

Cross sections for cosmic rays

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"PAW'24 - Physics at Amber Int. Workshop "

CRs in the Galaxy

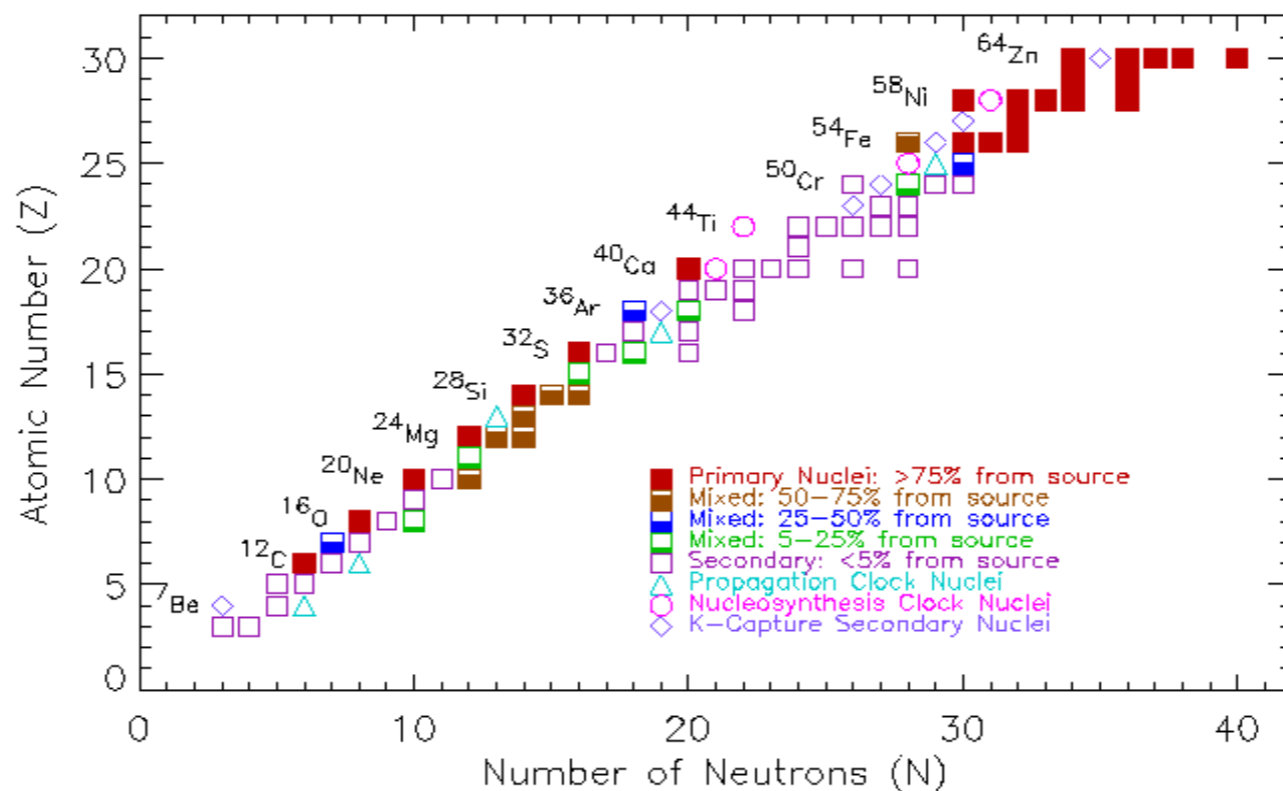
Primaries: produced in the sources (SNR and Pulsars)

H, He, CNO, Fe; e^- , e^+ ; possibly e^+ , p^- , d^- from Dark Matter annihilation

Secondaries: produced by spallation of primary CRs (p, He, C, O, Fe) on the interstellar medium (ISM): Li, Be, B, sub-Fe, [...], (radioactive) isotopes ; e^+ , p^- , d^-

Primaries = present in sources:
 Nuclei: H, He, CNO, Fe; e^- , (e^+) in SNR (& pulsars)
 e^+ , p^+ , d^+ from Dark Matter annihilation

Secondaries = NOT present in sources, thus produced by
 spallation of primary CRs (p, He, C, O, Fe) on ISM
 Nuclei: LiBeB, sub-Fe, ... ;
 e^+ , p^+ , d^+ ; ... from inelastic scatterings



Propagation equation

$$\frac{\partial \psi}{\partial t} - \nabla \cdot \{D(E) \nabla \psi\} + \frac{\partial}{\partial E} \left\{ \frac{dE}{dt} \psi \right\} = Q(E, \mathbf{x}, t)$$

diffusion

en. losses

source spectrum

Diffusion: $D(x, R)$ a priori

usually assumed isotropic in the Galaxy: $D \sim D_0 R^\delta$

D_0 and δ usually fixed by B/C

Energy losses: Nuclei: ionisation, Coulomb

Leptons: Synchrotron on the galactic $B \sim 3.6 \mu\text{G}$

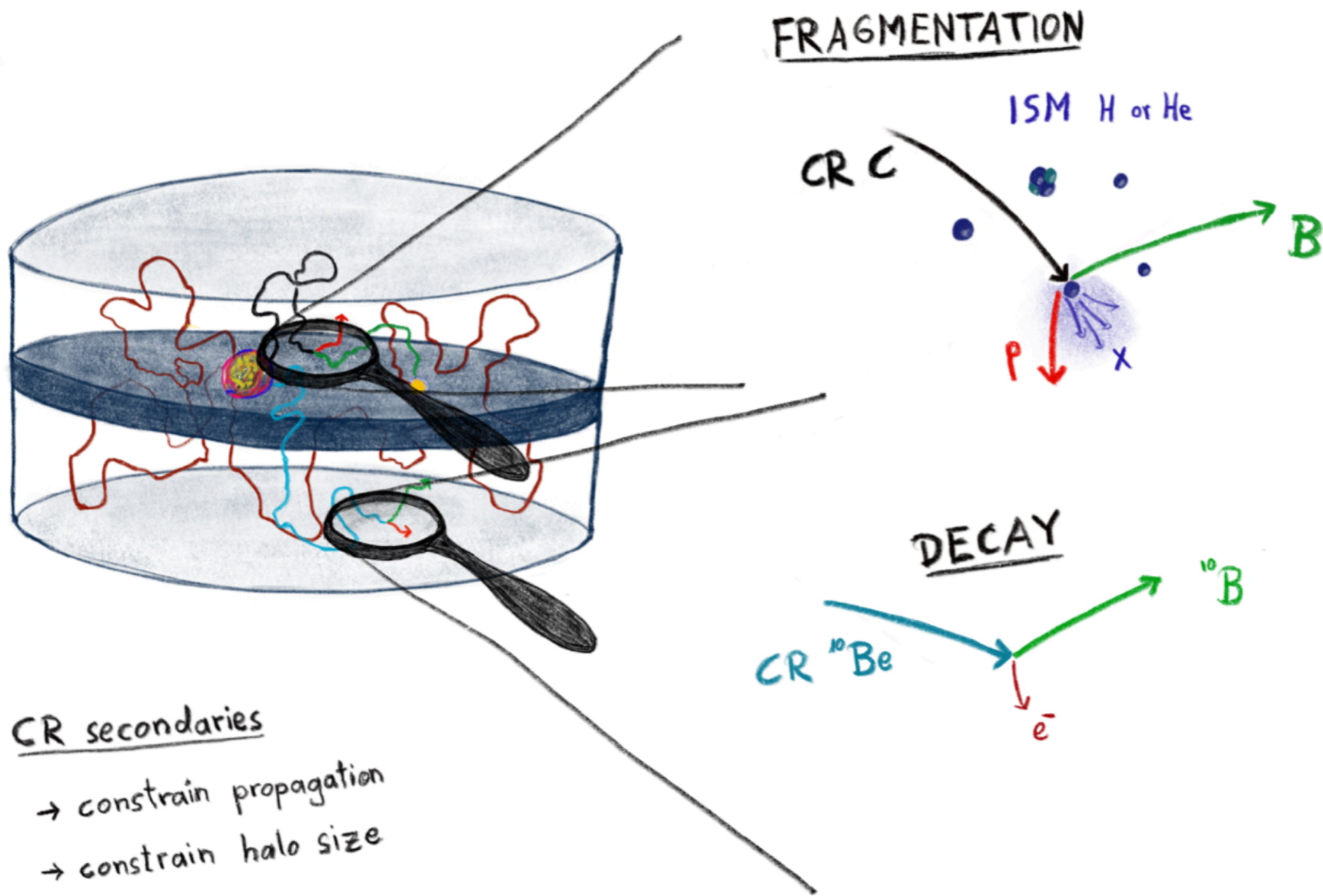
Inverse Compton on photon fields (stellar, CMB, UV, IR)

Sources: Supernova Remnants, $Q(E) \propto E^{-\gamma}$

Nuclear fragmentation, $Q_j(E) \propto n_{\text{ISM}} \sigma_{ij} \psi_i$

Dark Matter annihilation, see below

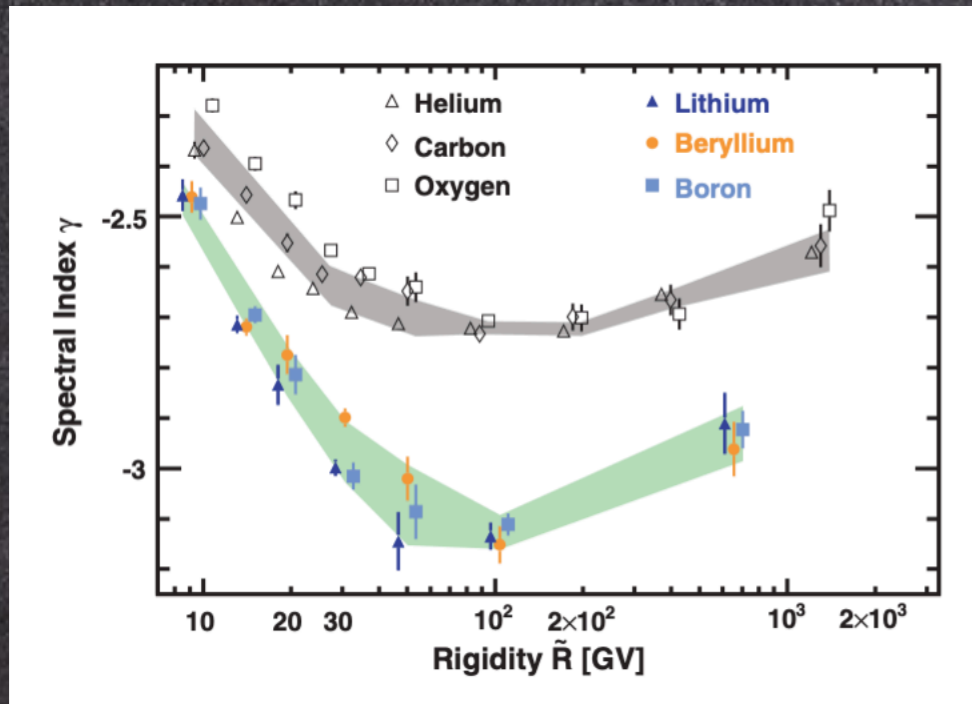
Interactions and decays in the Galaxy



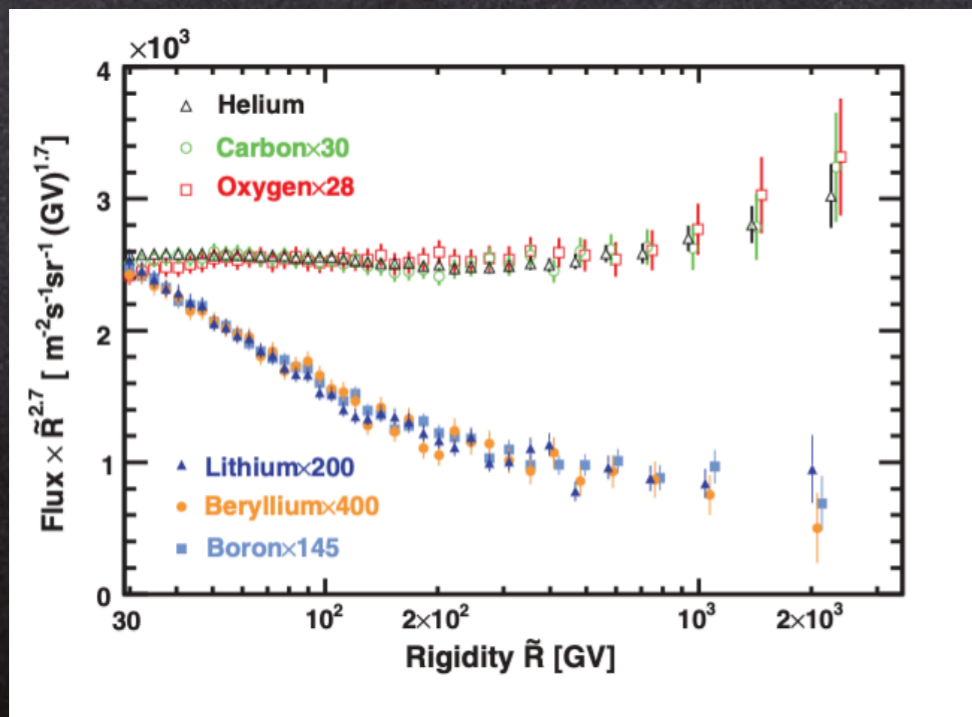
Courtesy of M. Korsmeier

The spectrum of secondary fluxes

See talk by A. Kounine



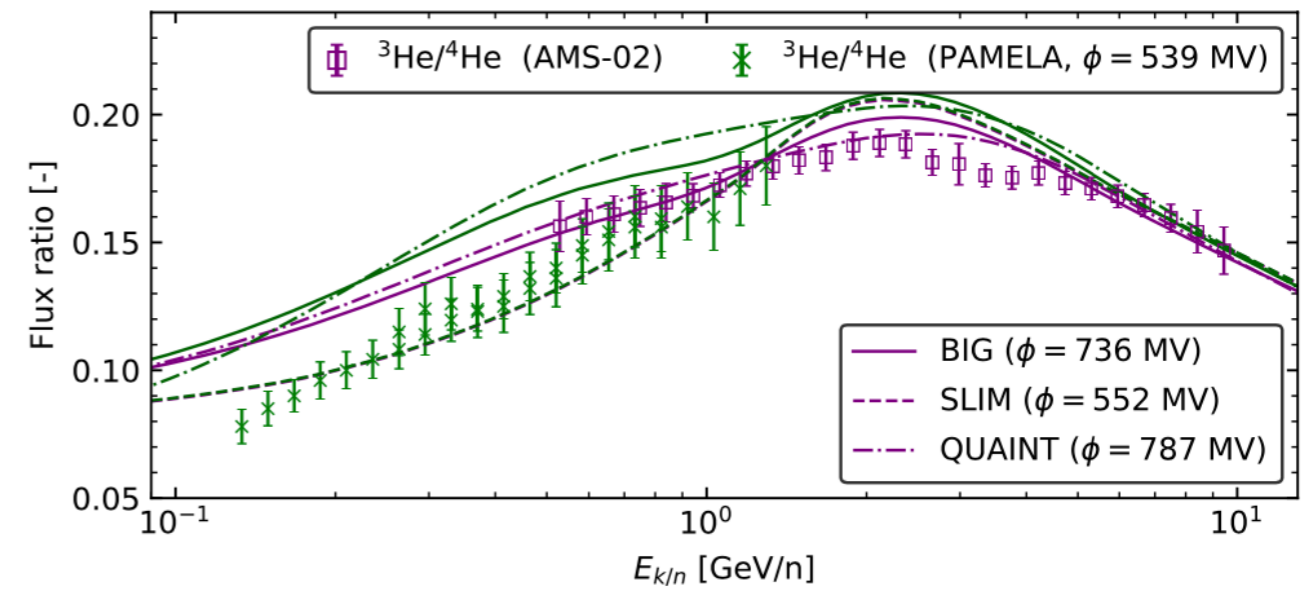
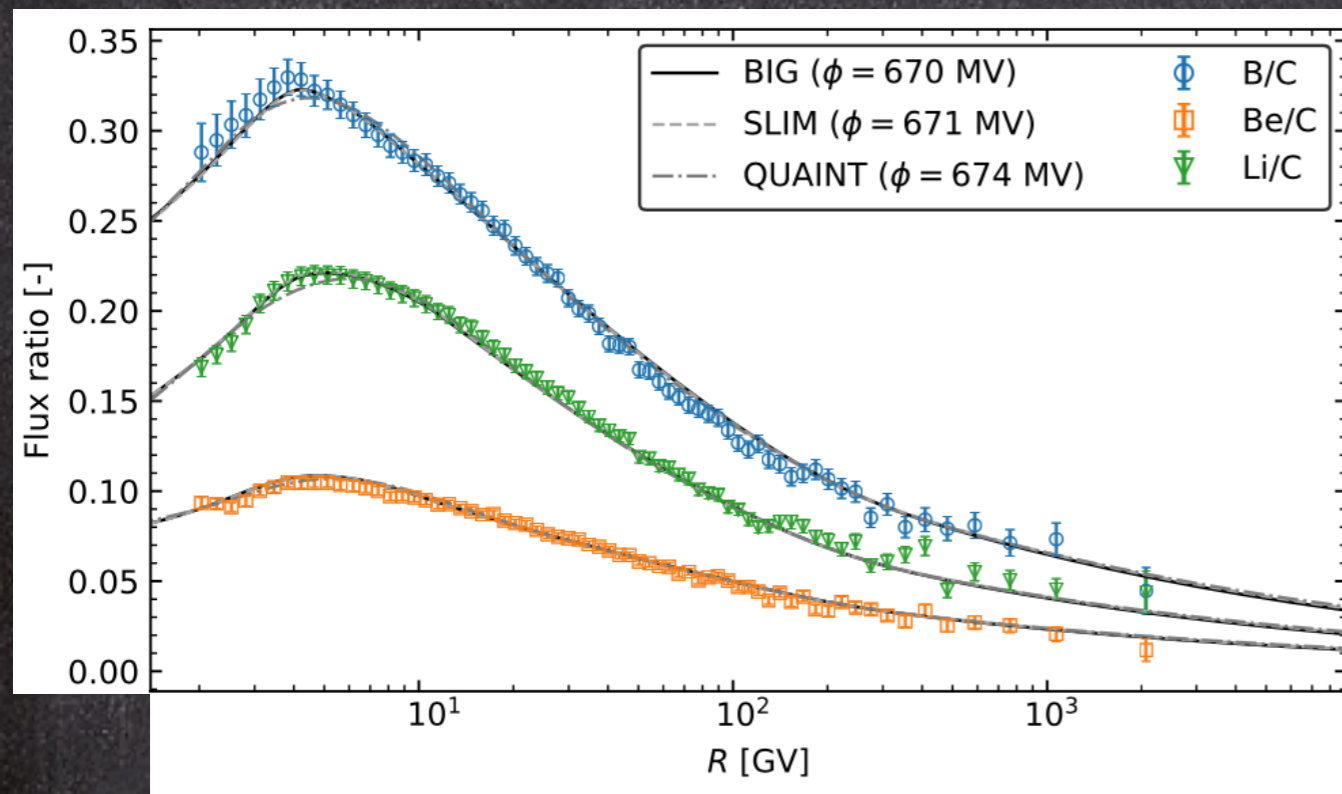
The rigidity dependence of Li, Be and B are nearly identical, but different from the primary He, C and O (and also p).



Li, Be, B fluxes measured by Pamela and AMS show an identical hardening w.r.t. energy above 200 GV. The spectral index of secondaries hardens 0.13 ± 0.03 more than for primaries

Propagation models vs data

Weinrich+ A&A 2020



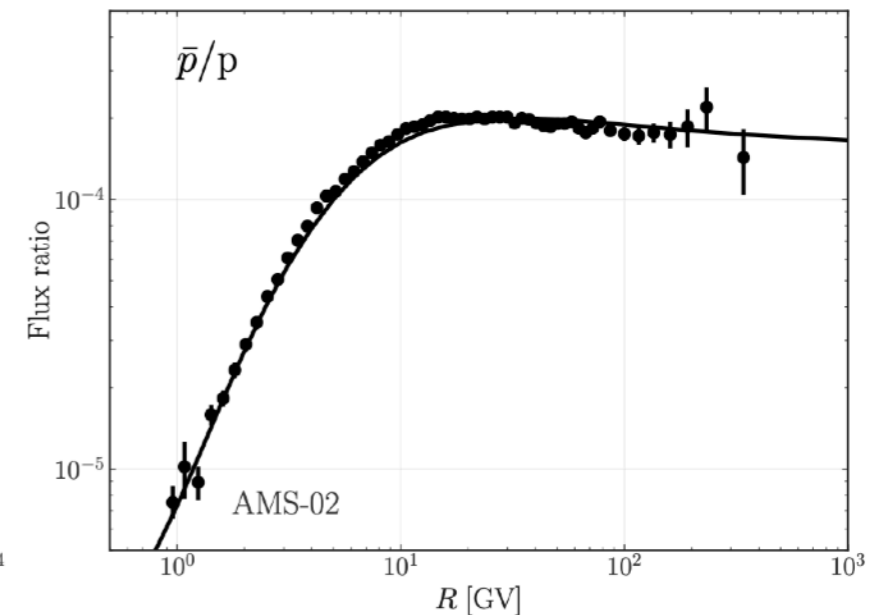
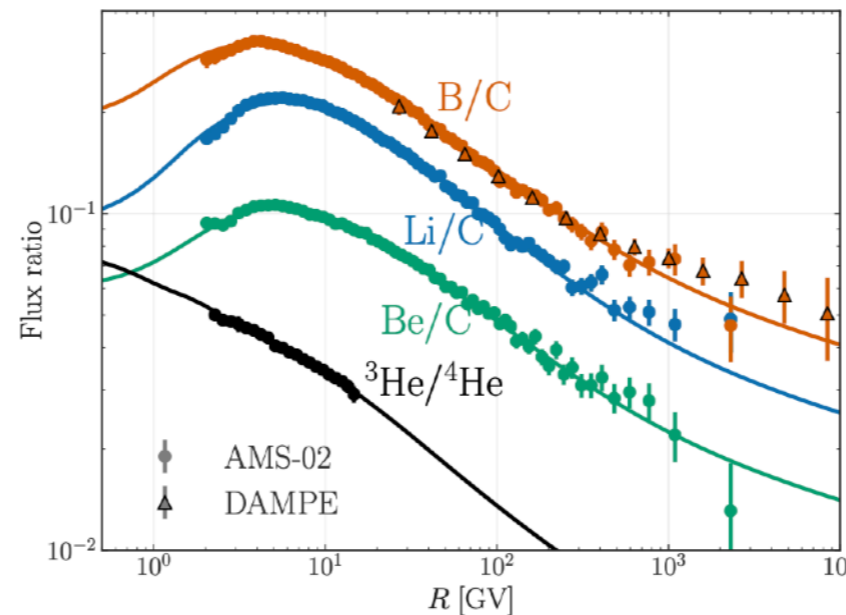
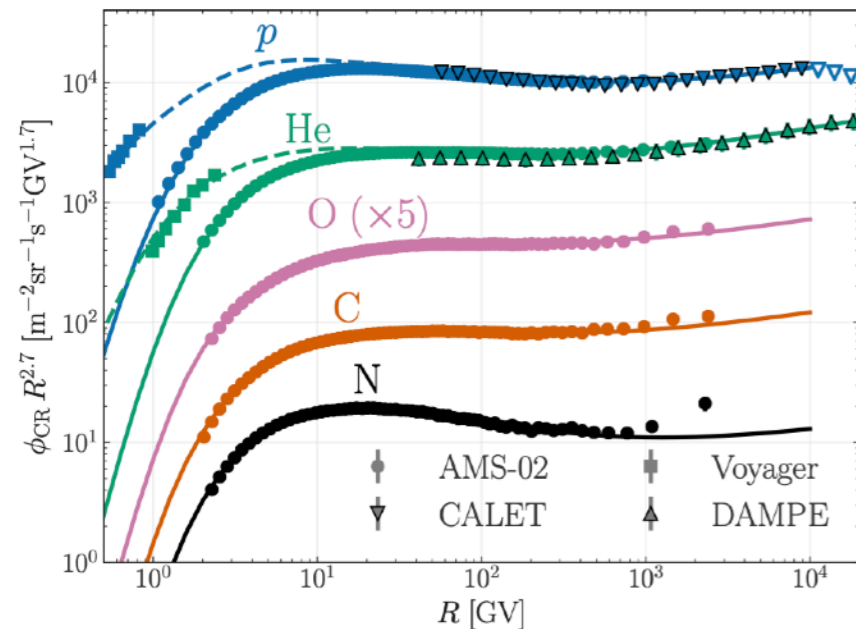
Data on secondary/primary species are well described by propagation model with diffusion coefficient power index $\delta = 0.50 \pm 0.03$.

Convection + reaccelerating, or pure diffusion both work.

Propagation models vs data

Several propagation models are tested

Di Mauro, Korsmeier, Cuoco 2311.17150

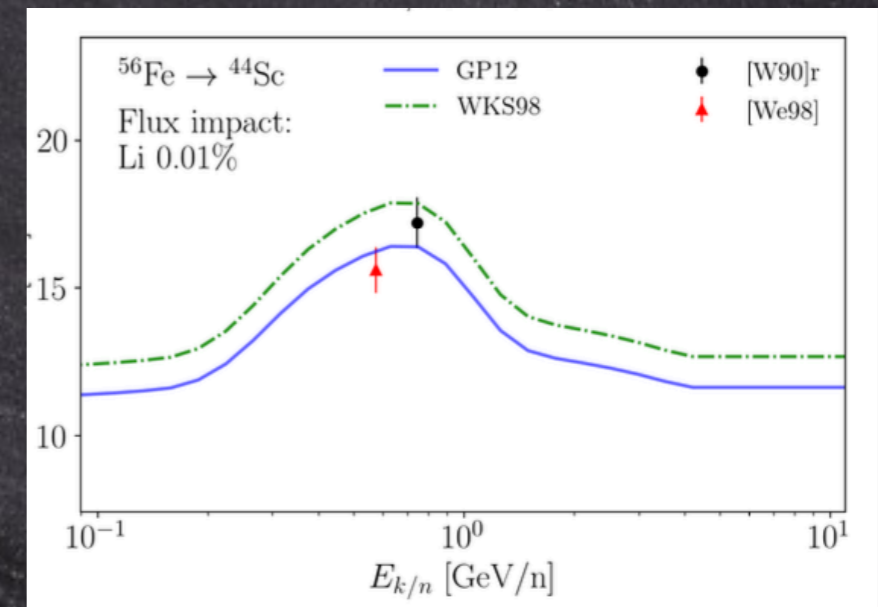
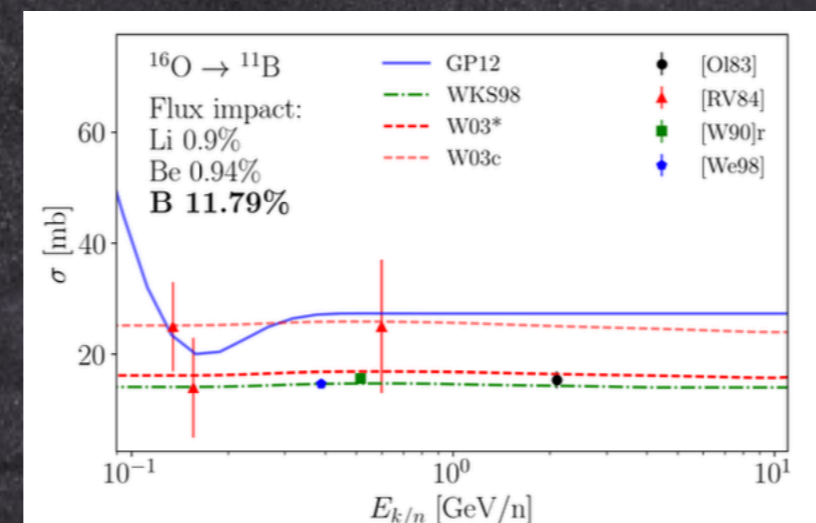
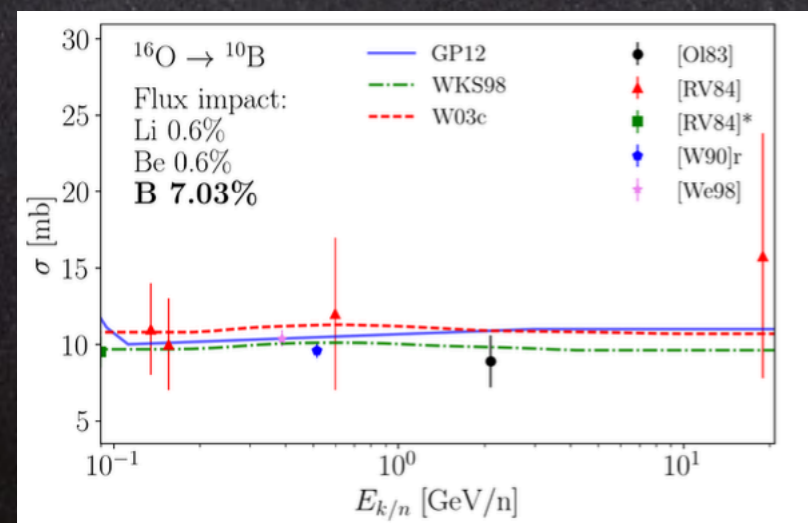
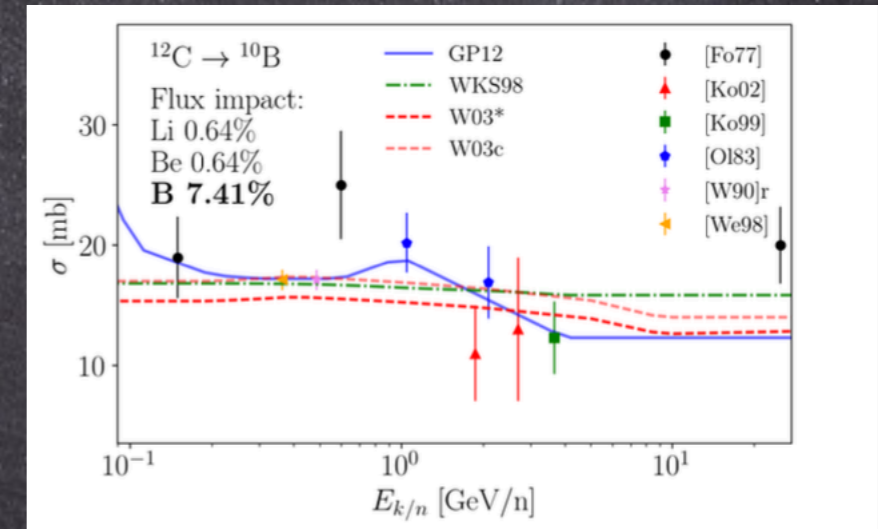
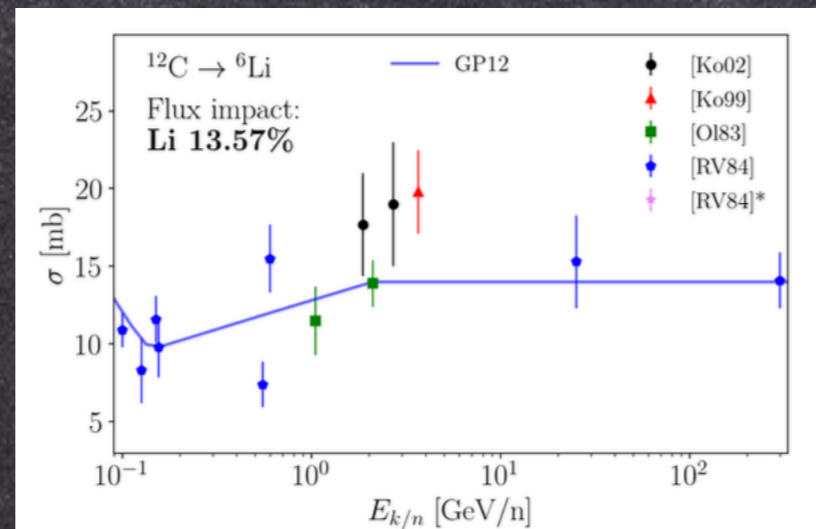
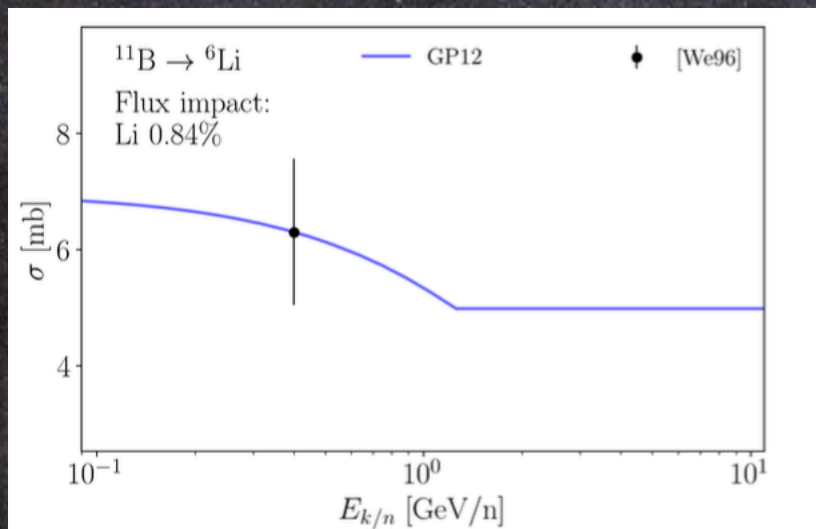


Fragmentation cross section uncertainties currently prevent a better understanding of CR propagation

Cross sections for Galactic cosmic rays

Data driven parameterizations (Silberberg & Tsao), semi-empirical formulae (Webber+), parametric formulae/direct fit to the data (Galprop), MonteCarlo codes (Fluka, Geant, ...)

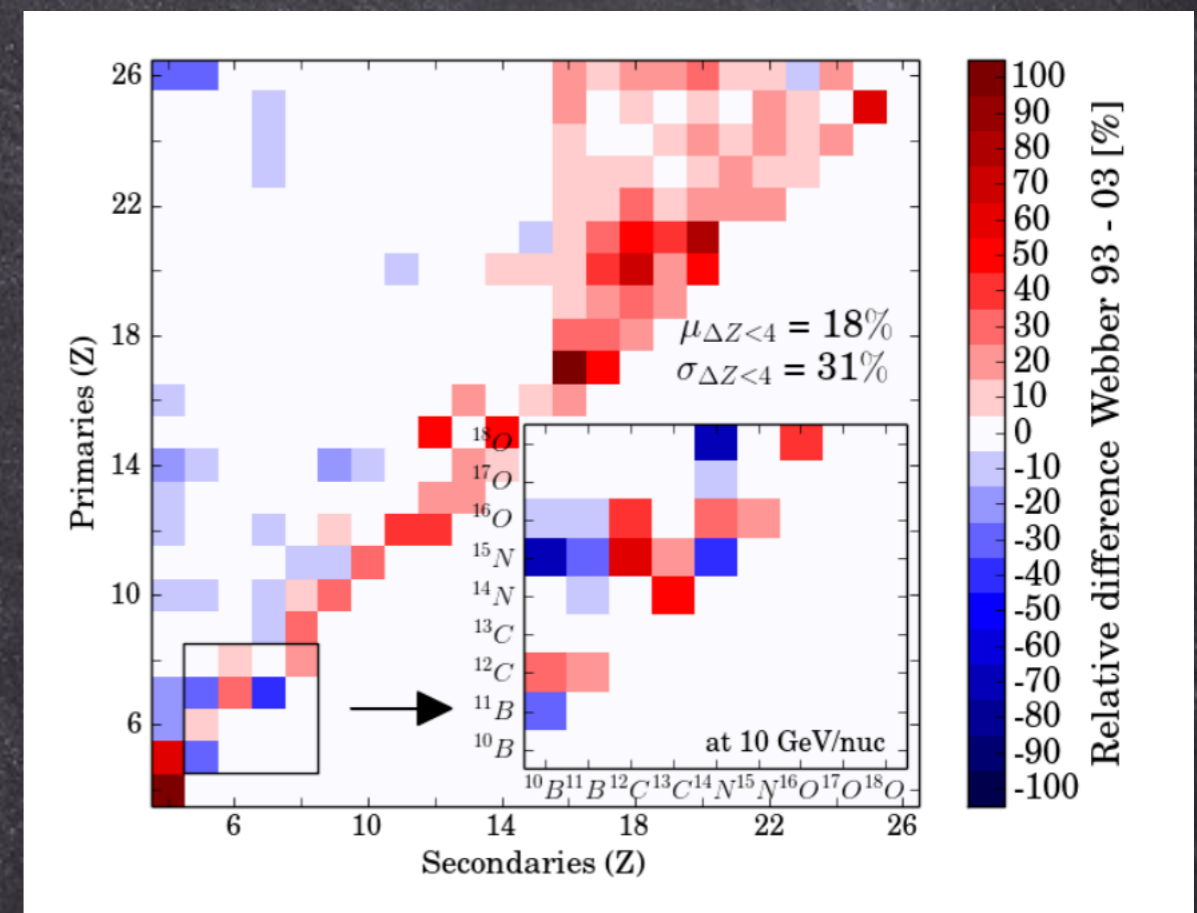
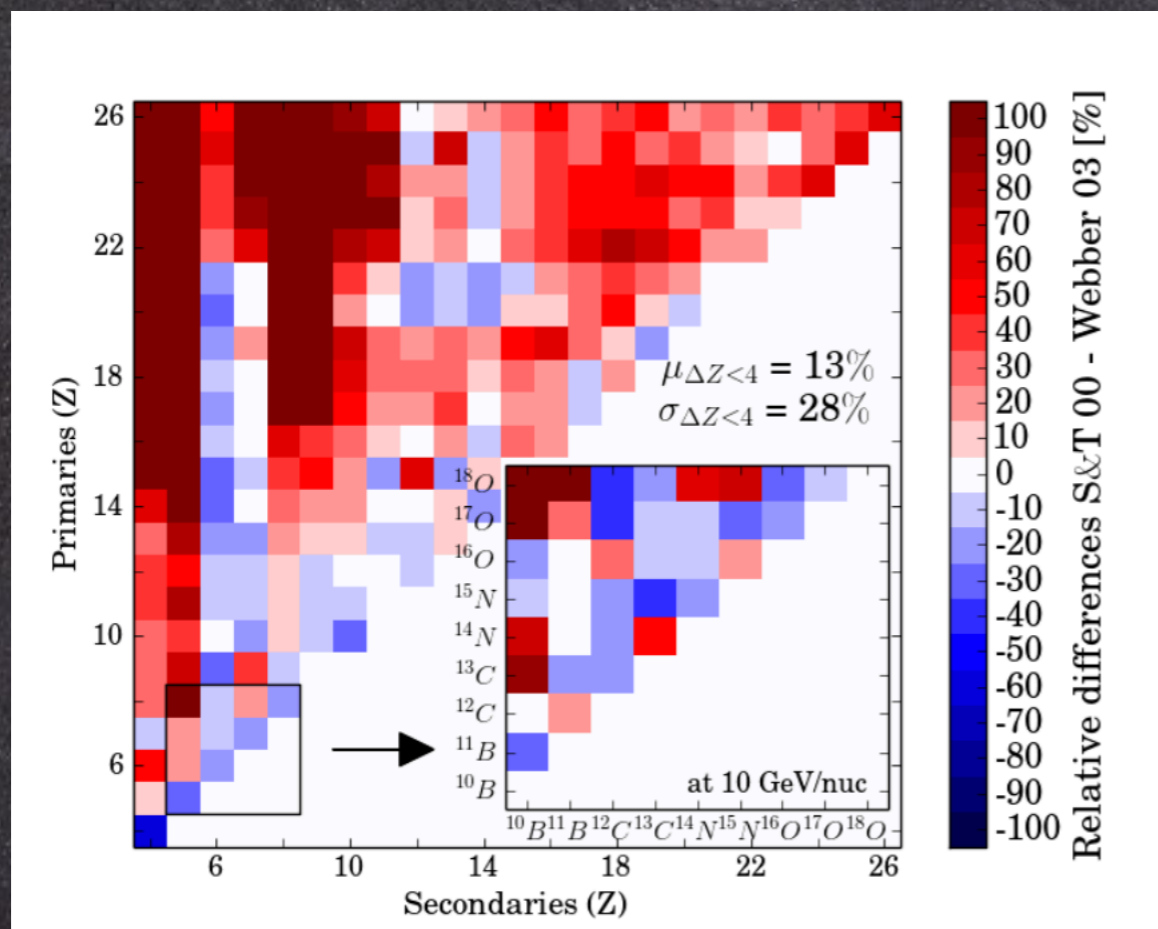
Genolini, Moskalenko, Maurin, Unger PRC 2018



Differences in the XS parameterizations

Genolini, Putze, Salati, Serpico A&A 2015

Differences in one parameterization wrt a benchmark model

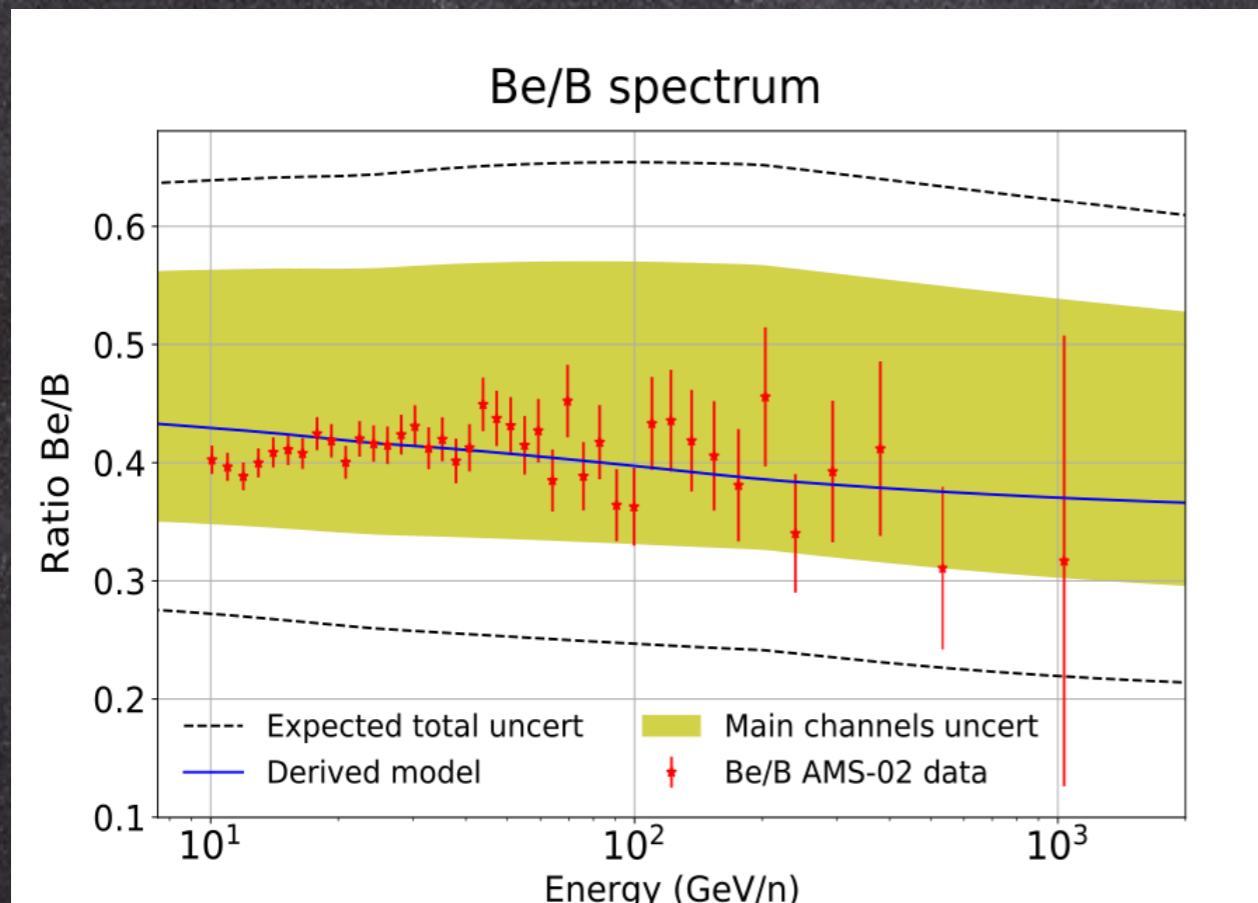


Even with the same, although scarce data, interpretation may be different

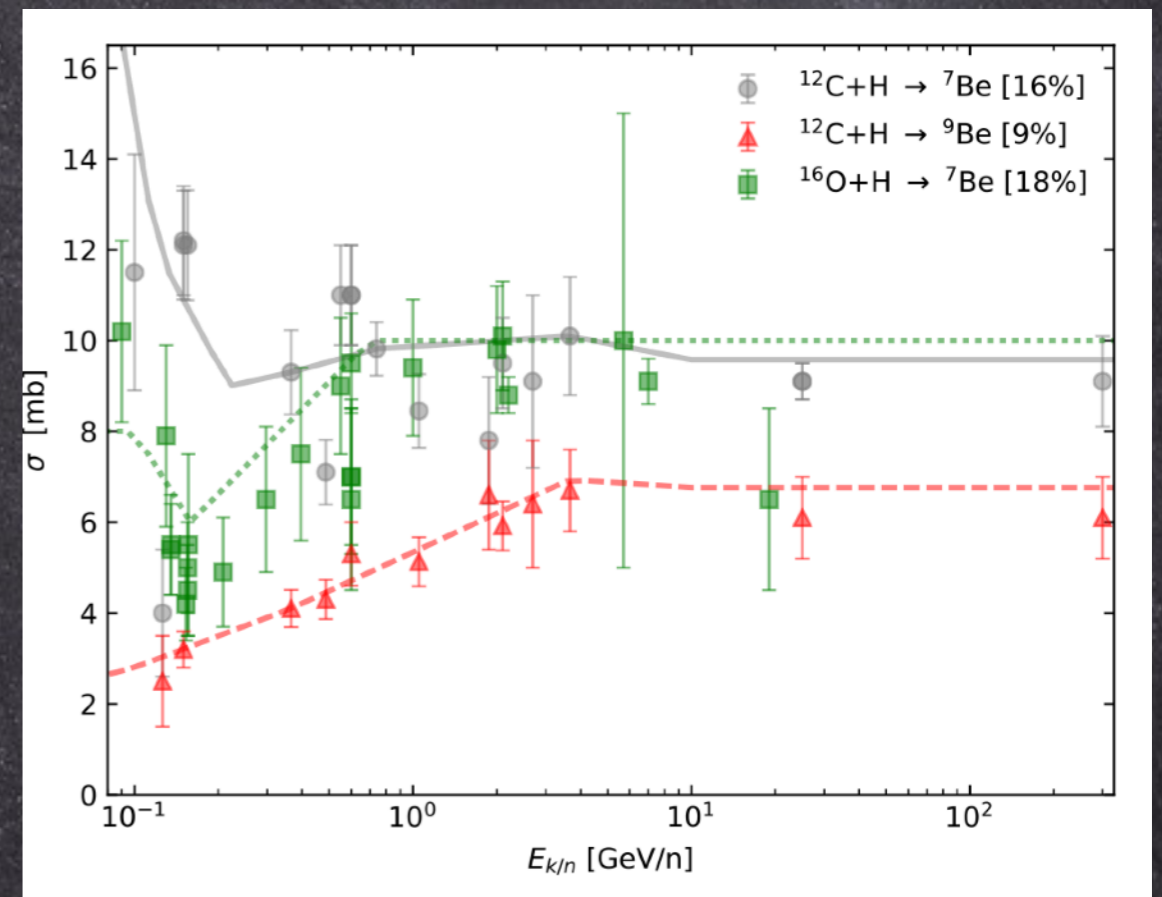
Fragmentation cross sections

They matter in both directions: as a loss term for progenitors,
as a source term for daughters

De La Torre Luquet+ JCAP 2021



Weinrich+ A&A 2021

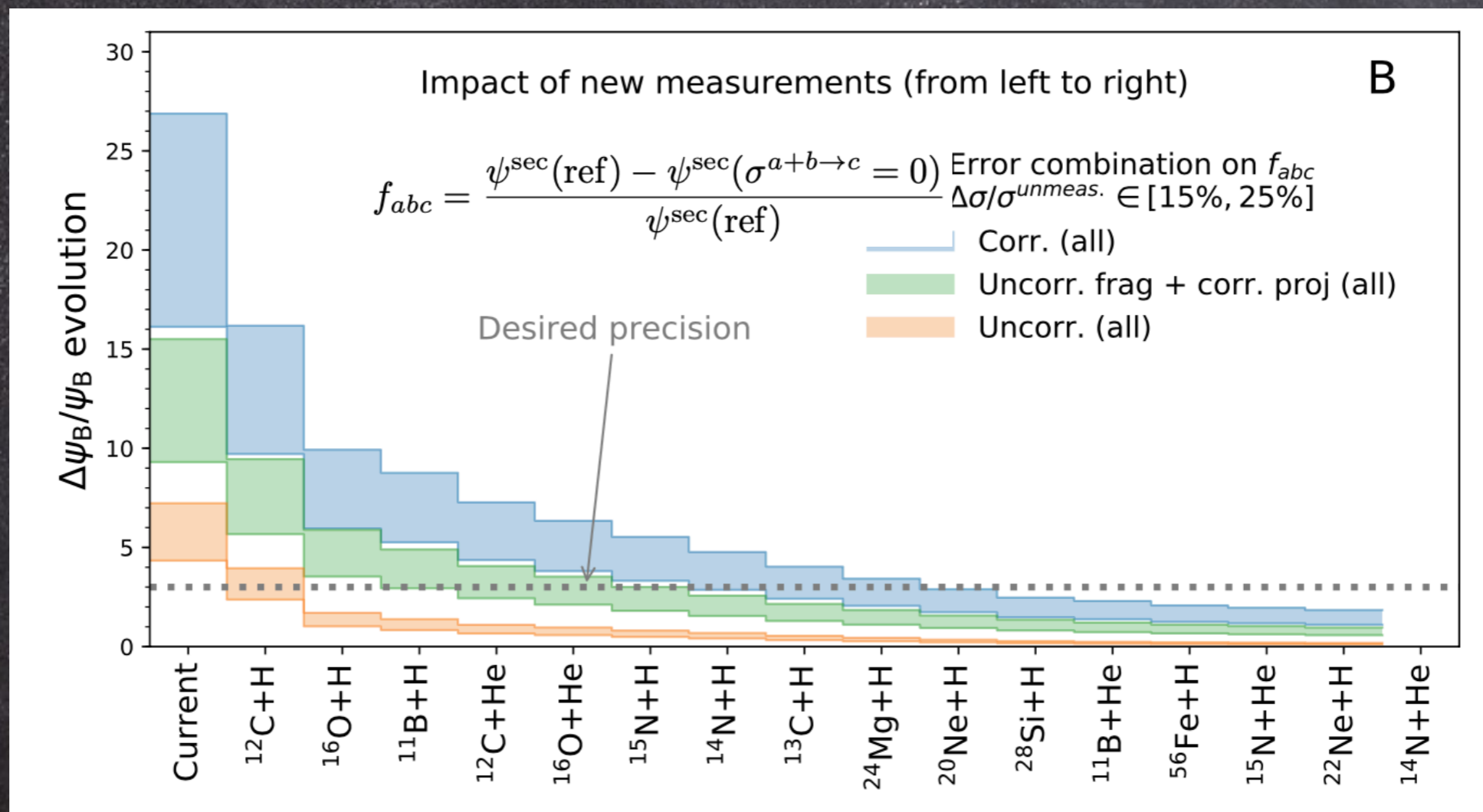


Probably the most limiting aspect now
Dedicated campaigns are needed (LHC, NA61, Amber, ...)

Most relevant physics cases

Improve Boron production cross sections

Genolini, Moskalenko, Maurin, Unger PRC 2018

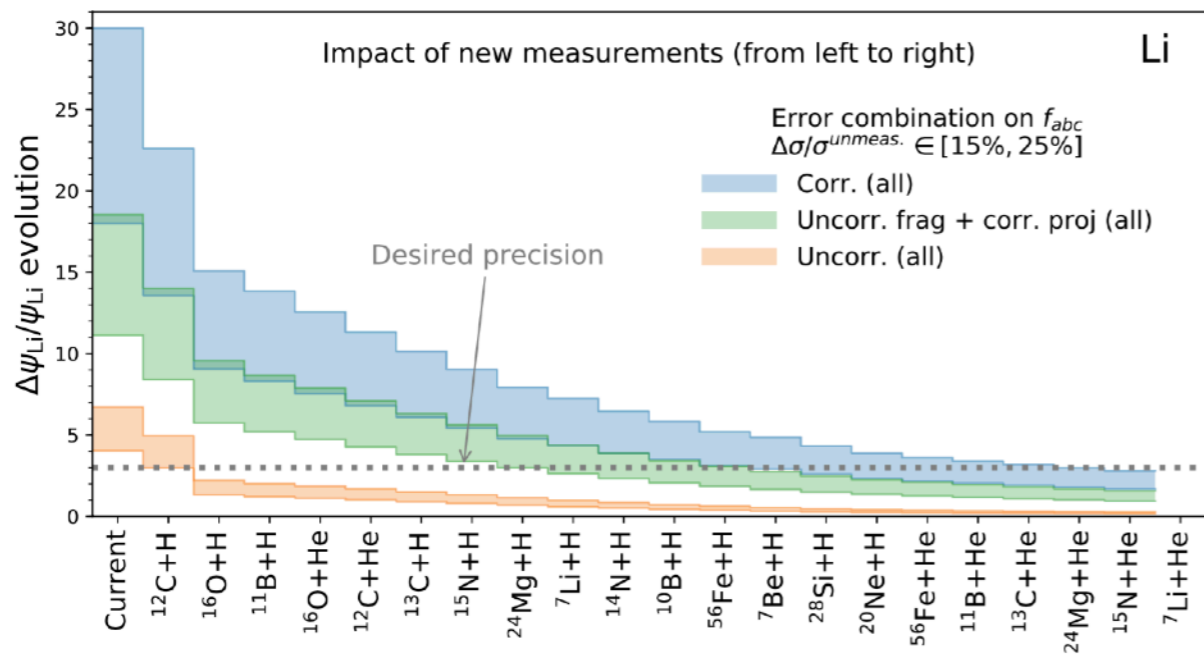


The evolution of error on the calculated B flux as if new reactions were measured with perfect accuracy ($^{12}\text{C}+\text{H}$ is the most critical one)

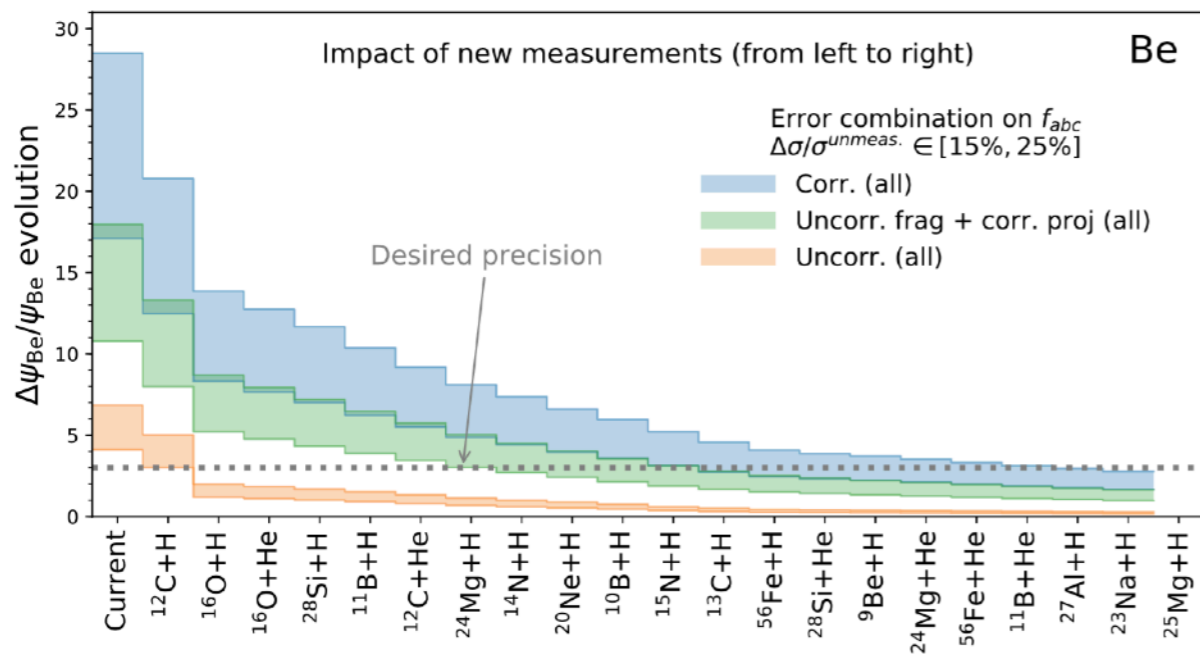
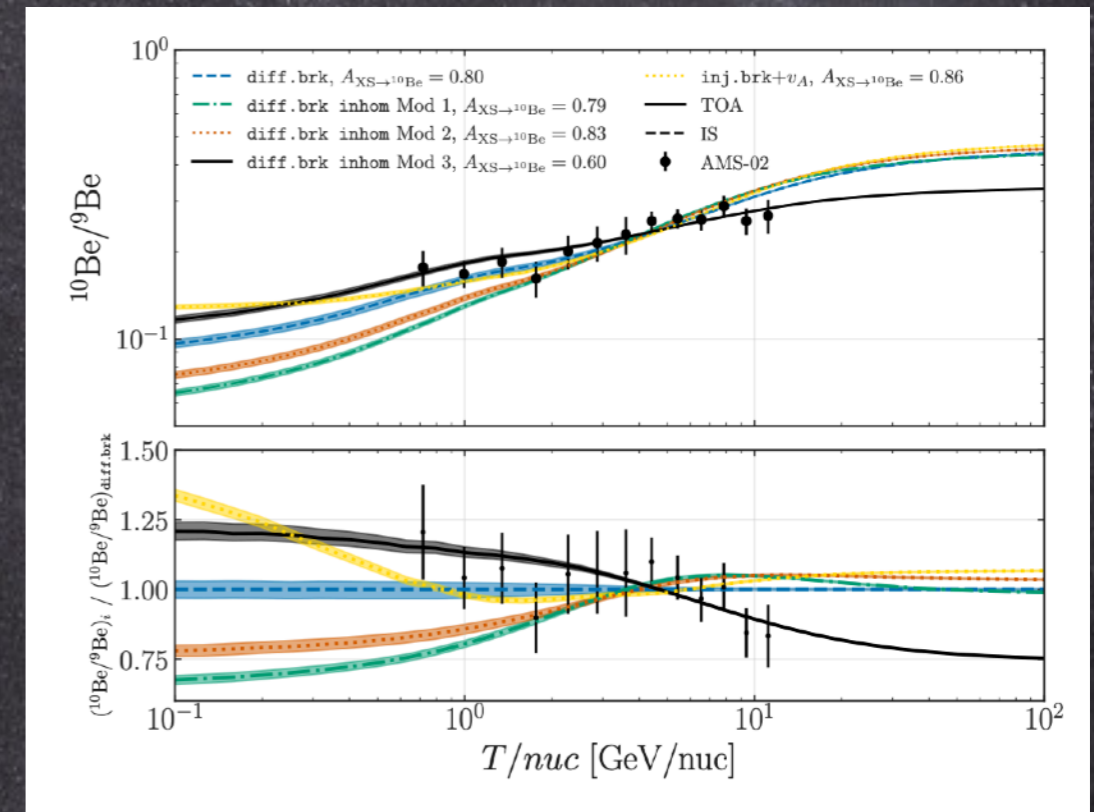
Most relevant physics cases

Improve Lithium and Berillium production cross sections

Genolini, Moskalenko, Maurin, Unger PRC 2018



Evoli, Morlino, Blasi, Aloisio PRD 2020
 Di Mauro, Korsmeier, Cuoco 2311.17150

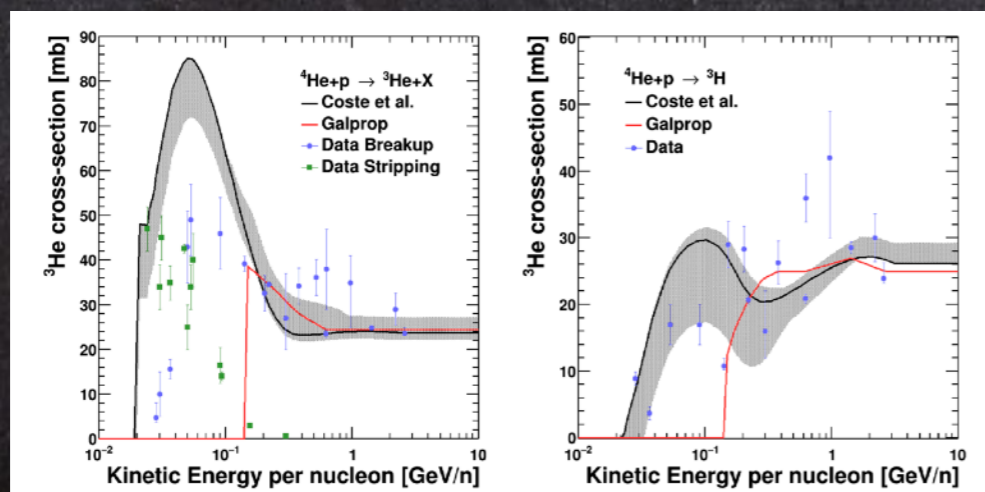
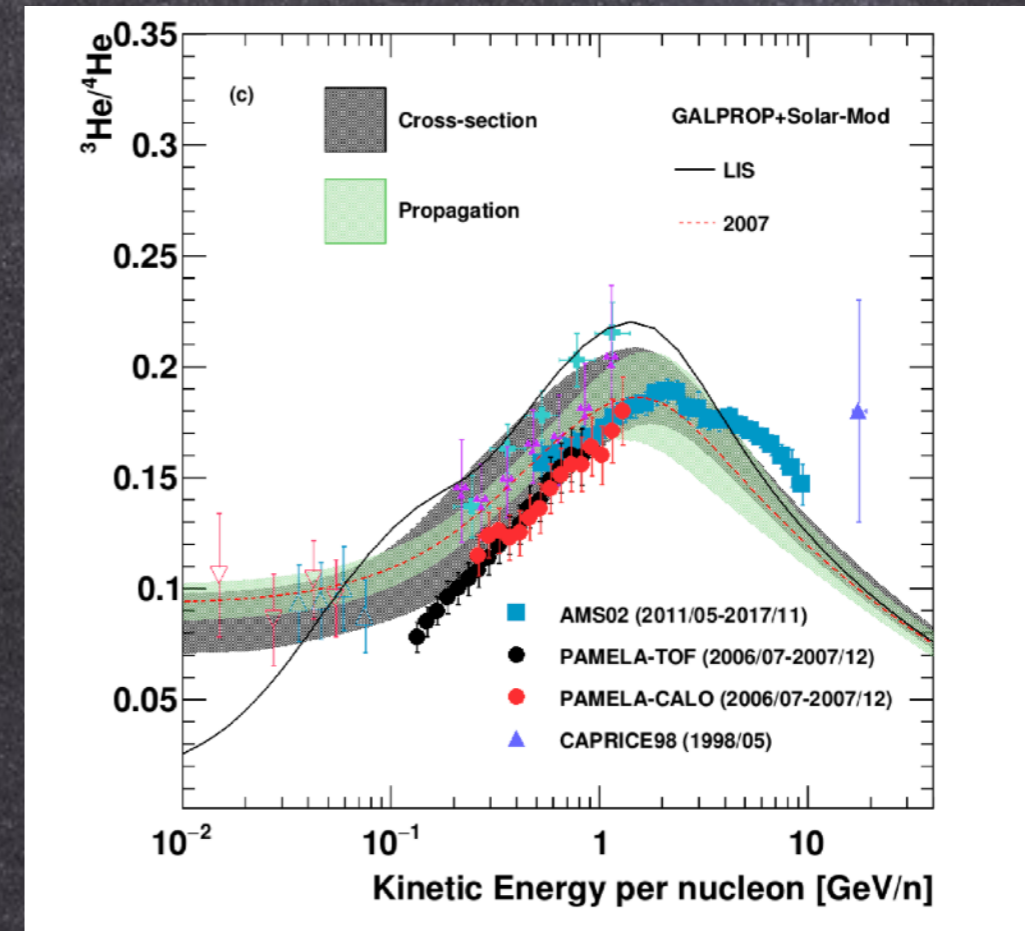
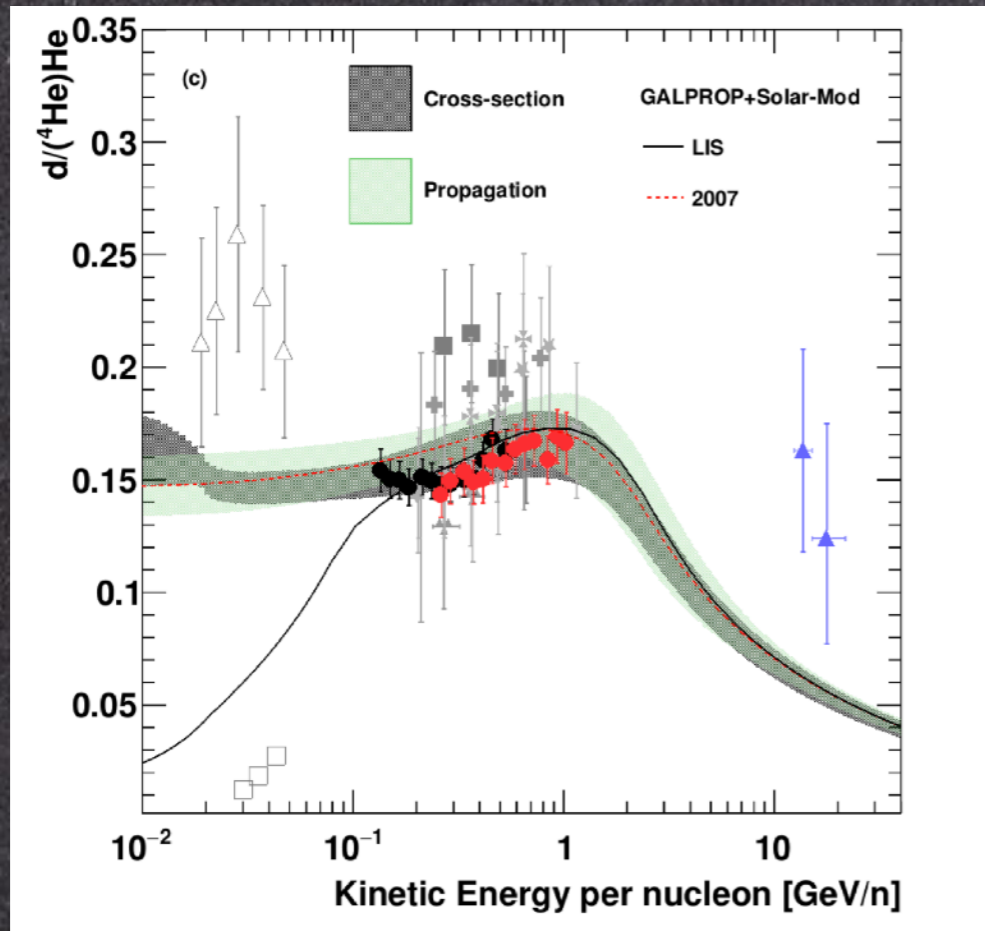


Berillium case very relevant

H and He isotopes

Coste+ A&A 2012

Gomez-Coral+ PRD 2024



Modelization of cross sections relies on poor data

Antimatter

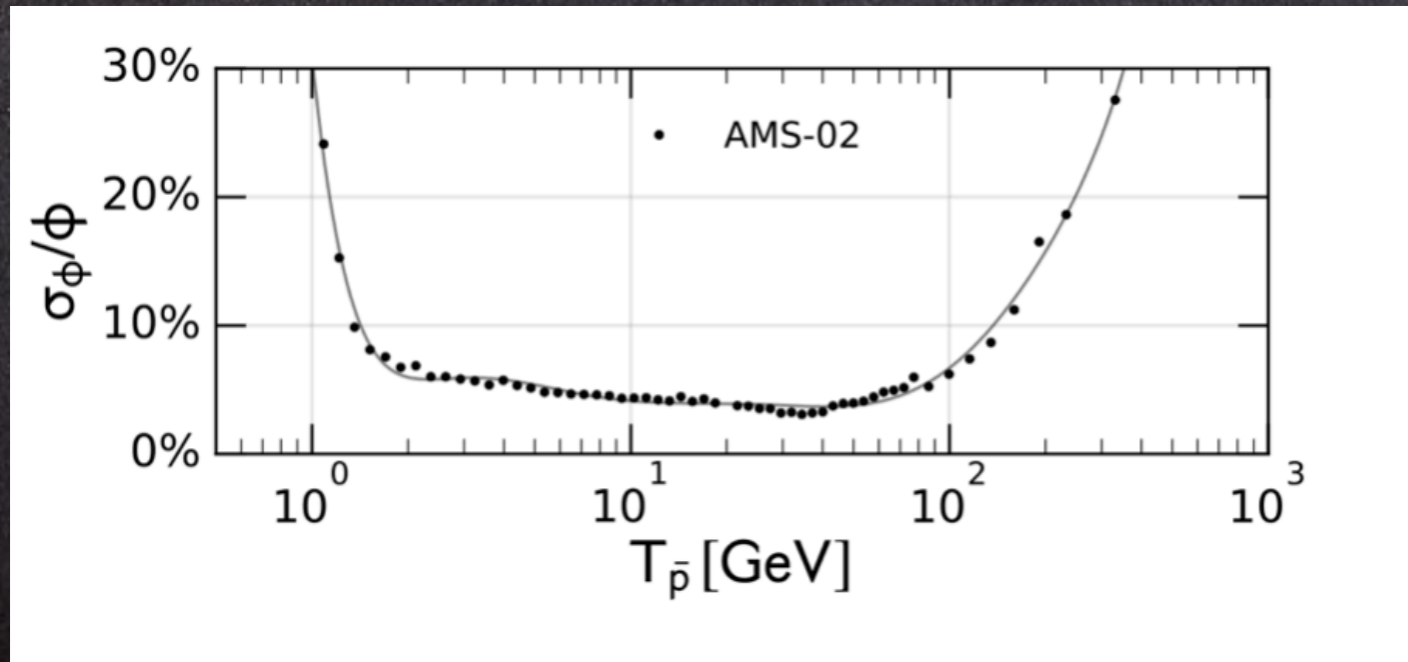
in the Galaxy

Antiproton production by inelastic scatterings

Korsmeier, FD, Di Mauro PRD 2018

$$q_{ij}(T_{\bar{p}}) = \int_{T_{\text{th}}}^{\infty} dT_i 4\pi n_{\text{ISM},j} \phi_i(T_i) \frac{d\sigma_{ij}}{dT_{\bar{p}}}(T_i, T_{\bar{p}}). \quad \frac{d\sigma_{ij}}{dT_{\bar{p}}}(T, T_{\bar{p}}) = p_{\bar{p}} \int d\Omega \sigma_{\text{inv}}^{(ij)}(T_i, T_{\bar{p}}, \theta).$$

Data from space are very precise

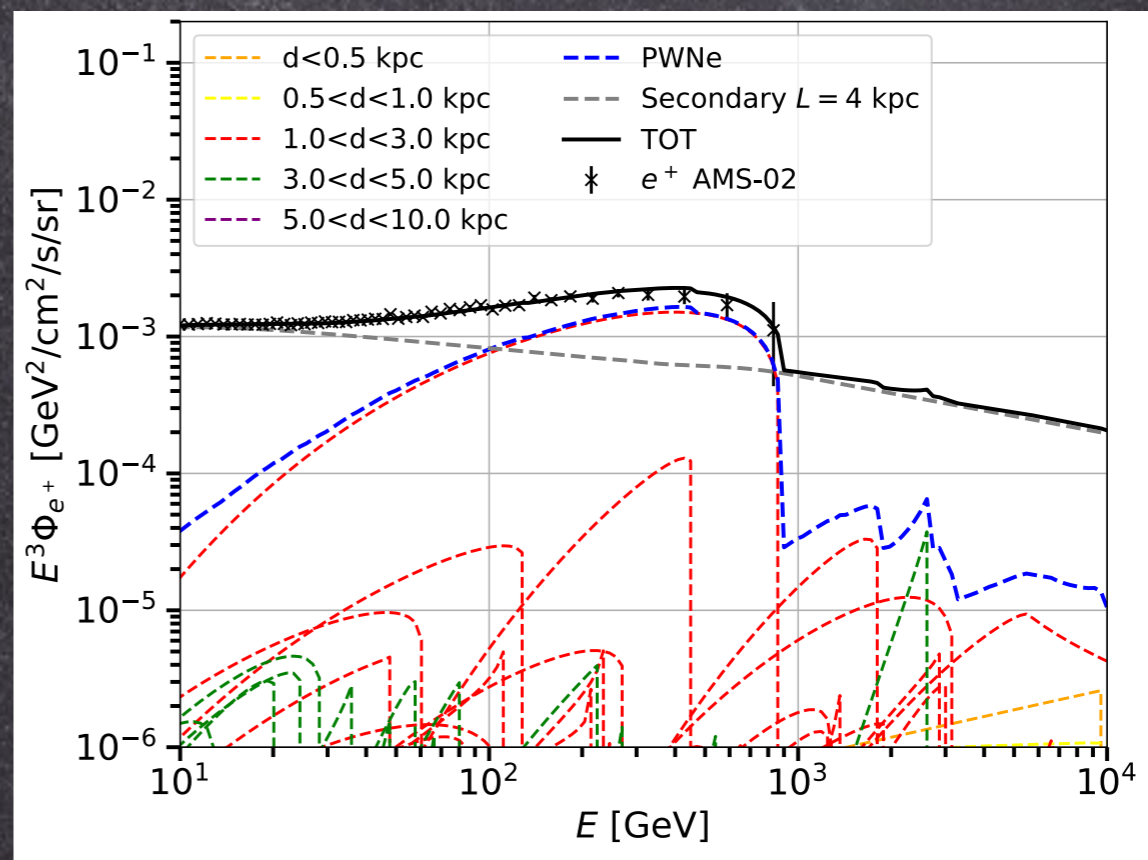
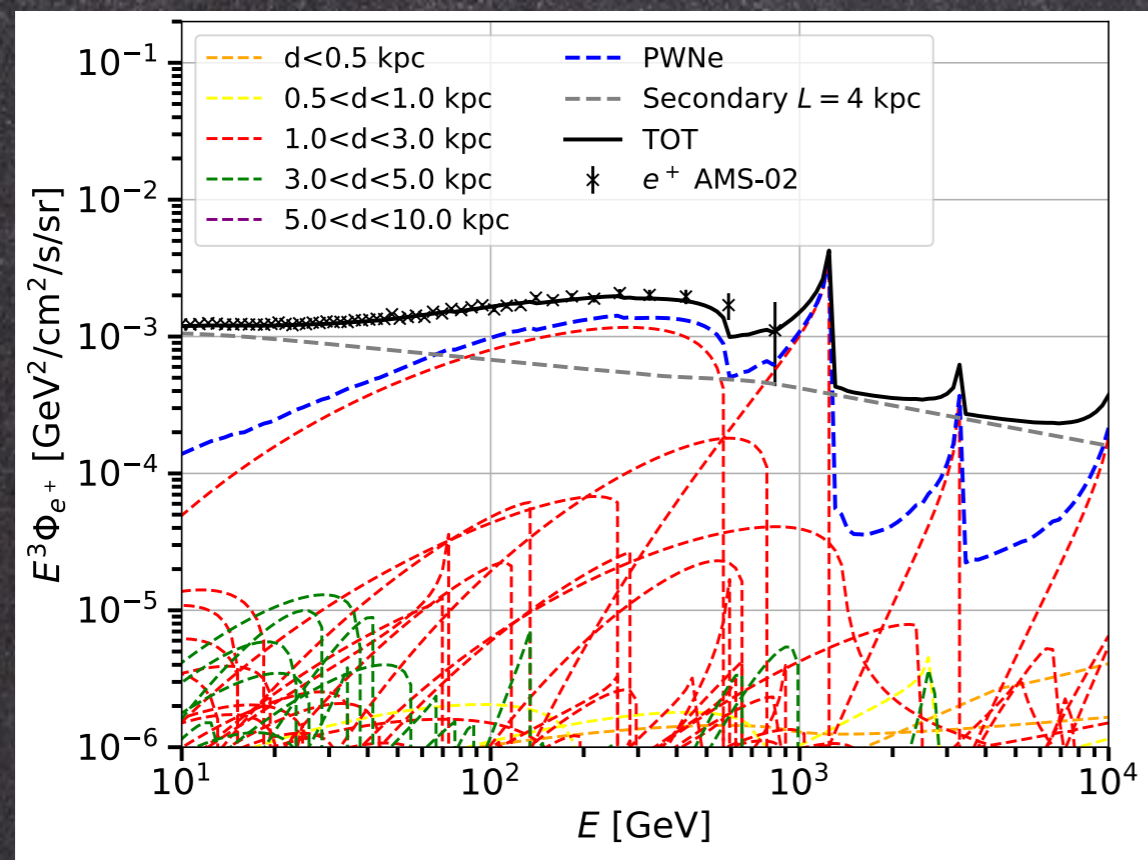


See talk by
M. Di Mauro

We need cross sections at $<3\%$

Fit of Galactic pulsar populations to AMS-02 e^+ data

Orusa, Di Mauro, FD, Manconi JCAP 2021



The contribution of pulsars to e^+ is dominant above 100 GeV and may have different features.

$E > 1$ TeV: unconstrained by data.

Secondaries forbid evidence of sharp cut-off.

No need for Dark Matter, indeed

The e^\pm production chain from π^\pm production

$$\frac{d\sigma_{ij}}{dT_{e^+}}(T_i, T_{e^+}) = \int dT_{\pi^+} \frac{d\sigma_{ij}}{dT_{\pi^+}}(T_i, T_{\pi^+}) P(T_{\pi^+}, T_{e^+})$$

Integral over the pion production cross section convolved with the probability density function P

$$\frac{d\sigma_{ij}}{dT_{\pi^+}}(T_i, T_{\pi^+}) = p_{\pi^+} \int d\Omega \sigma_{\text{inv}}^{(ij)}(T_i, T_{\pi^+}, \theta)$$

The pion production cross section is the integral of the Lorentz Invariant cross section over scattering angle (or p_T)

$$\sigma_{\text{inv}}^{(ij)} = E_{\pi^+} \frac{d^3\sigma_{ij}}{dp_{\pi^+}^3}$$

← data

A fit is performed on the σ_{inv} data

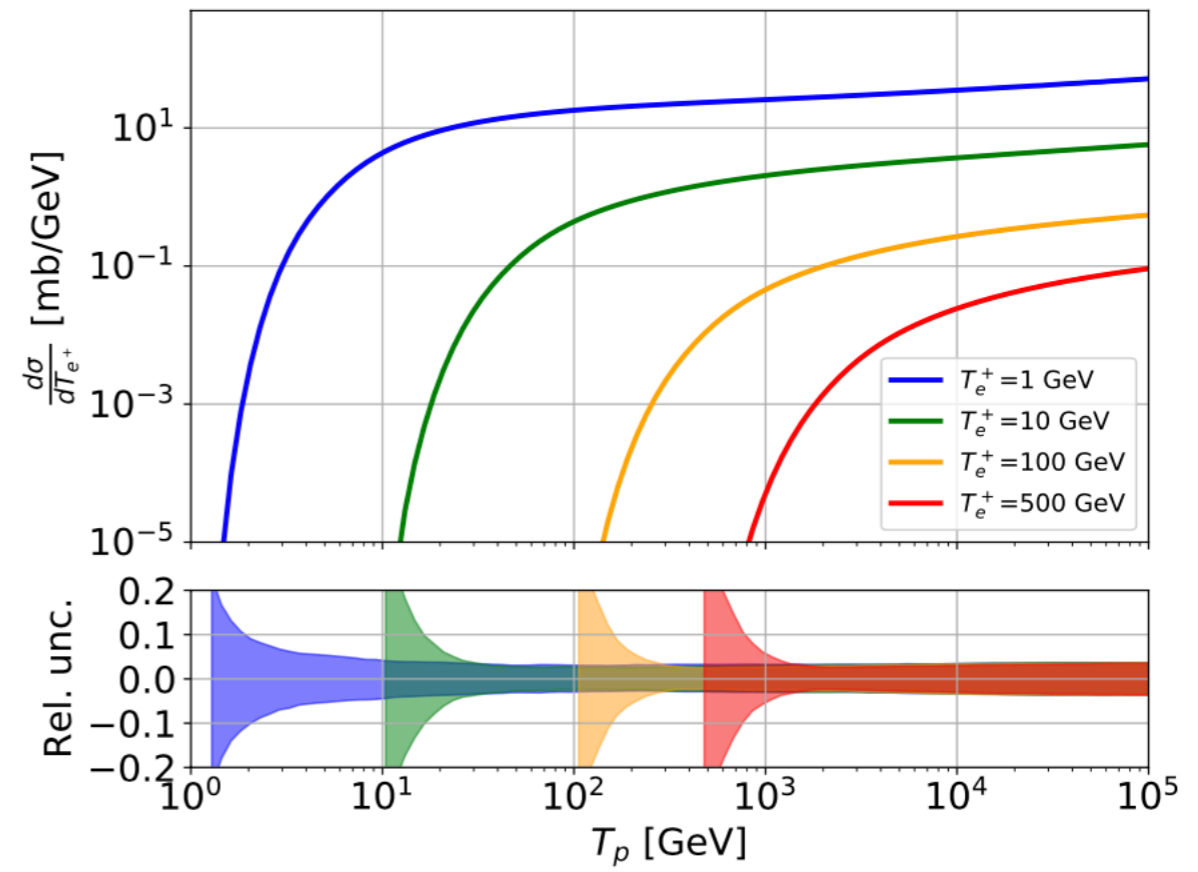
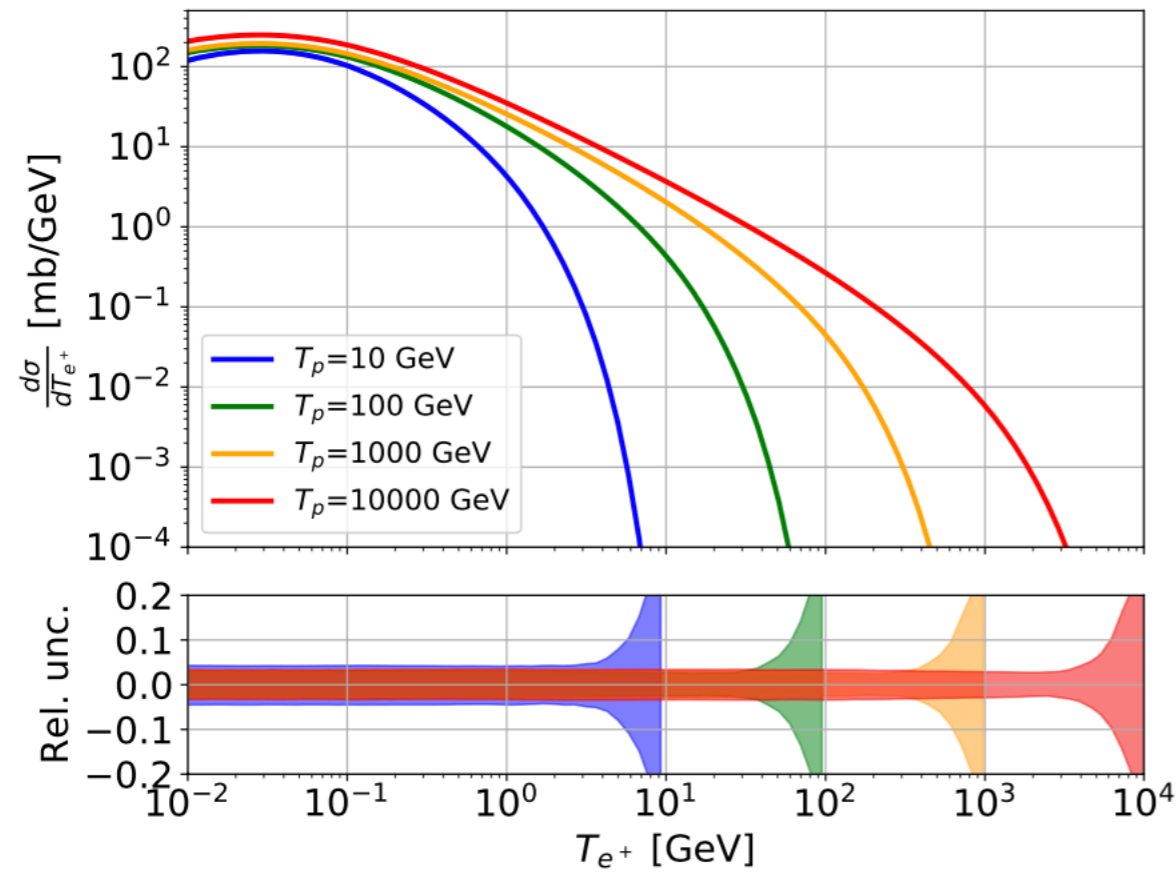
L. Orusa, M. Di Mauro, FD, M. Korsmeier PRD 2022

Experiment	\sqrt{s} [GeV]	σ_{inv}	n	Ref.
NA49	17.3	×	×	[22]
ALICE	900	×	-	[23]
CMS	900, 2760, 7000, 13000	×	-	[24, 25]
Antinucci	π^+ (3.0, 3.5, 4.9, 5.0, 6.1, 6.8)	-	×	[26]
	π^- (3.0, 3.5, 4.9, 5.0, 6.1, 6.8)	-	×	[26]
	K^+ (2.8, 3.0, 3.2, 5.0, 6.1, 6.8)	-	×	[26]
	K^- (4.9, 5.0, 6.1, 6.8)	-	×	[26]
NA61	6.3, 7.7, 8.8, 12.3, 17.3	-	×	[21]

We use data on σ_{inv} , the multiplicity n or both.

Results on the σ_{inv} for π^+ production

L. Orusa, M. Di Mauro, FD, M. Korsmeier PRD 2022

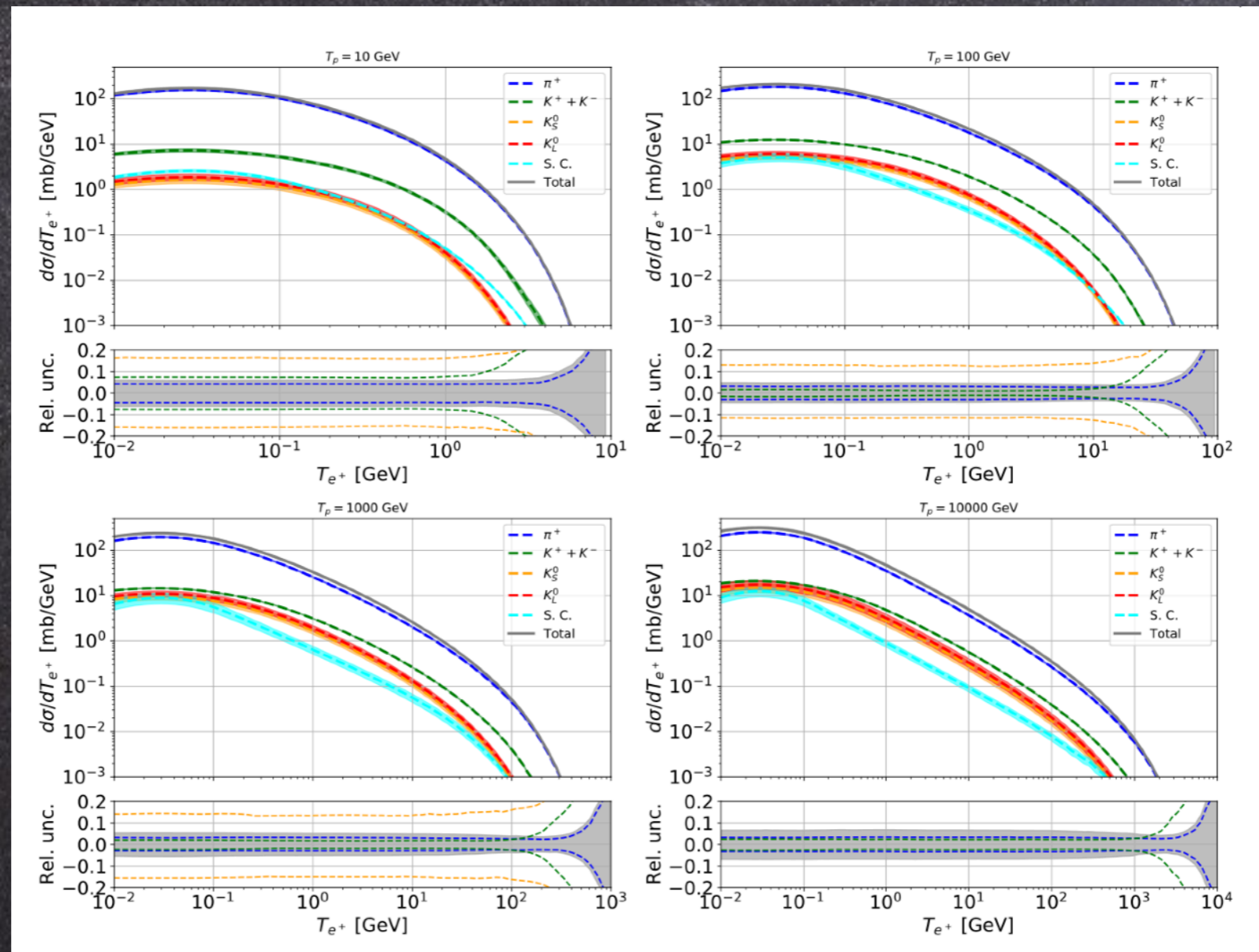


Data are fitted with very small uncertainties

Our parameterizations result appropriate, data are very precise

Total cross section from $pp \rightarrow e^+ + X$

L. Orusa, M. Di Mauro, FD, M. Korsmeier PRD 2022

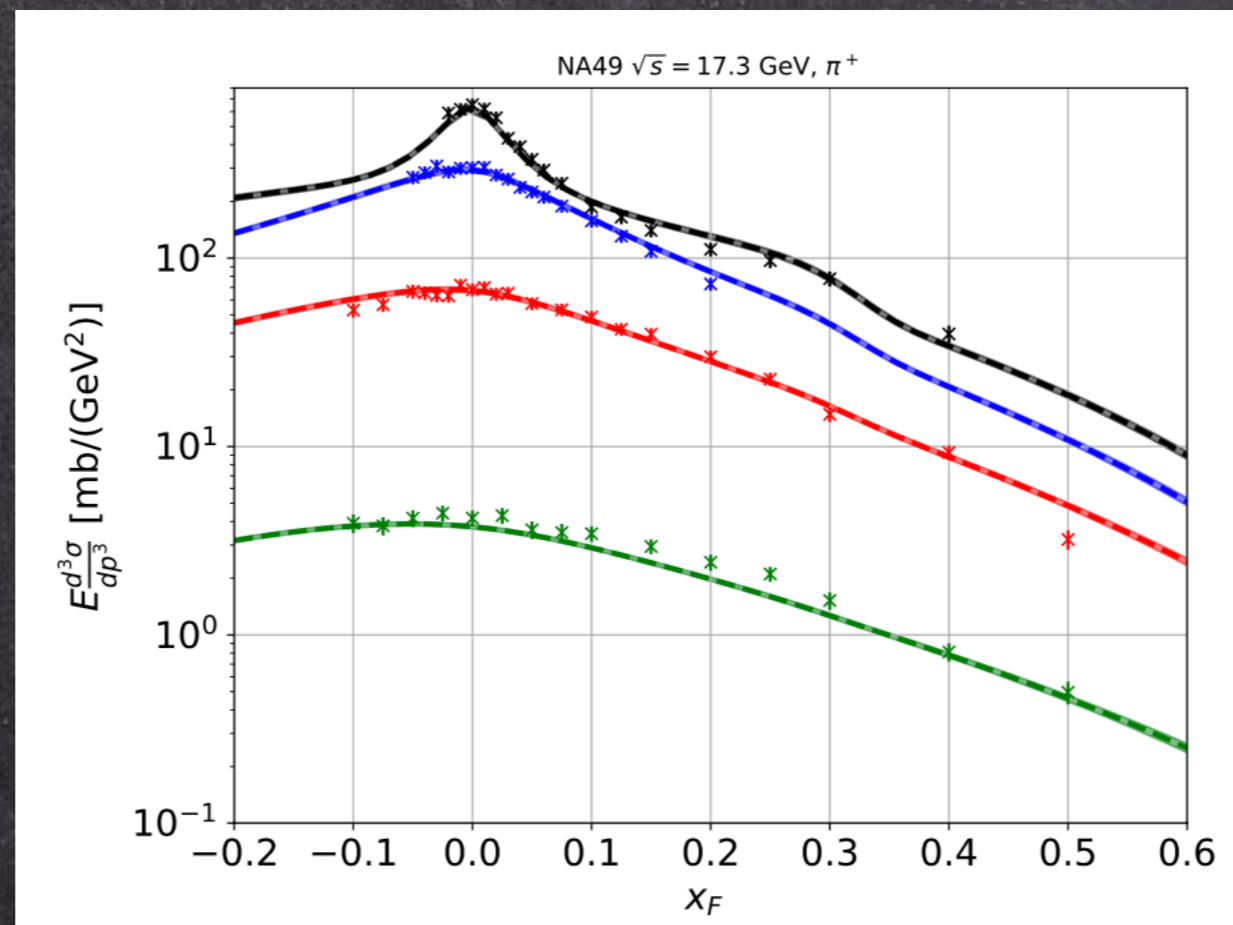


All channels contributing $>0.5\%$ are included.
Uncertainty globally contained to $<10\%$

Effect of scattering off nuclei

L. Orusa, M. Di Mauro, FD, M. Korsmeier PRD 2022

We need a model for the scattering involving He.
No data are there. We rely on NA49 $p+C \rightarrow e^++X$ data



Uncertainty is small, but very likely is not the true one

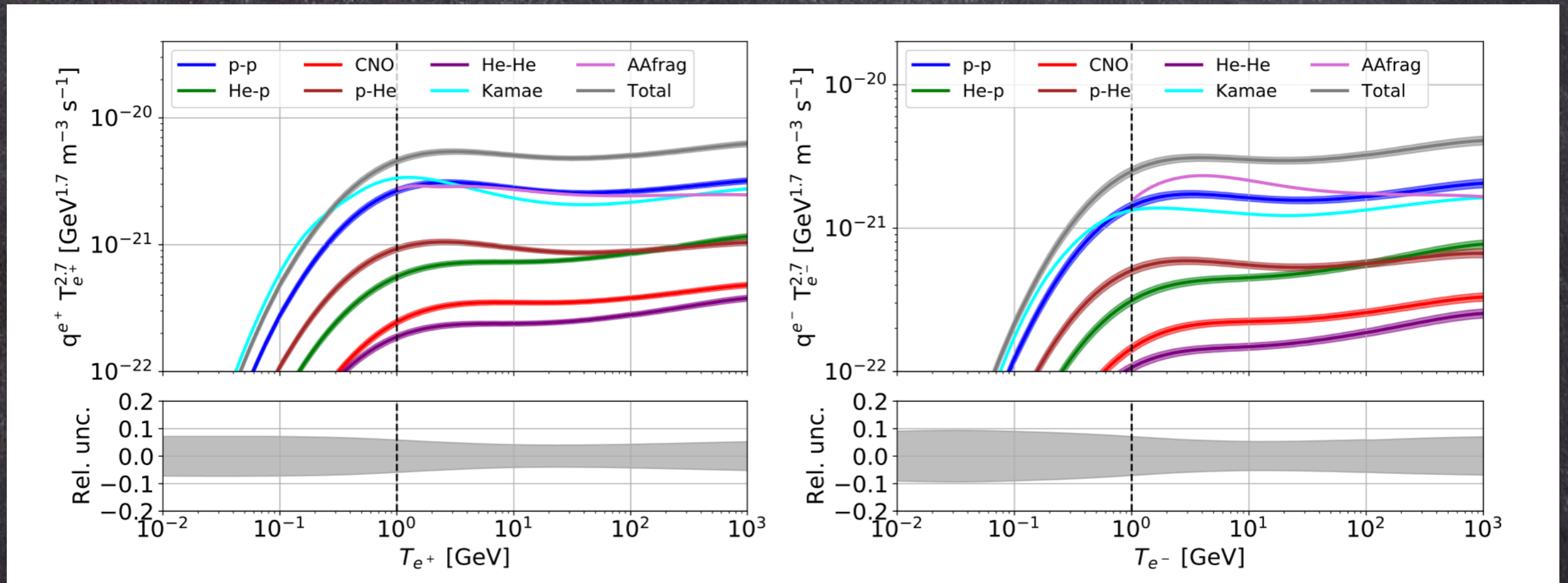
Data on He are necessary

Final results on e^+ cross section

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Positrons

Electrons

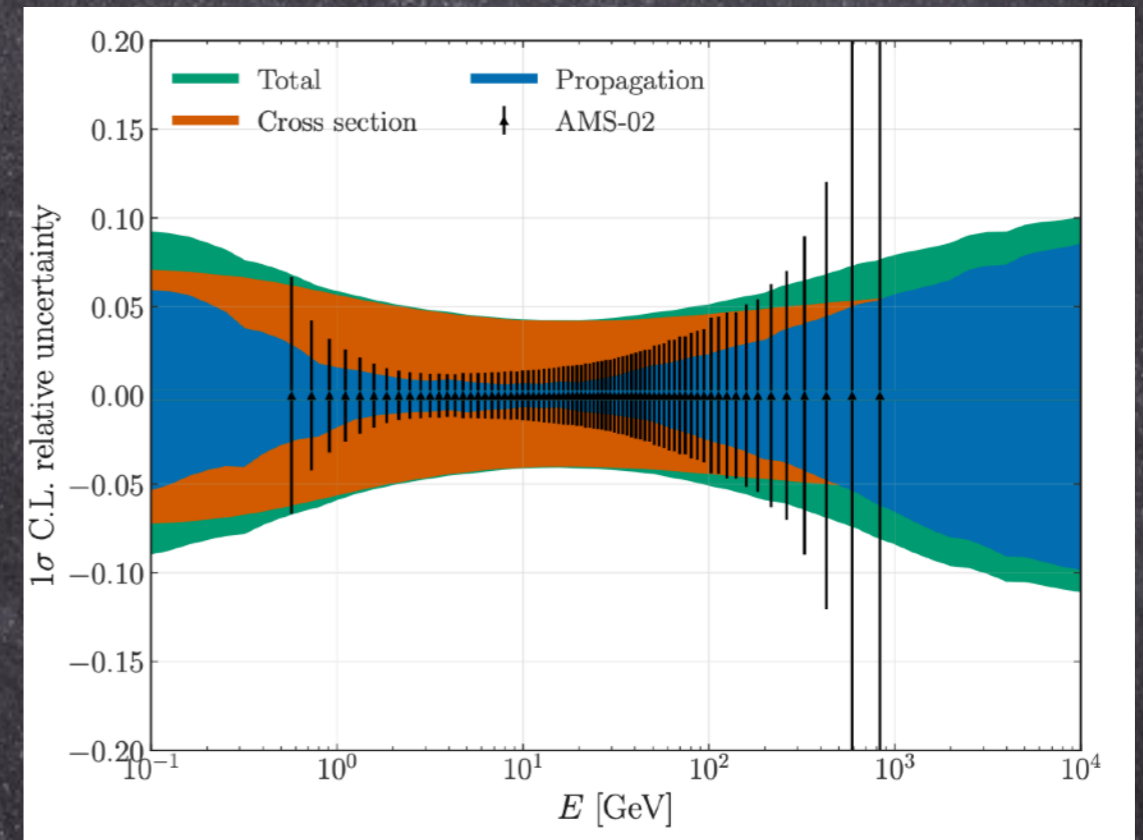
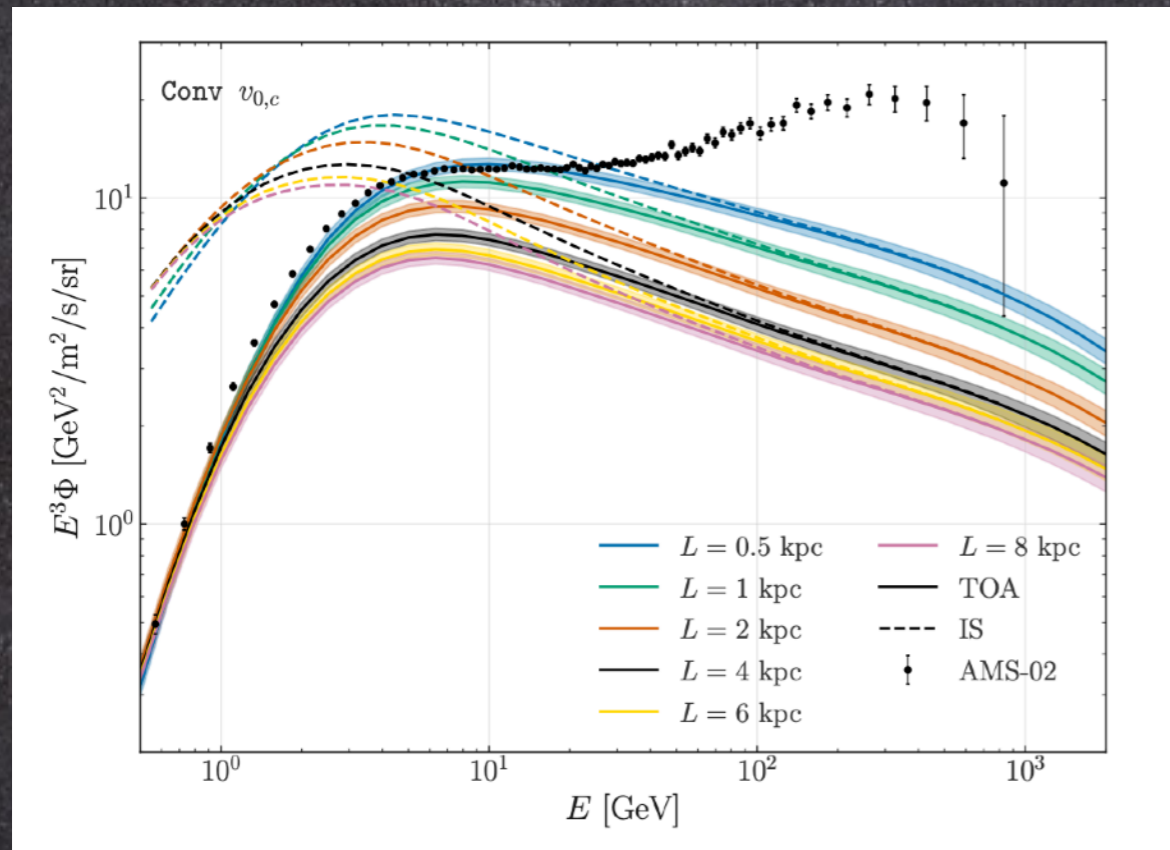


Production cross section is now known why 7-8% uncertainty above 1 GeV. Below we extrapolate.

Comparison with MonteCarlo computations is done for p-p. Similar results for e-.

The role of e^\pm secondaries

Di Mauro, FD, Korsmeier, manconi, Orusa, PRD 2023

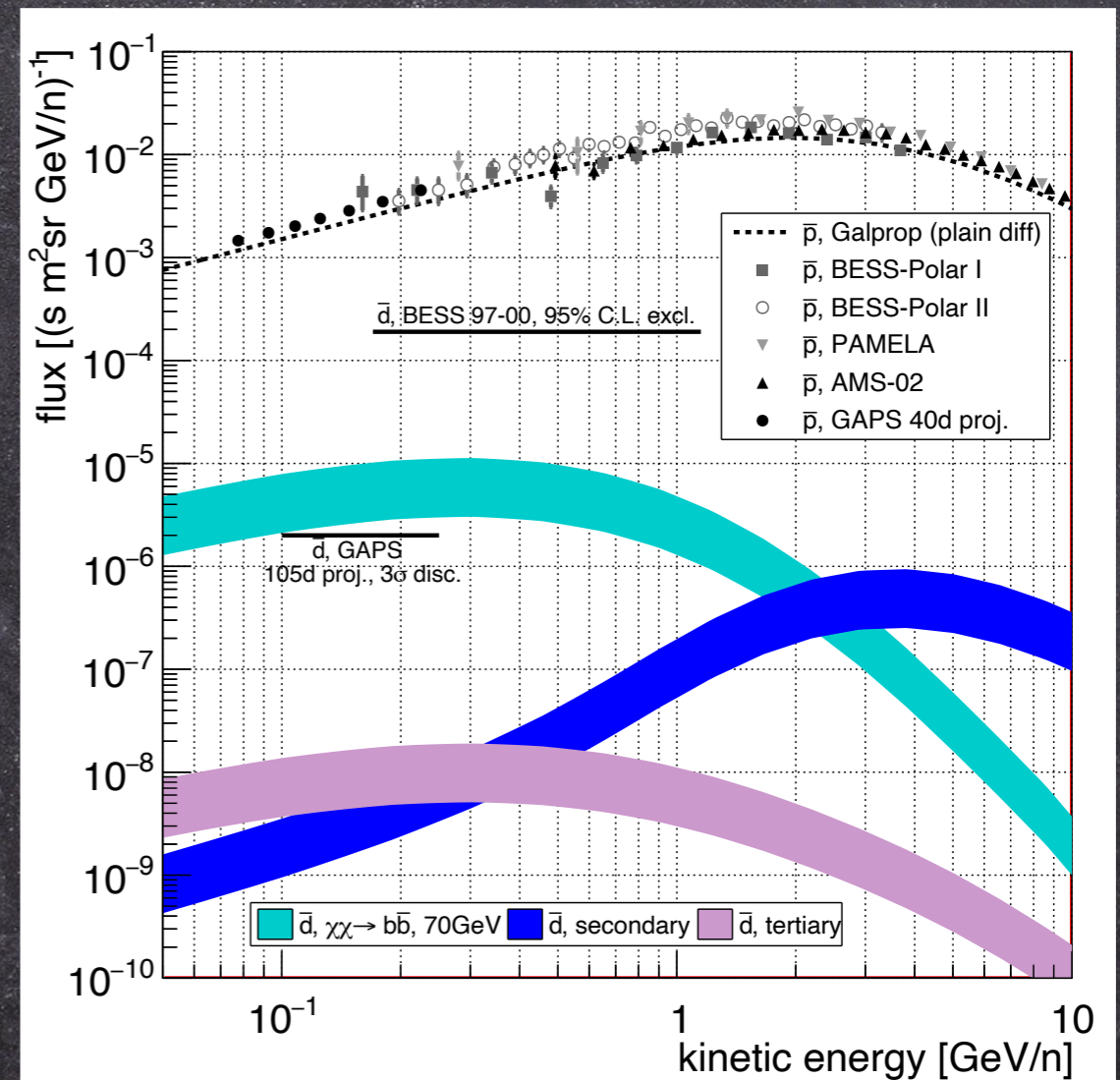
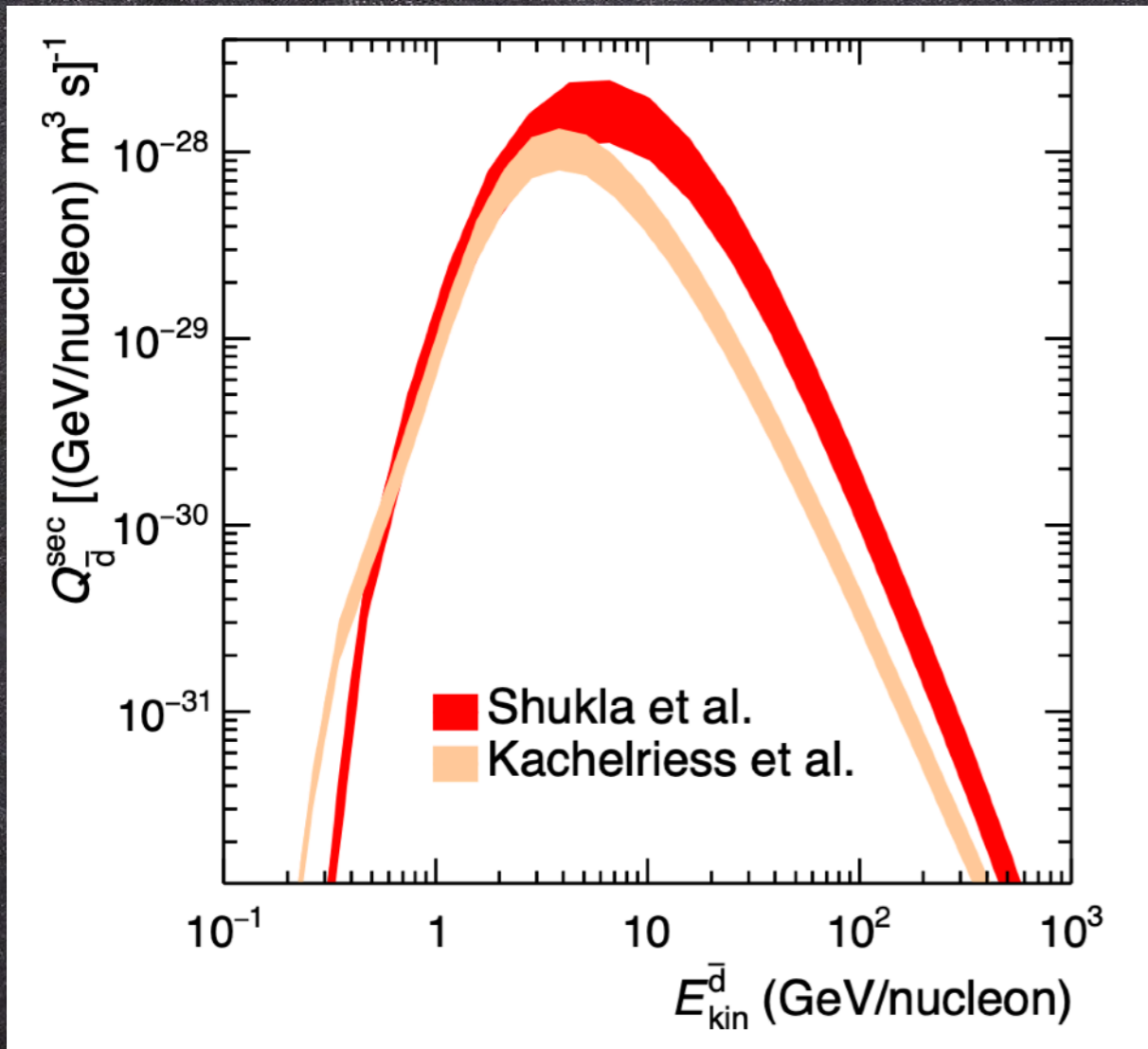


e^\pm secondaries contribute significantly to shape the spectrum at Earth below few GeV

Cross section uncertainties at the same level or greater than propagation ones ($L = 4$ kpc)

Antideuterons perspectives

Serksnyte et al, PRD 2022



Low energy window is a discovery field

Uncertainty in cross sections (left)

Uncertainties on P_c is $\pm 70\%$, P_c^3 in the flux (right)

Conclusions

Great efforts to better understand nuclei and antinuclei in CRS:
theory models, data from space, data from colliders.

Data from space are actually hampered by lack of precise (<10%) cross
section: nuclei, isotopes, antimatter, γ s

Data from colliders are highly desirable.

A specific receipt can be provided by the astroparticle community

A dedicated workshop at CERN 16-18/10/2024

<https://indico.cern.ch/event/1377509/>

Wishes' List

Partial, and personal

1. Low energy ($0.1 < T_{p\bar{p}} < 10$ GeV) antiprotons from p-p
 2. Antideuteron fusion at low energies (p beam $\sim 10-10^2$ GeV)
 3. $p+\text{He} \rightarrow e^++X$ ($p+\text{He} \rightarrow \pi^++X$)
 4. $^{12}\text{C}+p \rightarrow \text{LiBeB}$ fragments with isotopes
- + many more!

Analytical formulae for $e\pm$ production XS

L. Orusa, M. Di Mauro, FD, M. Korsmeier PRD 2022

The procedure is fully data driven

$$\sigma_{\text{inv}} = \sigma_0(s) c_1 \left[F_p(s, p_T, x_R) + F_r(p_T, x_R) \right] A(s),$$

$$F_p(s, p_T, x_R) = (1 - x_R)^{c_2} \exp(-c_3 x_R) p_T^{c_4} \quad (8)$$
$$\times \exp \left[-c_5 \sqrt{s/s_0}^{c_6} \left(\sqrt{p_T^2 + m_\pi^2} - m_\pi \right)^{c_7 \sqrt{s/s_0}^{c_6}} \right],$$

$$F_r(p_T, x_R) = (1 - x_R)^{c_8}$$
$$\times \exp \left[-c_9 p_T - \left(\frac{|p_T - c_{10}|}{c_{11}} \right)^{c_{12}} \right]$$
$$\times \left[c_{13} \exp(-c_{14} p_T^{c_{15}} x_R) + \right.$$
$$\left. + c_{16} \exp \left(- \left(\frac{|x_R - c_{17}|}{c_{18}} \right)^{c_{19}} \right) \right]$$

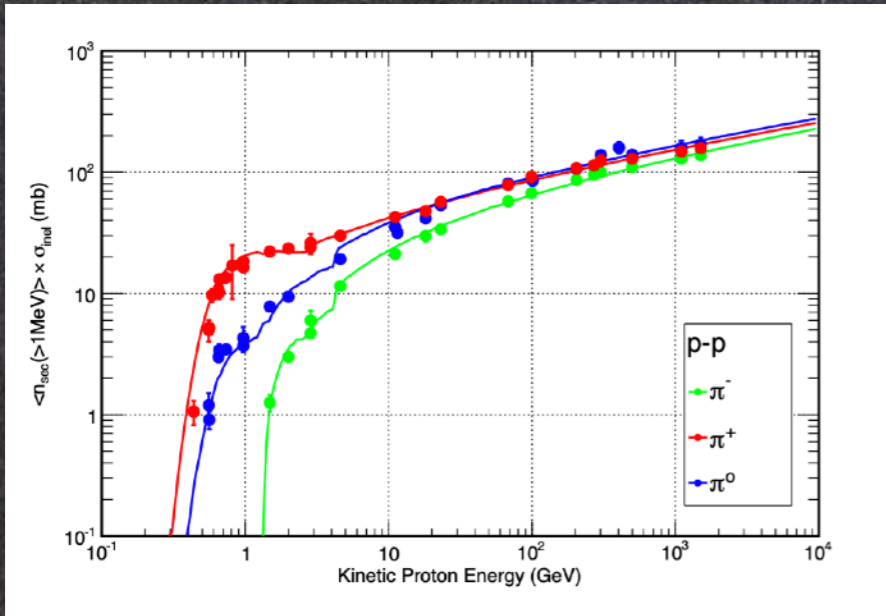
$$A(s) = \frac{1 + \left(\sqrt{s/c_{20}} \right)^{c_{21} - c_{22}}}{1 + \left(\sqrt{s_0/c_{20}} \right)^{c_{21} - c_{22}}} \left(\sqrt{\frac{s}{s_0}} \right)^{c_{22}}$$

F_s and F_r mainly driven by NA49 data

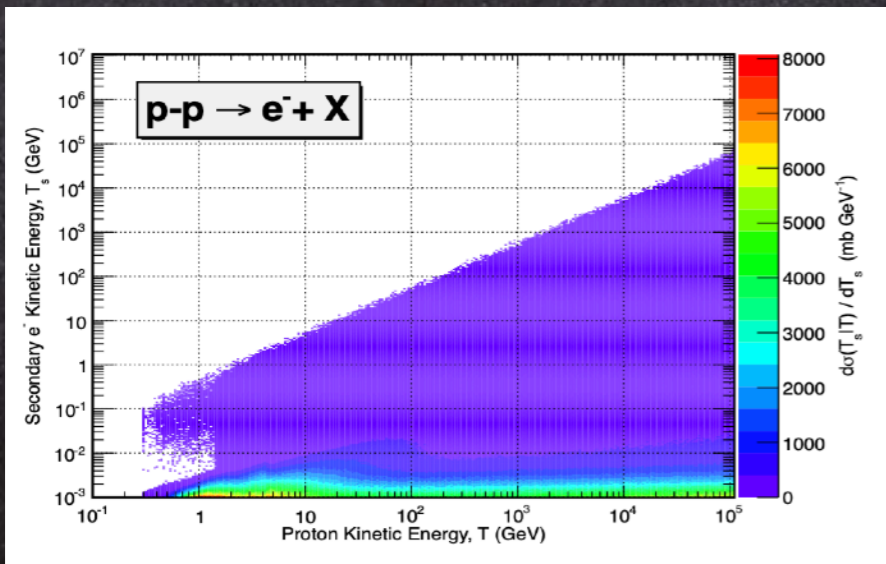
High energy behavior $A(s)$ tested on CMS and ALICE data

Fluka MC generator

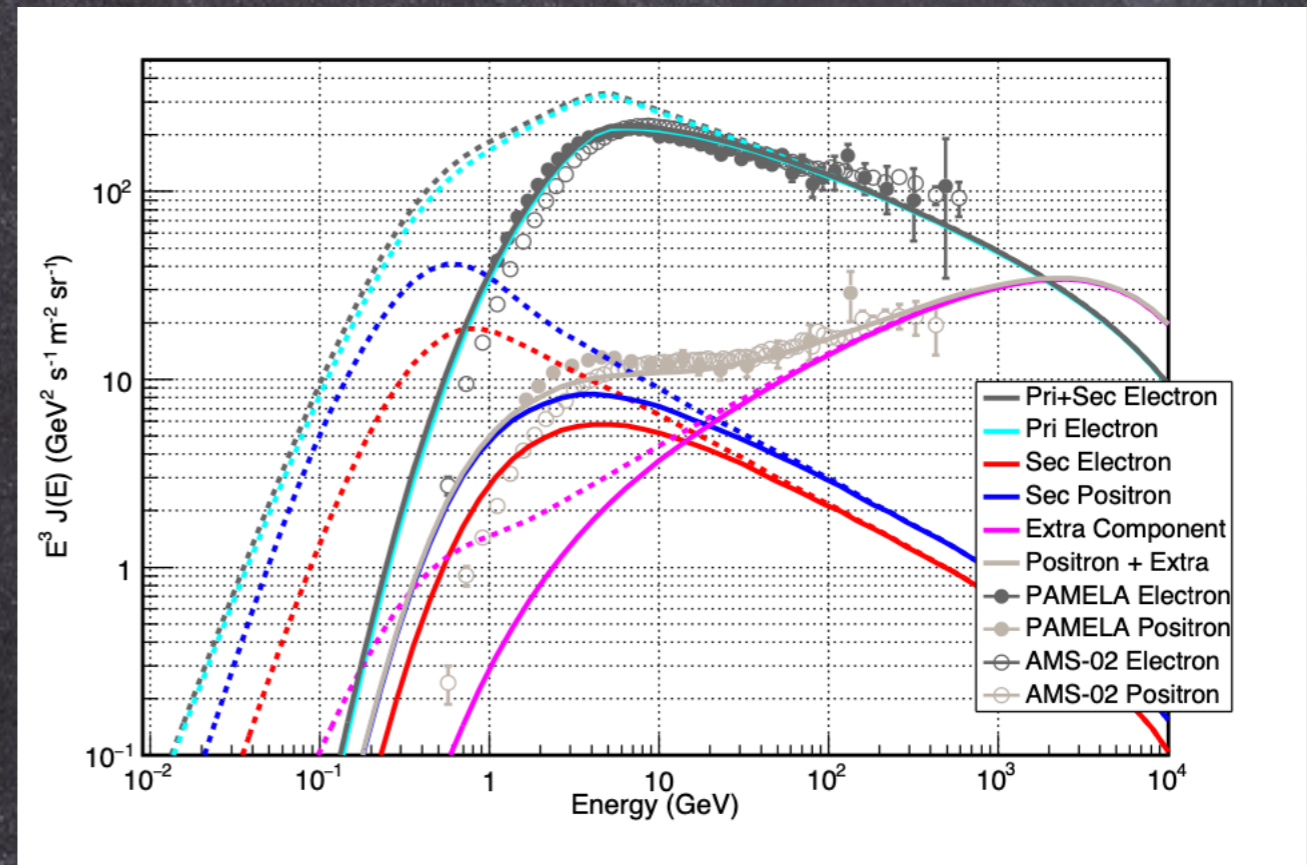
N. Mazziotta+, AP 2017



Points are from Dermer 1986



T_e is severely degraded from Projectile energy



Propagated e^+ and e^- w.r.t. data

Comparison with Monte Carlo generators

Koldobskiy et al., PRD 2021, 2110.00496

Results with Aafrag

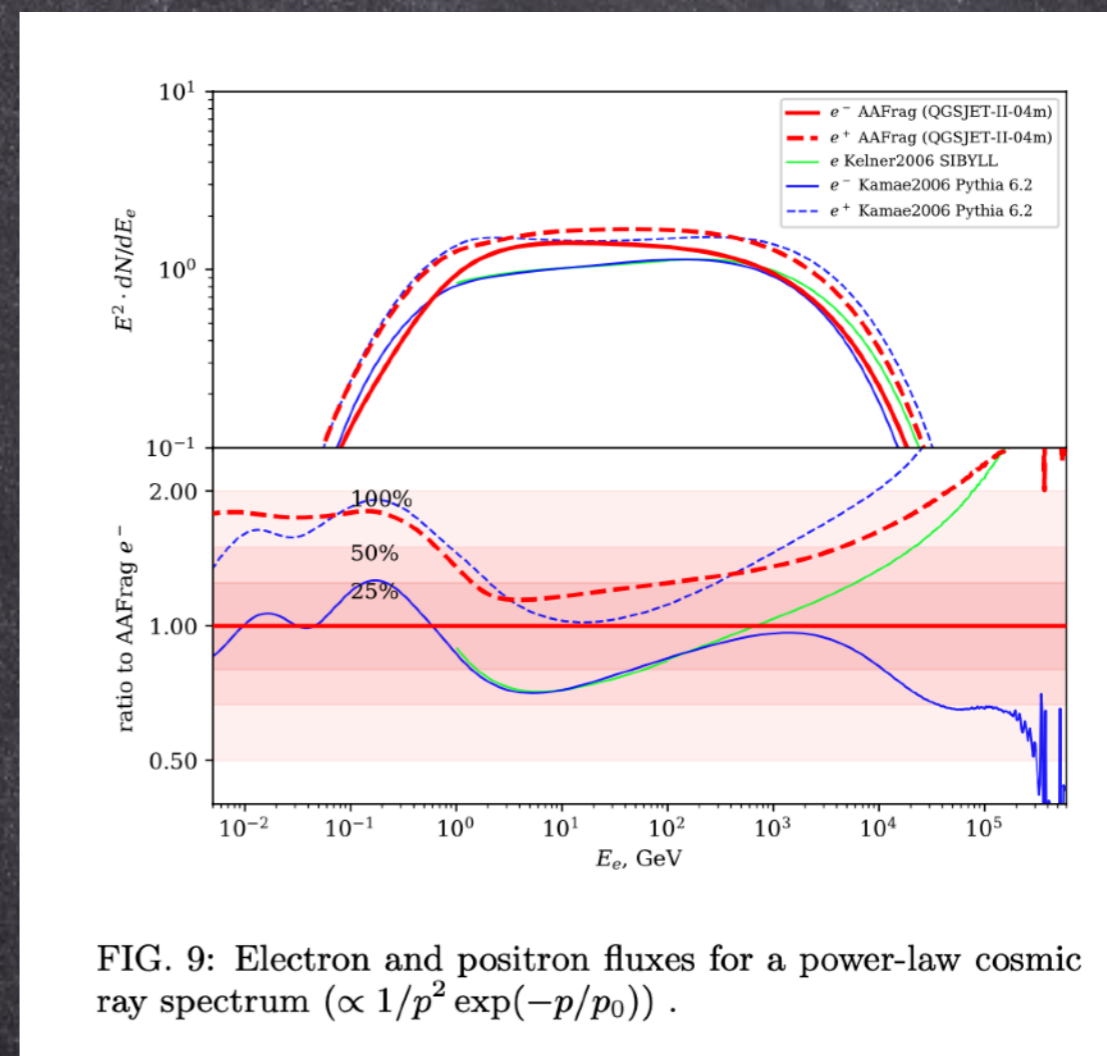
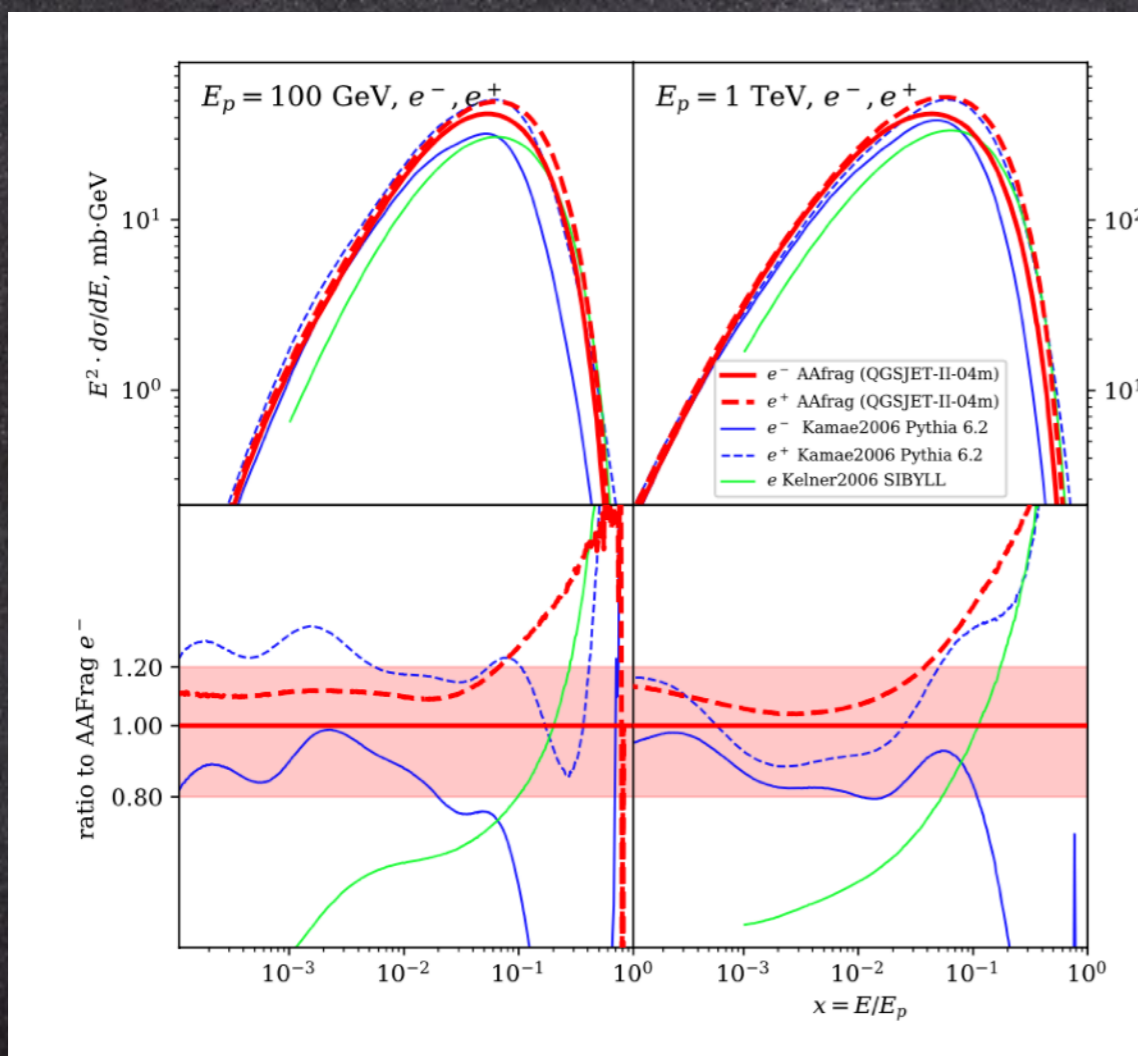


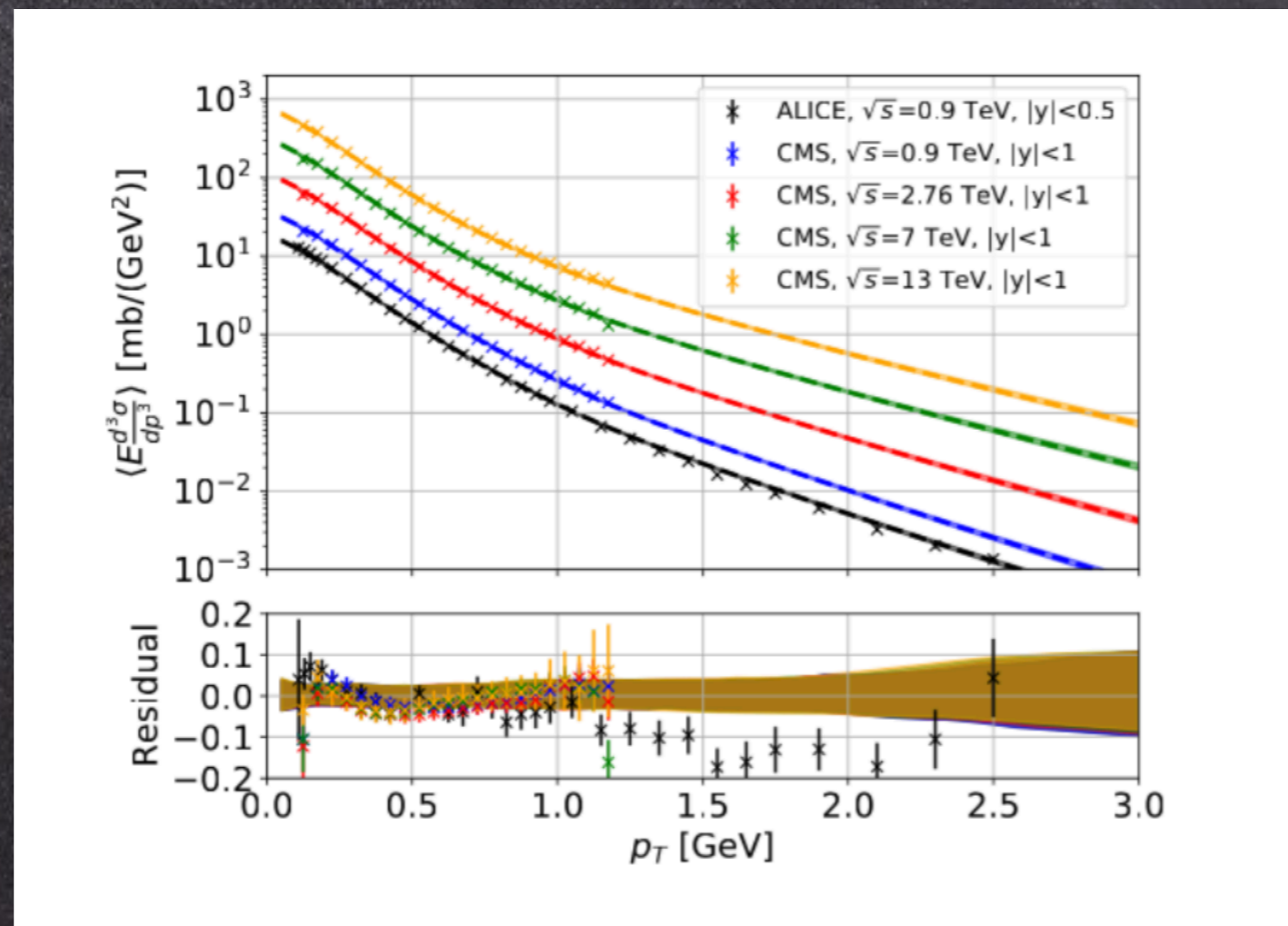
FIG. 9: Electron and positron fluxes for a power-law cosmic ray spectrum ($\propto 1/p^2 \exp(-p/p_0)$).

Different MC modelings lead to considerable differences in the Production cross section, and consequently on the source spectrum

Results at large \sqrt{s}

L. Orusa, M. Di Mauro, FD, M. Korsmeier PRD 2022

We use σ_{inv} or multiplicity



Uncertainties between 5% and 10% – most relevant is 5% at low p_T

Antimatter or γ -rays sources from DARK MATTER

Annihilation

$$Q_{\text{ann}}(\vec{x}, E) = \epsilon \left(\frac{\rho(\vec{x})}{m_{\text{DM}}} \right)^2 \sum_f \langle \sigma v \rangle_f \frac{dN_{e^\pm}^f}{dE}$$

Decay

$$Q_{\text{dec}}(\vec{x}, E) = \left(\frac{\rho(\vec{x})}{m_{\text{DM}}} \right) \sum_f \Gamma_f \frac{dN_{e^\pm}^f}{dE}$$

- ρ DM density in the halo of the MW
- m_{DM} DM mass
- $\langle \sigma v \rangle$ thermally averaged annihilation cross section in SM channel f
- Γ DM decay time
- e^+ , e^- energy spectrum generated in a single annihilation or decay event

Annihilations take place in the whole diffusive halo