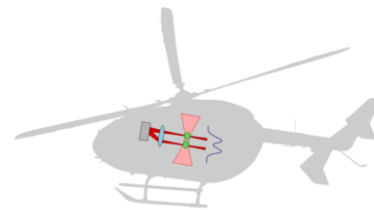


Matter-wave collimation to picokelvin energies with scattering length and potential shape control

Alexander Herbst

Leibniz University Hannover



**Guided matter-wave
interferometry group**

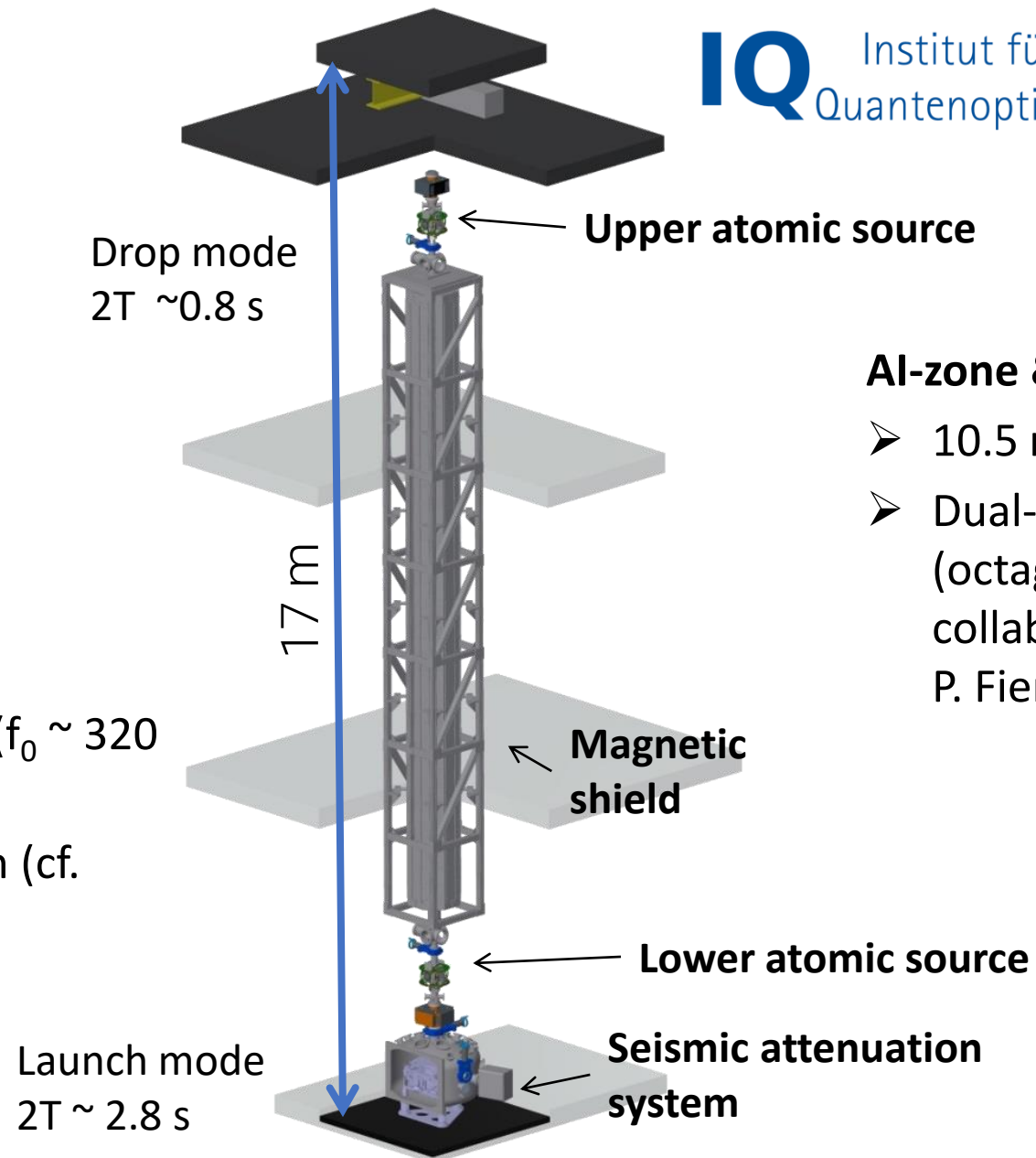
The VLBAI facility

Atomic sources

- Rubidium & ytterbium (quantum-degenerate)
- Drop/launch mode

Inertial reference

- Geometric anti-spring system ($f_0 \sim 320$ mHz) with active stabilization
- Rotation compensation system (cf. Kasevich, Müller groups)



Al-zone & magnetic shield

- 10.5 m CF 200 Al tube
- Dual-layer mu-metal (octagonal geometry; collaboration with P. Fierlinger, TU München)

The VLBAI facility

$$\Delta\phi = k_{\text{eff}} \cdot g \cdot T^2$$

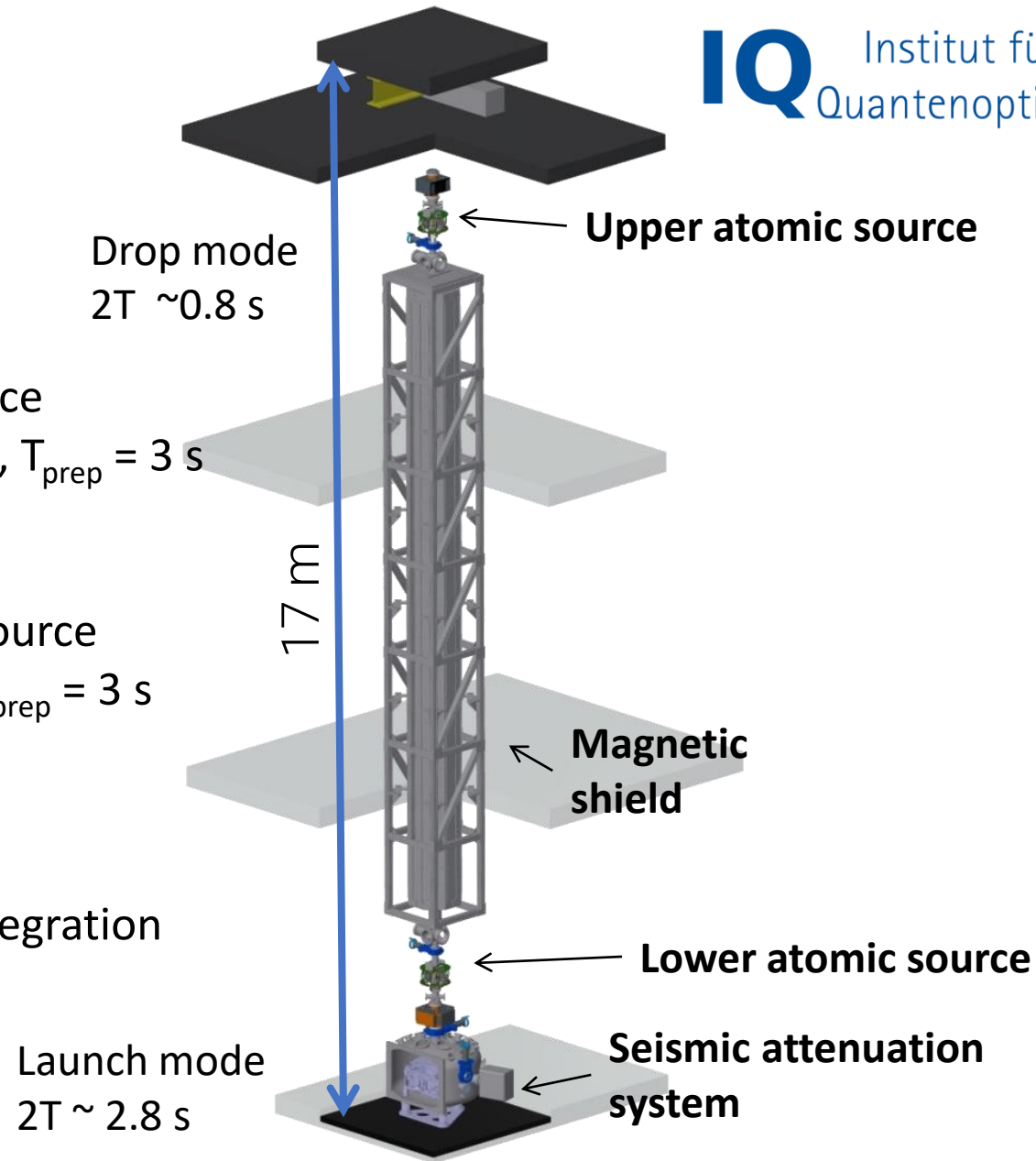
Drop mode – 1st generation source

2x10⁵ at, 2 photons, 2T = 800 ms, T_{prep} = 3 s
 → 1.7x10⁻⁹ m/s² @1s

Launch mode – 2nd generation source

1x10⁶ at, 4 photons, 2T = 2.8 s, T_{prep} = 3 s
 → 4x10⁻¹¹ m/s² @1s

➔ ~10⁻¹³ g after O(1 day) of integration



Gen. Rel. Grav. 43(7) 1931 (2011)

Wuhan (China)



J. Hogan, PhD thesis, Stanford U. (2010)

Stanford (USA)

+ ANU, Firenze, Berkeley, Stanford (II), UK, ...

[Schlippert et al., arXiv:1909.08524]

Long baseline challenges

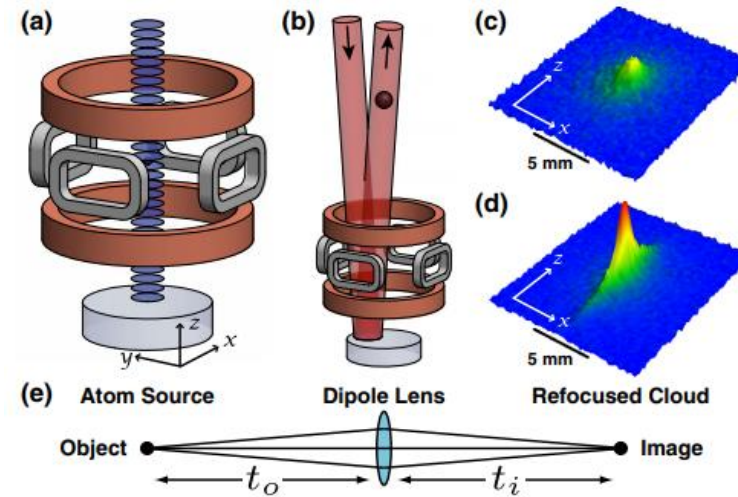


Long baselines usually demand large free apertures in-line (~10 cm or more)

- Challenging circumstances for strongly confining magnetic traps (large distance or beam clipping)
- Yb and Sr are non-magnetic
- Solution: All-optical setups

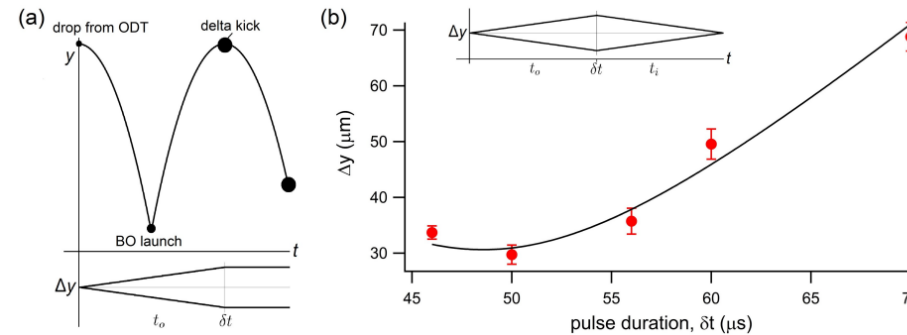
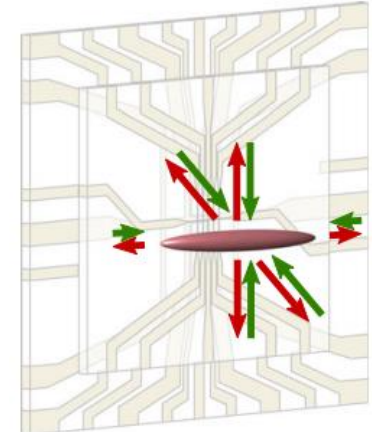
Matter-wave lensing

- Access to lowest kinetic energies
- Shortcuts of evaporation trajectories → higher flux
- Long initial time-of-flight can be experimentally challenging w/o micro-gravity

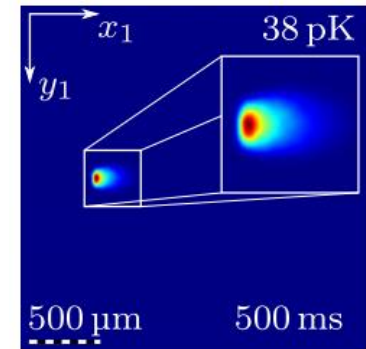


[Kovachy et al., PRL 2015]

QM-enh. DKC
3D collimation

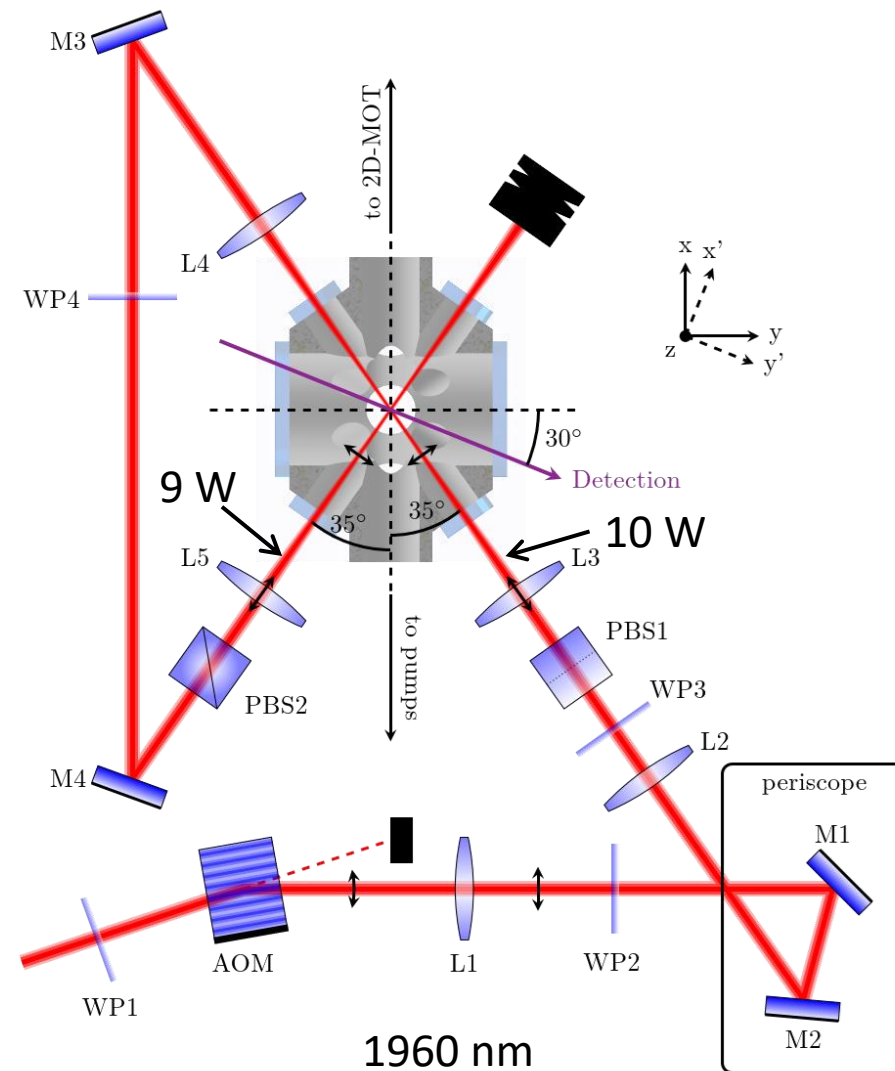
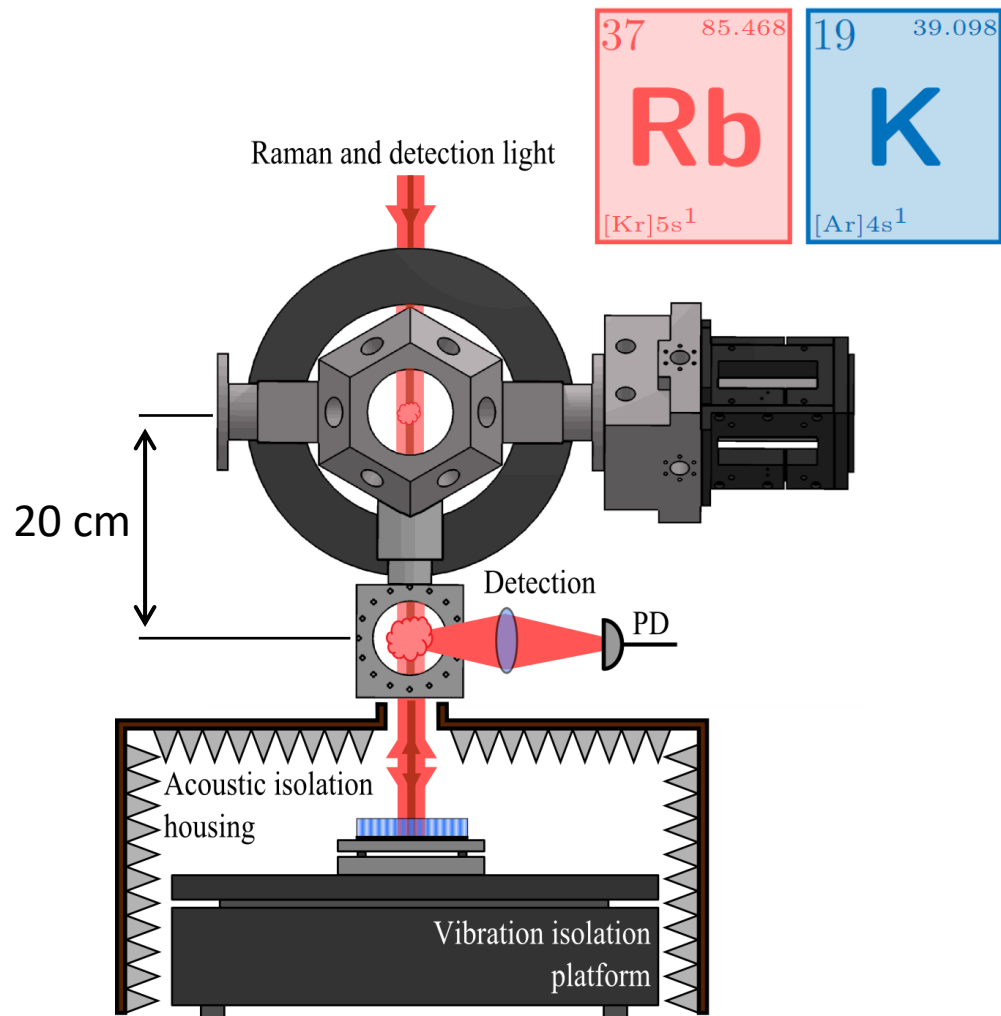


[Gochnauer et al., atoms 2021]

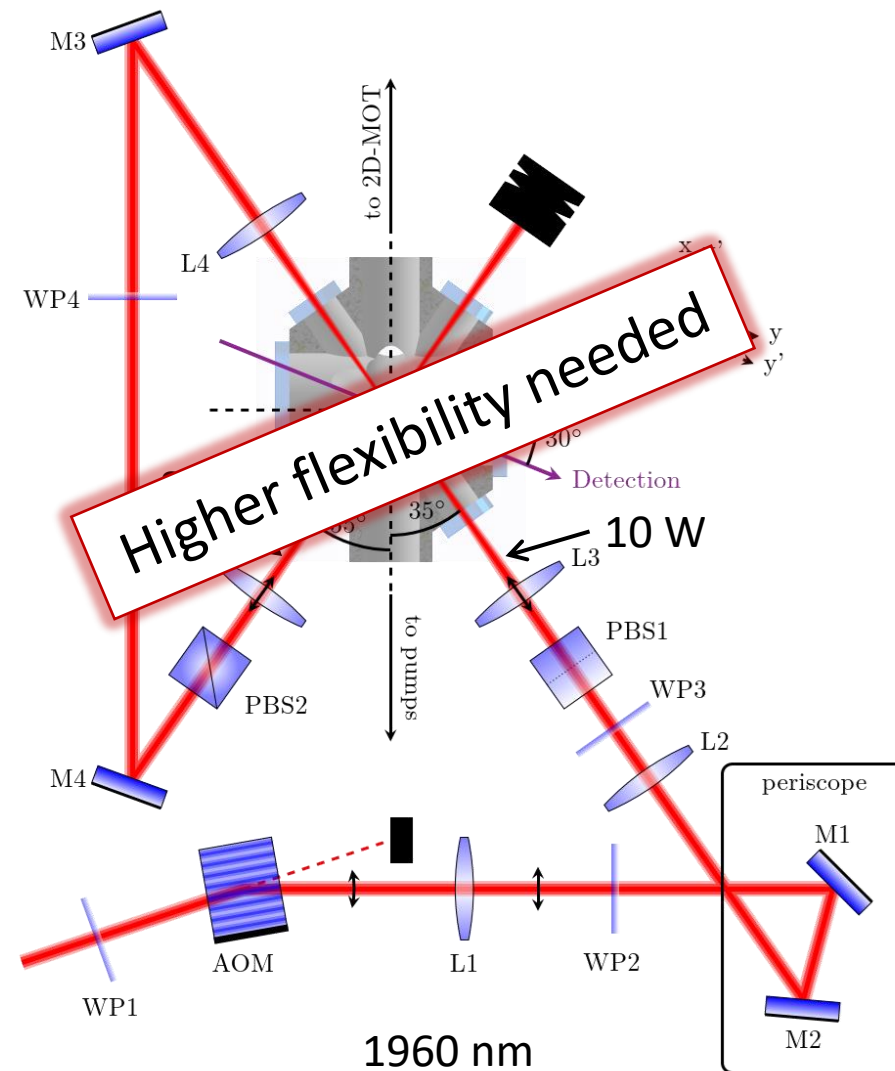
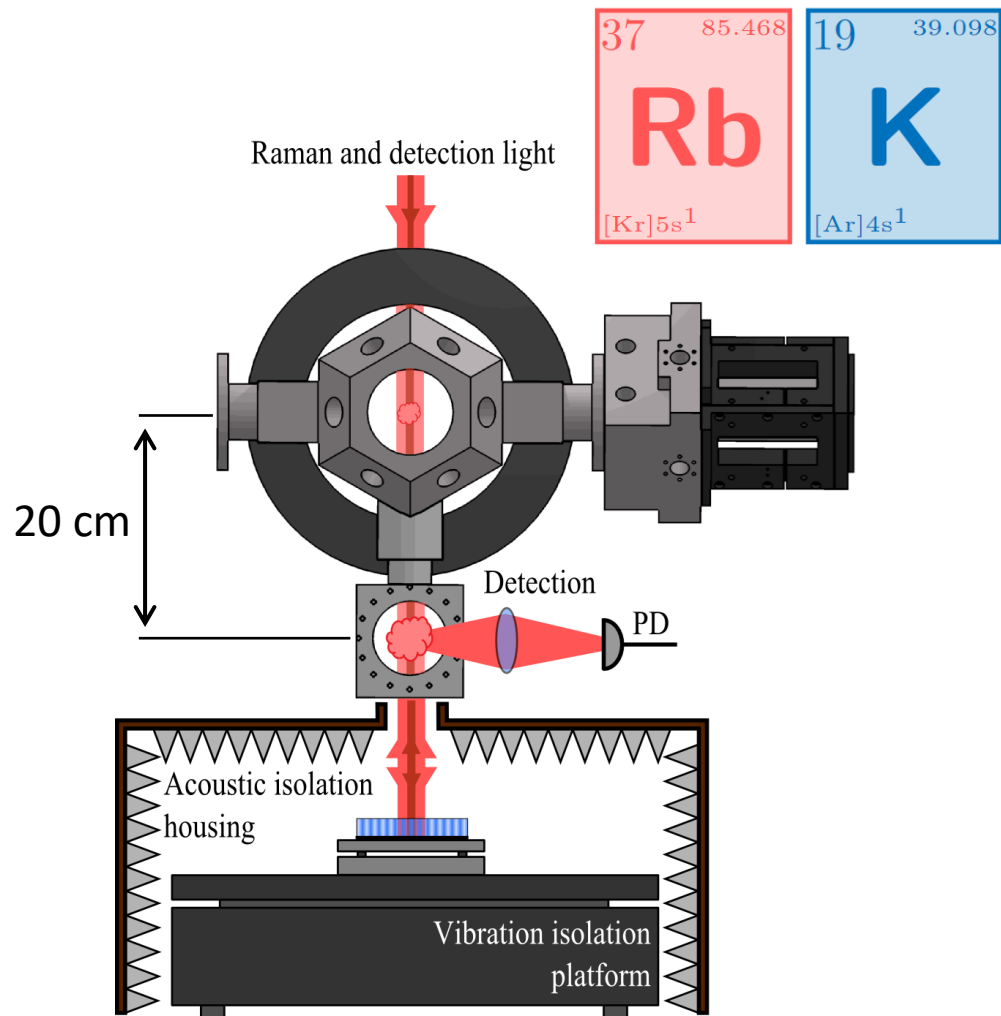


[Deppner et al., PRL 2021]

Experimental setup

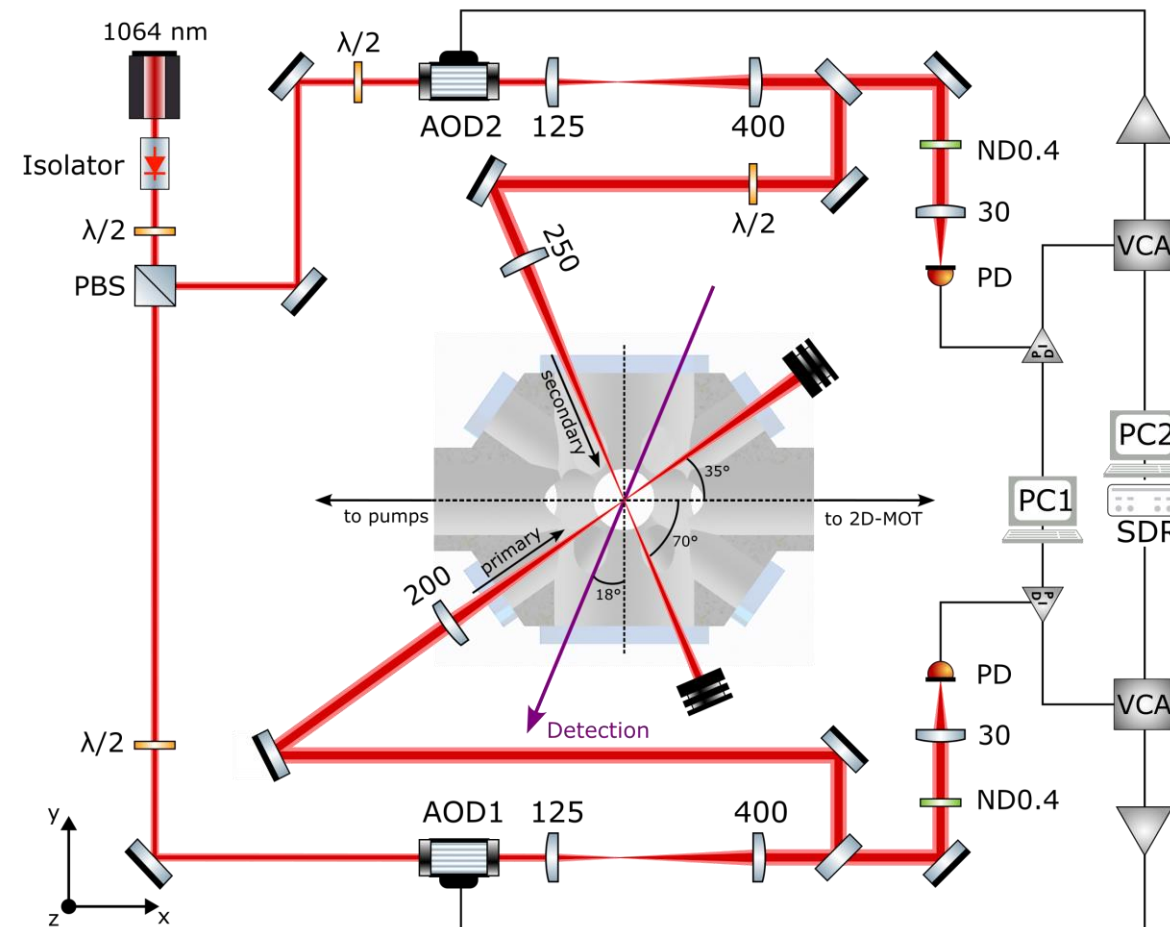


Experimental setup



Experimental setup

- Source systems for molasses cooled ^{39}K , ^{87}Rb
- Independent control of ODT beams:
 - Primary beam: 16 W with $w_0 = 24 \mu\text{m}$ waist
 - Secondary beam: 16 W with $w_0 = 30 \mu\text{m}$ waist
- 2D-AOD: 1.5 mm horizontal and vertical modulation stroke

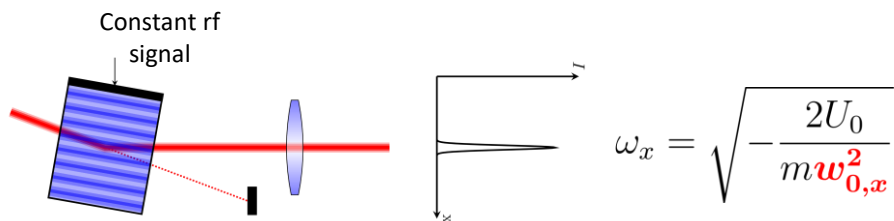


Experimental setup

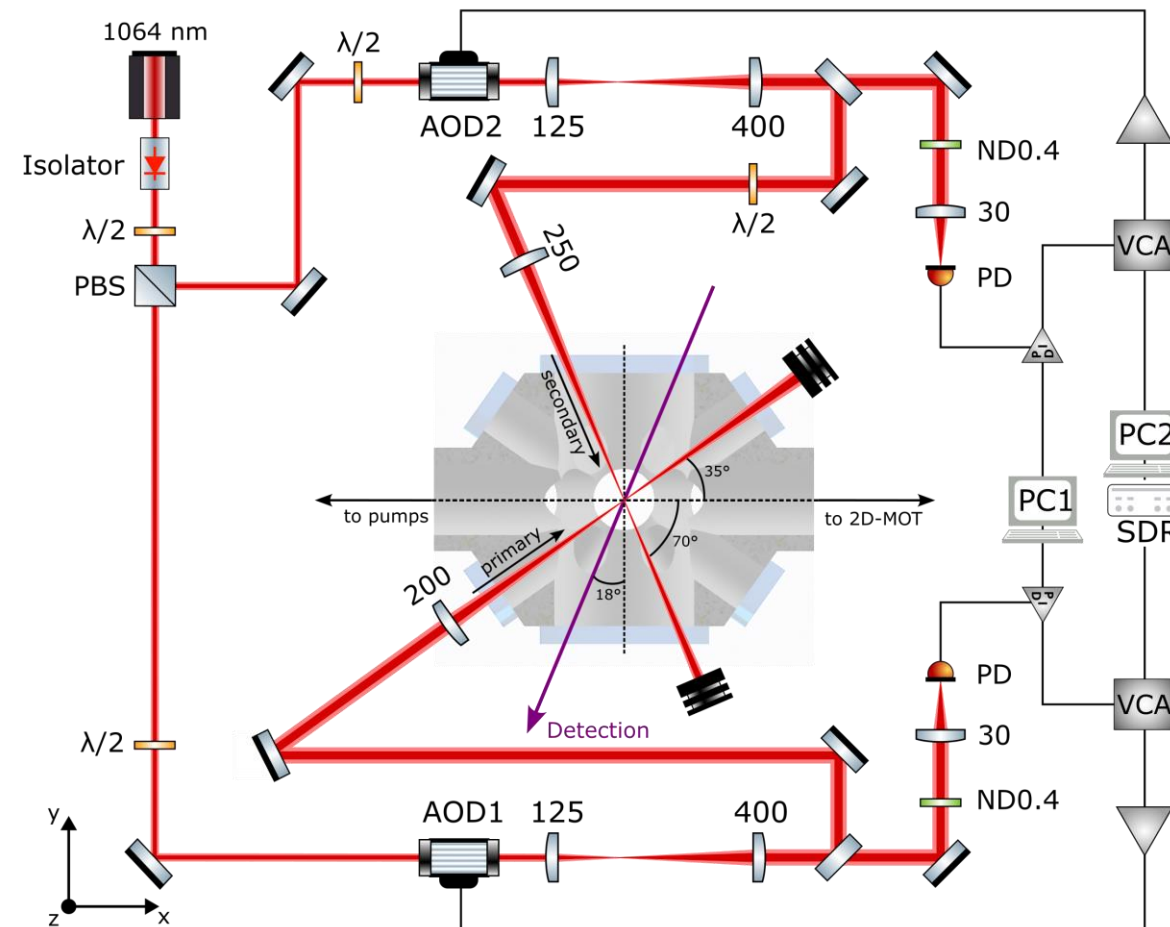
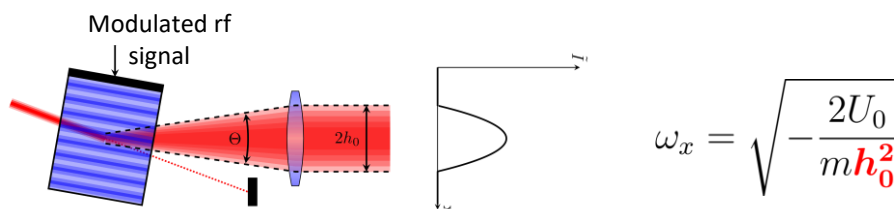
- Source systems for molasses cooled ^{39}K , ^{87}Rb
- Independent control of ODT beams:
 - Primary beam: 16 W with $w_0 = 24 \mu\text{m}$ waist
 - Secondary beam: 16 W with $w_0 = 30 \mu\text{m}$ waist
- 2D-AOD: 1.5 mm horizontal and vertical modulation stroke

Time-averaged optical potentials

Without center-position modulation:

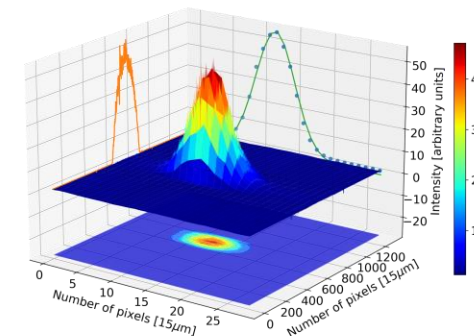
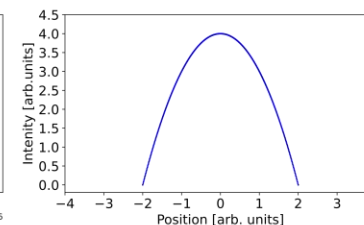
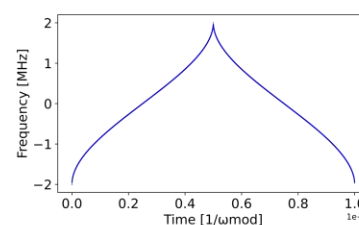


With center-position modulation:



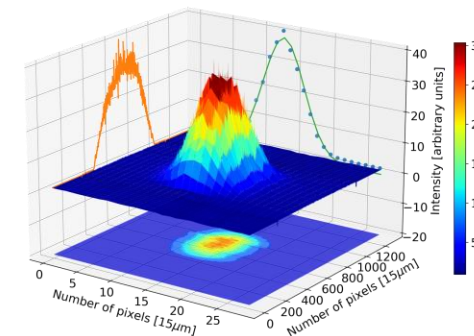
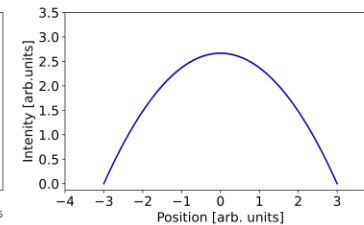
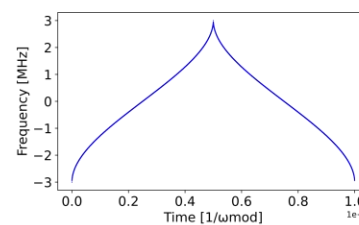
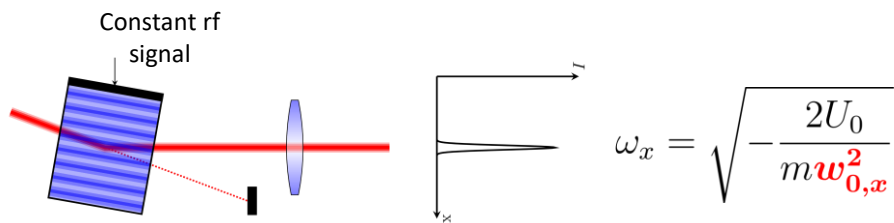
Experimental setup

- Source systems for molasses cooled ^{39}K , ^{87}Rb
- Independent control of ODT beams:
 - Primary beam: 16 W with $w_0 = 24 \mu\text{m}$ waist
 - Secondary beam: 16 W with $w_0 = 30 \mu\text{m}$ waist
- 2D-AOD: 1.5 mm horizontal and vertical modulation stroke

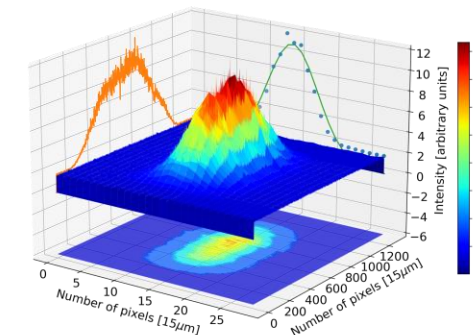
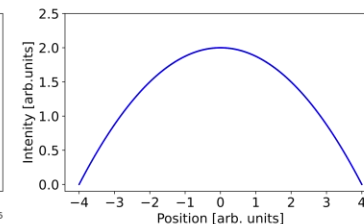
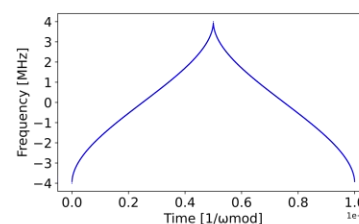
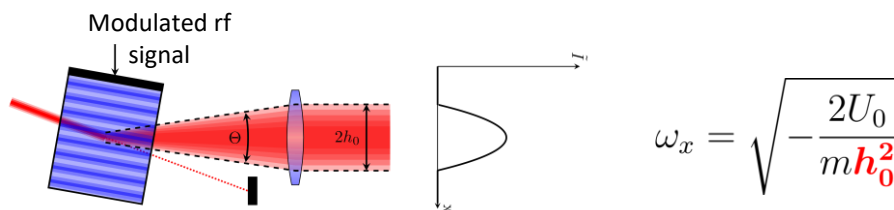


Time-averaged optical potentials

Without center-position modulation:



With center-position modulation:

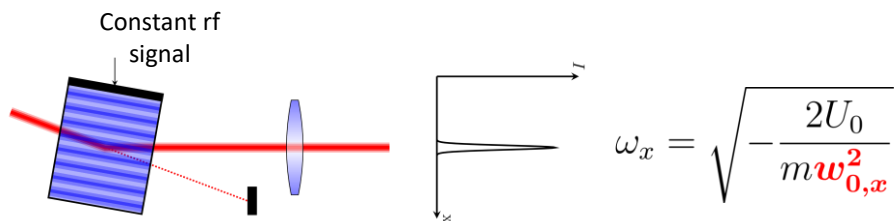


Experimental setup

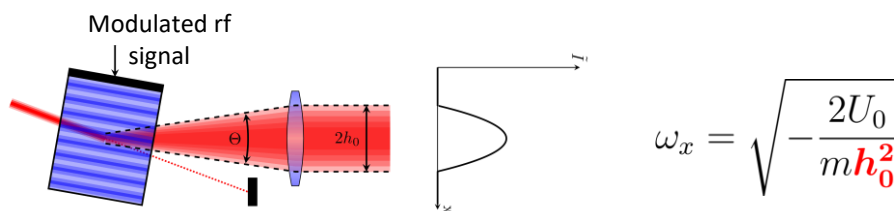
- Source systems for molasses cooled ^{39}K , ^{87}Rb
- Independent control of ODT beams:
 - Primary beam: 16 W with $w_0 = 24 \mu\text{m}$ waist
 - Secondary beam: 16 W with $w_0 = 30 \mu\text{m}$ waist
- 2D-AOD: 1.5 mm horizontal and vertical modulation stroke

Time-averaged optical potentials

Without center-position modulation:

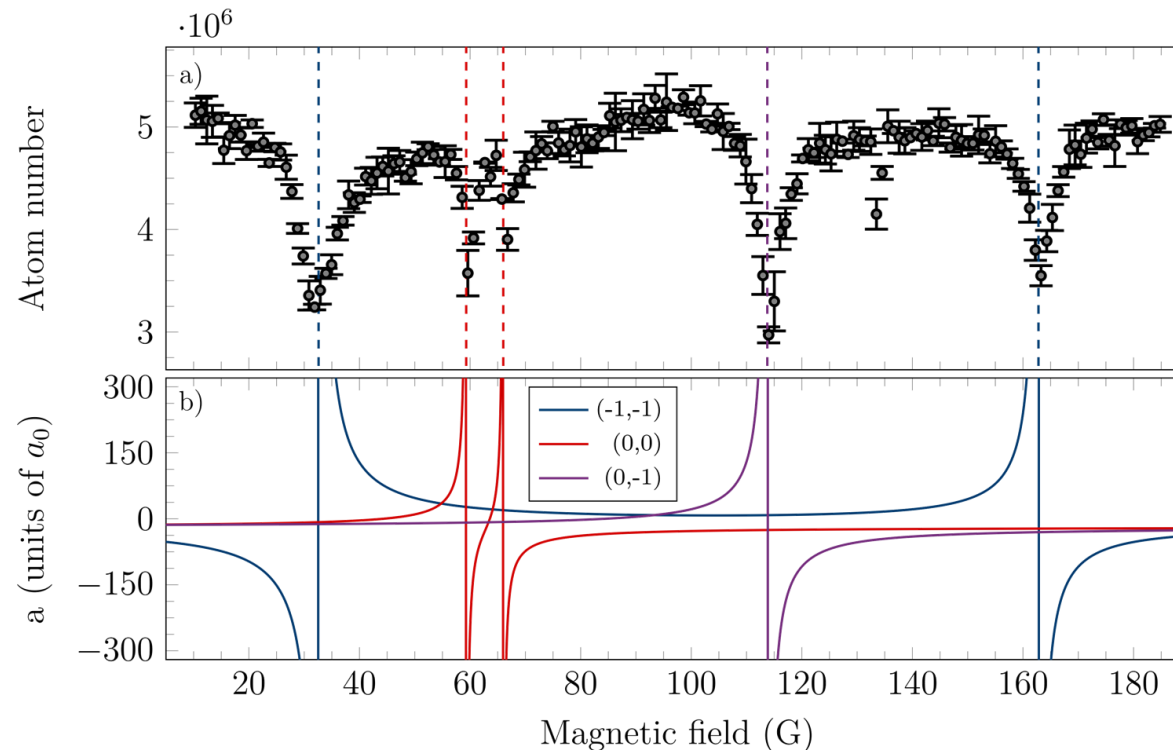


With center-position modulation:



Feshbach resonances

$$a(B) = a_{\text{bg}} \left(1 - \sum_i \frac{\Delta B_i}{B - B_{0i}} \right)$$



Evaporation optimization

Evaporative cooling dynamics:

$$\frac{dN}{dt} = - (\Gamma_{\text{ev}} + \Gamma_{3\text{b}} + \Gamma_{\text{bg}}) N$$

background gas collisions

$$\frac{dT}{dt} = - \left(\frac{\Gamma_{\text{ev}}}{3} (\eta + \alpha - 3) - \frac{\Gamma_{3\text{b}}}{3} - \frac{\dot{\bar{\omega}}}{\bar{\omega}} \right) T + \frac{\Gamma_{\text{sc}} E_r}{3k_b}$$

compression

off-resonant scattering

Evaporation rate:

$$\Gamma_{\text{ev}} \propto \frac{N \bar{\omega}^3 a^2}{T}$$

Three-body loss rate:

$$\Gamma_{3\text{B}} \propto \frac{N^2 \bar{\omega}^6 a^4}{T^3}$$

Evaporation efficiency:

$$R \propto \frac{T^2}{N \bar{\omega}^3 a^2}$$

Evaporation optimization

Evaporative cooling dynamics:

$$\frac{dN}{dt} = - (\Gamma_{ev} + \Gamma_{3b} + \Gamma_{bg}) N$$

background gas collisions

$$\frac{dT}{dt} = - \left(\frac{\Gamma_{ev}}{3} (\eta + \alpha - 3) - \frac{\Gamma_{3b}}{3} - \frac{\dot{\bar{\omega}}}{\bar{\omega}} \right) T + \frac{\Gamma_{sc} E_r}{3k_b}$$

compression

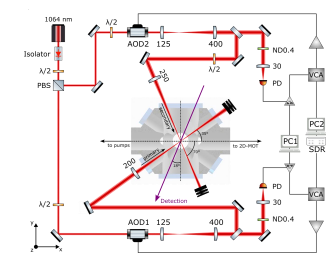
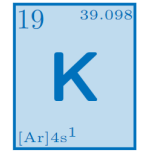
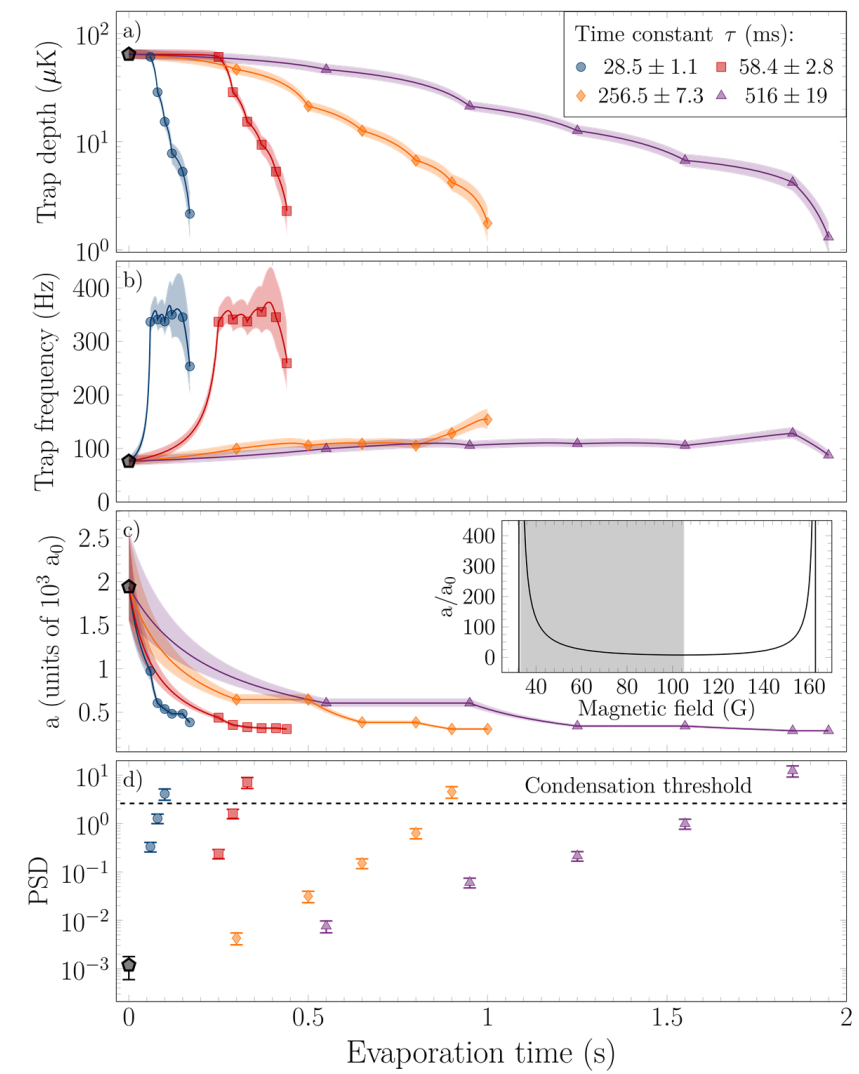
off-resonant scattering

Evaporation rate: $\Gamma_{ev} \propto \frac{N \bar{\omega}^3 a^2}{T}$

Three-body loss rate: $\Gamma_{3B} \propto \frac{N^2 \bar{\omega}^6 a^4}{T^3}$

Evaporation efficiency: $R \propto \frac{T^2}{N \bar{\omega}^3 a^2}$

- 6×10^4 condensed particles after 170 ms
- 6×10^5 condensed particles after 2 s

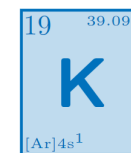
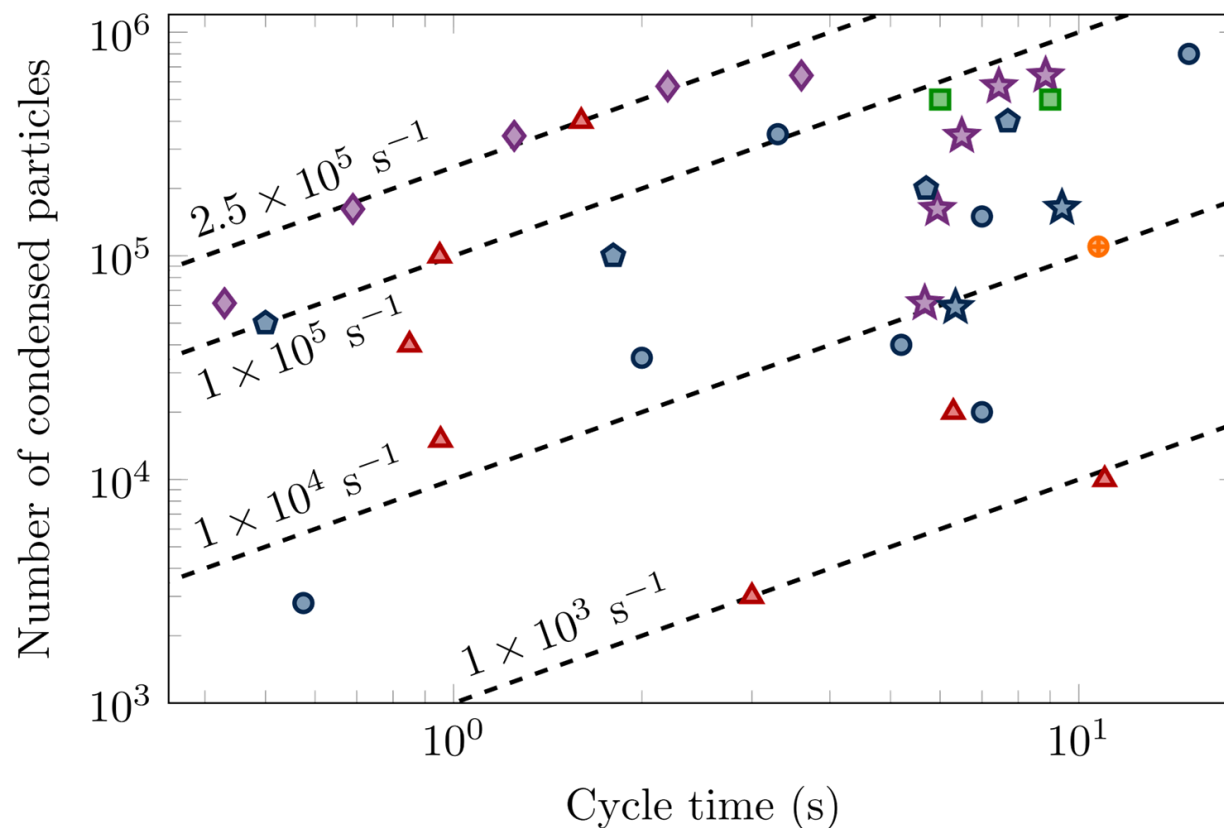


[Herbst et al., Phys. Rev. Research 6, 013139 (2024)]

Performance comparison

Current setup limited by MOT loading time

- Combination of ODTs with TAPs and tunable interactions allows to match the performance of chip traps

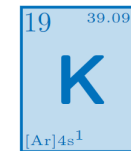
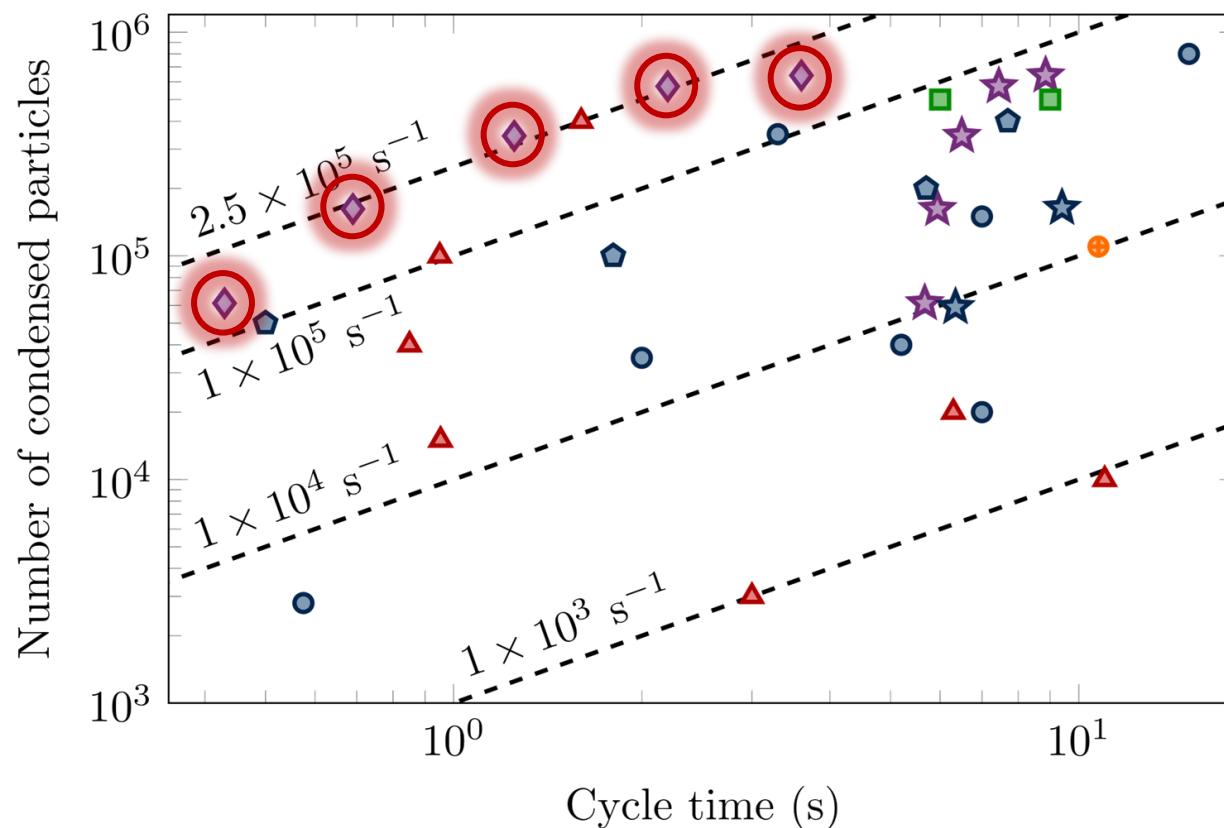


This work:	◆ - Adv. sce.	★ - 1 μm trap	★ - 2 μm trap
Others:	● - ODT	⬠ - TAP	⊕ - Laser cooled
	▲ - Chip	■ - Hybrid	

Performance comparison

Current setup limited by MOT loading time

- Combination of ODTs with TAPs and tunable interactions allows to match the performance of chip traps



This work:	◆ - Adv. sce.	☆ - 1 μm trap	★ - 2 μm trap
Others:	● - ODT	⬠ - TAP	⊕ - Laser cooled
	▲ - Chip	■ - Hybrid	

Performance comparison

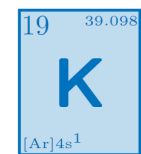
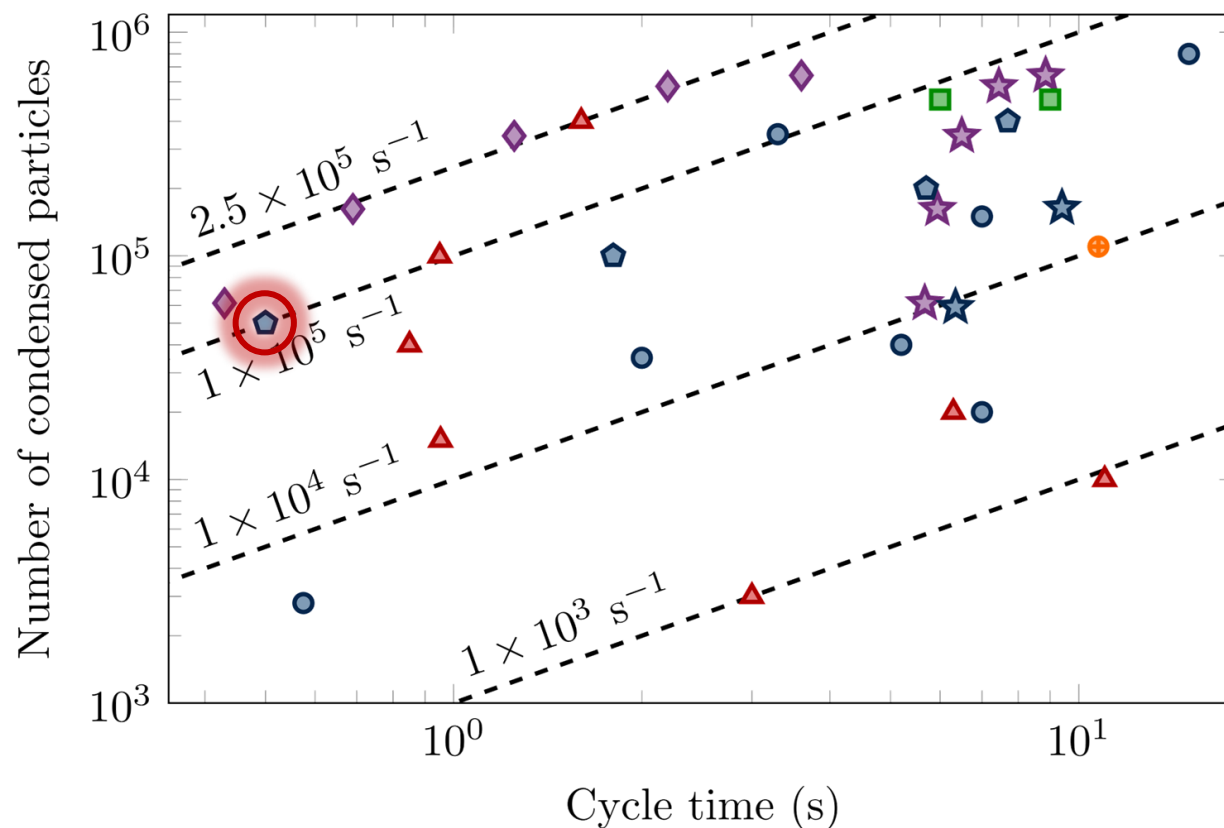
Current setup limited by MOT loading time

- Combination of ODTs with TAPs and tunable interactions allows to match the performance of chip traps

Similar results achieved with ^{87}Rb by further increasing the trap frequencies, using smaller beam waists:

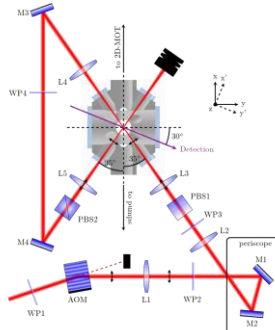
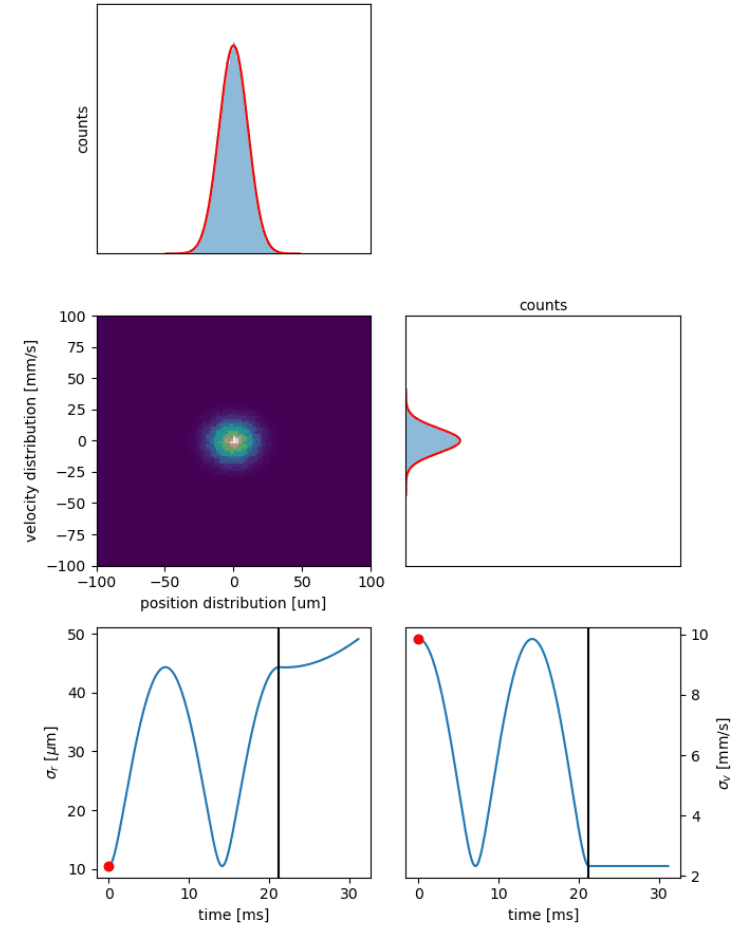
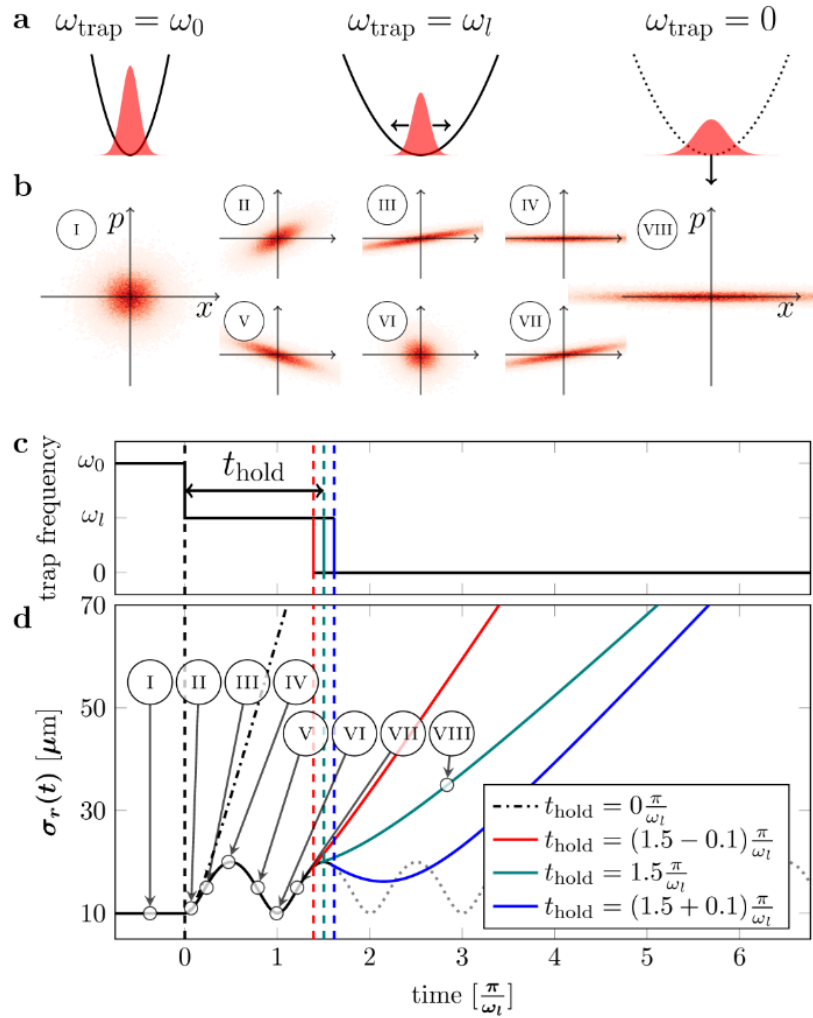
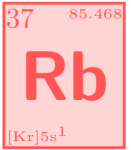
$$\Gamma_{\text{ev}} \propto \frac{N\bar{\omega}^3}{T} \quad R \propto \frac{T^2}{N\bar{\omega}^3}$$

[Hetzl, Dissertation, 2023, Leibniz University Hannover]



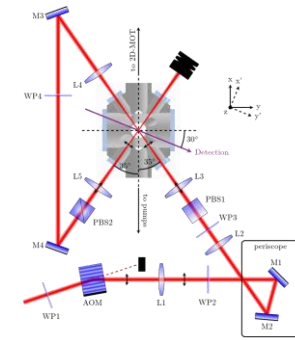
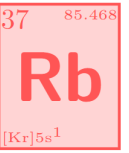
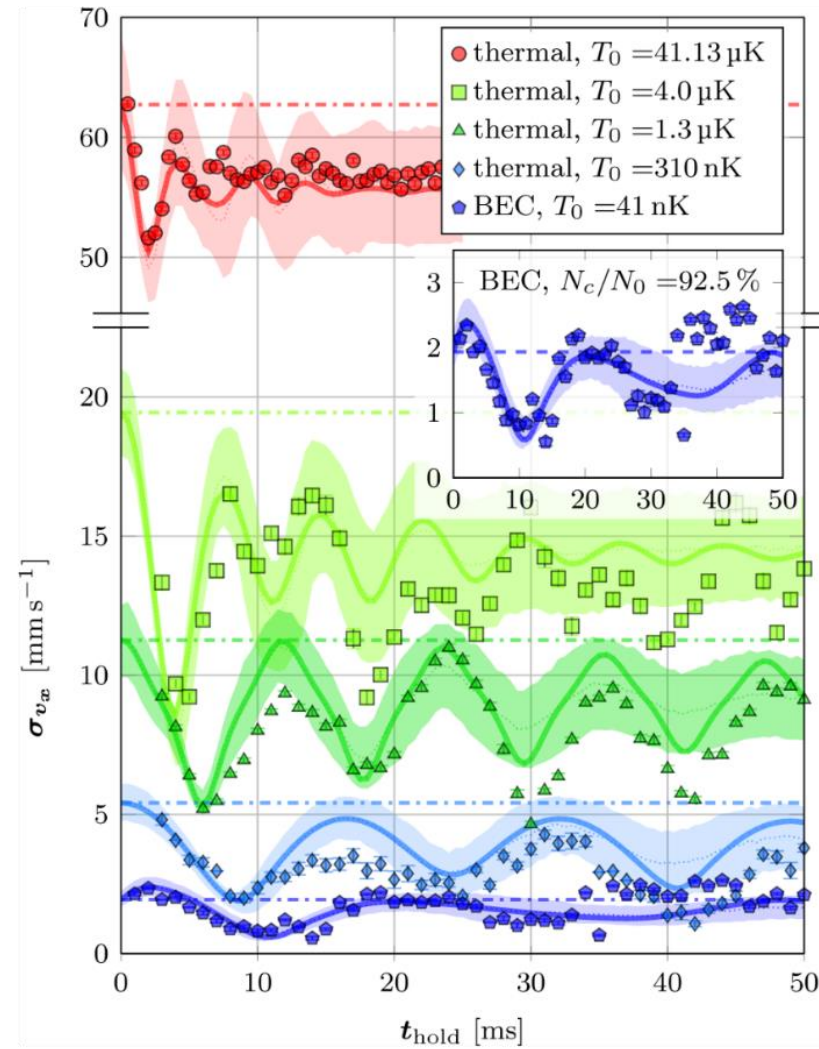
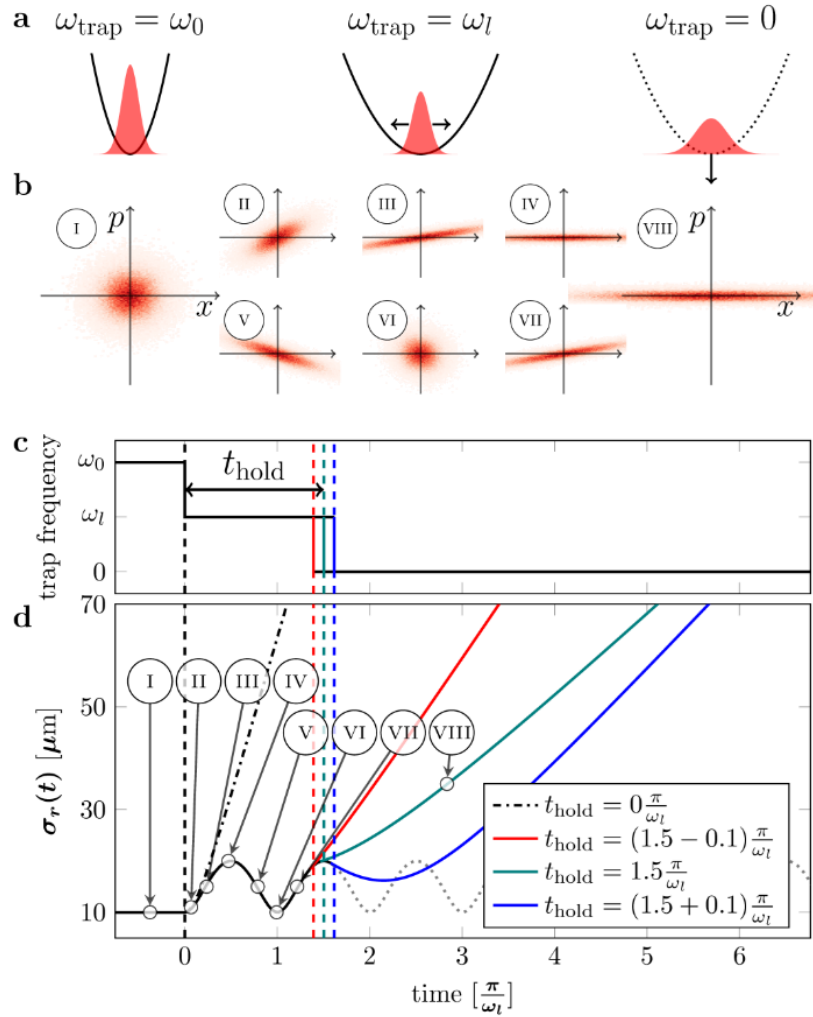
This work:	◆ - Adv. sce.	★ - 1 μm trap	★ - 2 μm trap
Others:	● - ODT	⬠ - TAP	⊕ - Laser cooled
	▲ - Chip	■ - Hybrid	

All-optical matter-wave lens



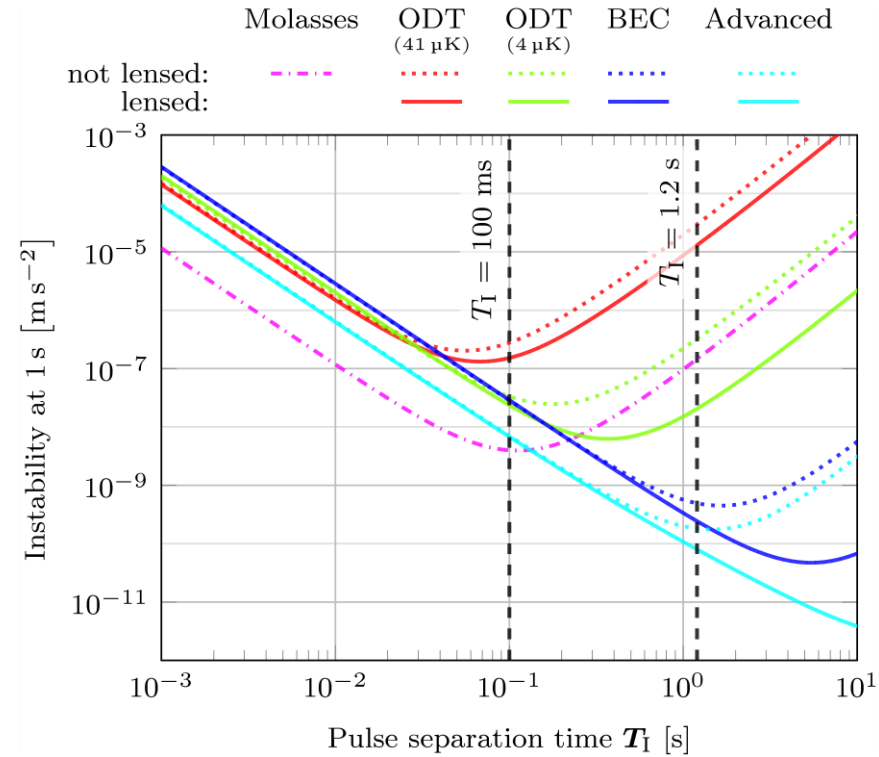
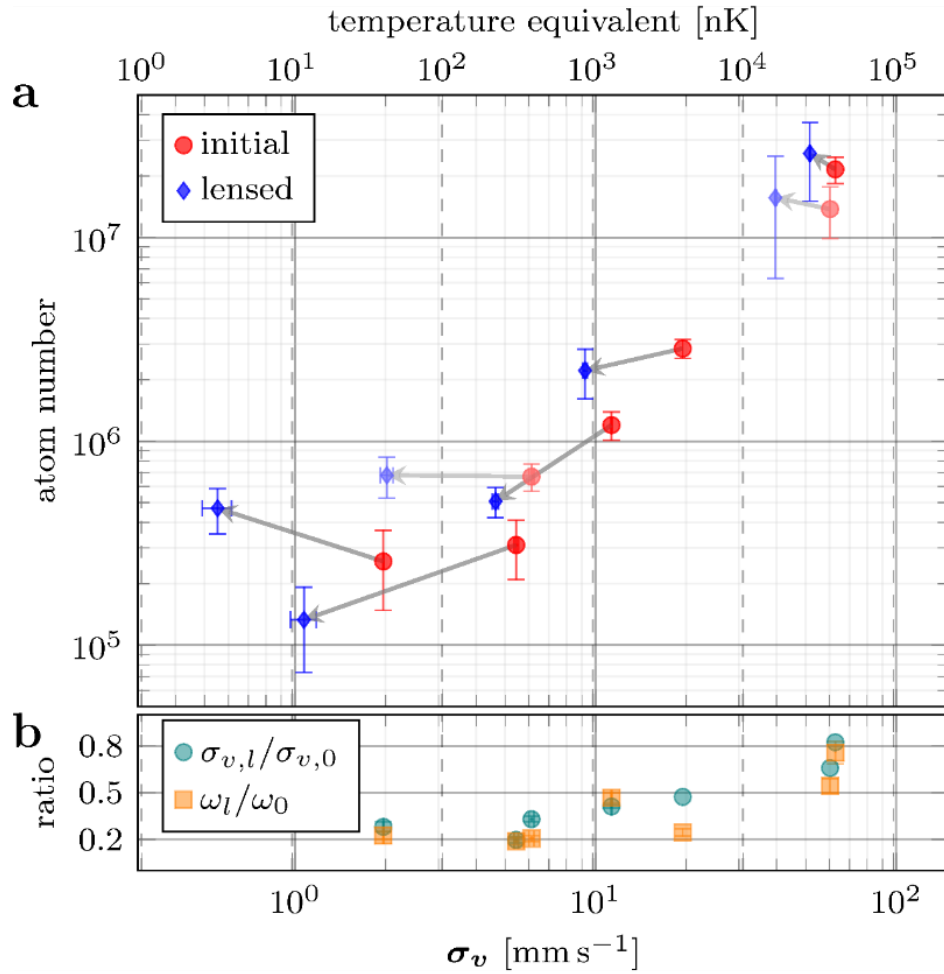
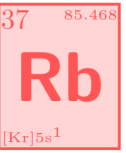
[Albers et al., Commun. Phys. 5, 60 (2022)]

All-optical matter-wave lens



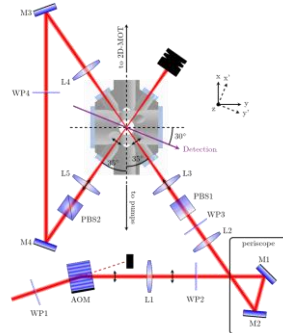
[Albers et al., Commun. Phys. 5, 60 (2022)]

All-optical matter-wave lens



Shortcuts of evaporative cooling possible

$$\sigma_a(\tau) = \frac{1}{C\sqrt{N}nk_{\text{eff}}T_I^2} \cdot \sqrt{\frac{t_{\text{cycle}}}{\tau}}$$



Matter-wave lens with ^{39}K

- Expansion fundamentally bound by uncertainty principle:

$$\Delta x \cdot \Delta p \geq \frac{\hbar}{2}$$

- Initial Δx depends on trap geometry/frequencies
- Repulsive interactions increase momentum:

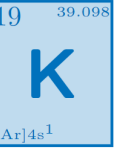
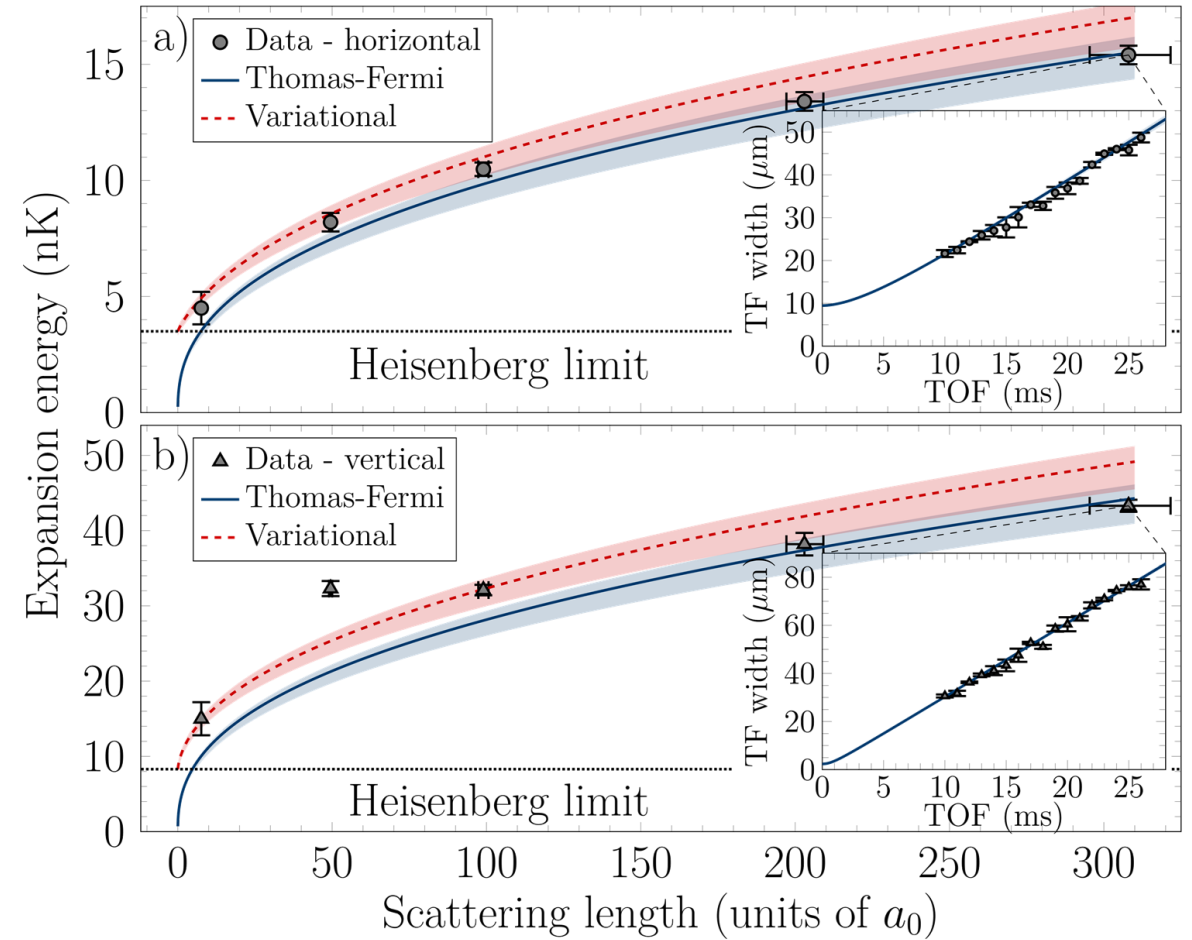
$$i\hbar \frac{\partial \Psi(\vec{r}, t)}{\partial t} = \left(-\frac{\hbar^2}{2m} \Delta + V(\vec{r}) + g|\Psi(\vec{r}, t)|^2 \right) \Psi(\vec{r}, t)$$

Thomas-Fermi approximation ($a \geq 50 - 150 a_0$):

- neglects kinetic energy term
- Thomas-Fermi distribution

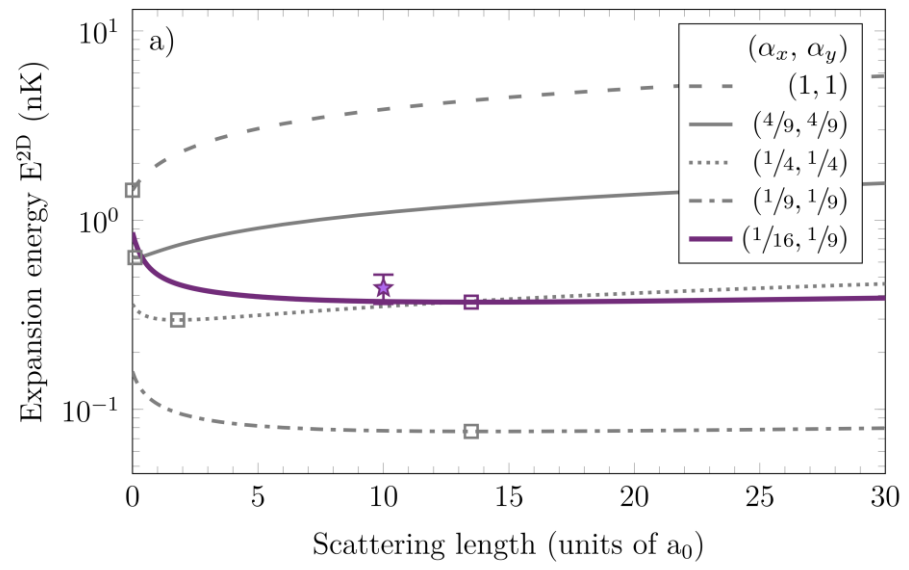
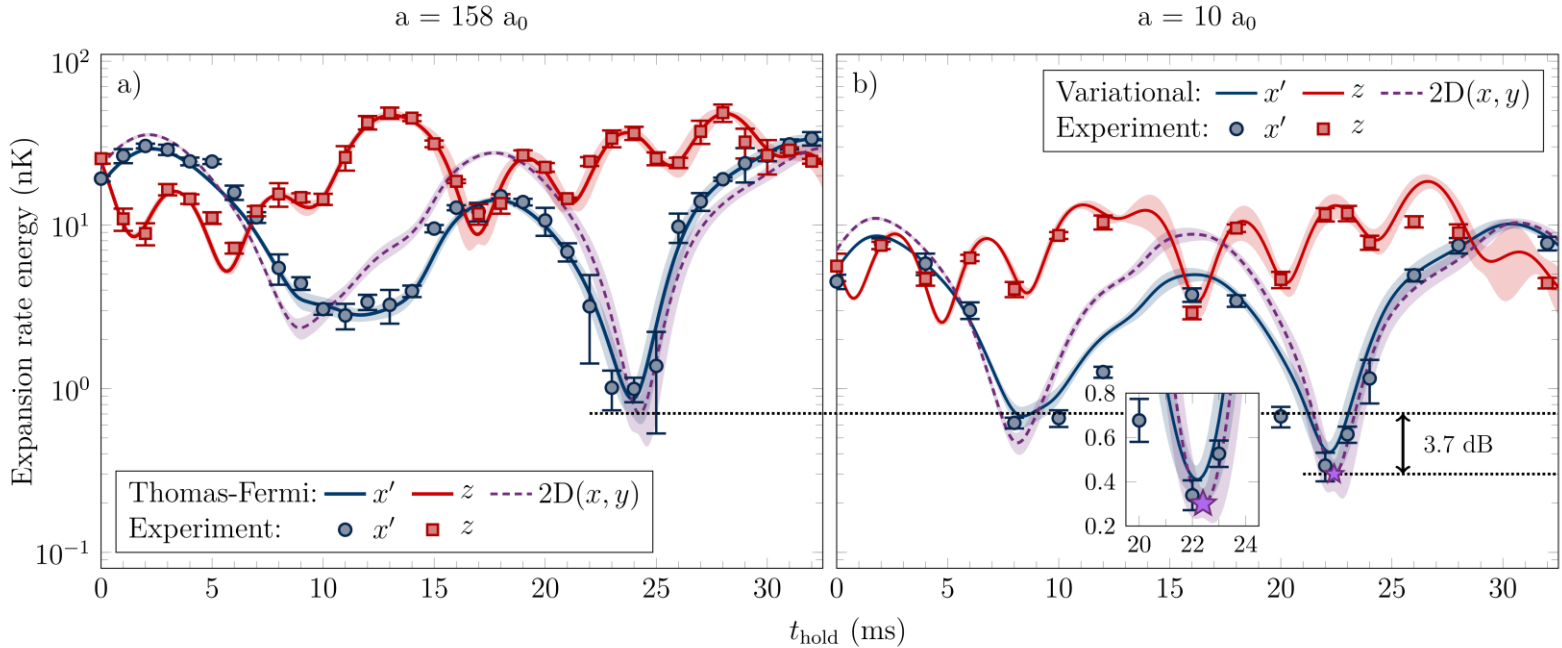
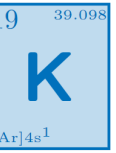
Variational approach ($a < 50 - 150 a_0$):

- includes kinetic energy term
- Gaussian distribution

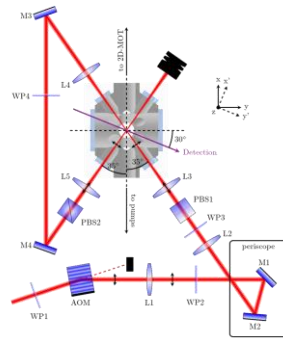


Matter-wave lens with ^{39}K

- Application of same lensing scheme with ^{39}K
- Dedicated tuning of scattering length allows to reduce mean-field energy prior to release

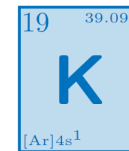


Results at 10 a_0 :
1D: 340 ± 12 pK
2D: 438 ± 77 pK

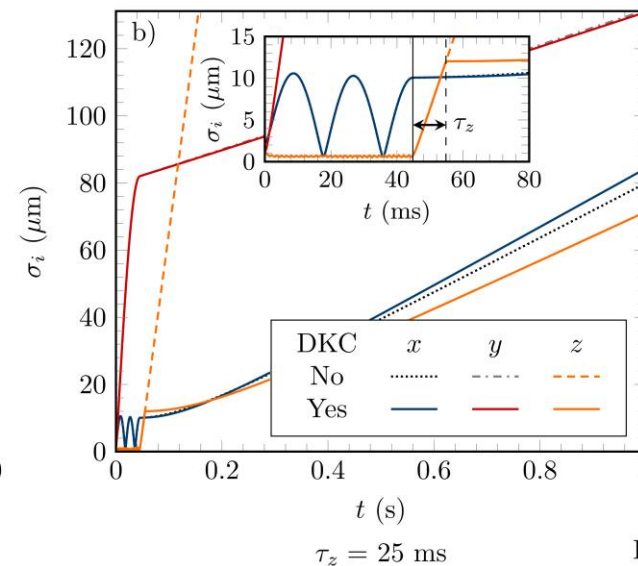
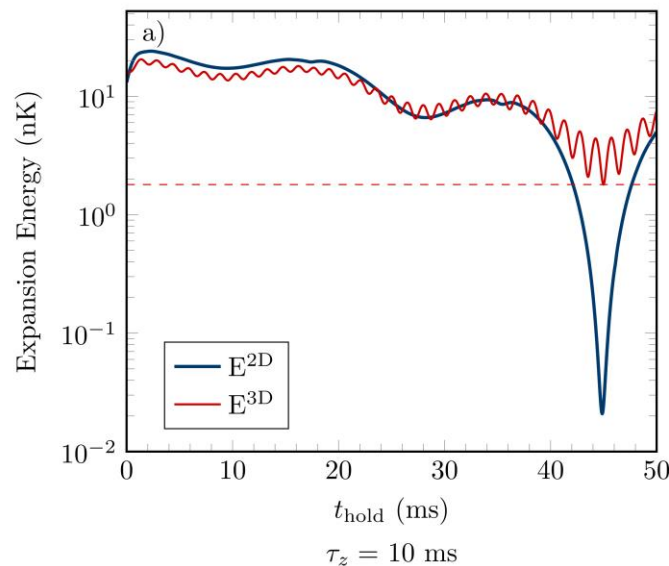


[Herbst et al., arXiv:2310.04383 (2024)]

Outlook

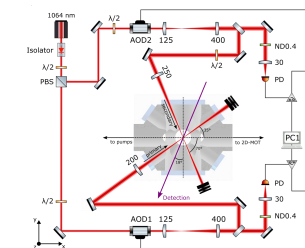
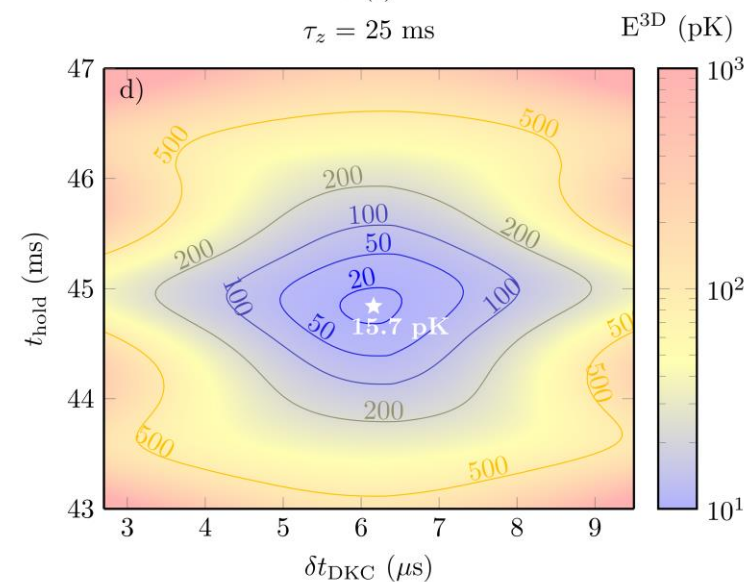
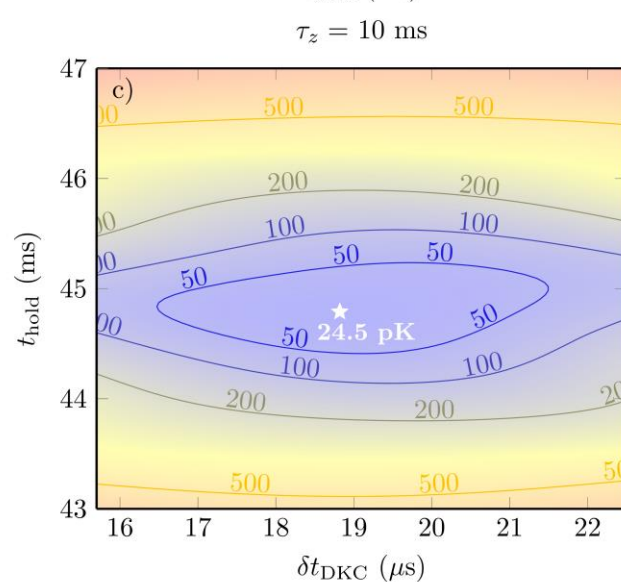


- Combination of continuous matter-wave lens with pulsed delta-kick collimation after release
- Co-moving ODT (via AOD) allows for up to 30 ms of free-fall



3D simulations:

- For ^{39}K below 20 pK with tunable interactions
- For ^{87}Rb below 100 pK w/o tunable interactions
- Directly applicable to Sr and Yb



[Herbst et al., arXiv:2310.04383 (2024)]

Thank you for your attention!

