On constraning cosmology at the level of the field with a bootstrap approach

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A summary of the problem

The context of LSS cosmology:

- We want to extract the most information, test new models and discover/constrain new physics
- Cosmological inference on LSS is based on correlators of the galaxy density field $\delta_g(\vec{x})$ (P+B)
- Correlators suffer from many degeneracies (bias relation) and become impractical beyond 2-point (covariances, higher order correlators)
- The alternative: use the field itself, it contains all the information and has minimal amounts of degeneracies

REFERENCES:

Babić, Schmidt, Tucci 24 [<u>2407.01524v1]</u> Nguyen+ 24 [<u>2403.03220v3]</u> Stadler, Schmidt, Reinecke 23 [<u>2303.09876v1]</u>

This work

Outline:

- Data is the dark matter density field from N-body simulations + Poissonian shot noise as a proxy for galaxy density
- Perturbative forward model (from initial conditions to late time) within the "LSS bootstrap" approach
- Define a likelihood (gaussian in first approximation), and sample a posterior through MCMC

REFERENCES: Kostić+22 [2212.07875v2] Schmidt 20 [2009.14176v3] Cabass 20 [2007.14988v2] Schmidt+ 20 [2004.06707v2]

Forward model: bootstrap

Non-linearities (kernels) of the field constrained by symmetries and conservation laws (MODEL INDEPENDENT):

- Rotational and translational invariance
- Mass conservation (bias relation when dropped)
- Scale independence (no physical scales)
- Extended Galilean Invariance

Extra degrees of freedom in the kernels (get fixed when choosing EOMs)

REFERENCES: Marinucci, Pardede, Pietroni 24 [2405.08413v1] D'Amico+ 21 [2109.09573v2] Fujita, Vlah 20 [2003.10114v2]

Forward model: implementation

GridSPT (Taruya, Nishimichi, Jeong 21 [2109.06734v1])

$$\begin{pmatrix} \delta_n \\ \theta_n \end{pmatrix} = \frac{2}{(2n+3)(n-1)} \begin{pmatrix} n+\frac{1}{2} & 1 \\ \frac{3}{2} & n \end{pmatrix} \sum_{m=1}^{n-1} \begin{pmatrix} \vec{\nabla} \delta_m \cdot \vec{u}_{n-m} + \delta_m \theta_{n-m} \\ \frac{1}{2} \Delta(\vec{u}_m \cdot \vec{u}_{n-m}) \end{pmatrix};$$

N-body

GridSPT

difference



3rd order, 1 Gpc box, Ng=128, z=1, kmax=0.12

Real space likelihood and results

GridSPT + 2nd order bootstrap + 3rd order UV counterterms:

$$\begin{pmatrix} \delta_2 \\ \theta_2 \end{pmatrix} = \begin{pmatrix} \frac{a_{\gamma}^{(2)}}{2} & 1 - \frac{a_{\gamma}^{(2)}}{2} \\ \frac{d_{\gamma}^{(2)}}{2} & 1 - \frac{d_{\gamma}^{(2)}}{2} \end{pmatrix} \begin{pmatrix} \vec{\nabla} \delta_1 \cdot \vec{u}_1 + \delta_1 \theta_1 \\ \frac{1}{2} \Delta u_1^2 \end{pmatrix}$$
$$\begin{pmatrix} \delta_3 \\ \theta_3 \end{pmatrix} = \text{GridSPT}|_{n=3} + c_2 \begin{pmatrix} 1 \\ 0 \end{pmatrix} \Delta \delta_1$$

Likelihood:

$$\mathcal{L}(\delta \,|\, \vec{\alpha}) \propto \exp\left[-\frac{\bar{n}}{2L^3} \sum_{k < k_{\max}} \left(\delta_{\vec{k}} - \delta_{\vec{k}}^{\mathrm{PT}}(\vec{\alpha})\right)^2\right]$$



Redshift space results

Redshift space GridSPT (bootstrap+UV):

$$\begin{split} \delta_1^{(S)} &= D_1 \\ \delta_2^{(S)} &= D_2 + f \tilde{\nabla}_z \Big(D_1 u_{z,1} \Big), \\ \delta_3^{(S)} &= D_3 + f \tilde{\nabla}_z \Big(D_1 u_{z,2} + D_2 u_{z,1} \Big) + \frac{f^2}{2!} \tilde{\nabla}_z^2 \Big(D_1 u_{z,1}^2 \Big) \end{split}$$

+ 3rd order angle dependent counterterms: $c_{z,2}f(1 + \varepsilon_f)\mu^2\Delta\delta_1 \qquad c_{z,4}f^4(1 + \varepsilon_f)^4\mu^4\left(1 + f(1 + \varepsilon_f)\mu^2\right)\Delta^2\delta_1$ $(f \equiv f_{\Lambda \text{CDM}})$



Future developments

- Understand the running and find ways to correct for it
- Also see what this implies for the cosmological parameters
- Perform MCMC with much finer grids (possible in the case of bootstrap)
- Hamiltonian MC (Julia GridSPT) to marginalize over ICs
- Compare with P + B

About running on bootstrap and counterterm (real space, WIP):

The running can actually be computed explicitly for bootstrap parameters and counterterms: they act similarly as bias parameters in the perturbative expansions (factors x fields).

Take the derivative of the likelihood wrt the Λ -dependent parameter to 0 (to trace the maximum), $a_{\gamma}^{(2)}$ and c_2 will be co-dependent and so will run together as Λ changes

#	ng	ag2	c_2
	64	1.2212762932584806	20.28538993651349
	128	1.3571428571428283	0.024750000000000022
	256	1.3892006054156556	2.5069942207809537
	512	1.4281154551674966	-0.07416138676181436

About running on bootstrap and counterterm (real space, WIP):

If the instead of using the N-body as data in the likelihood, we perform the same calculations with PT with a large Λ , we find that $a_{\gamma}^{(2)}$ doesn't run, but c_2 does.

#	ng	ag2	c_2
	64	1.421875	0.02438672
	128	1.4296875	0.02424805
	256	1.42849312	0.17251358
	512	1.42857143	0.