begin{pmatrix} t_L \\ b_L \end{pmatrix}^\alpha$$

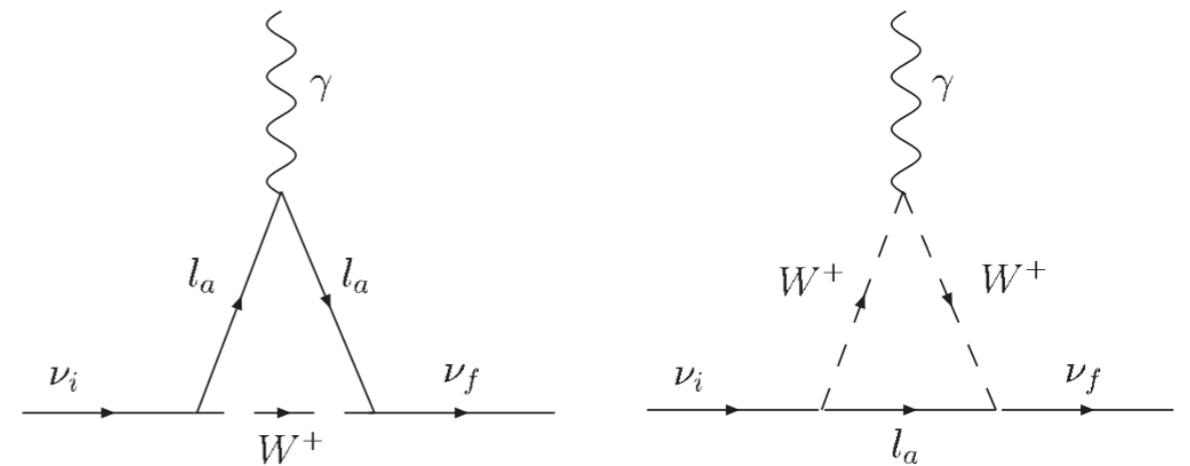
$$Q_R^i \left(0, \frac{1}{2}, \frac{1}{3} \right): \begin{pmatrix} u_R \\ d_R \end{pmatrix}^\alpha, \quad \begin{pmatrix} c_R \\ s_R \end{pmatrix}^\alpha, \quad \begin{pmatrix} t_R \\ b_R \end{pmatrix}^\alpha$$

$$\Psi_L^a \left(\frac{1}{2}, 0, -1 \right): \begin{pmatrix} \nu_{eL} \\ e_L^- \end{pmatrix}, \quad \begin{pmatrix} \nu_{\mu L} \\ \mu_L^- \end{pmatrix}, \quad \begin{pmatrix} \nu_{\tau L} \\ \tau_L^- \end{pmatrix},$$

$$\Psi_R^a \left(0, \frac{1}{2}, -1 \right): \begin{pmatrix} N_{eR} \\ e_R^- \end{pmatrix}, \quad \begin{pmatrix} N_{\mu R} \\ \mu_R^- \end{pmatrix}, \quad \begin{pmatrix} N_{\tau R} \\ \tau_R^- \end{pmatrix}$$

type-II seesaw dominance

$$U = \begin{pmatrix} U_L & 0_{3 \times 3} \\ 0_{3 \times 3} & U_R^* \end{pmatrix}$$



$$\mathcal{L}_{\text{Yukawa}}^L = \{ \bar{L}_{Li} (h_{ij} \phi + \tilde{h}_{ij} \tilde{\phi}) L_{Rj} + \text{h.c.} \}$$

$$+ \{ \bar{L}_{Ri}^c (h_M)_{ij} \Sigma_L L_{Lj} + \bar{L}_{Li}^c (h_M)_{ij} \Sigma_R L_{Rj} + \text{h.c.} \}$$

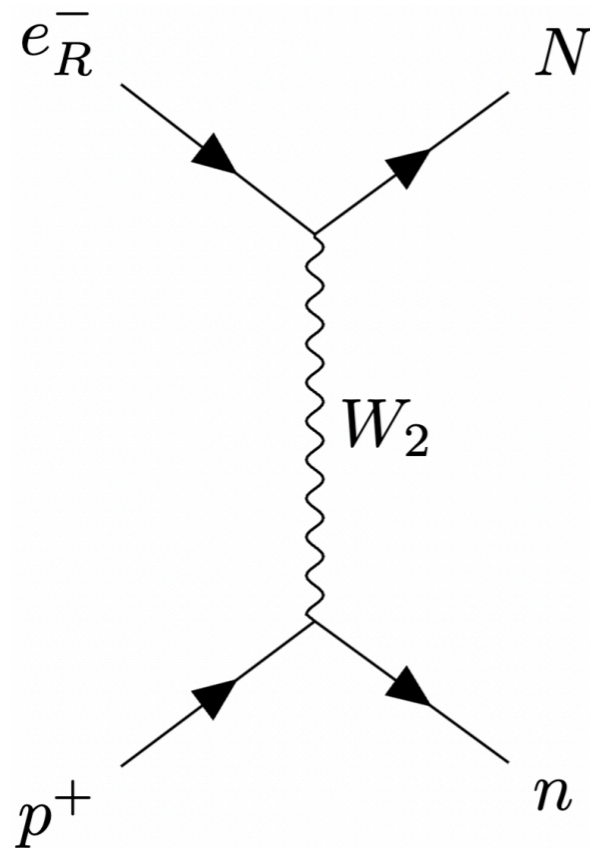
bi-doublet

triplet

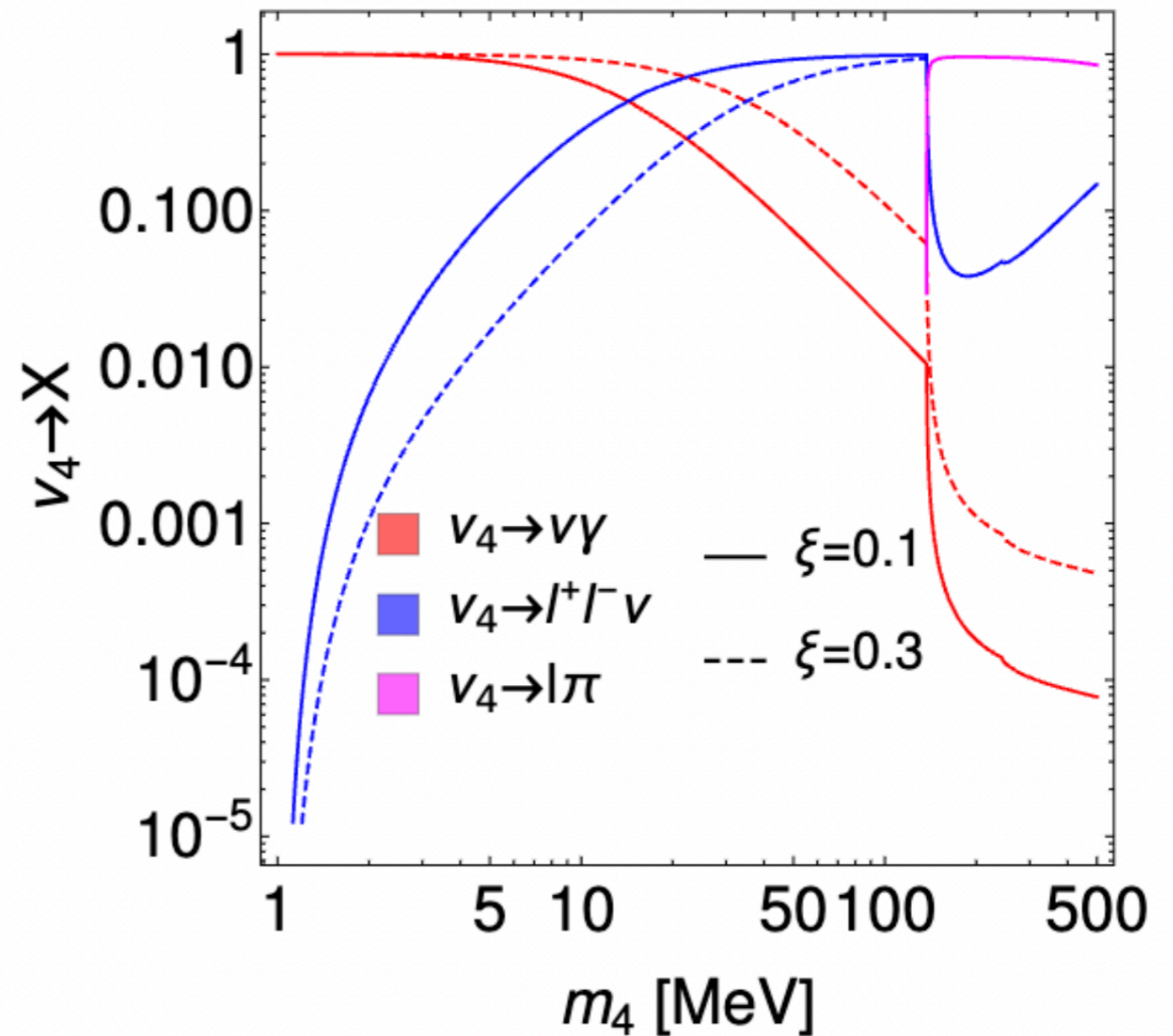
$$\mu_{ji} = \frac{eG_F}{2\sqrt{2}\pi^2} \sin 2\zeta \sum_{\alpha=1}^3 m_{l\alpha} \text{Im} [(PU^*)_{\alpha j} (P_s U^*)_{\alpha i}]$$

Left - Right symmetric model

Tree level production in SN



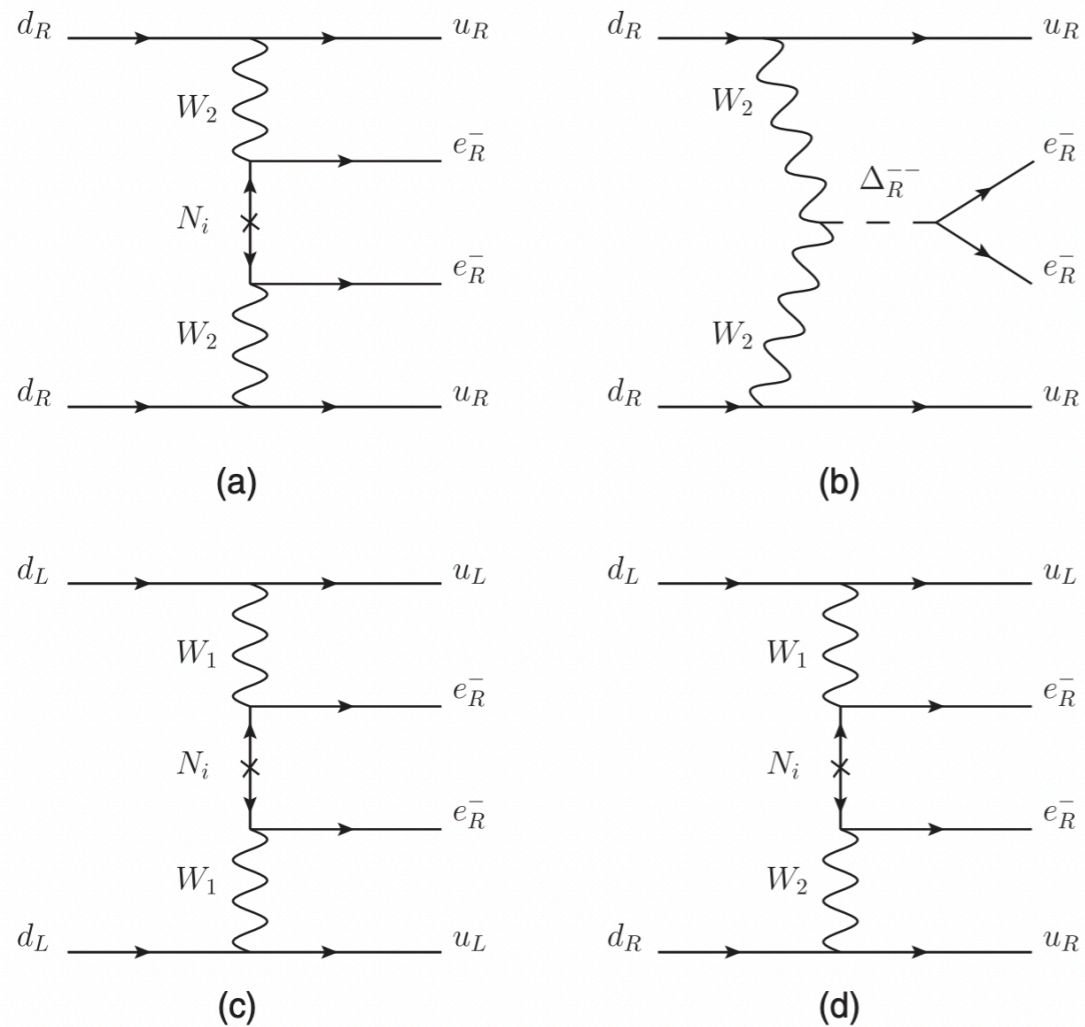
including three-body decay channel



[G. Li, YYL, Y.-L. Zhou, in preparation]

Left - Right symmetric model

Majorana Case

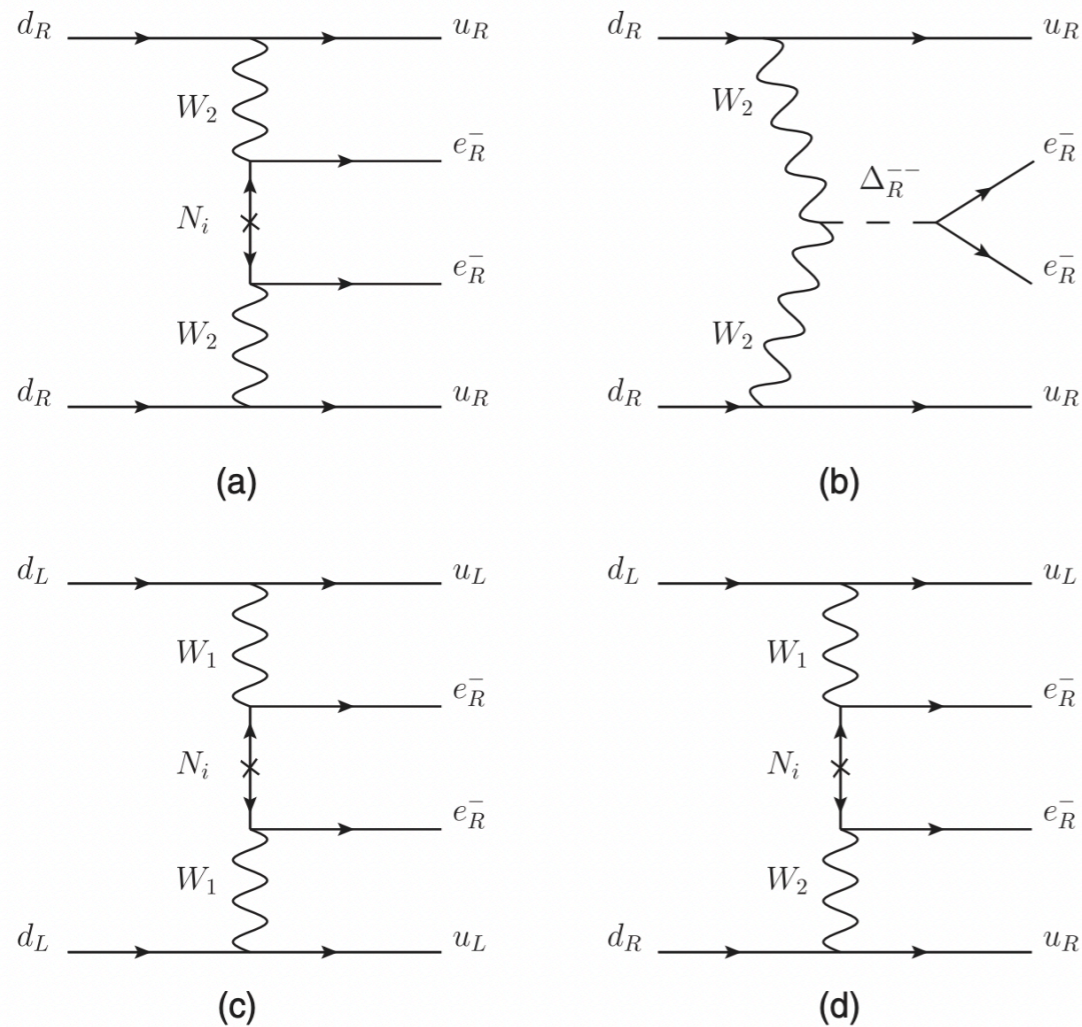


next generation of $0\nu\beta\beta$ decay
searches with ton-scale detectors

[G. Li et al., arXiv:2009.01257]

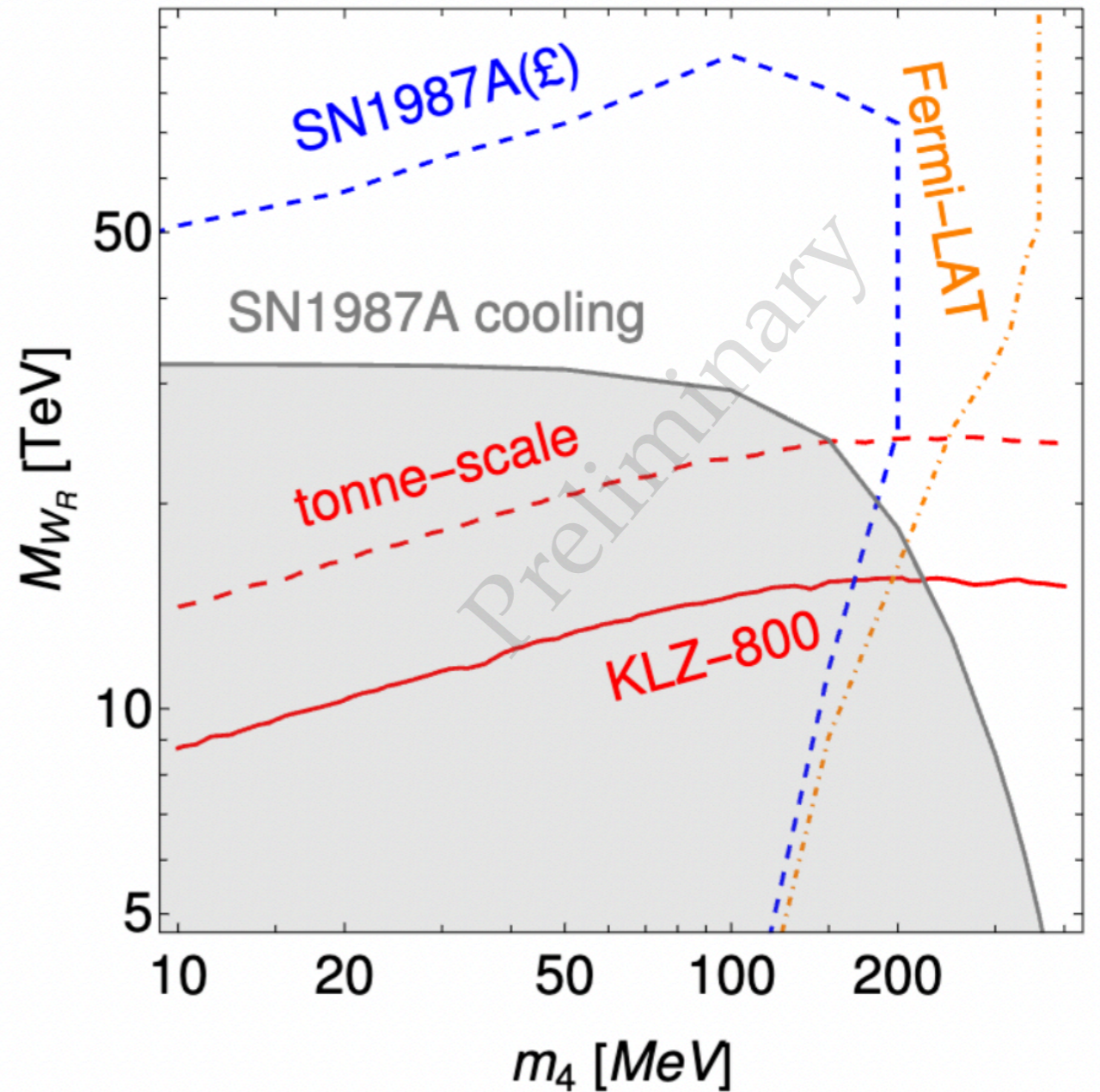
Left - Right symmetric model

Majorana Case



next generation of $0\nu\beta\beta$ decay searches with ton-scale detectors

[G. Li et al., arXiv:2009.01257]



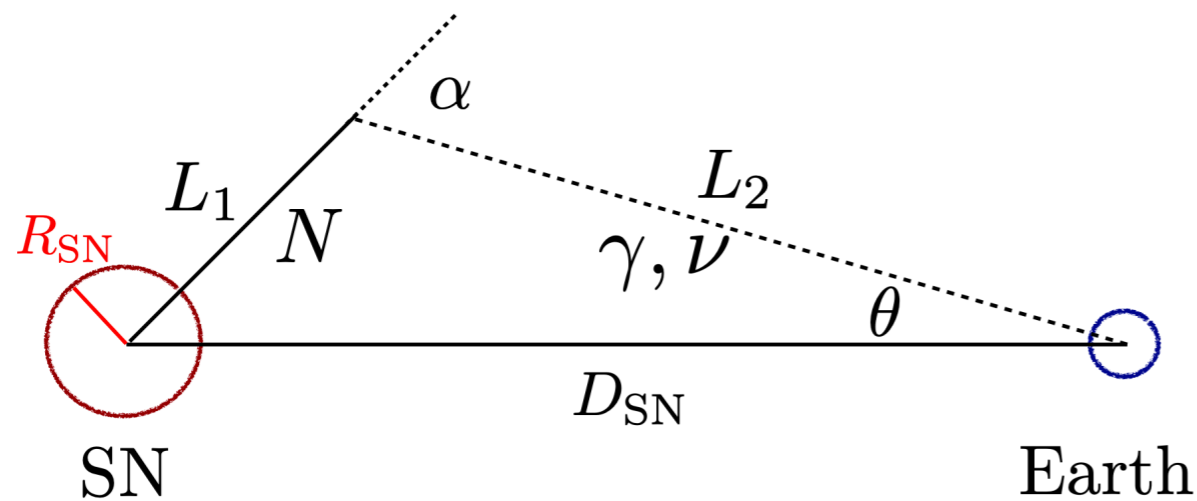
[G. Li, YYL, Y.-L. Zhou, in preparation]

Conclusion

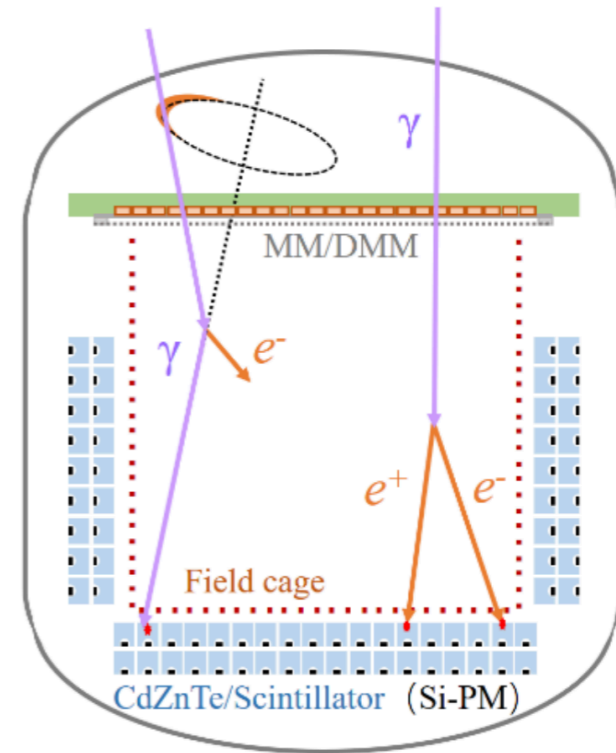
Multi-messenger Signals: Photon + Neutrino

Neutrino Timing Information

Outlooks

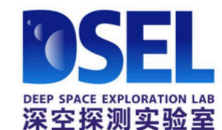


photon polarization



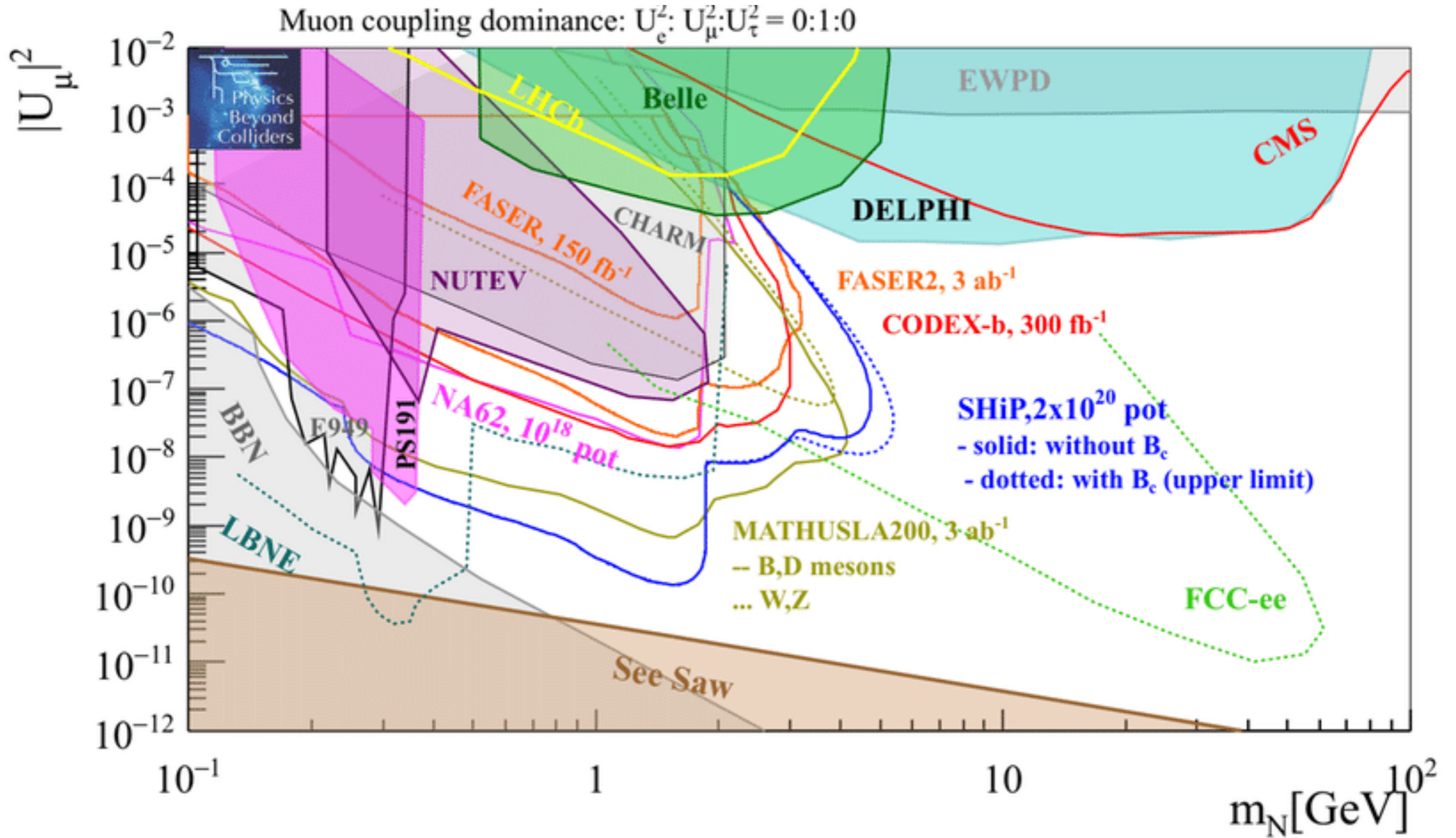
CP violation in neutral lepton transition dipole moment?

Photon signals from sterile neutrino decay or axion decay?

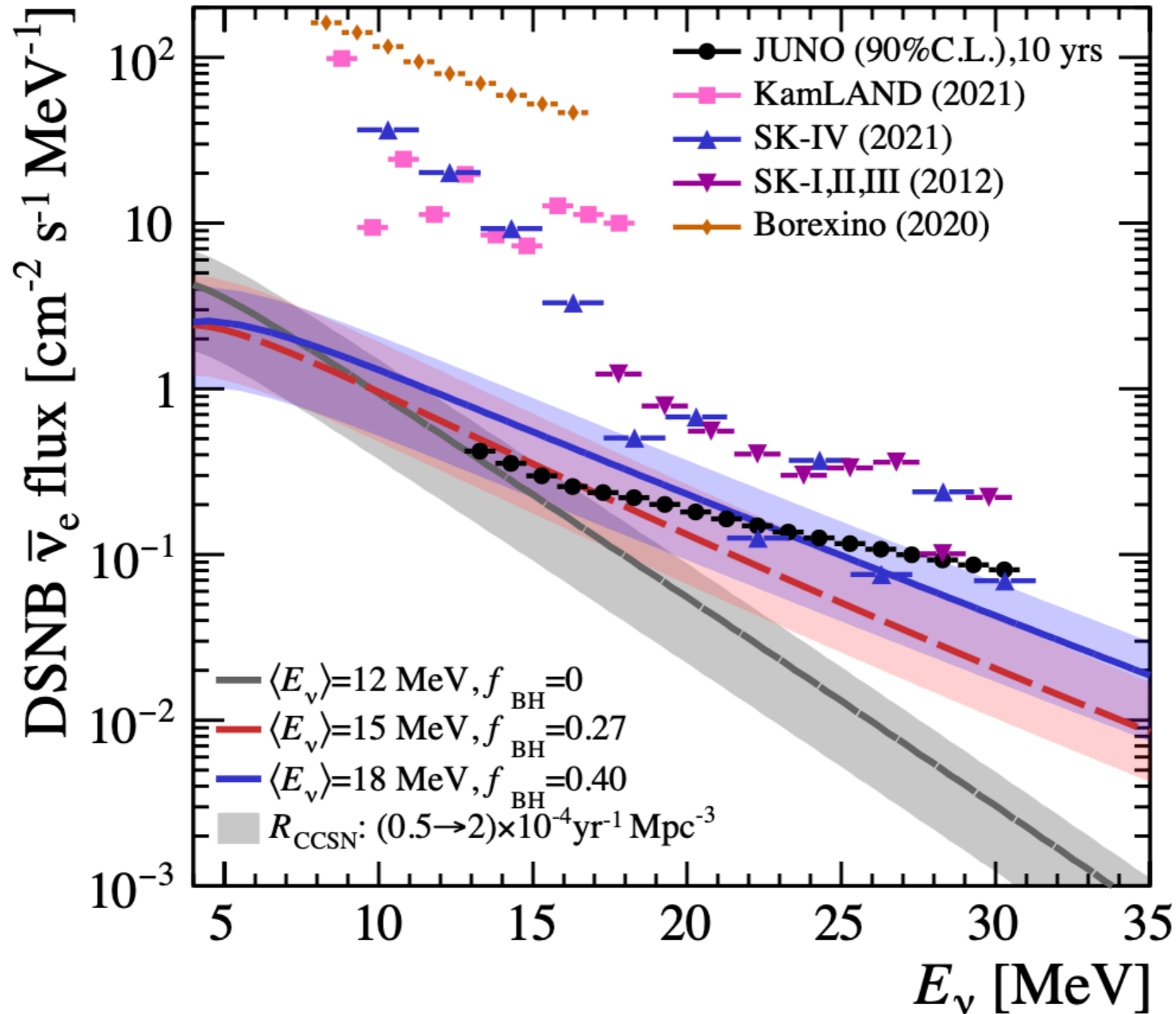


Thank you

BACK UP



Multimessenger Signals : diffused BSM Photon and neutrino background



[A. Abusleme, arXiv:2205.08830]