

Dark photon pair production via off-shell dark Higgs at FASER

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arXiv:2406.17760

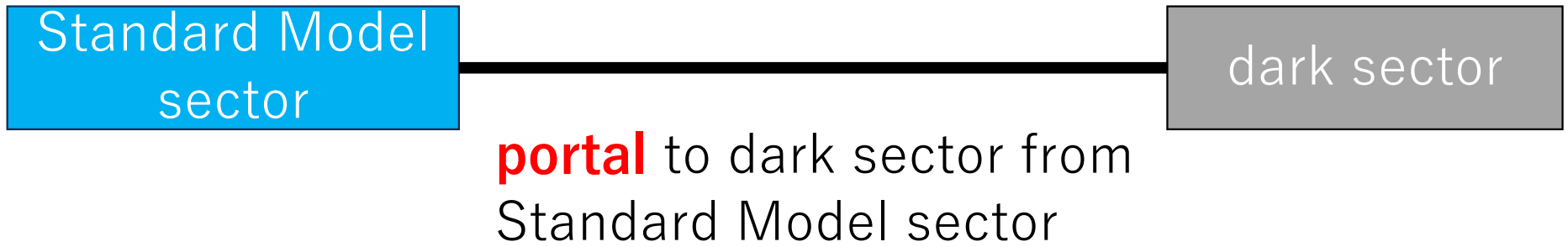
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14th workshop of the Long-Lived Particle Community

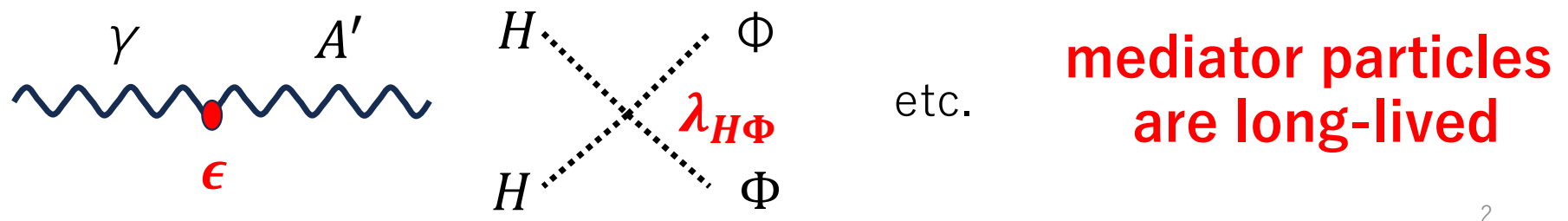
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Introduction

- Non-observation of dark matter implies dark sector idea.
- Portal mediates Standard Model sector and dark sector.



example : kinetic mixing between photon and dark photon,
mixing between Higgs and dark Higgs



origin of dark photon mass

origin of dark photon mass

- **spontaneously symmetry breaking of dark Higgs(our case)**
- Stuckelberg mechanism etc

dark sector Lagrangian

$$L_{DS} = -\frac{1}{4} F'_{\mu\nu} F'^{\mu\nu} + |D_\mu \Phi|^2 - V(\Phi), \quad F'_{\mu\nu} = \partial_\mu A'_\nu - \partial_\nu A'_\mu,$$

$$D_\mu = \partial_\mu - i g' A'_\mu, \quad \Phi = \langle \Phi \rangle + \phi/\sqrt{2}$$

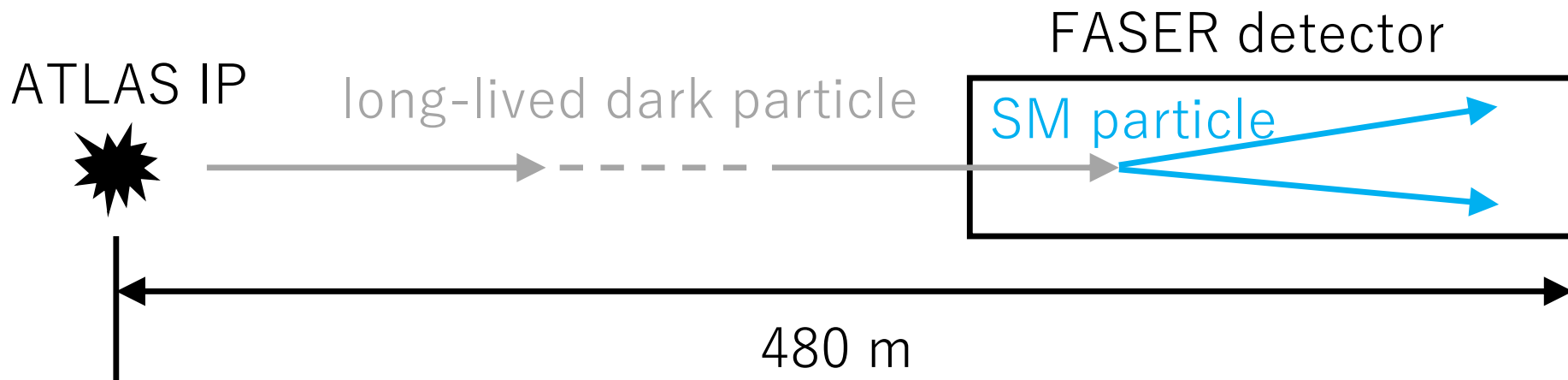
g' : extra U(1) gauge coupling, A' : dark photon,

ϕ : dark Higgs as origin of dark photon mass

FASER(started from 2022)

FASER detector is 480 meters away from ATLAS interaction point(IP) along with proton beam axis.

→ **FASER is suitable for long-lived particle search.**



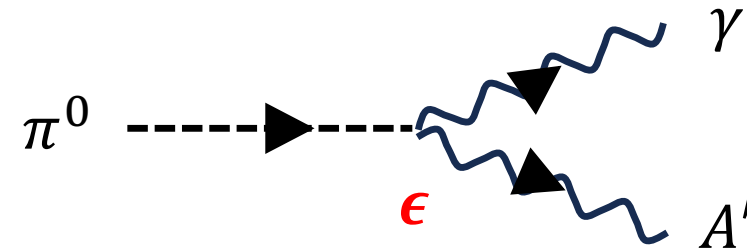
FASER is planned to upgrade to FASER2 at HL-LHC.

Dark photon and dark Higgs have been well studied at FASER.

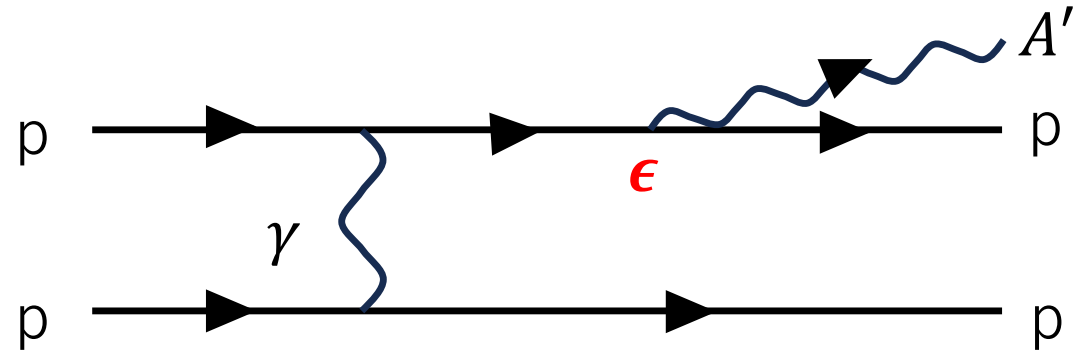
Felix Kling et al. Phys. Rev. D 97 (2018) 5, 055034, A. Ariga et al. Phys. Rev. D 99 (2019) 9, 095011

Main dark photon production process at FASER

- ① meson decay (mainly π^0)
 pp collision \rightarrow meson,
 meson decay \rightarrow dark photon

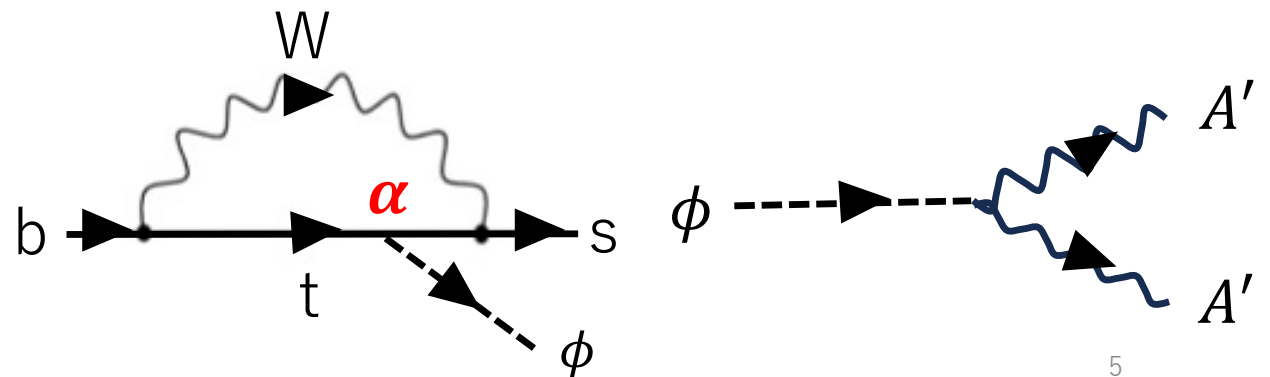


- ② dark bremsstrahlung

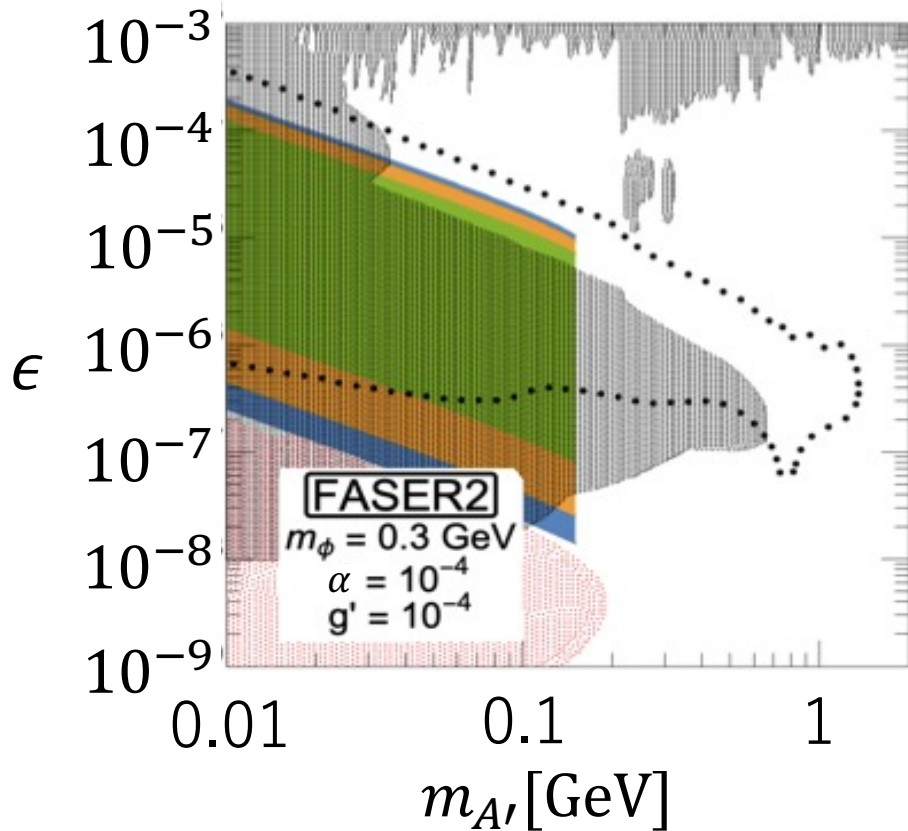


- ③ **dark Higgs decay**

- pp collision \rightarrow meson,
 meson decay (mainly B)
 \rightarrow dark Higgs,
 dark Higgs decay
 \rightarrow dark photon



Dark photon sensitivity at FASER2



dotted line
from meson decay and dark
bremsstrahlung(①, ②)

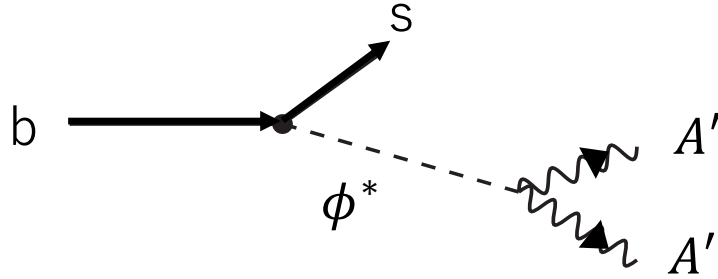
A. Ariga et al. Phys. Rev. D 99, 095011 (2019)

colored region 
from dark Higgs decay(③)

T. Araki et al. JHEP03(2021)072

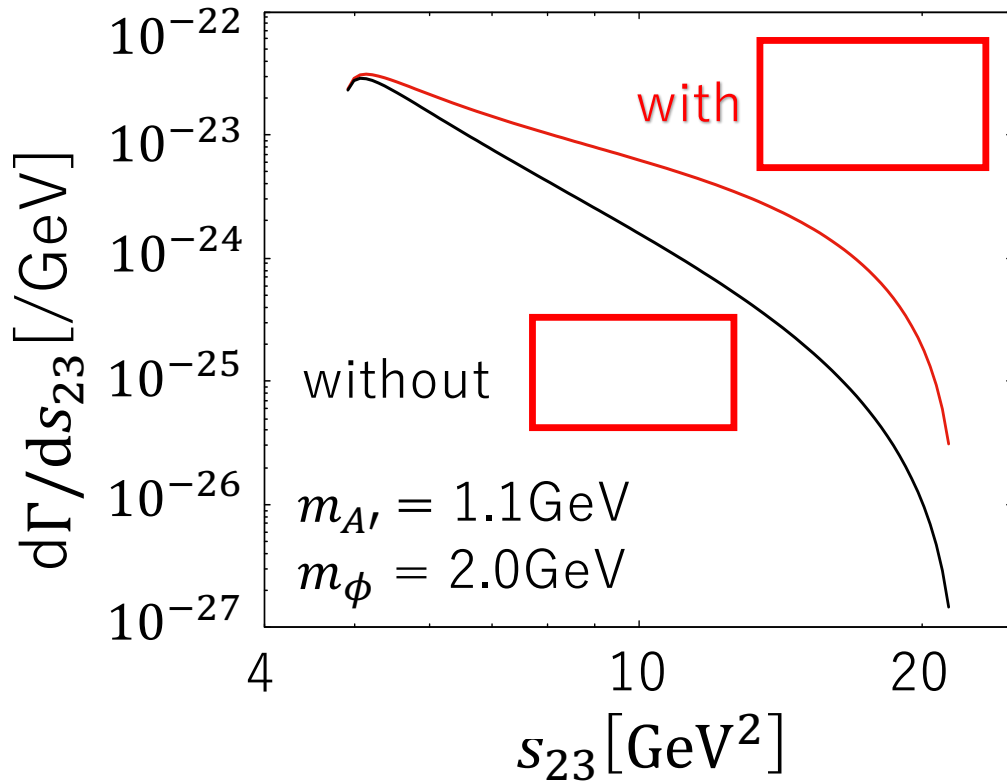
**Sensitivity region closes at $m'_{A'} = m_\phi/2$.
→ contribution of off-shell dark Higgs?**

Dark photon production from **off-shell** dark Higgs decay



$$k^\mu = (E, 0, 0, k),$$

$$\epsilon_L^\mu(k) = \left(\frac{k}{m_{A'}}, 0, 0, \frac{E}{m_{A'}} \right)$$



$$\frac{d\Gamma}{ds_{23}} (b \rightarrow s A' A')$$

$$\propto \frac{1}{(s_{23} - m_\phi^2)^2 + m_\phi^2 \Gamma_\phi^2} \left[2 + \left(\frac{s_{23}}{2m_{A'}^2} - 1 \right)^2 \right]$$

R .S. Chivukula and A. V. Manohar

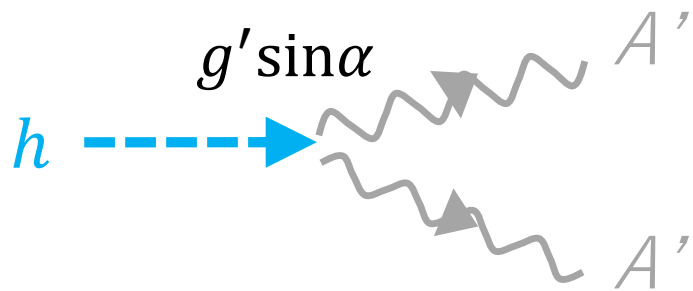
**longitudinal
enhancement**

s_{23} : square of dark Higgs momentum

$$4m_{A'}^2 < s_{23} < (m_b - m_s)^2$$

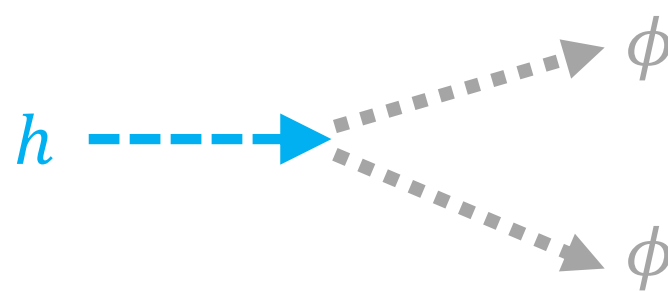
Constraints from Higgs \rightarrow invisible ①

- SM Higgs can decay into dark photons or dark Higgs.
- **Dark Higgs and dark photon are invisible at ATLAS and so on due to their long lifetime.**



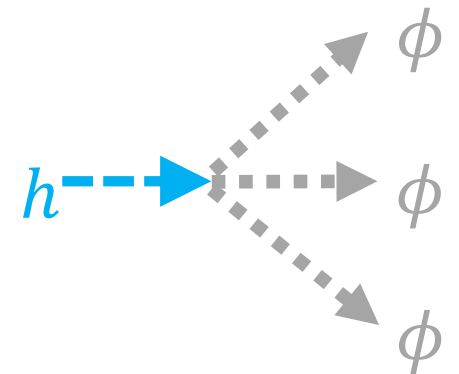
$$\Gamma(h \rightarrow A'A') \propto \frac{m_h^3}{m_{A'}^2} g'^2 \sin^2 \alpha$$

longitudinal enhancement



$$\Gamma(h \rightarrow \phi\phi) \propto \frac{m_h^3}{m_{A'}^2} g'^2 \sin^2 \alpha$$

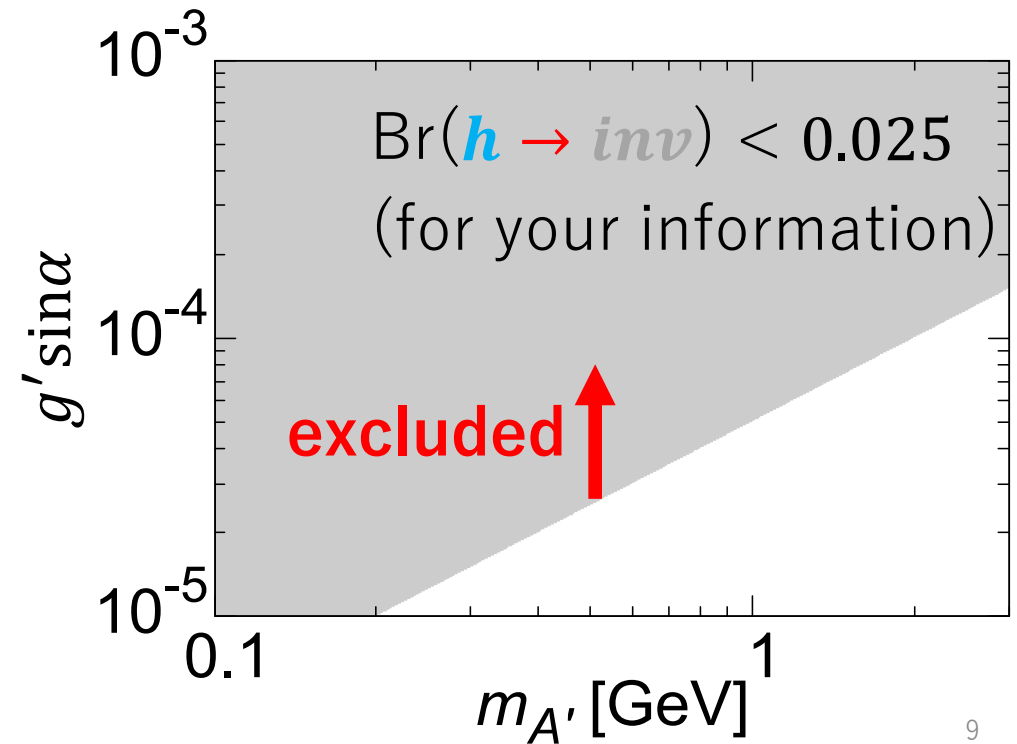
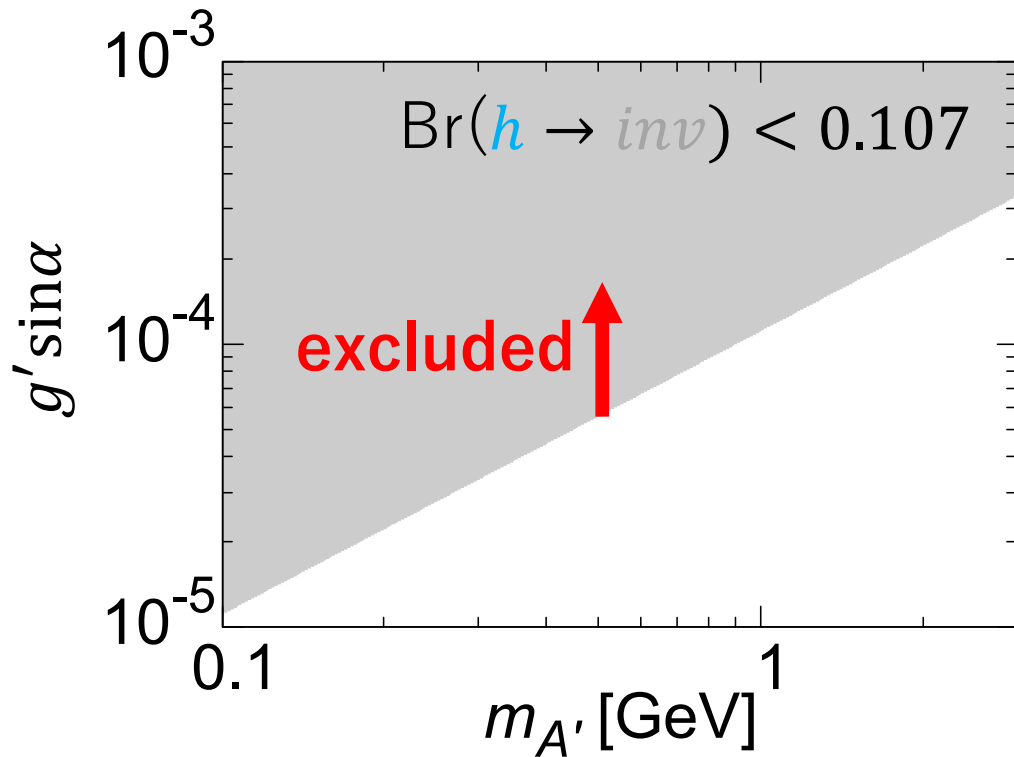
$$h\phi\phi \text{ coupling} \propto \frac{m_h^2}{m_{A'}} g' \sin \alpha$$



phase space
suppression

Constraints from Higgs \rightarrow invisible ②

- From $\text{Br}(h \rightarrow \text{inv}) < 0.107$ (PDG), $g' \sin\alpha$ is constrained.
- Upper bound on $\text{Br}(h \rightarrow \text{inv})$ is expected to be $\text{Br}(h \rightarrow \text{inv}) < 0.025$ at HL-LHC. “Snowmass White Paper Contribution: Physics with the Phase-2 ATLAS and CMS Detectors”, (2022)



Constraints from perturbative unitarity①

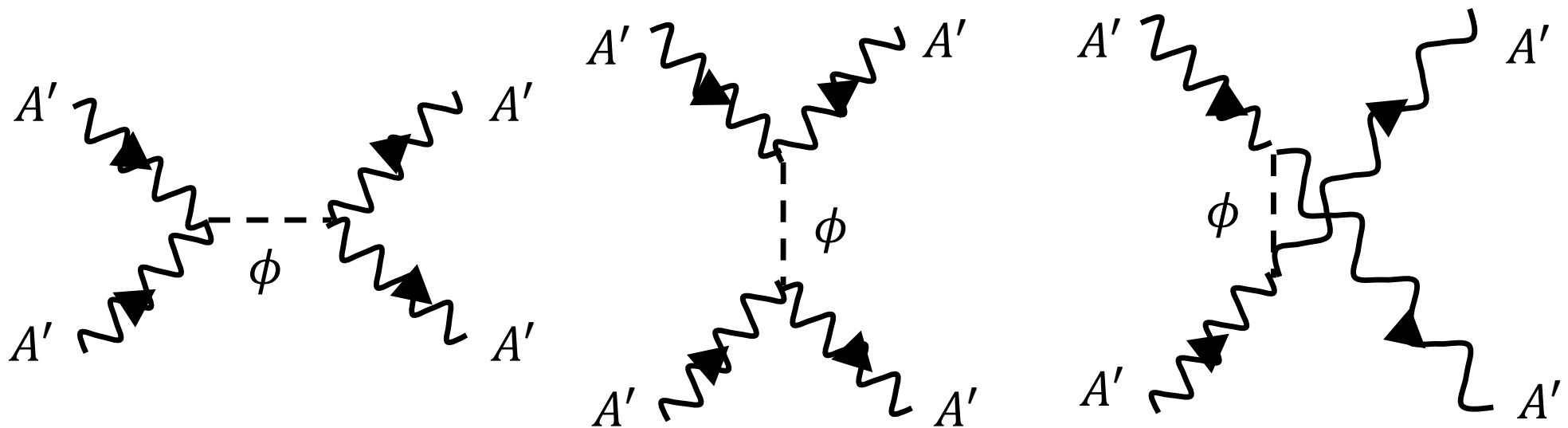
$$\frac{d\Gamma}{ds_{23}} (b \rightarrow sA'A') \propto \frac{s_{23}^2}{m_{A'}^4}$$

For too light $m_{A'}$, is decay width too enhanced ?

→ check by considering perturbative unitarity B, Lee et al. Phys. Rev. D40, 1145, (1977)

calculate two-body scattering amplitude in dark sector at tree level

example : $A'A' \rightarrow A'A'$ diagrams (unitary gauge)

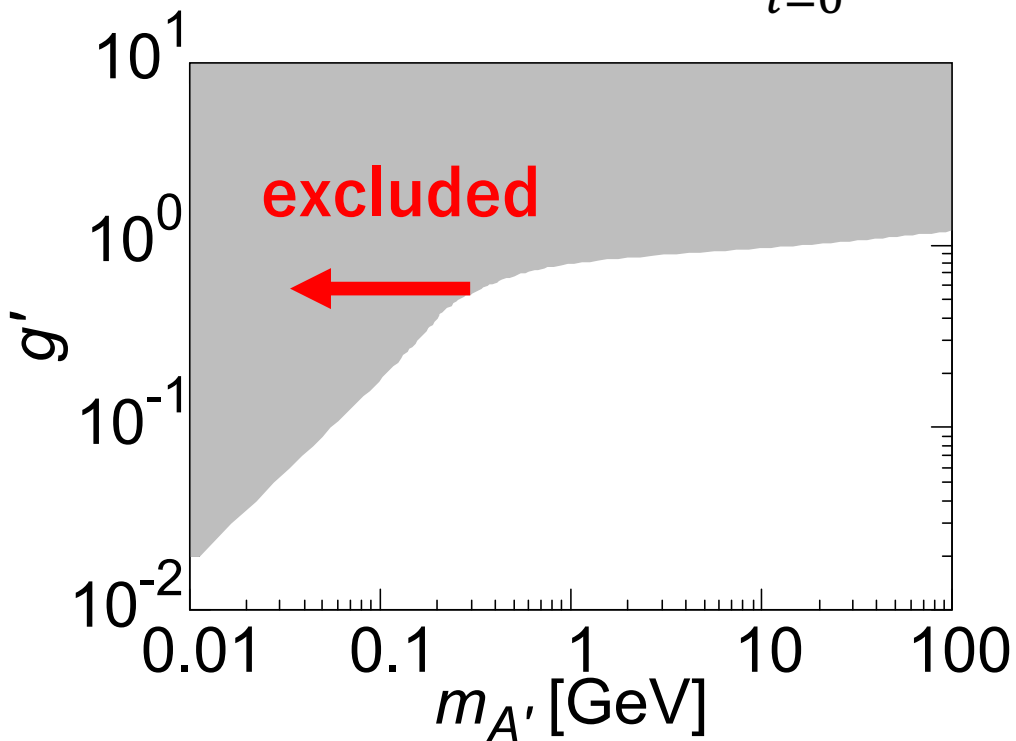


Constraints from perturbative unitarity②

$$\frac{d\Gamma}{ds_{23}}(b \rightarrow sA'A') \propto \frac{s_{23}^2}{m_{A'}^4}$$

expand $M(A'A' \rightarrow A'A')$ and so on by Legendre polynomial

$$M = 16\pi \sum_{l=0}^{\infty} (2l+1) a_l(|\vec{p}|) P_l(\cos\theta)$$



perturbative unitarity bound

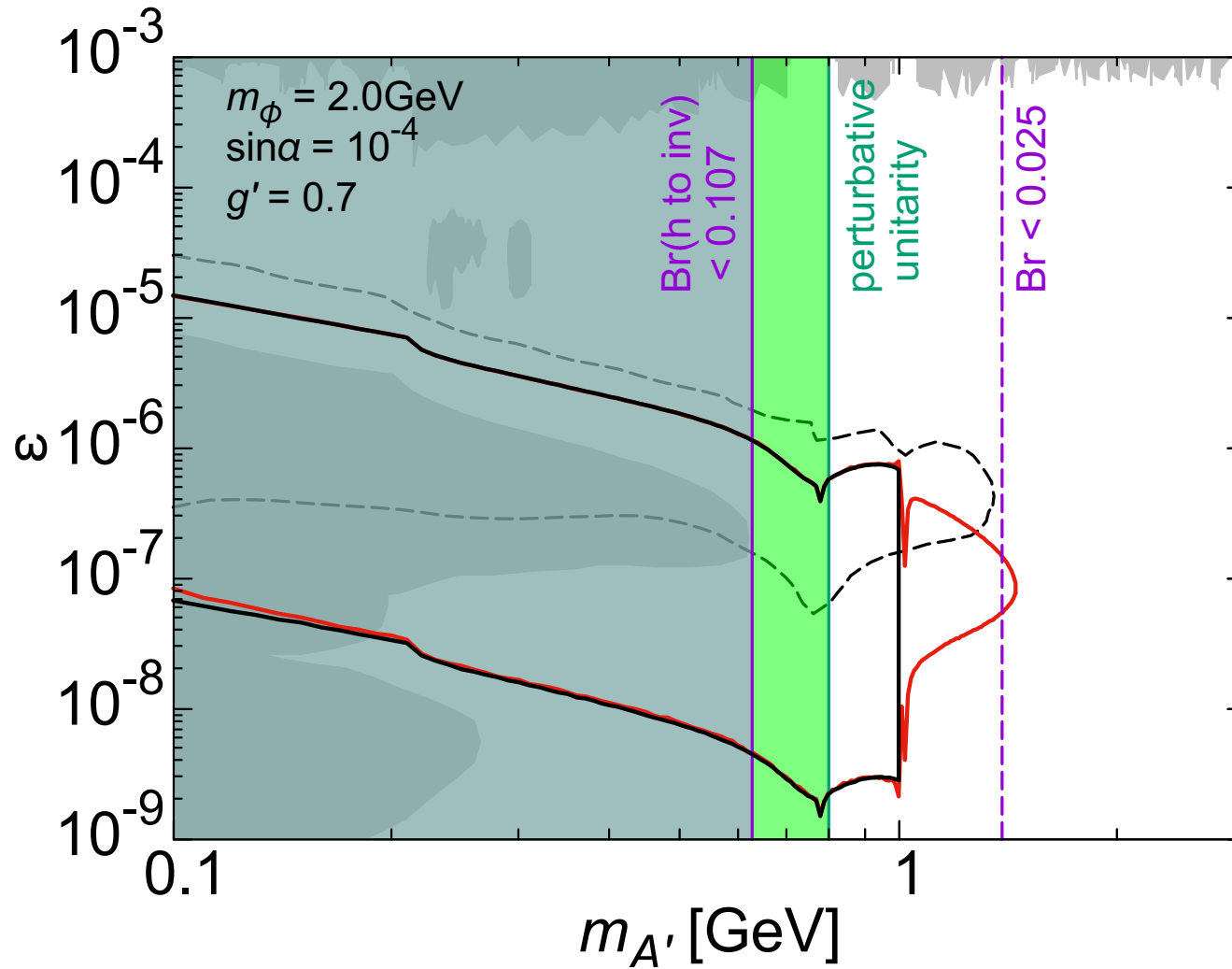
$$|\text{Re } a_0| < 1/2$$



$$1 - \frac{2m_\phi^2}{m_{A'}^2} - 2\log\left(\frac{s}{m_{A'}^2}\right) \leq \frac{8\pi}{g'^2} \text{ etc}$$

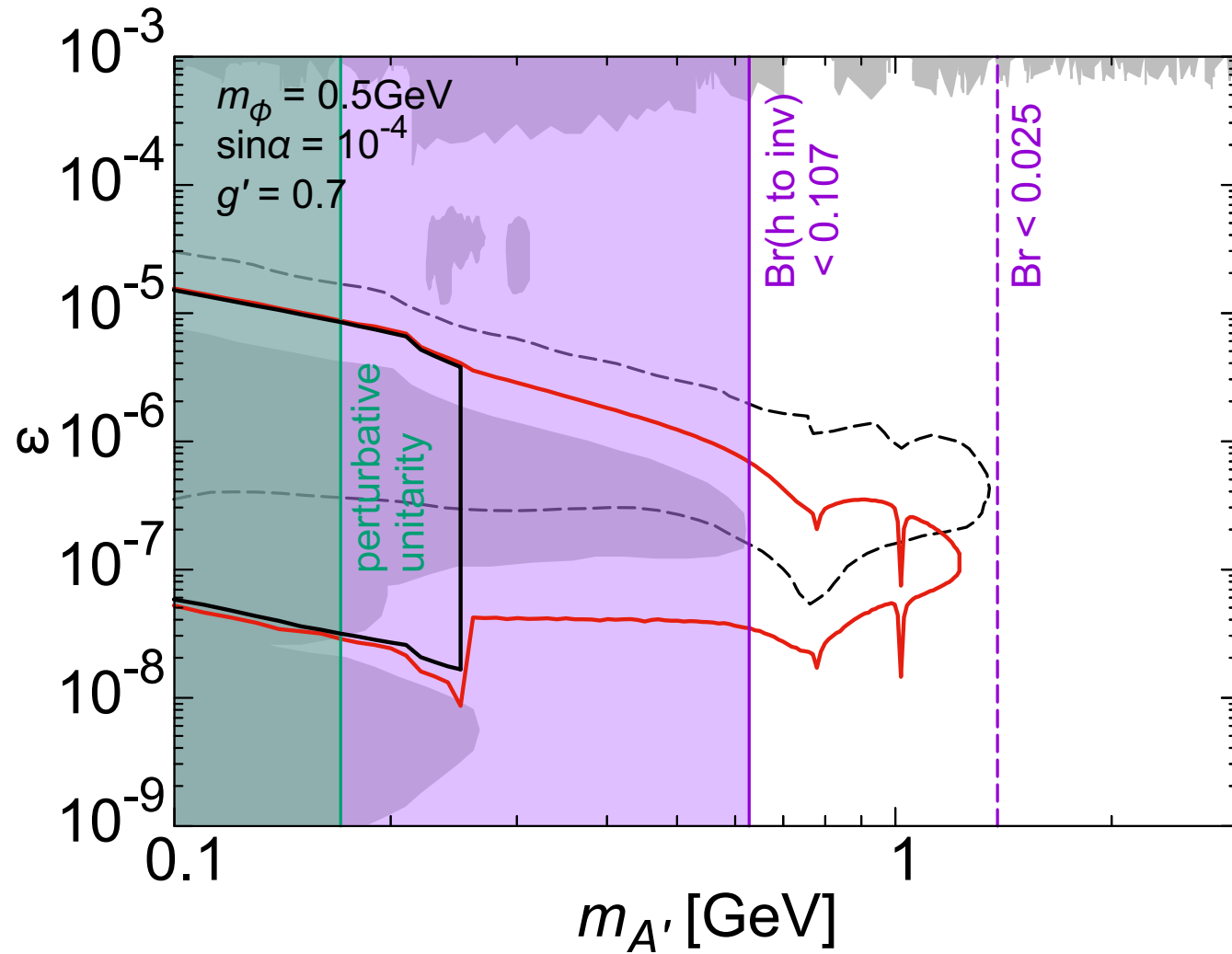
There is lower bound on $m_{A'}$.

Results of **off-shell** dark Higgs contribution ①



Sensitivity region is extended!

Results of **off-shell** dark Higgs contribution②



**Sensitivity region
is extended!**

Summary

- We have studied the possibility of dark photon production from off-shell dark higgs decay at FASER2.
- Dark photon production process is enhanced by its longitudinal component.
- The scalar mixing and extra U(1) gauge coupling constant are constrained by $h \rightarrow inv$.
- We can not use perturbative calculation for too light dark photon and large gauge coupling.
- **Sensitivity region is extended by contribution of off-shell dark higgs.**

appendix

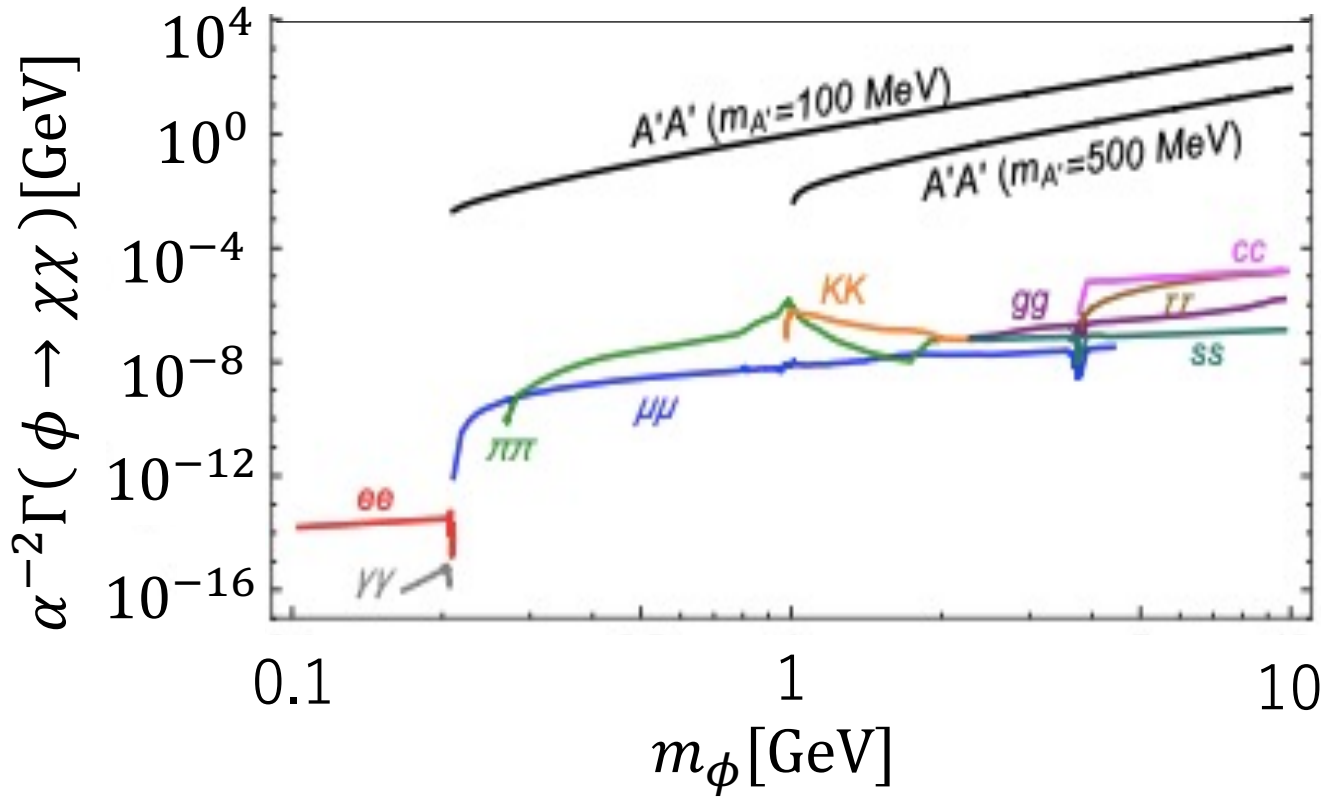
FASER2

	Lmin(m)	Lmax(m)	R(m)	L (ab ⁻¹)
FASER	478.5	480	0.1	0.15
FASER2	475	480	1.0	3.0

Decay width of dark Higgs

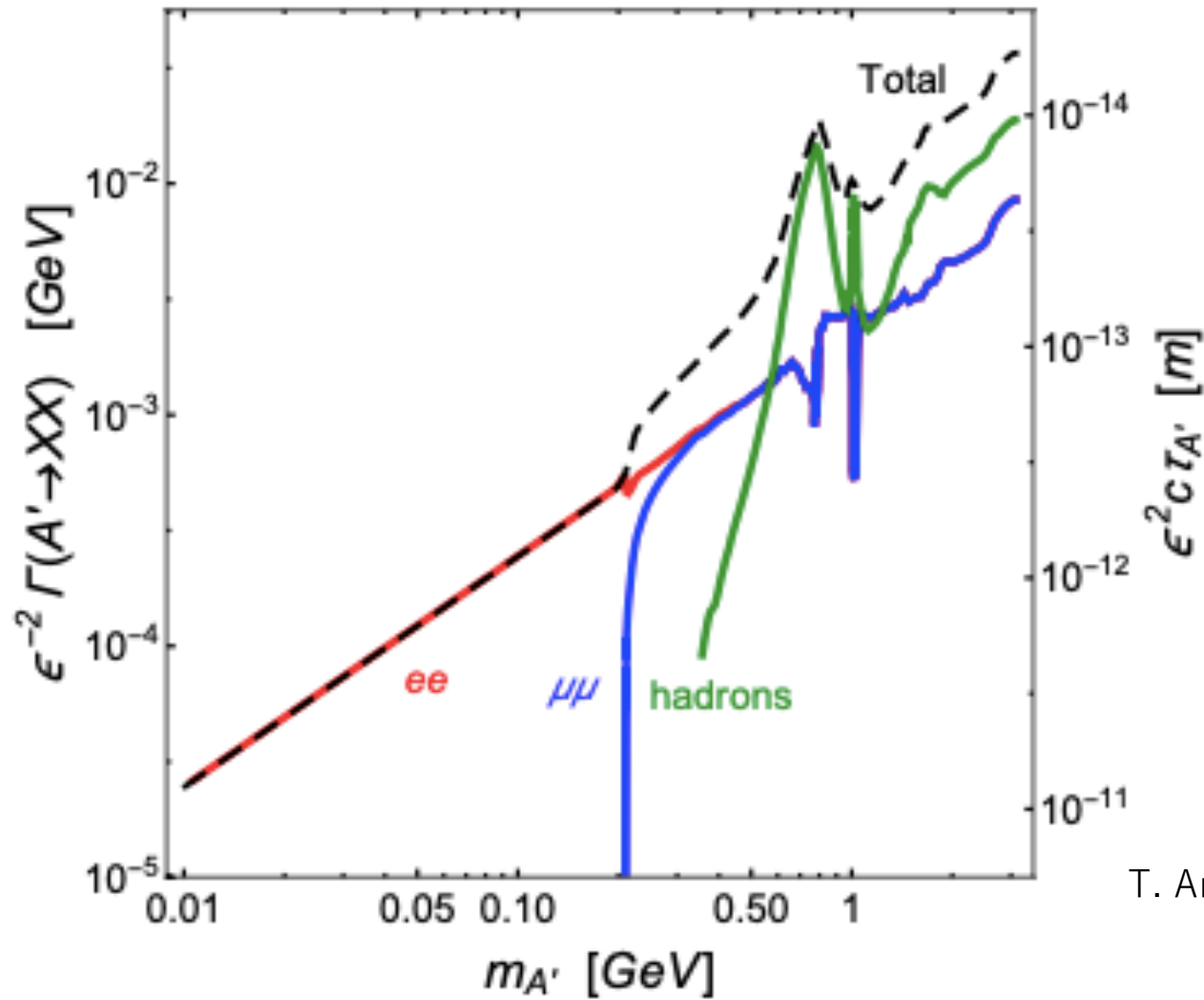
decay width of dark Higgs to dark photon

$$\Gamma(\phi \rightarrow A'A') = \frac{g'}{8\pi} \frac{m_{A'}^2}{m_\phi} \beta_\phi(A') \left(2 + \frac{m_\phi^4}{4m_{A'}^4} \left(1 - \frac{2m_{A'}^2}{m_\phi^2} \right)^2 \right) \text{ (enhanced by } \frac{m_\phi^4}{m_{A'}^4} \text{)}$$



The Decay to dark photon is dominant.

Decay width of dark photon



T. Araki et al. JHEP03(2021)072

Full $\frac{d}{ds_{23}} \Gamma(b \rightarrow s + A'A')$

$$\begin{aligned}
 \frac{d}{ds_{23}} \Gamma(b \rightarrow s + A'A') &= \frac{1}{2m_b} \frac{1}{256\pi^3} \left(1 - \frac{s_{23}}{m_b^2}\right) \sqrt{1 - \frac{4m_A'^2}{s_{23}}} \\
 &\times \frac{9\alpha^2}{32^2\pi^2 \sin^4 \theta_W} |V_{tb}|^2 \frac{m_t^4}{m_W^4} |V_{ts}|^2 \frac{s^2}{v^2} m_b^2 g'^2 m_{A'}^2 c^2 \\
 &\times 2(m_b^2 - s_{23}) \frac{1}{(s_{23} - m_\phi^2)^2 + m_\phi^2 \Gamma_\phi^2} \left[2 + \left(\frac{s_{23}}{2m_{A'}^2} - 1 \right)^2 \right]
 \end{aligned}$$

Event number of $B \rightarrow X_s, \phi^* \rightarrow X_s, A', A'$

$$\begin{aligned}
 N &= L \times \int dp_B d\theta_B ds_{23} d\widehat{\theta}_{A'} d\widehat{\phi}_{A'} d\theta'_s d\phi'_s \frac{d\sigma(pp \rightarrow X, B)}{dp_B d\theta_B} \\
 &\quad \times \frac{dB(B \rightarrow X_s, A', A')}{ds_{23} d\widehat{\theta}_{A'} d\widehat{\phi}_{A'} d\theta'_s d\phi'_s} \times (P_{A'_1} + P_{A'_2} + P_{A'_1} P_{A'_2}) \\
 P_{A_i} &= \left(e^{-\frac{L_{\min}}{d_{A_i}}} - e^{-\frac{L_{\max}}{d_{A_i}}} \right) \times \Theta(R - L_{\max} \tan\theta_{A_i})
 \end{aligned}$$

$\widehat{}$ is the value of dark Higgs rest frame, \prime is the value of B meson rest frame. This is calculated by Monte Carlo method.

Constrain for light dark Higgs A. Fradette, Phys.RevD99(2019)7, 075004

