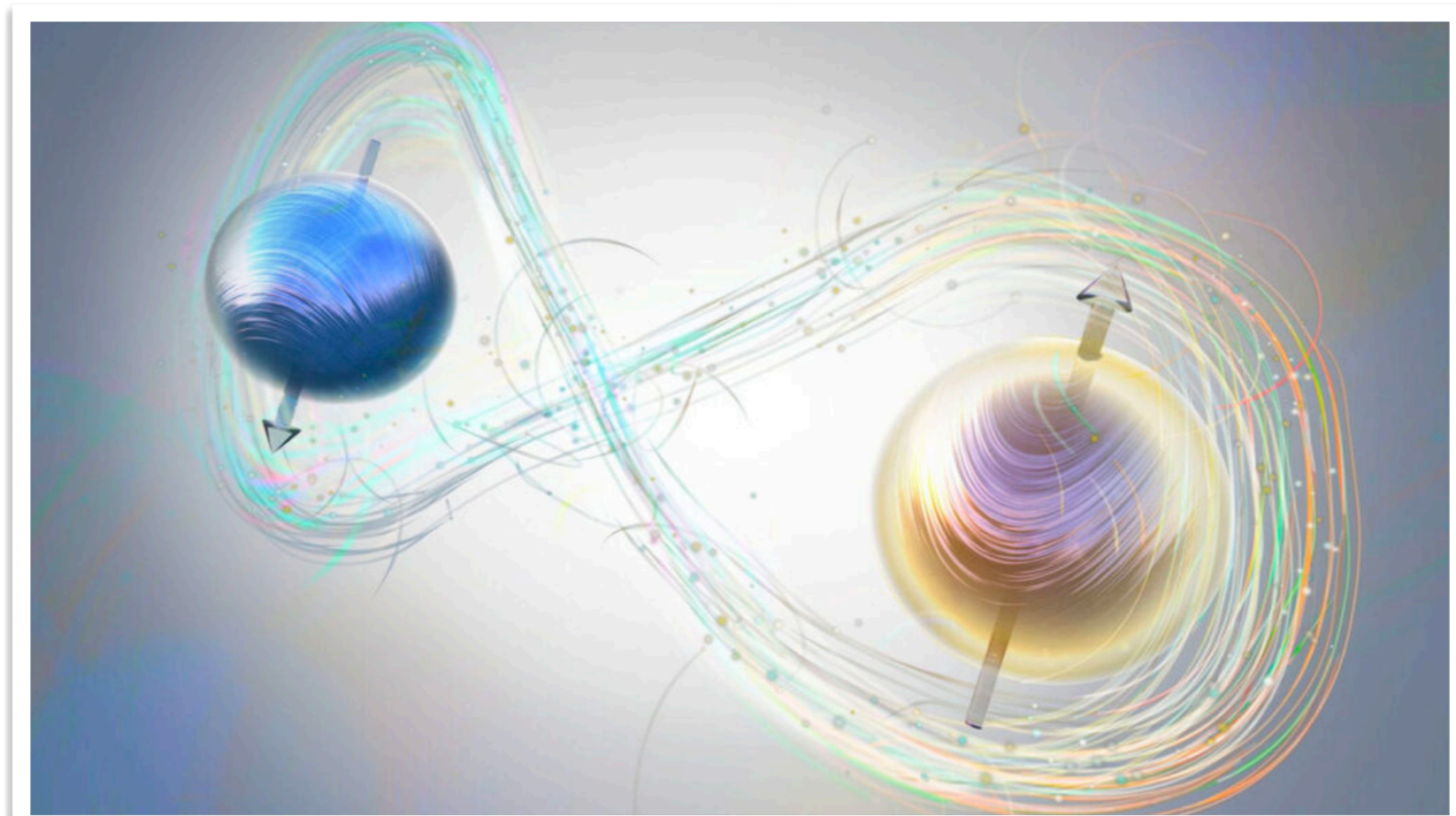


Hadronic Top Quark Polarimetry

Top 2024, Saint-Malo, Sep 25, 2024

Dorival Gonçalves 



Dong, DG, Kong, Navarro '23

Dong, DG, Kong, Larkoski, Navarro '24

Dong, DG, Kong, Larkoski, Navarro '24

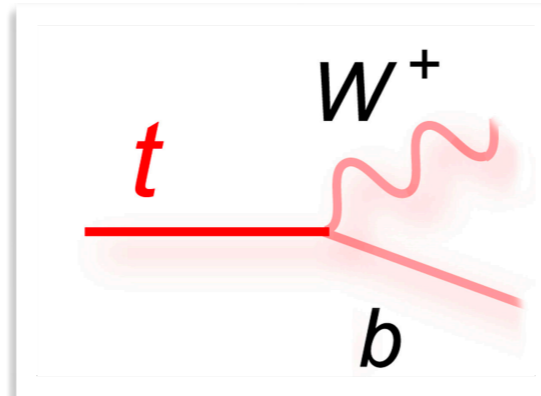
Top Quark is Unique

- Decays before it hadronizes or its spin flips

$$\tau_{top} \approx 5 \times 10^{-25} s$$

$$\tau_{had} \sim 1/\Lambda_{QCD} \sim 10^{-24} s$$

$$\tau_{flip} \sim m_t/\Lambda_{QCD}^2 \sim 10^{-21} s$$

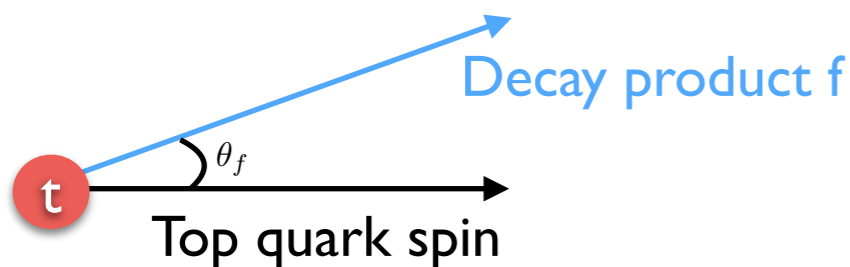


- Top polarization directly observable via angular distributions of its decay products

$$\frac{1}{\Gamma} \frac{d\Gamma}{d \cos \theta_f} = \frac{1}{2} (1 + \beta_f \cos \theta_f)$$

	l^+, \bar{d}	b	$\bar{\nu}, u$
β_f	1	-0.4	-0.3

Spin analyzing power: maximum for charged leptons



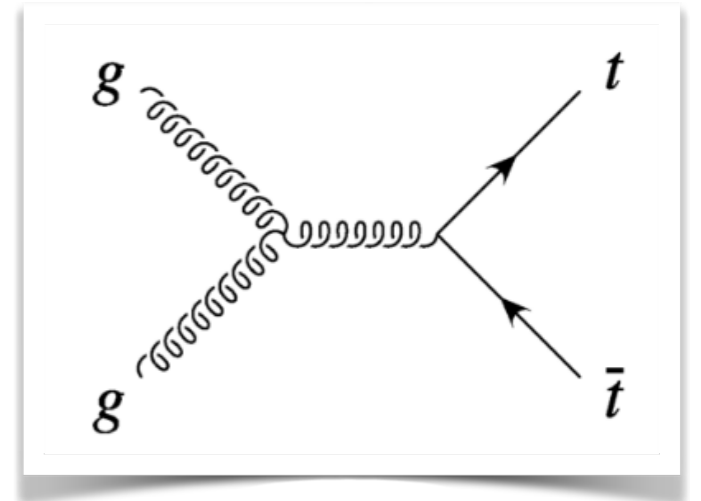
$$\beta_{l^+}^{\text{NLO}} = 0.998, \beta_{\bar{d}}^{\text{NLO}} = 0.966, \beta_b^{\text{NLO}} = -0.393$$

Brandenburg, Si, Uwer (2002)

Top quark pair production as a two qubit system

- The most general two-qubit system can be represented by

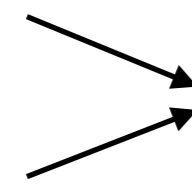
$$\rho = \frac{\mathbb{I} \otimes \mathbb{I} + (B_i \sigma_i \otimes \mathbb{I} + \bar{B}_i \mathbb{I} \otimes \sigma_i) + C_{ij} \sigma_i \otimes \sigma_j}{4}$$



- Characterized by 15 parameters: B_i , \bar{B}_i , and C_{ij}

$$B_i = \langle \sigma_i \otimes \mathbb{I} \rangle$$

$$\bar{B}_i = \langle \mathbb{I} \otimes \sigma_i \rangle$$


Polarizations

$$C_{ij} = \langle \sigma_i \otimes \sigma_j \rangle \longrightarrow \text{Spin correlations}$$

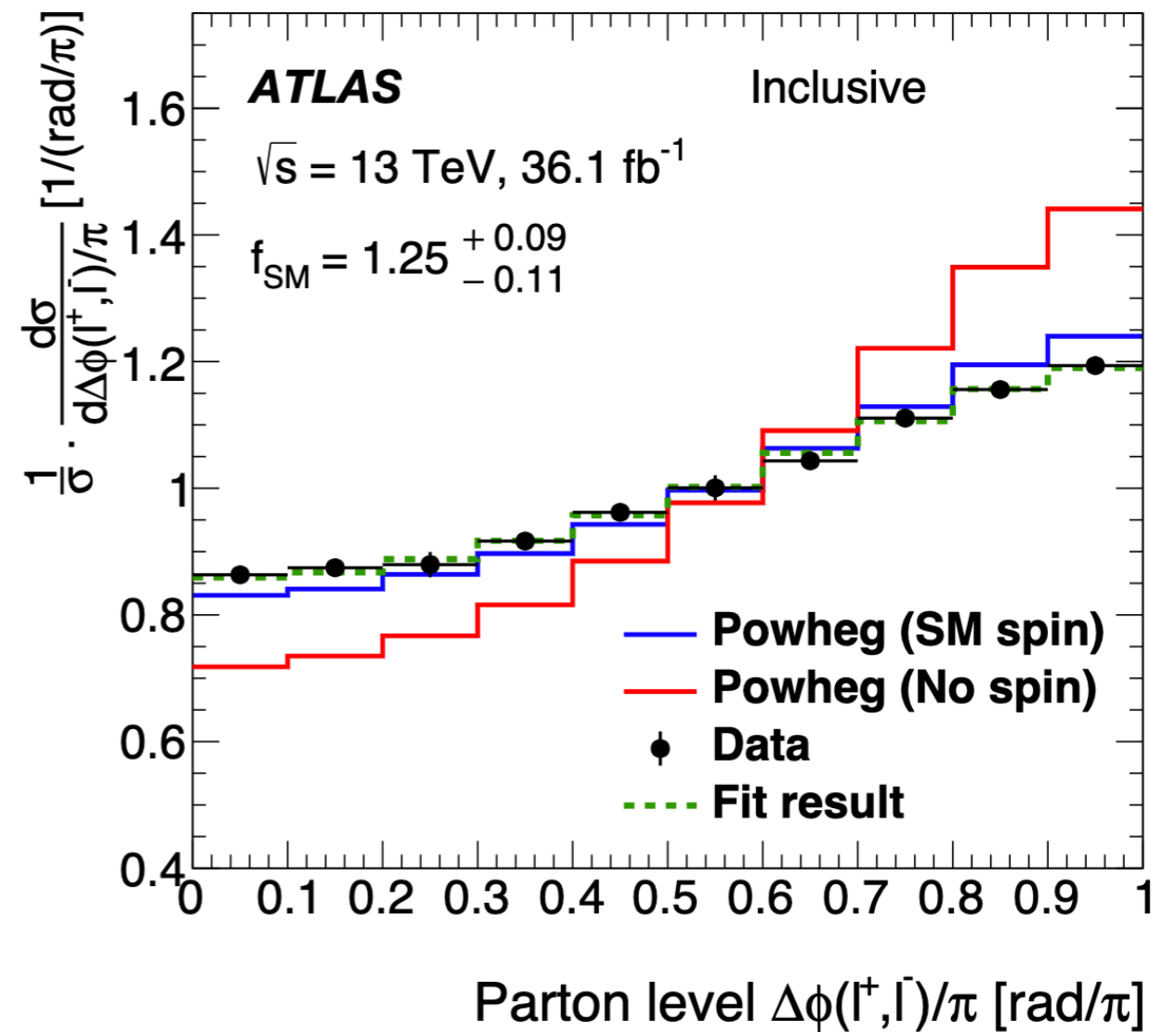
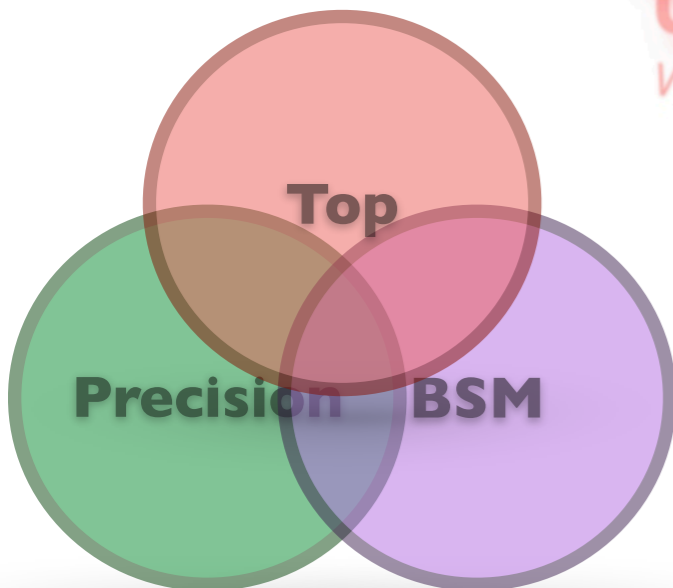
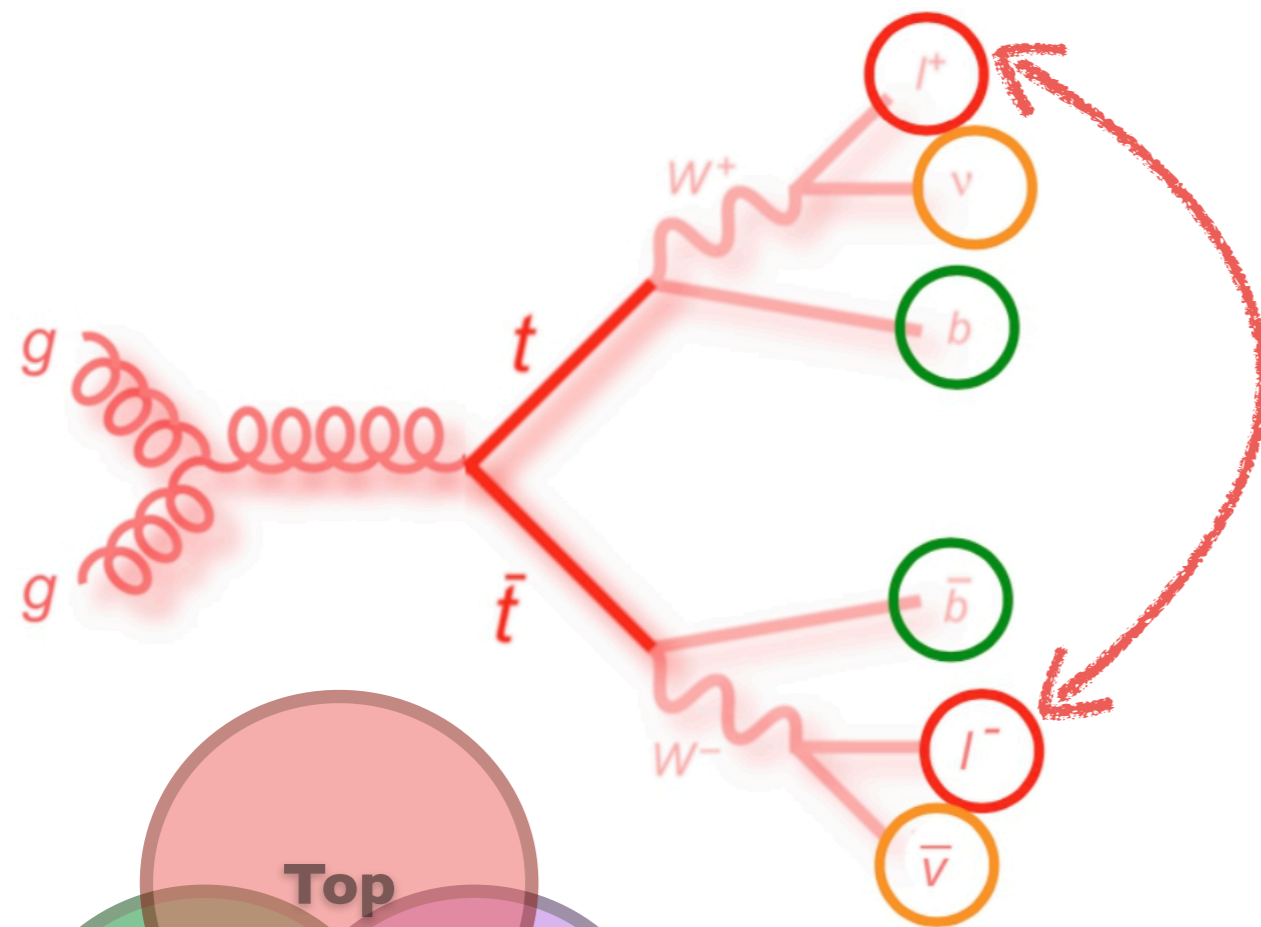
- P and CP invariance under $t\bar{t}$ production $\rightarrow B_i = \bar{B}_i = 0$ and $C_{ij} = C_{ji}$

- Things further simplify in the helicity basis: only non-vanishing parameters are the diagonal terms C_{ii} and one off-diagonal term $C_{12} \simeq C_{21}$

Bernreuther, Heisler, Si '15
 Czakon, Mitov, Poncelet '20
 Frederix, Tsinikos, Vitis '21

Top quark pair production as a two qubit system

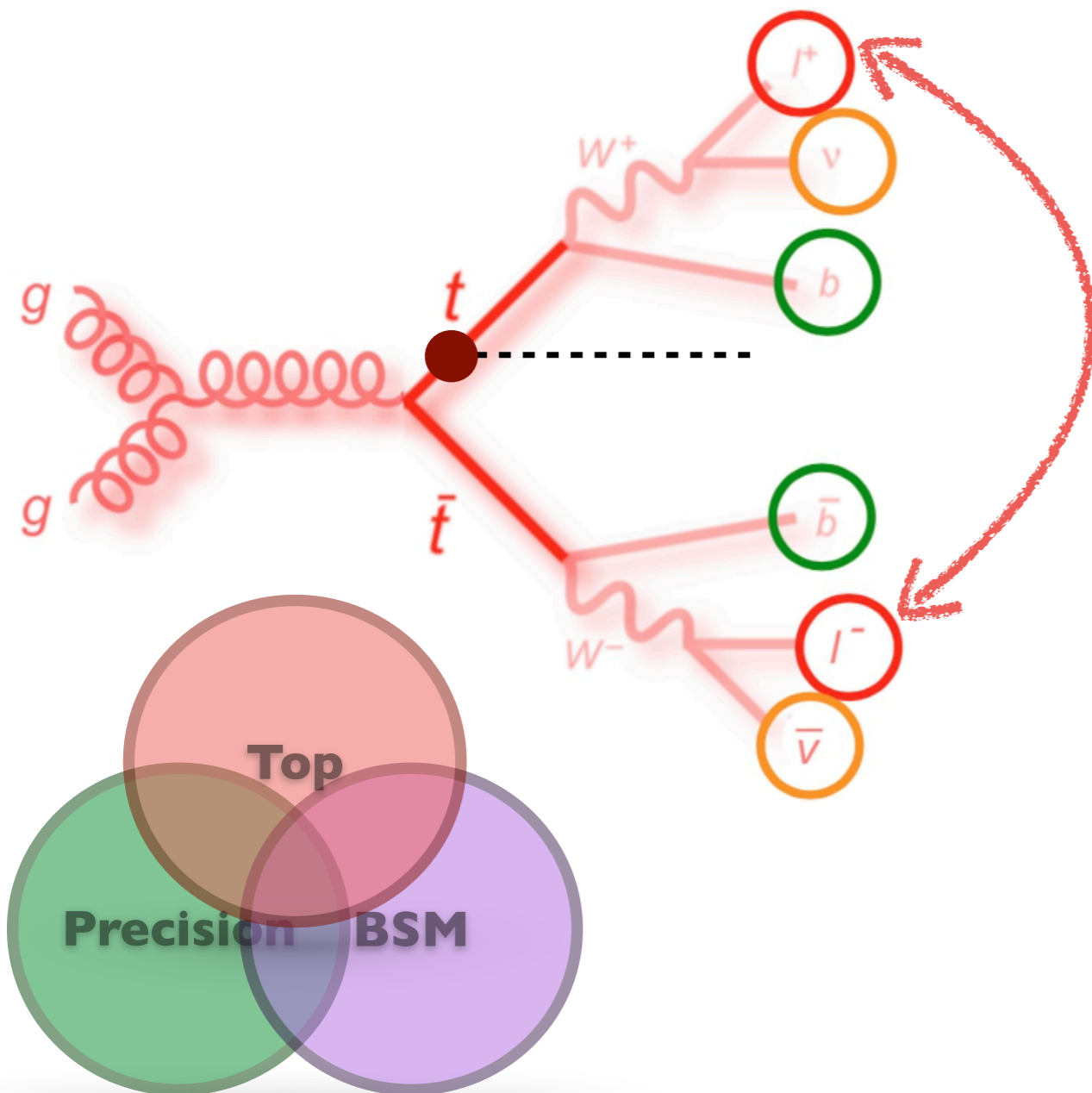
- We are probing the spin correlations of top and anti-top since the Tevatron era and continue to study them at the LHC in different forms: relevant precision observable



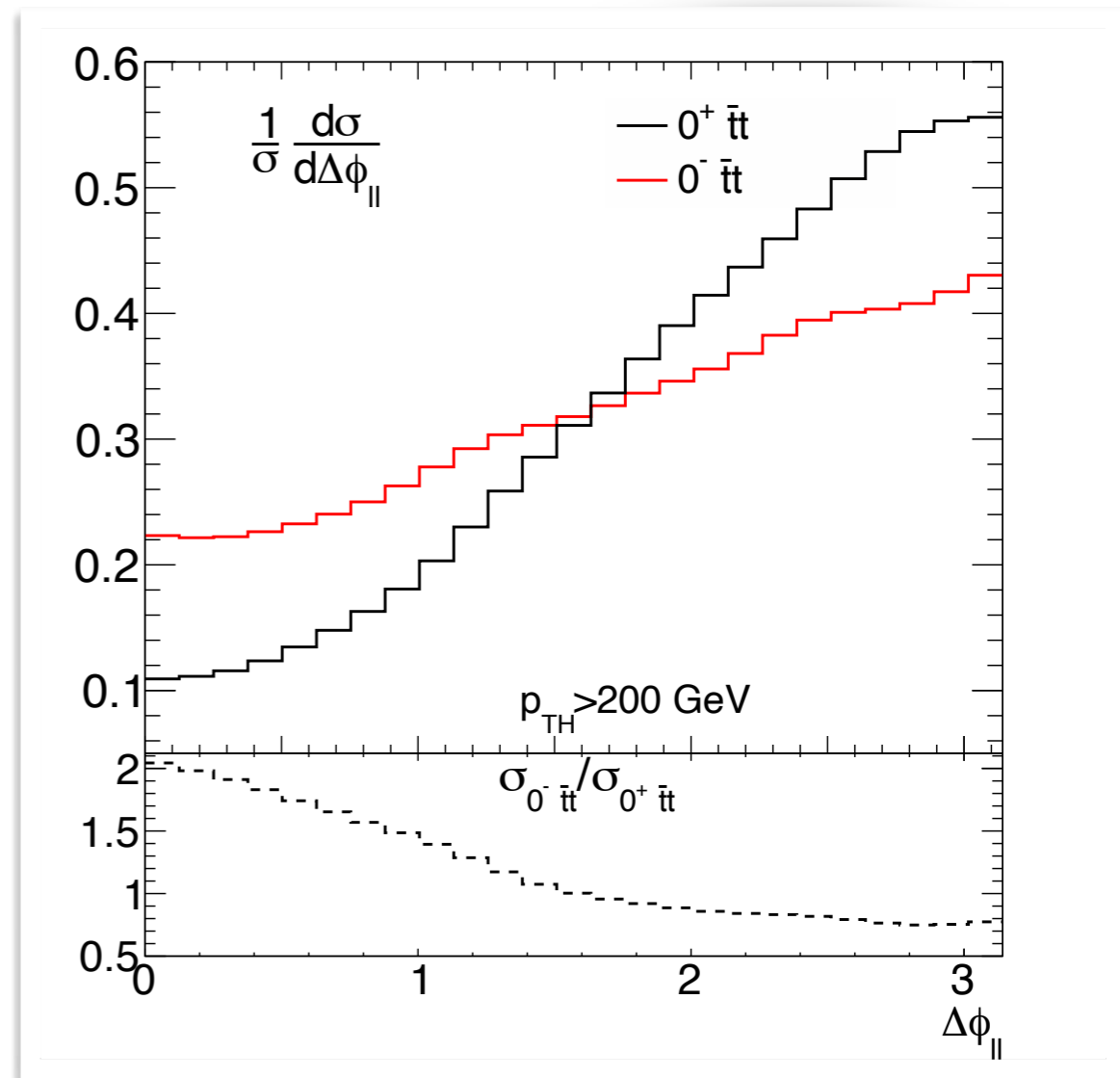
Parke, Mahlon '95

Top quark pair production as a two qubit system

- We are probing the spin correlations of top and anti-top since the Tevatron era and continue to study them at the LHC in different forms: relevant precision observable



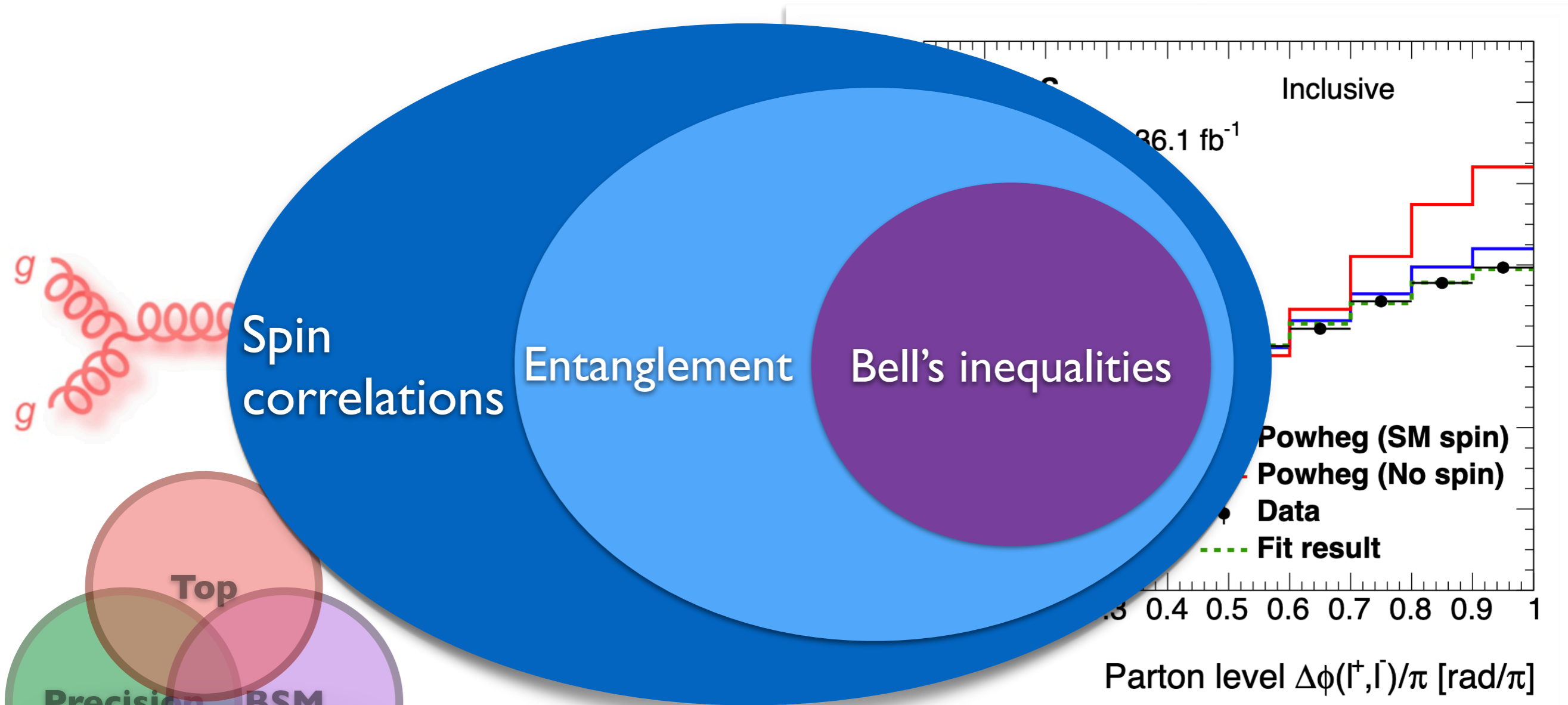
$$\mathcal{L} \supseteq -\frac{m_t}{v} K \bar{t} (\cos \alpha + i \gamma_5 \sin \alpha) t H$$



Buckley, DG (PRL '15)

Top quark pair production as a two qubit system

- We are probing the spin correlations of top and anti-top since the Tevatron era and continue to study them at the LHC in different forms: relevant precision observable

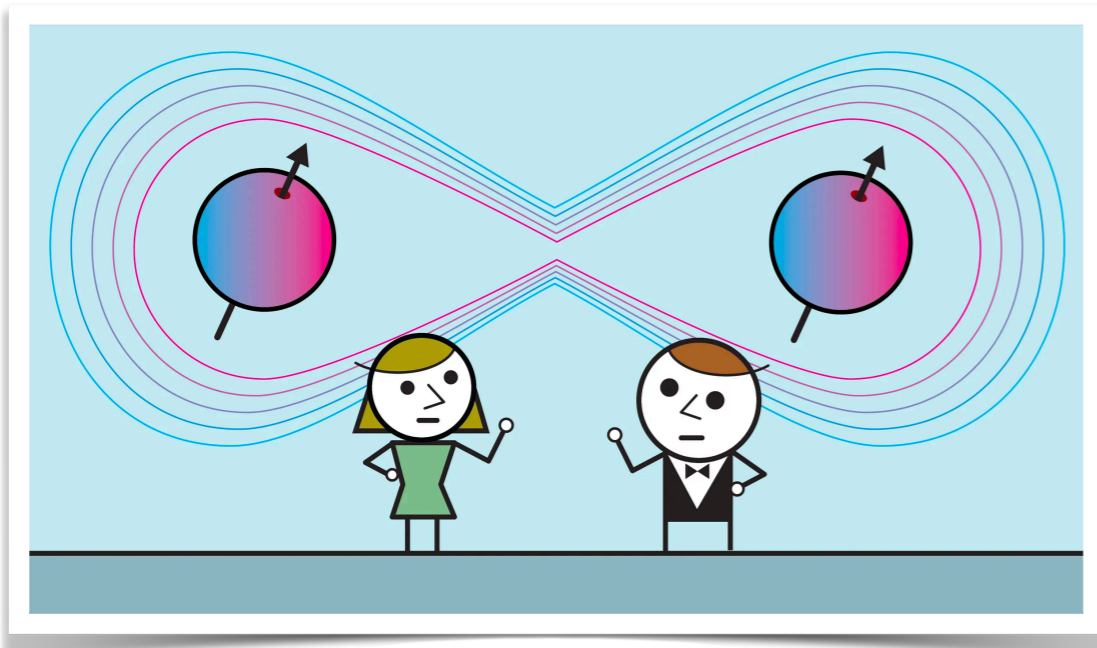


$$\rho = \frac{\mathbb{I} \otimes \mathbb{I} + (B_i \sigma_i \otimes \mathbb{I} + \bar{B}_i \mathbb{I} \otimes \sigma_i) + C_{ij} \sigma_i \otimes \sigma_j}{4}$$

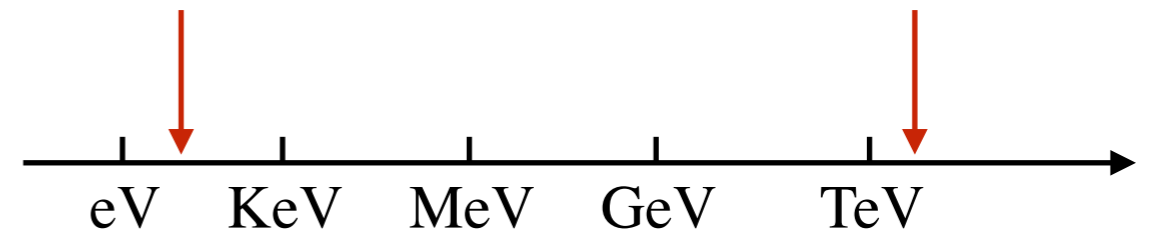
Parke, Mahlon '95

Entanglement and Bell Inequalities with Top Quarks

- LHC can provide a unique environment to study entanglement and violation of Bell's inequalities at the highest energy available to date



Typical entanglement
experiment with photons



- Top quark pair production is an optimal candidate for these studies

Afik, Nova '20; Fabbrichesi, Floreanini, Panizzo '21; Severi, Boschi, Maltoni, Sioli '21

Saavedra, Casas '22; Severi, Vryonidou 22

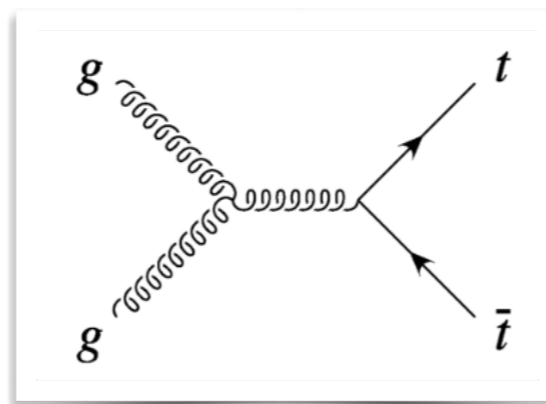
Dong, DG, Kong, Navarro '23

Han, Low, Wu '23

ATLAS Nature vol 633, 542-547 (2024)

CMS 2406.03976

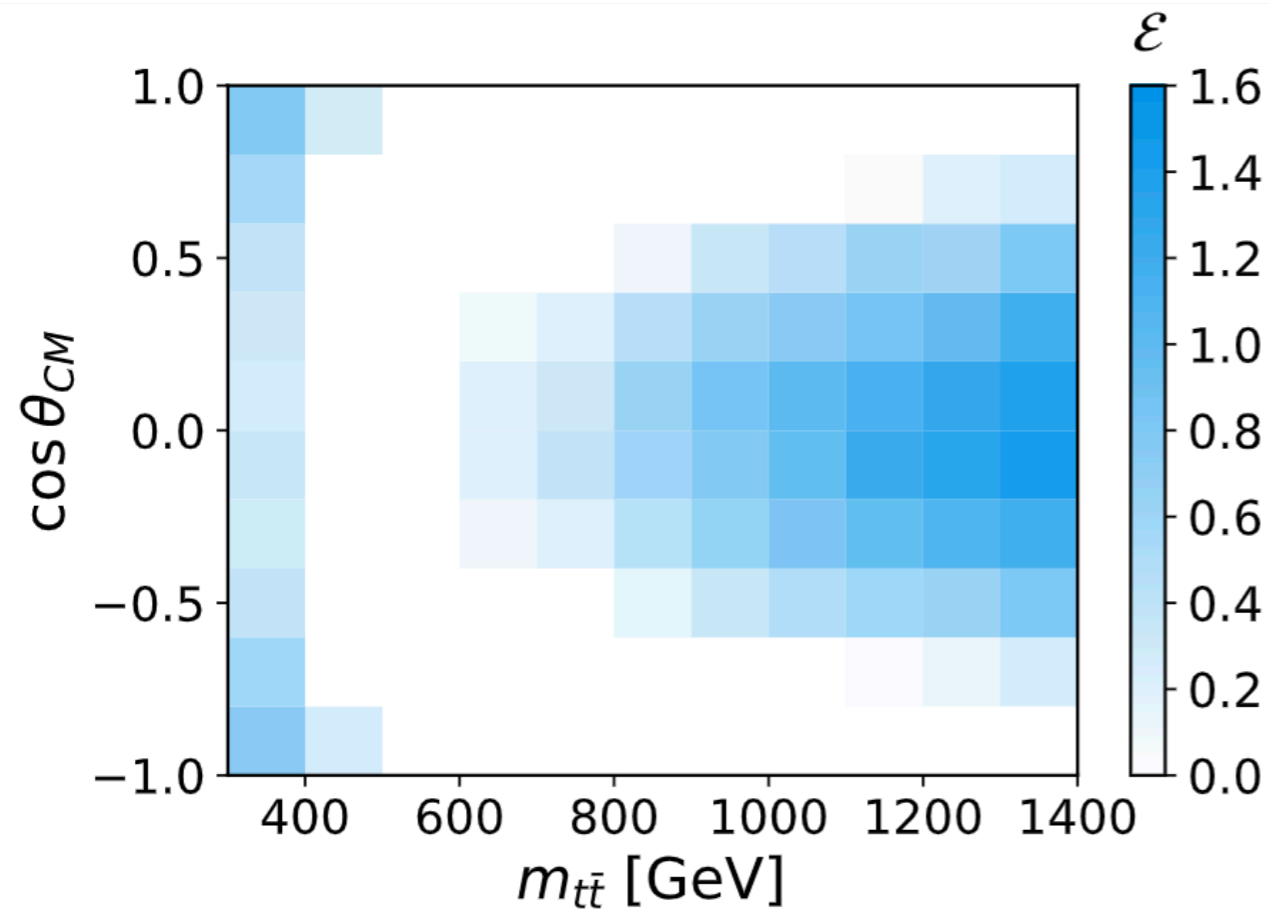
CMS 2409.11067



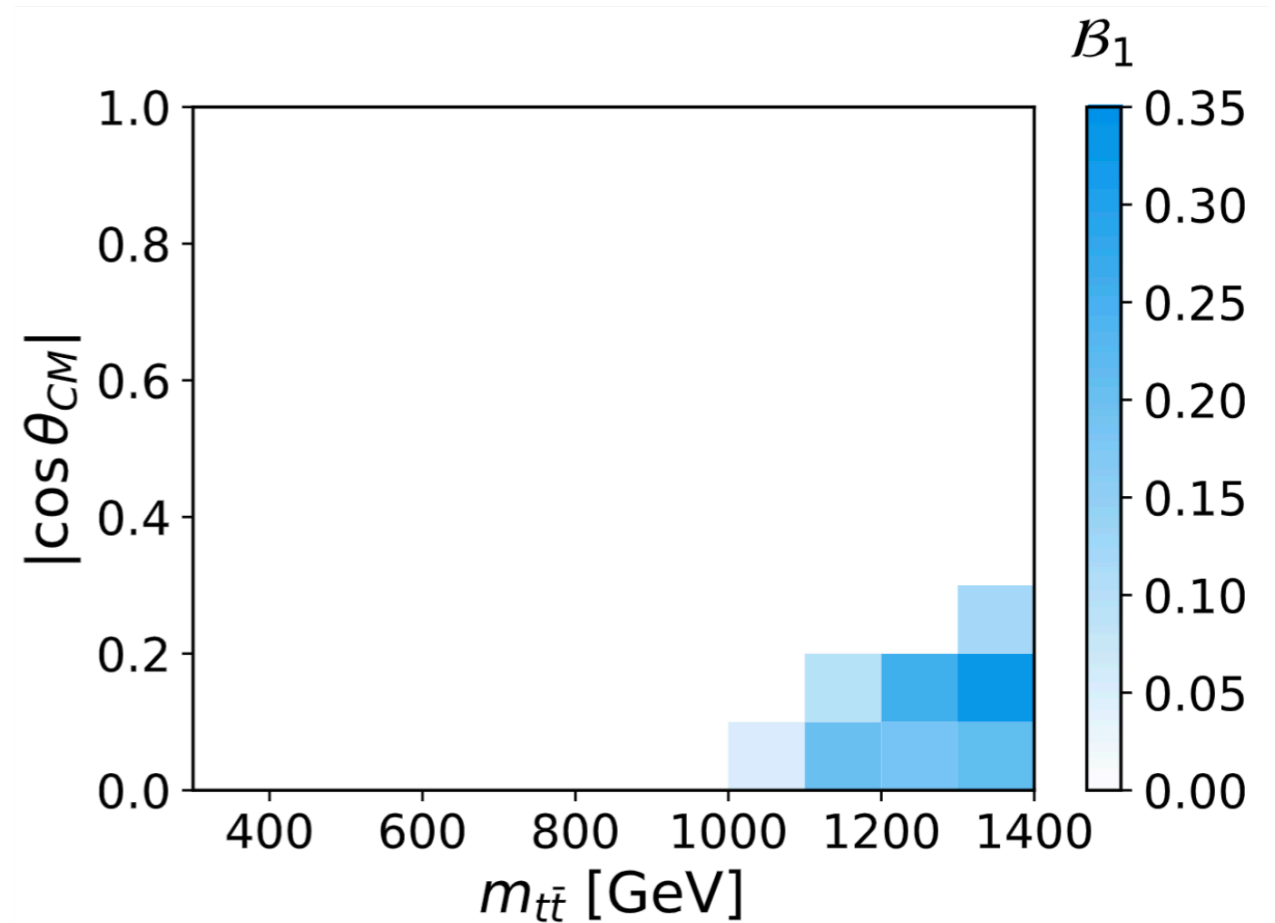
See talks by Eleni Vryonidou, Chris White,
Juan Antonio A. Saavedra, Giulia Negro,
Roman Lysak

Quantum Correlations

Entanglement



Bell/CHSH inequalities



Dong, DG, Kong, Navarro '23

Afik, Nova '20

Severi, Boschi, Maltoni, Sioli '21

Saavedra, Casas '22

➡ Entanglement and Bell/CHSH violation studies well match **boosted** top pair searches

Semi-leptonic top pair

Pros:

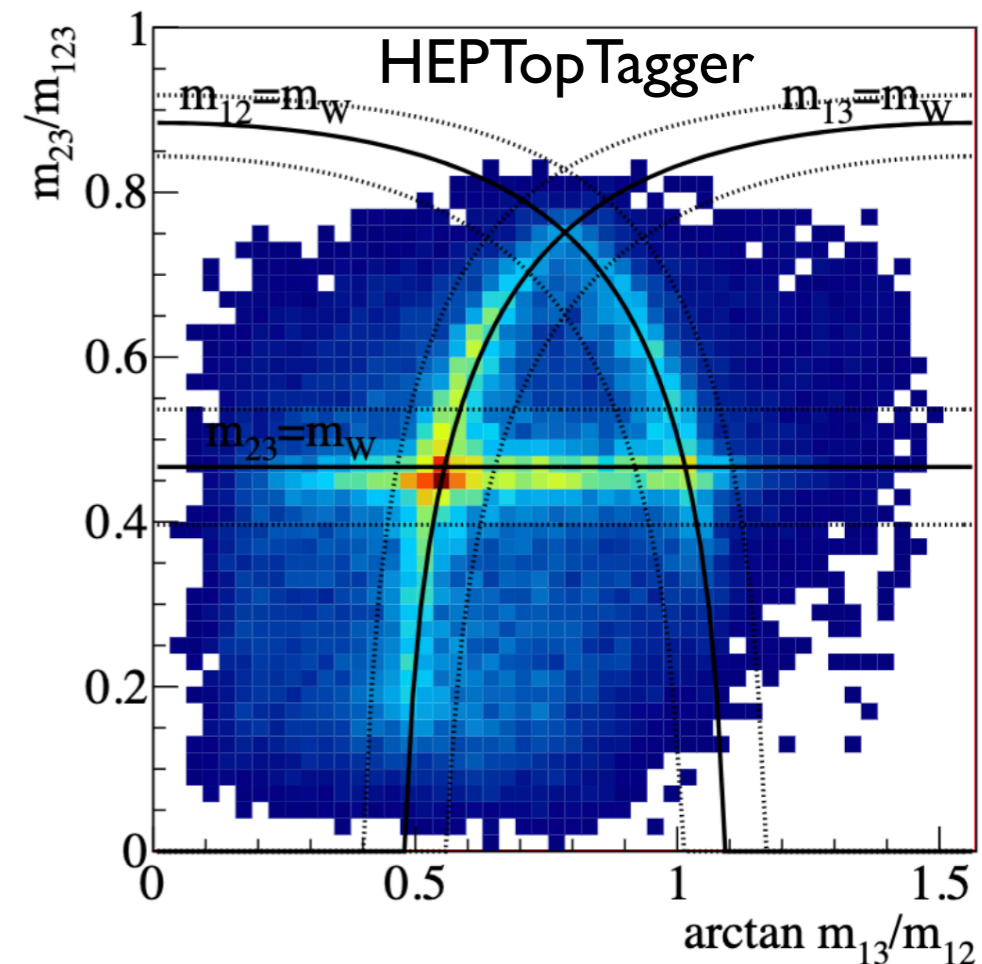
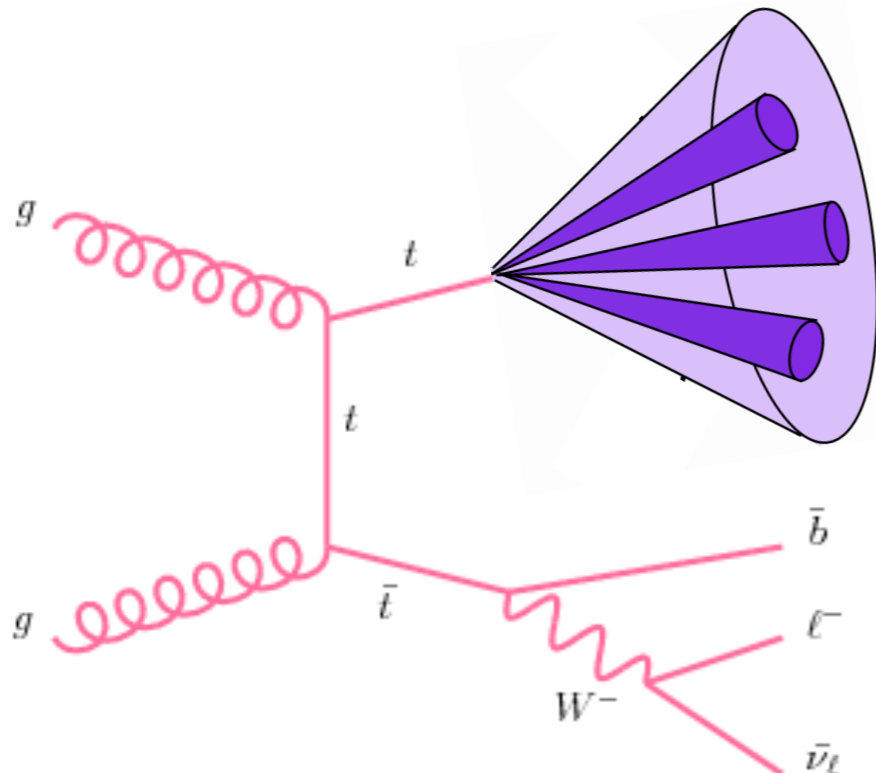
- 6 times higher event rate than dileptonic case
- It can more effectively probe the boosted regime
- Easier reconstruction

Cons:

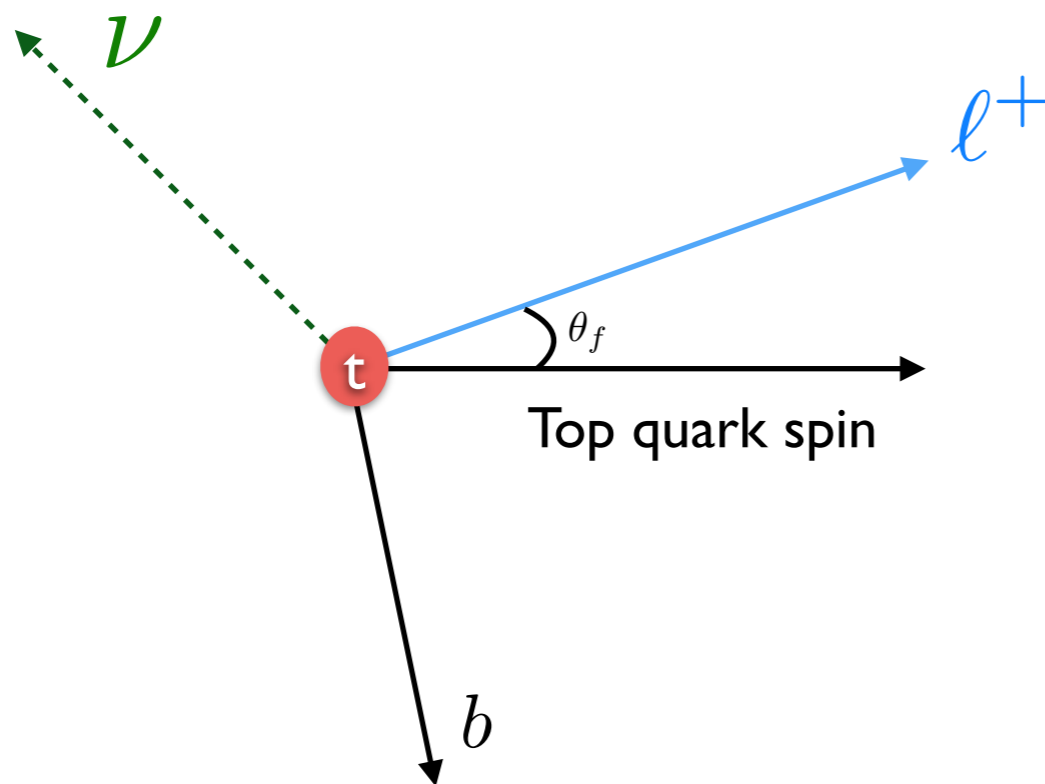
- down-type quark is best polarimeter, but tagging it in a collider environment is challenging

Solution:

- Boosted top tagging to aid subjet and light quark matching
- Proxy for down-quark: optimal hadronic polarimeter

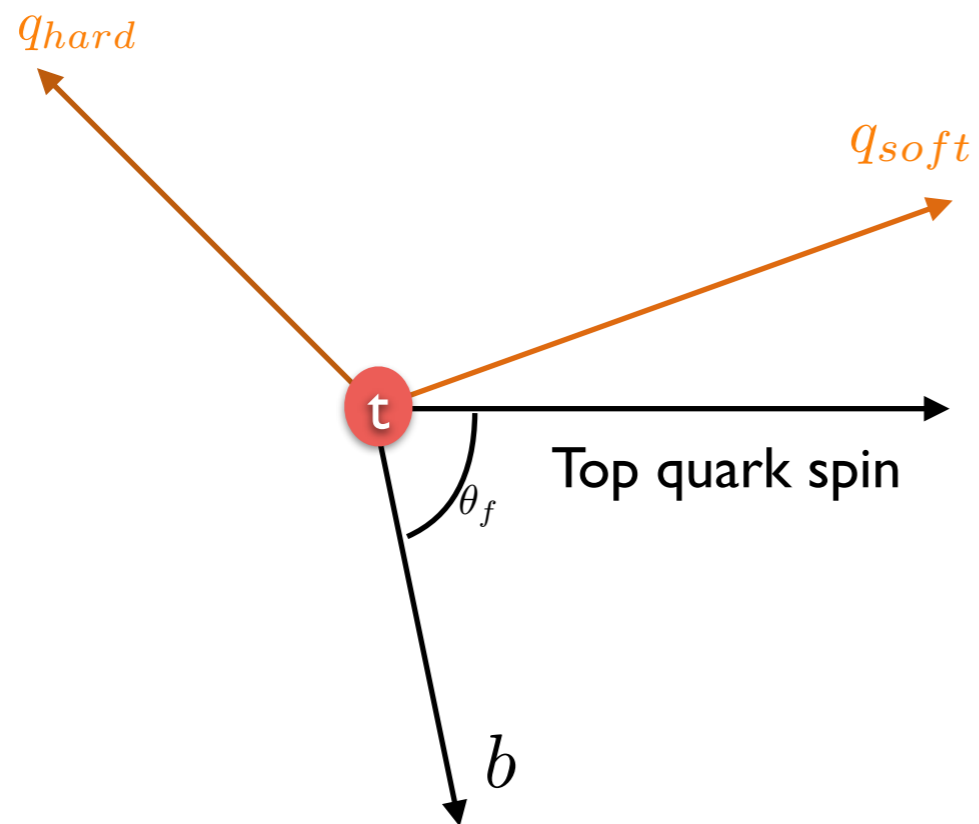


Hadronic top quark polarimetry



$$\frac{1}{\Gamma_f} \frac{d\Gamma_f}{d \cos \theta_f} = \frac{1}{2} (1 + \mathbf{1.0} \cos \theta_f)$$

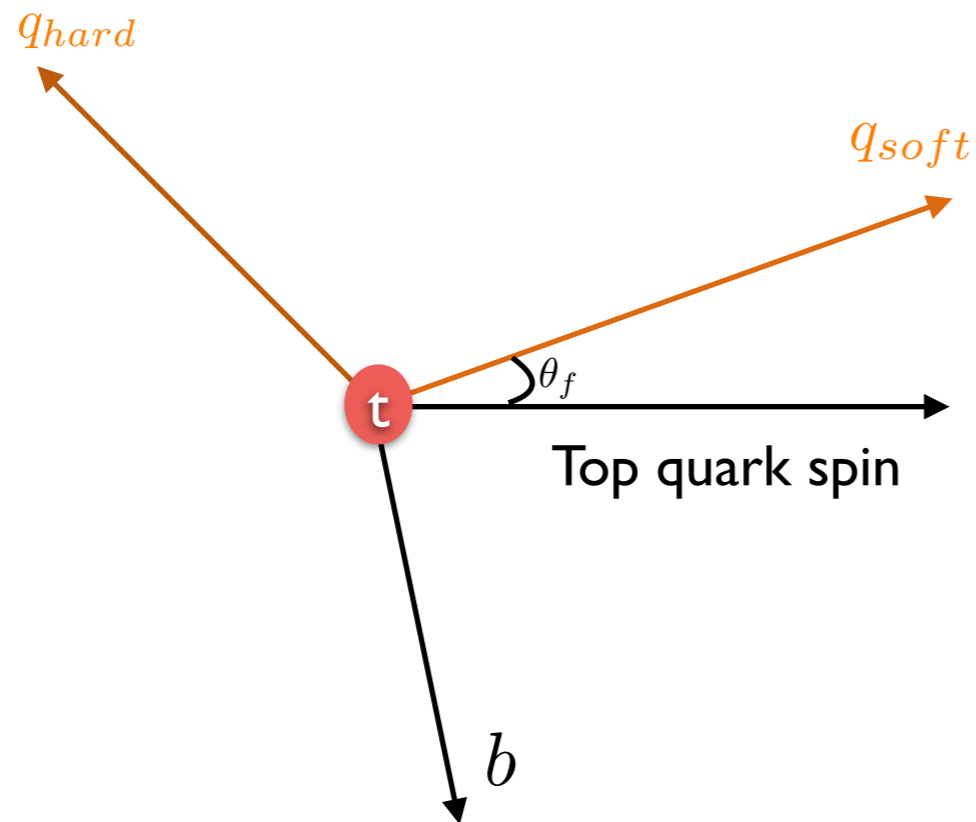
Hadronic top quark polarimetry



$$\frac{1}{\Gamma_f} \frac{d\Gamma_f}{d \cos \theta_f} = \frac{1}{2} (1 - 0.4 \cos \theta_f)$$

Hadronic top quark polarimetry

- In the high-energy regime, boosted techniques can be employed to tag the hadronic top and efficiently identify a spin analyzer

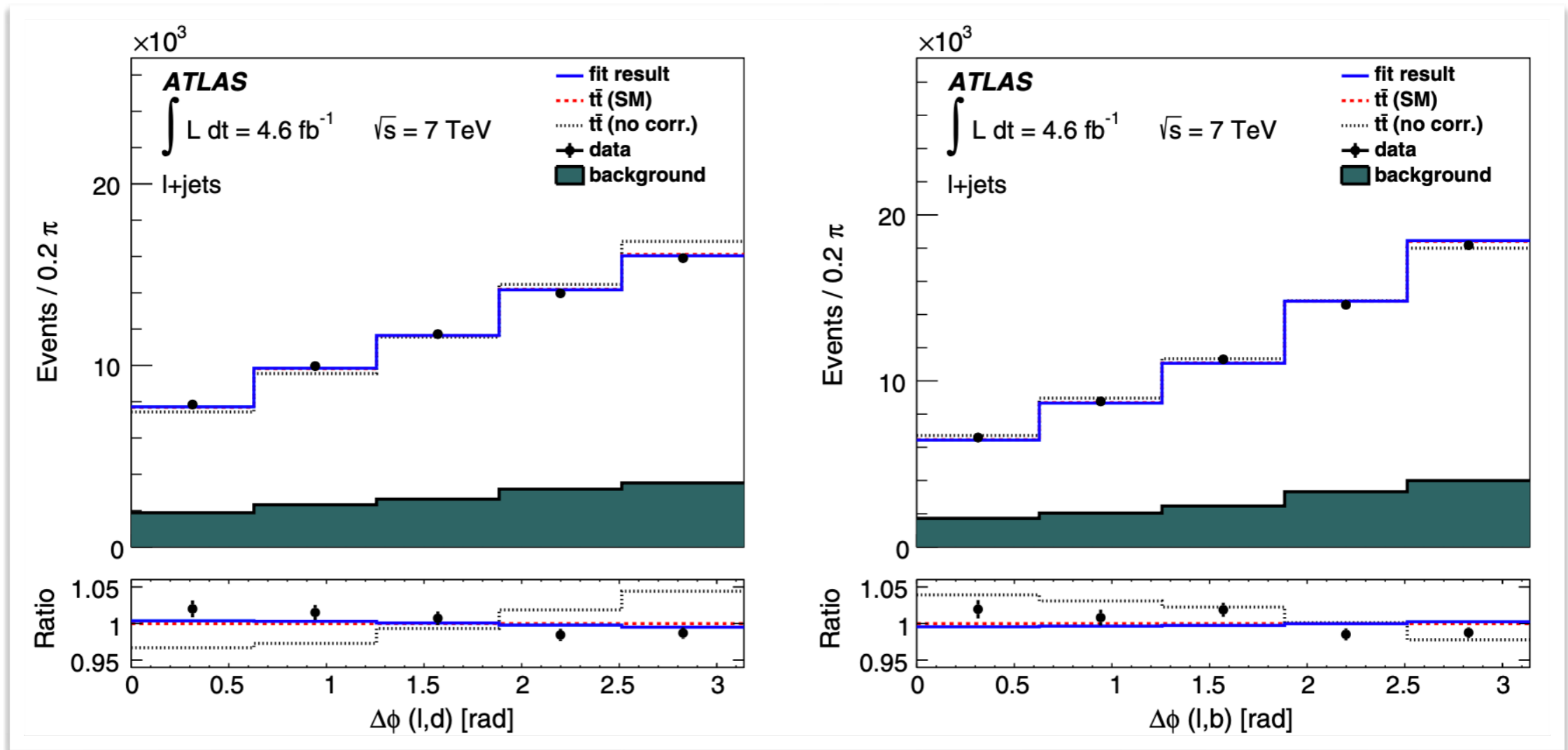


$$\frac{1}{\Gamma_f} \frac{d\Gamma_f}{d \cos \theta_f} = \frac{1}{2} (1 + \mathbf{0.5} \cos \theta_f)$$

Jezabek '94

Hadronic top quark polarimetry

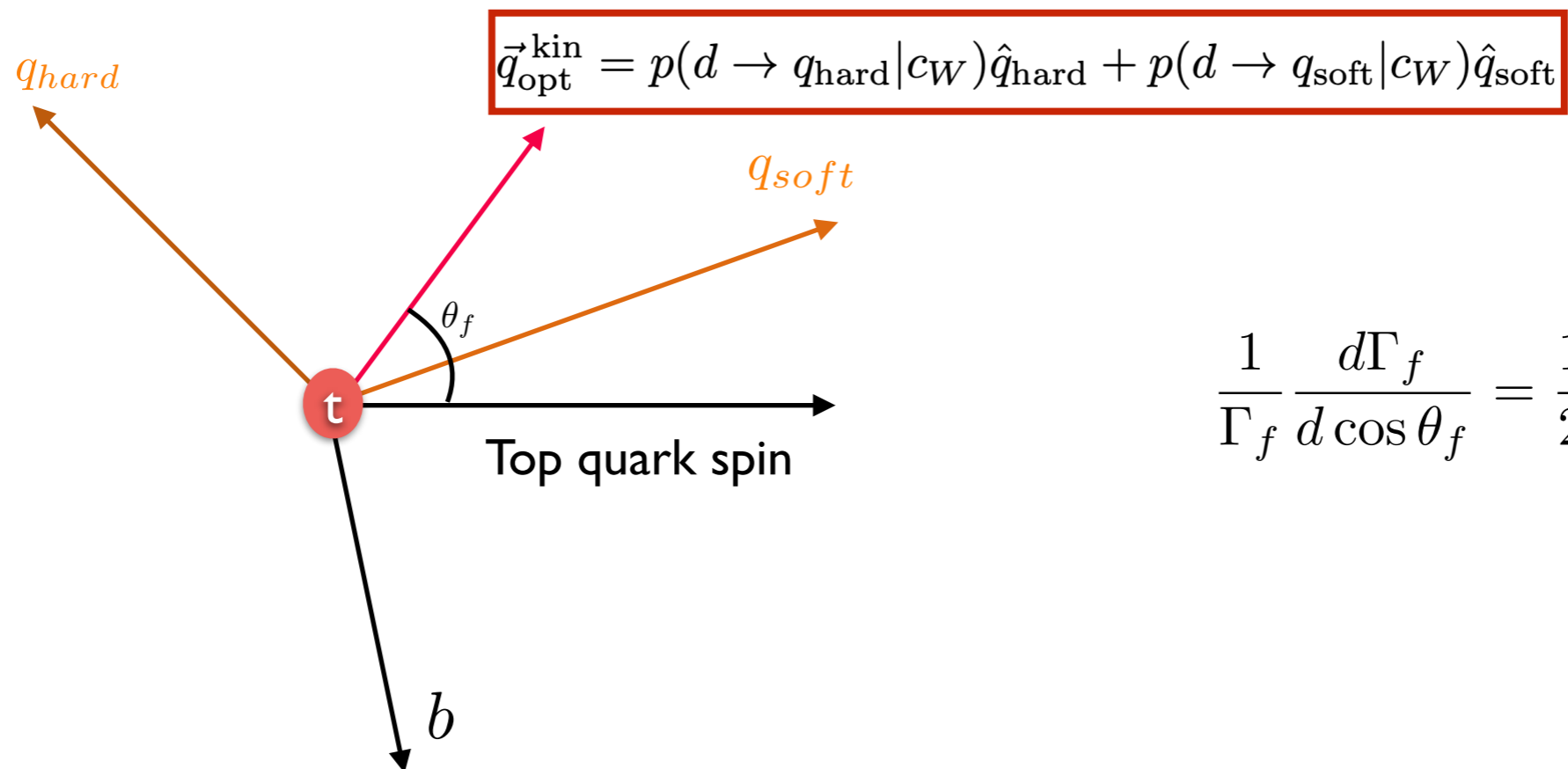
In the high-energy regime, boosted techniques can be employed to tag the hadronic top and efficiently identify a spin analyzer



PHYSICAL REVIEW D 90, 112016 (2014)

Hadronic top quark polarimetry

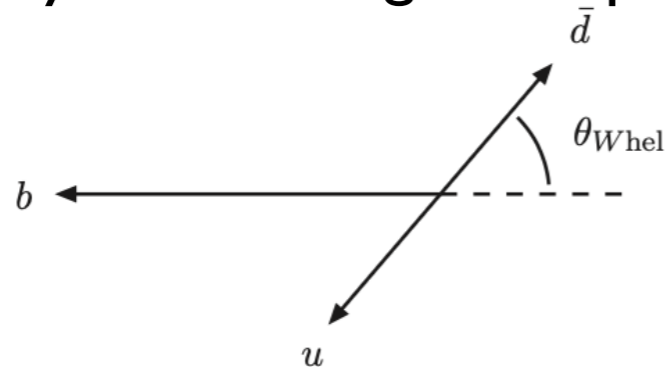
- In the high-energy regime, boosted techniques can be employed to tag the hadronic top and efficiently identify a spin analyzer



$$\frac{1}{\Gamma_f} \frac{d\Gamma_f}{d \cos \theta_f} = \frac{1}{2} (1 + 0.64 \cos \theta_f)$$

Hadronic top quark polarimetry

- We begin by considering the top quark decay in the W boson rest frame



W^\pm is given by

$$\rho(c_{W_{\text{hel}}}) = \frac{3}{8} f_R (1 \pm c_{W_{\text{hel}}})^2 + \frac{3}{4} f_0 (1 - c_{W_{\text{hel}}}^2) + \frac{3}{8} f_L (1 \mp c_{W_{\text{hel}}})^2.$$

- If we do not distinguish the d-subjet and u-subjet, we need to identify $c_{W_{\text{hel}}} \leftrightarrow -c_{W_{\text{hel}}}$
- However,

$$p(d \rightarrow q_{\text{hard}}) = \frac{\rho(|c_{W_{\text{hel}}}|)}{\rho(|c_{W_{\text{hel}}}|) + \rho(-|c_{W_{\text{hel}}}|)}$$

$$p(d \rightarrow q_{\text{soft}}) = \frac{\rho(-|c_{W_{\text{hel}}}|)}{\rho(|c_{W_{\text{hel}}}|) + \rho(-|c_{W_{\text{hel}}}|)}$$

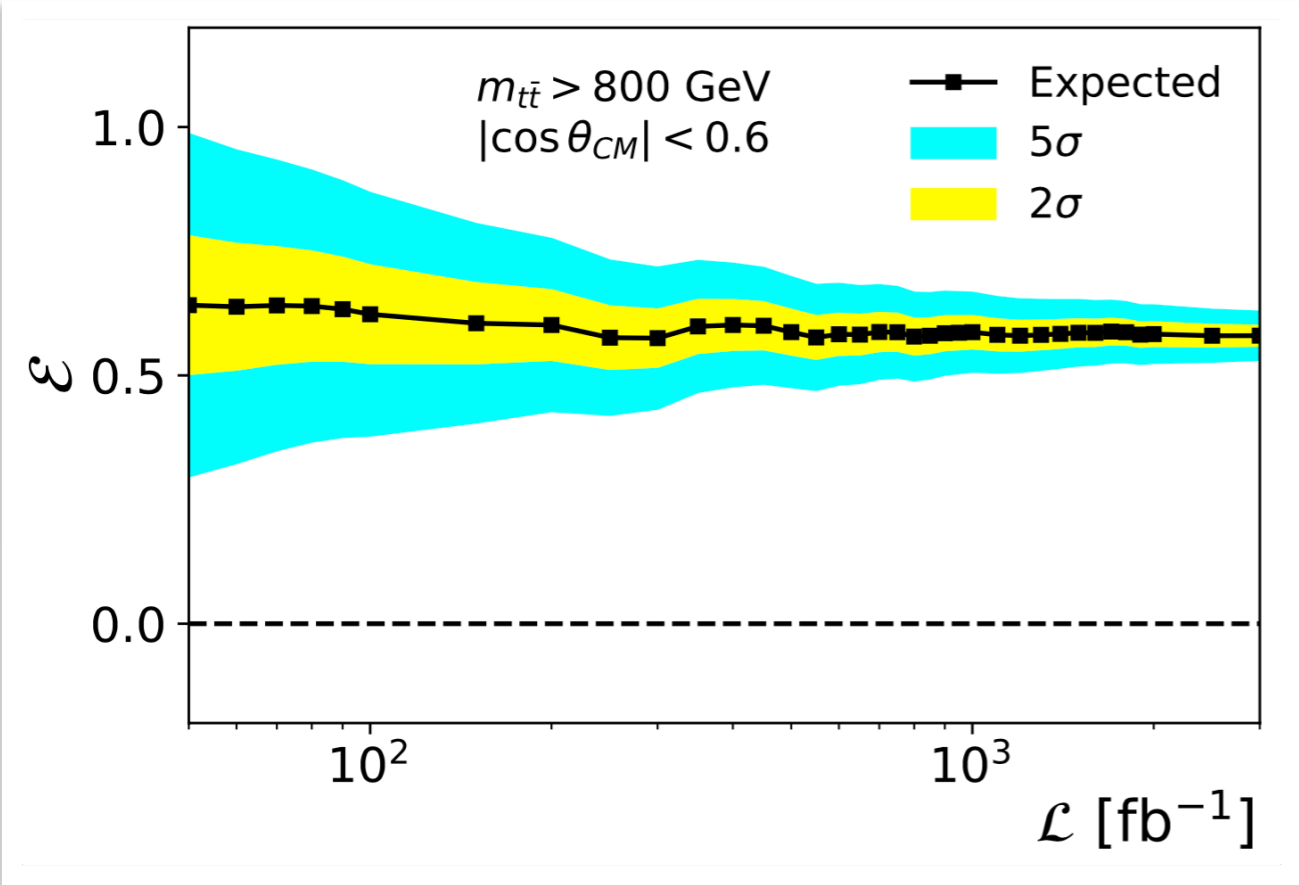
→ quark emitted in forward direction in W rest frame will be harder and more separated from b-quark in top rest frame

→ quark emitted in backward direction in W rest frame will be softer and more aligned with b-quark in top rest frame

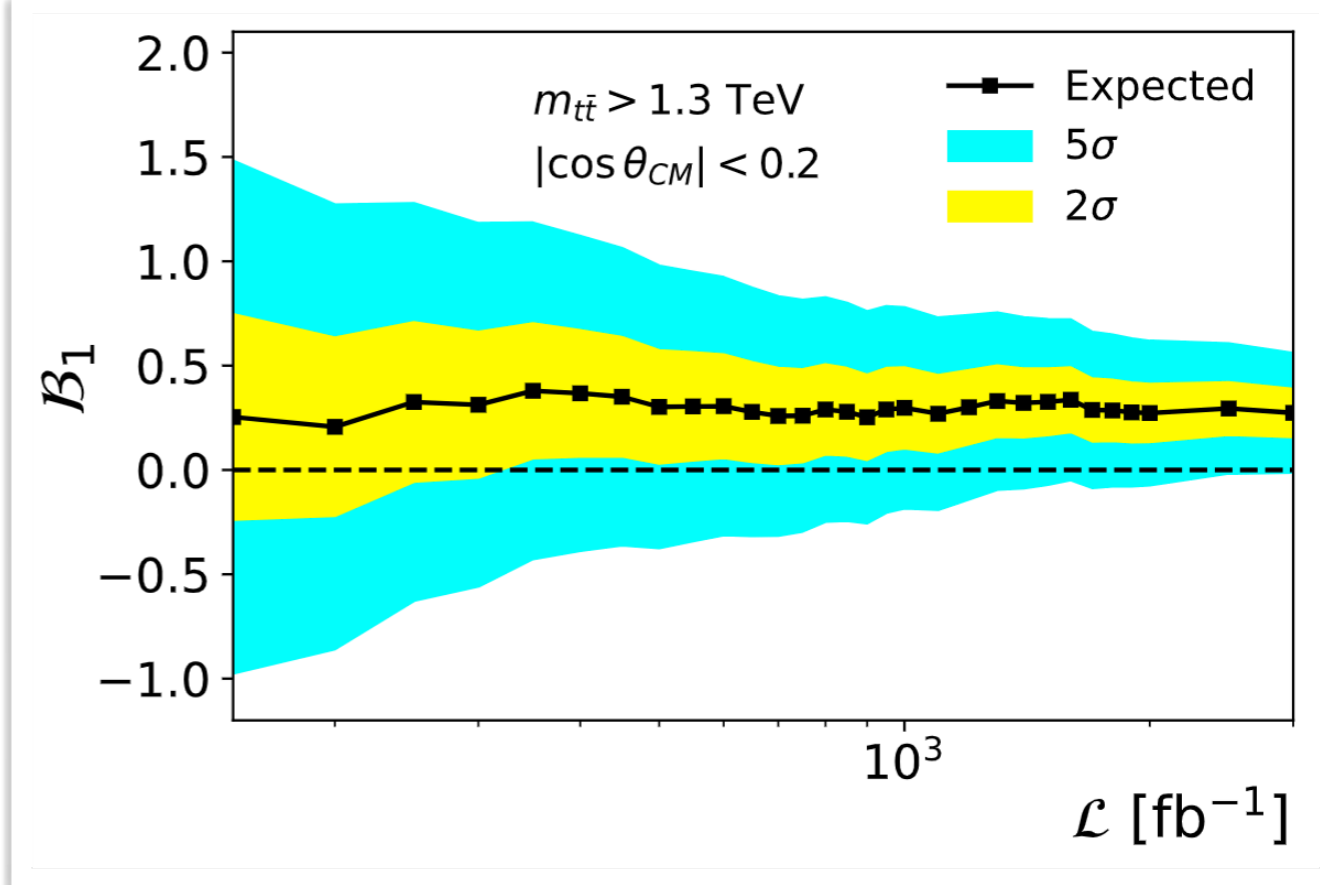
$$\vec{q}_{\text{opt}}^{\text{kin}} = p(d \rightarrow q_{\text{hard}} | c_W) \hat{q}_{\text{hard}} + p(d \rightarrow q_{\text{soft}} | c_W) \hat{q}_{\text{soft}}$$

LHC Projections

Entanglement



Bell/CHSH inequalities:



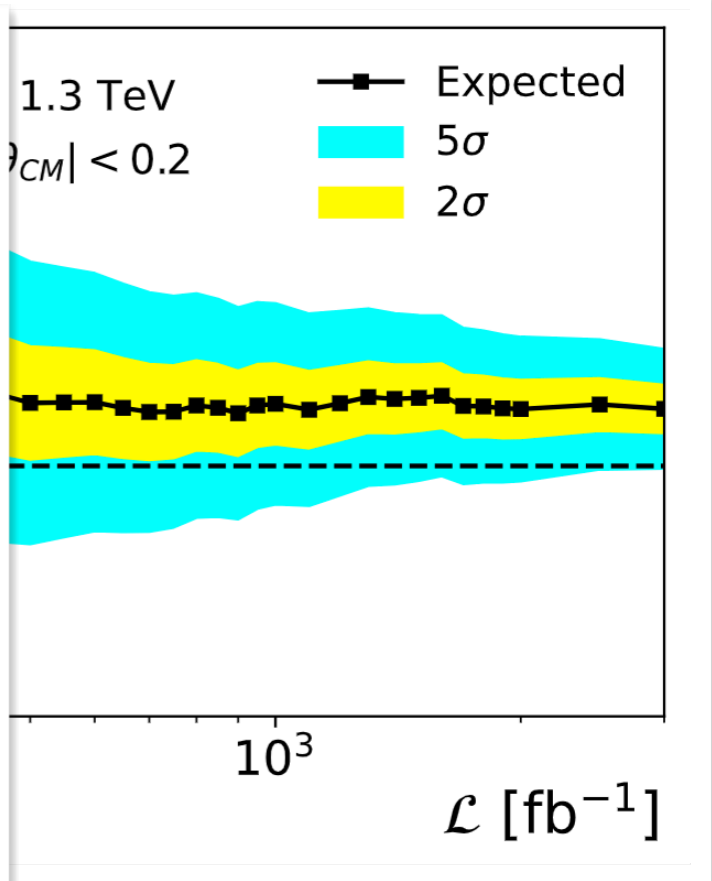
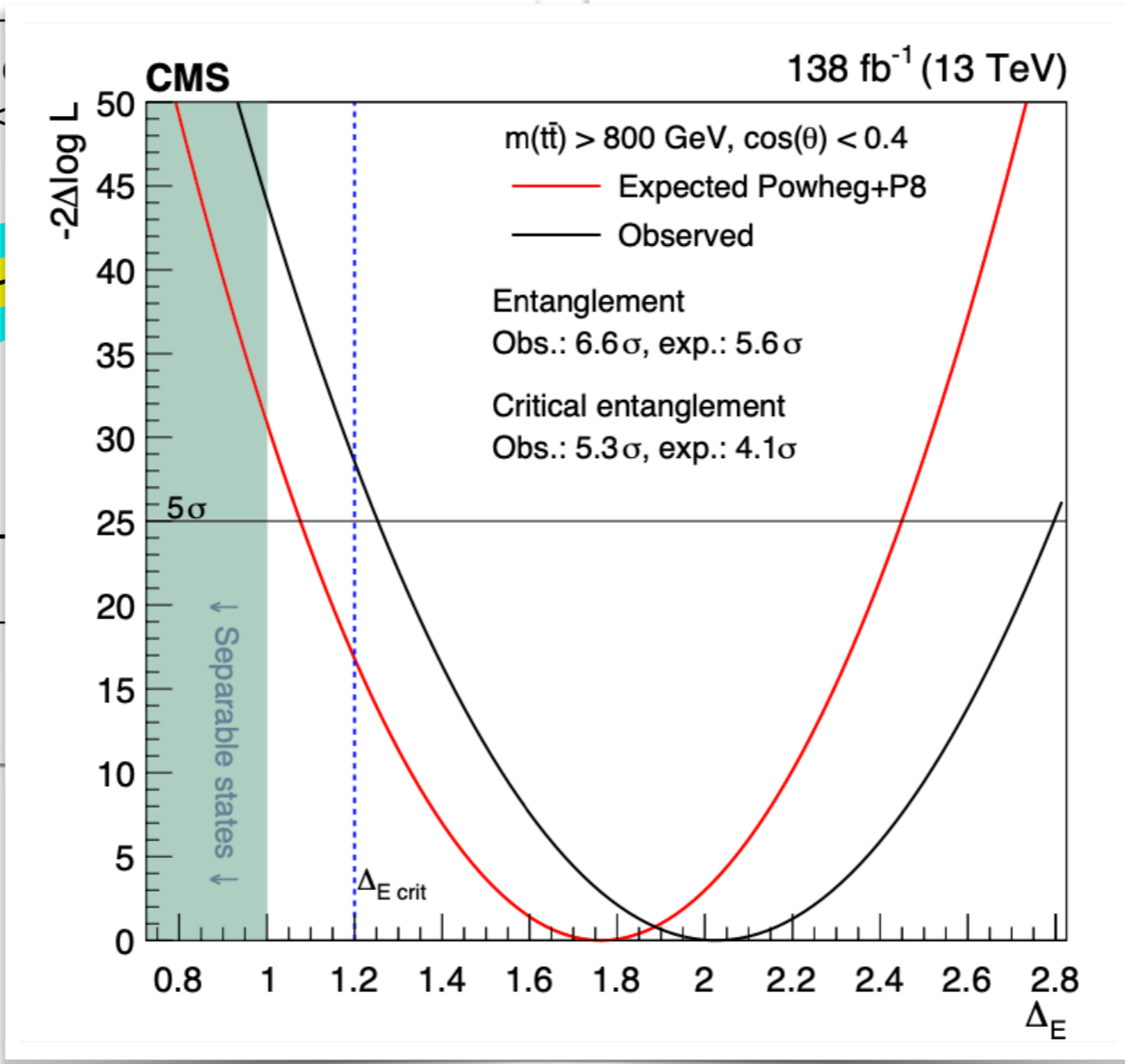
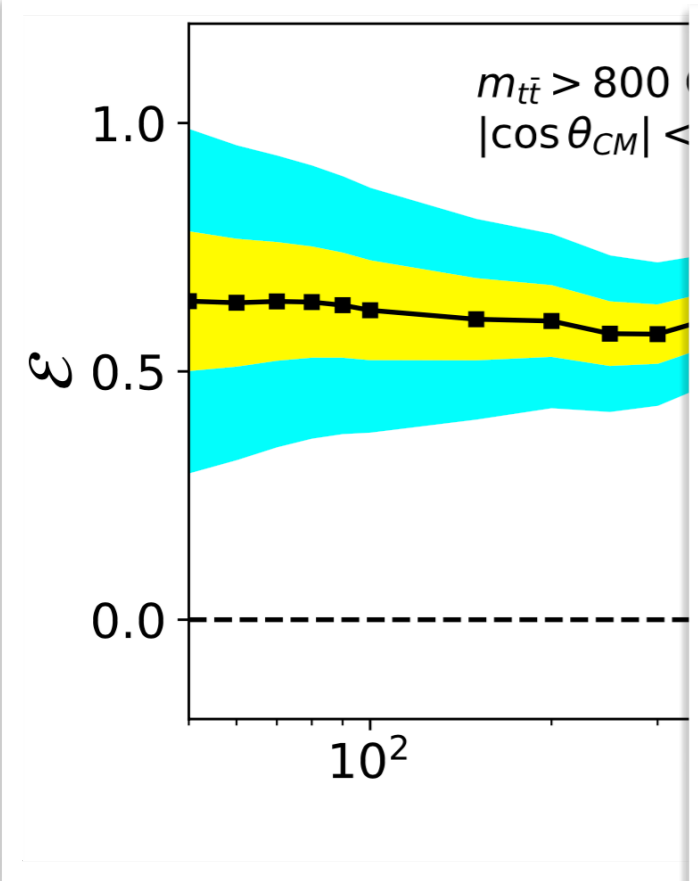
Dong, DG, Kong, Navarro '23

See talks by Eleni Vryonidou, Chris White, Juan Antonio A. Saavedra, Giulia Negro, Roman Lysak

LHC Projections

 Entanglement

 Bell/CHSH inequalities:

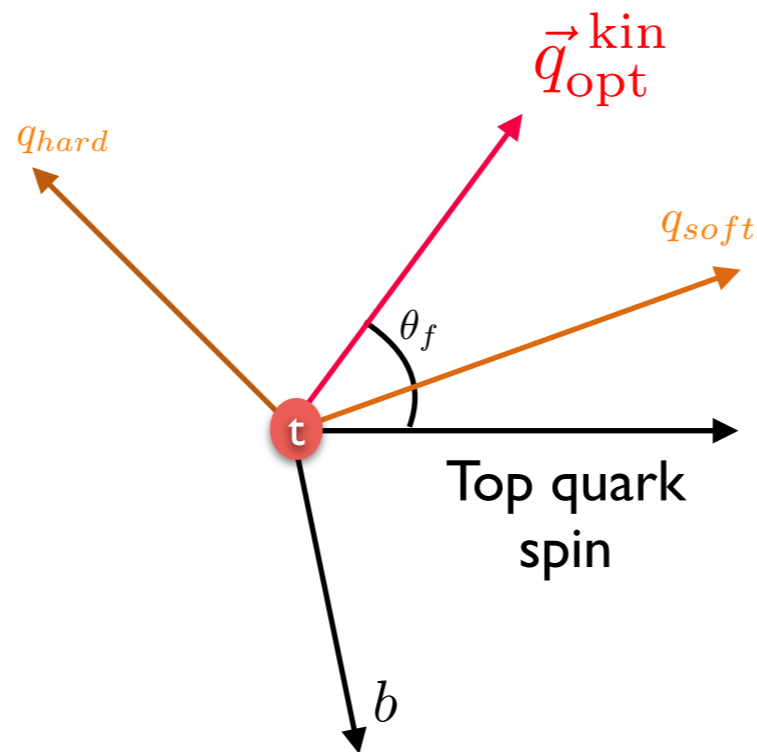


ong, Navarro '23

Eleni Vryonidou, Chris White,
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CMS 2409.11067

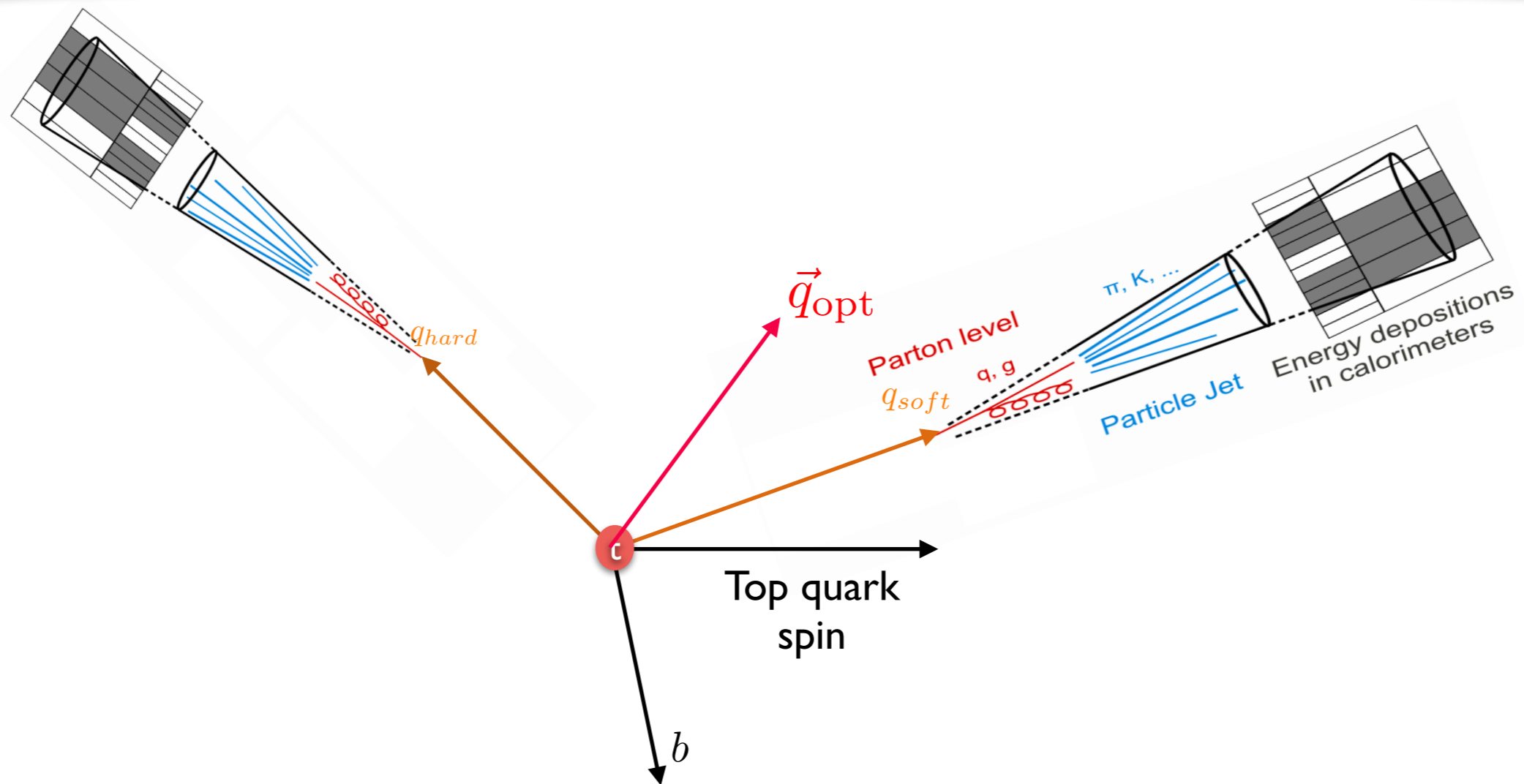
Hadronic Top Quark Polarimetry with ParticleNet



$$\vec{q}_{opt}^{kin} = p(d \rightarrow q_{hard} | c_W) \hat{q}_{hard} + p(d \rightarrow q_{soft} | c_W) \hat{q}_{soft}$$

Dong, DG, Kong, Larkoski, Navarro '24
Dong, DG, Kong, Larkoski, Navarro '24

Hadronic Top Quark Polarimetry with ParticleNet



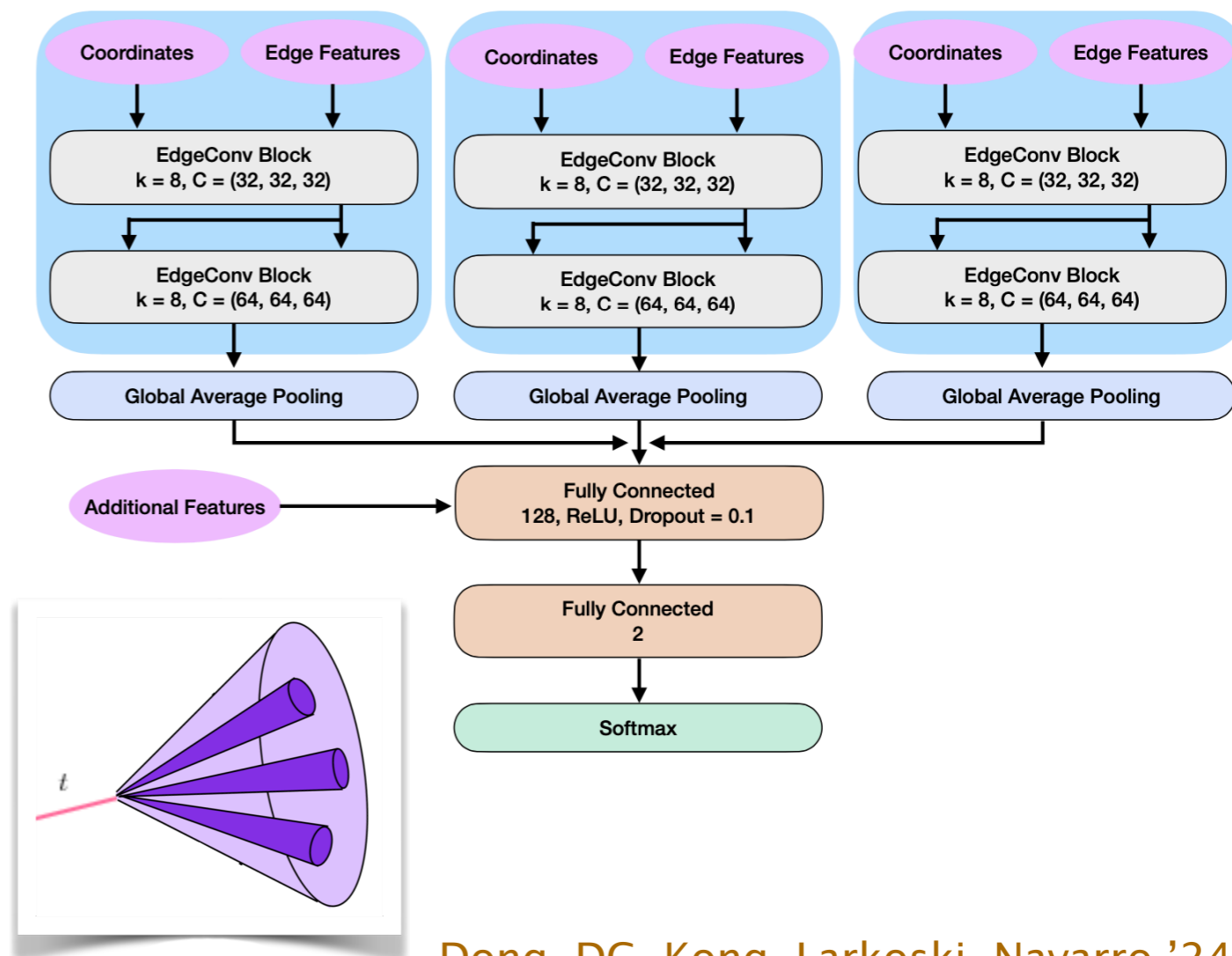
$$\vec{q}_{\text{opt}} = p(d \rightarrow q_{\text{hard}} | c_W, \{\mathcal{O}\}) \hat{q}_{\text{hard}} + p(d \rightarrow q_{\text{soft}} | c_W, \{\mathcal{O}\}) \hat{q}_{\text{soft}}$$

Dong, DG, Kong, Larkoski, Navarro '24
 Dong, DG, Kong, Larkoski, Navarro '24

Hadronic Top Quark Polarimetry with ParticleNet

- Instead of viewing a jet as an ordered set of constituents, it is treated as an unordered set of its constituent particles or a “particle cloud”. This approach efficiently captures low-level jet details while respecting permutation symmetry [Qu, Gouskos 2019](#)

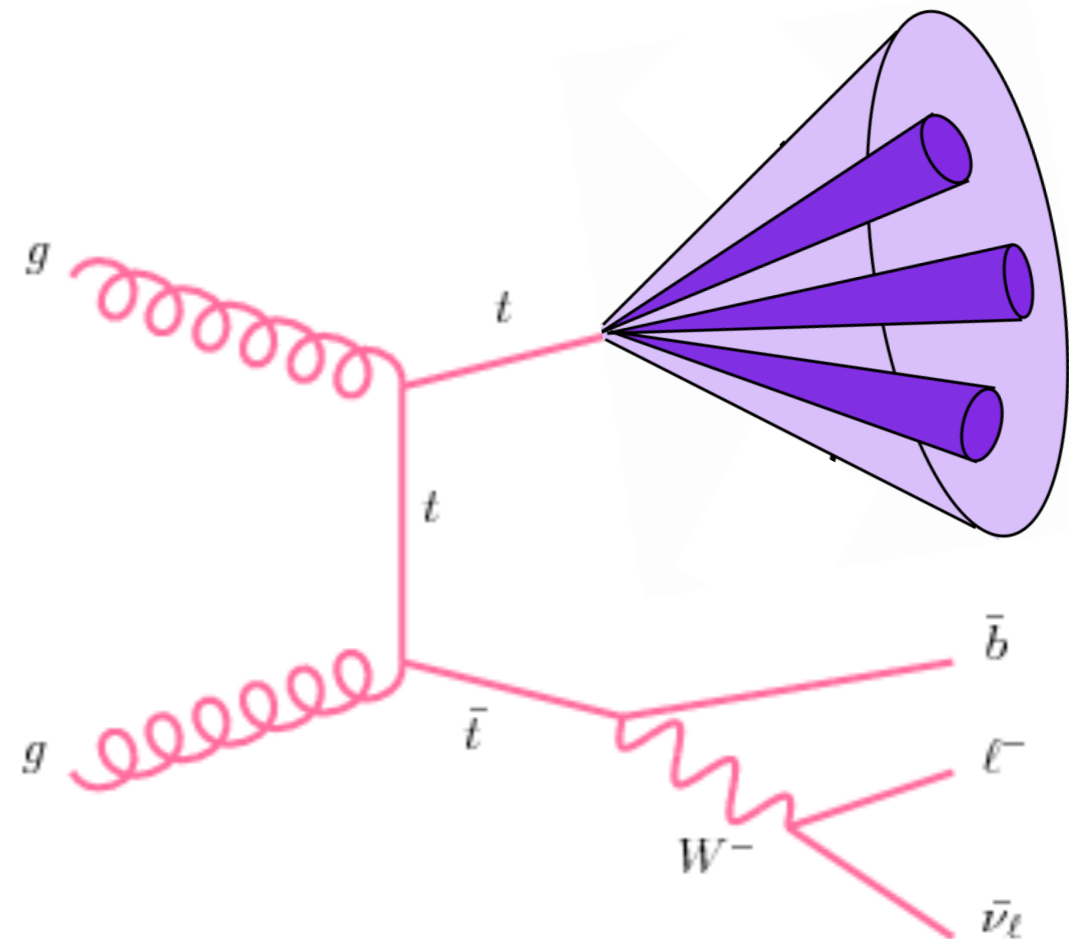
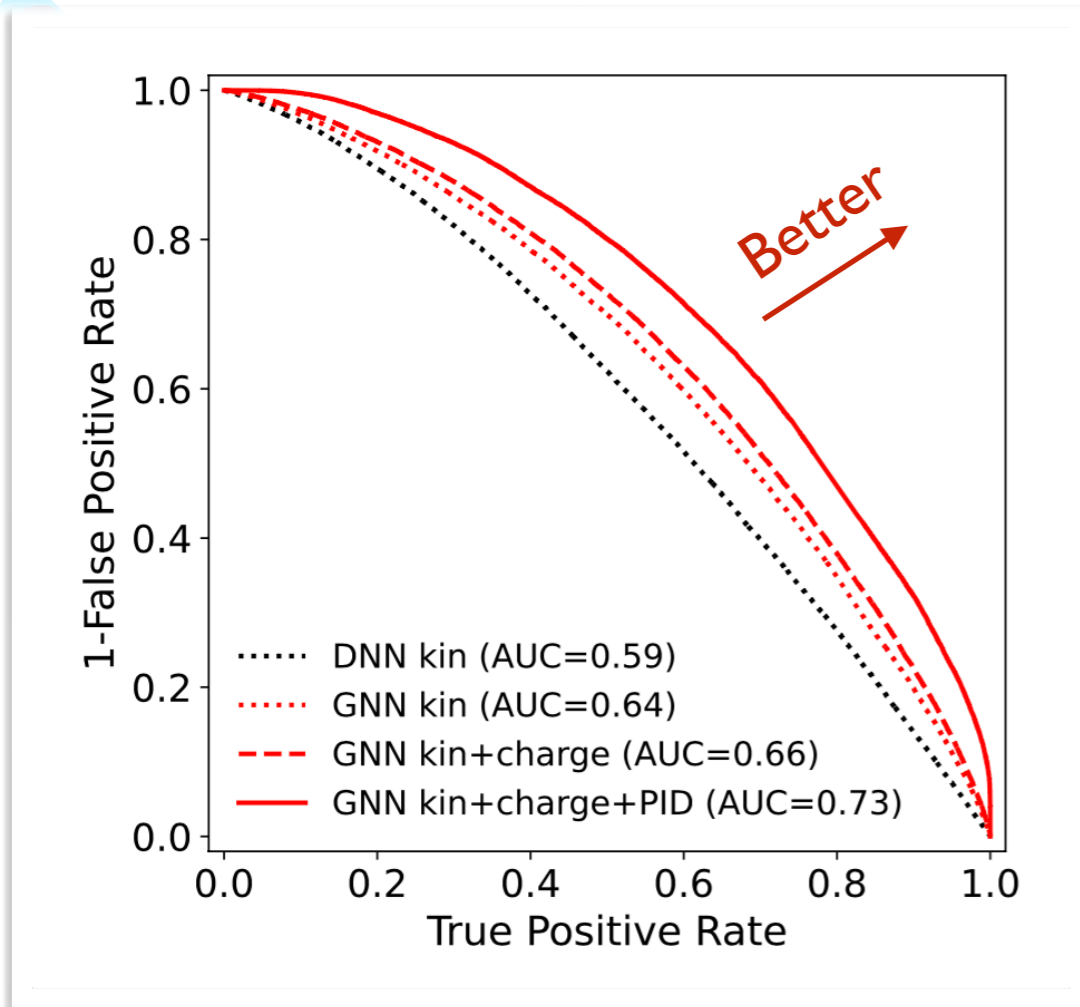
Variable	Definition
$\Delta\eta_t$	difference in pseudorapidity between the particle and the top jet axis
$\Delta\phi_t$	difference in azimuthal angle between the particle and the top jet axis
$\Delta\eta_j$	difference in pseudorapidity between the particle and the subjet axis
$\Delta\phi_j$	difference in azimuthal angle between the particle and the subjet axis
$\log p_T$	logarithm of the particle's p_T
$\log E$	logarithm of the particle's Energy
q	electric charge of the particle
isElectron	if the particle is an electron
isMuon	if the particle is a muon
isPhoton	if the particle is a photon
isChargedHadron	if the particle is a charged hadron
isNeutralHadron	if the particle is a neutral hadron



[Dong, DG, Kong, Larkoski, Navarro '24](#)
[Dong, DG, Kong, Larkoski, Navarro '24](#)

Hadronic Top Quark Polarimetry with ParticleNet

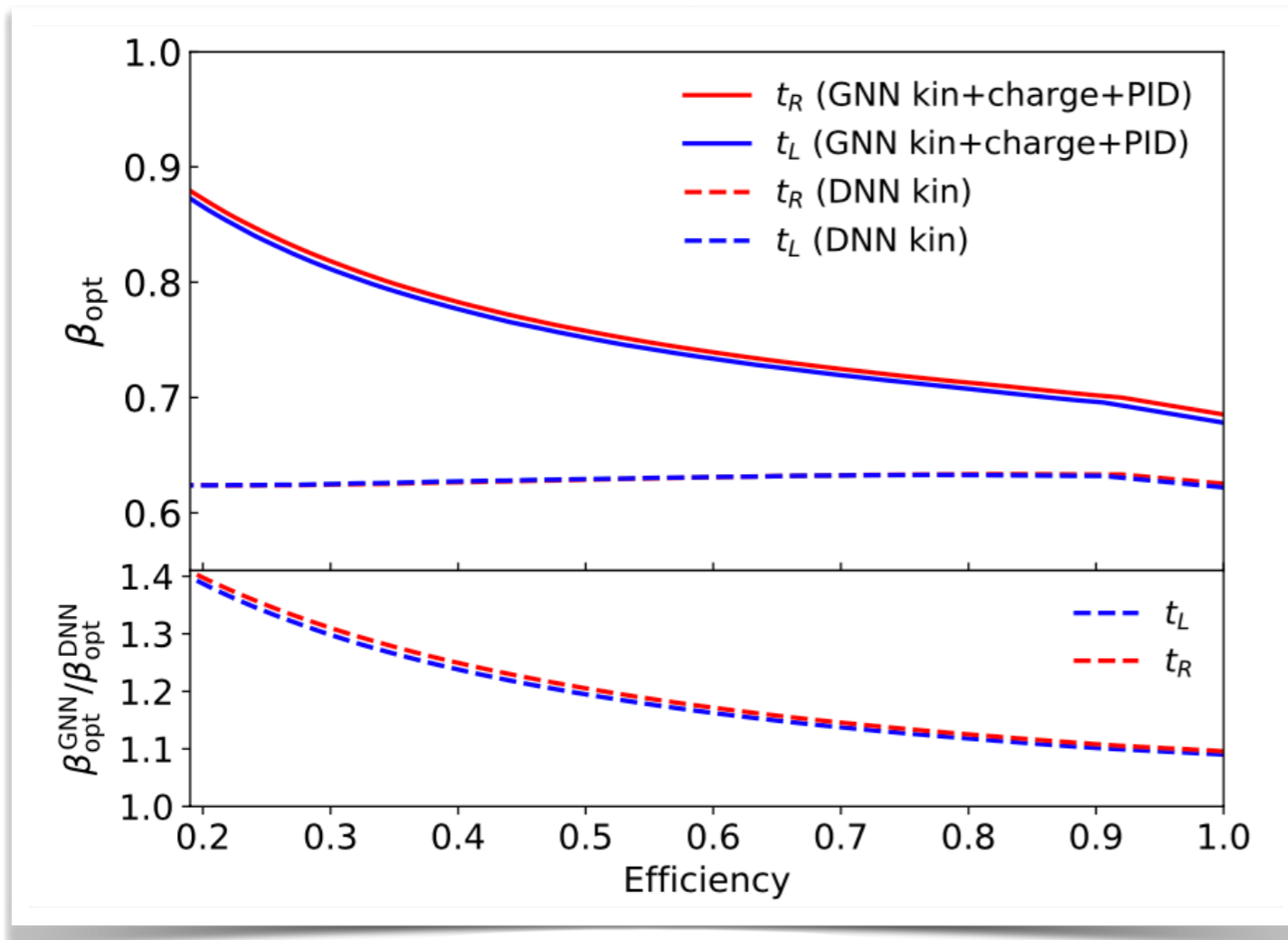
Down-type subjet flavor identification within fatjet



Dong, DG, Kong, Larkoski, Navarro '24

Dong, DG, Kong, Larkoski, Navarro '24

Hadronic Top Quark Polarimetry with ParticleNet



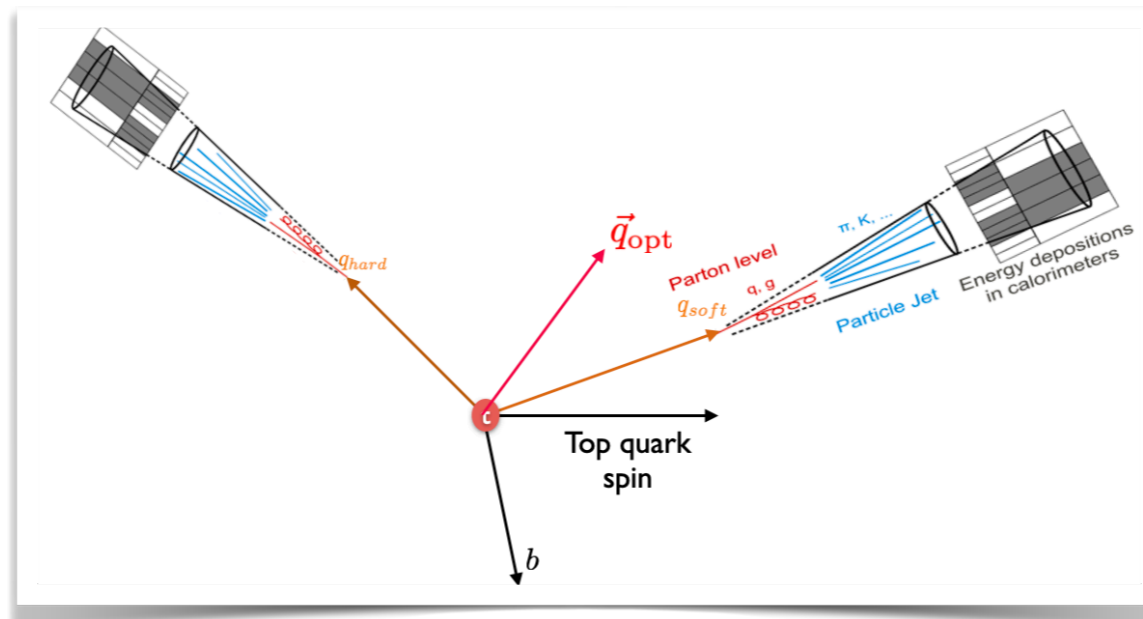
	$\beta_{\text{opt}}^{t_L}$	$\beta_{\text{opt}}^{t_R}$
DNN _{Eff=100%}	0.622	0.625
GNN _{Eff=100%}	0.678	0.685
GNN _{Eff=50%}	0.751	0.758
GNN _{Eff=20%}	0.863	0.869

Dong, DG, Kong, Larkoski, Navarro '24

Dong, DG, Kong, Larkoski, Navarro '24

Summary

- Optimal hadronic polarimeter: crucial ingredient to access higher event rate from hadronic final states
- Down-type Jet flavor tagging significantly improve spin analyzing power in hadronic decays, going beyond exclusive kinematic information
- Imposing an efficiency of 50% (20%) can increase the spin analyzing power of the new artificial direction to $\beta_{\text{opt}}=0.75$ (0.86)
- These improvements can potentially boost spin correlations studies for precision physics, BSM searches, entanglement, Bell inequalities,...



Quantum Entanglement

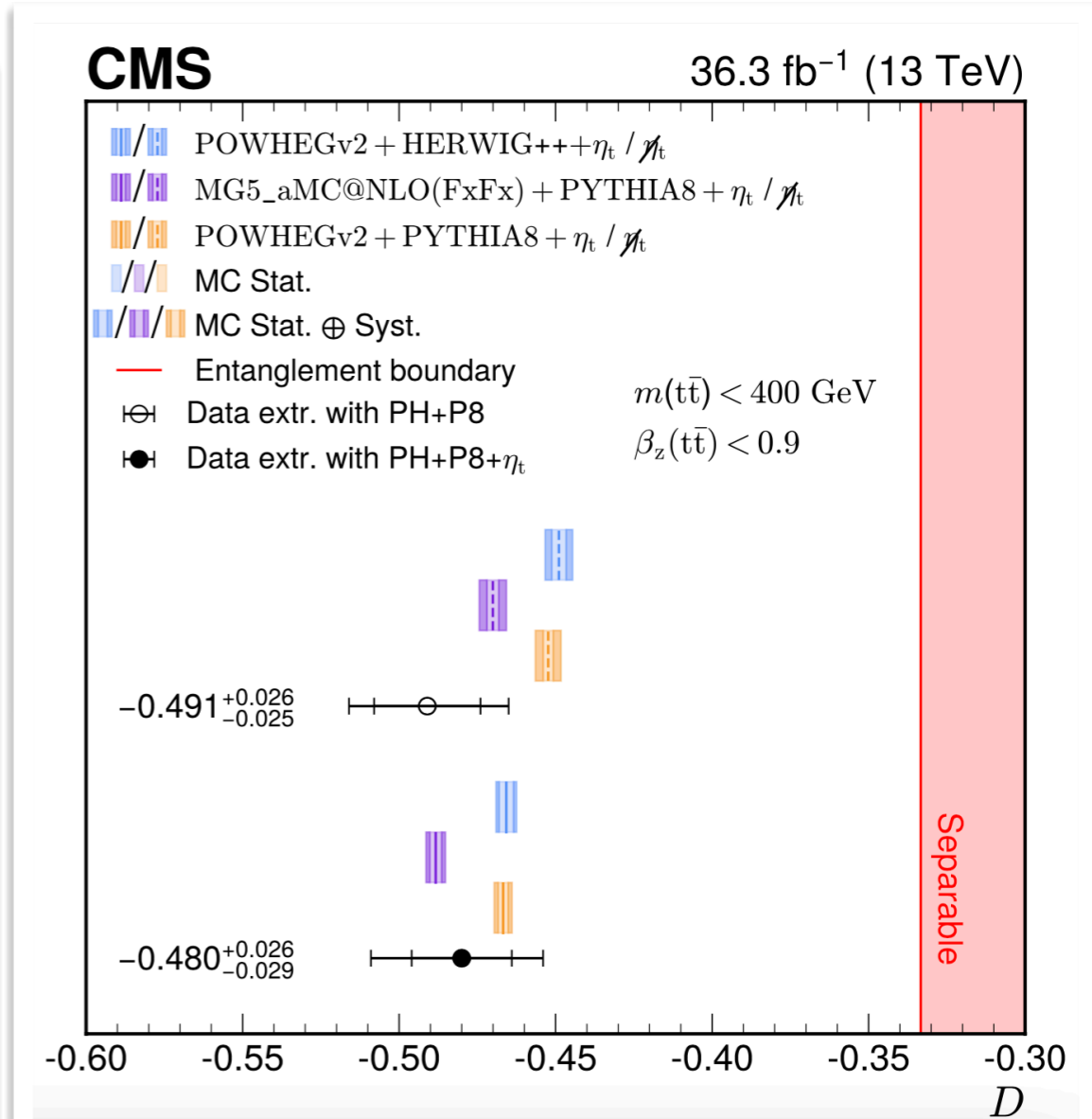
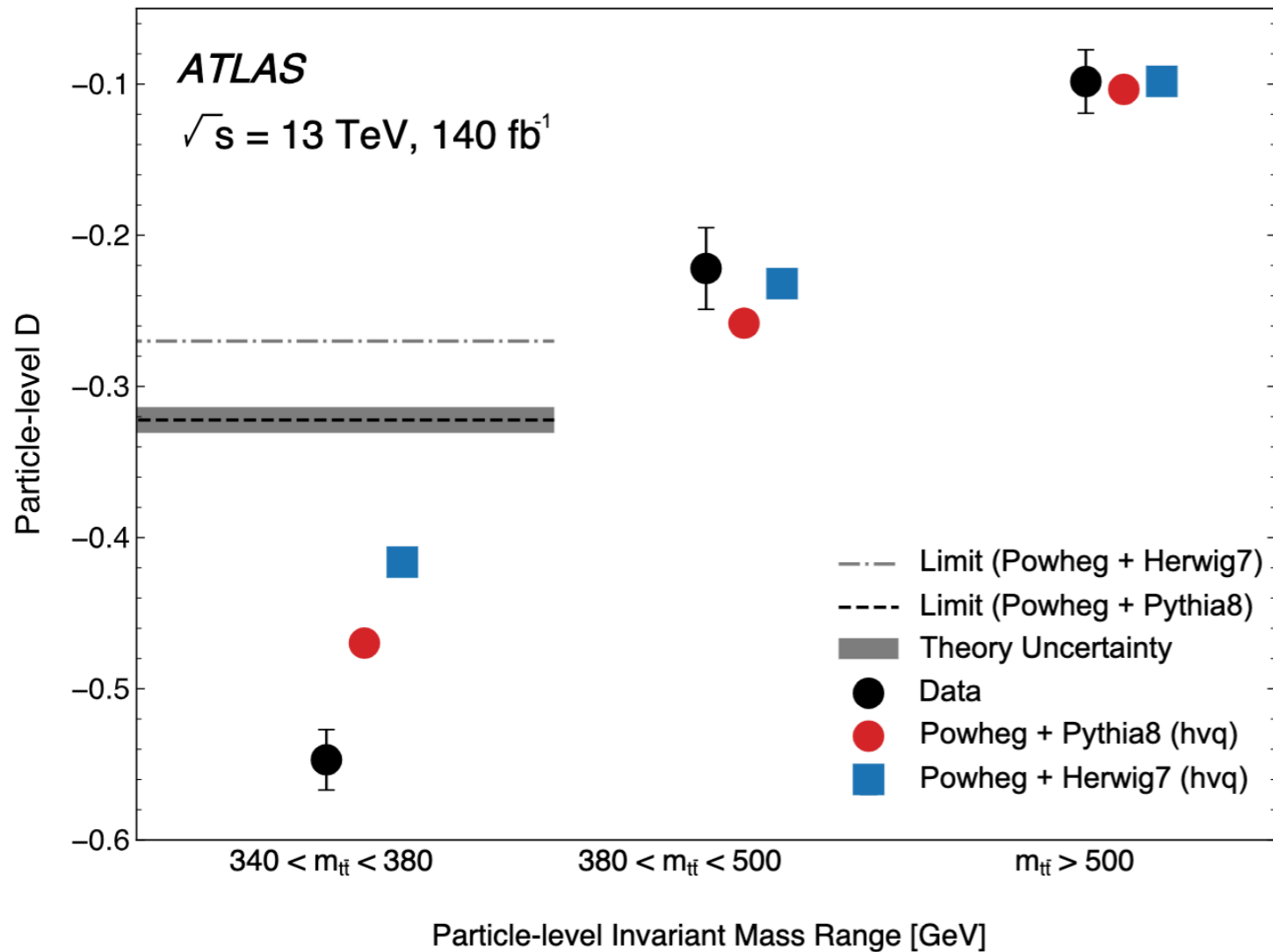
Entanglement stress test our current understanding

- Theoretical modeling
- Experimental uncertainties

ATLAS Nature vol 633, 542-547 (2024)

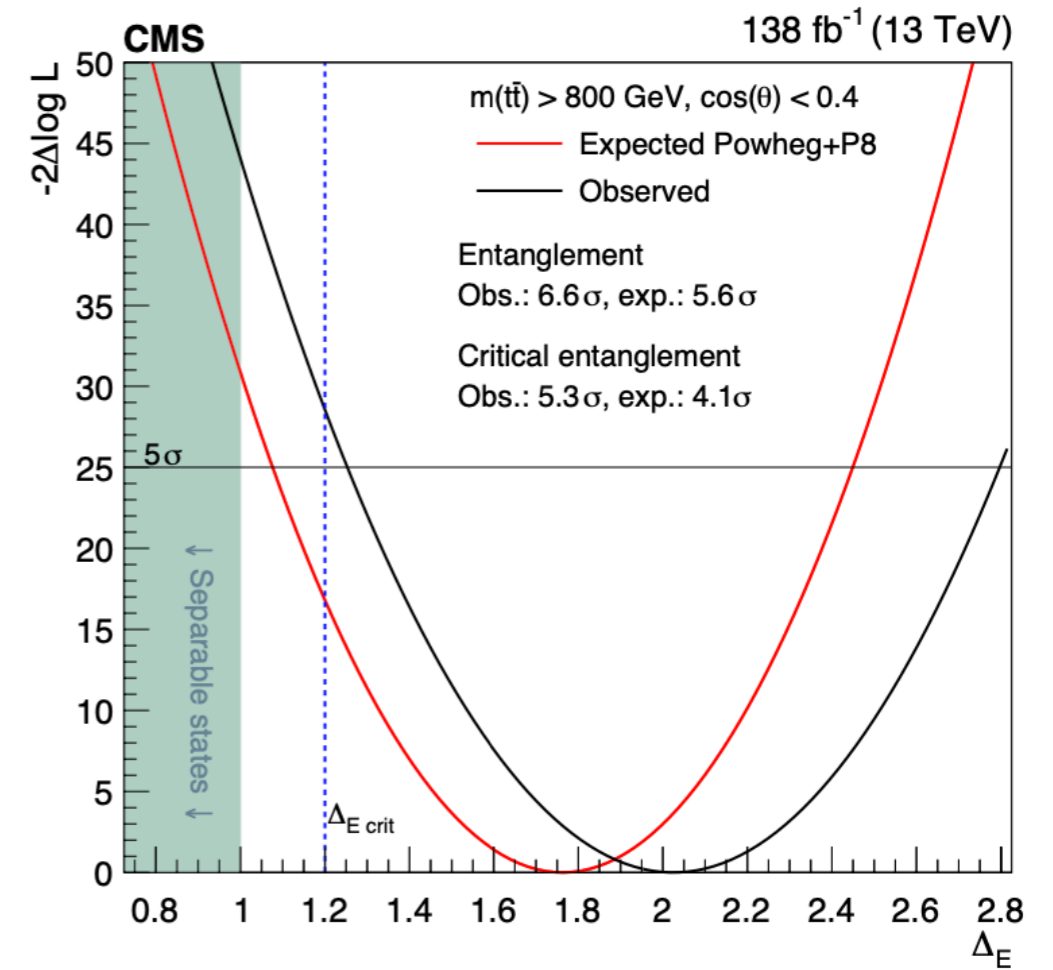
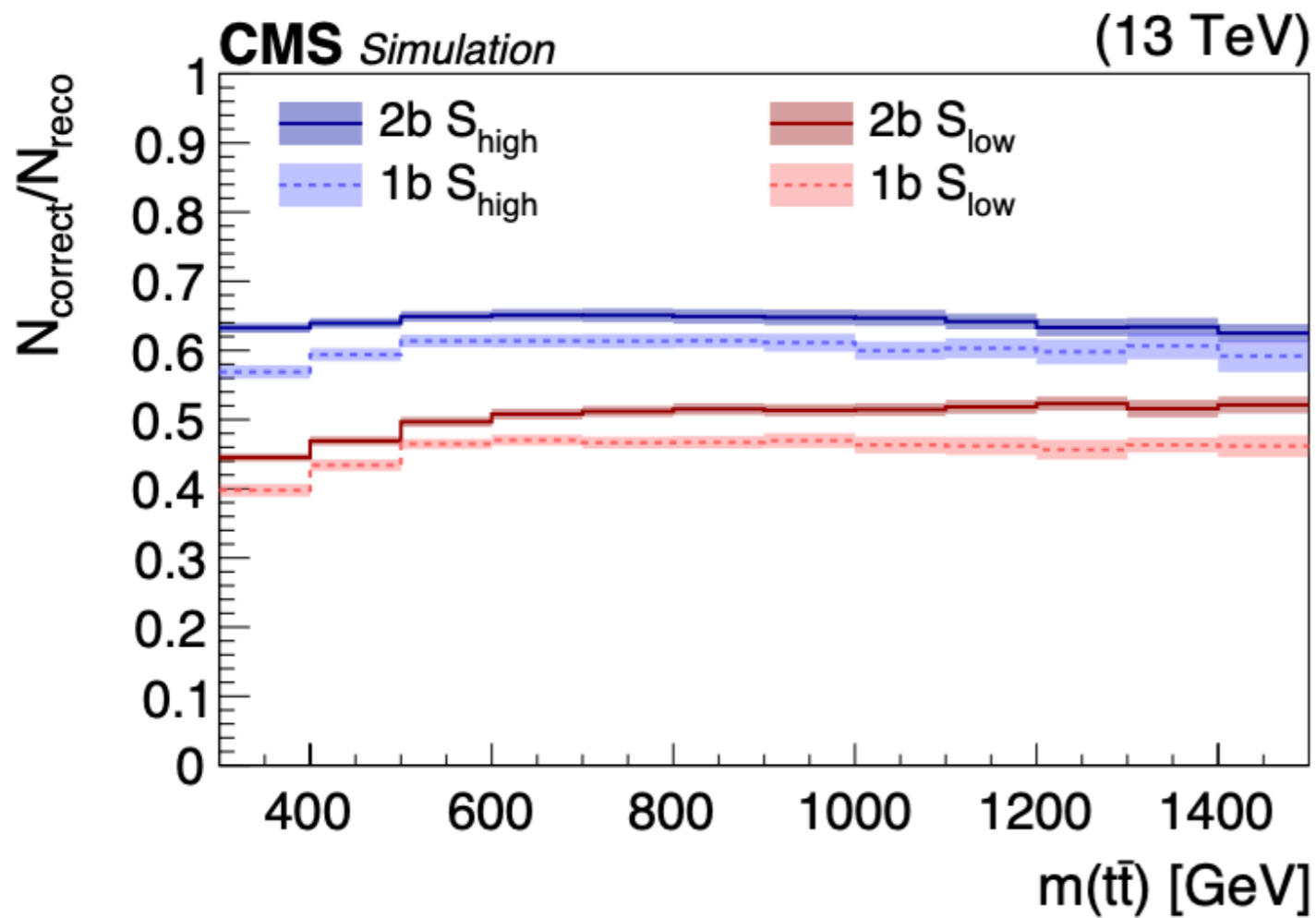
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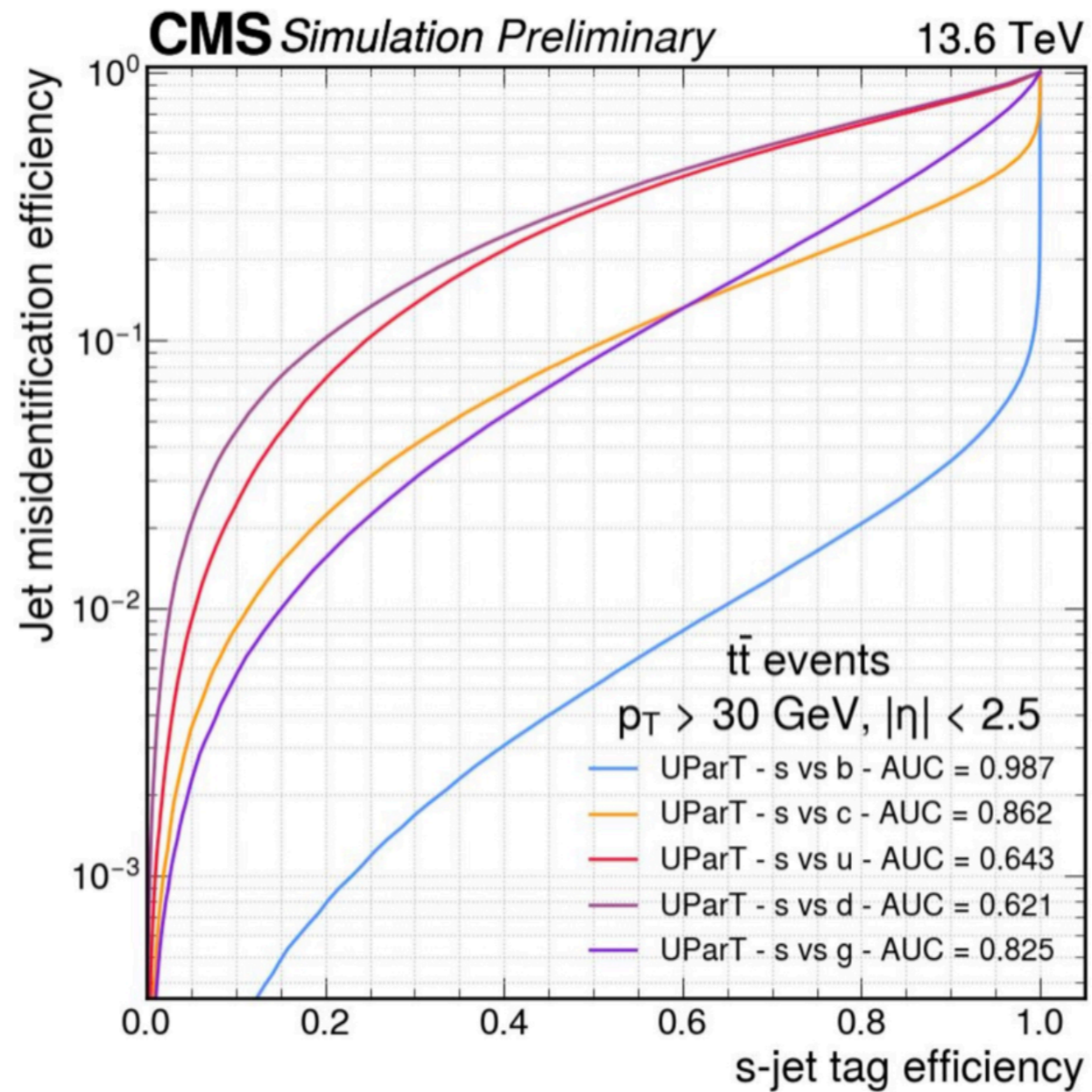


Quantum Entanglement

Semi-leptonic final state



s-tagging performance



See talk by Jieun Choi

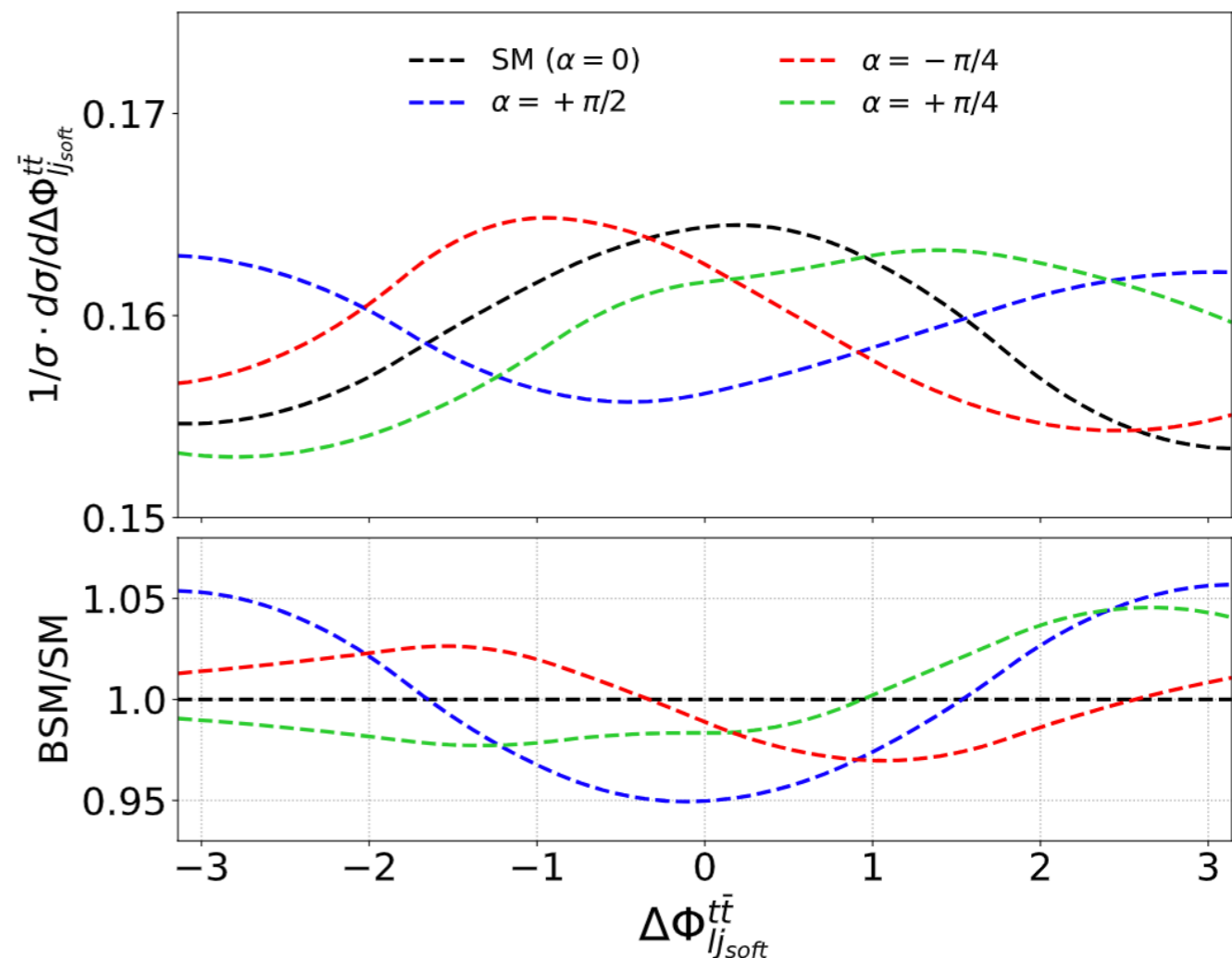
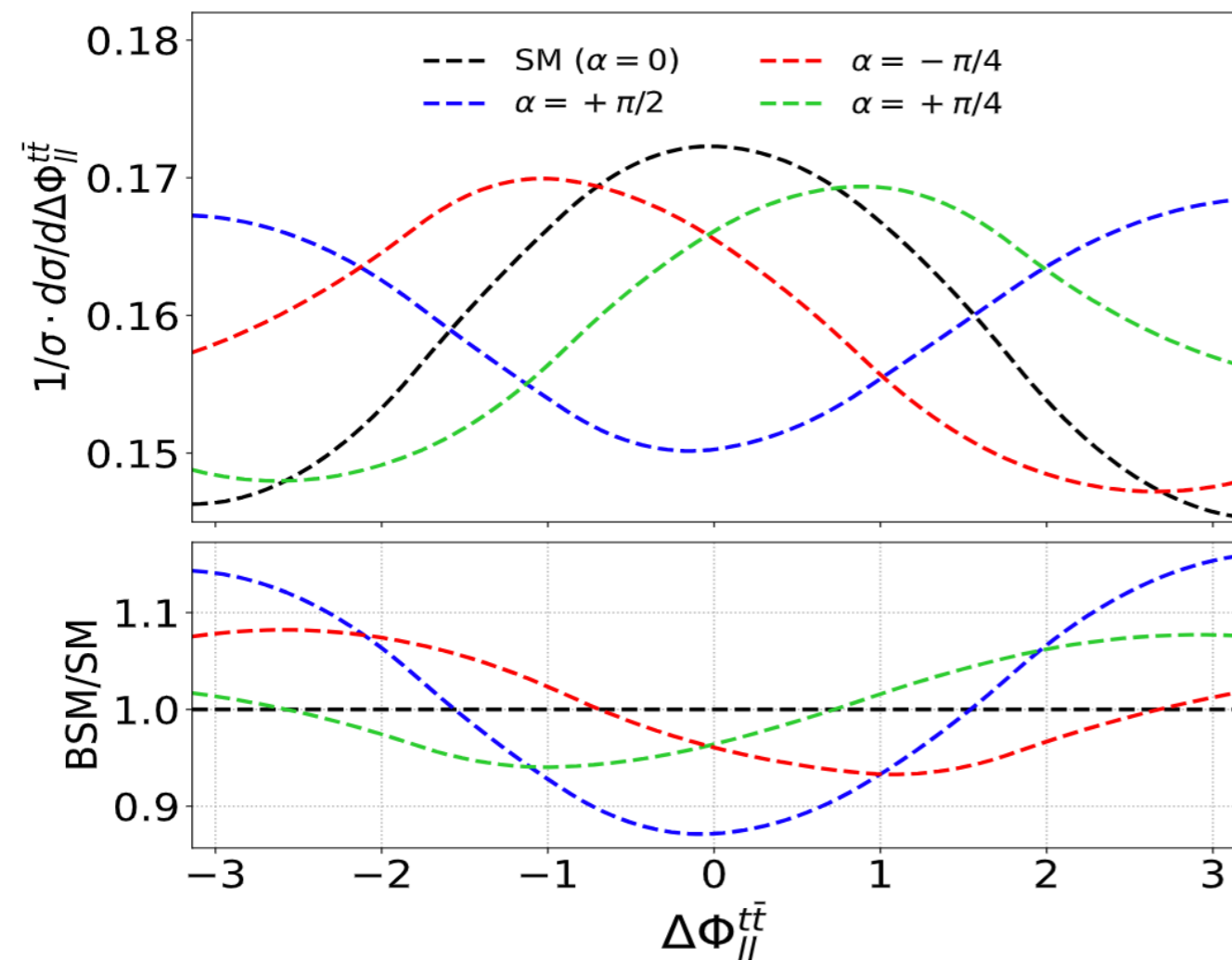
CP sensitive observables

CPV observables best defined at the top pair rest frame:

$$d\sigma(gg \rightarrow t(n_t)\bar{t}(n_{\bar{t}})H) = \sin^2 \alpha f_1(p_i \cdot p_j) + \cos^2 \alpha f_2(p_i \cdot p_j) + \sin \alpha \cos \alpha \sum_l g(p_i \cdot p_j) \epsilon_l$$

$$\epsilon_{\mu\nu\rho\sigma} p_a^\mu p_b^\nu p_c^\rho p_d^\sigma = E_a \vec{p}_b \cdot (\vec{p}_c \times \vec{p}_d) + E_c \vec{p}_d \cdot (\vec{p}_a \times \vec{p}_b) - E_b \vec{p}_c \cdot (\vec{p}_d \times \vec{p}_a) - E_d \vec{p}_a \cdot (\vec{p}_b \times \vec{p}_c)$$

$$\epsilon(p_t, p_{\bar{t}}, p_{\ell^+}, p_{\ell^-})|_{t\bar{t} \text{ CM}} \propto p_t \cdot (p_{\ell^+} \times p_{\ell^-})$$



DG, Kong, Kim '18 & '21, Barman, DG, Kling '21