

Real-time dynamics of pseudo-Goldstone and critical mode in QCD

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Collaborate with Yong-rui Chen & Shi Yin & Chuang Huang & Wei-jie Fu & Wei-jia Li.

Based on: **arXiv:2403.03503**

arXiv:24xx.xxxxx in preparation

July 19, 2024

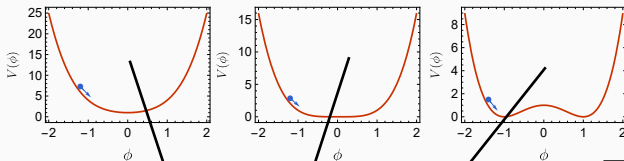
Dalian University of Technology



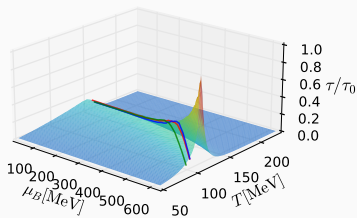
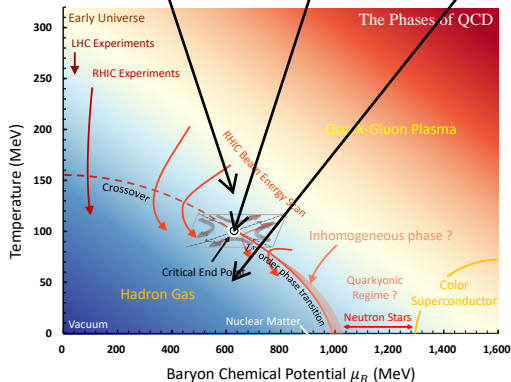
1. Introduction
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4. Summary

Introduction

QCD Phase Diagram



Critical slowing down
with the relaxation time:
 $\tau = \xi^z f(k\xi)$
 z : dynamic critical exponent



Tan, Yin, Chen, Huang, WF, in preparation

- Nonperturbative approach
- Real-time description

Functional renormalization group approach

$$Z_k[J] = \int \mathcal{D}\varphi \exp(-S[\varphi] + J \cdot \varphi - \Delta S_k[\varphi])$$

where $\Delta S_k[\varphi] = \frac{1}{2} \int \frac{d^d q}{(2\pi)^d} \varphi(-q) R_k(q) \varphi(q)$

The scale dependent effective action:

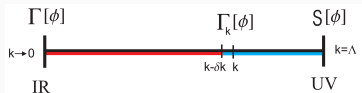
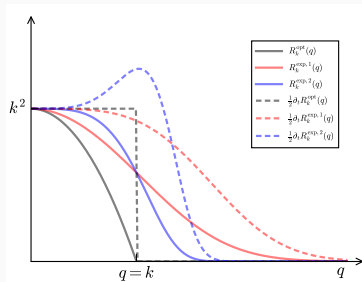
$$\Gamma_k[\phi] = -\ln Z_k[J] + \int d^d x J(x) \phi(x) - \Delta S_k[\phi]$$

Wetterich Equation :

$$k \partial_k \Gamma_k[\phi] = \frac{1}{2} \text{STr} \left[\left(\Gamma_k^{(2)}[\phi] + R_k \right)^{-1} k \partial_k R_k \right]$$

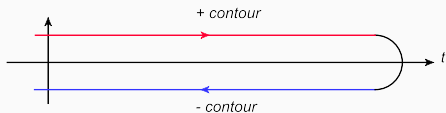
The scale evolution equation for n- point function (one loop exact)

$$\partial_t \Gamma_k^{(n)}[\phi] = \tilde{\partial}_t \left(\begin{array}{l} \text{all one-loop correction} \\ \text{diagrams of } \Gamma_k^{(n)}[\phi] \end{array} \right)$$



- Derivative expansion
- Vertex expansion

fRG in Keldysh path integral



$$\varphi_c = \frac{1}{\sqrt{2}} (\varphi_+ + \varphi_-), \quad \varphi_q = \frac{1}{\sqrt{2}} (\varphi_+ - \varphi_-)$$

The flow equation in the closed time path :

$$\begin{aligned} \partial_\tau \Gamma_k[\Phi] &= \frac{i}{2} \text{Tr} \left[(\partial_\tau R_k^*) (\Gamma_k^{(2)}[\Phi] + R_k)^{-1} \right] \\ &\equiv \frac{i}{2} \text{Tr} [(\partial_\tau R_k^*) G_k] \end{aligned}$$

- Propagator

$$G_k = \begin{pmatrix} G_k^K & G_k^R \\ G_k^A & 0 \end{pmatrix}$$

- Regulator

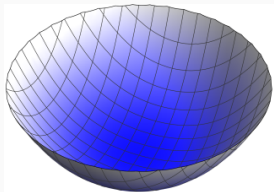
$$R_k = \begin{pmatrix} 0 & R_k^A \\ R_k^R & R_k^F \end{pmatrix}$$

$G^R = (G^A)^*$ and the Regulator $R_k^F = 0$. We use frequency-independent flat regulator to maintain the causal structure:

$$R_k^R(\omega, \mathbf{q}) = R_k^A(\omega, \mathbf{q}) = -Z_k(k^2 - \mathbf{q}^2)\theta(k^2 - \mathbf{q}^2)$$

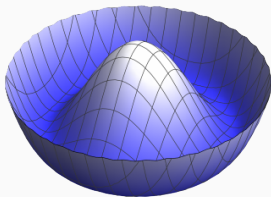
Damping of pseudo-Goldstone near critical points

Symmetry breaking of chiral phase transition



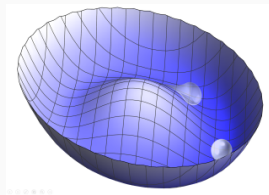
$$SU(2)_L \times SU(2)_R \simeq O(4)$$

Symmetric phase
Two-flavor QCD with
 $m_q \rightarrow 0$



$$O(4) \rightarrow O(3) \times Z_2 (\vec{\pi}, \sigma)$$

Spontaneous symmetry
breaking
 π - Goldstone bosons
with $m_\pi = 0$



Approximate symmetry
spontaneous breaking
 m_q is small but $\neq 0$
or $c \neq 0$
 π - pseudo-Goldstone
bosons with $m_\pi \neq 0$

$O(N)$ Effective action on the S-K contour

$$\Gamma[\phi_c, \phi_q] = \int d^4x \left(Z_a^{(t)} \phi_{a,q} \partial_t \phi_{a,c} - Z_a^{(i)} \phi_{a,q} \partial_i^2 \phi_{a,c} + V'(\rho_c) \phi_{a,q} \phi_{a,c} - 2 Z_a^{(t)} T \phi_{a,q}^2 - \sqrt{2} c \sigma_q \right)$$

Universal damping of pseudo-Goldstone

The real-time dispersion relation of pseudo-Goldstone:

$$\omega(\mathbf{p}) = \pm c_s \sqrt{m_\varphi^2 + \mathbf{p}^2} - \frac{i}{2} [\Omega_\varphi + (D_Q + D_\varphi) \mathbf{p}^2] + \dots$$

Ω_φ : relaxation/damping rate, D_φ : Goldstone diffusivity,

D_Q : charge diffusivity, m_φ : screening mass, c_s : speed of sound

$$T = 0.3 T_c$$

Universal relation for pseudo-Goldstone damping

$$\frac{\Omega_\varphi}{m_\varphi^2} \simeq D_\varphi + \mathcal{O}(m_\varphi^2/T^2)$$

Holographics:

Amoretti, Areán, Goutéraux, Musso, PRL 123 (2019) 211602;

Amoretti, Areán, Goutéraux, Musso, JHEP 10 (2019) 068;

Ammon et al., JHEP 03 (2022) 015;

Cao, Baggioli, Liu, Li, JHEP 12 (2022) 113

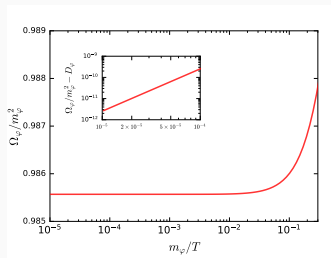
Hydrodynamics:

Delacrétaz, Goutéraux, Ziogas, PRL 128 (2022) 141601

EFT:

Baggioli, Phys. Rev. Res. 2 (2020) 022022;

Baggioli, Landry, SciPost Phys. 9 (2020) 062



$$\frac{\Omega_\varphi}{m_\varphi^2} - D_\varphi(T) \propto \left(\frac{m_\varphi}{T}\right)^\alpha$$

$$\alpha = 2.00021(2)$$

Universal damping of pseudo-Goldstone near the critical point

The relaxation/damping rate in Model A:

$$G_{\varphi\varphi}^R(\omega, q) = \frac{1}{-iZ_{\varphi}^{(t)}\omega + Z_{\varphi}^{(i)}(q^2 + m_{\varphi}^2)} \rightarrow \frac{\Omega_{\varphi}}{m_{\varphi}^2} = \frac{Z_{\varphi}^{(i)}}{Z_{\varphi}^{(t)}}$$

$$\frac{\Omega_{\varphi}}{m_{\varphi}^2} \propto m_{\varphi}^{\Delta_{\eta}}, \quad \Delta_{\eta} \equiv \eta_t - \eta > 0$$

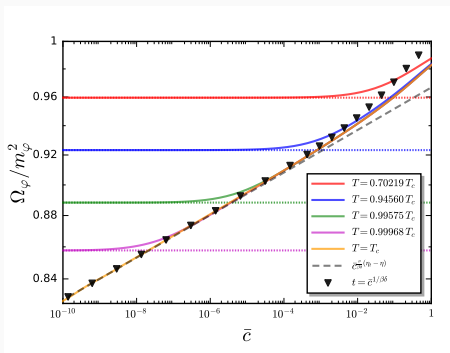
$$\Delta_{\eta} \simeq 0.0172$$

GMOR also breaks down with (only relates to the static properties):

$$m_{\varphi}^2 \propto c^{\frac{2\nu}{\beta\delta}} \simeq c^{0.806}$$

mean-field theory with:

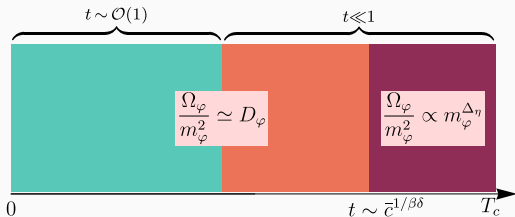
$$m_{\varphi}^2 \simeq c^{2/3}$$



Critical region of chiral phase transition:

$$m_{\pi 0} \lesssim 0.1 \sim 1 \text{ MeV}$$

Universal damping of pseudo-Goldstone near the critical point

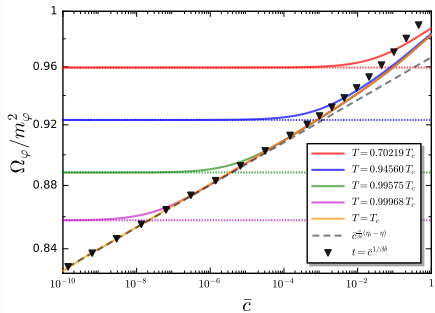


$$\frac{\Omega_\varphi}{m_\varphi^2} \propto m_\varphi^{\Delta_\eta}, \quad \Delta_\eta \equiv \eta_t - \eta > 0$$

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Large N limit

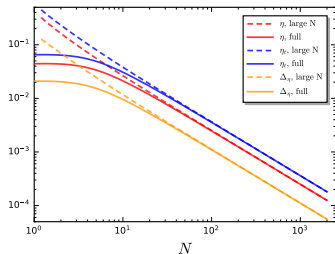
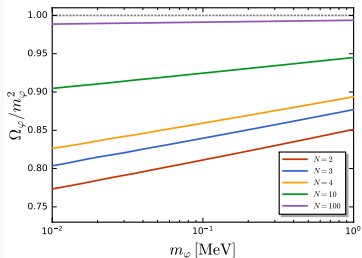
Neglect sigma mode in the fRG flow equation we have:

$$\eta = \frac{5}{N-1} \frac{(1+\eta)(1-2\eta)^2}{(5-\eta)(2-\eta)^2}$$

$$\eta_t = \frac{1}{9(N-1)} \frac{(1-2\eta)^2 (13+15\eta-2\eta^3)}{(2-\eta)^2}$$

In the large N limit

$\eta \propto 1/N, \eta_t \propto 1/N$ and $\Delta_\eta \propto 1/N$



Relaxation dynamics of critical mode near CEP

QCD-assisted Langevin transport model

Effective action on the S-K contour

$$\Gamma[\phi_c, \phi_q] = \int d^4x \left(Z_\phi^{(t)} \sigma_q \partial_t \sigma_c - Z_\phi^{(i)} \sigma_q \partial_i^2 \sigma_c + 2U'(\sigma_c) \sigma_q - 2Z_\phi^{(t)} T \sigma_q^2 \right)$$

Langevin equation:

$$Z_\phi^{(t)} \partial_t \sigma - Z_\phi^{(i)} \partial_i^2 \sigma + U'(\sigma) = \xi$$

with classical noise:

$$\langle \xi(t, \mathbf{x}) \xi(t', \mathbf{x}') \rangle = 2Z_\phi^{(t)} T \delta(t-t') \delta(\mathbf{x}-\mathbf{x}')$$

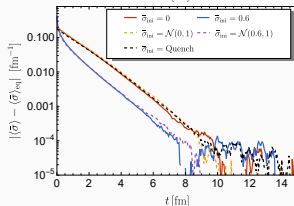
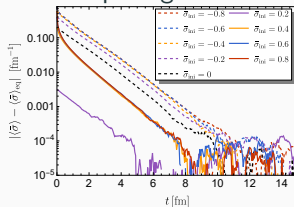
relaxation ansatz:

$$\langle \bar{\sigma}(t) \rangle = A e^{-\frac{t}{\tau}} + \langle \bar{\sigma}(t) \rangle_{\text{eq}}$$

Langevin simulations are performed on the GPUs using Julia with CUDA.jl and DifferentialEquations.jl packages.

Lattice size: $32 \times 32 \times 32$

Lattice spacing : 0.1fm



Inputs from first principal fRG-QCD

Flow of fRG-QCD (static properties):

$$\partial_t \Gamma_k[\Phi] = \frac{1}{2} \left(\text{orange loop} - \text{dotted loop} - \text{solid loop} + \frac{1}{2} \text{blue loop} \right)$$

Fu, Pawłowski, Rennecke,
PRD.101.054032

$$U'(\sigma) = \left. \frac{\delta \Gamma_{\text{QCD}}[\Phi]}{\delta \sigma} \right|_{\substack{\sigma(x)=\sigma \\ \tilde{\Phi}=\tilde{\Phi}_{\text{EoM}}}}$$

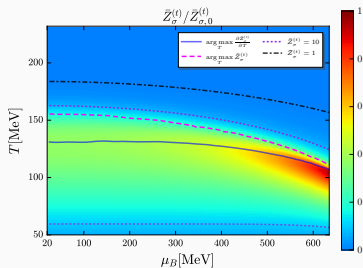
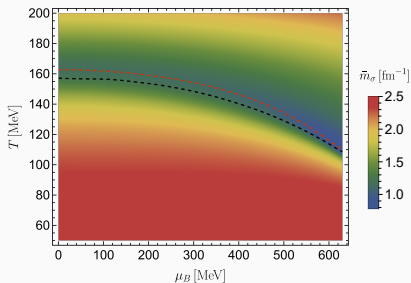
$$\Gamma_{\sigma\sigma}^{(2)} = \left. \frac{\delta^2 \Gamma_{\text{QCD}}[\Phi]}{\delta \sigma \delta \sigma} \right|_{\Phi=\Phi_{\text{EoM}}}, \quad Z_{\phi}^{(i)} = \left. \frac{\partial \Gamma_{\sigma\sigma}^{(2)}(p_0, \mathbf{p})}{\partial p^2} \right|_{\substack{p_0=0 \\ \mathbf{p}=0}}$$

Flow of two point function:

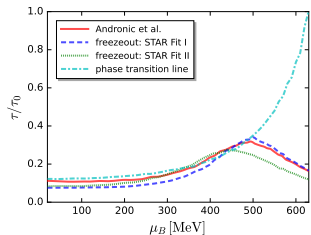
$$\partial_t \text{blob} = \tilde{\partial}_t \left(- \text{loop} + \frac{1}{2} \text{blue loop} + \frac{1}{2} \text{blue loop} \right)$$

$$\Gamma_{\sigma\sigma, \text{R}}^{(2)}(\omega, \mathbf{p}) = \lim_{\epsilon \rightarrow 0^+} \Gamma_{\sigma\sigma}^{(2)}(p_0 = -i(\omega + i\epsilon), \mathbf{p})$$

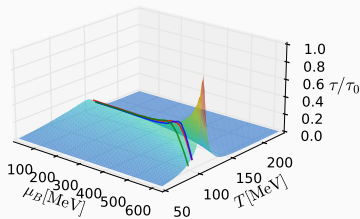
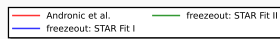
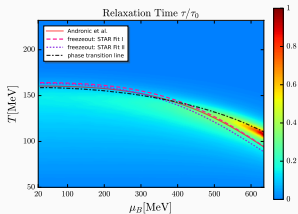
$$Z_{\phi}^{(t)} = \lim_{|\mathbf{p}| \rightarrow 0} \lim_{\omega \rightarrow 0} \frac{\partial}{\partial \omega} \text{Im} \Gamma_{\sigma\sigma, \text{R}}^{(2)}(\omega, \mathbf{p})$$



Relaxation time on the QCD phase diagram

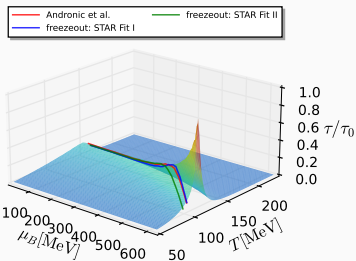
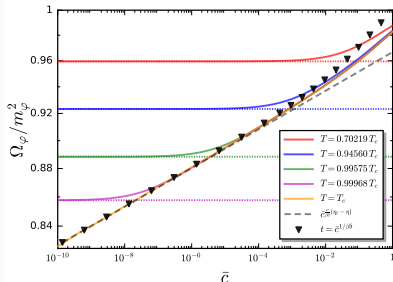


- Once away from the CEP the relaxation time quickly drops which indicates a small critical region around the CEP.
- The critical slowing down and nonequilibrium fluctuations have limited effects on the observables in the freezeout lines.



Summary

Summary



- The pseudo-Goldstone damping has a universal relation $\frac{\Omega_\varphi}{m_\varphi^2} \simeq D_\varphi + \mathcal{O}(m_\varphi^2/T^2)$ in the hydrodynamic regime while breaks down to $\frac{\Omega_\varphi}{m_\varphi^2} \propto m_\varphi^{\Delta_\eta}$, $\Delta_\eta > 0$ in the critical regime.
- The relaxation time quickly drops when away from the CEP and the relaxation time is relatively small in the freeze-out line thus the effect of critical slowing down is mild on the chemical freeze-out curve.