

CBS experiment @ **CTF2** (ARTI)

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The CTF2 gun









CLINXBT_0001

S-band standing wave 1.5 cell RF-gun intended for the new AWAKE injector.

Prototype constructed by INFN-Frascati and commissioned at CTF2.

Fabricated with brazing free technology [1].

[1] D. Alesini, et al., PRAB, vol. 21, n. 11, November 2018



Experimental set-up





Experimental match with simulations



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Simulations of CTF2 injector

A model of the CTF2 gun was implemented in RF-Track. The model comprised the cathode, gun, and solenoid.

Optimisation goals:

- Maximise flux
- Minimise background

0 – 20 mm/c 30 – 200 mm/c

Evolution of the phase space from cathode to screen



Cathode	Gun	Solenoid
Bunch charge (laser intensity)	RF phase	B field
Laser spot size (Gaussian)	RF gradient	



Flux optimisation (cathode → screen)





Small spot working point



Small spot obtained with a -10 deg RF phase and 200 µm laser spot on the cathode.

An electron RMS spot of 140 µm for a bunch charge of 200 pC.

The beam halo inducing a non-Gaussian beam profile accounts for 77% of charge.

Why the larger measured beam size?



Offsets in the cathode and solenoid alignment



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Novel solenoid offset inference method

Can infer the solenoid offset from tracking the electron beam size and position through a solenoid scan.

Offsets lead to an emittance increase \rightarrow larger beam size at focus. The beam displacement trajectory depends on the alignment of the solenoid with respect to the gun.





Novel solenoid offset inference method



Laser input	Value	RF input	Value
Pulse length	327 fs RMS	Bunch charge	53 pC
RMS spot X	160 um	Gradient	106.7 MV/m
RMS spot Y	230 um	Phase	-20 deg

Offset input	Value
Solenoid offset X	2.66 mm
Solenoid offset Y	3.19 mm
Solenoid pitch	0.625 mrad
Solenoid yaw	0.770 mrad
Cathode offset x	2.30 mm
Cathode offset y	1.33 mm



Use the inferred offsets to beam sizes measured for solenoid scans with various RF phases and laser spot sizes.



Electron beam size dependence on RF phase

200

80

-40

-20

-10

RF phase [deg]

10

20

0

-30

0.32

-30

-20

-10

RF phase [deg]

0

10



0.32

-30

-20

-10

RF phase [deg]

0

10

60

-40

-30

-20

-10

RF phase [deg]

0

20

10

Experiment: electron beam size from projection

Electron beam size dependence on the laser spot



Experiment



400 pC

1500

1000

500

70

0.2

0.25

0.35

0.3

RMS laser spot [mm]

0.4

0.45

200 pC

Gun solenoid field [T]



RF-Track





70

0.2

0.25

0.3

RMS laser spot [mm]

0.35

0.4

0.45

Cathode laser and solenoid alignment



Centering the laser on the gun axis

RF focusing was used to have beam with the solenoid off. To achieve this:

- Low beam energy (0.7 MeV)
- Low bunch charge (5 pC)

Followed the beam drift across a phase scan for two different laser position coordinates (transverse RF fields). Gun centre is where the two drifts intersect.

Beam centre actually outside the camera aperture -> brought it as close as possible.

Corrected cathode offset was: X = 160 um Y = 940 um









Centering the solenoid

The solenoid was aligned by increasing the solenoid strength (0.02 T) and bringing the drifted electron beam back to its initial position.

Adjusted solenoid movers according to trajectories seen in simulations.

Corrected solenoid offsets X: 0.818 mm Y: 1.625 mm









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Impact on beam size from solenoid and gun centering



A 40% reduction in beam size was achieved after correcting the solenoid and laser cathode position.

Minimum beam in plot reduced from 120 um to 72 um.

Offsets experimentally corrected (not final):

Cathode X: 0.160 mm Y: 0.940 mm Solenoid X: 0.818 mm Y: 1.625 mm



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Detection chamber vacuum test





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Vacuum test: Detection chamber

Vacuum test of the detection chamber showed the MCP detector is working as expected. Dark counts (field emission, cosmic muons, residual gas ionisation, local discharge) of 3 per second were observed.





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Conclusions & Next steps

- Simulations of the gun at CTF2 showed that it can be used to generate water window X-rays for a proof of principle experiment.
- The cathode and solenoid offsets have been adjusted, which led to a 40% reduction in beam size.
- The detection chamber has been assembled and tested, installation in the next few weeks.
- On track to have the experiment in the first half of this year.

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Compton backscattering

= The scattering of a low energy photon from an EM field to a high-energy photon (X-ray or gamma ray) during the interaction with a charged particle.

$$N_{\gamma} = \sigma_{c} \frac{N_{e} N_{laser} \cos(\phi/2)}{2\pi \sigma_{\gamma,y} \sqrt{\sigma_{\gamma,x}^{2} \cos^{2}(\phi/2) + \sigma_{\gamma,z}^{2} \sin^{2}(\phi/2)}} \qquad \qquad \mathcal{B} = \frac{\mathcal{F}}{4\pi^{2} \sigma_{\gamma,x} \sqrt{\epsilon_{x}/\beta_{x}} \sigma_{\gamma,y} \sqrt{\epsilon_{y}/\beta_{y}}}$$

$$\mathbf{Total flux} \qquad \qquad \mathbf{Average brilliance}$$

$$\frac{\sigma_{E_{\gamma}}}{E_{\gamma}} = \sqrt{\left(\frac{\sigma_{E_{\theta}}}{E_{\theta}}\right)^{2} + \left(2\frac{\sigma_{E_{e}}}{E_{e}}\right)^{2} + \left(\frac{\sigma_{E_{laser}}}{E_{laser}}\right)^{2} + \left(\frac{\sigma_{E_{\epsilon}}}{E_{\epsilon}}\right)^{2}} \qquad \qquad E_{X-ray} = 2\gamma^{2} E_{laser} \frac{1 + \cos \phi}{1 + \gamma^{2} \theta^{2}}$$

$$\mathbf{Photon bandwidth} \qquad \qquad \mathbf{Photon energy}$$

Water Window X-ray source

- Water window is a region in the electromagnetic spectrum where water is invisible.
- The region spans the K-edge of Carbon (282 eV) to Oxygen (533 eV).
- Only microscopy allowing for 10 nm range 3D imaging of cellular samples in their near-native state → quantitative characterisation of (sub)cellular organisation in single cells and cell-cell interaction [1].
- Most water window microscopes have been implemented as part of synchrotrons.

[1] V. Weinhardt, J.-H. Chen, A. Ekman, G. McDermott, M. A. Le Gros, and C. Larabell, "Imaging cell morphology and physiology using x-rays," Biochem. Soc. Trans. 47, 489–508 (2019).

Fig. 9. 2D cryo imaging with laboratory soft XRM. (a) Healthy and adhered HEK 293 T cell, 30 s exposure time; (b) slightly starved and rounded HEK 293 T cell, 20 s exposure time; (c) THP-1 cells with 5 min exposure time. Images (a) and (b) from the Stockholm microscope and (c) from the Berlin microscope.

https://doi.org/10.1364/OPTICA.393014

Alignment of laser spot on cathode wrt the solenoid

Method: scan the laser position on the cathode, and for each point, made a solenoid scan and determined the beam drift \rightarrow select laser coordinates with smallest drift.

Stable beam region set by laser spot position within 160 μ m.

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Alignment of laser spot on cathode wrt solenoid

Stable beam for a large range of solenoid strengths was obtained.

Beam size measurement

Three methods used to compute the beam size:

- 1. From the gaussian fit of projected pixels in X and Y
- 2. From the gaussian fit of a slice in x and y passing through the beam centre
- 3. From the weighted mean of the pixel intensities

Pixel projection is standard, however it tends to overestimate the beam size in the case of beam halo.

The weighted mean of the pixel intensities allows for computing the beam size with no fit. However, it is very sensitive to background.

To mitigate this, a pixel slice was used to compute the beam profile. The direct cut showed that beam size was up to 4 times smaller than the on predicted from the projected pixels.

