



Collinear laser spectroscopy at COLLAPS and analysis of ¹⁵⁵⁻¹⁷⁵Tm

STFC summer school 2024

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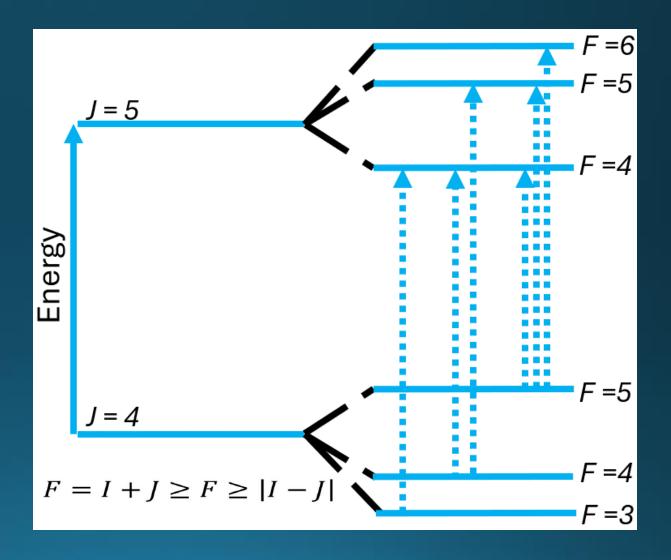




Laser spectroscopy gives a window into the Hyperfine structure.

Hyperfine structure arises in the atomic spectra from the fine structure of a nucleus due to interactions between electrons and the nucleus.

With enough resolution, a single electronic transition can be seen to split into various transitions between different levels.







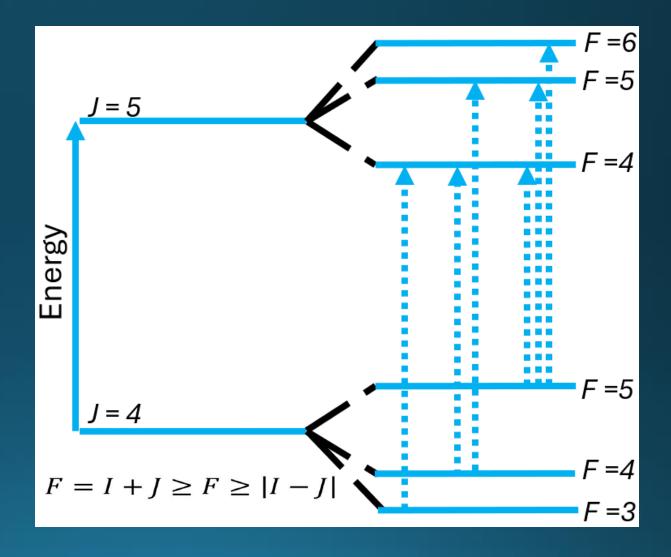
With enough resolution, a single electronic transition can be seen to split into various transitions between different levels.

This splitting is given by:

$$\frac{\Delta E}{h} = \frac{K}{2}A + \frac{3K(K+1) - 4I(I+1)J(J+1)}{8I(2I-1)J(2J-1)}B$$

A and B are the main parameters of note, as they relate to the nuclear properties:

$$A=\frac{\mu_I B_e(0)}{IJ}$$
 , $B=eQ_s\,\langle \frac{\partial^2 V_e}{\partial z^2}\rangle$







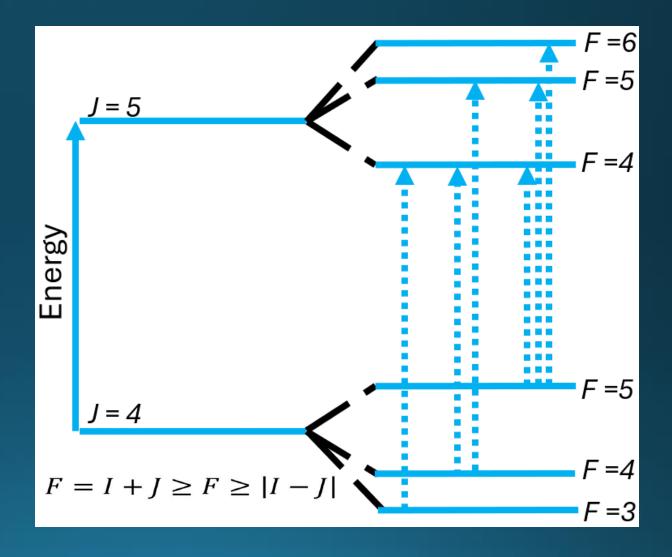
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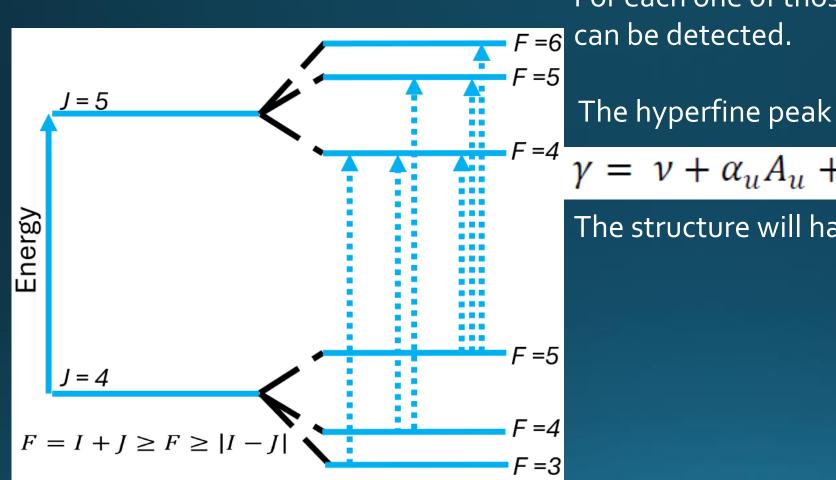
A and B are the main parameters of note, as they relate to the nuclear properties:

$$A = \underbrace{ \stackrel{\mu_I B_e(0)}{I J}}_{I J}, B = e \underbrace{ Q_s }_{\partial z^2} \frac{\partial^2 V_e}{\partial z^2} \rangle$$









For each one of those transitions. A hyperfine peak

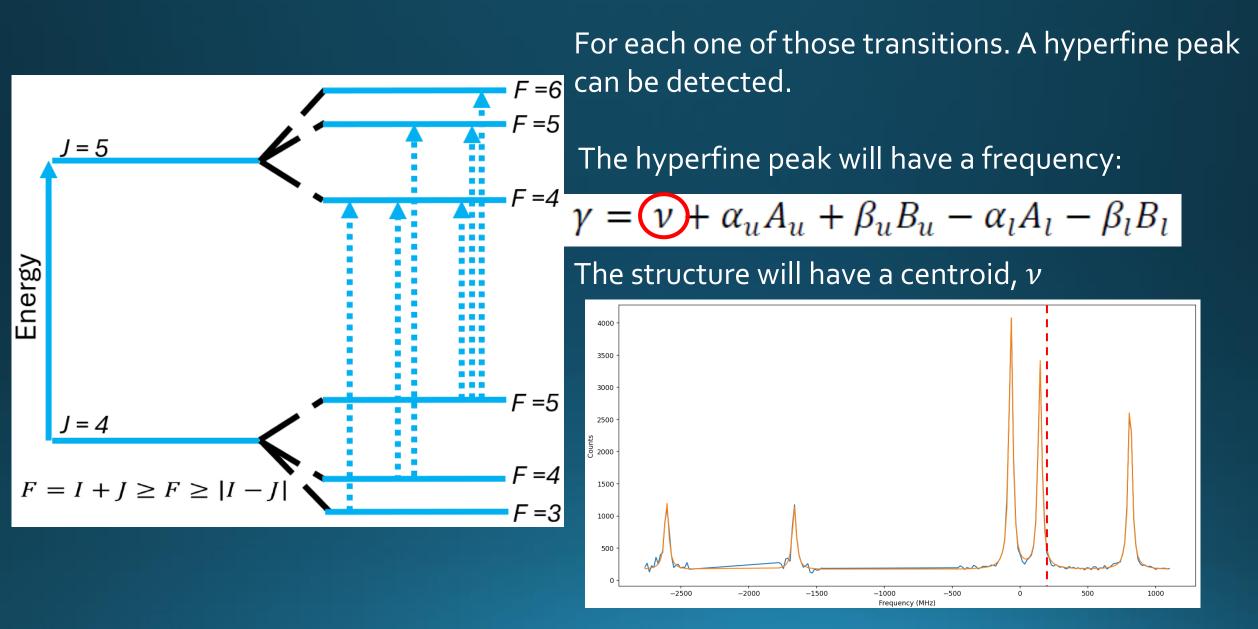
The hyperfine peak will have a frequency:

$$\nu = \nu + \alpha_u A_u + \beta_u B_u - \alpha_l A_l - \beta_l B_l$$

The structure will have a centroid, ν

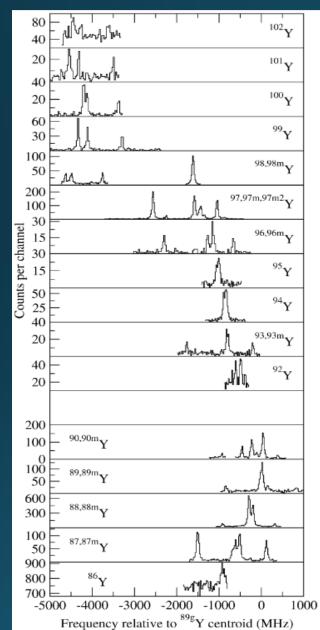












For each one of those transitions. A hyperfine peak can be detected.

The hyperfine peak will have a frequency:

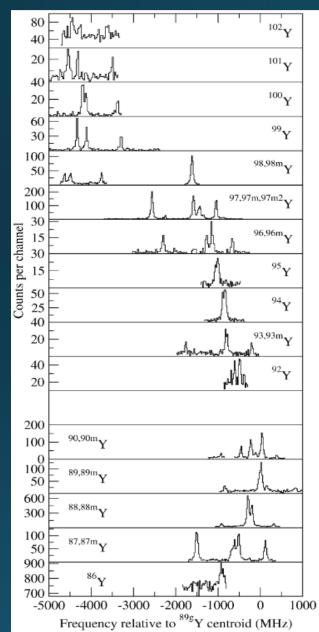
$$\gamma = \nu + \alpha_u A_u + \beta_u B_u - \alpha_l A_l - \beta_l B_l$$

The structure will have a centroid, ν

This centroid is seen to change from isotope to isotope, and is known as the isotope shift.







$$v = v + \alpha_u A_u + \beta_u B_u - \alpha_l A_l - \beta_l B_l$$

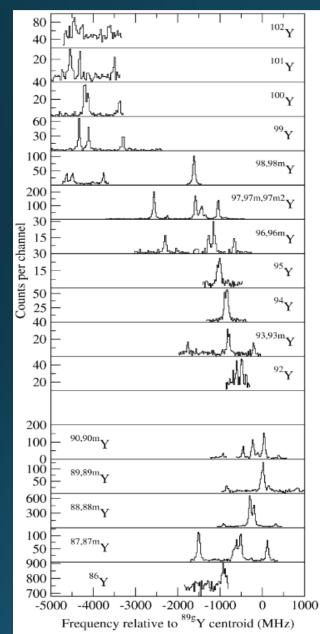
This centroid is seen to change from isotope to isotope, and is known as the isotope shift.

From this, we can deduce The change in mean square charge radius between 2 isotopes (A and A'):

$$\nu_{A'} - \nu_A = \delta \nu^{AA'} = M \frac{A' - A}{AA'} + F \delta \langle r^2 \rangle^{AA'}$$







$$v = v + \alpha_u A_u + \beta_u B_u - \alpha_l A_l - \beta_l B_l$$

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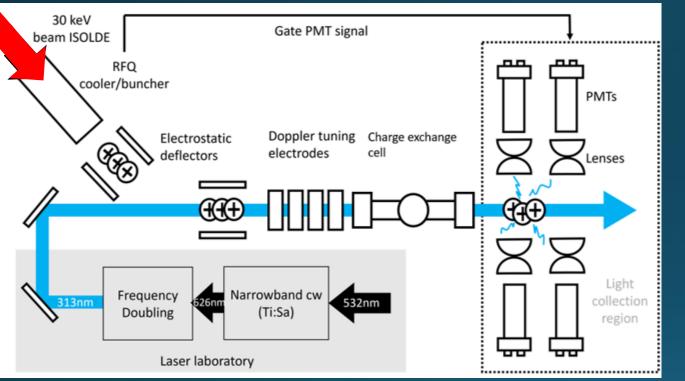
$$\nu_{A'} - \nu_A = \delta \nu^{AA'} = M \frac{A' - A}{AA'} + K \delta \langle r^2 \rangle^{AA'}$$

This property gives insight into the size and deformation of the nucleus









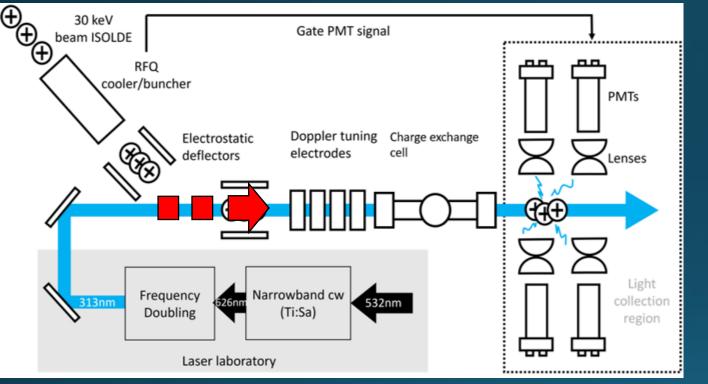
Radioactive ion beam from ISOLDE arrives into the cooler/buncher.

Cooler buncher converts the continuous beam into pulsed bunches.



COLLAPS





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Cooler buncher converts the continuous beam into pulsed bunches.

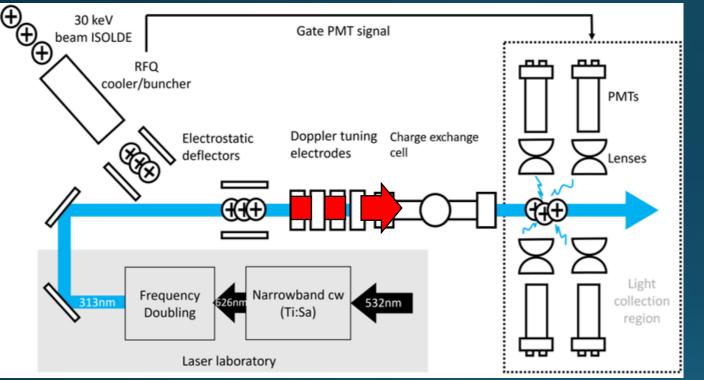
Beam accumulates in the buncher through the use of a potential.

Pulses are then accelerated and released from the buncher.



COLLAPS





Pulses are then accelerated and released from the buncher.

Ion pulses are deflected into the path of the laser.

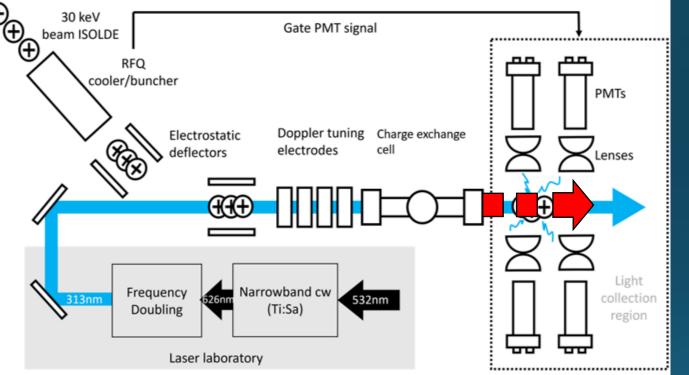
The speed of the ions is adjusted by the tuning electrodes.

This allows the frequency to be scanned using the Doppler shift









This allows the frequency to be scanned using the Doppler shift

Ions travel through the CEC and arrive at the light collection region

Photons which are re-emitted after resonant excitation can then be detected.

Pulsed beams allow for substantial background suppression.

From this, a hyperfine spectrum of counts against frequency can be obtained.



Previous work



A	I^{π}	$\tau_{1/2}$	μ ($\mu_{\rm N}$)	Q_s (b)	$\delta \langle \mathbf{r}^2 \rangle^{(169,A)} (\mathrm{fm}^2)$	Ref
175	$(1/2^+)$	15.2m	?	-	?	-
174	(4)-	$5.4\mathrm{m}$?	?	?	-
174m	0+	2.29s	-	-	?	-
173	$(1/2^+)$	8.24h	?	-	?	-
172	2-	63.6h	?	?	+0.154(30)	[1]
171	$1/2^+$	1.9y	-0.230(4)	-	+0.161(6)	[1]
170	1-	128.6d	+0.247(4)	+0.74(2)	+0.070(14)	[1]
169	$1/2^+$	stable	-0.2310(15)	-	0	[1]
168	3^{+}	93.1d	+0.226(11)	+3.23(7)	-0.084(4)	[1]
167	$1/2^+$	9.25d	-0.197(2)	-	-0.126(3)	[1]
166	2+	7.7h	+0.092(1)	+2.14(3)	-0.209(3)	[1]
166m	(6 ⁻)	$340 \mathrm{ms}$?	?	?	-
165	$1/2^{+}$	30.06h	-0.139(2)	-	-0.250(2)	[1]
164	1+	$2.0\mathrm{m}$	+2.37(3)	+0.706(51)	-0.347(6)	[1]
164m	6-	$5.1 \mathrm{m}$?	?	?	-
163	$1/2^+$	1.810h	-0.082(1)	-	-0.404(2)	[1]
162	1-	21.70m	+0.068(8)	+0.69(3)	-0.537(5)	[1]
162m	5+	21.70m	?	?	?	-
161	$7/2^+$	30.2m	2.39(2)	+2.90(7)	-0.632(3)	[1]
160	1-	9.4m	+0.156(18)	+0.582(44)	-0.741(4)	[1]
$160 \mathrm{m}$	5?	74.5s	?	?	?	-
159	$5/2^{+}$	9.13m	+3.408(34)	+1.93(7)	-0.850(4)	[1]
158	2-	3.98m	+0.042(17)	+0.74(11)	-1.002(7)	[1]
158m	(5^+)	≈20s	?	?	?	-
157	$1/2^{+}$	3.63m	+0.475(15)	-	-1.093(8)	[1]
156	2-	83.8s	?	?	?	-
155	$11/2^{-}$	21.6s	?	?	?	-
155m	$1/2^+$	45s	?	?	?	-
154	(2)	8.1s	-1.14(2)	+0.4(9)	-1.486(19)	[2]
154m	9+	3.30s	+5.91(5)	-0.2(4)	-1.522(15)	[2]
153	$(11/2^{-})$	1.48s	+6.93(11)	+0.5(10)	-1.615(31)	[2]
153m	$(1/2^+)$	2.5s	?	?	?	-
152	(2-)	88	?	?	?	-
152m	(9^+)	5.2s	?	?	?	-
151	$(11/2^{-})$	4.17s	?	?	?	-
$151 \mathrm{m}$	$(1/2^+)$	6.6s	?	?	?	-
150	(6 ⁻)	2.20s	?	?	?	-
149	$(11/2^{-})$	0.9s	?	?	?	-
148	(10^+)	0.7s	?	?	?	-
147	$11/2^{-1}$	0.58s	?	?	?	-

Previous laser spectroscopy performed on isotopes ¹⁵³Tm - ¹⁷²Tm in the work by G. D. Alkhazov *et al* [1] and A. E. Barzakh et al [2].

Previous and current work



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165	$1/2^{+}$	30.06h	-0.139(2)	-	-0.250(2)	[1]
164	1+	2.0m	+2.37(3)	+0.706(51)	-0.347(6)	[1]
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The eventual goal is to reach the heavily neutron deficient side of thulium to reach ¹⁴⁷Tm. This is because ¹⁴⁷Tm is a proton emitter with a 15% proton emission branch.

In august, 3 days of beam time were allocated to primarily test the yields of neutron deficient thulium and to find a suitable spectroscopic transition with sensitivity to the nuclear properties



New data



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148	(10^+)	0.7s	?	?	?	-
147	$11/2^{-1}$	0.58s	?	?	?	-

The transition of the ionic 313nm J=4 \rightarrow J=5 line from the ground state was tested and was very promising.

During these 3 days, despite various issues, lots of isotopes were measured with ¹⁵⁵Tm, ¹⁵⁶Tm, ¹⁷²Tm, ¹⁷³Tm, ¹⁷⁴Tm, ¹⁷⁵Tm being measured, via laser spectroscopy, for the first time. Several isomeric states were measured for the first time as well.

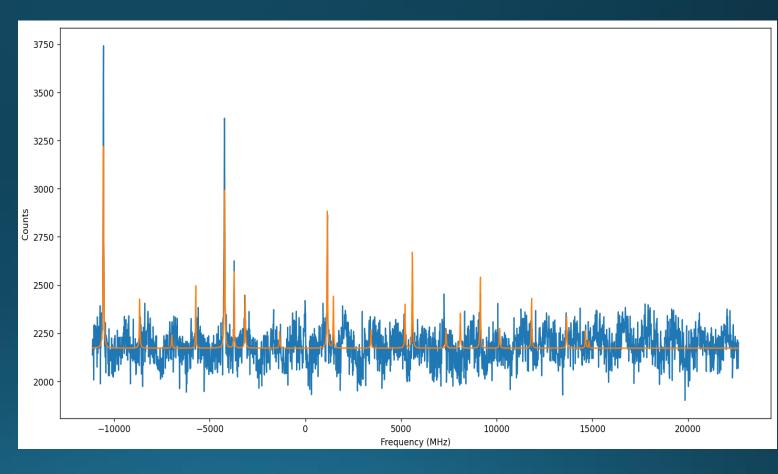


Issues in ¹⁵⁵Tm



¹⁵⁵Tm was the limit of the study, a poor signal to noise ratio was caused by isobaric contamination in the beam.

Despite this, there are clearly enough peaks visible to be able to fit and analyse the data.





Isomer in ¹⁶¹Tm

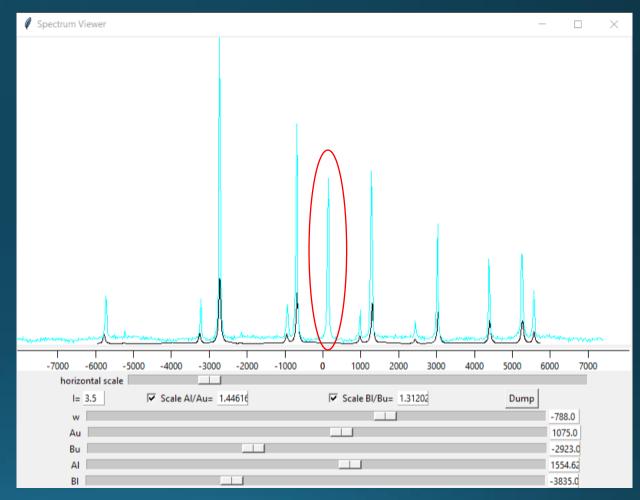


¹⁶¹Tm was more complex due to a single peak that could not be accounted for by the ground state alone.

There is no listed isomeric state on Nudat3 or the IAEA nuclear moments.

As it is an odd-even nucleus, an isomeric state spin cannot be an integer, and so a spin of 1/2 is the lowest possible spin state for this peak.

Having only the one peak, however, suggests that the isomeric peaks have collapsed into this single peak.



Initial estimate of the ground state ¹⁶¹Tm (black) against converted data from run 63 (blue). The single unaccounted peak has been circled in red

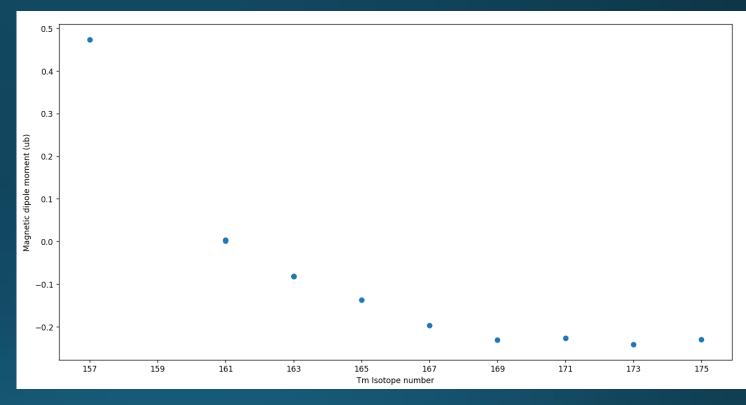


Isomer in ¹⁶¹Tm



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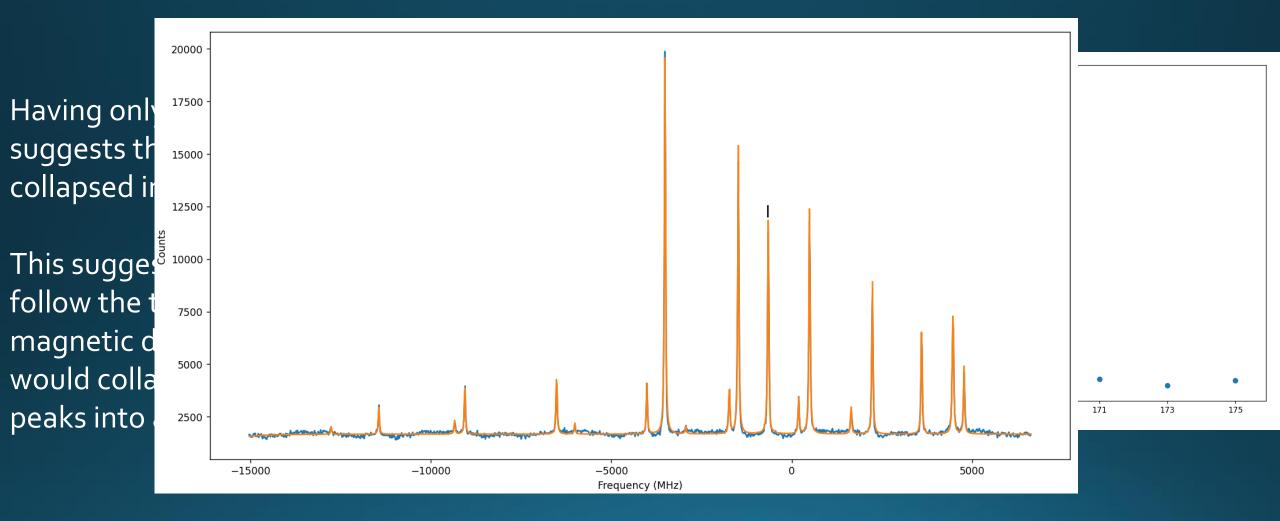
This suggestion is supported as it would follow the trend seen here, placing the magnetic dipole moment at 0, which would collapse the isomeric spin 1/2 peaks into a single peak





Isomer in ¹⁶¹Tm

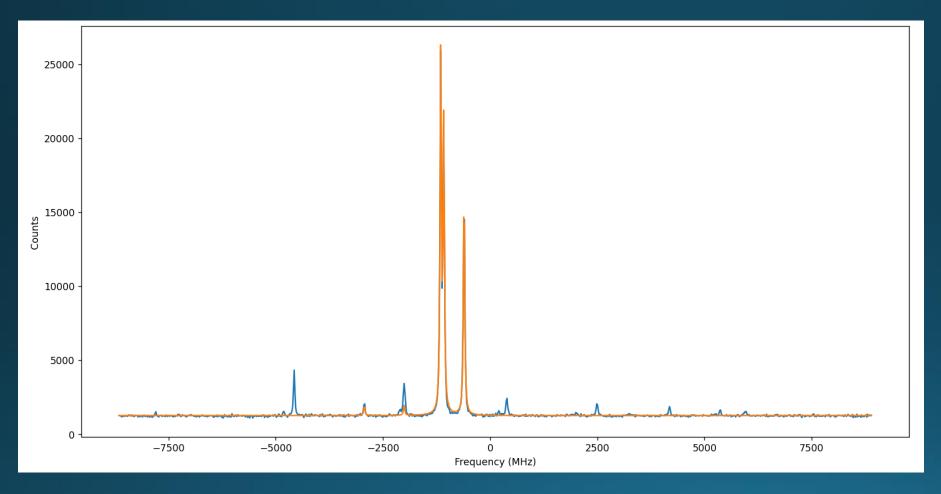






Isomer spin in ¹⁶⁰Tm





The ground state of ¹⁶⁰Tm was simple to fit, as all peaks could be accounted for.

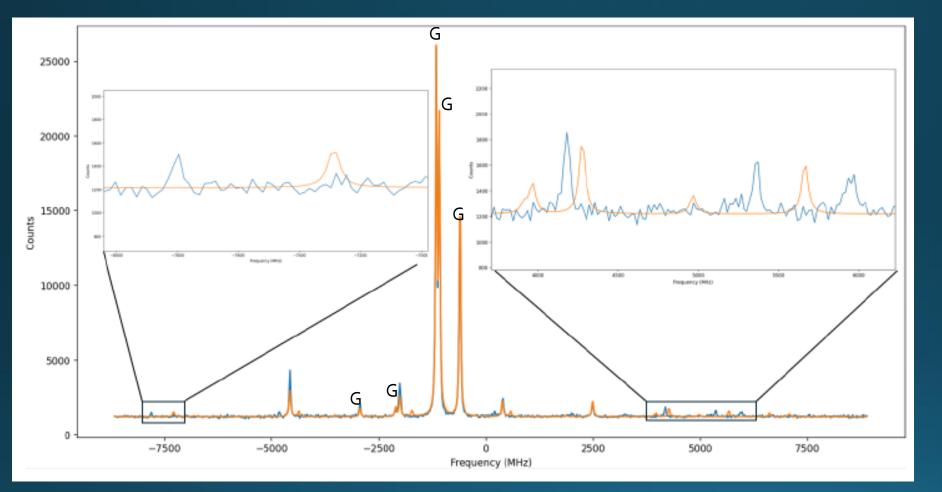
However, the Isomeric state was significantly harder to find an initial estimate for

Fit of ground state peaks (orange) against data (blue) for ¹⁶⁰Tm.

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Isomer spin in ¹⁶⁰Tm





Fit of ground and isomeric state peaks (orange) against data (blue) with an isomeric spin assignment of 5

Even with a good estimate, multiple discrepancies can be found.

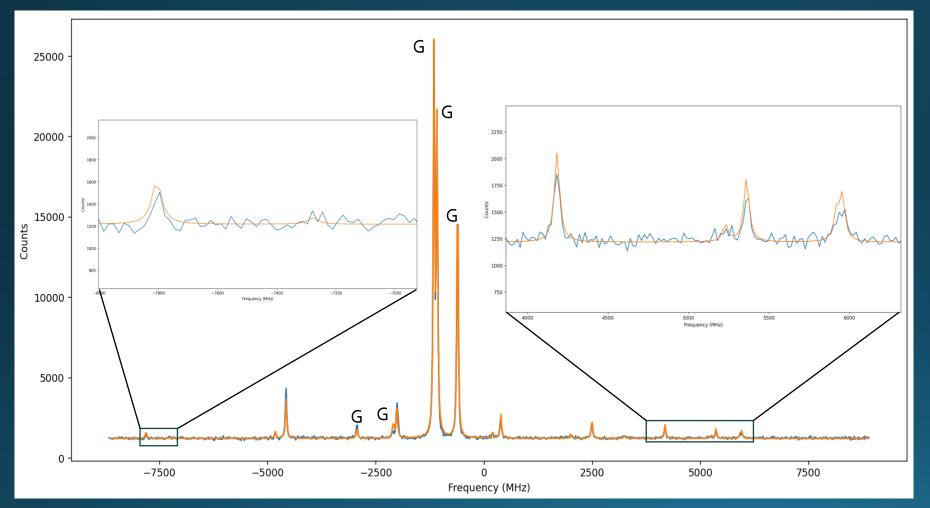
This led to some questions regarding the spin assignment of the isomeric state of ¹⁶⁰Tm

As such, different spins were investigated to see if other spins would produce a better fit



Isomer spin in ¹⁶⁰Tm





As such, different spins were investigated to see if other spins would produce a better fit

An isomeric state with a spin of 4 can be seen to be much more agreeable.

Fit of ground and isomeric state peaks (orange) against data (blue) with an isomeric spin assignment of 4



Results



\mathbf{Tm}	Ι	μ (μ_b)	Q_s (b)	$\delta \langle r^2 \rangle ~(\text{fm}^2)$
155	5.5	6.740(44)	-1.050(30)	-1341(28)
155m	0.5	-0.653(13)	-	-2.030(37)
156	2	-0.524(4)	-0.1691(85)	-1.269(26)
157	0.5	0.4735(31)	-	-1.100(23)
158	2	-0.079(1)	-0.400(13)	-1.006(21)
159	2.5	3.400(20)	1.865(50)	-0.842(19)
160	1	0.150(10)	0.602(16)	-0.744(17)
$160 \mathrm{m}$	4	3.005(20)	2.855(79)	-0.774(17)
161	3.5	2.392(16)	2.870(78)	-0.628(14)
161m	0.5	2.392(16)	-	-0.585(14)
162	1	0.0659(5)	0.709(19)	-0.539(13)
163	0.5	-0.082(1)	-	-0.403(10)
164	1	2.337(15)	0.473(17)	-0.3450(86)
164m	6	2.960(19)	4.36(12)	-0.4046(90)
165	0.5	-0.1377(9)	-	-0.2581(67)
166	2	0.0898(6)	2.139(58)	-0.2089(51)
167	0.5	-0.1964(13)	-	-0.1254(51)
168	3	0.2239(15)	3.262(88)	-0.0852(18)
169	0.5	-0.231(1)	-	0
170	1	0.2457(16)	0.74(2)	0.0443(16)
171	0.5	-0.2265(15)	-	0.1264(48)
172	2	-0.11204(81)	2.305(62)	0.1678(48)
173	0.5	-0.2416(16)	-	0.2220(64)
174	4	0.3191(25)	3.83(10)	0.2518(78)
175	0.5	-0.2303(25)	-	0.3117(93)

New nuclear moments for various isotopes and isomers across the thulium chain (bold).

Very strong evidence for a spin reassignment, as no other spin value can reproduce the frequency pattern of all the peaks simultaneously.

Very strong evidence for new isomeric state(s) as well.

Still a long way to go to get to ¹⁴⁷Tm.





Thank you!

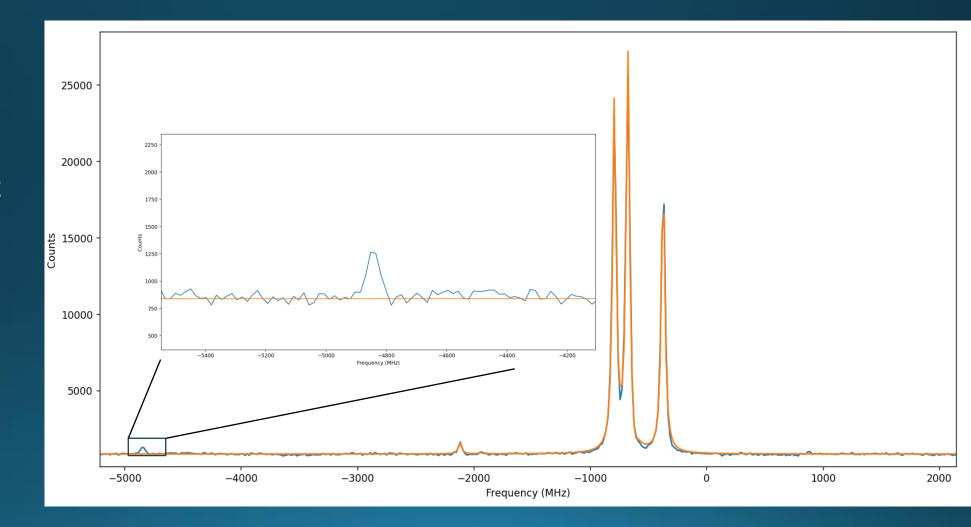
Presentation by Jack Hughes sgjhugh3@liverpool.ac.uk





¹⁶²Tm ground state fit

Small peaks above background at -5000 and 1000 MHz

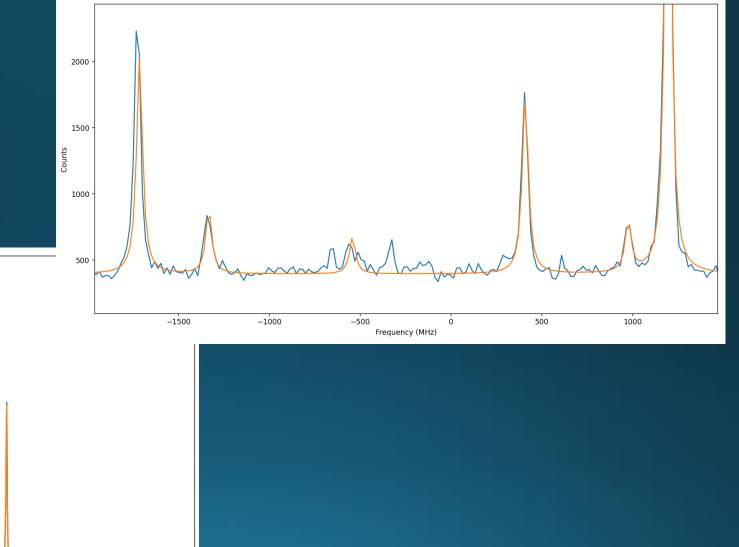


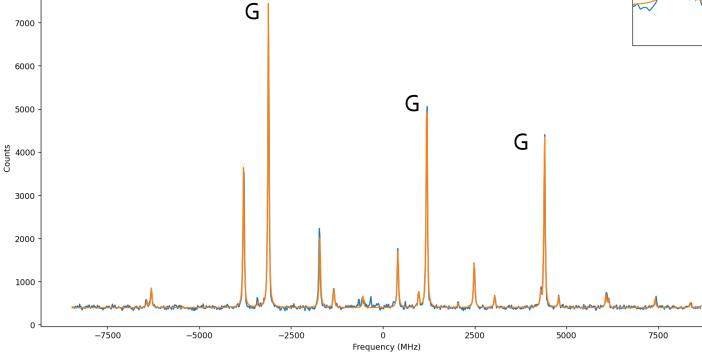




¹⁶⁴Tm ground and isomer fit

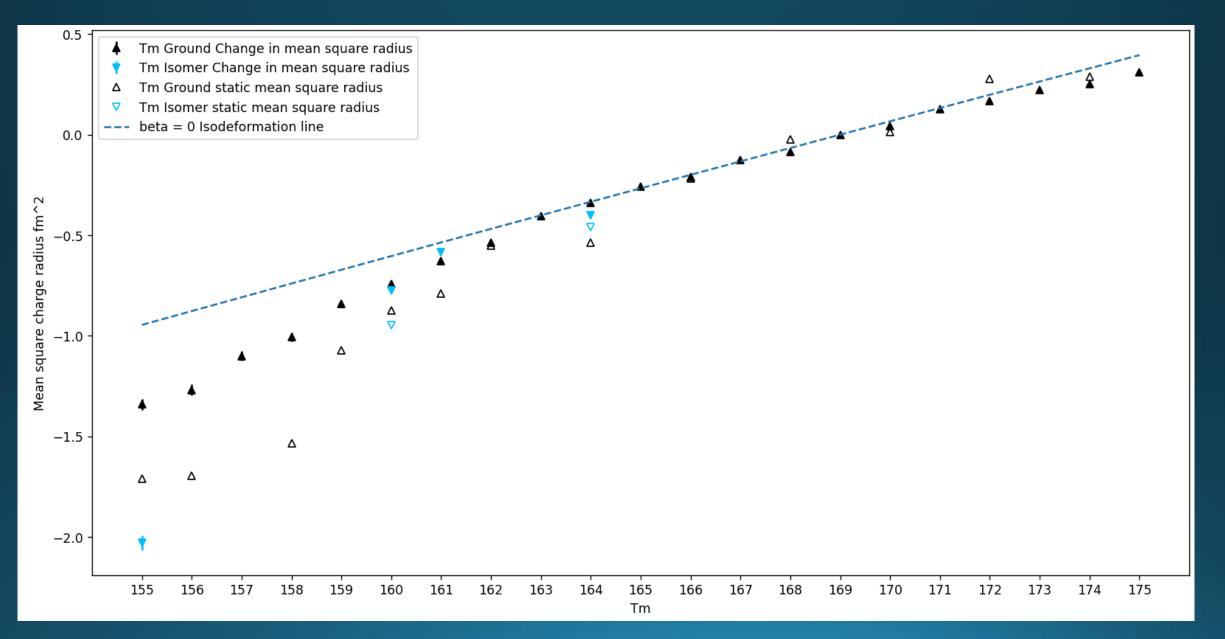
Something weird at approx. -500 MHz





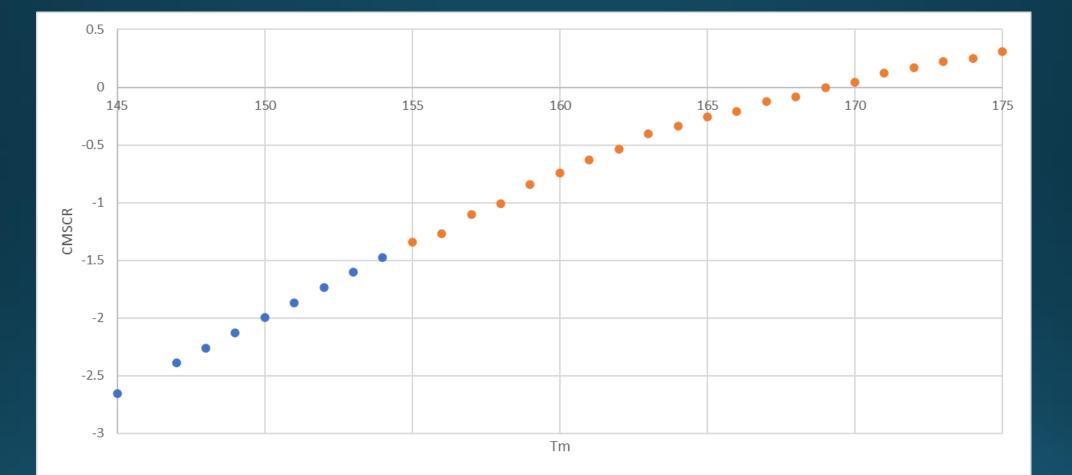














¹⁵⁵Tm raw data

Background increase at high line voltage in PMT 1

