

Bulk quantum corrections for non-spatial holographic entanglement

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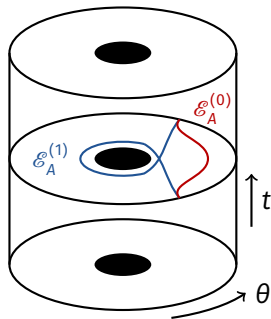
Vrije Universiteit Brussel

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Joint work with Dongming He to appear soon

Introduction

What is the field theory dual to the area of extremal surfaces \mathcal{E}_A that are topologically distinct from RT surfaces?

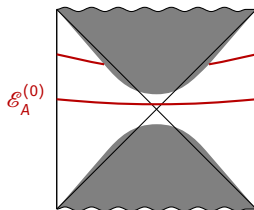
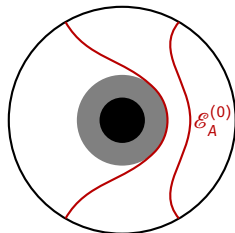


$$\frac{\text{Area}(\mathcal{E}_A^{(w)})}{4G_N} \stackrel{?}{=} S_A^{(w)} = -\text{Tr}(\rho_A^{(w)} \log \rho_A^{(w)})$$

→ Entanglement between different fields as well as between spatial DoF

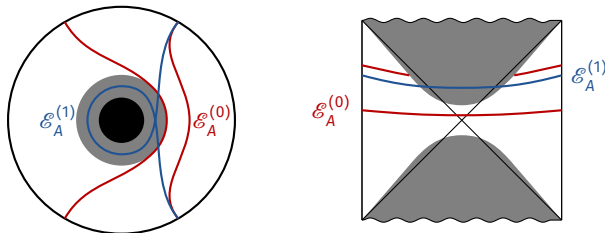
Motivation

- “entanglement builds geometry” [Swingle (2009), Van Raamsdonk (2010)]: bulk geometry encoded in terms of entanglement
- entanglement shadows [Freivogel, Jefferson, Kabir, Mosk, Yang (2014)]: RT surfaces don't probe all of the spacetime



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- Non-minimal extremal surfaces can probe closer to singularities/horizons
 - resolve entanglement shadows
- Expect large quantum corrections close to singularities
 - breakdown of “entanglement builds geometry”?

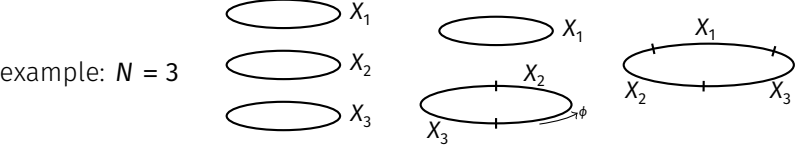
Entwinement: definition

[Balasubramanian, Chowdhury, Czech, de Boer (2014); Balasubramanian, Bernamonti, Craps, de Jonckheere, Galli (2016); MG (2021)]

- bottom-up AdS_3/CFT_2 setup: S_N orbifold CFT of generic seed CFT (e.g. D1/D5 system or $AdS_3 \times X$ with pure NS-NS flux [Eberhardt (2021); Knighton, Sriprachyakul (2024)])

$$S[X] = \sum_{i=1}^N S_{\text{seed}}[X_i] + \text{marginal deformations}$$

- field content: N indistinguishable copies of seed fields X_i ;
- Hilbert space contains states with long strands joining together multiple fields



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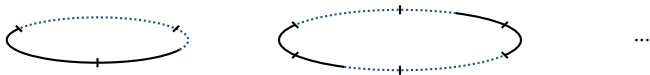
[Balasubramanian, Chowdhury, Czech, de Boer (2014); Balasubramanian, Bernamonti, Craps, de Jonckheere, Galli (2016); MG (2021)]

- ordinary entanglement entropy: consider one or multiple intervals of length L on long strands of arbitrary size



- entwinement: subsystem parametrized by two integers $w < m$ and interval length L

→ consider one or multiple intervals of length $w + L$ on long strands of length $m\mathbb{Z}$



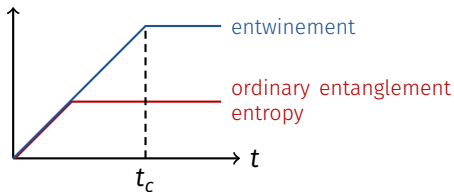
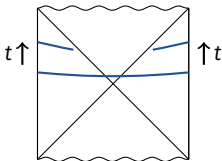
Entwinement at large N : bulk dual

- Explicit computations using large N techniques give

$$S_A^{(w,m)} = \frac{1}{m} \frac{\text{Length}[\mathcal{E}_A^{(w)}]}{4G_N} + O(1)$$

for conical defects [Balasubramanian, Chowdhury, Czech, de Boer (2014)]
and BTZ black holes [MG (2021)]

- Two-sided black hole: phase transition at $t_c \sim w$



- Winding number limit $w < N \sim 1/G_N$: cannot probe arbitrarily close to the horizon/singularity
- How reliable is the entanglement/geometry connection for finite w and finite N ?

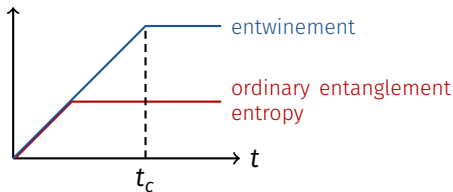
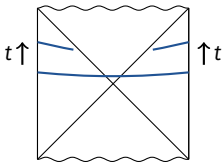
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1/N corrections: formal procedure

- Entwinement \approx ordinary entanglement entropy in “covering theory” on m -times larger torus
- Bulk quantum corrections for entwinement follow from [Faulkner, Lewkowycz, Maldacena (2013)] applied to the covering theory
- DoF for entwinement are located on a subset of all possible collections of long strands (subset \mathcal{H}_m of the Hilbert space)
- QFT with spectrum \mathcal{H}_m : $S_{N/m}$ orbifold of the same seed CFT where only spins $m\mathbb{Z}$ are allowed,

$$Z(\tau) = \text{Tr}_{\mathcal{H}_m} [q^{L_0} \bar{q}^{\bar{L}_0}] = \sum_{h-\bar{h} \in m\mathbb{Z}} q^{h-\frac{c}{24m}} \bar{q}^{\bar{h}-\frac{c}{24m}}$$

- can be rewritten as partition function of covering theory,

$$Z(\tau) = \sum a_{h,\bar{h}} \chi_h^{(c/m)}(-m/\tau) \chi_{\bar{h}}^{(c/m)}(-m/\bar{\tau})$$

- Entwinement = ordinary entanglement entropy for interval of length $(w + L)/m$ and state $\rho \propto \sum_{h-\bar{h} \in m\mathbb{Z}} q^{h-\frac{c}{24m}} \bar{q}^{\bar{h}-\frac{c}{24m}} |h, \bar{h}\rangle \langle h, \bar{h}|$

1/N corrections: explicit computation

- Direct computation of 1/N correction to entanglement at finite temperature gives (using techniques from [Barrella, Dong, Hartnoll, Martin (2013)]) valid for $m, w = O(1)$)

$$S_A^{(w,m)} = \frac{Nc_{\text{seed}}}{3m} \log \left[\frac{\beta}{2\pi^2 \epsilon_{UV}} \sinh \left(\frac{2\pi^2(L+w)}{\beta} \right) \right] \\ + e^{-2m \frac{2\pi}{\beta}} \left[8 - \frac{16\pi^2(L+w)}{\beta} \coth \left(\frac{2\pi^2(L+w)}{\beta} \right) \right] + O(e^{-4m \frac{2\pi}{\beta}}) \\ + O(1/N).$$

- Corrections increase \sim linearly with w but decrease exponentially with m
- ⇒ For small $O(1)$ winding numbers: quantum corrections are tiny, entanglement builds geometry is robust
- large $O(N)$ winding numbers: no universal result for 1/N corrections, depends on seed theory in question

$1/N$ corrections: comments

What about the exactly marginal deformations of the S_N orbifold?

- long strands exist as before but spectrum and OPE coefficients changes
- generalization of FLM: covering theory depends on coupling constant
- explicit computations for small winding numbers: result shown is the universal piece from $1/N$ corrections to the dominant conformal block
- further non-universal $1/N$ corrections may appear depending on the OPE coefficients and spectrum

Summary

- Entanglement between internal DoF (entwinement) is essential for understanding the entanglement/geometry connection in AdS/CFT
- FLM formula applies to entwinement, given the right choice of covering theory and state
- Obstructions to bulk geometry reconstruction from entwinement due to bulk quantum corrections can generically appear for large $O(N) = O(1/G_N)$ winding numbers

Open questions:

- Entwinement for large winding numbers?
- Entanglement between internal DoF in higher dimensions?
- String theoretic AdS/CFT constructions based on $\text{AdS}_{d+1} \times$ **internal space**. Probe geometry of the internal space using entanglement? [Mollabashi, Shiba, Takayanagi (2014); Karch, Uhlemann (2015); Taylor (2015); Das, Kaushal, Mandal, Nanda, Radwan, Trivedi (2022)]