

# Critical QFTs with spontaneous breaking of scale symmetry

Daniel Litim

CERN TH

&

US

UNIVERSITY  
OF SUSSEX

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# 3d critical theories

many applications in condensed matter physics

Dirac materials

Wilson-Fisher fixed points

mass generation and symmetry breaking

AdS/CFT conjecture

3d critical bosons and critical fermions

relate to higher-spin gauge theories on AdS4

Klebanov, Polyakov '02  
Sezgin, Sundell '03

Giombi, Yin '12

CFT vs higher spin symmetry

Maldacena, Zhiboedov '11, '12

Chern-Simons-matter dualities

Aharony, Giombi, Gur-Ari, Maldacena, Yacoby, '12  
Seiberg, Senthil, Wang, Witten, '16

# today:

3d QFTs with strongly interacting FPs  
and *spontaneous* scale symmetry breaking

**scalars**

**$O(N)$**

**fermions**

**Gross-Neveu**

**Yukawa**

**Gross-Neveu — Yukawa**

based on [2207.10115](#), [2212.06815](#), [2311.16246](#)  
and ongoing work with **Charlie Cresswell-Hogg**

# recap: $O(N)$ symmetric scalars

3d: **super**-renormalisable  $(\phi^* \phi)_{3d}^3$

free UV fixed point

Wilson-Fisher IR fixed point

exactly solvable at infinite N

**main tool:** functional RG

Wilson '71  
Polchinski '84  
Wetterich '92

$$\partial_t \Gamma_k = \frac{1}{2} \text{tr} \left\{ \left[ \Gamma_k^{(2)} + R_k \right]^{-1} \cdot \partial_t R_k \right\}$$

$$t = \ln k$$

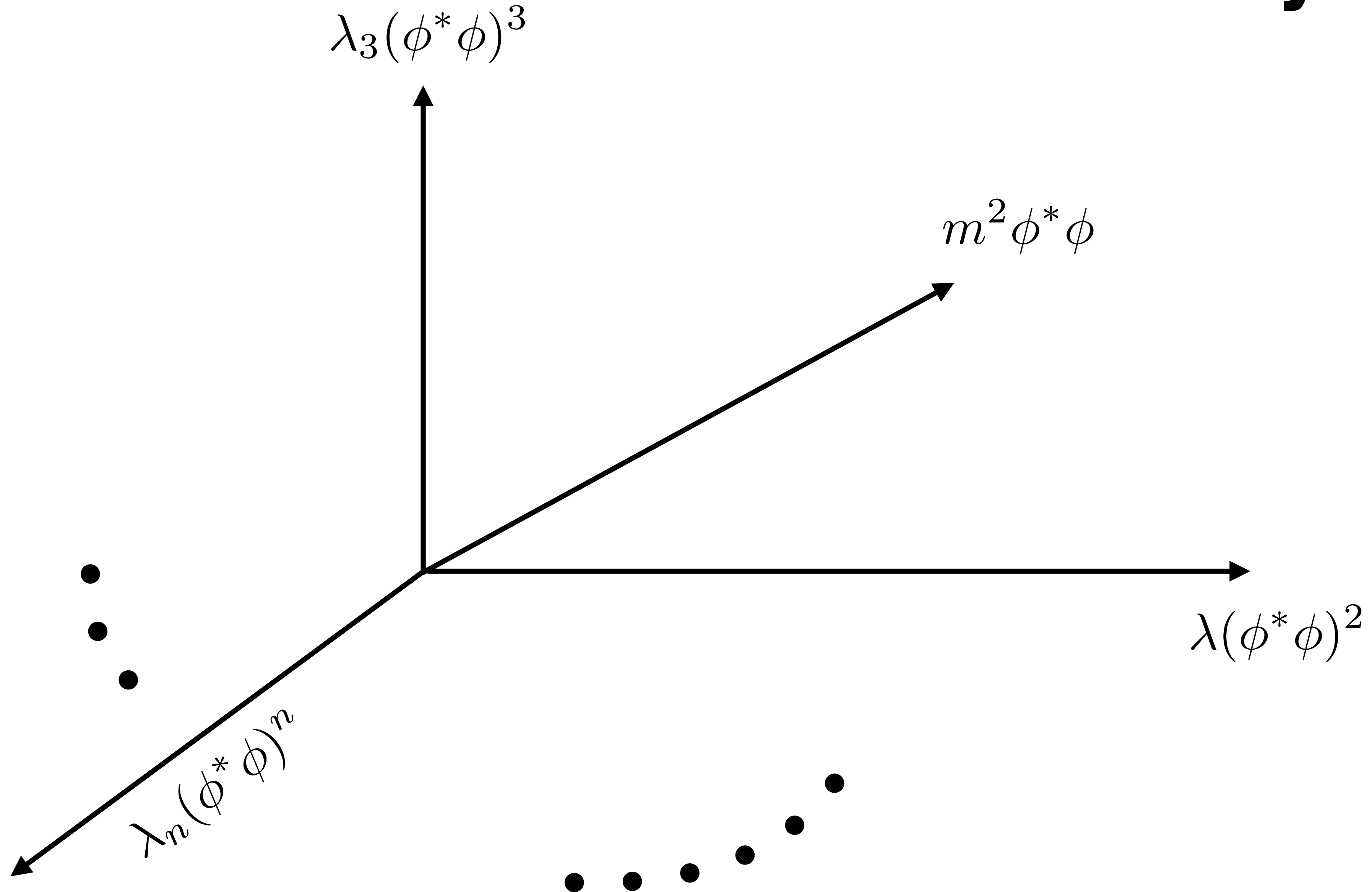
**UV**  
 $k = \Lambda$



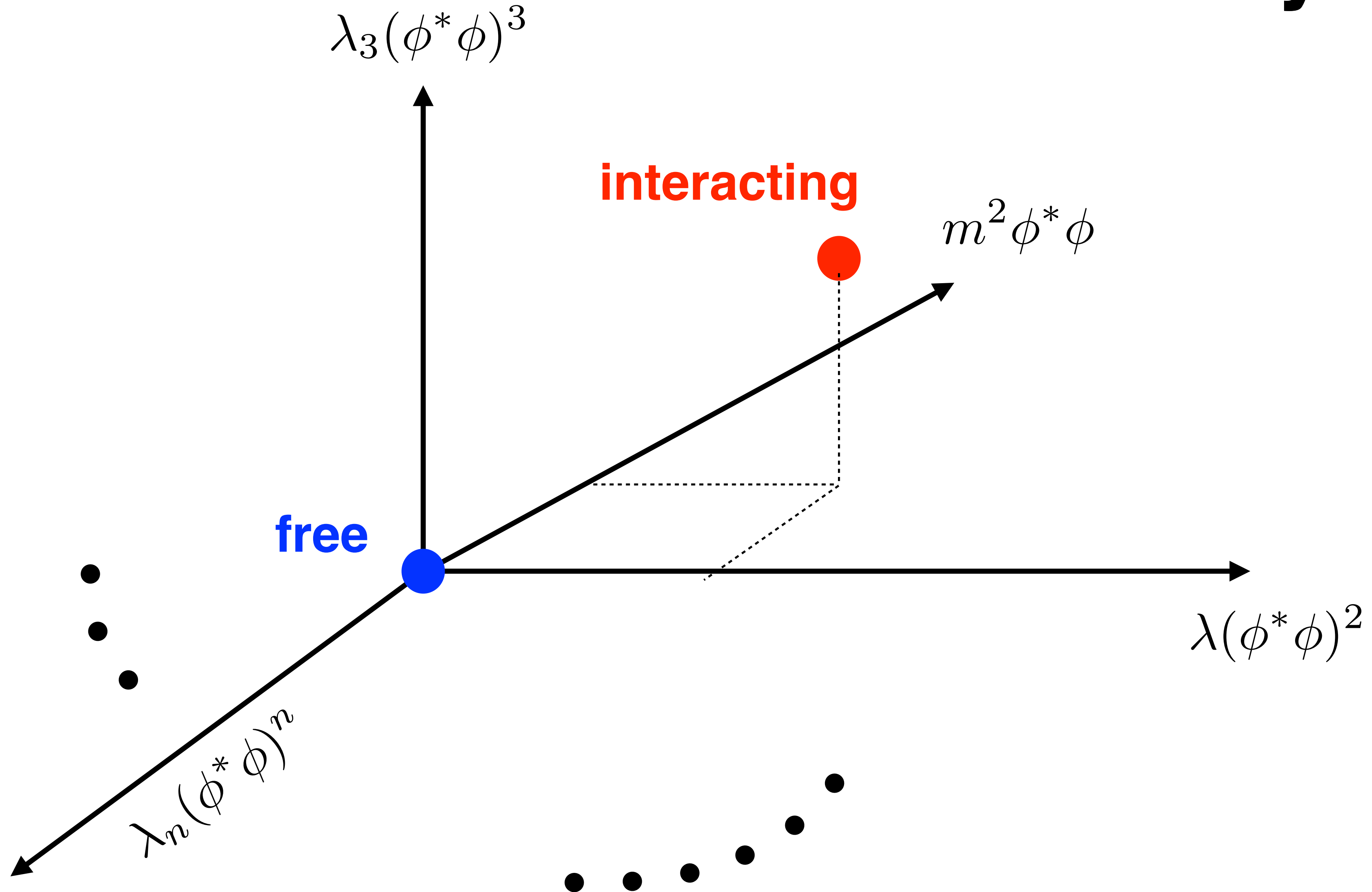
**IR**  
 $k = 0$



# “theory space”



# “theory space”



# “theory space”

## interactions

classical

$$m^2 \phi^* \phi$$

relevant

$$\lambda(\phi^* \phi)^2$$

relevant

---

$$\lambda_3(\phi^* \phi)^3$$

marginal

⋮

$$\lambda_n(\phi^* \phi)^n$$

irrelevant

⋮

UV

# “theory space”

## interactions

$$m^2 \phi^* \phi$$

classical

quantum

relevant

relevant

$$\lambda(\phi^* \phi)^2$$

relevant

$$\lambda_3(\phi^* \phi)^3$$

marginal

irrelevant

⋮

$$\lambda_n(\phi^* \phi)^n$$

irrelevant

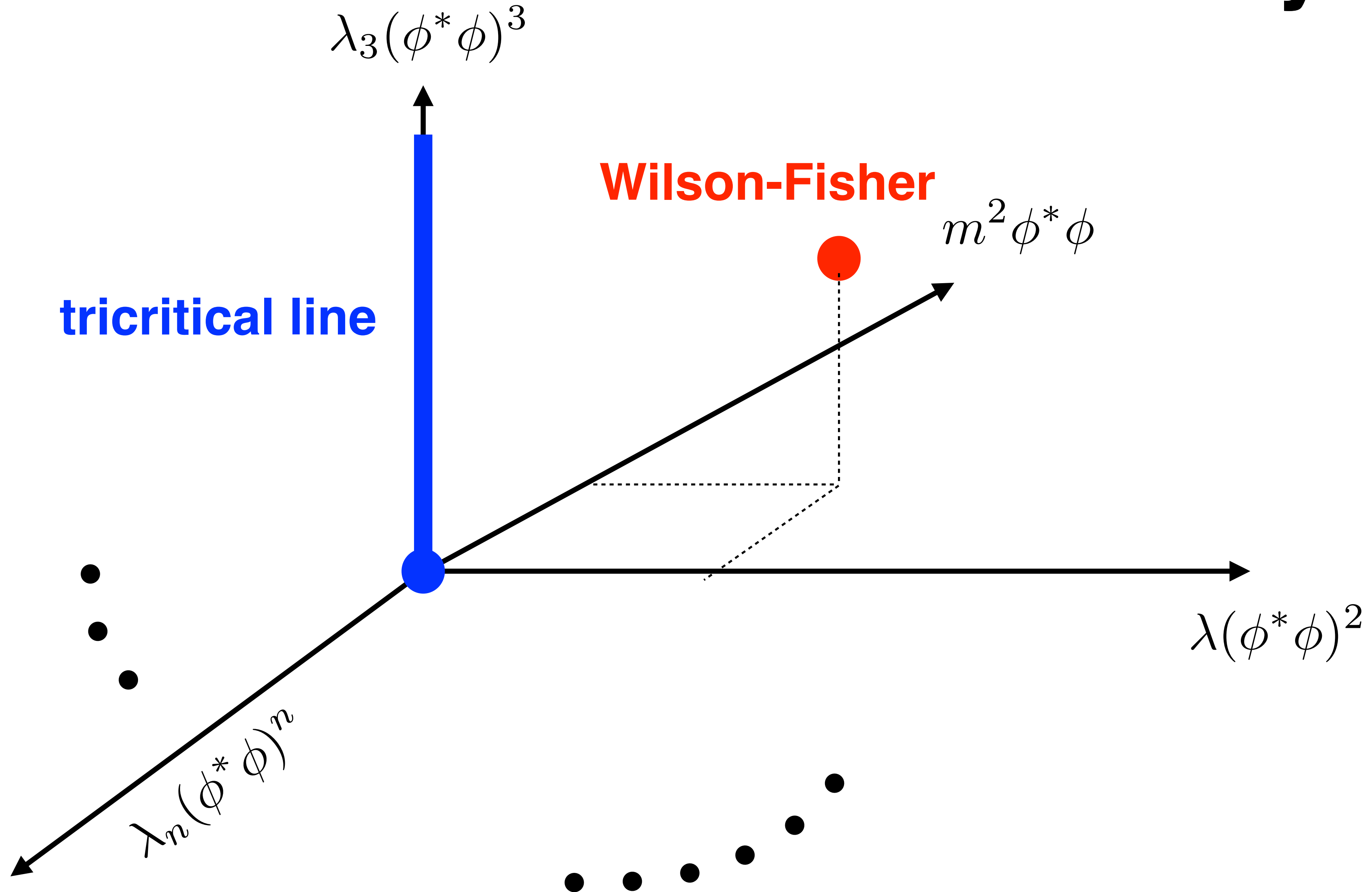
⋮

UV



IR

# “theory space”



# functional RG study

**Polchinski:  
UV cutoff**

$$P_{\text{UV}} = \frac{K(q^2/k^2)}{q^2}$$

Polchinski '84

**exact map  
R=R<sub>opt</sub>**

$$K(q^2/k^2) = \frac{R_k(q^2)}{q^2 + R_k(q^2)}$$

$$P_{\text{UV}} + P_{\text{IR}} = \frac{1}{q^2}$$

**Wetterich:  
IR cutoff**

$$P_{\text{IR}} = \frac{1}{q^2 + R_k(q^2)}$$

Wetterich '92

# functional RG study

**Polchinski:  
UV cutoff**

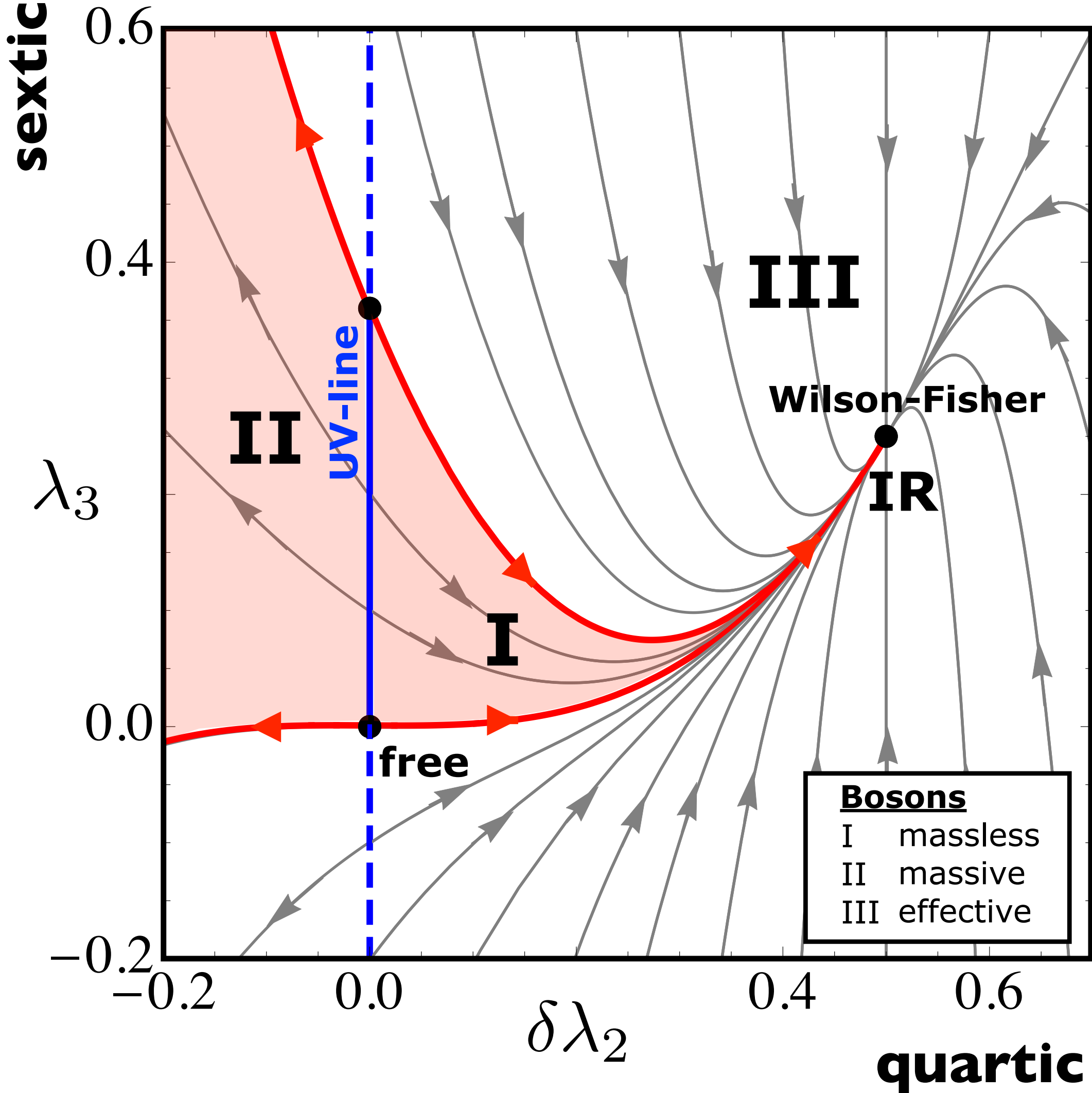
**exact map  
R=R<sub>opt</sub>**

$$\begin{cases} \partial_t u = -du + (d-2)\rho u' + 2\rho(u')^2 - (N-1)u' - (u' + 2\rho u'') \\ \partial_t w = -dw + (d-2)zw' + (N-1) \left( \frac{1}{1+w'} - 1 \right) + \left( \frac{1}{1+w' + 2zw''} - 1 \right) \end{cases}$$

**Wetterich:  
IR cutoff**

# functional RG

phase diagram





# functional RG

**BMB phenomenon**

Bardeen, Moshe, Bander '84  
David, Kessler, Neuberger '84

**spontaneous scale symmetry breaking**  
breaking of hyperscaling

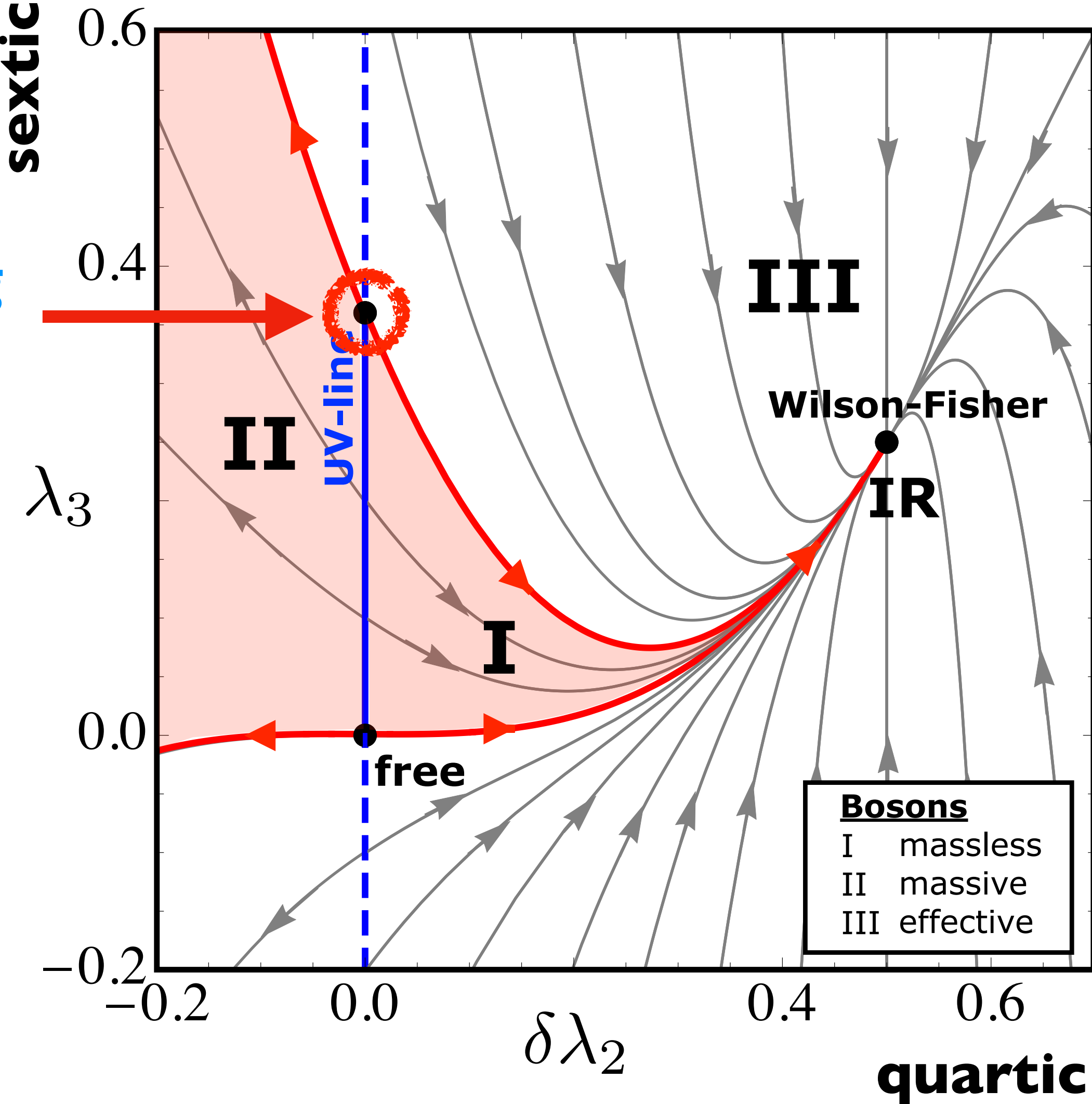


**compact conformal manifold**  
physical mass = free parameter



**non-perturbative**  
**infinite-order in local couplings**

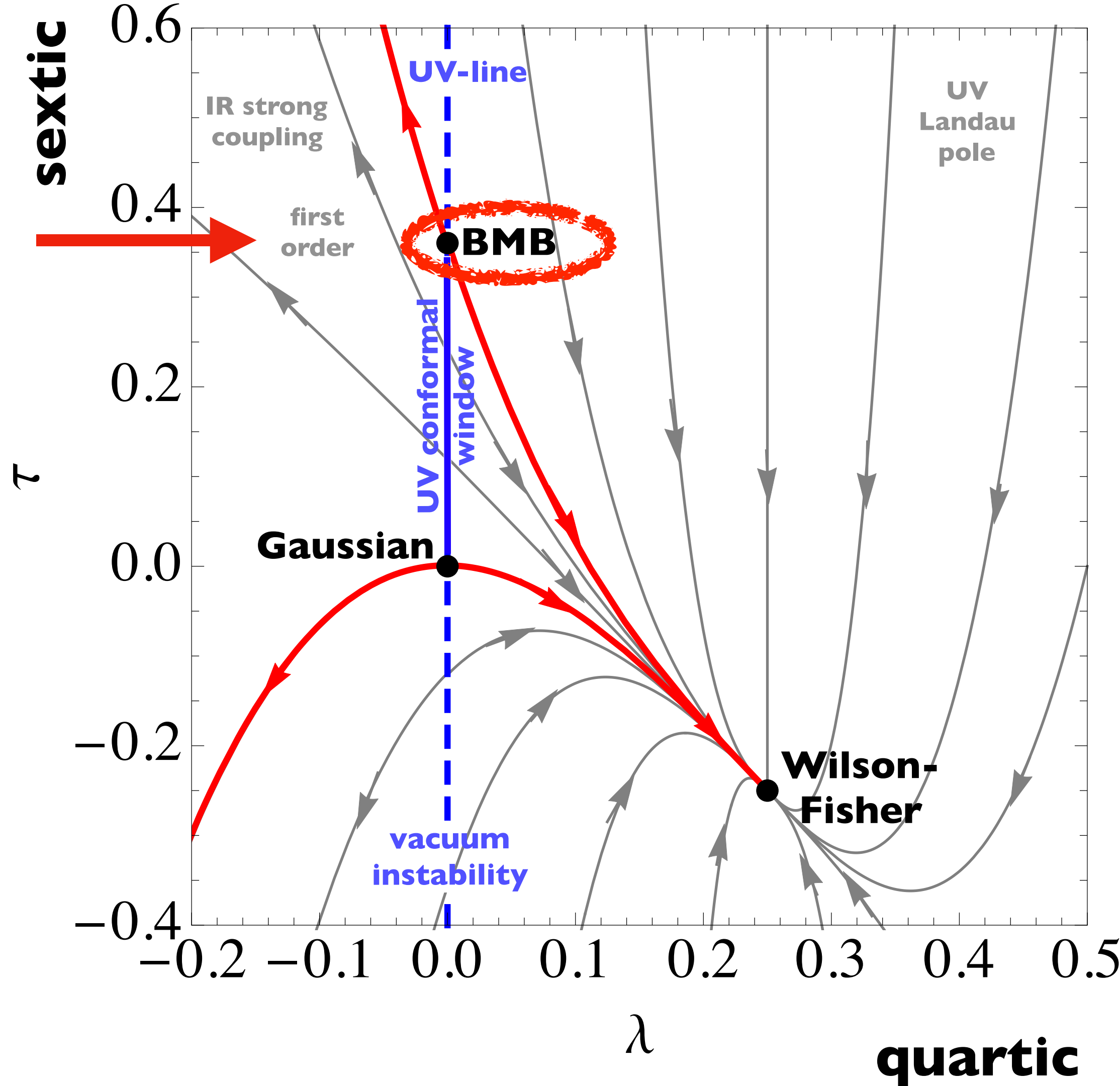
**phase diagram**



# Polchinski RG

spontaneous scale symmetry breaking  
breaking of hyperscaling

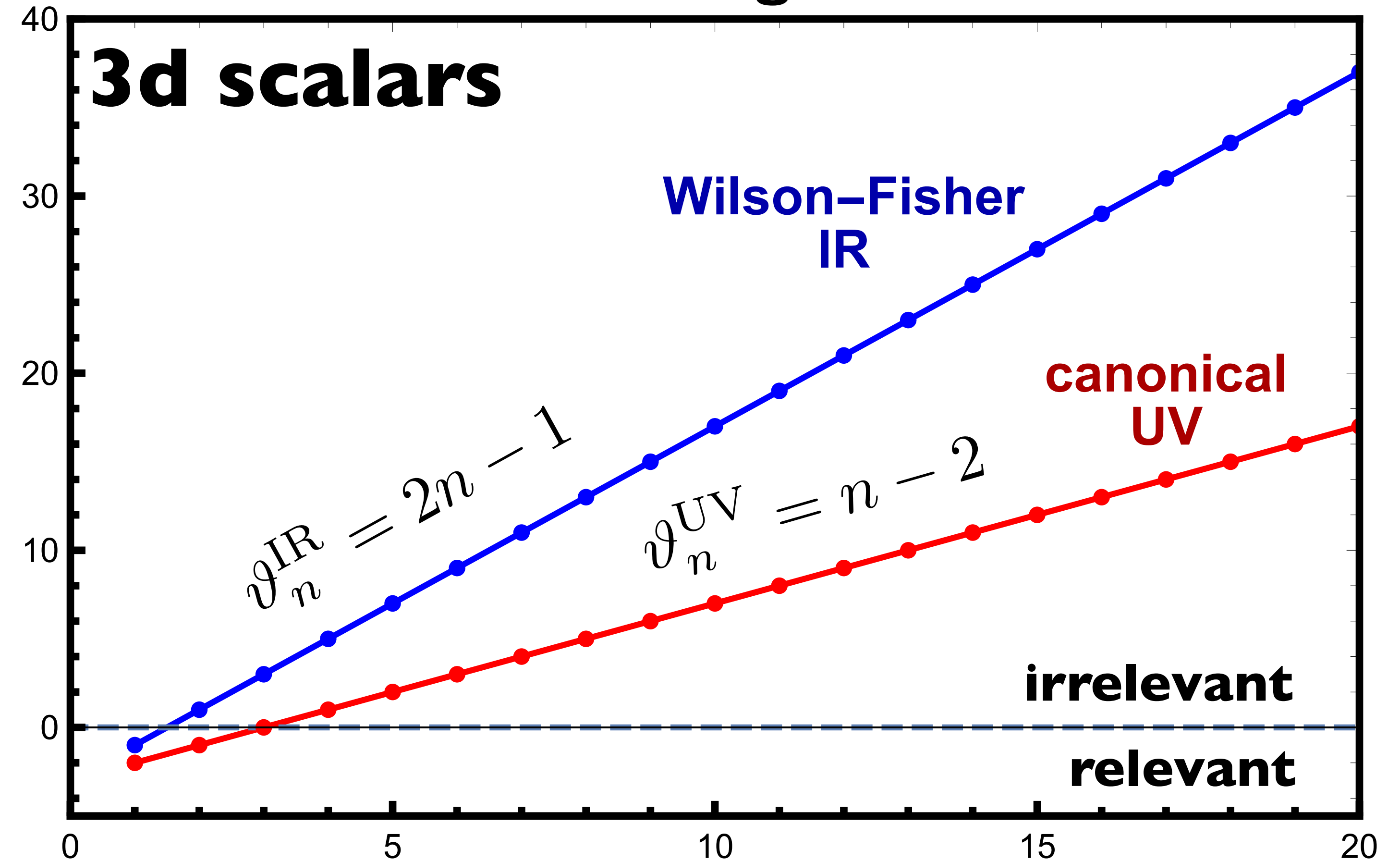
phase diagram



large quantum effects

$$\frac{\nu_n^{(\text{IR})} - \nu_n^{(\text{UV})}}{\nu_n^{(\text{IR})}} = \frac{n+1}{2n-1}$$

## universal scaling dimensions



# Gross-Neveu

U(N) symmetric fermions

4-fermion interactions  $G(\bar{\psi}\psi)^2$

Gross, Neveu '74

chiral symmetry  $\psi \rightarrow \gamma^5 \psi \quad \bar{\psi} \rightarrow -\bar{\psi} \gamma^5$

**3d:** perturbatively non-renormalisable...  $[G] = 2 - d$

**...yet** non-perturbatively renormalisable  
**interacting UV fixed point**

functional RG

Gawedzki, Kupiainen '85  
Rosenstein, War, Park '89  
de Calan, Faria da Veiga, Magnen, de Seneor '91

Jakovac, Patkos '13, '14  
Cresswell-Hogg, Litim, '22, '23 and in prep '24


# Gross-Neveu+

relax chiral symmetry

$$S = \int d^d x \left\{ \bar{\psi}_a \not{\partial} \psi_a + \frac{1}{2} G (\bar{\psi}_a \psi_a)^2 + \frac{1}{3!} H (\bar{\psi}_a \psi_a)^3 \right\}.$$

mass term permitted  $m \bar{\psi} \psi$

6-fermion interactions permitted

$$[H] = 3 - 2d$$


functional RG

exactly solvable at infinite Nf

# “theory space”

**interactions**

classical

$$\lambda_1 \bar{\psi}\psi$$

**relevant**

---

$$\lambda_2 (\bar{\psi}\psi)^2$$

**irrelevant**

$$\lambda_3 (\bar{\psi}\psi)^3$$

▪

▪

$$\lambda_n (\bar{\psi}\psi)^n$$

▪

▪

▪

▪

**IR**

# “theory space”

## interactions

classical

quantum

$$\lambda_1 \bar{\psi}\psi$$

relevant

relevant

---

$$\lambda_2 (\bar{\psi}\psi)^2$$

irrelevant

relevant

$$\lambda_3 (\bar{\psi}\psi)^3$$

·

marginal

$$\lambda_n (\bar{\psi}\psi)^n$$

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irrelevant

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IR

·

·

UV

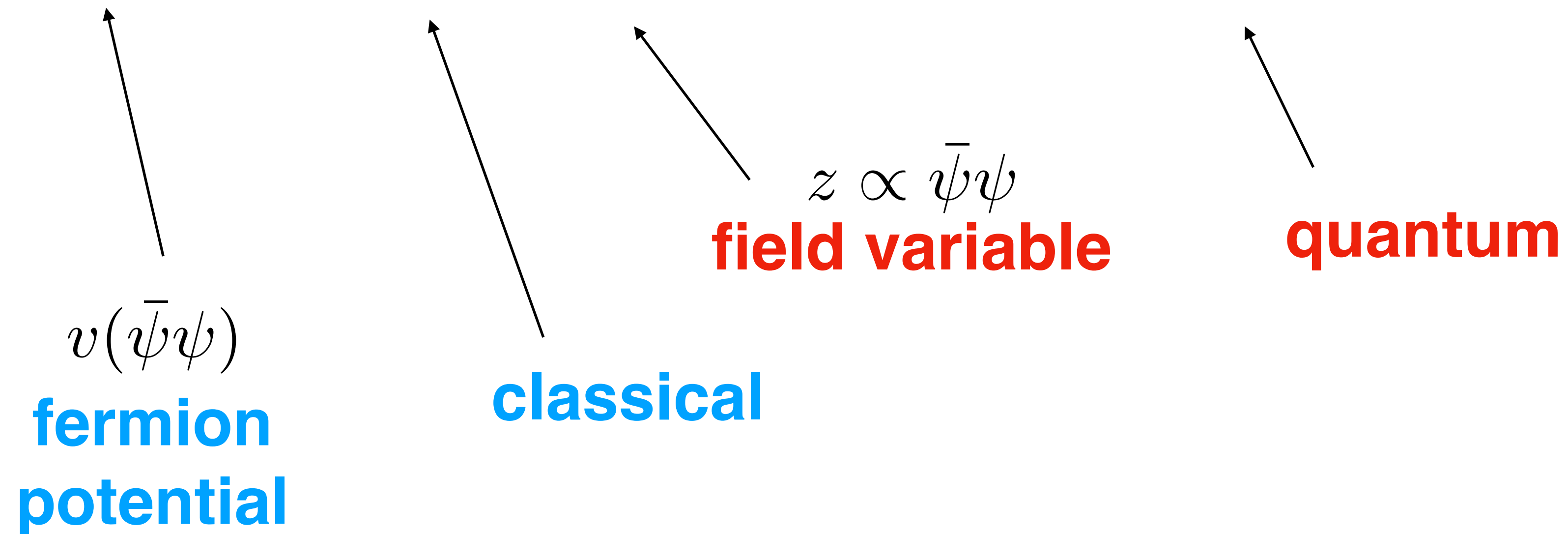
# Gross-Neveu+

functional RG

$$\Gamma_k = \int d^d x \left( \sum_{i=1}^{N_f} \bar{\Psi}_i \not{\partial} \Psi_i + V_k(\Psi, \bar{\Psi}) \right)$$

local potential

$$\partial_t v = -dv + (d-1)zv' - (4N_f + 1) \ell_d[(v')^2] + \ell_d[v' \cdot (v' + 2zv'')],$$





# Gross-Neveu+

large Nf:

$$v(z) = \sum_n \frac{\lambda_n}{n!} z^n$$

**mass**  $\beta_1 = -\lambda_1 + \frac{2\lambda_1\lambda_2}{(1 + \lambda_1^2)^2}$

**4F**  $\beta_2 = (d - 2)\lambda_2 + \frac{2\lambda_1\lambda_3}{(1 + \lambda_1^2)^2} + \frac{(2 - 6\lambda_1^2)\lambda_2^2}{(1 + \lambda_1^2)^3}$

**6F**  $\beta_3 = (2d - 3)\lambda_3 + \frac{2\lambda_1\lambda_4}{(1 + \lambda_1^2)^2} + \frac{6\lambda_2\lambda_3(1 - 3\lambda_1^2)}{(1 + \lambda_1^2)^3} + \frac{24\lambda_1\lambda_2^3(\lambda_1^2 - 1)}{(1 + \lambda_1^2)^4}$ .

**mass=0** is an exact RG fixed point

generation of mass

# Gross-Neveu+

**mass=0:**

**4F fixed point**

**4F**  $\tilde{\beta}_2 = (d - 2 + 2\lambda_2)\lambda_2,$

**6F**  $\tilde{\beta}_3 = (2d - 3 + 6\lambda_2)\lambda_3,$

**mass = 0** renders 4F and 6F betas homogeneous

# Gross-Neveu+

**mass=0:**

**4F fixed point**

**4F**  $\tilde{\beta}_2 = (d - 2 + 2\lambda_2)\lambda_2,$

**6F**  $\tilde{\beta}_3 = (2d - 3 + 6\lambda_2)\lambda_3,$

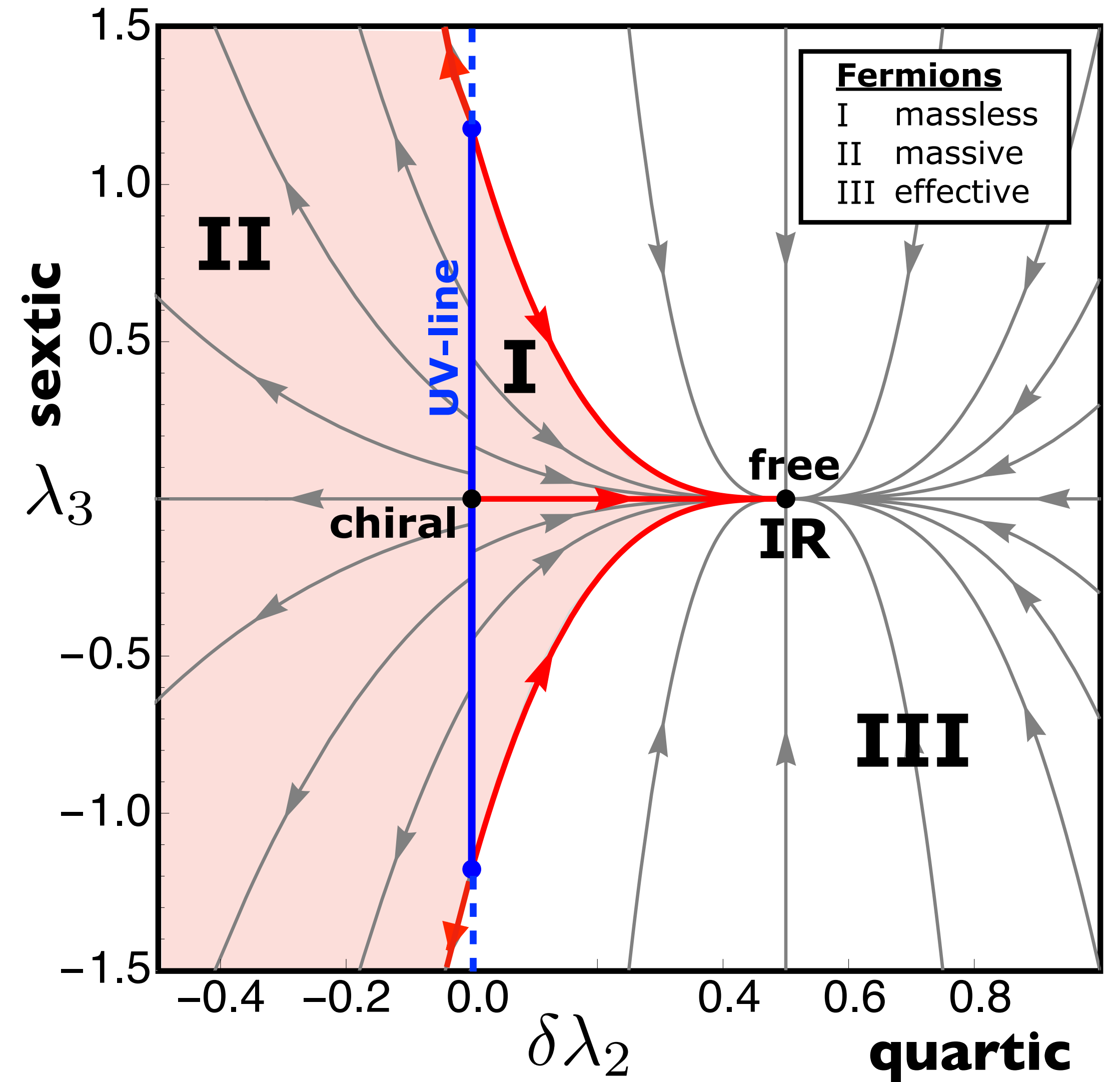
**3d:** **4F fixed point** renders 6F coupling **exactly marginal**

scheme independent

# Gross-Neveu+

UV-IR  
connecting  
trajectories

exactly marginal  
sextic coupling

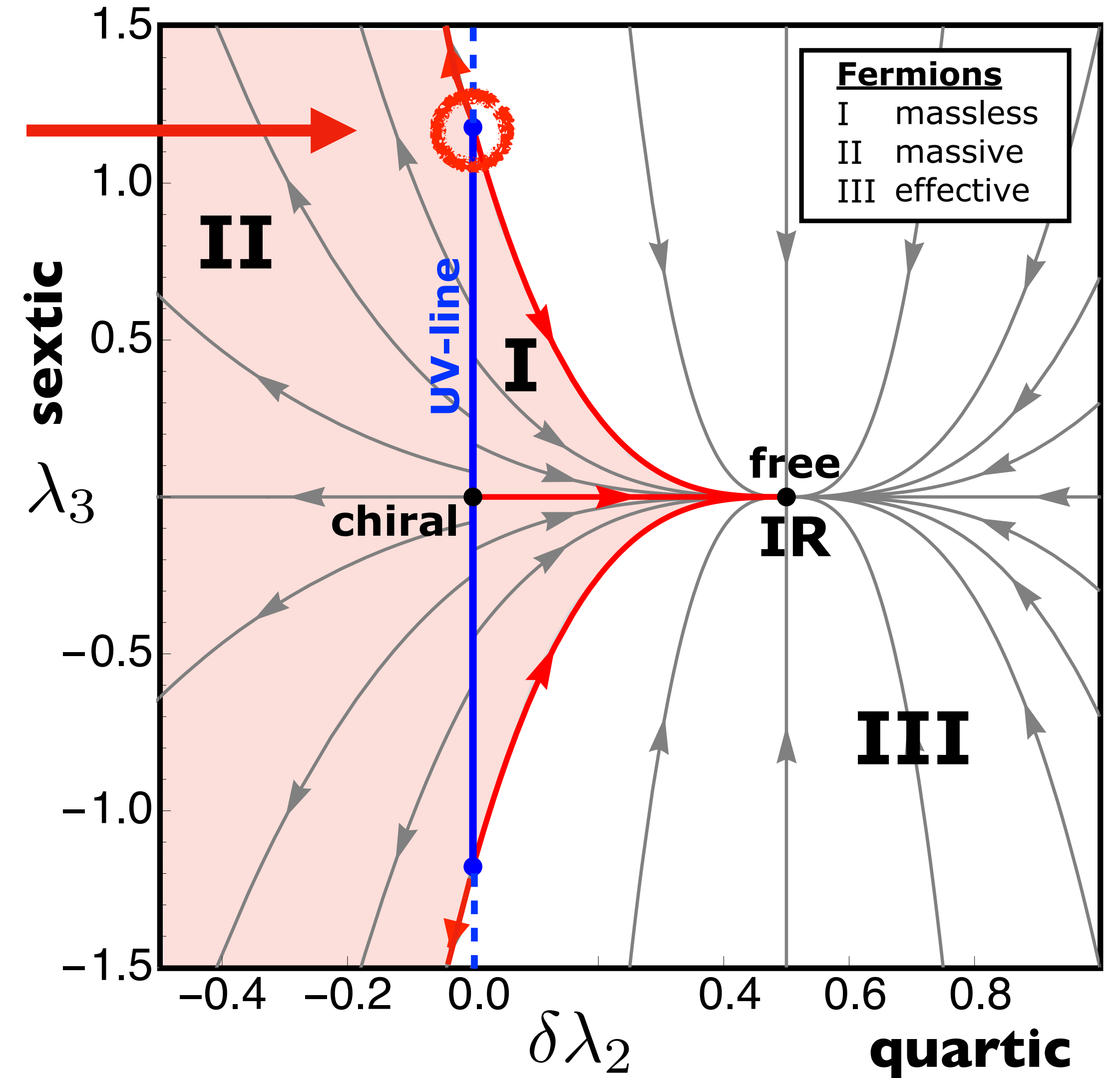


# Gross-Neveu+

spontaneous scale symmetry breaking  
breaking of hyperscaling

compact conformal manifold  
physical mass = free parameter

non-perturbative  
infinite-order in local couplings





# Gross-Neveu+

global fixed points

$$v' = v'(z) \text{ for all } z$$

$$z=0: v' = 0$$

**mass=0**

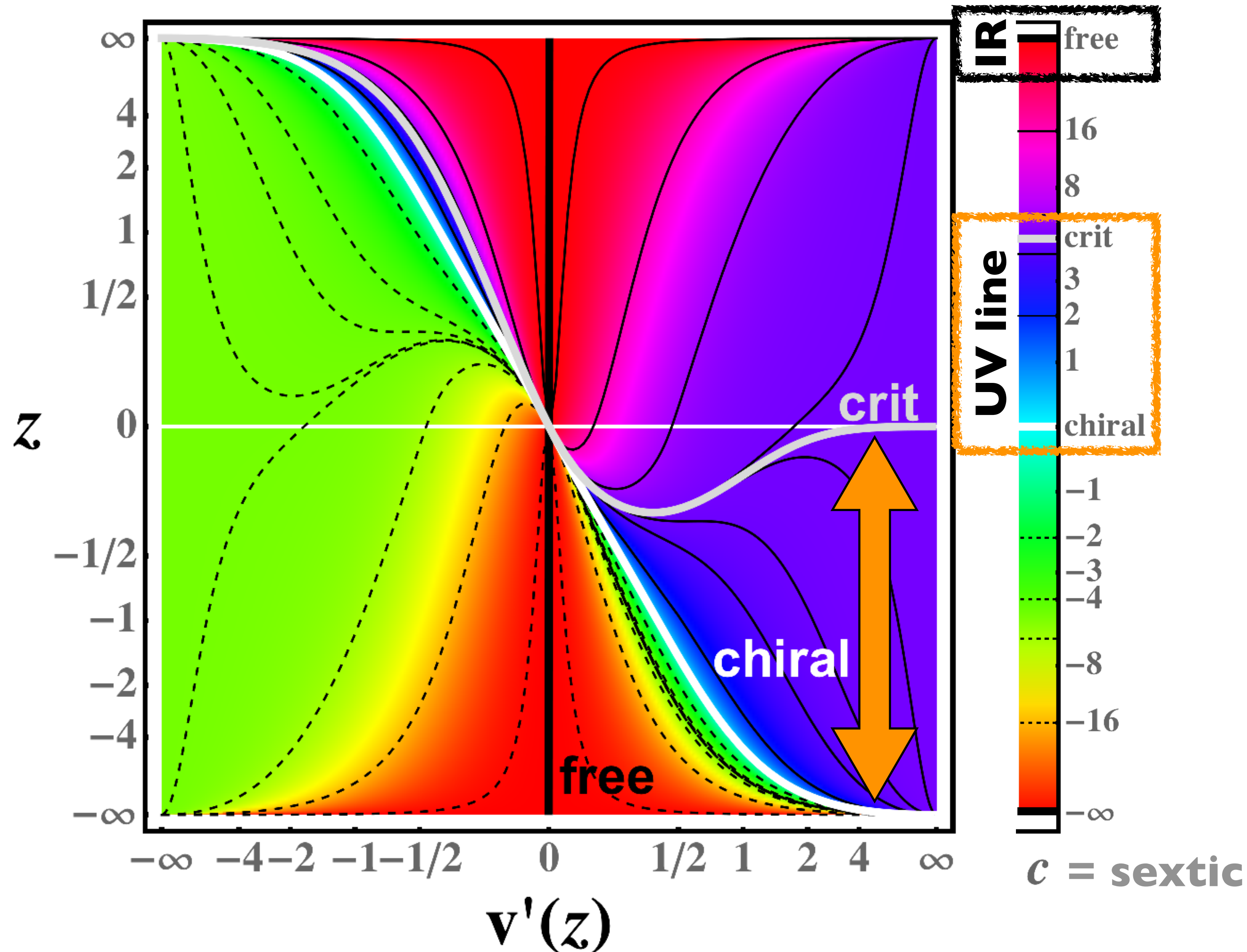
$$v'' = -1/2$$

**4F**

$$v''' = c$$

**6F**

finite UV conformal manifold



# Gross-Neveu+

spontaneous generation of mass

$$M = \lim_{k \rightarrow 0} k \cdot v'(0)$$

gap equation

$$\left[ c - \frac{3\pi}{2} \operatorname{sgn}(M) \right] M^2 = 0$$

two solutions

$$[\dots] = 0$$

M = free parameter

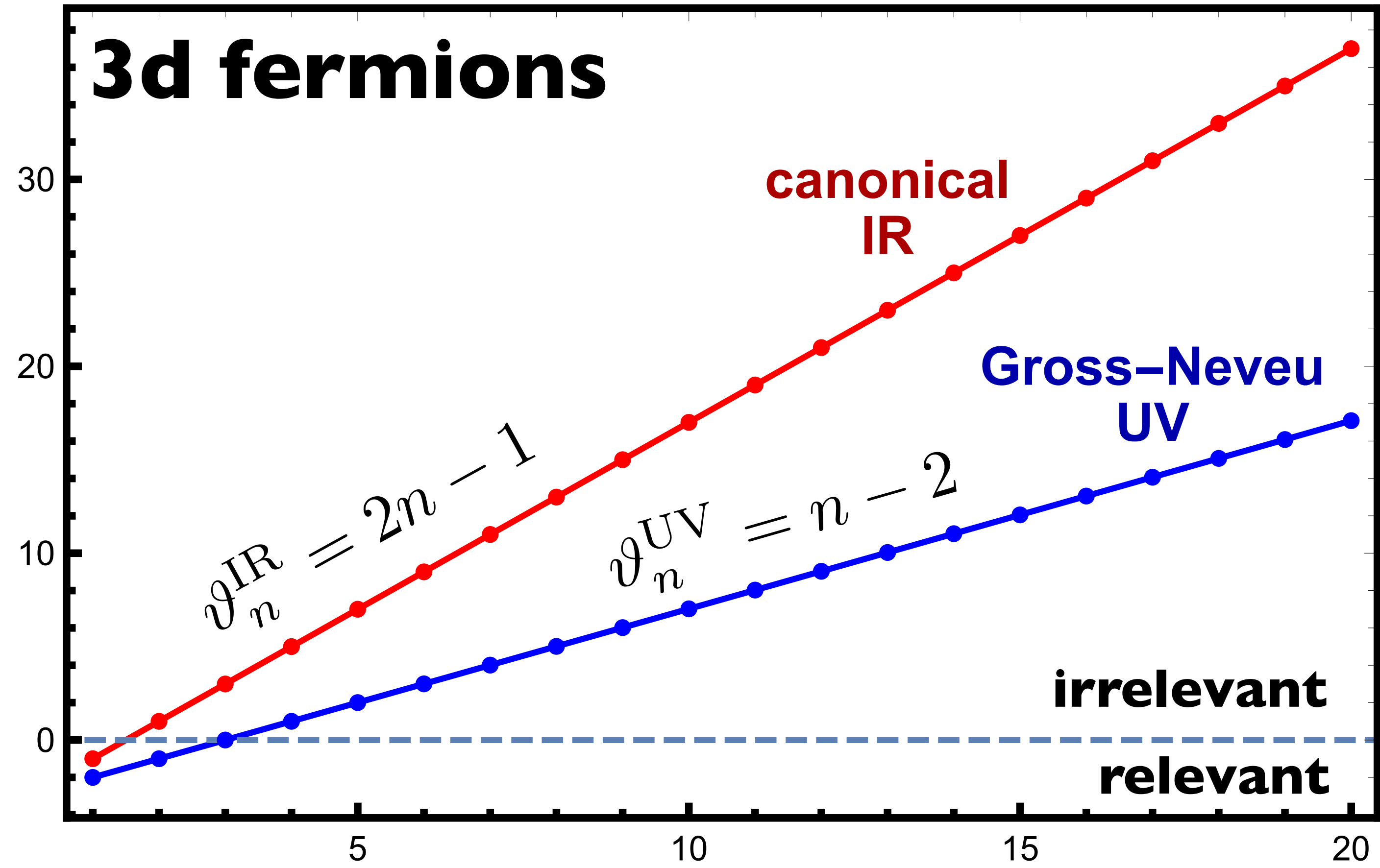
$$M = 0$$



scale symmetry broken **spontaneously**

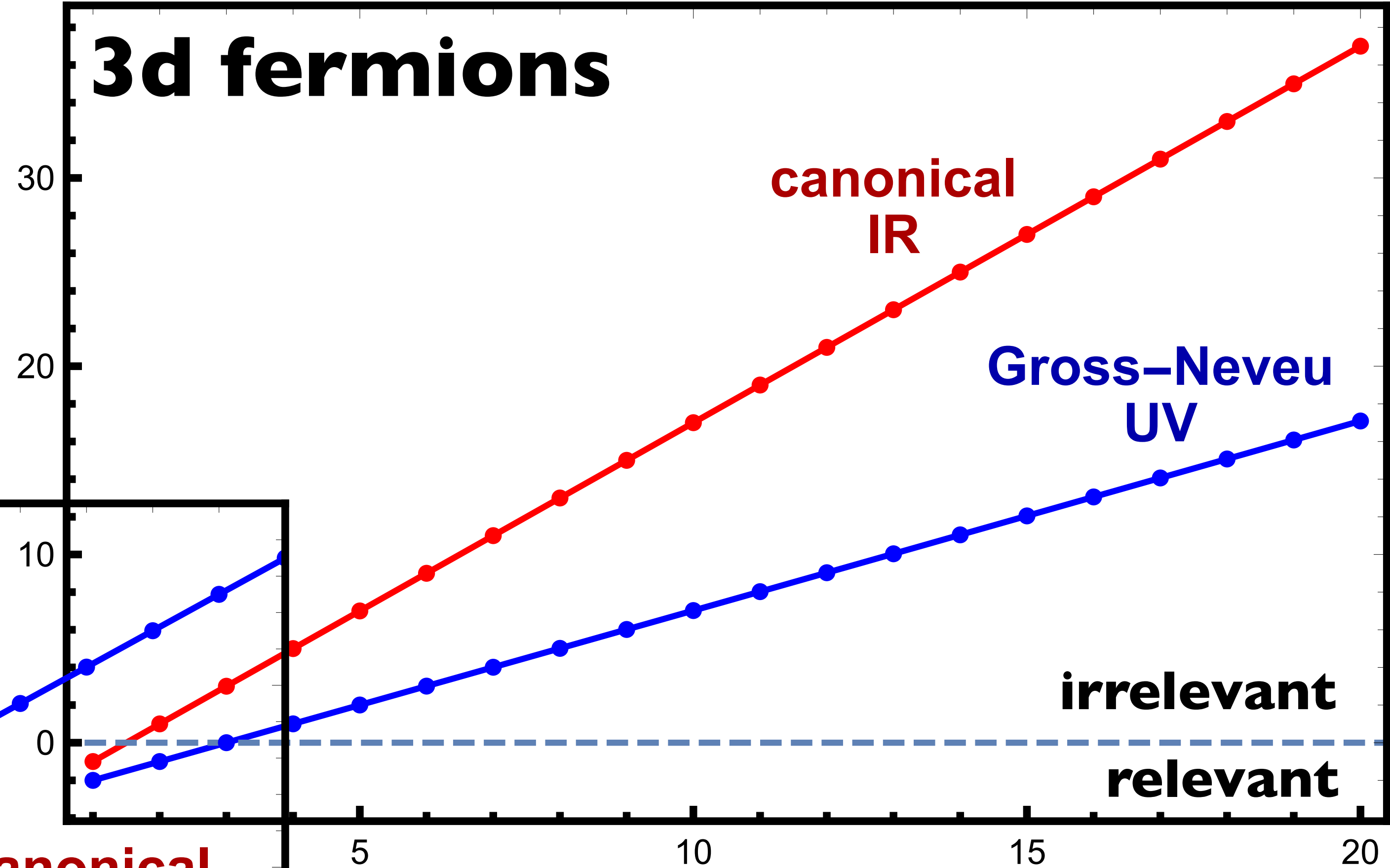
prerequisite: **6F interactions**, hence **no** chiral symmetry breaking

# universal scaling dimensions

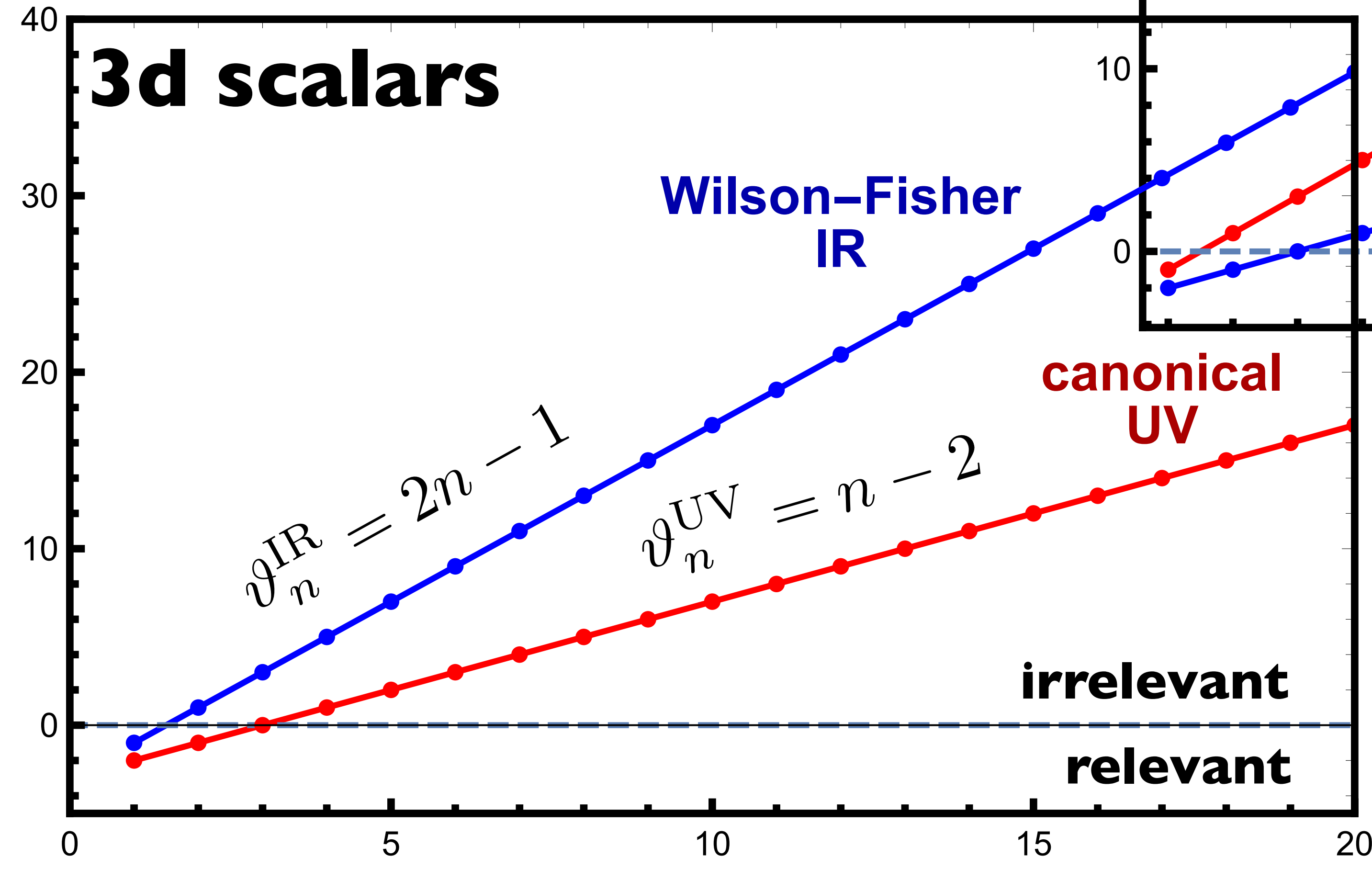




# 3d fermions



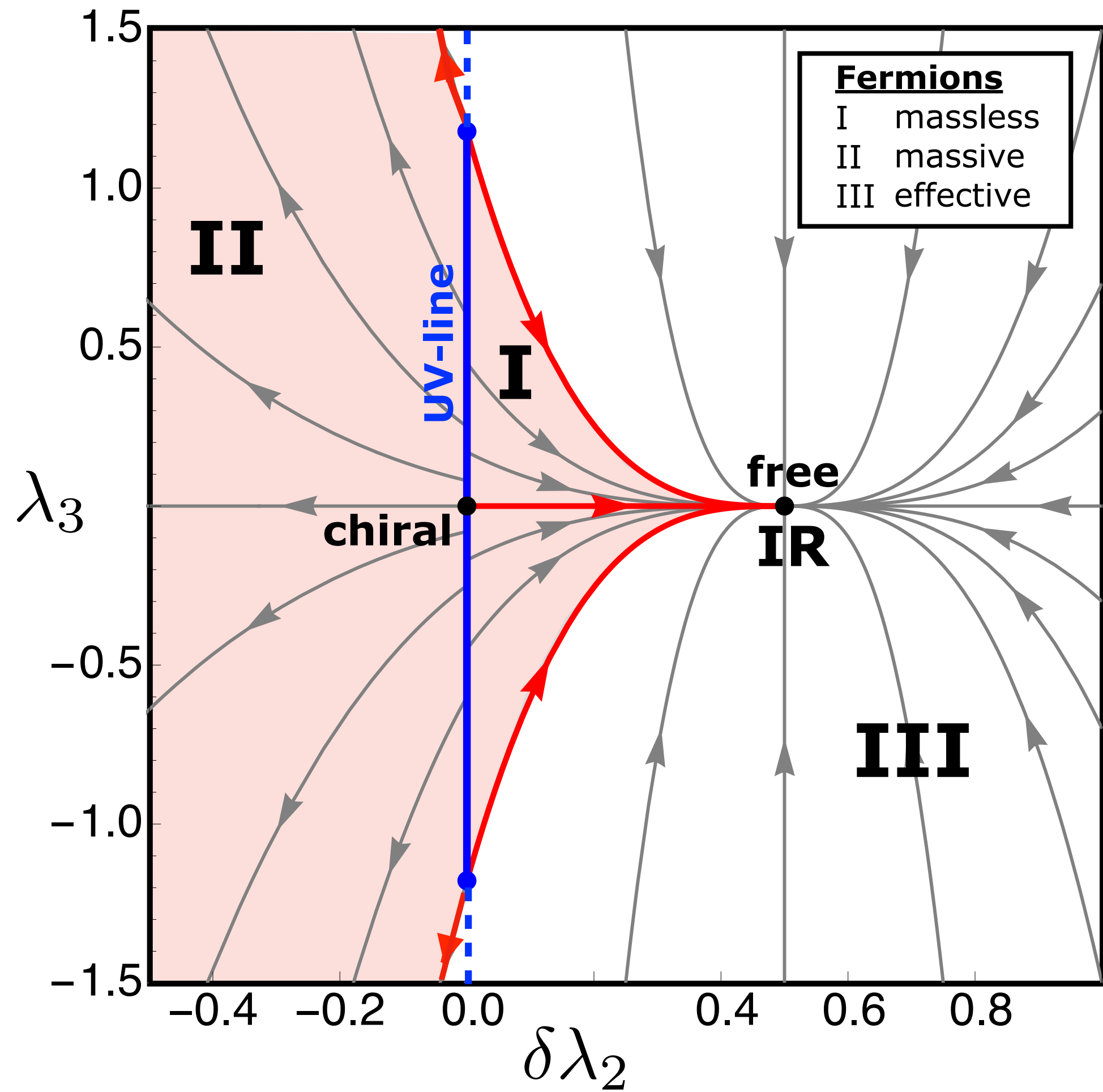
# 3d scalars



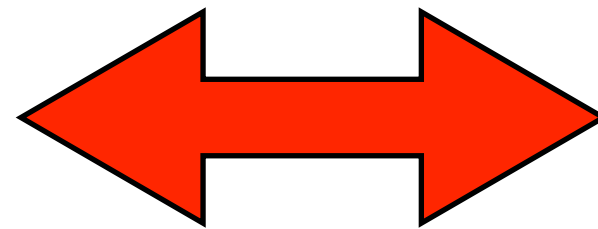
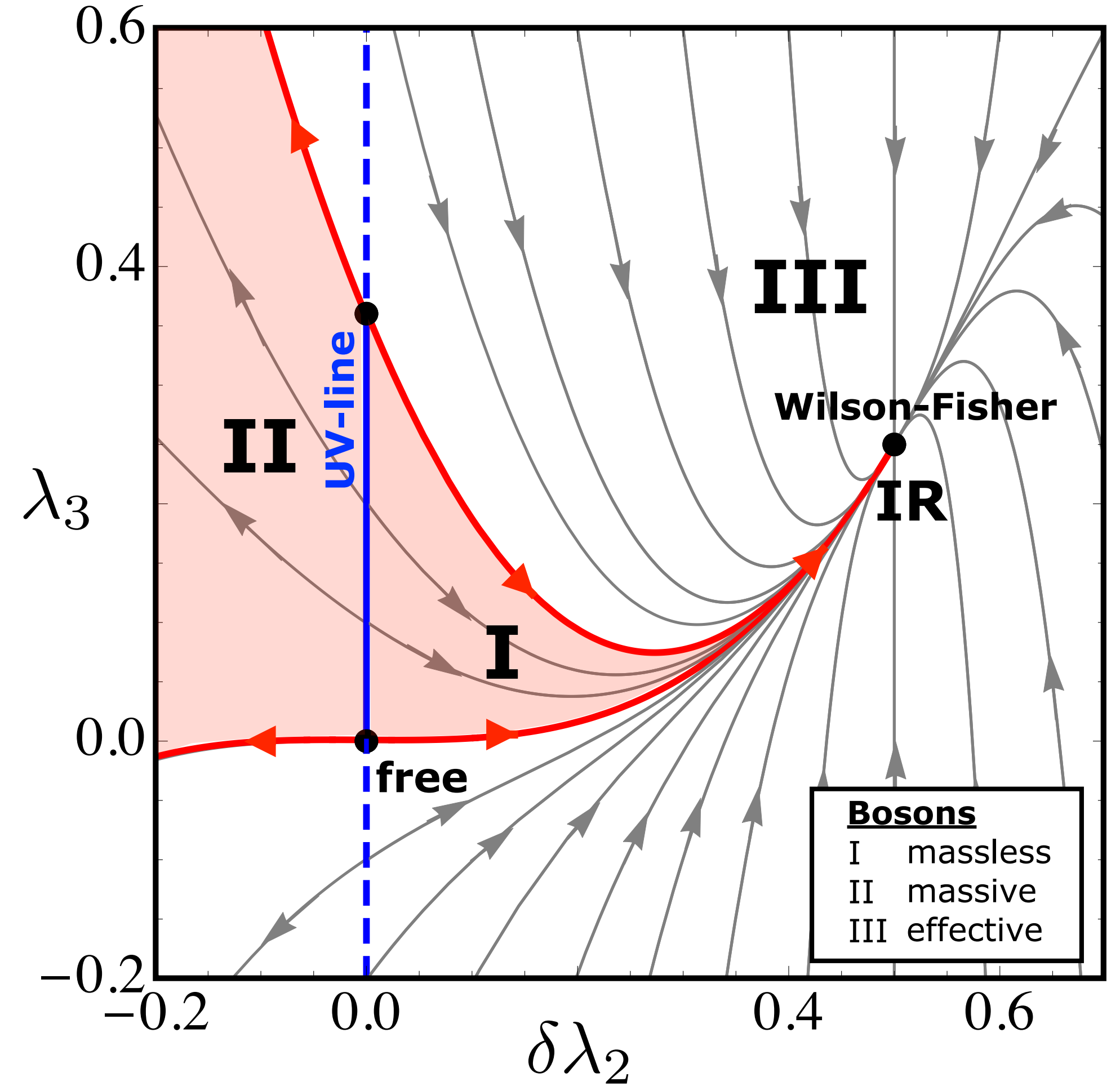
**irrelevant**  
**relevant**

**irrelevant**  
**relevant**

# Fermions



# Bosons



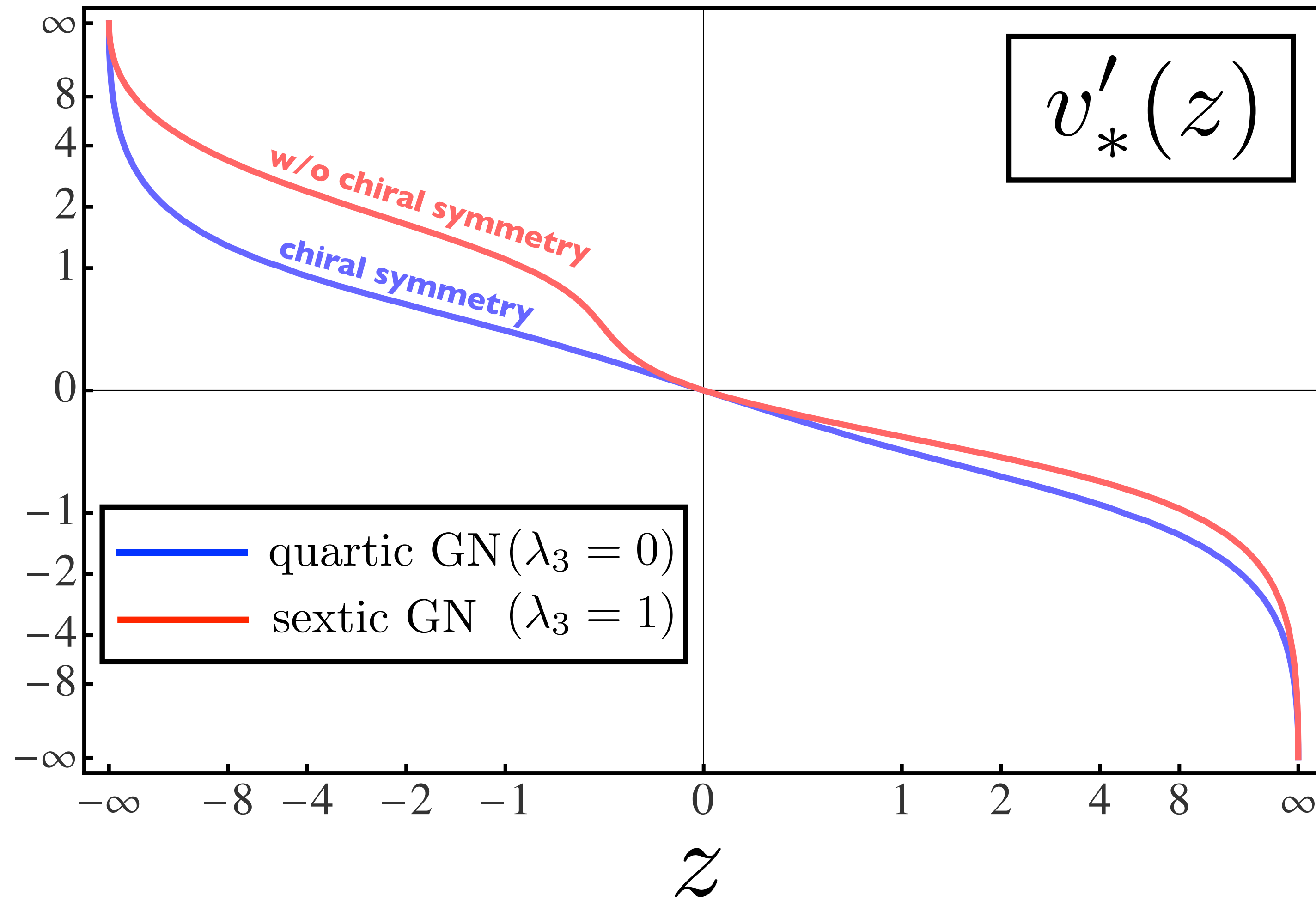
**AdS/CFT**  
**higher spin gauge theories**

**Chern-Simons-Matter**

Klebanov, Polyakov, hep-th/0210114, Szegin, Sundell, hep-th/0305040  
 Maldacena, Zhiboedov, 1112.1016, 1204.3882  
 Giombi, Zin, 1208.4036

Aharony, Giombi, Gur-Ari, Maldacena, Yacoby, 1211.4843  
 Seiberg, Senthil, Wang, Witten, 1606.01989

# non-chiral **CFT** duals for higher-spin GTs on AdS4



# Gross-Neveu—Yukawa

$$S_{\text{GNY}} = \int_x \left\{ \bar{\psi}_a \not{\partial} \psi_a + \frac{1}{2} (\partial\phi)^2 + H\phi \bar{\psi}_a \psi_a + U(\phi) \right\}$$

Yukawa

even / odd

perturbatively renormalisable

“integrate-out the scalar”

PT, large Nf

functional RG

Gracey '90  
Moshe Moshe, Zinn-Justin '03

Gies, MM Scherer '10  
Braun, Gies, DD Scherer '12

**relax chiral symmetry**

odd phi powers + fermion mass term permitted

functional RG

exactly solvable at infinite Nf

# Gross-Neveu — Yukawa

large  $N_f$

IR fixed point

**mass**

$$\partial_t m_F = -m_F \left( 1 + \frac{2h^2}{(1 + m_F^2)^2 \lambda_2} \right)$$

**Yukawa**

$$\partial_t h = -\frac{1}{2} (1 - \eta_\phi) h$$

$$\eta_\phi = \frac{5}{2} h^2$$

**cubic**

$$\partial_t \lambda_3 = -\frac{3}{2} (1 - \eta_\phi) \lambda_3$$

# Gross-Neveu—Yukawa

large Nf

IR fixed point

**mass = 0** is an exact fixed point

**mass**

$$\partial_t m_F = -m_F \left( 1 + \frac{2h^2}{(1 + m_F^2)^2 \lambda_2} \right)$$

**Yukawa**

$$\partial_t h = -\frac{1}{2} (1 - \eta_\phi) h$$

$$\eta_\phi = \frac{5}{2} h^2$$

**cubic**

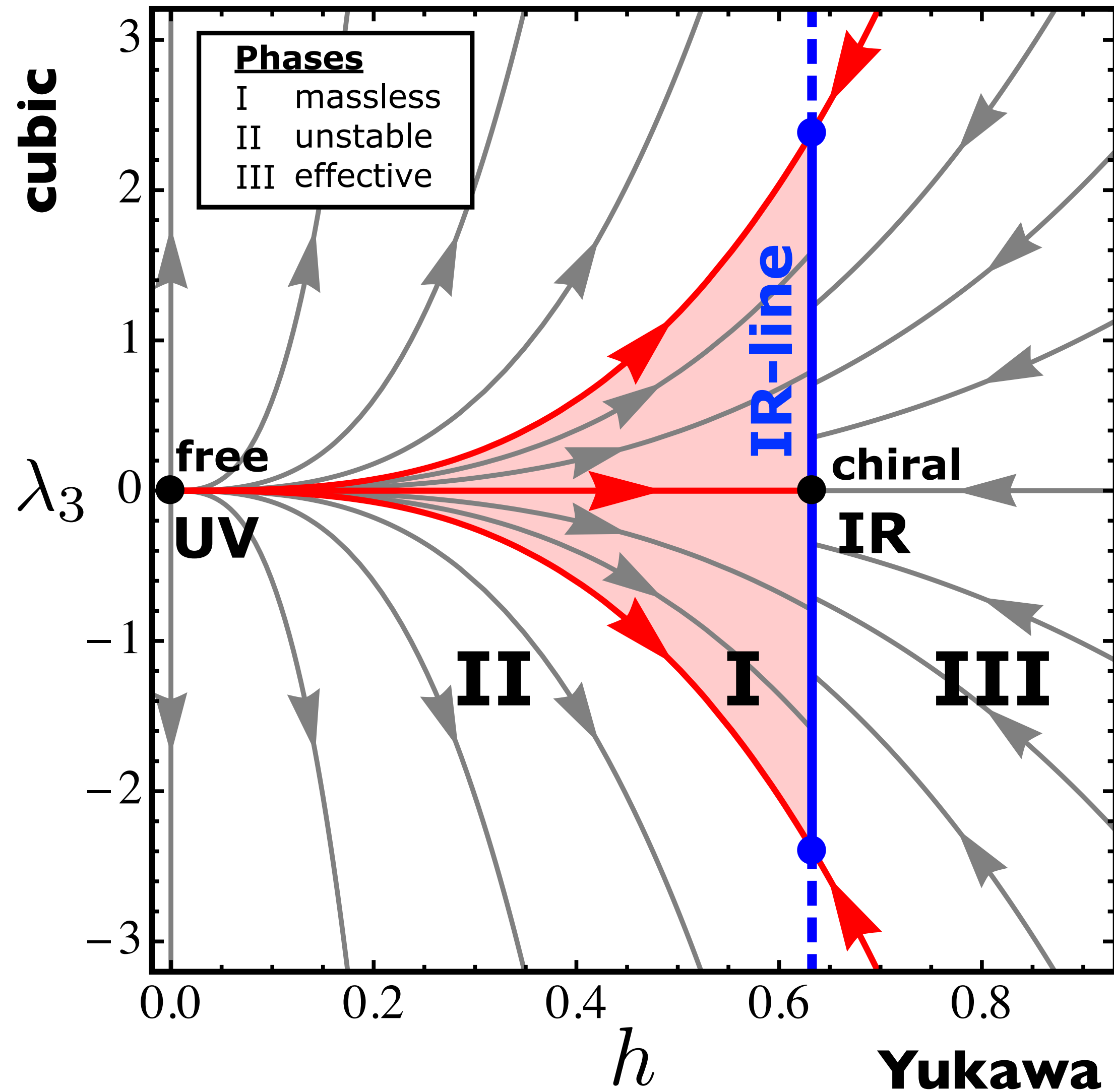
$$\partial_t \lambda_3 = -\frac{3}{2} (1 - \eta_\phi) \lambda_3$$

Yukawa fixes **eta = 1**

**eta = 1** renders the cubic **exactly marginal**

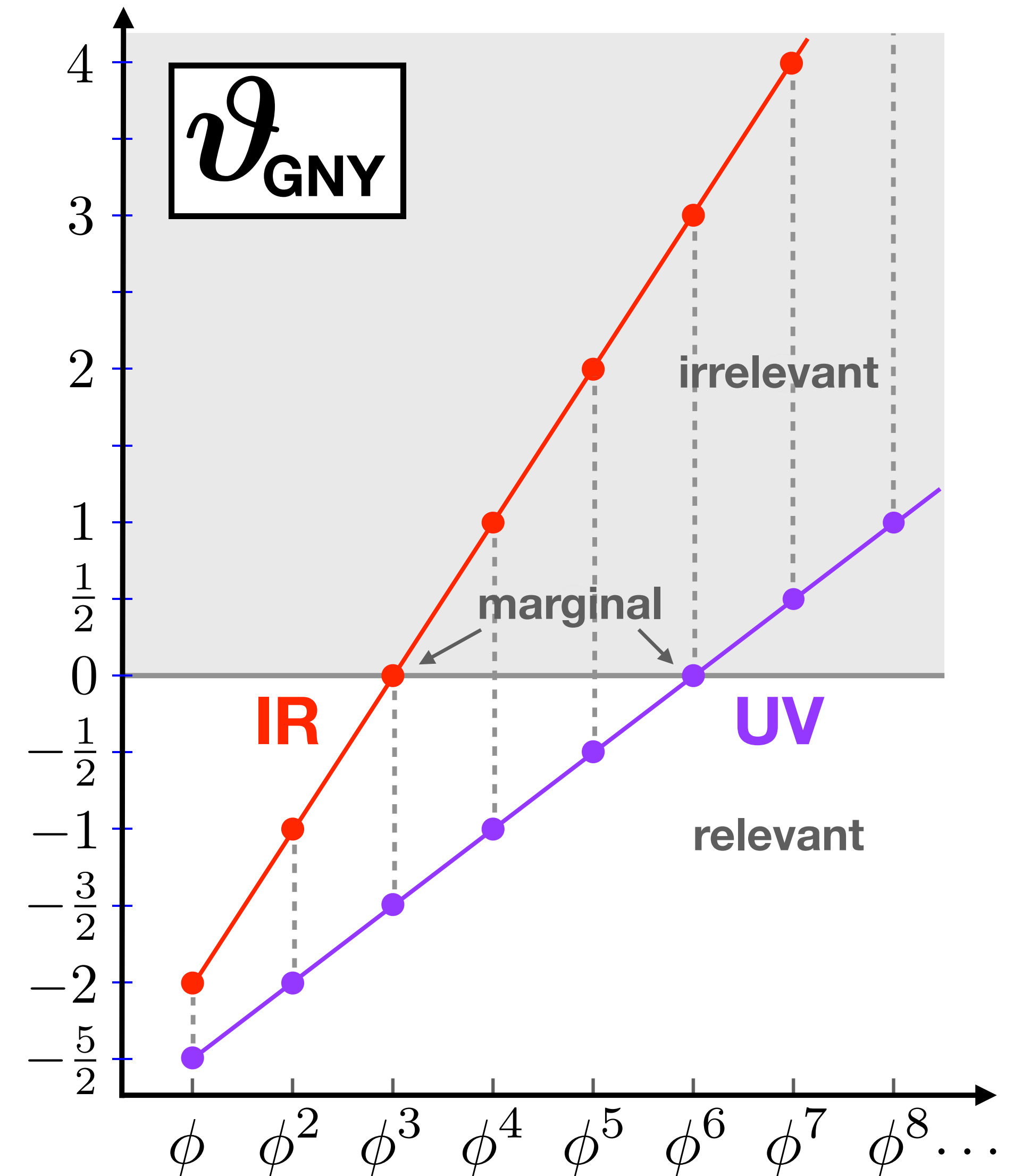
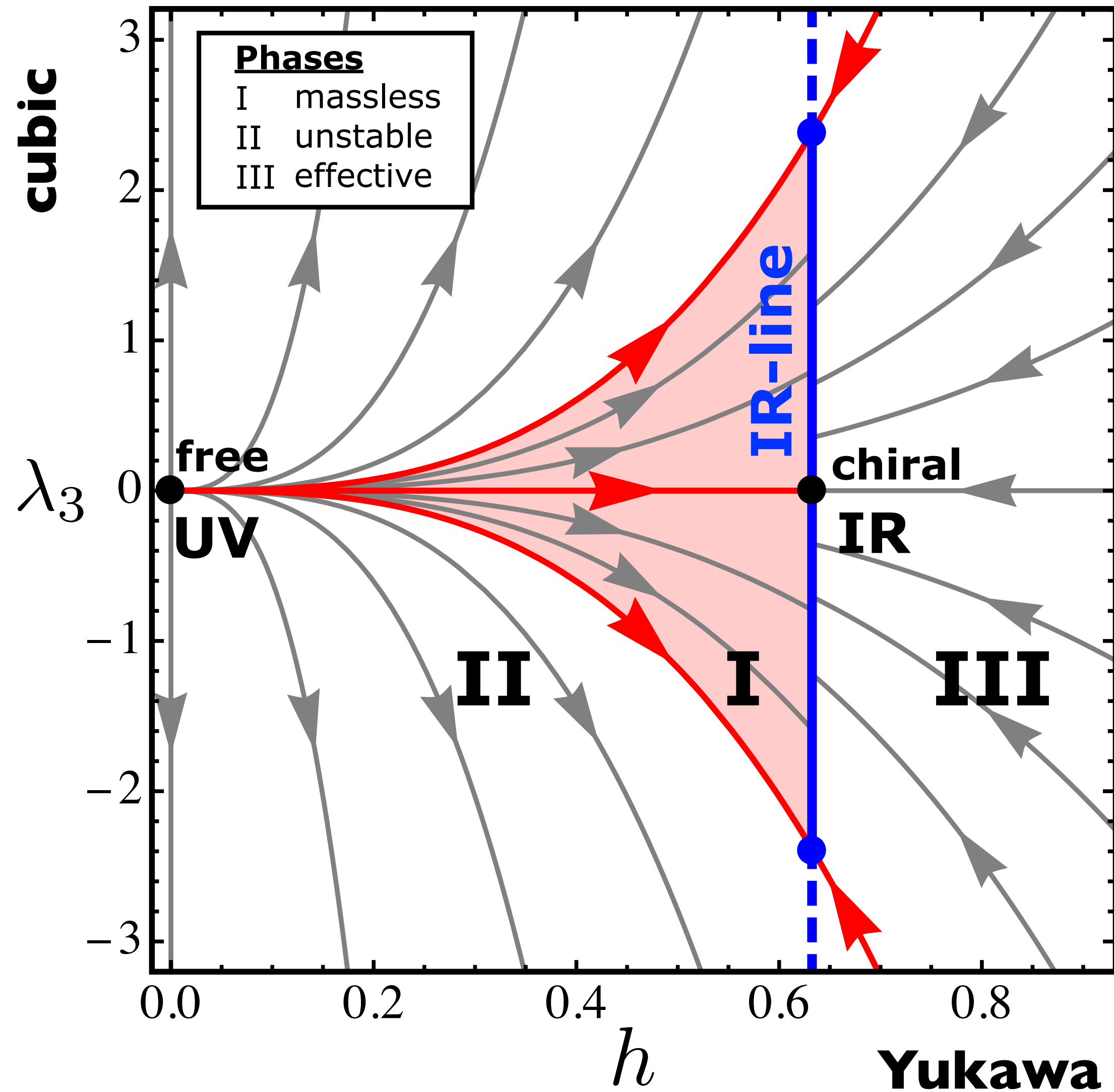
**→** conformal manifold

# Gross-Neveu – Yukawa



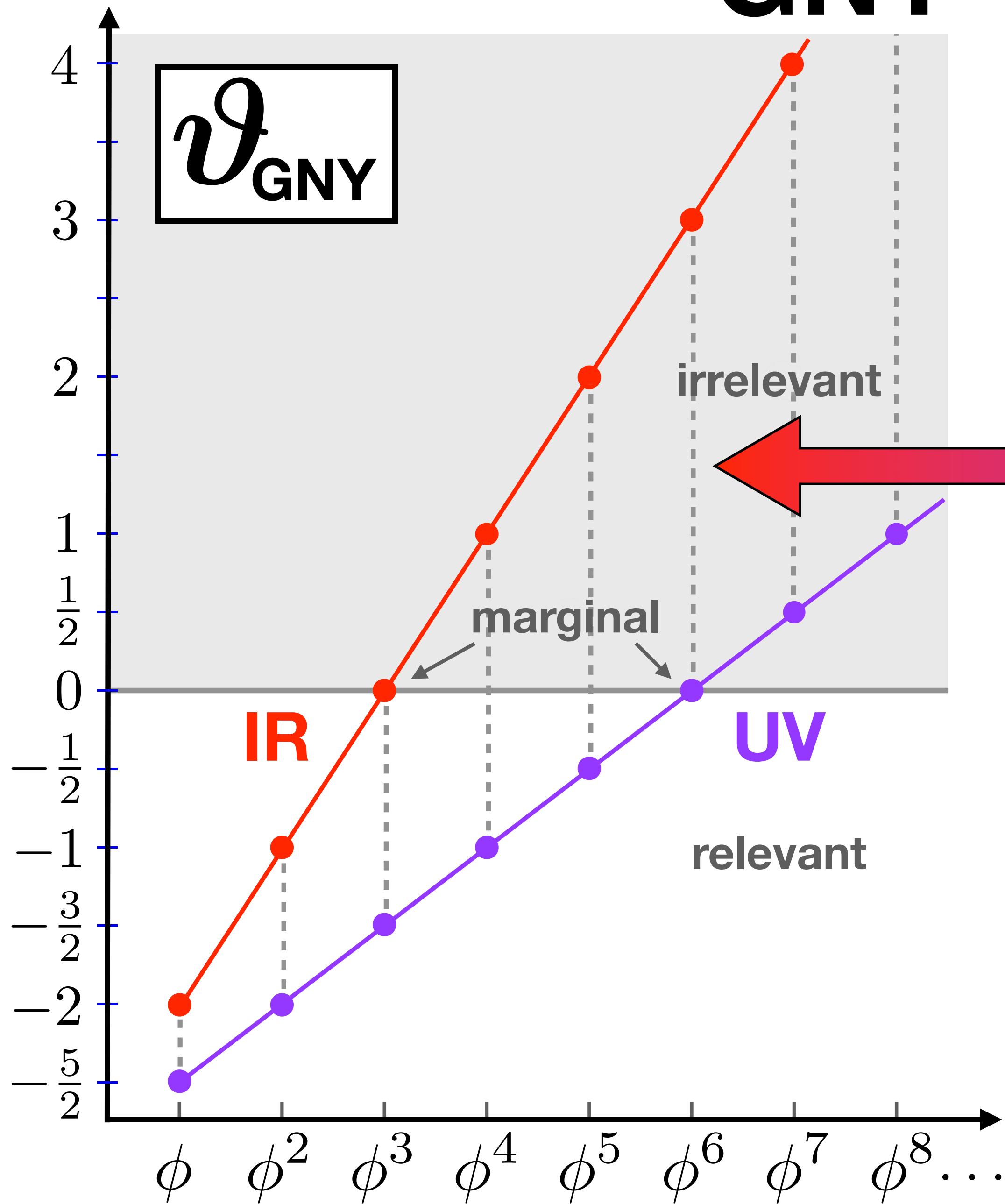


# Gross-Neveu – Yukawa

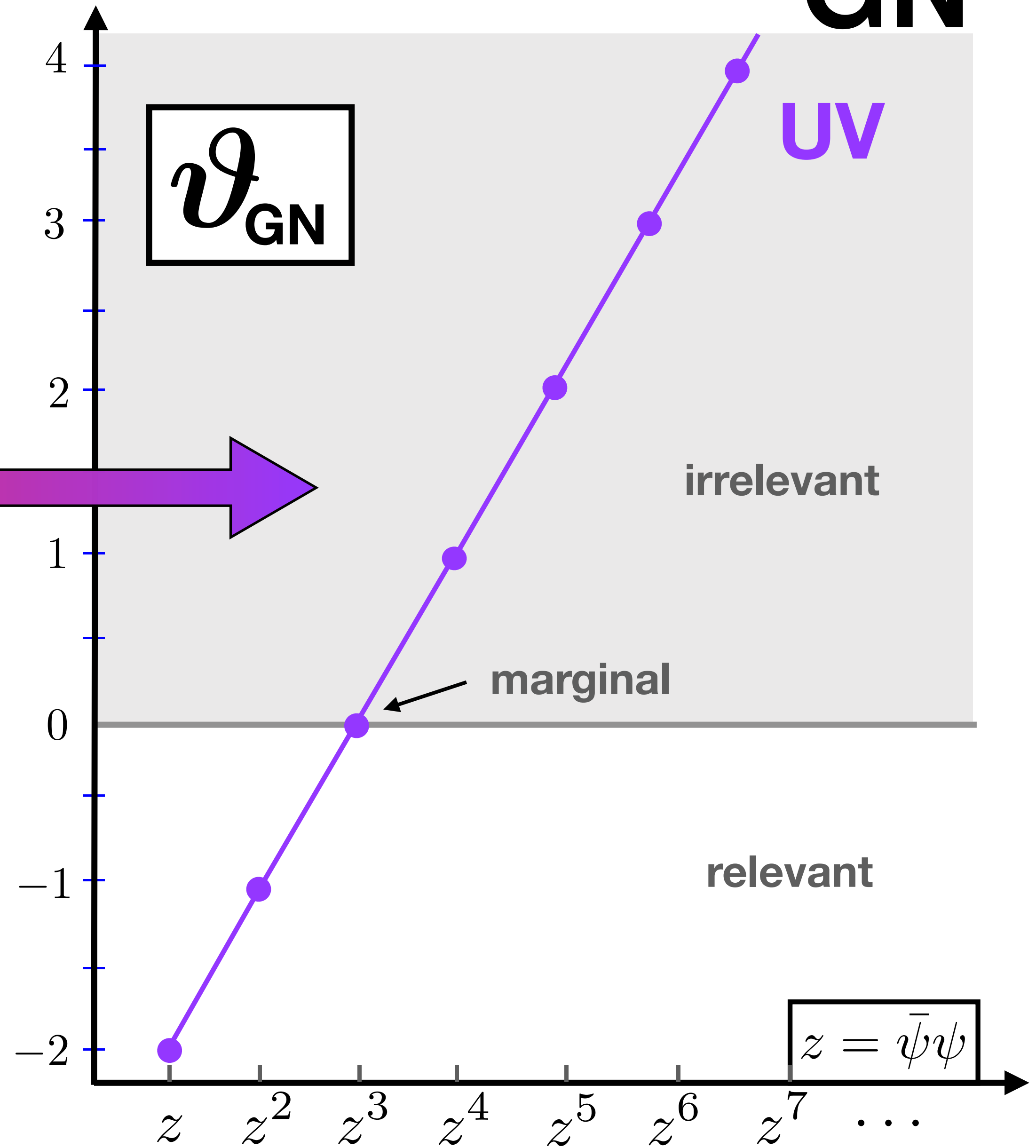




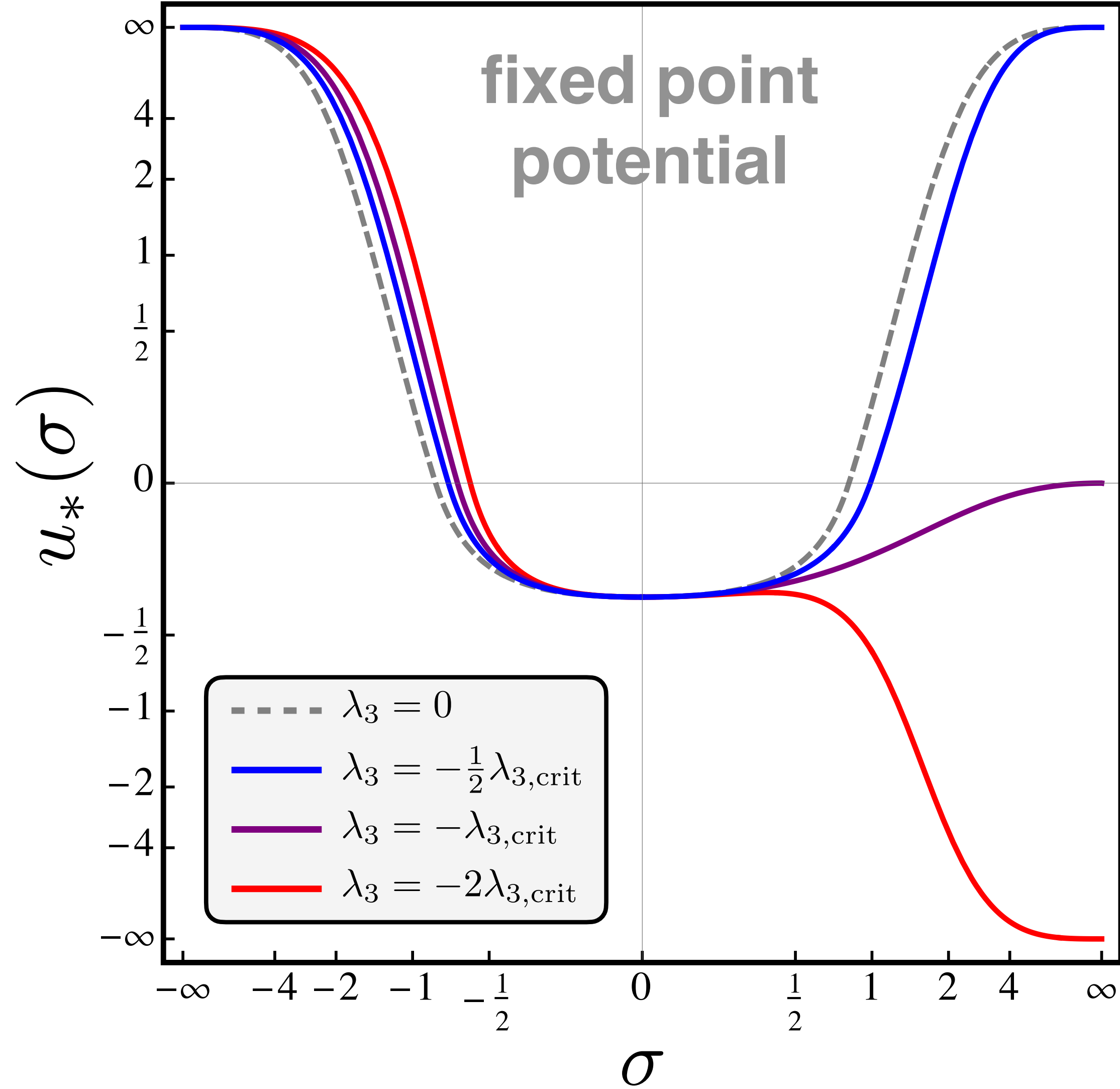
# GNY



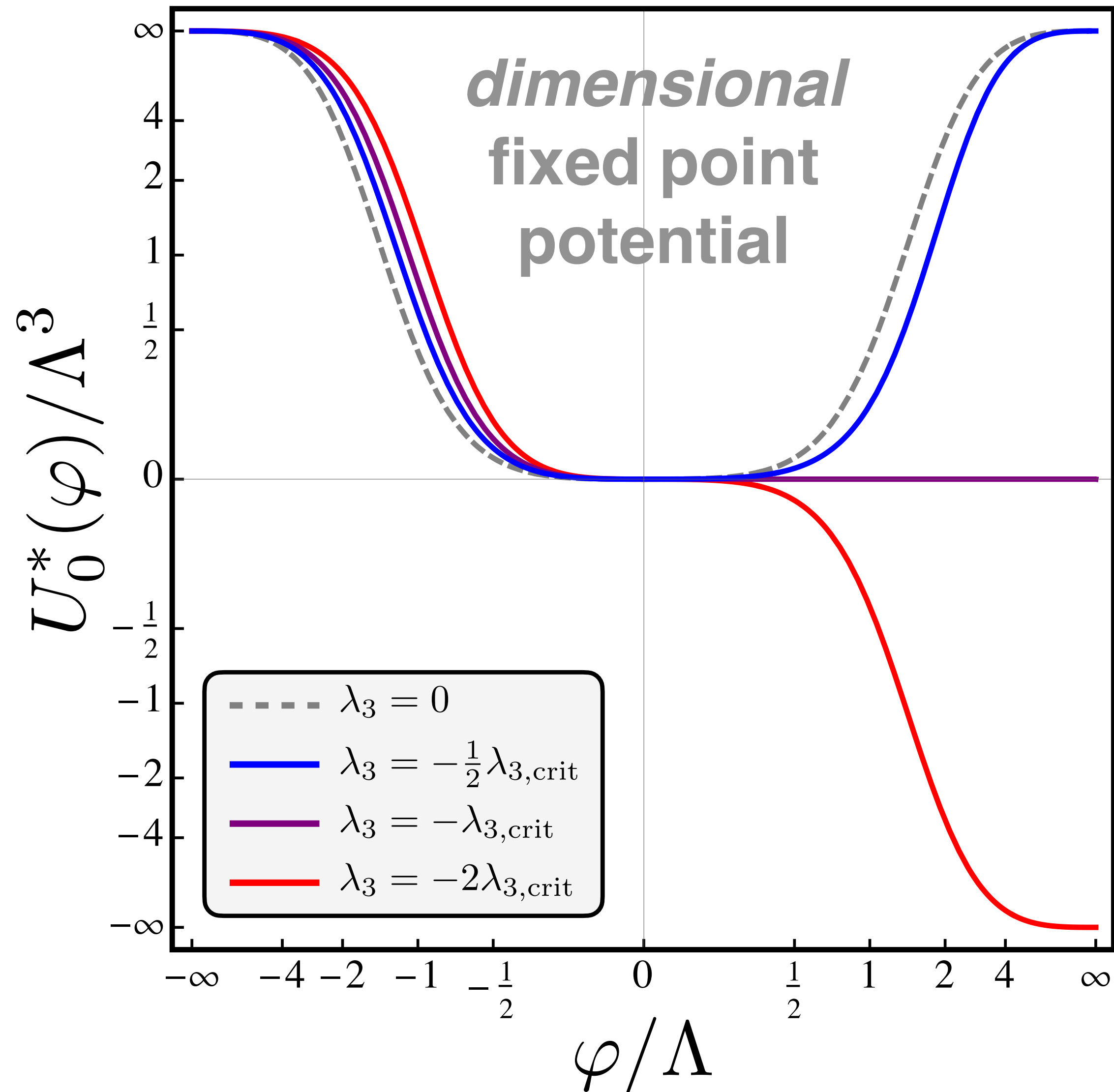
# GN



# Gross-Neveu — Yukawa



# Gross-Neveu—Yukawa



$$\propto [\lambda_3^* + \lambda_3^{\text{crit}} \text{sgn}(\phi)] \phi^3$$

conformal window  $|\lambda_3^*| \leq \lambda_3^{\text{crit}}$

scale symmetry broken **spontaneously**

$$|\lambda_3^*| = \lambda_3^{\text{crit}}$$



$\langle \phi \rangle = \text{free parameter}$



$M_F = H_* \langle \phi \rangle = \text{free parameter}$

$$M_S = 0$$

# bosonisation duality

large N fermionic theory

$$\int_x \bar{\psi}_i \not{\partial} \psi_i + F[\bar{\psi}_i \psi_i]$$

equivalent to

$$\int_x \left\{ \bar{\psi}_i \not{\partial} \psi_i + \sigma \bar{\psi}_i \psi_i \right\} + G[\phi]$$

functional Legendre transform

# bosonisation duality

$$\bar{\psi}_a \not{\partial} \psi_a + V_k(\bar{\psi}_a \psi_a)$$

GN

$$\bar{\psi}_a \not{\partial} \psi_a + \frac{1}{2} Z_\phi (\partial \phi)^2 + H_k \phi \bar{\psi}_a \psi_a + U_k(\phi)$$

GNY

map

$$V_k(\bar{\psi}_a \psi_a) = H_k \phi \bar{\psi}_a \psi_a + U_k(\phi)$$

$$U'_k(\phi) = -H_k \phi \bar{\psi}_a \psi_a$$

$$\lambda_{2\text{F}} = h \sigma_0 ,$$

$$\lambda_{4\text{F}} = -h^2 / \lambda_2 ,$$

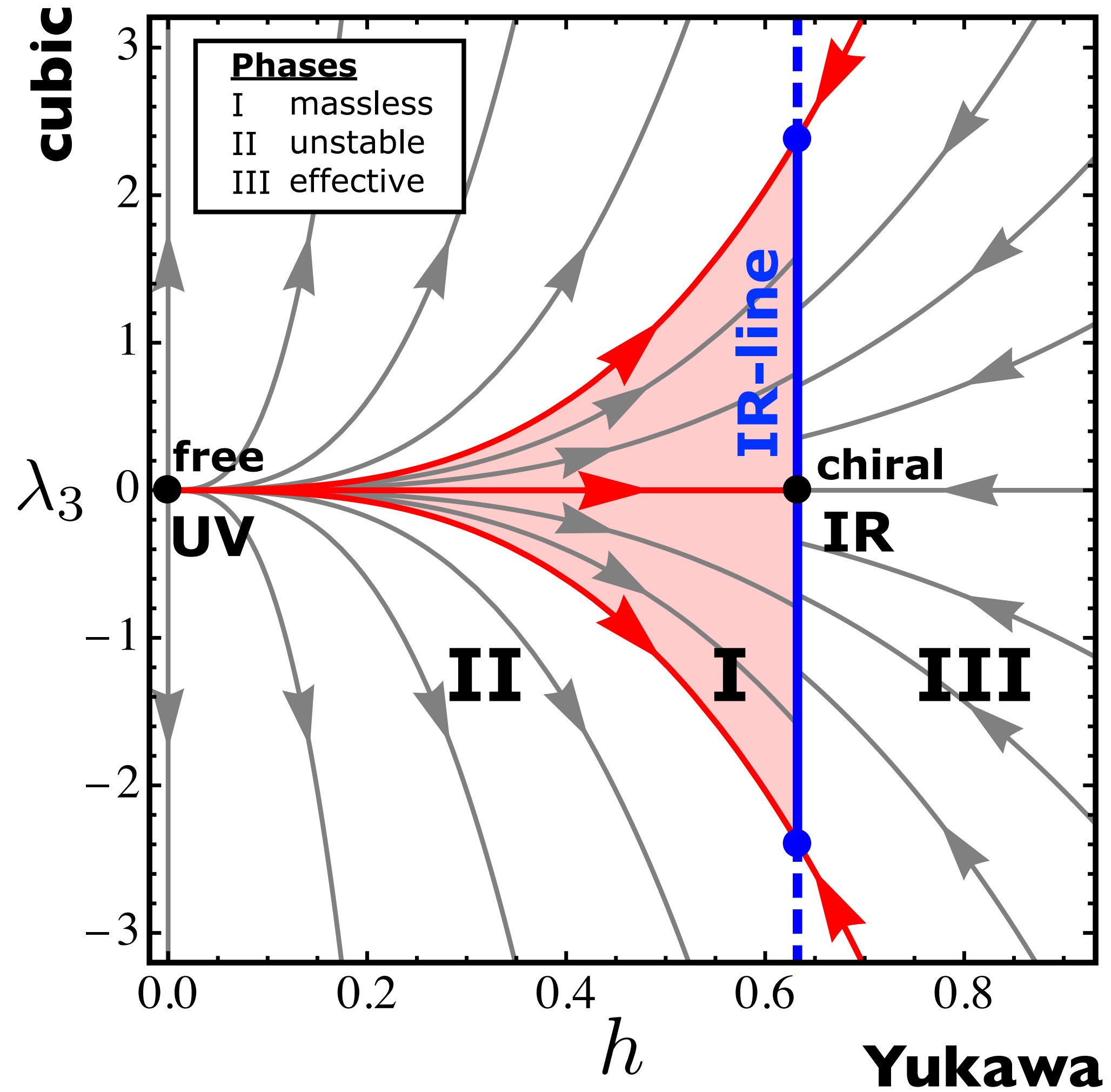
$$\lambda_{6\text{F}} = -h^3 \lambda_3 / \lambda_2^3 ,$$

$$\lambda_{8\text{F}} = h^4 (\lambda_2 \lambda_4 - 3 \lambda_3^2) / \lambda_2^5 ,$$

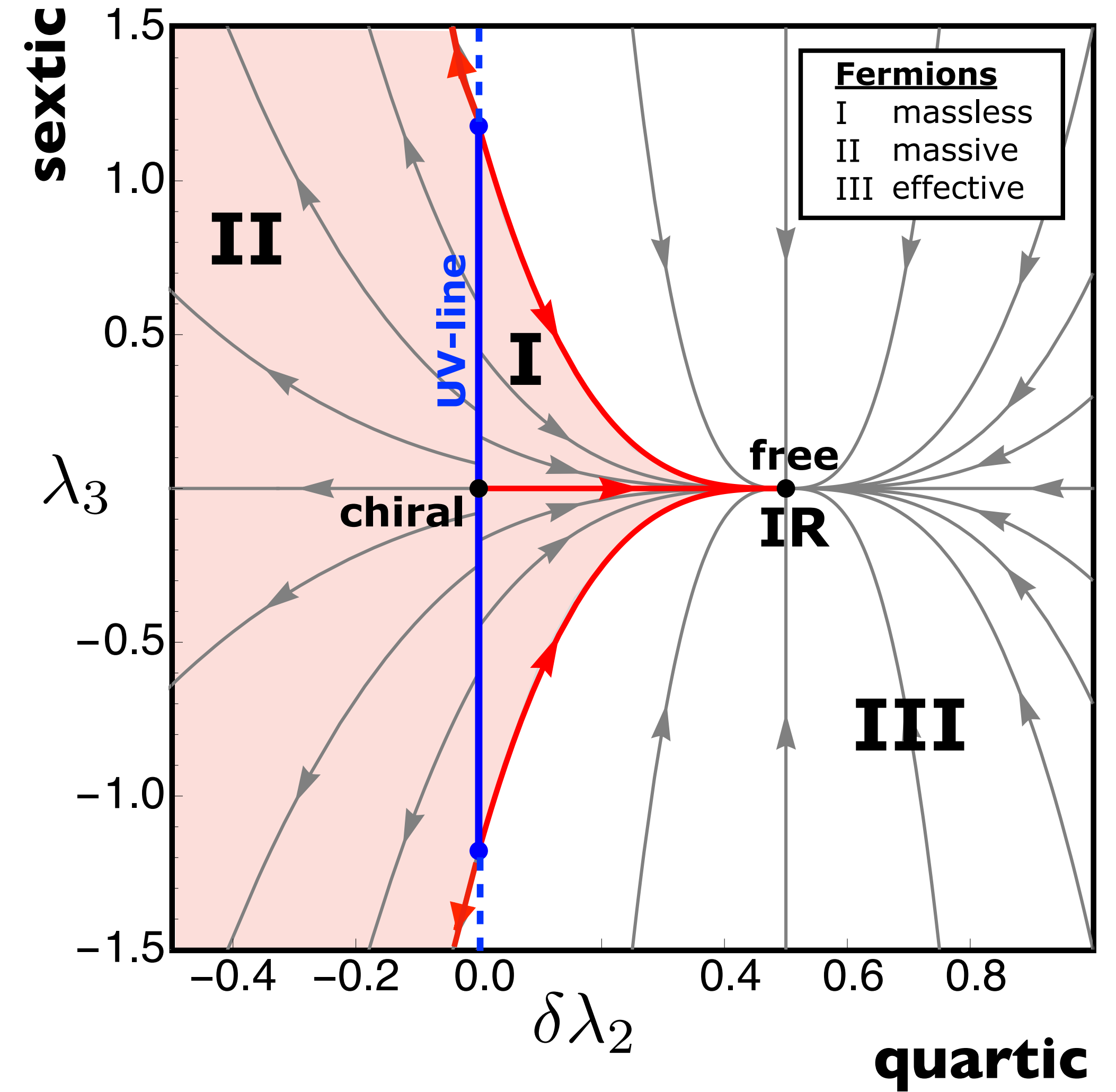
⋮

valid along RG trajectories

# GNY



# GN





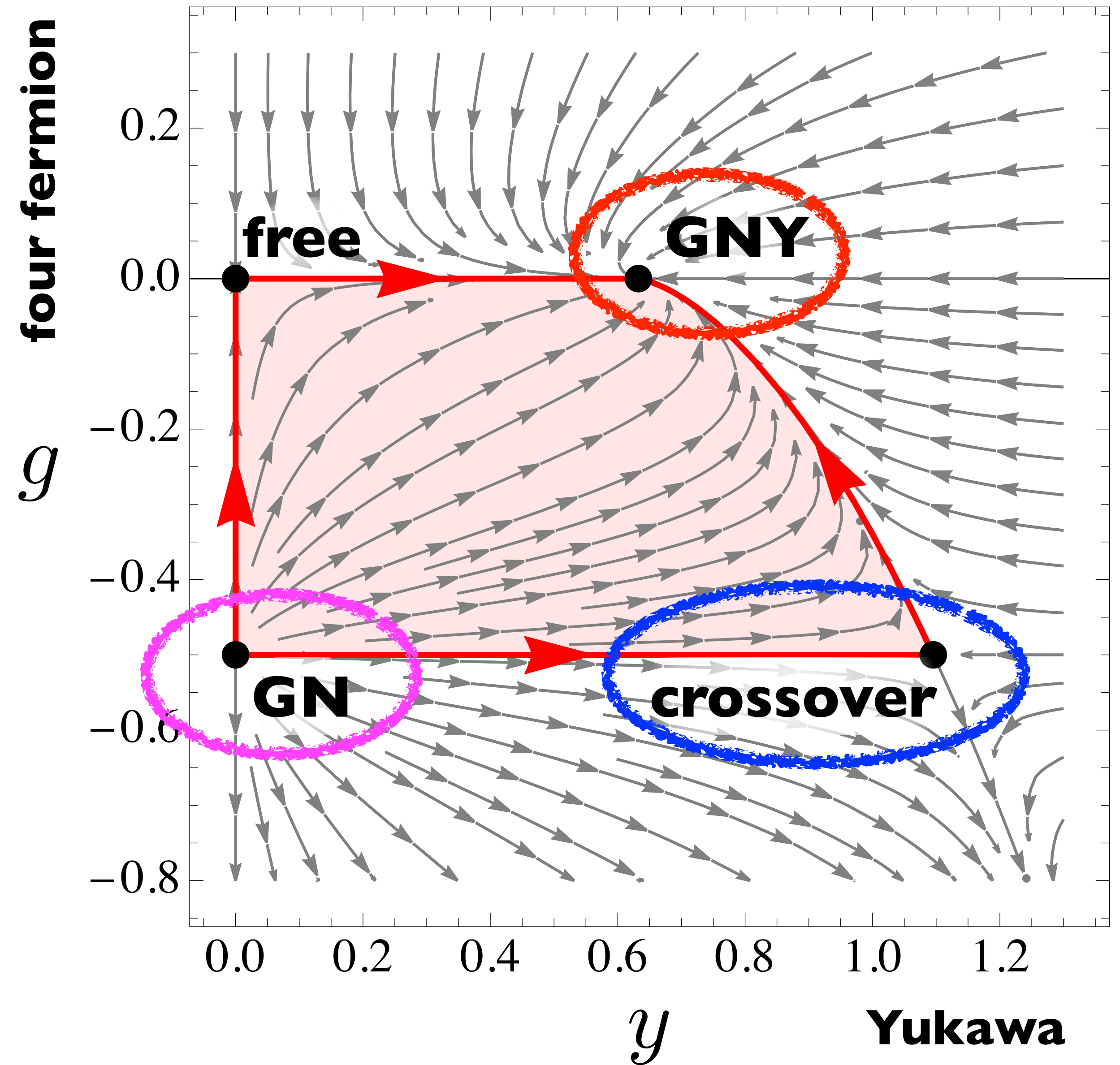
# outlook I

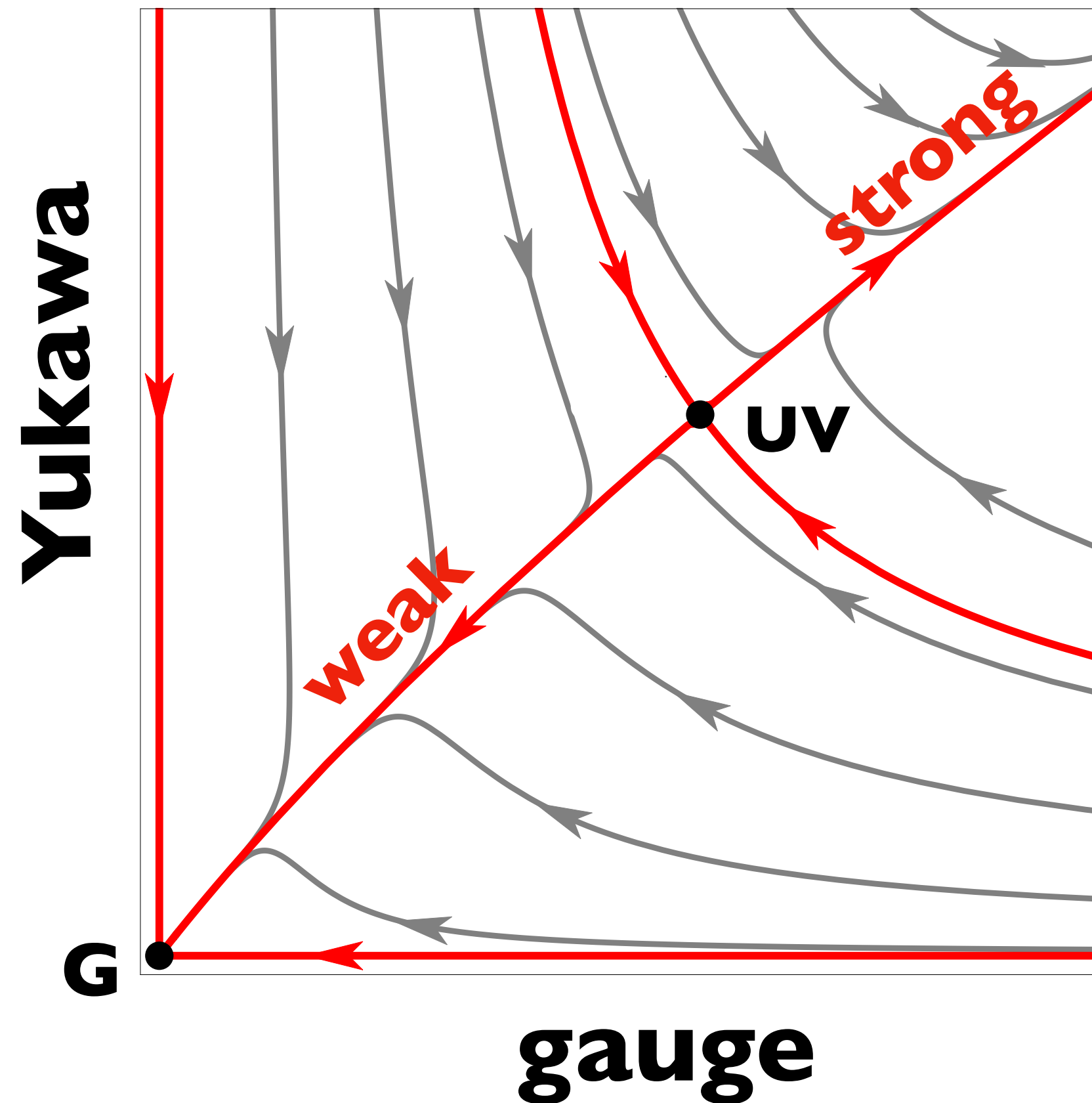
action

$$\bar{\psi}_i \not{\partial} \psi_i + \frac{1}{2} (\partial \phi)^2 + W_k(\phi, \bar{\psi}_i \psi_i)$$

e.g.

$$W_k = \frac{1}{2} G (\bar{\psi}_i \psi_i)^2 + Y_k \phi \bar{\psi}_i \psi_i$$





**SU(N) + Diracs**  
+ mesons

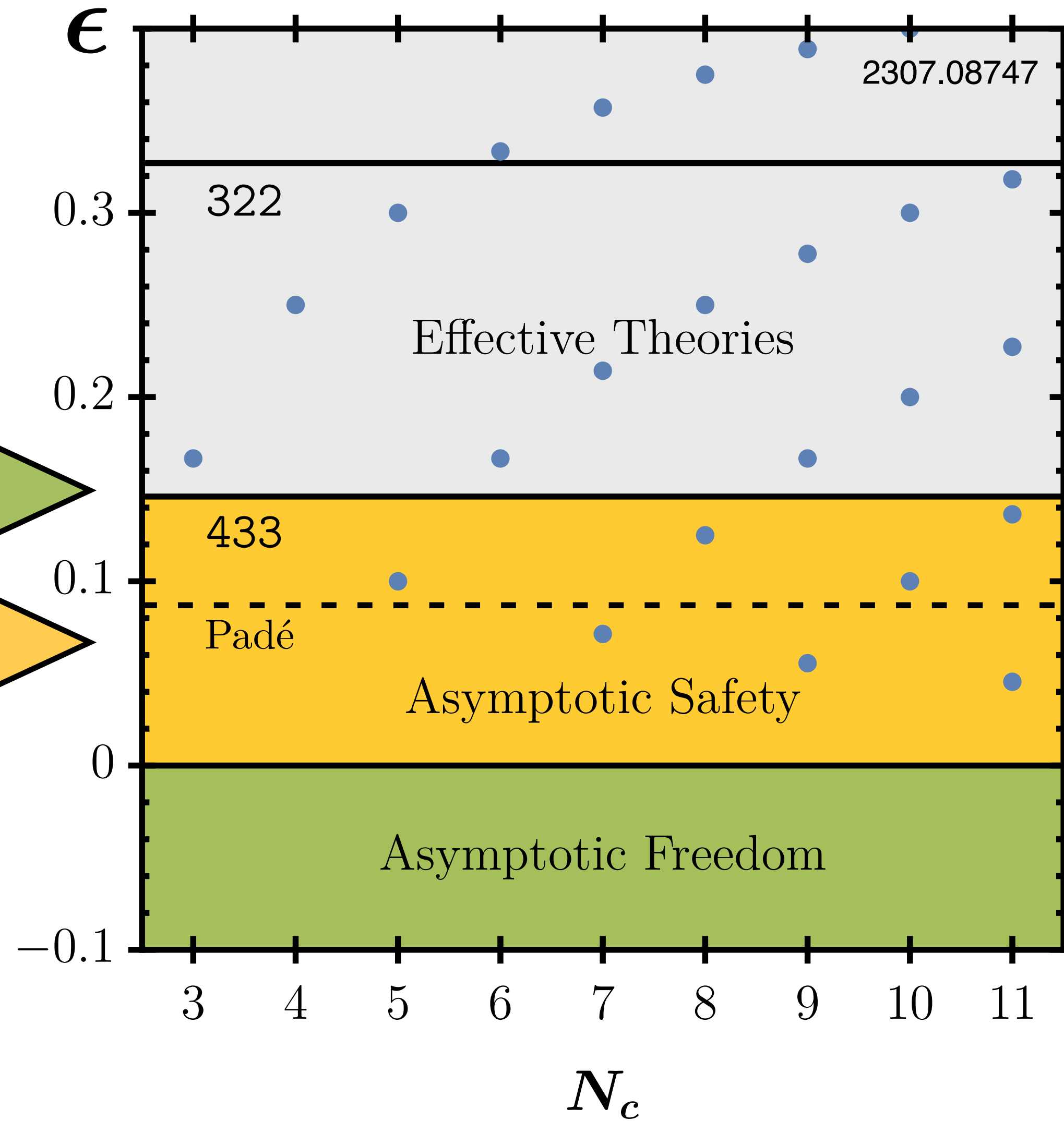
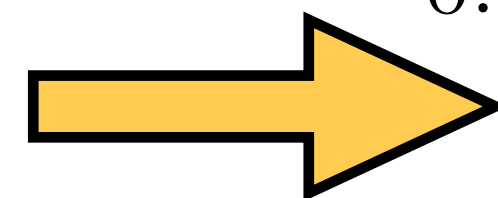
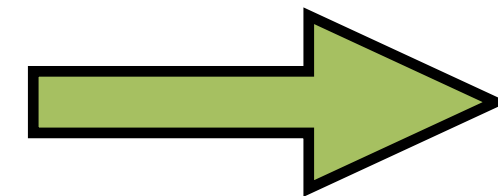
**SO(N) + Majoranas**  
+ mesons

**Sp(N) + Majoranas**  
+ mesons



loss of vacuum stability?  
merger?

conformal  
window



3 Loop

4 Loop

$$\epsilon = \frac{N_f}{N_c} - \frac{11}{2}$$

AD Bond, DF Litim, G Medina Vazquez, T Steudtner, **Conformal window for asymptotic safety**, 1710.07615 (PRD)

AD Bond, DF Litim, G Medina Vazquez, **Conformal windows beyond asymptotic freedom**, 2107.13020 (PRD)

DF Litim, N Riyaz, E Stamou, T Steudtner, **Asymptotic safety guaranteed at four loop**, 2307.08747 (PRD)

# Summary

## critical points in 3d QFTs

fermionic theories from first principles, UV / IR fixed points, mass generation & **spontaneous scale symmetry breaking**

**prerequisite:** interactions break chiral symmetry

links with CFTs and AdS/CFT

## large N dualities

maps between seemingly different 3d / 4d QFTs

new critical points, **fermions vs bosons**, **GN vs GNY**

## what's next?

4F interactions **beyond GN**, more dualities, spontaneous scale symmetry breaking in 4d, gauge fields, finite  $N$  ...

**Thank you!**