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# *Non-Hermitian control of Hermitian waveguide arrays*

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# Introduction

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✓ Dynamical control of wave propagation in optical lattices execution  
of optical operations 

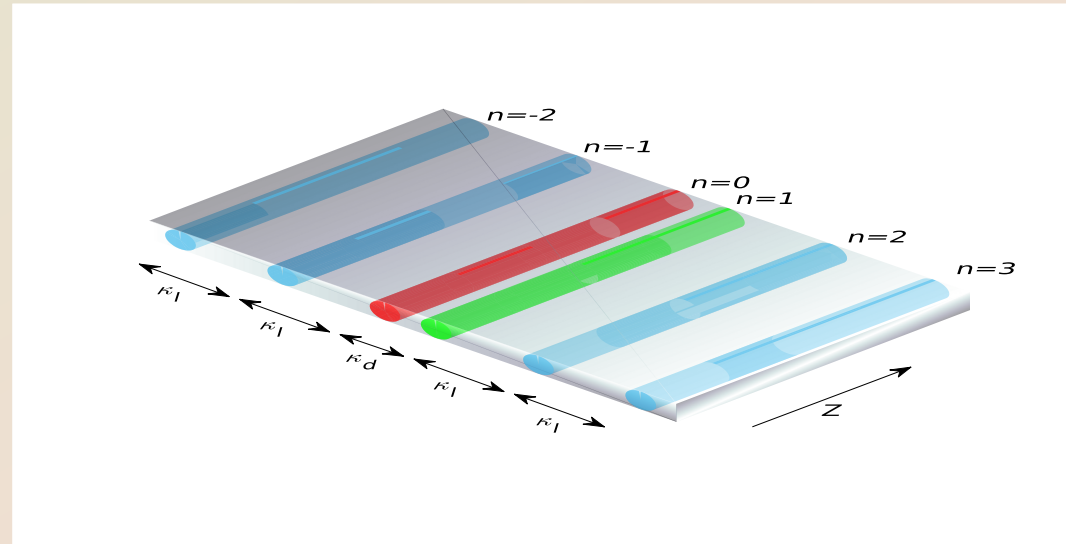
➤ Optical waveguides' arrays beam propagates under the influence of a spatially periodic, refractive index distribution

➤ in our case, wave dynamics approximated considering only amplitude/phase of each individual waveguide, simulating dynamical behavior, characteristic of discrete systems



# Objective

Device under investigation:  
array of optical waveguides  
with embedded, active,  
nonlinear elements



Analysis performed by following 3  
primary  
steps:

- theoretical background of wave propagation in individual elements
- verification through numerical analysis in the aforementioned elements
- numerical analysis on the device under investigation

# Coupled mode theory

## Modelling wave propagation in coupled waveguides (CMT)

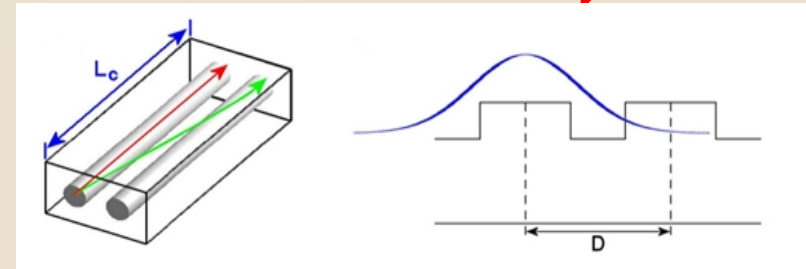


### Single-mode waveguide fields

$$\vec{E}_n = \sum_r E_{nr} \hat{r} = \sum_r e_{nr}(x, y) e^{ik_{nr}z} \hat{r}$$

$$\vec{H}_n = \sum_r H_{nr} \hat{r} = \sum_r h_{nr}(x, y) e^{ik_{nr}z} \hat{r}$$

general  
Maxwell  
equations



### Perturbed fields

$$\vec{E}' = \sum_{m,r} u_{mr}(z) e_{mr}(x, y) e^{ik_{mr}z} \hat{r}$$

$$\vec{H}' = \sum_{m,r} u'_{mr}(z) h_{mr}(x, y) e^{ik_{mr}z} \hat{r}$$



$$\nabla \times \vec{E}_n = i\omega\mu_0 \vec{H}_n$$

$$\nabla \times \vec{H}_n = -i\omega\epsilon_n \vec{E}_n$$

$$\nabla \times \vec{E}' = i\omega\mu_0 \vec{H}'$$

$$\nabla \times \vec{H}' = -i\omega(\epsilon' \vec{E}' + \vec{P}'^{NL})$$

nonlinear term -  
Kerr effect

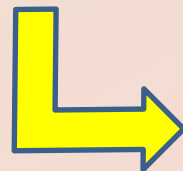
### General coupled mode equations:

$$i \frac{dA_{nx}}{dz} + k_{nx} A_{nx} + i \sum_{m \neq n} c_{mnxy} \frac{dA_{mx}}{dz} + \sum_{m \neq n} c_{mnxy} k_{mx} A_{mx}$$

$$+ \chi_{nx} A_{nx} + \sum_{m \neq n} \kappa_{nm} A_{mx} + \gamma_{nx} A_{nx} |A_{nx}|^2 = 0$$

Discrete nonlinear  
Schrodinger equation (DNLS)

assumptions



$$i \frac{dA_n}{dz} + \beta A_n + \kappa(A_{n+1} + A_{n-1}) + \gamma_n |A_n|^2 A_n = 0$$

# Hermitian lattice dynamics

## ➤ Linear waveguide array ( $\gamma_n=0$ )

- plane wave:

$$E_n(z) = \varphi_n e^{ink_x D} e^{ik_z z}$$



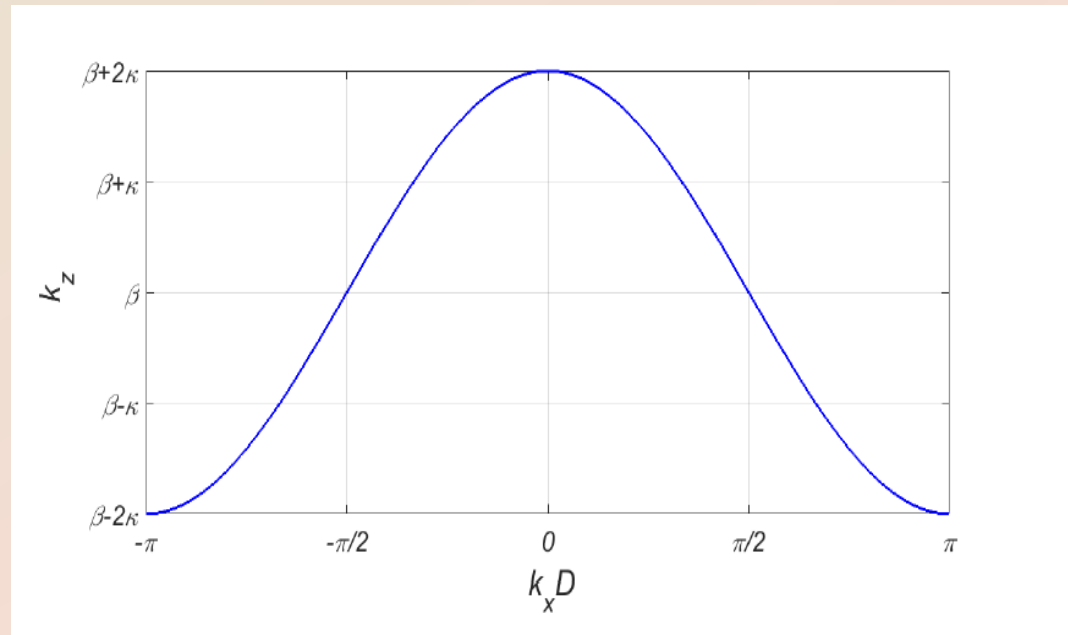
DNLS

$$k_z = \beta + 2\kappa \cos(k_x D)$$

1D waveguide array  
dispersion relation



e. phonon-like behaviour



# non-Hermitian dimer dynamics

➤ Asymmetric, nonlinear, active coupler:

$$\begin{aligned}\frac{dA_1}{dz} &= i(\beta_1 + i\alpha_1)A_1 + i\gamma(|A_1|^2 + \sigma|A_2|^2)A_1 + i\kappa A_2 \\ \frac{dA_2}{dz} &= i(\beta_2 + i\alpha_2)A_2 + i\gamma(|A_2|^2 + \sigma|A_1|^2)A_2 + i\kappa A_1\end{aligned}$$



Steady states:  
Nonlinear Supermodes (NS)

2 supported NS:

$$A_j = |A_j| \exp(i(bz + \varphi_j))$$

with:

$$|A_1|^2 = \frac{\beta_1 \alpha (\beta - 1) + \kappa \sqrt{\alpha} (\alpha - 1) \cos \varphi}{\gamma (1 - \sigma) (\alpha - 1)}$$

$$|A_2|^2 = \frac{1}{\alpha} |A_1|^2$$

where:

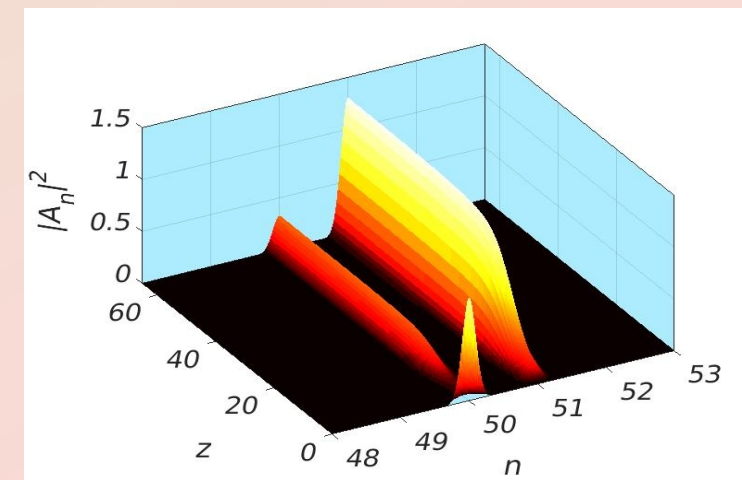
$$b = \frac{1}{2} \left[ \beta_1 (\beta + 1) + \kappa \frac{\alpha + 1}{\sqrt{\alpha}} \cos \varphi + \gamma (\sigma + 1) \frac{\alpha + 1}{\alpha} |A_1|^2 \right]$$

$$\sin \varphi = \frac{|\alpha_1| \sqrt{\alpha}}{\kappa}$$

$$\alpha \equiv -\alpha_2 / \alpha_1$$

$$\beta \equiv \beta_2 / \beta_1$$

e.g.  $\alpha_1=1$ ,  $\alpha_2=-0.25$ ,  $\beta_1=1$ ,  $\beta_2=2$ ,  
 $\gamma=-1$ ,  $\sigma=0$ ,  $\kappa=0.5$



# non-Hermitian dimer dynamics

How dimer dynamics determine array's behaviour

Stability of zero (ground state solution) is on/off

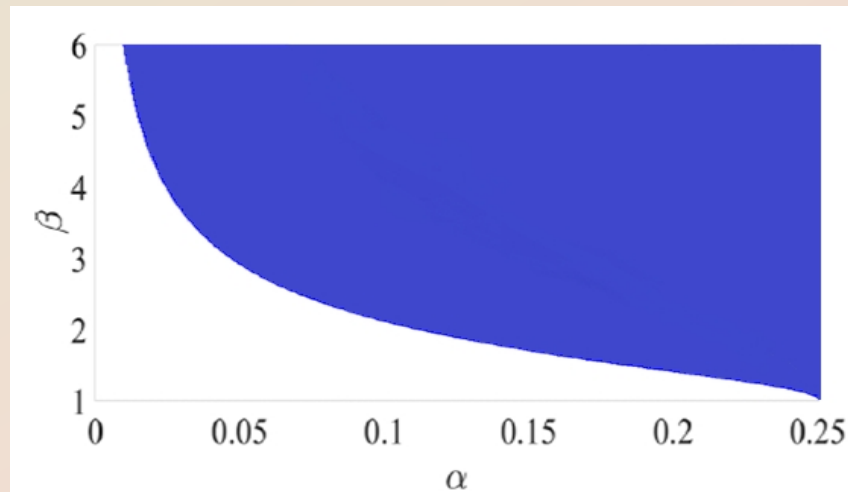
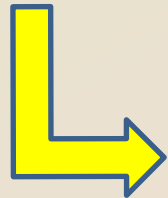
whether di 

Stability of the NS in ing perturbations



$$u_{1,2} = (A_{1,2} + \varepsilon_{1,2}) e^{ibz}$$

$$\varepsilon_{1,2} = c_{1,2} e^{-i\lambda z} + d_{1,2}^* e^{i\lambda^* z}$$



simulating noise in realistic applications

stability region, i.e. eigenvalues of the system with  $\text{Im}(\lambda) < 0$

$$\gamma = -1, \sigma = 0, \kappa_d = 0.5$$



# Model

➤ Coupled mode equations under the assumption of nearest neighbour coupling:

$$\frac{dA_n}{dz} = i \left( (\beta_n + i\alpha_n)A_n + \kappa_{n+1}A_{n+1} + \kappa_{n-1}A_{n-1} + \gamma_n A_n |A_n|^2 \right)$$

- discrete array of linear, passive waveguides  
( $\gamma_n=0$ ,  $\alpha_n=0$ , without loss of generality  $\beta_n=0$ )
- asymmetric, active dimer, with self-defocusing nonlinearity ( $\gamma_{N/2}=\gamma_{N/2+1}=-1$ )
- we investigate wave propagation by varying array's parameters,  $\alpha_1$ ,  $\alpha_2$ ,  $\beta_1$ ,  $\beta_2$ ,  $\kappa_l$ ,  $\kappa_d$
- evaluating array's properties, with the dimer acting as a scatterer, an input/output device and a detector
- scattered beam      discrete gaussian with initial amplitude profile:



$$f(n, z = 0) = \frac{1}{\sigma\sqrt{2\pi}} \exp \left( ik(n - n_i) - \frac{(n - n_i)^2}{2\sigma^2} \right)$$



# Scattering dynamics

Plane wave  $\rightarrow$  good approximation for a sufficiently wide gaussian beam

$$A_n = \exp(ikn - i\beta z) + R \exp(-ikn - i\beta z), \quad n \leq 0$$

$$A_n = T \exp(ikn - i\beta z), \quad n \geq 1$$

reflection/transmission coefficients for the linear case (laser switched off)

DNLS for dispersionless beam  
( $k = \pi/2$ )

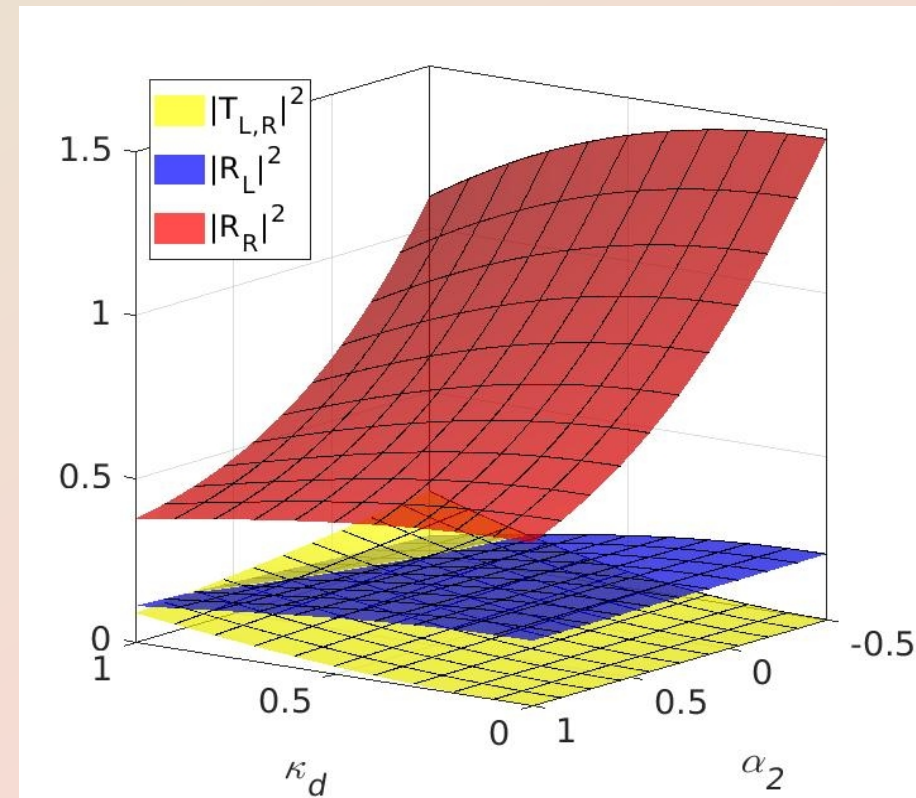
$$|T_L|^2 = \frac{4\bar{\kappa}^2}{D}$$

$$|R_L|^2 = \frac{\{\bar{\kappa}^2 + \bar{\kappa}(\bar{\alpha}_2 - \bar{\alpha}_1) - 1 + \bar{b}_1\bar{b}_2 - \bar{\alpha}_1\bar{\alpha}_2\}^2 + \{\bar{\kappa}(\bar{b}_2 - \bar{b}_1) - (\bar{\alpha}_1\bar{b}_2 + \bar{b}_1\bar{\alpha}_2)\}^2}{D}$$

where:

$$D = \{\bar{\kappa}^2 + \bar{\kappa}(\bar{\alpha}_1 + \bar{\alpha}_2) + 1 - \bar{b}_1\bar{b}_2 + \bar{\alpha}_1\bar{\alpha}_2\}^2 + \{\bar{\kappa}(\bar{b}_1 + \bar{b}_2) + (\bar{\alpha}_1\bar{b}_2 + \bar{b}_1\bar{\alpha}_2)\}^2$$

$$\bar{\kappa} = \frac{\kappa_l}{\kappa_d}, \quad \bar{\alpha}_n = \frac{\alpha_n}{\kappa_d}, \quad \bar{b}_n = \frac{b_n}{\kappa_d}$$



$\alpha_1=1, b_1=1,$   
 $b_2=2, \kappa_l=1$

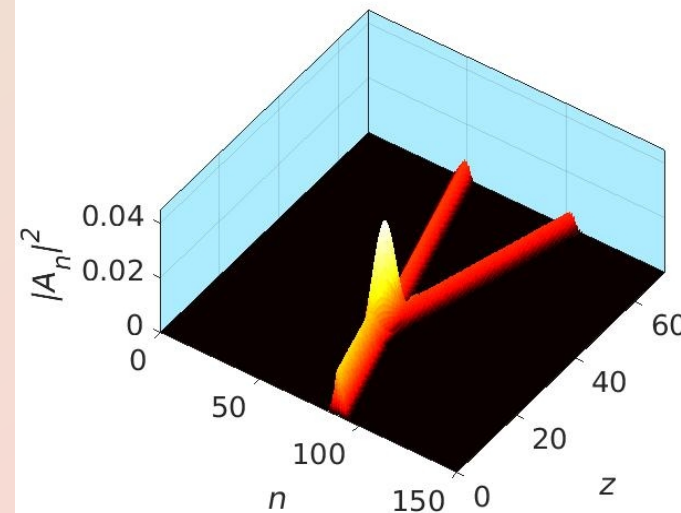
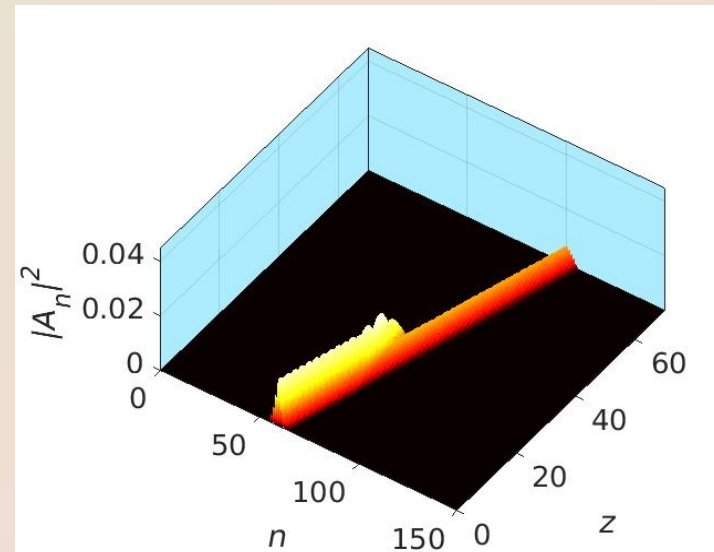
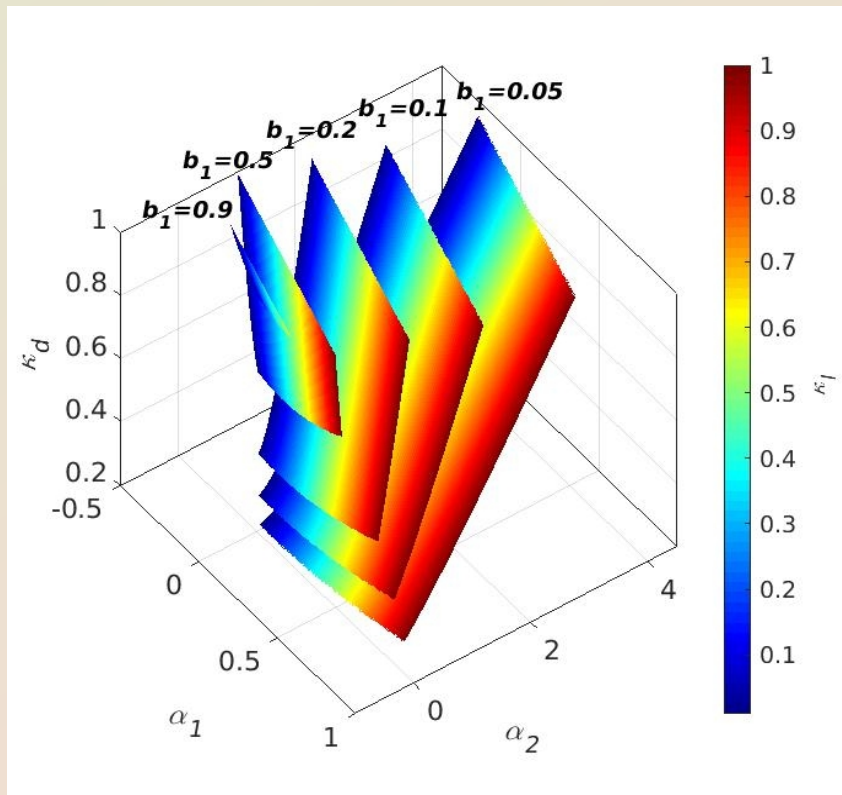
$R_R, T_R \rightarrow$  changing  $\bar{\alpha}_1, \bar{b}_1$  with  $\bar{\alpha}_2, \bar{b}_2$

# Scattering dynamics

➤ Zero reflection (geometrically dissimilar dimer ( $b_1/b_2=1$ )):

gaussian wavepacket simulation

PW calculated parameters

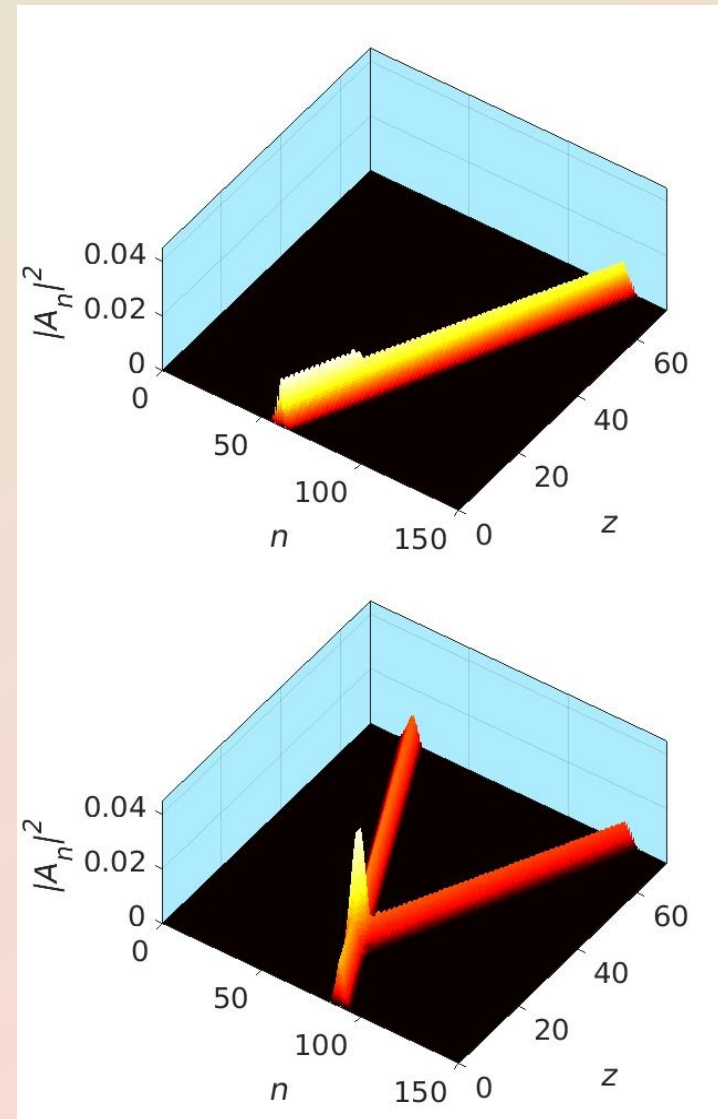
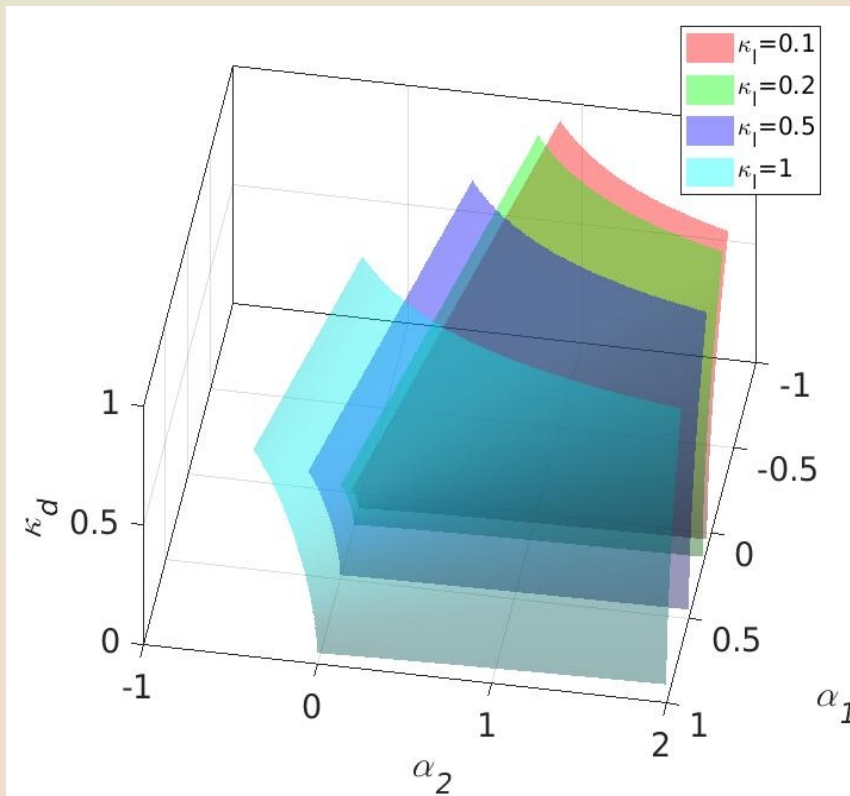


$a_1=-0.124$   
 $a_2=0.548$   
 $\kappa_l=0.300$   
 $\kappa_d=0.927$   
 $b_1=0.5$   $b_2=1$

# Scattering dynamics

➤ Zero reflection (geometrically identical waveguides ( $b_n=0$ )): gaussian wavepacket simulation

PW calculated parameters

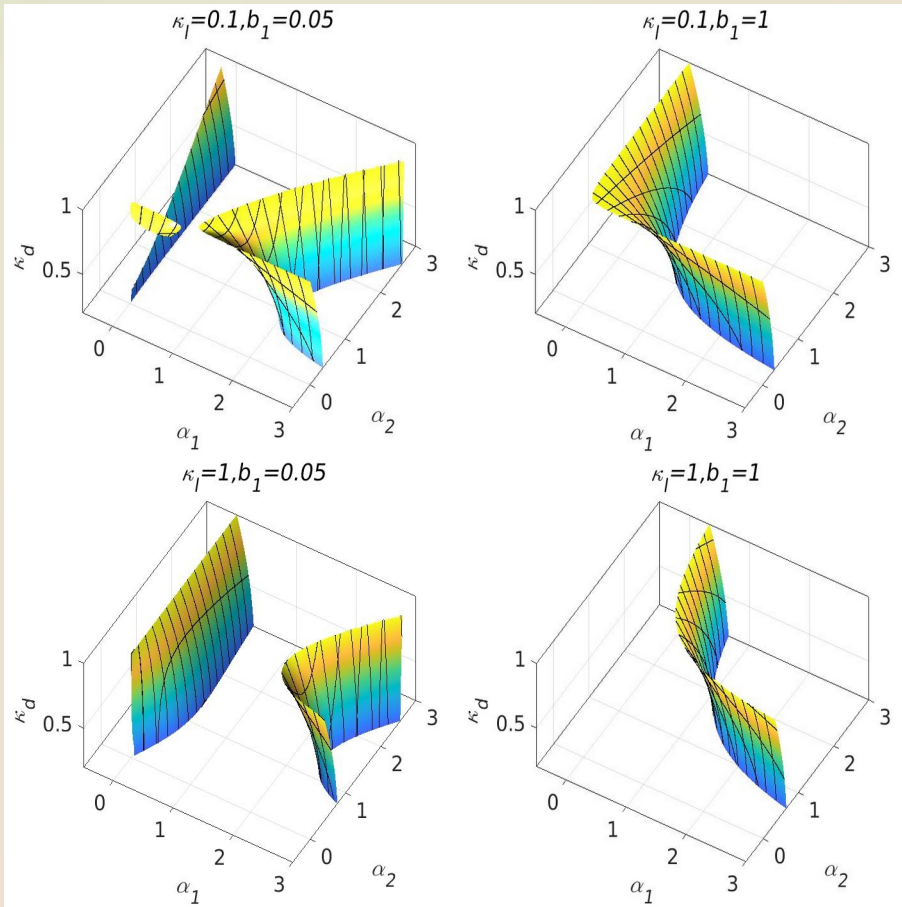


$a_1=-0.301$   
 $a_2=0.472$   
 $\kappa_l=0.500$   
 $\kappa_d=0.882$   
 $b_1=0$   
 $b_2=0$

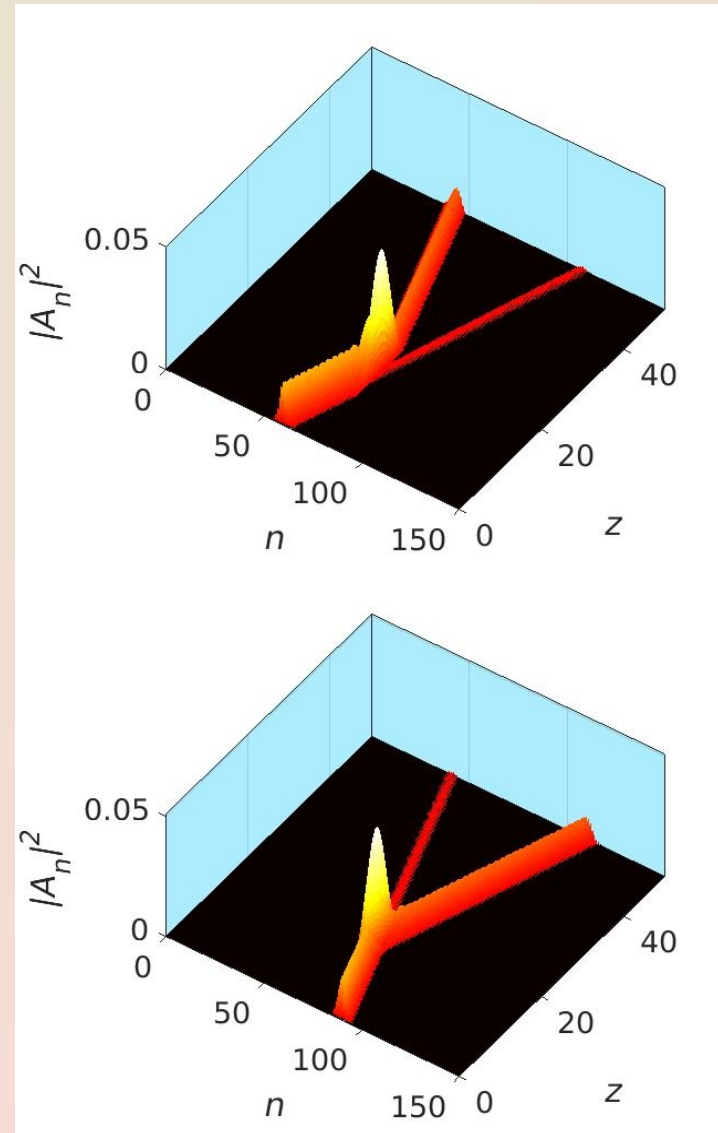
# Scattering dynamics

## Symmetrical reflection:

### PW calculated parameters



## gaussian wavepacket simulation



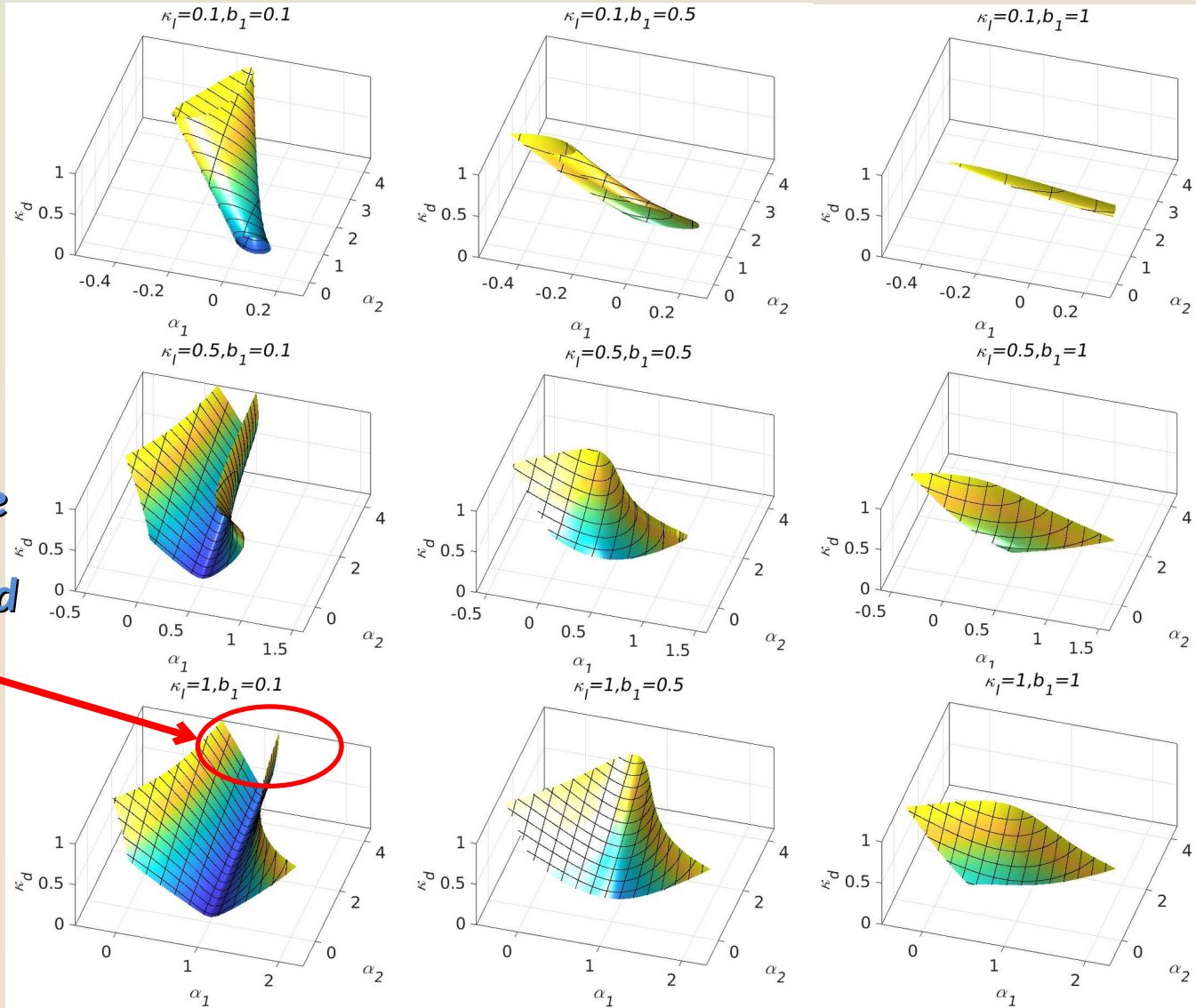
$a_1=-0.100$   
 $a_2=3.43$   
 $\kappa_l=0.500$   
 $\kappa_d=0.795$   
 $b_1=0.1$   
 $b_2=1$



# Scattering dynamics

## Reflection equal to transmission:

### PW calculated parameters

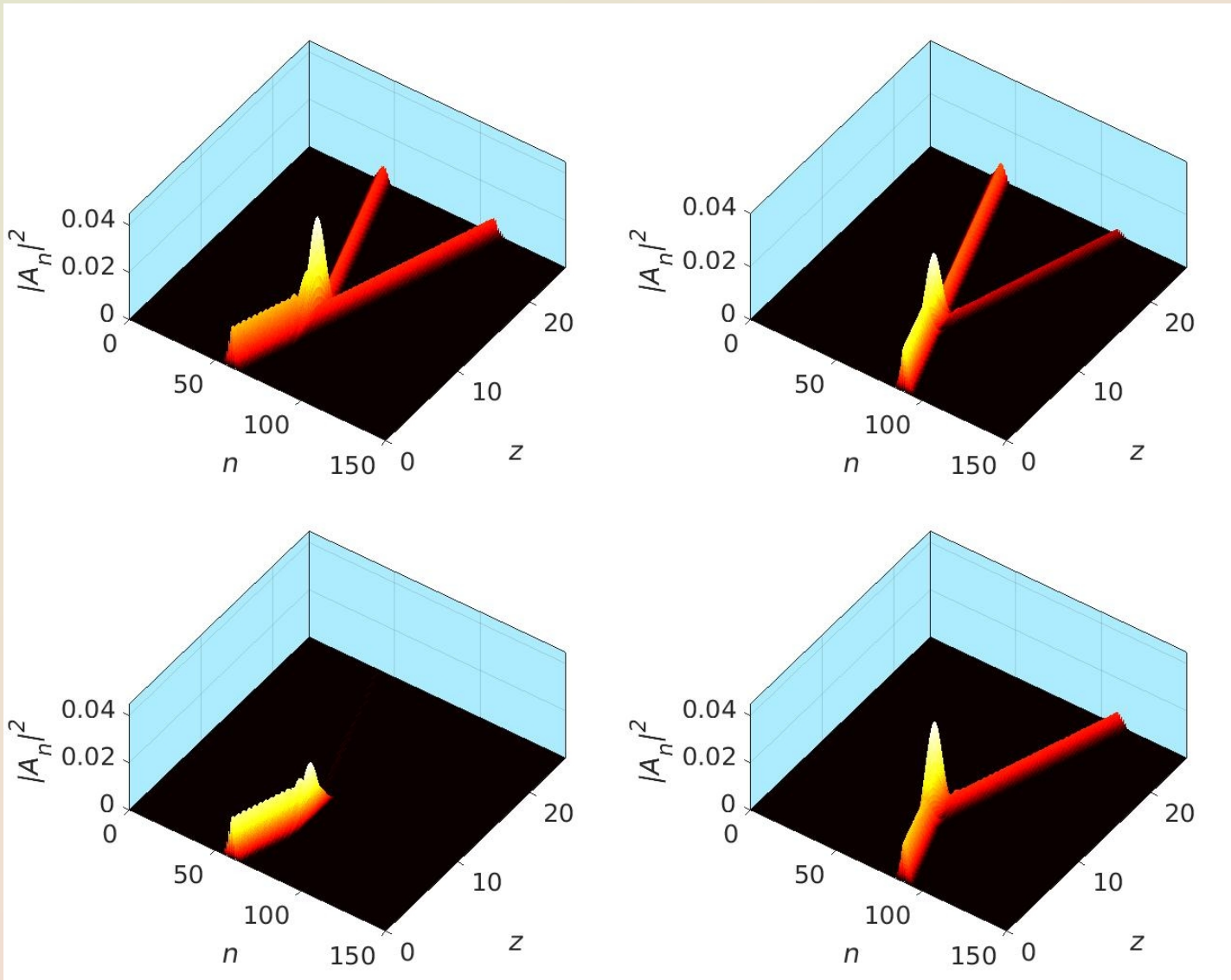


areas with large  $\alpha_1, \alpha_2$ , additionally lead to  $R, T \approx 0$ .

# Scattering dynamics

## Reflection equal to transmission:

gaussian wavepacket simulation



$a_1 = -0.144$   
 $a_2 = 0.868$   
 $\kappa_1 = 1$   
 $\kappa_d = 0.877$   
 $b_1 = 0.1$   
 $b_2 = 1$

$a_1 = 1.177$   
 $a_2 = 4.097$   
 $\kappa_1 = 1$   
 $\kappa_d = 0.799$   
 $b_1 = 0.1$   
 $b_2 = 1$

# Scattering dynamics

## Nonlinear scattering (dimer on):

gaussian wavepacket simulation

$$\begin{aligned} a_1 &= 1, a_2 = -0.05 \\ \kappa_d &= 0.45, \kappa_l = 0.5 \\ b_1 &= 1, b_2 = 6 \end{aligned}$$

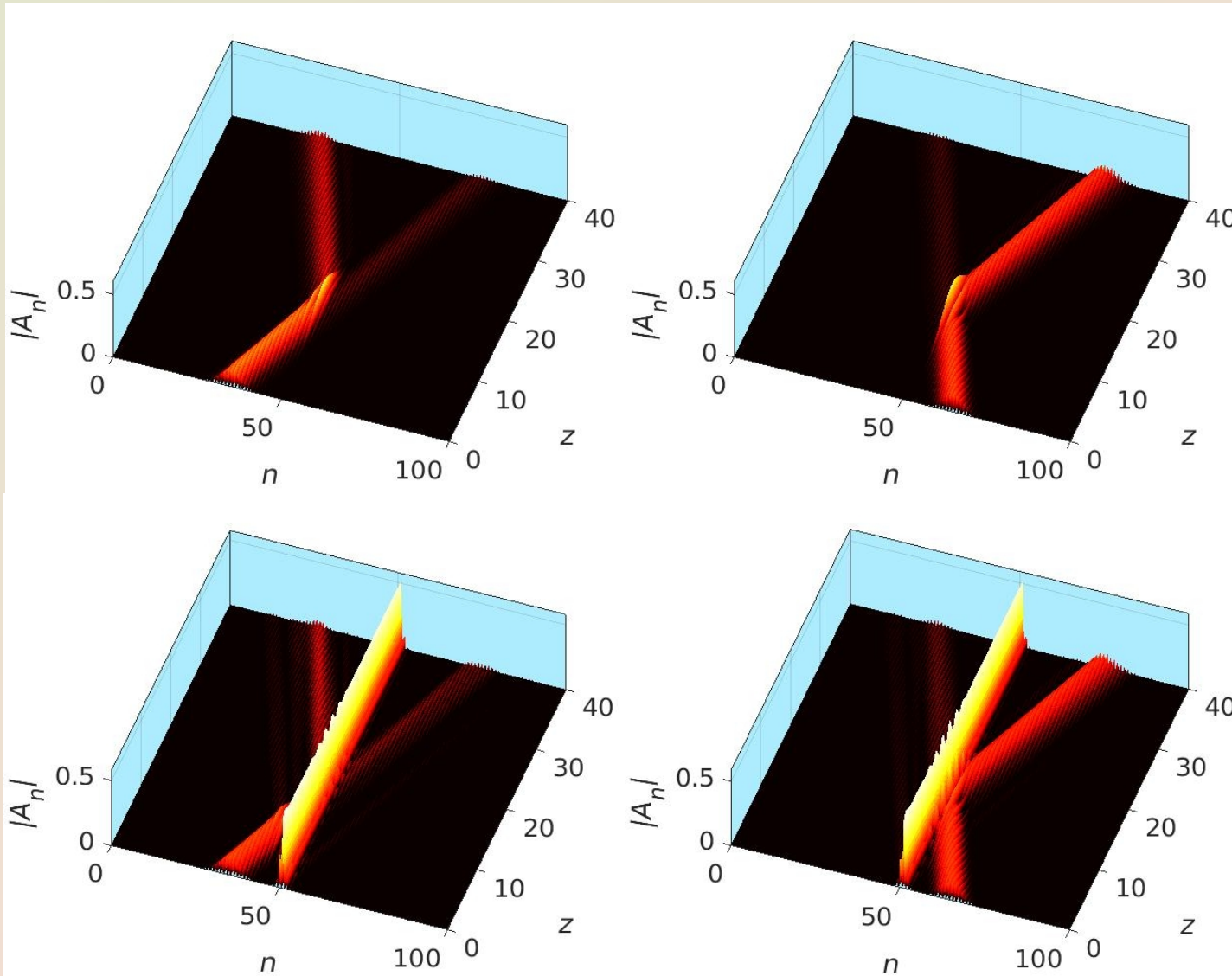
Dimer off

$$\begin{aligned} |R_L|^2 &= 0.25 \\ |T_{L,R}|^2 &= 0.06 \\ |R_R|^2 &= 0.92 \end{aligned}$$

no qualitative difference

Dimer on

$$\begin{aligned} |R_L|^2 &= 0.21 \\ |T_L|^2 &= 0.10 \\ |R_R|^2 &= 0.86 \\ |T_R|^2 &= 0.12 \end{aligned}$$





# Embedded dimer dynamics

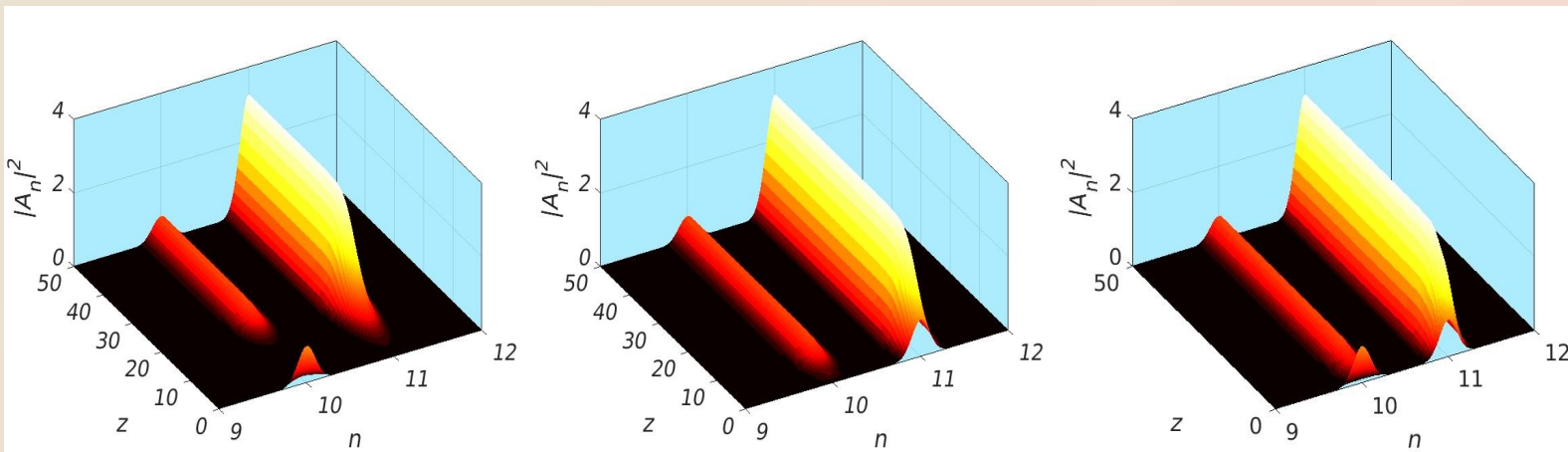
## Input/output device:

Isolated dimer:



- test for different initial conditions (IC) i.e. input in loss/gain/both dimer's waveguides
- taking advantage of different convergence speed to final state (NS)

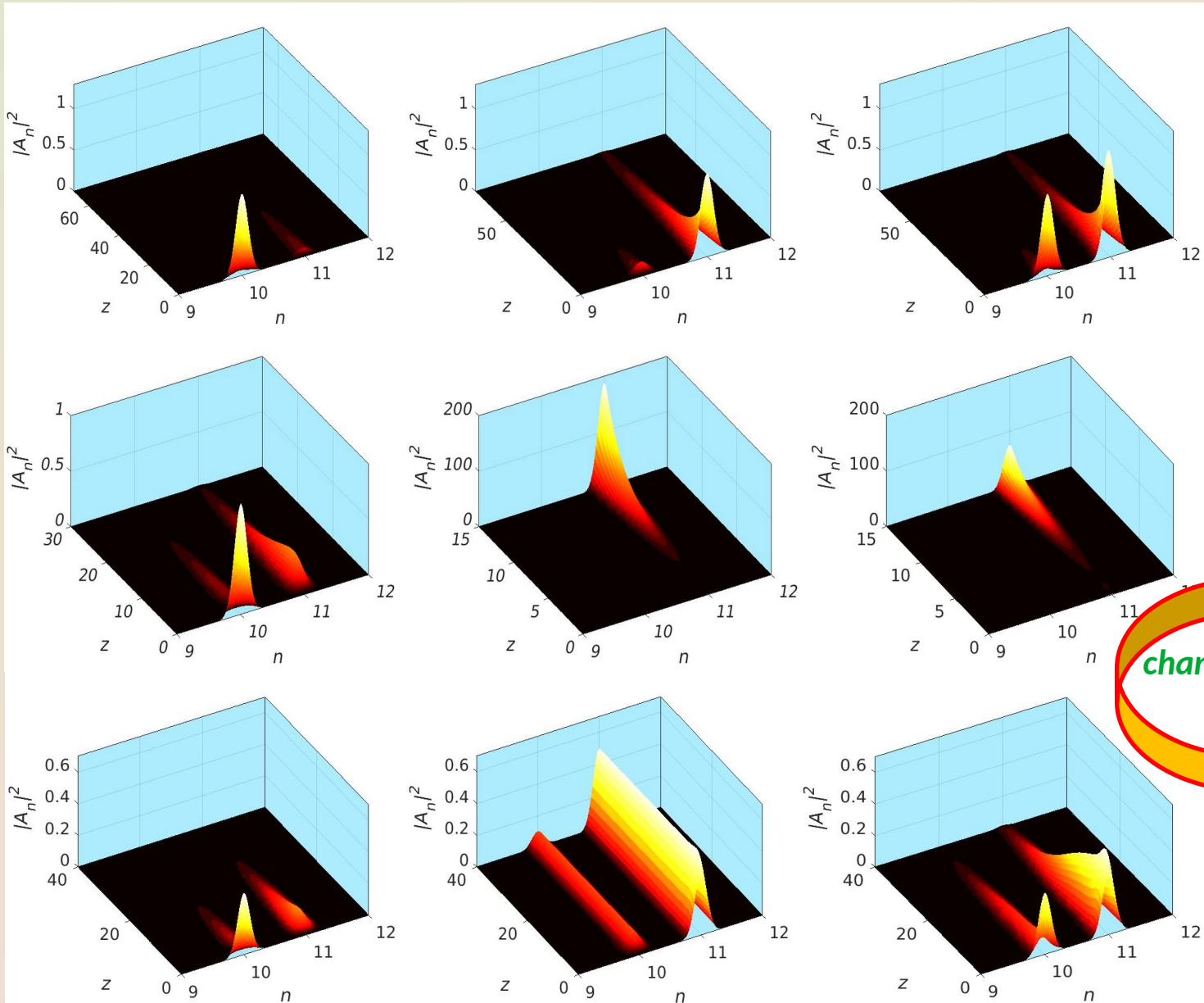
identical output irrespectively of IC  
(with  $(|A_2|/|A_1|)^2 = -a_1/a_2$ )



$a_1=1, a_2=-0.2, b_1=1,$   
 $b_2=4, \kappa_d=0.5, P_0=1$

# Embedded dimer dynamics

But...



$a_1=1, a_2=-0.05, b_1=1, b_2=2.5, \kappa_d=0.5, P_0=1$   
zero output irrespectively of IC

$a_1=1, a_2=-0.2, b_1=1, b_2=1, \kappa_d=0.5, P_0=1$   
infinite/zero output

$a_1=1, a_2=-0.2, b_1=1, b_2=1, \kappa_d=0.5, P_0=0.6256$   
finite/zero output

changing input power

# Embedded dimer dynamics

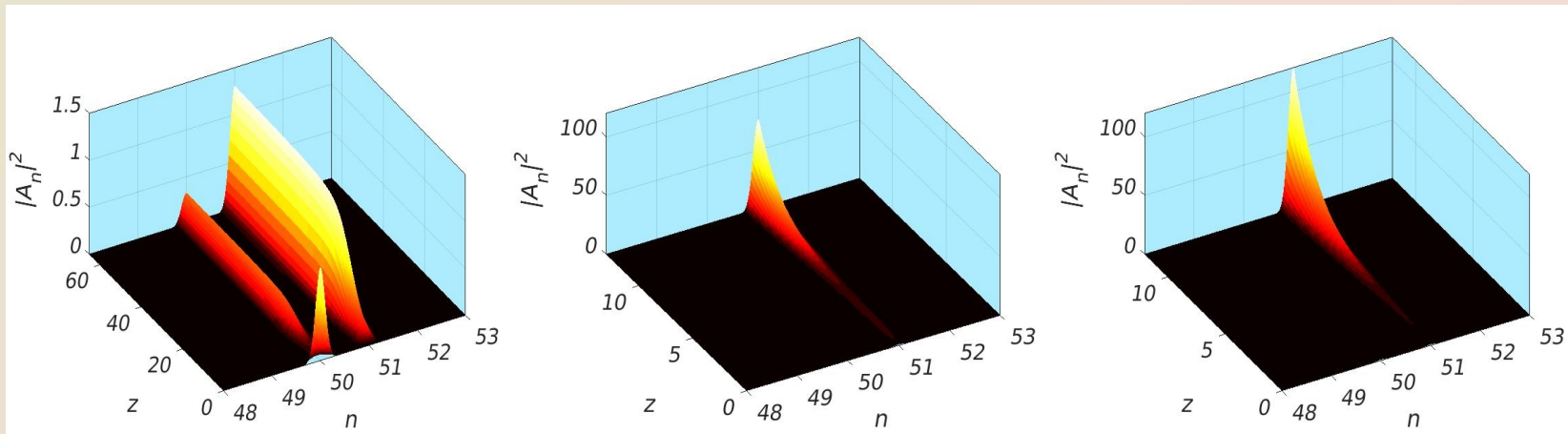
## Input/output device:

- Treating infinities:



stabilizing through coupling with the linear lattice ( $\kappa_l \neq 0$ )

e.g.  $a_1=1$ ,  $a_2=-0.25$ ,  $b_1=1$ ,  $b_2=2$ ,  $\kappa_d=0.5$ ,  $P_0=1$

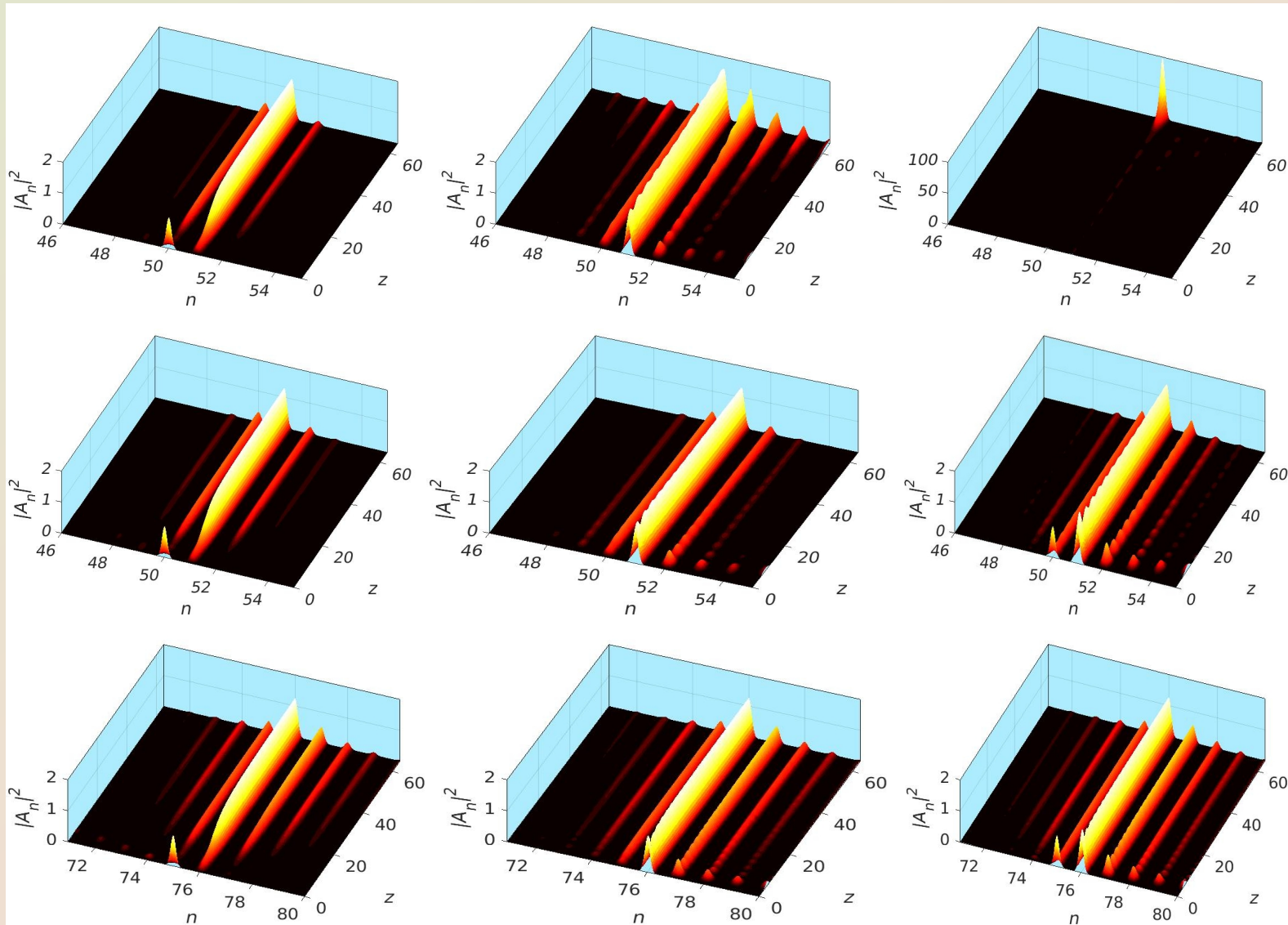


$\kappa_l = 0$

uncoupled dimer  
(only loss input  
stable)

# Embedded dimer dynamics

But...



$\kappa_I = 0.225$   
gain input  
stabilized

$\kappa_I = 0.3$   
all inputs  
stabilized

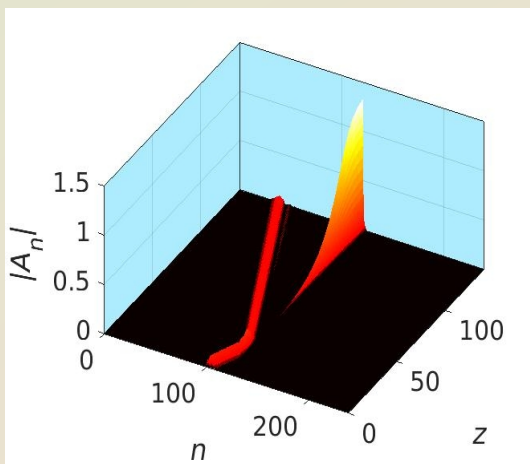
$\kappa_I = 0.5$



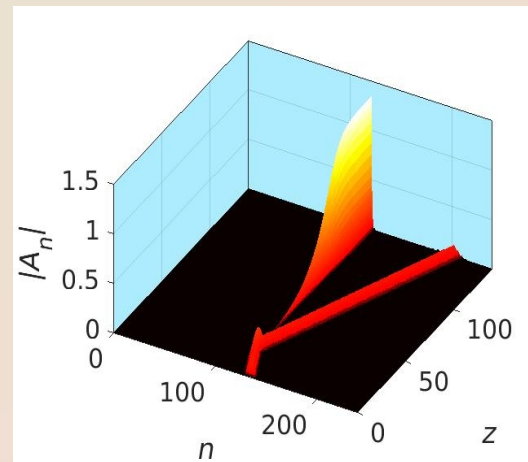
# Embedded dimer dynamics

## Array as a detector:

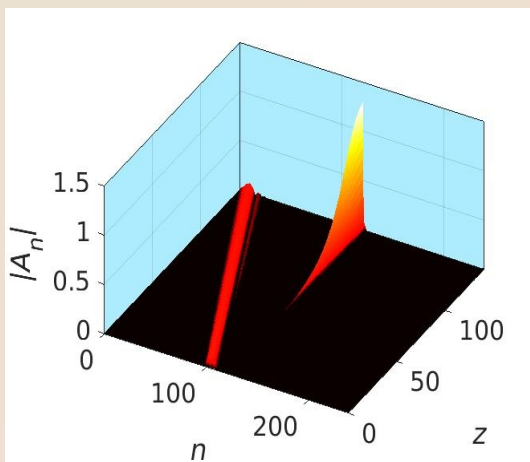
e.g.  $a_1=1$ ,  $a_2=-0.08$ ,  $b_1=1$ ,  $b_2=3$ ,  $\kappa_l=0.35$ ,  $\kappa_d=0.3$



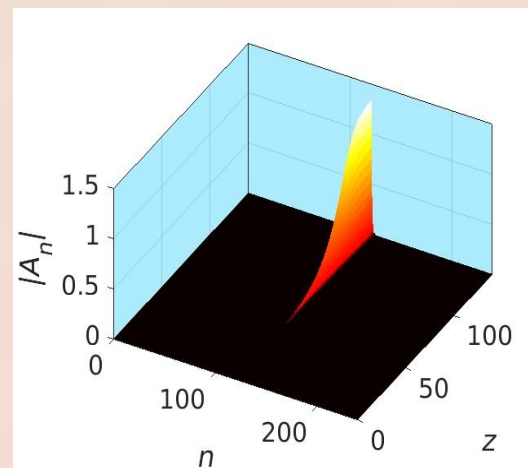
beam scattered by  
'loss' waveguide



beam scattered by  
'gain' waveguide



non-interacting  
beam



ring input

# Conclusions

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*In total, embedding an active, nonlinear, asymmetric dimer in a waveguide array:*

- i. offers the ability to adjust asymmetric scattering of beams*
- ii. offers the ability to adjust power output during dimer excitation by controlling IC, gain/loss and  $\kappa_d$*
- iii. has a stabilizing effect on propagating waveforms*
- iv. can be used for detection of propagating waves*



*Dimer controls scattering properties*

*Lattice modifies dimer dynamics*

*The End*

*Thank you for your attention!*