Double Higgs Boson Production in Association of Jets in the Triplet Higgs Model at the LHC

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Extended Scalars form All Angles

CERN, Oct. 22, 2024



Single and Double Higgs Boson Production in Association of Jets using Herwig 7 in the SM and in BSM

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References and Thanks to Collaborators...

Collaborators: Simon Platzer, Peter Schichtel, Michael Rauch, Malin Sjodahl, Francisco Campanario, Tinghua Chen, and Bathiya Samarakoon.

- https://arxiv.org/abs/2109.03730
- https://arxiv.org/abs/1802.09955
- https://arxiv.org/abs/1610.07922
- https://arxiv.org/abs/1308.2932
- https://arxiv.org/abs/0710.5621



Eur. Phys. J. C (2022) 82: 704 https://doi.org/10.1140/epjc/s10052-022-10671-9 Special Article - Tools for Experiment and Theory

NLO multijet merging for Higgs production beyond the VBF approximation

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In a full talk I would discuss the history of calculations.



More References and Collaborators

- <u>Anomalous Higgs boson couplings in vector boson fusion</u> (Ph.D. Thesis 2023, Tinghua Chen). Related paper in progress.
- <u>Azimuthal angle correlations for Higgs boson plus multi-jet</u> <u>events</u> (J. Andersen, K. Arnold, D. Zeppenfeld)
- In Progress: Double Higgs via VBF in BSM scenarios.
- Collaborators: Tinghua Chen and Bathiya Samarakoon



WHAT IS VBF?

The second most copious Higgs production mode after gluon-gluon fusion











Vector Boson Fusion (+Jet)









Simulation Tools and Matrix Elements

- Herwig 7 Event Generator (<u>https://herwig.hepforge.org</u>)
 - General purpose event generator,
 - parton showers, hadronization, MPI modeling
- HJETS++ (<u>https://hjets.hepforge.org</u>)
 - Add-on module to Herwig 7
- VBFNLO (<u>https://www.itp.kit.edu/vbfnlo</u>)
 - Interfaced to Herwig 7 via the Binoth One-Loop Accord





F. Campario, T. M. Figy, S. Platzer, and M. Sjodahl, PRL 111, 211802

- Matchbox [S. Platzer and S. Gieseke, arXiv:1109.6256]
 - Catani-Seymour Dipole subtraction [hep-ph/9605323]
 - Subtractive and POWHEG style matching to parton shower
 - ColorFull [M. Sjodahl, arXiv:1211.2099, http://colorfull.hepforge.org]
- Tensorial Reduction [F. Capanario, arXiv:1105.0920]
- Scalar Loop Integrals: OneLOop [A. van Hameren arXiv:1007.4716]









box lines





pen tagons











Real Corrections





VBFNLO



Project Description

VBFNLO is a fully flexible parton level Monte Carlo program for the simulation of vector boson fusion, double and triple vector boson production in hadronic collisions at next to leading order in the strong coupling constant. VBFNLO includes Higgs and vector boson decays with full spin correlations and all off-shell effects. In addition, VBFNLO implements CP-even and CP-odd Higgs boson via gluon fusion, associated with two jets, at the leading order one loop level with the full top-quark and bottom-quark mass dependence in a generic two Higgs doublet model.

A variety of effects arising from beyond the Standard Model physics are implemented for selected processes. This includes anomalous couplings of Higgs and vector bosons and a Warped Higgsless extra dimension model. The program offers the possibility to generate Les Houches Accord event files for all processes available at leading order.

All implemented processes can be found here.

The list of people involved in VBFNLO is here.



Herwig 7

Bootstrap script pulling in all dependencies. Tested on a large number of platforms.

ightarrow ./herwig-bootstrap /where/to/install

Documentation re-written from scratch: "living" sphinx sites replacing old wiki pages.

 \rightarrow Check out <code>herwig.hepforge.org</code>

Update of detailed physics & manual will follow in due course.

Usage can be done as before, though lots of parallelization added:

- Separate building, grid adaption, and event generation
 - \rightarrow Cheaper parameter variations.
- Grid adaption parallelized in separate jobs (no IPC required)
 - \rightarrow Herwig build --maxjobs=6 LHC-Matchbox.in
 - \rightarrow Herwig integrate --jobid=3 LHC-Matchbox.run ...
- Multicore capabilities
 - \rightarrow Herwig run --jobs=24 LHC-Matchbox.run





.

cd /Herwig/EventHandlers set EventHandler:LuminosityFunction:Energy 13000*GeV set EventHandler:Weighted On

Note that event generation may fail if no matching matrix element has ## been found. Coupling orders are with respect to the Born process, ## i.e. NLO QCD does not require an additional power of alphase.

Model assumptions
read Matchbox/StandardModelLike.in
read Matchbox/DiagonalCKM.in

Set the order of the couplings
cd /Herwig/MatrixElements/Matchbox
set Factory:OrderInAlphaEU 3
#set Factory:AlphaParameter 1.0
Select the process
You may use identifiers such as p, pba, j, l, mu+, h0 etc.
do Factory:Process p -> h0 j j j

The next line can switch of hadronization # # and MPI modelling. Use with care!! read Matchbox/PQCDLevel.in

Special settings required for on-shell production of unstable particles
enable for on-shell top production

h+3 pt3





insert JetCuts:JetRegions 0 FirstJet insert JetCuts:JetRegions 1 SecondJet insert JetCuts:JetRegions 2 ThirdJet

#read Matchbox/MCatNLO-DefaultShower.in # read Matchbox/Powheg-DefaultShower.in ## use for strict LO/NLO comparisons # read Matchbox/MCatLO-DefaultShower.in ## use for improved LO showering # read Matchbox/LO-DefaultShower.in

read Matchbox/MCatNLO-DipoleShower.in
read Matchbox/Powheg-DipoleShower.in
use for strict LO/NLO comparisons
read Matchbox/MCatLO-DipoleShower.in
use for improved LO showering
#read Matchbox/LO-DipoleShower.in

#read Matchbox/NLO-NoShower.in
#read Matchbox/LO-NoShower.in

read Matchbox/MuDown.in
read Matchbox/MuUp.in

Various setup:

- 1. Fixed Order
- 2. NLO/LO matching
 - 1. Subtractive
 - 2. Powheg





NLO Matching in Herwig 7

Parton showers vs fixed-order calculations



 $\textbf{Fixed-order calculations} \rightarrow \textbf{hard jets}$

- $\bullet\,$ reliable at high scales if no large scale hierarchies are present
- accurate predictions for limited number of legs (+ loops)
- determines **perturbative accuracy** (LO, NLO, NNLO, ...)
- $\textbf{Showers} \rightarrow \mathsf{jet} \ \mathsf{substructure}$
 - reliable in soft/collinear regions if large scale hierarchies are present
 - approximate predictions for many particles
 - determines logarithmic accuracy (LL, NLL, NNLL, ...)
- \Rightarrow largely complementary, so ideally combine them!

See

https://arxiv.org/pdf/1605.07851.pdf https://arxiv.org/pdf/1912.06509.pdf https://arxiv.org/pdf/1109.6256.pdf







NLO Merging in Herwig 7

See https://arxiv.org/pdf/1705.06700.pdf



Dijet Production at the LHC



Definitions

• The rapidity y is defined as

$$y = \frac{1}{2} \ln \left(\frac{E + p_z}{E - p_z} \right),$$

• The separation of rapidity between two jets is

$$\Delta y_{j_1 j_2} = |y_{j_1} - y_{j_2}|,$$

• The transverse momentum, p_T is defined by

$$p_T = \sqrt{p_x^2 + p_y^2},$$

• invariant mass for two jets, defined as

$$m_{j_1j_2} = \sqrt{(p_{j_1} + p_{j_2})^2},$$





Input Parameters and Event Selection

- CMS Energy of Collider and Beam Type: pp LHC (13 TeV).
- PDF set (LHAPDF): PDF4LHC15_nnlo_100_pdfas.
- Factorization and renormalization scales set to HT(jets) scale.
- Anti- k_T Jet clustering (R=0.4) using FastJet, at least two jets with $p_{T,jets} > 25$ GeV.
- Matched results use a MC@NLO type matching. No hadronization or MPI have been included in the simulations.
- The LOOSE selection cuts is defined as

$$m_{j_1 j_2} > 200 \text{ GeV}, \quad \Delta y_{j_1 j_2} > 1.$$
 (34)

• The TIGHT selection cuts is defined as

$$m_{j_1j_2} > 600 \,\,{
m GeV}, \quad \Delta y_{j_1j_2} > 4.5, \quad y_{j_1} \cdot y_{j_2} < 0.$$
 (35)



Rapidity Gap and Invariant Mass



NLO+PS does a good job overall for 2 jet observables.

8 9

 $\Delta y_{j_1 j_2}$



Inclusive Jet and Gap Jet Cross Section





Using HT(jets) for the scale.

Gap jets reside between yj1 and yj2.

Beyond 2 jets you need NLO 3 jet.



Gap Jet Activity





Gap Jets are quite soft.

Again, you need NLO for more than 2 jets



Transverse Momentum Spectrum (fixed scale: mh)



VBFNLO vs HJETS – Tight Selection Cuts

Figure: The distribution of invariant mass (left) and rapidity gap (right) of the two tagging jets with TIGHT cuts.

Anomalous Higgs Couplings

The effective Lagrangian can be written as the following equation:

$$\mathcal{L} = \frac{g_{5e}^{HWW}}{\Lambda_{5e}} HW^{+}_{\mu\nu}W^{-\mu\nu} + \frac{g_{5o}^{HWW}}{\Lambda_{5o}} H\tilde{W}^{+}_{\mu\nu}W^{-\mu\nu} + \frac{g_{5e}^{HZZ}}{2\Lambda_{5e}} HZ_{\mu\nu}Z^{\mu\nu} + \frac{g_{5o}^{HZZ}}{2\Lambda_{5o}} H\tilde{Z}_{\mu\nu}Z^{\mu\nu} .$$
(36)

The general tensor structure of the HVV vertex in the massless quark limit is

$$T^{\mu\nu}(q_1, q_2) = a_1(q_1, q_2)g^{\mu\nu} + a_2(q_1, q_2)[q_1 \cdot q_2g^{\mu\nu} - q_2^{\mu}q_1^{\nu}] + a_3(q_1, q_2)\epsilon^{\mu\nu\rho\sigma}q_{1\rho}, q_{2\sigma}$$
(37)

if $a_2 = a_3 = 0$, then a_1 represents the SM case.

Anomalous Higgs Couplings

The form factors a_2 and a_3 can be derived from the effective Lagrangian of Eq. (36),

$$a_2(q_1, q_2) = -\frac{2}{\Lambda_{5e}} g_{5e}^{HWW}$$
, $a_3(q_1, q_2) = \frac{2}{\Lambda_{5o}} g_{5o}^{HWW}$ (38)

for the HWW vertex, and

$$a_2(q_1, q_2) = -\frac{2}{\Lambda_{5e}} g_{5e}^{HZZ}$$
, $a_3(q_1, q_2) = \frac{2}{\Lambda_{5o}} g_{5o}^{HZZ}$ (39)

Anomalous Higgs Couplings

pure CP-even: g^{HZZ}_{5e} = g^{HZZ}_{5o} = 0.5,
pure CP-odd: g^{HWW}_{5o} = g^{HWW}_{5o} = 0.5,
CP-mixed: g^{HWW}_{5e} = g^{HWW}_{5o} = g^{HZZ}_{5o} = g^{HZZ}_{5o} = 0.5.

CP Sensitive Observables

 $|\Delta \phi_{j_1 j_2}|$ is defined as

$$|\Delta \phi_{j_1 j_2}| = |\phi_{j_1} - \phi_{j_2}|$$

 $\Delta \phi_{j_f j_b}$ is define as

$$\Delta \phi_{j_f j_b} = \begin{cases} \phi_{j_1} - \phi_{j_2}, & y_{j_1} > y_{j_2} \\ \phi_{j_2} - \phi_{j_1}, & y_{j_1} < y_{j_2} \end{cases}$$

 $\Delta \phi_{j_f j_b}$ also satisfies

$$\Delta \phi_{j_f j_b} = \begin{cases} \Delta \phi_{j_f j_b} + 2\pi, & \Delta \phi_{j_f j_b} < -\pi \\ \Delta \phi_{j_f j_b} - 2\pi, & \Delta \phi_{j_f j_b} > \pi \end{cases}$$

 ϕ_2 is defined as

$$\phi_2 = \angle (\mathbf{q}_{a\perp}, \mathbf{q}_{b\perp}),$$

where

$$q_a = \sum_{j \in \{ j \in s: y_j < y_h \}} p_j, \quad q_b = \sum_{j \in \{ j \in s: y_j > y_h \}} p_j,$$

Anomalous Higgs in VBFNLO

- Anomalous Higgs Couplings were implemented into VBFNLO since 2004 for Higgs + 2 jets via VBF.
- In 2012 there was an interest in using the PowHEG implementation of Higgs + 2 jets via VBF and a reweighting technique to generate predictions include parton shower and hadronization effects.
- Today you can use VBFNLO with Herwig 7 to simulate Higgs Production with Anom Couplings (See Release Note—VBFNLO 3.0 for details, <u>https://arxiv.org/abs/2405.06990</u>)

H+2 jets at Fixed Order (2004, arXiv: 0403297v1)

Results using VBFNLO interfaced to Herwig 7

Azimuthal Angle Correlations

Figure: The distribution for of $|\Delta \phi_{j_1 j_2}|$ (left) and $\Delta \phi_{j_f j_b}$ (right) for matched setup $h(2^*) \bigoplus PS$ using TIGHT selection cuts.

Azimuthal Angle Correlations

Figure: The distribution for of $|\Delta \phi_{j_1 j_2}|$ (left) and $\Delta \phi_{j_f j_b}$ (right) for merged setup h(2^{*}, 3^{*}, 4) using TIGHT selection cuts.

Azimuthal Angle Correlations

Figure: Distributions of ϕ_2 for $h(2^*) \bigoplus PS$ (left) and $h(2^*, 3^*, 4)$ (right) using TIGHT cuts.

Invariant Mass of Di-Tag System

Figure: The distribution of $m_{j_1j_2}$ with INCL cuts (left) and LOOSE cuts (right) for CP-mixed anomalous couplings. The ratio plot compared $h(2^*, 3^*, 4)$, $h(2^*, 3)$, $h(2^*) \bigoplus PS$, and $h(2^*)$ fixed order, with the merged setup $h(2^*, 3^*, 4)$ as the reference.

Rapidity Gap

Figure: The distribution of $\Delta y_{j_1j_2}$ with INCL cuts (left) and TIGHT cuts (right) for CP-mixed anomalous couplings. The ratio plot compared $h(2^*, 3^*, 4)$, $h(2^*, 3)$, $h(2^*) \bigoplus PS$, and $h(2^*)$ fixed order, with the merged setup $h(2^*, 3^*, 4)$ as the reference.

Rapidity Gap and Rapidity

Figure: The distribution of $\Delta y_{j_1j_2}$ (left) and y_{j_1} (right) for $h(2^*) \bigoplus PS$ with INCL cuts.

Rapidity Gap and Rapidity

Figure: The distribution of $\Delta y_{j_1j_2}$ (left) and y_{j_1} (right) for $h(2^*, 3^*, 4)$ with INCL cuts.

Transverse Momentum

Figure: The distribution of P_{T,j_1} for $h(2^*) \bigoplus PS$ (left) and $h(2^*, 3^*, 4)$ (right) with INCL cuts.

Higgs Triplet Model (Type II Seesaw Model)

arXiv:1105.1925

• The motivation for the Higgs Triplet Model stems from the observation that two doublets can be decomposed into a triplet and a singlet representation $(2 \otimes 2 = 3 \oplus 1)$.

$$\Delta = \begin{pmatrix} \delta^+ / \sqrt{2} & \delta^{++} \\ \delta^0 & -\delta^+ / \sqrt{2} \end{pmatrix}$$

 $V(\Phi, \Delta) = -\mu^2 H^{\dagger} H + \frac{\lambda}{4} \left(H^{\dagger} H \right)^2 + M_{\Delta}^2 Tr(\Delta^{\dagger} \Delta) + \left[\mu \left(H^T i \sigma \Delta^{\dagger} H \right) + h.c \right] + \lambda_1 \left(H^{\dagger} H \right) Tr(\Delta^{\dagger} \Delta) + \lambda_2 \left(Tr(\Delta^{\dagger} \Delta) \right)^2 + \lambda_3 \left(Tr(\Delta^{\dagger} \Delta)^2 \right) + \lambda_4 H^{\dagger} \Delta \Delta^{\dagger} H$

• There are 7 physical Higgs states: h^0 , H^0 , A^0 , H^{\pm} , $H^{\pm\pm}$ }

Double Higgs Production via VBF

Double Higgs Production via VBF

W W → h h

 $Z Z \rightarrow h h$

T2 P2 N3

T3 P1 N0

T4 P1 N9

Double Higgs via VBF Matched to Parton Showers (Type II Seesaw Model)

Auxiliary Slides

Parameter Space

• We can establish a parameter space as follows:

 $P = \{M_{h^0} = 125 \ GeV, M_{H^0}, M_{A^0}, M_{H^{\pm}}, M_{H^{\pm\pm}}v_{\Delta}, v, \cos(\alpha)\}$

ν	246.2 <i>GeV</i>
v_{Δ}	5.78 GeV
cosα	0.9
M_{h^0}	125.000 GeV
М _Н о	758.695 GeV
M_A o	758.695 GeV
$M_{H^{\pm}}$	570.63 GeV
$M_{H^{\pm\pm}}$	275.0 <i>GeV</i>

Work on Muon Colliders

Double Higgs boson production via photon fusion at muon colliders within the triplet Higgs model

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Published in: *Phys.Rev.D* 109 (2024) 7, 075015 e-Print: <u>2312.12594</u> [hep-ph] DOI: <u>10.1103/PhysRevD.109.075015</u> (publication)

