



The case for measuring the gravitational field of the LHC beam

Daniel Braun

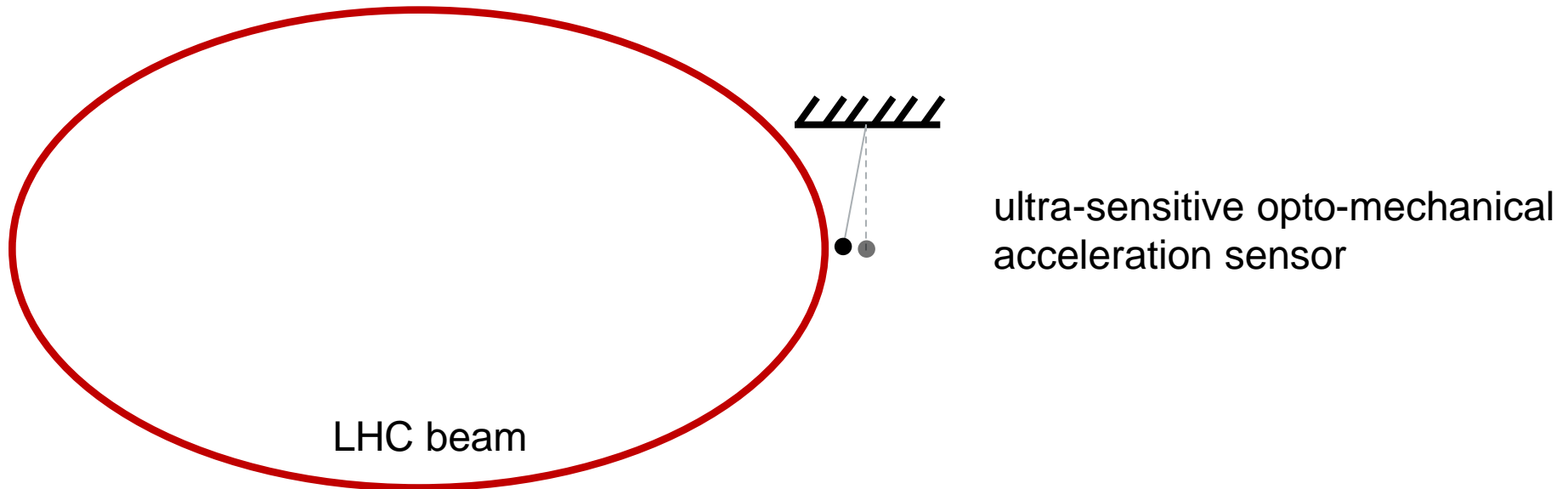
Alessio Belenchia, Daniel Carney,
Maurice-Garcia Sciveres, Dennis Rätzel,
Hendrik Ulbricht

Physics beyond Collider Workshop
CERN, March 27, 2024



Proposal

To measure the **gravitational near field** of the LHC proton beam



1. The science case
2. Why LHC?
3. How?
4. Technical challenges
5. Links to other proposals



The science case

- Lab-scale test of GR in entirely new parameter regime
 - **ultrarelativistic source of gravity** (in the sense of special relativity)

$$T_{\mu\nu} \simeq \mathcal{E} \begin{pmatrix} 1 & 0 & 0 & -1 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ -1 & 0 & 0 & 1 \end{pmatrix}$$

- Non-relativistic sensor
 - sensitive to T_{00} component
 - Newtonian gravity, but **source of gravity = almost pure kinetic energy** rather than mass
 - $1/r$ decay of gravitational force rather than $1/r^2$
 - **mm scale distance**
- Tests of GR so far:
 - on astrophysical scales (planetary orbits, BH or NS mergers, grav. lensing)
 - gravitational redshift experiments
 - precision tests of equivalence principle
 - Yukawa-type deviations from $1/r$ potential on small distances (down to $56\mu\text{m}$)?

... **all with mass dominating as source**



The science case

- Test alternative theories against experiment
 - e.g. Brans-Dicke theory
 - theories with free energy as source
 - in general, parametrized post-Newtonian theories, 10 parameters
- Medium- and long term perspectives:
 - cooling and squeezing of (ion?) beam => non-classical source of gravity?



Why LHC?

- Alternatives:
 - high power lasers (source = e.m. energy)
 - other accelerators
- What counts is **average power** and distance from beam:

Radial acceleration of non-relativistic probe mass:

$$a_\rho = -\frac{4GP_{\text{av}}}{c^3} \frac{1}{\rho} \quad \simeq -1.9 \cdot 10^{-20} \text{ m/s}^2 \text{ at } \rho = 2 \text{ mm}$$

	P_p	T_p^{cav}	P_{avg}	w_B	
laser {	pulses in cavity	$3 \cdot 10^{14} \text{ W}$	$100 \text{ fs} \cdot 10 \text{ kHz} \frac{8 \cdot 10^5}{\omega_0}$	$2 \cdot 10^{10} \text{ W}$	$< 100 \mu\text{m}$
	cw cavity	$2 \cdot 10^{11} \text{ W}$	$\frac{\pi}{\omega_0}$	$1 \cdot 10^{11} \text{ W}$	$< 100 \mu\text{m}$
	LHC	10^{14} W	10^{-9} s	$3.8 \cdot 10^{12} \text{ W}$	$16 \mu\text{m}$

Table 1: Comparison of LHC beam and laser-based sources: P_p pulse power, P_{avg} power averaged over time, w_B waist of beam. ω_0 = desired modulation frequency, T_p^{cav} effective pulse length (lasers: in cavity)



Detectors: mg-scale, monolithic pendulum

Featured in Physics



Demonstration of Displacement Sensing of a mg-Scale Pendulum for mm- and mg-Scale Gravity Measurements

Nobuyuki Matsumoto, Seth B. Cataño-Lopez, Masakazu Sugawara, Seiya Suzuki, Naofumi Abe, Kentaro Komori, Yuta Michimura, Yoichi Aso, and Keiichi Edamatsu

Phys. Rev. Lett. **122**, 071101 – Published 19 February 2019

Catano-Lopez et al. 2020

- $Q_m = 10^5$, $\omega_m = 2\pi \times 4.4$ Hz, improved version with $Q_m = 2 \times 10^6$
- Optical spring => resonance frequency $400 \text{ Hz} < \frac{\omega_0}{2\pi} < 1800 \text{ Hz}$
Feedback cooling compensates heating from optical spring, reduces Q

- Starting from room temperature, **demonstrated**

$$\omega_0 = 2\pi \cdot 280 \text{ Hz} \quad Q_{fb} = 250 \quad T_{fb} \simeq \text{mK}$$

$$\text{Displacement sensitivity } 3 \cdot 10^{-14} \text{ m}/\sqrt{\text{Hz}}$$

- Expect displacement from LHC: $\sim 10^{-24} \text{ m}$ (distance 2mm, full resonant amplitude)

But gain ~ 1000 from week-long measurement, ~ 1000 from nK cooling of pendulum mode in noise reduction



Optimization of pendulum setup

Total signal/noise ratio:

$$S/N = x_{\text{grav}} (1 - \exp(-\omega_0(\tau_{\text{tot}} - \tau_m)/(2Q_{\text{fb}}))) \frac{\sqrt{\tau_m}}{\sqrt{S_{\text{xx,tot}}}}$$

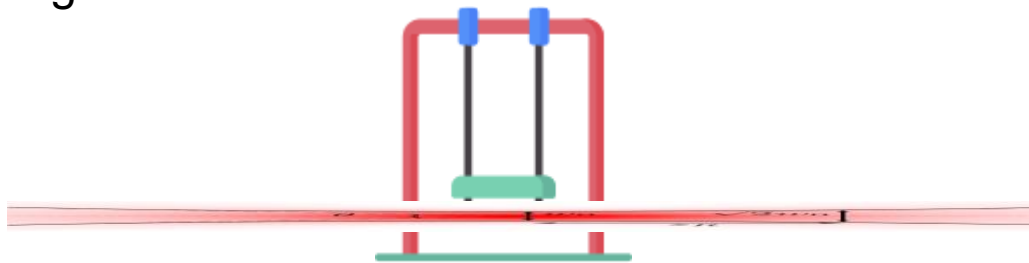
$$\simeq 0.01 \frac{(1 - e^{((\tau_m - \tau_{\text{tot}}) \frac{\omega_0}{2Q_{\text{fb}}})}) \sqrt{Q_{\text{fb}} m \tau_m}}{\omega_0 \sqrt{1 + \coth \frac{4 \cdot 10^{-12} \omega_0}{T_{\text{fb}}}}}$$

- Thermal + quantum noise
- SI units for all quantities
- $\rho_{\text{min}} = 100 \mu\text{m}$ from LHC beam line center

$$\tau_{\text{tot}} = \tau_r + \tau_m \quad \text{rise time + measurement time}$$

cylindrical pendulum mass, up to 50cm long, max mass to maintain ρ_{min}

=> $m = 33 \text{ mg}$ for Si





Optimization

- Max S/N over τ_m , ω_0 , Q_{fb} , T_{fb} with $\tau_{tot}=1$ week, $1 \leq \omega_0 \leq 10^4/s$, $1 \leq Q_{fb} \leq 10^8$, $1 \text{ nK} \leq T_{fb}$

$$\Rightarrow S/N \approx 0.6$$

$$\text{for } \tau_m = 3 \times 10^5 \text{ s, } \omega_0 = 2\pi \times 0.16\text{Hz, } T_{fb} = 1\text{nK, } Q_{fb} = 1.2 \times 10^5$$

relatively flat max, can accommodate technical constraints

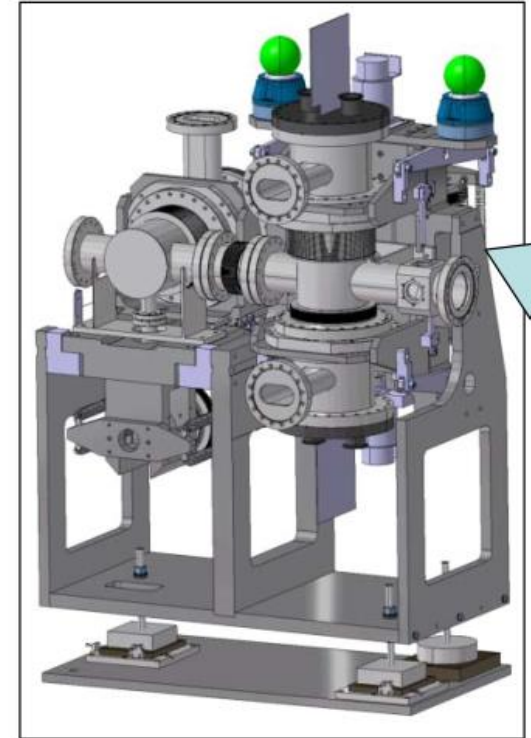
- Further improvements, but $\rho_{min} = 2 \text{ mm}$:
 - High-luminosity LHC \Rightarrow factor 2 in average power
 - Tungsten as material of pendulum body ($\rightarrow m=98 \text{ g}$)

$$\Rightarrow S/N \approx 6.1$$



Technical Challenges

1. Getting close to beam
 - avoid beam exposure
 - Roman pot design (secondary vacuum chamber in main vacuum)
 - convex pendulum design?
2. Cooling of detector:
 - cryostat (50mK) + electronic feedback cooling
 - optical spring
3. Shielding of e.m. forces from beam
4. Radioactive environment, beam halo
-> first discussion with Stefano Redaelli
5. Beam modulation (~kHz)



M. Orionno et al. EPAC 2006



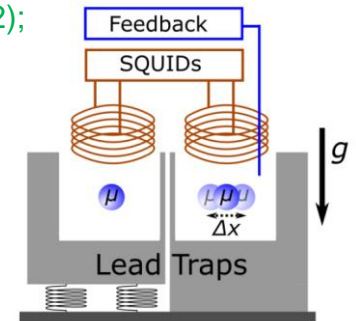
Alternative detectors – connection to DRD5 / RDq technology development

- Network of sensors
 - larger distance from beam, but many along the beam?
 - coherent averaging, $S/N \sim N_{\text{sensor}}$ DB, S. Popescu (2014); JME Fraise, DB (2015)
(+ new theory work ongoing with Maria M. Marchese, Stefan Nimmrichter, DB, D. Rätzel)
 - squeezed quantum noise ?
 - digital filtering, precisely known time-dependent signal?

- Suspended (nano-?)magnets in superconducting trap
 - projected acceleration sensitivity $10^{-14} \text{ g/Hz}^{1/2}$
 - UHV
 - superconducting electronics, SQUID read-out of displacements
 - natural electric field shielding, radiation protection ?

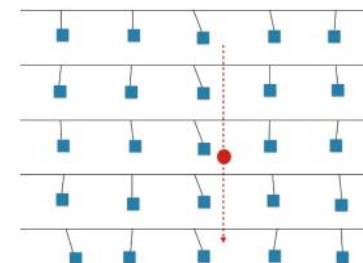
Horizon 2020 project

C. Timberlake et al. arXiv 2110.02263;
A. Vinante, C. Timberlake, H. Ulbricht
review (2022);



- Connection to wind-chime dark-matter detector

D. Carney et al. arXiv 1903.00492





Conclusion

- Measurement of gravitational near-field of LHC beam with optomechanical detector(s) appears to be possible in principle in near future.
- Would allow test of GR and alternative gravity theories in a completely new parameter regime on short distances in a lab experiment: kinetic energy as source!
- Substantial technological hurdles need to be overcome – let's join forces !
- Long-term perspective: cooled and squeezed particle beam as quantum source of gravity
 - Felix Spengler, Dennis Rätzel, DB, *New J. Phys.* 24 (2022) 053021
 - Fabienne Schneider, Dennis Rätzel, DB, *Class. Quant. Grav.* 35, 195007 (2018) and 36, 205007 (2019)



Marie Skłodowska-Curie Action IF program –
Project-Name "Phononic Quantum Sensors for Gravity"
(PhoQuS-G) – Grant-Number 832250

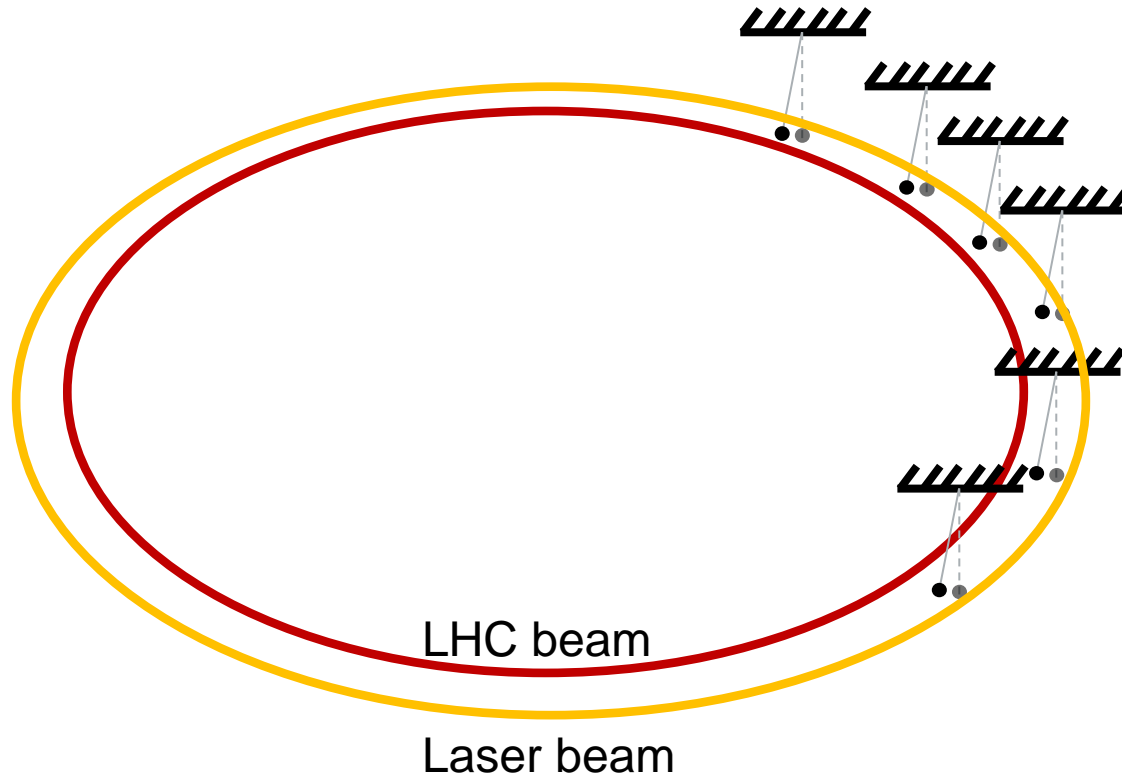


Alexander von Humboldt
Stiftung/Foundation

Feodor Lynen Research Fellowship



Backup



- Phase coherent addition of signal from N detectors by coupling to a common oscillator, e.g. mode of a ring laser.
- Expect S/N improvement $\sim N$ rather than $\sim N^{1/2}$

new theory work ongoing with Maria M. Marchese, Stefan Nimmrichter, DB, D. Rätzel



Sources 1: Modulated cw lasers + cavity

- Source high-power cw laser into high-finesse cavity

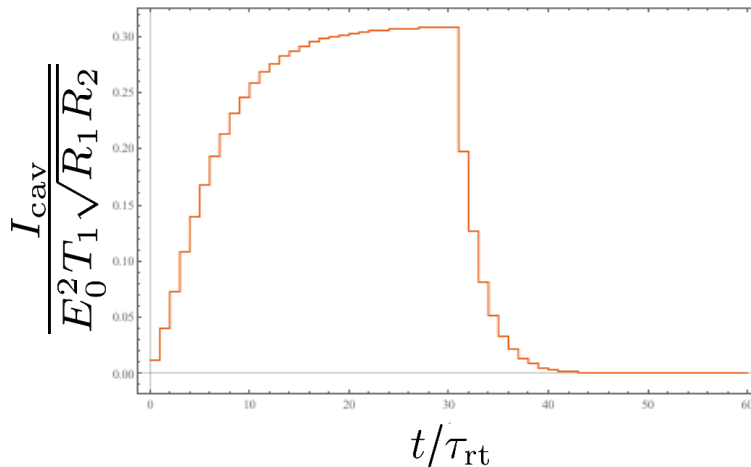
$$P_{\text{cav}}^{\text{max}} = \frac{2F}{\pi} P_{\text{pump}}$$

$$\tau_{\text{mod}} \gg \tau_{\text{L}}$$

modulation period loss time



- Slow modulation => periodic build up and decay of circulating power in cavity
- Commercially available:
 - IR 500 kW cw multi-mode; $F \sim 10^6 \Rightarrow P_{\text{av}} \sim 100 \text{ GW}$
 - 20 kW single mode => focus down to $\sim 1 \mu\text{m}$ waist



$$\tau_p / \tau_{\text{rt}} = 30, R_1 = R_2 = 0.8$$



Sources 1: Modulated cw lasers + cavity

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$$\tau_{mod} \gg \tau_L$$

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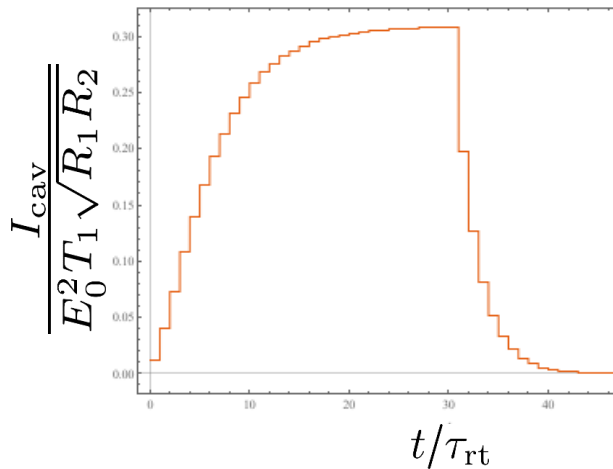
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➤ 20 kW single mode => focus down to $\sim 1 \mu m$ waist

- m^2 size mirrors to avoid thermal damage, stabilize thermal modes?



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YLS, 1-120+ kW

Ytterbium CW Laser Systeme

YLS Ytterbium Faserlasersysteme haben Leistungen von 1 bis 120 kW und arbeiten im CW oder moduliertem Betrieb von bis zu 5 kHz mit einem Steckdosenwirkungsgrad über 40 %. Die Leistung kann frei zwischen 10% - 100% der Nennleistung gewählt werden, ohne Änderung der Strahlqualität oder Strahleigenschaften. Daher eignen sich diese Laser ideal zum Schweißen, Bohren und Präzisionsschneiden. Es werden Bearbeitungsoptiken mit langer Brennweite für eine verbesserte Tiefenschärfe und minimale Verschmutzung der Schutzgläser verwendet. Die Laser können mit Faserlängen von bis zu 100 Metern und einer Vielzahl an Multi-Port Strahlschaltern, Strahlkopplern und Scannern ausgestattet werden.

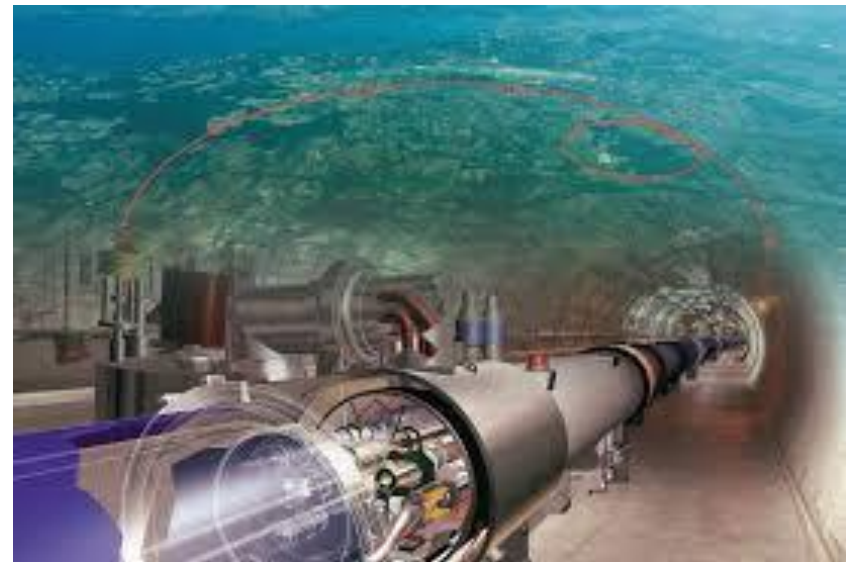


[Überblick](#) | [YLS bis zu 120 kW](#) | [Anwendungen](#) | [Strahlführungen](#) | [Kühler](#) | [QCW 2x Peak Power](#)



Sources 2: LHC particle beam

- 6.5 TeV energy per proton
 - ⇒ rest mass negligible
 - ⇒ almost identical energy-momentum tensor as light beam
- 2808 bunches with $\sim 10^{11}$ protons each over ring of 26,659 m length
 - ⇒ bunch rate 31.2Mhz
 - ⇒ lower frequencies via modulation?
- $P_{av} \sim 10^{12}$ W, beam waist $16 \mu\text{m}$





Signals versus sensitivities

	rod		liquid helium		pendulum
ω_0	$2\pi \cdot 10^3 \text{ Hz}$	$2\pi \cdot 10^9 \text{ Hz}$	$2\pi \cdot 2.8 \cdot 10^3 \text{ Hz}$		$2\pi \cdot 280 \text{ Hz}$
sensitivity	$1 \cdot 10^{-17} \frac{\text{m}}{\sqrt{\text{Hz}}}$	$4 \cdot 10^{-17} \frac{\text{m}}{\sqrt{\text{Hz}}}$	$2 \cdot 10^{-17} \frac{\text{N}}{\sqrt{\text{Hz}}}$	$1 \cdot 10^{-12} \frac{\text{m}}{\sqrt{\text{Hz}}}$	$3 \cdot 10^{-14} \frac{\text{m}}{\sqrt{\text{Hz}}}$
limiting factor	thermal noise	SQL	thermal noise		thermal noise
expected amplitude					
laser pulses	$2 \cdot 10^{-25} \text{ m}$	$2 \cdot 10^{-34} \text{ m}$	$2 \cdot 10^{-25} \text{ N}$	$1 \cdot 10^{-20} \text{ m}$	$3 \cdot 10^{-26} \text{ m}$
cw cavity	$4 \cdot 10^{-25} \text{ m}$	$4 \cdot 10^{-34} \text{ m}$	$3 \cdot 10^{-25} \text{ N}$	$2 \cdot 10^{-20} \text{ m}$	$5 \cdot 10^{-26} \text{ m}$
LHC beam *	$8 \cdot 10^{-24} \text{ m}$	$9 \cdot 10^{-32} \text{ m}$	$7 \cdot 10^{-24} \text{ N}$	$4 \cdot 10^{-19} \text{ m}$	$1 \cdot 10^{-24} \text{ m}$

* assumption of modulation of LHC beam while maintaining average power

- Full resonant amplitudes, require long build-up times for large Q
- $T=5\text{mK}$ for thermal noise, 2mm distance for pendulum
- What can help: long integration time, lower T , get closer to beam