Lessons from light-cone formulation for physics at null infinity

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Flat-space holography





Flat-space holography



Focus: This region of spacetime

This talk

A different setup with two intersecting null hypersurfaces \rightarrow Light-cone formulation

Light-cone formulation

"Forms of relativistic dynamics" [Dirac '49]

- Choice of 'time' for Hamiltonian dynamics
 - (a) Instant form: time x^0

Spatial foliations (initial data on $x^0 = 0$)

(b) Front form: time $x^+ = \frac{x^0 + x^3}{\sqrt{2}}$

Null foliations (initial data on $x^+ = 0$)

• Different IVP in GR: Cauchy vs. Characteristics



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	Instant Form	
	Cartesian coords: $(x^0, x^a), a = 1, 2, 3$	
Kinematical	$\mathcal{K} = \{P_a, M_{ab}\}$: 6-dim	
Dynamical	$\mathcal{D} = \{P_0, M_{0a}\}: 4\text{-dim}$	



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(a) Instant form (b) Front form

Poincaré algebra $(P_{\mu}, M_{\mu\nu})$

	Instant Form Cartesian coords: $(x^0, x^a), a = 1, 2, 3$	Front Form Light-cone coords, $(x^+, x^- x^i), i = 1, 2$
Kinematical	$\mathcal{K} = \{P_a, M_{ab}\}$: 6-dim	$\mathcal{K} = \{P, P_i, M_{-i}, M_{ij}, M_{-+}\} : \mathbf{7-dim}$
Dynamical	$\mathcal{D} = \{\textit{P}_0,\textit{M}_{0a}\}$: 4-dim	$\mathcal{D} = \{\textit{P}_+,\textit{M}_{+i}\}$: 3-dim

Largest Kinematical subgroup possible: Easier to derive interacting actions

Why light-cone?

- Reason 1: Largest kinematical subgroup
- Reason 2: Non-relativistic or Galilean aspects [Weinberg '66; Susskind'68]

 $\mathsf{Light}\mathsf{-}\mathsf{cone}\;\mathsf{Physics}\longleftrightarrow\mathsf{non-relativistic}\;\mathsf{Galilean}\;\mathsf{invariance}$

 \rightarrow 3D Galilei subgroup within 4D light-cone Poincaré

$$P_{\mu}P^{\mu} = 2P_{+}P_{-} - P_{i}P^{i} = 0$$

$$\Rightarrow \text{ Hamiltonian } P_{+} = \frac{P^{i}P_{i}}{2P}$$

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Reason 3: Gauge constraints often solvable → only physical d.o.f.

"Constrained Hamiltonian Systems" [Dirac 1959; Bergmann 1959]

$$S_{\mathcal{H}}[\phi, \pi_{\phi}, \lambda_i] = \int dt \int d^3x \left(\pi_{\phi} \dot{\phi} - \mathcal{H} - \lambda_i \mathcal{G}^i \right) , \quad \mathcal{G}^i$$
 gauge constraints

Closer to on-shell physics: Scattering amplitudes, Helicity states, GW waves...

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Disclaimer: Non-covariant and ugly!

Brief Outline

- Light-cone formulation of QFT
- LC approach to Asymptotic symmetries
- Some recent works:
 - a) Null-front canonical analysis
 - b) Links to Carrollian Physics



Gauge-fixing the Maxwell action

Light-cone gauge : $A_{-} = -A^{+} = -\frac{A^{0} + A^{3}}{\sqrt{2}} = 0$

Maxwell equations: $\partial_{\mu}F^{\mu\nu} = 0$

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$$\nu = +$$
): $\partial_{-}^{2}A^{-} + \partial_{i}\partial_{-}A^{i} = 0$

$$\mathbf{A}^{-} = -\frac{\partial_i \mathbf{A}^{i}}{\partial_{-}} + \mathbf{a}_1(\mathbf{x}^{+}, \mathbf{x}^{i}) \mathbf{x}^{-} + \mathbf{a}_0(\mathbf{x}^{+}, \mathbf{x}^{i})$$

"Inverse derivative" [Mandelstam '83, Leibbrandt '83]

$$\frac{1}{\Theta_{-}}g(x^{-}) = -\int \epsilon(x^{-} - y^{-}) g(y^{-}) dy^{-} + "const."$$

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• Trivial equation ($\nu = -$) relates a_0 and $a_1 \Rightarrow$ one arbitrary constant Let's set them to zero

Dynamical equations (ν = i)

 $(2\partial_{-}\partial_{+} - \partial_{i}\partial^{j})A^{j} = \Box_{lc}A^{j} = 0 \quad \rightarrow \quad \text{two propagating modes}$

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Light-cone action for Electromagnetism

Complexify the $A^i \rightarrow (A, \overline{A})$: ± 1 helicity states of the photon

$$\mathcal{S}_{\mathit{lc}}[A, \bar{A}] = \int d^4x \; \bar{A} \left(\partial_+ \partial_- - \partial \bar{\partial} \right) A \; \rightarrow \; \mathit{lc}_2 \; \mathsf{formalism}$$

Boundary conditions are sneaky!

Boundary conditions in light-cone formalism

Constraint:
$$A^- = -\frac{\partial_i A^i}{\partial_-} + a_1(x^+, x^i) x^- + a_0(x^+, x^i)$$

 $\triangle a_0 = \partial_+ a_1 = \triangle \Phi$; $\triangle = 2\partial \overline{\partial}$
Fall-off at large x^- : $A^i = \partial^i \Phi + \frac{A^i_{(0)}}{(x^-)} + \frac{A^i_{(1)}}{(x^-)^2} + \dots$



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Large gauge transformations $A_{\mu} \rightarrow A_{\mu} + \partial_{\mu} \epsilon(x)$

Case I: Zero modes a₀, a₁ set to zero

- Fall-off: $A^i = O\left(\frac{1}{x^-}\right)$
- LC action:

$$S_{lc}[A,\bar{A}] = \int d^4x \,\bar{A} \left(\partial_+\partial_- - \partial\bar{\partial}\right) A$$

• LGTs : $\epsilon(x, \bar{x}) = f(x) + \bar{f}(\bar{x}) \rightarrow \text{Constrained}$

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Case II: Zero modes $a_0, a_1 \neq 0$

- Fall-off: $A^i = \mathcal{O}(1)$
- LC action:

$$\mathcal{S}[A, \bar{A}, \Phi] = \mathcal{S}_{lc}[A, \bar{A}] + \int_{\partial \Sigma} (\Phi\text{-term})$$

[SM (2022)]

Key lesson: Boundary conditions, zero modes in x^- subtle: handle with care!

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- Light-cone formulation of QFT
- LC approach to Asymptotic symmetries
- Some recent works:
 - a) Null-front canonical analysis

[Barnich, SM, Speziale, Tan, arXiv: 2401.14873]

b) Links to Carrollian Physics

[SM, arXiv: 2406.10353]



Back to basics: Chiral Bosons in 2D

Issue: Boundary conditions, zero modes in x^- delicate

How to account for boundary conditions or zero modes correctly?

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How to account for boundary conditions or zero modes correctly?

The problem: 2D chiral boson Partition function

Partition function $Z(\beta, \alpha) = Tr \ e^{-\beta \hat{H} + i\alpha \hat{P}}$

Time x^0 , Periodic boundary conditions $x^1 \rightarrow x^1 + L$

$$Z(au,ar{ au})=rac{1}{\sqrt{ au_2}|\eta(au)|^2}$$

au modular parameter

[Di Francesco-Mathieu-Sènèchal or your favourite CFT textbook]

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Scalar field Lagrangian

$$S = \frac{1}{2} \int d^2 x \, \partial_\mu \phi \partial^\mu \phi = \int dx^+ dx^- \, \partial_- \phi \partial_+ \phi \quad \rightarrow \quad \text{EOM: } \phi = \phi_+(x^+) + \phi_-(x^-)$$

Left and Right movers

Goal: To reproduce $Z(\tau, \bar{\tau})$ using light-cone qunatization

2D Chiral bosons in light-cone approach

From Hamiltonian analysis and IVP

$$S_{H}[\phi,\pi^{+},\lambda^{+}] = \int dx^{+} \int dx^{-} [\pi^{+}\partial_{+}\phi - \lambda^{+}(\pi^{+} - \partial_{-}\phi)]$$

General Solution:

$$\phi(x^{+}, x^{-}) = \int_{c^{+}}^{x^{+}} dy^{+} \frac{\bar{\lambda}^{+}(y^{+})}{\Delta t} + \int_{c^{-}}^{x^{-}} dy^{-} \frac{\pi^{+}(y^{-})}{\pi^{+}(y^{-})} + \frac{\phi(c^{+}, c^{-})}{\Phi(c^{+}, c^{-})}$$

at $x^{-} = c^{-}$ at $x^{+} = c^{+}$ matching conditions



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Key results [Barnich, SM, Speziale, Tan (2024)]

- Lagrange multiplier λ^+ carries part of initial data
- Infinite tower of global shift symmetries
- Matching conditions crucial for zero modes
- Must quantize on two intersecting light fronts
 - \longrightarrow treat both x^+ and x^- as time

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at $x^- = c^-$ at $x^+ = c^+$ matching conditions

x^{+} (c^{+}, c^{-}) $x^{+} = c^{+}$

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"Two notions of light-cone time"



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Does the LC Poincaré algebra know about this?

 \downarrow

Carrollian Physics

Light-cone and Carrollian Physics

Two notions of time: Newtonian and Carrollian [Duval, Gibbons, Hovarthy, Zhang 2014]

Subgroups of 4D light-cone Poincaré

 x^+ Newtonian, x^- Carrollian

$$g_{+} \qquad b_{+} \qquad c_{+}$$

$$P_{+} \qquad P_{i} \qquad P_{-}$$

$$M_{+i} \qquad M_{ij} \qquad M_{-i} \qquad M_{+-}$$

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Two copies of 3D Carroll c_{\pm} , Bargmann b_{\pm} , Galilei g_{\pm} [SM, arXiv: 2406.10353; Bagchi, Nachiketh, Soni]

$$(\mathfrak{g}_+,\mathfrak{b}_+,\mathfrak{c}_+)\xleftarrow{x^+\leftrightarrow x^-}(\mathfrak{g}_-,\mathfrak{b}_-,\mathfrak{c}_-)$$

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Physical relevance

- 4D Light-cone physics \leftrightarrow 3D Galilean invariance: $\mathfrak{g}_+, \mathfrak{b}_+$ [Susskind]
- $\mathfrak{c}_+ \rightarrow \mathfrak{stability}$ group of light front at $x^+ = \mathfrak{constant}$
- $\mathfrak{c}_- \rightarrow$ stability group of light front at $x^- = \text{constant}$



Carrollian aspects of light-cone field theories

Scalar field action:

$$S = \int dx^{+} dx^{-} d^{d-1} x \left(\partial_{+} \phi \partial_{-} \phi - \frac{1}{2} \partial_{i} \phi \partial^{i} \phi \right)$$

Conjugate momenta:

$$\pi = \frac{\delta \mathcal{L}}{\delta(\partial_+ \phi)} = \partial_- \phi \quad \Rightarrow \quad \text{Constraint} : \chi = \pi - \partial_- \phi$$

Hamiltonian density

$$\mathcal{H}^{lc} = \pi \partial_+ \phi - \mathcal{L} = \frac{1}{2} \partial_i \phi \partial^i \phi \quad \rightarrow \quad \text{No } \partial_+ \text{ or } \pi \text{ terms in } \mathcal{H}$$

Poisson bracket algebra (or more precisely, Dirac bracket)

 $[\mathcal{H}(x), \mathcal{H}(y)] = 0$ [Henneaux (1979)]

→ Light-cone Hamiltonians are of the magnetic Carroll type!

[SM, arXiv: 2406.10353

Shortcut to obtaining Carrollian actions from Lorentzian ones

Some concluding remarks

Lessons

- Boundary conditions in LC subtle: Consider both fronts + matching conditions
- LC theories exhibit both Carrollian and Galilean features
 - E.g., magnetic Carroll nature of LC Hamiltonians

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To-do list

- Dictionary between symmetries in LC approach and asymptotic symmetries at *I* [Barnich, Ciambelli, Gonzalez, Arxiv: 2405.17722]
- Connections to scattering amplitudes, double copy, Self-dual YM and GR, ...
- LC BMS symmetries as conformal Carroll symmetries, ...
- Ambitious goal: Explore flat-space holography

Two intersecting null surfaces, matching conditions, Carrollian aspects, ... [Bekaert, Raj, Arxiv: 2407.17860]

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Modest goals

Learn about LC theories from a Carrollian perspective

AND/ OR

Learn about Carrollian field theories from LC theories



Thank you for your attention!