

Developing an amplitude level parton shower — CVolver

Fernando Torre González

With Jeff Forshaw and Simon Plätzer



The University of Manchester

Updating on progress made since the last
publication:
M. de Angelis, J. Forshaw, S. Plätzer (2020)
[2007.09648]



PSR24

Described in detail in Angeles Martinez, de Angelis, Forshaw, Plätzer, Seymour [1802.08531]

The soft evolution algorithm

We dress the hard process density matrix with iterative real and virtual operators:

$$\begin{aligned}
 \sigma_0 &= \text{Tr} \left(\mathbf{V}_{\mu,Q} \mathbf{H}(Q) \mathbf{V}_{\mu,Q}^\dagger \right) \equiv \text{Tr} \mathbf{A}_0(\mu) \\
 d\sigma_1 &= \text{Tr} \left(\mathbf{V}_{\mu,E_1} \mathbf{D}_1^\mu \mathbf{V}_{E_1,Q} \mathbf{H}(Q) \mathbf{V}_{E_1,Q}^\dagger \mathbf{D}_{1\mu}^\dagger \mathbf{V}_{\mu,E_1}^\dagger \right) d\Pi_1 \\
 &\equiv \text{Tr} \mathbf{A}_1(\mu) d\Pi_1, \\
 d\sigma_2 &= \text{Tr} \left(\mathbf{V}_{\mu,E_2} \mathbf{D}_2^\nu \mathbf{V}_{E_2,E_1} \mathbf{D}_1^\mu \mathbf{V}_{E_1,Q} \mathbf{H}(Q) \mathbf{V}_{E_1,Q}^\dagger \mathbf{D}_{1\mu}^\dagger \mathbf{V}_{E_2,E_1}^\dagger \mathbf{D}_{2\nu}^\dagger \mathbf{V}_{\mu,E_2}^\dagger \right) d\Pi_1 d\Pi_2 \\
 &\equiv \text{Tr} \mathbf{A}_2(\mu) d\Pi_1 d\Pi_2 \\
 &\vdots \\
 d\sigma_n &= \text{Tr} \mathbf{A}_n(\mu) \prod_{i=1}^n d\Pi_i
 \end{aligned}$$

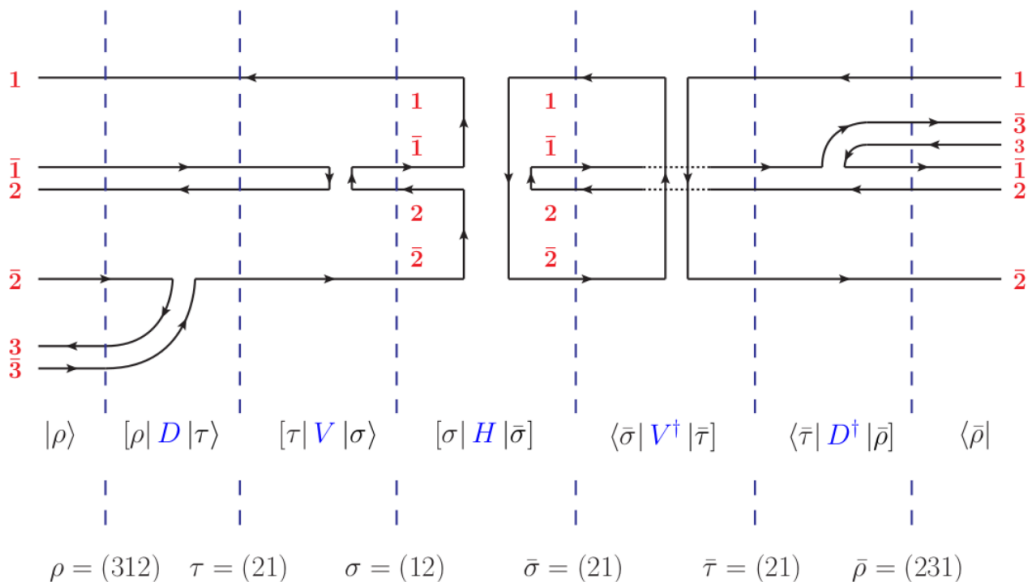
$$\mathbf{D}_i^\mu = \sum_j \mathbf{T}_j \frac{n_j^\mu}{n_j \cdot n_i}$$

$$\mathbf{V}_{a,b} = \exp \left(-\frac{\alpha_s}{\pi} \ln \left(\frac{b}{a} \right) \sum_{i < j} (-\mathbf{T}_i \cdot \mathbf{T}_j) \int \frac{d\Omega_k}{4\pi} \omega_{ij}(\hat{k}) \right)$$

By working in the colour flow basis, we can expand the exponentiated anomalous dimension as a series in $1/N_c$ Plätzer [1312.2448]:

$$[\tau | \mathbf{V}_{E,E'} | \sigma \rangle \simeq \delta_{\tau\sigma} R(\{\sigma\}) + \sum_{l=1}^d \left(-\frac{1}{N_c} \right)^l \sum_{\{\sigma_0, \dots, \sigma_l\}} \delta_{\tau\sigma_0} \delta_{\sigma_l, \sigma} \left(\prod_{\alpha=0}^{l-1} \Sigma_{\sigma_{\alpha+1}, \sigma_\alpha} \right) R(\{\sigma_0, \dots, \sigma_l\})$$

Keeping track of colour



At each step in the evolution, the colour state after the action of the real and virtual operators is sampled. This defines each event as a trajectory in colour space.

We count every factor of $1/N_c$ included at each step. There are four possible sources: the reals, the virtuals, the scalar product of the final colour states, and the hard process.

Thus, we can expand the cross-section in this way:

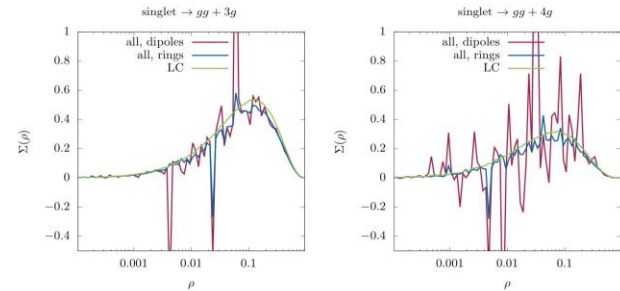
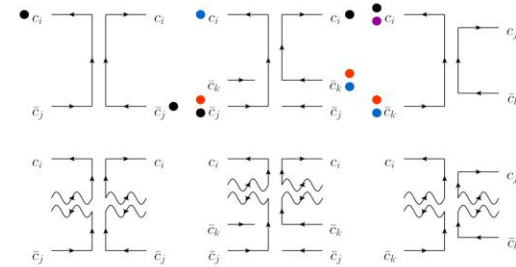
$$\sum_m N_c^m g_m \left(C_{0,m} + \frac{C_{1,m}}{N_c} + \frac{C_{2,m}}{N_c^2} + \dots \right)$$

Significant developments

- Implemented "rings and strings" in the calculation of real emission matrix elements, which exploit the collinear-finiteness of subleading colour.

See Forshaw, Holguin, Plätzer [2112.13124]

- Switched from an angular collinear cutoff to a rapidity-type cutoff.
- Added functionality to write event data in HDF5 format, allowing for reanalysis of data and reducing computing time significantly.
- A lot of different improvements to the sampling - notably the ability to count colour orders and specify the desired colour accuracy of the evolution.

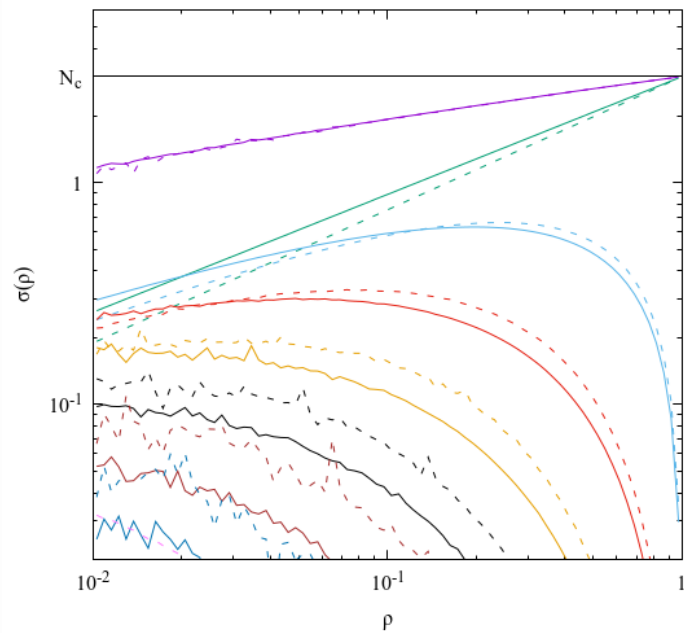


$$\sum_m N_c^m g_m \left(C_{0,m} + \frac{C_{1,m}}{N_c} + \frac{C_{2,m}}{N_c^2} + \dots \right)$$

We could force the evolution to stop here, for example

$e^+e^- \rightarrow q\bar{q}$ | gaps between jets

Collinear cutoff agreement



Full cross-section

0 emissions

1 emission

2 emissions

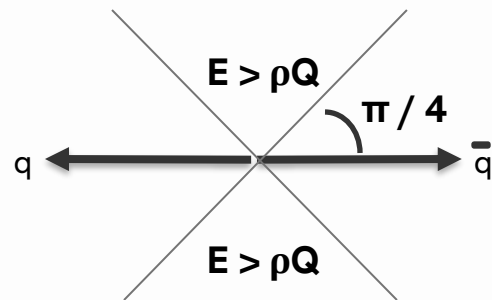
3 emissions

4 emissions

5 emissions

6 emissions

8 emissions

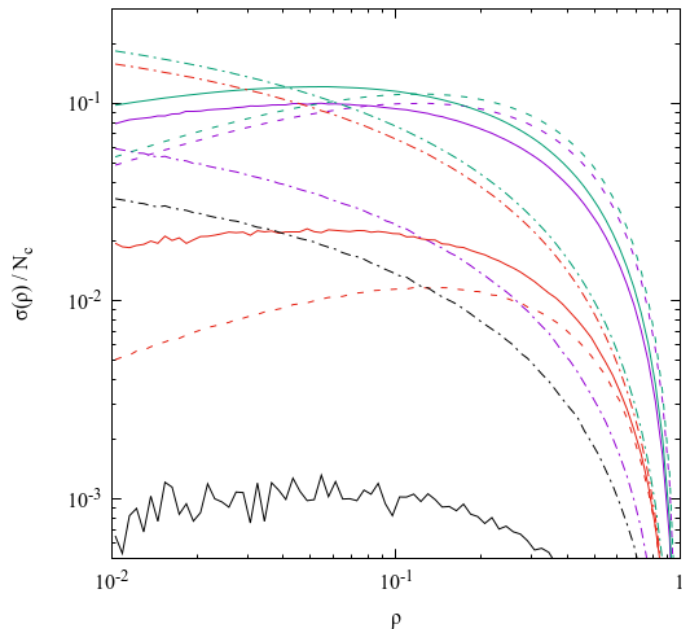


Collinear cutoff: 0.01

Collinear cutoff: 0.005

$e^+e^- \rightarrow q\bar{q}$ | gaps between jets

Breakdown in powers of $1/N_c$ for a specific multiplicity



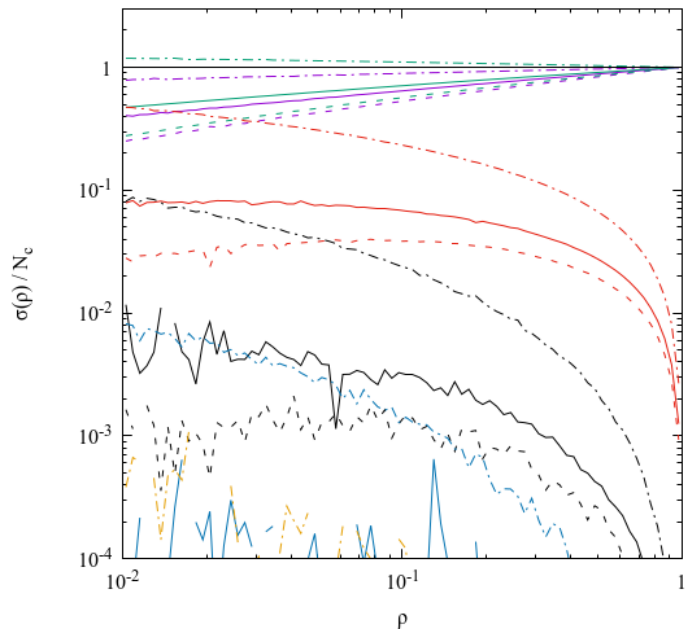
Full colour
 LC'
 $-(1/N_c^2)$: NLC'
 $(1/N_c^4)$: NNLC'

$N_c = 3$ ———
 $N_c = 4$ - - -
 $N_c = \text{sqrt}(2)$ - . - . -

$$N_c^2 g_2 \left(C_{0,2} + \frac{C_{1,2}}{N_c^2} + \frac{C_{2,2}}{N_c^4} \right)$$

$e^+e^- \rightarrow q\bar{q}$ | gaps between jets

Breakdown in powers of $1/N_c$ for the whole cross-section

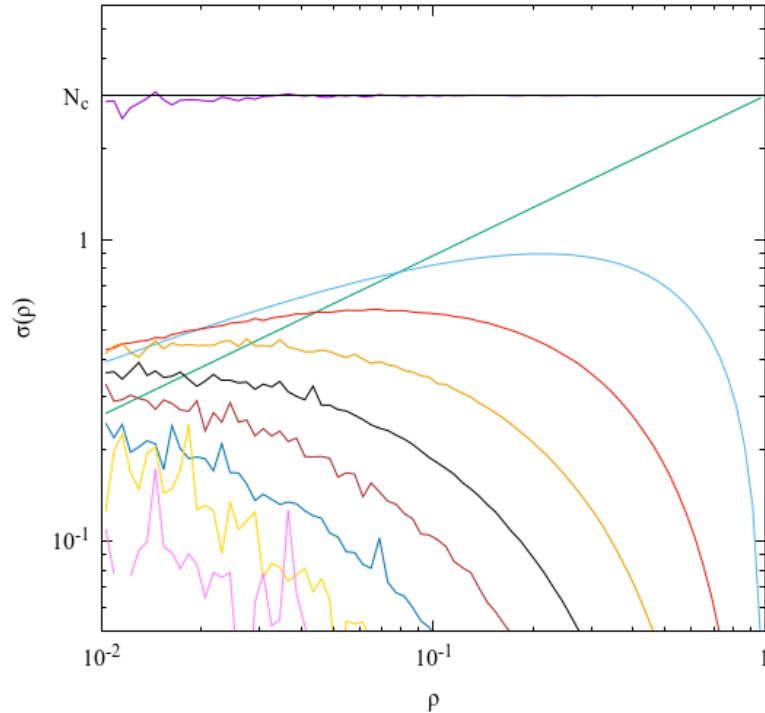


- Full colour
- LC'
- $-(1/N_c^2)$: NLC'
- $(1/N_c^4)$: NNLC'
- $-(1/N_c^6)$: N3LC'
- $(1/N_c^8)$: N4LC'

$$\sum_m N_c^m g_m \left(C_{0,m} + \frac{C_{1,m}}{N_c^2} + \frac{C_{2,m}}{N_c^4} + \frac{C_{3,m}}{N_c^6} + \frac{C_{4,m}}{N_c^8} + \dots \right)$$

- $N_c = 3$ —————
- $N_c = 4$ - - - - -
- $N_c = \text{sqrt}(2)$ - . - . - .

$e^+e^- \rightarrow q\bar{q}$ | unitarity check



Full cross-section

0 emissions

1 emission

2 emissions

3 emissions

4 emissions

5 emissions

6 emissions

7 emissions

8 emissions

We can run the evolution fully inclusive, accepting every emission.

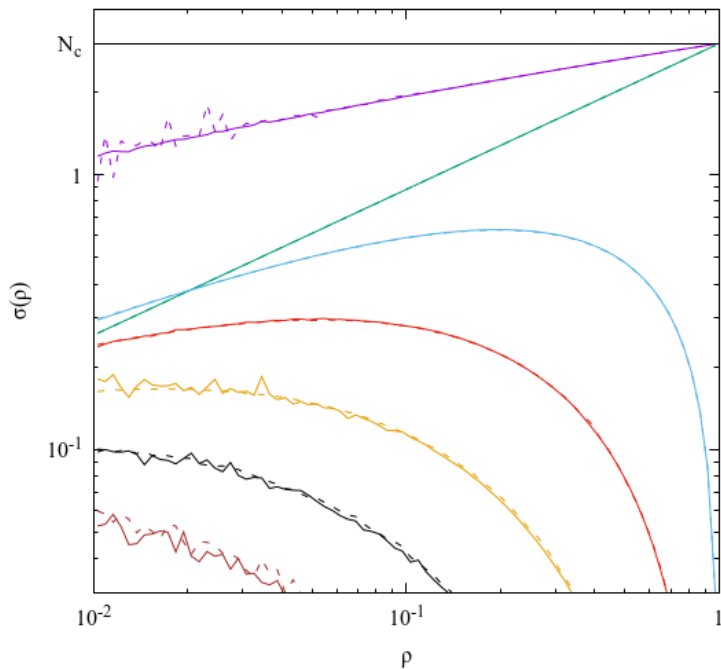
We expect all the virtual and real emissions to cancel, which serves as a sanity check

No gap...



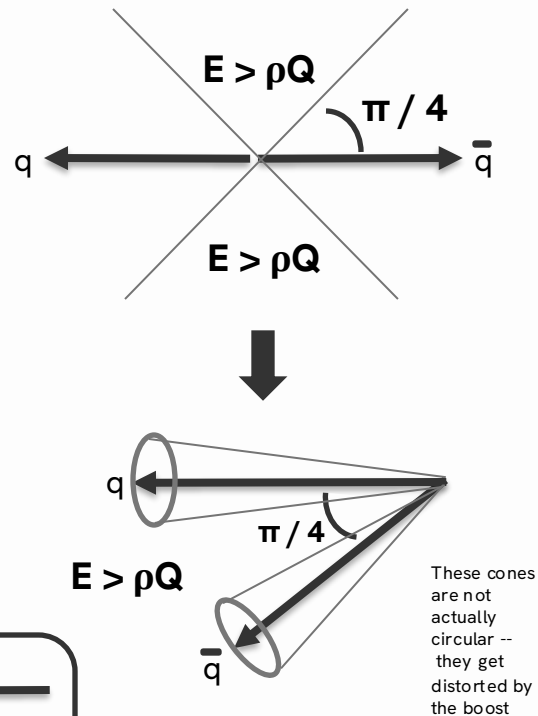
$e^+e^- \rightarrow q\bar{q}$ | gaps between jets

Lorentz invariance



- Full cross-section
- 0 emissions
- 1 emission
- 2 emissions
- 3 emissions
- 4 emissions
- 5 emissions

- Back-to-back frame
collinear cutoff = 0.01
- Boosted frame
collinear cutoff = 0.00146

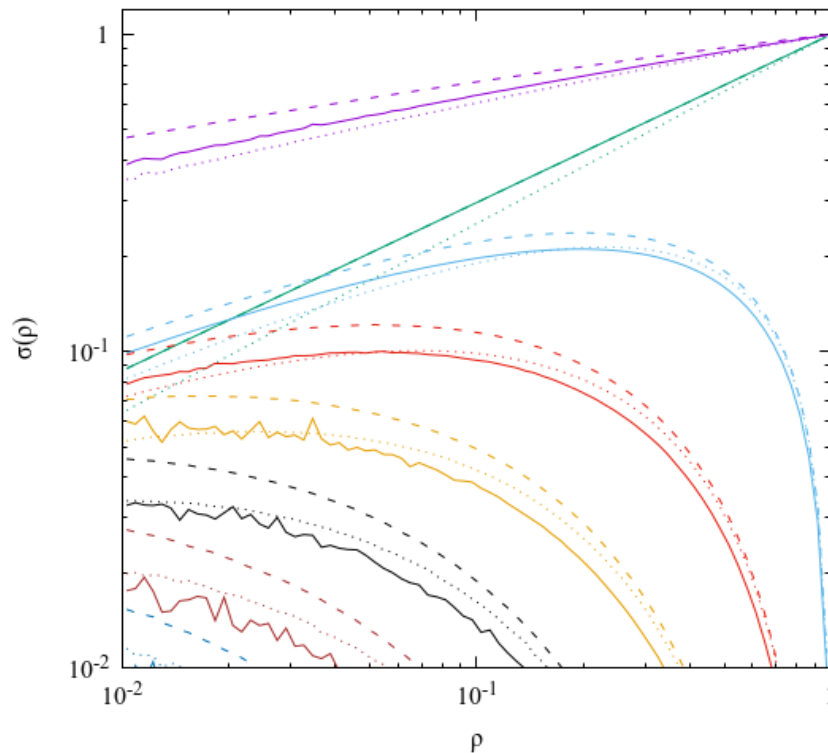
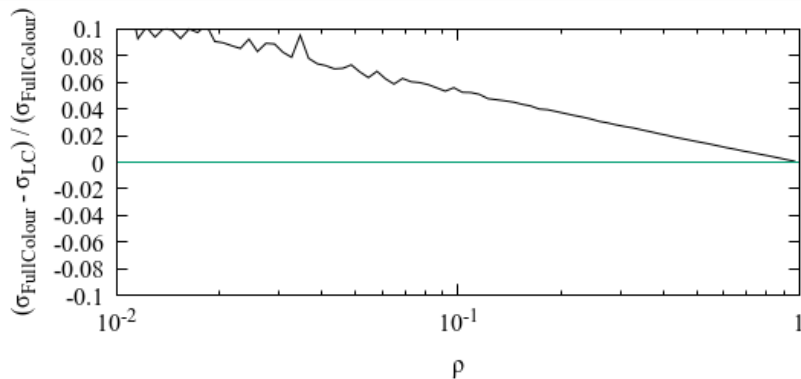


$e^+e^- \rightarrow q\bar{q}$ | gaps between jets

Effect of the subleading colour

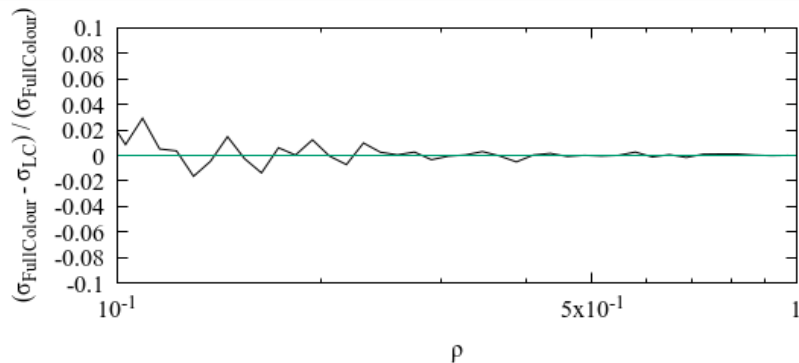
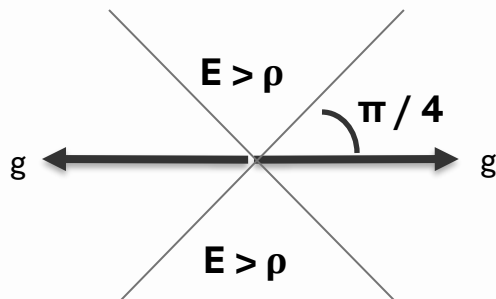
Full cross-section 0 emissions 1 emission
2 emissions 3 emissions 4 emissions 5 emissions

Full colour ———
LC' - - -
LC ·····



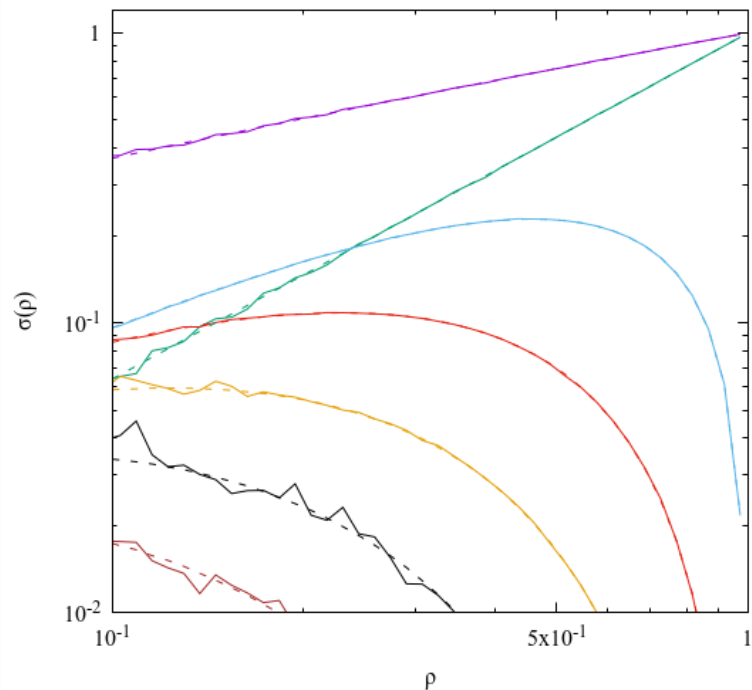
$e^+e^- \rightarrow gg$ | gaps between jets

Effect of the subleading colour



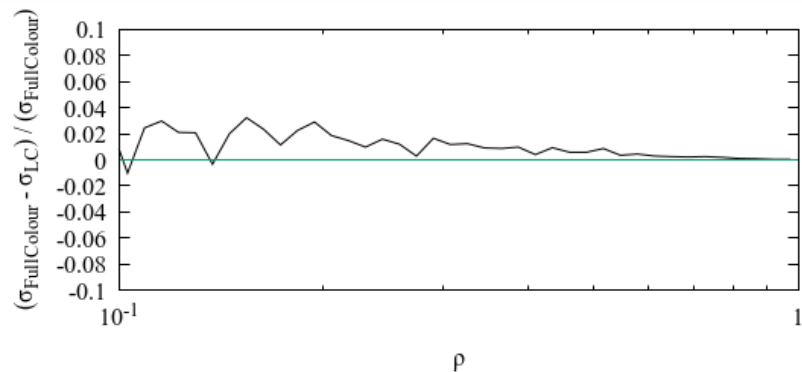
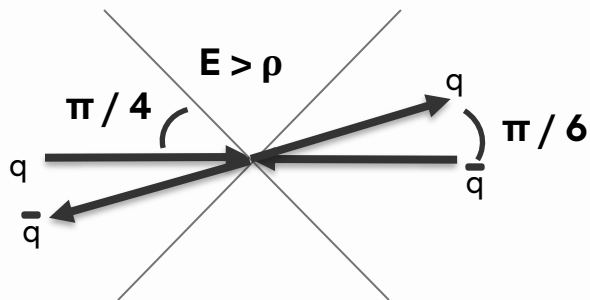
Full colour
 LC

Full cross-section 0 emissions 1 emission
 2 emissions 3 emissions 4 emissions 5 emissions



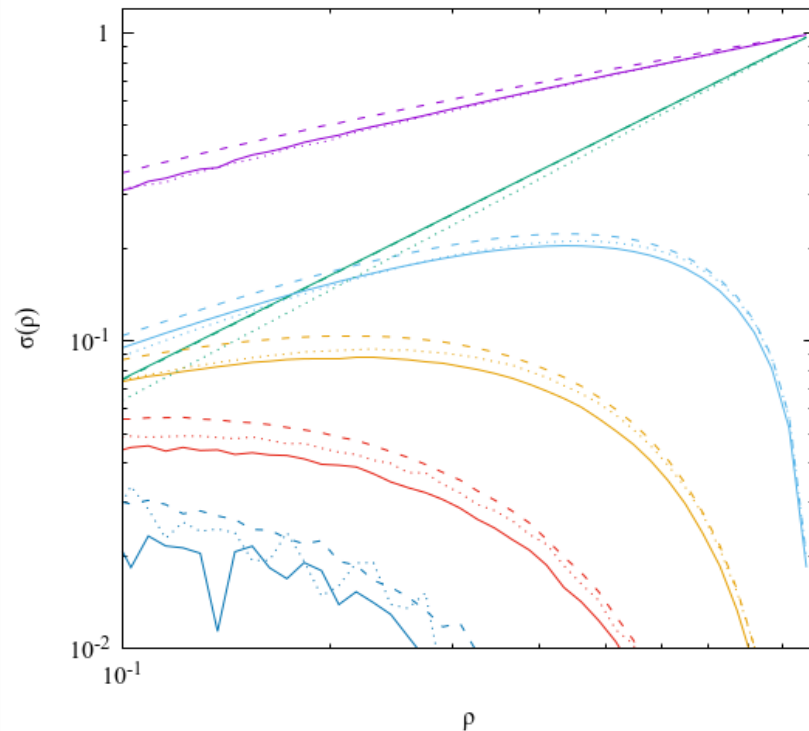
$q\bar{q} \rightarrow q\bar{q}$ | gaps between jets

Back-to-back configuration



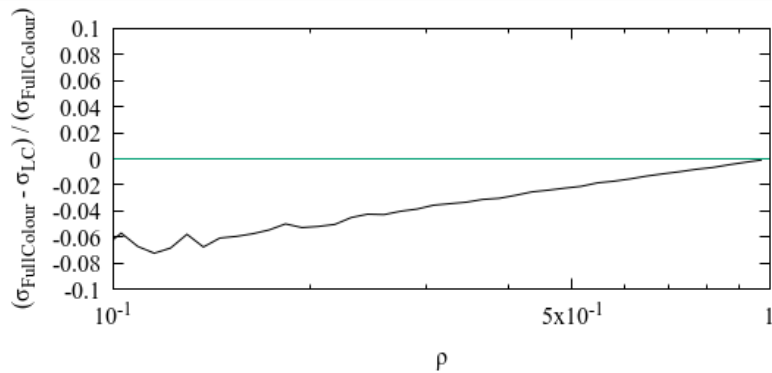
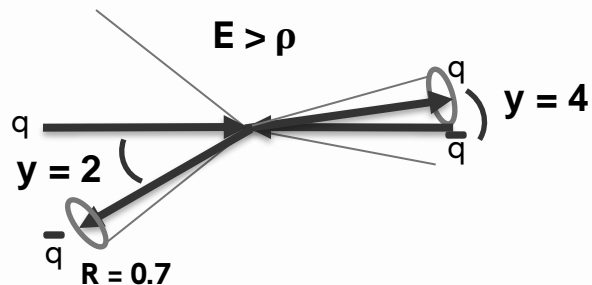
Full colour ———
 LC' - - -
 LC ·····

Full cross-section 0 emissions 1 emission
 2 emissions 3 emissions 4 emissions 5 emissions



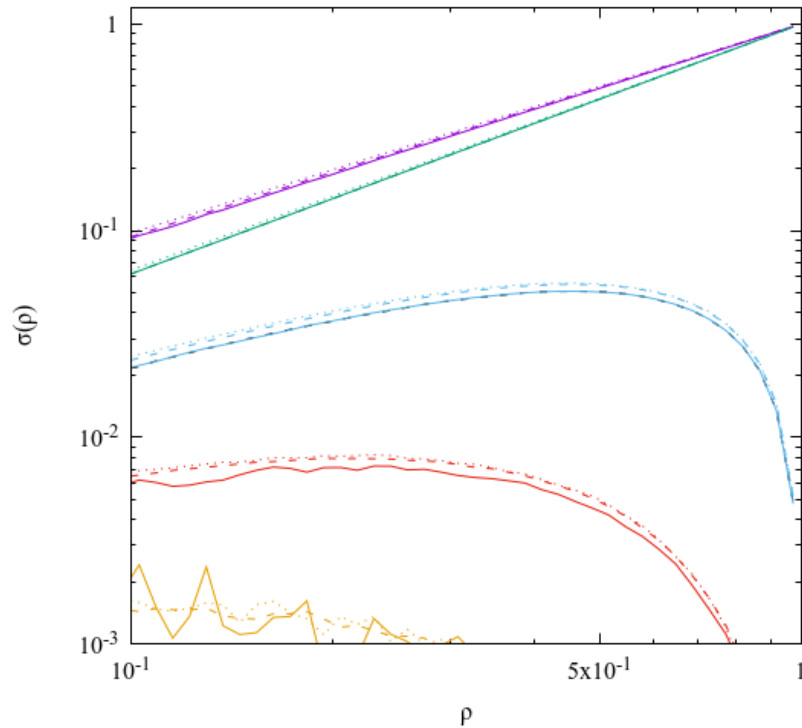
$q\bar{q} \rightarrow q\bar{q}$ | gaps between jets

Asymmetric final jet rapidities



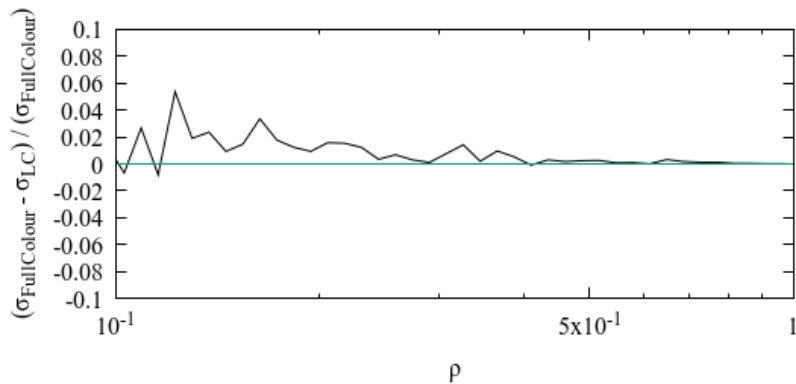
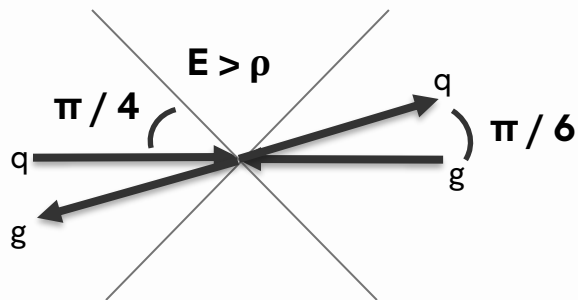
Full colour ———
 LC' - - - -
 LC
 (Note: The legend indicates that the solid line represents Full colour, the dashed line represents LC', and the dotted line represents LC.)

Full cross-section 0 emissions 1 emission
 2 emissions 3 emissions 4 emissions 5 emissions
 (Note: The legend indicates that different colors represent different numbers of emissions: purple for 0, green for 1, red for 2, orange for 3, blue for 4, and yellow for 5.)



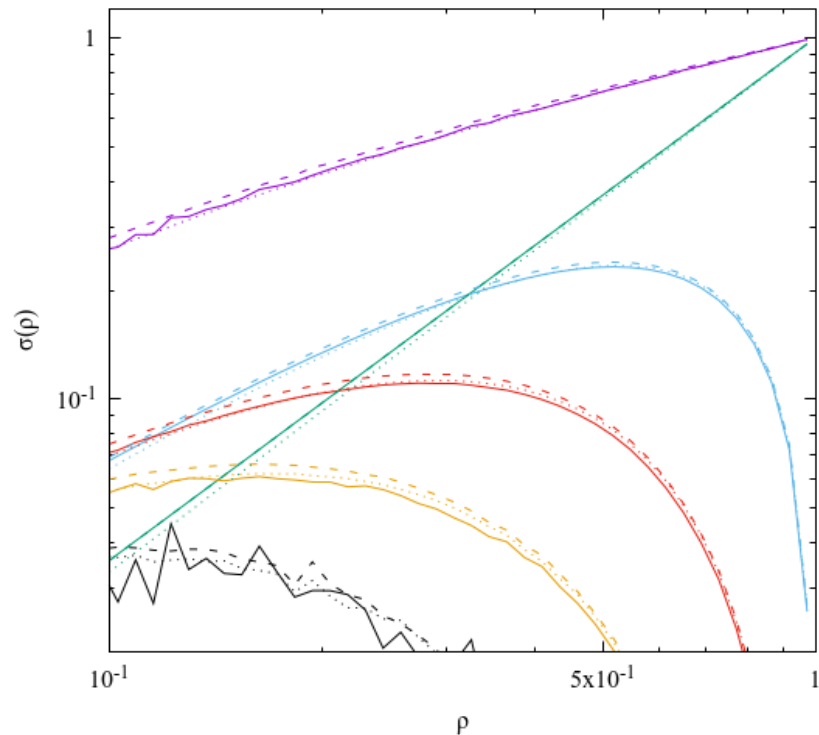
qg → qg | gaps between jets

Back-to-back configuration



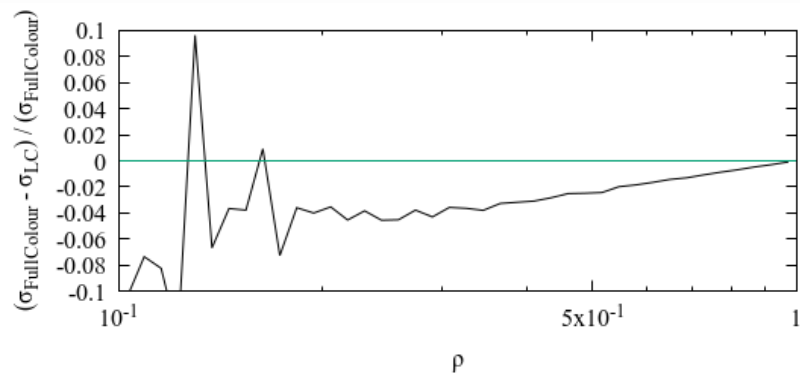
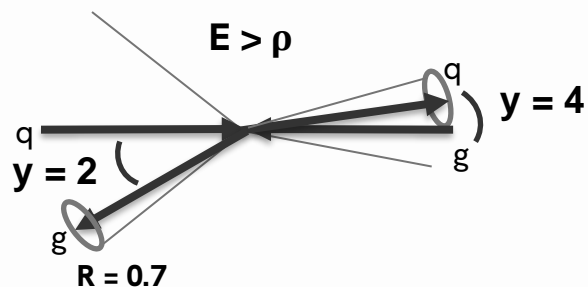
Full colour ———
 LC' - - -
 LC ·····

Full cross-section
 0 emissions (purple)
 1 emission (green)
 2 emissions (red)
 3 emissions (orange)
 4 emissions (blue)
 5 emissions (black)



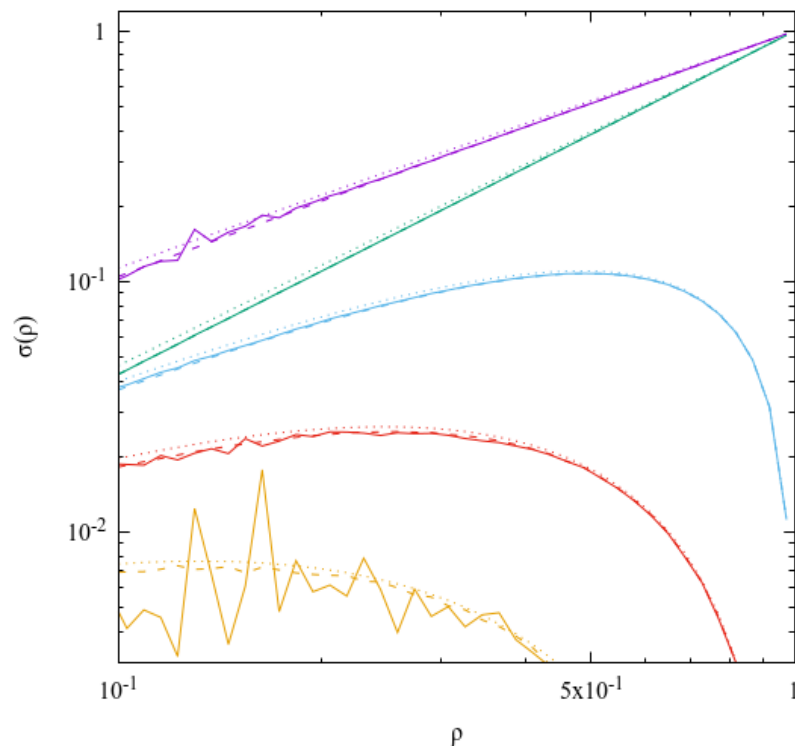
$qg \rightarrow qg$ | gaps between jets

Asymmetric final jet rapidities



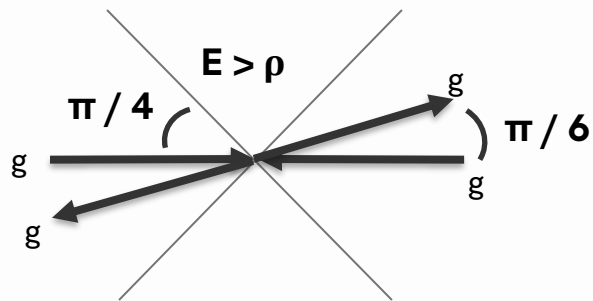
Full colour ———
 LC' — — — — —
 LC ·······

Full cross-section 0 emissions 1 emission
 2 emissions 3 emissions 4 emissions 5 emissions



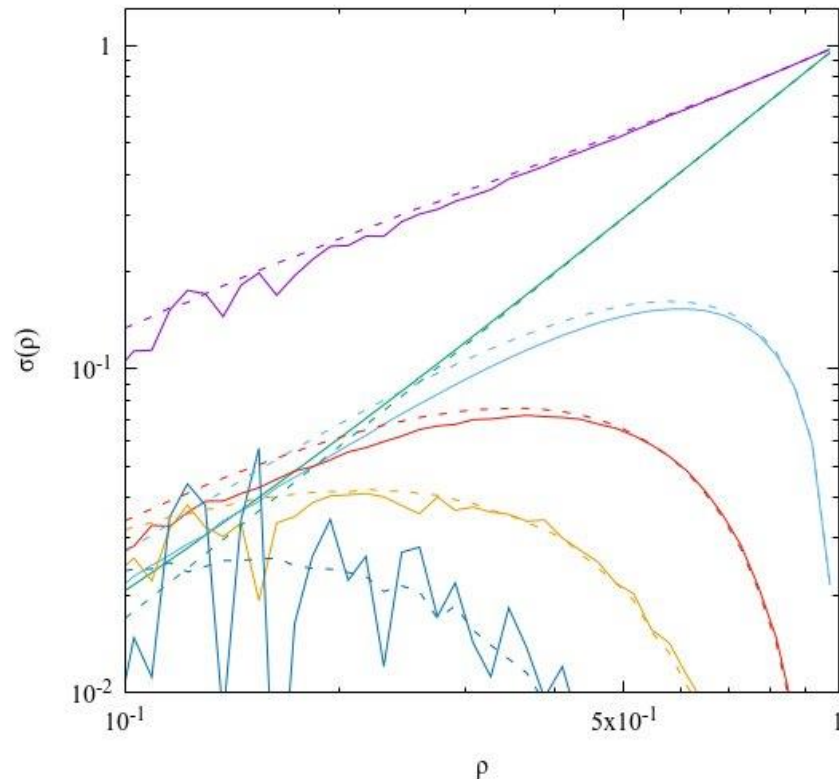
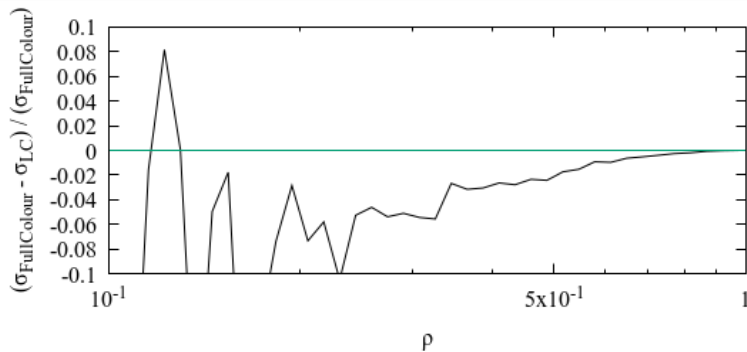
gg → gg | gaps between jets

Back-to-back configuration



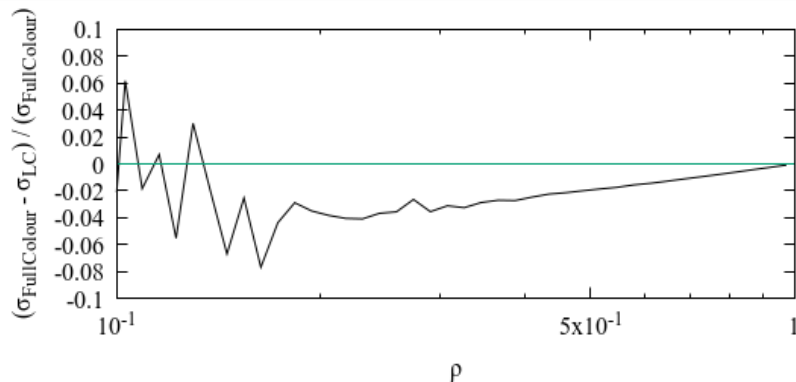
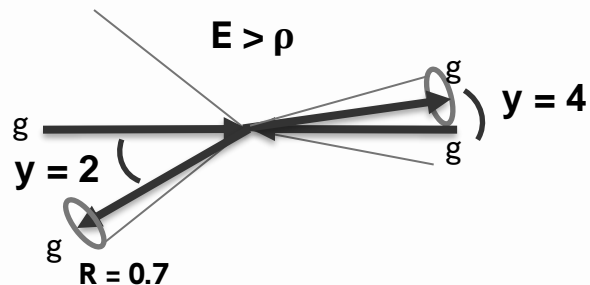
Full colour
 LC

Full cross-section 0 emissions 1 emission
 2 emissions 3 emissions 4 emissions 5 emissions



$gg \rightarrow gg$ | gaps between jets

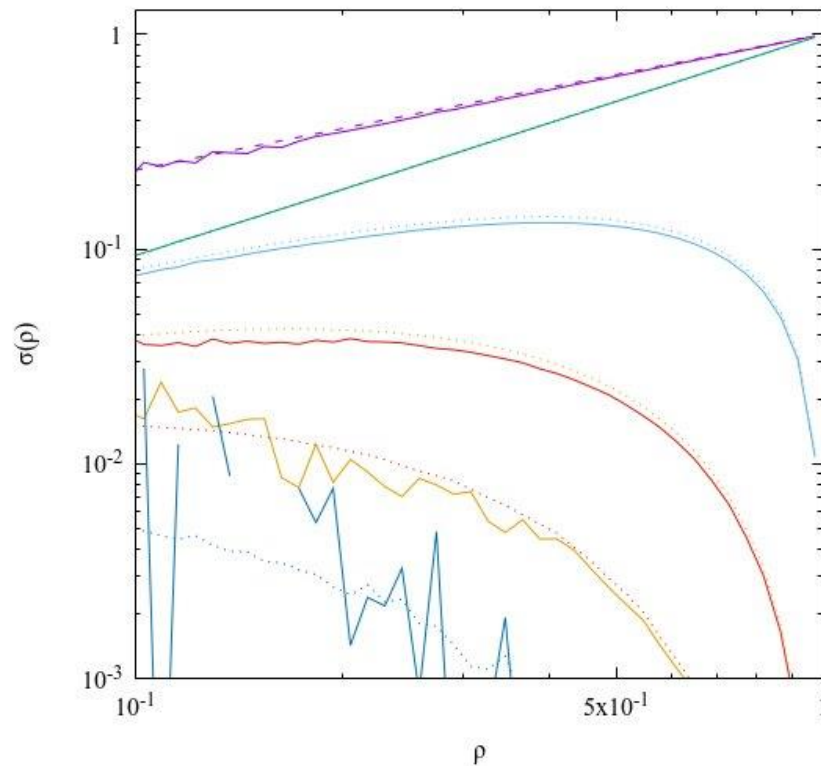
Asymmetric final jet rapidities



Full colour
 LC

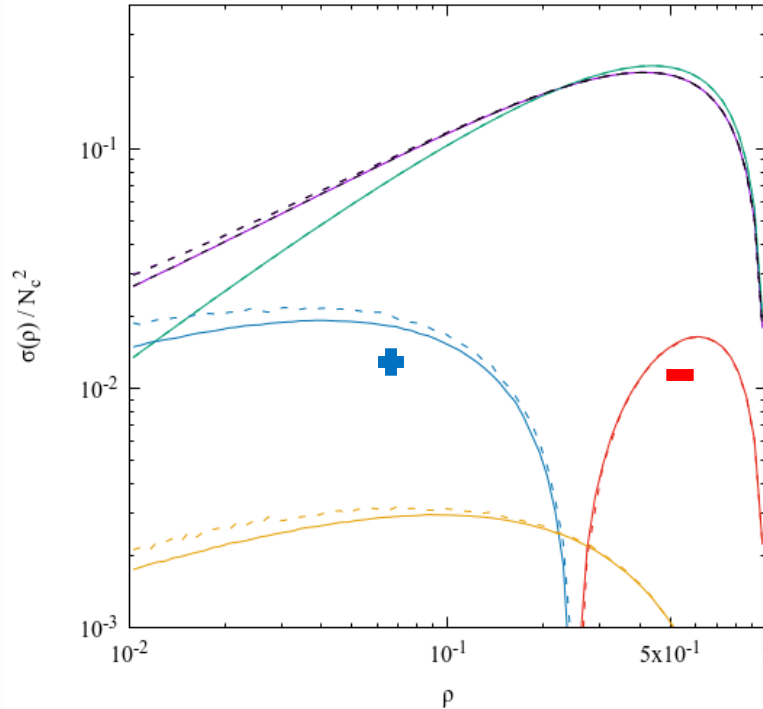
Full cross-section 0 emissions 1 emission

2 emissions 3 emissions 4 emissions 5 emissions



$q\bar{q} \rightarrow q\bar{q}$ | Coulomb exchanges

1 emission, broken down in powers of $1/N_c$



- Full colour
- LC'
- $-(1/N_c^2)$: NLC'
- $+(1/N_c^2)$: NLC'
- $+(1/N_c^4)$: NNLC'

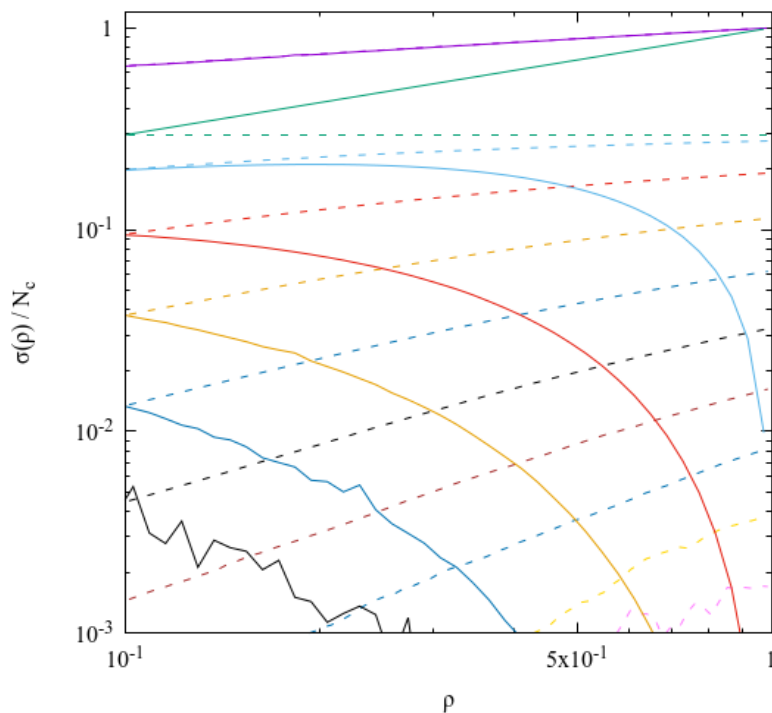
- No coulomb gluons
- With coulomb gluons

We can simply insert the Coulomb/Glauber phases in the anomalous dimension, and the algorithm is able to resum them.

Perfect agreement with independent calculation.

$$V_{a,b} = \exp \left(-\frac{\alpha_s}{\pi} \ln \left(\frac{b}{a} \right) \sum_{i < j} (-\mathbf{T}_i \cdot \mathbf{T}_j) \left(\int \frac{d\Omega_k}{4\pi} \omega_{ij}(\hat{k}) - i\pi \tilde{\delta}_{ij} \right) \right)$$

$e^+e^- \rightarrow \bar{q}q$ | event generator mode



Full cross-section

0 emissions

1 emission

2 emissions

3 emissions

4 emissions

5 emissions

6 emissions

7 emissions

8 emissions

Until now we have only discussed using CVolver as a resummation tool for gaps between jets.

But it can also be run as an event generator -- evolving all events down to some infra-red scale μ , and then applying the measurement function to generate a differential cross-section.

Resummation



Event generation

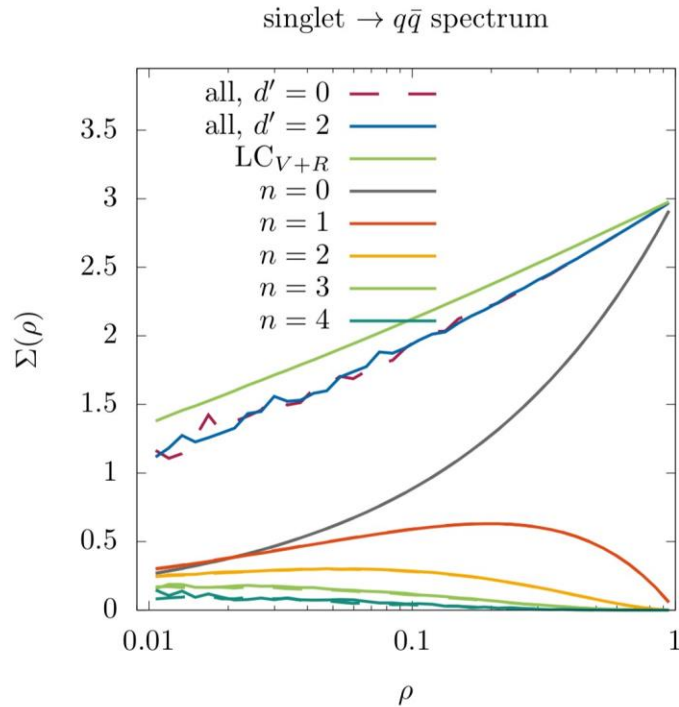


Summary and next steps

- We have implemented a full colour soft evolution algorithm that keeps track of every source of subleading colour in a systematic way.
- We have performed every cross-check we can think of.
- It has provided plenty of evidence that there is significant non-trivial subleading colour coming in at phenomenologically relevant scales.
- It is ready to produce lots of other interesting physics -- for example we can study differential observables in event generator mode (like colour reconnection) and Glauber effects.
- The algorithm design is such that we can "plug in" new features, which are under development:
 - Full kinematics parton shower, with hard process elements.
 - k_t ordering.
 - Hard-collinear physics
See Forshaw, Holguin, Plätzer [1905.08686]



Appendix: PRL plot for qqbar



$e^+e^- \rightarrow W^+W^- \rightarrow q\bar{q}q\bar{q}$

Colour reconnection

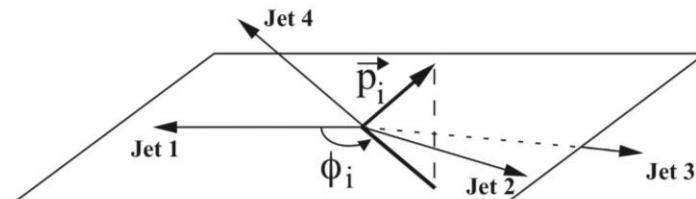
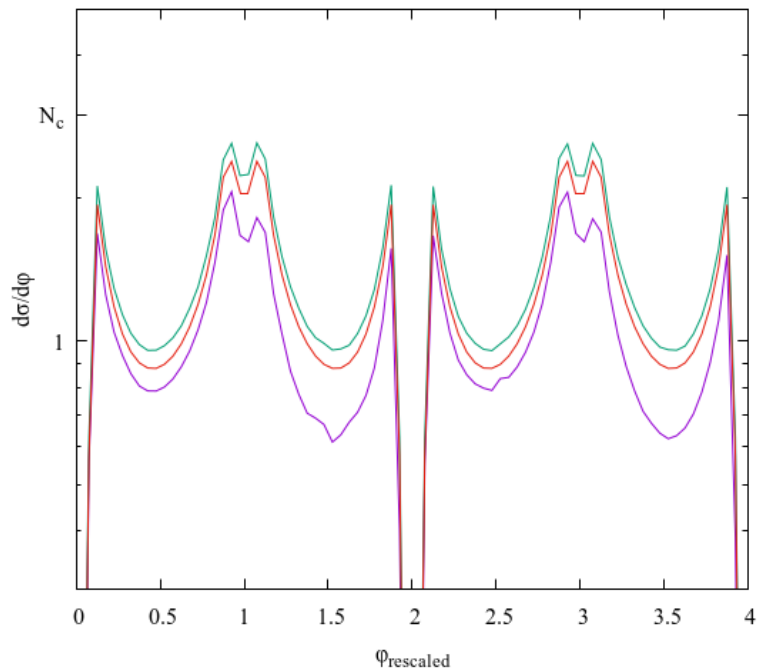


Figure taken from the L3 collaboration [hep-ex/0303042]

- Full colour
- LC'
- LC