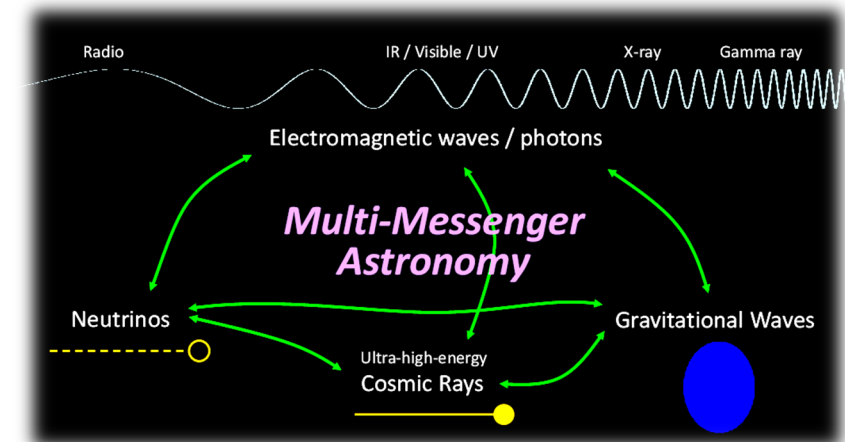
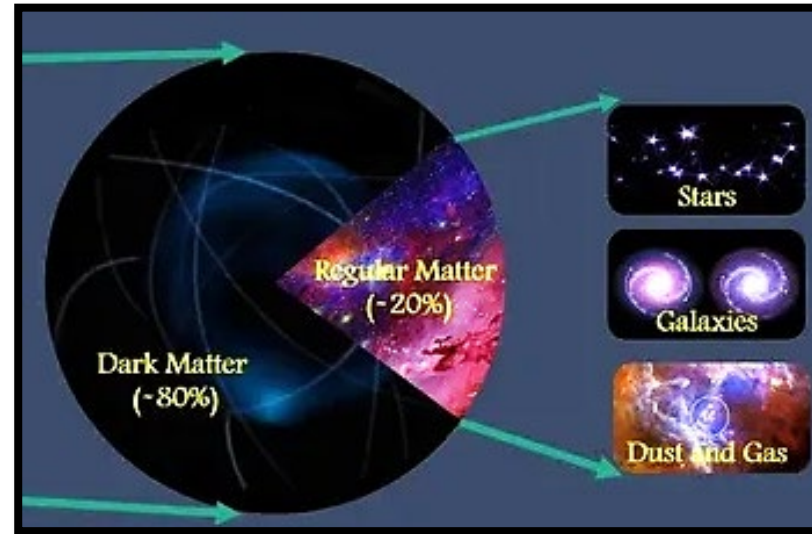


Inclusion of bosonic dark matter in neutron star to satisfy the observable features induced by nuclear matter models



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*The 60th Karpacz Winter School on Theoretical Physics and WE-Heraeus Physics School
16-25 May 2024*



D.R. K, S. Shakeri, V. Sagun, O. Ivanytskyi,
Bosonic dark matter in neutron stars and its effect on gravitational wave signal
Phys. Rev. D 105, 023001 (2022), [arXiv:2109.03801v2]

S. Shakeri, D.R. K,
Bosonic Dark Matter in Light of the NICER Precise Mass-Radius Measurements
Phys. Rev. D 109, 043029 (2024), [arXiv:2210.17308v2]

D.R. K, M. ShahrbaF, S. Shakeri, S. Typel
Exploring the distribution and impact of bosonic dark matter in neutron stars
Particles 7 (2024) 1, 201-213, [arXiv:2402.18696]

M. ShahrbaF, D.R. K, S. Typel
Constraints on the mass of a bosonic dark matter candidate within the DD2Y-T model
arXiv:2402.18686

D.R. K, S. Shakeri, V. Sagun, O. Ivanytskyi,
Tidal deformability as a probe of dark matter in neutron stars
MG16 Proceedings, (**World Scientific pp. 3713-3731 (2023)**), [arXiv:2112.14231]



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Mahboubeh ShahrbaF

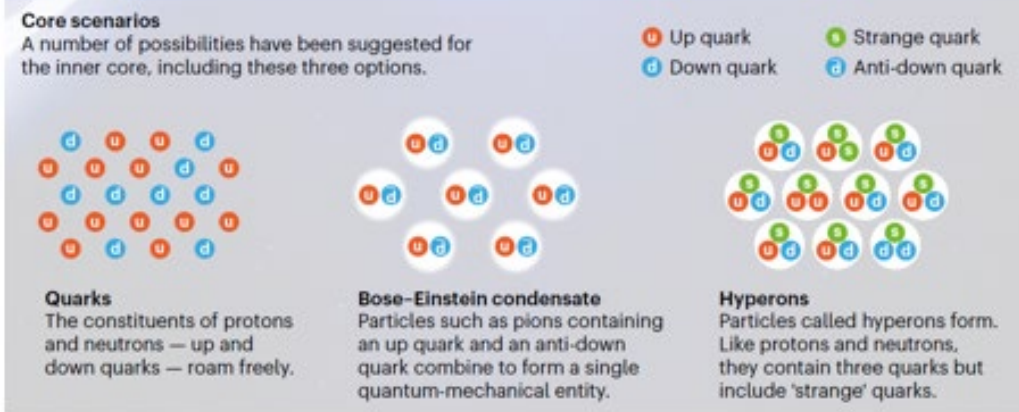
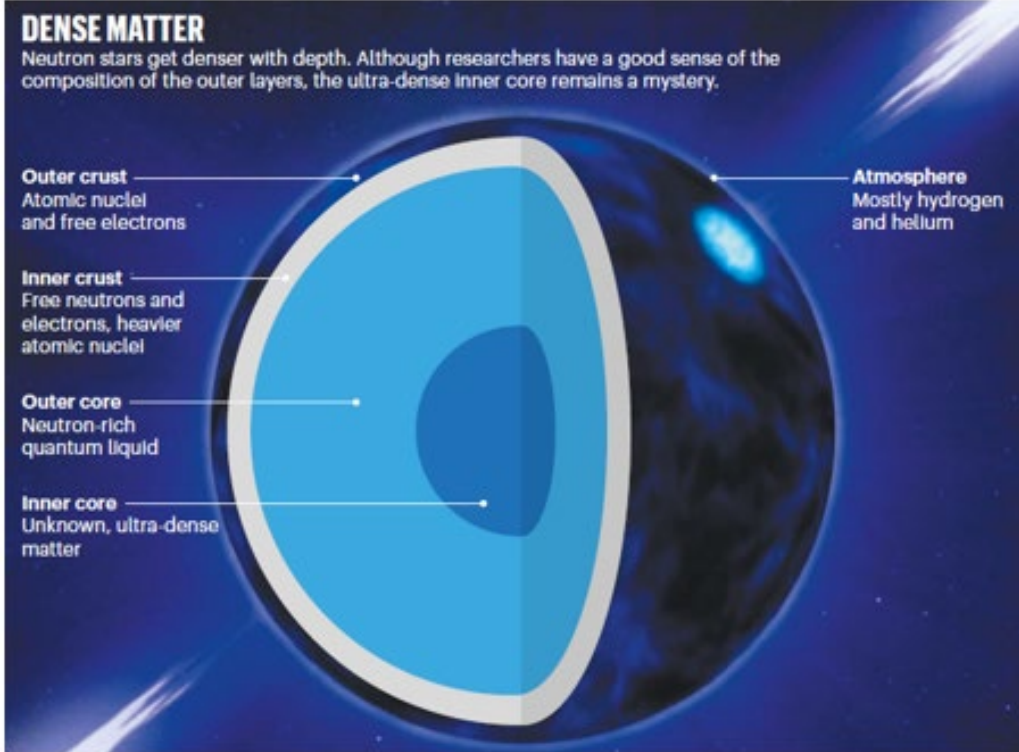
University of Wroclaw, Poland



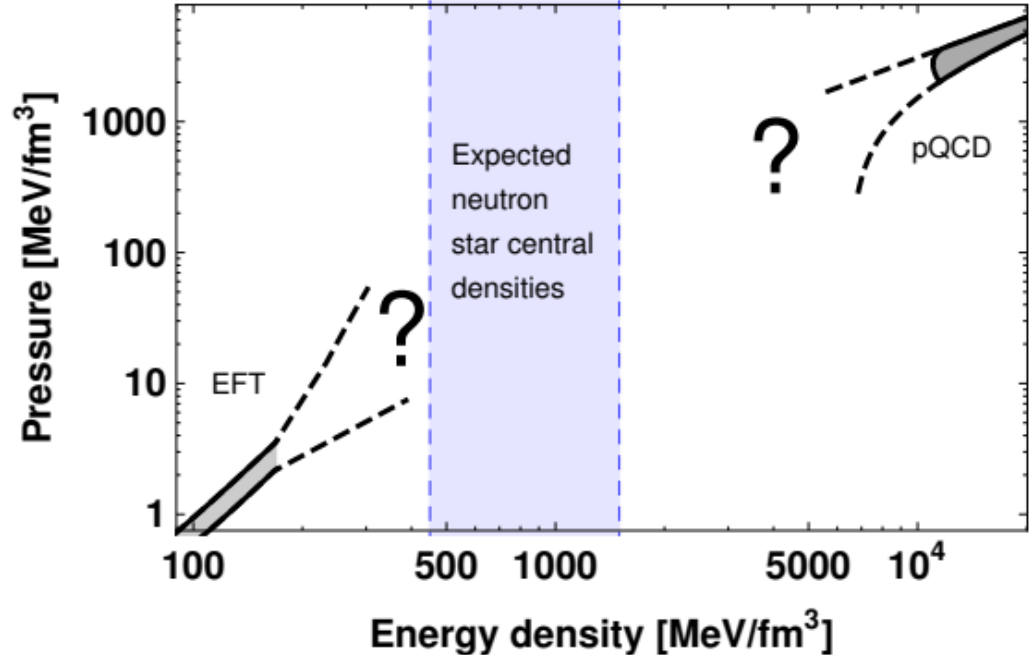
Stefan Typel

Darmstadt University and GSI institute, Germany

Neutron stars as a natural laboratory for high density matter



A. Mann, *Nature (London)* 579, 20 (2020)



Matti Järvinen, *Eur. Phys. J. C* (2022) 82:282

$$\frac{dP(r)}{dr} = -\frac{GM(r)\epsilon(r)}{c^2 r^2} \left[1 + \frac{P(r)}{\epsilon(r)} \right] \left[1 + \frac{4\pi r^3 P(r)}{M(r)c^2} \right] \left[1 - \frac{2GM(r)}{c^2 r} \right]^{-1}$$

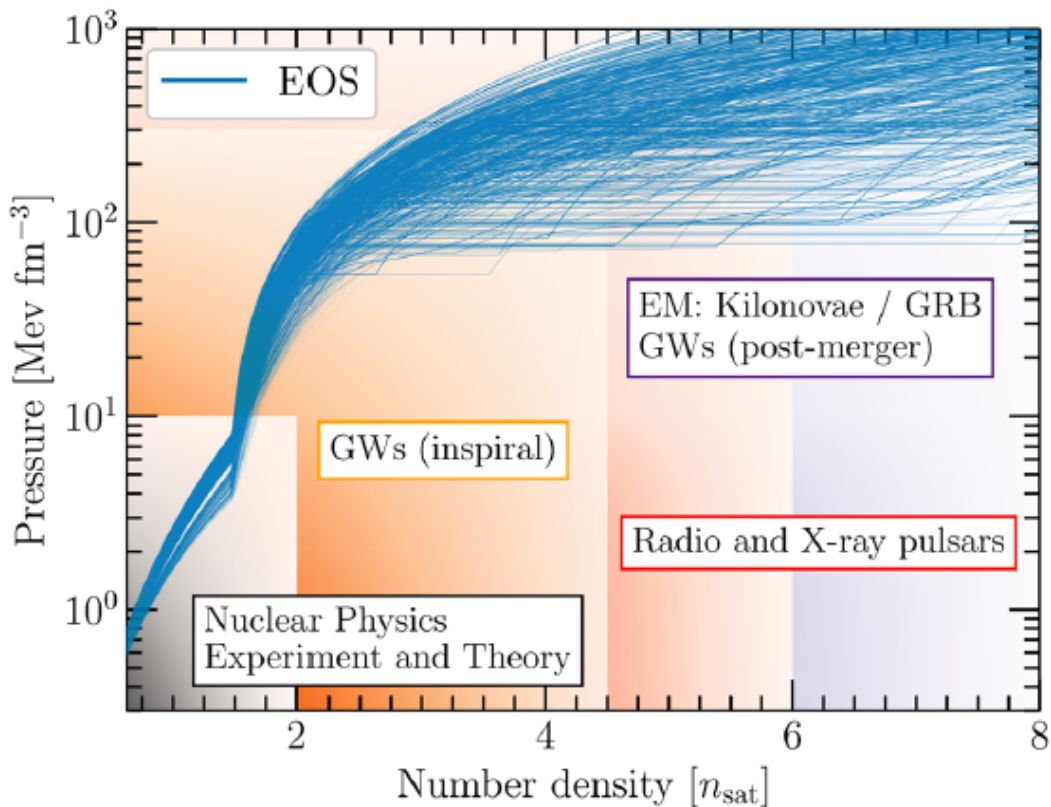
$$\frac{dM(r)}{dr} = \frac{4\pi r^2 \epsilon(r)}{c^2}$$

Tolman-Oppenheimer-Volkoff
(TOV) equations

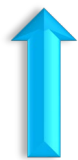
R. C. Tolman, *Phys. Rev.* 55, 364 (1939).

J. R. Oppenheimer and G. M. Volko, *Phys. Rev.* 55,374 (1939).

Multi-messenger observations of neutron stars



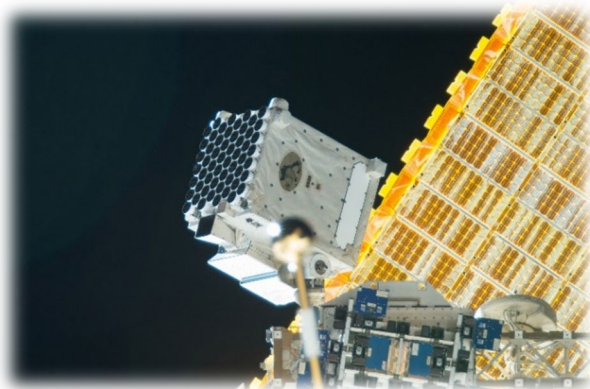
Peter T. H. Pang, et al. , *Nature Commun.* 14 (2023) 1, 8352



Radio telescopes



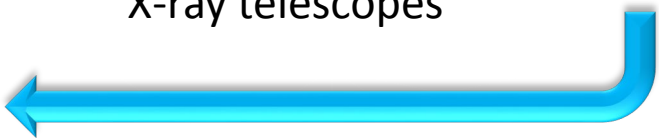
Optical telescopes



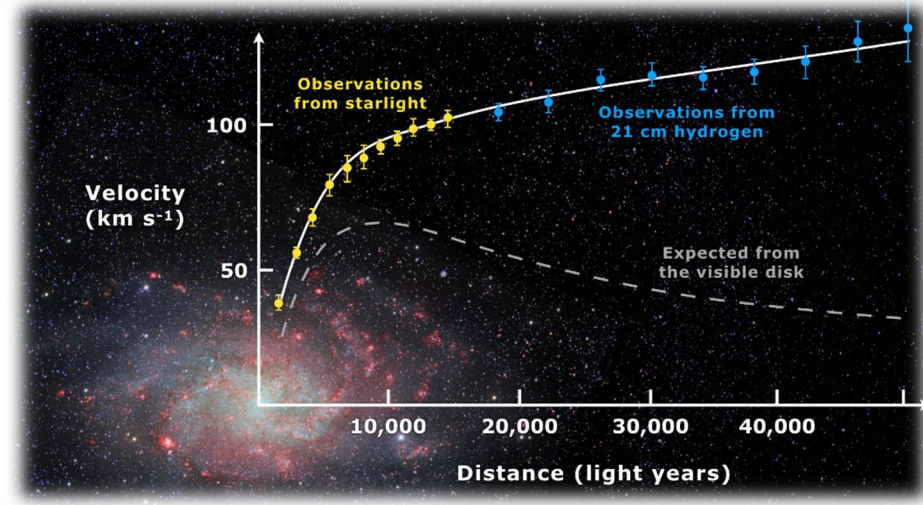
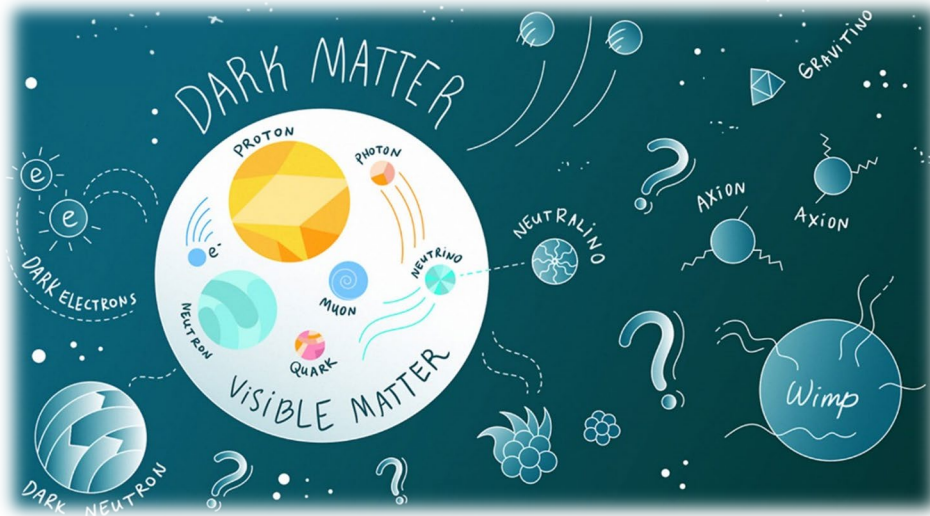
X-ray telescopes



Gravitational-Wave detectors



Observable features of neutron stars such as mass, radius and tidal deformability



Gravitationally stable astrophysical objects composed of dark matter

Fermionic or Bosonic dark matter

Dark Star
Boson or Fermion star

Dark matter admixed
neutron star

Andrea Maselli, et al. [PRD 96, 023005 \(2017\)](#)
 Joshua Eby, et al. [JHEP 02 \(2016\) 028](#)
 G. Narain, J. Schaffner-Bielich, et al. [PRD 74, 063003 \(2006\)](#)
 Chris Kouvaris, et al. [PRD 92 \(2015\) 6, 063526](#)
 P.A.Seoane, J.Barranco, A.Bernal, L. Rezzolla, [JCAP 11 \(2010\) 002](#)
 S. L. Liebling, C. Palenzuela, [Living Rev.Rel. 26 \(2023\) 1](#)
 Luca Visinelli, [Int.J.Mod.Phys.D 30 \(2021\) 15, 2130006](#)

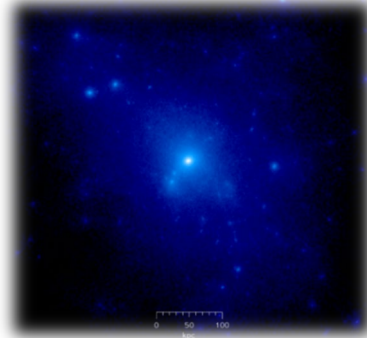
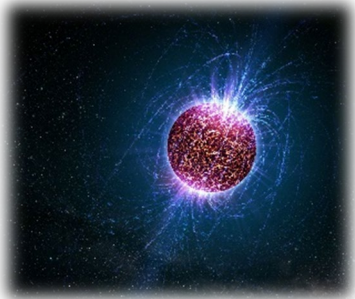
E. Giangrandi, V. Sugun, O. Ivanytskyi, C. Providência, T. Dietrich
[Astrophys.J. 953 \(2023\) 1, 115](#)
 H.M Liu, J.B. Wei, Z.H. Li, G.F. Burgio, H.-J. Schulze, [Phys.Dark Univ. 42 \(2023\) 101338](#)
 Harish Chandra Das, et al. [Mon.Not.Roy.Astron.Soc. 495 \(2020\) 4893-4903](#)
 A. Nelson, S. Reddy, D. Zhou, [JCAP07\(2019\)012](#)
 John Ellis, et al. [PRD 97, 123007 \(2018\)](#)
 Y.Dengler, J. Schaffner-Bielich, L. Tolos, [PRD 105 \(2022\) 4, 043013](#)

Dark matter (DM) admixed neutron star (NS)

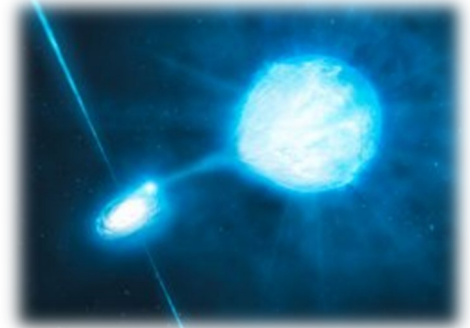
Ask Oleksii

Accumulation of DM by a star or a NS during its life time

A) Progenitor B) Main sequence (MS) star, C) Supernova explosion & formation of a proto-NS D) Equilibrated NS



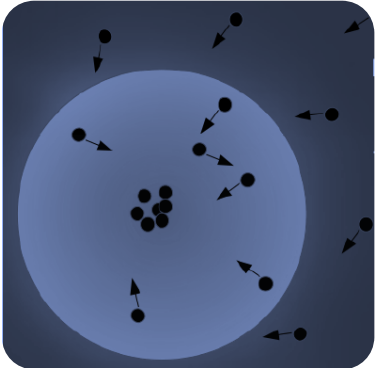
NS exists in a dense halo or region of DM or passes through it (Near the center of galaxy)



Dark star as an accretion center of baryonic matter

DM production in the NS matter or supernova explosions

Accretion of DM into a NS



Neutron decay anomaly

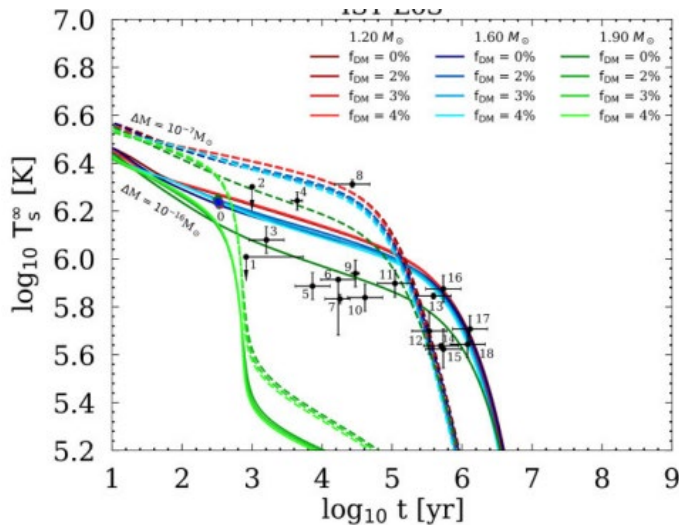
DM capture by NS in a binary system including Dark star or Dark star – NS merger



A. Nelson, S. Reddy, D. Zhou, *JCAP*07(2019)012
 John Ellis, et al. *PRD* 97, 123007 (2018)
 A. Del Popolo, et al. *Universe* 6 (2020) 12, 222

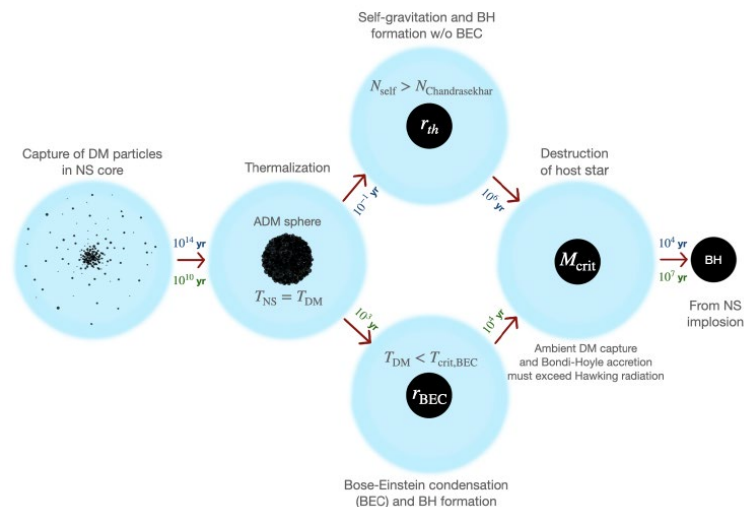
D.R. Karkevandi, S. Shakeri, V. Sagun, O. Ivanytskyi, *PRD* 105, 023001 (2022) P. Ciarcelluti & F. Sandin. *Phys.Lett.* B695:19-21,2011
 O. Ivanytskyi, V. Sagun, I. Lopes. *PRD* 102, 063028 (2020) S. Shirke, S. Ghosh, D. Chatterjee, L. Sagunski, J. Schaffner-Bielich, *JCAP* 12 (2023) 008
 Raul Ciancarella, et al. *Phys.Dark Univ.* 32 (2021) 100796 M. Deliyergiyev, A. Del Popolo, M. Le Delliou, *Mon.Not.Roy.Astron.Soc.* 527 (2023) 3

Cooling and Heating of NSs



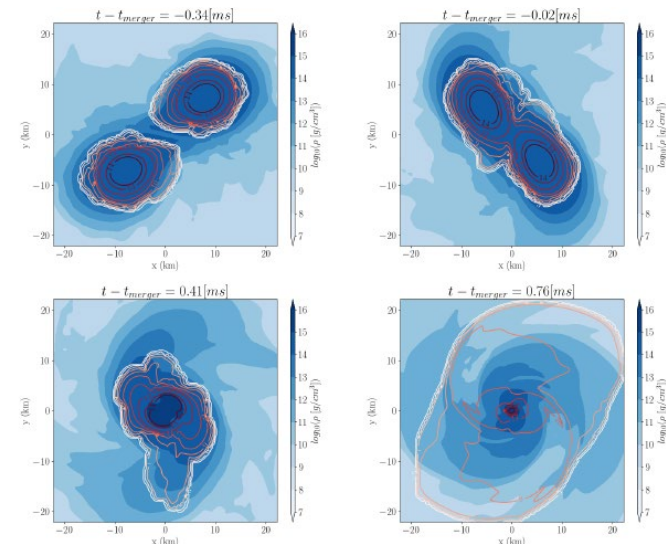
Ávila, E. Giangrandi, V. sagun, O. Ivanytskyi,
C. Providência, *Mon.Not.Roy.Astron.Soc.* **528** (2024) 4
Armen Sedrakian, *Phys.Rev.D* **93** (2016) 6, 065044
Chris Kouvaris, *Phys.Rev.D* **77** (2008) 023006

Black hole formation inside NSs



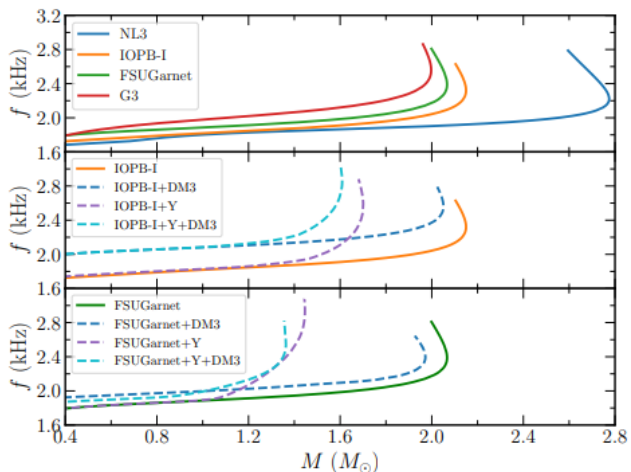
D. Singh, A. Gupta, E. Berti, S. Reddy, B. S. Sathyaprakash,
Phys.Rev.D **107** (2023) 8, 083037
N. F. Bell, A. Melatos, K. Petraki, *Phys.Rev.D* **87** (2013) 12, 123507

Numerical simulation of compact objects



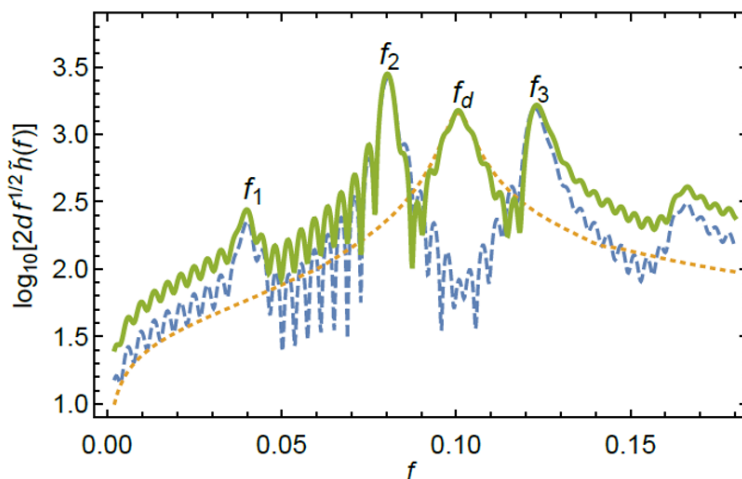
M. Emma, F. Schianchi, F. Pannarale, V. Sagun,
T. Dietrich, *Particles* **5** (2022) 3, 273-286
Andreas Bauswein, et al. *Phys.Rev.D* **107** (2023) 8, 083002
H. R. Rüter, V. Sagun, W. Tichy, T. Dietrich, *PRD* **108**, 124080

Different modes oscillation



Harish Chandra Das, et al. *Phys.Rev.D* **104** (2021) 12, 123006
S. Shirke, S. Ghosh, D. Chatterjee, L. Sagunski, J. Schaffner-Bielich,
JCAP **12** (2023) 008

Gravitational waves signals



John Ellis, et al. *Phys.Lett.B* **781** (2018) 607-610
Harish Chandra Das, et al. *Mon.Not.Roy.Astron.Soc.* **57** 4053 (2021)

Mass-Radius profile,
Tidal deformability and
moment of inertia

Pulse profile modeling

Z. Miao, Y. Zhu, Ang Li, F. Huang, *Ap.J.* **936** (2022) 1, 69
S. Shakeri, *D.R. K.*, *Phys. Rev. D* **109**, 043029 (2024)

Modeling of a DM admixed NS

Asymmetric DM

Single fluid DM admixed NS

An equation of state (EoS) by considering
DM-Baryonic matter (BM) interaction

M. ShahrbaF, D. Blaschke, S. Typel, et al. [Phys. Rev. D 105, 103005 \(2022\)](#)

Harish Chandra Das, [arXiv:2305.02065](#)

G. Panotopoulos and I. Lopes, [Phys.Rev.D 96 \(2017\) 8, 083004](#)

D. E. Alvarez-Castillo, M. Marczenko, [Phys.Polon.Supp. 15 \(2022\) 3, 28](#)

Direct impacts on nuclear models

Ask Harish, Mahboubeh

See Mahboubeh's poster

Two-fluid DM admixed NS

DM and BM interact only
through gravitational force

EoS for BM and EoS for DM

Our considered model

*Impacts on the observable features
resulting from nuclear models*

DM and Baryonic matter EoSs for two-fluid DM admixed NSs

VOLUME 57, NUMBER 20

PHYSICAL REVIEW LETTERS

17 NOVEMBER 1986

Boson Stars: Gravitational Equilibria of Self-Interacting Scalar Fields

Monica Colpi,^(a) Stuart L. Shapiro, and Ira Wasserman

DM: Self-interacting complex scalar field
Bosonic DM with **repulsive** self-interaction
 $V(\phi) = \frac{1}{4}\lambda|\phi|^4$ Leads to stellar mass Boson star

$$\mathcal{L} = \frac{1}{2}\partial_\mu\phi^*\partial^\mu\phi - \frac{m_\chi^2}{2}\phi^*\phi - \frac{\lambda}{4}(\phi^*\phi)^2.$$

Free parameters of the DM model
boson mass (m_χ), coupling constant (λ)

Strong coupling regime (Perfect fluid approximation)

In locally flat space-time by mean-field approximation

D.R. Karkevandi, S. Shakeri, V. Sagun, O. Ivanytskyi, PRD 105, 023001 (2022)

$$P = \frac{m_\chi^4}{9\lambda} \left(\sqrt{1 + \frac{3\lambda}{m_\chi^4}\rho} - 1 \right)^2.$$

DD2, a widely used and well-known nuclear matter equation of state (EoS)

Stiff EoS for which the tidal deformability is not consistent with the observational constraints

S. Typel and H.H. Wolter, Nucl.Phys.A 656 (1999) 331-364

S. Typel, Phys.Rev.C 71 (2005) 064301

Two-fluid DM admixed NS

BM and DM fluids interact only gravitationally



$$G_{\mu\nu} = R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = 8\pi(T_{DM}^{\mu\nu} + T_{BM}^{\mu\nu})$$

Energy-momentum tensors are conserved separately

Two-fluid Tolman-Oppenheimer-Volkof equation

F. Sandin & P. Ciarcelluti. Astropart.Phys.32:278-284,2009.

P. Ciarcelluti & F. Sandin. Phys.Lett. B695:19-21,2011.

$$\frac{dp_B}{dr} = - (p_B + \varepsilon_B) \frac{m + 4\pi r^3 p}{r(r - 2m)}$$

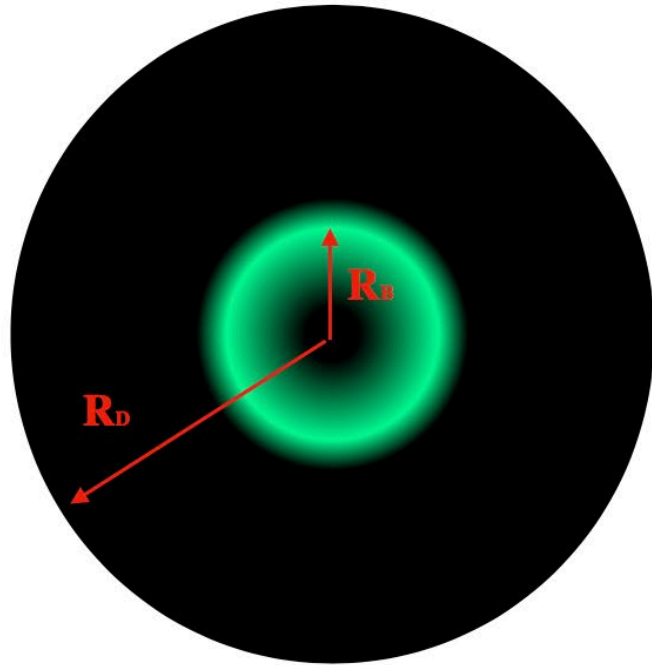
$$\frac{dp_D}{dr} = - (p_D + \varepsilon_D) \frac{m + 4\pi r^3 p}{r(r - 2m)}$$

$$m(r) = \underbrace{\int_0^r 4\pi r^2 \varepsilon_B}_{m_B(r)} + \underbrace{\int_0^r 4\pi r^2 \varepsilon_D}_{m_D(r)}$$

$$p(r) = p_B(r) + p_D(r)$$

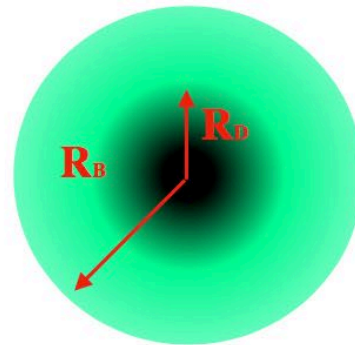
Three Possible DM distributions within NSs

DM halo



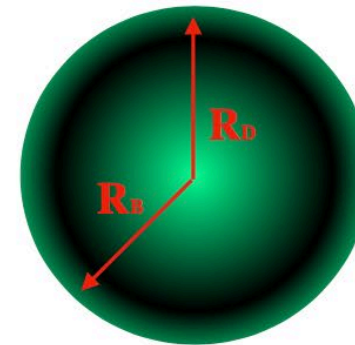
$$R_D > R_B$$

DM Core



$$R_B > R_D$$

DM distributed in entire NS



$$R_B \approx R_D$$

Core of a DM admixed NS is composed of both of the fluids

Green : BM
Black : DM

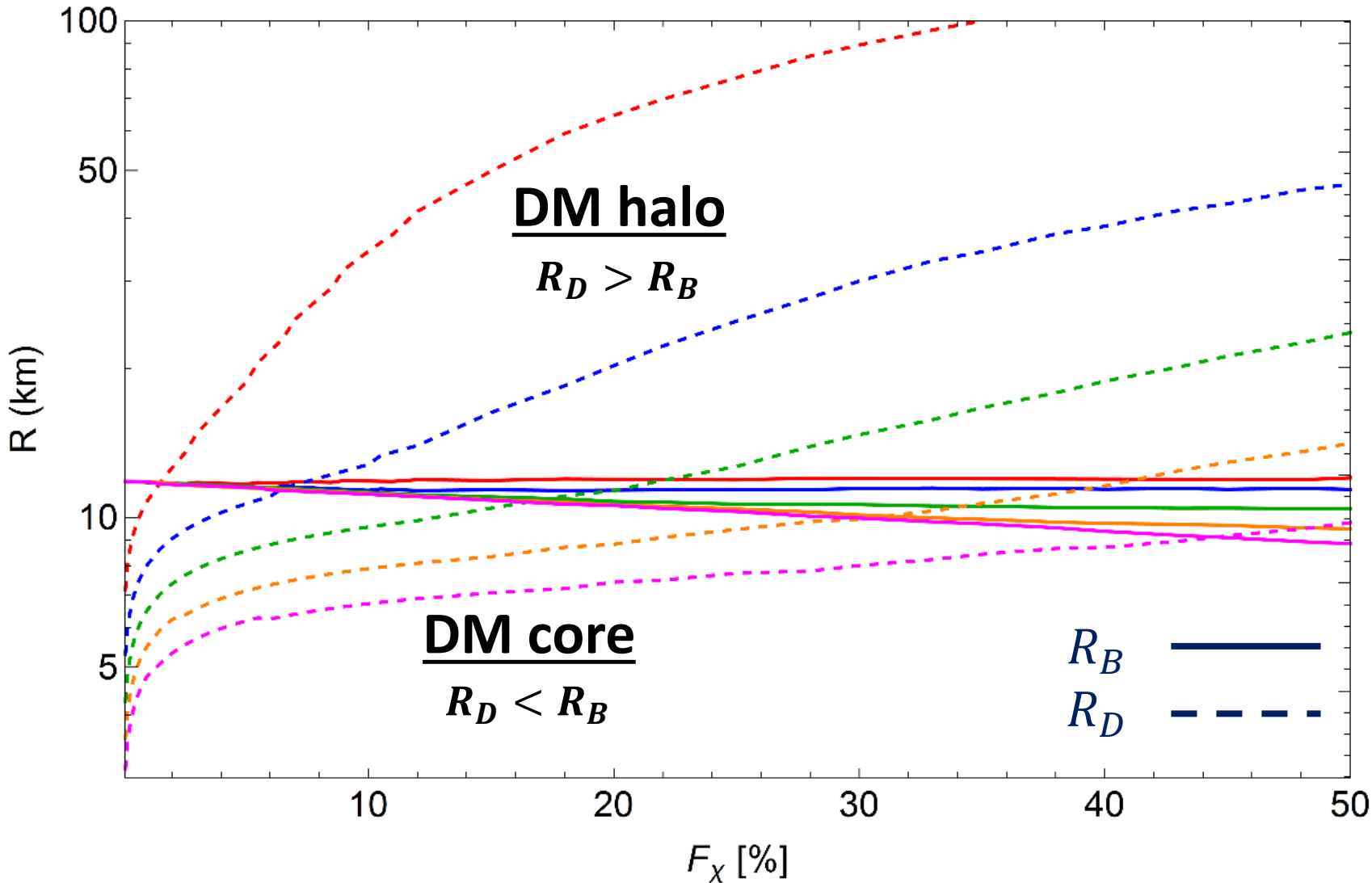
$$M_T = M_B(R_B) + M_D(R_D)$$

$$F_\chi = \frac{M_D(R_D)}{M_T}, \text{ DM Fraction}$$

R_B is the visible radius

Variation of BM radius (*solid lines*) and DM radius (*dashed lines*) in DM admixed NSs

By increasing the DM fraction, a transition can be seen from DM core to DM halo formation

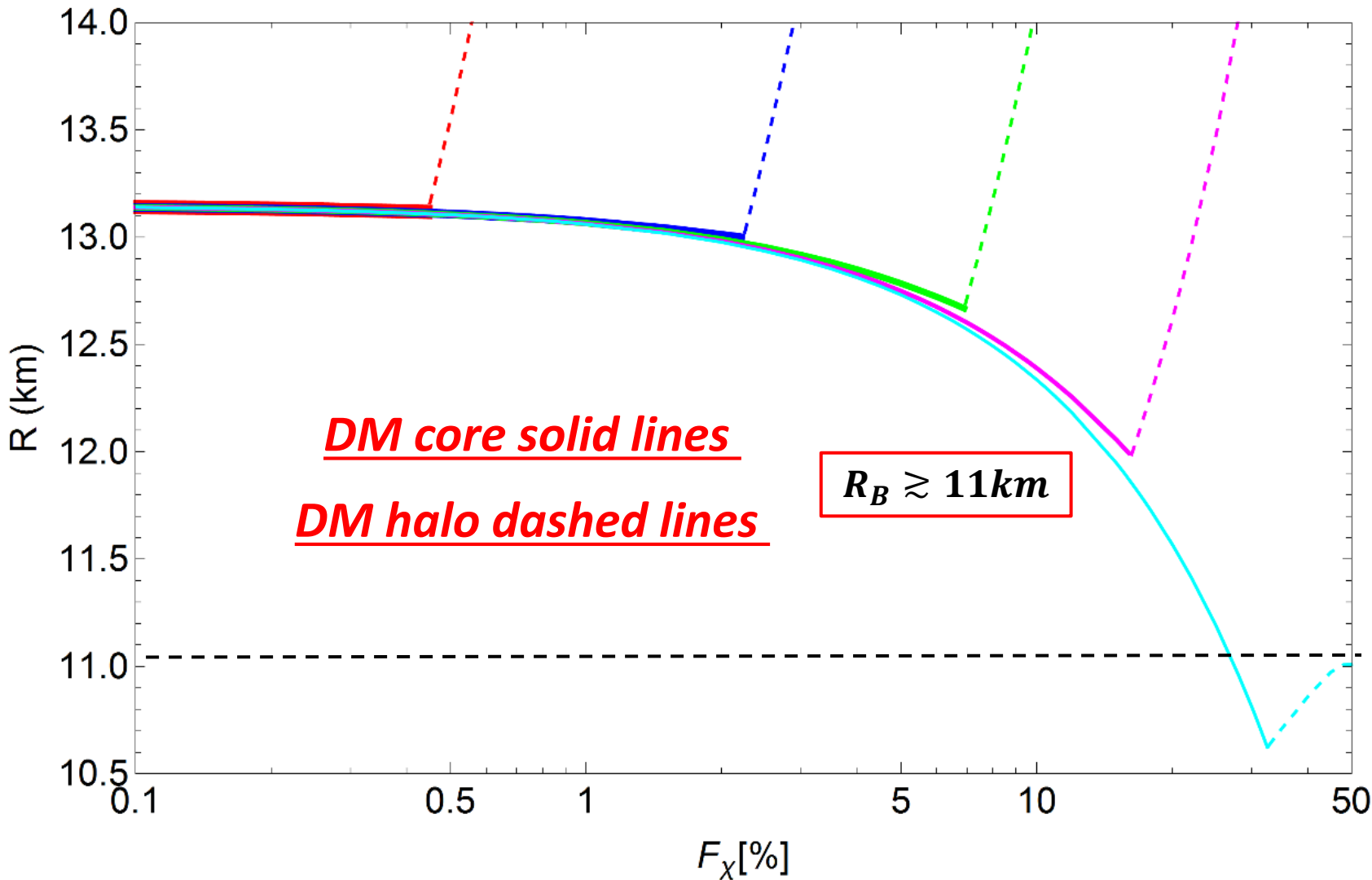


- $\lambda = \pi$
- 100 MeV
 - 150 MeV
 - 200 MeV
 - 250 MeV
 - 300 MeV

Even in low DM fractions a DM halo can be formed for light bosons. However, Heavy bosons lead mainly to DM core formation.

Variation of outermost radius of DM admixed NSs

DM core: R_B (BM radius) is the outermost radius **DM halo:** R_D (DM radius) is the outermost radius



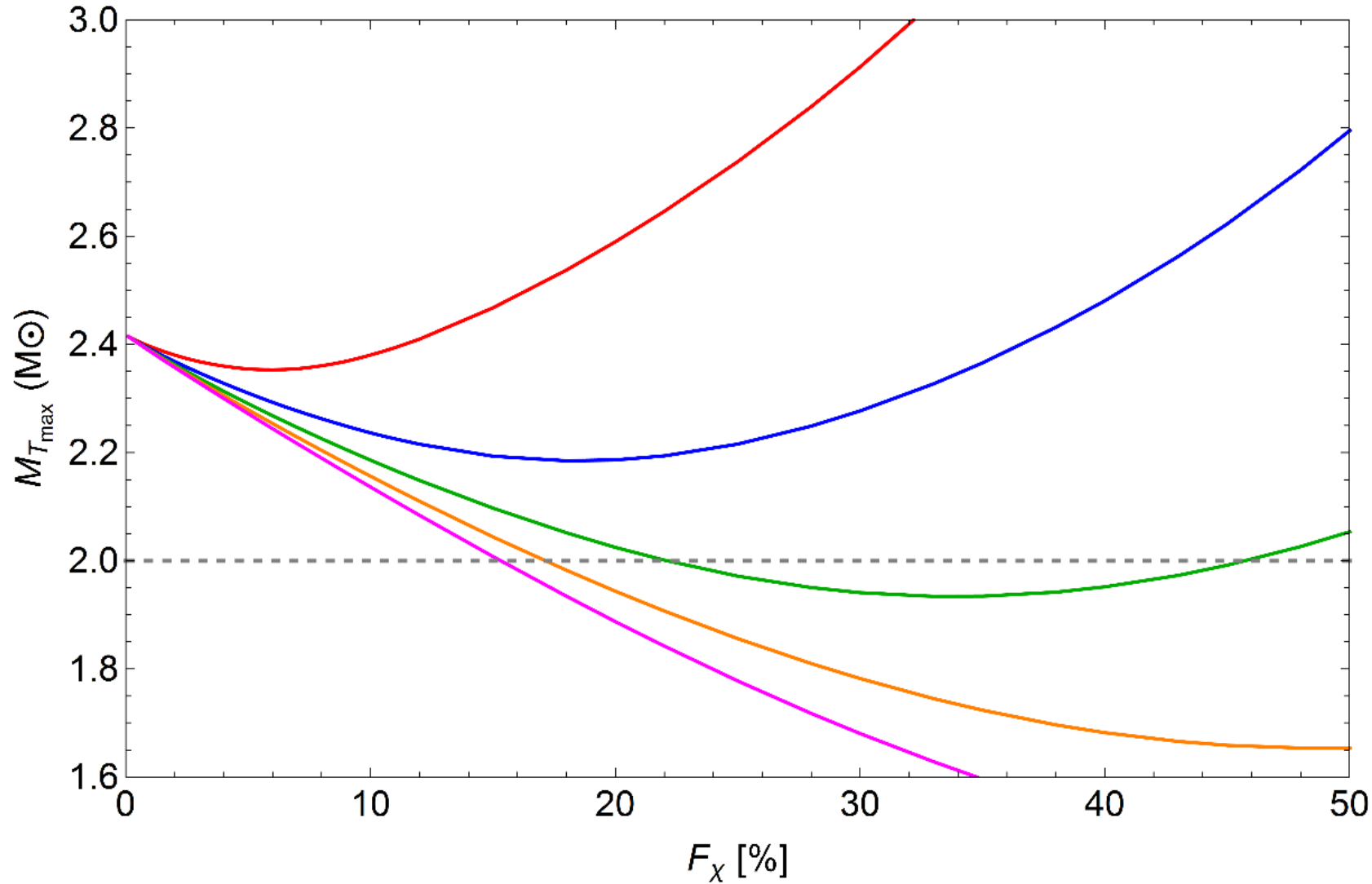
- $\lambda = \pi$
- 100 MeV
 - 150 MeV
 - 200 MeV
 - 250 MeV
 - 300 MeV

DM core formation
, solid lines,
decreases
the outermost radius.

DM halo formation
, dashed lines,
increases
the outermost radius.

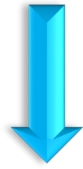
Variation of the total maximum mass of DM admixed NSs

For light bosons the maximum mass constraint is satisfied for the whole range of DM fractions while for heavy bosons the $2M_{\odot}$ limit will be violated for some DM fractions.



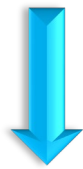
- $\lambda = \pi$
- 100 MeV
 - 150 MeV
 - 200 MeV
 - 250 MeV
 - 300 MeV

DM core



Decrease
in the
maximum mass

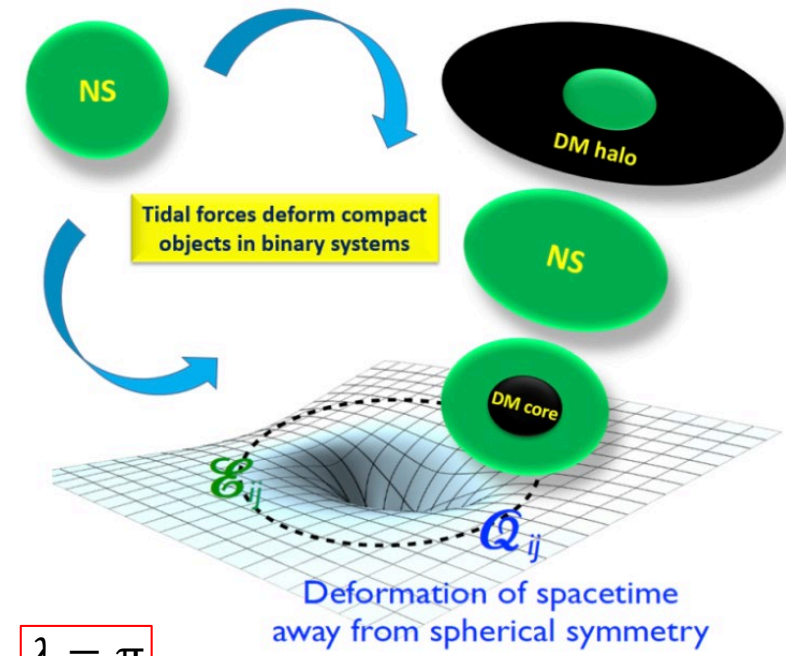
DM halo



Increase
in the
maximum mass

Variation of the tidal deformability of DM admixed NSs

Heavy bosons mainly reside as a core inside NS,
while light bosons form a large halo around the NS.



Tidal forces deform compact objects in binary systems

Deformation of spacetime away from spherical symmetry

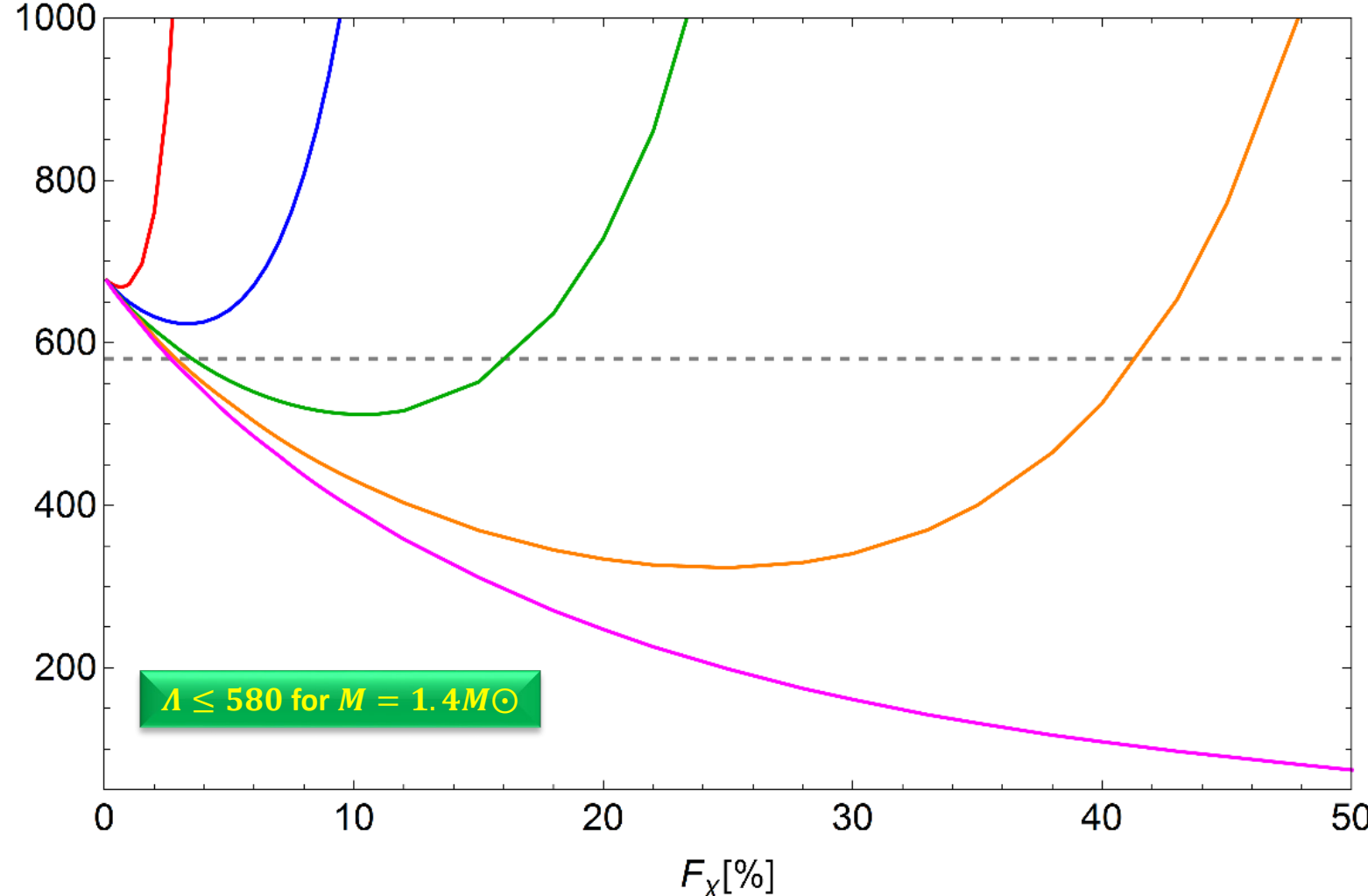
$$\lambda = \pi$$

- 100 MeV
- 150 MeV
- 200 MeV
- 250 MeV
- 300 MeV

$$\Lambda = \frac{\lambda_t}{M^5} = \frac{2}{3} k_2 \left(\frac{R}{M}\right)^5$$

DM core decreases tidal deformability

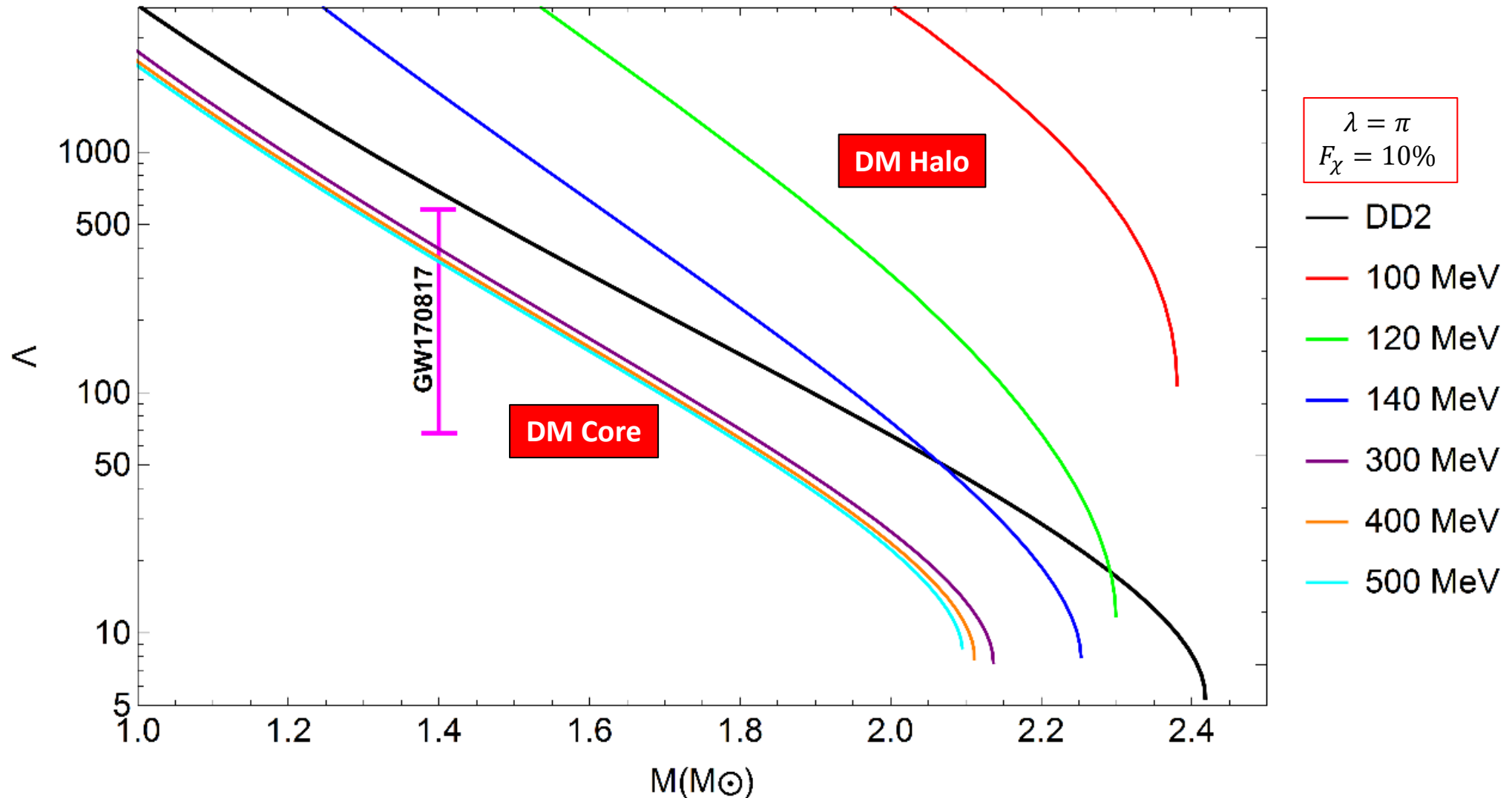
DM halo increases tidal deformability



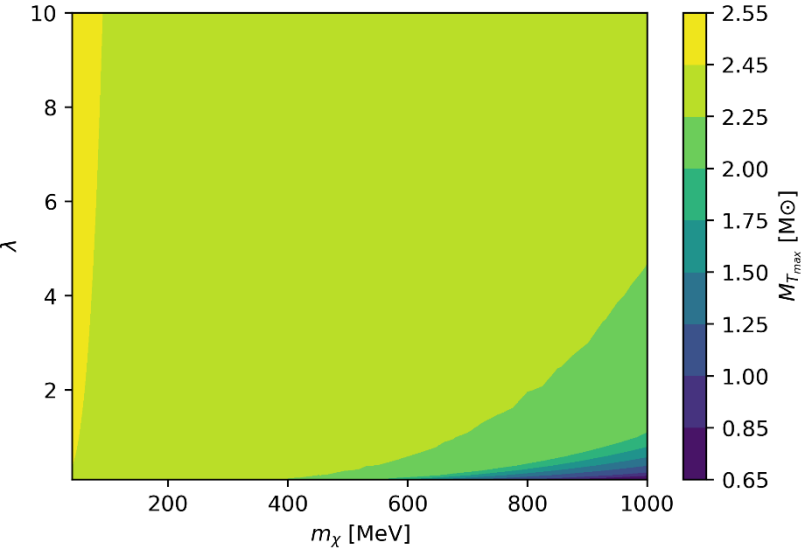
$\Lambda \leq 580$ for $M = 1.4 M_{\odot}$

Tidal deformability of DD2 EoS will be **modified** due to the presence of bosonic DM in the core of NSs.

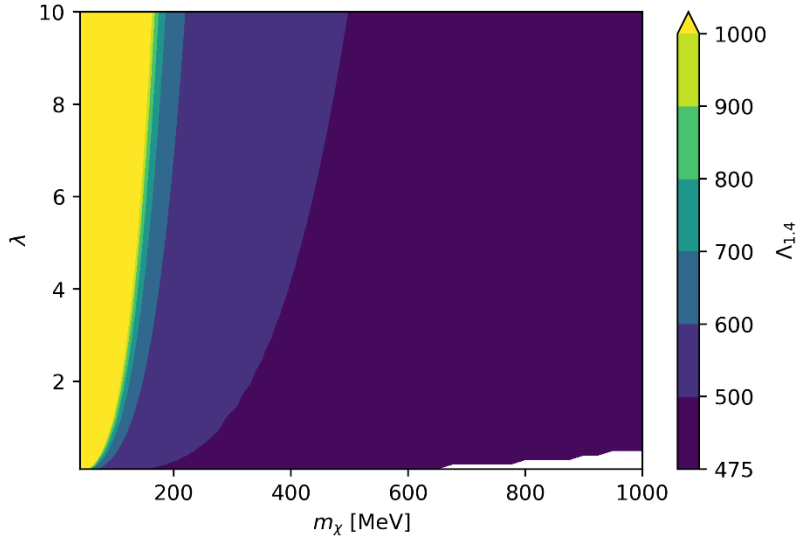
Black solid lines corresponds to DD2 EoS without DM



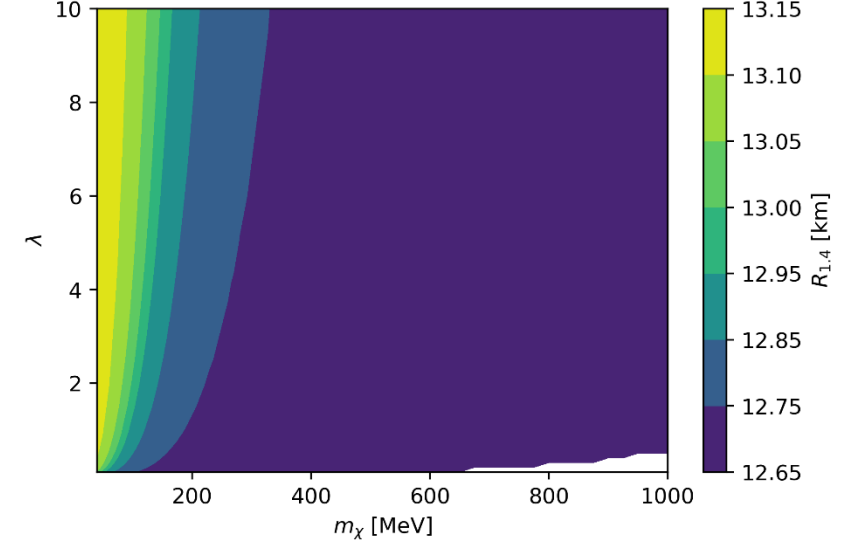
Maximum mass



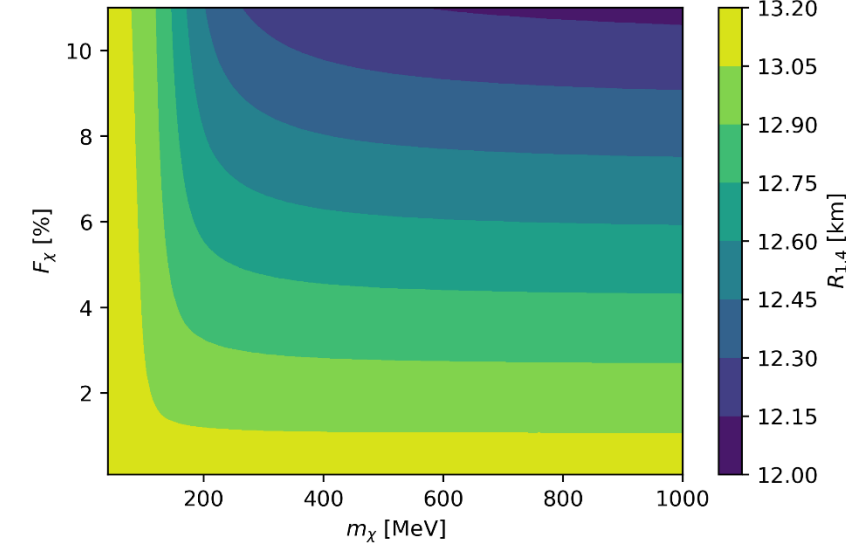
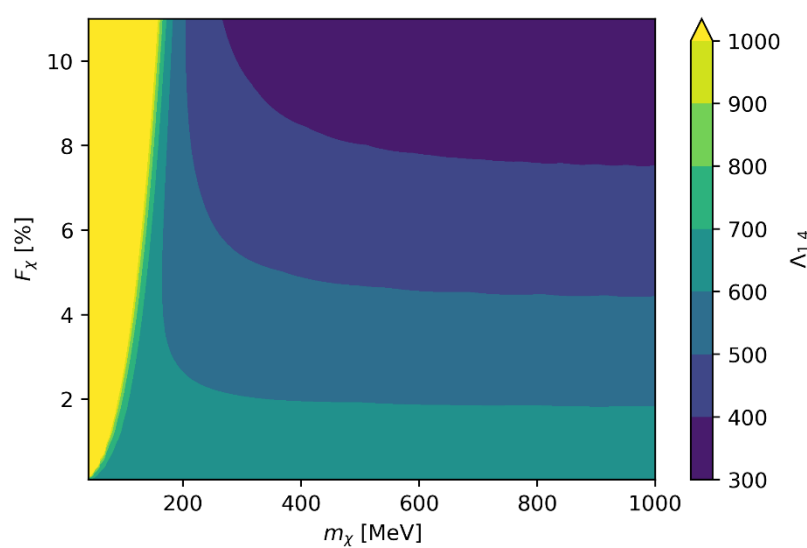
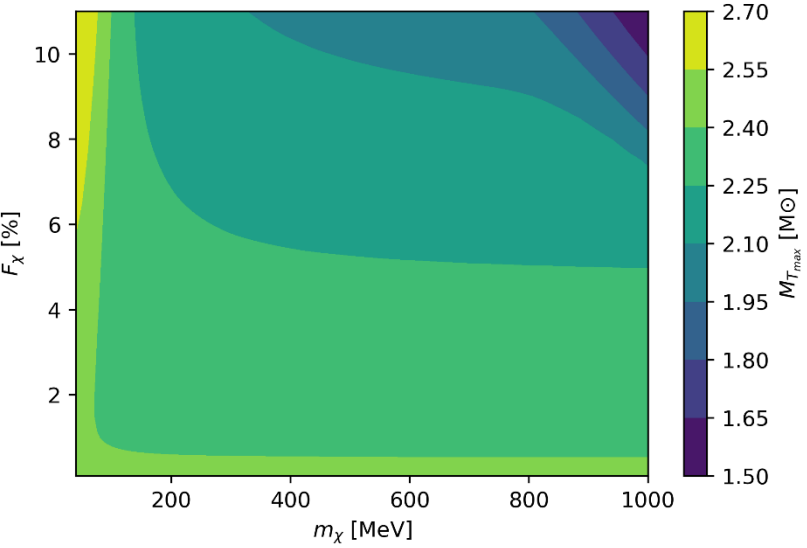
Tidal deformability



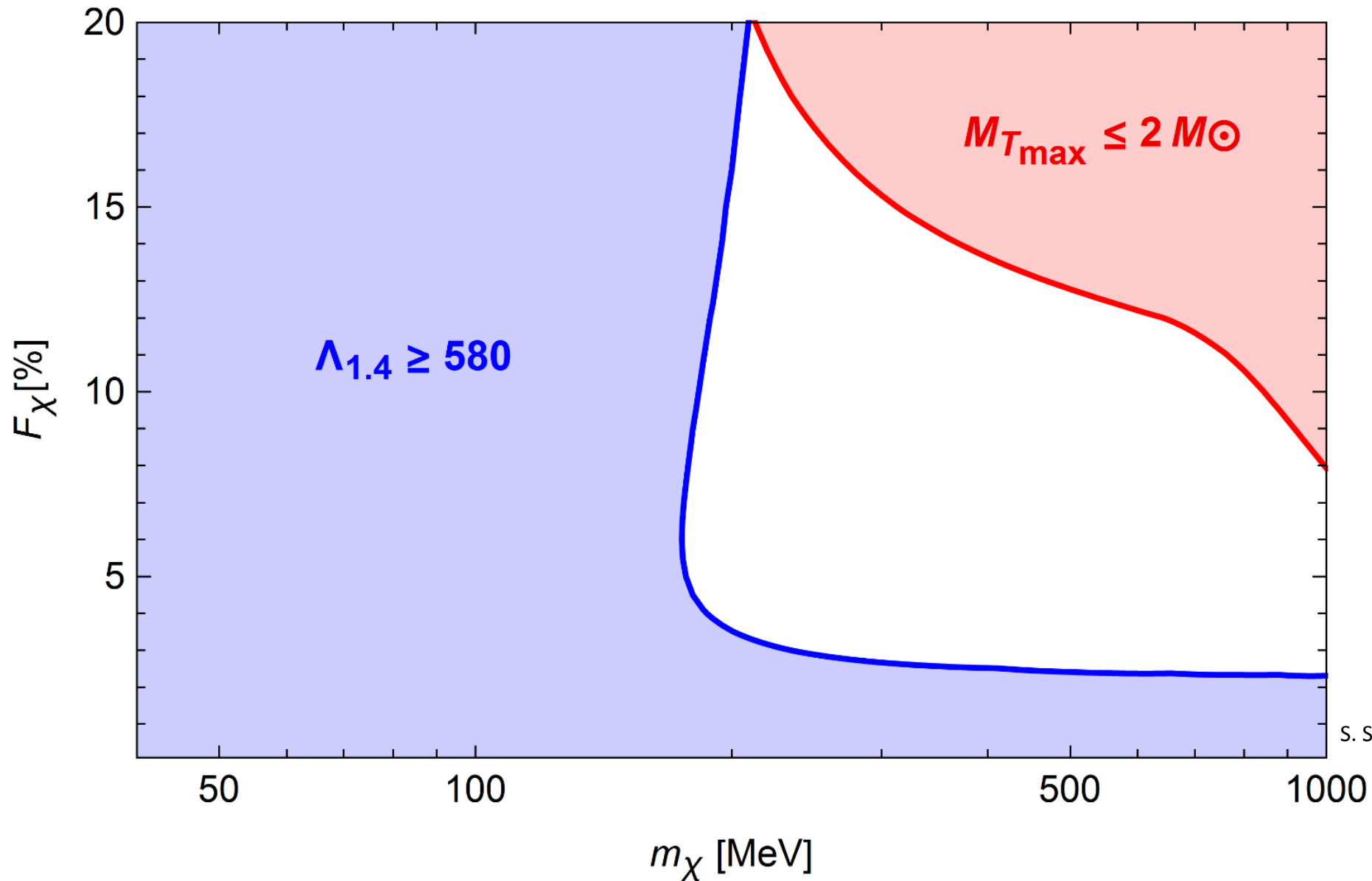
Visible radius (R_B)



Scan over the parameter space of the bosonic DM model, **Top:** $\lambda - m_\chi$ and **Bottom:** $F_\chi - m_\chi$



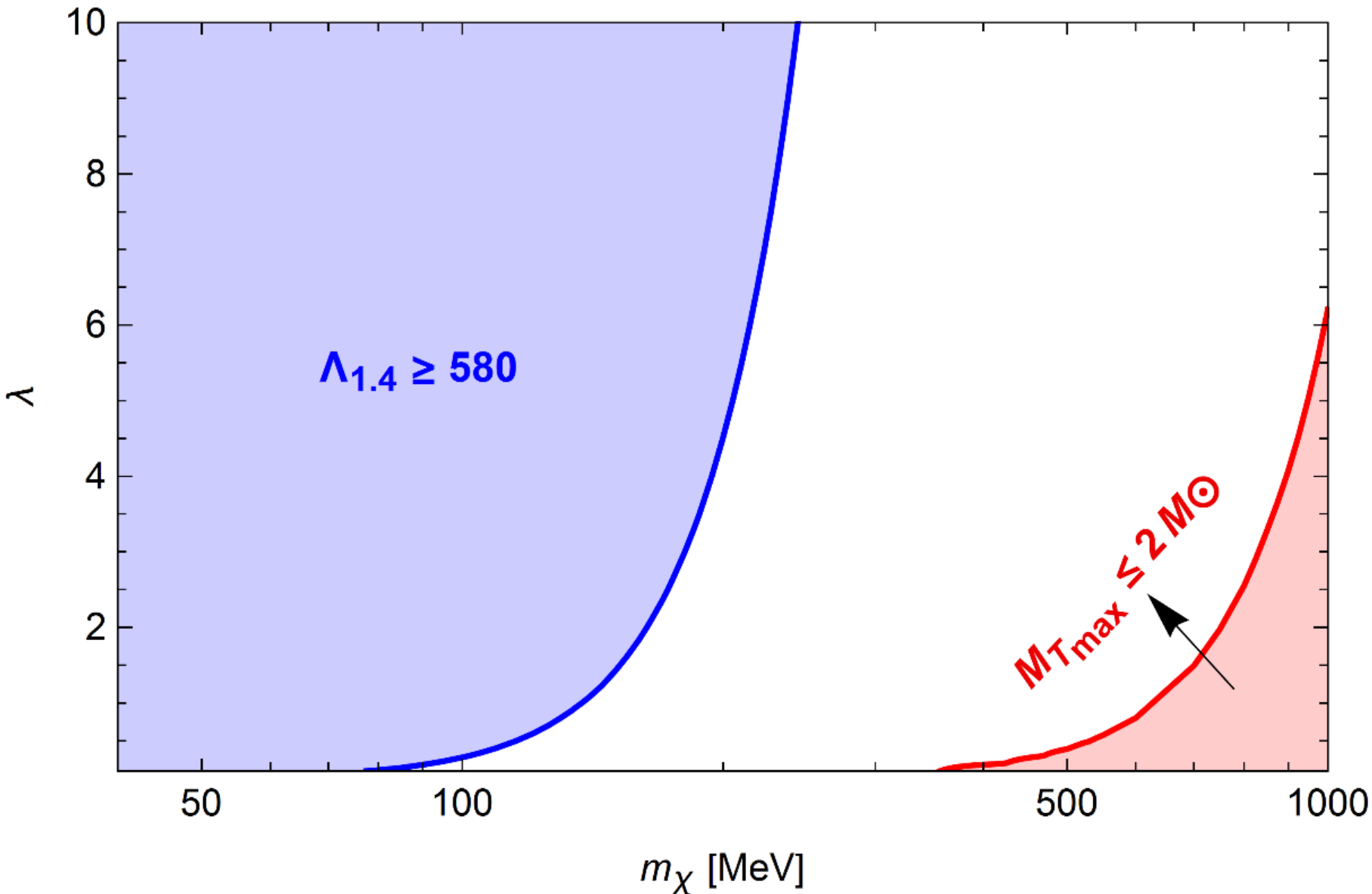
Scan over the $F_\chi - m_\chi$ parameter space of DM admixed NSs for $\lambda = \pi$



DM fraction is limited between 2% and 20% for which all the considered observable features are satisfied.

Scan over the $\lambda - m_\chi$ parameter space for $F_\chi = 10\%$

Light bosonic DM particles ($m_\chi \lesssim 200 \text{ MeV}$) are excluded by tidal deformability parameter for the whole range of λ .

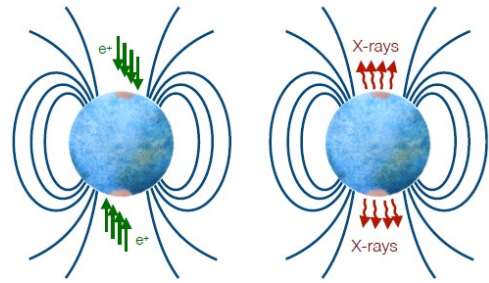


For heavy bosons ($m_\chi \gtrsim 500 \text{ MeV}$) and low coupling constants, the maximum mass is not consistent with observational constraint.

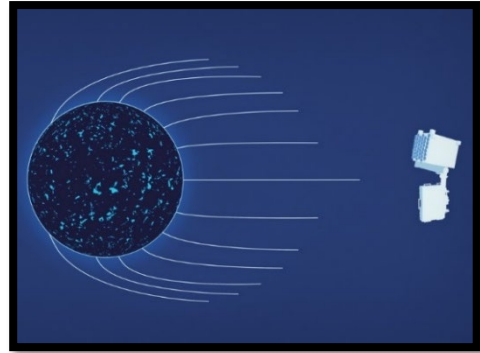
The allowed region shows where all the given observable properties of the NS, induced from DD2 EoS, is in agreement with the astrophysical bounds.

Pulse profile modeling as a novel probe for DM halo formation around NSs

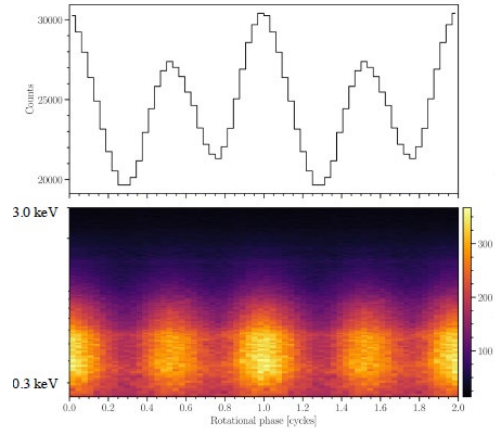
X-ray hot spots



NICER and light bending due to the curved of space time



X-ray Pulse profile



DM halo



Visible surface is increasing with the compactness

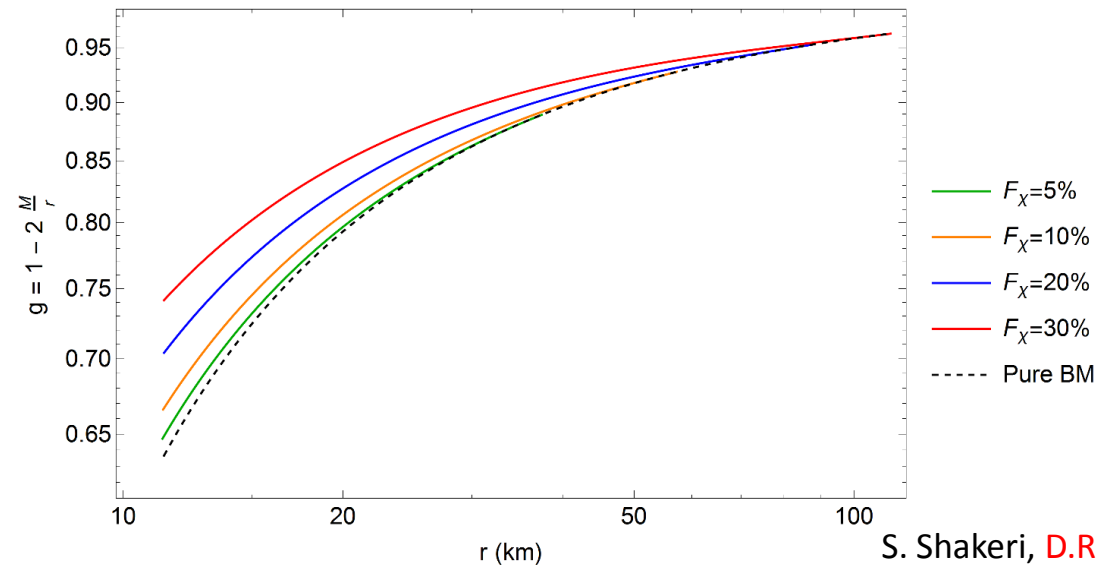
$C = \frac{M_T(R_B)}{R_B}$

C_1 , C_2 , C_3

$M_T = 1.4M_\odot$ For all cases

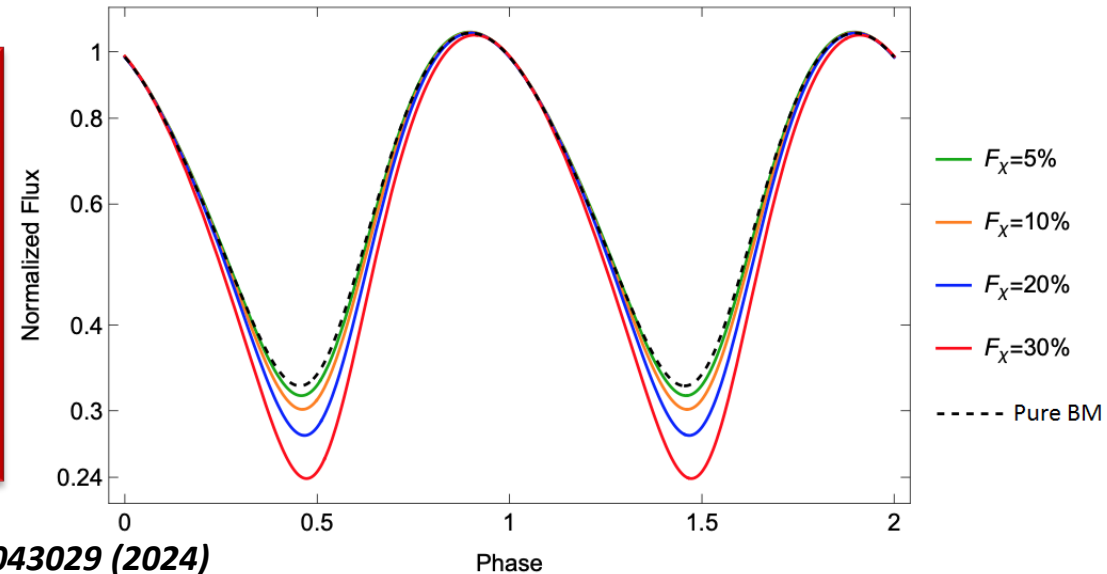
$C_1 > C_2 > C_3$

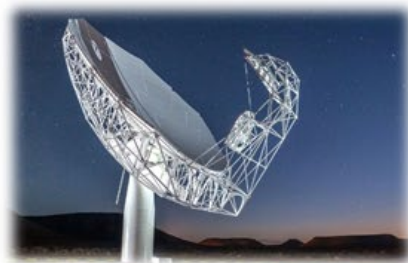
Variation of metric function due to the DM halo



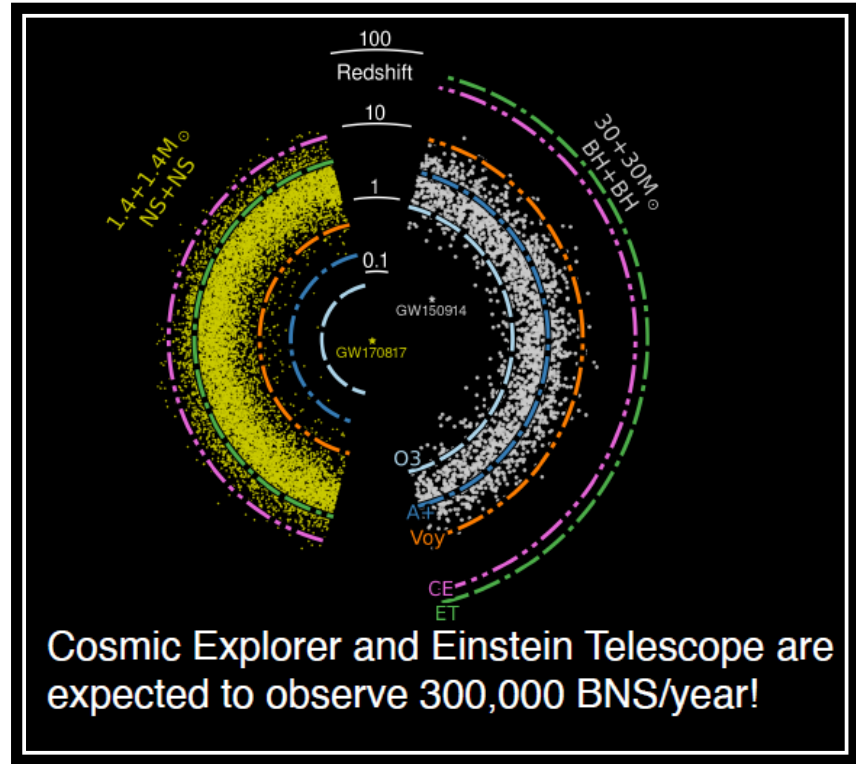
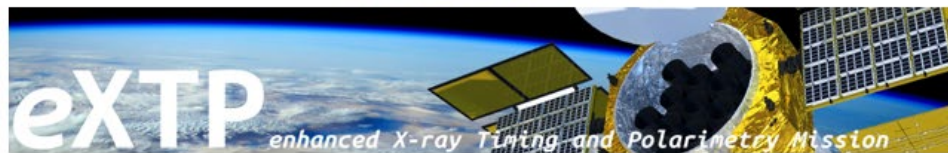
The deviation of the minimums of the fluxes compare to the pure NS is a remarkable signature of the DM halo.

X-ray pulse profile in the presence of DM halo





MeerKAT



[arXiv:2109.09882](https://arxiv.org/abs/2109.09882)

Exotic measurements

THE ASTROPHYSICAL JOURNAL LETTERS, 896:L44 (20pp), 2020 June 20
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<https://doi.org/10.3847/2041-8213/ab960f>
OPEN ACCESS
GW190814: Gravitational Waves from the Coalescence of a 23 Solar Mass Black Hole with a 2.6 Solar Mass Compact Object

Article | Published: 24 October 2022
A strangely light neutron star within a supernova remnant $M = 0.77M_{\odot}, R = 10.4\text{km}$
 Victor Doroshenko, Valery Suleimanov, Gerd Pühlhofer & Andrea Santangelo
Nature Astronomy 6, 1444–1451 (2022) | [Cite this article](#)

PRD. 105, 063005
 APJ. 922 (2021) 242
 PRD. 104, 063028 (2021)
 Astrophys. J. 958, 49 (2023)
 arXiv:2307.12748



Isfahan miasto Polskich dzieci

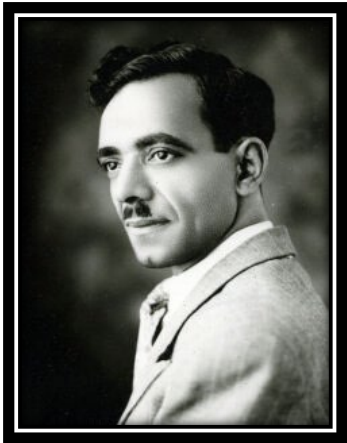
Isfahan - the City of Polish Children

1942-1945

2590 Polish children mainly below 9

Khajoo bridge, Isfahan-Iran



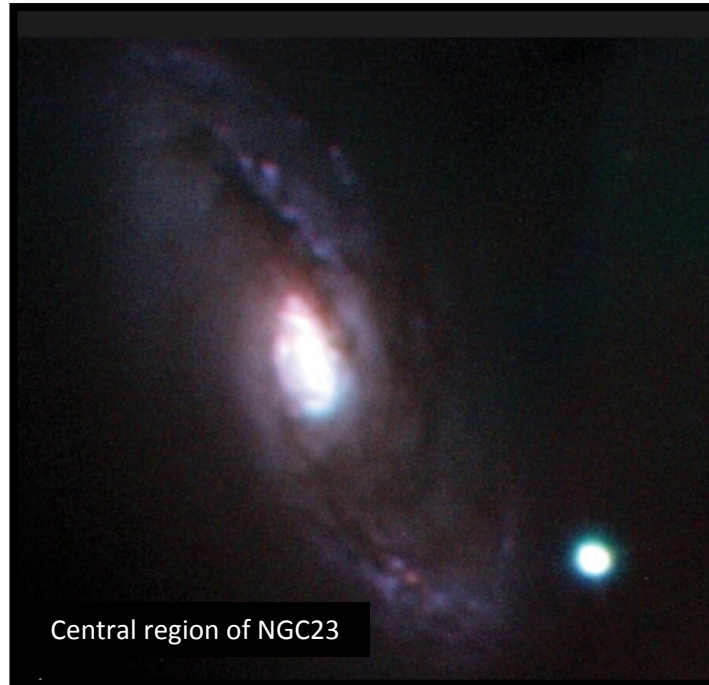
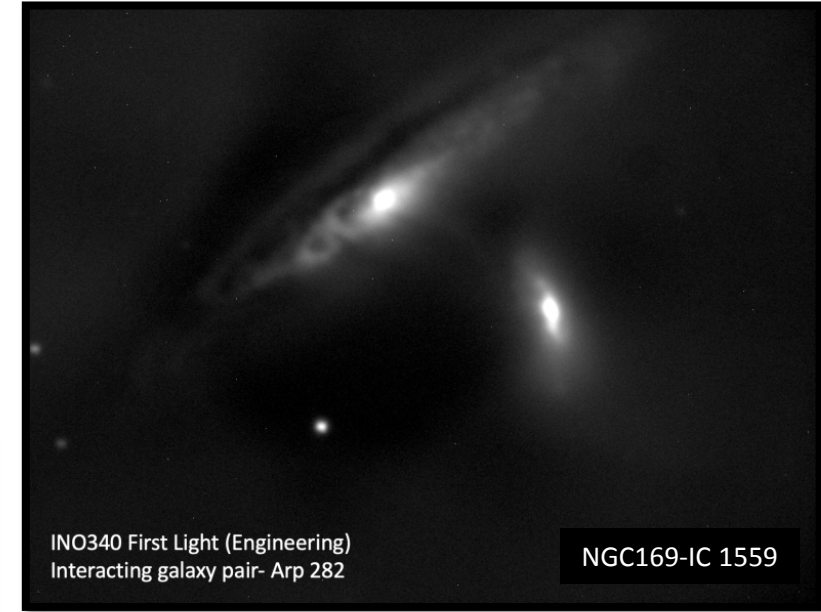


The Photographer
Abolqasem Jala



Optical observations of black-widow pulsars

Iranian National Observatory (INO), the largest home-grown scientific facility project, has recorded the first light image of its 3.4m optical telescope on October 2022.



'The door is open': Iranian astronomers seek collaborations for their new, world-class telescope.

Iranian National Observatory (INO)
3.4 meter optical telescope
3600m above the sea level

Science
AAAS



Seventeenth Marcel Grossmann Meeting

7–12 Jul 2024

The 'Gabriele d'Annunzio' University, ICRANet and Aurum

Europe/Rome timezone

Bosonic dark matter and/in neutron stars

DM5 

Chairpersons: Soroush Shakeri and Davood Rafiei Karkevandi



Among different Dark Matter (DM) candidates the scalar or pseudoscalar bosons are of great interest from various aspects in astrophysics and cosmology. Generally, bosonic DM could form gravitationally stable configurations or be substantially accumulated in compact objects such as Neutron Stars (NSs) through different scenarios. In this regard, the advent of multi-messenger observations via gravitational and electromagnetic waves provide a unique opportunity to probe the existence of dense astrophysical objects made entirely or fractionally by bosonic DM.

In this section, we will focus on both theoretical and observational aspects of boson stars, fermion-boson stars, DM admixed NSs and other exotic type of compact stars. Recently, the GW detections by LIGO-Virgo-KAGRA and X-ray observations by NICER telescopes have opened a new window towards understanding the structure of compact objects and may shed light on the nature of DM through exploring exotic results.