



# Quantum sensing with ultracold atoms in phase modulated optical lattices

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3 April 2024

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Infleqtion



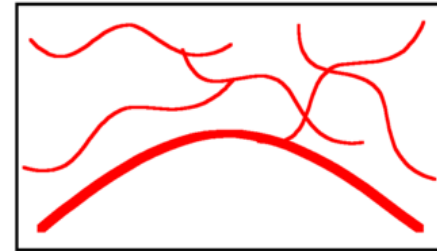
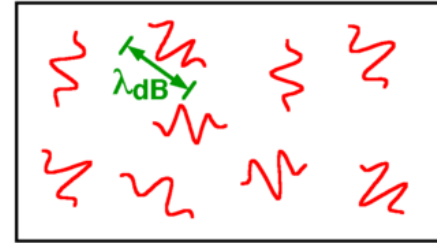
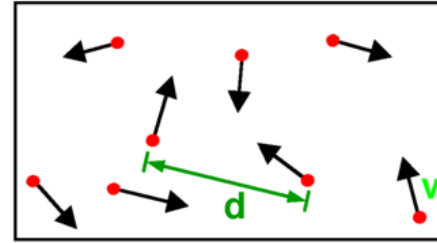
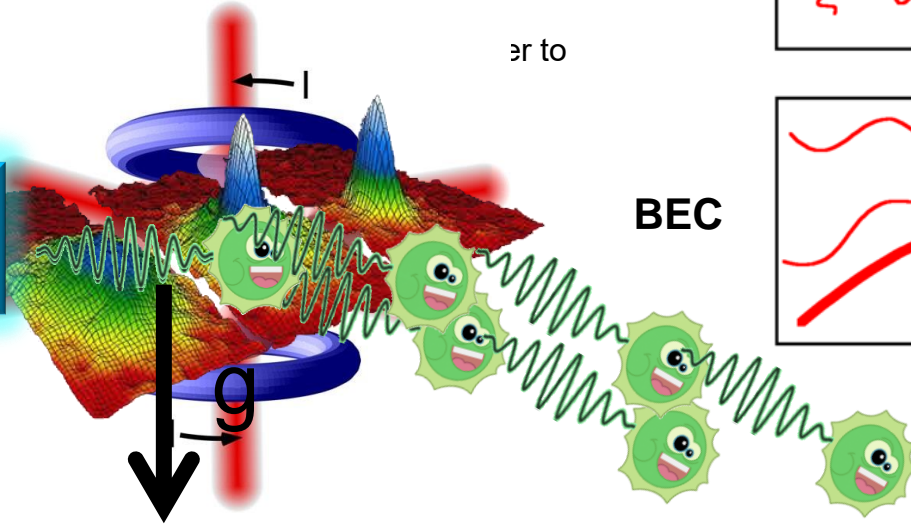
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# The preliminaries

# Making bosonic atoms very cold

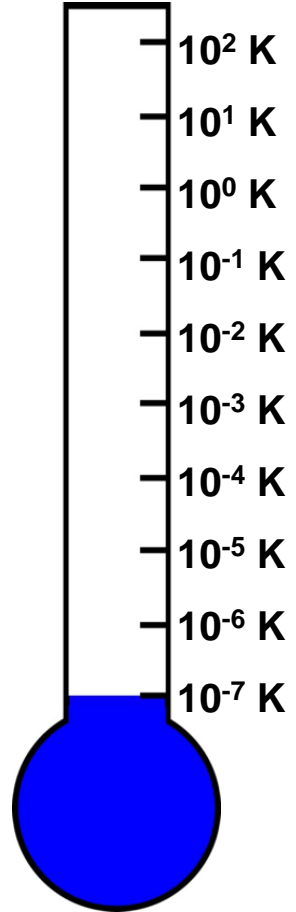
- If you get bosonic atoms cold enough ( $\approx 100$  nK), their deBroglie wavelength is on the order of the inter-particle spacing
  - This is known as a Bose-Einstein condensate, or BEC.
  - The atoms become mutually coherent, like the photons:
  - Our “atom control”

**ATOM LASER**



**MOT**

**BEC**



# Trapping atoms with light

- Light induces a dipole moment in an atom

$$\vec{d} = \alpha \vec{E}$$

- This gives rise to a force and potential

$$\vec{F} = -\nabla U \propto -\alpha(\omega) I(\vec{r})$$

- For red-detuned light ( $\Delta = \omega_{light} - \omega_{atom} < 0$ ) this potential is attractive, and the atoms move towards the intensity maxima.

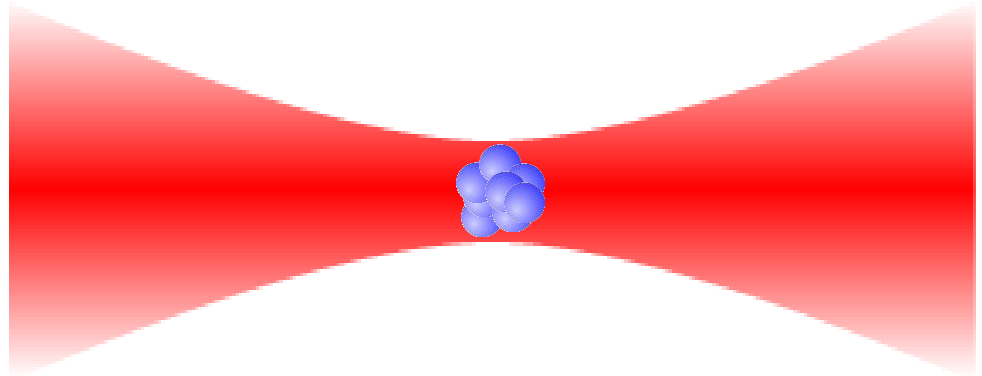
- Depth

$$U \propto I(\vec{r})/\Delta$$

- Scattering rate

$$\Gamma_{sc} \propto I(\vec{r})/\Delta^2$$

- We want a lot of power from a laser far-detuned from resonance!



# The optical lattice: an egg carton for atoms

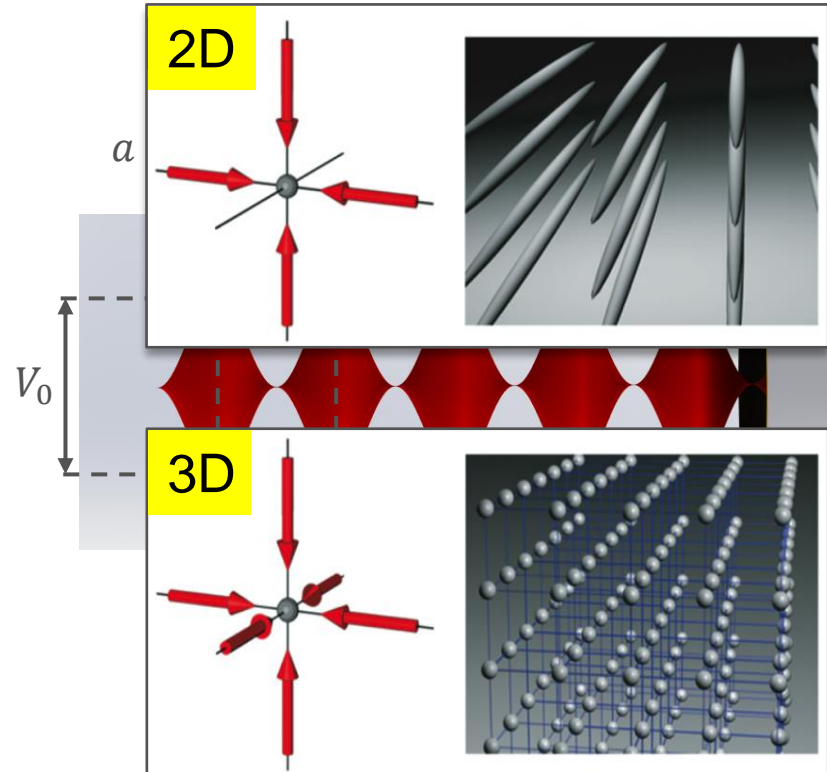
- Reflect a dipole laser back on itself to create a sinusoidally-varying potential

$$V(x) = V_0 \cos(2k_L x)$$

- Depth typically expressed in recoils

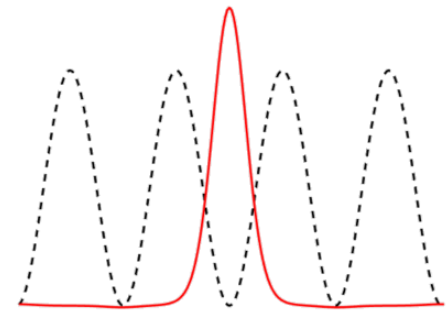
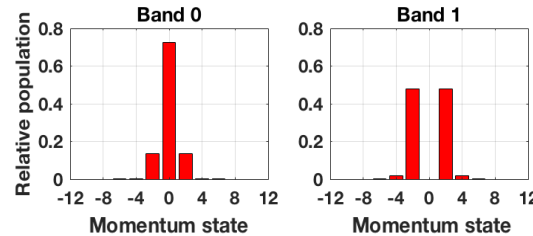
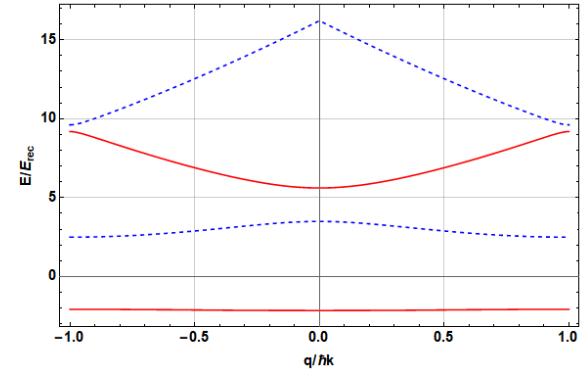
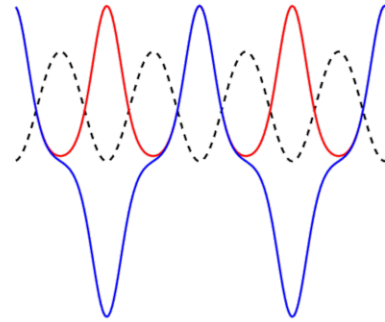
$$V_0 = sE_R = s \frac{\hbar^2 k_L^2}{2m}$$

- Can work in (up to) three dimensions!



# How do we describe the atom wavefunction in a lattice?

- Two (equivalent) bases are commonly used
- Bloch functions
  - Atoms delocalized in position, localized in momentum space
  - Gives rise to band structure within a Brillouin zone
- Wannier functions
  - Atoms localized in position space (to a single lattice site), delocalized in momentum space
  - Composed of sums of Bloch functions in a given band
- Localized or delocalized? It depends on the lattice depth (and the problem).
  - Deeper lattices: more localized atoms
  - SLI uses shallow lattices—**we control the momentum states of the atoms!**



# Inertial sensing with ultracold atoms trapped in phase-modulated optical lattices [PRL **120**, 263201, (2018)]

## Experimental Demonstration of Shaken-Lattice Interferometry

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(Received 28 January 2018; published 27 June 2018)

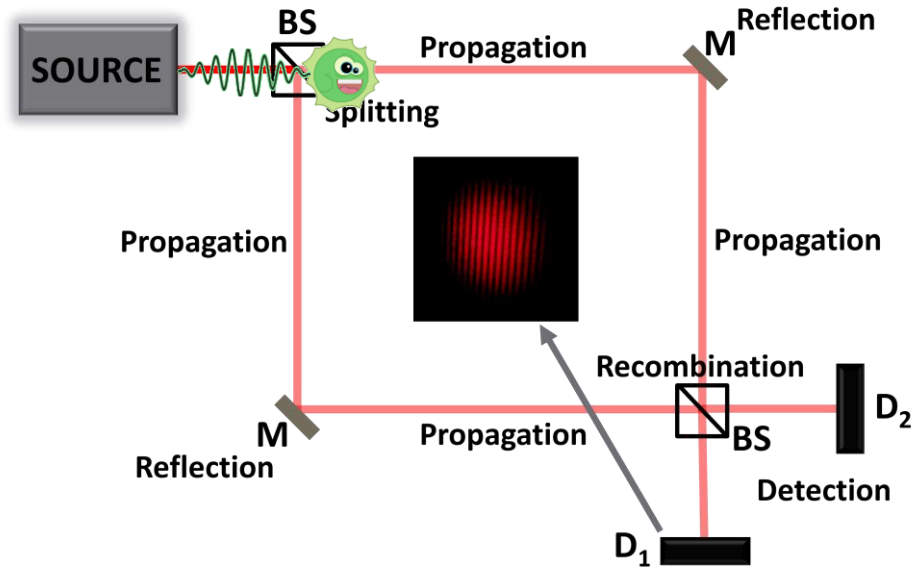
# Shaken lattice interferometry: building a sensor with atoms in optical lattices

The recipe:

- Take your favourite atom, and make it very cold
- Load it into the ground state of a **shallow** optical lattice potential
- Modulate the lattice to implement the atom-optical elements of an interferometer

$$V(x, t) = V_0 \cos(2kx + \phi(t))$$

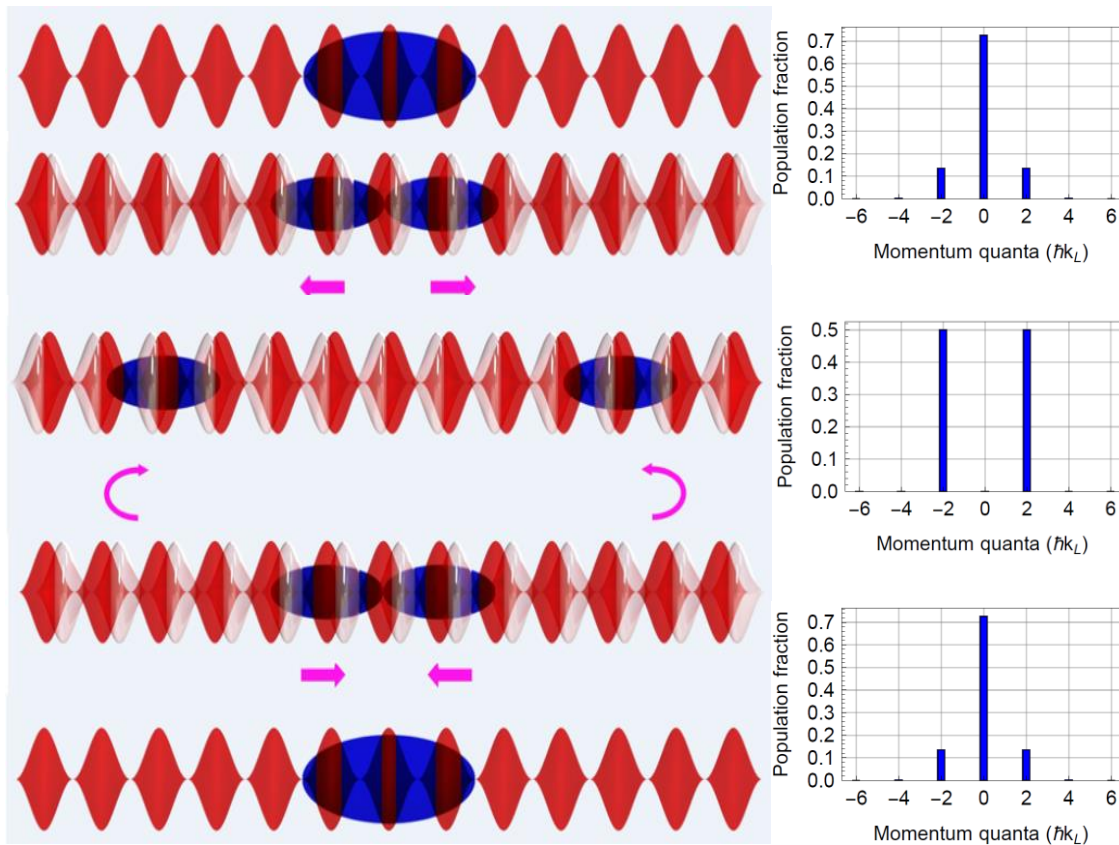
What we control!





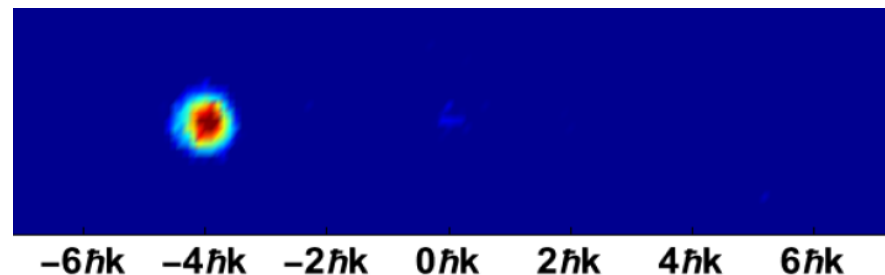
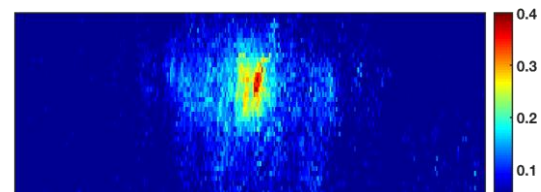
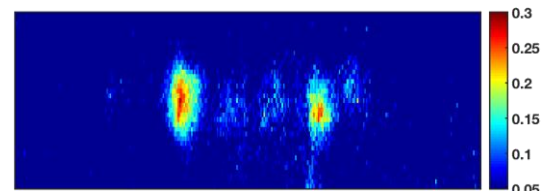
# Building a shaken lattice interferometer

- Work in the Bloch basis: atoms delocalized in position, localized in momentum
- Starting with atoms in the ground state of the lattice potential, we implement:
  - Splitting
  - Propagation
  - Reflection
  - Reverse propagation
  - Recombination back into the ground state
- The best shaking function  $\phi(t)$  is determined via optimal control



# Building a shaken lattice interferometer

- Measurement: relative population in the atoms' momentum states
  - Define a vector  $\vec{P}$  with elements  $\{P_n\}$  containing the relative population in the  $2n\hbar k$  state
  - We do not have access to phase information!
- Once the shaking function is known, it is fixed.
  - Can then calibrate the system's response to a signal (acceleration  $a$ )
  - Scale sensitivity by changing the total interrogation (shaking) time  $T$



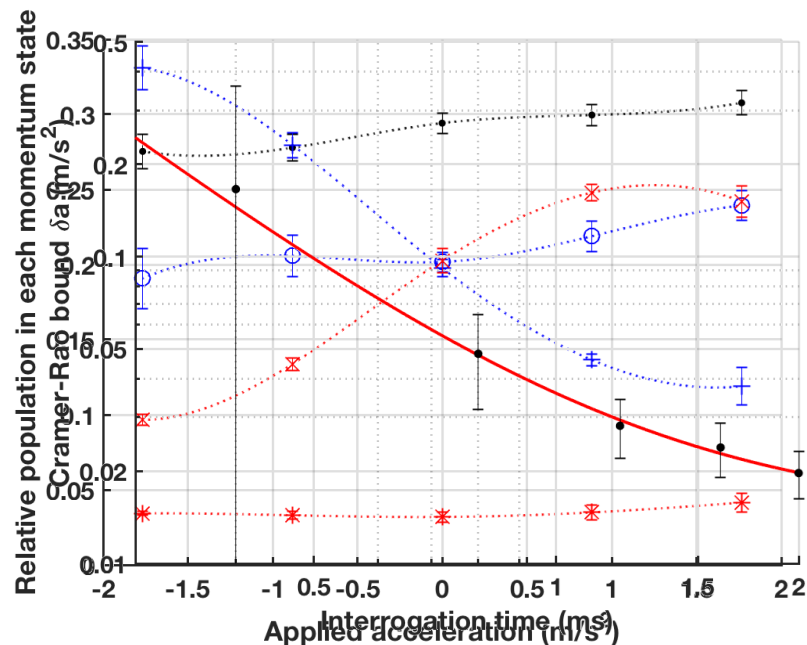
# But is it a sensor? Adding a signal

- We determine a signal by measuring how the atom momentum populations change with the applied signal
- The magnitude and direction of a signal is easily determined here, due to symmetry breaking as the lattice begins to shake
- Use the classical Fisher information  $F_c$  to define a minimum detectable acceleration  $\delta a = 1/\sqrt{F_c}$  given the momentum population vectors  $\vec{P}$  that we measure.

- CFI:

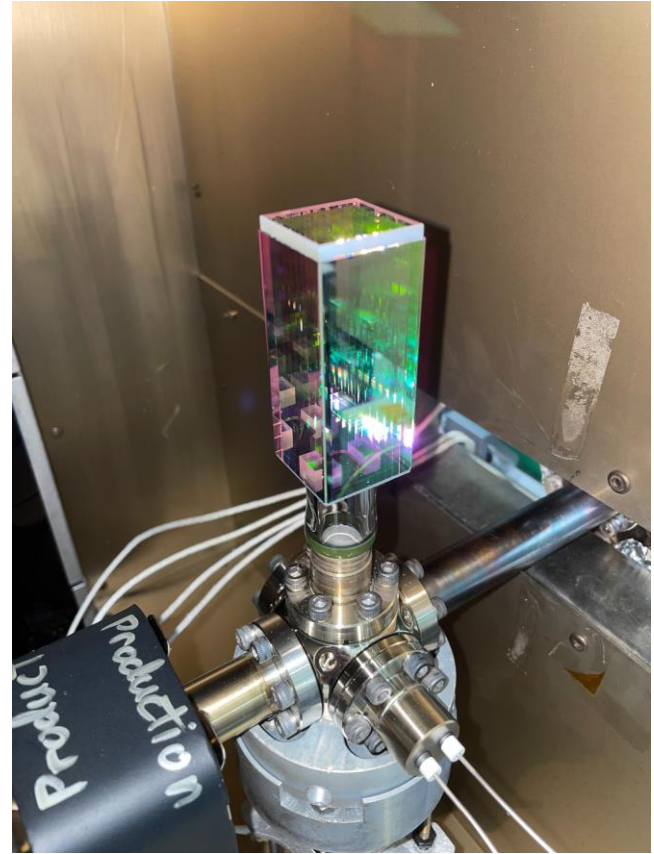
$$F_c(a) = N_{at} \sum_{n=-N}^N \frac{(\partial P_{a,n}/\partial a)^2}{P_{a,n}}$$

- Use this to find how  $\delta a$  scales with  $T$
- Simulations (experiments) give  $n = 2.21 \pm 0.31$  ( $1.96 \pm 0.13$ ) consistent with typical atom interferometers where  $n = 2$ .



# So what's next?

- Build a 3D lattice system in Bristol
- Demonstrate a multi-axis inertial sensor (3 axes of acceleration, 3 axes of rotation)



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- Open question #1: What is the best scaling with  $T$  that we can get?

## Abstract

Motivated by recent results using shaken optical lattices to perform atom interferometry, we explore the splitting of an atom cloud trapped in a phase-modulated ('shaken') optical lattice. Using a simple analytic model we are able to show that we can obtain the simplest case of  $\pm 2\hbar k_L$  splitting via single-frequency shaking. This is confirmed both via simulation and experiment. Furthermore, we are able to split with a relative phase  $\theta$  between the two split arms of 0 or  $\pi$  depending on our shaking frequency. Addressing higher-order splitting, we determine that  $\pm 6\hbar k_L$  splitting is sufficient to be able to accelerate the atoms in counterpropagating lattices. Finally, we show that we can use a genetic algorithm to optimize  $\pm 4\hbar k_L$  and  $\pm 6\hbar k_L$  splitting to within  $\approx 0.1\%$  by restricting our optimization to the resonance frequencies corresponding to single- and two-photon transitions between Bloch bands. As a proof-of-principle, an experimental demonstration of simplified optimization of  $4\hbar k_L$  splitting is presented.

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### PAPER

Simplified landscapes for optimization of shaken lattice interferometry

C A Weidner and D Z Anderson<sup>1</sup>

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- Open question #1: What is the best scaling with  $T$  that we can get?
- Open question #2: How robust is this method in the real world?

## Statistically characterizing robustness and fidelity of quantum controls and quantum control algorithms

Irtaza Khalid,<sup>1,\*</sup> Carrie A. Weidner<sup>2,\*</sup> Edmond A. Jonckheere,<sup>3,†</sup> Sophie G. Shermer,<sup>4,§</sup> and Frank C. Langbein<sup>1,‡</sup>

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IEEE CONTROL SYSTEMS LETTERS, VOL. 7, 2023

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## Analyzing and Unifying Robustness Measures for Excitation Transfer Control in Spin Networks

S. P. O'Neil<sup>Ⓞ</sup>, Member, IEEE, I. Khalid, A. A. Rompokos, Student Member, IEEE, C. A. Weidner<sup>Ⓞ</sup>, Member, IEEE, F. C. Langbein<sup>Ⓞ</sup>, Member, IEEE, S. Schirmer<sup>Ⓞ</sup>, Member, IEEE, and E. A. Jonckheere<sup>Ⓞ</sup>, Life Fellow, IEEE



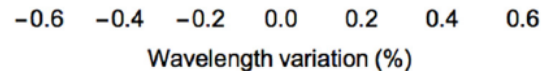
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IEEE TRANSACTIONS ON AUTOMATIC CONTROL, VOL. 69, NO. 4, APRIL 2024



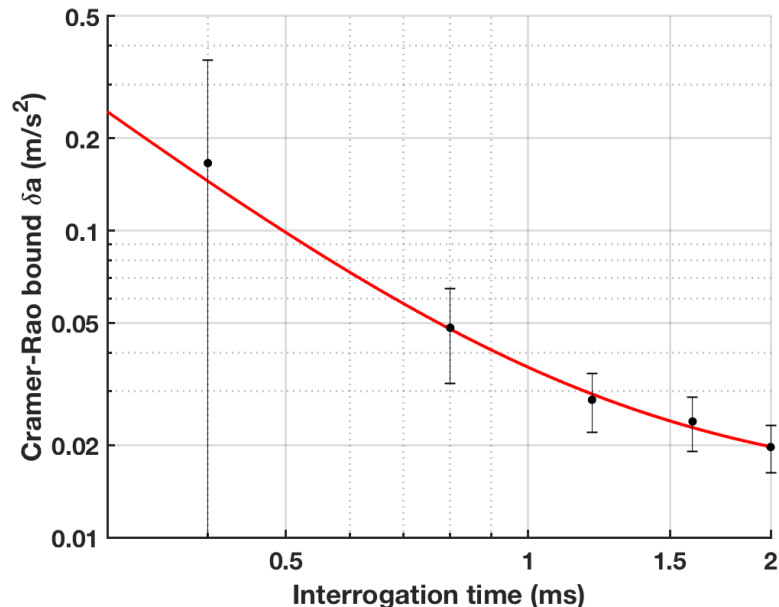
## Time-Domain Sensitivity of the Tracking Error

Sean O'Neil<sup>Ⓞ</sup>, Member, IEEE, Sophie Schirmer<sup>Ⓞ</sup>, Member, IEEE, Frank C. Langbein<sup>Ⓞ</sup>, Member, IEEE, Carrie A. Weidner<sup>Ⓞ</sup>, Member, IEEE, and Edmond A. Jonckheere<sup>Ⓞ</sup>, Life Fellow, IEEE



# So what's next?

- Build a 3D lattice system in Bristol
- Demonstrate a multi-axis inertial sensor (3 axes of acceleration, 3 axes of rotation)
- Open question #1: What is the best scaling with  $T$  that we can get?
- Open question #2: How robust is this method in the real world?
- Open question #3: What are the fundamental limitations of shaken lattice interferometry?



Thanks to:

--The Bristol team: Dr. Vineet Bharti (experiment), Dr. Dhritiman Chakraborty (theory), Harry Kendell (both experiment and theory)

--Collaborators Prof. Sophie Schirmer (Swansea), Prof. Frank Langbein (Cardiff), Prof. Edmond Jonckheere and Sean O'Neil (USC)

# Thank you for listening!

