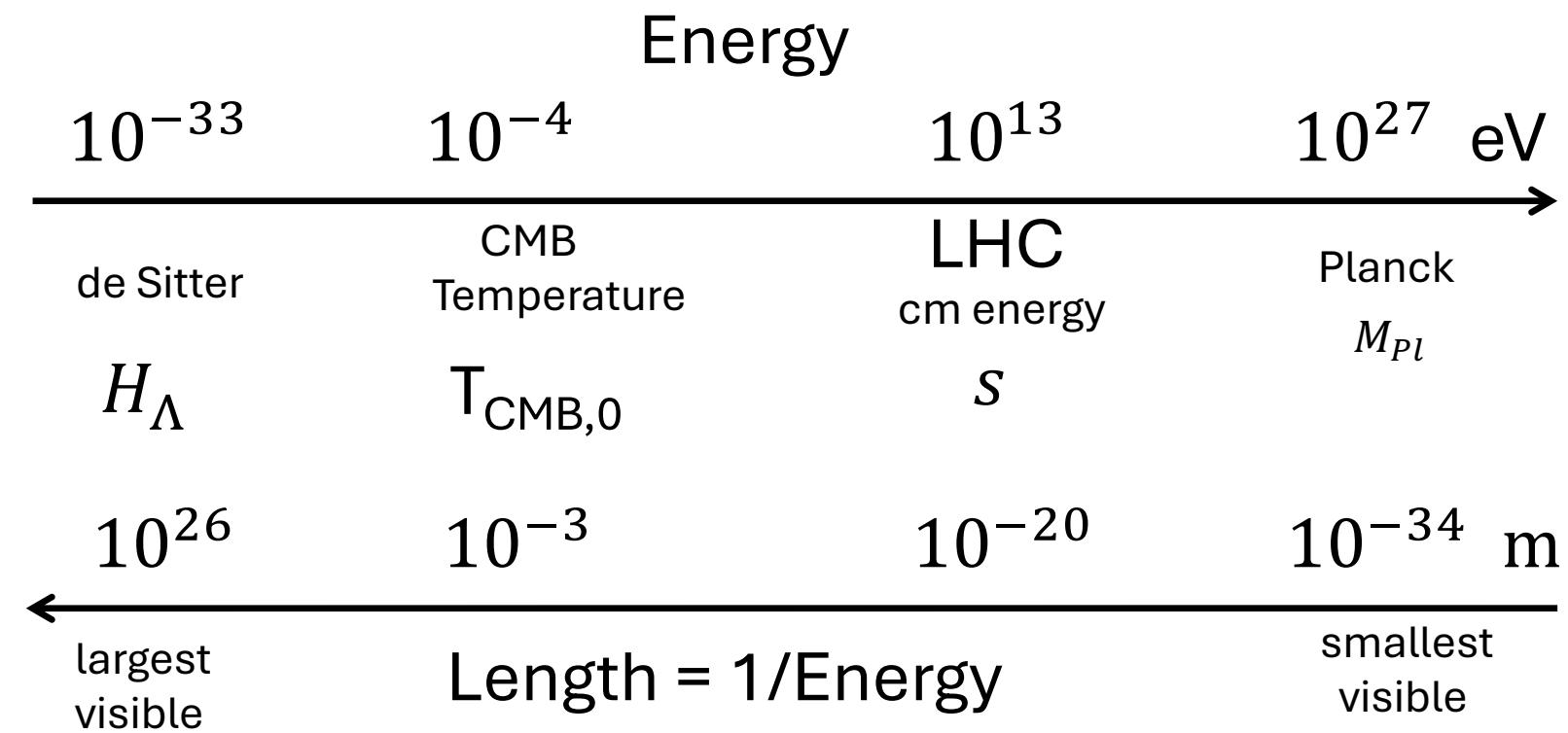


## Neil Turok Lecture 1

the universe is an amazing natural laboratory for fundamental physics



# FLRW metric: homogeneous and isotropic 3-spac

$$ds^2 = -dt^2 + a(t)^2 \gamma_{ij} dx^i dx^j; \quad a(t_0) \equiv 1$$

proper time

comoving coordinates: for a maximally symmetric space

Purely kinematical effects

$$R^{(3)} = 6\kappa \begin{cases} > 0 \text{ sphere} \\ = 0 \text{ torus} \\ < 0 \text{ hyperboloid} \end{cases}$$

1. Redshifting of wavelengths and momenta

e.g. photons

$$\lambda \propto a(t) \Rightarrow \frac{\lambda_{obs} - \lambda_{em}}{\lambda_{em}} = \frac{a_{obs}}{a_{em}} - 1 \equiv Z_{em} \quad \text{redshift} \quad a_{em} \propto 1/(1 + Z)$$

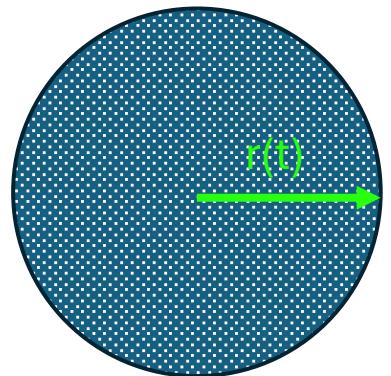
likewise the de Broglie wavelength and momentum of massive particles

2. Distances: objects which comove with expansion obey

$$d \propto a(t) \Rightarrow \frac{\dot{d}}{d} = \frac{\dot{a}}{a} = H(t): \text{recession velocity } \dot{d} \propto \text{distance } d \text{ Hubble's Law}$$

# Friedmann-Lemaître-Roberston-Walker Dynamics

Newtonian derivation



$$\frac{1}{2}\dot{r}^2 - \frac{GM}{r} = \frac{E}{m} = \text{const}; \quad M = \frac{4}{3}\pi r^3 \rho = \text{const}$$
$$\Rightarrow r \propto a(t); \quad \frac{\dot{a}^2}{a^2} = \frac{8\pi G}{3}\rho + \frac{\text{const}}{a^2}$$

Einstein's derivation:  $G_{\mu\nu} \equiv R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = 8\pi G T_{\mu\nu}$

where, for a perfect fluid at rest in comoving coords,  $T_{\mu}^{\nu} = \begin{pmatrix} -\rho & 0 & 0 & 0 \\ 0 & 0 & 0 & P(\rho)\delta_j^i \\ 0 & 0 & P(\rho)\delta_j^i & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$

↑  
Equation of state

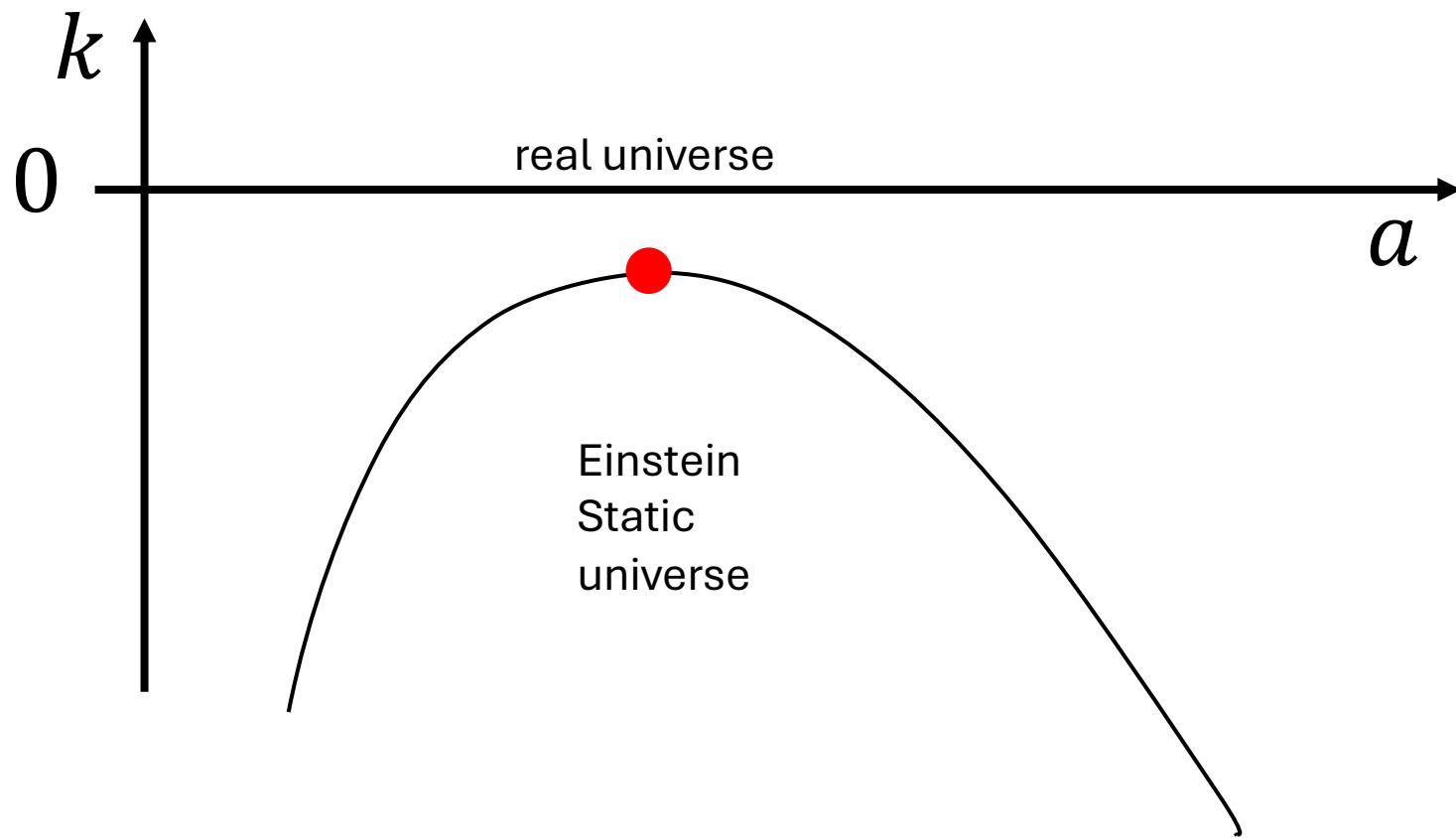
$G_{00}$  equation is a *constraint*:  $3\frac{\dot{a}^2}{a^2} - 3\frac{\kappa}{a^2} = 8\pi G\rho$  Friedmann  
 (involves only first time derivatives)

Conservation of energy-momentum  $\nabla_\mu T^{\mu\nu} = 0 \Rightarrow \dot{\rho} = -3\frac{\dot{a}}{a}(P + \rho)$   
 Types of energy:  $(dE = -P dV)$

- |                             |   |
|-----------------------------|---|
| 1. cosmological constant    | $T_{\mu\nu} = -\Lambda g_{\mu\nu} \Rightarrow P = -\rho$  |
| 2. matter (nonrelativistic) | $P = 0 \Rightarrow \rho \propto a^{-3}$ (approx.)   |
| 3. Radiation (relativistic) | $P = \frac{1}{3}\rho$ (or, more fundamentally, $T_\mu^\mu = 0$ )<br>$\Rightarrow \rho \propto a^{-4}$ <small>conformal symmetry</small><br><i>i.e.</i> , energy per photon $E_\gamma \propto \lambda^{-1} \propto a^{-1}$ |

So Friedmann reads  $\left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3}(\rho_\Lambda - \frac{\kappa}{a^2} + \rho_m + \rho_r) = H_0^2 \left( \Omega_\Lambda + \frac{\Omega_\kappa}{a^2} + \frac{\Omega_m}{a^3} + \frac{\Omega_r}{a^4} \right)$

i.e.,  $\dot{a}^2 - \lambda a^2 - ma^{-1} - ra^{-2} = k$ , with  $\lambda, m, r, k$  const



# FLRW metric: homogeneous and isotropic 3-spac

$$ds^2 = -dt^2 + a(t)^2 \gamma_{ij} dx^i dx^j; \quad a(t_0) \equiv 1$$

Friedmann equation

$$\left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3} \rho_{tot} - \frac{\kappa}{a^2}$$

$$= H_0^2 \left( \Omega_\Lambda + \frac{\Omega_\kappa}{a^2} + \frac{\Omega_m}{a^3} + \frac{\Omega_r}{a^4} \right)$$

Lambda      space      matter      radiation  
curvature

$$(\Omega_\Lambda + \Omega_\kappa + \Omega_m + \Omega_r \equiv 1)$$

$$R^{(3)} = 6\kappa \begin{cases} > 0 \text{ sphere} \\ = 0 \text{ torus} \\ < 0 \text{ hyperboloid} \end{cases}$$

Today,

$$H_0 \equiv 100h \text{ km s}^{-1} \text{ Mpc}^{-1}; h = 0.7 \pm 0.05; 1\text{pc}=3.26 \text{ light years}$$

Hubble time:  $H_0^{-1} = 9.78 \times 10^9 h^{-1} \text{ years}$

Hubble radius:  $H_0^{-1}c = 2998 h^{-1} \text{ Mpc}$

Critical density:  $\rho_c = \frac{3}{8\pi G} H_0^2 = 2.8 \times 10^{11} h^2 M_\odot \text{ Mpc}^{-3}$

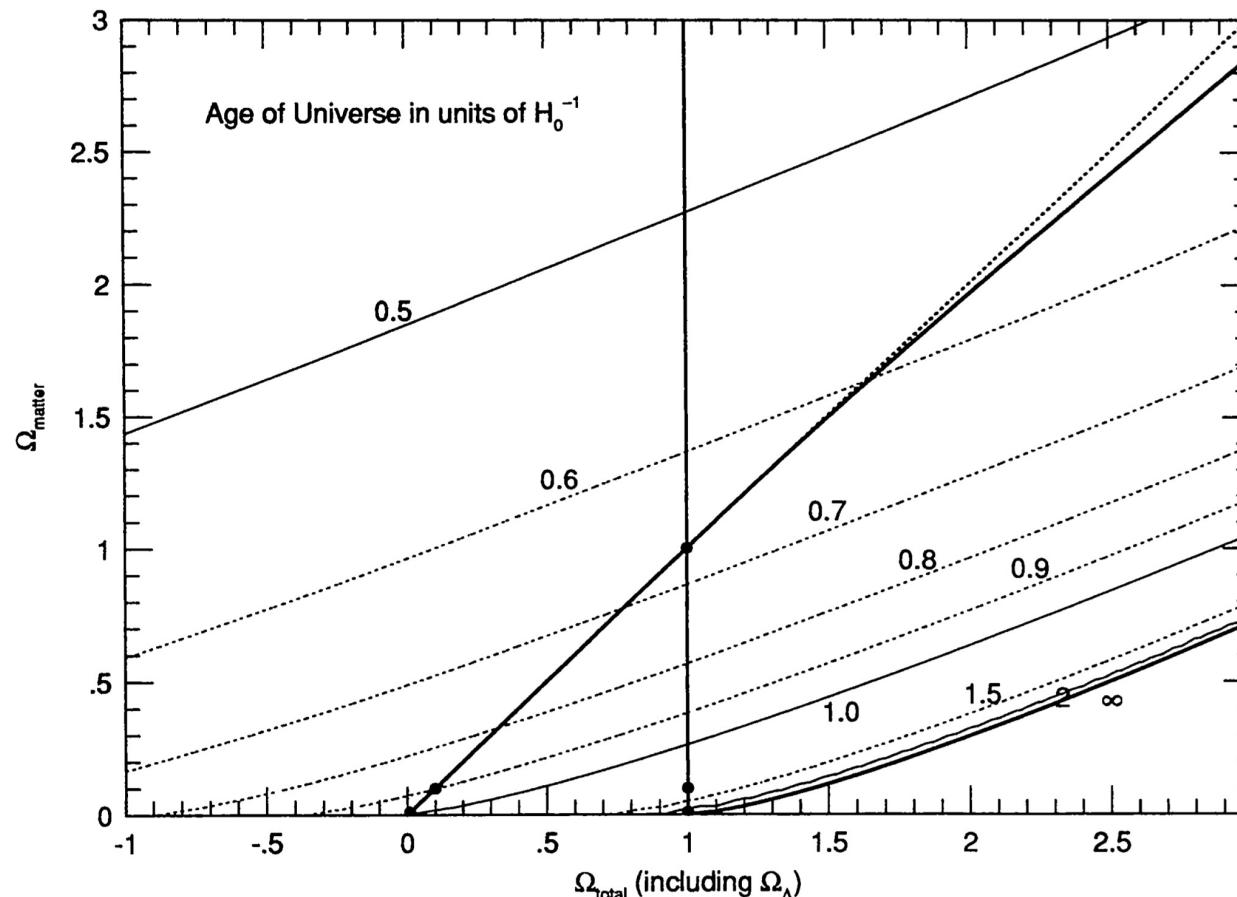
Radiation density  $\rho_{r,0} = 2\frac{\pi^2}{30} T_0^4 (1 + \frac{7}{8} \left(\frac{4}{11}\right)^{4/3} N_\nu) \Rightarrow \Omega_r = 8.5 \times 10^{-5} h_{.7}^{-2}$   
0.68 for 3 light ν's

(reheating of photons due to  $e^\pm$  annihilation)

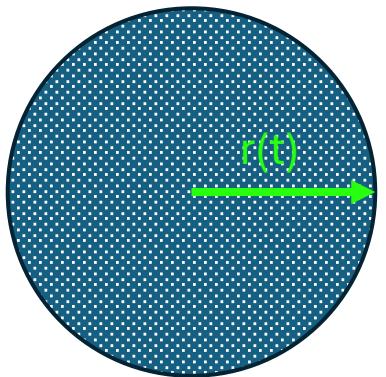
# Exact solution for LCDM model (neglecting radiation and spatial curvature)

$$a(t) = \left(\frac{1-\Omega_\Lambda}{\Omega_\Lambda}\right)^{1/3} \left(\sinh \frac{3}{2}H_0\Omega_\Lambda^{1/2}t\right)^{2/3} \Rightarrow t_0 = \frac{2}{3}H_0^{-1} \frac{1}{\sqrt{\Omega_\Lambda}} \sinh^{-1}(\sqrt{\Omega_\Lambda/(1-\Omega_\Lambda)})$$

(recommended exercise!)



# Cosmological perturbations (“Newtonian” derivation)



$$\frac{1}{2}\dot{r}^2 - \frac{GM}{r} - \frac{4\pi G}{3}r^2\rho_{\Lambda} = const$$

$$1. \ r \rightarrow r + \delta r \Rightarrow \dot{r}\delta r + \frac{GM}{r^2}\delta r - \frac{8\pi G}{3}r\delta r\rho_{\Lambda} = \delta k$$

$$2. \frac{d}{dt} \Rightarrow \ddot{r} = -\frac{GM}{r^2} + \frac{8\pi G}{3}\rho_{\Lambda}r \Rightarrow \frac{d}{dt}\left(\frac{\delta r}{r}\right) = \frac{\ddot{r}}{\dot{r}}\frac{\delta r}{r} - \frac{\dot{r}\delta r}{r^2} + \frac{\delta k}{r\dot{r}}$$

(recommended exercise!)

$$\delta_m \propto \frac{5}{2} \Omega_m \sqrt{\Omega_\Lambda + \Omega_m a^{-3}} \int_0^a \frac{da}{a^3 (\Omega_\Lambda + \Omega_m a^{-3})^{\frac{3}{2}}} \approx a, \quad a \ll (\frac{\Omega_m}{\Omega_\Lambda})^{1/3},$$

$$\approx 1.43 (\frac{\Omega_m}{\Omega_\Lambda})^{1/3}, \quad a \gg (\frac{\Omega_m}{\Omega_\Lambda})^{1/3}$$

Newtonian potential  $a^{-2} \nabla^2 \Phi = 4\pi G \rho_m \delta_m \Rightarrow \Phi \sim \text{const}, a^{-1}$  respectively

LCDM provides a remarkably good fit to the large-scale universe, with just 5 fundamental physics parameters

the energy content

1.  $\rho_\Lambda = (2.3 \text{ meV})^4 (\pm 1\%)$
2.  $\rho_{DM}/\rho_B = 5.36 (\pm 1\%)$
3.  $n_B/n_\gamma = 6 \times 10^{-10} (\pm 1\%)$

the perturbations

$$\langle \Phi^2 \rangle = \int \frac{dk}{k} A_\Phi \left( \frac{k}{k_*} \right)^{n_s - 1}$$

4. amplitude  $A_\Phi \approx 7.6 \pm 0.1 \times 10^{-10}$
5. “tilt”  $n_s - 1 \approx -0.041 \pm 0.0056$

many parameters so far consistent with zero:  
tensor and “isocurvature” perturbations,  
spatial curvature  $\kappa$ , non-Gaussianity...

