

Hidden Simplicity of Cosmological Correlators

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Based on:

2312.13803 with C. Chowdhury, J. Mei, I. Sachs, P. Vanhove

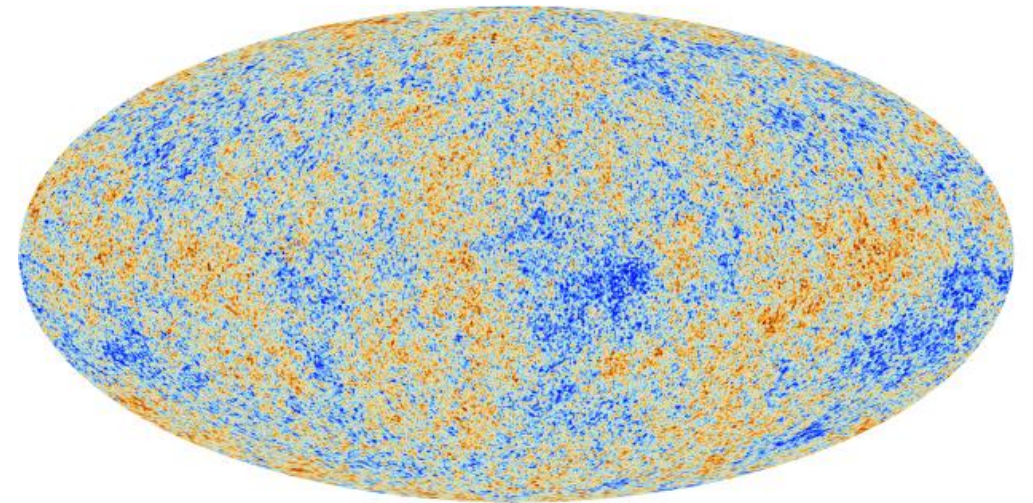
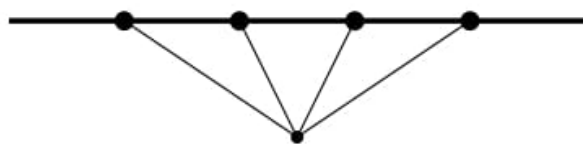
24xx.xxxxx with C. Chowdhury, J. Mei, J. Marshall, I. Sachs

Cosmological Observables

- Inflation: early Universe approximately described by dS_4 . CMB comes from correlations on future boundary

$$ds^2 = \frac{-d\eta^2 + (dx^i)^2}{\eta^2}, \quad -\infty < \eta < 0$$

$\eta=0$



Cosmological Wavefunction

- In-in correlators (Maldacena, Weinberg):

$$\langle \phi(\vec{k}_1) \dots \phi(\vec{k}_n) \rangle = \frac{\int \mathcal{D}\phi \phi(\vec{k}_1) \dots \phi(\vec{k}_n) |\Psi[\phi]|^2}{\int \mathcal{D}\phi |\Psi[\phi]|^2}$$

- Wavefunction:

$$\ln \Psi[\phi] = - \sum_{n=2}^{\infty} \frac{1}{n!} \int \prod_{i=1}^n \frac{d^d k_i}{(2\pi)^d} \psi_n(\vec{k}_1, \dots, \vec{k}_n) \phi(\vec{k}_1) \dots \phi(\vec{k}_n)$$

- ψ_n can be computed from analytically continued Feynman diagrams in EAdS (Maldacena, Pimentel, McFadden, Skenderis)

AdS Momentum Space

- Poincaré patch of Euclidean AdS₄ with unit radius:

$$ds^2 = \frac{dz^2 + (dx^i)^2}{z^2}$$

- Wick rotation $z = \pm i\eta$ gives dS₄ metric

- Boundary momenta: $\sum_{i=1}^n \vec{k}_i = 0$

- Energy: $k_{12\dots n} = \sum_{i=1}^n k_i, \quad k_i = |\vec{k}_i|$

Shadow Formalism

- In-in correlators can be computed from Witten diagrams in EAdS using the following Lagrangian:

$$iS_c = \int_0^\infty \frac{dz d^d x}{z^{d+1}} \left[\sin \left(\pi \left(\Delta_+ - \frac{d}{2} \right) \right) \left((\partial \phi_+)^2 - m^2 \phi_+^2 \right) \right. \\ \left. + \sin \left(\pi \left(\Delta_- - \frac{d}{2} \right) \right) \left((\partial \phi_-)^2 - m^2 \phi_-^2 \right) \right. \\ \left. + e^{i\pi \frac{d-1}{2}} V \left(e^{-i\frac{\pi}{2} \Delta_+} \phi_+ + e^{-i\frac{\pi}{2} \Delta_-} \phi_- \right) + e^{-i\pi \frac{d-1}{2}} V \left(e^{i\frac{\pi}{2} \Delta_+} \phi_+ + e^{i\frac{\pi}{2} \Delta_-} \phi_- \right) \right]$$

(Sleight, Taronna, Di Pietro, Gorbenko, Komatsu)

- ϕ_+ is a ghost and only appears on internal lines

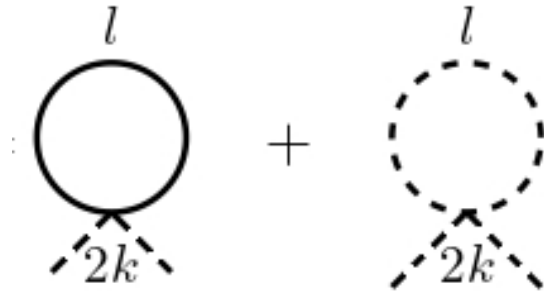
Conformally coupled ϕ^4

- Choose $d=3$ and $\Delta_- = d - \Delta_+ = 1$. $V(\phi_+, \phi_-) = \frac{1}{2} \frac{\lambda}{4!} (\phi_+^4 - 6\phi_+^2\phi_-^2 + \phi_-^4)$
- Feynman rules: (ϕ_- dotted, ϕ_+ solid)

$$\begin{array}{c} \vec{k} \\ \text{---} z_1 \quad \text{---} z_2 \end{array} = \frac{1}{\pi} \int_{-\infty}^{\infty} \frac{dp}{p^2 + k^2} \cos(pz_1) \cos(pz_2) \qquad \begin{array}{c} \vec{k} \\ \text{---} z_1 \quad \text{---} z_2 \end{array} = -\frac{1}{\pi} \int_{-\infty}^{\infty} \frac{dp}{p^2 + \vec{k}^2} \sin(pz_1) \sin(pz_2)$$

$$\begin{array}{c} \vec{k} \\ \text{---} \end{array} = e^{-kz} \qquad \begin{array}{c} \text{---} \\ \text{---} \end{array} = \frac{\lambda}{2}, \qquad \begin{array}{c} \text{---} \\ \text{---} \end{array} = -3\lambda, \qquad \begin{array}{c} \text{---} \\ \text{---} \end{array} = \frac{\lambda}{2}$$

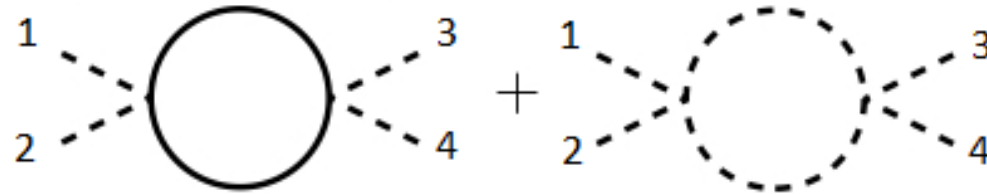
1-loop 2-point



$$\begin{aligned} &= \lambda \int \frac{d^3 l dp}{(2\pi)^4} \frac{1}{p^2 + l^2} \int_0^\infty dz e^{-2kz} (\sin^2(pz) + \cos^2(pz)) \\ &= \frac{\lambda}{k} \int \frac{d^4 L}{(2\pi)^4} \frac{1}{L^2}, \quad L^\mu = (p, \vec{l}) \end{aligned}$$

- 4d Lorentz invariance is restored
- Doesn't happen for wavefunction

1-loop 4-point



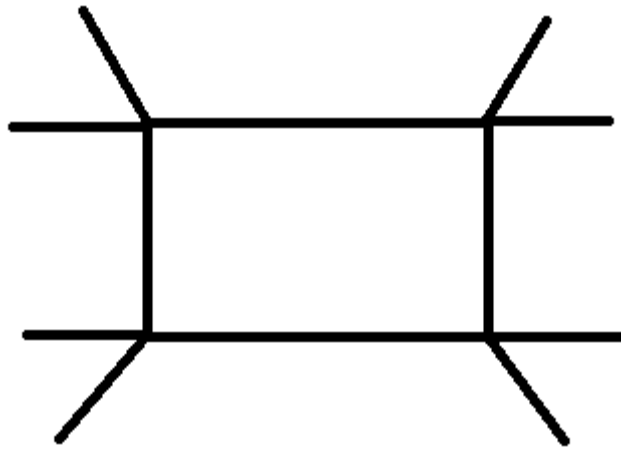
$$= \left(\frac{\lambda}{2}\right)^2 k_{12} k_{34} \int \frac{dp}{(2\pi)^5} \frac{1}{(p^2 + k_{12}^2)(p^2 + k_{34}^2)} \int \frac{d^4 L}{L^2(L-P)^2}, \quad P^\mu = (p, \vec{k}_{12})$$

$$= \frac{1}{(2\pi)^2} \frac{\lambda^2}{32(k_{12} + k_{34})} \left[\ln \left(\frac{(k_{12} + |\vec{k}_{12}|)(k_{34} + |\vec{k}_{12}|)}{4\Lambda^2} \right) + \frac{k_{12} + k_{34}}{k_{12} - k_{34}} \ln \left(\frac{k_{34} + |\vec{k}_{12}|}{k_{12} + |\vec{k}_{12}|} \right) \right], \quad \vec{k}_{ij} = \vec{k}_i + \vec{k}_j$$

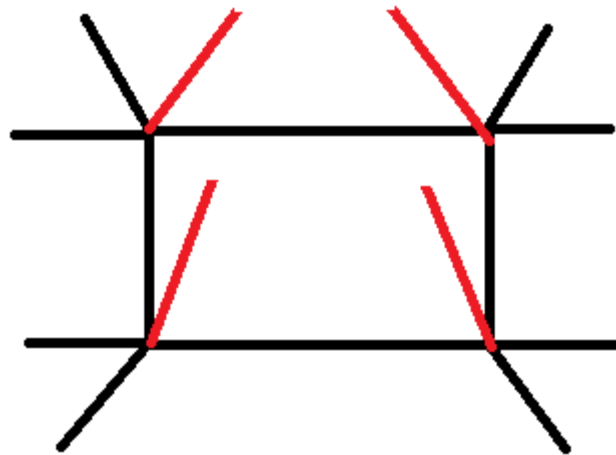
- Same analytic structure as scattering amplitude
- In contrast, wavefunction is a dilog ([Albayrak, Chowdhury, Kharel](#))

Question: How do we make the simplicity of in-in correlators manifest?

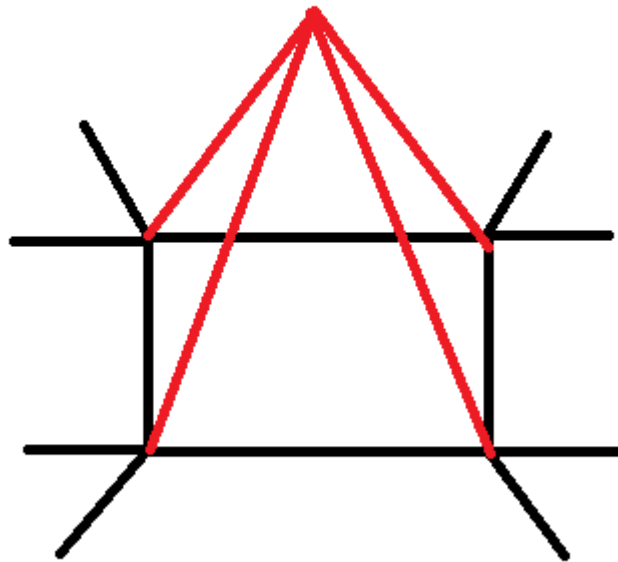
- Start with a flat space Feynman diagram



- Start with a flat space Feynman diagram
- Add an auxiliary propagator at each vertex

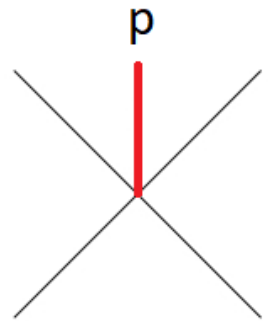


- Start with a flat space Feynman diagram
- Add auxiliary propagator at each vertex
- Connect auxiliary propagators to a common point
- Impose internal energy conservation at each vertex
- Integrate over auxiliary energies



Dressing Rules

- For conformally coupled ϕ^4 , the auxiliary propagator is:

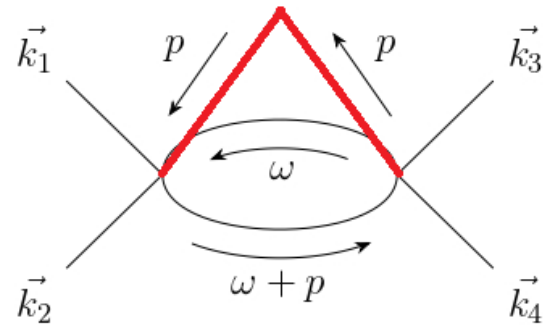


A Feynman diagram showing a central vertex with four external legs. One leg is highlighted in red and labeled 'p'. To the right of the diagram is the mathematical expression for the auxiliary propagator:

$$\frac{k_{ext}}{p^2 + k_{ext}^2}$$

- k_{ext} = total energy of external legs attached to vertex
- p = total energy of internal legs attached to vertex
- Similar rules can be derived for conformally coupled ϕ^3 , massless ϕ^4 , and massless ϕ^3

Example:



$$\begin{aligned}
 &= \left(\frac{\lambda}{2}\right)^2 \int \frac{1}{(2\pi)^5} dp d\omega d^3l \left(\frac{k_{12}}{p^2 + k_{12}^2}\right) \left(\frac{k_{34}}{p^2 + k_{34}^2}\right) \left(\frac{1}{\omega^2 + \vec{l}^2}\right) \left(\frac{1}{(\omega + p)^2 + (\vec{l} + \vec{k}_{12})^2}\right) \\
 &= \left(\frac{\lambda}{2}\right)^2 k_{12} k_{34} \int \frac{dp}{(2\pi)^5} \frac{1}{(p^2 + k_{12}^2)(p^2 + k_{34}^2)} \int \frac{d^4L}{L^2(L - P)^2},
 \end{aligned}$$

Conclusion and Future

- In-in correlators are simpler than wavefunctions
- Can be obtained by dressing flat space Feynman diagrams
- Discontinuities? (work in progress with [Chowdhury, Jazayeri, Mei, Marshall, Sachs](#))
- Renormalisation? (see recent work by [Bzowski, McFadden, Skenderis](#))
- General backgrounds?