



Alexander von
HUMBOLDT
STIFTUNG

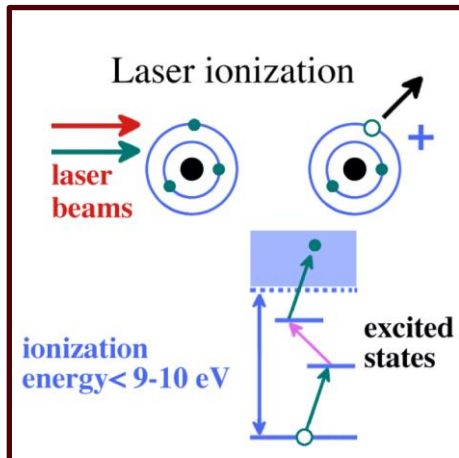


TECHNISCHE
UNIVERSITÄT
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FUTURE UPDATES AND HYPERNUCLEI PRODUCTION AT THE PUMA EXPERIMENTAL SETUP

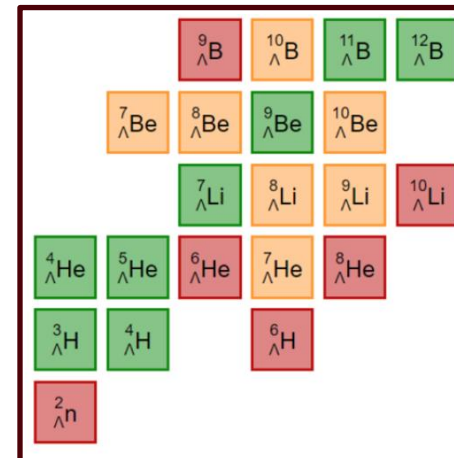
Moritz Schlaich
for the PUMA Collaboration
09.04.2024





Upgrade 1

- Upgrade ion source setup with laser ablation source
- Enable ionization and *investigation of metal ions*

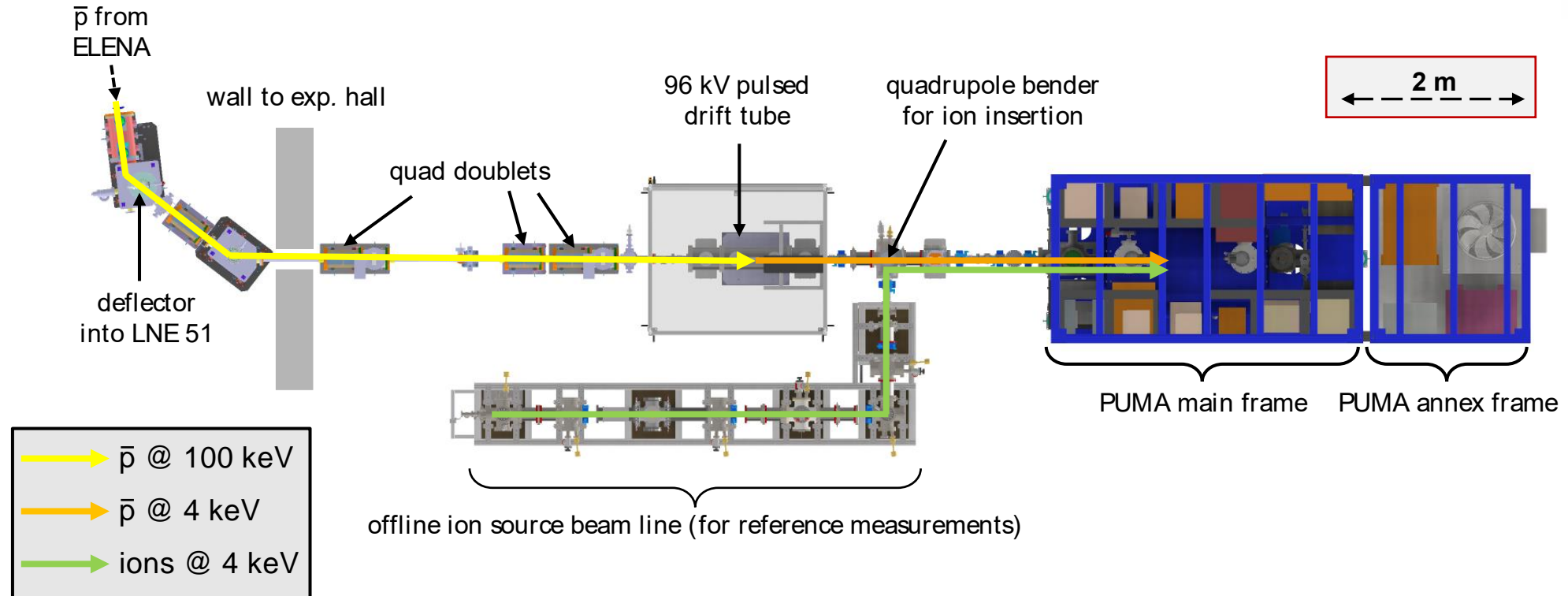


Upgrade 2

- Modify trap infrastructure with new detection system
- Use antiprotonic atoms to *produce new hypernuclei*



INTRODUCTION

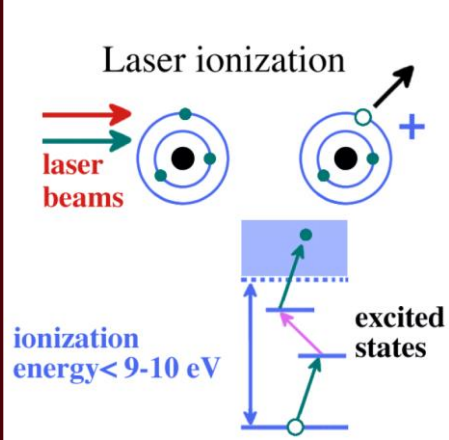


experimental campaign 2024 / 2025

- Characterization of detection system: p, d
- Study final state interactions with nucleon number: ${}^4\text{He}, {}^{16}\text{O}, {}^{20}\text{Ne}, {}^{40}\text{Ar}, {}^{132}\text{Xe}$
- Investigate isospin dependence along isotopic chains: ${}^{78-86}\text{Kr}, {}^{124-136}\text{Xe}$



Laser ionization



The diagram illustrates the laser ionization process. At the top, two atoms are shown. The first is a neutral atom with a nucleus (black dot) and two electrons (green dots). Two laser beams, represented by red and green arrows, strike it. The second atom is an ion with a nucleus and one electron, and a plus sign (+) next to it. Below this, a blue rectangular block represents a laser source. A blue arrow points from the source to the ground state of an atom. A green arrow points from the ground state to an excited state. A pink arrow points from the excited state to the ionization energy level. A blue arrow points from the ionization energy level to the ionization energy $< 9-10 \text{ eV}$.

ionization energy $< 9-10 \text{ eV}$

excited states

Upgrade 1

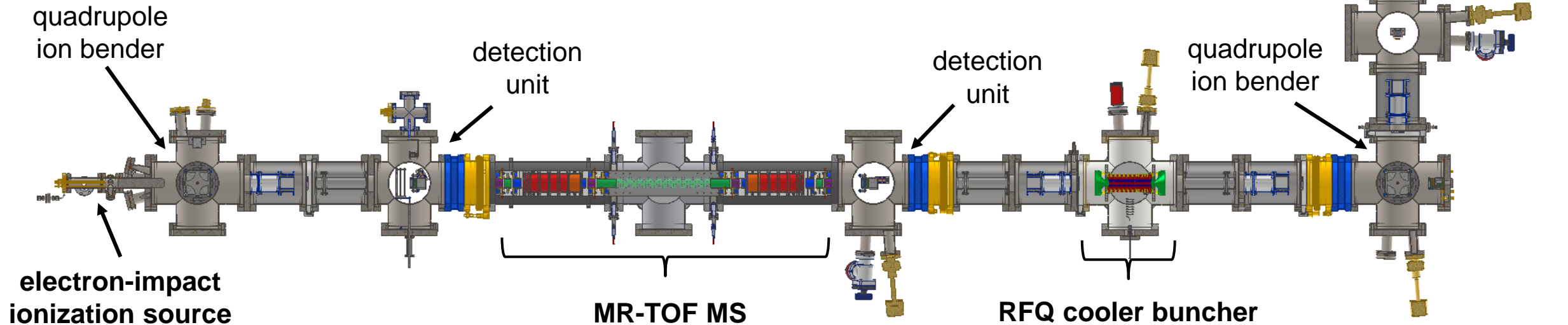
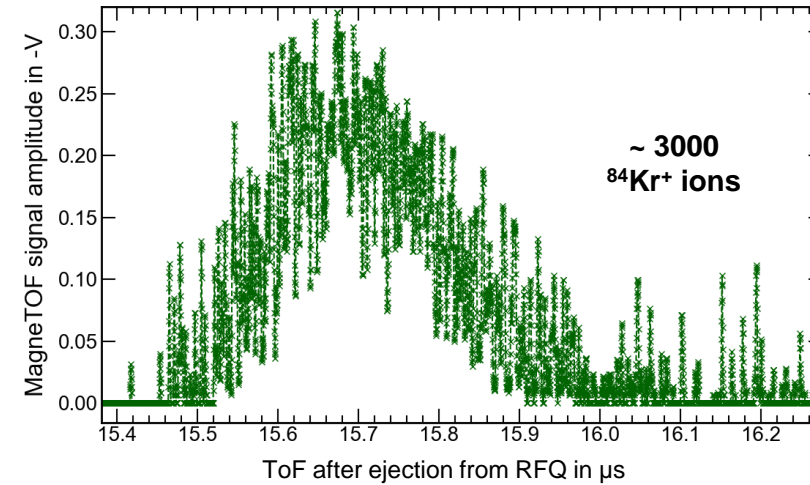
- Upgrade ion source setup with laser ablation source
- Enable ionization and *investigation of metal ions*



LASER ABLATION SOURCE – INTRODUCTION

PUMA ion source setup

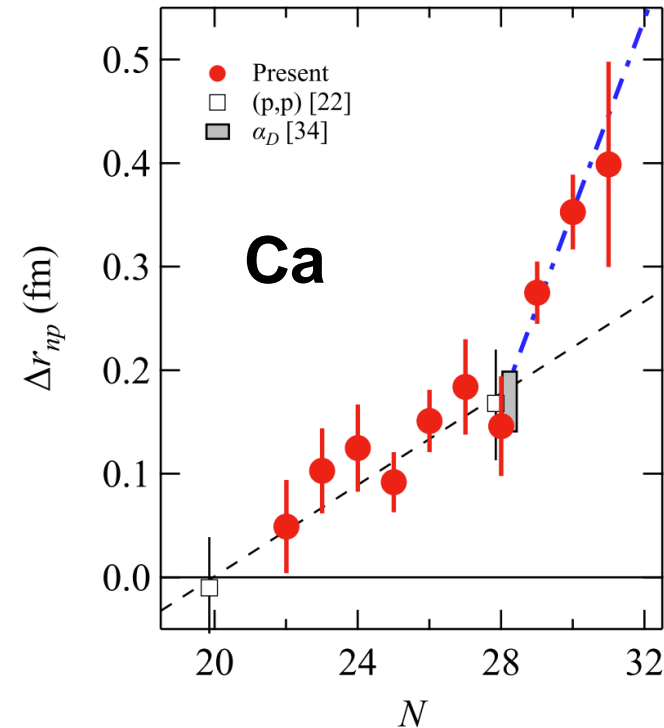
- Produces isotopically pure ion bunches
- Provides up to 10^4 ions / bunch
- Current setup lacks versatility
→ *Metal ions cannot be produced*



LASER ABLATION SOURCE – MOTIVATION



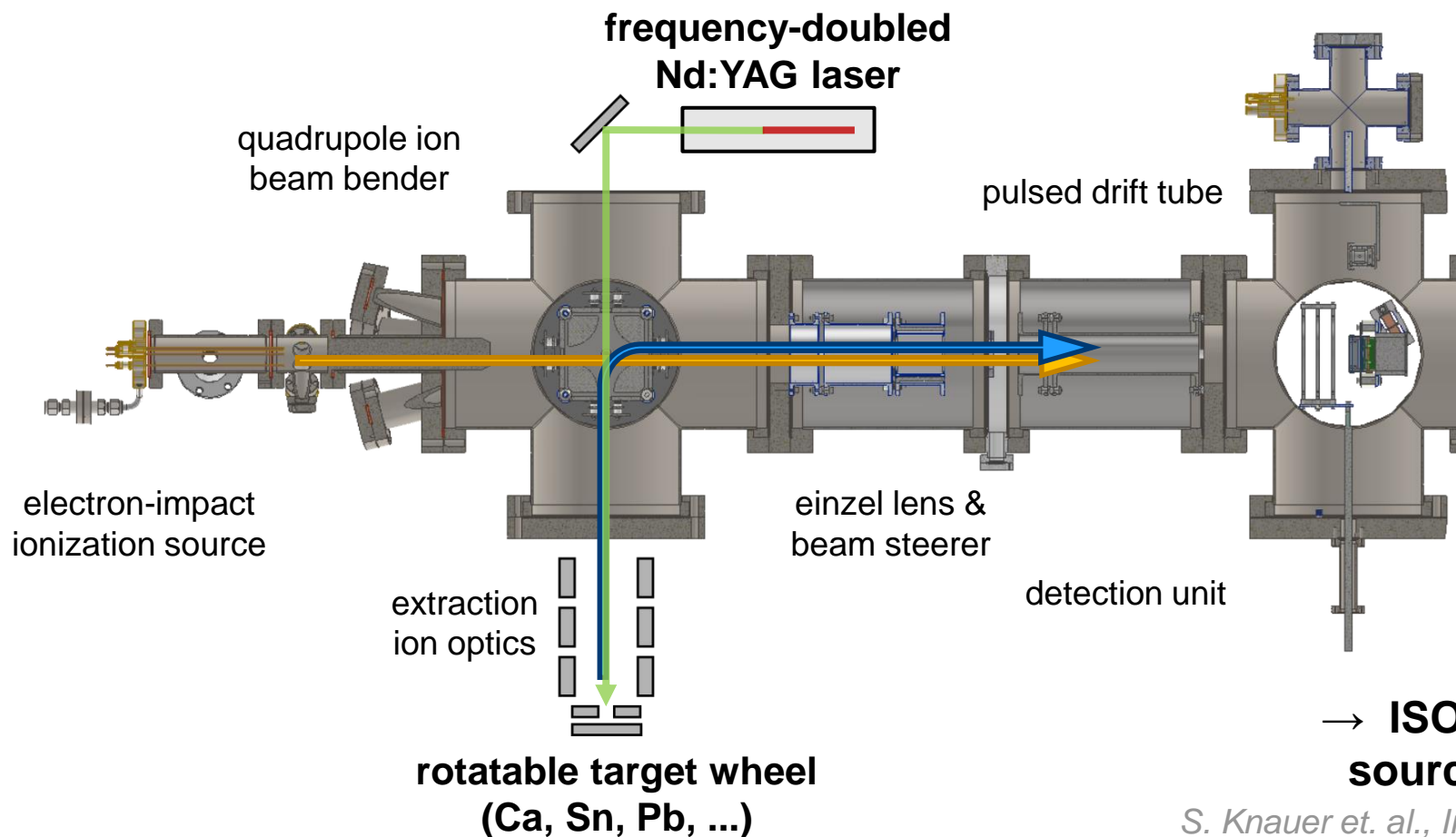
- Equation of state (EoS) of infinite nuclear matter: $\frac{E}{A}(\rho, \delta) = \frac{E}{A}(\rho, \delta = 0) + S(\rho) \cdot \delta^2 + \dots$
 with $S(\rho) = S(\rho_0) + L \cdot \left(\frac{\rho - \rho_0}{3\rho_0}\right) + \dots$
 and $\rho = \rho_n + \rho_p, \quad \delta = \rho_n - \rho_p$
- Neutron skin thickness correlated with slope parameter L of EoS
- Medium-mass closed shell nuclei with neutron excess ideal to study neutron skins
 → $^{48}\text{Ca}, ^{208}\text{Pb}$
- Study proton-closed shell isotopic chains
 → $^{40-48}\text{Ca}, ^{112-124}\text{Sn}$



M. Tanaka et. al., *PRL* **124**, 102501 (2020)



LASER ABLATION SOURCE



laser properties

- 532 nm wavelength
- 10 Hz repetition rate
- 330 mJ per pulse
- 1 mm² spot size

→ ISOLTRAP's laser-ablation
source serves as reference

S. Knauer et. al., Int. J. Mass Spectrom. 446, 116189 (2019)



		${}^9_{\Lambda}\text{B}$	${}^{10}_{\Lambda}\text{B}$	${}^{11}_{\Lambda}\text{B}$	${}^{12}_{\Lambda}\text{B}$	
	${}^7_{\Lambda}\text{Be}$	${}^8_{\Lambda}\text{Be}$	${}^9_{\Lambda}\text{Be}$	${}^{10}_{\Lambda}\text{Be}$		
		${}^7_{\Lambda}\text{Li}$	${}^8_{\Lambda}\text{Li}$	${}^9_{\Lambda}\text{Li}$	${}^{10}_{\Lambda}\text{Li}$	
${}^4_{\Lambda}\text{He}$	${}^5_{\Lambda}\text{He}$	${}^6_{\Lambda}\text{He}$	${}^7_{\Lambda}\text{He}$	${}^8_{\Lambda}\text{He}$		
${}^3_{\Lambda}\text{H}$	${}^4_{\Lambda}\text{H}$		${}^6_{\Lambda}\text{H}$			
${}^2_{\Lambda}\text{n}$						

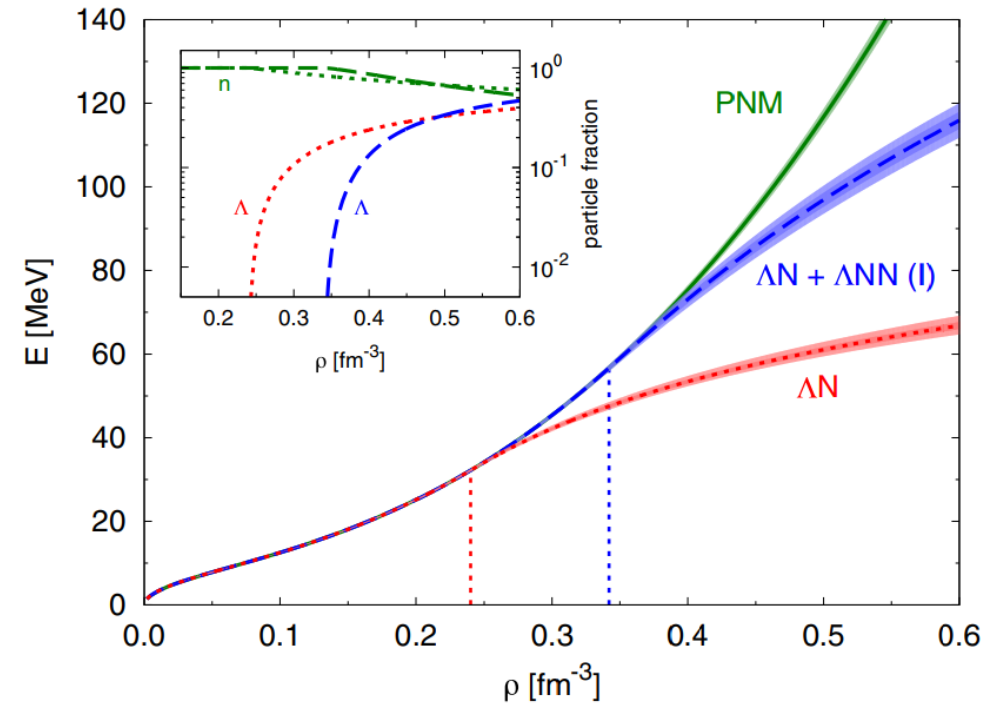
Upgrade 2

- Modify trap infrastructure with new detection system
- Use antiprotonic atoms to *produce new hypernuclei*



HYPERNUCLEI – MOTIVATION

- Strangeness in modern nuclear physics largely unexplored
 - *~ 500,000 bound Λ hypernuclei ($\Lambda < 70, Z < 120$) predicted*
E. Khan et al., Phys. Rev. C 92, 044313 (2015)
 - but only *~ 40* observed so far
- How does strangeness content evolve with baryonic density?
 - *Investigate influence on the nuclear EoS?*
- Increase experimental activities involving hypernuclei
- Antiprotonic atoms as a tool for hypernuclei production
 - *Use PUMA infrastructure*
 - *Redesign trap and detection system*



D. Lonardoni et al., *Phys. Rev. Lett.* **114**, 092301 (2015)

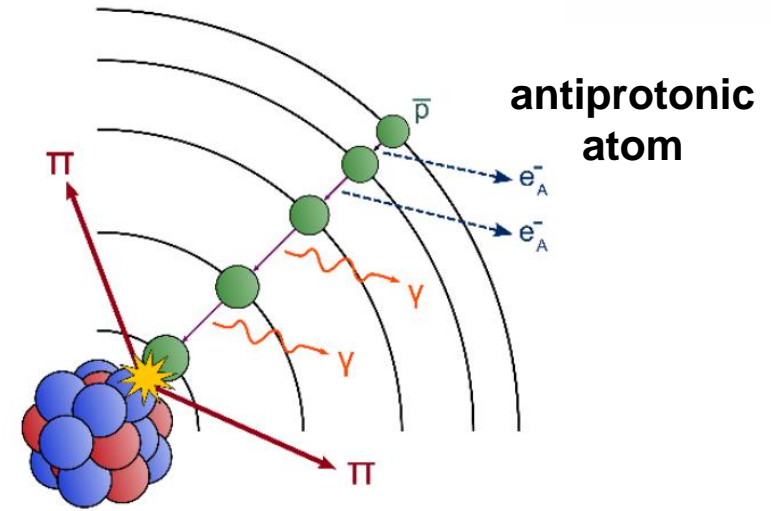


HYPERNUCLEI – PRODUCTION MECHANISM

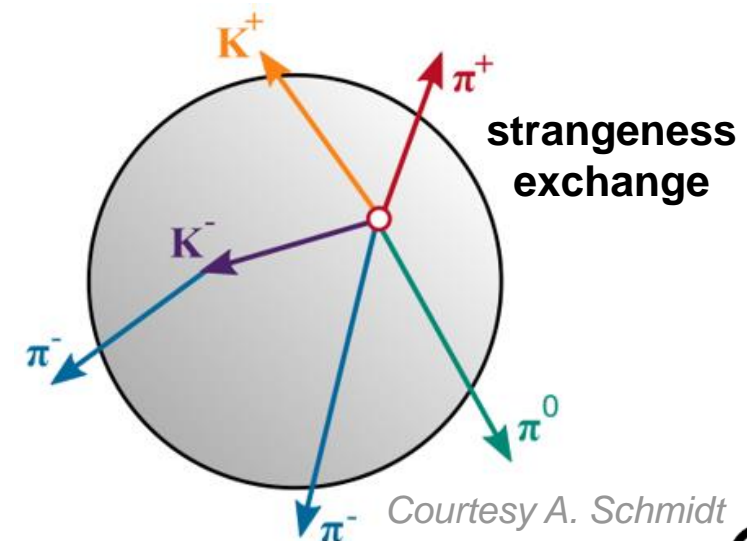
- Antiprotonic atom formed with target nucleus
- Meson emission after annihilation (dominated by pions)
- Kaons emitted in ~ 5% of all annihilations
- Strangeness production allows formation of Λ particles

two routes for hyperon generation

- Final state interactions of K^- or K^0 with nucleons
 $K^- n \rightarrow \Lambda \pi^-$, $K^- p \rightarrow \Lambda \pi^0$
 \rightarrow *Strangeness exchange*
- Final state interactions of pions with nucleons
 $\pi n \rightarrow \Lambda K \pi$, $\pi p \rightarrow \Lambda K \pi$
 \rightarrow *Strangeness pair production*



Courtesy A. Schmidt

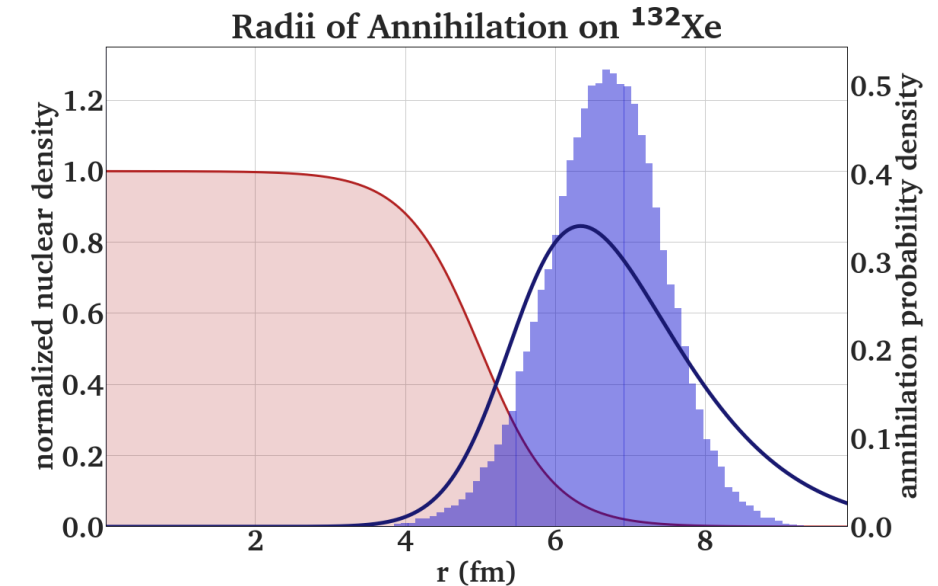


Courtesy A. Schmidt



HYPERNUCLEI – SIMULATIONS

- **1. STEP:** Simulation with GiBUU (Gießen Boltzmann-Uehling-Uhlenbeck) code
 - Antiprotons ($E_{\text{kin}} = 10$ eV) collide with nucleus at rest
 - Annihilation profile obtained from data on antiprotonic atoms
 - Propagation of annihilation products & interaction with residual nucleus → Λ production
 - Definition of excited hypernuclei based on phase space coalescence: $\rho_B > 0.01 \rho_0$

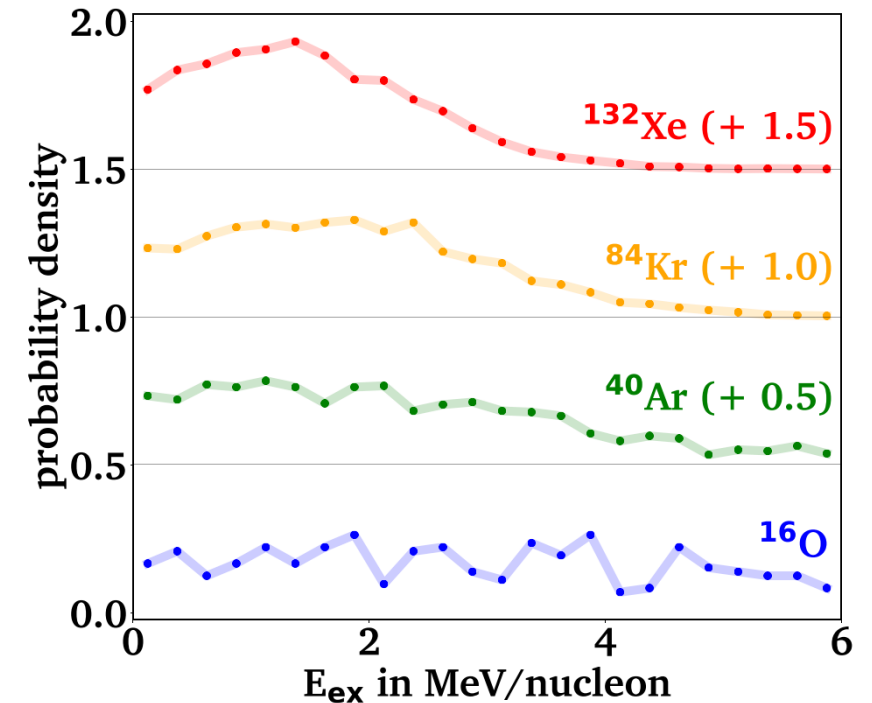


A. Schmidt et al., *Eur. Phys. J. A* **60**, 55 (2024)



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- **2. STEP:** Deexcitation of hypernuclei simulated in ABLA++
 - Typical excitation energies: 0 – 6 MeV / nucleon
 - Deexcitation via fission or evaporation of n, p, light particles
 - Derivation of yields of different hypernuclei



A. Schmidt et al., *Eur. Phys. J. A* **60**, 55 (2024)



HYPERNUCLEI – PRODUCTION RATES

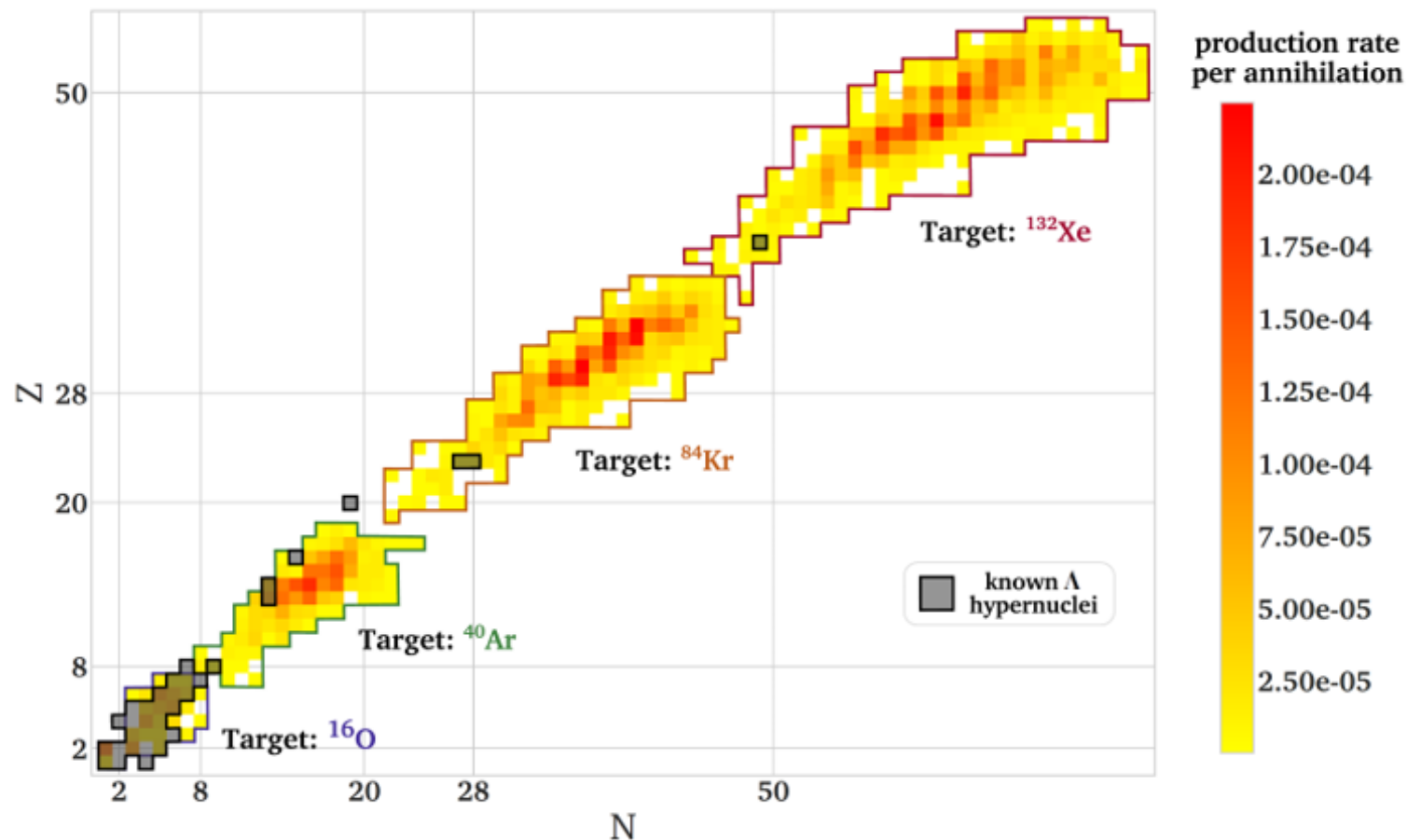


production rate per annihilation

^{16}O		^{40}Ar	
nucleus	p in 10^{-5}	nucleus	p in 10^{-5}
$^{10}_{\Lambda}\text{B}$	9 ± 2	$^{29}_{\Lambda}\text{Al}$	19 ± 3
$^{7}_{\Lambda}\text{Li}$	8 ± 2	$^{32}_{\Lambda}\text{Si}$	15 ± 3
$^{10}_{\Lambda}\text{Be}$	7 ± 2	$^{28}_{\Lambda}\text{Al}$	15 ± 3
$^{11}_{\Lambda}\text{B}$	6 ± 2	$^{31}_{\Lambda}\text{Al}$	13 ± 2
$^{9}_{\Lambda}\text{Be}$	5 ± 2	$^{28}_{\Lambda}\text{Mg}$	13 ± 2

^{84}Kr		^{132}Xe	
nucleus	p in 10^{-5}	nucleus	p in 10^{-5}
$^{65}_{\Lambda}\text{Cu}$	28 ± 3	$^{109}_{\Lambda}\text{Ag}$	21 ± 3
$^{72}_{\Lambda}\text{Ge}$	24 ± 3	$^{113}_{\Lambda}\text{In}$	19 ± 3
$^{71}_{\Lambda}\text{Ga}$	21 ± 3	$^{104}_{\Lambda}\text{Pd}$	18 ± 3
$^{69}_{\Lambda}\text{Ga}$	20 ± 3	$^{106}_{\Lambda}\text{Pd}$	16 ± 3
$^{62}_{\Lambda}\text{Ni}$	19 ± 3	$^{107}_{\Lambda}\text{Ag}$	16 ± 3

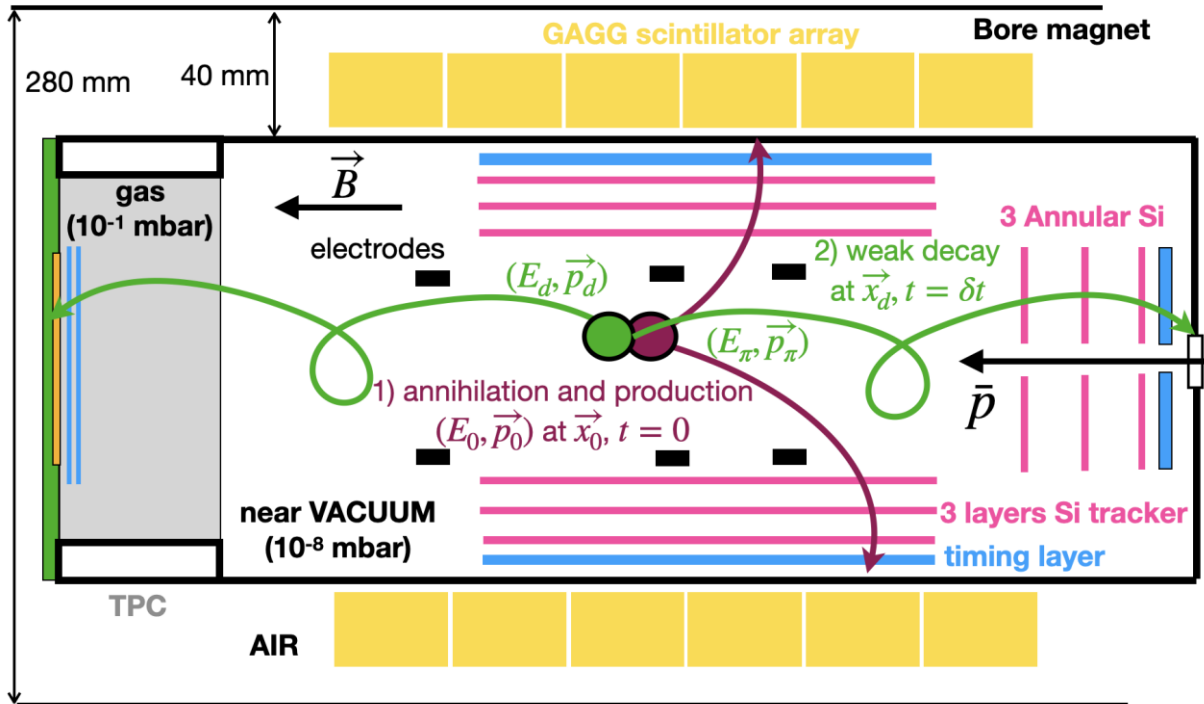
A. Schmidt et al., Eur. Phys. J. A **60**, 55 (2024)



A. Schmidt et al., Eur. Phys. J. A **60**, 55 (2024)



HYPERNUCLEI – EXPERIMENTAL SETUP



Courtesy A. Obertelli

Residual nucleus identification → TPC
 Weak decay pion identification → silicon trackers
 Hypernuclei identification → invariant mass method

- Antiprotons decelerated to 50-80 keV with PDT and focused onto separation window (~ 10 mm)
- Antiprotons trapped in Penning trap to form antiprotonic atoms with target gas atoms
- Tune annihilation rate with target gas pressure
 → Approx. 10^{-8} mbar
- Hypernuclei production tagged by time difference between annihilation and weak decay (~ 200 ps)

determine invariant mass of hypernucleus

$$M_Y^2 = m_\pi^2 + M^2 + 2\sqrt{m_\pi^2 + p_\pi^2 c^2} \sqrt{M^2 + p_r^2 c^2} - 2\vec{p}_\pi \cdot 2\vec{p}_r$$



HYPERNUCLEI – EFFICIENCY ESTIMATION



- Antiprotons per bunch: $N_{\bar{p}} = 3 \cdot 10^6$ (every two minutes)
- Number of bunches: $N_{\text{spill}} = 7800$ (15 days beam time)
- Trapping efficiency: $x_{\text{cap}} = 30 \%$ (conservative assumption)
- Y production rate: $f(\Lambda Y) = 0.01 \%$ (based on simulations)
- Annihilation detection efficiency: $\epsilon_a = 78 \%$ (based on simulations)
- Mesonic decay branching ratio: $BR = 5 \%$ (unfavorable case)
- Recoil fragment reaching TPC: $\epsilon_f = 50 \%$ (geometrical acceptance)
- Weak-decay-event timing selection: $\epsilon_\omega = 50 \%$ (conservative assumption)
- Recoil fragment entering TPC: $\epsilon_{\text{TPC}} = 60 \%$ (conservative assumption)
- Tracking & identification efficiency: $\epsilon_{\text{track}} = 60 \%$ (conservative assumption)

$$N(\Lambda Y) = N_{\bar{p}} \cdot N_{\text{spill}} \cdot x_{\text{cap}} \cdot f(\Lambda Y) \cdot \epsilon_a \cdot BR \cdot \epsilon_f \cdot \epsilon_\omega \cdot \epsilon_{\text{TPC}} \cdot \epsilon_{\text{track}} = \mathbf{1500 \text{ identified hypernuclei}}$$



SUMMARY



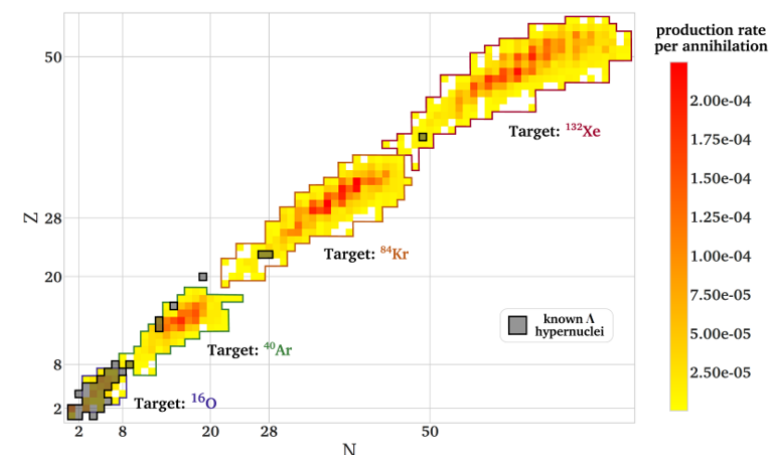
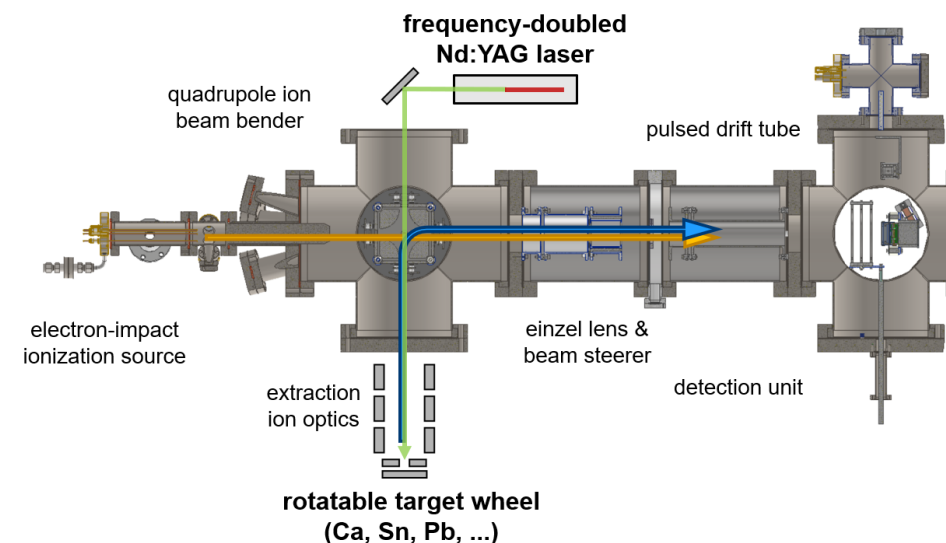
- Two upgrades for PUMA experiment planned

Upgrade 1

- Develop laser ablation ion source
 - Allow production of metal ions
- Study medium-mass closed shell nuclei
- *Commissioning planned for CERN's Long Shutdown 3*

Upgrade 2

- Develop new detection system and modify Penning trap
 - Produce and detect new hypernuclei
- Simulations suggest accessibility of broad hyperisotope range
- *Experimental activity planned in post-PUMA future (2027-2028?)*



THE PUMA COLLABORATION

T. Aumann, N. Azaryan, W. Bartmann, A. Bouvard, O. Boine-Frankenheim, A. Broche, F. Butin, D. Calvet, J. Carbonell, P. Chiggiato, H. De Gersem, R. De Oliveira, T. Dobers, F. Ehm, J. Ferreira Somoza, J. Fischer, M. Fraser, E. Friedrich, M. Gomez-Ramos, J.-L. Grenard, G. Hupin, K. Johnston, C. Klink, M. Kowalska, Y. Kubota, P. Indelicato, R. Lazauskas, S. Malbrunot-Ettenauer, N. Marsic, W. Müller, S. Naimi, N. Nakatsuka, R. Necca, D. Neidherr, A. Obertelli, Y. Ono, S. Pasinelli, N. Paul, E. C. Pollacco, L. Riik, D. Rossi, H. Scheit, M. Schlaich, R. Seki, A. Schmidt, L. Schweikhard, S. Sels, E. Siesling, T. Uesaka, M. Wada, F. Wienholtz, S. Wycech, C. Xanthopoulou, S. Zacarias



THANK YOU!



APPENDIX – BRANCHING RATIOS

Antiproton - Proton	
Final State	Probability in %
$\rho^+ \rho^-$	3.37
$\pi^+ \pi^- \pi^0$	2.34
$\pi^+ \pi^- \rho^0$	2.02
$\pi^+ \pi^0 \rho^-$	2.02
$\pi^- \pi^0 \rho^+$	2.02
$\pi^+ \pi^- \omega$	3.03
$\pi^+ \pi^- \pi^0 \omega$	2.84
$\pi^+ \pi^+ \pi^- \pi^-$	2.74
$\pi^+ \pi^- \pi^0 \pi^0$	3.89
$\pi^+ \pi^+ \pi^- \rho^-$	2.58
$\pi^+ \pi^- \pi^- \rho^+$	2.58
$\pi^+ \pi^- \pi^0 \rho^0$	6.29
$\pi^+ \pi^0 \pi^0 \rho^-$	5.05
$\pi^- \pi^0 \pi^0 \rho^+$	5.05

Antiproton - Proton	
Final State	Probability in %
$\pi^+ \pi^+ \pi^- \pi^- \pi^0$	2.61
$\pi^+ \pi^- \pi^0 \pi^0 \omega$	2.58
$\pi^+ \pi^+ \pi^+ \pi^- \pi^- \pi^-$	2.83
$\pi^+ \pi^+ \pi^- \pi^- \pi^0 \pi^0$	9.76
$\pi^+ \pi^- \pi^0 \pi^0 \pi^0 \pi^0$	2.68
$K^{*+} K^{*-}$	0.225
$K^{*0} \bar{K}^{*0}$	0.225
$K^0 \bar{K}^0 \omega$	0.232
$K^+ K^- \omega$	0.232
$K^0 \bar{K}^0 \rho^0$	0.202
$K^+ K^- \rho^0$	0.202
$K^0 K^- \rho^+$	0.234
$\bar{K}^0 K^+ \rho^-$	0.234
$K^{*+} \bar{K}^0 \pi^-$	0.23
$K^{*-} K^0 \pi^+$	0.23

A. Schmidt et al., Eur. Phys. J. A **60**, 55 (2024)

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APPENDIX – BRANCHING RATIOS

Antiproton - Neutron	
Final State	Probability in %
$\rho^- \rho^0$	3.51
$\rho^- \eta$	2.27
$\rho^- \omega$	3.51
$\pi^+ \pi^- \pi^-$	2.86
$\pi^+ \pi^- \rho^-$	3.62
$\pi^- \pi^0 \rho^0$	5.61
$\pi^0 \pi^0 \rho^-$	3.51
$\pi^- \rho^+ \rho^-$	2.09
$\pi^- \pi^0 \omega$	5.05
$\pi^+ \pi^- \pi^- \omega$	10.52
$\pi^+ \pi^- \pi^- \pi^0$	5.51

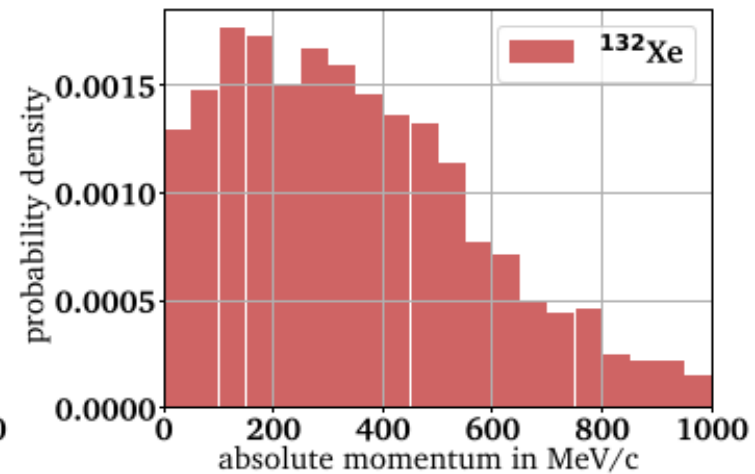
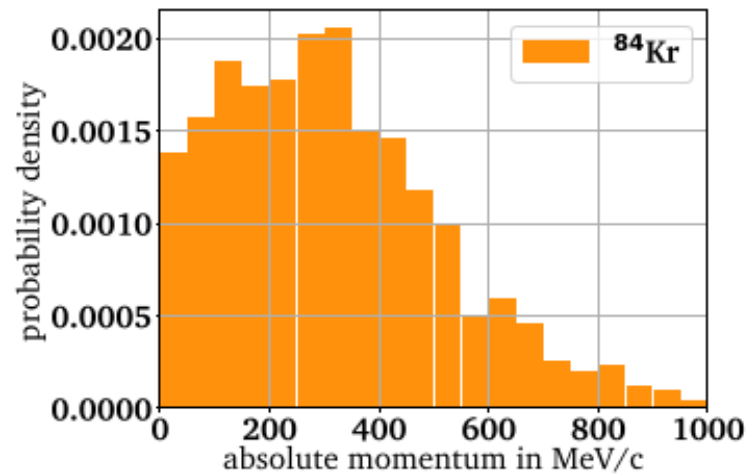
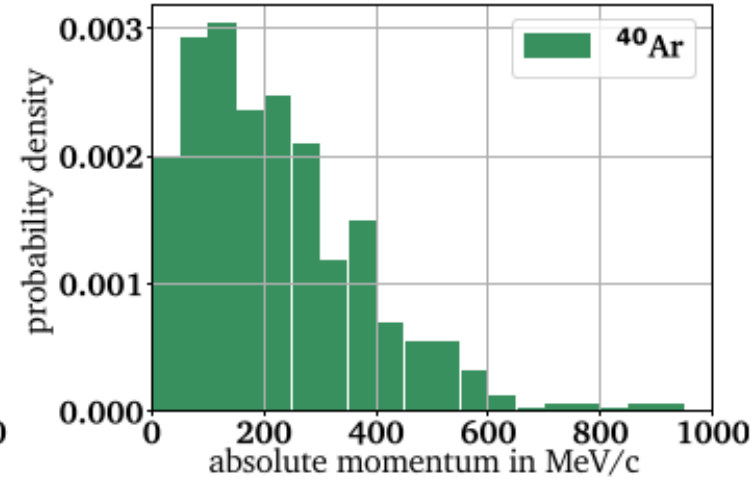
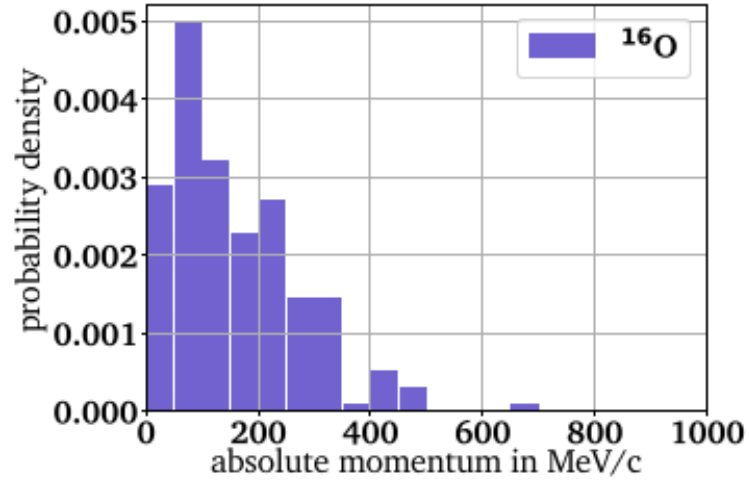
Antiproton - Neutron	
Final State	Probability in %
$\pi^+ \pi^- \pi^- \pi^0 \pi^0$	2.72
$\pi^+ \pi^+ \pi^- \pi^- \pi^- \pi^0$	8.33
$\pi^+ \pi^- \pi^- \pi^0 \pi^0 \pi^0$	6.67
$K^0 K^- \pi^0$	0.316
$K^0 \bar{K}^0 \pi^-$	0.432
$K^+ K^- \pi^-$	0.513
$K^0 K^- \omega$	0.35
$K^0 \bar{K}^0 \rho^-$	0.77
$K^+ K^- \rho^-$	0.77
$K^{*-} K^0 \pi^0$	0.245
$K^{*0} K^- \pi^0$	0.245

A. Schmidt et al., *Eur. Phys. J. A* **60**, 55 (2024)

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APPENDIX – MOMENTUM DISTRIBUTIONS



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