

Top-Quark Decay at Next-to-Next-to-Next-to-Leading Order in QCD

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In collaboration with Long Chen, Xin Guan, Yan-Qing Ma
Based on: [arXiv: 2309.01937v1]



1. Introduction: Motivation and Background

2. Computational Techniques

3. Numerical Results

4. Summary and Outlook

Top Quark

➤ Standard Model  New Physics?

➤ The heaviest fundamental particle in SM: **Top Quark** $m_t = 172.69 \pm 0.30 \text{ GeV}$

Particle Data Group. (2022)

➤ Relevant physical problems:

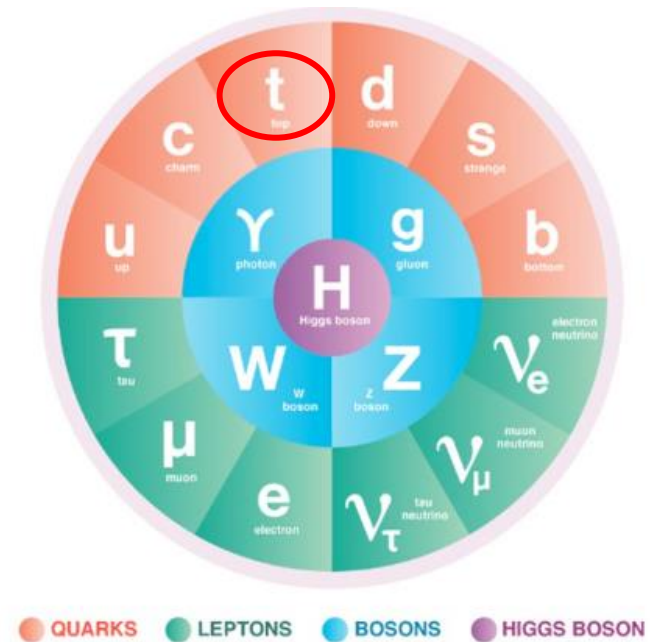
- The origin of Higgs mass
- The quark mixing
- CP violation...

➤ **Precision test of SM**, and probes for BSM physics

➤ $m_t \gg \Lambda_{QCD}$, allows us to use perturbative QCD

➤ Unique decay phenomenon:

- Heavy mass: **dominantly decay to W boson and b -quark**
- Before hadronization: the **information about its spin state** is preserved in distributions of top-quark decay products



Experimental Measurements for Mass and Decay Width

- Top quark mass m_t (PDG average)

$$m_t = 172.69 \pm 0.30 \text{ GeV}$$

Particle Data Group. (2022)

- Current best measurement for Γ_t :

$$\Gamma_t = 1.36 \pm 0.02(\text{stat.})_{-0.11}^{+0.14}(\text{syst.}) \text{ GeV} \quad \text{CMS. (2014)}$$

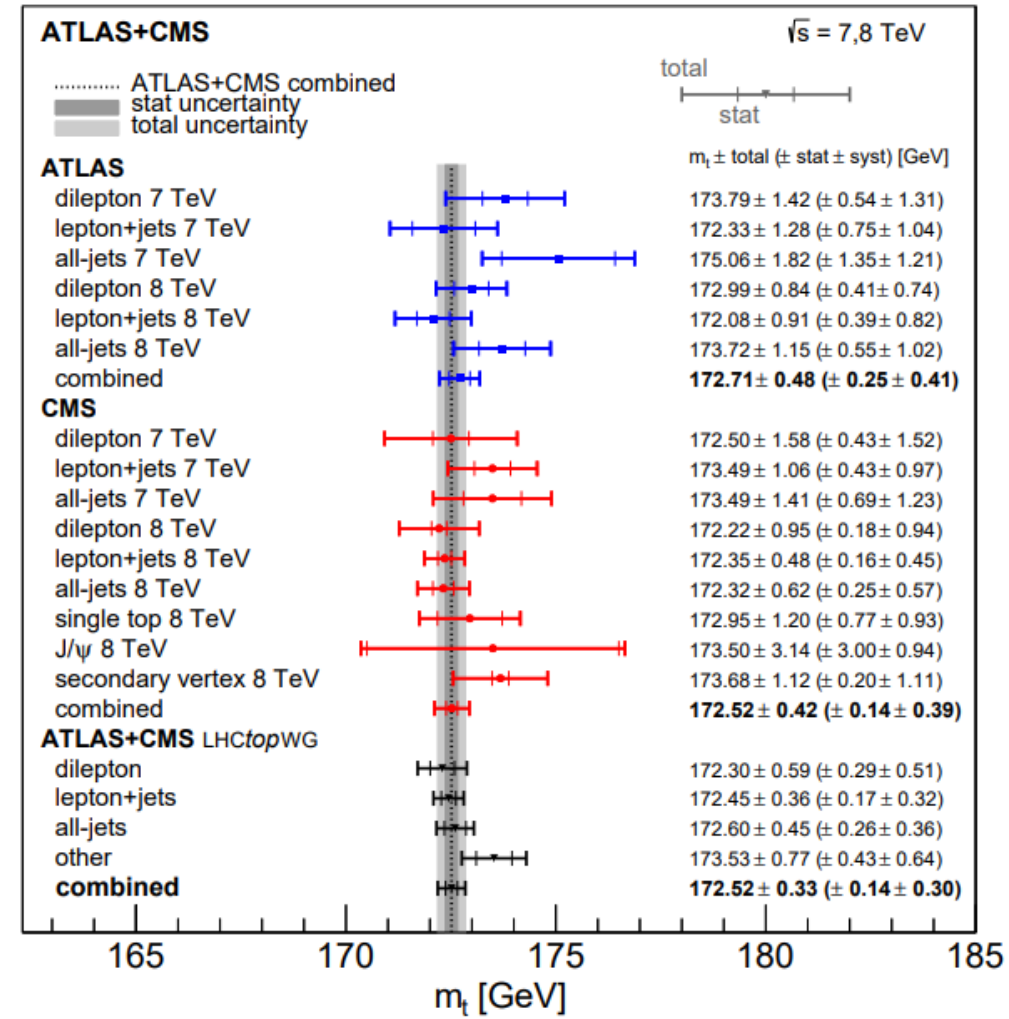
- Expected experimental precision at future colliders:

20~26 MeV

CLIC. (2019)

ILC. (2013)

CEPC. (2023)



From ATLAS, CMS. (2024)

Experimental Measurements for W -Helicity Fractions

➤ The W boson from decay process $t \rightarrow W + b$ is **polarized** even if the top quark is unpolarized.

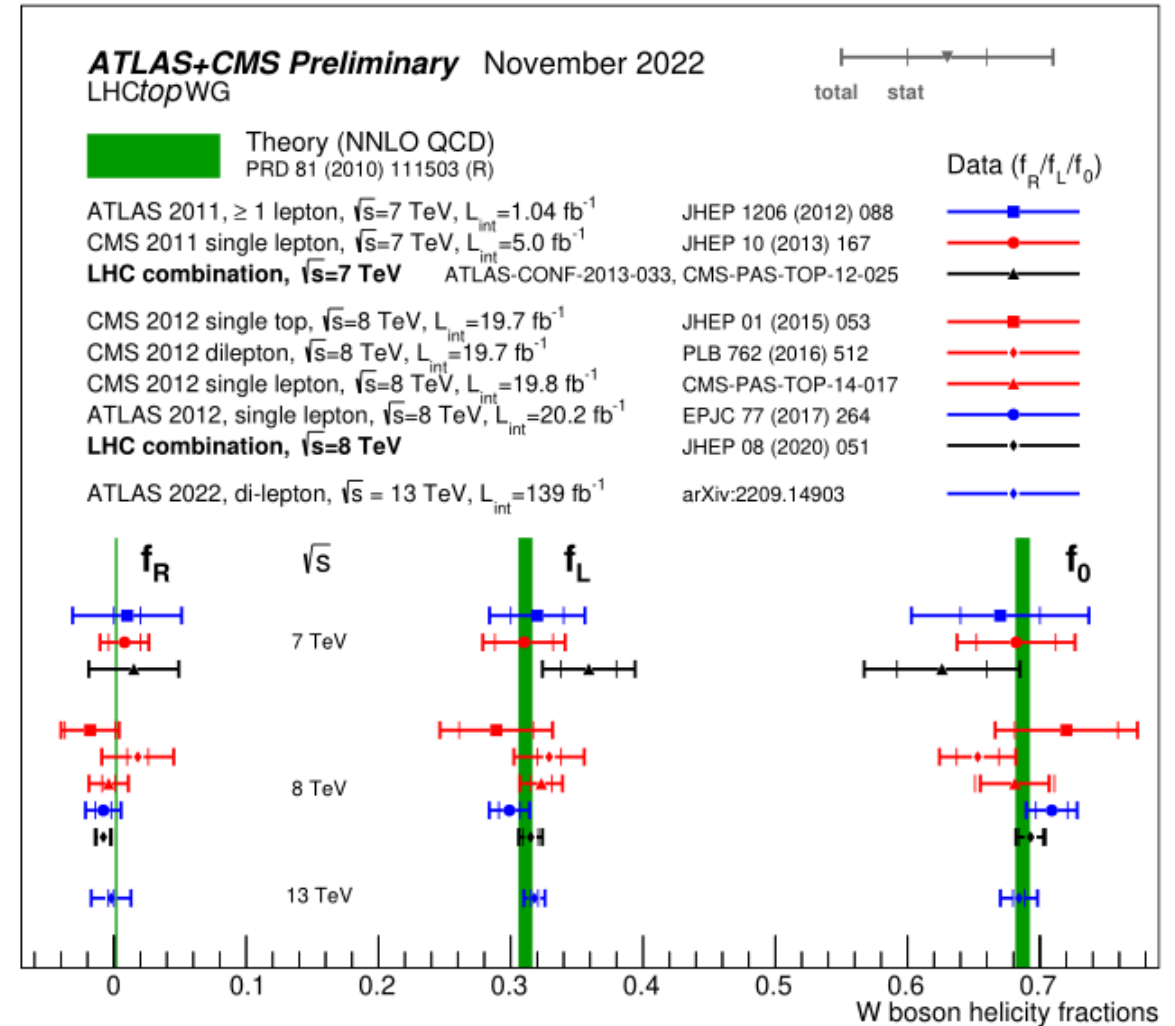
➤ The current best measurements for $f_\lambda (= \frac{\Gamma_\lambda}{\Gamma_t})$

$$f_0 = 0.684 \pm 0.005(\text{stat.}) \pm 0.014(\text{syst.})$$

$$f_L = 0.318 \pm 0.003(\text{stat.}) \pm 0.008(\text{syst.})$$

$$f_R = -0.002 \pm 0.002(\text{stat.}) \pm 0.014(\text{syst.})$$

ATLAS. (2022)



Theoretical Efforts in Literature

➤ Improvement of **experimental accuracy**



Need more **precise theoretical prediction**

➤ **Total decay width Γ_t**

- NLO(QCD)

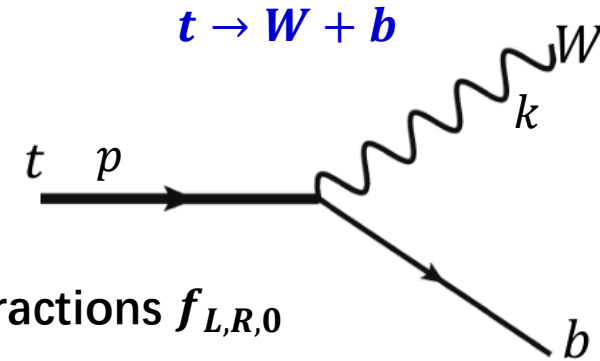
M. Jezabek, J. H. Kuhn. (1988)
A. Czarnecki. (1990)
C.-S. Li, R. J. Oakes, T. C. Yuan. (1990)

- NNLO(QCD)

A. Czarnecki, K. Melnikov. (1998)
K. G. Chetyrkin, R. Harlander, T. Seidensticker, M. Steinhauser. (1999)
M. Fischer, S. Groote, J. G. Korner, M. C. Mauser. (2001)
I. Blokland, A. Czarnecki, M. Slusarczyk, F. Tkachov. (2004)
I. Blokland, A. Czarnecki, M. Slusarczyk, F. Tkachov. (2005)
A. Czarnecki, J. G. Korner, J. H. Piclum. (2010)
J. Gao, C.-S. Li, H.-X. Zhu. (2012)
M. Brucherseifer, F. Caola, K. Melnikov. (2013)
R.-Q. Meng, S.-Q. Wang, T. Sun, C.-Q. Luo, J.-M. Shen, X.-G. Wu. (2022)
L.-C. Chen, H.-T. Li, J. Wang Y.-F. Wang. (2022)

- NLO(EW)

A. Denner, T. Sack. (1990)



➤ **W-helicity fractions $f_{L,R,0}$**

- NNLO(QCD)

A. Czarnecki, J. G. Korner, J. H. Piclum. (2010)
J. Gao, C.-S. Li, H.-X. Zhu. (2012)
M. Brucherseifer, F. Caola, K. Melnikov. (2013)

- NLO(EW)

H. S. Do, S. Groote, J. G. Korner, M. C. Mauser. (2002)

➤ **Differential distributions**

- NLO(QCD)

M. Fischer, S. Groote, J. G. Korner, M. C. Mauser. (2001)

- NNLO(QCD)

J. Gao, C.-S. Li, H.-X. Zhu. (2012)
M. Brucherseifer, F. Caola, K. Melnikov. (2013)

Theoretical uncertainty of Γ_t is about 1%(13MeV) at NNLO!

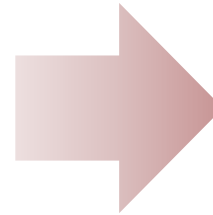
From L.-C. Chen, H.-T. Li, J. Wang Y.-F. Wang. (2022)

Full computation of NNNLO QCD correction is still indispensable!

Challenges for Perturbation Theory

➤ **Challenges** in high-order perturbative calculations:

- The number of Feynman diagrams increases rapidly
 - ✓ Loops
 - ✓ External legs
- The number of loops of the Feynman integrals increases
 - ✓ 1-loop: has been solved analytically
 - ✓ Beyond 1-loop: no systematic and analytic method for all situations
- The number of mass scales increases
 - ✓ Mass of particles in SM (W, Z, Higgs, etc)
 - ✓ Mandelstam variables



- Require **systematic** approach
- Difficult to calculate analytically, calculate **numerically**

Equipped with the-state-of-art techniques, the high-order corrections are **available now!!**

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Work-flow

➤ Work-flow for high-order computation:

Generating Integrands
(\mathcal{O} (few minutes to hours))

$$\mathcal{A} = \sum_{i=1}^{\mathcal{O}(10^4)} f_i \times I_i$$

Feynman diagrams



Integrals Reduction
(\mathcal{O} (hours to days))

$$I_i = \sum_{j=1}^{\mathcal{O}(10^2)} c_{ij} M_j$$

Integrate-By-Part (IBP) method



Evaluating master integrals
(\mathcal{O} (days) typically in frontier)

$$I_i(p_i \cdot p_j, m_i^2) = \int \frac{d^D l_1 \cdots d^D l_n}{D_1^{\nu_1} D_2^{\nu_2} \cdots D_N^{\nu_N}}$$

Differential equations

Formula

➤ To calculate the **semi-inclusive** distribution, Γ_t is expressed as

$$\Gamma_t = \frac{1}{2m_t} \int \frac{d^{D-1}k}{(2\pi)^{D-1}2E} \mathcal{W}_{tb}^{\mu\nu} \sum_{\lambda}^{L,R,0} \varepsilon_{\mu}^*(k, \lambda) \varepsilon_{\nu}(k, \lambda)$$

- Hadronic tensor $\mathcal{W}_{tb}^{\mu\nu}$ can be decomposed to 5 linearly-independent Lorentz-tensor structures:

$$\mathcal{W}_{tb}^{\mu\nu}(p, k) = W_1 g^{\mu\nu} + W_2 p^{\mu} p^{\nu} + W_3 k^{\mu} k^{\nu} + W_4 (p^{\mu} k^{\nu} + k^{\mu} p^{\nu}) + W_5 i \epsilon^{\mu\nu\rho\sigma} p_{\rho} k_{\sigma}$$

- $\mathcal{W}_{tb}^{\mu\nu}$
 - directly \Rightarrow The energy distribution of **W boson**
 - $\varepsilon_{\mu}^*(k, \lambda) \varepsilon_{\nu}(k, \lambda) \rightarrow \mathbb{P}_{\lambda}^{\mu\nu}$, and $\int dE \dots$ manually \Rightarrow Total decay width Γ_t and partial width $\Gamma_{\lambda} (\lambda = L, R, 0)$

- Projectors for polarized W boson

$$\begin{cases} \mathbb{P}^{\mu\nu} = -g^{\mu\nu} + \frac{k^{\mu} k^{\nu}}{m_W^2} \\ \mathbb{P}_0^{\mu\nu} = \frac{m_W^2}{m_t^2} \frac{1}{|\vec{k}|^2} (p^{\mu} + \frac{p \cdot k}{m_W^2} k^{\mu})(p^{\nu} - \frac{p \cdot k}{m_W^2} k^{\nu}) \\ \mathbb{P}_F^{\mu\nu} = -\frac{1}{m_t} \frac{1}{|\vec{k}|} i \epsilon^{\mu\nu\rho\sigma} p_{\rho} k_{\sigma} \end{cases} \Rightarrow \begin{cases} \mathbb{P}_L^{\mu\nu} = (\mathbb{P}^{\mu\nu} - \mathbb{P}_0^{\mu\nu} - \mathbb{P}_F^{\mu\nu})/2 \\ \mathbb{P}_R^{\mu\nu} = (\mathbb{P}^{\mu\nu} - \mathbb{P}_0^{\mu\nu} + \mathbb{P}_F^{\mu\nu})/2 \end{cases}$$

K. G. Chetyrkin, A. Retey. (2000)

Computational Details for $\mathcal{W}_{tb}^{\mu\nu}$

➤ Generating integrands:

- Phase-space integrals  Loop integrals

$$\delta(x) \rightarrow \frac{1}{2\pi i} \left(\frac{1}{x - i0} - \frac{1}{x + i0} \right) \quad \text{C. Anastasiou, K. Melnikov. (2002)}$$

- Complexity of Feynman amplitudes:
 - ✓ The number of integral families: 245
 - ✓ The number of Feynman integrals: $\sim 7 \times 10^4$

➤ Integrals reduction:

- Package: **Blade** (Integration-By-Parts relations + **Block-Triangular form**)
- The number of master integrals: 2988
- Computing resources: $\sim 1.1 \times 10^4 \text{CPU} \cdot h$

K. G. Chetyrkin, F. V. Tkachov. (1981)
S. Laporta. (2000)
A. V. Manteuffel, R. M. Schabinger. (2015)

X. Guan, X. Liu, Y.-Q. Ma. (2020)

➤ Evaluating master integrals:

- Package: **AMFlow** (Numerical differential equations + **Auxiliary mass flow method**)
- Expressed as deeply expanded power-log series
- Computing resources: $\sim 1.5 \times 10^4 \text{CPU} \cdot h$

R. N. Lee, V. A. Smirnov. (2018)
R. Bonciani, G. Degrossi, P. P. Giardino, R. Grober. (2019)
H. Frellesvig, M. Hidding, L. Maestri, F. Moriello, G. Salvatori. (2020)
L. Chen, M. Czakon, M. Niggetiedt. (2021)
Martijin Hidding. (2021)
X. Liu, Y.-Q. Ma. (2021)
M. Fael, F. Lange, K. Schonwald, M. Steinhauser. (2022)

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Total Decay Width

- The QCD corrections to Γ_t can be parameterized as

$$\Gamma_t = \Gamma_0 \left(c_0 + \frac{\alpha_s}{\pi} c_1 + \left(\frac{\alpha_s}{\pi} \right)^2 c_2 + \left(\frac{\alpha_s}{\pi} \right)^3 c_3 \right)$$

where $\Gamma_0 = \frac{G_F m_W^2 m_t |V_{tb}|^2}{12\sqrt{2}}$.

- Input values: $G_F = 1.166379 \times 10^{-5} \text{GeV}^{-2}$, $\alpha_s(\mu = m_t/2) = 0.1189$, $m_t = 172.69 \text{GeV}$, $m_W = 80.377 \text{GeV}$, $\mu = m_t/2$
 motivated by **kinematic energy** $m_t - m_W - m_b$

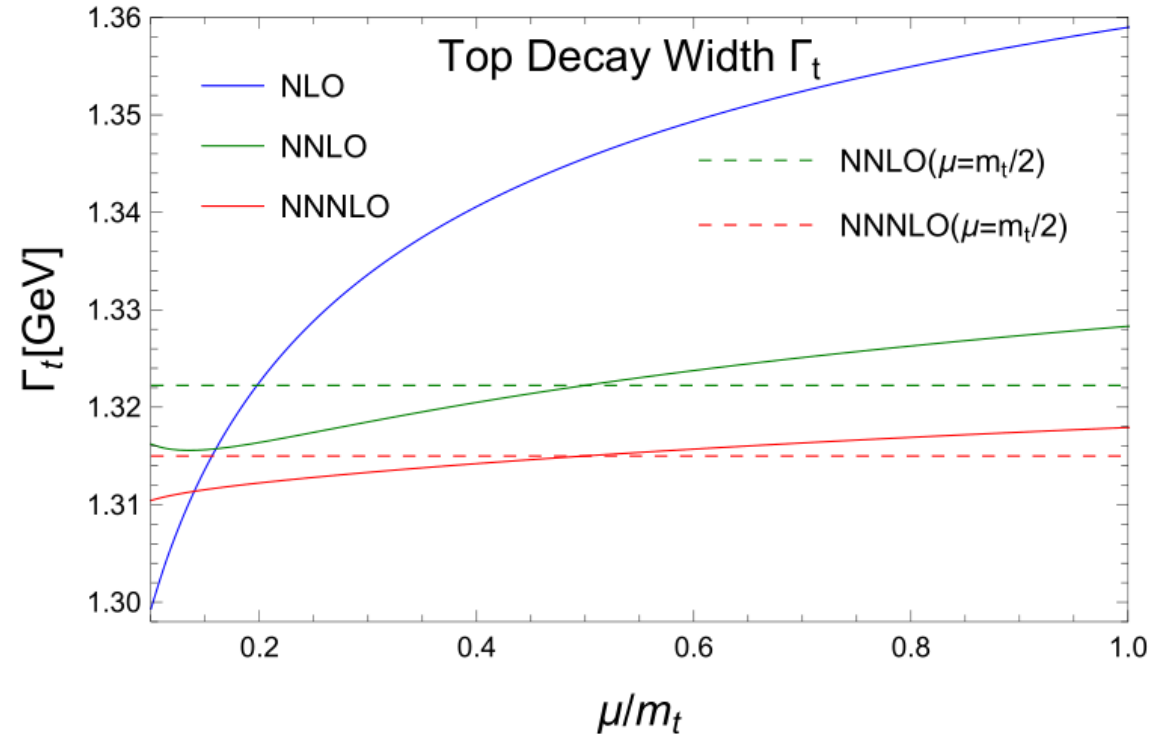
$$\Gamma_t = 1.48642 - 0.140877 - 0.023306 - 0.007240 \text{ GeV}$$

$$= \mathbf{1.31500 \text{ GeV}}$$

- The **pure $\mathcal{O}(\alpha_s^3)$** correction reduces the NNLO results of Γ_t by about 10MeV (**exceeding NNLO scale variation**).
- The **NNLO** green curve can **never** cover the **NNNLO** red curve at any scales less than $\frac{\mu}{m_t} = 0.6$.



Studying the μ dependence at NNLO \Rightarrow Underestimate the theoretical error!!



Other Effects

➤ Similar perturbative parametrization: $\Gamma_t = \Gamma_0(c_0 + \frac{\alpha_s}{\pi}c_1 + (\frac{\alpha_s}{\pi})^2c_2 + (\frac{\alpha_s}{\pi})^3c_3)$

- **Off-shell W** effects (**effective mass of W** is the **invariant-mass of the lepton pair** produced from W -decay)

$$\frac{\tilde{c}_0 - c_0}{c_0} = -1.54\%, \quad \frac{\tilde{c}_1 - c_1}{c_1} = -1.53\%, \quad \frac{\tilde{c}_2 - c_2}{c_2} = -1.39\%, \quad \frac{\tilde{c}_3 - c_3}{c_3} = -1.23\%$$

- **Finite b -quark mass** effects, $m_b = 4.78\text{GeV}$

$$\frac{c_1^{m_b} - c_1}{c_1} \approx \frac{c_2^{m_b} - c_2}{c_2} \approx -1.47\%$$

- **NLO electroweak corrections:** $\delta_{EW} = 1.68\%$

➤ Final results for decay width Γ_t :

$$\Gamma_t = 1.3148_{-0.005}^{+0.003} \times |V_{tb}|^2 + 0.027(m_t - 172.69)\text{GeV}$$

- Conservative estimate of the QCD scale uncertainty ($\frac{\mu}{m_t} \in [0.1,1]$)
- Input t -quark mass value $172.69 \pm 0.30\text{GeV}$

➤ The error of which **meets the request by future colliders.**

W-Helicity Fractions

- Partial width for a **polarized W**:

$$\Gamma_\lambda = \frac{1}{2m_t} \int \frac{d^{D-1}k}{(2\pi)^{D-1}2E} \mathcal{W}_{tb}^{\mu\nu} \varepsilon_\mu^*(k, \lambda) \varepsilon_\nu(k, \lambda)$$

- W-helicity fractions: $f_\lambda^{[n]} = \frac{\sum_{i=0}^n \Gamma_\lambda^{[n]}}{\sum_{i=0}^n \Gamma_t^{[n]}}$

$$f_0^{[3]} = 0.697706 - 0.008401 - 0.001954 - 0.000613 \\ = 0.686737$$

$$f_L^{[3]} = 0.302294 + 0.007254 + 0.001799 + 0.000586 \\ = 0.311933$$

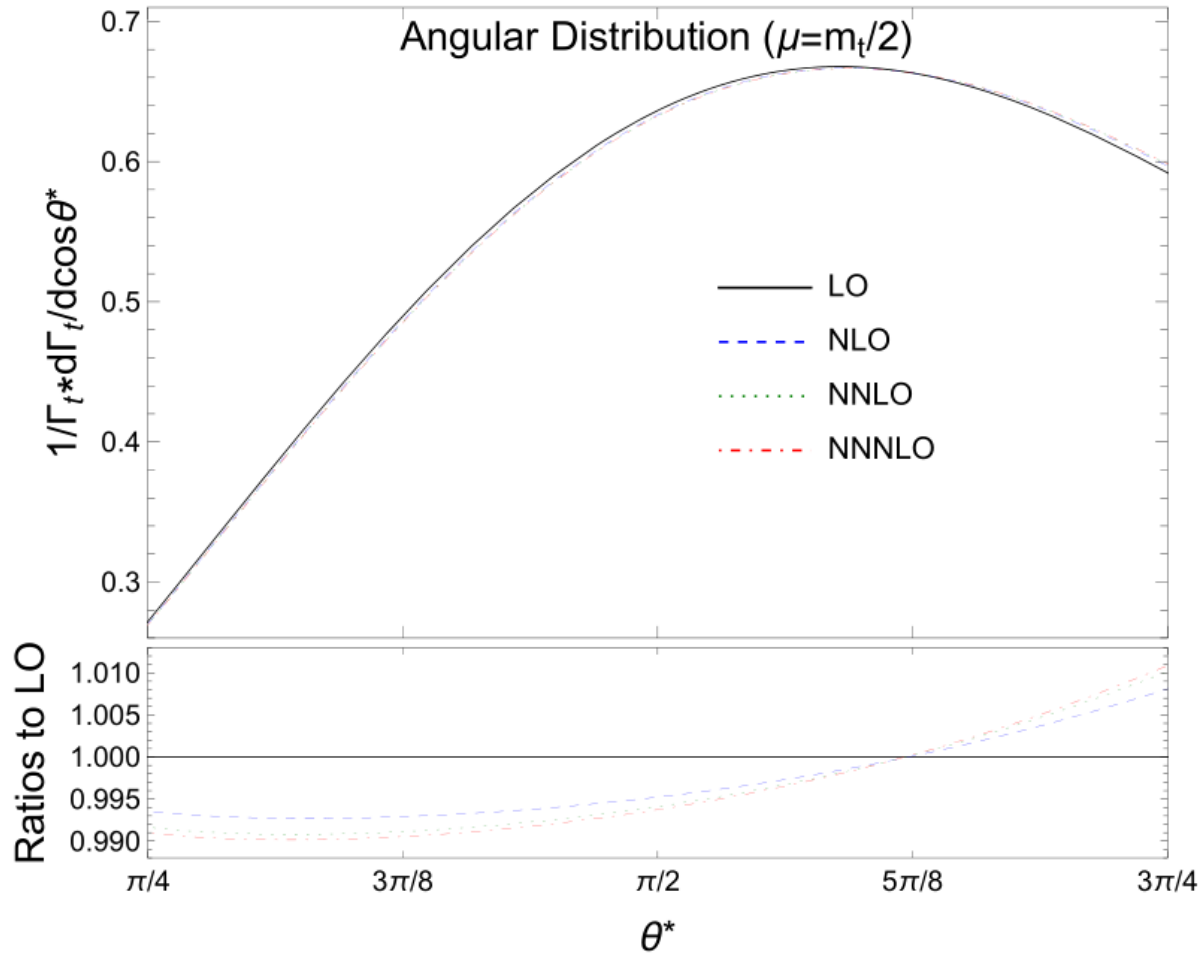
$$f_R^{[3]} = 0. + 0.001147 + 0.000155 + 0.000027 \\ = 0.001330$$

- After incorporating the **NLO electroweak** correction and **m_b effects**, the final results:

$$f_0^{[3]} = 0.686_{-0.003}^{+0.002}, \quad f_L^{[3]} = 0.312_{-0.002}^{+0.001}, \quad f_R^{[3]} = 0.00157_{-0.00002}^{+0.00002}$$

$\cos \theta^*$ Angular Distribution

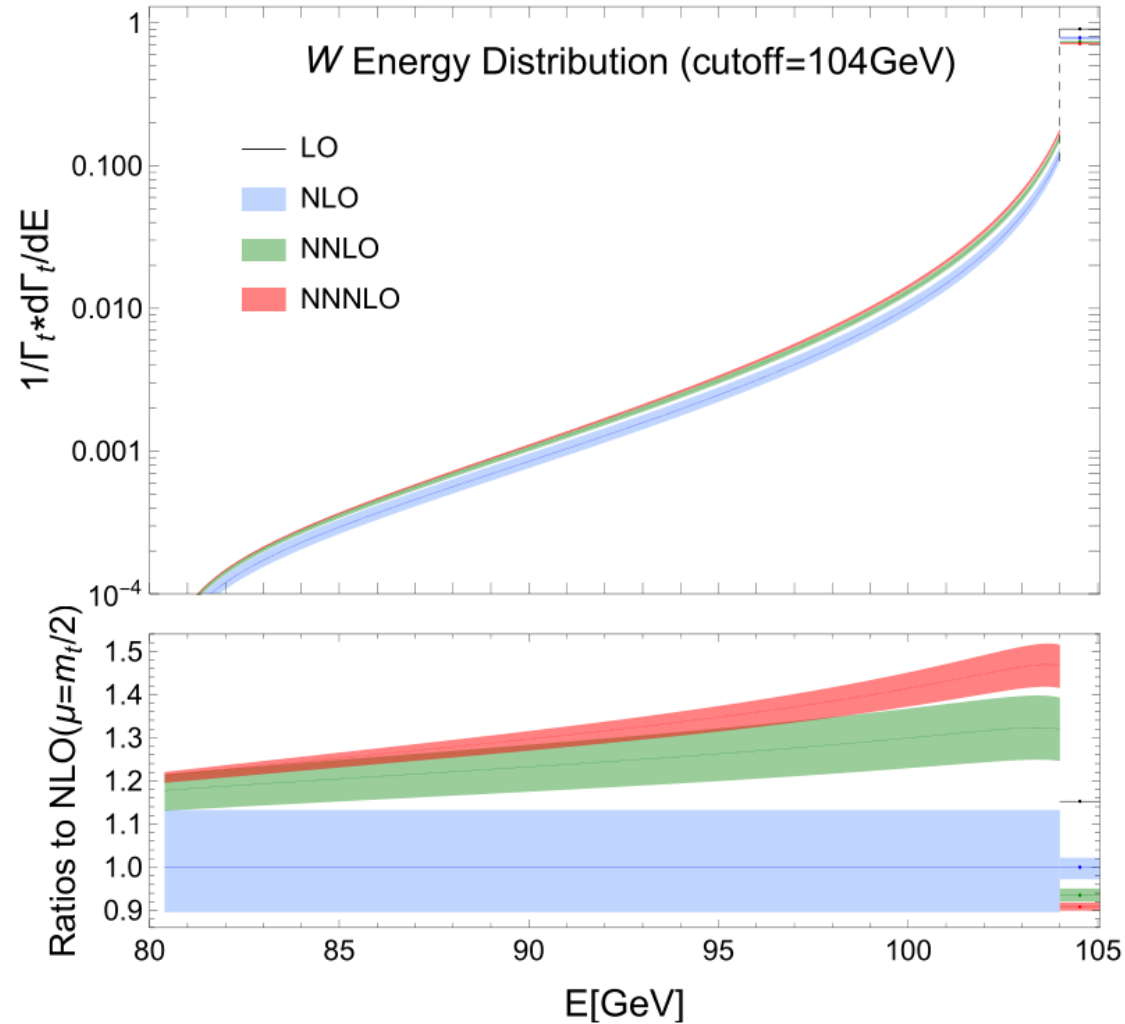
- $\cos \theta^*$ angular distribution (θ^* is the angle between the **momentum direction of the charged lepton** from the W -decay in the W -rest frame and the **W -momentum direction** in the t -quark rest frame)



$$\frac{1}{\Gamma_t} \frac{d\Gamma_t}{d\cos\theta^*} = \frac{3}{4} (\sin\theta^*)^2 f_0 + \frac{3}{8} (1 - \cos\theta^*)^2 f_L + \frac{3}{8} (1 + \cos\theta^*)^2 f_R$$

W-Energy Distribution

➤ W-energy distribution



- In the continuous region, the QCD corrections are **positive and quite sizable**. The **pure $\mathcal{O}(\alpha_s^3)$** correction modifies the lowest order result by about **7~14%** for $E \in [94, 104]$ GeV.
- In the rightmost end, the QCD corrections up to **$\mathcal{O}(\alpha_s^3)$** **decrease** the Born-level result.

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Summary and Outlook

- We have provided the to-date most-precise high-precision theoretical prediction for top-quark

decay width, the error of which meets the request by future colliders.

$$\Gamma_t = 1.3148_{-0.005}^{+0.003} \times |V_{tb}|^2 + 0.027(m_t - 172.69)\text{GeV}$$

- We determined W -helicity fractions, $\cos \theta^*$ angular distribution and W -energy distribution at $\mathcal{O}(\alpha_s^3)$

for the first time.

- To achieve a real sub-percent level, we also need the NNLO mixed QCD-EW corrections.

- The approach can be readily applied to the decay of polarized t -quark at $\mathcal{O}(\alpha_s^3)$.

Thanks for your attention!