

Mapping the SMEFT at High-Energy colliders

Based on SMEFIT3.0 [arXiv:2404.12809]

with T. Giani, J. ter Hoeve, L. Mantani, J. Rojo, A. N. Rossia, M. Thomas, E. Vryonidou

Top 2024,

Saint-Malo, Brittany, France

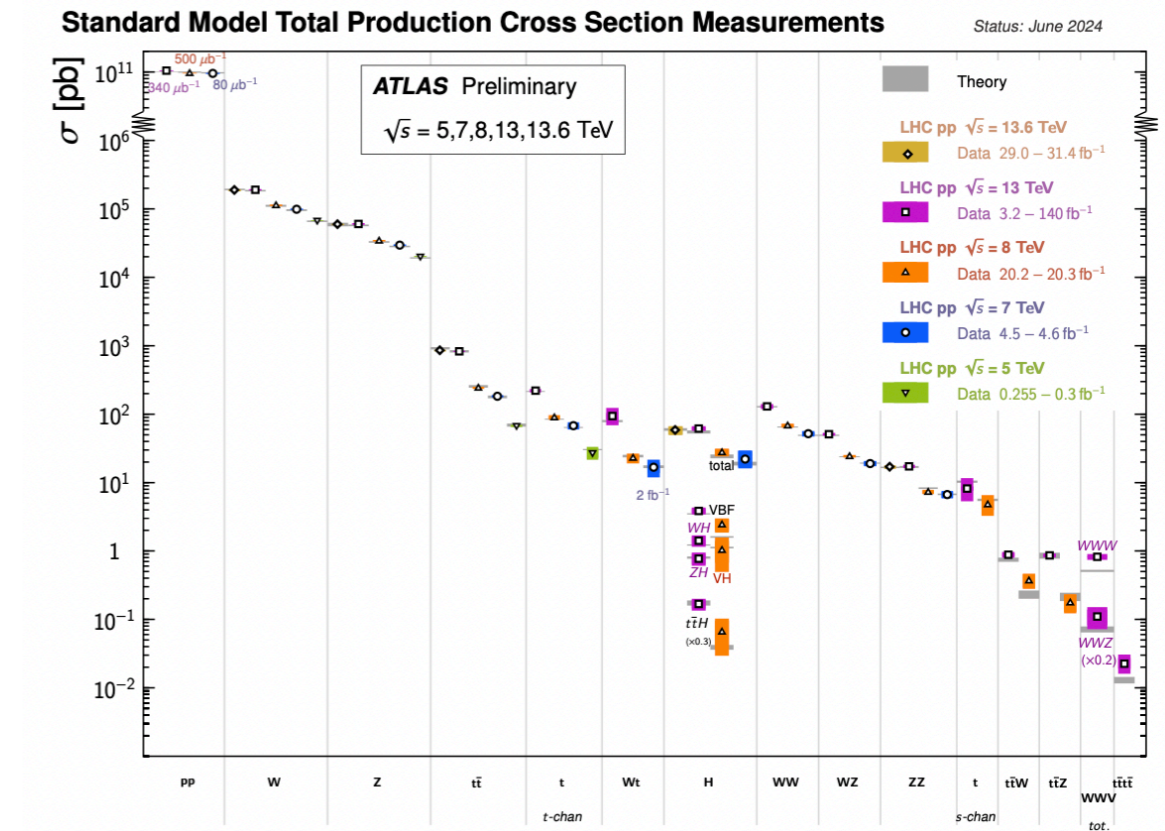
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Eugenia Celada
University of Manchester



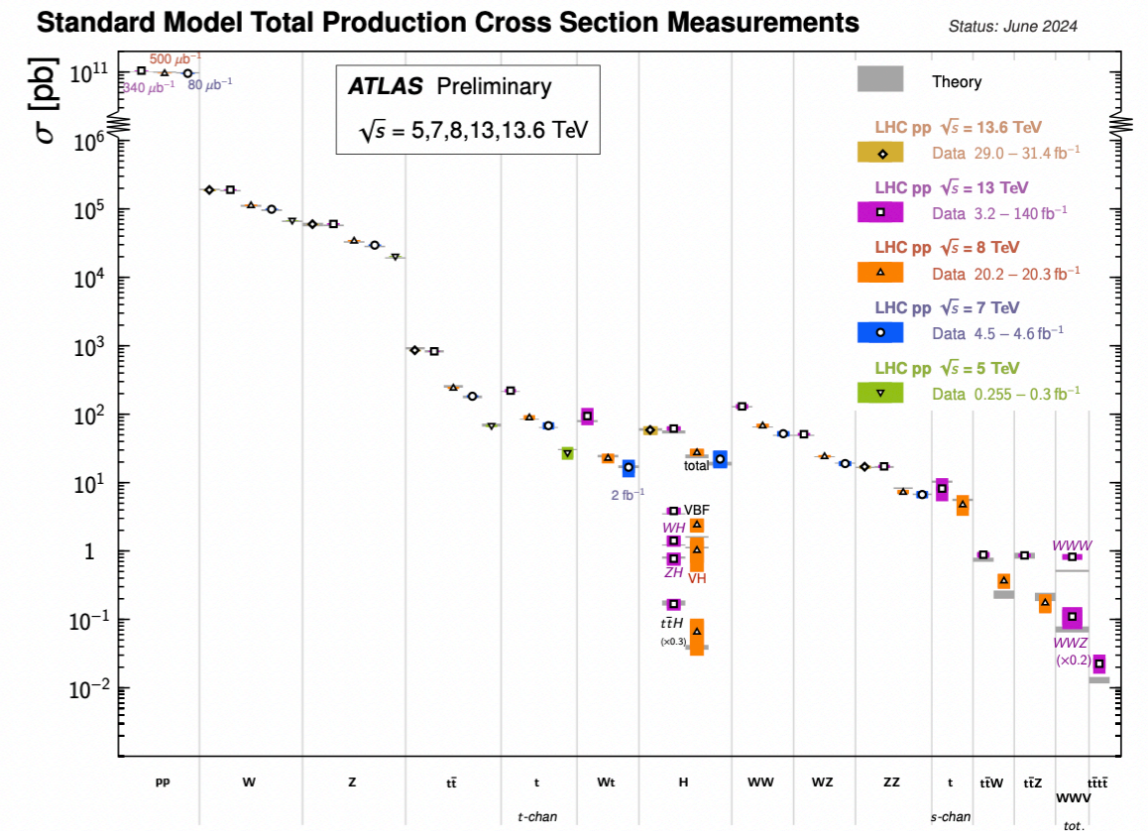
The special role of top

- Lots of top measurements, but no deviation with the SM detected so far
- New physics searches require a **general approach**

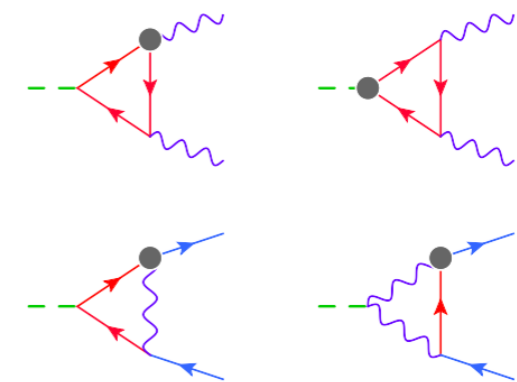


The special role of top

- Lots of top measurements, but no deviation with the SM detected so far
- New physics searches require a **general approach**



- **Top plays a special role in the SM**
 - largest fermion mass
 - largest Yukawa coupling
 - top-loop effects are crucial in Higgs and EW processes



[arXiv: 1804.09766]

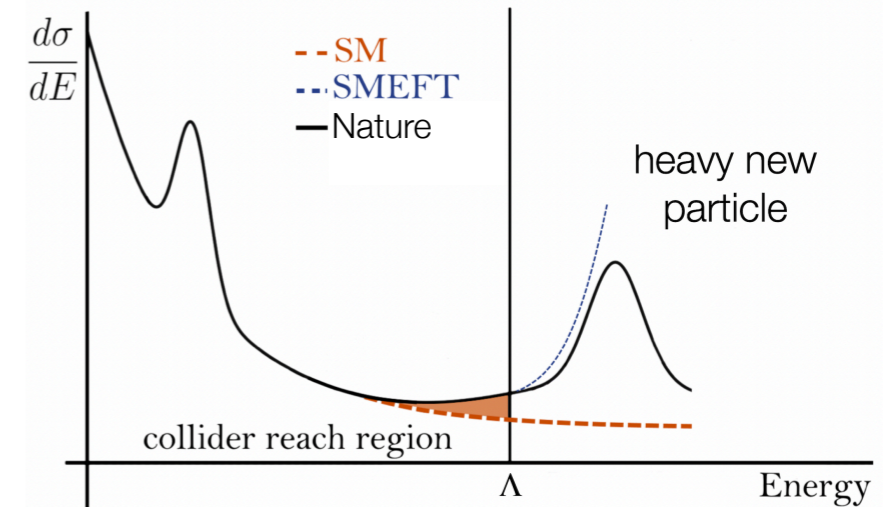
- **Cross-talk between different sectors** (Higgs, top, EW) requires a global analysis

SMEFT global fits

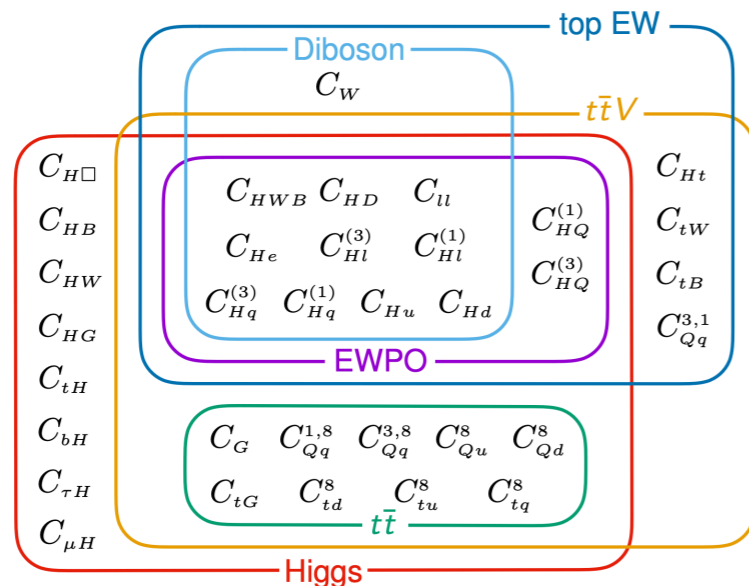
- The SMEFT reveals high energy physics effects through precise measurements at low energy

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_i \frac{c_i^{(6)}}{\Lambda^2} \mathcal{O}_i^{(6)} + \mathcal{O}(\Lambda^{-3})$$

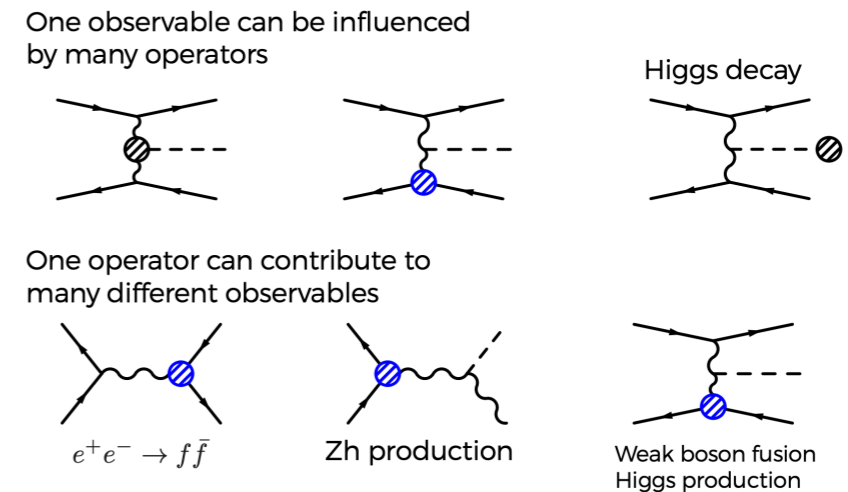
$$\sigma = |\mathcal{M}_{\text{SM}}|^2 + \frac{1}{\Lambda^2} \left(\sum c^{(6)} 2\text{Re}[\mathcal{M}_{\text{SM}}^* \mathcal{M}_{\text{EFT}}^{(6)}] \right) + \frac{1}{\Lambda^4} \left(\sum c^{(6)} \mathcal{M}_{\text{EFT}}^{(6)} \right)^2$$



- Ideal framework to look for overall pattern in global fits
- Large number of operator coefficients, many datasets needed to break degeneracies

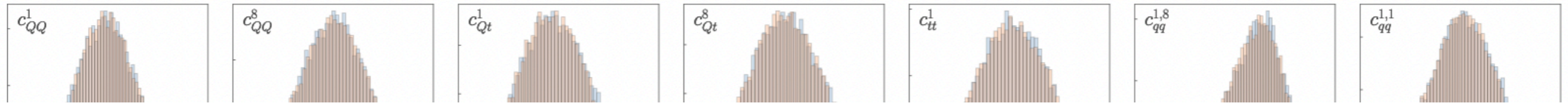


Fitmaker Collaboration [arXiv:2012.02779]



Anke Biekötter - HET seminar Brookhaven

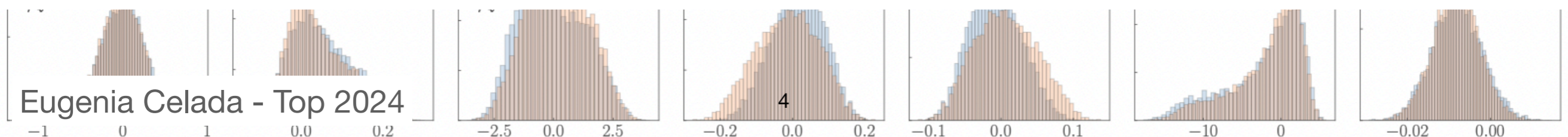
SMEFIT3.0: a summary



- Extension of SMEFIT2.0 with recent **LHC Run-II datasets** on **top, diboson and Higgs production**
[arXiv:2105.00006]
- Exact treatment of LEP and SLD Electroweak Precision Observables (EWPOs)
- **Projections for HL-LHC pseudodata** extrapolated from Run-II data
- **FCC-ee and CEPC pseudodata** from Snowmass predictions updated with FCC midterm Feasibility Report
[arXiv:2206.08326]
[CERN/3789/RA]
- Results for LHC Run-II and future colliders in terms of Wilson coefficients and **UV-complete models**
[arXiv:2309.04523]
- **Public code, data and theory:** results are fully reproducible



[lhcfithef.github.io/smefit_release/]



The SMEFiT setup

Experimental data

445 data points from Higgs, top, diboson (LHC) & EWPOs (LEP)

Experimental **uncertainties + correlations** as provided by experiments



Giani, Magni, Rojo, arXiv:2302.06660

The SMEFiT setup

Theory

SM: **(N)NLO QCD + NLO EW**

EFT: **NLO QCD**, linear and quadratics, with
SMEFT@NLO

NNPDF4.0 no top

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Methodology

Linear fit: **Analytical solution**

Quadratic fit: **Nested sampling**
(Bayesian inference)

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Output

Automatised fit report with **bounds** on
coefficients, **posterior** distributions,
PCA, Fisher information...

Experimental input

Extension of SMEFiT2.0 with recent **LHC Run-II datasets** on **top, diboson and Higgs production**

SMEFiT3.0: EC, Giani, ter Hoeve, Mantani, Rojo, Rossia, Thomas, Vryonidou [arXiv:2404.12809]

SMEFiT2.0: Ethier, Maltoni, Mantani, Nocera, Rojo, Slade, Vryonidou, Zhang [arXiv:2105.00006]

Category	Processes	n_{dat}	
		SMEFiT2.0	SMEFiT3.0
Top quark production	$t\bar{t} + X$	94	115
	$t\bar{t}Z, t\bar{t}W$	14	21
	$t\bar{t}\gamma$	-	2
	single top (inclusive)	27	28
	tZ, tW	9	13
	$t\bar{t}t\bar{t}, t\bar{t}b\bar{b}$	6	12
	Total	150	191
Higgs production and decay	Run I signal strengths	22	22
	Run II signal strengths	40	36 (*)
	Run II, differential distributions & STXS	35	71
	Total	97	129
Diboson production	LEP-2	40	40
	LHC	30	41
	Total	70	81
EWPOs	LEP-2	-	44
Baseline dataset	Total	317	445

+118 data points

Theory input: operators

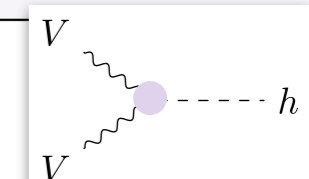
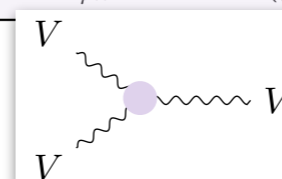
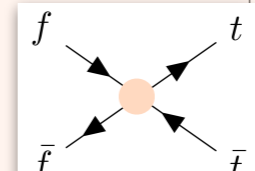
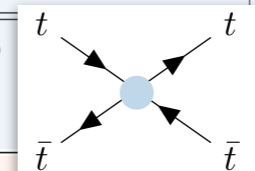
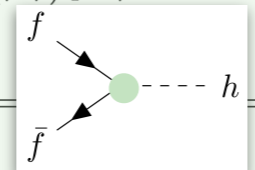
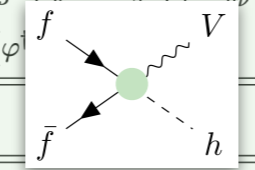
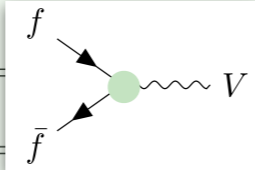
- Warsaw basis of dim-6 SMEFT operators
- Flavour symmetry:
 $U(2)_q \times U(3)_d \times U(2)_u \times (U(1)_l \times U(1)_e)^3$
 + Yukawa of bottom, charm and tau

Simultaneous fit of 50
Wilson coefficients

Operator	Coefficient	Definition	Operator	Coefficient	Definition
3rd generation quarks					
$\mathcal{O}_{\varphi Q}^{(1)}$	$c_{\varphi Q}^{(1)}$ (*)	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{Q} \gamma^\mu Q)$	\mathcal{O}_{tW}	c_{tW}	$i(\bar{Q} \tau^{\mu\nu} \tau_I t) \tilde{\varphi} W_{\mu\nu}^I + \text{h.c.}$
$\mathcal{O}_{\varphi Q}^{(3)}$	$c_{\varphi Q}^{(3)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \tau_I \varphi)(\bar{Q} \gamma^\mu \tau^I Q)$	\mathcal{O}_{tB}	c_{tB} (*)	$i(\bar{Q} \tau^{\mu\nu} t) \tilde{\varphi} B_{\mu\nu} + \text{h.c.}$
$\mathcal{O}_{\varphi t}$	$c_{\varphi t}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{t} \gamma^\mu t)$	\mathcal{O}_{tG}	c_{tG}	$ig_s(\bar{Q} \tau^{\mu\nu} T_A t) \tilde{\varphi} G_{\mu\nu}^A + \text{h.c.}$
$\mathcal{O}_{t\varphi}$	$c_{t\varphi}$	$(\varphi^\dagger \varphi) \bar{Q} t \tilde{\varphi} + \text{h.c.}$	$\mathcal{O}_{b\varphi}$	$c_{b\varphi}$	$(\varphi^\dagger \varphi) \bar{b} t \tilde{\varphi} + \text{h.c.}$
1st, 2nd generation quarks					
$\mathcal{O}_{\varphi q}^{(1)}$	$c_{\varphi q}^{(1)}$ (*)	$\sum_{i=1,2} i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{q}_i \gamma^\mu q_i)$	$\mathcal{O}_{\varphi d}$	$c_{\varphi d}$	$\sum_{i=1,2,3} i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{d}_i \gamma^\mu d_i)$
$\mathcal{O}_{\varphi q}^{(3)}$	$c_{\varphi q}^{(3)}$	$\sum_{i=1,2} i(\varphi^\dagger \overleftrightarrow{D}_\mu \tau_I \varphi)(\bar{q}_i \gamma^\mu \tau^I q_i)$	$\mathcal{O}_{c\varphi}$	$c_{c\varphi}$	$(\varphi^\dagger \varphi) \bar{q}_2 c \tilde{\varphi} + \text{h.c.}$
$\mathcal{O}_{\varphi u}$	$c_{\varphi u}$	$\sum_{i=1,2} i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{u}_i \gamma^\mu u_i)$			
two-leptons					
$\mathcal{O}_{\varphi l_i}$	$c_{\varphi l_i}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{l}_i \gamma^\mu l_i)$	$\mathcal{O}_{\varphi \mu}$	$c_{\varphi \mu}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{\mu} \gamma^\mu \mu)$
$\mathcal{O}_{\varphi l_i}^{(3)}$	$c_{\varphi l_i}^{(3)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \tau_I \varphi)(\bar{l}_i \gamma^\mu \tau^I l_i)$	$\mathcal{O}_{\varphi \tau}$	$c_{\varphi \tau}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{\tau} \gamma^\mu \tau)$
$\mathcal{O}_{\varphi e}$	$c_{\varphi e}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{e} \gamma^\mu e)$	$\mathcal{O}_{\tau\varphi}$	$c_{\tau\varphi}$	$(\varphi^\dagger \varphi) \bar{l}_3 \tau \varphi + \text{h.c.}$
four-leptons					
$\mathcal{O}_{\ell\ell}$	$c_{\ell\ell}$	$(\bar{l}_1 \gamma_\mu l_2)(\bar{l}_2 \gamma^\mu l_1)$			

DoF	Definition (in Warsaw basis notation)	DoF	Definition (in Warsaw basis notation)
c_{QQ}^1	$2c_{qq}^{1(3333)} - \frac{2}{3}c_{qq}^{3(3333)}$	c_{QQ}^8	$8c_{qq}^{3(3333)}$
c_{Qt}^1	$c_{qu}^{1(3333)}$	c_{Qt}^8	$8c_{qu}^{3(3333)}$
$c_{Qq}^{1,8}$	$c_{qq}^{1(i33i)} + 3c_{qq}^{3(i33i)}$	$c_{Qq}^{1,1}$	$c_{qq}^{1(ii33)} + \frac{1}{6}c_{qq}^{1(i33i)} + \frac{1}{2}c_{qq}^{3(i33i)}$
$c_{Qq}^{3,8}$	$c_{qq}^{1(i33i)} - c_{qq}^{3(i33i)}$	$c_{Qq}^{3,1}$	$c_{qq}^{3(ii33)} + \frac{1}{6}(c_{qq}^{1(i33i)} - c_{qq}^{3(i33i)})$
c_{tq}^8	$c_{qu}^{8(ii33)}$	c_{tq}^1	$c_{qu}^{1(ii33)}$
c_{tu}^8	$2c_{uu}^{(ii33)}$	c_{tu}^1	$c_{uu}^{(ii33)} + \frac{1}{3}c_{uu}^{(i33i)}$
c_{Qu}^8	$c_{qu}^{8(33ii)}$	c_{Qu}^1	$c_{qu}^{1(33ii)}$
c_{td}^8	$c_{ud}^{8(33jj)}$	c_{td}^1	$c_{ud}^{1(33jj)}$
c_{Qd}^8	$c_{qd}^{8(33jj)}$	c_{Qd}^1	$c_{qd}^{1(33jj)}$

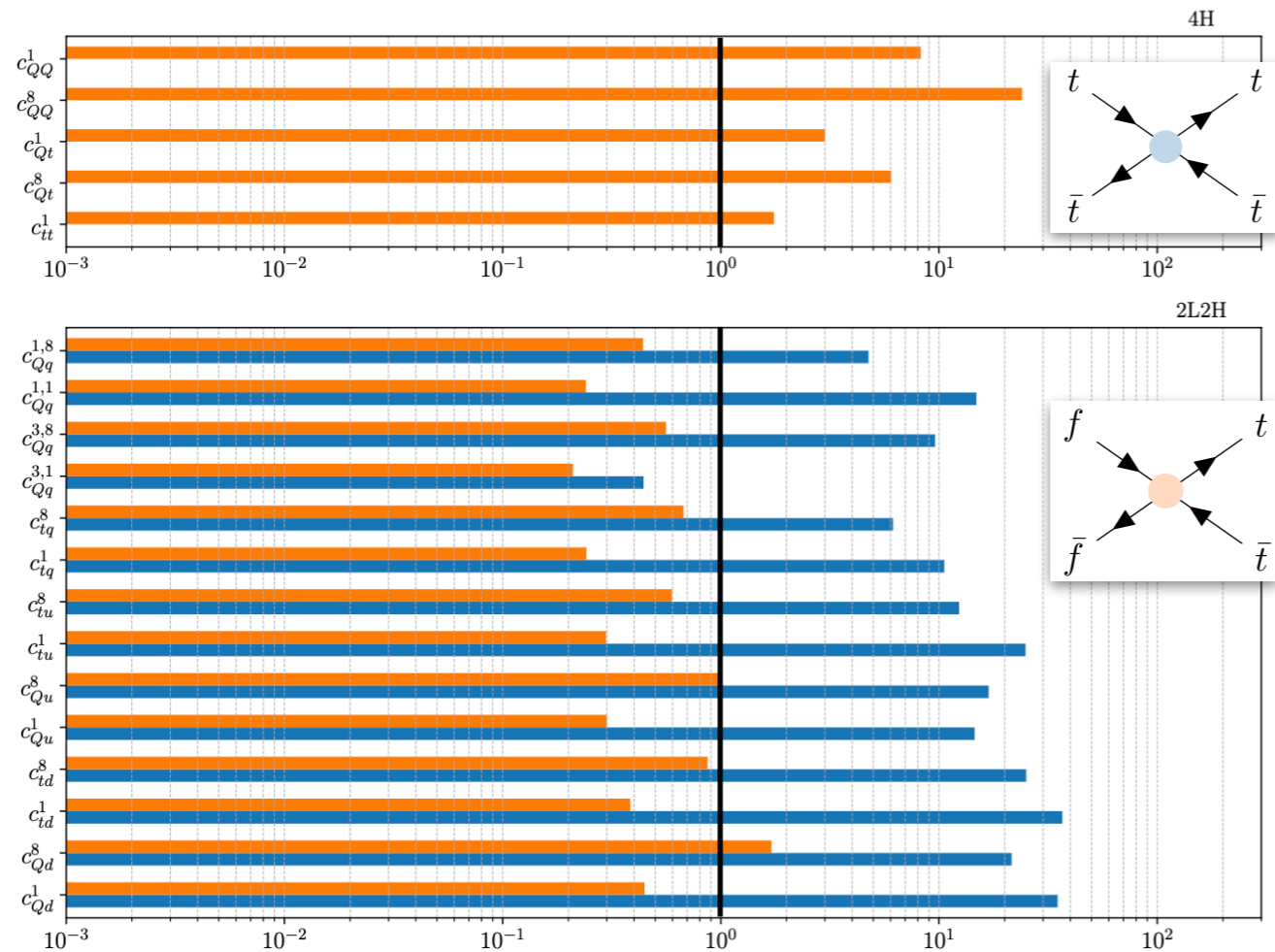
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$\mathcal{O}_{\varphi G}$	$c_{\varphi G}$	$(\varphi^\dagger \varphi) G_A^{\mu\nu} G_{\mu\nu}^A$	$\mathcal{O}_{\varphi \square}$	$c_{\varphi \square}$	$\partial_\mu(\varphi^\dagger \varphi) \partial^\mu(\varphi^\dagger \varphi)$
$\mathcal{O}_{\varphi B}$	$c_{\varphi B}$	$(\varphi^\dagger \varphi) B^{\mu\nu} B_{\mu\nu}$	$\mathcal{O}_{\varphi D}$	$c_{\varphi D}$	$(\varphi^\dagger D^\mu \varphi)^\dagger (\varphi^\dagger D_\mu \varphi)$
$\mathcal{O}_{\varphi W}$	$c_{\varphi W}$	$(\varphi^\dagger \varphi) W_I^{\mu\nu} W_{\mu\nu}^I$	\mathcal{O}_W	c_{WWW}	$\epsilon_{IJK} W_{\mu\nu}^I W^{\nu\rho J} W_\rho^{K,\mu}$
$\mathcal{O}_{\varphi WB}$	$c_{\varphi WB}$	$(\varphi^\dagger \tau_I \varphi) B^{\mu\nu} W_{\mu\nu}^I$			



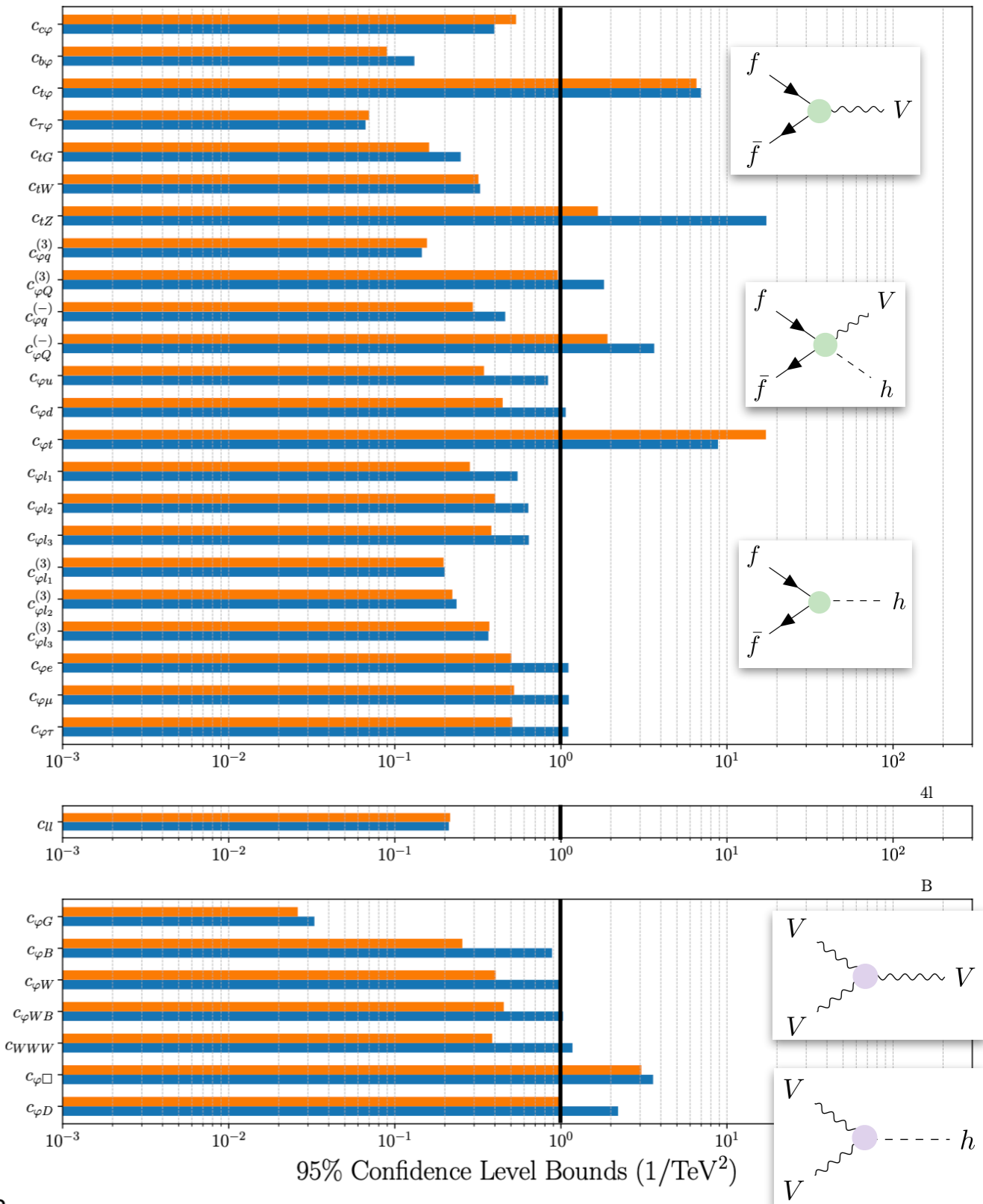
Current global fit results

Marginalised **linear** and **quadratic**

■ NLO $\mathcal{O}(\Lambda^{-2})$ ■ NLO $\mathcal{O}(\Lambda^{-4})$



- Most bounds below 1 for $\Lambda = 1$ TeV



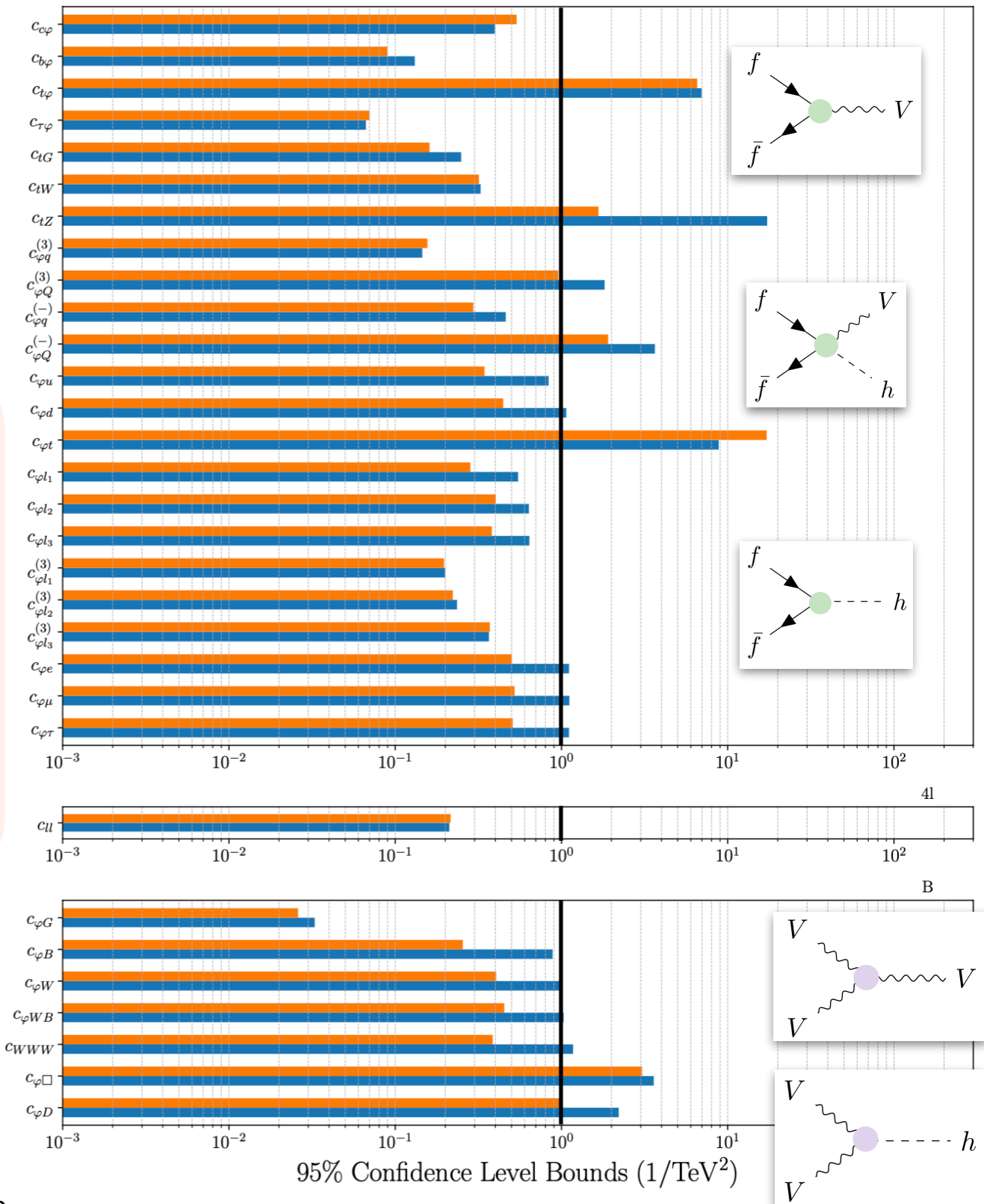
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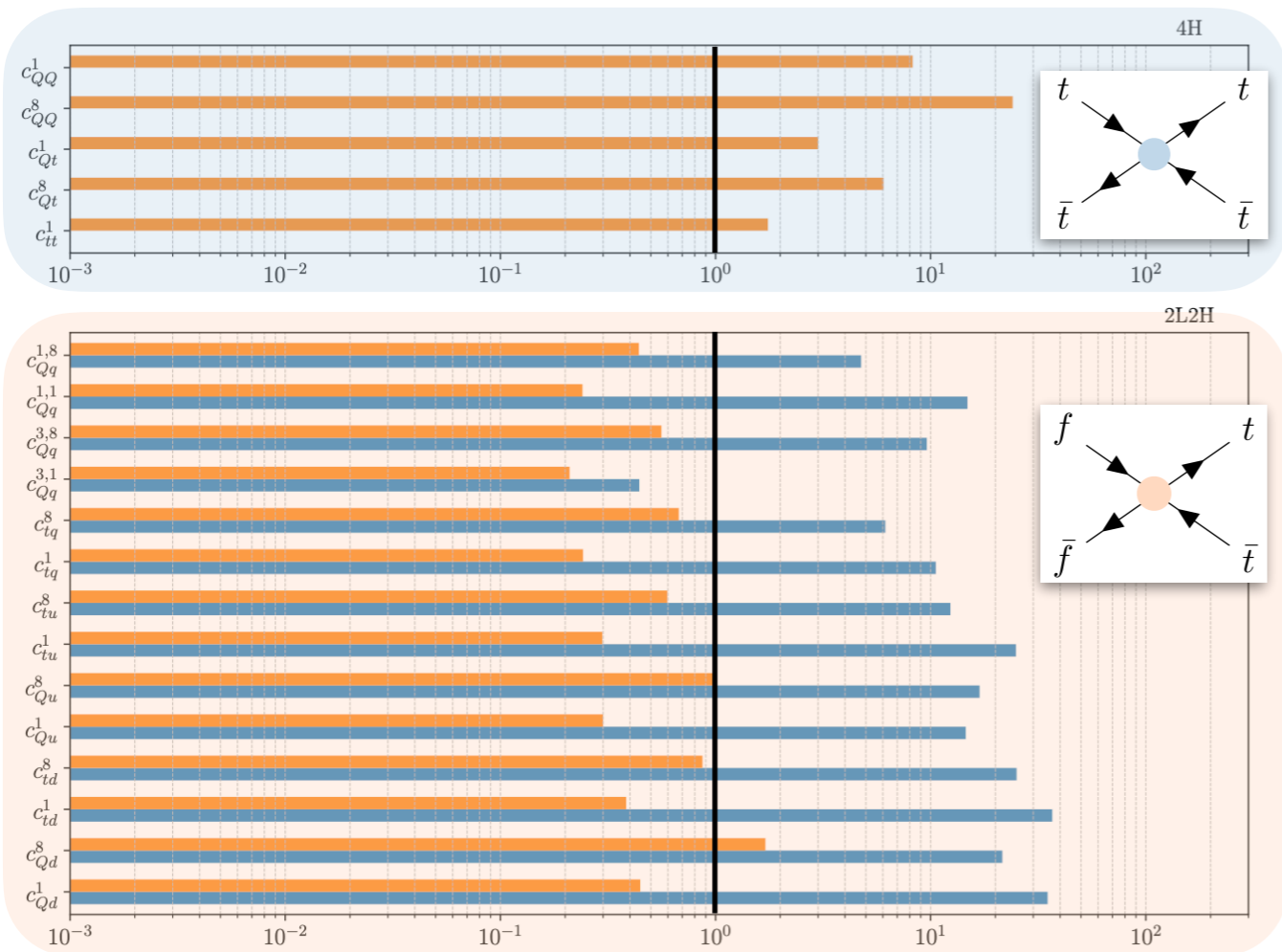
- Most bounds below 1 for $\Lambda = 1$ TeV
- Quadratic terms important (mostly **2-light-2-heavy**)



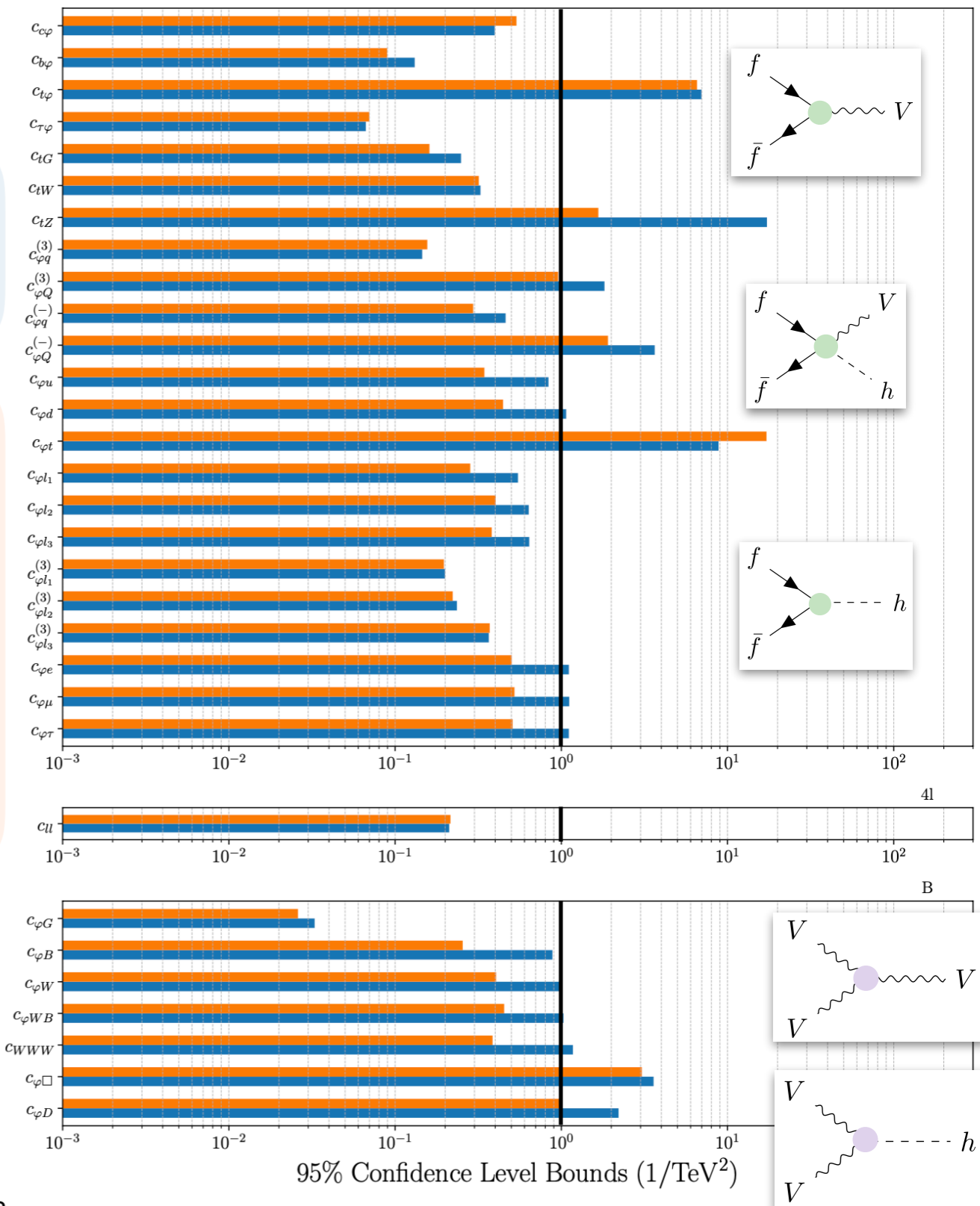
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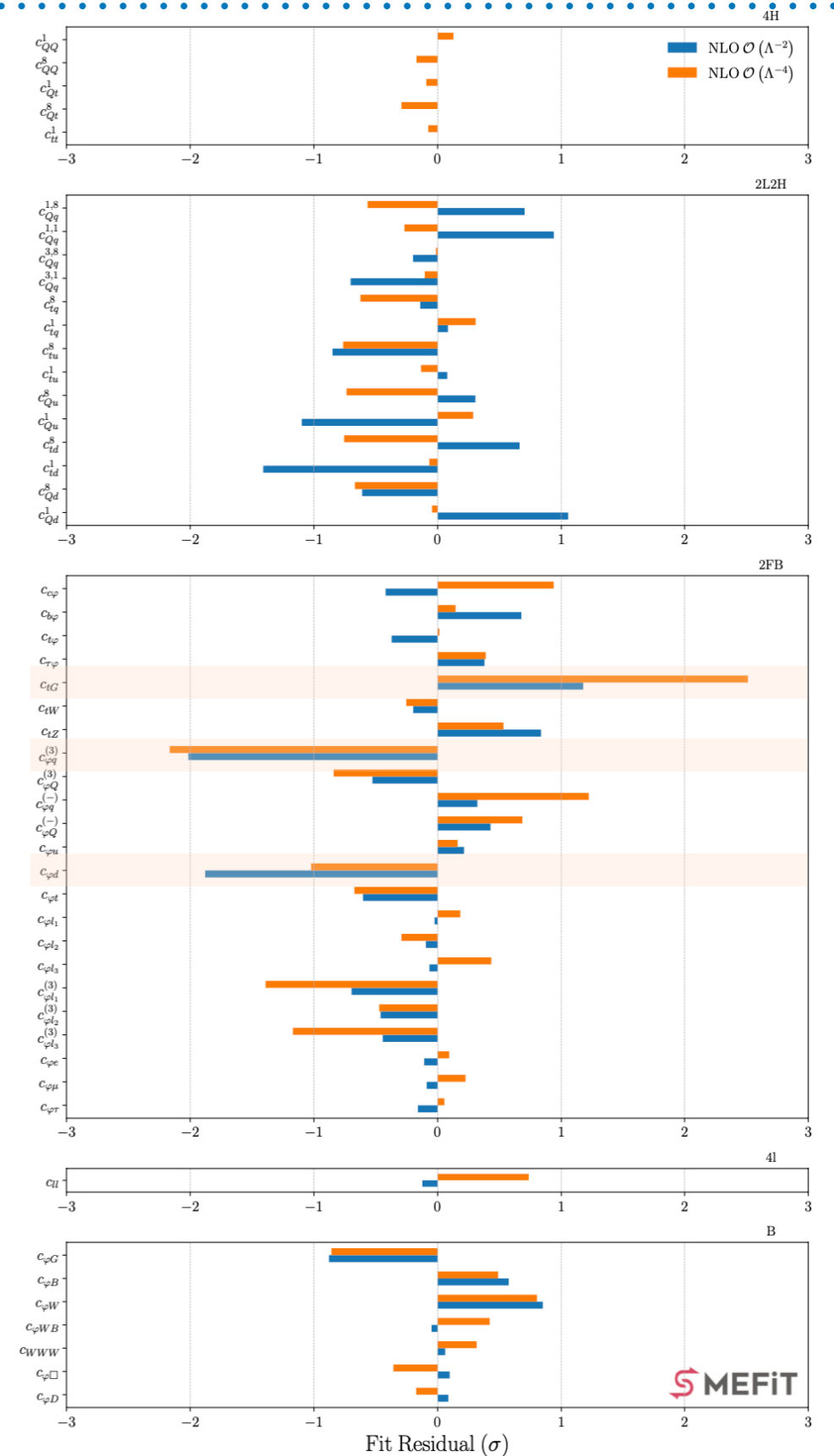
- Most bounds below 1 for $\Lambda = 1$ TeV
- Quadratic terms important (mostly **2-light-2-heavy**)
- Least constrained are **4-heavy operators**



Current global fit results

- Fit residuals (pulls) largely consistent with the SM

$$P_i \equiv 2 \left(\frac{\langle c_i \rangle - c_i^{(\text{SM})}}{[c_i^{\text{min}}, c_i^{\text{max}}]^{68\% \text{ CI}}} \right)$$

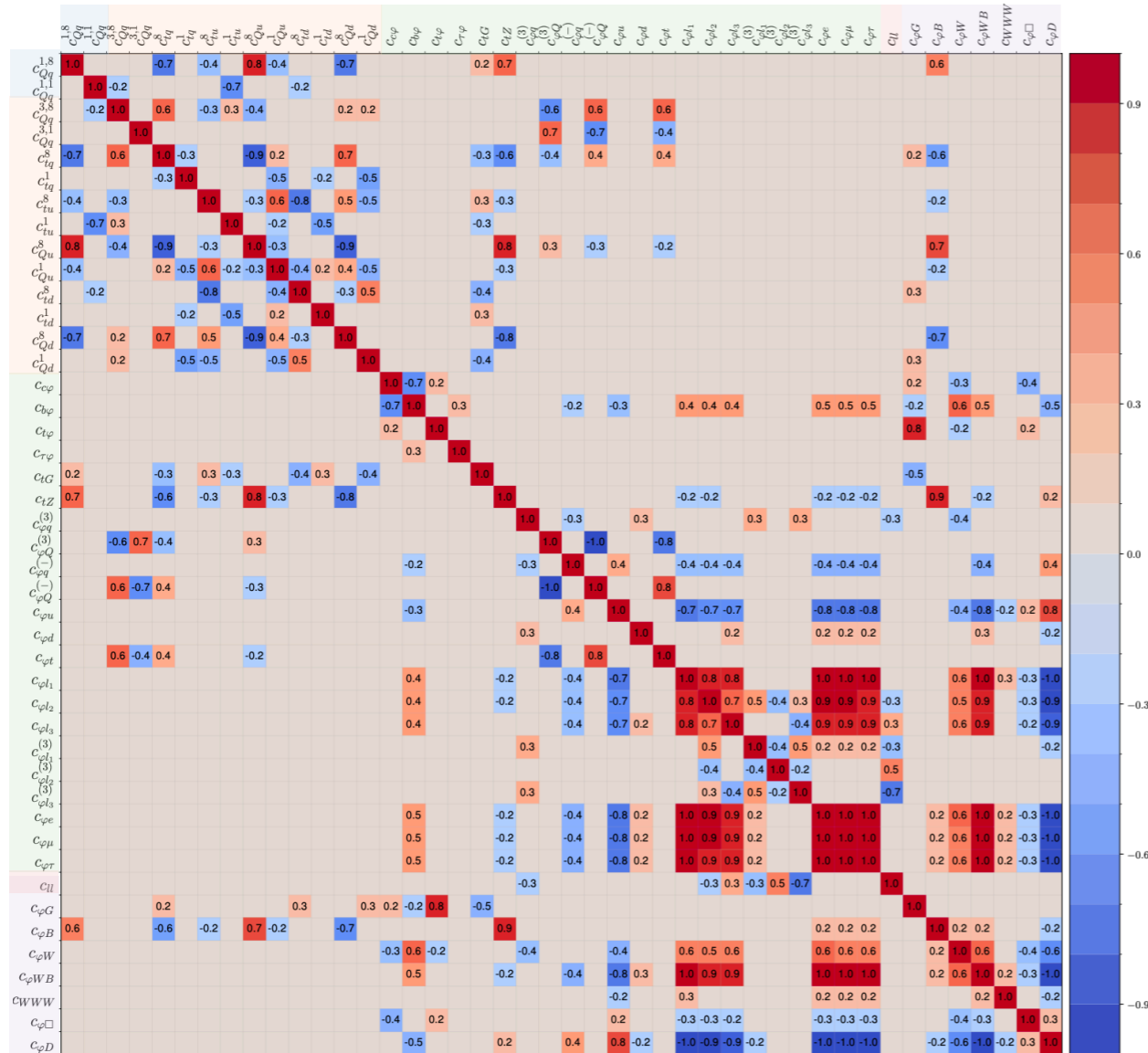


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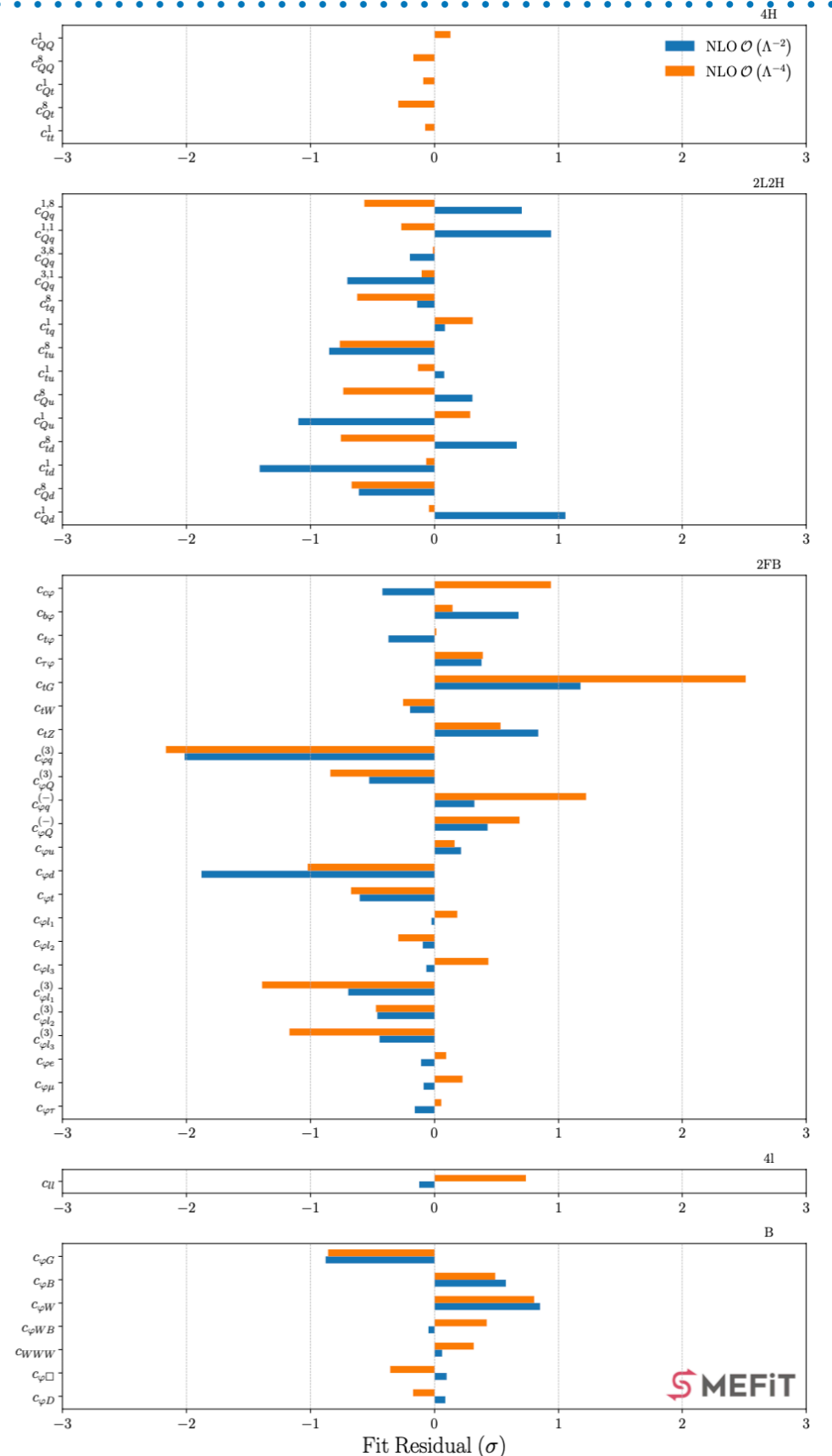
- Fit residuals (pulls) largely consistent with the SM

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- Large correlations in the linear fit are lifted in the quadratic fit



Correlation: NLO $\mathcal{O}(\Lambda^{-2})$

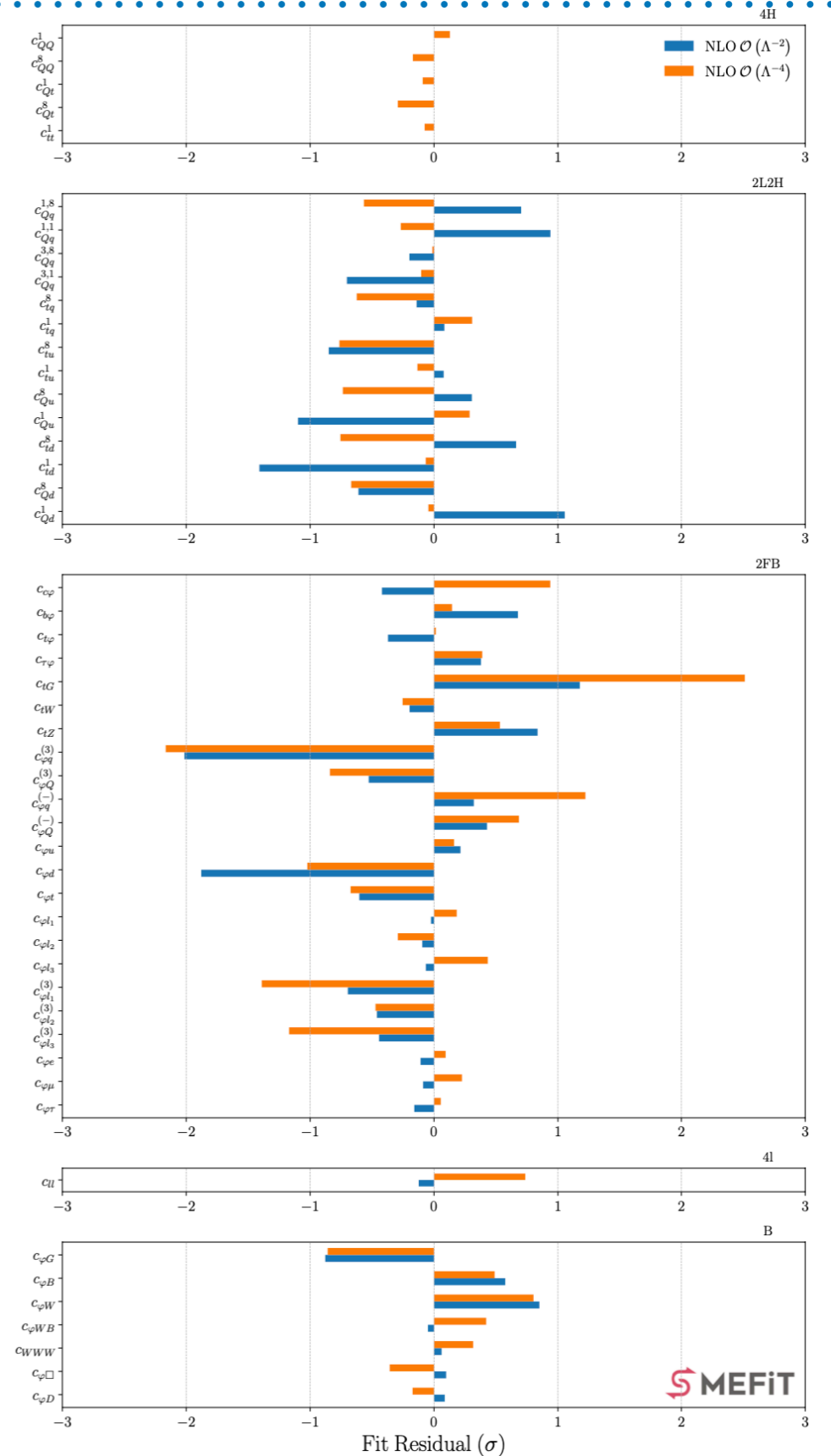
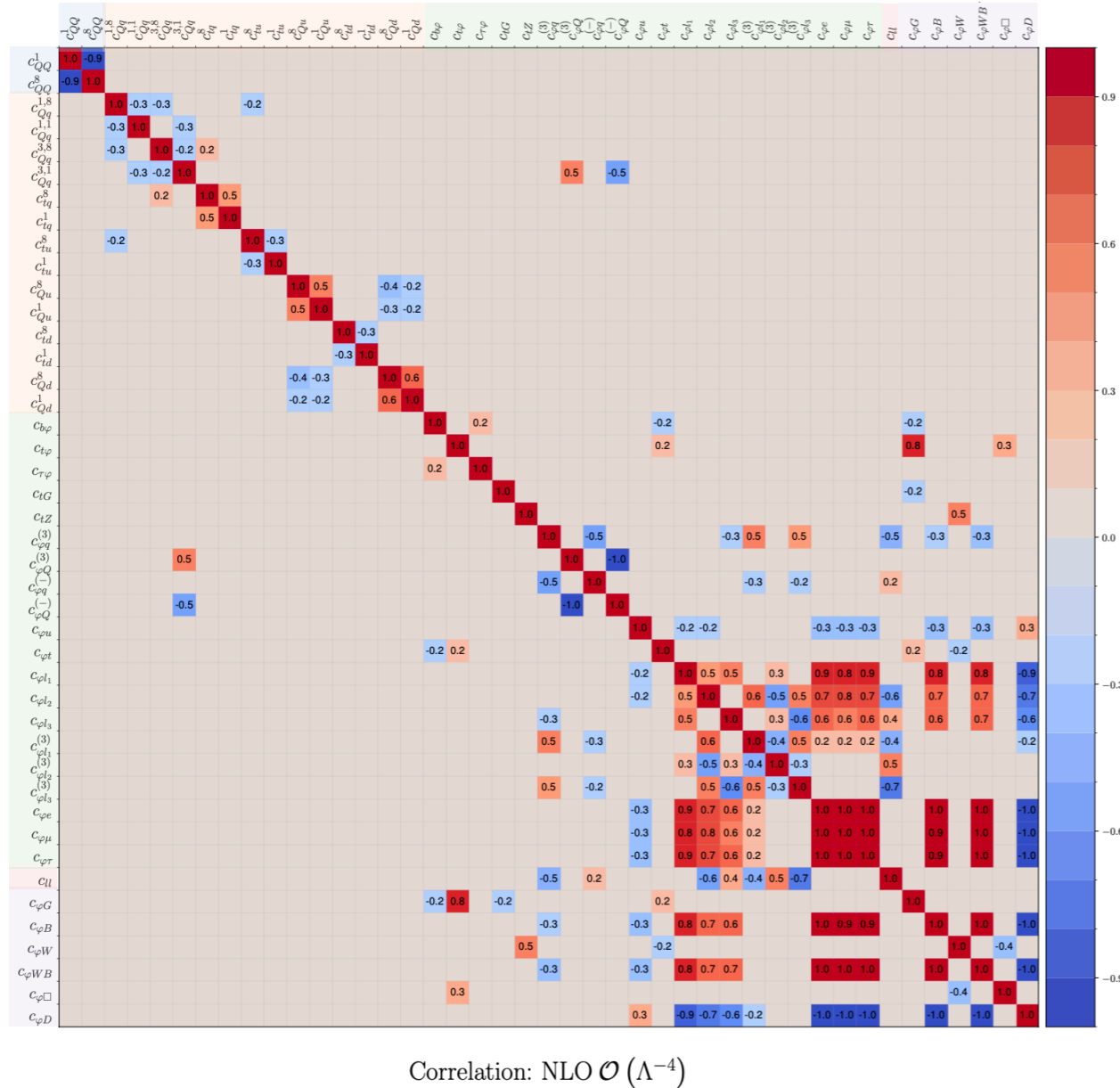


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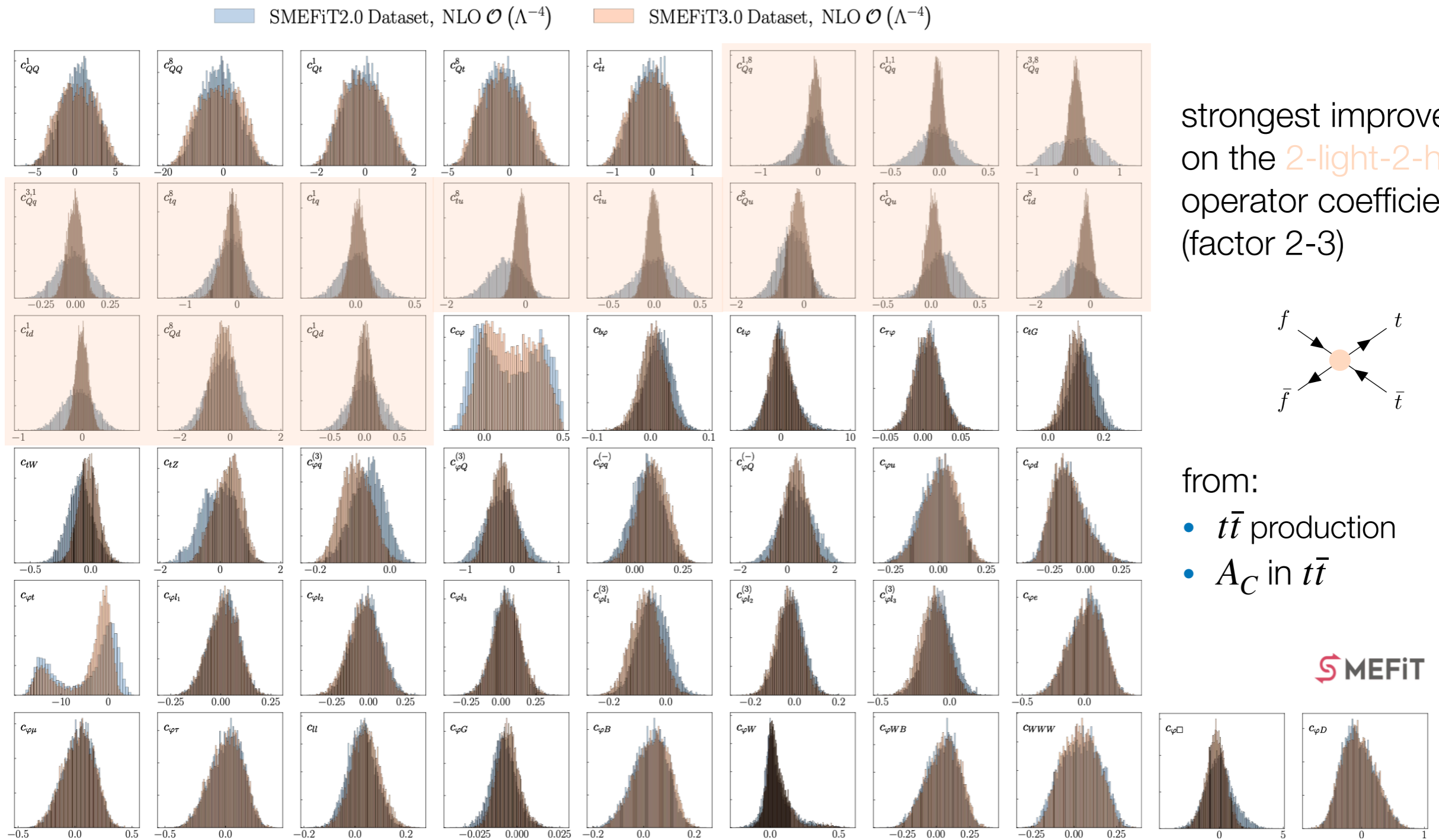
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- Large correlations in the linear fit are lifted in the **quadratic** fit



Impact of new datasets



Prospects at future colliders



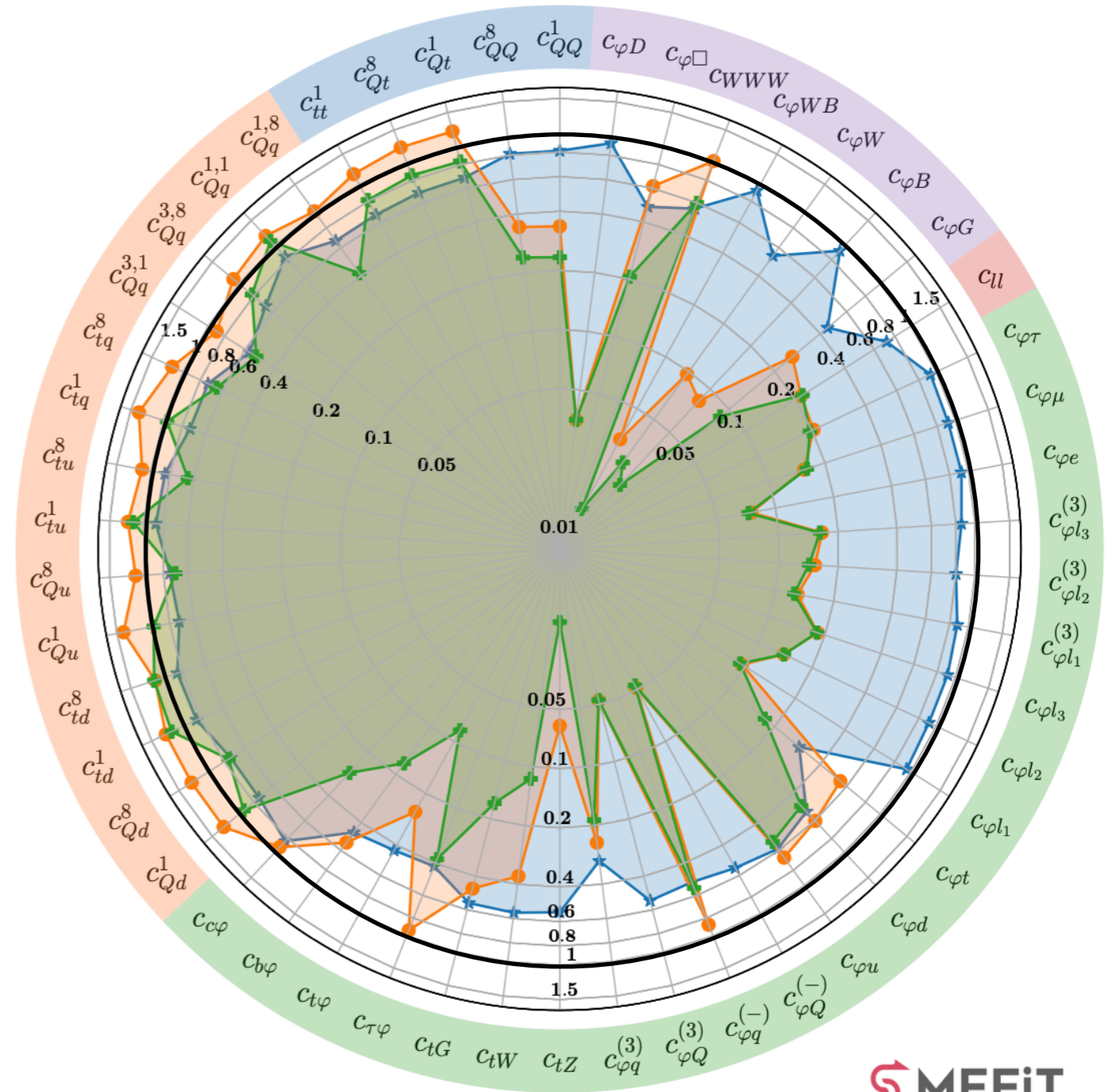
Implemented in SMEFiT3.0



HL-LHC results

$$R_{\delta c_i} = \frac{[c_i^{\min}, c_i^{\max}]^{95\% \text{ CL}} (\text{baseline} + \text{HL-LHC})}{[c_i^{\min}, c_i^{\max}]^{95\% \text{ CL}} (\text{baseline})}$$

Ratio of Uncertainties to SMEFiT3.0 Baseline, $\mathcal{O}(\Lambda^{-4})$, Marginalised



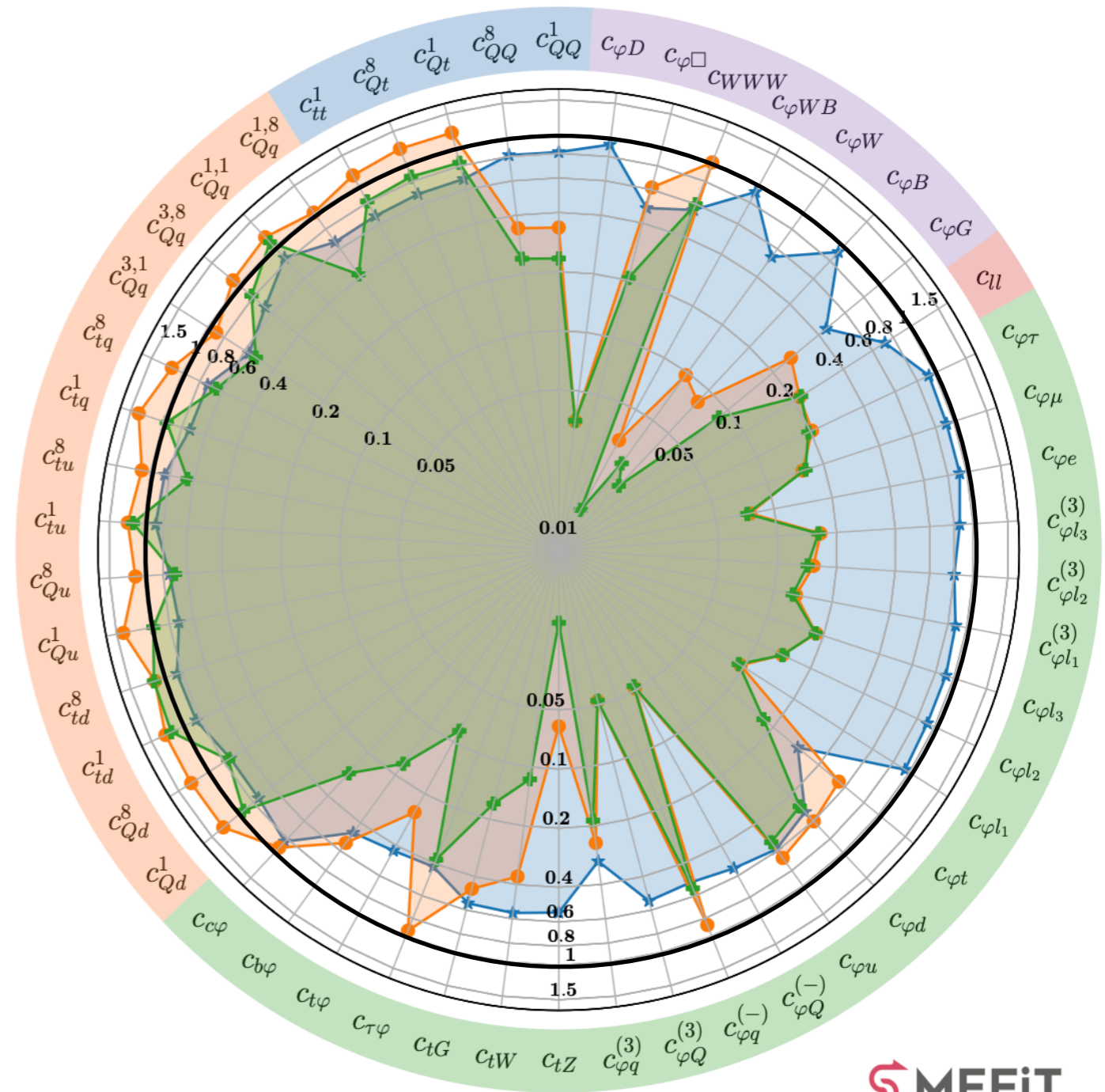
- ★ HL-LHC
- SMEFiT3.0, individual
- ✚ HL-LHC, individual
- Baseline: LHC, marginalised, quadratic

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- L1 projections of **Run II datasets**
 - pseudodata fluctuated around SM theory
 - statistics rescaled by luminosity
 - systematics reduced by a factor 2



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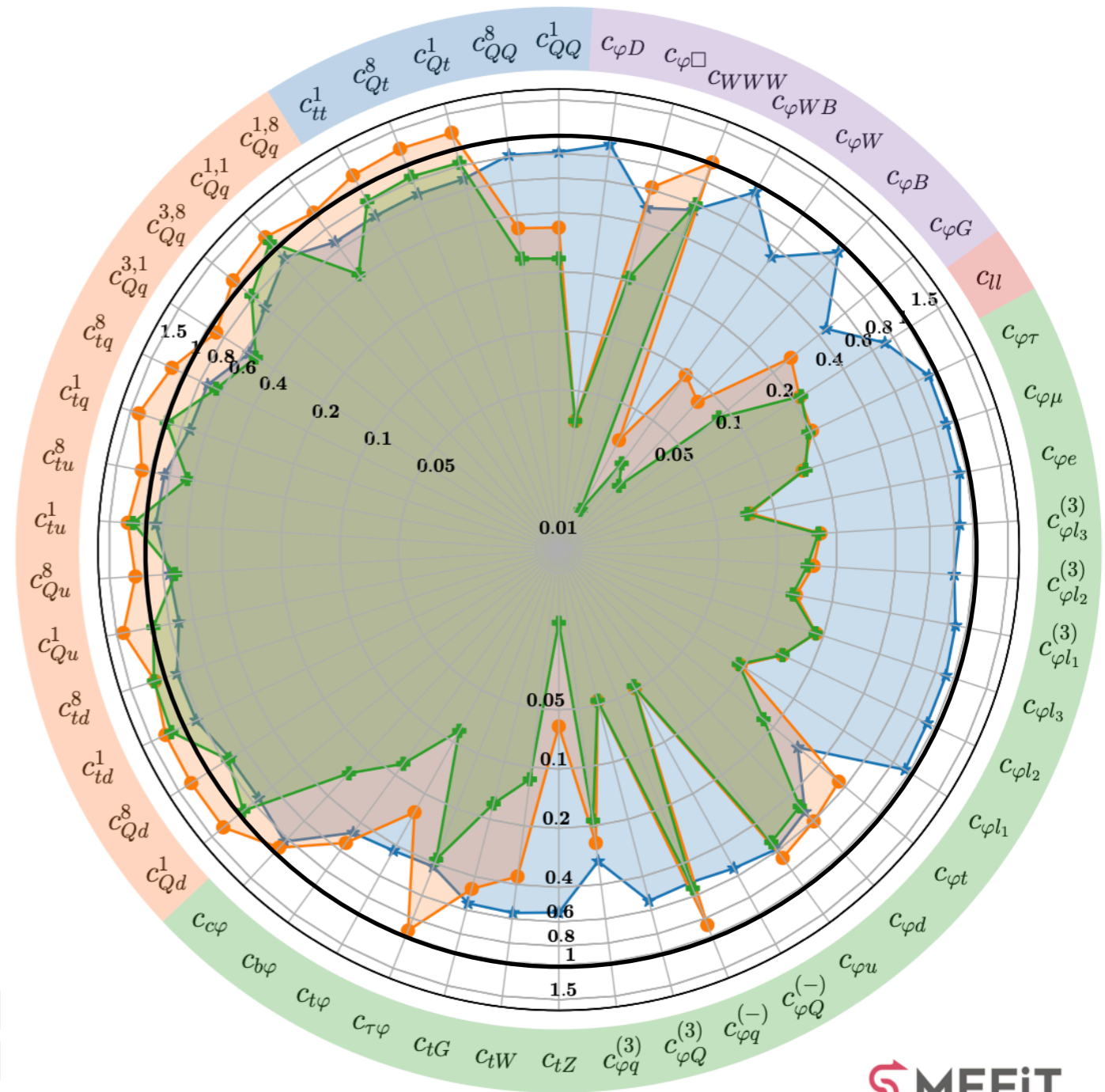
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- L1 projections of **Run II datasets**
 - pseudodata fluctuated around SM theory
 - statistics rescaled by luminosity
 - systematics reduced by a factor 2
- **marginalised bounds improve by a factor 1.2-3**
- individual bounds are overly optimistic
- **2-light-2-heavy** improved by 30% (further improvement of a factor 2 expected with a dedicated binning)

[arXiv:2206.08326]
[arXiv:2205.02140]

Ratio of Uncertainties to SMEFiT3.0 Baseline, $\mathcal{O}(\Lambda^{-4})$, Marginalised



- ★ HL-LHC
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Future circular lepton colliders



- Dataset input
 - EWPOs at the Z pole
 - light fermion pair production
 - Higgs production (Higgstrahlung and VBF)
 - diboson (WW) production
 - top pair production
- Uncertainty projections from Snowmass study updated with FCC midterm Feasibility Report

[[arXiv:2206.08326](https://arxiv.org/abs/2206.08326)]
[[CERN/3789/RA](https://cds.cern.ch/record/3789/RA)]

Energy (\sqrt{s})	\mathcal{L}_{int} (Run time)		$\mathcal{L}_{\text{FCC-ee}}/\mathcal{L}_{\text{CEPC}}$
	FCC-ee (4 IPs)	CEPC (2 IPs)	
91 GeV (Z -pole)	300 ab^{-1} (4 years)	100 ab^{-1} (2 years)	3
161 GeV ($2m_W$)	20 ab^{-1} (2 years)	6 ab^{-1} (1 year)	3.3
240 GeV	10 ab^{-1} (3 years)	20 ab^{-1} (10 years)	0.5
350 GeV	0.4 ab^{-1} (1 year)	0.2 ab^{-1}	2
365 GeV ($2m_t$)	3 ab^{-1} (4 years)	1 ab^{-1} (5 years)	3

FCC-ee results

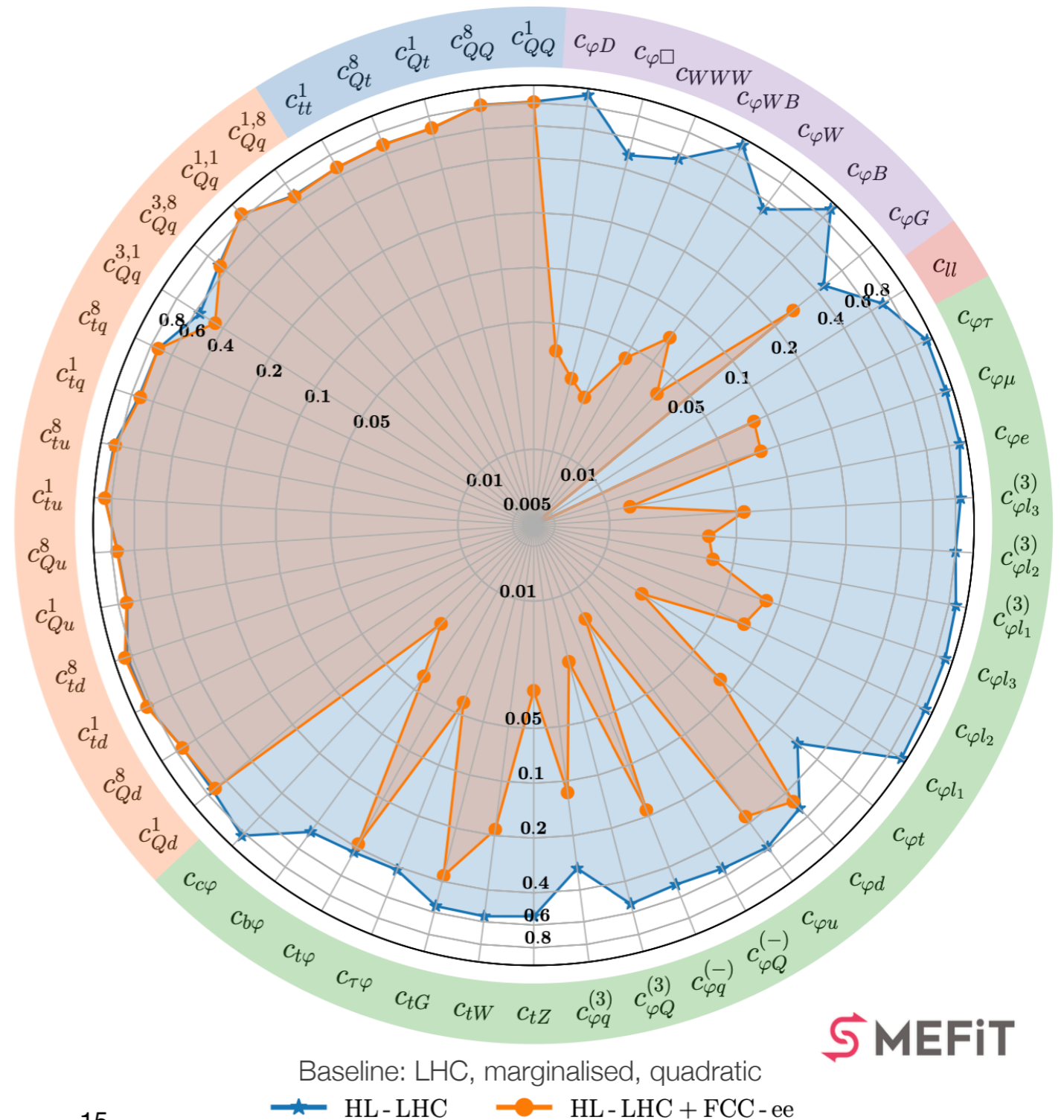
$$R_{\delta c_i} = \frac{[c_i^{\min}, c_i^{\max}]^{95\% \text{ CL}} (\text{baseline} + \text{HL-LHC})}{[c_i^{\min}, c_i^{\max}]^{95\% \text{ CL}} (\text{baseline})}$$

- gauge operators improve of up to a factor 30
- 2-fermion operators improve of up to a factor 50
- no sensitivity to 4-quark operators: improvement only from marginalisation

further improvement expected from one loop corrections

[see for instance [arXiv: 1809.03520](https://arxiv.org/abs/1809.03520),
[arXiv: 2409.11466](https://arxiv.org/abs/2409.11466)]

Ratio of Uncertainties to SMEFiT3.0 Baseline, $\mathcal{O}(\Lambda^{-4})$, Marginalised

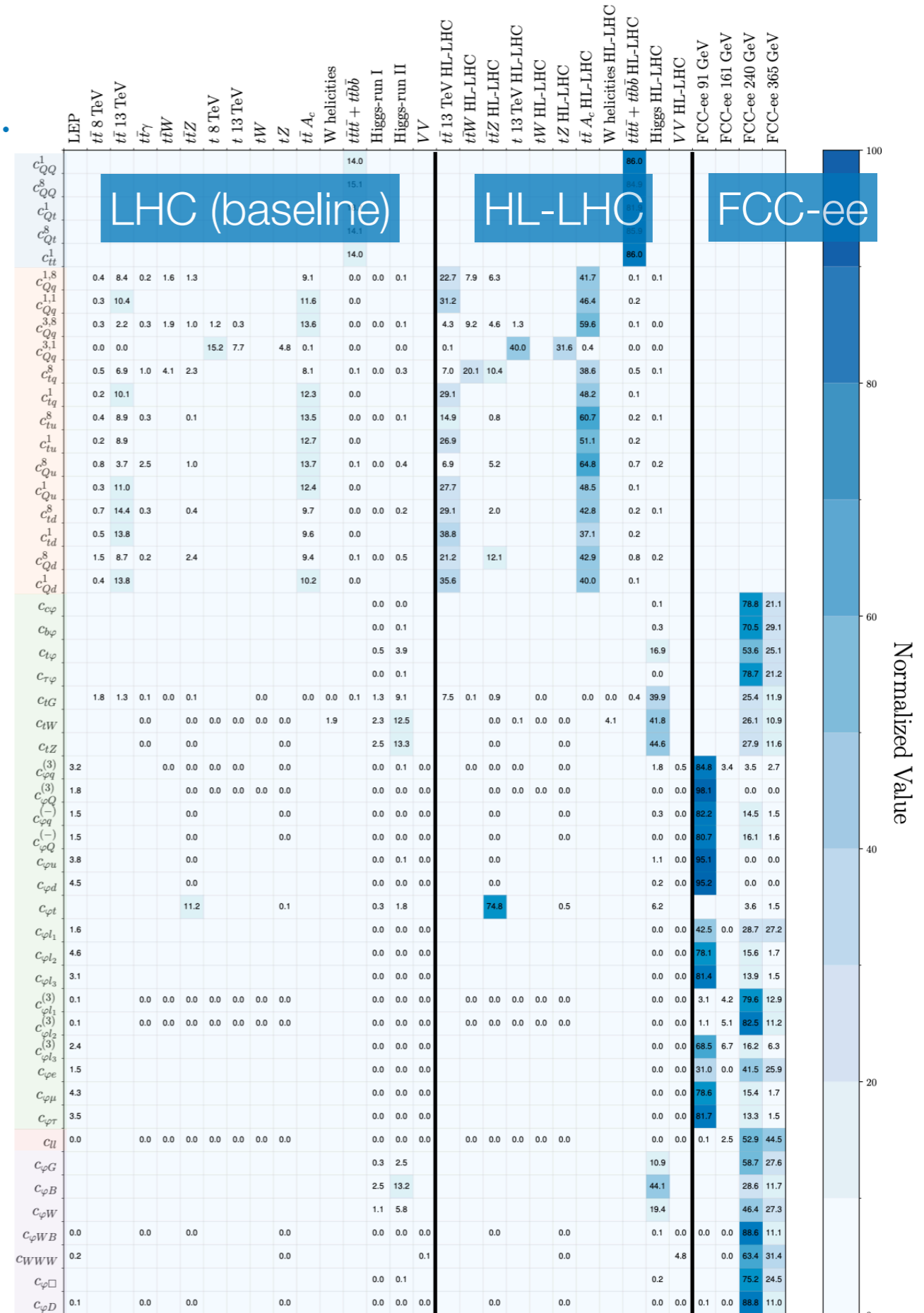


Fisher information


- The Fisher information matrix quantifies **which dataset is most sensitive** to a EFT parameter (each row is normalised by 100)

$$I_{ij} = \sum_{m=1}^{n_{\text{dat}}} \frac{\sigma_{m,i}^{(\text{eft})} \sigma_{m,j}^{(\text{eft})}}{\delta_{\text{exp},m}^2}, \quad i, j = 1, \dots, n_{\text{eft}}$$

- neglects correlations and quadratic effects
- HL-LHC dominates in the **2-light-2-heavy** and **4-heavy** sectors
- FCC-ee dominates in the **bosonic** and **2-fermion** sector
- FCC-ee run at 161 GeV is the less useful for SMEFT



Conclusions & outlook

- SMEFT is a consistent way to look for new interactions
- Global fits need to combine all the processes available at LHC
- **SMEFIT3.0 in the biggest global SMEFT analysis to date: 50 Wilson coefficients and 445 datapoints**
- **Significant improvement** in New Physics reach **expected at HL-LHC and at future circular lepton colliders**
- **Extension to other future colliders** (ILC, CLIC, muon collider) and **RG running** in progress
- **Code, data and theory are public** and available at  lhcfithef.github.io/smeFIT_release/

Backup

EWPOs exact implementation

- In the SMEFT, Z and W fermionic coupling receive contributions from dim-6 operators

- SMEFIT2.0:** assume LEP/SLD measurements precise enough to set coupling shift to 0

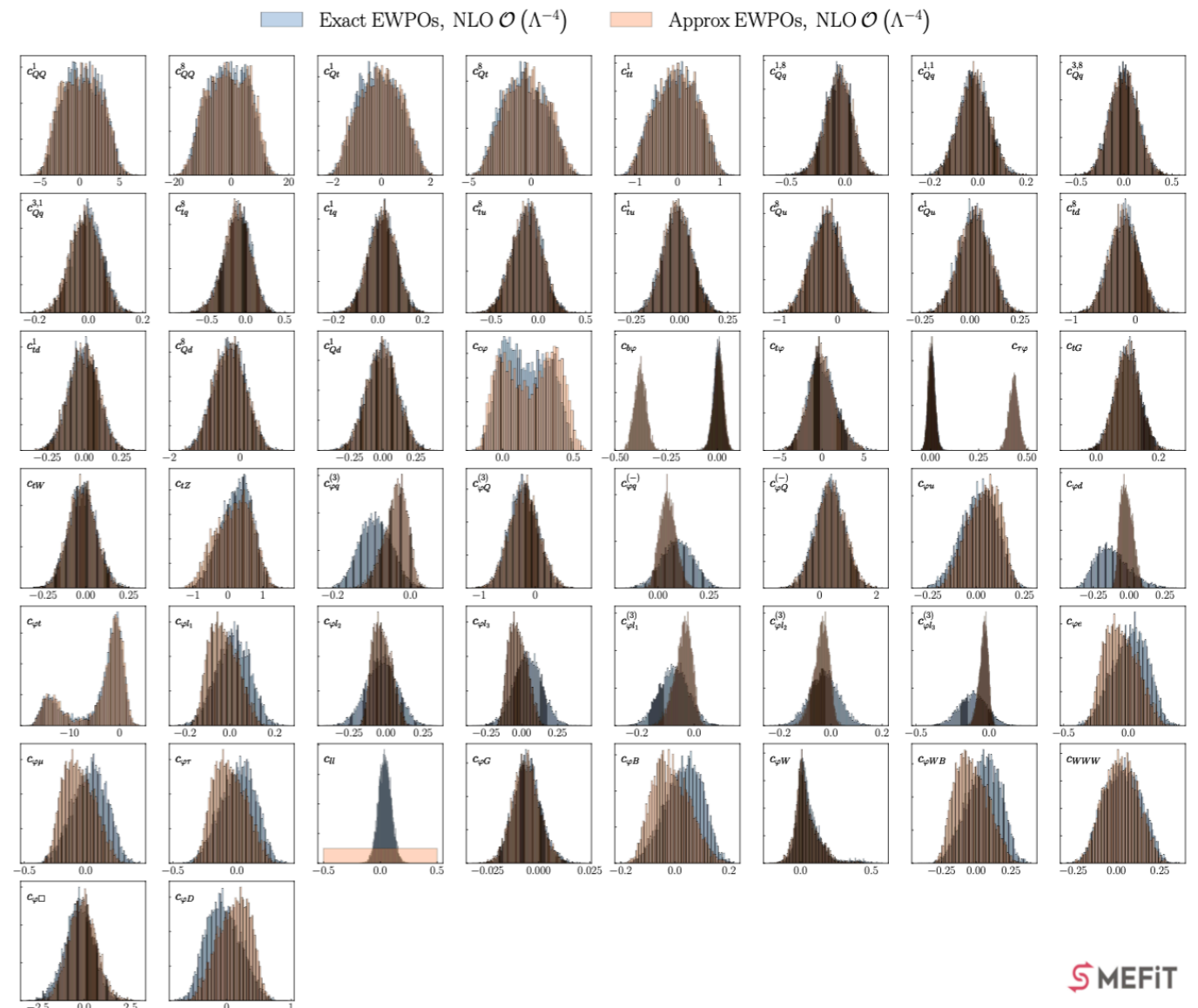
➔ only 2 independent Wilson coefficients

$$\begin{pmatrix} c_{\phi l_i}^{(3)} \\ c_{\phi l_i}^{(1)} \\ c_{\phi e/\mu/\tau}^{(-)} \\ c_{\phi q}^{(-)} \\ c_{\phi q}^{(3)} \\ c_{\phi u} \\ c_{\phi d} \\ c_{ll} \end{pmatrix} = \begin{pmatrix} -\frac{1}{t_W} & -\frac{1}{4t_W^2} \\ 0 & -\frac{1}{4} \\ 0 & -\frac{1}{2} \\ \frac{1}{t_W} & \frac{1}{4s_W^2} - \frac{1}{6} \\ -\frac{1}{t_W} & -\frac{1}{4t_W^2} \\ 0 & \frac{1}{3} \\ 0 & -\frac{1}{6} \\ 0 & 0 \end{pmatrix} \begin{pmatrix} c_{\phi WB} \\ c_{\phi D} \end{pmatrix}$$

- SMEFIT3.0:** EWPOs implemented like the rest of LHC data

➔ 16 independent Wilson coefficients

no major difference, but the approximate implementation resulted in too stringent constraints on some EW operators



New LHC datasets

Dataset	\sqrt{s} (TeV)	\mathcal{L} (fb $^{-1}$)	Info	Observables	n_{dat}	ref.
ATLAS-STXS-RunII-13TeV-2022	13	139	$ggF, \text{VBF}, Vh, t\bar{t}h, th$	$d\sigma/dp_T^h$ $d\sigma/dm_{jj}$ $d\sigma/dp_T^V$	36	[82]
CMS-WZ-pTZ-13TeV-2022	13	137	WZ, fully leptonic	$1/\sigma d\sigma/dp_T^Z$	10	[83]
CMS-tt-13TeV-ljets_inc	13	137	$\ell + \text{jets}$	$\sigma(t\bar{t})$	1	[84]
CMS-tt-13TeV-Mtt	13	137	$\ell + \text{jets}$	$1/\sigma d\sigma/dm_{t\bar{t}}$	14	[84]
CMS-tt-13TeV-asy	13	138	$\ell + \text{jets}$	A_C	3	[85]
ATLAS-tt-13TeV-asy-2022	13	139	$\ell + \text{jets}$	A_C	5	[86]
ATLAS-Whel-13TeV	13	139	W-helicity fraction	F_0, F_L	2	[87]
ATLAS-ttZ-13TeV-pTZ	13	139	$t\bar{t}Z$	$d\sigma/dp_T^Z$	7	[88]
ATLAS-tta-8TeV	8	20.2	Inclusive	$\sigma(t\bar{t}\gamma)$	1	[89]
CMS-tta-8TeV	8	19.7	Inclusive	$\sigma(t\bar{t}\gamma)$	1	[90]
ATLAS-tttt-13TeV-slep_inc	13	139	single-lepton	$\sigma_{\text{tot}}(t\bar{t}t\bar{t})$	1	[91]
CMS-tttt-13TeV-slep_inc	13	35.8	single-lepton	$\sigma_{\text{tot}}(t\bar{t}t\bar{t})$	1	[92]
ATLAS-tttt-13TeV-2023	13	139	multi-lepton	$\sigma_{\text{tot}}(t\bar{t}t\bar{t})$	1	[93]
CMS-tttt-13TeV-2023	13	139	same-sign or multi-lepton	$\sigma_{\text{tot}}(t\bar{t}t\bar{t})$	1	[94]
CMS-ttbb-13TeV-dilepton_inc	13	35.9	dilepton	$\sigma_{\text{tot}}(t\bar{t}b\bar{b})$	1	[95]
CMS-ttbb-13TeV-ljets_inc	13	35.9	$\ell + \text{jets}$	$\sigma_{\text{tot}}(t\bar{t}b\bar{b})$	1	[95]
ATLAS-t_sch-13TeV_inc	13	139	s-channel	$\sigma_{\text{tot}}(t + \bar{t})$	1	[96]
CMS-tZ-13TeV-pTt	13	138	dilepton	$d\sigma_{\text{fid}}(tZj)/dp_T^t$	3	[97]
CMS-tW-13TeV-slep_inc	13	36	single-lepton	$\sigma_{\text{tot}}(tW)$	1	[98]

L0 vs L1 projections

- **Level-0:** no fluctuation of the pseudo-data

$$\mathcal{O}_i^{(\text{exp})} = \mathcal{O}_i^{(\text{th})}, \quad i = 1, \dots, n_{\text{bin}}$$

- **Level-1:**

- pseudodata fluctuated around SM

$$\mathcal{O}_i^{(\text{exp})} = \mathcal{O}_i^{(\text{th})} \left(1 + r_i \delta_i^{(\text{stat})} + \sum_{k=1}^{n_{\text{sys}}} r_{k,i} \delta_{k,i}^{(\text{sys})} \right), \quad i = 1, \dots, n_{\text{bin}}$$

- statistics rescaled by luminosity

$$\delta_i^{(\text{stat})} = \tilde{\delta}_i^{(\text{stat})} \sqrt{\frac{\mathcal{L}_{\text{Run2}}}{\mathcal{L}_{\text{HLLHC}}}}$$

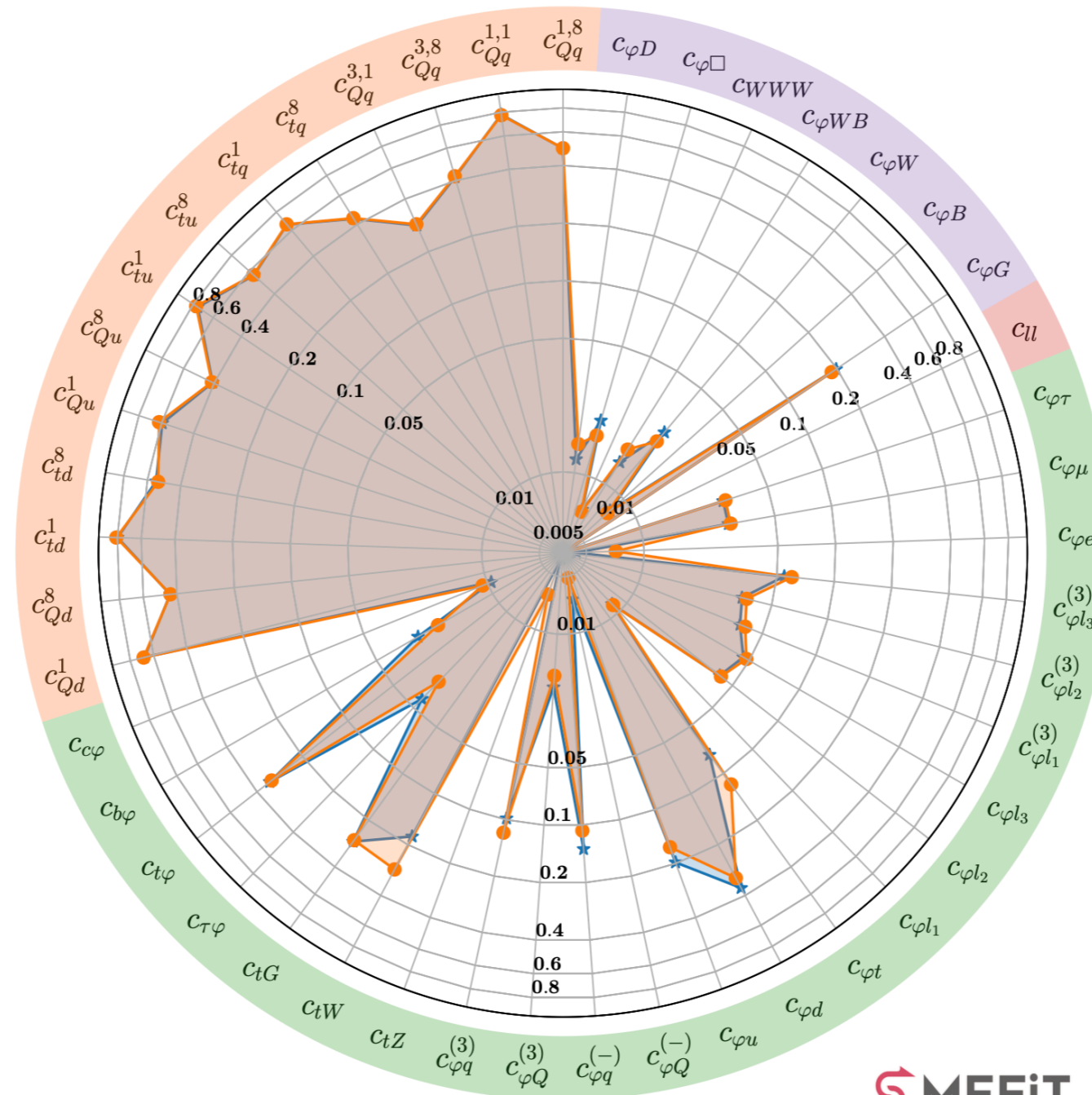
- systematics reduced by a factor $f_{\text{red}}^{(k)} = 1/2$

$$\delta_{k,i}^{(\text{sys})} = \tilde{\delta}_{k,i}^{(\text{sys})} \times f_{\text{red}}^{(k)}, \quad i = 1, \dots, n_{\text{bin}}, \quad k = 1, \dots, n_{\text{sys}}$$

Good agreement found between the two approaches.

FCC-ee and CEPC

Ratio of Uncertainties to SMEFIT3.0 Baseline, $\mathcal{O} (\Lambda^{-2})$, Marginalised



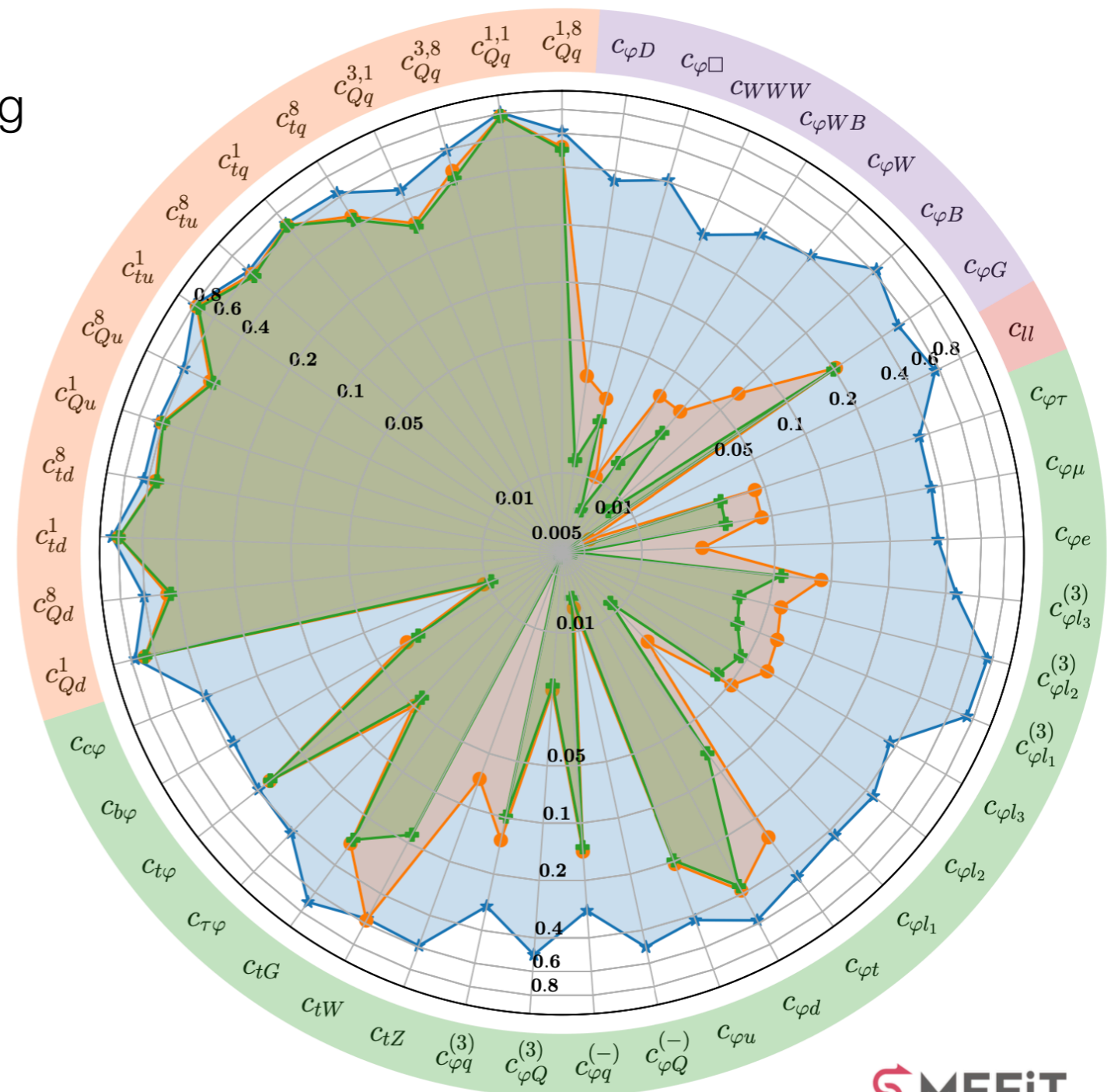
★ HL - LHC + FCC - ee ● HL - LHC + CEPC

FCC-ee energy breakdown

Ratio of Uncertainties to SMEFiT3.0 Baseline, $\mathcal{O}(\Lambda^{-2})$, Marginalised

Study the impact of sequentially adding different runs

- largest impact from 91 + 240 GeV combination
- 365 GeV also relevant



- HL - LHC + FCC - ee (91 GeV)
- HL - LHC + FCC - ee (91 + 240 GeV)
- HL - LHC + FCC - ee (91 + 161 + 240 + 365 GeV)

