

Preliminary Investigation of a Higgs Factory based on Proton-Driven Plasma Wakefield Acceleration

ALEGRO

Lissabon, 19-22 March 2024

J. Farmer, A. Caldwell, and A. Pukhov



Introduction

A Higgs factory is a high priority for particle physics.

Plasma wakefield acceleration offers the advantage of high acceleration gradients

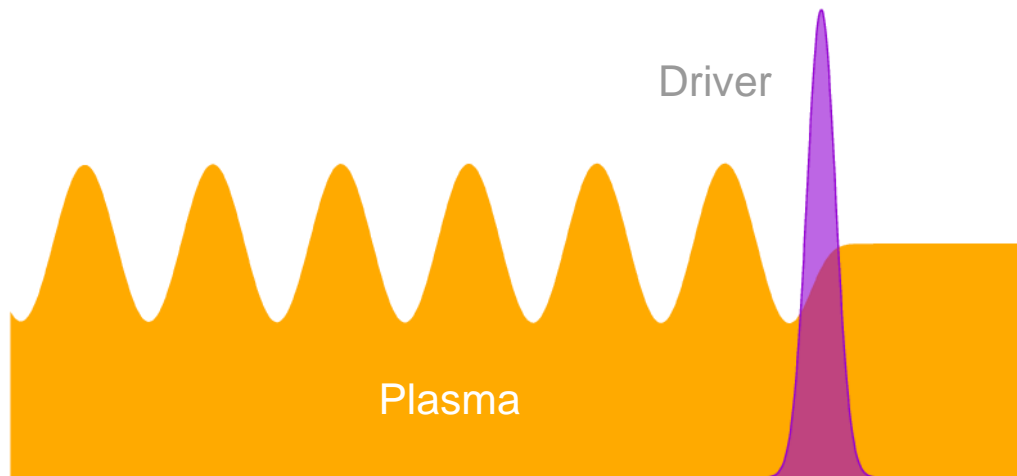
Proton-driven plasma wakefield acceleration would allow high energy gain in a single stage

- potentially higher average gradients
- positron acceleration

Motivation: plasma wakefields

Damage threshold not a concern for plasma

- already “broken”
- gradients $> 1\text{GeV/m}$



Wakeboarding

Plasma wakefield acceleration

Energy gain limited by choice of driver

- laser pulse
- electron bunch
- proton bunch

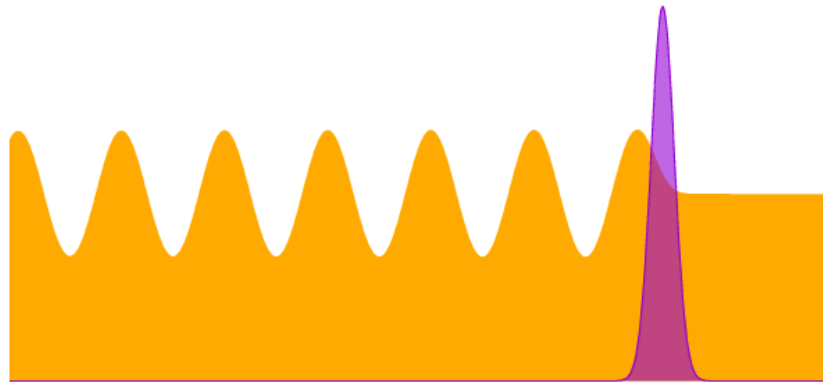
Not enough energy in laser/electron driver to reach high witness energy. Solutions include

- structured driver (stability)
- staging (alignment, average gradient)

Proton-driven PWFA

Proton beams have plenty of energy

BUT available beams “too long” to efficiently drive a wake



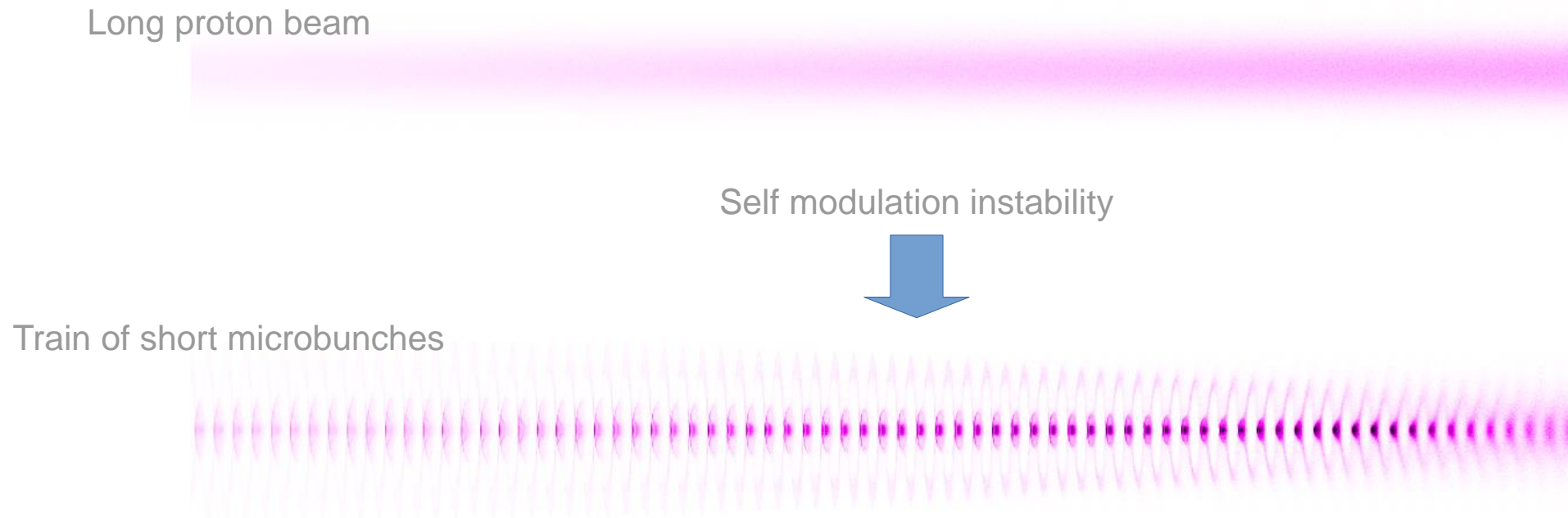
Short driver efficiently excites wakefield



Long driver suppresses its own wake

Proton-driven PWFA

Focussing/defocussing fields in plasma



Resulting train of microbunches can drive large wakefields

Proton-driven PWFA

Focussing/defocussing fields in plasma



Resulting train of microbunches can drive large wakefields

Short proton drivers revisited



It's worth revisiting short proton drivers.

Pros:

Higher gradients

Higher efficiency

Cons:

Such drivers ($L \sim 150 \mu\text{m}$)
don't exist

Short proton drivers revisited

nature
physics

ARTICLES

PUBLISHED ONLINE: 12 APRIL 2009; CORRECTED ONLINE: 24 APRIL 2009 | DOI: 10.1038/NPHYS1248

Proton-driven plasma-wakefield acceleration

Allen Caldwell^{1*}, Konstantin Lotov^{2,3}, Alexander Pukhov⁴ and Frank Simon^{1,5}

[Caldwell et al. \(2009\)](#)

A short proton wakefield driver is not a new idea (2009).
Predates AWAKE! So why now?

Short proton drivers revisited

IEEE TRANSACTIONS ON APPLIED SUPERCONDUCTIVITY, VOL. 32, NO. 6, SEPTEMBER 2022

4100404

Record High Ramping Rates in HTS Based Superconducting Accelerator Magnet

H. Piekarz ^{ID}, *Senior Member, IEEE*, S. Hays, B. Claypool, M. Kufer ^{ID}, and V. Shiltsev, *Fellow, IEEE*

[Piekarz et al. \(2022\)](#)

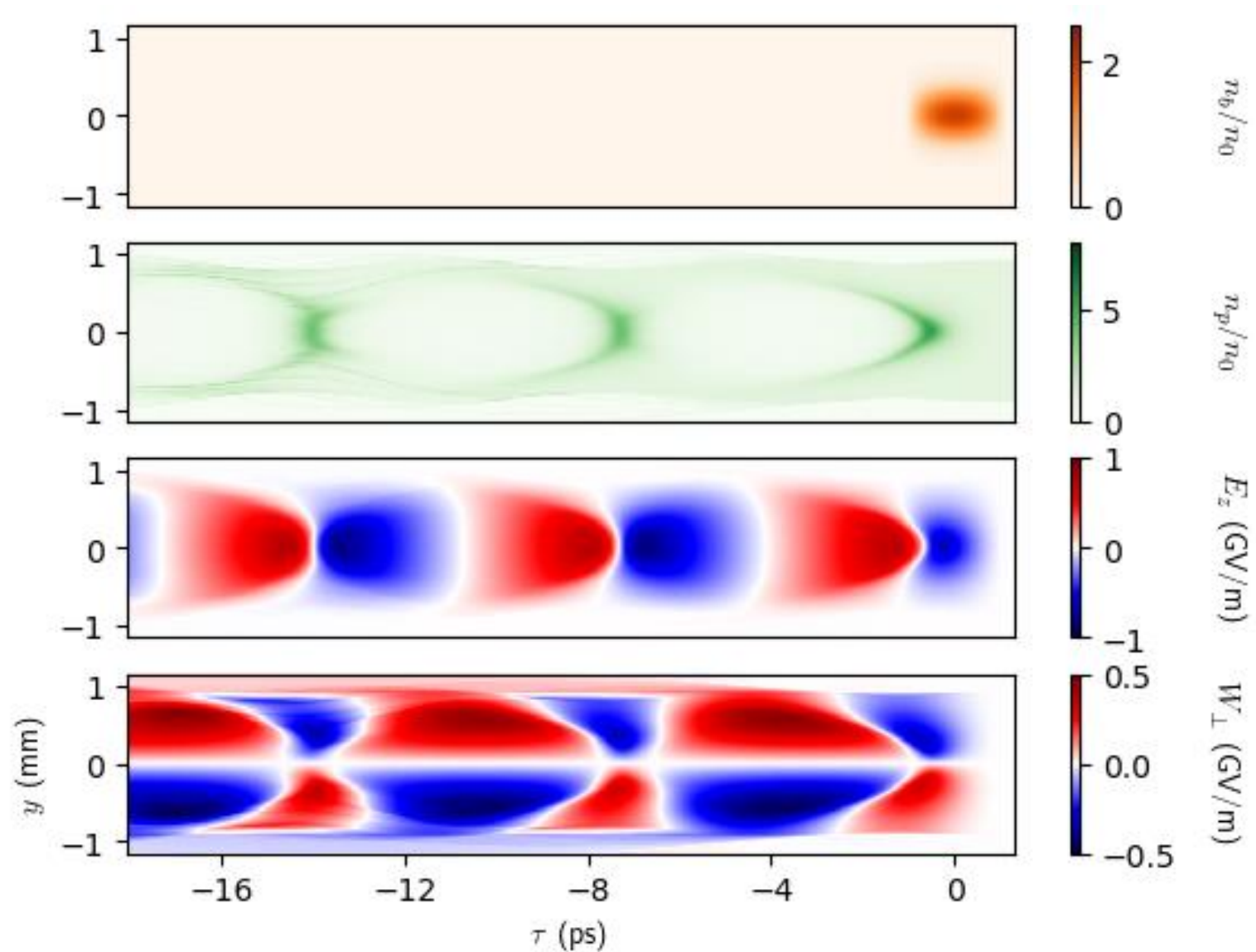
Developments in fast-ramping magnets would allow rapid-cycling (~5 Hz) synchrotrons.

Would allow for competitive luminosities for a proton-driven Higgs factory *if* bunch length can be achieved.

Configuration

Assume a suitably short proton driver can be generated

Moderately nonlinear wakefield allows acceleration of both electrons and positrons



Configuration

Higgs–Z threshold is 216.4 GeV

Add some margin → collision energy of 250 GeV

- 125 GeV e^- colliding with 125 GeV e^+

- HALHF considers 500 GeV e^- colliding with 31.3 GeV e^+

We need to demonstrate

- efficiency

- stability

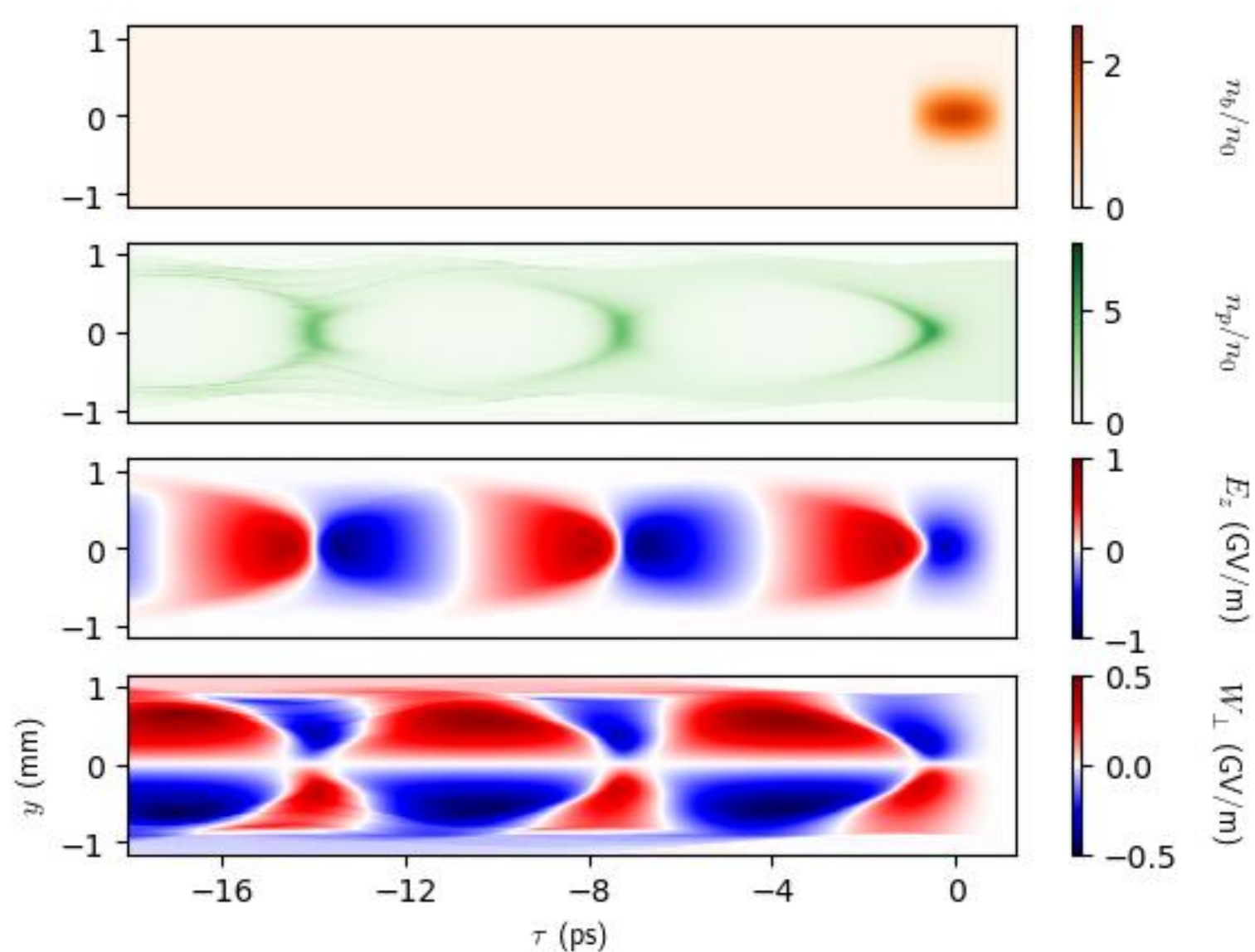
Picking the driver

Assume a suitably short proton driver can be generated

Need to pick driver parameters

- High efficiency

- Moderately nonlinear wake for positron acceleration

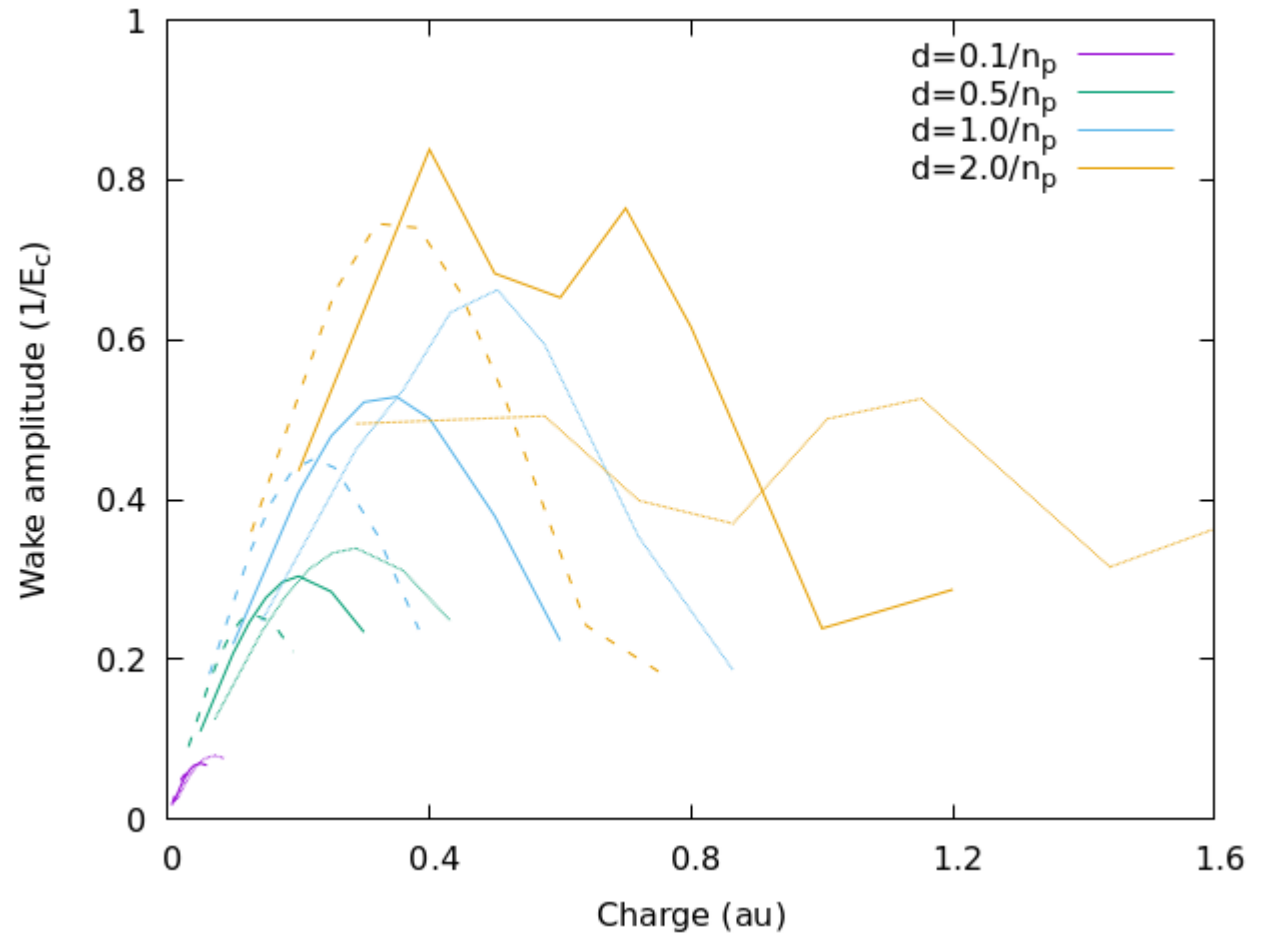


Picking the driver: efficiency

Scans with different driver length, density and radius

Optimal driver length
~driver charge density

Too-high current leads to highly-nonlinear wake



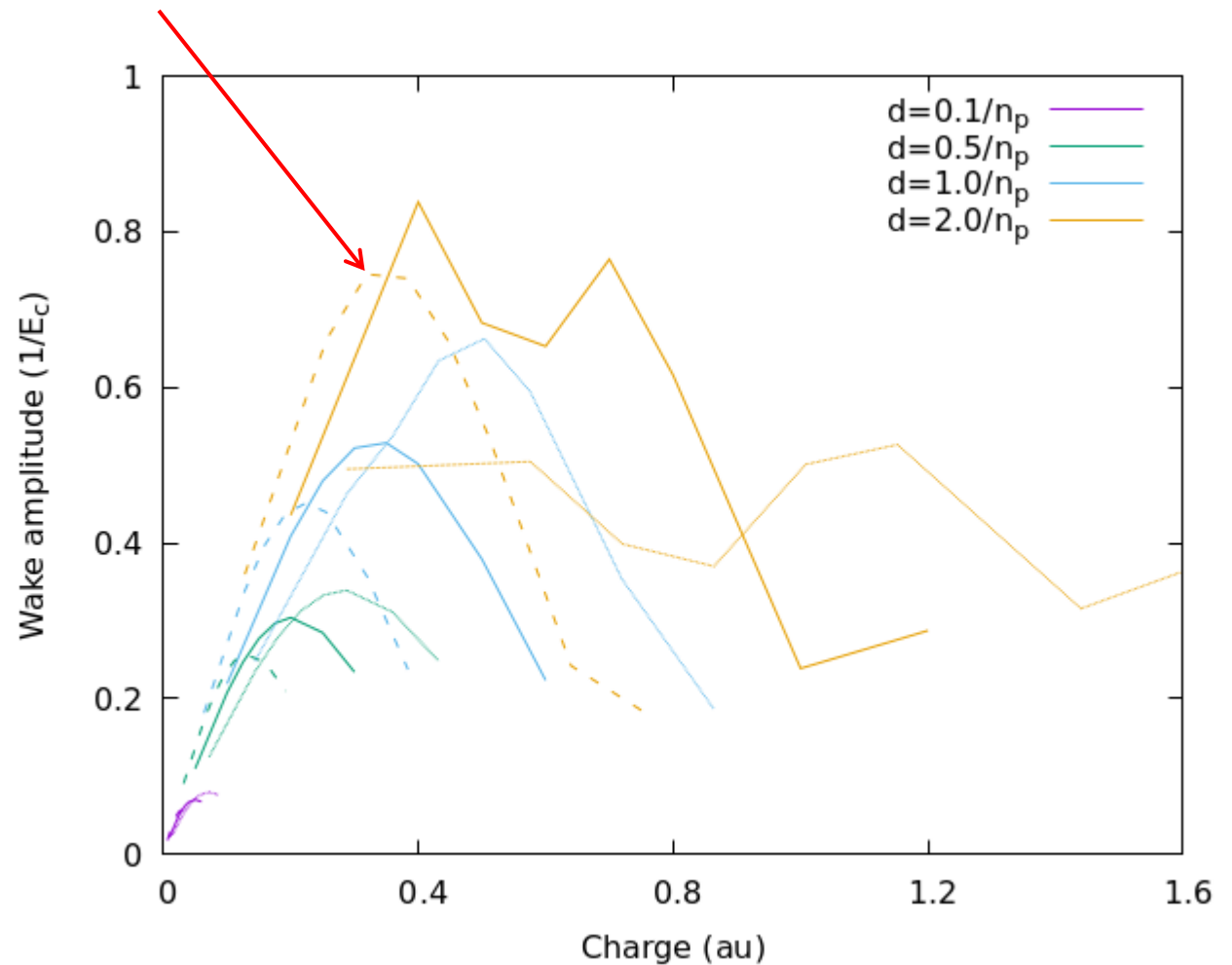
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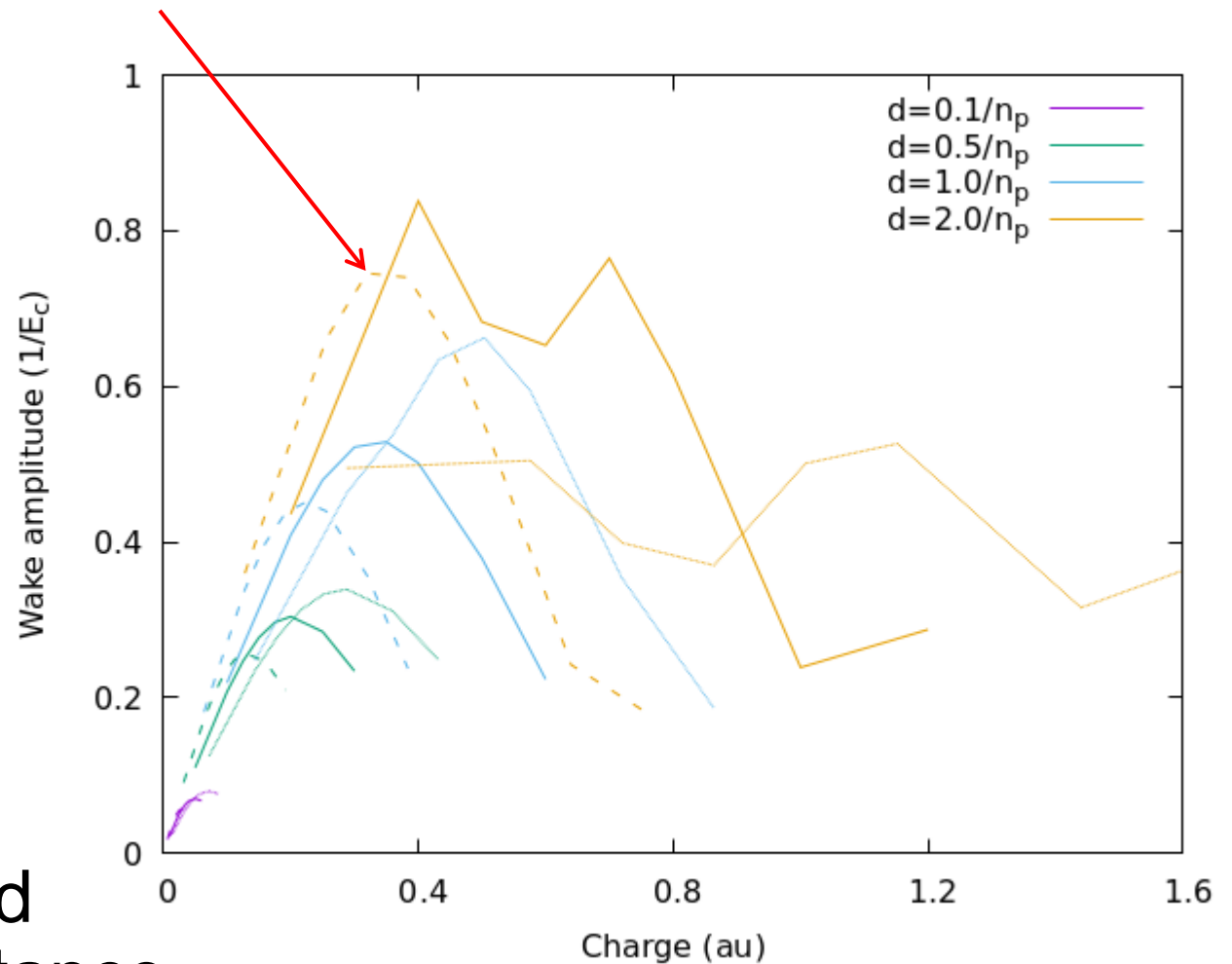
Picking the driver: efficiency

Everything scales with plasma frequency

1×10^{11} protons gives

- plasma density $3 \times 10^{14} \text{ cm}^{-3}$
- driver length $150 \mu\text{m}$
- Initial wakefields $\sim 0.8 \text{ GV/m}$

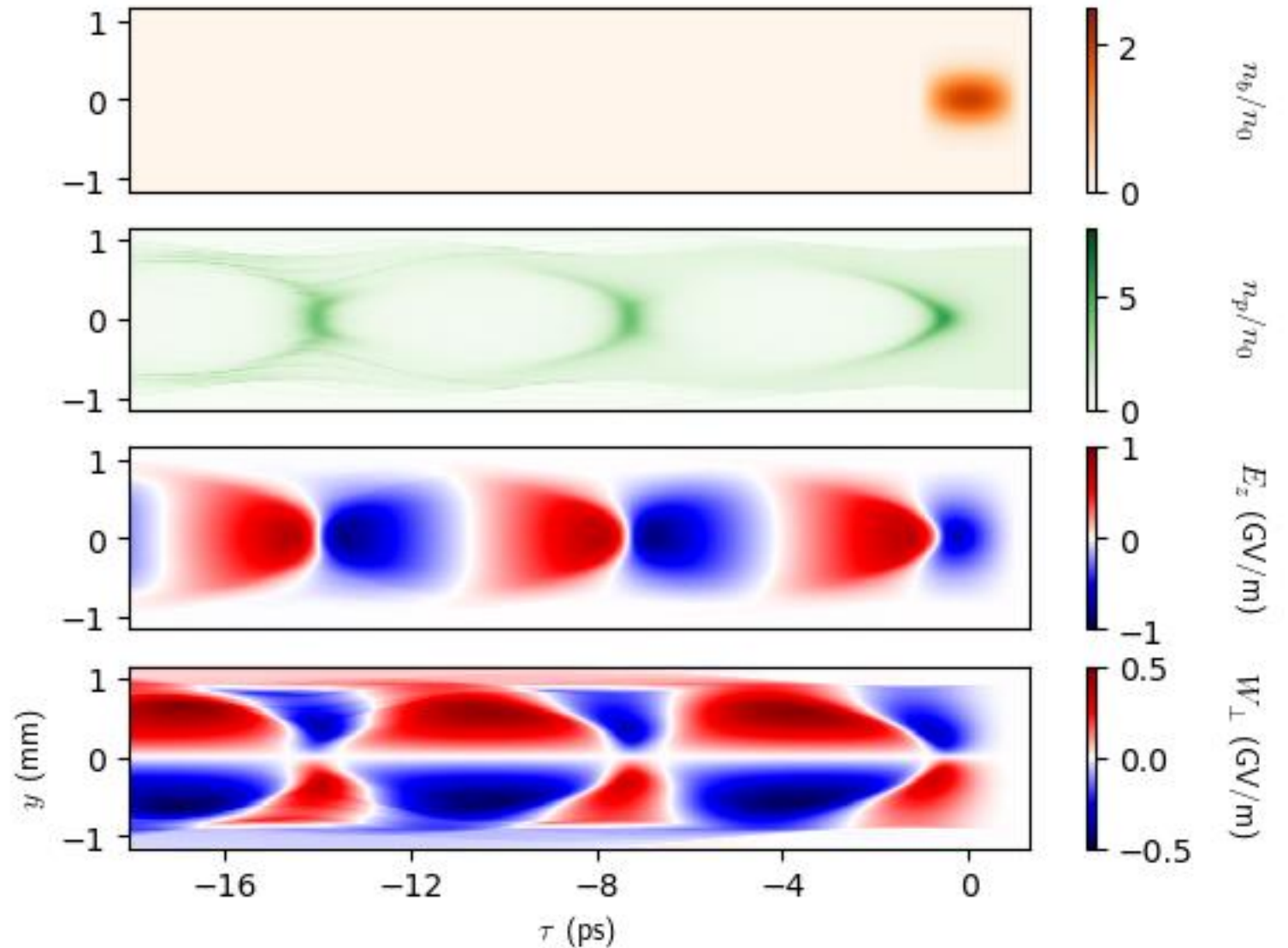
Pick 10% driver energy spread for “realistic” longitudinal emittance



Dashed line: $k_{pr} = 0.8$
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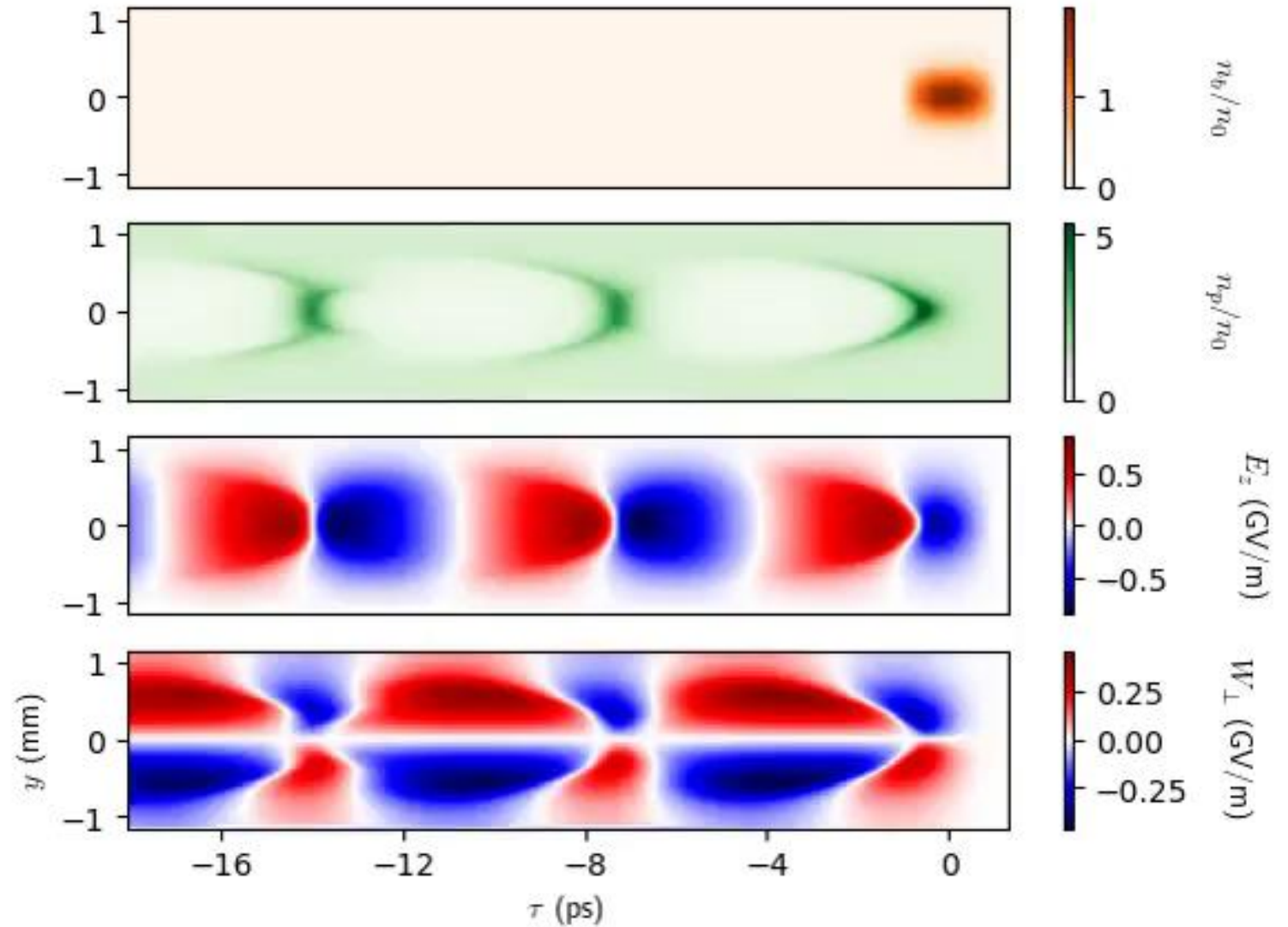
Picking the driver: stability

Initial proton driver
chosen to generate
suitable wakefields



Picking the driver: stability

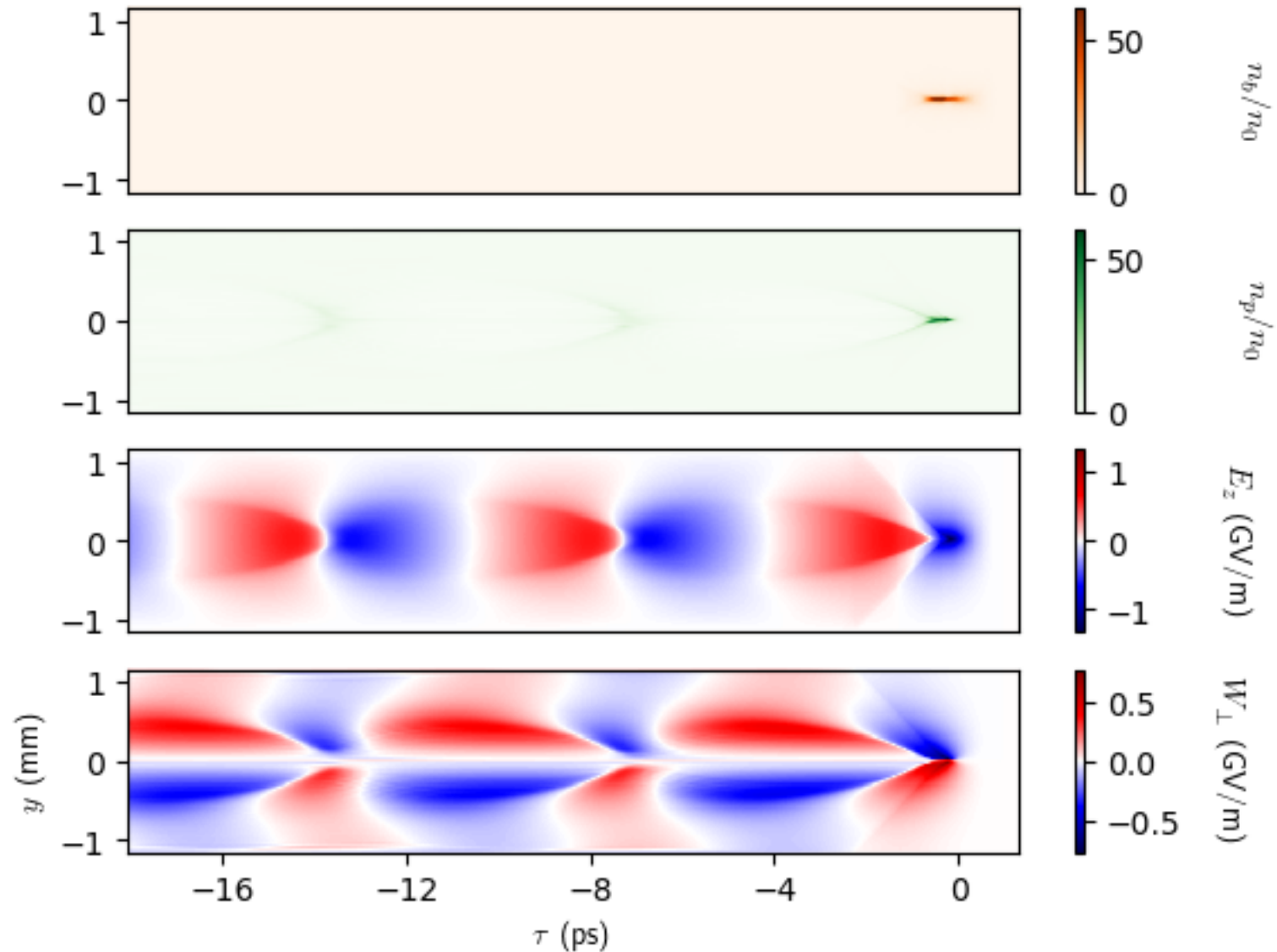
Initial proton driver
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Picking the driver: stability

Initial proton driver
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Driver rapidly pinches



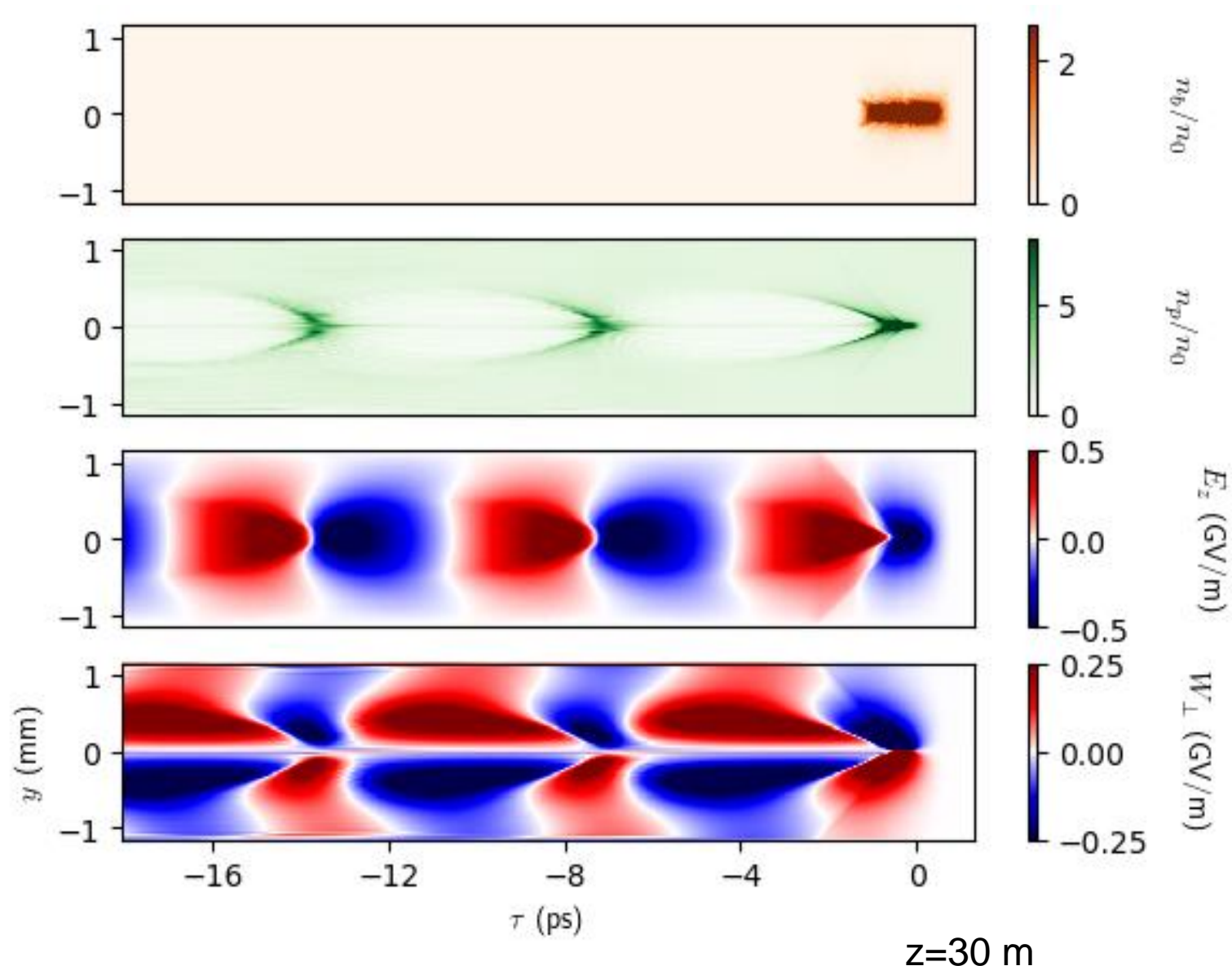
$z=30$ m

Picking the driver: stability

Initial proton driver
chosen to generate
suitable wakefields

Driver rapidly pinches

• Highly nonlinear
wakefield not suitable
for positron
acceleration



Picking the driver: stability

Good initial wakefields not sufficient:

- driver needs to evolve slowly
- counteract strong focussing wakefields

Picking the driver: stability

$$\sigma_z = 150 \mu\text{m}$$

$$\sigma_r = 240 \mu\text{m}$$

$$n_b = 1 \times 10^{11}$$

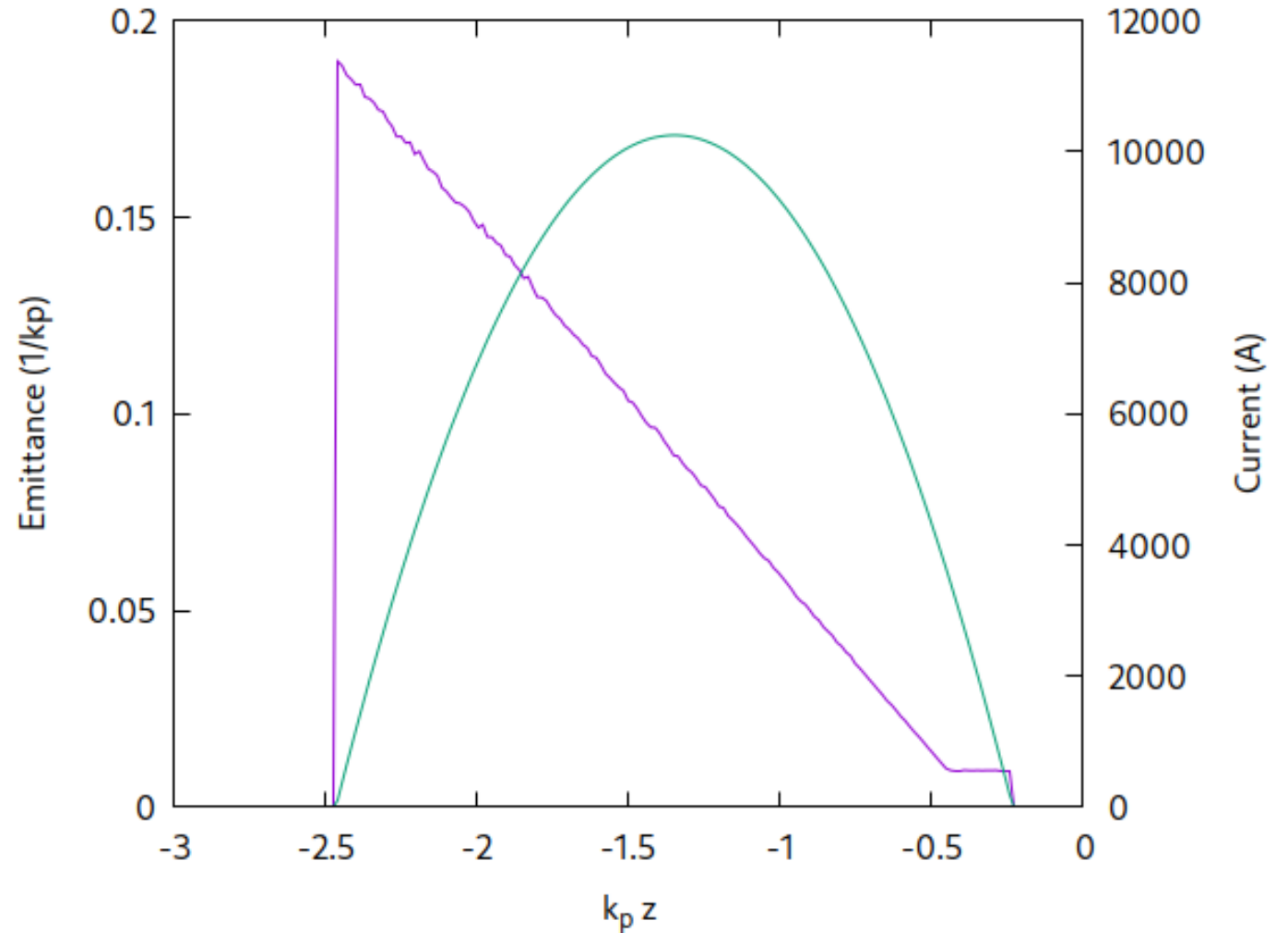
$$E = 400 \text{ GeV}$$

$$\epsilon_N = \textit{tailored}$$

- 3 μm at head

- initially constant

- rises linearly to 75 μm

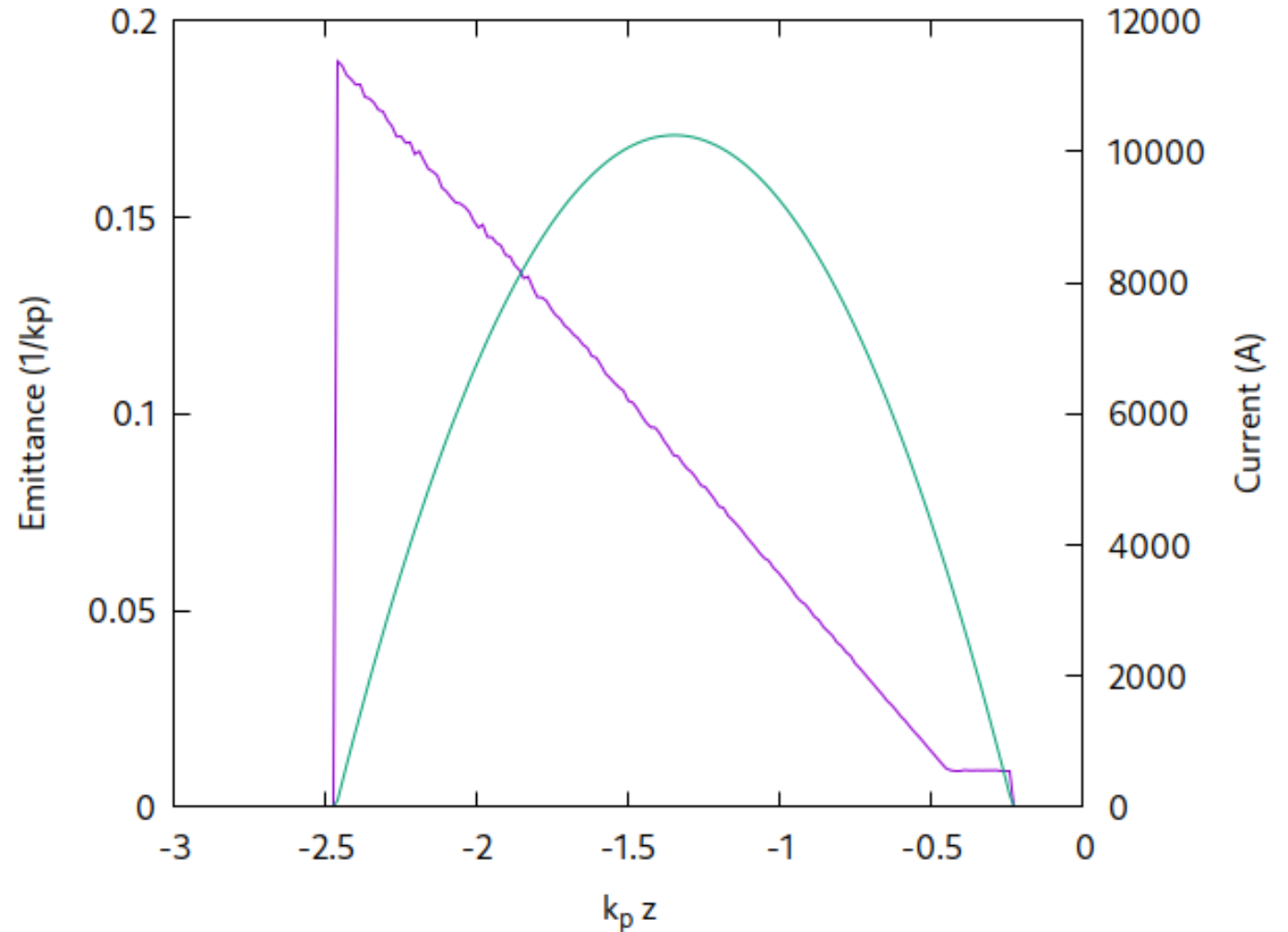


Picking the driver: stability

How can we generate a tailored emittance profile?

Most likely:
with difficulty

BUT emittance is initially constant before growing monotonically

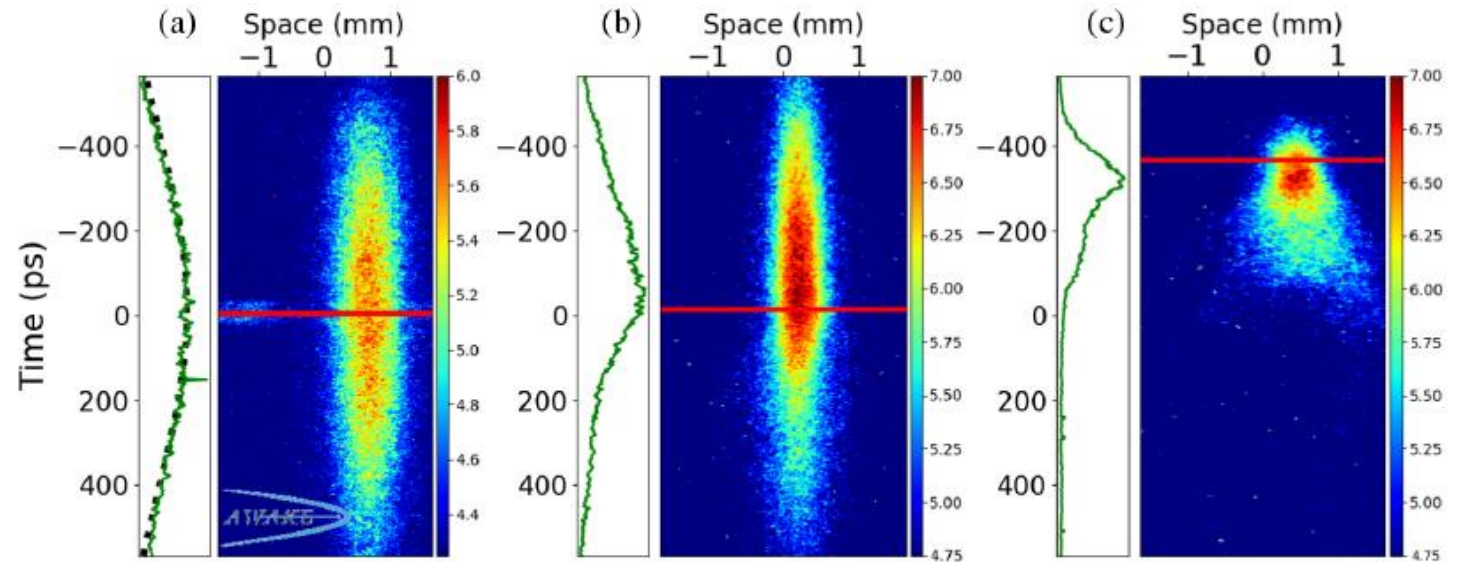


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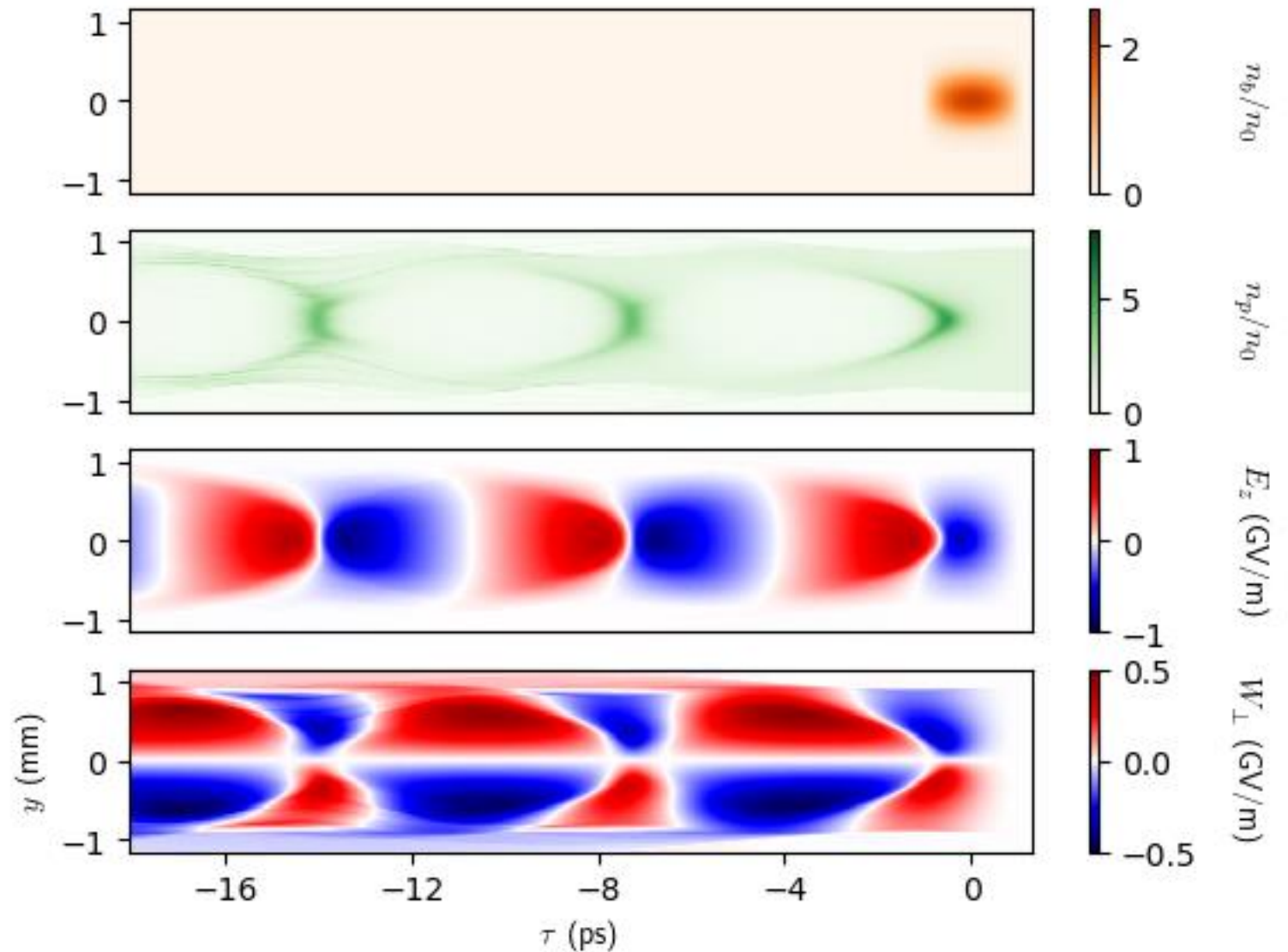


[AWAKE Collaboration, PRL \(2019\)](#)

Harness plasma instabilities?

Acceleration: dephasing

Initial proton driver
chosen to generate
suitable wakefields

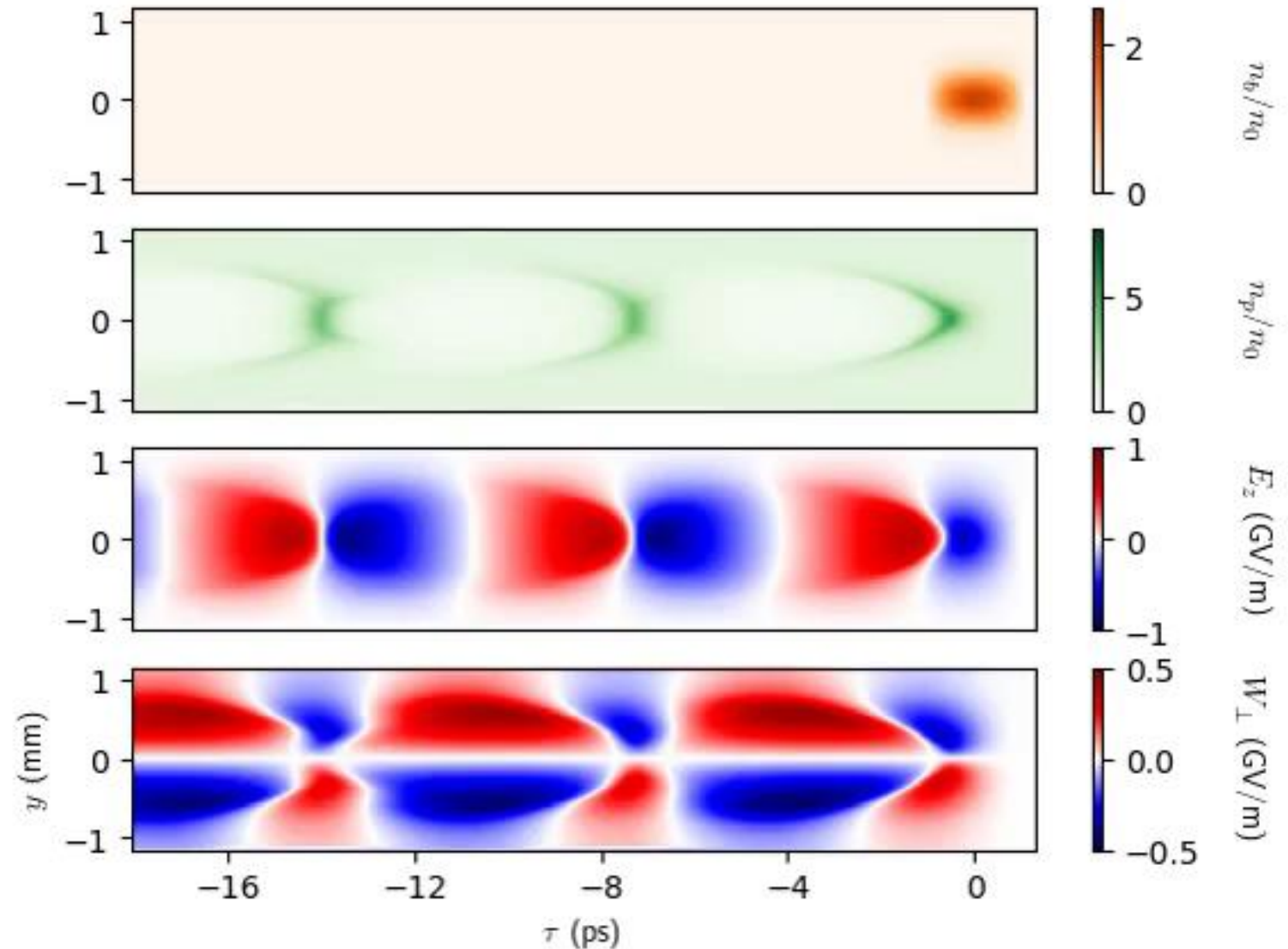


Acceleration: dephasing

Initial proton driver chosen to generate suitable wakefields

Tailored emittance profile stops the bunch from pinching

BUT:
protons “fall back”
in the light frame

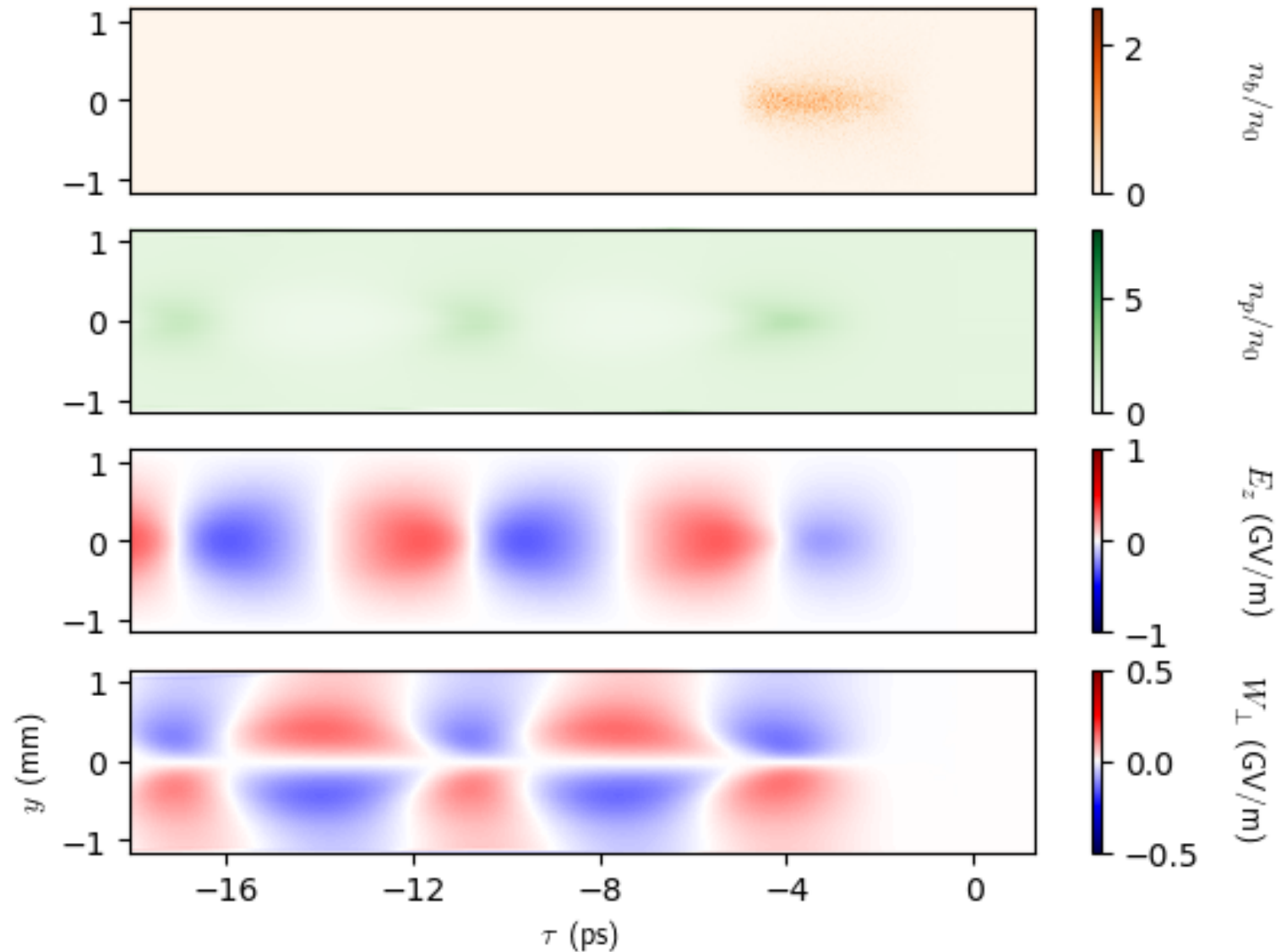


Acceleration: dephasing

Initial proton driver
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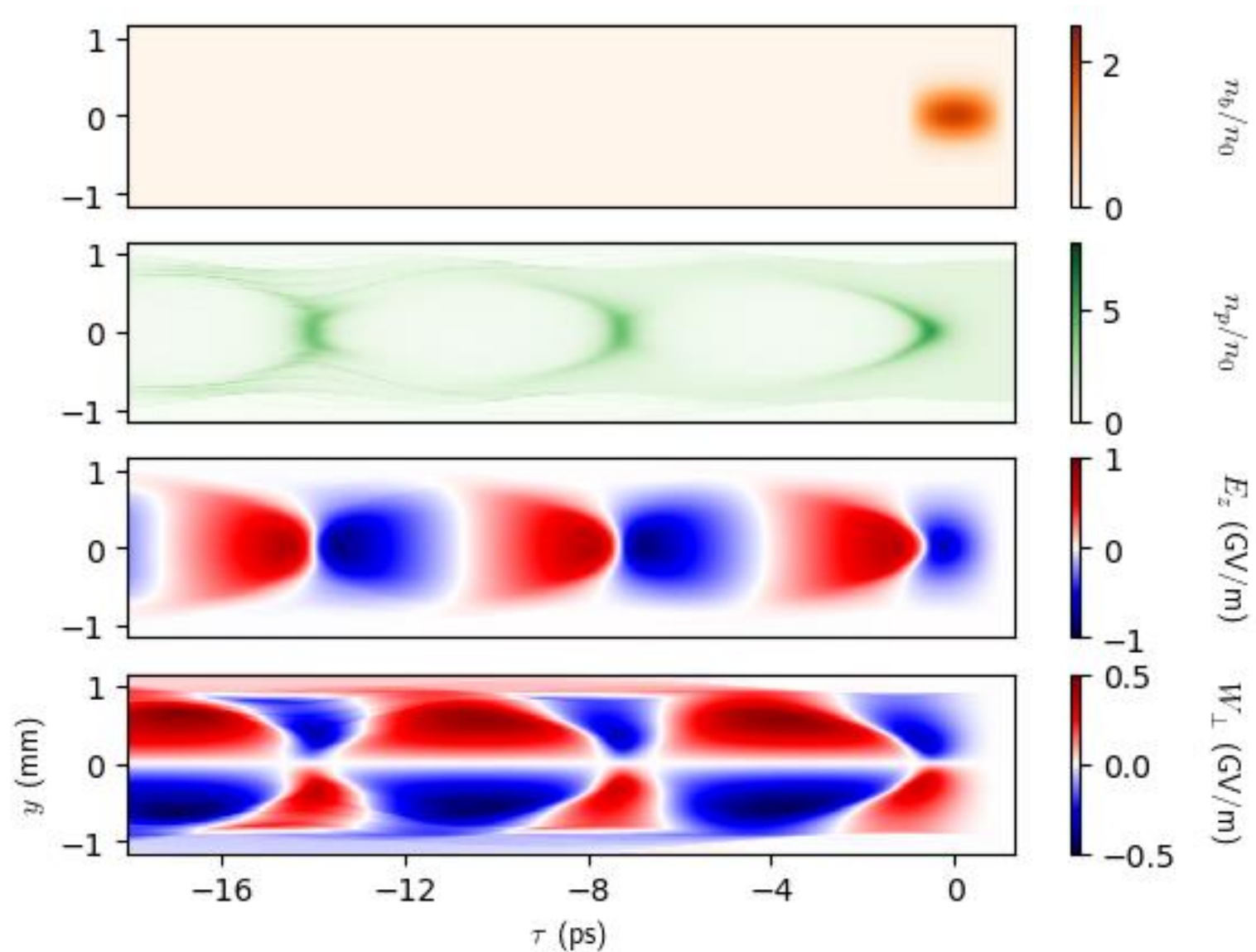
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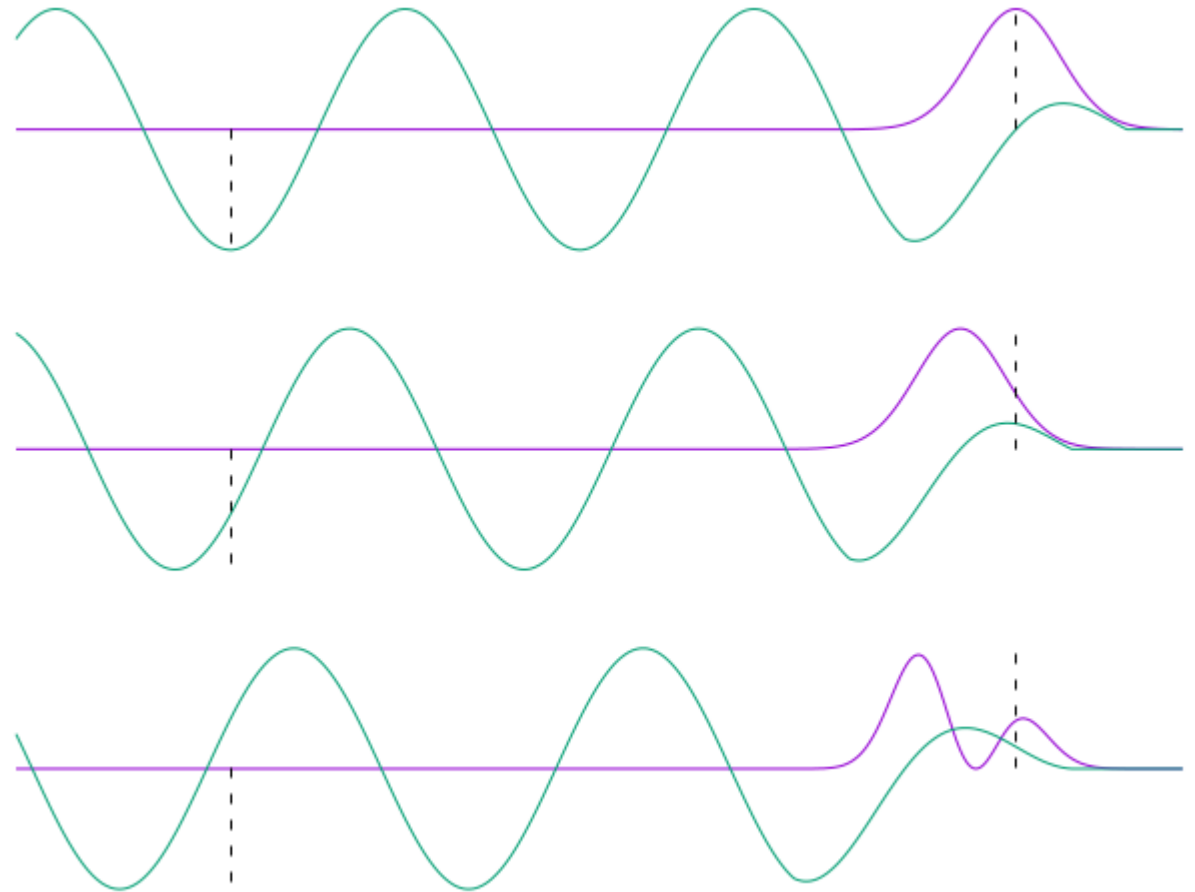


Acceleration: dephasing

Protons are fast,
but not that fast.

Driver evolution
will also modify
wakefield phase.

Witness will “catch up”
with the driver.

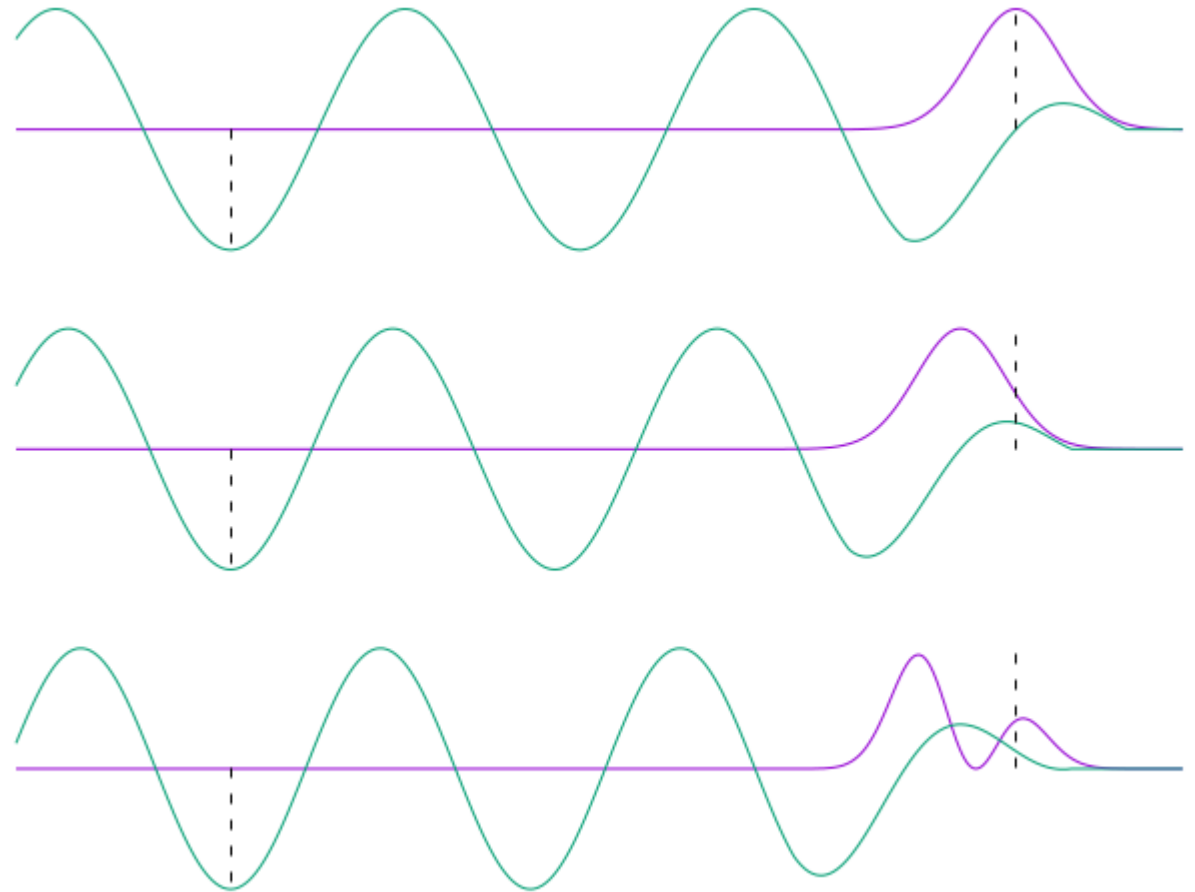


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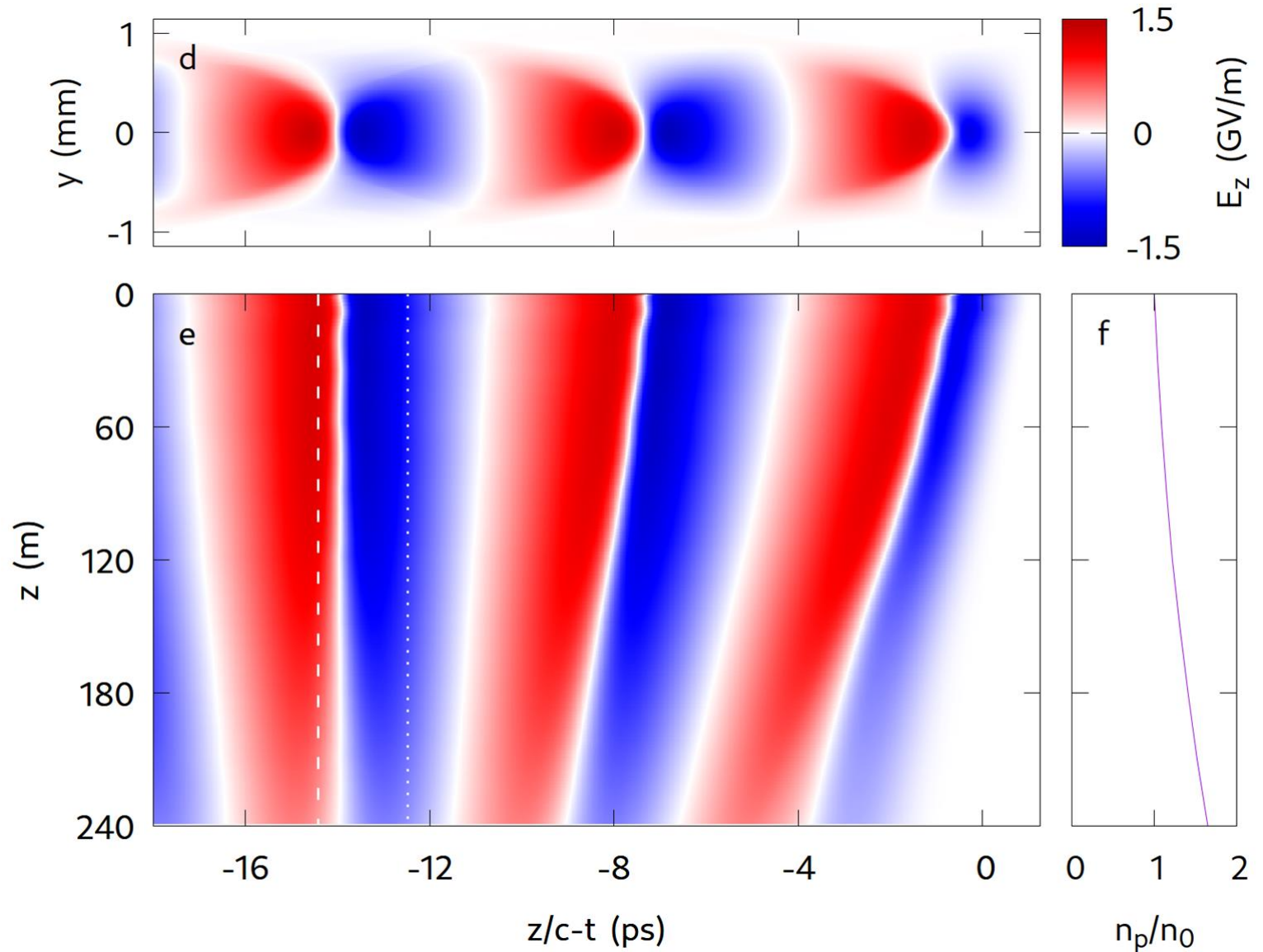
Witness will “catch up”
with the driver.



Change plasma density to keep phase constant

Acceleration: dephasing

Change plasma density to keep phase constant



Acceleration

We now have all the building blocks for Higgs factory

- Large accelerating wakefields
- Regions suitable for electron and positron acceleration
- Stable accelerating phase

Just (!) need to simulate acceleration

Acceleration

Large focussing fields mean extremely tiny witness radius at high energy

- ~1 μm for electrons

- ~100 nm for positrons

Simulations don't resolve this

- Nonphysical emittance growth

- Nonphysical energy spread growth

However, simulations do provide proof-of-concept for energy gain

Acceleration

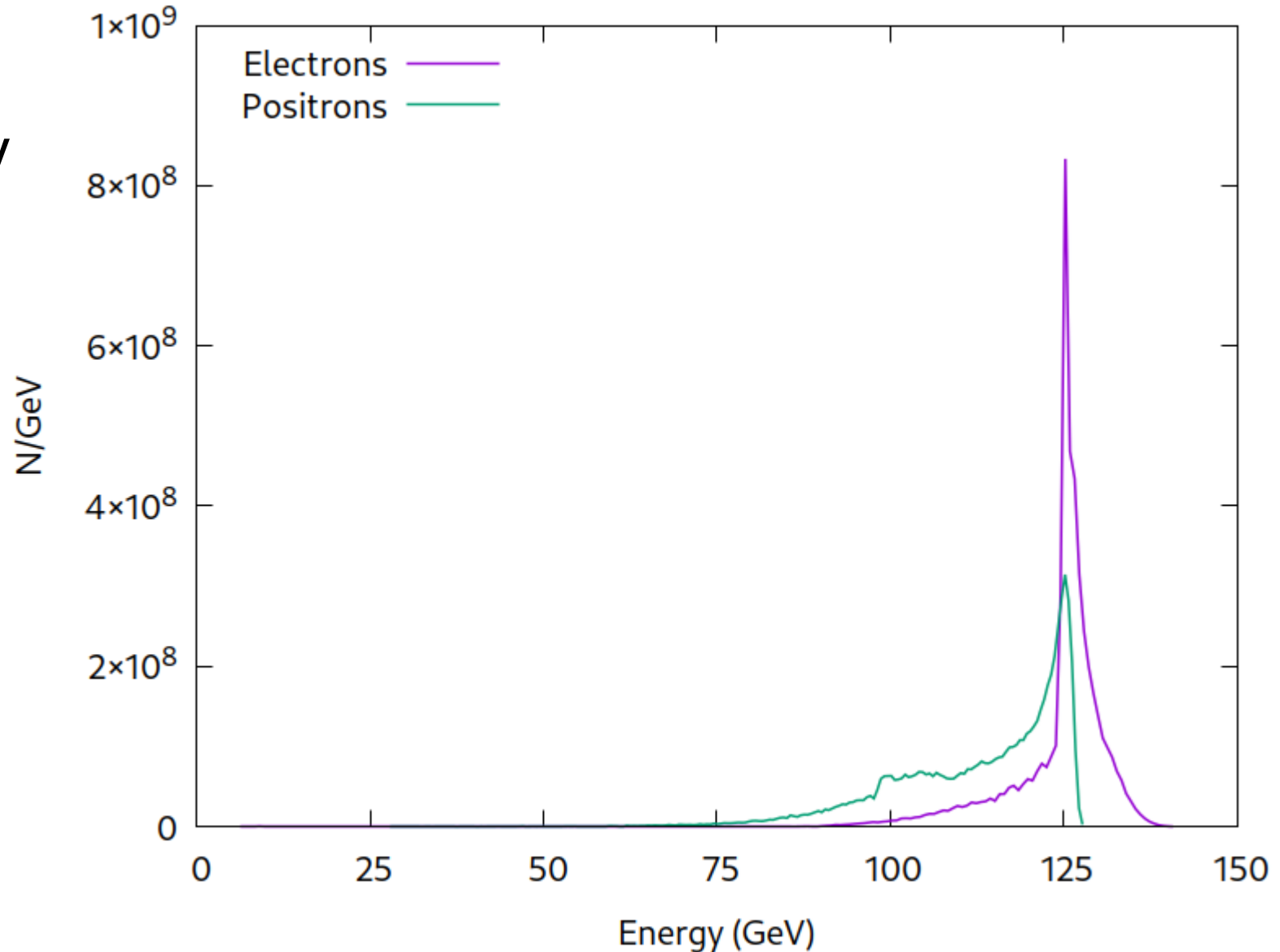
Driver:

10^{11} protons at 400 GeV

Witness:

10^{10} electrons/positrons
injected at 1 GeV

Energy spread is
nonphysical!



Acceleration

Simulations show proof-of-concept for acceleration

.10% of driver charge accelerated to 125 GeV in ~180 m

With tailored witness current profile, 20% should be possible

S. van der Meer, CLIC Note (1985)

Requires long plasma stage:

.CERN is already investing in this technology

[N. E. Torrado et al \(2023\)](#)

[B Buttenschön, N Fahrenkamp and O Grulke \(2019\)](#)

Luminosity

Simulations don't resolve emittance or energy spread, so we assume that these can be controlled.

For electrons: AWAKE Run 2c

[Olsen et al. \(2019\)](#), [Farmer et al \(arXiv\)](#)

For positrons: simulation study by Hue *et al.*

[C. S. Hue et al. \(2021\)](#)

Luminosity

Combine everything:

- Assume proton beams at 5 Hz, with 1000 bunches per beam

- Assume witness beams with 20% driver charge, 100 nm emittance*, ILC optics, and negligible energy spread

*Flat beams should be investigated

and this scheme is competitive:

$$1.7 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$$

Proton Accelerator Parameter	Symbol	Unit	Value
Proton energy	E_p	GeV	400
Refill Time	τ	s	0.2
Bunch population	N_p	10^{10}	10
Number of bunches	n		1000
Longitudinal RMS	σ_z	μm	150
Transverse RMS	$\sigma_{x,y}$	μm	240
Normalized transverse emittance	$\epsilon_{T,p}$	μm	3 – 75 μm
Power Usage	P	MW	150
Plasma Parameters	Symbol	Unit	Value
e^- cell Length	L_{e^-}	m	240
e^+ cell Length	L_{e^+}	m	240
density - upstream	n_p	10^{14} cm^{-3}	3.2
density - downstream	n_p	10^{14} cm^{-3}	5.2
e^\pm Bunch Parameters	Symbol	Unit	Value
Injection Energy	$E_{e,in}$	GeV	1
Final Energy	E_e	GeV	125
Bunch population	N_{e^\pm}	10^{10}	2
Normalized transverse emittance	$\epsilon_{T,e}$	nm	100
Hor. beta fn.	β_x^*	mm	13
Ver. beta fn.	β_y^*	mm	0.41
Hor. IP size.	σ_x^*	nm	73
Ver. IP size.	σ_y^*	nm	13
$e^- e^+$ Collider Parameter	Symbol	Unit	Value
Center-of-Mass Energy	E_{cm}	GeV	250
Average Collision Rate	f	kHz	5
Luminosity	\mathcal{L}	$\text{cm}^{-2} \text{ s}^{-1}$	1.7×10^{34}

Luminosity

Combine everything:

• As ... at 5 Hz,

• with ...
 • based on Proton-Driven Plasma Wakefield Acceleration
 • J. Farmer, A. Caldwell, and A. Pukhov
[Now on arXiv](#)

and neg ...

*Flat beams should be investigated

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Proton Accelerator Parameter	Symbol	Unit	Value
Proton energy	E_p	GeV	400
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Power Usage	P	MW	150

Plasma Parameters	Symbol	Unit	Value
e^- cell Length	L_{e^-}	m	240
cell Length	L_{e^+}	m	240
	n_p	10^{14} cm^{-3}	3.2
		10^{14} cm^{-3}	5.2

	Unit	Value
	GeV	1
	GeV	125
	10^{10}	2
	nm	100
	mm	13
	mm	0.41
	nm	73
	nm	13

	Symbol	Unit	Value
Center-of-Mass Energy	E_{cm}	GeV	250
Average Collision Rate	f	kHz	5
Luminosity	\mathcal{L}	$\text{cm}^{-2} \text{ s}^{-1}$	1.7×10^{34}

Outlook

Key challenges:

- rapid-cycling proton synchrotron with short bunches
- demonstration of drive bunches with tailored emittance
- experimental demonstration of emittance control for PWFA acceleration of positron bunches
- acceleration of flat witness bunches
- long plasma stages at high rep rate

~ **Fin** ~

~ Backups ~

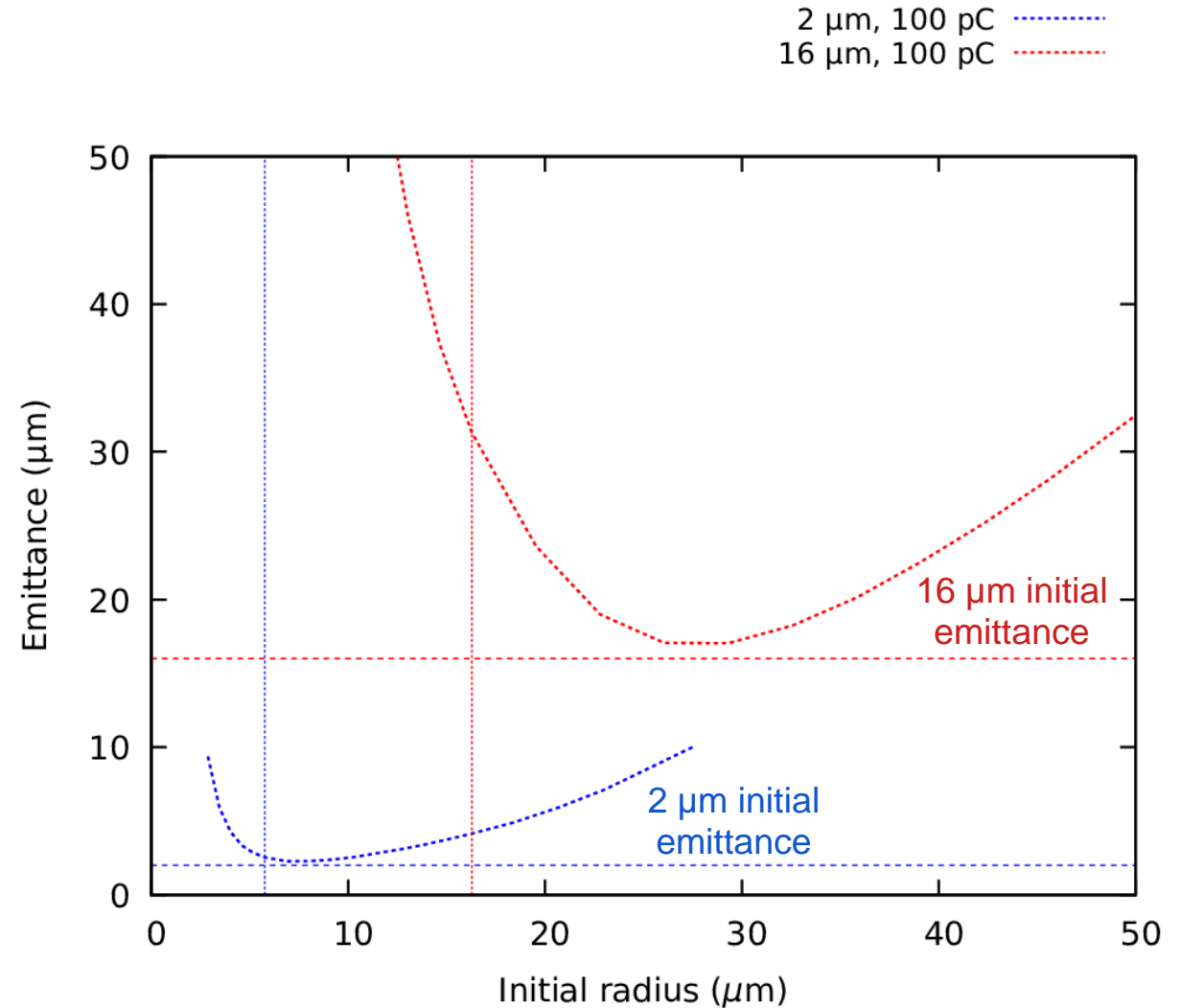
Footprint

Fits on the Fermilab site
(P5 review)

Emittance

Simulations for AWAKE Run 2c
show that a blowout is *not*
required to have emittance
control.

Obviously a big gap between
this and positron beams with
sub- μm emittance!



Adapted from [Farmer et al \(arXiv\)](#)

Extension to HALHF

Acceleration limited by dispersion, not energy depletion:
scales as $1/\gamma^2$.

Use a 800 GeV driver, accelerate to ~500 GeV in ~1 km.

Cooling

Witness with 10% driver charge
absorbs ~20% of wakefield energy

Witness with 20% driver charge
absorbs ~40% of wakefield energy

Assume acceleration over 240m,
gives required cooling as 12.5 kW/m

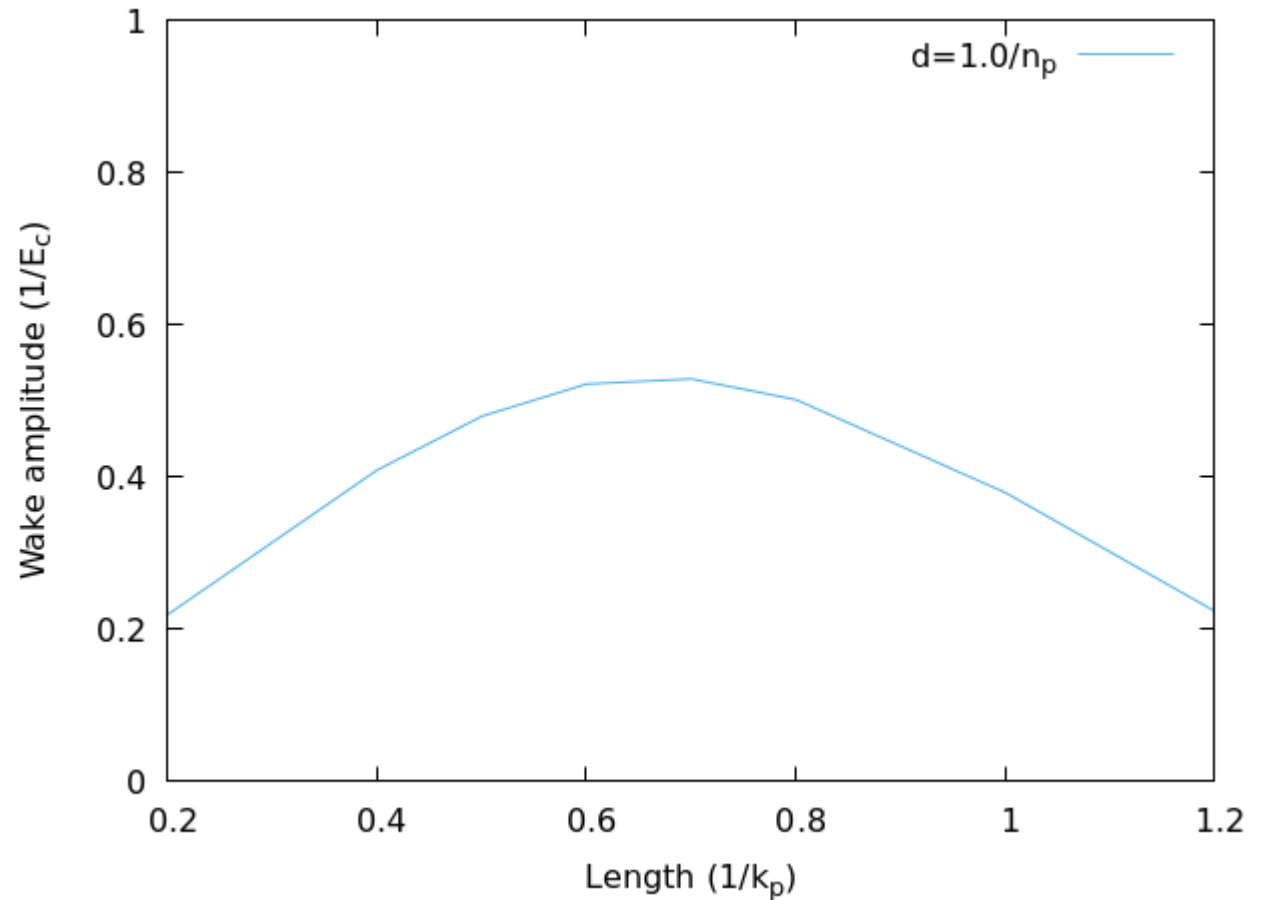
Cooling

Moderately nonlinear wakefields retain their structure after loading.

Could use a second witness bunch to “mop up” excess wakefield

Picking the driver: efficiency

Scan driver length.



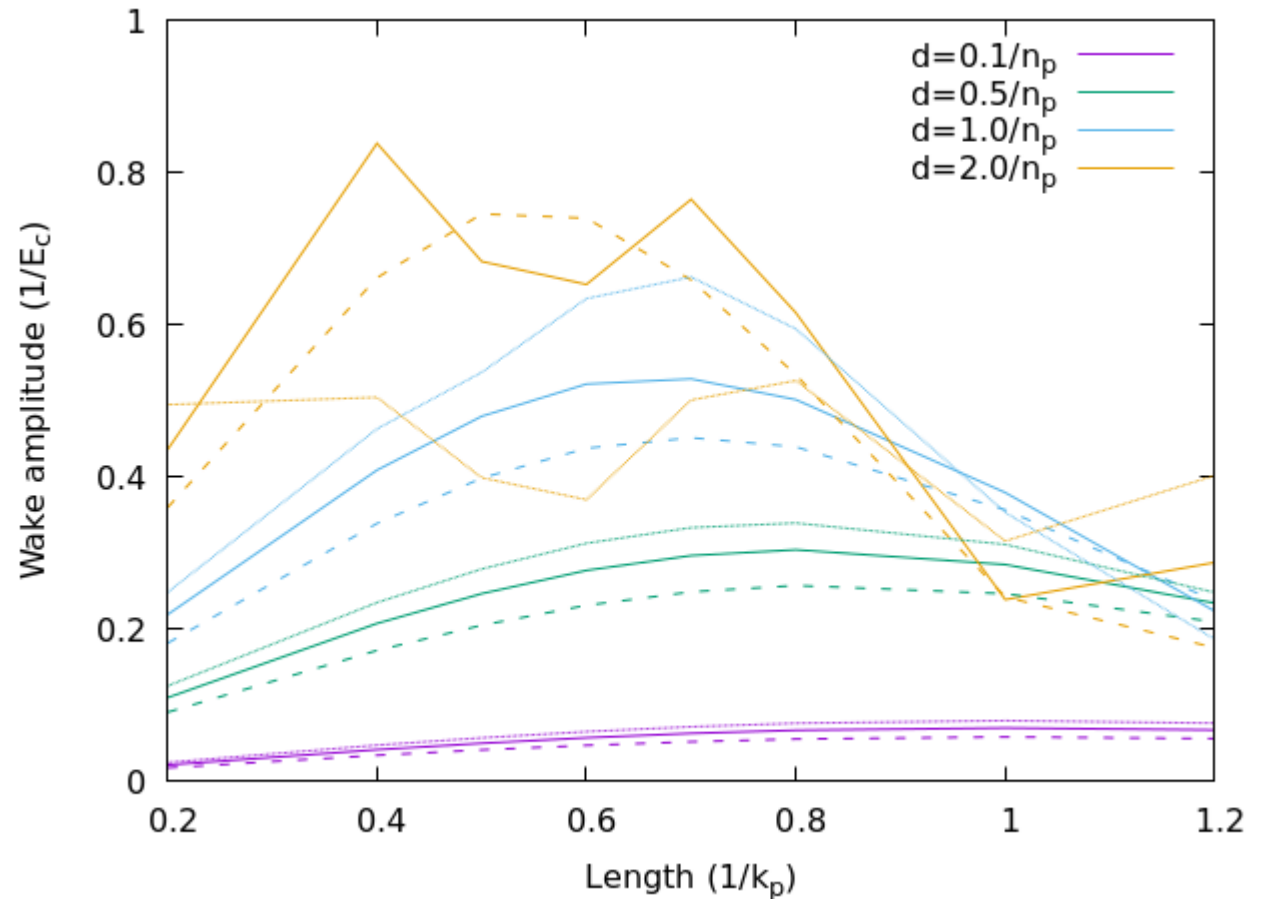
Solid line: $k_p r = 1.0$

Picking the driver: efficiency

Scan driver length.

Optimal length depends on beam charge density.

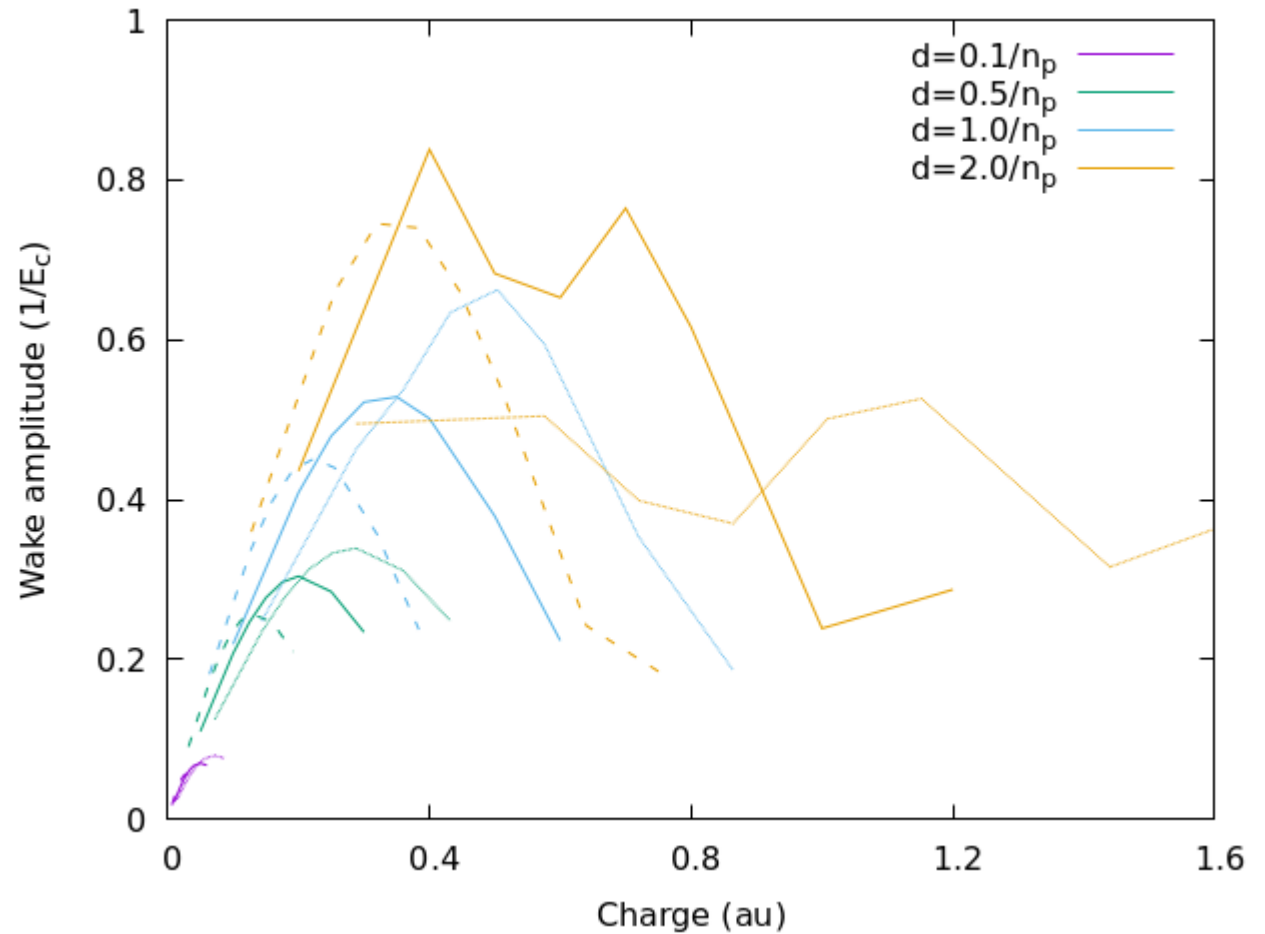
High current leads to nonlinear plasma response.



Dashed line: $k_p = 0.8$
Solid line: $k_p = 1.0$
Dotted line: $k_p = 1.2$

Picking the driver: efficiency

Replot wakefield as a function of beam charge.

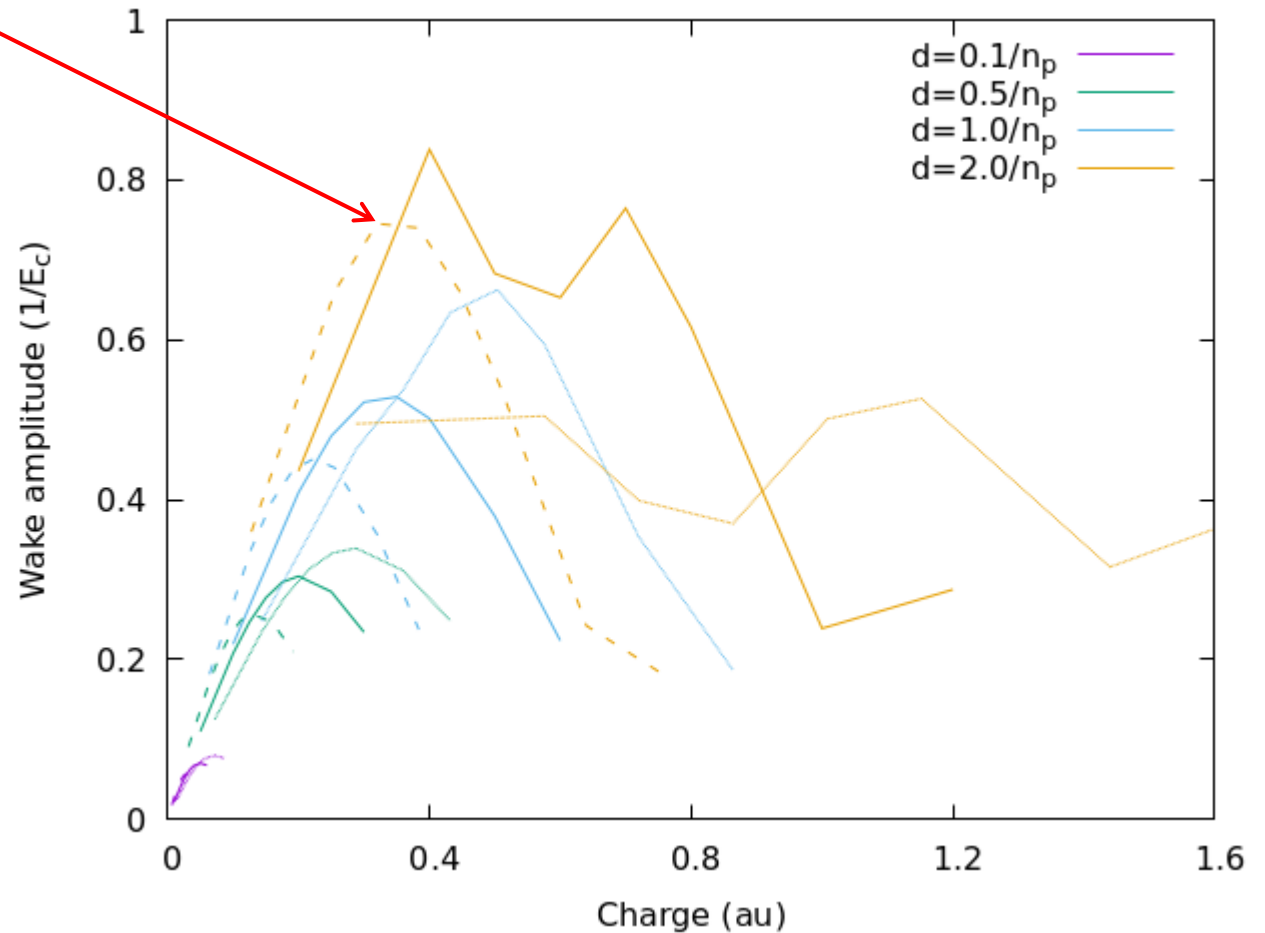


Dashed line: $k_{pr} = 0.8$
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Picking the driver: efficiency

Replot wakefield as a function of beam charge.

Optimal is best gradient per unit charge.



Dashed line: $k_{pr} = 0.8$
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