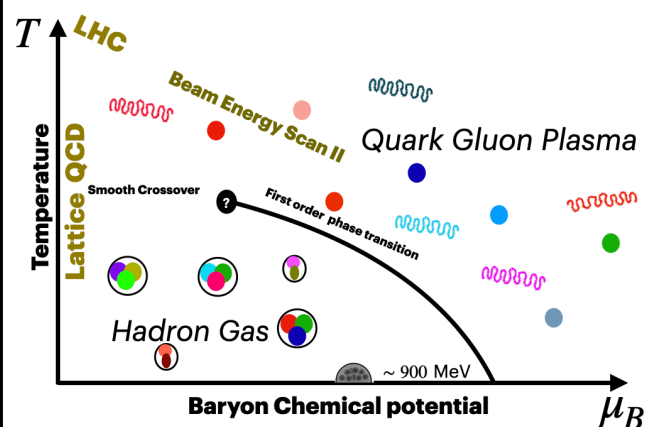


## Motivation

- Hydrodynamics simulations of heavy ion collisions require an **equation of state (EoS)** to comprehend experimental results.

**Goal: To provide an equation of state (EoS) for large  $\mu_B$  with a critical point, that matches lattice results.**

## Phase diagram



- At  $\mu_B = 0$  QCD phase diagram is well established.
- At finite  $\mu_B$  a critical point is expected but has not yet been observed.

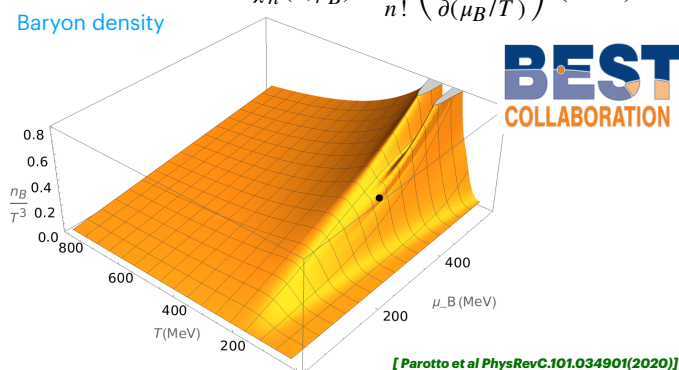
## Taylor Expansion

### Previous Approach:

- The EoS with a critical point from Taylor expansion is available and works well for low  $\mu_B$ .

$$n_B(T, \mu_B) = T^3 \sum_{n=1}^2 \frac{1}{(2n-1)!} \chi_{2n}^{\text{non-Ising}}(T) \left(\frac{\mu_B}{T}\right)^{2n-1} + \frac{T^4}{T} n_B^{\text{Ising}}(T, \mu_B)$$

$$\chi_n^B(T, \mu_B) = \frac{1}{n!} \left( \frac{\partial}{\partial (\mu_B/T)} \right)^n (P/T^4)$$



- For  $\mu_B > 450$  MeV, thermodynamics observables have **unphysical wiggles** due to limitations of truncated Taylor expansion, hindering critical phenomena studies

## T'-Expansion Scheme

- As a solution, the Wuppertal-Budapest collaboration developed a T'-expansion scheme that exhibits smooth behavior at high  $\mu_B$  and copes well with QCD transition temperature.

$$T \frac{\chi_1^B(T, \mu_B)}{\mu_B} = \chi_2^B(T', 0)$$

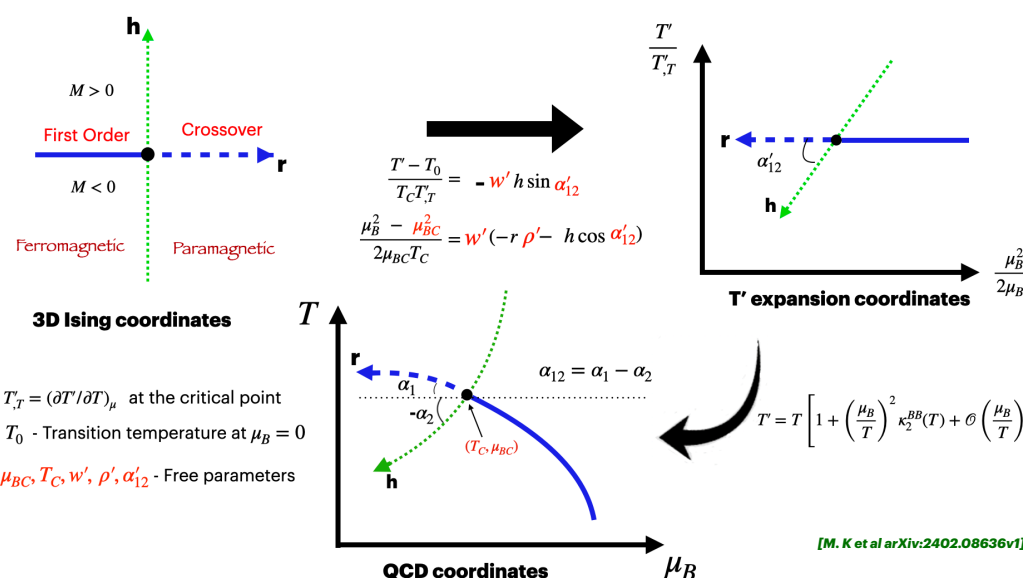
$$T'(T, \mu_B) = T \left[ 1 + \kappa_2^{BB}(T) \left(\frac{\mu_B}{T}\right)^2 + \kappa_4^{BB}(T) \left(\frac{\mu_B}{T}\right)^4 + \mathcal{O}\left(\frac{\mu_B}{T}\right)^6 \right]$$

[Borsányi, S et al PRL. 108(1), 101.034901(2021)]

- From the T'-expansion scheme, as long as  $T\chi_1^B/\mu_B$  is smooth, then finite density physics, such as the critical point can be encoded in T'.

## Mapping 3D-Ising to QCD

- If the critical point in QCD exists, then it must be in 3D-Ising model Universality class



- This mapping allows us to transport any physical observable from 3D-Ising to QCD with free parameters determined by the current physics knowledge.

## Merging 3D-Ising with T'-Expansion

### Methodology:

$$\frac{n_B(T, \mu_B)}{T^3} = \chi_1^B(T, \mu_B) = \left(\frac{\mu_B}{T}\right) \chi_{2,lat}^B(T', 0)$$

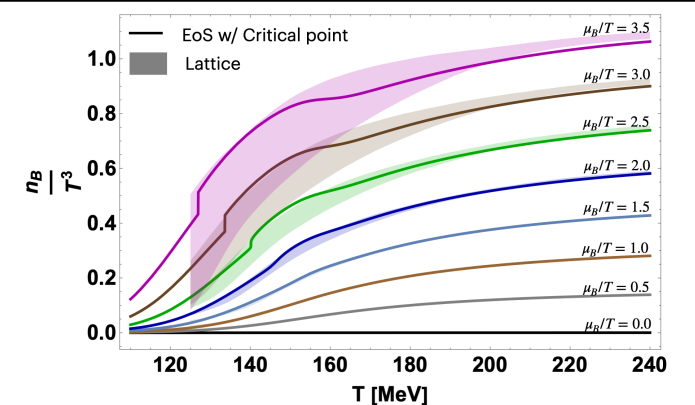
- We introduce the critical point in T' by separating into critical  $T'_{crit}$  and non-critical parts

$$T' = \underbrace{T'_{lat}(T, \mu_B)}_{\text{lower order in } (\mu_B/T)} + \underbrace{T'_{crit}(T, \mu_B) - Taylor[T'_{crit}(T, \mu_B)]}_{\text{higher orders in } (\mu_B/T)}$$

$$T'_{crit}(T, \mu_B) \approx \left( \frac{\partial \chi_{2,lat}^B(T, 0)}{\partial T} \Big|_{T=T_0} \right)^{-1} \frac{n_B^{crit}(T, \mu_B)/T^3}{(\mu_B/T)}$$

[M. K et al arXiv:2402.08636v1]

## Results

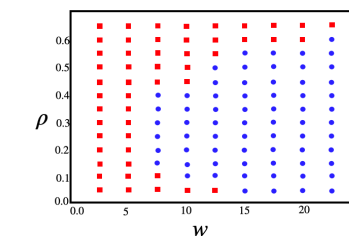


$\mu_{BC} = 350$  MeV,  $\alpha_{12} = 90$ ,  $w = 2$ ,  $\rho = 2$  [M. K et al arXiv:2402.08636v1]

- Our EoS with the critical point for some choice of parameters is within the error band of Extrapolated lattice QCD results.

## Constraints

- Transition line:** Choosing  $\mu_{BC}$  fixes  $T_C$  and  $\alpha_{12}$
- Physical quark masses:**  $\alpha_{12} = \alpha_1$
- Stability and causality:** fix  $w$  and  $\rho$

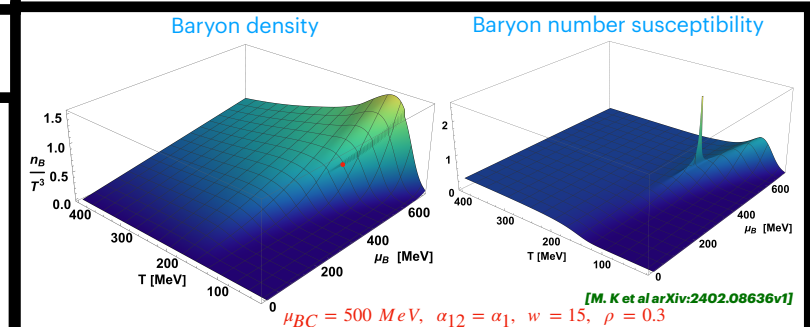


$\mu_{BC} = 500$  MeV,  $\alpha_{12} = \alpha_1$

- Blue point -> acceptable EoS
- Red point -> unacceptable EoS

[M. K et al arXiv:2402.08636v1]

## Thermodynamics



[M. K et al arXiv:2402.08636v1]

## Summary

### Disclaimer! : We don't predict the location of the critical point

- We offer a family of EoS with an enhance coverage ranging from  $\mu_B = 0$  to 700 MeV with adjustable parameters
- Our EoS can be tuned to reproduce physical quark masses