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- Computer codes `MB.m` and `MBresolve.m`

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- Simple one-loop examples
- General prescriptions for resolving singularities in ϵ in multiple Mellin-Barnes integrals. Two strategies
- Computer codes `MB.m` and `MBresolve.m`
- Various examples

Mellin transformation, Mellin integrals as a tool for Feynman integrals:

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Systematic evaluation of dimensionally regularized Feynman integrals (in particular, systematic resolution of the singularities in ϵ)
[V.A. Smirnov'99, J.B. Tausk'99]

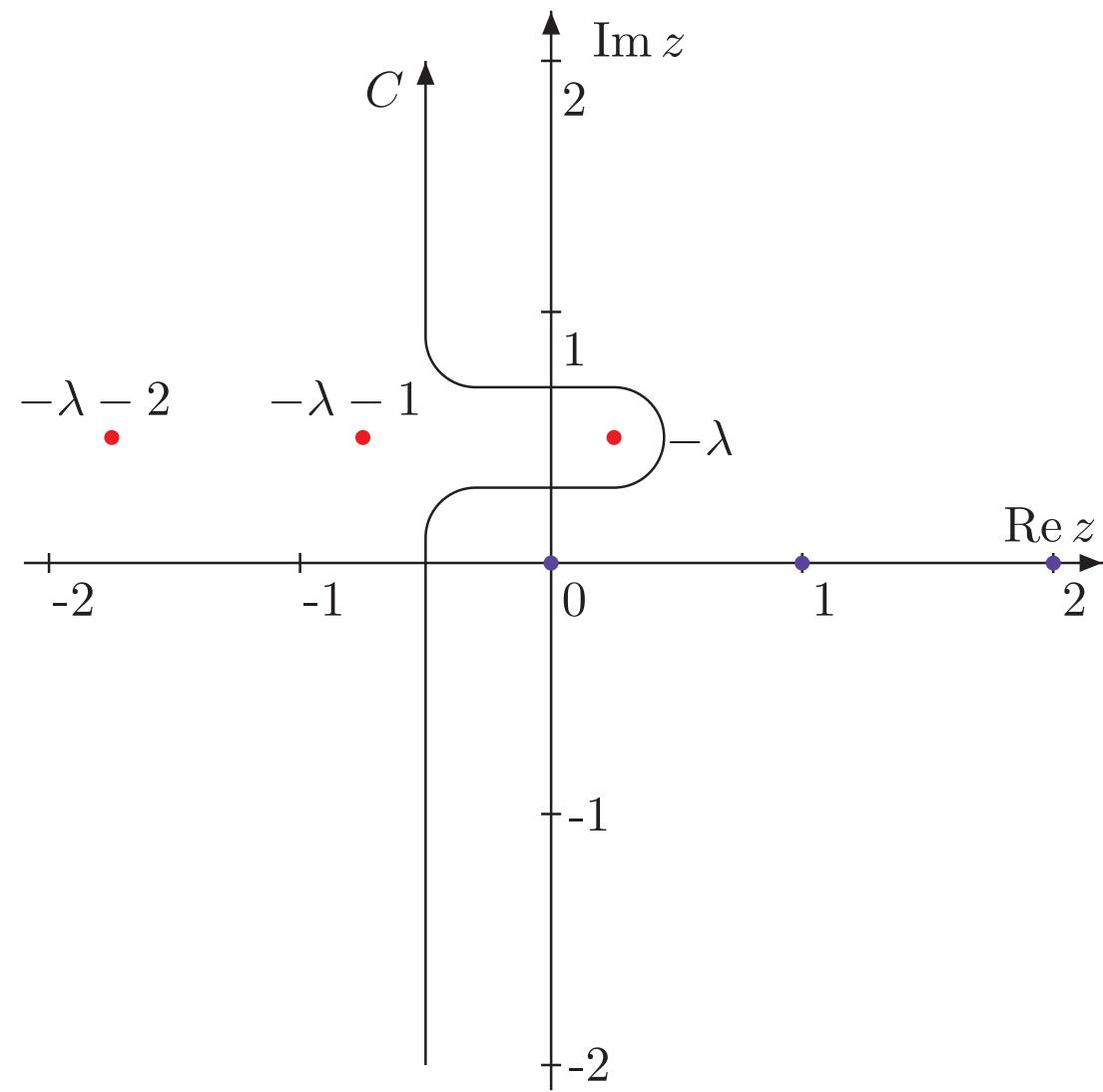
The basic formula:

$$\frac{1}{(X+Y)^\lambda} = \frac{1}{\Gamma(\lambda)} \frac{1}{2\pi i} \int_{-i\infty}^{+i\infty} dz \frac{Y^z}{X^{\lambda+z}} \Gamma(\lambda + z) \Gamma(-z).$$

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The poles with a $\Gamma(\dots + z)$ dependence are to the left of the contour and the poles with a $\Gamma(\dots - z)$ dependence are to the right



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- Evaluate expanded MB integrals

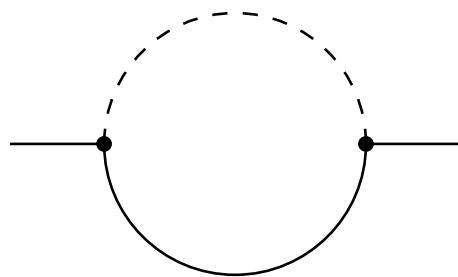
The simplest possibility:

$$\frac{1}{(m^2 - k^2)^\lambda} = \frac{1}{\Gamma(\lambda)} \frac{1}{2\pi i} \int_{-i\infty}^{+i\infty} dz \frac{(m^2)^z}{(-k^2)^{\lambda+z}} \Gamma(\lambda + z) \Gamma(-z)$$

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Example 1



$$F_\Gamma(q^2, m^2; a_1, a_2, d) = \int \frac{d^d k}{(m^2 - k^2)^{a_1} ((q - k)^2)^{a_2}}$$

$$F_\Gamma = \frac{1}{\Gamma(a_1)} \frac{1}{2\pi i} \int_{-\text{i}\infty}^{+\text{i}\infty} dz (m^2)^z \Gamma(a_1 + z) \Gamma(-z) \\ \times \int \frac{\mathbf{d}^d k}{(-k^2)^{a_1+z}(-(q-k)^2)^{a_2}}$$

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$$\int \frac{d^d k}{(-k^2)^{a_1+z} [-(q-k)^2]^{a_2}} = i\pi^{d/2} \frac{G(a_1+z, a_2)}{(-q^2)^{a_1+a_2+\epsilon-2+z}} ,$$

$$G(a_1, a_2) = \frac{\Gamma(a_1 + a_2 + \epsilon - 2)\Gamma(2 - \epsilon - a_1)\Gamma(2 - \epsilon - a_2)}{\Gamma(a_1)\Gamma(a_2)\Gamma(4 - a_1 - a_2 - 2\epsilon)}$$

$$\begin{aligned}
F_\Gamma(q^2, m^2; a_1, a_2, d) &= \frac{i\pi^{d/2} \Gamma(2 - \epsilon - a_2)}{\Gamma(a_1) \Gamma(a_2) (-q^2)^{a_1 + a_2 + \epsilon - 2}} \\
&\times \frac{1}{2\pi i} \int_{-i\infty}^{+i\infty} dz \left(\frac{m^2}{-q^2} \right)^z \Gamma(a_1 + a_2 + \epsilon - 2 + z) \\
&\times \frac{\Gamma(2 - \epsilon - a_1 - z) \Gamma(-z)}{\Gamma(4 - 2\epsilon - a_1 - a_2 - z)}
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\end{aligned}$$

Unambiguous prescriptions for contours:
the poles with a $\Gamma(\dots + z)$ dependence are to the left and
the poles with a $\Gamma(\dots - z)$ dependence are to the right of a
contour

Strategy A

[V.A. Smirnov'99]

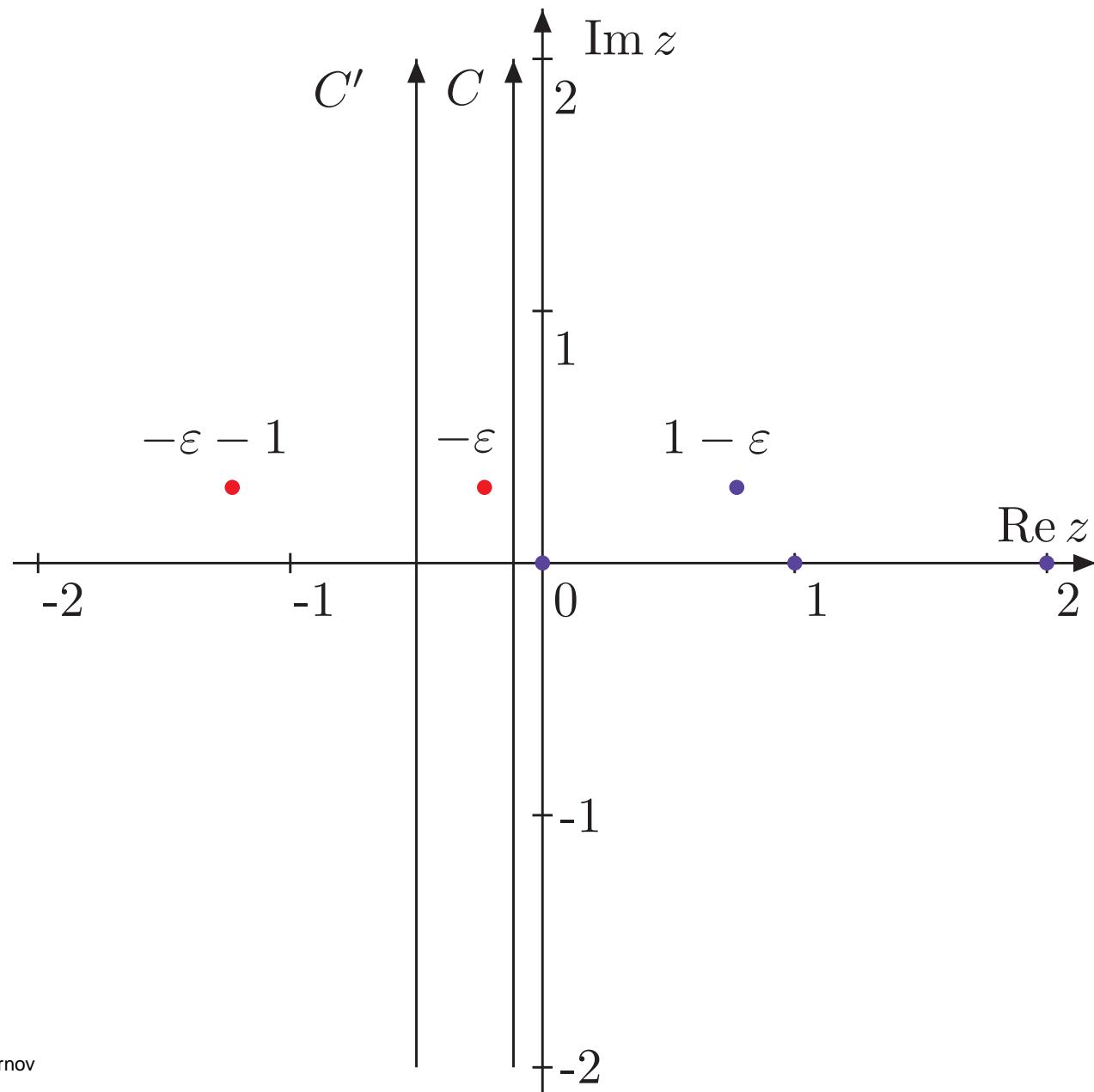
$$\begin{aligned} F_\Gamma(q^2, m^2; 1, 1, d) &= \frac{i\pi^{d/2}\Gamma(1-\epsilon)}{(-q^2)^\epsilon} \\ &\times \frac{1}{2\pi i} \int_C dz \left(\frac{m^2}{-q^2} \right)^z \frac{\Gamma(\epsilon+z)\Gamma(-z)\Gamma(1-\epsilon-z)}{\Gamma(2-2\epsilon-z)} \end{aligned}$$

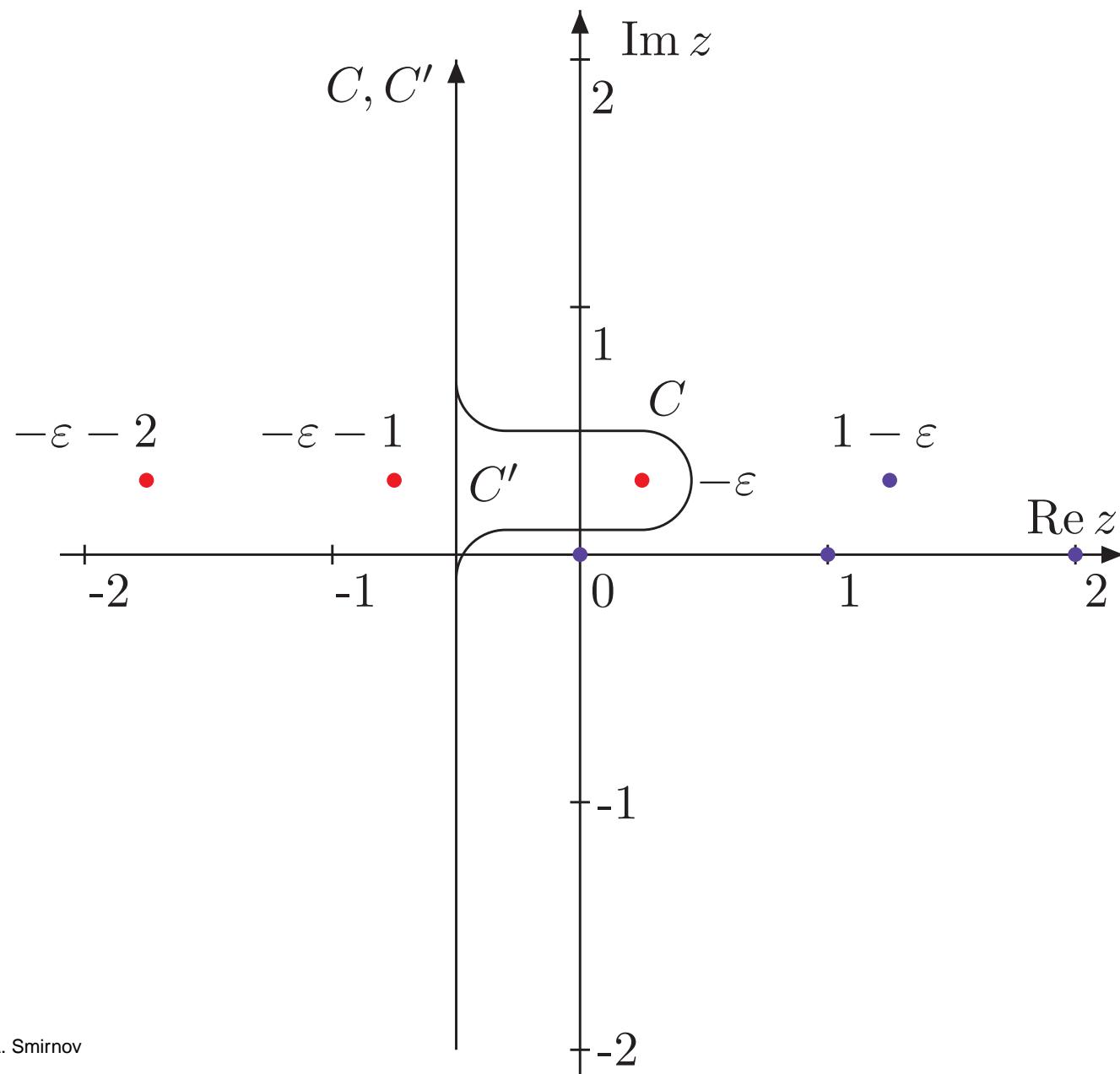
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$\Gamma(\epsilon+z)\Gamma(-z) \rightarrow$ a singularity in ϵ





Take a residue at $z = -\epsilon$:

$$i\pi^2 \frac{\Gamma(\epsilon)}{(m^2)^\epsilon(1-\epsilon)}$$

and shift the contour:

$$\frac{i\pi^{d/2}\Gamma(1-\epsilon)}{(-q^2)^\epsilon} \frac{1}{2\pi i} \int_{C'} dz \left(\frac{m^2}{-q^2} \right)^z \frac{\Gamma(\epsilon+z)\Gamma(-z)\Gamma(1-\epsilon-z)}{\Gamma(2-2\epsilon-z)}$$

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NB:

$$\Gamma(\epsilon + z)\Gamma(1 - \epsilon - z) = -\Gamma(1 + \epsilon + z)\Gamma(-\epsilon - z)$$

The integral can be expanded in ϵ , e.g., the value at $\epsilon = 0$ is

$$\begin{aligned}\frac{1}{2\pi i} \int_{C'} f(z) dz &= - \sum_{n=0} \operatorname{res}_{z=n} f(z) \\ &= + \sum_{n=1} \operatorname{res}_{z=n} f(z) \\ &= 1 - \left(1 - \frac{m^2}{q^2}\right) \ln \left(1 - \frac{q^2}{m^2}\right)\end{aligned}$$

where

$$f(z) = \left(\frac{m^2}{-q^2}\right)^z \frac{\Gamma(z)\Gamma(-z)\Gamma(1-z)}{\Gamma(2-z)} = \left(\frac{m^2}{-q^2}\right)^z \frac{\Gamma(z)\Gamma(-z)}{(1-z)}$$

Strategy A in a modified form

[A.V. Smirnov and V.A. Smirnov'09]

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Choose a straight contour C_0 for which the gamma functions in the numerator of the integrand are spoiled at $\epsilon = 0$ in a minimal way, i.e. the initial rules for choosing a contour are changed in a minimal way

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Let us choose C_0 with $\operatorname{Re} z = -1/4$.

Then $\Gamma(\epsilon + z)$ which transforms into $\Gamma(z)$ at $\epsilon = 0$ is spoiled.

$$\Gamma(\epsilon + z) \rightarrow \Gamma^{(1)}(\epsilon + z)$$

$\Gamma^{(1)}(\epsilon + z)$ means that the rule $\operatorname{Re}(\epsilon + z) > 0$ when crossing the real axis is changed to $-1 < \operatorname{Re}(\epsilon + z) < 0$

We do not need to spoil it more, e.g., by

$\Gamma(\epsilon + z) \rightarrow \Gamma^{(2)}(\epsilon + z)$ with the rule
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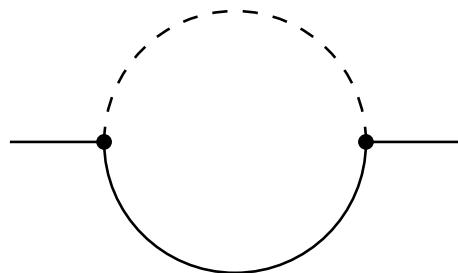
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Then

$$\begin{aligned} \frac{1}{2\pi i} \int_C f(z, \epsilon) dz &= \frac{1}{2\pi i} \int_{C_0} f(z, \epsilon) dz \\ &\quad + \left(\frac{1}{2\pi i} \int_C f(z, \epsilon) dz - \frac{1}{2\pi i} \int_{C_0} f(z, \epsilon) dz \right) \\ &= \frac{1}{2\pi i} \int_{C_0} f(z, \epsilon) dz + \operatorname{res}_{z=\epsilon} f(z, \epsilon) \end{aligned}$$

Strategy B Example 1

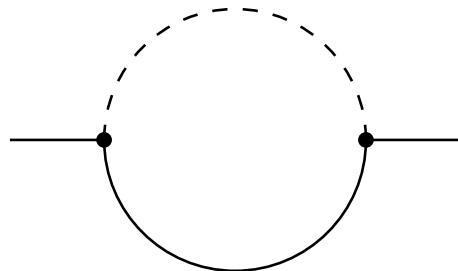
[J.B. Tausk'99, C. Anastasiou & A. Daleo'05, Czakon'05]



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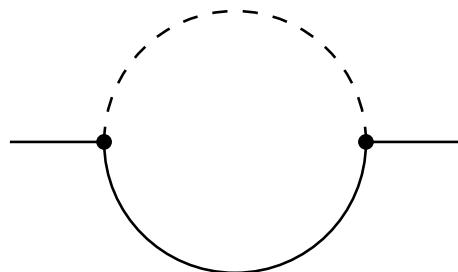


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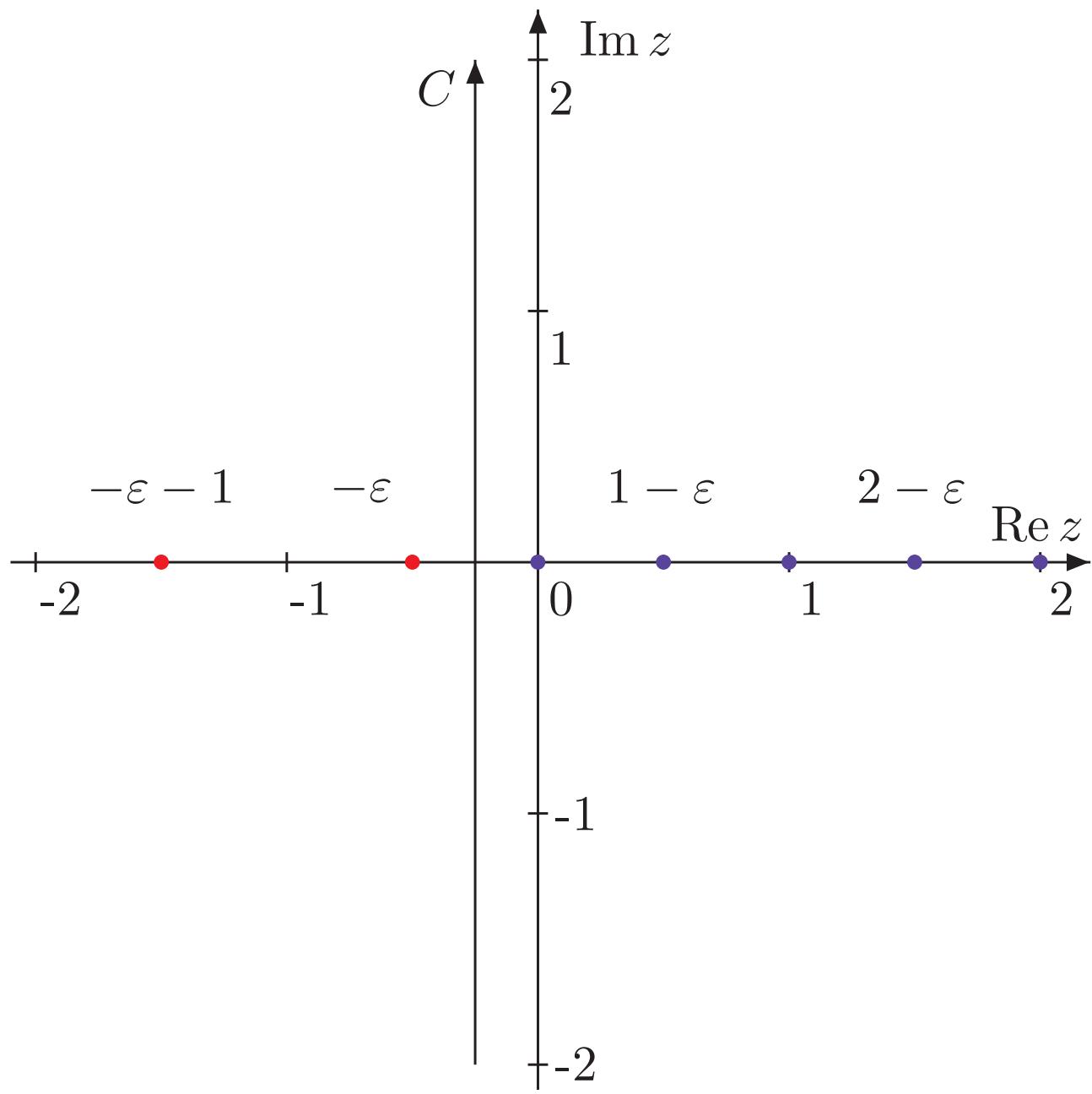
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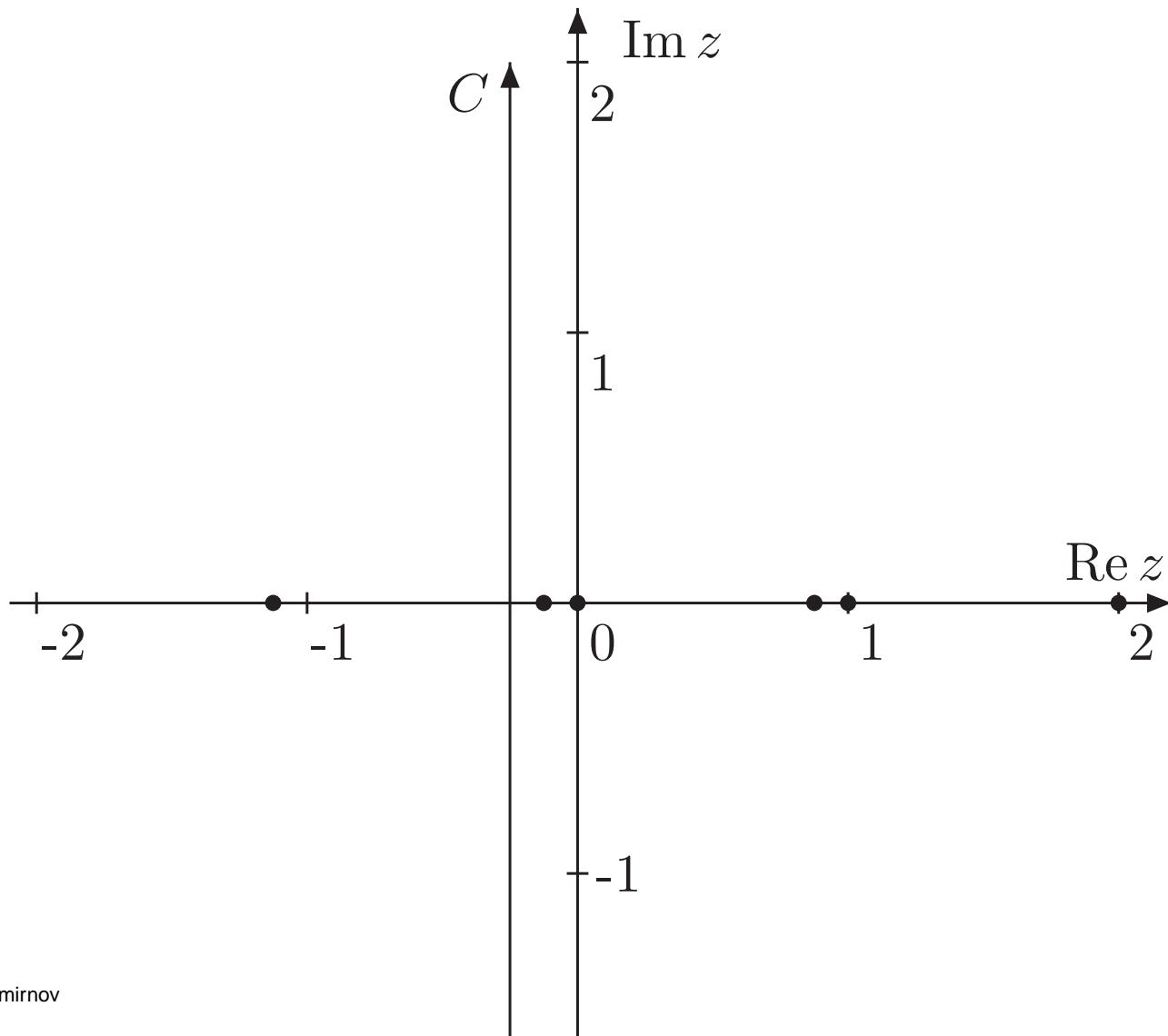
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For example, take $\epsilon = 1/2$, $\operatorname{Re} z = -1/4$. The contour is kept fixed. Tend ϵ to zero.



Whenever a pole of some gamma function is crossed add a residue and tend ϵ to zero further



$$\frac{1}{2\pi i} \int_C f(z, \epsilon) dz = \frac{1}{2\pi i} \int_{\Re z = -1/4} f(z, \epsilon) dz + \text{res}_{z=\epsilon} f(z, \epsilon)$$

General recipes for resolving the singularity structure in ϵ .

$$\frac{1}{(2\pi i)^n} \int_{-i\infty}^{+i\infty} \dots \int_{-i\infty}^{+i\infty} \frac{\prod_i \Gamma(a_i + b_i \epsilon + \sum_j c_{ij} z_j)}{\prod_i \Gamma(a'_i + b'_i \epsilon + \sum_j c'_{ij} z_j)} \prod_k x_k^{d_k} \prod_{l=1}^n dz_l$$

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The goal is to represent a given MB integral as a sum of integrals where a Laurent expansion in ϵ becomes possible.

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Two strategies: Strategy A and Strategy B



Strategy B

[J.B. Tausk'99, C. Anastasiou & A. Daleo'05, M. Czakon'05]

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Let $\epsilon \rightarrow 0$. Whenever a pole of some gamma function is crossed, take into account the corresponding residue.

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For every resulting residue, which involves one integration less, apply a similar procedure, etc.

Two algorithmic descriptions [C. Anastasiou & A. Daleo'05, M. Czakon'05]

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The Czakon's version **MB.m** implemented in Mathematica
is public.

<http://projects.hepforge.org/mbtools/>

- Strategy A in a modified form

[A.V. Smirnov & V.A. Smirnov'09]

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Strategy B: straight contours in the beginning

Strategy A: straight contours in the end

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Let $\prod \Gamma(A_i)$ with $A_i = a_i + b_i\epsilon + \sum_j c_{ij}z_j$

be the numerator of a multiple MB integral

Let $\sigma(x) = [(1 - x)_+]$ where $[\dots]$ is the integer part of a number and $x_+ = x$ for $x > 0$ and 0 otherwise.

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Choose contours, i.e. $\operatorname{Re} z_i$, for which

$$\sum_i \sigma(\operatorname{Re} A_i|_{\epsilon=0}) \equiv \sum_i \sigma \left(a_i + \sum_j c_{ij} \operatorname{Re} z_j \right)$$

is minimal

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Proceed iteratively: every residue is considered from the scratch, i.e. treated in the same way as the initial MB integral

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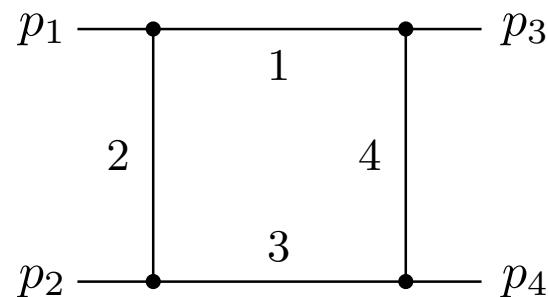
[MBresolve.m](#)

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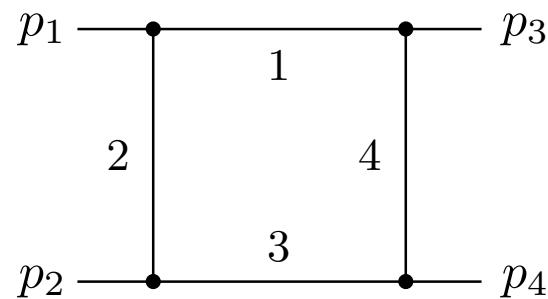
How to derive MB representations

Example 2. The massless on-shell box diagram, i.e. with
 $p_i^2 = 0, i = 1, 2, 3, 4$



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$$F_\Gamma(s, t; a_1, a_2, a_3, a_4, d)$$

$$= \int \frac{d^d k}{(-k^2)^{a_1} [-(k+p_1)^2]^{a_2} [-(k+p_1+p_2)^2]^{a_3} [-(k-p_3)^2]^{a_4}} ,$$

where $s = (p_1 + p_2)^2$ and $t = (p_1 + p_3)^2$

$$\mathcal{U} = \alpha_1 + \alpha_2 + \alpha_3 + \alpha_4 , \quad \mathcal{V} = t\alpha_1\alpha_3 + s\alpha_2\alpha_4 .$$

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$$\begin{aligned} F_{\Gamma}(s, t; a_1, a_2, a_3, a_4, d) &= i\pi^{d/2} \frac{\Gamma(a + \epsilon - 2)}{\prod \Gamma(a_l)} \\ &\times \int_0^\infty \dots \int_0^\infty \frac{\delta \left(\sum_{l=1}^4 \alpha_l - 1 \right)}{(-t\alpha_1\alpha_3 - s\alpha_2\alpha_4)^{a+\epsilon-2}} \prod_l \alpha_l^{a_l-1} d\alpha_1 \dots d\alpha_4 \ , \end{aligned}$$

$$a=a_1+\ldots+a_4.$$

$$\mathcal{U} = \alpha_1 + \alpha_2 + \alpha_3 + \alpha_4 , \quad \mathcal{V} = t\alpha_1\alpha_3 + s\alpha_2\alpha_4 .$$

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$$\times \int_0^\infty \dots \int_0^\infty \frac{\delta \left(\sum_{l=1}^4 \alpha_l - 1 \right)}{(-t\alpha_1\alpha_3 - s\alpha_2\alpha_4)^{a+\epsilon-2}} \prod_l \alpha_l^{a_l-1} d\alpha_1 \dots d\alpha_4 ,$$

$$a = a_1 + \dots + a_4.$$

Introduce new variables by $\alpha_1 = \eta_1\xi_1$, $\alpha_2 = \eta_1(1 - \xi_1)$, $\alpha_3 = \eta_2\xi_2$, $\alpha_4 = \eta_2(1 - \xi_2)$, **with the Jacobian** $\eta_1\eta_2$

$$\begin{aligned}
& F_{\Gamma}(s, t; a_1, a_2, a_3, a_4, d) \\
&= i\pi^{d/2} \frac{\Gamma(a + \epsilon - 2)\Gamma(2 - \epsilon - a_1 - a_2)\Gamma(2 - \epsilon - a_3 - a_4)}{\Gamma(4 - 2\epsilon - a) \prod \Gamma(a_l)} \\
&\times \int_0^1 \int_0^1 \frac{\xi_1^{a_1-1}(1-\xi_1)^{a_2-1}\xi_2^{a_3-1}(1-\xi_2)^{a_4-1}}{[-s\xi_1\xi_2 - t(1-\xi_1)(1-\xi_2) - i0]^{a+\epsilon-2}} d\xi_1 d\xi_2
\end{aligned}$$

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& F_{\Gamma}(s, t; a_1, a_2, a_3, a_4, d) \\
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&\quad \times \int_0^1 \int_0^1 \frac{\xi_1^{a_1-1}(1-\xi_1)^{a_2-1}\xi_2^{a_3-1}(1-\xi_2)^{a_4-1}}{[-s\xi_1\xi_2 - t(1-\xi_1)(1-\xi_2) - i0]^{a+\epsilon-2}} d\xi_1 d\xi_2
\end{aligned}$$

Apply the basic formula to separate
 $-s\xi_1\xi_2$ and $-t(1-\xi_1)(1-\xi_2)$ in the denominator

$$\begin{aligned}
& F_\Gamma(s, t; a_1, a_2, a_3, a_4, d) \\
&= i\pi^{d/2} \frac{\Gamma(a + \epsilon - 2)\Gamma(2 - \epsilon - a_1 - a_2)\Gamma(2 - \epsilon - a_3 - a_4)}{\Gamma(4 - 2\epsilon - a) \prod \Gamma(a_l)} \\
&\quad \times \int_0^1 \int_0^1 \frac{\xi_1^{a_1-1}(1-\xi_1)^{a_2-1}\xi_2^{a_3-1}(1-\xi_2)^{a_4-1}}{[-s\xi_1\xi_2 - t(1-\xi_1)(1-\xi_2) - i0]^{a+\epsilon-2}} d\xi_1 d\xi_2
\end{aligned}$$

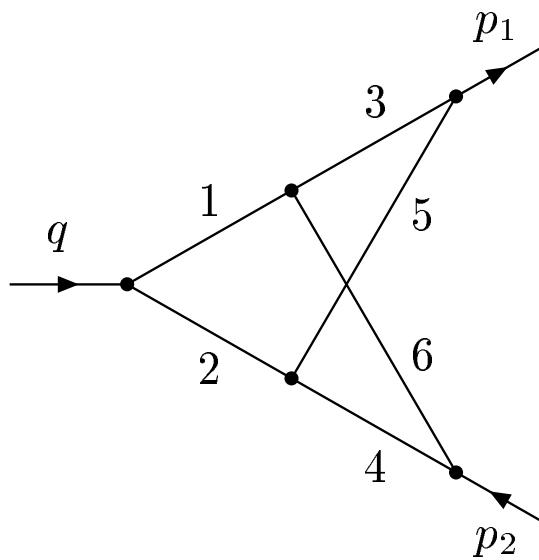
Apply the basic formula to separate
 $-s\xi_1\xi_2$ and $-t(1-\xi_1)(1-\xi_2)$ in the denominator

Change the order of integration over z and ξ -parameters,
evaluate parametric integrals in terms of gamma functions

$$\begin{aligned}
F_\Gamma(s, t; a_1, a_2, a_3, a_4, d) &= \frac{i\pi^{d/2}}{\Gamma(4 - 2\epsilon - a) \prod \Gamma(a_l)(-s)^{a+\epsilon-2}} \\
&\times \frac{1}{2\pi i} \int_{-i\infty}^{+i\infty} dz \left(\frac{t}{s}\right)^z \Gamma(a + \epsilon - 2 + z) \Gamma(a_2 + z) \Gamma(a_4 + z) \Gamma(-z) \\
&\times \Gamma(2 - a_1 - a_2 - a_4 - \epsilon - z) \Gamma(2 - a_2 - a_3 - a_4 - \epsilon - z)
\end{aligned}$$

$$\begin{aligned}
F_{\Gamma}(s, t; a_1, a_2, a_3, a_4, d) &= \frac{i\pi^{d/2}}{\Gamma(4 - 2\epsilon - a) \prod \Gamma(a_l)(-s)^{a+\epsilon-2}} \\
&\times \frac{1}{2\pi i} \int_{-i\infty}^{+i\infty} dz \left(\frac{t}{s}\right)^z \Gamma(a + \epsilon - 2 + z) \Gamma(a_2 + z) \Gamma(a_4 + z) \Gamma(-z) \\
&\times \Gamma(2 - a_1 - a_2 - a_4 - \epsilon - z) \Gamma(2 - a_2 - a_3 - a_4 - \epsilon - z) \\
F_{\Gamma}(s, t; 1, 1, 1, 1, d) &= \frac{i\pi^{d/2}}{\Gamma(-2\epsilon)(-s)^{2+\epsilon}} \\
&\times \frac{1}{2\pi i} \int_{-i\infty}^{+i\infty} dz \left(\frac{t}{s}\right)^z \Gamma(2 + \epsilon + z) \Gamma(1 + \cancel{z})^2 \Gamma(-1 - \epsilon - \cancel{z})^2 \Gamma(-z)
\end{aligned}$$

Example 3. Non-planar two-loop massless vertex diagram
 with $p_1^2 = p_2^2 = 0$, $Q^2 = -(p_1 - p_2)^2 = 2p_1 \cdot p_2$



$$\begin{aligned}
 F_\Gamma(Q^2; a_1, \dots, a_6, d) &= \int \int \frac{\mathbf{d}^d k \mathbf{d}^d l}{[(k + l)^2 - 2p_1 \cdot (k + l)]^{a_1}} \\
 &\times \frac{1}{[(k + l)^2 - 2p_2 \cdot (k + l)]^{a_2} (k^2 - 2p_1 \cdot k)^{a_3} (l^2 - 2p_2 \cdot l)^{a_4} (k^2)^{a_5} (l^2)^{a_6}}
 \end{aligned}$$

$$\frac{1}{(k^2 - 2p_1 \cdot k)^{a_3} (k^2)^{a_5}} = \frac{(-1)^{a_3+a_5} \Gamma(a_3 + a_5)}{\Gamma(a_3) \Gamma(a_5)} \\ \times \int_0^1 \frac{d\xi_1 \xi_1^{a_3-1} (1-\xi_1)^{a_5-1}}{[-(k - \xi_1 p_1)^2 - i0]^{a_3+a_5}}$$

and, similarly, for the second pair, with the replacements

$$\xi_1 \rightarrow \xi_2, \quad p_1 \rightarrow p_2, \quad k \rightarrow l, \quad a_3 \rightarrow a_4, \quad a_5 \rightarrow a_6$$

Change the integration variable $l \rightarrow r = k + l$ and integrate over k by means of our massless one-loop formula

$$\int \frac{dk}{[-(k - \xi_1 p_1)^2]^{a_3+a_5} [-(r - \xi_2 p_2 - k)^2]^{a_4+a_6}} \\ = i\pi^{d/2} \frac{G(a_3 + a_5, a_4 + a_6)}{[-(r - \xi_1 p_1 - \xi_2 p_2)^2]^{a_3+a_4+a_5+a_6+\epsilon-2}}$$

Apply Feynman parametric formula to the propagators 1 and 2 and the propagator arising from the previous integration, with a resulting integral over r evaluated in terms of gamma functions:

$$\int \frac{d^d r}{[-(r^2 - Q^2 A(\xi_1, \xi_2, \xi_3, \xi_4))]^{a+\epsilon-2}} \\ = i\pi^{d/2} \frac{\Gamma(a + 2\epsilon - 4)}{\Gamma(a + \epsilon - 2)} \frac{1}{(Q^2)^{a+2\epsilon-4} A(\xi_1, \xi_2, \xi_3, \xi_4)^{a+2\epsilon-4}}$$

where $a = a_1 + \dots + a_6$ **and**

$$A(\xi_1, \xi_2, \xi_3, \xi_4) = \xi_3 \xi_4 + (1 - \xi_3 - \xi_4)[\xi_2 \xi_3 (1 - \xi_1) + \xi_1 \xi_4 (1 - \xi_2)]$$

Gonsalves'83:

$$\begin{aligned} F_\Gamma(Q^2; a_1, \dots, a_6, d) &= \frac{(-1)^a \left(i\pi^{d/2}\right)^2 \Gamma(2 - \epsilon - a_{35}) \Gamma(2 - \epsilon - a_{46})}{(Q^2)^{a+2\epsilon-4} \prod \Gamma(a_l) \Gamma(4 - 2\epsilon - a_{3456})} \\ &\times \Gamma(a + 2\epsilon - 4) \int_0^1 d\xi_1 \dots \int_0^1 d\xi_4 \xi_1^{a_3-1} (1 - \xi_1)^{a_5-1} \xi_2^{a_4-1} (1 - \xi_2)^{a_6-1} \\ &\times \xi_3^{a_1-1} \xi_4^{a_2-1} (1 - \xi_3 - \xi_4)_+^{a_{3456}+\epsilon-3} A(\xi_1, \xi_2, \xi_3, \xi_4)^{4-2\epsilon-a} \end{aligned}$$

$$\begin{aligned}
& \frac{\Gamma(a + 2\epsilon - 4)}{[\eta\xi(1 - \xi) + (1 - \eta)(\xi\xi_2(1 - \xi_1) + (1 - \xi)\xi_1(1 - \xi_2))]^{a+2\epsilon-4}} \\
&= \frac{1}{2\pi i} \int_{-i\infty}^{+i\infty} \frac{dz_1 \Gamma(-z_1) \eta^{z_1} \xi^{z_1} (1 - \xi)^{z_1}}{(1 - \eta)^{a+2\epsilon-4+z_1}} \\
&\times \frac{\Gamma(a + 2\epsilon - 4 + z_1)}{[\xi\xi_2(1 - \xi_1) + (1 - \xi)\xi_1(1 - \xi_2)]^{a+2\epsilon-4+z_1}}
\end{aligned}$$

The last line →

$$\frac{1}{2\pi i} \int_{-i\infty}^{+i\infty} \frac{dz_2 \Gamma(a + 2\epsilon - 4 + z_1 + z_2) \Gamma(-z_2) \xi^{z_2} \xi_2^{z_2} (1 - \xi_1)^{z_2}}{(1 - \xi)^{a+2\epsilon-4+z_1+z_2} \xi_1^{a+2\epsilon-4+z_1+z_2} (1 - \xi_2)^{a+2\epsilon-4+z_1+z_2}}$$

$$\begin{aligned}
F_\Gamma(Q^2; a_1, \dots, a_6, d) &= \frac{(-1)^a \left(i\pi^{d/2} \right)^2 \Gamma(2 - \epsilon - a_{35})}{(Q^2)^{a+2\epsilon-4} \Gamma(6 - 3\epsilon - a) \prod \Gamma(a_l)} \\
&\times \frac{\Gamma(2 - \epsilon - a_{46})}{\Gamma(4 - 2\epsilon - a_{3456})} \frac{1}{(2\pi i)^2} \int_{-i\infty}^{+i\infty} \int_{-i\infty}^{+i\infty} dz_1 dz_2 \Gamma(a + 2\epsilon - 4 + z_1 + z_2) \\
&\quad \times \Gamma(-z_1) \Gamma(-z_2) \Gamma(a_4 + z_2) \Gamma(a_5 + z_2) \Gamma(a_1 + z_1 + z_2) \\
&\quad \times \frac{\Gamma(2 - \epsilon - a_{12} - z_1) \Gamma(4 - 2\epsilon + a_2 - a - z_2)}{\Gamma(4 - 2\epsilon - a_{1235} - z_1) \Gamma(4 - 2\epsilon - a_{1246} - z_1)} \\
&\quad \times \Gamma(4 - 2\epsilon + a_3 - a - z_1 - z_2) \Gamma(4 - 2\epsilon + a_6 - a - z_1 - z_2) ,
\end{aligned}$$

where $a_{3456} = a_3 + a_4 + a_5 + a_6$, etc.

The massless box diagram with two legs on shell,
 $p_3^2 = p_4^2 = 0$, and two legs off shell, $p_1^2, p_2^2 \neq 0$

$$\begin{aligned}
B_{1100} &= i\pi^{d/2} \frac{\Gamma(a + \epsilon - 2)}{\prod \Gamma(a_l)} \\
&\times \int_0^\infty \dots \int_0^\infty \left(\prod_{l=1}^4 \alpha_l^{a_l-1} d\alpha_l \right) \delta \left(\sum_{l=1}^4 \alpha_l - 1 \right) \\
&\times (-s\alpha_1\alpha_3 - t\alpha_2\alpha_4 - p_1^2\alpha_1\alpha_2 - p_2^2\alpha_2\alpha_3 - i0)^{2-a-\epsilon}
\end{aligned}$$

Apply

$$\frac{1}{(X_1 + \dots + X_n)^\lambda} = \frac{1}{\Gamma(\lambda)} \frac{1}{(2\pi i)^{n-1}} \int_{-i\infty}^{+i\infty} \dots \int_{-i\infty}^{+i\infty} dz_2 \dots dz_n \prod_{i=2}^n X_i^{z_i}$$
$$\times X_1^{-\lambda - z_2 - \dots - z_n} \Gamma(\lambda + z_2 + \dots + z_n) \prod_{i=2}^n \Gamma(-z_i)$$

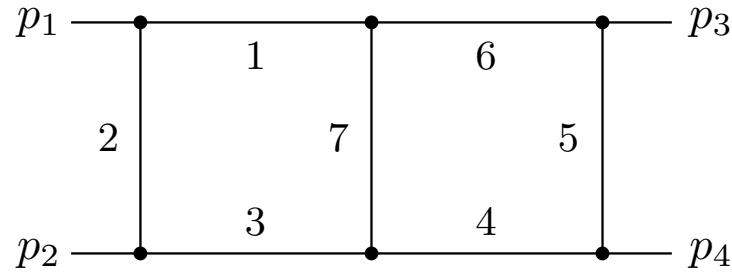
Separate terms with p_1^2 and p_2^2 , turn to new variables by

$$\alpha_1 = \eta_1 \xi_1, \quad \alpha_2 = \eta_1 (1 - \xi_1), \quad \alpha_3 = \eta_2 \xi_2, \quad \alpha_4 = \eta_2 (1 - \xi_2)$$

and evaluate integrals over parameters to obtain a three fold MB representation

$$\begin{aligned}
B_{1100} &= \frac{i\pi^{d/2}}{\Gamma(4 - 2\epsilon - a) \prod \Gamma(a_l)(-s)^{a+\epsilon-2}} \\
&\times \frac{1}{(2\pi i)^3} \int_{-i\infty}^{+i\infty} \int_{-i\infty}^{+i\infty} \int_{-i\infty}^{+i\infty} dz_2 dz_3 dz_4 \frac{(-p_1^2)^{z_2} (-p_2^2)^{z_3} (-t)^{z_4}}{(-s)^{z_2+z_3+z_4}} \\
&\times \Gamma(a + \epsilon - 2 + z_2 + z_3 + z_4) \Gamma(a_2 + z_2 + z_3 + z_4) \Gamma(a_4 + z_4) \\
&\times \Gamma(2 - \epsilon - a_{234} - z_3 - z_4) \Gamma(2 - \epsilon - a_{124} - z_2 - z_4) \\
&\times \Gamma(-z_2) \Gamma(-z_3) \Gamma(-z_4) .
\end{aligned}$$

Double box with irreducible numerator $(k + p_1 + p_2 + p_4)^2$



$$\begin{aligned}
 B_2(s, t; a_1, \dots, a_8, \epsilon) &= \int \int \frac{\mathbf{d}^d k \mathbf{d}^d l}{(k^2)^{a_1} [(k + p_1)^2]^{a_2} [(k + p_1 + p_2)^2]^{a_3}} \\
 &\times \frac{[(k + p_1 + p_2 + p_4)^2]^{-a_8}}{[(l + p_1 + p_2)^2]^{a_4} [(l + p_1 + p_2 + p_4)^2]^{a_5} (l^2)^{a_6} [(k - l)^2]^{a_7}}
 \end{aligned}$$

$$\begin{aligned}
B_2(s, t; a_1, \dots, a_8, \epsilon) &= \int \frac{\mathbf{d}^d k [(k + p_1 + p_2 + p_4)^2]^{-a_8}}{(k^2)^{a_1} [(k + p_1)^2]^{a_2} [(k + p_1 + p_2)^2]^{a_3}} \\
&\times B_{1100}(s, (k + p_1 + p_2 + p_4)^2, k^2, (k + p_1 + p_2)^2; a_6, a_7, a_4, a_5, d)
\end{aligned}$$

$$B_2(s, t; a_1, \dots, a_8, \epsilon) = \int \frac{\mathbf{d}^d k [(k + p_1 + p_2 + p_4)^2]^{-a_8}}{(k^2)^{a_1} [(k + p_1)^2]^{a_2} [(k + p_1 + p_2)^2]^{a_3}} \\ \times B_{1100}(s, (k + p_1 + p_2 + p_4)^2, k^2, (k + p_1 + p_2)^2; a_6, a_7, a_4, a_5, d)$$

After using the threefold MB representation for B_{1100} and changing the order of integration we obtain an on-shell box integral with indices shifted by z -variables. Apply then the onefold representation for the this box.

$$B_2(s, t; a_1, \dots, a_8, \epsilon) = \int \frac{\mathbf{d}^d k [(k + p_1 + p_2 + p_4)^2]^{-a_8}}{(k^2)^{a_1} [(k + p_1)^2]^{a_2} [(k + p_1 + p_2)^2]^{a_3}} \\ \times B_{1100}(s, (k + p_1 + p_2 + p_4)^2, k^2, (k + p_1 + p_2)^2; a_6, a_7, a_4, a_5, d)$$

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The loop by loop derivation of MB representations.

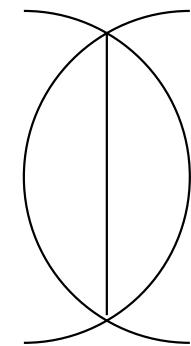
AMBRE

[J. Gluza, K. Kajda & T. Riemann'07]

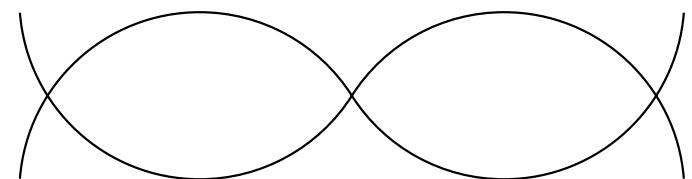
How to check a given MB representation

	a_1		a_6	
a_2		a_7		a_5
	a_3		a_4	

$$a_1, a_3, a_4, a_6 \rightarrow 0$$



$$a_2, a_5, a_7 \rightarrow 0$$



How to evaluate MB integrals after expanding in ϵ

The first Barnes lemma

$$\begin{aligned} & \frac{1}{2\pi i} \int_{-i\infty}^{+i\infty} dz \Gamma(\lambda_1 + z) \Gamma(\lambda_2 + z) \Gamma(\lambda_3 - z) \Gamma(\lambda_4 - z) \\ &= \frac{\Gamma(\lambda_1 + \lambda_3) \Gamma(\lambda_1 + \lambda_4) \Gamma(\lambda_2 + \lambda_3) \Gamma(\lambda_2 + \lambda_4)}{\Gamma(\lambda_1 + \lambda_2 + \lambda_3 + \lambda_4)} \end{aligned}$$

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Multiple corollaries, e.g.,

$$\begin{aligned} & \frac{1}{2\pi i} \int_{-i\infty}^{+i\infty} dz \Gamma(\lambda_1 + z) \Gamma^*(\lambda_2 + z) \Gamma(-\lambda_2 - z) \Gamma(\lambda_3 - z) \\ &= \Gamma(\lambda_1 - \lambda_2) \Gamma(\lambda_2 + \lambda_3) [\psi(\lambda_1 - \lambda_2) - \psi(\lambda_1 + \lambda_3)] \end{aligned}$$

$$\begin{aligned}
& \frac{1}{2\pi i} \int_{-i\infty}^{+i\infty} \frac{dz}{z} \Gamma(\lambda_1 + z) \Gamma(\lambda_2 + z) \Gamma(\lambda_3 - z) \Gamma(\lambda_4 - z) \\
&= \frac{\Gamma(2 - \lambda_1 - \lambda_3) \Gamma(1 - \lambda_2 - \lambda_3) \Gamma(\lambda_1 + \lambda_3 - 1) \Gamma(\lambda_2 + \lambda_3)}{\Gamma(1 - \lambda_1) \Gamma(1 - \lambda_2)} \\
&\times [\Gamma(1 - \lambda_1) \Gamma(1 - \lambda_2) - \Gamma(2 - \lambda_1 - \lambda_2 - \lambda_3) \Gamma(\lambda_3)]
\end{aligned}$$

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&\quad \times [\Gamma(1 - \lambda_1) \Gamma(1 - \lambda_2) - \Gamma(2 - \lambda_1 - \lambda_2 - \lambda_3) \Gamma(\lambda_3)]
\end{aligned}$$

The second Barnes lemma

$$\begin{aligned}
& \frac{1}{2\pi i} \int_{-i\infty}^{+i\infty} dz \frac{\Gamma(\lambda_1 + z) \Gamma(\lambda_2 + z) \Gamma(\lambda_3 + z) \Gamma(\lambda_4 - z) \Gamma(\lambda_5 - z)}{\Gamma(\lambda_6 + z)} \\
&= \frac{\Gamma(\lambda_1 + \lambda_4) \Gamma(\lambda_2 + \lambda_4) \Gamma(\lambda_3 + \lambda_4) \Gamma(\lambda_1 + \lambda_5)}{\Gamma(\lambda_1 + \lambda_2 + \lambda_4 + \lambda_5) \Gamma(\lambda_1 + \lambda_3 + \lambda_4 + \lambda_5)} \\
&\quad \times \frac{\Gamma(\lambda_2 + \lambda_5) \Gamma(\lambda_3 + \lambda_5)}{\Gamma(\lambda_2 + \lambda_3 + \lambda_4 + \lambda_5)}, \quad \lambda_6 = \lambda_1 + \lambda_2 + \lambda_3 + \lambda_4 + \lambda_5
\end{aligned}$$

Transform a given multiple MB integral originating after expanding in ϵ into multiple series.

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Summing up series with nested sums

$$S_i(n) = \sum_{j=1}^n \frac{1}{j^i}, \quad S_{ik}(n) = \sum_{j=1}^n \frac{S_k(j)}{j^i},$$

$$S_{ikl}(n) = \sum_{j=1}^n \frac{S_{kl}(j)}{j^i}, \quad S_{iklm}(n) = \sum_{j=1}^n \frac{S_{klm}(j)}{j^i}$$

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For example, with one index:

$$\begin{aligned}\psi(n) &= S_1(n-1) - \gamma_E, \\ \psi^{(k)}(n) &= (-1)^k k! (S_{k+1}(n-1) - \zeta(k+1)), \quad k = 1, 2, \dots,\end{aligned}$$

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SUMMER

[J.A.M. Vermaseren'00]

Harmonic polylogarithms (HPL)

$H_{a_1, a_2, \dots, a_n}(x) \equiv H(a_1, a_2, \dots, a_n; x)$, with $a_i = 1, 0, -1$

[E. Remiddi & J.A.M. Vermaseren'00]

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[E. Remiddi & J.A.M. Vermaseren'00]

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$$H(a_1, a_2, \dots, a_n; x) = \int_0^x f(a_1; t) H(a_2, \dots, a_n; t) dt,$$

where $f(\pm 1; t) = 1/(1 \mp t)$, $f(0; t) = 1/t$,

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are generalizations of the usual polylogarithms $\text{Li}_a(z)$ and
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Harmonic polylogarithms (HPL)

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HPL implemented in Mathematica

[D. Maitre'06]

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MB representations for non-planar diagrams?

The loop by loop strategy is hardly applicable for non-planar
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Poles in ϵ can arise not only locally but also from an
integration over large z .

$$\frac{1}{2\pi i} \int_C \frac{\Gamma(1+2\epsilon+z)\Gamma(-z)}{1+\epsilon+z} (-1)^z dz$$

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$$\Gamma(x \pm iy) \sim \sqrt{2\pi} e^{\pm i\frac{\pi}{4}(2x-1)} e^{\pm iy(\ln y - 1)} e^{-\frac{\pi}{2}y}$$

when $y \rightarrow +\infty$

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The integrand behaves like

$$2\pi \frac{1}{y^{1-2\epsilon}}$$

MB tools at <http://projects.hepforge.org/mbtools/>:

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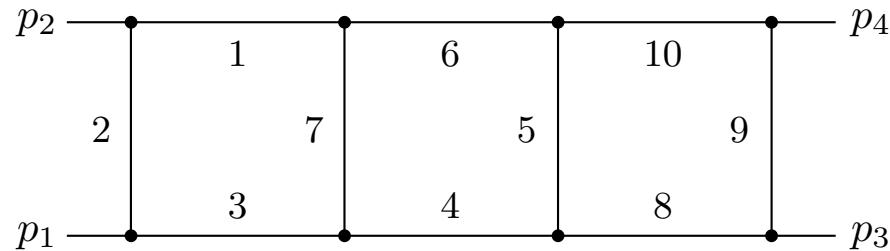
MBresolve.m [A. Smirnov'09]

MBasymptotics.m [M. Czakon'09]

barnesroutines.m [D. Kosower'08]

(applying Barnes lemmas automatically)

backup slides



The general planar triple box Feynman integral

$$\begin{aligned}
 T(a_1, \dots, a_{10}; s, t; \epsilon) &= \int \int \int \frac{\mathbf{d}^d k \mathbf{d}^d l \mathbf{d}^d r}{[k^2]^{a_1} [(k + p_2)^2]^{a_2}} \\
 &\times \frac{1}{[(k + p_1 + p_2)^2]^{a_3} [(l + p_1 + p_2)^2]^{a_4} [(r - l)^2]^{a_5} [l^2]^{a_6} [(k - l)^2]^{a_7}} \\
 &\times \frac{1}{[(r + p_1 + p_2)^2]^{a_8} [(r + p_1 + p_2 + p_3)^2]^{a_9} [r^2]^{a_{10}}}
 \end{aligned}$$

General 7fold MB representation:

$$\begin{aligned}
T(a_1, \dots, a_{10}; s, t, m^2; \epsilon) &= \frac{\left(i\pi^{d/2}\right)^3 (-1)^a}{\prod_{j=2,5,7,8,9,10} \Gamma(a_j) \Gamma(4 - a_{589(10)} - 2\epsilon) (-s)^{a-6+3\epsilon}} \\
&\times \frac{1}{(2\pi i)^7} \int_{-i\infty}^{+i\infty} dw \prod_{j=2}^7 dz_j \left(\frac{t}{s}\right)^w \frac{\Gamma(a_2 + w)\Gamma(-w)\Gamma(z_2 + z_4)\Gamma(z_3 + z_4)}{\Gamma(a_1 + z_3 + z_4)\Gamma(a_3 + z_2 + z_4)} \\
&\times \frac{\Gamma(2 - a_1 - a_2 - \epsilon + z_2)\Gamma(2 - a_2 - a_3 - \epsilon + z_3)\Gamma(a_7 + w - z_4)}{\Gamma(4 - a_1 - a_2 - a_3 - 2\epsilon + w - z_4)\Gamma(a_6 - z_5)\Gamma(a_4 - z_6)} \\
&\times \Gamma(+a_1 + a_2 + a_3 - 2 + \epsilon + z_4)\Gamma(w + z_2 + z_3 + z_4 - z_7)\Gamma(-z_5)\Gamma(-z_6) \\
&\times \Gamma(2 - a_5 - a_9 - a_{10} - \epsilon - z_5 - z_7)\Gamma(2 - a_5 - a_8 - a_9 - \epsilon - z_6 - z_7) \\
&\times \Gamma(a_4 + a_6 + a_7 - 2 + \epsilon + w - z_4 - z_5 - z_6 - z_7)\Gamma(a_9 + z_7) \\
&\times \Gamma(4 - a_4 - a_6 - a_7 - 2\epsilon + z_5 + z_6 + z_7) \\
&\times \Gamma(2 - a_6 - a_7 - \epsilon - w - z_2 + z_5 + z_7)\Gamma(2 - a_4 - a_7 - \epsilon - w - z_3 + z_6 + z_7) \\
&\times \Gamma(a_5 + z_5 + z_6 + z_7)\Gamma(a_5 + a_8 + a_9 + a_{10} - 2 + \epsilon + z_5 + z_6 + z_7),
\end{aligned}$$

General 7fold MB representation:

$$\begin{aligned}
T(a_1, \dots, a_{10}; s, t, m^2; \epsilon) &= \frac{\left(i\pi^{d/2}\right)^3 (-1)^a}{\prod_{j=2,5,7,8,9,10} \Gamma(a_j) \Gamma(4 - a_{589(10)} - 2\epsilon) (-s)^{a-6+3\epsilon}} \\
&\times \frac{1}{(2\pi i)^7} \int_{-i\infty}^{+i\infty} dw \prod_{j=2}^7 dz_j \left(\frac{t}{s}\right)^w \frac{\Gamma(a_2 + w)\Gamma(-w)\Gamma(z_2 + z_4)\Gamma(z_3 + z_4)}{\Gamma(a_1 + z_3 + z_4)\Gamma(a_3 + z_2 + z_4)} \\
&\times \frac{\Gamma(2 - a_1 - a_2 - \epsilon + z_2)\Gamma(2 - a_2 - a_3 - \epsilon + z_3)\Gamma(a_7 + w - z_4)}{\Gamma(4 - a_1 - a_2 - a_3 - 2\epsilon + w - z_4)\Gamma(a_6 - z_5)\Gamma(a_4 - z_6)} \\
&\times \Gamma(+a_1 + a_2 + a_3 - 2 + \epsilon + z_4)\Gamma(w + z_2 + z_3 + z_4 - z_7)\Gamma(-z_5)\Gamma(-z_6) \\
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&\times \Gamma(2 - a_6 - a_7 - \epsilon - w - z_2 + z_5 + z_7)\Gamma(2 - a_4 - a_7 - \epsilon - w - z_3 + z_6 + z_7) \\
&\times \Gamma(a_5 + z_5 + z_6 + z_7)\Gamma(a_5 + a_8 + a_9 + a_{10} - 2 + \epsilon + z_5 + z_6 + z_7),
\end{aligned}$$

$$\begin{aligned}
& T(1, 1, \dots, 1; s, t; \epsilon) \\
&= \frac{(i\pi^{d/2})^3}{\Gamma(-2\epsilon)(-s)^{4+3\epsilon}} \frac{1}{(2\pi i)^7} \int_{-i\infty}^{+i\infty} dw \prod_{j=2}^7 dz_j \left(\frac{t}{s}\right)^w \frac{\Gamma(1+w)\Gamma(-w)}{\Gamma(1-2\epsilon+w-z_4)} \\
&\times \frac{\Gamma(-\epsilon+z_2)\Gamma(-\epsilon+z_3)\Gamma(1+w-z_4)\Gamma(-z_2-z_3-z_4)\Gamma(1+\epsilon+z_4)}{\Gamma(1+z_2+z_4)\Gamma(1+z_3+z_4)} \\
&\times \frac{\Gamma(z_2+z_4)\Gamma(z_3+z_4)\Gamma(-z_5)\Gamma(-z_6)\Gamma(w+z_2+z_3+z_4-z_7)}{\Gamma(1-z_5)\Gamma(1-z_6)\Gamma(1-2\epsilon+z_5+z_6+z_7)} \\
&\times \Gamma(-1-\epsilon-z_5-z_7)\Gamma(-1-\epsilon-z_6-z_7)\Gamma(1+z_7) \\
&\times \Gamma(1+\epsilon+w-z_4-z_5-z_6-z_7)\Gamma(-\epsilon-w-z_2+z_5+z_7) \\
&\times \Gamma(-\epsilon-w-z_3+z_6+z_7)\Gamma(1+z_5+z_6+z_7)\Gamma(2+\epsilon+z_5+z_6+z_7)
\end{aligned}$$

Result

[V.A. Smirnov'03]

$$T(1, 1, \dots, 1; s, t; \epsilon) = -\frac{\left(i\pi^{d/2}e^{-\gamma_E\epsilon}\right)^3}{s^3(-t)^{1+3\epsilon}} \sum_{j=0}^6 \frac{c_j(x, L)}{\epsilon^j},$$

where $x = -t/s$, $L = \ln(s/t)$, and

$$c_6 = \frac{16}{9}, \quad c_5 = -\frac{5}{3}L, \quad c_4 = -\frac{3}{2}\pi^2,$$

$$c_3 = 3(H_{0,0,1}(x) + LH_{0,1}(x)) + \frac{3}{2}(L^2 + \pi^2)H_1(x) - \frac{11}{12}\pi^2L - \frac{131}{9}\zeta_3,$$

$$c_2 = -3(17H_{0,0,0,1}(x) + H_{0,0,1,1}(x) + H_{0,1,0,1}(x) + H_{1,0,0,1}(x))$$

$$-L(37H_{0,0,1}(x) + 3H_{0,1,1}(x) + 3H_{1,0,1}(x)) - \frac{3}{2}(L^2 + \pi^2)H_{1,1}(x)$$

$$-\left(\frac{23}{2}L^2 + 8\pi^2\right)H_{0,1}(x) - \left(\frac{3}{2}L^3 + \pi^2L - 3\zeta_3\right)H_1(x) + \frac{49}{3}\zeta_3L - \frac{1411}{1080}\pi^4,$$

$$\begin{aligned}
c_1 = & 3(81H_{0,0,0,0,1}(x) + 41H_{0,0,0,1,1}(x) + 37H_{0,0,1,0,1}(x) + H_{0,0,1,1,1}(x) \\
& + 33H_{0,1,0,0,1}(x) + H_{0,1,0,1,1}(x) + H_{0,1,1,0,1}(x) + 29H_{1,0,0,0,1}(x) \\
& + H_{1,0,0,1,1}(x) + H_{1,0,1,0,1}(x) + H_{1,1,0,0,1}(x)) + L(177H_{0,0,0,1}(x) + 85H_{0,0,1,1}(x) \\
& + 73H_{0,1,0,1}(x) + 3H_{0,1,1,1}(x) + 61H_{1,0,0,1}(x) + 3H_{1,0,1,1}(x) + 3H_{1,1,0,1}(x)) \\
& + \left(\frac{119}{2}L^2 + \frac{139}{12}\pi^2\right)H_{0,0,1}(x) + \left(\frac{47}{2}L^2 + 20\pi^2\right)H_{0,1,1}(x) \\
& + \left(\frac{35}{2}L^2 + 14\pi^2\right)H_{1,0,1}(x) + \frac{3}{2}(L^2 + \pi^2)H_{1,1,1}(x) \\
& + \left(\frac{23}{2}L^3 + \frac{83}{12}\pi^2L - 96\zeta_3\right)H_{0,1}(x) + \left(\frac{3}{2}L^3 + \pi^2L - 3\zeta_3\right)H_{1,1}(x) \\
& + \left(\frac{9}{8}L^4 + \frac{25}{8}\pi^2L^2 - 58\zeta_3L + \frac{13}{8}\pi^4\right)H_1(x) - \frac{503}{1440}\pi^4L + \frac{73}{4}\pi^2\zeta_3 - \frac{301}{15}\zeta_5,
\end{aligned}$$

$$\begin{aligned}
c_0 = & - (951H_{0,0,0,0,0,1}(x) + 819H_{0,0,0,0,1,1}(x) + 699H_{0,0,0,1,0,1}(x) + 195H_{0,0,0,1,1,1}(x) \\
& + 547H_{0,0,1,0,0,1}(x) + 231H_{0,0,1,0,1,1}(x) + 159H_{0,0,1,1,0,1}(x) + 3H_{0,0,1,1,1,1}(x) \\
& + 363H_{0,1,0,0,0,1}(x) + 267H_{0,1,0,0,1,1}(x) + 195H_{0,1,0,1,0,1}(x) + 3H_{0,1,0,1,1,1}(x) \\
& + 123H_{0,1,1,0,0,1}(x) + 3H_{0,1,1,0,1,1}(x) + 3H_{0,1,1,1,0,1}(x) + 147H_{1,0,0,0,0,1}(x) \\
& + 303H_{1,0,0,0,1,1}(x) + 231H_{1,0,0,1,0,1}(x) + 3H_{1,0,0,1,1,1}(x) + 159H_{1,0,1,0,0,1}(x) \\
& + 3H_{1,0,1,0,1,1}(x) + 3H_{1,0,1,1,0,1}(x) + 87H_{1,1,0,0,0,1}(x) + 3H_{1,1,0,0,1,1}(x) \\
& + 3H_{1,1,0,1,0,1}(x) + 3H_{1,1,1,0,0,1}(x)) \\
& - L (729H_{0,0,0,0,1}(x) + 537H_{0,0,0,1,1}(x) + 445H_{0,0,1,0,1}(x) + 133H_{0,0,1,1,1}(x) \\
& + 321H_{0,1,0,0,1}(x) + 169H_{0,1,0,1,1}(x) + 97H_{0,1,1,0,1}(x) + 3H_{0,1,1,1,1}(x) \\
& + 165H_{1,0,0,0,1}(x) + 205H_{1,0,0,1,1}(x) + 133H_{1,0,1,0,1}(x) + 3H_{1,0,1,1,1}(x) \\
& + 61H_{1,1,0,0,1}(x) + 3H_{1,1,0,1,1}(x) + 3H_{1,1,1,0,1}(x)) \\
& - \left(\frac{531}{2}L^2 + \frac{89}{4}\pi^2 \right) H_{0,0,0,1}(x) - \left(\frac{311}{2}L^2 + \frac{619}{12}\pi^2 \right) H_{0,0,1,1}(x) \\
& - \left(\frac{247}{2}L^2 + \frac{307}{12}\pi^2 \right) H_{0,1,0,1}(x) - \left(\frac{71}{2}L^2 + 32\pi^2 \right) H_{0,1,1,1}(x)
\end{aligned}$$

$$\begin{aligned}
& - \left(\frac{151}{2}L^2 - \frac{197}{12}\pi^2 \right) H_{1,0,0,1}(x) - \left(\frac{107}{2}L^2 + 50\pi^2 \right) H_{1,0,1,1}(x) \\
& - \left(\frac{35}{2}L^2 + 14\pi^2 \right) H_{1,1,0,1}(x) - \frac{3}{2} (L^2 + \pi^2) H_{1,1,1,1}(x) \\
& - \left(\frac{119}{2}L^3 + \frac{317}{12}\pi^2 L - 455\zeta_3 \right) H_{0,0,1}(x) - \left(\frac{47}{2}L^3 + \frac{179}{12}\pi^2 L \right. \\
& \quad \left. - 120\zeta_3 \right) H_{0,1,1}(x) - \left(\frac{35}{2}L^3 + \frac{35}{12}\pi^2 L - 156\zeta_3 \right) H_{1,0,1}(x) - \left(\frac{3}{2}L^3 + \pi^2 L \right. \\
& \quad \left. - 3\zeta_3 \right) H_{1,1,1}(x) - \left(\frac{69}{8}L^4 + \frac{101}{8}\pi^2 L^2 - 291\zeta_3 L + \frac{559}{90}\pi^4 \right) H_{0,1}(x) \\
& - \left(\frac{9}{8}L^4 + \frac{25}{8}\pi^2 L^2 - 58\zeta_3 L + \frac{13}{8}\pi^4 \right) H_{1,1}(x) \\
& - \left(\frac{27}{40}L^5 + \frac{25}{8}\pi^2 L^3 - \frac{183}{2}\zeta_3 L^2 + \frac{131}{60}\pi^4 L - \frac{37}{12}\pi^2 \zeta_3 + 57\zeta_5 \right) H_1(x) \\
& + \left(\frac{223}{12}\pi^2 \zeta_3 + 149\zeta_5 \right) L + \frac{167}{9}\zeta_3^2 - \frac{624607}{544320}\pi^6.
\end{aligned}$$

‘Inverse Feynman parameters’

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Integrating over a MB variable (not by a Barnes lemma)

$$\begin{aligned} & \frac{1}{2\pi i} \int_C \Gamma(a+z)\Gamma(b+z)\Gamma(c-z)\Gamma(-z)x^z dz \\ &= \frac{\Gamma(a)\Gamma(b)\Gamma(a+c)\Gamma(b+c)}{\Gamma(a+b+c)} {}_2F_1(a; b; a+b+c; 1-x) \\ &= \Gamma(a)\Gamma(b+c) \int_0^1 t^{b-1}(1-t)^{a+c-1}(1-t+tx)^{-a} dt \end{aligned}$$



Strategy A

[V.A. Smirnov'99]

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Analysis of the integrand. Think of integrations over z -variables in various orders.

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For example, the product $\Gamma(1+z)\Gamma(-1-\epsilon-z)$ generates a pole of the type $\Gamma(-\epsilon)$ where $-\epsilon = (1+z) + (-1-\epsilon-z)$

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The general rule: $\Gamma(a+z)\Gamma(b-z)$, where a and b depend on the rest of the variables, generates a pole of the type $\Gamma(a+b)$

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The general rule: $\Gamma(a+z)\Gamma(b-z)$, where a and b depend on the rest of the variables, generates a pole of the type $\Gamma(a+b)$

Identifying key gamma functions (responsible for the generation of poles in ϵ).

Let $\Gamma(A_i)$ with $A_i = a_i + b_i\epsilon + \sum_j c_{ij}z_j$
be one of the key gamma functions. Consider ϵ real.

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'Changing the nature' of these key gamma functions (i.e.
changing rules for the contours)

$$\begin{aligned}\operatorname{Re} A_i > 0 &\rightarrow -1 < \operatorname{Re} A_i < 0 \\ \Gamma(A_i) &\rightarrow \Gamma^{(1)}(A_i)\end{aligned}$$

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Changing more:

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Taking residues and shifting contours.

For each resulting residue, which involves one integration
less, apply a similar procedure, etc.