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Enhanced Di-Higgs Production in Extended Higgs Models within Photonic Colliders

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International Symposium on High Energy Physics (ISHEP-2024) October 19, 2024 $\begin{array}{l} \mbox{Introduction} & \mbox{Introduction} \\ \mbox{The Inet Doublet Plus Two Active Doublets Model } (1+2)\mbox{HDM} \\ & \mbox{The process of } \gamma \gamma \rightarrow hh \mbox{ in the SM} \\ \mbox{The process of } \gamma \gamma \rightarrow hh \mbox{ in the } (1+2)\mbox{HDM} \\ & \mbox{Conclusion} \end{array}$

Outlines



- The Inet Doublet Plus Two Active Doublets Model I(1+2)HDM
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Higgs Boson July 04, 2012 : Discovery BSM



Opening wide windows in this field for searching on the Higgs properties either on Standard model or Beyond.

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I(1+2)HDM in a nutshell

- The I(1+2)HDM is an extension of the 2HDM, comprising the scalar doublets Φ_1 and Φ_2 , by an inert scalar doublet η which can provide a stable dark matter candidate.
- The model exhibits a $\mathbb{Z}_2 \times \mathbb{Z}'_2$ symmetry, with the first factor being the inert-doublet \mathbb{Z}_2 , where only the field η transforms as $\eta \to -\eta$, while all other fields remain neutral.

$$\Phi_{i} = \begin{pmatrix} \phi_{i}^{+} \\ (\mathbf{v}_{i} + \eta_{i} + i\mathbf{z}_{i})/\sqrt{2} \end{pmatrix}, (i = 1, 2); \quad \eta = \begin{pmatrix} \chi^{+} \\ (\chi + i\chi_{a})/\sqrt{2} \end{pmatrix}$$
(1)

• The most general scalar potential consistent with the gauge invariant and CP invariant is given by :

$$V(\Phi_1, \Phi_2, \eta) = V(\Phi_1, \Phi_2) + V(\eta) + V(\Phi_1, \Phi_2, \eta),$$
(2)

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I(1+2)HDM in a nutshell

where

$$V(\Phi_{1}, \Phi_{2}) = -\frac{1}{2} \left\{ m_{11}^{2} \Phi_{1}^{\dagger} \Phi_{1} + m_{22}^{2} \Phi_{2}^{\dagger} \Phi_{2} + \left[m_{12}^{2} \Phi_{1}^{\dagger} \Phi_{2} + \text{h.c.} \right] \right\} \\ + \frac{\rho_{1}}{2} (\Phi_{1}^{\dagger} \Phi_{1})^{2} + \frac{\rho_{2}}{2} (\Phi_{2}^{\dagger} \Phi_{2})^{2} + \rho_{3} (\Phi_{1}^{\dagger} \Phi_{1}) (\Phi_{2}^{\dagger} \Phi_{2}) \\ + \rho_{4} (\Phi_{1}^{\dagger} \Phi_{2}) (\Phi_{2}^{\dagger} \Phi_{1}) + \frac{1}{2} \left[\rho_{5} (\Phi_{1}^{\dagger} \Phi_{2})^{2} + \text{h.c.} \right],$$
(3)

the potential of the inert sector is written as :

$$V(\eta) = m_{\eta}^2 \eta^{\dagger} \eta + \frac{\rho_{\eta}}{2} (\eta^{\dagger} \eta)^2, \qquad (4)$$

and the mixing terms are defined as follows :

$$V(\Phi_{1}, \Phi_{2}, \eta) = \rho_{1133}(\Phi_{1}^{\dagger}\Phi_{1})(\eta^{\dagger}\eta) + \rho_{2233}(\Phi_{2}^{\dagger}\Phi_{2})(\eta^{\dagger}\eta) + \rho_{1331}(\Phi_{1}^{\dagger}\eta)(\eta^{\dagger}\Phi_{1}) + \rho_{2332}(\Phi_{2}^{\dagger}\eta)(\eta^{\dagger}\Phi_{2}) + \frac{1}{2} \left[\rho_{1313}(\Phi_{1}^{\dagger}\eta)^{2} + \text{h.c.}\right] + \frac{1}{2} \left[\rho_{2323}(\Phi_{2}^{\dagger}\eta)^{2} + \text{h.c.}\right].$$
(5)

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I(1+2)HDM in a nutshell

 The spectrum of the I(1+2)HDM containes 9 scalar particles : two CP-even, H and h which is the SM Higgs boson, a CP-odd A and a pair of charged scalars H[±], equivalent to pure 2HDM and for inert sector contains three scalars : χ, χ_a and χ[±]. Their masses are given by :

$$m_{\chi^{\pm}}^2 = m_{\eta}^2 + \frac{1}{2} \rho_a v^2,$$
 (6)

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$$m_{\chi}^{2} = m_{\chi^{\pm}}^{2} + \frac{1}{2}(\rho_{b} + \rho_{c})v^{2}, \quad m_{\chi_{a}}^{2} = m_{\chi^{\pm}}^{2} + \frac{1}{2}(\rho_{b} - \rho_{c})v^{2}.$$
(7)

wher

$$\rho_a \equiv \rho_{1133} = \rho_{2233}, \quad \rho_b \equiv \rho_{1331} = \rho_{2332}, \quad \rho_c \equiv \rho_{1313} = \rho_{2323}, \quad (8)$$

• The final independent parameters for the model are 12 which we choose as follows :

$$\Omega = \left\{ m_{h}, m_{A}, m_{H}, m_{H^{\pm}}, m_{12}^{2}, \beta, \alpha, m_{\chi}, m_{\chi_{a}}, m_{\chi^{\pm}}, m_{\eta}^{2}, \rho_{\eta} \right\}, \qquad (9)$$

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Theoretical and Experimental constraints

- Theoretical constraints : Perturbativity, Vacuum Stability, Charge-breaking minima, Unitarity..
- Experimental constraints :
 - * Contraints from Higgs data at LHC :
 - * Direct search from LEP :
 - $m_A, m_H > 80 90 \, GeV$
 - $m_{H^\pm}, m_{\chi^\pm} >$ 80 100 GeV
 - * electroweak precision through S, T, U .
 - * DM relic density, direct, indirect and collider searche.
 - * tools used : HiggsBounds-5.10.2, HiggsSignals-2.6.2, MicrOmegas-5.2.

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Feynman diagram at one-loop



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Total cross section for $\gamma\gamma \rightarrow hh$ and $e^+e^- \rightarrow \gamma\gamma \rightarrow hh$ as a function of the collision energy.



the box loop contributions almost entirely dominate the triangle ones for the partonic cross-section.

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Feynman diagram Results and discussion

Vertex type diagrams contribution at one-loop



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It can be clearly seen that box contribution has either the quartic vertex $hh\chi^+\chi^-$ and/or hH^+H^- or twice the triple vertex $hH^+_{-}H^-_{-}$ and/or $h\chi^+_{-}\chi^-_{-}$

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total cross section $\sigma(e^+e^- \rightarrow \gamma\gamma \rightarrow hh)$



• The threshold effect is observed for each BPs when $E_{\gamma\gamma} \approx 2m_{\chi^{\pm}}(2m_{H^{\pm}})$, corresponding to the opening of the inert charged Higgs/chagred Higgs pair channel, $\gamma\gamma \rightarrow \chi^{+}\chi^{-}$; $H^{+}H^{-}$.

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total cross section $\sigma(e^+e^- \rightarrow \gamma\gamma \rightarrow hh)$

- The cross-section is enhanced significantly for small $m_{\chi^{\pm}}$ when $m_{\chi^{\pm}} < m_{H^{\pm}}$, and the production channel $\gamma\gamma \rightarrow \chi^{+}\chi^{-}$ opened firstly, wihle for the BP4 when $m_{H^{\pm}} < m_{\chi^{\pm}}$ the cross-section decreases compared to the other BPs and the production channel $\gamma\gamma \rightarrow \chi^{+}\chi^{-}$ opened secondly.
- The total cross-section is now principally determined by the contributions of the boxes, both at low and high center-of-mass energy. This is because the triple couplings $h\chi^+\chi^-$, hH^+H^- and quartic couplings $hh\chi^+\chi^-$, hhH^+H^- , which contribute to the boxes, are relatively large.

BPs	BP1	BP2	BP3	BP4
$\lambda_{hhh}^{\rm I(1+2)HDM}\diagup\lambda_{hhh}^{\rm SM}\approx$	0.024	1.55	0.89	1.03
$\lambda_{Hhh}^{\rm I(1+2)HDM}\diagup\lambda_{hhh}^{\rm SM}\approx$	0.63	-0.15	0.48	-0.12
$\lambda_{hH^+H^-}^{\rm I(1+2)HDM}\diagup\lambda_{hhh}^{\rm SM}\approx$	10.25	11.03	4.43	3.22
$\lambda_{HH^+H^-}^{\rm I(1+2)HDM}\diagup\lambda_{hhh}^{\rm SM}\approx$	8.94	9.18	5.33	2.45
$\lambda_{h\chi^+\chi^-}^{\rm I(1+2)HDM} \diagup \lambda_{hhh}^{\rm SM} \approx$	12.84	10.52	14.54	10.17
$\lambda_{H\chi^+\chi^-}^{\rm I(1+2)HDM} \diagup \lambda_{hhh}^{\rm SM} \approx$	8.12	4.68	6.36	3.31
$\lambda_{H^+H^-hh}^{\rm I(1+2)HDM} \diagup \lambda_{hhhh}^{\rm SM} \approx$	9.62	10.35	5.76	2.98
$\lambda^{\rm I(1+2)HDM}_{\chi^+\chi^-hh}\diagup\lambda^{\rm SM}_{hhhh}\approx$	15.49	11.73	16.18	10.90

The trilpe and quartic Higgs couplings in I(1+2)HDM normalized by SM ,

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Conclusion

The cross section for production of Di-Higgs boson via photon- fusion at photonic collider within the framework of I(1+2)HDM can be enhanced more than the corresponding SM after all constraints.

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