

On Exploring Leptogenesis in an Extension of the Scotogenic Model

PPC 2024 IIT Hyderabad

Parallel Session Talk

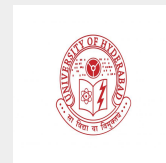
Based on: **JCAP 02, 041 (2024)**

In collaboration with

Devabrat Mahanta & Surender Verma

Labh Singh

Central University of Himachal Pradesh



• Introduction

- Why do all the structures in the Universe consist up of ordinary matter

Matter-Antimatter Asymmetry

$$\eta = \frac{n_B - n_{\bar{B}}}{n_\gamma} = 6.1 \times 10^{-10}$$

PDG, PTEP, 2022

- Observation of BBN ($T \sim 1 \text{ MeV}$) and CMB ($T \sim 1 \text{ eV}$) agree with each other
- What is the theoretical origin of matter-antimatter asymmetry?

• Introduction

- Why do all the structures in the Universe consist up of ordinary matter

Matter-Antimatter Asymmetry

$$\eta = \frac{n_B - n_{\bar{B}}}{n_\gamma} = 6.1 \times 10^{-10}$$

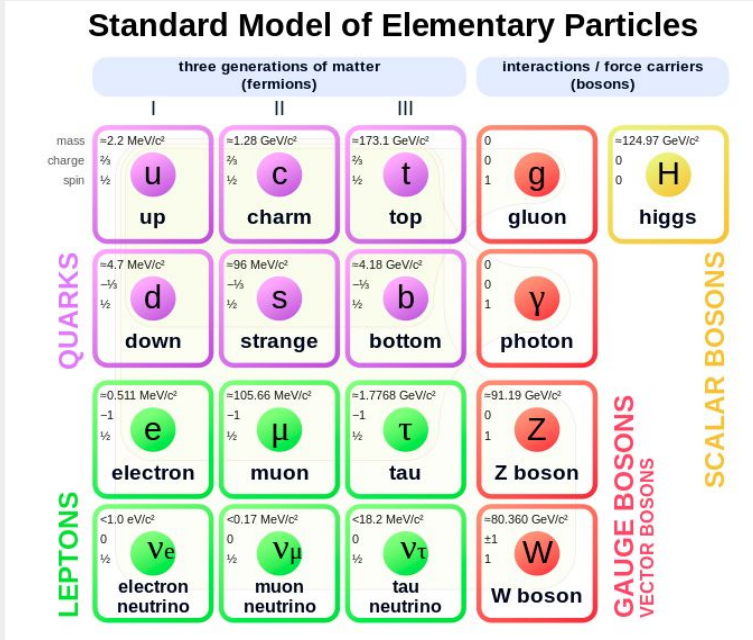
PDG, PTEP, 2022

- Observation of BBN ($T \sim 1 \text{ MeV}$) and CMB ($T \sim 1 \text{ eV}$) agree with each other
- What is the theoretical origin of matter-antimatter asymmetry?

Leptogenesis

Leptogenesis

Requirement of Leptogenesis



+

3 (or atleast 2) Heavy Majorana fermions (N_R)

Leptogenesis ← SM+3 RHN

→ Type-I Seesaw

- Type-I seesaw + Leptogenesis

$$\mathcal{L} = Y_\nu \nu_L \tilde{\Phi} \bar{N}_R + M_N \bar{N}_R^c N_R + h.c. \quad \text{Fukugita and Yanagida 1986}$$

• Type-I seesaw + Leptogenesis

Dirac-Yukawa, If Y
complex \rightarrow CP
Violation

$$\mathcal{L} = Y_\nu \nu_L \tilde{\Phi} \bar{N}_R + M_N \bar{N}_R^c N_R + h.c.$$

Fukugita and Yanagida 1986

- Type-I seesaw + Leptogenesis Lepton no. violating Majorana mass term $\Delta L = 2$

Dirac-Yukawa, If Y
complex \rightarrow CP
Violation

$$\mathcal{L} = Y_\nu \nu_L \tilde{\Phi} \bar{N}_R + M_N \bar{N}_R^c N_R + h.c.$$

Fukugita and Yanagida 1986

• Type-I seesaw + Leptogenesis Lepton no. violating Majorana mass term $\Delta L = 2$

Dirac-Yukawa, If Y
complex \rightarrow CP
Violation

$$\mathcal{L} = Y_\nu \nu_L \tilde{\Phi} \bar{N}_R + M_N \bar{N}_R^c N_R + h.c.$$

Fukugita and Yanagida 1986

$$\begin{pmatrix} 0 & m_D \\ m_D^T & M_N \end{pmatrix}$$

• Type-I seesaw + Leptogenesis Lepton no. violating Majorana mass term $\Delta L = 2$

Dirac-Yukawa, If Y complex \rightarrow CP Violation

$$\mathcal{L} = Y_\nu \nu_L \tilde{\Phi} \bar{N}_R + M_N \bar{N}_R^c N_R + h.c.$$

Fukugita and Yanagida 1986

$$\begin{pmatrix} 0 & m_D \\ m_D^T & M_N \end{pmatrix}$$

$$m_\nu = \frac{Y_\nu^2 v^2}{M_N} \approx 0.1 eV$$

• Type-I seesaw + Leptogenesis Lepton no. violating Majorana mass term $\Delta L = 2$

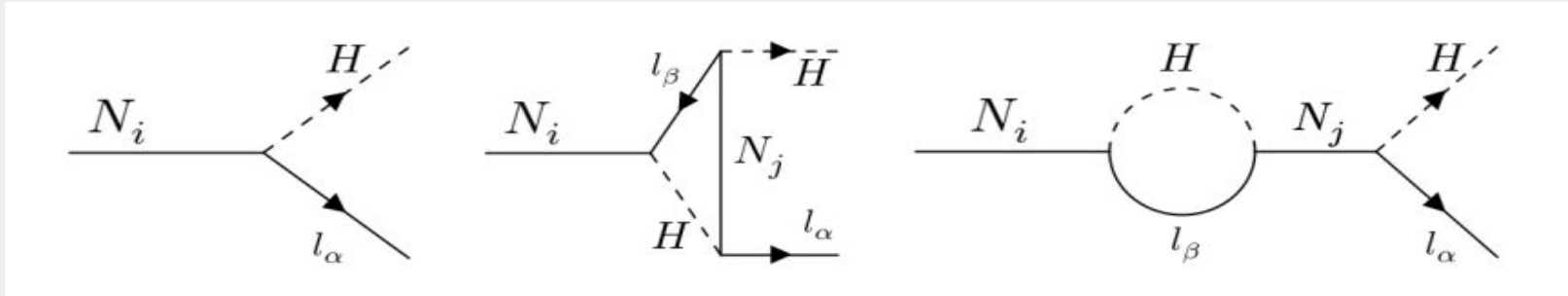
Dirac-Yukawa, If Y complex \rightarrow CP Violation

$$\mathcal{L} = Y_\nu \nu_L \tilde{\Phi} \bar{N}_R + M_N \bar{N}_R^c N_R + h.c.$$

Fukugita and Yanagida 1986

$$\begin{pmatrix} 0 & m_D \\ m_D^T & M_N \end{pmatrix}$$

$$m_\nu = \frac{Y_\nu^2 v^2}{M_N} \approx 0.1 eV$$



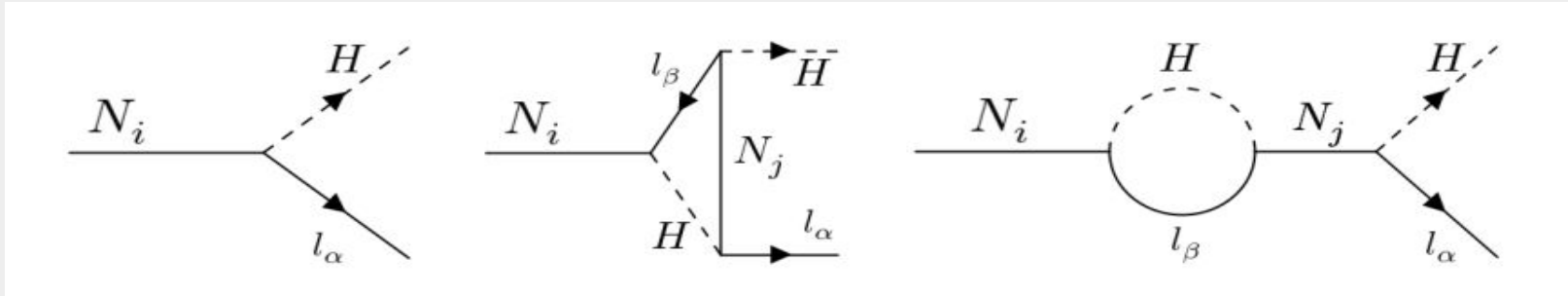
• Type-I seesaw + Leptogenesis Lepton no. violating Majorana mass term $\Delta L = 2$

Dirac-Yukawa, If Y complex \rightarrow CP Violation

$$\mathcal{L} = Y_\nu \nu_L \tilde{\Phi} \bar{N}_R + M_N \bar{N}_R^c N_R + h.c.$$

Fukugita and Yanagida 1986

Generates Neutrino mass and matter-antimatter asymmetry: Elegant



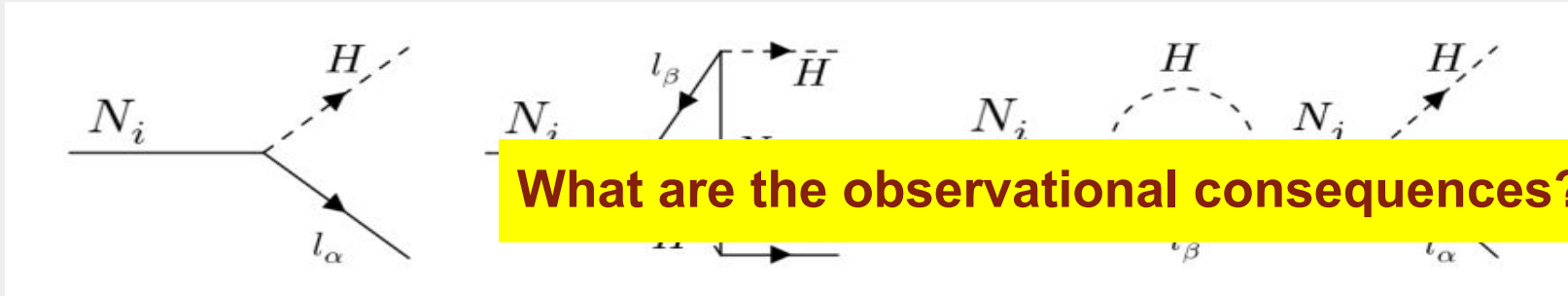
• Type-I seesaw + Leptogenesis Lepton no. violating Majorana mass term $\Delta L = 2$

Dirac-Yukawa, If Y complex \rightarrow CP Violation

$$\mathcal{L} = Y_\nu \nu_L \tilde{\Phi} \bar{N}_R + M_N \bar{N}_R^c N_R + h.c.$$

Fukugita and Yanagida 1986

Generates Neutrino mass and matter-antimatter asymmetry: Elegant



What are the observational consequences?

• Type-1 seesaw + Leptogenesis

$$\mathcal{L} = Y_\nu \nu_L \tilde{\Phi} \bar{N}_R + M_N \bar{N}_R^c N_R + h.c.$$

$$m_\nu = \frac{Y_\nu^2 v^2}{M_N} \approx 0.1 eV \quad \text{For } Y_\nu \approx \mathcal{O}(1), M_N \text{ have to be } 10^{13} - 10^{14} \text{ GeV.}$$

• Type-1 seesaw + Leptogenesis

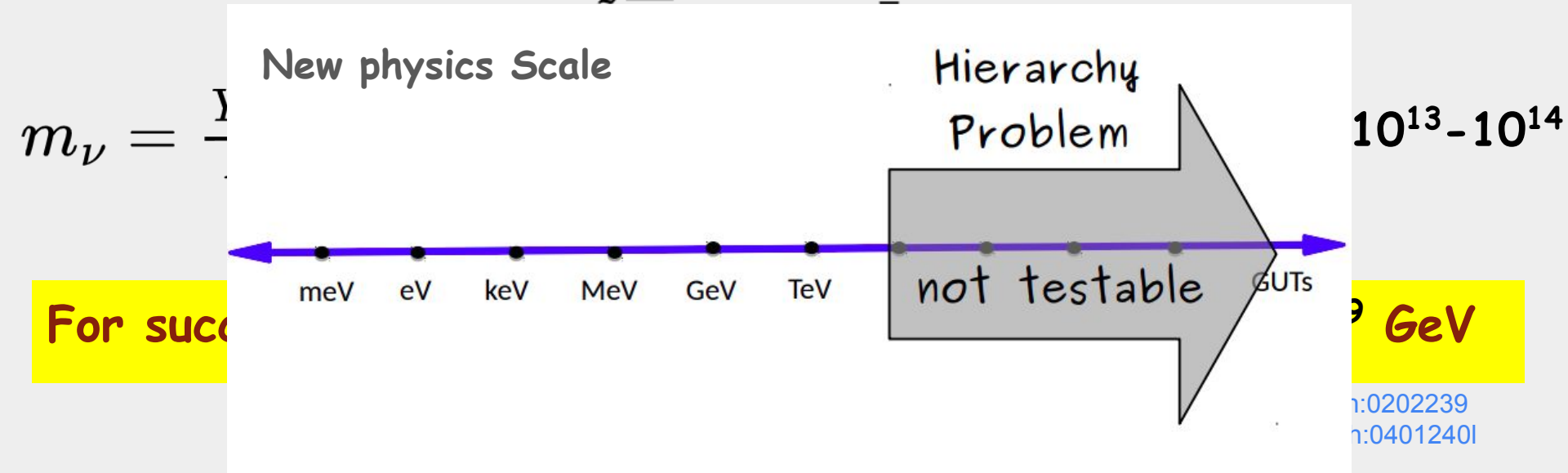
$$\mathcal{L} = Y_\nu \nu_L \tilde{\Phi} \bar{N}_R + M_N \bar{N}_R^c N_R + h.c.$$

$$m_\nu = \frac{Y_\nu^2 v^2}{M_N} \approx 0.1 eV \quad \text{For } Y_\nu \approx \mathcal{O}(1), M_N \text{ have to be } 10^{13} - 10^{14} \text{ GeV.}$$

For successful leptogenesis and correct ν mass, $M_N \geq 10^9 \text{ GeV}$

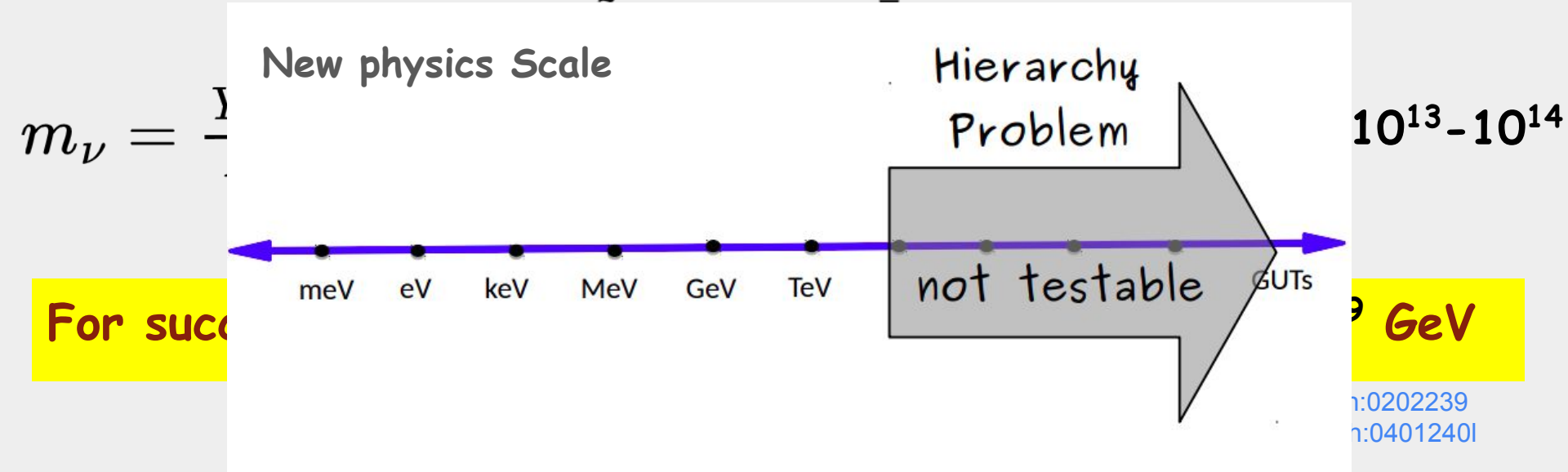
Davidson, Ibarra, hep-ph:0202239
Buchmuller, et. al, hep-ph:0401240

Type-1 seesaw + Leptogenesis



Inaccessible to conventional laboratory experiments!!

Type-1 seesaw + Leptogenesis



Inaccessible to conventional laboratory experiments!!

Therefore, we are behind low scale leptogenesis scenarios

Singlet-Triplet Scotogenic Model

Particle Content	Generations	Symmetry
		$(SU(2)_L \times U(1)_Y \times Z_2)$
L	3	$(2, -\frac{1}{2}, +)$
ϕ	1	$(2, \frac{1}{2}, +)$
N	1	$(1, 0, -)$
Σ	1	$(3, 0, -)$
Ω	1	$(3, 0, +)$
η	1	$(2, \frac{1}{2}, -)$

Table: Particle content of the singlet-triplet scotogenic model

Singlet-Triplet Scotogenic Model

Beyond SM
Symmetry

Particle Content	Generations	Symmetry
		$(SU(2)_L \times U(1)_Y \times Z_2)$
SM particles	L	3
	ϕ	1
Beyond SM particle	N	1
	Σ	1
	Ω	1
	η	1

Table: Particle content of the singlet-triplet scotogenic model

Singlet-Triplet Scoto-genic Model

Beyond SM Symmetry

Contribute to ν mass generation and Leptogenesis

Particle Content	Generations	Symmetry
		$(SU(2)_L \times U(1)_Y \times Z_2)$
L	3	$(2, -\frac{1}{2}, +)$
ϕ	1	$(2, \frac{1}{2}, +)$
N	1	$(1, 0, -)$
Σ	1	$(3, 0, -)$
Ω	1	$(3, 0, +)$
η	1	$(2, \frac{1}{2}, -)$

SM particles

Beyond SM particle

Table: Particle content of the singlet-triplet scotogenic model

Singlet-Triplet Scoto-genic Model

Beyond SM Symmetry

Particle Content	Generations	Symmetry
		$(SU(2)_L \times U(1)_Y \times Z_2)$
<i>L</i>	3	$(2, -\frac{1}{2}, +)$
ϕ	1	$(2, \frac{1}{2}, +)$
<i>N</i>	1	$(1, 0, -)$
Σ	1	$(3, 0, -)$
Ω	1	$(3, 0, +)$
η	1	$(2, \frac{1}{2}, -)$

Contribute to ν mass generation and Leptogenesis

Contribute to the W-mass anomaly

Table: Particle content of the singlet-triplet scotogenic model

Singlet-Triplet Scoto-genic Model

Beyond SM Symmetry

Particle Content	Generations	Symmetry
		$(SU(2)_L \times U(1)_Y \times Z_2)$
<i>L</i>	3	$(2, -\frac{1}{2}, +)$
ϕ	1	$(2, \frac{1}{2}, +)$
<i>N</i>	1	$(1, 0, -)$
Σ	1	$(3, 0, -)$
Ω	1	$(3, 0, +)$
η	1	$(2, \frac{1}{2}, -)$

SM particles

Beyond SM particle

Contribute to ν mass generation and Leptogenesis

Contribute to the W-mass anomaly

Dark-matter candidate

Table: Particle content of the singlet-triplet scotogenic model

The relevant Yukawa Lagrangian of the model is given by

$$\begin{aligned} \mathcal{L} = & Y_L^{\alpha\beta} \bar{L}_\alpha \phi \ell_\beta + Y_N^\alpha \bar{L}_\alpha i\sigma_2 \eta N + Y_\Sigma^\alpha \bar{L}_\alpha C \Sigma^\dagger i\sigma_2 \eta + Y_\Omega \text{Tr}(\bar{\Sigma} \Omega) N \\ & + \frac{1}{2} M_N \bar{N}^c N + \frac{1}{2} M_\Sigma \text{Tr}(\bar{\Sigma}^c \Sigma) + h.c. \end{aligned}$$

The scalar potential of the model is given by

$$\begin{aligned} V = & -m_\phi^2 \phi^\dagger \phi + m_\eta^2 \eta^\dagger \eta + \frac{\lambda_1}{2} (\phi^\dagger \phi)^2 + \frac{\lambda_2}{2} (\eta^\dagger \eta)^2 + \frac{m_\Omega^2}{2} \text{Tr}(\Omega^\dagger \Omega) + \lambda_3 (\phi^\dagger \phi) (\eta^\dagger \eta) \\ & + \lambda_4 (\phi^\dagger \eta) (\eta^\dagger \phi) \frac{\lambda_5}{2} [(\phi^\dagger \eta)^2 + h.c.] + \frac{\lambda^\eta}{2} (\eta^\dagger \eta) \text{Tr}(\Omega^\dagger \Omega) \\ & + \frac{\lambda_1^\Omega}{2} (\phi^\dagger \phi) \text{Tr}(\Omega^\dagger \Omega) + \frac{\lambda_2^\Omega}{4} \text{Tr}(\Omega^\dagger \Omega)^2 + \mu_1 \phi^\dagger \Omega \phi + \mu_2 \eta^\dagger \Omega \eta. \end{aligned}$$

After EWSB, the masses of scalar sector

$$m_{\eta^\pm}^2 = m_\eta^2 + \frac{1}{2}\lambda_3 v_\phi^2 + \frac{1}{2}\lambda^\eta v_\Omega^2 + \frac{1}{\sqrt{2}}v_\Omega \mu_2,$$

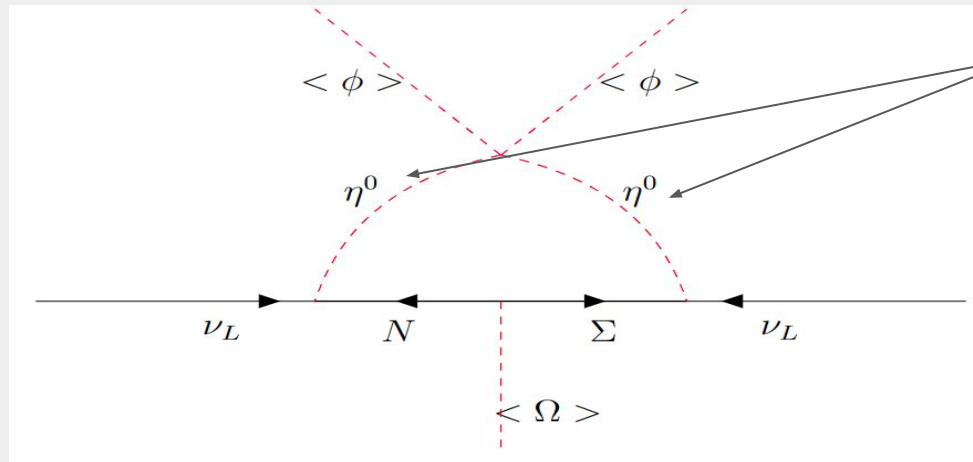
$$m_{\eta^R}^2 = m_\eta^2 + \frac{1}{2}\lambda_3 v_\phi^2 + \frac{1}{2}(\lambda_4 + \lambda_5)v_\phi^2 - \frac{1}{\sqrt{2}}v_\Omega \mu_2 + \frac{1}{2}\lambda^\eta v_\Omega^2,$$

$$m_{\eta^I}^2 = m_\eta^2 + \frac{1}{2}\lambda_3 v_\phi^2 + \frac{1}{2}(\lambda_4 - \lambda_5)v_\phi^2 - \frac{1}{\sqrt{2}}v_\Omega \mu_2 + \frac{1}{2}\lambda^\eta v_\Omega^2.$$

And masses of fermion sector are

$$m_{\chi^\pm} = M_\Sigma,$$
$$m_{\chi_{1,2}^0} = \frac{1}{2} \left(M_\Sigma + M_N \mp \sqrt{(M_\Sigma - M_N)^2 + 4(v_\Omega Y_\Omega)^2} \right),$$

The neutrino mass generated at one-loop level



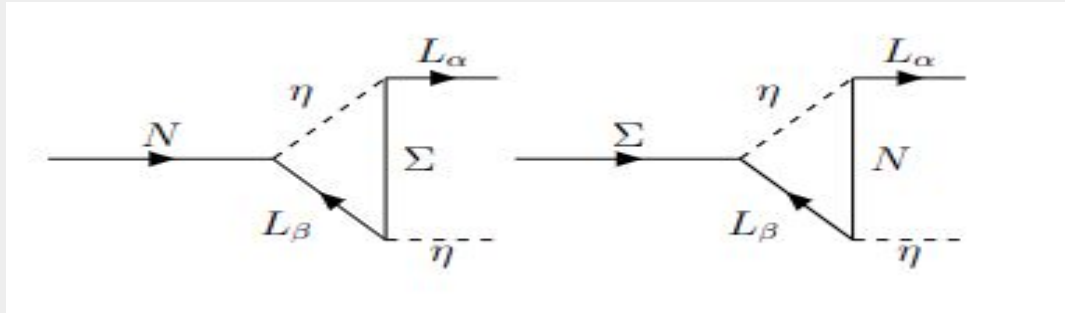
Dark-Matter contribute to neutrino mass generation

M. Hirsch, et. al, JHEP, 2013

$$(\mathcal{M}_\nu)_{\alpha\beta} = \sum_{i=1}^2 \frac{h_{\alpha i} h_{\beta i} m_{\chi_i^0}}{2(4\pi)^2} \left[\frac{m_{\eta R}^2 \ln\left(\frac{m_i^2}{m_{\eta 0}^2}\right)}{m_{\chi_i^0}^2 - m_{\eta R}^2} - \frac{m_{\eta I}^2 \ln\left(\frac{m_{\chi_i^0}^2}{m_{\eta I}^2}\right)}{m_{\chi_i^0}^2 - m_{\eta I}^2} \right]$$

For low scale leptogenesis in our model, we considered both the heavy fermions to be moderately hierarchical.

$$M_N \approx M_\Sigma$$



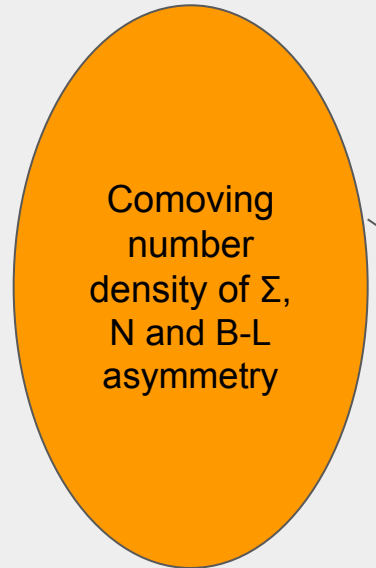
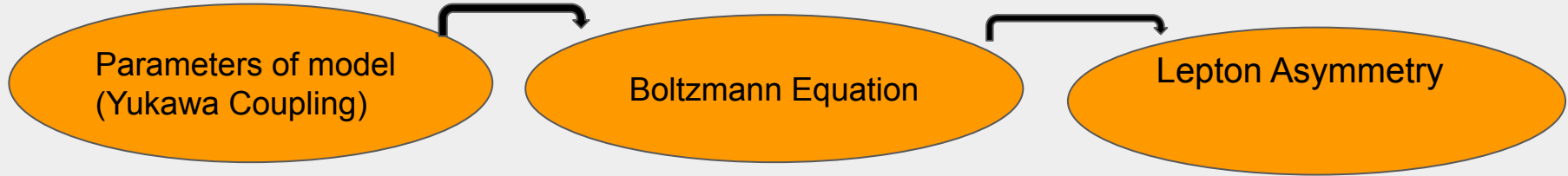
LS, D. Mahanta, S. Verma, JCAP, 2024

CP-Asymmetry parameter:

$$\epsilon^{F(\Sigma)} = \frac{1}{8\pi (Y_{F(\Sigma)}^\dagger Y_{F(\Sigma)})} \sum_{j \neq i} \text{Im} \left[\left(Y_{F(\Sigma)}^\dagger Y_{\Sigma(F)} \right)_{ij}^2 \right] \mathcal{F} \left(\frac{M_{\Sigma(F)}^2}{M_{F(\Sigma)}^2} \right)$$

where
$$\mathcal{F}(x) = \sqrt{x} \left[1 + \frac{1}{1-x} + (1+x) \ln \left(\frac{x}{1+x} \right) \right]$$

$$\eta = \frac{n_{B^-} - n_{\bar{B}}}{n_\gamma} = 6.1 \times 10^{-10}$$



$$\begin{aligned} \frac{dn_\Sigma}{dz} &= -D_\Sigma(n_\Sigma - n_\Sigma^{eq}) - S_A(n_\Sigma^2 - (n_\Sigma^{eq})^2), \\ \frac{dn_N}{dz} &= -D_N(n_N - n_N^{eq}), \\ \frac{dn_{B-L}}{dz} &= -\epsilon^\Sigma D_\Sigma(n_\Sigma - n_\Sigma^{eq}) - \epsilon^N D_N(n_N - n_N^{eq}) - W_{ID}^\Sigma n_{B-L} - W_{ID}^N n_{B-L} \\ &\quad - W_{\Delta L} n_{B-L}. \end{aligned}$$

LS, D. Mahanta, S. Verma, JCAP, 2024

Due to CP- violations

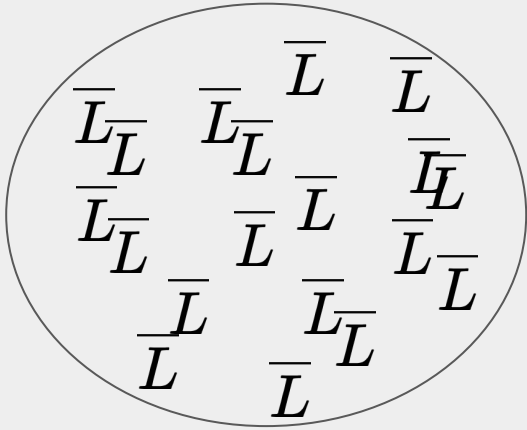


More **Anti-Leptons** than **Leptons**

Due to CP- violations



More **Anti-Leptons** than **Leptons**



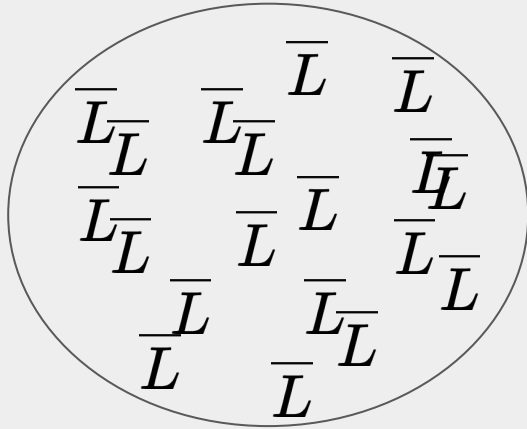
Asymmetry in the Lepton sector

Due to CP-violations



More **Anti-Leptons** than **Leptons**

$L=-1, B=0$
 $B-L=1, B+L=-1$

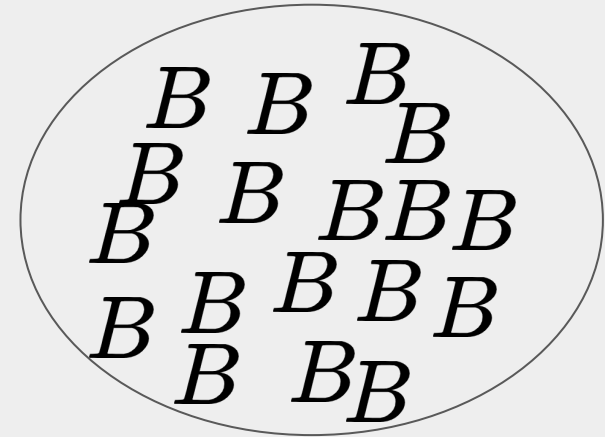


Asymmetry in the Lepton sector

**Sphaleron
Process**

(Converts
Lepton
Asymmetry into
Baryon
Asymmetry
before EWPT)

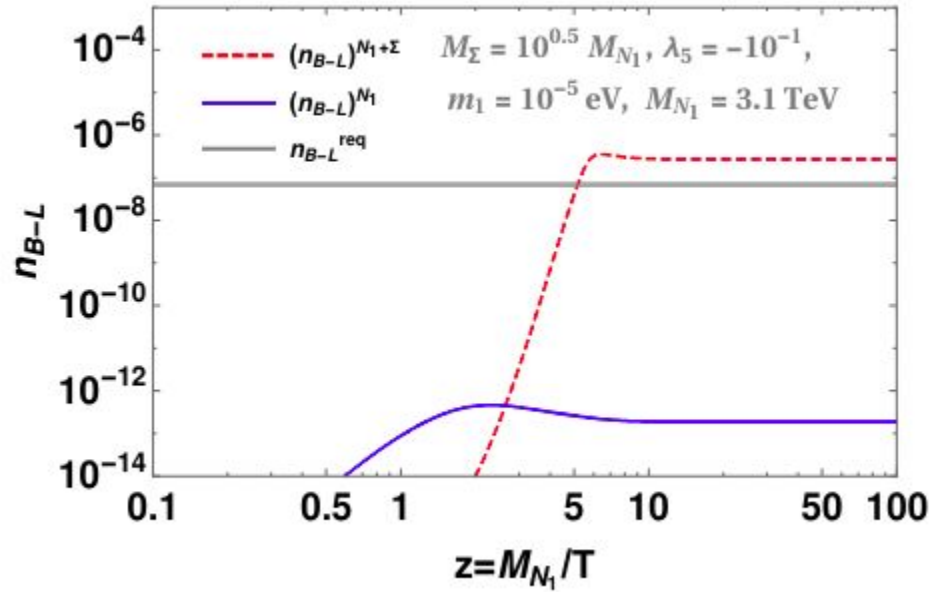
$L=0, B=1$
 $B-L=1, B+L=1$



Asymmetry in the Baryon sector

$$\eta_{B-L} = 9.2 \times 10^{-3} \eta$$

Results



Results

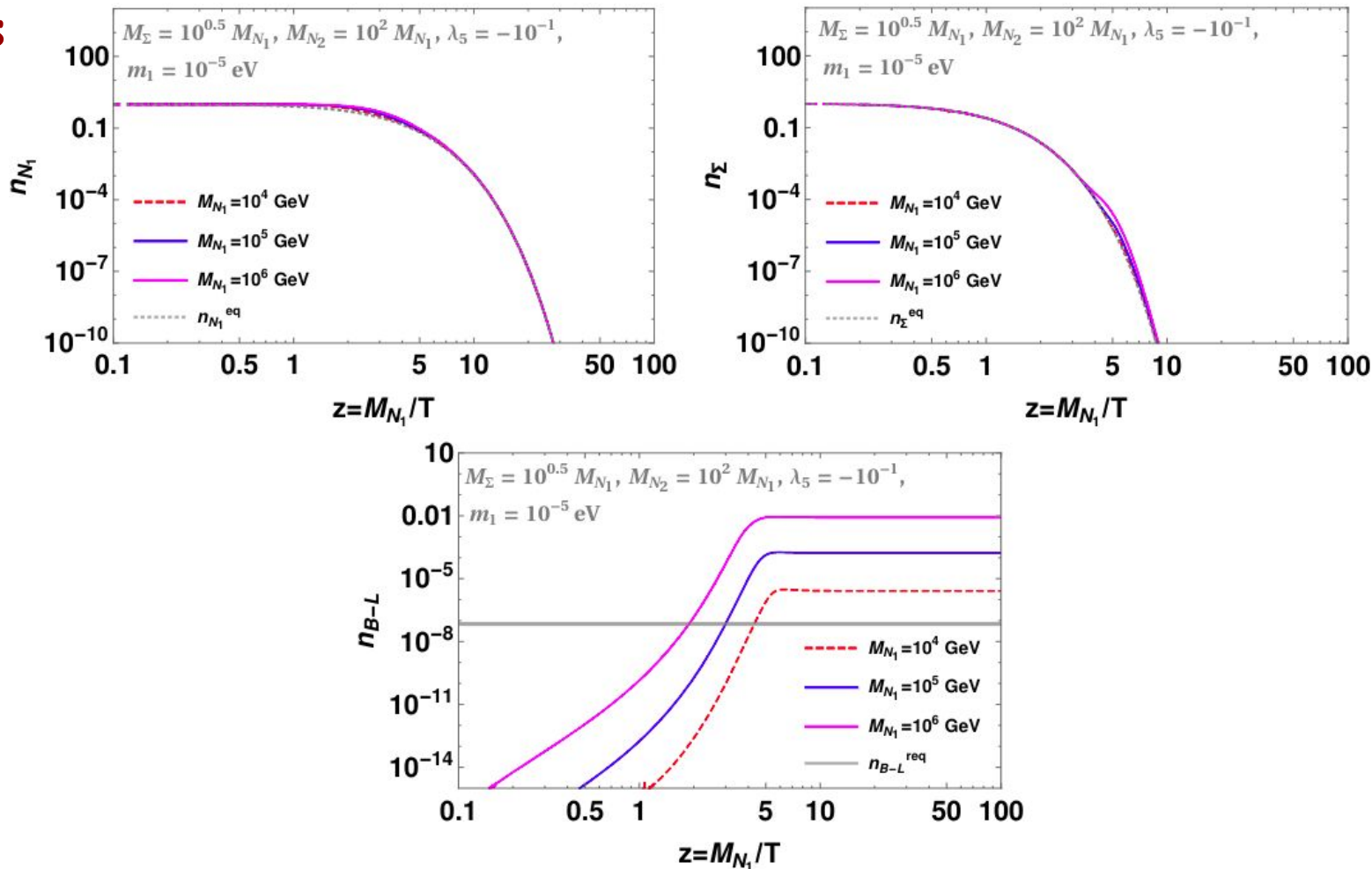
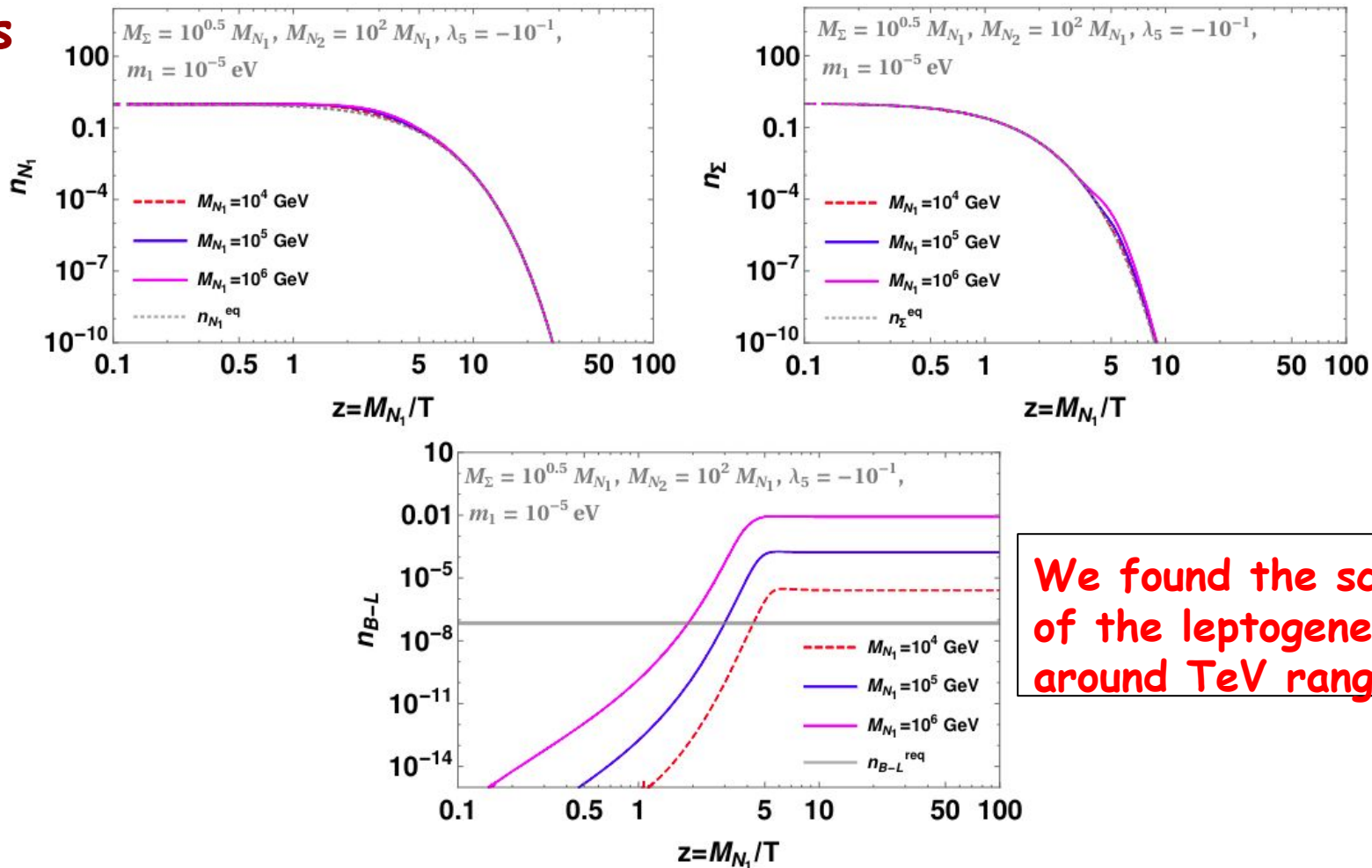


Figure 2: Evolution of comoving number density w.r.t Z .

Results



We found the scale of the leptogenesis around TeV range

Figure 2: Evolution of comoving number density w.r.t Z .

Summary

- In this work, we have realised low-scale leptogenesis in singlet-triplet scotogenic model.
- We obtained the scale of the leptogenesis in the TeV range.
- Such a low scale can be probed in the future collider experiments.

Summary

- In this work, we have realised low-scale leptogenesis in singlet-triplet scotogenic model.
- We obtained the scale of the leptogenesis in the TeV range.
- Such a low scale can be probed in the future collider experiments.

Thank You