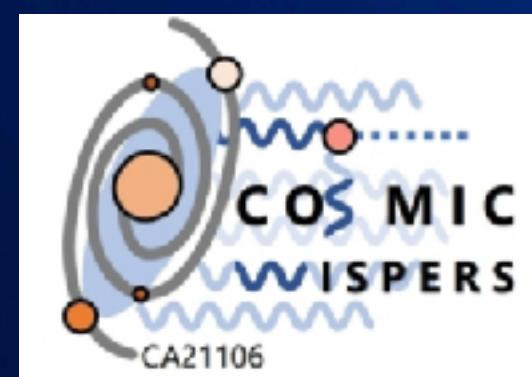


Diffuse Axion Background

Joshua Eby
Oskar Klein Centre
Stockholm University

Athens symposium on Exploring the Universe
ATHEXIS
2024/06/10



Based on

Eby, Shirai, Stadnik, Takhistov (2106.14893)
Arakawa, **Eby**, Safronova, Takhistov, Zaheer (2306.16468)
Arakawa, Zaheer, **Eby**, Takhistov, Safronova (2402.06736)
Eby, Takhistov (2402.00100)



Stockholms
universitet

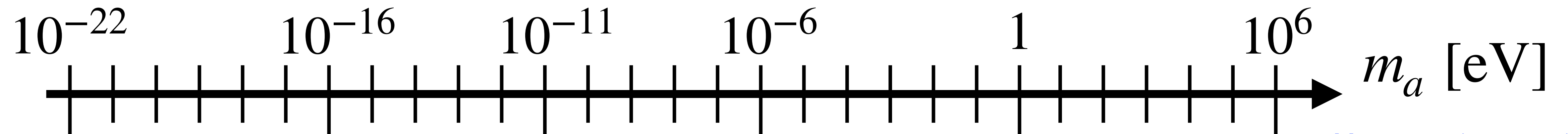
Axions / Axion-like Particles*

*in this talk, same thing

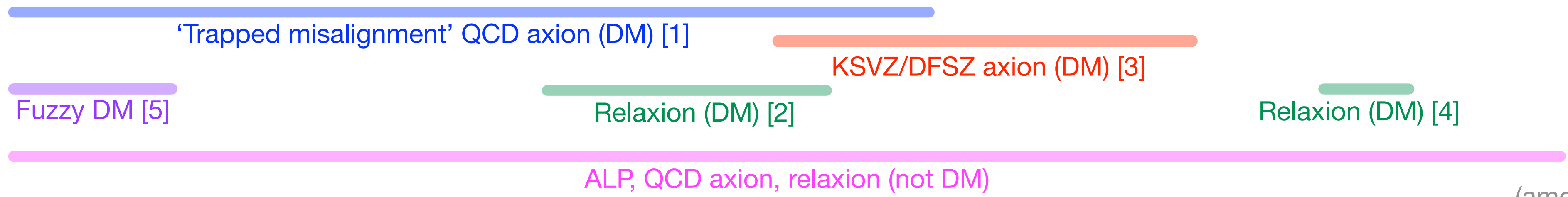
(1) Light or ultralight

(2) Scalar or pseudoscalar

(3) Dark matter or not



- [1] Di Luzio++ (2102.01082)
- [2] Banerjee++ (1810.01889)
- [3] di Cortona++ (1511.02867)
(recent review)
- [4] Fonseca++(1809.04534)
- [5] Hu++ (astro-ph/0003365)



(among many others...)

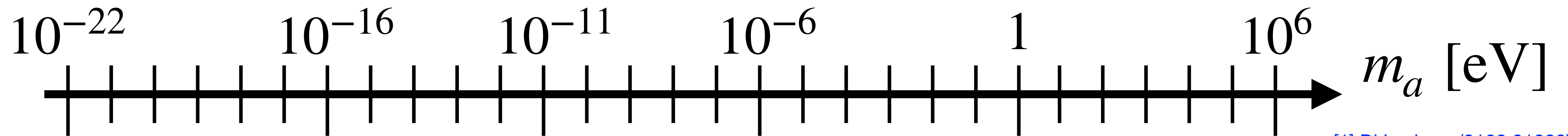
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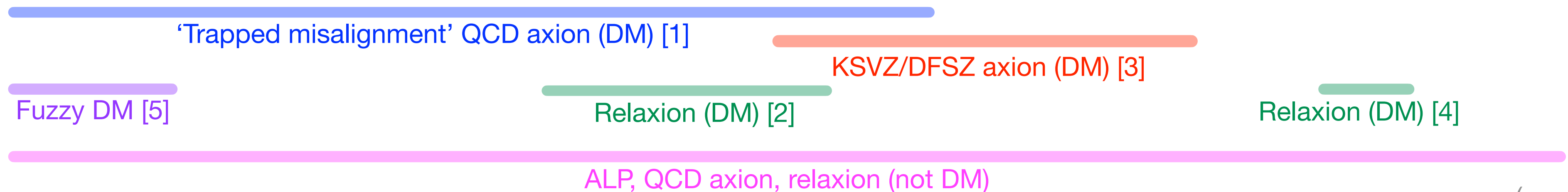
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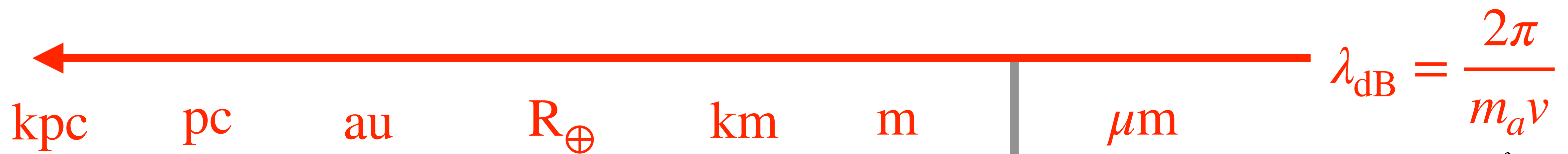
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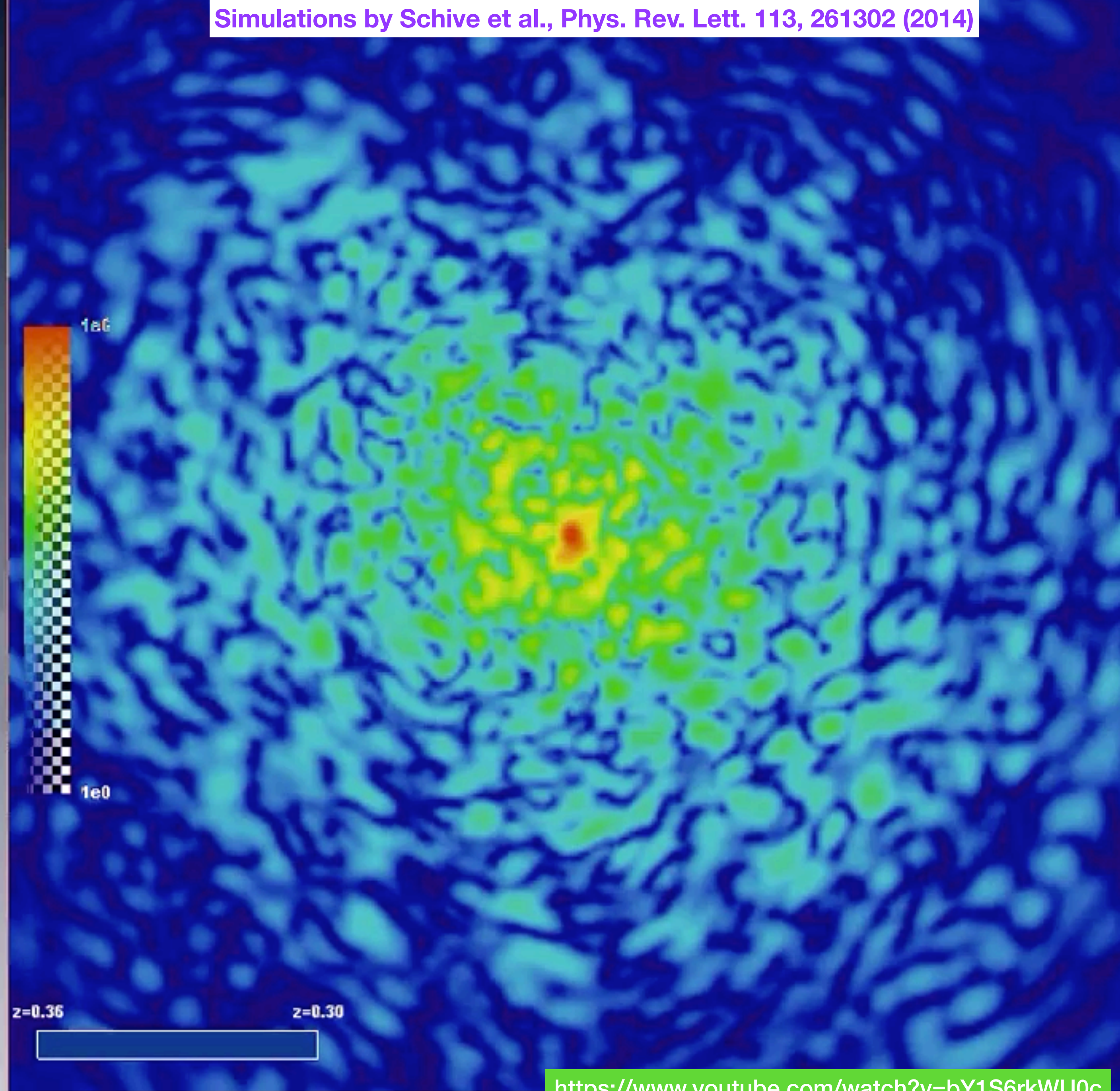


Solutions for the axion field

$$a(\vec{x}, t) \simeq a_0 \cos \left(m_a t + m_a \vec{v}_{\text{dm}} \cdot \vec{x} + \delta \right)$$

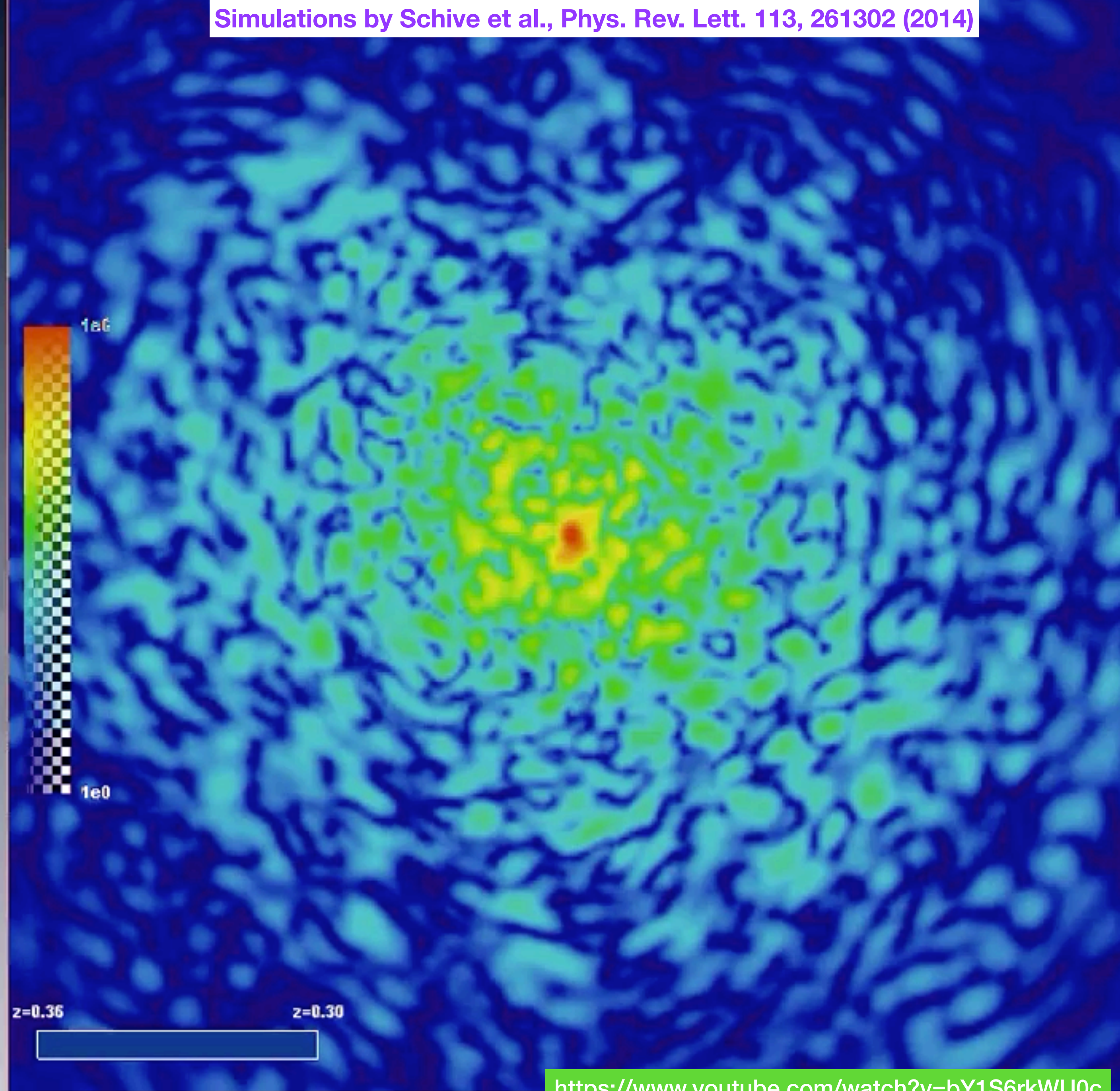
amplitude $a_0 \simeq \sqrt{2\rho_{\text{dm}}}/m_a$ DM velocity arbitrary phase

Density field
 $\rho(\vec{x}, t) \simeq m_a^2 a(\vec{x}, t)^2$



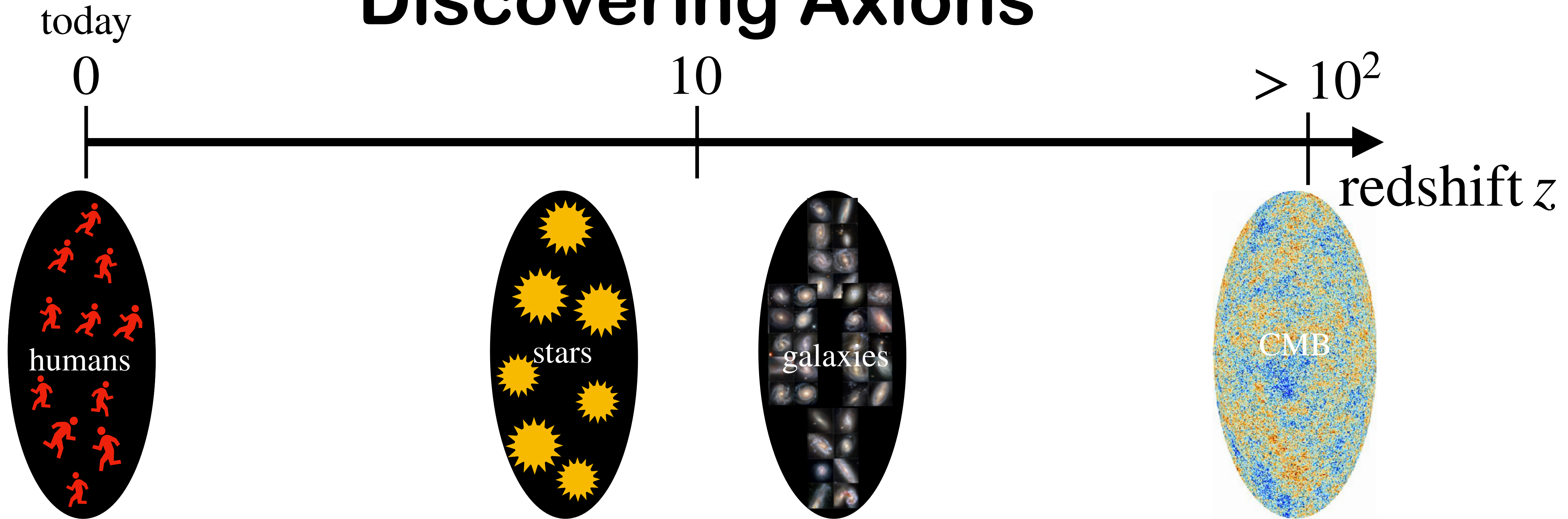
<https://www.youtube.com/watch?v=bY1S6rkWU0c>

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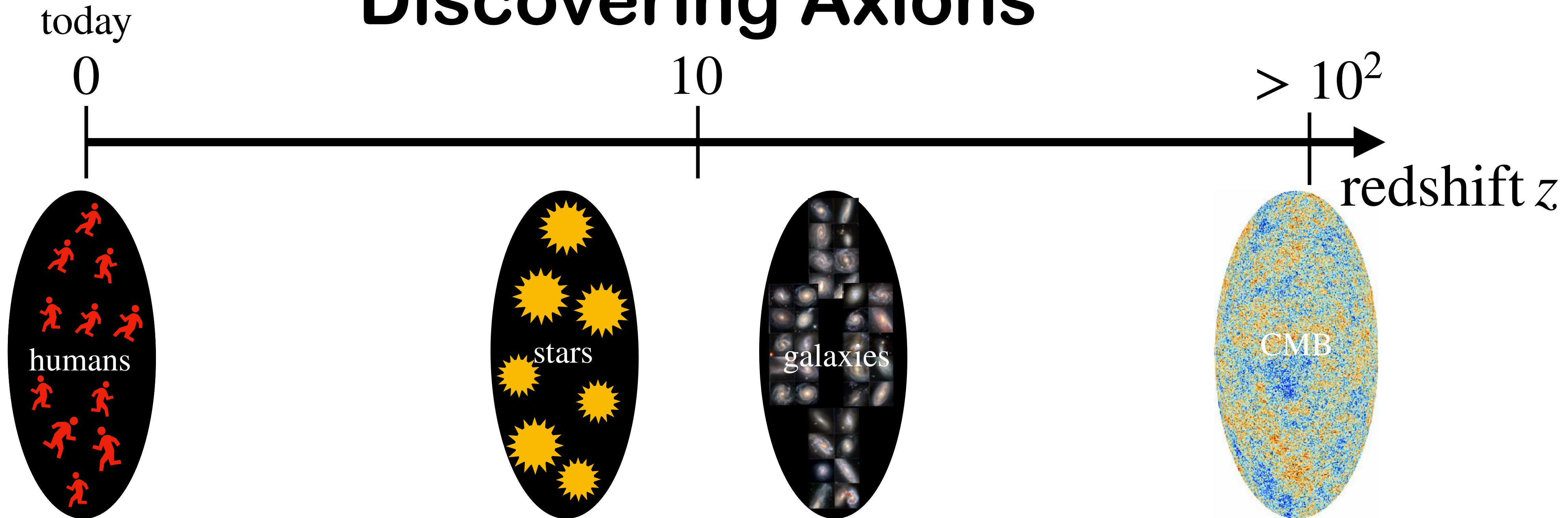


<https://www.youtube.com/watch?v=bY1S6rkWU0c>

Discovering Axions



Discovering Axions



traditional DM searches

$$z \sim 0$$

- cold, $v_{\text{dm}} \sim 10^{-3}c$
- direct detection: local DM with $\rho_{\text{dm}} \simeq 0.4 \text{ GeV}/\text{cm}^3$
- indirect detection: annihilation flux from e.g. galactic center

Discovering Axions

today

0

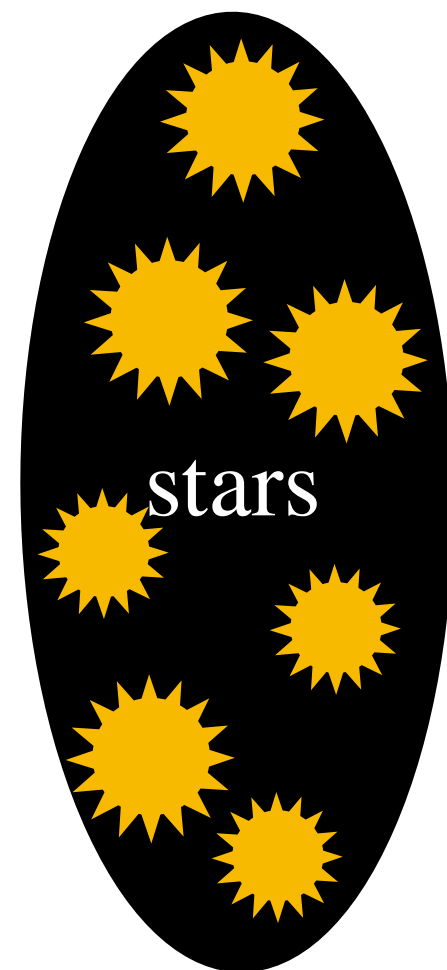
10

$> 10^2$

redshift z



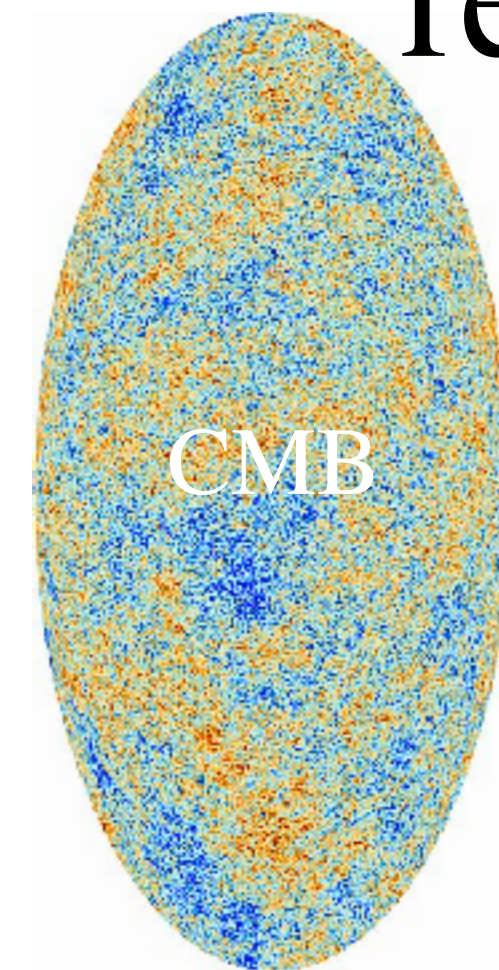
humans



stars



galaxies



CMB

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cosmic axion background

$$z \gg 30$$

- “Hot”, $v \sim c$
- Relativistic population of axions from cosmological sources

Conlon and Marsh (1304.1804, 1305.3603)
Dror, Murayama, Rodd (2101.09287)

Discovering Axions

today

0

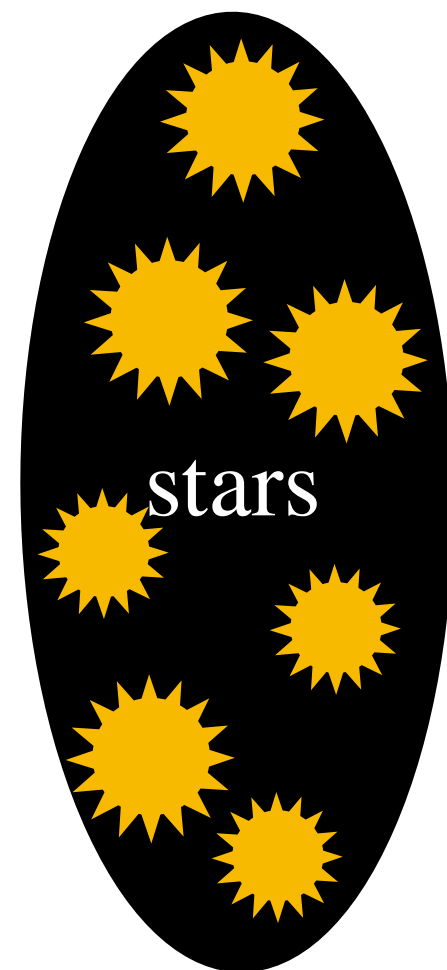
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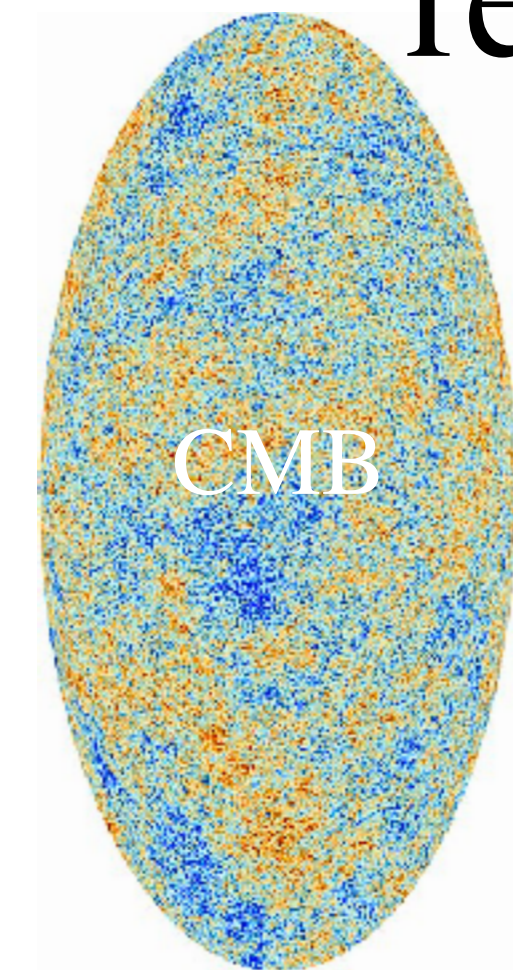
humans



stars



galaxies



CMB

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- indirect detection: annihilation flux from e.g. galactic center

transient searches

$z \sim 0 - \text{few}$

- Relativistic burst passes Earth, leaving detectable signal
- Galactic or extragalactic

Eby, Takhistov,
with Shirai, Stadnik (2106.14893)
with Arakawa, Safronova, Zaheer,
(2306.16468, 2402.06736)

diffuse axion background

$z \sim \text{few} - 30$

- build-up of large population of relativistic axions originating in astrophysical bursts
- direct and indirect signals

Eby, Takhistov (2402.00100)

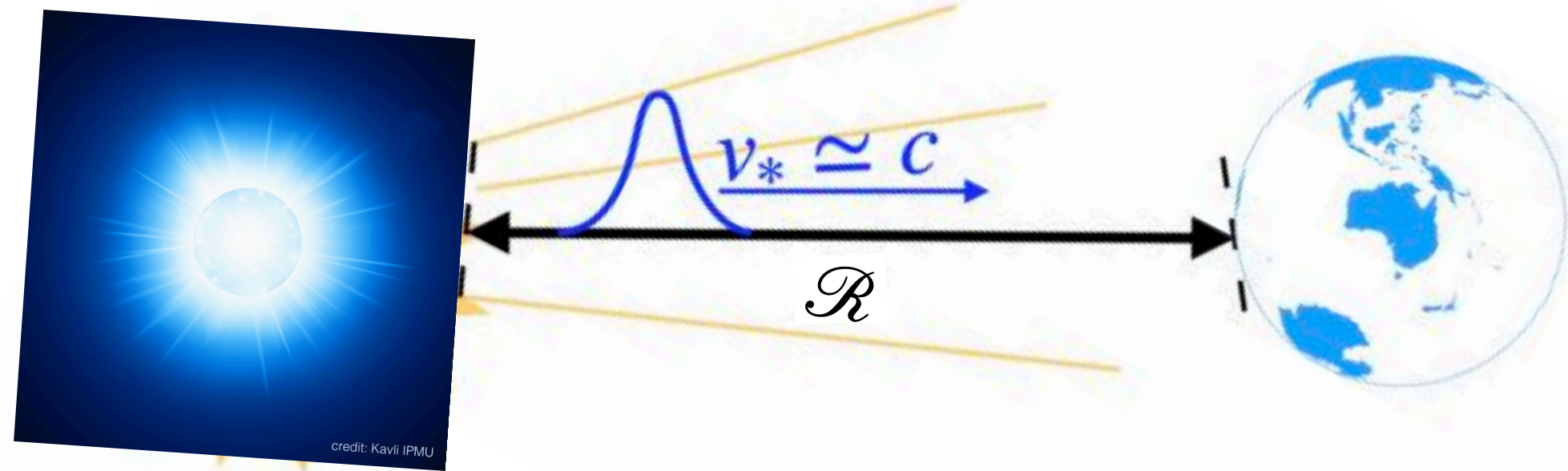
cosmic axion background

$z \gg 30$

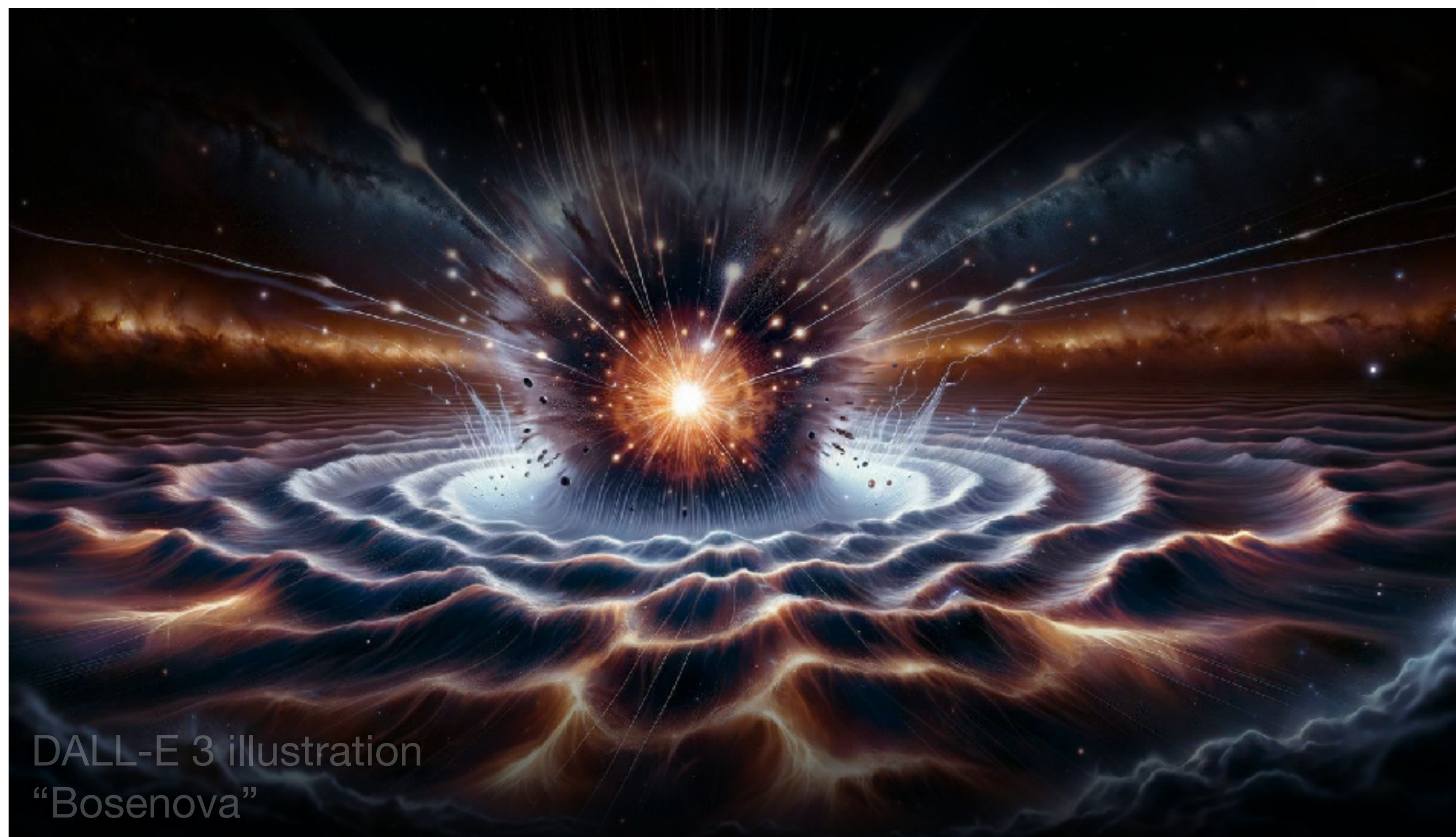
- “Hot”, $v \sim c$
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Conlon and Marsh (1304.1804, 1305.3603)
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Outline



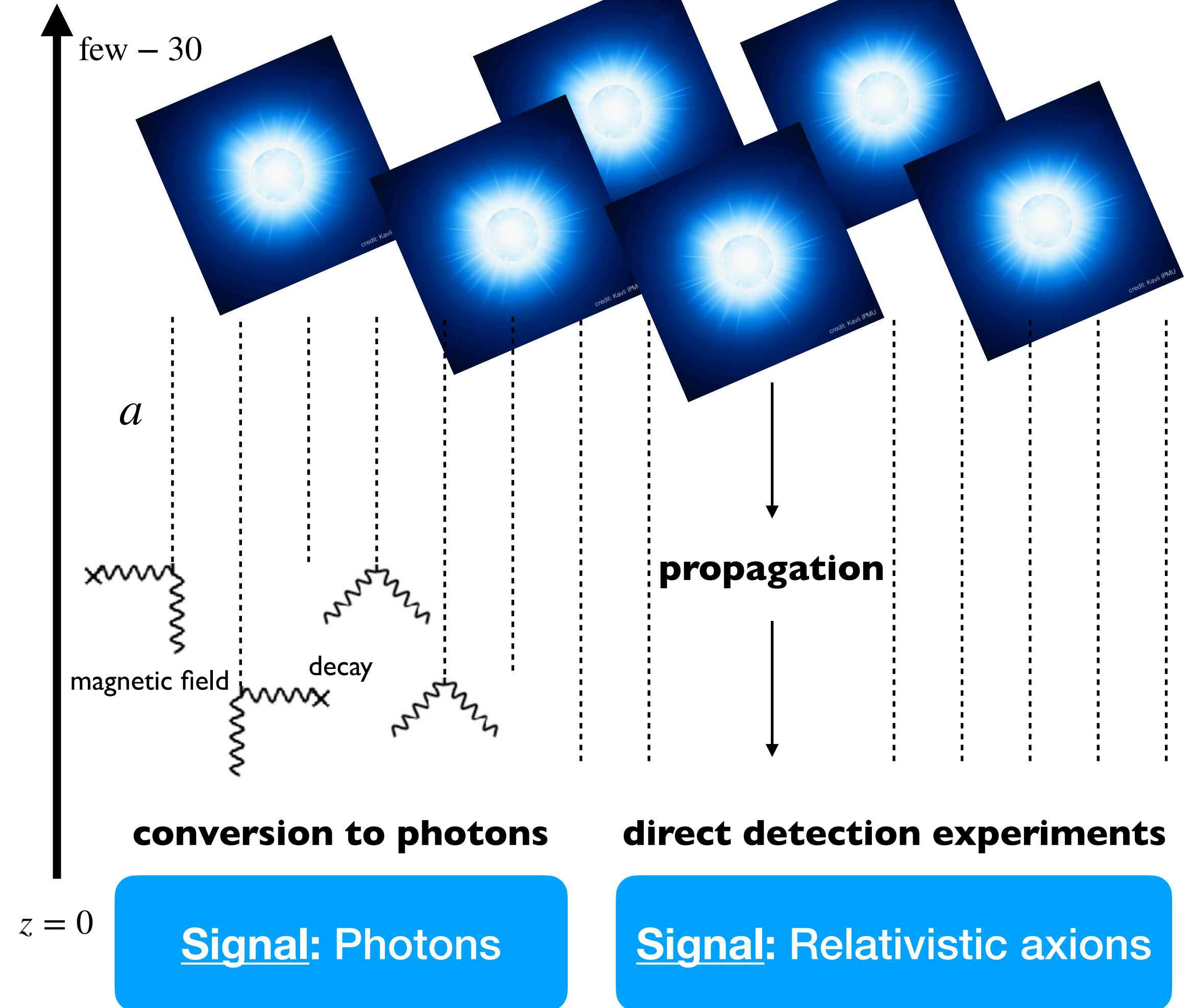
1. Direct searches for transient signals from relativistic axion bursts



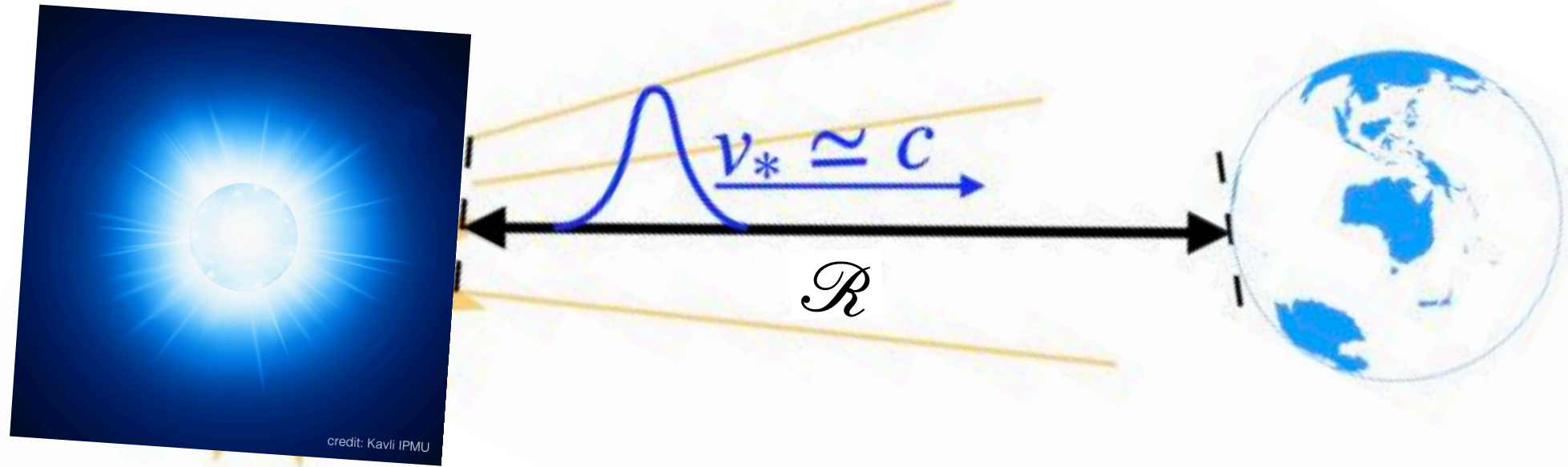
DALL-E 3 illustration
"Bosenova"

redshift z

relativistic 'bursts' of axions



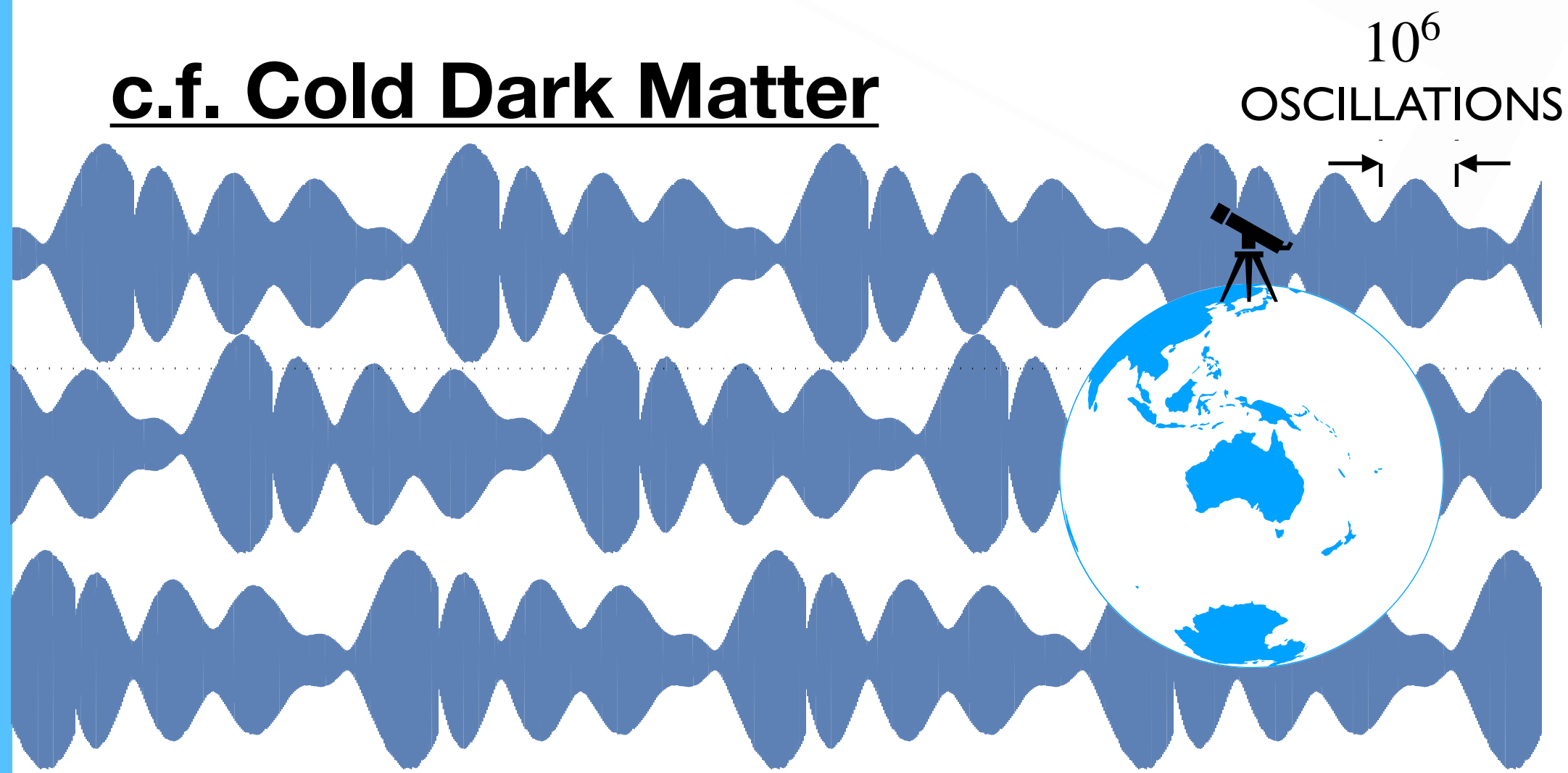
2. Diffuse Axion Background



1. Direct searches for transient signals from relativistic axion bursts

Key takeaway: DM search experiments can discover axions by detecting bursts

c.f. Cold Dark Matter



Local DM density:

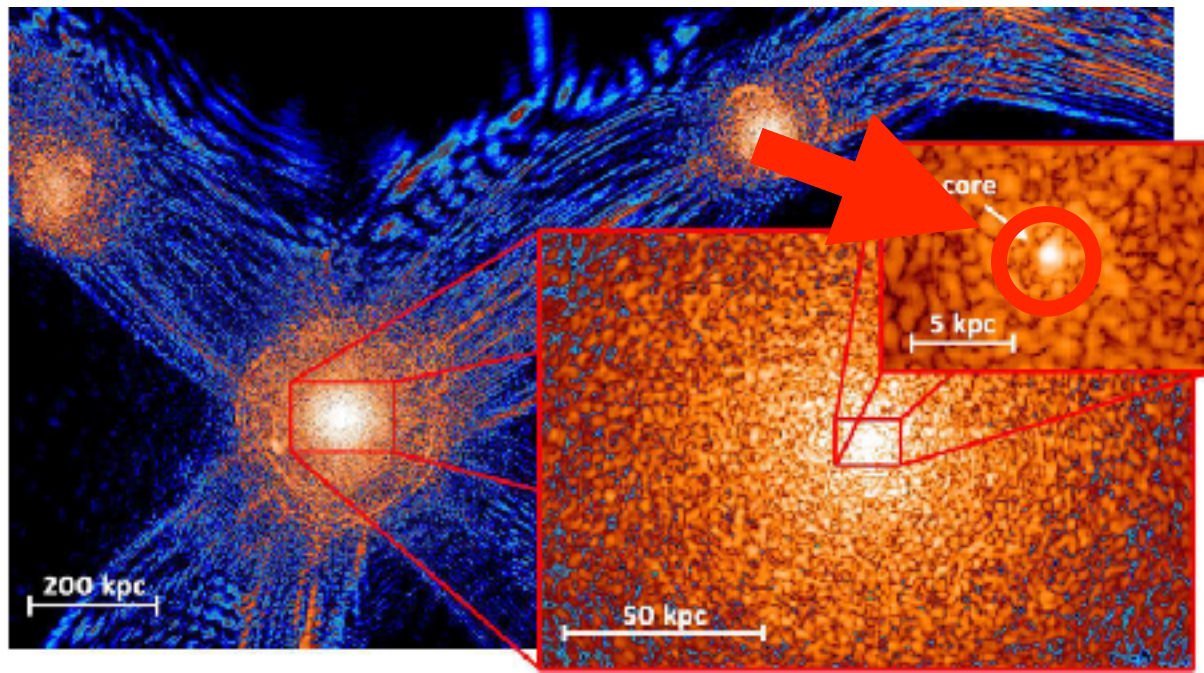
$$\rho_{\text{dm}} \simeq 0.4 \text{ GeV/cm}^3$$

Ultralight scalars are

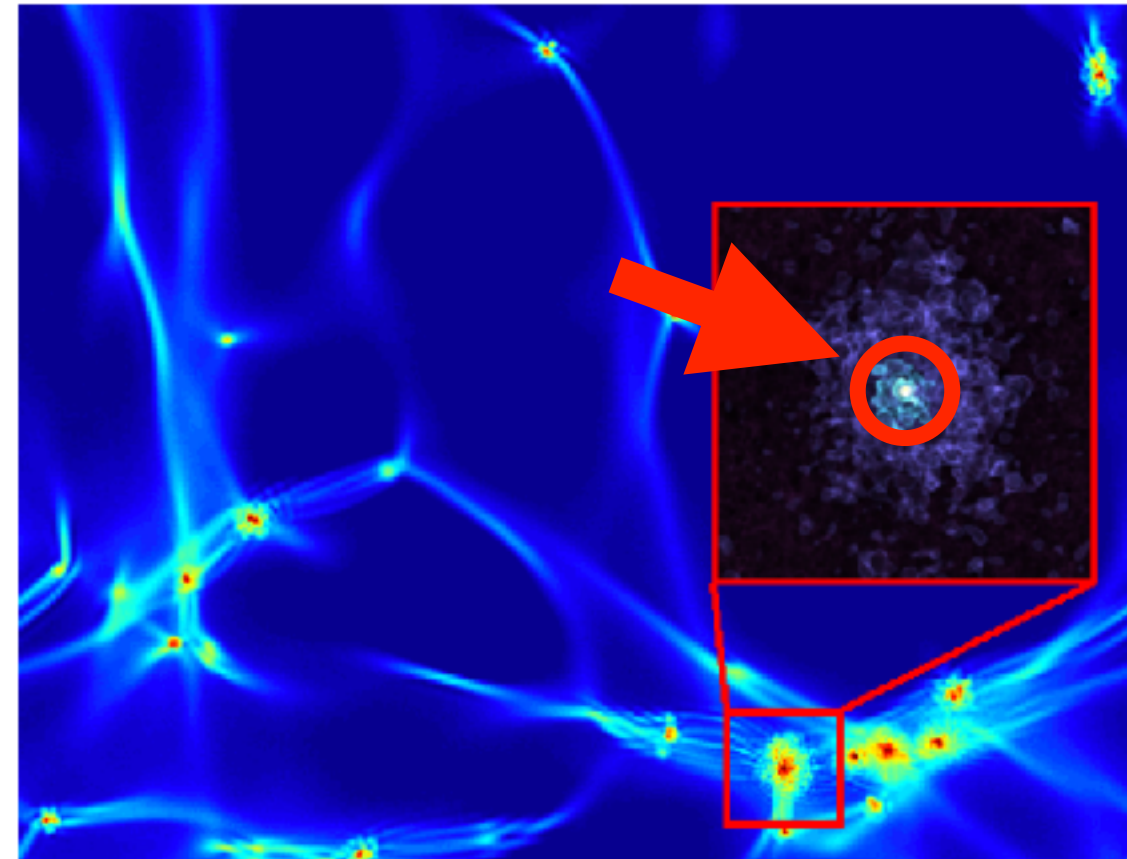
high 'quality' oscillators: $v_{\text{dm}} \simeq 10^{-3}c$

$$\tau_{\text{dm}} \sim 2\pi v_{\text{dm}}^{-2} m_a^{-1} \sim 10^6 / m_a$$

Burst Source of Interest: Boson Stars

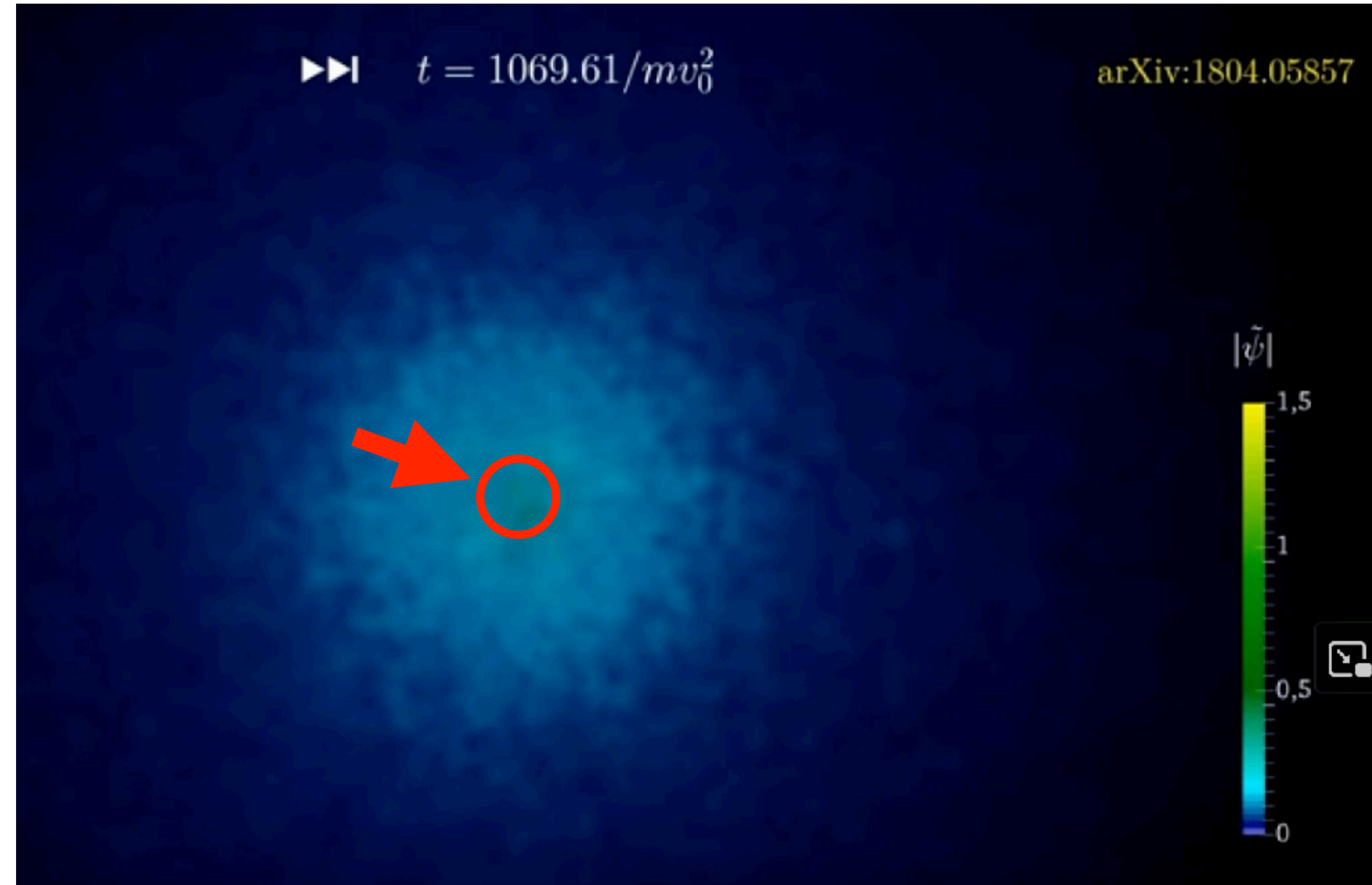


Schive++ (1406.6586)

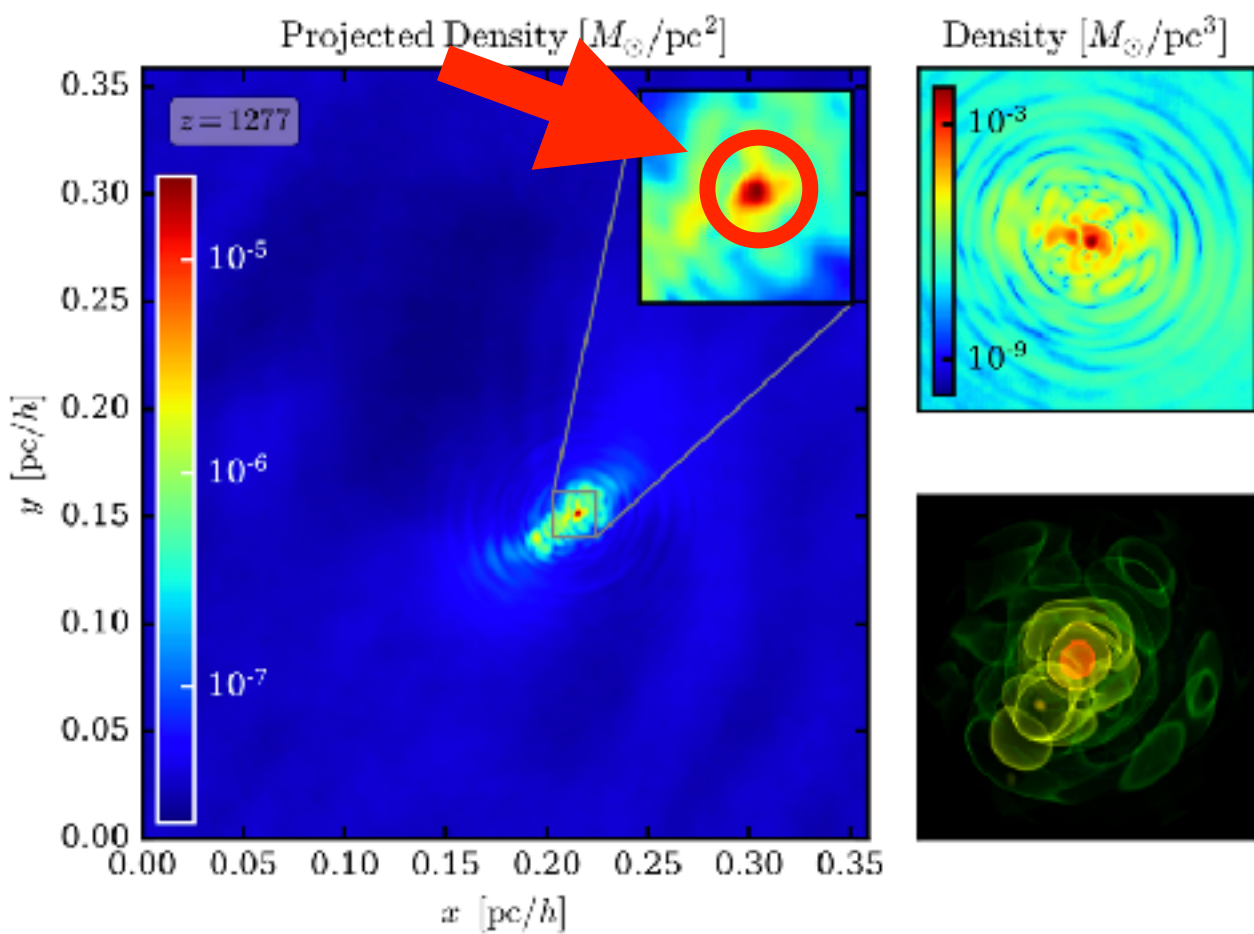


Mocz++ (1705.05845)

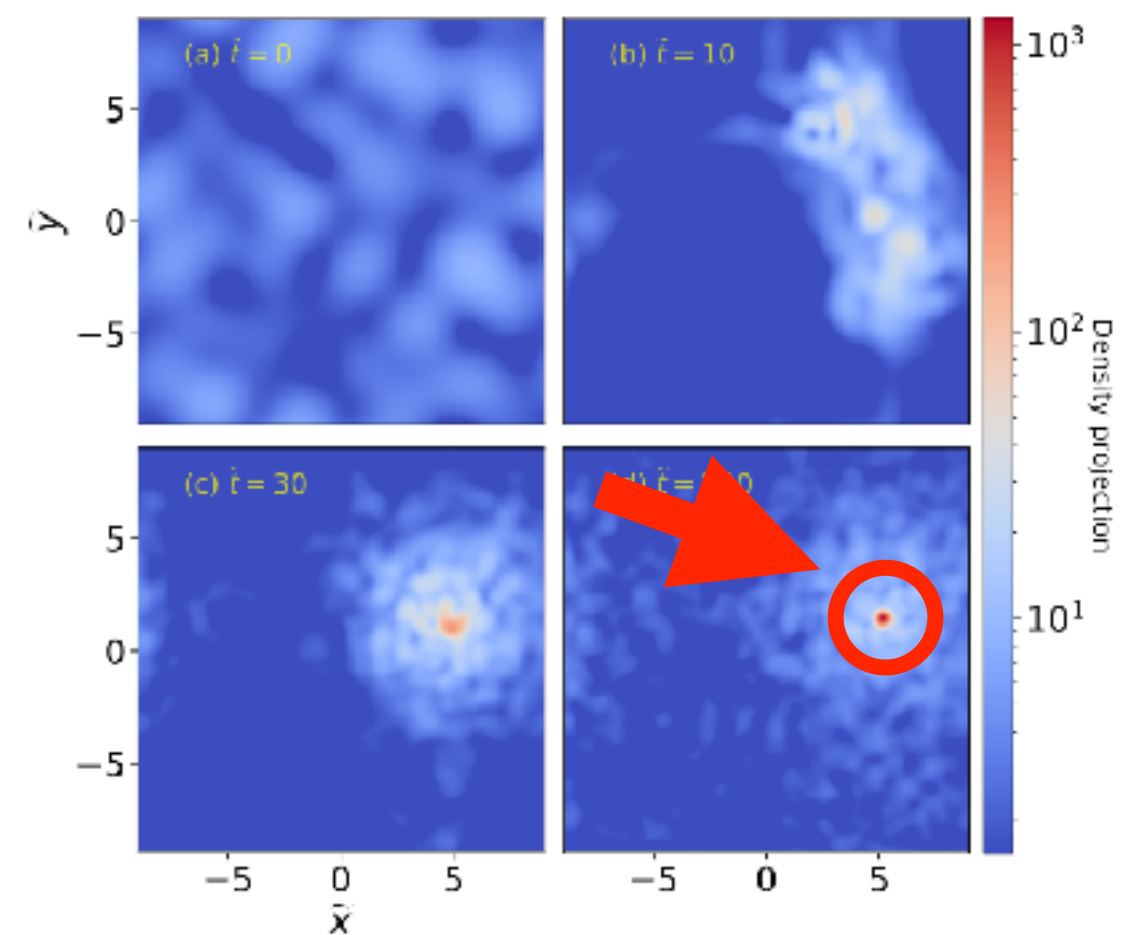
These dense configurations form
in astrophysical DM halos



Levkov, Panin, Tkachev (1804.05857)
Video via Alexander Panin on YouTube



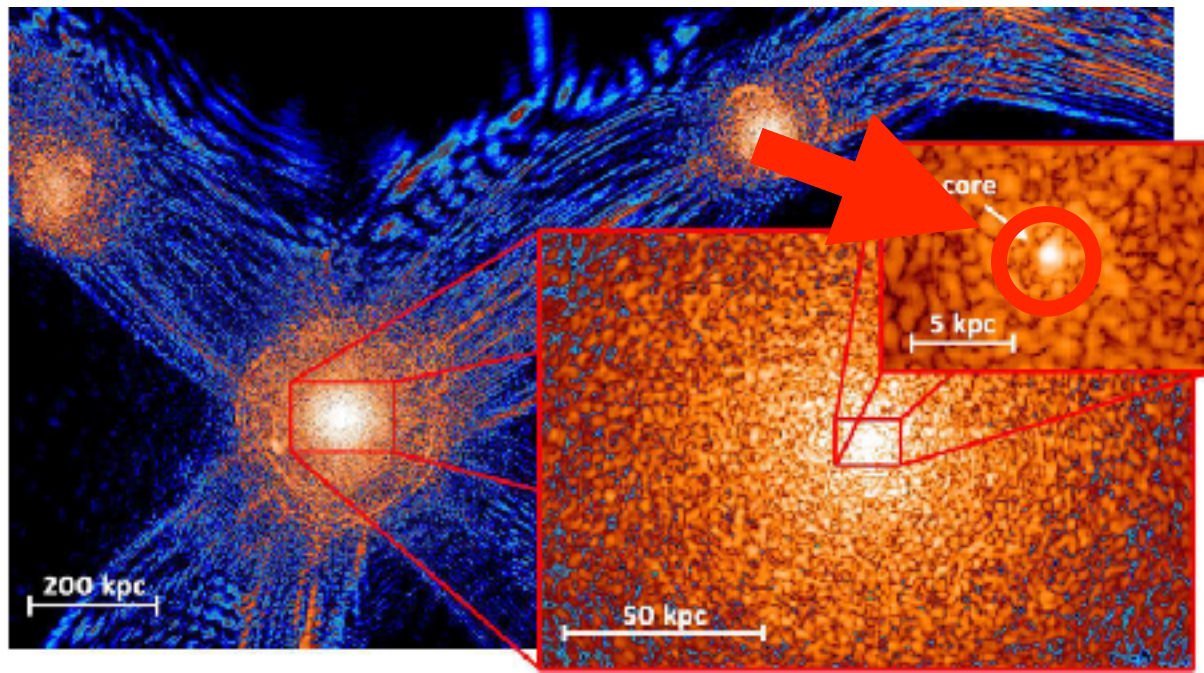
Eggemeier and Niemeyer (1906.01348)



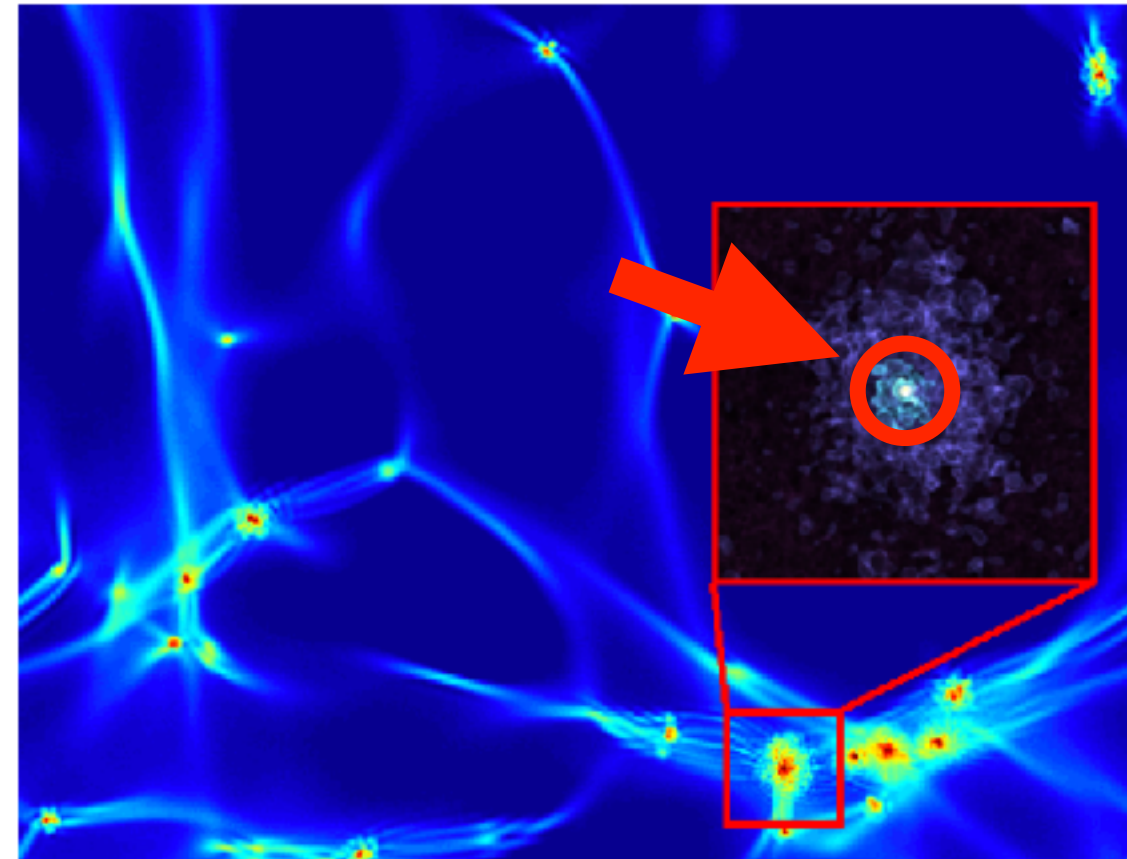
Chen++ (2011.01333)

(among others!)

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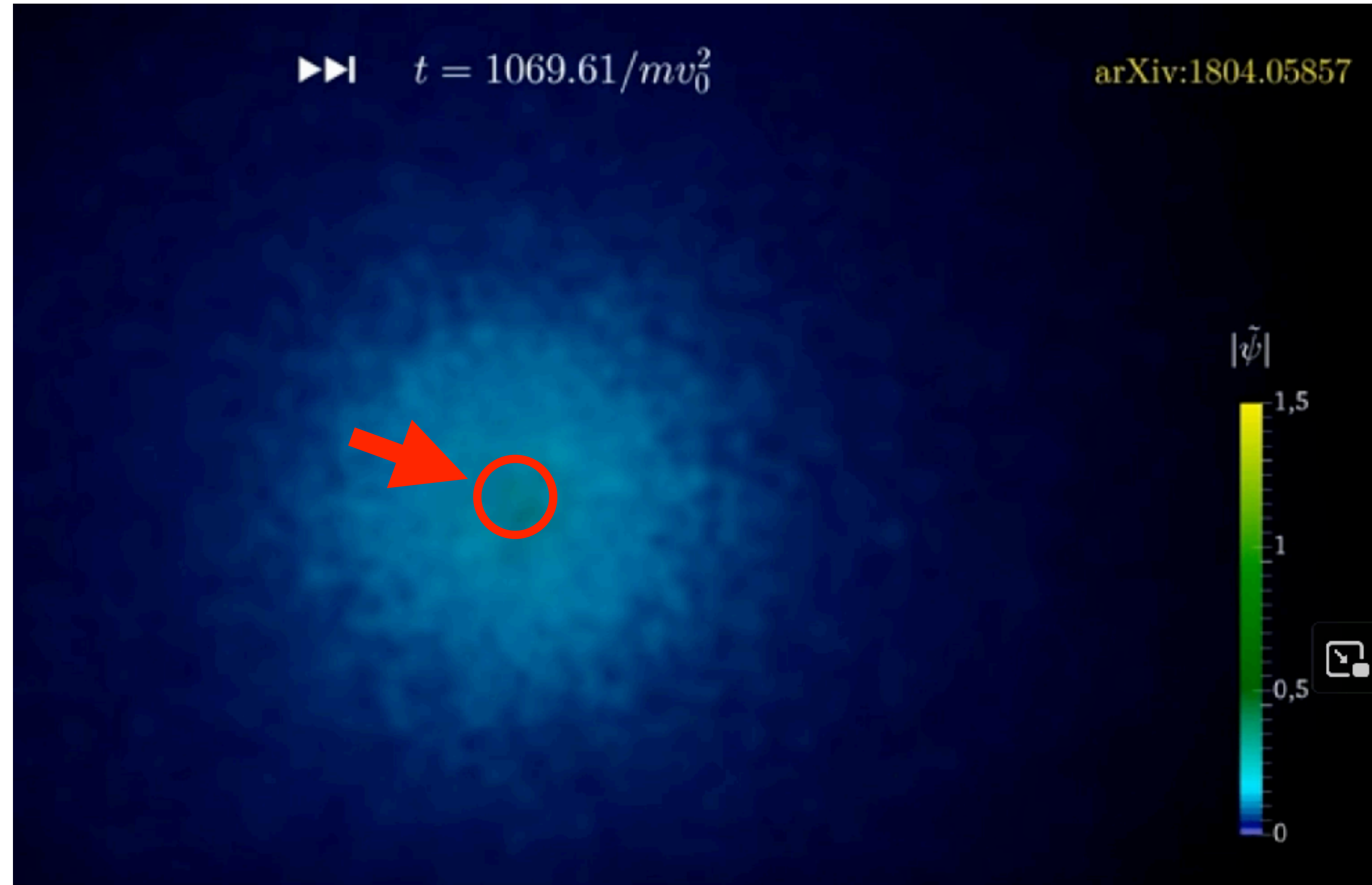


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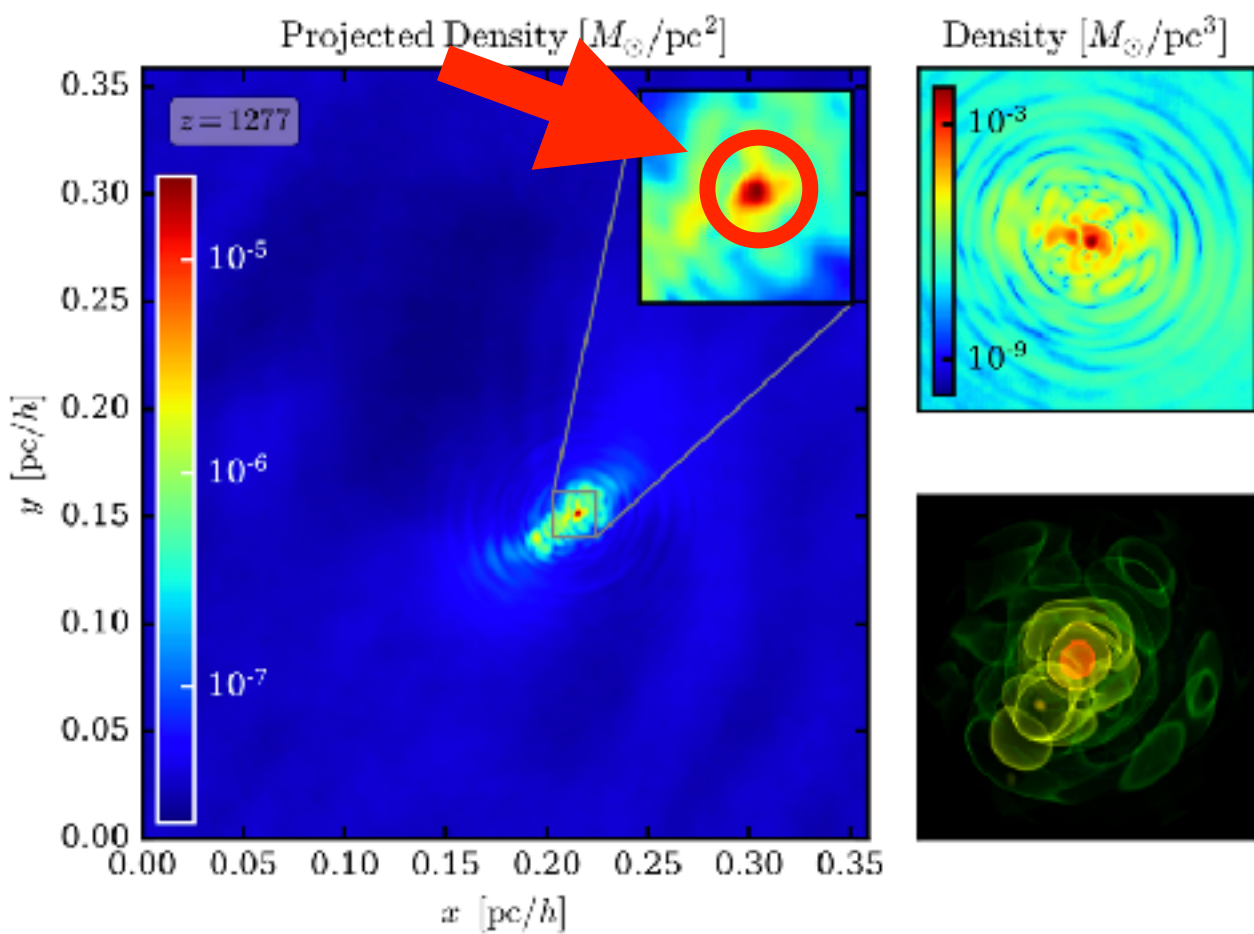


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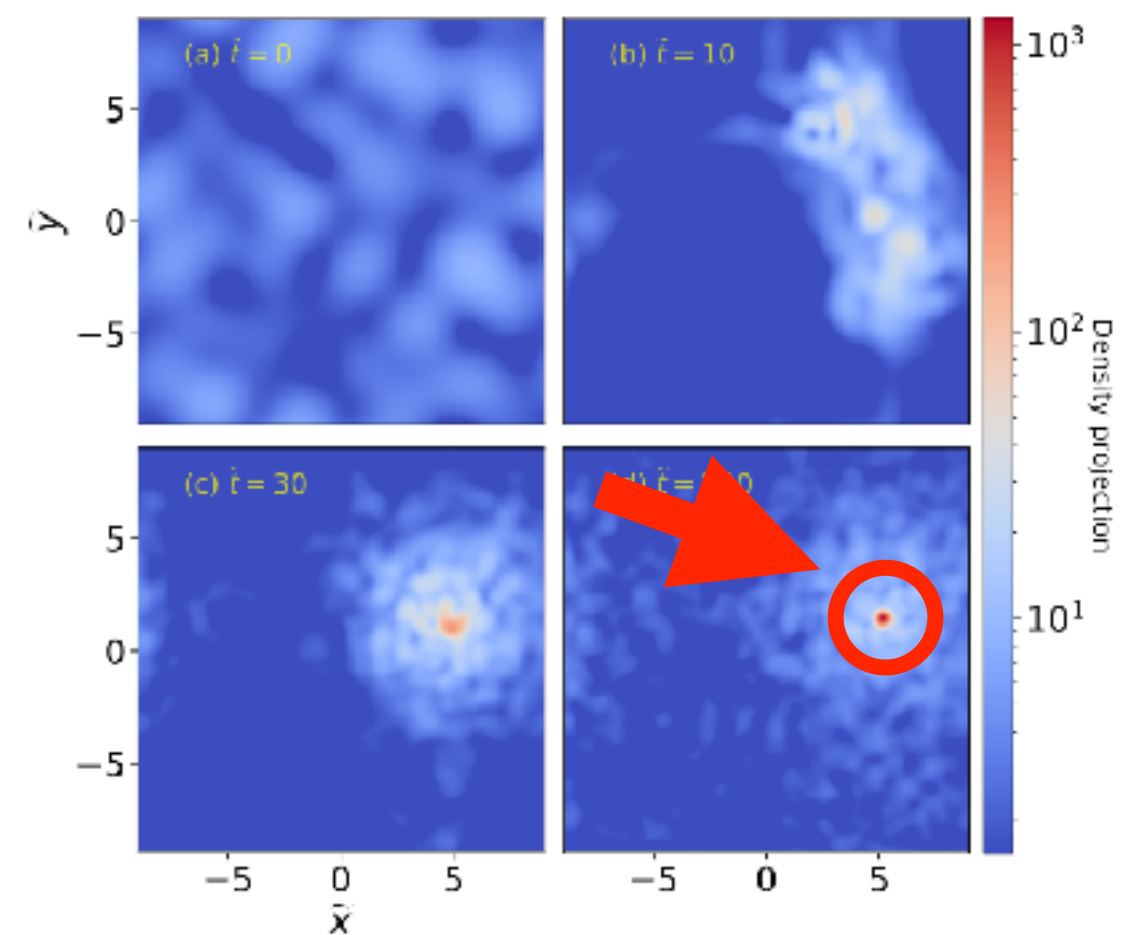
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(among others!)

Equations of Motion

$$\mathcal{L}_a \supset \frac{1}{2} \partial_\mu a \partial^\mu a - \frac{m_a^2}{2} a^2 + \frac{\lambda}{4!} a^4 - \dots$$

- Axions are
- non-relativistic
 - classical field
 - gravitating

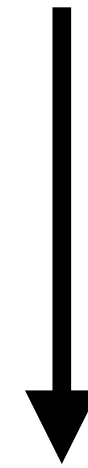


$$i \frac{\partial \psi}{\partial t} = \left[-\frac{\nabla^2}{2m_a} + V_g(|\psi|^2) + V_{int}(|\psi|^2) \right] \psi$$

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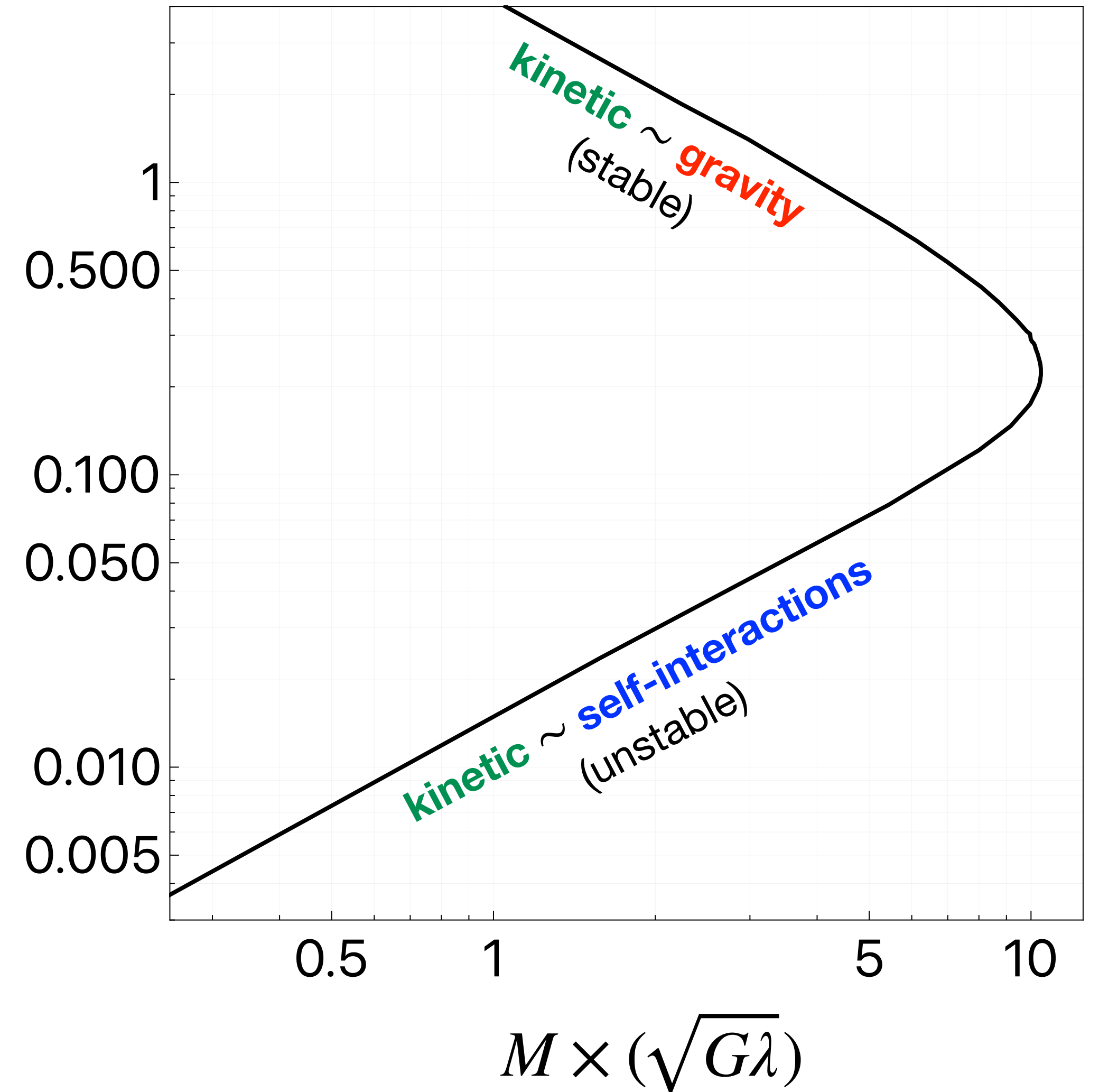
Boson star:
ground state of
Schrödinger eq.
with self-gravity

kinetic
 $\propto \frac{1}{m_a R^2}$

gravity
 $\propto -\frac{G m_a M}{R}$

self-interactions
 $\propto -\frac{|\lambda| M}{m_a^3 R^3} + \dots$

$R \times (m_a^2 \sqrt{G/\lambda})$



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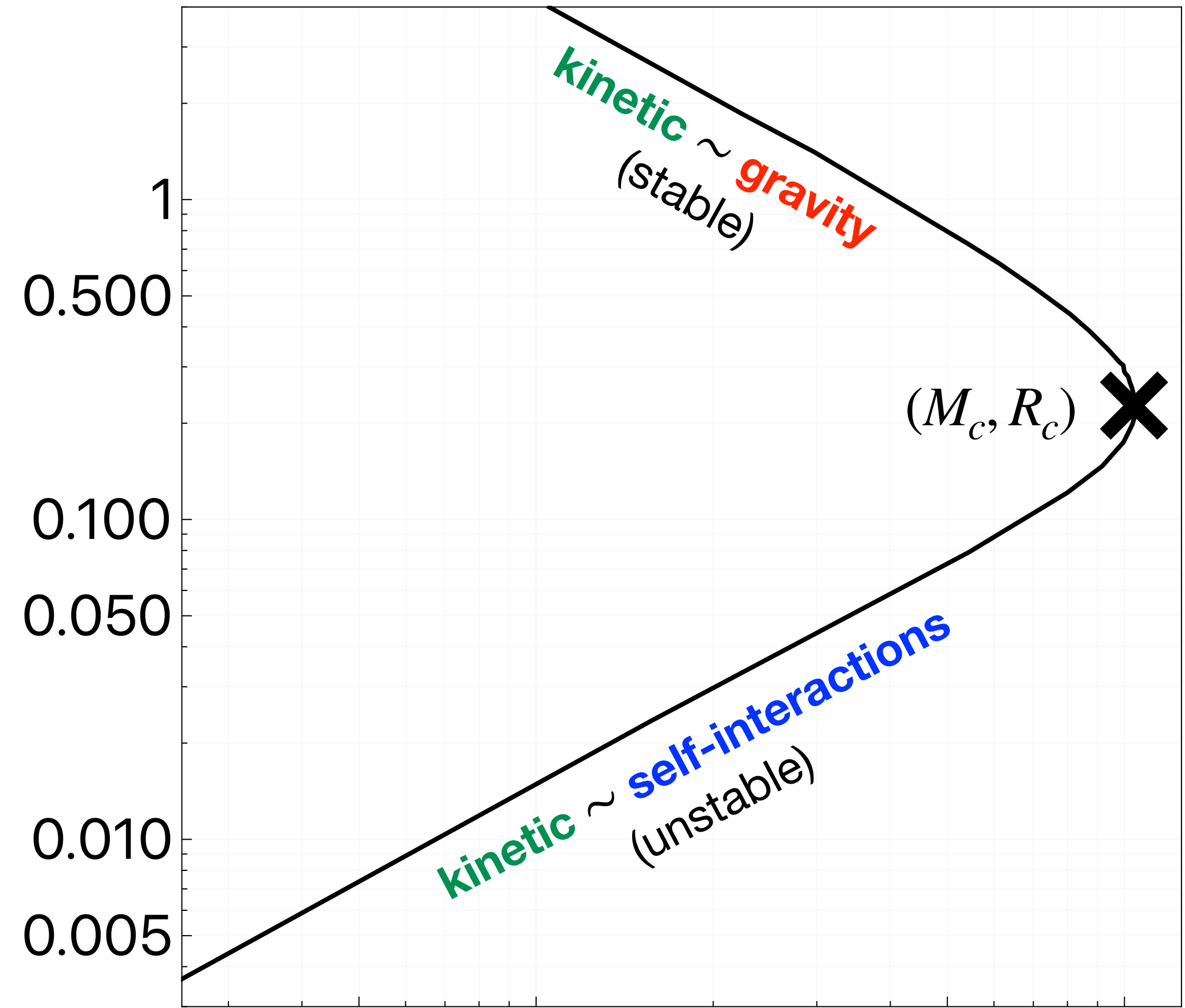


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kinetic	gravity	self-interactions
$\propto \frac{1}{m_a R^2}$	$\propto -\frac{G m_a M}{R}$	$\propto -\frac{ \lambda M}{m_a^3 R^3} + \dots$

$R \times (m_a^2 \sqrt{G/\lambda})$

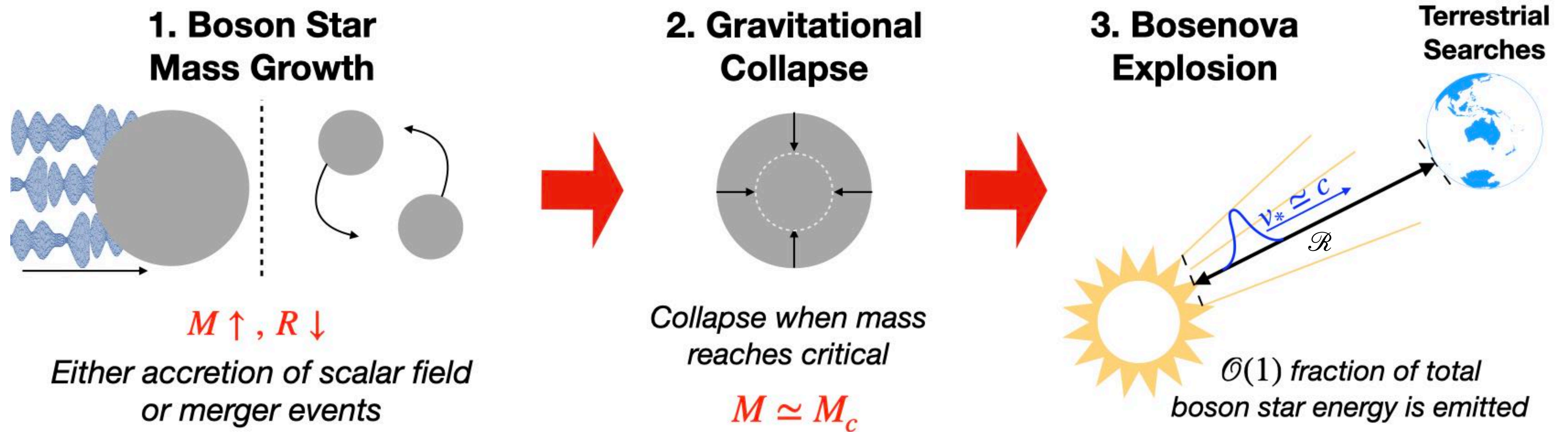


$$M_c \simeq \frac{10}{\sqrt{G\lambda}} \simeq 10^3 M_\odot \left(\frac{10^{-80}}{\lambda} \right)^{1/2}$$

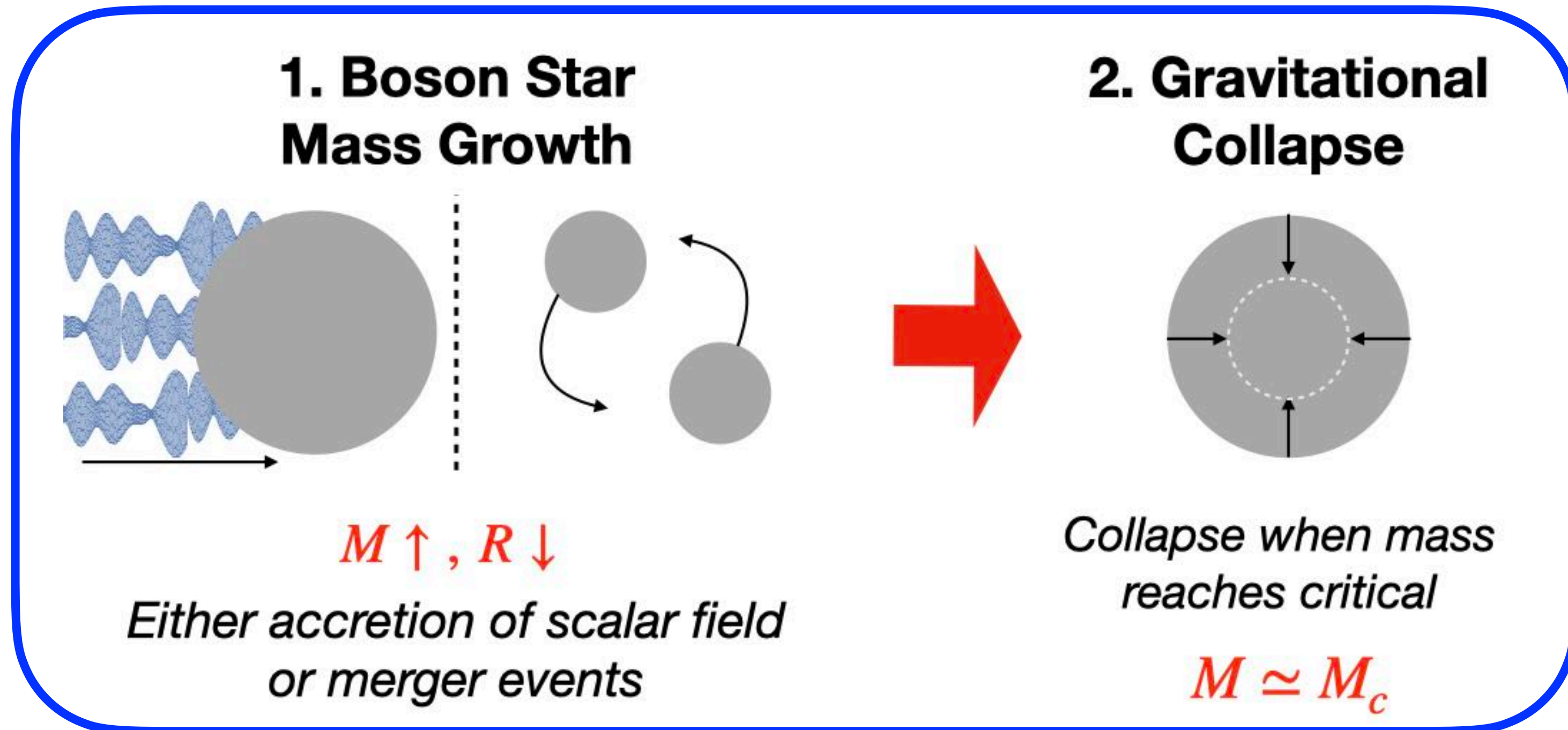
$$R_c \simeq \frac{0.2}{m_a^2} \sqrt{\frac{\lambda}{G}} \simeq 70 R_\odot \left(\frac{10^{-15} \text{ eV}}{m_a} \right)^2 \left(\frac{\lambda}{10^{-80}} \right)^{1/2}$$

$M \times (\sqrt{G\lambda})$

Life Cycle of a Boson Star



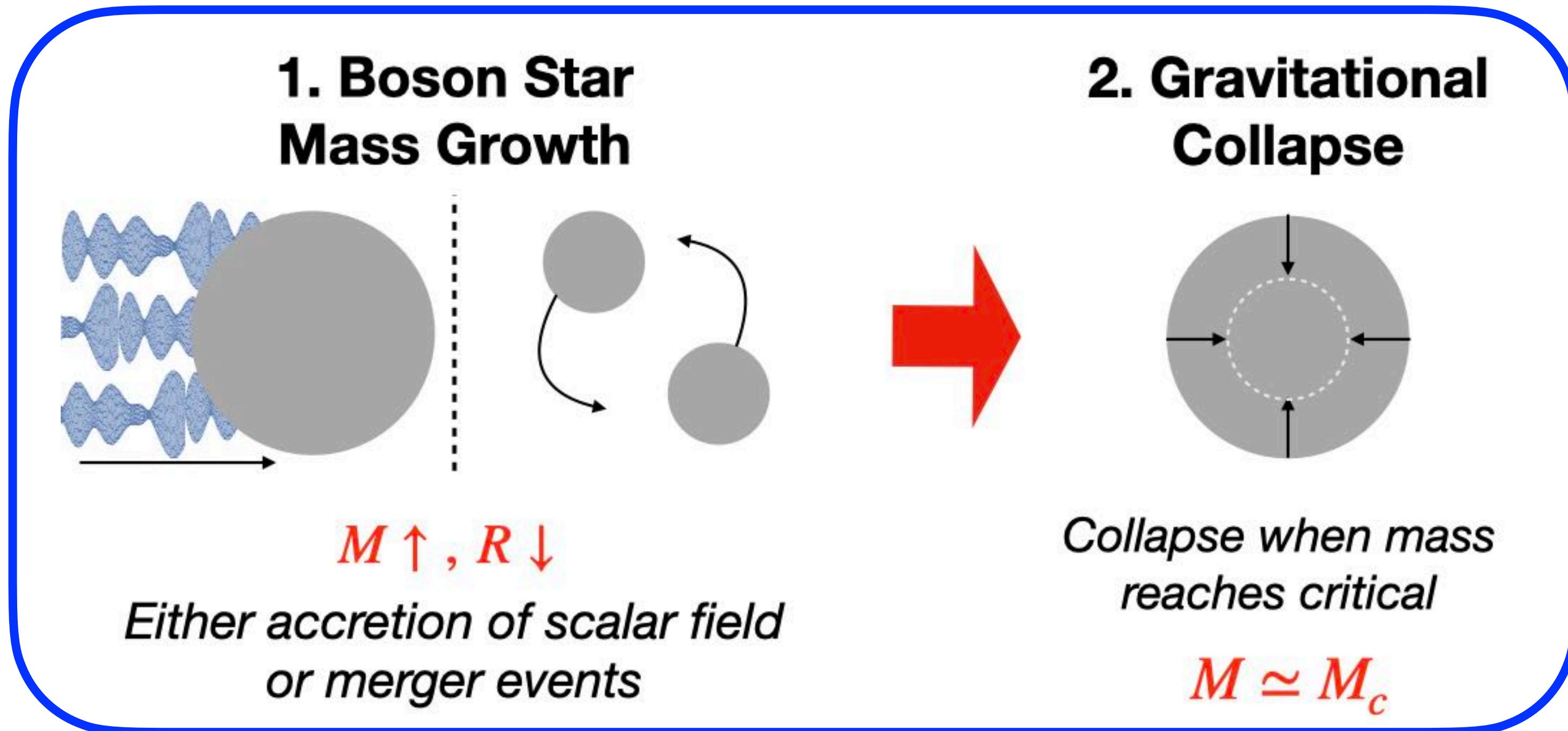
Event Rate vs Burst Flux



Highly model-dependent!

- *Cosmology (formation history, mass fraction)*
- *Astrophysics (mass growth, merger history)*
- *Particle physics (interactions \rightarrow relaxation rate)*

Event Rate vs Burst Flux



DISTRIBUTIONS AND COLLISION RATES OF ALP STARS IN THE MILKY WAY

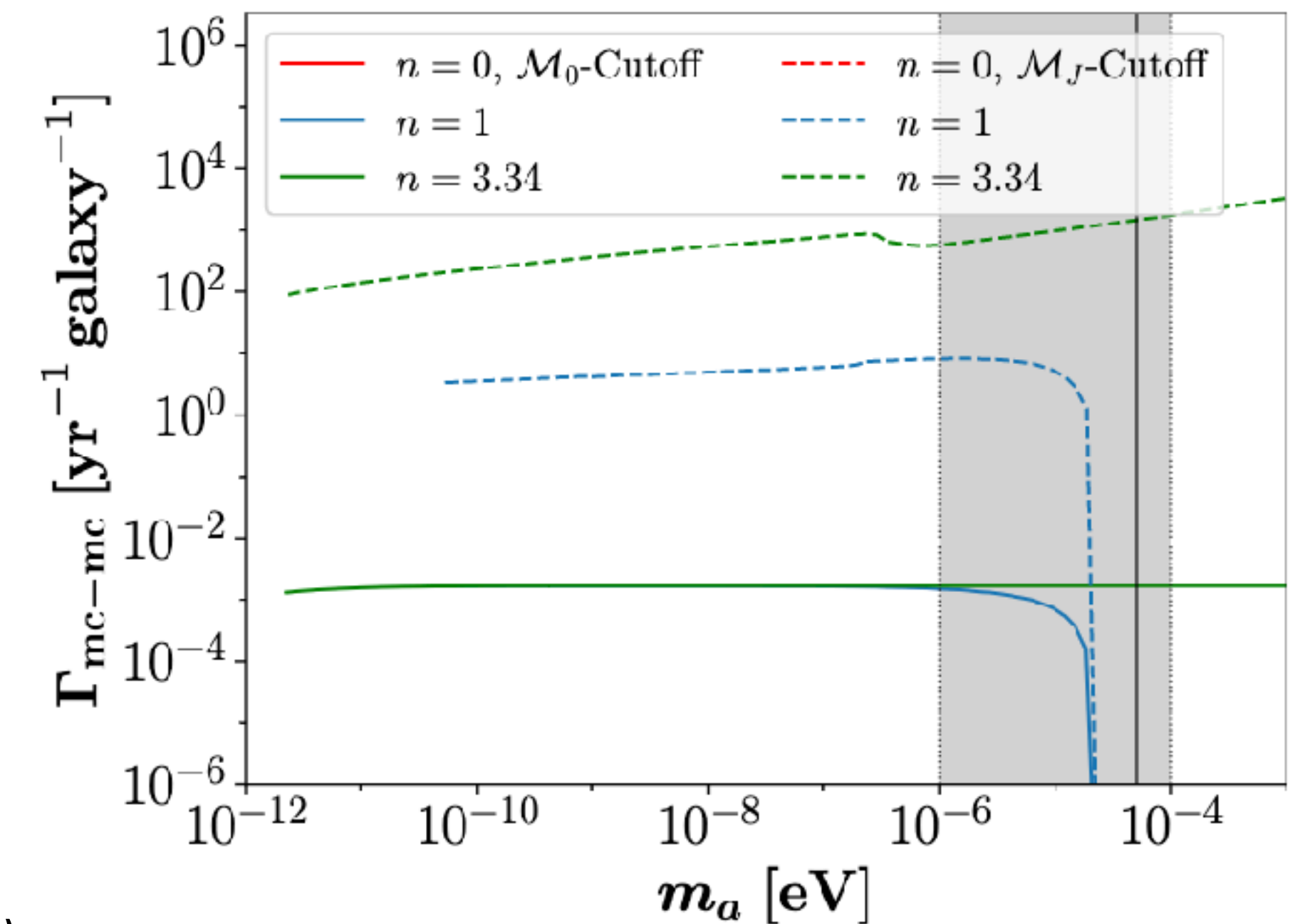
Dennis Maseizik
 II. Institut for theoretical Physics
 Hamburg University
 dennis.maseizik@desy.de

Günter Sigl
 II. Institut for theoretical Physics
 Hamburg University
 guenter.sigl@desy.de

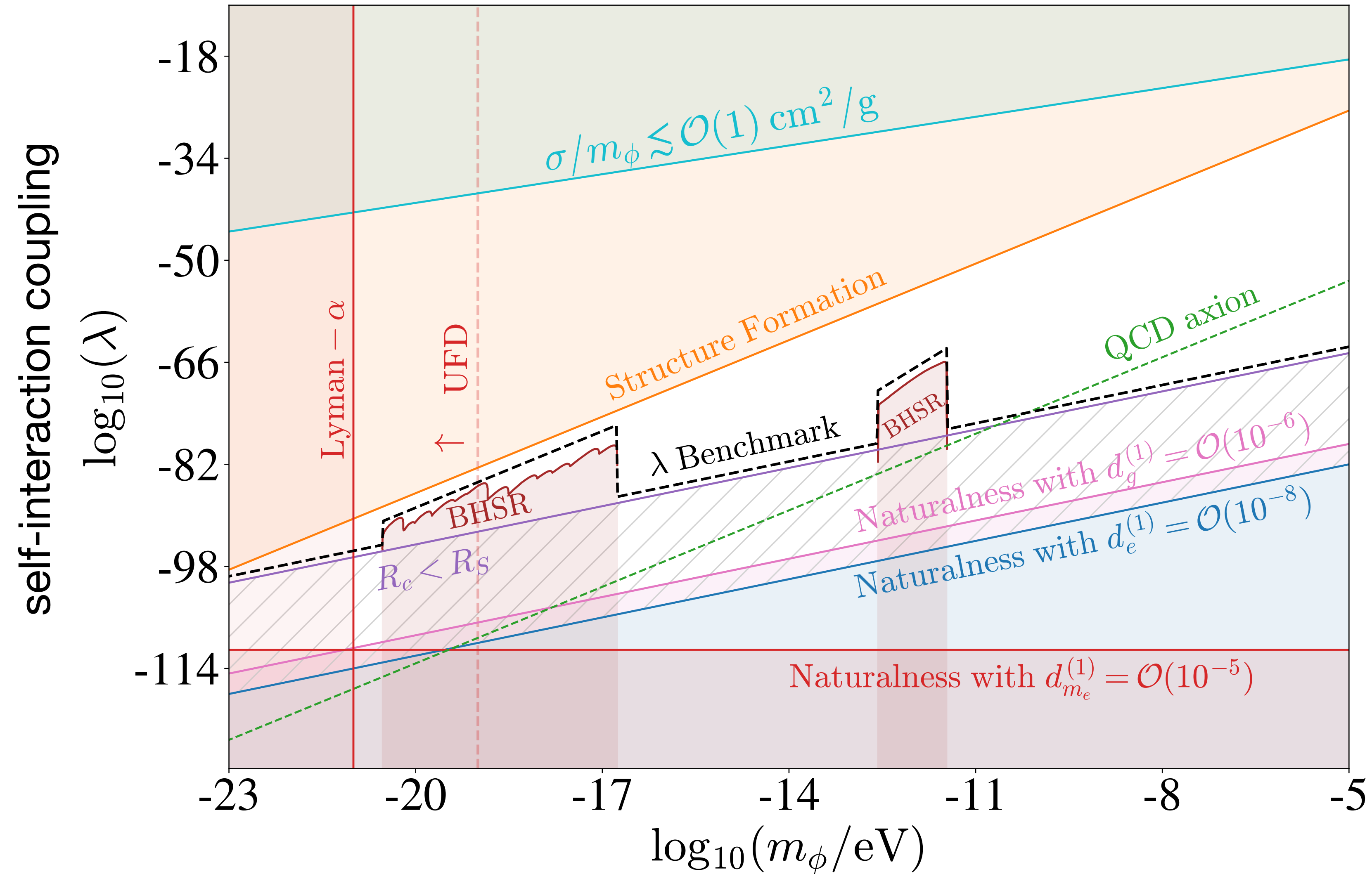
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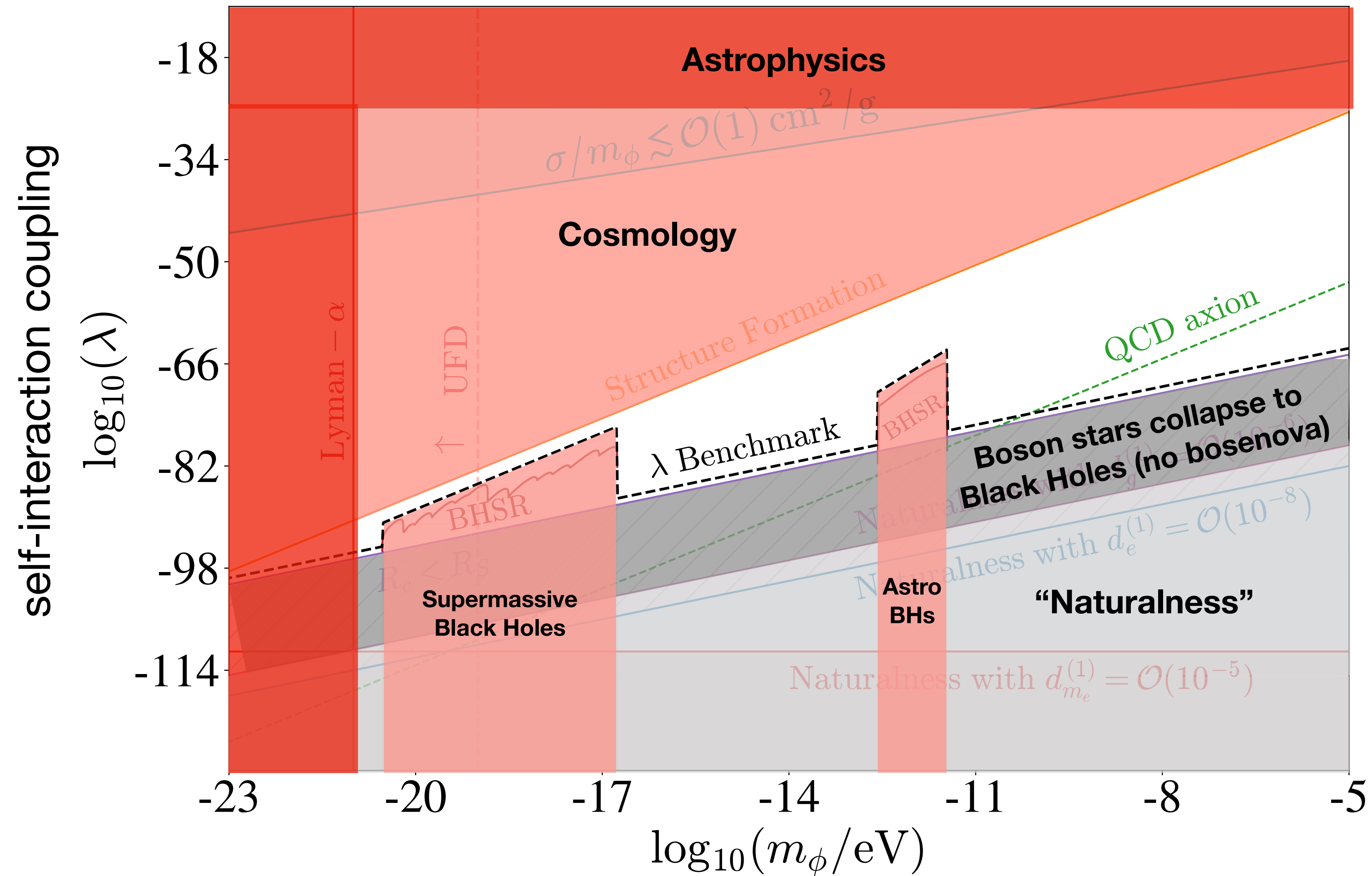
Maseizik, Sigl (2404.07908)



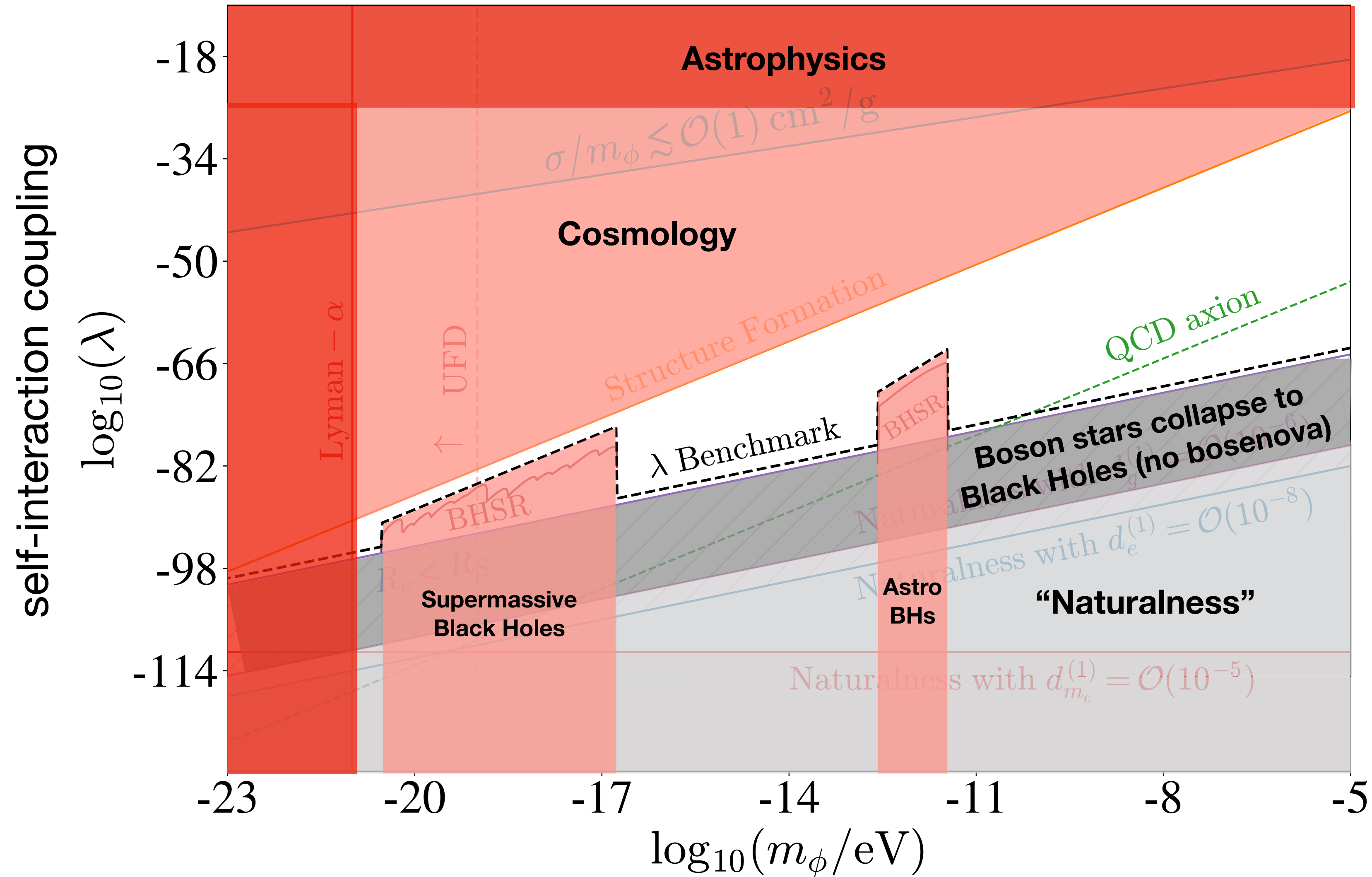
Event Rate vs Burst Flux



Event Rate vs Burst Flux



Event Rate vs Burst Flux



Smaller M_c (signal)
Larger Γ (rate)

Recall that

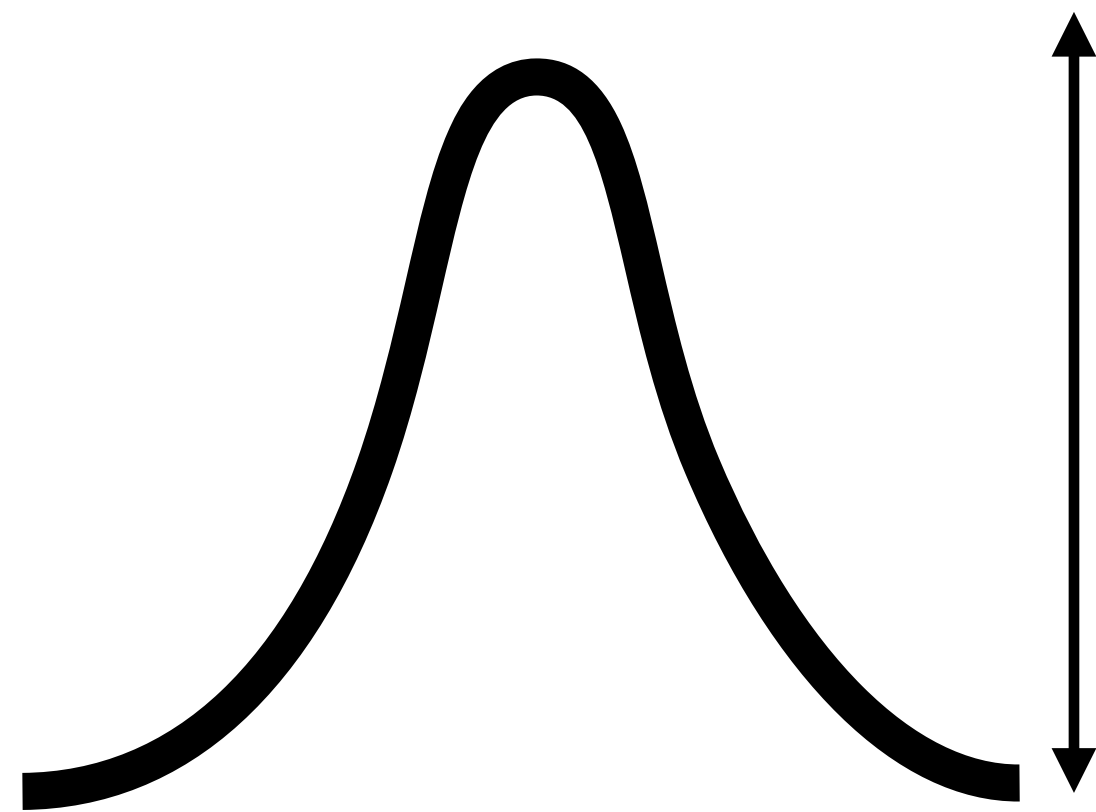
$$M_c \simeq \frac{10}{\sqrt{G\lambda}} \simeq 10^3 M_\odot \left(\frac{10^{-80}}{\lambda} \right)^{1/2}$$

Larger M_c (signal)
Smaller Γ (rate)

Future work is needed
to fully estimate the rate
across range of models

Event Rate vs Burst Flux

Burst properties:



momentum spectrum of emission:

peaked at $k_0 \sim \text{few} \times m_a$

with width $\delta k \sim m_a$

total emitted energy,

$$\sim M_c = \frac{10M_{\text{Pl}}}{\sqrt{\lambda}}$$

$$\rho_* \sim \frac{E_{\text{tot}}}{4\pi R^2 \delta x}$$

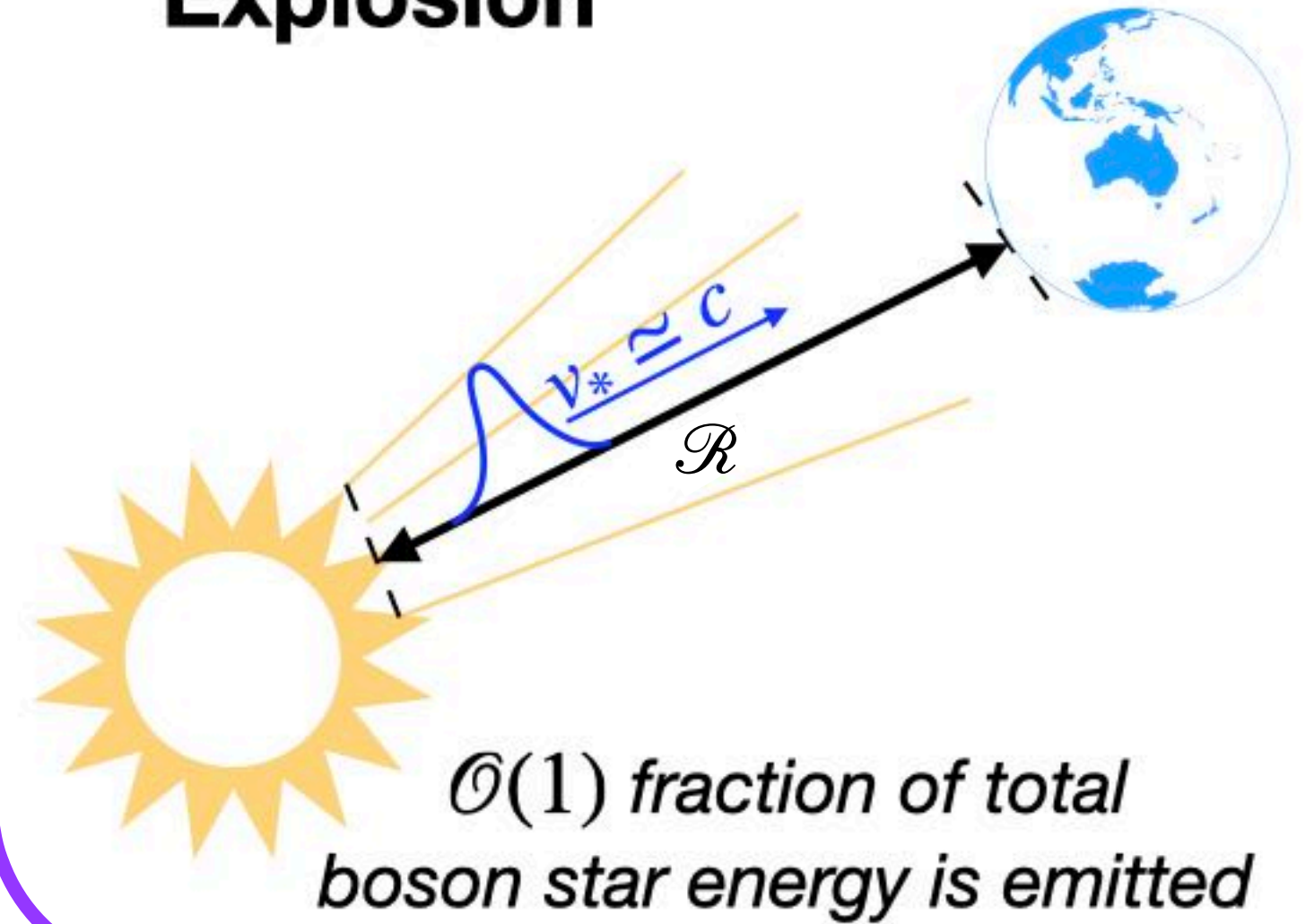
distance to burst

burst 'size'
(duration)

**Energy density
in burst**

3. Bosenova Explosion

Terrestrial
Searches



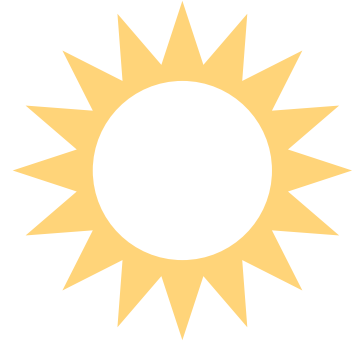
Prediction of Bosenova

**Eby, Leembruggen,
Suranyi, Wijewardhana (1608.06911)**

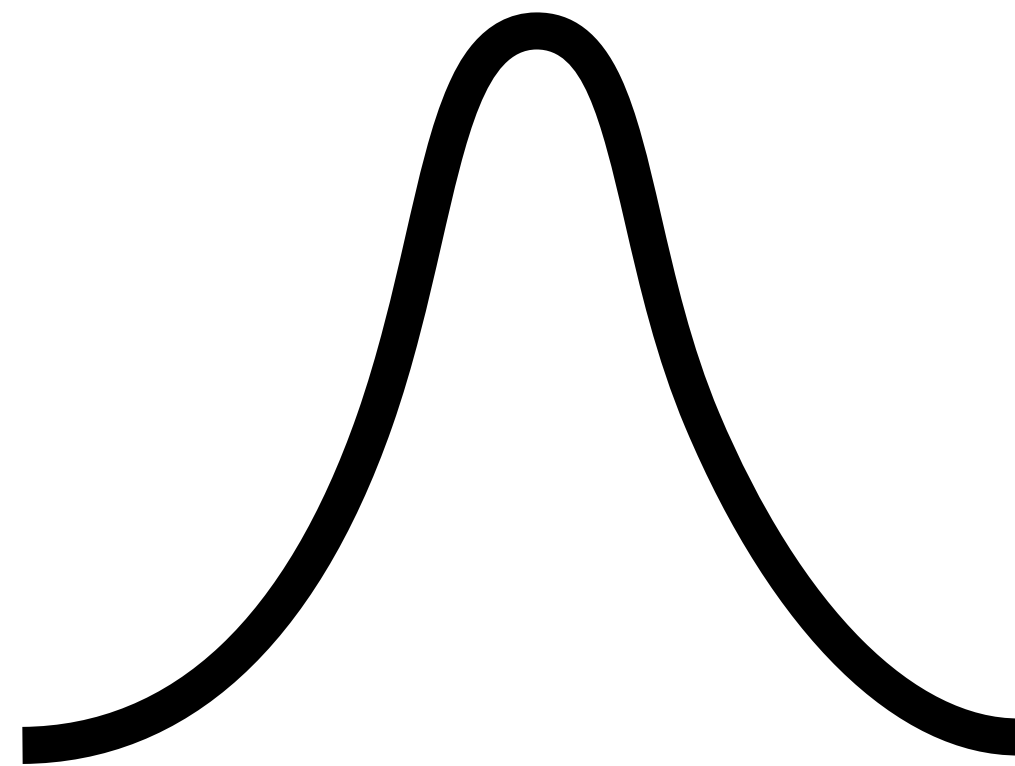
**Numerical simulations of
collapse+Bosenova process**

Levkov, Panin, Tkachev (1609.03611)

Wave Spreading in Flight



At Source

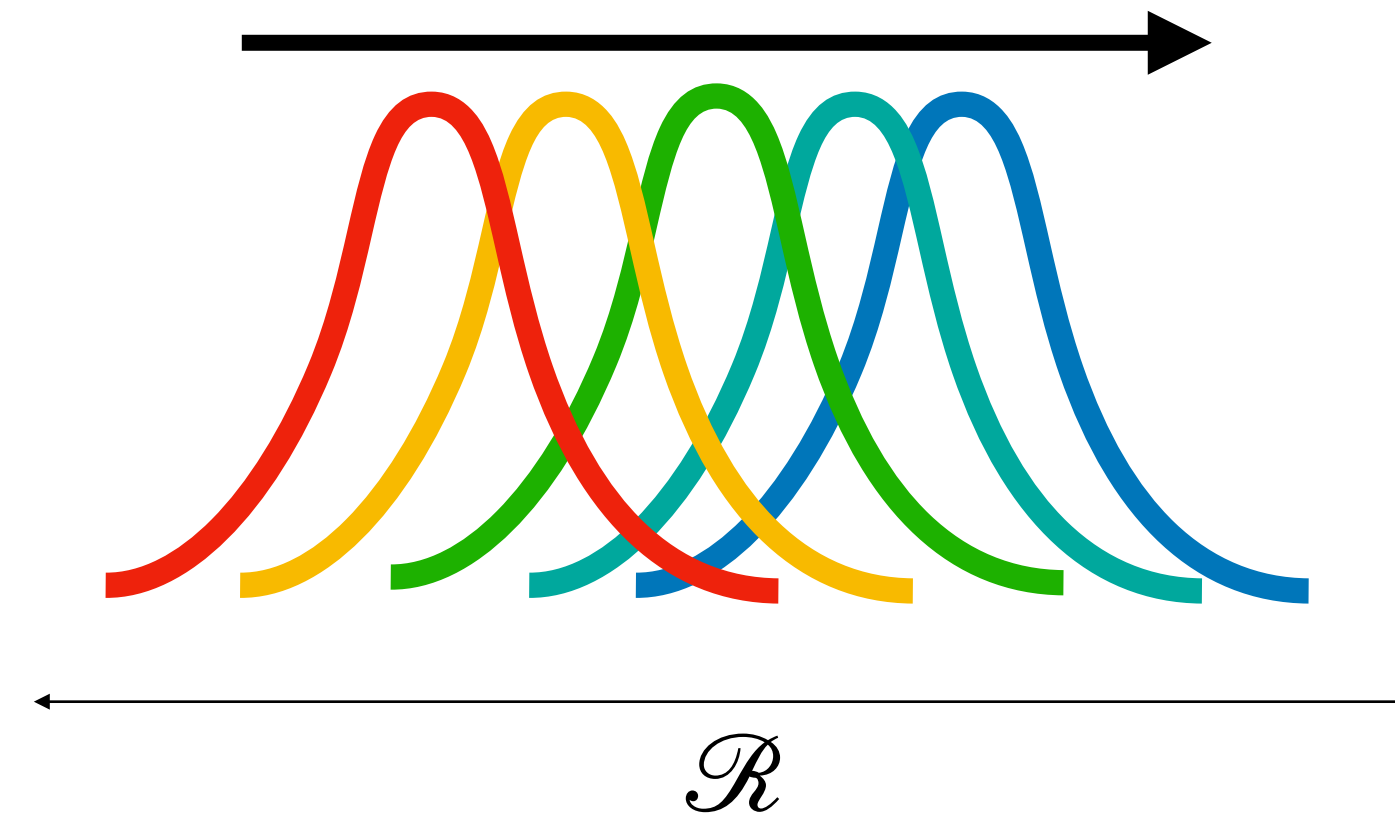


momentum spectrum of emission:
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with width $\delta k \sim m_a$

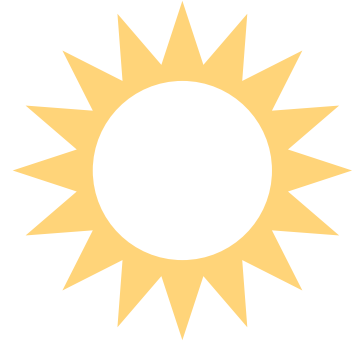
imagine discrete momentum modes
 k_1, k_2, k_3, \dots

Fastest momentum modes
“escape” from slower ones
during propagation

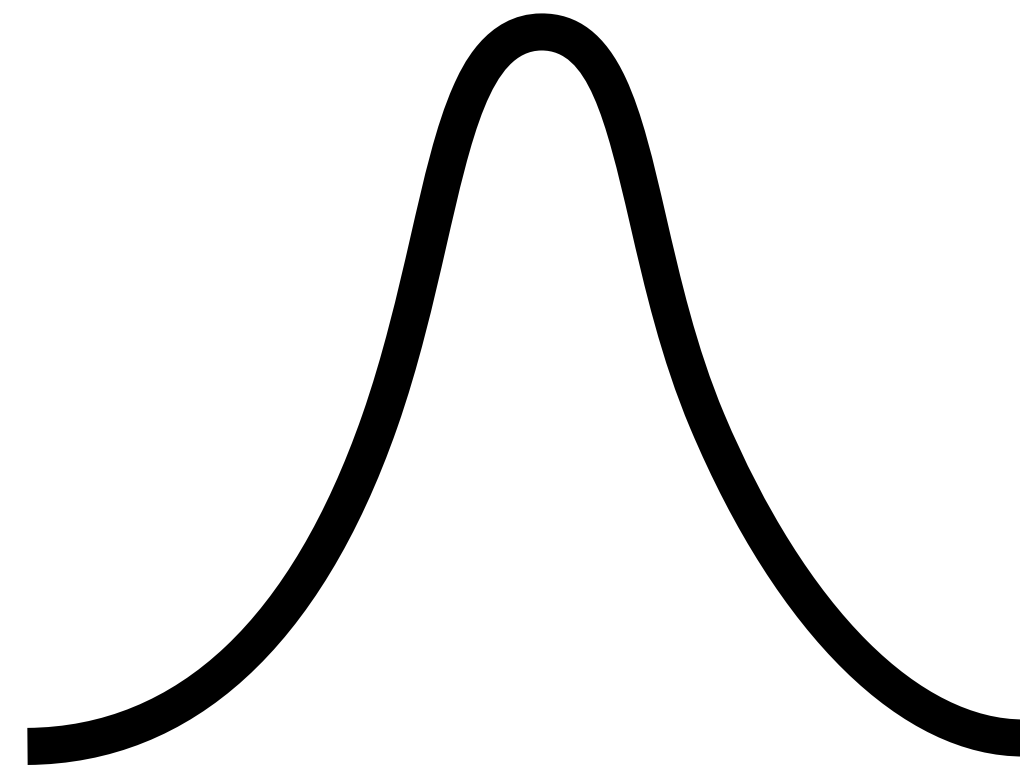
$$k_1 < k_2 < k_3, \dots$$



Wave Spreading in Flight



At Source

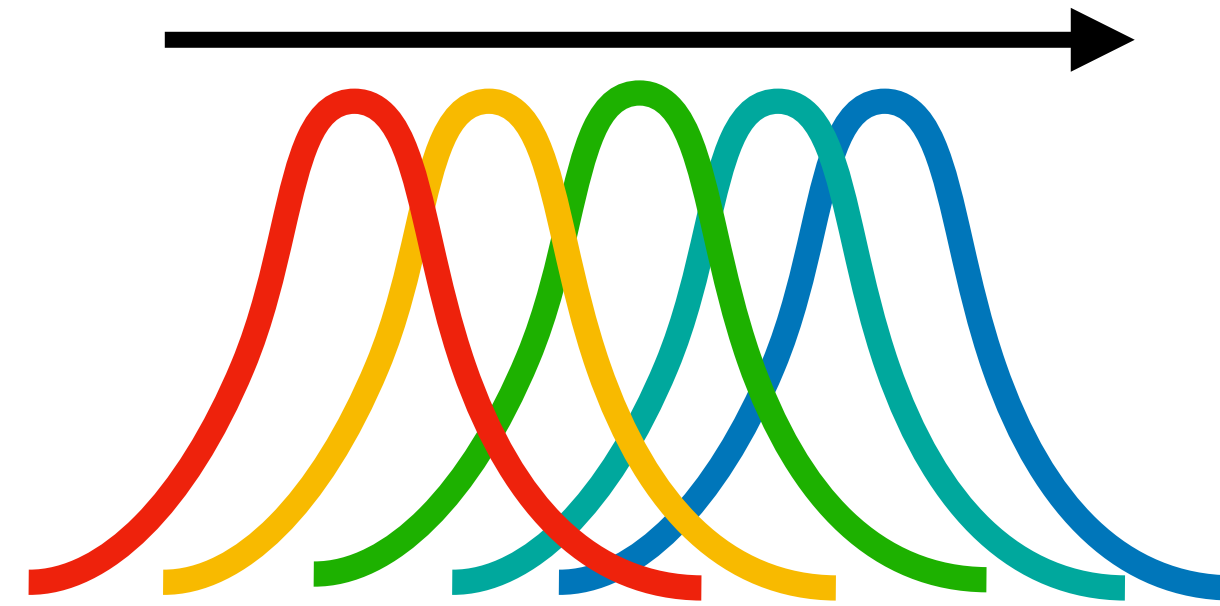


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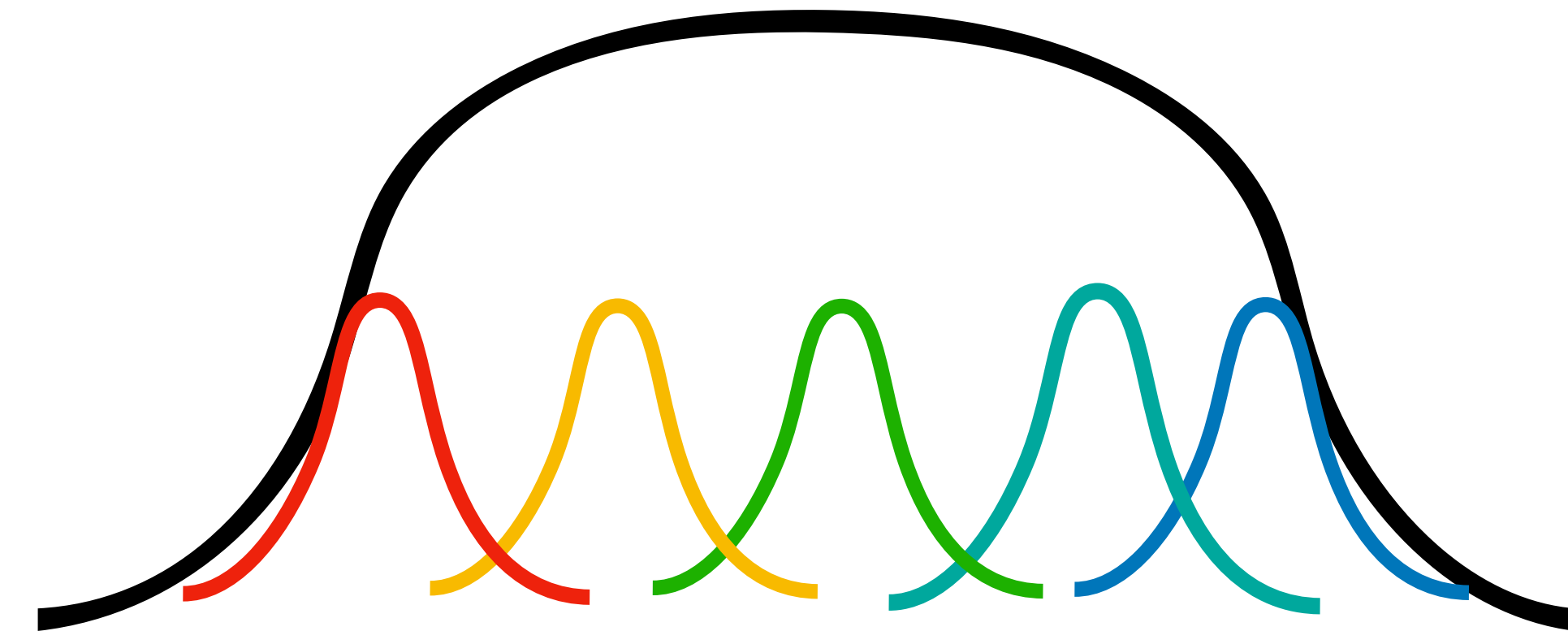
$$k_1 < k_2 < k_3, \dots$$



\mathcal{R}



At Detector



1. Long duration $\delta t \sim \text{month} \left(\frac{\mathcal{R}}{\text{pc}} \right)$

2. At any given moment in the detector,
one sees a **narrow** distribution
of momentum / energy

\Rightarrow “effective coherence time” $\tau_* \sim 10^{-2} \mathcal{R}$

Sensitivity to Transient Signal

Sensitivity Ratio
to SM coupling d_i at
fixed frequency ω_0

$$\frac{\text{(burst)} \quad d_{i,*}(\omega_0)}{\text{(DM)} \quad d_i(\omega_0)} \sim \left(\frac{\rho_{\text{dm}}}{\rho_*} \right)^n \frac{t_{\text{int}}^{1/4} \min(\tau_{\text{dm}}^{1/4}, t_{\text{int}}^{1/4})}{\min(\delta t^{1/4}, t_{\text{int}}^{1/4}) \min(\tau_*^{1/4}, t_{\text{int}}^{1/4})}$$

Timescales:

Interrogation time [†]	t_{int}
DM coherence time	τ_{dm}
Burst coherence time	τ_*
Burst duration	δt

“Is a given DM experiment equally/more sensitive to relativistic bursts compared to cold DM search?”

[†]laser coherence, natural linewidth, ...

Sensitivity to Transient Signal

Sensitivity Ratio
to SM coupling d_i at
fixed frequency ω_0

usual scaling of DM signal

$$\frac{d_{i,*}(\omega_0)}{d_i(\omega_0)} \sim \left(\frac{\rho_{\text{dm}}}{\rho_*} \right)^n \frac{t_{\text{int}}^{1/4} \min(\tau_{\text{dm}}^{1/4}, t_{\text{int}}^{1/4})}{\min(\delta t^{1/4}, t_{\text{int}}^{1/4}) \min(\tau_*^{1/4}, t_{\text{int}}^{1/4})}$$

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Sensitivity to Transient Signal

**Sensitivity Ratio
to SM coupling d_i at
fixed frequency ω_0**

depends on axion-SM coupling:

$n = 1/2$ (linear coupling)

$n = 1$ (quadratic coupling)

usual scaling of DM signal

$$\frac{d_{i,*}(\omega_0)}{d_i(\omega_0)} \sim \left(\frac{\rho_{\text{dm}}}{\rho_*} \right)^n \frac{t_{\text{int}}^{1/4} \min(\tau_{\text{dm}}^{1/4}, t_{\text{int}}^{1/4})}{\min(\delta t^{1/4}, t_{\text{int}}^{1/4}) \min(\tau_*^{1/4}, t_{\text{int}}^{1/4})}$$

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Sensitivity to Transient Signal

**Sensitivity Ratio
to SM coupling d_i at
fixed frequency ω_0**

$$\frac{d_{i,*}(\omega_0)}{d_i(\omega_0)} \sim$$

depends on axion-SM coupling:

$$n = 1/2 \text{ (linear coupling)}$$

$$n = 1 \text{ (quadratic coupling)}$$

usual scaling of DM signal

$$\left(\frac{\rho_{\text{dm}}}{\rho_*} \right)^n t_{\text{int}}^{1/4} \min(\tau_{\text{dm}}^{1/4}, t_{\text{int}}^{1/4})$$

$$\min(\delta t^{1/4}, t_{\text{int}}^{1/4})$$

$$\min(\tau_*^{1/4}, t_{\text{int}}^{1/4})$$

energy
density in
the burst

“did we catch all/most
of the signal?”

“is the signal coherent for
the whole integration time?”

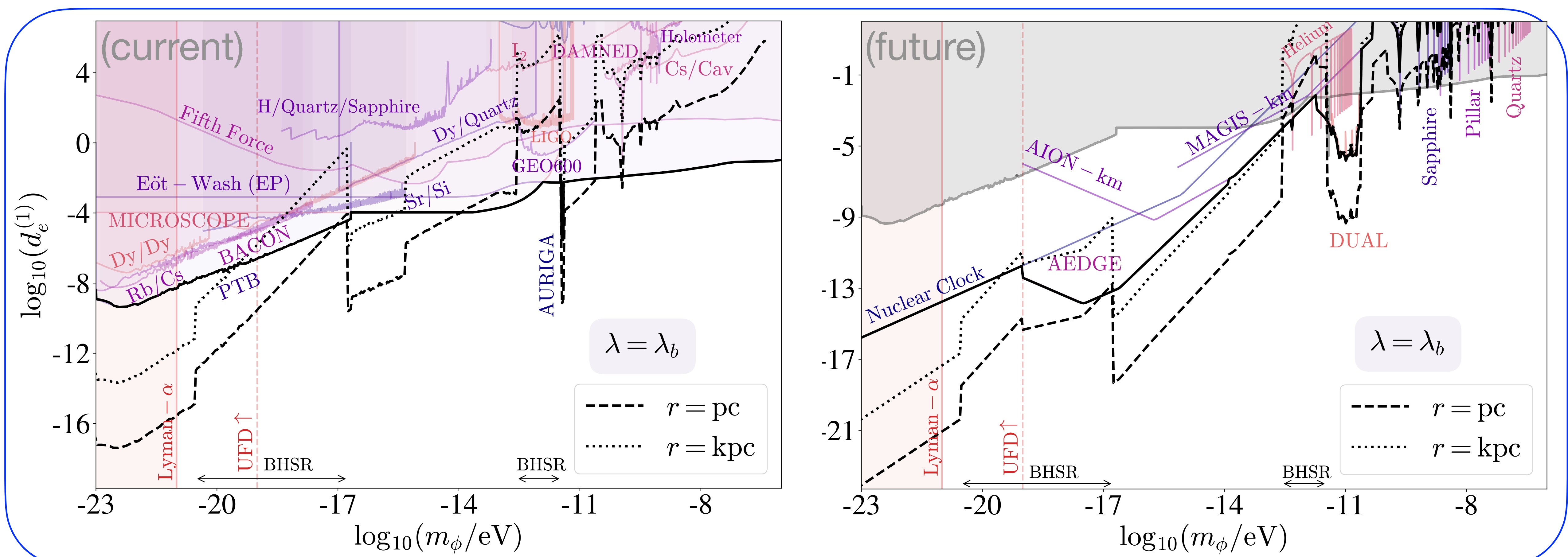
**“Is a given DM experiment equally/more sensitive
to relativistic bursts compared to cold DM search?”**

Timescales:

- Interrogation time[†] t_{int}
- DM coherence time τ_{dm}
- Burst coherence time τ_*
- Burst duration δt

[†]laser coherence, natural linewidth, ...

coupling:	linear, psuedoscalar	linear, scalar	quadratic, scalar
Lagrangian:	$\mathcal{L} \supset g_{a\gamma} a F^{\mu\nu} \tilde{F}_{\mu\nu}$	$\mathcal{L} \supset d_e^{(1)} \frac{a}{2M_{\text{Pl}}} F^{\mu\nu} F_{\mu\nu}$	$\mathcal{L} \supset d_e^{(2)} \left(\frac{a}{2M_{\text{Pl}}} \right)^2 F^{\mu\nu} F_{\mu\nu}$
example field:	QCD axion	relaxion	QCD axion or relaxion
experiments:	ABRACADABRA, DM-Radio, ...	Quantum sensors	Quantum sensors
reference:	Eby , Shirai, Stadnik, Takhistov (2106.14893)	Arakawa, Eby , Safronova, Takhistov, Zaheer (2306.16468)	Arakawa, Zaheer, Eby , Takhistov, Safronova (2402.06736)



coupling:

linear, pseudoscalar

Lagrangian:

$$\mathcal{L} \supset g_{a\gamma} a F^{\mu\nu} \tilde{F}_{\mu\nu}$$

example field:

QCD axion

experiments:

ABRACADABRA, DM-Radio, ...

reference:

Eby, Shirai, Stadnik,
Takhistov (2106.14893)

linear, scalar

$$\mathcal{L} \supset d_e^{(1)} \frac{a}{2M_{\text{Pl}}} F^{\mu\nu} F_{\mu\nu}$$

relaxion

Quantum sensors

Arakawa, **Eby**, Safronova, Takhistov,
Zaheer (2306.16468)

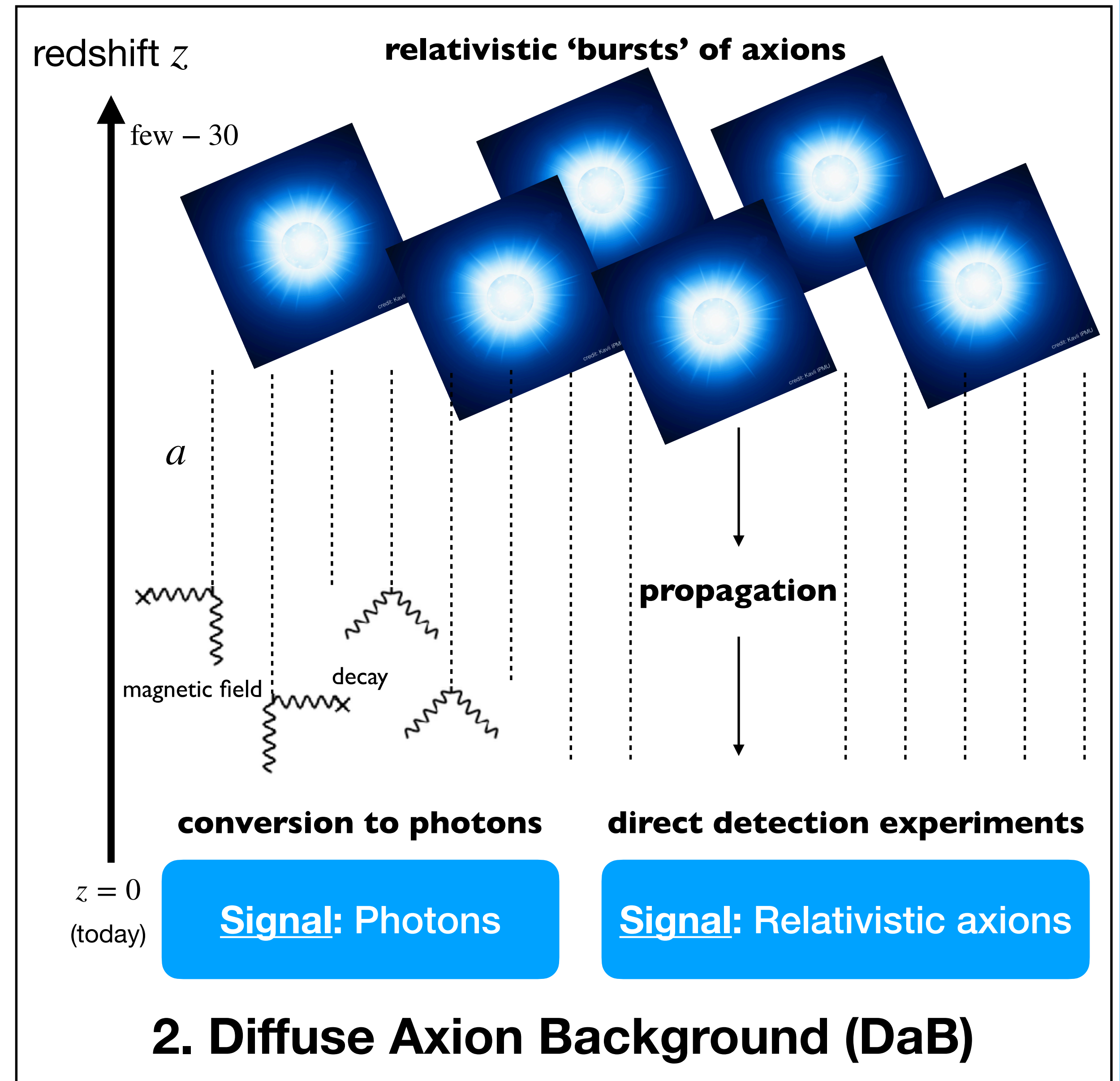
quadratic, scalar

$$\mathcal{L} \supset d_e^{(2)} \left(\frac{a}{2M_{\text{Pl}}} \right)^2 F^{\mu\nu} F_{\mu\nu}$$

QCD axion or relaxion

Quantum sensors

Arakawa, Zaheer, **Eby**, Takhistov,
Safronova (2402.06736)



Astrophysical bursts of relativistic axions

generally characterised by

flux

rate

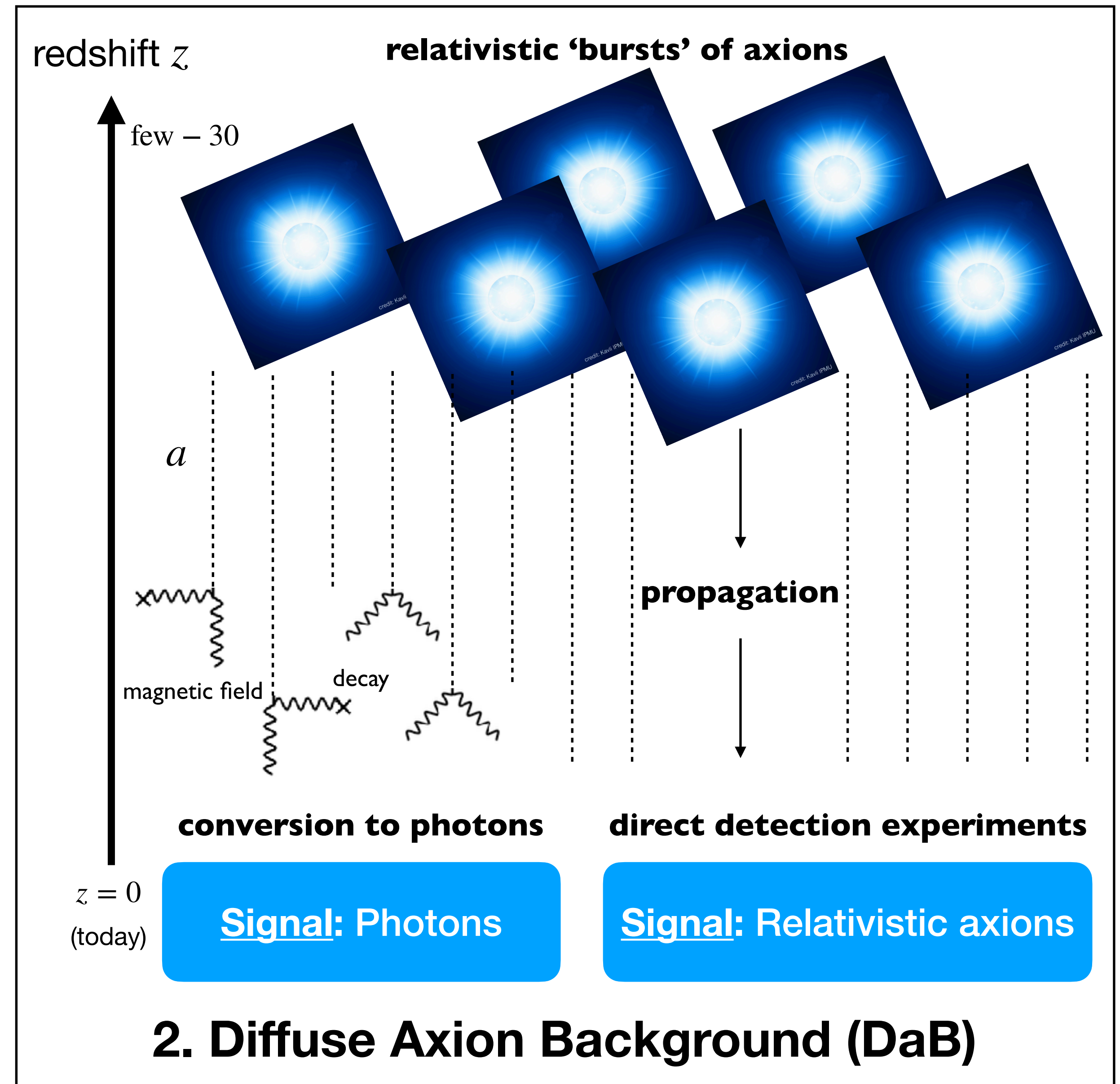
$$\frac{dN_a}{d\omega}(\omega)$$

$$R_{\text{burst}}(z)$$

DaB flux* in present day

$$\frac{d\phi}{d\omega}(\omega) = \int_0^\infty dz \frac{dN_a(\omega(1+z))}{d\omega} \frac{R_{\text{burst}}(z)}{H(z)}$$

*note: flux $\frac{dN_a}{d\omega}$ (# of particles) vs flux $\frac{d\phi}{d\omega}$ (# per area per time)

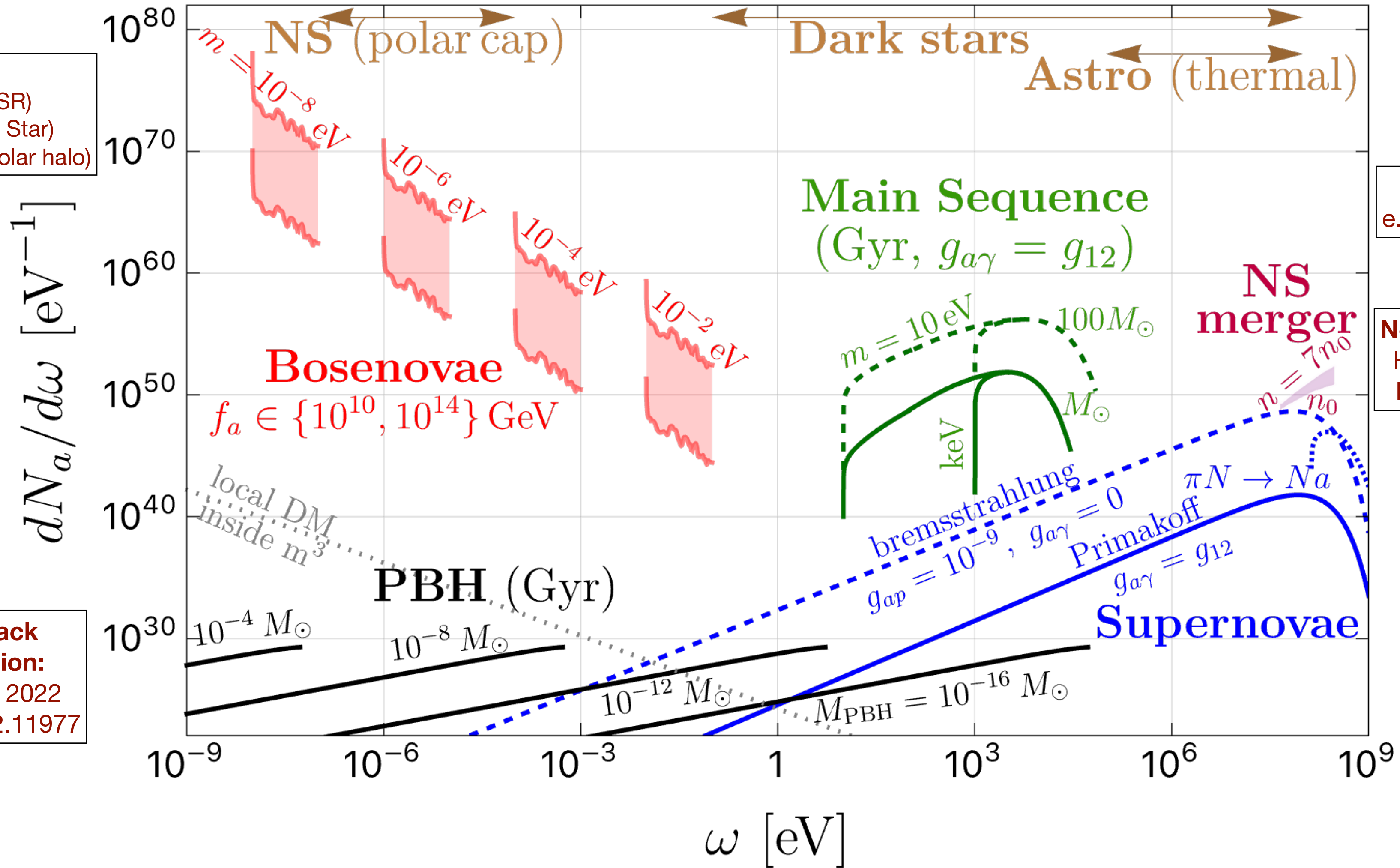


Bursts Abound

NS polar cap:
Prabhu, PRD 2021
Noordhuis++, PRL 2023, arXiv:2307.11811

Dark stars:
Maselli++, PRD 2017
Curtin and Setford, PRL 2019
Hippert++, PRD 2021

Bosenovae:
Yoshino++, PTP 2012 (BHSR)
Levkov++, PRL 2016 (Boson Star)
Budker, **J Eby**++, JCAP 2023 (Solar halo)



Main sequence stars:
e.g. Nguyen++, JCAP 2023

Neutron star mergers:
Harris++, JCAP 2020
Fiorillo++, PRD 2022

Primordial black hole evaporation:
Agashe++, PRD 2022
Jho++, arXiv:2212.11977

Supernovae:
fits via
Carenza++, JCAP 2019
Calore++, PRD 2020
(First DaB proposal!)

Eby, Takhistov (2402.00100)

Parameterization: DaB

$$\frac{dN_a}{d\omega}(\omega) \propto \frac{E_{\text{tot}}}{m_a^2} \frac{\exp\left(-\frac{(\omega - \bar{\omega})^2}{\delta\omega^2}\right)}{\delta\omega/m_a}$$

$$R_{\text{burst}}(z) \propto \frac{\mathcal{F} \bar{\rho}_U H_0}{E_{\text{tot}}} f(z)$$

\mathcal{F} : total DM fraction converted to DaB

$f(z)$: dimensionless rate of bursts

$\bar{\omega}$: peak burst energy per particle

$\delta\omega$: spread in burst energy per particle

E_{tot} : energy emitted per burst

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input parameters

particle physics

$$m_a, f_a, g_{a\gamma}, \dots$$

burst parameterisation

$$\mathcal{F}, f(z), \bar{\omega}, \delta\omega, E_{\text{tot}}$$

(cosmology) (individual bursts) cancels in product

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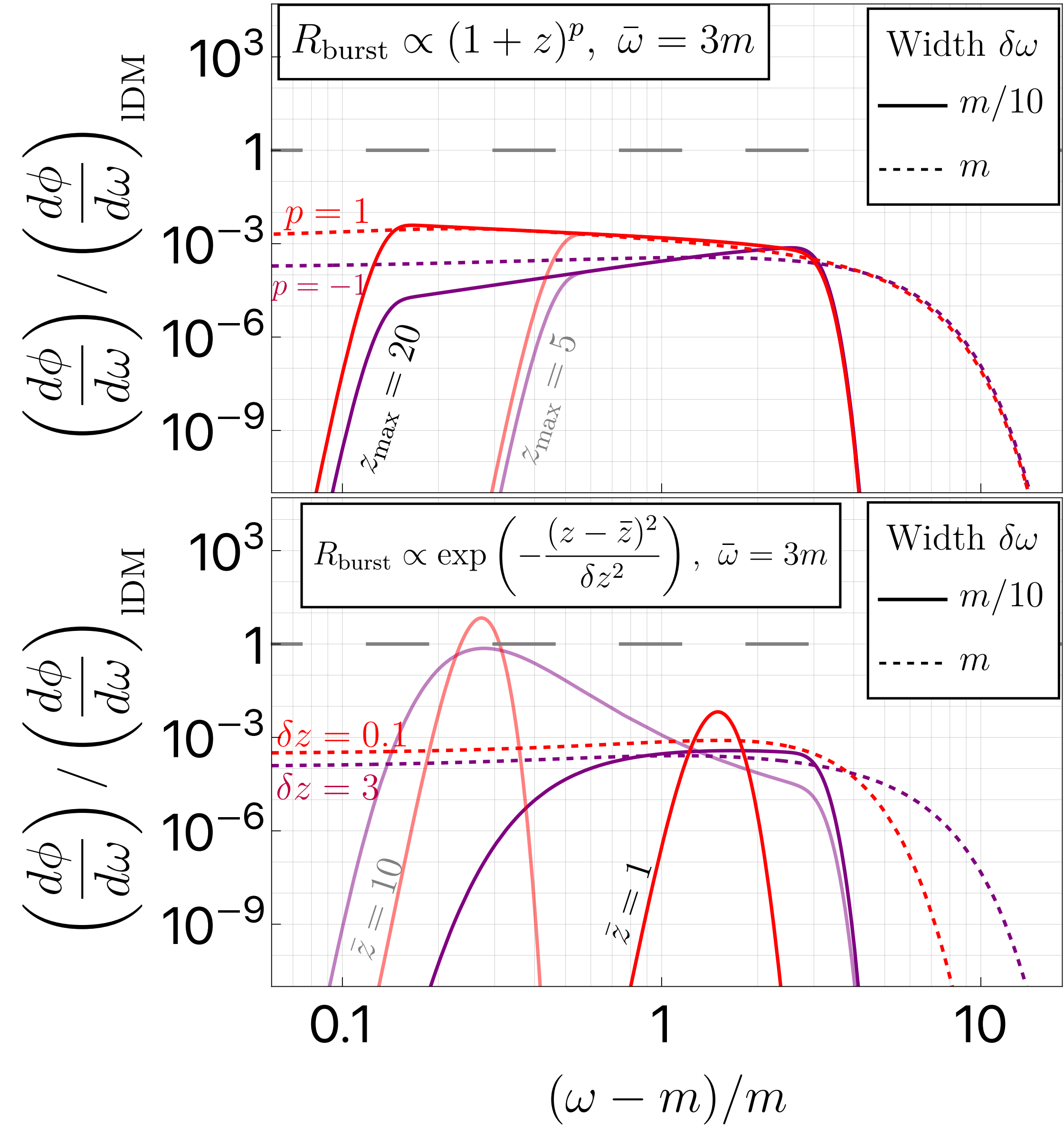
How to search for DaB: (1) direct detection, (2) photon signals, [more to come]

DaB Flux vs DM Flux

Locally, $\left(\frac{d\phi}{d\omega}\right)_{\text{local DM}} \simeq \frac{n_a v_{\text{dm}}}{m_a} \simeq \frac{\rho_{\text{dm}}}{m_a^2} v_{\text{dm}}$

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DaB Flux vs DM Flux

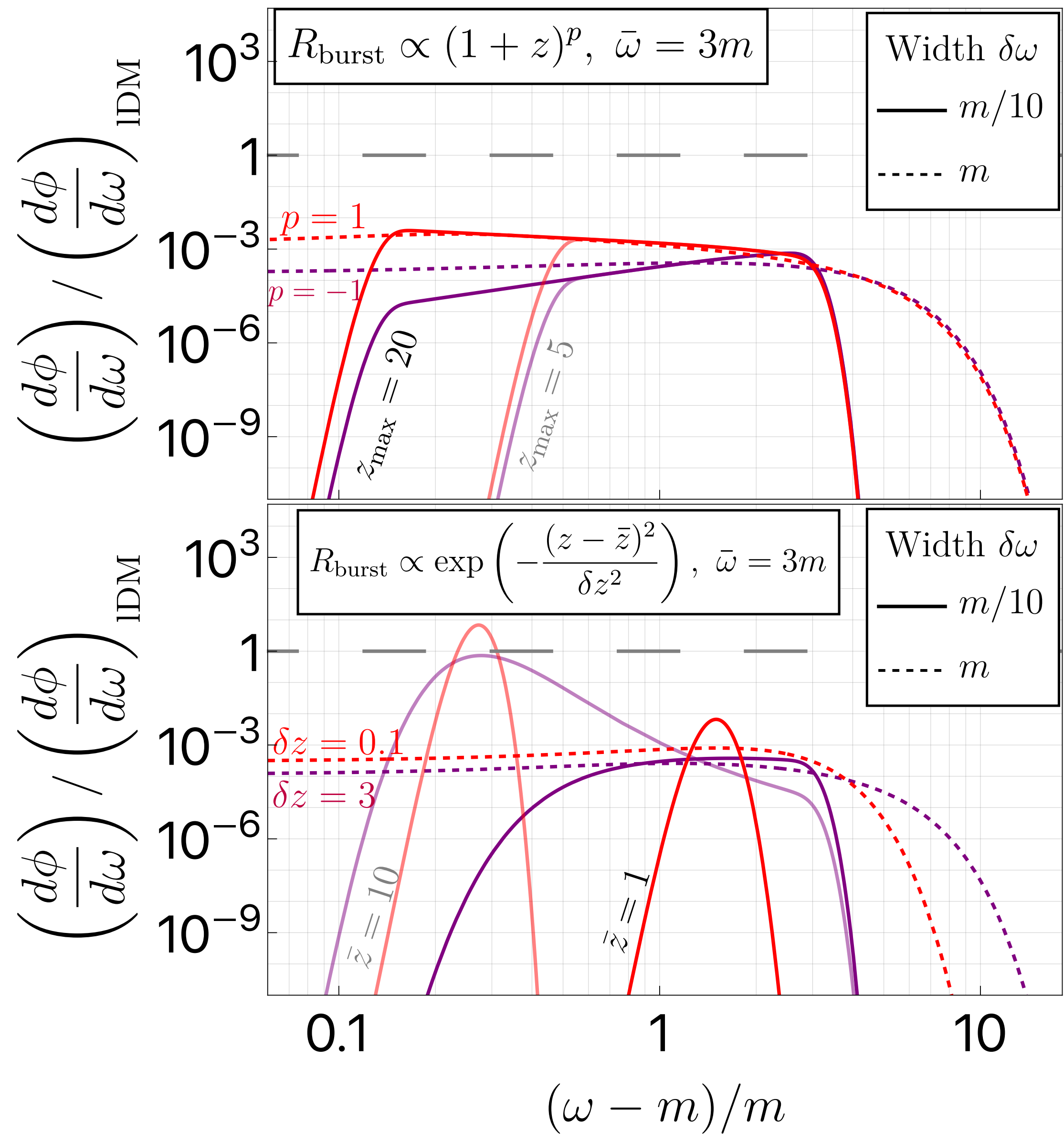
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Parameterise flux and rate

$$\sim \frac{\mathcal{F} \bar{\rho}_U}{m_a \delta\omega} \int dz f(z) \frac{H_0}{H(z)} \exp \left[- \left(\frac{(\omega(1+z) - \bar{\omega})}{\delta\omega} \right)^2 \right]$$



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Locally, $\left(\frac{d\phi}{d\omega}\right)_{\text{local DM}} \simeq \frac{n_a v_{\text{dm}}}{m_a} \simeq \frac{\rho_{\text{dm}}}{m_a^2} v_{\text{dm}}$

DaB flux in present day

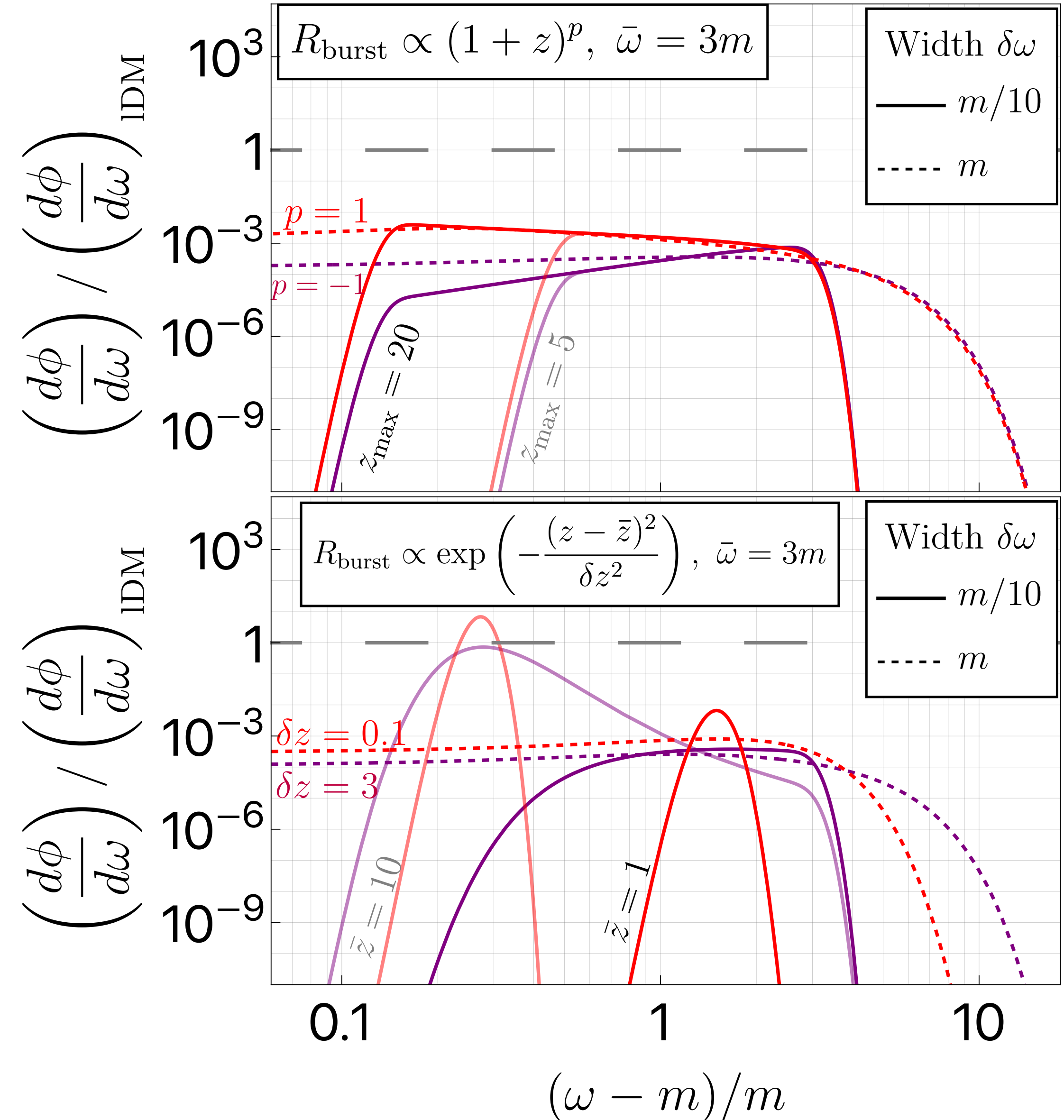
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narrow: $\frac{\delta\omega}{\omega} \rightarrow 0$
recent: $z \sim 0$

$$\sim \frac{\mathcal{F} \bar{\rho}_U}{\bar{\omega}^2}$$



DaB Flux vs DM Flux

Locally, $\left(\frac{d\phi}{d\omega}\right)_{\text{local DM}} \simeq \frac{n_a v_{\text{dm}}}{m_a} \simeq \frac{\rho_{\text{dm}}}{m_a^2} v_{\text{dm}}$

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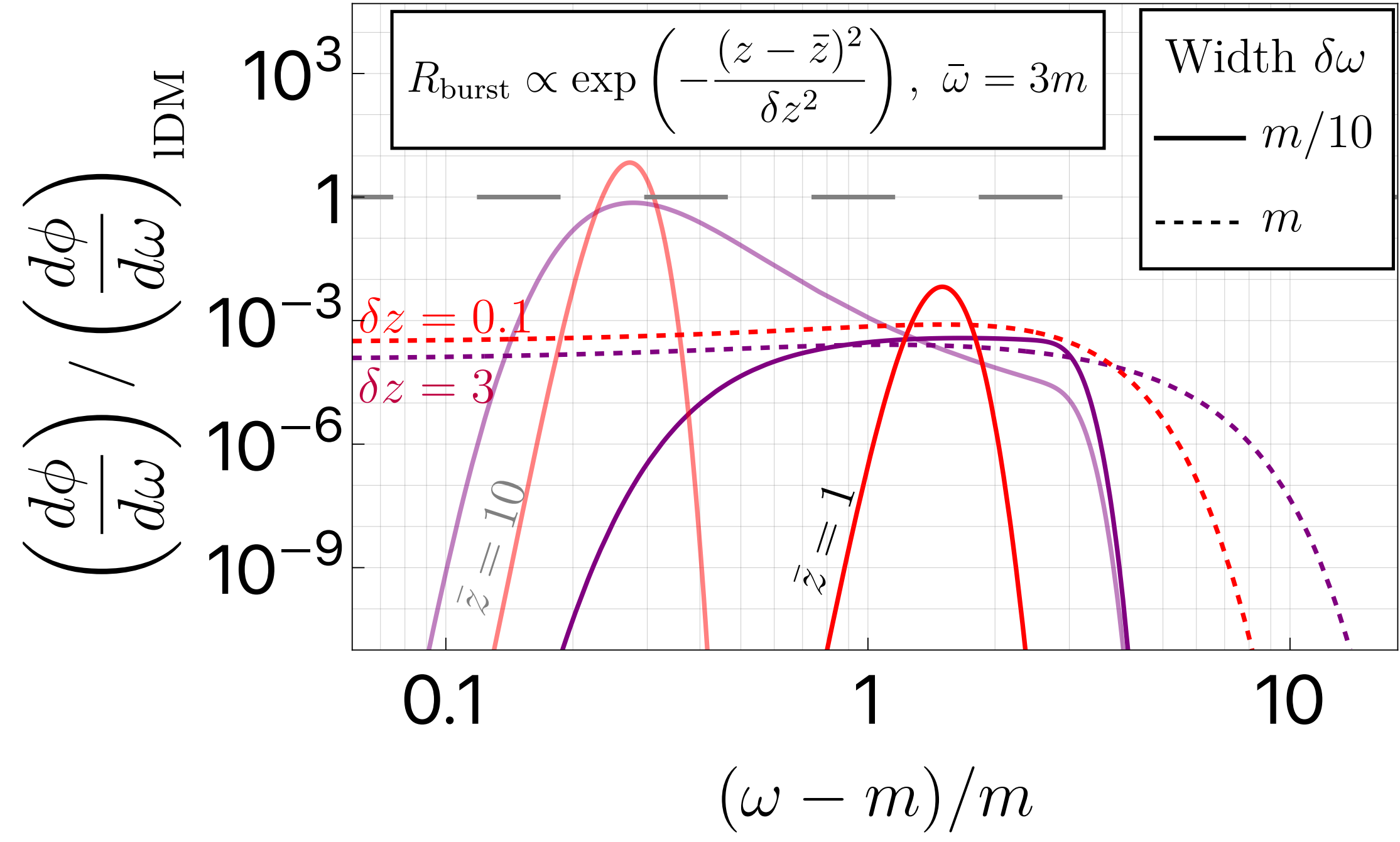
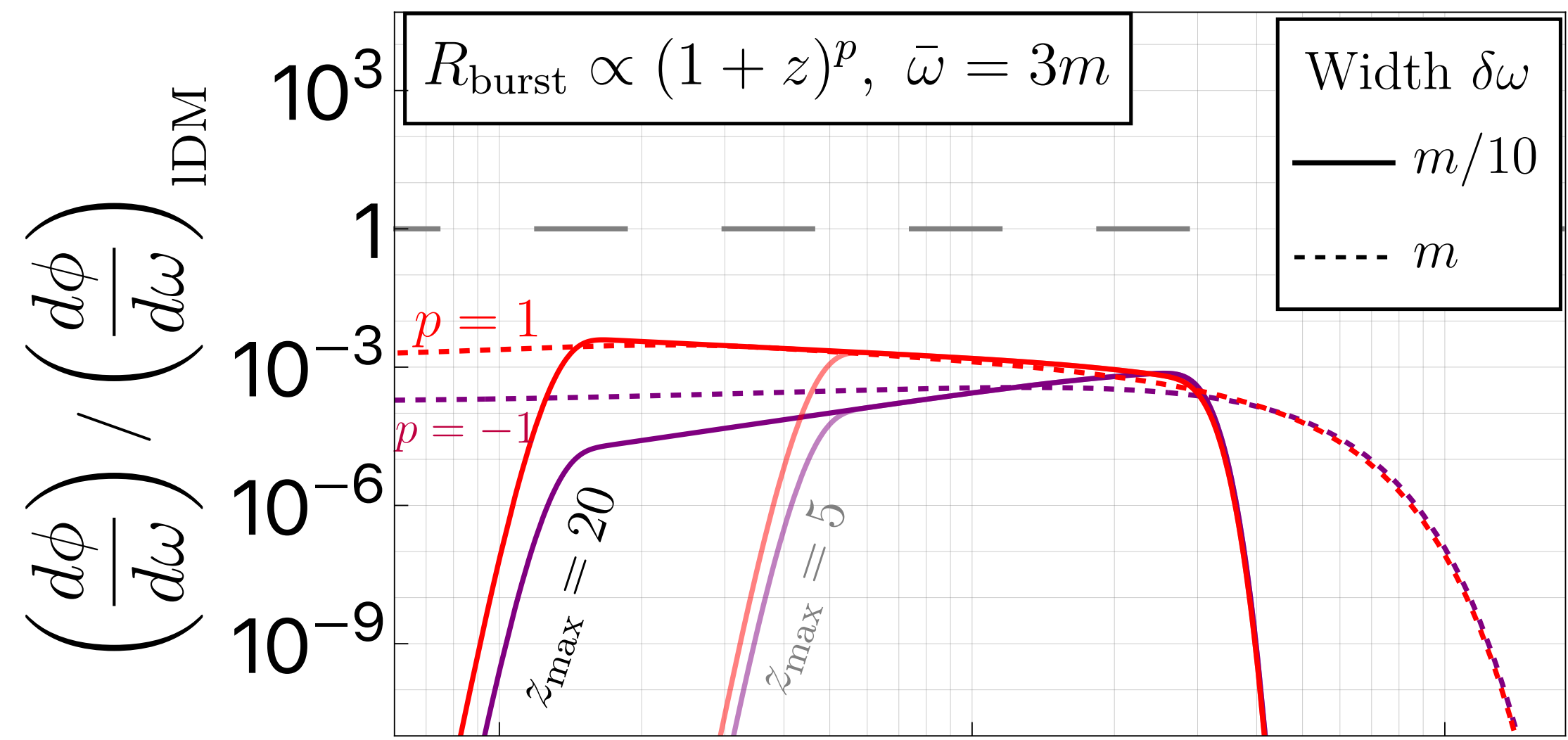
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narrow: $\frac{\delta\omega}{\omega} \rightarrow 0$
recent: $z \sim 0$

$$\frac{d\phi/d\omega}{(d\phi/d\omega)_{\text{IDM}}} \simeq \left(\frac{1}{v_{\text{dm}}}\right) \left(\frac{m_a}{\bar{\omega}}\right)^2 \left(\frac{\mathcal{F} \bar{\rho}_U}{\rho_{\text{dm}}}\right) \simeq 3 \cdot 10^{-3} \mathcal{F} \left(\frac{m_a}{\bar{\omega}}\right)^2$$

↑ (large) ↓ (small) ↓ (small)



Direct Detection

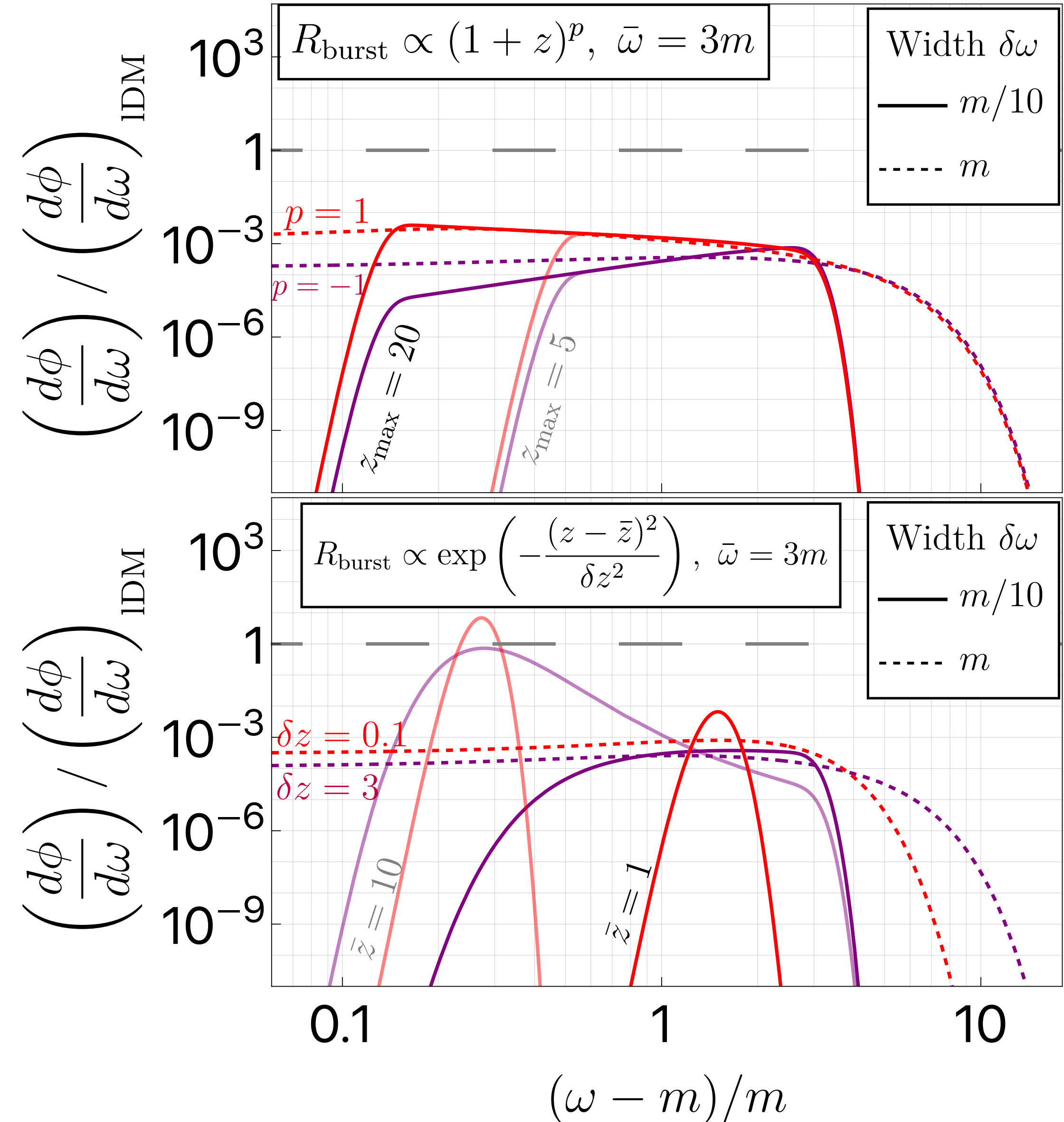
Likely challenging!

- DaB flux generally \lesssim local DM flux
- Signal likely much less coherent than local DM

$$\tau_{\text{coh}} \simeq \frac{2\pi}{m_a v^2}, \quad v_{\text{dm}} \sim 10^{-3} \text{ vs } v_{\text{DaB}} \sim 1$$

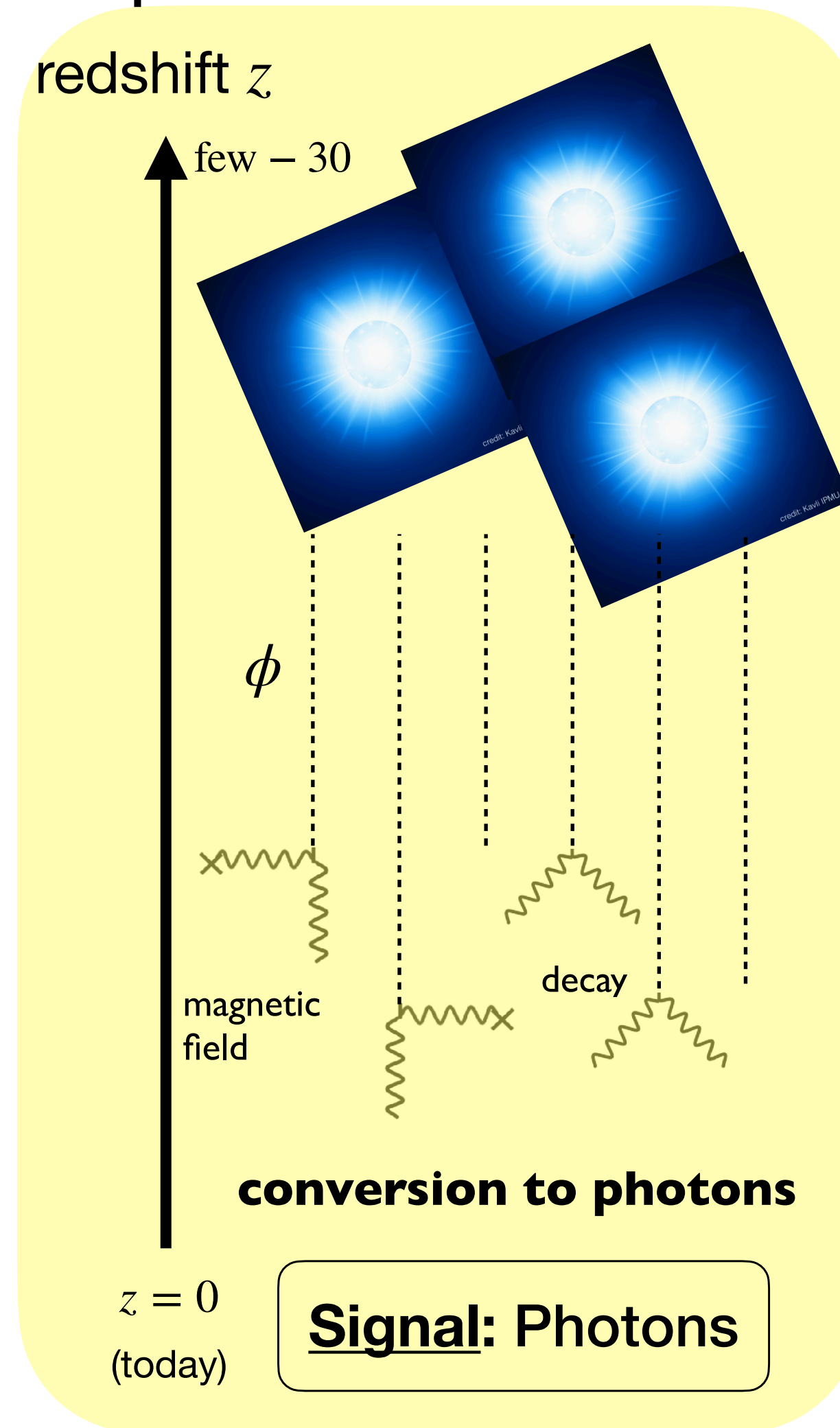
Worth investigating!

- Nontrivial energy distribution encodes cosmological evolution and source properties
- Can also encode information about fundamental axion potential, e.g. self-interactions



Photon Signals from DaB

$$\mathcal{L} \supset \frac{1}{4} g_{a\gamma} a F_{\mu\nu} \tilde{F}^{\mu\nu} = g_{a\gamma} a \mathbf{E} \cdot \mathbf{B}$$

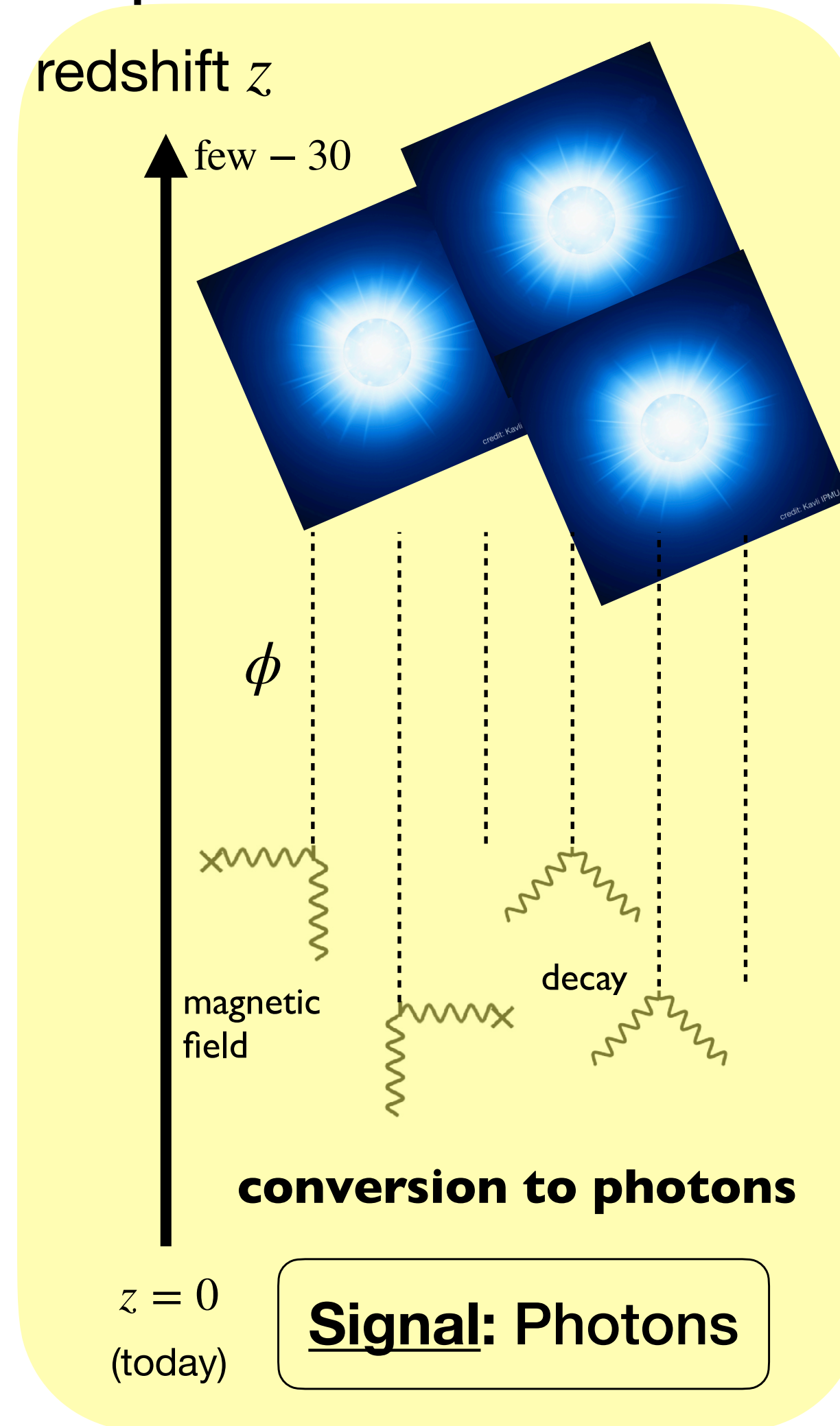
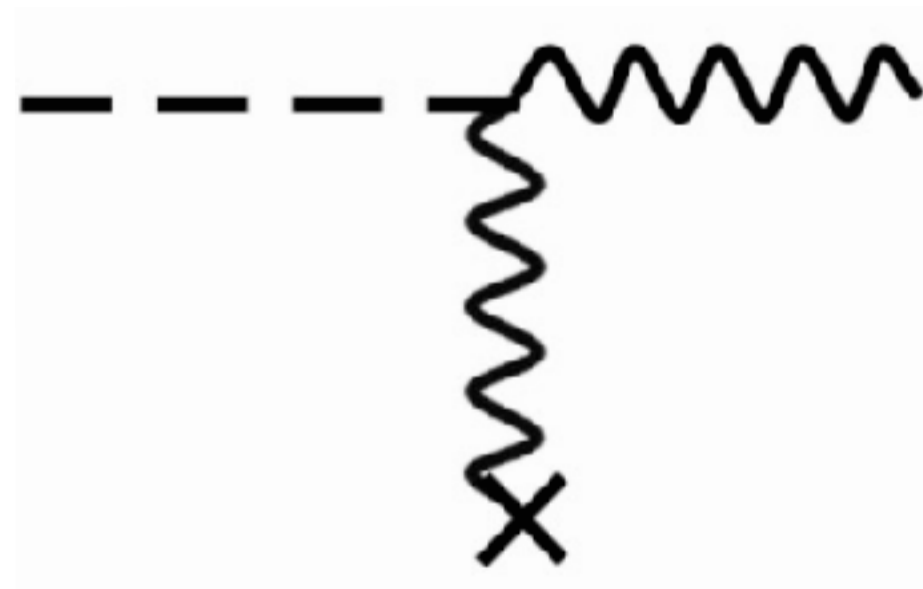


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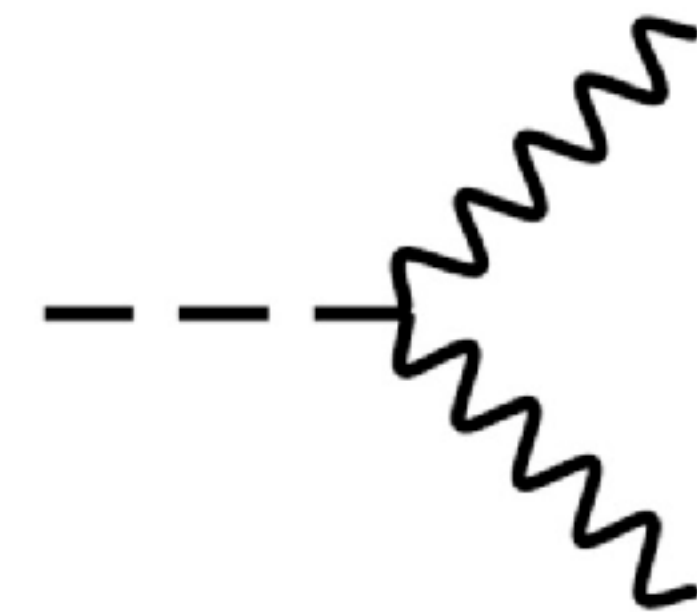
Magnetic field conversion

$$\left. \frac{d\phi_\gamma}{d\omega} \right|_{B\text{-field}} = P_{\gamma \rightarrow a} \frac{d\phi}{d\omega}$$



Axion decay to photons

$$\left. \frac{d\phi_\gamma}{d\omega} \right|_{\text{decay}} \simeq P_{\text{decay}} \frac{d\phi}{d\omega}$$

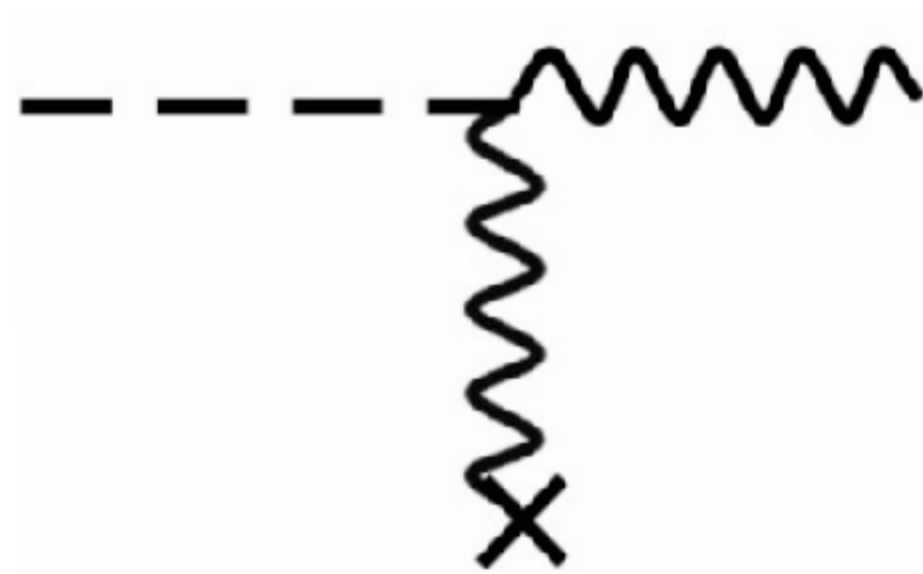


Photon Signals from DaB

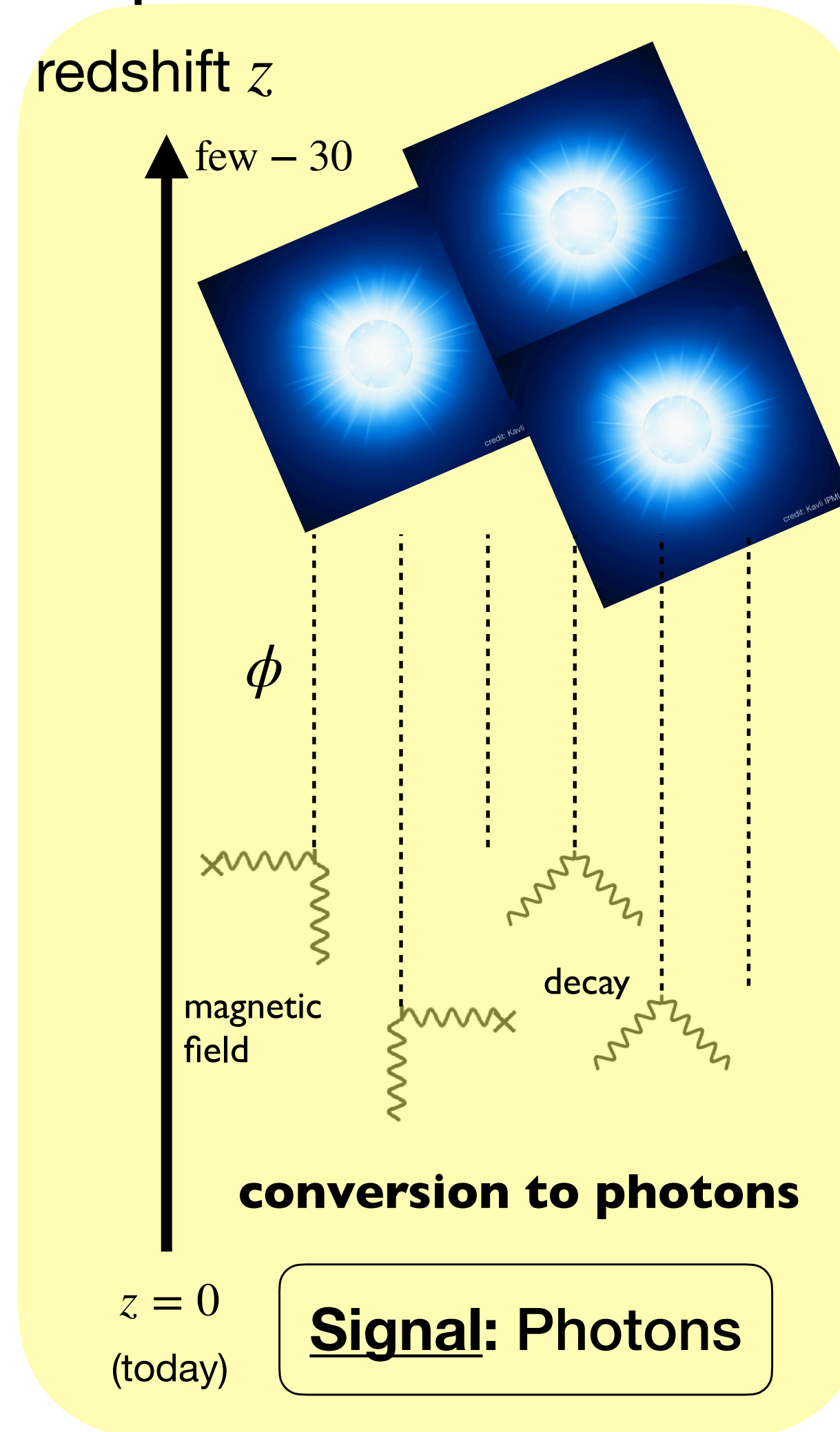
$$\mathcal{L} \supset \frac{1}{4} g_{a\gamma} a F_{\mu\nu} \tilde{F}^{\mu\nu} = g_{a\gamma} a \mathbf{E} \cdot \mathbf{B}$$

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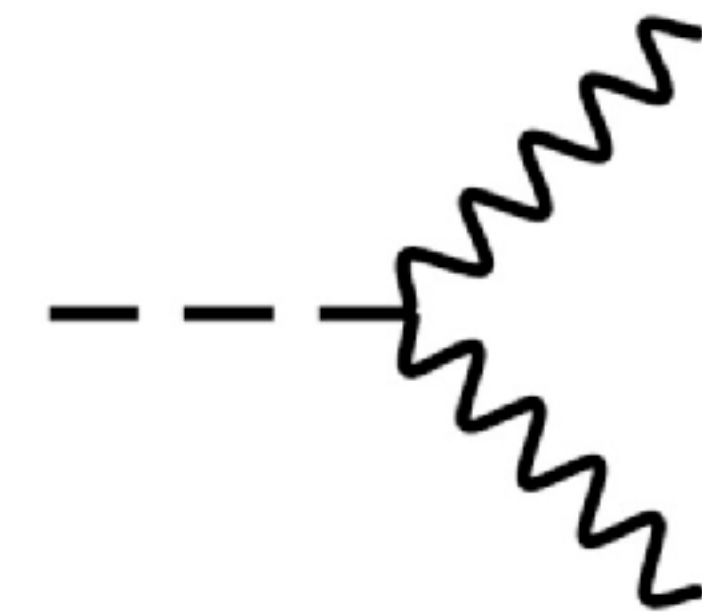


- Galactic magnetic fields of $\sim \mu\text{G}$ dominate (typical distances $\sim \text{kpc} - \text{Mpc}$)
- $P_{\gamma \rightarrow a}$ grows with large ω and small m_a
 \Rightarrow largest when $\omega \gg m_a$ with small m_a



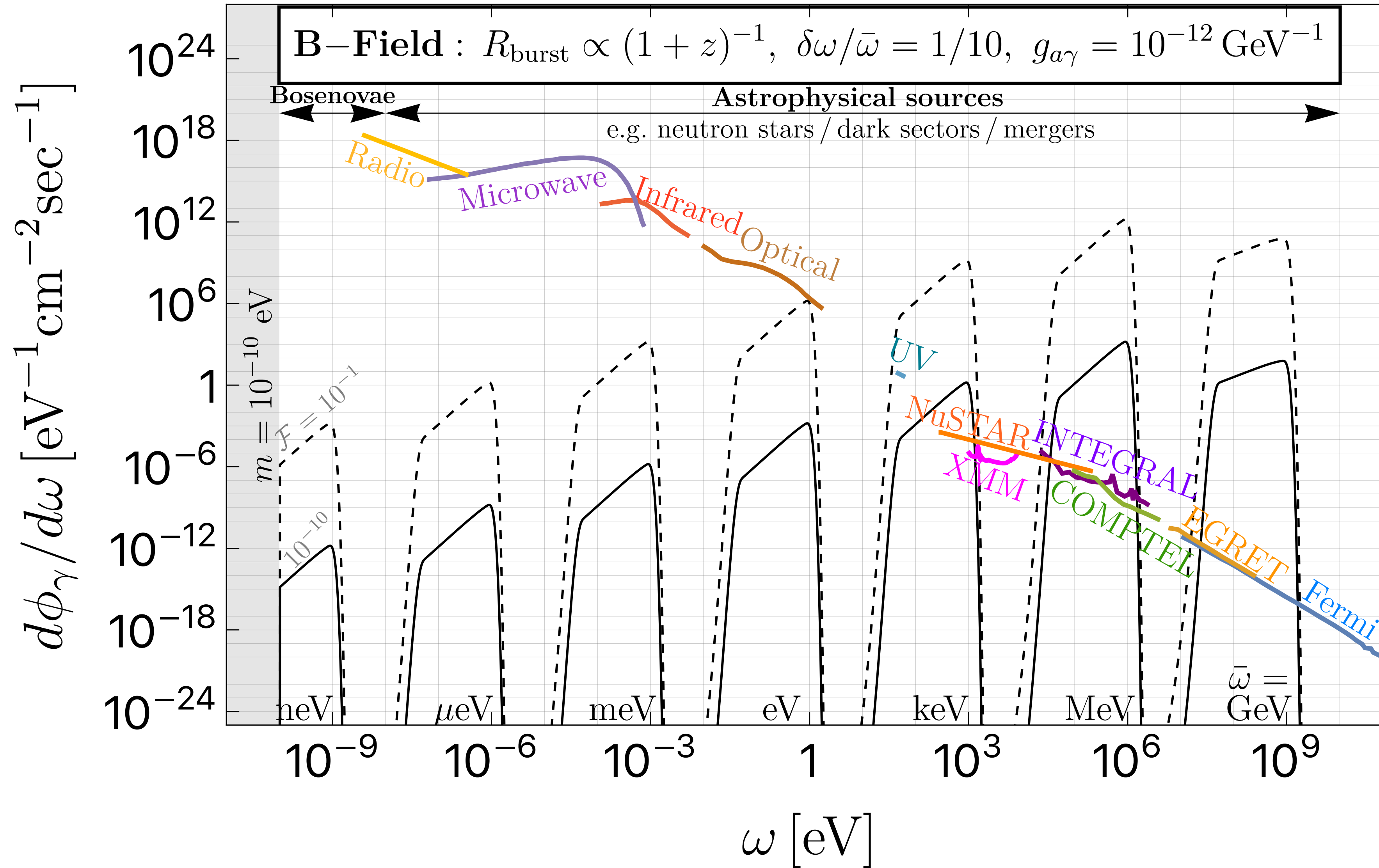
Axion decay to photons

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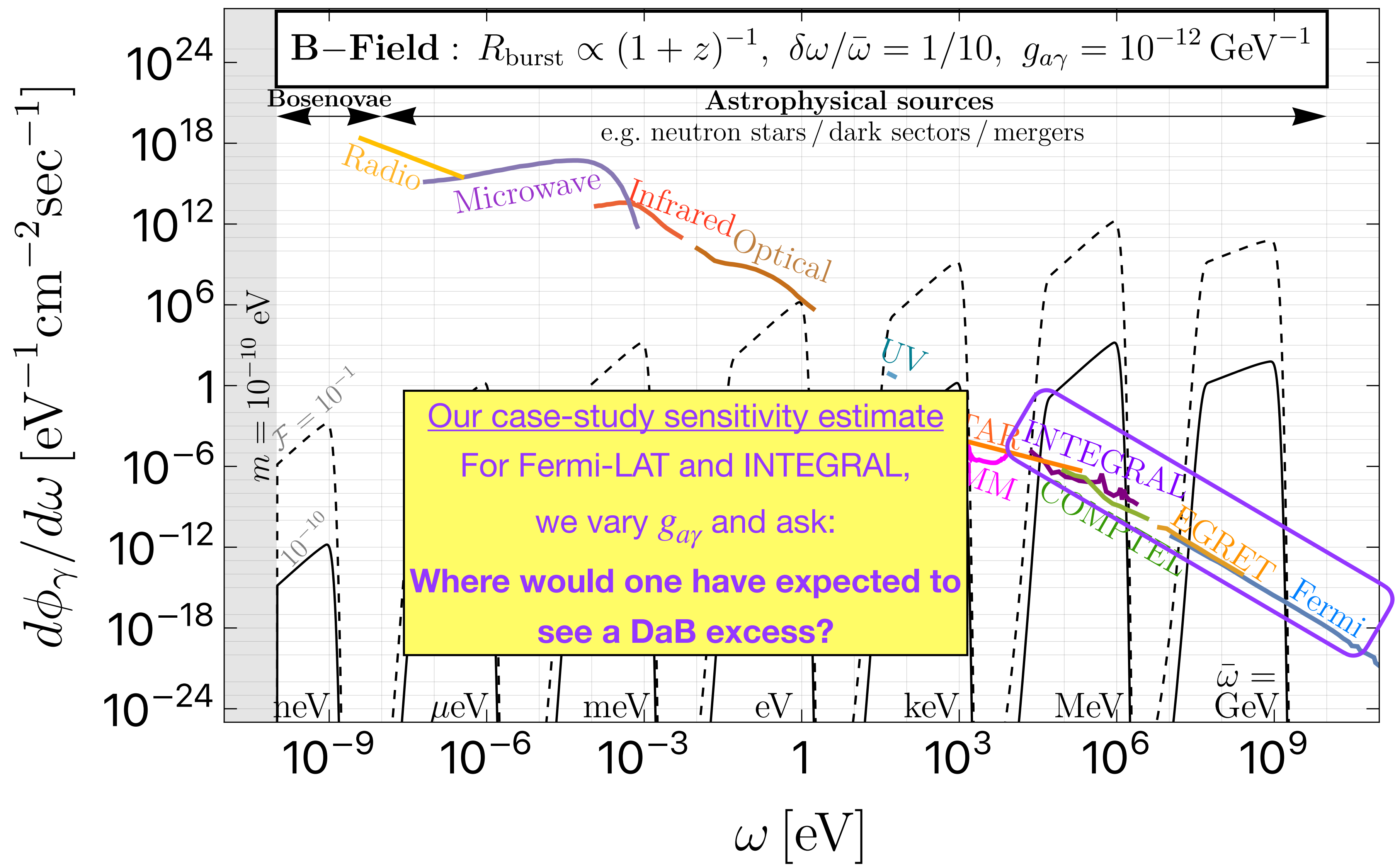


- Decay can occur anywhere in space (typical distances $\sim \text{Gpc}$)
- P_{decay} grows with small ω and large m_a
 \Rightarrow largest when $\omega \gtrsim m_a$ with large m_a

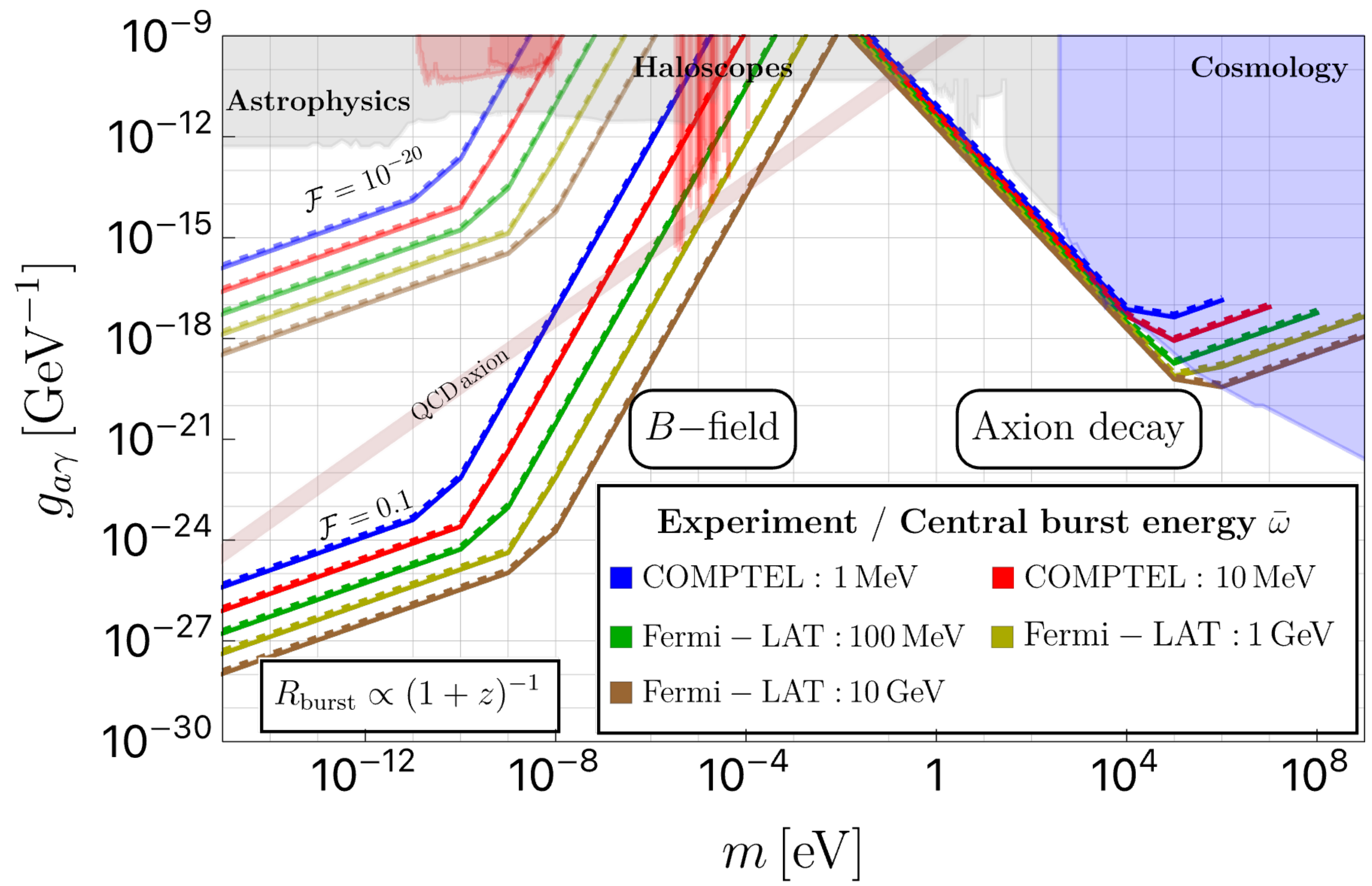
Where to Search: Today



Where to Search: Today



Searches for DaB Gamma-Rays

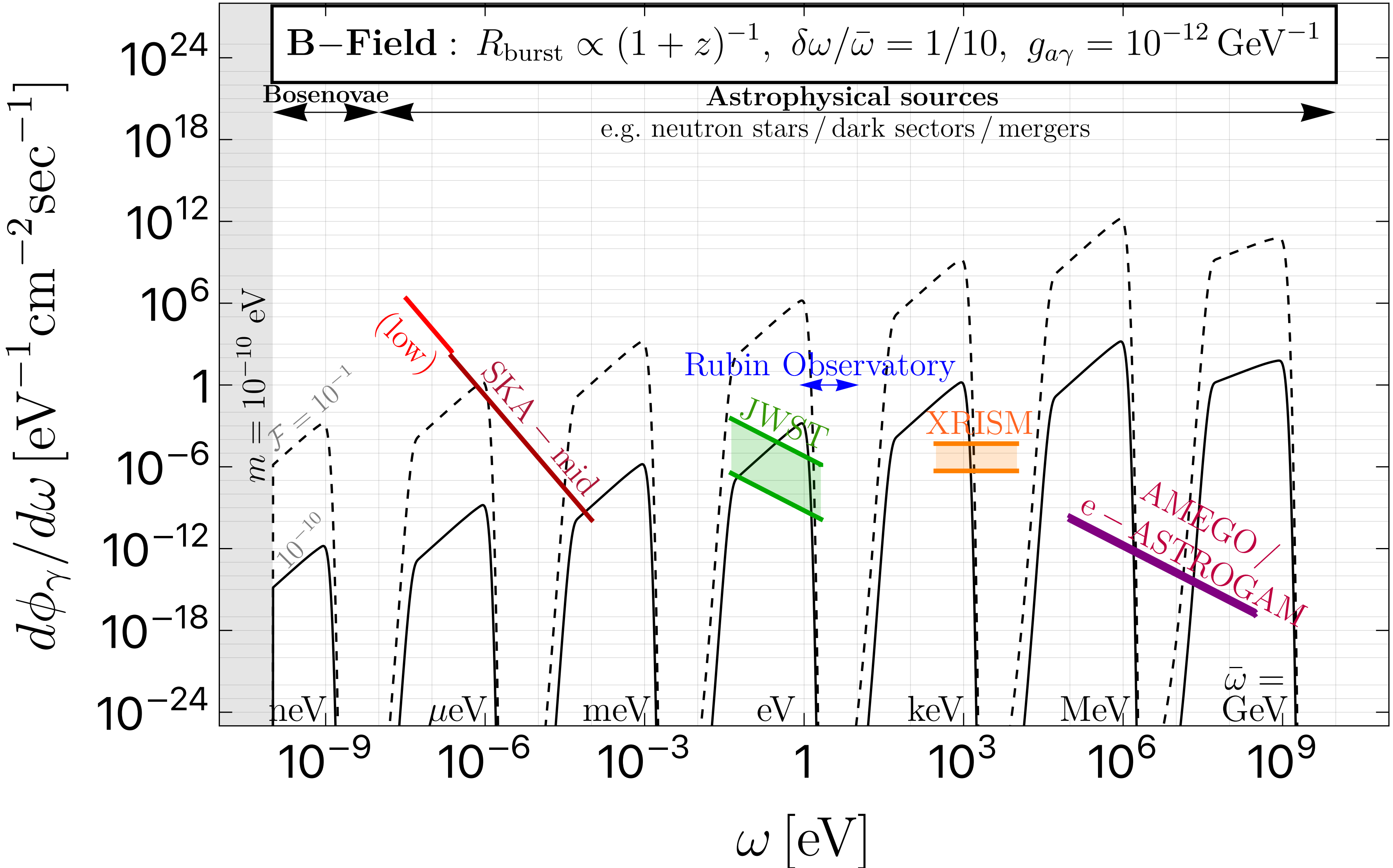


A very tiny energy fraction in DaB can give rise to striking signals!

Best sensitivity when $\bar{\omega} \gg m_a$

What about low-energy signals?

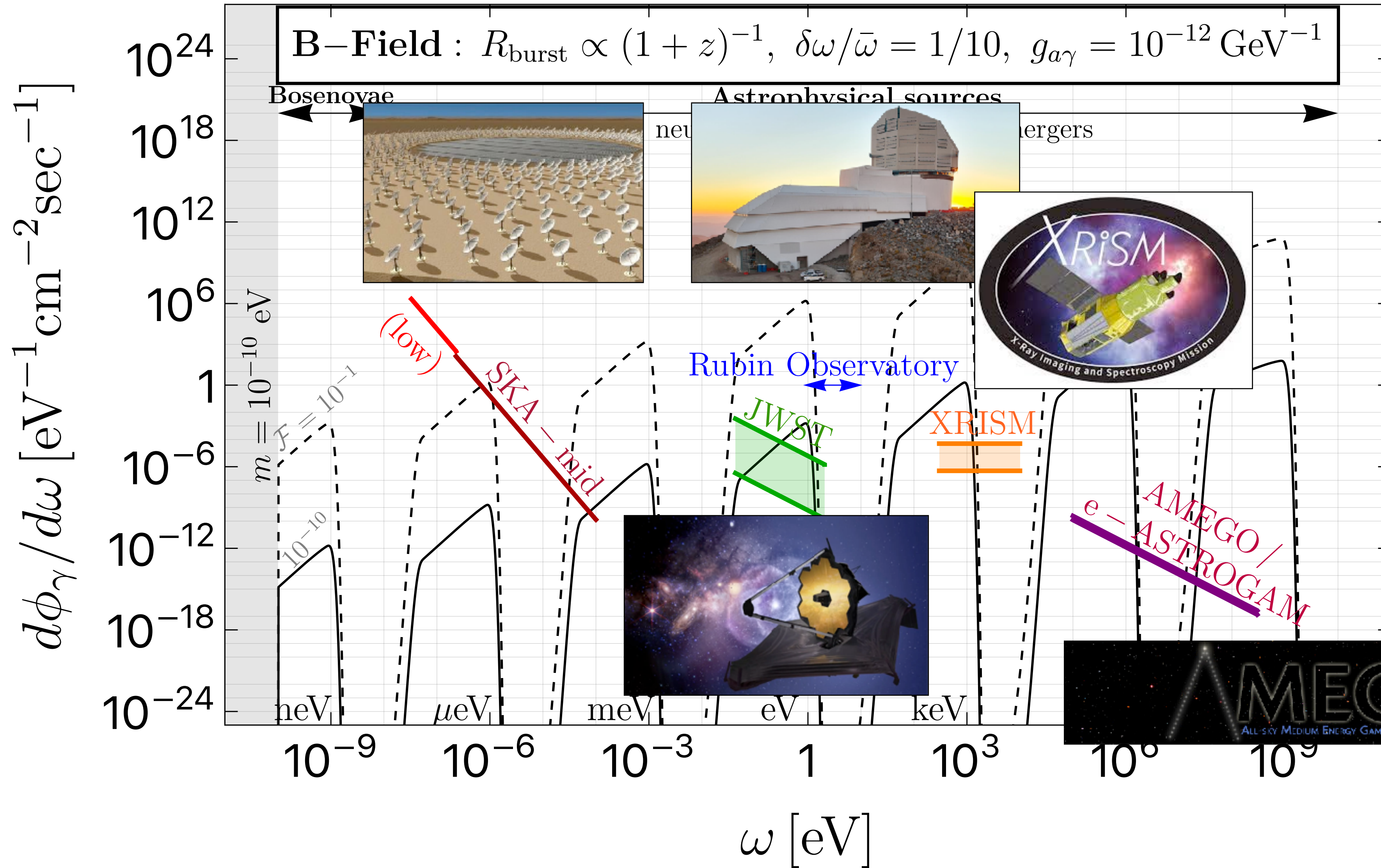
Where to Search: Future



Next-Gen searches will see huge improvements, especially at low energy

Future: search even for semi-relativistic burst sources of DaB, like bosenovae!

Where to Search: Future



Next-Gen searches will see huge improvements, especially at low energy

Future: search even for semi-relativistic burst sources of DaB, like bosenovae!

Conclusions

- [Relativistic bursts of axions](#) commonly originate in astrophysical processes, both in SM / dark sectors; give rise to a [diffuse axion background \(DaB\)](#)
- Axions emitted in recent transient bursts, e.g. from *bozenovae* in boson star collapse, [lead to direct-detection signals which can exceed sensitivity of local DM searches](#). At present, **detection viable but exclusion difficult**; need stronger prediction of burst rate
- The DaB encodes novel information about cosmology and burst sources, [implying complementarity with existing DM searches](#). Direct detection difficult but promising!
- [Existing photon-search experiments \(e.g. Fermi\) can already constrain DaB](#) using photon couplings. In progress: investigate signals from other couplings, e.g. electrons and quarks
- **Further characterization of burst sources is worthy of dedicated study, strong discovery potential!**

Thank you for your attention!



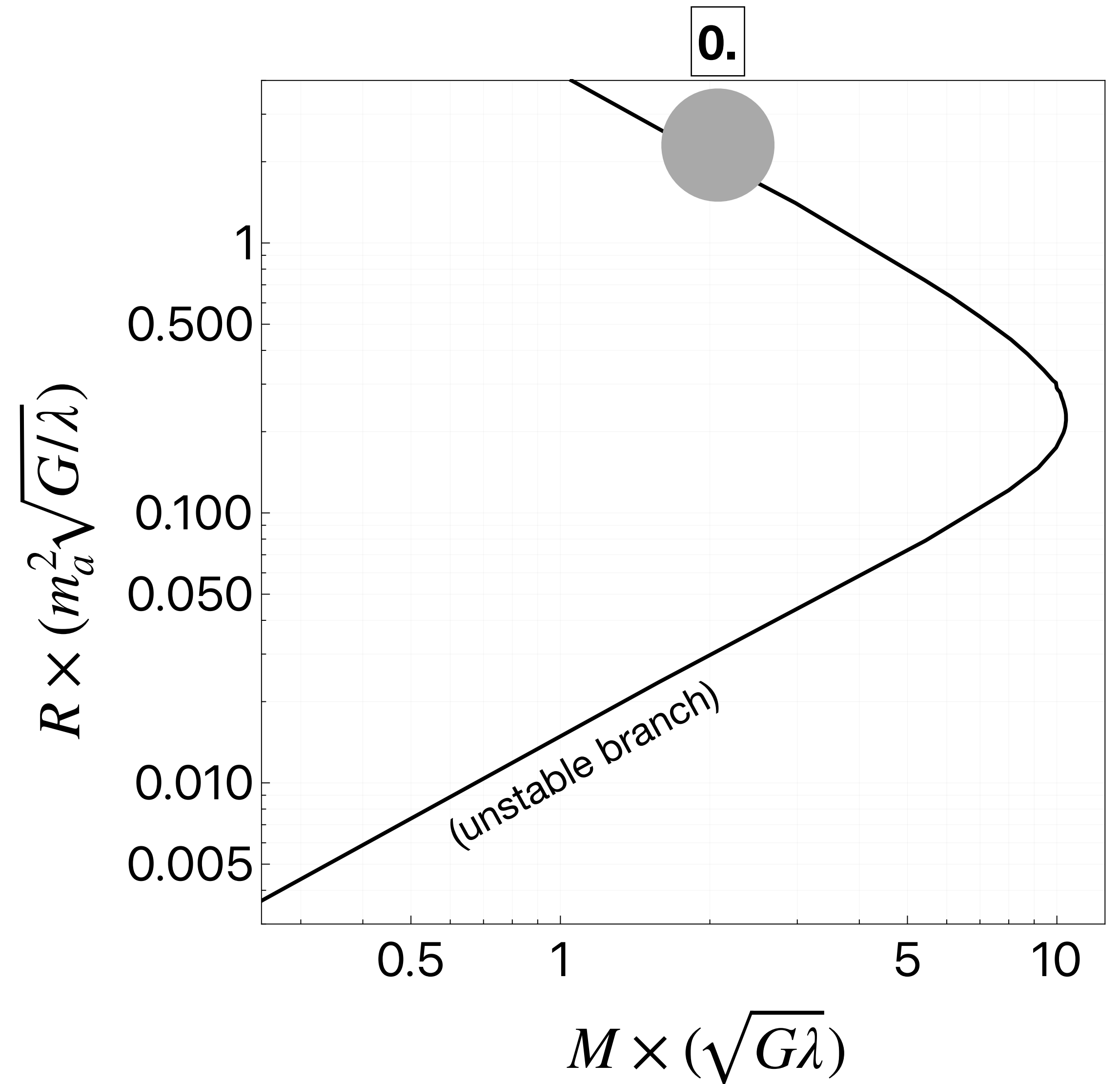
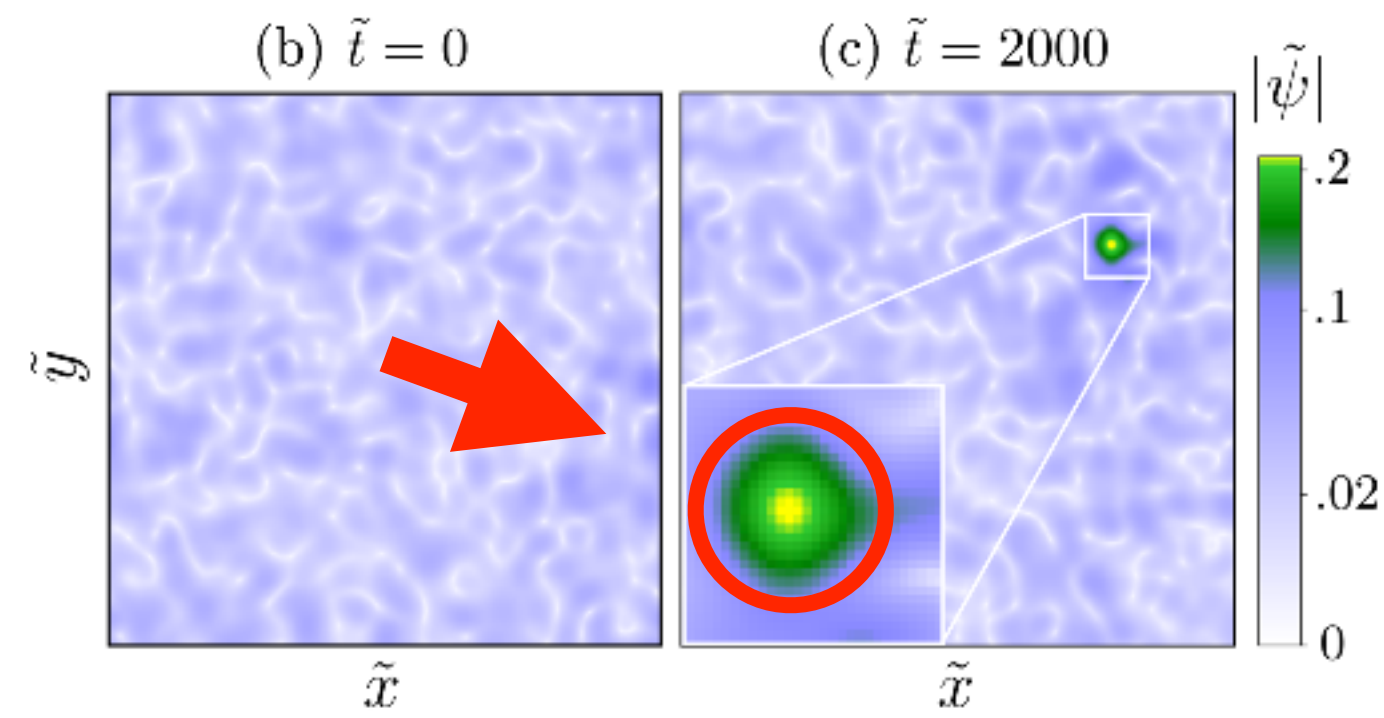
DALL-E 3 illustration:
"Athens symposium on
Exploring the Universe"

Backup Slides



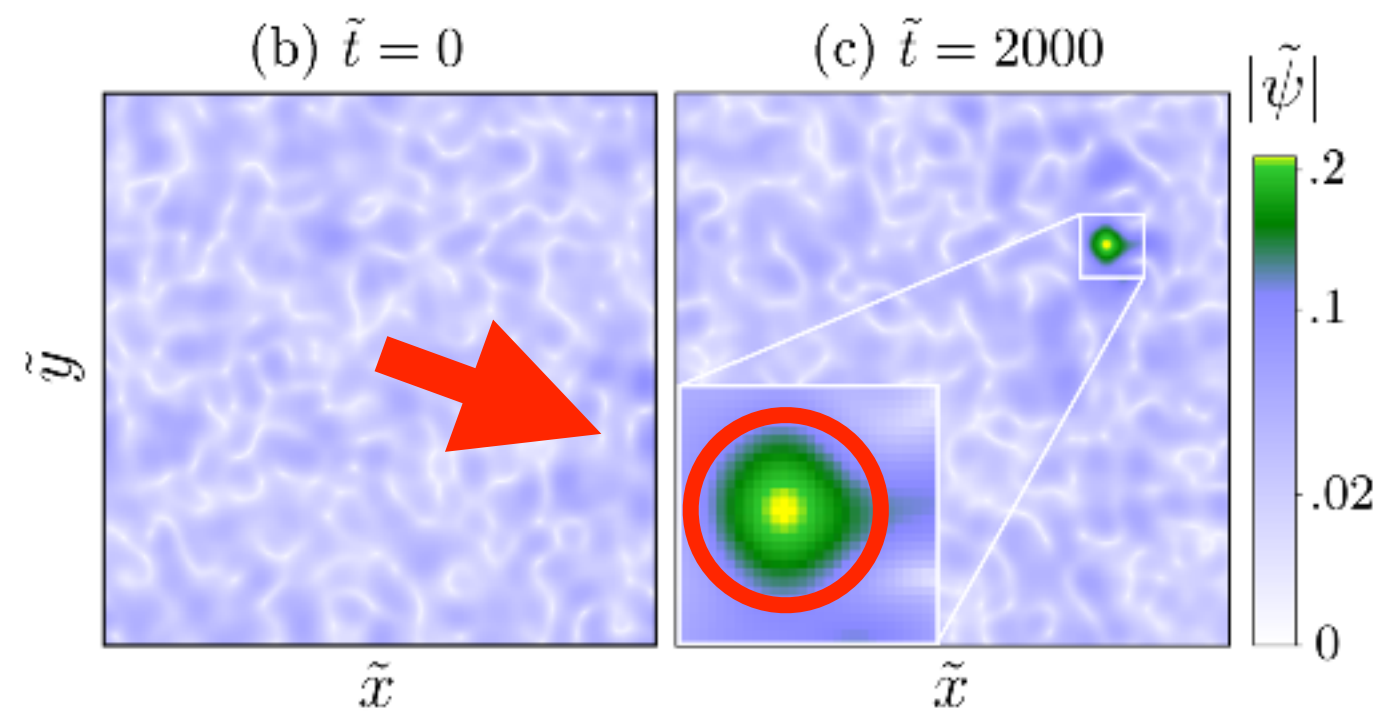
Life Cycle of a Boson Star

0. Birth

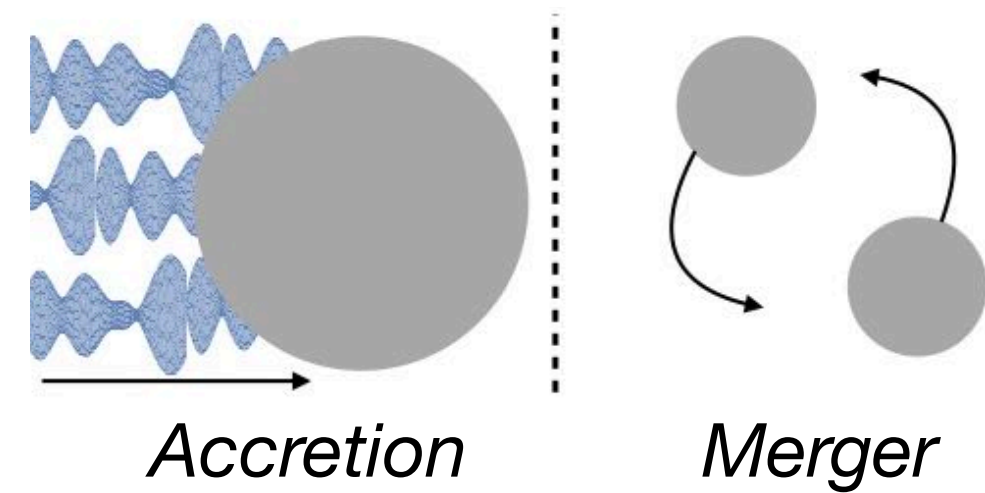


Life Cycle of a Boson Star

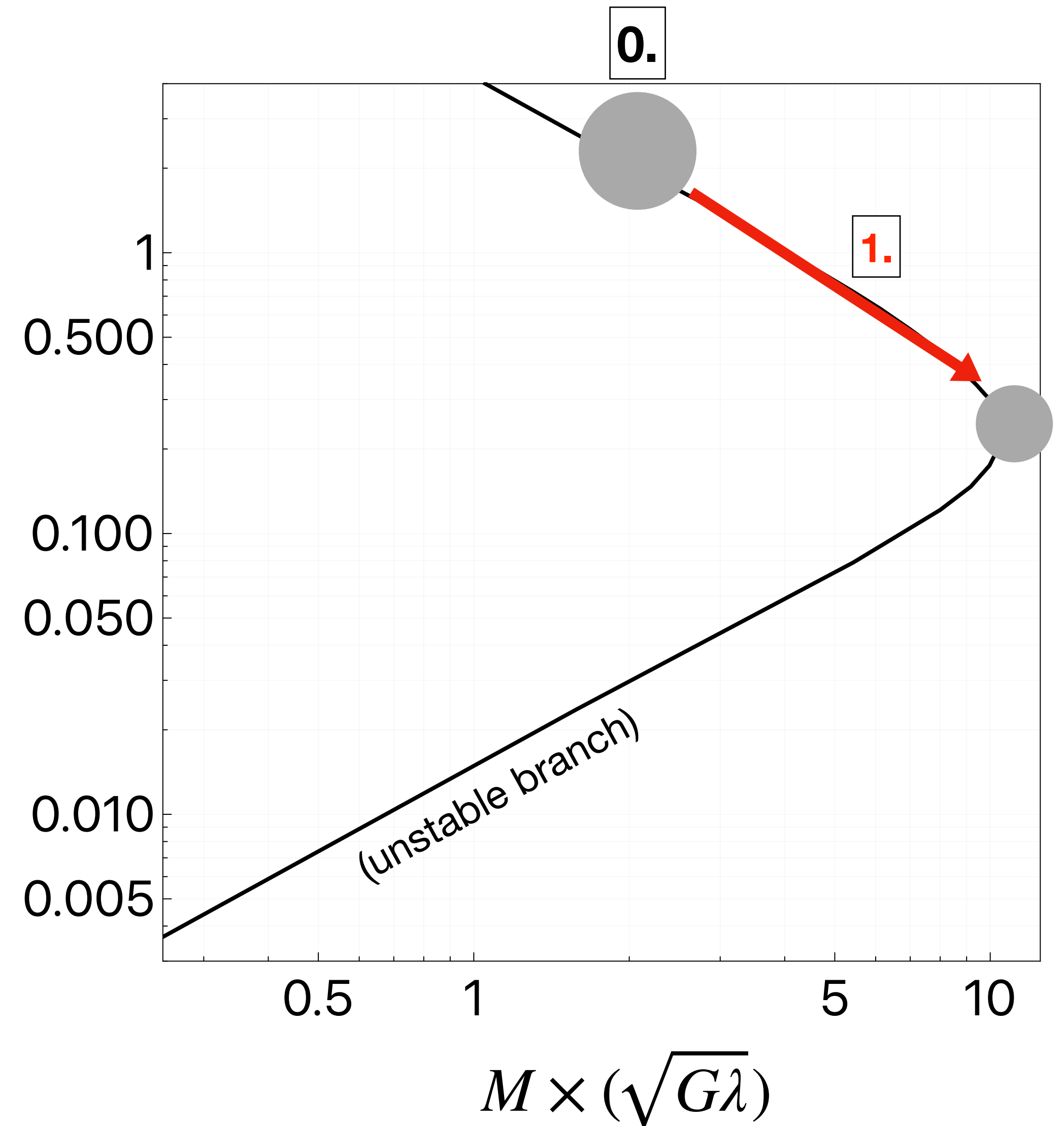
0. Birth



1. Growing up

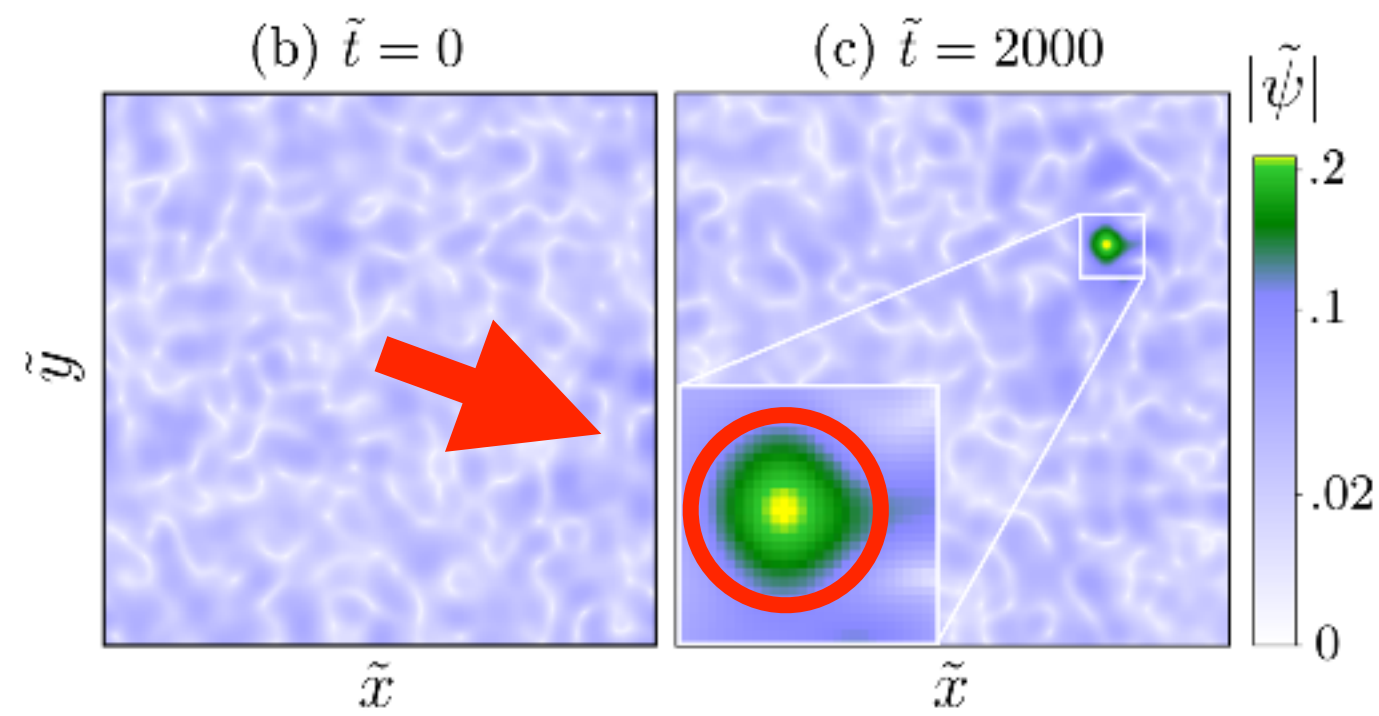


$R \times (m_a^2 \sqrt{G/\lambda})$

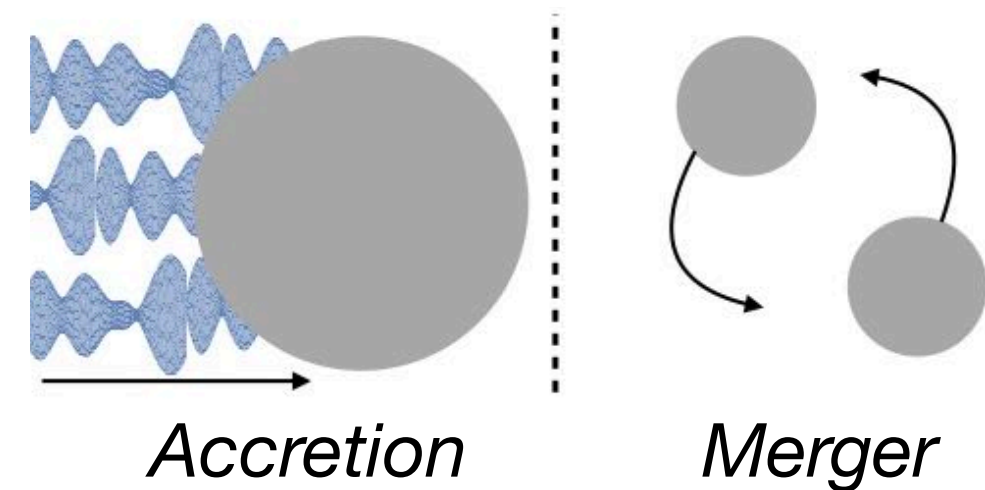


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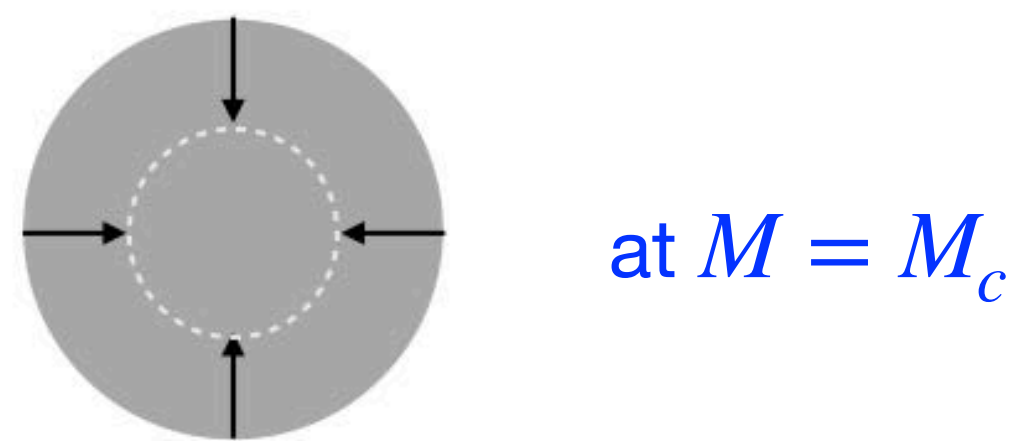
0. Birth



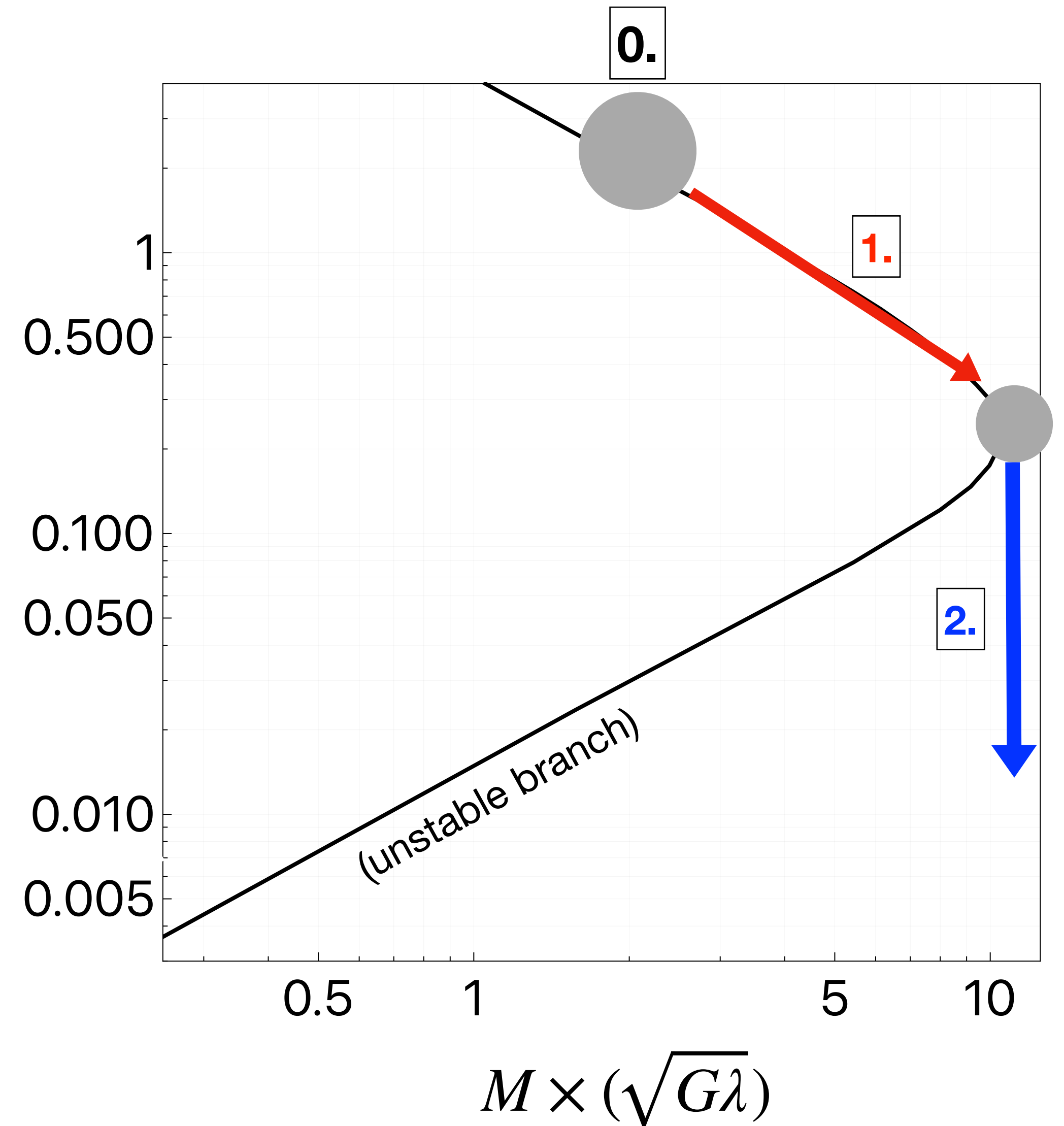
1. Growing up



2. Gravitational collapse

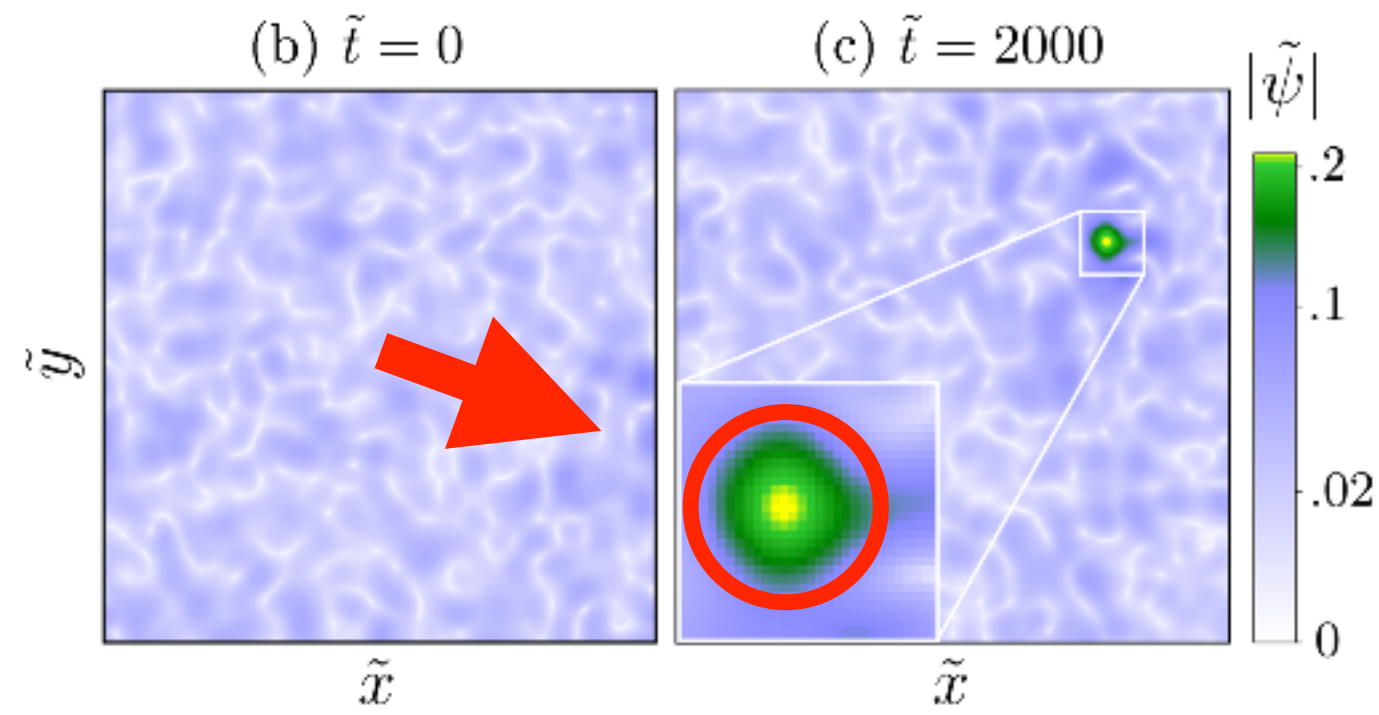


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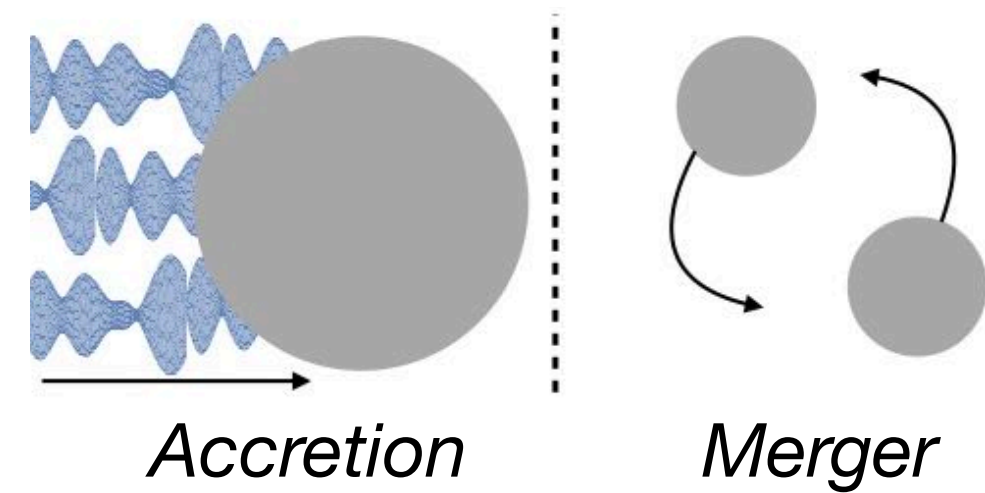


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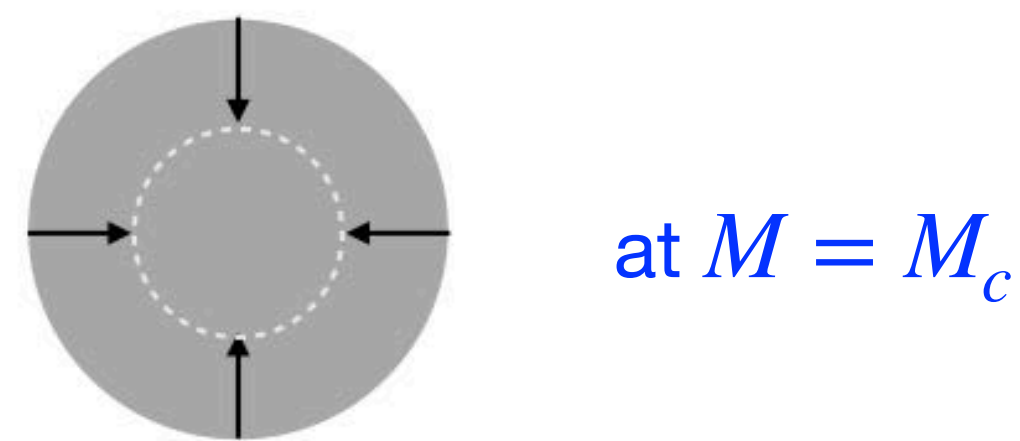
0. Birth



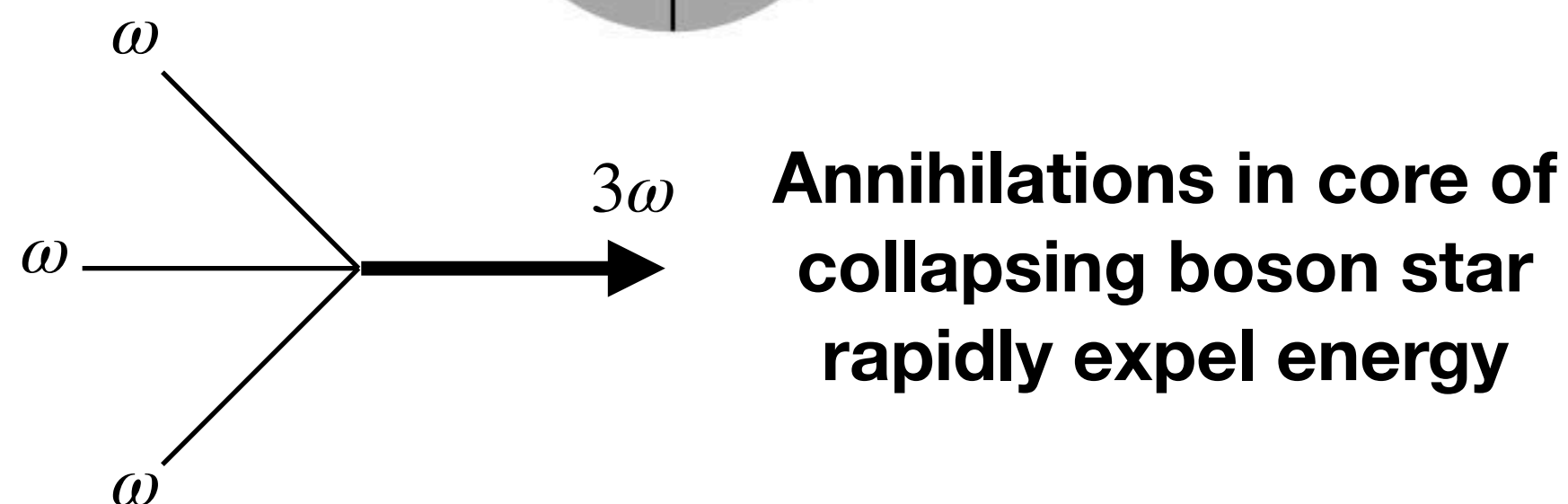
1. Growing up



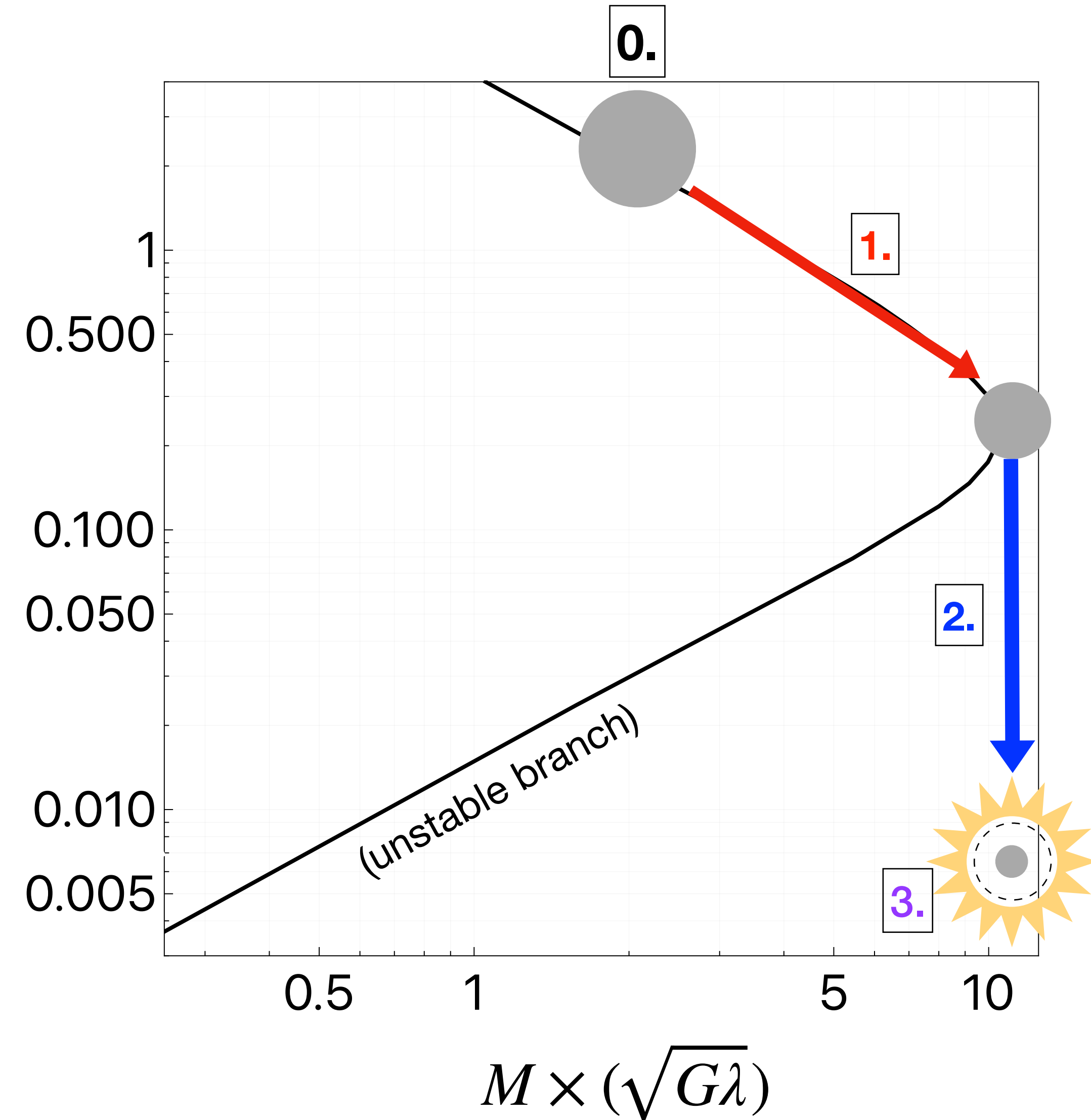
2. Gravitational collapse



3. "Bosenova"

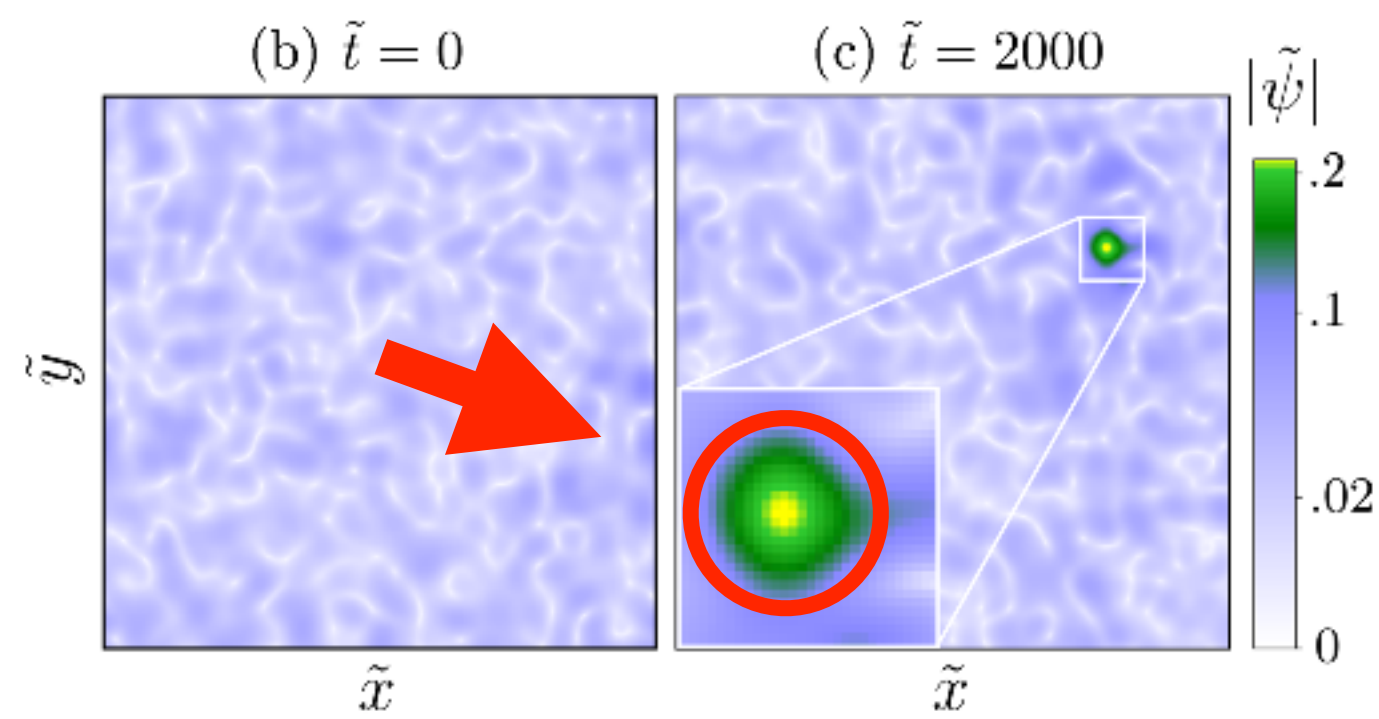


$R \times (m_a^2 \sqrt{G/\lambda})$

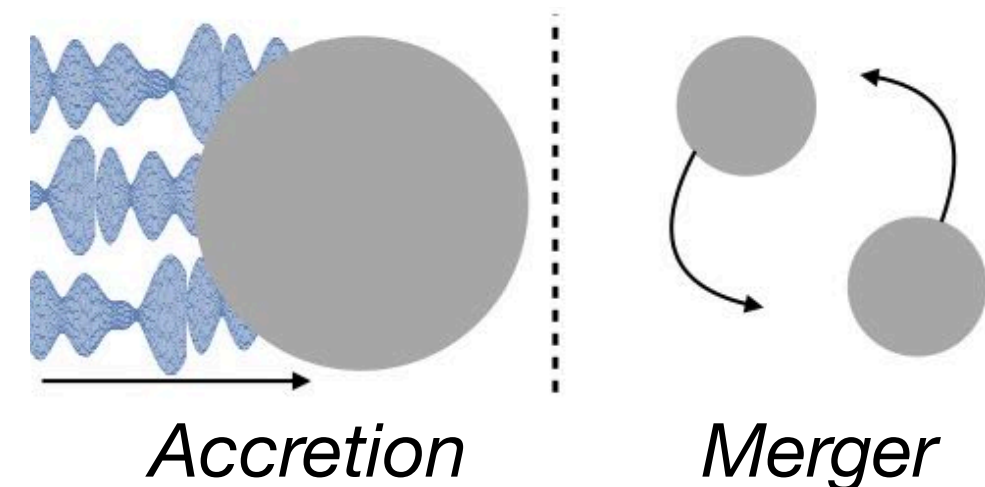


Life Cycle of a Boson Star

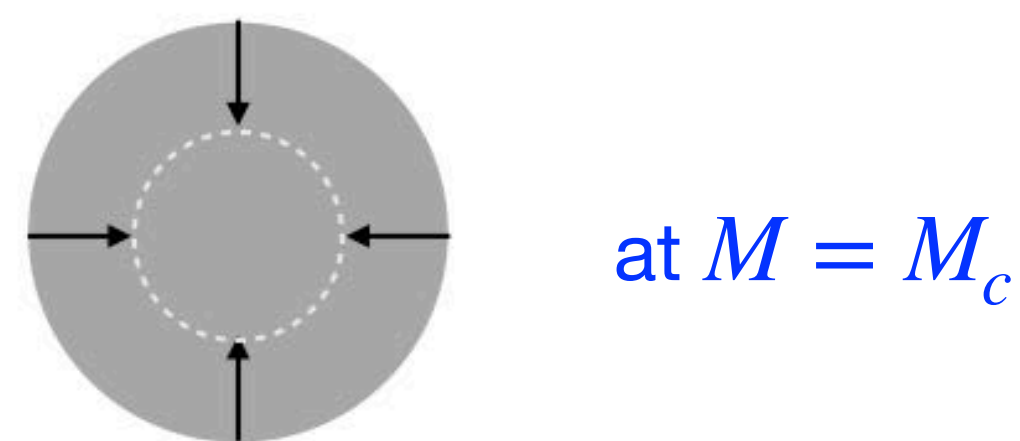
0. Birth



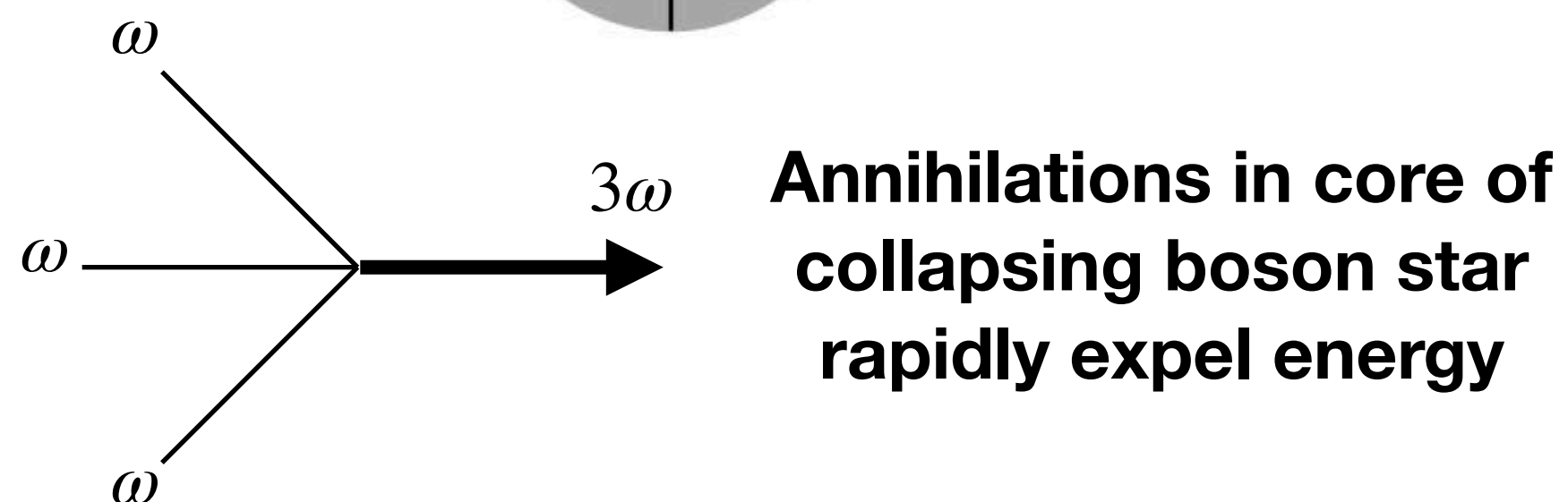
1. Growing up



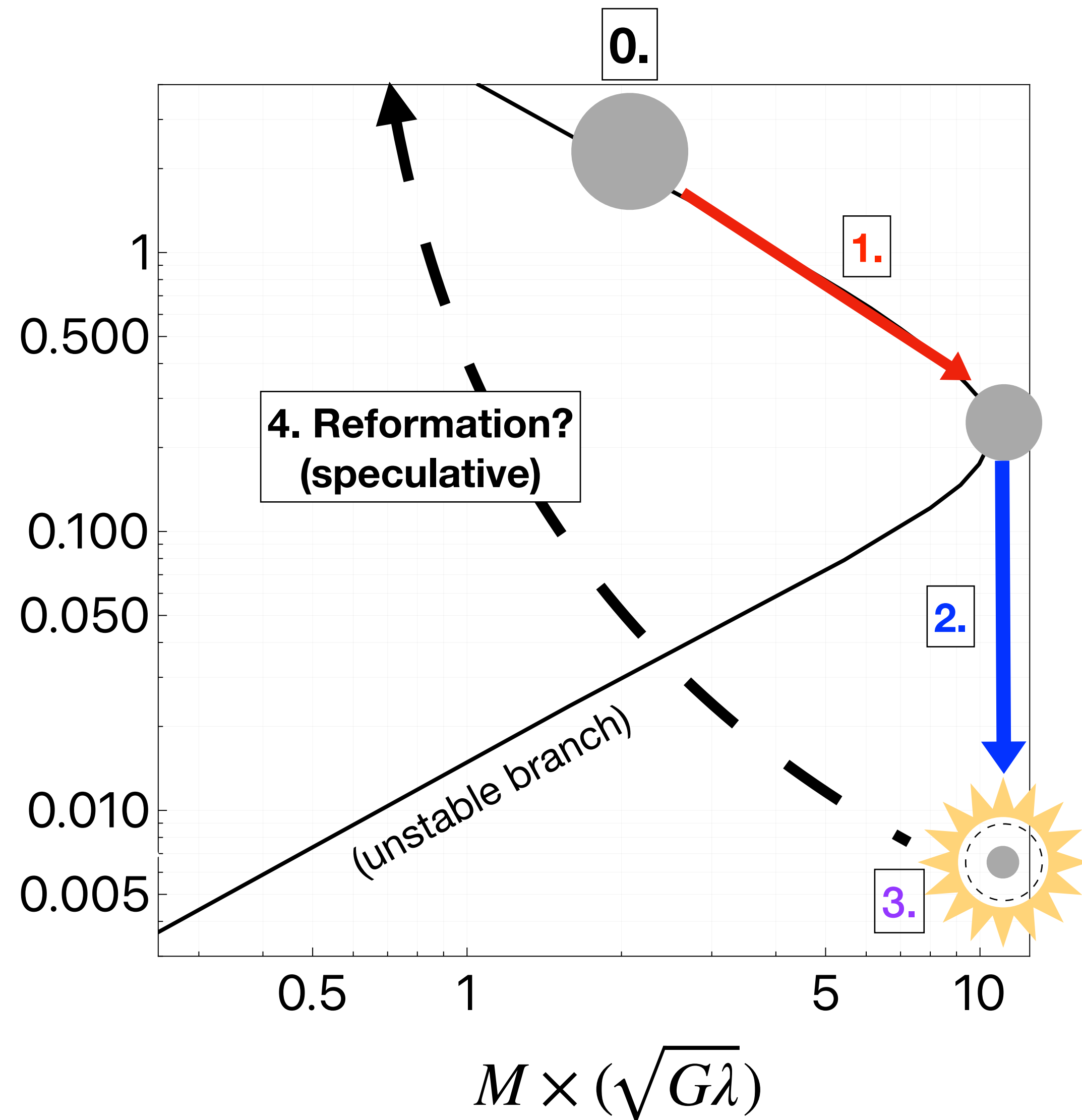
2. Gravitational collapse



3. "Bosenova"

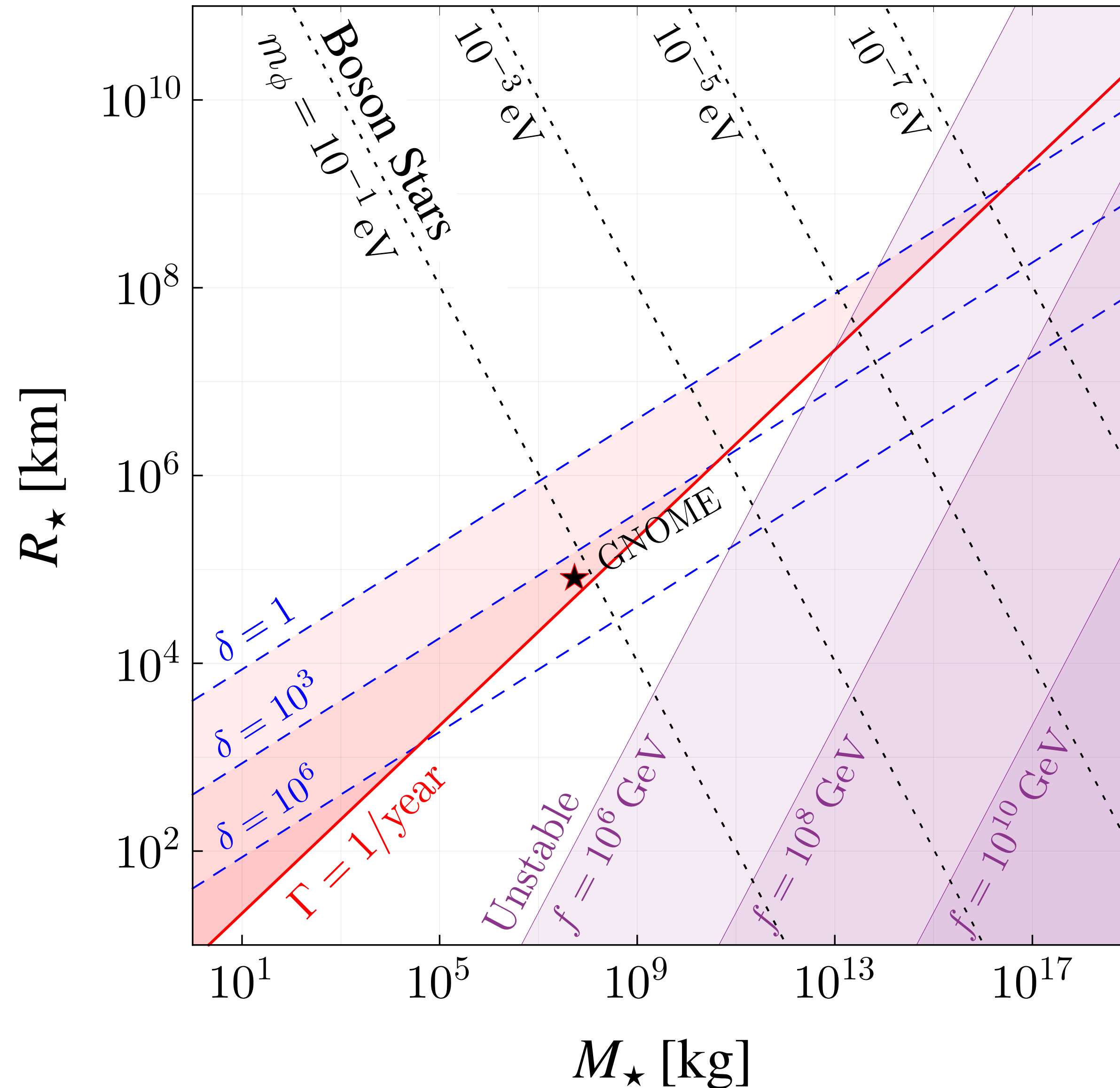


$R \times (m_a^2 \sqrt{G/\lambda})$



Boson Star Encounters

Budker, Banerjee, Eby, Kim,
Perez (1902.08212)



Overdensity

$$\delta \equiv \frac{\rho_\star}{\rho_{\text{dm}}} \propto \rho_{\text{local}}^{-1} R_\star^{-4} m_\phi^{-2}$$

Encounter rate

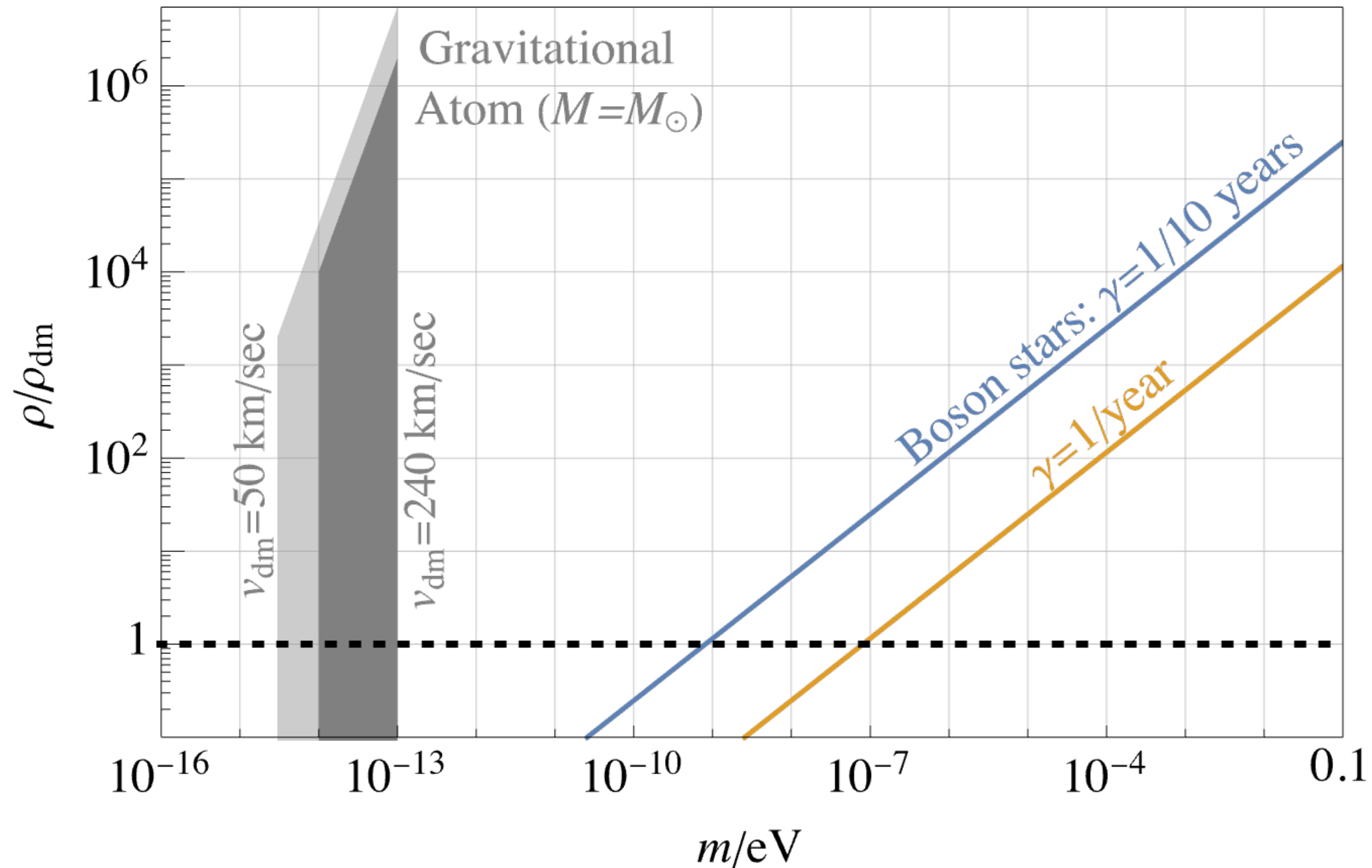
$$\Gamma \equiv \frac{\rho_{\text{dm}}}{M_\star} \sigma_\star v_\star \propto \rho_{\text{local}} R_\star^3 m_\phi^2$$

Both parameters are only significant when

$$m_\phi \gtrsim 10^{-7} \text{ eV}$$

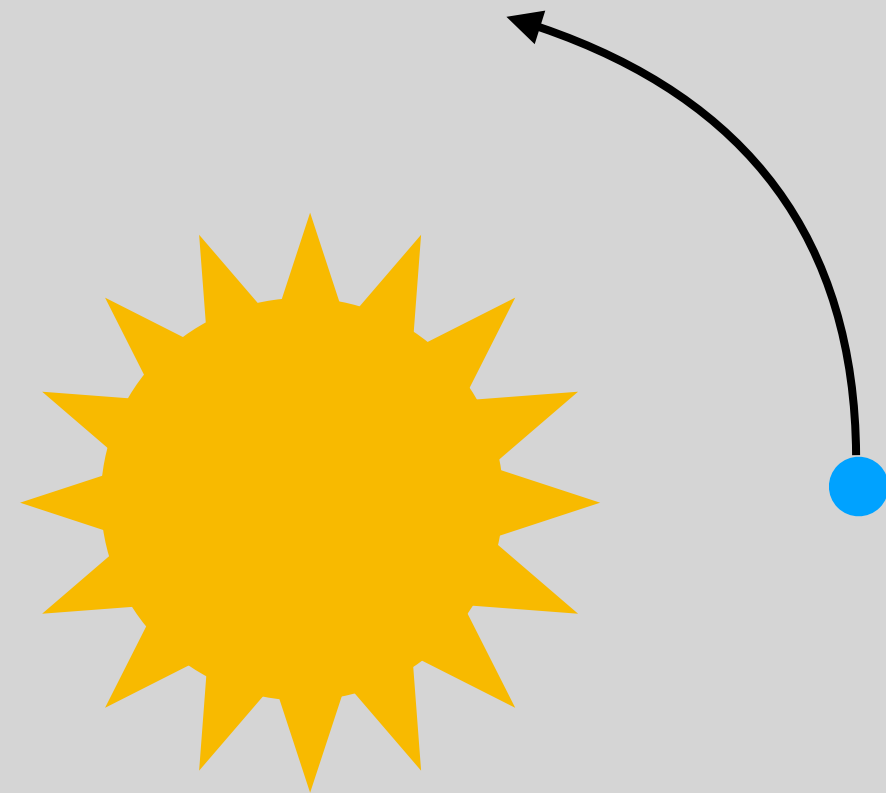
Boson Star Encounters (2)

Budker, Eby, Gorghetto, Jiang,
Perez (2306.12477)



The Very Local DM Density

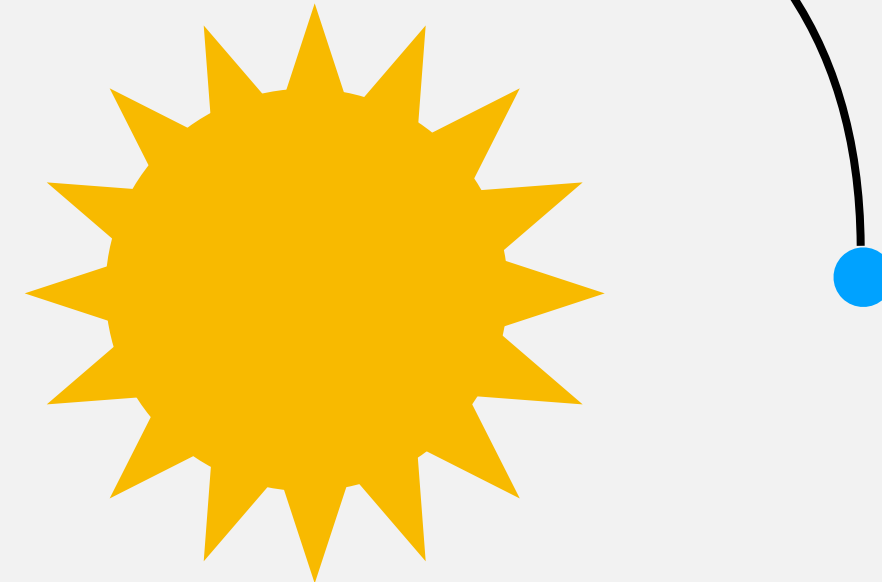
Standard Scenario



No small-scale overdensities

$$\rho = \rho_{\text{dm}}$$

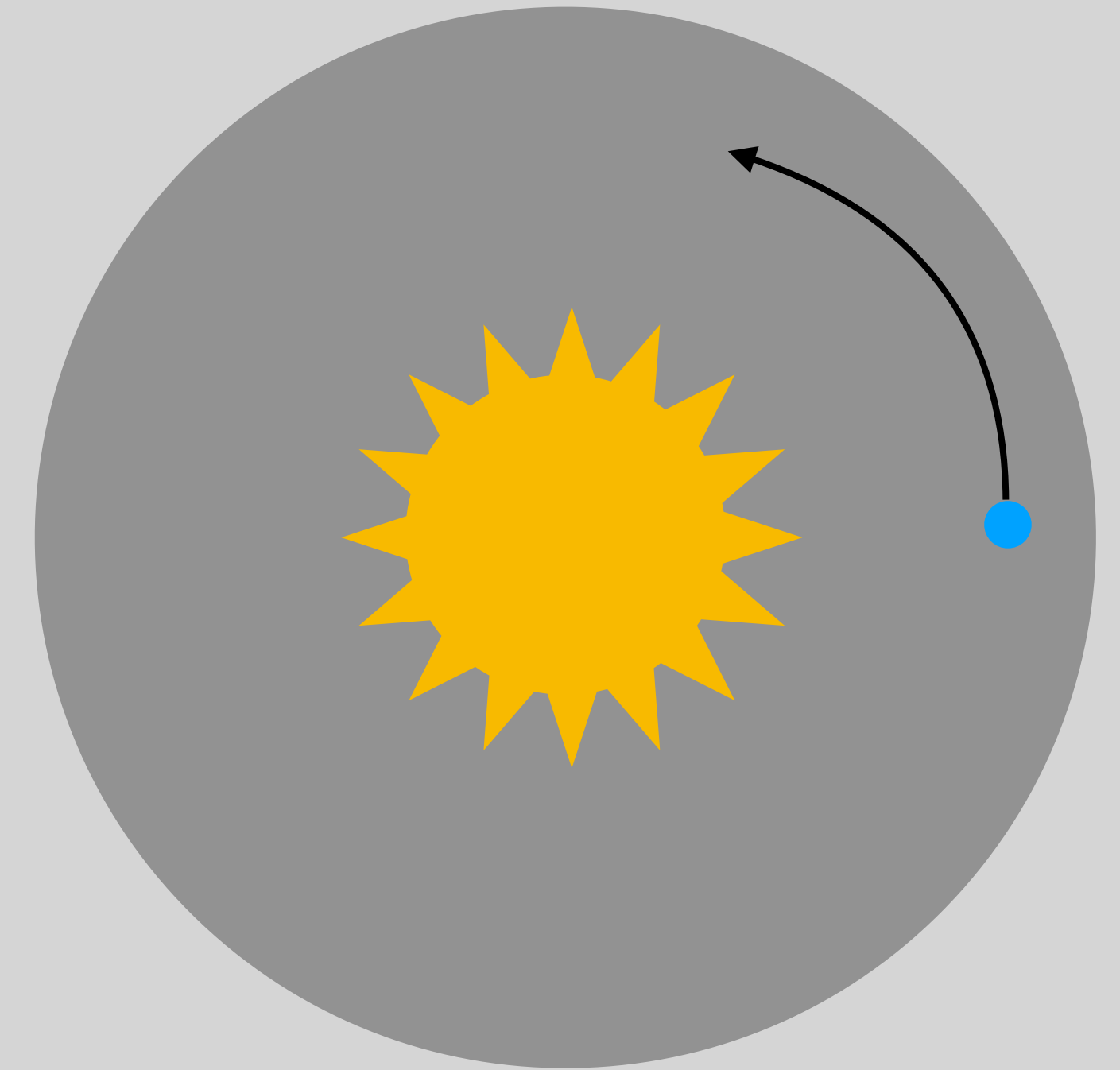
Boson Stars



DM bound in clumps

$$\rho \ll \rho_{\text{dm}}$$

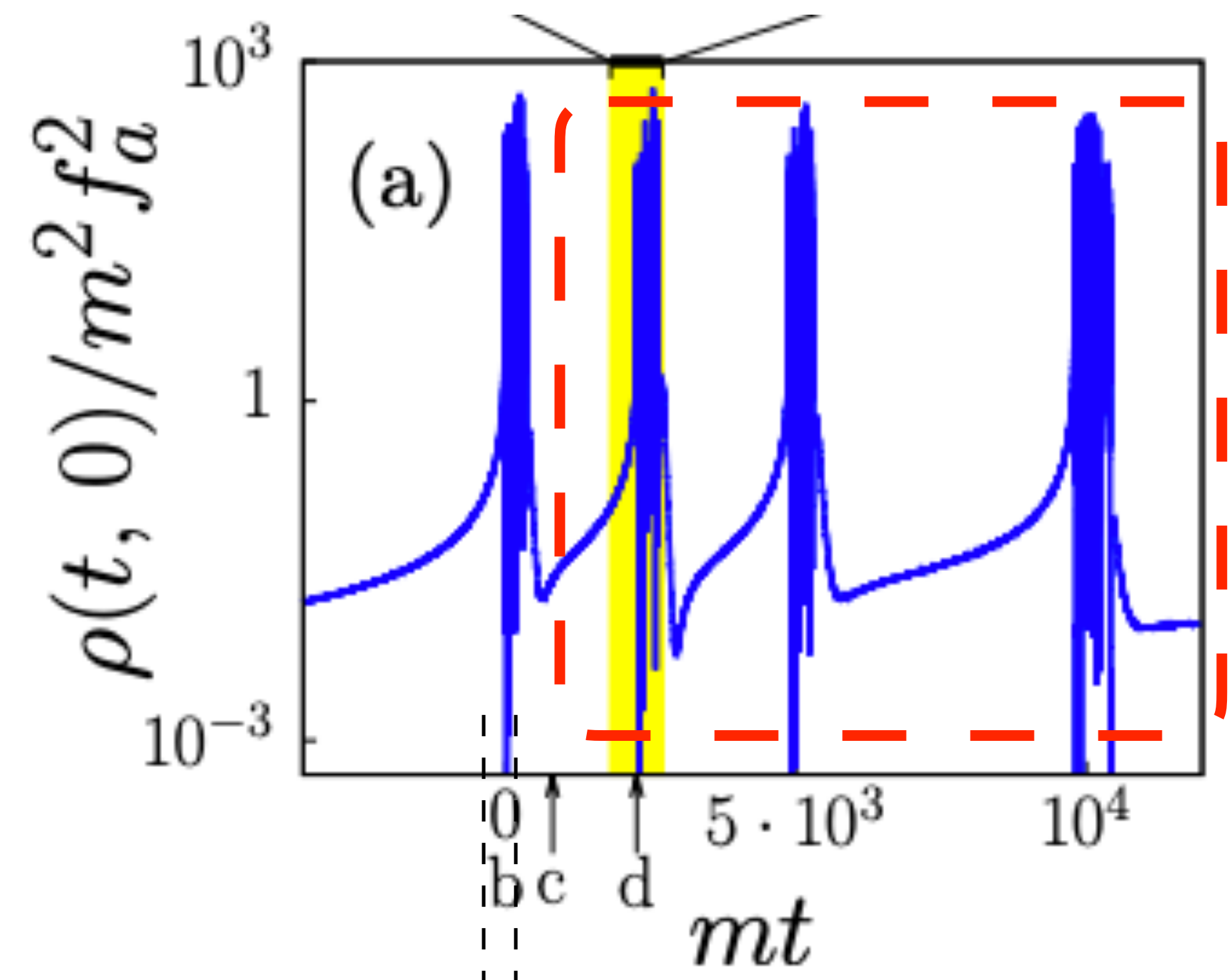
Gravitational Atoms



Overdensities inside the Solar System

$$\rho \gg \rho_{\text{dm}}$$

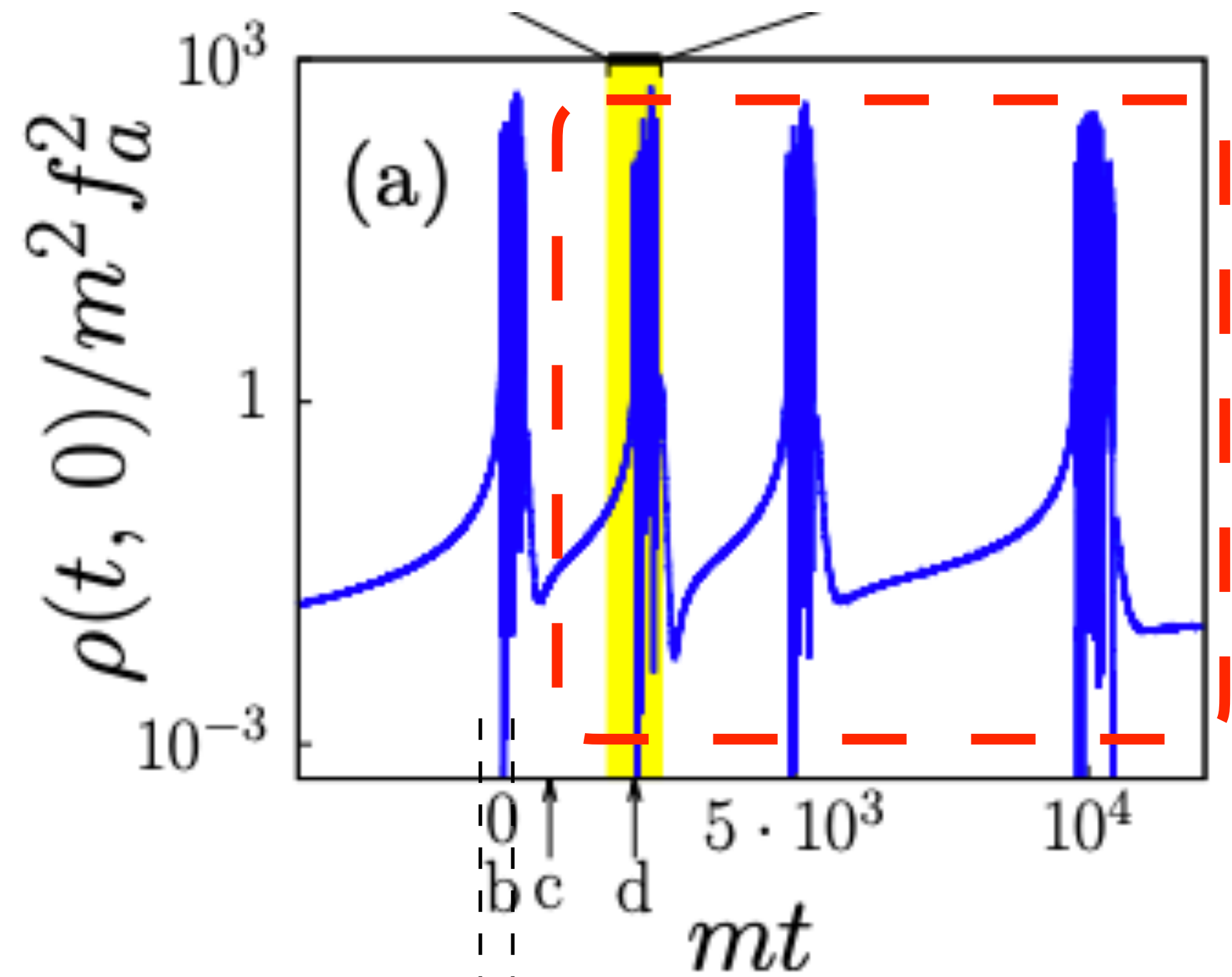
Bosenova Simulation Results



Multiple explosions
per collapse!
(Ignored here)

Short duration $\delta t_{\text{burst}} \sim \mathcal{O}(400)/m$

Bosenova Simulation Results

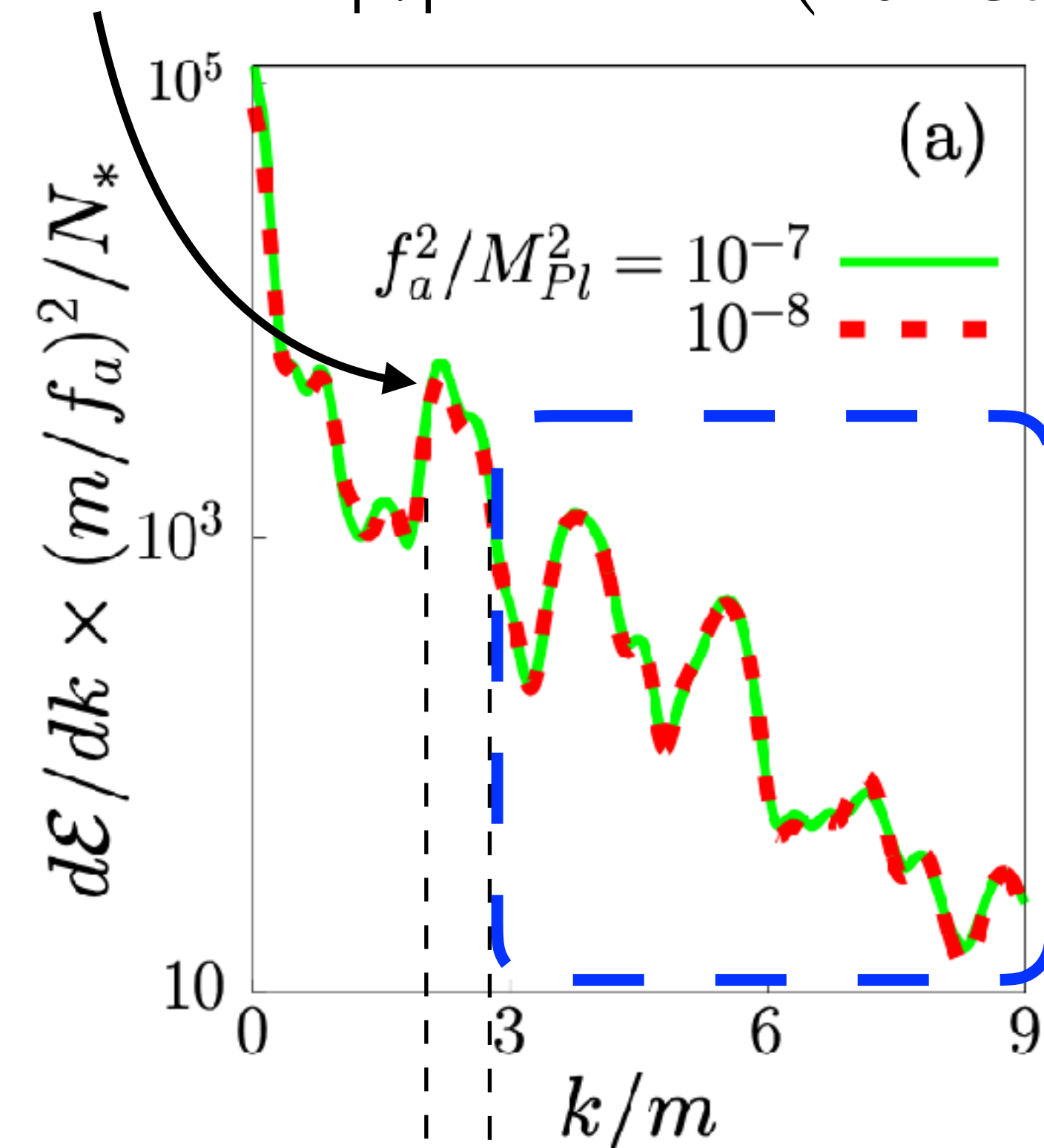


Multiple explosions per collapse!
(Ignored here)

Short duration $\delta t_{\text{burst}} \sim \mathcal{O}(400)/m$

Large integrated energy! In first peak alone,

$$\mathcal{E}_{\text{peak}} \approx 3400 \frac{m}{|\lambda|} \simeq 10^2 M_{\odot} \left(\frac{f}{10^{16} \text{ GeV}} \right)^2 \frac{10^{-15} \text{ eV}}{m}$$

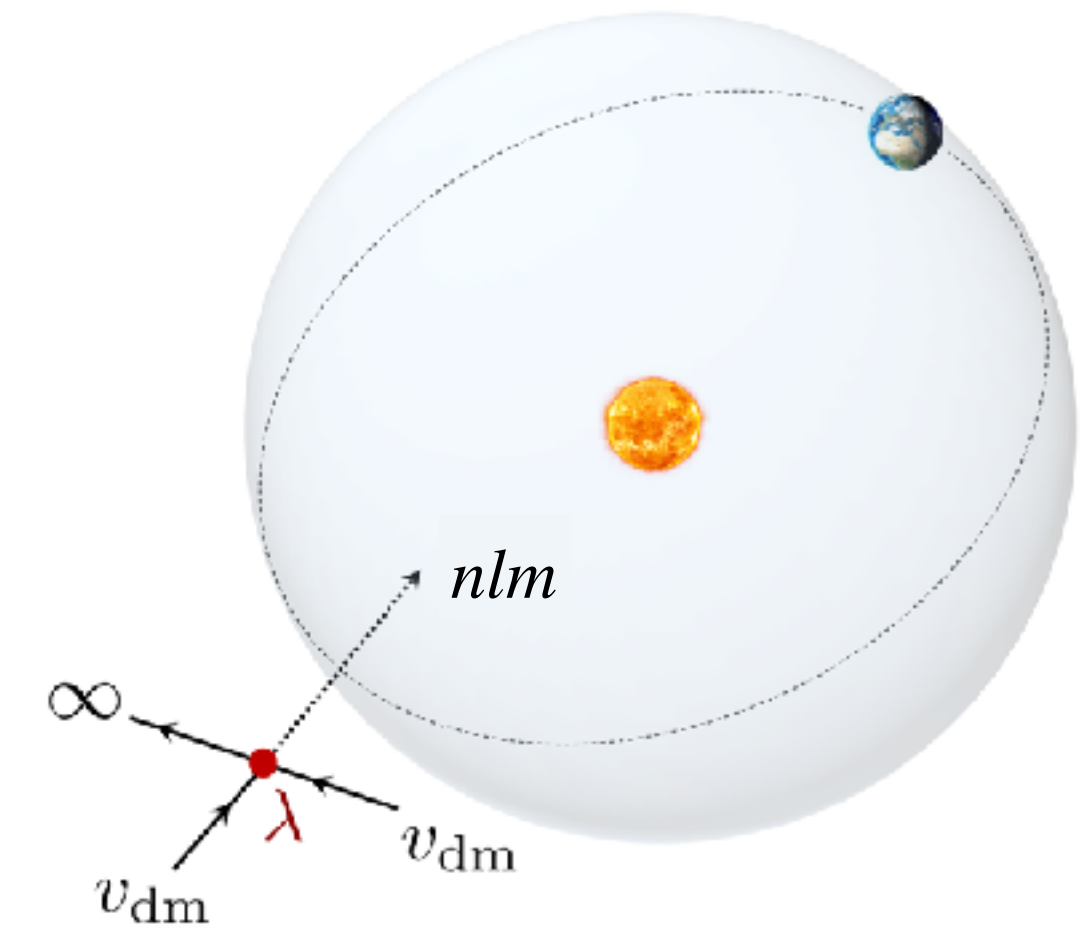


Higher-energy peaks!
(Ignored here)

Momentum peak $k_0 \simeq 2.4 m$
with spread of $\delta k \sim m$

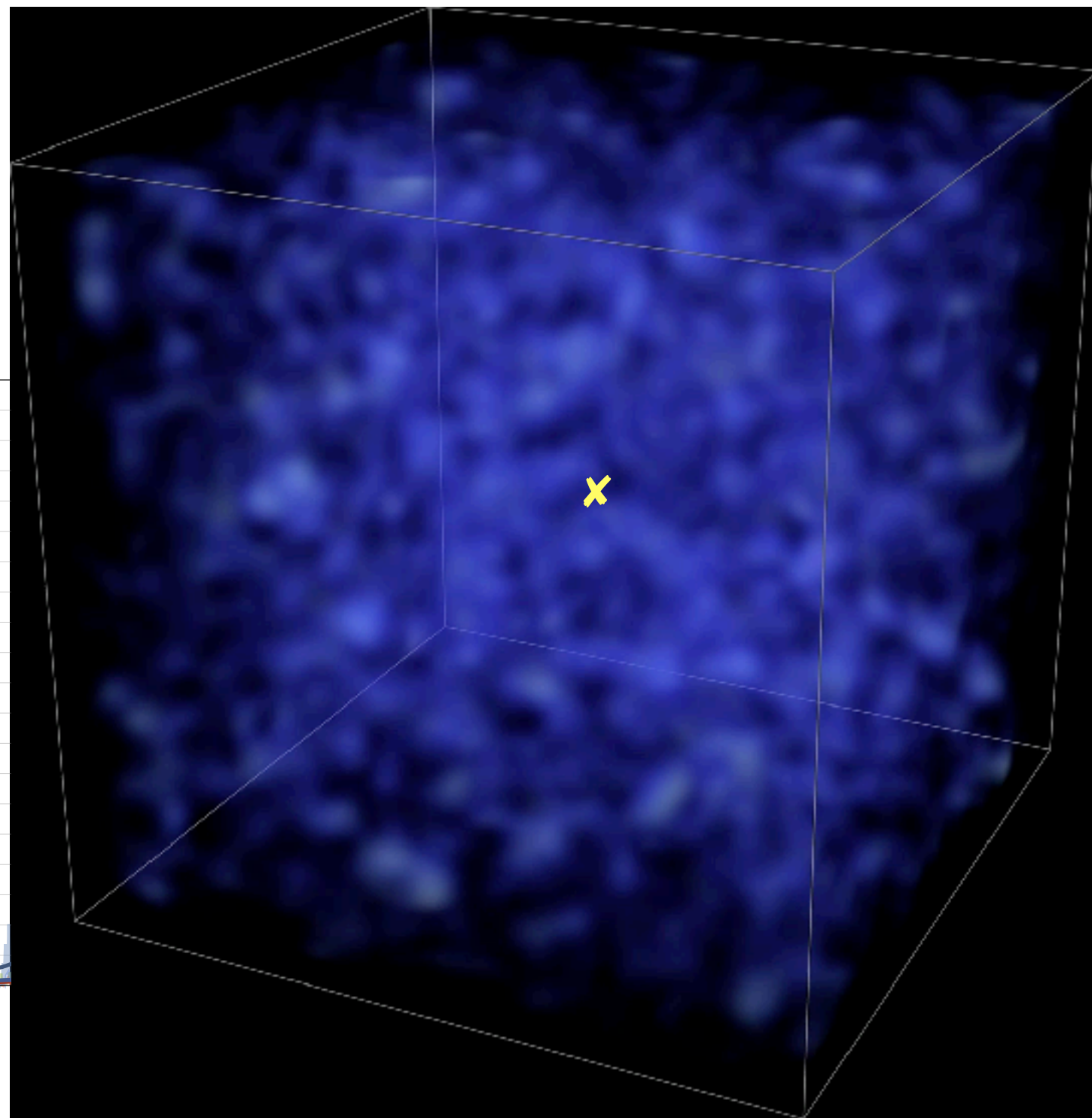
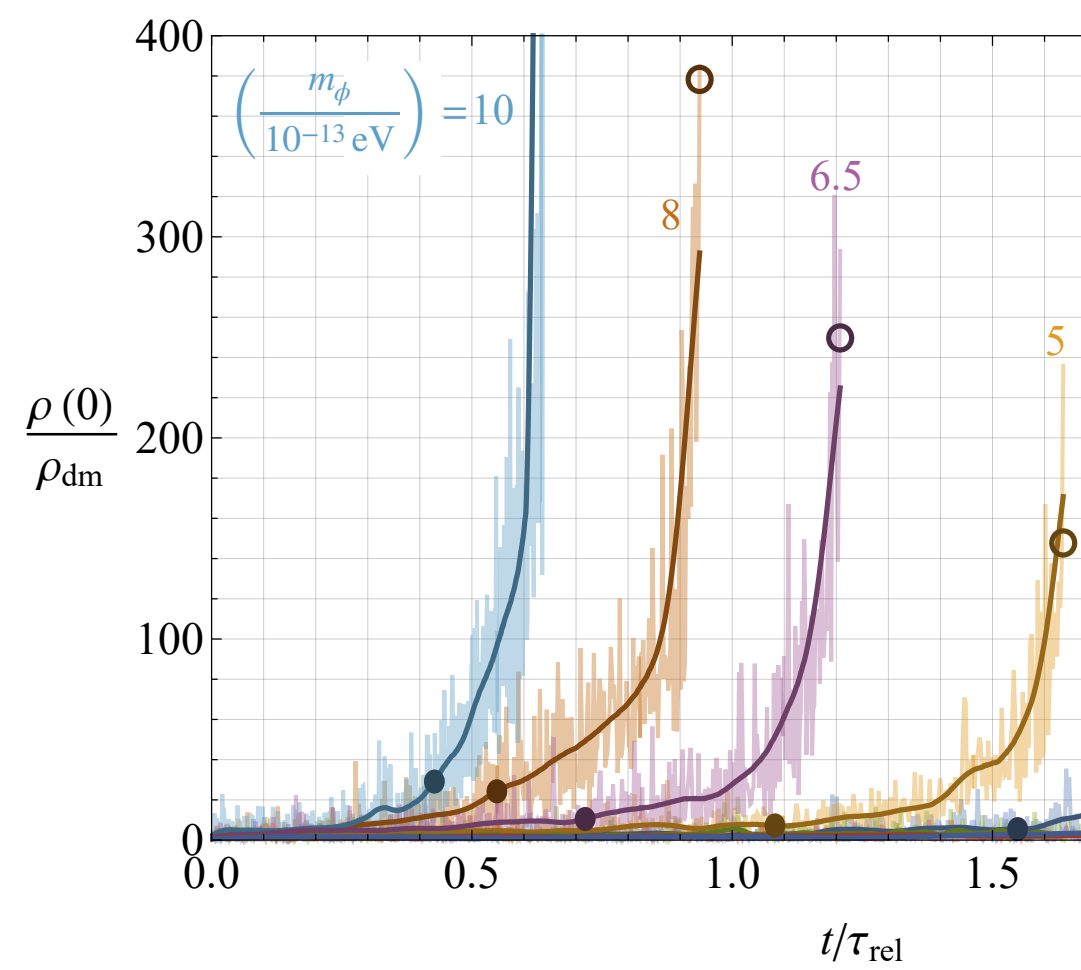
Other sources of Bosonovae

Gravitational Atoms



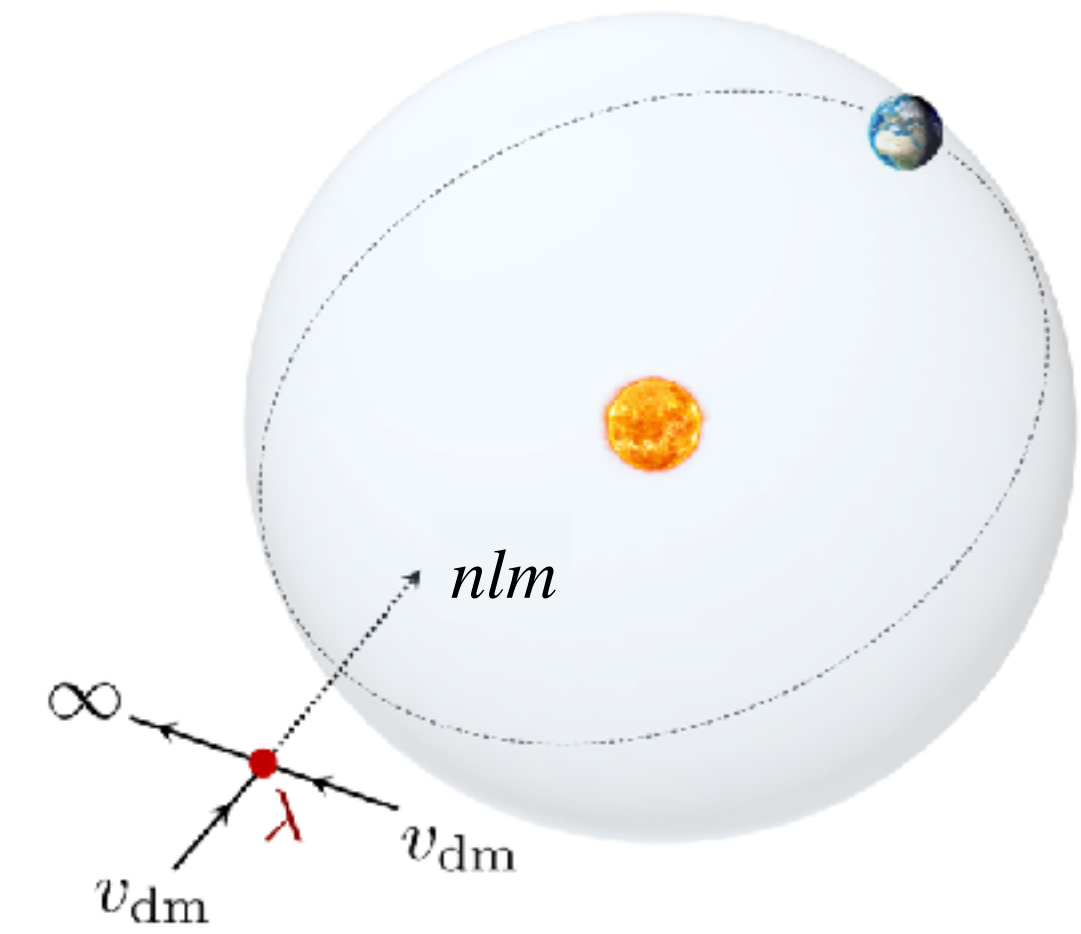
Budker, Banerjee, **J Eby**, Kim,
Perez, Comm Phys 2020
with Matsedonskyi and Flambaum, JCAP 2020

Budker, **J Eby**, Gorghetto, Jiang,
Perez, JCAP 2023



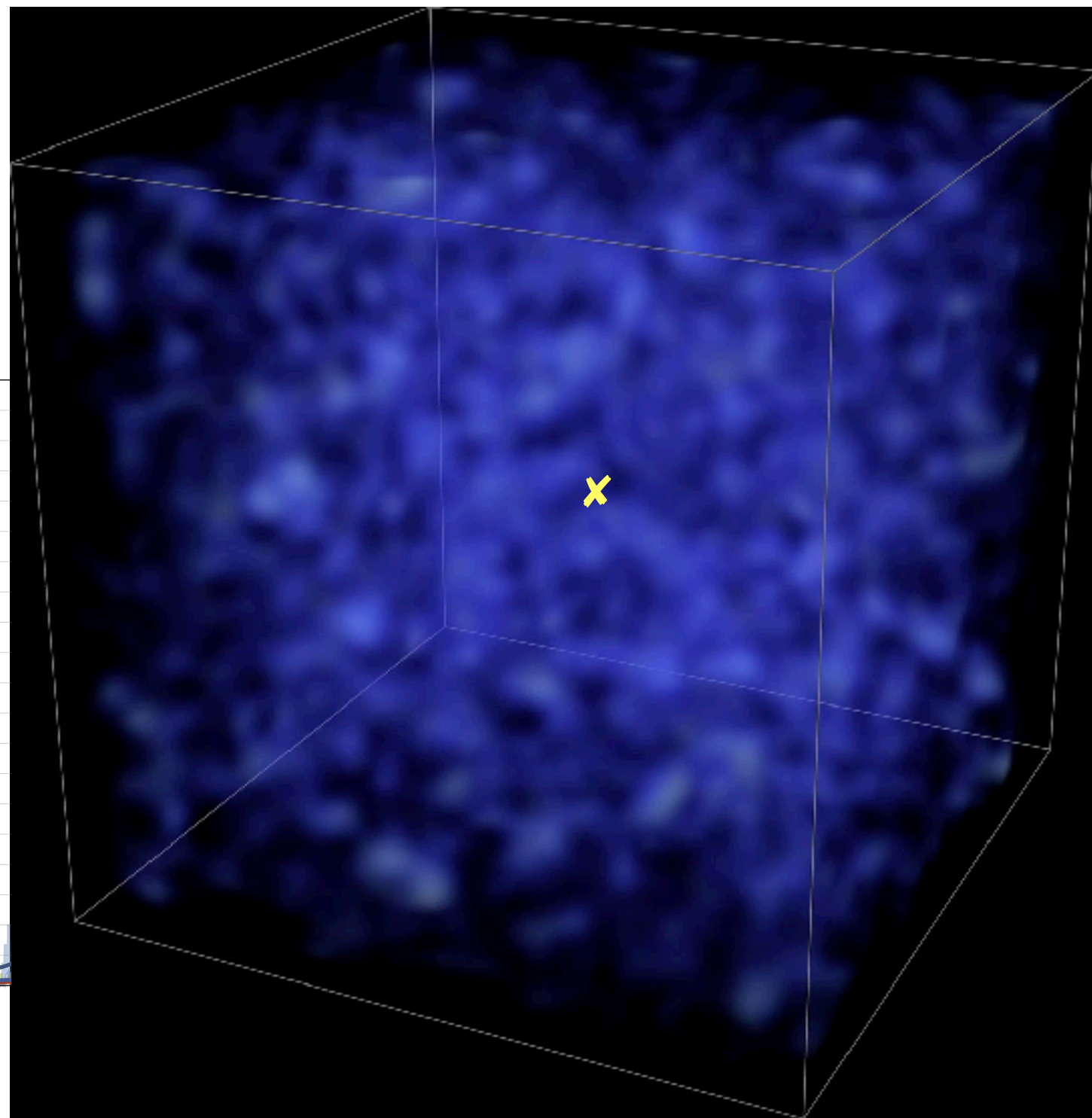
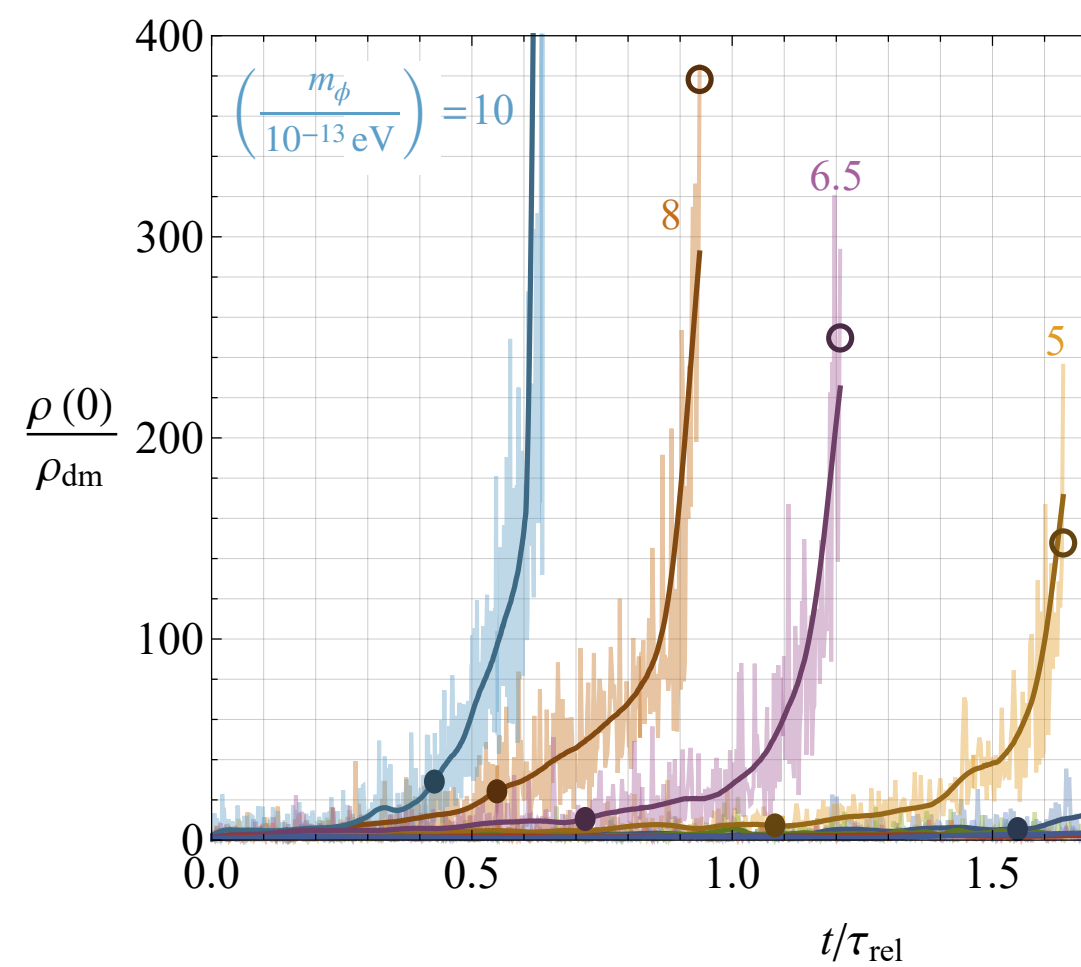
Other sources of Bosonovae

Gravitational Atoms



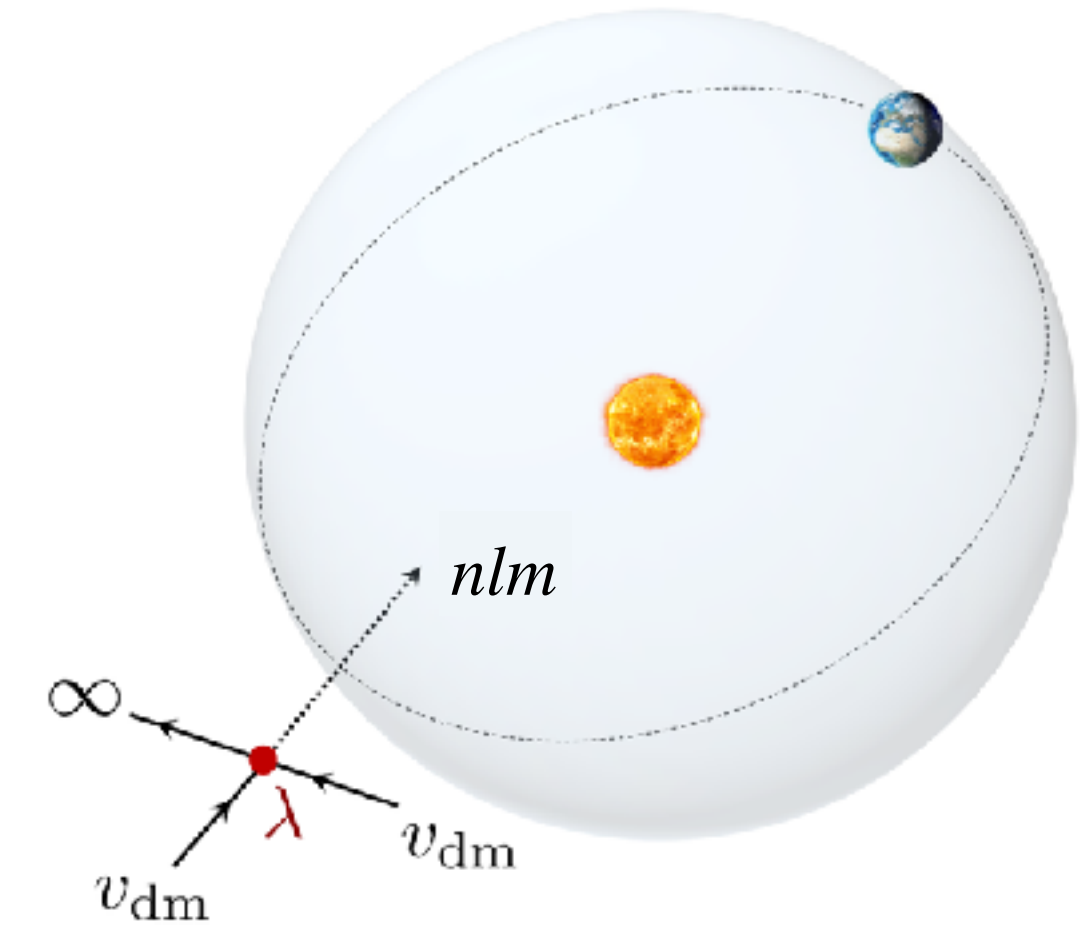
Budker, Banerjee, **J Eby**, Kim,
Perez, Comm Phys 2020
with Matsedonskyi and Flambaum, JCAP 2020

Budker, **J Eby**, Gorghetto, Jiang,
Perez, JCAP 2023



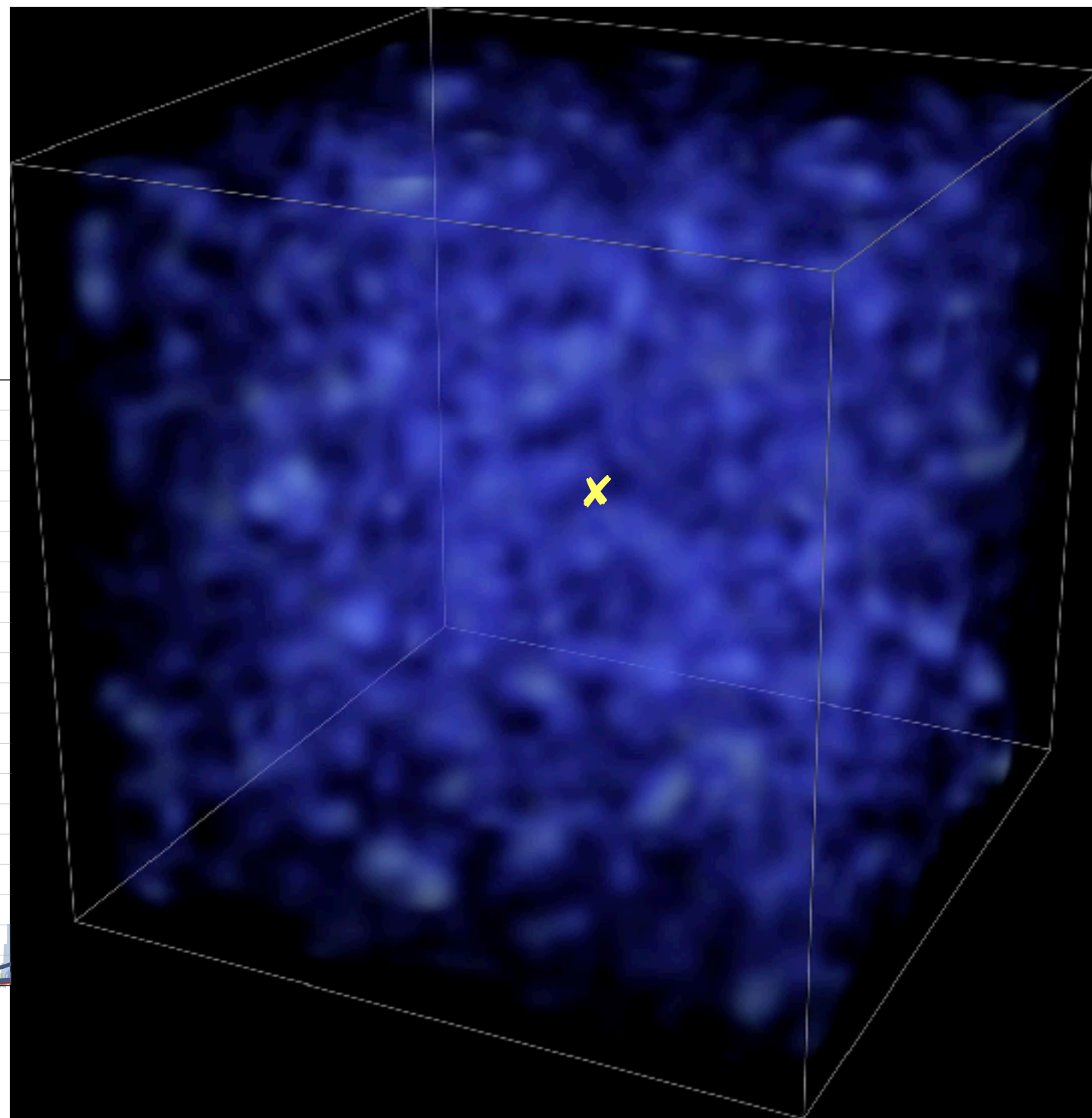
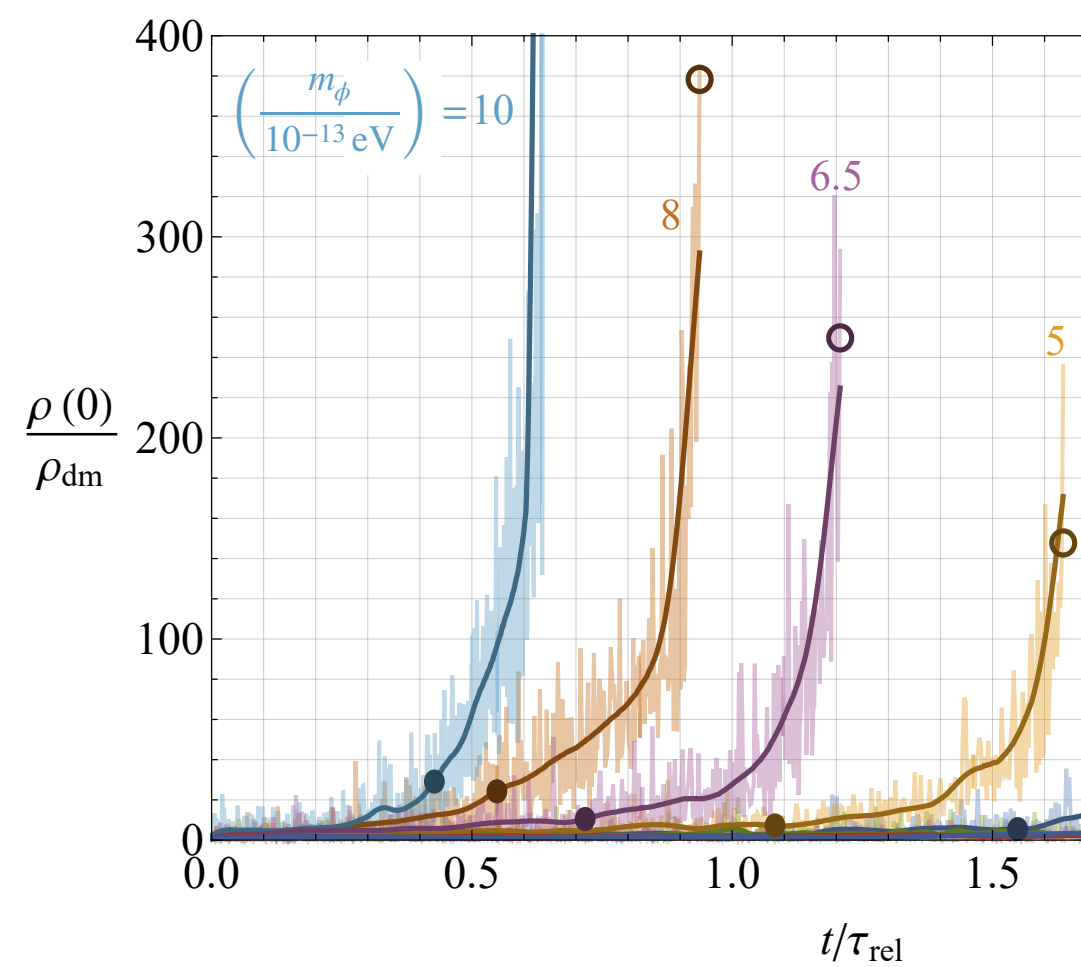
Other sources of Bosenovae

Gravitational Atoms



Budker, Banerjee, **J Eby**, Kim, Perez, Comm Phys 2020
with Matsedonskyi and Flambaum, JCAP 2020

Budker, **J Eby**, Gorghetto, Jiang, Perez, JCAP 2023



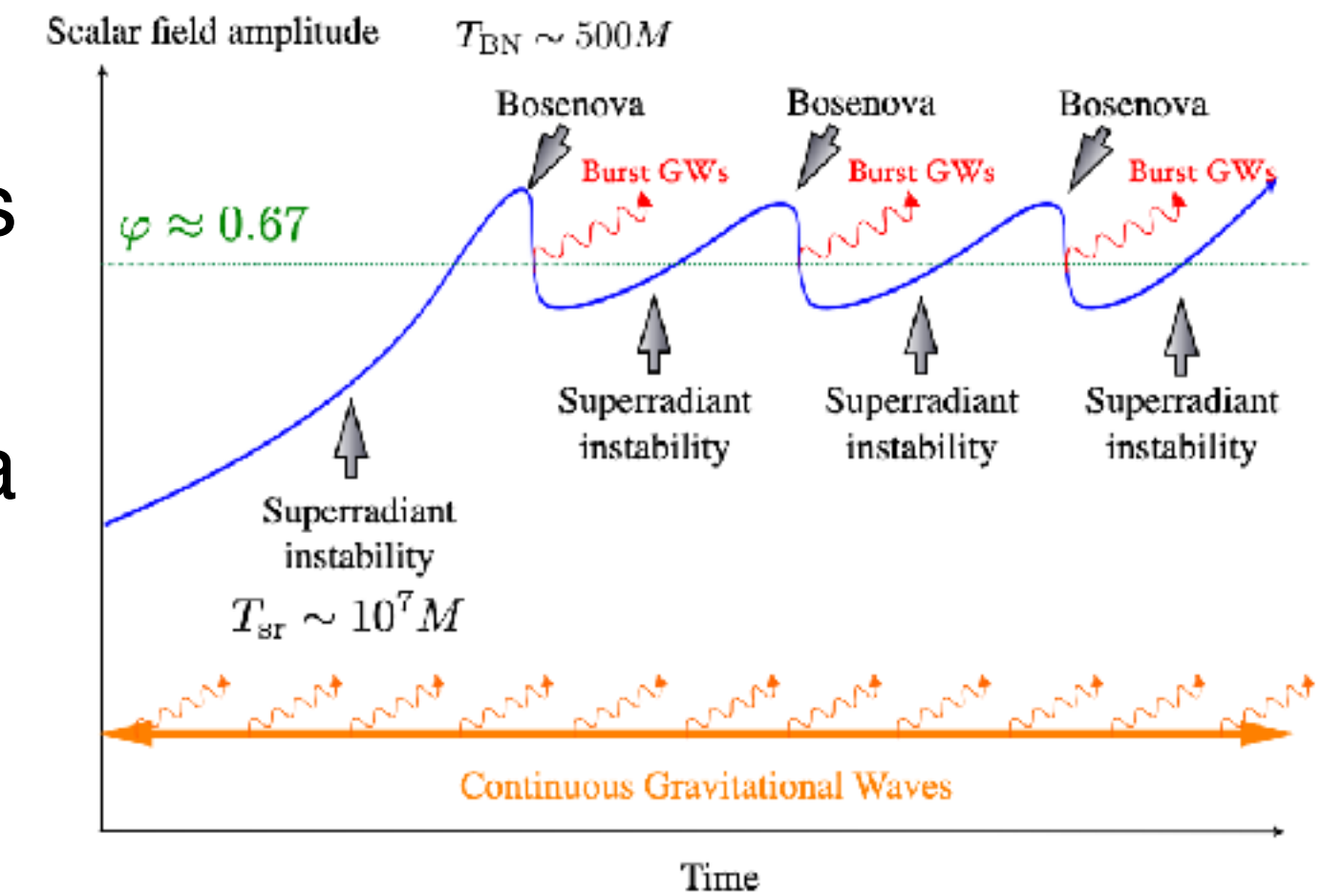
Black Hole Superradiance

Bosenova predicted in Arvanitaki and Dubovsky, PRD 2011

Weak self-interactions

→ bosenova

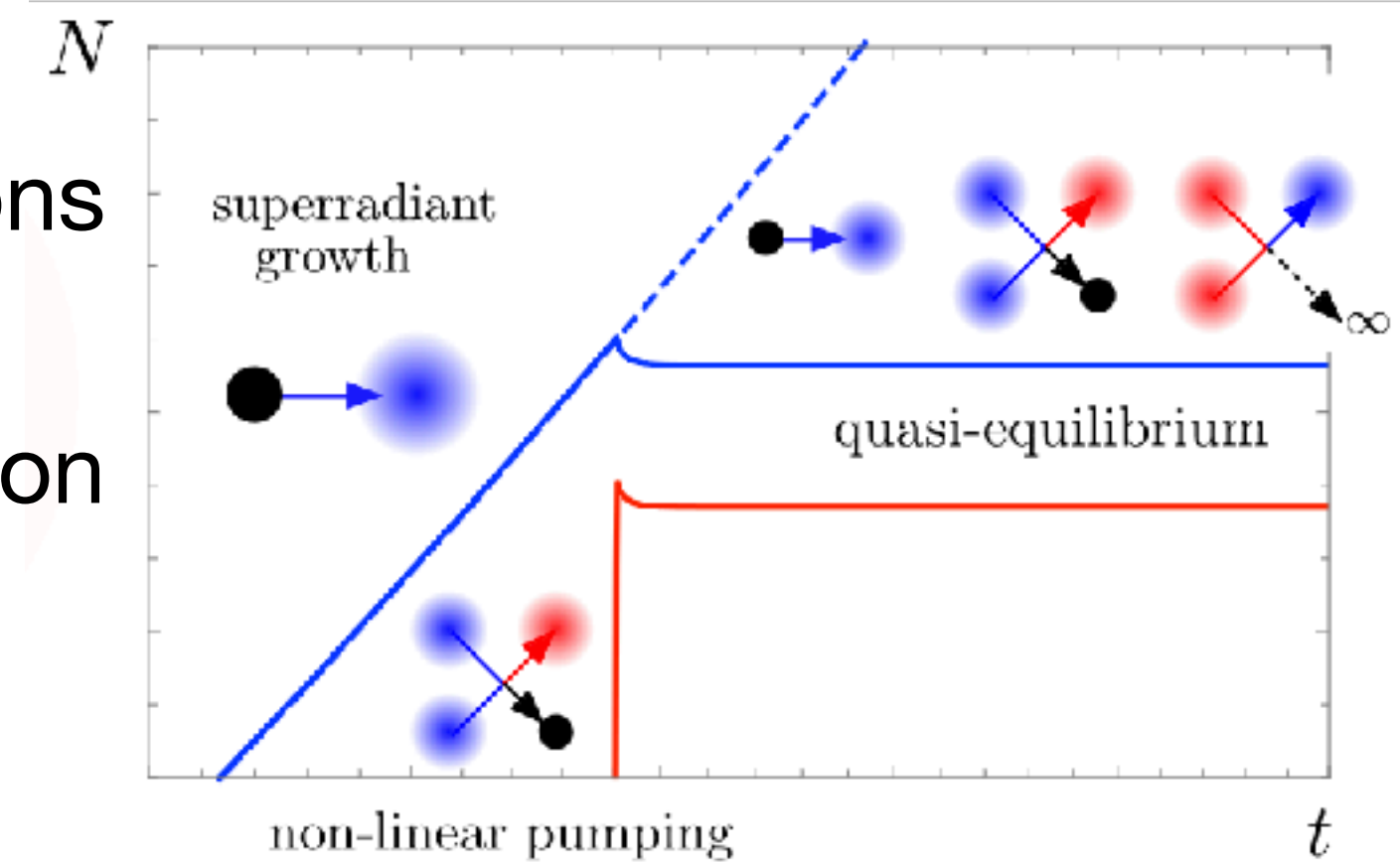
Yoshino and Kodama, PTP 2012, CQG 2015



Large self-interactions

→ axion wave emission

Baryakhtar++, PRD 2021

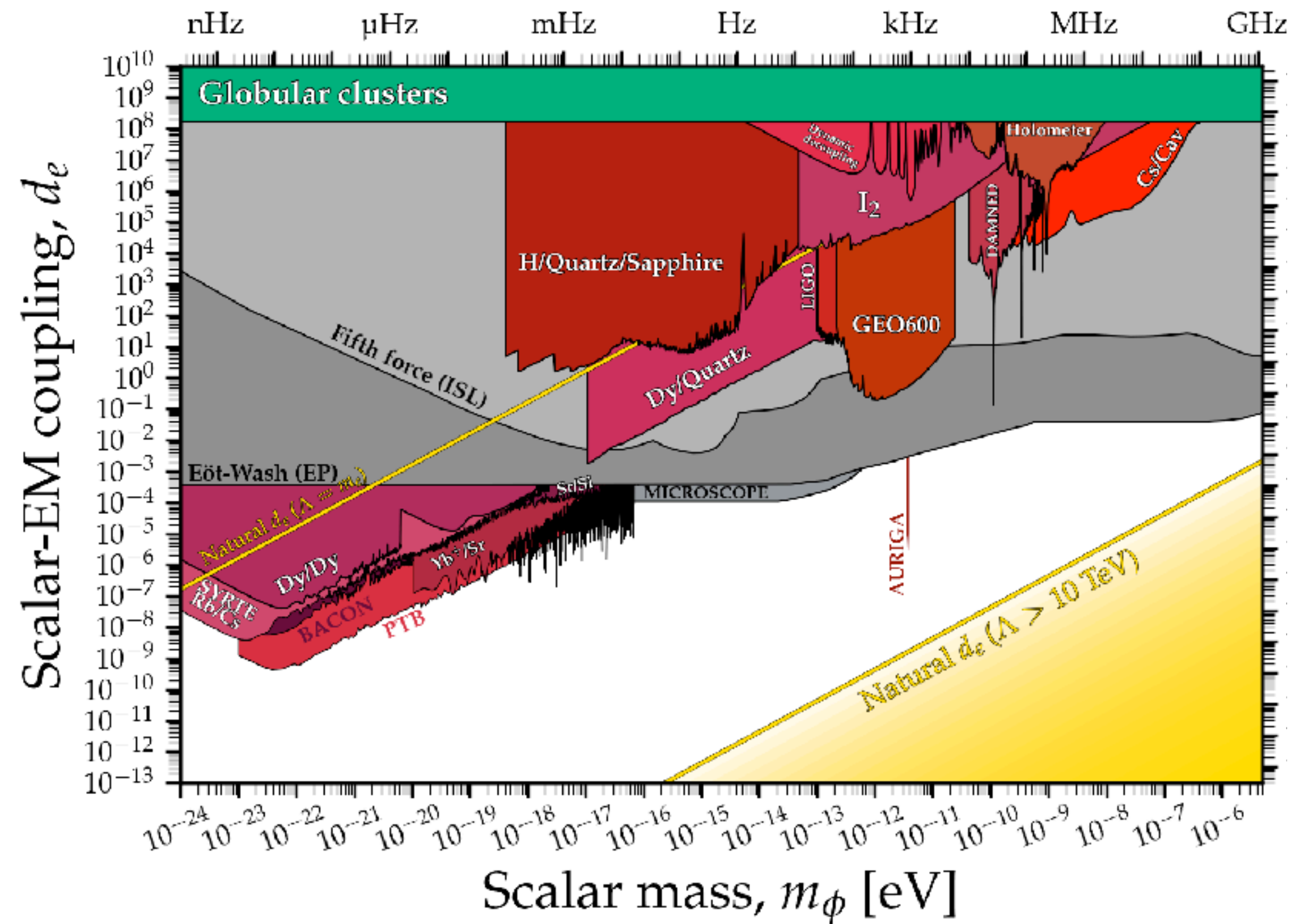


Other Couplings

Coupling to photons

Oscillation of
fine-structure constant

$$\mathcal{L} \supset d_e \frac{\phi}{2M_{\text{Pl}}} F^{\mu\nu} F_{\mu\nu} \longrightarrow \alpha \rightarrow \alpha + \delta\alpha(d_e, \rho, m_\phi)$$



Other Couplings

Coupling to photons

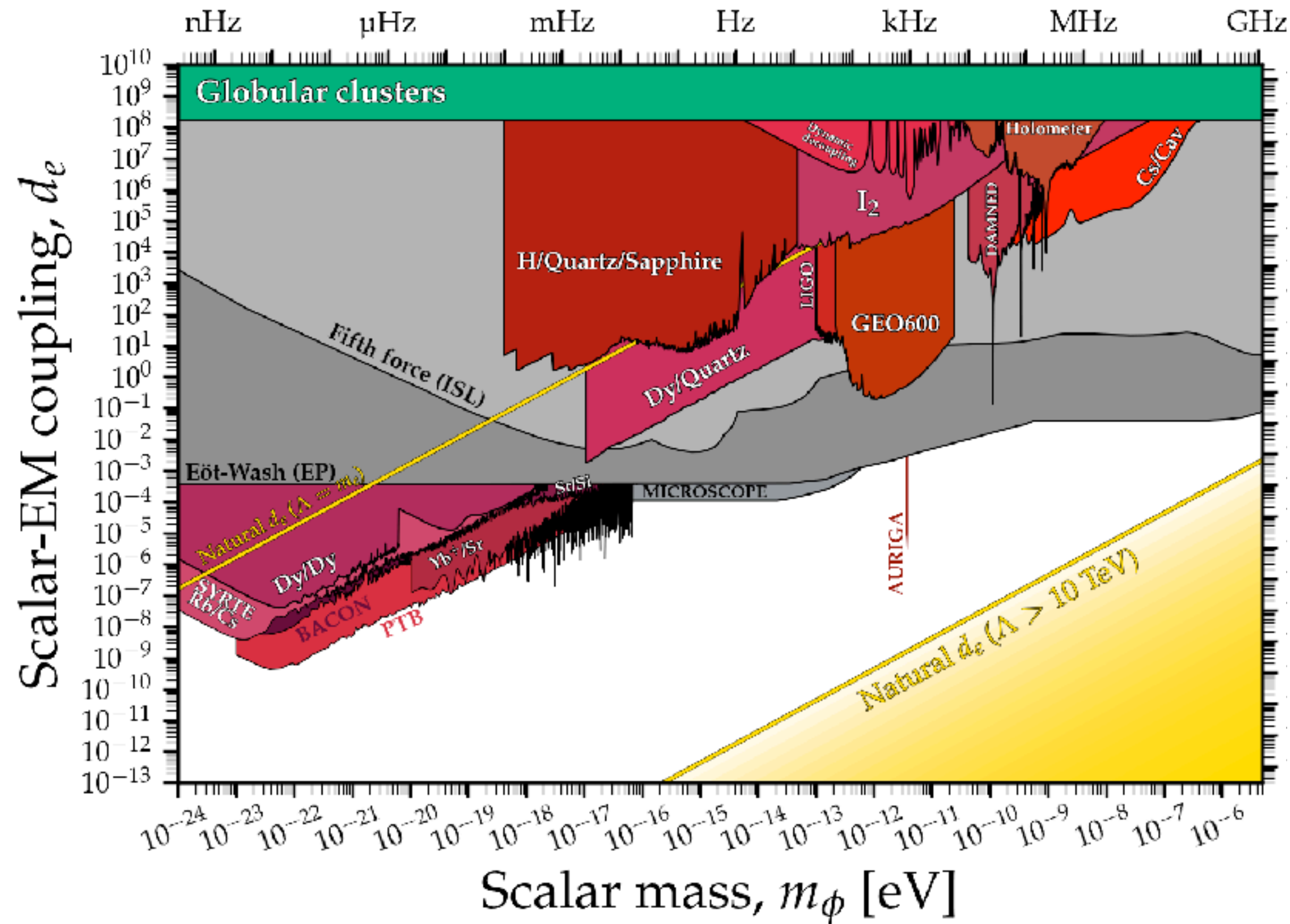
$$\mathcal{L} \supset d_e \frac{\phi}{2M_{\text{Pl}}} F^{\mu\nu} F_{\mu\nu} \longrightarrow \alpha \rightarrow \alpha + \delta\alpha(d_e, \rho, m_\phi)$$

Oscillation of fine-structure constant

Coupling to electrons

$$\mathcal{L} \supset d_{m_e} \frac{\phi}{2M_{\text{Pl}}} \bar{e}e \longrightarrow \frac{m_e}{m_p} \rightarrow \frac{m_e}{m_p} + \frac{\delta m_e(d_{m_e}, \rho, m_\phi)}{m_p}$$

Oscillation of electron-proton mass ratio



Other Couplings

Coupling to photons

$$\mathcal{L} \supset d_e \frac{\phi}{2M_{\text{Pl}}} F^{\mu\nu} F_{\mu\nu} \longrightarrow \alpha \rightarrow \alpha + \delta\alpha(d_e, \rho, m_\phi)$$

Oscillation of fine-structure constant

Coupling to electrons

$$\mathcal{L} \supset d_{m_e} \frac{\phi}{2M_{\text{Pl}}} \bar{e}e \longrightarrow \frac{m_e}{m_p} \rightarrow \frac{m_e}{m_p} + \frac{\delta m_e(d_{m_e}, \rho, m_\phi)}{m_p}$$

Oscillation of electron-proton mass ratio

Coupling to gluons

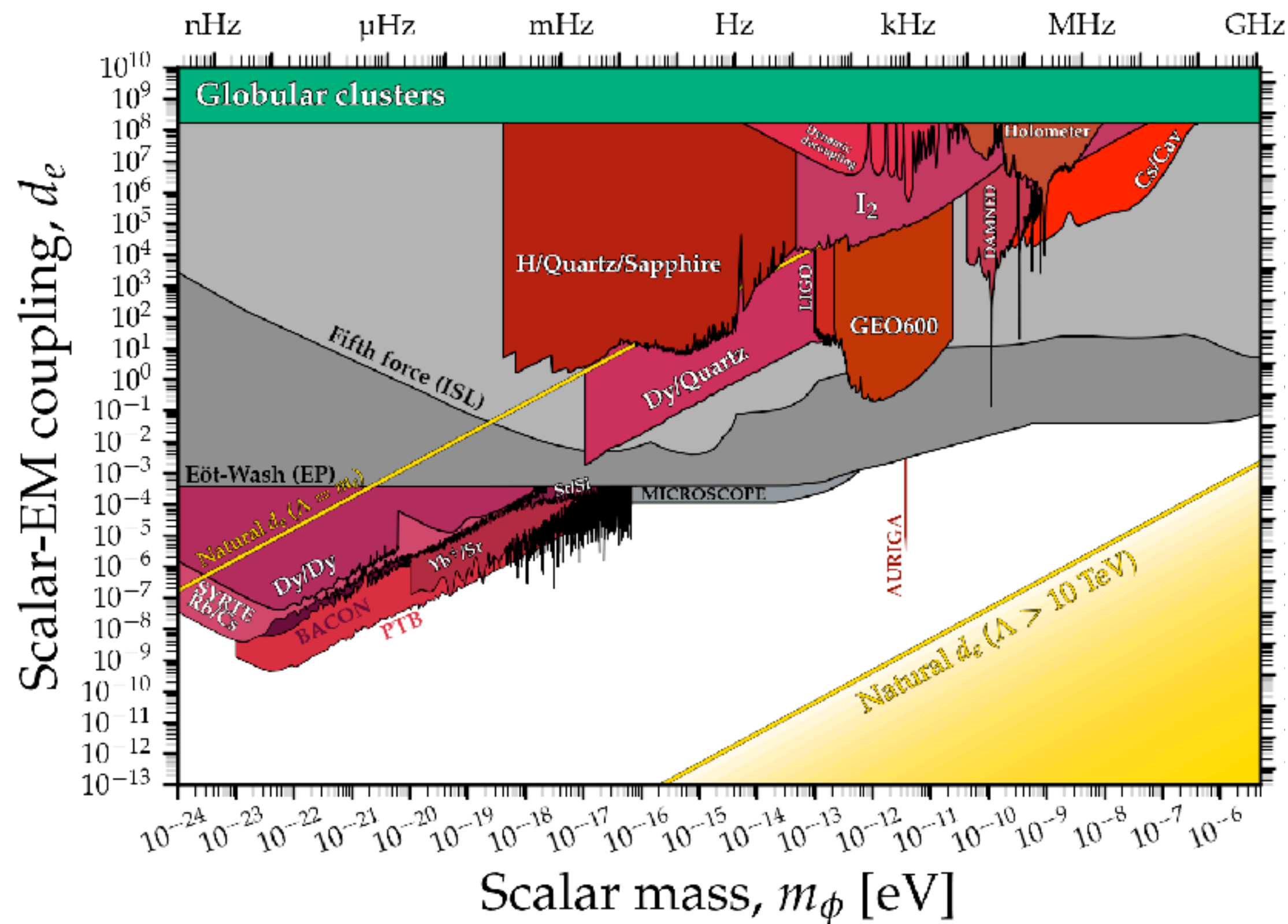
$$\mathcal{L} \supset d_g \frac{\phi}{2M_{\text{Pl}}} G^{a\mu\nu} G_{\mu\nu}^a$$

Coupling to quarks

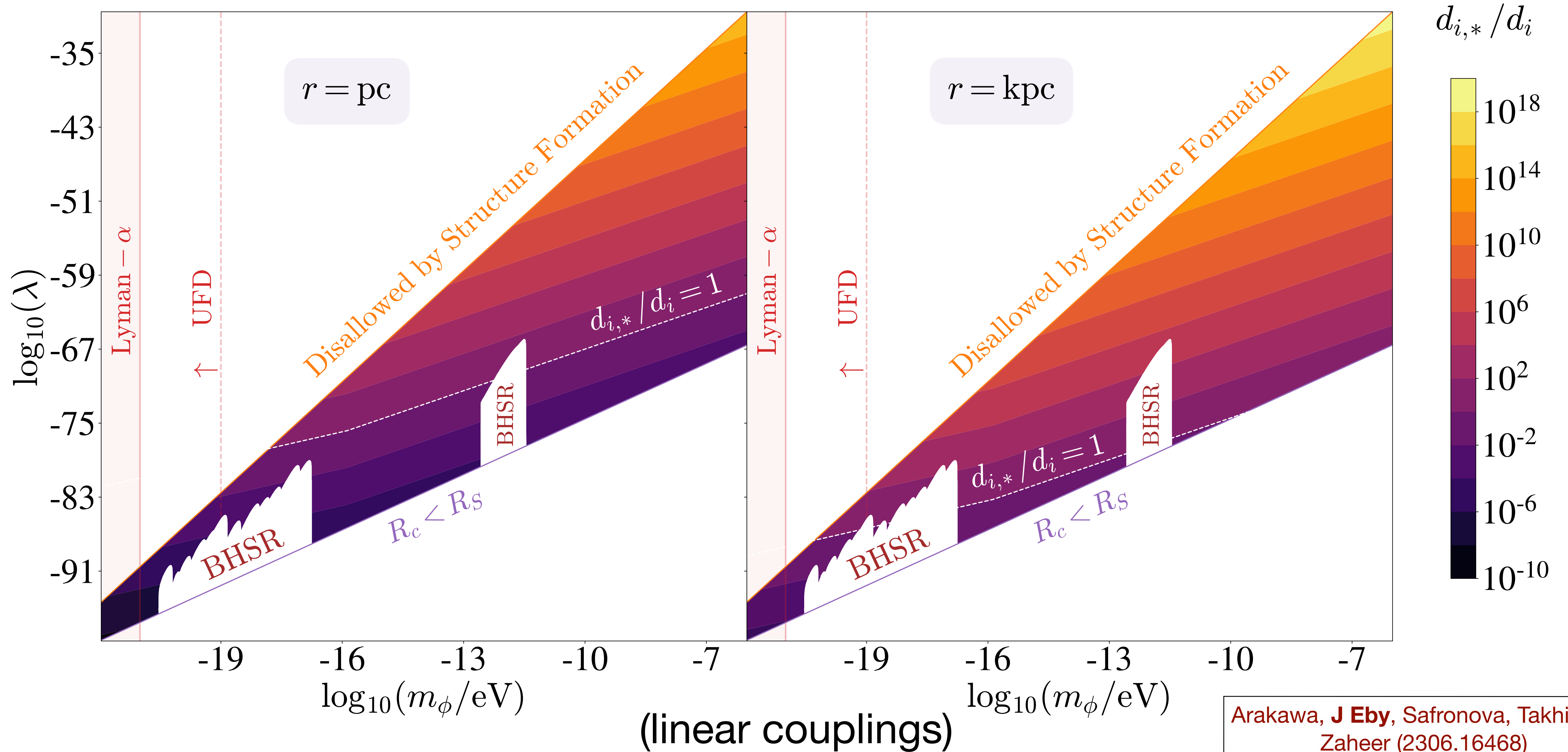
$$\mathcal{L} \supset d_{m_q} \frac{\phi}{2M_{\text{Pl}}} \bar{q}q$$

$$\frac{m_q}{\Lambda_{\text{QCD}}} \rightarrow \frac{m_q}{\Lambda_{\text{QCD}}} + \delta \left(\frac{m_q}{\Lambda_{\text{QCD}}} \right) (d_{m_q}, d_g, \rho, m_\phi)$$

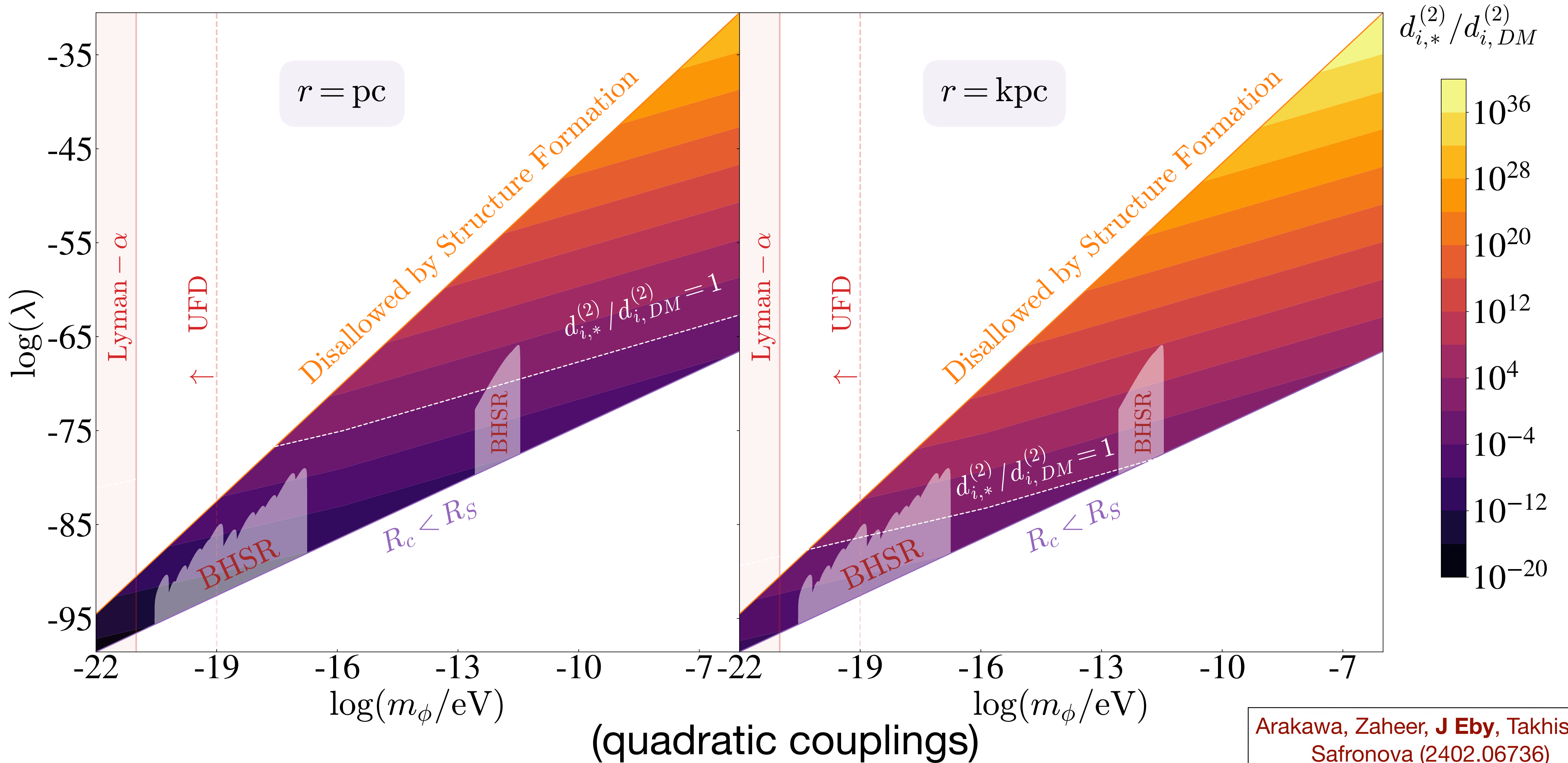
Oscillation of ratio of quark mass to QCD scale



“Experiment-Independent” Sensitivity



“Experiment-Independent” Sensitivity (2)



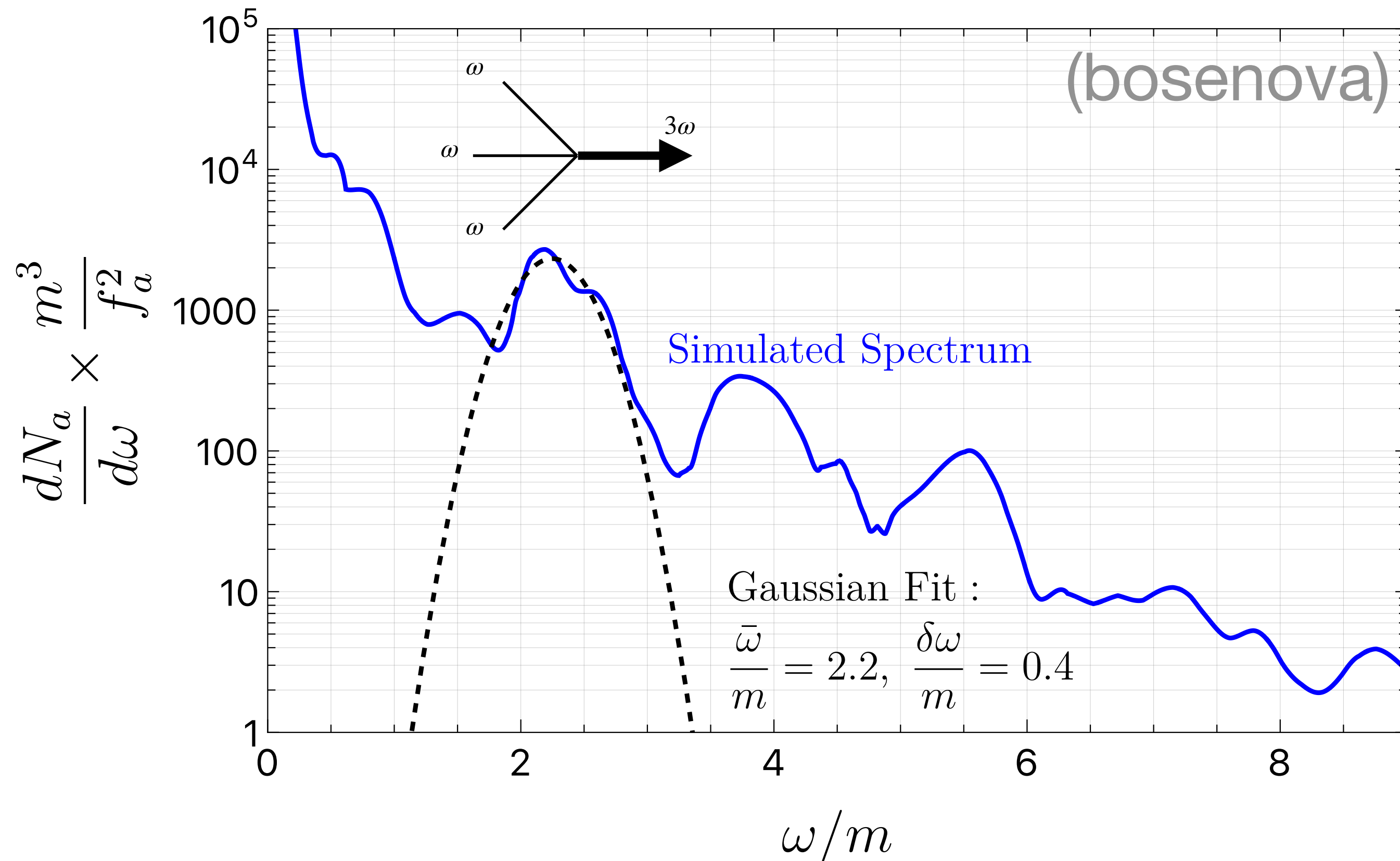
Parameterization: Flux

$$\frac{dN_a}{d\omega}(\omega) \propto \frac{E_{\text{tot}}}{m_a^2} \frac{\exp\left(-\frac{(\omega - \bar{\omega})^2}{\delta\omega^2}\right)}{\delta\omega/m_a}$$

E_{tot} : total energy emitted
in single burst

$\bar{\omega}$: peak energy

$\delta\omega$: energy width



- easily captures peaked distribution
- computationally simple
- sum of Gaussians can be used for asymmetric distributions, e.g. power-law

Parameterization: Rate

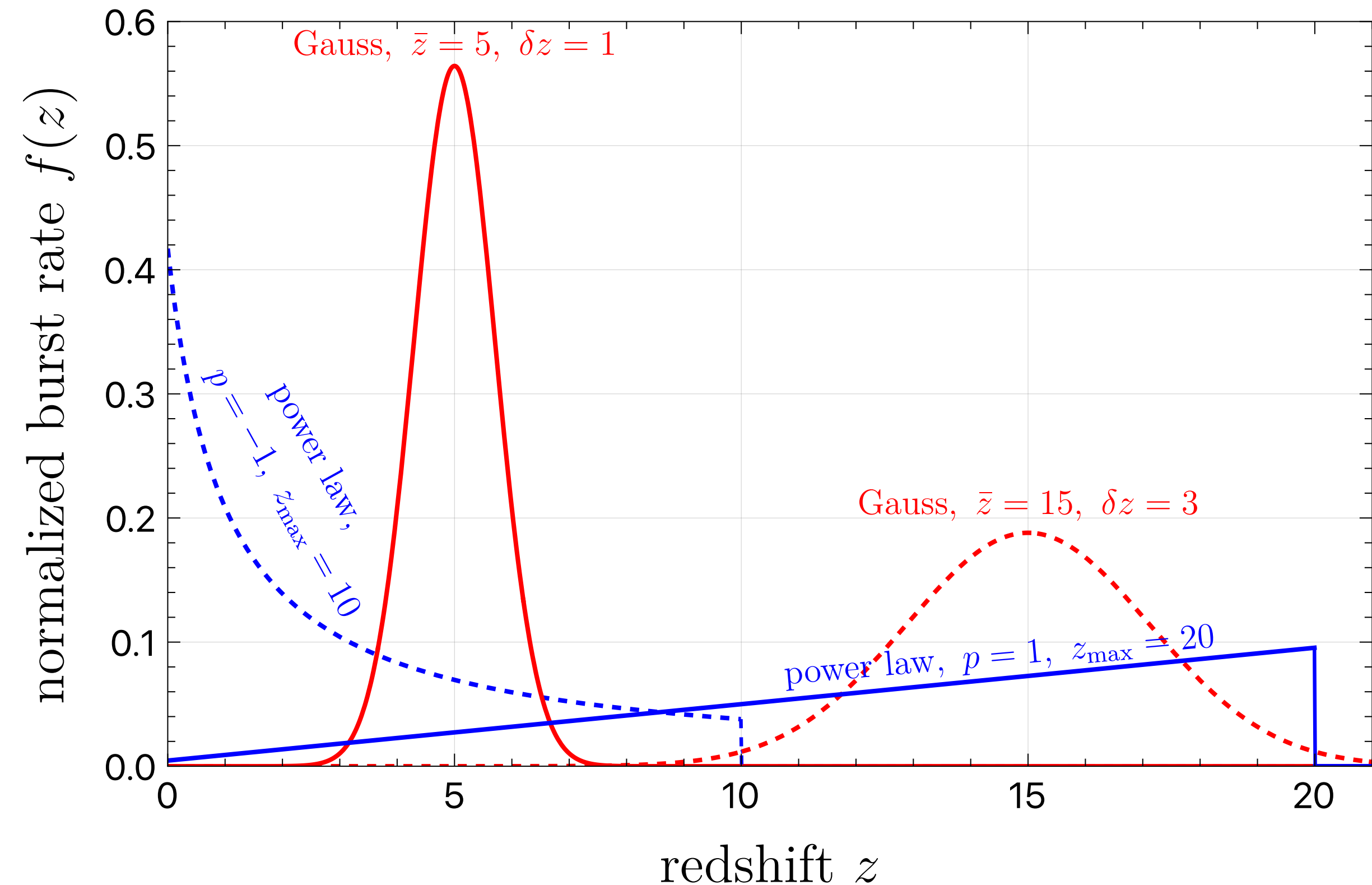
$$f(z) = (1 + z)^p \Theta(z - z_{\max}) \text{ for power-law} \quad \longleftarrow \quad R_{\text{burst}}(z) \propto \frac{\rho_{\text{loss}} H_0}{E_{\text{tot}}(z)} f(z)$$

$$f(z) = \exp\left(-\frac{(z - \bar{z})^2}{\delta z^2}\right) \text{ for Gaussian}$$

ρ_{loss} : total relativistic energy
density emitted across all z

Convenient normalisation:

$$\rho_{\text{loss}} \equiv \mathcal{F} \bar{\rho}_U \quad \text{with } \bar{\rho}_U \simeq 10^{-6} \text{ GeV/cm}^3$$



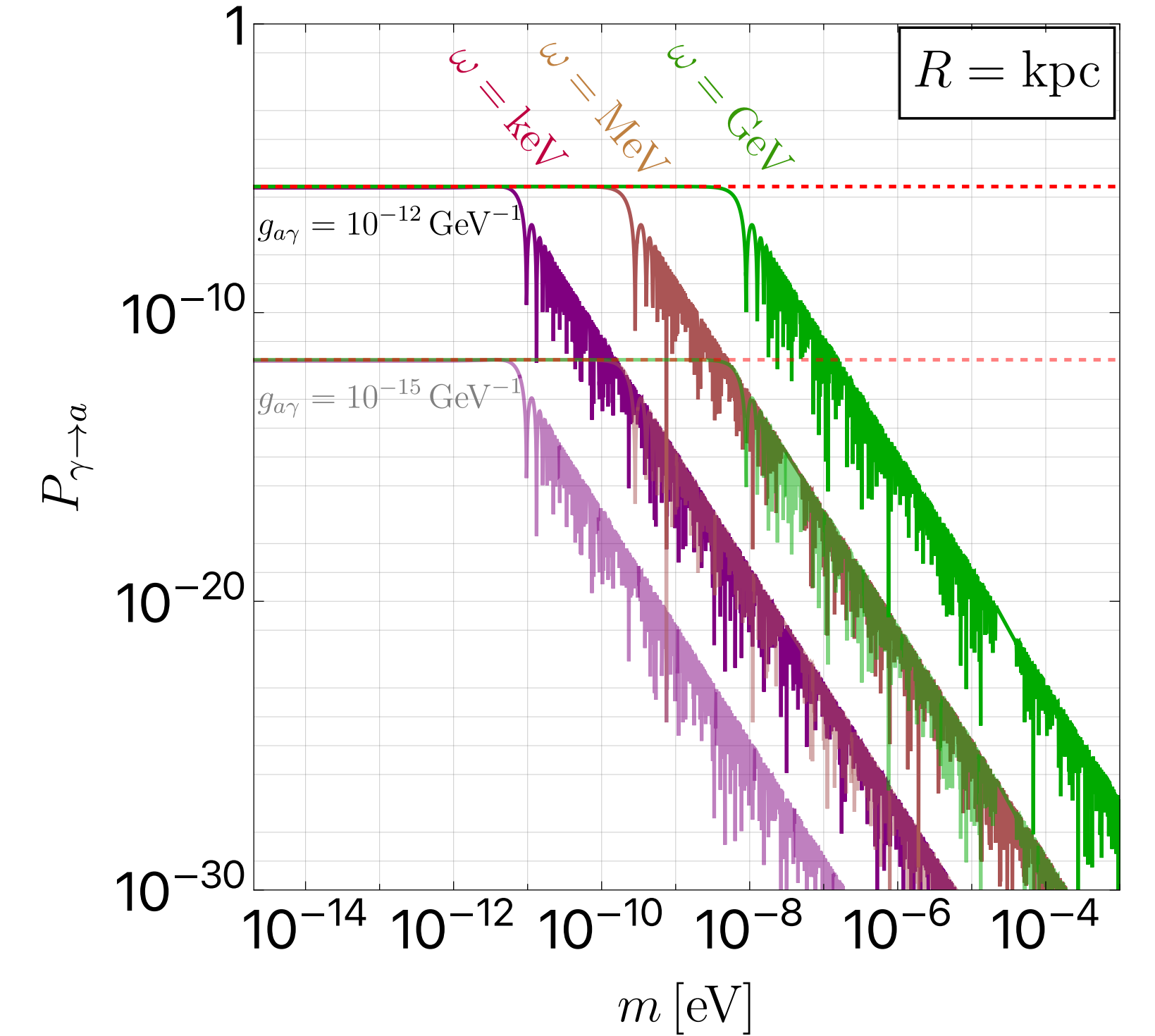
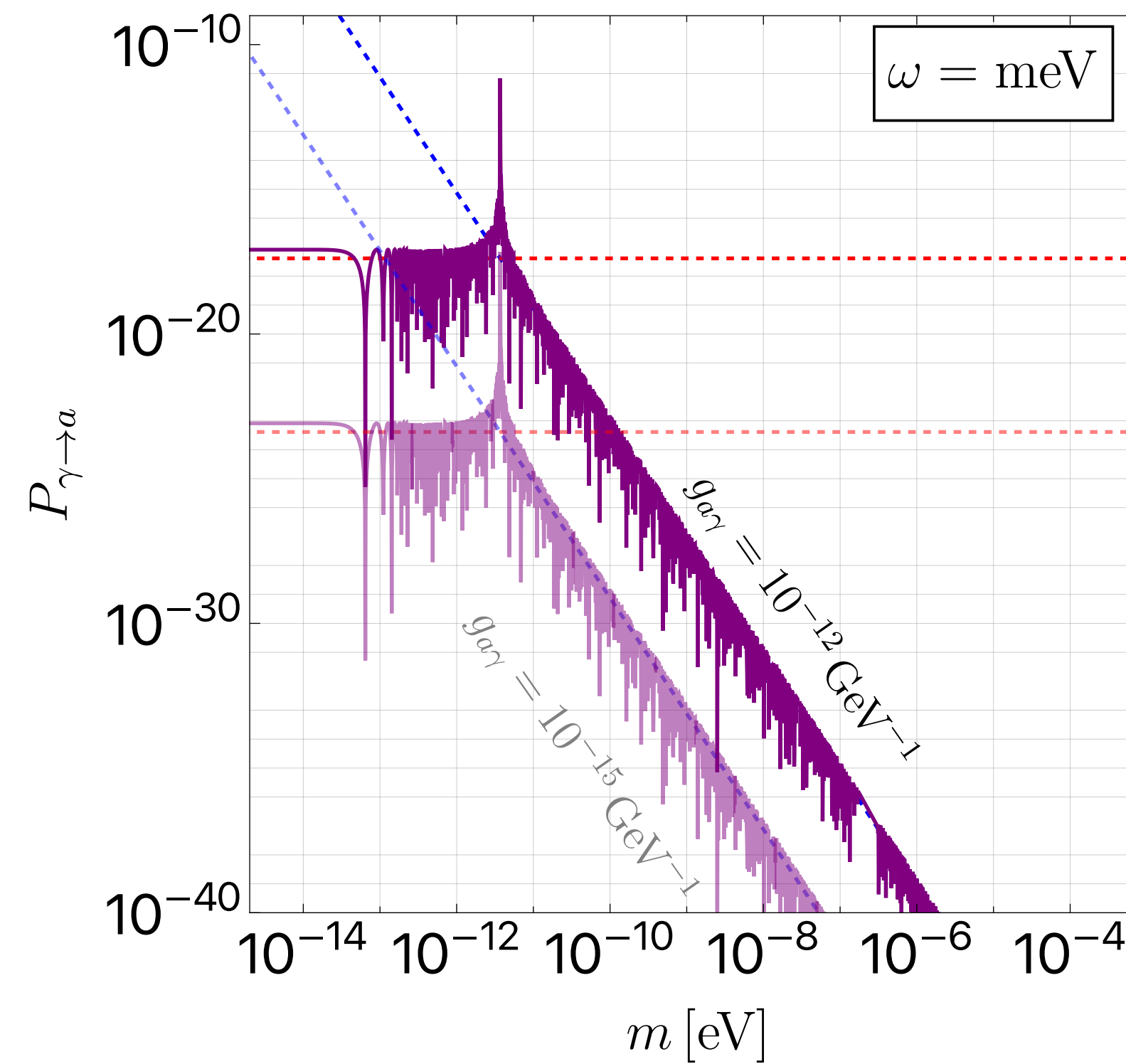
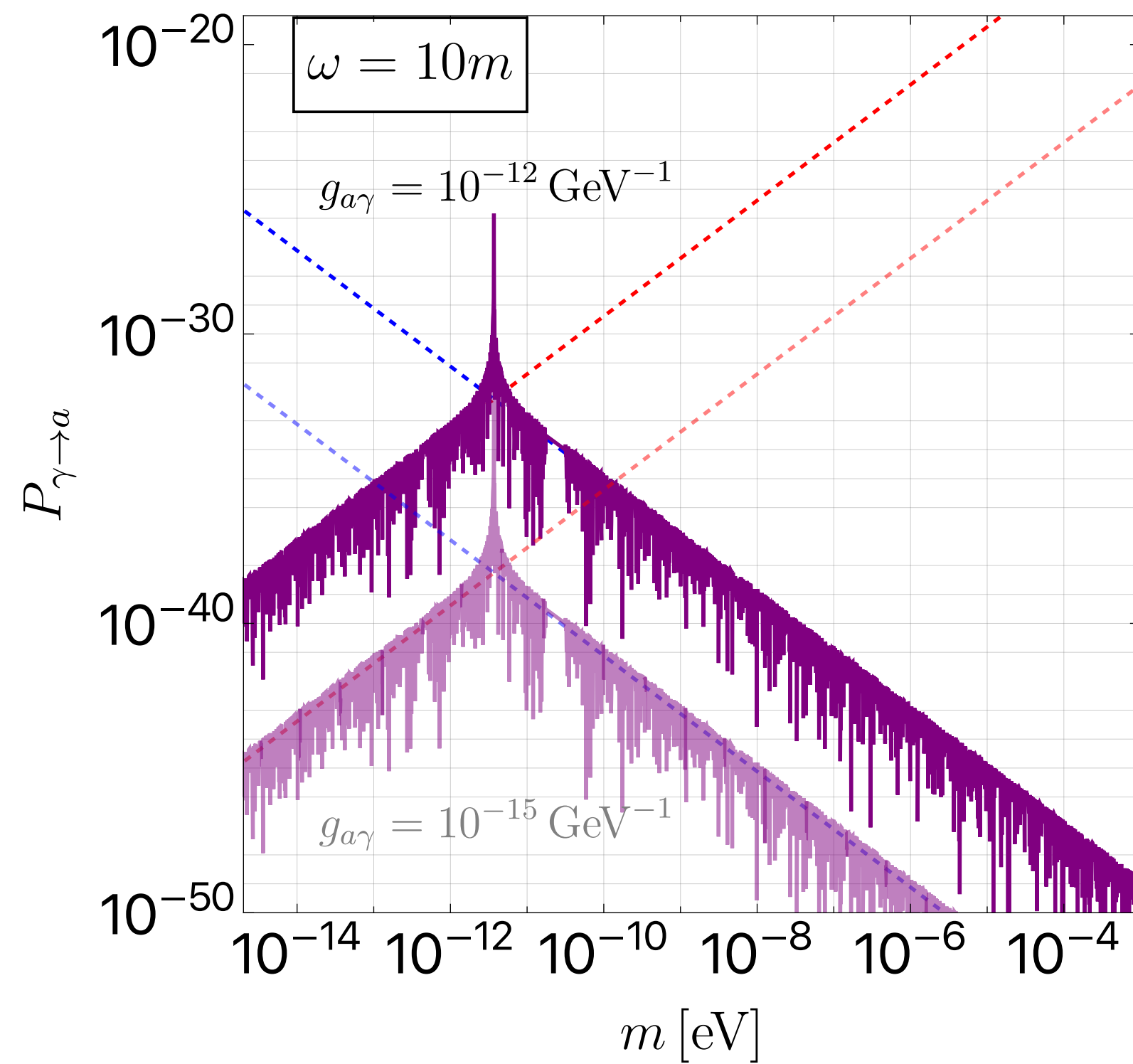
B-Field Conversion Probability

$$P_{\gamma \rightarrow a} = (\Delta_{a\gamma} R)^2 \frac{\sin^2(\Delta_{\text{osc}} R/2)}{(\Delta_{\text{osc}} R/2)^2}$$

$$\Delta_{a\gamma} \equiv \frac{g_{a\gamma} B_T}{2} \simeq 1.5 \cdot 10^{-4} \left(\frac{g_{a\gamma}}{10^{-12} \text{ GeV}^{-1}} \right) \left(\frac{B_T}{\mu\text{G}} \right) \text{ kpc}^{-1},$$

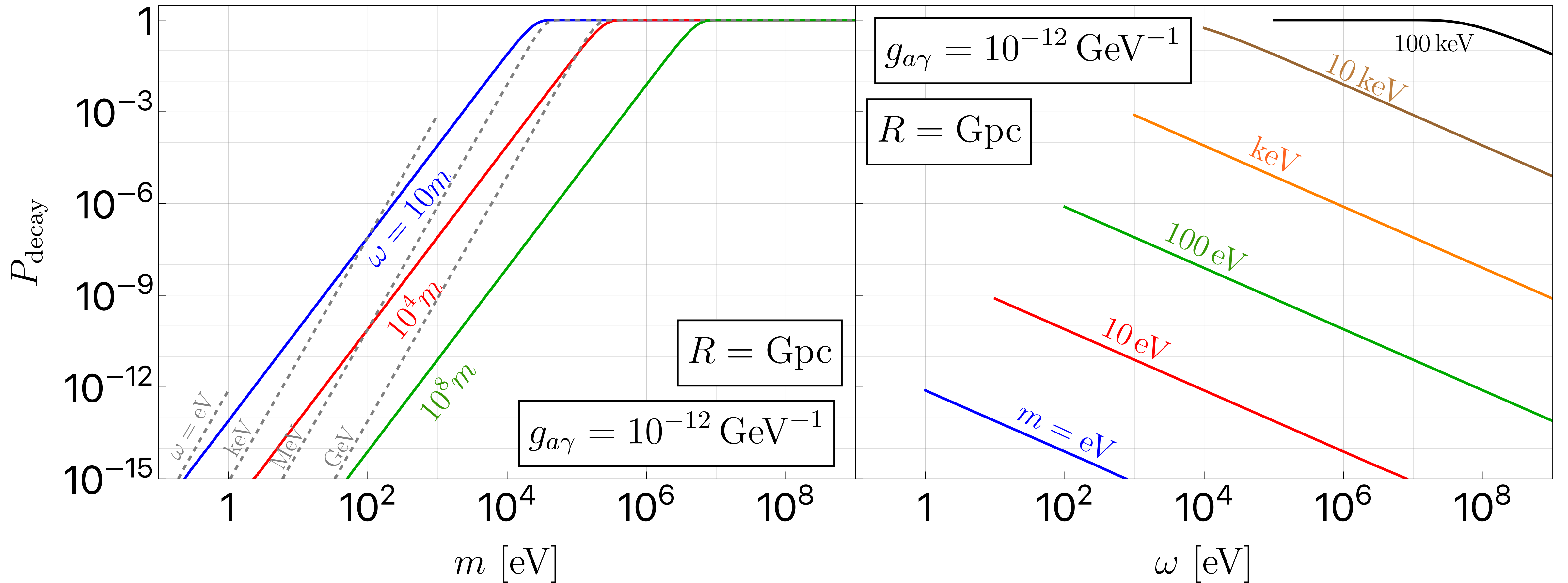
$$\Delta_a \equiv -\frac{m^2}{2\omega} \simeq -7.8 \cdot 10^{13} \left(\frac{m}{10^{-11} \text{ eV}} \right)^2 \left(\frac{10^{-10} \text{ eV}}{\omega} \right) \text{ kpc}^{-1},$$

$$\Delta_{\text{pl}} \equiv -\frac{\omega_{\text{pl}}^2}{2\omega} \simeq -1.1 \cdot 10^{13} \left(\frac{n_e}{10^{-2} \text{ cm}^{-3}} \right) \left(\frac{10^{-10} \text{ eV}}{\omega} \right) \text{ kpc}^{-1},$$



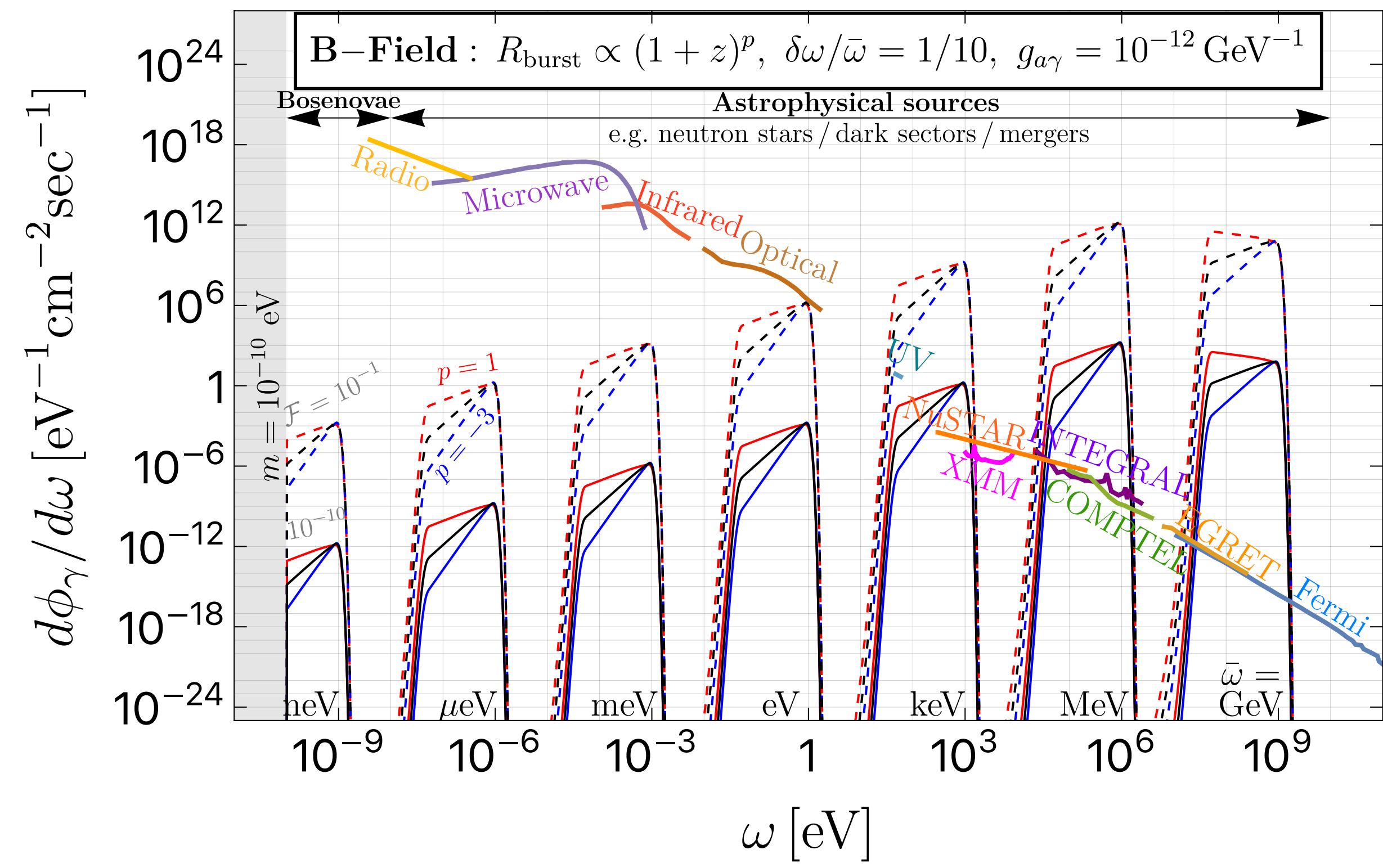
Decay Probability

$$l(\omega) \simeq \frac{\gamma v_{\text{burst}}}{\Gamma_{\gamma\gamma}} \simeq \left(\frac{\omega}{m}\right) \frac{64\pi}{g_{a\gamma}^2 m^3} \simeq \text{Mpc} \left(\frac{\omega}{\text{MeV}}\right) \left(\frac{100 \text{ keV}}{m}\right)^4 \left(\frac{10^{-12} \text{ GeV}^{-1}}{g_{a\gamma}}\right)^2$$

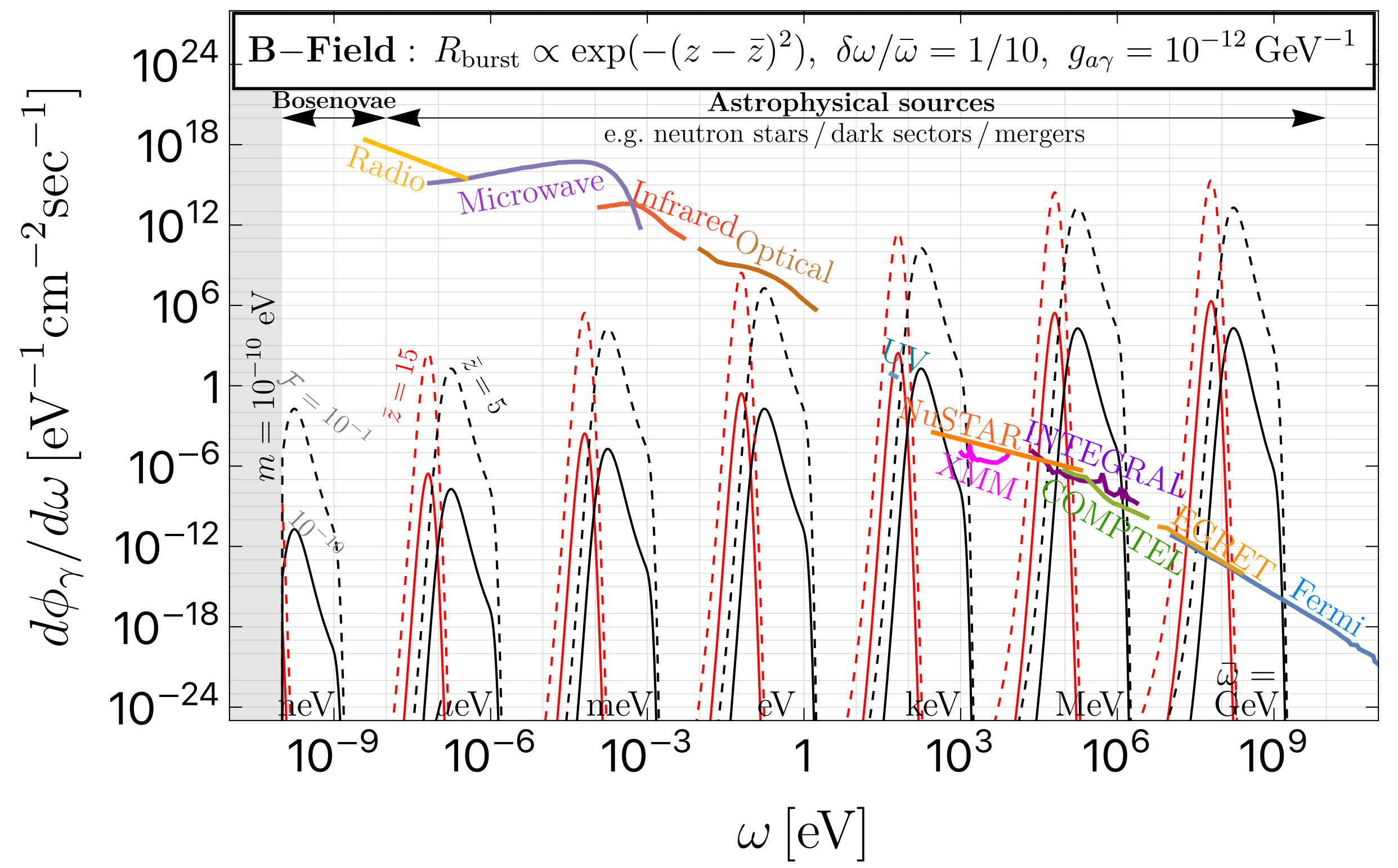


DaB Flux: Other $f(z)$

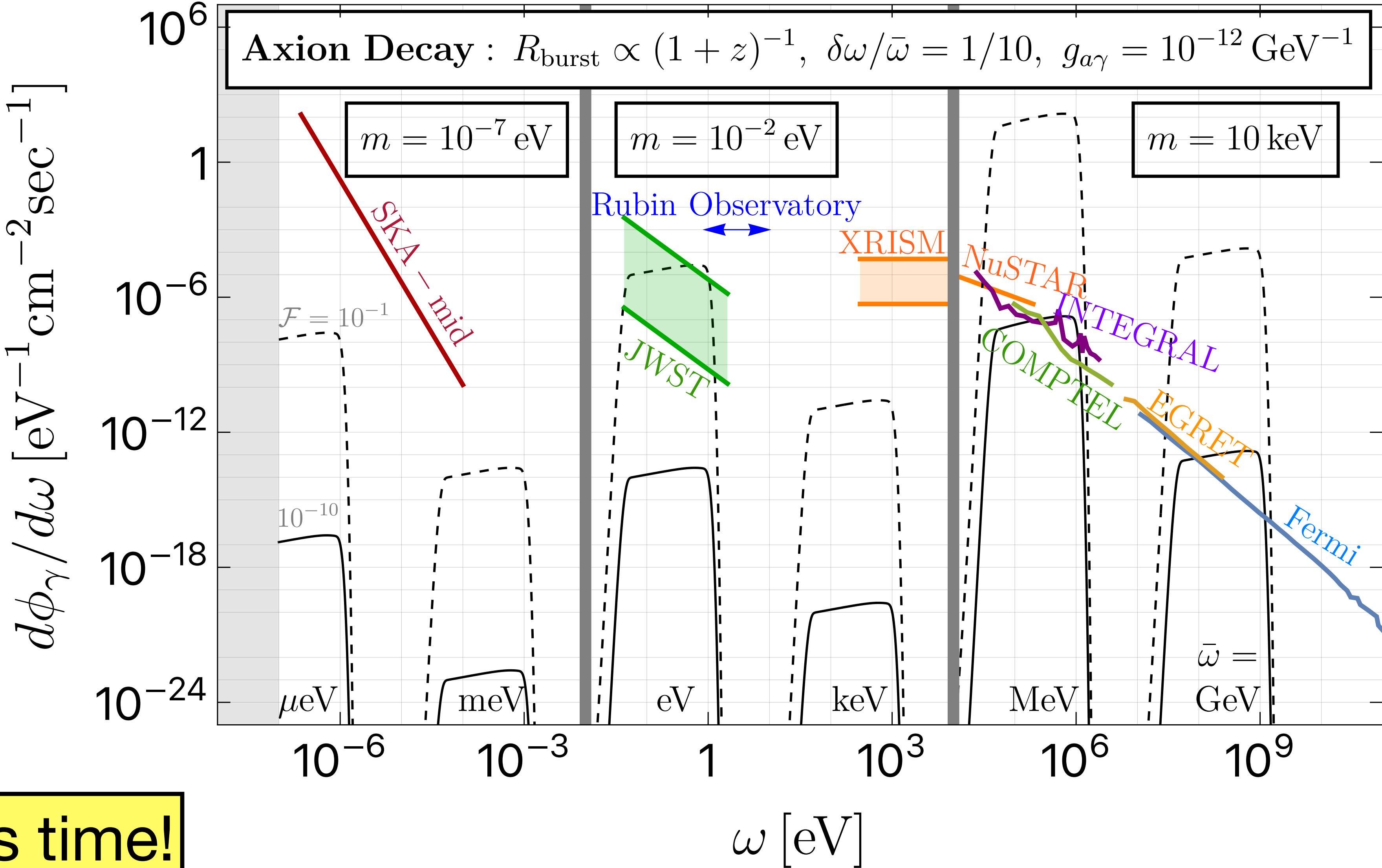
Power law



Gaussian



DaB Flux from Decay



If there's time!

Case Study: DaB From Bosenovae in SKA

