

# A Quantum Description of Wave Dark Matter

w/ Dhong Yeon Cheong & Lian-Tao Wang



# Motivation

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Establish a more rigorous description of wave DM and the wave-particle boundary

# Outline

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1. What is the density matrix of dark matter?
2. A rigorous definition of the coherence time
3. A single calculation across the wave-particle boundary

Part I

# The Density Matrix of Dark Matter



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Recall, coherent states defined by  $\hat{a} |\alpha\rangle = \alpha |\alpha\rangle$  are complete (but not orthogonal), so can decompose density matrix as

$$\hat{\rho} = \int d^2\alpha P(\alpha) |\alpha\rangle\langle\alpha|$$

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Properties of  $P(\alpha)$ :

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**NB:**  $P(\alpha)$  is not a probability distribution,  $P(\alpha) < 0$  allowed

# The Density Matrix of Dark Matter

[Glauber 1963]:  $P(\alpha)$  obeys the central limit theorem  
So generally expect (e.g. thermal radiation) that

$$\hat{\rho}_{\mathbf{k}} = \int d^2\alpha_{\mathbf{k}} \underbrace{\left( \frac{1}{\pi \langle N_{\mathbf{k}} \rangle} \exp \left[ -\frac{|\alpha_{\mathbf{k}}|^2}{\langle N_{\mathbf{k}} \rangle} \right] \right)}_{P(\alpha_{\mathbf{k}})} |\alpha_{\mathbf{k}}\rangle \langle \alpha_{\mathbf{k}}|$$

**k: mode of the field**

Cf. Coherent state:  
 $P(\alpha) = \delta^{(2)}(\alpha - \beta)$

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$$\langle N_{\mathbf{k}} \rangle = \frac{\text{density of particles}}{\text{density of states}} \simeq \frac{(2\pi\hbar)^3}{g_s} \bar{n} p(\mathbf{k})$$

# density

e.g. Standard Halo Model

Axion:  $g_s = 1$   
 Dark photon:  $g_s = 3$

≈ for local DM

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$\langle N_{\mathbf{k}} \rangle \simeq \bar{n} \times V_{\text{coherence}} \simeq \#$  of indistinguishable particles

Defines wave-particle boundary (given  $\rho_{\text{DM}}$  etc)  
Axions:  $m \simeq 14.4$  eV  
Dark photons:  $m \simeq 11.0$  eV

# Scalar Field Statistics

Let's determine the implications for a scalar field

$$\hat{\phi}(t, \mathbf{x}) = \sum_{\mathbf{k}} \frac{1}{\sqrt{2V\omega_{\mathbf{k}}}} \left( \hat{a}_{\mathbf{k}} e^{-ik \cdot x} + \hat{a}_{\mathbf{k}}^{\dagger} e^{ik \cdot x} \right)$$



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As usual,  $\langle \hat{\mathcal{O}} \rangle = \text{Tr}[\hat{\rho} \hat{\mathcal{O}}]$ , but if  $[\hat{a}, \hat{a}^{\dagger}] = 0$ , set  $\hat{a}_{\mathbf{k}}^{(\dagger)} = \alpha_{\mathbf{k}}^{(*)}$

$[\hat{a}, \hat{a}^{\dagger}] \neq 0$   
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$$\Rightarrow \phi(t, \mathbf{x}) = \sum_{\mathbf{k}} \sqrt{\frac{2}{V\omega_{\mathbf{k}}}} \text{Re} [\alpha_{\mathbf{k}} e^{-ik \cdot x}]$$

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$\Rightarrow \phi$  is a Gaussian random field, with

$$\langle \phi(t, \mathbf{x}) \rangle = 0 \quad \& \quad \langle \phi^2(t, \mathbf{x}) \rangle \simeq \frac{\rho}{m^2}$$

Also  $\partial_t \phi \sim \text{Im}[\alpha]$  is  
an independent  
Gaussian random field

# $P(\alpha)$ Experimentally Testable

Key assumption: Gaussian  $P(\alpha)$

May not be true, e.g. coherent state or Bose-Einstein  
Condensate

BEC: e.g. [Sikivie, Yang 2009]  
[Erken, Sikivie, Tam, Yang 2012]

Could resolve with experiment (post discovery of DM):  
look for non-Gaussianities in the fluctuations of  $\phi$

Part II

# The Coherence Time

# Autocorrelation function

Having understood  $\langle \phi^n(t, \mathbf{x}) \rangle$ , natural to next consider

$$\Gamma(\tau, \mathbf{d}) = \langle \phi(t, \mathbf{x}) \phi(t + \tau, \mathbf{x} + \mathbf{d}) \rangle$$

Assume stationary/homogeneous  
 $\Rightarrow \langle \mathcal{O} \rangle$  independent of  $(t, \mathbf{x})$

Intuition: how much does knowledge of the field at one point tell you about it at another?

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If stationary\* can derive (with  $\mathbf{d} = 0$ )

$$\Gamma(\tau) = \frac{\rho}{\bar{\omega}} \int d\omega \frac{p(\omega)}{\omega} \cos(\omega\tau)$$

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 Also have results for  $\mathbf{d} \neq 0$   
 Cf. [Derevianko 2018]



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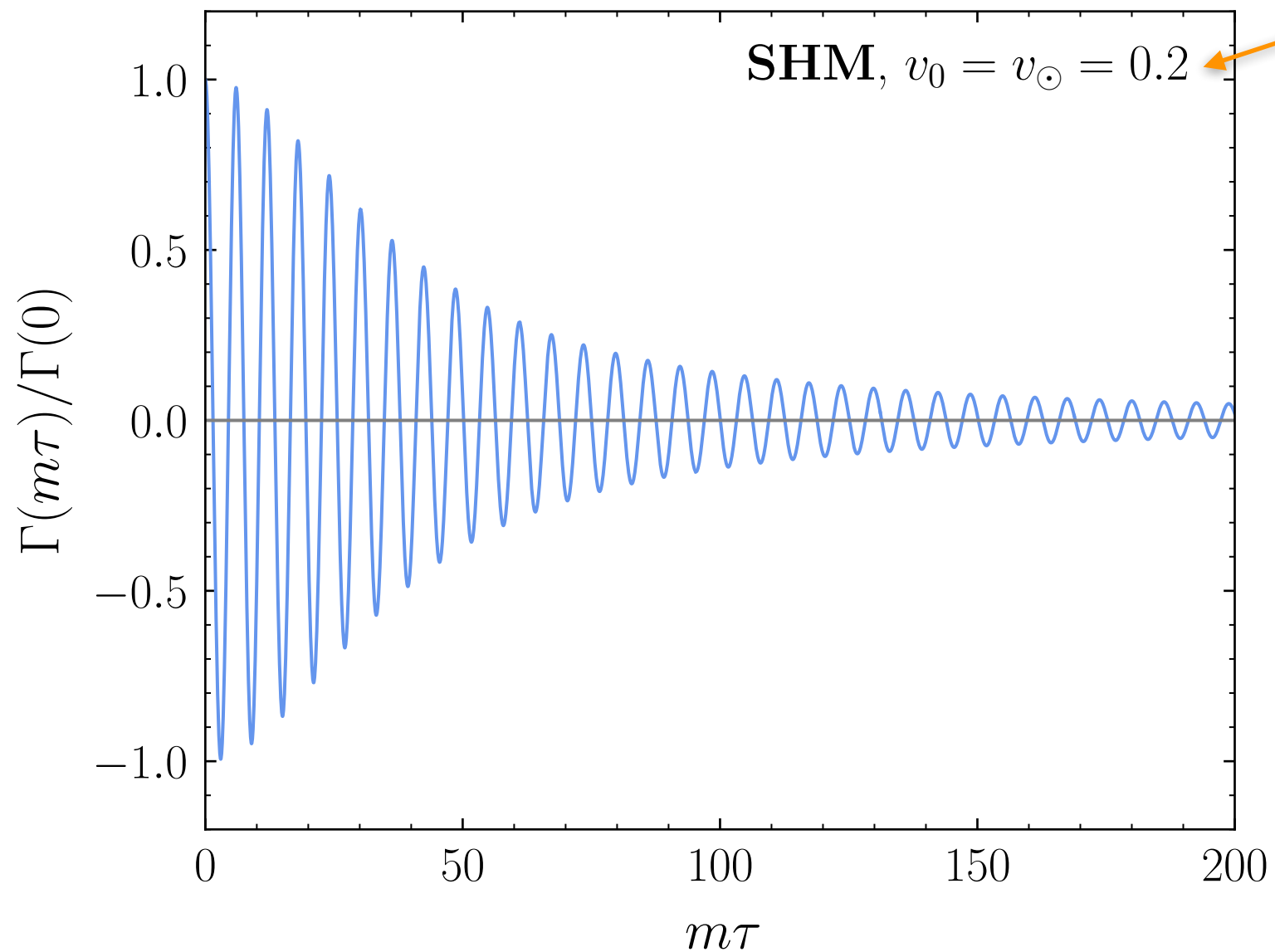
For DM,  $\omega \simeq m + \frac{1}{2}mv^2$ , with  $v$  set by e.g.

$$f(\mathbf{v}) = \frac{1}{\pi^{3/2} v_0^3} e^{-(\mathbf{v} + \mathbf{v}_\odot)^2 / v_0^2}$$

Standard Halo Model

# Autocorrelation function

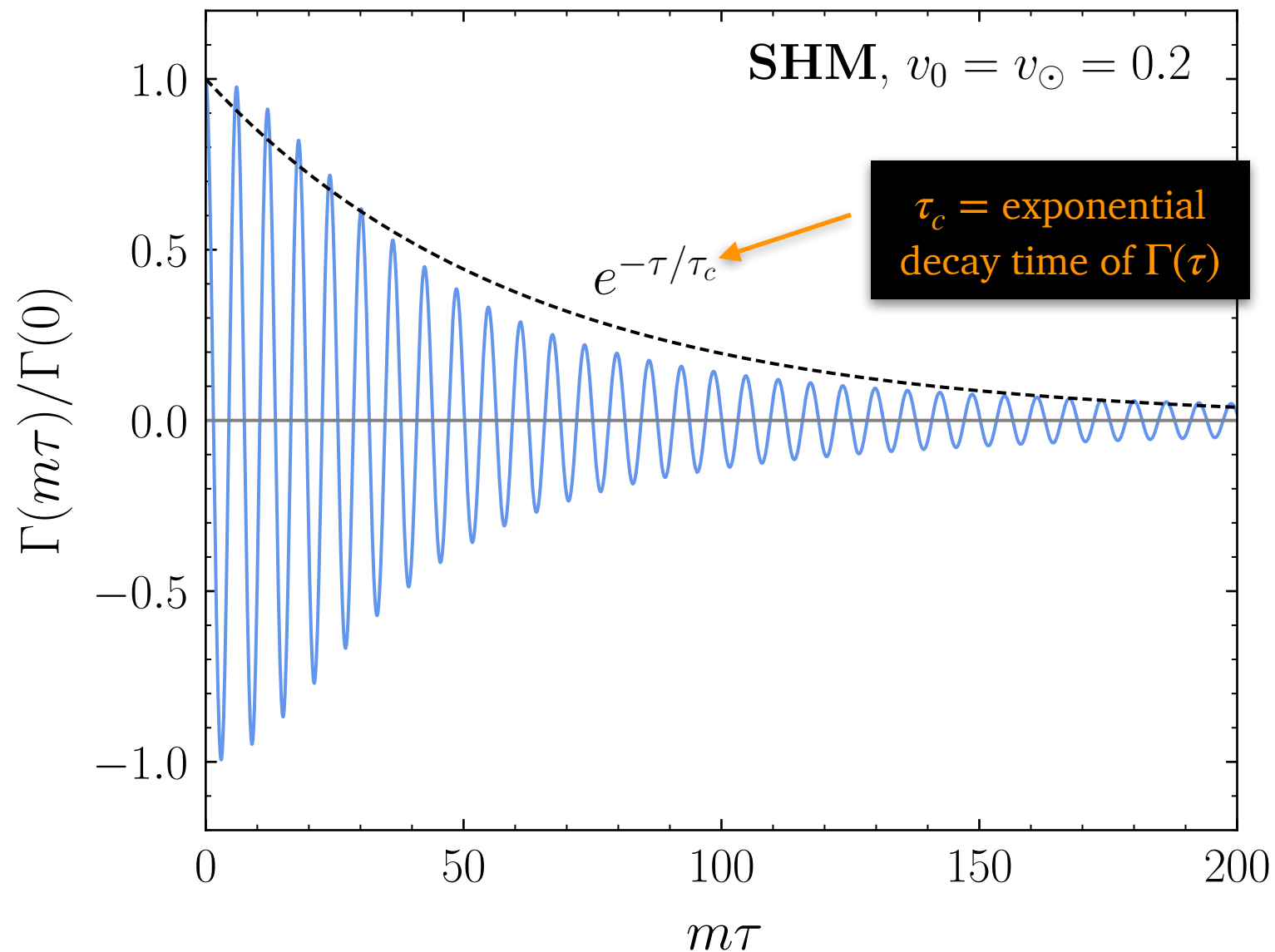
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In reality,  
 $v_0 \sim v_{\odot} \sim 10^{-3}$

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# Coherence Time

$$\text{Define: } \tau_c = \int_{-\infty}^{\infty} d\tau \left| \frac{\Gamma(\tau)}{\Gamma(0)} \right|^2$$

Common def. in quantum optics,  
e.g. [Mandel & Wolf, “Optical  
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**Example 2:** DM with the SHM

$$\tau_c = \frac{\sqrt{2\pi}\text{Erf}\left[\sqrt{2}v_{\odot}/v_0\right]}{mv_0v_{\odot}} \left( 1 + \frac{3v_0^2}{4} - \frac{v_0v_{\odot}e^{-2v_{\odot}^2/v_0^2}}{\sqrt{2\pi}\text{Erf}\left[\sqrt{2}v_{\odot}/v_0\right]} + \mathcal{O}(v^4) \right)$$

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$$\simeq 2.8 \text{ s} \left( \frac{1 \text{ neV}}{m} \right)$$

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By the Wiener-Khinchin theorem,

$$S(\omega) = \int_{-\infty}^{\infty} d\tau \Gamma(\tau) e^{i\omega\tau}$$

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Further, width of  $S(\omega)$  is  $\Delta\omega = 1/\tau_c$

Intuition:  $\tau_c$  measures how long  
 $\phi(t) = \phi_0 \cos(mt)$  is a good approximation  
 See also [Dror, Gori, Leedom, NLR 2023]

Part III

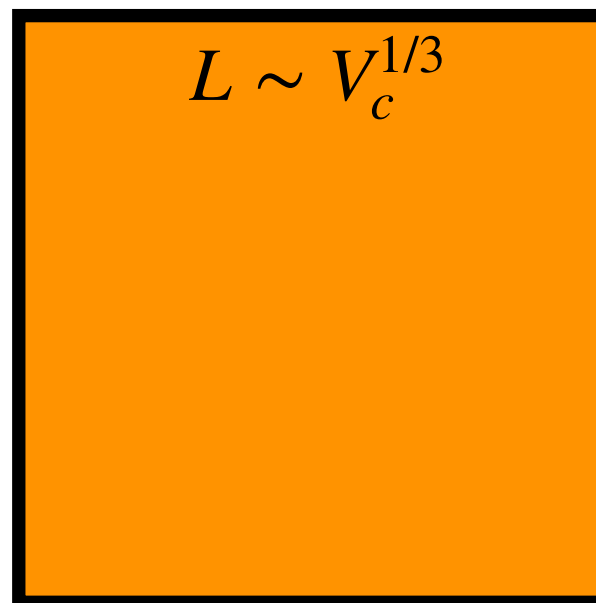
# Wave-Particle Boundary

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So far  $[\hat{a}, \hat{a}^\dagger] \simeq 0$  (justify by  $N \gg 1$ )

Now  $[\hat{a}, \hat{a}^\dagger] = 1$ , but for simplicity take a single mode  
( $\omega = m$ )

**Question:** what is the energy in a box of volume  $V_c$ ?



Similar result holds for  
calculation in a finite  
physical volume

# Wave-Particle Boundary

Rewrite Gaussian  $\hat{\rho}$  in the number basis

$$\begin{aligned}\hat{\rho} &= \int d^2\alpha \frac{e^{-|\alpha|^2/\langle N \rangle}}{\pi\langle N \rangle} |\alpha\rangle\langle\alpha| \\ &= \frac{1}{1 + \langle N \rangle} \sum_{k=0}^{\infty} \left( \frac{\langle N \rangle}{1 + \langle N \rangle} \right)^k |k\rangle\langle k|\end{aligned}$$

Here  $k \in \mathbb{N}$ , not wavevector!

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Can use to show  
 $\text{Tr}[\hat{\rho}^2] = (1 + 2\langle N \rangle)^{-1}$

Probability of seeing  $k$  quanta in  $V_c$  is

$$p(k) = \frac{1}{1 + \langle N \rangle} \left( \frac{\langle N \rangle}{1 + \langle N \rangle} \right)^k$$

For a single mode:  $E = m \times k$ , so we can just study  $k$

# Wave-Particle Boundary

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The mean and standard deviation of  $k$ :

$$\mu_k = \langle k \rangle = \langle N \rangle$$
$$\sigma_k^2 = \langle k^2 \rangle - \langle k \rangle^2 = \langle N \rangle(1 + \langle N \rangle)$$

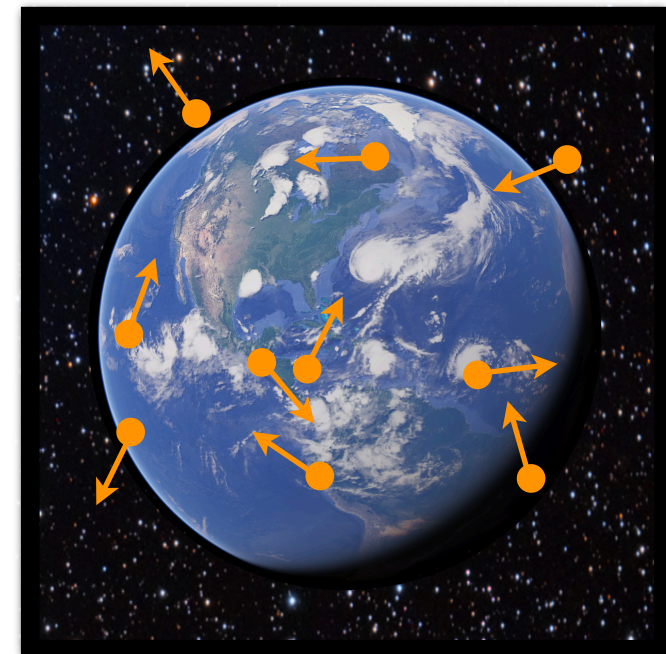


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For  $\langle N \rangle \ll 1$ ,  $\sigma_k^2 = \mu_k$   
Poisson distributed



Holds for all higher moments

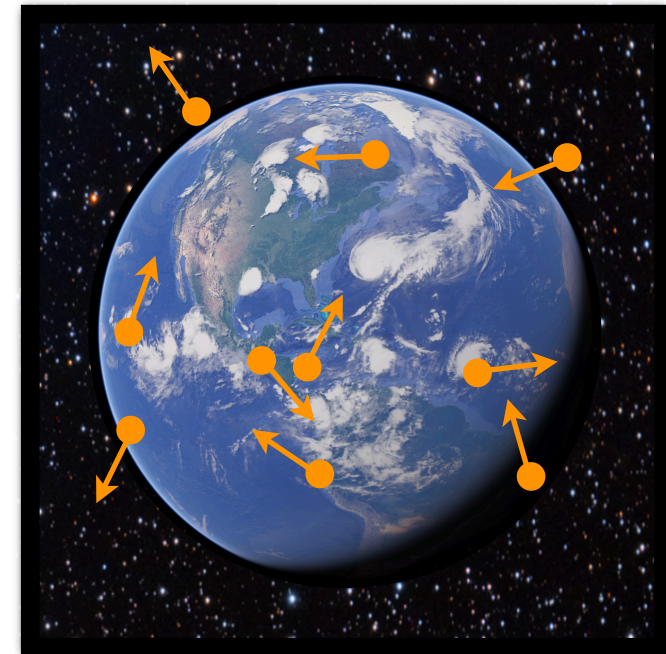
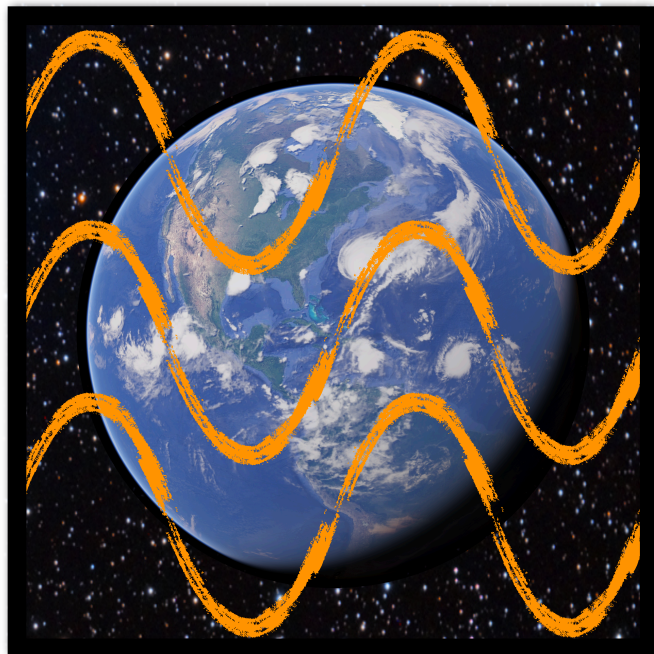
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For  $\langle N \rangle \gg 1$ ,  $\sigma_k^2 = \mu_k^2$   
Exponentially distributed

For  $\langle N \rangle \ll 1$ ,  $\sigma_k^2 = \mu_k$   
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For  $\langle N \rangle \sim 1$  neither Poisson nor exponential



# Conclusion

The quantum approach opens a path to a rigorous description of wave dark matter

Open questions:

- Determine the exact  $P(\alpha)$  of DM
- Interface with experiment (quantum measurement theory)
- Resolve the distribution of polarizations for dark photons
- ...

