

# Full Run 2 overview of electro-weak multiboson production in CMS

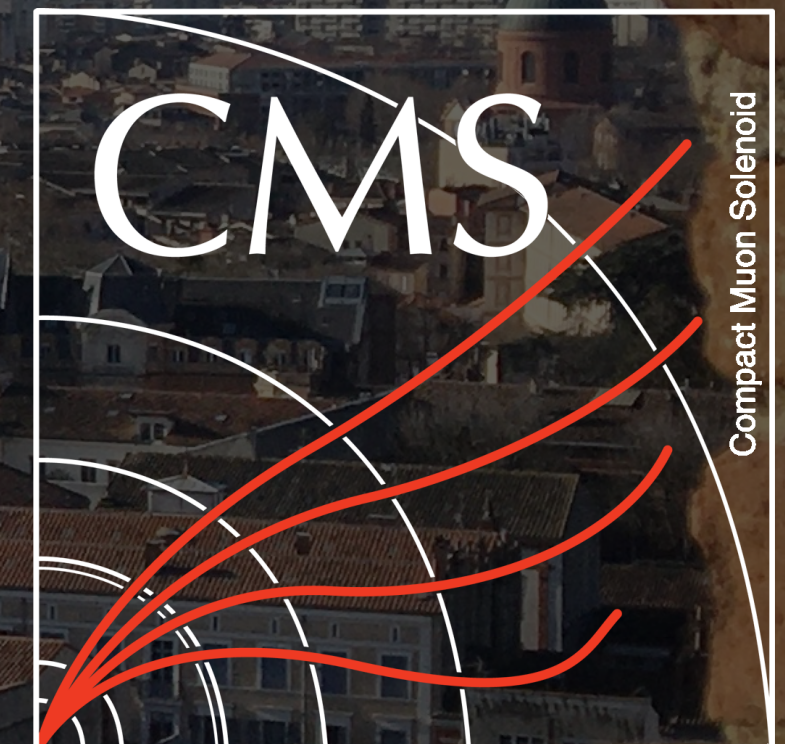
**SMU**



**Alexander von Humboldt**  
Stiftung/Foundation



**WAYNE STATE**  
UNIVERSITY



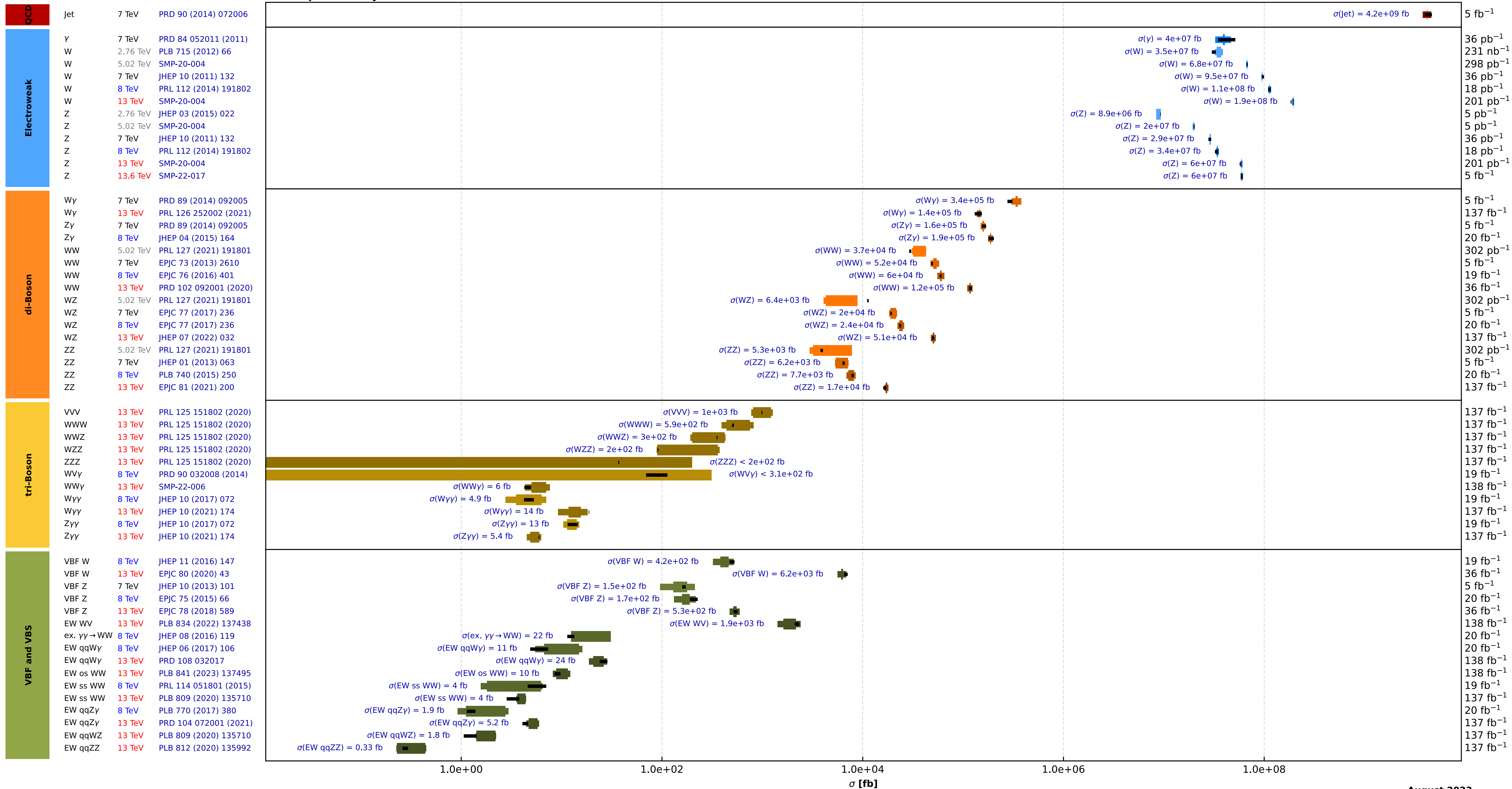
**Saptaparna Bhattacharya**  
**Multiboson Interactions**

September 25-28th, 2024

# Overview of CMS cross section results

CMS preliminary

$3 \mu\text{b}^{-1} - 138 \text{fb}^{-1}$  (2.76, 5.02, 7, 8, 13, 13.6 TeV)



Measured cross sections and exclusion limits at 95% C.L. See here for all cross section summary plots

Inner colored bars statistical uncertainty, outer narrow bars statistical+systematic uncertainty Light to Dark colored bars: 2.76, 5.02, 7, 8, 13, 13.6 TeV, Black bars: theory prediction

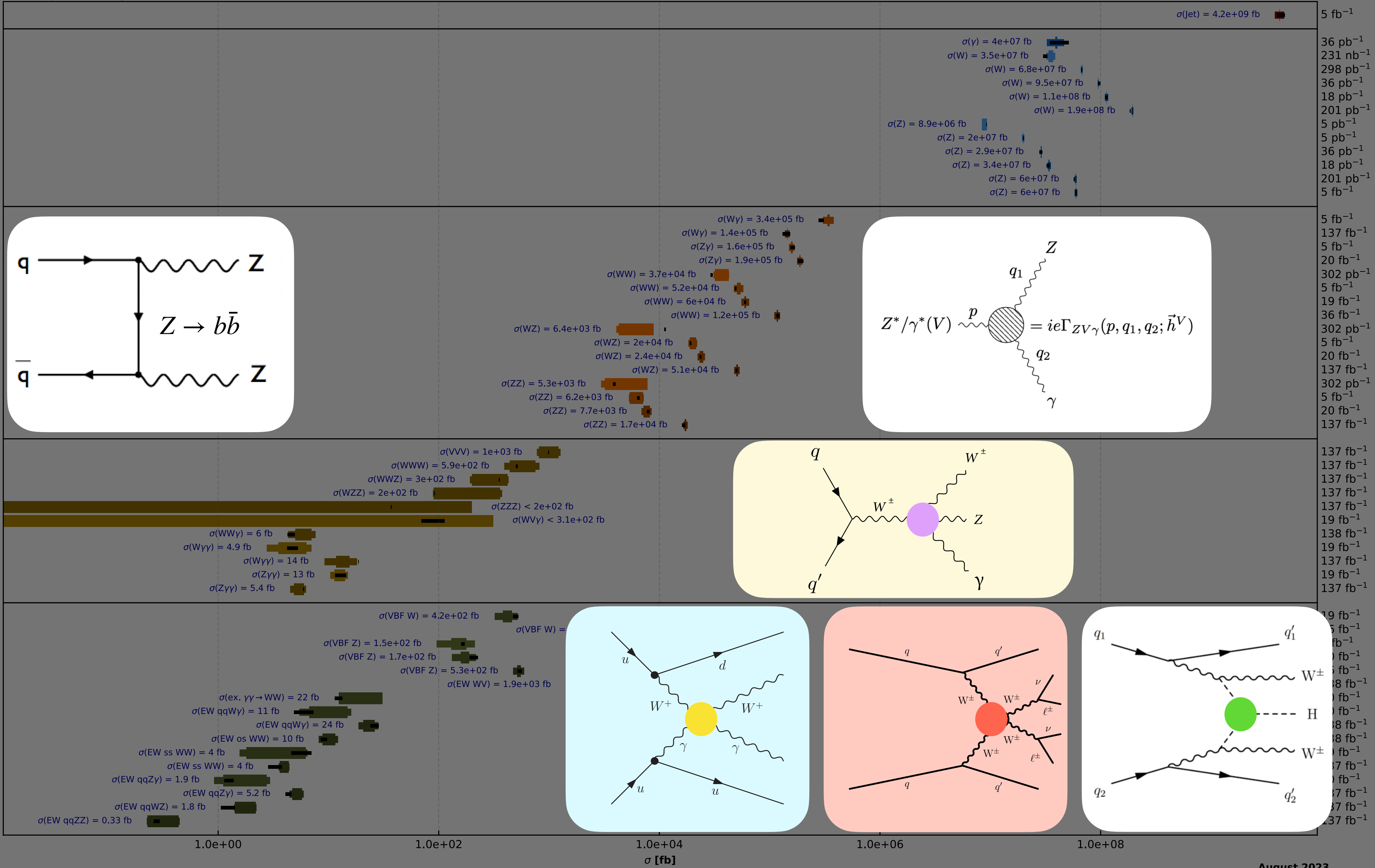
August 2023

# Overview of CMS cross section results

CMS preliminary

$3 \mu\text{b}^{-1} - 138 \text{fb}^{-1}$  (2.76, 5.02, 7, 8, 13, 13.6 TeV)

Category	Process	Energy	Reference
QCD	Jet	7 TeV	PRD 90 (2014) 072006
Electroweak	$\gamma$	7 TeV	PRD 84 052011 (2011)
	W	2.76 TeV	PLB 715 (2012) 66
	W	5.02 TeV	SMP-20-004
	W	7 TeV	JHEP 10 (2011) 132
	W	8 TeV	PRL 112 (2014) 191802
	W	13 TeV	SMP-20-004
	Z	2.76 TeV	JHEP 03 (2015) 022
	Z	5.02 TeV	SMP-20-004
	Z	7 TeV	JHEP 10 (2011) 132
	Z	8 TeV	PRL 112 (2014) 191802
Z	13 TeV	SMP-20-004	
Z	13.6 TeV	SMP-22-017	
di-Boson	W $\gamma$	7 TeV	PRD 89 (2014) 092005
	W $\gamma$	13 TeV	PRL 126 252002 (2021)
	Z $\gamma$	7 TeV	PRD 89 (2014) 092005
	Z $\gamma$	8 TeV	JHEP 04 (2015) 164
	WW	5.02 TeV	PRL 127 (2021) 191801
	WW	7 TeV	EPJC 73 (2013) 2610
	WW	8 TeV	EPJC 76 (2016) 401
	WW	13 TeV	PRD 102 092001 (2020)
	WZ	5.02 TeV	PRL 127 (2021) 191801
	WZ	7 TeV	EPJC 77 (2017) 236
	WZ	8 TeV	EPJC 77 (2017) 236
	WZ	13 TeV	JHEP 07 (2022) 032
	ZZ	5.02 TeV	PRL 127 (2021) 191801
	ZZ	7 TeV	JHEP 01 (2013) 063
	ZZ	8 TeV	PLB 740 (2015) 250
ZZ	13 TeV	EPJC 81 (2021) 200	
tri-Boson	VVV	13 TeV	PRL 125 151802 (2020)
	WWW	13 TeV	PRL 125 151802 (2020)
	WWZ	13 TeV	PRL 125 151802 (2020)
	WZZ	13 TeV	PRL 125 151802 (2020)
	ZZZ	13 TeV	PRL 125 151802 (2020)
	WW $\gamma$	8 TeV	PRD 90 032008 (2014)
	WW $\gamma$	13 TeV	SMP-22-006
	W $\gamma\gamma$	8 TeV	JHEP 10 (2017) 072
	W $\gamma\gamma$	13 TeV	JHEP 10 (2021) 174
	Z $\gamma\gamma$	8 TeV	JHEP 10 (2017) 072
Z $\gamma\gamma$	13 TeV	JHEP 10 (2021) 174	
VBF and VBS	VBF W	8 TeV	JHEP 11 (2016) 147
	VBF W	13 TeV	EPJC 80 (2020) 43
	VBF Z	7 TeV	JHEP 10 (2013) 101
	VBF Z	8 TeV	EPJC 75 (2015) 66
	VBF Z	13 TeV	EPJC 78 (2018) 589
	EW WW	13 TeV	PLB 834 (2022) 137438
	ex. $\gamma\gamma \rightarrow WW$	8 TeV	JHEP 08 (2016) 119
	EW qqW $\gamma$	8 TeV	JHEP 06 (2017) 106
	EW qqW $\gamma$	13 TeV	PRD 108 032017
	EW os WW	13 TeV	PLB 841 (2023) 137495
	EW ss WW	8 TeV	PRL 114 051801 (2015)
	EW ss WW	13 TeV	PLB 809 (2020) 135710
	EW qqZ $\gamma$	8 TeV	PLB 770 (2017) 380
	EW qqZ $\gamma$	13 TeV	PRD 104 072001 (2021)
	EW qqWZ	13 TeV	PLB 809 (2020) 135710
EW qqZZ	13 TeV	PLB 812 (2020) 135992	

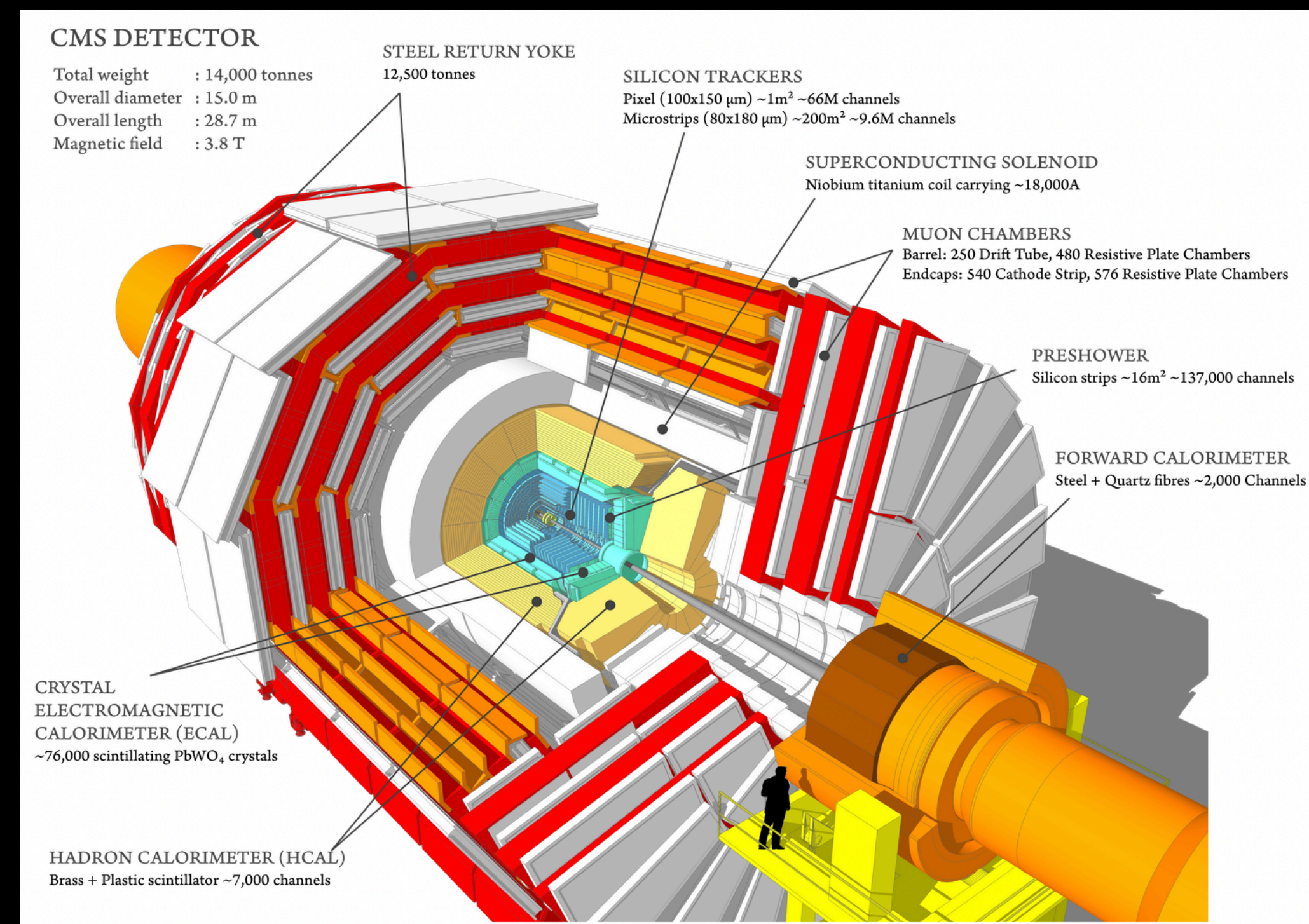


Measured cross sections and exclusion limits at 95% C.L. See here for all cross section summary plots

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# Outline

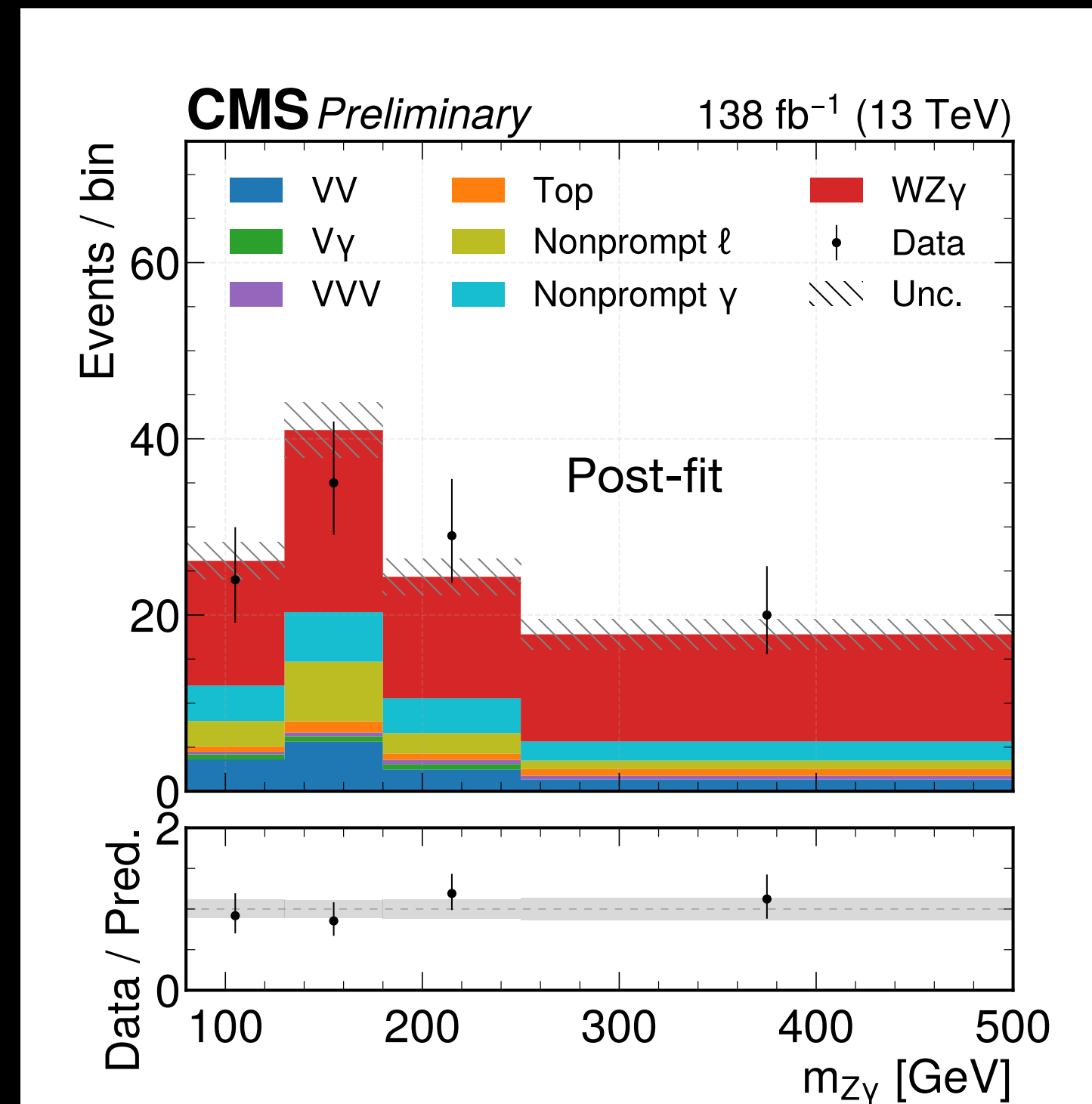
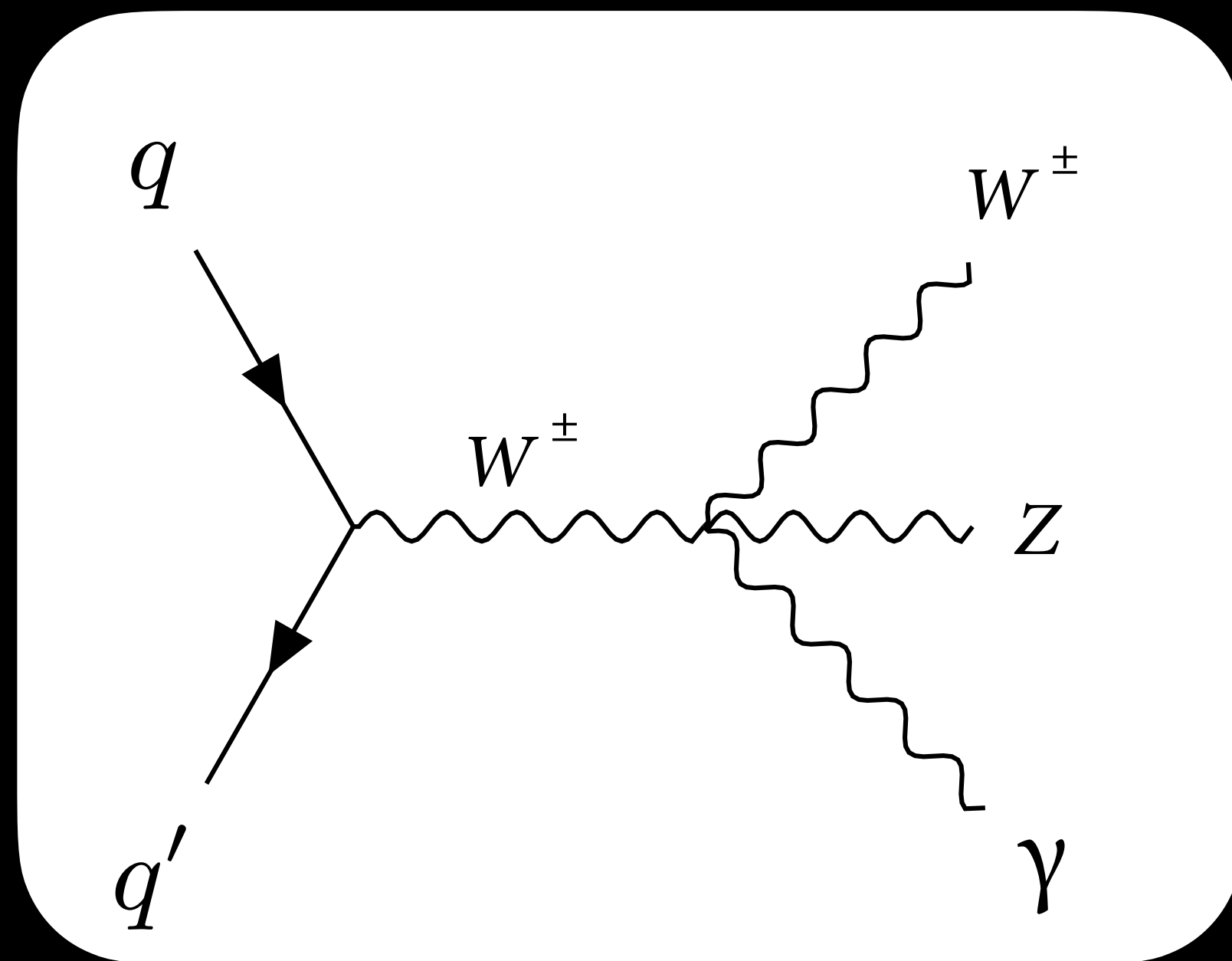
- Observation of  $WZ\gamma$  and  $WW\gamma$
- Exploring the implication of new physics modifying the quartic coupling in the Standard Model — effective field theory (EFT) interpretation
- The  $\gamma\gamma \rightarrow \tau\tau$  process and measuring  $g_\tau - 2$
- Tagging the protons with the Precision Proton Spectrometer
- Evolution of vector boson scattering (VBS) measurements
  - Same-signed  $WW$  with a hadronic  $\tau$
  - $W\gamma$  with 2 jets
  - $HHWW$  coupling
  - Future prospects
- Dibosons
  - Neutral couplings and  $ZZ/HZ, Z \rightarrow b\bar{b}, H \rightarrow b\bar{b}$



Brass + Plastic scintillator  $\sim 7,000$  channels  
HADRON CALORIMETER (HCAL)

# Observation of $WZ\gamma$ production at $\sqrt{s} = 13$ TeV

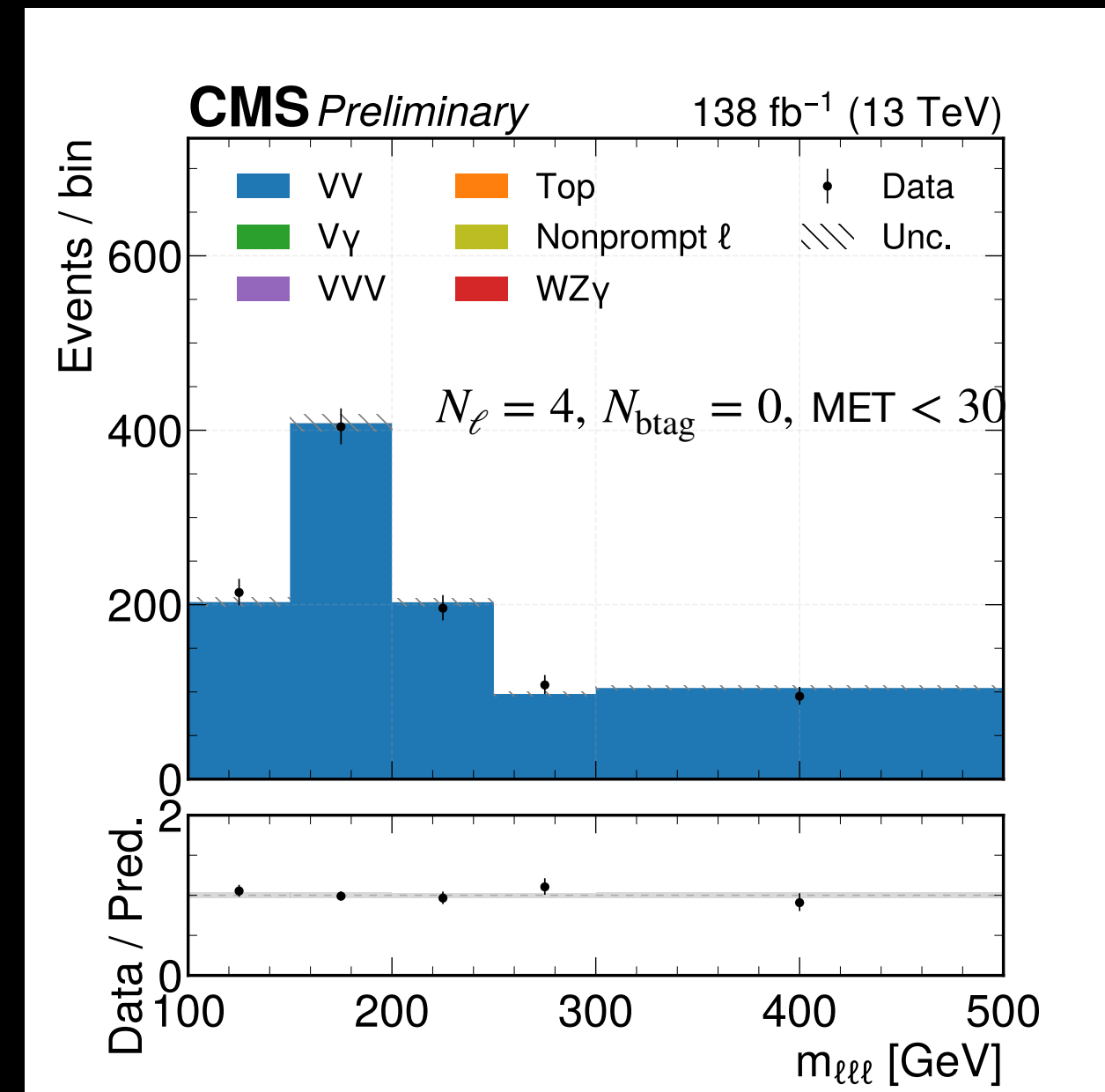
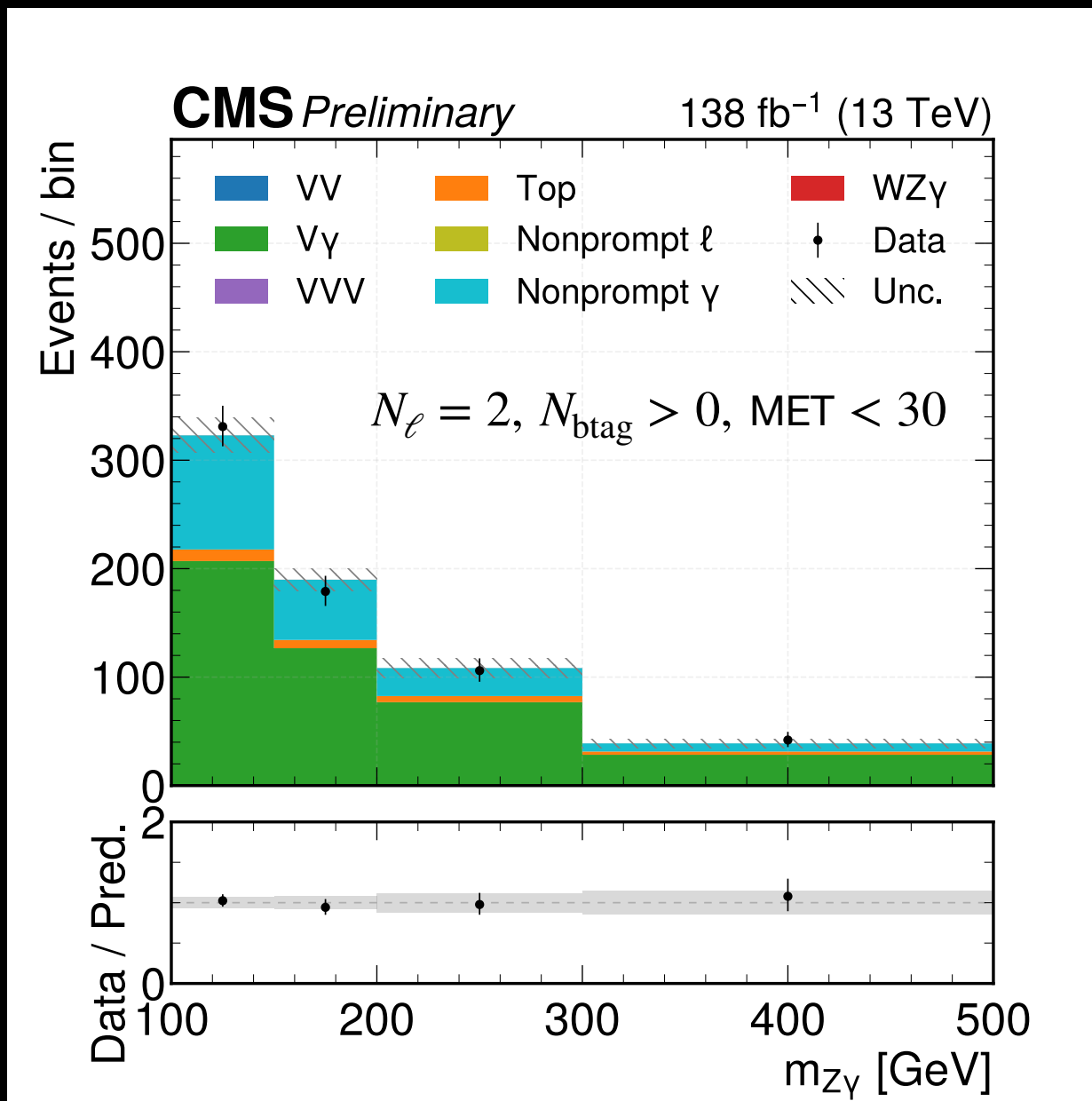
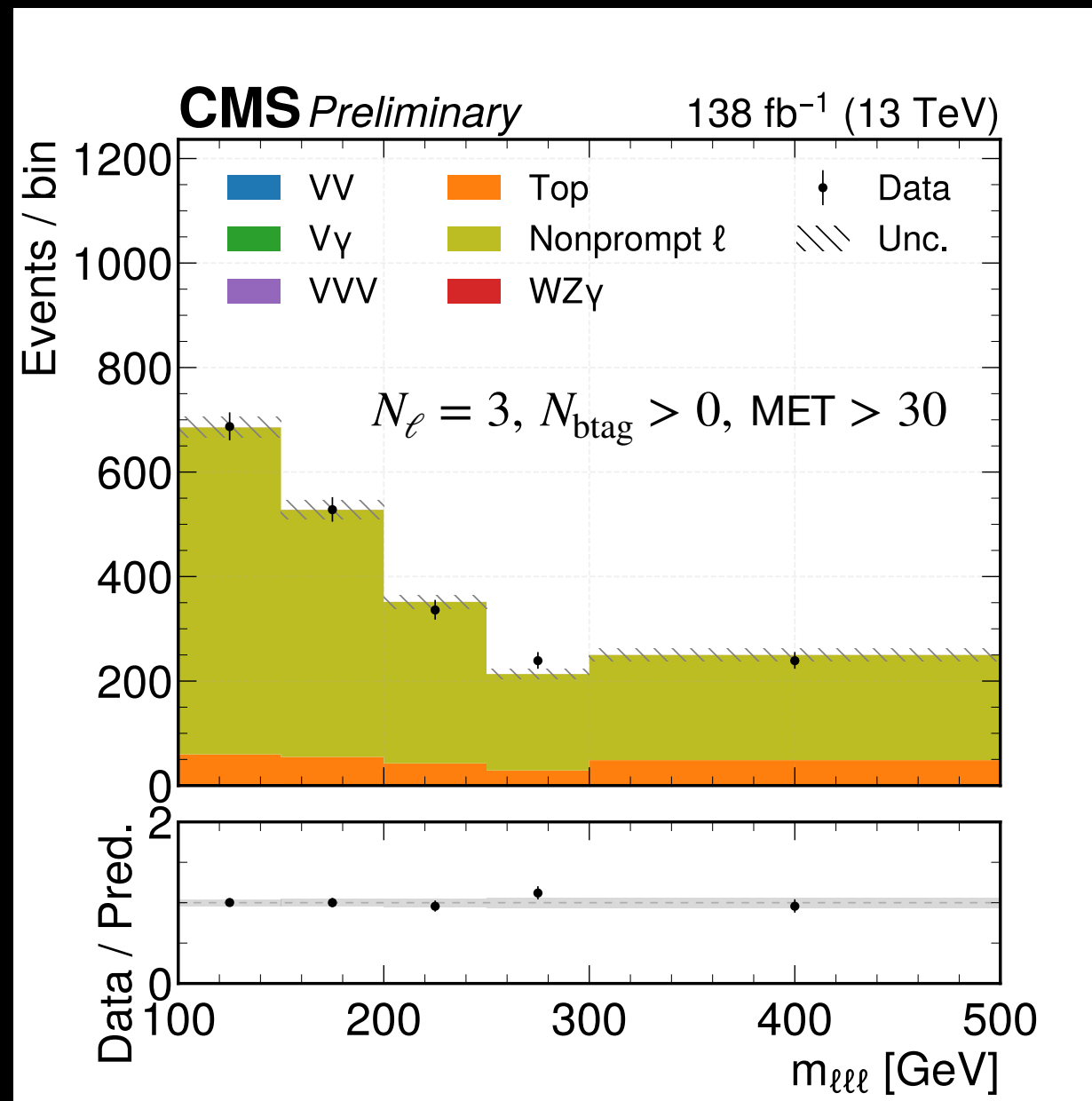
- $WZ\gamma$  process observed (expected) with a significance of 5.4 (3.8)  $\sigma$
- Final state defined by requiring three charged leptons ( $WZ \rightarrow \ell\nu\ell\ell$ ) and a photon
- Fiducial cross section measured as:  $5.48 \pm 1.11$  fb
- Several new physics scenarios explored including axions-like particles



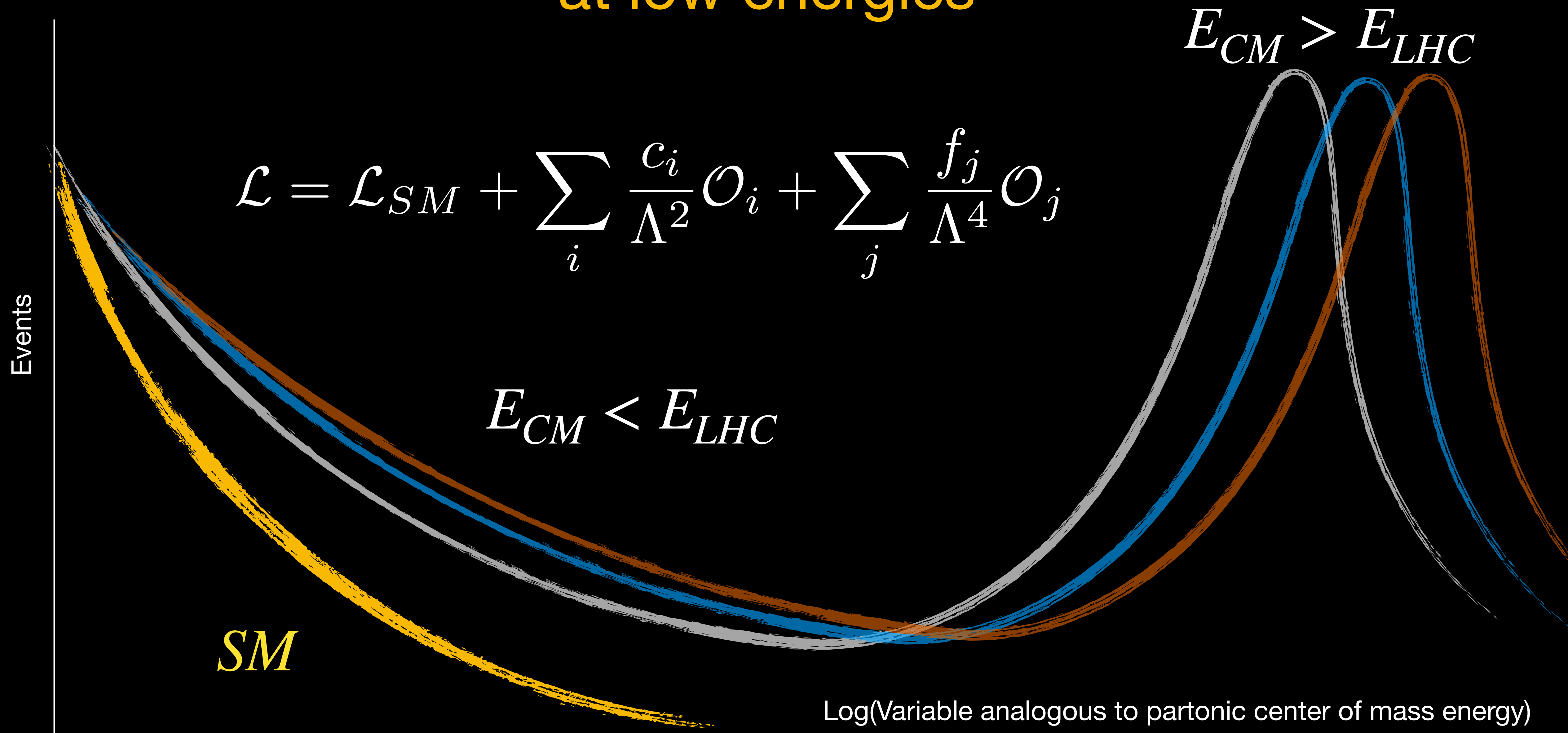
# Observation of $WZ\gamma$ production at $\sqrt{s} = 13$ TeV

- Simultaneous fit of signal and background regions
- Major background sources
  - Non-prompt photon and leptons
  - Prompt ZZ with an ISR photon

Process	SR	Nonprompt $\ell$ CR	Nonprompt $\gamma$ CR	ZZ CR
VV	$13.0 \pm 0.3$	$1.86 \pm 0.12$	$0.16 \pm 0.02$	$1016 \pm 12$
VVV	$0.69 \pm 0.05$	$0.36 \pm 0.11$	$0.01 \pm 0.01$	$0.10 \pm 0.04$
V $\gamma$	$1.38 \pm 0.76$	$4.66 \pm 2.05$	$438 \pm 27$	$0.01 \pm 0.01$
Top	$3.34 \pm 0.55$	$227 \pm 15$	$27.0 \pm 5.9$	$0.30 \pm 0.04$
Nonprompt $\ell$	$12.9 \pm 2.8$	$1792 \pm 34$	$<0.1$	$<0.1$
Nonprompt $\gamma$	$15.8 \pm 2.2$	$<0.1$	$195 \pm 19$	$<0.1$
WZG signal	$60.8 \pm 3.5$	$0.66 \pm 0.01$	$0.20 \pm 0.04$	$0.02 \pm 0.01$
Total background	$48.5 \pm 3.7$	$2027 \pm 33$	$660 \pm 21$	$1016 \pm 12$
Total prediction	$109 \pm 5$	$2027 \pm 33$	$660 \pm 21$	$1016 \pm 12$
Observed	108	2029	658	1017



# New physics at high energies manifesting at low energies



# BSM implication of $WZ\gamma$

<https://arxiv.org/pdf/2004.05174>

EFT interpretation of dimension-8 operators featuring  $SU(2)_L$  and  $U(1)_Y$  field strength

Operators	Observed limits [ TeV <sup>-4</sup> ]	Expected limits [ TeV <sup>-4</sup> ]	Unitarity bound [ TeV ]
$F_{T,0}/\Lambda^4$	[-2.60, 2.60]	[-2.52, 2.52]	1.32
$F_{T,1}/\Lambda^4$	[-3.28, 3.24]	[-3.18, 3.14]	1.48
$F_{T,2}/\Lambda^4$	[-7.15, 7.05]	[-6.95, 6.85]	1.35
$F_{T,5}/\Lambda^4$	[-2.54, 2.56]	[-2.46, 2.50]	1.55
$F_{T,6}/\Lambda^4$	[-3.18, 3.22]	[-3.08, 3.14]	1.61
$F_{T,7}/\Lambda^4$	[-6.85, 7.05]	[-6.65, 6.85]	1.71

YITP-SB-2020-8

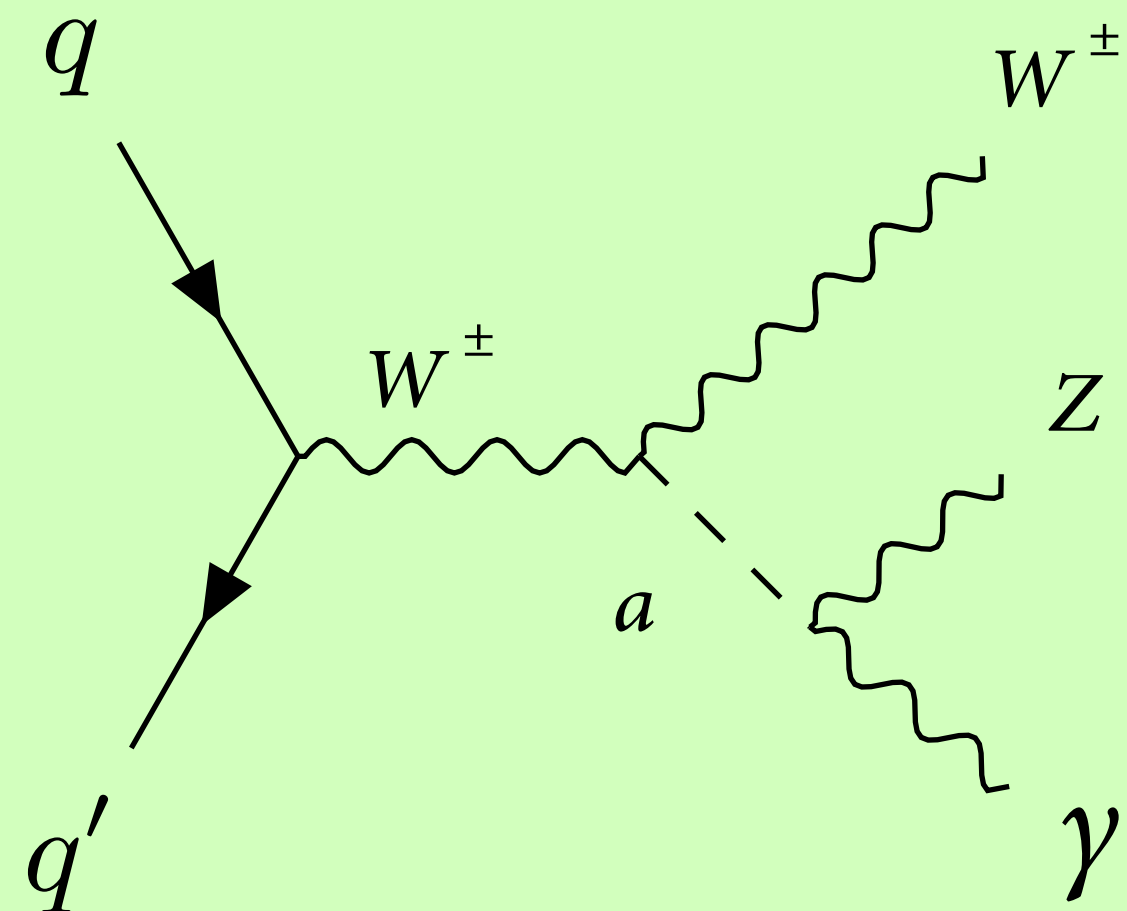
**Unitarity Constraints on Anomalous Quartic Couplings**

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*Instituto de Física, Universidade de São Paulo, São Paulo - SP, Brazil.*

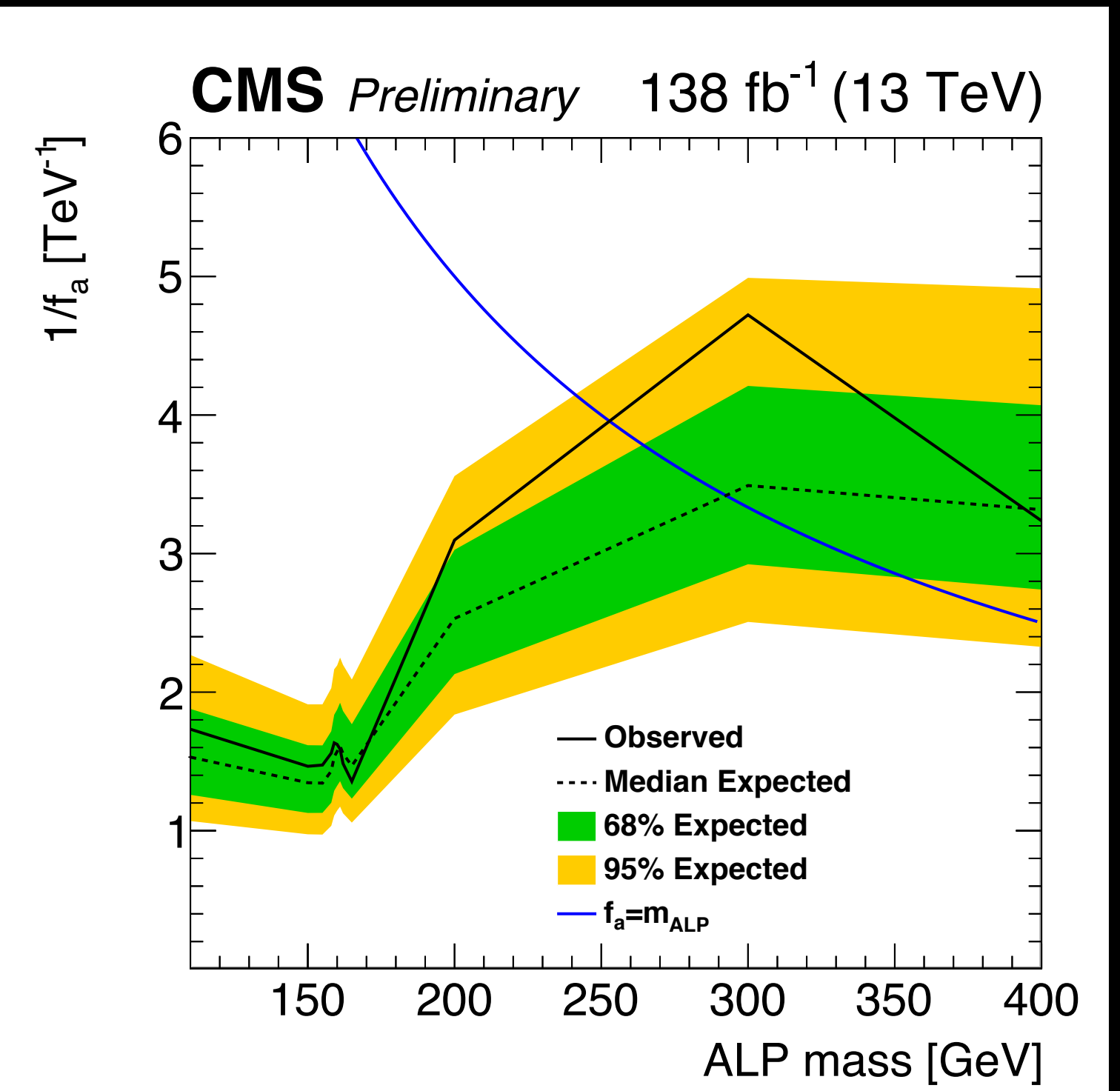
M. C. Gonzalez-Garcia<sup>3</sup>  
*Institució Catalana de Recerca i Estudis Avançats (ICREA),  
 Departament d'Estructura i Constituents de la Matèria,  
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 C.N. Yang Institute for Theoretical Physics, SUNY at Stony Brook, Stony Brook, NY 11794-3840, USA*

We obtain the partial-wave unitarity constraints on the lowest-dimension effective operators which generate anomalous quartic gauge couplings but leave the triple gauge couplings unaffected. We consider operator expansions with linear and nonlinear realizations of the electroweak symmetry and explore the multidimensional parameter space of the coefficients of the relevant operators: 20 dimension-eight operators in the linear expansion and 5  $O(p^4)$  operators in the derivative expansion. We study two-to-two scattering of electroweak gauge bosons and Higgs bosons taking into account all coupled channels and all possible helicity amplitudes for the  $J = 0, 1$  partial waves. In general, the bounds degrade by factors of a few when several operator coefficients are considered to be nonvanishing simultaneously. However, this requires considering constraints from both  $J = 0$  and  $J = 1$  partial waves for some sets of operators.

$$\mathcal{L} = \mathcal{L}_{SM} + \sum_i \frac{c_i}{\Lambda^2} \mathcal{O}_i + \sum_j \frac{f_j}{\Lambda^4} \mathcal{O}_j$$

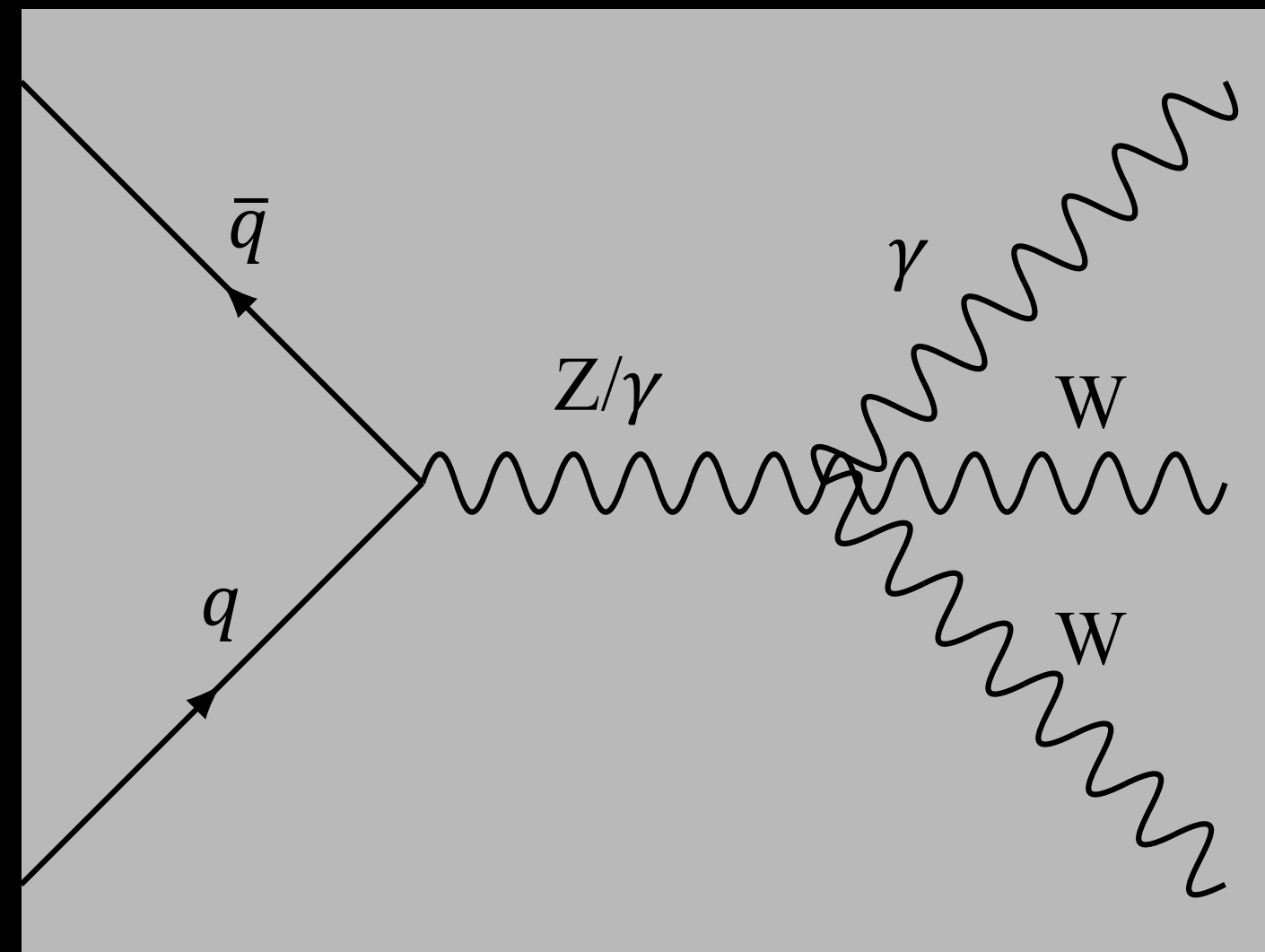


The  $WZ\gamma$  process is also sensitive to mediation by axion-like particles





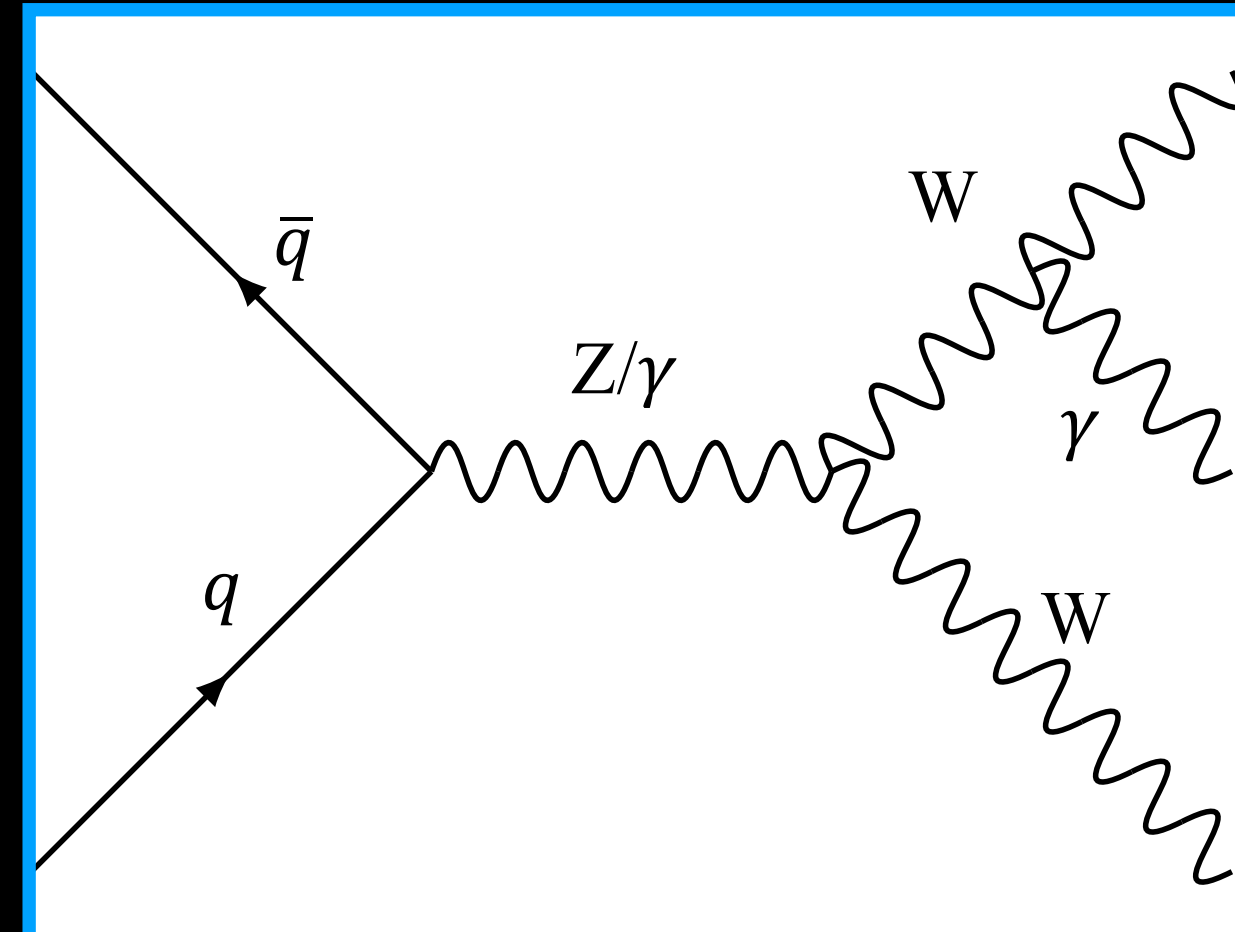
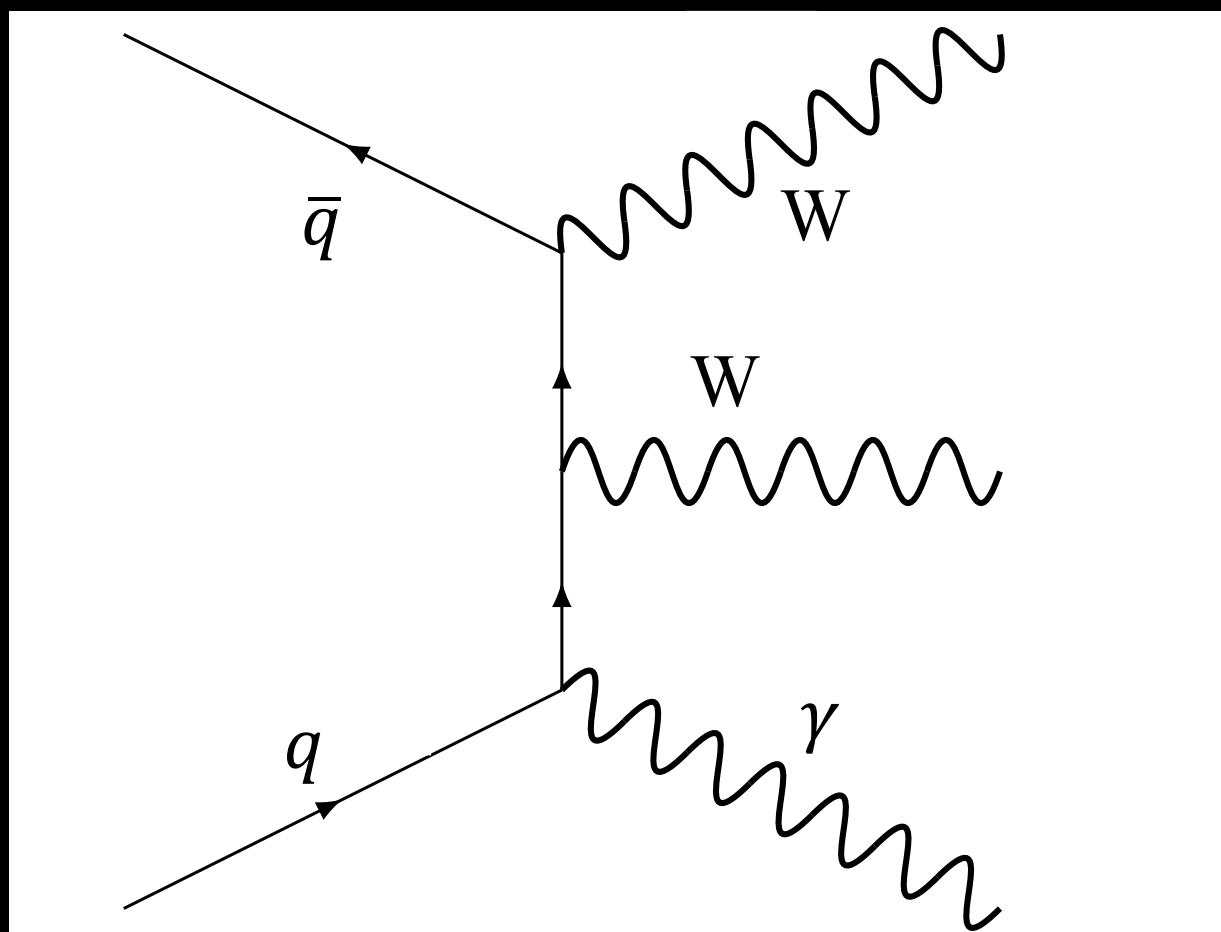
# Observation of $WW\gamma$ production at $\sqrt{s} = 13$ TeV



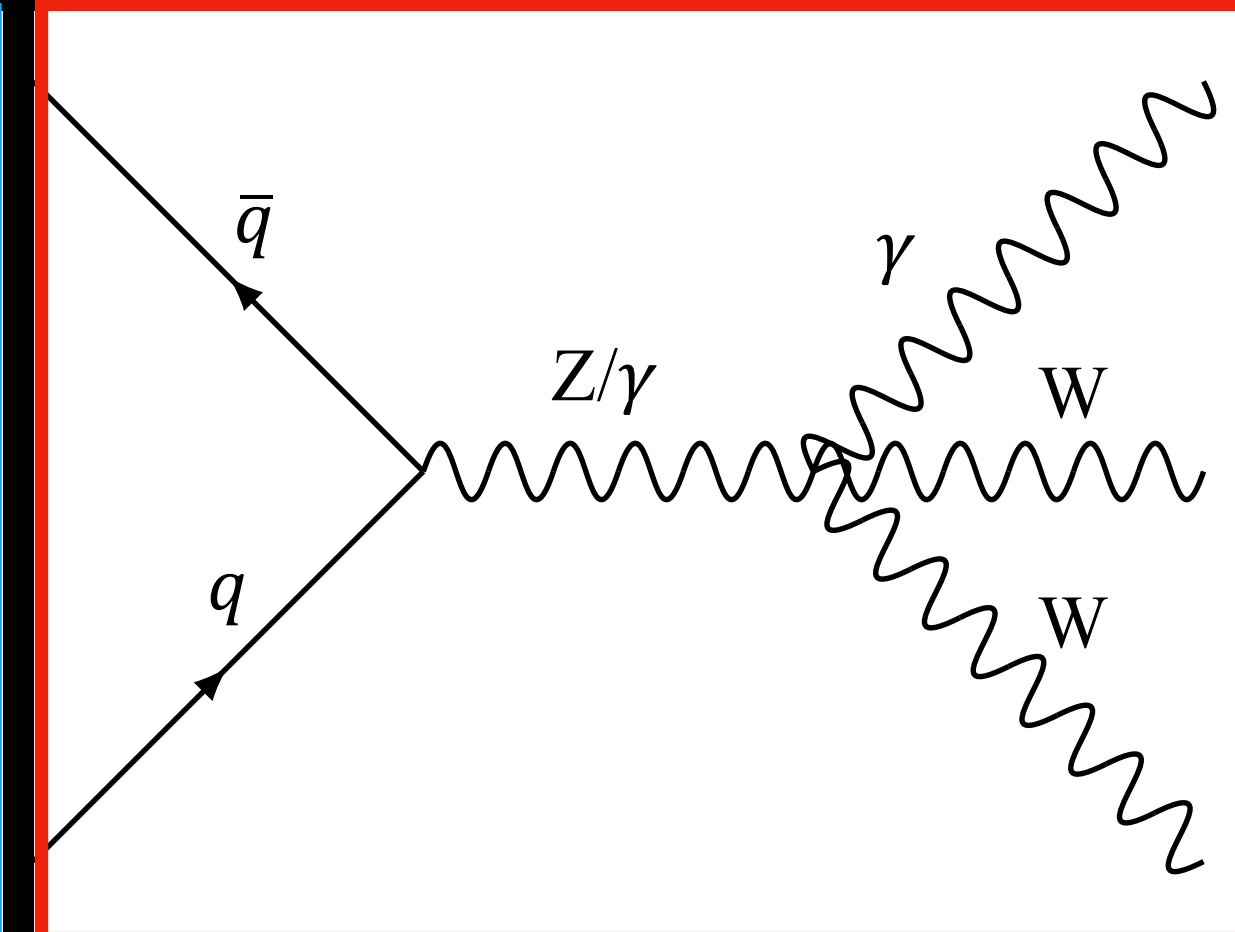
# Observation of $WW\gamma$ production at $\sqrt{s} = 13$ TeV

- $WW\gamma$  process observed (expected) with a significance of 5.6 (5.1)  $\sigma$
- Fiducial cross section measured as:  $5.9 \pm 0.8$  (stat.)  $\pm 0.8$  (syst.)  $\pm 0.7$  (modeling\*) fb
- Associated search for H with a photon explored  $\rightarrow$  generated by coupling of the Higgs boson to light quarks

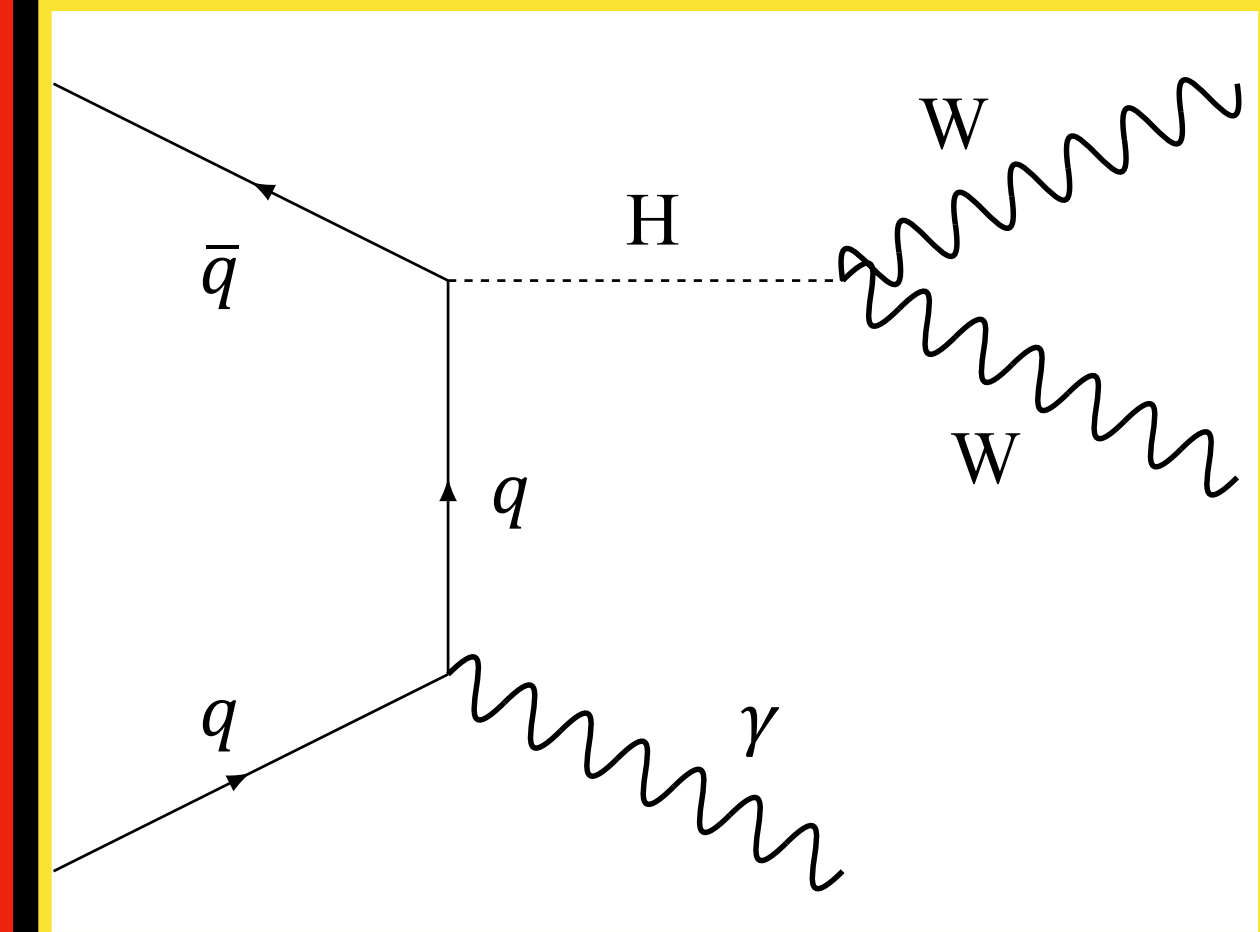
Phys. Rev. Lett. 132 (2024) 121901



Trilinear coupling



Quartic coupling



Associated production ( $H\gamma$ )

\*The theoretical modeling uncertainties include the renormalization and factorization of QCD scales, PDFs, and parton shower modeling

# Observation of $WW\gamma$ production at $\sqrt{s} = 13$ TeV

- $WW\gamma \rightarrow e^+\nu_e\mu^-\bar{\nu}_\mu\gamma$  or  $\mu^+\nu_\mu e^-\bar{\nu}_e\gamma$  final state

- Events with b-jets vetoed

- Additional loose leptons vetoed

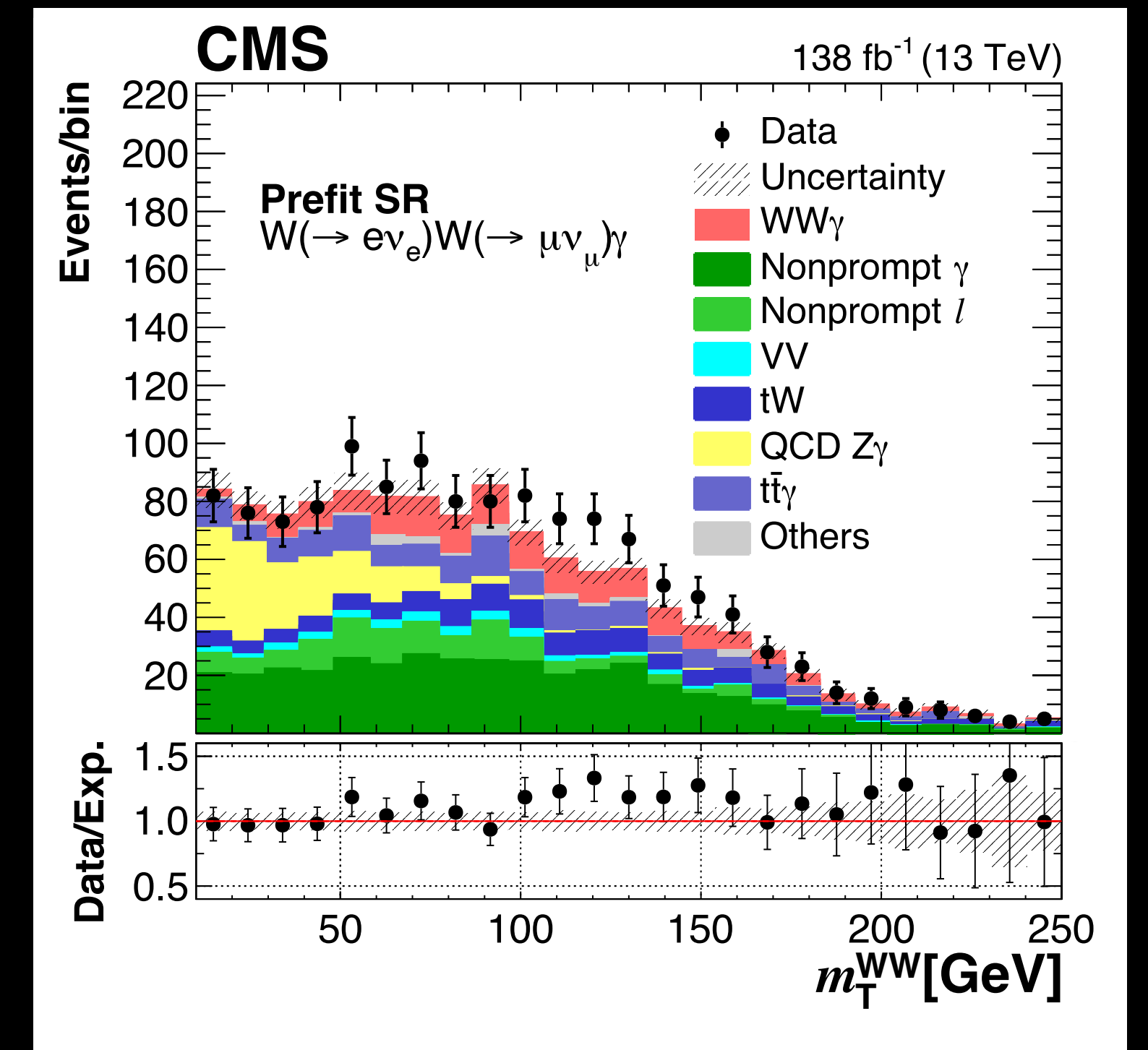
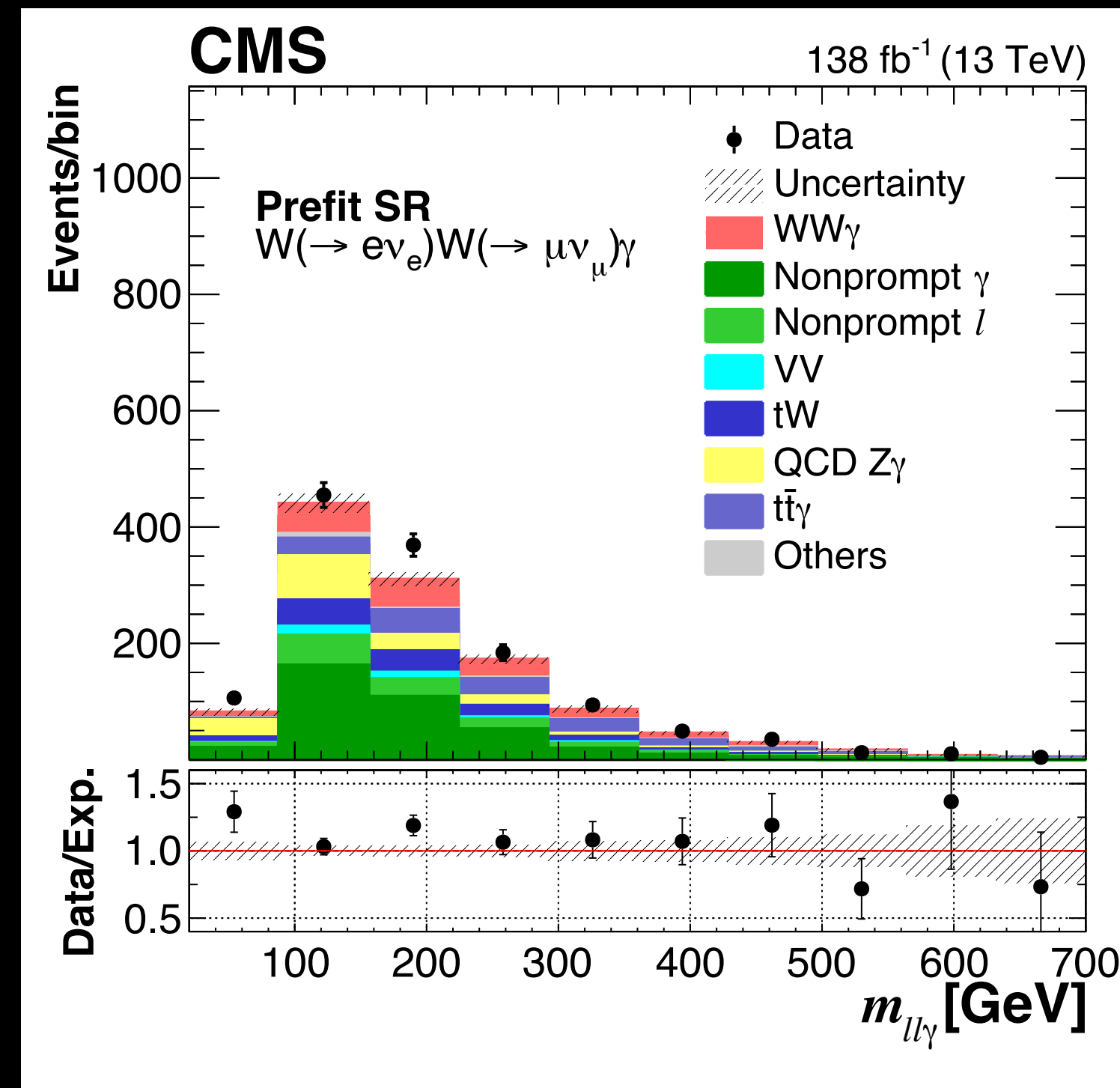
- Backgrounds suppressed by

- $M_{\ell\ell} > 10$  GeV

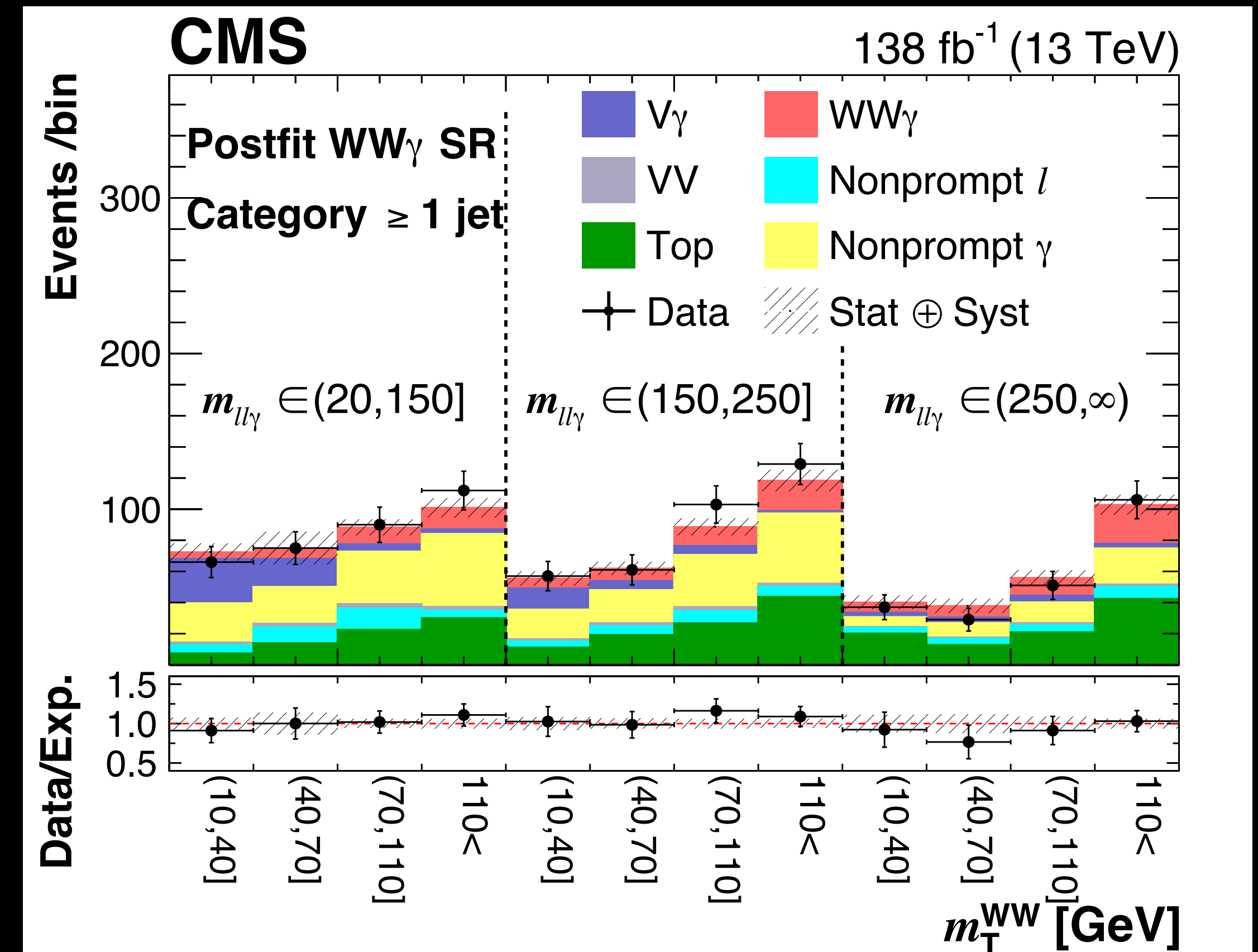
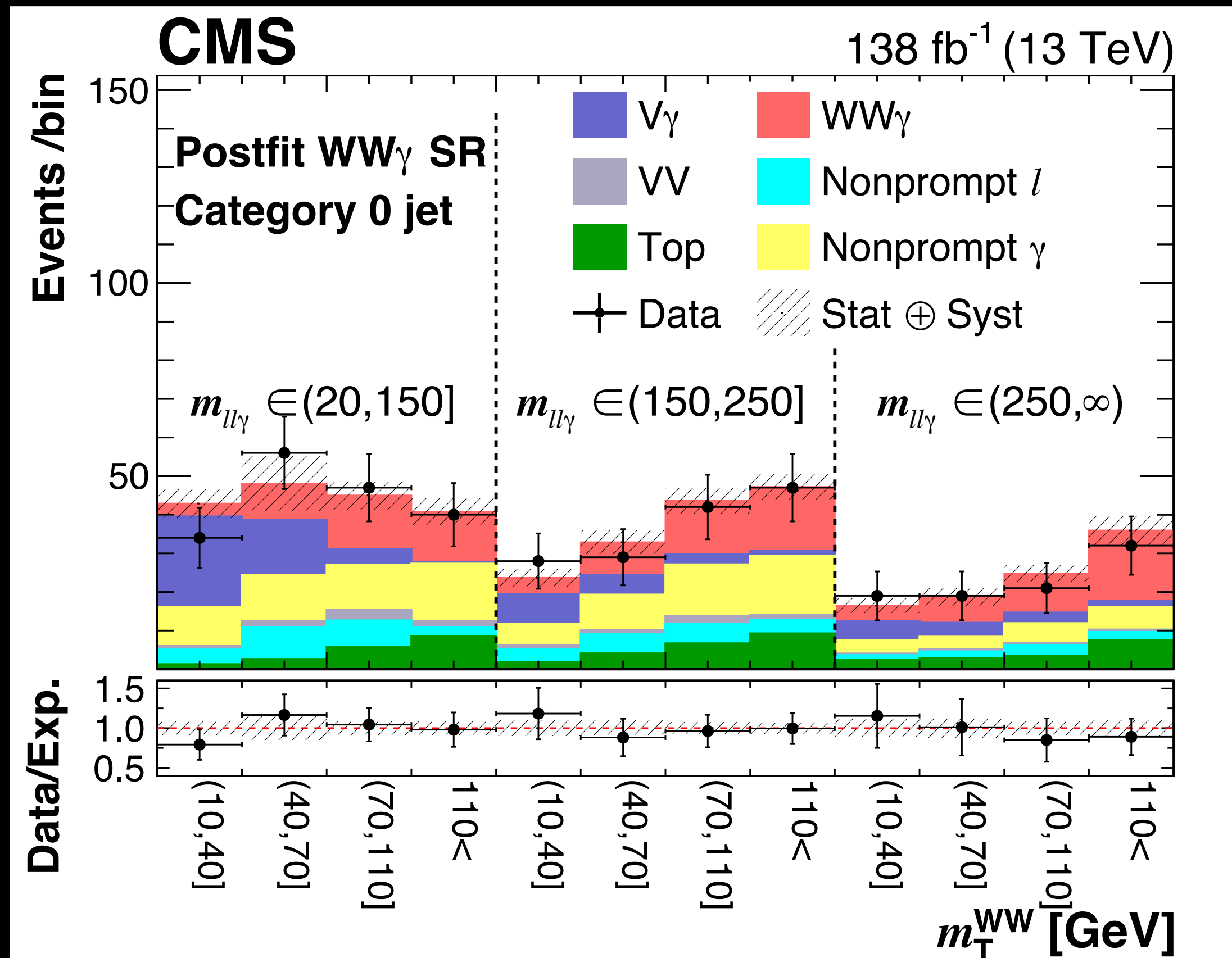
- $p_T^{\ell\ell} > 15$  GeV

- $m_T^{WW} > 10$  GeV

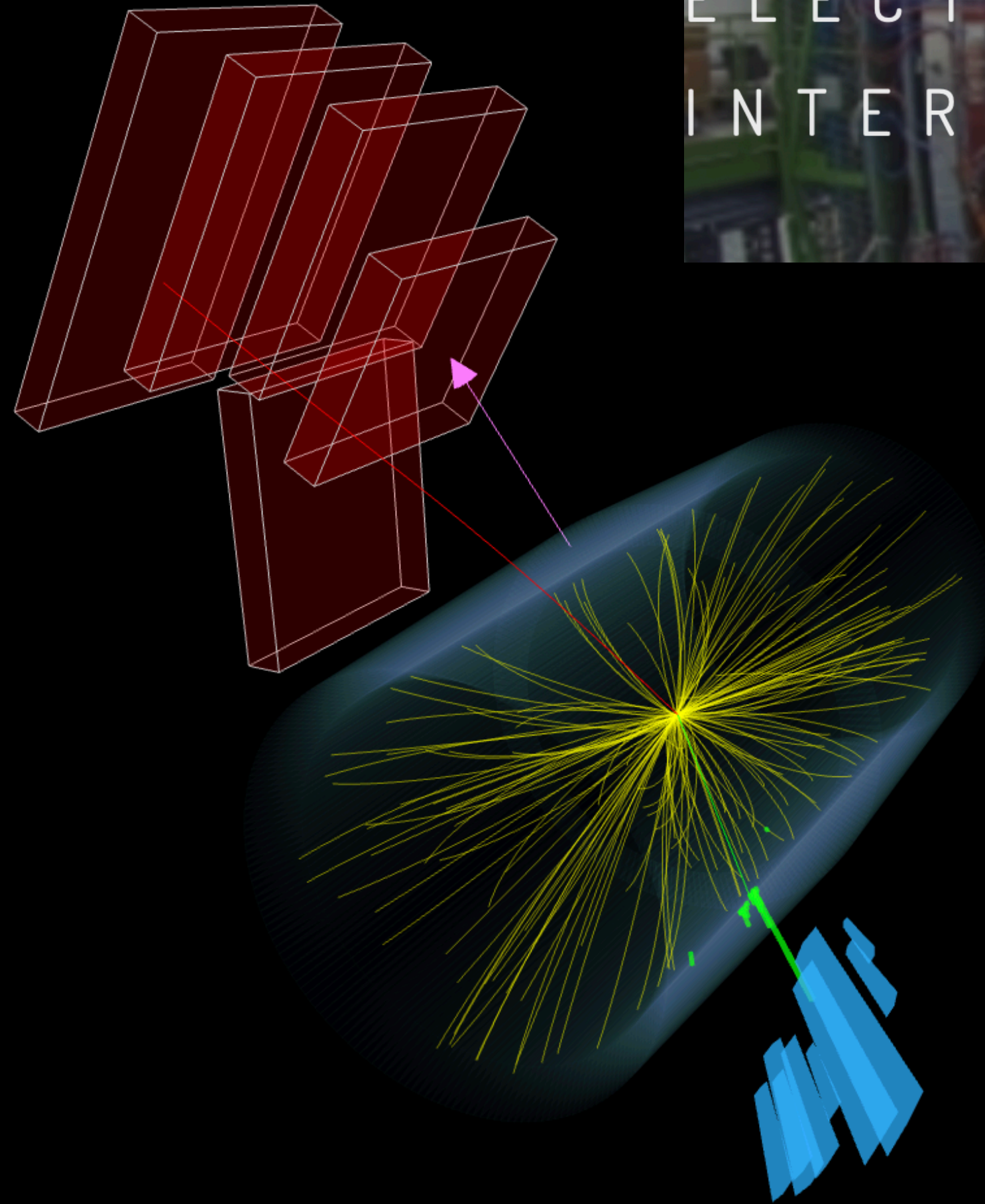
$$m_T^{WW} = \sqrt{2p_T^{\ell\ell} p_T^{\text{miss}} [1 - \cos \Delta\phi(\vec{p}_T^{\ell\ell}, \vec{p}_T^{\text{miss}})]}$$



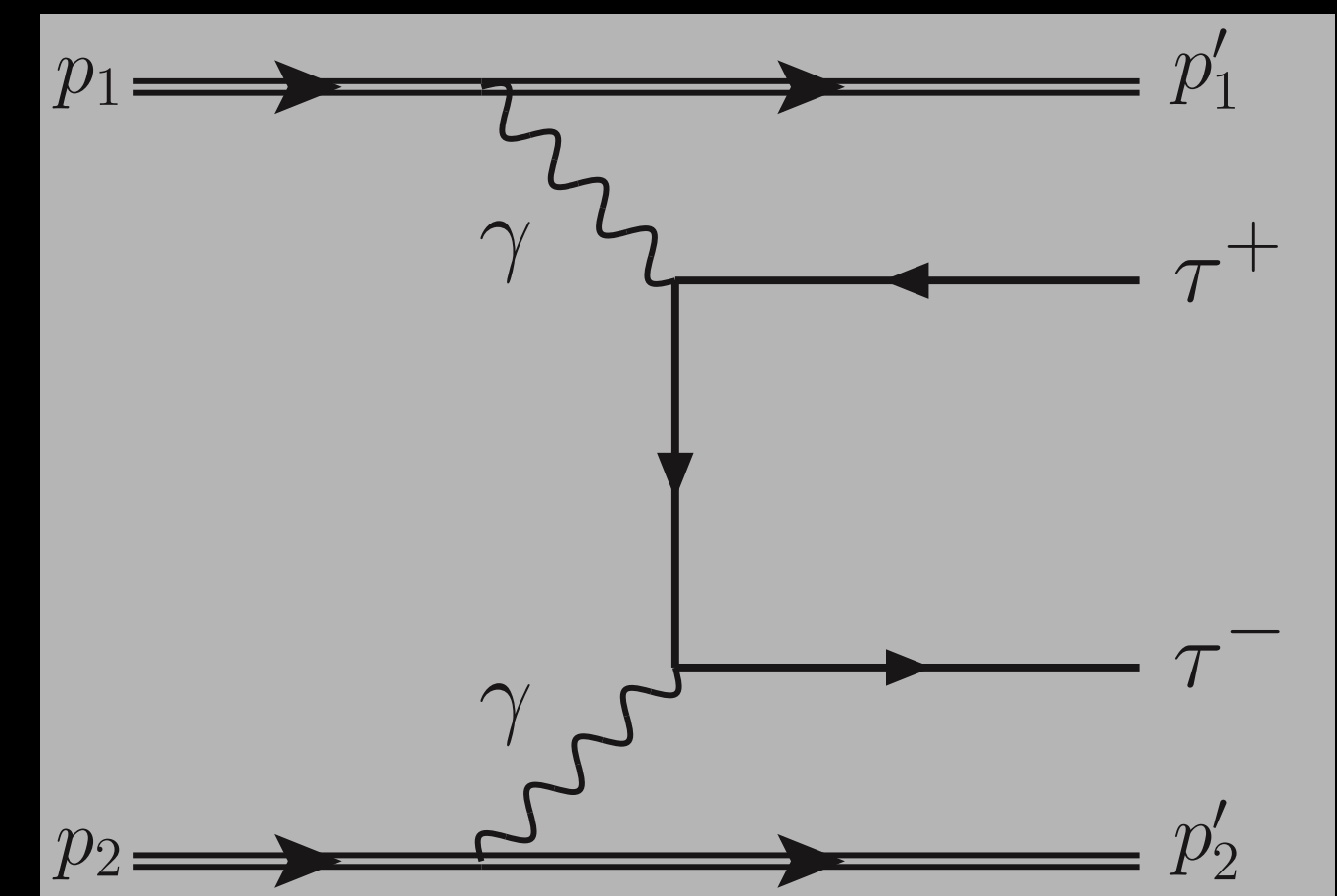
# Observation of $WW\gamma$ production at $\sqrt{s} = 13$ TeV



- Signal extracted from a binned maximum likelihood fit using two dimensional distributions in  $m_T^{WW}$  and  $m_{ll\gamma}$  (product of the Poisson probability mass functions for each bin forms the likelihood function)



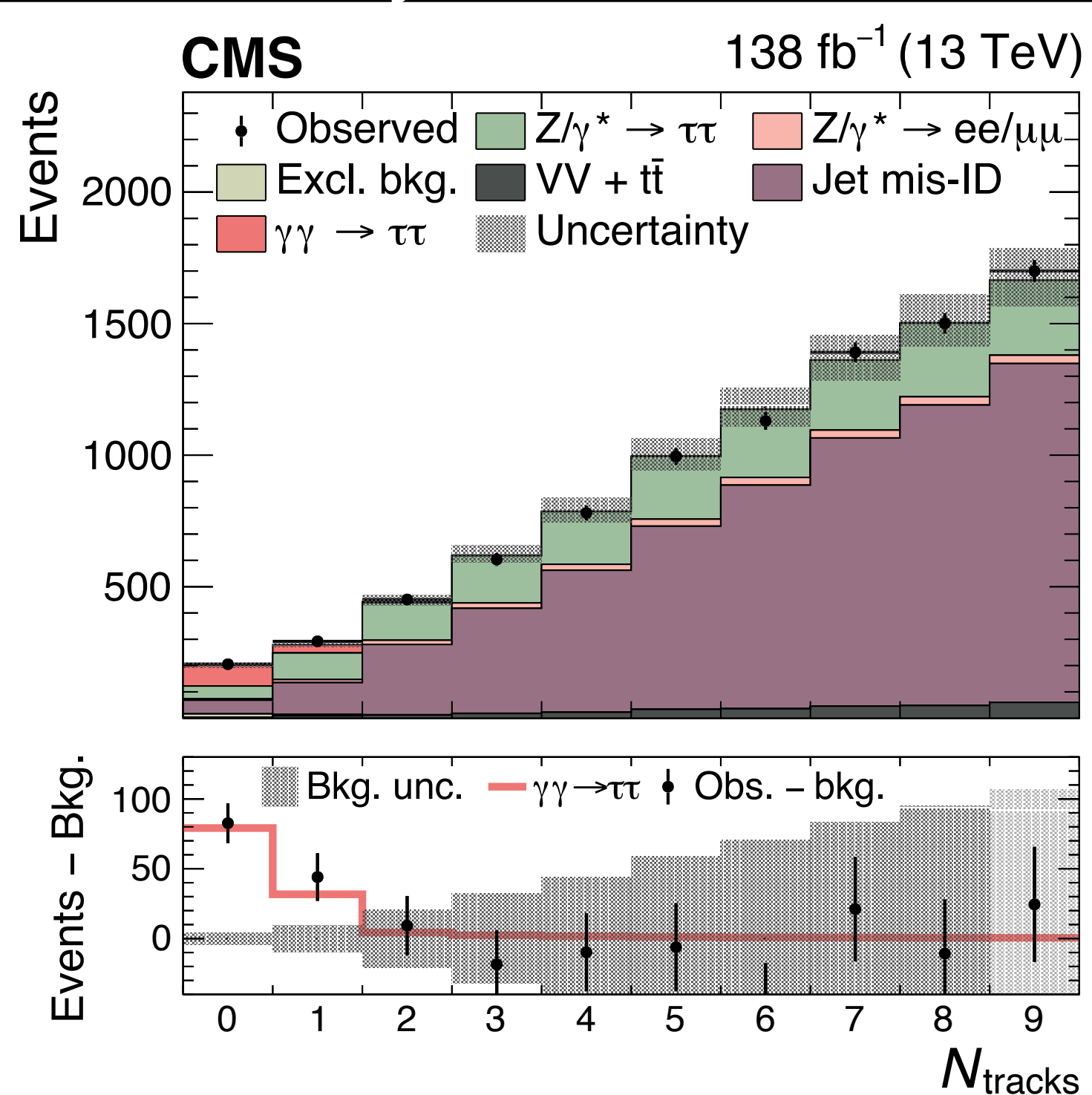
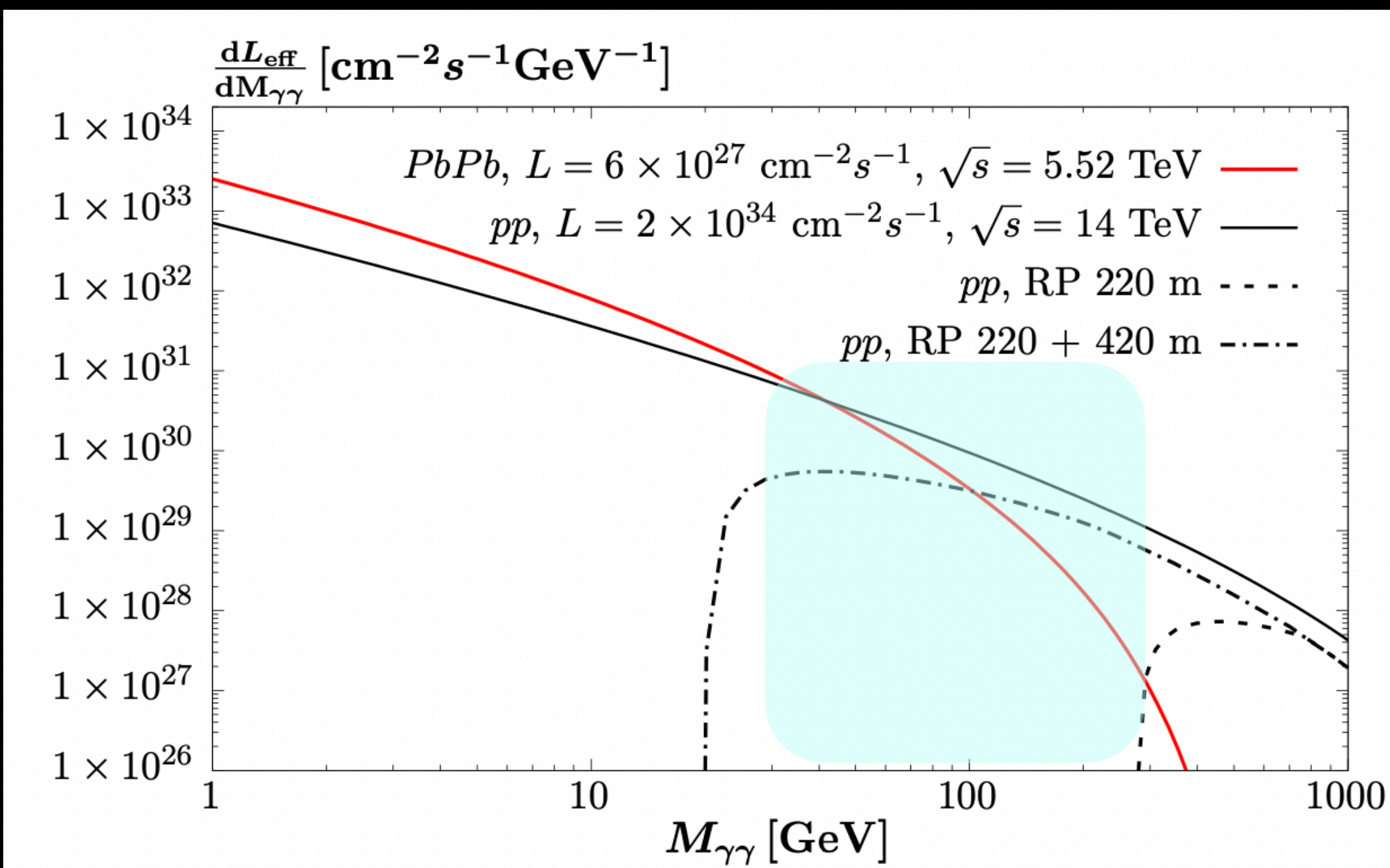
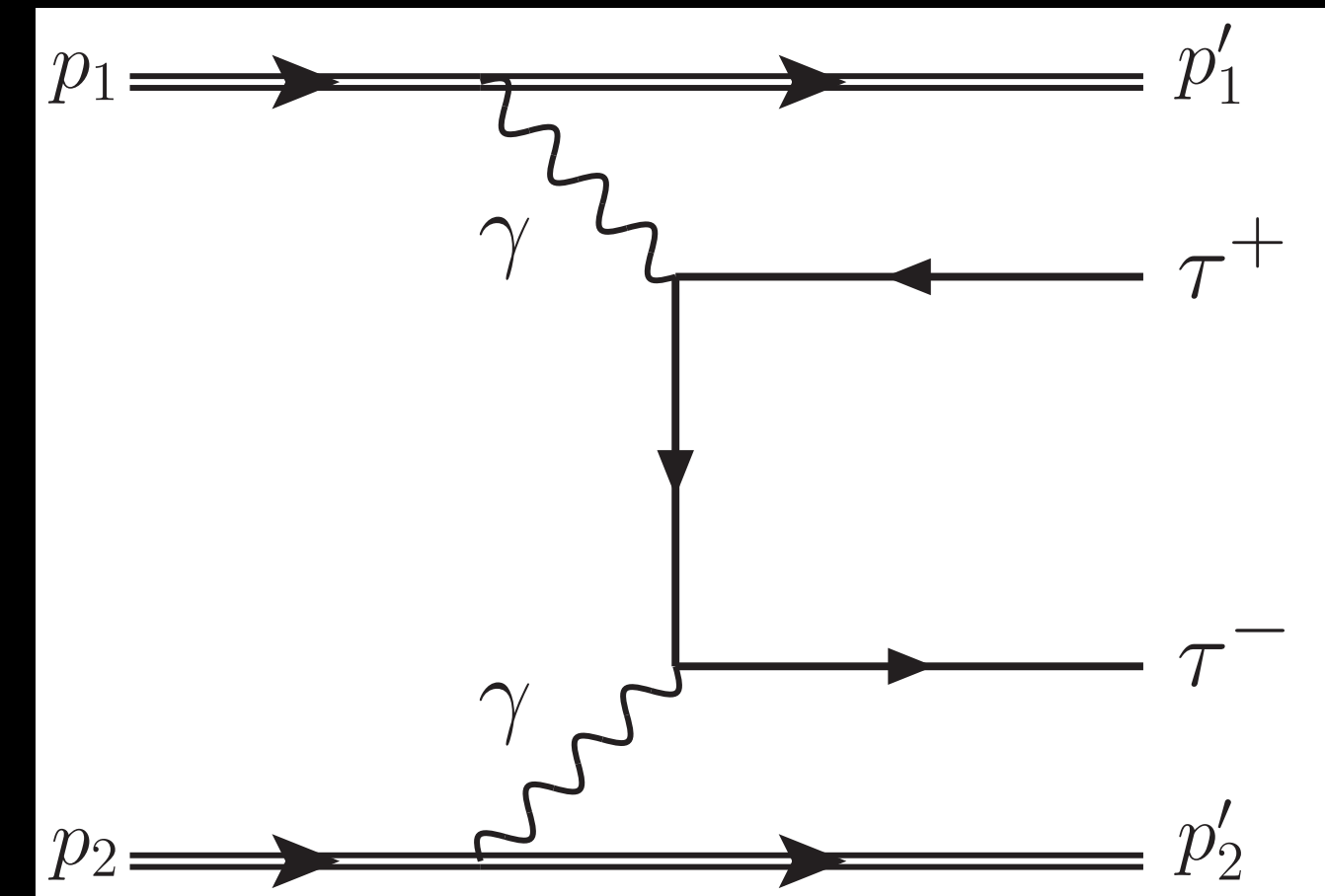
DISCLOSING QUANTUM  
CORRECTIONS TO  
ELECTROMAGNETIC  
INTERACTIONS OF TAU LEPTONS



Candidate  $\gamma\gamma \rightarrow \tau\tau$  event measured in proton-proton collisions by CMS. The event is reconstructed as having a leptonic  $\tau$  decay,  $\tau \rightarrow \mu\nu\nu$ , with the  $\mu$  track indicated in red, and a hadronic  $\tau$  decay,  $\tau \rightarrow \pi\pi\pi\nu$ , with the 3 charged pions indicated by the yellow tracks and by the energy deposits in the ECAL (green) and HCAL (cyan)

# The LHC as a photon collider

- When protons pass each other at relativistic velocities
  - generate intense electromagnetic fields
    - photon-photon collisions occur
- First observation of  $\gamma\gamma \rightarrow \tau\tau$  in proton-proton collisions
  - $5.3 \sigma$  observed,  $6.5 \sigma$  expected
- Non-dissociative protons → low hadronic activity



Elastic signal events characterized by low track multiplicity

# Measurement of the $g_\tau - 2$

Most general form of the QED vertex:

$$\Gamma^\mu = \gamma^\mu F_1(q^2) + \frac{\sigma^{\mu\nu} q_\nu}{2m} [iF_2(q^2) + F_3(q^2)\gamma_5]$$

$$F_2(0) = a_\ell \equiv (g_\ell - 2)/2$$

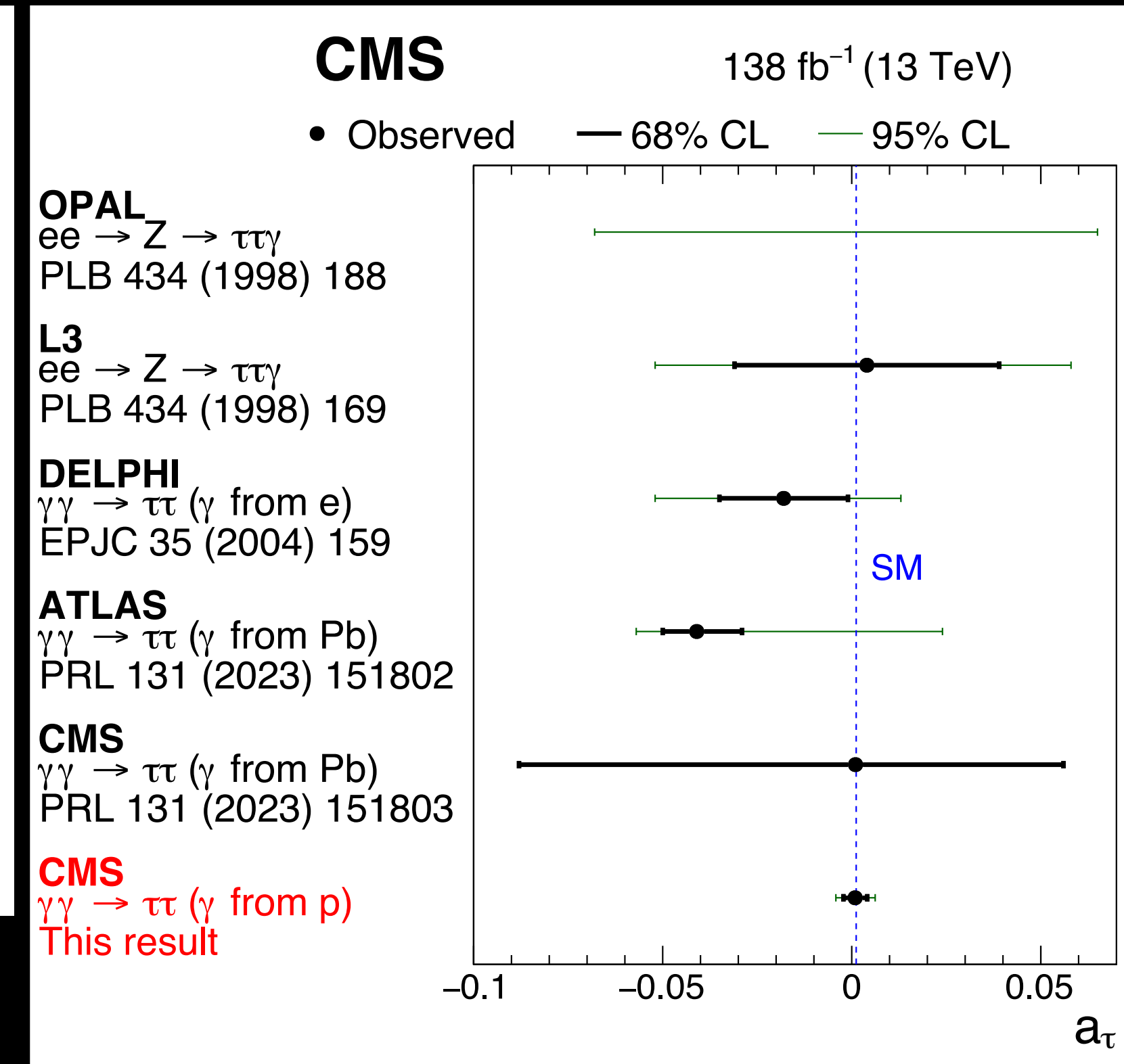
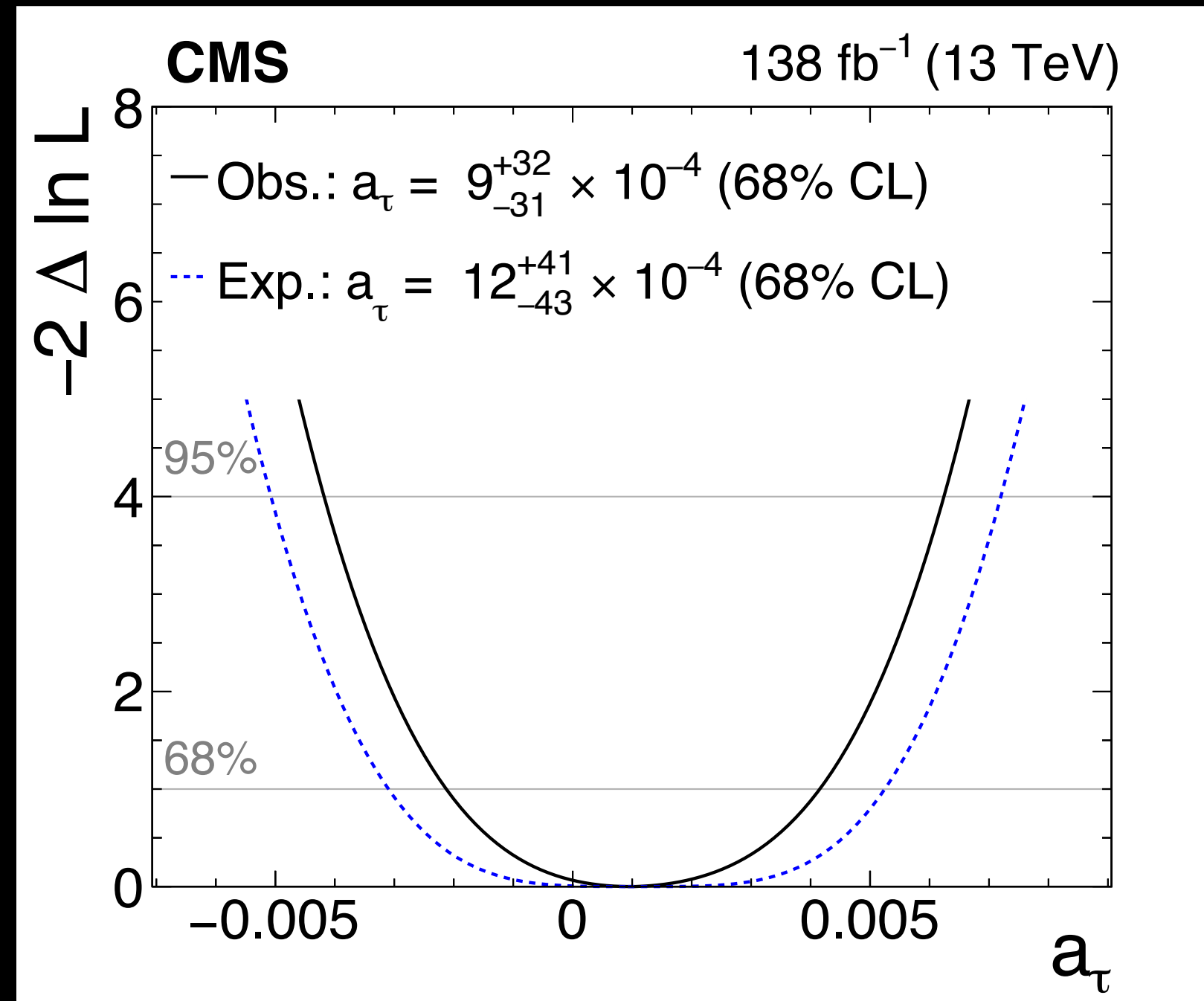
$$F_3(0) = -\frac{2m}{e} d_\ell$$

$g_\ell$  (gyromagnetic ratio)  $\rightarrow$  relates the magnetic moment to the spin of the lepton

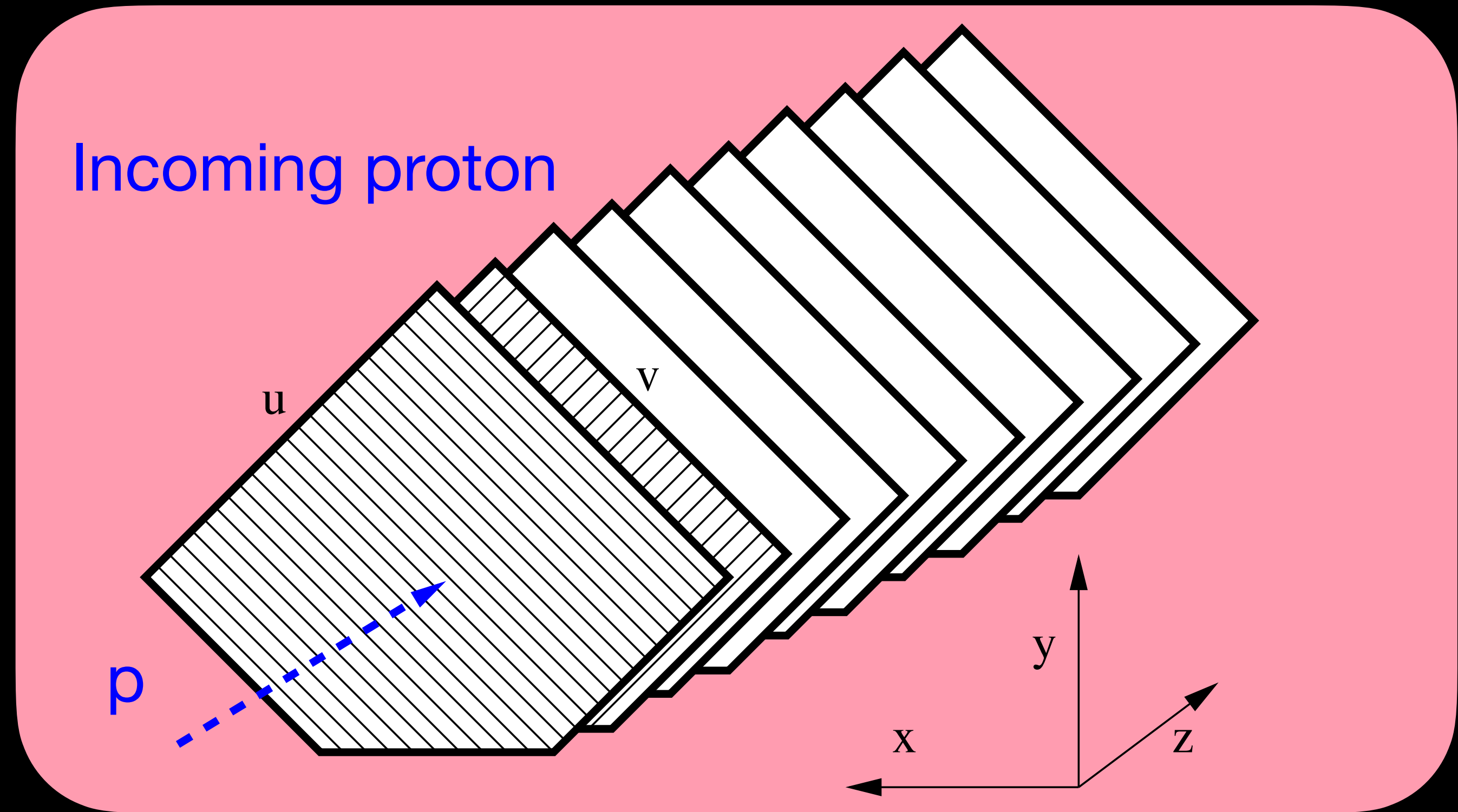
**Schwinger term:**

$$a_\ell = \frac{\alpha}{2\pi} \simeq 0.00116$$

- Extremely sensitive analysis
- Limits set to  $\sim 3$  X Schwinger term



**Milestone in collider physics!**



The TOTEM Roman pot detectors

- reconstruct transverse momentum of scattered protons
- estimate the transverse location of the primary interaction

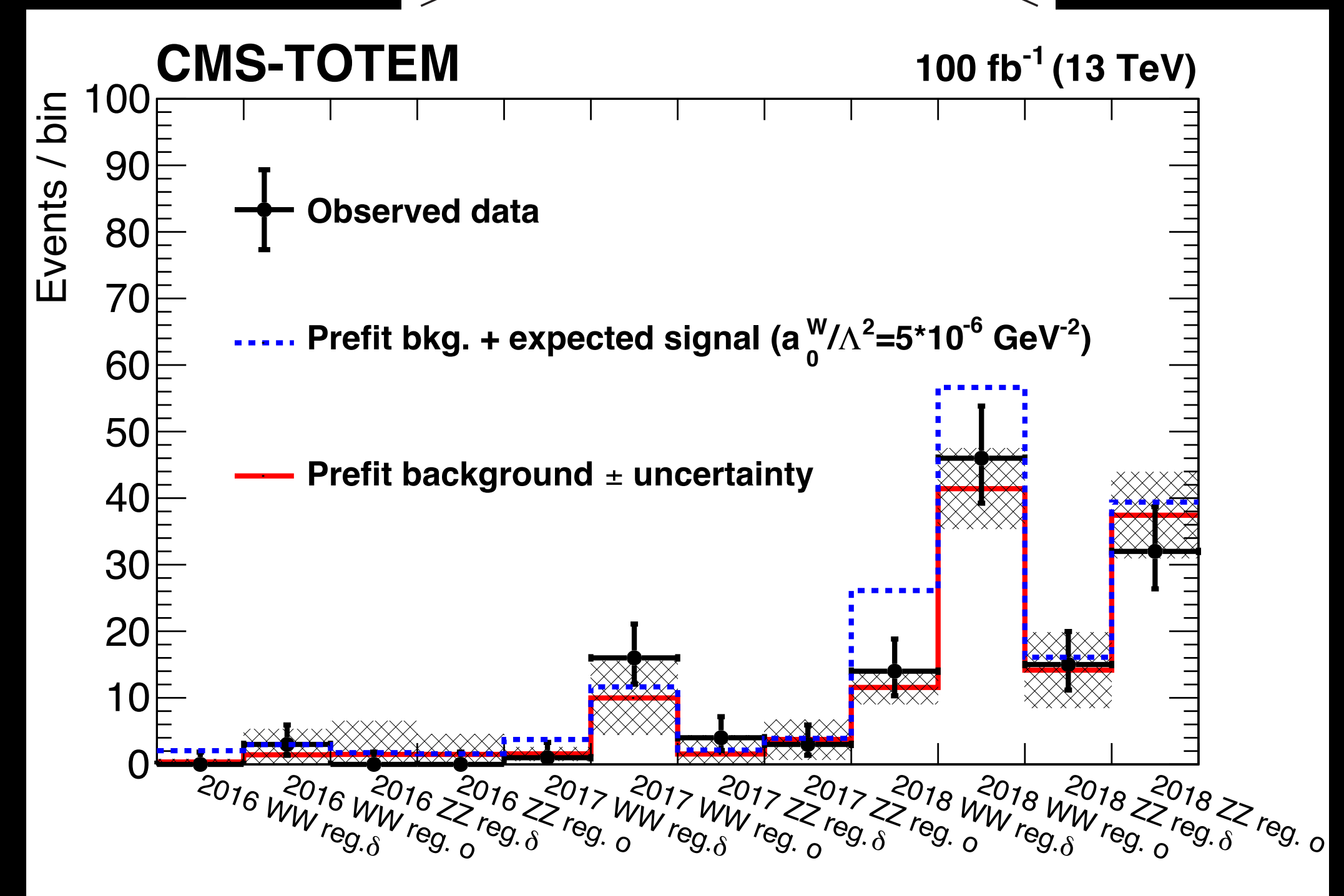
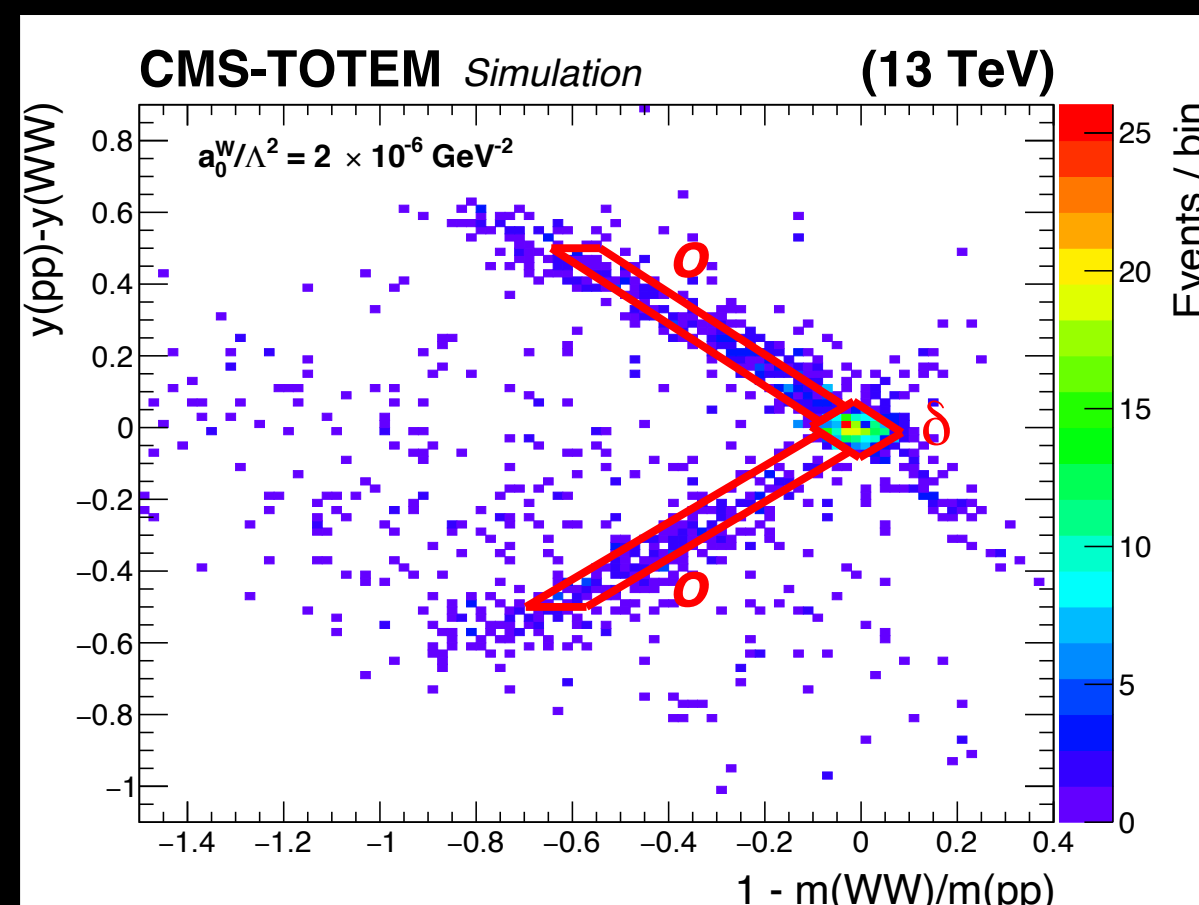
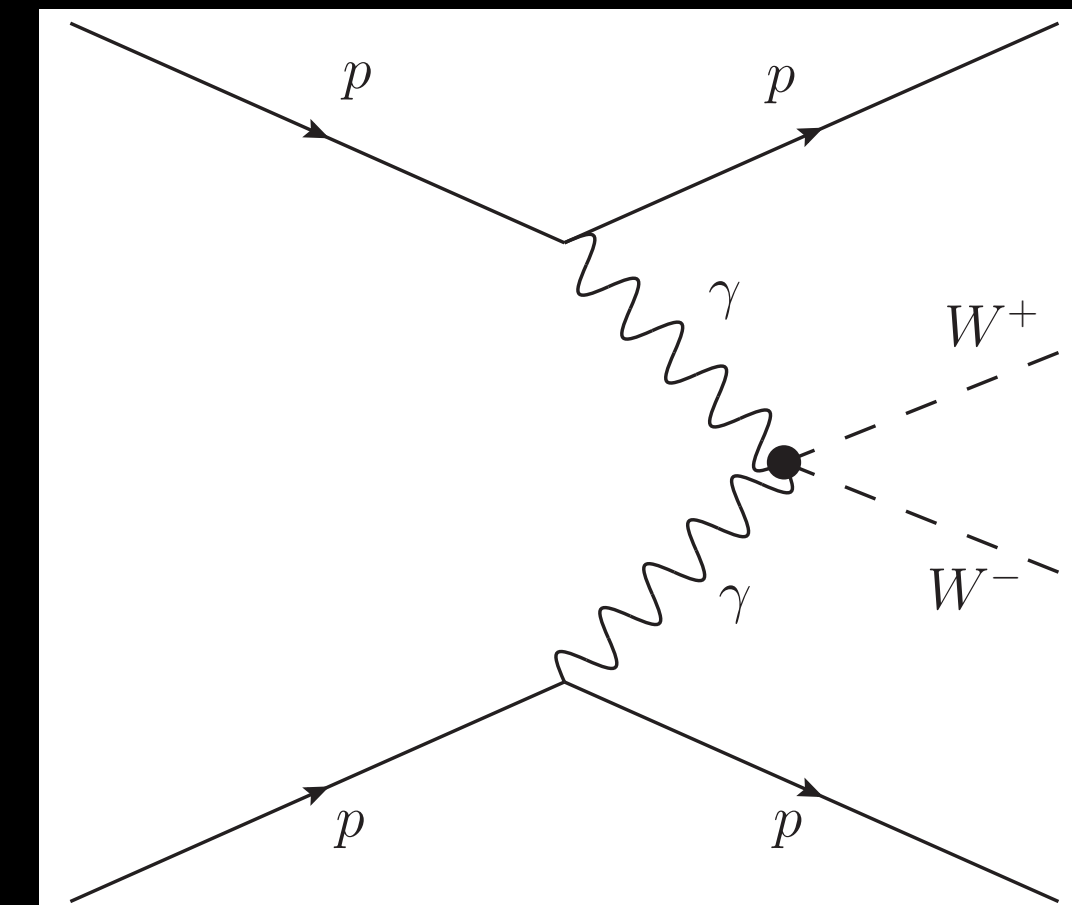
New studies on high- $\beta^*$  period in 2018 expected to enable several physics analyses



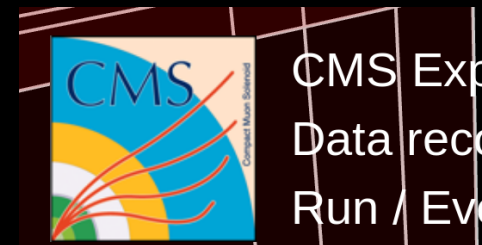


# Search for exclusive $\gamma\gamma \rightarrow WW$ and $\gamma\gamma \rightarrow ZZ$ production in final states with jets and forward protons

- Both protons tagged by the precision proton spectrometer (PPS)
- The  $\gamma\gamma \rightarrow WW$  process allows the study of the quartic coupling
- Events selected based on properties of jets, the protons and their correlation
- First search for anomalous high-mass  $\gamma\gamma \rightarrow WW$  and  $\gamma\gamma \rightarrow ZZ$  using reconstructed forward protons
  - Limits 15-20x more stringent than previous results



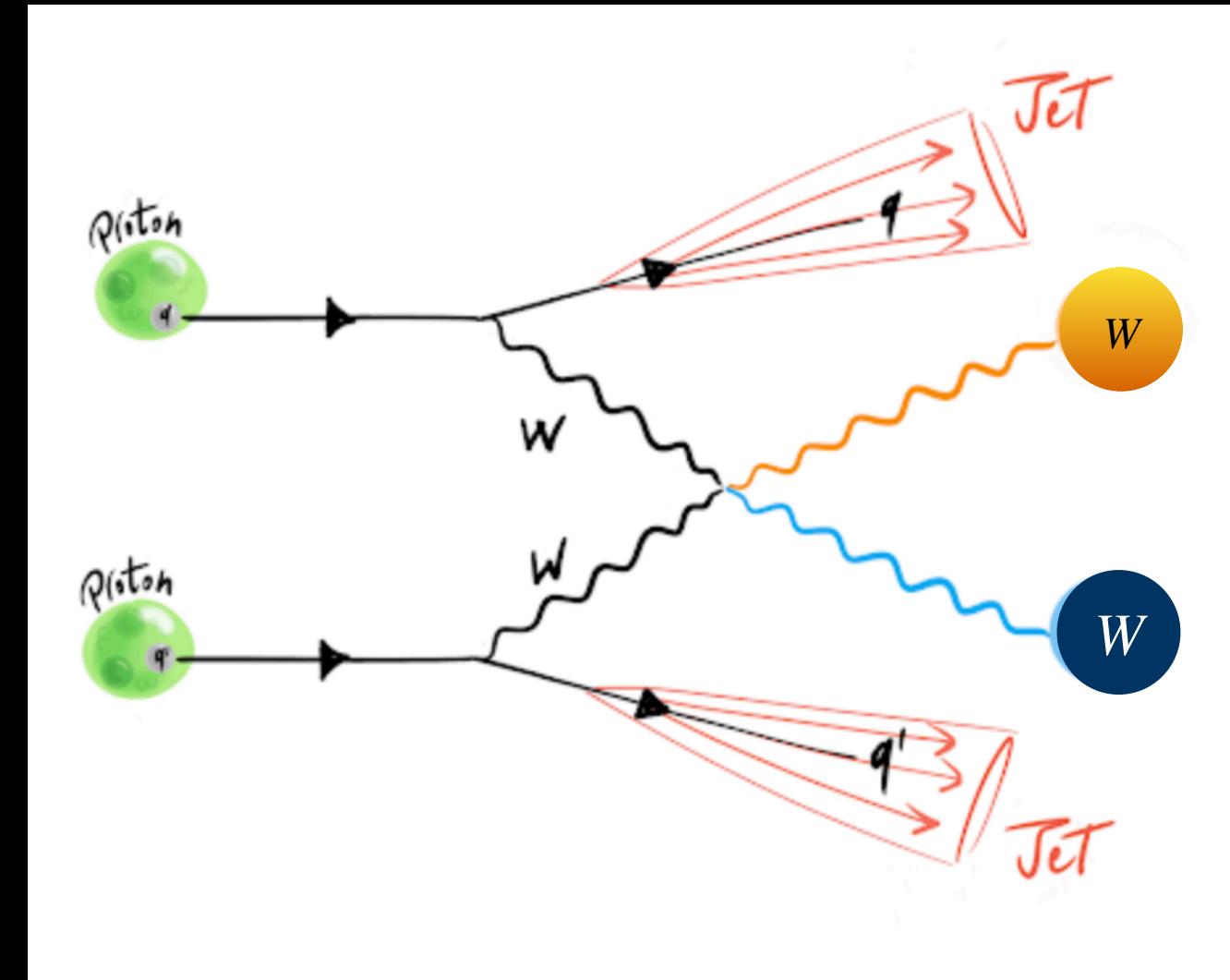
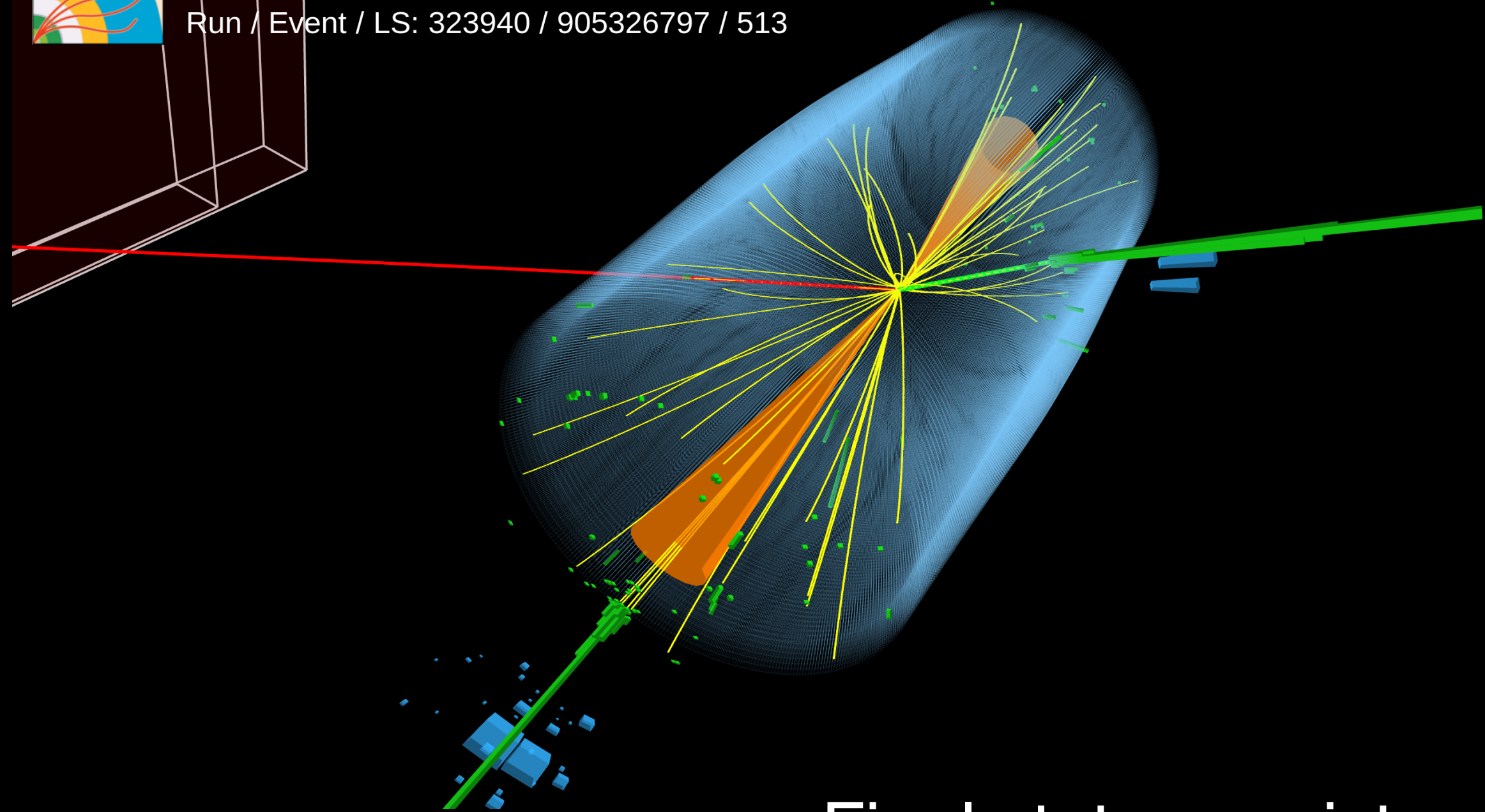
# Vector Boson Scattering



CMS Experiment at the LHC, CERN

Data recorded: 2018-Oct-03 04:13:04.188416 GMT

Run / Event / LS: 323940 / 905326797 / 513

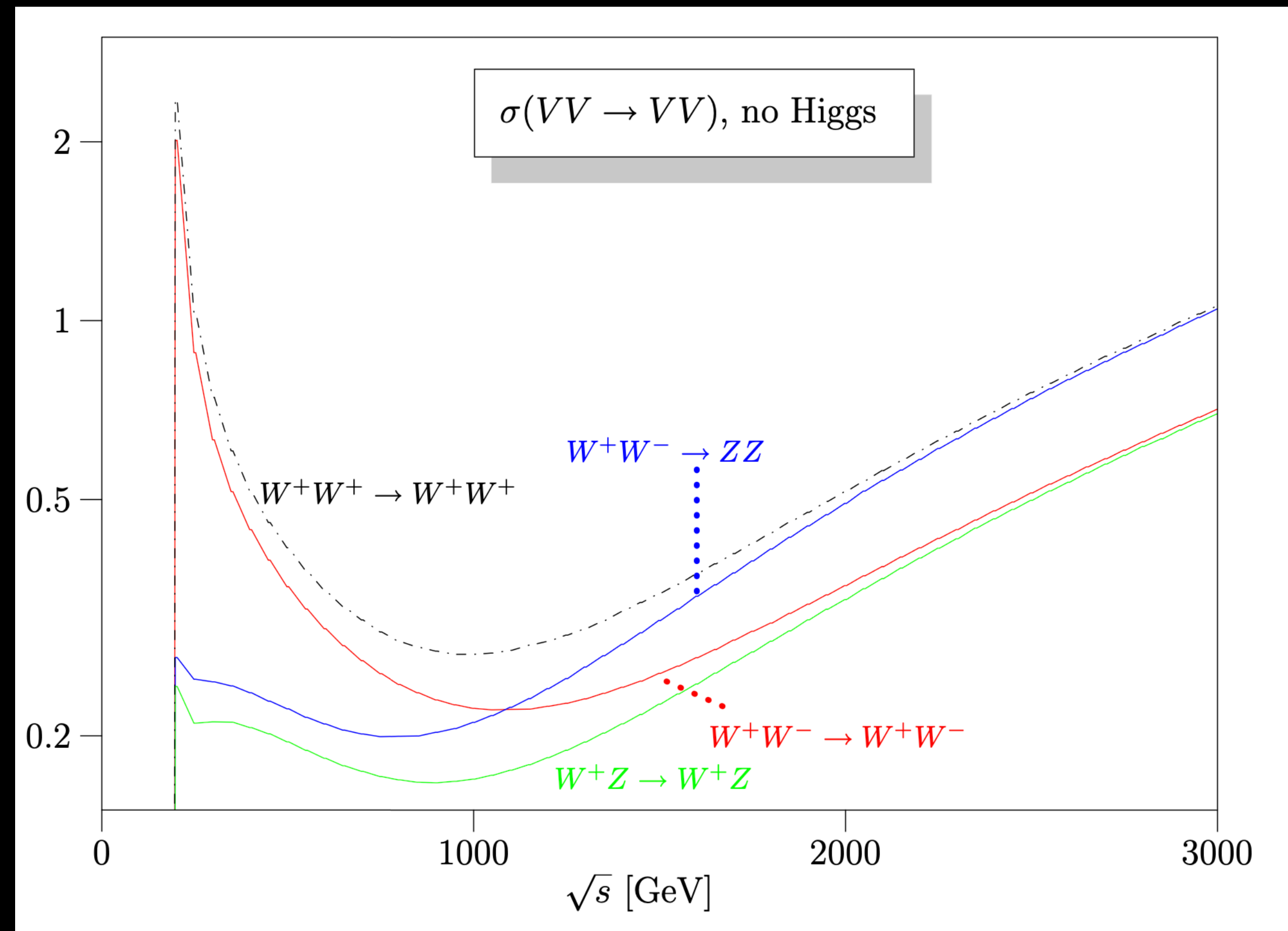


- Pure electro-weak interactions of order at least  $\alpha_{EW}^4$

- Final state consists of two high transverse momenta ( $p_T$ ) jets
- Rapidity gap between jets  $\rightarrow$  no color flow
- High mass of the two jets ( $M_{jj}$ )
- Decay product of the gauge bosons: central with respect to jets

# Why Vector Boson Scattering?

Studying the Higgs mechanism *Higgslessly!*



arXiv:0806.4145

ChatGPT

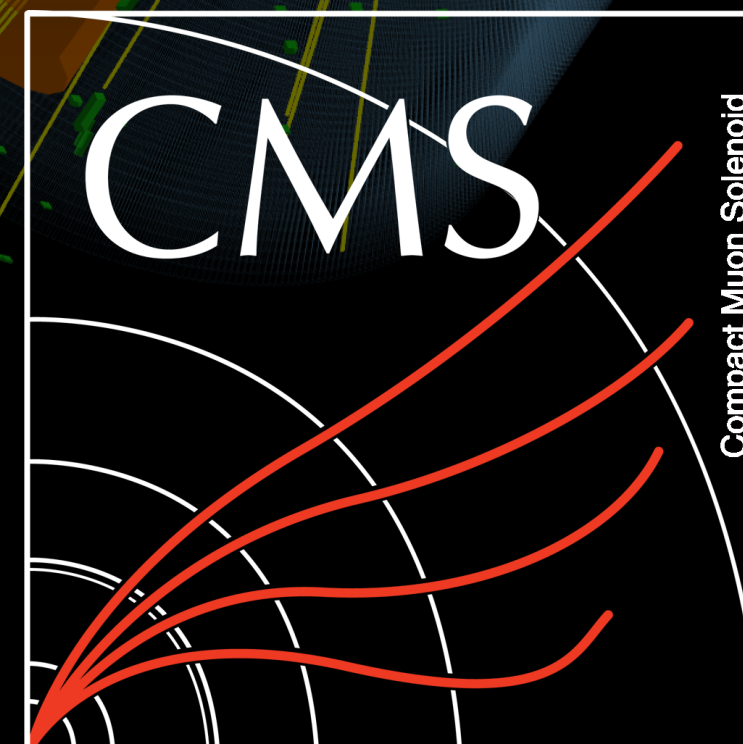
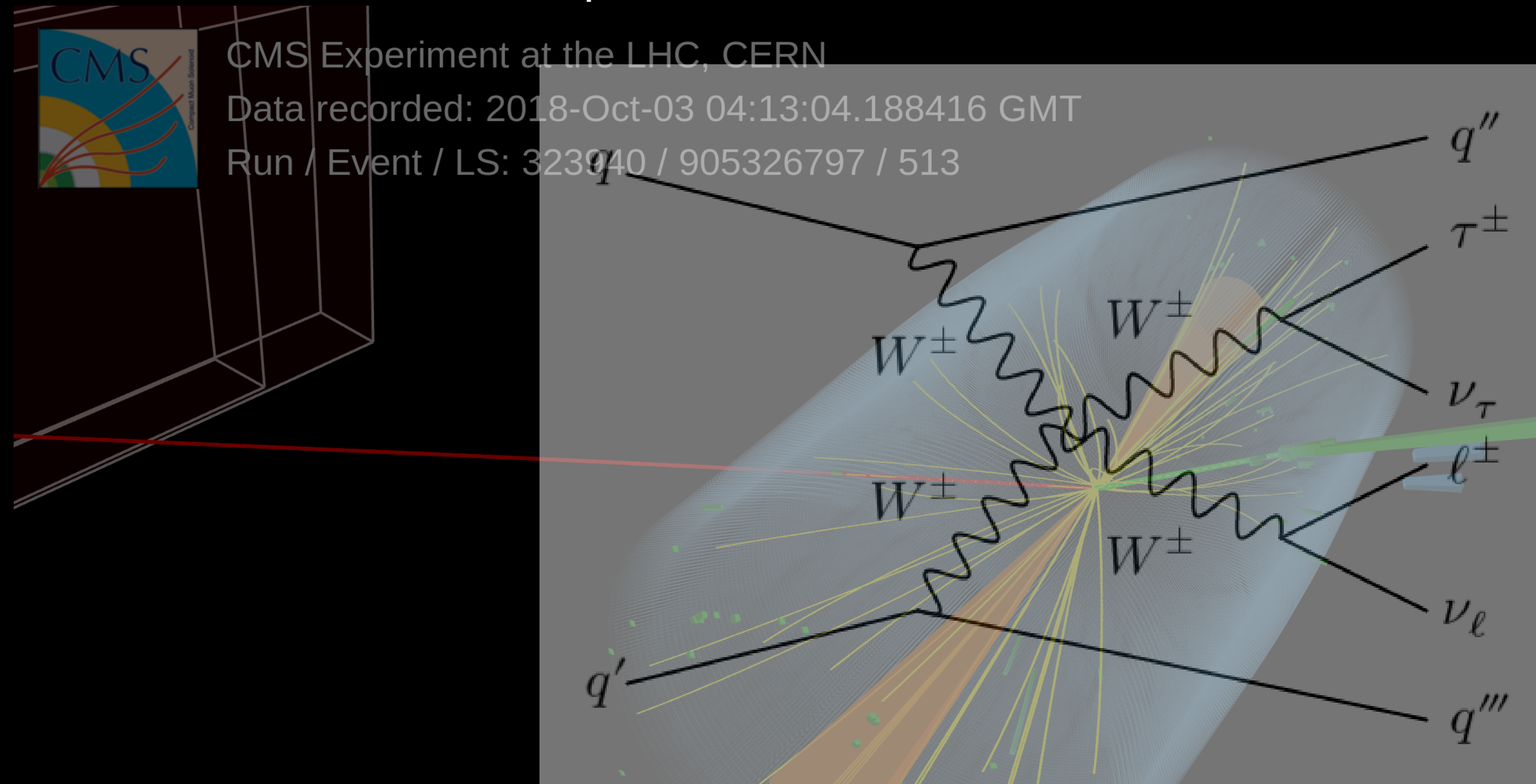
“However, it is important to note that the precise role of the Higgs boson in unitarizing  $WW$  scattering, in particular, is still an active area of research and there are ongoing studies to refine our understanding of this process.”

# Vector Boson Scattering – evolution of knowledge

$\sqrt{s}$	Process	Reference	Significance/Status
8 TeV	EW $W^\pm W^\pm jj$ ( $2\ell 2\nu jj$ )	<a href="#">PhysRevLett.114.051801</a>	$2\sigma$
	EW $Z\gamma jj$ ( $\nu\nu/\ell\ell\gamma jj$ )	<a href="#">PhysLettB770(2017)380-402</a>	$3\sigma$
	EW $W^\pm\gamma jj$ ( $\ell\nu\gamma jj$ )	<a href="#">JHEP06(2017)106</a>	$2.7\sigma$
	EW $W^\pm Z jj$ ( $3\ell\nu jj$ )	<a href="#">PhysRevLett.114.051801</a>	$2\sigma$
13 TeV	EW $W^\pm W^\pm jj$ ( $2\ell 2\nu jj$ )	<a href="#">PhysLettB809(2020)</a>	2016: $5.5\sigma$ , Run II $\gg 5.0\sigma$
	EW $ZZ jj$ ( $4\ell jj$ )	<a href="#">PhysLettB812(2021)135992</a>	2016: $2.7\sigma$ , Run II: $4.0\sigma$
	EW $W^\pm Z jj$ ( $3\ell jj$ )	<a href="#">PhysLettB809(2020)135710</a>	$6.8\sigma$
	EW $Z\gamma jj$ ( $\ell\ell\gamma jj$ )	<a href="#">PhysRevD.104.072001</a>	$4.7\sigma$ , Run II $\gg 5.0\sigma$
	EW $W^\pm\gamma jj$ ( $\ell\nu\gamma jj$ )	<a href="#">PhysLettB811(2020)135988</a>	2016: $5.3\sigma$
	EW $W^\pm V jj$ ( $\ell\nu jjjj$ )	<a href="#">arXiv:2112.05259</a>	$4.4\sigma$
	EW $W^\pm W^\mp jj$ ( $2\ell 2\nu jj$ )	<a href="#">arXiv:2205.05711</a>	$5.6\sigma$
	EW $VVpp$ ( $4jpp$ )	<a href="#">arXiv:2211.16320</a>	CMS with TOTEM
	EW $W^\pm\gamma jj$ ( $\ell\nu\gamma jj$ )	<a href="#">arXiv:2212.12592</a>	Cross section measurement with full Run II
	$W^\pm W^\pm jj$ ( $\ell\tau_{\text{had}} 2\nu jj$ )	<a href="#">SMP-22-008</a>	$2.9\sigma$

# VBS Same Sign $WW$ with hadronic $\tau$

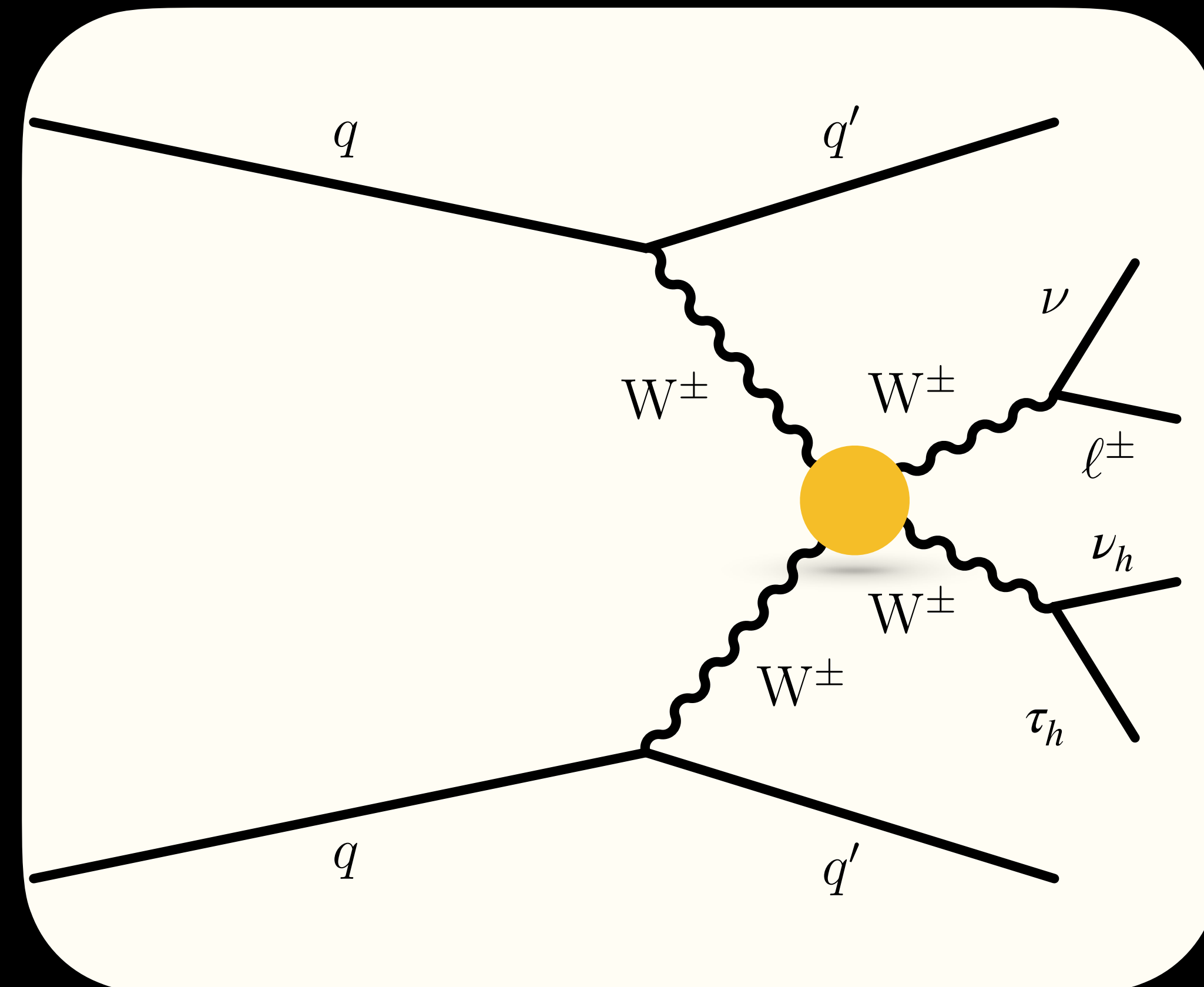
<http://cds.cern.ch/record/2867989?ln=en>



# Vector Boson Scattering (VBS)

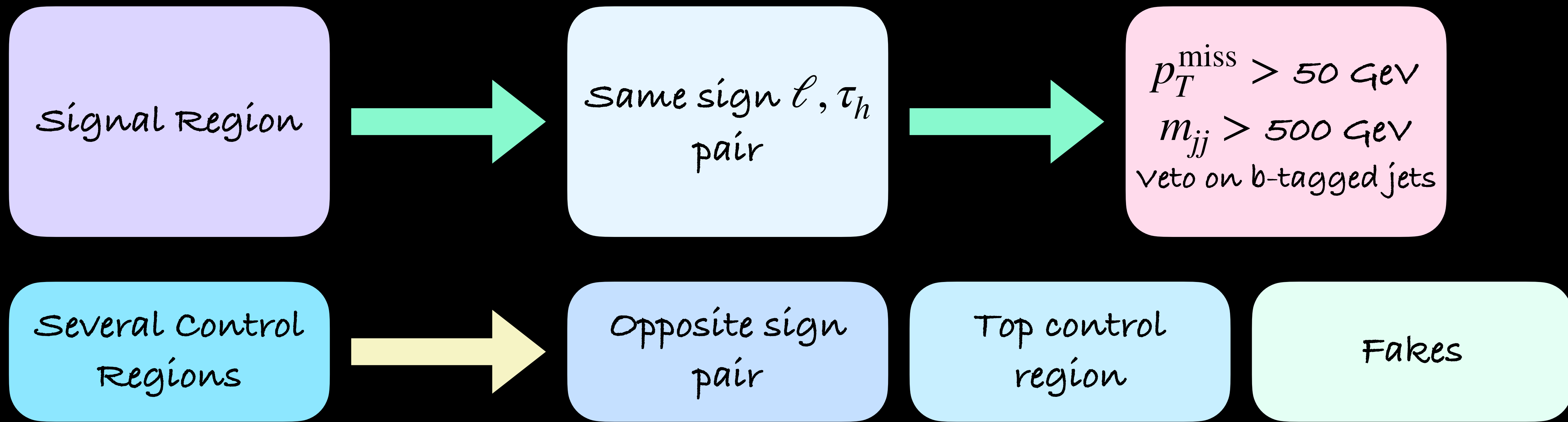
- Study of vector boson scattering processes  
→ essential to probe nature of **electroweak symmetry breaking**
- Constraints on various **effective field theory operators** evaluated
- Several kinematic observables sensitive to the effect of new physics

Pure electro-weak interactions of order  $\alpha_{EW}^6$



# VBS Same Sign $WW$ with hadronic $\tau$

- Final state where one of the two same-signed W-bosons decays to a hadronic  $\tau$ 
  - Signature:  $\tau_h \nu_\tau \ell \nu_\ell (\ell = e, \mu)$
- Evidence of SM process at  $2.7 \sigma$ , signal strength:  $1.44^{+0.63}_{-0.56}$
- Public since August 2023



# Exploring the implication of new physics modifying the quartic coupling in the Standard Model

$$\mathcal{L} = \mathcal{L}_{SM} + \sum_i \frac{c_i}{\Lambda^2} \mathcal{O}_i + \sum_j \frac{f_j}{\Lambda^4} \mathcal{O}_j$$

- In the past, diboson analyses used to extract sensitivity to dimension-6 operators and vector boson scattering primarily used to extract sensitivity to dimension-8 operators
- Now explore vector boson scattering for both dimension-6 and dimension-8 exploration



# Exploring the implication of new physics modifying the quartic coupling in the Standard Model

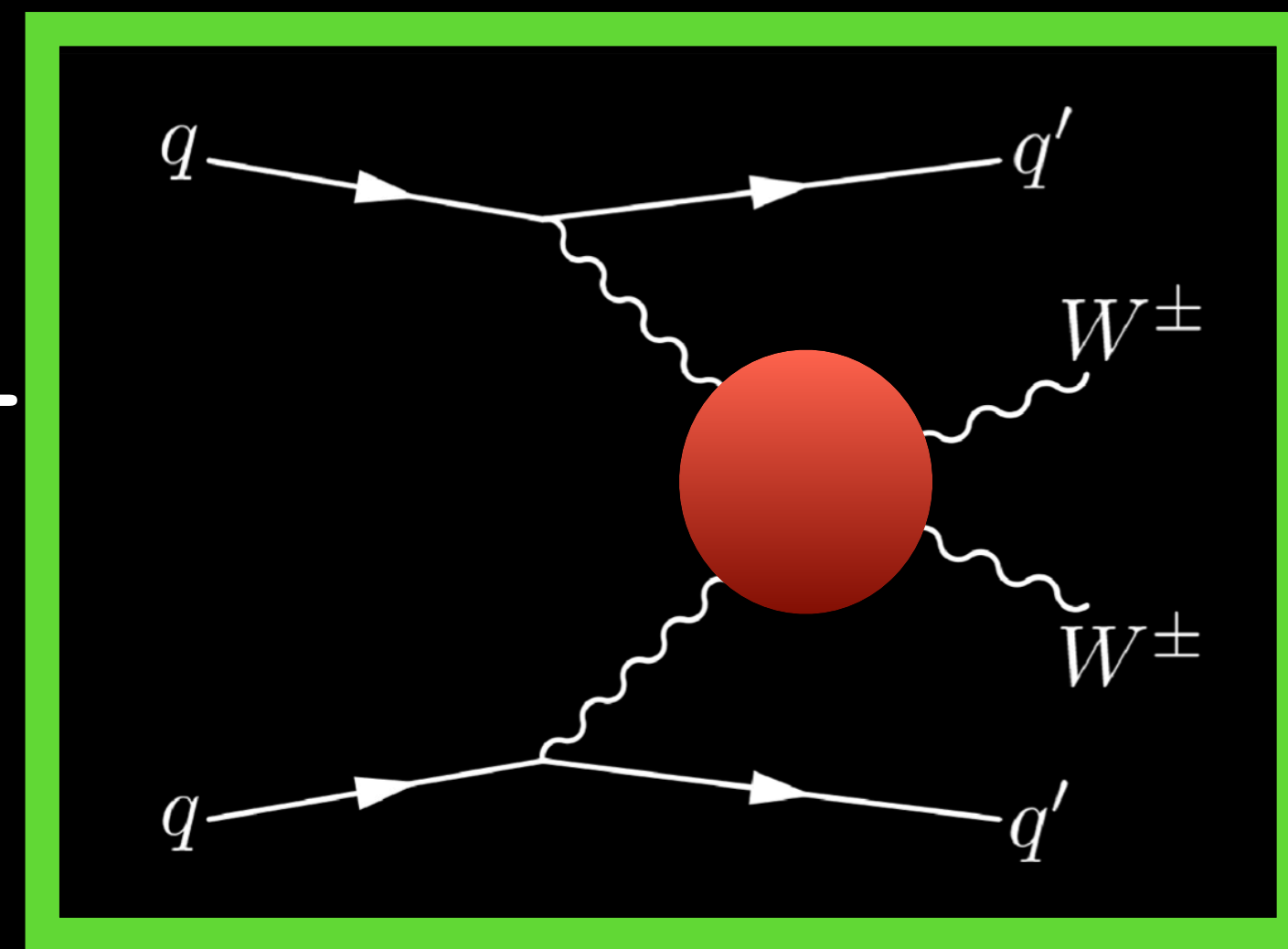
$$\mathcal{L} = \mathcal{L}_{SM} + \text{[Green Box Diagram]} + \text{[Cyan Box Diagram]}$$

The diagram illustrates the Lagrangian  $\mathcal{L}$  as the sum of the Standard Model Lagrangian  $\mathcal{L}_{SM}$  and two new physics contributions. The first contribution (green box) shows a vertex  $g$  where two fermions  $f_1$  and  $f_2$  meet, emitting a  $Z$  boson. This  $Z$  boson then interacts with a red circle representing a new physics operator, producing two  $W^\pm$  bosons. The second contribution (cyan box) shows a vertex where two quarks  $q$  meet, emitting a  $Z$  boson. This  $Z$  boson then interacts with a red circle representing a new physics operator, producing two quarks  $q'$  and two  $W^\pm$  bosons.

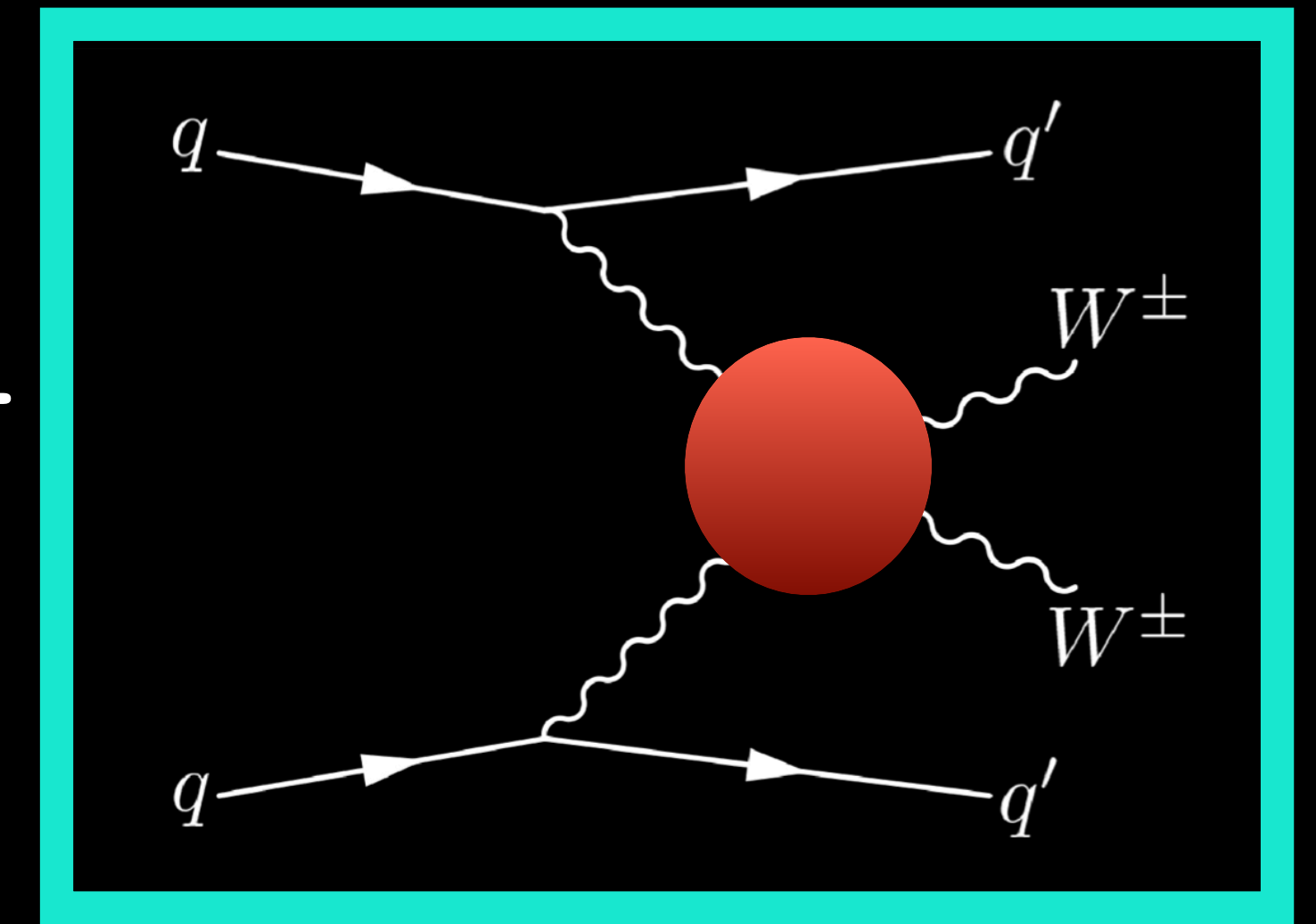
- In the past, diboson analyses used to extract sensitivity to dimension-6 operators and vector boson scattering primarily used to extract sensitivity to dimension-8 operators

# Exploring the implication of new physics modifying the quartic coupling in the Standard Model

$$\mathcal{L} = \mathcal{L}_{SM} +$$



+



- Now explore vector boson scattering for both dimension-6 and dimension-8 exploration

# VBS Same Sign $WW$ with hadronic $\tau$

- Deep Neural Network (DNNs) trained with used to gain sensitivity to BSM signals

extracted with dim-6 DNN output distributions	Wilson coefficient	68% CL interval(s)		95% CL interval	
		Expected	Observed	Expected	Observed
dim-6	$c_{ll}^{(1)}$	$[-12.9, -8.03] \cup [-2.95, 1.91]$	$[-11.6, 0.045]$	$[-14.6, 3.53]$	$[-13.5, 2.11]$
	$c_{qq}^{(1)}$	$[-0.501, 0.576]$	$[-0.341, 0.416]$	$[-0.742, 0.818]$	$[-0.605, 0.681]$
	$c_W$	$[-0.681, 0.669]$	$[-0.513, 0.481]$	$[-0.987, 0.974]$	$[-0.842, 0.818]$
	$c_{HW}$	$[-7.00, 6.09]$	$[-5.48, 4.31]$	$[-9.99, 9.05]$	$[-8.68, 7.60]$
	$c_{HWB}$	$[-41.7, 69.6]$	$[30.7, 89.2]$	$[-66.6, 96.4]$	$[-49.7, 110]$
	$c_H$	$[-16.6, 18.1]$	$[-12.0, 14.0]$	$[-24.7, 26.3]$	$[-20.9, 22.7]$
	$c_{HD}$	$[-24.6, 34.7]$	$[-15.3, 31.5]$	$[-38.2, 48.8]$	$[-31.4, 45.5]$
	$c_{Hl}^{(1)}$	$[-28.8, 29.9]$	$[-38.2, 39.5]$	$[-49.4, 49.7]$	$[-69.3, 68.3]$
	$c_{Hl}^{(3)}$	$[-1.43, 2.23] \cup [5.88, 9.54]$	$[-0.045, 8.58]$	$[-2.64, 10.8]$	$[-1.59, 9.94]$
	$c_{Hq}^{(1)}$	$[-4.53, 4.42]$	$[-3.27, 3.44]$	$[-6.56, 6.44]$	$[-5.55, 5.60]$
	$c_{Hq}^{(3)}$	$[-2.39, 1.37]$	$[-1.88, 0.705]$	$[-3.24, 2.16]$	$[-2.82, 1.61]$
	extracted with dim-8 DNN output distributions	$f_{T0}$	$[-1.02, 1.08]$	$[-0.774, 0.842]$	$[-1.52, 1.58]$
$f_{T1}$		$[-0.426, 0.480]$	$[-0.319, 0.381]$	$[-0.640, 0.695]$	$[-0.552, 0.613]$
$f_{T2}$		$[-1.15, 1.37]$	$[-0.851, 1.12]$	$[-1.75, 1.98]$	$[-1.51, 1.76]$
$f_{M0}$		$[-9.89, 9.74]$	$[-8.07, 7.70]$	$[-14.6, 14.5]$	$[-13.1, 12.8]$
$f_{M1}$		$[-12.5, 13.3]$	$[-9.54, 11.15]$	$[-18.7, 19.6]$	$[-16.4, 17.7]$
$f_{M7}$		$[-20.3, 19.2]$	$[-17.6, 15.3]$	$[-29.9, 28.8]$	$[-27.6, 25.8]$
$f_{S0}$		$[-11.6, 12.0]$	$[-9.60, 9.82]$	$[-17.4, 17.9]$	$[-15.9, 16.1]$
$f_{S1}$		$[-37.4, 38.8]$	$[-40.9, 41.3]$	$[-57.2, 58.6]$	$[-60.9, 61.8]$
dim-8	$f_{S2}$	$[-37.4, 38.8]$	$[-40.9, 41.3]$	$[-57.2, 58.6]$	$[-60.9, 61.8]$

# VBS Same Sign $WW$ with hadronic $\tau$

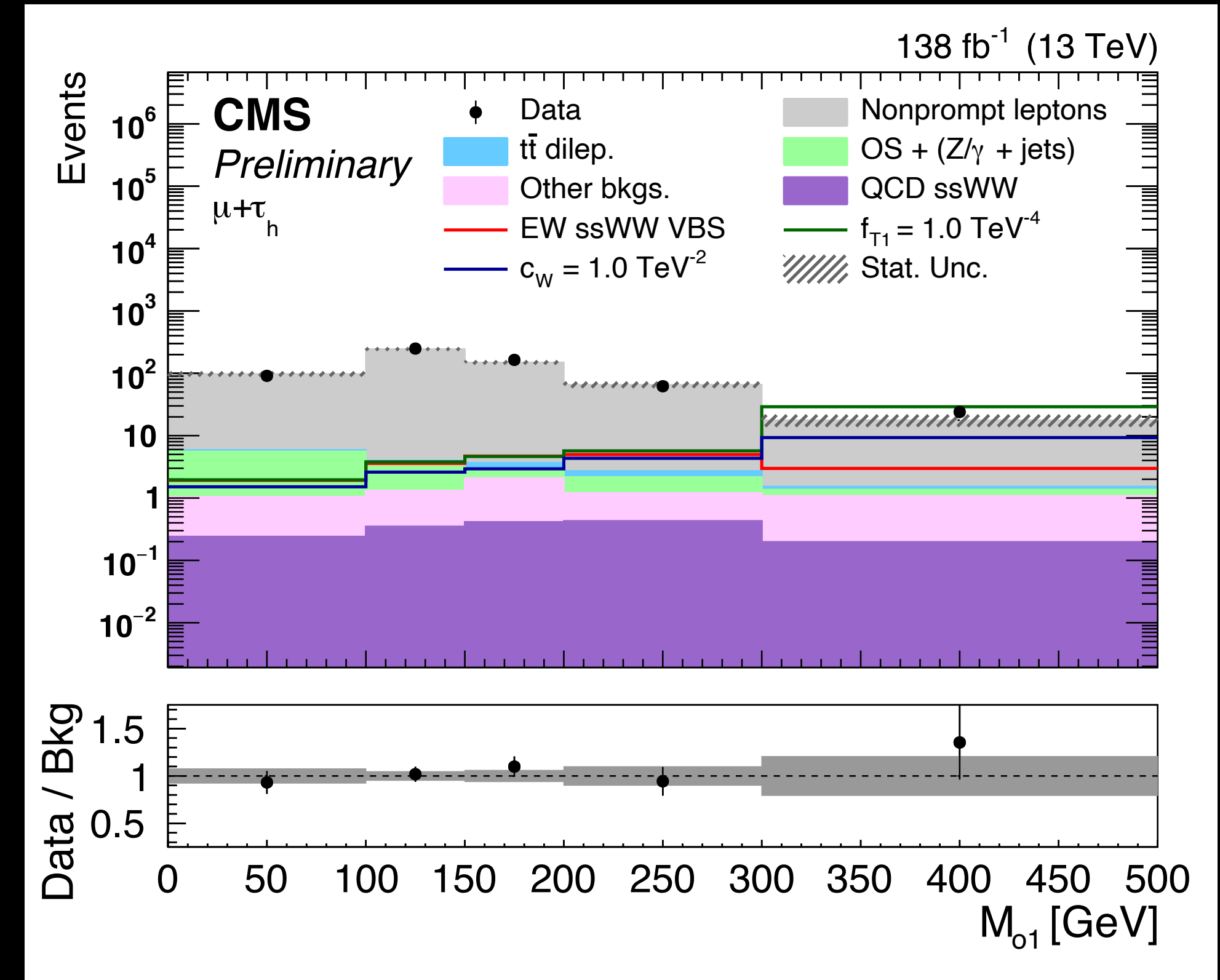
- First simultaneous extraction of dim-6 and dim-8 constraints
- Transverse mass ( $M_{01}$ ) used as the variable of interest for 2D constraints

$$N \propto |\mathcal{A}|^2 = |\mathcal{A}_{SM}|^2 + \sum_{\alpha} \frac{C_{\alpha}}{\Lambda^2} \cdot 2\text{Re}(\mathcal{A}_{SM} \mathcal{A}_{Q\alpha}^{(6)\dagger}) + \sum_{\alpha, \beta} \frac{C_{\alpha} C_{\beta}}{\Lambda^4} \cdot (\mathcal{A}_{Q\alpha}^{(6)} \mathcal{A}_{Q\beta}^{(6)\dagger}) +$$

Dim6 including linear, BSM and mixed contributions

$$\sum_k \left[ \frac{f_k}{\Lambda^4} \cdot 2\text{Re}(\mathcal{A}_{SM} \mathcal{A}_{Qk}^{(8)\dagger}) \right] + \sum_k \frac{f_k^2}{\Lambda^8} \cdot (\mathcal{A}_{Qk}^{(8)} \mathcal{A}_{Qk}^{(8)\dagger})$$

Dim8 including linear and BSM contributions

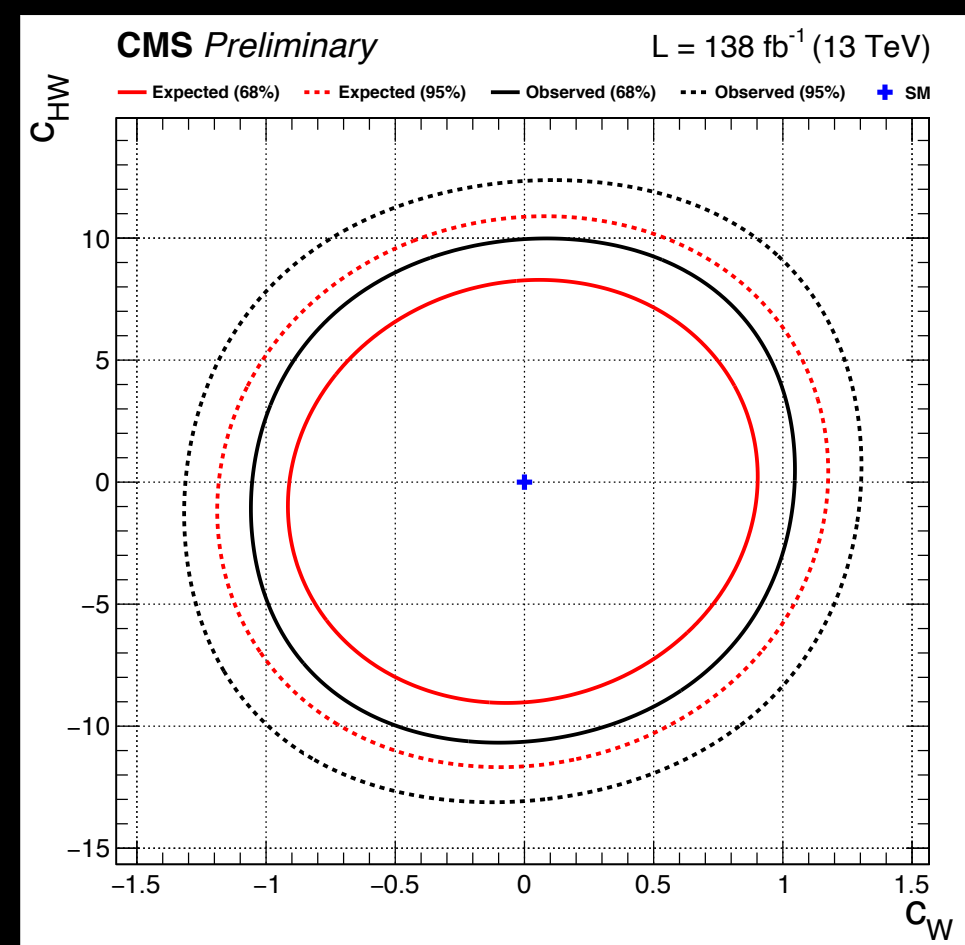
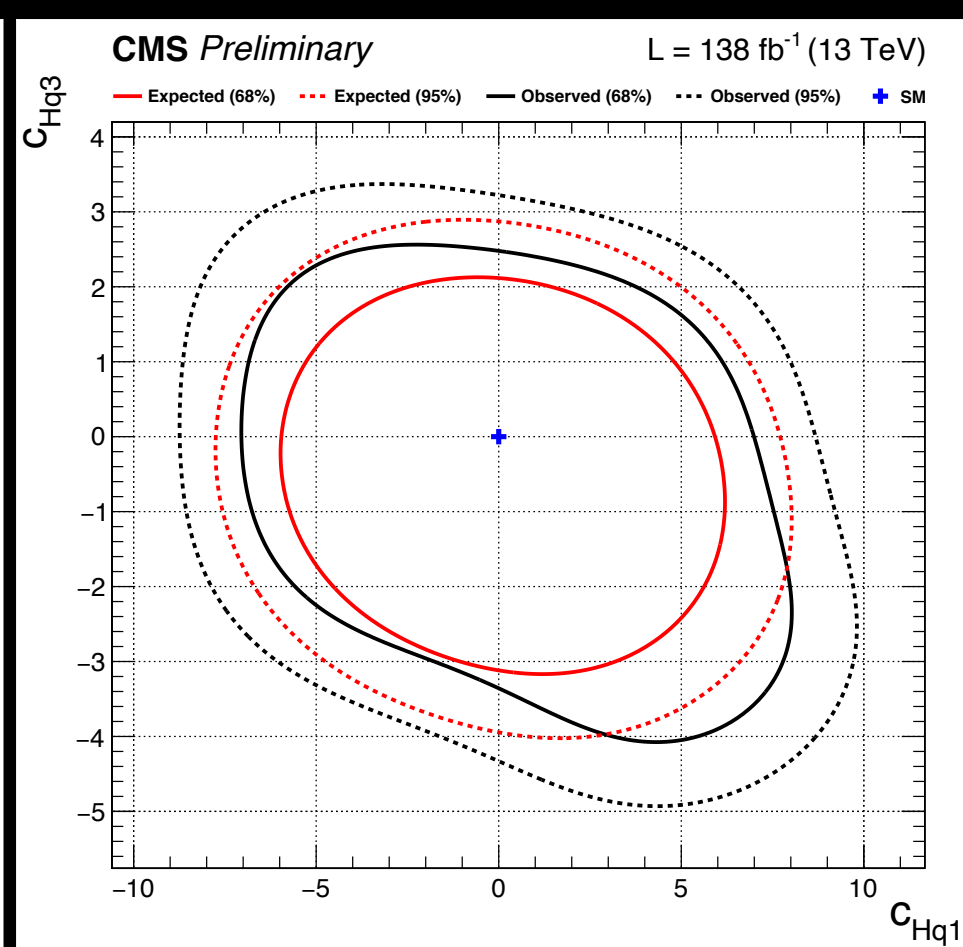
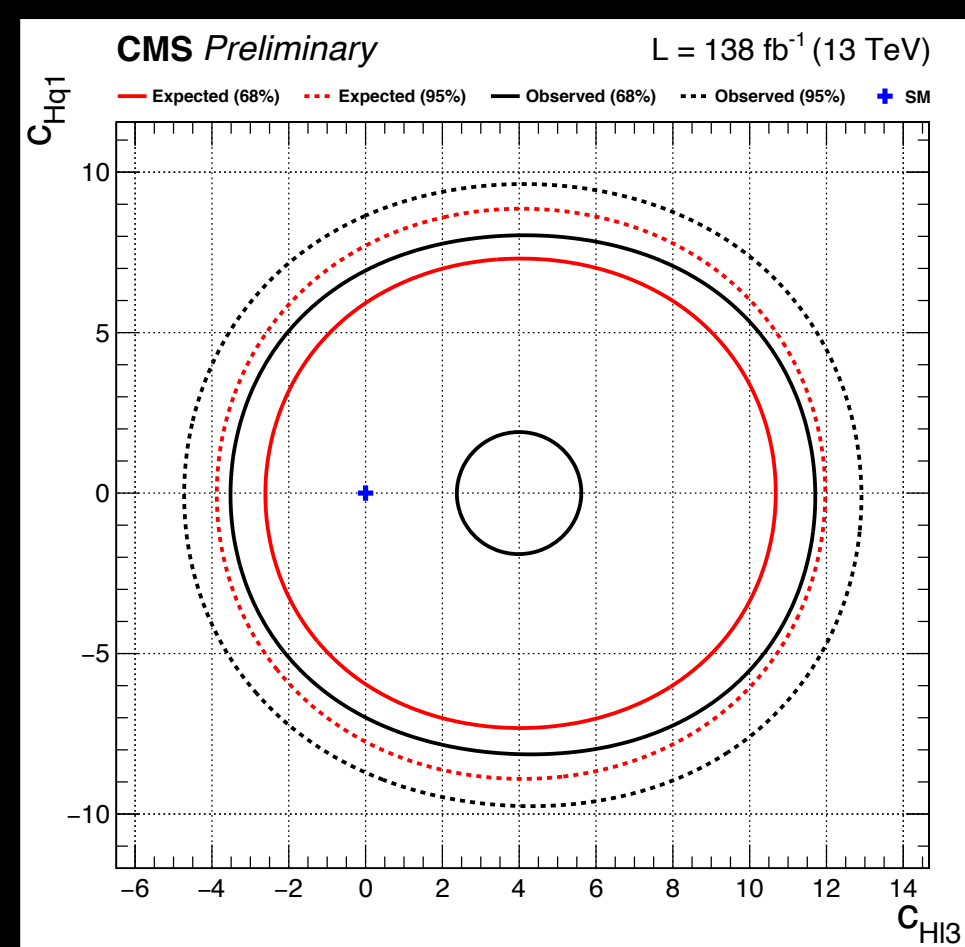


$$M_{01} = (p_{T\tau} + p_l + \vec{p}_T^{miss})^2 + |p_{T\tau} + \vec{p}_T^{miss}|^2$$

# VBS Same Sign $WW$ with hadronic $\tau$

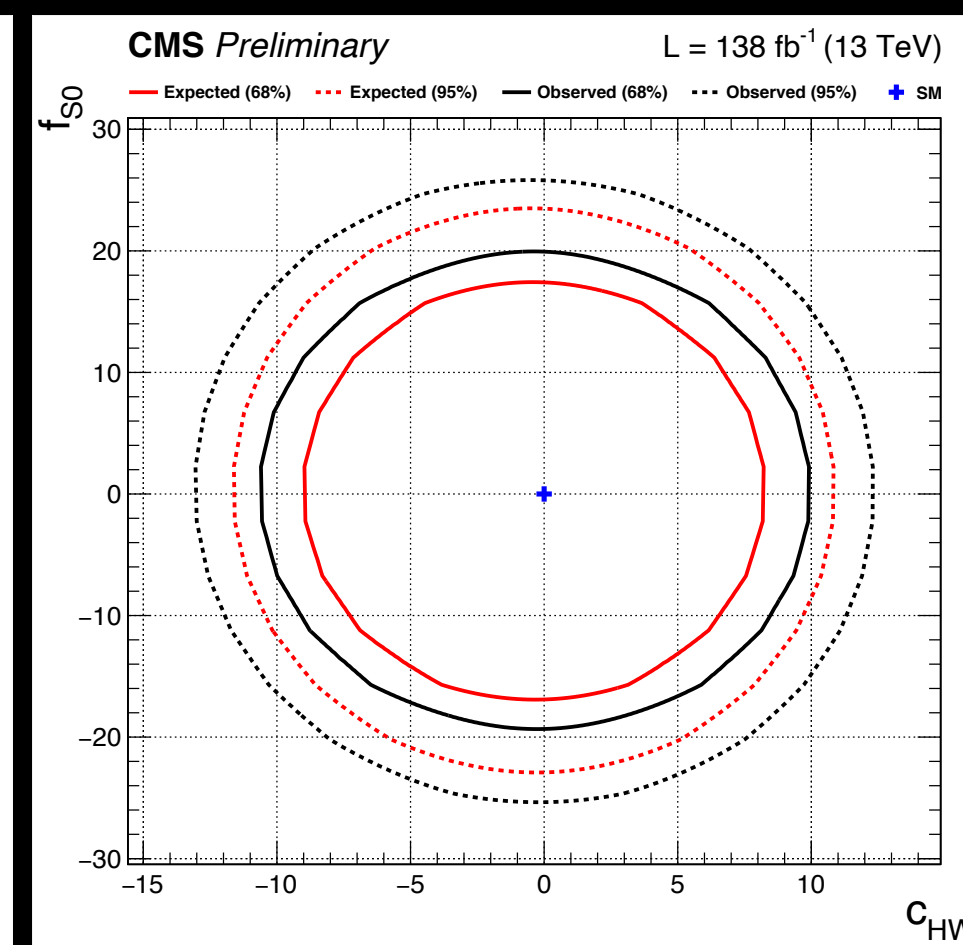
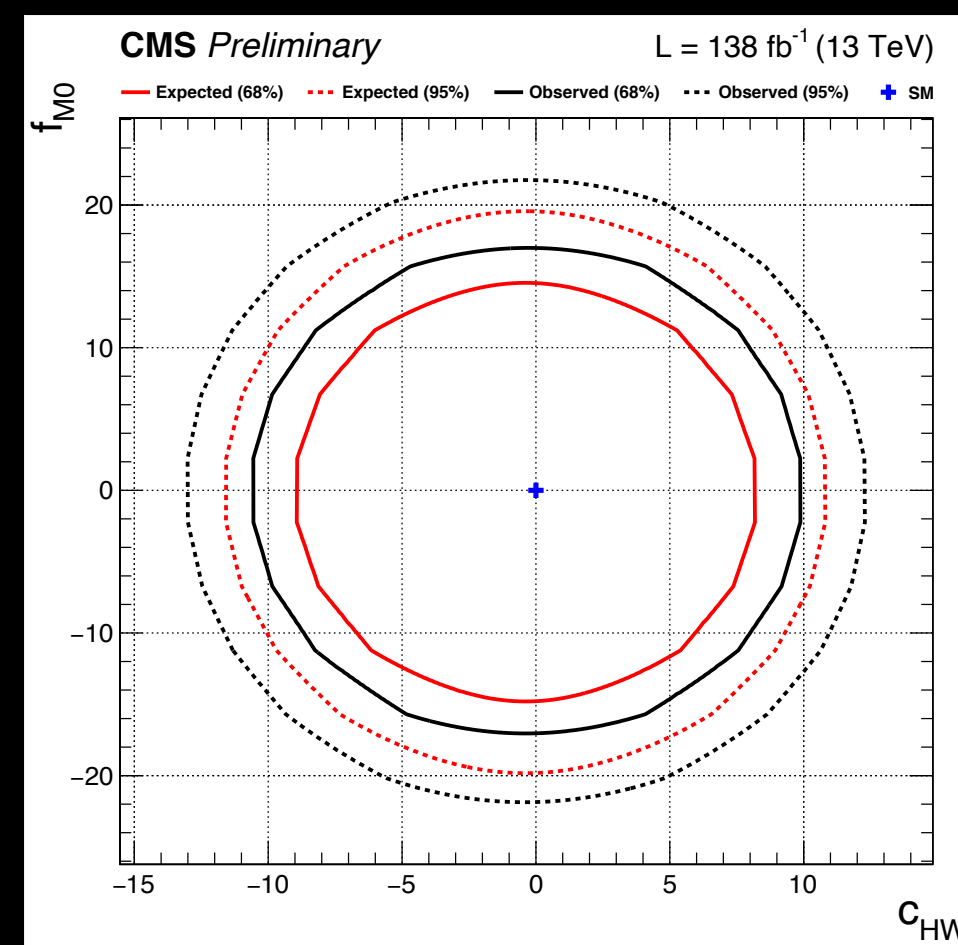
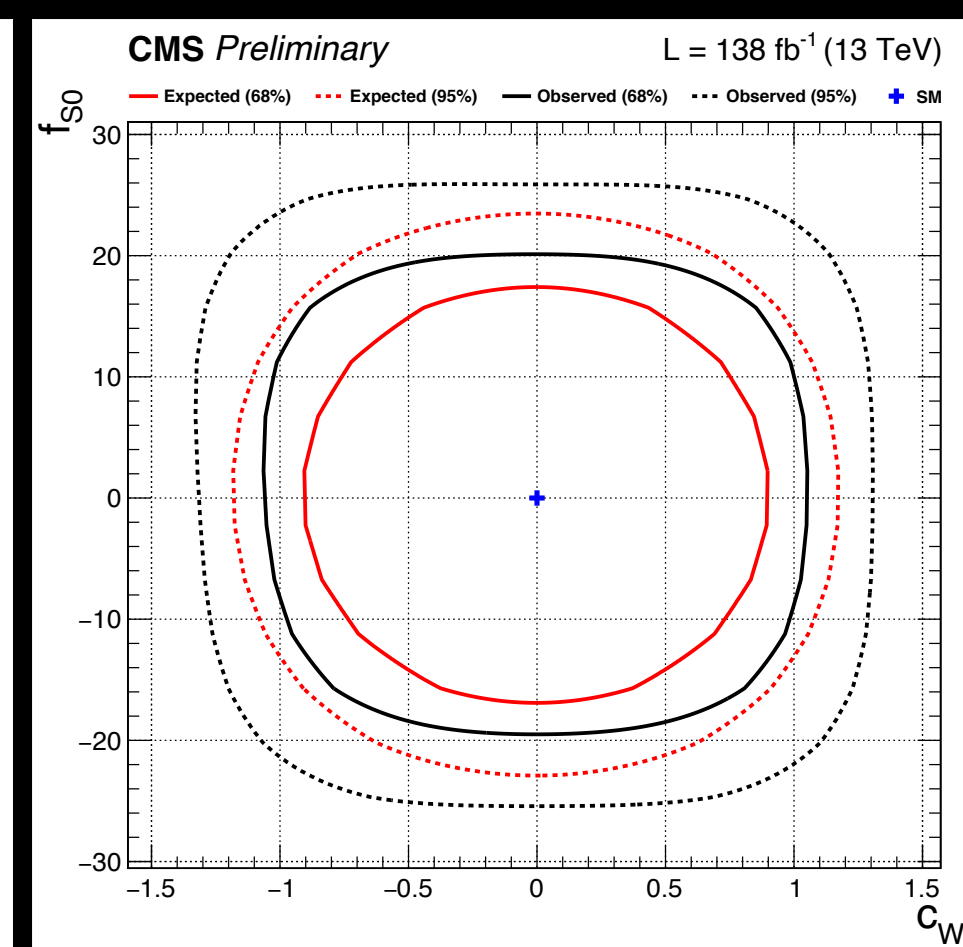
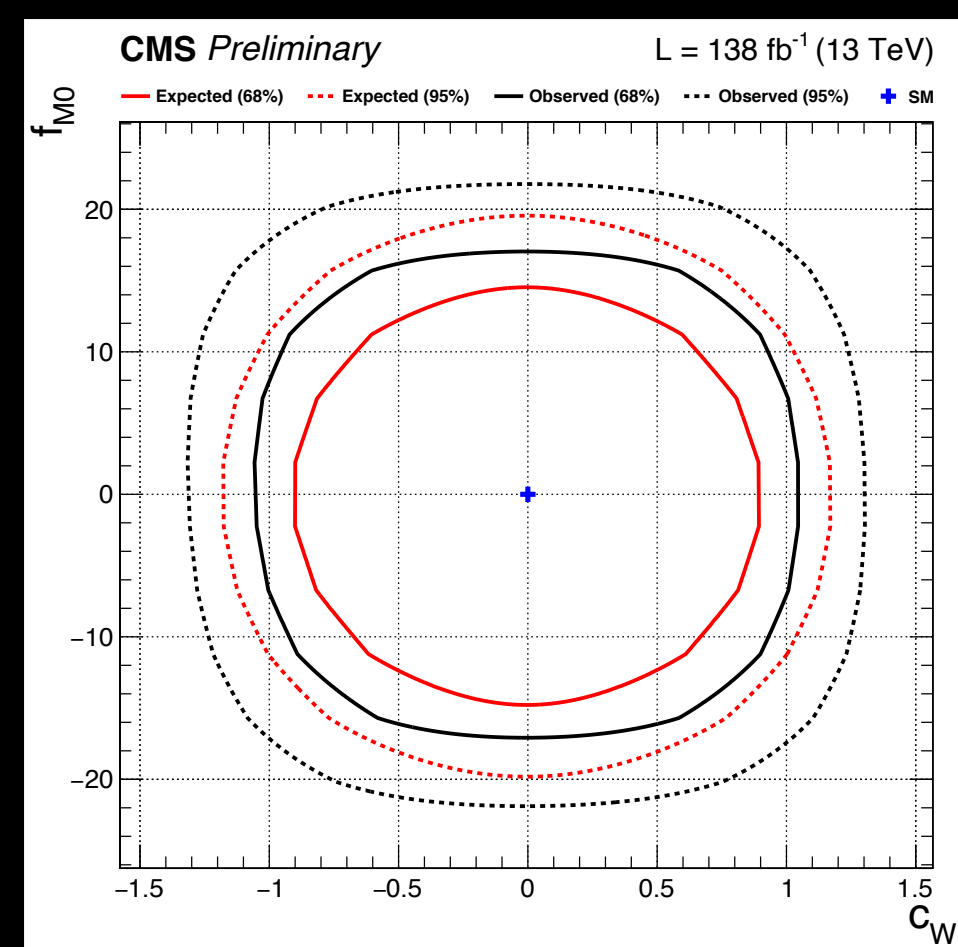
$$M_{01} = (p_{T\tau} + p_l + \vec{p}_T^{miss})^2 + |p_{T\tau} + \vec{p}_T^{miss}|^2$$

- Simultaneous extraction of dim-6 and dim-8 constraints

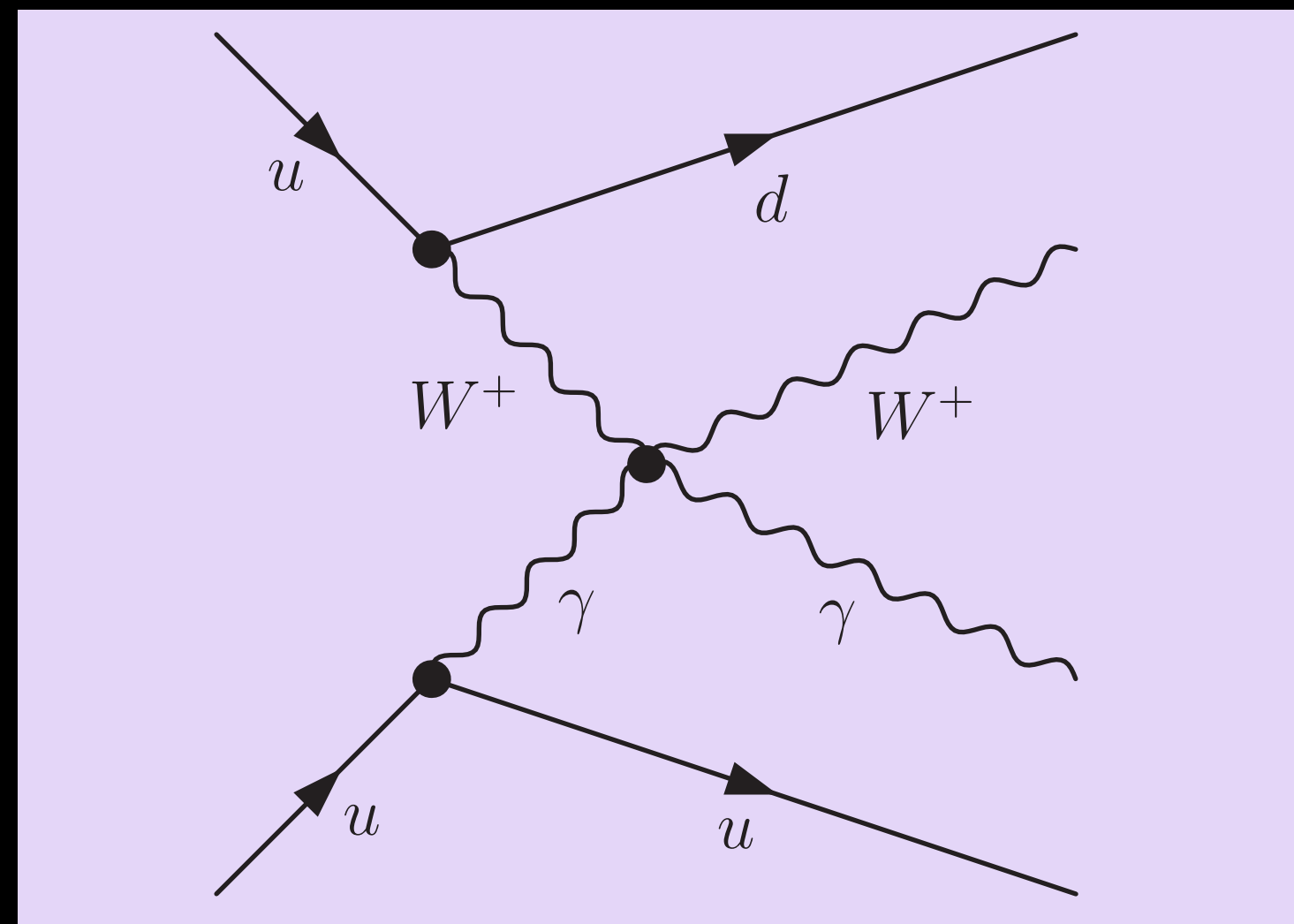


Sensitive to correlation between Higgs-fermion operators

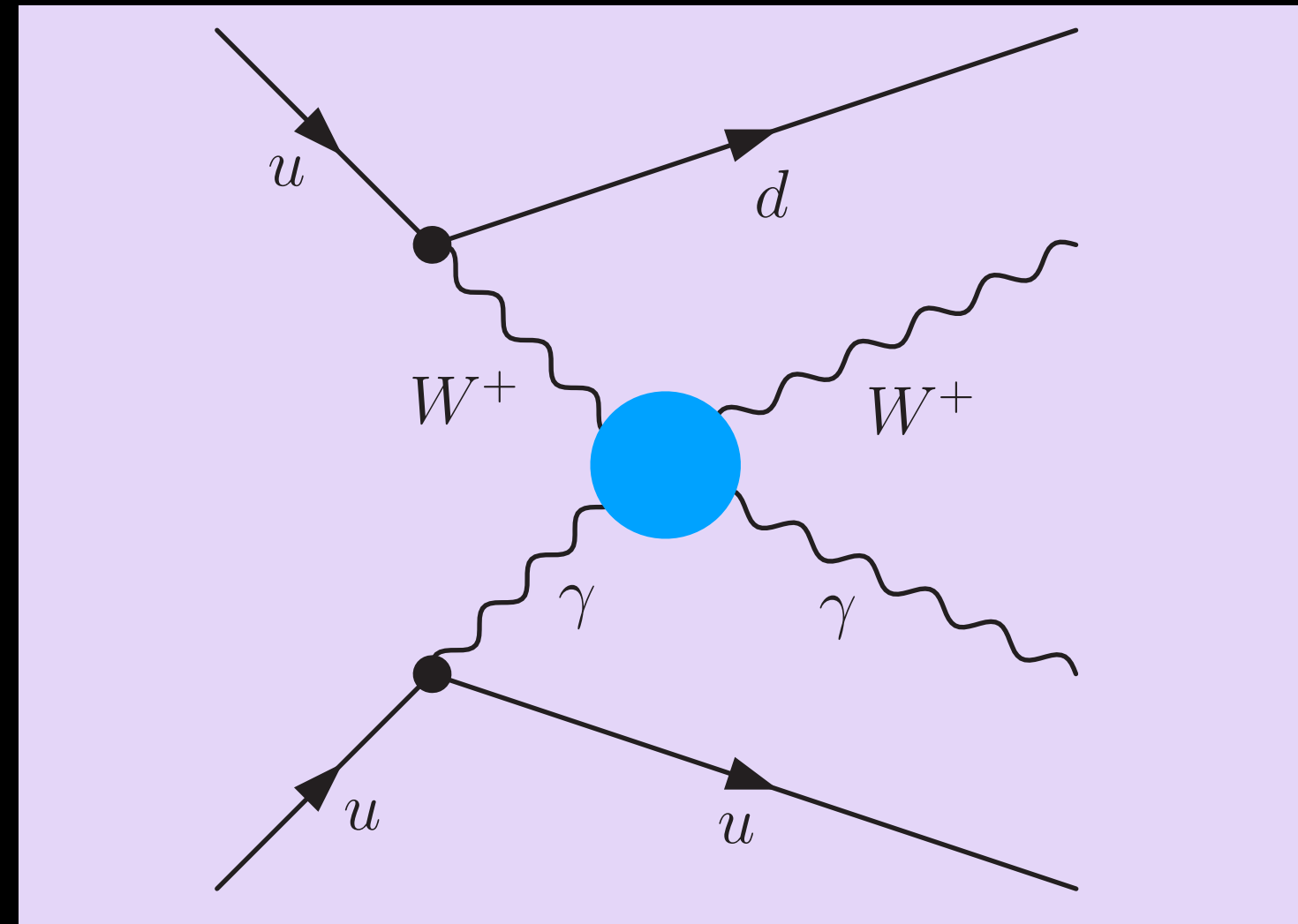
Fits between boson operators of dim-6 and dim-8



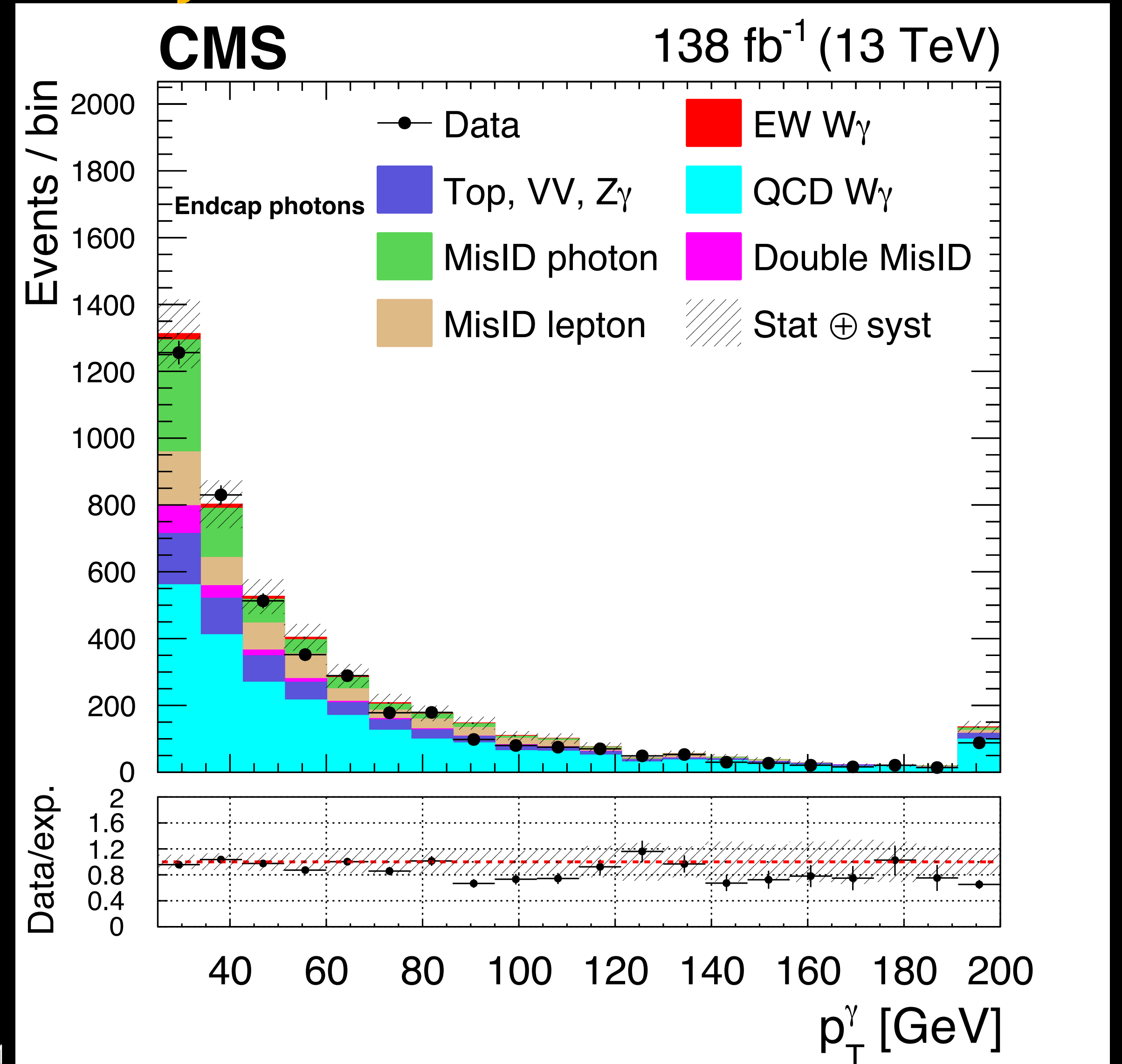
# Measurement of the electroweak production of $W\gamma$ with two jets



# Measurement of the electroweak production of $W\gamma$ with two jets



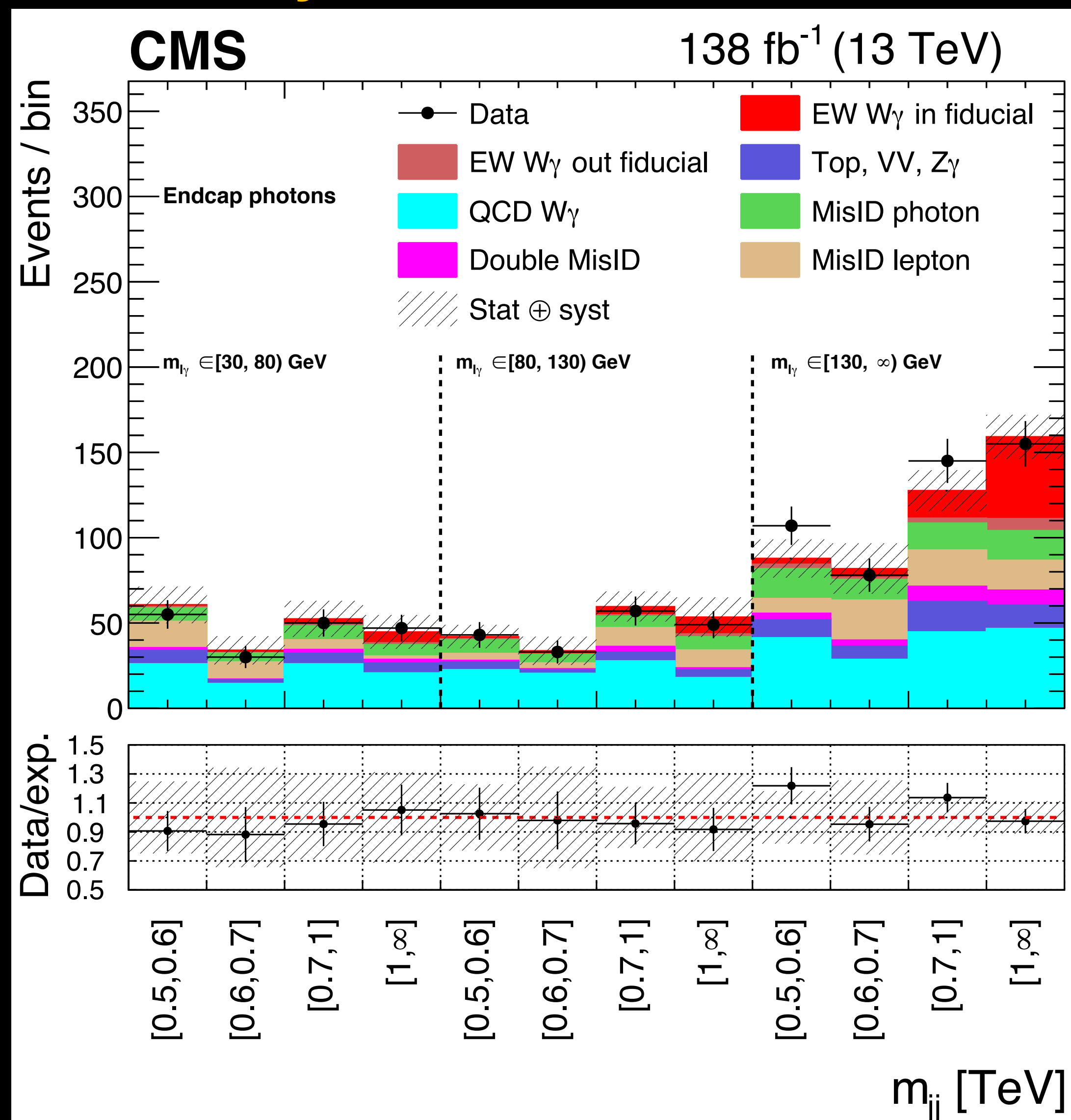
- First observation of the  $W\gamma + 2$  jets process with observed (expected) significance of 5.3 (4.8)  $\sigma$  with  $35.9 \text{ fb}^{-1}$
- Extensive measurement now possible with full Run II dataset
- Invariant mass of the  $W\gamma$  system is sensitive to presence of dim-8 operators



# Measurement of the electroweak production of $W\gamma$ with two jets

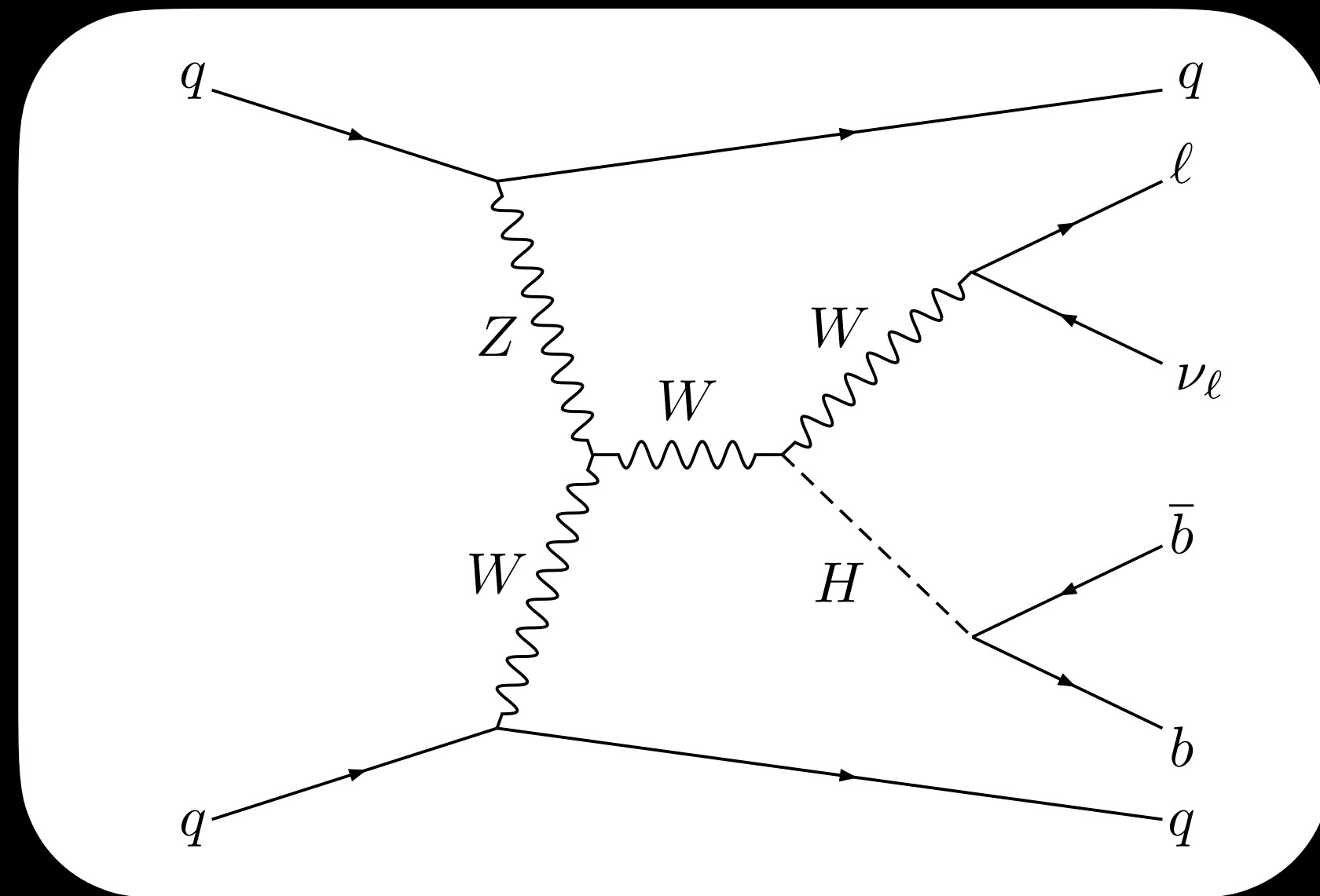
- Major backgrounds from  $W$ +jets and  $t\bar{t}$  processes where the jet constituents is misidentified as a photon

	Barrel	Endcap
EW $W\gamma$ in fiducial region	$316 \pm 16$	$90.2 \pm 5.5$
EW $W\gamma$ out of fiducial region	$64.7 \pm 2.0$	$20.4 \pm 1.0$
QCD $W\gamma$	$1301 \pm 28$	$362 \pm 13$
top, $VV, Z\gamma$	$402 \pm 14$	$93.3 \pm 7.2$
Nonprompt photon	$434 \pm 13$	$120.2 \pm 5.7$
Nonprompt muon	$134 \pm 27$	$45 \pm 11$
Nonprompt electron	$189 \pm 20$	$86 \pm 13$
Nonprompt photon, nonprompt muon	$43.0 \pm 7.0$	$14.6 \pm 3.4$
Nonprompt photon, nonprompt electron	$75.5 \pm 5.5$	$25.0 \pm 2.0$
Total prediction	$2960 \pm 43$	$856 \pm 21$
Data	$2959 \pm 57$	$849 \pm 32$

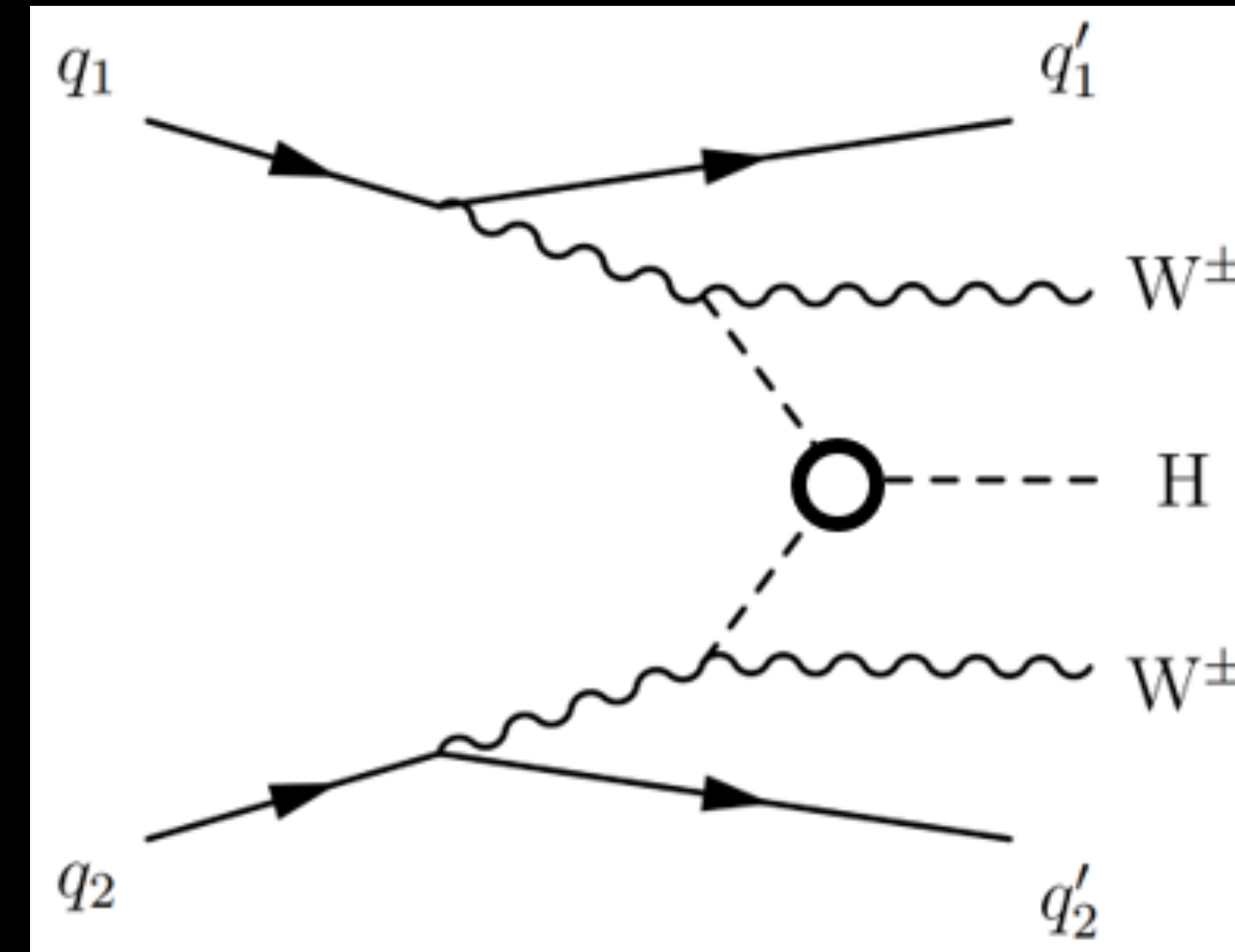




# Search for $HWW/HHWW$ couplings in the VBS production of $W^\pm W^\pm H$ with $H \rightarrow b\bar{b}$ decays



<https://arxiv.org/pdf/2405.16566>



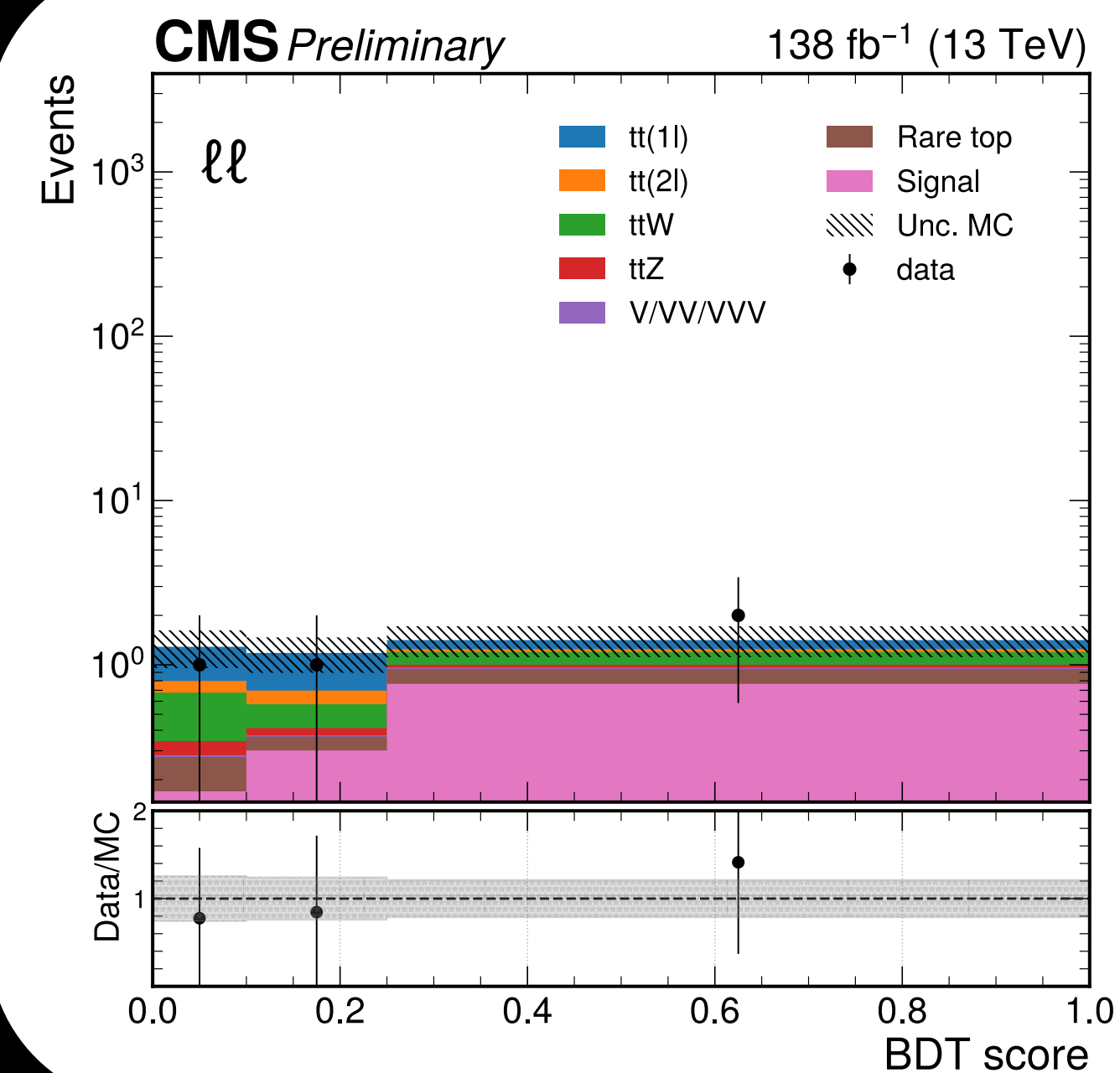
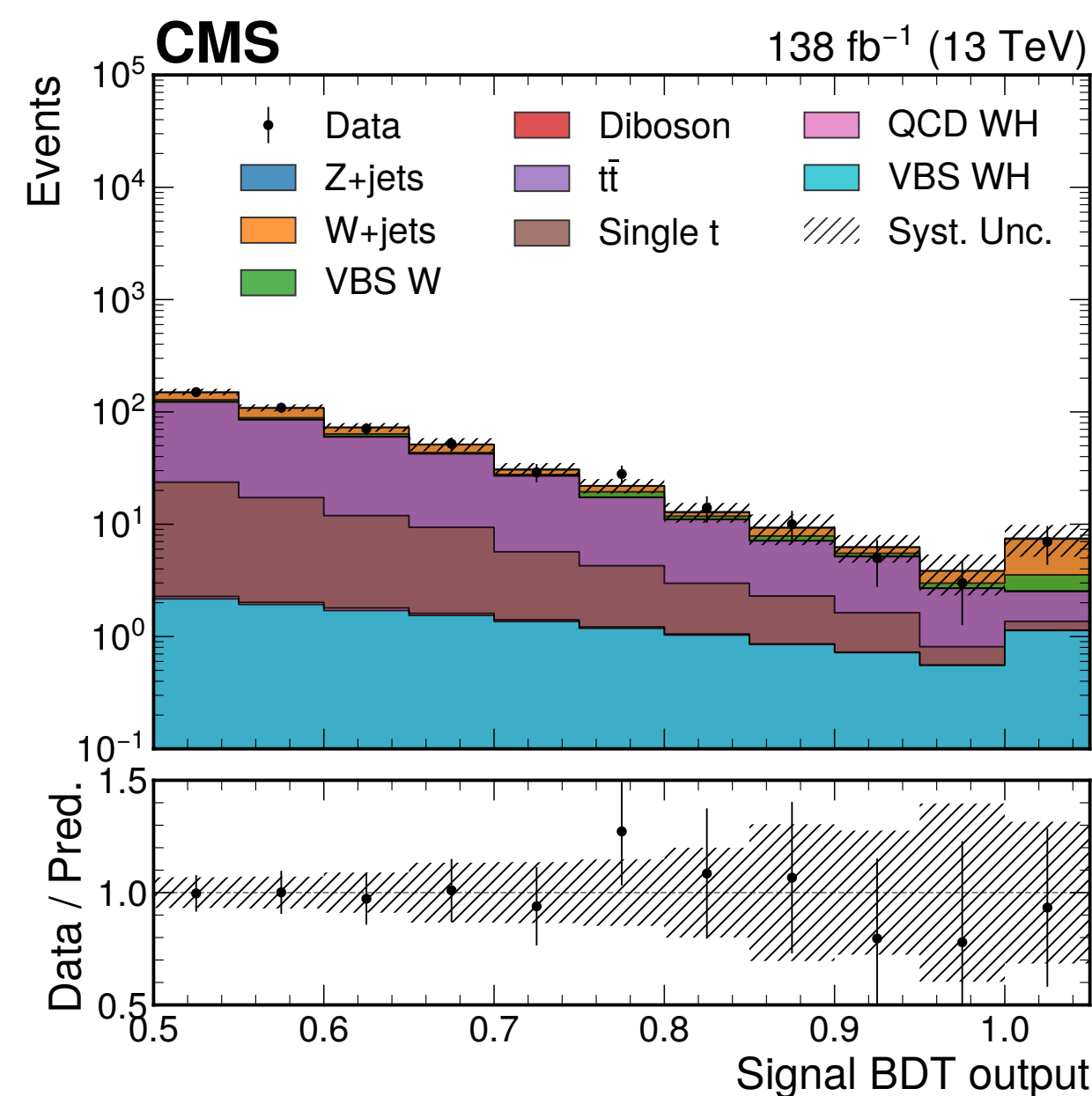
<https://cds.cern.ch/record/2905615>

# Search for $HHWW$ couplings in the VBS production of $W^\pm W^\pm H$ with $H \rightarrow b\bar{b}$ decays

- Extraction of sign of W and Z coupling to Higgs performed with W and Higgs in final state
- Boosted decision tree trained against major backgrounds
- All scenarios with  $\lambda_{WZ} = \frac{k_W}{k_Z} < 0$  excluded
- For  $W^\pm W^\pm H$  low signal yields  $\rightarrow$  need more data to set tight constraints on couplings

## $W^\pm W^\pm H$

Shorthand	Description
$\eta_J$	$\eta$ of the leading merged jet
$p_{T,J}$	$p_T$ of the leading merged jet
$p_{T,jj}$	$p_T$ of the VBS-jet system
$P_{j_0}$	magnitude of the three-momentum of the leading VBS jet
$P_{j_1}$	magnitude of the three-momentum of the subleading VBS jet
$M_{\ell\ell}$	invariant mass of the SS dilepton system
$p_{T,\ell_0}$	$p_T$ of the leading lepton
$p_{T,\ell_1}$	$p_T$ of the subleading lepton
$E_T^{\text{miss}}$	missing transverse energy
$L_T$	scalar sum of $p_{T,\ell_0}, p_{T,\ell_1}$ , and $E_T^{\text{miss}}$
$S_T$	scalar sum of $p_{T,J}$ and $L_T$



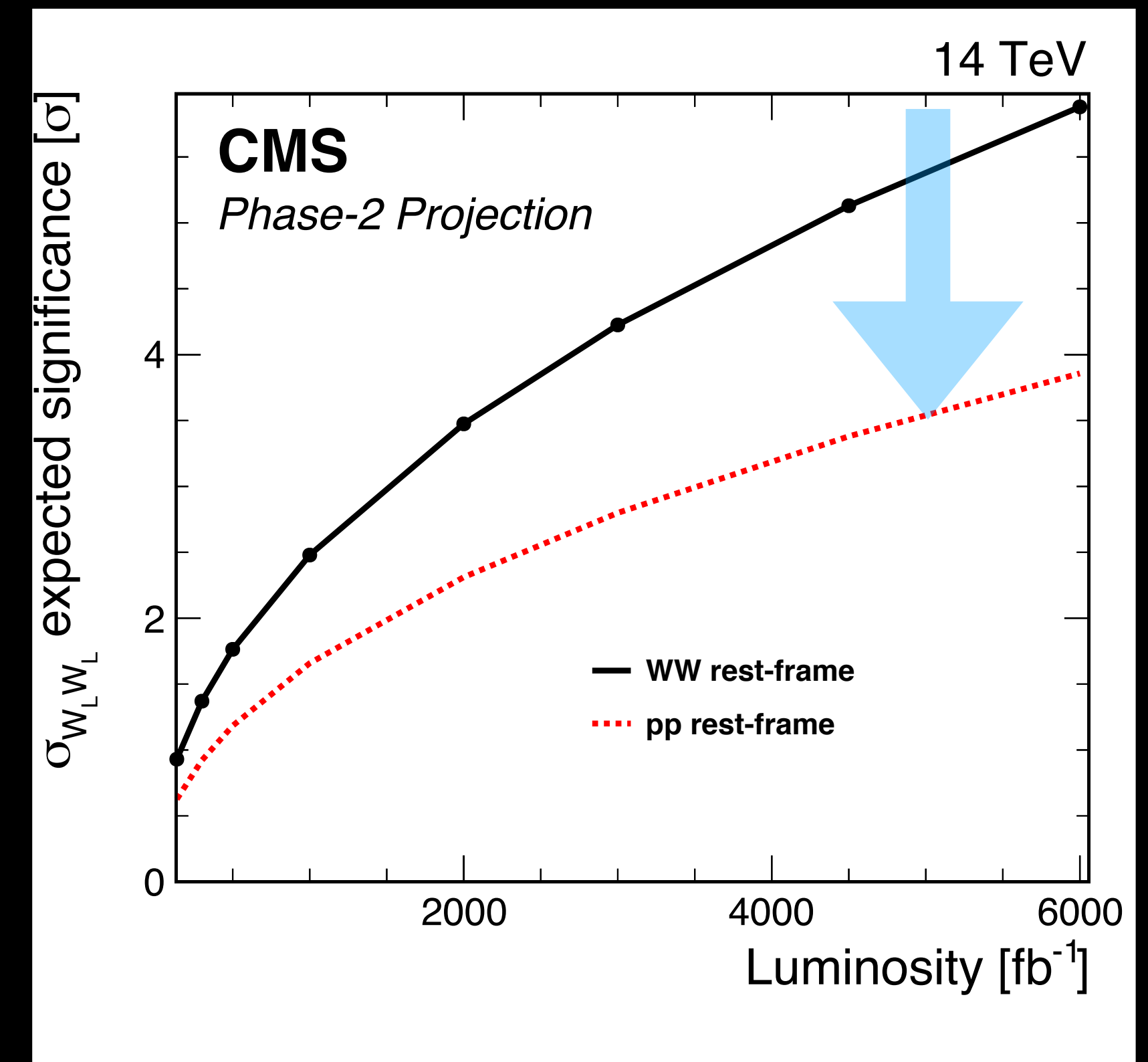
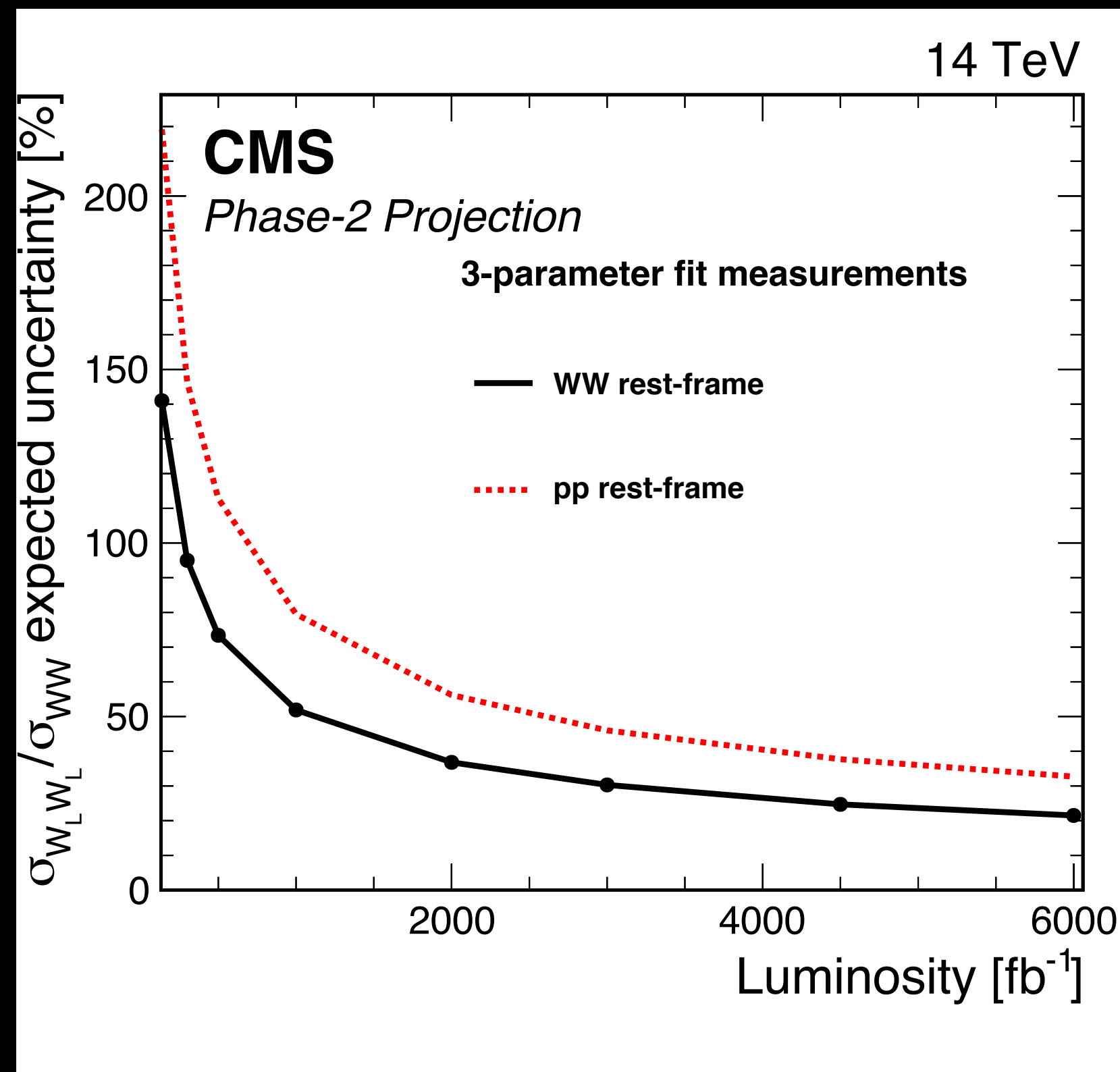
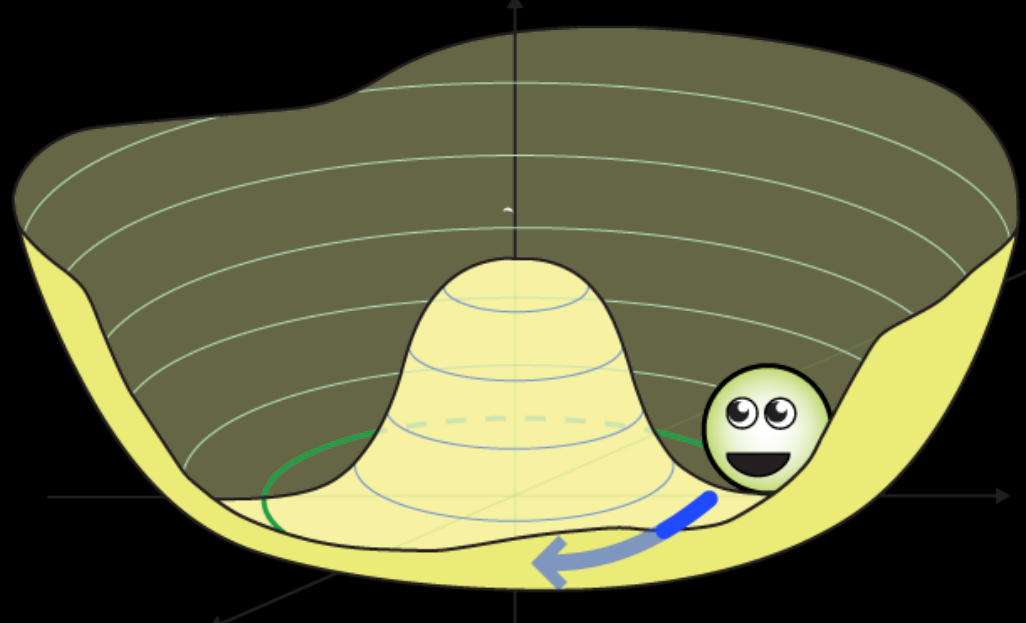
# Going forward... (practicalities)

- Study of vector boson scattering processes is an exciting area of research with many new analyses in the last few years
- Run III dataset provides additional opportunities to study these processes in depth
- As the precision program of the LHC is realized, studying generator modeling is crucial
  - Many different generators studied with various configurations
    - Features discussion of choice of dipole recoil scheme
    - Synchronize sample production between ATLAS and CMS

Sample name	Generator	$\mu$ -scale	Shower	Tune	PDF	further settings
Sherpa (ATLAS)	SHERPA v2.2.2	dynamic scale, $m_{WW}$	internal	internal	NNPDF3.0-NLO	multileg-LO, exactly six EW vertices with one additional parton at LO accuracy in QCD
PW+Py8 (ATLAS)	POWHEG v2, VBS approx.	fixed scale, $m_W$	PYTHIA 8.212	AZNLO	NNPDF3.0-NLO	NLO
PW+Py8 dipole-recoil (ATLAS)	POWHEG v2	fixed scale, $m_W$	PYTHIA 8.235	AZNLO	NNPDF3.0-NLO	Dipole Recoil [6]
MG5+Py8 dipole-recoil (ATLAS)	MG5_AMCNLO v2.6.2	dynamic scale, $\sqrt{p_T^{\text{jet1}} p_T^{\text{jet1}}}$	PYTHIA 8.235	A14	NNPDF3.0-NLO	LO, Dipole Recoil [6]
MG5+Py8 (CMS)	MG5_AMCNLO v2.3.3	dynamic scale, using a 2→2 topology from the clustered external state	PYTHIA 8.212	CUETP8M1 [7]	NNPDF3.0-LO	LO, exactly six EW vertices
PW+Py8 (VBSscan)	POWHEG v2	dynamic scale, $\sqrt{p_T^{\text{jet1}} p_T^{\text{jet2}}}$	PYTHIA 8.230	Monash	NNPDF3.0-NLO	NLO

# Going forward...

- Understanding electroweak symmetry breaking → crucial part of LHC physics program
- Longitudinally polarized scattering of W and Z complementary to direct measurements of the Higgs coupling to gauge bosons
- Analysis projected from Run II to 3000 fb<sup>-1</sup>



# Anomalous neutral gauge couplings in the mono photon channel

The diagram shows a central shaded circular vertex. An incoming wavy line from the left is labeled  $Z^*/\gamma^*(V)$  and has momentum  $p$ . Two outgoing wavy lines to the right are labeled  $Z$  (top) and  $\gamma$  (bottom), with momenta  $q_1$  and  $q_2$  respectively.

$$Z^*/\gamma^*(V) \sim p \text{ --- } \text{shaded circle} \text{ --- } \begin{matrix} q_1 \\ Z \\ q_2 \\ \gamma \end{matrix} = ie\Gamma_{ZV\gamma}(p, q_1, q_2; \vec{h}^V)$$

# Anomalous neutral gauge couplings in the mono photon channel

- Generalized theory of forbidden neutral gauge couplings
  - Parametrized with 8 parameters ( $h_1$  to  $h_4$ )
  - Vertex functions describe the most general Lorentz and  $U_Y(1)$  invariant interactions of the incoming off-shell Z or  $\gamma$  and outgoing on-shell Z and  $\gamma$
  - CP violating:  $h_1^V, h_2^V$ ; CP conserving:  $h_3^V, h_4^V$

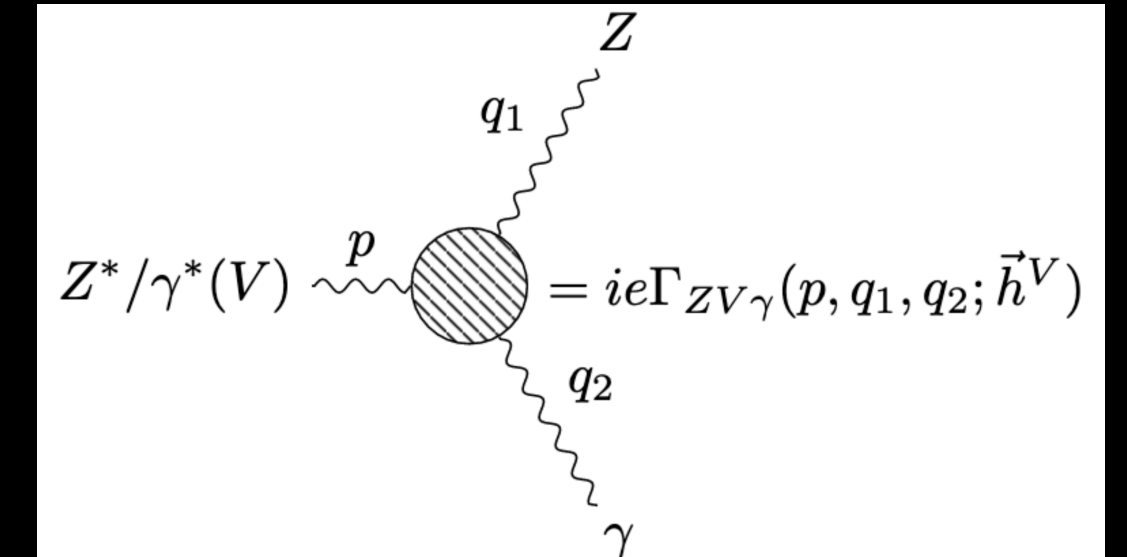
$$\Gamma_{ZZ\gamma}^{\alpha\beta\mu} = \frac{p^2 - q_1^2}{m_Z^2} \left[ h_1^Z (q_2^\mu g^{\alpha\beta} - q_2^\alpha g^{\mu\beta}) + \frac{h_2^Z}{m_Z^2} p^\alpha [(p \cdot q_2) g^{\mu\beta} - q_2^\mu p^\beta] + h_3^Z \epsilon^{\mu\alpha\beta\rho} q_{2\rho} + \frac{h_4^Z}{m_Z^2} p^\alpha \epsilon^{\mu\beta\rho\sigma} p_\rho q_{2\sigma} \right]$$

$\Gamma_{Z\gamma\gamma}^{\alpha\beta\mu}$  obtained with some simple substitutions to the above equation

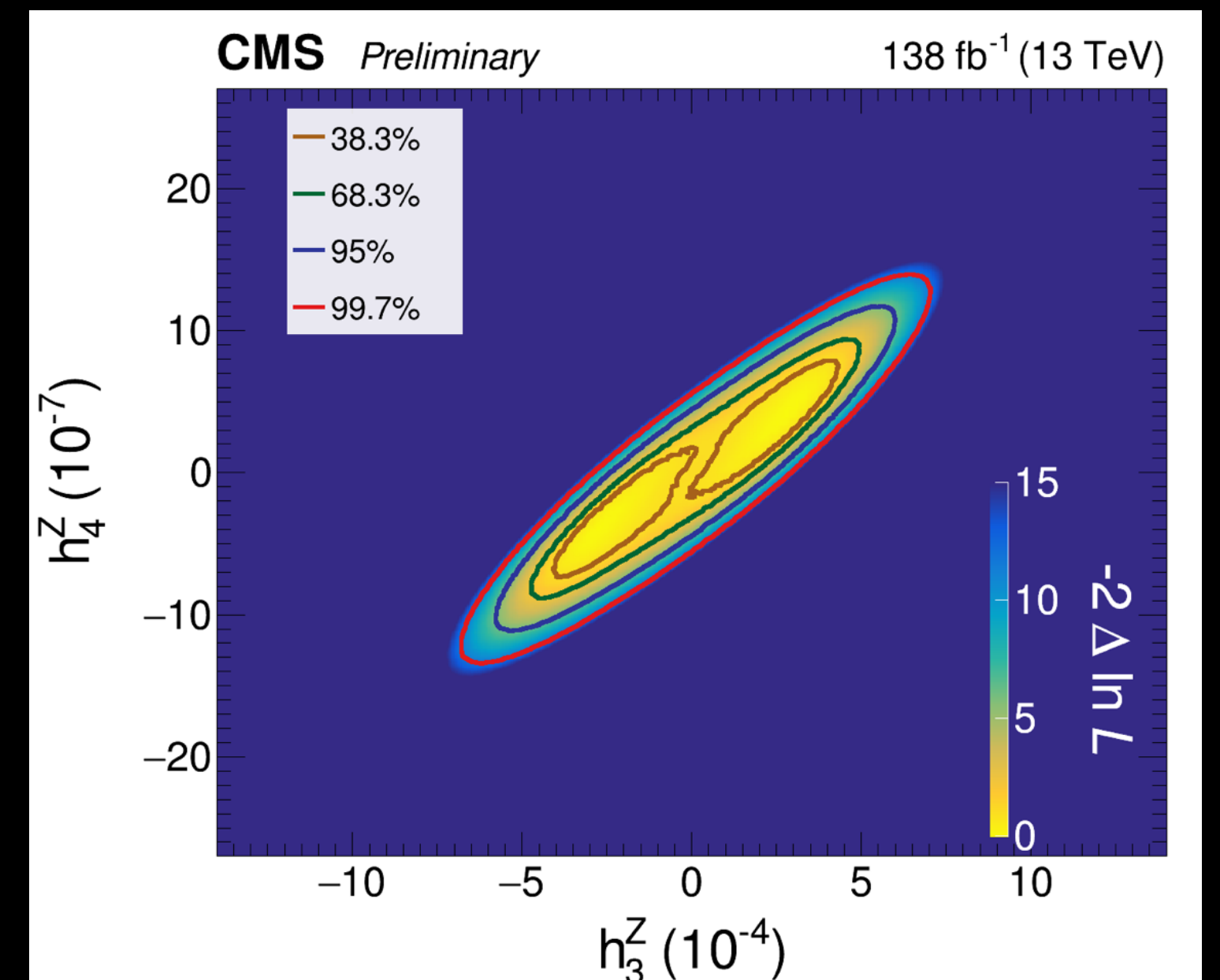
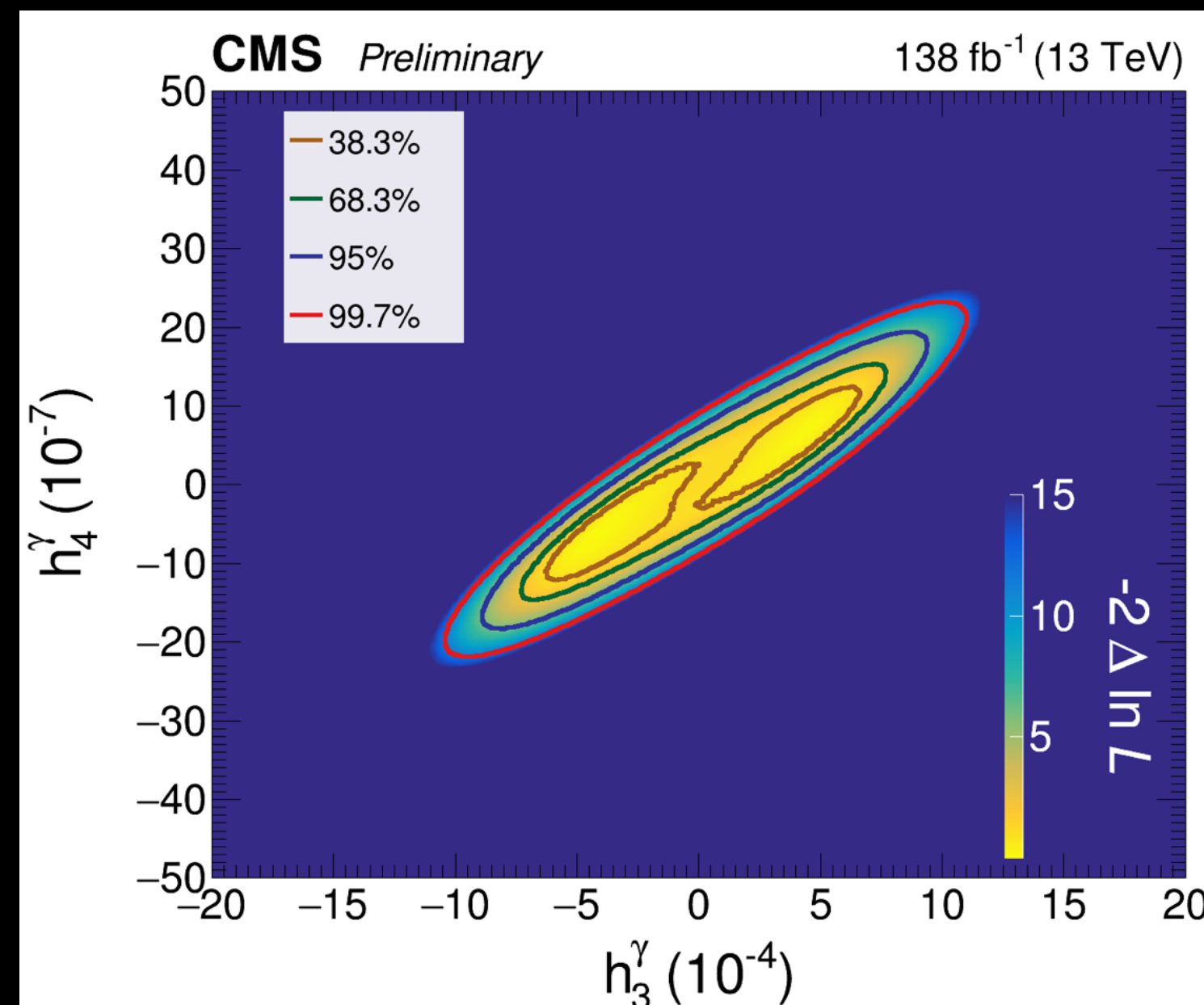
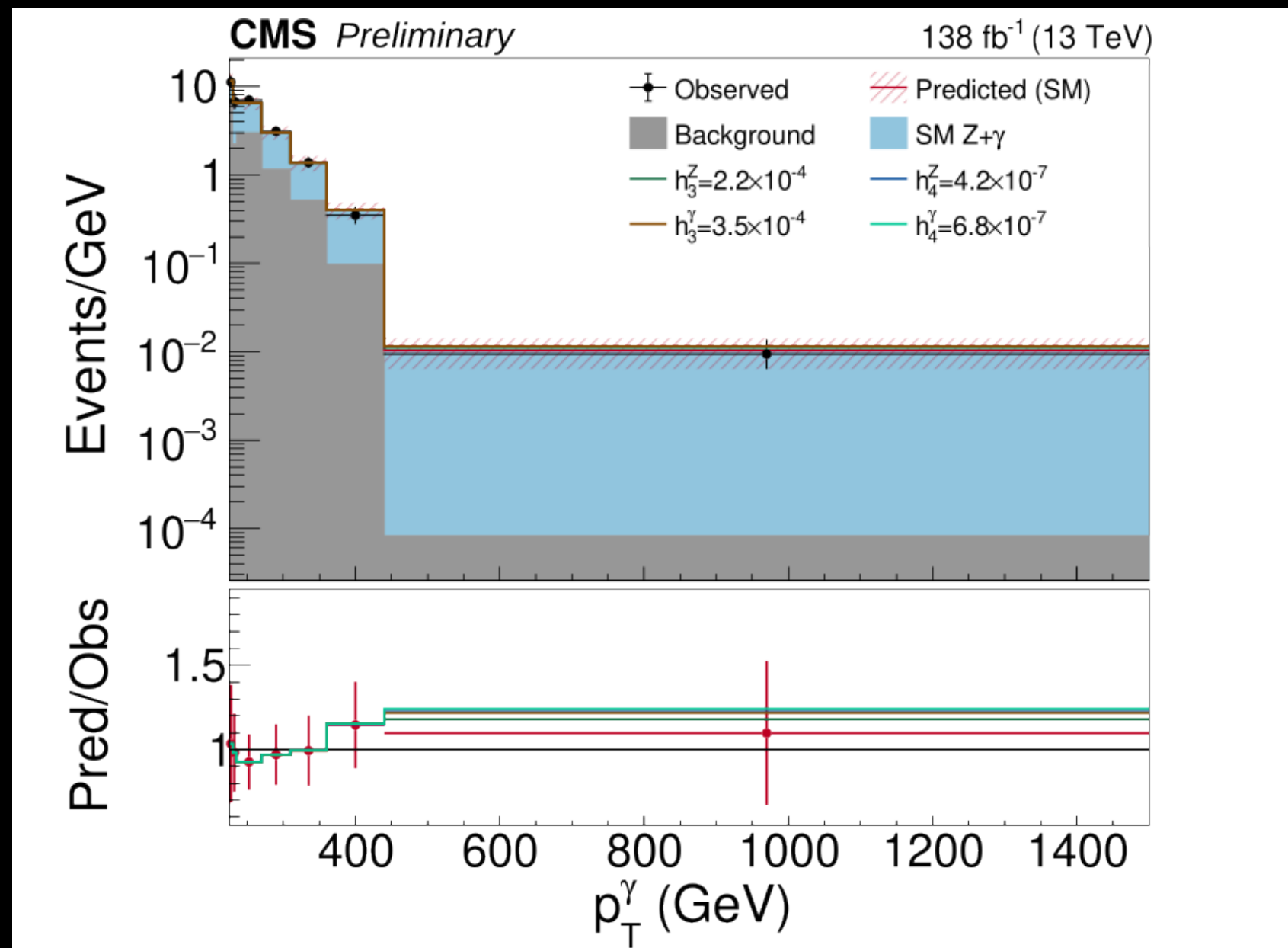
Ellis et al:

[https://indico.cern.ch/event/1330671/contributions/5601599/attachments/2735050/4755768/Offshell+CPV\\_Dimension8.pdf](https://indico.cern.ch/event/1330671/contributions/5601599/attachments/2735050/4755768/Offshell+CPV_Dimension8.pdf)

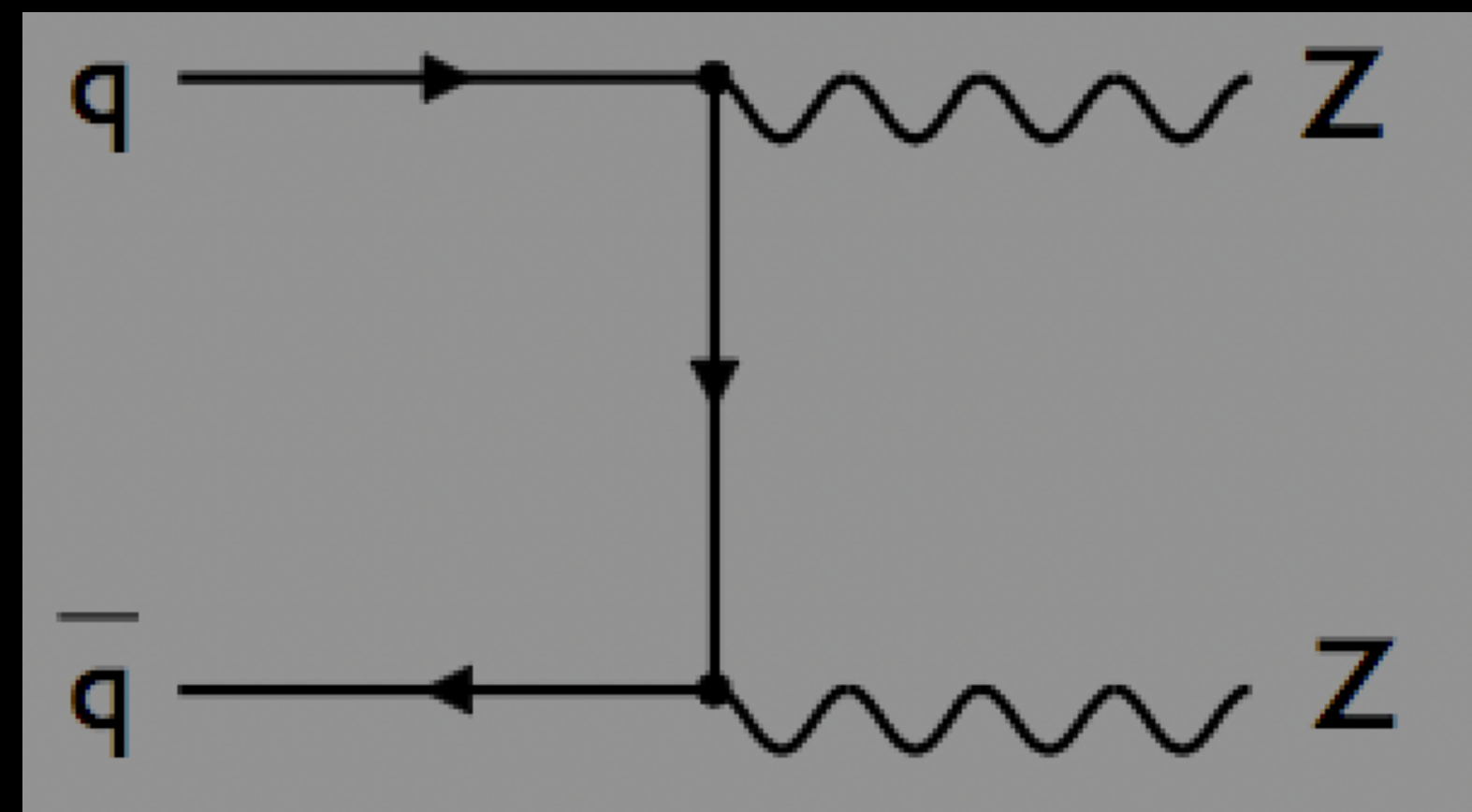
# Anomalous neutral gauge couplings in the mono photon channel



- New physics most likely to show up as deviation at high  $p_T$ 
  - Photon  $p_T$  is the variable of interest
- Backgrounds range from  $W + \gamma$ ,  $V + \text{jets}$ ,  $\gamma + \text{jets}$  and **dibosons** estimated using **data-driven methods** and **simulations**



Search for  $ZZ \rightarrow b\bar{b}$  and  $ZH \rightarrow b\bar{b}$



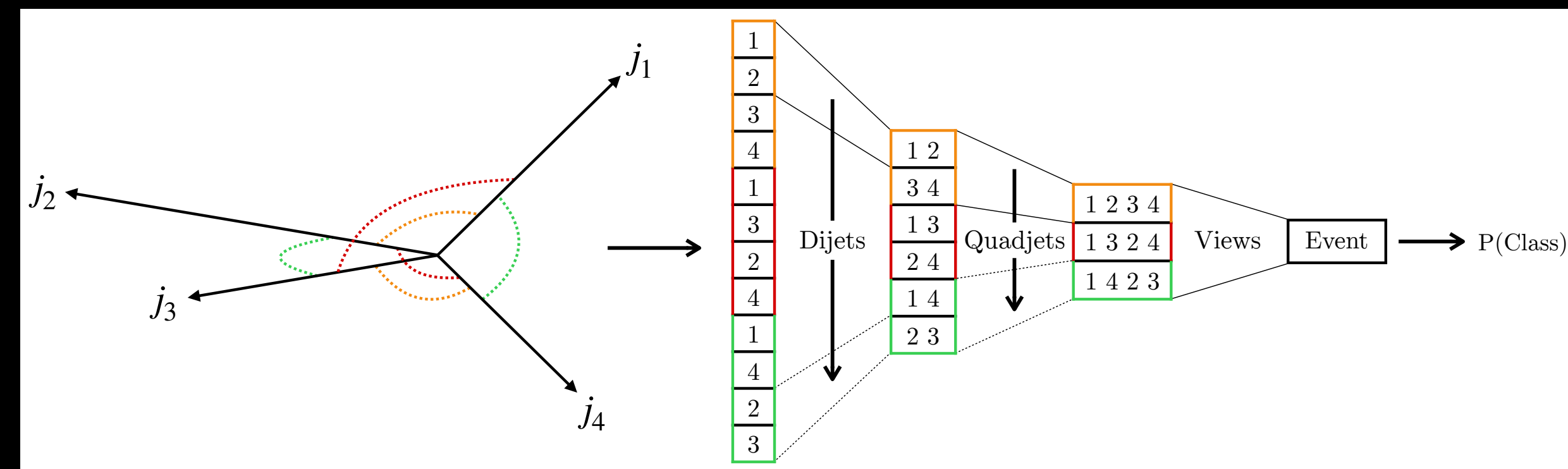
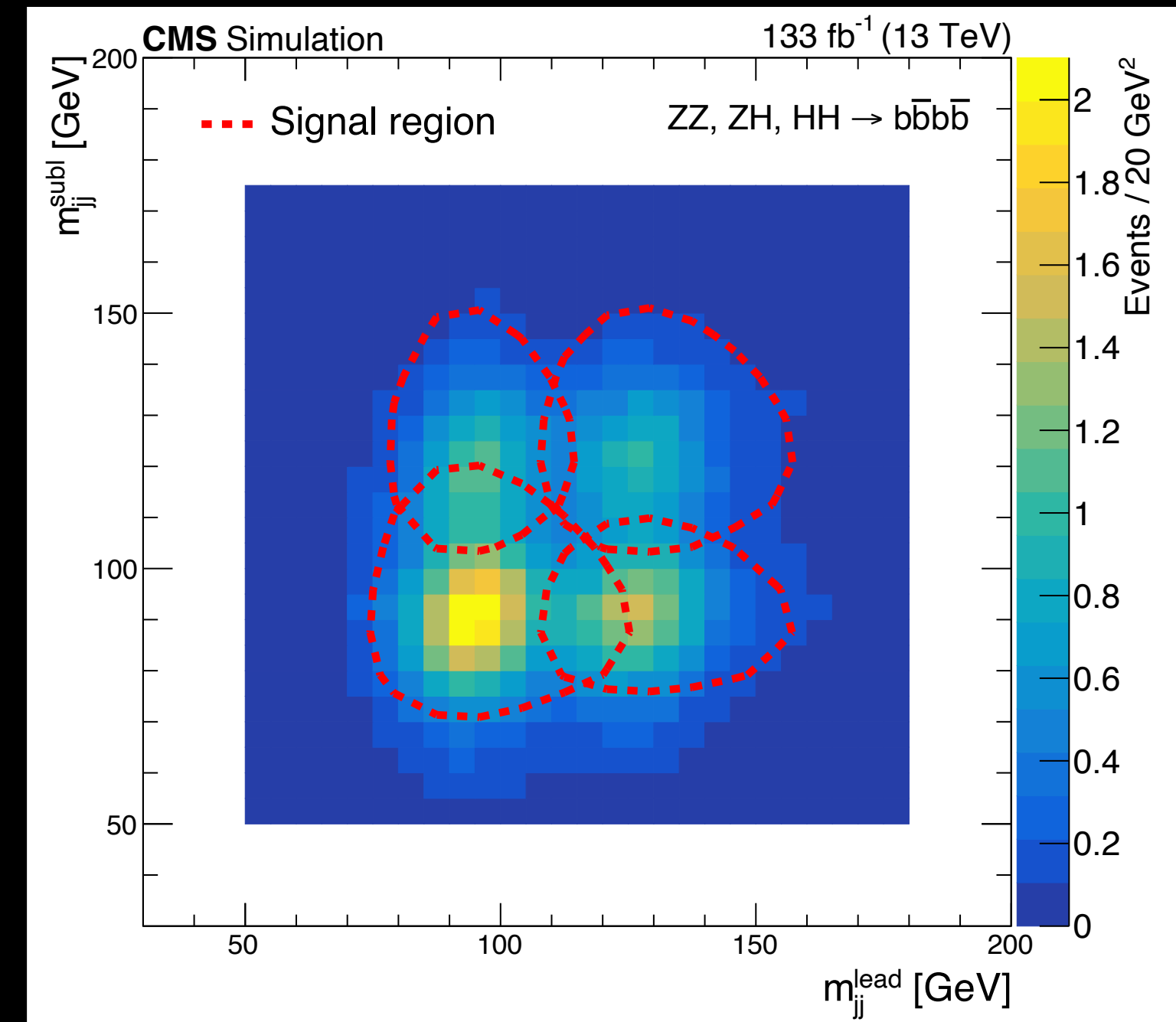


# Search for $ZZ \rightarrow b\bar{b}$ and $ZH \rightarrow b\bar{b}$

- Diboson  $ZZ$  process is interesting to study in its own right
  - Additionally, major background for DiHiggs measurement
- Define signal regions based on:

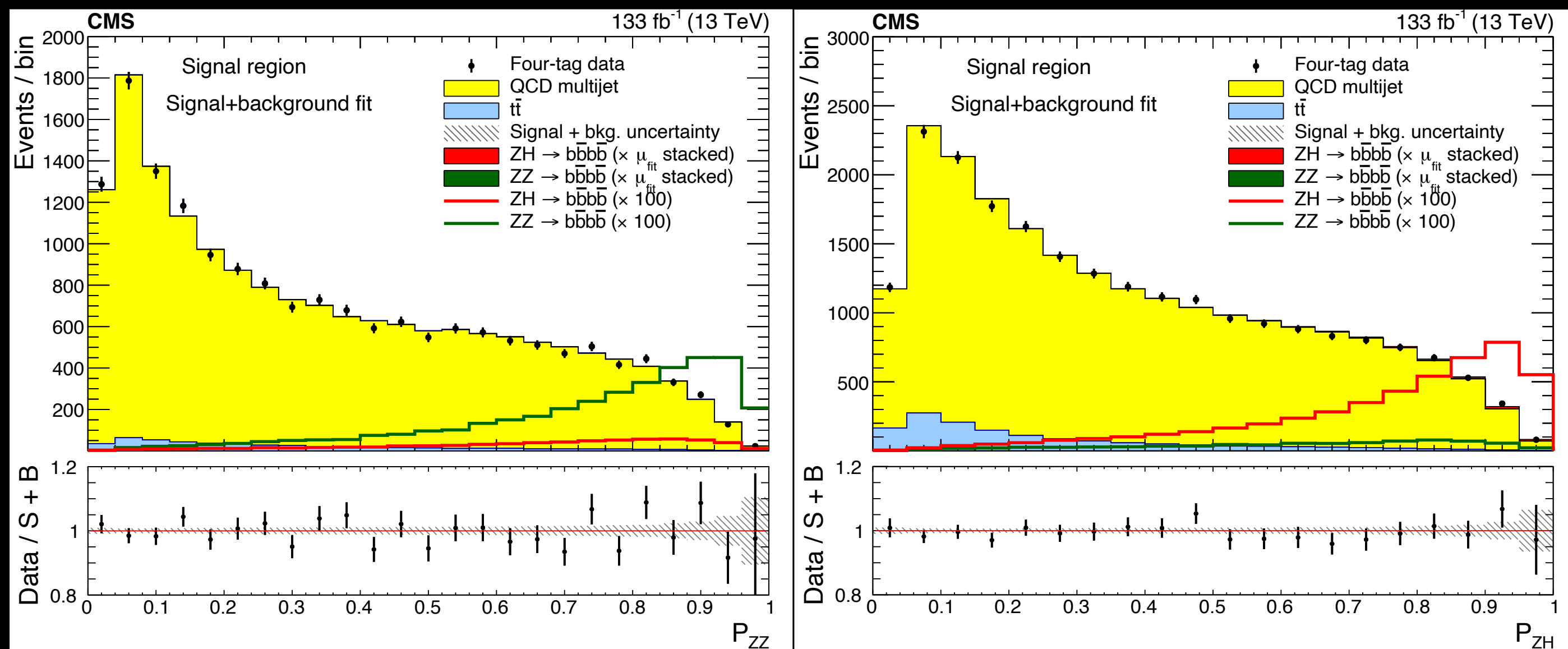
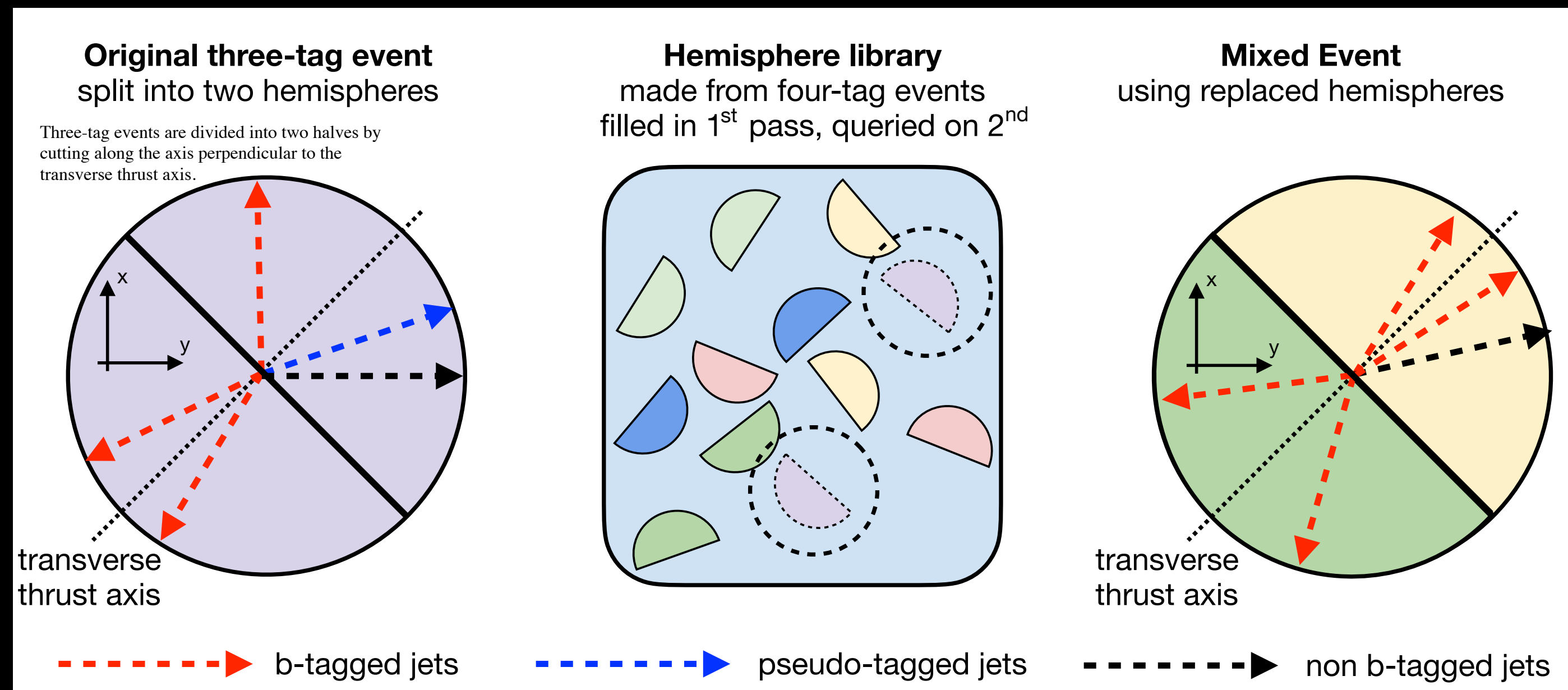
$$X_{B_1 B_2} = \sqrt{\left(\frac{m_{jj}^{\text{lead}} - m_{B_1}}{\sigma_{m_{jj}^{\text{lead}}}}\right)^2 + \left(\frac{m_{jj}^{\text{sublead}} - m_{B_2}}{\sigma_{m_{jj}^{\text{sublead}}}}\right)^2}$$

- Hierarchical combinatorial residual network trained for 4-jet diboson topology
  - Jet image formed from pixels representing jet 4-vector
  - Copies of jet pixels arranged to form 1-D image  $\rightarrow$  pairs of adjacent pixels represent the three possible jet pairings
  - Second layer  $\rightarrow$  six dijet pixels to form a three-pixel quadjet image
  - Final processing leads to event probabilities that determine event class
- Multijet background modeled similarly after reversing b-tagging criteria



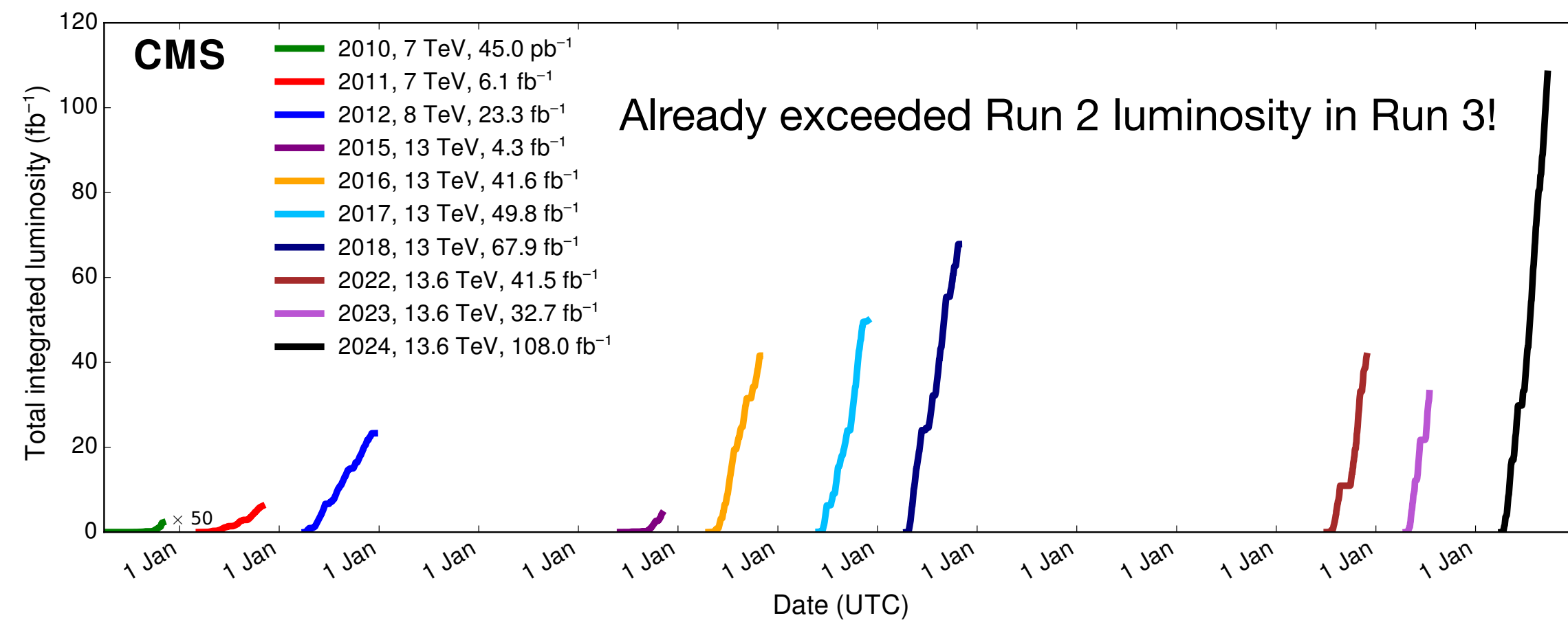
# Search for $ZZ \rightarrow b\bar{b}$ and $ZH \rightarrow b\bar{b}$

- Background modeling validated by high-event-count proxy for the four-tag background
- Synthetic data generated by using the hemisphere technique:
  - Event split into hemispheres and library generated
  - Corrections introduced to account for  $t\bar{t}$  background contamination
- Mixed models allow testing for biases in the background model that can mimic a signal
  - Check with and without unconstrained signal template
- $ZZ \rightarrow b\bar{b}$  observed with  $3.8 \sigma$  significance
- $ZH \rightarrow b\bar{b}$  observed with  $5.0 \sigma$  significance

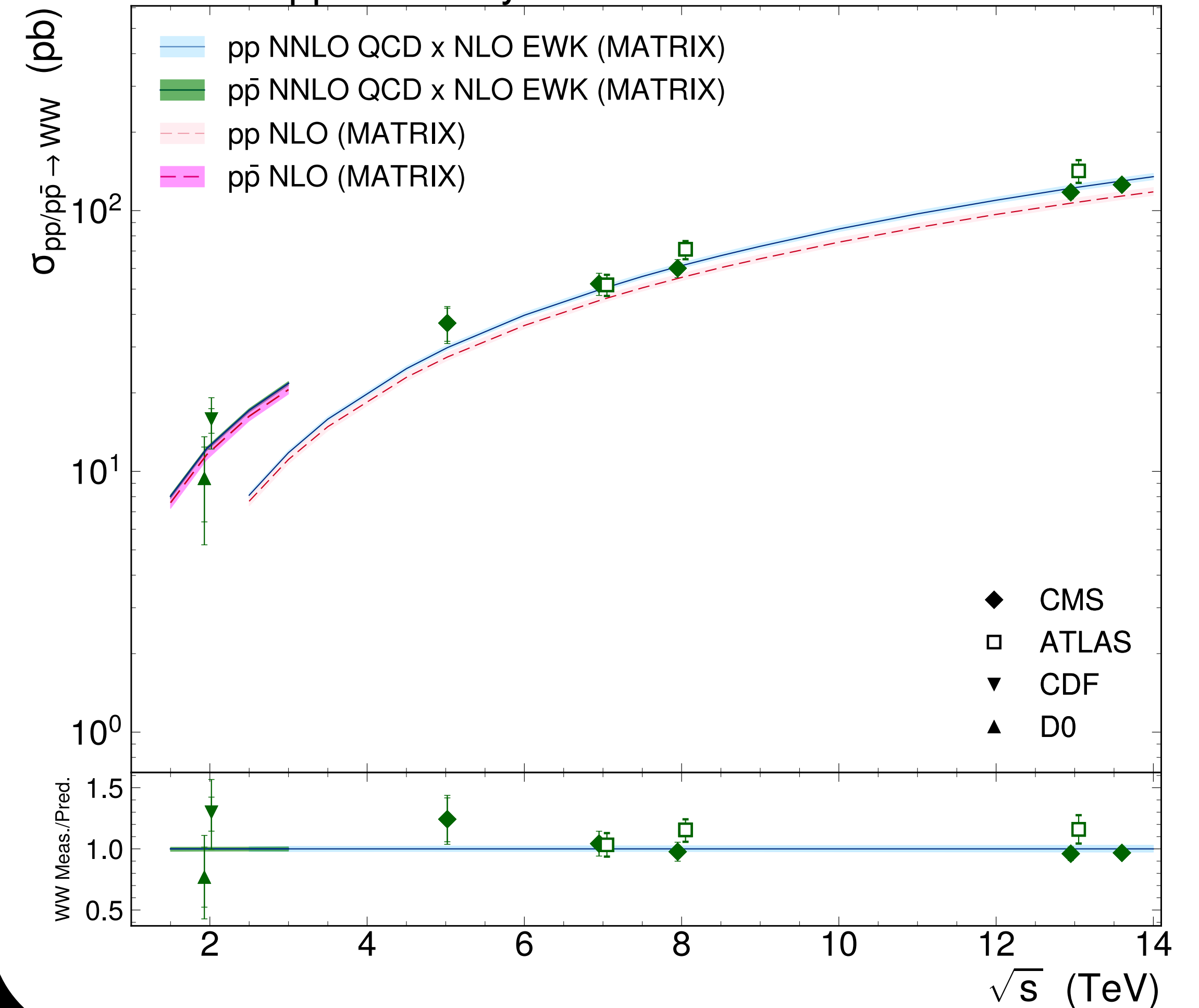


# Stress testing the Standard Model

- The full Run 2 dataset allowed us to observe rare processes predicted by the SM
- Precision tests of multiboson processes and improving background predictions for Higgs self coupling measurements were possible
- Run 3 luminosity already exceeds Run 2



CMS Supplementary



# Additional Material

# Introducing the dramatis personae

## The Standard Model of Particle Physics

### Quarks

$u$	$c$	$t$
$d$	$s$	$b$

### Leptons

$e$	$\mu$	$\tau$
$\nu_e$	$\nu_\mu$	$\nu_\tau$

### Forces

HIGGS

$\nu$

$W^\pm$	$g$
$Z^0$	$\gamma$

### Focus of my talk:

- Explore the gauge boson sector → interplay between  $W$ ,  $Z$ , and  $\gamma$
- Enables fundamental tests of the Standard Model

# Observation of $WW\gamma$ production at $\sqrt{s} = 13$ TeV

Control region	Same sign $WW\gamma$ control region	Top $\gamma$ control region
Definition	Require charged leptons of identical charge	Flip b-veto and remove cut on transverse mass of the WW-system
Target processes	Validate non-prompt lepton background modeling	Validate top-quark and non-prompt photon background modeling

# Observation of $WW\gamma$ production at $\sqrt{s} = 13$ TeV

Process	SR (0 jet)	SR ( $\geq 1$ jet)	SR (total)	SSWW $\gamma$ CR	Top $\gamma$ CR
WW $\gamma$	$122 \pm 23$	$132 \pm 27$	$254 \pm 47$	$1.0 \pm 0.2$	$12.8 \pm 2.7$
QCD V $\gamma$	$72.0 \pm 6.4$	$94.7 \pm 9.3$	$167 \pm 14$	$12.2 \pm 2.2$	$12.6 \pm 1.2$
VV	$15.1 \pm 1.4$	$21.6 \pm 2.4$	$36.7 \pm 3.5$	$24.9 \pm 1.7$	$2.0 \pm 0.3$
Top	$56.6 \pm 6.5$	$271 \pm 26$	$328 \pm 32$	$2.4 \pm 0.6$	$2434 \pm 85$
Nonprompt $\ell$	$45.7 \pm 4.0$	$77.2 \pm 6.5$	$122.9 \pm 9.7$	$197 \pm 14$	$40 \pm 11$
Nonprompt $\gamma$	$109.1 \pm 9.0$	$301 \pm 24$	$410 \pm 32$	$19.9 \pm 1.6$	$793 \pm 62$
Total	$420 \pm 20$	$898 \pm 29$	$1318 \pm 43$	$257 \pm 14$	$3294 \pm 57$
Data	414	916	1330	259	3287

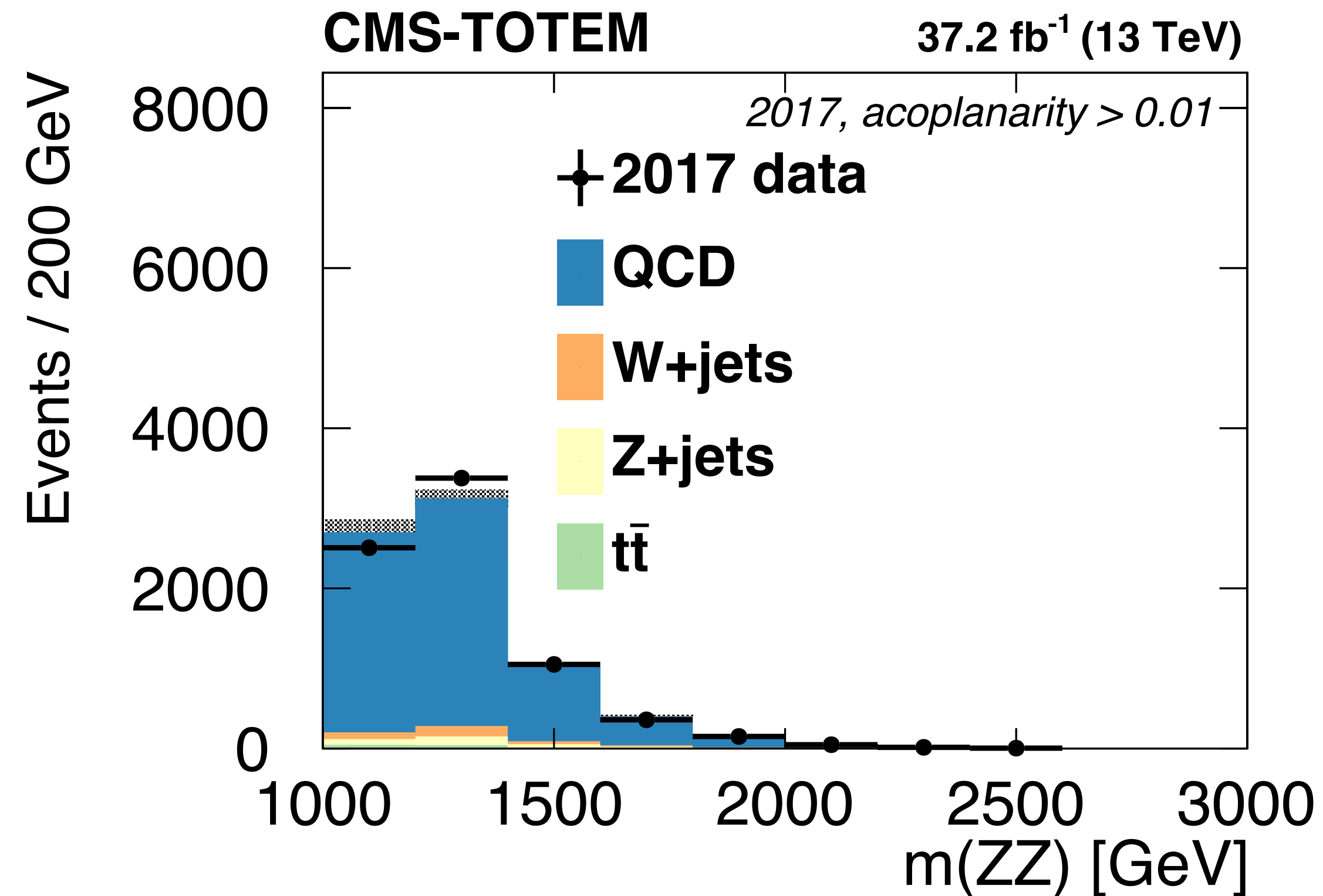
Background  
dominated

- Simultaneous extraction of signal and control regions



# Search for exclusive $\gamma\gamma \rightarrow WW$ and $\gamma\gamma \rightarrow ZZ$ production in final states with jets and forward protons

Number of events	region	$N_{\text{evt}}$ (2016)	$N_{\text{evt}}$ (2017)	$N_{\text{evt}}$ (2018)
Anti-acoplanarity sideband	$\delta$	$1.5 \pm 1.1$	$1.6 \pm 0.8$	$14.2 \pm 3.0$
Anti-pruned mass sideband	$\delta$	$0.4 \pm 0.2$	$0.9 \pm 0.2$	$9.9 \pm 0.9$
Event mixing	$\delta$	$0.5 (< 2.1)$	$1.5 (< 3.6)$	$11.6 \pm 9.4$
Expected signal ( $a_0^Z / \Lambda^2 = 1 \times 10^{-5} \text{ GeV}^{-2}$ )	$\delta$	1.3	1.4	9.0
Anti-acoplanarity sideband	$o$	$1.5 \pm 1.1$	$3.7 \pm 1.5$	$37.4 \pm 5.6$
Anti-pruned mass sideband	$o$	$2.1 \pm 0.8$	$5.4 \pm 1.3$	$41.7 \pm 3.1$
Event mixing	$o$	$2.0 \pm 1.8$	$6.3 \pm 5.1$	$42 \pm 16$
Expected signal ( $a_0^Z / \Lambda^2 = 1 \times 10^{-5} \text{ GeV}^{-2}$ )	$o$	1.0	1.6	12.8





# Effective Field Theory interpretation in $\gamma\gamma \rightarrow \tau\tau$

$$\mathcal{L}_{BSM} = \frac{C_{\tau B}}{\Lambda^2} \bar{L}_L \sigma^{\mu\nu} \tau_R H B_{\mu\nu} + \frac{C_{\tau W}}{\Lambda^2} \bar{L}_L \sigma^{\mu\nu} \tau_R \sigma^i H W_{\mu\nu}^i + h.c.$$

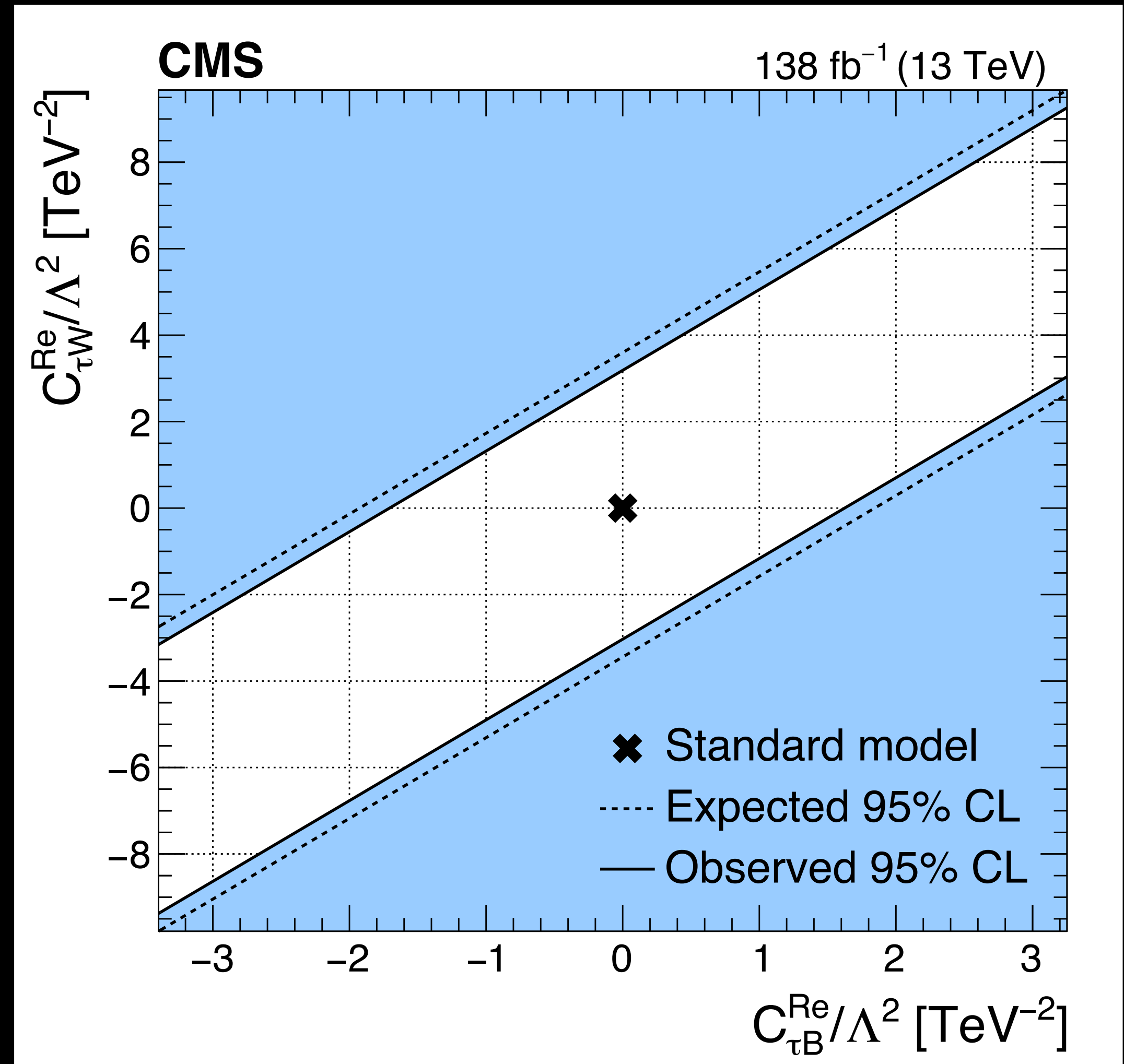
$\tau\tau\gamma$  vertex parametrized as:

$$V_{\tau\tau\gamma} = ie\gamma^\mu = \frac{v\sqrt{2}}{\Lambda^2} [\text{Re}[C_{\tau\gamma}] + \text{Im}[C_{\tau\gamma}]i\gamma_5] \sigma^{\mu\nu} q_\nu$$

$$C_{\tau\gamma} = (\cos\theta_W C_{\tau B} - \sin\theta_W C_{\tau W})$$

BSM contribution to  $a_\tau$ :

$$\delta a_\tau = \frac{2m_\tau \sqrt{2}v}{e \Lambda^2} \text{Re}[C_{\tau\gamma}]$$



# Measurement of the electroweak production of $W\gamma$ with two jets

This paper presents a measurement of the EW  $W\gamma jj$  production at  $\sqrt{s} = 13$  TeV based on the complete Run 2 data collected during 2016–2018, superseding the previous CMS result [4]. A complete set of tabulated results of this analysis is available in the HEPData database [5]. In addition to increased integrated luminosity, our new results include: (i) an updated fiducial region requiring jets with  $p_T > 50$  GeV; (ii) the removal of the missing transverse momentum requirement from the fiducial region definition; (iii) the treatment of the interference term between the EW- and quantum chromodynamics (QCD) induced processes as a background component; (iv) and the treatment of the out-of-fiducial signal contribution as a background component.

## 9 Fiducial cross section measurement

The fiducial cross section measurement for the EW  $W\gamma$  production at 13 TeV is extracted with the same 2D  $m_{jj}-m_{\ell\gamma}$  binning used for the signal significance. The fiducial region is defined based on the particle-level (for leptons, photons, jets) quantities: one lepton  $p_T^\ell > 35$  GeV and

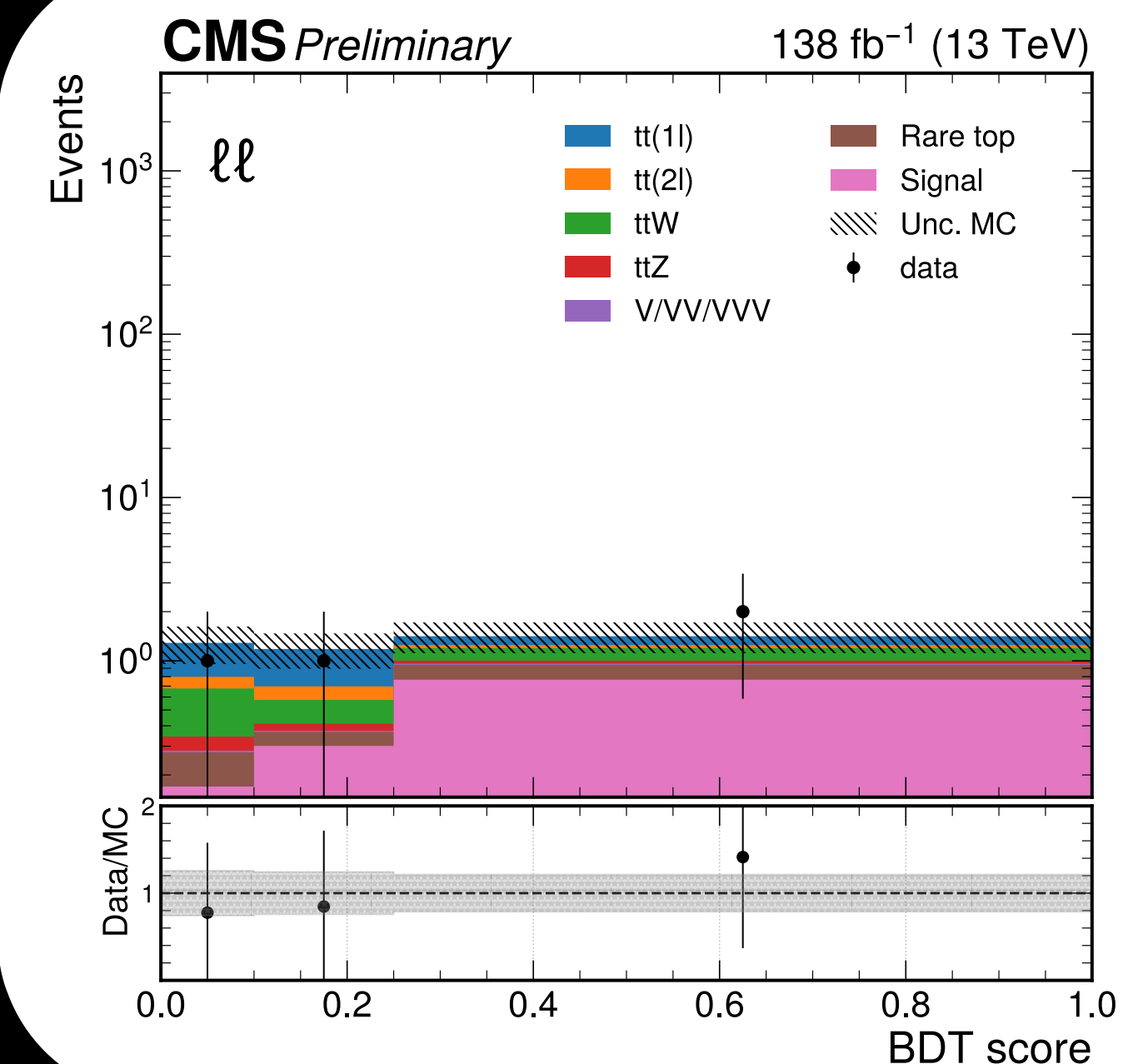
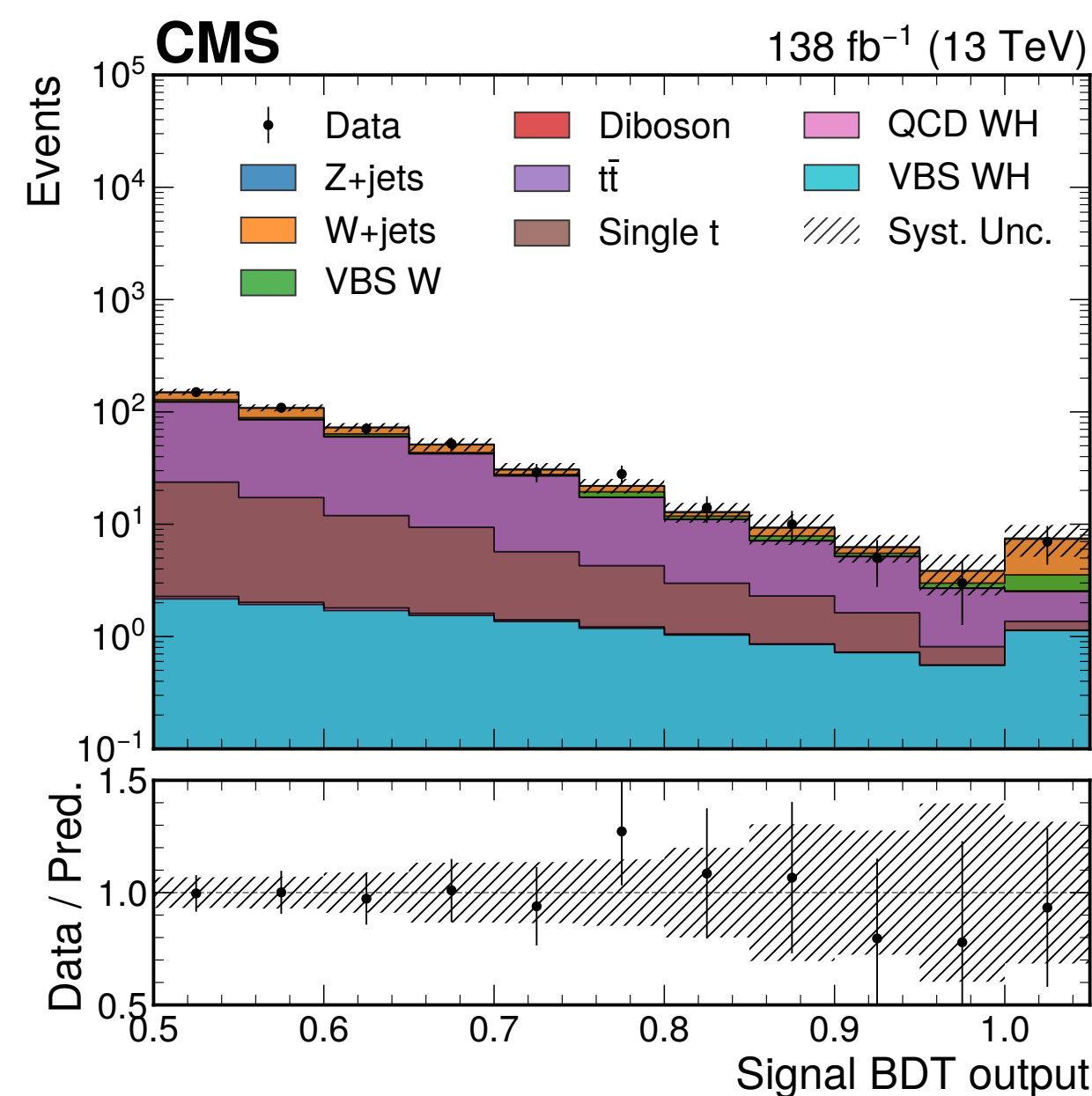
$|\eta_\ell| < 2.4$ ,  $p_T^{\text{miss}} > 30$  GeV,  $p_T^\gamma > 25$  GeV,  $|\eta_\gamma| < 1.444$  or  $1.566 < |\eta_\gamma| < 2.5$ ,  $\Delta R_{\ell\gamma} > 0.5$ ,  $m_T^W > 30$  GeV, and two jets with  $p_T^{j1(2)} > 50$  GeV,  $|\eta_j| < 4.7$ ,  $m_{jj} > 500$  GeV,  $\Delta R_{jj} > 0.5$ ,  $\Delta R_{j\ell} > 0.5$ ,  $\Delta R_{j\gamma} > 0.5$ , and  $|\Delta\eta_{jj}| > 2.5$ . The leptons are reconstructed at the particle level with fully recovered final-state radiation. The acceptance is defined as the fraction of the signal events passing the fiducial region selection, and is estimated using MG5. The theoretical uncertainty in the extrapolation between the fiducial and SR is negligible ( $< 1\%$ ). We define the cross section as  $\sigma^{\text{fid}} = \sigma_g \hat{\mu} \alpha_{\text{gf}}$ , where the cross section for the signal events is  $\sigma_g = 0.776$  pb calculated with MG5 at LO in QCD [12], the observed signal strength parameter  $\hat{\mu} = 0.88_{-0.18}^{+0.19}$ , and the

# Search for $HHWW$ couplings in the VBS production of $W^\pm W^\pm H$ with $H \rightarrow b\bar{b}$ decays

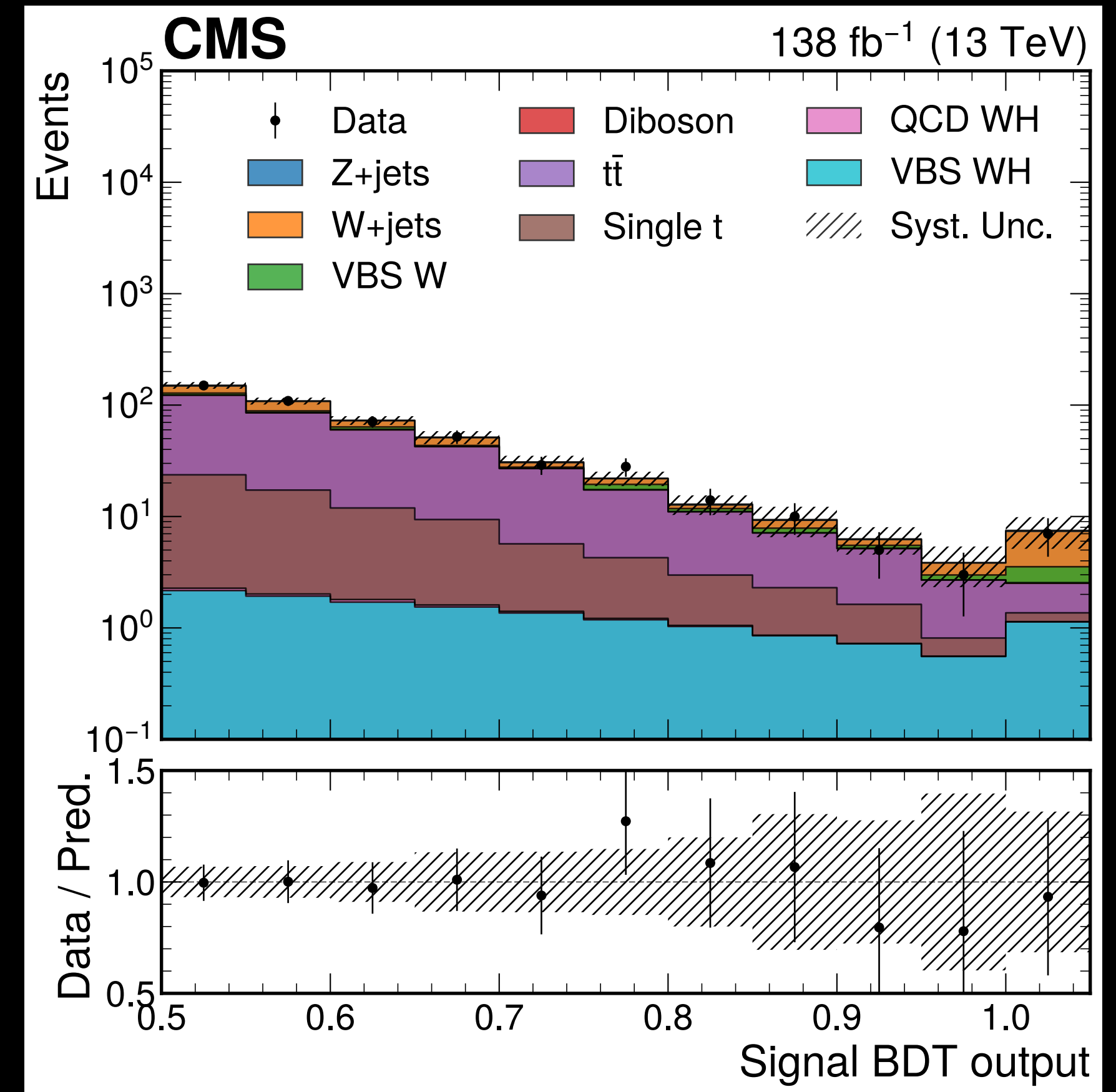
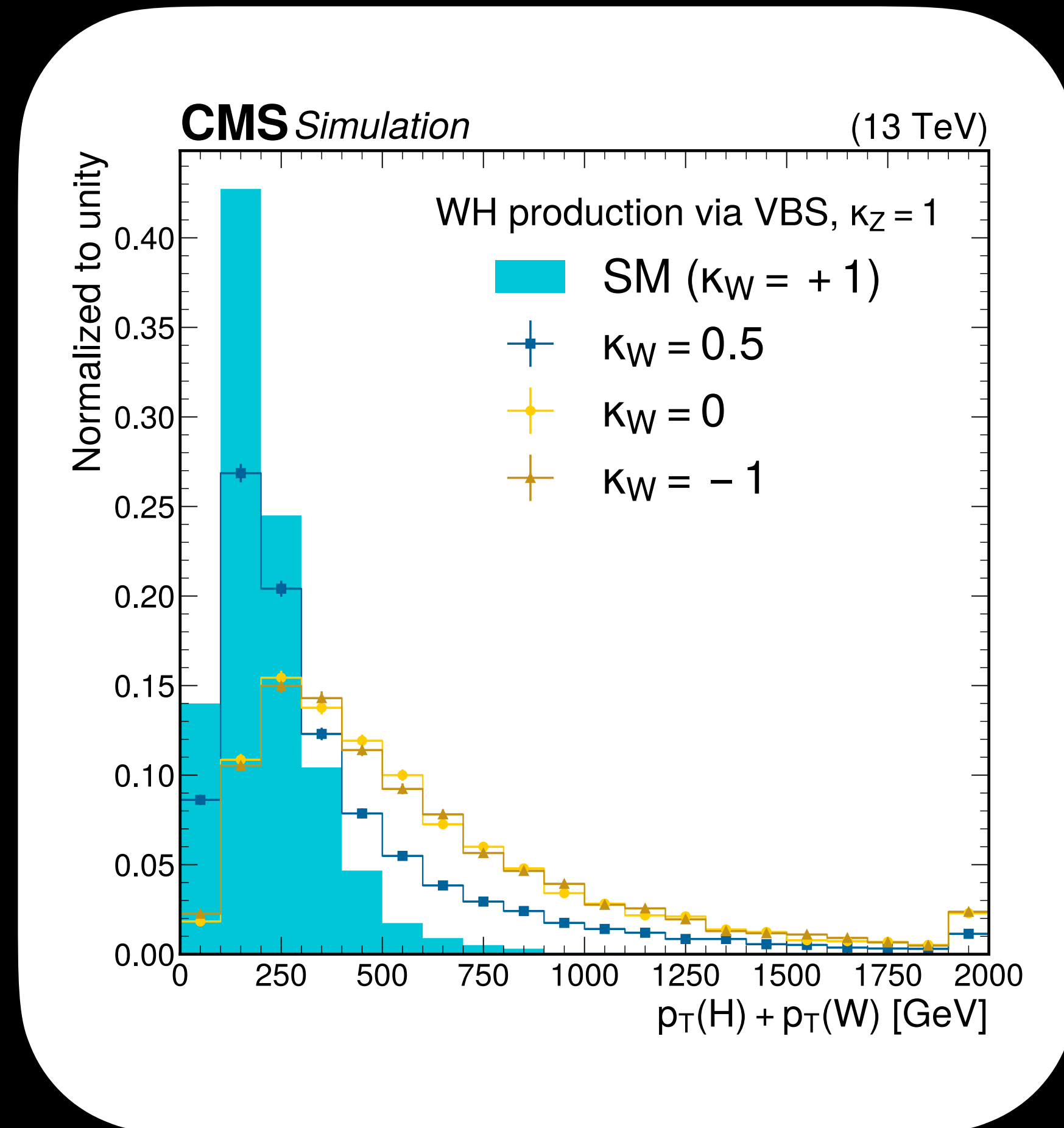
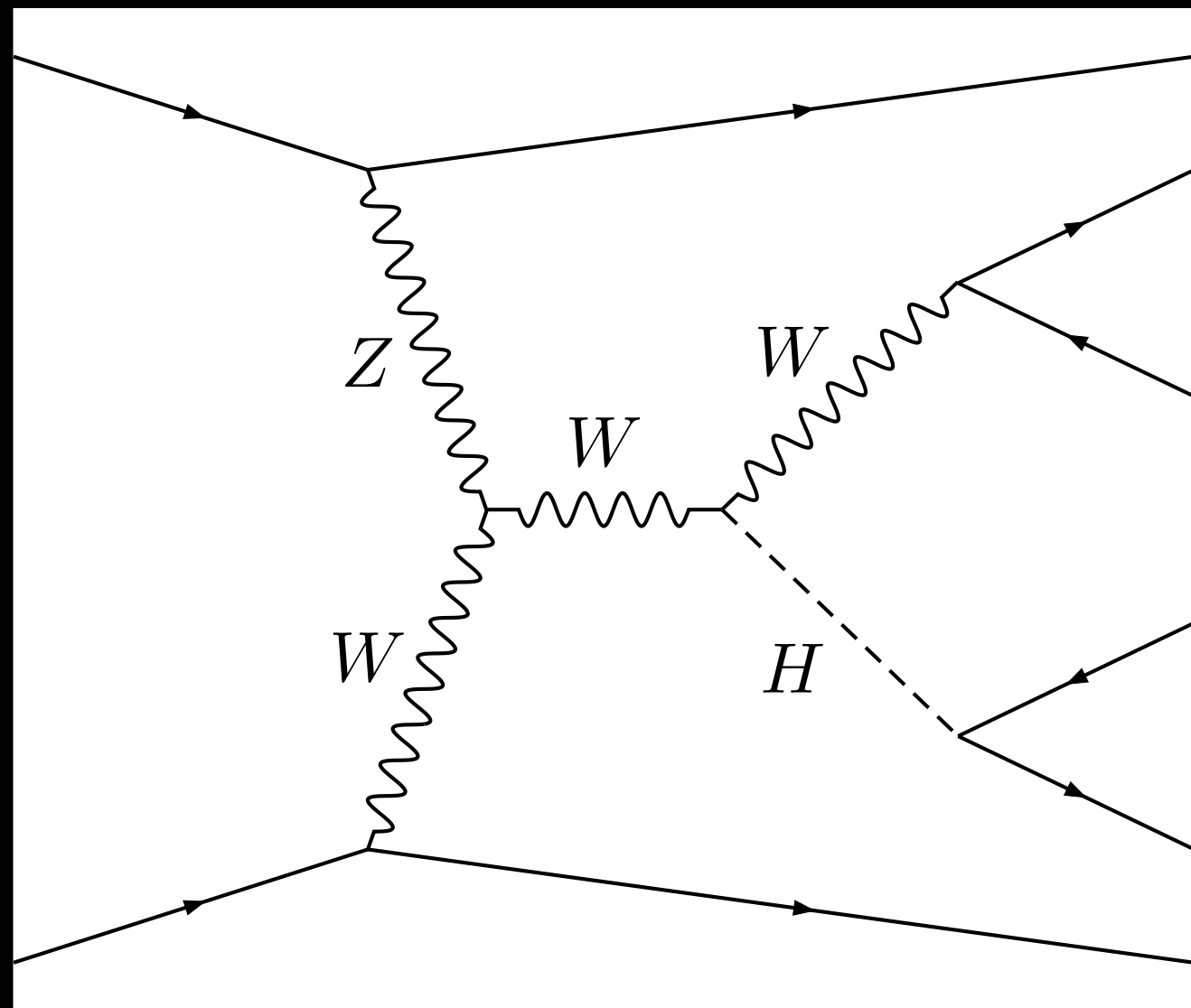
- Extraction of sign of W and Z coupling to Higgs performed with single W and Higgs in final state
- Final state: 2 same-signed leptons (one  $\tau$  lepton permitted), a boosted jet, non-VBS contribution removed by requiring  $m_{jj} > 100$  GeV
- Boosted decision tree trained against major backgrounds
- Low signal yields  $\rightarrow$  need more data to set tight constraints on couplings

## $W^\pm W^\pm H$

Shorthand	Description
$\eta_J$	$\eta$ of the leading merged jet
$p_{T,J}$	$p_T$ of the leading merged jet
$p_{T,jj}$	$p_T$ of the VBS-jet system
$P_{j_0}$	magnitude of the three-momentum of the leading VBS jet
$P_{j_1}$	magnitude of the three-momentum of the subleading VBS jet
$M_{\ell\ell}$	invariant mass of the SS dilepton system
$p_{T,\ell_0}$	$p_T$ of the leading lepton
$p_{T,\ell_1}$	$p_T$ of the subleading lepton
$E_T^{\text{miss}}$	missing transverse energy
$L_T$	scalar sum of $p_{T,\ell_0}, p_{T,\ell_1}$ , and $E_T^{\text{miss}}$
$S_T$	scalar sum of $p_{T,J}$ and $L_T$



# Study of $WH$ production and extraction of relative sign of the $W$ and $Z$ couplings



# Search for $ZZ \rightarrow b\bar{b}$ and $ZH \rightarrow b\bar{b}$

Source	ZZ	ZH
Statistical uncertainty	75	77
Total systematic uncertainty	67	64
Background model	61	56
(Variance)	(46)	(46)
(Extrapolation)	(40)	(33)
b tagging	9	17
Jet energy scale and resolution	9	5
Others	24	24

	ZZ	ZH
Signal strength expected (stat. only)	$1.0^{+1.9}_{-1.7}$ ( $1.0^{+1.4}_{-1.3}$ )	$1.0^{+1.5}_{-1.4}$ ( $1.0^{+1.1}_{-1.1}$ )
Signal strength observed	$0.0^{+2.0}_{-1.7}$	$2.2^{+0.9}_{-0.8}$
Expected upper limit at 95% CL (stat. only)	3.8 (2.8)	2.9 (2.3)
Observed upper limit at 95% CL	3.8	5.0

# Search for $ZZ \rightarrow b\bar{b}$ and $ZH \rightarrow b\bar{b}$

The mixing algorithm is based on the technique developed in a previous CMS  $HH \rightarrow 4b$  analysis [20]. The first step involves creating a collection of hemispheres (hemisphere library) from events in the four-tag data set. Each event is split into two hemispheres using the plane orthogonal to the transverse thrust axis [20], which is chosen based on the assumption that it is a good proxy for the initial gluon directions in the underlying scattering process. Jets on one side of the plane are assigned to one hemisphere, those on the other side are assigned to the other hemisphere. Four variables are computed using the sum of the four-vectors of all the jets in that hemisphere: the invariant mass, the longitudinal momentum, and the transverse momentum perpendicular and parallel to the transverse thrust axis. The jet and b jet multiplicities are also computed for each hemisphere. The library is created with events that pass the jet kinematic requirements but before the dijet invariant mass requirement given in Eq. (1) is applied.