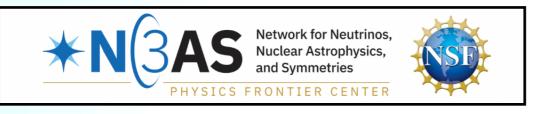
Detecting Rare Species of Dark Matter with Terrestrial Detectors

- i) Phys. Rev. Lett. 131, 011005 (2023) [arXiv: 2303.03416]
- ii) JCAP 01 029 (2024) [arXiv: 2309.10032]
- iii) arXiv: 2402.03431

Anupam Ray

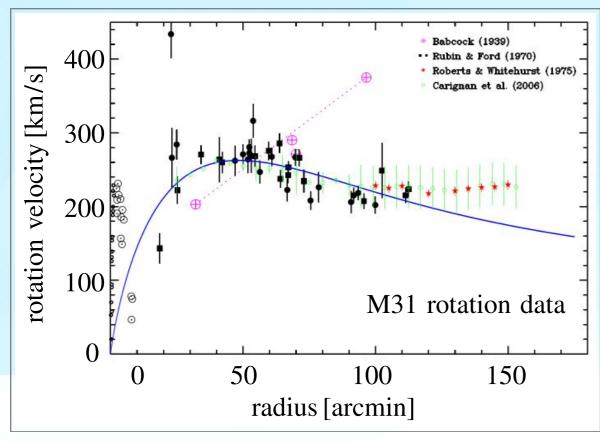
N3AS Fellow, UC Berkeley & University of Minnesota

Mitchell Conference, 2024 05.25.2024

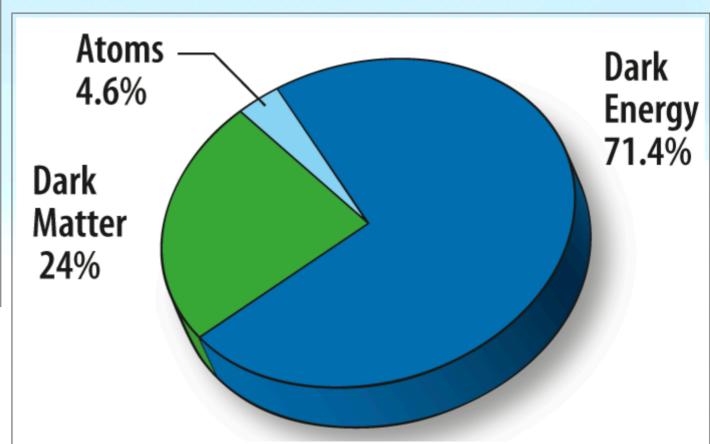




Dark Matter (DM)



From: Bertone and Hooper, Rev. Mod. Physics (2016)



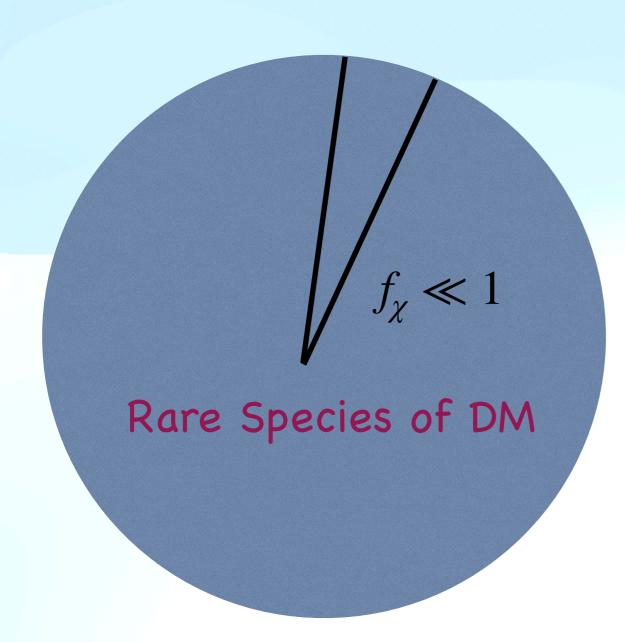
https://wmap.gsfc.nasa.gov/universe/uni_matter.html

• DM mass?

DM interactions with baryons?

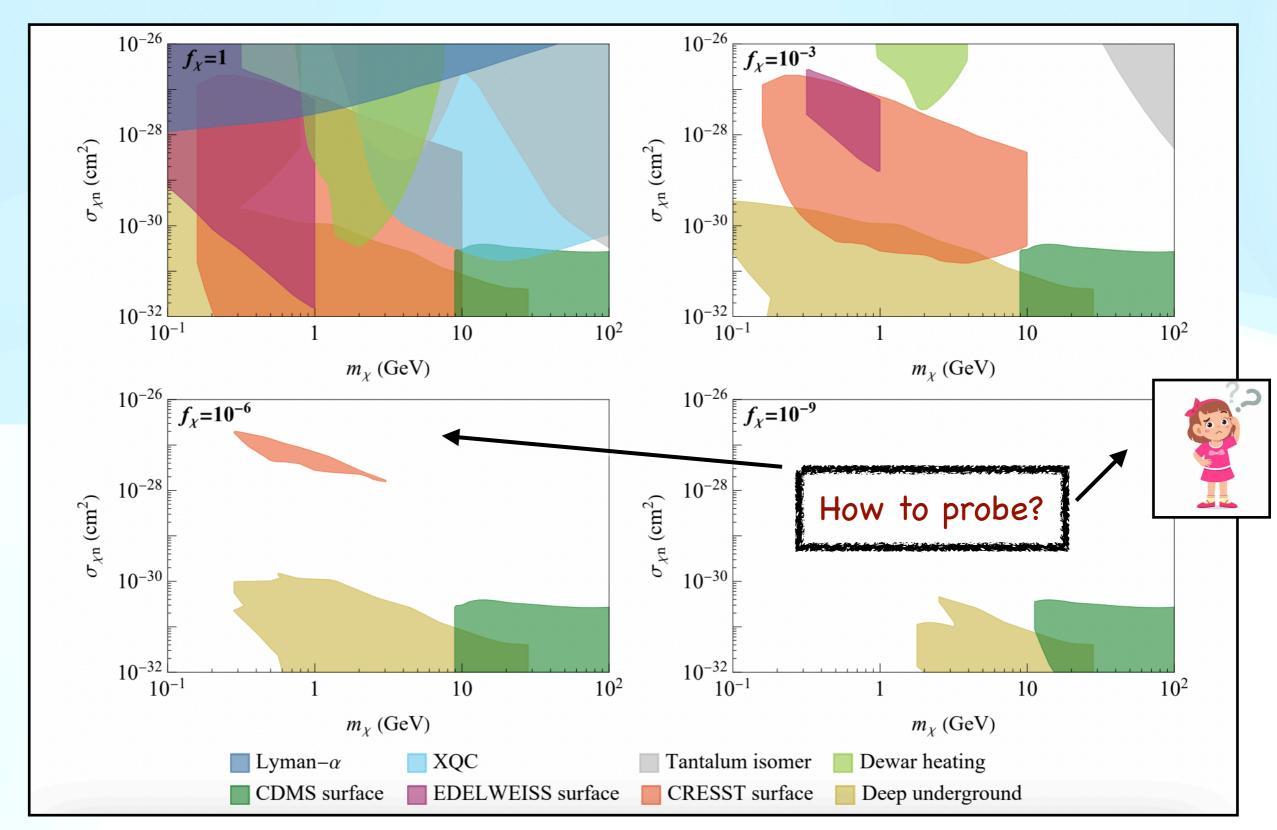
Strongly-interacting DM Component

A sub-component of DM can be strongly interacting.



 χ makes up a sub-component of the total DM energy budget.

Strongly-interacting DM Component



Take Away

• "Earth-bound" DM provides a novel powerful probe.



The density of "Earth-bound DM" can be huge.

Annihilating DM

 Local annihilation inside any large-volume neutrino detectors (such as Super-Kamiokande)

Ray, (with Mckeen, Morissey, Pospelov, Ramani) [PRL, 2023]

 Neutrinos from annihilation of Earth-bound DM.

Pospelov & Ray [JCAP, 2024]

Non-Annihilating DM

 Earth-bound DM can be up-scattered by fast neutrons inside the nuclear reactors, and subsequently detected.

(similar scheme as $CE\nu NS$)

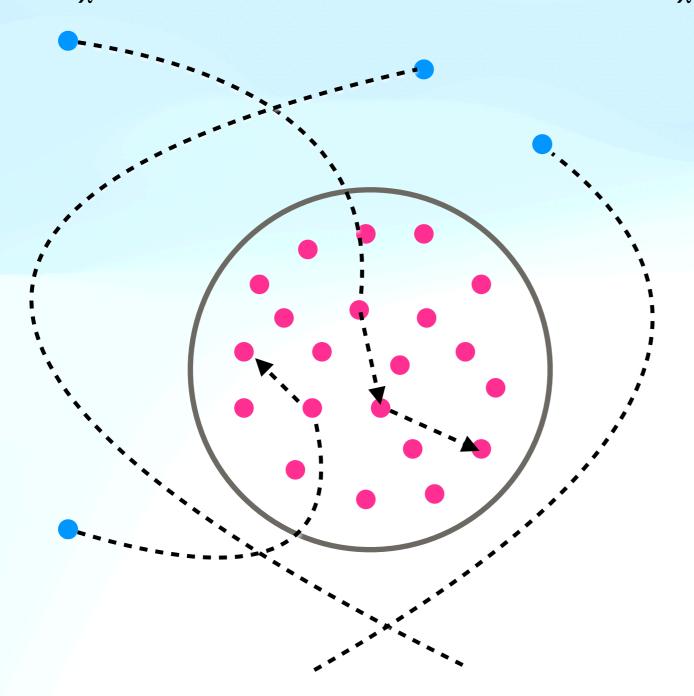
Ray, (with Ema, Pospelov)
[2402.03431]

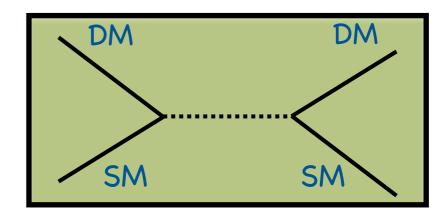
Earth-Bound DM

Press & Spergel (1985, ApJ), Gould (1987, ApJ),...

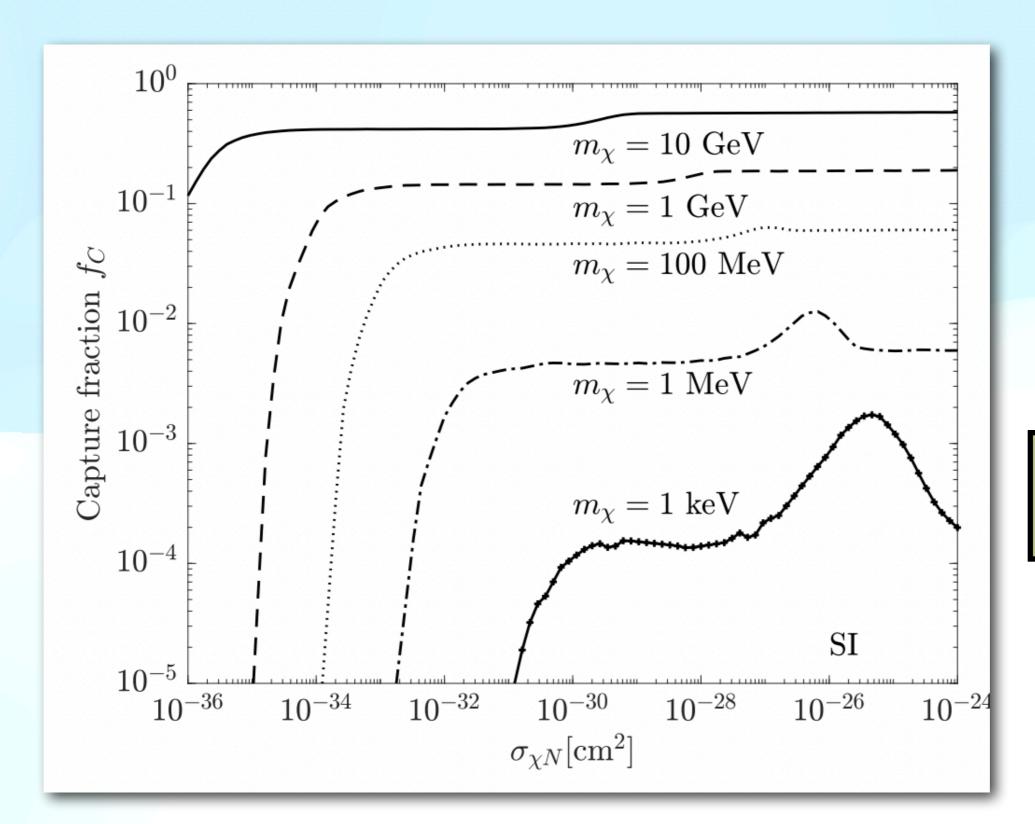
Small $\sigma_{\chi n} \to \text{ single collision,}$

large $\sigma_{\!\chi n} o \,$ multiple collisions.





Earth-Bound DM



$$f_c\left(\sigma_{\chi n}, m_{\chi}\right)$$

Earth-Bound DM

Lets do some estimate:

For DM mass of 1 GeV and
$$\sigma_{\!\chi n}=10^{-28}\,\mathrm{cm}^2$$

$$C_{\rm geo} = 1.3 \times 10^{25} \, {\rm s}^{-1}$$
 and $f_c \sim 0.1$ $f_{\chi} = 1$

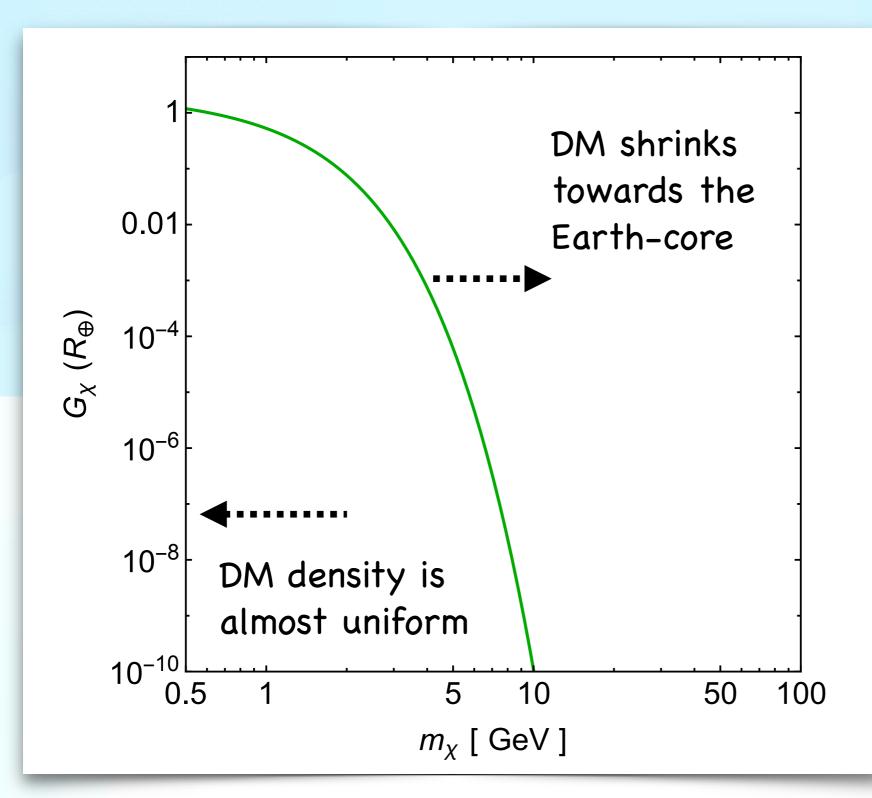
DM density (assuming they uniformly distribute over the Earth-volume)

$$\rho_{\chi} = m_{\chi} \frac{f_c \times C_{\text{geo}} \times t_{\oplus}}{V_{\oplus}} \sim 3 \times 10^{14} \,\text{GeV/cm}^3$$

$$f_{\chi} = 1$$

15 orders of magnitude larger than the Galactic DM density!

DM Distribution in Stellar Objects



 Dimensionless profile function:

$$G_{\chi}(R_{\oplus}) = \frac{n_{\chi}(R_{\oplus})V_{\oplus}}{N_{\chi}}$$

 For uniform DM density:

$$G_{\chi}(R_{\oplus}) = 1$$

Signal at Super-K

 Earth-bound DM, of mass GeV scale have an enormously large surface density.

• Their detection via scattering is almost impossible as they acquire very little amount kinetic energy (0.03 eV).

How to detect them?

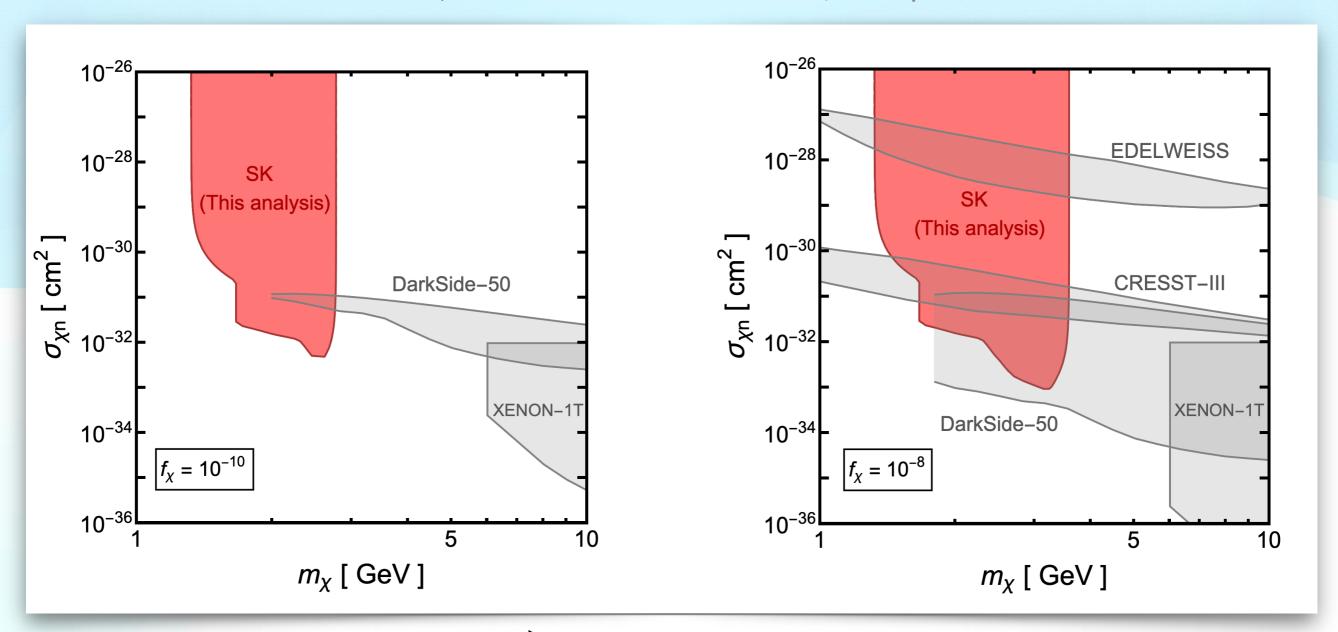
Ray, (with Mckeen, Morissey, Pospelov, Ramani) [PRL, 2023]

Our proposal: simply look at their annihilation signature inside large-volume detectors (annihilation is not limited to the tiny kinetic energy)!

Results

• Using existing di-nucleon annihilation searches at Super-K

Ray, (with Mckeen, Morissey, Pospelov, Ramani) [PRL, 2023]

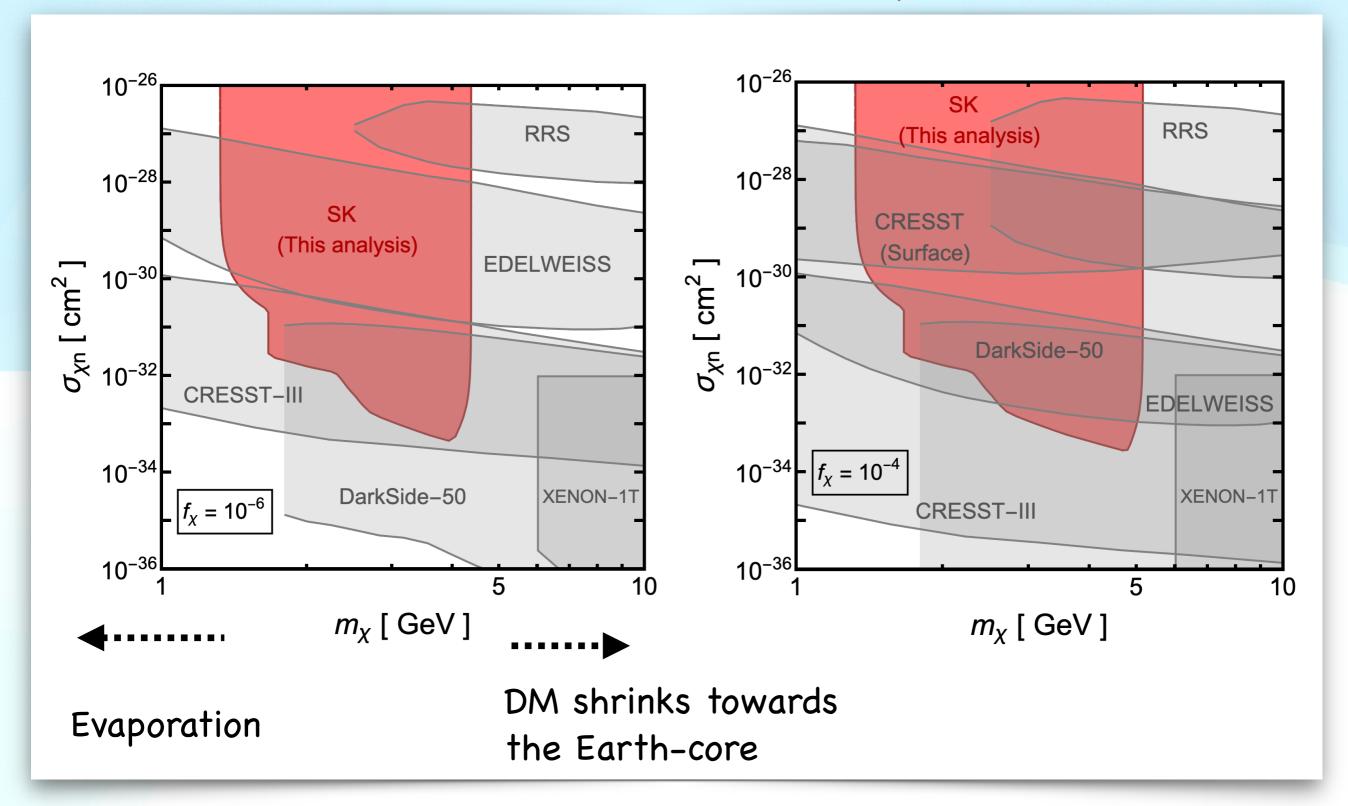


Evaporation

DM shrinks towards the Earth-core Up to $f_{\chi} = 10^{-10}$

Results

Ray, (with Mckeen, Morissey, Pospelov, Ramani) [PRL, 2023]



Model

 Let's illustrate our result in a concrete phenomenological model.

$$\mathcal{L} = -\frac{1}{4} \left(F'_{\mu\nu} \right)^2 - \frac{\epsilon}{2} F'_{\mu\nu} F^{\mu\nu} + \frac{1}{2} m_{A'}^2 \left(A'_{\mu} \right)^2 + \bar{\chi} (i \gamma^{\mu} D_{\mu} - m_{\chi}) \chi$$

 χ : Dirac fermion which can couple to a dark photon A'

• The perturbative cross section for χ to scatter on a nucleus (Z, A) is related to the model parameters

$$\sigma_{\chi A} = \frac{16\pi Z^2 \alpha \alpha_d \epsilon^2 \mu_{\chi A}^2}{m_{A'}^4}$$

Model

· We are interested in the following channel

$$\chi \bar{\chi} \rightarrow A'A'$$
 with $A' \rightarrow SM + SM$ (say $e^+ + e^-$)

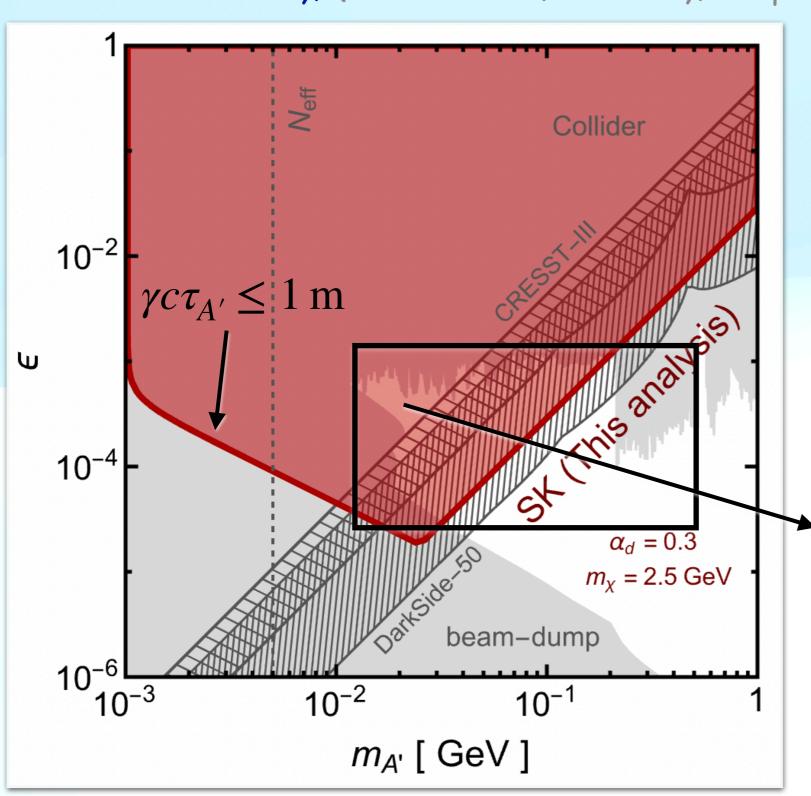
$$\langle \sigma v \rangle_{\text{ann}} = \frac{\pi \alpha_d^2}{m_\chi^2} \frac{\left(1 - m_{A'}^2 / m_\chi^2\right)^{3/2}}{\left(1 - m_{A'}^2 / 4 m_\chi^2\right)^2}$$

$$\Gamma_{A'} = \frac{1}{3} \alpha \epsilon^2 m_{A'} \left(1 + \frac{2m_e^2}{m_{A'}^2}\right) \left(1 - \frac{4m_e^2}{m_{A'}^2}\right)^{1/2}$$

• To ensure the decay within the Super-K fiducial volume, we restrict the decay length $\gamma c \tau_{A'} \leq 1 \, \mathrm{m}$.

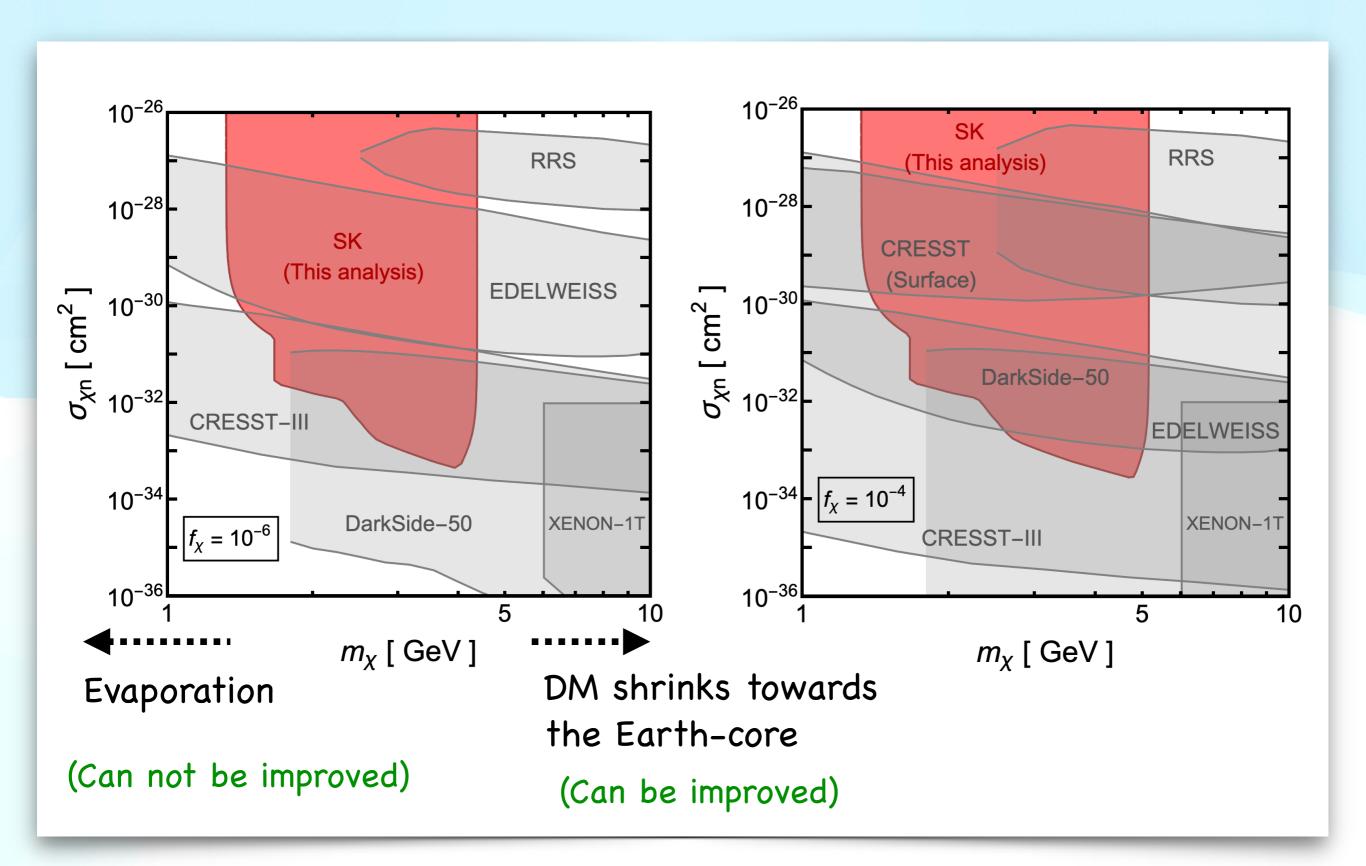
Results

Ray, (with Mckeen, Morissey, Pospelov, Ramani) [PRL, 2023]



Unprecedented sensitivity on parts of the parameter space.

What about heavy DM?



Neutrino Signal

 Earth-bound DM if sufficiently heavy, shrinks towards the core, leading to a negligible surface density.

gravity dominates over the diffusion processes

 Annihilation to neutrinos can occur at the Earth-core, if Earth-bound DM if sufficiently heavy. Since the number density is huge, annihilation rate is also fairly large.

 Neutrinos, because of their feeble interactions, can reach detectors like Super-K, IceCube-DeepCore, and searching these annihilated neutrinos can provide sensitivity to DM interactions.

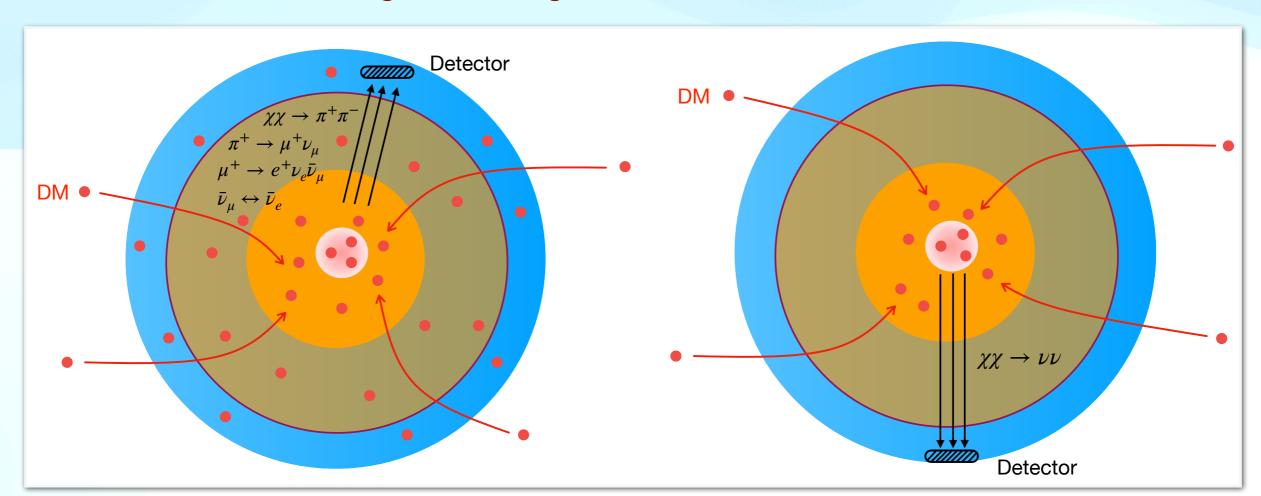
Pospelov & Ray [JCAP, 2024]

Neutrino Signal

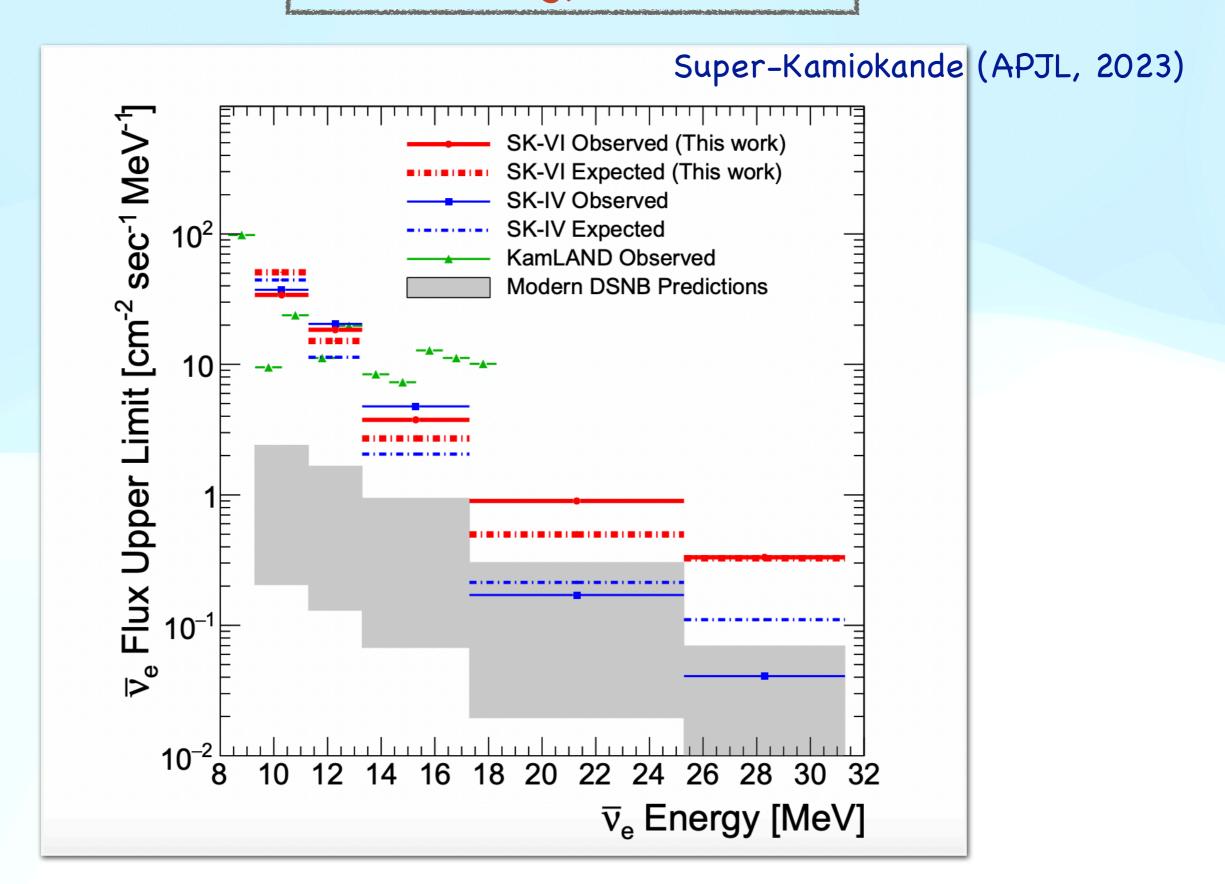
• We consider two phenological scenarios:

Lower energy neutrinos from the stopped pion decay

Higher energy neutrino lines from direct annihilation

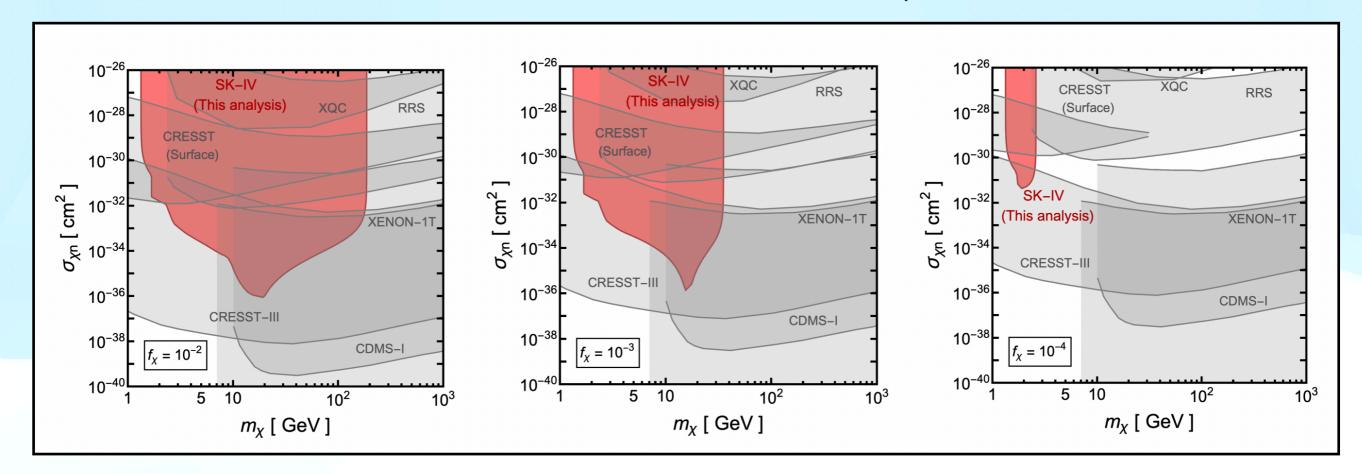


Low Energy Neutrinos



Low Energy Neutrinos

Pospelov & Ray [JCAP, 2024]

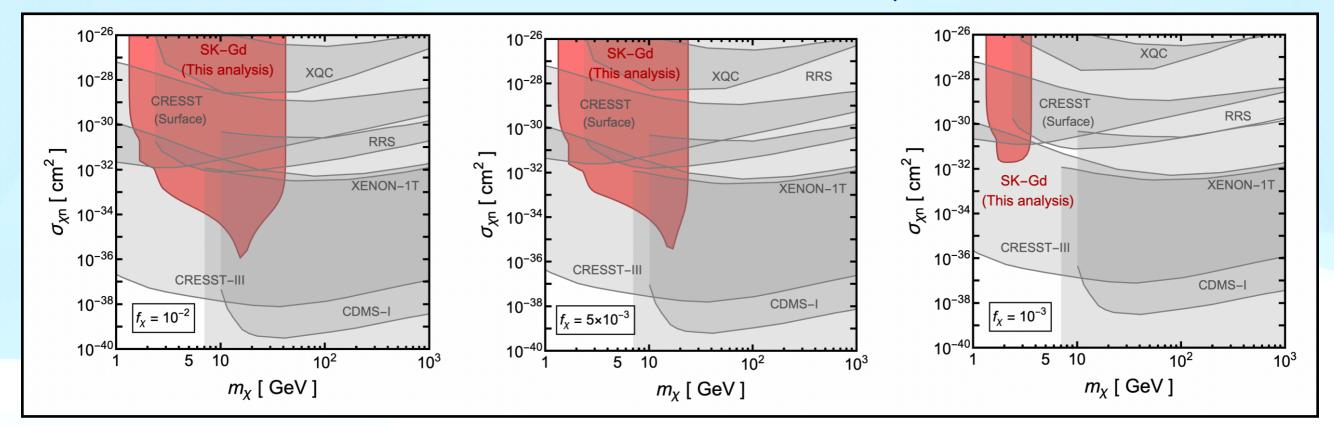


We use the Super-K DSNB search result with pure-water $(22.5 \text{ kton} \times 2970 \text{ days})$ to derive the exclusion limits.

Super-Kamiokande (PRD, 2021)

Low Energy Neutrinos

Pospelov & Ray [JCAP, 2024]



We use the Super-K DSNB search result with 0.01 wt% gadolinium loaded water (22.5 kton \times 552.2 days) to derive the exclusion limits

Super-Kamiokande (APJL, 2023)

*Gd-loaded water gives competitive limit (as compared to the pure-water limits) although the data is 5 times less.

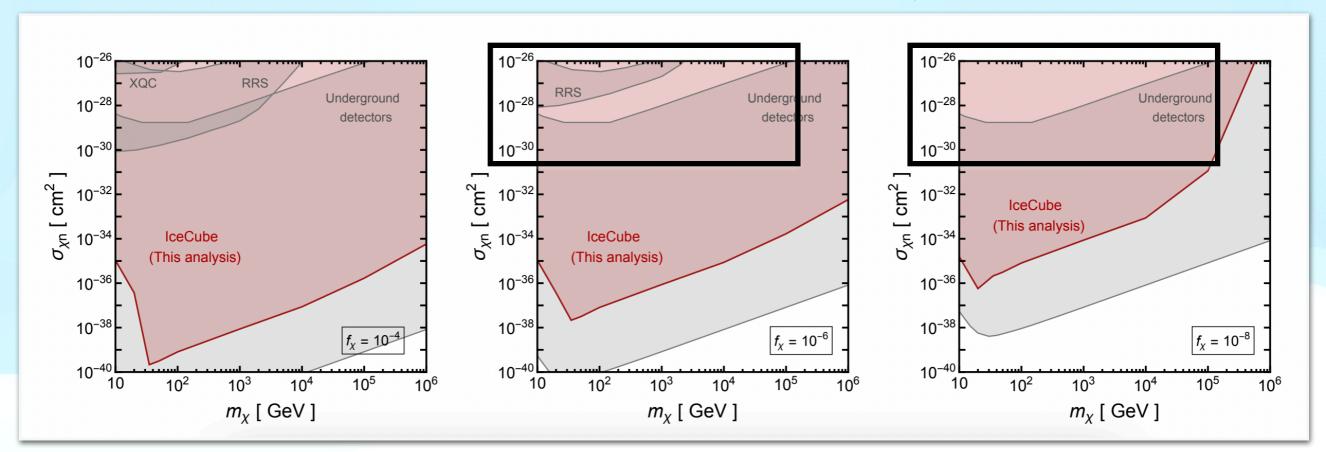
High Energy Neutrinos

- DM annihilation directly to neutrinos yields a line at $E_{\nu}=m_{\chi}$ high-energy neutrinos can also come from $\chi\chi\to W^+W^-, b\bar{b}, \tau\bar{\tau}$, giving a continuum spectra up to $E_{\nu}=m_{\chi}$ (or $\chi\chi\to A'A'\to 4\nu$).
 - We search the "neutrino-line" signature in the IceCube DeepCore data with a total live-time of 6.75 years.
- We use the null-detection of the neutrino-line signature in the IceCube DeepCore data to derive the exclusions

			Marie Commission Commi
	$\mathbf{b}\mathbf{ar{b}}$	$ auar{ au}$	$lackbr{\ell}$ $ uar{ u}$
Mass (GeV)	$\Gamma_{ m ann}~[{ m s}^{-1}]~ imes 10^{23}$	$\Gamma_{\mathrm{ann}} \ [\mathrm{s}^{-1}] \ imes 10^{23}$	$\Gamma_{ m ann}~[m s^{-1}]~ imes 10^{23}$
5	139	139.3	· ·
10	396	7.0	1.37
20	29.7	0.97	0.27
35	7.41	0.22	0.09
50	3.51	0.096	0.05
100	1.39	0.038	0.027

High Energy Neutrinos

Pospelov & Ray [JCAP, 2024]



We probe up to $f_{\gamma} \ge 10^{-8}$ for significantly heavy Earth-bound DM.

Earth as the most optimal detector

• Earth accumulates fewer number of DM particles as compared to the Sun. (by a factor of $\sim R_{\oplus}^2/R_{\odot}^2$)

$$\Gamma_{\text{cap}} = f_c \frac{\rho_{\chi}}{m_{\chi}} \pi R^2 \int \frac{f(u)du}{u} (u^2 + v_{\text{esc}}^2)$$

• But, for Earth-bound DM, distance to the detector is far less.

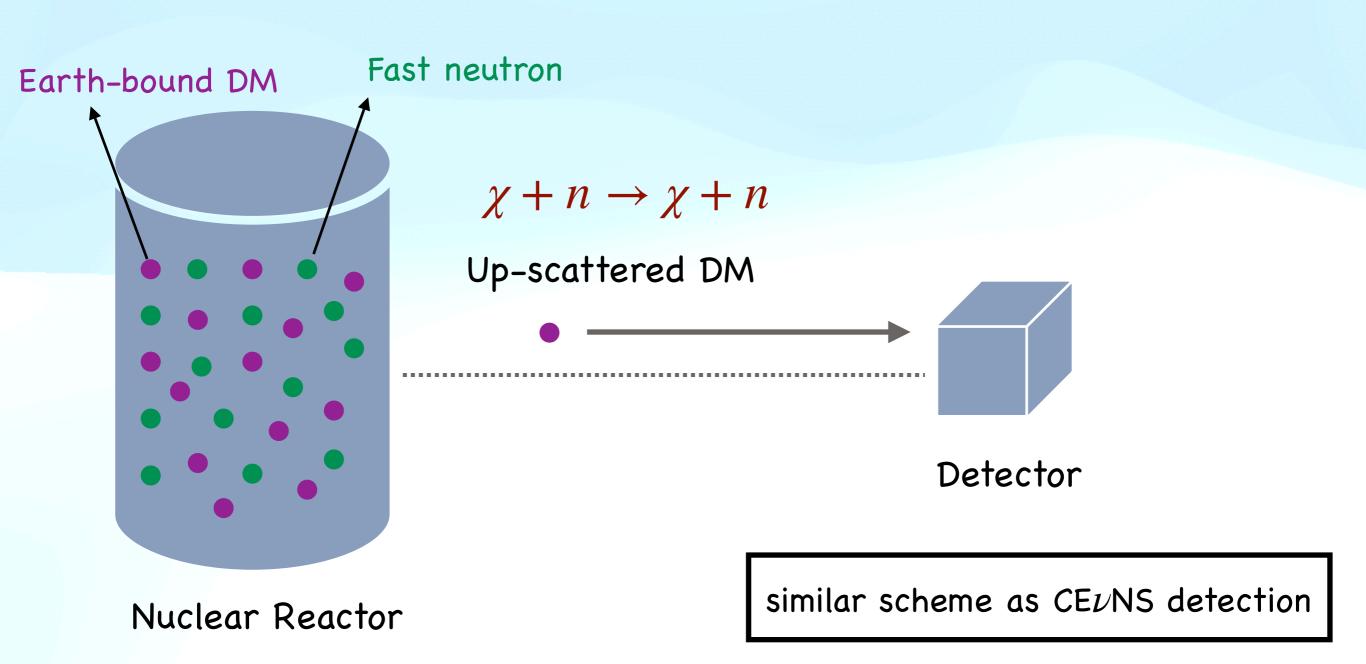
$$\phi_{\oplus} \sim \frac{\Gamma_{\mathrm{cap}}}{4\pi R_{\oplus}^2}$$
 and $\phi_{\odot} \sim \frac{\Gamma_{\mathrm{cap}}}{4\pi D^2}$

Flux for Earth-bound DM is ~ 4000 larger than the neutrino flux from Sun.

This is quite different from standard weakly-interacting paradigm where Sun is the most-optimal detector, and hence, has been studied over the past few decades.

Non-Annihilating DM

 Nuclear Reactors act as powerful probe of Earth-bound DM detection.



Ray, (with Ema, Pospelov) [2402.03431]

Non-Annihilating DM

· Accumulation of Earth-bound DM.



• Distribution of Earth-bound DM.



 Up-scattering of Earth-bound DM inside Nuclear Reactors by fast neutrons (typically of MeV energy).

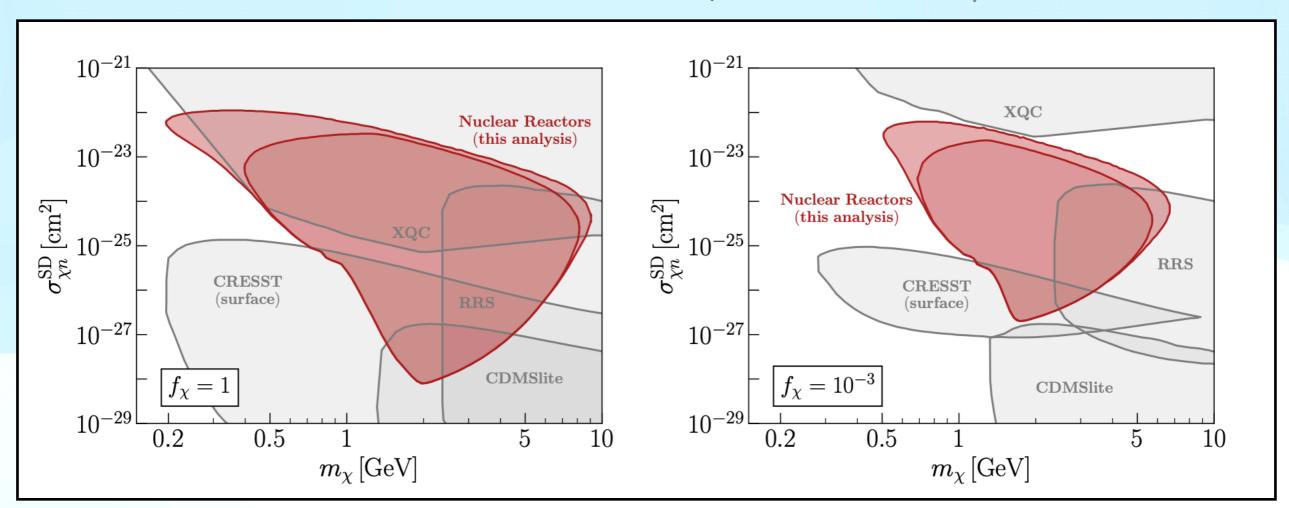
We use CONUS experiment setup for our analysis.

 Subsequent propagation through shielding and detection via scattering.

We use MC simulations for the propagation along with provide an analytical recipe.

Results

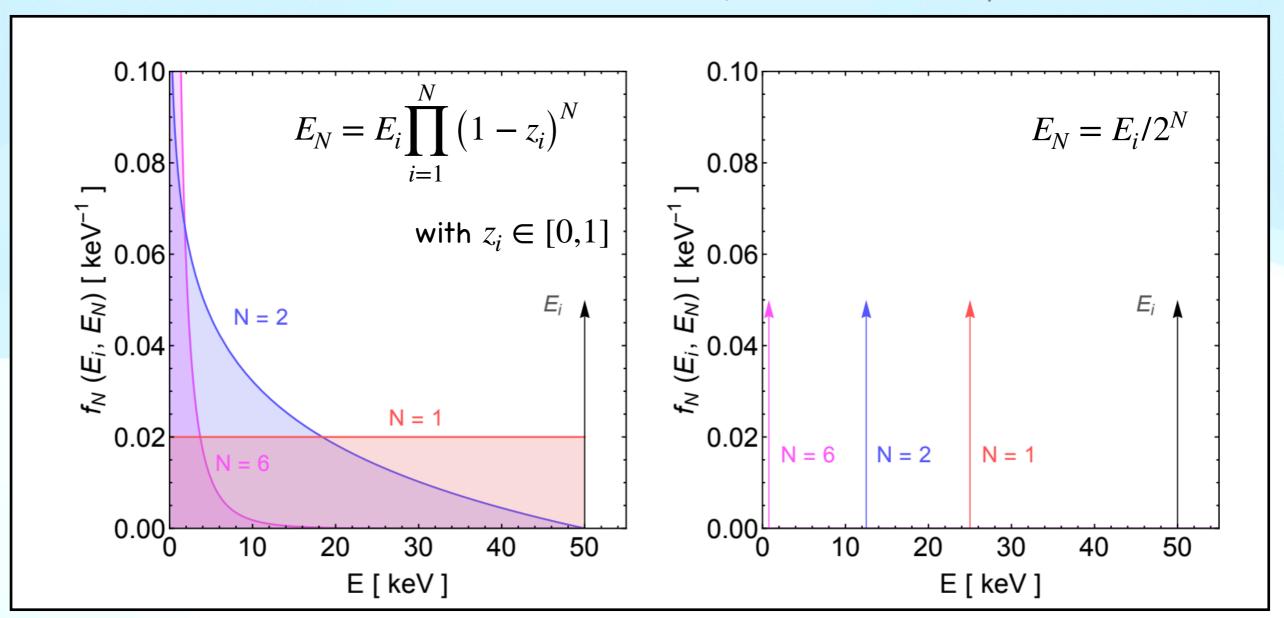
Ray, (with Ema, Pospelov) [2402.03431]



Smaller regions: includes only the DM particles which do not experience any collisions. Bigger regions: includes the full multiple-scattering contributions.

Propagation of Up-scattered DM

Ray, (with Ema, Pospelov) [2402.03431]



Tail of the distribution is utterly important. Many previous studies (e.g., Bramante et al [PRD, 2017], Leane et al. [JCAP, 2022] etc) neglect this simple yet important point.

Summary

- Earth accumulates significant number of DM particles from the Galactic halo, leading to a DM density 15 orders of magnitude larger than the Galactic DM density!
- Despite their prodigious abundance, their detection is extremely challenging as they acquire tiny amount of kinetic energy.
- Annihilation of such Earth-bound DM at large-volume neutrino detectors, provides a novel way for their detection and can be used to probe strongly-interacting DM component.
- If Earth-bound DM do not annihilate among themselves, upscattering them inside nuclear reactors provides a powerful probe of their detection.

Conclusion

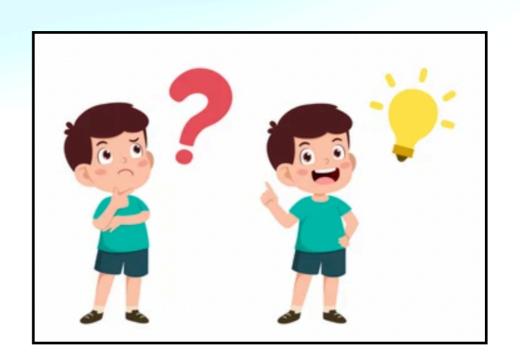


How to detect rare species of DM?





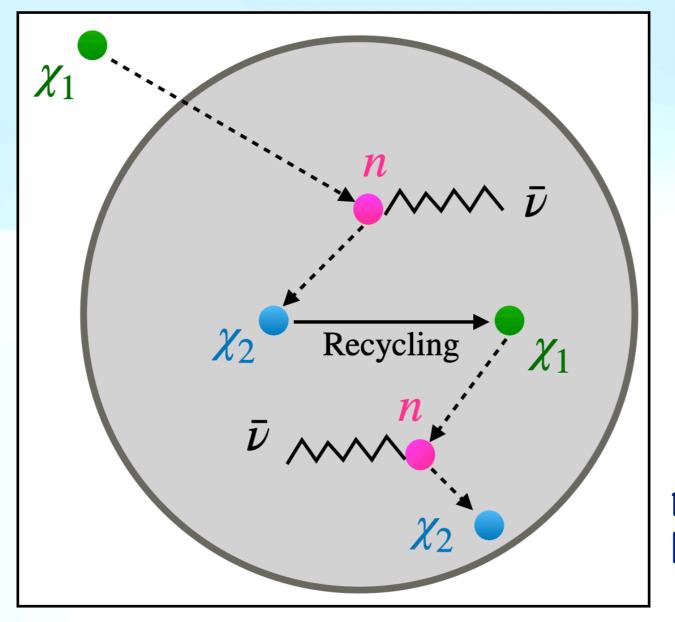
Look at the Earth-bound DM!



Questions & Comments: anupam.ray@berkeley.edu

What about the Major Portion?

 Major portion of the DM could be weakly-interacting, and can have baryon-number-violating interactions.



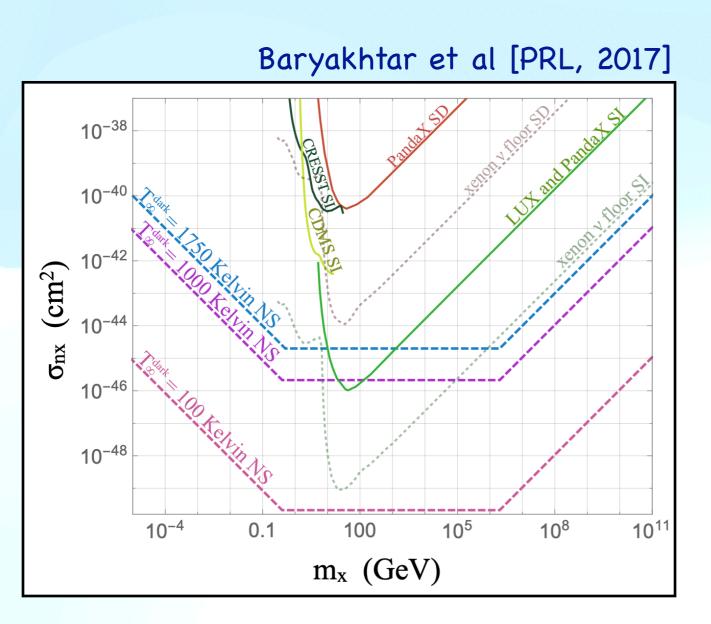
How to probe?

Ray, (with Ema, McGhee, Pospelov) [in prep.]

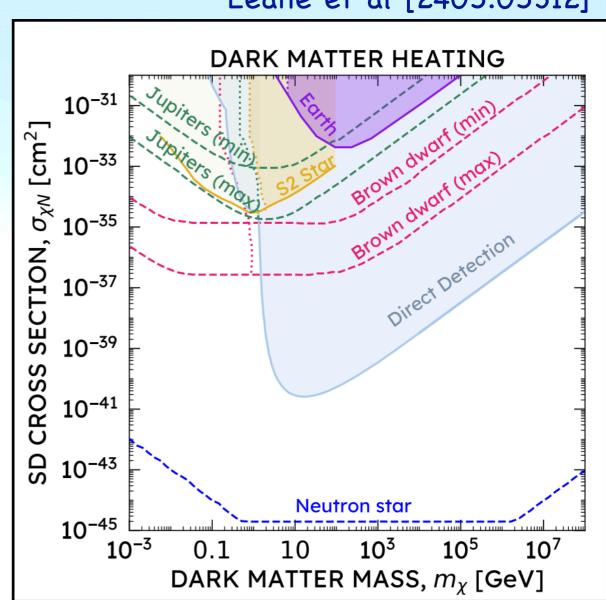
Leads to anomalous heating of cold Neutron Stars.

What about the Major Portion?

 Anomalous neutron star heating via captured DM annihilation (or via kinetic energy transfer).



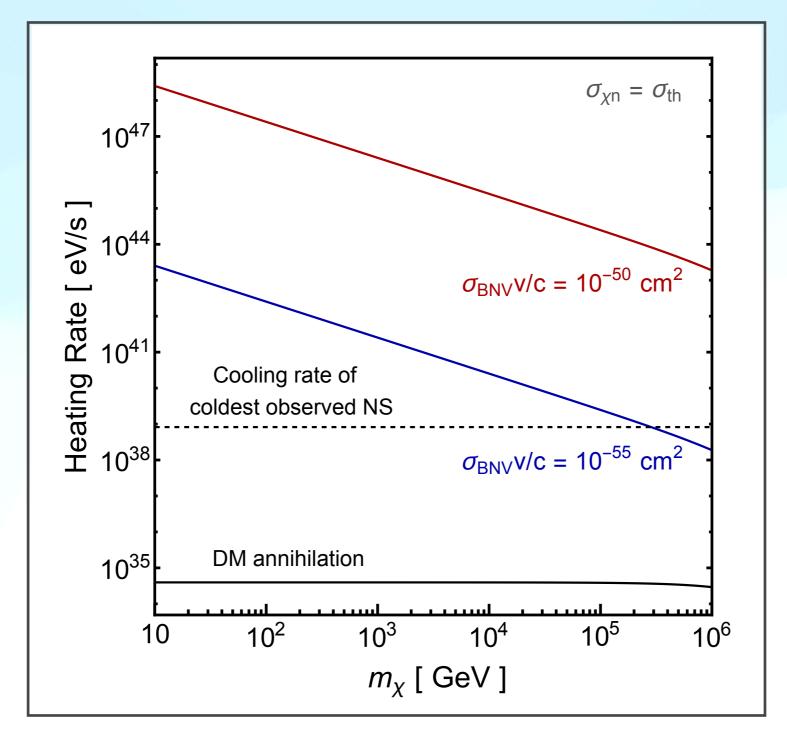
Leane et al [2405.05312]



Coldest NS so far seen has $T \sim 40,000\,K$ (significantly larger than these estimates!)

DM-induced BNV

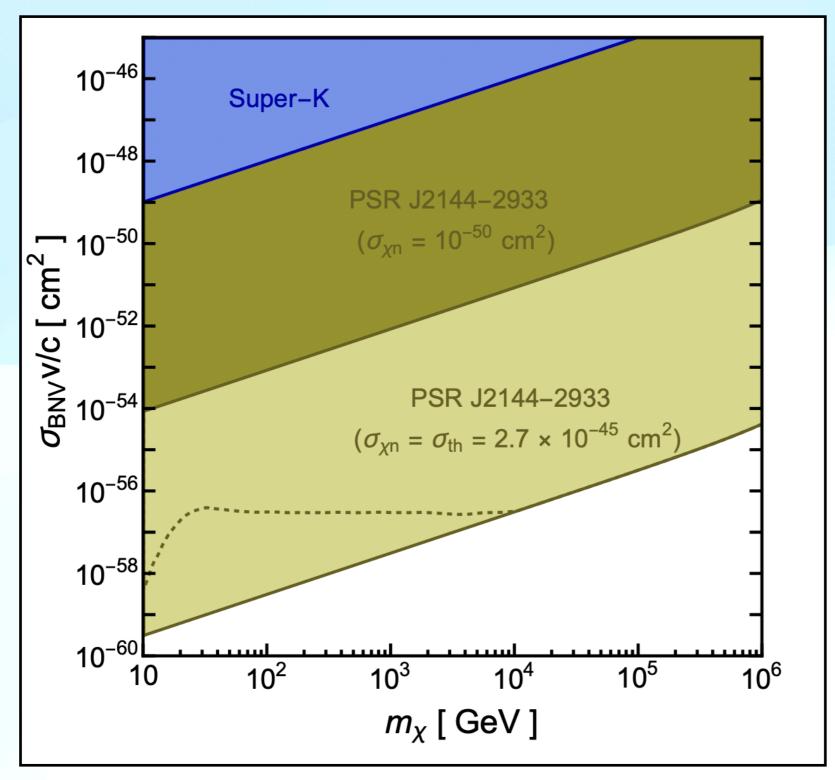
 Total amount of DM captured inside a neutron star is minuscule as compared to the total neutron mass.



Ray, (with Ema, McGhee, Pospelov) [in prep.]

DM-induced BNV

How strong can DM-induced BNV interactions be?



Ray, (with Ema, McGhee, Pospelov) [in prep.]