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The Inert Doublet Model: Insights into Dark Matter

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Outline

- Cosmological Standard Model
- Inert Doublet Model (IDM)
- Bridge between particle physics and cosmology
- Extensions of IDM

From Gravitation to Cosmology





 $\Omega_{total} = \Omega_m^{\approx 0.31} + \Omega_r^{\approx 5 \times 10^{-5}} + \Omega_\Lambda^{\approx 0.69} + \Omega_k^{\approx 0} \approx 1$

- Galactic Rotation Curves: The observed rotation speed of stars in galaxies remains constant far from the center, contrary to predictions based on visible matter alone.
- Gravitational Lensing: Light from distant galaxies is bent more than can be explained by visible matter, indicating the presence of unseen mass.
- Cosmic Microwave Background (CMB): The patterns in the CMB require dark matter to explain the observed density fluctuations.
- Galaxy Cluster Dynamics: The motion of galaxies within clusters shows more mass than can be accounted for by visible matter.
 Structure Formation: The formation of galaxies and large-scale structures in the universe is best explained by cold dark matter models.

Good DM candidate

- Dark matter is a form of matter that does not emit, absorb, or reflect electromagnetic radiation, making it invisible to telescopes. Its presence is inferred through its gravitational effects on visible matter, and on the cosmic microwave background (CMB).
- It constitutes approximately 84% of the matter in the universe and 27% of the total mass-energy content of the universe, while ordinary (baryonic) matter makes up about 5%.
- Dark matter is described as being cold (velocity is non-relativistic at the epoch of radiation-matter equality),
- Non-baryonic (consisting of matter other than protons and neutrons),
- Dissipationless (cannot cool by radiating photons)

- Collisionless (i.e., the dark matter particles interact with each other and other particles only through gravity and possibly the weak force).
- Evidence suggests it does not interact strongly with itself or baryonic matter, apart from weak interactions. (because such interactions would produce observable effects that are inconsistent with astrophysical and cosmological observations.). Strong selfinteractions would lead to core collapse in dark matter halos or the formation of dense cores, which are not consistent with observations.
- Stable: If dark matter were unstable and decayed, Decay products would be observable; and alter the anisotropies in the CMB, which are well-matched to current models of stable dark matter.

Call for Particle Physics

Looking in the Standard Model

- Quarks carry <u>color charges</u> (red, blue, or green) so they participate in strong interactions.
- All particles have <u>electric charge</u> participate in electromagnetic interaction.
- Quarks and leptons carry <u>flavor charges</u> and participate in weak nuclear interactions.

The Inert Doublet Model

- One of the simplest and most studied BSM scenario is the inert doublet model (IDM) where the Higgs sector of the SM is extended by an additional electroweak doublet with a stabilizing discrete Z₂ symmetry.
- The lightest neutral component of the new doublet serves as a DM candidate.
- In this scenario, DM interacts with the SM sector via its interaction with the Higgs boson, and weak gauge interactions.

$$\Phi_{h} = \begin{pmatrix} \Phi_{h}^{+} \\ \Phi_{h}^{0} \end{pmatrix} = \begin{pmatrix} G^{+} \\ \frac{1}{\sqrt{2}}(v+h+iG^{0}) \end{pmatrix}$$

This Z₂ symmetry ensures that:

- Φ_{inert} does not participate in electroweak symmetry breaking,
- does not couple to the SM fermion
- the lightest neutral component of Φ_{inert} is stable and can serve as a dark matter candidate.

The scalar sector of the model is described by the potential

 $V(\Phi_h, \Phi_{inert}) = \left\{ \begin{array}{l} -\mu_h^2 |\Phi_h|^2 + \lambda_h |\Phi_h|^4 + \mu_{inert}^2 |\Phi_{inert}|^2 + \lambda_{inert} |\Phi_{inert}|^4 \\ + \lambda_3 |\Phi_h|^2 |\Phi_{inert}|^2 + \lambda_4 |\Phi_h^{\dagger} \Phi_{inert}|^2 + \frac{\lambda_5}{2} \left[(\Phi_h^{\dagger} \Phi_{inert})^2 + h.c. \right] \end{array} \right\}$

 $\Phi_{inert} = \begin{pmatrix} \Phi_{inert}^{+} \\ \Phi_{inert}^{0} \end{pmatrix} = \begin{pmatrix} H^{+} \\ \frac{1}{\sqrt{2}}(H^{0} + iA^{0}) \end{pmatrix}$

 Z_2 symmetry

 $\Phi_h \to \Phi_h$, $\Phi_{inert} \to -\Phi_{inert}$

 H^0 and A^0 are lighter than H^{\pm}

 H^0 the lightest particle in the inert doublet

$$V(\Phi_{h}, \Phi_{inert}) = \begin{cases} \mu_{1}^{2} |\Phi_{h}|^{2} + \mu_{2}^{2} |\Phi_{inert}|^{2} + \frac{1}{2} \lambda_{1} |\Phi_{h}|^{4} + \frac{1}{2} \lambda_{2} |\Phi_{inert}|^{4} \\ + \lambda_{3} |\Phi_{h}|^{2} |\Phi_{inert}|^{2} + \lambda_{4} |\Phi_{h}^{\dagger} \Phi_{inert}|^{2} + \frac{\lambda_{5}}{2} [(\Phi_{h}^{\dagger} \Phi_{inert})^{2} + h.c.] \end{cases}$$

Mass Spectrum

$$M_{h}^{2} = 2\lambda_{h}v^{2} = 2\mu_{h}^{2}$$
$$M_{H^{0}}^{2} = \mu_{inert}^{2} + \frac{1}{2}(\lambda_{3} + \lambda_{4} + \lambda_{5})v^{2}$$
$$M_{A^{0}}^{2} = \mu_{inert}^{2} + \frac{1}{2}(\lambda_{3} + \lambda_{4} - \lambda_{5})v^{2}$$
$$M_{H^{\pm}}^{2} = \mu_{inert}^{2} + \frac{1}{2}\lambda_{3}v^{2}$$

Theoretical constraints

Regarding the theoretical constraints, we incorporate considerations of perturbativity, stability, and tree-level unitarity in the scalar sector

vacuum stability

$$\begin{split} \lambda_1 &> 0, \\ \lambda_2 &> 0, \\ \lambda_3 &> -\sqrt{\lambda_1 \lambda_2}, \\ \lambda_3 &+ \lambda_4 - |\lambda_5| > -\sqrt{\lambda_1 \lambda_2}, \end{split}$$

unitarity

$$\begin{split} & \beta(\lambda_{1} + \lambda_{2}) + \sqrt{9(\lambda_{1} - \lambda_{2})^{2} + (2\lambda_{3} + \lambda_{4})^{2}} < 16\pi, \\ & \lambda_{1} + \lambda_{2} + \sqrt{(\lambda_{1} - \lambda_{2})^{2} + 4\lambda_{4}^{2}} < 16\pi, \\ & \lambda_{1} + \lambda_{2} + \sqrt{(\lambda_{1} - \lambda_{2})^{2} + 4\lambda_{5}^{2}} < 16\pi, \\ & |\lambda_{3} + 2\lambda_{4} \pm 3\lambda_{5}| < 8\pi, \\ & |\lambda_{3} \pm \lambda_{4}| < 8\pi, \\ & |\lambda_{3} \pm \lambda_{5}| < 8\pi. \end{split}$$

perturbativity

 $|\lambda_i| < 4\pi$

Experimental constraints

Production of inert scalars at the LEP: No signals have been reported by LEP for processes like :

 $Z \rightarrow (H^0 A^0, H^+ H^-)$ $W^{\pm} \rightarrow (H^0 H^{\pm}, A^0 H^{\pm})$

This leads to the conditions

$$m_{H^0,A^0} + m_{H^{\pm}} > m_W$$

 $m_{A^0} + m_{H^0} > m_Z$
 $2m_{H^{\pm}} > m_Z$

https://arxiv.org/pdf/2410.00638

Higgs invisible decays: When the inert scalars have masses lower than half of Higgs masse, certain decay modes become kinematically possible. These include:

DM relic density? How?

From particle physics to cosmology

The Boltzmann equation connects microphysical interactions (annihilation) rates) with cosmological dynamics (expansion of the Universe).

$$\frac{dn_{\chi}}{dt} = -3Hn_{\chi} - \langle \sigma_{eff} v \rangle [n_{\chi}^2 - (n_{\chi}^{eq})^2]$$
Dilution of the number
density due to the
expansion of the
Universe
$$\int e^{-\frac{E}{K_B T}} d^3v$$

The product ov describes the rate of particle annihilation per unit pair of particles.

it ensures that all possible velocities and energy distributions in the thermal bath are considered.

 $\frac{E}{K_B T} d^3 v$

The effective annihilation cross-section for a system with multiple processes is a sum over all possible final states: $H^0H^0 \rightarrow W^+W^-$

$$\sigma_{eff}v = \sum_{i} \sigma_{i}v_{rel;i}$$

$$\begin{split} H^{0}H^{0} &\to W^{+}W^{-}, ZZ, hh, f\bar{f}, \dots \\ A^{0}A^{0} &\to W^{+}W^{-}, ZZ, hh, f\bar{f}, \dots \\ H^{+}H^{-} &\to W^{+}W^{-}, ZZ, hh, f\bar{f}, \dots \end{split}$$

► For co-annihilation, the effective σv becomes $H^0A^0 \rightarrow Zh$, $H^0H^+ \rightarrow W^+h$, $A^0H^+ \rightarrow W^+Z$,...

$$\langle \sigma v \rangle_{eff} = \sum_{ij} \langle \sigma_{ij} v \rangle \frac{n_i n_j}{(n_{total})^2}$$

Approximation for Non-Relativistic Dark Matter

$$\sigma v = a + bv^2 + O(v^4)$$

$$\sigma \sim \sigma_0 + \sigma_2 v^2$$

 $\sigma \sim \sigma_0$ cold dark matter

- To quantify these processes, we numerically solve the Boltzmann equation, accounting for all possible annihilation and co-annihilation contributions.
- The relic density of DM is required to align with the PLANCK measurement

$$\Omega_{DM}h^2 = 0.1200 \pm 0.0012$$

Aghanim, N., et al. "Erratum: Planck 2018 results: VI. Cosmological parameters (Astronomy and Astrophysics (2020) 641 (A6." Astronomy & Astrophysics 652 (2021): 1-3.

Problems of the Inert Doublet Model:

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- 1. Ensuring correct dark matter (DM) relic density requires tuning parameters like the mass of the DM candidate and scalar couplings.
- 2. Indirect detection faces issues in probing high-mass regions due to limited experimental sensitivity.
- 3. The model's ability to achieve a strong first-order **EWPT for** baryogenesis is sensitive to parameter choices.
- 4. The IDM does not inherently explain the smallness of **neutrino masses** or their observed mixing patterns.
- 5. The model lacks a built-in mechanism for explaining the baryon asymmetry of the universe.

Extensions of the Inert Doublet Model:

- Adding singlet or triplet scalars or fermions can address neutrino masses (via seesaw mechanisms) and baryogenesis.
- Adding right-handed neutrinos allows for neutrino mass generation via the seesaw mechanism and provides new interactions with the inert doublet.
- Enlarging the gauge group (e.g., adding U(1) symmetry) can stabilize dark matter and introduce new interactions that modify IDM phenomenology.
- Introducing additional discrete symmetries, to stabilize multi-component DM or allow for new decay channels.
- Embedding IDM within a supersymmetric framework (e.g., SUSY) can address finetuning and provide additional dark matter candidates.
- Adding new gauge bosons or interactions to the model to enhance DM annihilation channels or address indirect detection signals.
- Studying loop corrections to the scalar potential can provide more precise constraints and potentially resolve vacuum stability issues.
- Incorporating IDM into a framework with modifications to gravity, such as extra dimensions, may provide connections to cosmology.