Extended Scalar Sectors From All Angles @ CERN, October 21-26, 2024

Kennesaw State University, Georgia, USA

KENNESAW STATE UNIVERSITY

Dissecting Multi-Higgs Boson Production in New Physics Models

Symmetry Breaking

→ *h* **is the Higgs boson! (LHC, 2012)**

- → **the Higgs boson's self-interactions.**
- → **Determine shape of this potential by measuring:**
	-

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4 **Andreas Papaefstathiou**
Andreas Papaefstathiou [e.g. **AP**, Sakurai, **arXiv:1508.06524**, Fuks, Kim, Lee, **arXiv:1510.07697** & **arXiv:1704.04298**, **AP**, Tetlalmatzi-Xolocotzi, Zaro, **arXiv:1909.09166**, …]

- → **the Higgs boson's self-interactions.**
- → **Determine shape of this potential by measuring:**

4 **Andreas Papaefstathiou**
Andreas Papaefstathiou Fuks, Kim, Lee, **arXiv:1510.07697** & **arXiv:1704.04298**, **AP**, Tetlalmatzi-Xolocotzi, Zaro, **arXiv:1909.09166**, …]

Karkout, **AP**, Postma, Tetlalmatzi-Xolocotzi, van de Vis, du Pree, a**rXiv:2404.12425**, **AP**, Tetlalmatzi-Xolocotzi, **aXiv:2312.13562**] **SEE LATER!**

• \exists factor of $\mathcal{O}(10^{-3})$ each time you "draw" an extra Higgs boson @ pp colliders.

σ(*h*) ∼ 50 pb

SM, 14 TeV

• \exists factor of $\mathcal{O}(10^{-3})$ each time you "draw" an extra Higgs boson @ pp colliders.

σ(*h*) ∼ 50 pb $× 6(1)$

$σ(hh) \sim 40$ fb

SM, 14 TeV

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σ(*h*) ∼ 50 pb $×$ 0(1)

SM, 14 TeV

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~ ×60 increase in cross section as pp COM $14 TeV \rightarrow 100 TeV.$

• Cranking up the **pp COM energy** could help!

SM hhh Boson Production "Fun" Facts

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~ ×60 increase in cross section as pp COM 14 TeV → **100 TeV.**

SM hhh Boson Production "Fun" Facts

THE SECRET **iNGREDIENT** IS ALWAYS LOVE

THE SECRET iNGREDIENT IS ALWAYS New Physics

THE SECRET **INGREDIENT** IS ALWAYS New Physics

A. **hhh** in SM+2 singlet scalar fields, B. **hhh** with anomalous couplings.

Here:

A. **hhh** in SM+2 singlet scalar fields,

SM + Two Real Singlet Scalars [= TRSM]

- Add two real singlet scalar fields S , X .
- & impose discrete \mathcal{Z}_2 symmetries:
- *S* $2: S \to -S, X \to X$
- *X* $\frac{X}{2}: X \to -X, \ S \to S$

5^{4} + $\bullet S^{2}$ + $\bullet S^{4}$ + $\bullet X^{2}$ + $\bullet X^{4}$

 $2S^2 + 2\phi^2 + 2\phi^2$

⇒ TRSM scalar potential:

$$
\mathcal{V}(\phi, S, X) = \bullet |\phi|^2 + |\phi|
$$

+
$$
S^2 X^2
$$

+
$$
|\phi|^2 S^2 +
$$

SM + Two Real Singlet Scalars [= TRSM]

- Go through electroweak symmetry breaking…
- \Rightarrow Get **three** scalar bosons: $h_1, h_2, h_3 \rightarrow h_1 \approx$ SM-like Higgs boson.
- **hhh** detectable at the LHC! ⇒ **[AP**, Robens, Tetlalmatzi-Xolocotzi, **arXiv:2101.00037]** ✨**see Gilberto Tetlalmatzi-Xolocotzi's talk!**✨

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- ⇒ **hhh** detectable at the LHC! ^[AP, Robens, Tetlalmatzi-Xolocotzi, arXiv:2101.00037]

→ **"Double-Resonant Enhancement"!**

[AP, Robens, Tetlalmatzi-Xolocotzi, arXiv:2101.00037]

$$
through: pp \rightarrow h
$$

✨**see Gilberto Tetlalmatzi-Xolocotzi's talk!**✨

hhh in the TRSM [Karkout, AP, Postma, Tetlalmatzi-Xolocotzi, eq. (2.14), and determined the scales at which perturbativity gets violated and when we hhh in

van de Vis, du Pree, **arXiv:2404.12425**]

Viable points with $\sigma > 10 \times \sigma_{SM}(gg \rightarrow hhh)$ @13.6 TeV

hhh in the TRSM [Karkout, AP, Postma, Tetlalmatzi-Xolocotzi, eq. (2.14), and determined the scales at which perturbativity gets violated and when we hhh in

van de Vis, du Pree, **arXiv:2404.12425**]

\bullet Narrow-width $(\Gamma, \times M, \& \Gamma, \times M)$ \sim 130110W WIQUITIPLES \sim $m_2 \propto$ 13.8 m_3 tribution from *pp* Ñ *h*³ Ñ *h*2*h*¹ Ñ *h*1*h*1*h*1. Only points with a factor 10 enhancement or greater are shown. The density of points increases from the dark blue to yellow shade. two heavy scalars *h*³ and *h*2. \mathbf{v} and different singlets with intermediate states can contribute states can contribute t_{max} boson production $\mathbf{3}$ s mixing and be overcome by resonance e \mathbf{u} if the intermediate states if the intermediate states if the intermediate states if the intermediate states in the intermediate states in the intermediate states in the in are produced on-shell and above the threshold *M*³ ° *M*² ` *M*¹ and *M*² ° 2*M*1. In the • Narrow-width $(\Gamma_2 \ll M_2 \& \Gamma_3 \ll M_3)$ double-resonant production **suggests a "simplified" factorized approach**.

 \bigcap \bigcap • Coming soon! [AP, Tetlalmatzi-Xolocotzi, arXiv:24!?.!?!?] **15td** [**AP**, Tetlalmatzi-Xolocotzi, **arXiv:24**⁉**.**⁉⁉]

FO-EWPT and hhh in the TRSM

- related to **electro-weak baryogenesis**?
- and if so, will this lead to enhanced hhh at the LHC? If an de Vis, du Pree, arXiv:2404.12425]

• *Q*: Can there be a **First-Order Electroweak Phase Transition** in the **TRSM**,

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FO-EWPT and hhh in the TRSM

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• *Q*: Can there be a **First-Order Electroweak Phase Transition** in the **TRSM**,

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 $\nu(\boldsymbol{\phi})$

 $\boldsymbol{\phi}$

- ‣ **FO-EWPT** & **enhanced hhh** are **mutually exclusive**!
- ‣ "barrier" **not** generated **if both new scalars attain a VEV**,
- ‣ and **non-zero VEVs are necessary** for sufficient mixing!

 \rightarrow Removing the \mathscr{L}_{2} restrictions might help!

•**NO!** ✨**see Osama Karkout's talk!**✨

- Add **higher-dimensional operators** to the SM Lagrangian!
	- → Capture effects of new particles at Scales \gg Collision Energies.
-

• e.g. Add D=6 operators **relevant to multi-Higgs boson production**, of the form :

[see e.g. Goertz, AP, Yang, Zurita, $arXiv:1410.3471$ for $pp \rightarrow hh$ study in this framework]

$$
S_{M} \rightarrow \mathcal{L}_{h^n} \supset -\mu^2 |\phi|^2 - \lambda |\phi|^4 - \left(y_t \bar{Q}_L \phi^c t_R + y_b \bar{Q}_L \phi b_R + \text{h.c.}\right)
$$

+
$$
\frac{c_H}{2\Lambda^2} (\partial^\mu |\phi|^2)^2 - \frac{c_6}{\Lambda^2} \lambda_{SM} |\phi|^6 + \frac{\alpha_s c_g}{4\pi\Lambda^2} |\phi|^2 G_{\mu\nu}^a G_{a}^{\mu\nu}
$$

-
$$
\left(\frac{c_t}{\Lambda^2} y_t |\phi|^2 \bar{Q}_L \phi^c t_R + \frac{c_b}{\Lambda^2} y_b |\phi|^2 \bar{Q}_L \phi b_R + \text{h.c.}\right)
$$

For 1-loop computations see: **smeft@nlo**: [Degrande, Durieux, Maltoni, Mimasu, Vryonidou, Zhang, **arXiv:2008.11743**]

• Go through EWSB… ⇒ in terms of the physical scalar Higgs boson *h*:

$$
\mathcal{L}_{\text{D=6}} \supset -\frac{m_h^2}{2v} \left(1 + c_6\right) h^3 - \frac{m_h^2}{8v^2} \left(1 + 6c_6\right) h^4 \n+ \frac{\alpha_s c_g}{4\pi} \left(\frac{h}{v} + \frac{h^2}{2v^2}\right) G^a_{\mu\nu} G^{\mu\nu}_a \n- \left[\frac{m_t}{v} \left(1 + c_t\right) \bar{t}_L t_R h + \frac{m_b}{v} \left(1 + c_b\right) \bar{b}_L b_R h + \text{h.c.}\right] \n- \left[\frac{m_t}{v^2} \left(\frac{3c_t}{2}\right) \bar{t}_L t_R h^2 + \frac{m_b}{v^2} \left(\frac{3c_b}{2}\right) \bar{b}_L b_R h^2 + \text{h.c.}\right] \n- \left[\frac{m_t}{v^3} \left(\frac{c_t}{2}\right) \bar{t}_L t_R h^3 + \frac{m_b}{v^3} \left(\frac{c_b}{2}\right) \bar{b}_L b_R h^3 + \text{h.c.}\right],
$$

• Go through EWSB… ⇒ in terms of the physical scalar Higgs boson *h*:

$$
\frac{m_h^2}{2v} (1+c_6) h^3 - \frac{m_h^2}{8v^2} (1+6c_6) h^4
$$
\n
$$
+\frac{\alpha_s c_g}{4\pi} \left(\frac{h}{v} + \frac{h^2}{2v^2}\right) G^a_{\mu\nu} G^{\mu\nu}_a
$$
\n
$$
-\left[\frac{m_t}{v} (1+c_t) \bar{t}_L t_R h + \frac{m_b}{v} (1+c_b) \bar{b}_L b_R h + \text{h.c.}\right]
$$
\n
$$
-\left[\frac{m_t}{v^2} \left(\frac{3c_t}{2}\right) \bar{t}_L t_R h^2 + \frac{m_b}{v^2} \left(\frac{3c_b}{2}\right) \bar{b}_L b_R h^2 + \text{h.c.}\right]
$$
\n
$$
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$$

• Go through EWSB... \overline{up} \Rightarrow in terms of the physical scalar Higgs boson *h*:

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+
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\frac{\alpha_s c_g}{4\pi} \left(\frac{h}{v} + \frac{h^2}{2v^2}\right) G^a_{\mu\nu} G^{\mu\nu}_a
$$

-
$$
\left[\frac{m_t}{v} \left(1 + c_t\right) \bar{t}_L t_R h + \frac{m_b}{v} \left(1 + c_b\right) \bar{b}_L b_R h + \text{h.c.}\right]
$$

-
$$
\left(\frac{h}{v^2} \left(\frac{3c_t}{2}\right) \bar{t}_L t_R h^2 + \frac{m_b}{v^2} \left(\frac{3c_b}{2}\right) \bar{b}_L b_R h^2 + \text{h.c.}\right]
$$

-
$$
\left[\frac{m_t}{v^3} \left(\frac{c_t}{2}\right) \bar{t}_L t_R h^3 + \frac{m_b}{v^3} \left(\frac{c_b}{2}\right) \bar{b}_L b_R h^3 + \text{h.c.}\right],
$$

q

g

 \mathcal{S}

g

g

• A slightly more "general" picture is obtained by "**dissociating**" the operators as:

$$
\mathcal{L}_{\text{Pheno}} \supset -\frac{m_h^2}{2v} (1+d_3) h^3 - \frac{m_h^2}{8v^2} (1+d_4) h^4
$$

+
$$
\frac{\alpha_s}{4\pi} \left(c_{g1} \frac{h}{v} + c_{g2} \frac{h^2}{2v^2} \right) G_{\mu\nu}^a G_a^{\mu\nu}
$$

-
$$
\left[\frac{m_t}{v} (1+c_{t1}) \bar{t}_L t_R h + \frac{m_b}{v} (1+c_{b1}) \bar{b}_L b_R h + \text{h.c.} \right]
$$

-
$$
\left[\frac{m_t}{v^2} \left(\frac{3c_{t2}}{2} \right) \bar{t}_L t_R h^2 + \frac{m_b}{v^2} \left(\frac{3c_{b2}}{2} \right) \bar{b}_L b_R h^2 + \text{h.c.} \right]
$$

-
$$
\left[\frac{m_t}{v^3} \left(\frac{c_{t3}}{2} \right) \bar{t}_L t_R h^3 + \frac{m_b}{v^3} \left(\frac{c_{b3}}{2} \right) \bar{b}_L b_R h^3 + \text{h.c.} \right],
$$

Recover D=6 by setting:

\n
$$
d_{3} = c_{6},
$$
\n
$$
d_{4} = 6c_{6},
$$
\n
$$
c_{g1} = c_{g2} = c_{g},
$$
\n
$$
c_{f1} = c_{f2} = c_{f3} = c_{f}.
$$

[**AP**, Tetlalmatzi-Xolocotzi, **arXiv:2312.13562**]

*Note***: This can be also be motivated via the Electro-weak Chiral Lagrangian, [see e.g. Buchalla, Catá, Krause arXiv:1307.5017]**

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$$

+
$$
\frac{\alpha_s}{4\pi} \left(\frac{h}{c_g} \frac{h}{v} + \frac{h^2}{c_g} \frac{h^2}{2v^2} \right) G_{\mu\nu}^a G_a^{\mu\nu} \qquad \text{instead of } c_g
$$

-
$$
\left[\frac{m_t}{v} (1+c_{t1}) \bar{t}_L t_R h + \frac{m_b}{v} (1+c_{b1}) \bar{b}_L b_R h + \text{h.c.} \right]
$$

-
$$
\left[\frac{m_t}{v^2} \left(\frac{3c_{t2}}{2} \right) \bar{t}_L t_R h^2 + \frac{m_b}{v^2} \left(\frac{3c_{b2}}{2} \right) \bar{b}_L b_R h^2 + \text{h.c.} \right]
$$

-
$$
\left[\frac{m_t}{v^3} \left(\frac{c_{t3}}{2} \right) \bar{t}_L t_R h^3 + \frac{m_b}{v^3} \left(\frac{c_{b3}}{2} \right) \bar{b}_L b_R h^3 + \text{h.c.} \right],
$$

[**AP**, Tetlalmatzi-Xolocotzi, **arXiv:2312.13562**]

Recover D=6 by setting: $d_3 = c_6$, $d_4 = 6c_6$ $c_{g1} = c_{g2} = c_g$ $c_{f1} = c_{f2} = c_{f3} = c_f$.

*Note***: This can be also be motivated via the Electro-weak Chiral Lagrangian, [see e.g. Buchalla, Catá, Krause arXiv:1307.5017]**

D=6-Inspired Anomalous Couplings

• A slightly more "general" picture is obtained by "**dissociating**" the operators as:

$$
\mathcal{L}_{\text{Pheno}} \supset -\frac{m_h^2}{2v} (1+d_3) h^3 - \frac{m_h^2}{8v^2} (1+d_4) h^4 \n+ \frac{\alpha_s}{4\pi} \left(c_{g1} \frac{h}{v} + c_{g2} \frac{h^2}{2v^2} \right) G_{\mu\nu}^a G_{a}^{\mu\nu} \n- \left[\frac{m_t}{v} (1 + C_{t1}) \bar{t}_L t_R h + \frac{m_b}{v} (1 + c_{b1}) \bar{b}_L b_R h + \text{h.c.} \right] \qquad \text{Recover D=6 by setting:} \\ - \left[\frac{m_t}{v^2} \left(\frac{3C_0}{2} \right) \bar{t}_L t_R h^2 + \frac{m_b}{v^2} \left(\frac{3c_{b2}}{2} \right) \bar{b}_L b_R h^2 + \text{h.c.} \right] \qquad \qquad c_{g1} = c_{g2} = c_g, \\ - \left[\frac{m_t}{v^3} \left(\frac{C_0}{2} \right) \bar{t}_L t_R h^3 + \frac{m_b}{v^3} \left(\frac{c_{b3}}{2} \right) \bar{b}_L b_R h^3 + \text{h.c.} \right], \qquad \text{instead of } c_t
$$

[**AP**, Tetlalmatzi-Xolocotzi, **arXiv:2312.13562**]

*Note***: This can be also be motivated via the Electro-weak Chiral Lagrangian, [see e.g. Buchalla, Catá, Krause arXiv:1307.5017]**

 h^4

D=6-Inspired Anomalous Couplings

• **Further modify** to match more closely **LHC experiments' definitions**:

$$
\mathcal{L}_{\text{PhenoExp}} \supset -\lambda_{\text{SM}} v (1+d_3) h^3 - \frac{\lambda_{\text{SM}}}{4} (1+d_4) h^4
$$

+
$$
\frac{\alpha_s}{12\pi} \left(c_{g1} \frac{h}{v} - c_{g2} \frac{h^2}{2v^2} \right) G_{\mu\nu}^a G_a^{\mu\nu}
$$

-
$$
\left[\frac{m_t}{v} (1+c_{t1}) \bar{t}_L t_R h + \frac{m_b}{v} (1+c_{b1}) \bar{b} - \left[\frac{m_t}{v^2} c_{t2} \bar{t}_L t_R h^2 + \frac{m_b}{v^2} c_{b2} \bar{b}_L b_R h^2 + \text{h.c.} \right] - \left[\frac{m_t}{v^3} \left(\frac{c_{t3}}{2} \right) \bar{t}_L t_R h^3 + \frac{m_b}{v^3} \left(\frac{c_{b3}}{2} \right) \bar{b}_L \right]
$$

Obtain **CMS-like** parametrization by:

And **ATLAS-like** parametrization by:

$$
\kappa_{\lambda} = (1 + d_3),
$$

\n
$$
k_t = c_{t1},
$$

\n
$$
c_2 = c_{t2},
$$

\n
$$
c_g = c_{g1},
$$

\n
$$
c_{gg} = c_{2g}.
$$

$$
chhh = (1+d3),
$$

$$
cggh = 2cg1/3,
$$

$$
cgghh = -cg2/3.
$$

[**AP**, Tetlalmatzi-Xolocotzi, **arXiv:2312.13562**]

 $L^b_R h + h.c.$ + h.c. $\left(\frac{b^3}{2}\right)\bar{b}_L b_R h^3 + \text{h.c.}\right],$ Defined: $\lambda_{\text{SM}} = m_h^2 / 2v^2$.

Figure 8: Example Feynman diagrams with *one* EFT operator insertion contributing to Higgs boson triple production. The contribution of the production of the production of the production of the p
Higgs boson triple production. The production of the production of the production of the production of the pro

Monte Carlo Implementation of Anomalous Couplings • We have implemented a **MadGraph5_aMC@NLO "loop" model** for $\mathscr{L}_{\text{PhenoExp}}$. \mathbf{q} *g h g q h q h*

• Includes <u>Loop \times Tree Level interference</u> between diagrams.

[see: Hirschi, [https://cp3.irmp.ucl.ac.be/projects/madgraph/wiki/LoopInducedTimesTree](https://cp3.irmp.ucl.ac.be/projects/madgraph/wiki/LoopInducedTimesTree%5D)] <u>tt</u> *h*

• Model available at <u>[https://gitlab.com/apapaefs/multihiggs_loop_sm](https://gitlab.com/apapaefs/multihiggs_loop_sm%5D)</u>. [includes necessary patch for MG5_aMC].

ce h

• Most couplings **validated** vs. a **HERWIG** 7 $pp \rightarrow hh$ implementation, e.g.:

$\bar{t}h^3$

Model Validation [**AP**, Tetlalmatzi-Xolocotzi, **arXiv:2312.13562**]

• "New" non-trivial coupling that appears, $\alpha c_{t3} t\bar{t}h^3$ has been **validated** via an "EFT" limit, in the $t\bar{t} \rightarrow hhh$ process:

 0.5

 0.5

 -0.5

 0.5

 $0₀$

 -0.5

 0.5

 0.0

 0.5

 0.0

 -0.5

 -0.5

 0.0

 -0.5

 C_{91}

 C_{Q2}

 C_{t1}

hhh Cross Sections @ 13.6 TeV

 C_{b1}

 C_{t2}

 C_{b2}

 C_{t3}

 C_{b3}

 \mathcal{Q}_4

- Cross section as a multiple of the SM.
- \bullet n.b.: $\sigma_{\text{SM}} \sim 0.04$ fb at $LO@13.6$ TeV.
- Each 2D panel shown: **all other coefficients set to zero!**

Andreas Papaefstathiou $\mathbf{1}$

[**AP**, Tetlalmatzi-Xolocotzi, **arXiv:2312.13562**]

hhh Cross Sections @ 13.6 TeV

- Cross section as a multiple of the SM.
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Anomalous Couplings Constraints bottom quarks, with relevant coecients *ct*³ and *cb*3, as well as the anomalous modification to the Higgs boson's quartic interaction, related to the *d*⁴ coecient. While it is beyond the

-
- $f_{\rm eff}$ the two schools two schools that we examine $f_{\rm eff}$ the 13.6 μ Frojected conditionmed for od • **Projected constraints for other coefficients:**

• Other processes constrain (<u>at LO</u>) all coefficients **<u>except</u>** $\{c_{t3}, d_4\} \rightarrow$ only in hhh).

✨ For details, see: [**AP**, Tetlalmatzi-Xolocotzi, **arXiv:2312.13562**]✨!

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Anomalous Couplings Constraints

• Focusing on a model with **non-zero**: $\{c_{t2}, d_3, c_{t3}, d_4\}$:

[**AP**, Tetlalmatzi-Xolocotzi, **arXiv:2312.13562**]

[**AP**, Tetlalmatzi-Xolocotzi, **arXiv:2312.13562**]

h

constrained by *pp* → *hh*

[**AP**, Tetlalmatzi-Xolocotzi, **arXiv:2312.13562**] *h*

h

constrained by *pp* → *hh*

h

q

g

q

q

g

g

 \mathcal{C}

h

h

h

h q h

constrained by $pp \rightarrow hhh$

g

h

q

h

g

h

g

q

g

g

h

h

h

g

h

g

{

Anomalous Couplings Constraints

- Focusing on a model with only $\{c_{t2}, d_3, c_{t3}, d_4\}$,
- Using the **hhh** \rightarrow 6 **b-jet final state**, and marginalizing over $\{c_{t2}, d_3\}$ within projected constraints:

- **Anomalous Couplings Constraints** • Focusing on a model with only $\{c_{t2}, d_3, c_{t3}, d_4\}$, *d*⁴ r´28*.*0*,* 41*.*7s r´99*.*5*,* 152*.*9s r´24*.*9*,* 20*.*8s r´40*.*8*,* 23*.*1s **0us Couplings Constraints** LAP, Tetlalmatzi-Xolocot: [**AP**, Tetlalmatzi-Xolocotzi, **arXiv:2312.13562**]
- Using the **hhh** \rightarrow 6 **b-jet final state,** and **marginalizing** over $\{c_{t2}, d_3\}$ within projected constraints: $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$ and $\frac{1}{2}$ and $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$, $\frac{1}{2}$, over *ct*2, *d*³ and either *d*4, or *ct*3.

 \Rightarrow $\begin{aligned} c_{t3} &\sim \mathcal{O}(0.1 - 1) \\ d_4 &\sim \mathcal{O}(10) \end{aligned}$ either *d*4, or *ct*3. $d_4 \sim c_1$ to d_1 The situation is greatly improved, as expected, at the FCC-hh, where the range of *d*⁴ is *d*⁴ ∼ (10)

- **Multi-Higgs production processes** → crucial rôle in understanding EWSB.
	- **‣** e.g. **hhh** production → **Higgs quartic self-coupling**.
- **SM** hhh \rightarrow hopeless at the LHC,
- BUT: **Enhanced** in models with **extended scalar sectors** or **anomalous interactions**. **Could we see hints of hhh at the LHC?**
	- ‣ TRSM (**two singlets**) → through double-resonant process. **Related to baryogenesis?**
	- ‣ **Anomalous coupling picture** → an agnostic framework for **h**/**hh**/**hhh**.

MG5_aMC model: https://gitlab.com/apapaefs/multihiggs_loop_sm

• Questions merit investigation @ LHC & future colliders!

Summary & Outlook

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Summary & Outlook

Appendices

TRSM Monte Carlo Event Generation

- We have implemented a **MadGraph5_aMC@NLO (MG5_aMC) "loop" model for the TRSM:**
	- **MG5_aMC input parameters**: the three mixing angles, two masses/widths and **all** the scalar couplings (only 7 are independent in **TRSM**).
	- Comes with a **Python script** that:
		- allows conversion of M_2 , M_3 + <u>three</u> mixing angles + <u>two</u> VEVs to the **MG5_aMC** model input,
		- calculates several single-production cross sections, branching ratios, widths,
		- and writes associated **MG5_aMC** parameter card (**param_card.dat**) automatically.
	- **Get it at:** [https://gitlab.com/apapaefs/twosinglet.](https://gitlab.com/apapaefs/twosinglet)

[AP, Tania Robens, Gilberto Tetlalmatzi-Xolocotzi, arXiv:2101.00037]

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hhh with Anomalous Couplings

hhh: Final states

[AP, Sakurai, 1508.06524]

1510.07697.

 \rightarrow AP, Sakurai, 1508.06524, Chen, Yan, Zhao, Zhao, Zhong, 1510.04013, Fuks, Kim, Lee, → Fuks, Kim, Lee, 1510.07697, Fuks, Kim, Lee, 1704.04298. →Kilian, Sun, Yan, Zhao, Zhao, 1702.03554.

Assume: K-factor = 2.

[Maltoni, Vryonidou, Zaro, 1408.6542]

The 6b final state, analysis [AP, Gilberto Tetlalmatzi-Xolocotzi, Marco Zaro, **arXiv:1909.09166]**

- What can we learn about the anomalous couplings via **hhh** at 13.6 TeV?
- Begin by using the **6 b-jet final state!**
- **1. Require 6 tagged b-jets.**

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t_{max} is pseudo-rapidity coverage for the identified *b*-jets on \mathbf{r} on \mathbf{r} on \mathbf{r} on \mathbf{r} coupming the main at 16.6 possible.

3. For each pairing construct:

$$
\chi^2 = \sum_{qr \in \text{pairings } I} (M_{qr} - m_h^2)^2
$$

where *Mqr* is the invariant mass of the *b*-jet pairing *qr* in the **≡ sum of squared differences from Higgs mass (~125 GeV)**

30 Andreas Papaefstathiou Andreas Papaefstathiou

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 $\chi^2_{\rm n}$

⇒ **4. Pairing that gives minimum** *χ2* **determines "reconstructed Higgs boson".** where *Mqr* is the invariant mass of the *b*-jet pairing *qr* in the actements ictonsulatica inggs boson.

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$$

≡ sum of squared differences from Higgs mass (~125 GeV)

h^i_r *r* [→]**Higgs boson candidates**

$$
]\mathord{\mathrm{GeV}}, i=1,2,3
$$

 \langle [3.5*,*3*.5,*3*.5*], (*i, j*) = [(1*,*2)*,*(1*,3*)*,*(2*,3*)] $\leq 8, 8, 11$ GeV **the three terms in** χ^2 _{*min*}.

signal/backgrounds after analysis $\frac{1}{2}$ an estimate of the expected increase in cross section from leading order to next-to-leading order. The fourth column gives $\mathcal O$ and the final column gives the expected number of events at 20 ab1 of integrated luminosity at 20 ab1 of integrated number of events at 20 ab1 of integrated number of events at 1000 ab1 of integrated luminosity at

Reducible backgrounds (1000) to the "worst-case" of 80%, see Appendix $\mathcal{L}(\mathcal{L})$ of 80%, see Appendix $\mathcal{L}(\mathcal{L})$ d. A. Shows section of the starting cross section of the starting cross section of the starting cross section

 $-pjets)$ (pb)

⇒ Assuming perfect b-tagging + identical analysis efficiency to QCD 6b:
 $\frac{1}{2}$

to QCD 6 *b*-jet production. We do not consider process that Keducibie back

c.f. $\sigma_{GEN}(6b) = 26.15 \text{ pb}$ applied:

 α ^o α backgrounds are expected to contribute *O COMMORTON HOM TCAGEDIC*
O rounds, for perfect *b*-tagging to *Pb*!*^b* = 0*.*8, respectively. Therefore →**~10% contribution from reducible backgrounds.**

for $P(b-tagging) = 0.8$:

 $P_{a\rightarrow b}=0.1$

 θ th θ is the cross section of the cross section estimates, see the cross section estimates, see the cross section estimates, θ t_{S} →**~30% contribution.**

TRSM hhh → **6b analysis details**

Introduce two observables: $\chi^{2,(4)} = \sum$

qr∈*I* $\left(M_{qr}-M_1\right)$ 2 $\chi^{2,(6)} = \sum (M_{qr} - M_1)$ *qr*∈*J* 2

\rightarrow constructed from different pairings of 4 and 6 b-tagged jets, M_{qr} is the invariant mass of the pairing *qr*.

TRSM hhh → **6b analysis details**

Table 3. The optimised selection cuts for each of the benchmark points within **BP3** shown in table 2. The cuts not shown above are common for all points, as follows: $|\eta|_b < 2.35$, $\Delta m_{\text{min, med, max}} <$ $[15, 14, 20]$ GeV, $p_T(h_1^i) > [50, 50, 0]$ GeV, $\Delta R(h_1^i, h_1^j) < 3.5$ and $\Delta R_{bb}(h_1) < 3.5$. For some of the points a m_{4b}^{inv} cut is not given, as this was found to not have an impact when combined with the m_{6b}^{inv} cut.

TRSM hhh → **6b analysis details (Signal vs Bkg)** [15*,* 14*,* 20] GeV, *p^T* (*hⁱ* ¹) *>* [50*,* 50*,* 0] GeV, *R*(*hⁱ* 1*, h^j* ¹) *<* 3*.*5 and *Rbb*(*h*1) *<* 3*.*5. For some of the $ph \rightarrow 6b$ analysis details (Signal w

Table 4. The resulting selection efficiencies, $\varepsilon_{\text{Sig.}}$ and $\varepsilon_{\text{Bkg.}}$, number of events, *S* and *B* for the signal and background, respectively, and statistical significances for the sets of cuts presented in table 3. A *b*-tagging efficiency of 0.7 has been assumed. The number of signal and background events are provided at an integrated luminosity of 300 fb⁻¹. Results for 3000 fb⁻¹ are obtained via simple extrapolation. The significance is given at both values of the integrated luminosity excluding (including) systematic errors in the background according to Eq. (5.1) (or Eq. (5.2) with $\sigma_b = 0.1 \times B$).

TRSM BP3 Definition

Parameter

 v_X

TRSM BP3 Benchmark Point Info $\sqrt{ }$

Table 5. The total widths and new scalar branching ratios for the parameter points considered in the analysis. For the SM-like h_1 , we have $M_1 = 125 \,\text{GeV}$ and $\Gamma_1 = 3.8 \,\text{MeV}$ for all points considered. The other input parameters are specified in table 1. The on-shell channel $h_3 \to h_2 h_2$ is kinematically forbidden for all points considered here.

hhh in the TRSM [14 TeV]

• In BP3: All params fixed $\overline{\text{except }M_2, M_3!}$

B (263, 455) 50.36 **Cross section can be much higher than in the SM!** \rightarrow **c.f. SM:** $\sigma \sim 0.1$ fb @ 14 TeV.

• Focus on a particular family of benchmark points: **"Benchmark Plane 3"** = **"BP3"** in [Robens, Stefaniak, Wittbrodt, arXiv:1908.08554].

[**AP**, Robens, Tetlalmatzi-Xolocotzi, arXiv:2101.00037]

hhh in the TRSM "BP3" [14 TeV] H (300*,* 475) 26*.*5 15 20 352 500

- Search for **hhh** via: $pp \rightarrow (b\bar{b})(b\bar{b})(b\bar{b}) \rightarrow$ **6** b-jets. \bullet Jean en Fernard Via. μ ² (*UU*)(*UU*) - \bullet U⁻Jeis.
- About **20%** of the **hhh** final state! Δ Δ Δ out Δ Ω θ shown in table Δ Δ shown in table 2.5 shown in t The cuts not shown above are common for all points, as follows: *|*⌘*|^b <* 2*.*35, *m*min*,* med*,* max *<*
- Significances **large**, even when including systematic uncert.: ¹) *<* 3*.*5 and *Rbb*(*h*1) *<* 3*.*5. For some of the • Significances large, even when including systematic

¹) *>* [50*,* 50*,* 0] GeV, *R*(*hⁱ* 1*, h^j* n including systematic uncert.:

$\mathbf{2D2U}$ 17 \mathbf{A} $\mathbf{T_0}$ \mathbf{V} \mathbf{I} H (300*,* 475) 26*.*5 15 20 352 500 $\overline{1}\times\overline{1}\times\overline{1}$

J (280*,* 500) 34*.*0 10 40 454 525

hhh in the TRSM "BP3'' [14 TeV]

- **hhh** will (**probably?**) **not be a discovery channel,**
-

• but could be **important in determining the parameters of the model**, if scalars are discovered!

- Get the **MG5_aMC model** at: https://gitlab.com/apapaefs/multihiggs_loop_sm.
- [A patch to MG5_aMC to enable **Loop** \times Tree is included].
- Can generate events either at:
	- **SM**^2 + interference of $[M \times One-Insection diagrams]$, i.e.: |ℳ| 2 $= |\mathcal{M}_{\rm SM}|$ 2 + 2Re{ $\mathscr{M}_{\rm SI}^{*}$ $*_{\text{SM}}$ $M_{1-\text{ins.}} \} \propto 1 + c_i$

Monte Carlo Implementation of Anomalous Couplings

or

• **SM^2** + interference of $[M \times One]$ **Two** insertion diagrams + [One Insertion]^2, i.e.: |ℳ| 2 $= |\mathcal{M}_{\rm SM}|$ 2 + 2Re{ $\mathscr{M}_{\rm SI}^{*}$ $\mathcal{L}_{\text{SM}}^*$ M₁-ins.} + 2Re{ $\mathcal{M}_{\text{SM}}^*$ M_{2-ins.}} + | M_{1-ins.}] α 1 + c_i + c_jc_k + c_{ℓ}^2

2

