# **PRODUCTION CROSS SECTIONS OF** $e^{\pm}$ **AND** $\gamma$ **RAYS**

## Luca Orusa XSCRC2024: Cross sections for Cosmic Rays CERN, 17/10/2024 Luca Orusa, Mattia Di Mauro, Fiorenza Donato, Michael Korsmeier

arXiv:2203.13143, Phys. Rev. D 105 (2022), 123021 arXiv:2302.01943, Phys. Rev. D 107 (2023), 083031 <u>https://github.com/lucaorusa/</u>



### The positron excess (Orusa et al. 2024 arXiv:2410.10951)



#### Secondary positrons

•Secondary contribution:  $q(T_{e^+}) = \sum_{i,j} 4\pi n_{\text{ISM},j} \times \int dT_i \phi_i(T_i) \frac{d\sigma_{ij}}{dT_{e^+}}(T_i, T_{e^+})$ 

•Scattering of a relativistic cosmic ray nucleus i colliding with an interstellar nucleus j at rest:

$$\sigma_{\rm inv}^{(ij)} = E_{\pi^{\pm}} \frac{d^3 \sigma_{ij}}{dp_{\pi^{\pm}}^3} (\sqrt{s}, p_T, x_R)$$

• $d\sigma(p + H \rightarrow e^{\pm} + X)$  former predictions affected by a factor 2 of uncertainty.



# Fit to $\pi^+$ data





- Total uncertainties between 5 and 10%.  $\sqrt{s} = 5-50$  GeV relevant.
- Integrating  $\sigma_{inv}$  over the solid angle and combining the result with the  $\pi^+$  decay, we obtain the  $\frac{d\sigma_{ij}}{dT_{e^+}}(T_i, T_{e^+})$  from  $\pi^+$ .

Experiment	$\sqrt{s}~[{ m GeV}]$		$\sigma_{ m inv}$	n	Ref.
NA49	17.3	$(\pi^{\pm},K^{\pm})$		-	[67, 76]
ALICE	900	$(\pi^+, K^\pm)$		-	[77]
CMS	900, 2760, 7000, 13000	$(\pi^{\pm},K^{\pm})$		-	[78, 79]
Antinucci	3.0,  3.5,  4.9,  5.0,  6.1,  6.8	$(\pi^{\pm})$	-	$\checkmark$	[80]
	2.8,  3.0, 3.2,  5.3,  6.1,  6.8	$(K^+)$	-	$\checkmark$	[80]
	4.9,5.0,6.1,6.8	$(K^{-})$	-	$\checkmark$	[80]
NA61/SHINE	6.3, 7.7, 8.8, 12.3, 17.3	$(\pi^{\pm}, K^{\pm})$	-		[68]

### Other channels

We consider the  $\pi^+$  created from weak decays of strange particles:

- $K^{\pm}$ ,  $K_S^0$ : fit on available data.
- $K_L^0$ : rescaled contribution from  $\overline{K_S^0}$ .
- $\overline{\Lambda}$ ,  $\Sigma$  and  $\Xi$ : rescaled contribution from the  $\Lambda$ .

We also consider the contribution from  $\pi^0$  to the  $e^+$  yield by multiplying the  $\pi^+$  cross sections by a normalization factor connected to the multiplicity of  $\pi^+, \pi^0$ .



## Results on the $e^+$ production cross section

- The  $\pi^+$  channel dominates the total cross section (10 times higher than the  $K^+$  channel).
- $e^+$  production from  $K_S^0$ ,  $K_L^0$ , and subdominant channels contributes at a few % level.
- $d\sigma/dT_{e^+}$  uncertainty: at  $1\sigma$  is 5% to 8% at all  $T_p$  energies.



• Secondary  $e^+$  are produced in nuclei interactions (p + A, A + p, and A + A).

• We use data from NA49 for  $p + \mathbf{C} 
ightarrow \pi^+ + \mathbf{X}$  and NA61 for the other channels. 6 17/10/2024

## Results on the $e^+$ production cross section

- The q(E) is predicted with a remarkably small uncertainty, ranging from 5% to 8%.
- The channels involving He, constitute 30-40% of the total spectrum.
- The heavier primary CNO nuclei contribute a non negligible few percent at the AMS-02 energies.



- Similar analysis performed for  $e^-$  with similar conclusions.
- Future measurements of pion production in the p + He could help to improve the predictions for nuclei channels.
- For Monte Carlo predictions, see talk by De la Torre Luque (tomorrow).

- •Most of the  $\gamma$  rays detected by Fermi-LAT are produced by the Galactic diffuse emission.
- •It originates from the interaction of CRs with interstellar gas and radiation fields within our own Galaxy.





### Diffuse $\gamma$ -ray emission





γ+γ

 $\rightarrow \gamma + \gamma$ 

 $\rightarrow \gamma + \gamma$ 

 $\rightarrow \gamma + \gamma$ 

γ + γ

Di Mauro et al., Phys. Rev. D 103, 123005 (2021)

 $\sigma_{\text{inv}} = \sigma_0(s) c_{20} [G_{\pi^+}(p_T, x_R) + G_{\pi^-}(p_T, x_R)] A(s)$ 

- $\pi^0$ : total uncertainties between 7 and 20%.  $\sqrt{s}$  = 5-50 GeV relevant.
- $K^{\pm}$ ,  $K_S^0$ : fit on available data.
- $K_L^0$ : rescaled contribution from  $\overline{K_S^0}$ .

•  $\eta$ : fit on available data and rescaled contribution from  $\pi^0$ .



#### Results on the $\gamma$ -ray production cross section



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### Diffuse $\nu$ emission (Orusa et al. in prep)

•Unlike the multi-component fluxes of  $\gamma$  rays, the flux of galactic diffuse  $\nu$  uniquely originate from the decay of charged mesons, that are produced in hadronic interactions.

#### • $\nu$ production cross section are needed.



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- 1.  $p + \text{He} \rightarrow e^+ + X (p + \text{He} \rightarrow \pi^+ + X)$ (p beam ~ 10 - 200 GeV)
- 2.  $p + p \rightarrow \gamma + X(p + p \rightarrow \pi^0 + X)$ (p beam ~ 10 - 200 GeV)
- 3.  $p + \text{He} \rightarrow \gamma + X (p + \text{He} \rightarrow \pi^0 + X)$ (p beam ~ 10 - 200 GeV)

$$\sigma_{\text{inv}} = \sigma_0(s) c_1 \left[ F_p(s, p_T, x_R) + F_r(p_T, x_R) \right] A(s)$$

$$F_p(s, p_T, x_R) = (1 - x_R)^{c_2} \exp(-c_3 x_R) p_T^{c_4} \times \exp\left[ -c_5 \sqrt{s/s_0}^{c_6} \left( \sqrt{p_T^2 + m_\pi^2} - m_\pi \right)^{c_7 \sqrt{s/s_0}^{c_6}} \right]$$

$$F_r(p_T, x_R) = (1 - x_R)^{c_8} \times \exp\left[ -c_9 p_T - \left( \frac{|p_T - c_{10}|}{c_{11}} \right)^{c_{12}} \right] \times \left[ c_{13} \exp(-c_{14} p_T^{c_{15}} x_R) + c_{16} \exp\left( - \left( \frac{|x_R - c_{17}|}{c_{18}} \right)^{c_{19}} \right) \right]$$

$$A(s) = \frac{1 + \left( \sqrt{s/c_{20}} \right)^{c_{21} - c_{22}}}{1 + \left( \sqrt{s_0/c_{20}} \right)^{c_{21} - c_{22}}} \left( \sqrt{\frac{s}{s_0}} \right)^{c_{22}}$$

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	$\pi^+$	$\pi^-$	$K^+$	$K^{-}$
$\chi^2_{ m NA49}/ m d.o.f.$	338/263	287/290.	146/151	197/151
$\chi^2_n/{ m d.o.f.}$	189/129	169/96	160/102	135/100
$\chi^2_{ m ALICE}$	77(33)	-	42(27)	36(27)
$\chi^2_{ m CMS}$	100 (88)	154 (88)	77~(68)	54~(68)
$\chi^2_{ m NA61,Antinucci}$	10(12)	15(12)	39(11)	44 (9)
$\chi^2_{ m tot}/ m d.o.f.$	527/392	456/386	306/253	332/251

- The uncertainties are about 5% for almost all  $T_{e^{+.}}$
- The relative uncertainty increases above 20% when approaching the maximum energy, that has a negligible impact on the final uncertainty.
- The results of this Section already hint at the final result. The by far dominant contribution of  $e^+$  production in p + p collisions comes from  $\pi^+$ .























