
HETERODYNE DETECTION OF AXIONS & GWS IN SRF CAVITIES

The Cavity Variations

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JHEP 07 (2020) 088, [hep-ph/1912.11048](#)

A. Berlin, R. T. D'Agnolo, **SARE**, P. Schuster, N. Toro,
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PRD 104 (2021) 11, L111701, [hep-ph/2007.15656](#)

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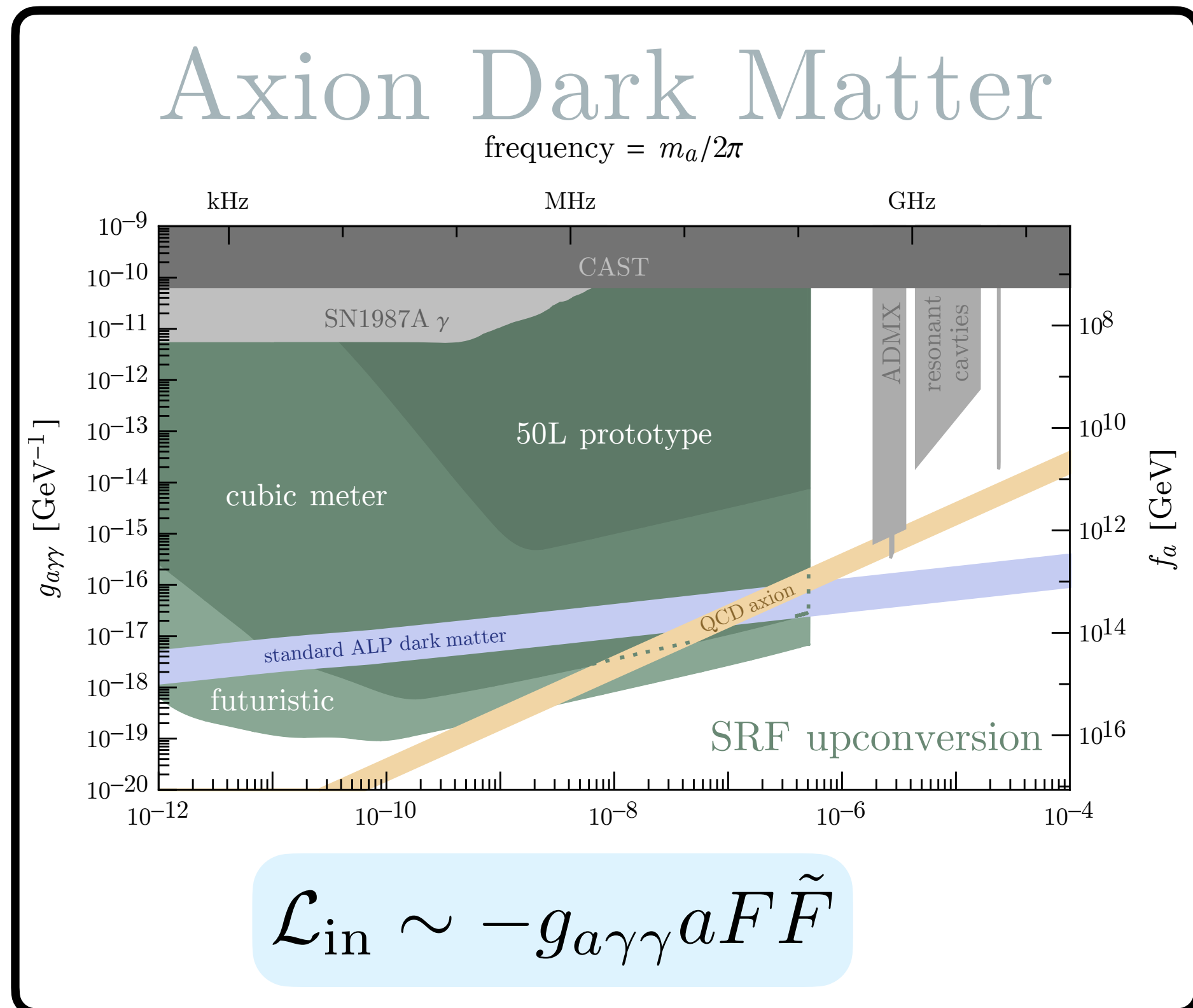
PRD 108 8, 084058 [hep-ph/2303.01518](#)

A. Berlin, D. Blas, R. T. D'Agnolo, **SARE**, R. Harnik, Y.
Kahn, J. Schütte-Engel, M. Wentzel

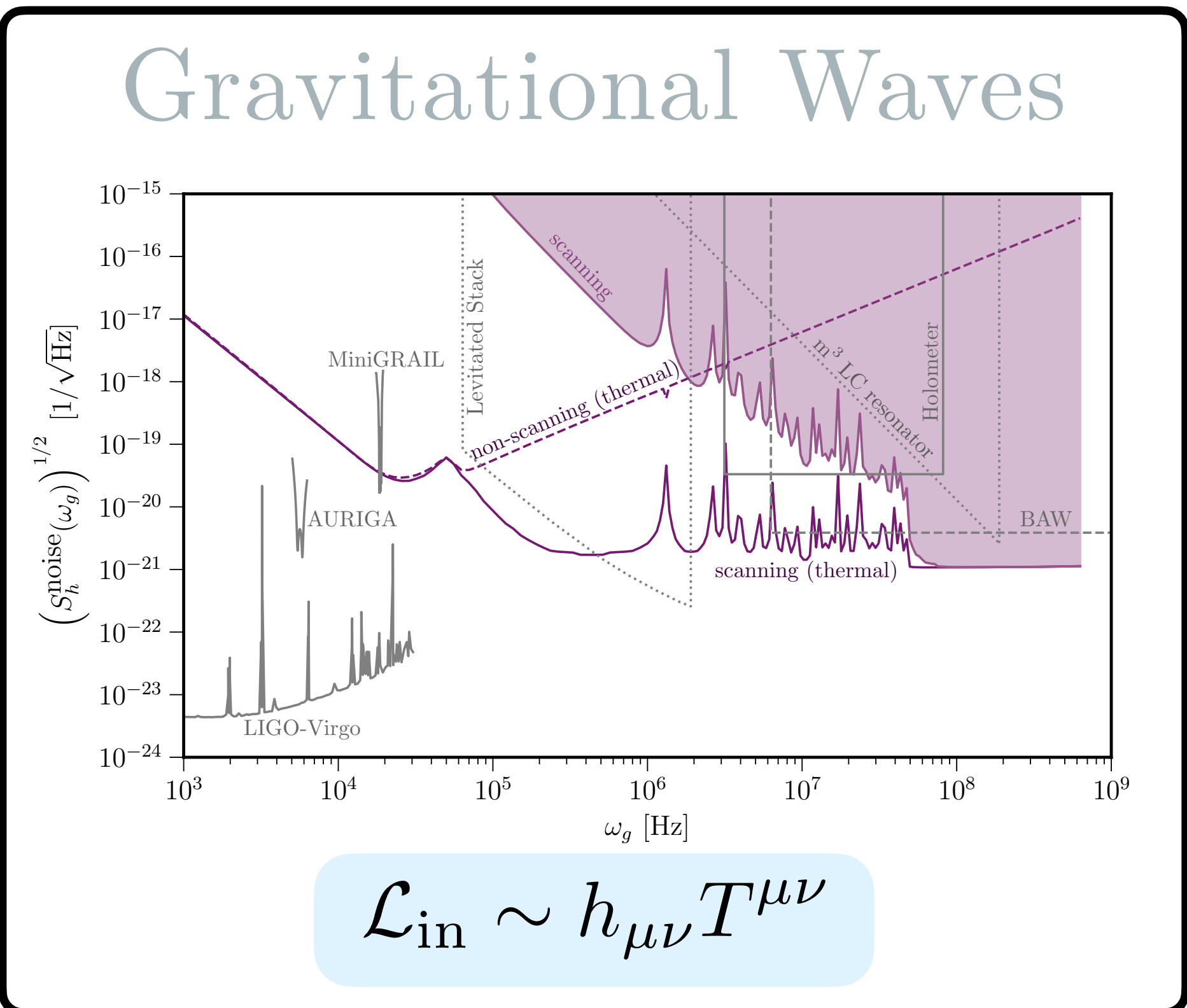
High-level Summary

Search for weakly-coupled signals in a *loaded* SRF cavity

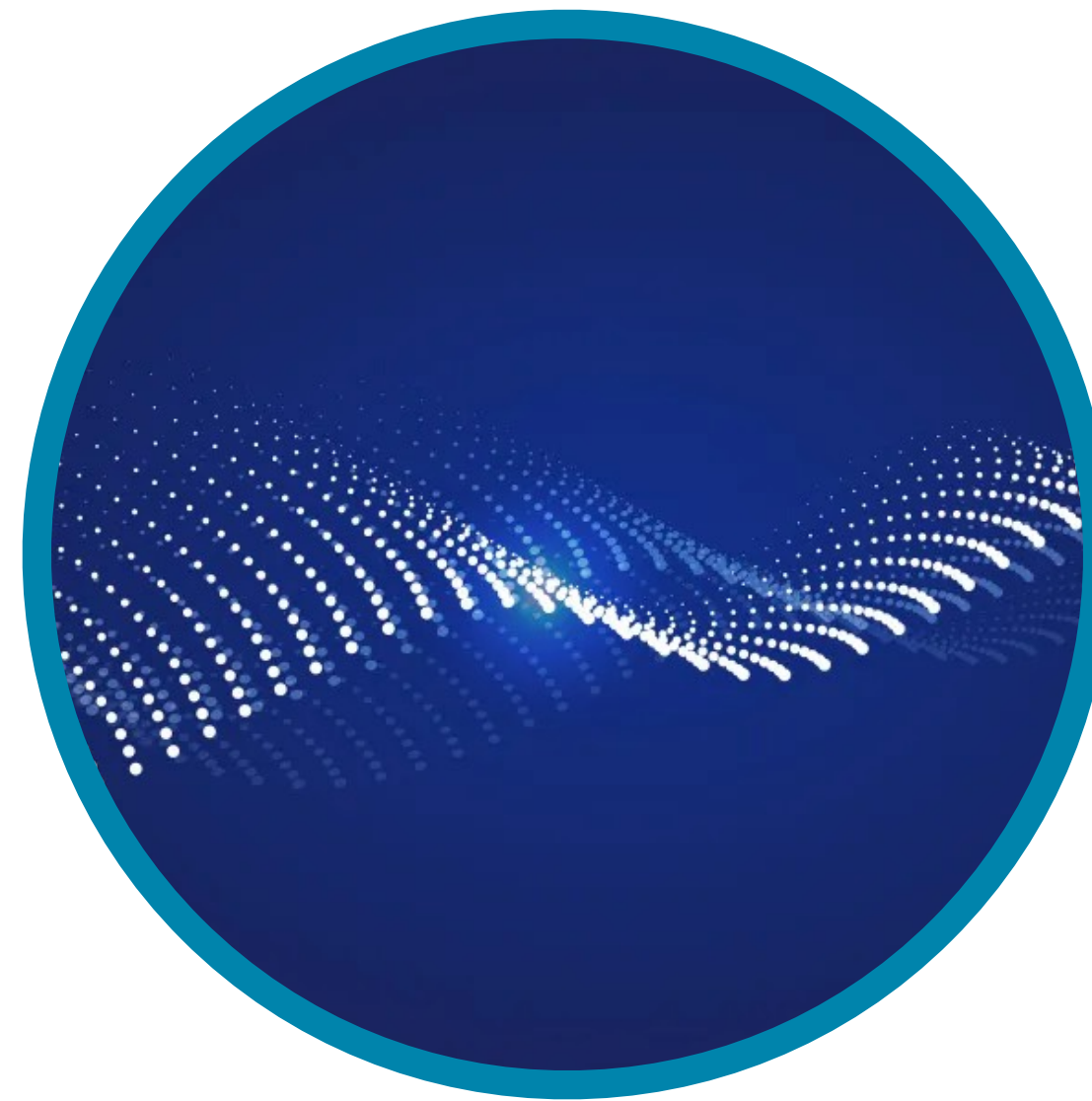
Oscillating background B-field: heterodyne **up-conversion** approach decouples input ω_{in} from V_{det}



$$\omega_{\text{sig}} = \omega_0 \pm \omega_{\text{in}}$$



Axion Dark Matter



Axion, ALPs and Axion Electrodynamics

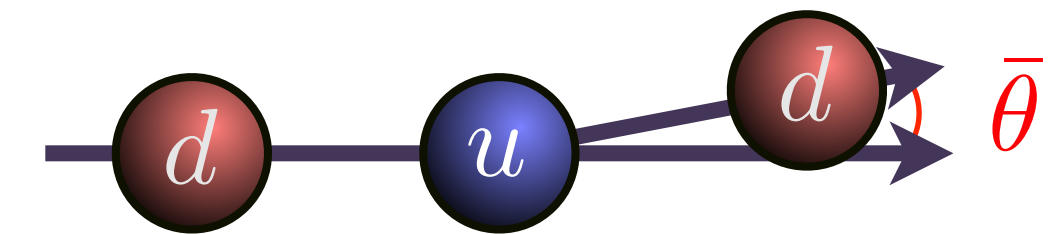
Axion introduced to solve strong CP problem

$$\mathcal{L} \supset \left(\frac{a}{f_a} + \bar{\theta} \right) \frac{g_s^2}{32\pi^2} G_{\mu\nu}^a \tilde{G}^{\mu\nu,a}$$

Peccei & Quinn (1977)
 Weinberg (1978)
 Wilczek (1978)

$$d_n \sim 10^{-16} \bar{\theta} \text{ e cm}$$

$$d_n^{\text{exp}} \lesssim 10^{-26} \text{ e cm}$$



Mixing w/ pion or from full theory:

$$\mathcal{L} \supset -\frac{g_{a\gamma\gamma}}{4} a F \tilde{F} = -g_{a\gamma\gamma} a \mathbf{E} \cdot \mathbf{B}$$

$$\nabla \cdot \mathbf{E} = \rho - g_{a\gamma\gamma} \mathbf{B} \cdot \nabla a$$

$$\nabla \times \mathbf{B} = \partial_t \mathbf{E} + \mathbf{J} - g_{a\gamma\gamma} (\mathbf{E} \times \nabla a - \mathbf{B} \partial_t a)$$

Maxwell's new and improved Equations

Axions and Cavities

FoM:

$$P_{\text{sig}} \sim \omega_{\text{sig}}^2 B_a^2 V \min \left(\frac{1}{\Delta\omega_r}, \frac{1}{\Delta\omega_a} \right)$$

and/or $\mathcal{R} \sim \frac{\Delta\omega_r}{t_{\text{int}}} \text{SNR}^2$

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Maximise: $\omega_{\text{sig}}, B_a, V$

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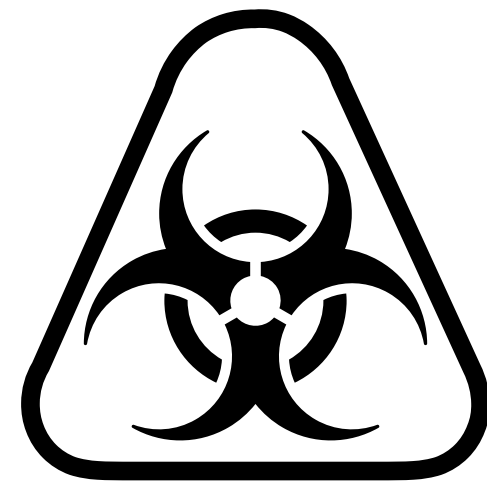
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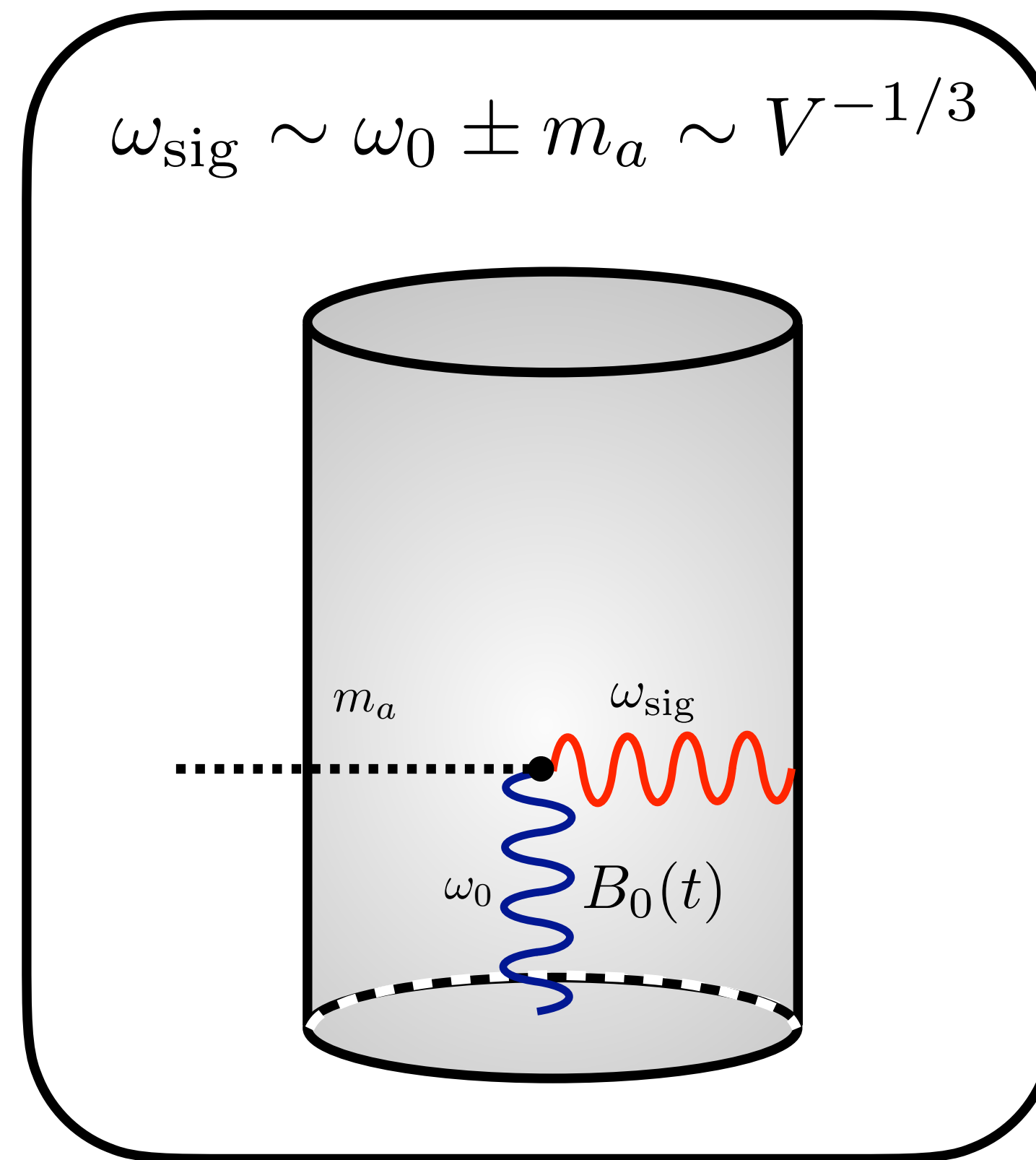


Maximise: $\omega_{\text{sig}}, B_a, V$

Minimise: P_n

New Approach

Heterodyne Superconducting Radio-Frequency Resonator:



Enhances ω_{sig}

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Decouples axion mass from detector volume

New Approach

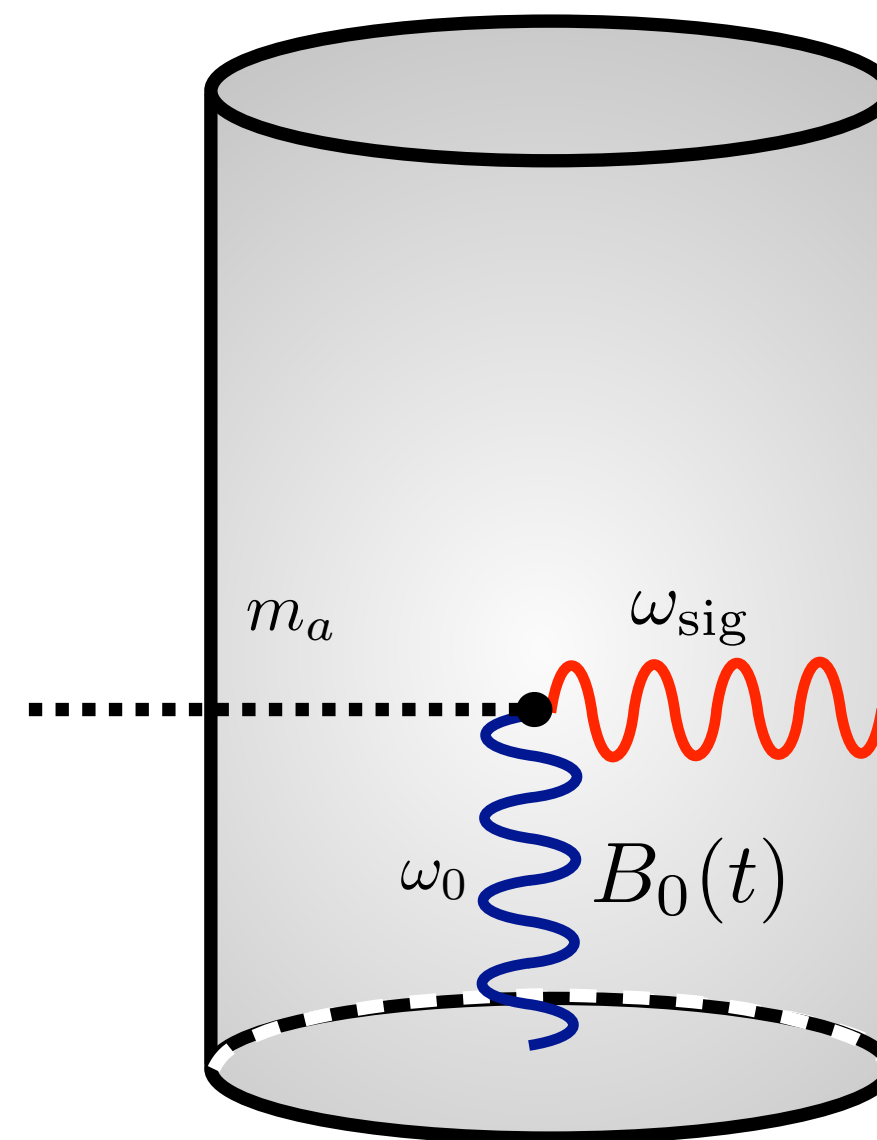
Heterodyne Superconducting Radio-Frequency Resonator:

Q-factor $\gtrsim 10^{10}$

Applications elsewhere,
e.g. Quantum Computing:
see e.g. [quant-ph/2304.09345](https://arxiv.org/abs/quant-ph/2304.09345)

Enhances B_a

$$\omega_{\text{sig}} \sim \omega_0 \pm m_a \sim V^{-1/3}$$



Enhances ω_{sig}

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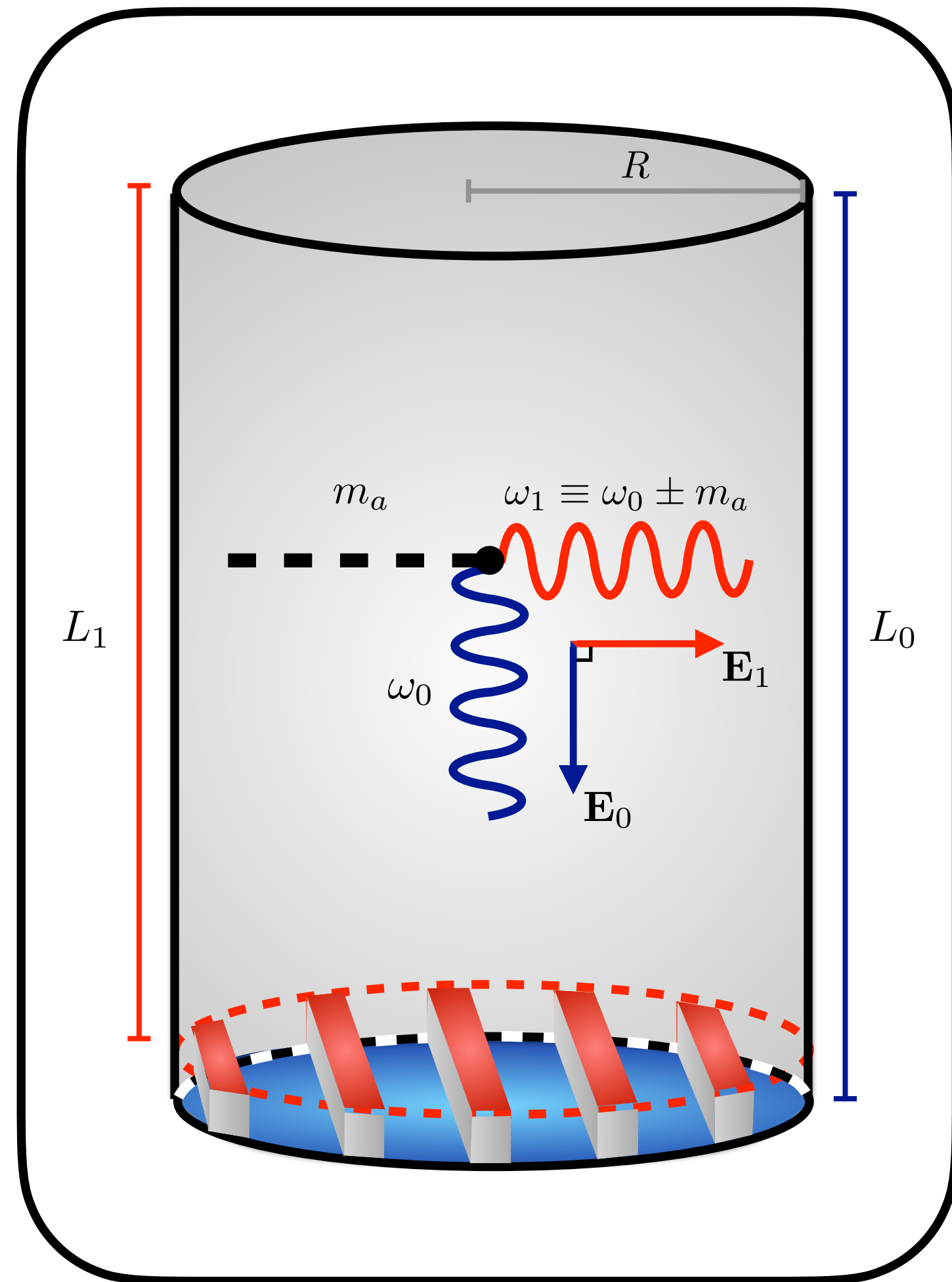
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Heterodyne Resonator

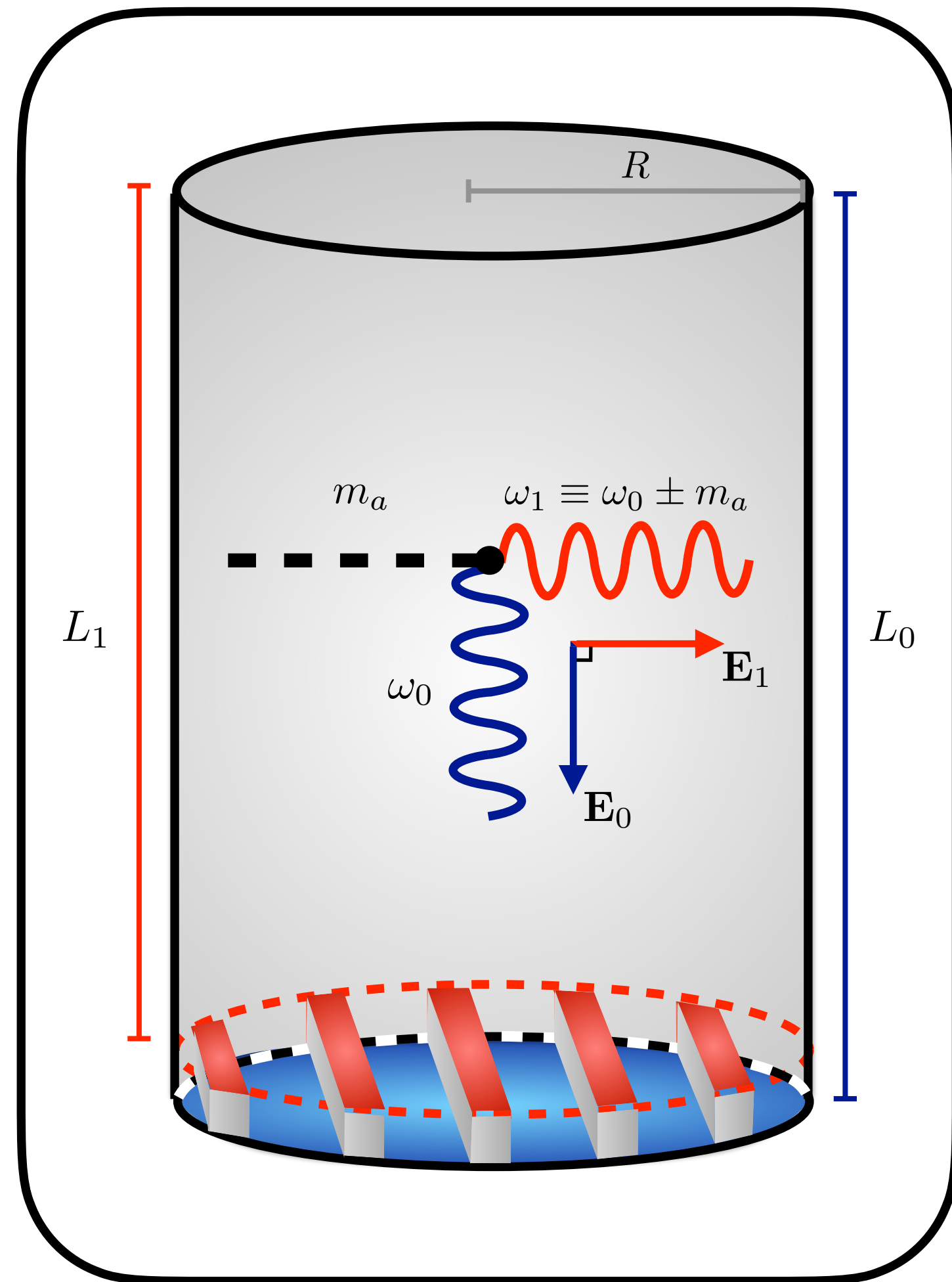


$$\underline{\omega_1 = \omega_0 + \delta\omega}$$

$$\underline{\omega_0 \sim \text{GHz}}$$

$$\underline{m_a \sim \delta\omega}$$

Heterodyne Resonator



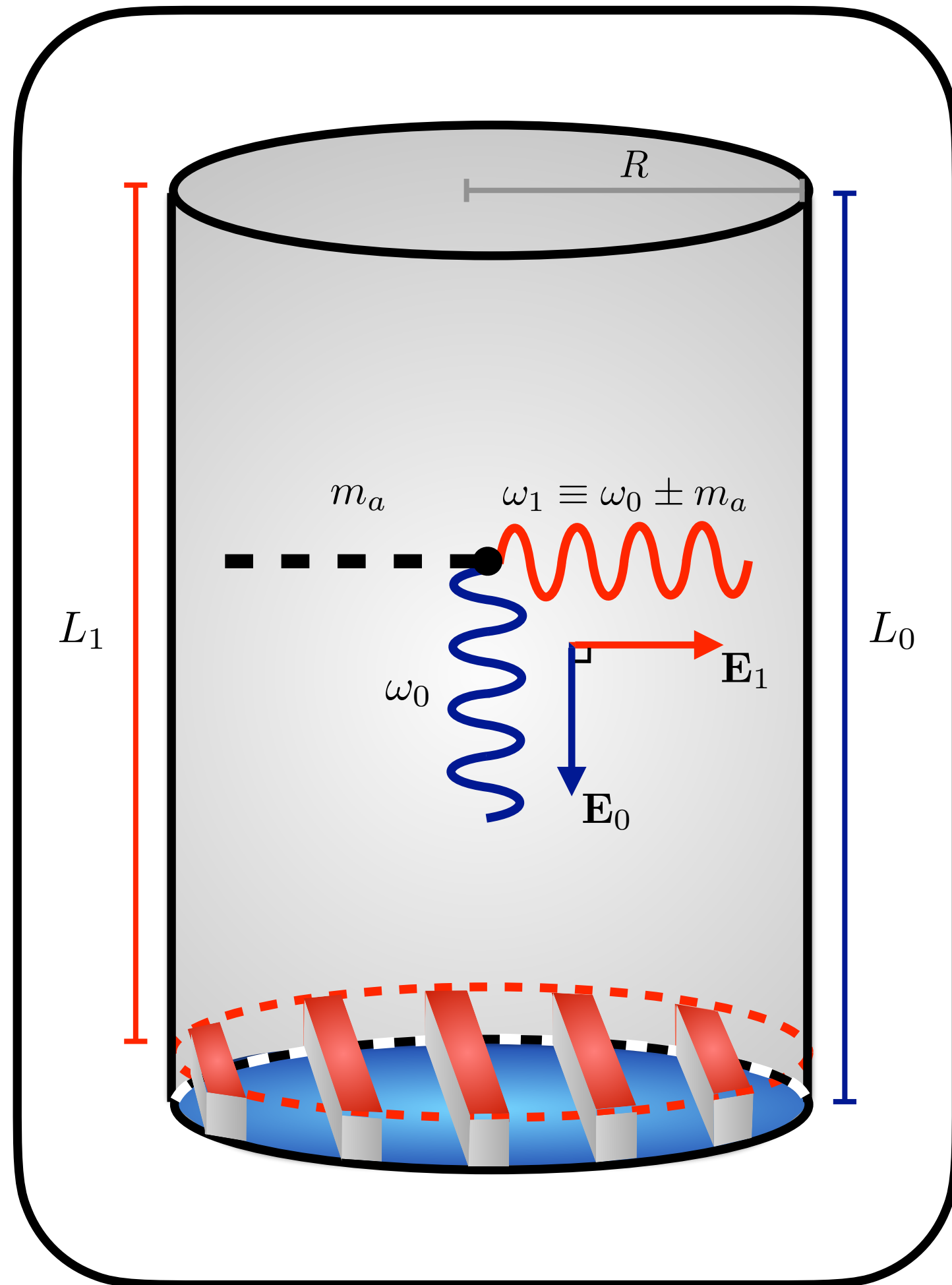
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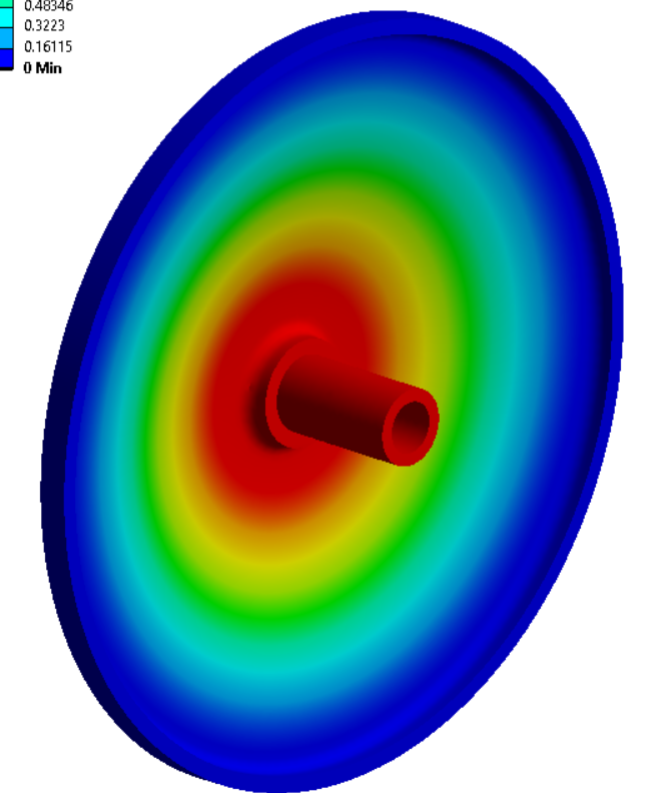
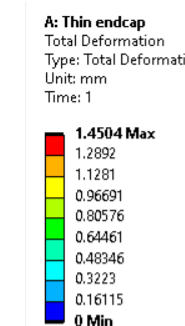
$$m_a \sim \delta\omega$$

Tunability

$$\delta\omega \ll \text{GHz}$$

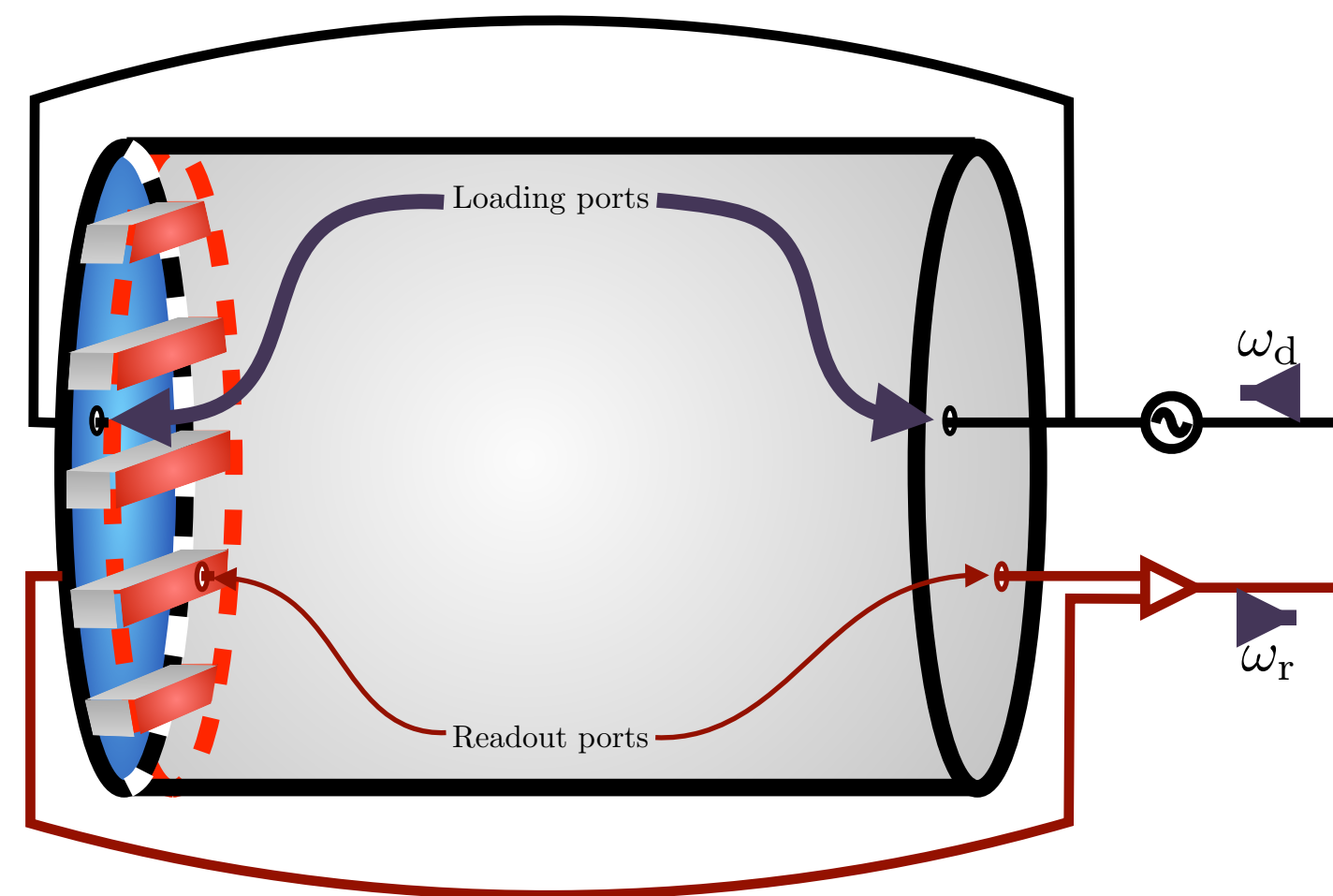
deformability:

~1mm



Courtesy: Marco Oriunno (SLAC)

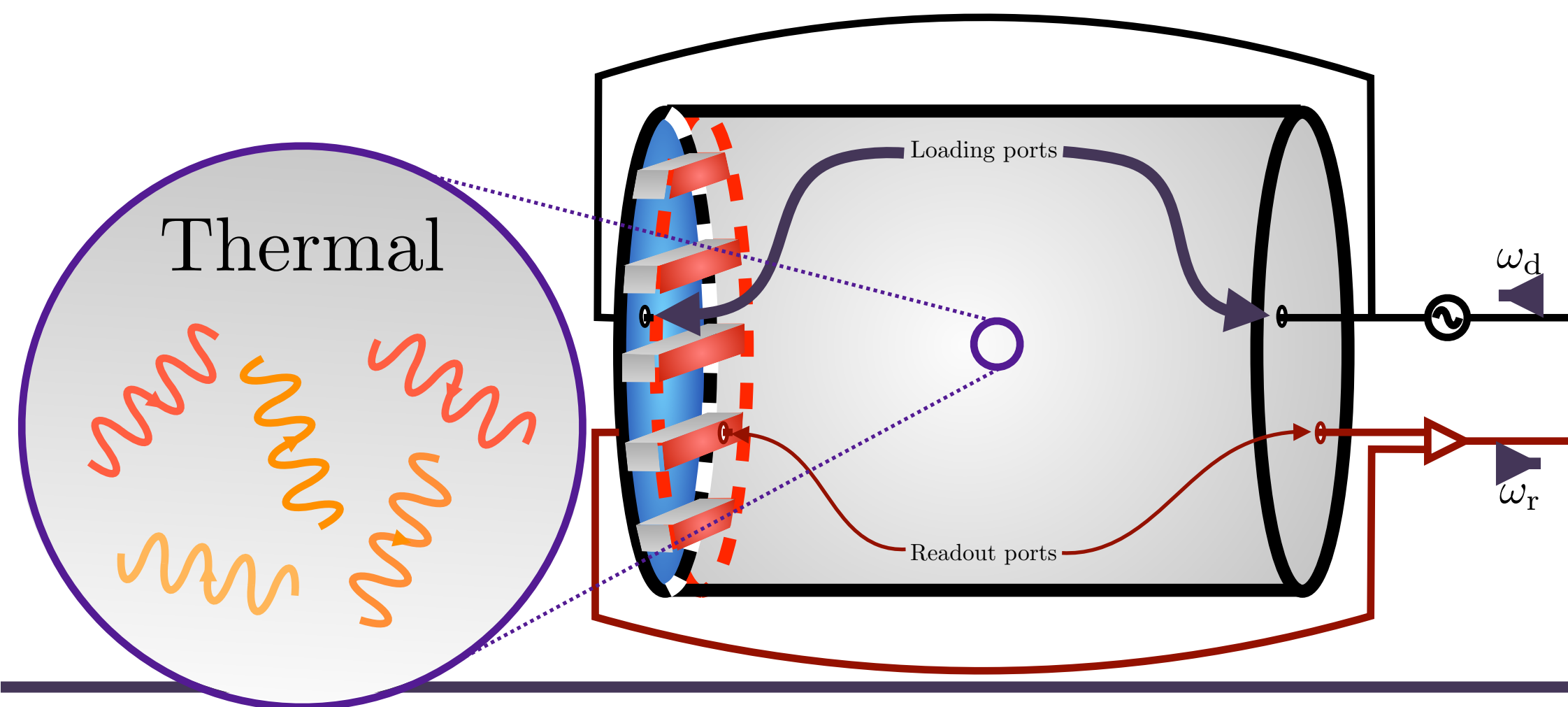
All Noise Sources



See backup for details on each noise source

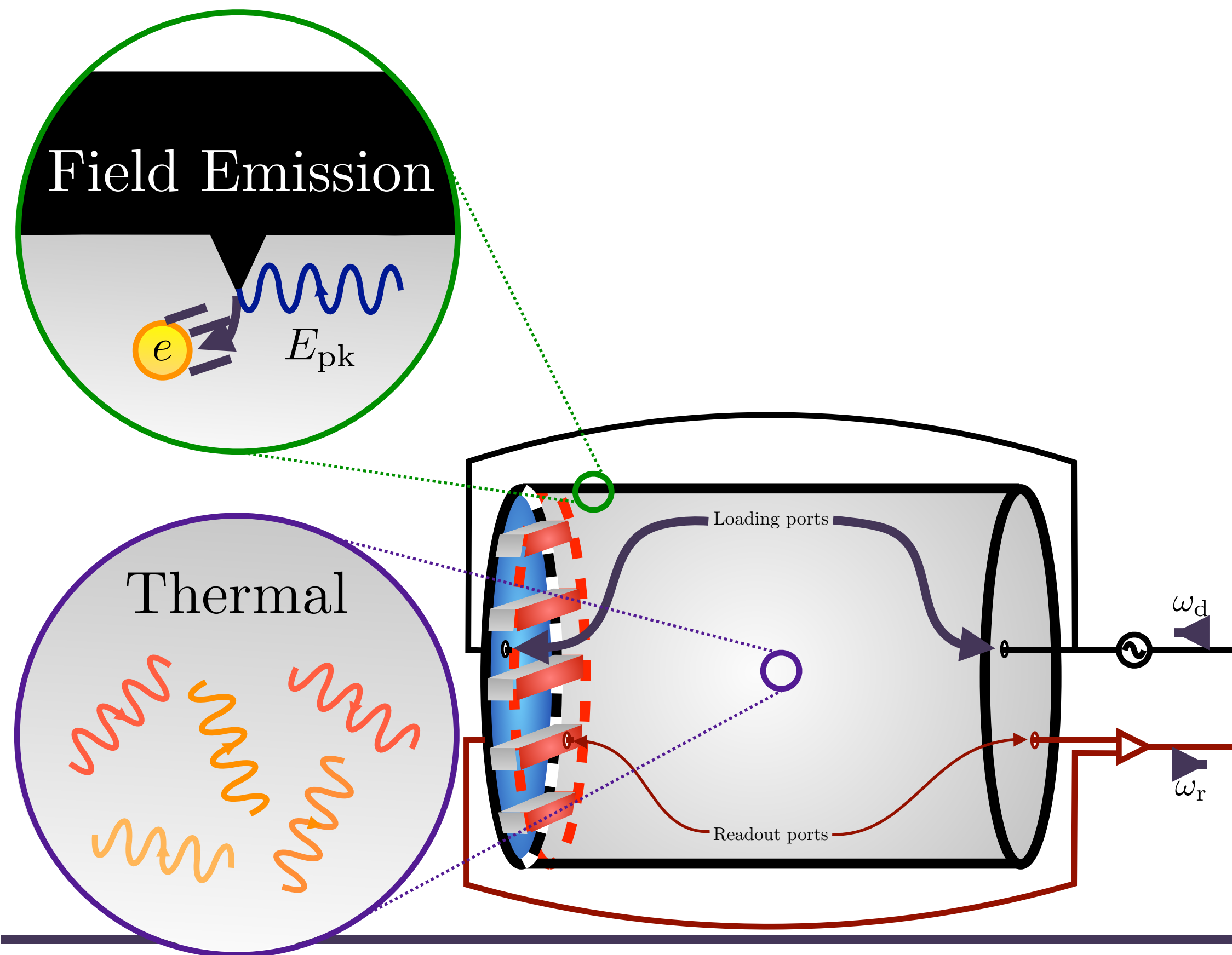
All Noise Sources

- *Thermal noise*: requires cryo



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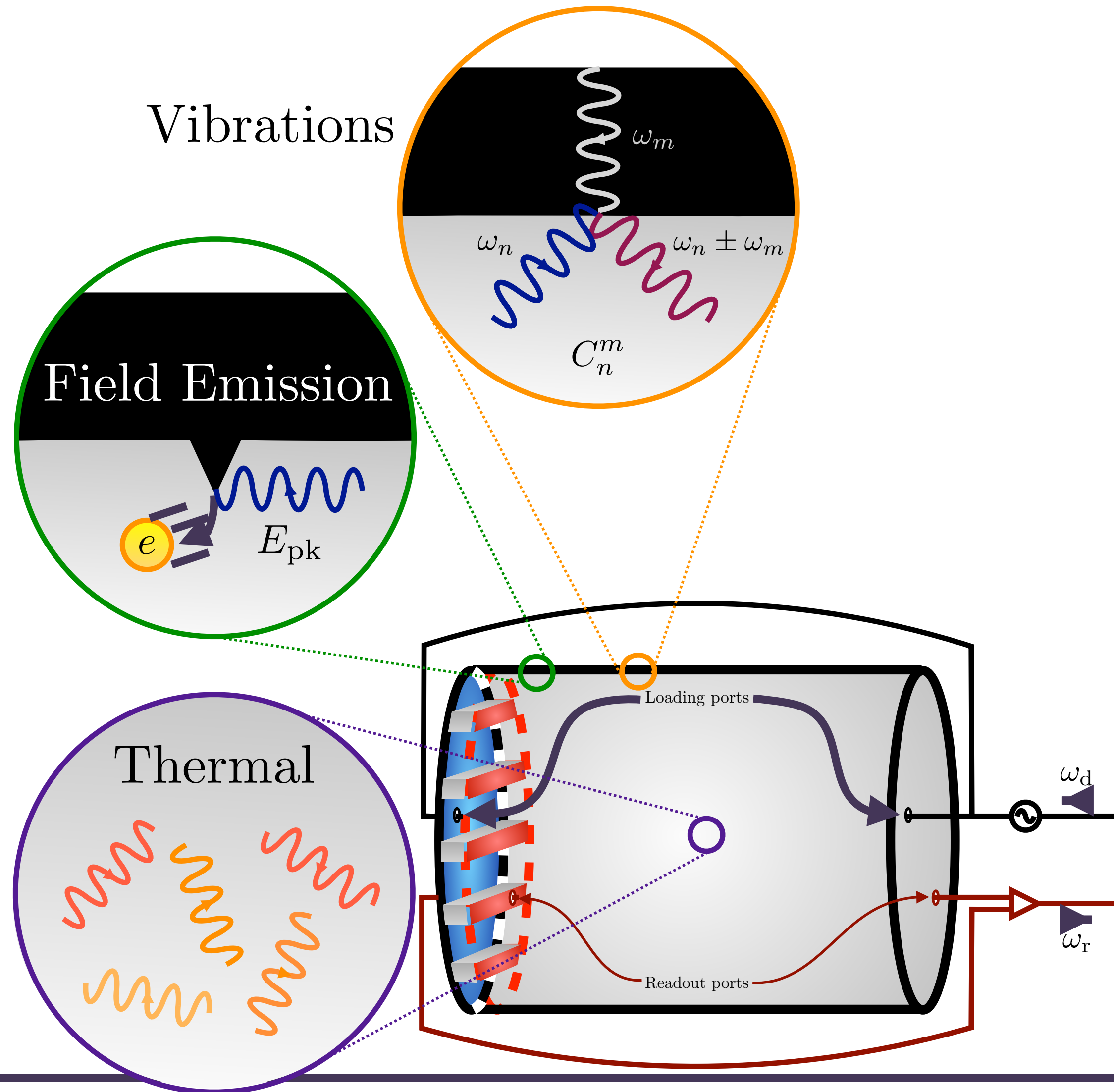
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- *Thermal noise*: requires cryo
- *Field Emission*: careful design & limits peak B-field

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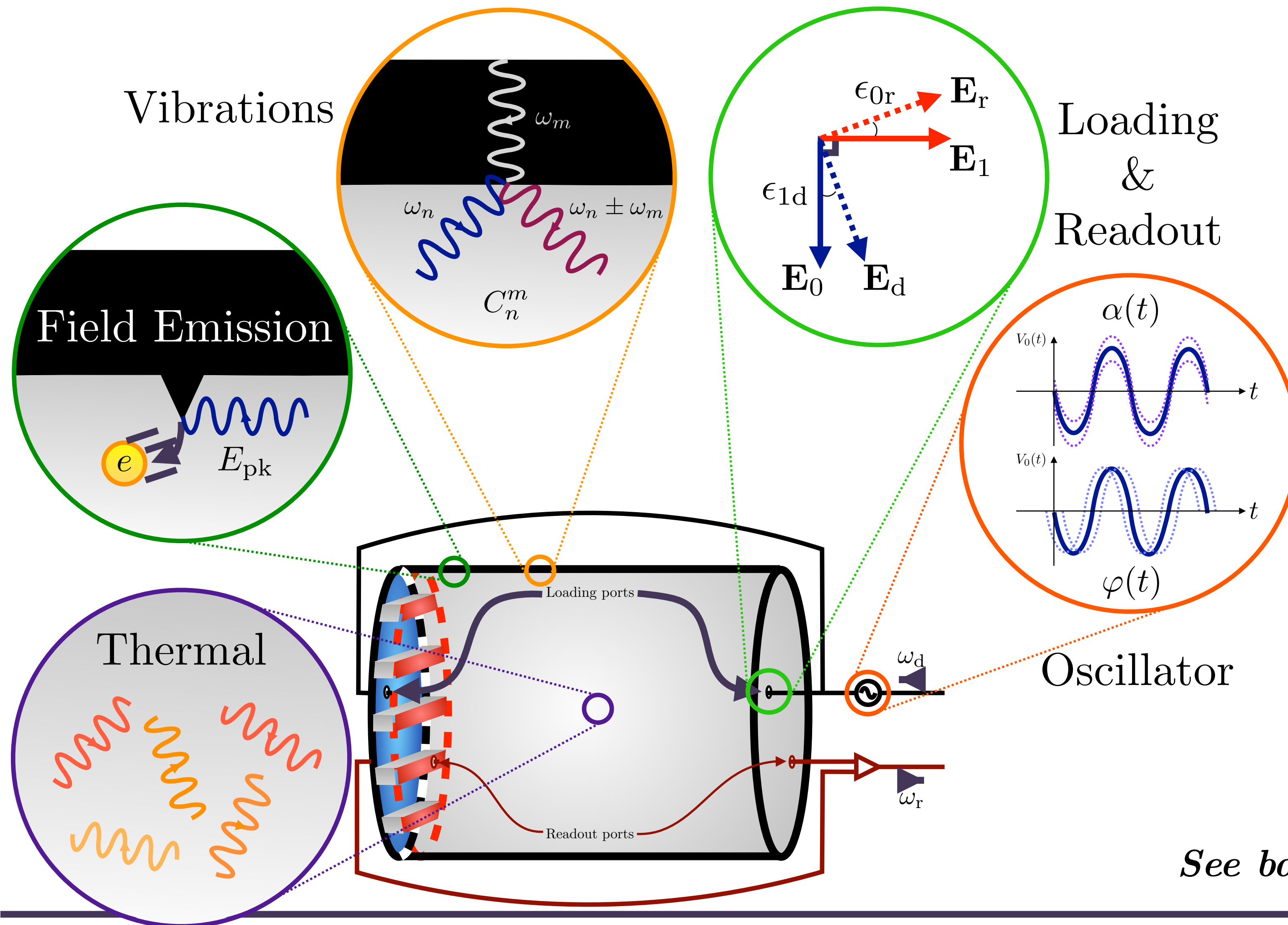
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- *Thermal noise*: requires cryo
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- *Vibrations*: design to reduce microphonics, isolation, cryo

See backup for details on each noise source

All Noise Sources



- **Thermal noise:** requires cryo
- **Field Emission:** careful design & limits peak B-field
- **Vibrations:** design to reduce microphonics, isolation, cryo
- **Loading/Readout & Phase:** design to improve coupling to pump & signal modes. Low phase-noise pump & readout electronics

See backup for details on each noise source

Signal to Noise

Signal to Noise

Thermal noise dominated:

$$\text{SNR} \sim \frac{\rho_{\text{DM}} V}{m_a \omega_1} (g_{a\gamma\gamma} \eta_{10} B_0)^2 \left(\frac{Q_a Q_{\text{int}} t_e}{T} \right)^{1/2}$$

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Comparison with LC resonator:

$$\frac{\text{SNR}}{\text{SNR}^{\text{LC}}} \sim \frac{\omega_0 \pm m_a}{m_a} \left(\frac{Q_{\text{int}}}{Q_{\text{LC}}} \right)^{1/2} \left(\frac{T_{\text{LC}}}{T} \right)^{1/2} \left(\frac{B_0}{B_{\text{LC}}} \right)^2$$

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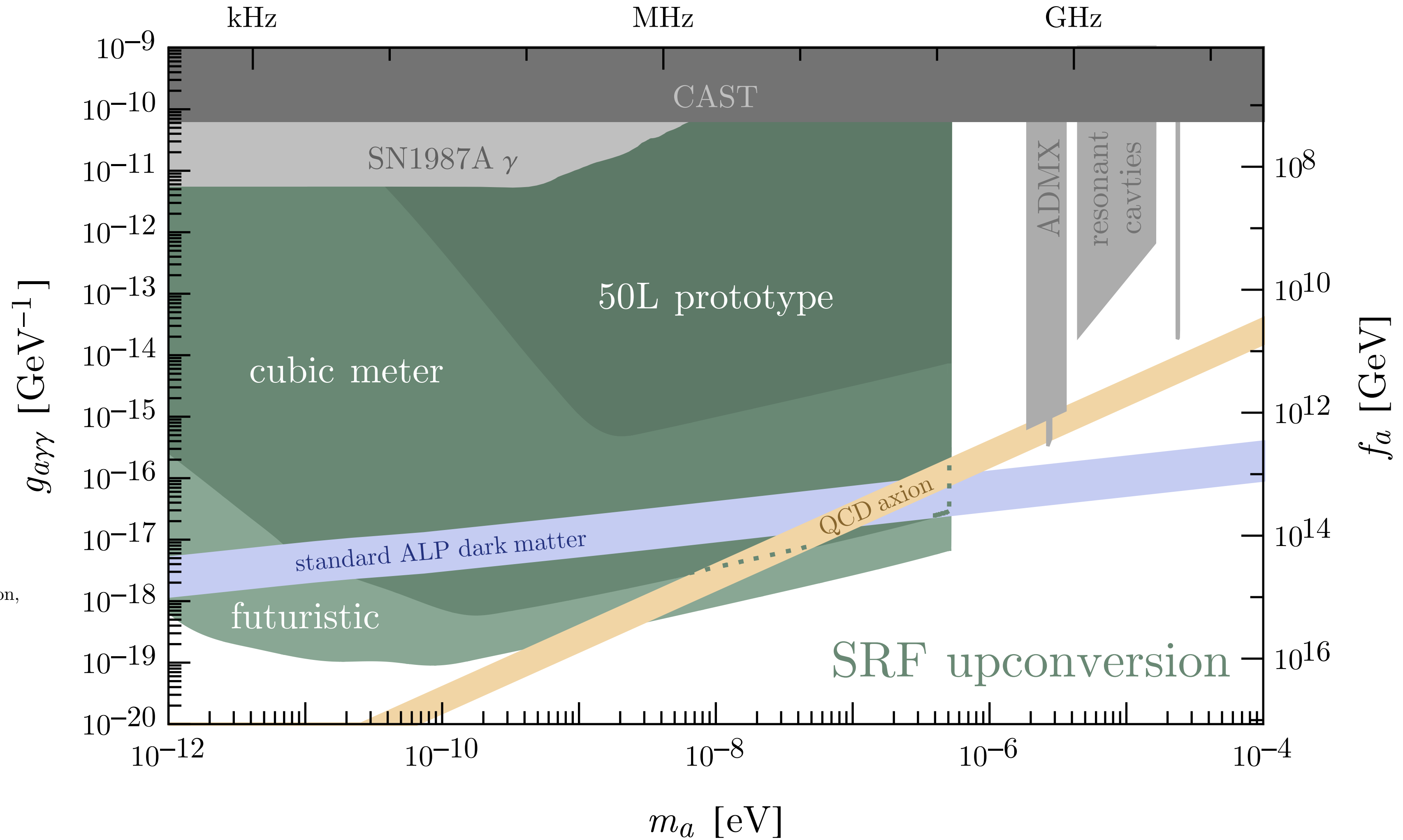
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Sensitivity

$$\text{frequency} = m_a / 2\pi$$



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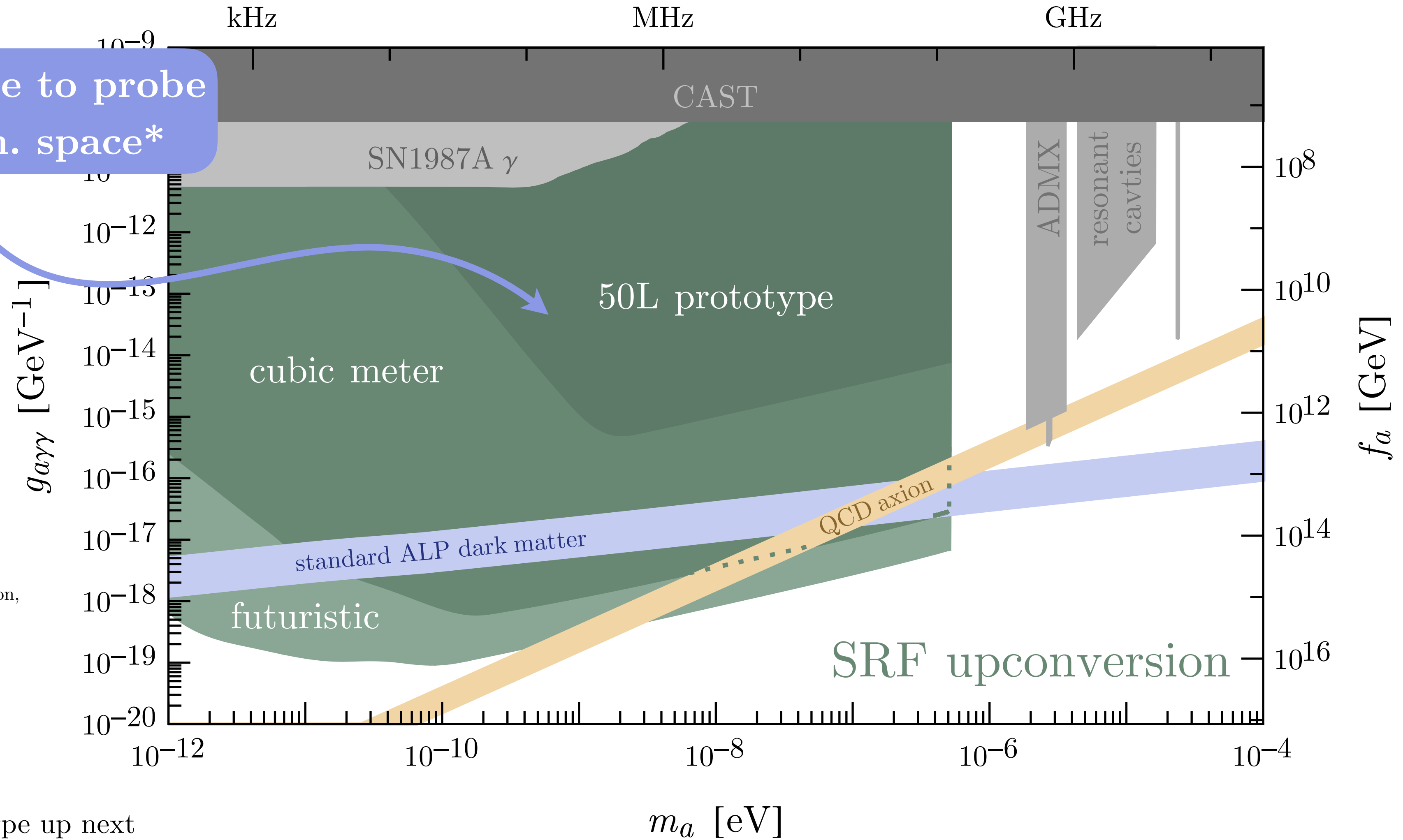
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SC Prototype to probe new param. space*



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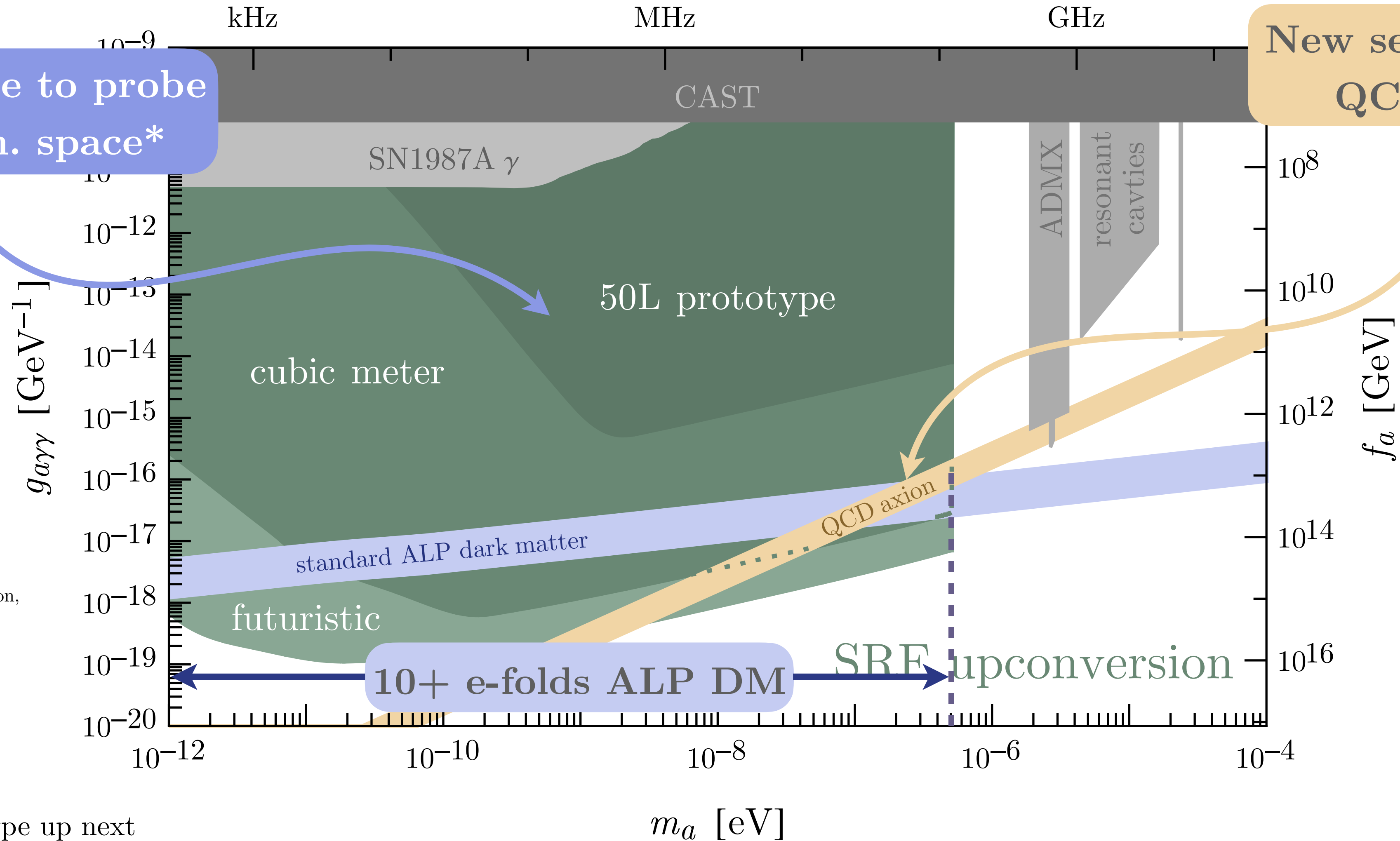
*Status of SC prototype up next

Sensitivity

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SC Prototype to probe new param. space*

New sensitivity to QCD axion



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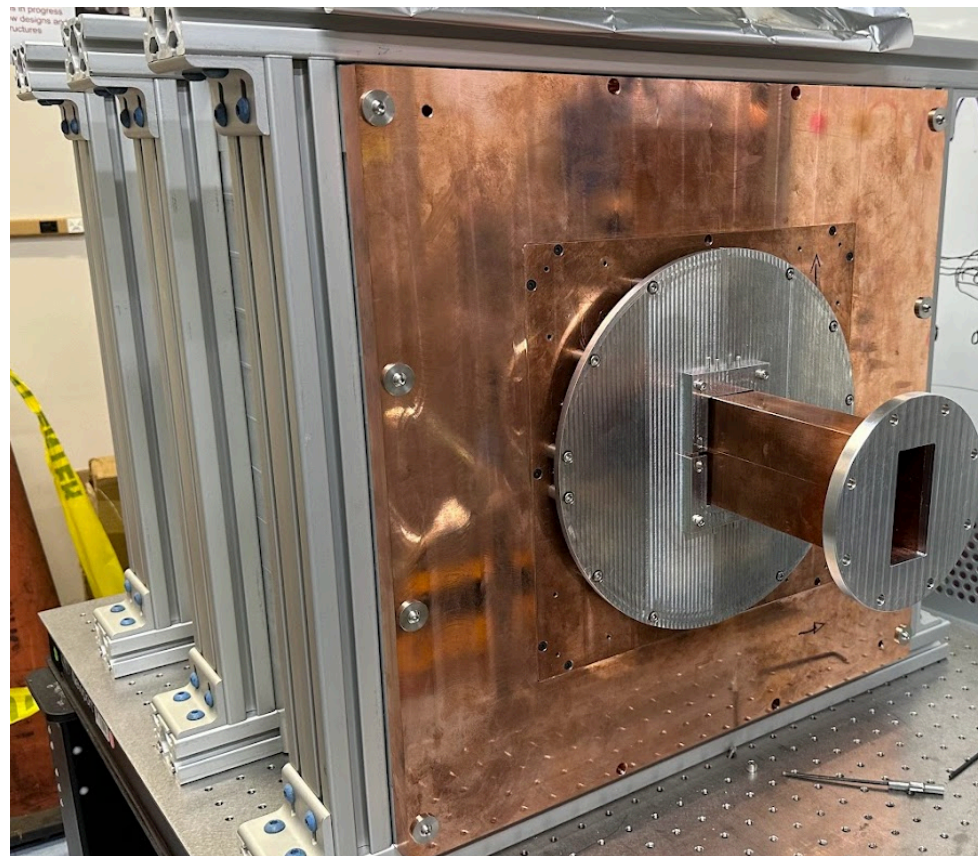
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Current Status

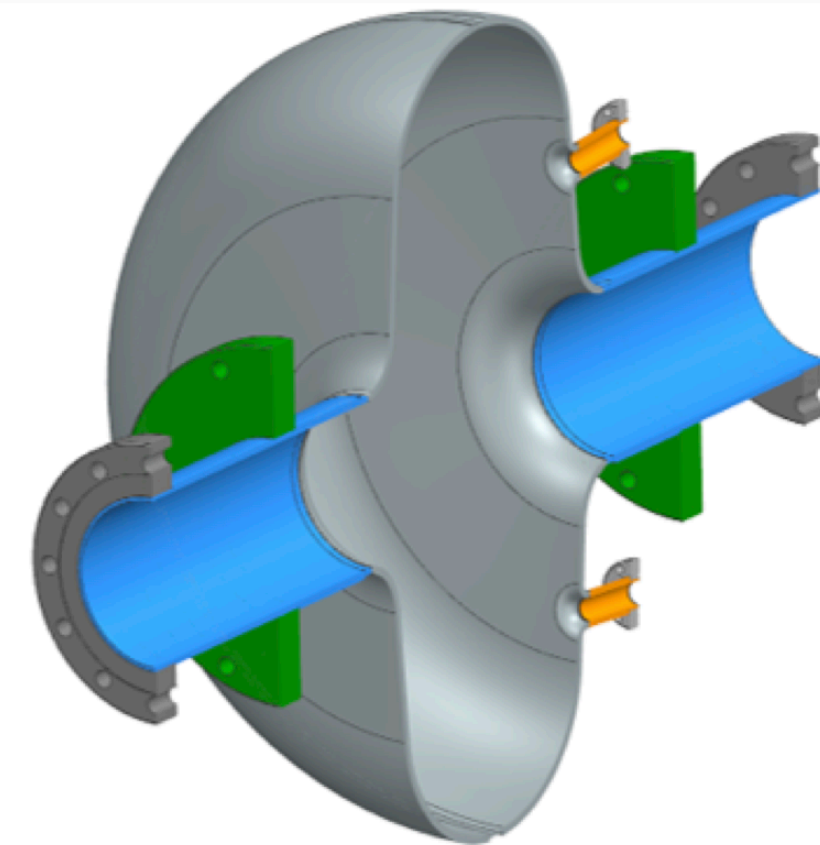
SLAC



LDRD led by S. Tantawi

- Copper prototype
- Input from PBC

FNAL



Giaccone et al [hep-ex/2207.11346]

- Single-cell SRF

CERN



SARE contact w/ A.MacP since 2019

- PBC since ~2021
- QTI since Jan. 2024

See Bianca's talk

Axion Summary

- ▶ Axion: good DM candidate & can solve strong CP
- ▶ Heterodyne detector enables low-mass searches
Ongoing efforts at SLAC, SQMS @FNAL & CERN QTI
- ▶ Noise level above Standard Quantum Limit
- ▶ Quantum techniques promise to improve scan rate

Gravitational Waves

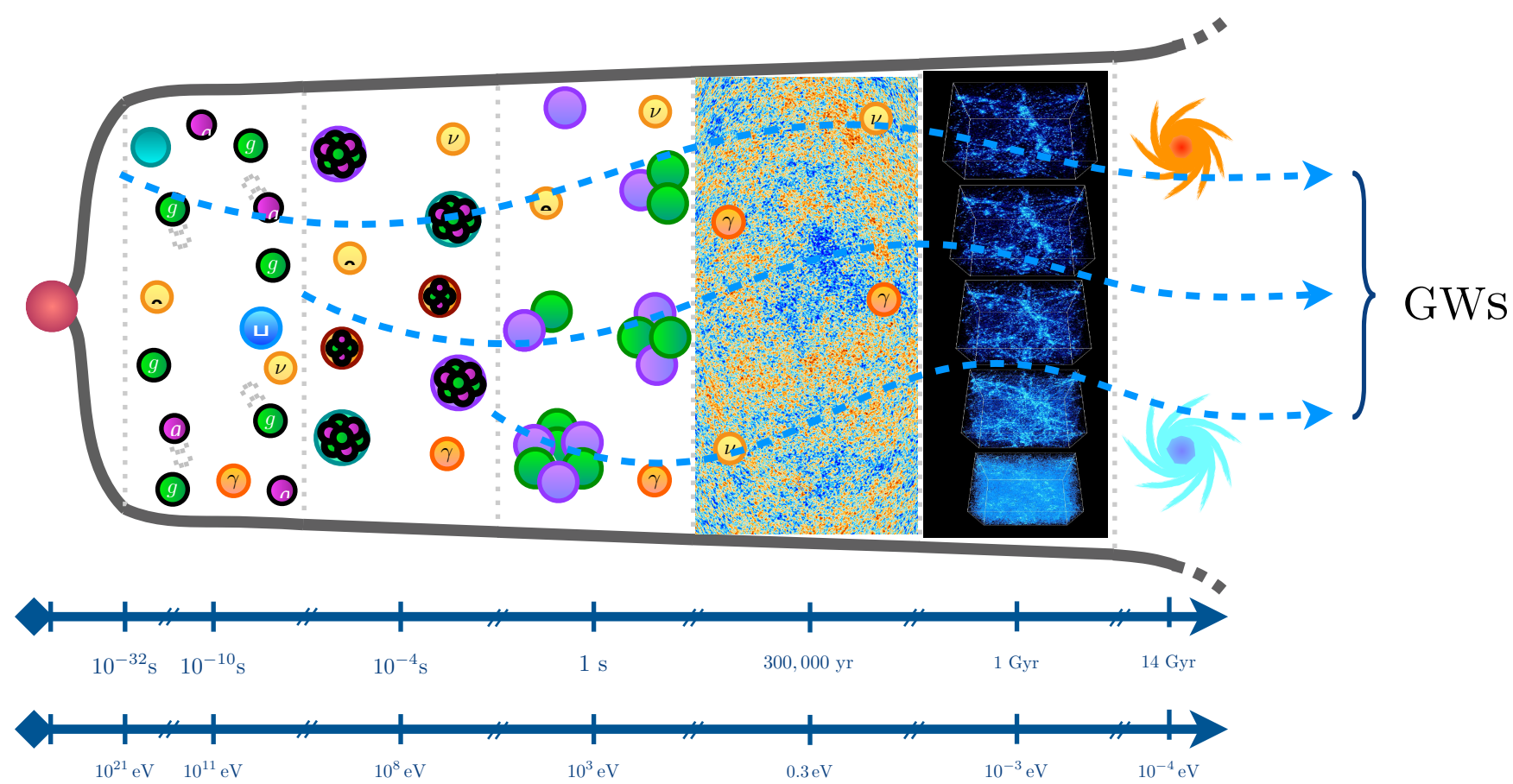


High-Frequency Gravitational Waves

High-Frequency Gravitational Waves

Cosmological GWs

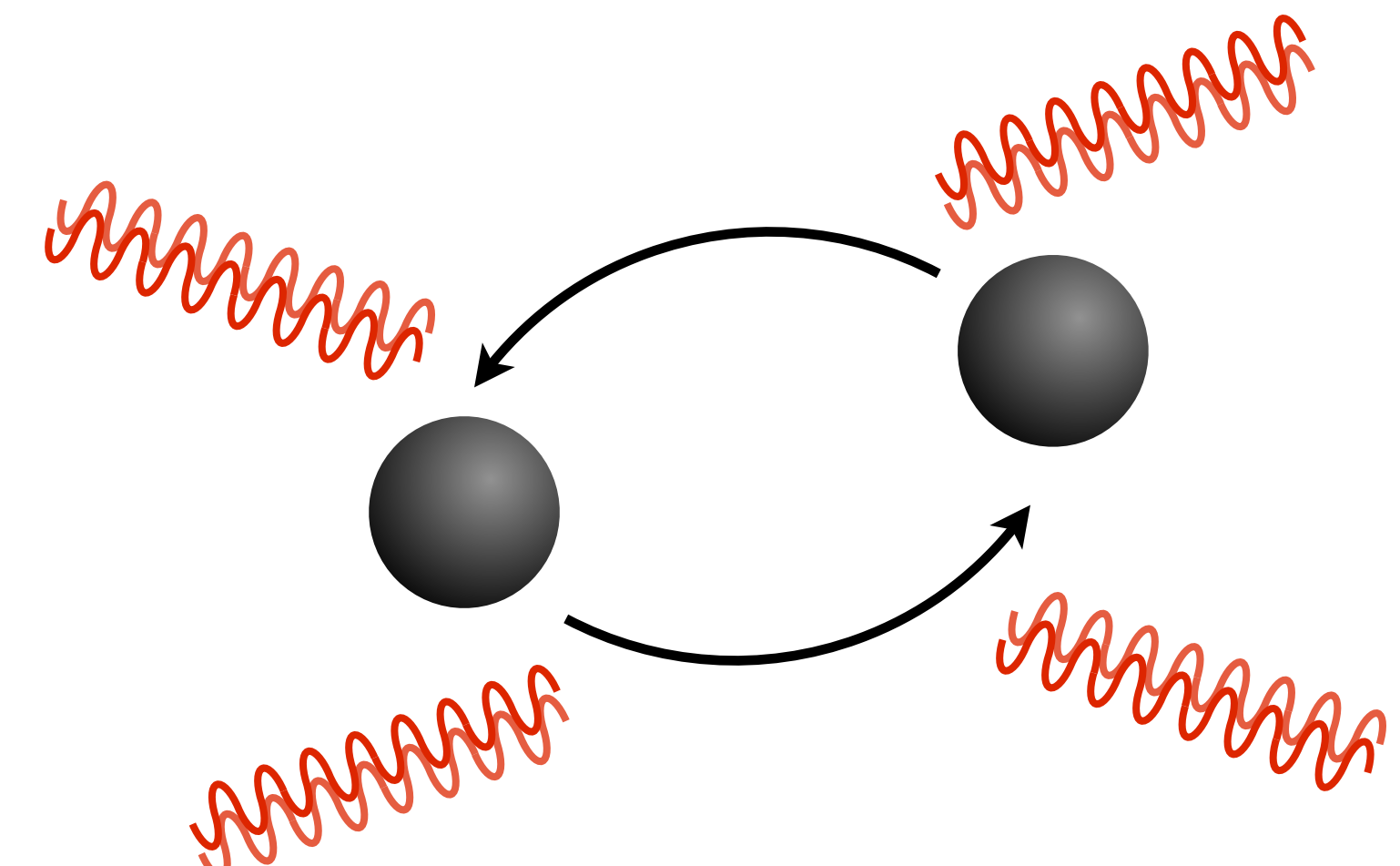
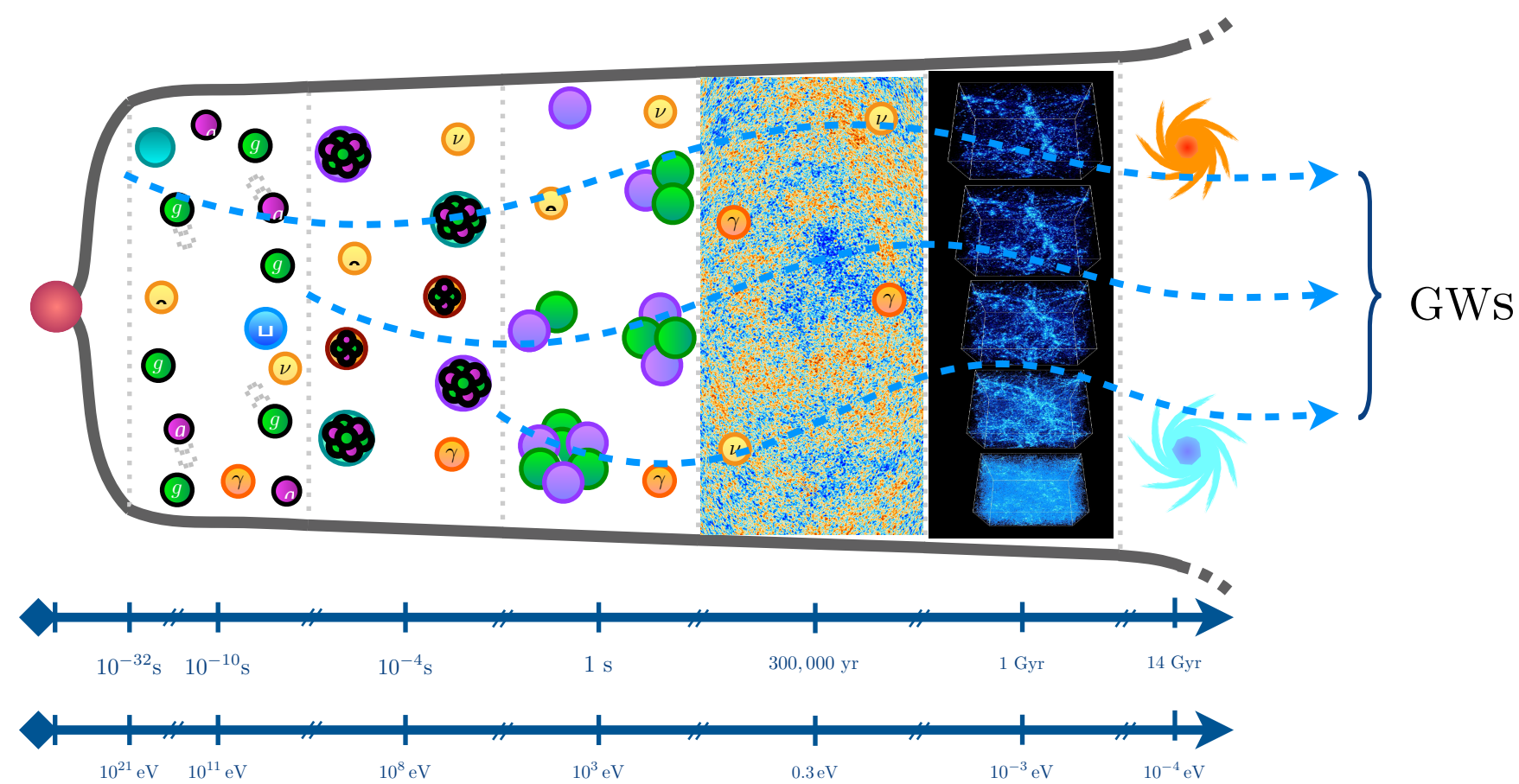
$$\omega_g \Leftrightarrow T_{\text{origin}}$$



High-Frequency Gravitational Waves

Cosmological GWs

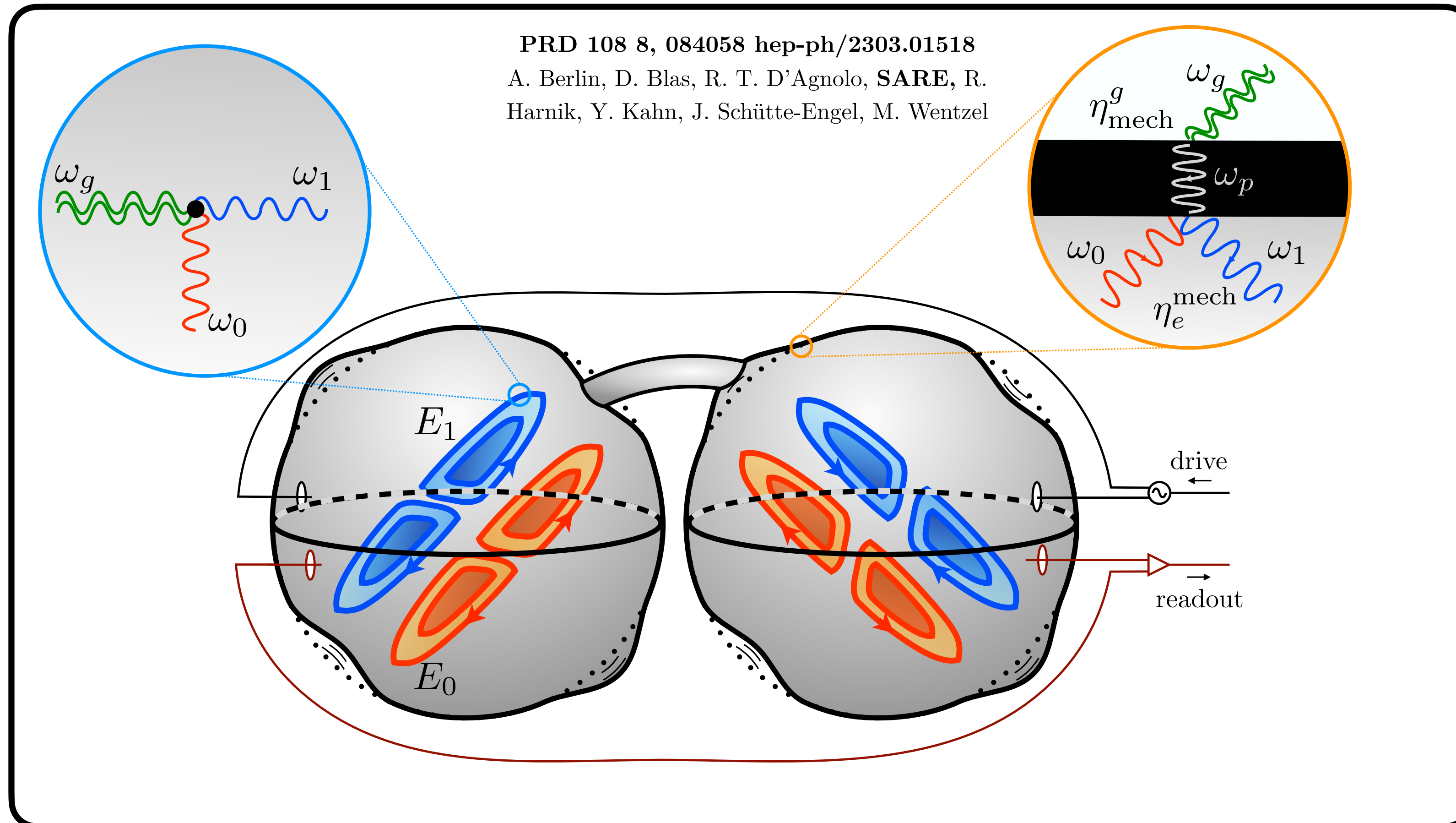
$$\omega_g \Leftrightarrow T_{\text{origin}}$$



Astrophysical GWs

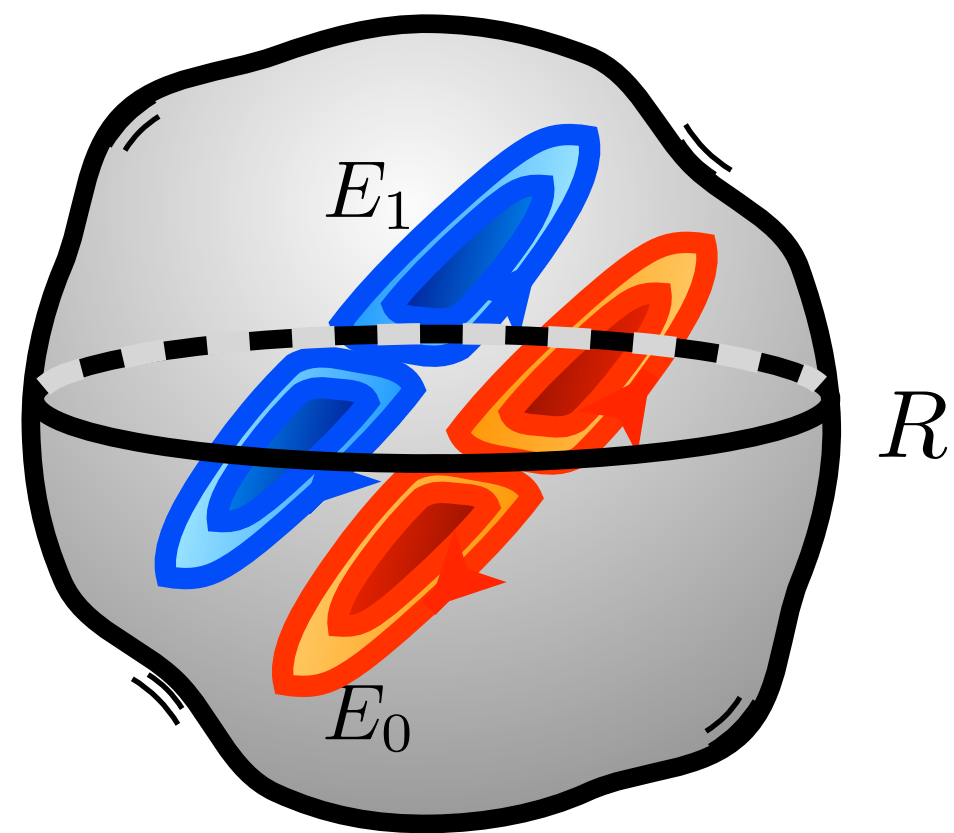
$$\omega_g \Leftrightarrow \Lambda_{\text{origin}}$$

MAGO 2.0

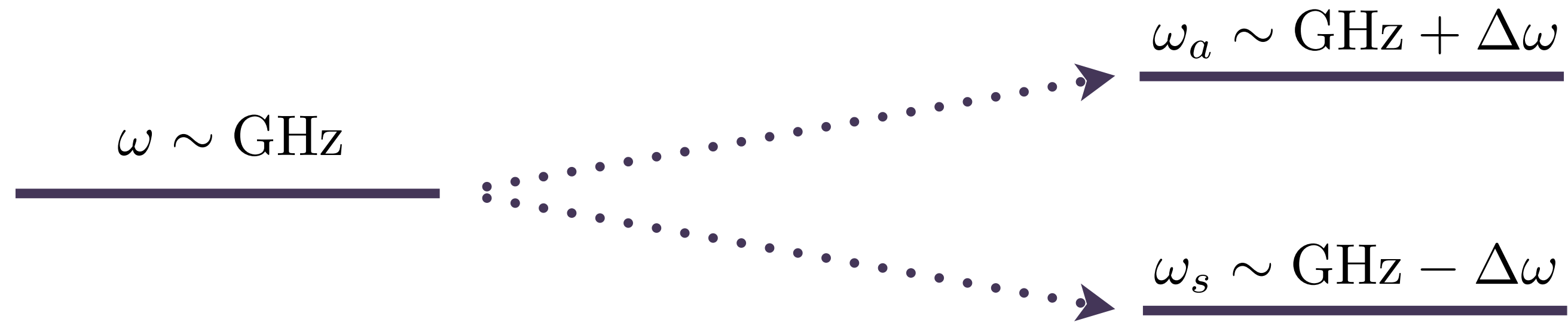


MAGO 2.0

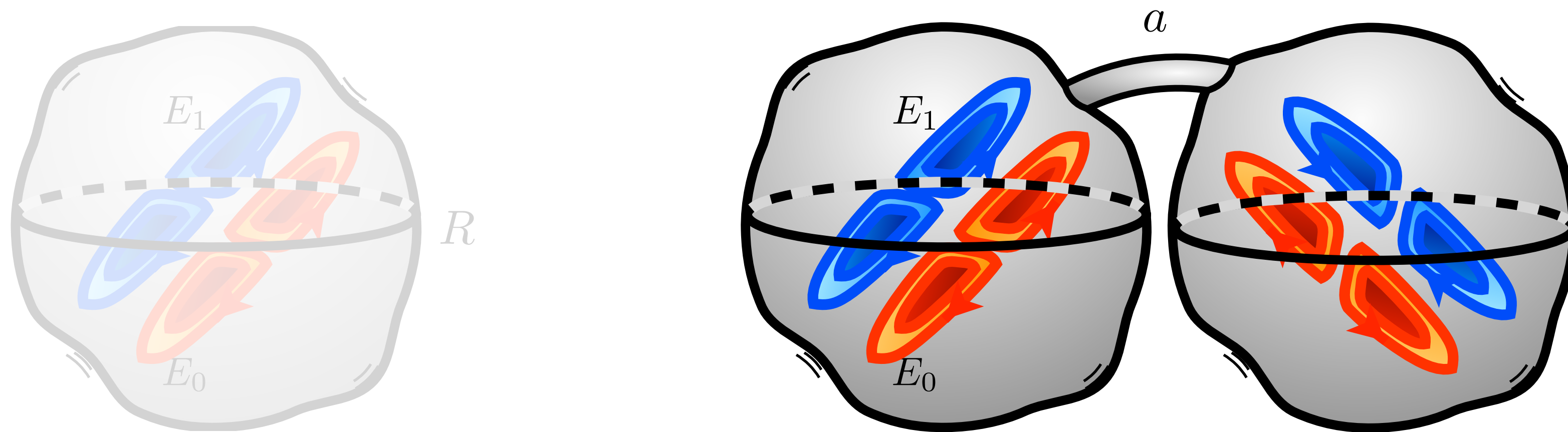
$\omega \sim \text{GHz}$



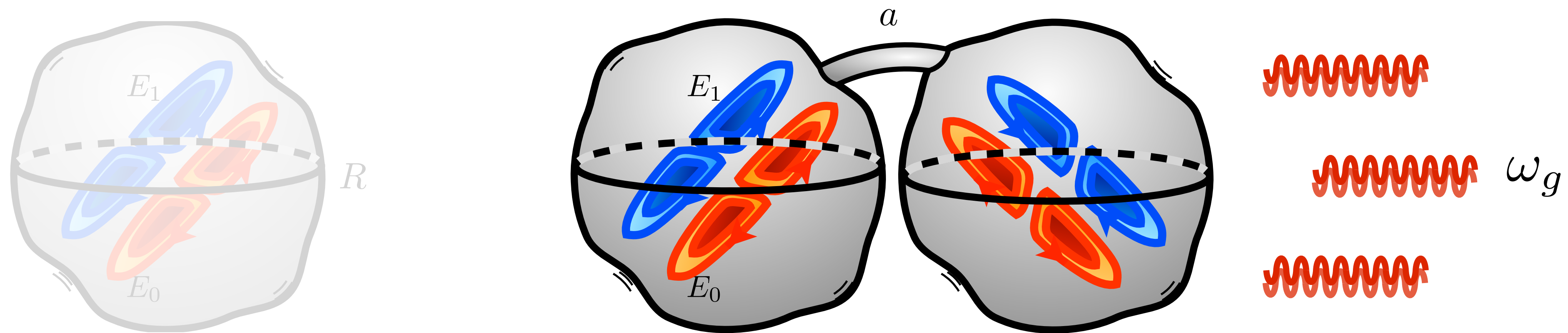
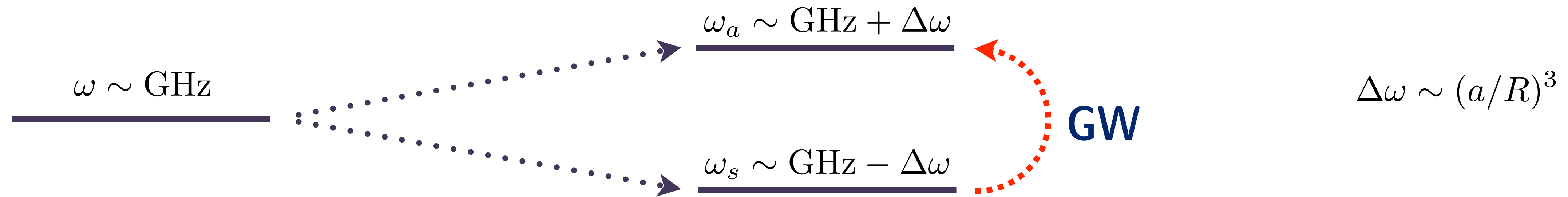
MAGO 2.0



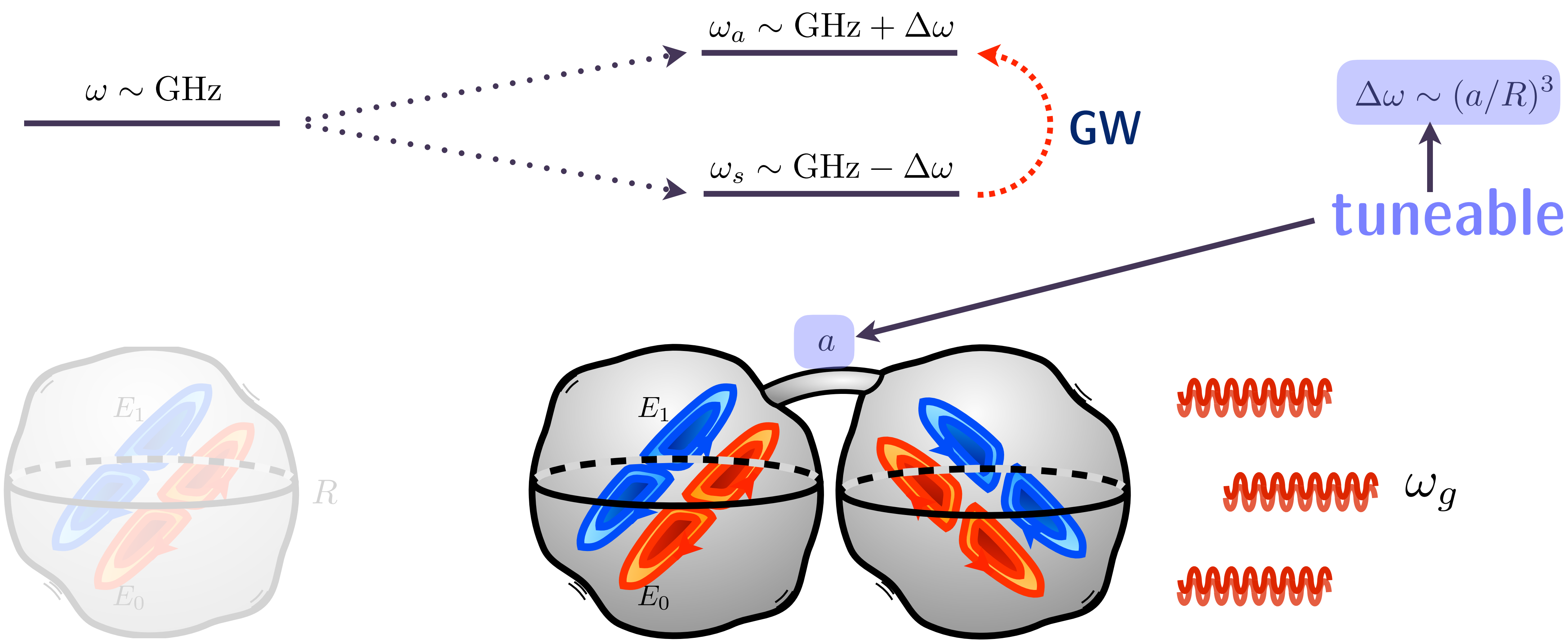
$$\Delta\omega \sim (a/R)^3$$



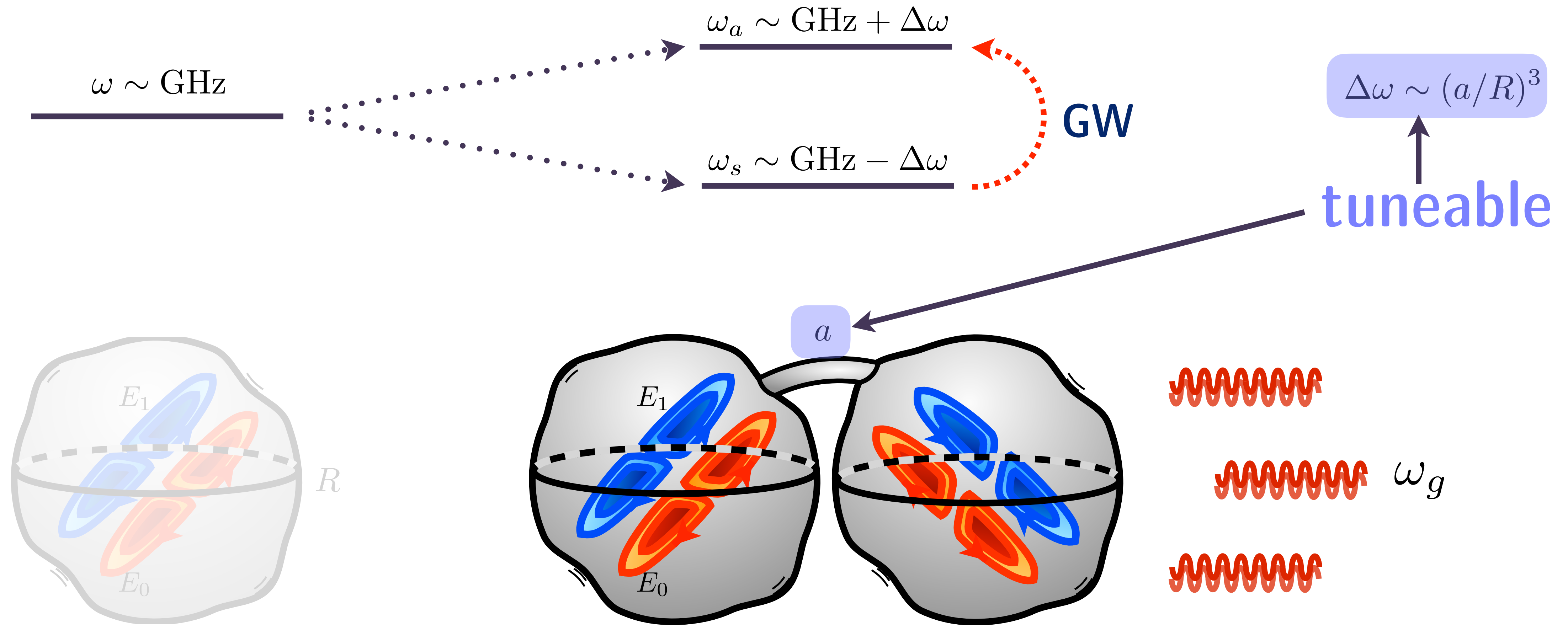
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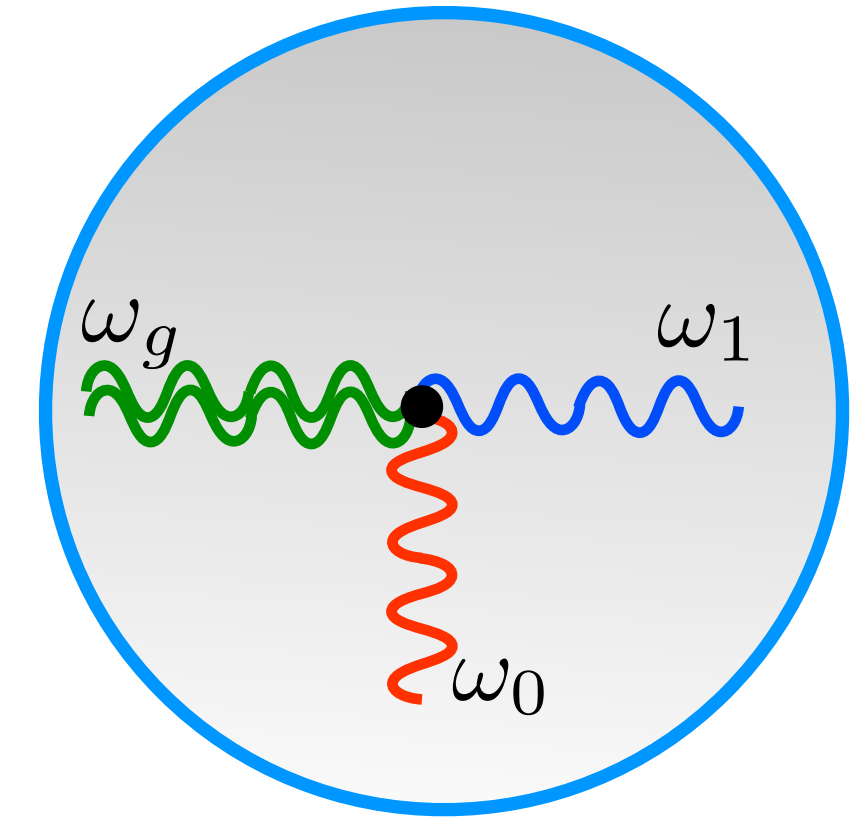


2. Resonant with *beat frequency* ✓

EM and Mechanical signals

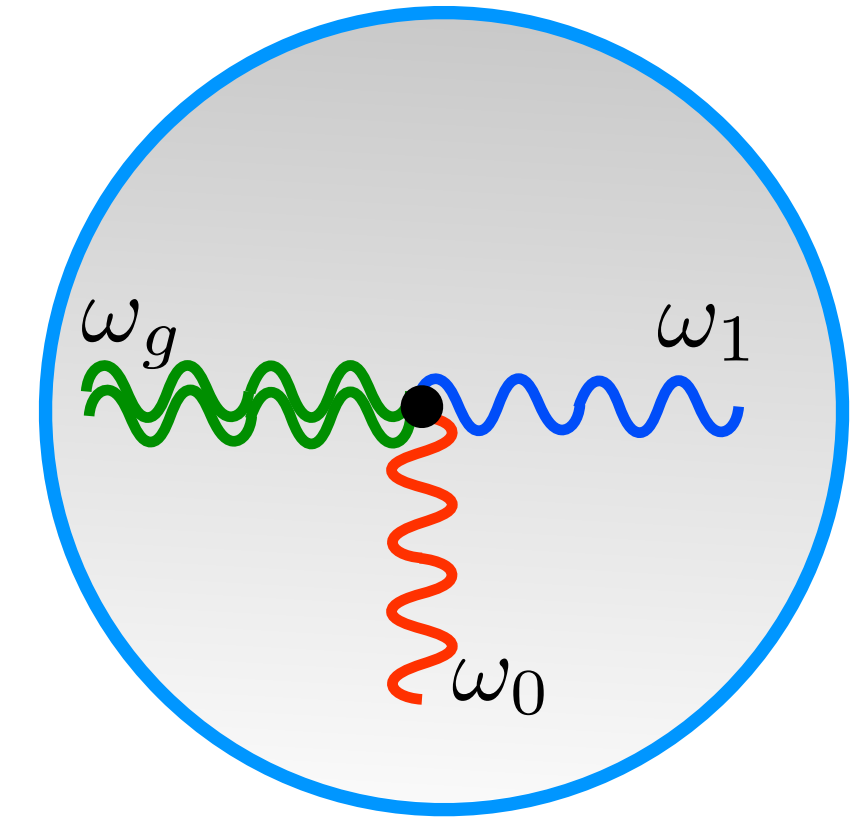
EM and Mechanical signals

Parametrics of the EM signal: $E_{\text{sig}}^{(\text{EM})} \sim Q_{\text{em}} (\omega_g L_{\text{cav}})^2 h^{\text{TT}} E_0$



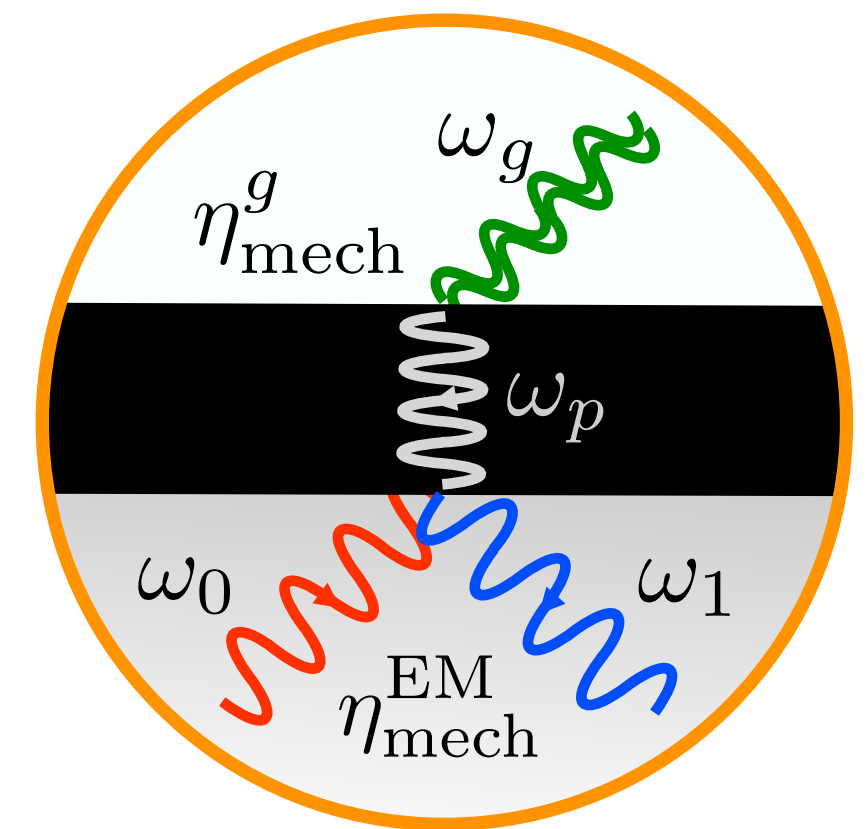
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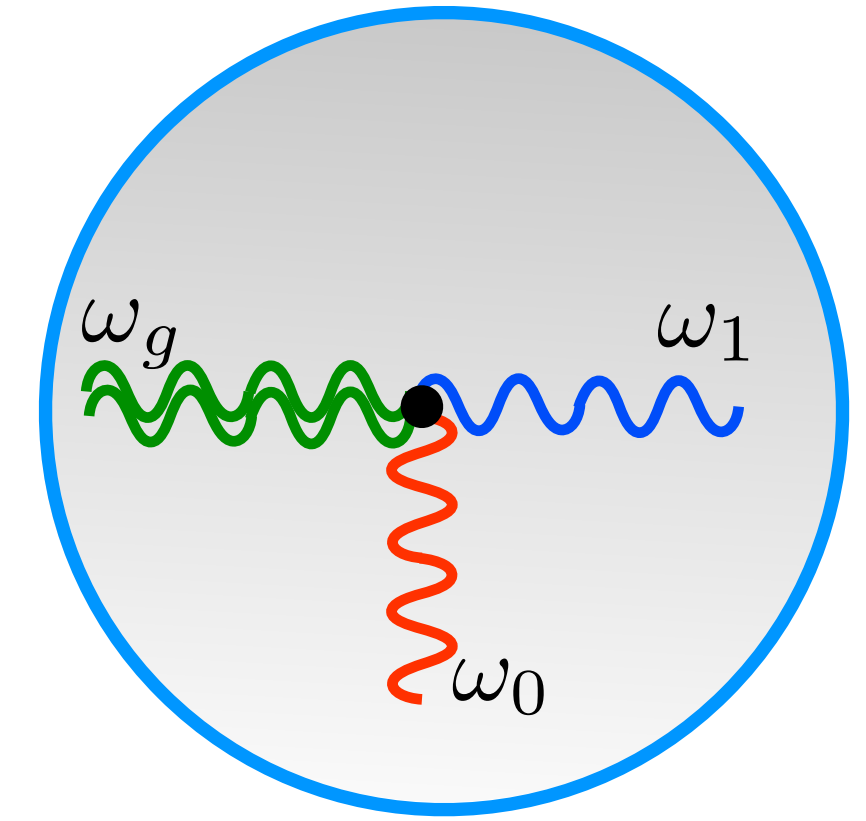
Mechanical signal:

$$E_{\text{sig}}^{(\text{mech})} \sim Q_{\text{em}} h^{\text{TT}} E_0 \min \left(1, \frac{\omega_g L_{\text{cav}}}{c_s} \right)^2$$



EM and Mechanical signals

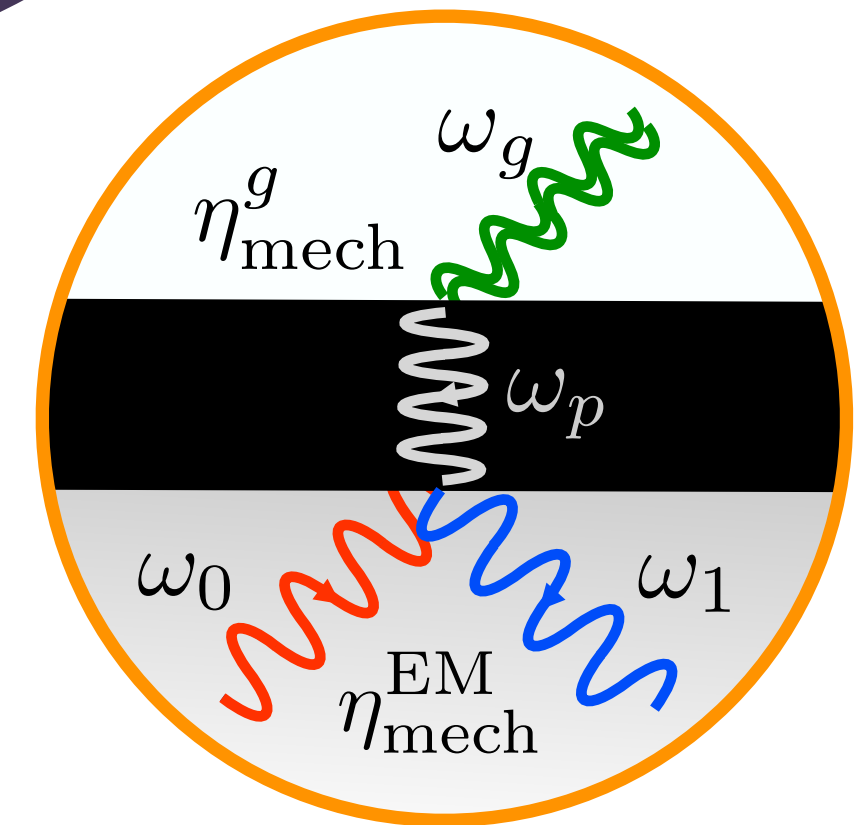
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Enhanced by $1/c_s^2 \gg 1$ (!)

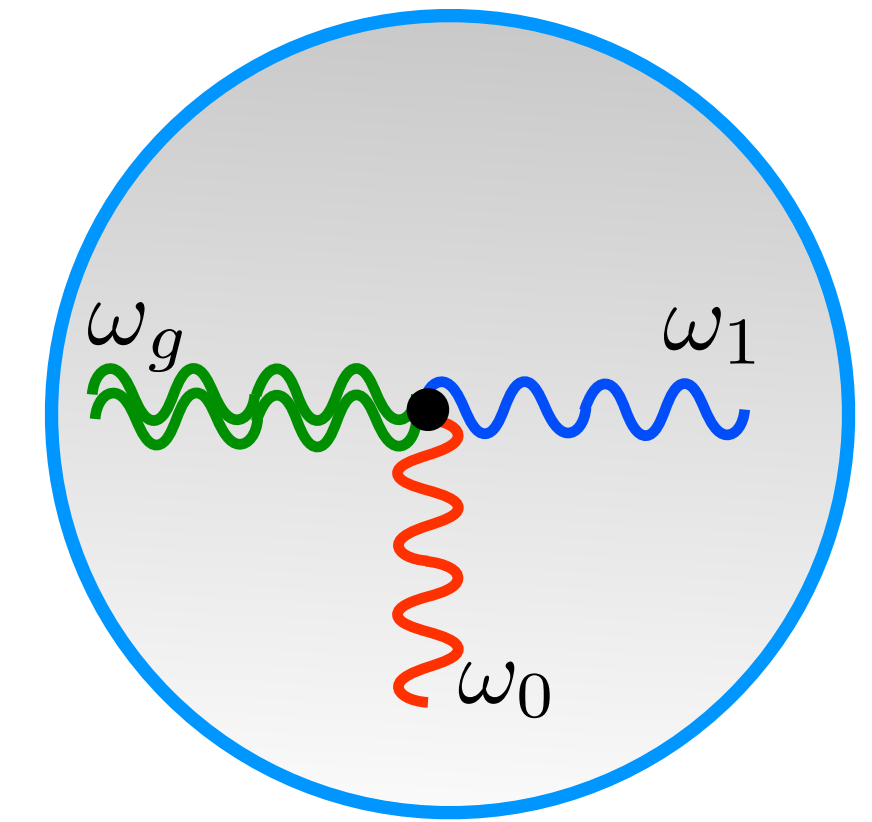
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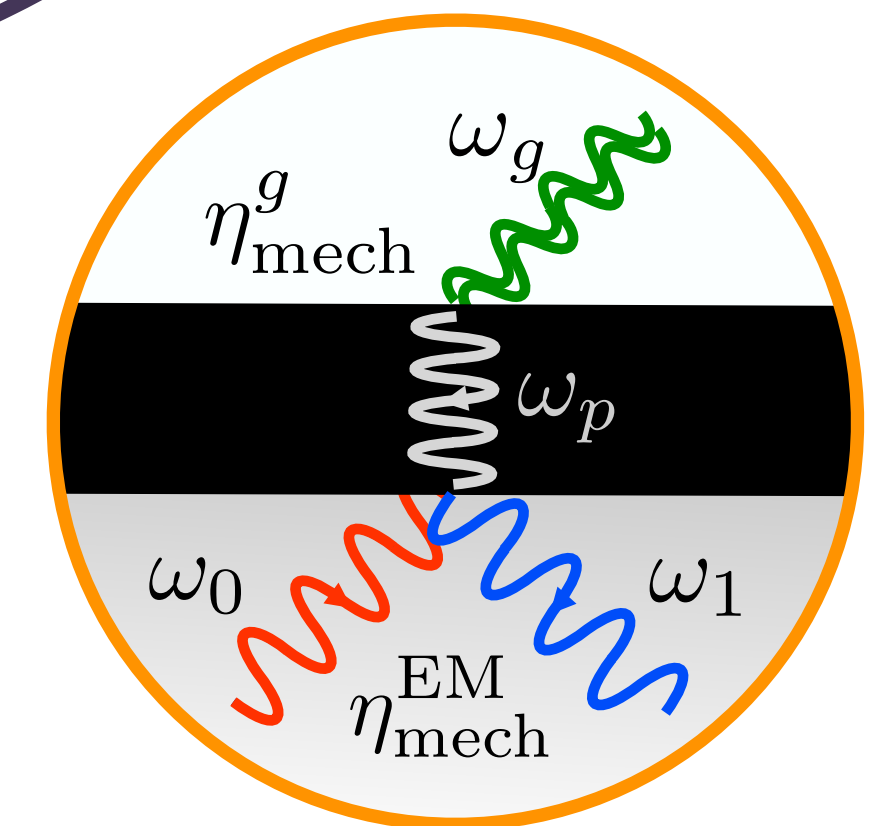


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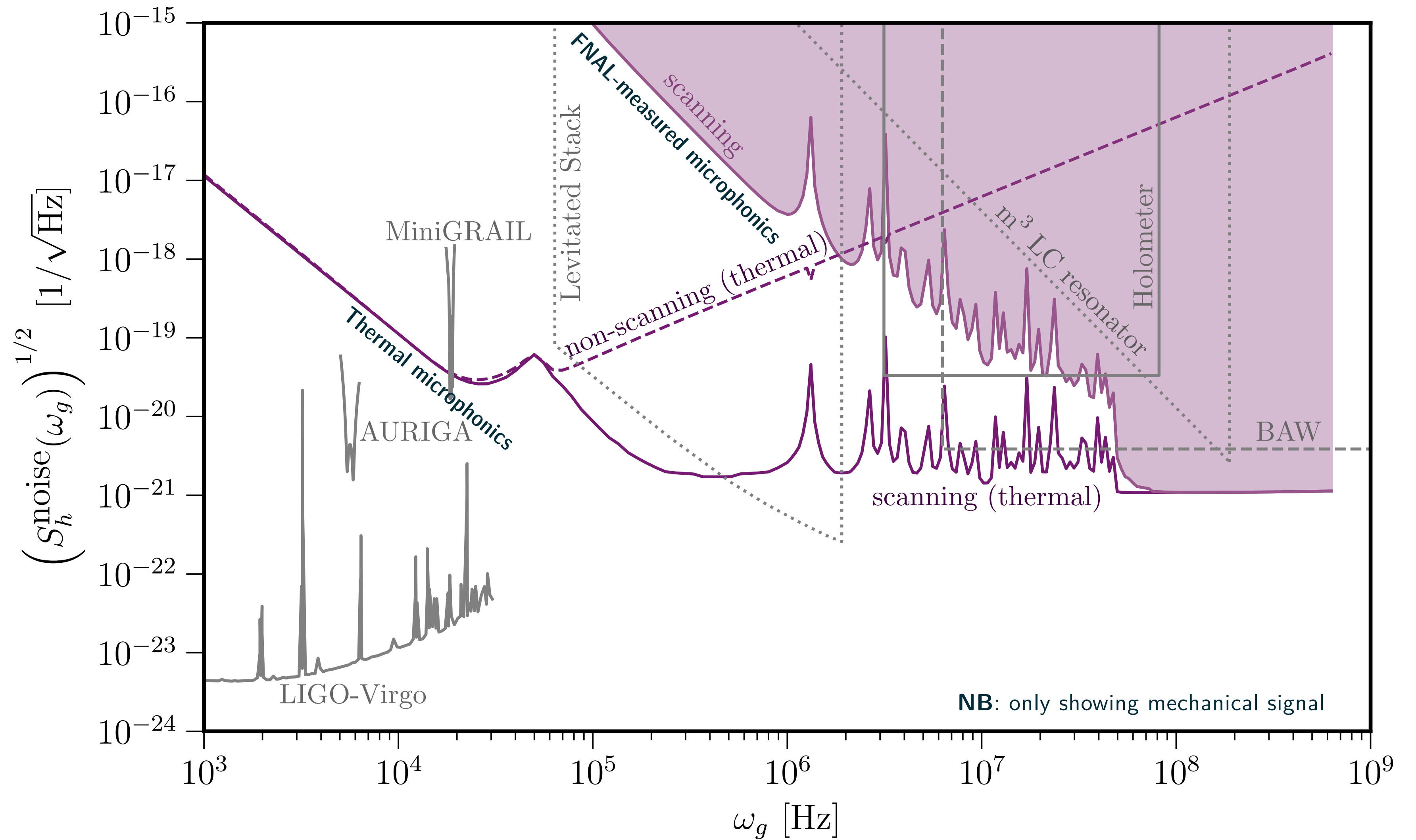
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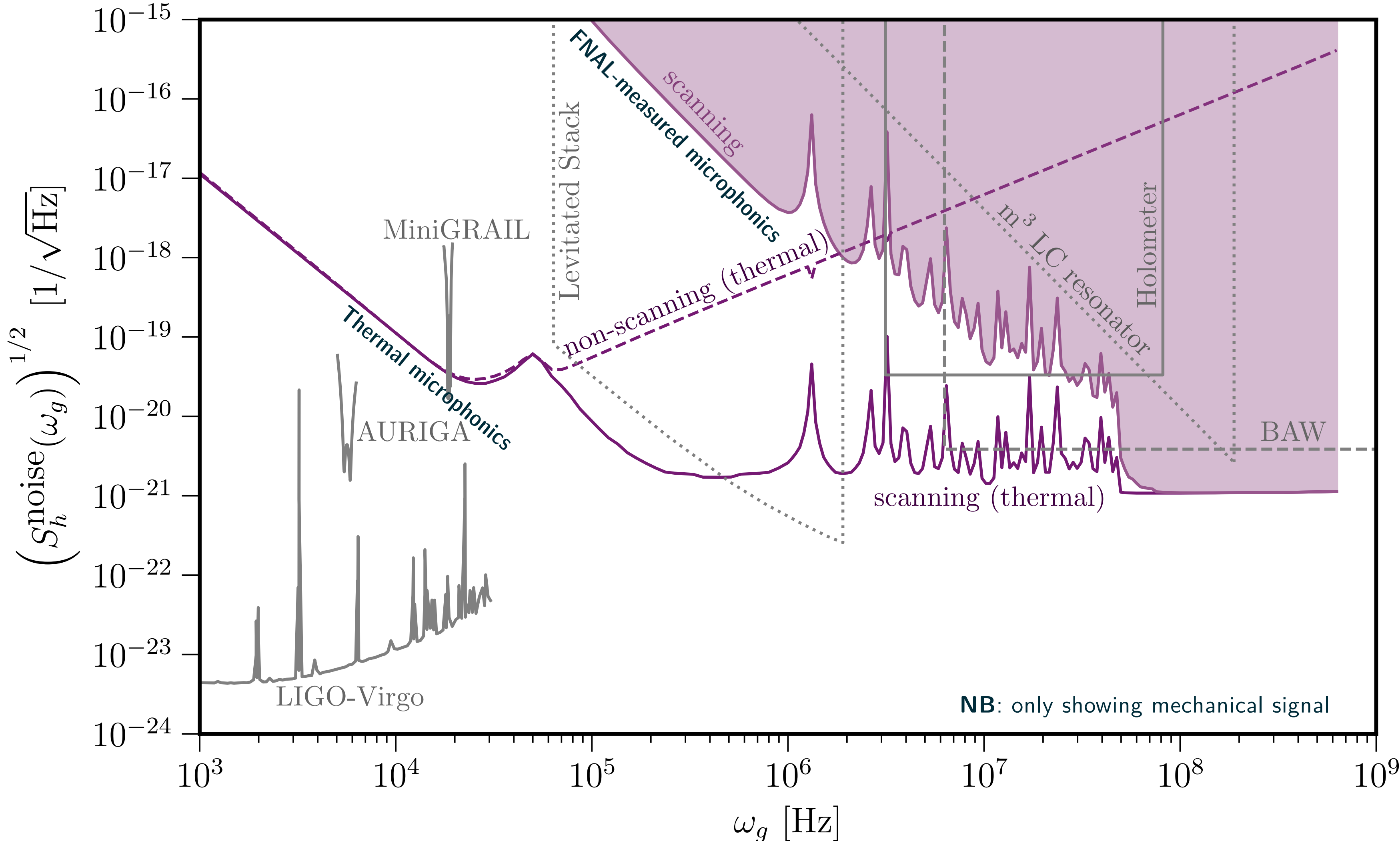
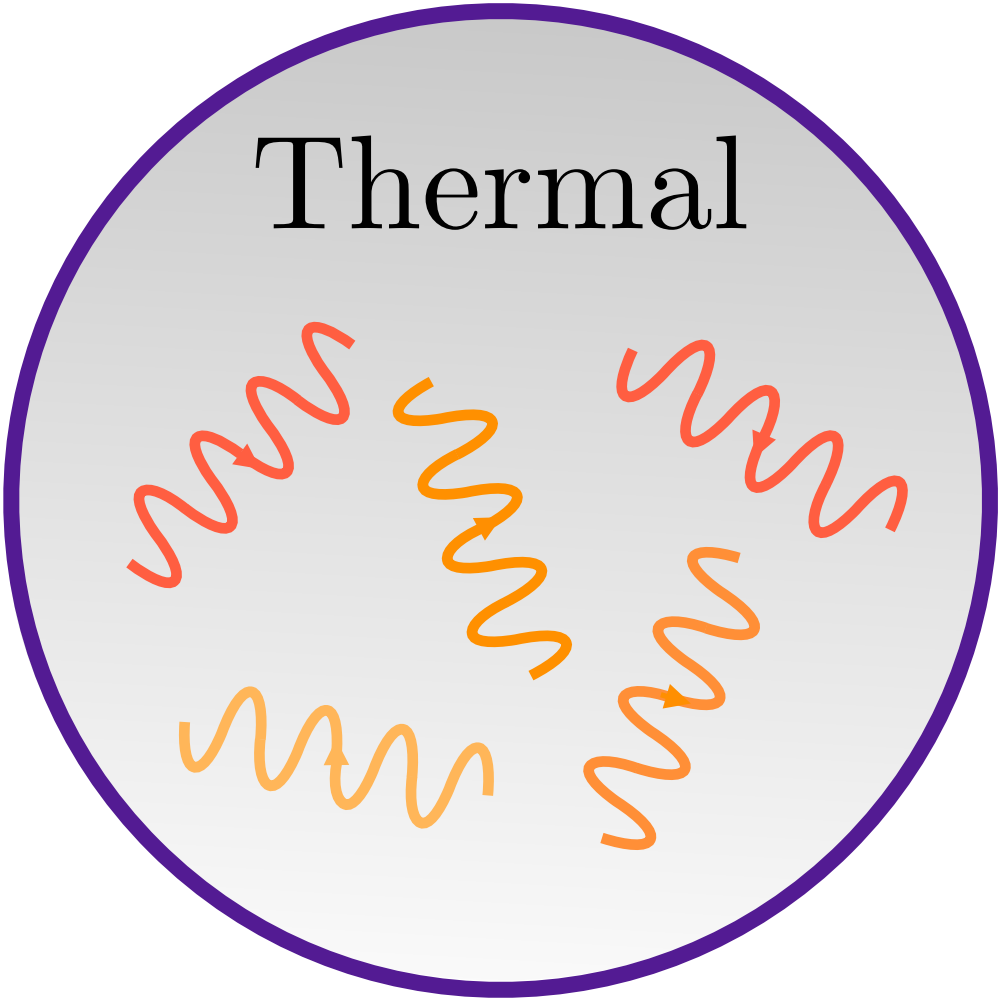
Mechanical modes less “rigid” than EM modes



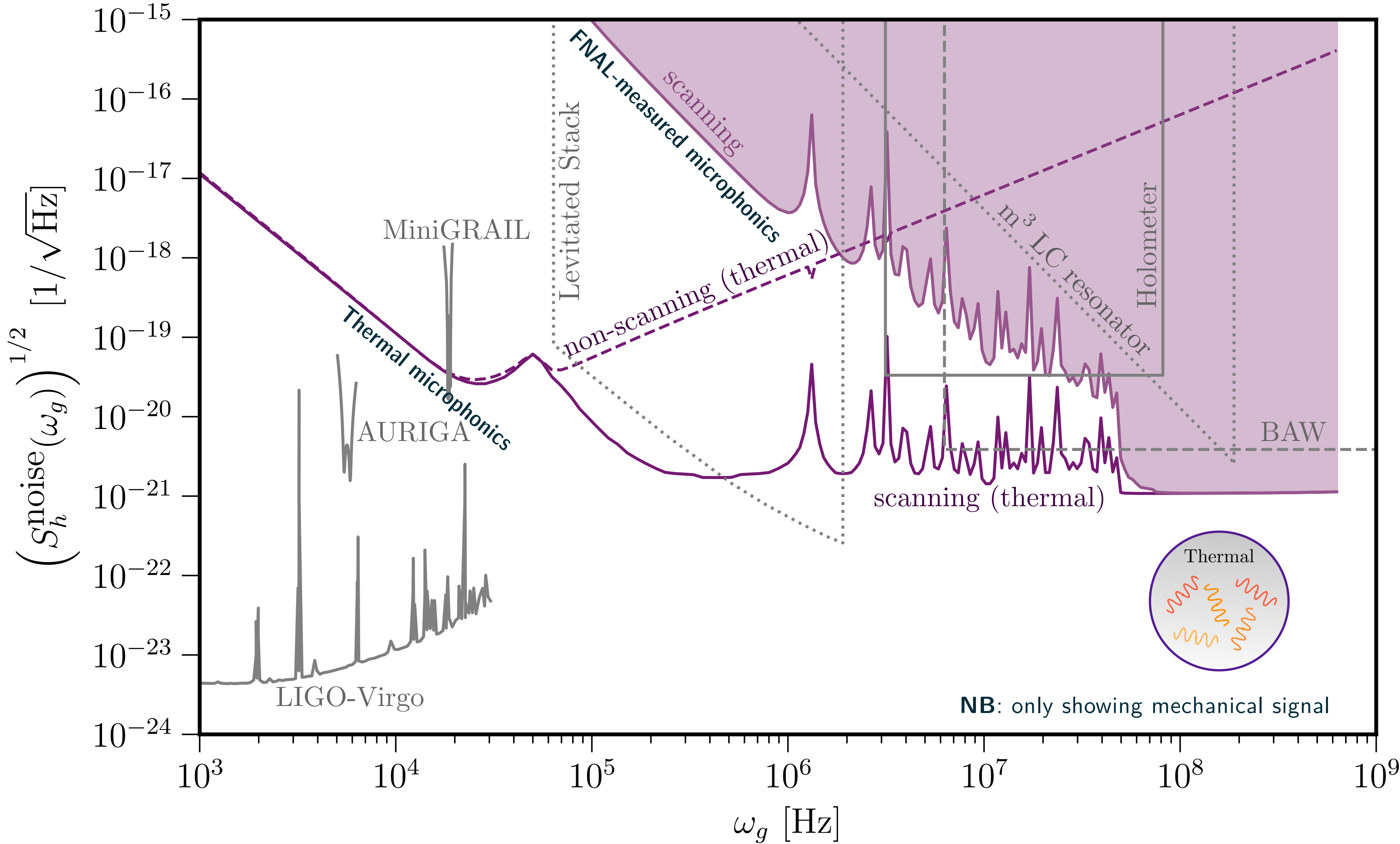
Noise in MAGO 2.0



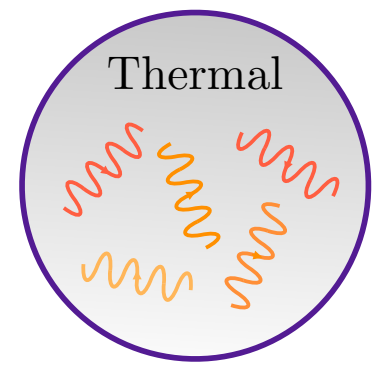
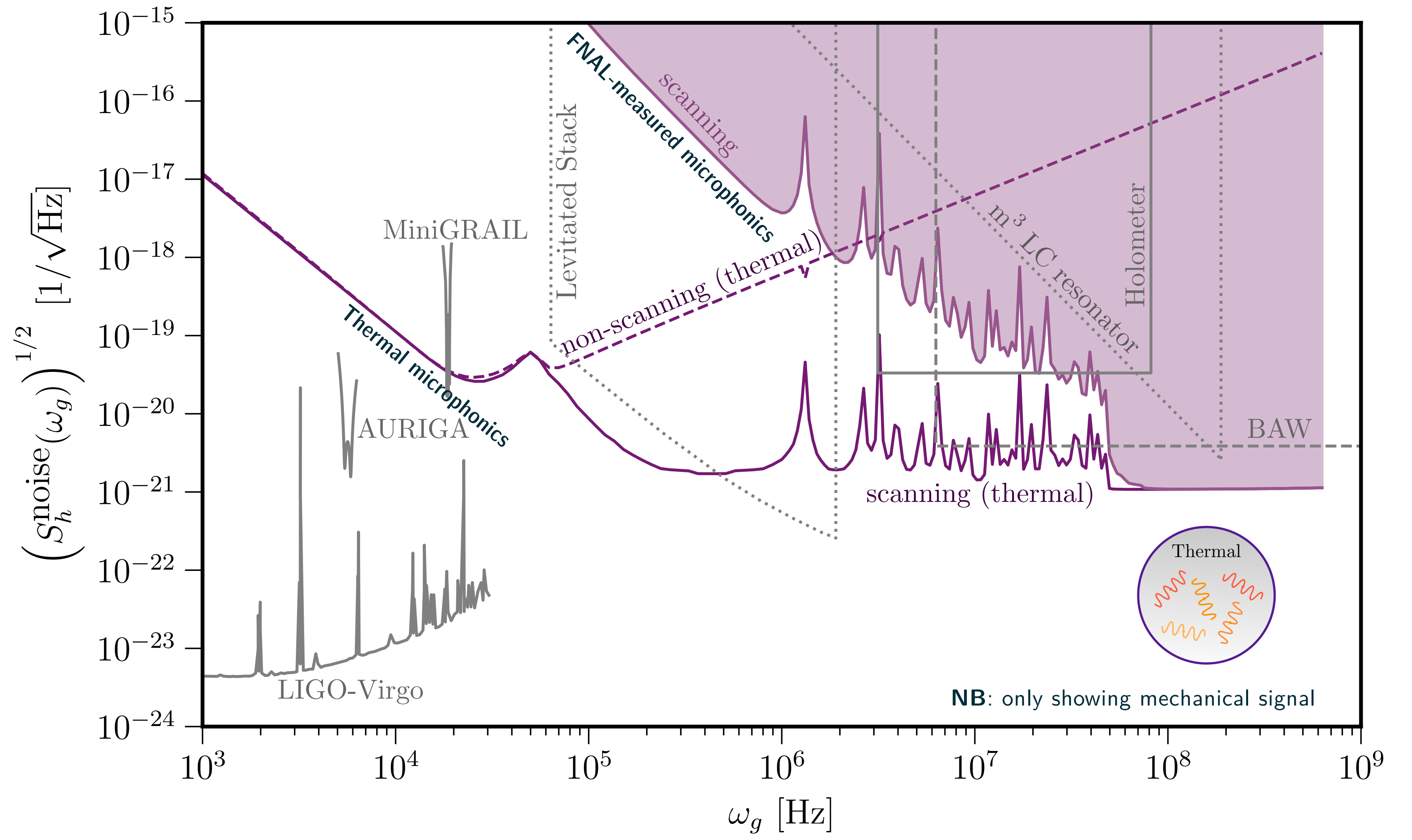
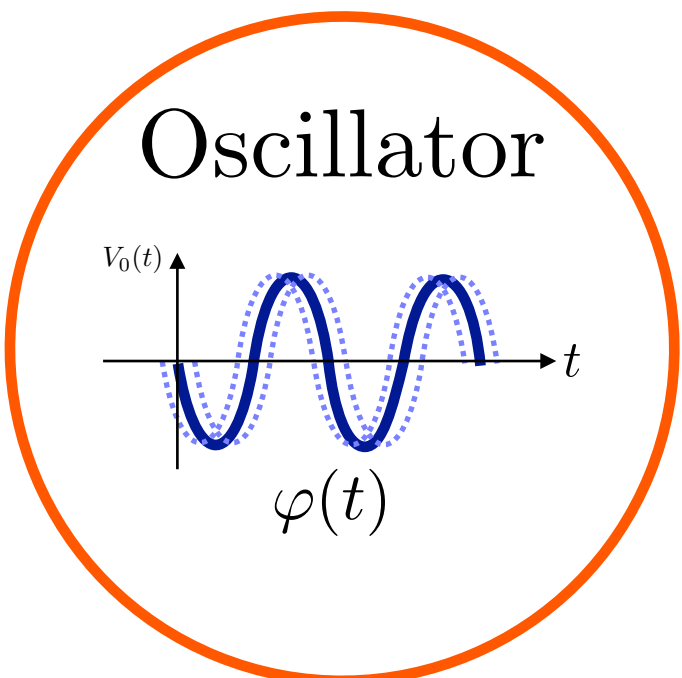
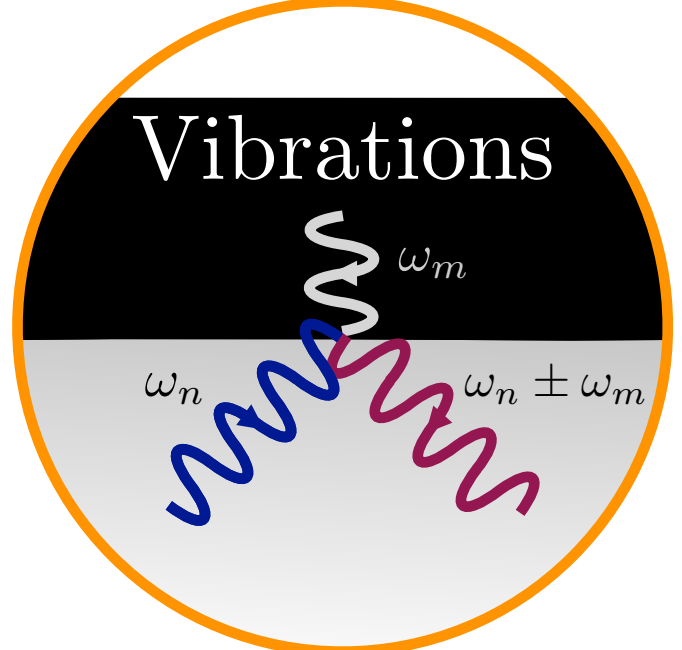
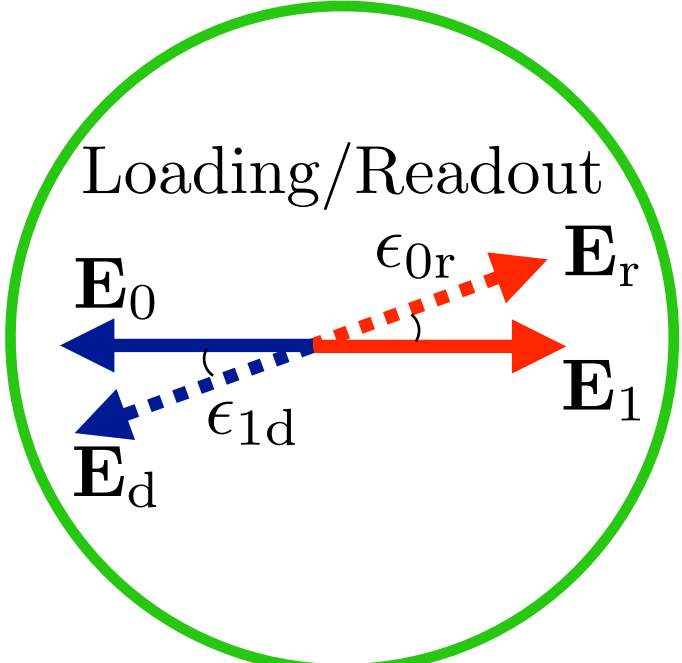
Noise in MAGO 2.0



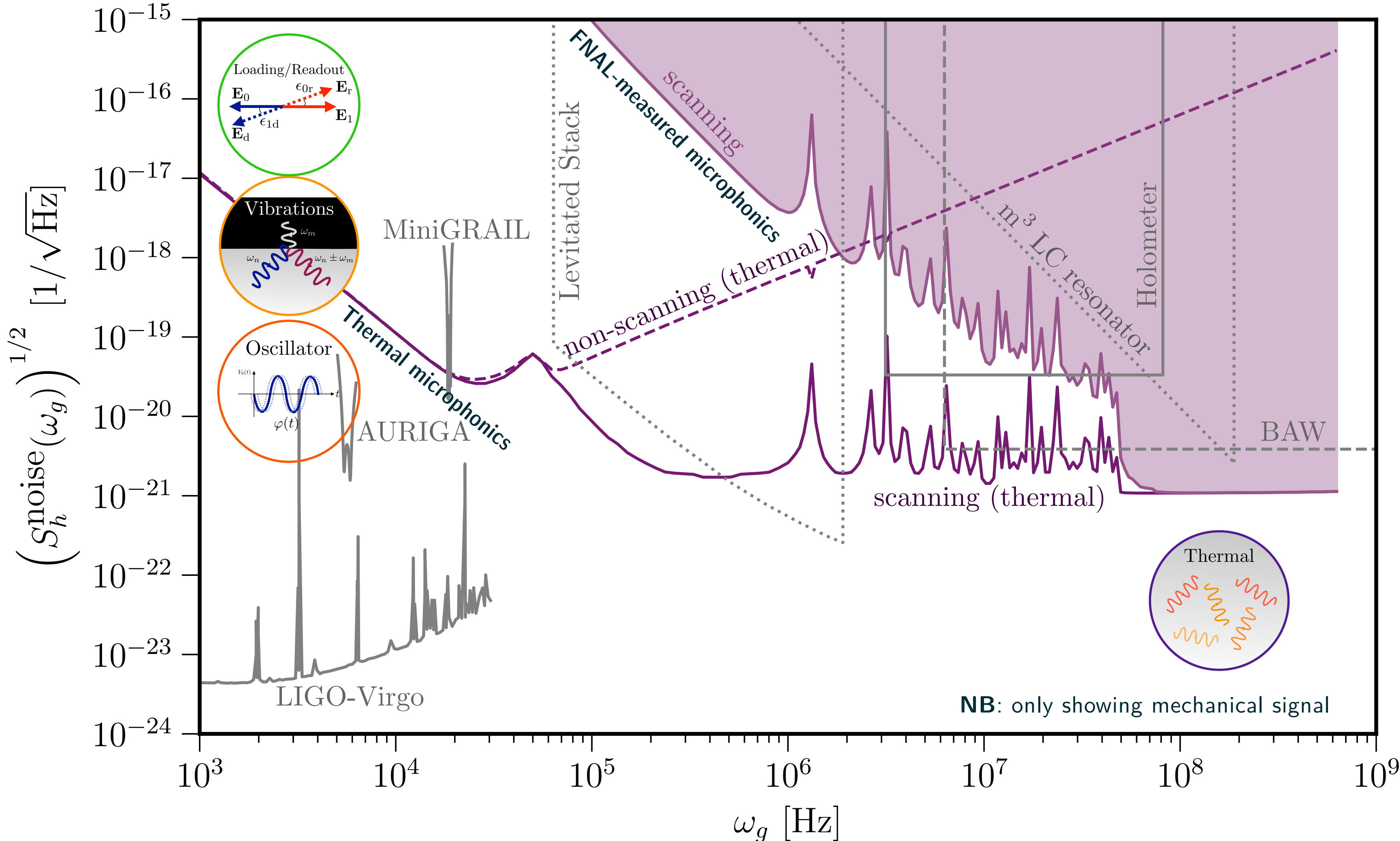
Noise in MAGO 2.0



Noise in MAGO 2.0

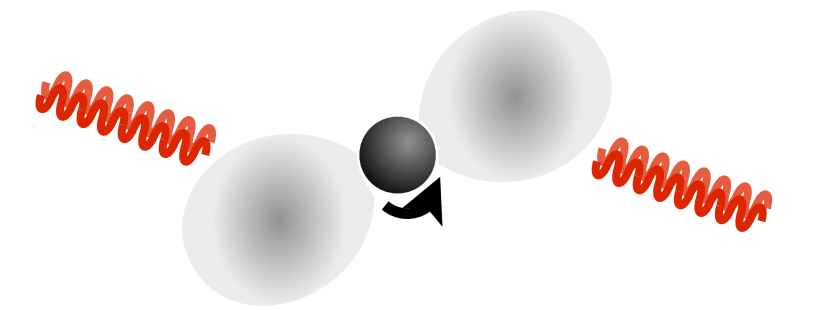
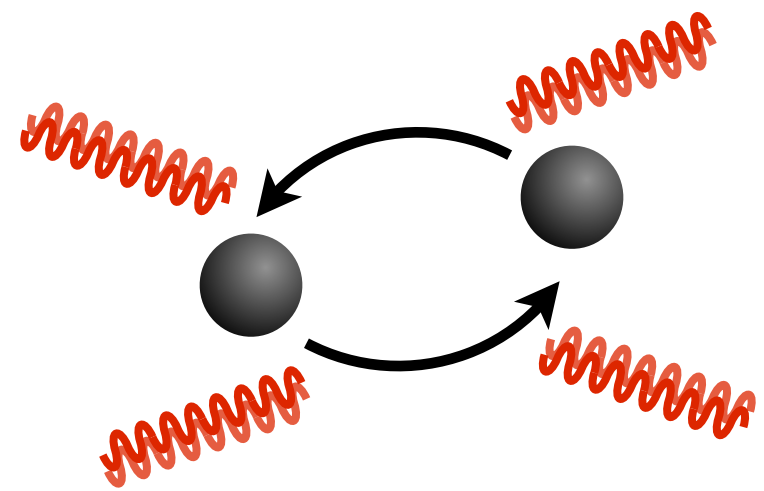


Noise in MAGO 2.0



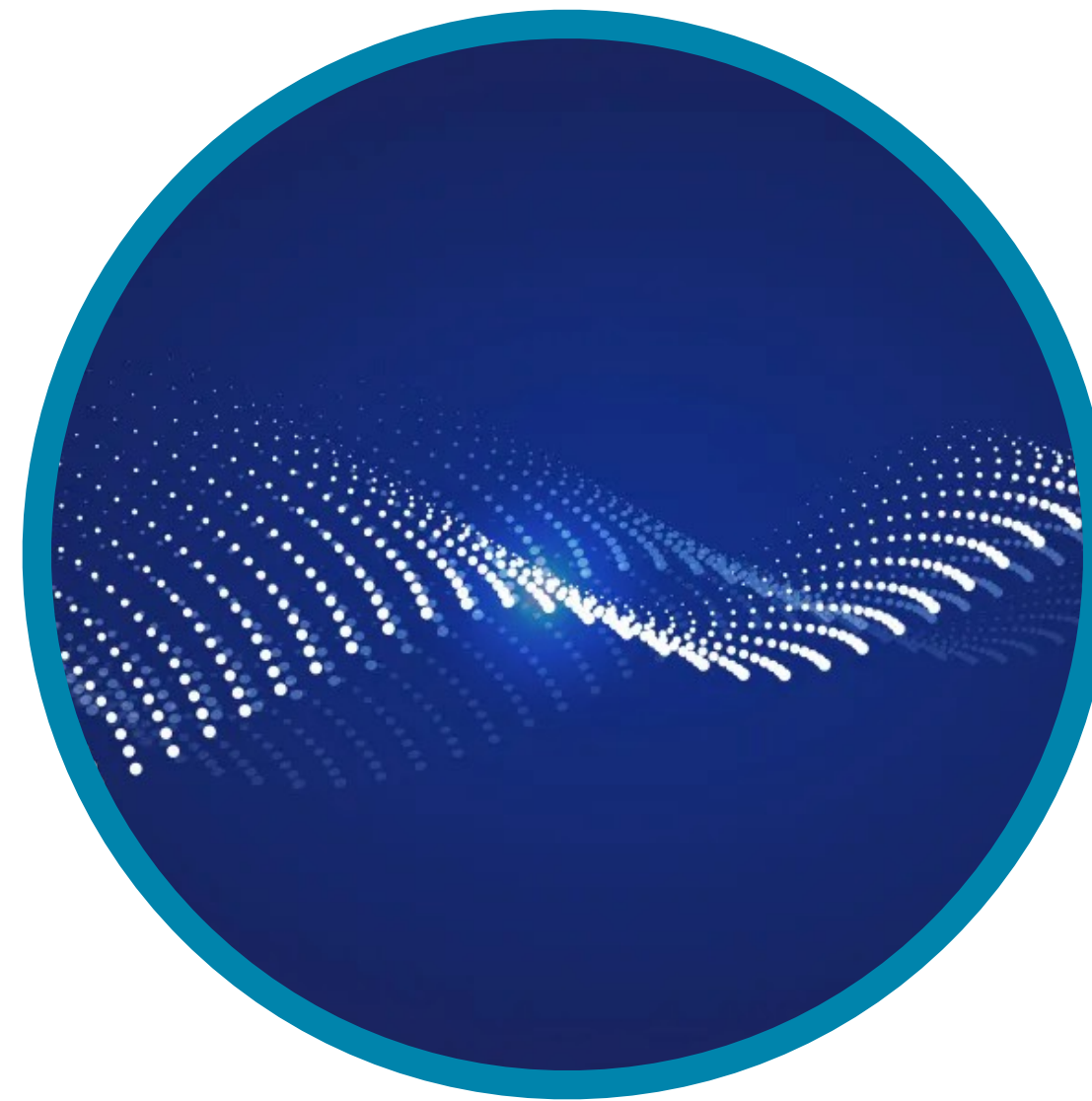
Gravitational Waves

- ▶ HFGWs provide insight into early universe/BSM
- ▶ Reaching anticipated stochastic sources difficult
- ▶ Broadband sensitivity of MAGO 2.0:
ongoing collaboration w/ DESY & SQMS @FNAL
See Bianca's talk
- ▶ Improve scan rate for resonant searches w/ Quantum techniques



BACKUP

Axion Dark Matter BACKUP



Comparison of approaches

	Static-field Haloscope	LC Resonator	RF Frequency Conversion
J_{eff}	$\propto B_0^{\text{static}} \cos(m_a t)$	$\propto B_0^{\text{static}} \cos(m_a t)$	$\propto B_0^{\text{RF}} \cos(\omega_0 \pm m_a)t$
\mathcal{E}_a	$\propto m_a/\omega_{\text{sig}} \sim 1$	$\propto m_a V^{1/3} \lesssim 1$	$\propto (\omega_0 \pm m_a)/\omega_{\text{sig}} \sim 1$
P_{sig}	$J_{\text{eff}}^2 V \min\left(\frac{Q_r}{m_a}, \frac{Q_a}{m_a}\right)$	$J_{\text{eff}}^2 m_a^2 V^{5/3} \min\left(\frac{Q_{\text{LC}}}{m_a}, \frac{Q_a}{m_a}\right)$	$J_{\text{eff}}^2 V \min\left(\frac{Q_{\text{SRF}}}{\omega_0 \pm m_a}, \frac{Q_a}{m_a}\right)$

Axion Signal

Signal Power Spectral Density (PSD):

$$S_{\text{sig}}(\omega) = \frac{\omega_1}{Q_1} (g_{a\gamma\gamma} \eta_{10} B_0)^2 V \frac{\omega^2}{(\omega^2 - \omega_1^2)^2 + (\omega \omega_1 / Q_1)^2} \int \frac{d\omega'}{(2\pi)^2} (\omega' - \omega)^2 S_{b_0}(\omega') S_a(\omega - \omega')$$

$$\text{Axion PSD: } \langle a(t)^2 \rangle = \frac{1}{(2\pi)^2} \int d\omega S_a(\omega) = \frac{\rho_{\text{DM}}}{m_a^2}$$

Background magnetic field PSD:

$$S_{b_i}(\omega) = \pi^2 \left(\delta(\omega - \omega_i) + \delta(\omega + \omega_i) \right) + S_{b_i}^{(\text{phase})} + S_{b_i}^{(\text{mech})}$$

$$\text{NB: } B_i \equiv \sqrt{\frac{1}{V_{\text{cav}}} \int_{V_{\text{cav}}} |\mathbf{B}_i(x)|^2} \quad \mathbf{B}_i(x, t) = \mathbf{B}_i(x) b_i(t)$$

Axion Signal

Signal Power Spectral Density (PSD):

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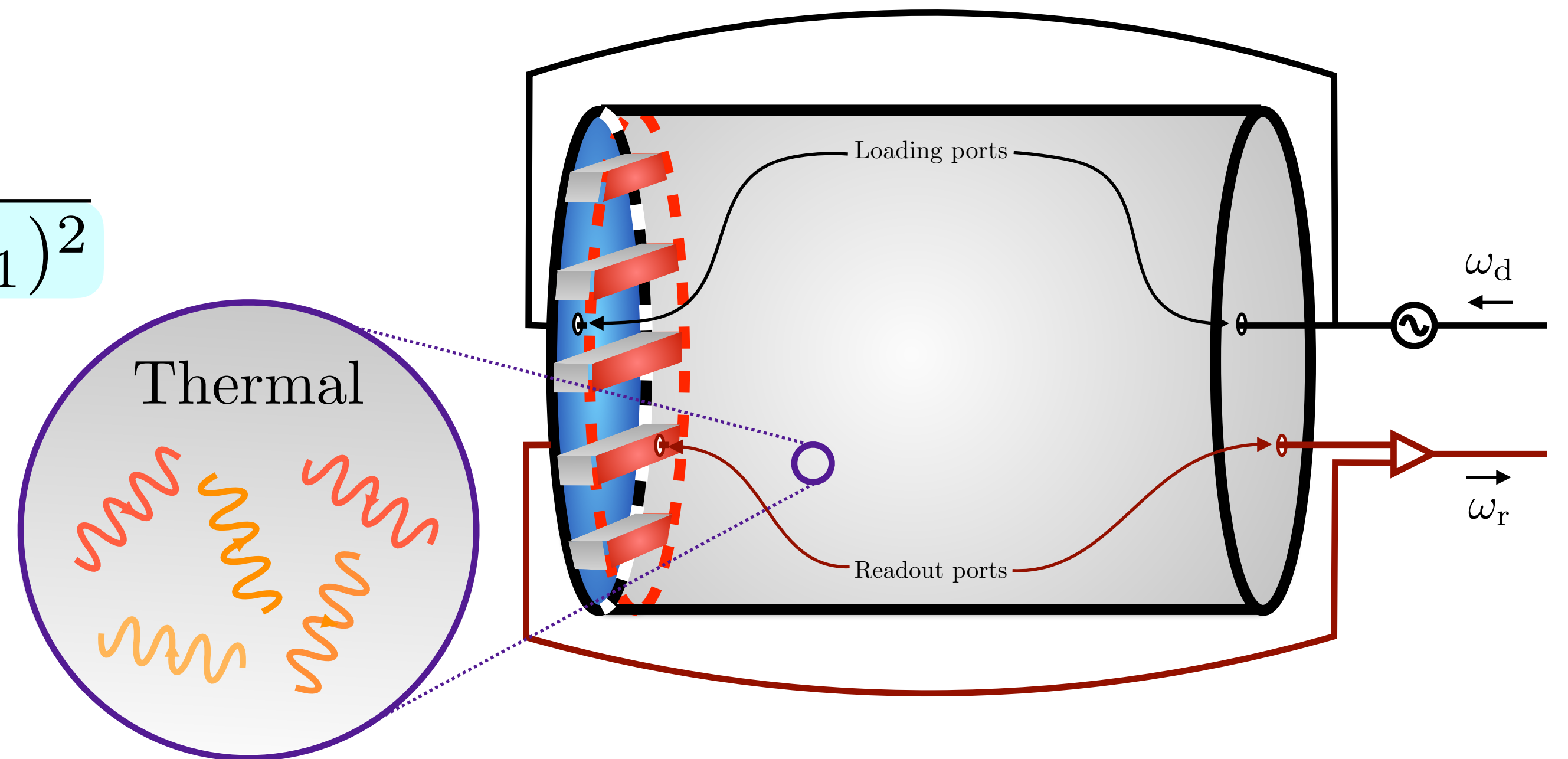
Signal Power (resonant):

$$P_{\text{sig}} \simeq \frac{1}{4} (g_{a\gamma\gamma} \eta_{10} B_0)^2 \rho_{\text{DM}} V \times \begin{cases} Q_1 / \omega_1 & \frac{m_a}{Q_a} \ll \frac{\omega_1}{Q_1} \\ \pi Q_a / m_a & \frac{m_a}{Q_a} \gg \frac{\omega_1}{Q_1} \end{cases},$$

Standard Noise Sources: Thermal Noise

Power Spectral Density:

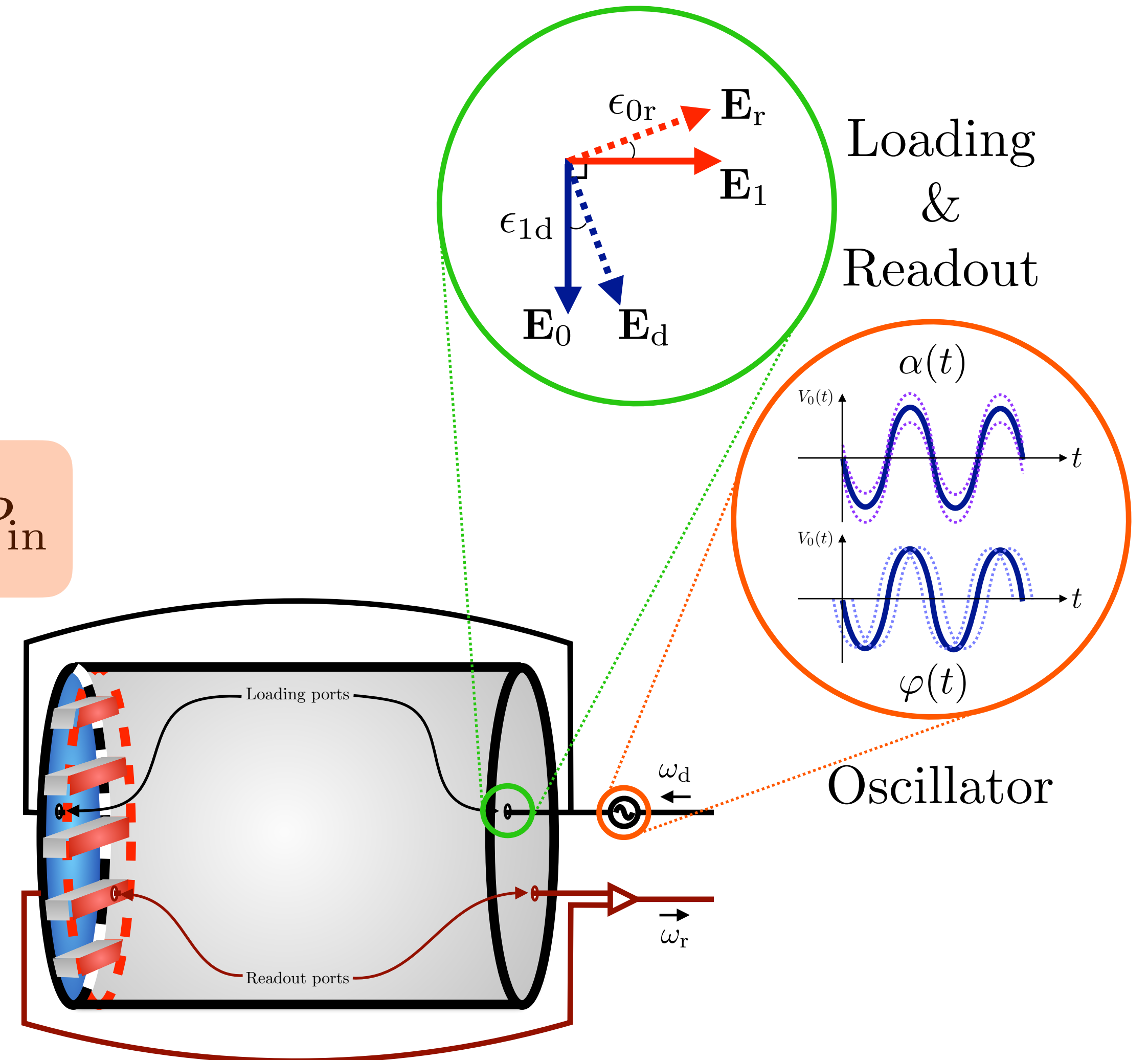
$$S_{\text{th}}(\omega) = \frac{Q_1}{Q_{\text{int}}} \frac{4\pi T (\omega \omega_1 / Q_1)^2}{(\omega^2 - \omega_1^2)^2 + (\omega \omega_1 / Q_1)^2}$$



Non-standard Noise Sources: Phase Noise

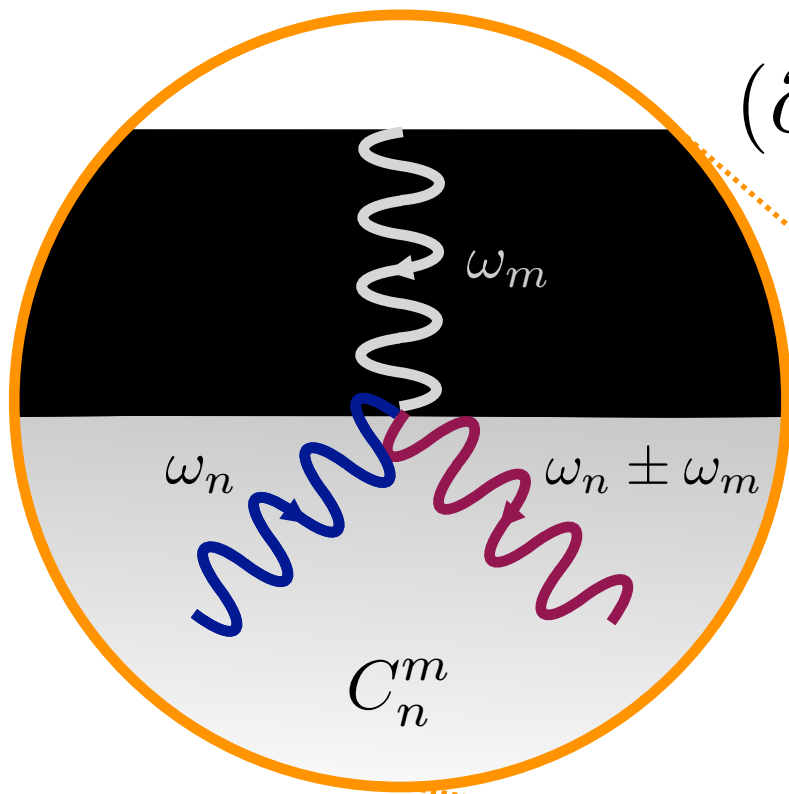
Power Spectral Density:

$$S_{\text{phase}}(\omega) \simeq \frac{1}{2} \epsilon_{1d}^2 S_{\varphi}(\omega - \omega_0) \times \frac{(\omega \omega_1 / Q_1)^2}{(\omega^2 - \omega_1^2)^2 + (\omega \omega_1 / Q_1)^2} \frac{\omega_0 Q_1}{\omega_1 Q_0} P_{\text{in}}$$



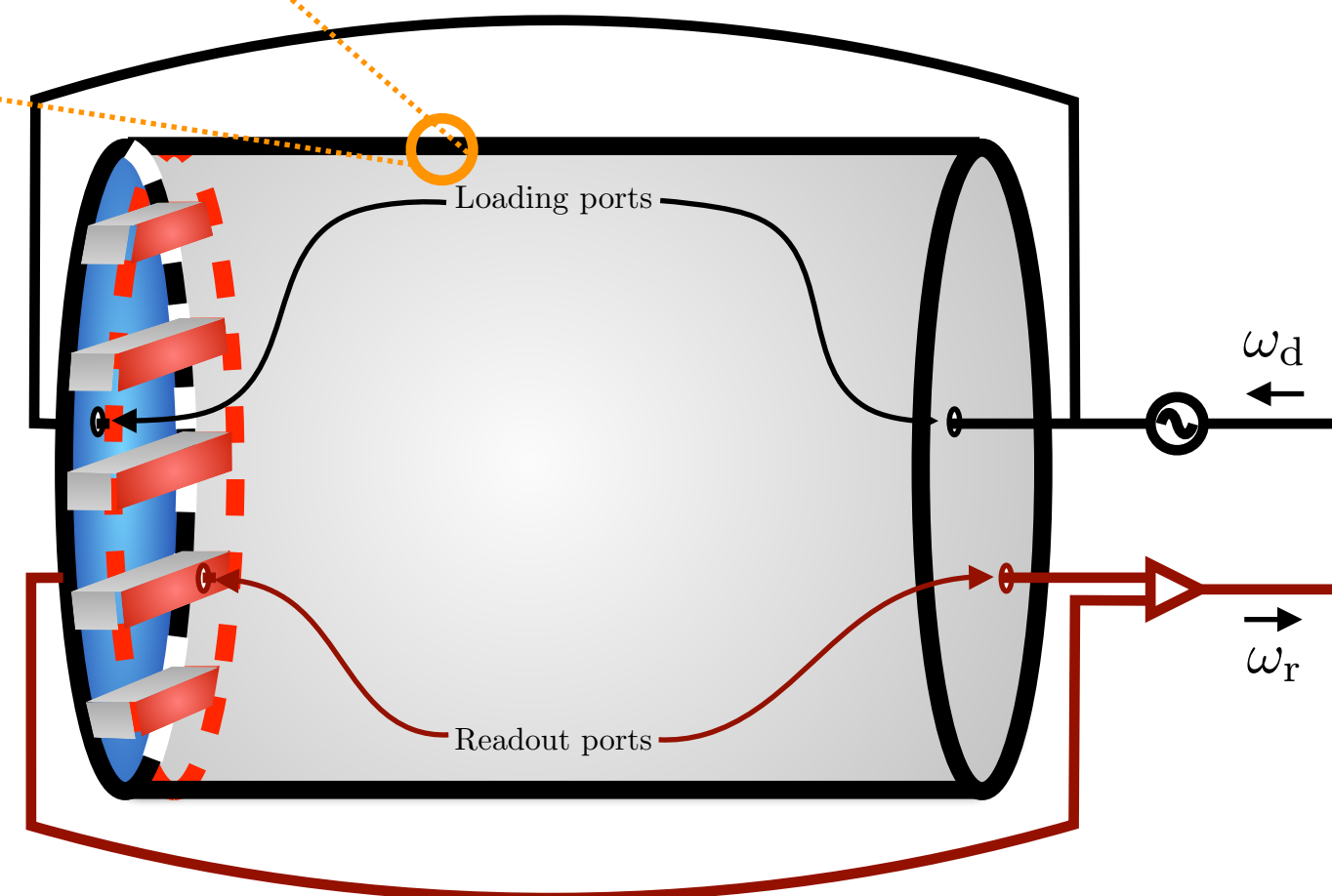
Non-standard Noise Sources: Vibrations

Vibrations



$$(\partial_t^2 + \omega_n^2) e_n \simeq \frac{2}{U_n} \sum_m e_m \begin{cases} \int_{\Delta V} d^3\mathbf{x} \left[\omega_n \omega_m \mathbf{B}_m \cdot \mathbf{B}_n^* - \frac{1}{2} (\omega_n^2 + \omega_m^2) \mathbf{E}_m \cdot \mathbf{E}_n^* \right] + \mathcal{O}(\Delta V^2) & (V' \subset V) \\ \int_{\Delta V} d^3\mathbf{x} \left[\omega_n \omega_m \mathbf{B}_m \cdot \mathbf{B}_n^* - \omega_n^2 \mathbf{E}_m \cdot \mathbf{E}_n^* \right] + \mathcal{O}(\Delta V^2) & (V \subset V') \end{cases}$$

Cavity perturbation theory, see e.g. Meidlinger (2009), Pozar, etc.

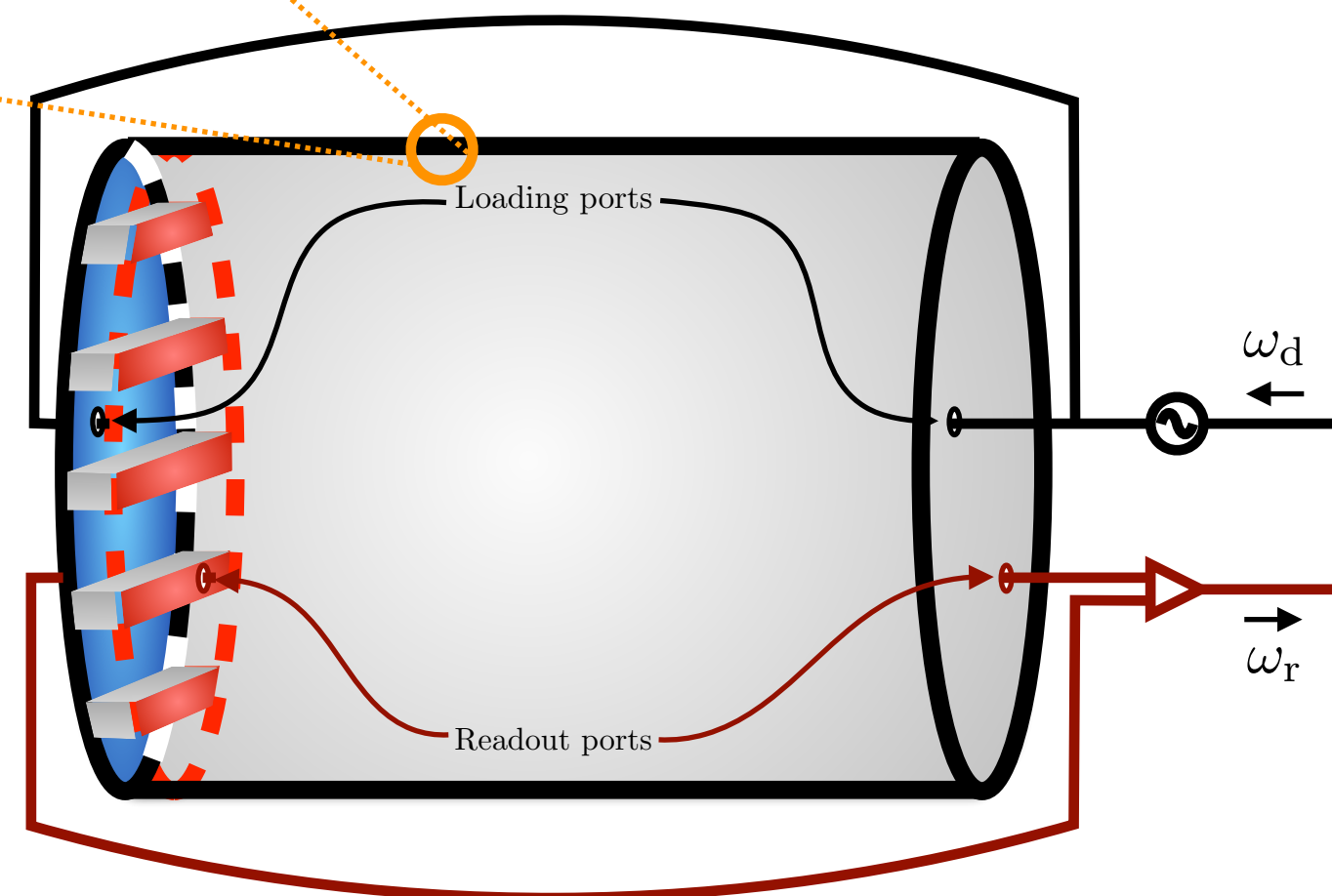
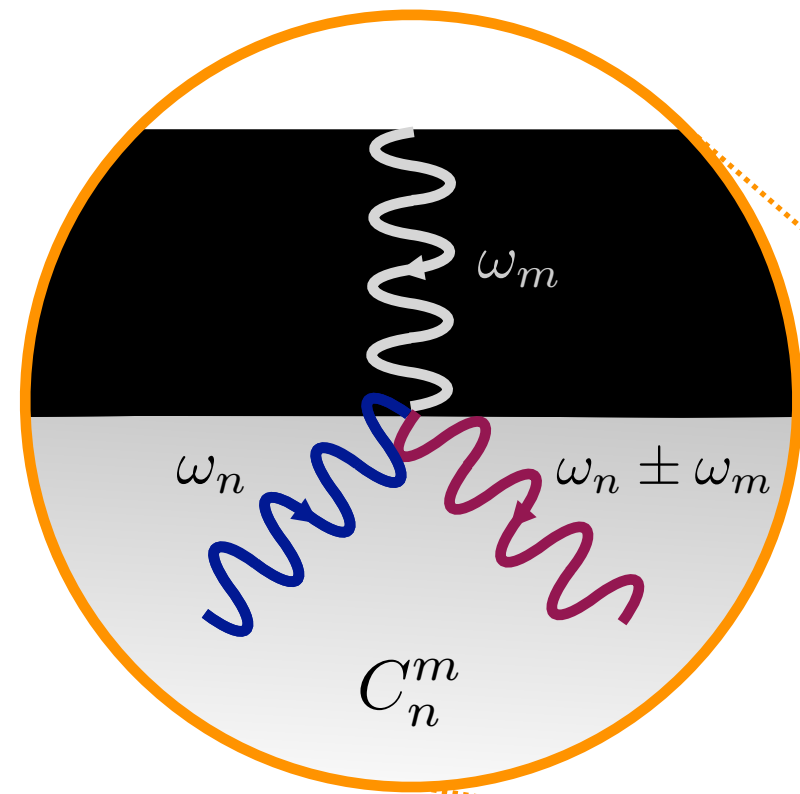


Non-standard Noise Sources: Vibrations

Vibrations

Power Spectral Density:

$$S_{\text{mix}} = \frac{\epsilon_{01}^2}{4} P_{\text{in}} \sum_{i=0,1} \frac{\omega_i^4}{(\omega_0^2 - \omega_i^2)^2 + (\omega_0 \omega_i / Q_i)^2} \frac{(\omega \omega_i)^2}{(\omega^2 - \omega_i^2)^2 + (\omega \omega_i / Q_i)^2} |C_i^m|^2 S_{q_m}(\omega - \omega_0)$$



Displacement PSD:

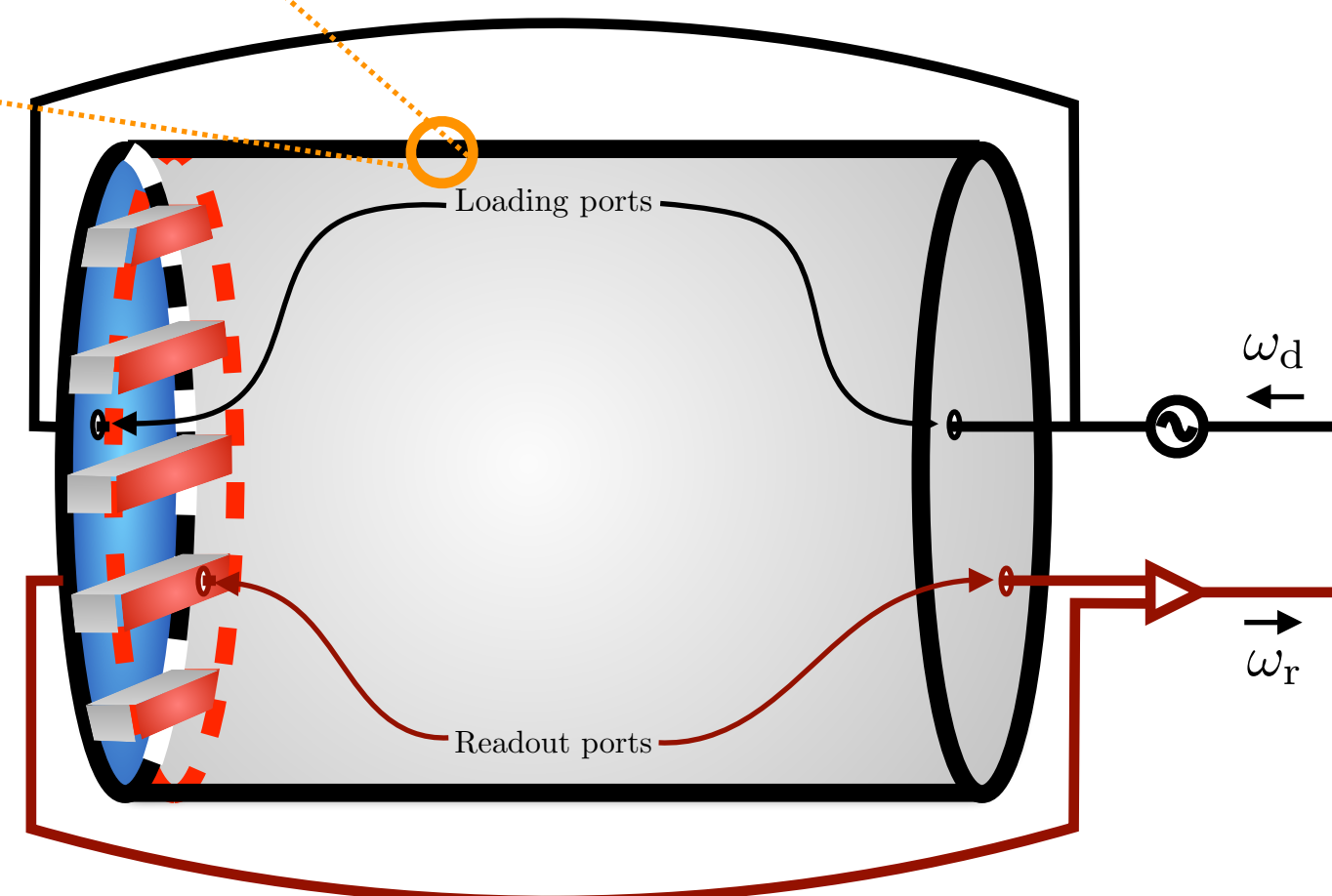
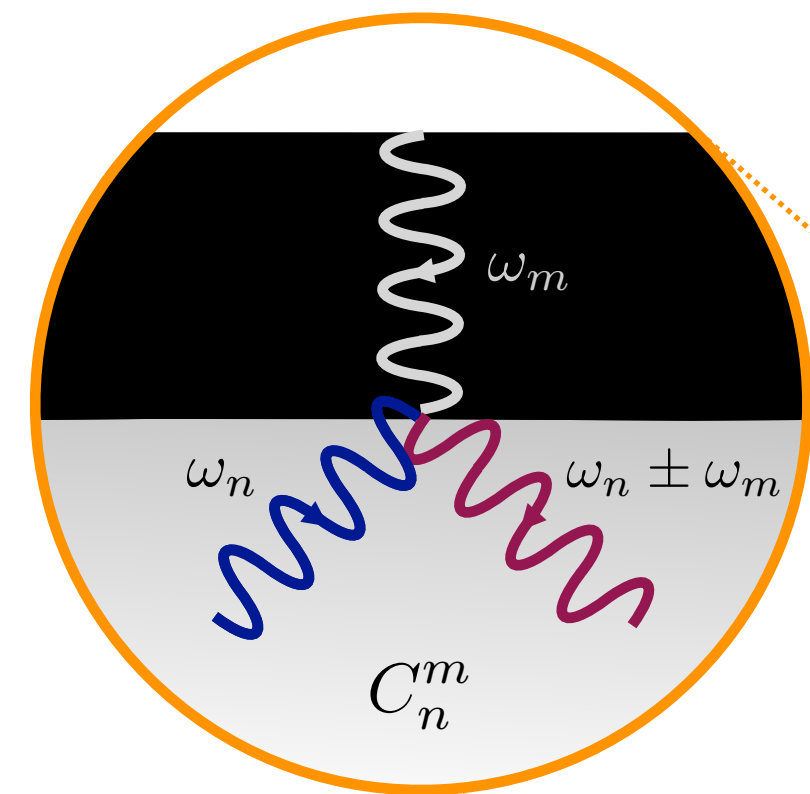
$$S_{q_m}(\omega) \simeq \frac{1}{M^2} \frac{S_{f_m}(\omega)}{(\omega^2 - \omega_m^2)^2 + (\omega_m \omega / Q_m)^2}$$

Non-standard Noise Sources: Vibrations mixing

Vibrations

Power Spectral Density:

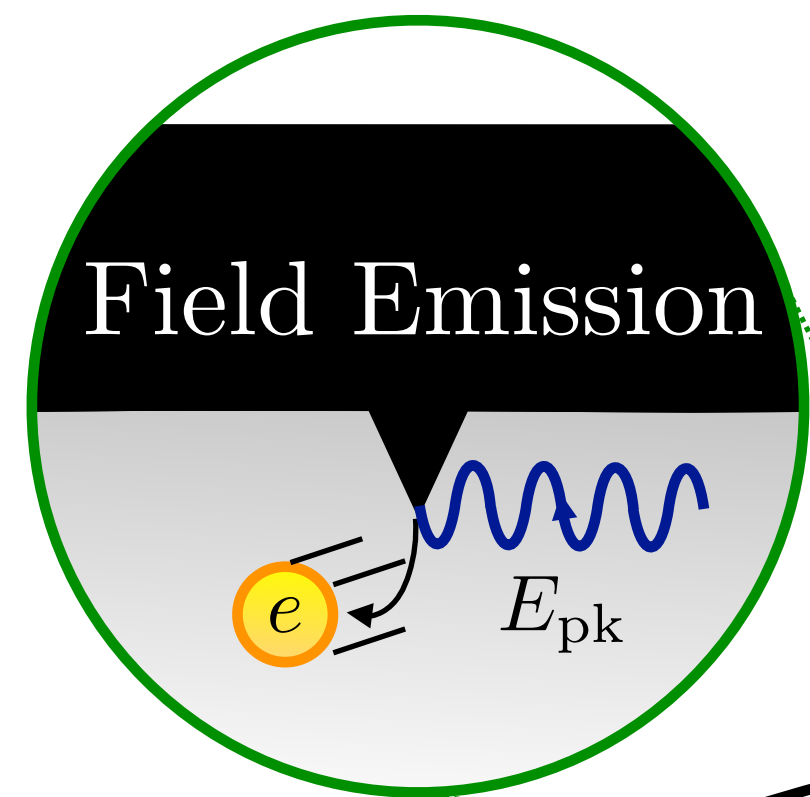
$$S_{\text{mix}} = \frac{\eta_m^2}{4} P_{\text{in}} \frac{\omega_0^4}{(\omega^2 - \omega_1^2)^2 + (\omega\omega_1/Q_1)^2} |C_1^m|^2 S_{q_m}(\omega - \omega_0)$$



Displacement PSD:

$$S_{q_m}(\omega) \simeq \frac{1}{M^2} \frac{S_{f_m}(\omega)}{(\omega^2 - \omega_m^2)^2 + (\omega_m\omega/Q_m)^2}$$

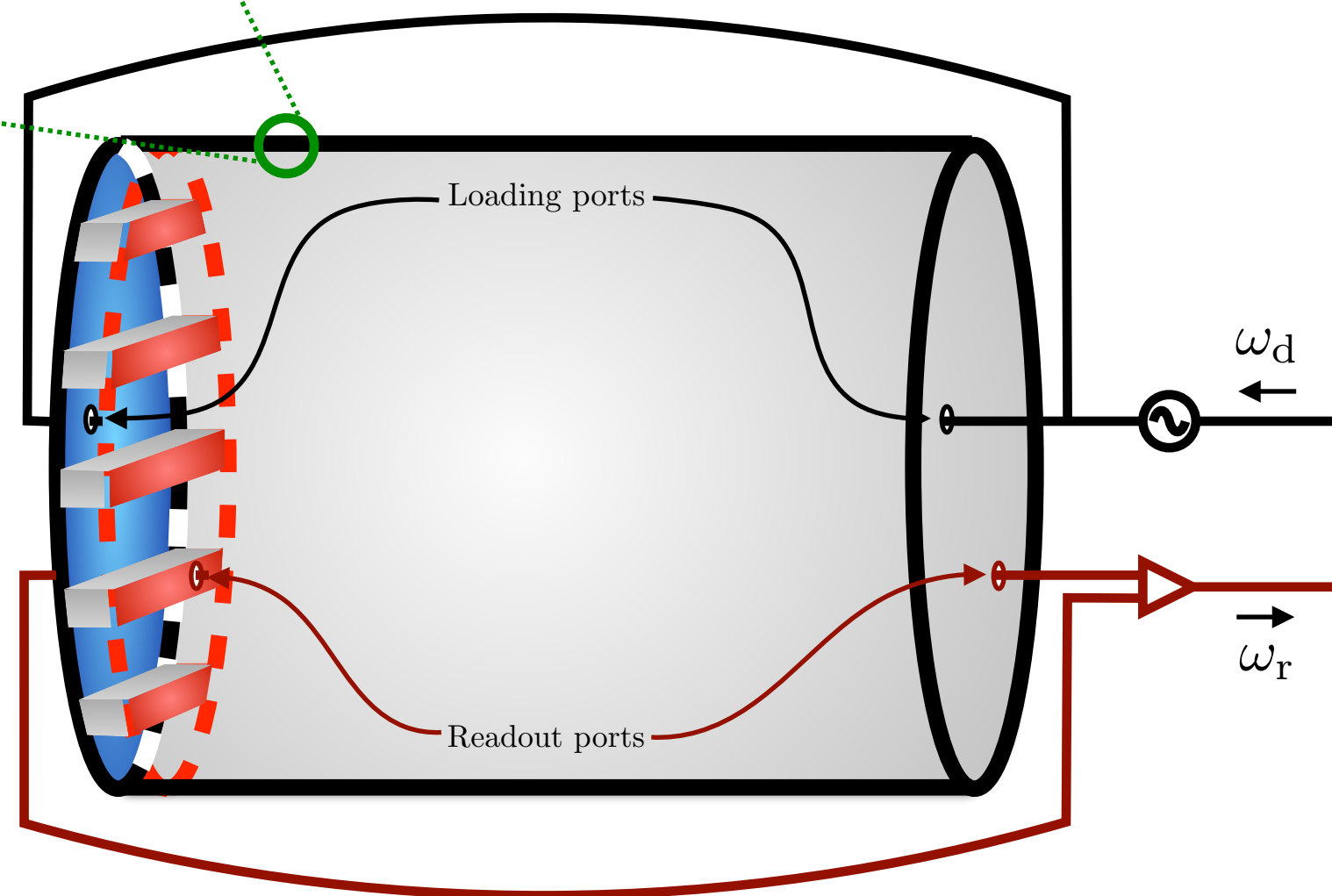
Non-standard Noise Sources: Field Emission



Power Spectral Density:

$$\frac{S(\omega_1)}{4\pi T} \sim \frac{P_{\text{tot}}}{0.1 \text{ W}} \times \begin{cases} 1 & \text{synchrotron} \\ 10^{-6} & \text{transition} \\ 10^{-5} & \text{Bremsstrahlung,} \end{cases}$$

Limits max B-field ~ 0.2 T



All Noise Sources

$B = 0.1 \text{ T}, \quad T = 2\text{K}, \quad \omega_0 = 2\pi \text{ GHz}, \quad V = 0.05 \text{ m}^3$

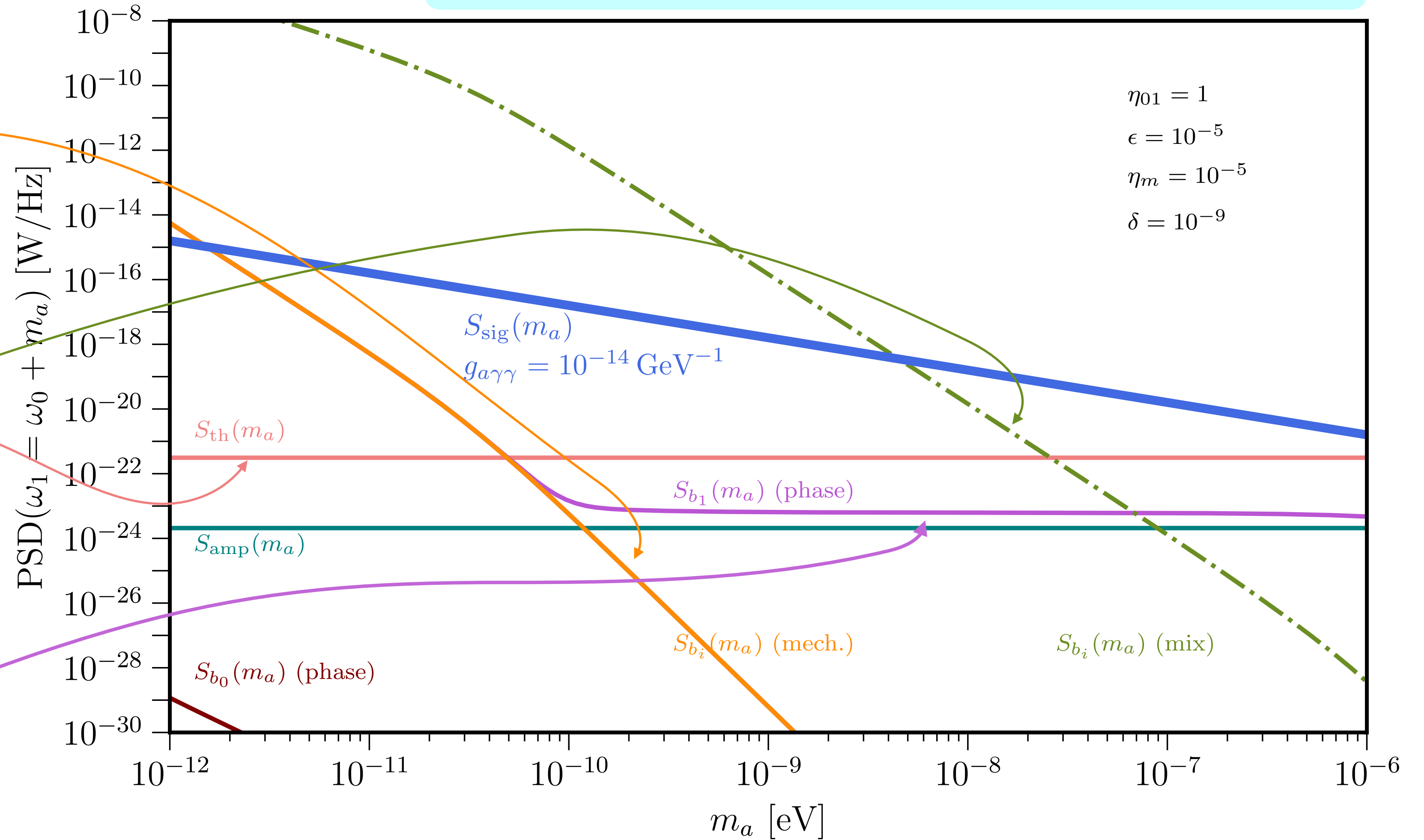
$$P_{\text{mech}} \sim \epsilon_{1d}^2 \delta^2 \frac{\omega_0^2 \omega_{\text{min}}^3}{m_a^5} P_{\text{in}}$$

fractional wall disp. δ

$$P_{\text{mix}} \sim \eta_m^2 \delta^2 \frac{Q_1 Q_m \omega_1 \omega_{\text{min}}^3}{m_a^4} P_{\text{in}}$$

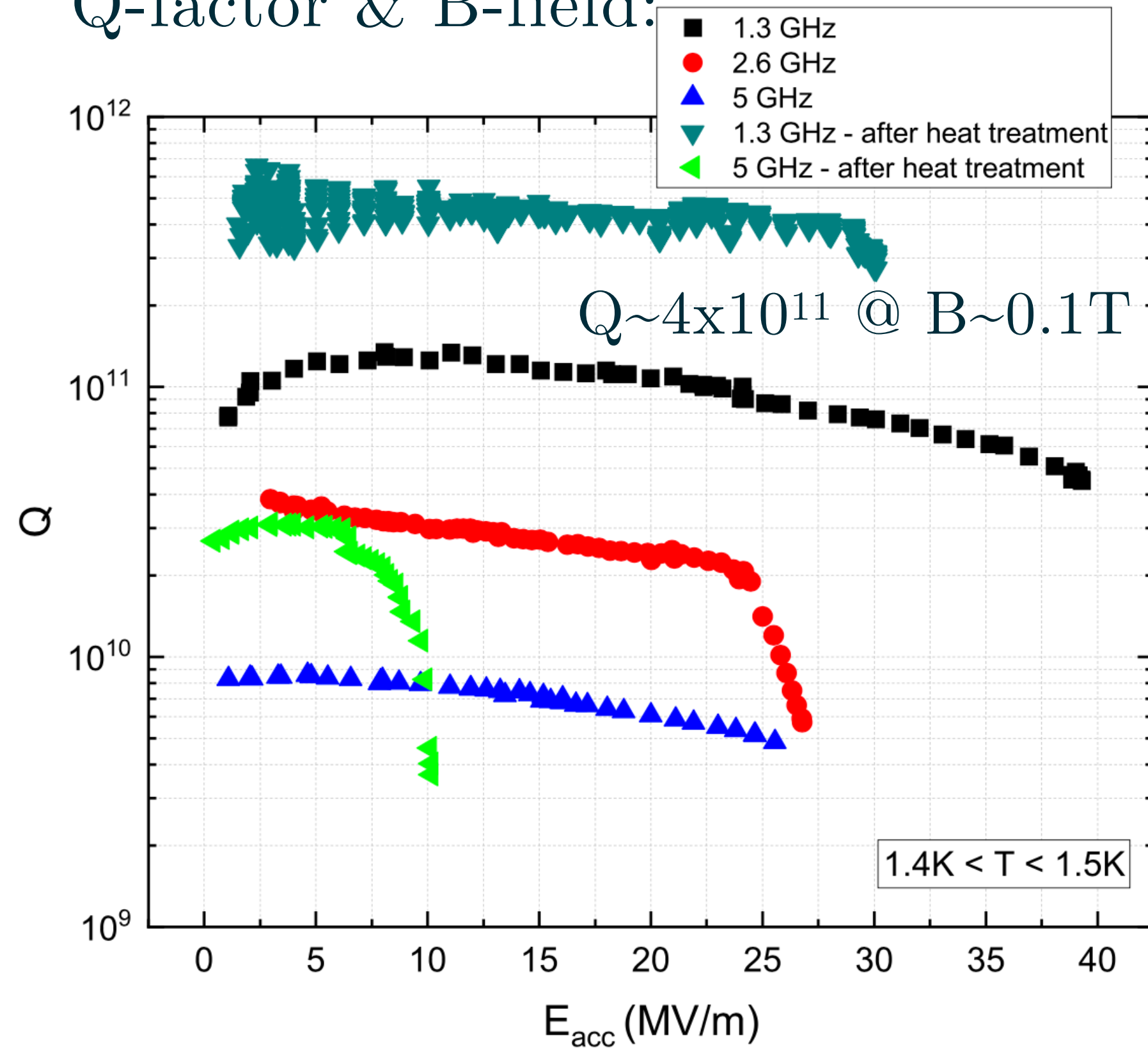
$$P_{\text{th}} \sim T \frac{\omega_1}{Q_{\text{int}}}$$

$$P_{\text{phase}} \sim \epsilon_{1d}^2 S_\varphi(m_a) \frac{\omega_1}{Q_{\text{int}}} P_{\text{in}}$$



Technology Requirements

Q-factor & B-field:

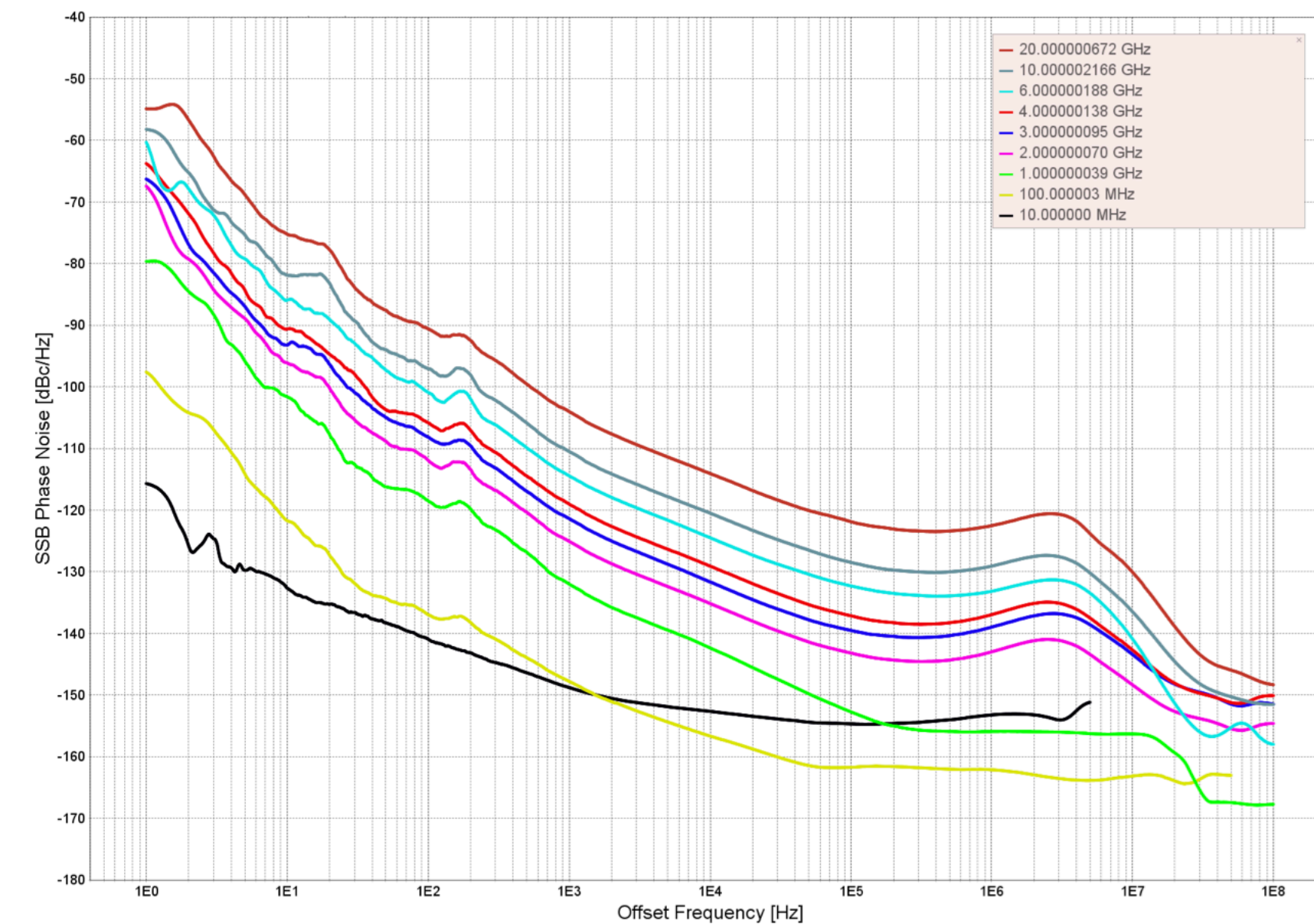


arXiv: 1810.03703 Romanenko et al.

Phase noise: *e.g.* BN865-M

TYPICAL PERFORMANCE CURVES

Phase Noise Performance with option LN



In total need ~ -240 dB/Hz @ GHz to reach thermal noise floor

Technology Requirements

Mode rejection:

$\epsilon = 10^{-7}$ achieved



gr-qc/0502054 Ballantini, ..., Calatroni et al
physics/0004031 Bernard, Gemme, Parodi, Picasso

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Original MAGO collaboration

gr-qc/0502054 Ballantini, ..., Calatroni et al
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Original MAGO collaboration

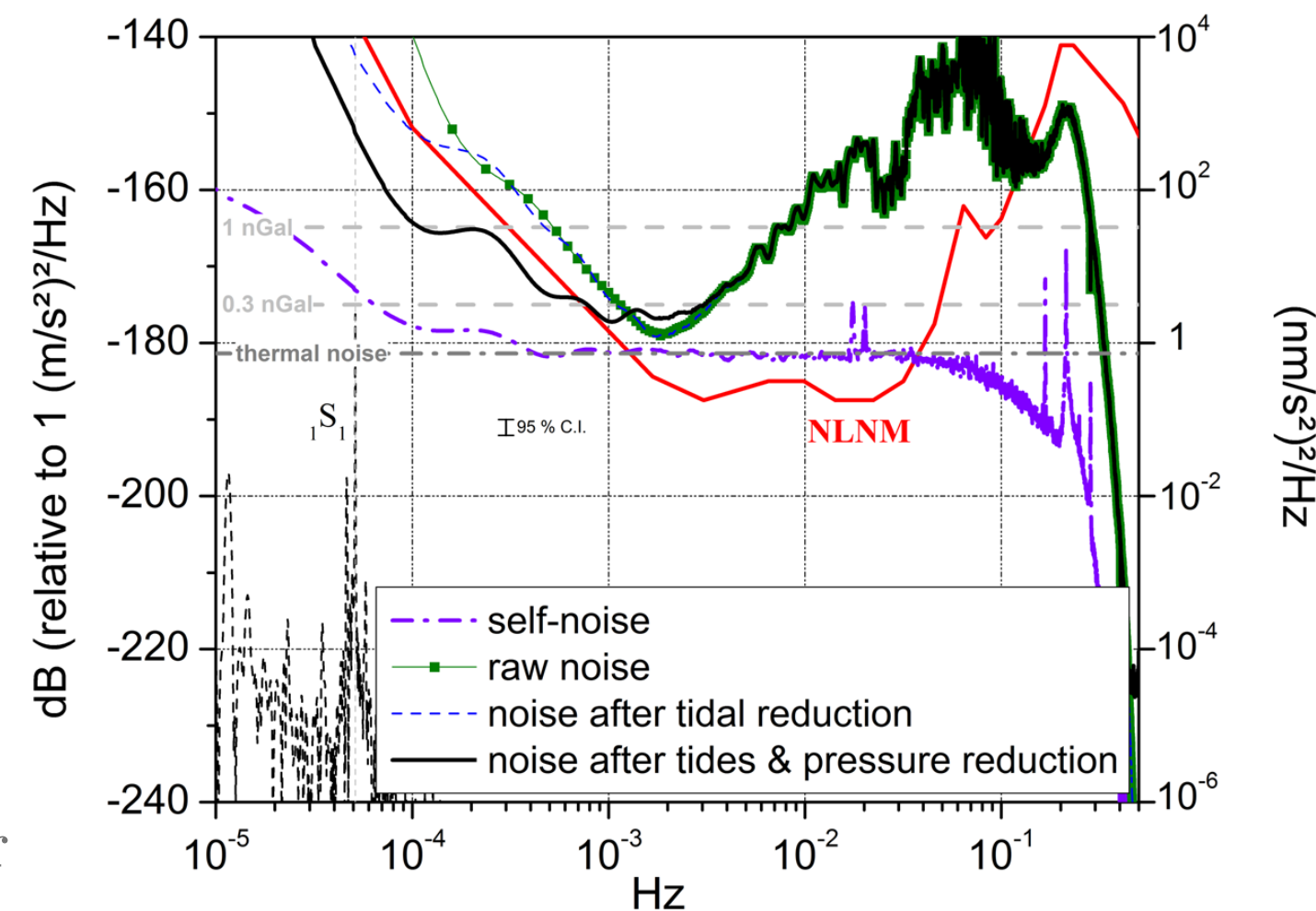
gr-qc/0502054 Ballantini, ..., Calatroni et al
 physics/0004031 Bernard, Gemme, Parodi, Picasso

Low-frequency
 seismic noise:

$$\Delta\omega/\omega \sim \delta \sim 10^{-10}$$

DarkSRF (2020)

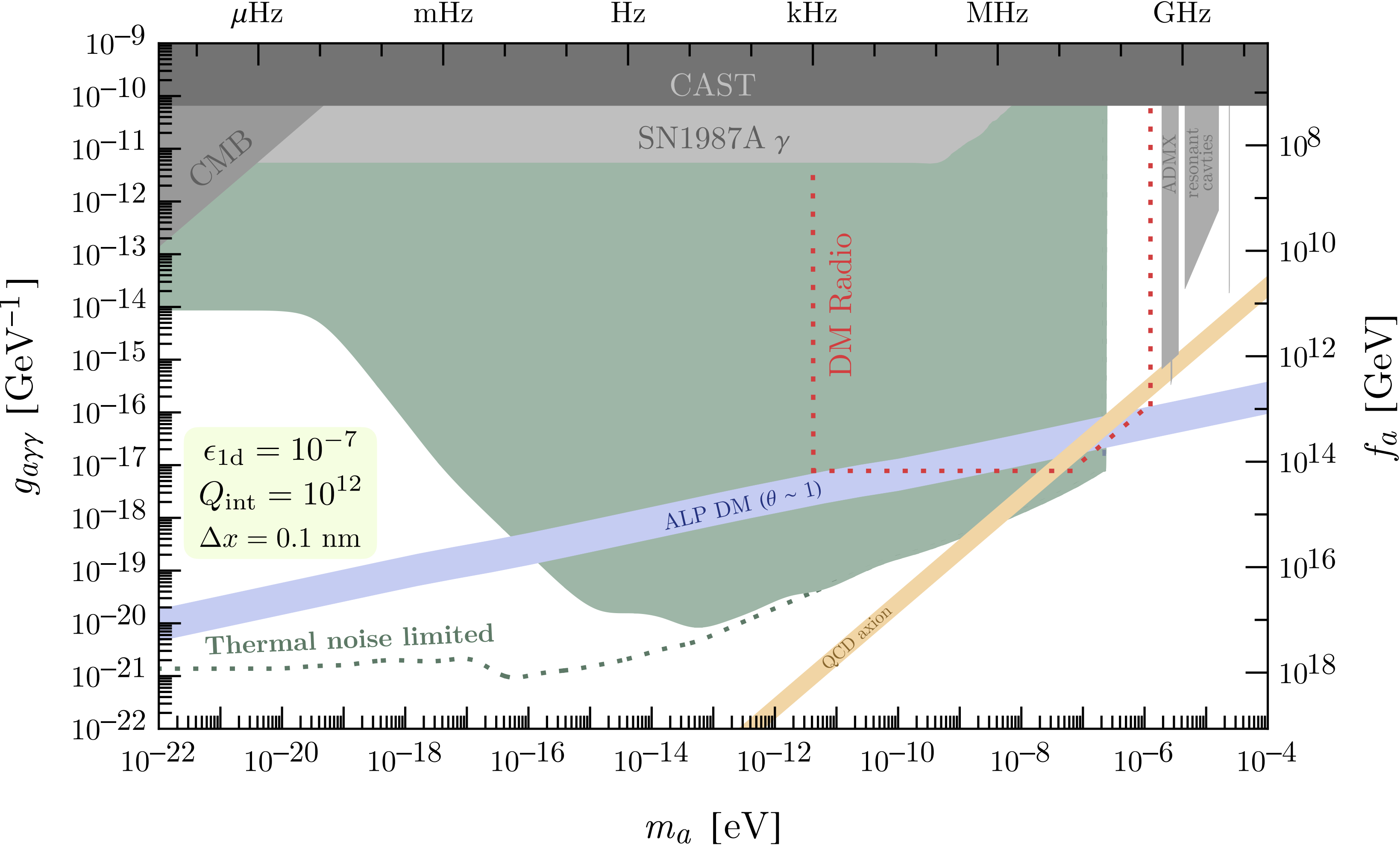
Scientific Reports 8, 15324 (2018) Rosat & Hinderer



Resonant Axion Resonant Frequency Conversion

$B = 0.2 \text{ T}, T = 2\text{K}, \omega_0 = 1 \text{ GHz}$

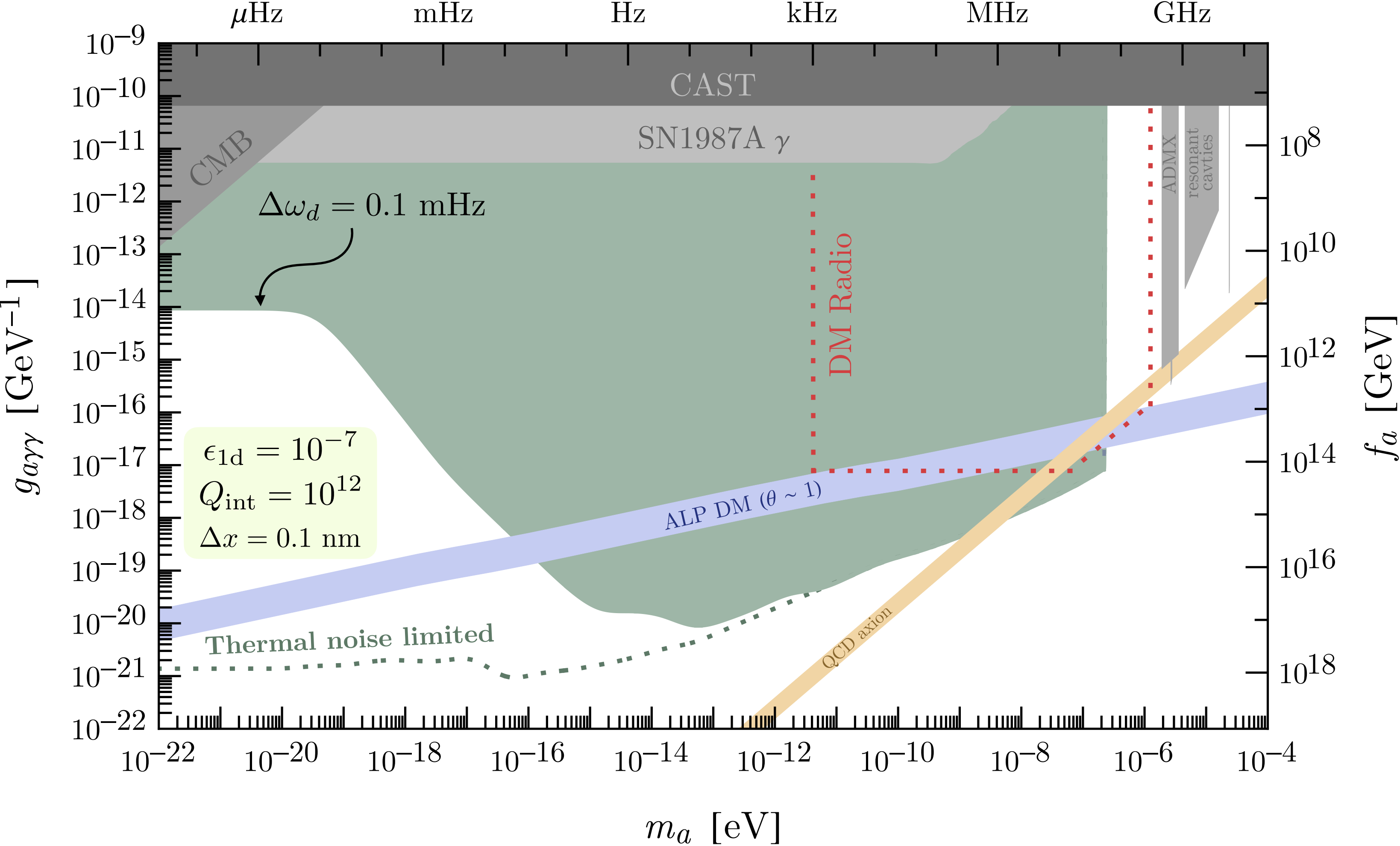
frequency = $m_a/2\pi$



Resonant Axion Resonant Frequency Conversion

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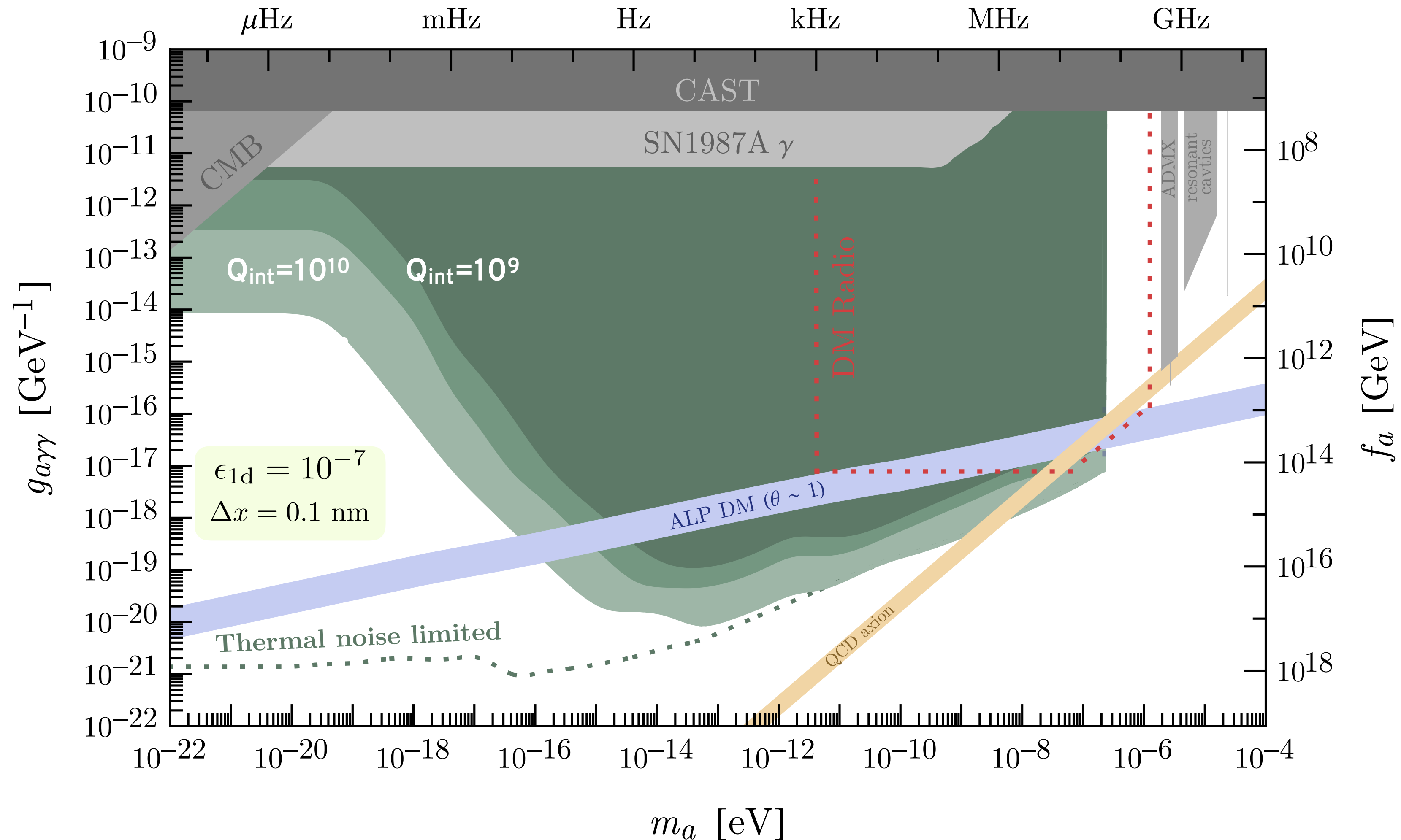
frequency = $m_a/2\pi$



Resonant parameter variations: *Q-factor*

$B = 0.2 \text{ T}, \quad T = 2\text{K}, \quad \omega_0 = 1 \text{ GHz}$

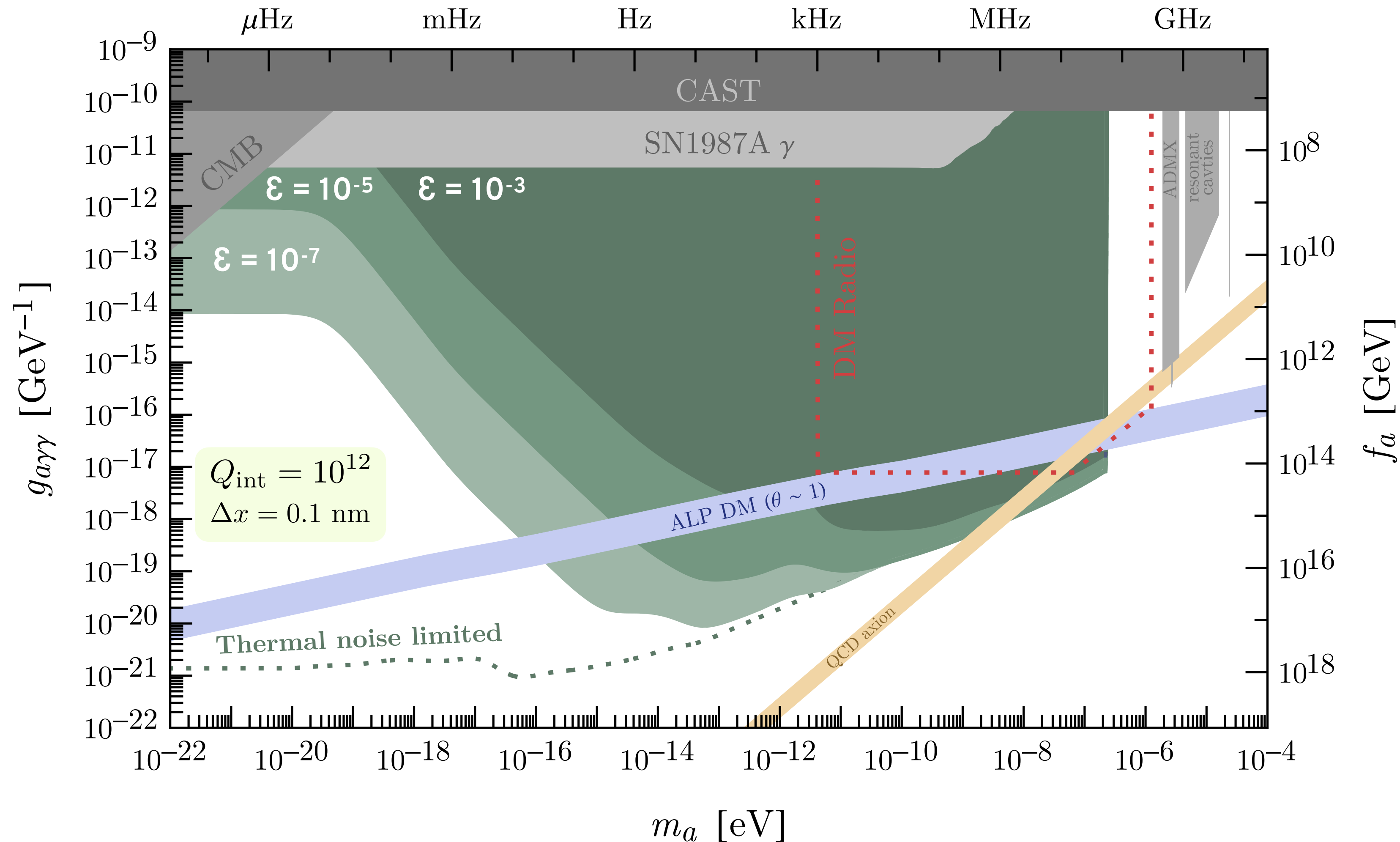
frequency = $m_a/2\pi$



Resonant parameter variations: *mode rejection*

$B = 0.2 \text{ T}, \quad T = 2\text{K}, \quad \omega_0 = 1 \text{ GHz}$

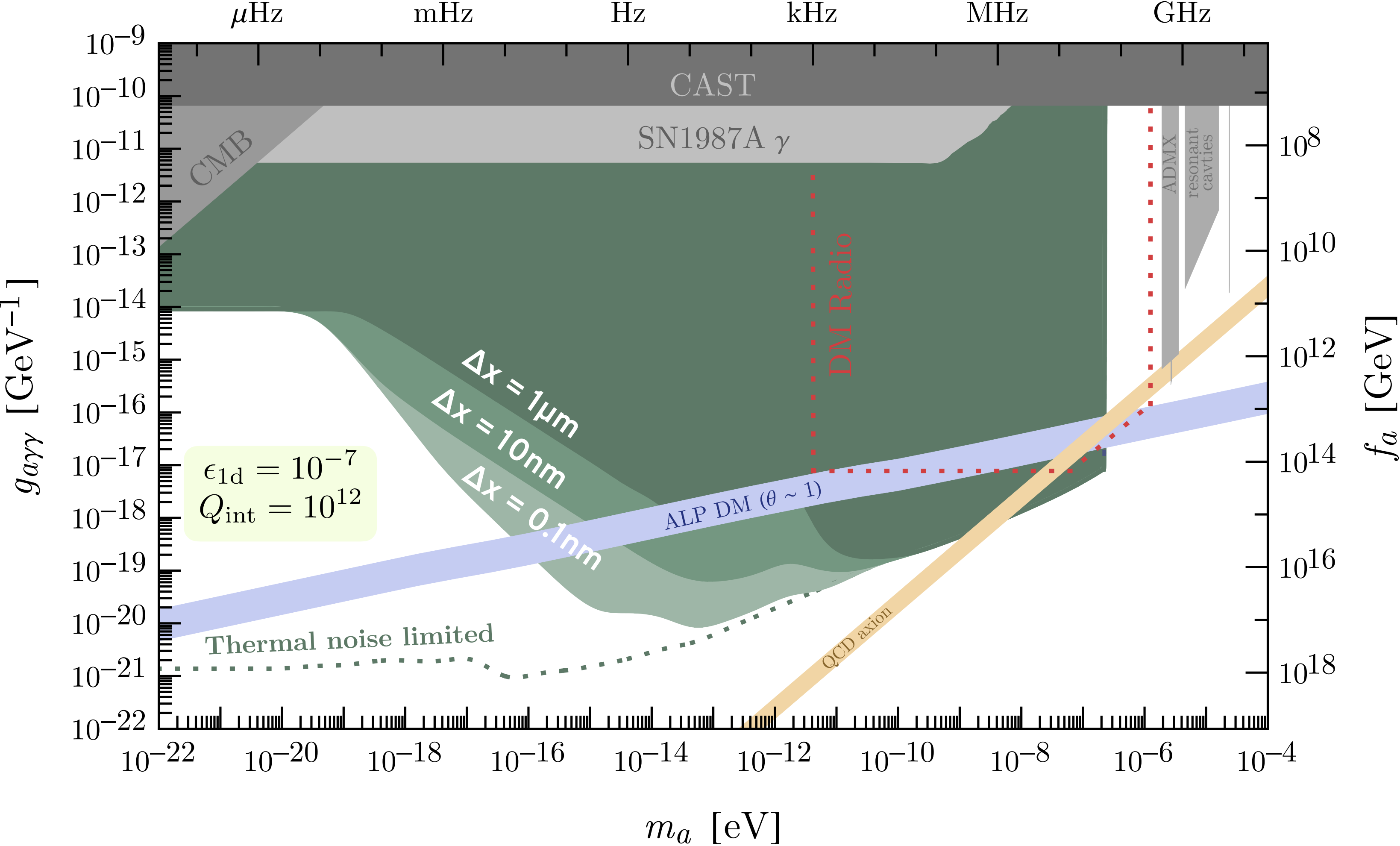
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$B = 0.2 \text{ T}, \quad T = 2\text{K}, \quad \omega_0 = 1 \text{ GHz}$

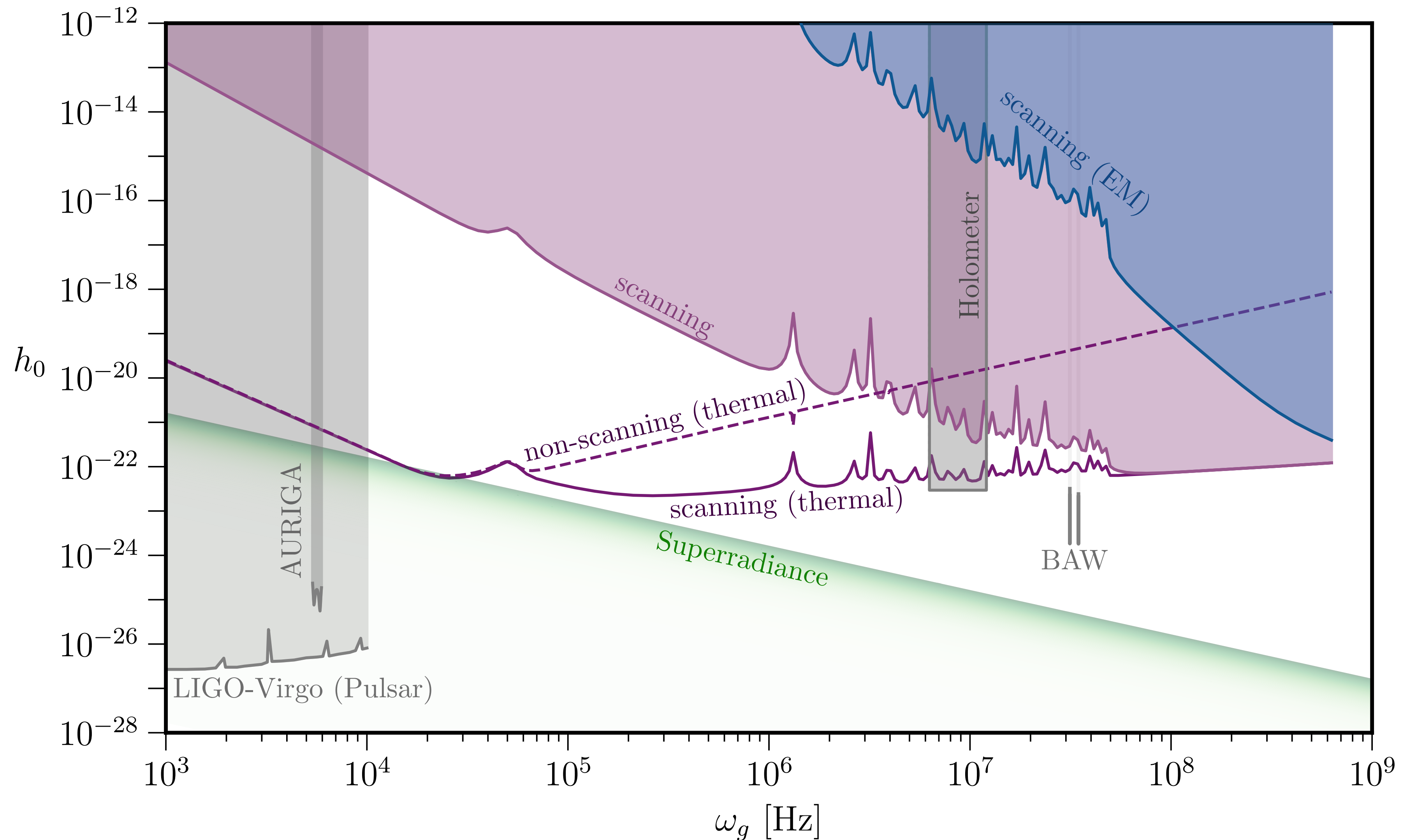
frequency = $m_a/2\pi$



Gravitational Waves BACKUP



MAGO 2.0 sensitivity to coherent GWs



Why SRF Cavity as a Weber Bar?

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Thermal mechanical noise smaller for larger mass $S_{\text{noise}}^{(\text{mech.})} \propto \frac{4T\omega_p}{Q_p M}$

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Thermal EM noise-limited reach scales as $S_h^{(\text{EM})} \propto \frac{1}{Q} \frac{T}{\omega_0 U_0}$

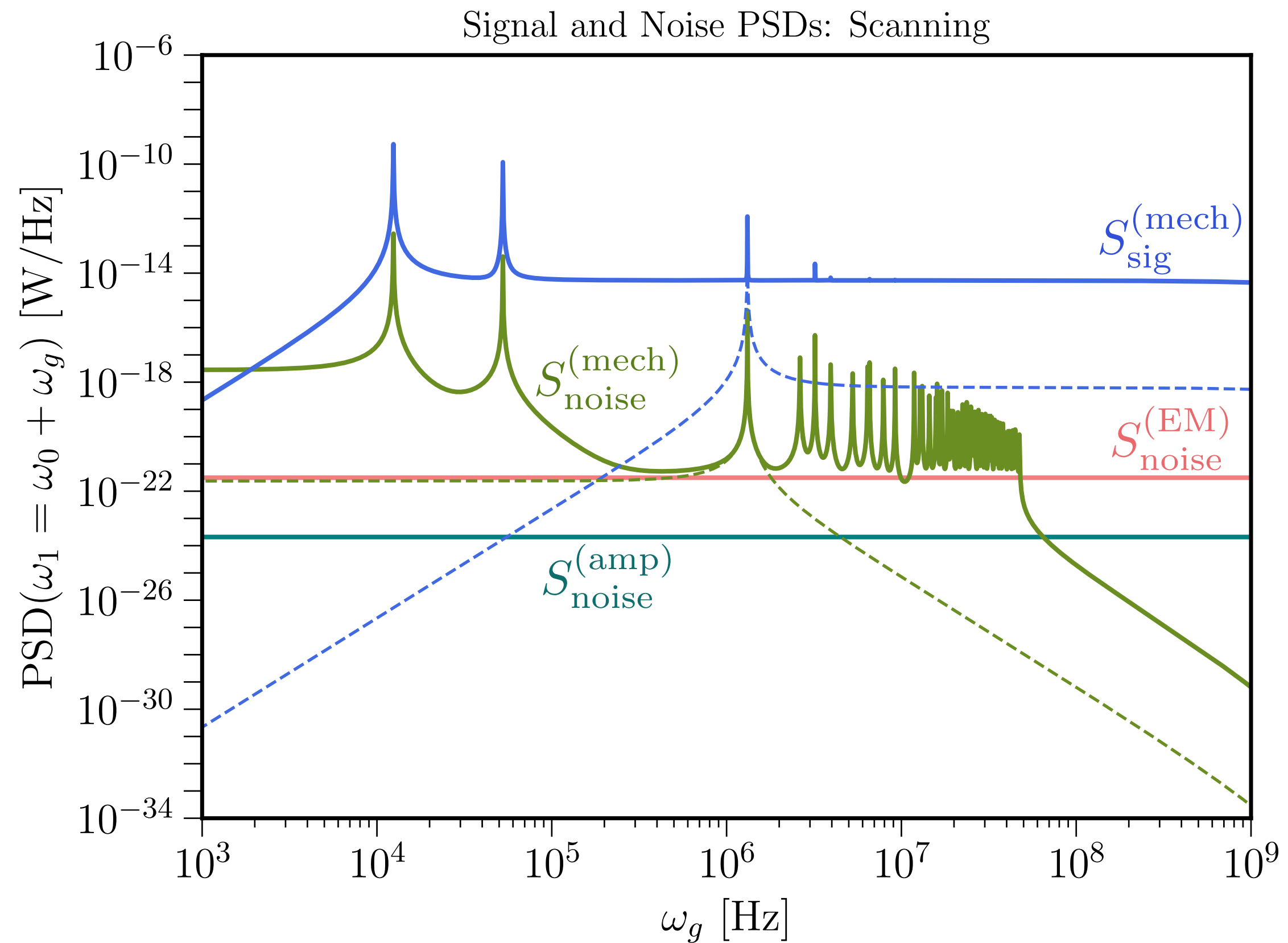
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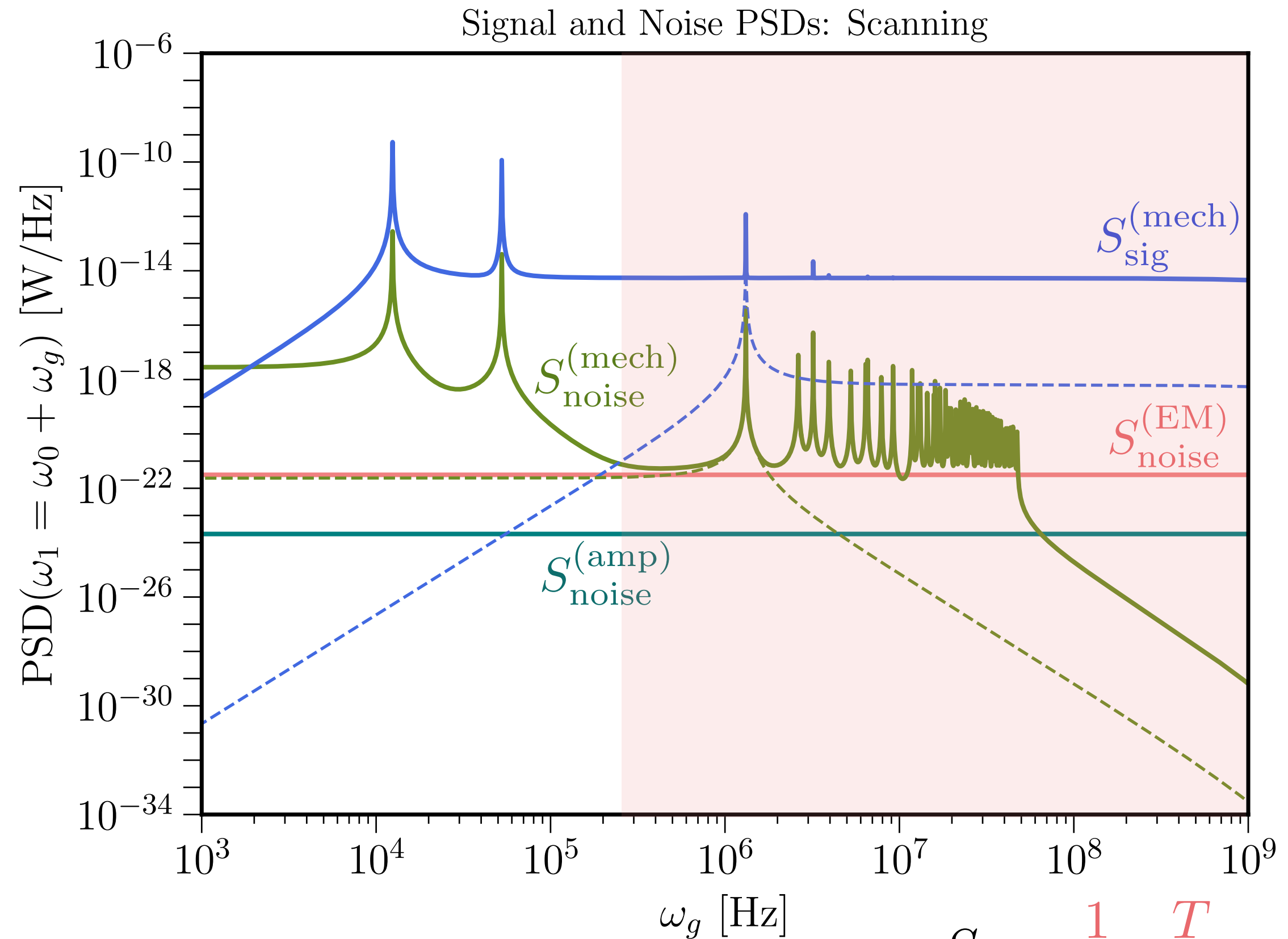
Thermal EM noise-limited reach scales as $S_h^{(\text{EM})} \propto \frac{1}{Q} \frac{T}{\omega_0 U_0}$

SRF Cavity w/ $Q \sim 10^{10}$ means better sensitivity

Noise in MAGO 2.0



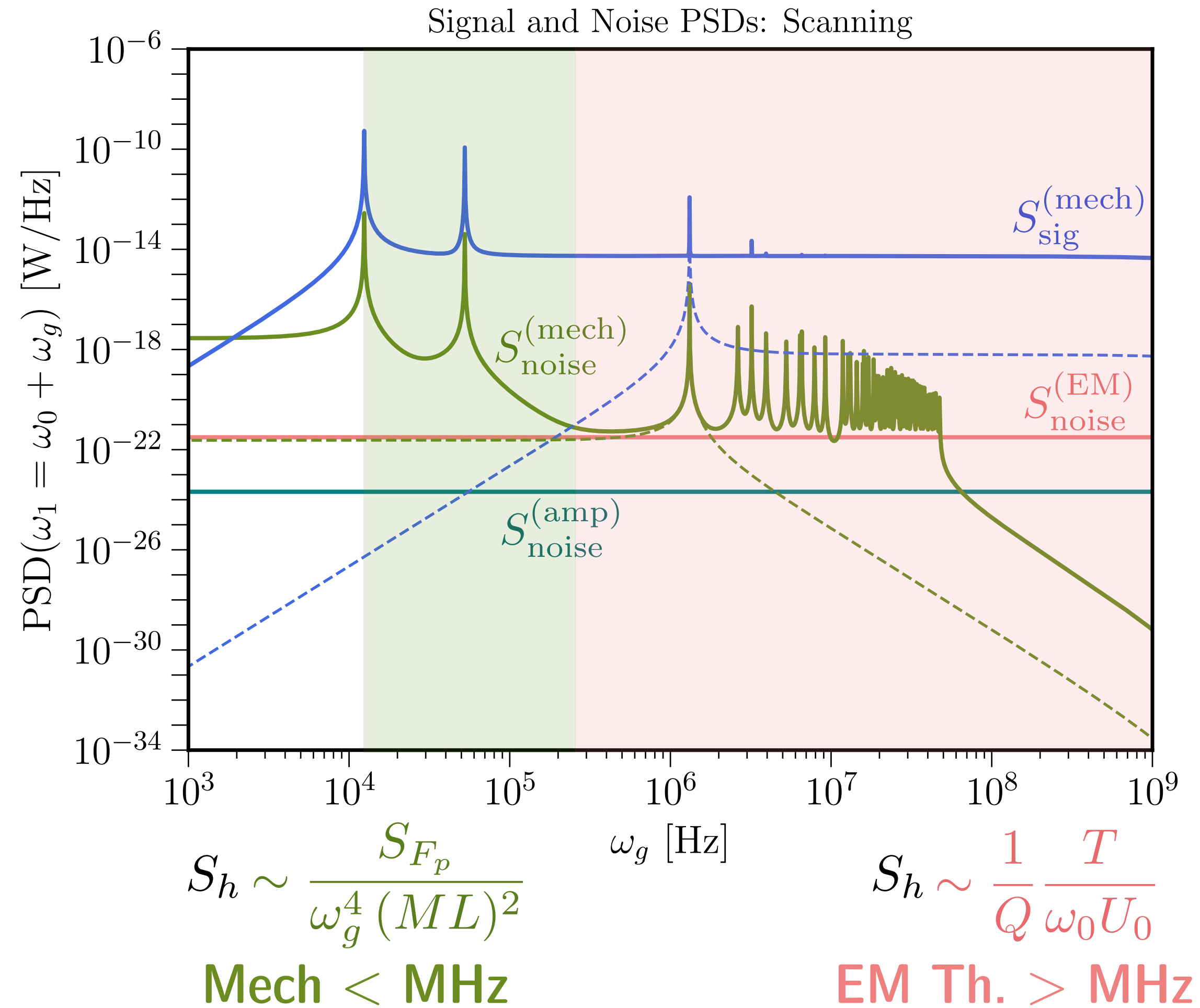
Noise in MAGO 2.0



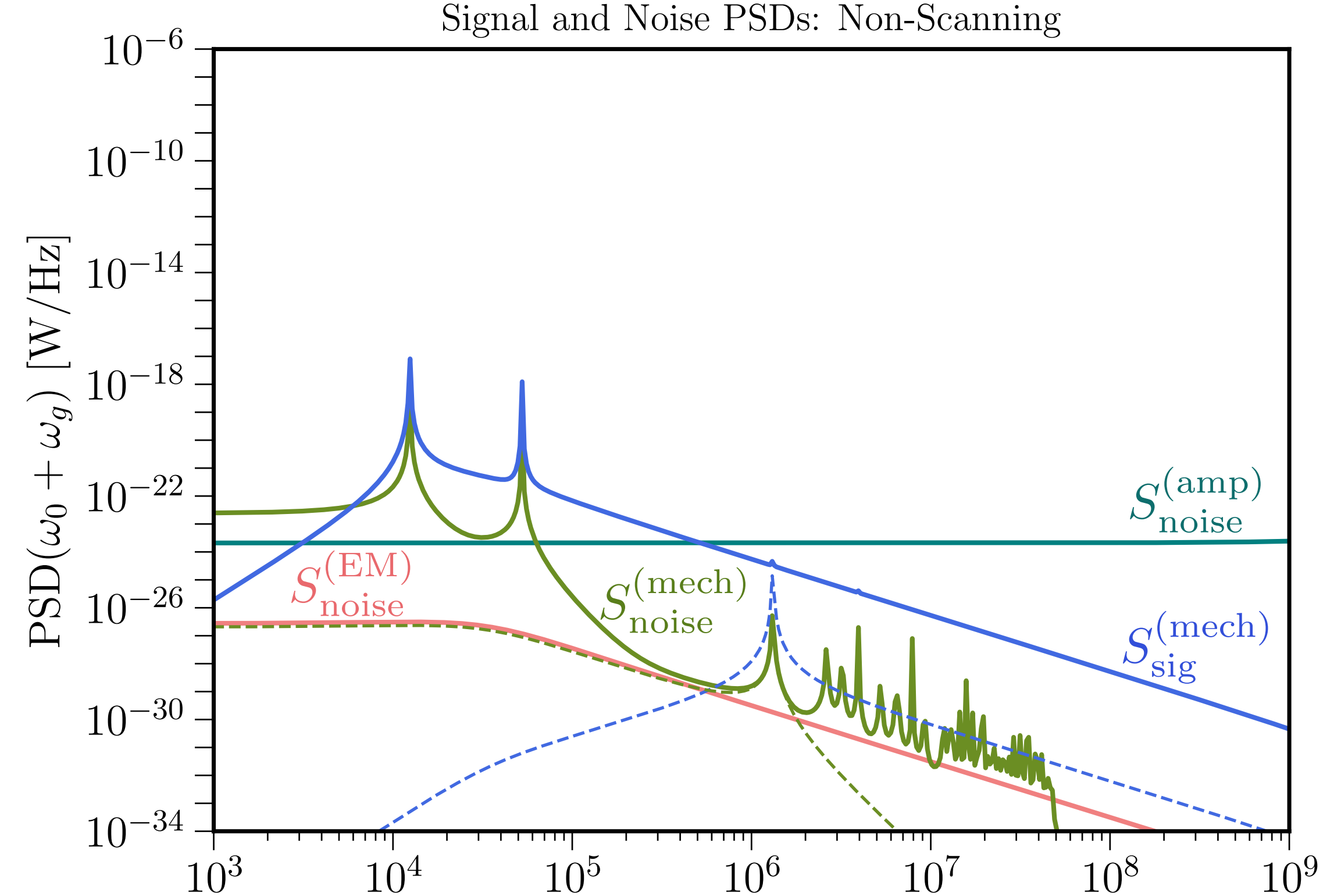
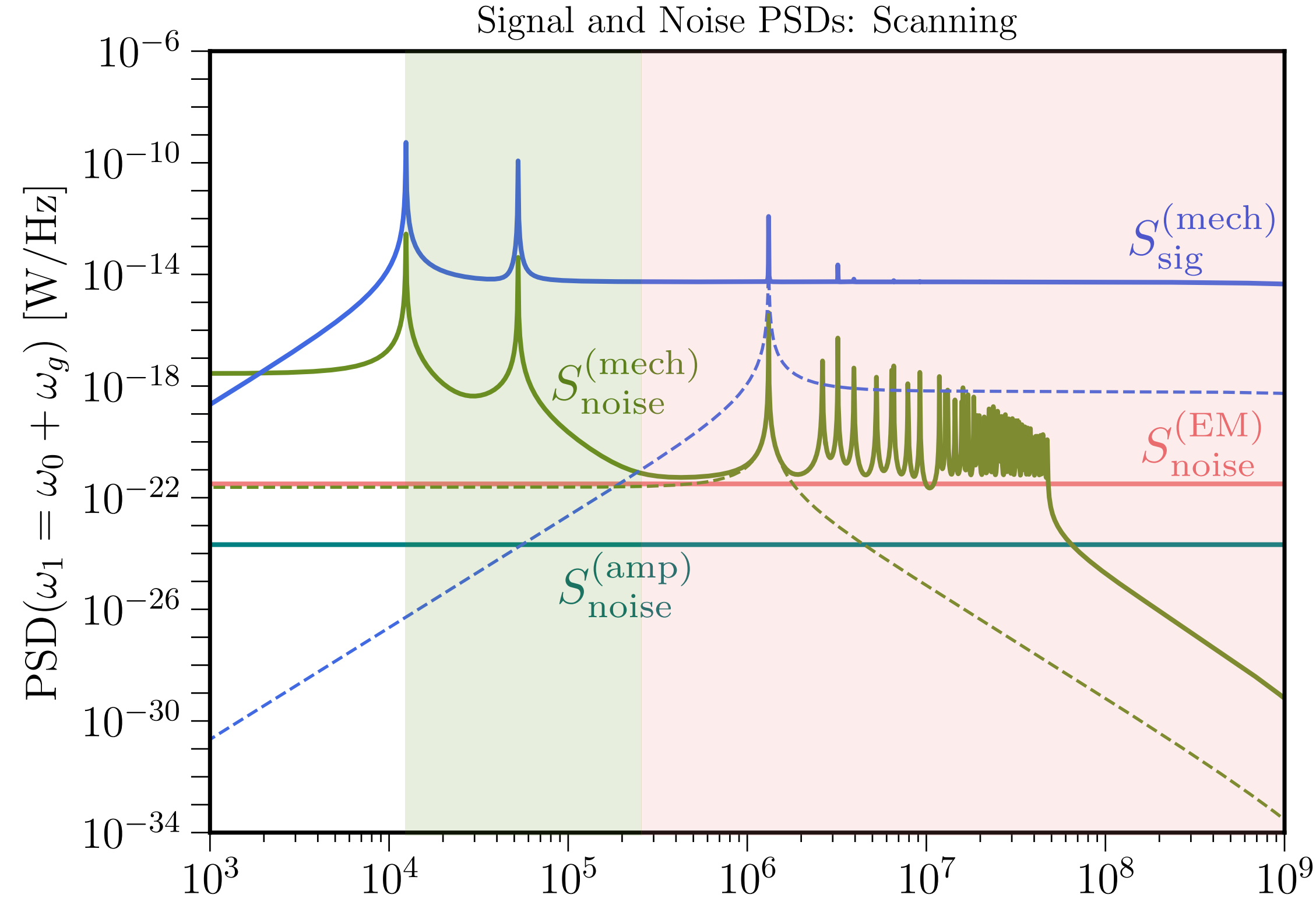
$$S_h \sim \frac{1}{Q} \frac{T}{\omega_0 U_0}$$

EM Th. > MHz

Noise in MAGO 2.0



Noise in MAGO 2.0



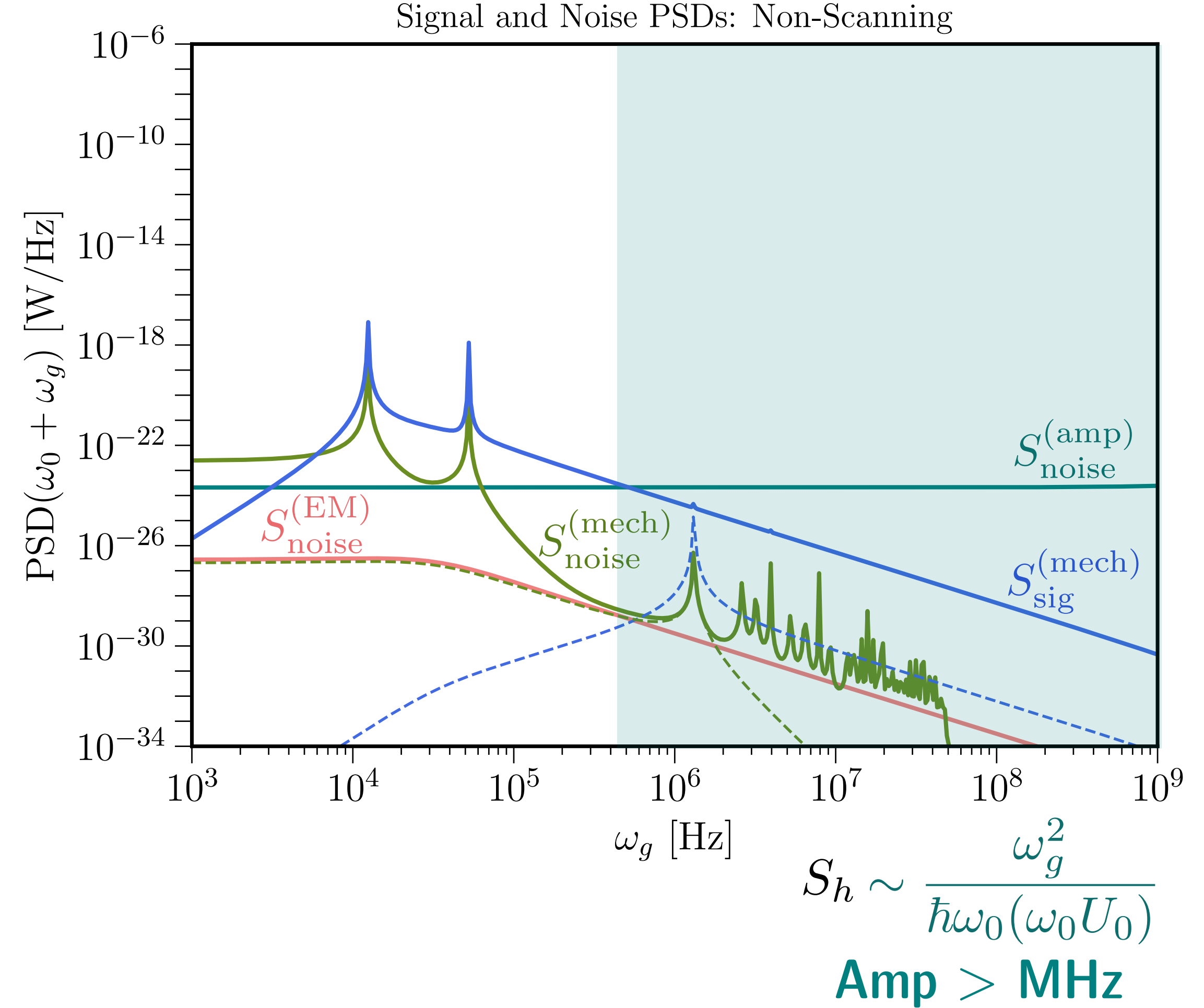
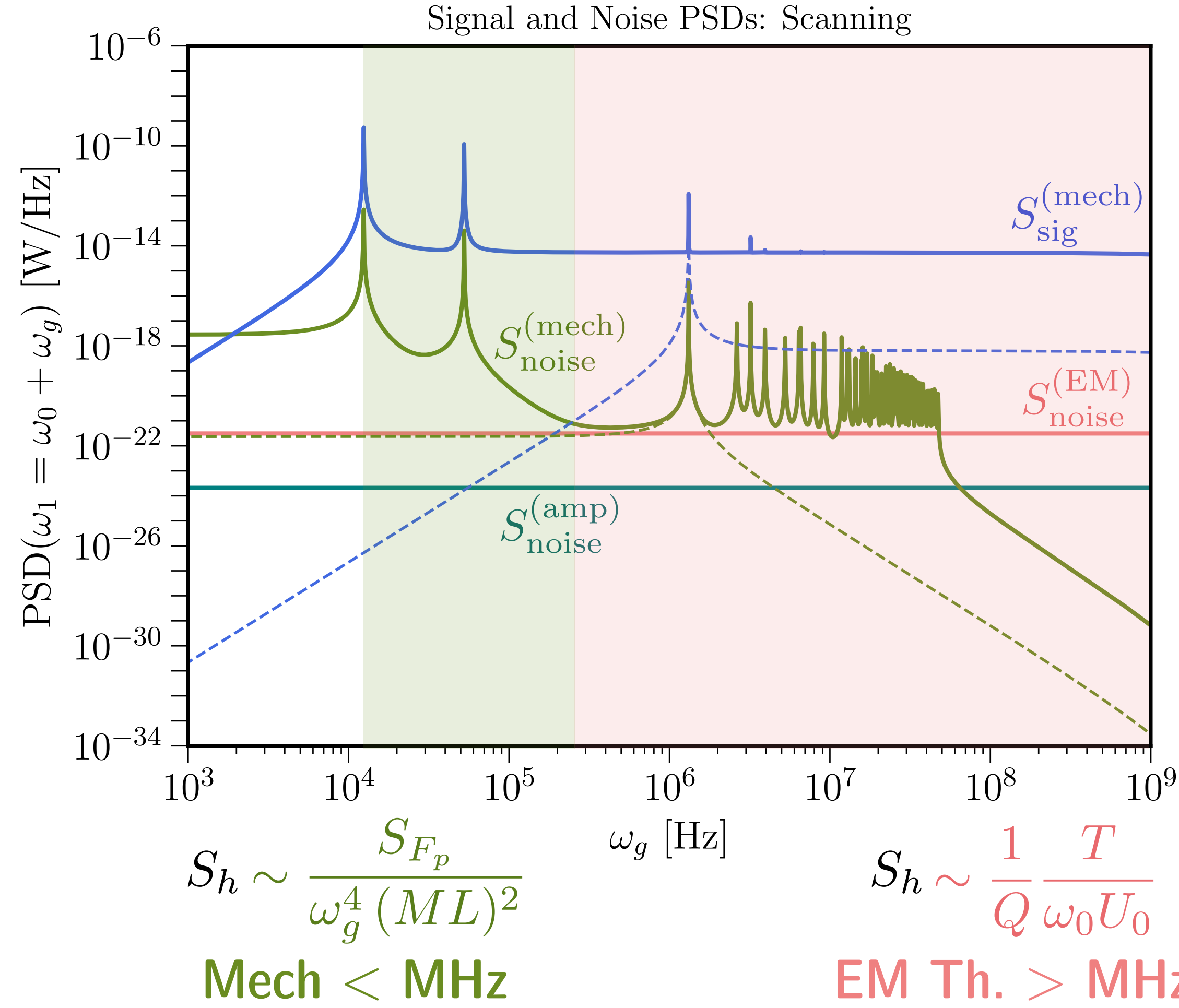
$$S_h \sim \frac{S_{F_p}}{\omega_g^4 (ML)^2}$$

Mech < MHz

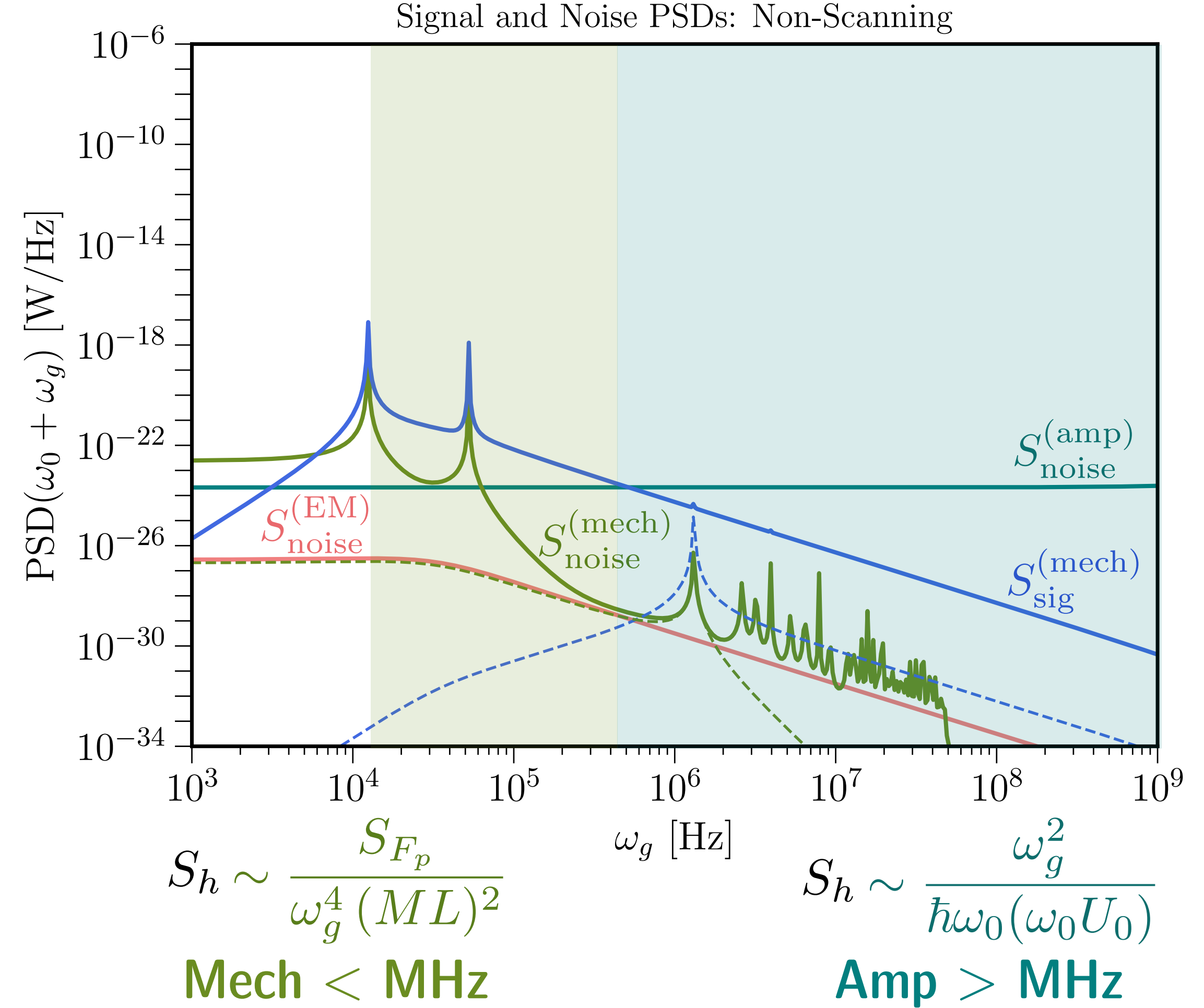
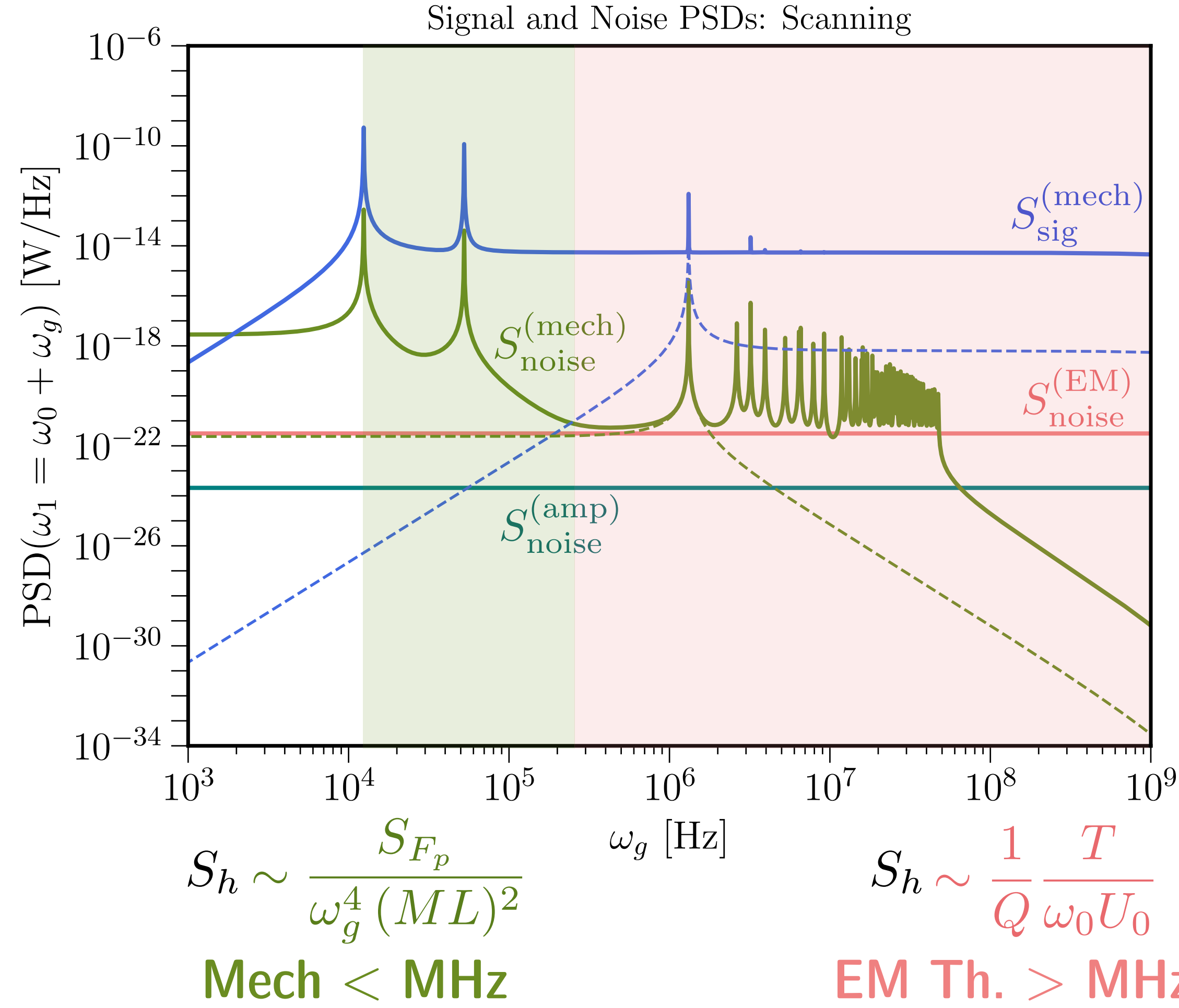
$$S_h \sim \frac{1}{Q} \frac{T}{\omega_0 U_0}$$

EM Th. > MHz

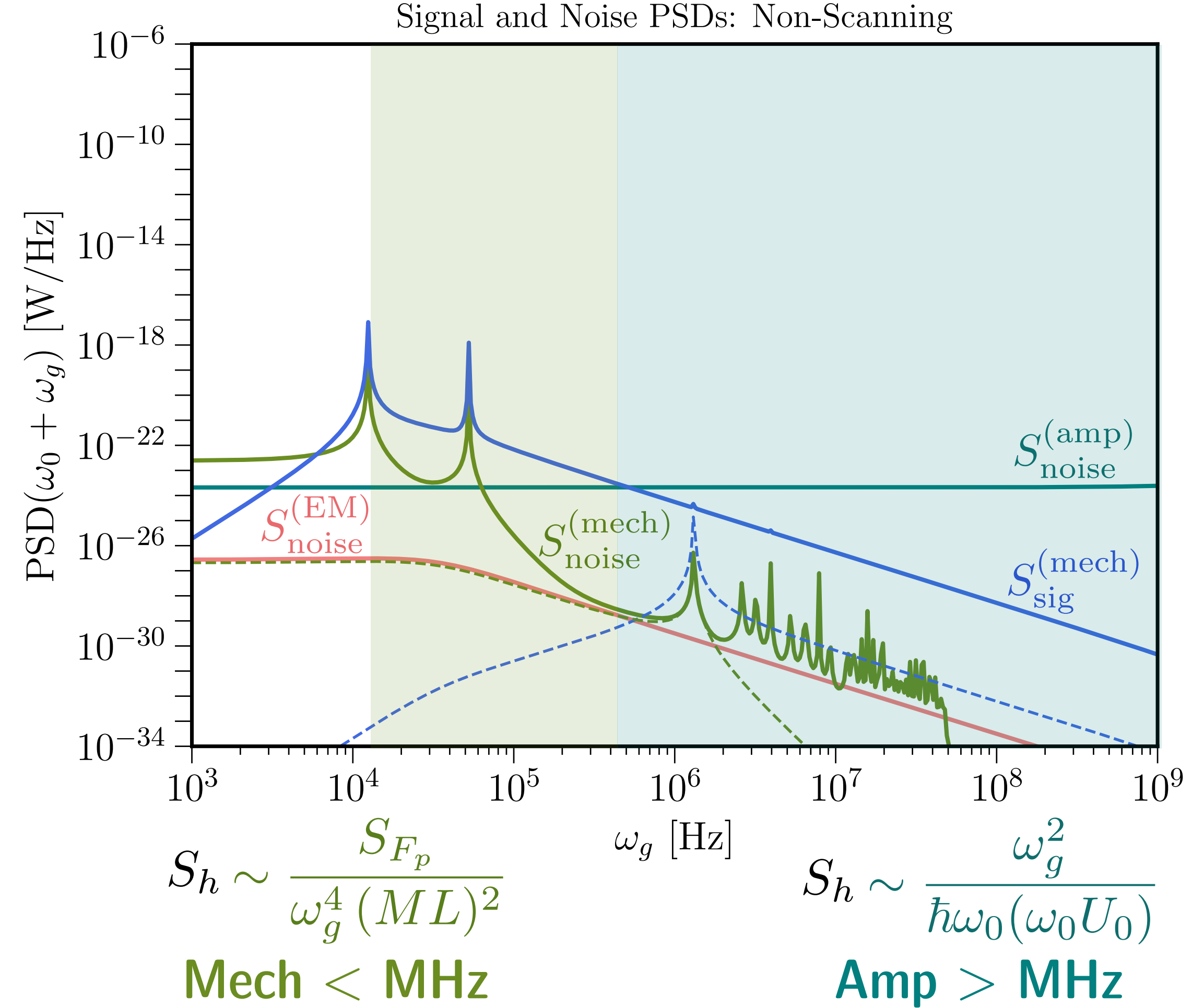
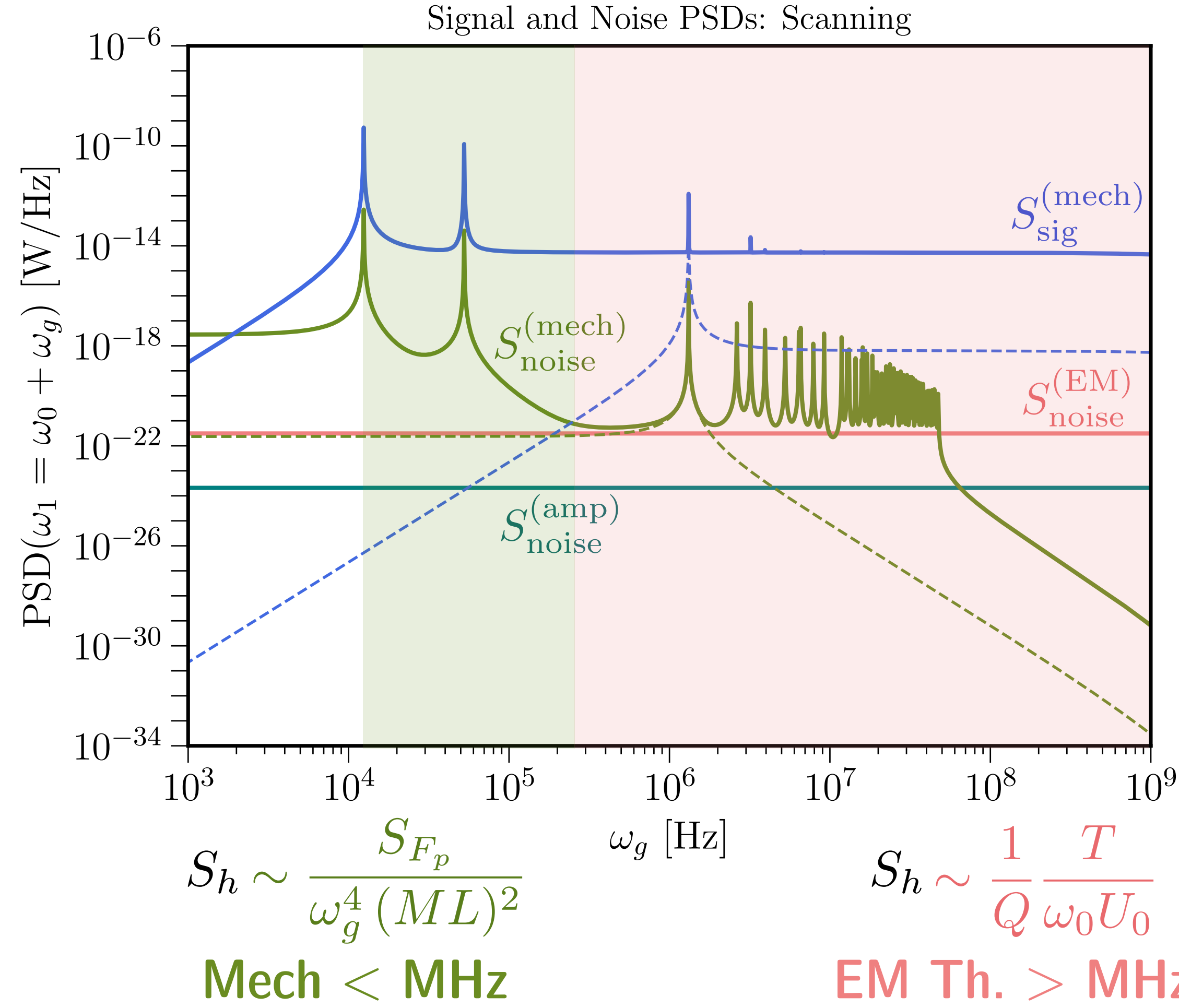
Noise in MAGO 2.0



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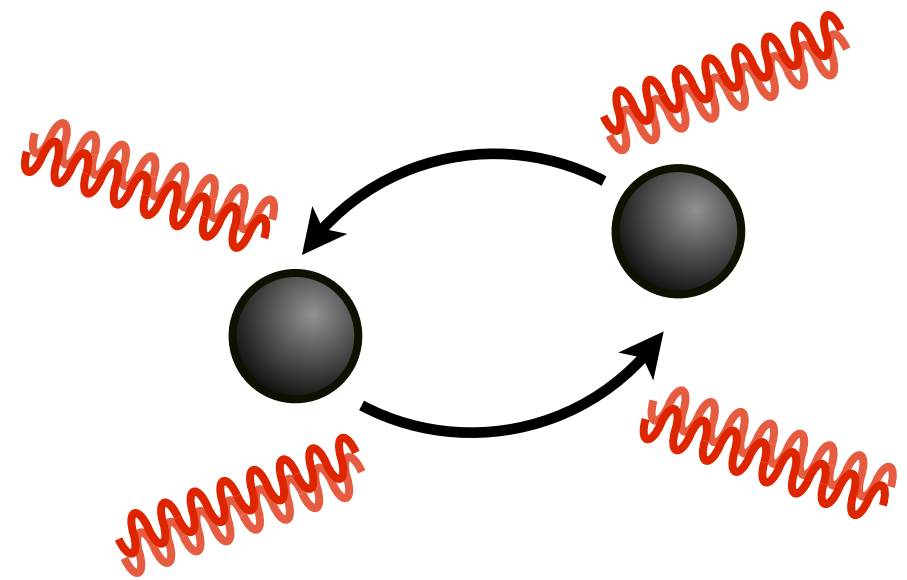
Noise in MAGO 2.0



NB: missing radiation damping effect studied in Löwenberg, Moortgat-Pick: 2307.14379

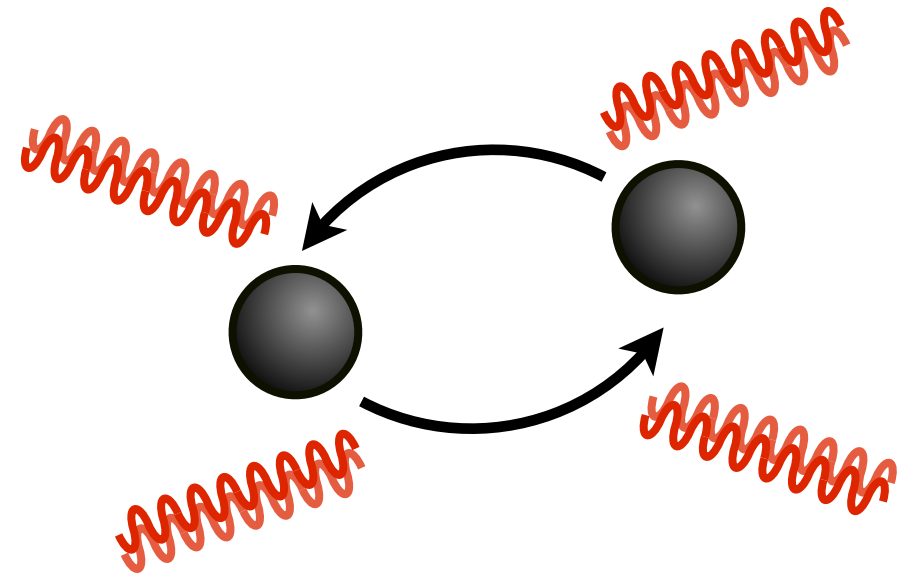
Is this good enough to see BSM sources?

Primordial Black Hole binary inspirals



Is this good enough to see BSM sources?

Primordial Black Hole binary inspirals

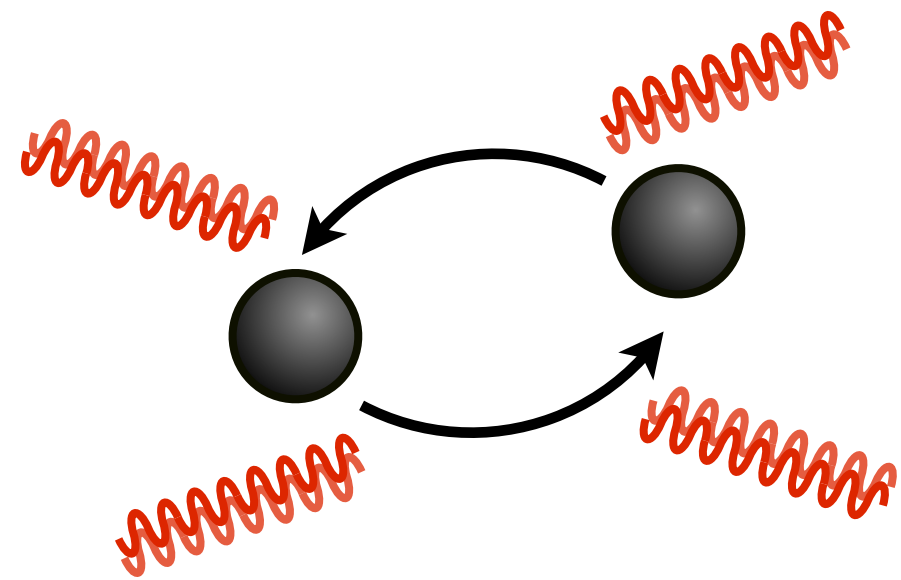


$$\omega_g \simeq 14 \text{ GHz} \times (10^{-6} M_{\odot}/M_b) (r_{\text{ISCO}}/r_b)^{3/2}$$

$$d\omega_g/dt \propto (M_b/r_b)^{11/6}$$

Is this good enough to see BSM sources?

Primordial Black Hole binary inspirals



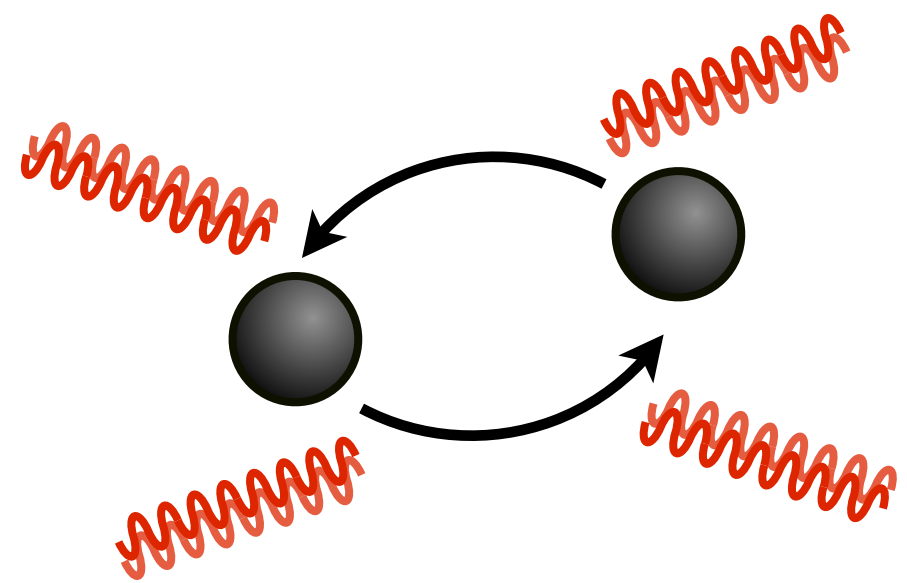
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$$h_0 \sim 10^{-29} \times \left(\frac{1 \text{ pc}}{D} \right) \left(\frac{M_b}{10^{-11} M_\odot} \right)^{5/3} \left(\frac{\omega_g}{1 \text{ GHz}} \right)^{2/3}$$

Is this good enough to see BSM sources?

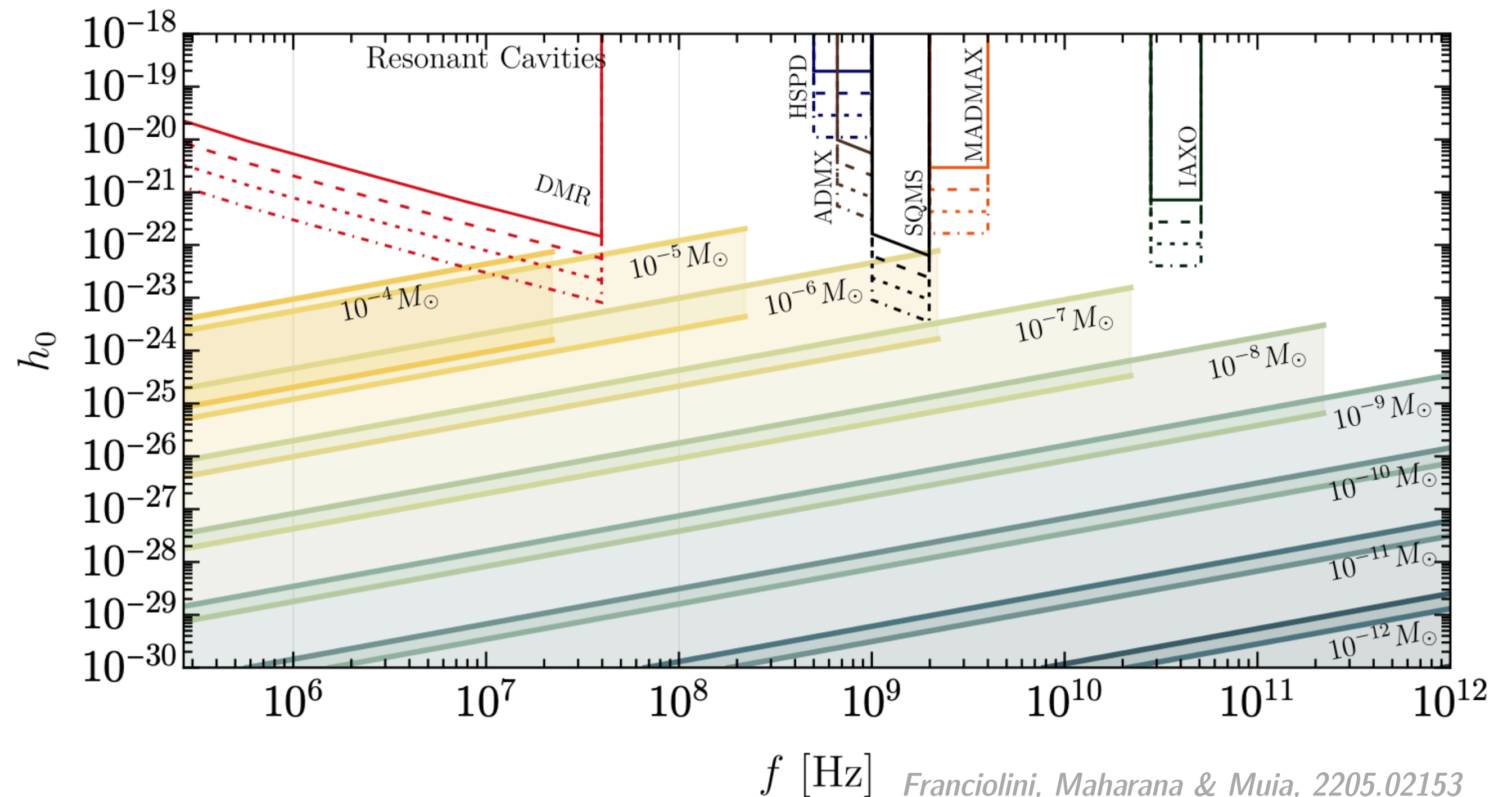
Primordial Black Hole binary inspirals



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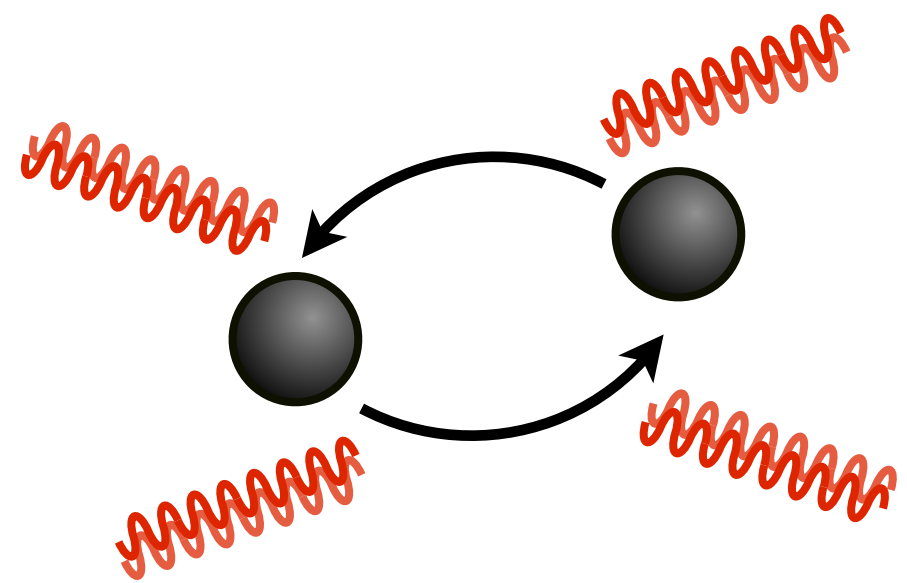
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Is this good enough to see BSM sources?

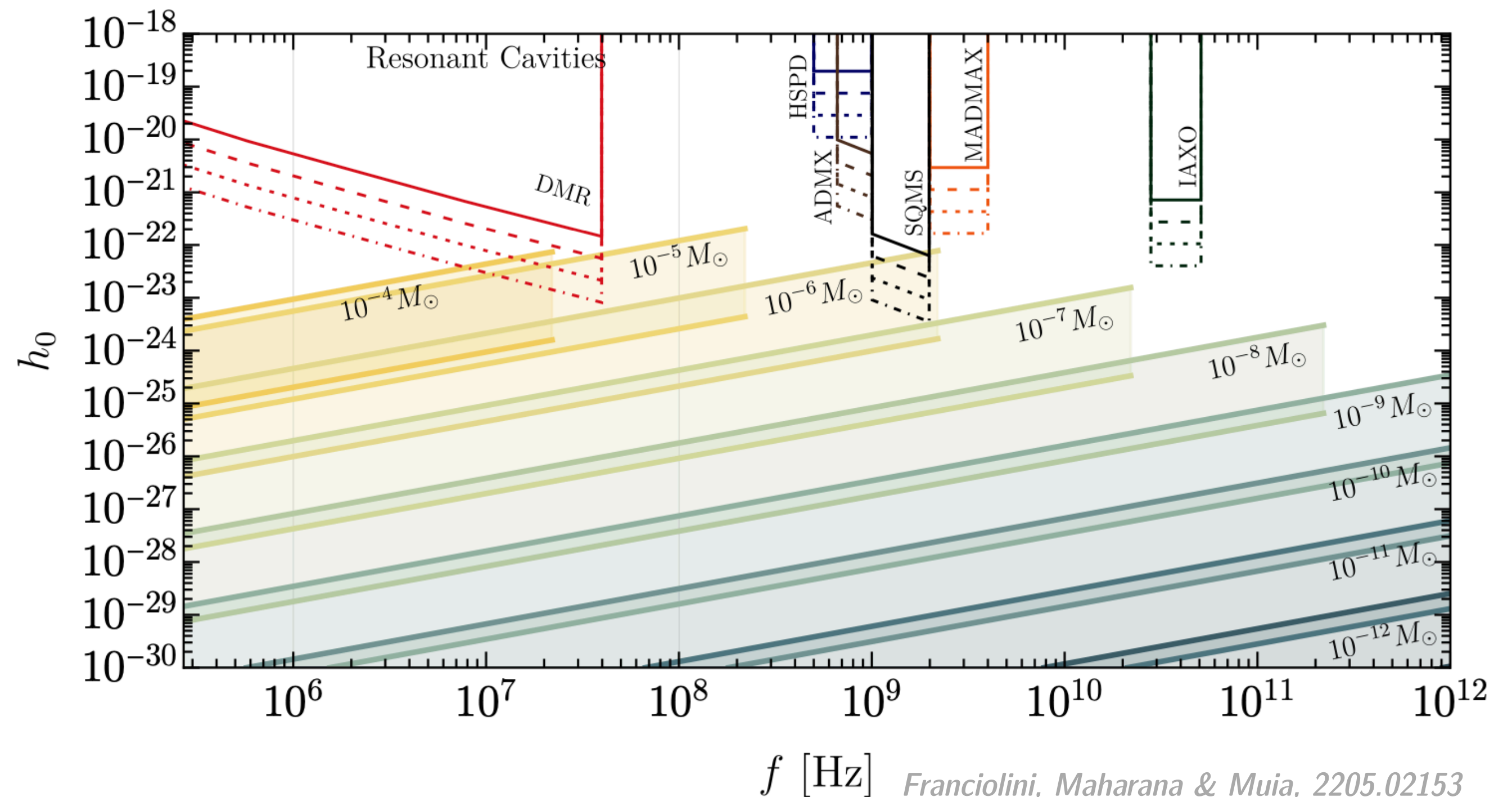
Primordial Black Hole binary inspirals



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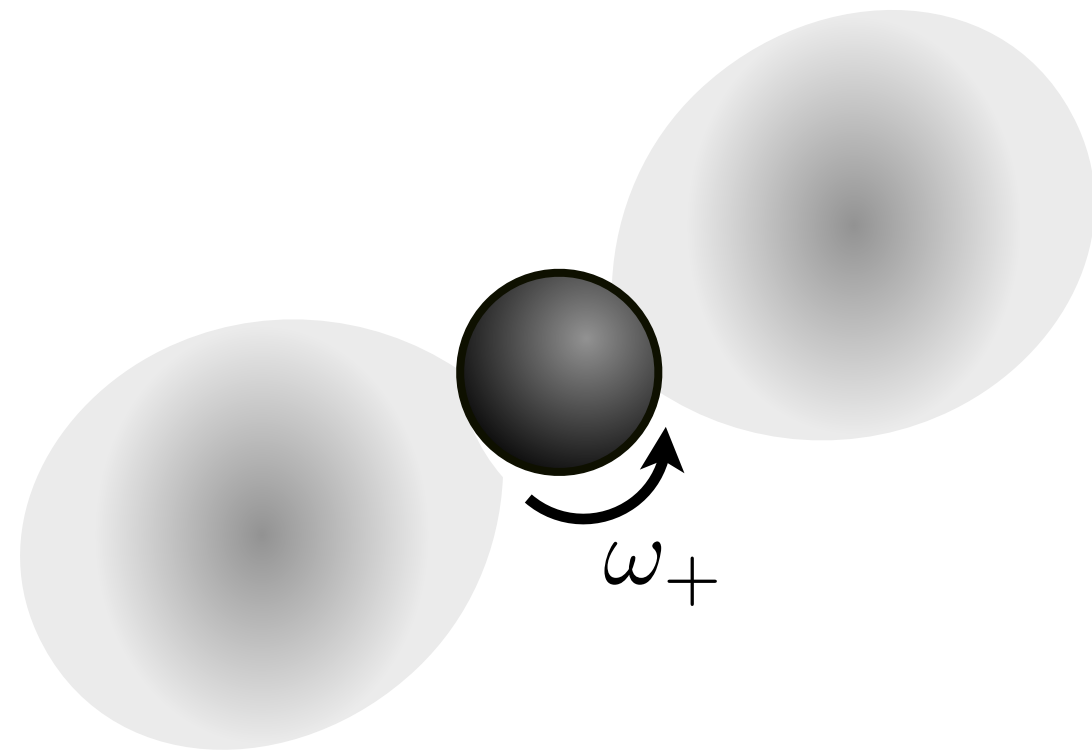
difficult at GHz, plausible at kHz

Is this good enough to see BSM sources?

Superradiance (BSM^2)

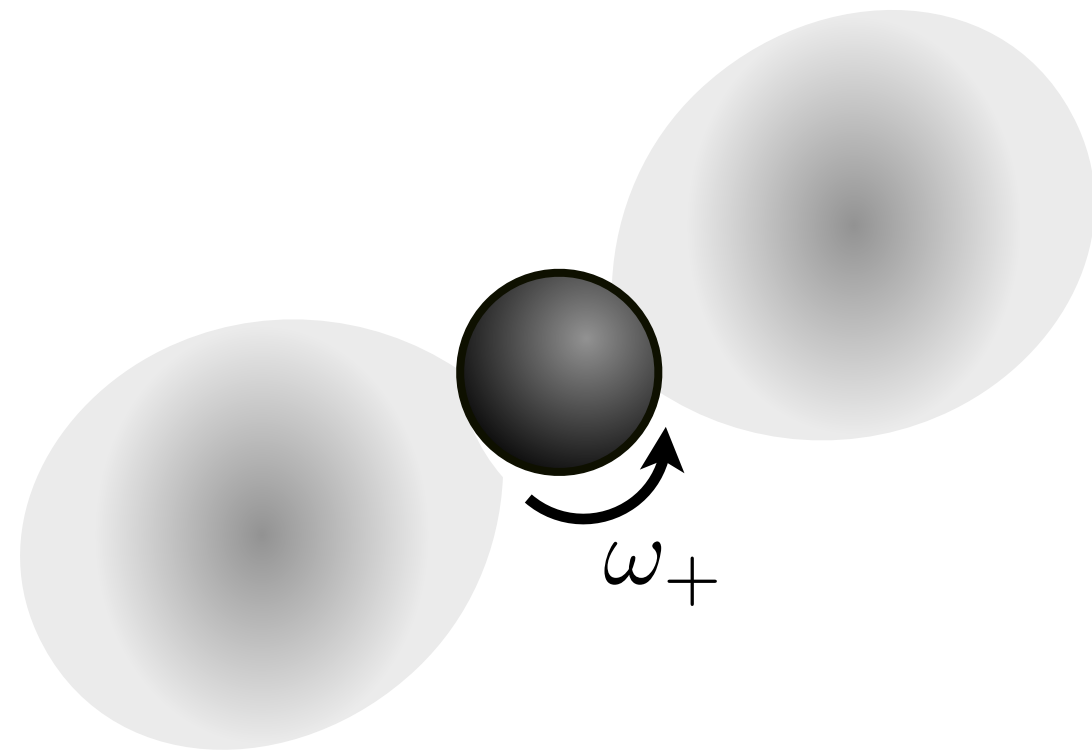
Is this good enough to see BSM sources?

Superradiance (BSM^2)



Is this good enough to see BSM sources?

Superradiance (BSM²)



Ternov et al (1978)

Zouros & Eardley (1979)

Arvanitaki et al (2009)

Arvanitaki & Dubovsky (2010)

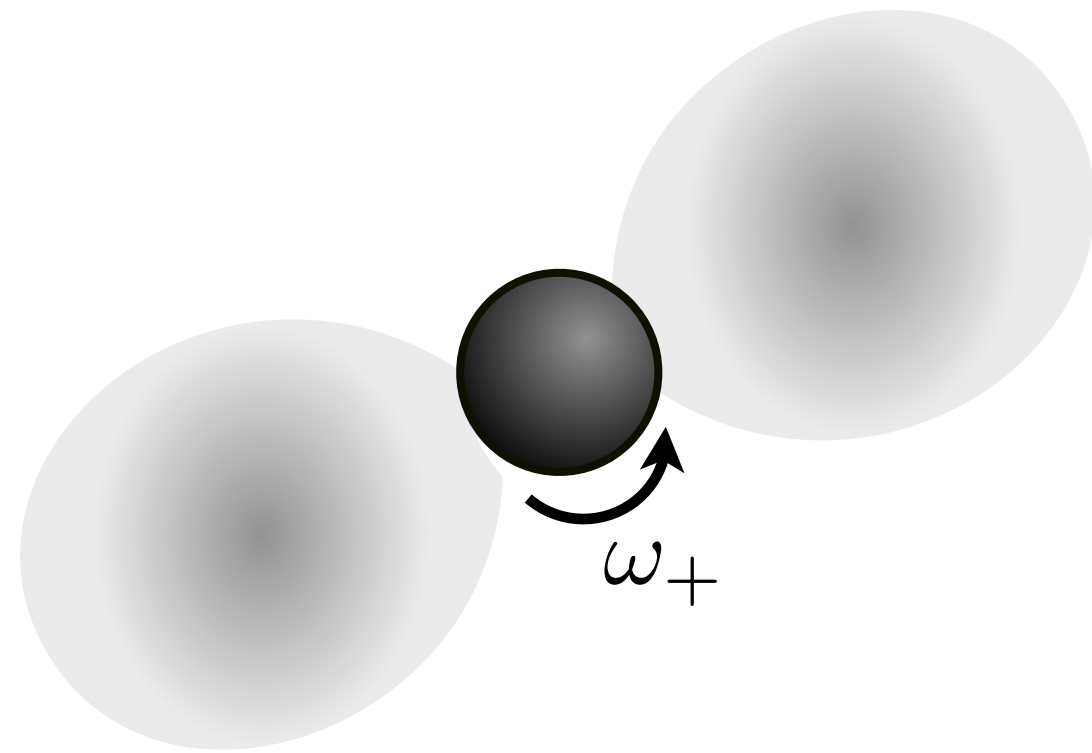
...

Is this good enough to see BSM sources?

Superradiance (BSM²)

$$0 < m_a \leq m\omega_+$$

$$\omega_+ = \frac{1}{4GM_{\text{BH}}} \frac{\chi_i}{1 + \sqrt{1 - \chi_i^2}}$$



Ternov et al (1978)

Zouros & Eardley (1979)

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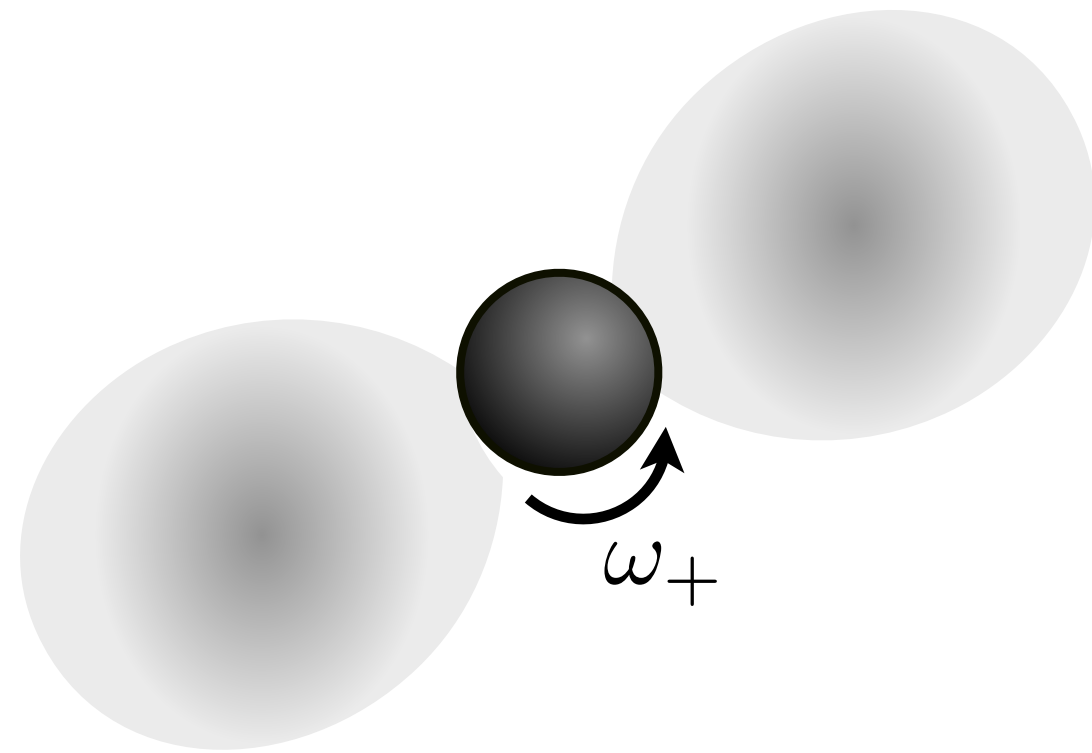
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$$m_a \sim \mu\text{eV} \times (10^{-4} M_{\odot}/M_{\text{PBH}})$$



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...

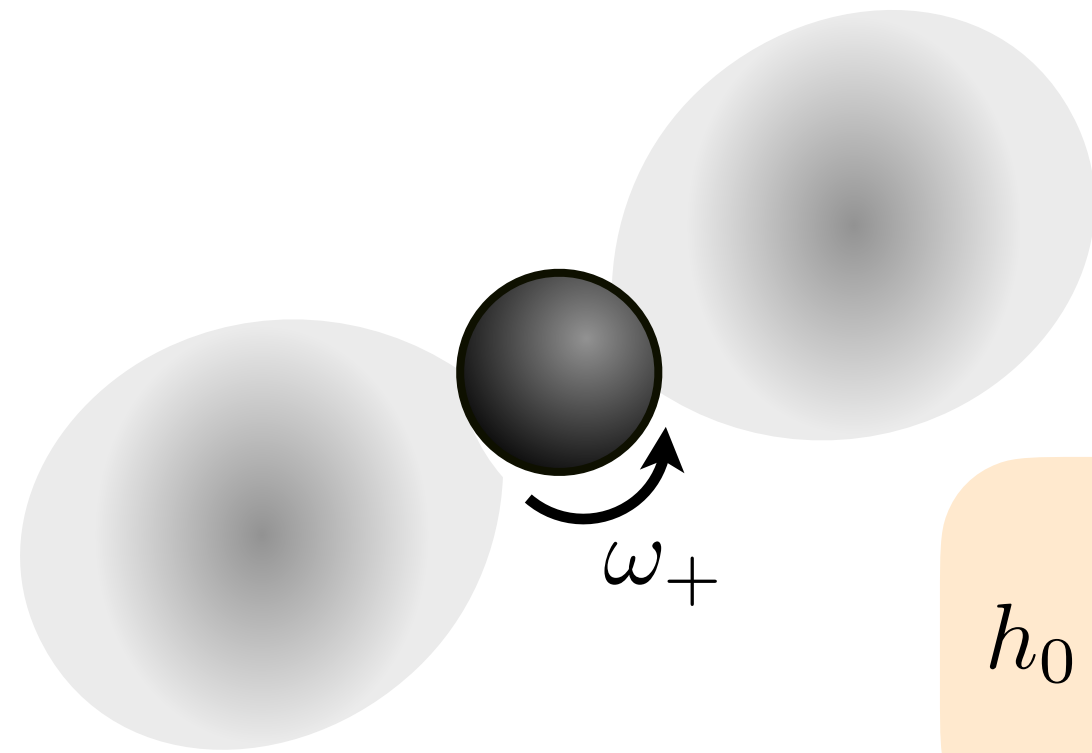
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$$h_0 \sim 10^{-27} \times \left(\frac{10 \text{ kpc}}{D} \right) \left(\frac{M_{\text{PBH}}}{10^{-4} M_{\odot}} \right)$$

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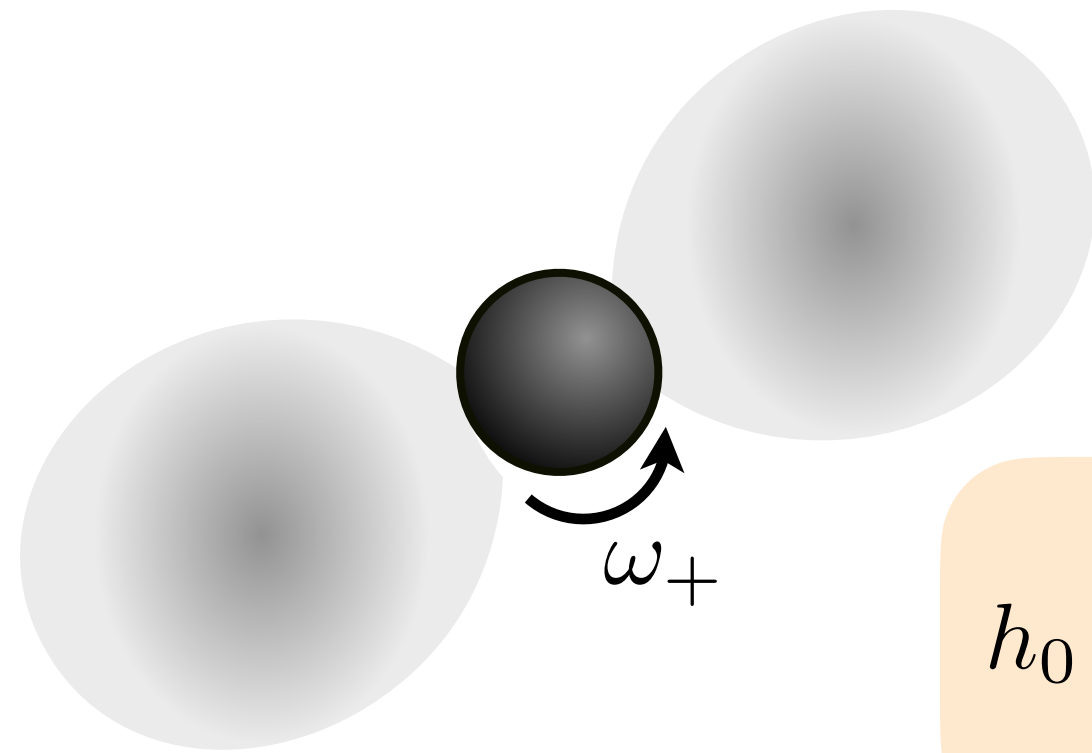
...

Is this good enough to see BSM sources?

Superradiance (BSM²)

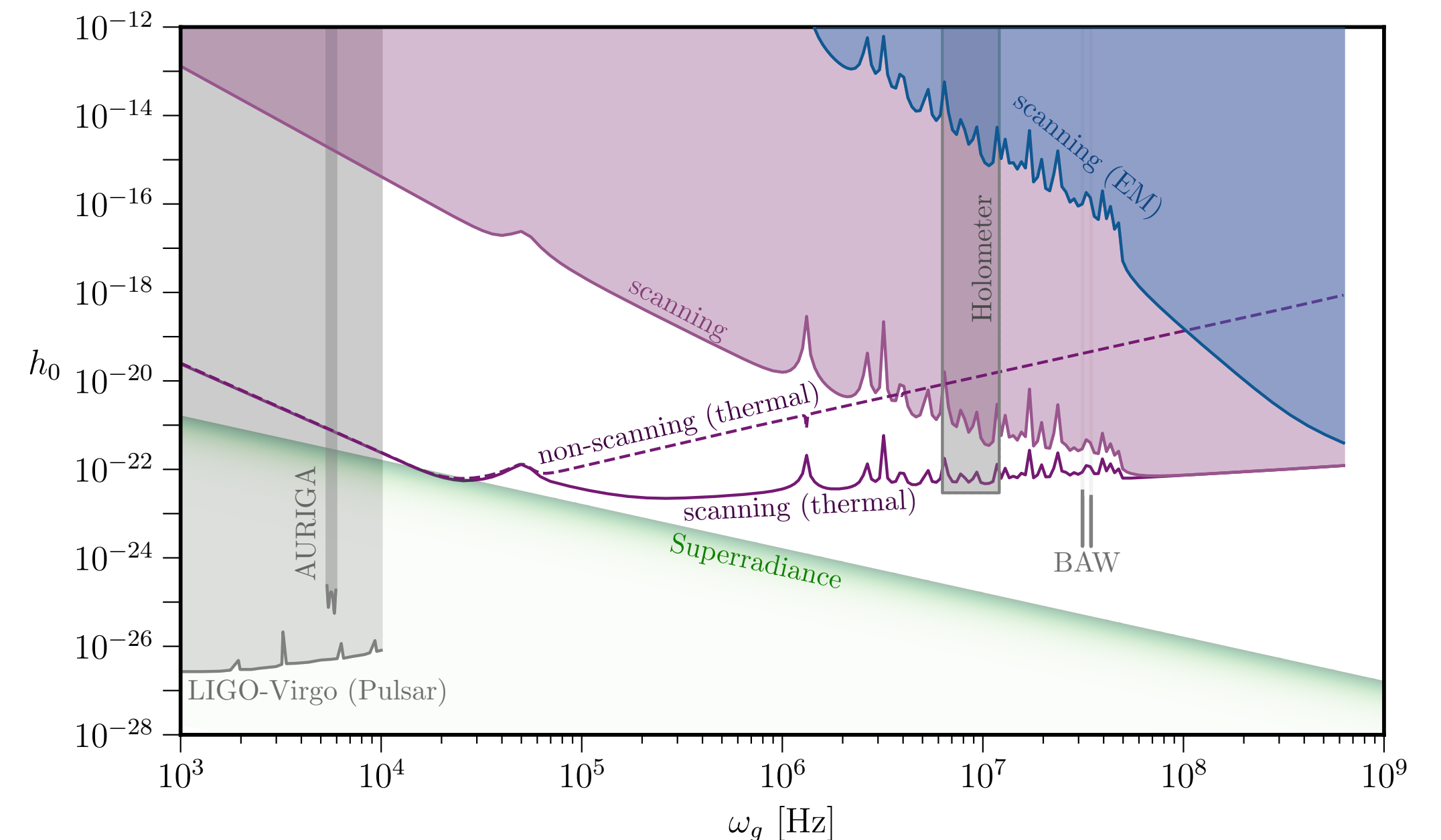
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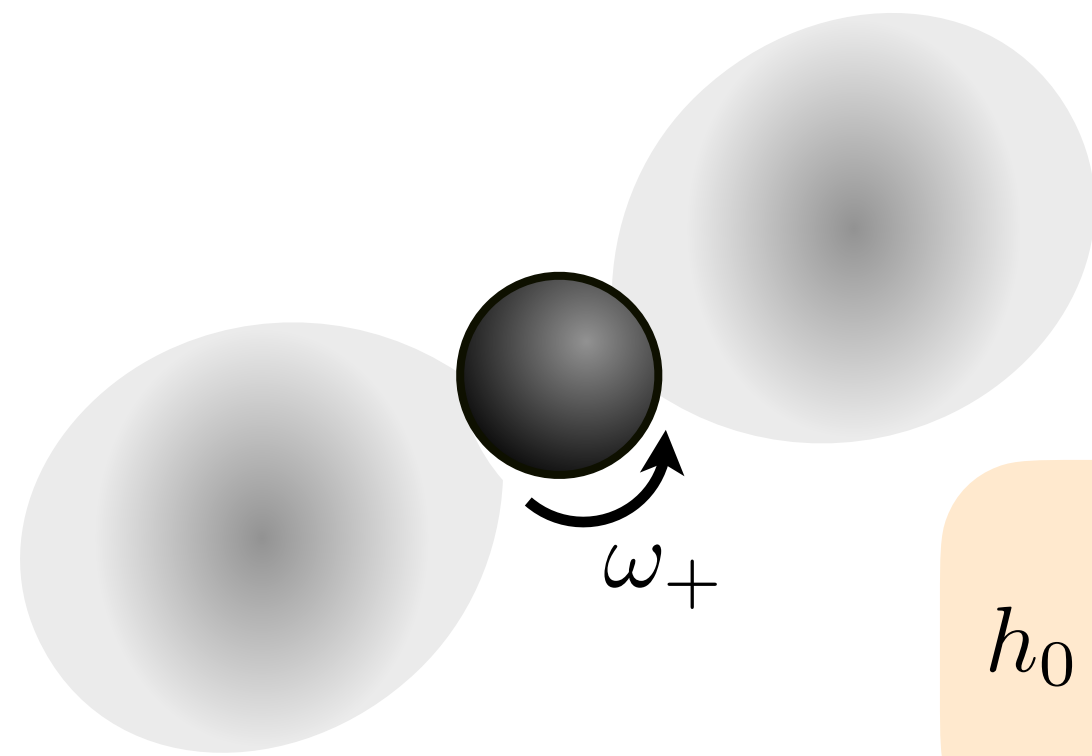


Is this good enough to see BSM sources?

Superradiance (BSM²)

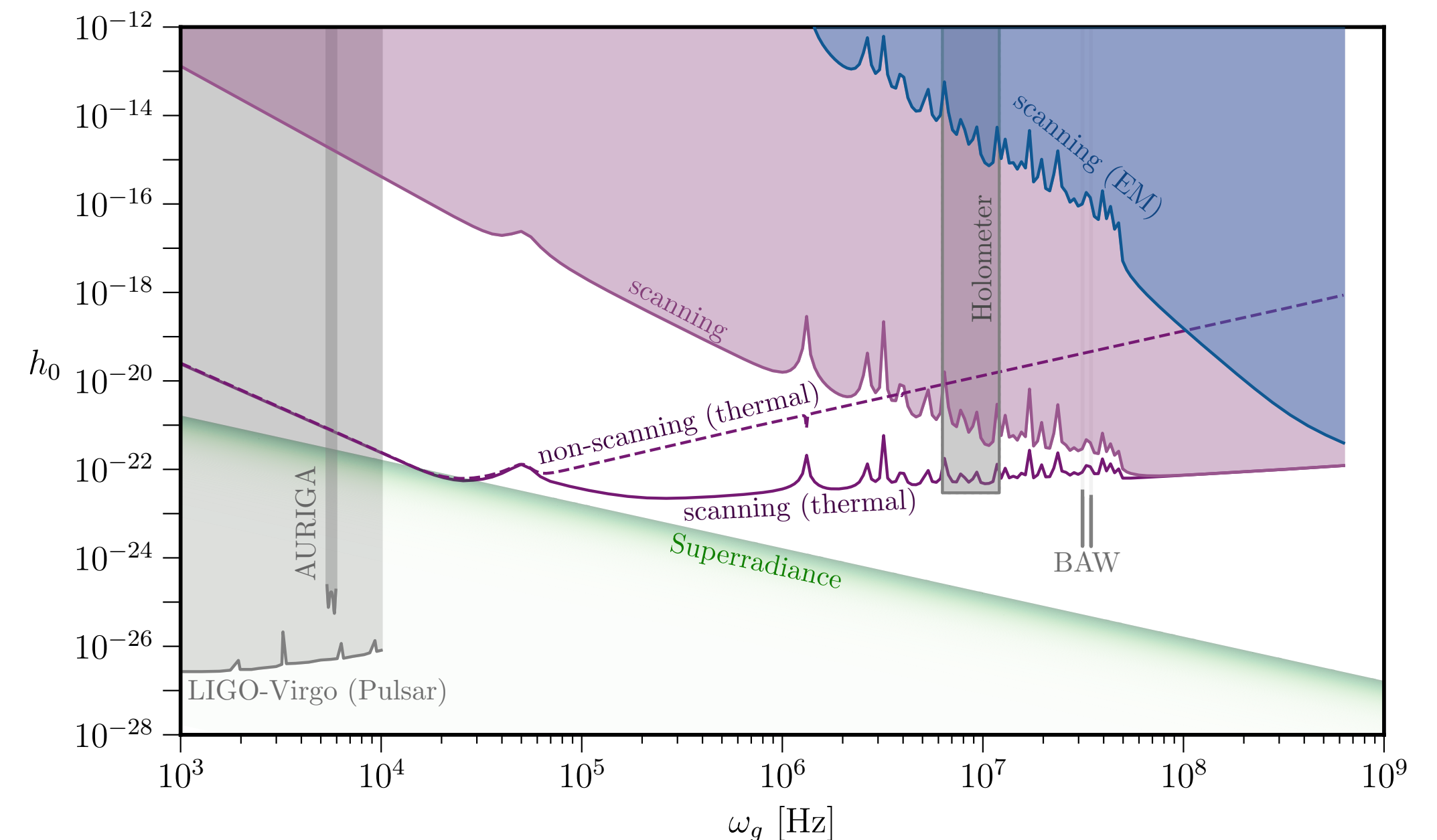
$$0 < m_a \leq m\omega_+ \quad \omega_+ = \frac{1}{4GM_{\text{BH}}} \frac{\chi_i}{1 + \sqrt{1 - \chi_i^2}}$$

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difficult at GHz, plausible at kHz