

# Which orders?

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# In SMEFT framework

$$|A|^2 = |A_{SM}|^2 + \frac{2\text{Re}(A_{SM}^* A_6)}{\Lambda^2} + \frac{1}{\Lambda^4} \left( |A_6|^2 + 2\text{Re}(A_{SM}^* A_8) \right) + \dots$$

interference piece,  
usually largest effect.  
State of the art  
SMEFT

'Higher order'  
 $\mathcal{O}(1/\Lambda^4)$   
corrections

Dual expansion: gotta match dimensions, so numerator  $\sim$  powers of

$$v, \partial_\mu \sim E$$

At high energy  $\left(\frac{E^n}{\Lambda^n}\right) > \left(\frac{v^n}{\Lambda^n}\right)$ : main advantage of SMEFT at LHC

# In SMEFT framework

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corrections

**larger expansion parameter = more sensitive to higher orders!**

- To know error on  $1/\Lambda^2$  piece, we should know next order!
- Additionally, there are circumstances where interference is suppressed. Then  $1/\Lambda^4$  is the leading SMEFT piece

# OK, so we'd like to include $\mathcal{O}(1/\Lambda^4)$ effects

**BUT!**

SMEFT Warsaw basis:  $\mathcal{O}(60)$  operators at dim-6  
(flavor universal, CP)  $\mathcal{O}(1000)$  operators at dim-8

Can't we just do  $|\text{dim} - 6|^2$ ?

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can be okay if nothing else, but **lots** of pitfalls

- $|\text{dim} - 6|^2$  is positive definite, total  $\mathcal{O}(1/\Lambda^4)$  need not be
- $|\text{dim} - 6|^2$  limited to dim - 6 operators...  
limited structure, some already bounded, small in some UV setups

**Can lead to wildly inaccurate estimates of  $\mathcal{O}(1/\Lambda^4)$  ...**

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**Can lead to wildly inaccurate estimates of  $\mathcal{O}(1/\Lambda^4)$  ...**

**Especially dangerous if  $|\text{dim} - 6|^2 > SM \times (\text{dim} - 6)$   
without a good reason!!**

# geoSMEFT-ist perspective

[2001.01453 Helset, AM, Trott]

**geoSMEFT** = re-organization of SMEFT that makes many key processes (for LHC SMEFT global fit) calculable  $\mathcal{O}(1/\Lambda^4)$  without needing 1000 operators. Clarifies E vs. v counting



Calculate away, forming a library of process to use as a laboratory to study ‘truncation error’.

Organize operators by the smallest vertex (# of particles that enter) they can impact at tree level: 2, 3,4, etc. Minimize the # of operators affecting 2, 3-particle vertices by strategically placing derivatives (IBP)

- $(H^\dagger D_\mu H)^*(H^\dagger D^\mu H) \supset v^2 (\partial_\mu h)^2$  contributes to 2-particle vertex

$(\psi^\dagger \psi)^2$  contributes to 4-particle vertex



- $\square (H^\dagger H) \square (H^\dagger H) \supset v^4 (\partial_\mu h)^2$  would contribute to

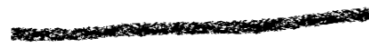
but can use IBP to manipulate to

$(D_\mu H^\dagger D^\mu H D_\nu H^\dagger D^\nu H)$  which only affects 4+ particle vertices



# At dimension-6, assuming B,L, flavor universal (59 total)

Min vertex:



Operator type:

[X = field strength,  
D = deriv]

$$\begin{aligned}
 &H^6 \\
 &H^4 D^2 \\
 &H^2 X^2 \\
 &\psi^2 H^3
 \end{aligned}$$

$$\begin{aligned}
 &X^3 \\
 &\psi^2 H X \\
 &\psi^2 H^2 D
 \end{aligned}$$

$$\psi^4$$

Number:

**14**

**20**

**25**

**0**

**At dimension-8, assuming B,L, flavor universal (993 total)**

Min vertex:



Operator type:

[X = field strength,  
D = deriv]

$$H^8$$

$$H^6 D^2$$

$$H^4 X^2$$

$$\psi^2 H^5$$

$$H^2 X^3$$

$$\psi^2 H^3 X$$

$$\psi^2 H^4 D$$

$$H^4 D^2 X$$

$$X^4$$

$$H^4 D^4$$

$$X^2 H^2 D^2$$

...



$$\psi^4 X$$

Number:

**19**

**47**

**927!**

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$$H^4 D^4$$

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...



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**927!**

If we also impose CP, U(3)<sup>5</sup> (remember, must interfere to enter 1/Λ<sup>4</sup>)

**8**

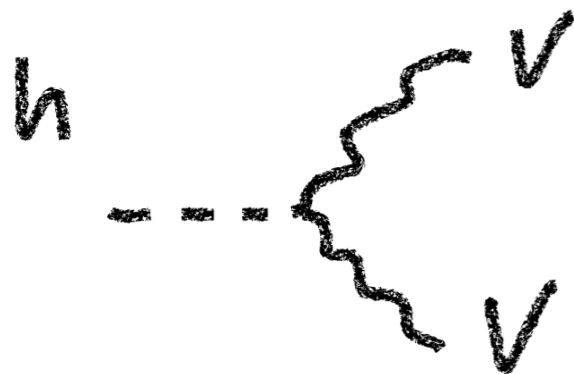
**22**

[trend continues to dim > 8 too!]

# Why is this a good idea?

- “Universal” corrections related to inputs  $\sim O(10)$  new operators.  
Simplest building block vertices  $\sim O(20)$  ops
- Bulk of operators pushed to more process-specific, 4+-particle interactions
- 2-, 3- particle interactions: going from dim-6 to dim-8 doesn't change kinematics — just added additional  $H^2$ ! Additional derivatives aren't possible, as all momentum products reduce to masses = constants.  
So the energy/vev scaling of these terms is set by whatever happens at dim-6

Ex.)



$$\sim \frac{E^2 v}{\Lambda^2} \text{ at dim-6}$$

$$H^\dagger H X_{\mu\nu} X^{\mu\nu}$$

$$\sim \frac{E^2 v}{\Lambda^2} \left( \frac{v^2}{\Lambda^2} \right) \text{ at dim-8}$$

$$(H^\dagger H)^2 X_{\mu\nu} X^{\mu\nu}$$

# Why is this a good idea?

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So, if we're hunting for energy enhanced effects  $\rightarrow$  energy enhanced  
dim-6 3-particle + dim-6, dim-8 contact

# Ok, what do I do with this?

1.) Simplest LHC processes: resonances, 2 -> 2 can be done 'fully' to  $\mathcal{O}(1/\Lambda^4)$  without an order of magnitude increase in operators

$$gg \rightarrow h \rightarrow \gamma\gamma, \gamma Z \quad \text{Z-pole, Drell-Yan} \quad pp \rightarrow V(\ell\ell)h$$
$$pp \rightarrow W(\ell\nu)\gamma$$

[Kim, AM 2203.11976 ] [Boughezal et al 2106.05337, 2207.01703

[Corbett, AM, Trott 2107.07470 ] [AM, Trott 2305.05879 ] [Hays, Helset, AM, Trott 2007.00565 ]

For these, can use  $\mathcal{O}(1/\Lambda^4)$  as an uncertainty on extraction of dim-6 operators [how to do this systematically?]

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For these, can use  $\mathcal{O}(1/\Lambda^4)$  as an uncertainty on extraction of dim-6 operators [how to do this systematically?]

2.) Initial step: focus on terms that grow with energy (fully, to  $\mathcal{O}(1/\Lambda^4)$ ). Assuming all WC are same size, these effects will be largest

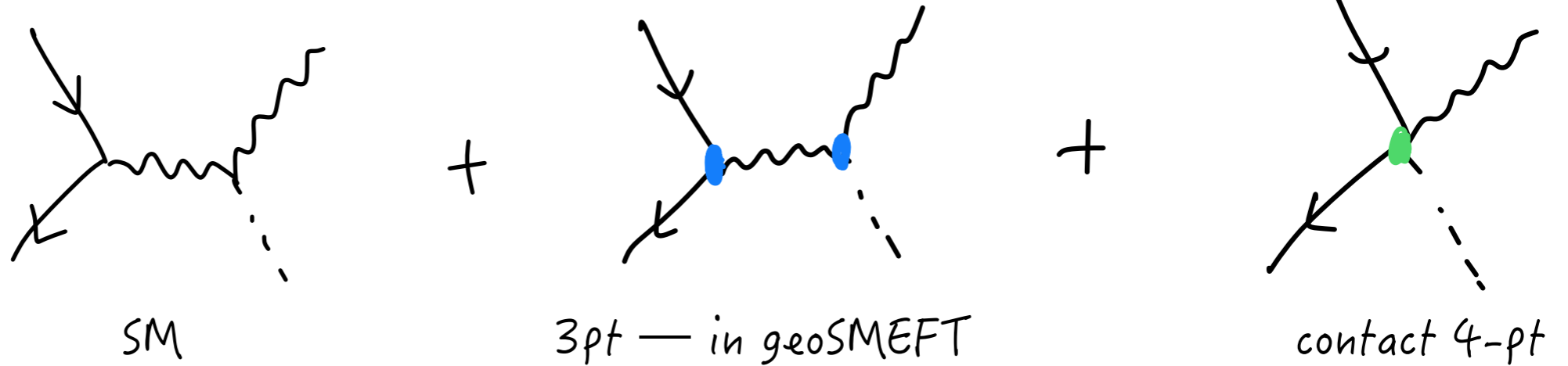
$$pp \rightarrow W^+W^-, W^\pm Z \quad \text{VBF } pp \rightarrow hjj$$

[2303.10493 Degrande]

[Assi,AM in prep]

# Do I gain something vs. using $|\text{dim-6}|^2$

Example: VH



Energy enhanced effects

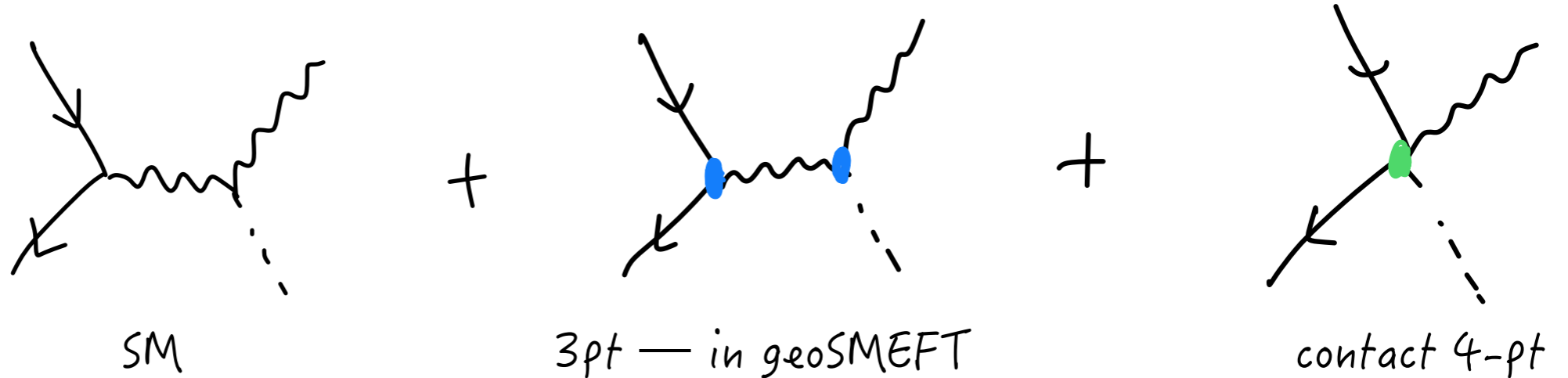
**dim-6:** vertex  $H^\dagger H W_{\mu\nu} W^{\mu\nu}$   
 contact  $(Q^\dagger \bar{\sigma}^\mu \tau^I Q) H^\dagger \overleftrightarrow{D}_I H$

**dim-8:** contact  $\psi^2 H^2 D^3 \supset (Q^\dagger \sigma^\mu D^\nu Q)(D^\mu H^\dagger D_\nu H)$   
 $\psi^2 H^2 X D \supset (Q^\dagger \bar{\sigma}^\mu Q) D^\nu (H^\dagger H) B_{\mu\nu}$



# Do I gain something vs. using $|\text{dim-6}|^2$

Example: VH



Energy enhanced effects

**dim-6:**

vertex

$$\cancel{H^\dagger H W_{\mu\nu} W^{\mu\nu}}$$

contact

$$(Q^\dagger \bar{\sigma}^\mu \tau^I Q) H^\dagger \overleftrightarrow{D}_I H$$

SM dominantly  $V_L$ , won't interfere with these!

**dim-8:**

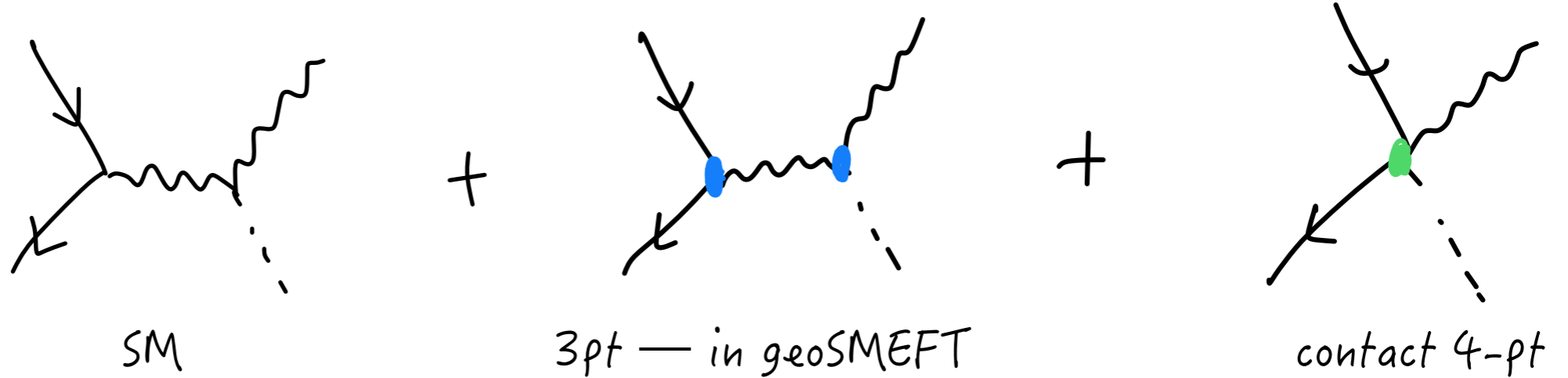
contact

$$\psi^2 H^2 D^3 \supset (Q^\dagger \sigma^\mu D^\nu Q)(D^\mu H^\dagger D_\nu H)$$

$$\cancel{\psi^2 H^2 X D \supset (Q^\dagger \bar{\sigma}^\mu Q) D^\nu (H^\dagger H) B_{\mu\nu}}$$

# Do I gain something vs. using |dim-6|^2

Example: VH



Energy enhanced effects

**dim-6:**

vertex

$$\cancel{H^\dagger H W_{\mu\nu} W^{\mu\nu}}$$

contact

$$(Q^\dagger \bar{\sigma}^\mu \tau^I Q) H^\dagger \overleftrightarrow{D}_I H$$

interference  $\sim g_{SM}^2 \frac{\hat{s}}{\Lambda^2}$

squared  $\sim \frac{\hat{s}^2}{\Lambda^4}$

**dim-8:**

contact

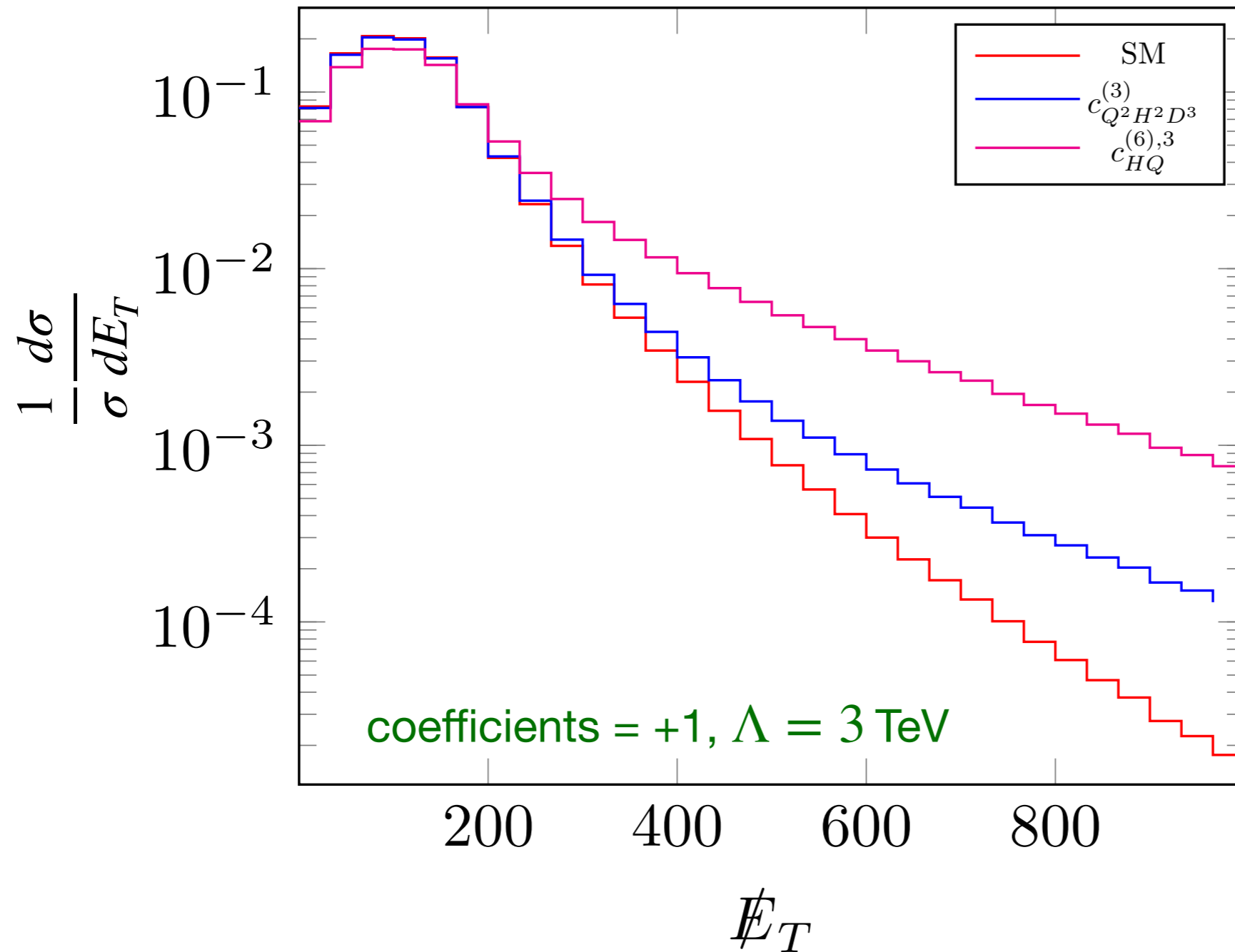
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interference  $g_{SM}^2 \frac{\hat{s}^2}{\Lambda^4}$

# Do I gain something vs. using |dim-6|^2

Effects at large  $\hat{s}$  controlled by:

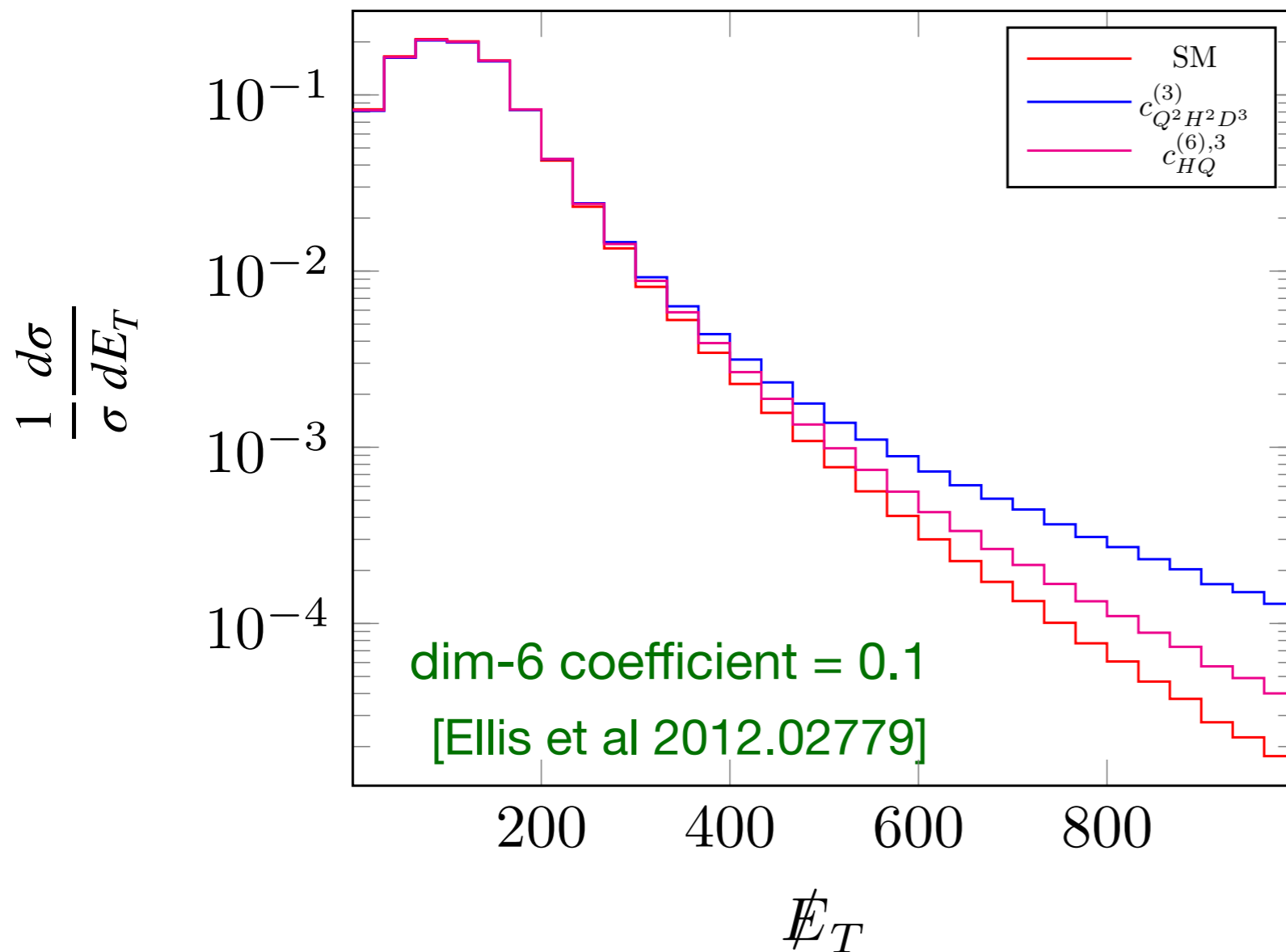


$$(Q^\dagger \bar{\sigma}^\mu \tau^I Q) H^\dagger \overleftrightarrow{D}_\mu \tau^I H$$

$$Q^\dagger \bar{\sigma}^\mu \tau^I D_\nu Q D^\mu H^\dagger \tau_I D_\nu H$$

# Do I gain something vs. using |dim-6|^2

But,  $Q^\dagger \bar{\sigma}^\mu \tau^I Q H^\dagger \overleftrightarrow{D}_I H$  etc.  $\supset Q^\dagger \bar{\sigma}^\mu Q Z_\mu$  are constrained by LEP, while  
 $Q^\dagger \bar{\sigma}^\mu \tau^I D_\nu Q D^\mu H^\dagger \tau_I D_\nu H$  are not ( $\not\supset Q^\dagger \bar{\sigma}^\mu Q Z_\mu$ )



complying with those constraints, dim-8 terms dominate in large  $\hat{s}$  regime

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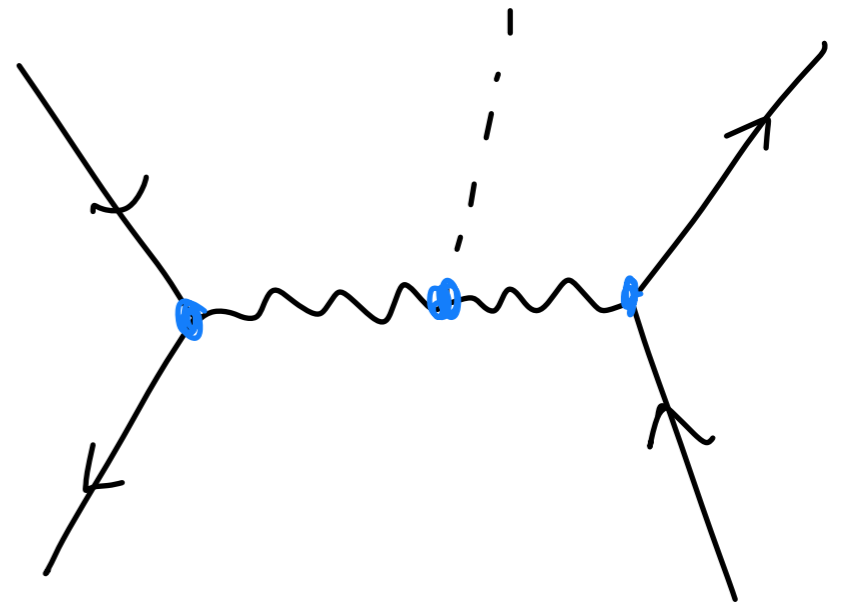
$\bar{q}q \rightarrow V(\bar{q}q)H$  + crossing  
 symmetry gets us VBF

Therefore, expect similar  
 operators to dominate, though  
 kinematics and cuts are  
 slightly different

i.e.  
 $\psi^4 H^2$



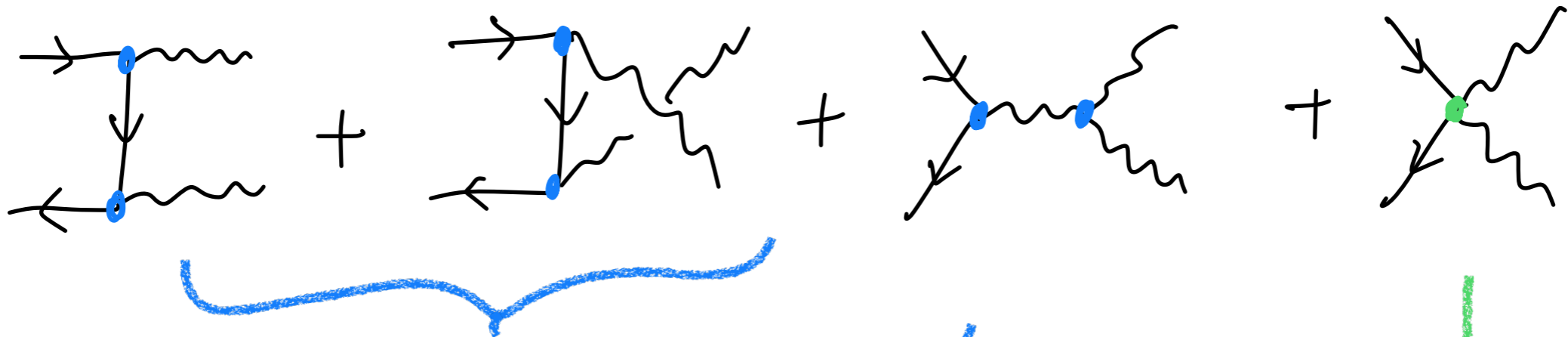
pp →  
 Hjj



pp → VH

# Do I gain something vs. using $|\text{dim}-6|^2$

Diboson



Assuming CP,  $U(3)^5$ , no energy enhancements

VW is energy enhanced ( $C_W W^3$ ).  
Important in global fit program, as  
first place triple gauge operators  
as appear.

Contact terms  
only show up at  
dim-8, ex. class  
 $\psi^2 X^2 D$

Example:  $\gamma W^\pm$ , organize calculation by the polarizations of the  $W, \gamma$

# Do I gain something vs. using |dim-6|<sup>2</sup>

Energy scaling of different polarization amplitudes

$\epsilon_\gamma \epsilon_W$	SM	dim-6 $C_W$
++	$\frac{v^2}{s}$	$\frac{s}{\Lambda^2}$
+-	1	0
+0	$\frac{v}{\sqrt{s}}$	$\frac{v\sqrt{s}}{\Lambda^2}$

$$|A_{SM}|^2 + \frac{2\text{Re}(A_{SM}^* A_6)}{\Lambda^2} + \frac{1}{\Lambda^4} |A_6|^2$$

with dim-6 alone, largest energy enhancement (to  $\mathcal{O}(1/\Lambda^4)$ ) comes from from

$$|\text{dim-6 } C_W|^2 \sim \frac{s^2}{\Lambda^4}$$

$$\hat{s} \gg m_W^2$$

# Do I gain something vs. using $|\text{dim-6}|^2$

[AM, 2312.09867]

		$W^3$	$\psi^2 W^2 D$
$\epsilon_\gamma \epsilon_W$	SM	dim-6 $C_W$	dim-8 contact
++	$\frac{v^2}{s}$	$\frac{s}{\Lambda^2}$	$\frac{s^2}{\Lambda^4}$
+-	1	0	$\frac{s^2}{\Lambda^4}$
+0	$\frac{v}{\sqrt{s}}$	$\frac{v\sqrt{s}}{\Lambda^2}$	$\frac{vs^{3/2}}{\Lambda^4}$

$$|A_{SM}|^2 + \frac{2\text{Re}(A_{SM}^* A_6)}{\Lambda^2} + \frac{1}{\Lambda^4} |A_6|^2$$

$$+ \frac{2\text{Re}(A_{SM}^* A_8)}{\Lambda^4}$$

**But:** dim 8

$$(Q^\dagger \bar{\sigma}^\mu \tau^I \overleftrightarrow{D}_\nu Q) W_{\mu\rho}^I B_{\rho\nu}$$

can interfere with dominant SM polarization

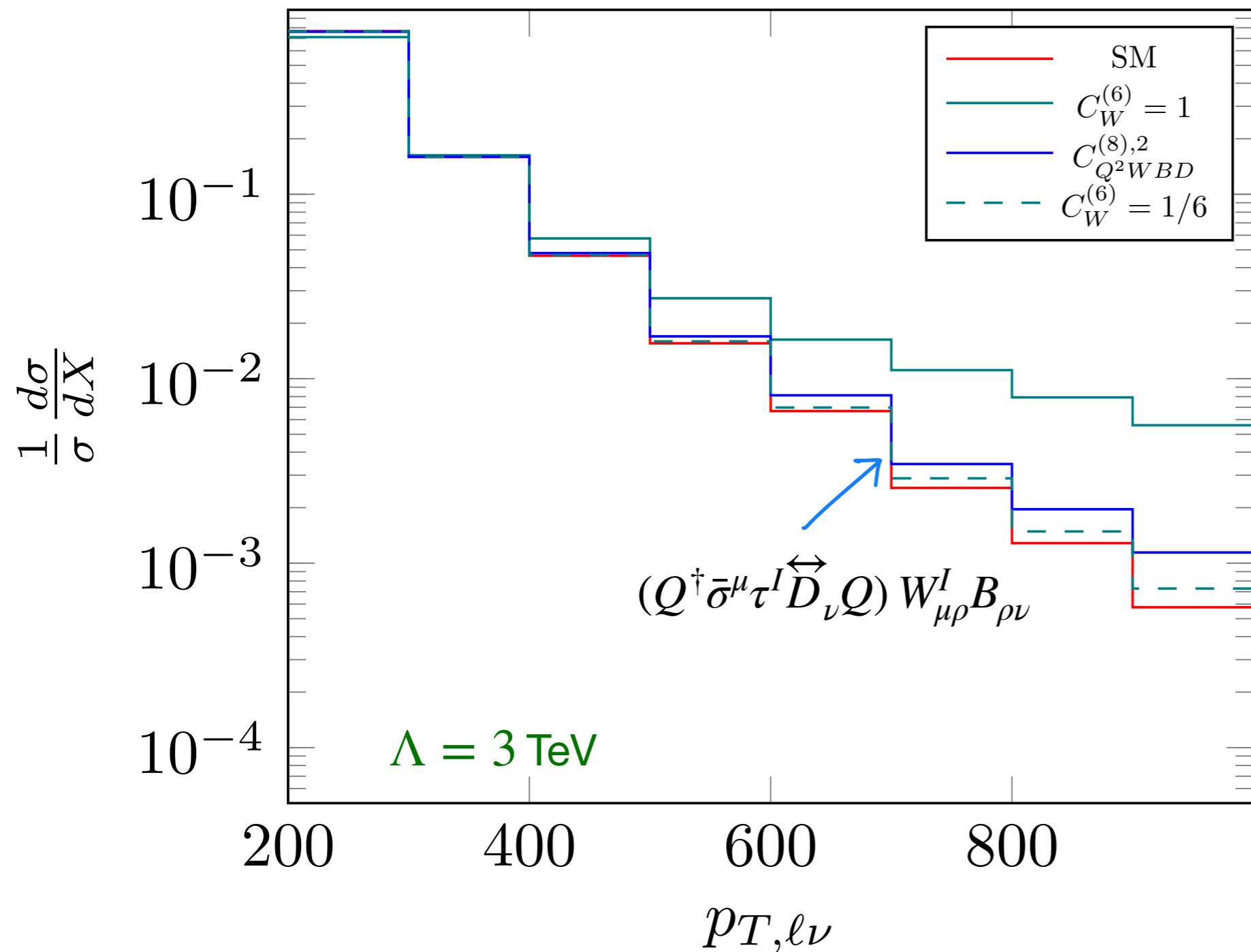
$$SM \times \text{dim-8} \sim \frac{s^2}{\Lambda^4}$$

$$\hat{s} \gg m_W^2$$

See also Degrande 2303.10493 (for WW, WZ)



# Do I gain something vs. using $|\text{dim-6}|^2$



As in VH, dim-8 effects non-negligible, even dominant

$$\varepsilon^{IJK} W_\mu^{I\nu} W_\nu^{J\rho} W_\rho^{K\mu}$$

Motivates polarization studies, 'taggers'

# Takeaways

To take advantage of ‘energy frontier’ at LHC, need to know next order SMEFT corrections.

$|\text{dim-6}|^2$  is an unreliable estimate at best! (And  $|\text{dim-6}|^2 > \text{dim-6} \times \text{SM}$  without good reason I don't trust at all)

geoSMEFT organization: minimizes operators that enter smallest (& most universal) vertices. Pushes new energy-enhanced effects to process-specific 4+ particle vertices

Facilitates full  $\mathcal{O}(1/\Lambda^4)$  calculations. Several key processes relevant for global SMEFT program worked out.

Easy energy vs. vev counting: as first step, focus on energy enhanced terms to  $\mathcal{O}(1/\Lambda^4)$ . Assuming all WC are the same size, these will dominate kinematic tails

# Takeaways

Facilitates full  $\mathcal{O}(1/\Lambda^4)$  calculations. Several key processes relevant for global SMEFT program worked out.

Easy energy vs. vev counting: as first step, focus on energy enhanced terms. Assuming all WC are the same size, these will dominate kinematic tails

In either scenario, how do you actually include  $\mathcal{O}(1/\Lambda^4)$  as uncertainty? At least from examples worked out so far, impact of  $\mathcal{O}(1/\Lambda^4)$  strongly depends on process and kinematic regime...

What's the best format for  $\mathcal{O}(1/\Lambda^4)$  theory calculations?