Exploring cosmological gravitational wave backgrounds through the synergy of LISA and ET

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LISA and ET



Schematic configurations of LISA¹ (left) and ET² (right).

¹Image from: J. Baker *et al.*, arXiv: 1907.06482 [astro-ph.IM] (2019).

²Image from: T. Regimbau *et al.*, Phys. Rev. D **86**, 122001 (2012).

Nominal sensitivity curves³





³R. Flauger *et al.*, JCAP **01**, 059 (2021); S. Hild *et al.*, Class. Quant. Grav. **28**, 094013 (2011).

SGWB with LISA and ET

Combined SNR

Our aim is to estimate the optimal signal-to-noise ratio (SNR) for detecting a stochastic GW signal with the two instruments LISA and ET working together.

The square of the total SNR is the sum of the squares of the individual SNRs⁴:

$$SNR_{tot} = \sqrt{T \int_0^\infty df \left[\frac{\Omega_{GW}^2(f)}{\Sigma_{LISA}^2(f)} + \frac{\Omega_{GW}^2(f)}{\Sigma_{ET}^2(f)}\right]} \\ = \sqrt{SNR_{LISA}^2 + SNR_{ET}^2}.$$

The two detectors together can reach higher values of SNR.

We are interested in SGW spectra enhanced within the broad frequency band:

 $10^{-5} \,\mathrm{Hz} \le f \le 445 \,\mathrm{Hz}.$

⁴T. L. Smith and R. R. Caldwell, Phys. Rev. D 100, 104055 (2019); J. D. Romano and N. J. Cornish, Living Rev. Rel. 20, 2 (2017).

Power-law integrated sensitivity (PLS) curves

GW spectra of a power-law form are assumed:

$$\Omega_{_{\mathrm{GW}}}(f) = \Omega_{\beta} \left(\frac{f}{f_*}\right)^{\beta}.$$

For a set of β and some choice of f_* , the following amplitude is evaluated:

$$\Omega_{\beta} = \frac{\mathrm{SNR}}{\sqrt{2T}} \left[\int_{f_{\min}}^{f_{\max}} \mathrm{d}f \, \frac{(f/f_*)^{2\beta}}{\Sigma^2(f)} \right]^{-1/2}.$$

For each pair (β, Ω_{β}) , $\Omega_{_{GW}}(f)$ is plotted against f.

The resulting envelope of the family of such curves is the PLS curve⁵:

$$\Omega_{_{\rm GW}}^{^{_{\rm PLS}}}(f) = \max_{\beta} \left[\Omega_{\beta} \; \left(\frac{f}{f_*} \right)^{\beta} \right]. \label{eq:GW}$$

⁵E. Thrane and J. D. Romano, Phys. Rev. D 88, 124032 (2013).

Plotting the envelope



PLS for ET

Integrated sensitivity curves



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Power-law (PL) and broken power-law (BPL) signals



A BPL signal can be parameterized as⁶:

$$\Omega_{\rm GW}(f) = \Omega_{\star} \left(\frac{f}{f_{\star}}\right)^{n_1} \left[\frac{1}{2} + \frac{1}{2}\left(\frac{f}{f_{\star}}\right)^{\sigma}\right]^{\frac{n_2 - n_1}{\sigma}}.$$

⁶LISA Cosmology Working Group, JCAP **10**, 020 (2024).

Power-law (CS) Fisher analysis



By operating together, LISA and ET can measure these parameters with better accuracy.

Broken power-law (PT) Fisher analysis



By operating together, LISA and ET can measure these parameters with an accuracy of $\sim 10\%$.

Lognormal signal



A lognormal signal can be parameterized as⁷:

$$\Omega_{\rm GW} = \Omega_{\star} \exp\left[-\frac{\ln^2\left(f/f_{\star}\right)}{2\rho^2}\right].$$

⁷R. Namba *et al.*, JCAP **01**, 041 (2016).

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Lognormal (inflation) Fisher analysis



By operating together, LISA and ET can measure these parameters with better accuracy.

Summary

- LISA and ET will be operational approximately around the same time.
- We are interested in a stochastic gravitational wave background that is characterized by a spectrum spanning a wide range of frequencies.
- We have investigated the consequences of using synergies between LISA and ET for improved detection of such a GW background.
- We have illustrated this effect using several examples of early-universe scenarios that lead to GW signals with different frequency profiles.

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Thank you!