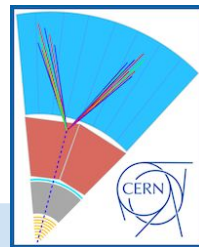


LLP2024 Workshop
Tokyo, 1-5 July 2024



Probing sterile neutrino magnetic moments with non-pointing photon searches @ LHC

based on [arXiv:2405.08877](https://arxiv.org/abs/2405.08877) with P. Bolton, F. Deppisch, C. Hati, M. Hirsch

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Motivation

Neutrino mass models often predict the existence of RH or **sterile neutrinos**.

Seesaw Type-I: $m_N, V_{\ell N} = f(m_\nu, m_N)$

Minimal scenario. Sterile neutrinos interact only through **neutrino mixing**:

$$\mathcal{L}_{\text{SM}+N_R} \supset \mathcal{L}_{N_R}^{\text{mass}} - \frac{g}{\sqrt{2}} V_{\ell N} (\bar{\ell} \gamma^\mu N_R^c) W_\mu - \frac{g}{2 \cos \theta_W} U_{\ell i} V_{\ell N}^* (\bar{N}_R^c \gamma^\mu \nu_{iL}) Z_\mu + \text{h.c.}$$

Independent parameters: $m_N, V_{\ell N}$

LLP searches @LHC have looked for sterile neutrinos in the minimal scenario.

Beyond the minimal scenario

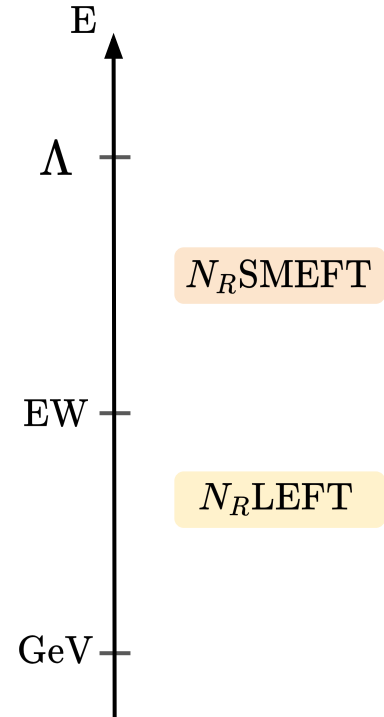
If sterile neutrinos couple to heavy new physics at Λ

$$\mathcal{L}_{\text{EFT}} = \mathcal{L}_{\text{SM}+N_R} + \sum_{d>4} C_i^{(d)} \mathcal{O}_i^{(d)} \quad C_i^{(d)} \propto \Lambda^{4-d}$$

Much more phenomenology to explore if we consider **new interactions** beyond the minimal scenario.



Neutrino magnetic moments



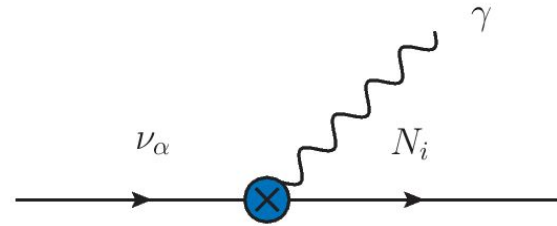
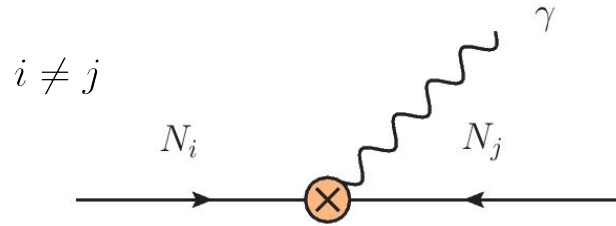
Sterile neutrino magnetic moments

Sterile neutrino magnetic moments

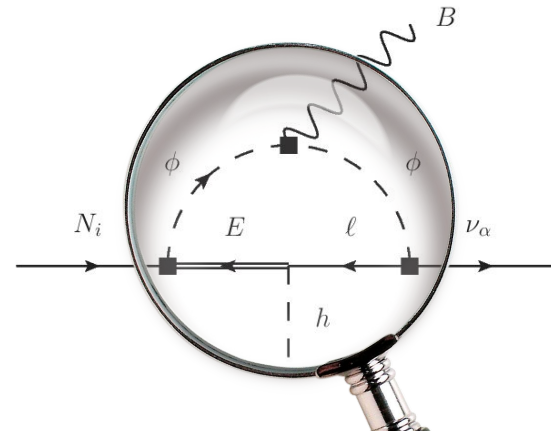
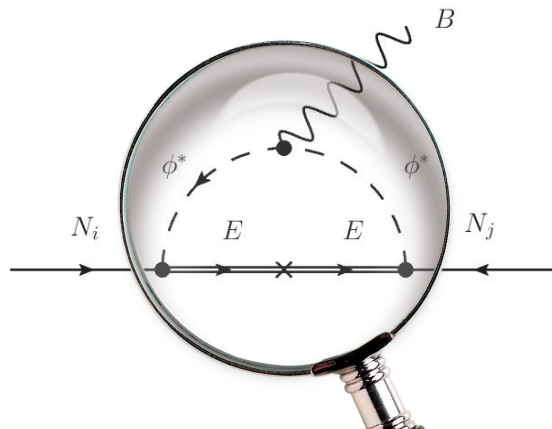
Sterile-to-sterile

Active-to-sterile

EFT description

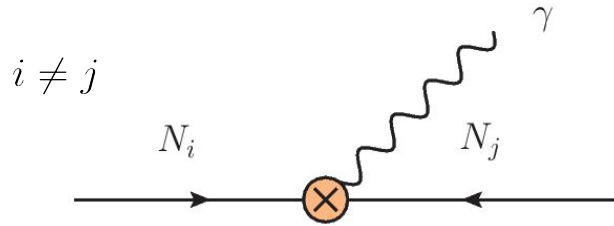


Specific UV model

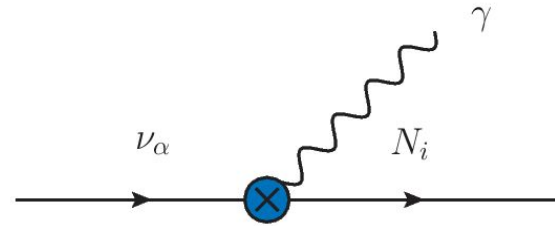


Effective description

Sterile-to-sterile



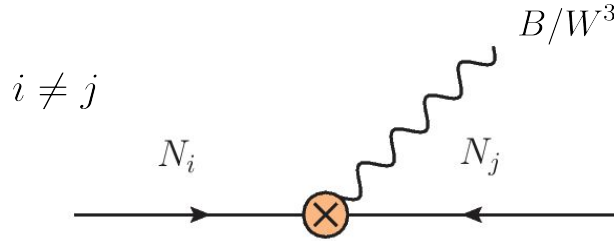
Active-to-sterile



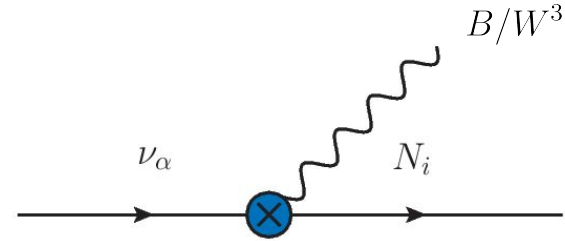
$$\mathcal{L}_{N\text{LEFT}} \supset d_{NN\gamma}^{ij} (\overline{N_{Ri}^c} \sigma^{\mu\nu} N_{Rj}) F_{\mu\nu} + d_{\nu N\gamma}^{\alpha i} (\overline{\nu_{L\alpha}} \sigma^{\mu\nu} N_{Ri}) F_{\mu\nu} + \text{h.c.}$$

Effective description

Sterile-to-sterile



Active-to-sterile



$$\mathcal{L}_{N\text{LEFT}} \supset d_{NN\gamma}^{ij} (\overline{N_{Ri}^c} \sigma^{\mu\nu} N_{Rj}) F_{\mu\nu} + d_{\nu N\gamma}^{\alpha i} (\overline{\nu_{L\alpha}} \sigma^{\mu\nu} N_{Ri}) F_{\mu\nu} + \text{h.c.}$$

Above EW scale



$$\begin{aligned} \mathcal{L}_{NSMEFT} \supset & C_{NNB}^{(5)ij} (\overline{N_{Ri}^c} \sigma^{\mu\nu} N_{Rj}) B_{\mu\nu} + \text{h.c.} \\ & + C_{NB}^{(6)\alpha i} (\overline{L_\alpha} \sigma^{\mu\nu} N_{Ri}) \tilde{H} B_{\mu\nu} + C_{NW}^{(6)\alpha i} (\overline{L_\alpha} \sigma^{\mu\nu} N_{Ri}) \tau^I \tilde{H} W_{\mu\nu}^I + \text{h.c.} \end{aligned}$$

Electroweak dipole moments

In NSMEFT, it is convenient to use rotated basis with electroweak gauge boson operators:

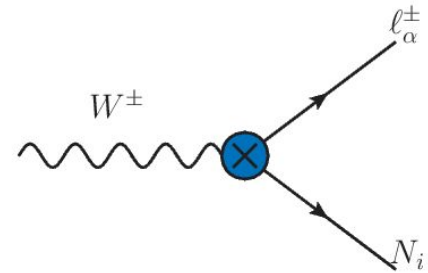
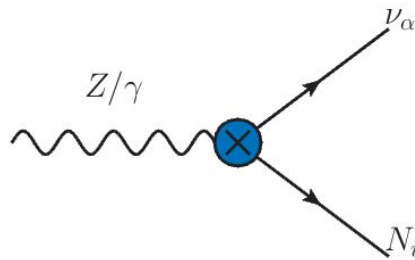
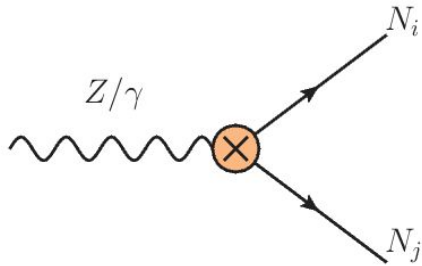
$$\mathcal{O}_{NN\gamma} = (\overline{N_R^c} \sigma_{\mu\nu} N_R) F^{\mu\nu}$$

$$\mathcal{O}_{\nu N\gamma} = (\overline{\nu_L} \sigma_{\mu\nu} N_R) F^{\mu\nu}$$

$$\mathcal{O}_{NNZ} = (\overline{N_R^c} \sigma_{\mu\nu} N_R) Z^{\mu\nu}$$

$$\mathcal{O}_{\nu NZ} = (\overline{\nu_L} \sigma_{\mu\nu} N_R) Z^{\mu\nu}$$

$$\mathcal{O}_{\ell NW} = (\overline{\ell_L} \sigma_{\mu\nu} N_R) W^{\mu\nu}$$



$$C_{NNB}^{(5)} \Leftrightarrow d_{NN\gamma}, d_{NNZ}$$

$$C_{NB}^{(6)}, C_{NW}^{(6)} \Leftrightarrow d_{\nu N\gamma}, d_{\nu NZ}, d_{\ell NW}$$

Dictionary

$$C_{NNB}^{(5)} \Leftrightarrow d_{NN\gamma}, d_{NNZ}$$

$$C_{NB}^{(6)}, C_{NW}^{(6)} \Leftrightarrow d_{\nu N\gamma}, d_{\nu NZ}, d_{\ell NW}$$

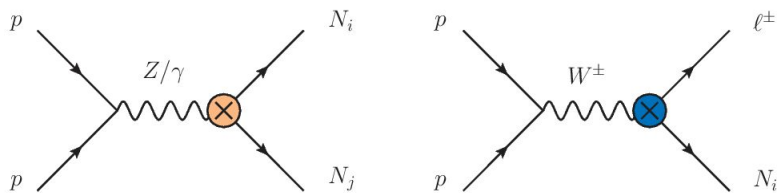
$$d_{NN\gamma}^{ij} = c_w C_{NNB}^{(5)ij}, \quad d_{NNZ}^{ij} = -s_w C_{NNB}^{(5)ij},$$

$$d_{\nu N\gamma}^{\alpha i} = \frac{v}{\sqrt{2}} \left(c_w C_{NB}^{(6)\alpha i} + \frac{s_w}{2} C_{NW}^{(6)\alpha i} \right), \quad d_{\nu NZ}^{\alpha i} = \frac{v}{\sqrt{2}} \left(-s_w C_{NB}^{(6)\alpha i} + \frac{c_w}{2} C_{NW}^{(6)\alpha i} \right)$$

$$d_{\ell NW}^{\alpha i} = \frac{v}{2} C_{NW}^{(6)\alpha i}$$

Phenomenology @ LHC

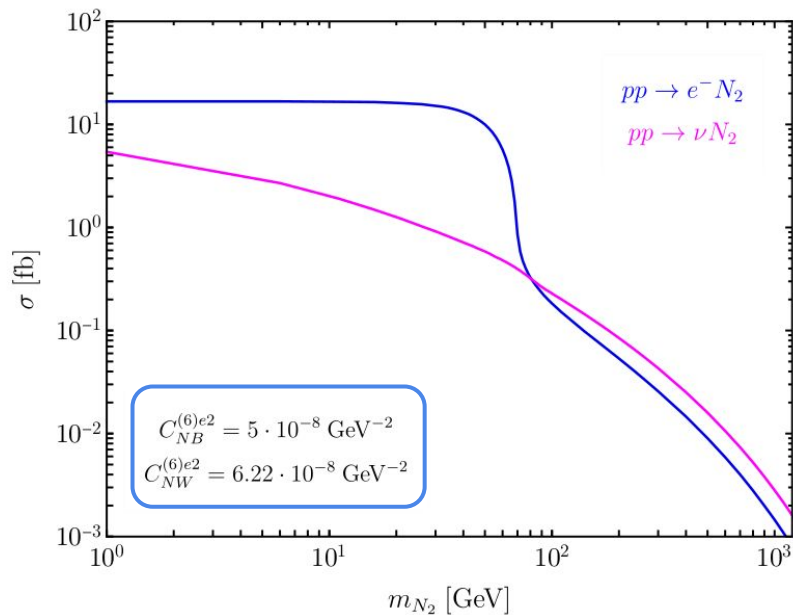
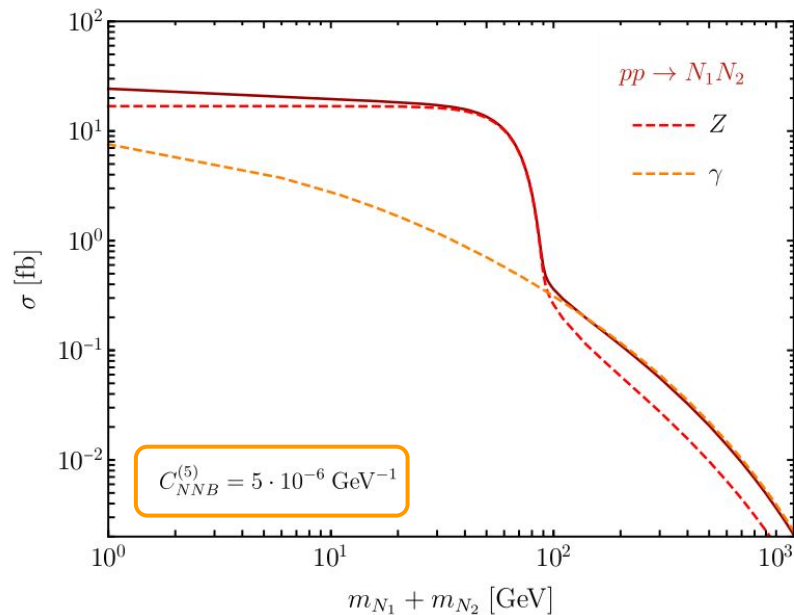
Sterile neutrino production



$$\sigma(pp \rightarrow N_1 N_2) \propto |C_{NNB}^{(5)}|^2$$

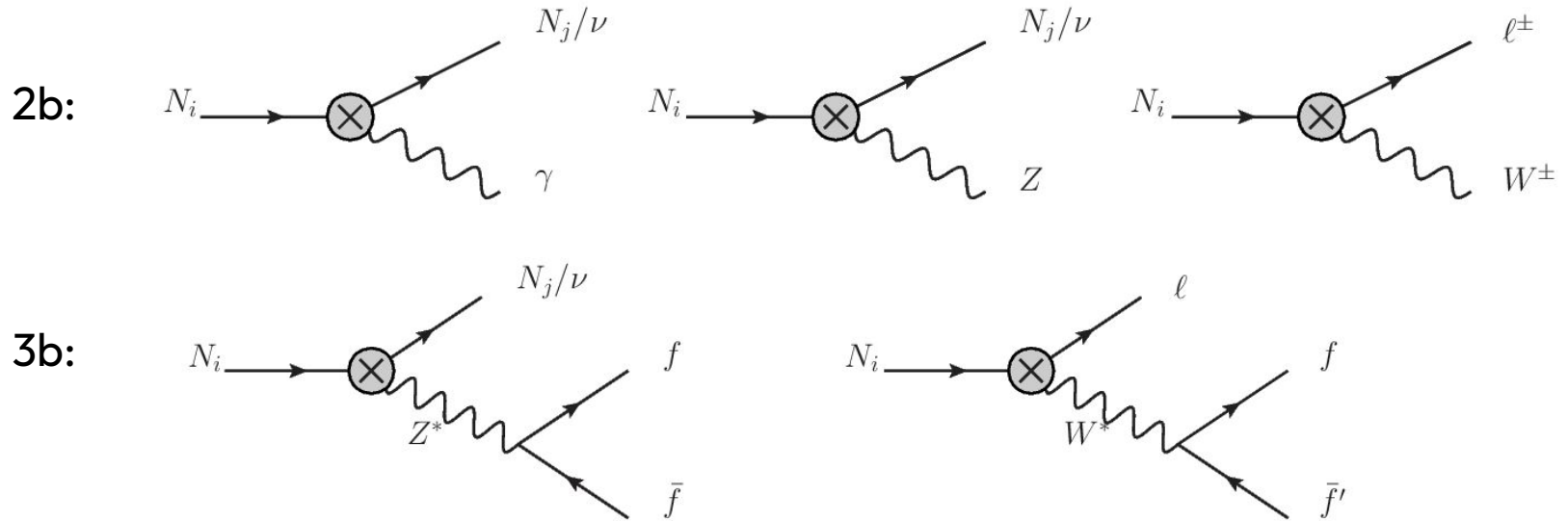
$$\sigma(pp \rightarrow \ell N_i) \propto |C_{NW}^{(6)}|^2$$

$\sqrt{s} = 14 \text{ TeV}$



Sterile neutrino decay

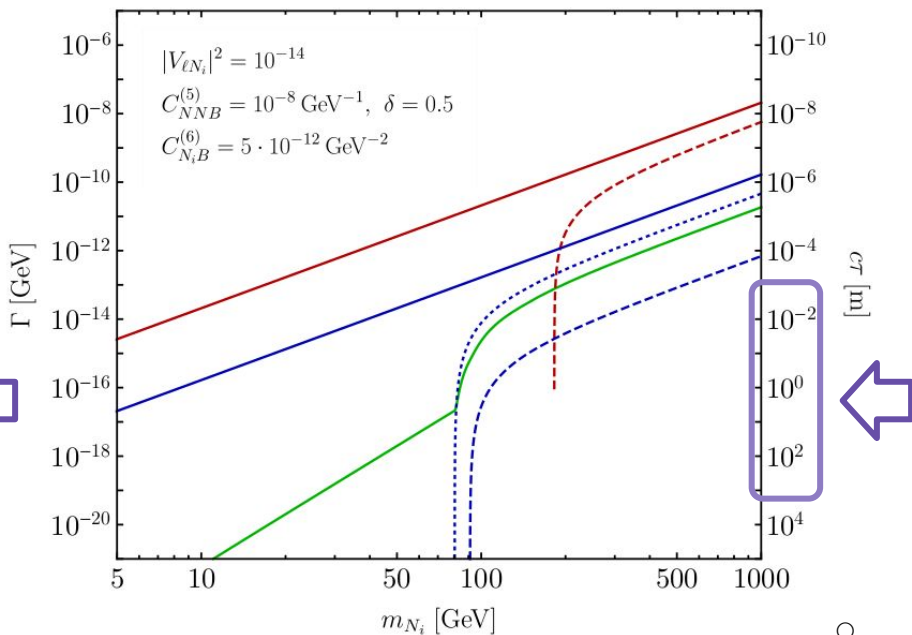
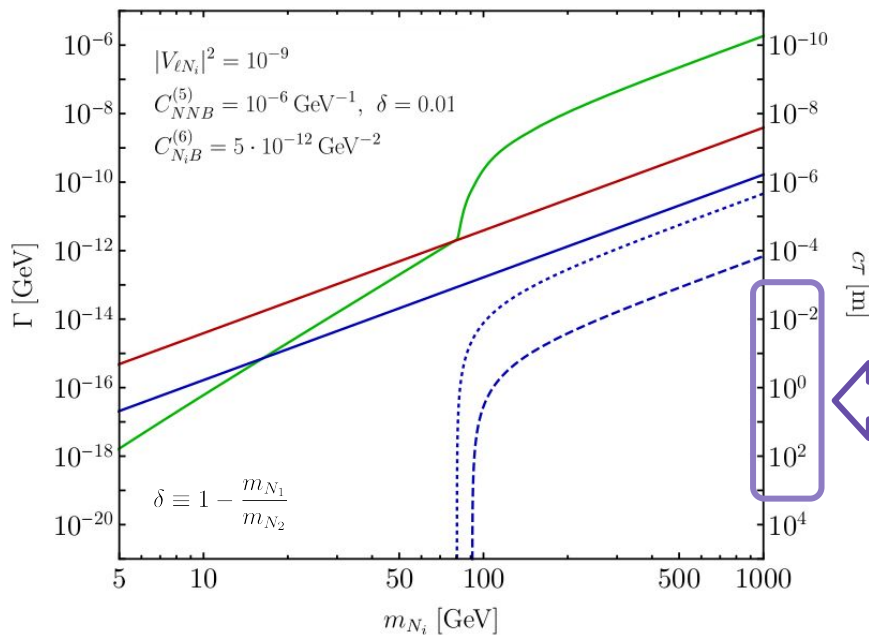
Sterile neutrinos can undergo 2-body or 3-body decays, depending on available phase space



Sterile neutrino decay

$$\left(m_{N_1}, m_{N_2}, C_{NNB}^{(5)}, C_{NB}^{(6)}, C_{NW}^{(6)}, V_{\ell N} \right)$$

- N_i mixing decay
- $N_i \rightarrow \nu + \gamma$
- $N_2 \rightarrow N_1 + \gamma$
- - - $N_i \rightarrow \nu + Z$
- - - $N_2 \rightarrow N_1 + Z$
- ⋯ $N_i \rightarrow e^- + W^-$



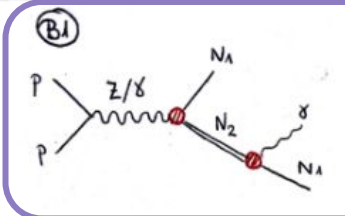
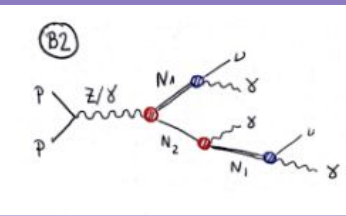
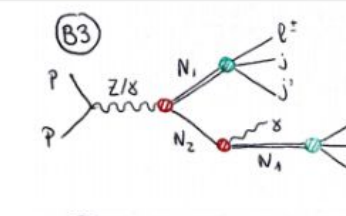
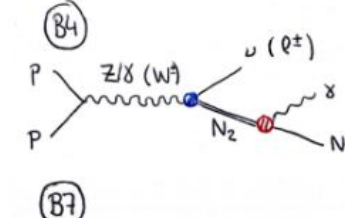
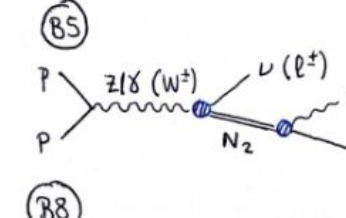
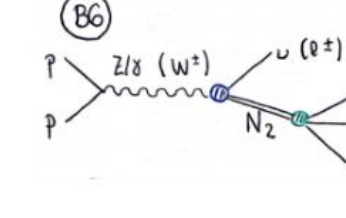
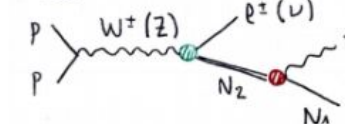
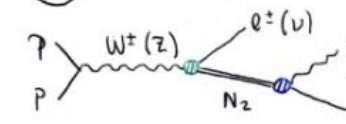
Possible scenarios

Each coupling can dominate production and decay (9 benchmarks)

Dec. \ Prod.	$C_{NNB}^{(5)}$	$C_{NB}^{(6)}, C_{NW}^{(6)}$	V_{eN}
$C_{NNB}^{(5)}$	<p>(B1)</p>	<p>(B2)</p>	<p>(B3)</p>
$C_{NB}^{(6)}, C_{NW}^{(6)}$	<p>(B4)</p>	<p>(B5)</p>	<p>(B6)</p>
V_{eN}	<p>(B7)</p>	<p>(B8)</p>	<p>Minimal scenario</p>

Possible scenarios

Each coupling can dominate production and decay (9 benchmarks)

Dec.	$C_{NNB}^{(5)}$	$C_{NB}^{(6)}, C_{NW}^{(6)}$	V_{eN}
$C_{NNB}^{(5)}$	 <p>(B1)</p>	 <p>(B2)</p>	 <p>(B3)</p>
$C_{NB}^{(6)}, C_{NW}^{(6)}$	 <p>(B4)</p>	 <p>(B5)</p>	 <p>(B6)</p>
V_{eN}	 <p>(B7)</p>	 <p>(B8)</p>	Minimal scenario



Non-pointing photon search

LLP signature: non-pointing photons

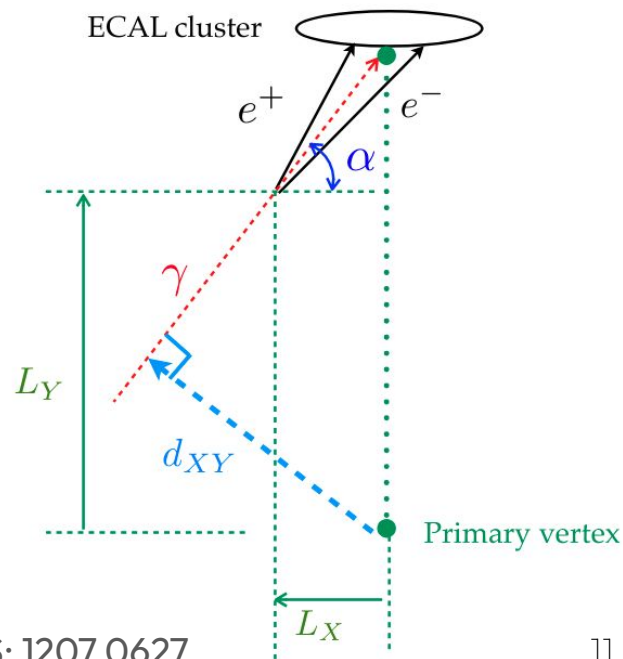
Non-pointing photons are emitted in the decay of LLPs:

- occur at a secondary vertex, displaced from PV
- leave a signal at ECAL
- displacement characterized by **Impact Parameter**

Minimal distance from photon trajectory to PV

$$d_{XY} = -L_X \cdot \sin \alpha + L_Y \cdot \cos \alpha$$
$$d_Z = L_Z - \frac{L_X \cdot p_X + L_Y \cdot p_Y}{p_T} \cdot \frac{p_Z}{p_T}$$

❗ requires knowing precisely PV location



Simulation process

We perform numerical study to estimate ATLAS sensitivity reach:

1. Model implemented in FeynRules
2. pp collisions at $\sqrt{s} = 14$ TeV in MadGraph5, 100k events at each grid point
3. Particle decays handled in MadSpin
4. Detection efficiencies obtained with Pythia8

↓
($m_{\text{LLP}}, C_{\text{decay}}$)

Scenario	Model parameters		Simulated decay
	Scan	Fixed	
B1	$m_{N_2}, C_{NNB}^{(5)}$	δ	$N_2 \rightarrow N_1 \gamma$
B2	$m_{N_1}, C_{N_1 X}^{(6)}$	$m_{N_2}, C_{NNB}^{(5)}$	$N_2 \rightarrow N_1 \gamma$ $N_1 \rightarrow \nu \gamma$
B3	$m_{N_1}, V_{eN_1} ^2$	$m_{N_2}, C_{NNB}^{(5)}$	$N_2 \rightarrow N_1 \gamma$ $N_1 \rightarrow e j j$

Efficiencies: selection cuts

Scenario	Signature	Selection cuts
B1	Non-pointing γ	$ p_T^\gamma > 10 \text{ GeV}$, $ \eta^\gamma < 2.47$
B2	Non-pointing γ ($\times 2$) (+ prompt γ)	$r_{\text{DV}} < 1450 \text{ mm}$, $ z_{\text{DV}} < 3450 \text{ mm}$ $ d_{XY}^\gamma > 6 \text{ mm}^*$
B3	Displaced Vertex ($\times 2$) (+ prompt γ)	$ p_T^e > 120 \text{ GeV}$, $ \eta^e < 2.47$ $4 \text{ mm} < r_{\text{DV}} < 300 \text{ mm}$, $ z_{\text{DV}} < 300 \text{ mm}$ 4 tracks with $ d_0 > 2 \text{ mm}$ $m_{\text{DV}} > 5 \text{ GeV}$

LLP decay before escaping ECAL

*original cut from CMS: 1207.0627
to reduce background

Number of signal events for each scenario:

$$N_{\text{sig.}}^{\text{B1}} = \sigma \cdot \mathcal{L} \cdot \mathcal{B}(N_2 \rightarrow N_1 \gamma) \cdot \epsilon_{\text{sel}}^{\text{B1}}$$

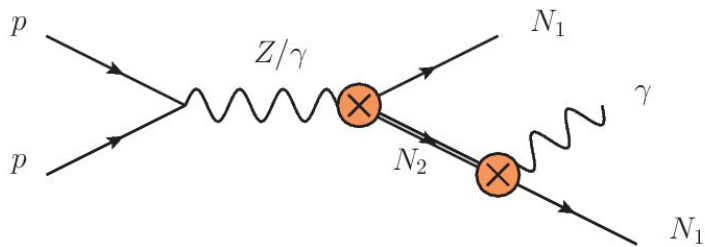
$$N_{\text{sig.}}^{\text{B2}} = \sigma \cdot \mathcal{L} \cdot \mathcal{B}(N_2 \rightarrow N_1 \gamma) \cdot 2 \cdot \mathcal{B}(N_1 \rightarrow \nu \gamma) \cdot \epsilon_{\text{sel}}^{\text{B2}}$$

$$N_{\text{sig.}}^{\text{B3}} = \sigma \cdot \mathcal{L} \cdot \mathcal{B}(N_2 \rightarrow N_1 \gamma) \cdot 2 \cdot \mathcal{B}(N_1 \rightarrow e j j) \cdot \epsilon_{\text{sel}}^{\text{B3}}$$

Sensitivity prospects: B1

$$pp \rightarrow N_1 N_2 : (d_{NN\gamma}, d_{NNZ})$$

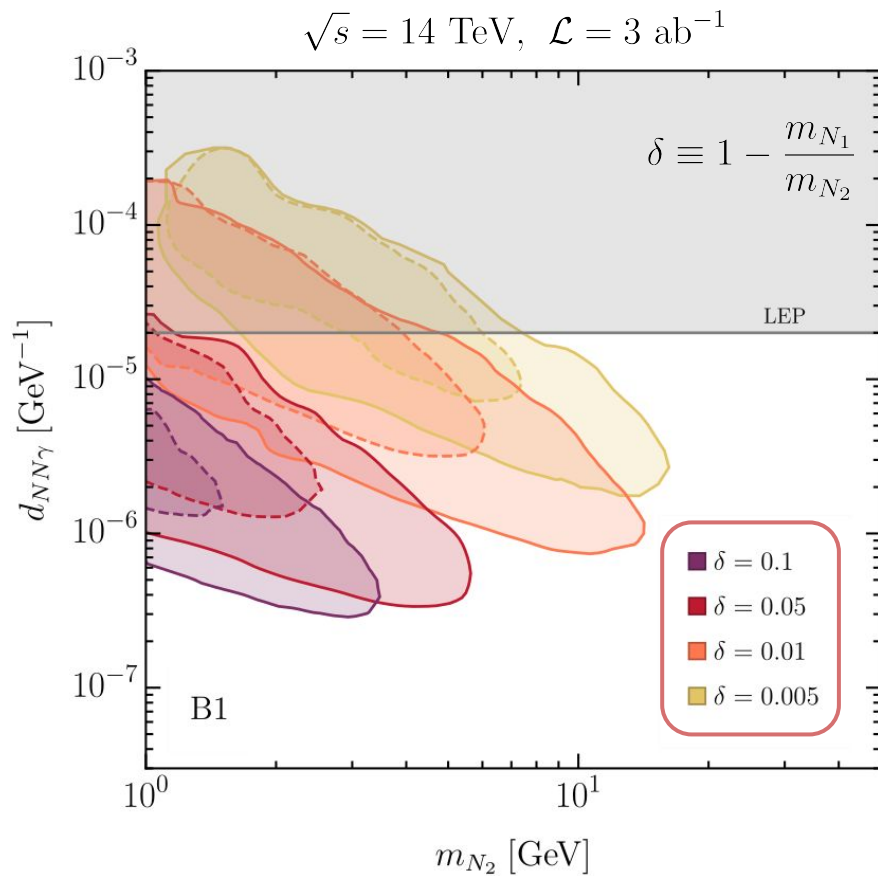
$$N_2 \rightarrow (N_1 \gamma)^{\text{LLP}} : (d_{NN\gamma}, \delta)$$



$$N_{\text{sig.}}^{\text{B1}} = \sigma \cdot \mathcal{L} \cdot \mathcal{B}(N_2 \rightarrow N_1 \gamma) \cdot \epsilon_{\text{sel}}^{\text{B1}}$$

solid: 3 events (95% C.L.)

dashed: 30 events

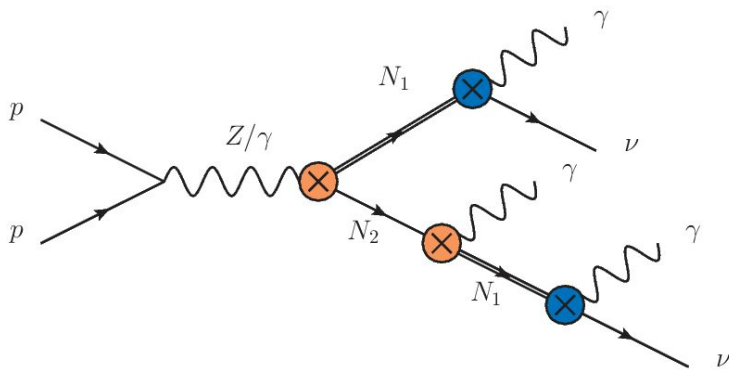


Sensitivity prospects: B2

$$pp \rightarrow N_1 N_2 : (d_{NN\gamma}, d_{NNZ})$$

$$N_1 \rightarrow (\nu\gamma)^{\text{LLP}} : (d_{\nu N\gamma})$$

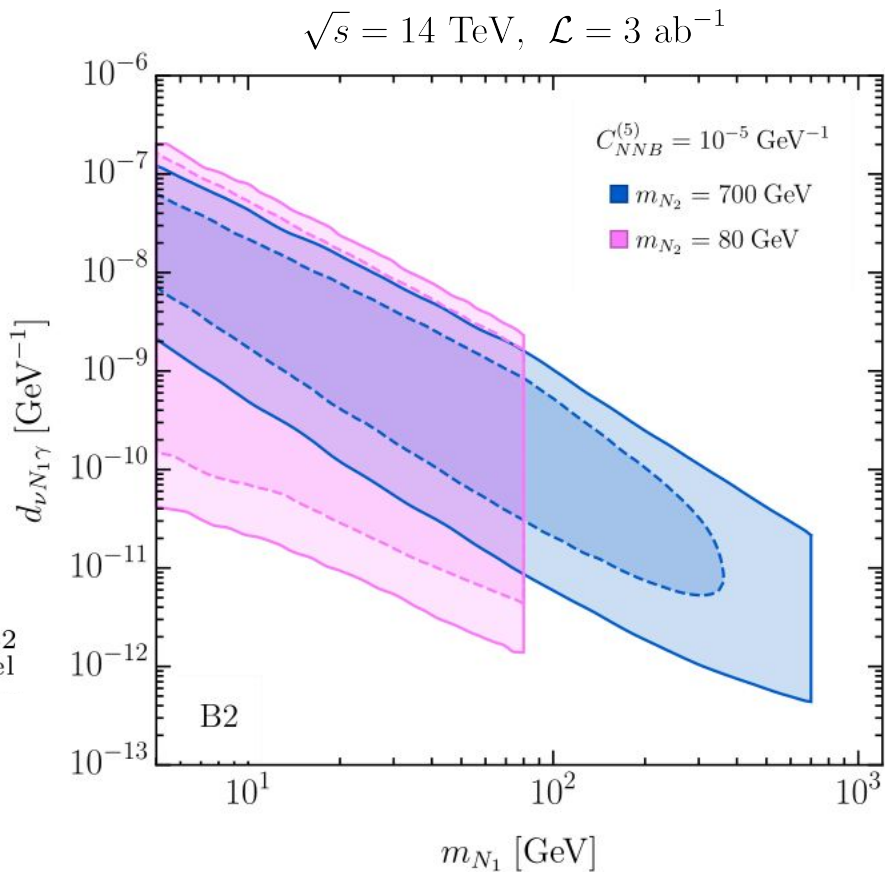
$$N_2 \rightarrow (N_1\gamma) \rightarrow (\nu\gamma)^{\text{LLP}} \gamma : (d_{\nu N\gamma})$$



$$N_{\text{sig.}}^{\text{B2}} = \sigma \cdot \mathcal{L} \cdot \mathcal{B}(N_2 \rightarrow N_1\gamma) \cdot 2 \cdot \mathcal{B}(N_1 \rightarrow \nu\gamma) \cdot \epsilon_{\text{sel}}^{\text{B2}}$$

solid: 3 events (95% C.L.)

dashed: 30 events



Conclusions

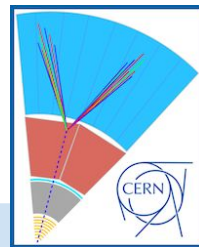
Rich phenomenology of heavy sterile neutrinos with **magnetic moments** @ LHC

↳ displaced vertices + non-pointing/delayed photons

LHC experiments can use these signatures to probe new sterile neutrino interactions (beyond minimal neutrino mixing)

↳ **Non-pointing photon searches** could be sensitive to new parts of parameter space for $d_{NN\gamma}, d_{\nu N\gamma}$

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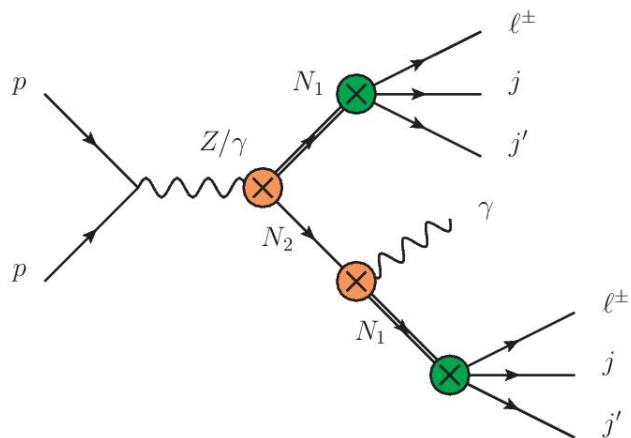


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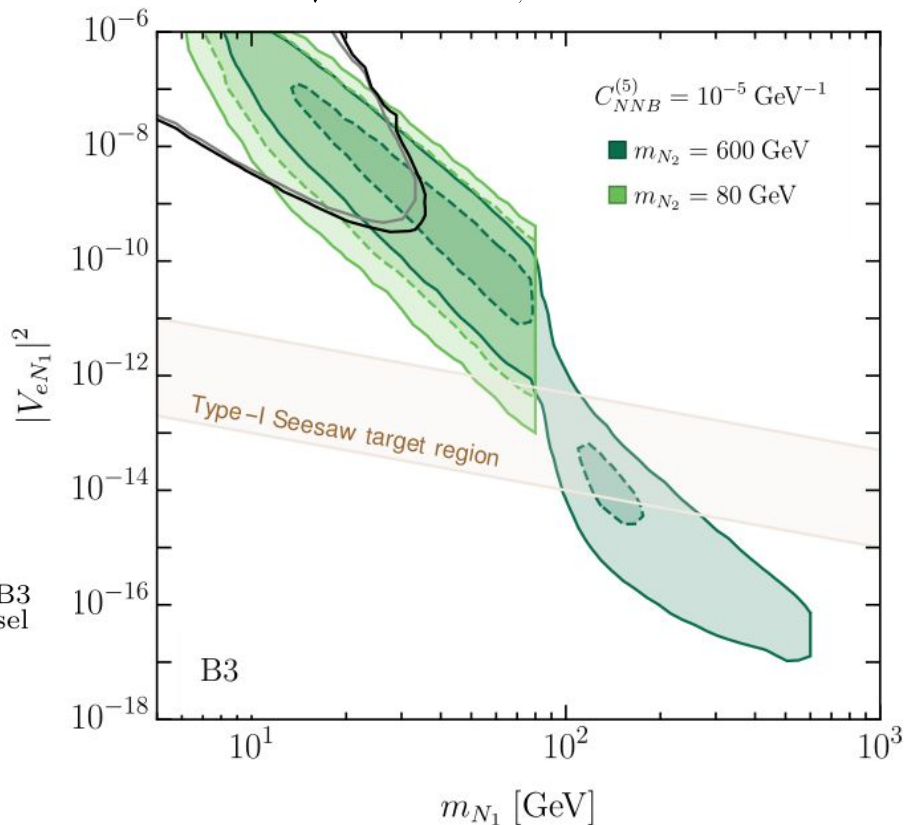
Sensitivity prospects: B3



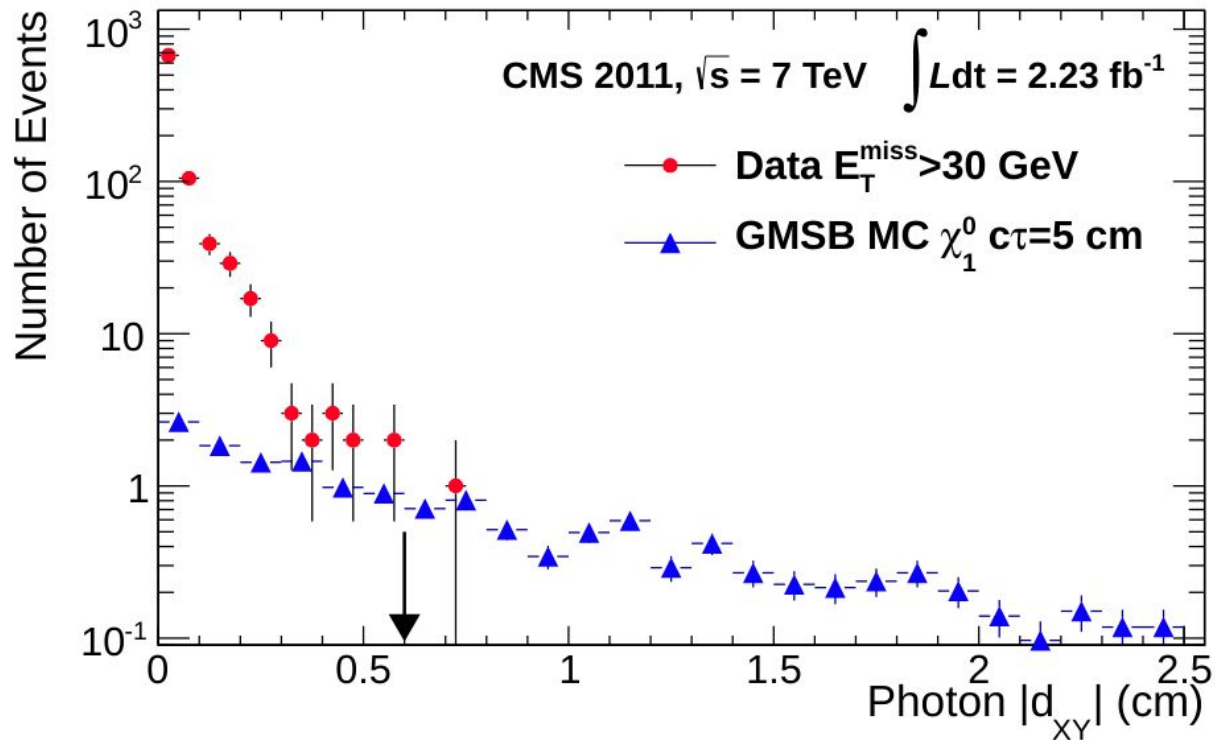
$$N_{\text{sig.}}^{\text{B3}} = \sigma \cdot \mathcal{L} \cdot \mathcal{B}(N_2 \rightarrow N_1 \gamma) \cdot 2 \cdot \mathcal{B}(N_1 \rightarrow e j j) \cdot \epsilon_{\text{sel}}^{\text{B3}}$$

solid: 3 events (95% C.L.)
dashed: 10 events

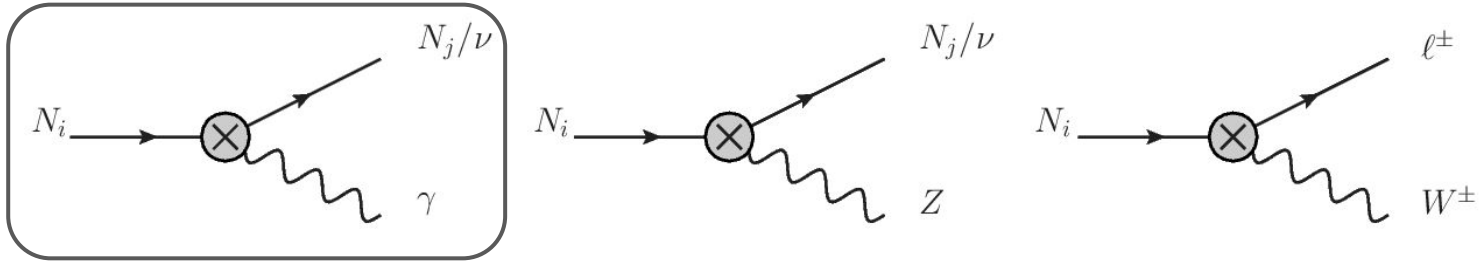
$\sqrt{s} = 14 \text{ TeV}, \mathcal{L} = 3 \text{ ab}^{-1}$



Non-pointing photon backgrounds



Sterile neutrino decay



$$\Gamma(N_2 \rightarrow N_1 \gamma) = \frac{2|d_{NN\gamma}^{12}|^2}{\pi} m_{N_2}^3 (2 - \delta)^3 \delta^3$$

$$\delta \equiv 1 - \frac{m_{N_1}}{m_{N_2}}$$

$$\Gamma(N_i \rightarrow \nu \gamma) = \frac{|d_{\nu N_i \gamma}|^2}{2\pi} m_{N_i}^3$$