

Applied Landau damping

Lecture notes available at <https://xbuffat.web.cern.ch/landaudampingCAS.pdf>



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Collective Effects and Impedances

CERN, Switzerland, Geneva

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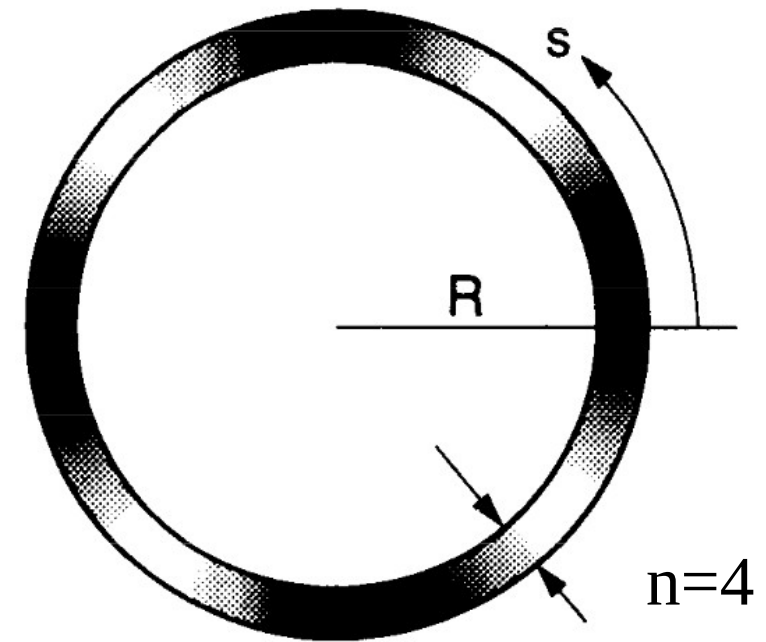
Recap

- Landau damping stems from the **interaction of single particles with waves**
 - A necessary condition for Landau damping is the a comparable velocity / frequency of the wave and the particles motion
- While collective forces such as wake fields or electron clouds tend to generate unstable modes of oscillation, Landau damping stabilises them **without emittance growth**
 - An external perturbation may also decay through a similar phenomenon, we rather talk about decoherence or filamentation. This mechanism leads to **emittance growth**
- Landau damping originates in the spread of oscillation frequencies of the particles in the beam
 - It is a **linear mechanism**, as in plasmas. However in accelerators the frequency spread often originates from **non-linear forces**

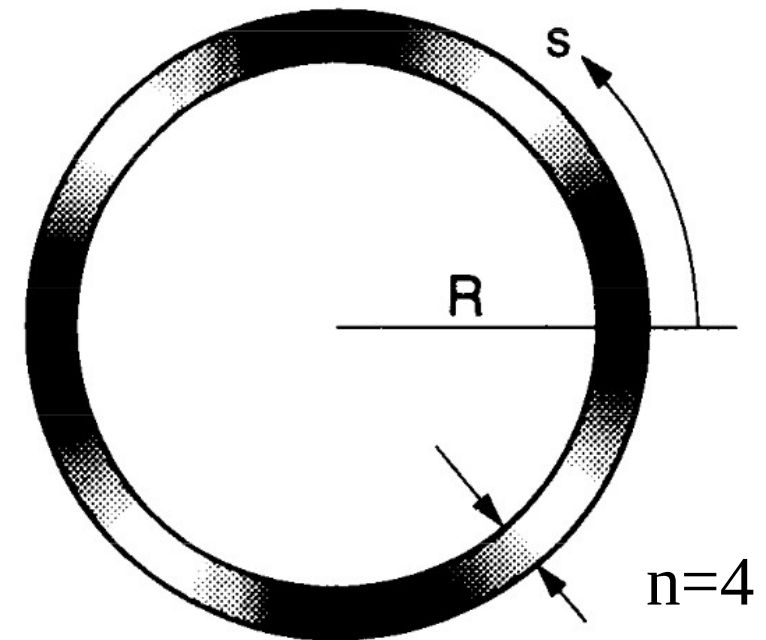
Content

- Part I (concept)
 - Wave – particle interaction
 - Van Kampen approach
 - Stability diagram and beam transfer function
- Part II (applications)
 - Longitudinal and transverse Landau damping in unbunched and bunched beams
 - Non-linear collective forces
 - Advanced Landau damping techniques

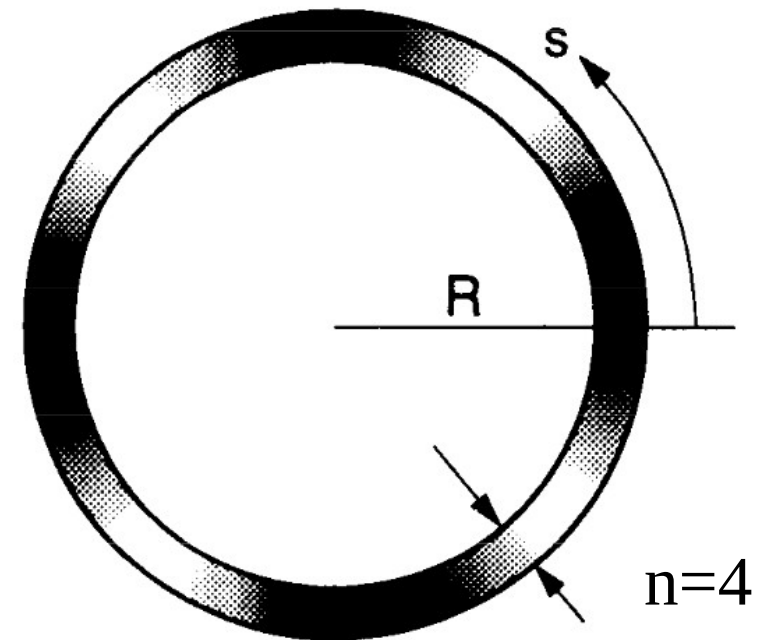
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 - The dispersion relation takes a **special form**:

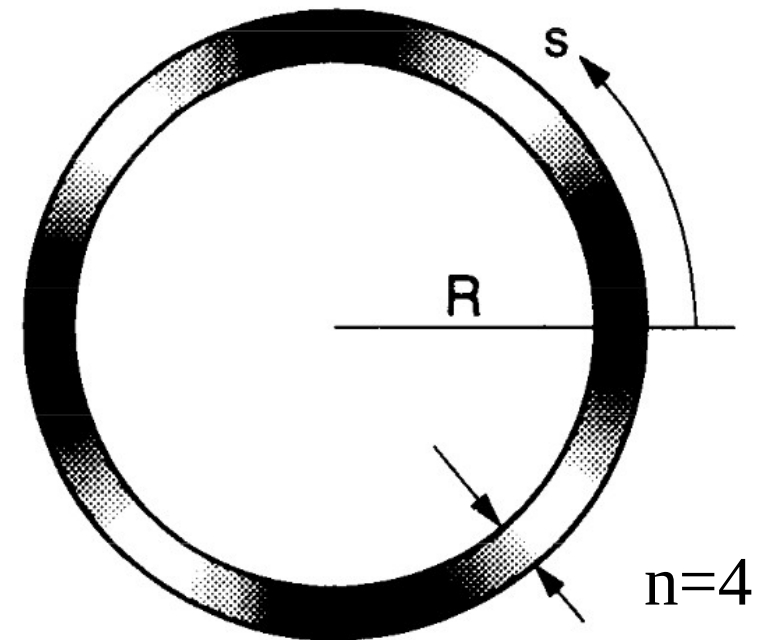


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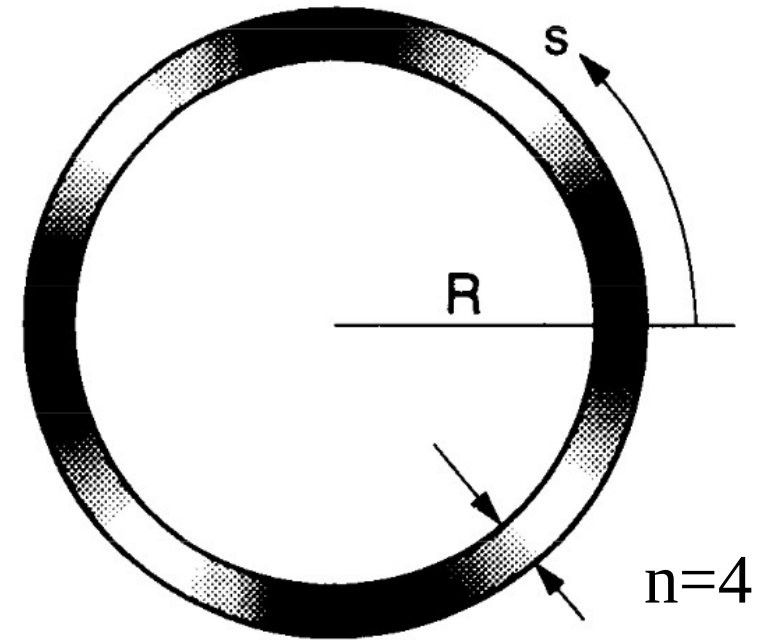
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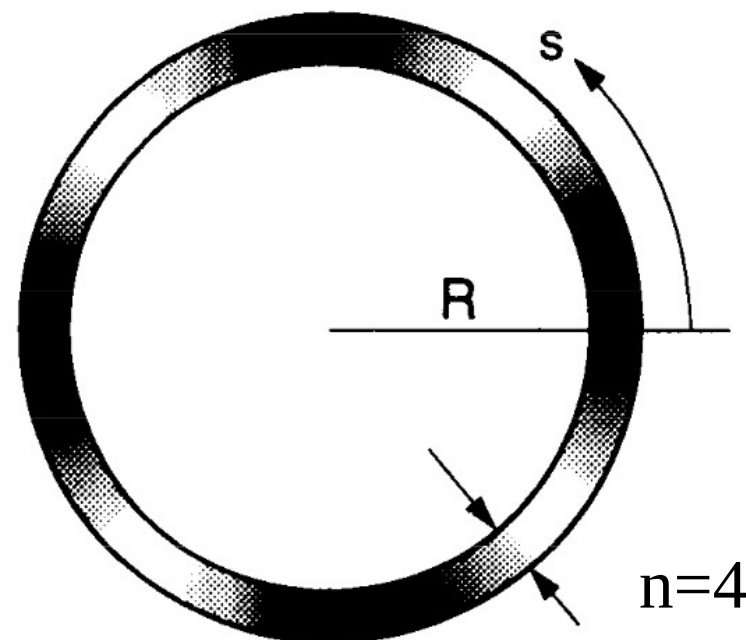


Mode frequency shift driven by wake fields (without Landau damping)

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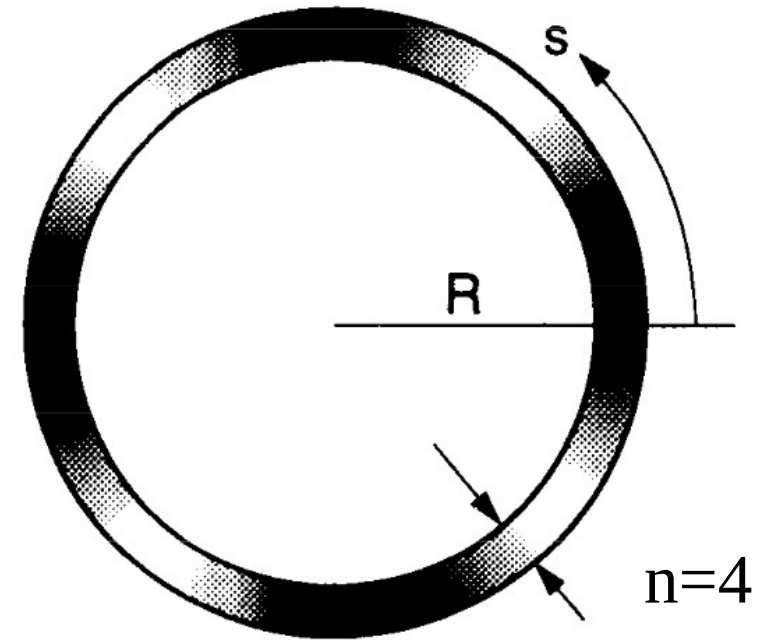
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Mode number \rightarrow (points to n in the denominator)

Distribution of revolution frequencies \rightarrow (points to $\rho(\omega)$)

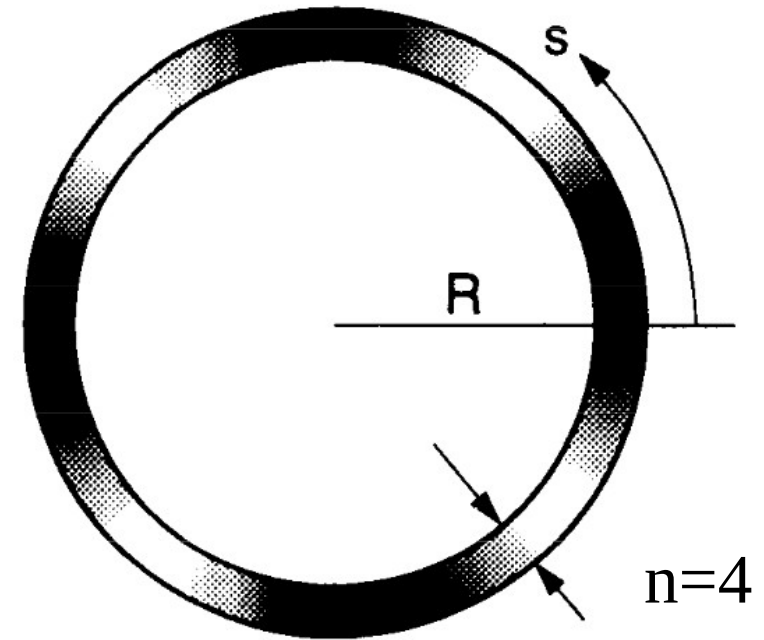
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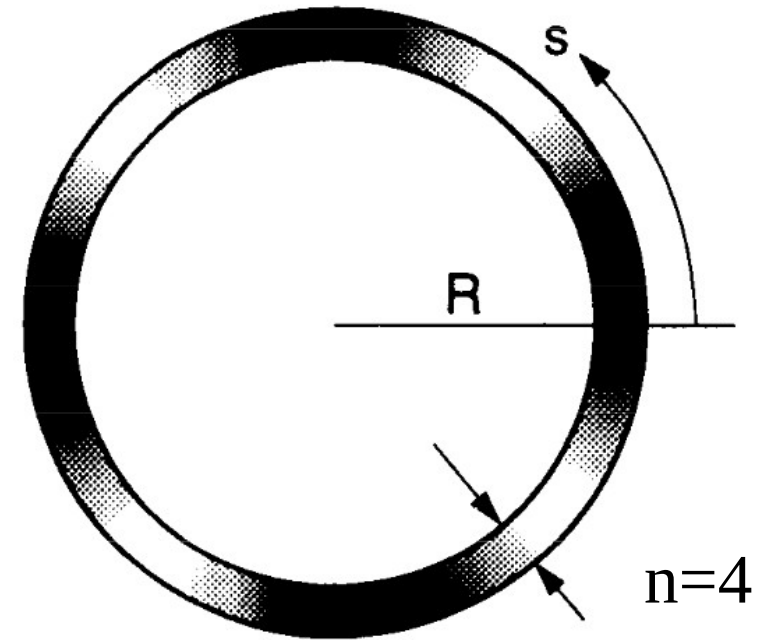
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Linked to the absence of focusing $\rightarrow (n\omega - \Omega)^2$

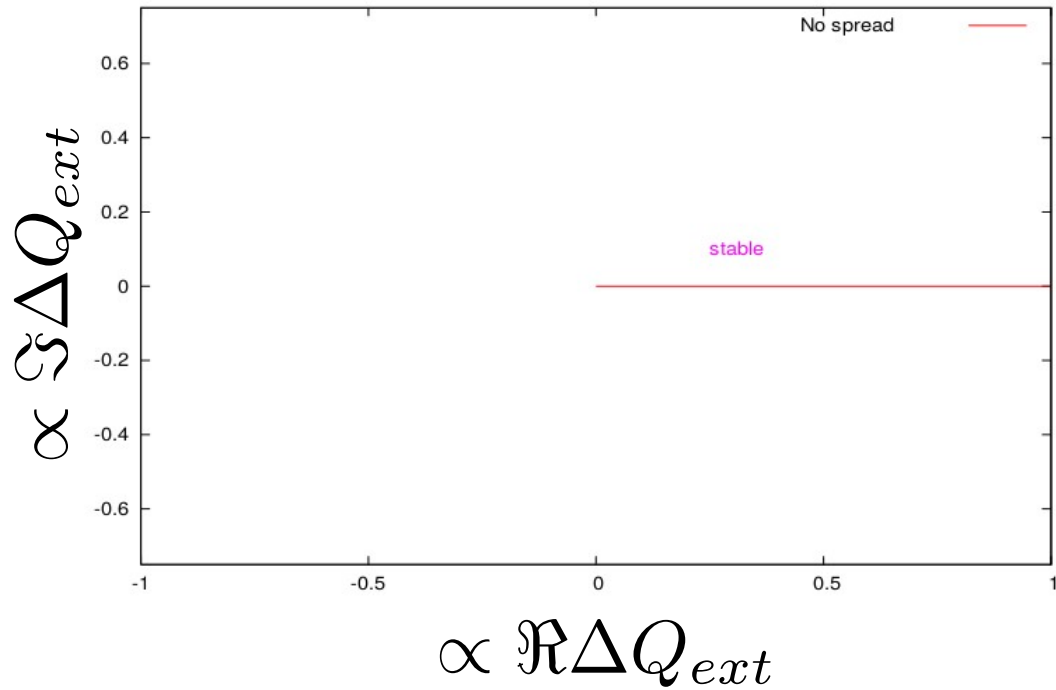
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Stability diagram – Keil-Schnell criterion

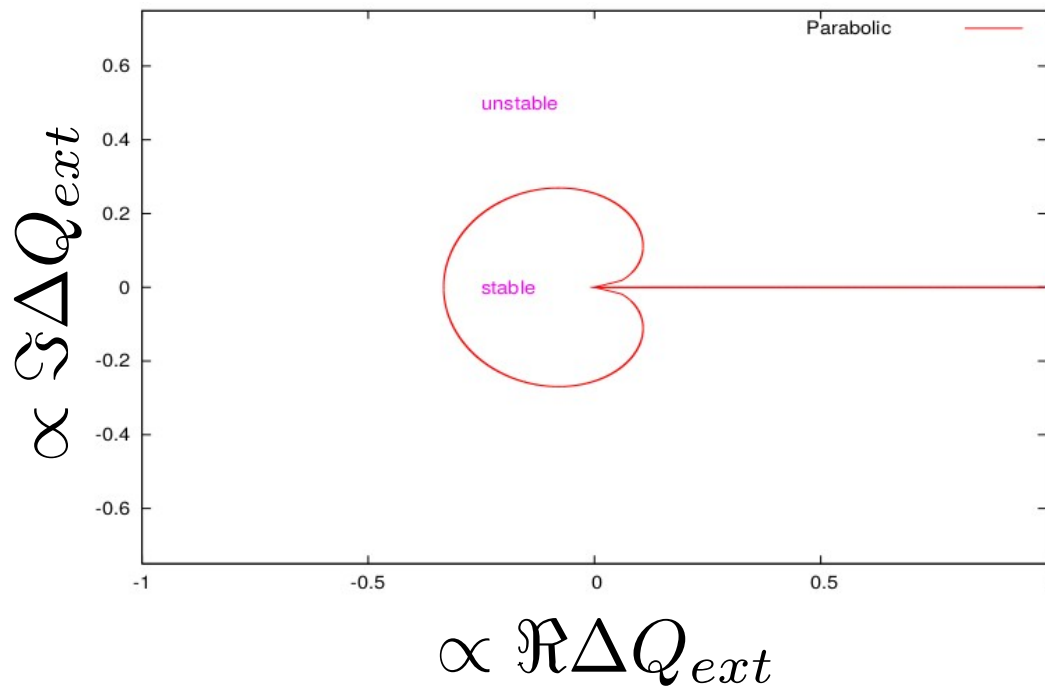
[Chao,
Herr]



- The beam is always unstable without energy spread

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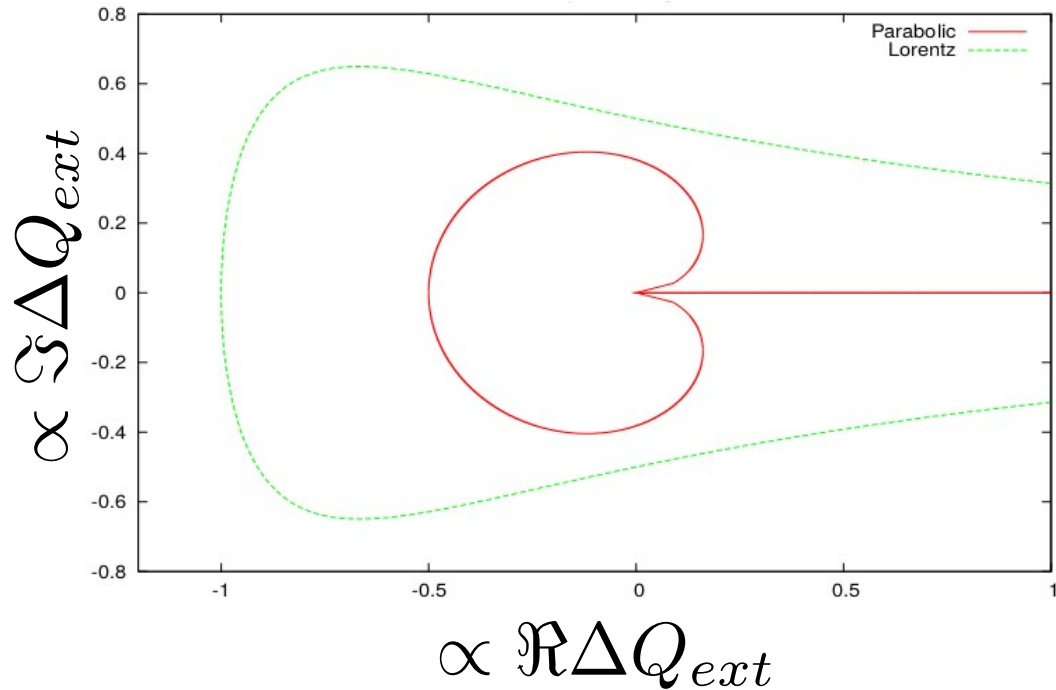
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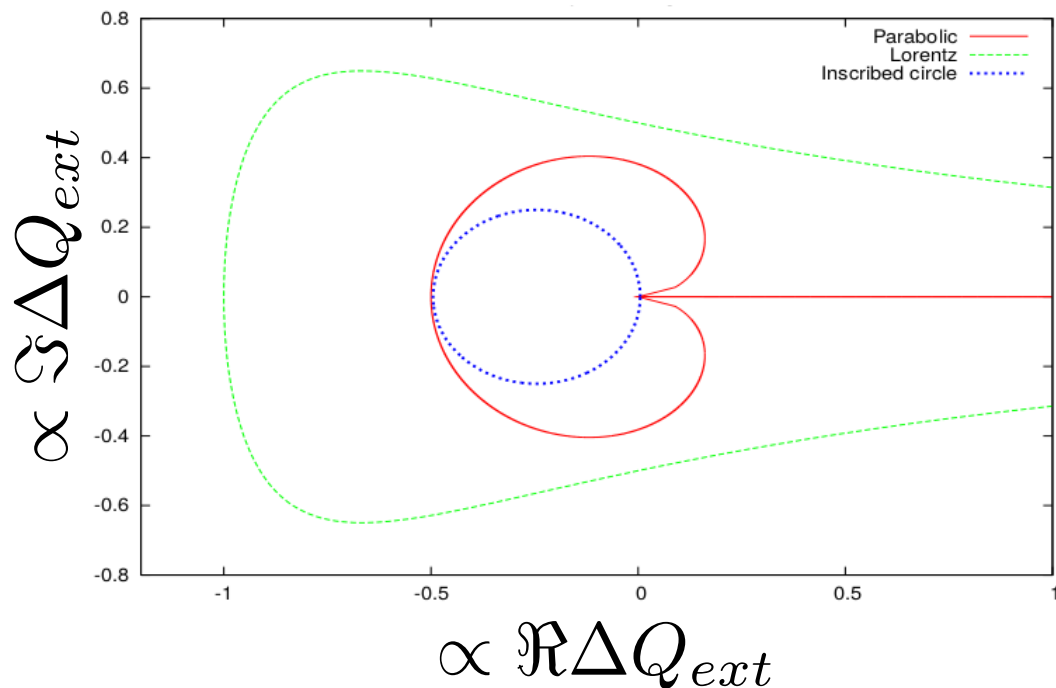
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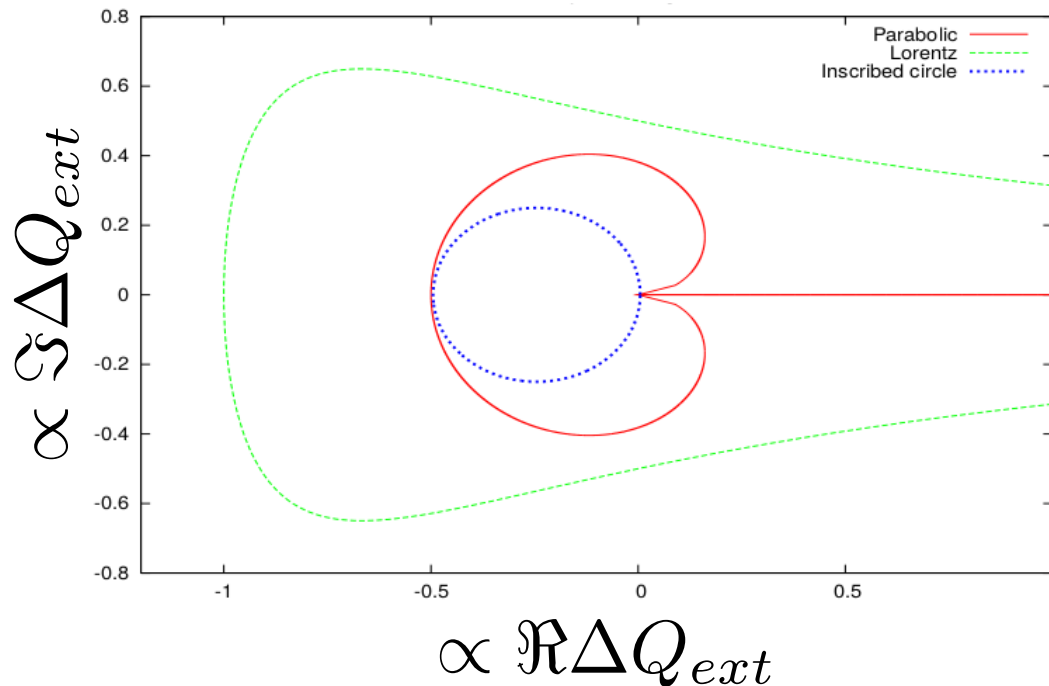
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- Usually the frequency distribution is poorly known, Keil-Schnell derived a conservative criterion based on the inscribed circle:

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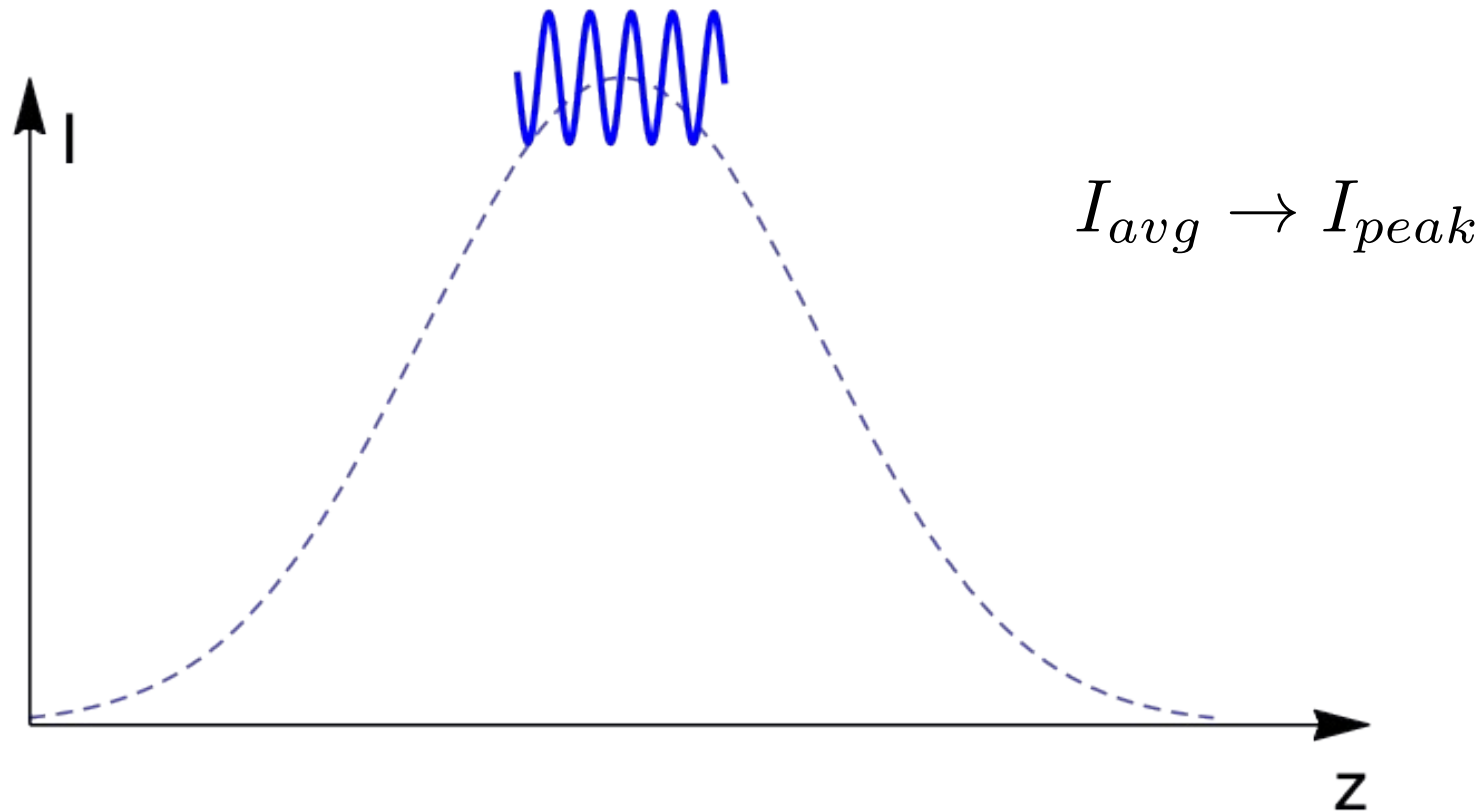
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- Revolution frequency spread:

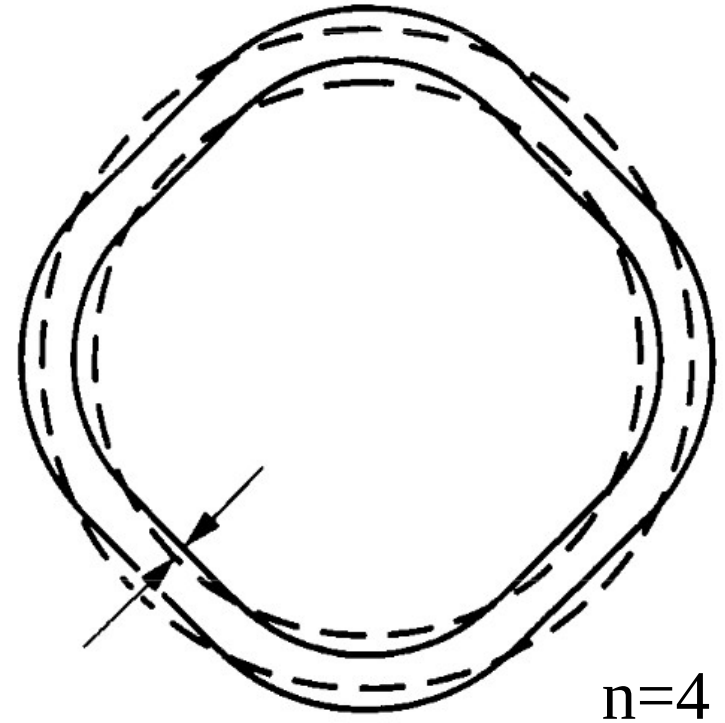
$$\Delta \omega \approx \omega_0 |\eta| \Delta \delta$$

Momentum spread

- The KS criterion also provides a good indication of the requirement to stabilise the **microwave instability** in bunched beams
 - Keil-Schnell-Boussard criterion

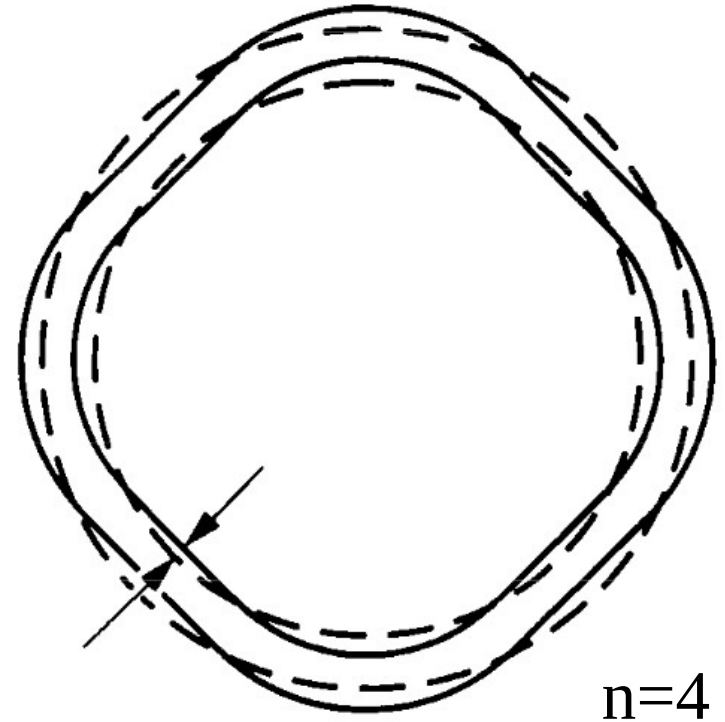


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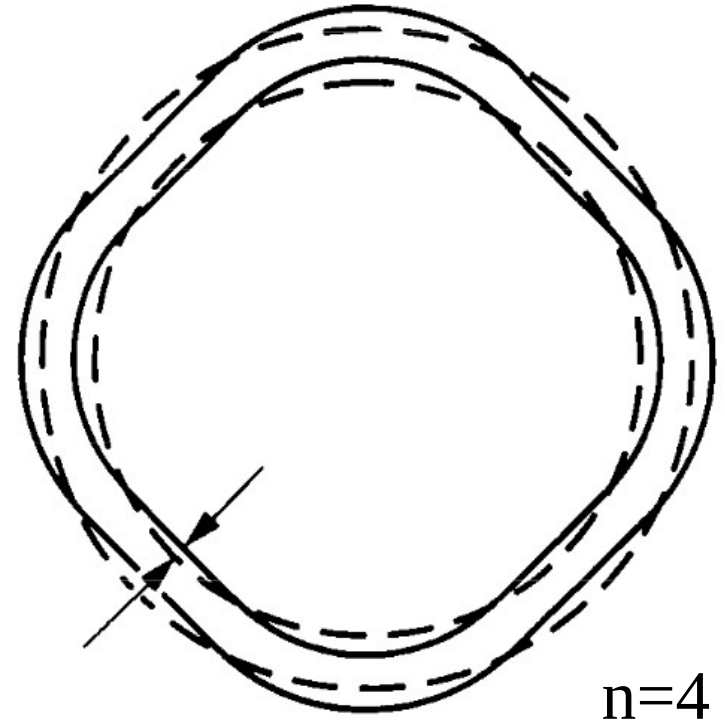
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Transverse frequency shift caused by the impedance, e.g. from perturbation theory:

$$\Delta\Omega_n = -i \frac{Nr_0 c^2 \eta}{2\gamma \omega_\beta T_0} Z^\perp (n\omega_0 + \omega_\beta)$$



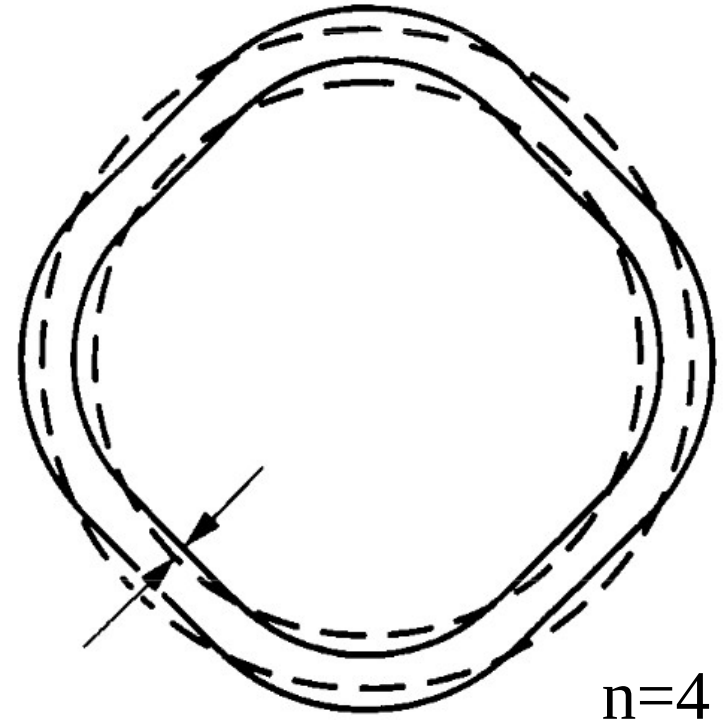
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Transverse frequency spread \rightarrow $\rho(\omega)$
 $\Delta\Omega_n$ \rightarrow

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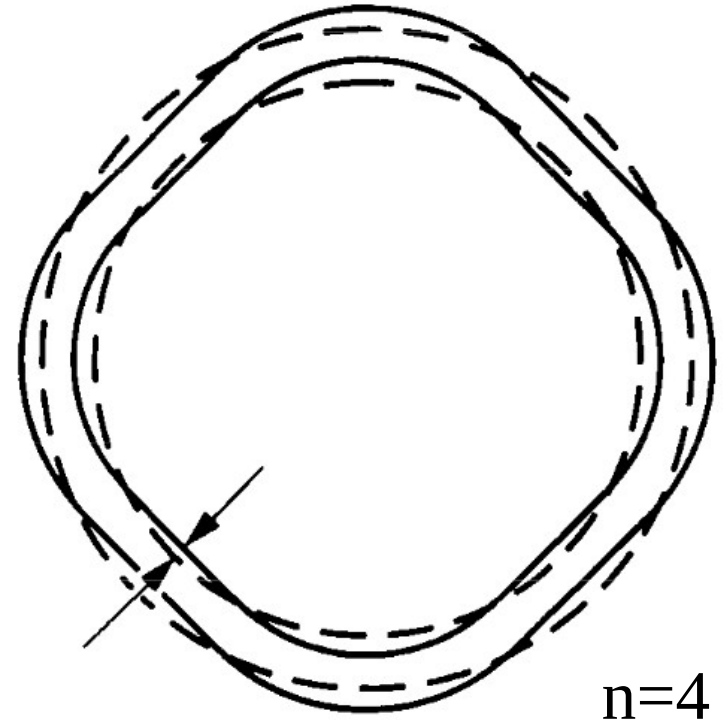
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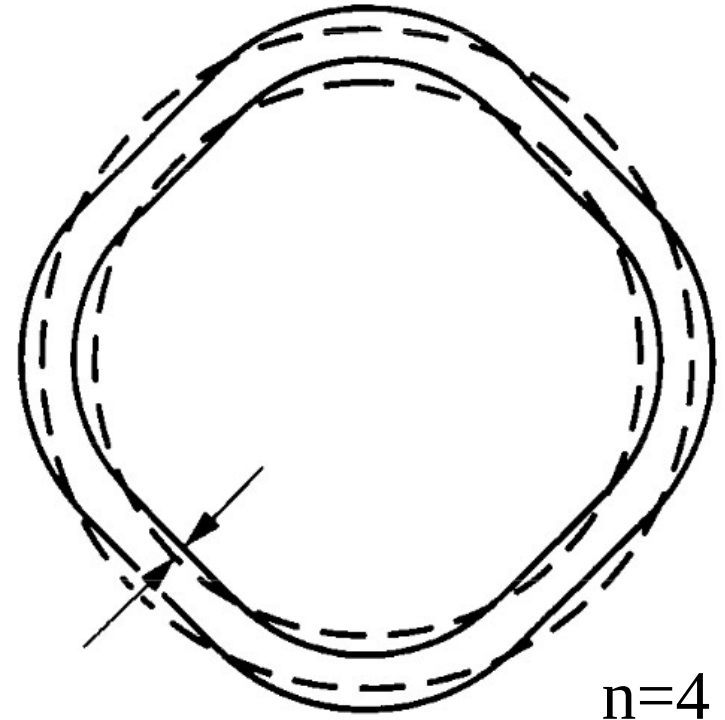
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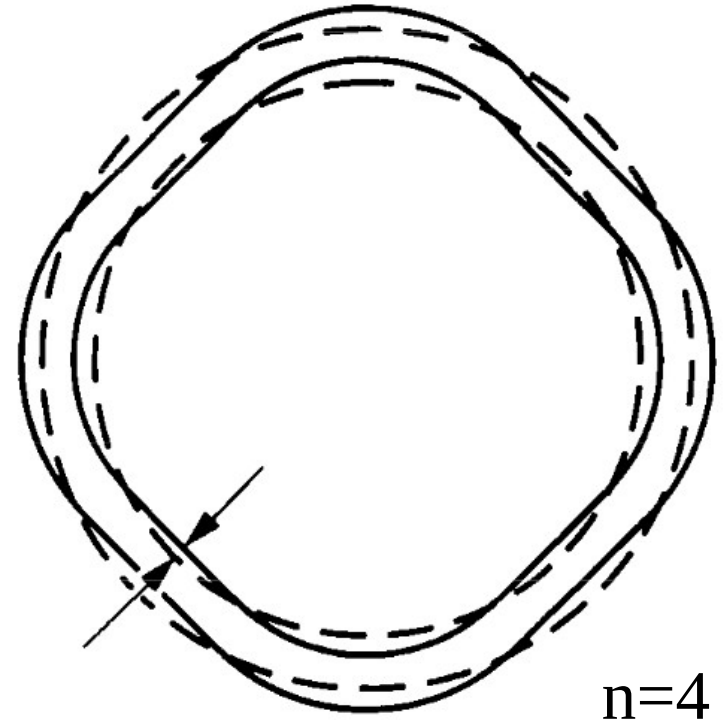
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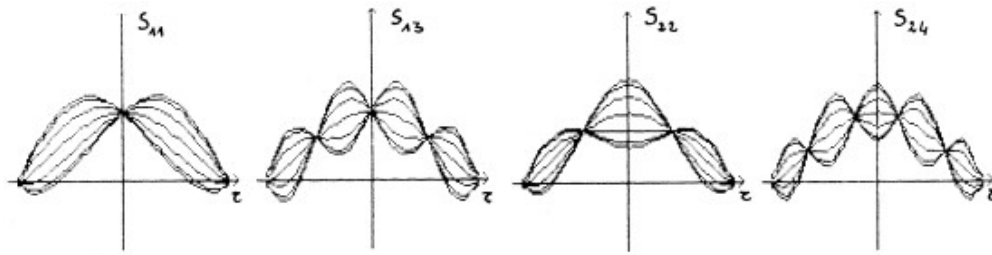
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 - Chromaticity (Q')

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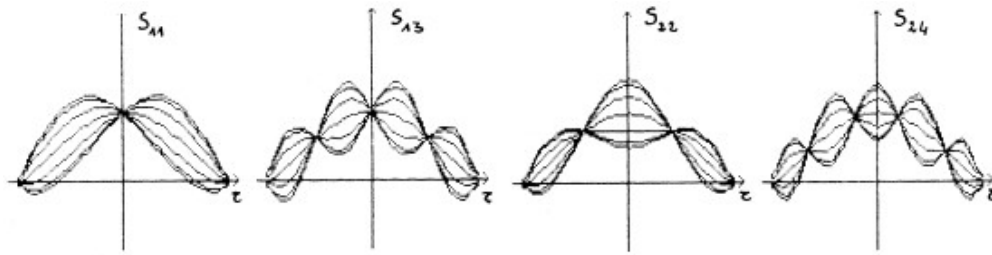
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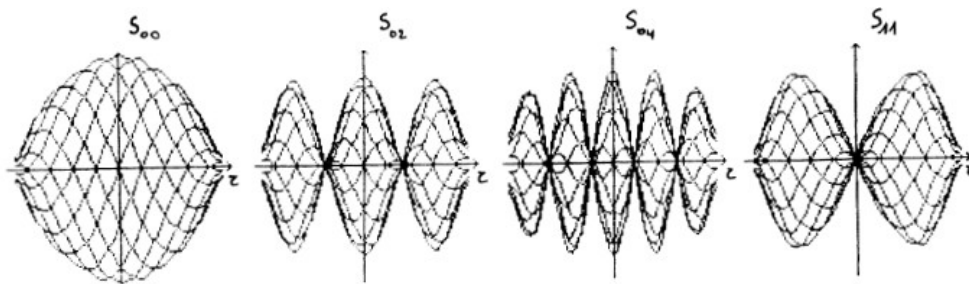


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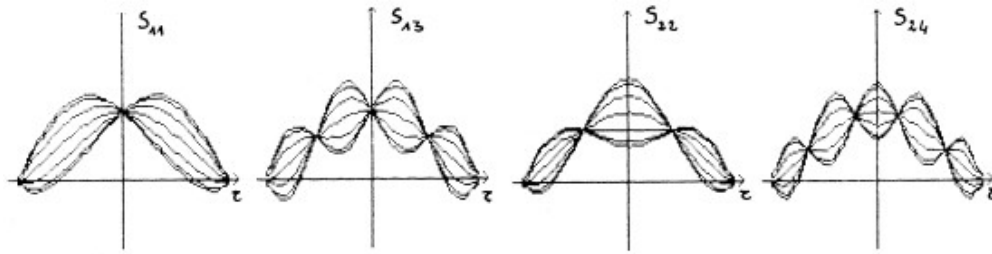
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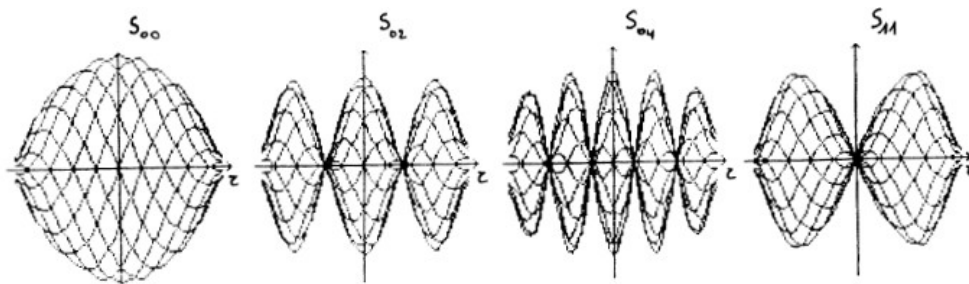


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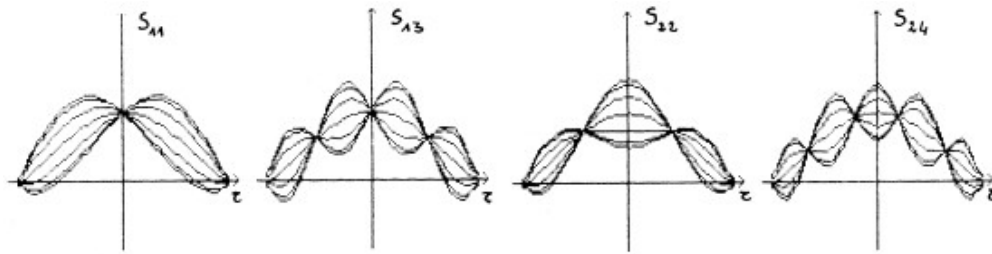


Mode frequency shift
(e.g. from Sacharer
formula)

$$\rightarrow |\Delta\Omega_n| \approx \Delta\omega$$



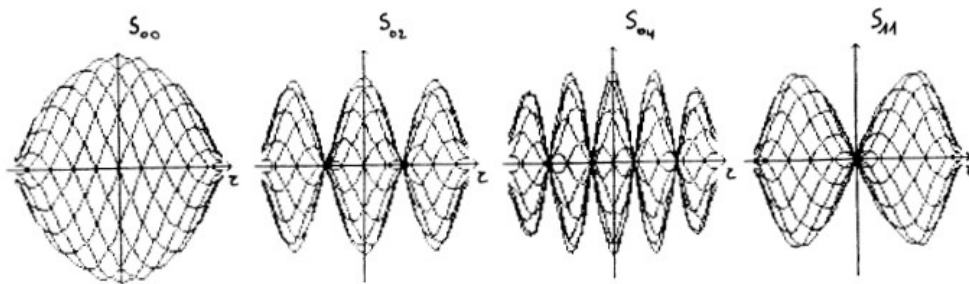
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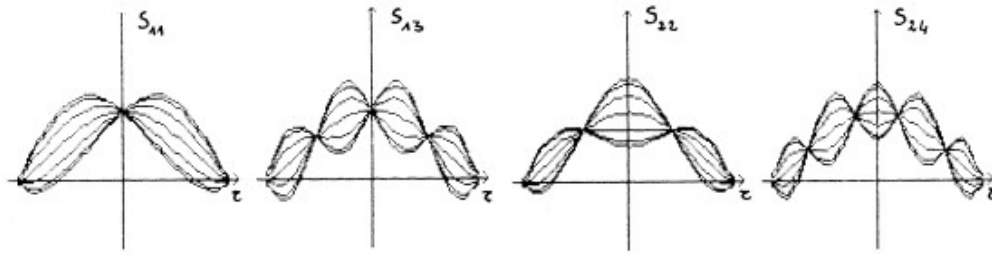
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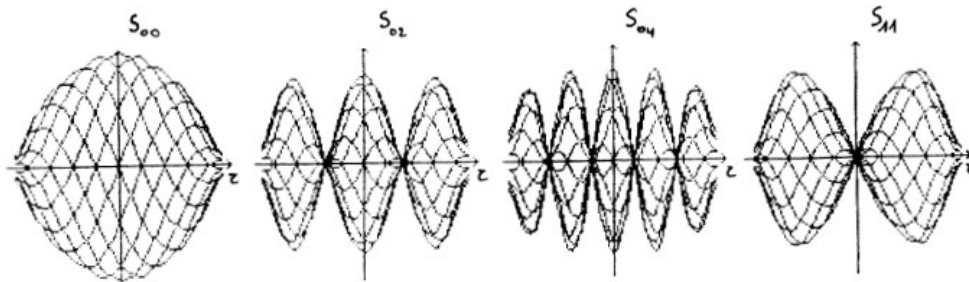
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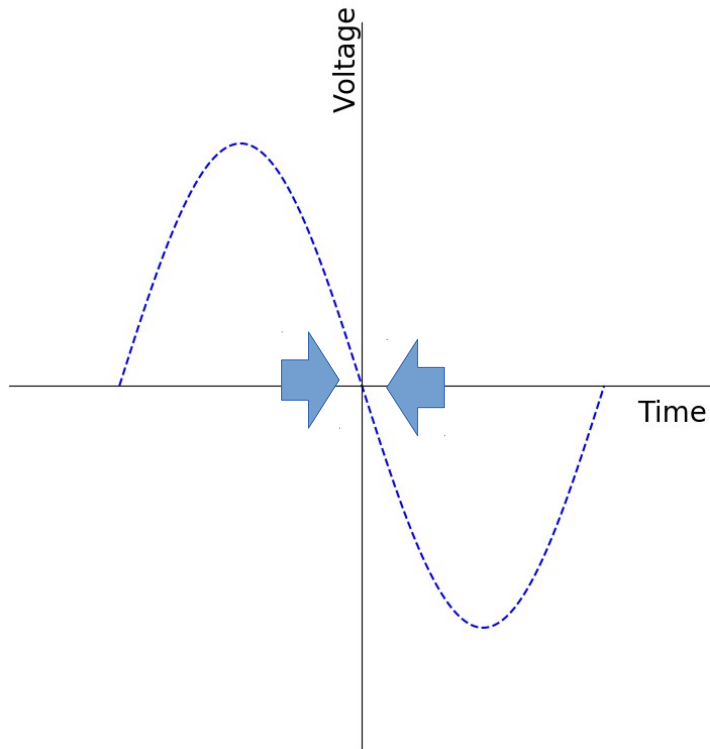
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 - Chromaticity
 - Non-linear forces

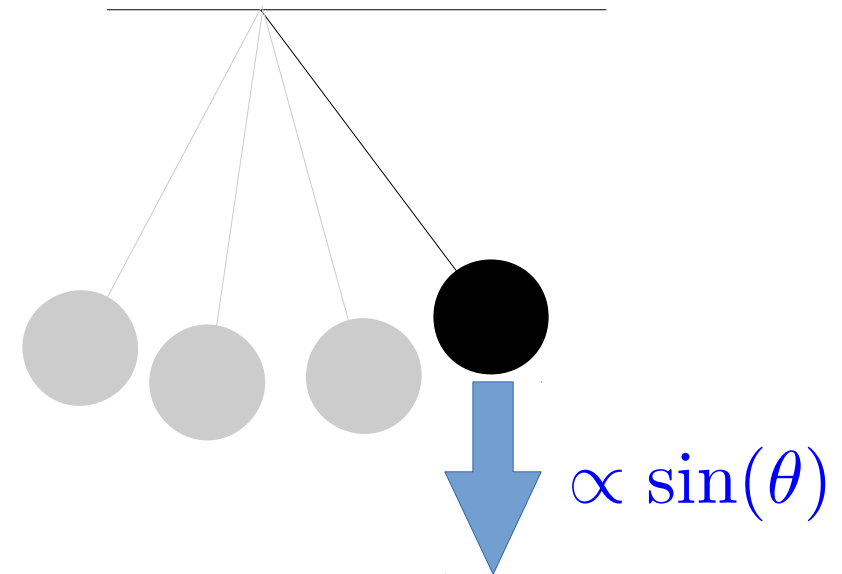
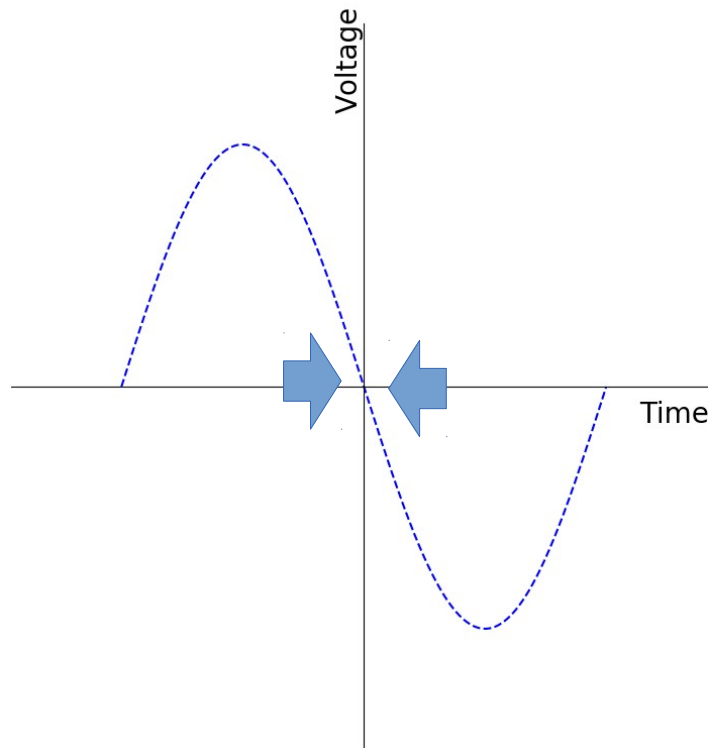
Longitudinal stability of bunched beams

- For bunched beams, the longitudinal focusing provokes oscillations around the fixed point with ω_s
 - As RF cavities function with sine wave, the focusing force is non-linear

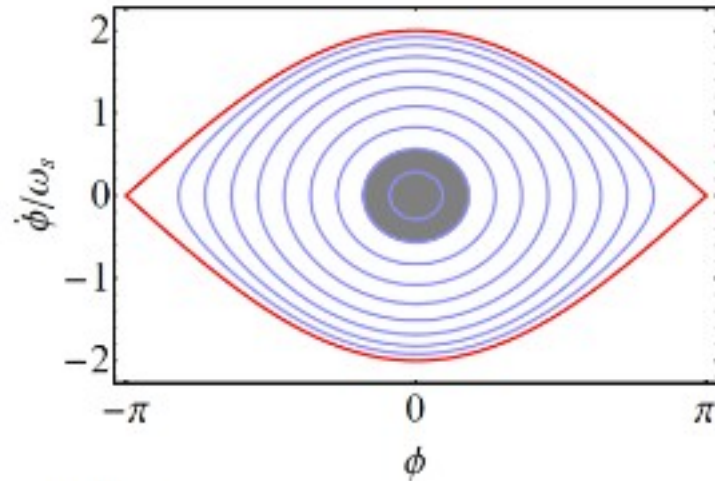


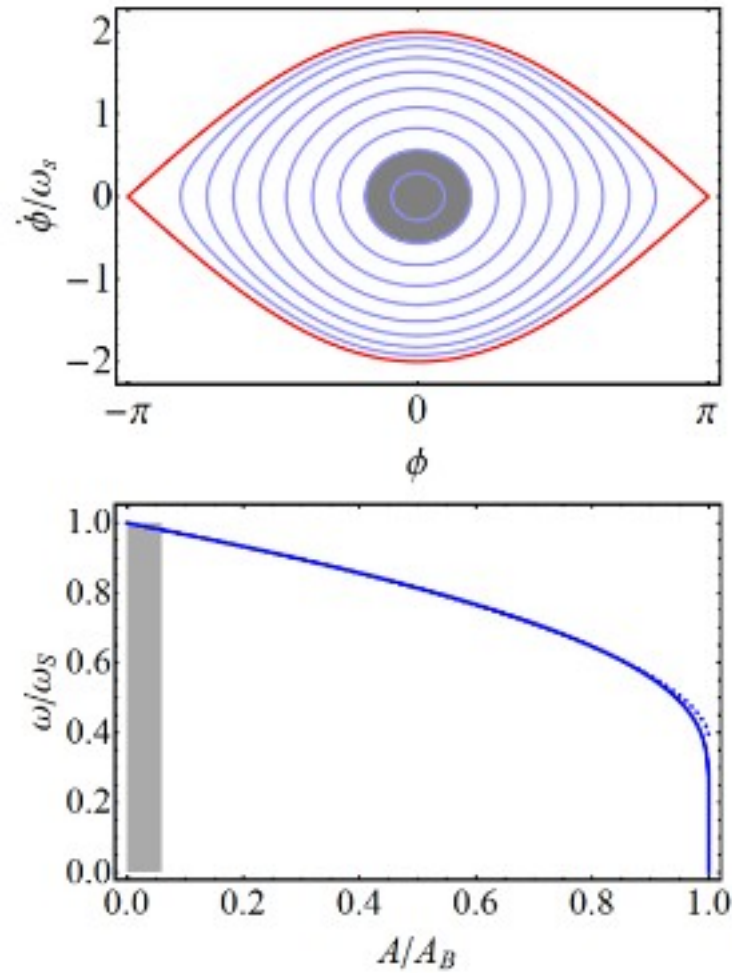
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- The behaviour is identical to the pendulum **without** the small angle approximations

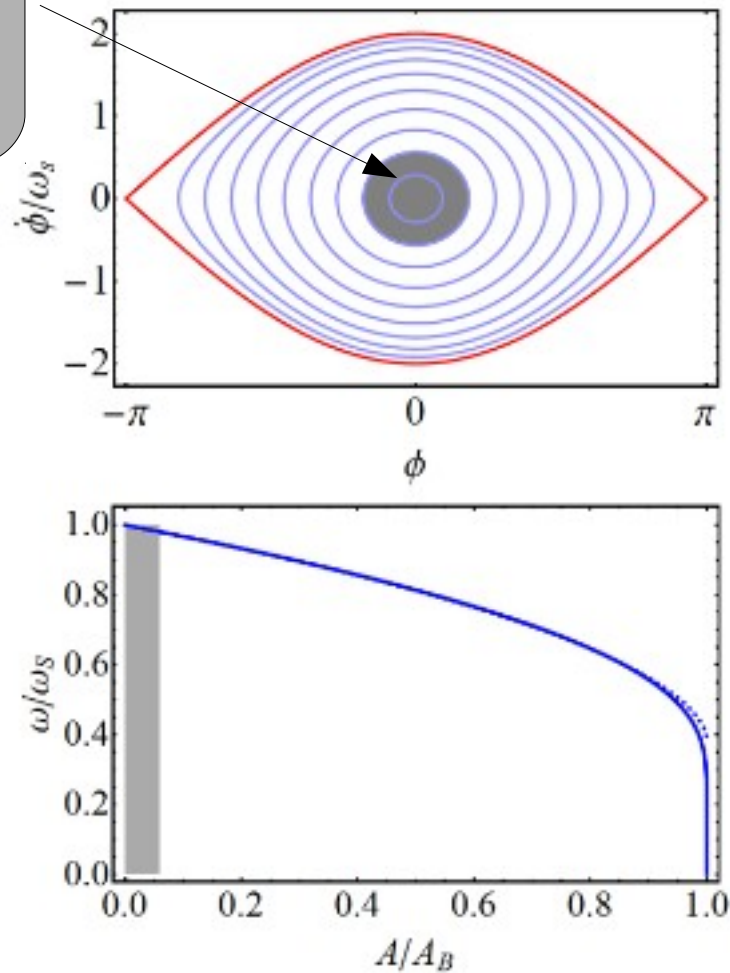




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[Damerau]

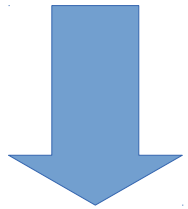
Small beam
wrt to
available
bucket



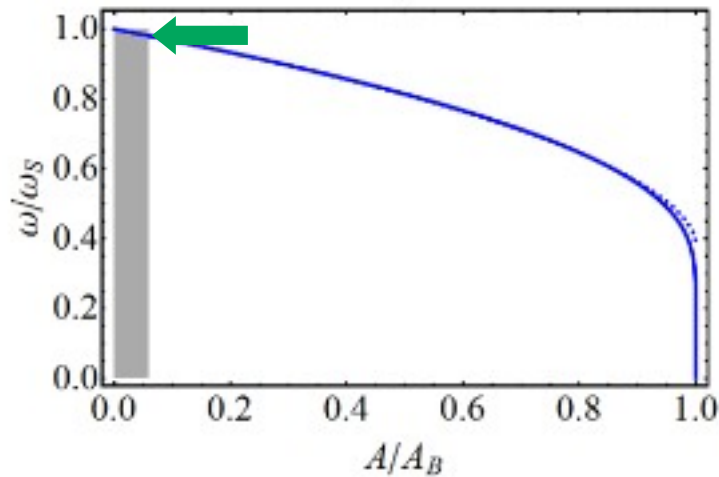
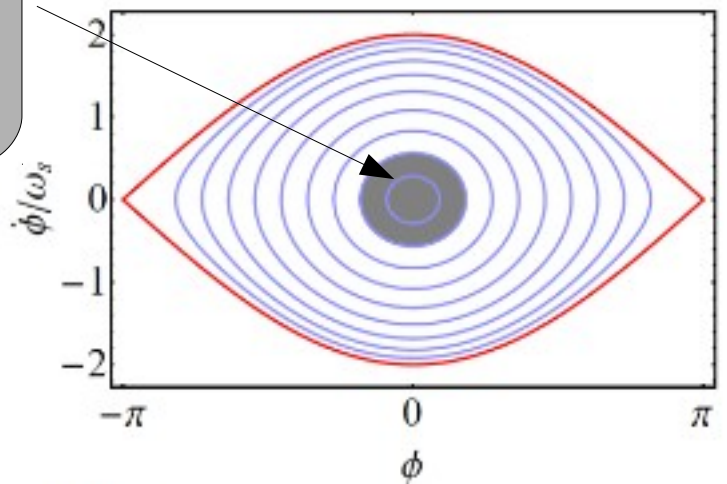
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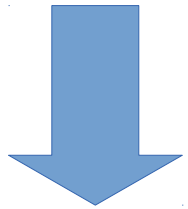
Small
frequency
spread



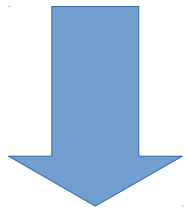
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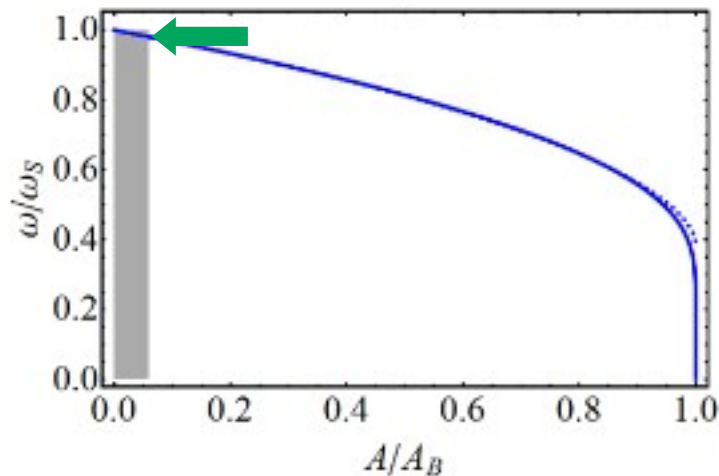
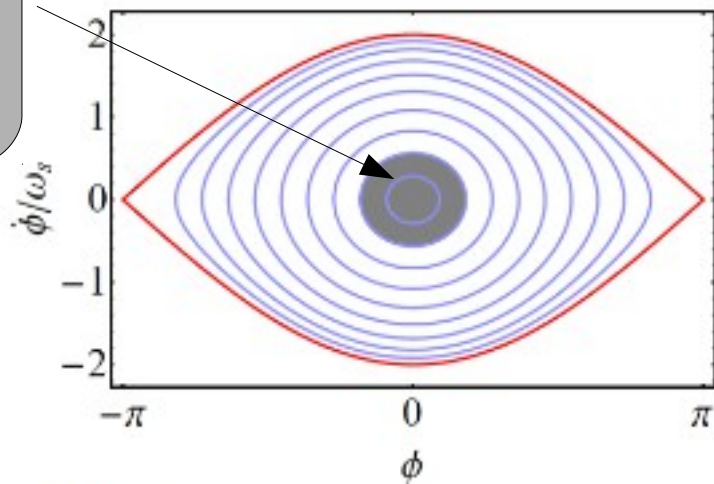
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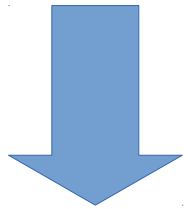
Weak
Landau
damping



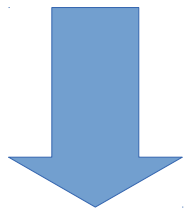
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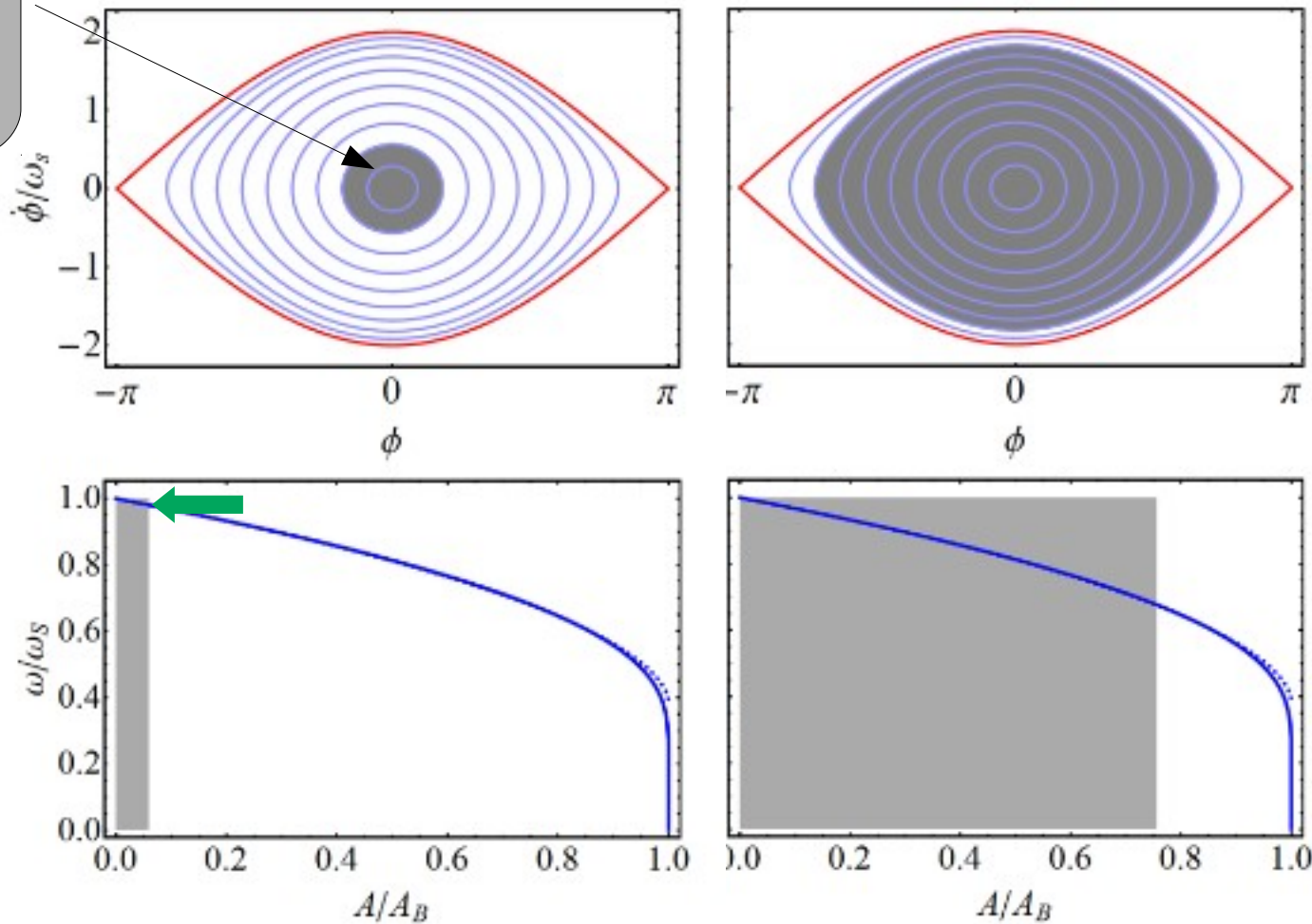
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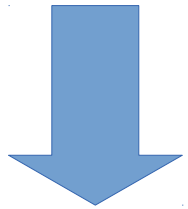
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Longitudinal stability of bunched beams

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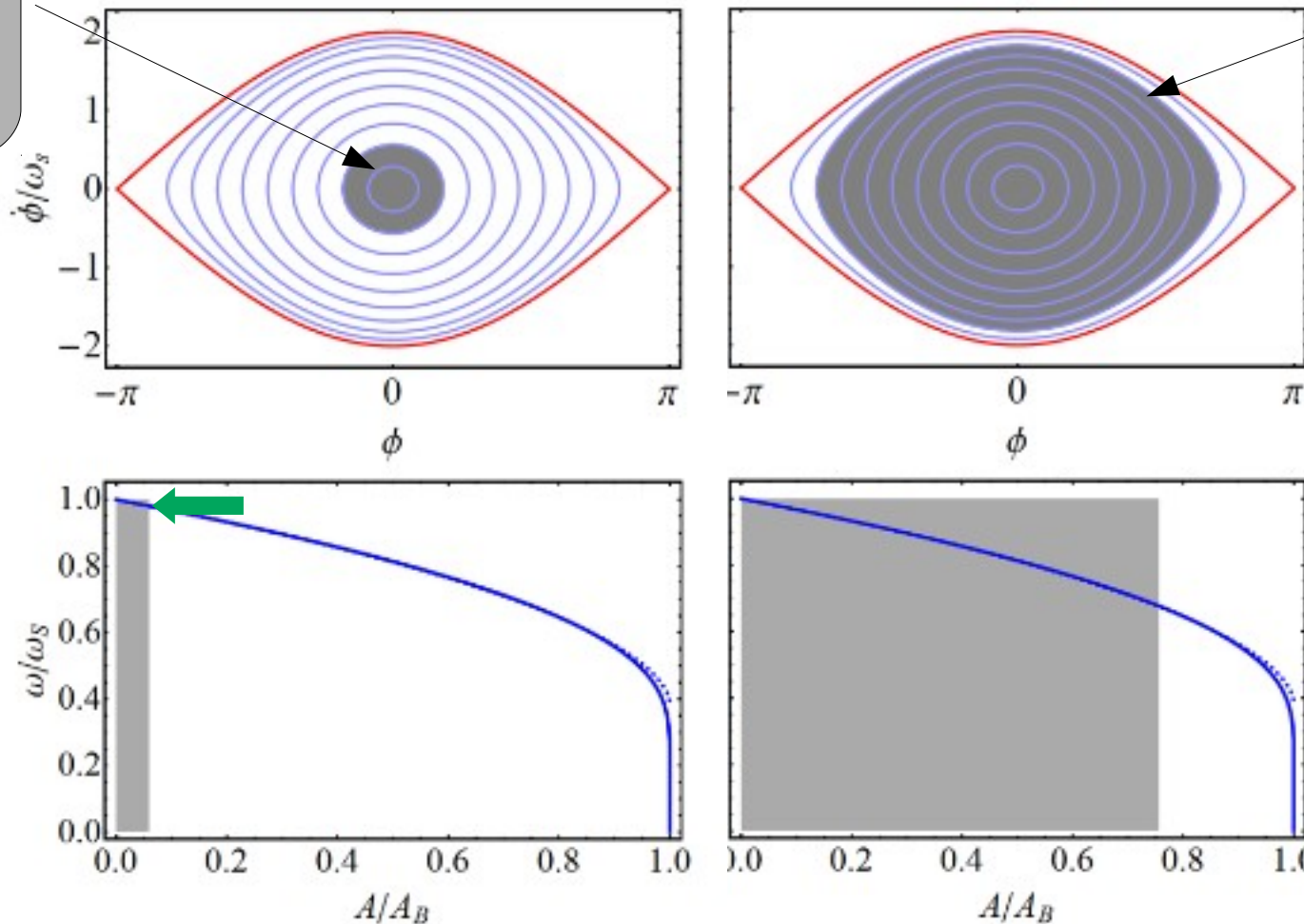
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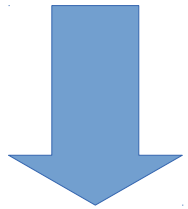


Bucket is
filled

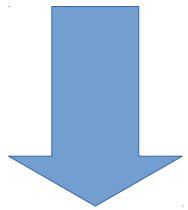
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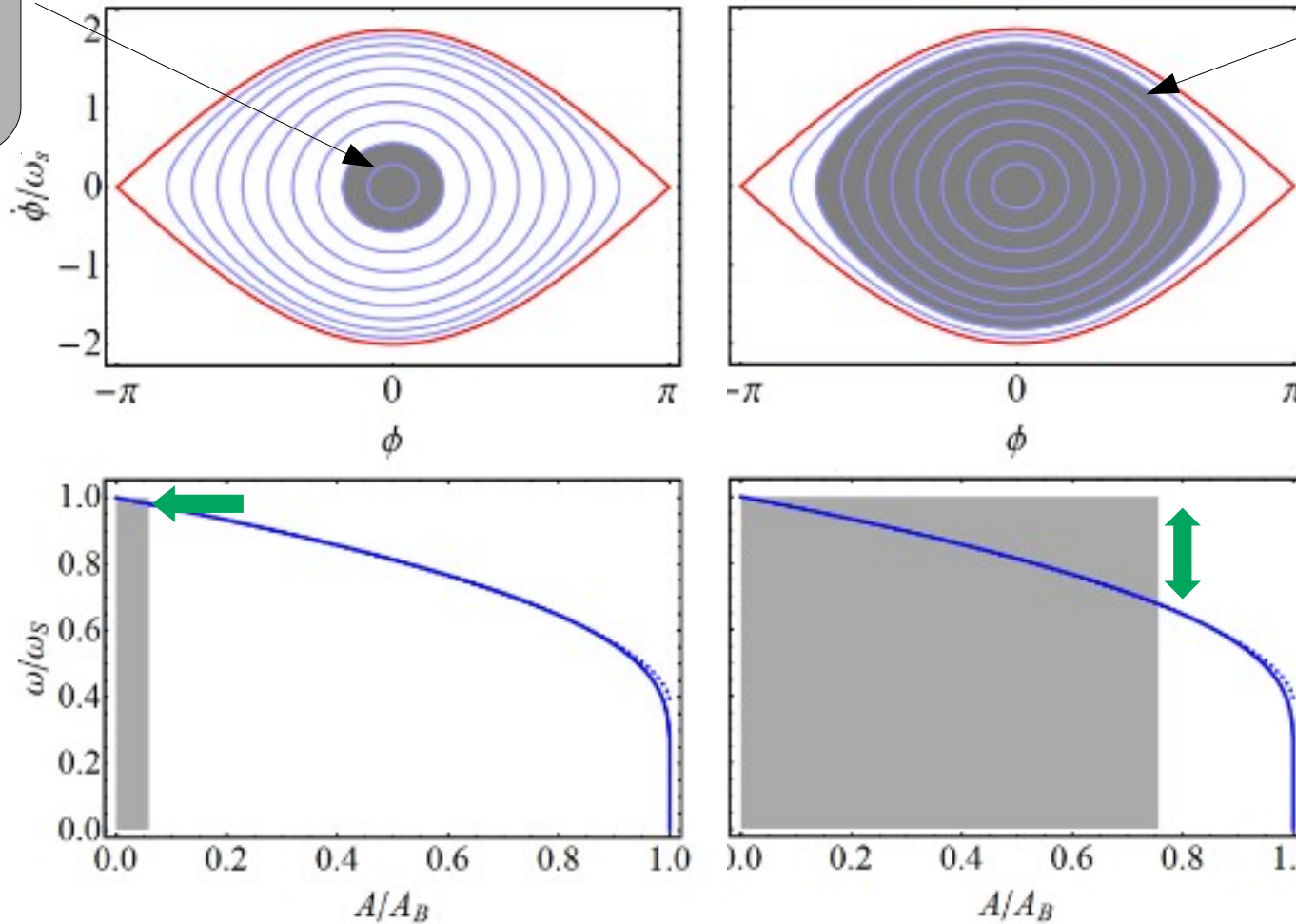
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available
bucket



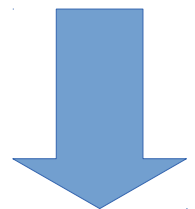
Small
frequency
spread



Weak
Landau
damping



Bucket is
filled

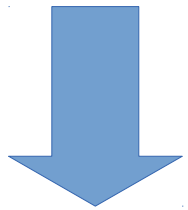


large
frequency
spread

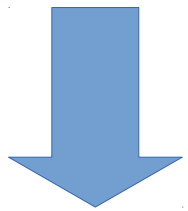
Longitudinal stability of bunched beams

[Damerou]

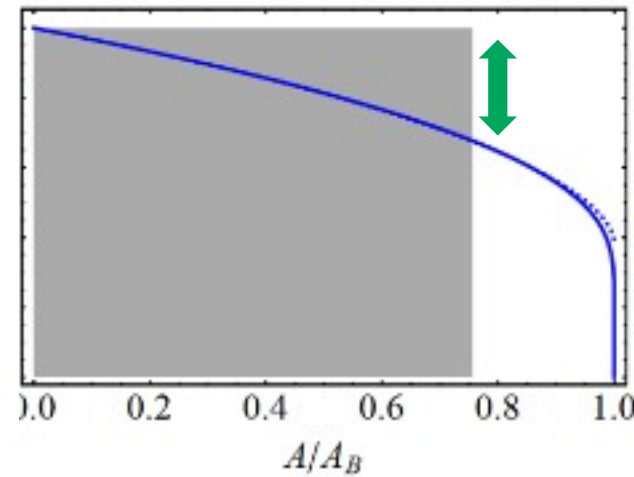
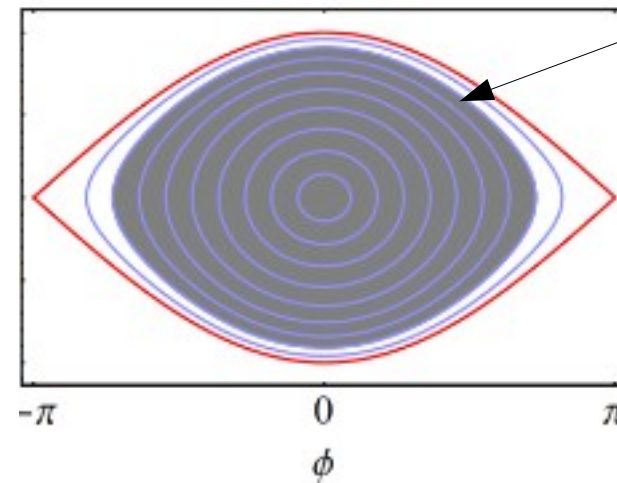
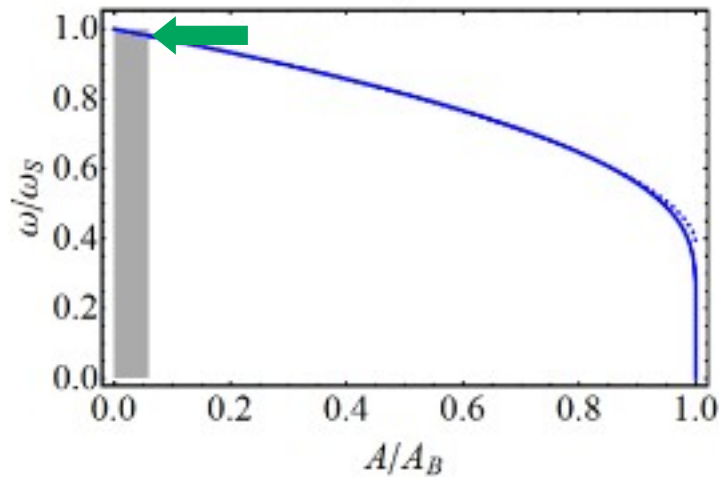
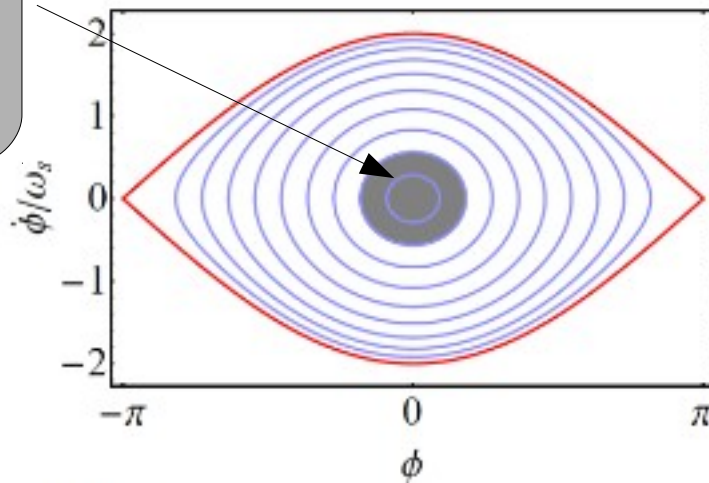
Small beam
wrt to
available
bucket



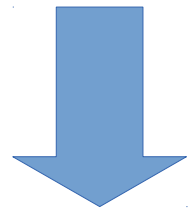
Small
frequency
spread



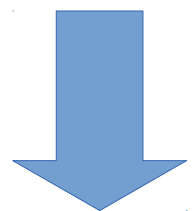
Weak
Landau
damping



Bucket is
filled



large
frequency
spread

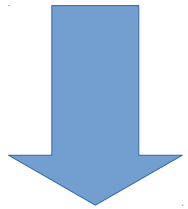


Strong
Landau
damping

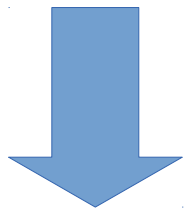
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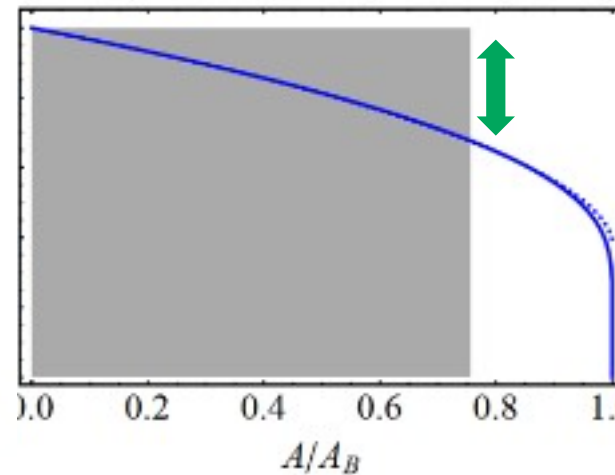
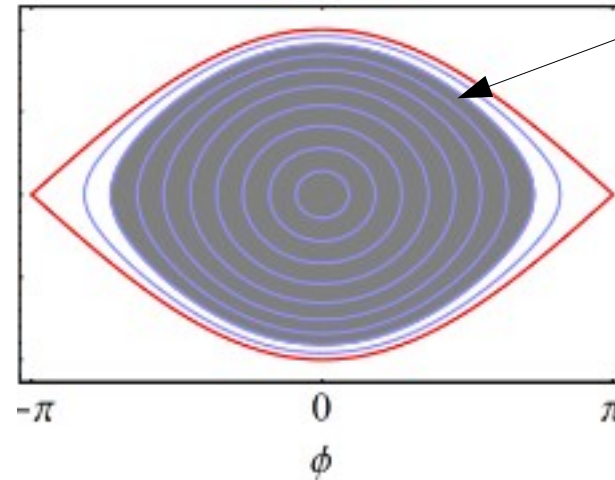
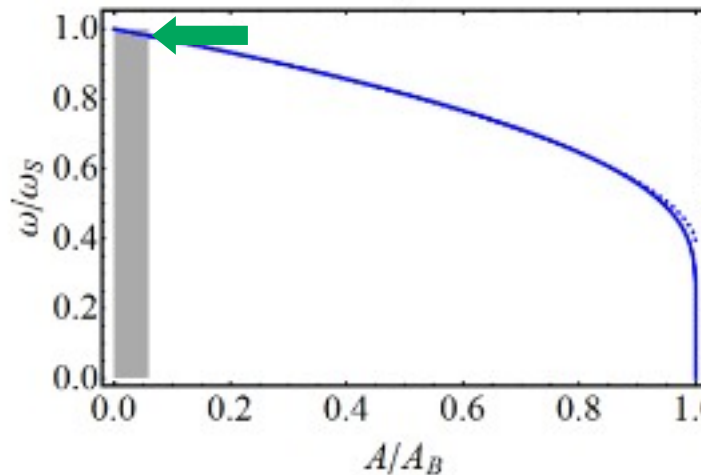
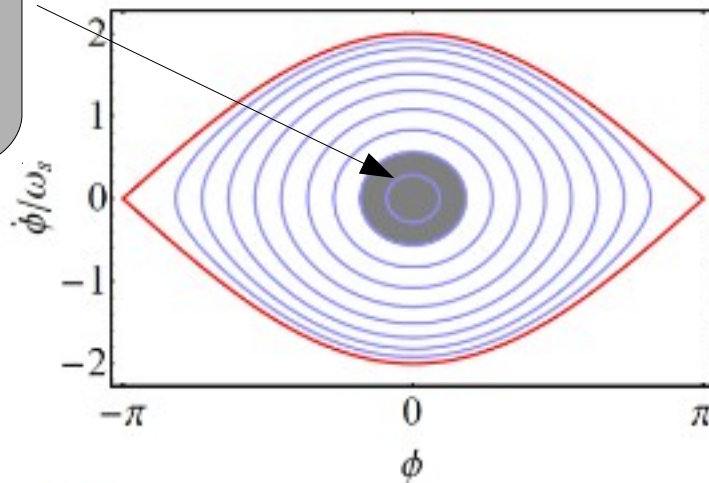
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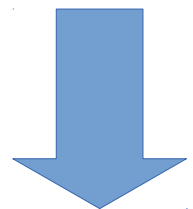
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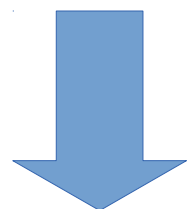
Weak
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large
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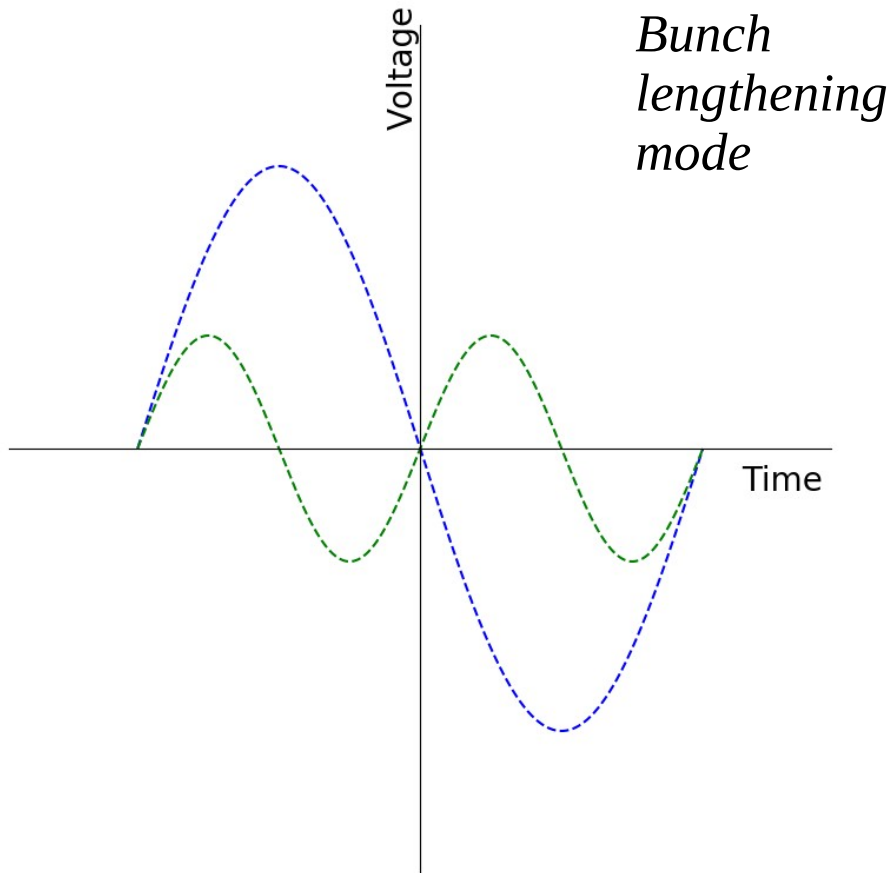


Strong
Landau
damping

- **Filling the available bucket** is key to maintain Landau damping in the longitudinal plane

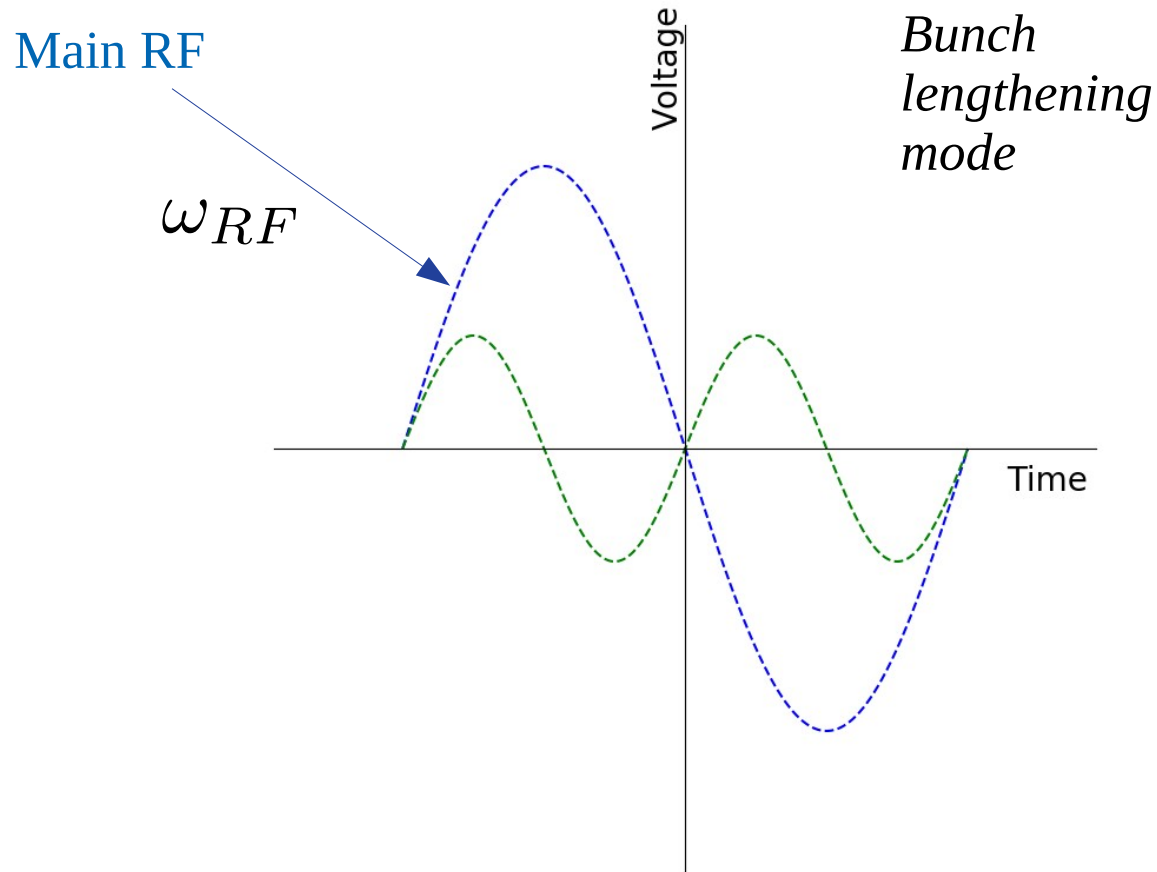
Double harmonic RF

- With a second harmonic RF (featuring a lower voltage) the total voltage becomes more non-linear



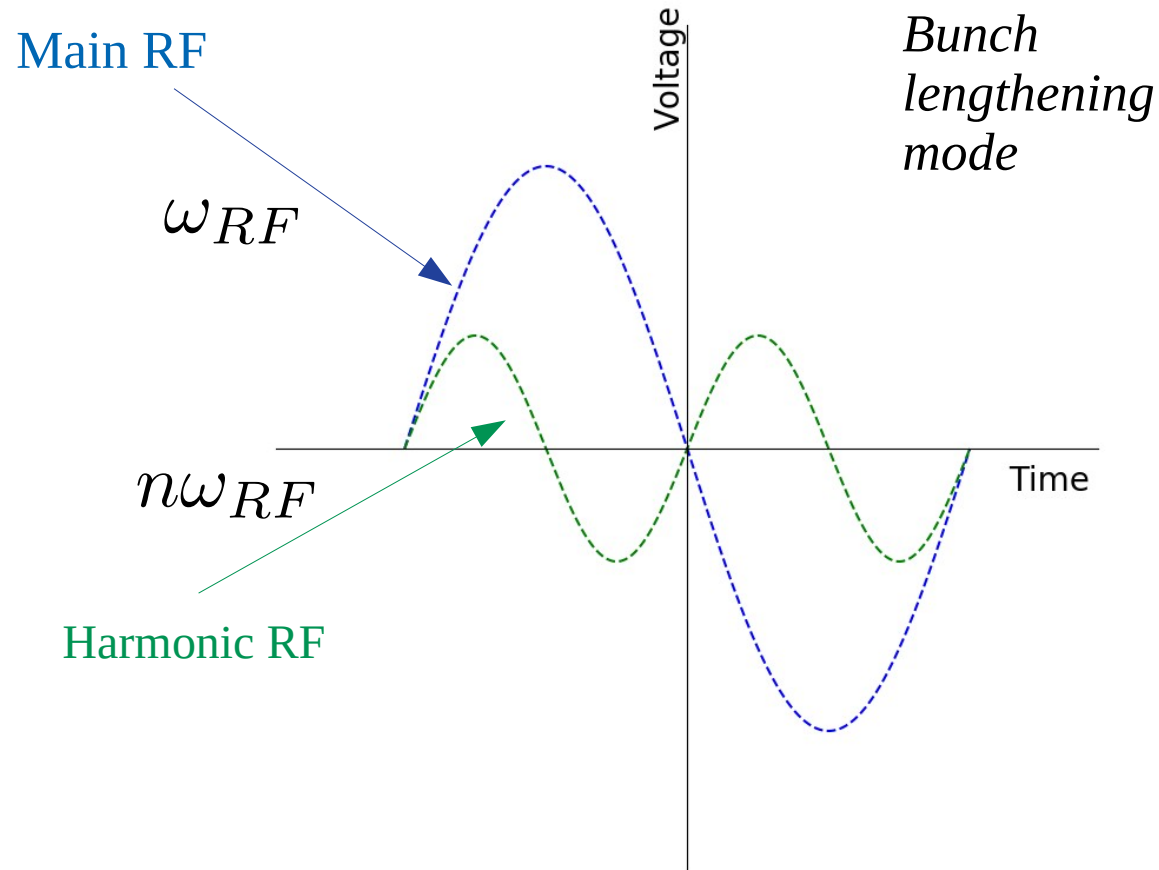
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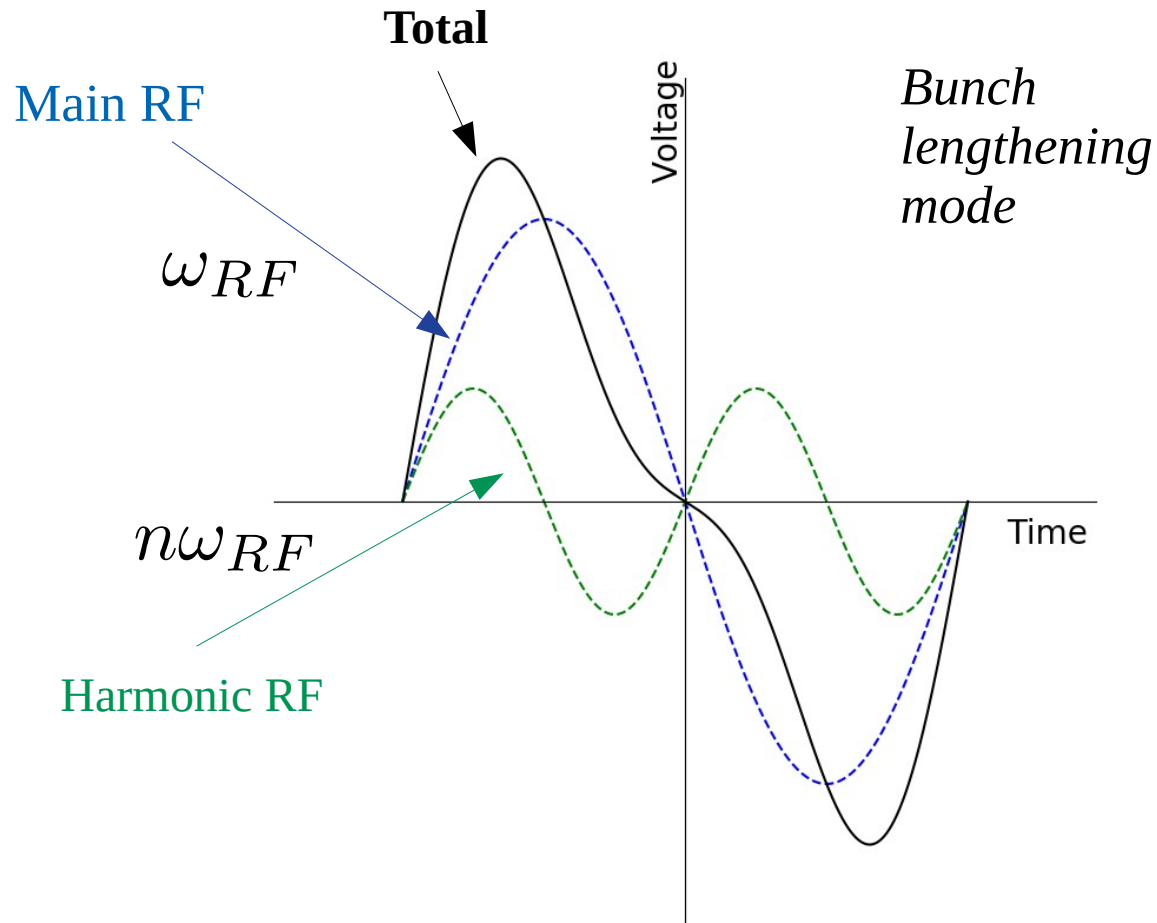
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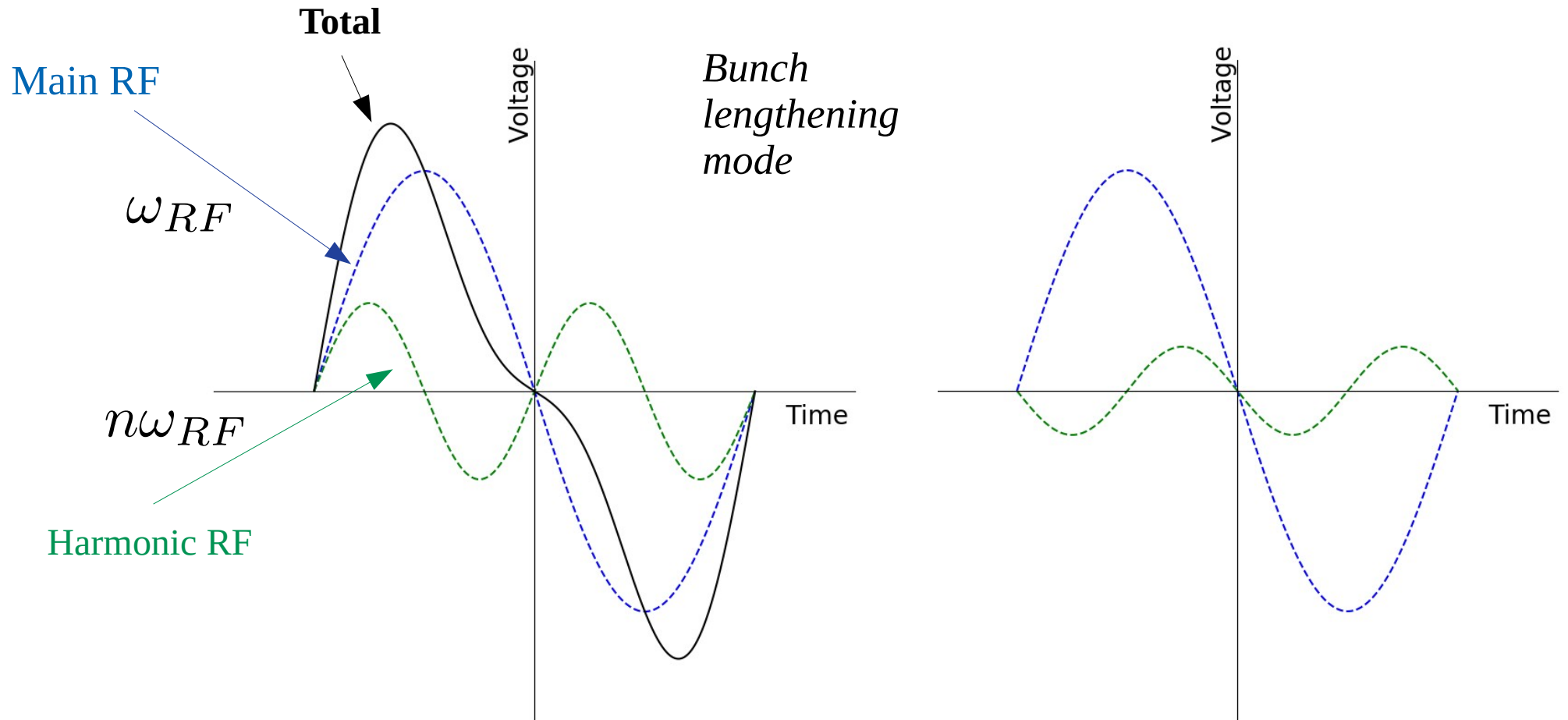
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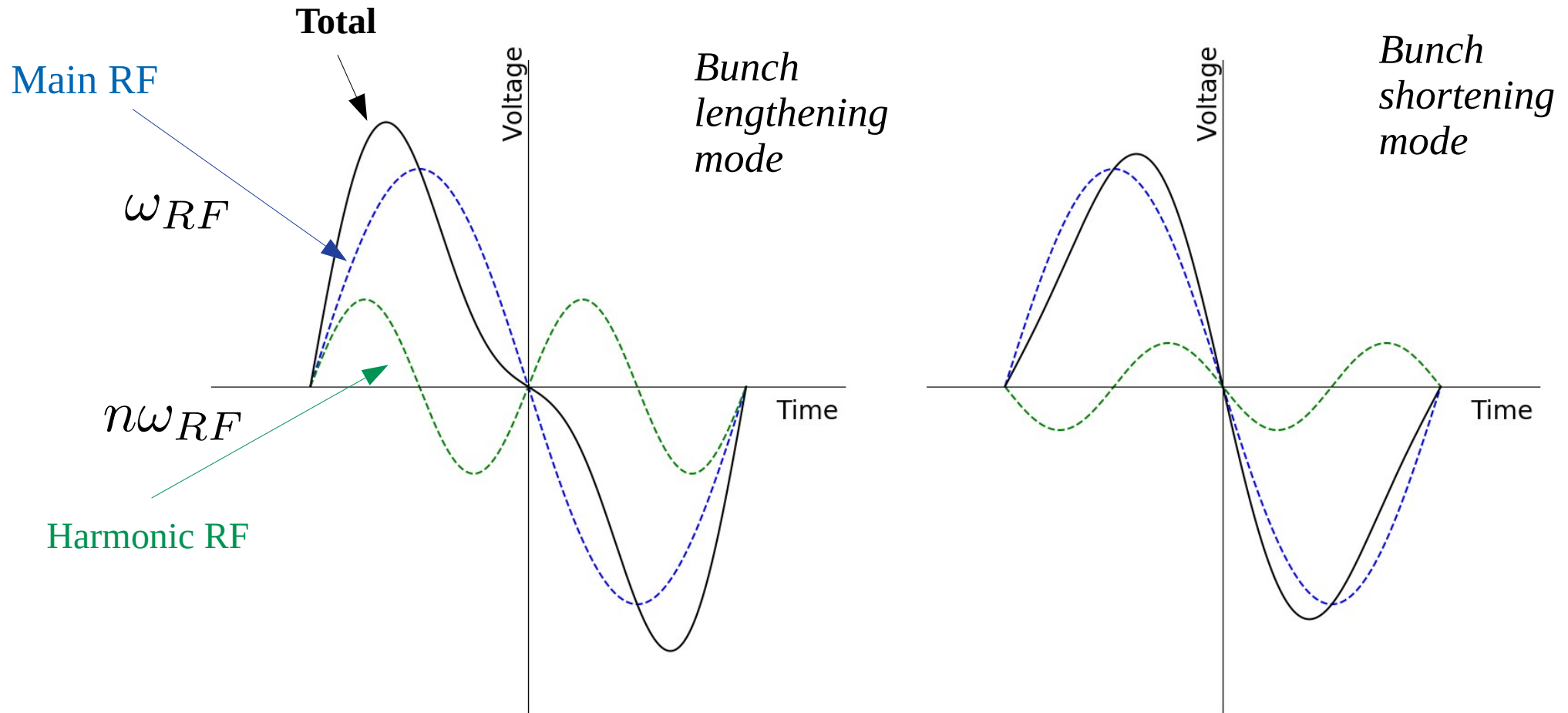
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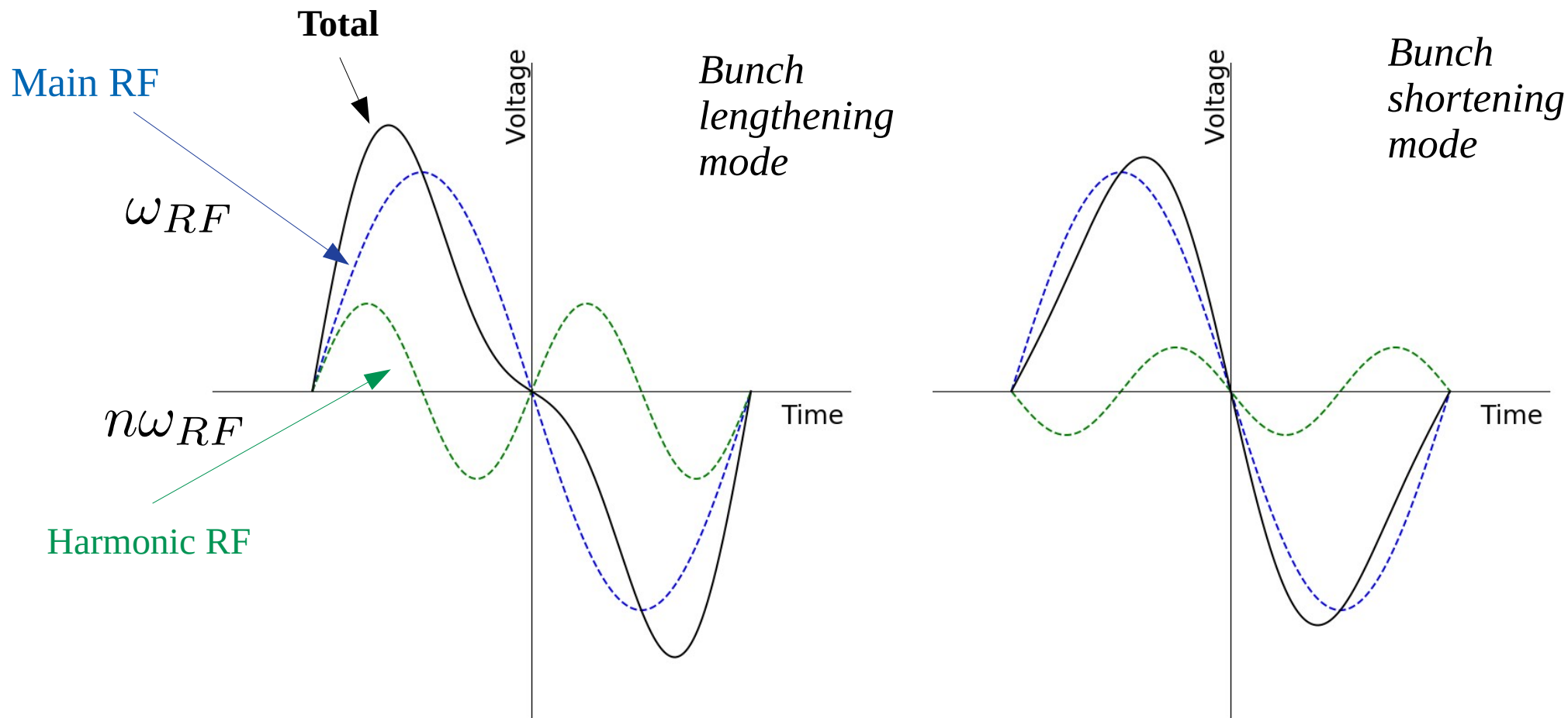
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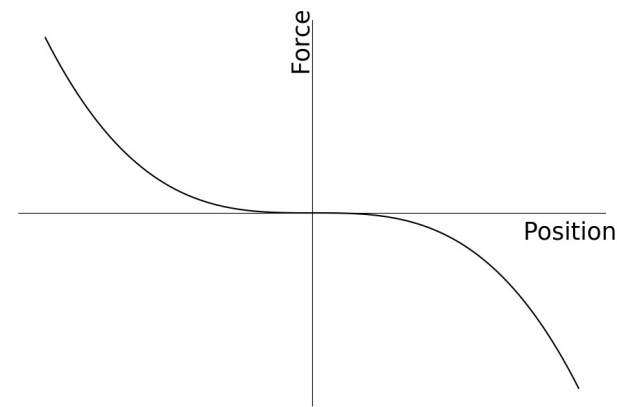
Double harmonic RF

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- The tune spread can be enhanced (or reduced) depending on the relative phase and voltage of the two RF systems
→ Improve / deteriorate Landau damping

- In high energy machines, the frequency spread linked to revolution frequency and the chromaticity is usually small
- Chromatic sextupole magnets are non-linear, yet to first order they don't contribute to the transverse tune spread
 - Dedicated octupole magnets (aka Landau octupoles)



$$\omega(J) = 2\pi(Q_0 + aJ)$$

Optics

Slippage factor
chromaticity correction

Practical aspects

Optics

Slippage factor
chromaticity correction

RF cavities

Frequency, voltage,
harmonic systems

Practical aspects

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Frequency, voltage,
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Magnets

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Impedance

Beam pipe dimensions,
Beam equipment designs
(e.g. instrumentation, collimator,
Vacuum valves),
Material choices, transitions

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Operation

Adiabatic damping during energy ramp
(RF voltage functions, longitudinal blowup)

Fuego's theoretical catch

$$\frac{-1}{\Delta\Omega_n} = \int dJ \frac{J \frac{df_0}{dJ}}{\Omega - \omega(J)}$$

?

$$\frac{-1}{\Delta\Omega_n} = \int d\omega \frac{\rho(\omega)}{\Omega - \omega}$$

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Fuego's theoretical catch

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\approx

- In plasmas, only the density of velocity matters
 - A treatment based on the frequency distribution remains a **reasonable approximation** for many applications in accelerator



Fuego's theoretical catch

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$$\frac{-1}{\Delta\Omega_n} = \int d\omega \frac{\rho(\omega)}{\Omega - \omega} \approx$$

- In plasmas, only the density of velocity matters
 - A treatment based on the frequency distribution remains a **reasonable approximation** for many applications in accelerator
- When the frequency spread arise from non-linear forces, the treatment is slightly different

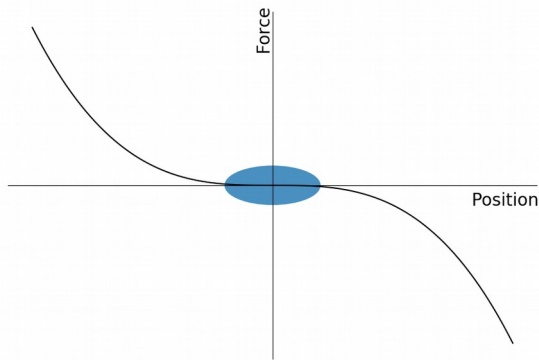
$$\omega(J) = \frac{\partial H}{\partial J}$$



Non-linear collective forces

- Some collective forces are non-linear, they have an impact on Landau damping
 - Due to their dynamic nature, they lead to different behaviours
 - Different dispersion relations

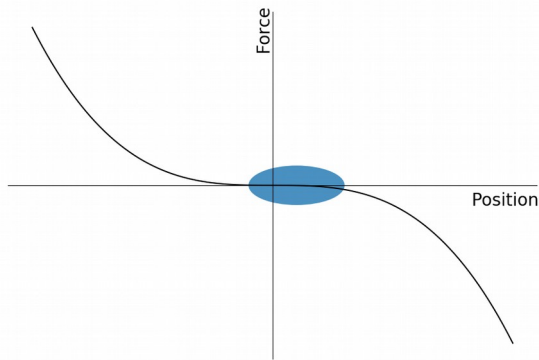
External forces



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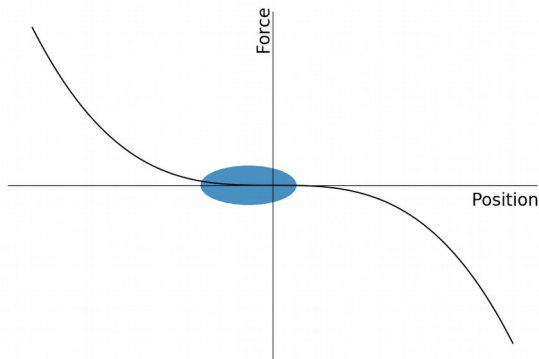
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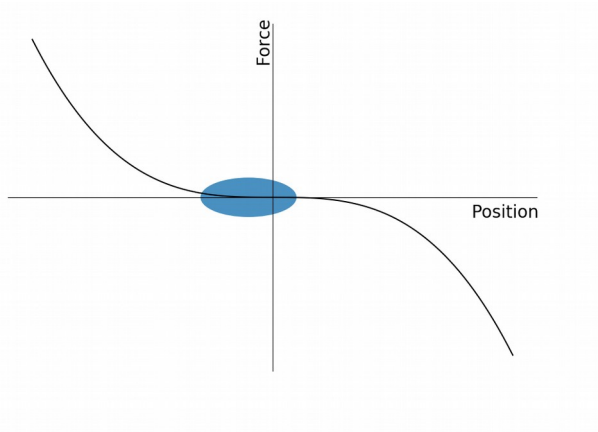
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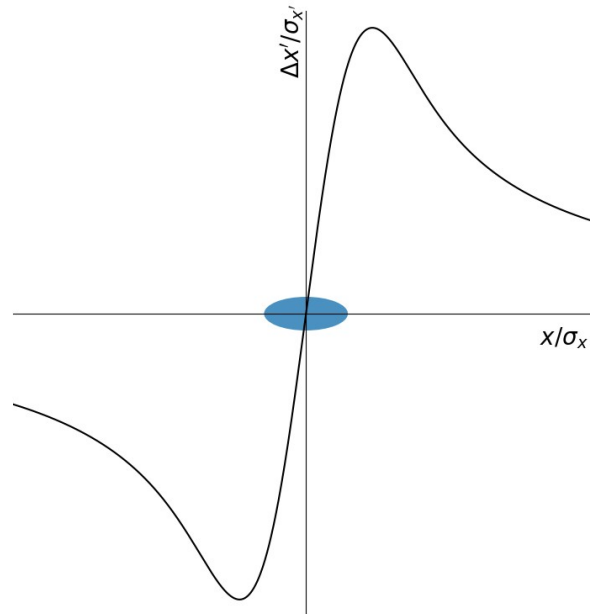
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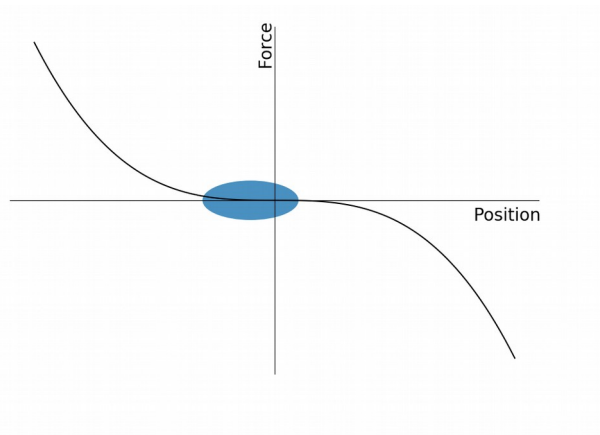
Space-charge



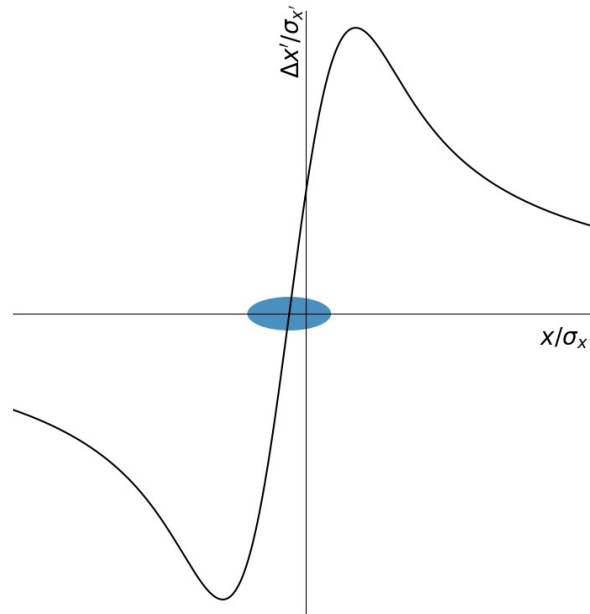
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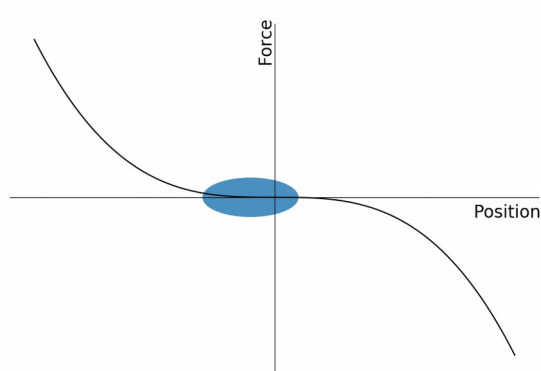
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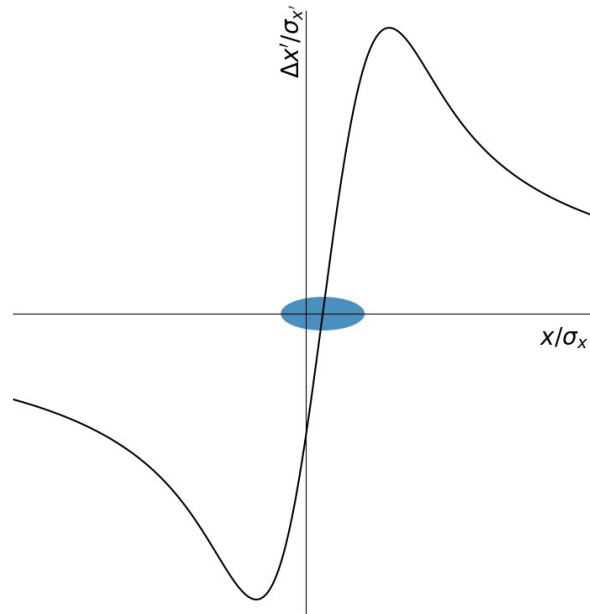
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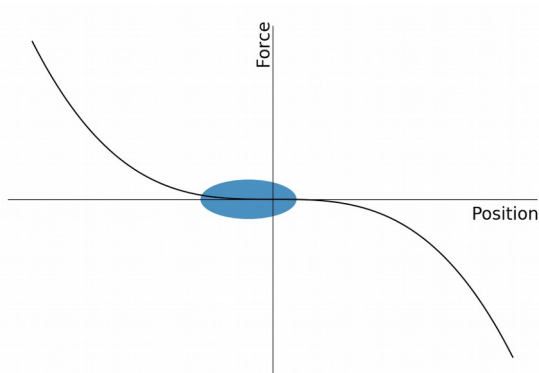
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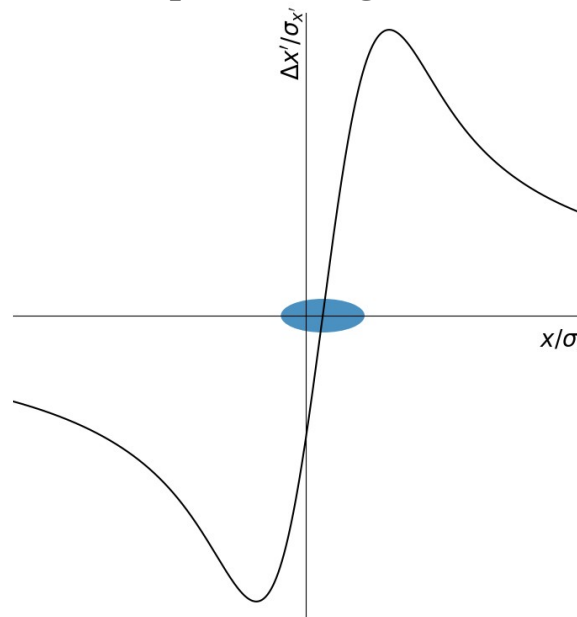
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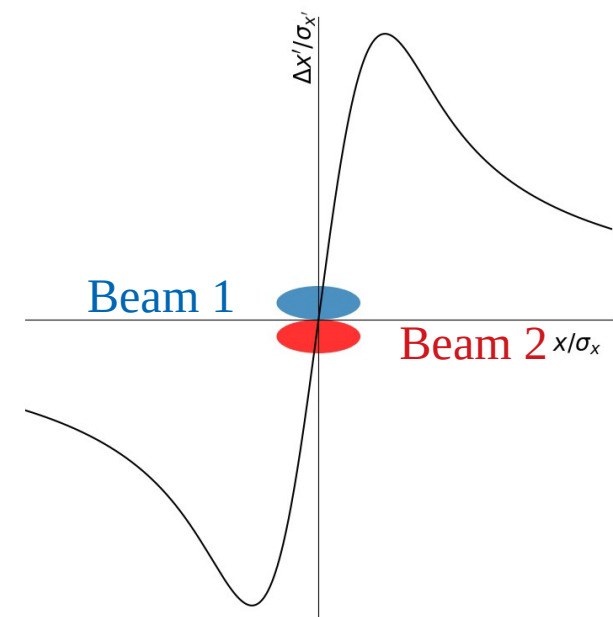
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Space-charge



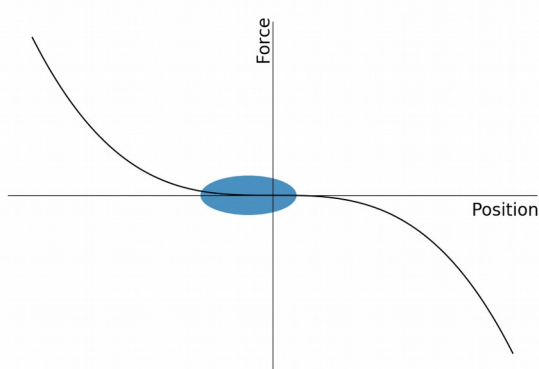
Beam-beam σ -mode



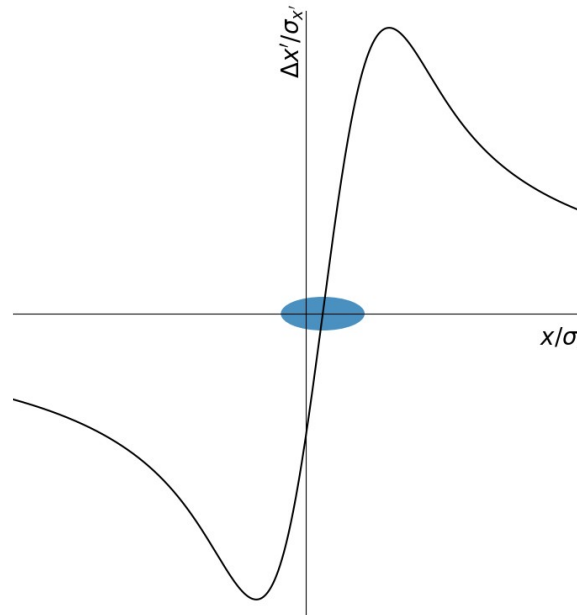
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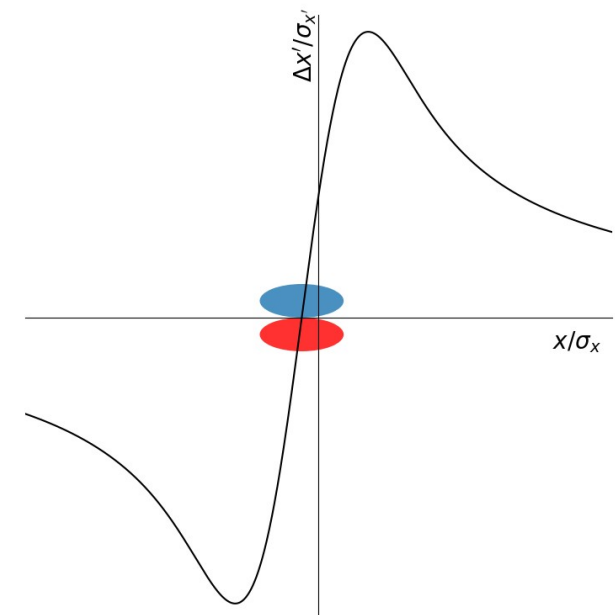
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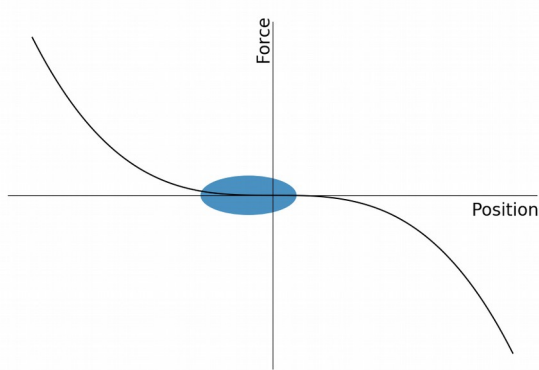
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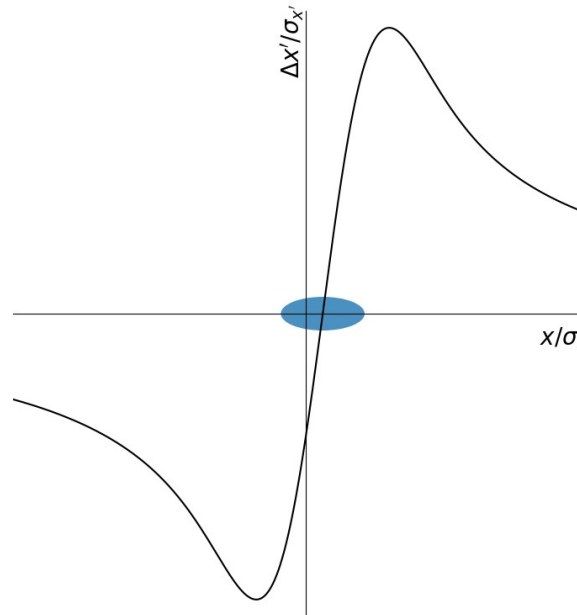
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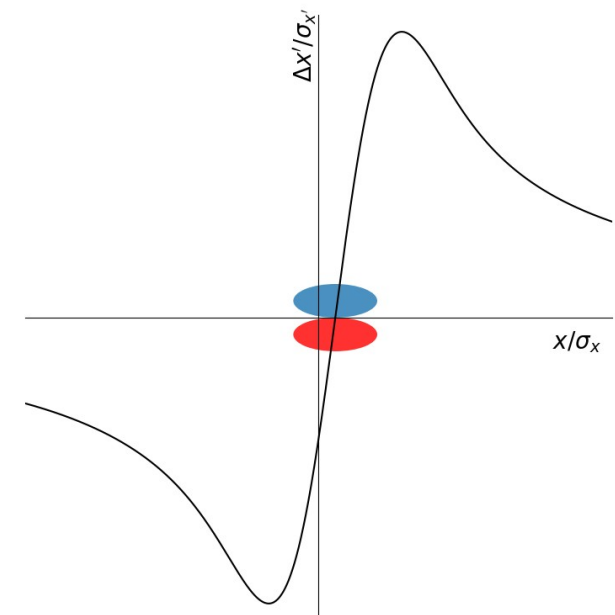
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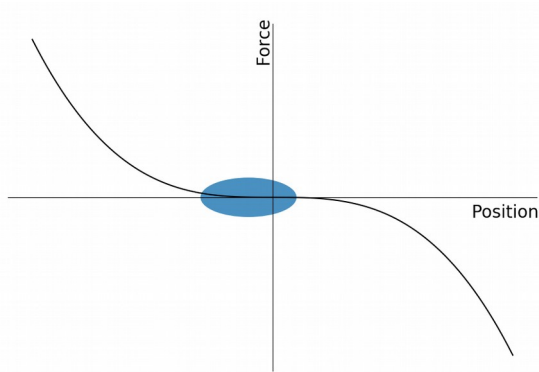
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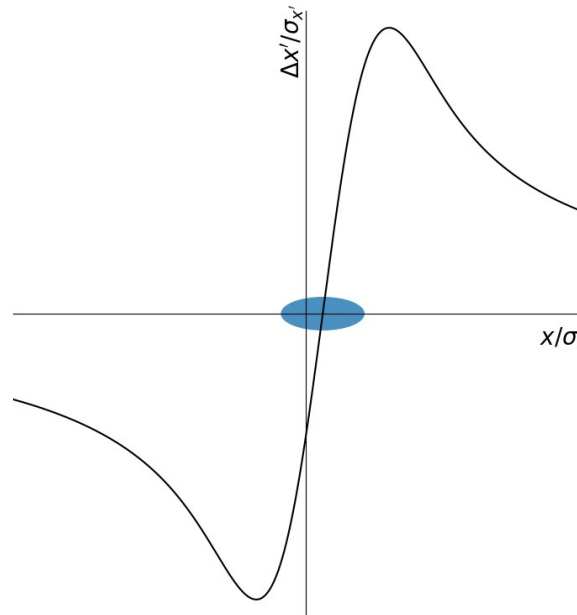
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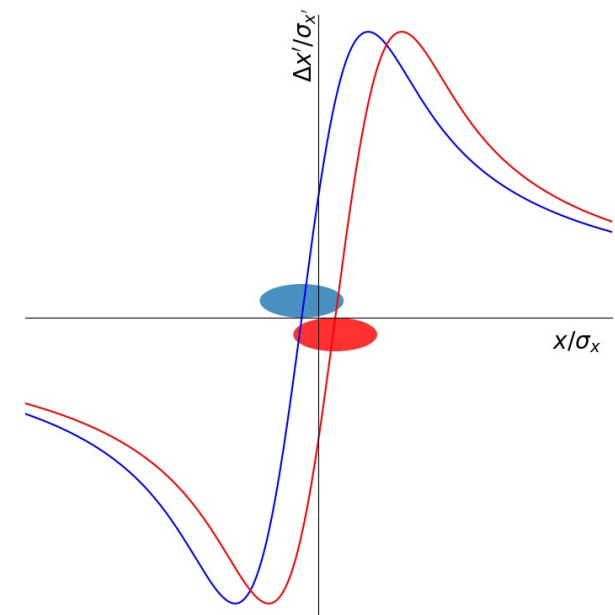
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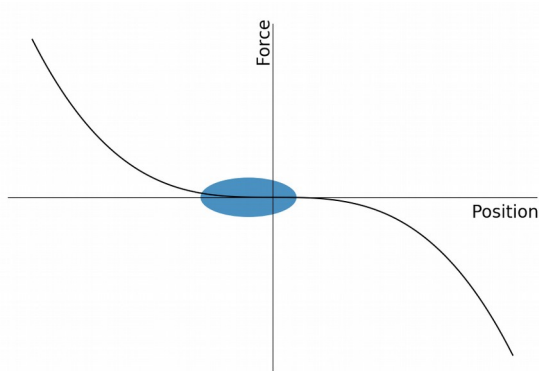
Beam-beam π -mode



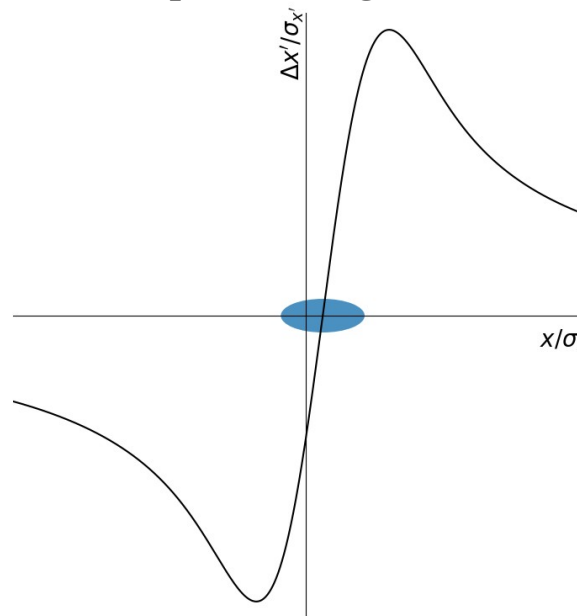
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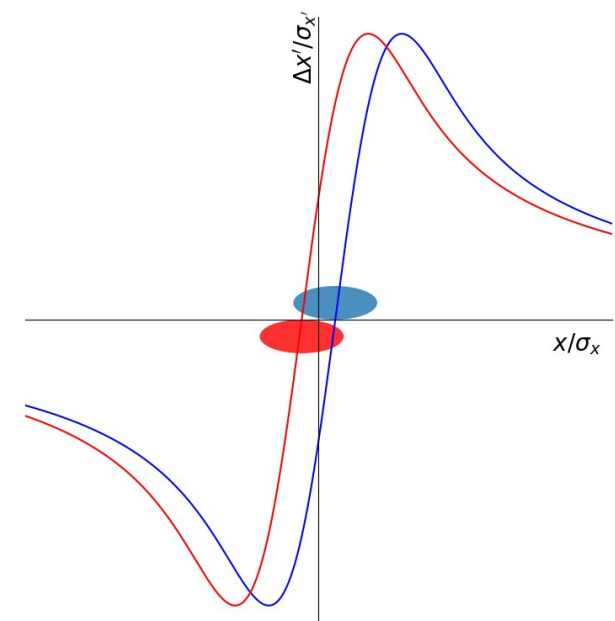
External forces



Space-charge



Beam-beam π -mode



$$\int dJ_x dJ_y \frac{J_x \frac{\partial f_0}{\partial J_x} (\Delta Q_n^x - \Delta Q_{SC}^x(J_x, J_y))}{Q^x - Q_0^x - \Delta Q^x(J_x, J_y) - \Delta Q_{SC}^x(J_x, J_y) - nQ_s} = -1$$

Tune shift of mode n

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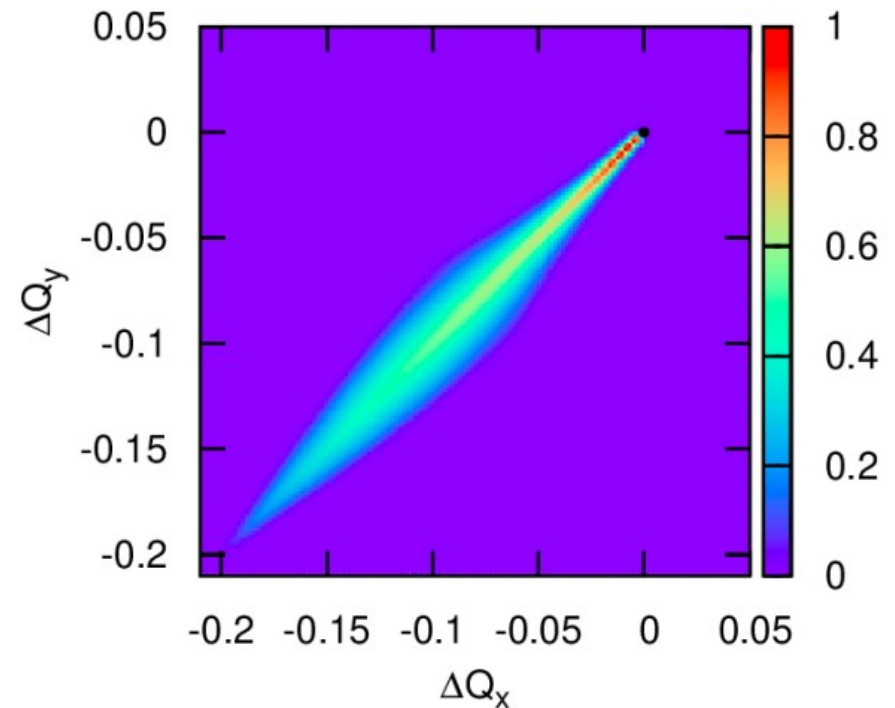
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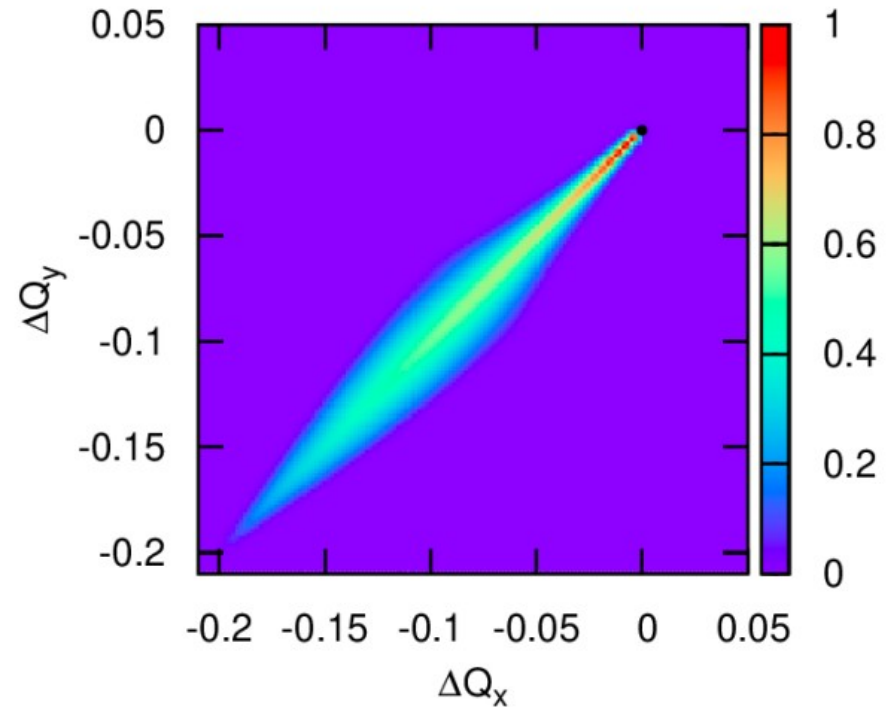
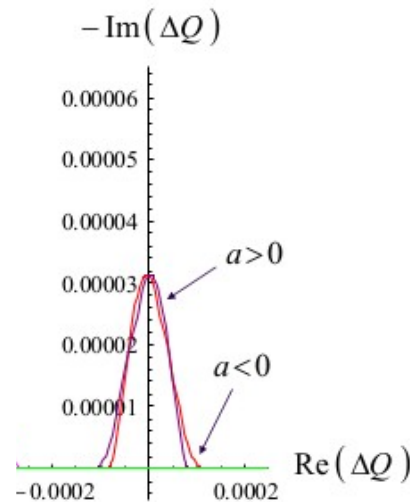
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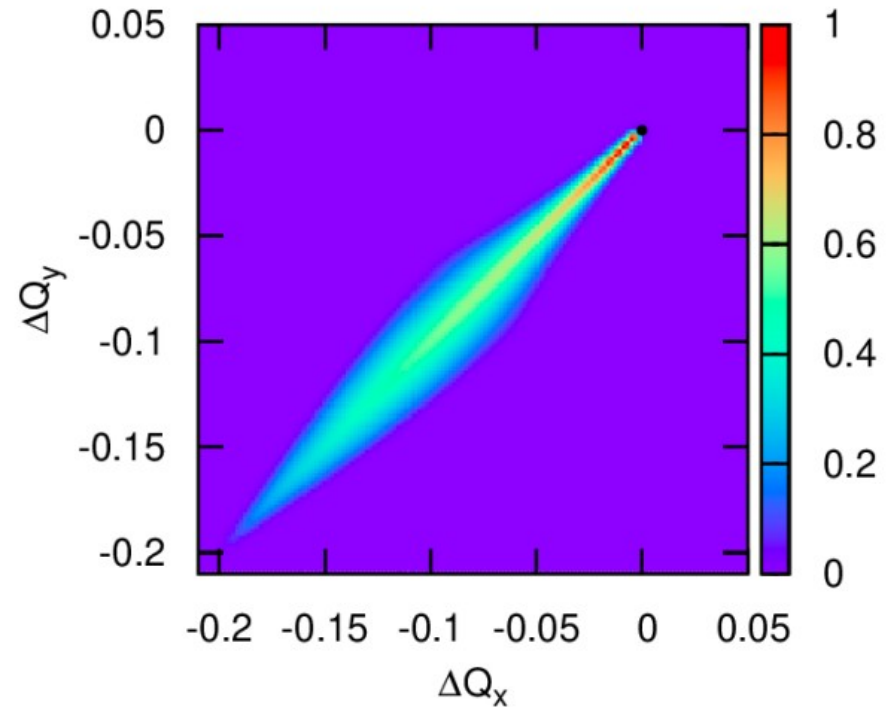
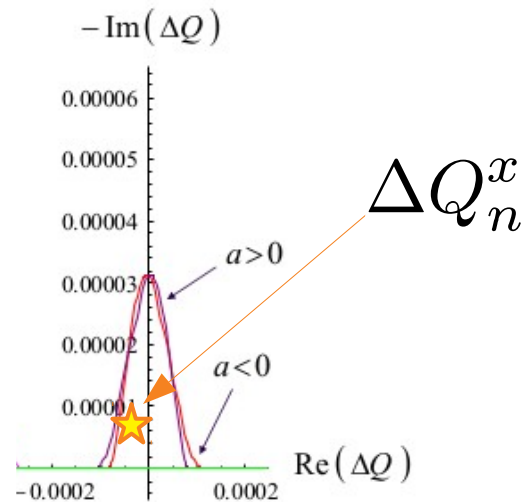


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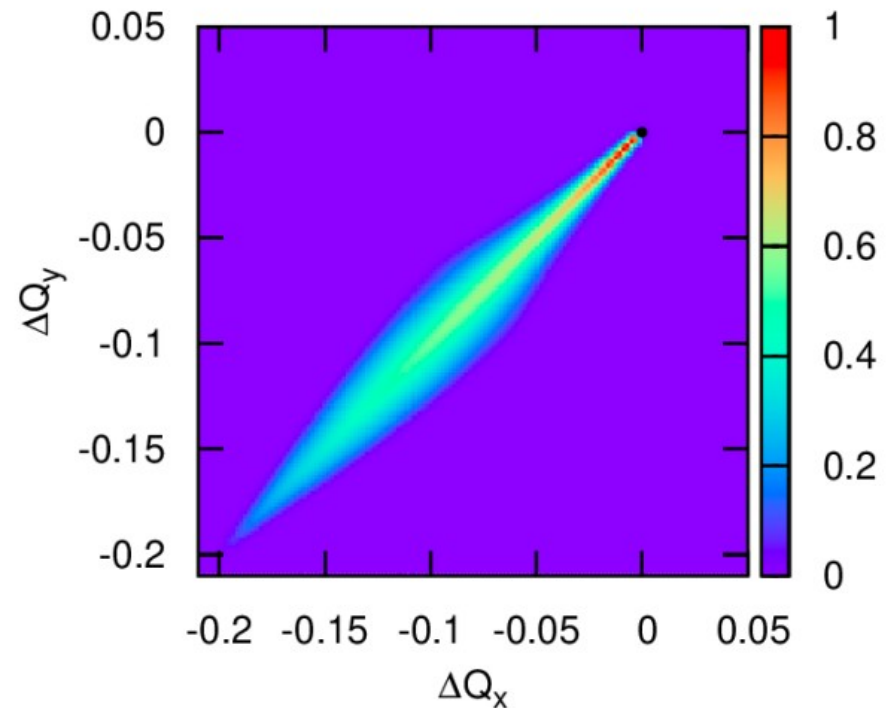
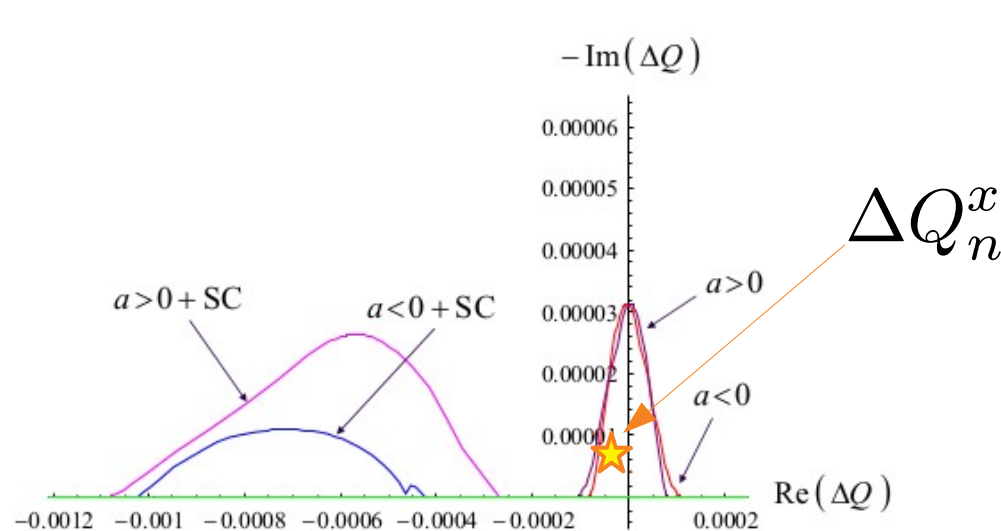
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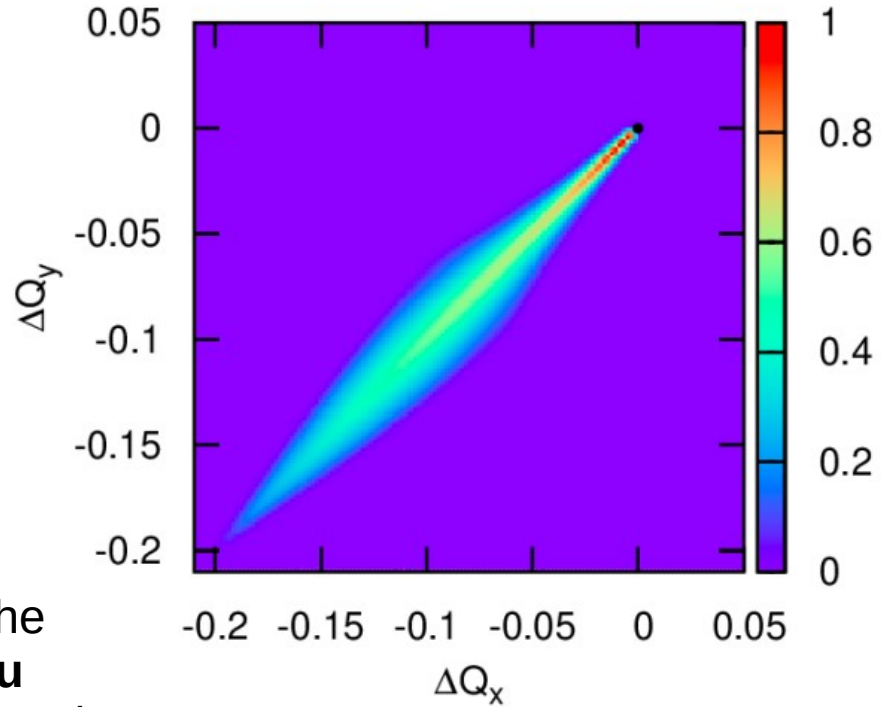
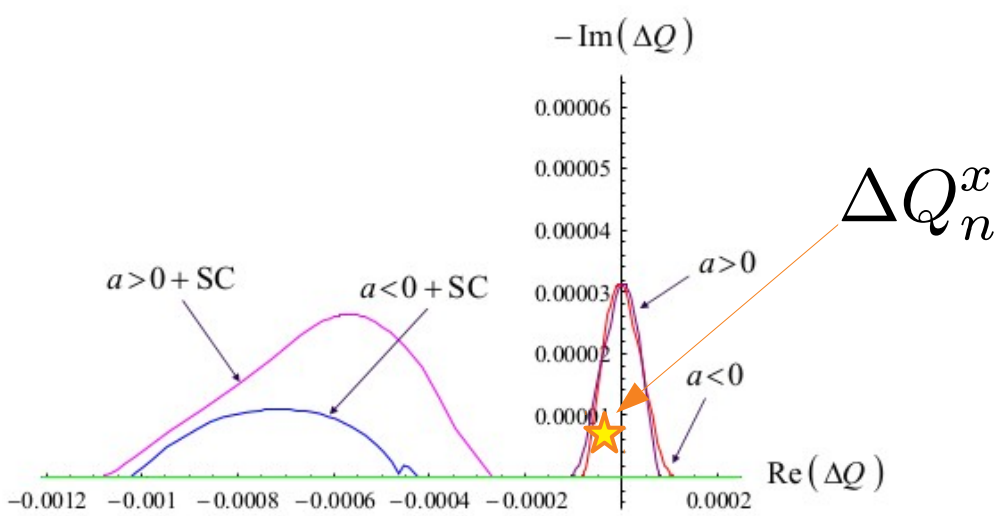


Stability of the rigid bunch mode with space-charge [Metral, Kornilov]

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$$\int dJ_x dJ_y \frac{J_x \frac{\partial f_0}{\partial J_x} (\Delta Q_n^x - \Delta Q_{SC}^x(J_x, J_y))}{Q^x - Q_0^x - \Delta Q^x(J_x, J_y) - \Delta Q_{SC}^x(J_x, J_y) - nQ_s} = -1$$

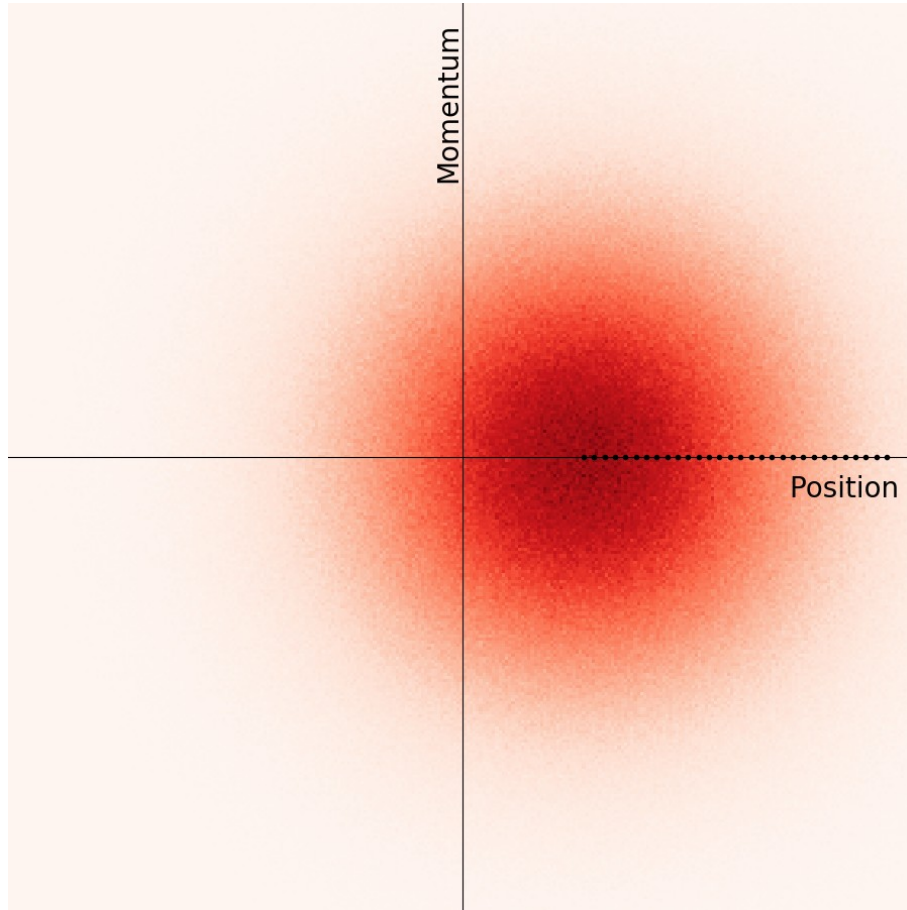
Tune shift of mode n
Tune of the mode including Landau damping
Amplitude detuning due to external non-linearities (e.g. octupole magnets)
Amplitude detuning due to space-charge



- By shifting the so-called incoherent spectrum from the coherent modes, space-charge can **remove Landau damping** for modes otherwise stabilised e.g. by octupoles

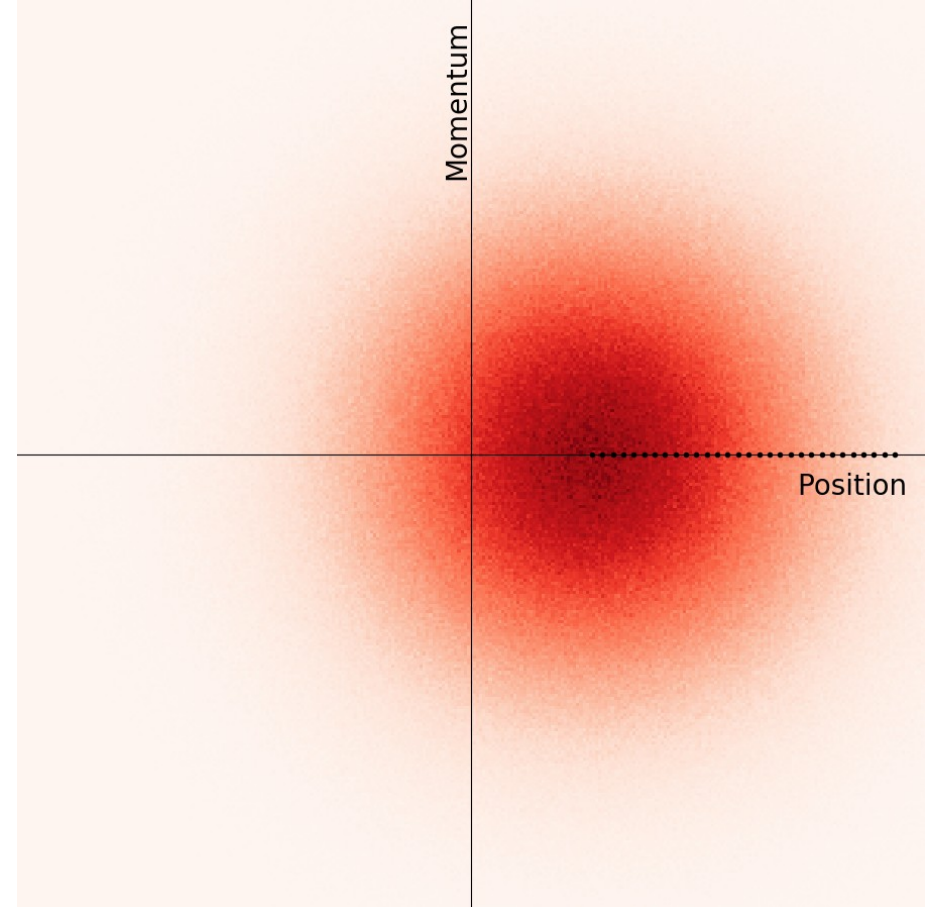
Decoherence of the rigid bunch mode with space-charge

Without space-charge



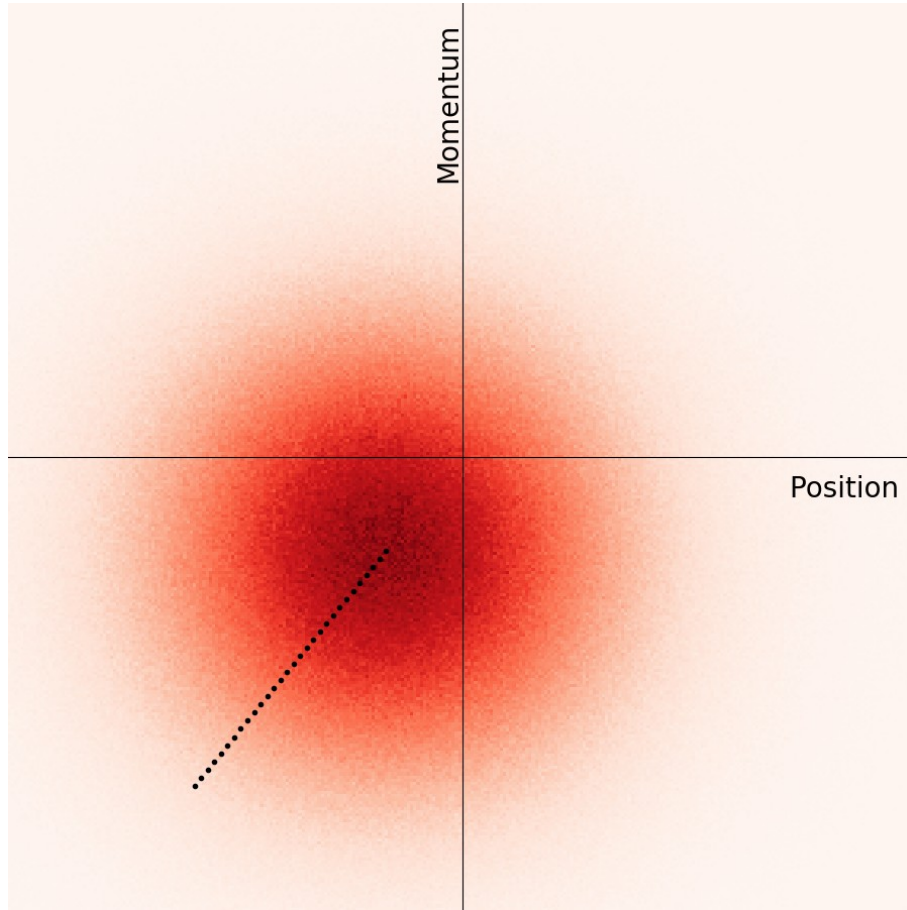
Turn
0

With space-charge



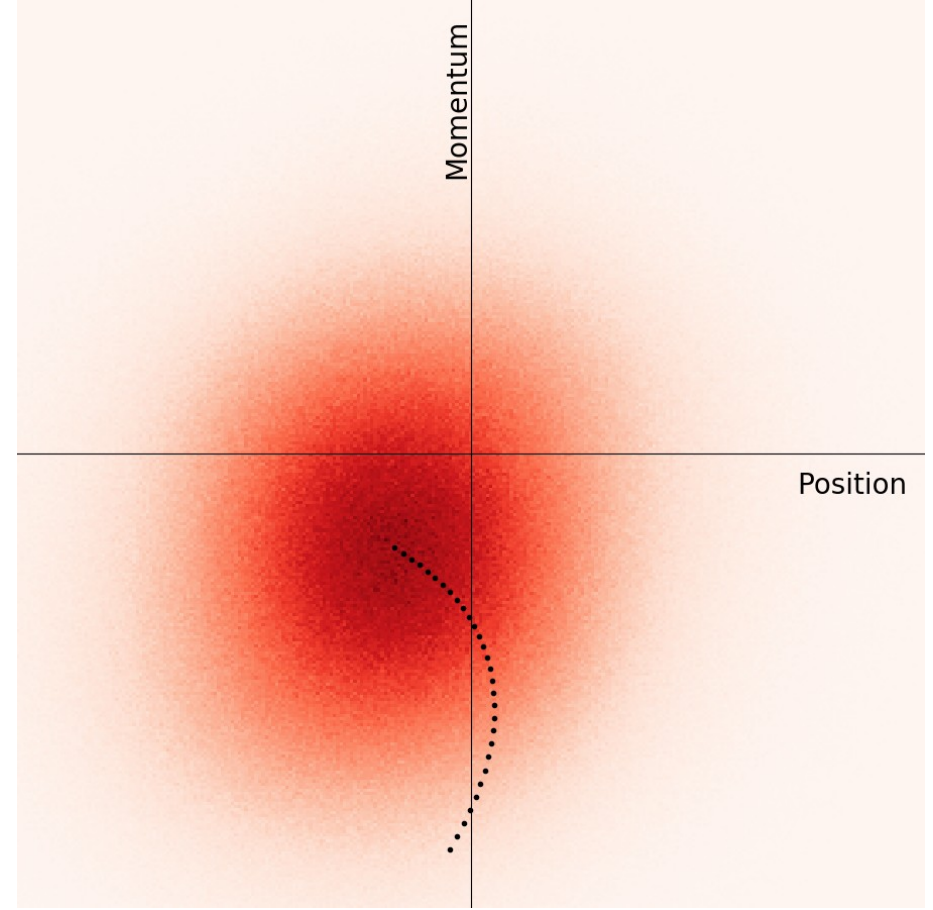
Decoherence of the rigid bunch mode with space-charge

Without space-charge



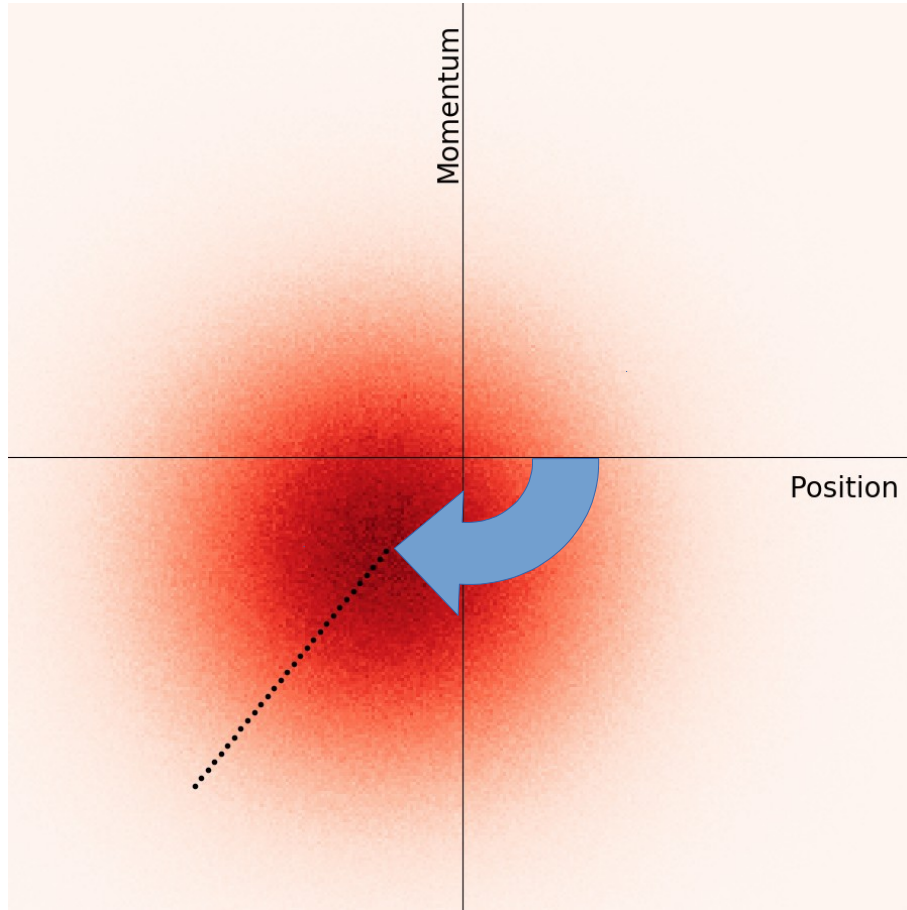
Turn
14

With space-charge



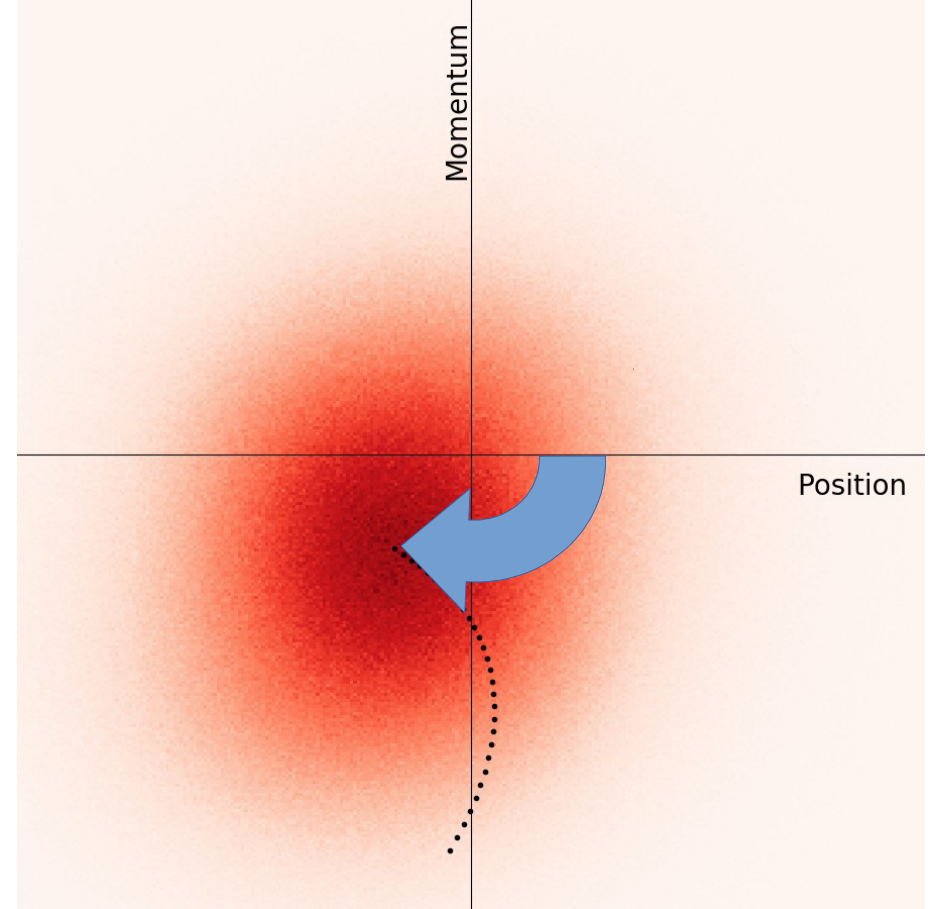
Decoherence of the rigid bunch mode with space-charge

Without space-charge



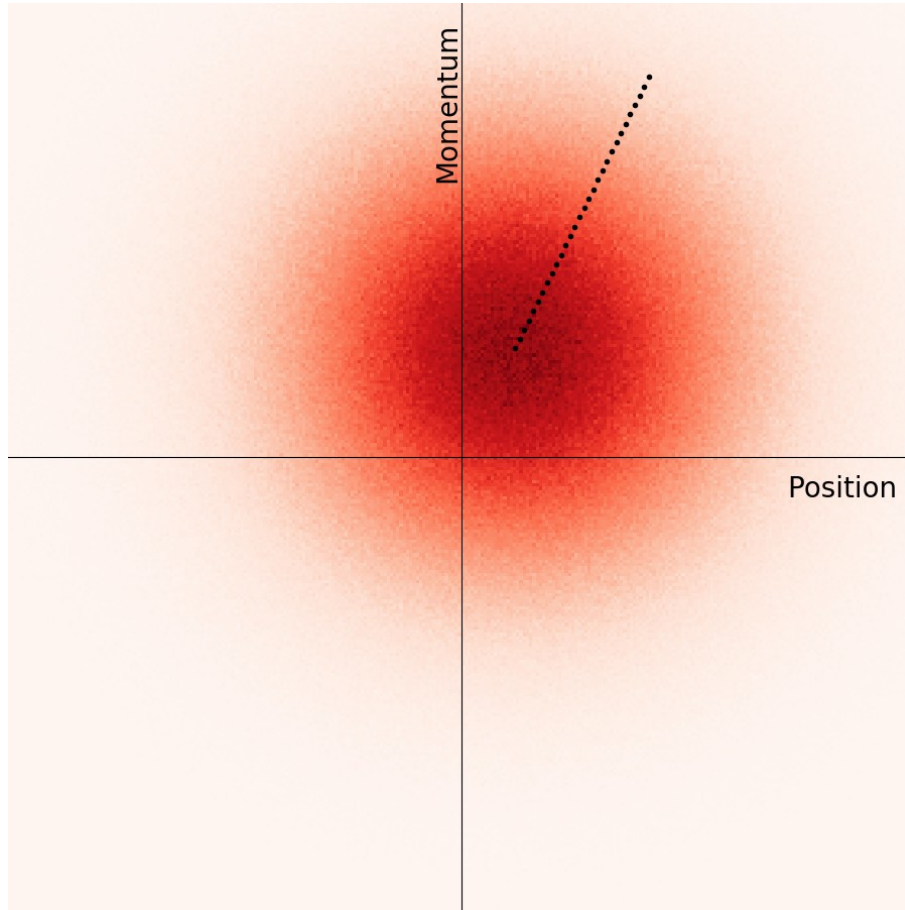
Turn
14

With space-charge



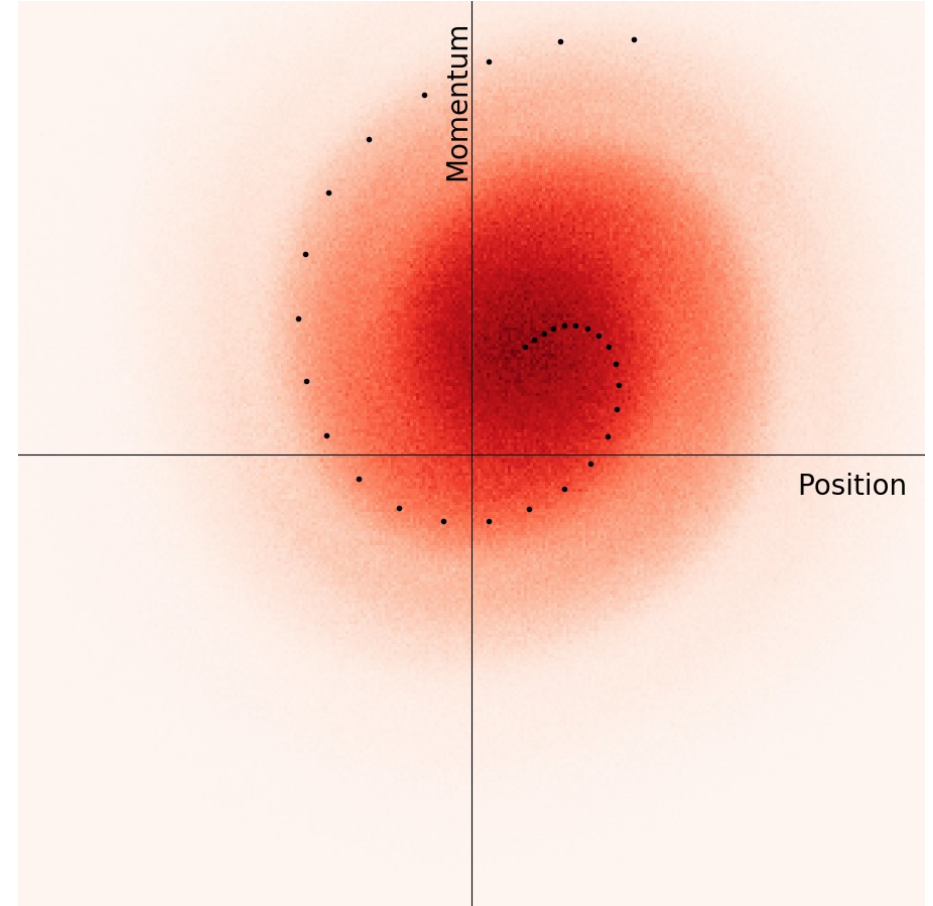
Decoherence of the rigid bunch mode with space-charge

Without space-charge



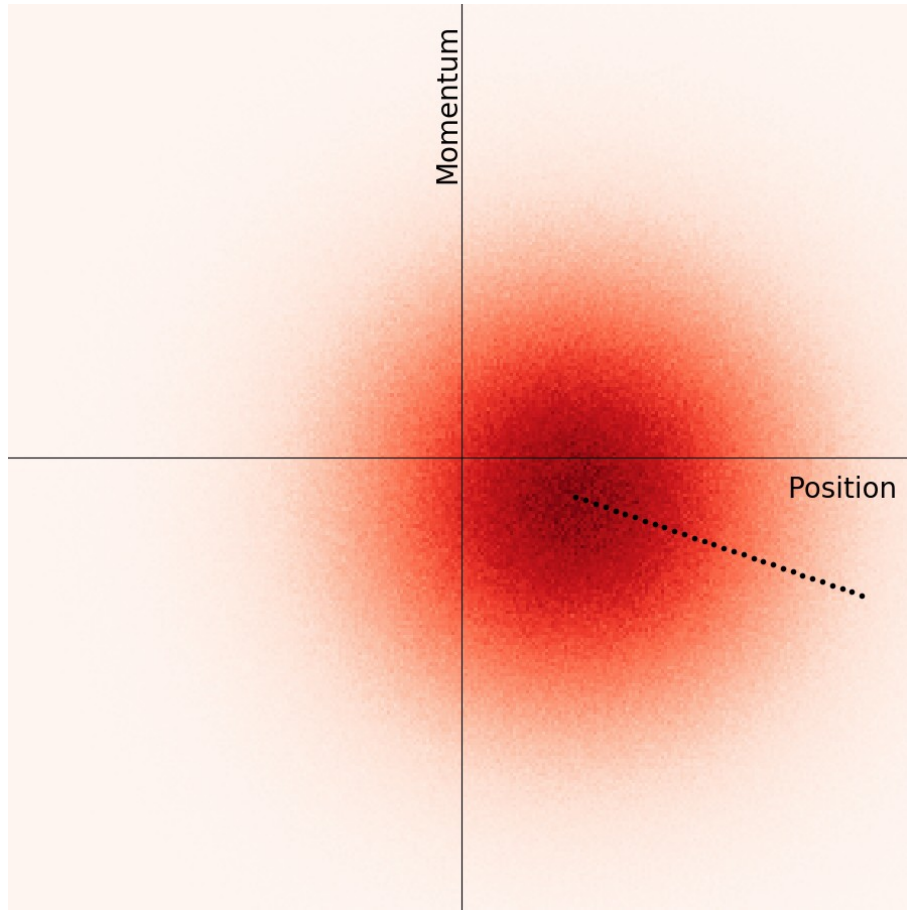
Turn
99

With space-charge



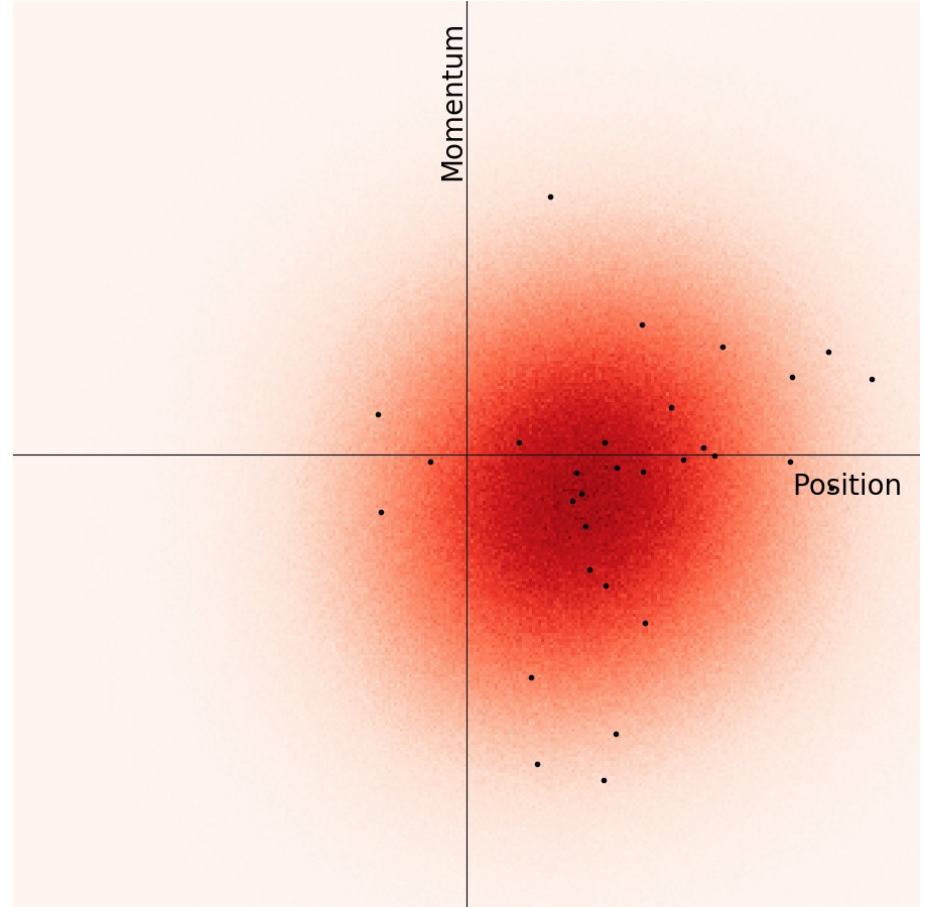
Decoherence of the rigid bunch mode with space-charge

Without space-charge

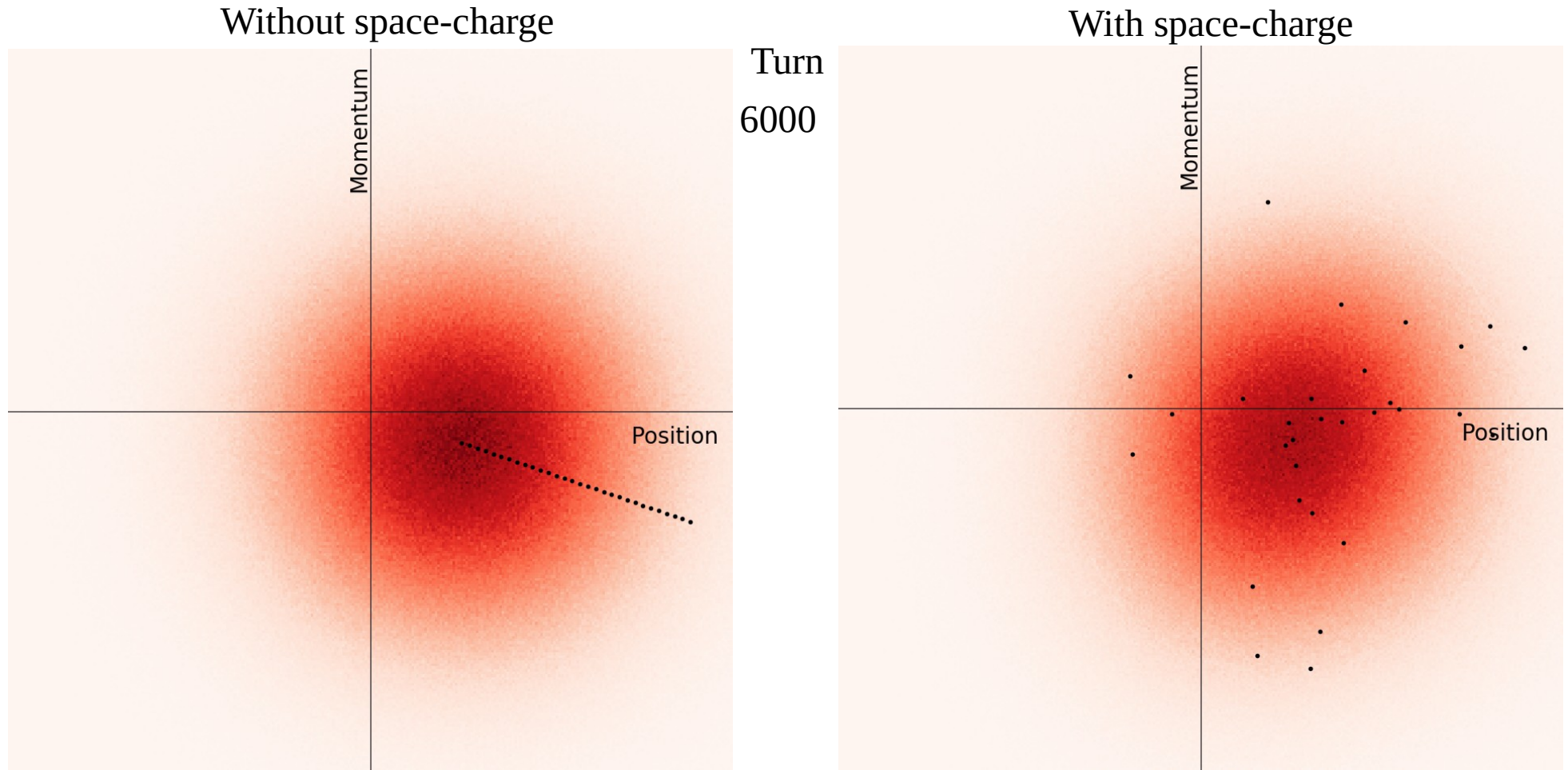


Turn
6000

With space-charge



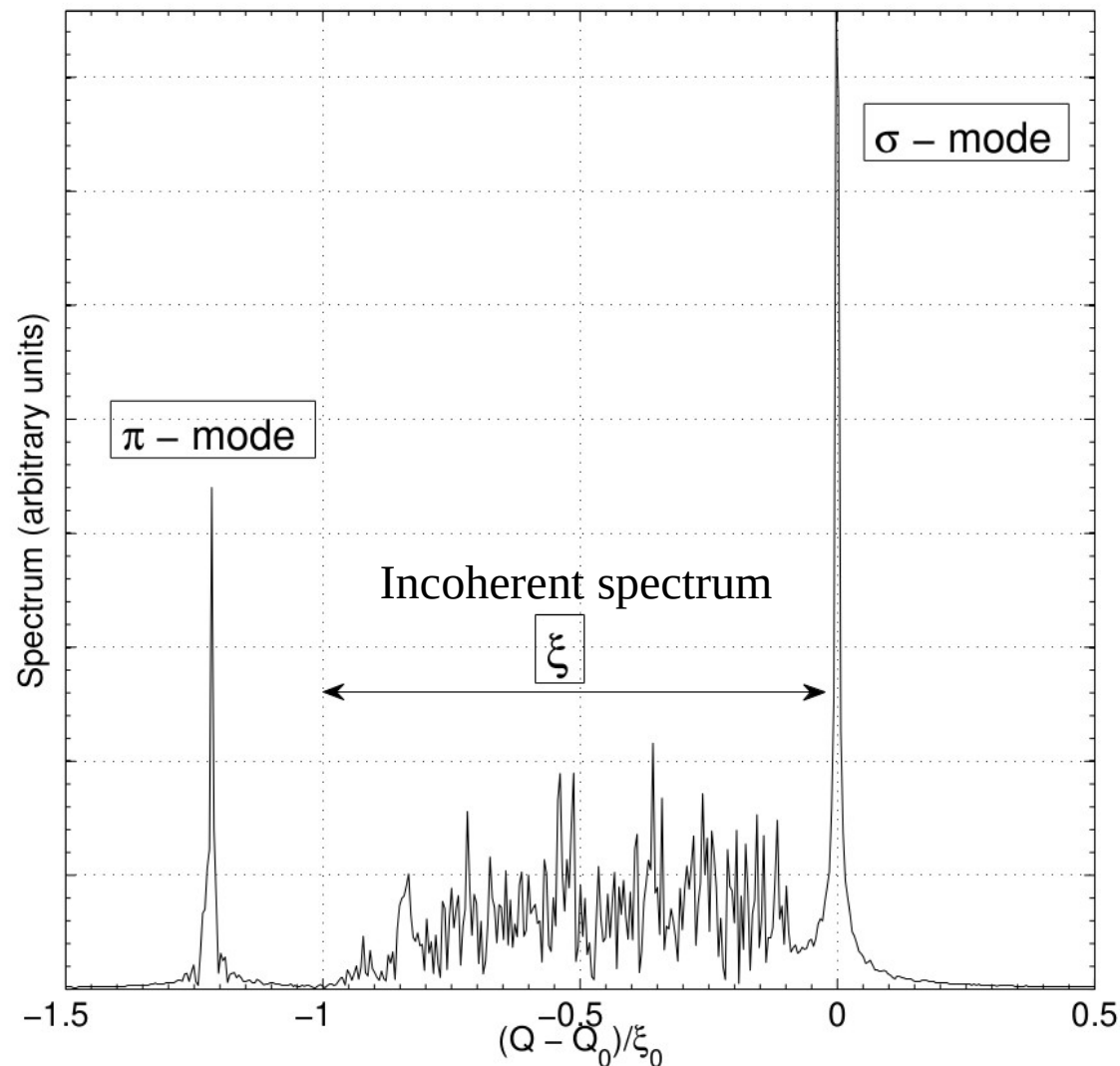
Decoherence of the rigid bunch mode with space-charge



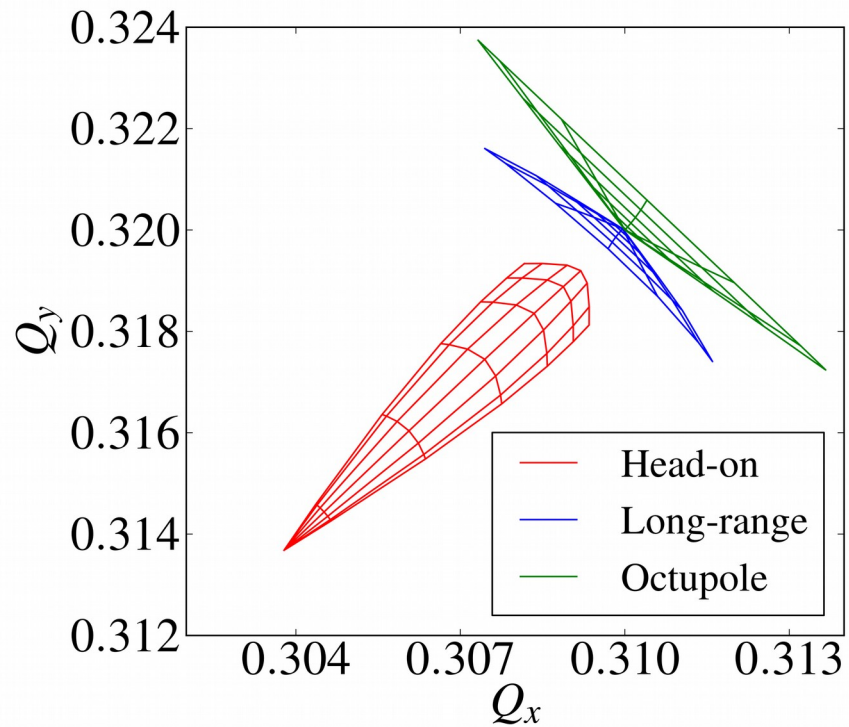
- The motion of the centroid is not affected by space-charge
 - Coherent mode
- The motion of single particles around the centroid is affected
 - Incoherent tune spread

Stability of the rigid bunch mode with beam-beam [Pieloni]

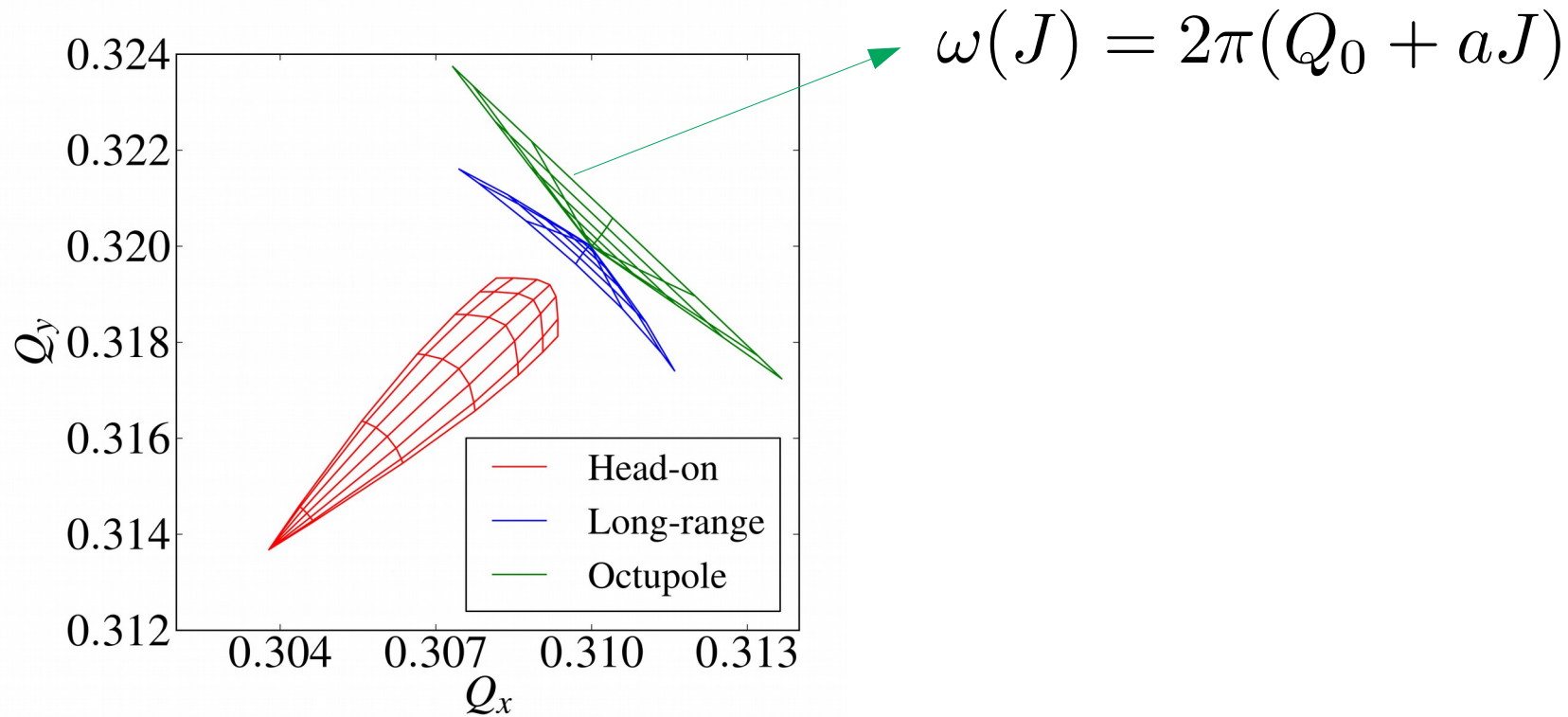
- A similar effect occurs with the coherent modes generated by beam-beam interactions
 - They are outside of the incoherent spectrum, Landau damping is lost



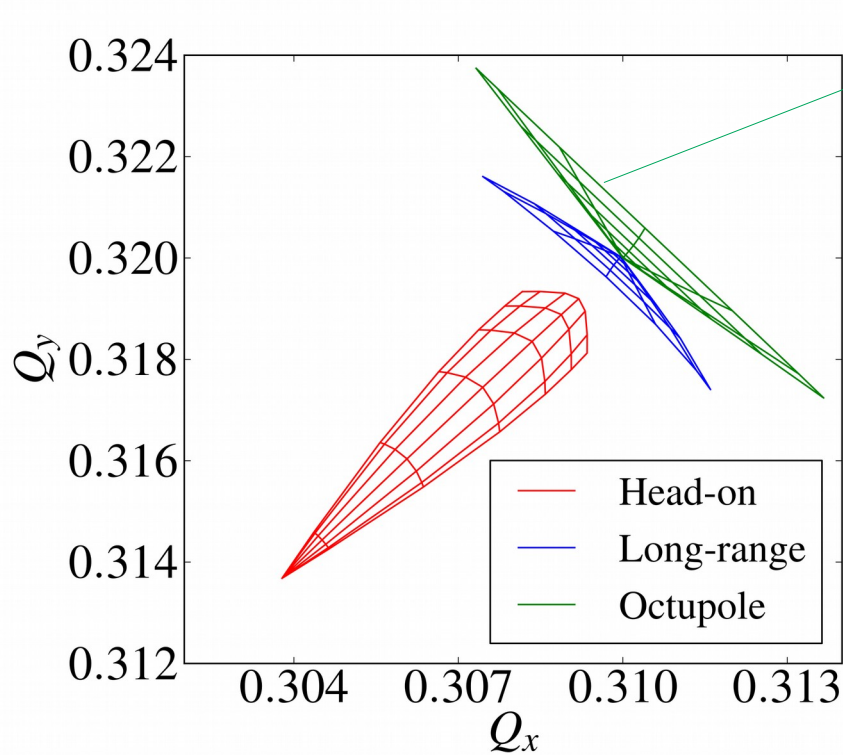
- If the coherent modes are suppressed (e.g. with an active feedback), the remaining tune spread can be beneficial for other modes



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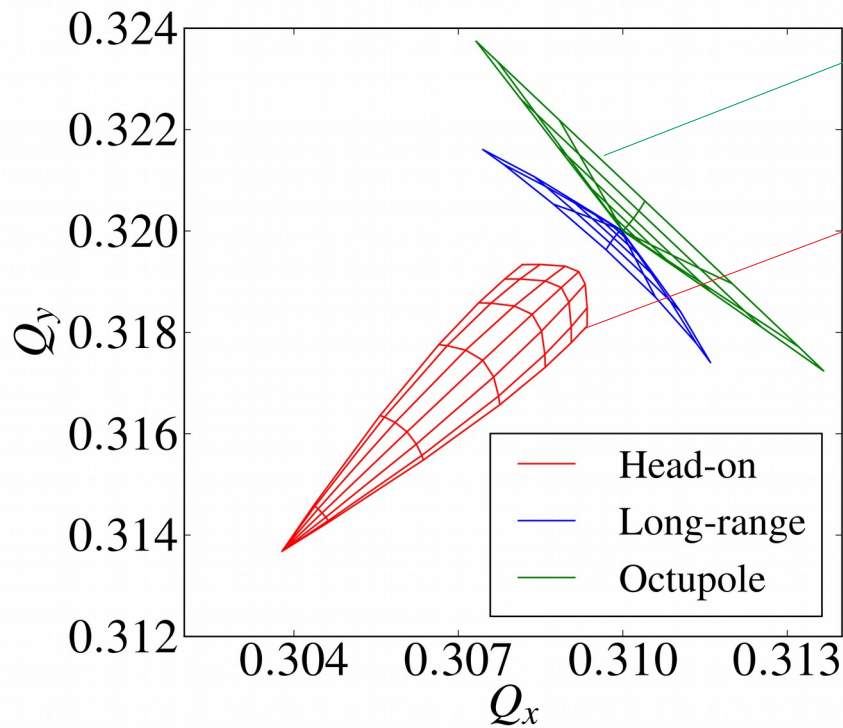
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$$\omega(J) = 2\pi(Q_0 + aJ)$$

$$\frac{-1}{\Delta\Omega_n} = \int dJ \frac{J \frac{df_0}{dJ}}{\Omega - \omega(J)}$$

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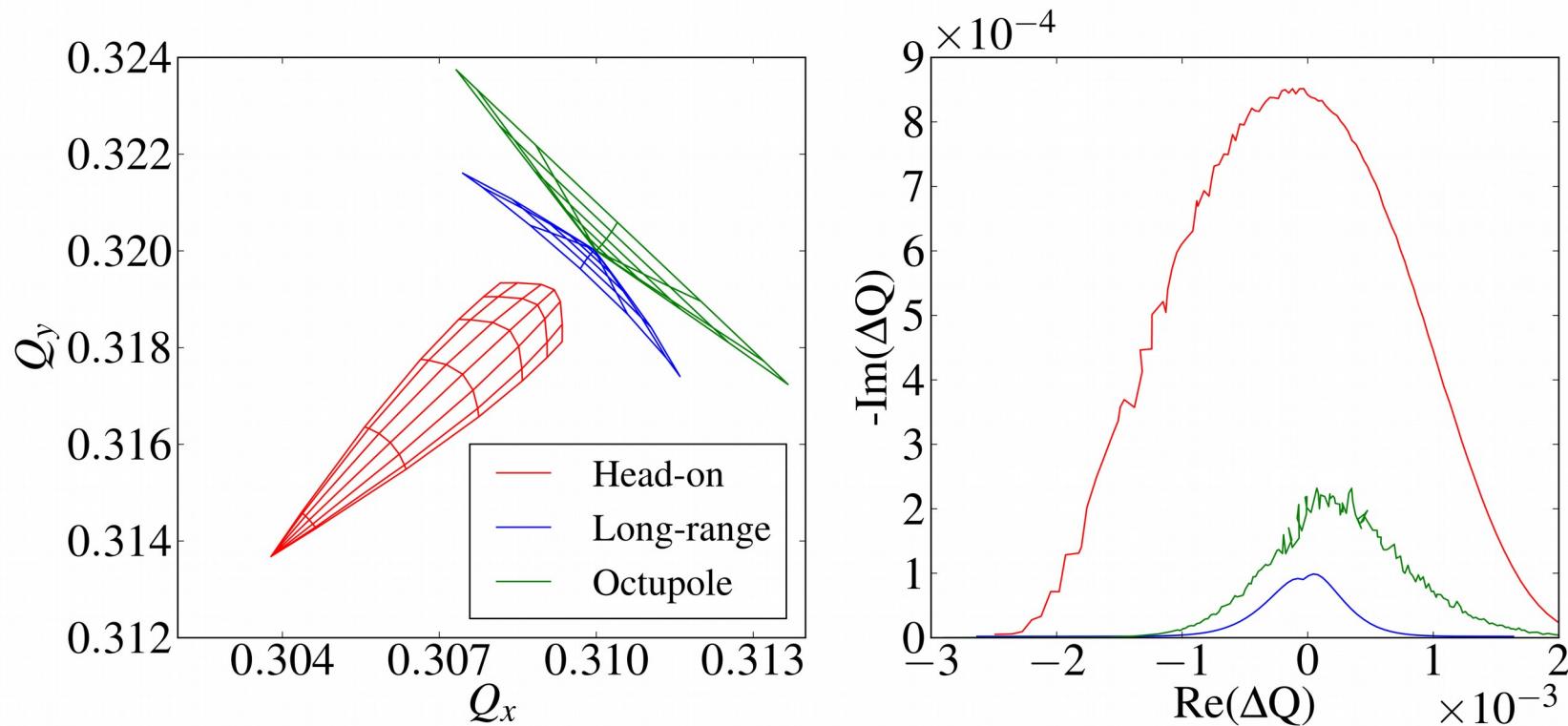


$$\omega(J) = 2\pi(Q_0 + aJ)$$

$$\omega(J) = 2\pi Q_0 - \frac{Nr_0}{\gamma J} \left(1 - e^{-\frac{J\beta}{2\sigma^2}} I_0 \left(\frac{J\beta}{2\sigma^2} \right) \right)$$

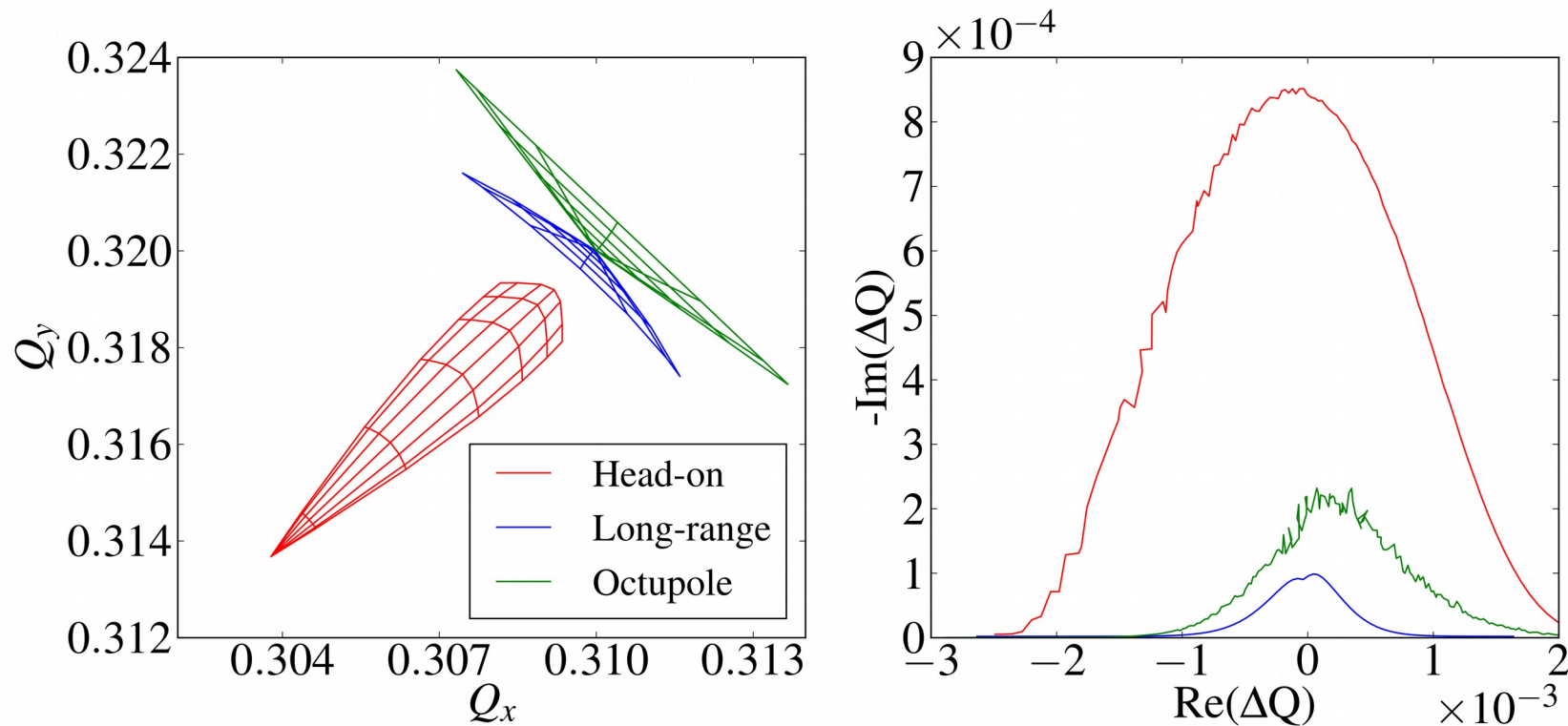
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- Due to its different dependence on the action, the amplitude detuning due to head-on beam-beam interactions is more efficient at producing Landau damping than octupoles!

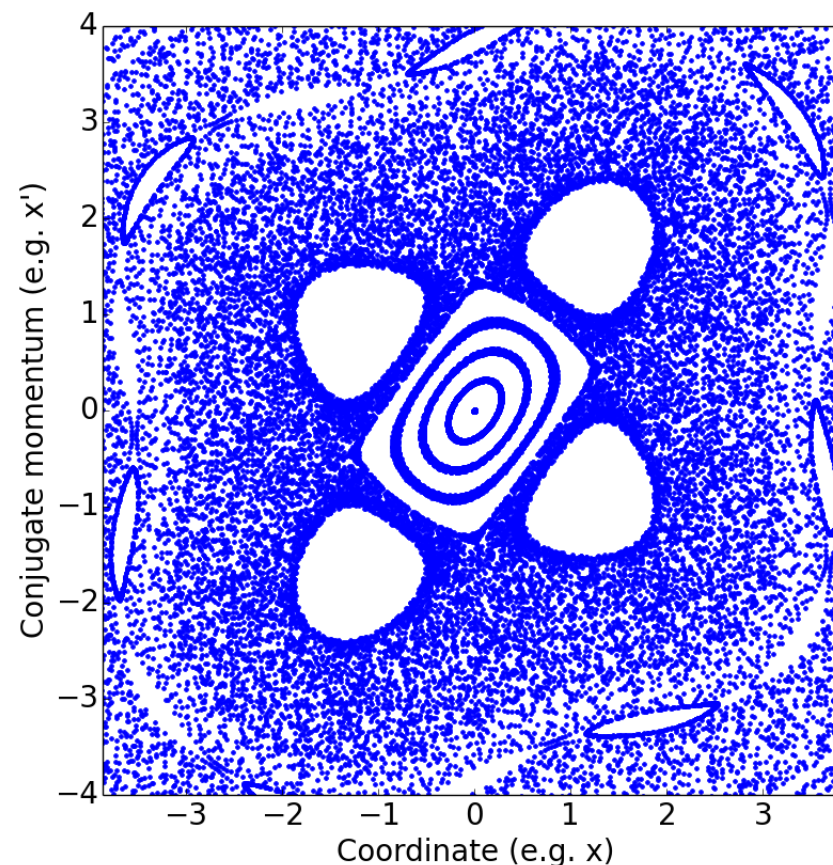
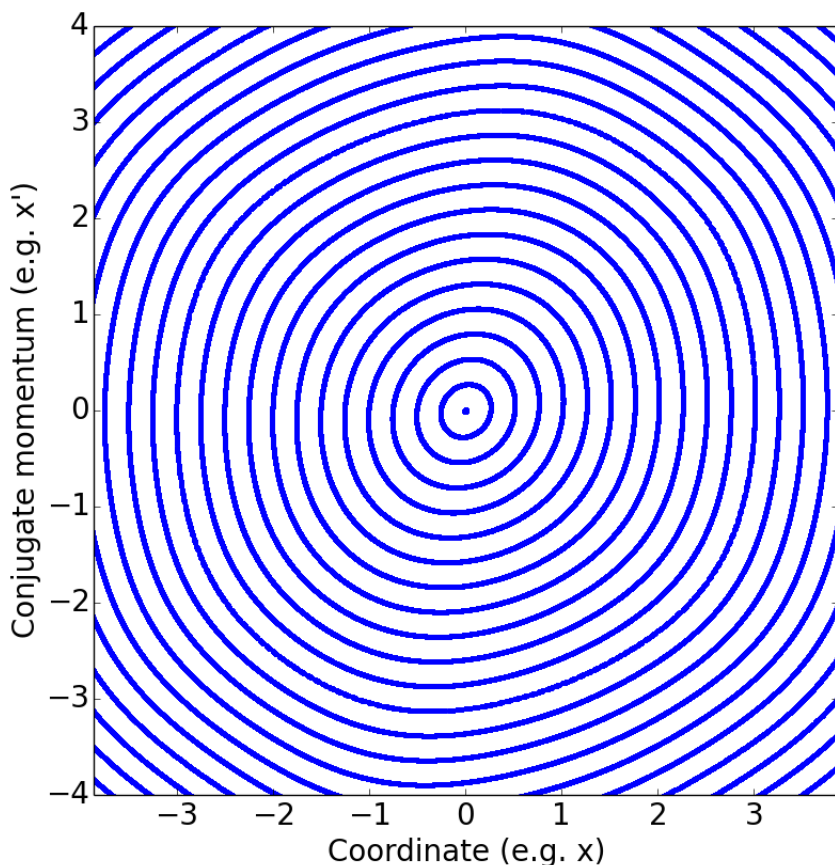
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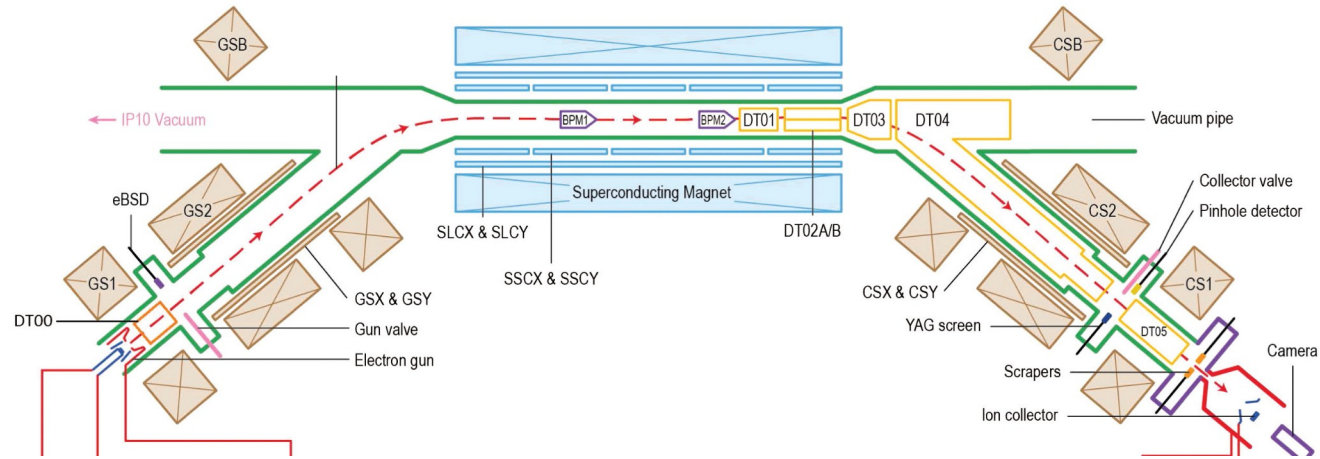
→ Maybe we should be inspired ?

The issue with non-linear forces



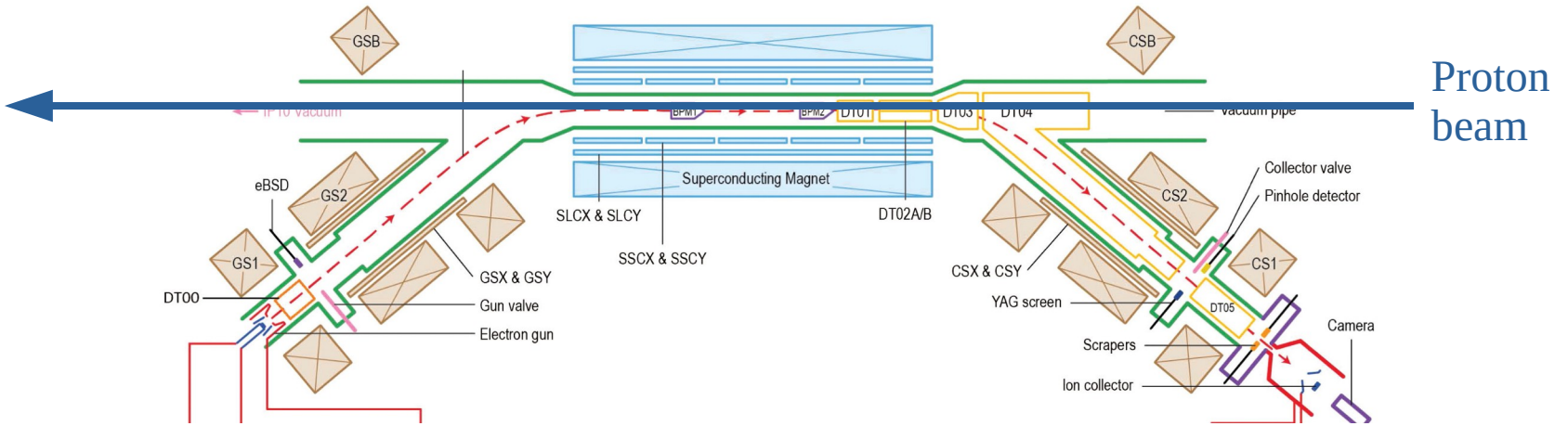
- Along with the tune spread required for Landau damping, non-linearities come with detrimental effect for the single particle trajectories:
 - Resonances, chaotic motion and eventually beam quality degradation (particles losses, emittance growth)
- The **amount of Landau** damping that can be obtained with octupoles is **limited** by their impact on beam losses

Electron lens

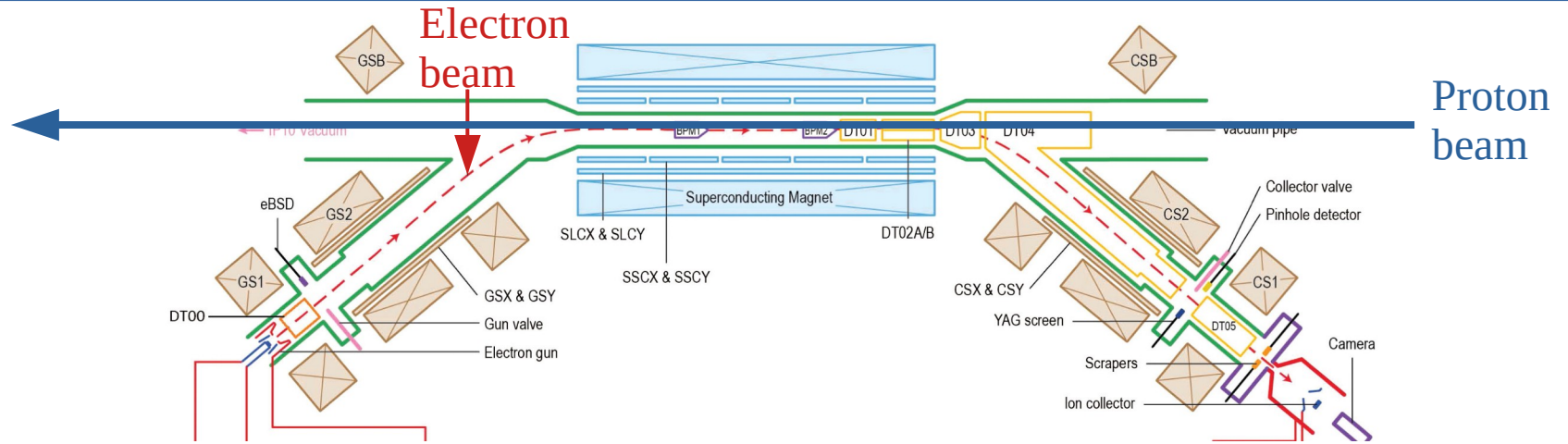


Electron lens

[RHIC,
elens]

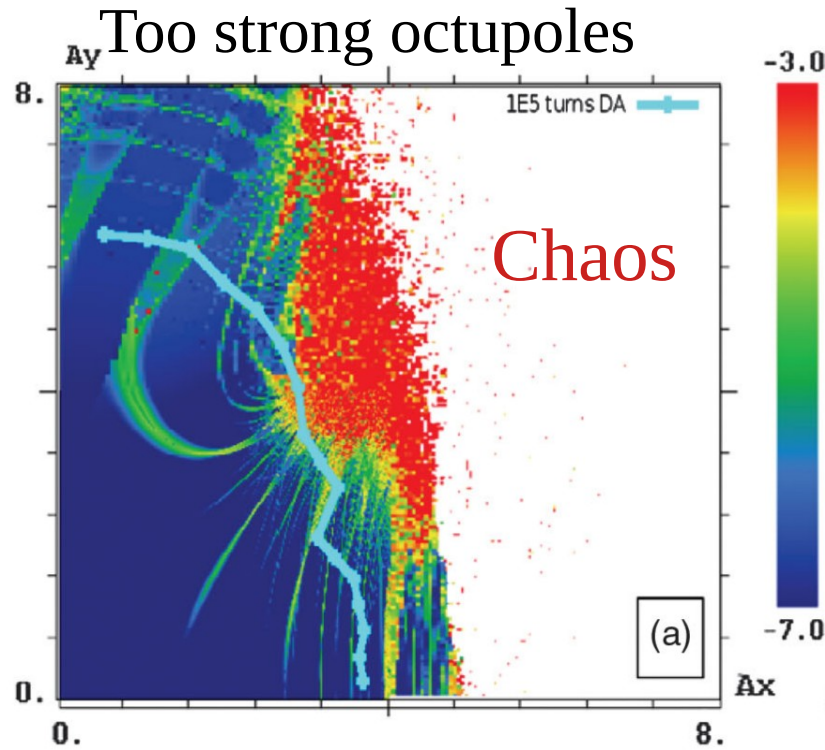
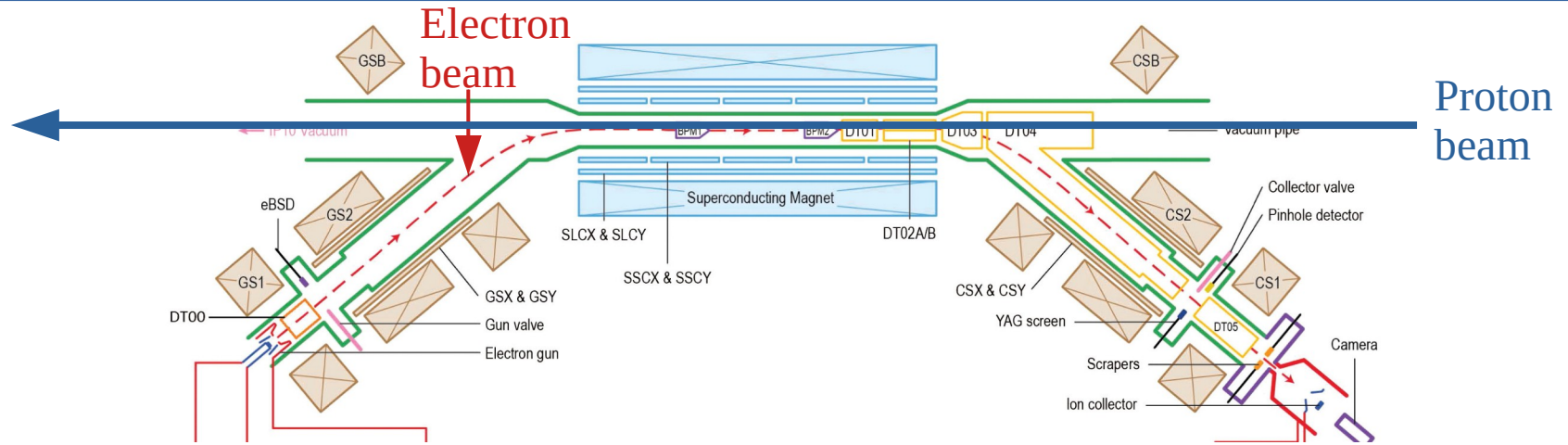


Electron lens

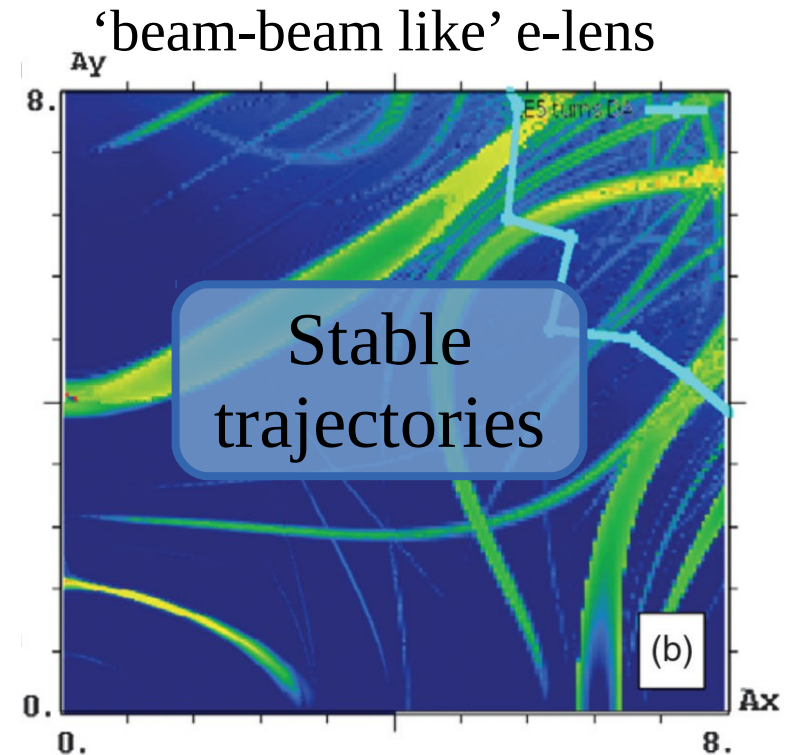
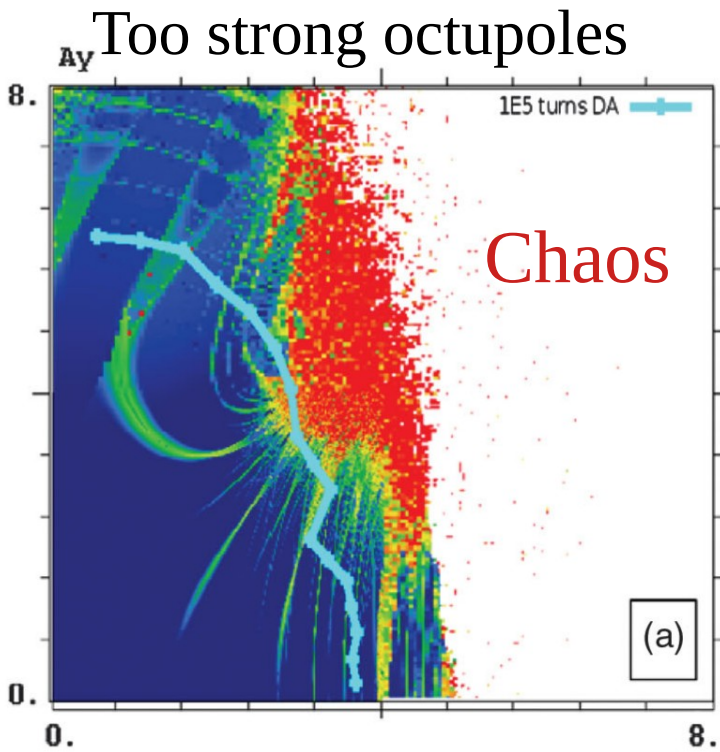
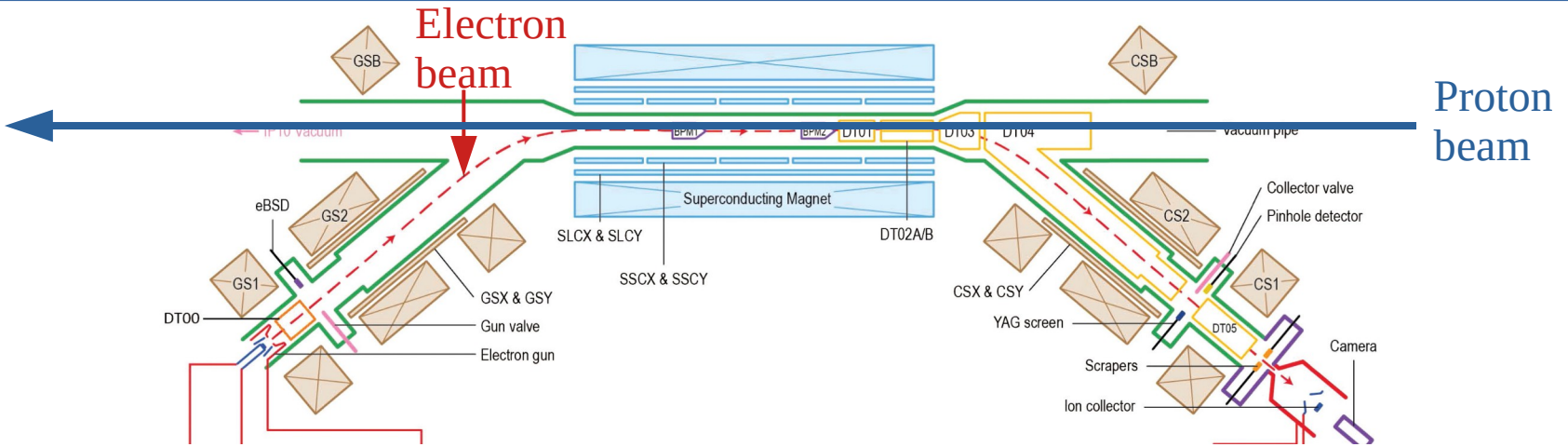


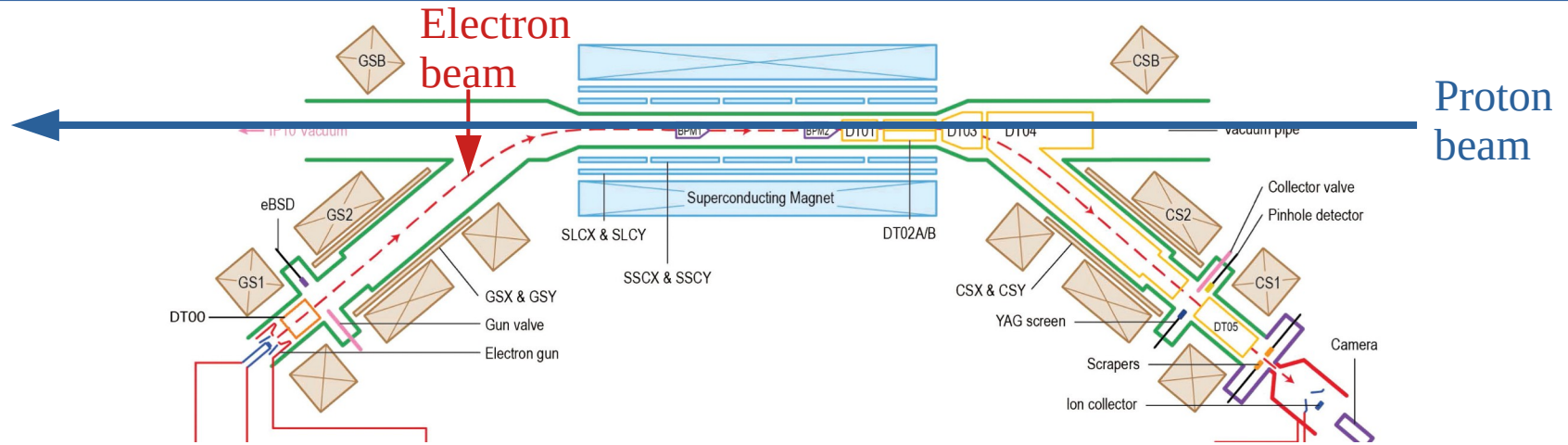
Electron lens

[RHIC,
elens]

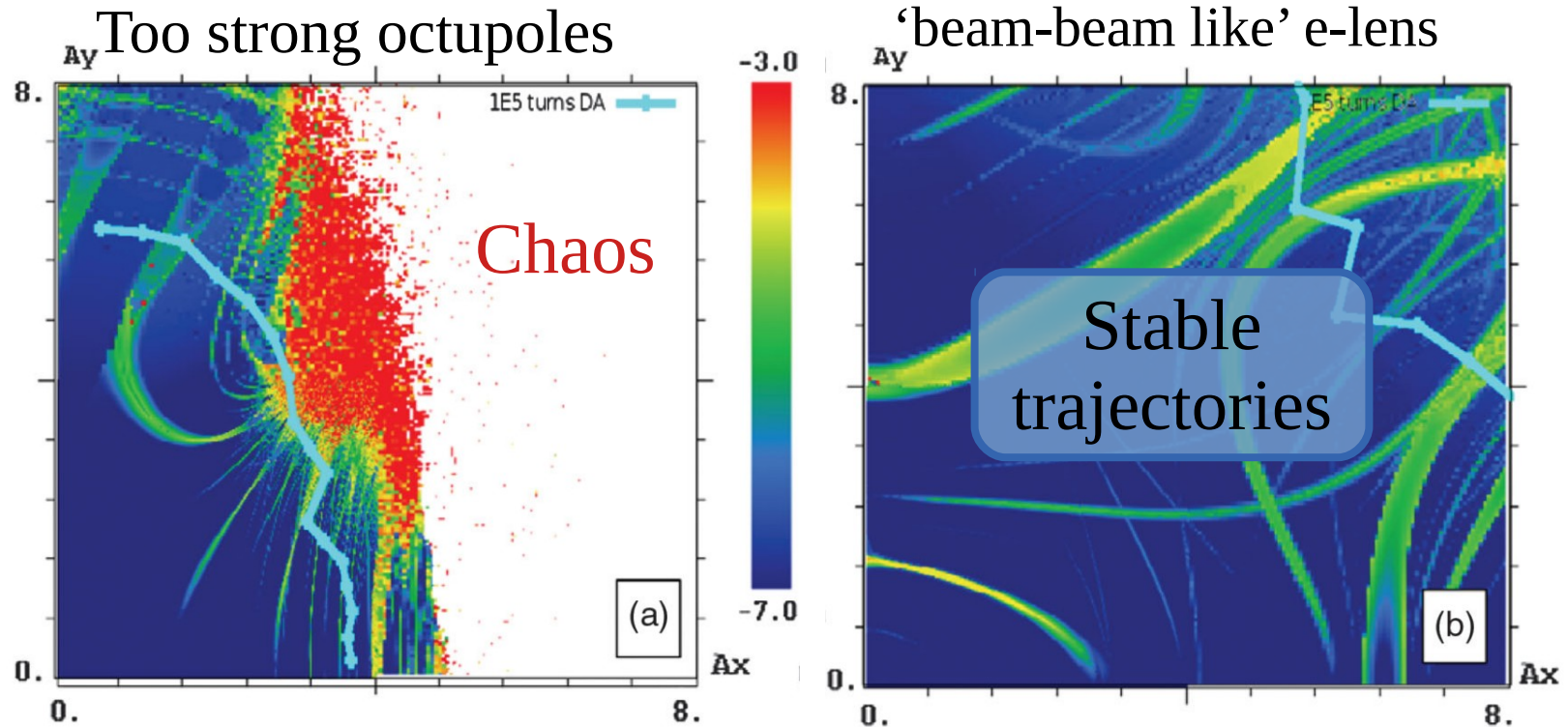


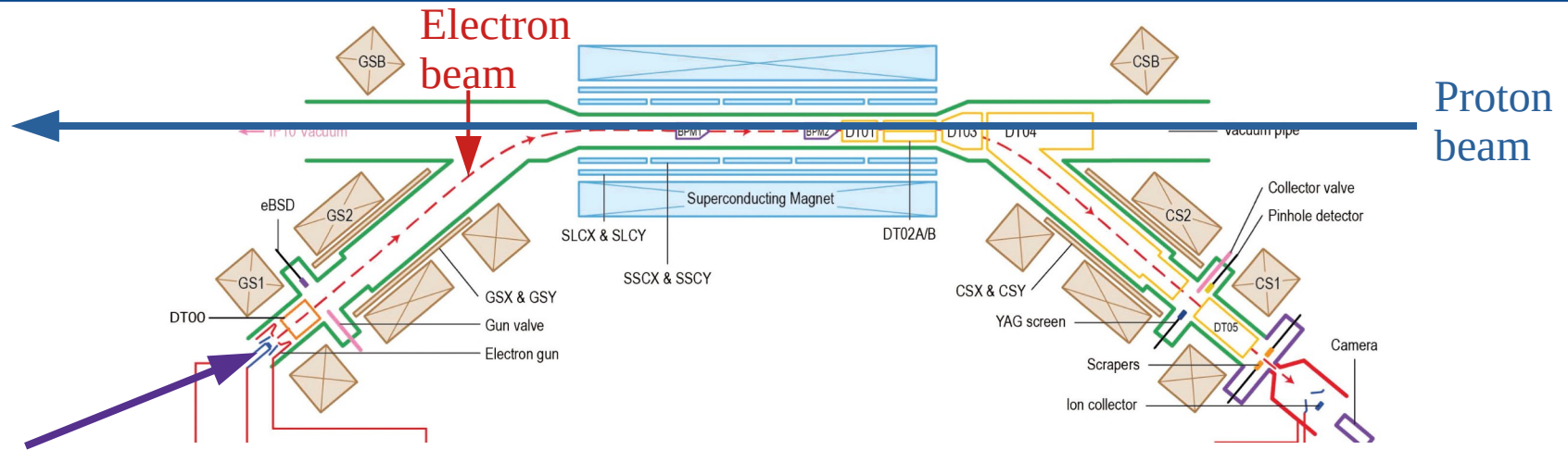
Electron lens



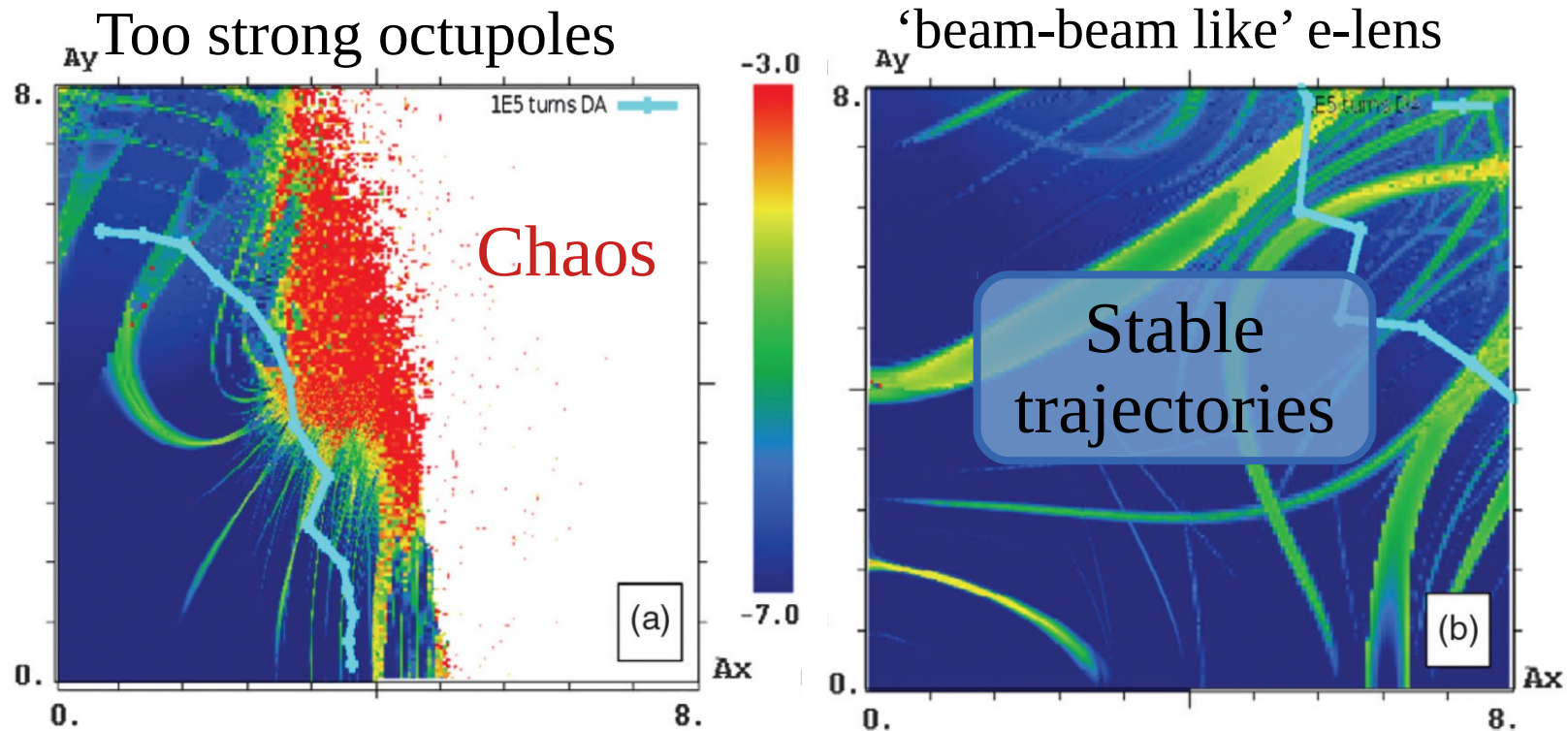


The gun design allows for various electron beam shapes

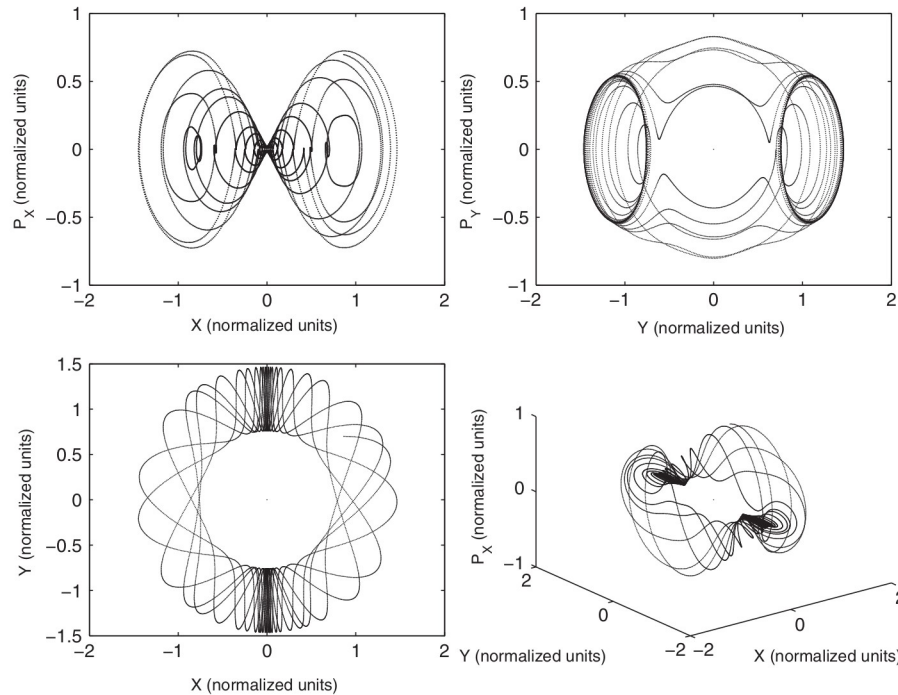




The gun design allows for various electron beam shapes
 → Optimise the force to maximise Landau damping with least impact on the beam quality

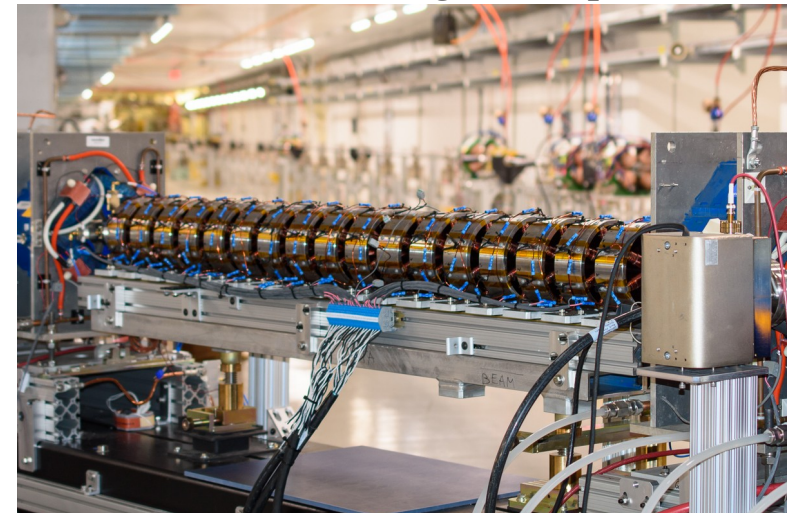
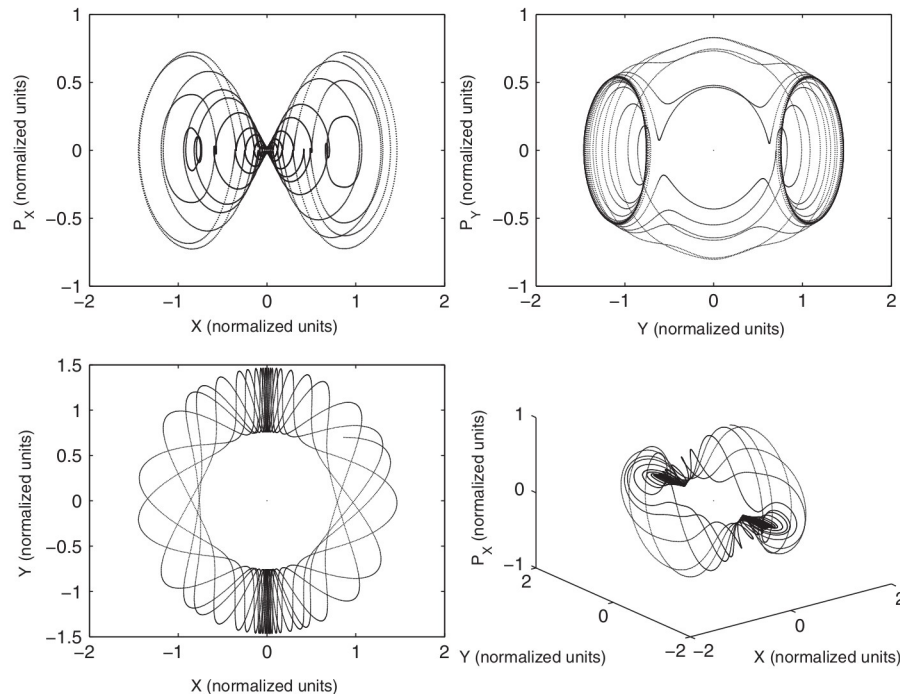


- It is possible to introduce ‘good’ non-linearities that generate a tune spread yet maintaining some invariants of motion
 - Possibly strong Landau damping without deterioration of the beam quality



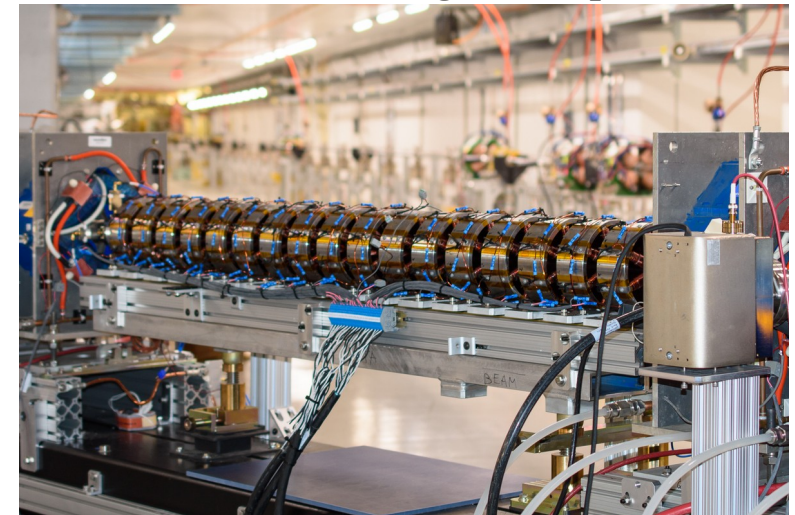
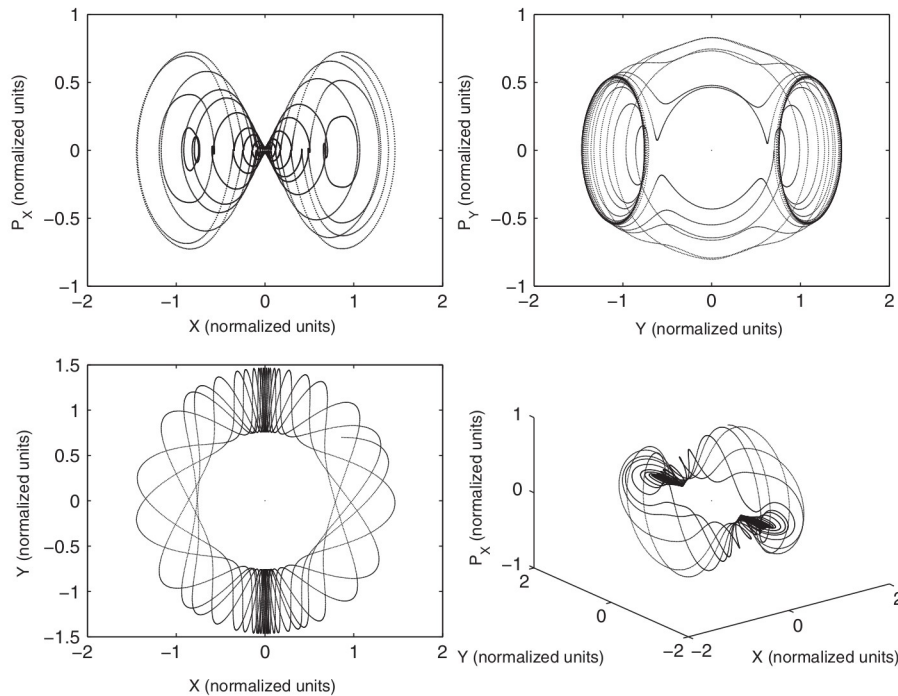
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A series of independently powered octupoles to generate a non-linear integrable optics at IOTA

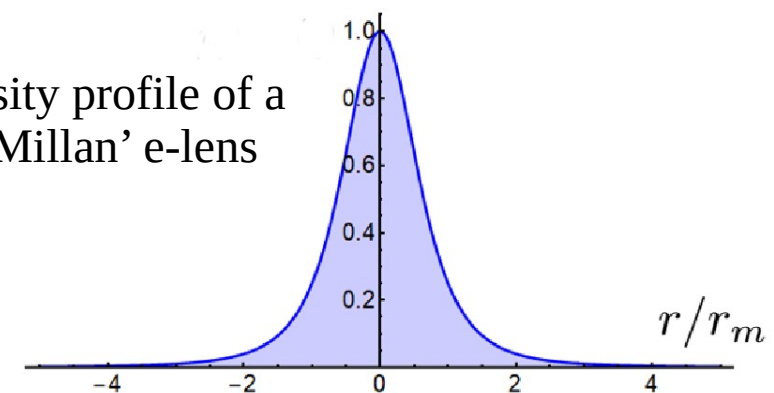


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Density profile of a ‘McMillan’ e-lens



$$\frac{\int dr r f_0(r) |H_l^k(r)|^2}{\Delta\Omega_{ext}^{l,k}} = \int dr \frac{r f_0(r) |H_l^k(r)|^2}{\Omega^{l,k} - \omega(r) - l\omega_s}$$

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Transverse frequency shift

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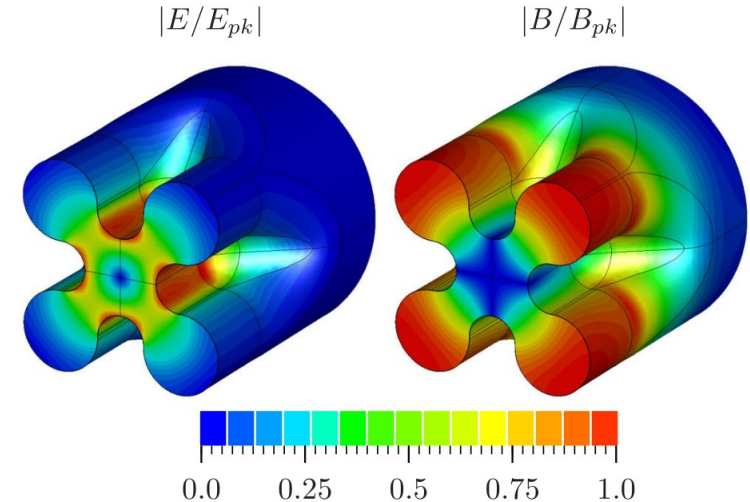
Longitudinal
oscillation
amplitude

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Transverse frequency shift

Longitudinal
oscillation
amplitude

- Transverse detuning with longitudinal amplitude can be achieved with
 - Dedicated optics (non-linear chromaticity)
 - RF quadrupoles



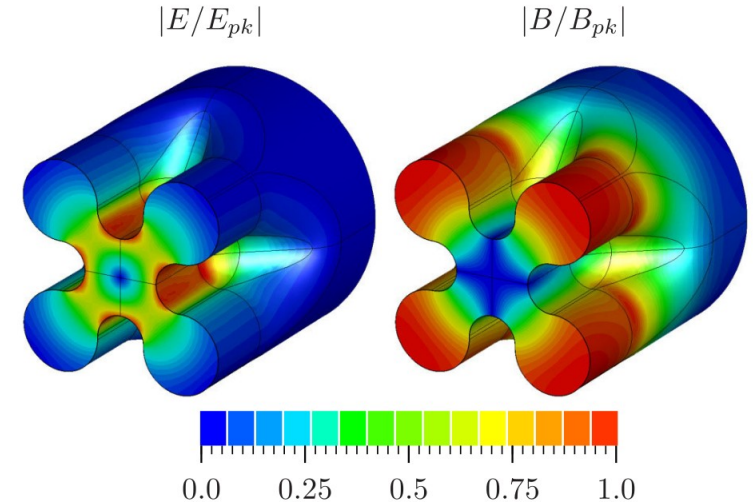
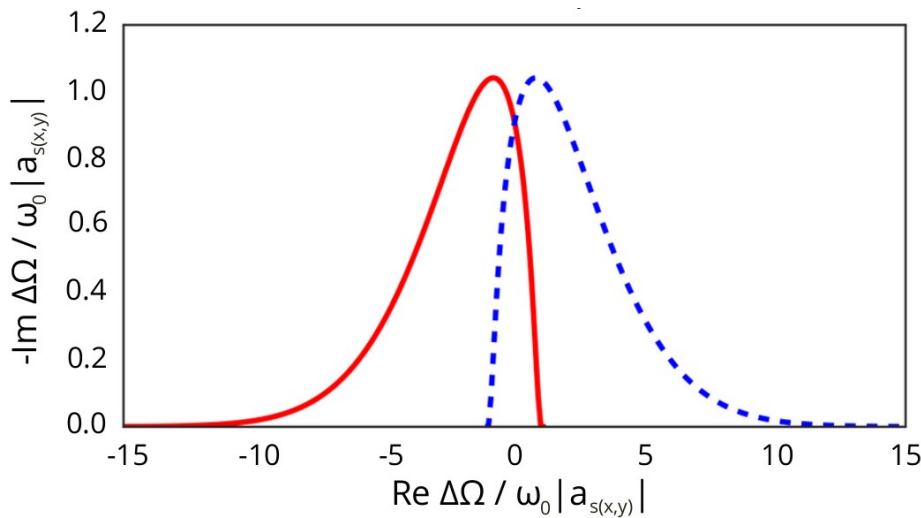
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Transverse frequency shift

Longitudinal oscillation amplitude

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Summary

- In some cases Landau damping arise naturally in accelerators
 - Momentum spread
 - Chromatic spread
 - Non-linearity of the longitudinal focusing (RF wave)
 - Non-linearity of collective forces (Space-charge, beam-beam)
 - **Watch out !** Due to their dynamic nature, the collective forces can lead to loss of Landau damping, by shifting the coherent mode frequencies w.r.t. the incoherent spectrum

Summary

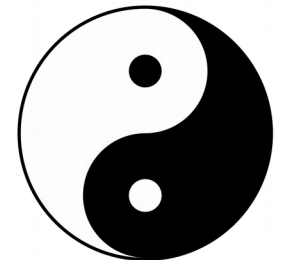
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 - More advanced tools (electron-lens, special magnets, RF quadrupoles)
- Several aspect of accelerator design are driven by the need for Landau damping (Beam parameters, optics, operation, ...)
- Landau damping is **beneficial** to maintain the beam quality, however the means to generate Landau damping can have a bad impact on the trajectories of single particles, leading to a **deterioration** of the beam quality



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