

Bootstrapping string and M-theory

Shai M. Chester
Imperial College London

Based on 2312.12576 and TBA with S. Pufu and R. Dempsey

Non-perturbative string theory

- Only non-perturbative def of string theory from holography. e.g.:
- IIB string theory on $AdS_5 \times S^5 \Leftrightarrow 4d \mathcal{N} = 4 SU(N)$ SYM.
 - String length ℓ_s and complex string coupling $\tau_s = \chi + i/g_s$ related to SYM $\tau \equiv \frac{4\pi i}{g_{YM}^2} + \frac{\theta}{2\pi}$ as $\tau_s = \tau$ and $L_{AdS}^4/\ell_s^4 = \lambda \equiv g_{YM}^2 N$.
- M-theory on $AdS_4 \times S^7/\mathbb{Z}_k \Leftrightarrow 3d U(N)_k \times U(N)_{-k}$ ABJM.
 - Planck length ℓ_{11} related as $L_{AdS}^9/\ell_{11}^9 = kN^{3/2}$.
- Graviton scattering on AdS \Leftrightarrow stress tensor correlator in CFT.

Non-perturbative string theory

- Only non-perturbative def of string theory from holography. e.g.:
- IIB string theory on $AdS_5 \times S^5 \Leftrightarrow 4d \mathcal{N} = 4 SU(N)$ SYM.
 - String length ℓ_s and complex string coupling $\tau_s = \chi + i/g_s$ related to SYM $\tau \equiv \frac{4\pi i}{g_{YM}^2} + \frac{\theta}{2\pi}$ as $\tau_s = \tau$ and $L_{AdS}^4/\ell_s^4 = \lambda \equiv g_{YM}^2 N$.
- M-theory on $AdS_4 \times S^7/\mathbb{Z}_k \Leftrightarrow 3d U(N)_k \times U(N)_{-k}$ ABJM.
 - Planck length ℓ_{11} related as $L_{AdS}^9/\ell_{11}^9 = kN^{3/2}$.
- Graviton scattering on AdS \Leftrightarrow stress tensor correlator in CFT.

Non-perturbative string theory

- Only non-perturbative def of string theory from holography. e.g.:
- IIB string theory on $AdS_5 \times S^5 \Leftrightarrow 4d \mathcal{N} = 4 SU(N)$ SYM.
 - String length ℓ_s and complex string coupling $\tau_s = \chi + i/g_s$ related to SYM $\tau \equiv \frac{4\pi i}{g_{YM}^2} + \frac{\theta}{2\pi}$ as $\tau_s = \tau$ and $L_{AdS}^4/\ell_s^4 = \lambda \equiv g_{YM}^2 N$.
- M-theory on $AdS_4 \times S^7/\mathbb{Z}_k \Leftrightarrow 3d U(N)_k \times U(N)_{-k}$ ABJM.
 - Planck length ℓ_{11} related as $L_{AdS}^9/\ell_{11}^9 = kN^{3/2}$.
- Graviton scattering on AdS \Leftrightarrow stress tensor correlator in CFT.

Non-perturbative string theory

- Only non-perturbative def of string theory from holography. e.g.:
- IIB string theory on $AdS_5 \times S^5 \Leftrightarrow 4d \mathcal{N} = 4 SU(N)$ SYM.
 - String length ℓ_s and complex string coupling $\tau_s = \chi + i/g_s$ related to SYM $\tau \equiv \frac{4\pi i}{g_{YM}^2} + \frac{\theta}{2\pi}$ as $\tau_s = \tau$ and $L_{AdS}^4/\ell_s^4 = \lambda \equiv g_{YM}^2 N$.
- M-theory on $AdS_4 \times S^7/\mathbb{Z}_k \Leftrightarrow 3d U(N)_k \times U(N)_{-k}$ ABJM.
 - Planck length ℓ_{11} related as $L_{AdS}^9/\ell_{11}^9 = kN^{3/2}$.
- Graviton scattering on AdS \Leftrightarrow stress tensor correlator in CFT.

Non-perturbative string theory

- Only non-perturbative def of string theory from holography. e.g.:
- IIB string theory on $AdS_5 \times S^5 \Leftrightarrow 4d \mathcal{N} = 4 SU(N)$ SYM.
 - String length ℓ_s and complex string coupling $\tau_s = \chi + i/g_s$ related to SYM $\tau \equiv \frac{4\pi i}{g_{YM}^2} + \frac{\theta}{2\pi}$ as $\tau_s = \tau$ and $L_{AdS}^4/\ell_s^4 = \lambda \equiv g_{YM}^2 N$.
- M-theory on $AdS_4 \times S^7/\mathbb{Z}_k \Leftrightarrow 3d U(N)_k \times U(N)_{-k}$ ABJM.
 - Planck length ℓ_{11} related as $L_{AdS}^9/\ell_{11}^9 = kN^{3/2}$.
- Graviton scattering on AdS \Leftrightarrow stress tensor correlator in CFT.

Non-perturbative string theory

- Only non-perturbative def of string theory from holography. e.g.:
- IIB string theory on $AdS_5 \times S^5 \Leftrightarrow 4d \mathcal{N} = 4 SU(N)$ SYM.
 - String length ℓ_s and complex string coupling $\tau_s = \chi + i/g_s$ related to SYM $\tau \equiv \frac{4\pi i}{g_{YM}^2} + \frac{\theta}{2\pi}$ as $\tau_s = \tau$ and $L_{AdS}^4/\ell_s^4 = \lambda \equiv g_{YM}^2 N$.
- M-theory on $AdS_4 \times S^7/\mathbb{Z}_k \Leftrightarrow 3d U(N)_k \times U(N)_{-k}$ ABJM.
 - Planck length ℓ_{11} related as $L_{AdS}^9/\ell_{11}^9 = kN^{3/2}$.
- Graviton scattering on AdS \Leftrightarrow stress tensor correlator in CFT.

Bootstrapping holography

- AdS/CFT has recast quantum gravity into easier problem of CFT, but CFTs still strongly coupled so cannot compute in practice.
 - At leading large N (i.e. classical string) can use integrability for finite λ , but cannot study quantum string (or M-theory).
 - For SYM can do weak coupling, but gravity at strong coupling.
- Today we combine three non-perturbative constraints to numerically study holographic CFTs (and thus quantum gravity):
 - 1 Unitarity: OPE coefficients real and Δ bounded.
 - 2 Crossing: combine with unitarity for numerical bootstrap algorithm.
 - 3 Localization: Input protected coupling-dependence to bootstrap.

Bootstrapping holography

- AdS/CFT has recast quantum gravity into easier problem of CFT, but CFTs still strongly coupled so cannot compute in practice.
 - At leading large N (i.e. classical string) can use integrability for finite λ , but cannot study quantum string (or M-theory).
 - For SYM can do weak coupling, but gravity at strong coupling.
- Today we combine three non-perturbative constraints to numerically study holographic CFTs (and thus quantum gravity):
 - 1 Unitarity: OPE coefficients real and Δ bounded.
 - 2 Crossing: combine with unitarity for numerical bootstrap algorithm.
 - 3 Localization: Input protected coupling-dependence to bootstrap.

Bootstrapping holography

- AdS/CFT has recast quantum gravity into easier problem of CFT, but CFTs still strongly coupled so cannot compute in practice.
 - At leading large N (i.e. classical string) can use integrability for finite λ , but cannot study quantum string (or M-theory).
 - For SYM can do weak coupling, but gravity at strong coupling.
- Today we combine three non-perturbative constraints to numerically study holographic CFTs (and thus quantum gravity):
 - 1 Unitarity: OPE coefficients real and Δ bounded.
 - 2 Crossing: combine with unitarity for numerical bootstrap algorithm.
 - 3 Localization: Input protected coupling-dependence to bootstrap.

Bootstrapping holography

- AdS/CFT has recast quantum gravity into easier problem of CFT, but CFTs still strongly coupled so cannot compute in practice.
 - At leading large N (i.e. classical string) can use integrability for finite λ , but cannot study quantum string (or M-theory).
 - For SYM can do weak coupling, but gravity at strong coupling.
- Today we combine three non-perturbative constraints to numerically study holographic CFTs (and thus quantum gravity):
 - 1 Unitarity: OPE coefficients real and Δ bounded.
 - 2 Crossing: combine with unitarity for numerical bootstrap algorithm.
 - 3 Localization: Input protected coupling-dependence to bootstrap.

Bootstrapping holography

- AdS/CFT has recast quantum gravity into easier problem of CFT, but CFTs still strongly coupled so cannot compute in practice.
 - At leading large N (i.e. classical string) can use integrability for finite λ , but cannot study quantum string (or M-theory).
 - For SYM can do weak coupling, but gravity at strong coupling.
- Today we combine three non-perturbative constraints to numerically study holographic CFTs (and thus quantum gravity):
 - 1 Unitarity: OPE coefficients real and Δ bounded.
 - 2 Crossing: combine with unitarity for numerical bootstrap algorithm.
 - 3 Localization: Input protected coupling-dependence to bootstrap.

Bootstrapping holography

- AdS/CFT has recast quantum gravity into easier problem of CFT, but CFTs still strongly coupled so cannot compute in practice.
 - At leading large N (i.e. classical string) can use integrability for finite λ , but cannot study quantum string (or M-theory).
 - For SYM can do weak coupling, but gravity at strong coupling.
- Today we combine three non-perturbative constraints to numerically study holographic CFTs (and thus quantum gravity):
 - 1 Unitarity: OPE coefficients real and Δ bounded.
 - 2 Crossing: combine with unitarity for numerical bootstrap algorithm.
 - 3 Localization: Input protected coupling-dependence to bootstrap.

Bootstrapping holography

- AdS/CFT has recast quantum gravity into easier problem of CFT, but CFTs still strongly coupled so cannot compute in practice.
 - At leading large N (i.e. classical string) can use integrability for finite λ , but cannot study quantum string (or M-theory).
 - For SYM can do weak coupling, but gravity at strong coupling.
- Today we combine three non-perturbative constraints to numerically study holographic CFTs (and thus quantum gravity):
 - 1 Unitarity: OPE coefficients real and Δ bounded.
 - 2 Crossing: combine with unitarity for numerical bootstrap algorithm.
 - 3 Localization: Input protected coupling-dependence to bootstrap.

4d $\mathcal{N} = 4$ SYM stress tensor correlator

- Expand stress tensor superprimary correlator in blocks:

$$\langle S(x_1)S(x_2)S(x_3)S(x_4) \rangle = G_{prot}(c, x_i) + \Theta_{kin}(x_i) \sum_{\Delta, \ell} \lambda_{\Delta, \ell}^2 g_{\Delta, \ell}(x_i)$$

- Susy ward identities solved by this expansion, where all dynamical info in λ, Δ of long multiplets in singlet of $SU(4)_R$.
- Theory specified by $c \sim N^2$ and τ .
 - $1/c$ enters linearly into $G_{prot}(c, x_i)$.
 - Both τ and c effect long multiplet CFT data nontrivially.

4d $\mathcal{N} = 4$ SYM stress tensor correlator

- Expand stress tensor superprimary correlator in blocks:

$$\langle S(x_1)S(x_2)S(x_3)S(x_4) \rangle = G_{prot}(c, x_i) + \Theta_{kin}(x_i) \sum_{\Delta, \ell} \lambda_{\Delta, \ell}^2 g_{\Delta, \ell}(x_i)$$

- Susy ward identities solved by this expansion, where all dynamical info in λ, Δ of long multiplets in singlet of $SU(4)_R$.
- Theory specified by $c \sim N^2$ and τ .
 - $1/c$ enters linearly into $G_{prot}(c, x_i)$.
 - Both τ and c effect long multiplet CFT data nontrivially.

4d $\mathcal{N} = 4$ SYM stress tensor correlator

- Expand stress tensor superprimary correlator in blocks:

$$\langle S(x_1)S(x_2)S(x_3)S(x_4) \rangle = G_{prot}(c, x_i) + \Theta_{kin}(x_i) \sum_{\Delta, \ell} \lambda_{\Delta, \ell}^2 g_{\Delta, \ell}(x_i)$$

- Susy ward identities solved by this expansion, where all dynamical info in λ, Δ of long multiplets in singlet of $SU(4)_R$.
- Theory specified by $c \sim N^2$ and τ .
 - $1/c$ enters linearly into $G_{prot}(c, x_i)$.
 - Both τ and c effect long multiplet CFT data nontrivially.

4d $\mathcal{N} = 4$ SYM stress tensor correlator

- Expand stress tensor superprimary correlator in blocks:

$$\langle S(x_1)S(x_2)S(x_3)S(x_4) \rangle = G_{prot}(c, x_i) + \Theta_{kin}(x_i) \sum_{\Delta, \ell} \lambda_{\Delta, \ell}^2 g_{\Delta, \ell}(x_i)$$

- Susy ward identities solved by this expansion, where all dynamical info in λ, Δ of long multiplets in singlet of $SU(4)_R$.
- Theory specified by $c \sim N^2$ and τ .
 - $1/c$ enters linearly into $G_{prot}(c, x_i)$.
 - Both τ and c effect long multiplet CFT data nontrivially.

4d $\mathcal{N} = 4$ SYM stress tensor correlator

- Expand stress tensor superprimary correlator in blocks:

$$\langle S(x_1)S(x_2)S(x_3)S(x_4) \rangle = G_{prot}(c, x_i) + \Theta_{kin}(x_i) \sum_{\Delta, \ell} \lambda_{\Delta, \ell}^2 g_{\Delta, \ell}(x_i)$$

- Susy ward identities solved by this expansion, where all dynamical info in λ, Δ of long multiplets in singlet of $SU(4)_R$.
- Theory specified by $c \sim N^2$ and τ .
 - $1/c$ enters linearly into $G_{prot}(c, x_i)$.
 - Both τ and c effect long multiplet CFT data nontrivially.

4d $\mathcal{N} = 4$ SYM stress tensor correlator

- Expand stress tensor superprimary correlator in blocks:

$$\langle S(x_1)S(x_2)S(x_3)S(x_4) \rangle = G_{prot}(c, x_i) + \Theta_{kin}(x_i) \sum_{\Delta, \ell} \lambda_{\Delta, \ell}^2 g_{\Delta, \ell}(x_i)$$

- Susy ward identities solved by this expansion, where all dynamical info in λ, Δ of long multiplets in singlet of $SU(4)_R$.
- Theory specified by $c \sim N^2$ and τ .
 - $1/c$ enters linearly into $G_{prot}(c, x_i)$.
 - Both τ and c effect long multiplet CFT data nontrivially.

Localization constraints

- Derivatives of free energy $F(m)$ deformed by hyper mass relate to S^4 integrals of correlator [Binder, SMC, Pufu, Wang '19; SMC, Pufu '20] :

$$\frac{1}{8c} \frac{\partial_m^2 \partial_\tau \partial_{\bar{\tau}} F}{\partial_\tau \partial_{\bar{\tau}} F} \Big|_{m=0} = I_2 \left[\sum_{\Delta, \ell} \lambda_{\Delta, \ell}^2 g_{\Delta, \ell}(x_i) - f_{\text{prot}}(c, x_i) \right],$$
$$-48\zeta(3)c^{-1} - c^{-2} \partial_m^4 F \Big|_{m=0} = I_4 \left[\sum_{\Delta, \ell} \lambda_{\Delta, \ell}^2 g_{\Delta, \ell}(x_i) - f_{\text{prot}}(c, x_i) \right].$$

- LHS can be computed as function of N and τ in terms of N dimensional matrix model from localization [Pestun '08].
 - LHS computed in closed form for any N and τ to high accuracy [Alday, SMC, Dorigoni, Green, Wen '23].
 - For any N , allows us to input τ into numerical bootstrap, as two extra linear constraints on CFT data [SMC, Dempsey, Pufu '21].

Localization constraints

- Derivatives of free energy $F(m)$ deformed by hyper mass relate to S^4 integrals of correlator [Binder, SMC, Pufu, Wang '19; SMC, Pufu '20] :

$$\frac{1}{8c} \frac{\partial_m^2 \partial_\tau \partial_{\bar{\tau}} F}{\partial_\tau \partial_{\bar{\tau}} F} \Big|_{m=0} = I_2 \left[\sum_{\Delta, \ell} \lambda_{\Delta, \ell}^2 g_{\Delta, \ell}(x_i) - f_{\text{prot}}(c, x_i) \right],$$
$$-48\zeta(3)c^{-1} - c^{-2} \partial_m^4 F \Big|_{m=0} = I_4 \left[\sum_{\Delta, \ell} \lambda_{\Delta, \ell}^2 g_{\Delta, \ell}(x_i) - f_{\text{prot}}(c, x_i) \right].$$

- LHS can be computed as function of N and τ in terms of N dimensional matrix model from localization [Pestun '08].
 - LHS computed in closed form for any N and τ to high accuracy [Alday, SMC, Dorigoni, Green, Wen '23].
 - For any N , allows us to input τ into numerical bootstrap, as two extra linear constraints on CFT data [SMC, Dempsey, Pufu '21].

Localization constraints

- Derivatives of free energy $F(m)$ deformed by hyper mass relate to S^4 integrals of correlator [Binder, SMC, Pufu, Wang '19; SMC, Pufu '20] :

$$\frac{1}{8c} \frac{\partial_m^2 \partial_\tau \partial_{\bar{\tau}} F}{\partial_\tau \partial_{\bar{\tau}} F} \Big|_{m=0} = I_2 \left[\sum_{\Delta, \ell} \lambda_{\Delta, \ell}^2 g_{\Delta, \ell}(x_i) - f_{\text{prot}}(c, x_i) \right],$$
$$-48\zeta(3)c^{-1} - c^{-2} \partial_m^4 F \Big|_{m=0} = I_4 \left[\sum_{\Delta, \ell} \lambda_{\Delta, \ell}^2 g_{\Delta, \ell}(x_i) - f_{\text{prot}}(c, x_i) \right].$$

- LHS can be computed as function of N and τ in terms of N dimensional matrix model from localization [Pestun '08].
 - LHS computed in closed form for any N and τ to high accuracy [Alday, SMC, Dorigoni, Green, Wen '23].
 - For any N , allows us to input τ into numerical bootstrap, as two extra linear constraints on CFT data [SMC, Dempsey, Pufu '21].

Localization constraints

- Derivatives of free energy $F(m)$ deformed by hyper mass relate to S^4 integrals of correlator [Binder, SMC, Pufu, Wang '19; SMC, Pufu '20] :

$$\frac{1}{8c} \frac{\partial_m^2 \partial_\tau \partial_{\bar{\tau}} F}{\partial_\tau \partial_{\bar{\tau}} F} \Big|_{m=0} = I_2 \left[\sum_{\Delta, \ell} \lambda_{\Delta, \ell}^2 g_{\Delta, \ell}(x_i) - f_{\text{prot}}(c, x_i) \right],$$
$$-48\zeta(3)c^{-1} - c^{-2} \partial_m^4 F \Big|_{m=0} = I_4 \left[\sum_{\Delta, \ell} \lambda_{\Delta, \ell}^2 g_{\Delta, \ell}(x_i) - f_{\text{prot}}(c, x_i) \right].$$

- LHS can be computed as function of N and τ in terms of N dimensional matrix model from localization [Pestun '08].
 - LHS computed in closed form for any N and τ to high accuracy [Alday, SMC, Dorigoni, Green, Wen '23].
 - For any N , allows us to input τ into numerical bootstrap, as two extra linear constraints on CFT data [SMC, Dempsey, Pufu '21].

Localization constraints

- Derivatives of free energy $F(m)$ deformed by hyper mass relate to S^4 integrals of correlator [Binder, SMC, Pufu, Wang '19; SMC, Pufu '20] :

$$\frac{1}{8c} \frac{\partial_m^2 \partial_\tau \partial_{\bar{\tau}} F}{\partial_\tau \partial_{\bar{\tau}} F} \Big|_{m=0} = I_2 \left[\sum_{\Delta, \ell} \lambda_{\Delta, \ell}^2 g_{\Delta, \ell}(x_i) - f_{\text{prot}}(c, x_i) \right],$$
$$-48\zeta(3)c^{-1} - c^{-2} \partial_m^4 F \Big|_{m=0} = I_4 \left[\sum_{\Delta, \ell} \lambda_{\Delta, \ell}^2 g_{\Delta, \ell}(x_i) - f_{\text{prot}}(c, x_i) \right].$$

- LHS can be computed as function of N and τ in terms of N dimensional matrix model from localization [Pestun '08].
 - LHS computed in closed form for any N and τ to high accuracy [Alday, SMC, Dorigoni, Green, Wen '23].
 - For any N , allows us to input τ into numerical bootstrap, as two extra linear constraints on CFT data [SMC, Dempsey, Pufu '21].

Perturbative results

- Small g_{YM} expansion at finite N , e.g. Konishi [Velizhanin '09]

$$\Delta = 2 + \frac{3\lambda}{4\pi^2} - \frac{3\lambda^2}{16\pi^4} + \frac{21\lambda^3}{256\pi^6} + \frac{\lambda^4 \left(-1440 \left(\frac{12}{N^2} + 1 \right) \zeta(5) + 576\zeta(3) - 2496 \right)}{65536\pi^8} + O(\lambda^5)$$

- Integrability gives leading large N at finite λ , indistinguishable from resummed weak coupling in regime we'll bootstrap.
- Large N finite τ from analytic bootstrap (i.e. crossing but NOT unitarity) + localization [SMC '19; SMC, Green, Pufu, Wang, Wen '19]:

$$\Delta = 4 - \frac{4}{c} + \frac{135}{7\sqrt{2}\pi^{3/2}c^{7/4}} E\left(\frac{3}{2}, \tau\right) + \frac{1199}{42c^2} - \frac{3825}{32\sqrt{2}\pi^{5/2}c^{9/4}} E\left(\frac{5}{2}, \tau\right) + \dots$$

Perturbative results

- Small g_{YM} expansion at finite N , e.g. Konishi [Velizhanin '09]

$$\Delta = 2 + \frac{3\lambda}{4\pi^2} - \frac{3\lambda^2}{16\pi^4} + \frac{21\lambda^3}{256\pi^6} + \frac{\lambda^4 \left(-1440 \left(\frac{12}{N^2} + 1 \right) \zeta(5) + 576\zeta(3) - 2496 \right)}{65536\pi^8} + O(\lambda^5)$$

- Integrability gives leading large N at finite λ , indistinguishable from resummed weak coupling in regime we'll bootstrap.
- Large N finite τ from analytic bootstrap (i.e. crossing but NOT unitarity) + localization [SMC '19; SMC, Green, Pufu, Wang, Wen '19] :

$$\Delta = 4 - \frac{4}{c} + \frac{135}{7\sqrt{2}\pi^{3/2}c^{7/4}} E\left(\frac{3}{2}, \tau\right) + \frac{1199}{42c^2} - \frac{3825}{32\sqrt{2}\pi^{5/2}c^{9/4}} E\left(\frac{5}{2}, \tau\right) + \dots$$

Perturbative results

- Small g_{YM} expansion at finite N , e.g. Konishi [Velizhanin '09]

$$\Delta = 2 + \frac{3\lambda}{4\pi^2} - \frac{3\lambda^2}{16\pi^4} + \frac{21\lambda^3}{256\pi^6} + \frac{\lambda^4 \left(-1440 \left(\frac{12}{N^2} + 1 \right) \zeta(5) + 576\zeta(3) - 2496 \right)}{65536\pi^8} + O(\lambda^5)$$

- Integrability gives leading large N at finite λ , indistinguishable from resummed weak coupling in regime we'll bootstrap.
- Large N finite τ from analytic bootstrap (i.e. crossing but NOT unitarity) + localization [SMC '19; SMC, Green, Pufu, Wang, Wen '19] :

$$\Delta = 4 - \frac{4}{c} + \frac{135}{7\sqrt{2}\pi^{3/2}c^{7/4}} E\left(\frac{3}{2}, \tau\right) + \frac{1199}{42c^2} - \frac{3825}{32\sqrt{2}\pi^{5/2}c^{9/4}} E\left(\frac{5}{2}, \tau\right) + \dots$$

Perturbative results

- Small g_{YM} expansion at finite N , e.g. Konishi [Velizhanin '09]

$$\Delta = 2 + \frac{3\lambda}{4\pi^2} - \frac{3\lambda^2}{16\pi^4} + \frac{21\lambda^3}{256\pi^6} + \frac{\lambda^4 \left(-1440 \left(\frac{12}{N^2} + 1 \right) \zeta(5) + 576\zeta(3) - 2496 \right)}{65536\pi^8} + O(\lambda^5)$$

- Integrability gives leading large N at finite λ , indistinguishable from resummed weak coupling in regime we'll bootstrap.
- Large N finite τ from analytic bootstrap (i.e. crossing but NOT unitarity) + localization [SMC '19; SMC, Green, Pufu, Wang, Wen '19]:

$$\Delta = 4 - \frac{4}{c} + \frac{135}{7\sqrt{2}\pi^{3/2}c^{7/4}} E\left(\frac{3}{2}, \tau\right) + \frac{1199}{42c^2} - \frac{3825}{32\sqrt{2}\pi^{5/2}c^{9/4}} E\left(\frac{5}{2}, \tau\right) + \dots$$

Perturbative results

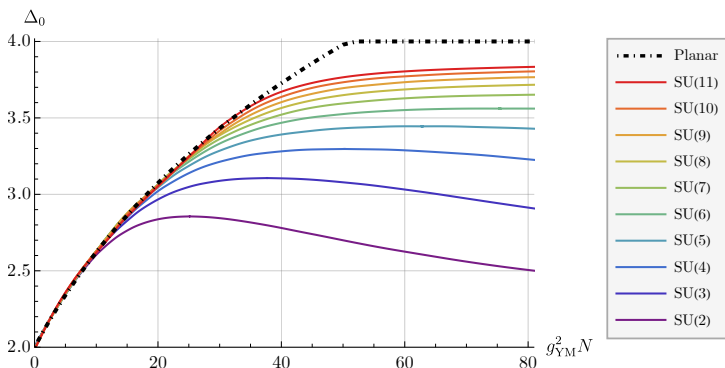
- Small g_{YM} expansion at finite N , e.g. Konishi [Velizhanin '09]

$$\Delta = 2 + \frac{3\lambda}{4\pi^2} - \frac{3\lambda^2}{16\pi^4} + \frac{21\lambda^3}{256\pi^6} + \frac{\lambda^4 \left(-1440 \left(\frac{12}{N^2} + 1 \right) \zeta(5) + 576\zeta(3) - 2496 \right)}{65536\pi^8} + O(\lambda^5)$$

- Integrability gives leading large N at finite λ , indistinguishable from resummed weak coupling in regime we'll bootstrap.
- Large N finite τ from analytic bootstrap (i.e. crossing but NOT unitarity) + localization [SMC '19; SMC, Green, Pufu, Wang, Wen '19] :

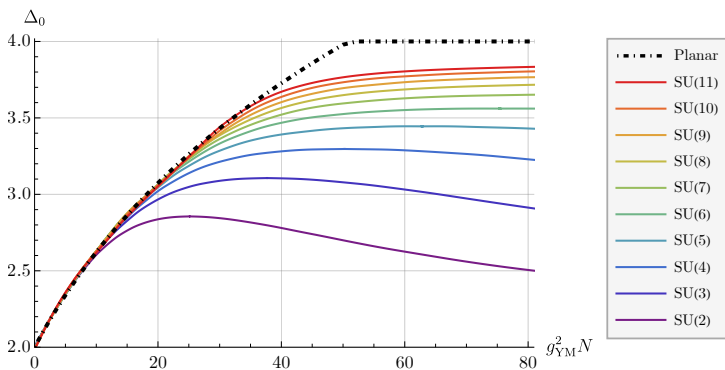
$$\Delta = 4 - \frac{4}{c} + \frac{135}{7\sqrt{2}\pi^{3/2}c^{7/4}} E\left(\frac{3}{2}, \tau\right) + \frac{1199}{42c^2} - \frac{3825}{32\sqrt{2}\pi^{5/2}c^{9/4}} E\left(\frac{5}{2}, \tau\right) + \dots$$

Bootstrap bounds on lowest Δ for various N



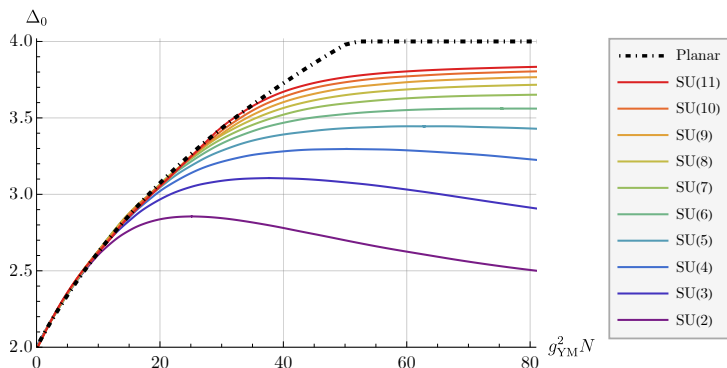
- Bounds are converging to Planar integrability spectrum (similar to Pade resummed 4-loop weak coupling in this regime).
- Planar integrability for double trace is trivially $4 = 2 + 2$.

Bootstrap bounds on lowest Δ for various N



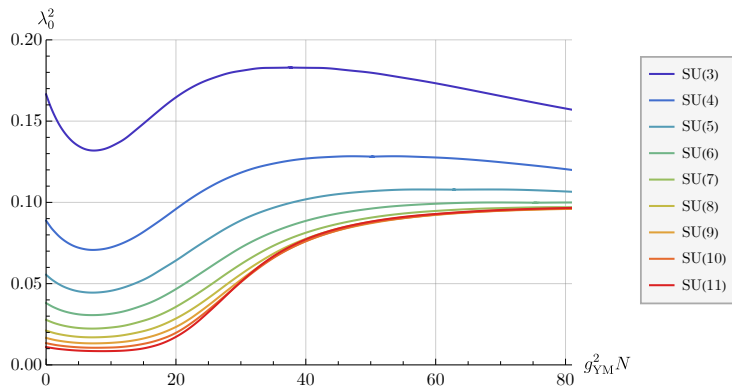
- Bounds are converging to Planar integrability spectrum (similar to Pade resummed 4-loop weak coupling in this regime).
- Planar integrability for double trace is trivially $4 = 2 + 2$.

Bootstrap bounds on lowest Δ for various N



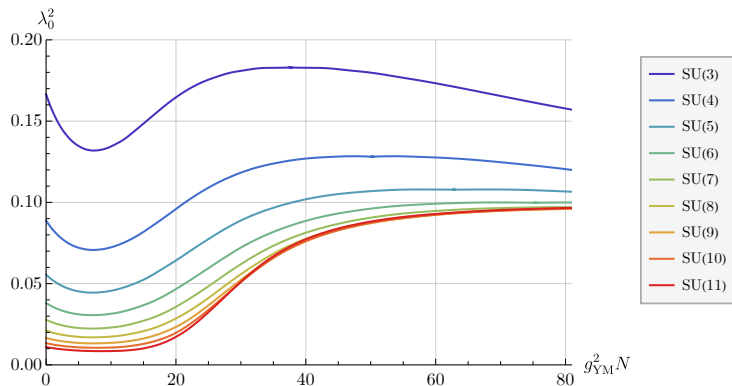
- Bounds are converging to Planar integrability spectrum (similar to Pade resummed 4-loop weak coupling in this regime).
- Planar integrability for double trace is trivially $4 = 2 + 2$.

Bounds on lowest λ^2 for various N



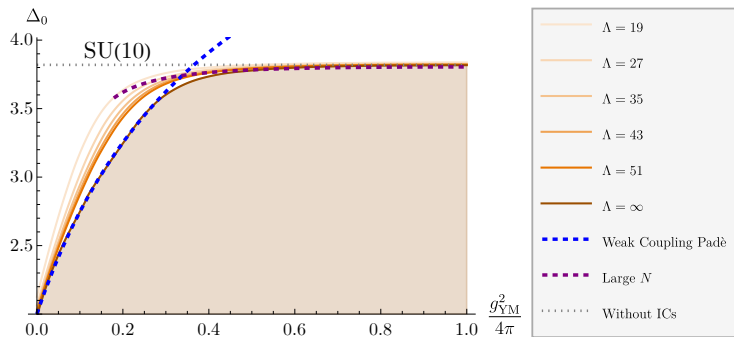
- No planar integrability results to compare to now.

Bounds on lowest λ^2 for various N



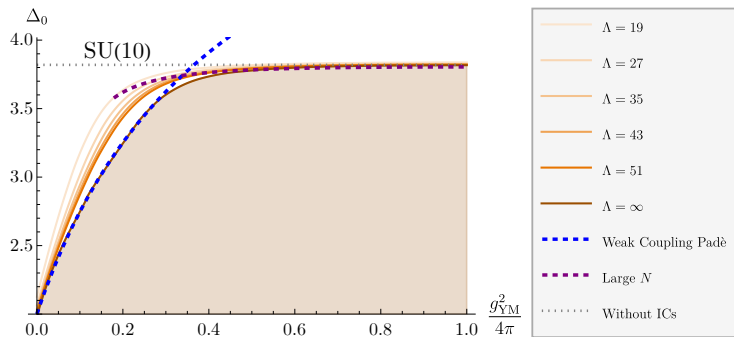
- No planar integrability results to compare to now.

Bounds: Lowest Δ for $SU(10)$



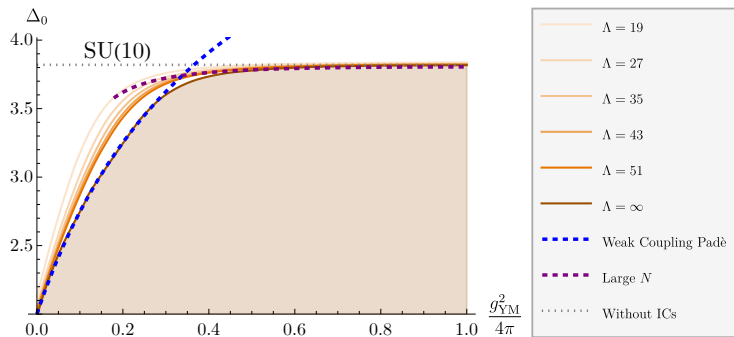
- Matches BOTH weak coupling and strong coupling expansions!
- Observe non-pert **level repulsion**, in between weak coupling for single trace and strong coupling for double trace.

Bounds: Lowest Δ for $SU(10)$



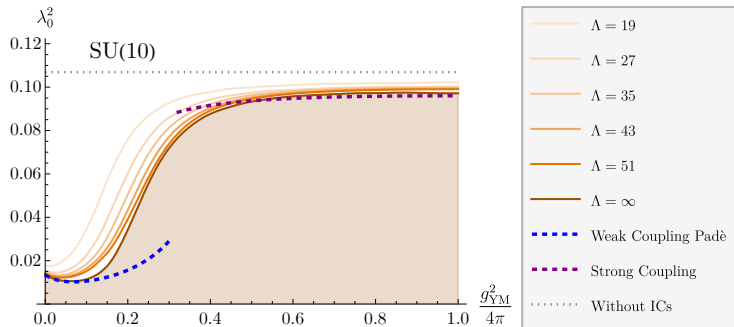
- Matches BOTH weak coupling and strong coupling expansions!
- Observe non-pert **level repulsion**, in between weak coupling for single trace and strong coupling for double trace.

Bounds: Lowest Δ for $SU(10)$



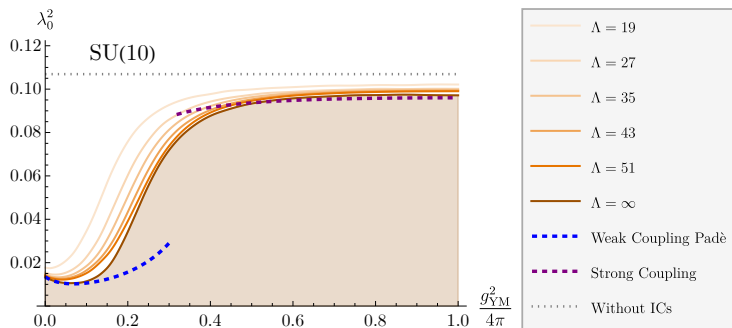
- Matches BOTH weak coupling and strong coupling expansions!
- Observe non-pert **level repulsion**, in between weak coupling for single trace and strong coupling for double trace.

Bounds: Lowest λ^2 for $SU(10)$



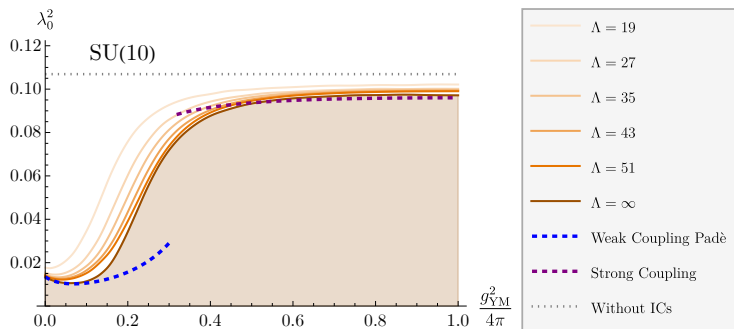
- Matches BOTH weak coupling and strong coupling expansions!
- Observe non-pert **level repulsion**, in between weak coupling for single trace and strong coupling for double trace.

Bounds: Lowest λ^2 for $SU(10)$



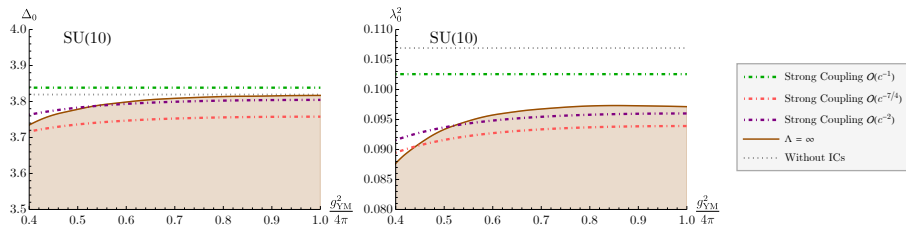
- Matches BOTH weak coupling and strong coupling expansions!
- Observe non-pert **level repulsion**, in between weak coupling for single trace and strong coupling for double trace.

Bounds: Lowest λ^2 for $SU(10)$



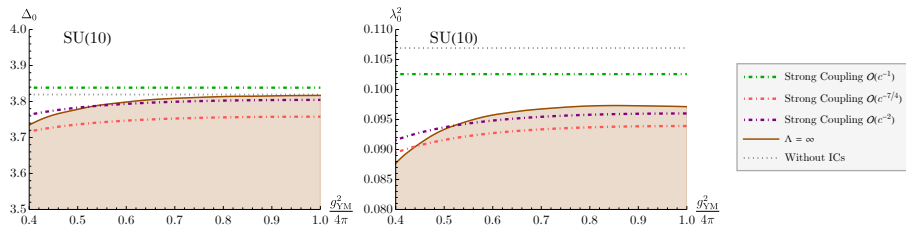
- Matches BOTH weak coupling and strong coupling expansions!
- Observe non-pert **level repulsion**, in between weak coupling for single trace and strong coupling for double trace.

Sensitivity to stringy corrections



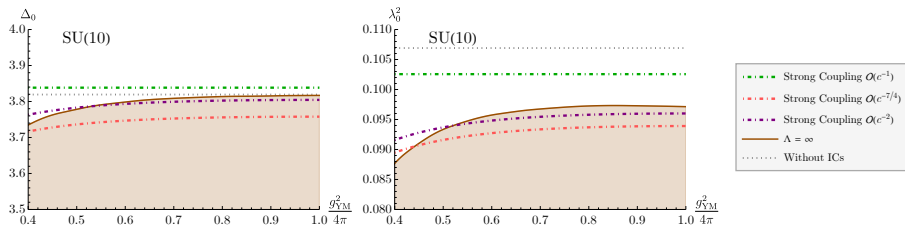
- For largish N (e.g. $SU(10)$), we see that analytic bootstrap result gets closer to bound as we include more $1/c$ corrections.
- $1/c$ is supergravity, $1/c^{7/4}$ is R^4 correction [SMC, Green, Pufu, Wang, Wen '19], $1/c^2$ is 1-loop correction [Alday, Bissi '17; Aprile, Drummond, Heslop, Paul '17] (which included contact term fixed from localization [SMC '19]).
- So bootstrap sensitive to stringy corrections!

Sensitivity to stringy corrections



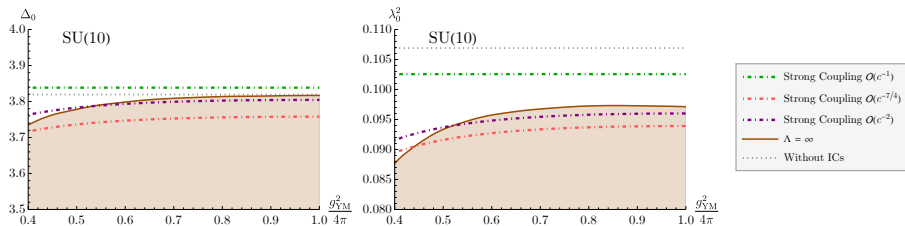
- For largish N (e.g. $SU(10)$), we see that analytic bootstrap result gets closer to bound as we include more $1/c$ corrections.
- $1/c$ is supergravity, $1/c^{7/4}$ is R^4 correction [SMC, Green, Pufu, Wang, Wen '19], $1/c^2$ is 1-loop correction [Alday, Bissi '17; Aprile, Drummond, Heslop, Paul '17] (which included contact term fixed from localization [SMC '19]).
- So bootstrap sensitive to stringy corrections!

Sensitivity to stringy corrections



- For largish N (e.g. $SU(10)$), we see that analytic bootstrap result gets closer to bound as we include more $1/c$ corrections.
- $1/c$ is supergravity, $1/c^{7/4}$ is R^4 correction [SMC, Green, Pufu, Wang, Wen '19], $1/c^2$ is 1-loop correction [Alday, Bissi '17; Aprile, Drummond, Heslop, Paul '17] (which included contact term fixed from localization [SMC '19]).
- So bootstrap sensitive to stringy corrections!

Sensitivity to stringy corrections



- For largish N (e.g. $SU(10)$), we see that analytic bootstrap result gets closer to bound as we include more $1/c$ corrections.
- $1/c$ is supergravity, $1/c^{7/4}$ is R^4 correction [SMC, Green, Pufu, Wang, Wen '19], $1/c^2$ is 1-loop correction [Alday, Bissi '17; Aprile, Drummond, Heslop, Paul '17] (which included contact term fixed from localization [SMC '19]).
- So bootstrap sensitive to stringy corrections!

3d ABJM stress tensor correlator

- Expand stress tensor superprimary correlator in superblocks:

$$\langle S^a(x_1) S^b(x_2) S^c(x_3) S^d(x_4) \rangle = \frac{1}{x_{12}^2 x_{34}^2} \left[\sum_{\Delta, \ell} \lambda_{\Delta, \ell}^2 G_{\Delta, \ell}^{abcd}(x_i) + \sum_M \lambda_M^2 G_M^{abcd}(x_i) \right]$$

- Susy ward identities solved by superblocks, written as linear combos of 3d blocks in terms of structures $\sum_{i=1}^6 T_i^{abcd} \frac{G^i(U, V)}{x_{12}^2 x_{34}^2}$.
- Protected multiplets are nontrivial, include $(A, 2)$ for odd spin, $(A, +)$ for even spin, $(B, 2)$, $(B, +)$ and stress tensor for spin zero. Can bootstrap both upper and lower bounds (i.e. islands)!
- Theory specified by $c \sim N^{3/2}$ and $k = 1, 2$.
 - For $k = 2$ two theories $U(N)_2 \times U(N)_{-2}$ ABJM and $U(N)_2 \times U(N+1)_{-2}$ ABJ.
 - For $k = 1$, theory has free and interacting parts.

3d ABJM stress tensor correlator

- Expand stress tensor superprimary correlator in superblocks:

$$\langle S^a(x_1) S^b(x_2) S^c(x_3) S^d(x_4) \rangle = \frac{1}{x_{12}^2 x_{34}^2} \left[\sum_{\Delta, \ell} \lambda_{\Delta, \ell}^2 G_{\Delta, \ell}^{abcd}(x_i) + \sum_M \lambda_M^2 G_M^{abcd}(x_i) \right]$$

- Susy ward identities solved by superblocks, written as linear combos of 3d blocks in terms of structures $\sum_{i=1}^6 T_i^{abcd} \frac{G^i(U, V)}{x_{12}^2 x_{34}^2}$.
- Protected multiplets are nontrivial, include $(A, 2)$ for odd spin, $(A, +)$ for even spin, $(B, 2)$, $(B, +)$ and stress tensor for spin zero. Can bootstrap both upper and lower bounds (i.e. islands)!
- Theory specified by $c \sim N^{3/2}$ and $k = 1, 2$.
 - For $k = 2$ two theories $U(N)_2 \times U(N)_{-2}$ ABJM and $U(N)_2 \times U(N+1)_{-2}$ ABJ.
 - For $k = 1$, theory has free and interacting parts.

3d ABJM stress tensor correlator

- Expand stress tensor superprimary correlator in superblocks:

$$\langle S^a(x_1) S^b(x_2) S^c(x_3) S^d(x_4) \rangle = \frac{1}{x_{12}^2 x_{34}^2} \left[\sum_{\Delta, \ell} \lambda_{\Delta, \ell}^2 G_{\Delta, \ell}^{abcd}(x_i) + \sum_M \lambda_M^2 G_M^{abcd}(x_i) \right]$$

- Susy ward identities solved by superblocks, written as linear combos of 3d blocks in terms of structures $\sum_{i=1}^6 T_i^{abcd} \frac{G^i(U, V)}{x_{12}^2 x_{34}^2}$.
- Protected multiplets are nontrivial, include $(A, 2)$ for odd spin, $(A, +)$ for even spin, $(B, 2)$, $(B, +)$ and stress tensor for spin zero. Can bootstrap both upper and lower bounds (i.e. islands)!
- Theory specified by $c \sim N^{3/2}$ and $k = 1, 2$.
 - For $k = 2$ two theories $U(N)_2 \times U(N)_{-2}$ ABJM and $U(N)_2 \times U(N+1)_{-2}$ ABJ.
 - For $k = 1$, theory has free and interacting parts.

3d ABJM stress tensor correlator

- Expand stress tensor superprimary correlator in superblocks:

$$\langle S^a(x_1) S^b(x_2) S^c(x_3) S^d(x_4) \rangle = \frac{1}{x_{12}^2 x_{34}^2} \left[\sum_{\Delta, \ell} \lambda_{\Delta, \ell}^2 G_{\Delta, \ell}^{abcd}(x_i) + \sum_M \lambda_M^2 G_M^{abcd}(x_i) \right]$$

- Susy ward identities solved by superblocks, written as linear combos of 3d blocks in terms of structures $\sum_{i=1}^6 T_i^{abcd} \frac{G^i(U, V)}{x_{12}^2 x_{34}^2}$.
- Protected multiplets are nontrivial, include $(A, 2)$ for odd spin, $(A, +)$ for even spin, $(B, 2)$, $(B, +)$ and stress tensor for spin zero. Can bootstrap both upper and lower bounds (i.e. islands)!
- Theory specified by $c \sim N^{3/2}$ and $k = 1, 2$.
 - For $k = 2$ two theories $U(N)_2 \times U(N)_{-2}$ ABJM and $U(N)_2 \times U(N+1)_{-2}$ ABJ.
 - For $k = 1$, theory has free and interacting parts.

3d ABJM stress tensor correlator

- Expand stress tensor superprimary correlator in superblocks:

$$\langle S^a(x_1) S^b(x_2) S^c(x_3) S^d(x_4) \rangle = \frac{1}{x_{12}^2 x_{34}^2} \left[\sum_{\Delta, \ell} \lambda_{\Delta, \ell}^2 G_{\Delta, \ell}^{abcd}(x_i) + \sum_M \lambda_M^2 G_M^{abcd}(x_i) \right]$$

- Susy ward identities solved by superblocks, written as linear combos of 3d blocks in terms of structures $\sum_{i=1}^6 T_i^{abcd} \frac{G^i(U, V)}{x_{12}^2 x_{34}^2}$.
- Protected multiplets are nontrivial, include $(A, 2)$ for odd spin, $(A, +)$ for even spin, $(B, 2)$, $(B, +)$ and stress tensor for spin zero. Can bootstrap both upper and lower bounds (i.e. islands)!
- Theory specified by $c \sim N^{3/2}$ and $k = 1, 2$.
 - For $k = 2$ two theories $U(N)_2 \times U(N)_{-2}$ ABJM and $U(N)_2 \times U(N+1)_{-2}$ ABJ.
 - For $k = 1$, theory has free and interacting parts.

3d ABJM stress tensor correlator

- Expand stress tensor superprimary correlator in superblocks:

$$\langle S^a(x_1) S^b(x_2) S^c(x_3) S^d(x_4) \rangle = \frac{1}{x_{12}^2 x_{34}^2} \left[\sum_{\Delta, \ell} \lambda_{\Delta, \ell}^2 G_{\Delta, \ell}^{abcd}(x_i) + \sum_M \lambda_M^2 G_M^{abcd}(x_i) \right]$$

- Susy ward identities solved by superblocks, written as linear combos of 3d blocks in terms of structures $\sum_{i=1}^6 T_i^{abcd} \frac{G^i(U, V)}{x_{12}^2 x_{34}^2}$.
- Protected multiplets are nontrivial, include $(A, 2)$ for odd spin, $(A, +)$ for even spin, $(B, 2)$, $(B, +)$ and stress tensor for spin zero. Can bootstrap both upper and lower bounds (i.e. islands)!
- Theory specified by $c \sim N^{3/2}$ and $k = 1, 2$.
 - For $k = 2$ two theories $U(N)_2 \times U(N)_{-2}$ ABJM and $U(N)_2 \times U(N+1)_{-2}$ ABJ.
 - For $k = 1$, theory has free and interacting parts.

3d ABJM stress tensor correlator

- Expand stress tensor superprimary correlator in superblocks:

$$\langle S^a(x_1) S^b(x_2) S^c(x_3) S^d(x_4) \rangle = \frac{1}{x_{12}^2 x_{34}^2} \left[\sum_{\Delta, \ell} \lambda_{\Delta, \ell}^2 G_{\Delta, \ell}^{abcd}(x_i) + \sum_M \lambda_M^2 G_M^{abcd}(x_i) \right]$$

- Susy ward identities solved by superblocks, written as linear combos of 3d blocks in terms of structures $\sum_{i=1}^6 T_i^{abcd} \frac{G^i(U, V)}{x_{12}^2 x_{34}^2}$.
- Protected multiplets are nontrivial, include $(A, 2)$ for odd spin, $(A, +)$ for even spin, $(B, 2)$, $(B, +)$ and stress tensor for spin zero. Can bootstrap both upper and lower bounds (i.e. islands)!
- Theory specified by $c \sim N^{3/2}$ and $k = 1, 2$.
 - For $k = 2$ two theories $U(N)_2 \times U(N)_{-2}$ ABJM and $U(N)_2 \times U(N+1)_{-2}$ ABJ.
 - For $k = 1$, theory has free and interacting parts.

Localization constraints

- Derivatives of $F(m_i)$ deformed by three masses m_i relate to S^3 integrals of correlator [Agmon, SMC, Pufu '17; Binder, SMC, Pufu '18]:

$$\partial_{m_{\pm}}^4 F|_{m=0} = -\frac{\pi^4 c_T^2}{2^{13}} [4G^1(1, 4) + 2G^2(4, 1) + 16G^4(1, 4) - G^4(4, 1)],$$

$$\partial_{m_+}^2 \partial_{m_-} F|_{m=0} = -\frac{\pi c_T^2}{2^{12}} \int d^3 \vec{x} \frac{G^1(|1 - \vec{x}|^4, |\vec{x}|^2)}{|1 - \vec{x}|^2 |\vec{x}|^2}, \quad m_{\pm} \equiv m_1 \pm m_2.$$

- Integral trivializes in first constraint due to 1d theory [Yacoby, SMC, Pufu '14], can be written in terms of $\lambda_{(B,2)}^2$, $\lambda_{(B,+)}^2$, and c_T .
- LHS can be computed as function of N and k in terms of N^2 dimensional matrix model from localization [Kapustin, Willett, Yaakov '09], solved for finite N, k [Nosaka '23].

Localization constraints

- Derivatives of $F(m_i)$ deformed by three masses m_i relate to S^3 integrals of correlator [Agmon, SMC, Pufu '17; Binder, SMC, Pufu '18]:

$$\partial_{m_{\pm}}^4 F|_{m=0} = -\frac{\pi^4 c_T^2}{2^{13}} [4G^1(1, 4) + 2G^2(4, 1) + 16G^4(1, 4) - G^4(4, 1)],$$

$$\partial_{m_+}^2 \partial_{m_-} F|_{m=0} = -\frac{\pi c_T^2}{2^{12}} \int d^3 \vec{x} \frac{G^1(|1 - \vec{x}|^4, |\vec{x}|^2)}{|1 - \vec{x}|^2 |\vec{x}|^2}, \quad m_{\pm} \equiv m_1 \pm m_2.$$

- Integral trivializes in first constraint due to 1d theory [Yacoby, SMC, Pufu '14], can be written in terms of $\lambda_{(B,2)}^2$, $\lambda_{(B,+)}^2$, and c_T .
- LHS can be computed as function of N and k in terms of N^2 dimensional matrix model from localization [Kapustin, Willett, Yaakov '09], solved for finite N, k [Nosaka '23].

Localization constraints

- Derivatives of $F(m_i)$ deformed by three masses m_i relate to S^3 integrals of correlator [Agmon, SMC, Pufu '17; Binder, SMC, Pufu '18]:

$$\partial_{m_{\pm}}^4 F|_{m=0} = -\frac{\pi^4 c_T^2}{2^{13}} [4G^1(1, 4) + 2G^2(4, 1) + 16G^4(1, 4) - G^4(4, 1)],$$

$$\partial_{m_+}^2 \partial_{m_-} F|_{m=0} = -\frac{\pi c_T^2}{2^{12}} \int d^3 \vec{x} \frac{G^1(|1 - \vec{x}|^4, |\vec{x}|^2)}{|1 - \vec{x}|^2 |\vec{x}|^2}, \quad m_{\pm} \equiv m_1 \pm m_2.$$

- Integral trivializes in first constraint due to 1d theory [Yacoby, SMC, Pufu '14], can be written in terms of $\lambda_{(B,2)}^2$, $\lambda_{(B,+)}^2$, and c_T .
- LHS can be computed as function of N and k in terms of N^2 dimensional matrix model from localization [Kapustin, Willett, Yaakov '09], solved for finite N, k [Nosaka '23].

Localization constraints

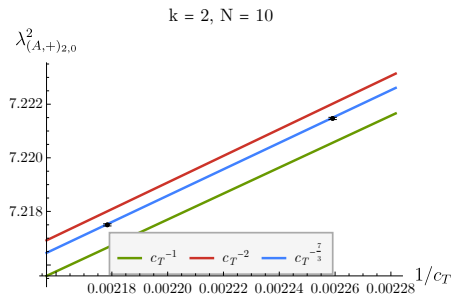
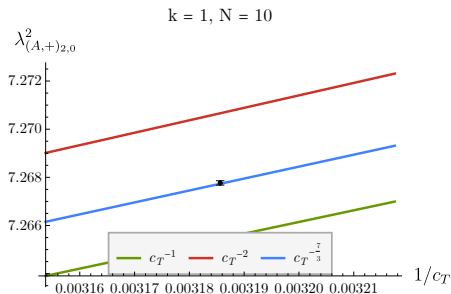
- Derivatives of $F(m_i)$ deformed by three masses m_i relate to S^3 integrals of correlator [Agmon, SMC, Pufu '17; Binder, SMC, Pufu '18]:

$$\partial_{m_{\pm}}^4 F|_{m=0} = -\frac{\pi^4 c_T^2}{2^{13}} [4G^1(1, 4) + 2G^2(4, 1) + 16G^4(1, 4) - G^4(4, 1)],$$

$$\partial_{m_+}^2 \partial_{m_-} F|_{m=0} = -\frac{\pi c_T^2}{2^{12}} \int d^3 \vec{x} \frac{G^1(|1 - \vec{x}|^4, |\vec{x}|^2)}{|1 - \vec{x}|^2 |\vec{x}|^2}, \quad m_{\pm} \equiv m_1 \pm m_2.$$

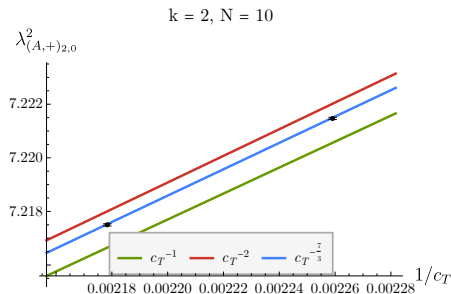
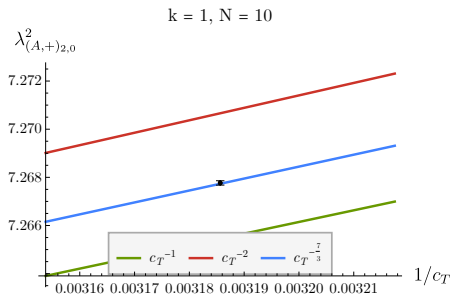
- Integral trivializes in first constraint due to 1d theory [Yacoby, SMC, Pufu '14], can be written in terms of $\lambda_{(B,2)}^2$, $\lambda_{(B,+)}^2$, and c_T .
- LHS can be computed as function of N and k in terms of N^2 dimensional matrix model from localization [Kapustin, Willett, Yaakov '09], solved for finite N, k [Nosaka '23].

Bootstrap islands: Sensitivity to M-theory corrections



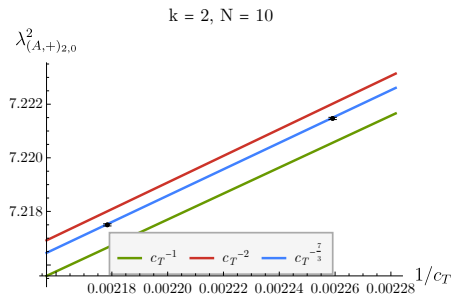
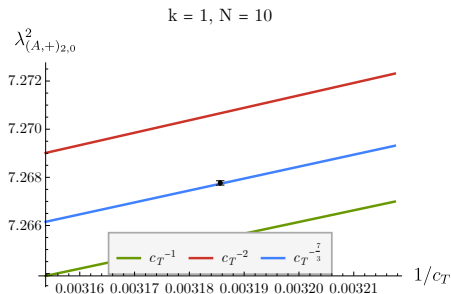
- For largish N , we see that analytic bootstrap result for $k = 1, 2$ gets closer to island as we include more $1/c_T$ corrections.
- $1/c_T$ is supergravity [SMC '18], $1/c_T^2$ is 1-loop correction [Alday, SMC, Raj '21, '22] (which included contact term fixed from localization), $1/c_T^{7/3}$ is $D^6 R^4$ correction [Binder, SMC, Pufu '18].
- So bootstrap sensitive to ALL protected M-theory corrections!

Bootstrap islands: Sensitivity to M-theory corrections



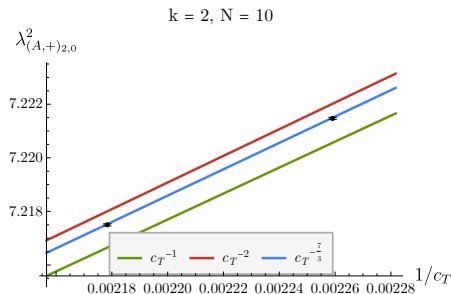
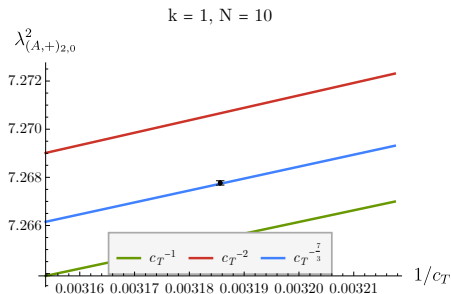
- For largish N , we see that analytic bootstrap result for $k = 1, 2$ gets closer to island as we include more $1/c_T$ corrections.
- $1/c_T$ is supergravity [SMC '18], $1/c_T^2$ is 1-loop correction [Alday, SMC, Raj '21, '22] (which included contact term fixed from localization), $1/c_T^{7/3}$ is $D^6 R^4$ correction [Binder, SMC, Pufu '18].
- So bootstrap sensitive to ALL protected M-theory corrections!

Bootstrap islands: Sensitivity to M-theory corrections



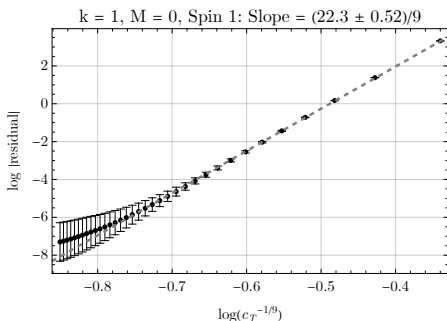
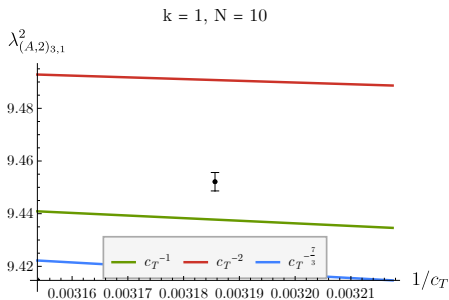
- For largish N , we see that analytic bootstrap result for $k = 1, 2$ gets closer to island as we include more $1/c_T$ corrections.
- $1/c_T$ is supergravity [SMC '18], $1/c_T^2$ is 1-loop correction [Alday, SMC, Raj '21,'22] (which included contact term fixed from localization), $1/c_T^{7/3}$ is $D^6 R^4$ correction [Binder, SMC, Pufu '18].
- So bootstrap sensitive to ALL protected M-theory corrections!

Bootstrap islands: Sensitivity to M-theory corrections



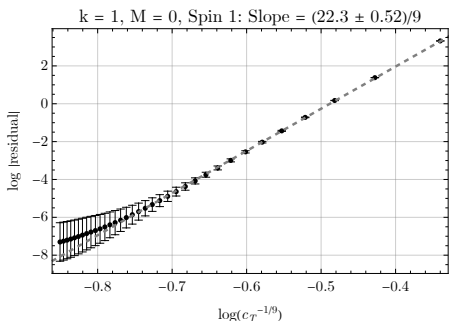
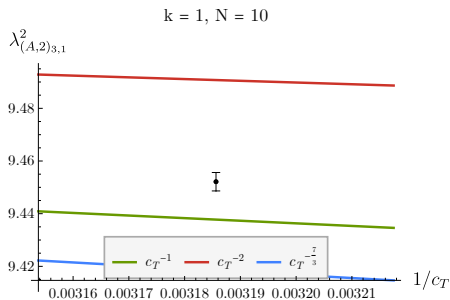
- For largish N , we see that analytic bootstrap result for $k = 1, 2$ gets closer to island as we include more $1/c_T$ corrections.
- $1/c_T$ is supergravity [SMC '18], $1/c_T^2$ is 1-loop correction [Alday, SMC, Raj '21,'22] (which included contact term fixed from localization), $1/c_T^{7/3}$ is $D^6 R^4$ correction [Binder, SMC, Pufu '18].
- So bootstrap sensitive to ALL protected M-theory corrections!

Bootstrap islands: Unprotected $D^8 R^4$ term



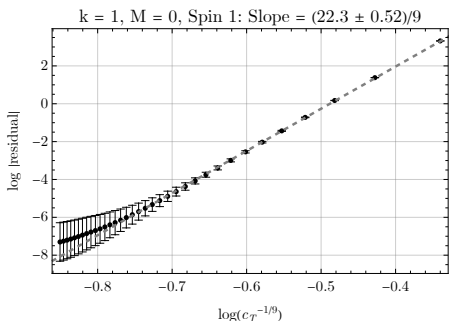
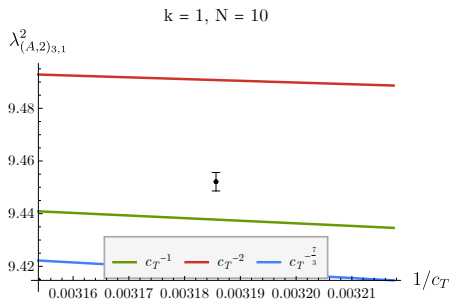
- Island is sensitive to corrections beyond $D^6 R^4$ at $c_T^{-7/3}$, i.e. the first unprotected correction $D^8 R^4$ at $c_T^{-23/9}$.
- The residual between bootstrap data and $O(c_T^{-7/3})$ expansion is consistent with $c_T^{-23/9}$ correction! Similar plots for $k = 2$.

Bootstrap islands: Unprotected $D^8 R^4$ term



- Island is sensitive to corrections beyond $D^6 R^4$ at $c_T^{-7/3}$, i.e. the first unprotected correction $D^8 R^4$ at $c_T^{-23/9}$.
- The residual between bootstrap data and $O(c_T^{-7/3})$ expansion is consistent with $c_T^{-23/9}$ correction! Similar plots for $k = 2$.

Bootstrap islands: Unprotected $D^8 R^4$ term



- Island is sensitive to corrections beyond $D^6 R^4$ at $c_T^{-7/3}$, i.e. the first unprotected correction $D^8 R^4$ at $c_T^{-23/9}$.
- The residual between bootstrap data and $O(c_T^{-7/3})$ expansion is consistent with $c_T^{-23/9}$ correction! Similar plots for $k = 2$.

Conclusion

- Bounds from bootstrap+localization give non-perturbative solution to holographic CFTs at finite N and coupling!
 - For $\mathcal{N} = 4$ SYM, evidence from comparing upper bound to both weak and strong coupling.
 - For ABJM, have rigorous upper and lower bounds (i.e. islands), and also matches strong coupling.
- Bounds accurate enough to read off lowest few protected corrections from string/M-theory to supergravity.
- For ABJM, can also read off prediction for unprotected $D^8 R^4$ correction to M-theory S-matrix!

Conclusion

- Bounds from bootstrap+localization give non-perturbative solution to holographic CFTs at finite N and coupling!
 - For $\mathcal{N} = 4$ SYM, evidence from comparing upper bound to both weak and strong coupling.
 - For ABJM, have rigorous upper and lower bounds (i.e. islands), and also matches strong coupling.
- Bounds accurate enough to read off lowest few protected corrections from string/M-theory to supergravity.
- For ABJM, can also read off prediction for unprotected $D^8 R^4$ correction to M-theory S-matrix!

Conclusion

- Bounds from bootstrap+localization give non-perturbative solution to holographic CFTs at finite N and coupling!
 - For $\mathcal{N} = 4$ SYM, evidence from comparing upper bound to both weak and strong coupling.
 - For ABJM, have rigorous upper and lower bounds (i.e. islands), and also matches strong coupling.
- Bounds accurate enough to read off lowest few protected corrections from string/M-theory to supergravity.
- For ABJM, can also read off prediction for unprotected $D^8 R^4$ correction to M-theory S-matrix!

Conclusion

- Bounds from bootstrap+localization give non-perturbative solution to holographic CFTs at finite N and coupling!
 - For $\mathcal{N} = 4$ SYM, evidence from comparing upper bound to both weak and strong coupling.
 - For ABJM, have rigorous upper and lower bounds (i.e. islands), and also matches strong coupling.
- Bounds accurate enough to read off lowest few protected corrections from string/M-theory to supergravity.
- For ABJM, can also read off prediction for unprotected $D^8 R^4$ correction to M-theory S-matrix!

Conclusion

- Bounds from bootstrap+localization give non-perturbative solution to holographic CFTs at finite N and coupling!
 - For $\mathcal{N} = 4$ SYM, evidence from comparing upper bound to both weak and strong coupling.
 - For ABJM, have rigorous upper and lower bounds (i.e. islands), and also matches strong coupling.
- Bounds accurate enough to read off lowest few protected corrections from string/M-theory to supergravity.
- For ABJM, can also read off prediction for unprotected $D^8 R^4$ correction to M-theory S-matrix!

Future directions

- More accurate bounds \Rightarrow more unprotected corrections.
 - Get greater accuracy from imposing more localization constraints (e.g. from the squashed sphere), or from mixing with lowest dimension long operator (which is also relevant).
- Combine localization + bootstrap to numerically solve ANY 3d $\mathcal{N} = 2$, 4d $\mathcal{N} = 2$, or 5d $\mathcal{N} = 1$ Lagrangian CFT, e.g.:
 - 4d $\mathcal{N} = 2$ dual to open strings [SMC '22; Behan, SMC, Ferrero '23] .
 - 3d $\mathcal{N} = 6$ ABJ(M) in string, M-theory, and higher spin regimes [Binder, SMC, Jerdee, Pufu '20] .
- Once we can access $\Delta \sim c$ higher twist operator, can study statistics of black hole states as function of c and coupling.

Future directions

- More accurate bounds \Rightarrow more unprotected corrections.
 - Get greater accuracy from imposing more localization constraints (e.g. from the squashed sphere), or from mixing with lowest dimension long operator (which is also relevant).
- Combine localization + bootstrap to numerically solve ANY 3d $\mathcal{N} = 2$, 4d $\mathcal{N} = 2$, or 5d $\mathcal{N} = 1$ Lagrangian CFT, e.g.:
 - 4d $\mathcal{N} = 2$ dual to open strings [SMC '22; Behan, SMC, Ferrero '23] .
 - 3d $\mathcal{N} = 6$ ABJ(M) in string, M-theory, and higher spin regimes [Binder, SMC, Jerdee, Pufu '20] .
- Once we can access $\Delta \sim c$ higher twist operator, can study statistics of black hole states as function of c and coupling.

Future directions

- More accurate bounds \Rightarrow more unprotected corrections.
 - Get greater accuracy from imposing more localization constraints (e.g. from the squashed sphere), or from mixing with lowest dimension long operator (which is also relevant).
- Combine localization + bootstrap to numerically solve ANY 3d $\mathcal{N} = 2$, 4d $\mathcal{N} = 2$, or 5d $\mathcal{N} = 1$ Lagrangian CFT, e.g.:
 - 4d $\mathcal{N} = 2$ dual to open strings [SMC '22; Behan, SMC, Ferrero '23] .
 - 3d $\mathcal{N} = 6$ ABJ(M) in string, M-theory, and higher spin regimes [Binder, SMC, Jerdee, Pufu '20] .
- Once we can access $\Delta \sim c$ higher twist operator, can study statistics of black hole states as function of c and coupling.

Future directions

- More accurate bounds \Rightarrow more unprotected corrections.
 - Get greater accuracy from imposing more localization constraints (e.g. from the squashed sphere), or from mixing with lowest dimension long operator (which is also relevant).
- Combine localization + bootstrap to numerically solve ANY 3d $\mathcal{N} = 2$, 4d $\mathcal{N} = 2$, or 5d $\mathcal{N} = 1$ Lagrangian CFT, e.g.:
 - 4d $\mathcal{N} = 2$ dual to open strings [SMC '22; Behan, SMC, Ferrero '23] .
 - 3d $\mathcal{N} = 6$ ABJ(M) in string, M-theory, and higher spin regimes [Binder, SMC, Jerdee, Pufu '20] .
- Once we can access $\Delta \sim c$ higher twist operator, can study statistics of black hole states as function of c and coupling.

Future directions

- More accurate bounds \Rightarrow more unprotected corrections.
 - Get greater accuracy from imposing more localization constraints (e.g. from the squashed sphere), or from mixing with lowest dimension long operator (which is also relevant).
- Combine localization + bootstrap to numerically solve ANY 3d $\mathcal{N} = 2$, 4d $\mathcal{N} = 2$, or 5d $\mathcal{N} = 1$ Lagrangian CFT, e.g.:
 - 4d $\mathcal{N} = 2$ dual to open strings [SMC '22; Behan, SMC, Ferrero '23] .
 - 3d $\mathcal{N} = 6$ ABJ(M) in string, M-theory, and higher spin regimes [Binder, SMC, Jerdee, Pufu '20] .
- Once we can access $\Delta \sim c$ higher twist operator, can study statistics of black hole states as function of c and coupling.

Future directions

- More accurate bounds \Rightarrow more unprotected corrections.
 - Get greater accuracy from imposing more localization constraints (e.g. from the squashed sphere), or from mixing with lowest dimension long operator (which is also relevant).
- Combine localization + bootstrap to numerically solve ANY 3d $\mathcal{N} = 2$, 4d $\mathcal{N} = 2$, or 5d $\mathcal{N} = 1$ Lagrangian CFT, e.g.:
 - 4d $\mathcal{N} = 2$ dual to open strings [SMC '22; Behan, SMC, Ferrero '23] .
 - 3d $\mathcal{N} = 6$ ABJ(M) in string, M-theory, and higher spin regimes [Binder, SMC, Jerdee, Pufu '20] .
- Once we can access $\Delta \sim c$ higher twist operator, can study statistics of black hole states as function of c and coupling.

See you in London!

Solving Holographic Theories

December 16-20, 2024
Imperial College London

<https://indico.cern.ch/e/sol-holography>

Speakers:

Luis Fernando Alday
Nikolay Bobev
Minjae Cho
Ross Dempsey
Lorenz Eberhardt
Barak Gabai
Rajesh Gopakumar
Michael Green
Nikolay Gromov
Sean Hartnoll
Manki Kim
Shota Komatsu
Henry Lin
Edward Mazenc
Sameer Murthy
Silviu Pufu
Mukund Rangamani
Ning Su
Arkady Tseytlin
Aron Wall
Congkao Wen
Zahra Zahraee
Zechuan Zheng

Organizers:

Shai Chester
Xi Yin
Deliang Zhong

