Golden period of particle physics



Golden period of particle physics

- 4π acceptance
- Trace a charged particle with good efficiency
- Calculate momentum
- Identify particle, e/π/μ/p/n
- Separate them out, spatial resolution (~8µm) with particle multiplicity up to 150
- Mass identification, identify mother
- Secondary vertex



Reconstruction of particles in present day HEP detector



Obviously large bias on examples to BELLE/CMS/INO detector

Measurement of track momentum

$$P_{T}[GeV/c] = 0.3 B[T] R[m]$$

$$\sin \frac{\theta}{2} = \frac{L}{2r}; \text{ for } r \gg L \quad \frac{\theta}{2} \approx \frac{L}{2r} \Rightarrow \theta \approx \frac{0.3BL}{P_{T}}$$

$$s = r - r \cos\left(\frac{\theta}{2}\right) \approx r \left[1 - \left(1 - \frac{1}{2}\frac{\theta^{2}}{4}\right)\right] = \frac{r\theta^{2}}{8} \approx \frac{0.3BL^{2}}{8P_{T}} = \frac{L^{2}}{8R}$$

$$s = x_{B} - \frac{x_{A} + x_{C}}{2}; \quad \therefore ds = dx_{B} - \frac{dx_{A} + dx_{C}}{2}; \quad (ds)^{2} = \sigma_{s}^{2} = (3/2)\sigma_{s}^{2}$$

$$\frac{dP_{T}}{P_{T}} = \frac{\sigma_{s}}{s} = \frac{\sqrt{3/2}}{s} \sigma_{x} = \sqrt{\frac{3}{2}} \sigma_{x} \frac{8P_{T}}{0.3BL^{2}}$$



For multiple coulomb scattering, average scattering angle

$$\theta_0 \approx \frac{13.6 MeV}{\beta Pc} Z \sqrt{\frac{L}{X_0}} \Rightarrow s_{ms} = \frac{L\theta_0}{4\sqrt{3}} \Rightarrow \frac{s_{ms}}{s} = \frac{dP}{P} \bigg|_{ms} \approx 0.05 \frac{1}{\beta B \sqrt{LX_0}}$$

Momentum resolution, σ_P/P varies

- Constant for multiple coulomb scattering (neglect $1/\beta$ term for $p > \sim GeV$)
- linearly with P_T
- **Inversely with** *BL*²

For higher momentum, detector size increases with \sqrt{L}

Interactions of charged particle and Track reconstruction

Туре	particles	fund. parameter	characteristics	effect
Multiple Scattering	all charged particle	radiation length X	almost gaussian average effect 0 depends $\sim 1/p$	deflects particles, increases measurement uncertainty
Ionisation loss	all charged particle	effective density $A/Z * ho$	small effect in tracker, small dependence on p	increases momentum uncertainty
Bremsstrahlung	all charged particle, dominant for e	radiation length X	highly non- gaussian, depends	introduces measurement bias
Hadronic Int.	all hadronic particles	nuclear interaction length Λ	destroys particle, rather constant effect in p	main source of track reconstruction inefficiency

Cylindrical Drift Chamber for collider experiment

Characteristics:

- Cylindrical symmetry
- Open drift cell geometry
- Require: Simple space-time relation
- given by E,B field and drift cell geometry





Single wire resolution ~ 50 μm (300μm) Double track resolution ~ 450 μm (7mm)







Introduction to track fitting

Tracking detector with cylindrical layers Input to track finding is all or parts of the measurements in the detector at a given instance

A successful track finder identifies a set of potential tracks as indicated in the figure Hits along these tracks are given to the track fitter for parameter estimation and final validation of track candidate



After the track fit one usually forgets about the hits and only cares about a compact representation of the tracks

Track finder

• Track fitting is involved with inversion of matrices and large dimension

- Large computing time to look all combinations of points
- Most of those combinations might not converse
- First identify combinations without CPU consuming fitting algorithm, which could be a simple pattern recognition
 - \rightarrow Track finding algorithm
- Track fitting algorithm : used hit from finder
 - Bubble chamber days
 - scanning team looked at photograph
 - recognition straight forward
 - Electronically read out detectors
 - less hits per track length
 - environment got more dense (more hits)
 - algorithms needs to replace 'look at'
 - Algorithm (time consuming)

Data Analysis Technique for High Energy Physics Ed : M. Regler and R. Frühwirth





Track finding algorithm

- Prerequisite:
 - Detector alignment
 - Hit reconstruction (space points with uncertainties)
- Depends on :
 - Tracker type (detector resolution, dominated in drift chamber, where contribution from electron statistics, diffusion, inhomogeneity of the drift field, electronic jitter, TDC binning, wire sagging....)
 - Accuracy in alignment (Shift of different layers)
 - Hits it produces (Inefficiency, # of hit in a layer, electronic noise)
 - Magnetic field (inhomogeneous, how much can it be linearised)
 - Event environment (Beam background, pileup, cosmic)
 - Kink for decay in flight



Some more concern of track finding algorithm



- Multiple scattering on sensitive material, support structure, beam pipe
 - Not only fluctuation, but also for dE/dx estimation

$$\theta_e = \frac{13.6MeV}{\beta pv} Z_1 \sqrt{\frac{L}{L_R}} \left[1 + 0.125 \log_{10} \left(\frac{L}{L_R} \right) \right]$$



Track finder : Track following

Outside in

- Start with seed at outer end of tracking volume
- Swim in general direction of the beamline
- Low occupancy outside, easy pattern reco, add hits moving in one knows already where to look
- For high multiplicity events, look for out-to-in.
- Certainly it can not find very low p_T tracks
- Inside out
 - Follow natural particle direction, least MS
 - -Detector cutoff in pseudorapidity has minimal effect
 - Track are built from inside to out because many particles will interact before crossing all layers of the tracker
 - Seeding difficult because of high occupancy



Track finder : Road Search

• First, last and a point at centre and look for other points (or points in inner and outer layers in collider along with interaction vertex as third point (computing time is prop to a factor between n^2 to n^3 , where *n* is the number of hit points. Slower than following methods, but suitable where redundancy in the co-ordinate measurement is very low.



- CMS uses pixel tracklet and outer TOB/TEC hit and build the road with a minimum $P_{\rm T}$
 - The standard seeding uses hits from the inner two (matched) layers of the TIP and TID for inner cood hits, and hits in the ou for outer seed hits.
 - The trajectory will k precision coordinate of
 - Initially combine laye layer
 - This trajectory is buil The smoothed trajector the beamline and the t again, this time addin higher occupancy layer



Track Reconstruction in collider (CMS)

• Clustering of silicon strips and pixels to find "hit" positions and errors

• Initial estimate of track parameters using a minimal number of hit

• Collection of the remaining hits associated to the particle trajectory

• Final estimation of the track parameters using the full set of associated hits

• Removal of tracks likely to be fakes



Steps of Track Reconstruction at collider (Track seeding)

- Default seeding is done in pixels layers
- Track are built from inside to out because many particles will interact before crossing all layers of the tracker
- Pixel hits are very precise. So two or three pixel hits can already give good initial estimation of the track parameters
 - single pixel area = $100 \mu m \times 150 \mu m = 1.5 \times 10^{-2} mm^2$
 - typical strip area = $10 \text{cm} \times 100 \mu \text{m} = 10 \text{mm}^2$

• This plot explains why CMS tracking is heavily pixel driven (seeding) : Low



Track follower at collider experiments



Hits in a subset of tracker layers are used to find trajectory seeds. Default seeds are made from pixels, the innermost layers of the tracker.



Seeds that point well outside the collision region are discarded



Triplets of hits (or pairs plus the beamspot) are combined to produce trajectory seeds whose directions are compatible with the beam collision region.



For the remaining seeds, each one is then propagated outward to collect more hits to find the full trajectory of the charged particle.

















- In the second level remove all hits belongs to these tracks
- Restart with remaining hits with loose criteria

Track finder

- Linear approximation, e.g., INO :
 - Triplet : Formation (allowed maximum gap of two layers), join them in a chain, sort out the best choice as a track candidate.
 - Need optimisation mainly for a gap



Track finder

- Linear approximation, e.g., INO :
 - Bending at the end of the track can not be included with this algorithm, but in the fitting time that can be taken care of.





Hough Transformation



Simulated signal in Tracker : (0,0,3.8T) field





pidin, momin, thein, phiin, posxin, posyin, poszin :// pdgId, three momentum vector and position of incident particle

- Radius[13] = {2.9, 6.8, 10.9, 16.0, 22, 29, 37, 45, 52, 59, 66, 73, 80}; // in cm //4 bit
- unsigned int detidTk = theTouchable->GetCopyNumber(0); //Layer number
- detidTk<<=13; //For layer number
- int nIThe = int((750.0 + localpos.z()) /0.2); //200micon strip length //13bit
- detidTk +=nIThe;
- detidTk<<=15;
- int nIPhi = int((acos(-1.0)/4. + localpos.phi()) *20000); // 50 micon strip width : 15bit
- detidTk +=nIPhi;
- simenrTk = (*TrkCollection)[ij]->GetEdep(); //in KeV
- unsigned int nsimhtTk; //#of simhit in tracker

Hough Transformation

- Tracks in uniform magnetic field along Z direction
- Straight line in (ρ,z) plane



Hough Transformation

• y = mx + c: points in straight line (ρ_i, z_i) in $(cot\theta, \rho_0)$ plan coincide in one point, whereas noise are not



Divide inverse space into cells and find maxima

- But, useless for collider experiment, mainly because of
 - Much precise measurement of points in r- ϕ plane (curvature)
 - No estimation of P_T of track (particularly not in trigger)

Hough transformation in helical path (P_T criteria)



Single stub with co-ordinate (r, ϕ) maps onto a straight line in track parameter space (q/P_T, ϕ_0)

The Hough Transform track finding algorithm. Left: Trajectory of a single particle in a quarter of the tracker barrel in the x-y plane. The numbers label the stubs produced by the passage of the particles through the tracking stations. Middle: The produced stubs draw straight lines in the HT space, where the axes correspond to the track parameters $(q/p_T/\phi_T)$. Intersection point identifies a track candidate. Right: The same Hough-space discretised with an array of 32×64 cells. Stubs fill a cell only if their local-p_T estimation is compatible with the bin momentum. [95]

Example of Hough transformation



No firm comments, but

- Looks like Hough transformation is able to reconstruct p_T , ϕ and θ
 - Measurements in non bending plane is much better than bending in plane
- Uncertainties due to
 - Position resolution
 - Energy loss/Multiple scattering

Exercises

- Hough transformation method
 - -Using silicon signals in r-z plane
 - Make plots of few events in $\cot\theta$ and intersection plane
 - Reconstruct the polar angle of muon
 - Use the points in 2D grid points as well as fit those data points with straight line to obtain those parameters
 - Using silicon signals in r-ø plane
 - Make plots of few events in $\boldsymbol{\phi}$ and \boldsymbol{p}_T plane
 - reconstruct the azimuthal angle and transverse momentum of muons
 - -Make the resolution plots of momentum, polar and azimuthal angle of muon
 - Compare these results with the inclusion of 1% noise hits in the silicon sensor
- Use road search method to find the hits in the muon trajectory and compare that with Hough transformation method