

The Automation of **SMEFT-Assisted** constraints on **UV-complete models**

Multi-Boson Interactions 2024

Toulouse, France

27 September 2024

Alejo N. Rossia

On behalf of the  **SMEFiT** Collaboration

Department of Physics and Astronomy

The University of Manchester

Based on:

[2309.04523] JHEP 01 (2024) 179 (w/ J. ter Hoeve, G. Magni, J. Rojo, and E. Vryonidou)

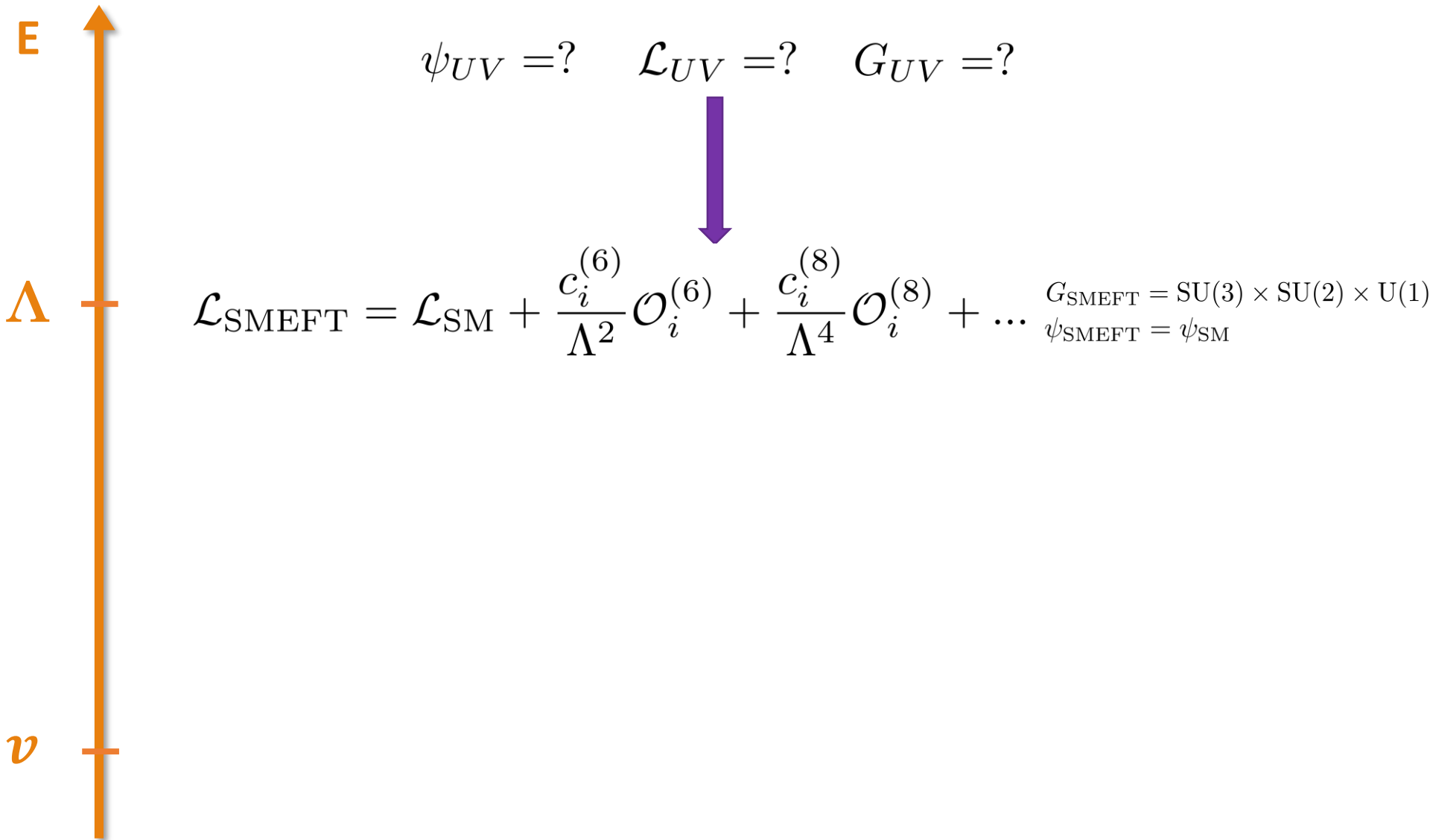
[2404.12809] JHEP 09 (2024) 091 (w/ E. Celada, T. Giani, J. ter Hoeve, L. Mantani, J. Rojo, M. Thomas and E. Vryonidou)

The Standard Model EFT (SMEFT)

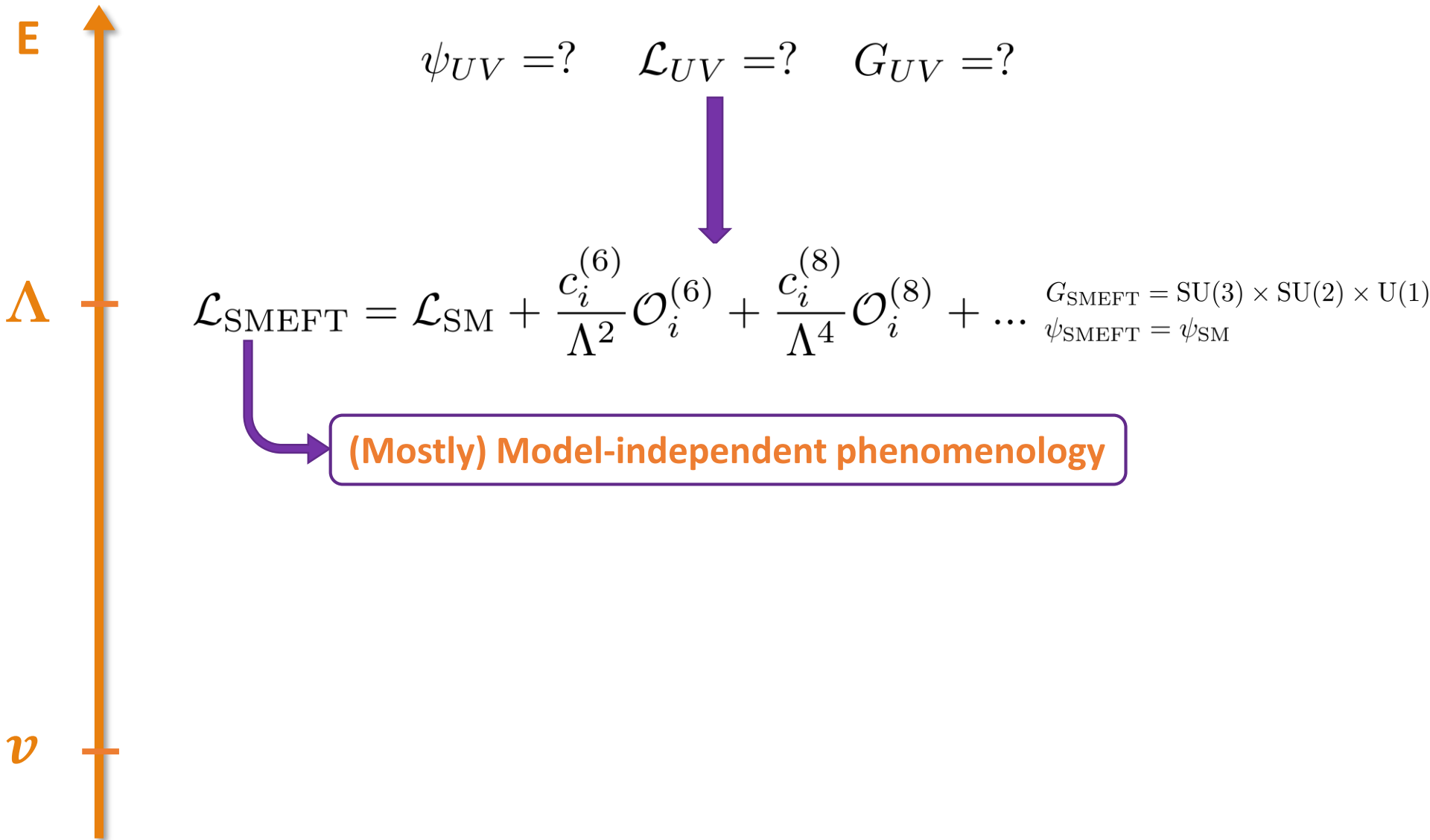


$$\psi_{UV} =? \quad \mathcal{L}_{UV} =? \quad G_{UV} =?$$

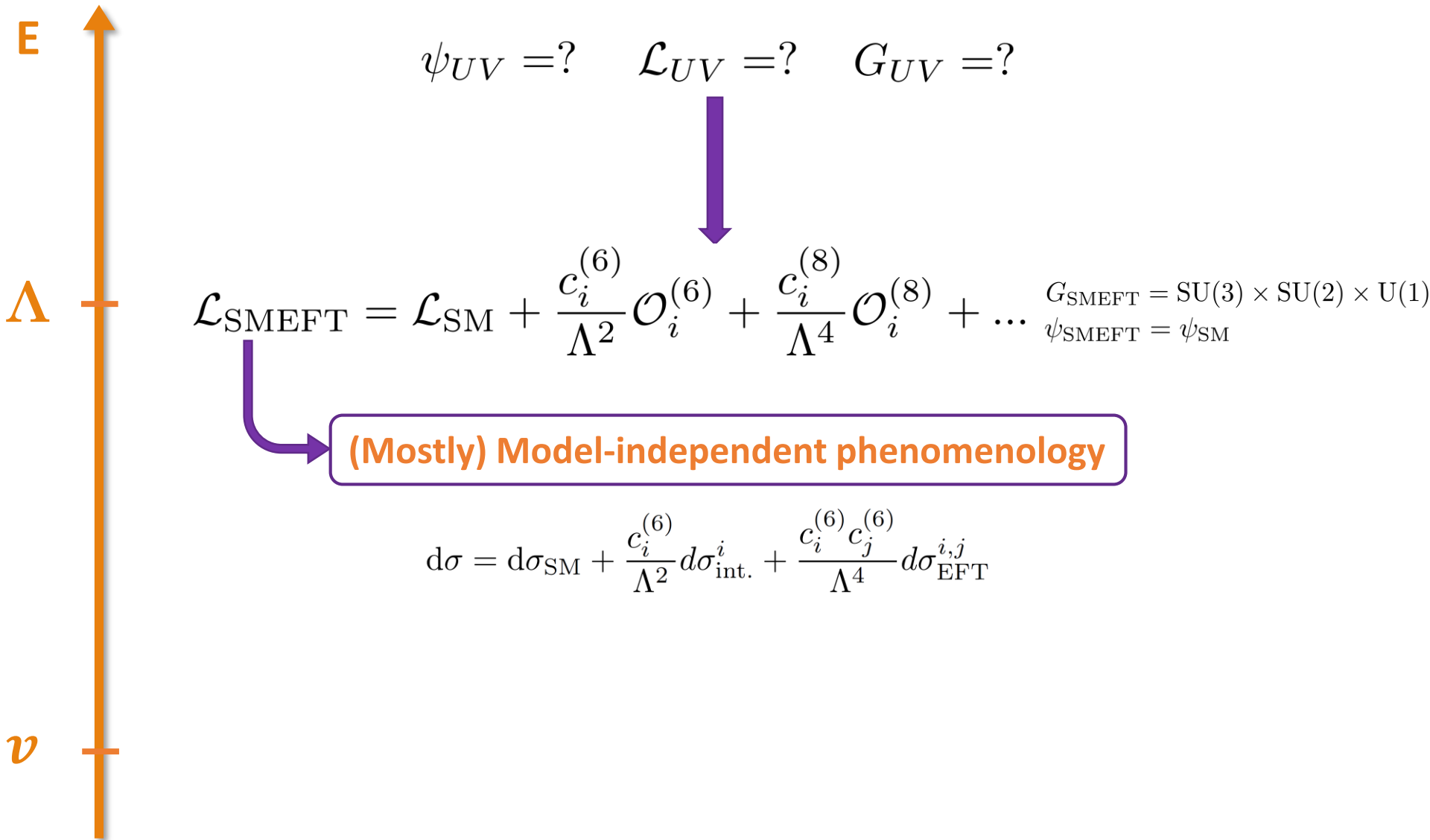
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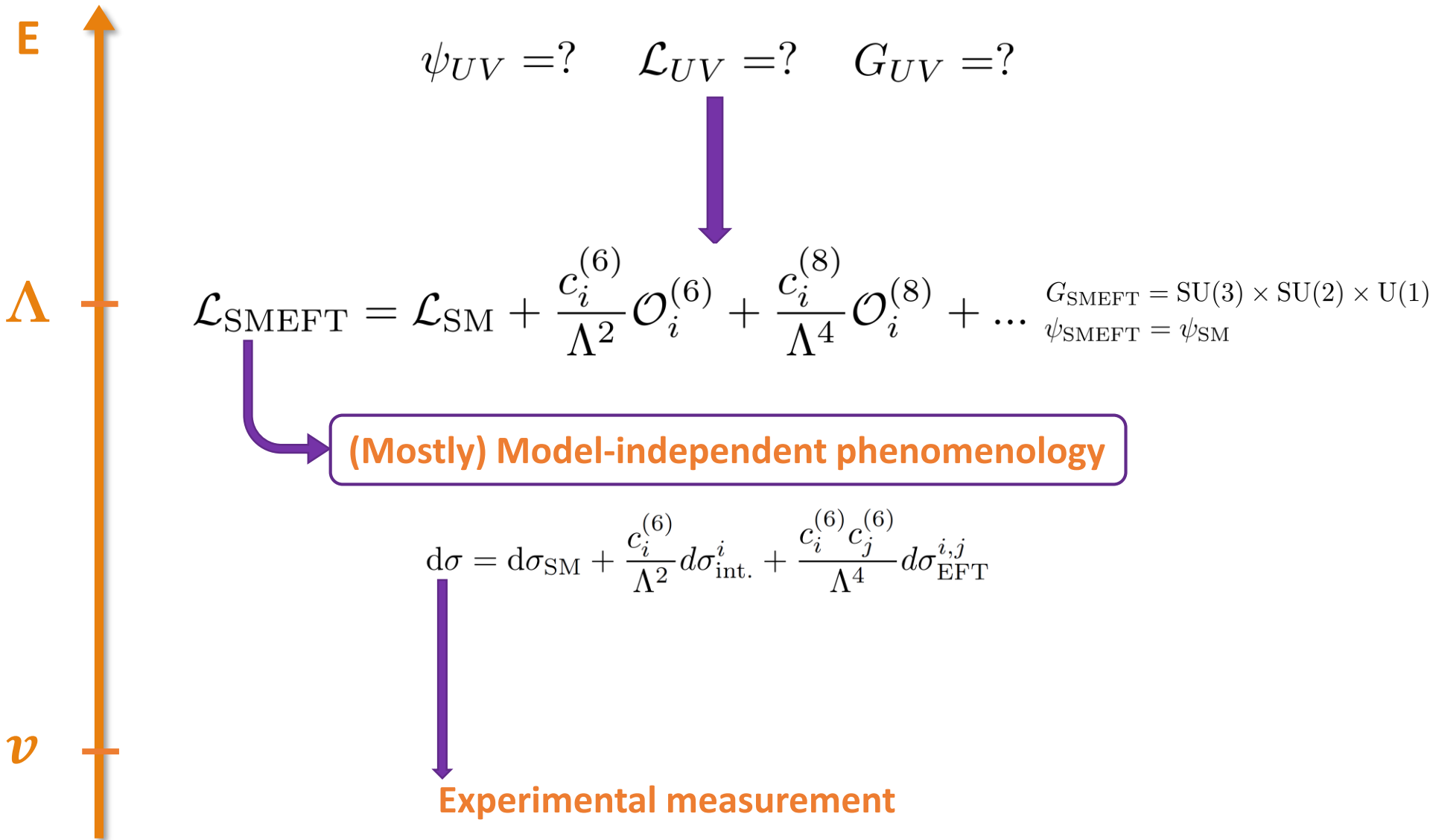
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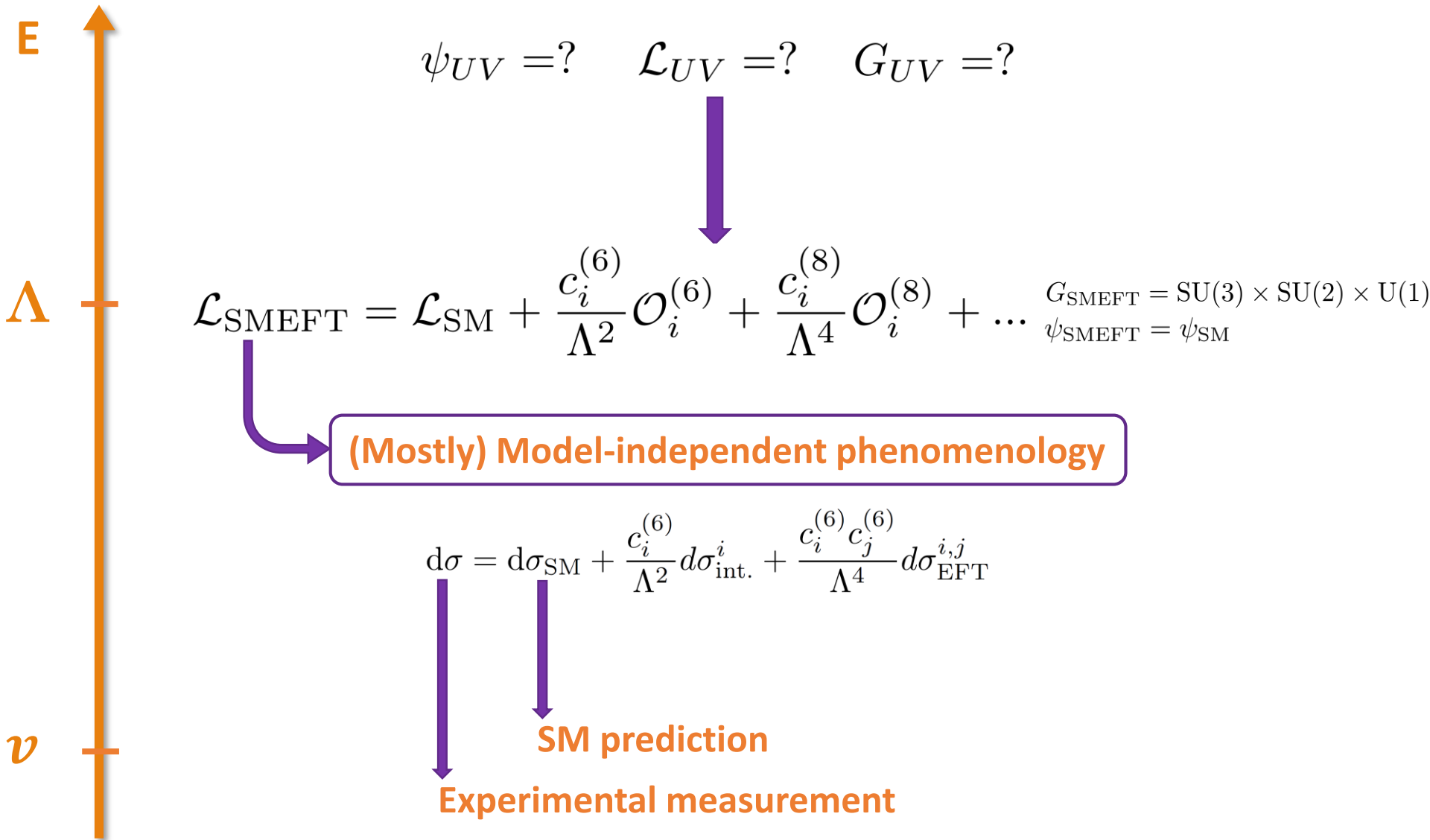
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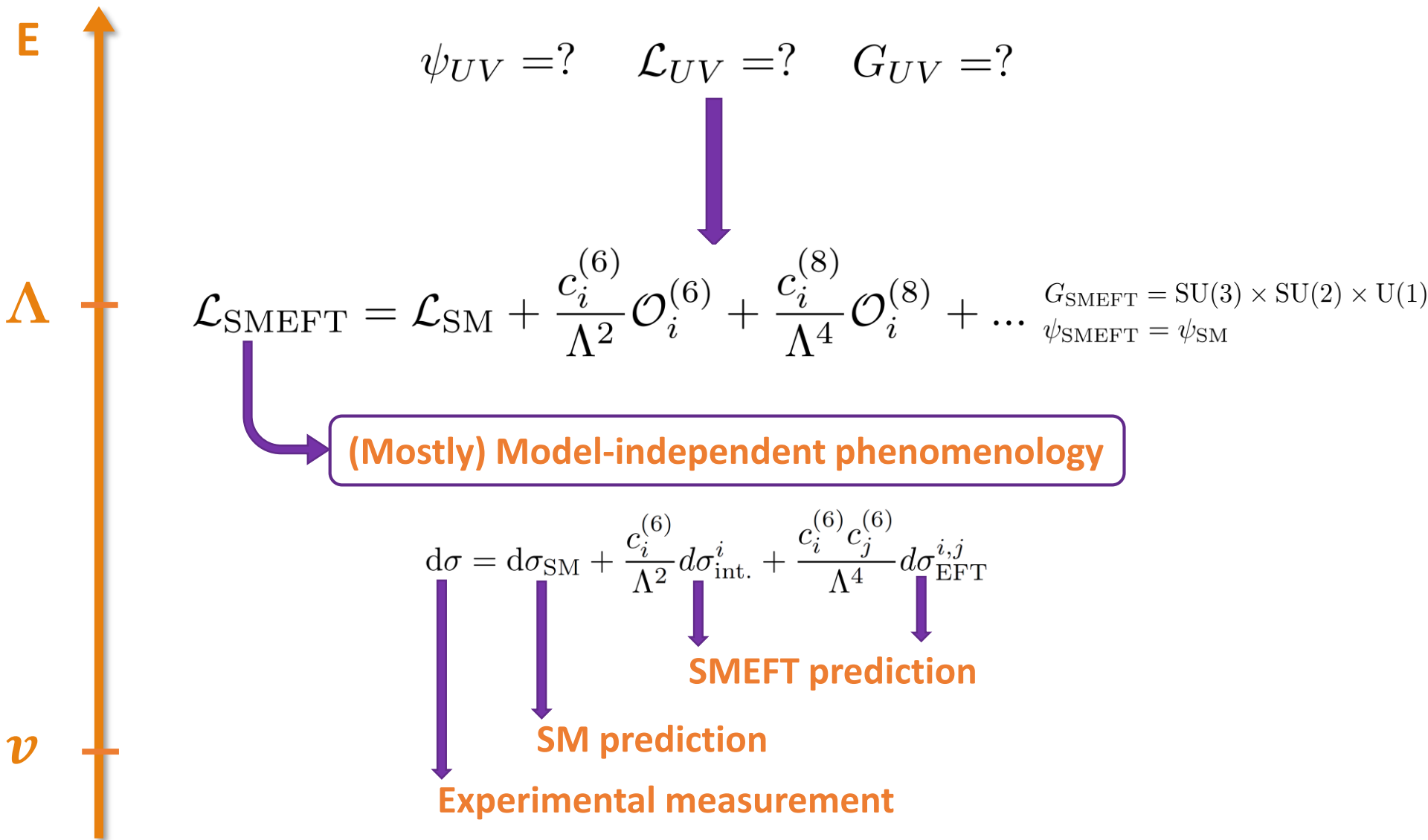
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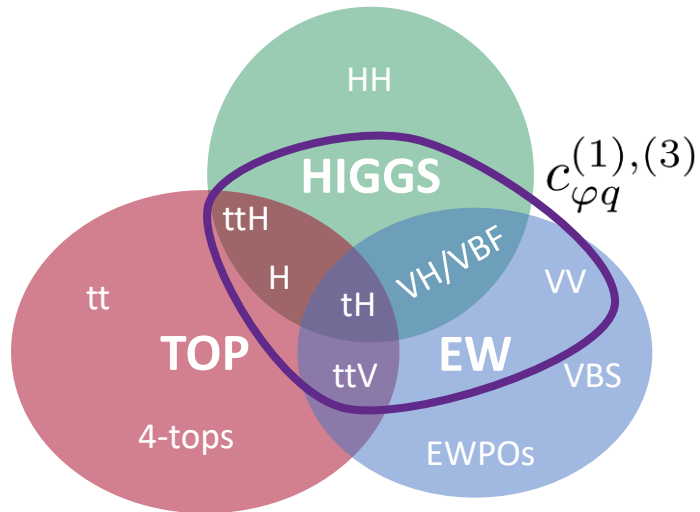


Global fits

Correlations, correlations everywhere...

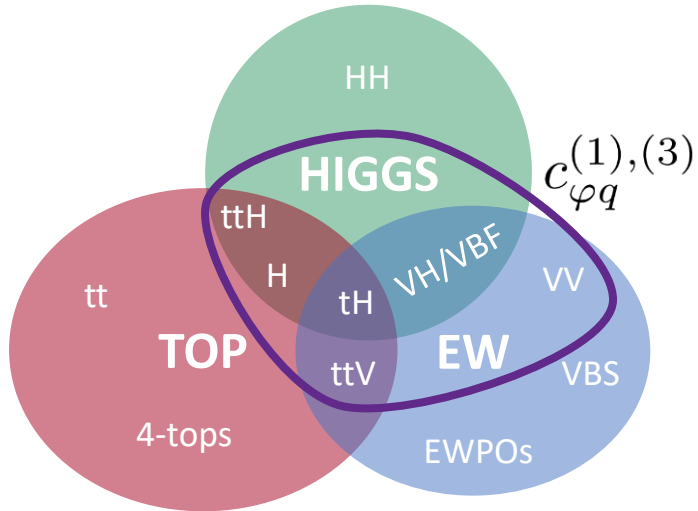
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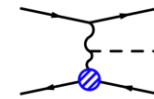
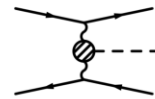


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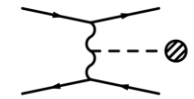
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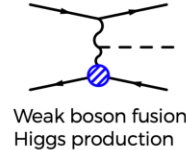
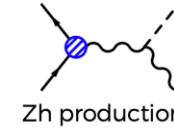
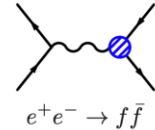
One observable can be influenced by many operators



Higgs decay



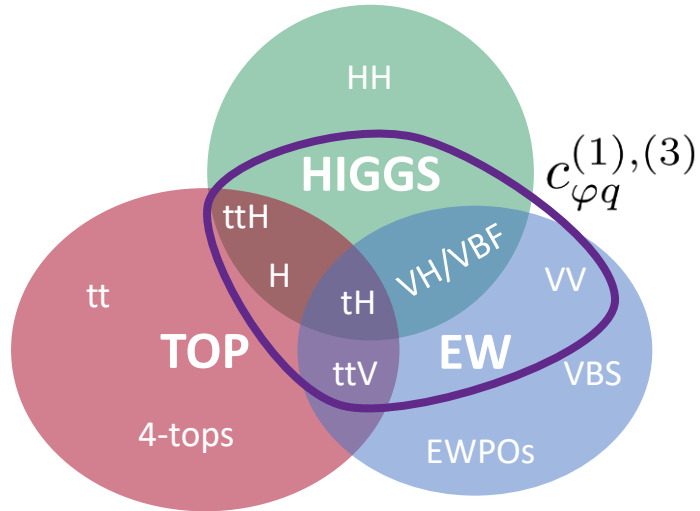
One operator can contribute to many different observables



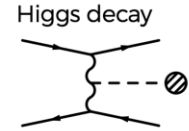
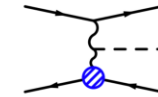
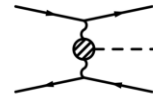
[A. Biekötter's seminar]

Global fits

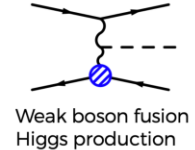
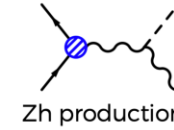
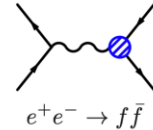
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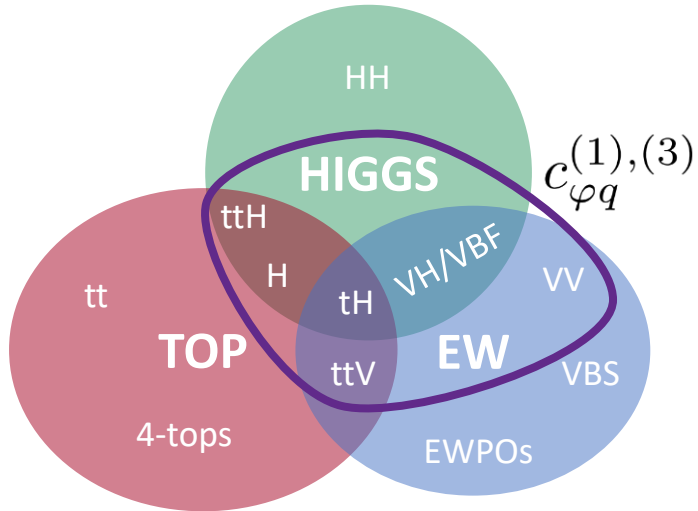


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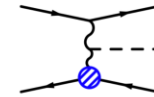
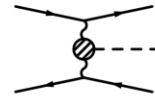
The choices in the fitter's way

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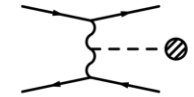
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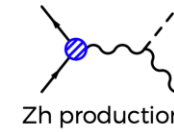
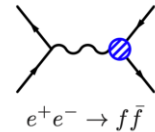
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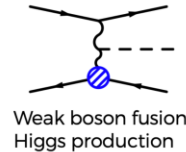
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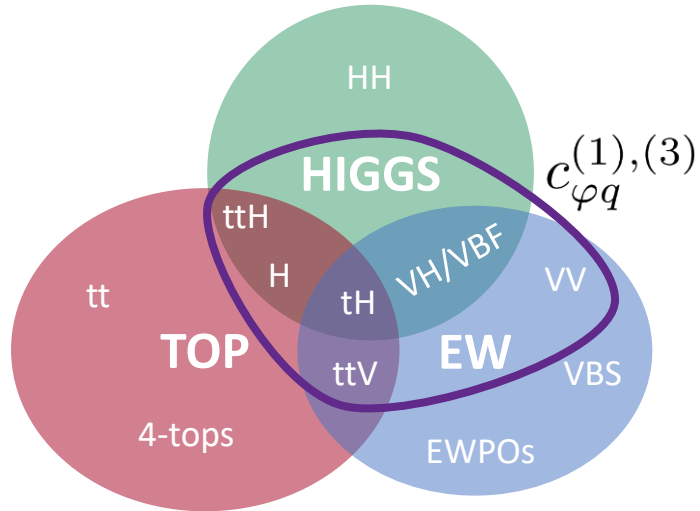


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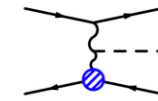
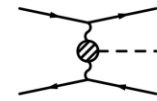
- Dataset

Global fits

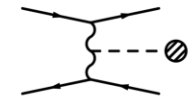
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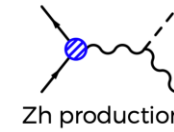
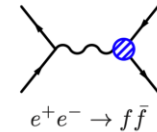
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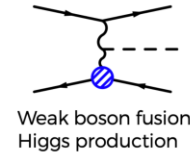
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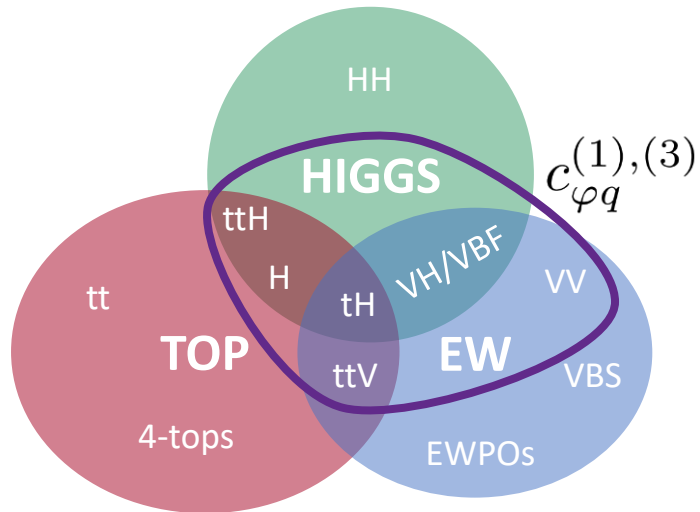


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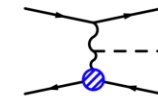
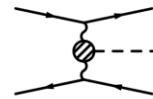
- Dataset
- Likelihoods

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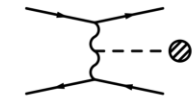
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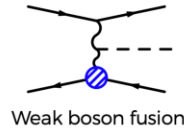
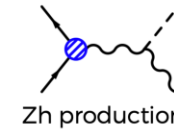
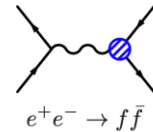


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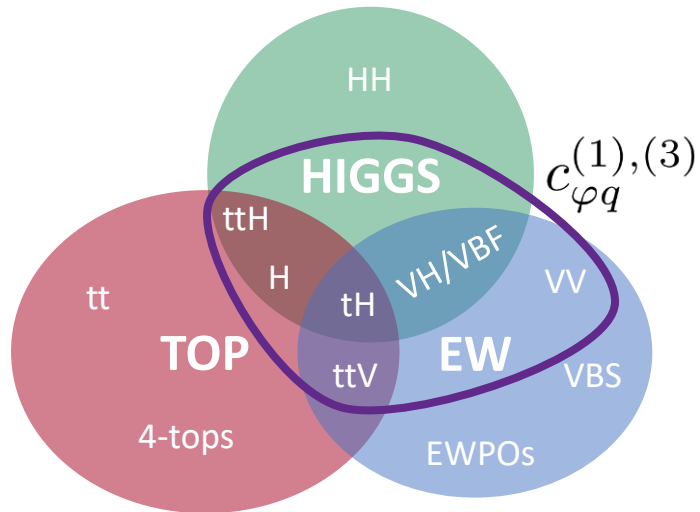
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Higgs production

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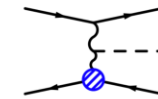
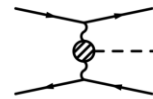
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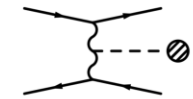
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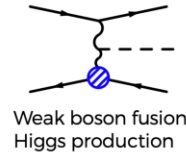
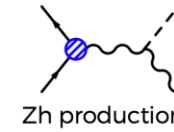
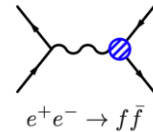


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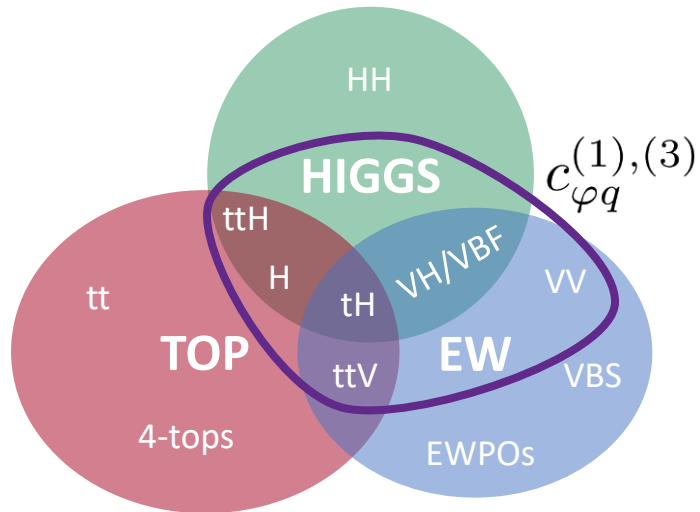


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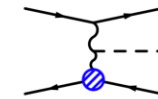
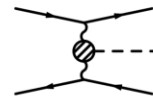
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Global fits

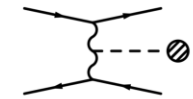
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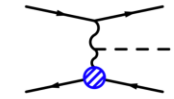
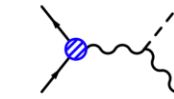
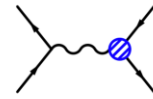


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$e^+e^- \rightarrow f\bar{f}$

Zh production

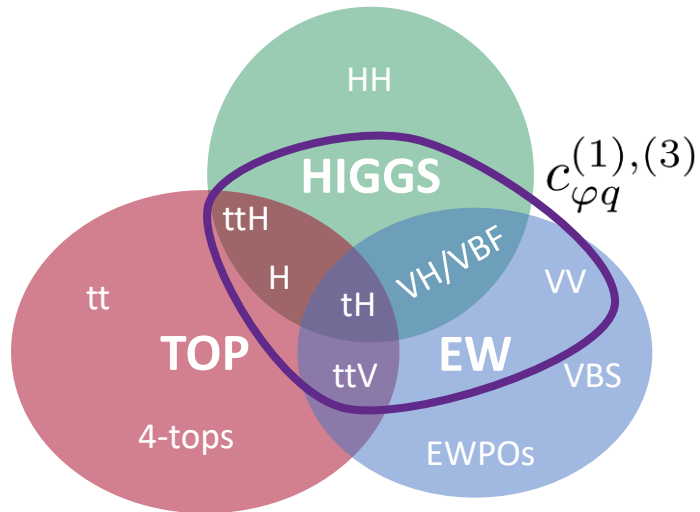
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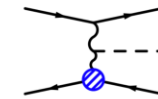
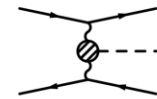
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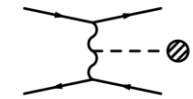
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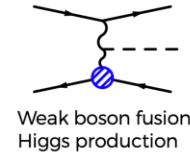
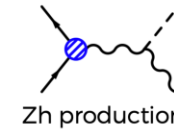
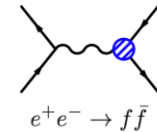
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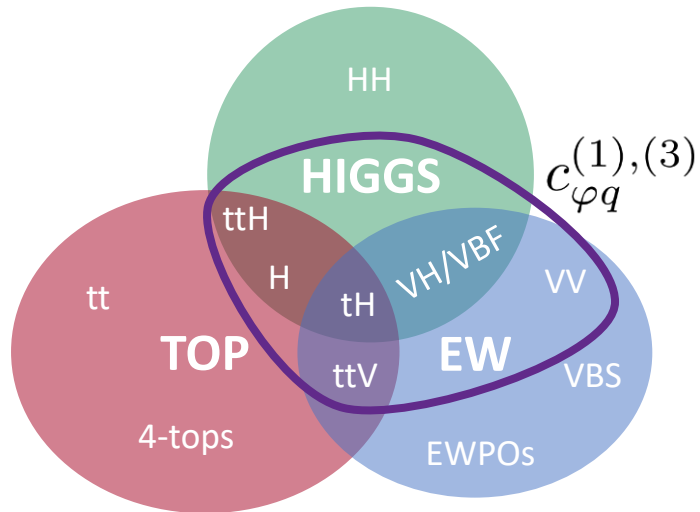
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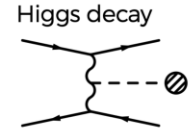
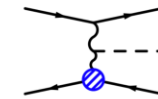
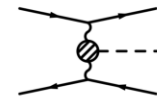
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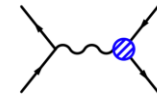


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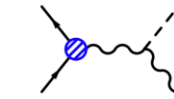


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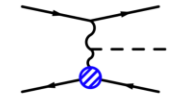
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Zh production



Weak boson fusion Higgs production

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The choices in the fitter's way

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Fitmaker

[2012. 02779]



[1910.14012]



[2302.06660]



[2208. 08454]

...



A Python software for global interpretation of particle physics data in SMEFT





A Python software for global interpretation of particle physics data in SMEFT

THEORY

SM: (N)NLO QCD + NLO EW
SMEFT: NLO QCD, quadratic in WCs
50 WCs!



SMEFiT



A Python software for global interpretation of particle physics data in SMEFT

THEORY

SM: (N)NLO QCD + NLO EW
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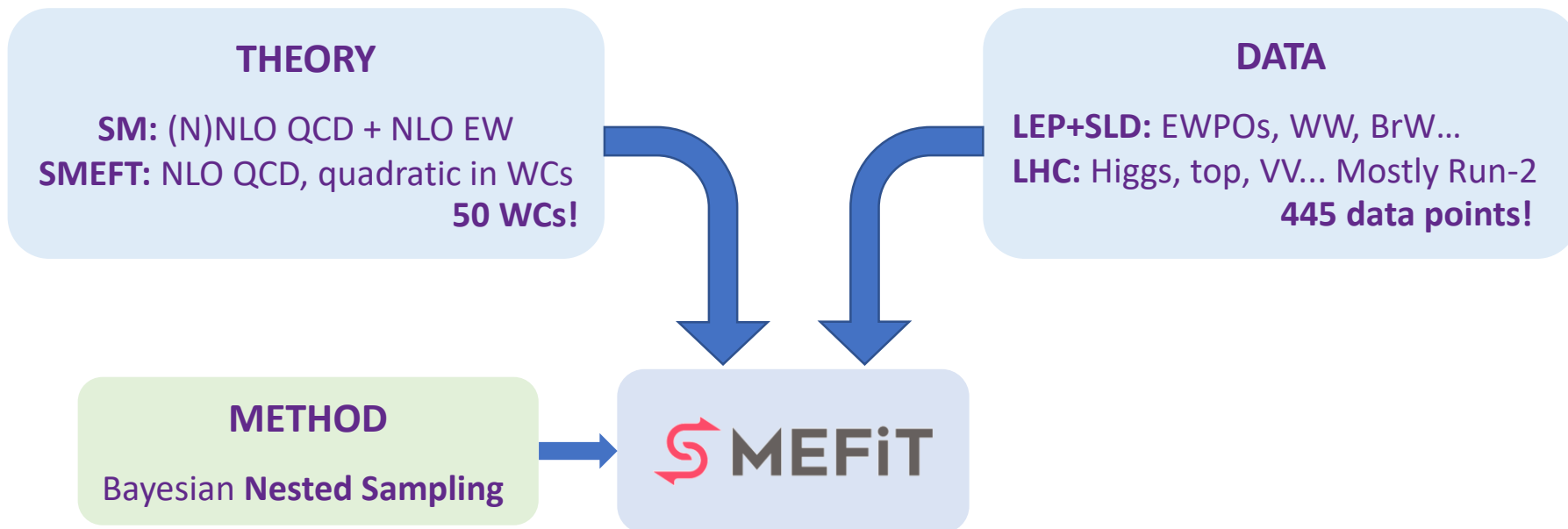
DATA

LEP+SLD: EWPOs, WW, BrW...
LHC: Higgs, top, VV... Mostly Run-2
445 data points!



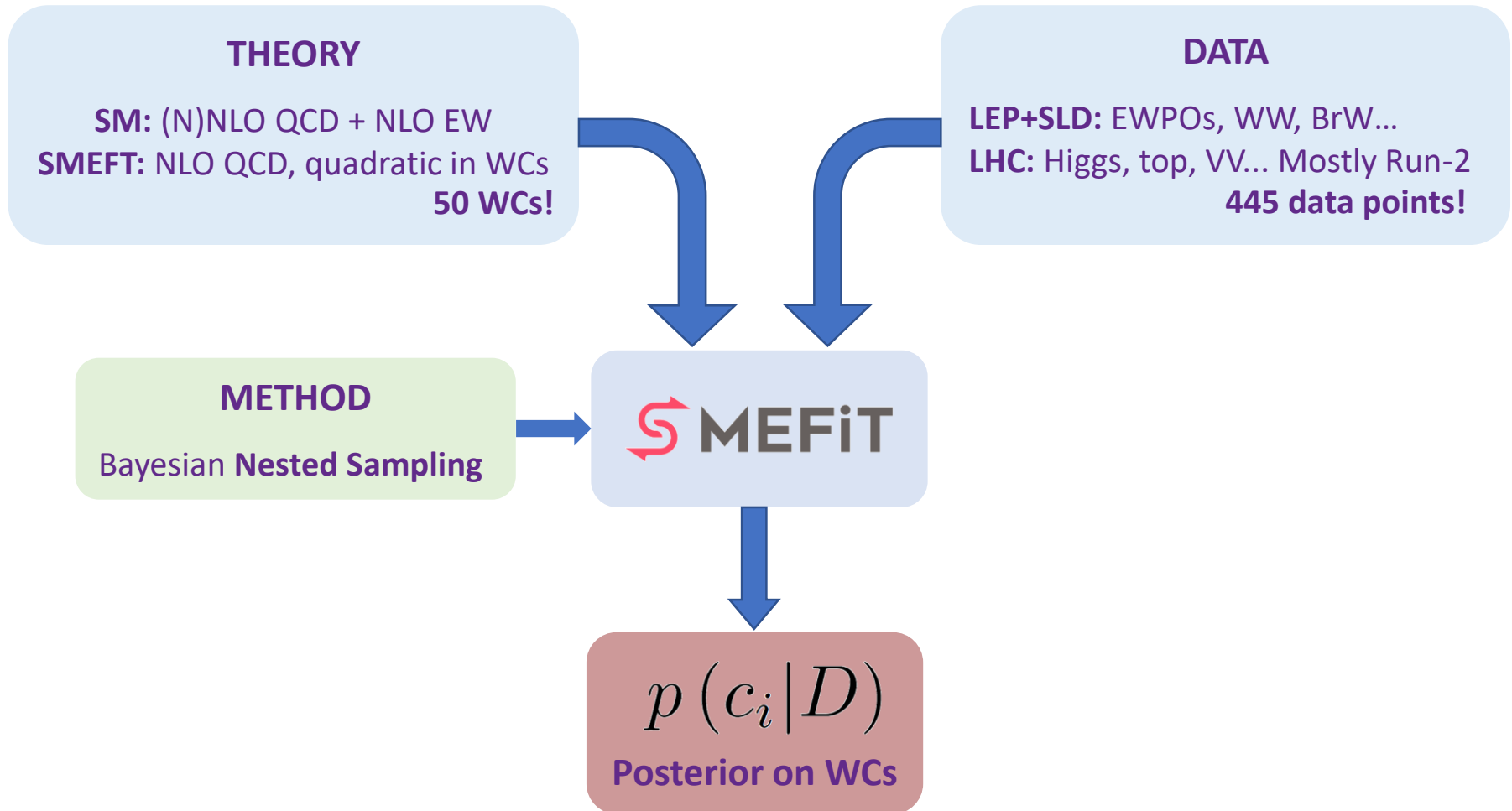


A Python software for global interpretation of particle physics data in SMEFT



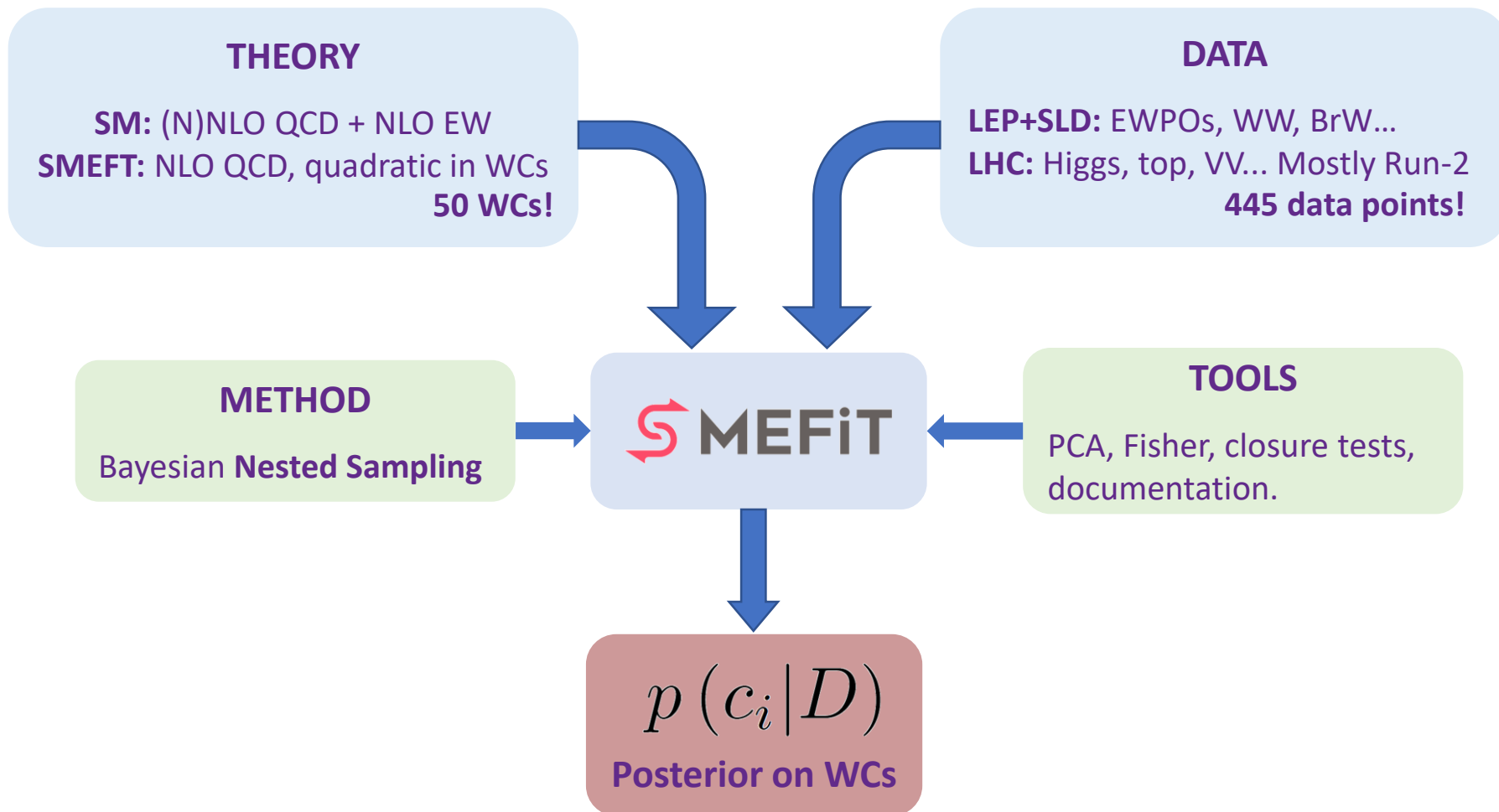


A Python software for global interpretation of particle physics data in SMEFT





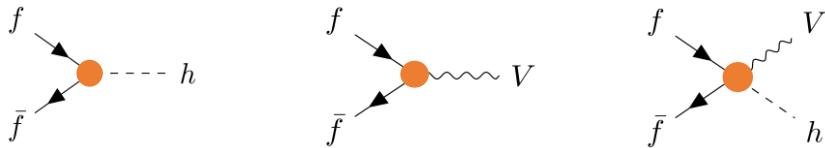
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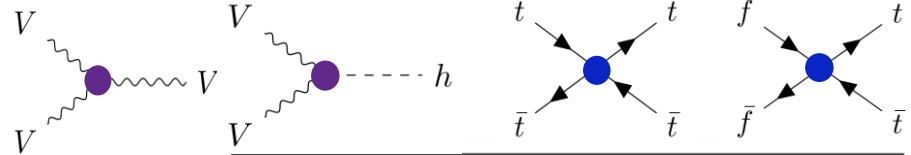
Operator basis

- Warsaw basis with rotations.
- Flavour sym.: $U(2)_q \times U(3)_d \times U(2)_u \times (U(1)_l \times U(1)_e)^3 + y_{b,c,\tau}^{SM} + c_{\varphi}(b,c,\tau)$

Operator	Coefficient	Definition	Operator	Coefficient	Definition
3rd generation quarks					
$\mathcal{O}_{\varphi Q}^{(1)}$	$c_{\varphi Q}^{(1)}$ (*)	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{Q} \gamma^\mu Q)$	\mathcal{O}_{tW}	c_{tW}	$i(\bar{Q} \tau^{\mu\nu} \tau_I t) \bar{\varphi} W_{\mu\nu}^I + \text{h.c.}$
$\mathcal{O}_{\varphi Q}^{(3)}$	$c_{\varphi Q}^{(3)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \tau_I \varphi)(\bar{Q} \gamma^\mu \tau^I Q)$	\mathcal{O}_{tB}	c_{tB} (*)	$i(\bar{Q} \tau^{\mu\nu} t) \bar{\varphi} B_{\mu\nu} + \text{h.c.}$
$\mathcal{O}_{\varphi t}$	$c_{\varphi t}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{t} \gamma^\mu t)$	\mathcal{O}_{tG}	c_{tG}	$igs(\bar{Q} \tau^{\mu\nu} T_A t) \bar{\varphi} G_{\mu\nu}^A + \text{h.c.}$
$\mathcal{O}_{t\varphi}$	$c_{t\varphi}$	$(\varphi^\dagger \varphi) \bar{Q} t \bar{\varphi} + \text{h.c.}$	$\mathcal{O}_{b\varphi}$	$c_{b\varphi}$	$(\varphi^\dagger \varphi) \bar{Q} b \varphi + \text{h.c.}$
1st, 2nd generation quarks					
$\mathcal{O}_{\varphi q}^{(1)}$	$c_{\varphi q}^{(1)}$ (*)	$\sum_{i=1,2} i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{q}_i \gamma^\mu q_i)$	$\mathcal{O}_{\varphi d}$	$c_{\varphi d}$	$\sum_{i=1,2,3} i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{d}_i \gamma^\mu d_i)$
$\mathcal{O}_{\varphi q}^{(3)}$	$c_{\varphi q}^{(3)}$	$\sum_{i=1,2} i(\varphi^\dagger \overleftrightarrow{D}_\mu \tau_I \varphi)(\bar{q}_i \gamma^\mu \tau^I q_i)$	$\mathcal{O}_{c\varphi}$	$c_{c\varphi}$	$(\varphi^\dagger \varphi) \bar{q}_2 c \bar{\varphi} + \text{h.c.}$
$\mathcal{O}_{\varphi u}$	$c_{\varphi u}$	$\sum_{i=1,2} i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{u}_i \gamma^\mu u_i)$			
two-leptons					
$\mathcal{O}_{\varphi \ell_i}$	$c_{\varphi \ell_i}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{\ell}_i \gamma^\mu \ell_i)$	$\mathcal{O}_{\varphi \mu}$	$c_{\varphi \mu}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{\mu} \gamma^\mu \mu)$
$\mathcal{O}_{\varphi \ell_i}^{(3)}$	$c_{\varphi \ell_i}^{(3)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \tau_I \varphi)(\bar{\ell}_i \gamma^\mu \tau^I \ell_i)$	$\mathcal{O}_{\varphi \tau}$	$c_{\varphi \tau}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{\tau} \gamma^\mu \tau)$
$\mathcal{O}_{\varphi e}$	$c_{\varphi e}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{e} \gamma^\mu e)$	$\mathcal{O}_{\tau\varphi}$	$c_{\tau\varphi}$	$(\varphi^\dagger \varphi) \bar{\ell}_3 \tau \varphi + \text{h.c.}$
four-leptons					
$\mathcal{O}_{\ell\ell}$	$c_{\ell\ell}$	$(\bar{\ell}_1 \gamma_\mu \ell_2)(\bar{\ell}_2 \gamma^\mu \ell_1)$			



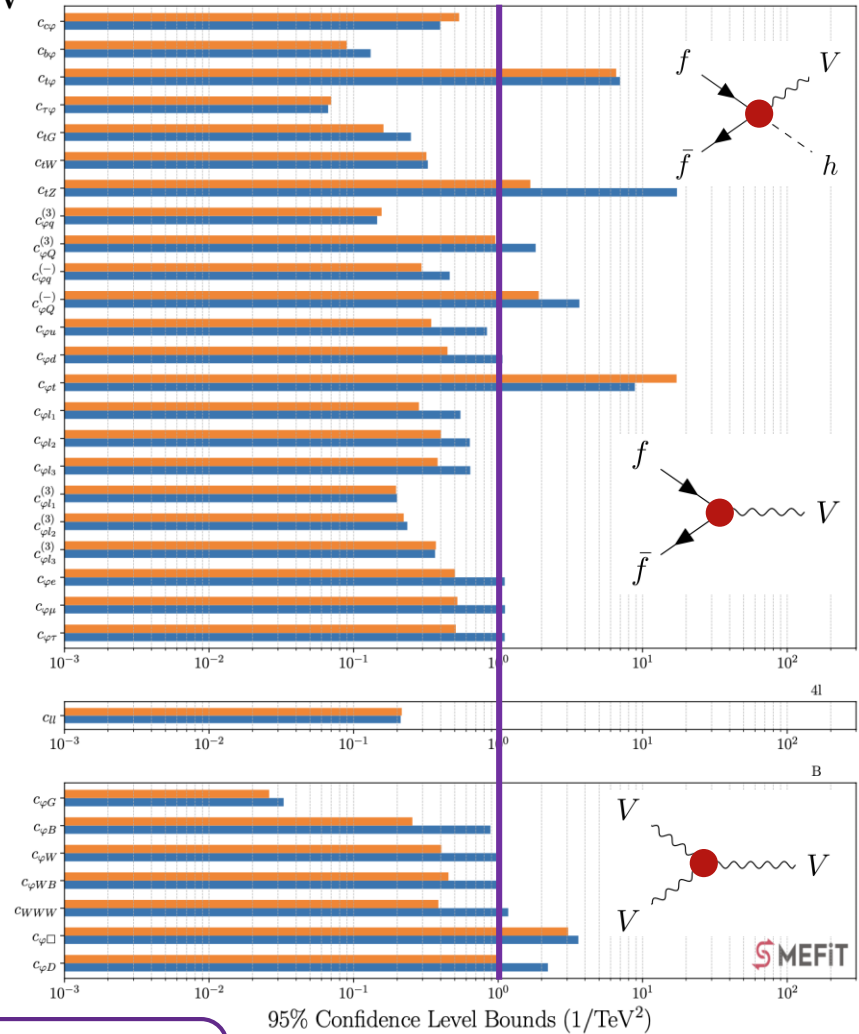
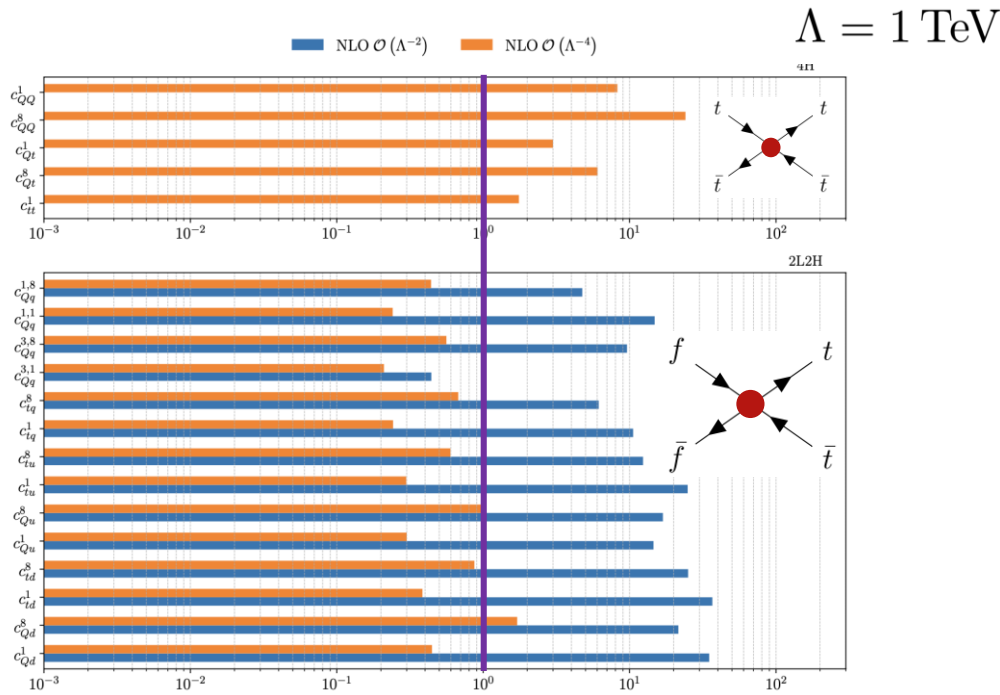
Operator	Coefficient	Definition	Operator	Coefficient	Definition
$\mathcal{O}_{\varphi G}$	$c_{\varphi G}$	$(\varphi^\dagger \varphi) G_A^{\mu\nu} G_{\mu\nu}^A$	$\mathcal{O}_{\varphi \square}$	$c_{\varphi \square}$	$\partial_\mu(\varphi^\dagger \varphi) \partial^\mu(\varphi^\dagger \varphi)$
$\mathcal{O}_{\varphi B}$	$c_{\varphi B}$	$(\varphi^\dagger \varphi) B^{\mu\nu} B_{\mu\nu}$	$\mathcal{O}_{\varphi D}$	$c_{\varphi D}$	$(\varphi^\dagger D^\mu \varphi)^\dagger (\varphi^\dagger D_\mu \varphi)$
$\mathcal{O}_{\varphi W}$	$c_{\varphi W}$	$(\varphi^\dagger \varphi) W_I^{\mu\nu} W_{\mu\nu}^I$	\mathcal{O}_W	c_{WWWW}	$\epsilon_{IJK} W_{\mu\nu}^I W^{J,\nu\rho} W_\rho^{K,\mu}$
$\mathcal{O}_{\varphi WB}$	$c_{\varphi WB}$	$(\varphi^\dagger \tau_I \varphi) B^{\mu\nu} W_{\mu\nu}^I$			



DoF	Definition (in Warsaw basis notation)	DoF	Definition (in Warsaw basis notation)
c_{QQ}^1	$2c_{qq}^{1(3333)} - \frac{2}{3}c_{qq}^{3(3333)}$	c_{QQ}^8	$8c_{qq}^{3(3333)}$
c_{Qt}^1	$c_{qu}^{1(3333)}$	c_{Qt}^8	$c_{qu}^{8(3333)}$
$c_{Qq}^{1,8}$	$\frac{1}{c_{qq}^{1(i33i)}} + 3c_{qq}^{3(i33i)}$	$c_{Qq}^{1,1}$	$\frac{1}{c_{qq}^{1(i33)}} + \frac{1}{6}c_{qq}^{1(i33i)} + \frac{1}{2}c_{qq}^{3(i33i)}$
$c_{Qq}^{3,8}$	$\frac{1}{c_{qq}^{1(i33i)}} - c_{qq}^{3(i33i)}$	$c_{Qq}^{3,1}$	$\frac{3}{c_{qq}^{1(i33)}} + \frac{1}{6}(c_{qq}^{1(i33i)} - c_{qq}^{3(i33i)})$
c_{tq}^8	$c_{qu}^{8(i333)}$	c_{tq}^1	$c_{qu}^{1(i333)}$
c_{tu}^8	$2c_{uu}^{(i33i)}$	c_{tu}^1	$\frac{(i333)}{c_{uu}} + \frac{1}{3}c_{uu}^{(i33i)}$
c_{Qu}^8	$c_{qu}^{8(33ii)}$	c_{Qu}^1	$c_{qu}^{1(33ii)}$
c_{ud}^8	$c_{ud}^{8(33jj)}$	c_{ud}^1	$c_{ud}^{1(33jj)}$
c_{Qd}^8	$c_{qd}^{8(33jj)}$	c_{Qd}^1	$c_{qd}^{1(33jj)}$

Fit of 45 (50) WCs at the linear (quadratic) level

MEFIT 3.0 results



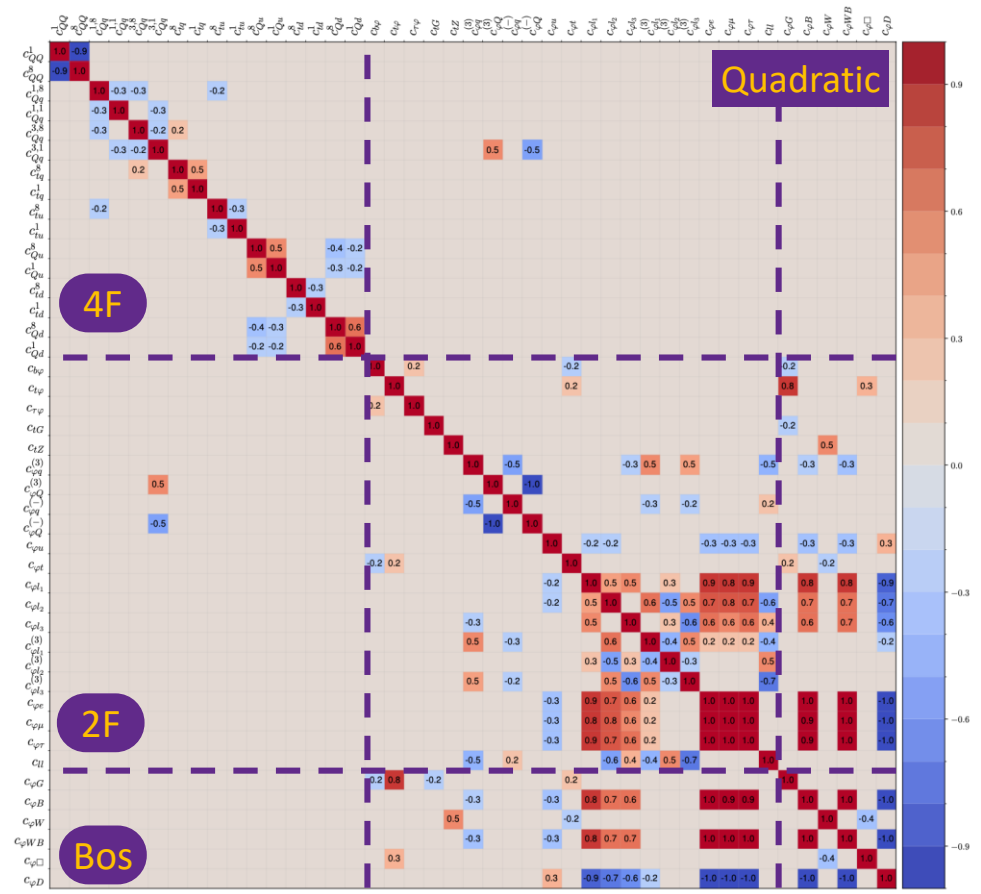
- Marginalised fit
- LEP + LHC
- Linear in EFT
- Quadratic in EFT

Many bounds dominated by quadratic contributions



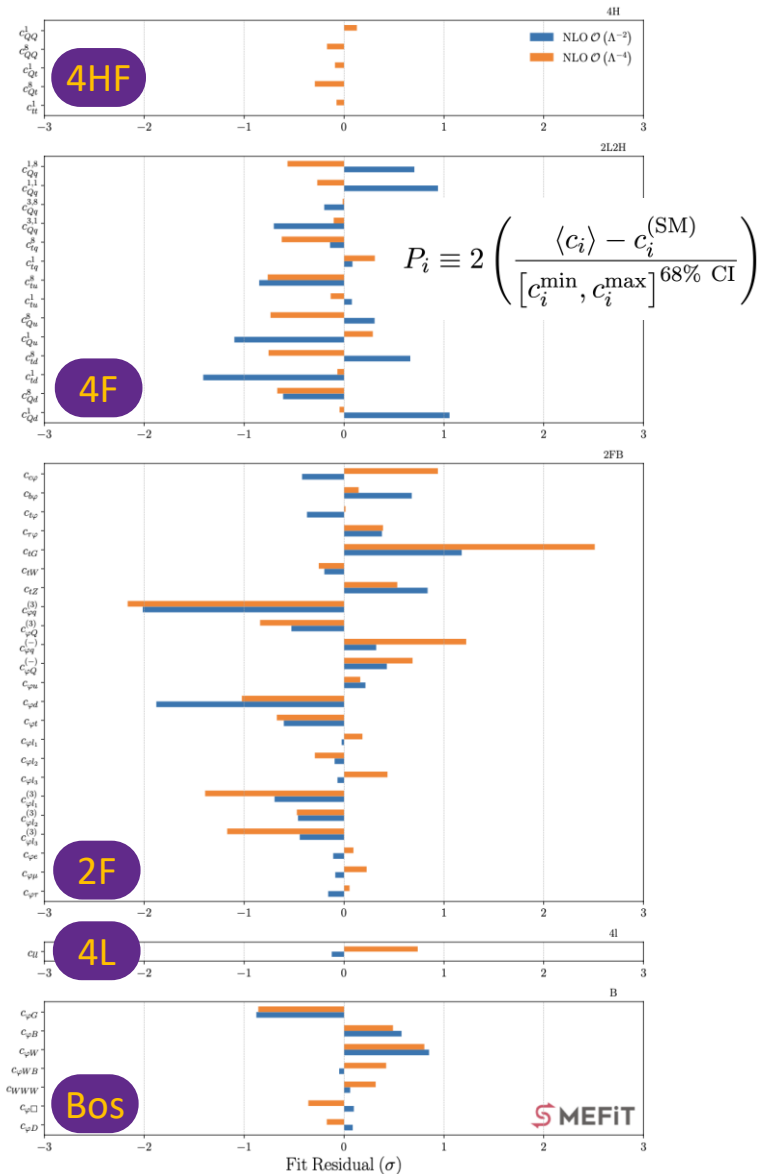
Pulls and correlations

No large correlations



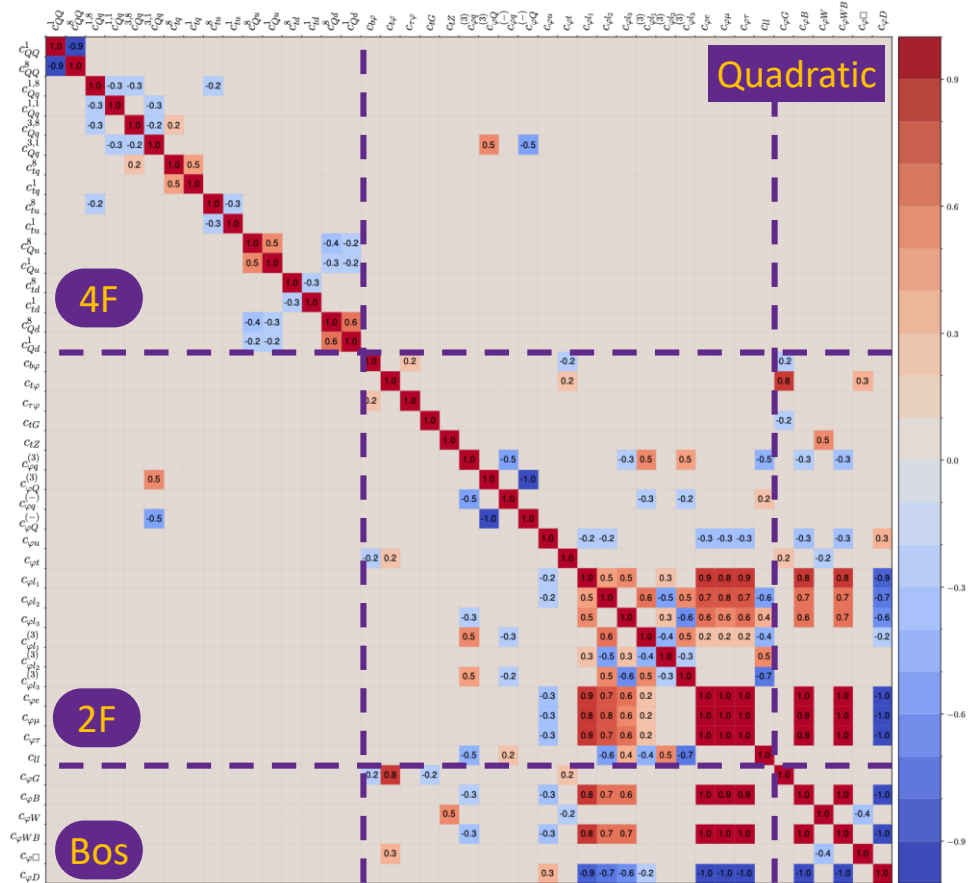
Correlation: NLO $\mathcal{O}(\Lambda^{-4})$

Pulls and correlations

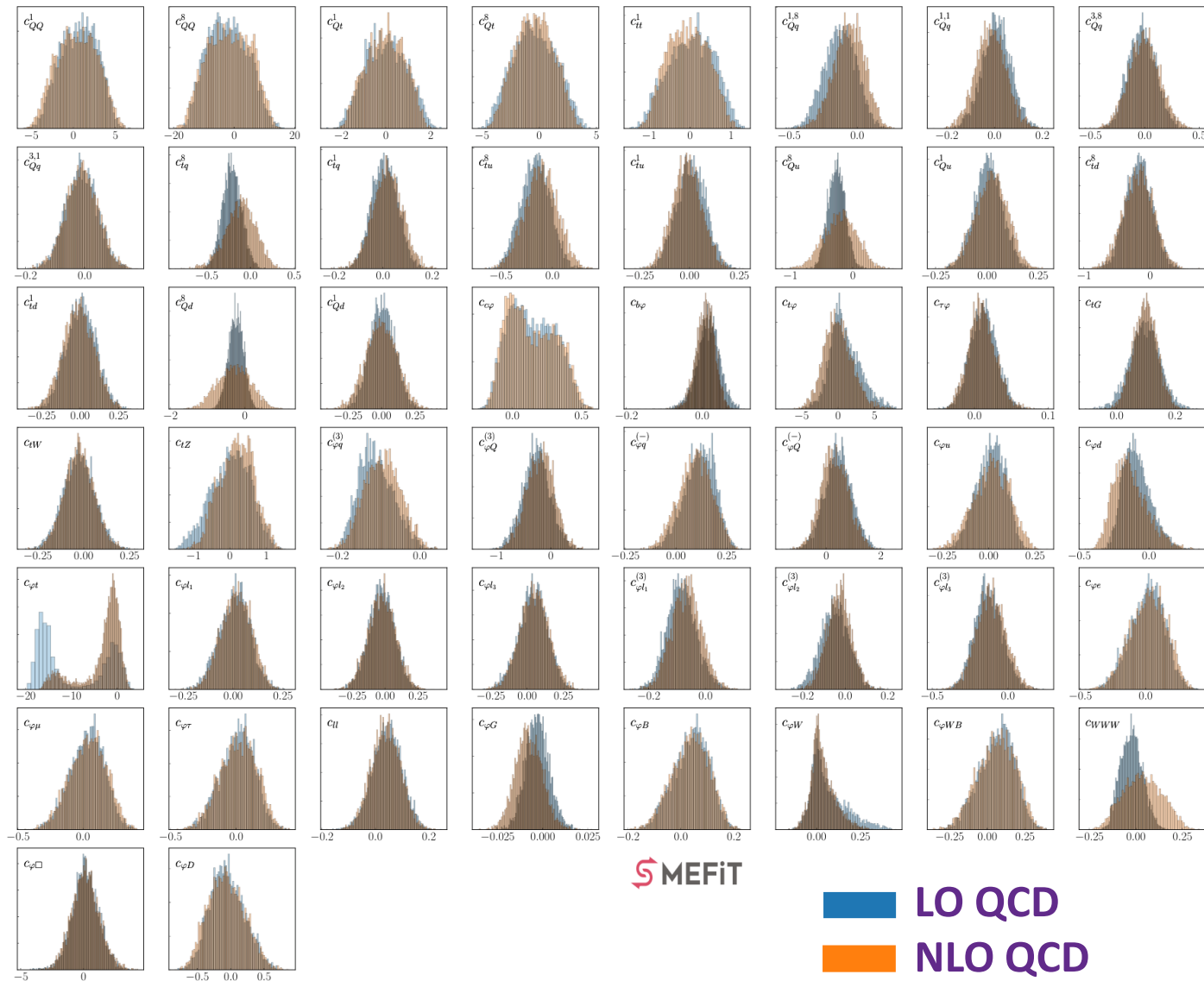


Fit residuals (pulls) largely SM compatible

No large correlations



NLO QCD in the EFT effects



An eye on the future

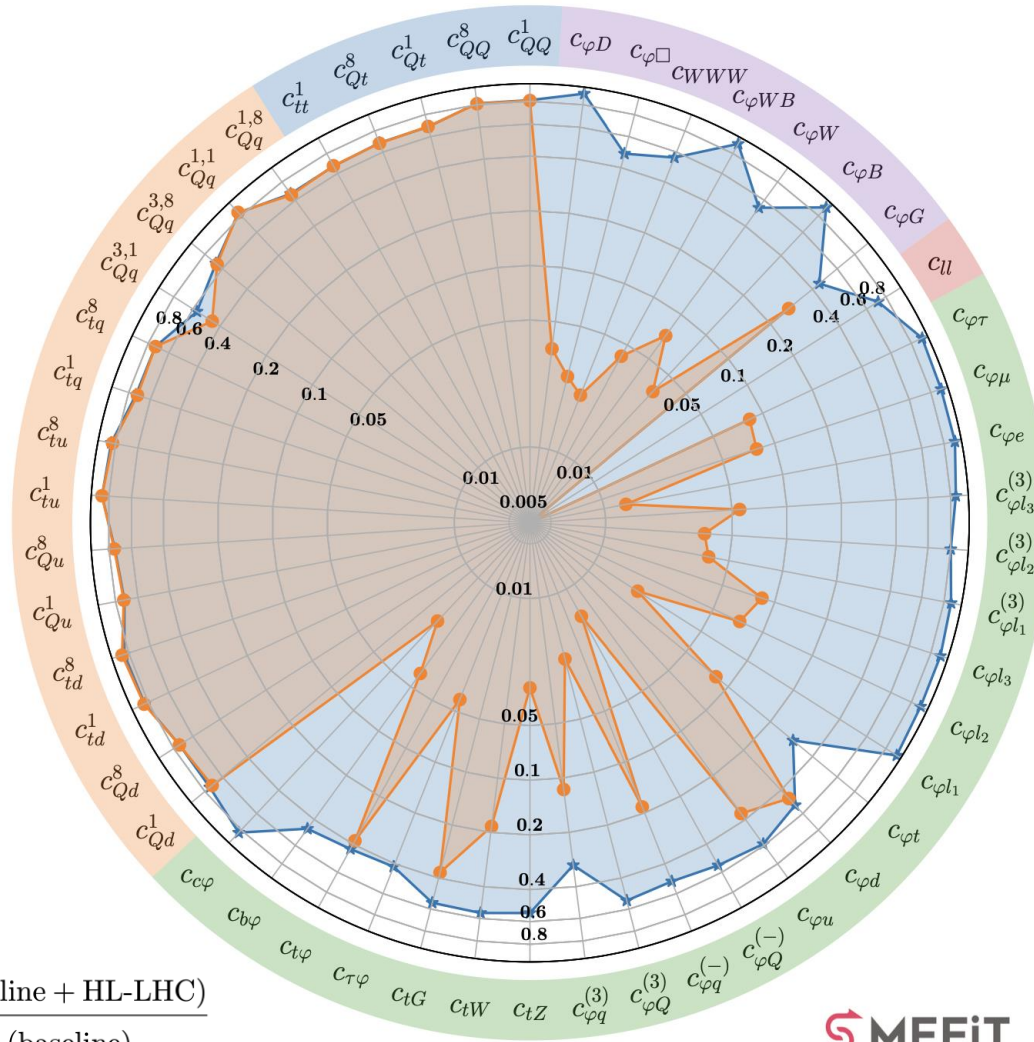


Implemented in SMEFiT



Improvements at HL-LHC and FCC-ee

Ratio of Uncertainties to SMEFiT3.0 Baseline, $\mathcal{O}(\Lambda^{-4})$, Marginalised



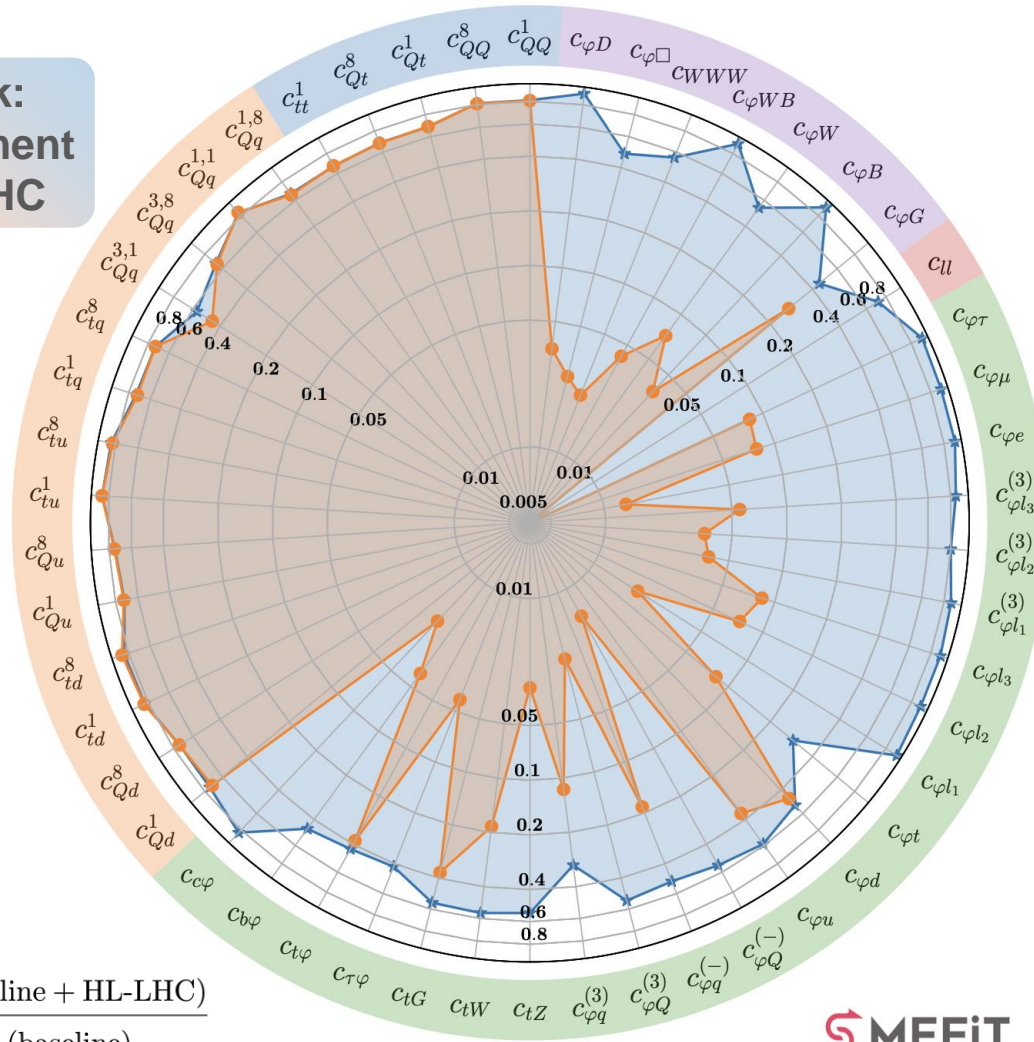
★ HL-LHC ● HL-LHC + FCC-ee

$$R_{\delta c_i} = \frac{[c_i^{\min}, c_i^{\max}]^{95\% \text{ CL}} (\text{baseline} + \text{HL-LHC})}{[c_i^{\min}, c_i^{\max}]^{95\% \text{ CL}} (\text{baseline})}$$

Improvements at HL-LHC and FCC-ee

Ratio of Uncertainties to SMEFiT3.0 Baseline, $\mathcal{O}(\Lambda^{-4})$, Marginalised

4-quark:
improvement
at HL-LHC



★ HL-LHC ● HL-LHC + FCC-ee

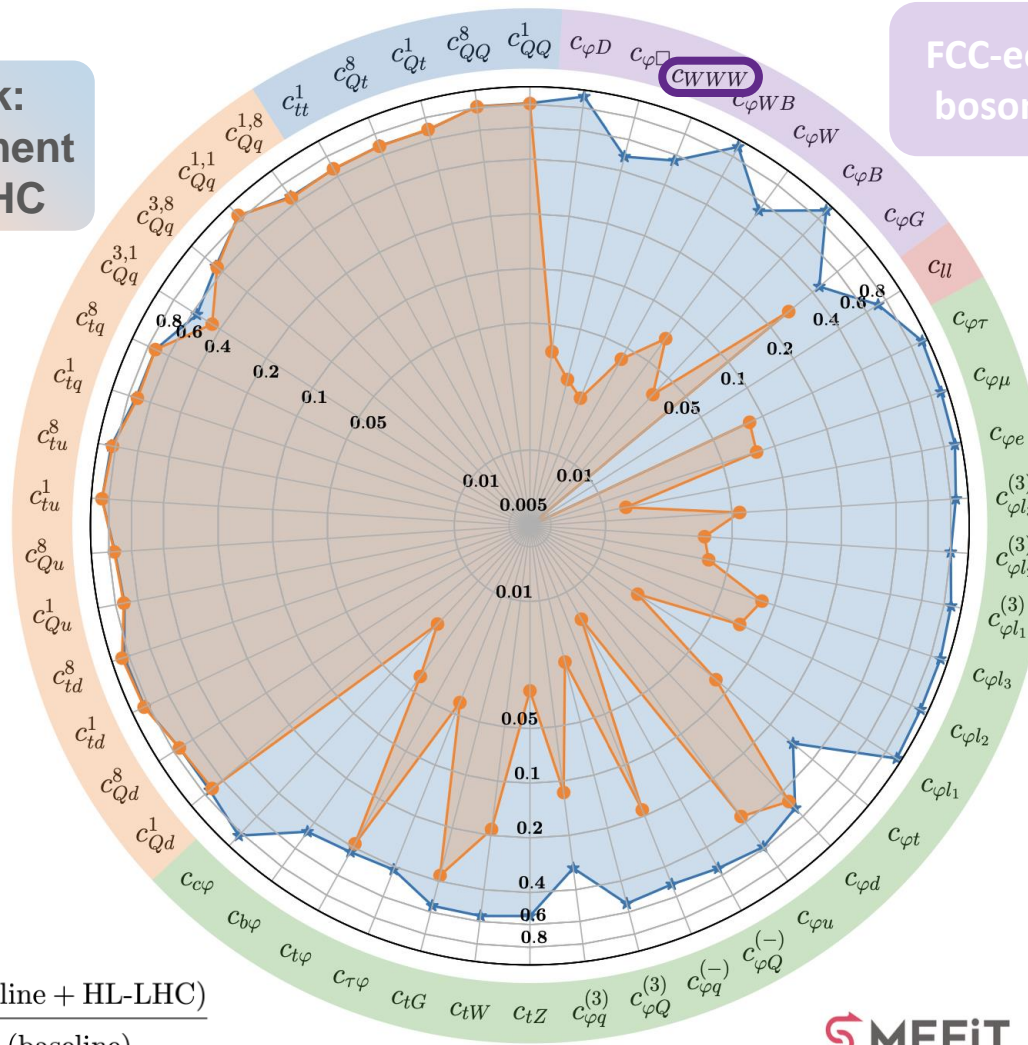
$$R_{\delta c_i} = \frac{[c_i^{\min}, c_i^{\max}]^{95\% \text{ CL}} (\text{baseline} + \text{HL-LHC})}{[c_i^{\min}, c_i^{\max}]^{95\% \text{ CL}} (\text{baseline})}$$

Improvements at HL-LHC and FCC-ee

Ratio of Uncertainties to SMEFiT3.0 Baseline, $\mathcal{O}(\Lambda^{-4})$, Marginalised

4-quark:
improvement
at HL-LHC

FCC-ee probes well
bosonic operators



$$R_{\delta c_i} = \frac{[c_i^{\min}, c_i^{\max}]^{95\% \text{ CL}} (\text{baseline} + \text{HL-LHC})}{[c_i^{\min}, c_i^{\max}]^{95\% \text{ CL}} (\text{baseline})}$$



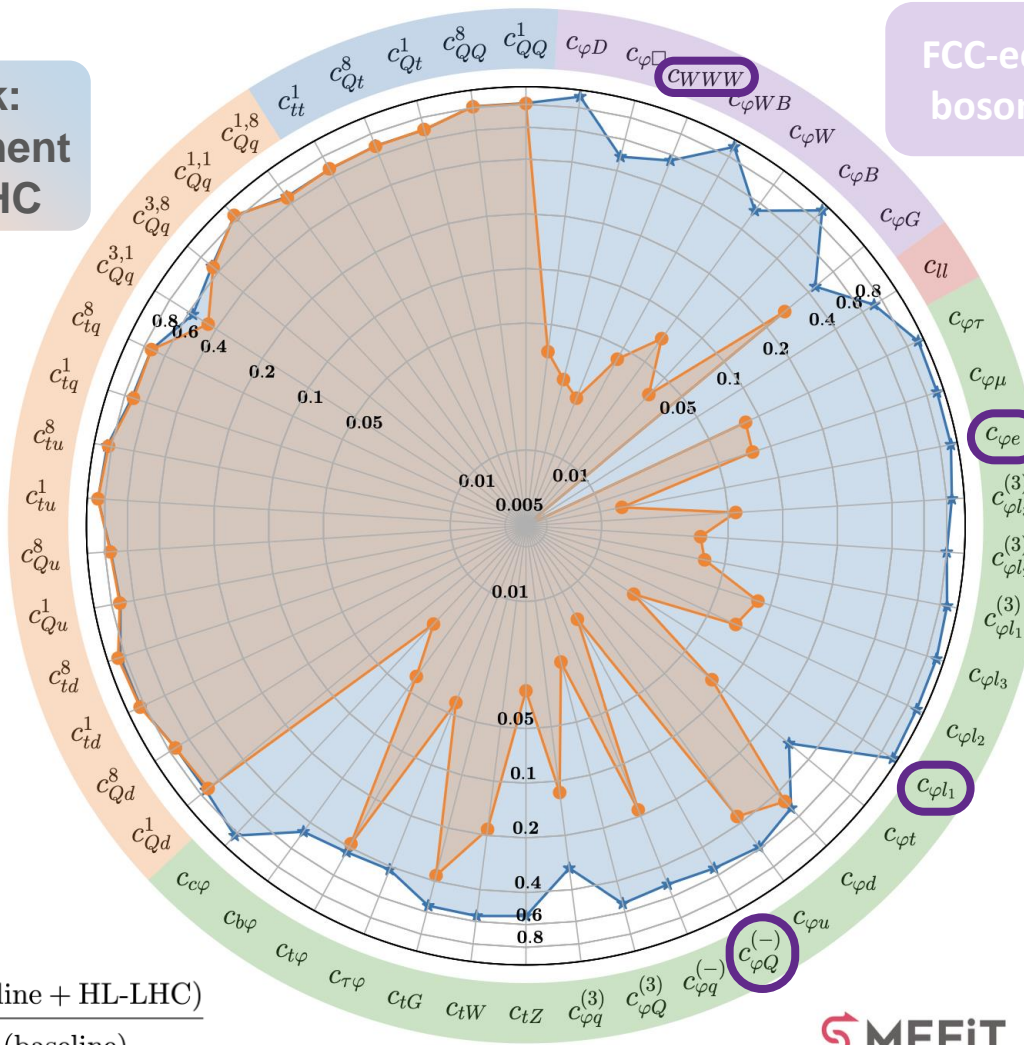
★ HL-LHC ● HL-LHC + FCC-ee

Improvements at HL-LHC and FCC-ee

Ratio of Uncertainties to SMEFiT3.0 Baseline, $\mathcal{O}(\Lambda^{-4})$, Marginalised

4-quark:
improvement
at HL-LHC

FCC-ee probes well
bosonic operators



And several 2-fermion operators!

$$R_{\delta c_i} = \frac{[c_i^{\min}, c_i^{\max}]^{95\% \text{ CL}} (\text{baseline} + \text{HL-LHC})}{[c_i^{\min}, c_i^{\max}]^{95\% \text{ CL}} (\text{baseline})}$$



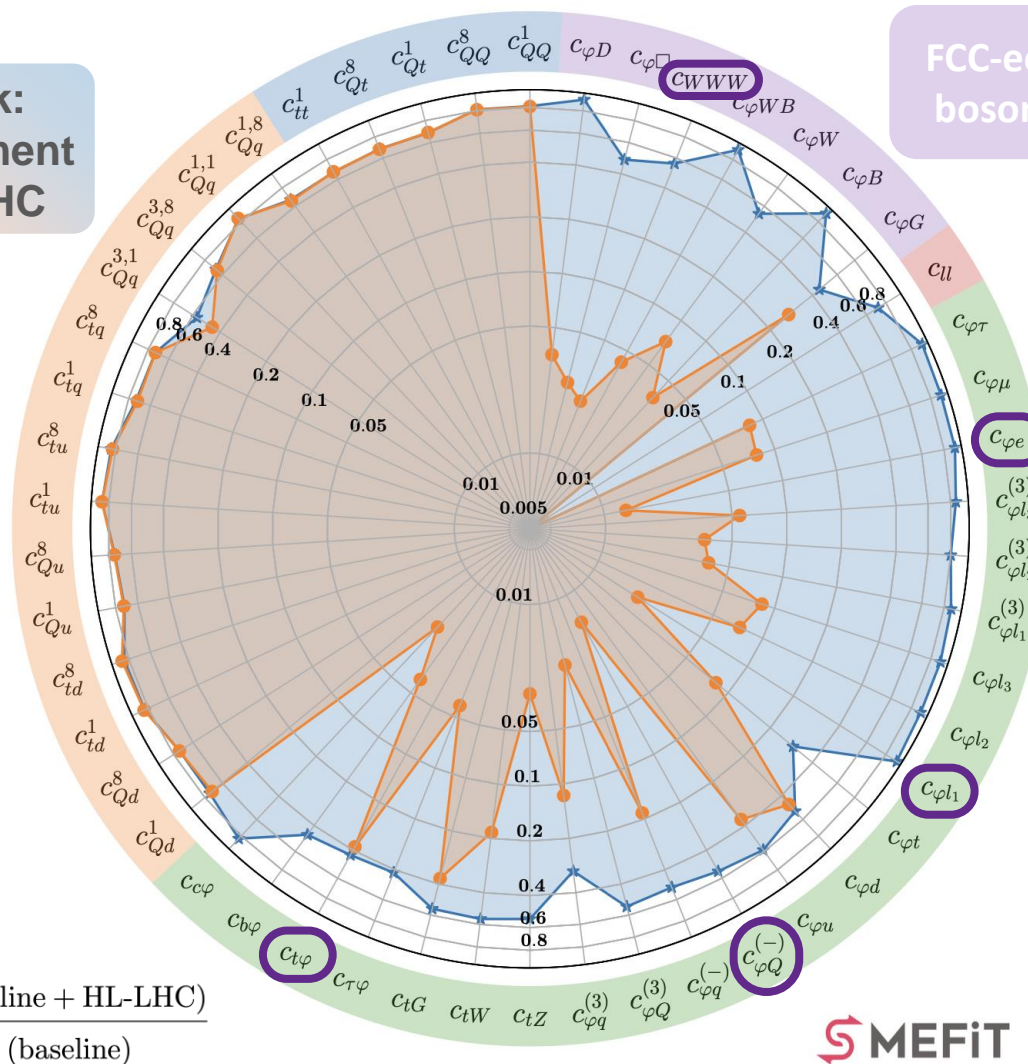
★ HL-LHC ● HL-LHC + FCC-ee

Improvements at HL-LHC and FCC-ee

Ratio of Uncertainties to SMEFiT3.0 Baseline, $\mathcal{O}(\Lambda^{-4})$, Marginalised

4-quark:
improvement
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FCC-ee probes well
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And several 2-fermion operators!

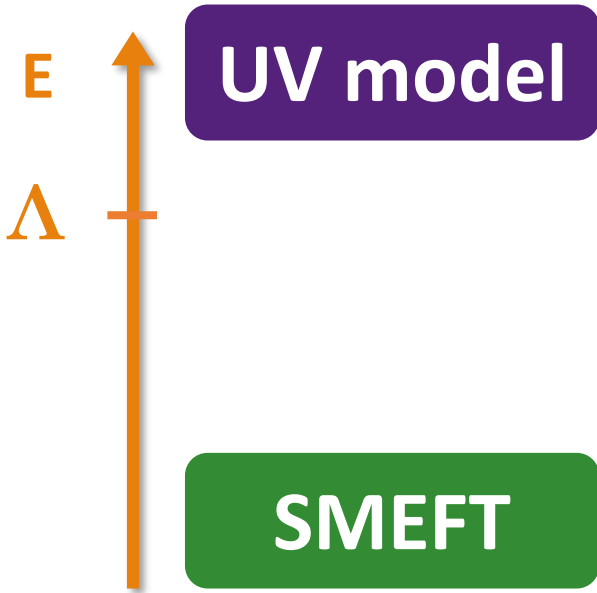
Top Yukawa only at HL-LHC

$$R_{\delta c_i} = \frac{[c_i^{\min}, c_i^{\max}]^{95\% \text{ CL}} (\text{baseline} + \text{HL-LHC})}{[c_i^{\min}, c_i^{\max}]^{95\% \text{ CL}} (\text{baseline})}$$

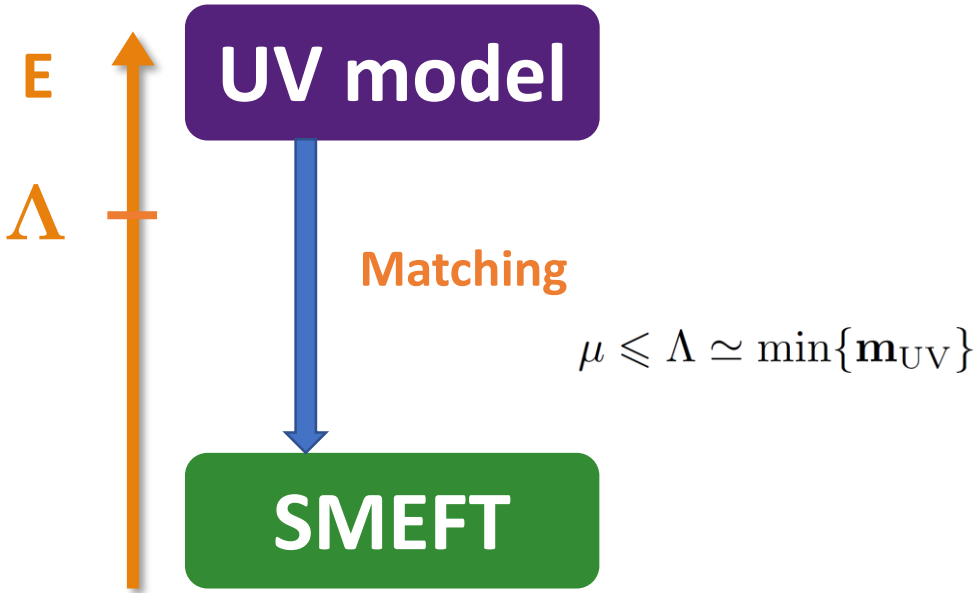


★ HL-LHC ● HL-LHC + FCC-ee

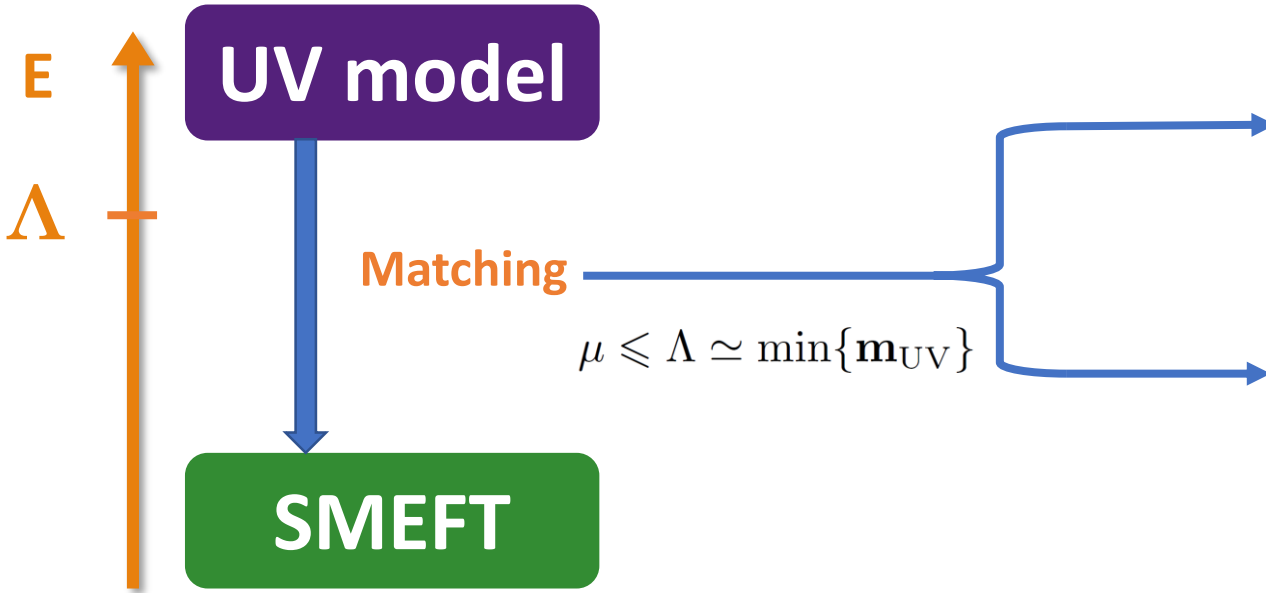
UV perspective



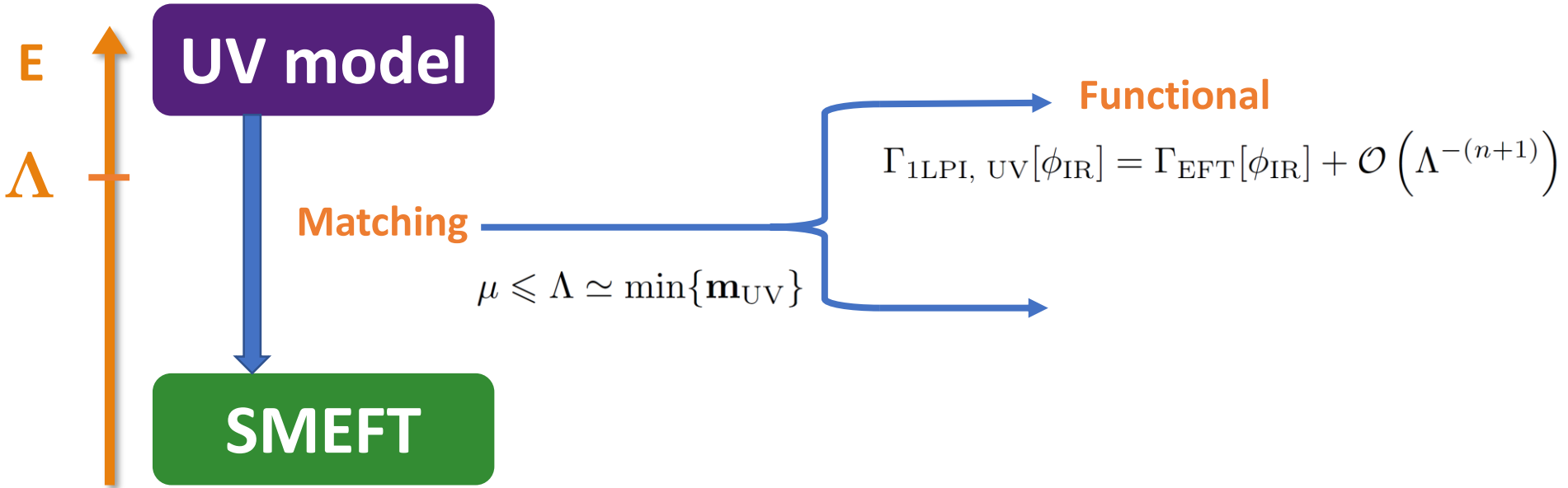
UV perspective



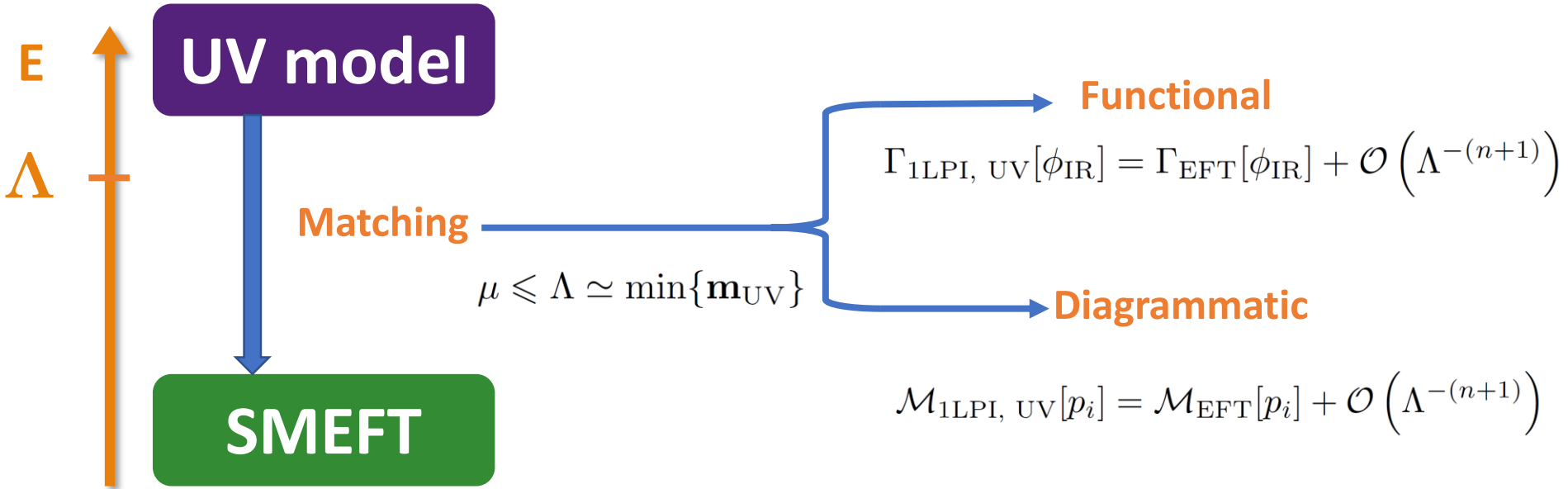
UV perspective



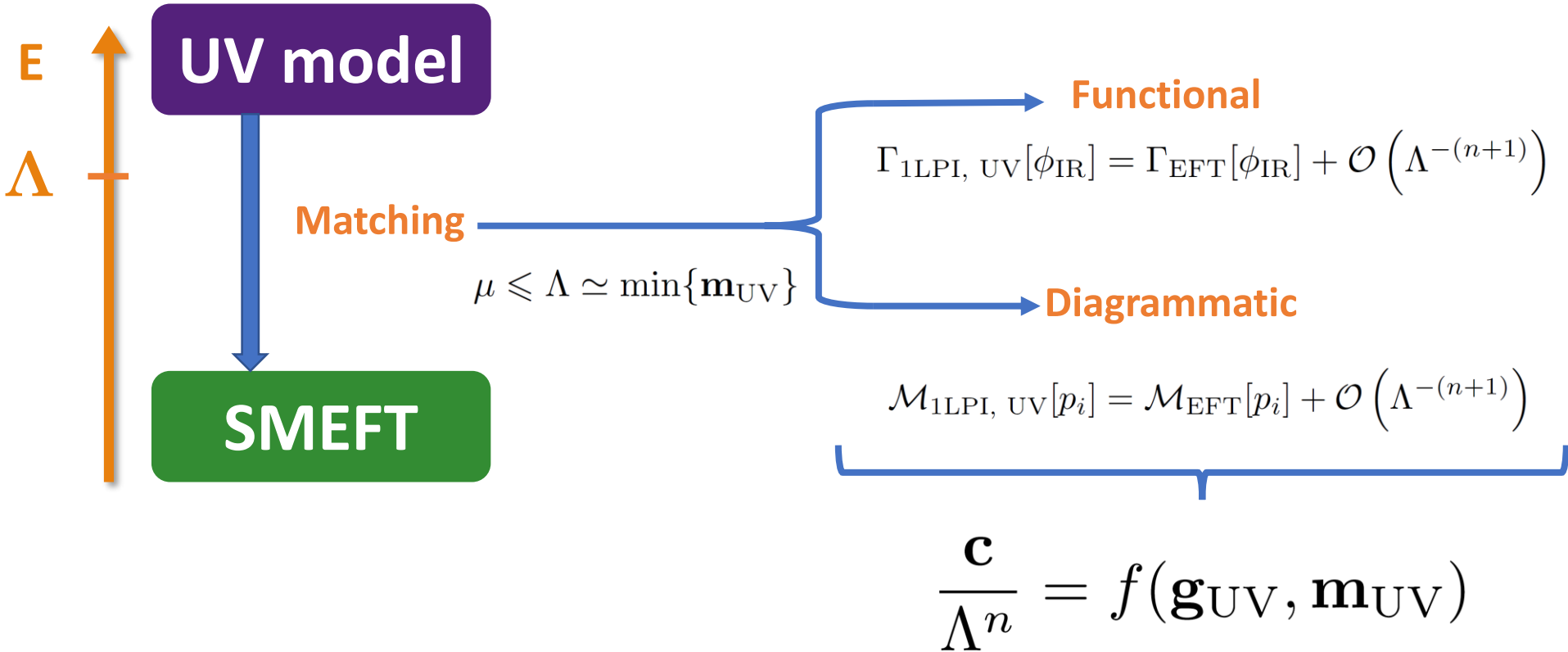
UV perspective



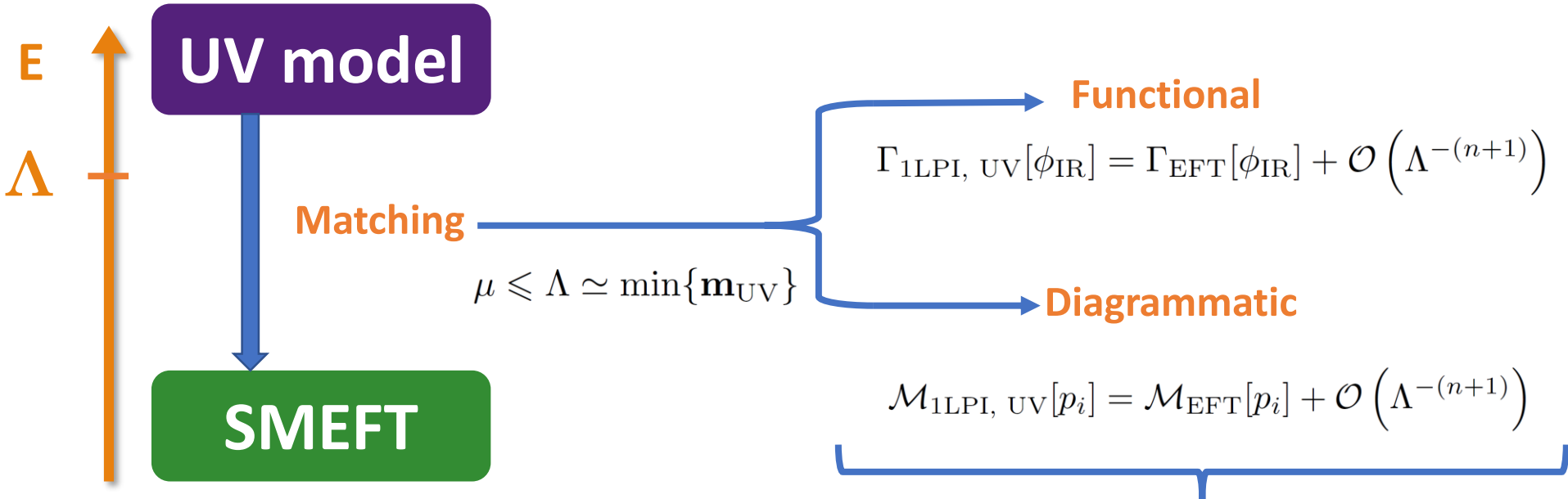
UV perspective



UV perspective



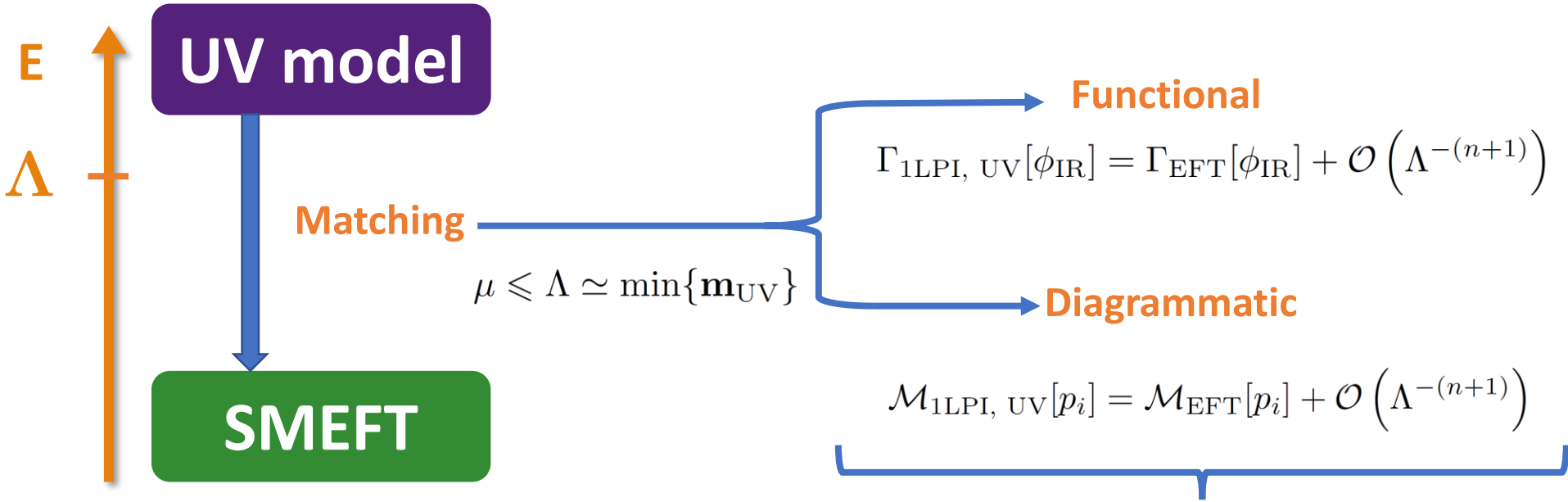
UV perspective



$$\frac{\mathbf{c}}{\Lambda^n} = f(\mathbf{g}_{\text{UV}}, \mathbf{m}_{\text{UV}})$$

$$\sigma(\mathbf{c}) \longrightarrow \sigma(\mathbf{g}_{\text{UV}})$$

UV perspective



- Less parameters
- Stronger correlations
- Model dependent
- Sharper interpretation

$$\frac{\mathbf{c}}{\Lambda^n} = f(\mathbf{g}_{\text{UV}}, \mathbf{m}_{\text{UV}})$$

$$\sigma(\mathbf{c}) \longrightarrow \sigma(\mathbf{g}_{\text{UV}})$$

The state of matching affairs

Automated
1-loop
matching*

***Only up to tree level for heavy spin-1 bosons.**

The state of matching affairs



*Only up to tree level for heavy spin-1 bosons.

The state of matching affairs

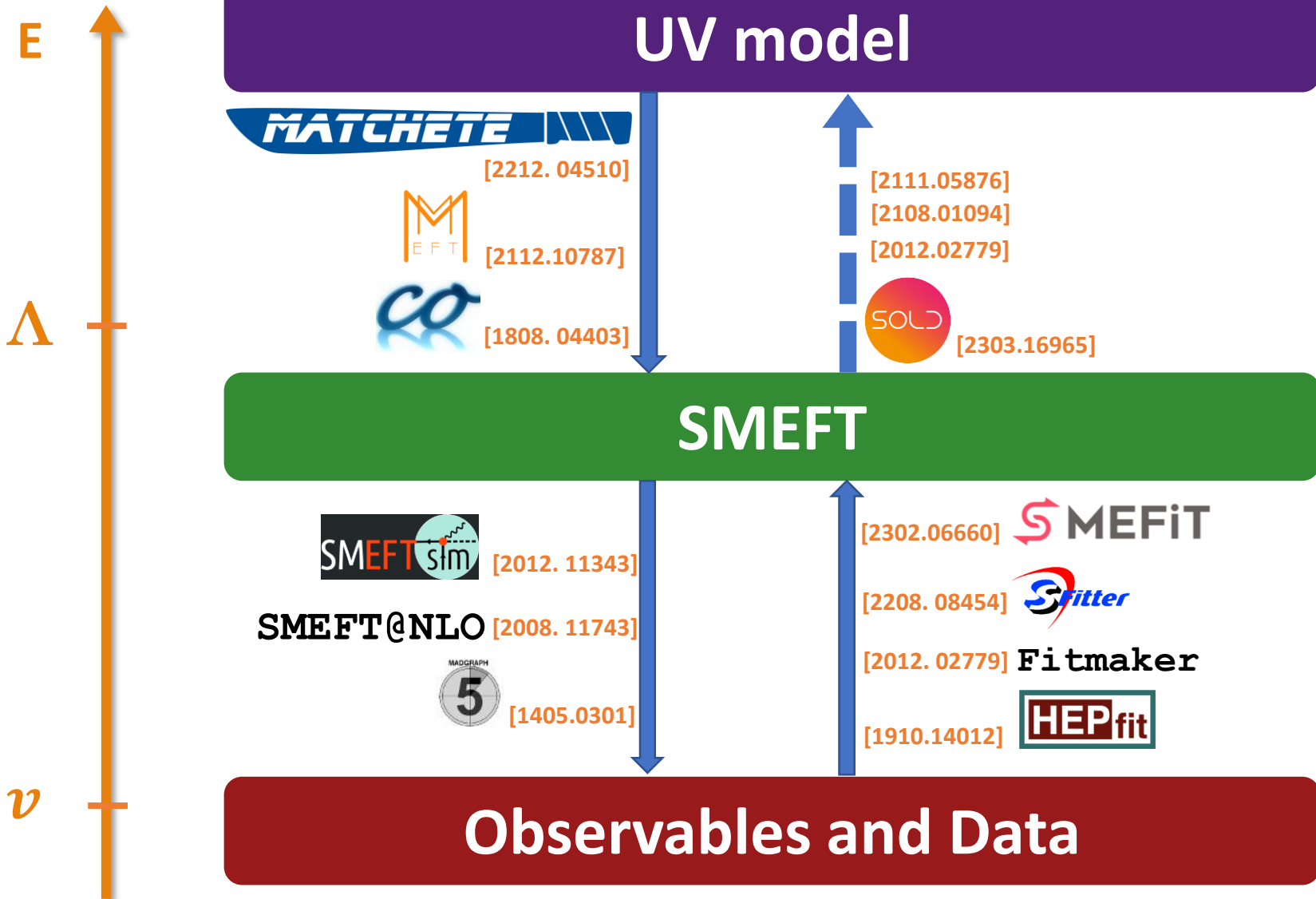


- On-shell matching techniques [2308.00035] [F. Vilches@Planck24] [J. López Miras@HEFT2024]
- Matching with Spontaneous Symmetry Breaking [2404.11640]
- Functional matching for dim. >6 [2306.09103] [2308.03849] [2311.12757]
- 1-loop dictionaries [2303.16965] [P. Olgoso@HEFT2024] [J. Gargalionis@Planck24]
- 2-loop matching [2311.13630] [J. Fuentes-Martín@HEFT2024]

Disclaimer: incomplete list.

*Only up to tree level for heavy spin-1 bosons.

Automation across scales



Apologies for not including all tools/codes due to space restrictions.

Bridging the gap



match2fit

Automating the reuse of SMEFT predictions
and global fits to bound UV models

Photo by Didier Descouens - Own work, CC BY-SA 4.0

Constraints on WCs from UV matching

Tree-level matching

$$\frac{\left(c_{qd}^{(1)}\right)_{3333}}{\Lambda^2} = -\frac{\left(y_\phi^d\right)_{33}^2}{6 m_\phi^2}, \quad \frac{\left(c_{qd}^{(8)}\right)_{3333}}{\Lambda^2} = -\frac{\left(y_\phi^d\right)_{33}^2}{m_\phi^2}, \quad \frac{\left(c_{d\varphi}\right)_{33}}{\Lambda^2} = \frac{\lambda_\phi \left(y_\phi^d\right)_{33}}{m_\phi^2}, \quad \frac{c_\varphi}{\Lambda^2} = \frac{\lambda_\phi^2}{m_\phi^2}$$

Constraints on WCs from UV matching

Tree-level matching

$$\frac{(c_{qd}^{(1)})_{3333}}{\Lambda^2} = -\frac{(y_\phi^d)_{33}^2}{6m_\phi^2}, \quad \frac{(c_{qd}^{(8)})_{3333}}{\Lambda^2} = -\frac{(y_\phi^d)_{33}^2}{m_\phi^2}, \quad \frac{(c_{d\varphi})_{33}}{\Lambda^2} = \frac{\lambda_\phi (y_\phi^d)_{33}}{m_\phi^2}, \quad \frac{c_\varphi}{\Lambda^2} = \frac{\lambda_\phi^2}{m_\phi^2}$$

One-loop level matching

$$\begin{aligned} \frac{c_{\varphi\Box}}{\Lambda^2} &= -\frac{g_1^4}{7680\pi^2} \frac{1}{m_\phi^2} - \frac{g_2^4}{2560\pi^2} \frac{1}{m_\phi^2} - \frac{3}{32\pi^2} \frac{\lambda_\phi^2}{m_\phi^2}, \\ \frac{c_{t\varphi}}{\Lambda^2} &= -\frac{\lambda_\phi (y_\phi^u)_{33}}{m_\phi^2} - \frac{g_2^4 y_t^{\text{SM}}}{3840\pi^2} \frac{1}{m_\phi^2} + \frac{y_t^{\text{SM}}}{16\pi^2} \frac{\lambda_\phi^2}{m_\phi^2} + \frac{(4(y_b^{\text{SM}})^2 - 13(y_t^{\text{SM}})^2) \lambda_\phi (y_\phi^u)_{33}}{64\pi^2 m_\phi^2} \\ &\quad - \left(12\lambda_\varphi^{\text{SM}} + (y_b^{\text{SM}})^2 - 11(y_t^{\text{SM}})^2\right) \frac{y_t^{\text{SM}} (y_\phi^u)_{33}^2}{64\pi^2 m_\phi^2} + \frac{3}{128\pi^2} \frac{\lambda_\phi (y_\phi^u)_{33}^3}{m_\phi^2}, \end{aligned}$$

Constrains on WCs from UV matching

Tree-level matching

$$\frac{(c_{qd}^{(1)})_{3333}}{\Lambda^2} = -\frac{(y_\phi^d)_{33}^2}{6m_\phi^2}, \quad \frac{(c_{qd}^{(8)})_{3333}}{\Lambda^2} = -\frac{(y_\phi^d)_{33}^2}{m_\phi^2}, \quad \frac{(c_{d\varphi})_{33}}{\Lambda^2} = \frac{\lambda_\phi (y_\phi^d)_{33}}{m_\phi^2}, \quad \frac{c_\varphi}{\Lambda^2} = \frac{\lambda_\phi^2}{m_\phi^2}$$

One-loop level matching

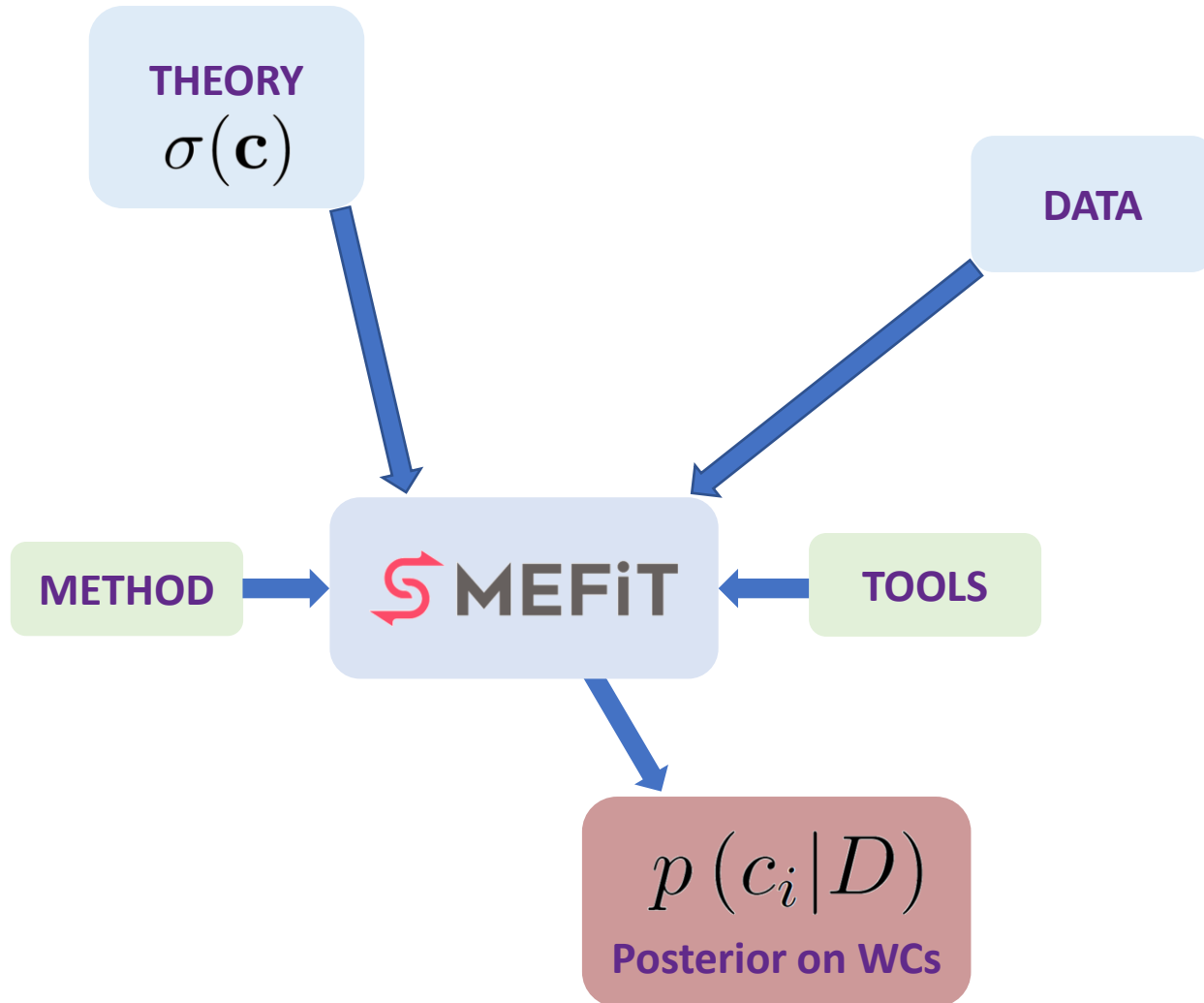
$$\frac{c_{\varphi\Box}}{\Lambda^2} = -\frac{g_1^4}{7680\pi^2} \frac{1}{m_\phi^2} - \frac{g_2^4}{2560\pi^2} \frac{1}{m_\phi^2} - \frac{3}{32\pi^2} \frac{\lambda_\phi^2}{m_\phi^2},$$

$$\frac{c_{t\varphi}}{\Lambda^2} = -\frac{\lambda_\phi (y_\phi^u)_{33}}{m_\phi^2} - \frac{g_2^4 y_t^{\text{SM}}}{3840\pi^2} \frac{1}{m_\phi^2} + \frac{y_t^{\text{SM}}}{16\pi^2} \frac{\lambda_\phi^2}{m_\phi^2} + \frac{(4 (y_b^{\text{SM}})^2 - 13 (y_t^{\text{SM}})^2) \lambda_\phi (y_\phi^u)_{33}}{64\pi^2 m_\phi^2}$$

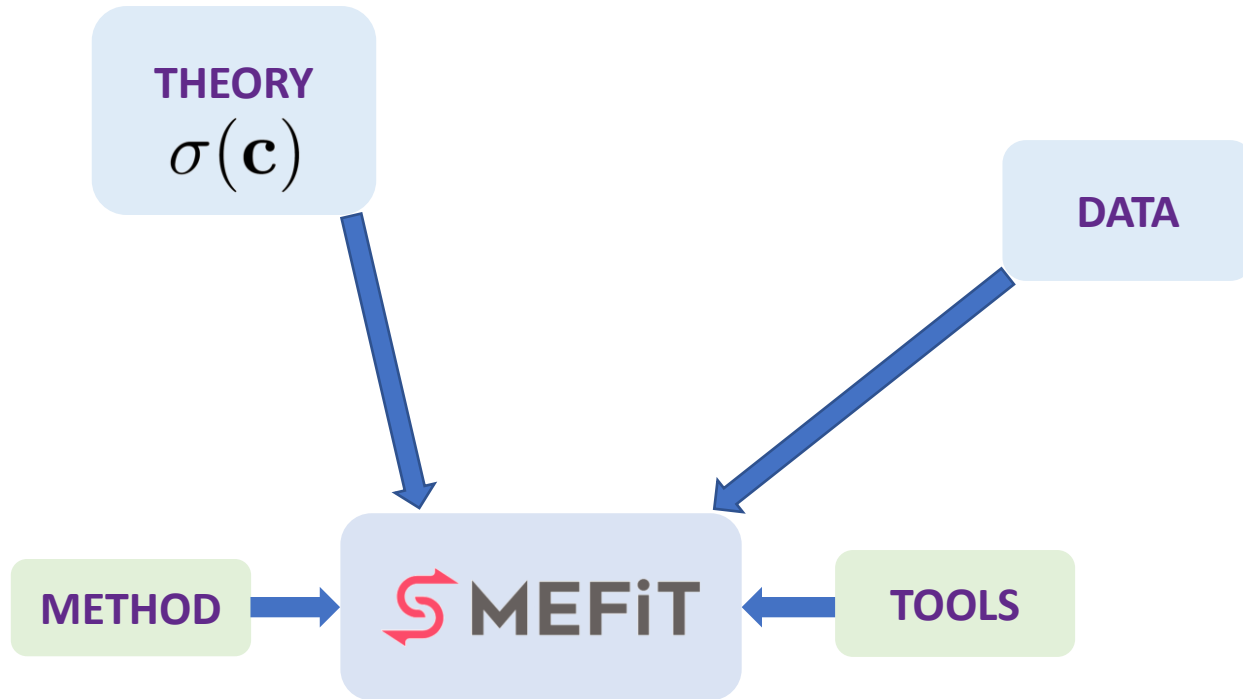
$$- \left(12\lambda_\phi^{\text{SM}} + (y_b^{\text{SM}})^2 - 11 (y_t^{\text{SM}})^2 \right) \frac{y_t^{\text{SM}} (y_\phi^u)_{33}^2}{64\pi^2 m_\phi^2} + \frac{3}{128\pi^2} \frac{\lambda_\phi (y_\phi^u)_{33}^3}{m_\phi^2},$$

UV constraints on the WC space are highly non trivial

Reusing EFT global fits for the UV

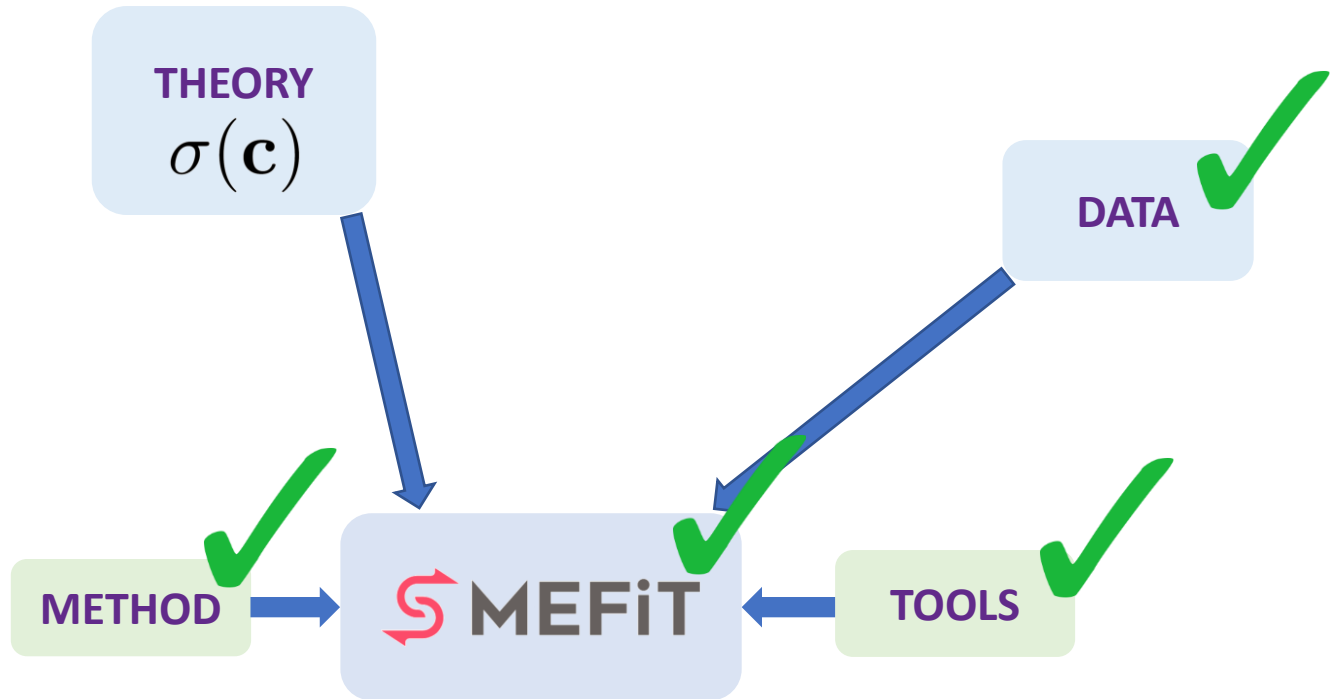


Reusing EFT global fits for the UV



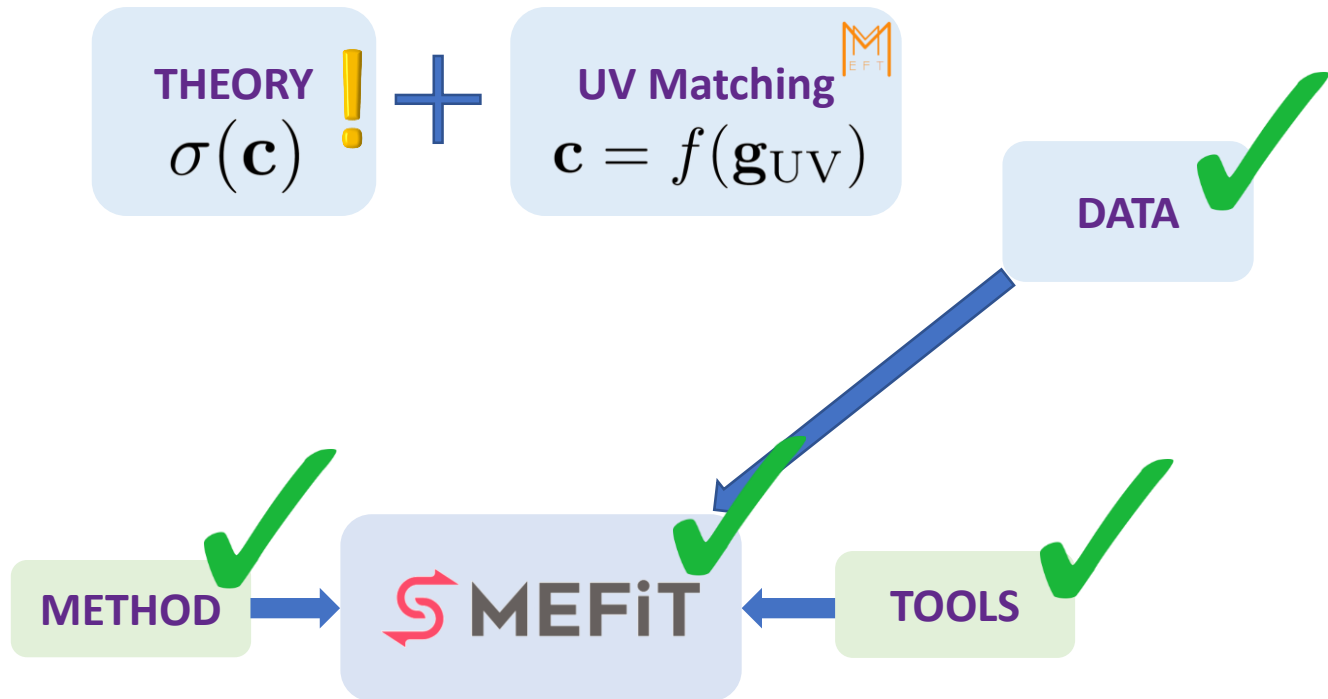
$p(c_i | D)$ ~~X~~
Posterior on WCs

Reusing EFT global fits for the UV



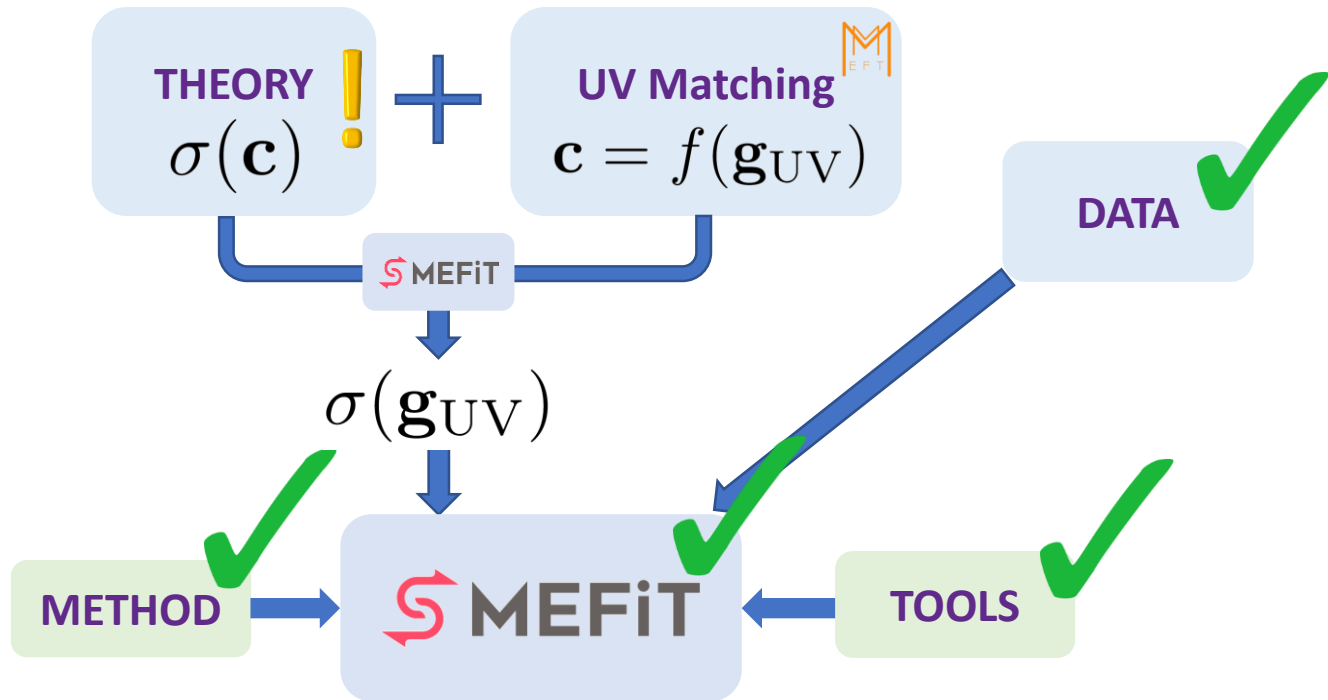
$p(c_i | D)$ ✗
Posterior on WCs

Reusing EFT global fits for the UV



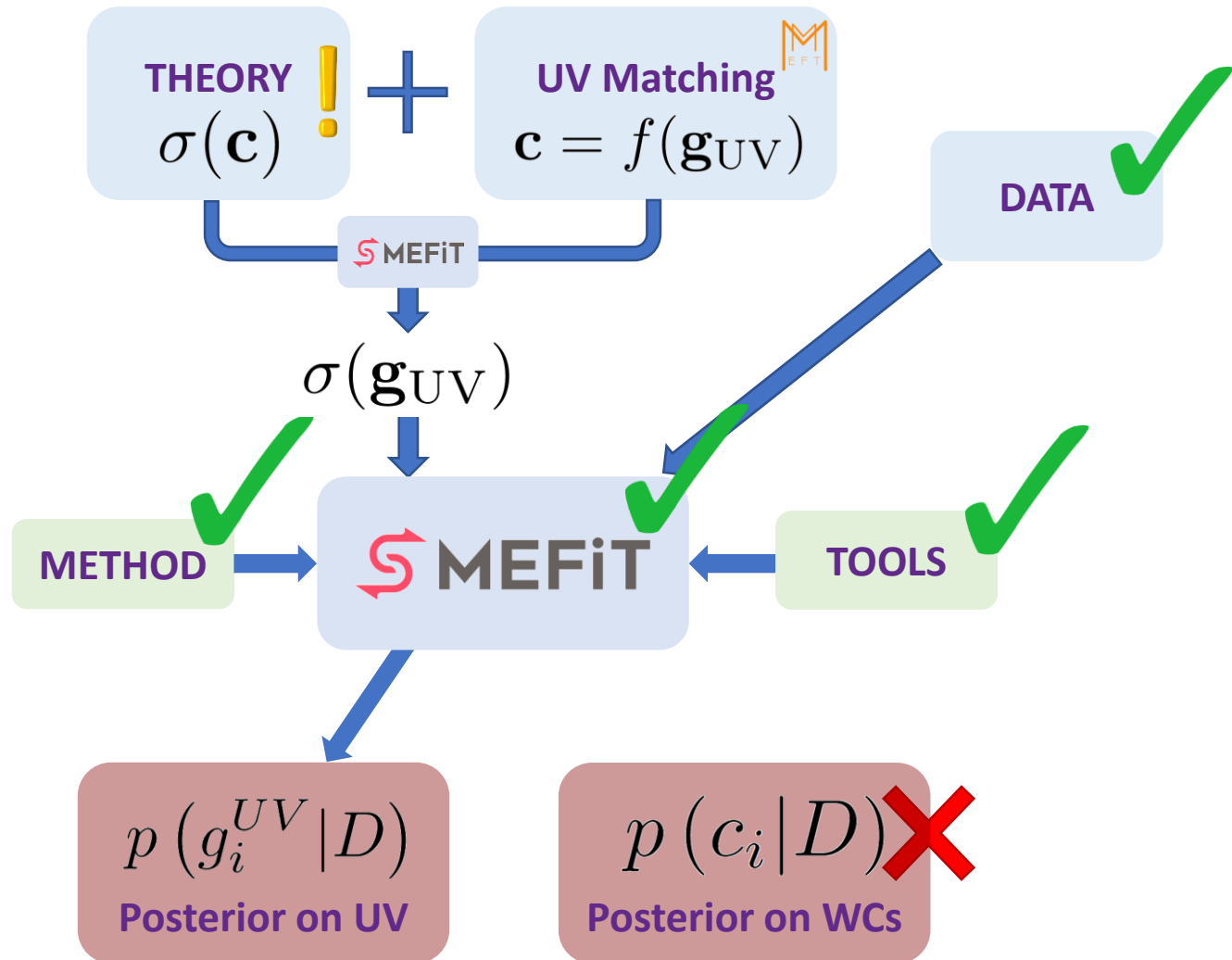
$p(c_i | D)$ ~~X~~
Posterior on WCs

Reusing EFT global fits for the UV



$p(c_i | D)$ ~~X~~
Posterior on WCs

Reusing EFT global fits for the UV





match2fit

- A Wolfram Mathematica™ package, fully documented.
- Reads results from `Matchmakereft` and produces run cards that can be fed into `smeFit` to perform a fit.
- Uses the same WC basis than SMEFiT.

$$U(2)_q \times U(3)_d \times U(2)_u \times (U(1)_\ell \times U(1)_e)^3 + c_{b\varphi}, c_{\tau\varphi}, c_{c\varphi}$$

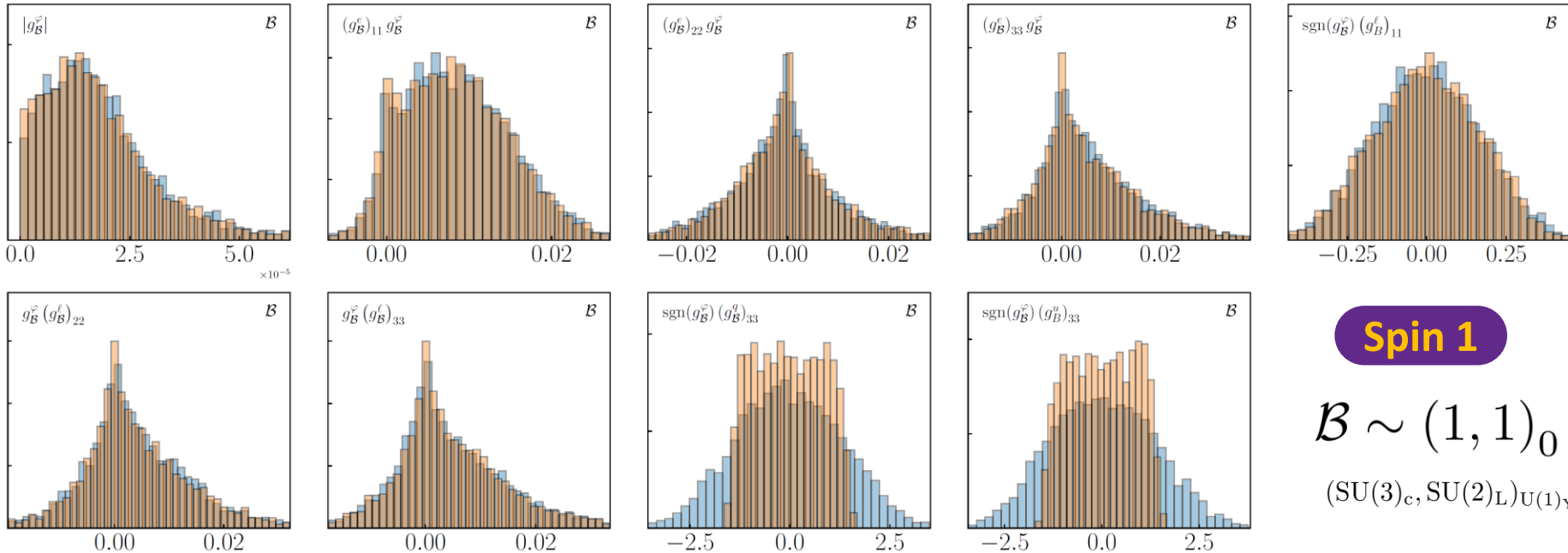
- It can impose UV flavor assumptions and evaluates the masses.
- It can run `Matchmakereft` to perform the matching and translation at once.

It supports 1-loop matching results.

One-part. models at tree level

■ NLO $\mathcal{O}(\Lambda^{-2})$
■ NLO $\mathcal{O}(\Lambda^{-4})$

$m_{\mathcal{B}} = 1 \text{ TeV}$



Spin 1

$\mathcal{B} \sim (1, 1)_0$

$(\text{SU}(3)_c, \text{SU}(2)_L)_{\text{U}(1)_Y}$

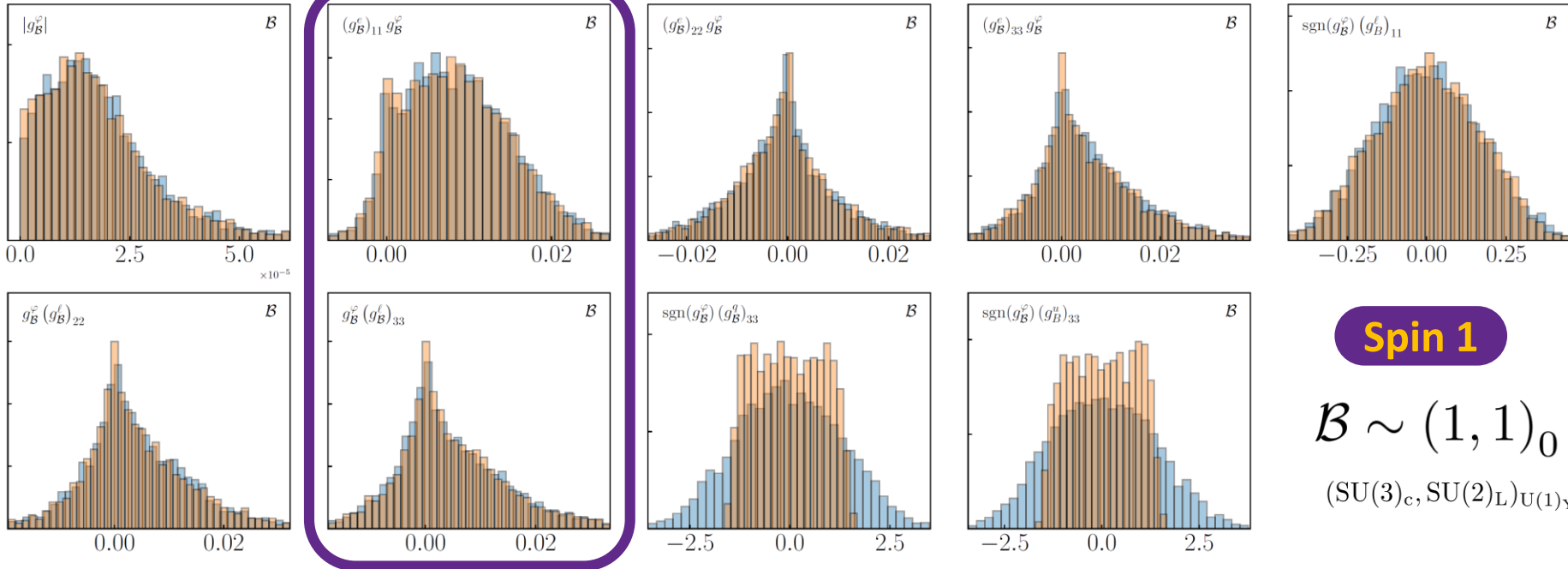
Dataset: SMEFIT 2.0 + EWPOs

Automated SMEFT-Assisted constraints on UV models | Alejo N. Rossia, 27 Sept 24

One-part. models at tree level

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Spin 1

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Sensitivity to the sign of UV couplings

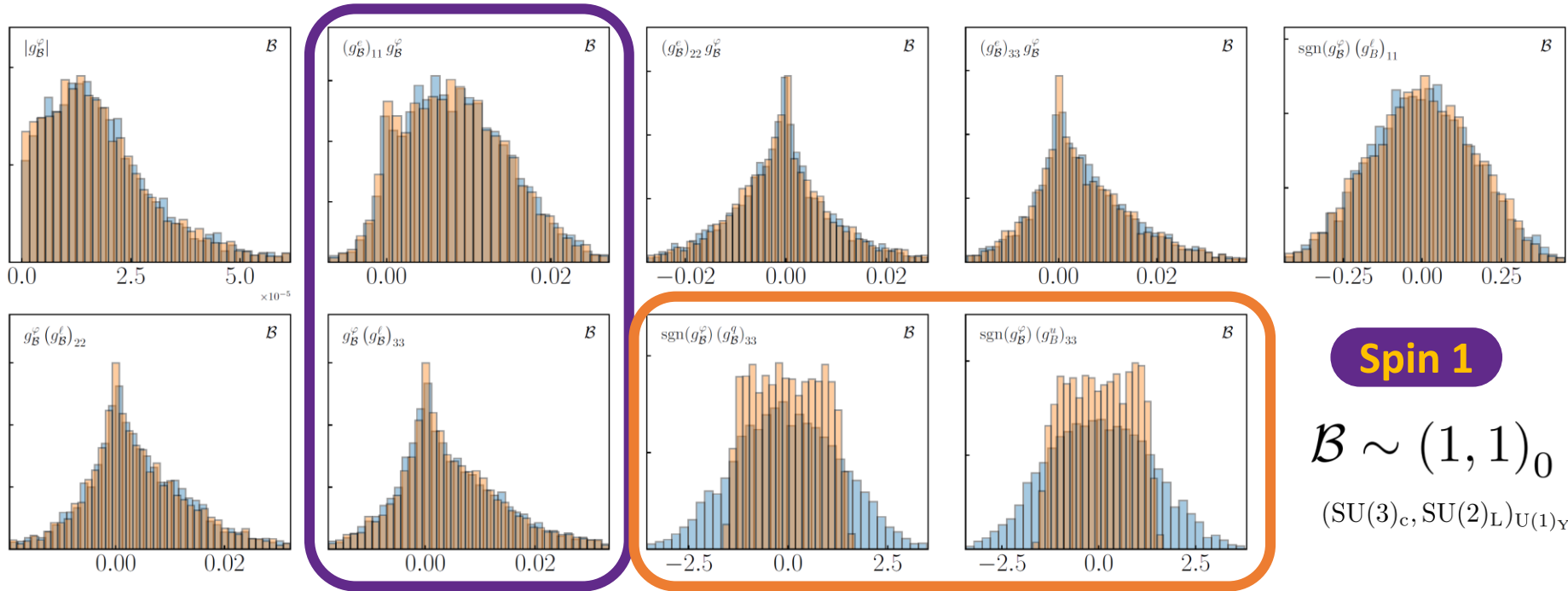
Dataset: SMEFiT 2.0 + EWPOs

Automated SMEFT-Assisted constraints on UV models | Alejo N. Rossia, 27 Sept 24

One-part. models at tree level

$m_{\mathcal{B}} = 1 \text{ TeV}$

■ NLO $\mathcal{O}(\Lambda^{-2})$ ■ NLO $\mathcal{O}(\Lambda^{-4})$



Spin 1

$\mathcal{B} \sim (1, 1)_0$

$(\text{SU}(3)_c, \text{SU}(2)_L)_{\text{U}(1)_Y}$

Sensitivity to the sign of UV couplings

Top couplings sensitive to quadratic-in-WCs pieces

Dataset: SMEFIT 2.0 + EWPOs

Automated SMEFT-Assisted constraints on UV models | Alejo N. Rossia, 27 Sept 24

One-loop matching makes a difference

$$\phi \sim (1, 2)_{1/2}$$

$$\mathcal{L}_{\text{UV}} \supset - (y_\phi^u)_{33} \phi^\dagger i \sigma_2 \bar{q}_L^T u_R^3 - \lambda_\phi \phi^\dagger H |H|^2 + \text{h.c.}$$

$$m_\phi = 1 \text{ TeV}$$

Dataset: SMEFiT 2.0 + EWPOs

Automated SMEFT-Assisted constraints on UV models | Alejo N. Rossia, 27 Sept 24

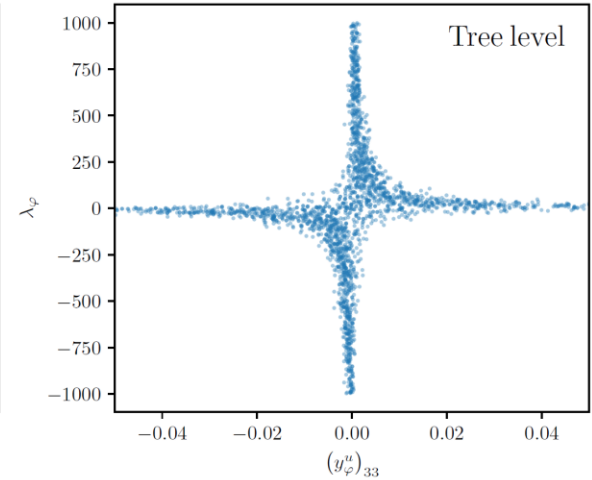
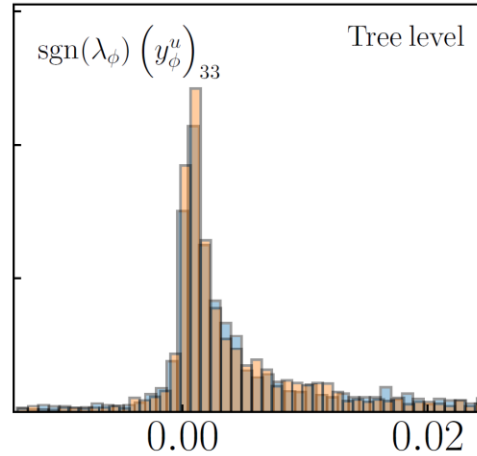
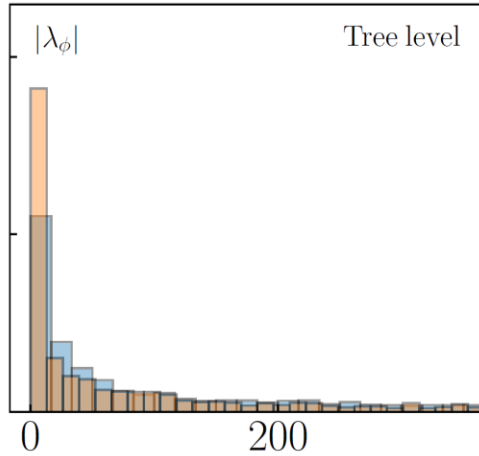
One-loop matching makes a difference

$$\phi \sim (1, 2)_{1/2}$$

$$\mathcal{L}_{\text{UV}} \supset - (y_\phi^u)_{33} \phi^\dagger i \sigma_2 \bar{q}_L^{T,3} u_R^3 - \lambda_\phi \phi^\dagger H |H|^2 + \text{h.c.}$$

$$m_\phi = 1 \text{ TeV}$$

■ NLO $\mathcal{O}(\Lambda^{-2})$
■ NLO $\mathcal{O}(\Lambda^{-4})$



Dataset: SMEFIT 2.0 + EWPOs

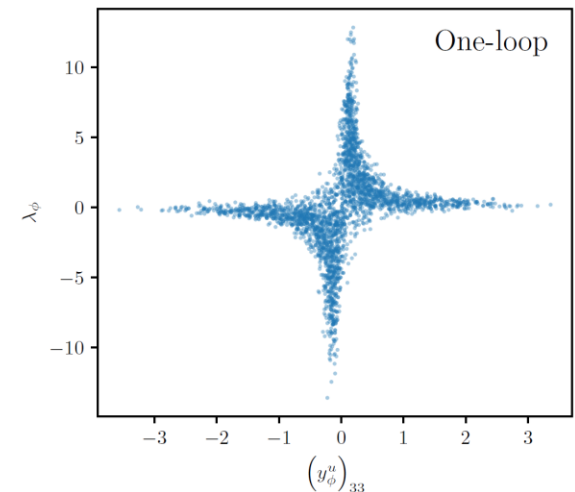
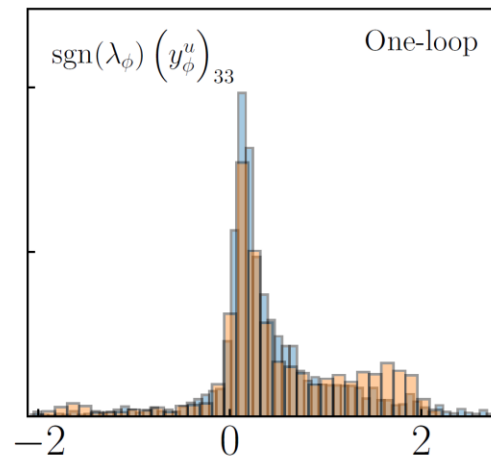
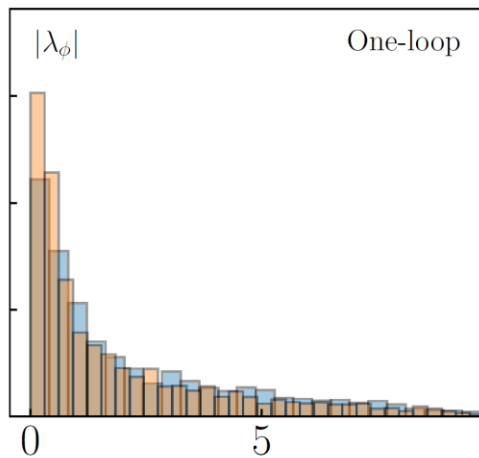
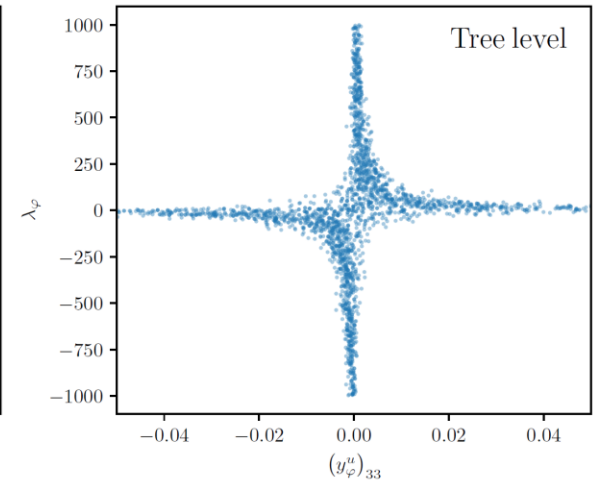
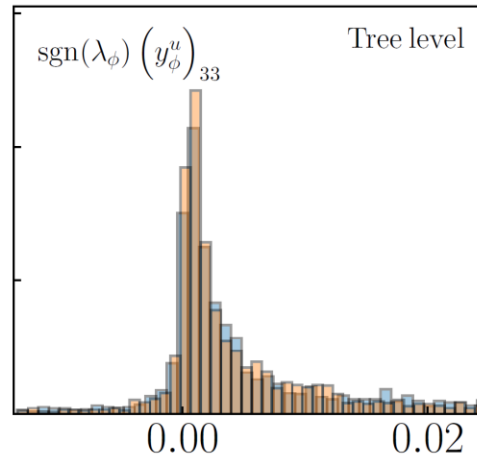
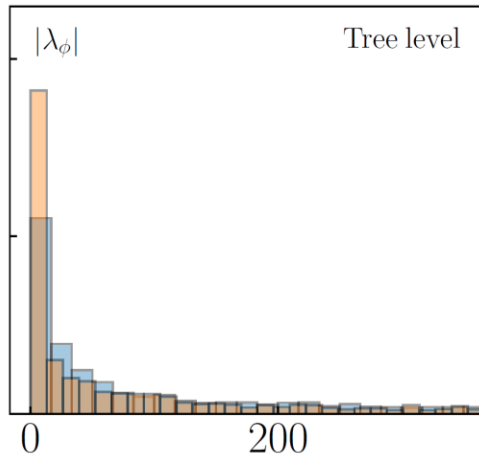
One-loop matching makes a difference

$$\phi \sim (1, 2)_{1/2}$$

$$\mathcal{L}_{\text{UV}} \supset - (y_\phi^u)_{33} \phi^\dagger i \sigma_2 \bar{q}_L^{T,3} u_R^3 - \lambda_\phi \phi^\dagger H |H|^2 + \text{h.c.}$$

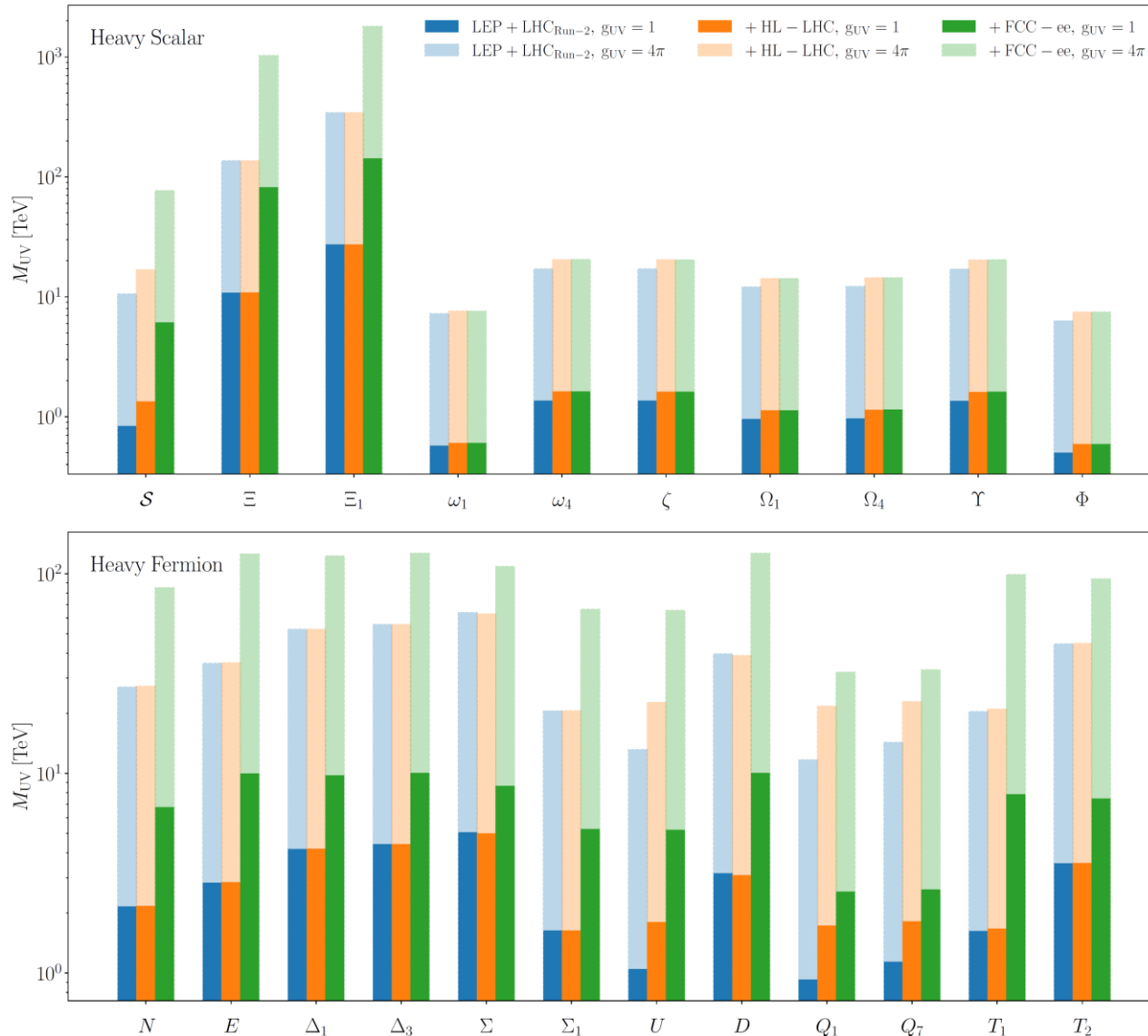
$$m_\phi = 1 \text{ TeV}$$

■ NLO $\mathcal{O}(\Lambda^{-2})$ ■ NLO $\mathcal{O}(\Lambda^{-4})$



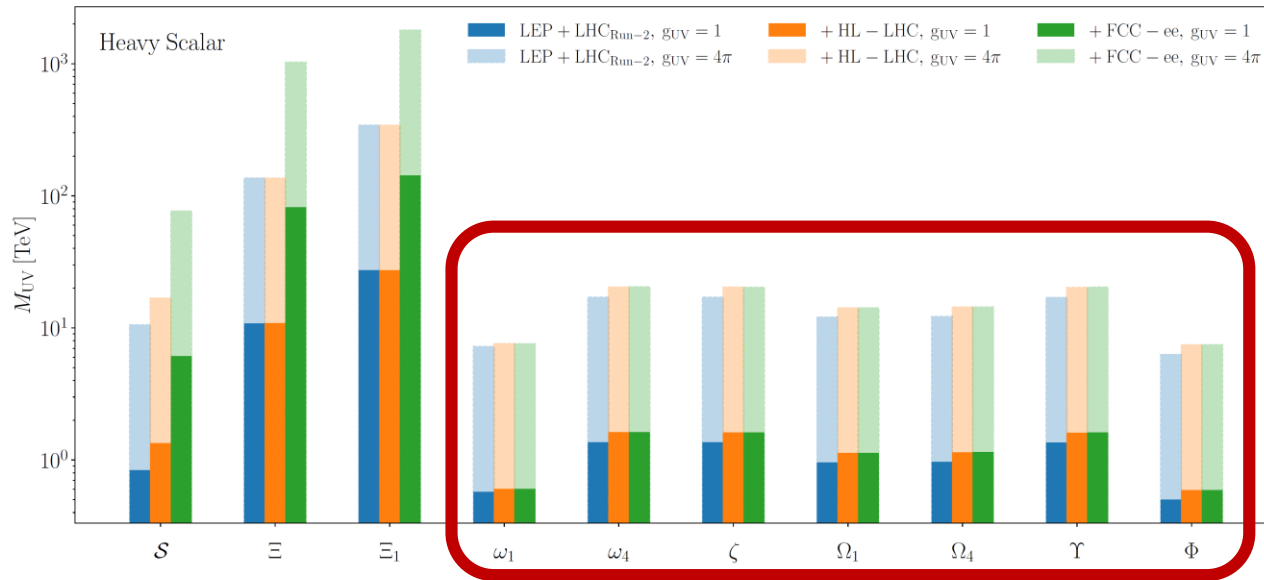
Dataset: SMEFIT 2.0 + EWPOs

Projections for the future: simple models



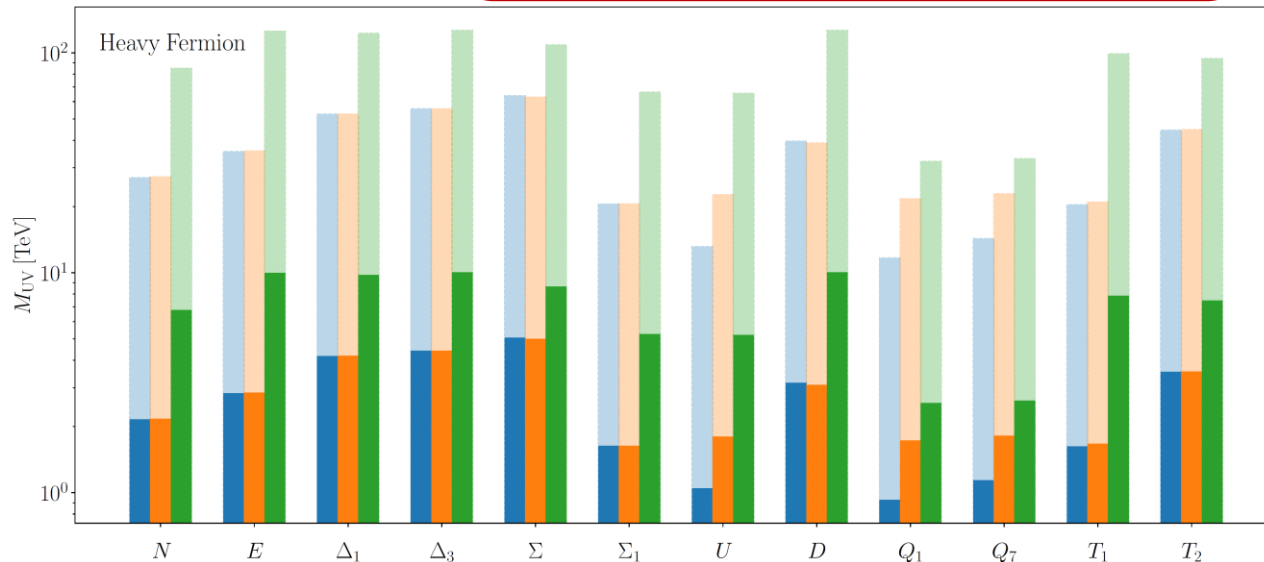
Tree-level matching

Projections for the future: simple models



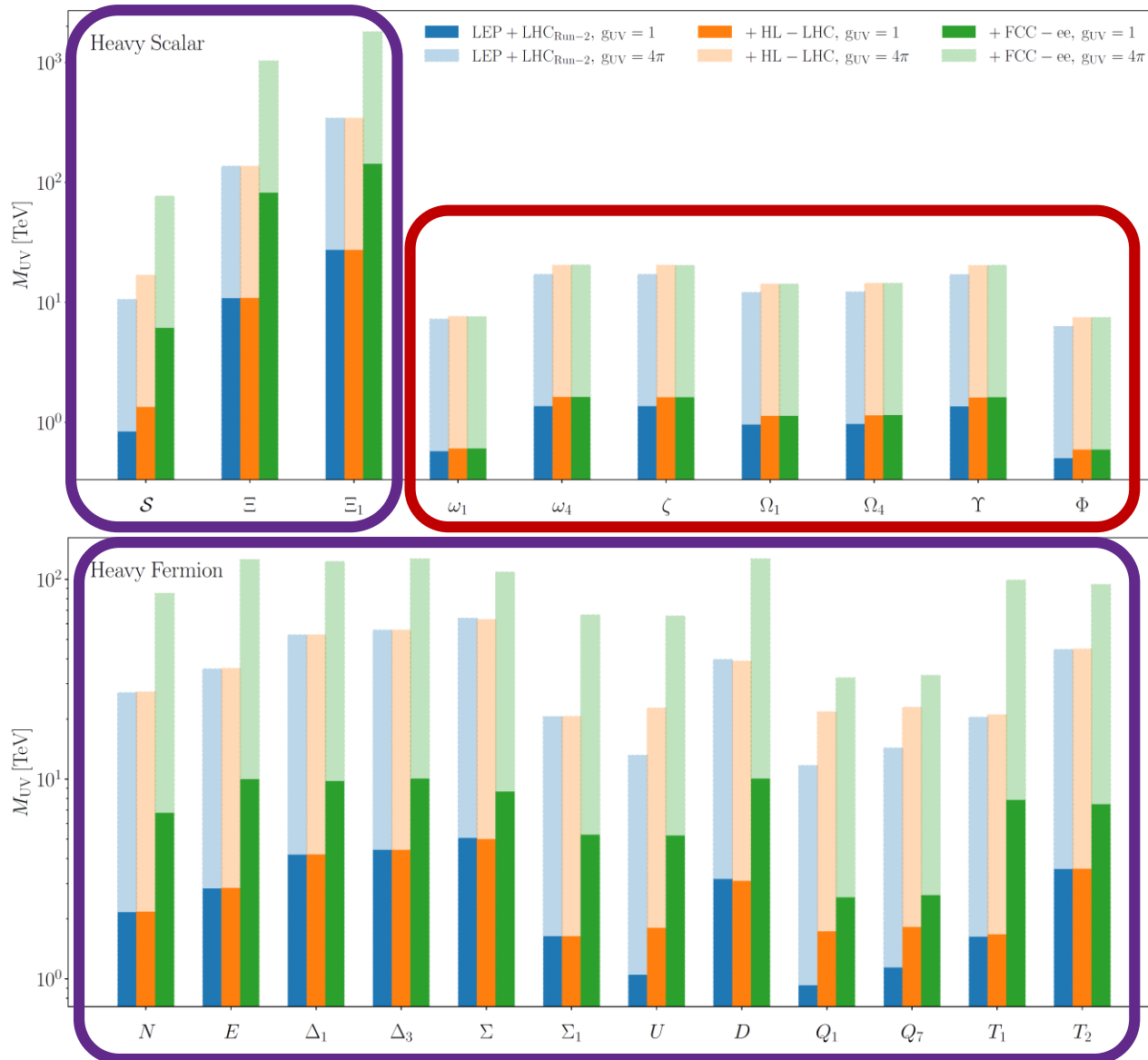
Driven by heavy 4-quark operators

Tree-level matching



Projections for the future: simple models

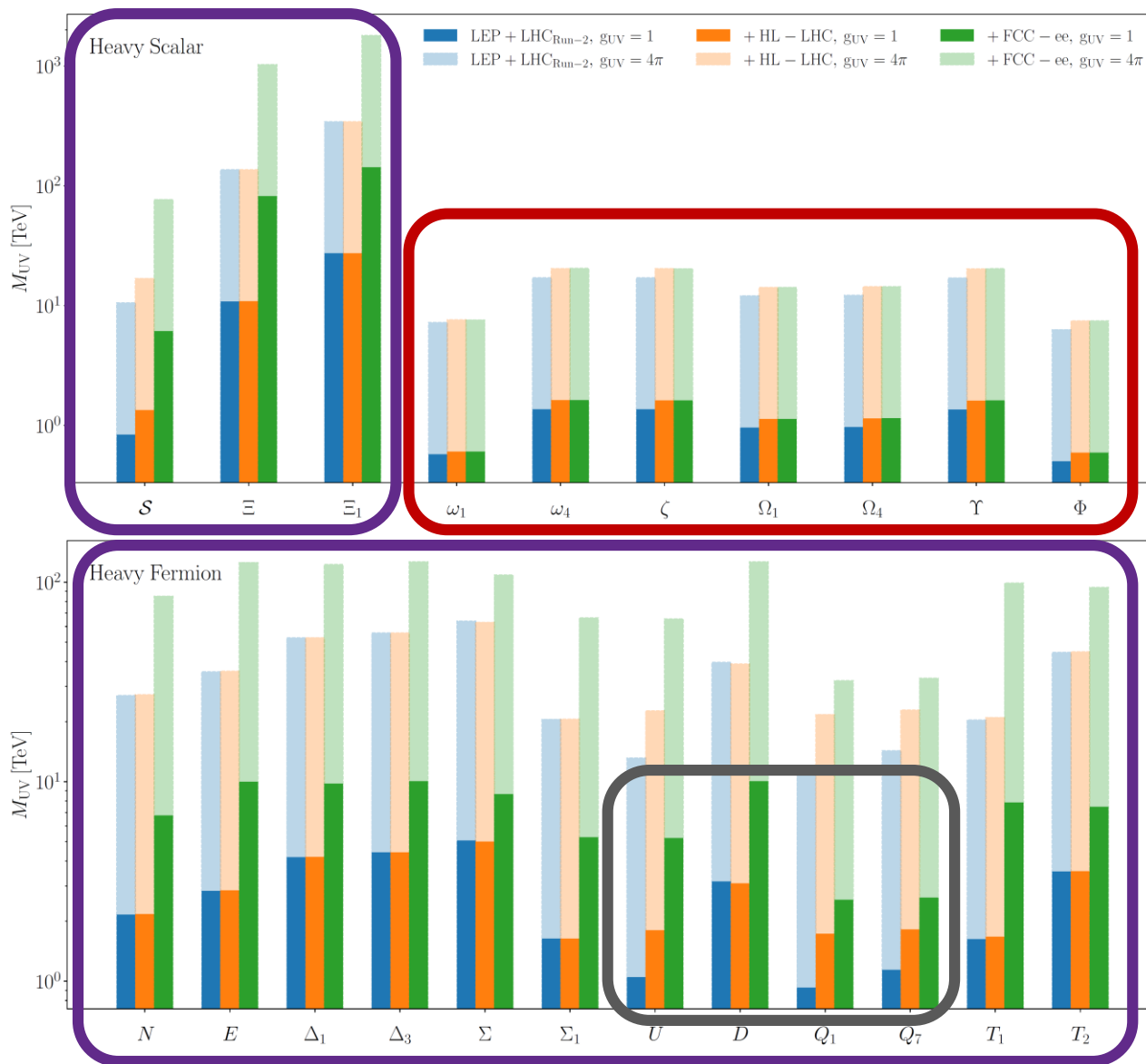
Driven by
EWPOs



Driven by
heavy
4-quark
operators

Tree-level
matching

Projections for the future: simple models



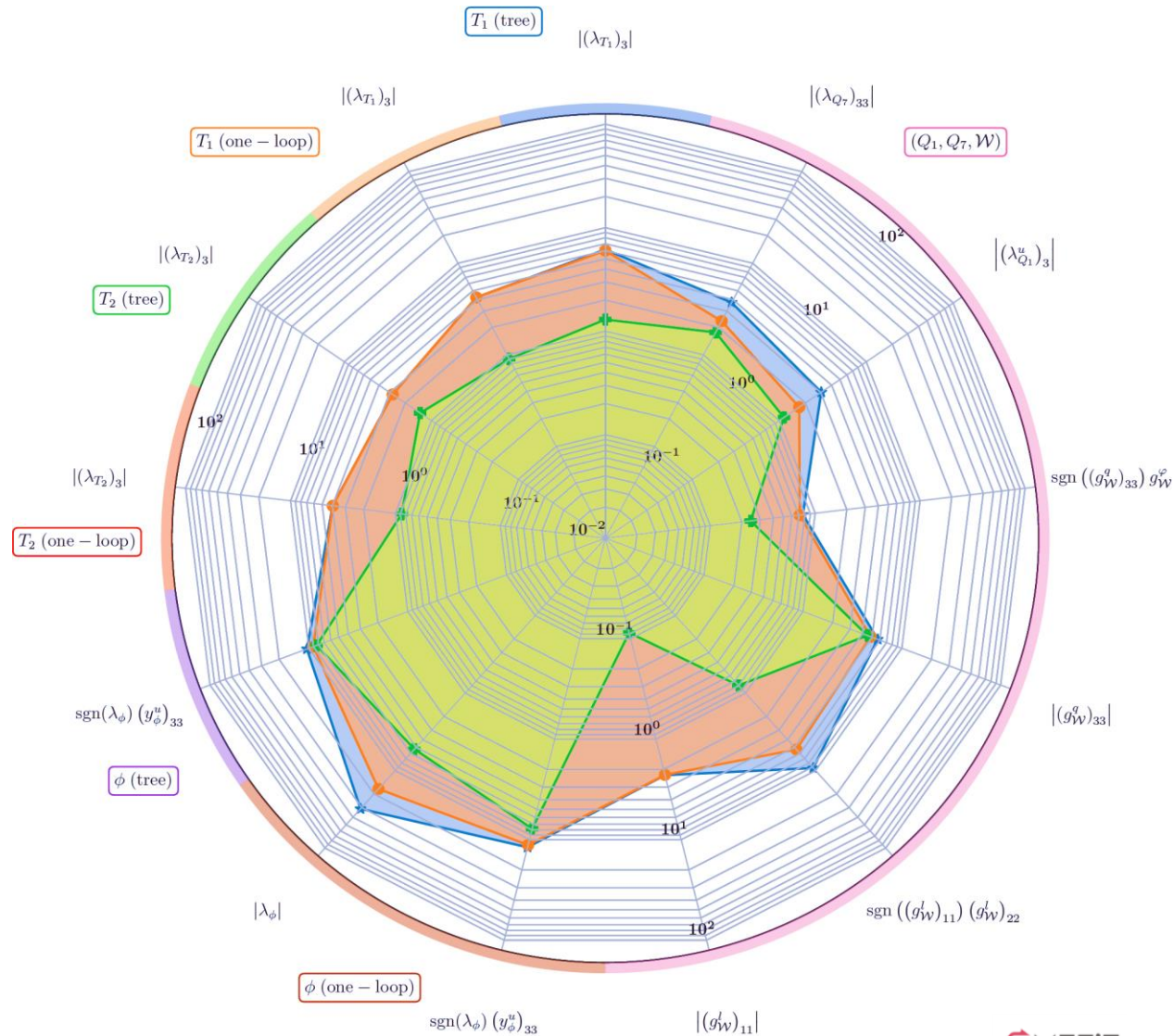
Driven by EWPOs

Top partners

Driven by heavy 4-quark operators

Tree-level matching

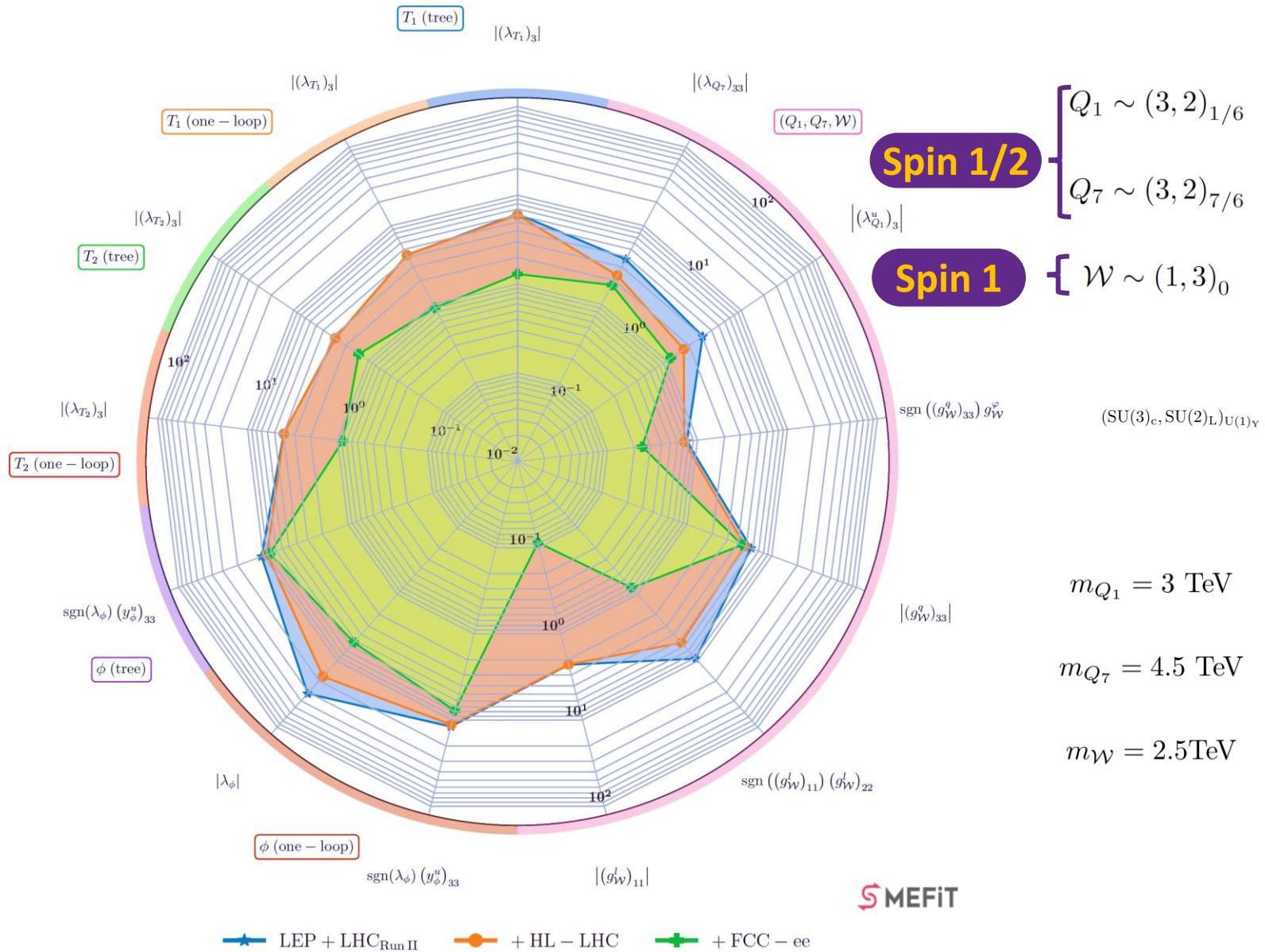
Projections for the future: complex models



★ LEP + LHC_{Run II} ● + HL-LHC + FCC-ee

Automated SMEFT-Assisted constraints on UV models | Alejo N. Rossia, 27 Sept 24

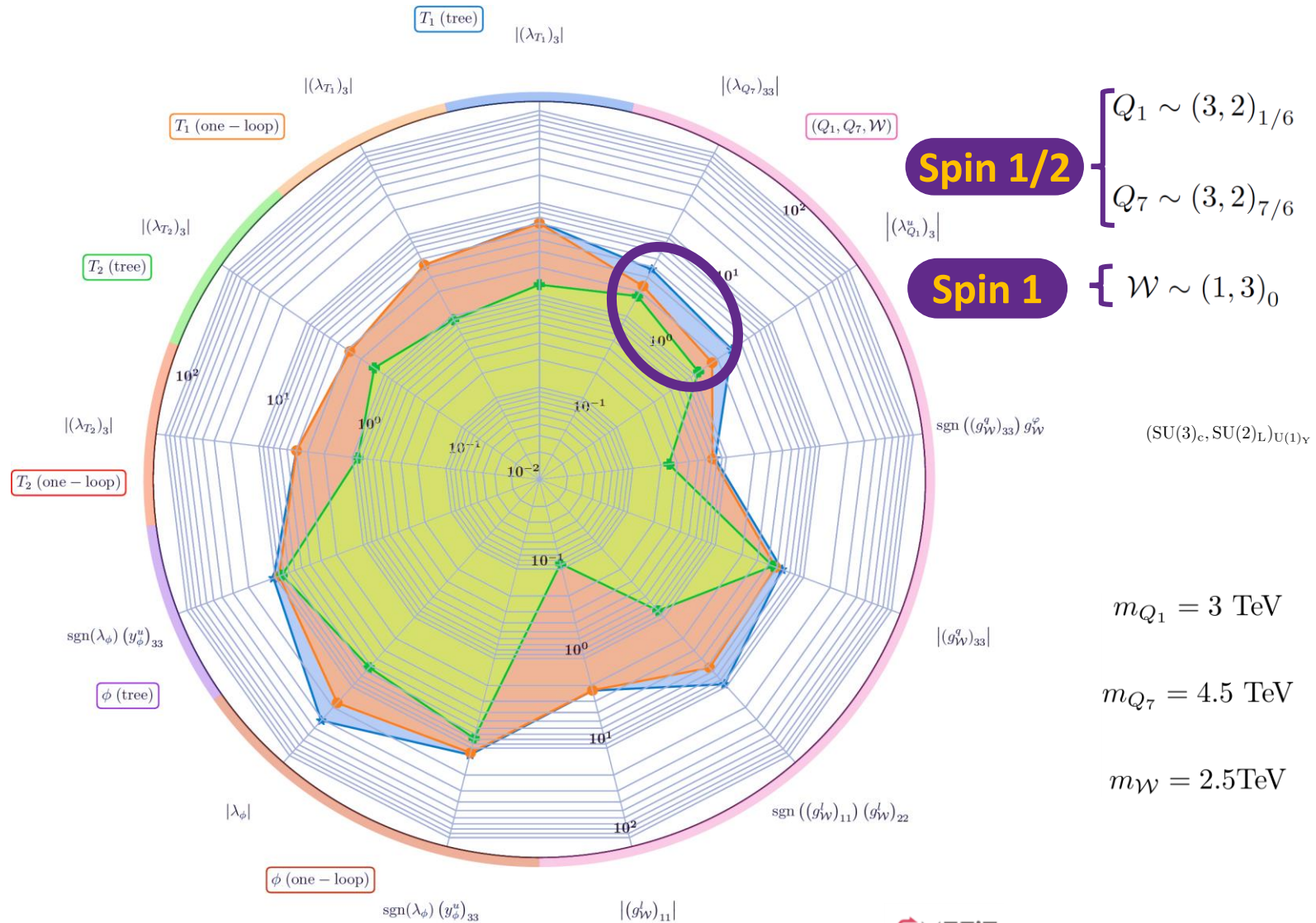
Projections for the future: complex models



SMEFIT

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Projections for the future: complex models

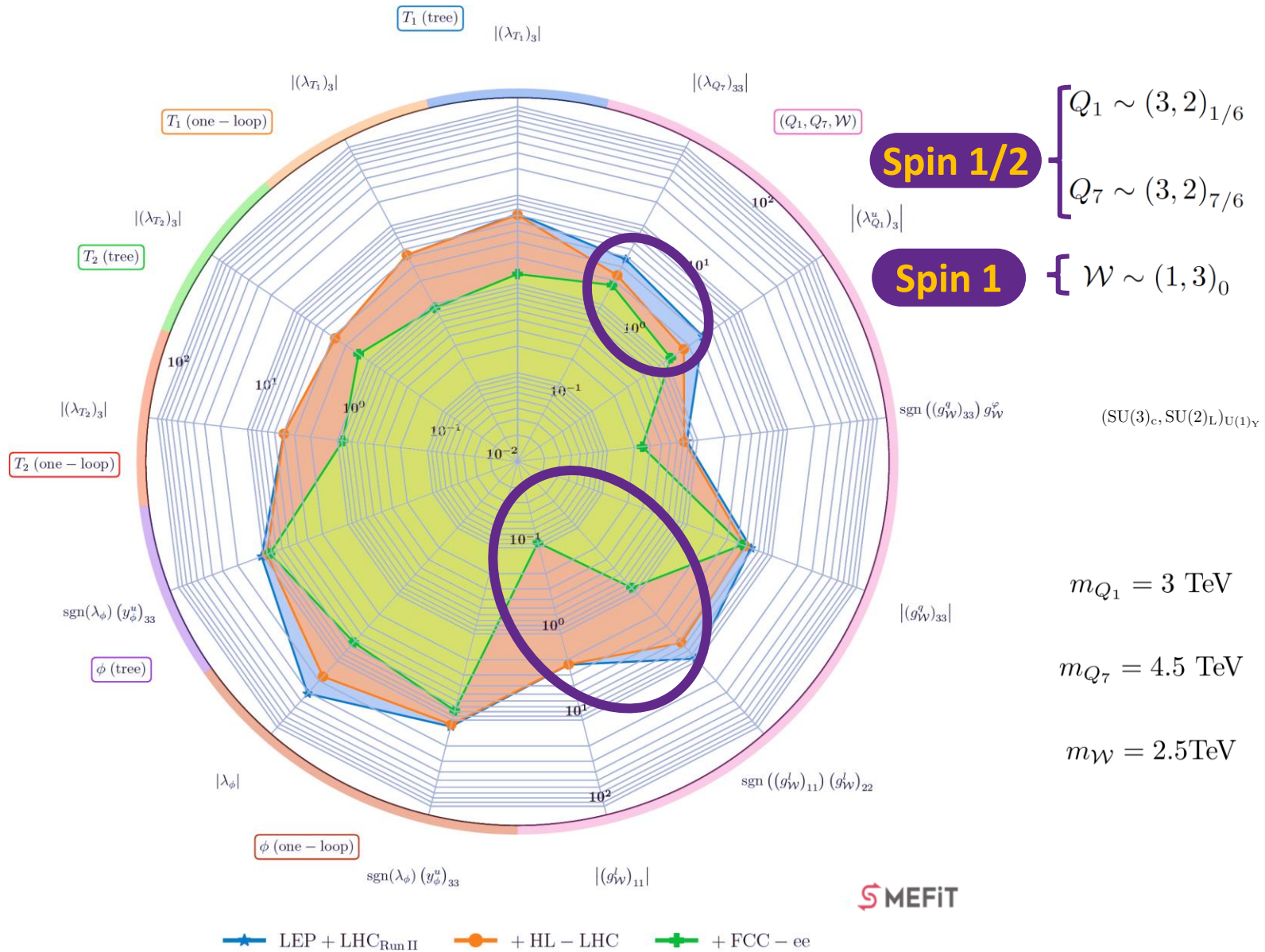


SMEFIT

★ LEP + LHC_{Run II} ● + HL - LHC + FCC - ee

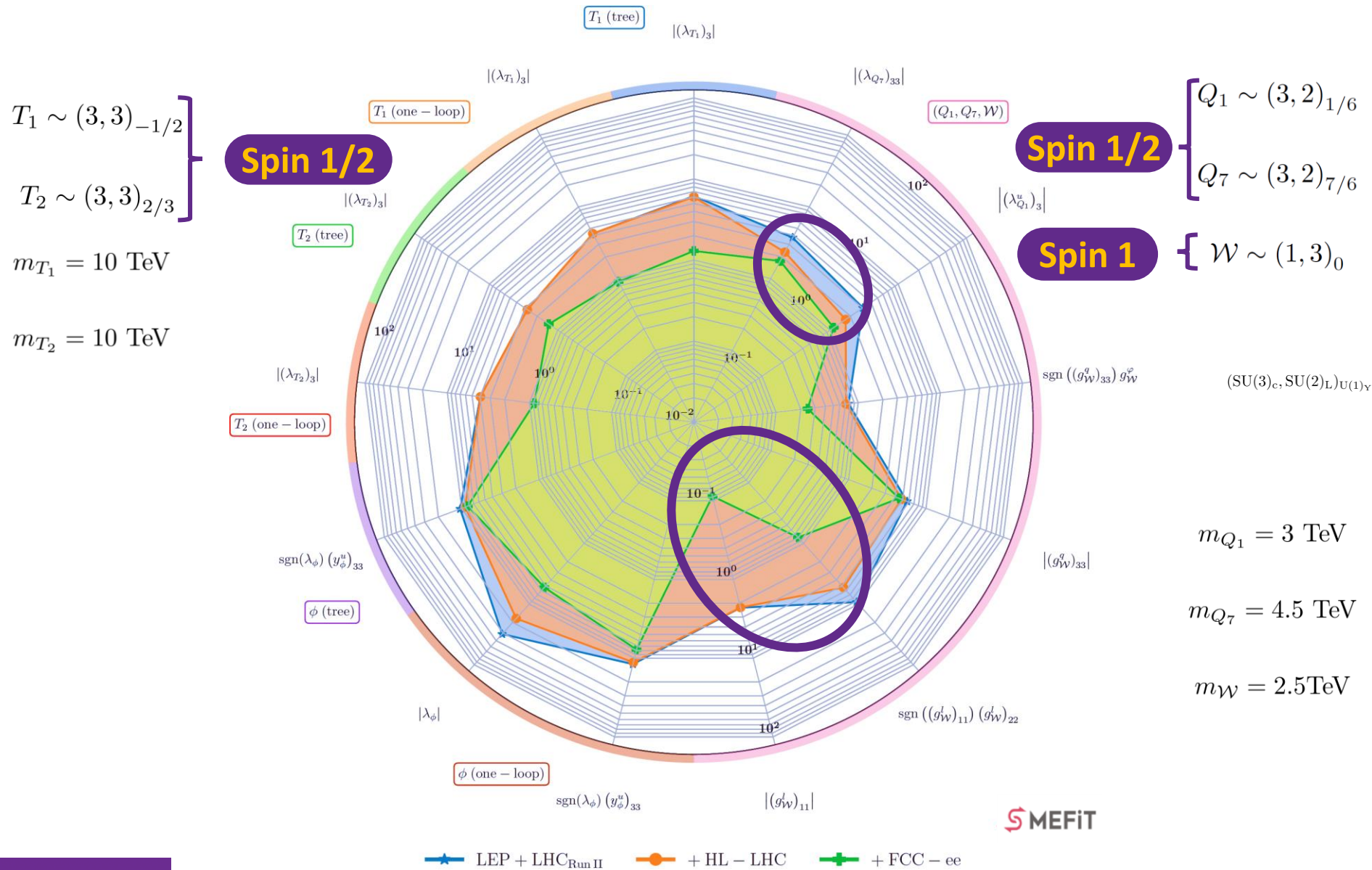
Automated SMEFT-Assisted constraints on UV models | Alejo N. Rossia, 27 Sept 24

Projections for the future: complex models

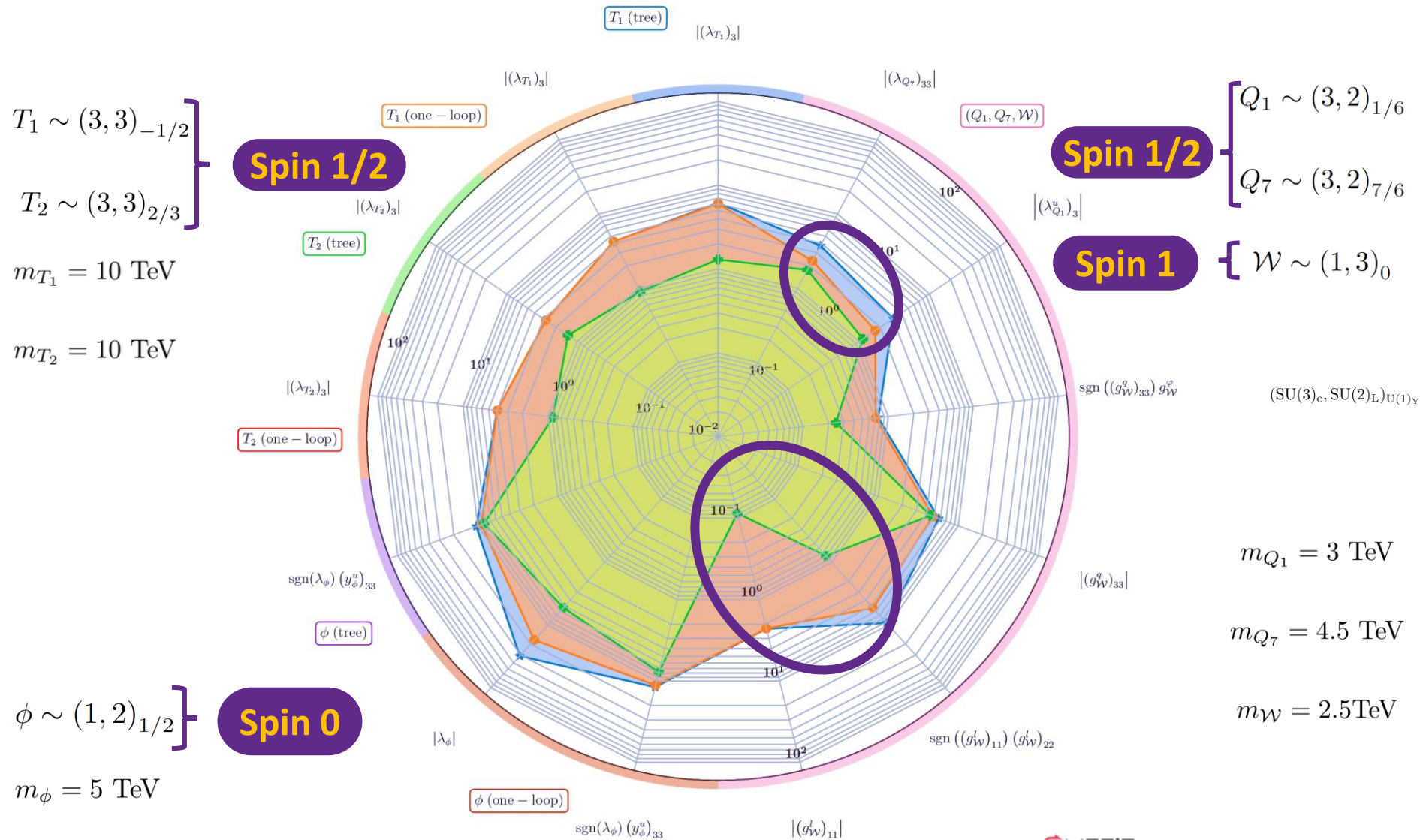


Automated SMEFT-Assisted constraints on UV models | Alejo N. Rossia, 27 Sept 24

Projections for the future: complex models



Projections for the future: complex models

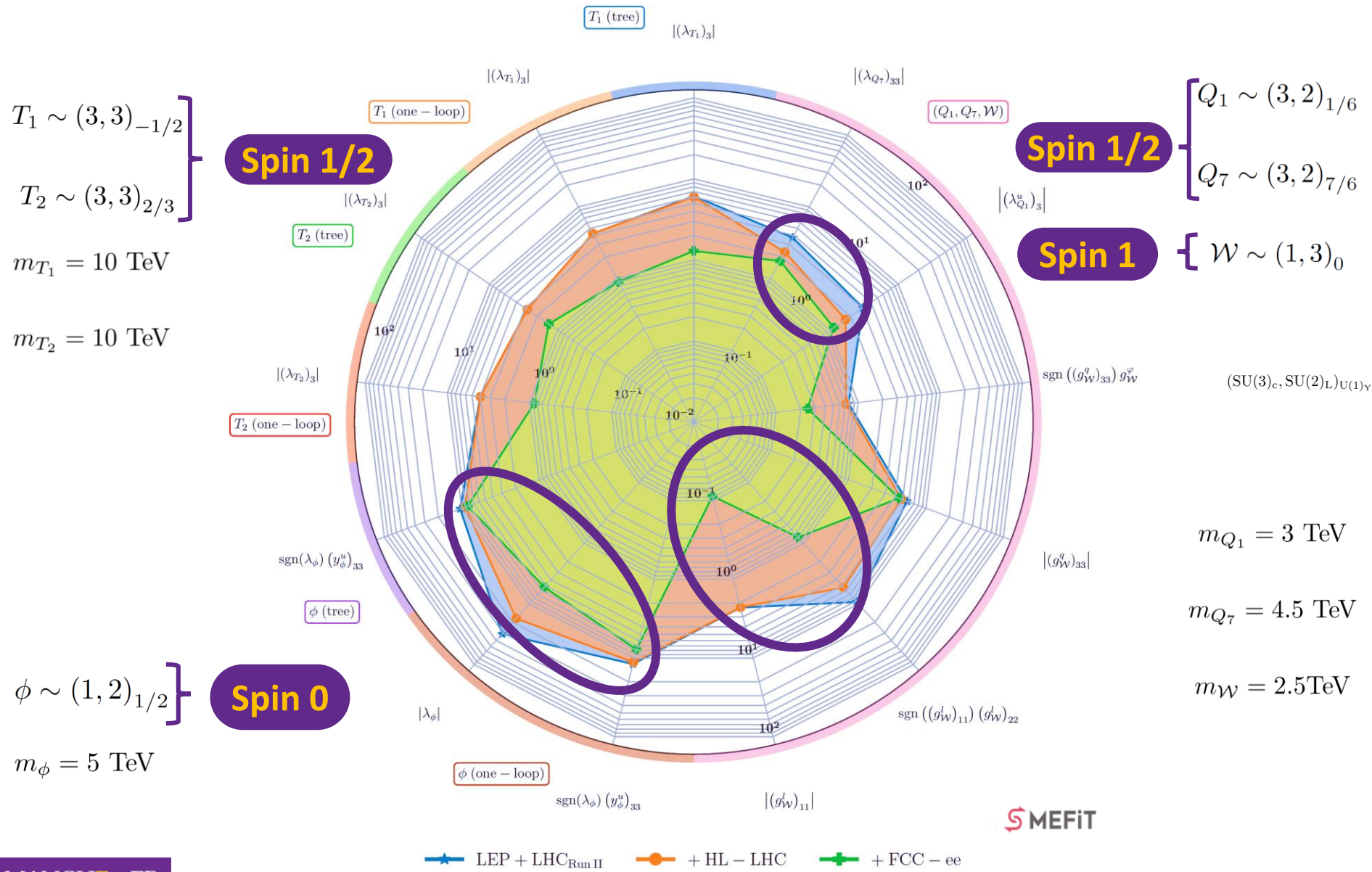


SMEFIT

★ LEP + LHC_{Run II} ● + HL - LHC + FCC - ee

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Projections for the future: complex models



SMEFIT

Automated SMEFT-Assisted constraints on UV models | Alejo N. Rossia, 27 Sept 24

Conclusions

- We have the tools for the full cycle of the EFT program for BSM Physics.
- SMEFiT allows to interpret LHC data at the EFT and UV model levels from one set of predictions.
- Match2fit provides a simple and flexible SMEFiT-MMEFT interface.
- LHC Run 2 data shows an impressive constrain power.
- We can understand the impact of future colliders at SMEFT and UV level.
- Several improvement possibilities: interfacing more codes, flavor data, RGE effects, more general flavor symmetries...

Thanks for your attention!

Contact:

Alejo N. Rossia

HEP Theory Group – Dept. Of Physics and Astronomy

E-mail: alejo dot rossia at manchester dot ac dot uk

Thanks to M. Thomas, E. Celada, V. Miralles and H. el Faham for ideas for the slides and discussions.

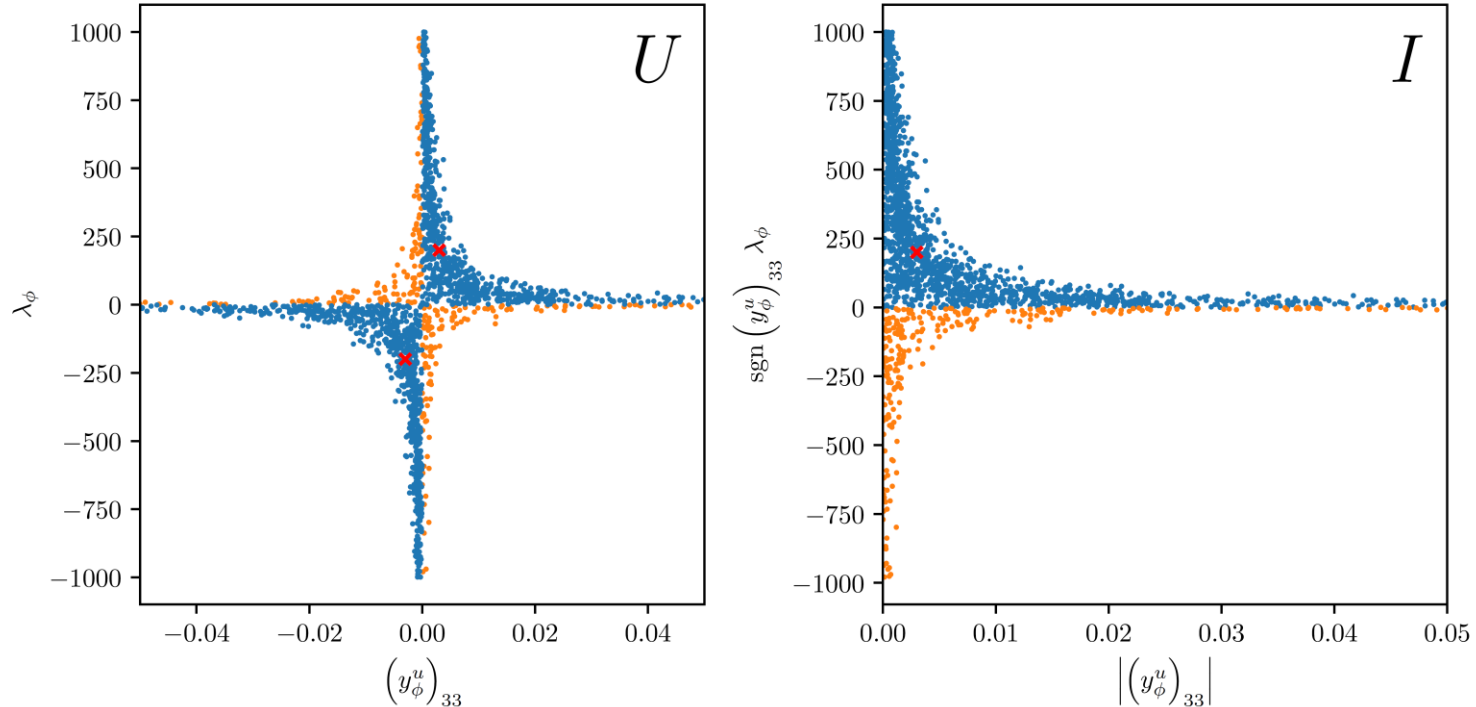
Appendix

UV invariants

We are sensitive only to combinations of UV couplings that enter the WCs.

$$h : U \rightarrow I \quad \text{“UV invariants”} \quad \mathbf{c} = f(\mathbf{g}_{\text{UV}})$$

$$f(h(g)) = f(h(g')) \iff h(g) = h(g') \quad \mathbf{c} = f(h(\mathbf{g}_{\text{UV}}))$$



Not necessary to do the fit, but useful to understand the results.

Restrictions from EFT flavor symmetry

- Your model produces an operator that should vanish and does not enter in any fitted process.
 - The bounds from the fit might be suboptimal with respect to bounds from other processes.
- Your model produces an operator that should vanish and enters some processes in the dataset.
 - The bounds from the fit might not be trustworthy and suboptimal.
- The symmetry assumes two WCs to be equal but your model produces them with different values.
 - Match2fit will take only one of those values and ignore the other. Unless the difference is small, the bounds from the fit are not trustworthy.

How to forecast



(3 ab⁻¹)



- Take SMEFIT 3.0 LHC datasets with highest int. luminosity
- Pseudodata fluctuated around SM
- Rescale uncertainties:
 - Statistical $\longrightarrow \mathcal{L}$
 - Systematics $\longrightarrow \frac{1}{2}$
- No HL-LHC optimization
- Snowmass + FCC midterm Feas. Rep.
- Z-pole+161+240+350/365 GeV
 - EWPOs
 - $f\bar{f}$ production
 - $ZH + \nu\bar{\nu}H + \text{all } H \text{ decays.}$
 - $W^+W^- + t\bar{t}$ with Optim. Obs.

Additional technicalities

SMEFIT supports relations among fit parameters like:

$$\sum_i a_i (c_1)^{n_{1,i}} \dots (c_N)^{n_{N,i}} = 0$$

The exponents can be rational numbers of any sign.
This imposes restrictions on the supported matching relations.

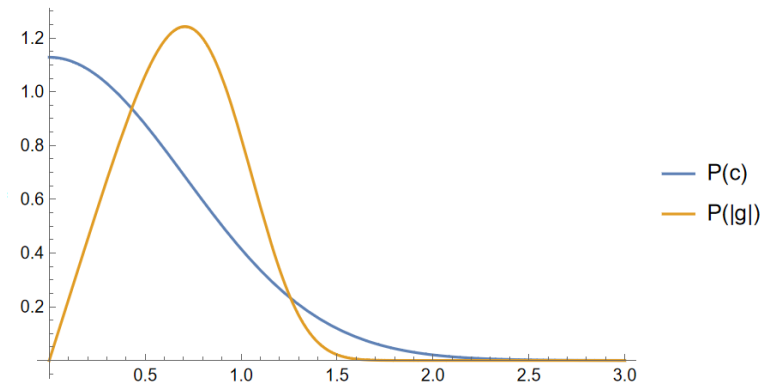
Probability in UV and WC spaces

The relation between PDFs in WC and UV space can be misleading.

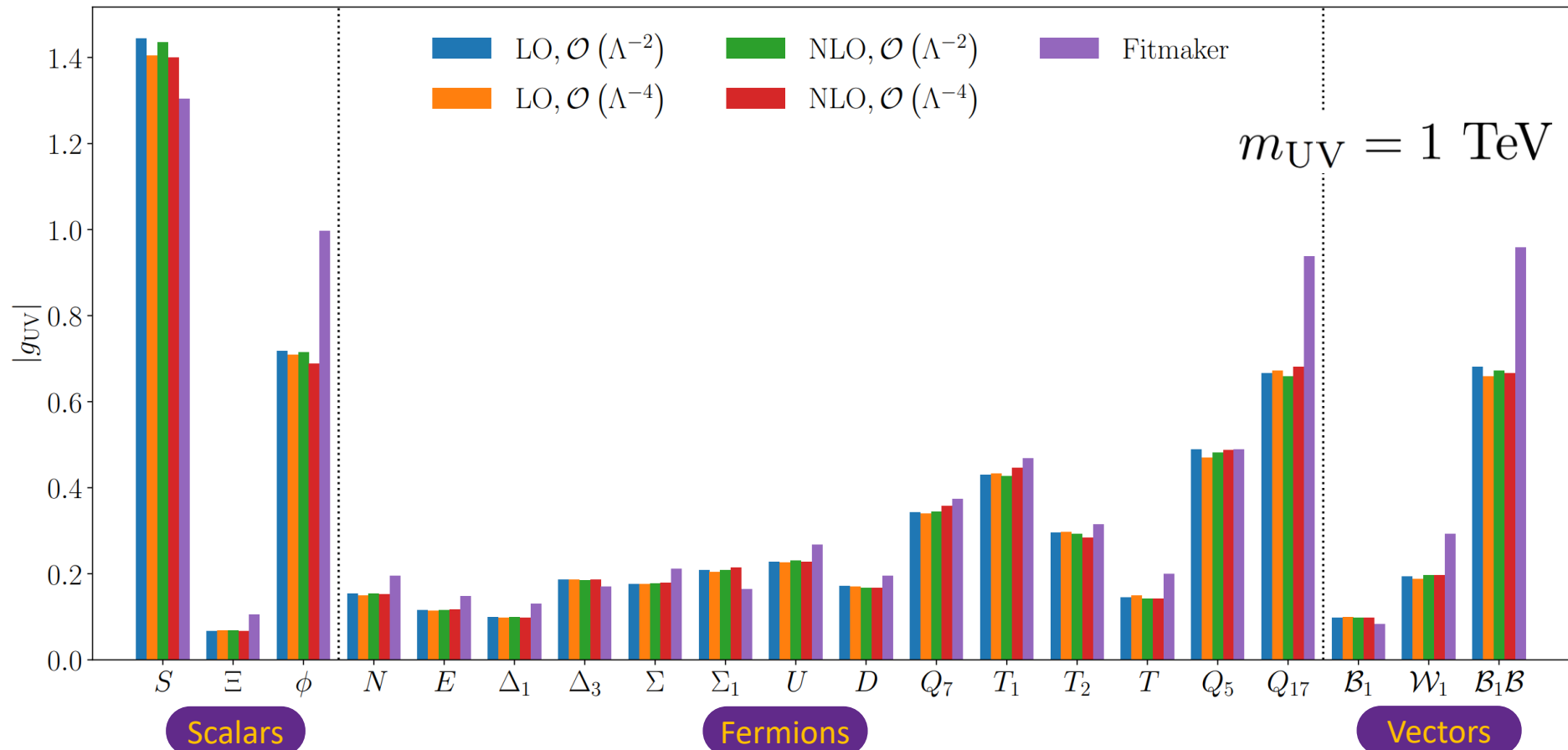
$$P(c) = \frac{2}{\sqrt{\pi}} e^{-c^2}, \quad \int_0^{\infty} dc P(c) = 1$$

$$c = g^2$$

$$P(|g|) = \frac{4}{\sqrt{\pi}} |g| e^{-|g|^4}, \quad \int_0^{\infty} d|g| P(|g|) = 1$$



One-part. models at tree level

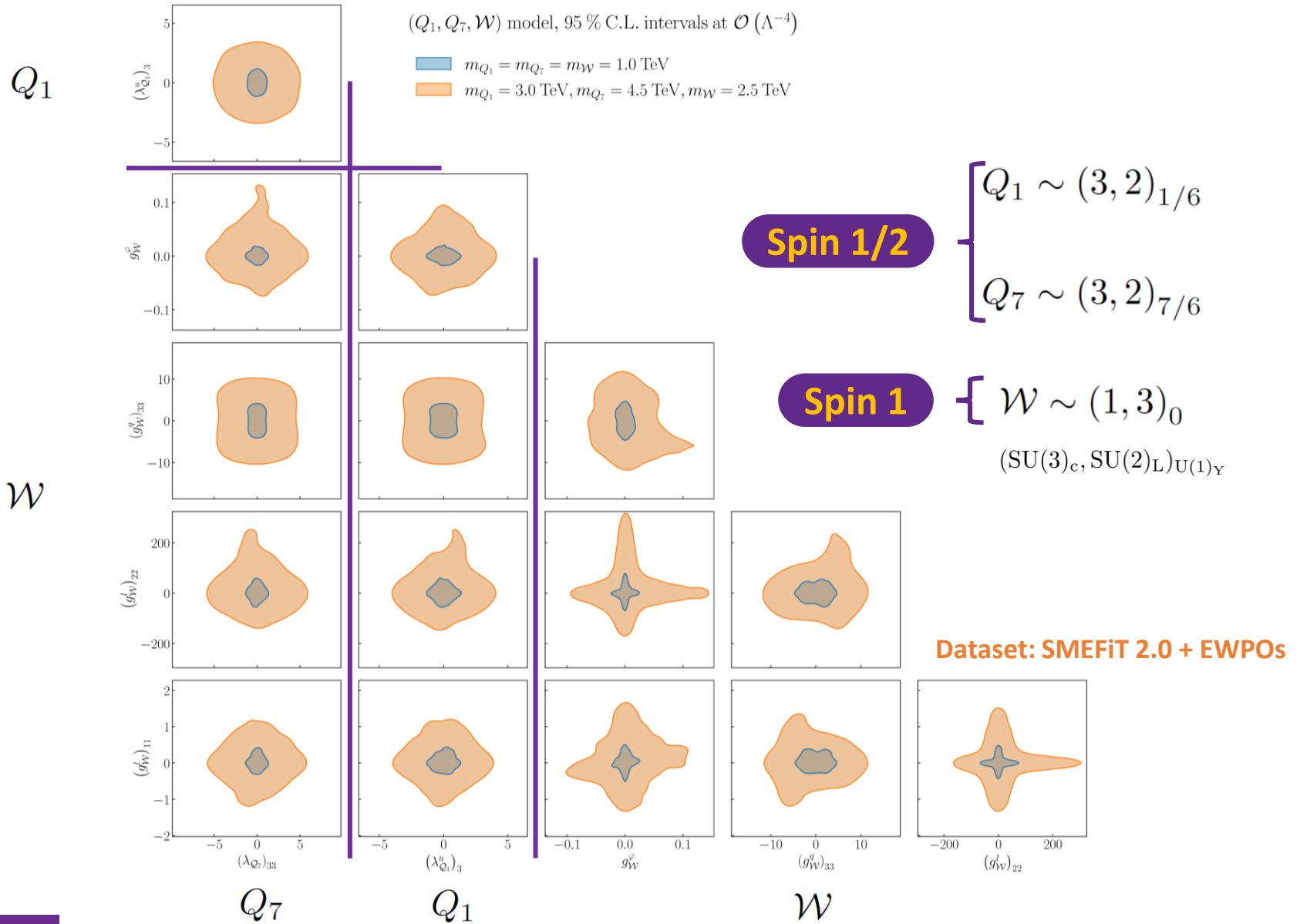


Good agreement with the Fitmaker results

Dataset: SMEFIT 2.0 + EWPOs

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Multi-particle models at tree level



List of models

Scalars		Fermions		Vectors	
Particle	Irrep	Particle	Irrep	Particle	Irrep
\mathcal{S}	$(1, 1)_0$	N	$(1, 1)_0$	\mathcal{B}	$(1, 1)_0$
\mathcal{S}_1	$(1, 1)_1$	E	$(1, 1)_{-1}$	\mathcal{B}_1	$(1, 1)_1$
ϕ	$(1, 2)_{1/2}$	Δ_1	$(1, 2)_{-1/2}$	\mathcal{W}	$(1, 3)_0$
Ξ	$(1, 3)_0$	Δ_3	$(1, 2)_{-3/2}$	\mathcal{W}_1	$(1, 3)_1$
Ξ_1	$(1, 3)_1$	Σ	$(1, 3)_0$	\mathcal{G}	$(8, 1)_0$
ω_1	$(3, 1)_{-1/3}$	Σ_1	$(1, 3)_{-1}$	\mathcal{H}	$(8, 3)_0$
ω_4	$(3, 1)_{-4/3}$	U	$(3, 1)_{2/3}$	\mathcal{Q}_5	$(8, 3)_0$
ζ	$(3, 3)_{-1/3}$	D	$(3, 1)_{-1/3}$	\mathcal{Y}_5	$(\bar{6}, 2)_{-5/6}$
Ω_1	$(6, 1)_{1/3}$	Q_1	$(3, 2)_{1/6}$		
Ω_4	$(6, 1)_{4/3}$	Q_7	$(3, 2)_{7/6}$		
Υ	$(6, 3)_{1/3}$	T_1	$(3, 3)_{-1/3}$		
Φ	$(8, 2)_{1/2}$	T_2	$(3, 3)_{2/3}$		
		Q_5	$(3, 2)_{-5/6}$		

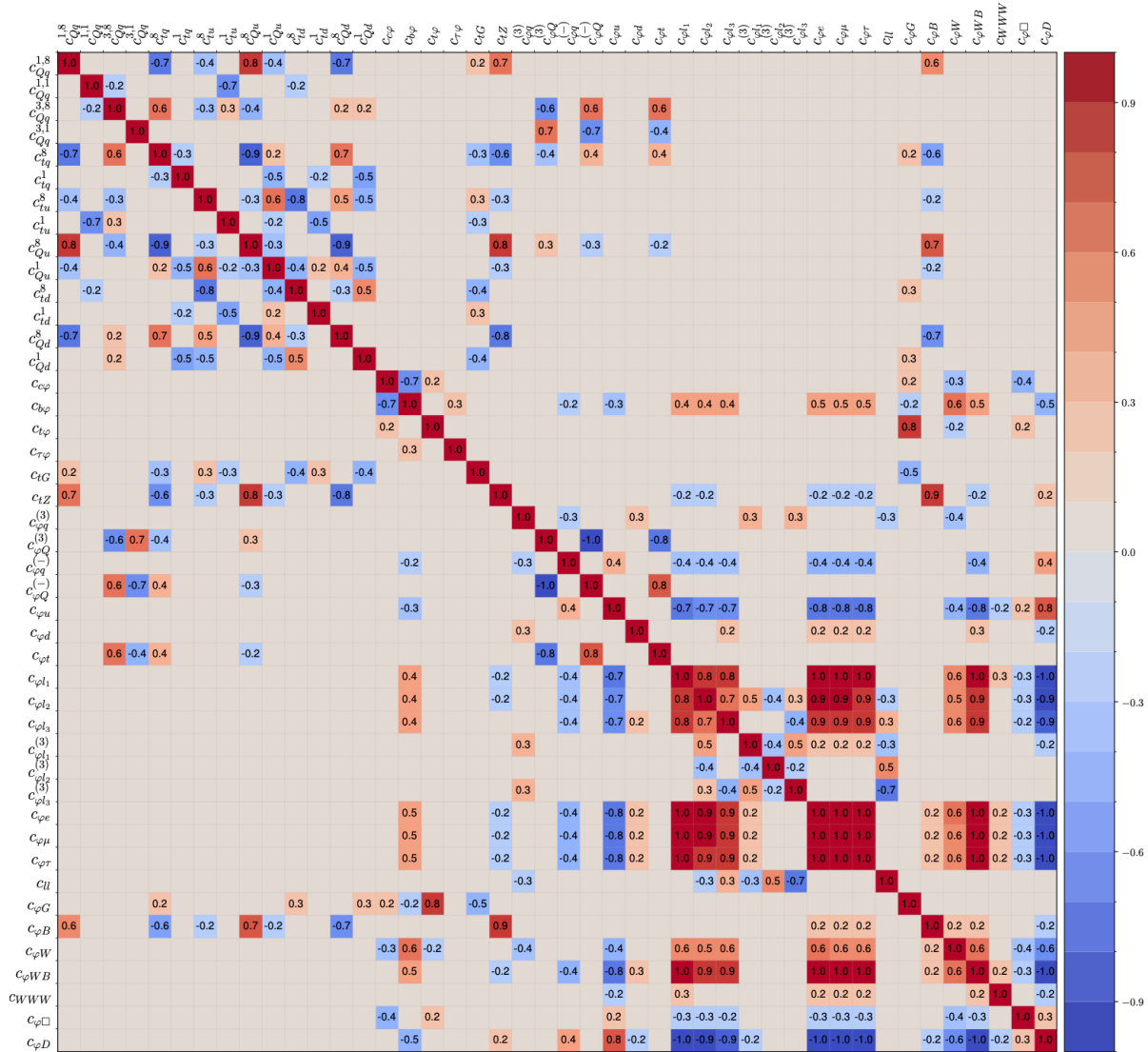
UV Couplings

Scalars		Fermions		Vectors	
Model	UV couplings	Model	UV couplings	Model	UV couplings
\mathcal{S}	$\kappa_{\mathcal{S}}$	N	$(\lambda_N^e)_3$	\mathcal{B}	$(g_B^u)_{33}, (g_B^q)_{33}, g_B^\varphi,$
ϕ	$\lambda_\phi, (y_\phi^u)_{33}$	E	$(\lambda_E)_3$		$(g_B^e)_{11}, (g_B^e)_{22}, (g_B^e)_{33},$
Ξ	κ_Ξ	Δ_1	$(\lambda_{\Delta_1})_3$		$(g_B^\ell)_{22}, (g_B^\ell)_{33}$
Ξ_1	κ_{Ξ_1}	Δ_3	$(\lambda_{\Delta_3})_3$	\mathcal{B}_1	$g_{\mathcal{B}_1}^\varphi$
ω_1	$(y_{\omega_1}^{qq})_{33}$	Σ	$(\lambda_\Sigma)_3$	\mathcal{W}	$(g_{\mathcal{W}}^l)_{11} = 2 (g_{\mathcal{W}}^l)_{22}, (g_{\mathcal{W}}^l)_{33}$
ω_4	$(y_{\omega_4}^{uu})_{33}$	Σ_1	$(\lambda_{\Sigma_1})_3$		$g_{\mathcal{W}}^\varphi, (g_{\mathcal{W}}^q)_{33}$
ζ	$(y_\zeta^{qq})_{33}$	U	$(\lambda_U)_3$	\mathcal{W}_1	$g_{\mathcal{W}_1}^\varphi$
Ω_1	$(y_{\Omega_1}^{qq})_{33}$	D	$(\lambda_D)_3$	\mathcal{G}	$(g_{\mathcal{G}}^q)_{33}, (g_{\mathcal{G}}^u)_{33}$
Ω_4	$(y_{\Omega_4})_{33}$	Q_1	$(\lambda_{Q_1}^u)_3$		
Υ	$(y_\Upsilon)_{33}$	Q_7	$(\lambda_{Q_7})_3$	\mathcal{H}	$(g_{\mathcal{H}})_{33}$
Φ	$(y_\Phi^{qu})_{33}$	T_1	$(\lambda_{T_1})_3$	\mathcal{Q}_5	$(g_{\mathcal{Q}_5}^{uq})_{33}$
		T_2	$(\lambda_{T_2})_3$	\mathcal{Y}_5	$(g_{\mathcal{Y}_5})_{33}$

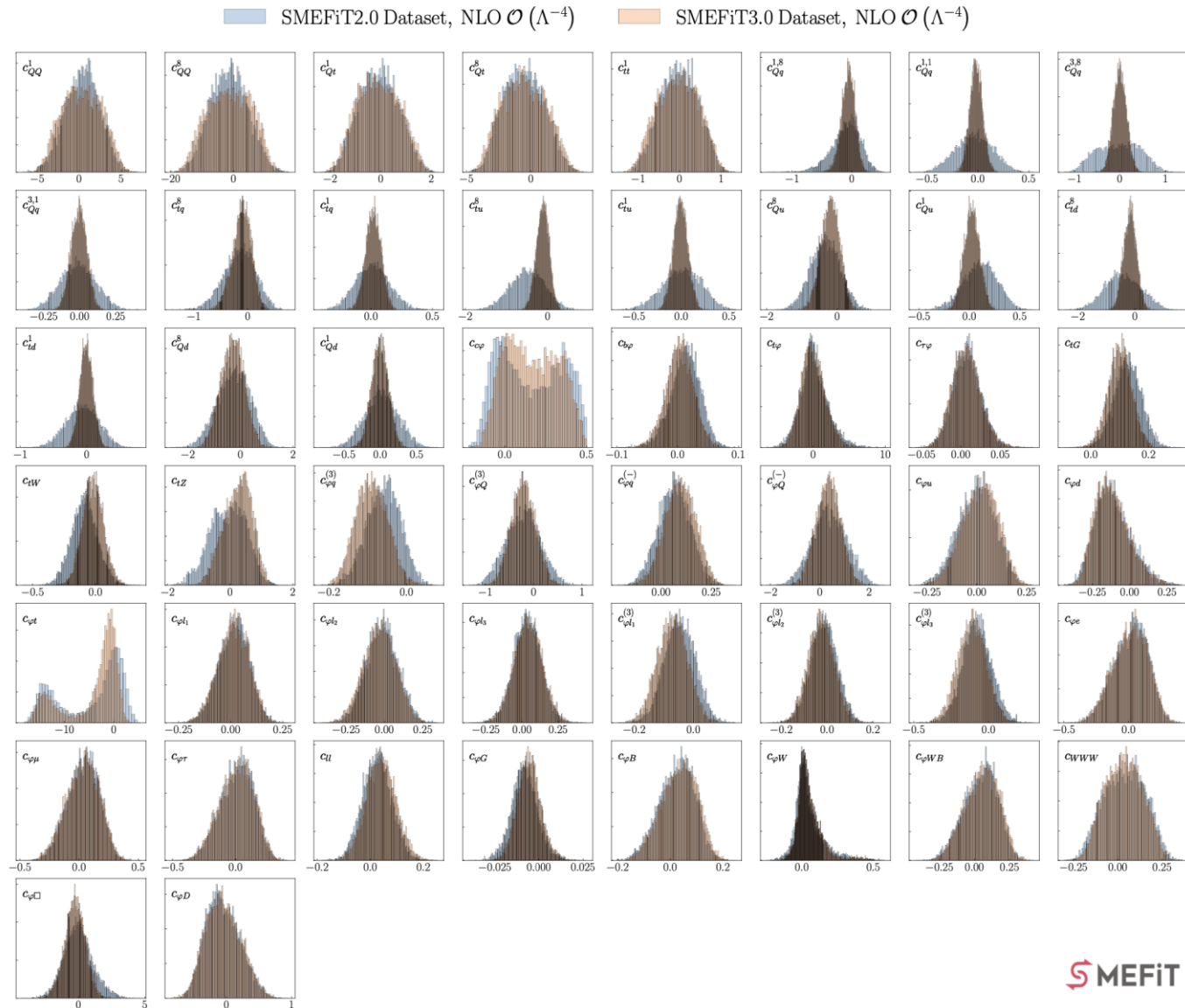
Dataset

Category	Processes	n_{dat}	
		SMEFT2.0	SMEFT3.0
Top quark production	$t\bar{t} + X$	94	115
	$t\bar{t}Z, t\bar{t}W$	14	21
	$t\bar{t}\gamma$	-	2
	single top (inclusive)	27	28
	tZ, tW	9	13
	$t\bar{t}t\bar{t}, t\bar{t}b\bar{b}$	6	12
	Total	150	191
Higgs production and decay	Run I signal strengths	22	22
	Run II signal strengths	40	36 (*)
	Run II, differential distributions & STXS	35	71
	Total	97	129
Diboson production	LEP-2	40	40
	LHC	30	41
	Total	70	81
EWPOs	LEP-2	-	44
Baseline dataset	Total	317	445

Correlations in linear fit



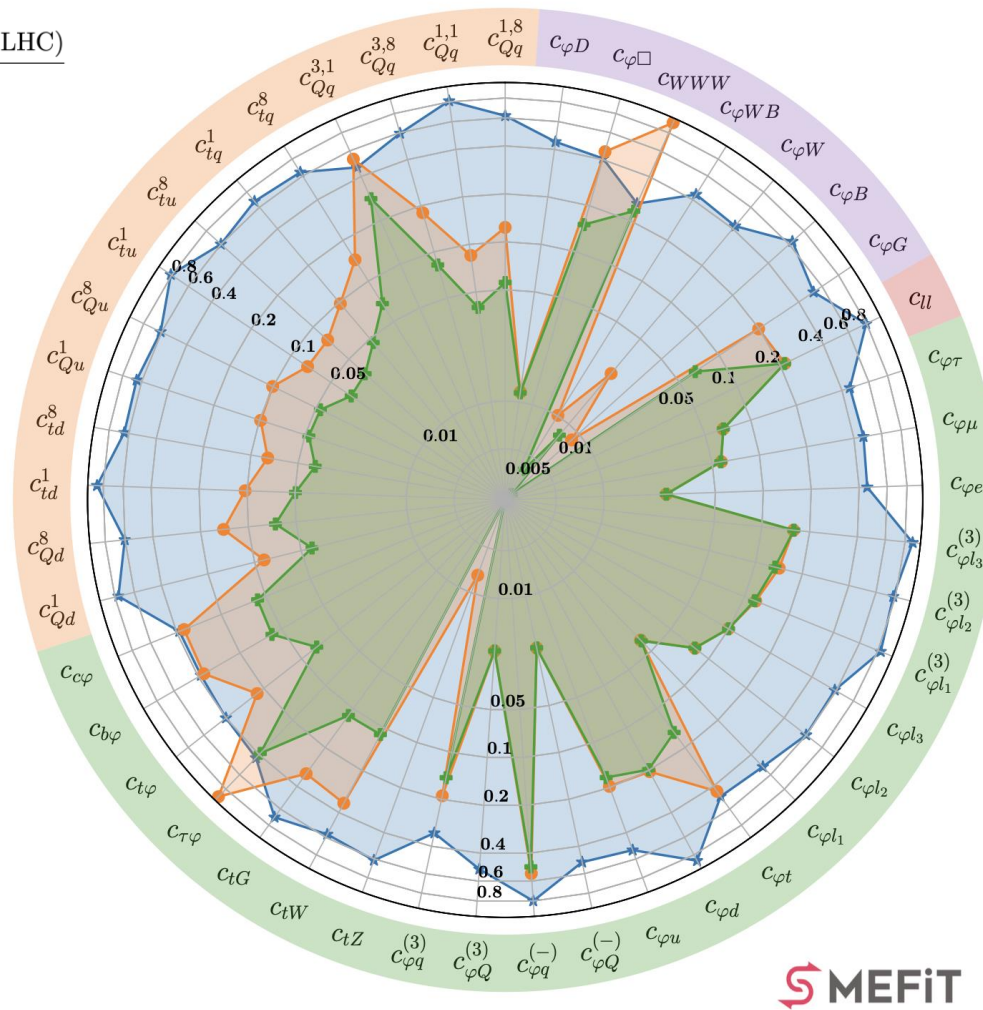
Impact of new LHC Run 2 data



HL-LHC impact in detail

Ratio of Uncertainties to SMEFiT3.0 Baseline, $\mathcal{O}(\Lambda^{-2})$, Marginalised

$$R_{\delta c_i} = \frac{[c_i^{\min}, c_i^{\max}]^{95\% \text{ CL}} (\text{baseline} + \text{HL-LHC})}{[c_i^{\min}, c_i^{\max}]^{95\% \text{ CL}} (\text{baseline})}$$

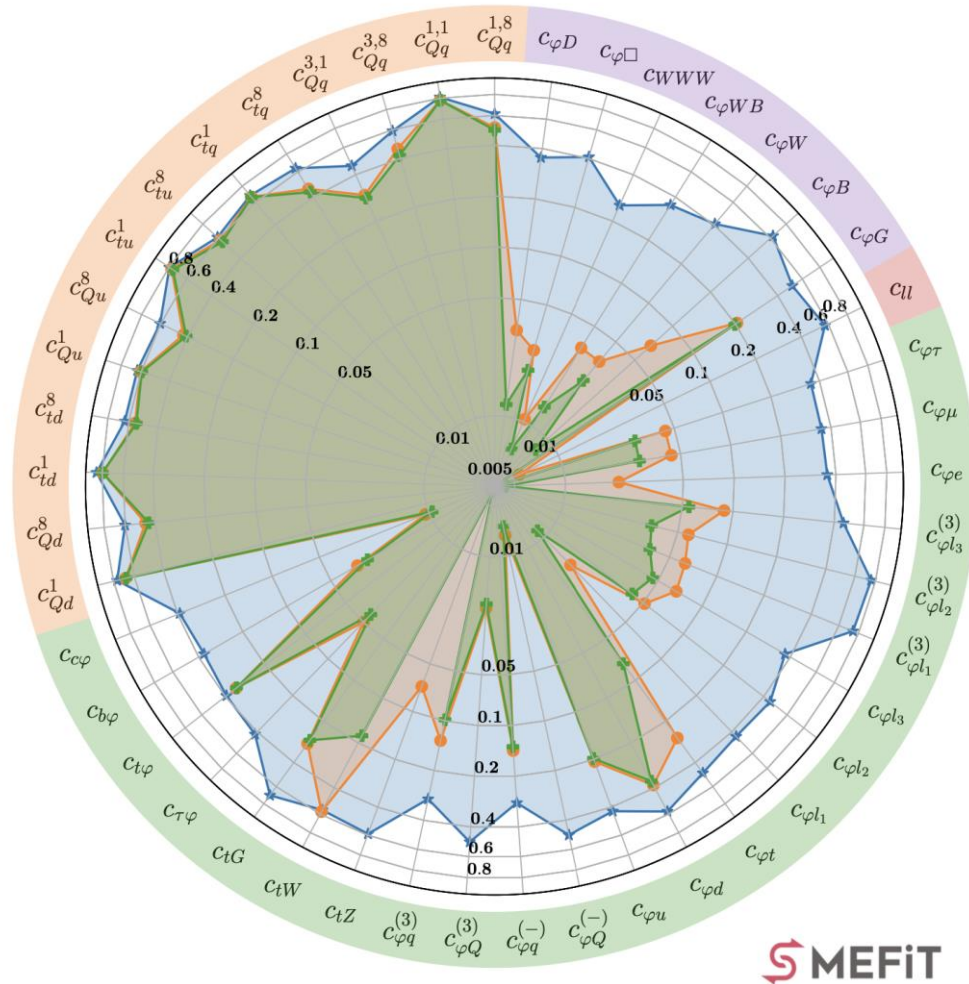


- ★ HL-LHC
- ✚ HL-LHC, individual
- SMEFiT3.0, individual



FCC-ee Energy runs

Ratio of Uncertainties to SMEFiT3.0 Baseline, $\mathcal{O}(\Lambda^{-2})$, Marginalised

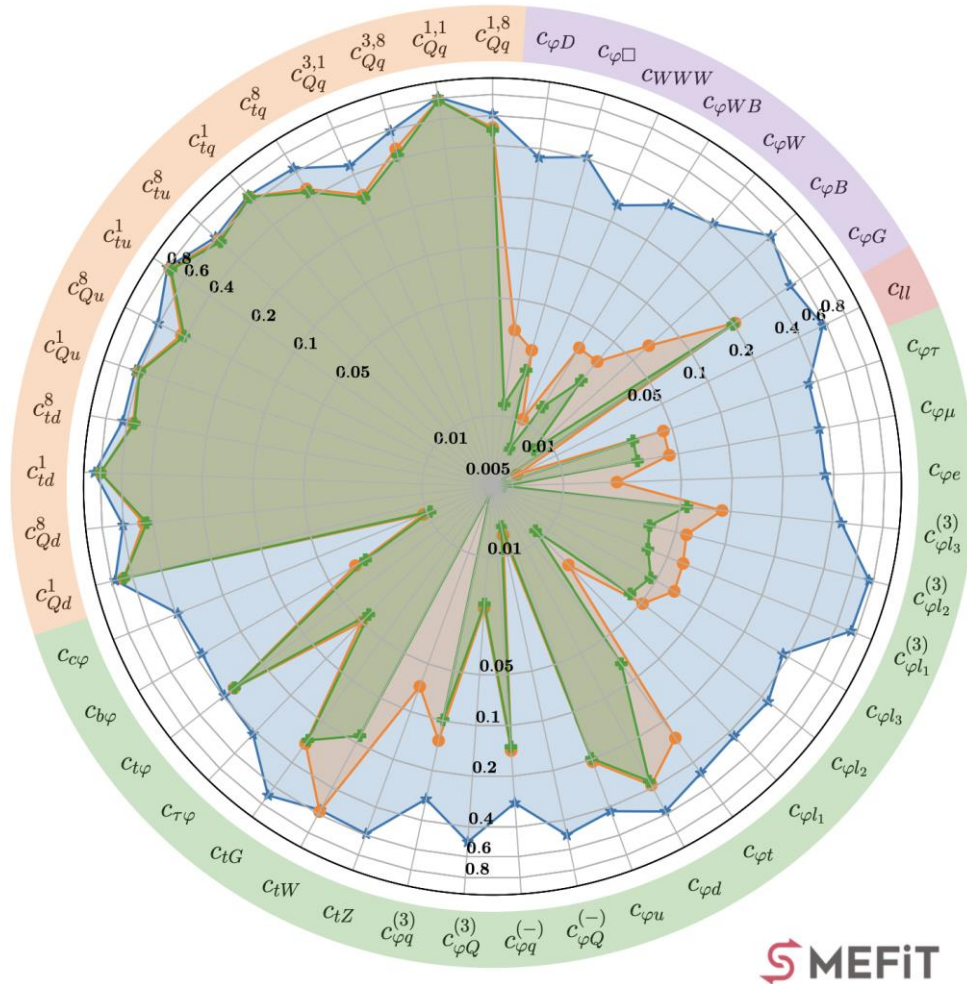


- ★ HL-LHC + FCC-ee (91 GeV)
- + HL-LHC + FCC-ee (91 + 161 + 240 + 365 GeV)
- HL-LHC + FCC-ee (91 + 240 GeV)



FCC-ee Energy runs

Ratio of Uncertainties to SMEFiT3.0 Baseline, $\mathcal{O}(\Lambda^{-2})$, Marginalised



SMEFiT

- ★ HL-LHC + FCC-ee (91 GeV)
- ✚ HL-LHC + FCC-ee (91 + 161 + 240 + 365 GeV)
- HL-LHC + FCC-ee (91 + 240 GeV)

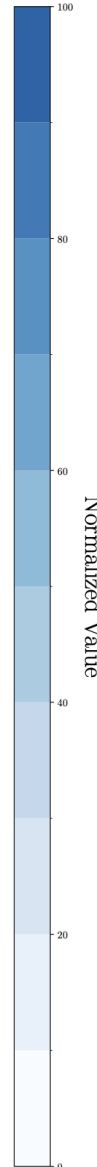
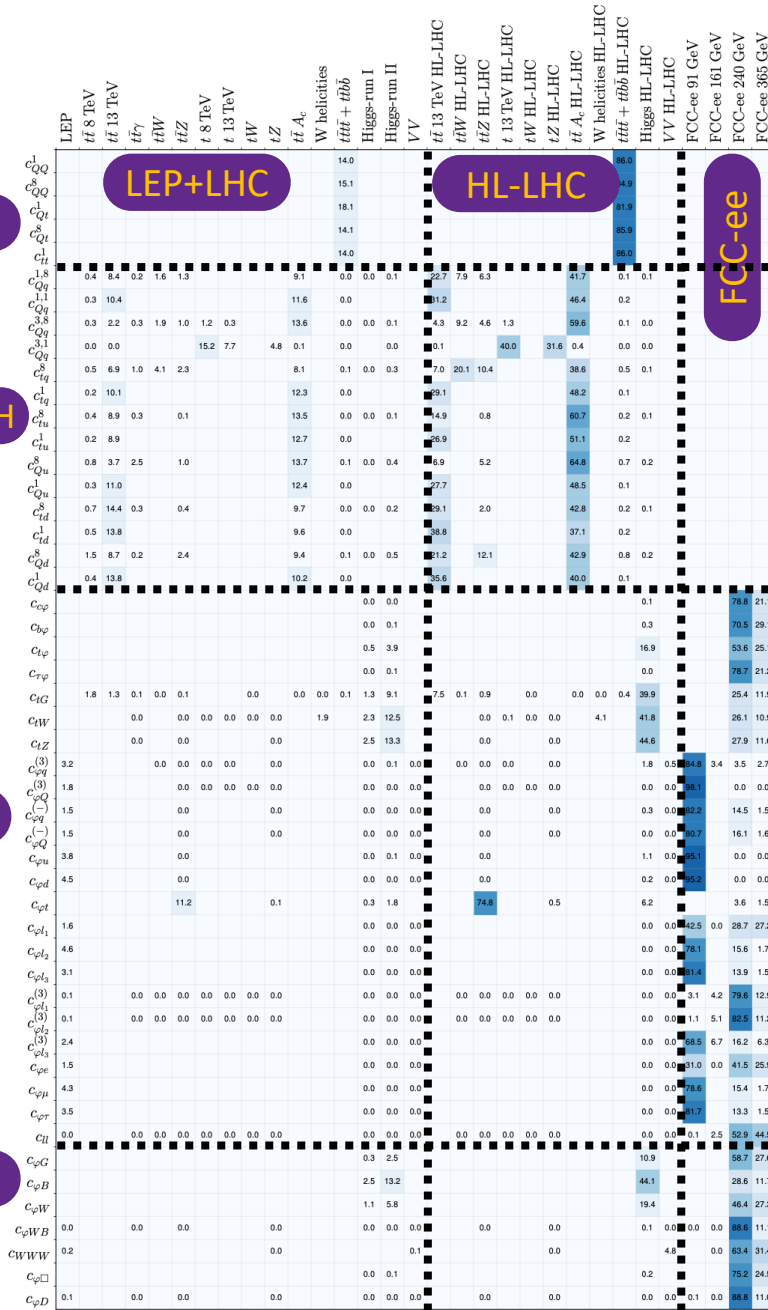
Fisher Information matrix

4H

2L2H

2F

Bos



- Quantifies which datasets have more sensitivity to given operator
- Proxy for linear individual fit
- FCC-ee dominates nearly all operators except 4-quark operators, only accessible in pp collisions (tree level)
- Combination of 91 GeV and 240 GeV runs important to pin down 2-fermion and gauge operators
- FCC-ee run at 161 GeV is the least useful for the SMEFT

$$I_{ij} = \sum_{m=1}^{n_{\text{dat}}} \frac{\sigma_{m,i}^{(\text{eft})} \sigma_{m,j}^{(\text{eft})}}{\delta_{\text{exp},m}^2}, \quad i, j = 1, \dots, n_{\text{eft}},$$

The power of multi(di)-boson

