

# Bootstrapping the three-loop hexagon

Johannes M. Henn



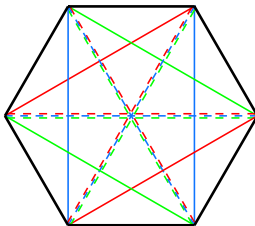
Humboldt-Universität zu Berlin

based on

L. J. Dixon, J. M. Drummond, J. M. H., arXiv:1108.4461 [hep-th]

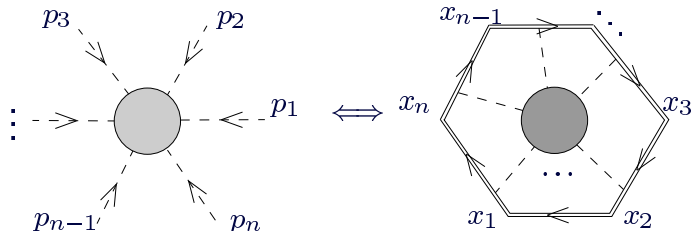
Ustroń 2011, September 13

## Bootstrapping the three-loop hexagon



- Cusped Wilson loops in  $\mathcal{N} = 4$  super Yang-Mills
- Iterated integrals and symbols
- ansatz and constraints for the three-loop hexagon
- prediction for Regge limit

# Duality between Wilson loops and scattering amplitudes



[Alday, Maldacena, 2007; Drummond, Korchemsky, Sokatchev, 2007; Brandhuber, Heslop, Travaglini, 2007]

- Checked by two-loop computations for  $n \leq 6$  points

[Drummond, J.M.H., Korchemsky, Sokatchev, 2007, 2008]

[Bern, Dixon, Kosower, Roiban, Spradlin, Vergu, Volovich, 2008]

- conformal symmetry of Wilson loops  $\rightarrow$  dual conformal symmetry of amplitudes

- **all-order** (dual) conformal Ward identities

[Drummond, J.M.H., Korchemsky, Sokatchev, 2007]

$\Rightarrow$  kinematical dependence of four- and five-point Wilson loops/scattering amplitudes **fixed to all orders in the coupling!**

# Conformal Ward identity and six-point Wilson loop

- known structure of UV divergences [Korchemsky, Radyushkin (1987); Korchemskaya, Korchemsky (1992)]

$$\log W_n = [\text{UV divergent}]_n + F_n^{\text{WL}}.$$

- solution to Ward identity at  $n = 6$  [Drummond, J.M.H., Korchemsky, Sokatchev, 2007]

$$F_6^{\text{WL}} = \gamma_K(a) F_{6,1\text{-loop}}^{\text{WL}} + R_6(a; u, v, w).$$

with  $a \equiv g^2 N_c / (8\pi^2)$ .  $\gamma_K(a)$  cusp anomalous dimension.

- loop expansion of remainder function

$$R_6(a; u, v, w) = \sum_{L=2}^{\infty} a^L R_6^{(L)}(u, v, w), .$$

Depends on three dual conformal cross ratios only,

$$u = \frac{x_{13}^2 x_{46}^2}{x_{14}^2 x_{36}^2}, \quad v = \frac{x_{24}^2 x_{15}^2}{x_{25}^2 x_{14}^2}, \quad w = \frac{x_{35}^2 x_{26}^2}{x_{36}^2 x_{25}^2}.$$

- $R_6^{(L)}$  expected to be expressible as sum of  $2L$ -fold iterated integrals

# Iterated integrals, pure functions and symbols

- we define a *pure* function of degree (or weight)  $k$  recursively,

$$df^{(k)} = \sum_r f_r^{(k-1)} d \log \phi_r .$$

$\phi_r$  are algebraic functions. Degree zero functions are constants.

The definition includes logarithms and classical polylogarithms, as well as other iterated integrals, such as harmonic polylogarithms of one or more variables.

- *symbol*  $\mathcal{S}(f)$  of a pure function  $f$  is defined recursively

$$\mathcal{S}(f^{(k)}) = \sum_r \mathcal{S}(f_r^{(k-1)}) \otimes \phi_r .$$

- $\mathcal{S}(f^{(k)})$  is an element of the  $k$ -fold tensor product of the space of algebraic functions,

$$\mathcal{S}(f^{(k)}) = \sum_{\vec{\alpha}} \phi_{\alpha_1} \otimes \dots \otimes \phi_{\alpha_k} ,$$

# Iterated integrals, pure functions and symbols

Important properties of symbols:

- property derived from  $\log(ab) = \log a + \log b$

$$\dots \otimes \phi_1 \phi_2 \otimes \dots = \dots \otimes \phi_1 \otimes \dots + \dots \otimes \phi_2 \otimes \dots$$

- **integrability condition**  $d^2 f^{(k)} = 0$  for any function implies relations among the different elements.
- **Branch cuts and discontinuities.** Given

$$\mathcal{S}(f^{(k)}) = \sum_{\vec{\alpha}} \phi_{\alpha_1} \otimes \dots \otimes \phi_{\alpha_k},$$

the degree  $k$  function  $f^{(k)}$  will have a branch cut starting at  $\phi_{\alpha_1} = 0$ . The discontinuity across this cut has the symbol

$$\mathcal{S}(\Delta_{\phi_{\alpha_1}} f^{(k)}) = \sum_{\vec{\alpha}} \phi_{\alpha_2} \otimes \dots \otimes \phi_{\alpha_k}.$$

- some examples:

$$\mathcal{S}(\log^2(u)) = 2u \otimes u$$

$$\mathcal{S}(-\text{Li}_2(u)) = (1 - u) \otimes u$$

$$\mathcal{S}(\log(u) \log(v)) = u \otimes v + v \otimes u$$

- Fine print:
  - Symbols do not know on which branch the functions are evaluated.
  - Related: symbols are defined only up to constants times lower degree functions.
- Nonetheless, **the symbol is an extremely useful tool**, especially in multi-variable cases. E.g., symbols were crucial to find a simple form for the two-loop remainder function  $R_6^{(2)}$ .

# Symbol of the two-loop remainder function

- $R_6^{(2)}$  known analytically

[Del Duca, Duhr, Smirnov (2009)]

can be expressed in terms of classical polylogarithms [Goncharov, Spradlin, Vergu, Volovich (2010)]

It has the symbol

$$\begin{aligned} \mathcal{S}(R_6^{(2)}) = & -\frac{1}{8} \left\{ \left[ u \otimes (1-u) \otimes \frac{u}{(1-u)^2} + 2(u \otimes v + v \otimes u) \otimes \frac{w}{1-v} \right. \right. \\ & \left. \left. + 2v \otimes \frac{w}{1-v} \otimes u \right] \otimes \frac{u}{1-u} \right. \\ & \left. + \left[ u \otimes (1-u) \otimes y_u y_v y_w - 2u \otimes v \otimes y_w \right] \otimes y_u y_v y_w \right\} \\ & + \text{permutations,} \end{aligned}$$

Here

$$y_u = \frac{u - z_+}{u - z_-}, \quad z_{\pm} = \frac{1}{2} \left( -1 + u + v + w \pm \sqrt{(1-u-v-w)^2 - 4uvw} \right).$$

- What is the symbol of  $R_6$  at higher loops?



# Ansatz for $\mathcal{S}(R_6^{(3)})$

- Ansatz: symbol built from the following letters:

$$\mathcal{A} = \{u, v, w, 1 - u, 1 - v, 1 - w, y_u, y_v, y_w\}.$$

i.e.

$$\mathcal{S}(R_6^{(3)}) = \sum_{\vec{\alpha}} c_{\vec{\alpha}} \times a_{\alpha_1} \otimes a_{\alpha_2} \otimes a_{\alpha_3} \otimes a_{\alpha_4} \otimes a_{\alpha_5} \otimes a_{\alpha_6},$$

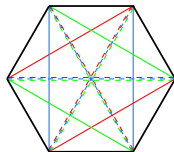
with  $a_{\alpha_i} \in \mathcal{A}$  and  $c_{\vec{\alpha}} \in \mathbb{R}$ .

- Motivation:

- explicit form of two-loop remainder function [Goncharov, Spradlin, Vergu, Volovich (2010)]
- experience with loop integrals appearing in MHV scattering amplitudes

[Drummond, J. .M. .H., Trnka (2010); Dixon, Drummond, J. .M. .H. (2011)]

- Is this ansatz consistent with all known constraints on  $\mathcal{S}(R_6^{(3)})$ ?



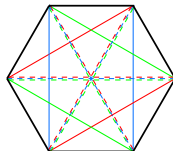
# Constraints on $\mathcal{S}(R_6^{(3)})$

What constraints should the symbol of the remainder function obey?

- should be *integrable*, i.e. symbol of a function
- *discontinuities* of loop integrals at  $x_{ij}^2 = 0 \longrightarrow$  first entry should be  $u, v$  or  $w$
- it should be completely symmetric in  $u, v, w$
- parity even  $\longrightarrow$  even number of  $y_u, y_v, y_w$  variables
- collinear limit:  $R(u, 1 - u, 0) = 0$
- constraints in multi-Regge kinematics

[Lipatov, Prygarin (2010), Bartels, Lipatov, Prygarin (2010)]

$$u \rightarrow 1, \quad \frac{v}{1-u} \rightarrow x, \quad \frac{w}{1-u} \rightarrow y.$$



- OPE (operator product expansion) for cusped Wilson loops

[Alday, Gaiotto, Maldacena, Sever, Vieira (2010)]

[Gaiotto, Maldacena, Sever, Vieira (2011)]

predicts (multiple) discontinuity  $\Delta_v^{L-1}$  at  $L$  loops from one-loop data

$$\Delta_v^{L-1} \mathcal{S}(R_6^{(L)}) \propto \mathcal{S} \left( \int \frac{dp}{2\pi} e^{-ip\sigma} \left( \sum_{m=1}^{\infty} \frac{[\gamma_{m+2}(p)]^{L-1} \cos(m\phi)}{p^2 + m^2} + \sum_{m=2}^{\infty} \frac{[\gamma_{m-2}(p)]^{L-1} \cos((m-2)\phi)}{p^2 + (m-2)^2} \right) C_m(p) \mathcal{F}_{m/2,p}(\tau) \right).$$

$\tau, \sigma, \phi$  related to  $u, v, w$

- we can test corollaries of this formula

$$\mathcal{D}_+ \mathcal{D}_- \Delta_v \Delta_v \mathcal{S}(R_6^{(3)}) = 0$$

and

$$\square \Delta_w \Delta_w \Delta_v \Delta_v \mathcal{S}(R_6^{(3)}) \propto \frac{w(1-u+v-w)}{(1-v)(1-w)}$$

- we find a solution consistent with OPE constraints, with 26 parameters  $\alpha_i$

$$\mathcal{S}(R_6^{(3)}) = \mathcal{S}(X) + \sum_{i=1}^{26} \alpha_i \mathcal{S}(f_i)$$

highly nontrivial that ansatz is consistent with all constraints!

- Regge limit imposes three more constraints
- additional constraint: final entry drawn only from set

$$\frac{u}{1-u}, \frac{v}{1-v}, \frac{w}{1-w}, y_u, y_v, y_w$$

can be motivated by supersymmetry  
leads to

$$\mathcal{S}(R_6^{(3)}) = \mathcal{S}(X) + \alpha_1 \mathcal{S}(f_1) + \alpha_2 \mathcal{S}(f_2)$$

[Caron-Huot (2011)]

# Prediction for multi-Regge limit

- multi-Regge limit

Analytically continue to physical branch  $u \rightarrow e^{-2\pi i} u$  and let

$$u \rightarrow 1, \quad \frac{v}{1-u} \rightarrow x, \quad \frac{w}{1-u} \rightarrow y.$$

Expect

$$R_6^{(3)} \rightarrow (2\pi i) \sum_{r=0}^2 \log^r(1-u) \left[ g_r^{(3)}(x, y) + 2\pi i h_r^{(3)}(x, y) \right]$$

- we find agreement with prediction for leading terms  $g_2^{(3)}$ ,  $h_{2,1}^{(3)}$  [Lipatov, Prygarin (2010)]
- we have **new predictions** for  $g_1^{(3)}$ ,  $g_0^{(3)}$  and  $h_0^{(3)}$

# Explicit prediction for multi-Regge limit

- explicit formula for  $g_1^{(3)}$

$$u \rightarrow 1, \quad \frac{v}{1-u} \rightarrow \frac{1}{(1+w)(1+w^*)}, \quad \frac{w}{1-u} \rightarrow \frac{ww^*}{(1+w)(1+w^*)}.$$

We find

$$\begin{aligned} g_1^{(3)}(w, w^*) &= \frac{1}{8} \left\{ \log |w|^2 \left[ \text{Li}_3 \left( \frac{w}{1+w} \right) + \text{Li}_3 \left( \frac{w^*}{1+w^*} \right) \right] \right. \\ &+ (5 \log |1+w|^2 - 2 \log |w|^2) \left[ \text{Li}_3(-w) + \text{Li}_3(-w^*) \right] \\ &- \frac{3}{2} \log |w|^2 \log \frac{|1+w|^4}{|w|^2} \left[ \text{Li}_2(-w) + \text{Li}_2(-w^*) \right] \\ &- \frac{1}{12} \log^2 |1+w|^2 \left[ \log |w|^2 (\log |w|^2 + 2 \log |1+w|^2) - 10 \log^2 \frac{|1+w|^2}{|w|^2} \right] \\ &+ \frac{1}{2} \log |w|^2 \log \frac{|1+w|^2}{|w|^2} \log(1+w) \log(1+w^*) - 2 \zeta_3 \log |1+w|^2 \left. \right\} \\ &+ \left( \frac{5}{2} + \gamma' \right) \zeta_2 g_1^{(2)}(w, w^*), \end{aligned}$$

# Summary and outlook

- **Summary**

- starting from our ansatz we fix the symbol of  $\mathcal{R}_6^{(3)}$  up to two constants

$$\mathcal{S}(\mathcal{R}_6^{(3)}) = \mathcal{S}(X) + \alpha_1 \mathcal{S}(f_1) + \alpha_2 \mathcal{S}(f_2)$$

without evaluating any loop integrals!

- prediction for multi-Regge limit, where  $\alpha_1$  and  $\alpha_2$  drop out
- new functions other than classical polylogarithms needed

- **Outlook**

- find functions  $X$  and  $f_2$  ( $f_1$  already determined)
  - constraints at function level are likely to fix  $\alpha_1$  and  $\alpha_2$ .
- analytical result for three-loop six-particle scattering process

