

# EFT Workshop at Notre Dame

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## Abstract

The LPC EFT workshop was held April 25-26, 2024 at the University of Notre Dame. The workshop was organized into five thematic sessions: *how far beyond linear* discusses issues of truncation and validity in interpretation of results with an eye towards practicality; *reconstruction-level results* visits the question of how best to design analyses directly targeting inference of EFT parameters; *logistics of combining likelihoods* addresses the challenges of bringing a diverse array of measurements into a cohesive whole; *unfolded results* tackles the question of designing fiducial measurements for later use in EFT interpretations, and the benefits and limitations of unfolding; and *building a sample library* addresses how best to generate simulation samples for use in data analysis. This document serves as a summary of presentations, subsequent discussions, and actionable items identified over the course of the workshop.



# 1 Introduction

The LPC<sup>1</sup> Effective Field Theory (EFT) workshop was held on April 25-26, 2024 at the University of Notre Dame<sup>2</sup>. It was the second in a series started at the LPC in September 2023, with a primary goal of growing the community of physicists interested in EFT interpretations of measurements and searches at the CERN Large Hadron Collider (LHC). The workshop included ample time for discussion centered around resolving outstanding practical questions towards performing large-scale EFT combinations within an experiment. The attendees were primarily experimentalists on the CMS experiment and theorists working on EFT-related topics. The desire is for the series to grow in breadth of attendees' backgrounds. Prior to the workshop, an introductory EFT tutorial and hackathon was held. The tutorial introduced EFT concepts and practical tools for new researchers, and the hackathon allowed attendees to work in a self-organized fashion on specific topics, such as generation of a particular sample of events.

This document serves to record what was discussed and what actionable items were identified over the course of the workshop. Speakers were invited to submit a short paragraph summarizing the most critical takeaway messages from their talk to be conveyed to a broader community, and the session conveners were charged with placing the summaries in context by capturing questions and comments raised in subsequent discussions. The sessions were organized into five themes which form the sections of this document: Section 2 discusses issues of truncation and validity in interpretation of results with an eye towards practicality; Section 3 visits the question of how best to design analyses directly targeting inference of EFT parameters; Section 4 addresses the challenges of bringing a diverse array of measurements into a cohesive whole; Section 5 tackles the question of designing fiducial measurements for later use in EFT interpretations, and the benefits and limitations of unfolding; and Section 6 addresses how best to generate simulation samples for use in data analysis.

## 2 How far beyond linear?

The goal of this session was to explore the considerations that should be kept in mind when choosing which orders of EFT operators to include in the modeling of EFT predictions for experimental analyses and to discuss the consequences that this choice has on interpretations. While experimental analyses often include just the linear and dimension-six squared pieces (for practical reasons), this is incomplete. Other contributions (e.g., interference between the SM and dimension-eight contributions, as well as interference between the SM and diagrams with double-insertions of dimension-six vertices) enter at the same  $1/\Lambda^4$  order as the dimension-six squared contributions. It is thus important to consider these contributions and work towards more proper methods of handling their effects.

Three main aspects of this topic were explored. Section 2.1 presents examples of cases where new physics could enter in a way that is not modeled by linear SMEFT contributions. Next, Section 2.2 describes an approach (known as “geoSMEFT”) for identifying and organizing the most relevant higher-order contributions. Section 2.3 discusses how these higher order contributions can be seen as an uncertainty on the linear piece. Finally, Section 2.4 summarizes the discussion.

### 2.1 SMEFT vs HEFT: When EFT for new physics is not linear

*Speaker:* Duarte Fontes

The Higgs Effective Field Theory (HEFT) can be used at the LHC and future colliders to parametrize possible deviations from the Standard Model. Eventual nonzero HEFT coefficients should then be converted into coefficients of specific UV models via a matching procedure. However, this procedure is not unambiguous. In fact, depending on the way in which the different parameters of a certain UV model are assumed to scale, different HEFT expansions (and hence matching relations) are obtained. This happens in such a way that, according to the process and region of the parameter space of the UV model, different matchings should be used to ensure a fast convergence of the HEFT expansion to the results of the UV model. This complicates the interpretation of HEFT measurements in terms of parameters of UV models.

### 2.2 Which orders: the geoSMEFT perspective

*Speaker:* Adam Martin

Using the geoSMEFT reorganization, many processes—including those that are central to the global SMEFT fit program such as the EPWO,  $h \rightarrow \gamma\gamma$ , and diboson processes—are fully calculable to  $O(1/\Lambda^4)$ . Using these

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<sup>2</sup> Agenda accessible at <https://indico.cern.ch/e/lpceft2024>

48 full higher order SMEFT calculations (which, should supersede  $|dim6|^2$  approximate calculations whenever they are  
 49 known), one can study truncation uncertainty. Dimension-8 operators are usually unconstrained by EWPO and can  
 50 contain different helicity/polarization structure, allowing them to interfere with the SM in scenarios where dimension-  
 51 6 operators cannot. As a result, they can lead to dramatic differences between the full  $O(1/\Lambda^4)$  SMEFT result and the  
 52  $|dim6|^2$  approximation. The geoSMEFT organization also simplifies the energy vs.  $vev^3$  scaling of higher dimensional  
 53 operators, as all new momentum dependence (novel energy enhanced vertices) are shuffled into process-specific 4+  
 54 particle vertices.

## 55 2.3 Truncation and validity: Treating higher order terms as uncertainties

56 *Speaker:* William Shepherd

57 Exploring specific new physics models and their realizations in EFT frameworks is an interesting exercise, but  
 58 it doesn't really reflect the state of our current experimental knowledge. A more honest approach to EFT searches  
 59 must remain model-agnostic, positing only that the SM as we've measured it is the low-energy approximation of some  
 60 higher-energy theory. This agnosticism suggests that we should not privilege some Wilson Coefficients over others or  
 61 assume that our framework has been configured such that the interesting effects of this new physics will generically  
 62 be captured by searches that specifically look for only a small subset of operators, but instead we must be aware that  
 63 all effects at a given order in the EFT perturbation theory deserve to be treated with the same seriousness.

64 Taking this perturbation theory seriously necessitates acknowledgement of the presence of higher order calculations  
 65 than those we can perform, and the treatment of some approximate of them as an error on our understanding of the  
 66 prediction of the EFT. Without this, we cannot be confident in any constraint we claim on EFT parameters. Luckily,  
 67 much like QCD offers us the scale dependence as a way to estimate the size of higher-order effects, the EFT also  
 68 gives us an easy-to-calculate quantity that we can use to get a generic understanding of the size of next-order effects.  
 69 In an analysis using only dimension-6 operators, this is the square of the amplitude term at order  $\frac{1}{\Lambda^2}$ . Using this,  
 70 and scaling to avoid accidental overconfidence (e.g. preventing strong would-be limits at dimension-6 from requiring  
 71 dimension-8 operators to be small) it is possible to honestly estimate the size of these uncertainties, and yield bounds  
 72 which are solid constraints on any future model of heavy new physics.

73 Further details on these ideas can be found in [1].

## 74 2.4 Discussion

75 It is evident that higher-order contributions should not be disregarded; however, in many cases, it is not yet practical  
 76 to fully account for dimension-eight effects. While there is not yet an overarching solution that would be feasible to  
 77 apply in all cases, we summarize the observations and recommendations that resulted from the discussion:

- 78 • The majority of experimental analyses currently include the linear and quadratic dimension-six pieces, without  
 79 higher order contributions. This approach tends to be utilized because of its practicality.
- 80 • If there are no other feasible alternatives available, these types of analyses are worthwhile; they provide a  
 81 method of testing the SM against an alternative hypothesis, so they do have discovery potential. However,  
 82 analyzers should be aware that the theoretical interpretation of these types of results are limited (not only in  
 83 the case of a possible sign of new physics, but also for confidence intervals placed on Wilson coefficients in the  
 84 absence of a signal).
- 85 • In the absence of a feasible method of properly including all  $1/\Lambda^4$  effects, it would be beneficial for analyses to  
 86 report linear-only results<sup>4</sup> in addition to the quadratic results reported. Many analyses already do this, but it  
 87 would be beneficial for this to become more standard.
- 88 • When the full  $1/\Lambda^4$  effects cannot be included easily, it is possible to treat the dimension-six quadratic piece  
 89 as an uncertainty to the linear term. While it may be challenging to incorporate this rigorously into a fit in an  
 90 experimental analysis (as the uncertainty would depend on the Wilson coefficient values, and would need to  
 91 be handled for each bin) it would be beneficial to invest effort in this area. Results which incorporate such an  
 92 uncertainty would be more theoretically meaningful than results which aggressively include linear and quadratic  
 93 dimension-six pieces (without any estimation of uncertainty from the missing  $1/\Lambda^4$  terms or the higher-order  
 94 corrections). Presented together, these conservative and aggressive limits together would help to provide a  
 95 more useful and interpretable picture of the Wilson coefficient space to be explored.

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<sup>3</sup> Higgs field vacuum expectation value

<sup>4</sup> It should be noted that there exist cases where a linear term is not present (when the EFT contribution does not interfere with the standard model).

- The geoSMEFT method of organizing dimension-eight contributions provides a promising method to properly handle the most important  $1/\Lambda^4$  effects. It would be beneficial for experimental analyses to explore this method, and to use it to begin incorporating proper  $1/\Lambda^4$  effects as soon as possible.

The points summarized here do not imply that aggressive analyses (which include linear and quadratic dimension-six contributions) should be abandoned; the discovery potential of these approaches is acknowledged, though the limited interpretability of these approaches is also recognized. The point to be emphasized is that the community should not become complacent in this limited approach. As improved ideas and techniques (such as geoSMEFT) become available, the community should strive to incorporate these into experimental analyses. Regardless of whether future EFT analyses will discover hints of new physics or if limits remain consistent with the standard model prediction, improvements in EFT modeling will strongly benefit the field by improving the interpretability and meaningfulness of LHC EFT results.

### 3 Reconstruction-level results

This session focused on the challenges and opportunities posed by performing a search for EFT effects at the reconstruction level (equivalently, detector level). Ideally there would be no loss of information going from truth level to detector level measurements, but to account for this, acceptance and efficiency are taken into account in such measurements but the implementation can be challenging. From designing and picking the most suitable observable for our analyses, to thinking about ways to better utilize and understand existing measurements and statistical methods to derive the results, there is a lot to consider in reconstruction level EFT measurements. To guide our discussion we will be focusing on the following

- Re-interpretability of detector-level measurements Section 3.1
- Entanglement and Bell's inequalities with boosted top quarks Section 3.2
- Observables for EFT measurements Section 3.3
- Simulation Based Inference Section 3.4

#### 3.1 Re-interpretability of detector-level measurements

*Speaker:* Sergio Sánchez Cruz

While measurements in dedicated regions of interest with reconstruction level observables have the advantage of reflecting detector effects and portraying a more accurate picture, reinterpreting the results can be challenging. Once a measurement is complete we often lose access to the full statistical model used in the analysis, and any statistical combination becomes impossible. Any changes (like testing the effect of a different EFT operator) or additions to the measurement would also require rerunning the whole analyses, which is not feasible in most cases. The problems posed can be solved if as a collaboration we make it a habit to release the full statistical model used in the analysis, which will make it easier to rerun the analyses or make statistical combinations. Using certain statistical techniques to conduct a post-generation reweighting of simulated samples would also help to study any new operator of interest.

#### 3.2 Entanglement and Bell's inequalities with boosted top quarks

*Speaker:* Dorival Goncalves

The LHC offers a unique opportunity to explore quantum correlations, such as entanglement and Bell inequality violation, at the highest energy scale available today. We discuss these quantum correlations using top quark pair production as a model for a two-qubit system, specifically focusing on the semi-leptonic top pair channel, which provides a sixfold increase in statistics and easier reconstruction compared to the dileptonic channel, which was previously used by ATLAS and CMS. While measuring the spin polarization of the hadronic top quark presents challenges, our study demonstrates the feasibility of reconstructing the spin density matrix of the two-qubit system using an optimized hadronic polarimeter. This involves employing jet substructure techniques and reconstruction methods inspired by neural networks to enhance the mapping between subjects and quarks. Our analysis reveals that entanglement can already be observed at a significance level exceeding  $5\sigma$  with existing data using this channel. Moreover, the violation of Bell inequalities may be probed at a significance level surpassing  $4\sigma$  at the HL-LHC with  $3 \text{ ab}^{-1}$  of data. Hence, the analysis of this channel can represent the next crucial step for ATLAS and CMS in their exploration of quantum correlations at the LHC.

### 143 3.3 Observables for EFT Measurements

144 *Speaker:* Andrei Gritsan

145 A crucial aspect when devising a detector-level EFT analysis is the proper definition of a set of quantities capable of  
146 capturing the peculiarities of the EFT contributions under investigation. The types of suitable observables range from  
147 the usual quantities calculated for the SM measurements to the EFT-sensitive, as well as to the optimized observables  
148 utilizing advanced matrix-element and machine-learning techniques. When building EFT-sensitive observables, it is  
149 worth considering that higher-dimension operators typically lead to enhancement at the higher values of four-momenta  
150 squared distributions of the particles appearing in the propagators. Therefore, observables based on such calculations  
151 or correlated with those quantities become sensitive probes of deviations from the SM. An example of such an  
152 observable could be the transverse momentum of reconstructed objects. At the same time, such generic probes may  
153 not be sensitive to distinguish multiple operators that all lead to the same enhancement at higher momenta. One  
154 example of such a situation is the study of CP-even and CP-odd operators, which may require special CP-sensitive  
155 observables to differentiate them. Moving to more complex variables, while the matrix-element calculations guarantee  
156 optimal performance from first principles, there are practical limitations to their applications. The most critical  
157 limitations are the transfer functions, which are difficult and time-consuming to model. Parton shower and detector  
158 effects may confuse and distort the input to matrix elements to such a degree that calculations become impractical.  
159 Machine Learning (ML) techniques may come to the rescue in such a case. Training of machine-learning algorithms is  
160 still based on MC samples utilizing the same matrix elements that would be used for optimal discriminants. However,  
161 these MC samples reflect the parton shower and detector effects and therefore allow the construction of optimized  
162 observables that incorporate these effects. When it comes to performing the training of such an ML model, there are  
163 two important aspects to take into account, which are selecting the observables to inject into the learning process  
164 and which samples should be used. The matrix-element approach provides answers to both questions, together with  
165 an insight into the process of constructing the optimized observables with ML. The first thing to consider is that the  
166 input observables should provide full information, which could be simply the four vectors of all particles involved,  
167 like in the matrix elements, or better-derived physics quantities that are equivalent to those. In addition, the optimal  
168 observables, like in the case of matrix-element-based calculations, can be separated into two categories, which are the  
169 set of quantities sensitive to quadratic terms in the amplitude, and the set of variables sensitive to interference terms.  
170 The first corresponds to the classic problem of differentiating between two models, and an ML algorithm is trained  
171 on two samples corresponding to two alternative models. Training an ML equivalent of the second type of matrix-  
172 element-based optimal observables is less obvious, as it requires isolating the interference component. A discriminant  
173 trained to differentiate the two models with maximal quantum-mechanical mixing is the best candidate for that goal,  
174 where the quadratic term may or may not be removed. Using the two observables jointly would guarantee optimal  
175 performance for any size of the contributing operator.

### 176 3.4 Simulation-based inference

177 *Speaker:* Harrison Prosper

178 Simulation-based inference (SBI) is in a mature state and high-quality tools exist to implement SBI, notably,  
179 the toolkit Madminer 4.2. While it is always helpful to implement SBI oneself, it is worth checking if an approach  
180 you wish to use is not already available. Given the mature state of SBI, there is an actual chance to move away  
181 from the traditional ML benchmark approach to testing new physics models. It is now computationally feasible to  
182 simulate events at a large number of parameter points and either directly construct inference summaries such as point  
183 estimates and confidence sets and intervals from simulations or approximate the statistical model  $p(X|\theta)$  directly,  
184 where  $X$  are potential observations and  $\theta$  the parameter space of the theoretical model of reference. It is important  
185 to assess the accuracy of the approximations, but if  $p(X|\theta)$  can be accurately modeled this has the virtue that all of  
186 the standard machinery of frequentist and Bayesian statistical inference can be deployed on a neural network model  
187 of  $p(X|\theta)$ .

### 188 3.5 Discussion

189 When considering observables for EFT measurements, it's crucial to prioritize selecting those that offer the most  
190 meaningful insights into the underlying physics. One recommendation is to incorporate a diverse range of observables  
191 that capture different aspects of the phenomenon under investigation. This ensures a comprehensive understanding  
192 of the system and increases the likelihood of detecting subtle effects indicative of new physics. In terms of tool  
193 selection, it's advisable to leverage versatile tools like MELA, which have evolved to accommodate various generators  
194 such as JHUGen, MCFM, and now Madgraph. This adaptability allows for more robust analyses and enhances the

195 compatibility of results across different platforms. When confronted with a multitude of observables, it's essential  
196 to exercise discretion in their selection. Instead of overwhelming analyses with a plethora of variables, focus on  
197 identifying a few key observables that are most sensitive to the parameters of interest. This targeted approach  
198 not only streamlines the analysis process but also facilitates a clearer interpretation of results. Furthermore, when  
199 constructing likelihood ratios, it's prudent to consider the Neyman-Pearson lemma carefully. While this lemma  
200 provides a framework for hypothesis testing, its applicability in multidimensional analyses may vary. It's advisable to  
201 assess the suitability of the hypothesis for the specific context of the analysis and to explore alternative approaches  
202 if necessary.

203 Navigating simulation-based inference techniques involves discerning when to employ matrix-element-driven op-  
204 timal observables, machine learning, or neither, depending on the analytical context.

- 205 • Prioritize analytical approximation of likelihood whenever possible. This approach fosters a sturdy statistical  
206 model and promotes the sharing of likelihoods among researchers. In cases where modeling the likelihood  
207 proves impracticable, consider the merits of unbinned analysis techniques. This method enables the direct  
208 computation of the probability density function or cumulative distribution function (CDF) of the targeted test  
209 statistic, furnishing valuable insights into the underlying statistical model.
- 210 • Tackle nuisance parameters by judiciously assigning believable priors and relaxing specific value requirements  
211 to facilitate averaging over the priors. However, it's vital to validate the efficacy of this strategy, particularly  
212 in higher-dimensional analyses where coverage may vary.
- 213 • Endeavor to release a CDF of the test statistic distribution, empowering further analysis and interpretation.  
214 However, determining the optimal approach between directly approximating the unbinned statistical model or  
215 compressing data into a test statistic remains an area requiring additional exploration.
- 216 • Although feasible, combining unbinned and binned analyses necessitates accurate consideration of dependen-  
217 cies between the two statistical models. This requires using compatible parameter spaces to ensure reliable  
218 outcomes.
- 219 • Recognize the considerable computational demands associated with simulation-based inference techniques,  
220 particularly when producing robust confidence sets. Strategically allocate resources and gradually populate  
221 parameter spaces near boundaries to minimize computational expenditures.
- 222 • Focus on identifying independent directions when exploring EFT parameter spaces. This targeted approach  
223 enables effective examination and learning along specific dimensions, enhancing analytical efficiency.

224 In the discussion about the re-interpretability of detector-level measurements, a pertinent question arises concern-  
225 ing the post-analysis re-visitation of full simulations. This practice, once completed, presents a formidable task, given  
226 the intricacies involved in re-calibration and re-evaluation. Such challenges underscore not only technical hurdles  
227 but also a broader “sociological problem,” wherein entrenched methodologies hinder the exploration of alternative  
228 approaches. However, within this landscape of challenges, avenues for progress emerge. Consider the proposition  
229 of presenting a comprehensive likelihood for theoretical scrutiny. While the magnitude of such an undertaking may  
230 appear daunting, precedent exists in the form of the published likelihoods in the context of Higgs studies, serving as  
231 a testament to the feasibility and benefits of such initiatives. Armed with this foundational framework, theorists gain  
232 access to a wealth of insights, fostering deeper exploration and collaborative endeavors. Amidst these discussions,  
233 questions arise regarding disparities observed between helicity-ignorant and helicity-aware reweighting techniques to  
234 implement a broad set of EFT scenarios within the same simulated samples. The integration of degrees of freedom  
235 and the preservation of spin correlation emerge as pivotal considerations in that respect. Delving deeper, nuances  
236 about the consideration of operators and the exploration of phase space come to the fore. A cautious approach is  
237 recommended, as drawing conclusive inferences necessitates a thorough understanding of underlying processes and  
238 methodological intricacies.

239 In the discourse surrounding top entanglement and Bell inequalities, several key themes emerge, each offering  
240 insights and avenues for further exploration. Firstly, the observation of improved agreement between CMS measure-  
241 ments of top-entanglement and theoretical predictions underscores the importance of refining modeling techniques,  
242 particularly in capturing spin correlations with precision. This highlights the need for ongoing efforts to enhance  
243 modeling methodologies to better align experimental observations with theoretical expectations. Concerns regarding  
244 systematic uncertainties prompt reflection on the choice of models for interpreting data, particularly in cases of dis-  
245 agreement between observed data and theoretical predictions. The imperative to ensure robustness and reliability in  
246 model selection underscores the critical role of methodological rigor in advancing our understanding of particle interac-  
247 tions. Considerations surrounding loopholes in Bell inequalities measurements shed light on the inherent challenges in  
248 experimental design and interpretation, though collider probes of entanglement can not be used for loophole-free tests

249 of Bell’s inequality. Explorations into the interplay between EFT contributions and top entanglement properties offer  
250 fertile ground for future research. Efforts to better exploit EFT frameworks in analyzing angular distributions and  
251 spin matrices hold promise for uncovering deeper insights into the underlying physics governing particle interactions.

## 252 **4 Logistics of combining likelihoods**

253 One of the central promises of EFT is that it provides a common language for describing physics beyond the standard  
254 model across different measurements and experiments. This common language facilitates global statistical combi-  
255 nations, which exploit the capability of different analyses to constrain orthogonal directions in Wilson coefficient  
256 space.

257 Statistical combinations in high energy physics have a long history, especially with regards to precision electroweak  
258 fits, the discovery of single top quark production, and the discovery of the Higgs boson. These combinations led to  
259 tremendous technical development of combination tools and procedures, but EFT combinations present some unique  
260 challenges, which will require additional development. It is crucial that these tools be specified in a human-readable  
261 fashion, so that they can be used far into the future even though programming languages and execution environments  
262 will change. It is also very important that the digital artifacts that become inputs to a combination have enough  
263 flexibility to permit changes in the underlying theoretical framework, for example the introduction of higher-dimension  
264 EFT operators or more precise calculation tools.

265 The time is ripe for this technical development, which will enable us to establish not only useful combination  
266 tools and procedures for EFT, but also to define the legacy of the HL-LHC and of future Higgs factories.

### 267 **4.1 Grand combinations at LEP and Tevatron**

268 *Speaker:* Tom Junk

269 Searches for new particles and interactions grew in complexity in the 1990s at LEP, with the use of neural  
270 networks and other multivariate analysis techniques. Statistical methods and software kept pace, being able to use  
271 distributions of discriminant variables for each event optimally, rather than just cutting and counting. Multivariate  
272 analysis techniques however are typically optimized to find a specific signal with an a priori model of the background.  
273 Searching for more general signals requires care to make sure the efficiencies stay high and are well understood, and  
274 that separation from backgrounds remains close to optimal.

275 The combinations of multivariate Higgs boson search results within each of the four LEP collaborations, ALEPH,  
276 DELPHI, L3, and OPAL, were necessary in order to produce publishable results. Combinations of the four collabo-  
277 rations’ results provided the most sensitivity to the sought-after signals. Preparing digital artifacts for inputs to the  
278 combination was organized by a working group that standardized the formats and set deadlines. The format was  
279 designed so that choices that had an impact on the combined results, except the choice of the final statistical method,  
280 were made by the experimental collaborations and not by the combiners. Issues of binning and interpolation methods  
281 were handled by exchanging Fortran code which could read in external data files as needed. Systematic uncertainties  
282 were provided by named source, where sources with the same name were considered 100% correlated and sources  
283 with different names were considered uncorrelated. Arbitrarily-correlated uncertainties could be decomposed into  
284 correlated and uncorrelated pieces.

285 Combinations at the Tevatron increased in complexity, mainly due to the need to handle larger systematic  
286 uncertainties and the shift to ROOT and C++. Single Top and Higgs combinations followed similar patterns to  
287 those used at LEP, with increased attention to shape uncertainties and profiling the likelihood in the calculation  
288 of the test statistic. As was the case with the LEP experiments, each collaboration provided inputs, but also each  
289 collaboration performed all of the combinations, reproducing their own and those of the other collaboration(s).  
290 Results could only be approved if they were in agreement among several combiners. The necessary checking of the  
291 inputs and reproduction of the results was a valuable tool for finding and correcting mistakes in the inputs or the  
292 interpretations of them. Better tools for preparing, exchanging, combining, and preserving results have been created  
293 for use at the LHC, expanding on experience gained at LEP and the Tevatron.

### 294 **4.2 Grand combinations at HL-LHC and FAIROS-HEP**

295 *Speaker:* Kyle Cranmer

296 The gold standard for EFT results should be based on statistical combinations similar to those used for the Higgs  
297 discovery, i.e., statistical models of the observations without unfolding. The infrastructure for these combinations is  
298 mature, but we lack an agreed upon standard for describing how the distribution of a specific process (e.g. ‘signal’)



299 depends on the EFT coefficients. Recent experience has shown that there are significant advantages to separating  
 300 the mathematical specification from the implementation, and that this facilitates efforts to publish these statistical  
 301 models. Such a specification could be a straight-forward extension of the specifications used for binned-template fits  
 302 such as the HistFactory / pyhf specification or what is used in the CMS Combine tool. It would be natural to have  
 303 a specification for both the linear and quadratic EFT expansion strategies. Finally, in order to extend the results  
 304 to new EFT operators, update background modeling, or incorporate the effect of EFT operators on backgrounds,  
 305 a RECAST-like service for EFTs would be valuable. In the case of new EFT operators, this can be achieved very  
 306 efficiently through reweighting, and the service can export a new statistical model.

### 307 4.3 SMEFT for top quarks at LHC and future Higgs factories

308 *Speaker:* Michael Peskin

309 One of the best opportunities for the discovery of Beyond-Standard-Model effects at colliders that we know how to  
 310 build comes in the precision study of the top quark. This motivates a detailed search for nonzero SMEFT coefficients  
 311 for one or more of the dimension 6 SMEFT operators associated with the top quark.

312 Different BSM models treat the top quark very differently. In some models, such as supersymmetry and two-  
 313 Higgs-doublet models, the top quark is an ordinary quark with an order-1 Yukawa coupling. Its form factors are  
 314 affected by electroweak perturbative corrections, of typical size  $(\alpha_w/4\pi)(v^2/M^2)$ , where  $\alpha_w$  is the weak interaction  
 315  $SU(2)$  coupling,  $v = 246$  GeV is the Higgs boson scale, and  $M$  is the mass scale of BSM physics. Other models — in  
 316 particular, models in which the Higgs boson is composite — require a stronger interaction with the top quark in order  
 317 to produce the large top quark mass. In these models, the BSM corrections are order-1 times  $v^2/M^2$ , and often with  
 318 a large numerical coefficient. In such models, the top quark is said to be “partially composite”. The phenomenology  
 319 of the top quark interactions with a composite Higgs sector presents real opportunities for discovery [2]. These ought  
 320 to be taken more seriously by LHC experimenters.

321 Even considering only dimension 6 operators, SMEFT introduces a large number of possible Lagrangian inter-  
 322 actions for the top quark. These divide into three classes: 4-fermion operators solely within the 3rd generation,  
 323 4-fermion operators that link the 3rd generation with lighter generations, and operators that couple 3rd-generation  
 324 fermion bilinears to bosonic operators. In the standard analysis for the LHC [3], the number of these operators  
 325 is  $11 + 14 + 9 = 34$ . The multiplicity of 4-fermion operators comes from the fact that individual operators are  
 326 current-current interactions with specific helicities in the initial and final states, e.g.  $(\bar{q}_L\gamma^\mu q_L)(\bar{t}_R\gamma_\mu t_R)$ , and possibly  
 327 with weak isospin or  $SU(3)$  color currents. At an  $e^+e^-$  Higgs factory operating above the top quark threshold, the  
 328 number of operators is also large [4], seven 4-fermion operators coupling the 3rd generation quarks to electrons, and  
 329 the above 9 fermion-boson operators. It will clearly be important to make joint fits to these two data sets. It should  
 330 be noted that the 4-fermion contact interactions involving first-generation quarks and those involving the electron  
 331 are distinct but should be similar in size. A global fit might include an explicit theoretical model to relate these  
 332 interactions.

333 The most important effects of composite Higgs physics should be from vector resonances coupling to  $t\bar{t}$  and from  
 334 mixing of the top quark with vectorlike top quark partners. The first class of effects should show up in 4-fermion  
 335 contact interactions, in resonant enhancement of the couplings of  $t$  to  $Z$  and  $W$  (associated with SMEFT operators  
 336 such as  $(\Phi^\dagger \overleftrightarrow{D}_\mu \Phi)(\bar{t}_{L,R}\gamma^\mu t_{L,R})$ ), and in resonant enhancement of the top quark Yukawa coupling (associated with  
 337 the SMEFT operator  $|\Phi|^2\bar{t}_L\tilde{\Phi}t_R$ ). The second effect is visible in decrements of the top quark couplings to  $Z$  and  $W$   
 338 and in loop effects decreasing the  $Hgg$  coupling and the overall scale of Higgs boson couplings.

339 The HL-LHC will offer a very large data set to explore for these effects, with approximately 3 billion top quark  
 340 pairs expected. For precision studies, the modeling of these events to extract angular distributions and top quark  
 341 spins will be challenging.

342  $e^+e^-$  colliders offer lower event samples but in a setting that makes it easier to distinguish the effects of the  
 343 various SMEFT operators. In particular, a linear collider can have high electron beam polarization. This makes it  
 344 possible to measure 6 independent cross section observables – the forward, central, and backward cross sections for  
 345 each of two polarization settings. These are expected to be measured at the parts-per-mil level of accuracy. Adding  
 346 positron polarization as in the ILC design enhances this separation and provides an additional tool for background  
 347 reduction [5]. An  $e^+e^-$  collider is particularly sensitive to the  $t\bar{t}Z$  couplings, since this is part of the top quark pair  
 348 production mechanism; see Fig. 1. One difficulty is that, at a fixed energy, the effect of the two  $t\bar{t}Z$  operators are  
 349 degenerate with the effects of 4-fermion operators. Disentangling these operators requires running at two different  
 350 center of mass energies above the top quark threshold (e.g., 550 GeV and 800 GeV).

351 I would like to add one simple comment on the archiving of LHC SMEFT fits. I am one of those conservative  
 352 people who does not believe in fits with dimension 6 operators including both linear and quadratic terms. The

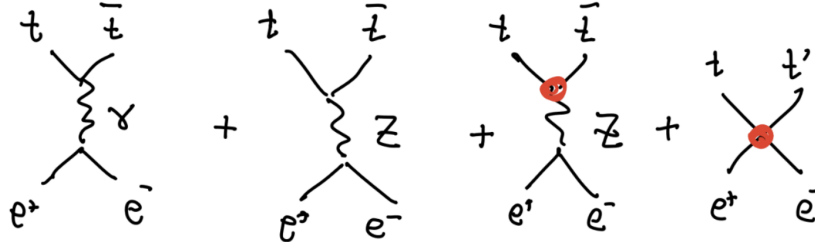


Fig. 1: Diagrams contributing to the cross section for  $e^+e^- \rightarrow t\bar{t}$  at the leading order.

quadratic terms, being positive while dimension 8 operators can give negative terms, lead to constraints on Wilson coefficients that have no physical basis. A better strategy is to fit with linear terms in the dimension 6 Wilson coefficients only and construct the inverse covariance matrix  $(C^{-1})_{ab}$ , where  $a, b$  run over the full set of Warsaw basis operators. This matrix has many zeros and, typically, is not invertible. Truncating this matrix to its positive subspace and then inverting gives the constraints on those operators that the data actually constrains. When additional data sets are added, the matrix  $C^{-1}$  will be invertible over a larger subspace, and so more directions will be constrained. It will be useful to quote two errors on  $C^{-1}$ , a direct experimental error and an error due to truncation of  $1/M^4$  terms (as proposed in Sec. 2.3).

#### 4.4 Discussion

The history of grand combinations in collider physics details:

- A tightly coupled dialogue between theorists and experimentalists that defines mathematical language in which we should summarize tests of wide classes of limits on new physics and corresponding measurements;
- The evolving complexity of statistical inference computational techniques to address experimental uncertainties in requisite detail; and
- The accessibility and ease of interpretation of these analysis data products for theorists, experimentalists, and scientists in other sub-fields of particle physics.

In this session we discussed the past, present, and future of these aspects from the perspective of both physics, software, and even sociology.

Beginning with the practical matter of exchanging and combining physics results, all three presentations provided unique insight on what needs to be achieved in a multifaceted analysis product like a grand EFT fit combination. The discussion of the history of combinations made a clear case about striving to maintain simplicity of components of the combination so that cross-checks are easy to perform and that the overall pieces of the combination each themselves make sense. This goal was easily attained in the era of FORTRAN based analyses of LEP where data interchange formats were necessarily simple and restricted through CERNLIB or custom text files. However, as technology evolved and resulted in the rise in popularity and utility of C++ during the Tevatron, this goal was challenged by the advent of bespoke file formats and increasing computing power since the number of viable methods to achieve the goal of a combination greatly increased. Accumulating experimental information became necessarily more complex and started to require standardization, which has persisted to the LHC experiments, and will be required for combinations with the output of future colliders. Special attention should be paid to making sure that analysis products, like the complete likelihoods and software to evaluate them, can be calculated as-is in future computing systems for combinations.

These grand combinations of the LHC data, much like the electroweak combinations of LEP, will stand as definitive results for multiple decades due to the time it takes to construct next-generation colliders and experiments, and then take sufficient data from them. As we are able to constrain more parameters with new data and the field approaches being able to invert the full matrix, the directions to pursue and subsequent machines to build will become better justified. Useful and guiding summaries like the inverse covariance matrix for all EFT parameters, and the ability to reproduce and combine them on demand to update our knowledge, will be essential tools in determining the course of the collider physics in the next century. While the inverse covariance matrix is certainly the most robust data product to changing software environments and compute techniques, it is the result that is least rich in experimental

392 information and hardest to precisely combine with new measurements as information about systematic uncertainties  
393 and correlations is lost. This means that we need to invest in analysis preservation infrastructure for EFT results,  
394 and that it will be critical to the use of LHC data products throughout the future of the field.

## 395 5 Unfolded results

396 Measured distributions in particle physics are distorted by the finite resolution and limited acceptance of the detectors.  
397 The transformation to the underlying true distribution, called unfolding, is crucial for estimating the EFT sensitivity  
398 of certain variables. Limitations can arise due to our unfolding capabilities and methods, which should be considered  
399 in the EFT interpretation. This session discusses the limitations and the unfolding methods used in physics analyses.  
400 In addition, the treatment of the fiducial measurements that can be used for EFT interpretations is highlighted. Both  
401 theoretical and experimental presentations were included and gave the current picture of unfolding measurements  
402 and treatment, emphasizing in the lessons learned so far and opening the discussion for a future Run3-combination.  
403 The main aspects of the topic were presented. Section 5.1 focus on the challenges in the EFT interpretation of the  
404 unfolded cross-section measurements, followed by section 5.2 on the Higgs and Simplified Template Cross Sections.  
405 Finally, on Section 5.3 we focus on SMEFT probes with LHC Drell-Yan data.

### 406 5.1 Challenges in the EFT interpretation of unfolded cross section measurements

407 *Speaker:* Andrew Gilbert

408 Unfolded differential cross-section measurements have traditionally been a crucial interface between theory and  
409 experiment, and can readily be used for reinterpretation with EFT. A key advantage is that such measurements do  
410 not require any particular new physics model at the time they are produced. As long as the definition of the fiducial  
411 selection and unfolded observables are preserved (e.g., via a RIVET plugin), a new interpretation can be produced  
412 potentially decades later. There are, however, some challenges. We are typically only able to unfold a small number  
413 of variables simultaneously, meaning we may not have optimal sensitivity to all operators. Backgrounds which may be  
414 EFT sensitive are assumed to follow standard model distributions and subtracted. We are often interested in the high  
415 energy tails of distributions, which are difficult to unfold when the expected number of events is small. Furthermore,  
416 there is an implicit assumption that the EFT effects do not modify the efficiency times acceptance within each bin.  
417 With the LHC experiments in the process of releasing the full statistical model information (from which the unfolded  
418 cross sections are derived), we have a more powerful tool to address some of these issues. Nonetheless, releasing  
419 differential cross section results remains important for comparisons to theory, and offers a fallback method for future  
420 reinterpretation.

### 421 5.2 Higgs and Simplified Template Cross Sections

422 *Speaker:* Jonathon Langford

423 The LHC Run 2 dataset has enabled the CMS and ATLAS experiments to go beyond inclusive measurements  
424 of Higgs boson production and decay, and begin measuring Higgs boson interactions differentially. One approach  
425 is within the Simplified Template Cross Section (STXS) framework, where Higgs boson events are first split by  
426 production mode and secondly by kinematic variables such as the transverse momentum of the Higgs boson or  
427 the number of additional jets. Both experiments have performed STXS measurements in a number of the major  
428 Higgs boson decay channels, thus developing a granular, kinematic description of Higgs boson production. One of  
429 the key advantages of STXS is that it provides a natural framework for BSM interpretations, including the use of  
430 Standard Model Effective Field Theory (SMEFT). The standard approach here is to parameterize the Higgs boson  
431 cross sections, at the granularity of the STXS, and decay rates as functions of the SMEFT Wilson coefficients, which  
432 enter the likelihood as signal-strength modifiers. In the talk, the caveats of this approach were discussed which arise  
433 from the fact that we cannot fully encapsulate all EFT effects within simple rate scaling functions. These caveats  
434 include acceptance corrections due to no fiducial selection on the Higgs boson decay products in the STXS, selection  
435 effects where EFT effects can vary significantly within a single STXS bin, and shape effects where the EFT can  
436 change the shape of the fitted observable used for signal extraction (e.g. a multivariate classifier score). In addition,  
437 the STXS binning choice is not optimized for SMEFT sensitivity. Following this, possible future improvements to  
438 the STXS were discussed, including adding a fiducial selection on decay products, updating the binning scheme with  
439 finer granularity and improving the tools used to derive the parametrization (such as standalone reweighting after  
440 the detector simulation). Finally the use of STXS measurements within global EFT fits was presented along with

441 a few points to be addressed. These include the choice of flavor scheme, the simultaneous parametrization of signal  
442 and background, and methods to ensure orthogonality between analyses.

### 443 **5.3 SMEFT probes with LHC Drell-Yan data**

444 *Speaker:* Frank Petriello

445 The experimental precision for a multitude of Drell-Yan observables is approaching the percent level. Future  
446 studies can take advantage of this high-precision data to search for subtle deviations from Standard Model (SM)  
447 predictions. A framework for heavy new physics searches in the absence of new particles is the SM effective field  
448 theory (SMEFT), which contains all operators consistent with SM symmetries and which assumes a mass gap to  
449 any new physics. Detailed studies of the existing Drell-Yan data and of simulated future HL-LHC data reveal a  
450 rich program of discovery in this channel. Observables such as the invariant mass and forward-backward asymmetry  
451 are measured precisely enough to reveal higher-order dimension-8 corrections in the SMEFT. Access to these terms  
452 can potentially discriminate between ultraviolet completions of the SMEFT even without direct observation of a  
453 new particle. This direction will take on a new dimension with high-precision transverse momentum distribution  
454 measurements at an HL-LHC. Angular distributions in the Drell-Yan process also provide a “smoking gun” signature  
455 of dimension-8 effects that are ripe for investigation. Future experimental analyses should take advantage of these  
456 vast possibilities in the Drell-Yan process.

### 457 **5.4 Discussion**

458 The approach of presenting unfolded results for later interpretation has several drawbacks. The primary ones pointed  
459 out by the speakers in this session include the necessity of using only a few observables (which cannot be sensitive to all  
460 SMEFT operators); the built-in assumptions regarding background processes (usually that they are SM); the impact  
461 of SMEFT on efficiency and acceptance (often neglected); and the challenge of defining bins such that SMEFT effects  
462 do not vary significantly within a single bin. Concerns were raised in the question session about producing unfolded  
463 results in multiple correlated observables. It was suggested that, since the same issue is faced in MC generator tuning,  
464 it could be possible to use Monte Carlo to understand these correlations. It was also emphasized that unfolded results  
465 must be produced in bins where the SMEFT expansion parameter remains small - this is something that should be  
466 made very explicit in all interpretations of unfolded results.

467 Considering these substantial drawbacks, the question was then raised about the utility of unfolded results for  
468 re-interpretation by theorists. As demonstrated in the final talk of the session (by Frank Petriello), unfolded cross  
469 sections are very useful for theorists to perform interpretations without diving deeply into the details of ATLAS and  
470 CMS. The consensus was that unfolded results should continue to be produced, but full (reco-level) likelihoods should  
471 also be published and preserved. One alternative to re-interpreting unfolded results would be for the experiments to  
472 provide recipes for forward-folding. This would reach a similar result, but would require more work and experimental  
473 expertise from the person performing the interpretation.

## 474 **6 Building a sample library**

475 In planning the session, a few discussion items that are particularly relevant to event sample production were identified:  
476 the choice of UFO<sup>5</sup> model and what diagram classes to consider (e.g. order in  $\alpha_s$ ); what theory systematic uncertainties  
477 are to be considered; the choice of starting point in model space to use for the initial event sample that may  
478 subsequently be re-weighted; and prospects for sharing samples across experiments. Speakers were invited to address  
479 any of the above points (or others), and a discussion followed.

### 480 **6.1 The interplay between PDF fits and heavy New Physics searches**

481 *Speaker:* Luca Mantani

482 The extraction of parton distribution functions (PDF) from data can potentially conceal effects of heavy new  
483 physics (NP), since the PDF parameterisation can be flexible enough to mimic the NP induced deviations in the tails  
484 of distributions. Moreover, expanding the kinematic coverage of the data at low-Q, for example by incorporating  
485 projected data from forward facilities, can help mitigate the NP “contamination” of the PDFs. This is because  
486 discrepancies between data at high-Q and low-Q would become apparent, prompting the exclusion of high-Q data

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<sup>5</sup> Universal FeynRules Output

487 from the global dataset. Ultimately, the capabilities of SIMUnet, a public framework designed for simultaneous  
488 global fits of PDFs and EFT Wilson Coefficients were presented. Through SMEFT PDF extraction, SIMUnet has  
489 the potential to disentangle NP effects in the data, even without the need for additional projected measurements.

## 490 6.2 Adventures in producing EFT samples

491 *Speaker:* Andrew Wightman

492 Monte Carlo reweighting is a powerful tool that enables exploring the high dimensional phase space of EFT  
493 parameters without needing to generate an unreasonable number of samples. A key part of this approach is the ability  
494 to extract quadratic parameterizations, which enables us to estimate event yields in any kinematic distribution and  
495 event selection for an arbitrary choice of EFT parameters.

496 There are already techniques and tools developed that implement much of the technical machinery needed to  
497 extract these quadratic parameterizations, such as WCFit, but these tools still need to be placed into CMSSW<sup>6</sup>  
498 proper so that they can be more widely accessed within the collaboration. This will also allow for concrete decisions  
499 to be made about conventional choices regarding accessing and storing these objects throughout the different stages  
500 of producing a EFT oriented MC sample. It is also important to consider the choice of starting point used to generate  
501 the MC samples. Some preliminary checks of how well the sample reweights to different parts of the phase space  
502 should be done in order to get a proper estimate for how many events will be needed to ensure reasonable statistical  
503 power for an analysis.

## 504 6.3 Building the STXS Parameterization

505 *Speaker:* Charlotte Knight

506 An EFT parameterization of the STXS is a big undertaking, requiring the use of multiple techniques and tools,  
507 e.g. both reweighting and generation-based approaches. Furthermore, post-generation tools are necessary to prop-  
508 erly account for acceptance effects, and to perform more validity checks (see Sec. 5.2). To reduce the amount of  
509 duplicated work, a common parameterization is being developed in collaboration with CMS, ATLAS and theorists.  
510 This parameterization will be published in a proposed common format and it is our hope that future publications  
511 will conform to this so that a library of parameterizations can be created. We would also like to encourage, where  
512 possible, the release of workflows to accompany these publications so parameterizations can be easily reproduced and  
513 altered if desired.

## 514 6.4 Discussion

515 In discussing the interplay between PDFs and SMEFT (Sec. 6.1), there was a question about the adequacy of existing  
516 methods for estimating PDF uncertainties in event samples produced for use with EFT fits in CMS, and how to fold  
517 the results of SIMUnet studies into experimental uncertainty estimates. It was understood that there will be a  
518 need for improved PDF measurements using high- $x$  low- $Q^2$  data such as that provided by forward physics facilities  
519 (e.g. FASER) that may better disentangle PDF degrees of freedom from EFT ones. SIMUnet could potentially be  
520 used to design a new PDF model for use in experimental fits that is more conservative in its extrapolations to high  
521  $Q^2$  in view of the sensitivity to NP effects in that region. A potential action item for CMS analysts is to measure  
522 the degree of correlation between PDF Hessian eigenvariations and EFT parameter effects on their analysis regions  
523 and compare the result with that provided by SIMUnet for the same event topology, to better understand if the  
524 uncertainty estimates can detect for possible absorption of NP discrepancies into the PDF model.

525 During the discussion on sample generation (Sec. 6.2), several questions revolved around what an optimal gener-  
526 ation strategy might look like. For example, it is unclear whether generating inclusive samples (e.g.  $t\bar{t}$ ) or exclusive  
527 decay mode samples with EFT weights is more resource-efficient. In the case of this analysis, samples were produced  
528 with exclusive decays generally. Should one use the same point in WC space for sampling events across different  
529 processes? Generally, no because it can be challenging to find a point in WC space that samples events that can  
530 be effectively reweighted to all interesting WC values. In the case of the analysis presented in 6.2, one sample was  
531 sufficient but in other cases there may be a need to combine several samples together that are generated from different  
532 starting points. Methods for combining samples can be constructed, but a thorough evaluation of the options should  
533 be an action item for CMS analysts. There is progress on integration of the code to compute the WC polynomial  
534 expansion into CMSSW, but it is not complete yet.

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<sup>6</sup> CMS offline software framework

535 In the discussion on building the STXS parameterization (Sec. 6.3), many in the audience were very interested  
536 in the post-generation reweighting technique presented. It was discussed how best to integrate the MadGraph-  
537 based reweighting module into vectorized python workflows and what common components are shared with the  
538 code for computing WC parameterization used in the previous talk that could be harmonized. The post-generation  
539 reweighting technique was identified as an effective method to determine the EFT parameterization of reducible  
540 backgrounds without having to re-generate potentially large simulation samples, provided it can be shown that the  
541 result is faithful to that of reweighting a dedicated sample. Even samples produced with a different generator could  
542 in principle be reweighted, but the results should be treated with caution, especially in the case where the considered  
543 diagrams are not common, and are likely not meaningful between LO and NLO QCD generators. It was highlighted  
544 that, in cases where the quadratic term is to be used mainly as an estimate of the missing higher order terms in  
545 the  $1/\Lambda$  expansion, that one can relax the requirements on the MC statistical uncertainty in determining those  
546 coefficients. The need for a schema for the JSON data format for bin parameterizations was also highlighted.

## 547 **7 Conclusion**

548 As evidenced by the extensive discussions, a dedicated workshop could easily have been arranged for each of the five  
549 themes. Although some consensus on the best approach to performing large-scale experimental SMEFT combinations  
550 is building, it is clear that more work will be needed. We believe that there are paths laid out here that are actionable  
551 in the short term, while still being flexible to future changes in requirements for (re-)interpretability. Future workshops  
552 will help refine a plan of action for this community.

## 553 **Acknowledgments**

554 We would like to thank the University of Notre Dame for making space available for this workshop. We would also  
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