



PAIR PRODUCTION OF HIGGS BOSONS AT NLO

Michael Spira (PSI)

I Introduction

II $gg \rightarrow HH$

III Conclusions

in collaboration with J. Baglio, A. Bhattacharya, F. Campanario, S. Carlotti, J. Chang, S. Glaus, J. Mazzitelli, J. Ronca, M. Mühlleitner and J. Schlenk

Extended Scalar Sectors From All Angles

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Angels

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Angels*

No name tags?

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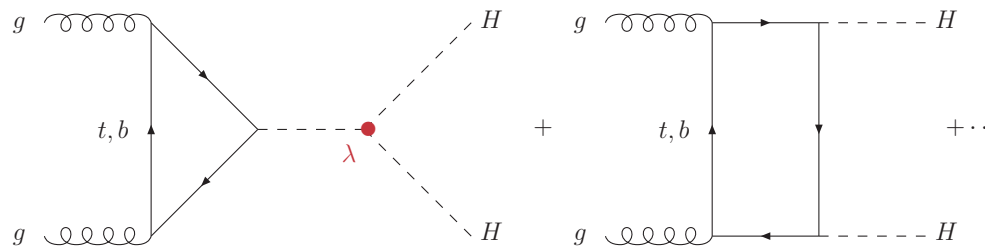
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Extended Scalar Sectors From All Angles
Angels

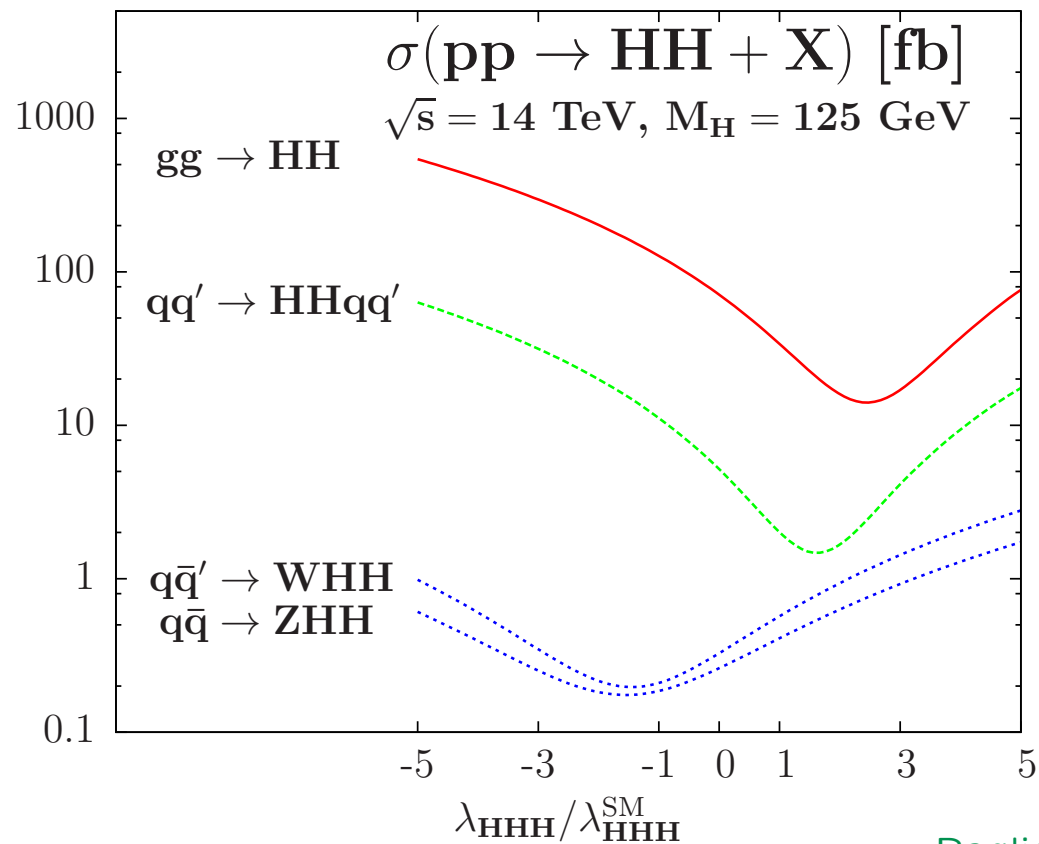
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Extended Calculations of
Extended Visions of
Extended Dreams of
Extended Extensions of

I INTRODUCTION



- third generation dominant: t (b)



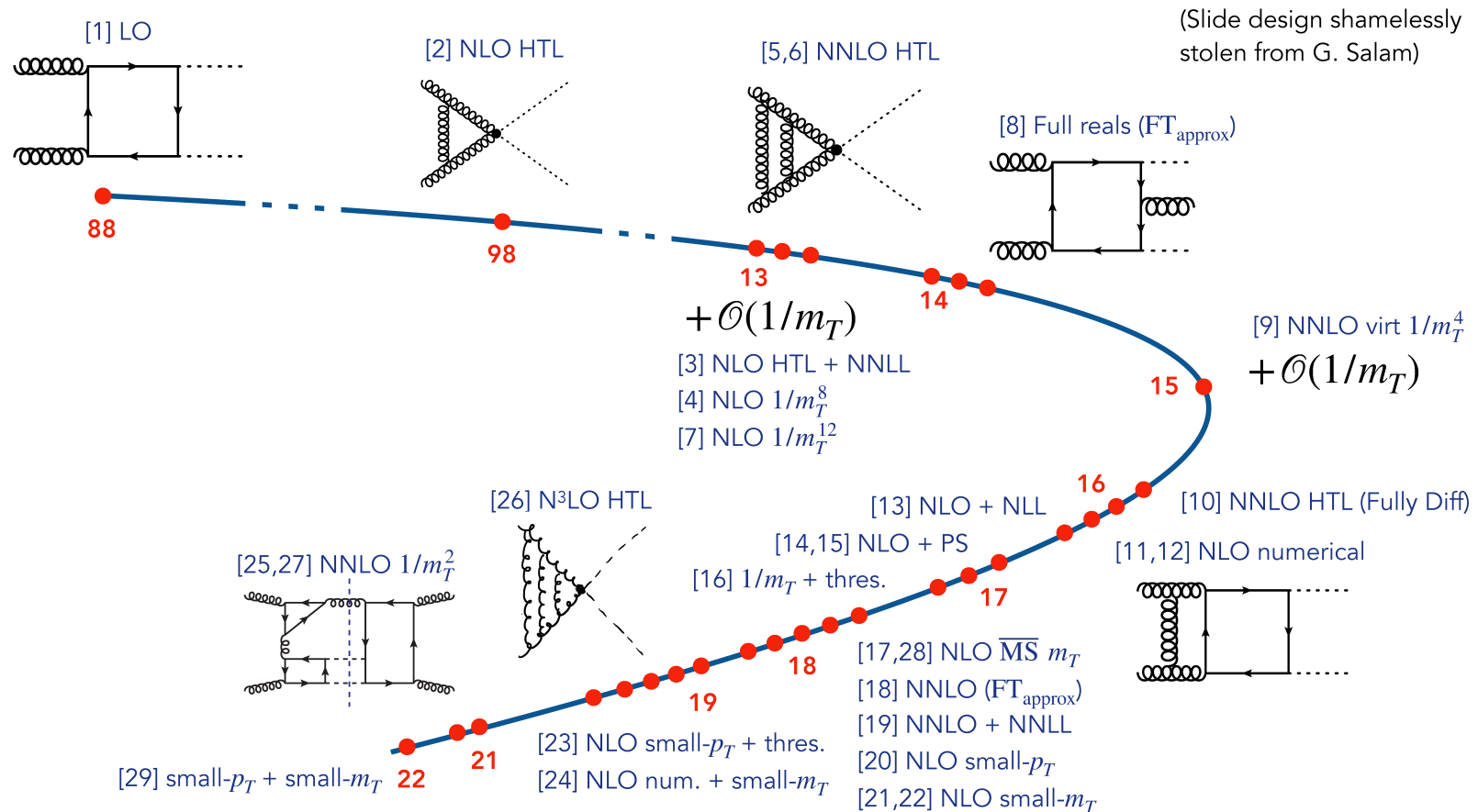
$gg \rightarrow HH :$

$$\frac{\Delta\sigma}{\sigma} \sim -\frac{\Delta\lambda}{\lambda}$$

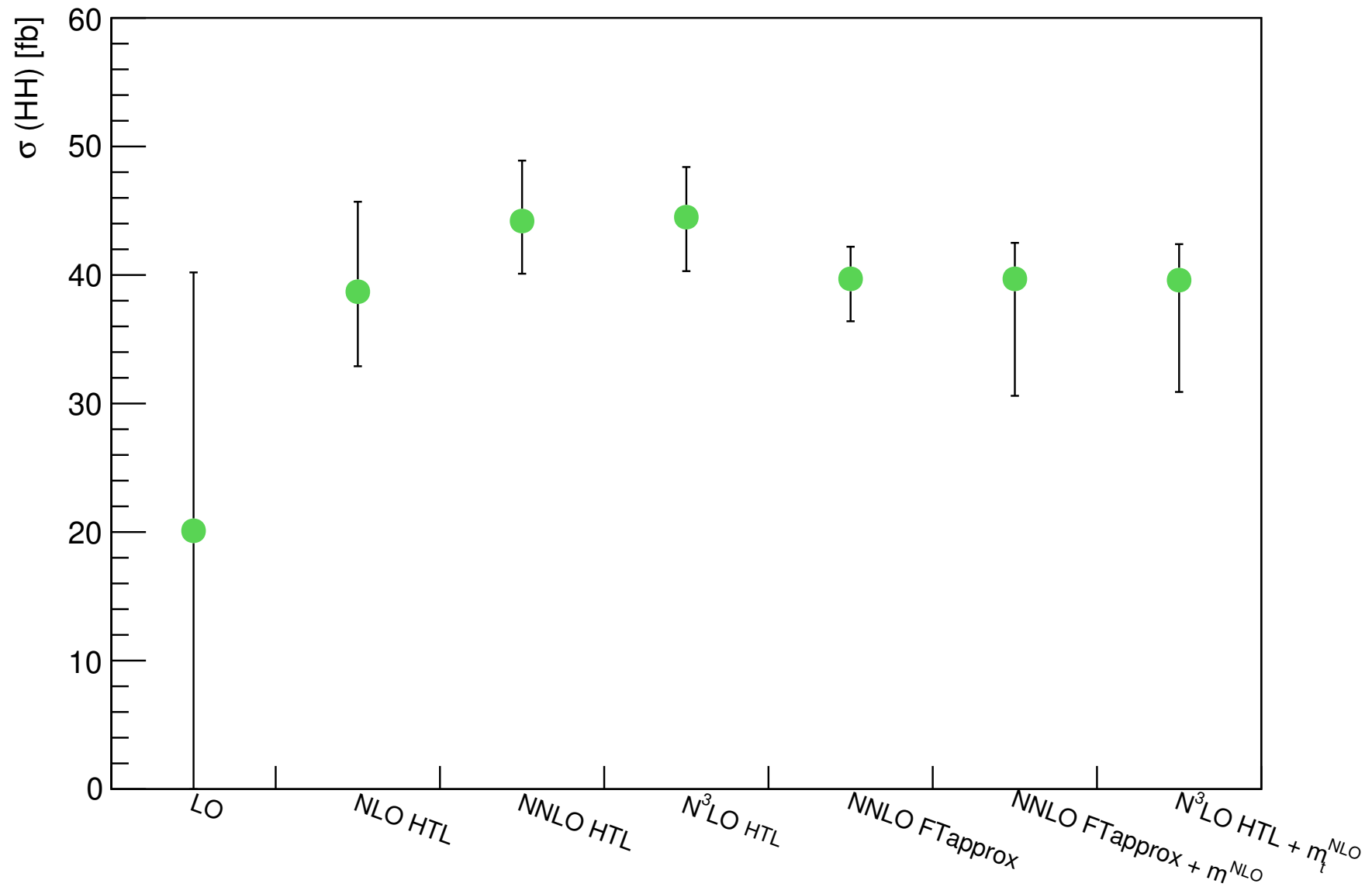
II $gg \rightarrow HH$

S. Jones

An approximate history (30 years in 30 seconds)



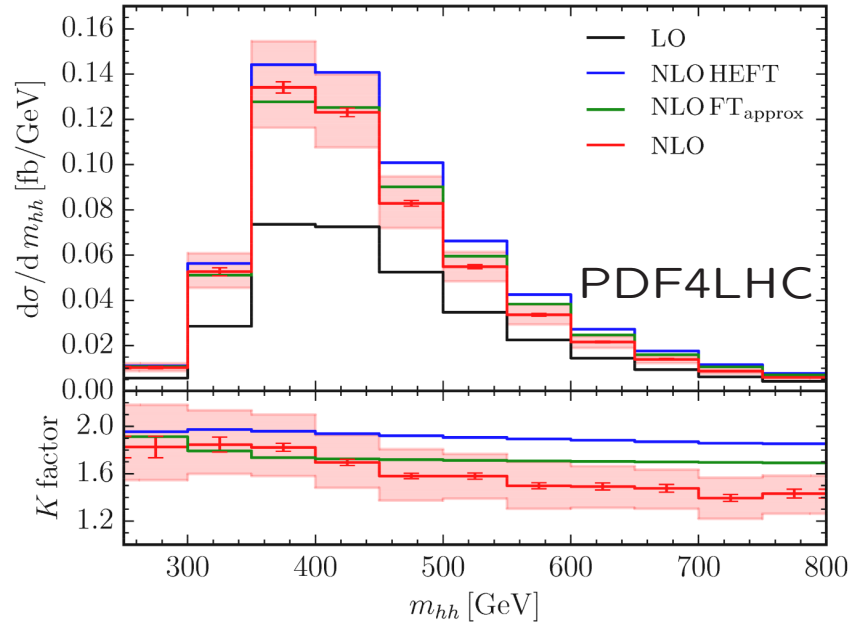
[1] Glover, van der Bij 88; [2] Dawson, Dittmaier, Spira 98; [3] Shao, Li, Li, Wang 13; [4] Grigo, Hoff, Melnikov, Steinhauser 13; [5] de Florian, Mazzitelli 13; [6] Grigo, Melnikov, Steinhauser 14; [7] Grigo, Hoff 14; [8] Maltoni, Vryonidou, Zaro 14; [9] Grigo, Hoff, Steinhauser 15; [10] de Florian, Grazzini, Hanga, Kallweit, Lindert, Maierhöfer, Mazzitelli, Rathlev 16; [11] Borowka, Greiner, Heinrich, SPJ, Kerner, Schlenk, Schubert, Zirke 16; [12] Borowka, Greiner, Heinrich, SPJ, Kerner, Schlenk, Zirke 16; [13] Ferrera, Pires 16; [14] Heinrich, SPJ, Kerner, Luisoni, Vryonidou 17; [15] SPJ, Kuttimalai 17; [16] Gröber, Maier, Rauh 17; [17] Baglio, Campanario, Glaus, Mühlleitner, Spira, Streicher 18; [18] Grazzini, Heinrich, SPJ, Kallweit, Kerner, Lindert, Mazzitelli 18; [19] de Florian, Mazzitelli 18; [20] Bonciani, Degrossi, Giardino, Gröber 18; [21] Davies, Mishima, Steinhauser, Wellmann 18, 18; [22] Mishima 18; [23] Gröber, Maier, Rauh 19; [24] Davies, Heinrich, SPJ, Kerner, Mishima, Steinhauser, David Wellmann 19; [25] Davies, Steinhauser 19; [26] Chen, Li, Shao, Wang 19, 19; [27] Davies, Herren, Mishima, Steinhauser 19, 21; [28] Baglio, Campanario, Glaus, Mühlleitner, Ronca, Spira 21; [29] Bellafronte, Degrossi, Giardino, Gröber, Vitti 22;



Full NLO calculation: top only, numerical integration

Borowka <i>et al.</i>	Baglio <i>et al.</i>
tensor reduction	no tensor reduction
sector decomposition	IR, end-point subtraction
contour deformation	IBP, Richardson extrapolation
$m_t = 173 \text{ GeV}$	$m_t = 172.5 \text{ GeV}$

Borowka, Greiner, Heinrich, Jones, Kerner, Schlenk, Schubert, Zirke
Baglio, Campanario, Glaus, Mühlleitner, Ronca, S., Streicher



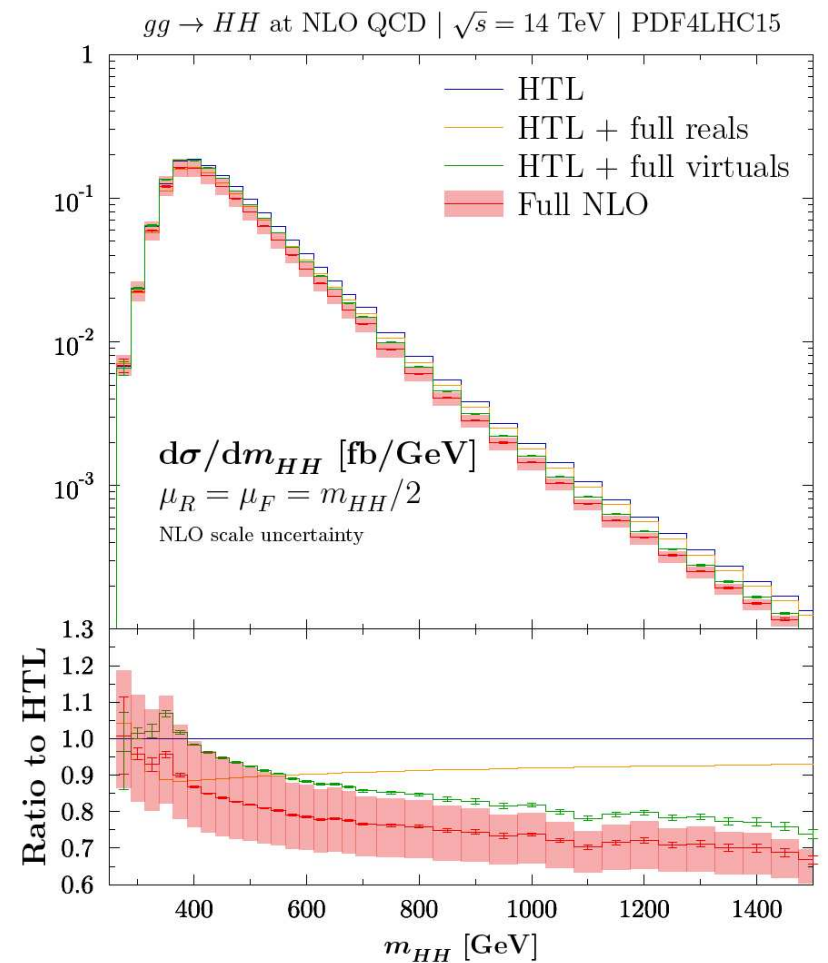
Borowka, Greiner, Heinrich, Jones, Kerner
Schlenk, Schubert, Zirke

$$\sigma_{NLO} = 32.91(10)_{-12.8\%}^{+13.8\%} \text{ fb}$$

$$\sigma_{NLO}^{HTL} = 38.75_{-15\%}^{+18\%} \text{ fb}$$

$$m_t = 173 \text{ GeV}$$

⇒ -15% mass effects on top of LO

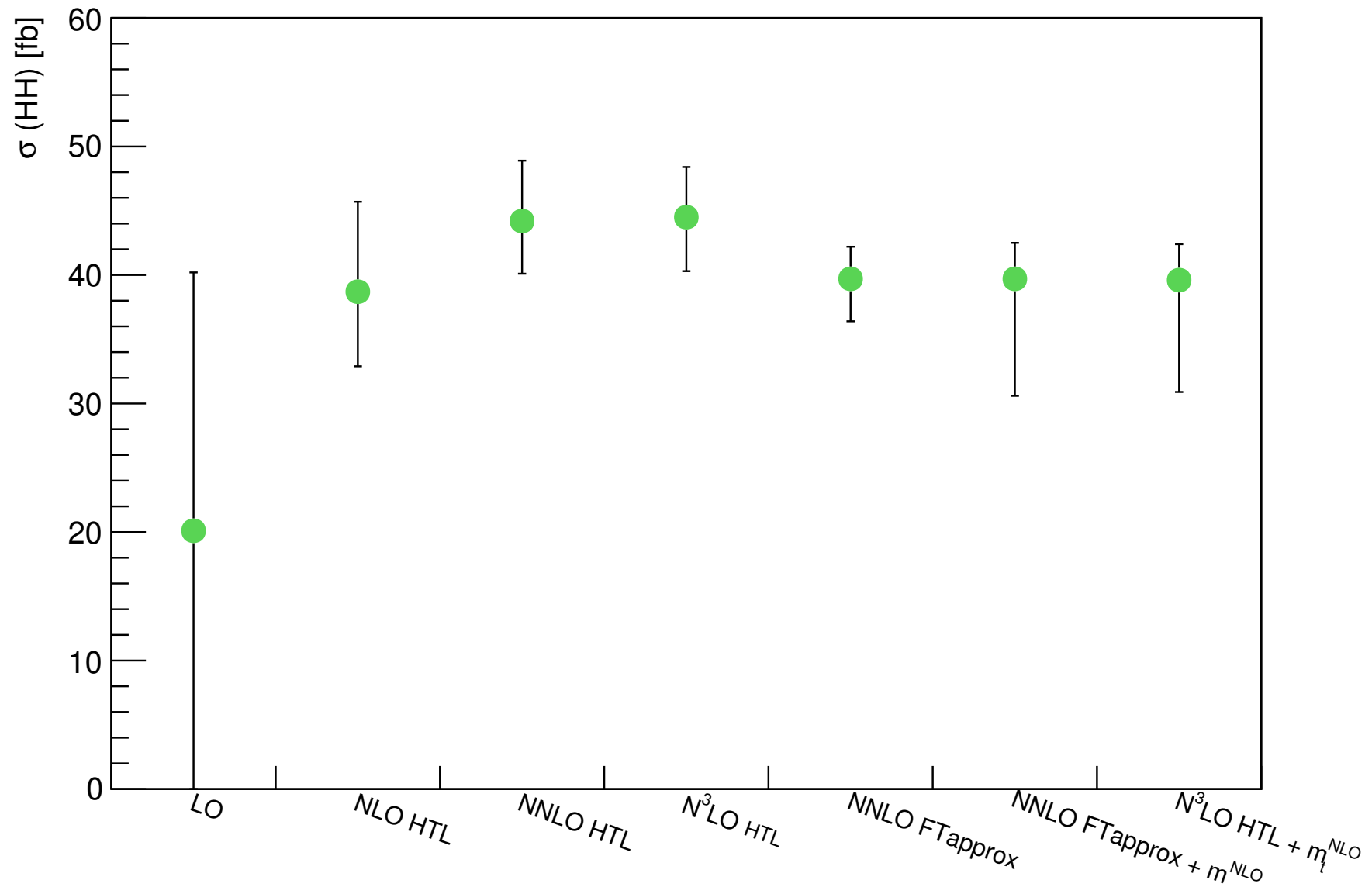


Baglio, Campanario, Glaus,
Mühlleitner, Ronca, S., Streicher

$$\sigma_{NLO} = 32.81(7)_{-12.5\%}^{+13.5\%} \text{ fb}$$

$$\sigma_{NLO}^{HTL} = 38.66_{-15\%}^{+18\%} \text{ fb}$$

$$m_t = 172.5 \text{ GeV}$$



uncertainties due to m_t

- use m_t , $\bar{m}_t(\bar{m}_t)$ and scan $Q/4 < \mu < Q \rightarrow$ uncertainty = envelope:

$$\frac{d\sigma(gg \rightarrow HH)}{dQ} \Big|_{Q=300 \text{ GeV}} = 0.02978(7)_{-34\%}^{+6\%} \text{ fb/GeV},$$

$$\frac{d\sigma(gg \rightarrow HH)}{dQ} \Big|_{Q=400 \text{ GeV}} = 0.1609(4)_{-13\%}^{+0\%} \text{ fb/GeV},$$

$$\frac{d\sigma(gg \rightarrow HH)}{dQ} \Big|_{Q=600 \text{ GeV}} = 0.03204(9)_{-30\%}^{+0\%} \text{ fb/GeV},$$

$$\frac{d\sigma(gg \rightarrow HH)}{dQ} \Big|_{Q=1200 \text{ GeV}} = 0.000435(4)_{-35\%}^{+0\%} \text{ fb/GeV}$$

- bin-by-bin interpolation:

$$\sigma(gg \rightarrow HH) = 32.81_{-18\%}^{+4\%} \text{ fb}$$

final combined ren./fac. scale and m_t scale/scheme unc. @ NNLO_{FTapprox}:

$\sqrt{s} = 13$ TeV :	$\sigma_{tot} = 31.05^{+6\%}_{-23\%}$ fb
$\sqrt{s} = 14$ TeV :	$\sigma_{tot} = 36.69^{+6\%}_{-23\%}$ fb
$\sqrt{s} = 27$ TeV :	$\sigma_{tot} = 139.9^{+5\%}_{-22\%}$ fb
$\sqrt{s} = 100$ TeV :	$\sigma_{tot} = 1224^{+4\%}_{-21\%}$ fb

- expansion methods → fully differential Monte Carlo

Boncianni, Degrassi, Giardino, Gröber
Bagnaschi, Degrassi, Gröber

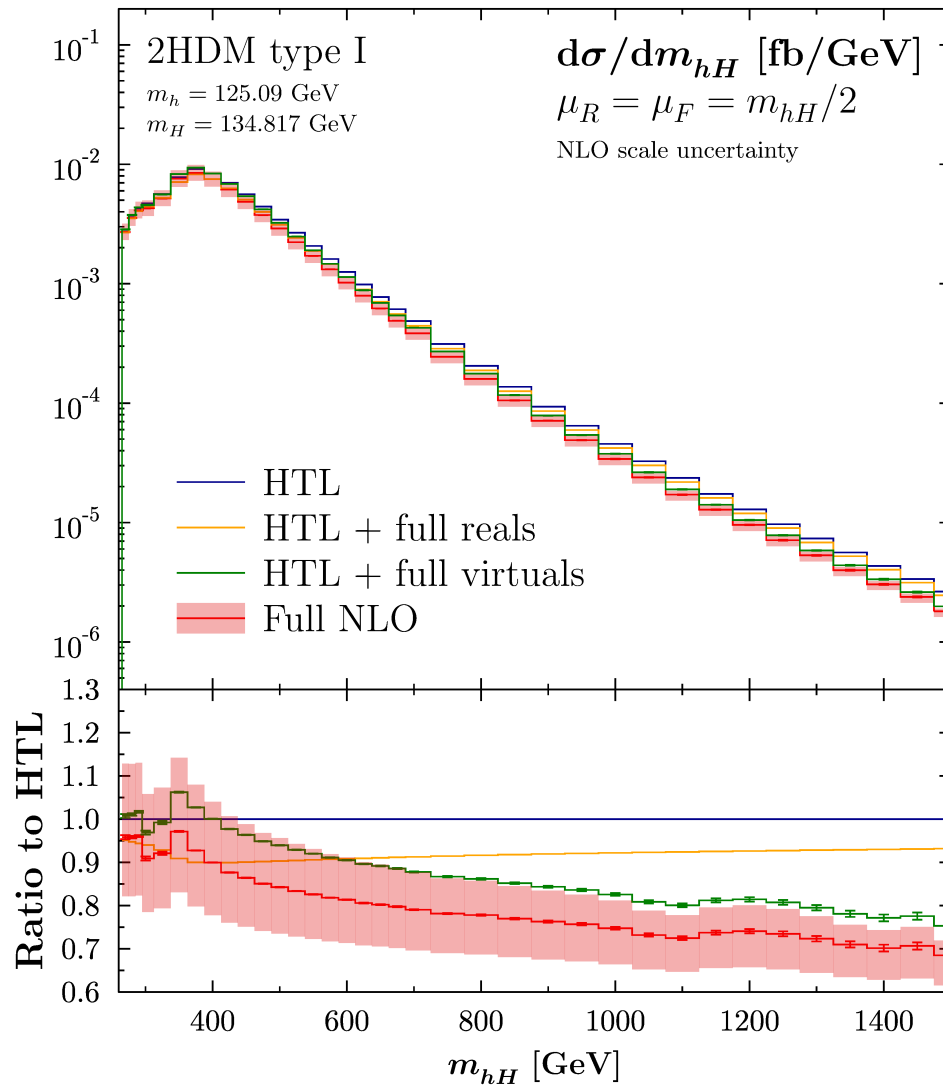
- 2HDM: $gg \rightarrow hh, hH, HH, AA$ available, too

2HDM [type I]: $gg \rightarrow hh, hH, HH, AA$ [no $hA, HA \rightarrow$ DY-like]

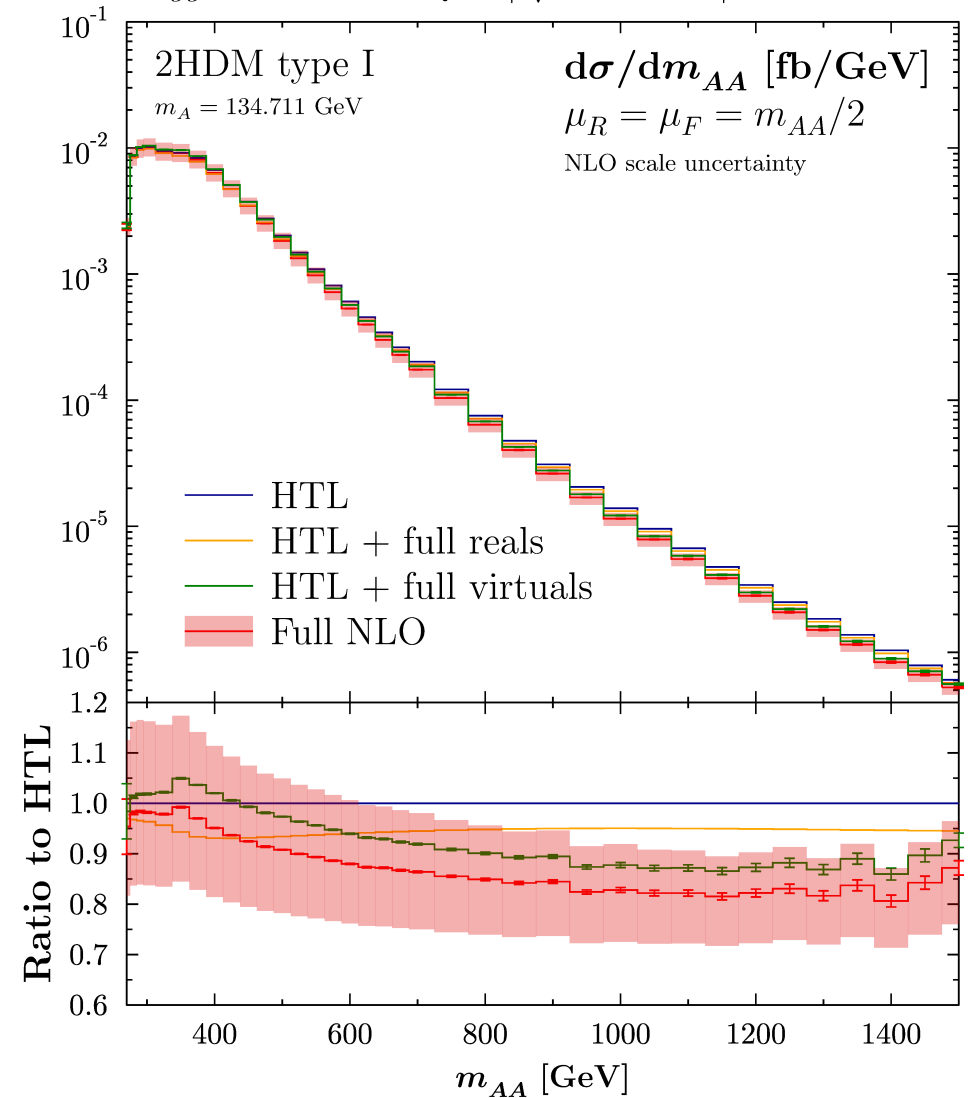
$M_h = 125.09$ GeV $M_H = 134.817$ GeV $M_A = 134.711$ GeV

$\tan\beta = 3.759$ $\alpha = -0.102$ $m_{12}^2 = 4305$ GeV² $\Rightarrow \cos(\beta - \alpha) = 0.157$

$gg \rightarrow hH$ at NLO QCD | $\sqrt{s} = 13$ TeV | PDF4LHC15



$gg \rightarrow AA$ at NLO QCD | $\sqrt{s} = 13$ TeV | PDF4LHC15



combined uncertainties

$$\begin{aligned} 13 \text{ TeV} : \quad & \sigma(gg \rightarrow hH) = 1.592(1)_{-24\%}^{+21\%} \text{ fb} \\ 14 \text{ TeV} : \quad & \sigma(gg \rightarrow hH) = 1.876(1)_{-24\%}^{+21\%} \text{ fb} \\ 27 \text{ TeV} : \quad & \sigma(gg \rightarrow hH) = 7.036(4)_{-23\%}^{+18\%} \text{ fb} \\ 100 \text{ TeV} : \quad & \sigma(gg \rightarrow hH) = 60.49(4)_{-25\%}^{+16\%} \text{ fb} \end{aligned}$$

$$\begin{aligned} 13 \text{ TeV} : \quad & \sigma(gg \rightarrow AA) = 1.643(1)_{-21\%}^{+26\%} \text{ fb} \\ 14 \text{ TeV} : \quad & \sigma(gg \rightarrow AA) = 1.927(1)_{-22\%}^{+26\%} \text{ fb} \\ 27 \text{ TeV} : \quad & \sigma(gg \rightarrow AA) = 7.012(4)_{-21\%}^{+23\%} \text{ fb} \\ 100 \text{ TeV} : \quad & \sigma(gg \rightarrow AA) = 58.12(3)_{-22\%}^{+22\%} \text{ fb} \end{aligned}$$

Is this everything?



Is this everything?

No...

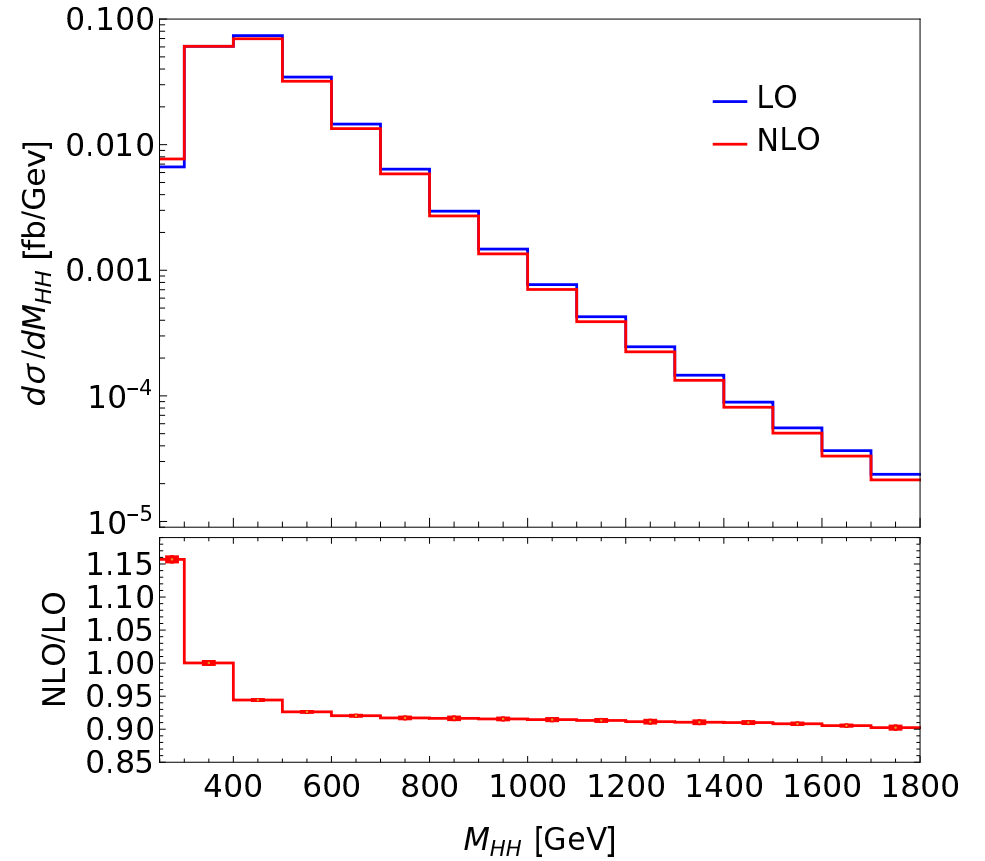
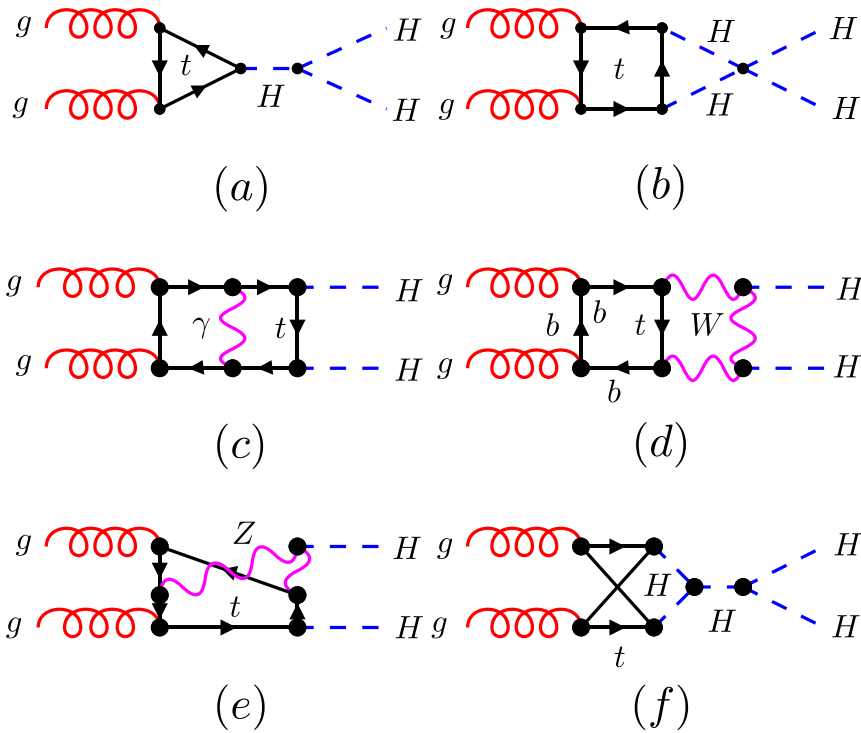
Is this everything?

No...

electroweak corrections...

- (i) y_t : HTL for $ggH(H)$ coupling + full corrections to HHH vertex
Mühlleitner, Schlenk, S.
- (ii) y_t : analytical results for $ggHH$ coupling in the HEL
Davies, Mishima, Schönwald, Steinhauser, Zhang
and close to the production threshold
Davies, Schönwald, Steinhauser, Zhang
- (iii) λ : elw. corrections due to the Higgs self-interactions
Borowka, Duhr, Maltoni, Pagani, Shivaji, Zhao
- (iv) g_t, λ : elw. corrections due to the top Yukawa and Higgs self-interactions [only Higgs exchange diagrams]
Heinrich, Jones, Kerner, Stone, Vestner
- (v) full elw. corrections (\leftarrow to be checked)
Bi, Huang, Huang, Ma, Yu

Full electroweak corrections

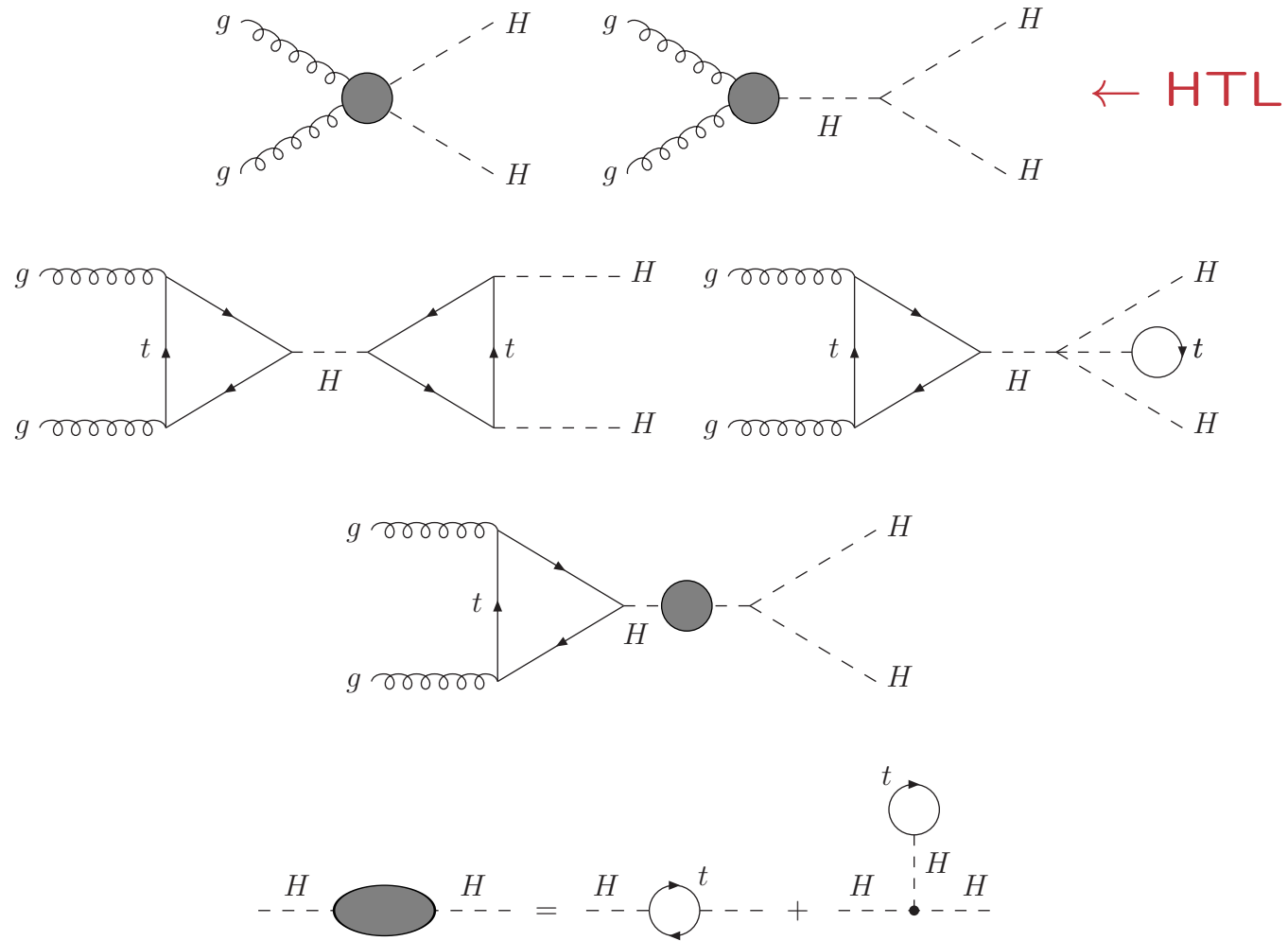


- -4.2% for total cross section

Bi, Huang, Huang, Ma, Yu

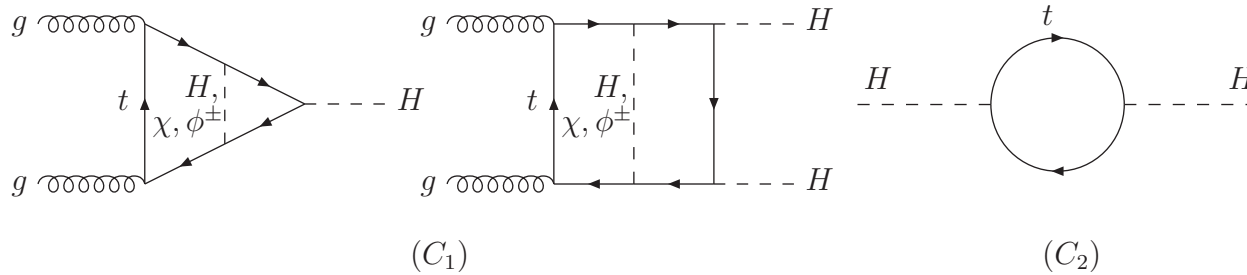
Top-Yukawa-induced elw. corrections

Mühlleitner, Schlenk, S.



(i) effective $ggH(H)$ couplings:

$$\mathcal{L}_{eff} = C_1 \frac{\alpha_s}{12\pi} G^{a\mu\nu} G_{\mu\nu}^a \log \left(1 + C_2 \frac{H}{v} \right)$$



- $C_1 = 1 - 3x_t$: genuine vertex corrections $[x_t = G_F m_t^2 / (8\sqrt{2}\pi^2)]$

Djoaudi, Gambino
Chetyrkin, Kniehl, Steinhauser

- $C_2 = 1 + 7x_t/2 [= 1 + \delta Z_H/2 - \delta v/v]$: universal corrections

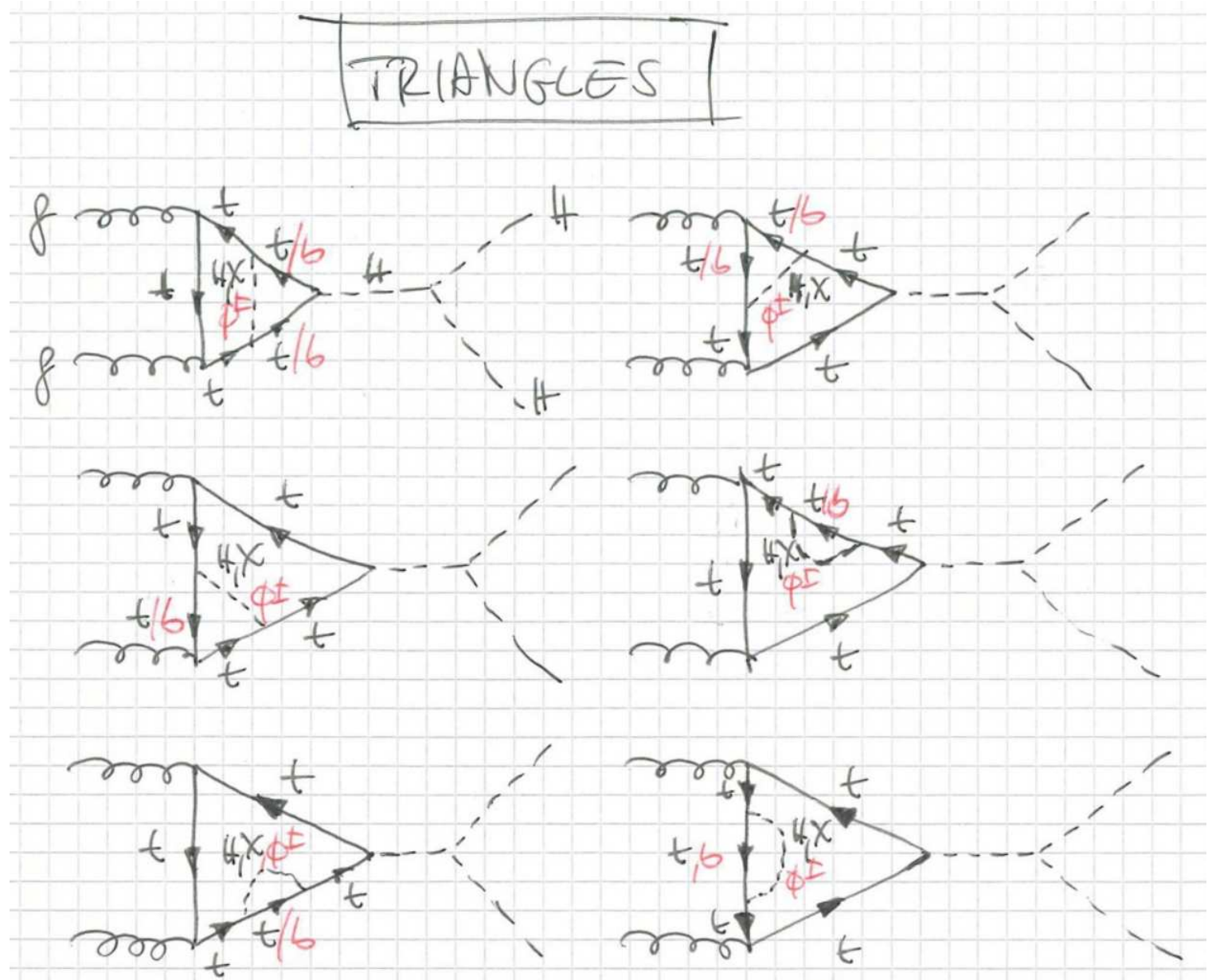
Kniehl, Spira
Kwiatkowski, Steinhauser

$$\mathcal{L}_{eff} = \frac{\alpha_s}{12\pi} G^{a\mu\nu} G_{\mu\nu}^a \left\{ (1 + \delta_1) \frac{H}{v} + (1 + \eta_1) \frac{H^2}{2v^2} + \mathcal{O}(H^3) \right\}$$

$$\delta_1 = \frac{x_t}{2} + \mathcal{O}(x_t^2) \qquad \eta_1 = 4x_t + \mathcal{O}(x_t^2)$$

INTERMEZZO

Full top-mass dependence (wave-function ren. adjusted appropriately)



elw. gaugeless limit + QCD = top-Yukawa model + QCD

$$\phi = \begin{pmatrix} G^+ \\ \frac{v + H + iG^0}{\sqrt{2}} \end{pmatrix}$$

$$\mathcal{L} = -\frac{1}{4}G^{a\mu\nu}G_{\mu\nu}^a + \bar{t}i\not{D}t + |\partial_\mu\phi|^2 - V(\phi) - g_t\bar{Q}_L\phi^c t_R$$

$$V(\phi) = -\mu^2|\phi|^2 + \frac{\lambda}{2}|\phi|^4$$

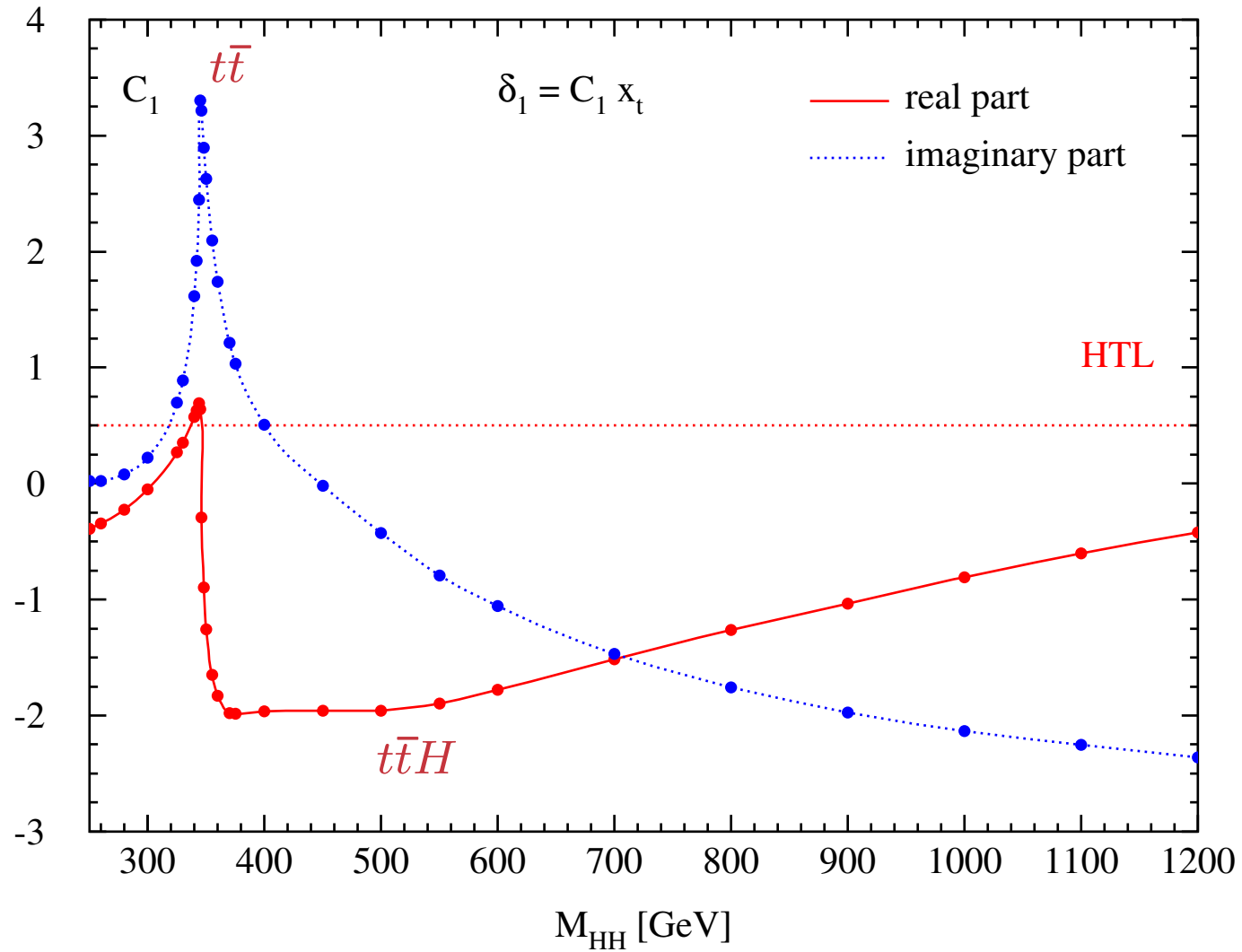
$$= -\frac{M_H^2}{8}v^2 + \frac{M_H^2}{2}H^2 + \frac{M_H^2}{v} \left[\frac{H^3}{2} + \frac{H}{2}(G^0)^2 + HG^+G^- \right]$$

$$+ \frac{M_H^2}{2v^2} \left[\frac{H^4}{4} + \frac{H^2}{2}(G^0)^2 + H^2G^+G^- + (G^+G^-)^2 + (G^0)^2G^+G^- + \frac{(G^0)^4}{4} \right]$$

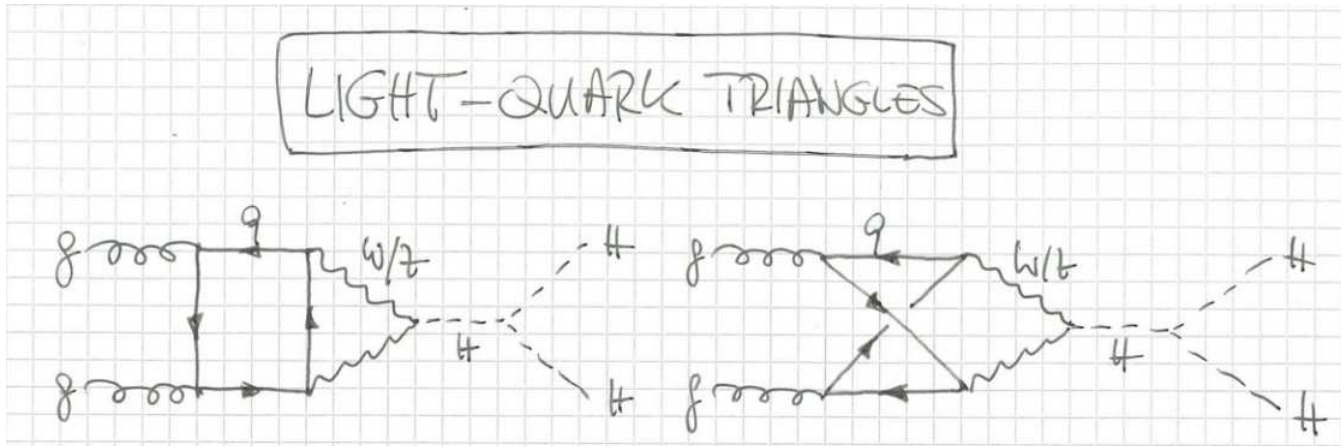
$$\Rightarrow M_{G^0} = M_{G^\pm} = 0$$

$$\delta_1 = \delta_{2loop} + \delta Z_H/2 - \delta v/v$$

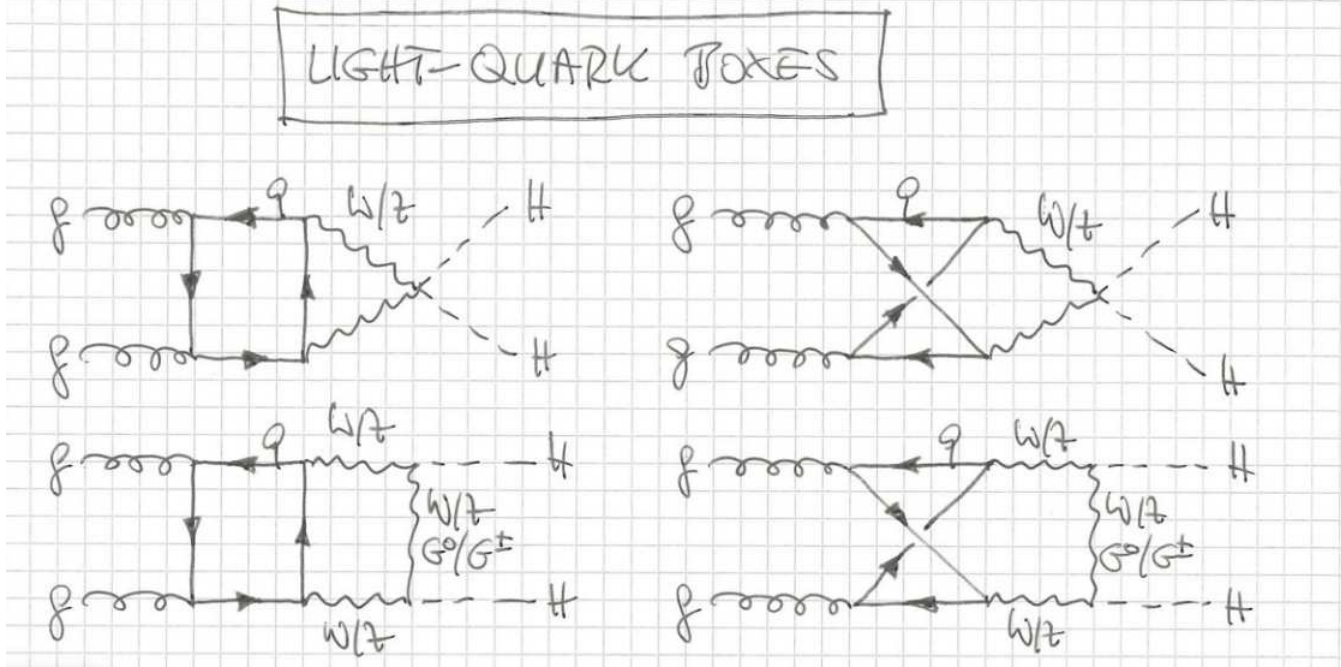
PRELIMINARY



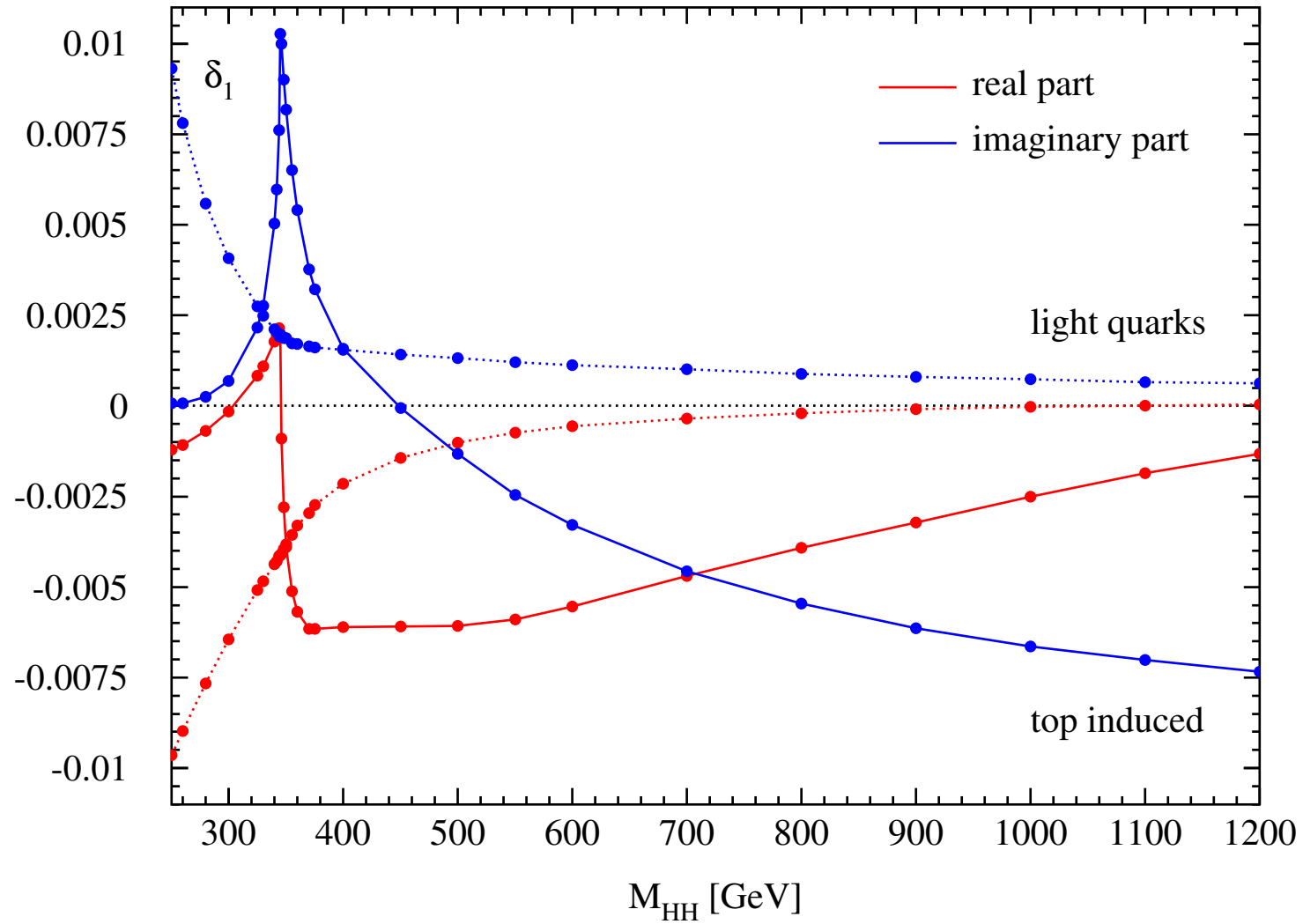
Light-quark loops



Aglietti, Bonciani, Degrassi, Vicini



PRELIMINARY



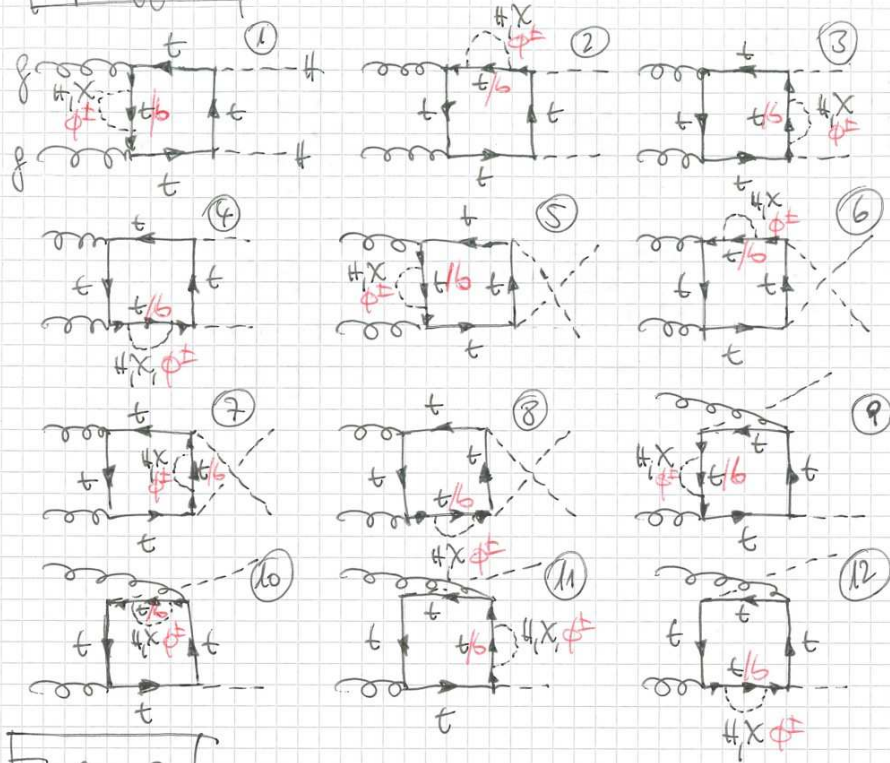
Bhattacharya, Campanario, Carlotti, Chang, Mazzitelli, Mühlleitner, Ronca, S.

- box diagrams in the making...

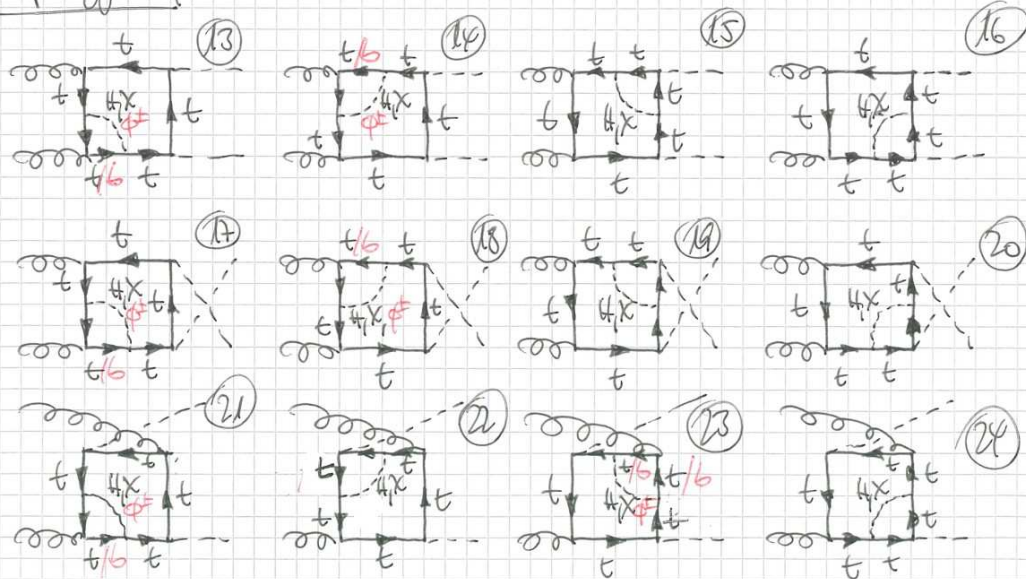
Topology 1

BOXES

1

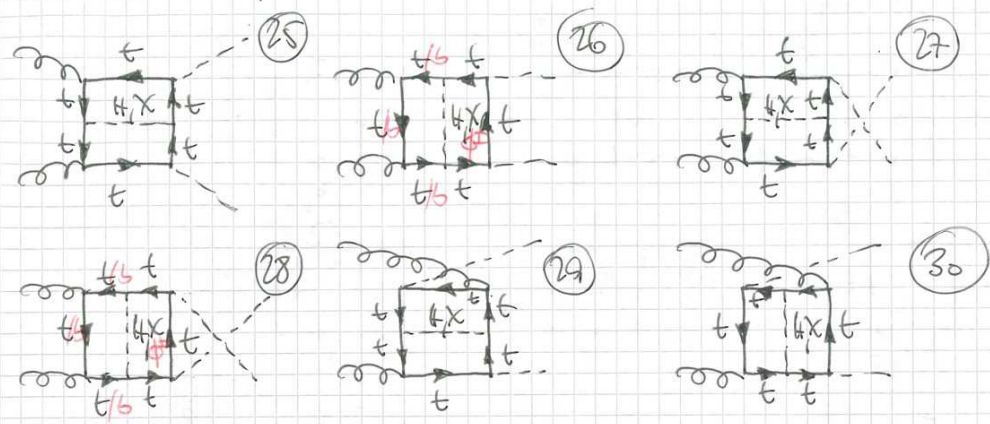


Topology 2



TOPOLOGY 3

2



(ii) effective $HHH(H)$ couplings:

- effective Higgs potential:

Coleman, Weinberg

$$V_{eff} = V_0 + V_1$$

$$V_0 = \mu_0^2 |\phi|^2 + \frac{\lambda_0}{2} |\phi|^4$$

$$V_1 = \frac{3\bar{m}_t^4}{16\pi^2} \Gamma(1 + \epsilon) (4\pi^2)^\epsilon \left(\frac{1}{\epsilon} + \log \frac{\bar{\mu}^2}{\bar{m}_t^2} + \frac{3}{2} \right)$$

$$\phi = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v + H \end{pmatrix} \quad \bar{m}_t = m_t \left(1 + \frac{H}{v} \right)$$

- after renormalization

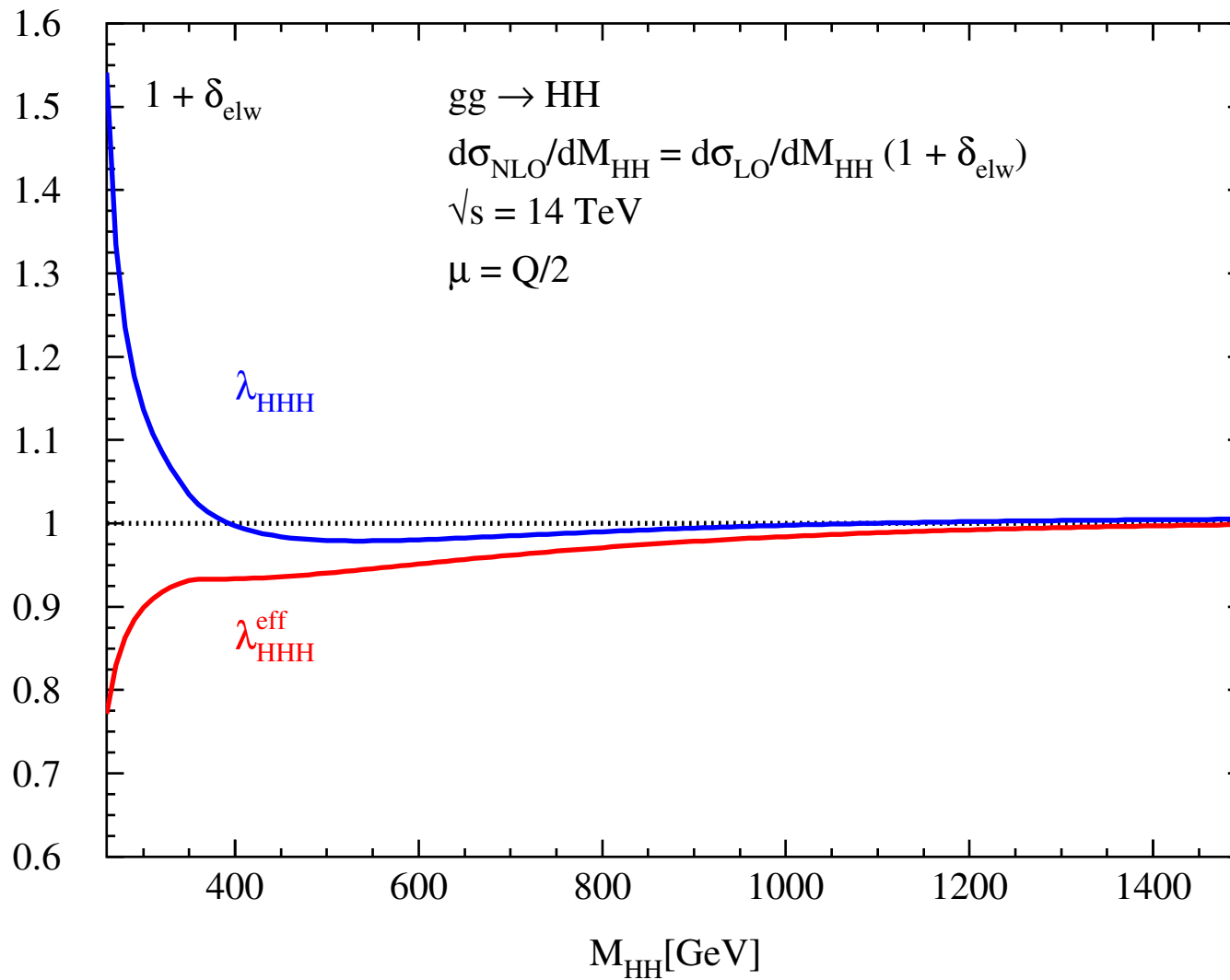
$$\lambda_{HHH}^{eff} = 3 \frac{M_H^2}{v} + \Delta\lambda_{HHH},$$

$$\lambda_{HHHH}^{eff} = 3 \frac{M_H^2}{v^2} + \Delta\lambda_{HHHH}$$

$$\Delta\lambda_{HHH} = -\frac{3m_t^4}{\pi^2 v^3},$$

$$\Delta\lambda_{HHHH} = -\frac{12m_t^4}{\pi^2 v^4}$$

$$\lambda_{HHH}^{eff} = 3 \frac{M_H^2}{v} - \frac{3m_t^4}{\pi^2 v^3} \approx 0.91 \times 3 \frac{M_H^2}{v}$$



$$\sigma = 1.002 \times \sigma_{\text{LO}} \quad (\lambda_{\text{HHH}})$$

$$\sigma = 0.938 \times \sigma_{\text{LO}} \quad (\lambda_{\text{HHH}}^{\text{eff}}) \leftarrow \text{disfavoured}$$

IV CONCLUSIONS

- scale and scheme uncertainties due to m_t relevant for large momenta
- Higgs pair production: m_t effects on top of LO $\sim -15\%$ for σ_{tot}
[larger for distributions]
- uncertainties due to factorization/renormalization scale and m_t scale/scheme choice @NNLO_{FTapprox} $\lesssim 25\%$
- combined uncertainties available for λ dependence, too.
- top-induced electroweak corrections: small for total cxn, larger for distributions
- effective radiatively corrected λ_{HHH}^{eff} disfavoured
[momentum dependence of same size]
- full elw. corrections $\sim 5 - 10\%$ (in absolute terms)

BACKUP SLIDES

II CALCULATION

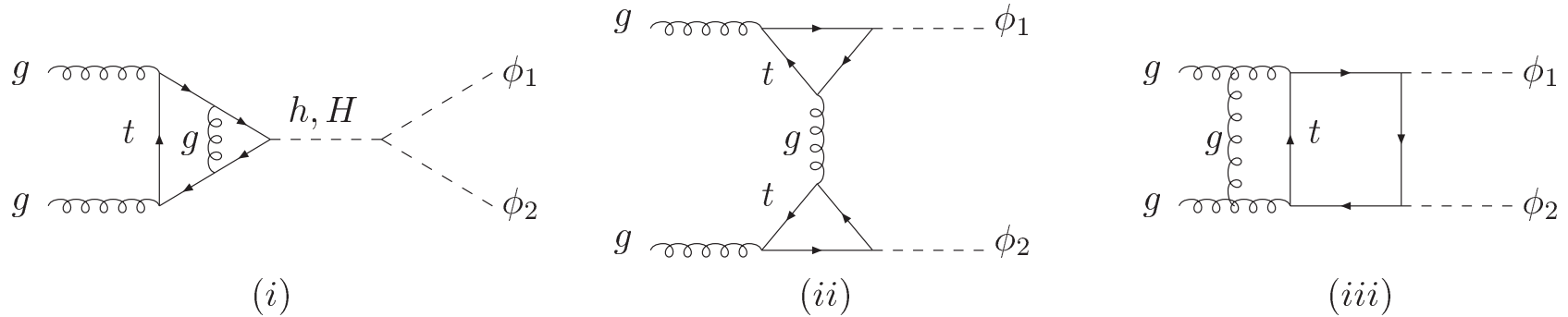
$$\sigma_{\text{NLO}}(pp \rightarrow HH + X) = \sigma_{\text{LO}} + \Delta\sigma_{\text{virt}} + \Delta\sigma_{gg} + \Delta\sigma_{gq} + \Delta\sigma_{q\bar{q}}$$

$$\begin{aligned} \sigma_{\text{LO}} &= \int_{\tau_0}^1 d\tau \frac{d\mathcal{L}^{gg}}{d\tau} \hat{\sigma}_{\text{LO}}(Q^2 = \tau s) \\ \Delta\sigma_{\text{virt}} &= \frac{\alpha_s(\mu)}{\pi} \int_{\tau_0}^1 d\tau \frac{d\mathcal{L}^{gg}}{d\tau} \hat{\sigma}_{\text{LO}}(Q^2 = \tau s) \ C \\ \Delta\sigma_{gg} &= \frac{\alpha_s(\mu)}{\pi} \int_{\tau_0}^1 d\tau \frac{d\mathcal{L}^{gg}}{d\tau} \int_{\tau_0/\tau}^1 \frac{dz}{z} \hat{\sigma}_{\text{LO}}(Q^2 = z\tau s) \left\{ -z P_{gg}(z) \log \frac{M^2}{\tau s} \right. \\ &\quad \left. + d_{gg}(z) + 6[1 + z^4 + (1 - z)^4] \left(\frac{\log(1 - z)}{1 - z} \right)_+ \right\} \\ \Delta\sigma_{gq} &= \frac{\alpha_s(\mu)}{\pi} \int_{\tau_0}^1 d\tau \sum_{q, \bar{q}} \frac{d\mathcal{L}^{gq}}{d\tau} \int_{\tau_0/\tau}^1 \frac{dz}{z} \hat{\sigma}_{\text{LO}}(Q^2 = z\tau s) \left\{ -\frac{z}{2} P_{gq}(z) \log \frac{M^2}{\tau s(1 - z)^2} + d_{gq}(z) \right\} \\ \Delta\sigma_{q\bar{q}} &= \frac{\alpha_s(\mu)}{\pi} \int_{\tau_0}^1 d\tau \sum_q \frac{d\mathcal{L}^{q\bar{q}}}{d\tau} \int_{\tau_0/\tau}^1 \frac{dz}{z} \hat{\sigma}_{\text{LO}}(Q^2 = z\tau s) \ d_{q\bar{q}}(z) \end{aligned}$$

$$C \rightarrow \pi^2 + \frac{11}{2} + C_{\Delta\Delta}, \quad d_{gg} \rightarrow -\frac{11}{2}(1 - z)^3, \quad d_{gq} \rightarrow \frac{2}{3}z^2 - (1 - z)^2, \quad d_{q\bar{q}} \rightarrow \frac{32}{27}(1 - z)^3$$

(i) virtual corrections

47 gen. box diags, 8 tria diags (\leftarrow single Higgs), 1PR ($\leftarrow H, A \rightarrow Z\gamma$)



- two formfactors:

$$\mathcal{A}^{\mu\nu} = F_1 T_1^{\mu\nu} + F_2 T_2^{\mu\nu} \quad F_1 = C_\Delta F_\Delta + F_\square \quad F_2 = G_\square$$

$$T_1^{\mu\nu} = g^{\mu\nu} - \frac{q_1^\nu q_2^\mu}{(q_1 q_2)},$$

$$T_2^{\mu\nu} = g^{\mu\nu} + \frac{M_H^2 q_1^\nu q_2^\mu}{p_T^2 (q_1 q_2)} - 2 \frac{(q_2 p_1) q_1^\nu p_1^\mu}{p_T^2 (q_1 q_2)} - 2 \frac{(q_1 p_1) p_1^\nu q_2^\mu}{p_T^2 (q_1 q_2)} + 2 \frac{p_1^\nu p_1^\mu}{p_T^2}$$

$$P_1^{\mu\nu} = \frac{(1 - \epsilon) T_1^{\mu\nu} + \epsilon T_2^{\mu\nu}}{2(1 - 2\epsilon)} \quad P_2^{\mu\nu} = \frac{\epsilon T_1^{\mu\nu} + (1 - \epsilon) T_2^{\mu\nu}}{2(1 - 2\epsilon)}$$

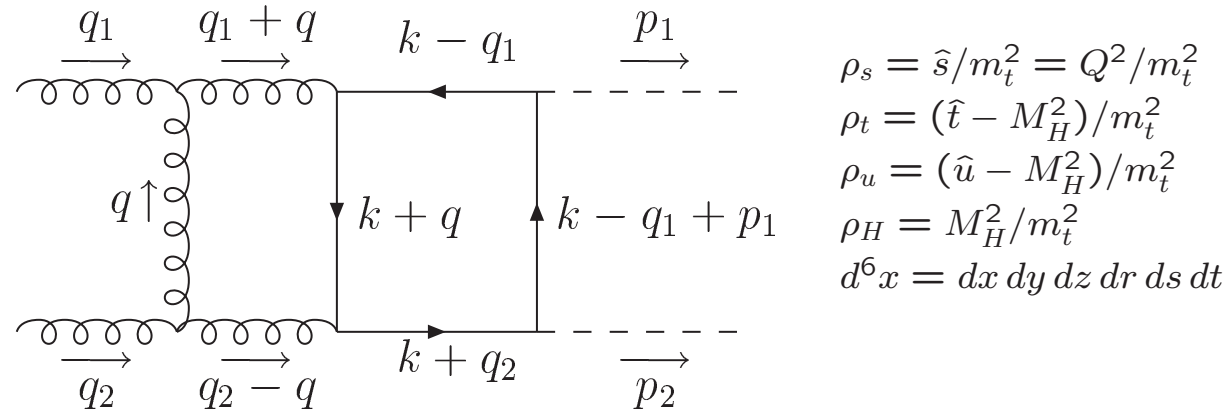
$$P_1^{\mu\nu} \mathcal{A}_{\mu\nu} = F_1 \quad P_2^{\mu\nu} \mathcal{A}_{\mu\nu} = F_2$$

- full diagram w/o tensor reduction \rightarrow 6-dim. Feynman integrals

- UV-singularities: end-point subtractions

$$\int_0^1 dx \frac{f(x)}{(1-x)^{1-\epsilon}} = \int_0^1 dx \frac{f(1)}{(1-x)^{1-\epsilon}} + \int_0^1 dx \frac{f(x) - f(1)}{(1-x)^{1-\epsilon}} = \frac{f(1)}{\epsilon} + \int_0^1 dx \frac{f(x) - f(1)}{1-x} + \mathcal{O}(\epsilon)$$

- IR-sing.: IR-subtraction (based on struc. of integr. and rel. to HTL)



$$\Delta F_i = \Gamma(1+2\epsilon) \left(\frac{4\pi\mu_0^2}{m_t^2} \right)^{2\epsilon} \int_0^1 d^6x \frac{x^{1+\epsilon}(1-x)^{\epsilon} r^{1+\epsilon} s^{-\epsilon} H_i(\vec{x})}{N^{3+2\epsilon}(\vec{x})}$$

$$N(\vec{x}) = ar^2 + br + c$$

$$a = x(1-x)ys \left[-\rho_s(1-y-t) + \rho_t yz - \rho_u z(1-y-t) + \rho_H yz^2 \right]$$

$$b = 1 - \rho_s x \left\{ xy(1-y) + (1-x)[(1-s)(1-y-t) + yst] \right\} - \rho_H xyz(1-xyz) \\ - \rho_t xyz[1-xy - (1-x)(1-s)] - \rho_u xyz[x(1-y) + (1-x)st]$$

$$c = -\rho_s x(1-x)(1-s)t$$

- subtract integrand with linear denominator

$$\Delta F_i = \frac{\alpha_s}{\pi} \Gamma(1 + 2\epsilon) \left(\frac{4\pi\mu_0^2}{m_t^2} \right)^{2\epsilon} (G_1 + G_2)$$

$$G_1 = \int_0^1 d^6x x^{1+\epsilon} (1-x)^\epsilon r^{1+\epsilon} s^{-\epsilon} \left\{ \frac{H_i(\vec{x})}{N^{3+2\epsilon}(\vec{x})} - \frac{H_i(\vec{x})|_{r=0}}{(br+c)^{3+2\epsilon}} \right\}$$

$$G_2 = \int_0^1 d^6x x^{1+\epsilon} (1-x)^\epsilon r^{1+\epsilon} s^{-\epsilon} \frac{H_i(\vec{x})|_{r=0}}{(br+c)^{3+2\epsilon}}$$

- $G_2 \rightarrow$ hypergeom. fct. after r-integration [arg \rightarrow 1/arg]
- thresholds: $Q^2 \geq 0, 4m_t^2 \rightarrow$ IBP \rightarrow reduction of power of denominator [$m_t^2 \rightarrow m_t^2(1 - ih)$]

$$\int_0^1 dx \frac{f(x)}{(a+bx)^3} = \frac{f(0)}{2a^2b} - \frac{f(1)}{2b(a+b)^2} + \int_0^1 dx \frac{f'(x)}{2b(a+bx)^2}$$

- extrapolation to NWA ($h \rightarrow 0$): Richardson extrapolation

1911

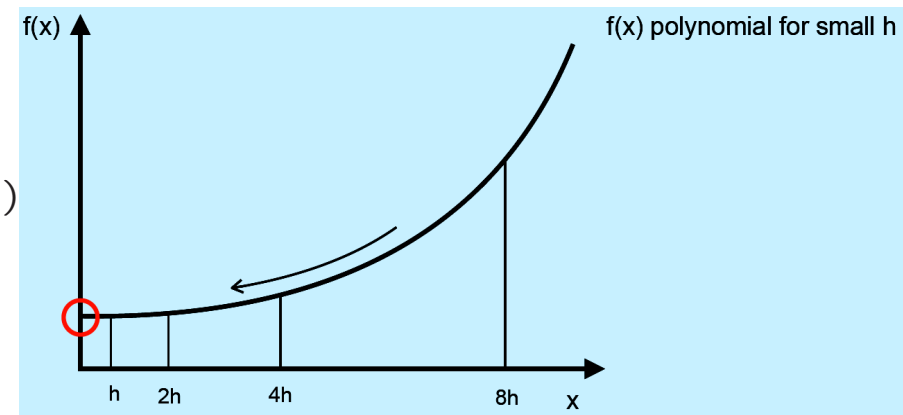
$$M_2 = 2f(h) - f(2h) = f(0) + \mathcal{O}(h^2)$$

$$M_4 = \{8f(h) - 6f(2h) + f(4h)\}/3 = f(0) + \mathcal{O}(h^3)$$

$$M_8 = \{64f(h) - 56f(2h) + 14f(4h) - f(8h)\}/21 = f(0) + \mathcal{O}(h^4)$$

etc.

$$[h \geq 0.05]$$



- renormalization: α_s : $\overline{\text{MS}}$, 5 flavours
 m_t : on-shell
- PS-integration \rightarrow 7-dim. integrals for $d\sigma/dQ^2$
- subtraction of HTL \rightarrow IR-finite mass effects
- add back HTL results \leftarrow HPAIR

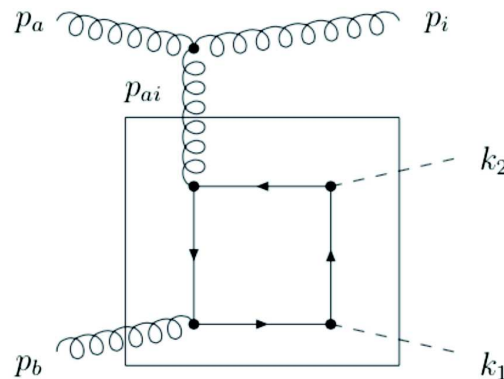
(ii) real corrections

- full matrix elements generated with FeynArts and FormCalc
- matrix elements in HTL involving full LO sub-matrix elements subtracted \rightarrow IR-, COLL-finite [adding back HTL results \leftarrow HPAIR]

$$\sum \overline{|\mathcal{M}_{gg}|^2} = \sum \overline{|\tilde{\mathcal{M}}_{LO}|^2} \frac{24\pi^2 \alpha_s}{Q^4 \pi} \left\{ \frac{s^4 + t^4 + u^4 + Q^8}{stu} - 4 \frac{\epsilon}{1-\epsilon} Q^2 \right\}$$

$$\sum \overline{|\mathcal{M}_{gq}|^2} = \sum \overline{|\tilde{\mathcal{M}}_{LO}|^2} \frac{32\pi^2 \alpha_s}{3Q^4 \pi} \left\{ \frac{s^2 + u^2}{-t} + \epsilon \frac{(s+u)^2}{t} \right\}$$

$$\sum \overline{|\mathcal{M}_{q\bar{q}}|^2} = \sum \overline{|\tilde{\mathcal{M}}_{LO}|^2} \frac{256\pi^2 \alpha_s}{9Q^4 \pi} (1-\epsilon) \left\{ \frac{t^2 + u^2}{s} - \epsilon \frac{(t+u)^2}{s} \right\}$$



- m_t scale/scheme uncertainties at LO:

$$\frac{d\sigma(gg \rightarrow HH)}{dQ} \Big|_{Q=300 \text{ GeV}} = 0.01656^{+62\%}_{-2.4\%} \text{ fb/GeV}$$

$$\frac{d\sigma(gg \rightarrow HH)}{dQ} \Big|_{Q=400 \text{ GeV}} = 0.09391^{+0\%}_{-20\%} \text{ fb/GeV}$$

$$\frac{d\sigma(gg \rightarrow HH)}{dQ} \Big|_{Q=600 \text{ GeV}} = 0.02132^{+0\%}_{-48\%} \text{ fb/GeV}$$

$$\frac{d\sigma(gg \rightarrow HH)}{dQ} \Big|_{Q=1200 \text{ GeV}} = 0.0003223^{+0\%}_{-56\%} \text{ fb/GeV}$$

$$F_i = F_{i,LO} + \Delta F_i$$

$$\Delta F_i = \Delta F_{i,HTL} + \Delta F_{i,mass}$$

- pole mass:

$$F_{1,LO} \rightarrow 4 \frac{m_t^2}{\hat{s}}$$

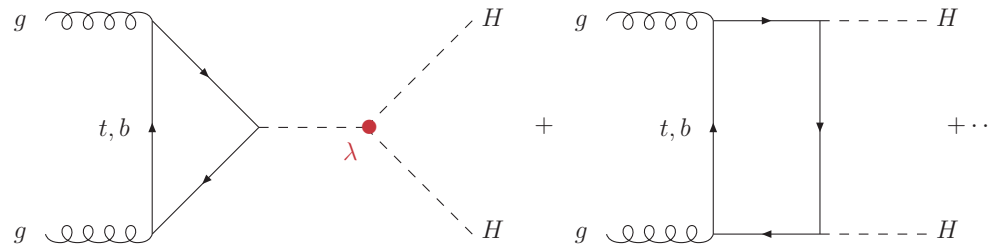
$$F_{2,LO} \rightarrow -\frac{m_t^2}{\hat{s}\hat{t}(\hat{s} + \hat{t})} \{(\hat{s} + \hat{t})^2 L_{1ts}^2 + \hat{t}^2 L_{ts}^2 + \pi^2 [(\hat{s} + \hat{t})^2 + \hat{t}^2]\}$$

- $\overline{\text{MS}}$ mass:

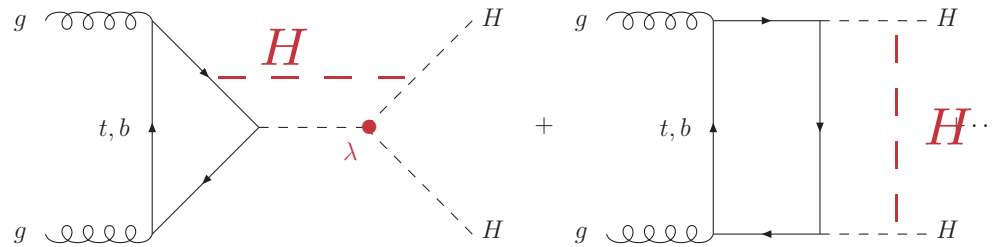
$$F_{1,LO} \rightarrow 4 \frac{\overline{m}_t^2(\mu_t)}{\hat{s}}$$

$$F_{2,LO} \rightarrow -\frac{\overline{m}_t^2(\mu_t)}{\hat{s}\hat{t}(\hat{s} + \hat{t})} \{(\hat{s} + \hat{t})^2 L_{1ts}^2 + \hat{t}^2 L_{ts}^2 + \pi^2 [(\hat{s} + \hat{t})^2 + \hat{t}^2]\}$$

- different scales for y_t in triangle (Q) and box (M_H) diagrams?
 → has to hold at all orders



- different scales for y_t in triangle (Q) and box (M_H) diagrams?
 → has to hold at all orders



elw. corrections

⇒ same scales in all diagrams

$$\sigma_{\text{LO}} = \int_{\tau_0}^1 d\tau \frac{d\mathcal{L}^{gg}}{d\tau} \hat{\sigma}_{\text{LO}}(Q^2 = \tau s)$$

$$\frac{d\mathcal{L}^{gg}}{d\tau} = \int_{\tau}^1 \frac{dx}{x} g(x, \mu_F) g\left(\frac{\tau}{x}, \mu_F\right)$$

$$\hat{\sigma}_{\text{LO}} = \frac{G_F^2 \alpha_s^2(\mu_R)}{512(2\pi)^3} \int_{\hat{t}_-}^{\hat{t}_+} d\hat{t} [|C_{\Delta} F_{\Delta} + F_{\square}|^2 + |G_{\square}|^2]$$

$$\hat{t}_{\pm} = -\frac{1}{2} \left[Q^2 - 2M_H^2 \mp Q^2 \sqrt{1 - 4\frac{M_H^2}{Q^2}} \right]$$

$$\lambda_{\text{HHH}} = 3 \frac{M_H^2}{v}$$

$$C_{\Delta} = \frac{\lambda_{\text{HHH}} v}{(Q^2 - M_H^2)}$$

$$\text{HTL: } F_{\Delta} \rightarrow 2/3, \quad F_{\square} \rightarrow -2/3, \quad G_{\square} \rightarrow 0$$

$$C_{\Delta} F_{\Delta} \rightarrow C_{\Delta} F_{\Delta} (1 + \Delta_{\Delta})$$

$$F_{\square} \rightarrow F_{\square} (1 + \Delta_{\square})$$

$$\Delta_{\Delta} = \delta_1 + \Delta_{\text{HHH}}$$

$$\Delta_{\square} = \eta_1$$

$$\Delta_{HHH} = \Delta_{vertex} + \Delta_{self} + \Delta_{CT}$$

$$\Delta_{vertex} = \frac{m_t^4}{v^2 M_H^2} \frac{8}{(4\pi)^2} \left\{ B_0(Q^2; m_t, m_t) + 2B_0(M_H^2; m_t, m_t) \right. \\ \left. + \left(4m_t^2 - \frac{Q^2 + 2M_H^2}{2} \right) C_0(Q^2, M_H^2, M_H^2; m_t, m_t, m_t) \right\} + \frac{T_1}{v M_H^2}$$

$$\Delta_{self} = \frac{\Sigma_H(Q^2)}{Q^2 - M_H^2} + \frac{1}{2} \Sigma'_H(M_H^2)$$

$$\Delta_{CT} = \frac{\delta M_H^2}{Q^2 - M_H^2} + \frac{\delta \lambda_{HHH}}{\lambda_{HHH}}$$

$$\Sigma_H(Q^2) = 3 \frac{T_1}{v} + 6 \frac{m_t^2}{(4\pi)^2 v^2} \left\{ 2A_0(m_t) + (4m_t^2 - Q^2) B_0(Q^2; m_t, m_t) \right\} + \mathcal{O}(m_t^0)$$

$$\Sigma'_H(Q^2) = 6 \frac{m_t^2}{(4\pi)^2 v^2} \left\{ (4m_t^2 - Q^2) B'_0(Q^2; m_t, m_t) - B_0(Q^2; m_t, m_t) \right\} + \mathcal{O}(m_t^0)$$

$$\frac{T_1}{v} = -12 \frac{m_t^2}{(4\pi)^2 v^2} A_0(m_t)$$

$$\frac{\delta \lambda_{HHH}}{\lambda_{HHH}} = \frac{\delta M_H^2}{M_H^2} + \frac{1}{2} \frac{\Sigma_W(0)}{M_W^2}$$

$$\frac{\Sigma_W(0)}{M_W^2} = 2 \frac{T_1}{v M_H^2} + \frac{2m_t^2}{(4\pi)^2 v^2} \left\{ B_0(0; m_t, 0) + 2B_0(0; m_t, m_t) + m_t^2 B'_0(0; m_t, 0) \right\} + \mathcal{O}(m_t^0)$$

$$\delta M_H^2 = -\Sigma_H(M_H^2)$$