

Higgs Self-Couplings in the MSSM

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in collaboration with M. Spira

Outline

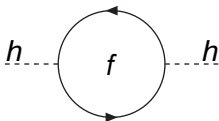
- 1 Motivations for the MSSM
- 2 The Higgs Sector in the MSSM
- 3 MSSM Higgs search
- 4 Radiative Corrections to Higgs Self-Couplings

Hierarchy Problem in the Standard Model

Embedding the SM in a grand unified theory (GUT)

→ introduce UV cutoff Λ

- Loop corrections to the Higgs mass generate quadratic divergence:



$$\Delta m_H^2 \sim \Lambda^2 \sim \mathcal{O}(M_{GUT}) \gg m_H^2$$

absorb in counterterm: $m_H^2 \rightarrow m_H^2 + \Delta m_H^2 - \delta m_H^2$

⇒ unnatural fine tuning ~ 28 digits

Motivations for the MSSM

- Supersymmetry: fermions \leftrightarrow bosons
- Quadratic divergencies are cancelled:



$$\Delta m_H^2 \sim (\tilde{m}^2 - m^2) \log \frac{\Lambda^2}{m^2} \Rightarrow \text{no fine-tuning for } \tilde{m} \lesssim \mathcal{O}(1 \text{ TeV})$$

\Rightarrow Hierarchy problem solved

- SUSY- $SU(5)$ -GUT predicts: $\sin^2 \Theta_W = 0.2334 \pm 0.0026$
LEP measures: $\sin^2 \Theta_W = 0.2317 \pm 0.0002$

Electro Weak Symmetry Breaking in the MSSM

Two complex Higgs doublets: H_1, H_2

→ 8 degrees of freedom (DoF)

Longitudinal polarizations of the W^\pm and Z bosons eat 3 DoF:

⇒ 5 physical Higgs bosons

- 2 scalar Higgs bosons h, H
- 1 pseudoscalar Higgs boson A
- 2 charged Higgs bosons H^+, H^-

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Higgs Boson Masses in the MSSM at Tree Level

The MSSM Higgs potential is

$$V_0 = (m_1^2 + \mu^2)|H_1|^2 + (m_2^2 + \mu^2)|H_2|^2 - B\mu\epsilon_{ij}(H_1^i H_2^j + h.c.) + \frac{g^2 + g'^2}{8}(|H_1|^2 - |H_2|^2)^2 + \frac{g^2}{2}|H_1^\dagger H_2|^2.$$

diagonalize mass matrix \rightarrow physical Higgs boson masses

$$m_{H^\pm}^2 = m_A^2 + m_W^2 \quad m_A^2 = \frac{2B\mu}{\sin(2\beta)}$$

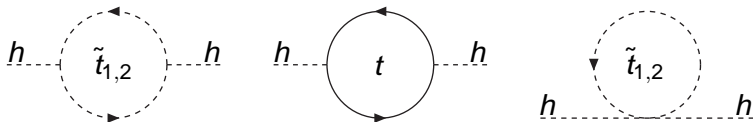
$$m_{H,h}^2 = \frac{1}{2} \left[m_A^2 + m_Z^2 \pm \sqrt{(m_A^2 + m_Z^2)^2 - 4m_A^2 m_Z^2 \cos^2(2\beta)} \right]$$

$$\Rightarrow m_h < m_Z$$

One-Loop Corrections to Higgs Masses

Top-Yukawa-coupling h_t is large

→ get large corrections from top/stop loops $\mathcal{O}(\alpha_t)$, $\alpha_t = \frac{h_t^2}{4\pi}$:



$$m_h^2 \leq m_Z^2 + \frac{3G_F}{\sqrt{2}\pi^2} \frac{m_t^4}{\sin^2(\beta)} \left[\log \frac{m_{\tilde{t}_1} m_{\tilde{t}_2}}{m_t^2} + \frac{X_t^2}{M_{SUSY}^2} \left(1 - \frac{X_t^2}{12M_{SUSY}^2} \right) \right]$$

$$X_t = A_t - \mu \cot(\beta)$$

Ellis, ... (1991)
 Haber, ... (1990)
 etc.

⇒ m_h is easily pushed beyond m_Z bound

Further Corrections to Higgs Masses

- QCD corrections are dominant at two-loop: $\mathcal{O}(\alpha_t \alpha_s)$

Hempfling, Hoang (1994)
Heinemeyer, . . . (1998)
Zhang (1999)
Slavich, . . . (2001)
etc.

- $\mathcal{O}(\alpha_t^2)$, $\mathcal{O}(\alpha_b \alpha_s)$ have been calculated.

Espinosa, Zhang (2000)
Slavich, . . . (2001)
Heinemeyer, . . . (2005)
etc.

- First 3-loop results: $\mathcal{O}(\alpha_t \alpha_s^2)$

Martin (2007)
Harlander, . . . (2010)

$$\Rightarrow m_h \lesssim 135 \text{ GeV}$$

Higgs couplings to SM particles

Modified couplings (tree level):

ϕ	g_u^ϕ	g_d^ϕ	g_V^ϕ
H_{SM}	1	1	1
h	$\cos(\alpha)/\sin(\beta)$	$-\sin(\alpha)/\cos(\beta)$	$\sin(\beta - \alpha)$
H	$\sin(\alpha)/\sin(\beta)$	$\cos(\alpha)/\cos(\beta)$	$\cos(\beta - \alpha)$
A	$\tan^{-1}(\beta)$	$\tan(\beta)$	0

mixing angle α :
$$\begin{pmatrix} h \\ H \end{pmatrix} = \begin{pmatrix} \cos(\alpha) & -\sin(\alpha) \\ \sin(\alpha) & \cos(\alpha) \end{pmatrix} \begin{pmatrix} H_1^0 \\ H_2^0 \end{pmatrix}$$

$$t_\beta \uparrow \Rightarrow g_u^\phi \downarrow \quad g_d^\phi \uparrow$$

Higgs Boson Decay

- low $\tan(\beta)$

mainly: $h \rightarrow \tau^+ \tau^- / b\bar{b}$

large m_h : $h \rightarrow gg / \gamma\gamma / WW / ZZ$

low m_A : $H, A \rightarrow \tau^+ \tau^- / b\bar{b}$

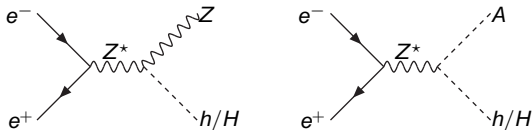
large m_A : $H, A \rightarrow t\bar{t}$ $H \rightarrow WW / ZZ / hh / \tilde{f}\tilde{f}$ $A \rightarrow Zh$

- large $\tan(\beta)$

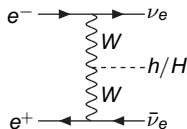
$h, H, A \rightarrow \tau^+ \tau^- / [\mu^+ \mu^-]$ $h, H, A \rightarrow b\bar{b}$

Higgs Production at LEP

- Higgs-strahlung and pair production

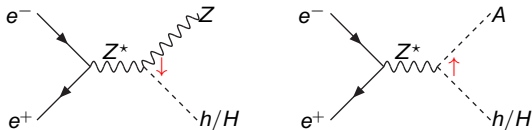


- Fusion process

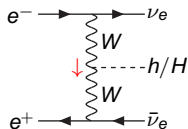


Higgs Production at LEP

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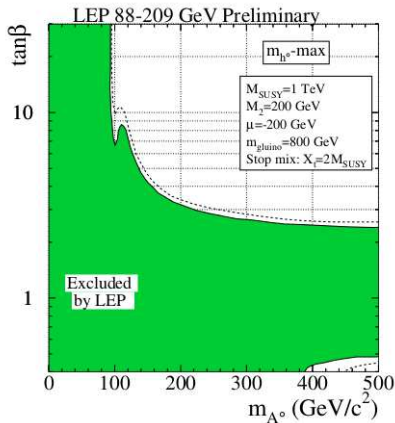


- Fusion process



$\tan(\beta) \uparrow$

Exclusion Limits from LEP



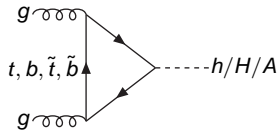
$$E_{\text{CMS}} = 209 \text{ GeV}$$

- $m_A < 93.4 \text{ GeV}$ and $m_{h,H} < 92.8 \text{ GeV}$ excluded
- $0.7 < \tan(\beta) < 2.4$ excluded

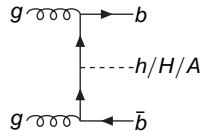
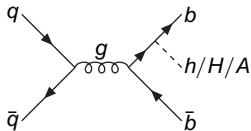
A linear collider with $E_{\text{CMS}} \gtrsim 250 \text{ GeV}$ could access the whole parameter space.

Higgs Production at the LHC

- Gluon fusion

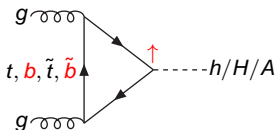


- Higgs bremsstrahlung off bottom quarks

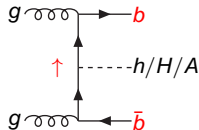
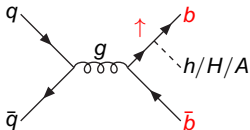


Higgs Production at the LHC

- Gluon fusion

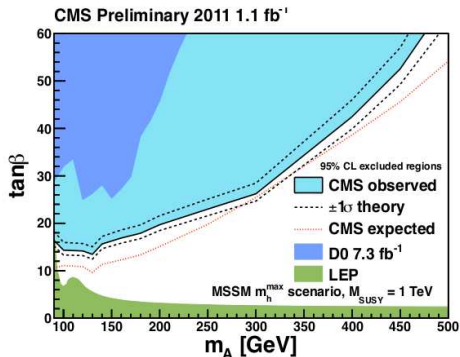


- Higgs bremsstrahlung off bottom quarks



$\tan(\beta) \uparrow$

Exclusion Limits from the LHC



Cross section is enhanced for

- large $\tan(\beta)$
- low m_A

High m_A and moderate $\tan(\beta) \rightarrow h = H_{SM}$
 \Rightarrow hard to detect because of huge $b\bar{b}$ background

Motivation for Calculating Two-Loop Corrections to the Couplings

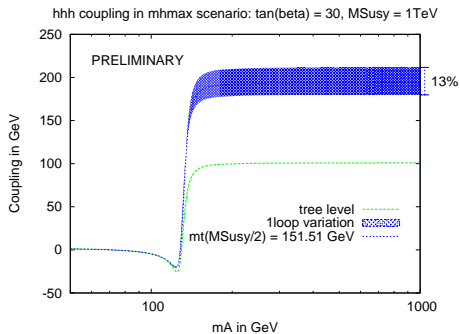
- Higgs self-couplings determine Higgs potential
- Higgs potential is responsible for Electro Weak Symmetry Breaking (EWSB)

⇒ need to measure Higgs self-interactions to understand EWSB

(very difficult at LHC, linear collider needed)

⇒ need high-precision predictions for trilinear couplings

Existing One-Loop Calculation



$$\lambda_{hhh}^{tree} = \frac{3m_Z^2}{v} c_{2\alpha} s_{\alpha+\beta}$$

- large corrections
- sizable uncertainties

⇒ two-loop calculation needed.

Barger, . . . (1992)

Effective Potential Method

Effective Potential V^{eff} :

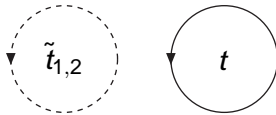
- Non-derivative part of the effective action
→ correct in the limit of **vanishing external momenta**
- Generating functional of 1PI Greens functions
with no external legs (**vacuum diagrams**)
- n -th derivative of V^{eff} :
sum of all 1PI diagrams with n external legs

$$\Rightarrow \lambda(h_1, h_2, h_3) = \left. \frac{\partial^3 V^{eff}}{\partial h_1 \partial h_2 \partial h_3} \right|_{min}$$

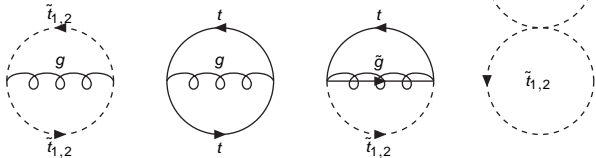
Computing the Effective Potential

First step: calculate ΔV^{α_t} and $\Delta V^{\alpha_t \alpha_s}$

1loop:
 $\mathcal{O}(\alpha_t)$



2loop:
 $\mathcal{O}(\alpha_t \alpha_s)$



Zhang, ... (1999)
 Slavich, ... (2001)

Renormalization

The fully renormalized coupling can be calculated by

$$\lambda_{\alpha_t \alpha_s}(h_1, h_2, h_3) = \left. \frac{\partial^3 (V_0 + \Delta V^{\alpha_t} + \Delta V^{\alpha_t \alpha_s})}{\partial h_1 \partial h_2 \partial h_3} \right|_{min} + \delta\lambda_{CT}.$$

The counterterm is obtained from derivatives

$$\delta\lambda_{CT}^{(2)} = \sum_i \frac{\partial \Delta\lambda_1}{\partial x_i} \delta x_i,$$

where $x_i = \{m_t^2, m_{\tilde{t}_1}^2, m_{\tilde{t}_2}^2, A_t\}$ are all parameters of the one-loop couplings that are renormalized at $\mathcal{O}(\alpha_s)$.

Cancellation of Divergences

For simplicity, start with \overline{DR} -scheme:

- \overline{DR} -counterterms $\delta^{\overline{DR}}$ are $\frac{1}{\epsilon}$ -divergences
- $\mathcal{O}(\epsilon)$ -terms in $\Delta V_1^{\alpha t}$ give finite contributions
- $\mathcal{O}(\epsilon^0)$ -terms in $\Delta V_1^{\alpha t}$ give $\frac{1}{\epsilon}$ poles

Non-trivial consistency check: All $\frac{1}{\epsilon^2}$ and $\frac{1}{\epsilon}$ poles cancel.

$\Rightarrow \lambda_{\alpha_t \alpha_s}$ is finite

Renormalization Scheme

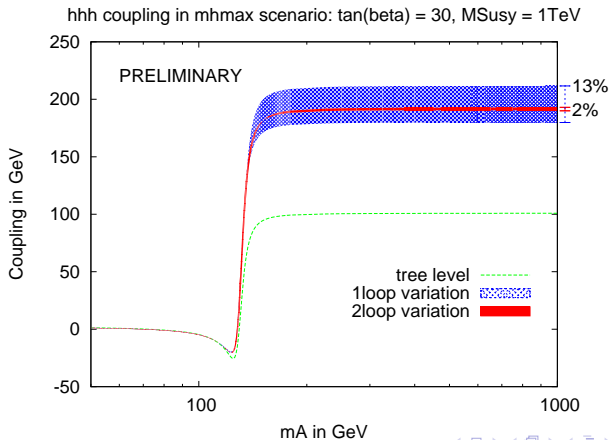
Can shift to any other scheme by adding finite counterterm.
e.g. on-shell-scheme:

$$\lambda_{\alpha_t \alpha_s}^{\text{OS}} = \lambda_{\alpha_t \alpha_s}^{\overline{\text{DR}}} + \Delta\lambda_{\text{CT}}^{\text{OS}}$$

- $\Delta\lambda_{\text{CT}}^{\text{OS}} = \sum_i \frac{\partial \Delta\lambda_1}{\partial x_i} \Delta^{\text{OS}} x_i$
- $\Delta^{\text{OS}} x_i$: finite part of on-shell counterterm

$\Rightarrow \lambda_{\alpha_t \alpha_s}^{\text{OS}}$ is independent of the 't Hooft scale Q .

Results: hhh



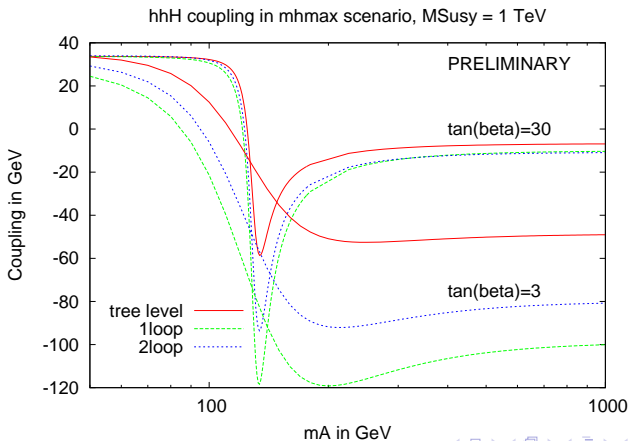
Summary

- The **effective potential method** provides an efficient way to calculate two-loop corrections to Higgs self-interactions.
- The $\mathcal{O}(\alpha_t\alpha_s)$ corrections to the hhh-coupling **are small at the central scale $M_{SUSY}/2$ and the theoretical uncertainty is reduced from $\sim 15\%$ to $\sim 2\%$.**

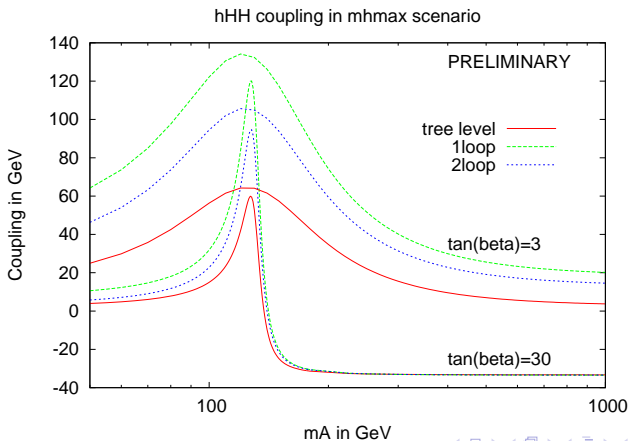
⇒ stabilization

- Outlook:
 - quartic couplings → done ✓
 - $\mathcal{O}(\alpha_t^2)$ corrections → in progress
 - analytic formulae, public code

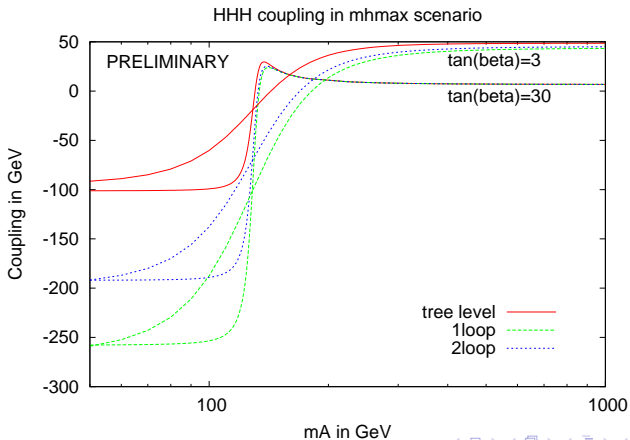
Results: hhH



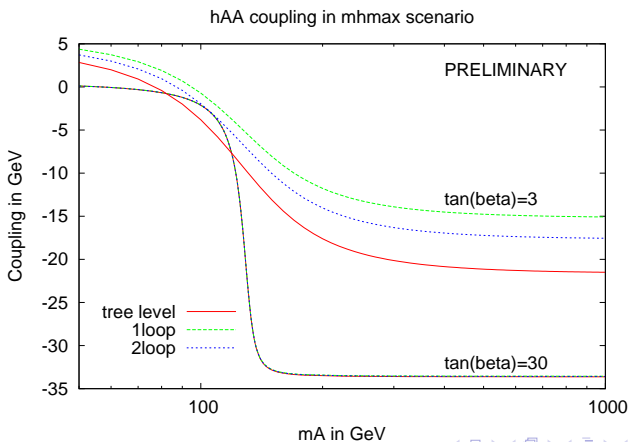
Results: hHH



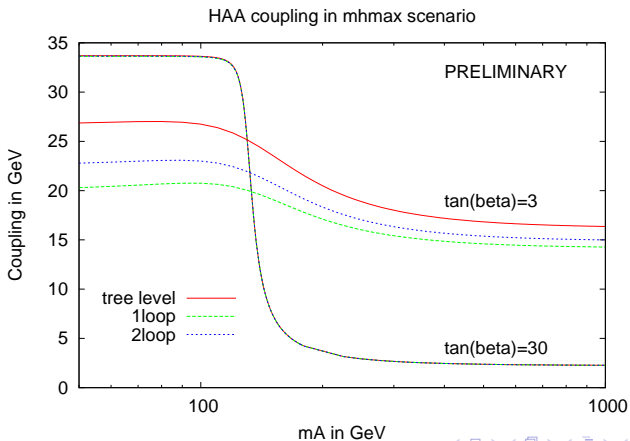
Results: HHH



Results: hAA



Results: HAA



Results: hhh

