

Future Circular Collider — Physics Potential —

2024 LHC Days Split



Christophe Grojean

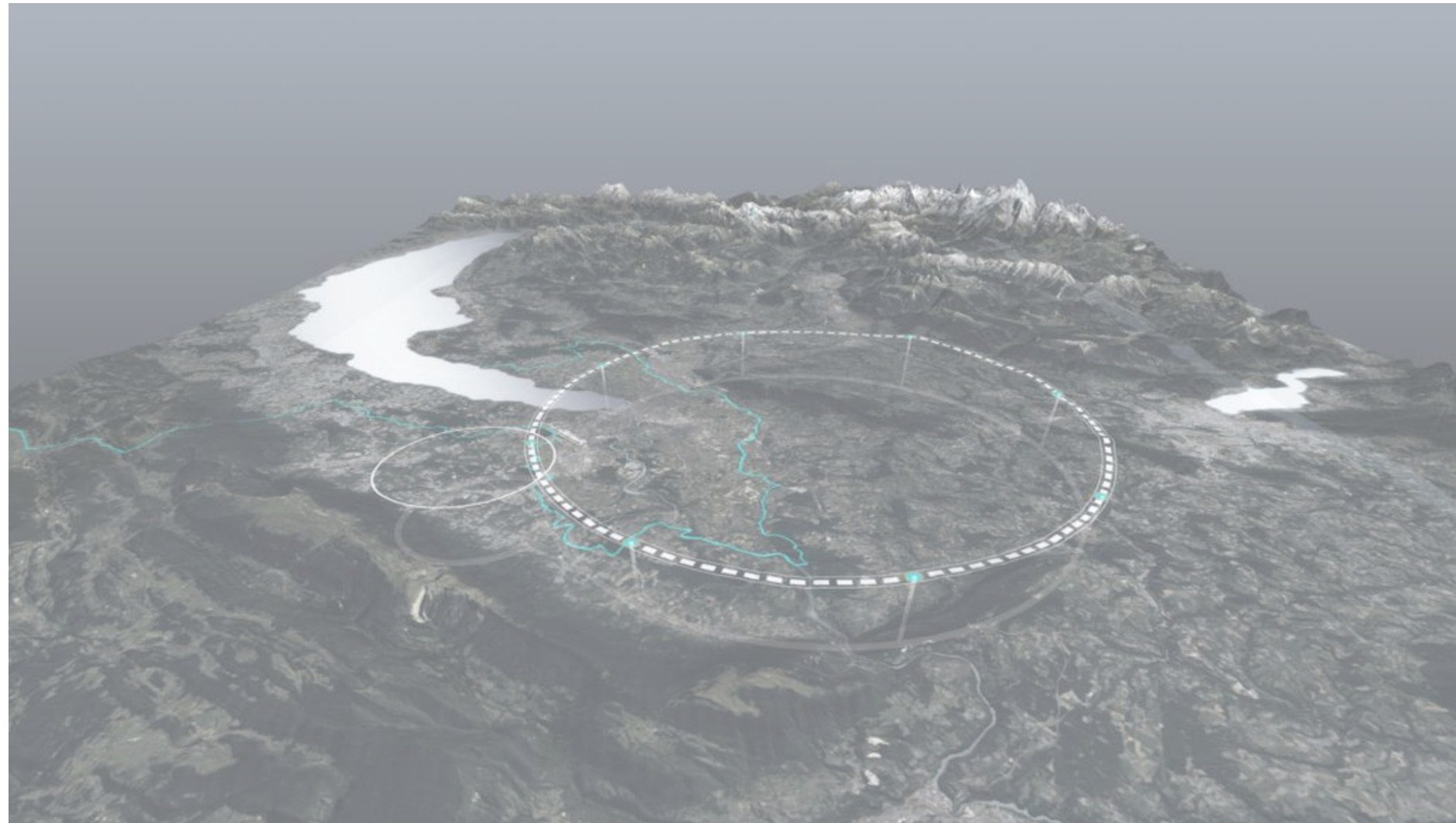
(christophe.grojean@desy.de)

— on behalf of the FCC team —



Future Circular Collider

- A versatile particle collider housed in a 91km underground ring around CERN.
- Implemented in several stages:
 - an e^+e^- “Higgs/EW/Flavour/top/QCD” factory running at 90-365 GeV  **FCC-ee**
 - followed by a high-energy pp collider reaching 100 TeV  **FCC-hh**



FCC on a Fast Track

After just over a decade of pioneering work, huge progress has been achieved:

- The first proposal of a high-luminosity e^+e^- circular collider to study the Higgs boson was made **thirteen years** ago (December 2011) [A. Blondel & F. Zimmermann following discussions with P. Janot at CERN cafeteria on a bright 2011 summer night speculating on the rumours of a Higgs at **140 GeV**];
- The Future Circular Collider collaboration was created **ten years** ago, towards the conceptual design study of a **100 TeV pp collider**, with an e^+e^- Higgs factory as a potential intermediate step;
- The **Conceptual Design Reports** of the FCC physics case, and of the FCC-ee and FCC-hh colliders, were published **five years** ago and submitted to the 2018-19 European Strategy Update;
- The CERN Council updated the European Strategy **three years** ago, stating that an e^+e^- Higgs factory would be the highest priority next collider, to be followed by a proton-proton collider at the highest achievable energy;
- **Two years** ago, the CERN Council consequently initiated and funded a **technical and financial feasibility** study for FCC with focus on an e^+e^- electroweak and Higgs factory as a first stage, study to be completed by the time of the next European Strategy Update;
- **Eleven months** ago, a 700+ pages **mid-term report** about the FCC feasibility was submitted to the CERN Council for a thorough review, with a conclusion expected at the beginning of 2025. Very positive feedback from CERN council in **Feb. 2.**

FCC-hh tunnel is great for FCC-ee

- **80-100 km is needed to accelerate pp up to 100 TeV**
- **80-100 km is also exactly what is needed**
 - to get enough luminosity (5 times more than in 27 km) to get sensitivity to the Higgs self coupling, the electron Yukawa coupling, or sterile neutrinos,
 - to make TeraZ a useful flavour factory,
 - for transverse polarisation to be available all the way to the WW threshold (allowing a precise W mass measurement)
 - for the top threshold to be reached and exceeded.

FCC-hh tunnel is great for FCC-ee

- 80-100 km is needed to accelerate pp up to 100 TeV
- 80-100 km is also exactly what is needed
 - to get enough luminosity (5 times more than in 27 km) to get sensitivity to the Higgs self coupling, the electron Yukawa coupling, or sterile neutrinos,
 - to make TeraZ a useful flavour factory,
 - for transverse polarisation to be available all the way to the WW threshold (allowing a precise W mass measurement)
 - for the top threshold to be reached and exceeded.

After the success of LHC, we need a broad, versatile and ambitious programme that

1. sharpens our knowledge of already discovered physics → **guaranteed deliverables**
2. pushes the frontiers of the unknown at high and low scales → **exploration**

— together FCC-ee & FCC-hh combine these 2 aspects —

more PRECISION and more ENERGY, for more SENSITIVITY to New Physics

Precision as a discovery tool

Many historical examples

- ▶ Uranus anomalous trajectory \rightsquigarrow Neptune
- ▶ Mercury perihelion \rightsquigarrow General Relativity
- ▶ Z/W interactions to quarks and leptons \rightsquigarrow Higgs boson
- ▶ ...

Precision as a discovery tool

Many historical examples

- ▶ Uranus anomalous trajectory \rightsquigarrow Neptune
- ▶ Mercury perihelion \rightsquigarrow General Relativity
- ▶ Z/W interactions to quarks and leptons \rightsquigarrow Higgs boson
- ▶ ...

Sometimes, these discoveries were expected based on theoretical arguments
(e.g. Rayleigh-Jeans UV catastrophe for QM, unitarity breakdown for the Higgs)
but precision gave valuable additional clues.

In any case, experimentalists shouldn't lean too heavily on theorist priors/prejudices
(remember discovery of P violation).

At times when we don't have a precise theoretical guidance, we need powerful experimental tools to make progress.

Precision as a discovery tool

Many historical examples

- ▶ Uranus anomalous trajectory \rightsquigarrow Neptune
- ▶ Mercury perihelion \rightsquigarrow General Relativity
- ▶ Z/W interactions to quarks and leptons \rightsquigarrow Higgs boson
- ▶ ...

Sometimes, these discoveries were expected based on theoretical arguments
(e.g. Rayleigh-Jeans UV catastrophe for QM, unitarity breakdown for the Higgs)
but precision gave valuable additional clues.

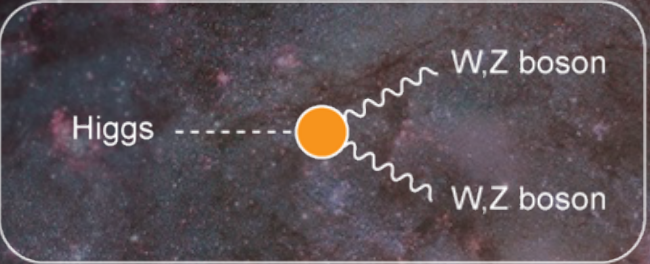
In any case, experimentalists shouldn't lean too heavily on theorist priors/prejudices
(remember discovery of P violation).

At times when we don't have a precise theoretical guidance, we need powerful experimental tools to make progress.

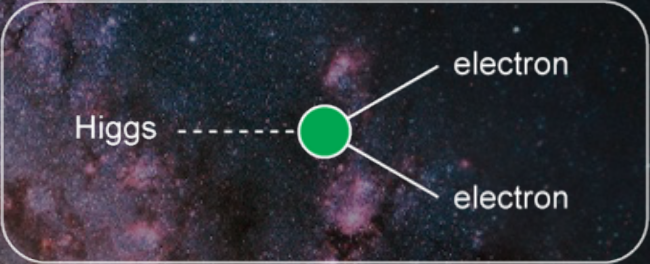
The FCC project offers unprecedented opportunities on many different fronts.
No LHC/SSC-like **no-lose theorem** but a **promise** of making significant
steps forward in our understanding of the fundamental laws of Nature.

Why a Higgs Factory?

Star life-time



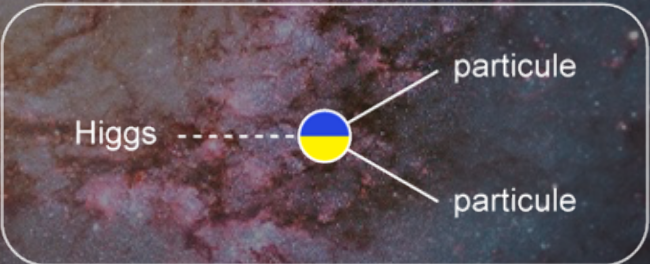
Size of atoms
Stability of nuclei/matter



Birth of vacuum



Matter/antimatter
dominance



Physics Overview

FCC-ee Run Plan.

LEP1 data accumulated in **every 2 mn.** Exciting & diverse programme with different priorities every few years.
(order of the different stages still subject to discussion/optimisation)

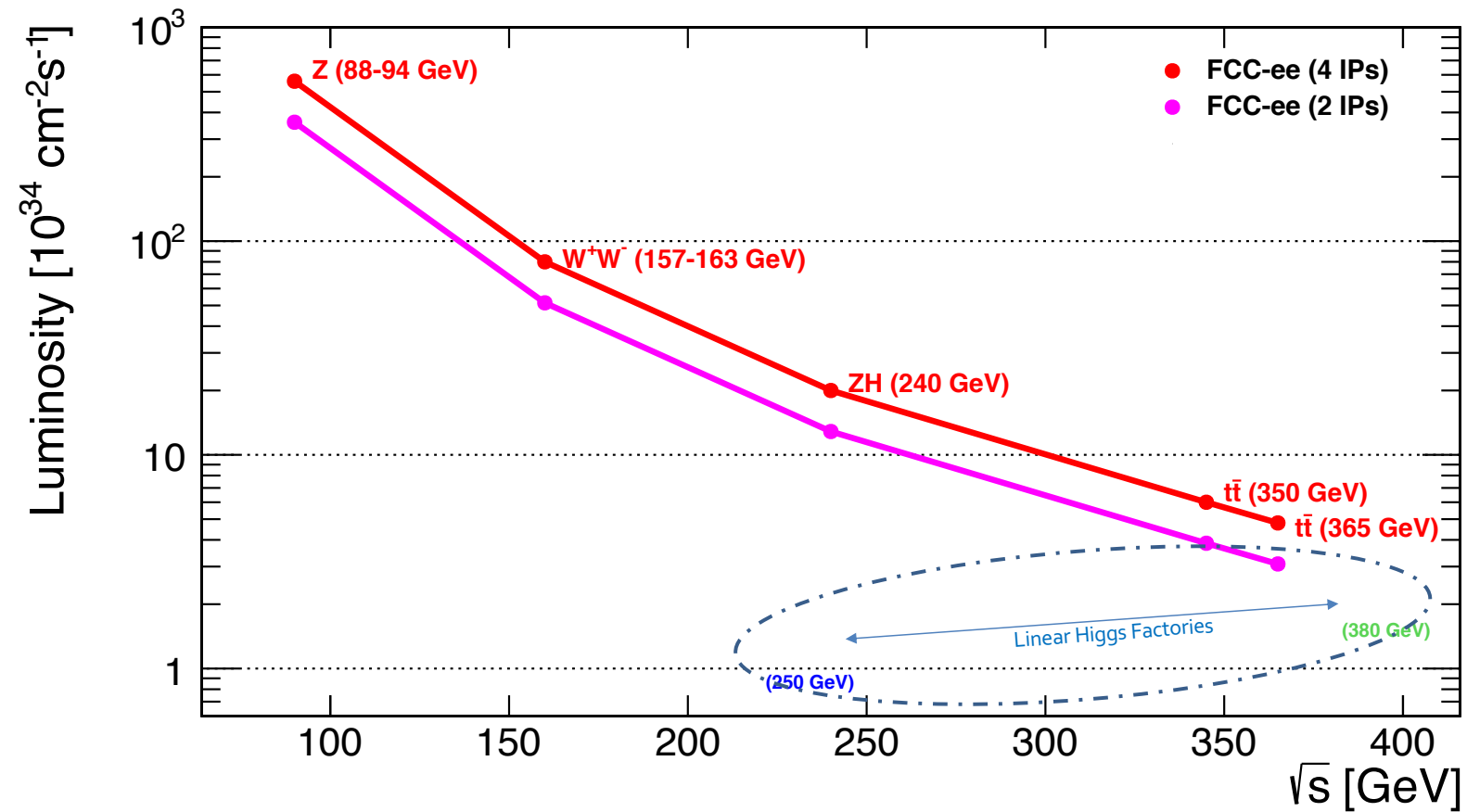


Fig. to be updated: new optics design (May 2024) gives 50% more lumi @ 240 GeV.

FCC-ee Run Plan.

LEP1 data accumulated in **every 2 mn.** Exciting & diverse programme with different priorities every few years.
 (order of the different stages still subject to discussion/optimisation)

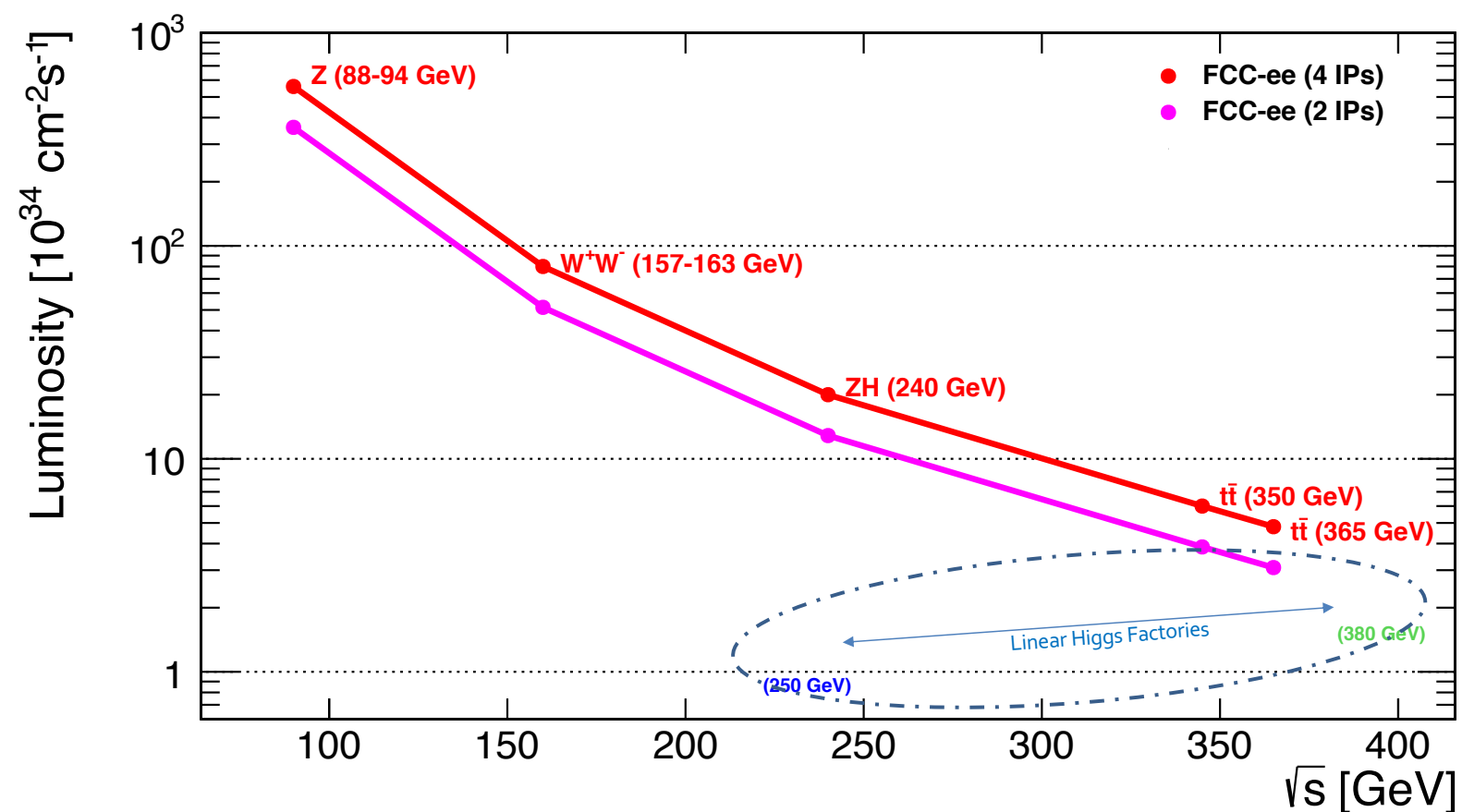
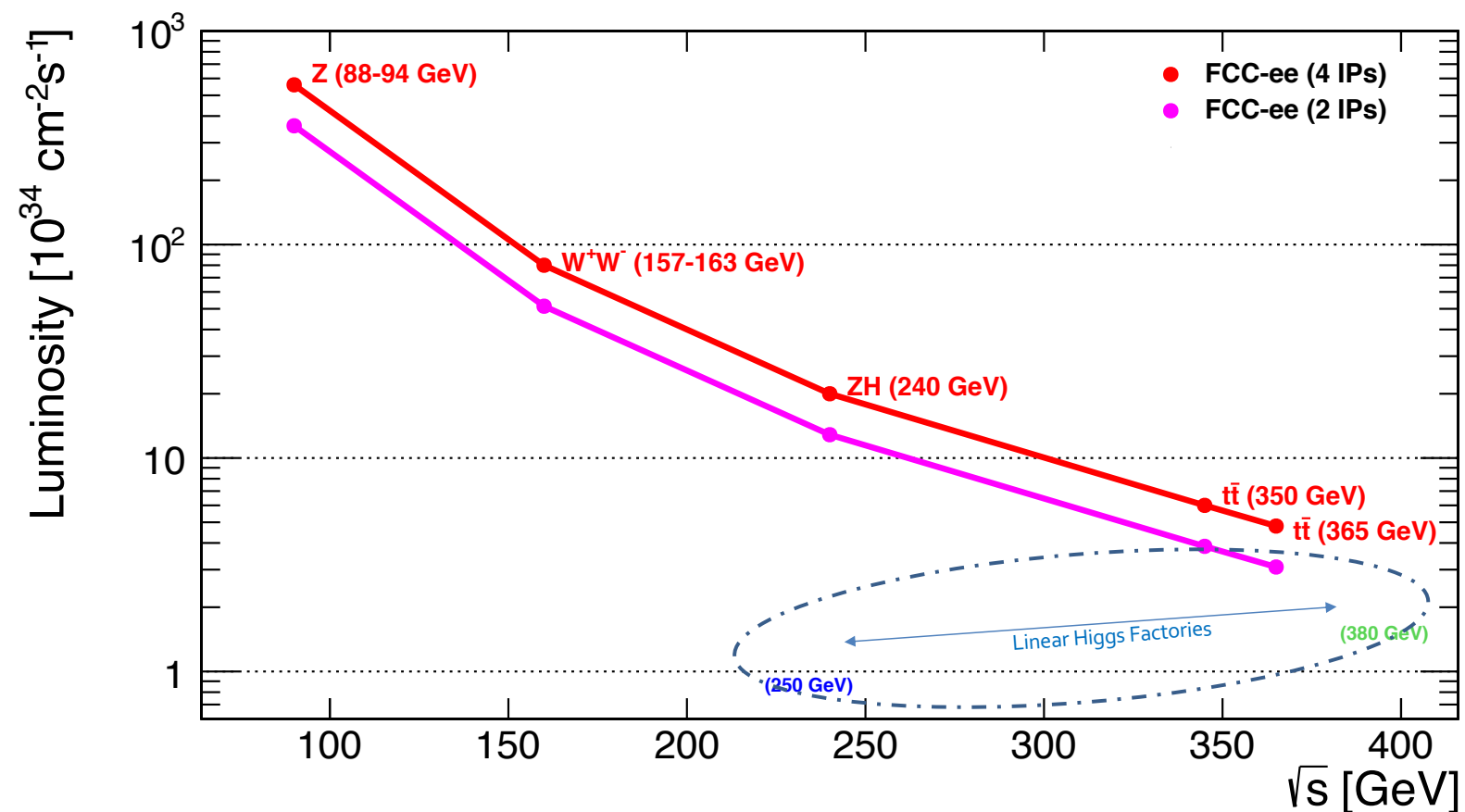


Fig. to be updated: new optics design (May 2024) gives 50% more lumi @ 240 GeV.

Working point	Z, years 1-2	Z, later	WW, years 1-2	WW, later	ZH	$t\bar{t}$
\sqrt{s} (GeV)	88, 91, 94		157, 163		240	340–350 365
Lumi/IP ($10^{34} \text{ cm}^{-2} \text{ s}^{-1}$)	70	140	10	20	5.0	0.75 1.20
Lumi/year (ab^{-1})	34	68	4.8	9.6	2.4	0.36 0.58
Run time (year)	2	2	2	–	3	1 4
Number of events	6×10^{12} Z		2.4×10^8 WW		1.45×10^6 ZH + 45k WW \rightarrow H	1.9×10^6 $t\bar{t}$ +330k ZH +80k WW \rightarrow H

FCC-ee Run Plan.

LEP1 data accumulated in **every 2 mn.** Exciting & diverse programme with different priorities every few years.
 (order of the different stages still subject to discussion/optimisation)



— Superb statistics achieved in only 15 years —

in each detector:
 10^5 Z/sec, 10^4 W/hour,
1500 Higgs/day, 1500 top/day

Fig. to be updated: new optics design (May 2024) gives 50% more lumi @ 240 GeV.

Working point	Z, years 1-2	Z, later	WW, years 1-2	WW, later	ZH	$t\bar{t}$
\sqrt{s} (GeV)	88, 91, 94		157, 163		240	340–350 365
Lumi/IP ($10^{34} \text{ cm}^{-2} \text{ s}^{-1}$)	70	140	10	20	5.0	0.75 1.20
Lumi/year (ab^{-1})	34	68	4.8	9.6	2.4	0.36 0.58
Run time (year)	2	2	2	–	3	1 4
Number of events	6×10^{12} Z		2.4×10^8 WW		1.45×10^6 ZH + 45k WW \rightarrow H	1.9×10^6 $t\bar{t}$ +330k ZH +80k WW \rightarrow H

FCC-ee Run Plan.

LEP1 data accumulated in **every 2 mn.** Exciting & diverse programme with different priorities every few years.
(order of the different stages still subject to discussion/optimisation)

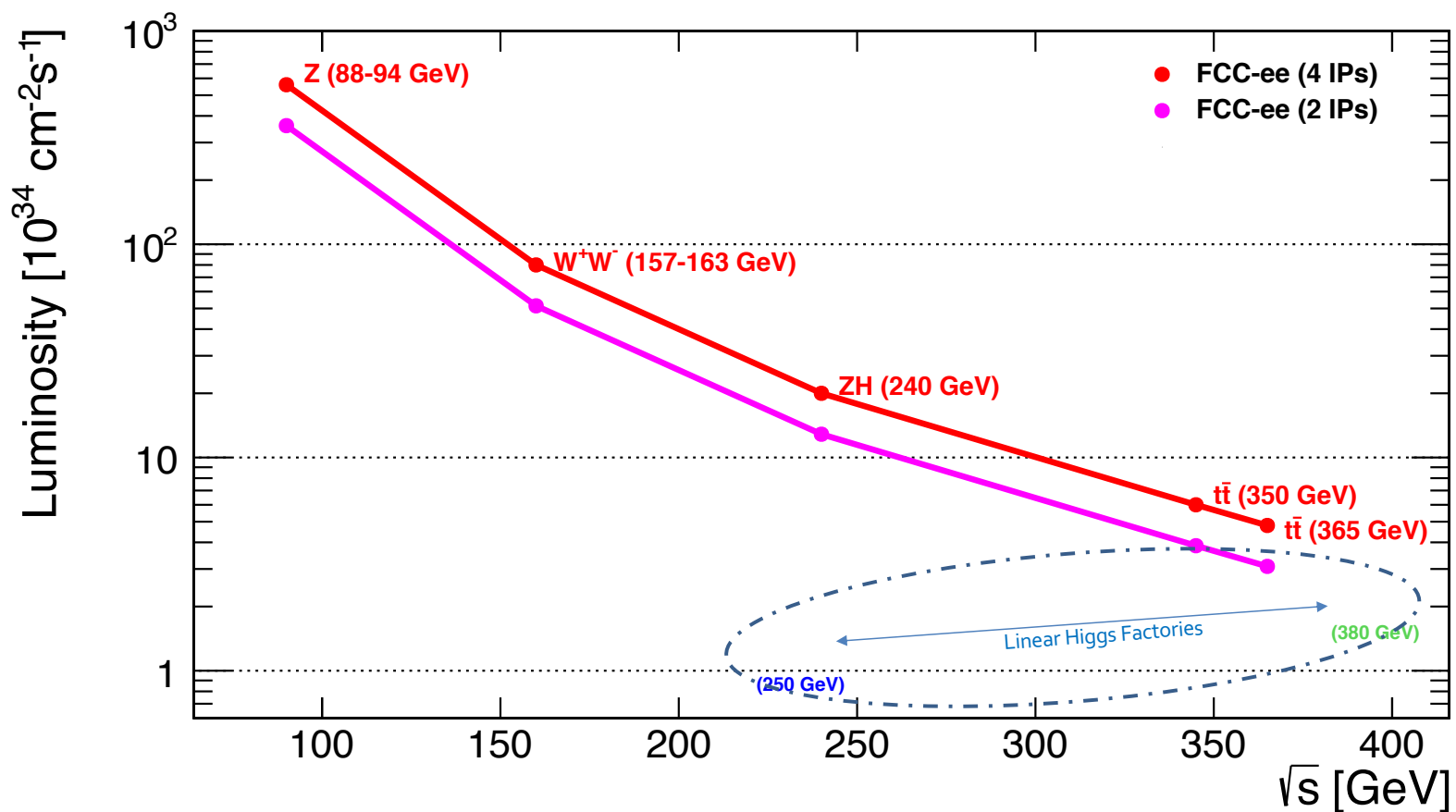
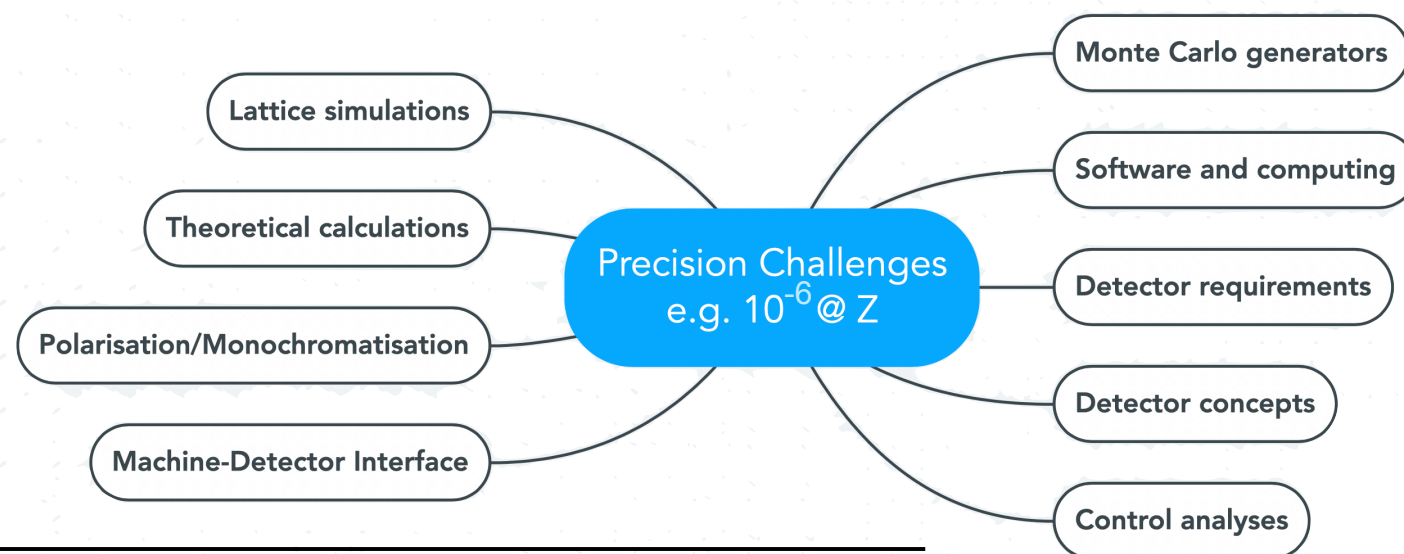


Fig. to be updated: new optics design (May 2024) gives 50% more lumi @ 240 GeV.

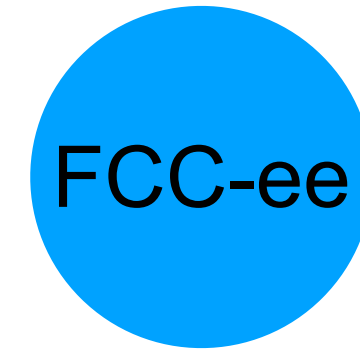
Working point	Z, years 1-2	Z, later	WW, years 1-2	WW, later	ZH	tt
\sqrt{s} (GeV)	88, 91, 94		157, 163		240	340–350 365
Lumi/IP ($10^{34} \text{ cm}^{-2} \text{ s}^{-1}$)	70	140	10	20	5.0	0.75 1.20
Lumi/year (ab^{-1})	34	68	4.8	9.6	2.4	0.36 0.58
Run time (year)	2	2	2	–	3	1 4
Number of events	6×10^{12} Z		2.4×10^8 WW		1.45×10^6 ZH + 45k WW \rightarrow H	1.9×10^6 tt +330k ZH +80k WW \rightarrow H

— Superb statistics achieved in only 15 years —

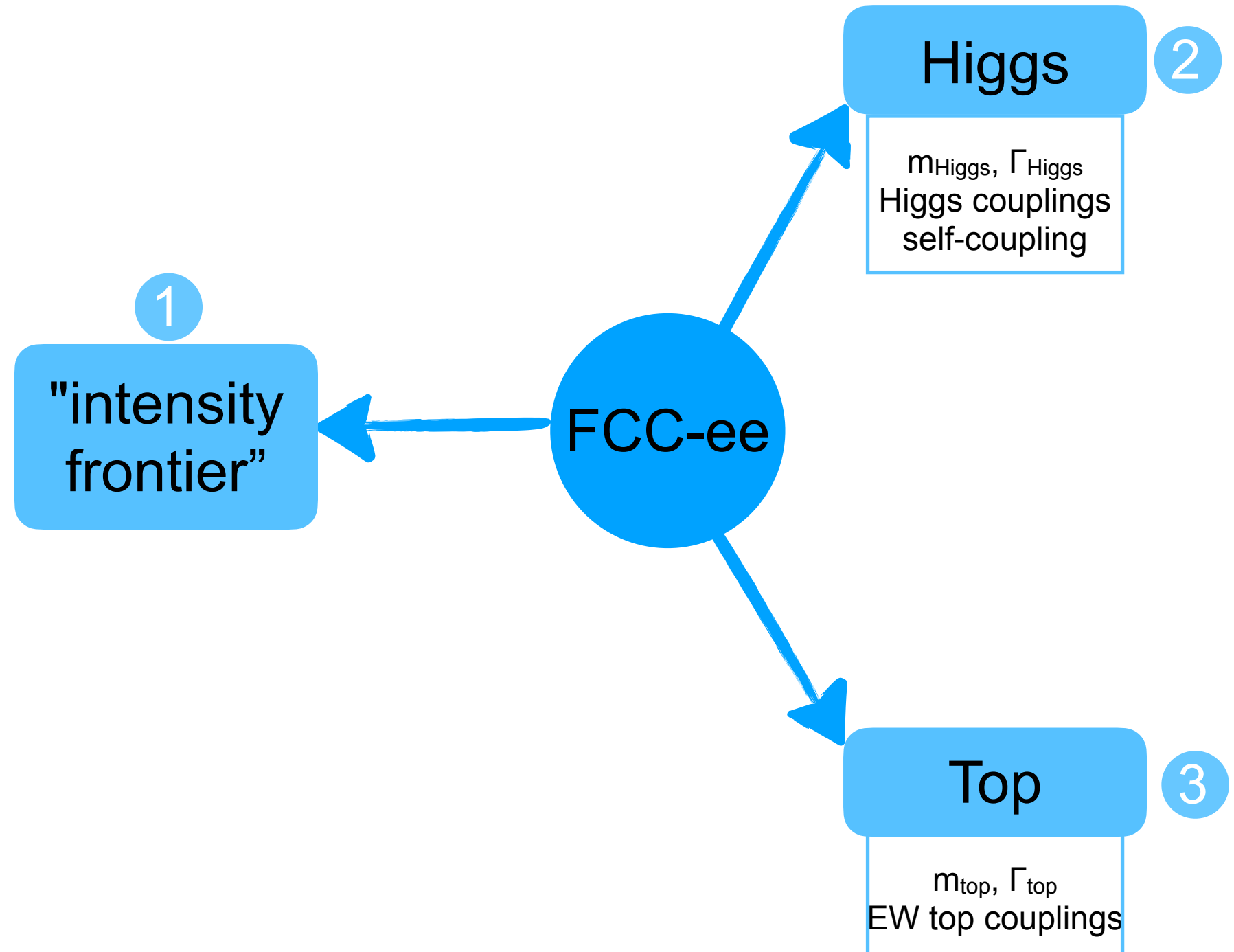
in each detector:
 **10^5 Z/sec, 10^4 W/hour,
 1500 Higgs/day, 1500 top/day**



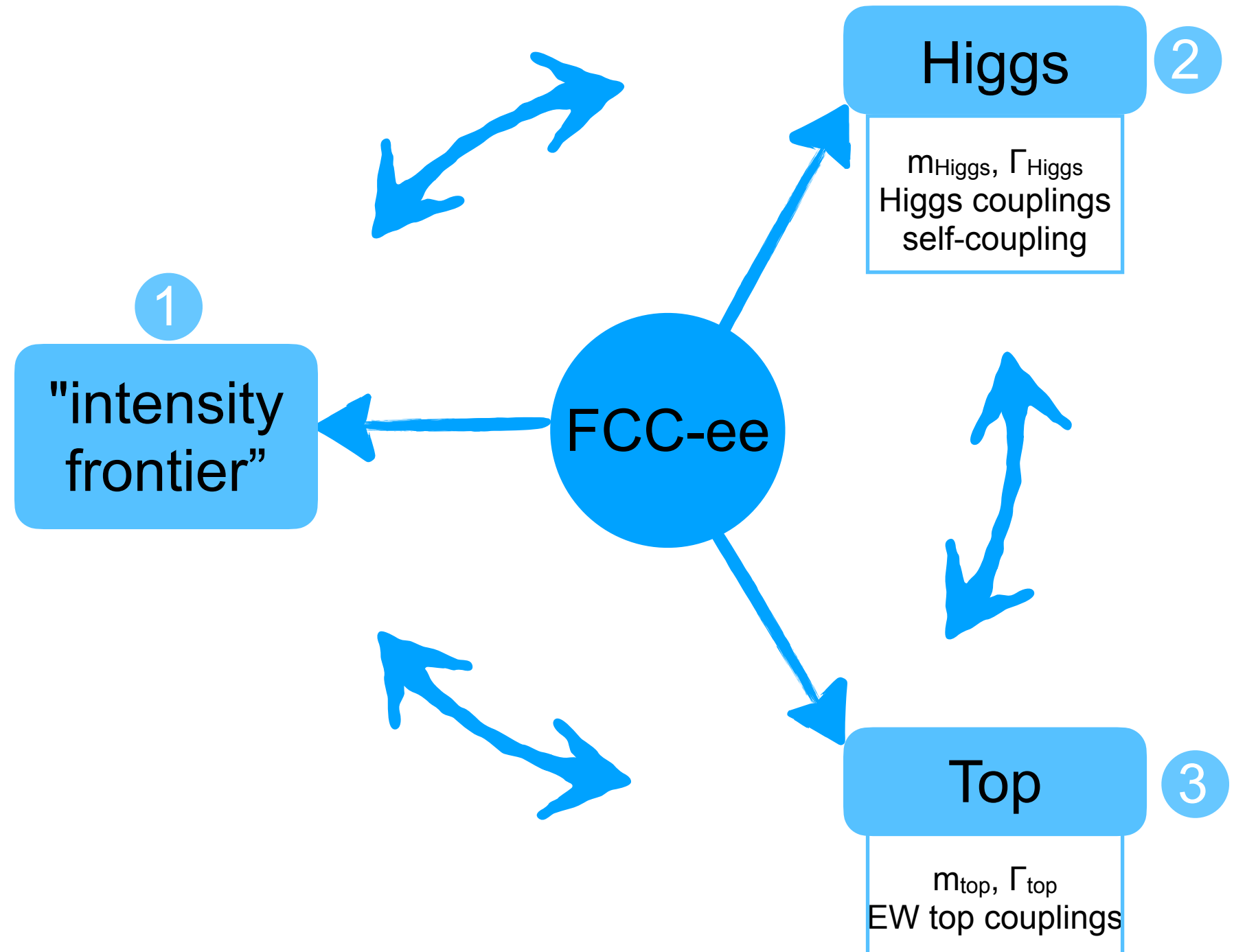
FCC-ee Physics Programme.



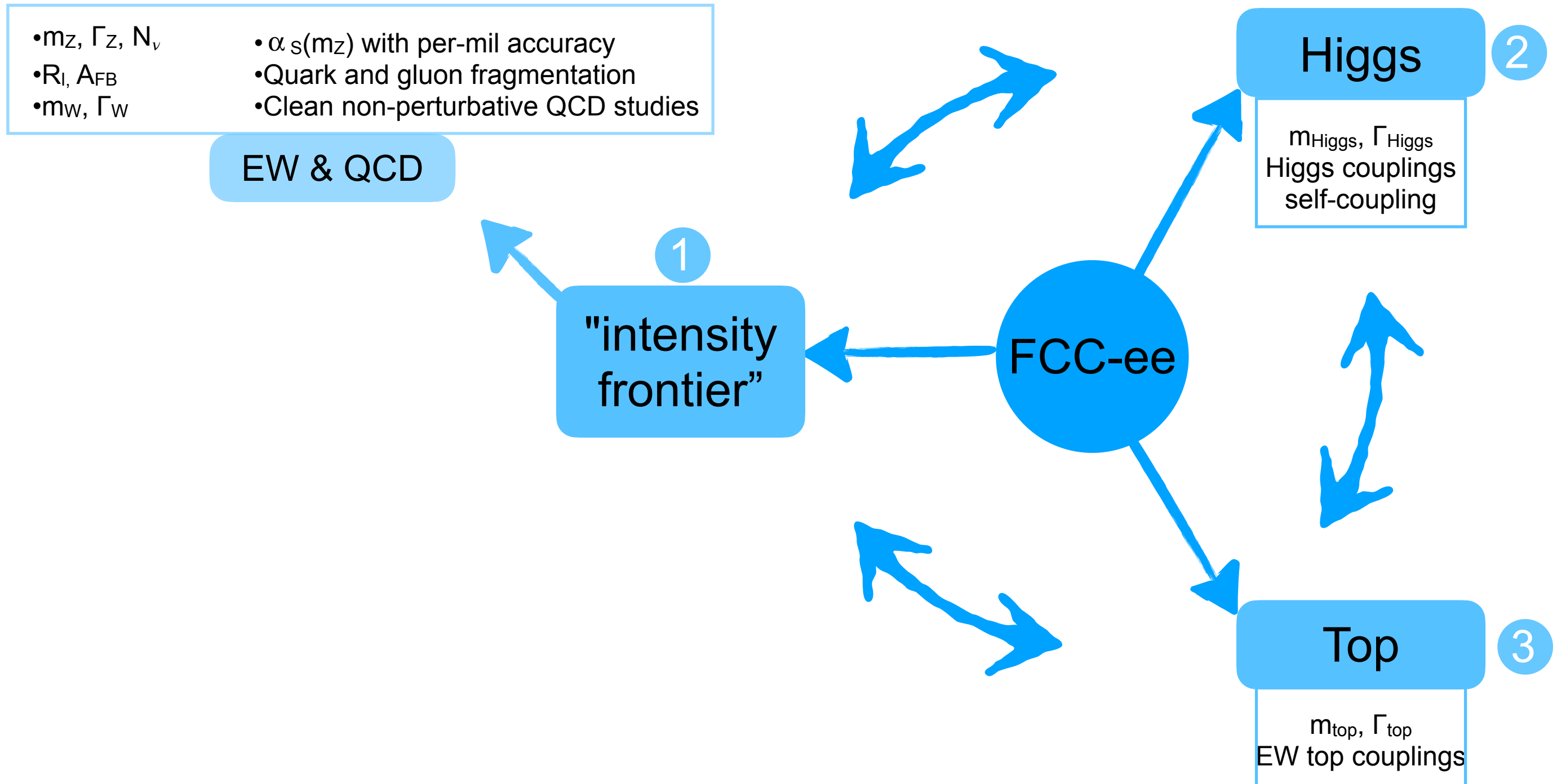
FCC-ee Physics Programme.



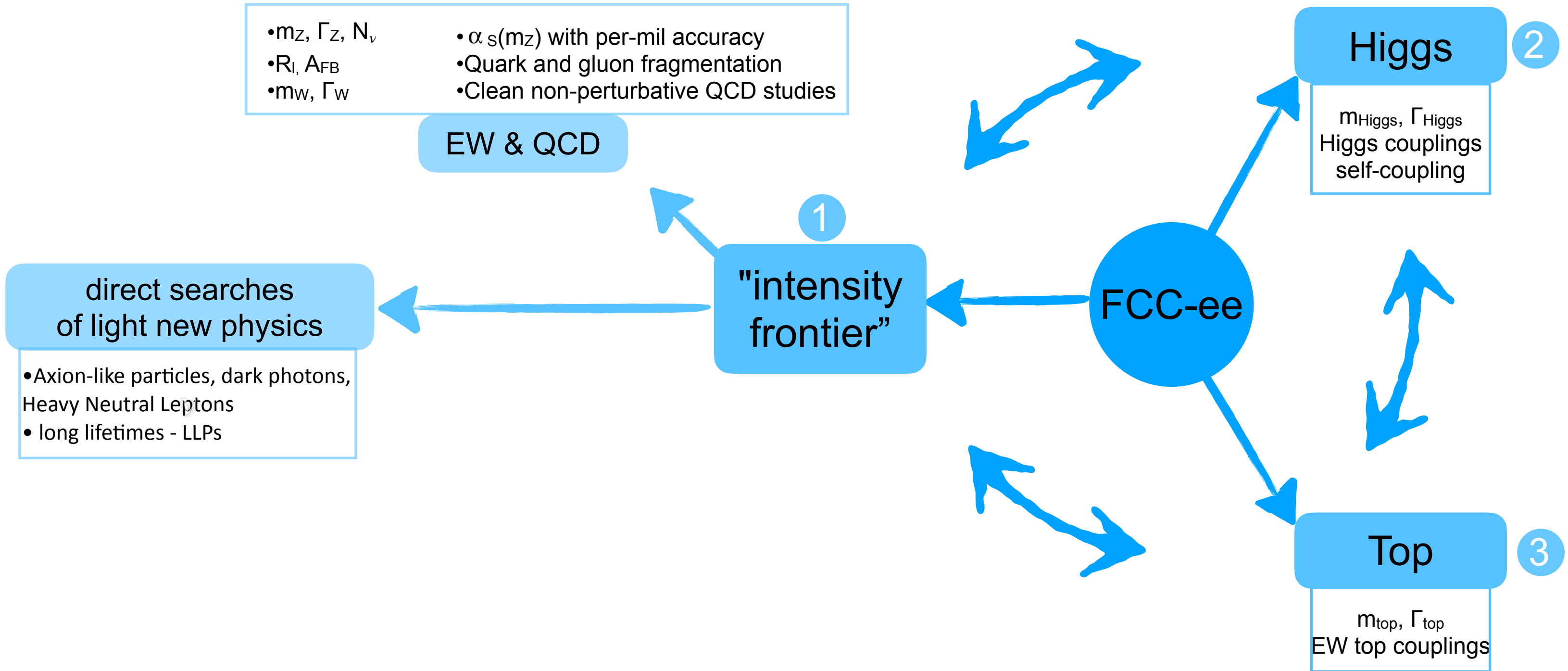
FCC-ee Physics Programme.



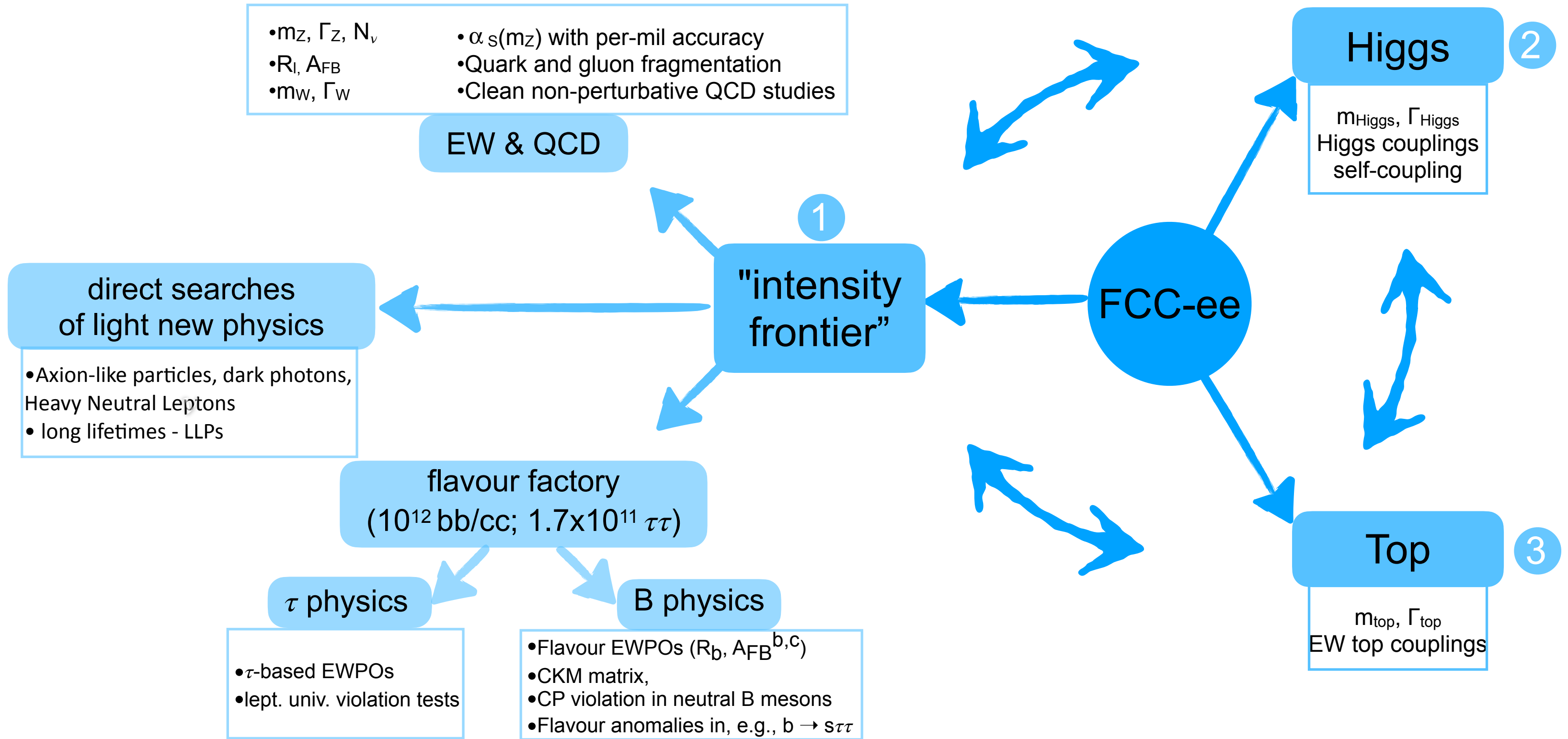
FCC-ee Physics Programme.



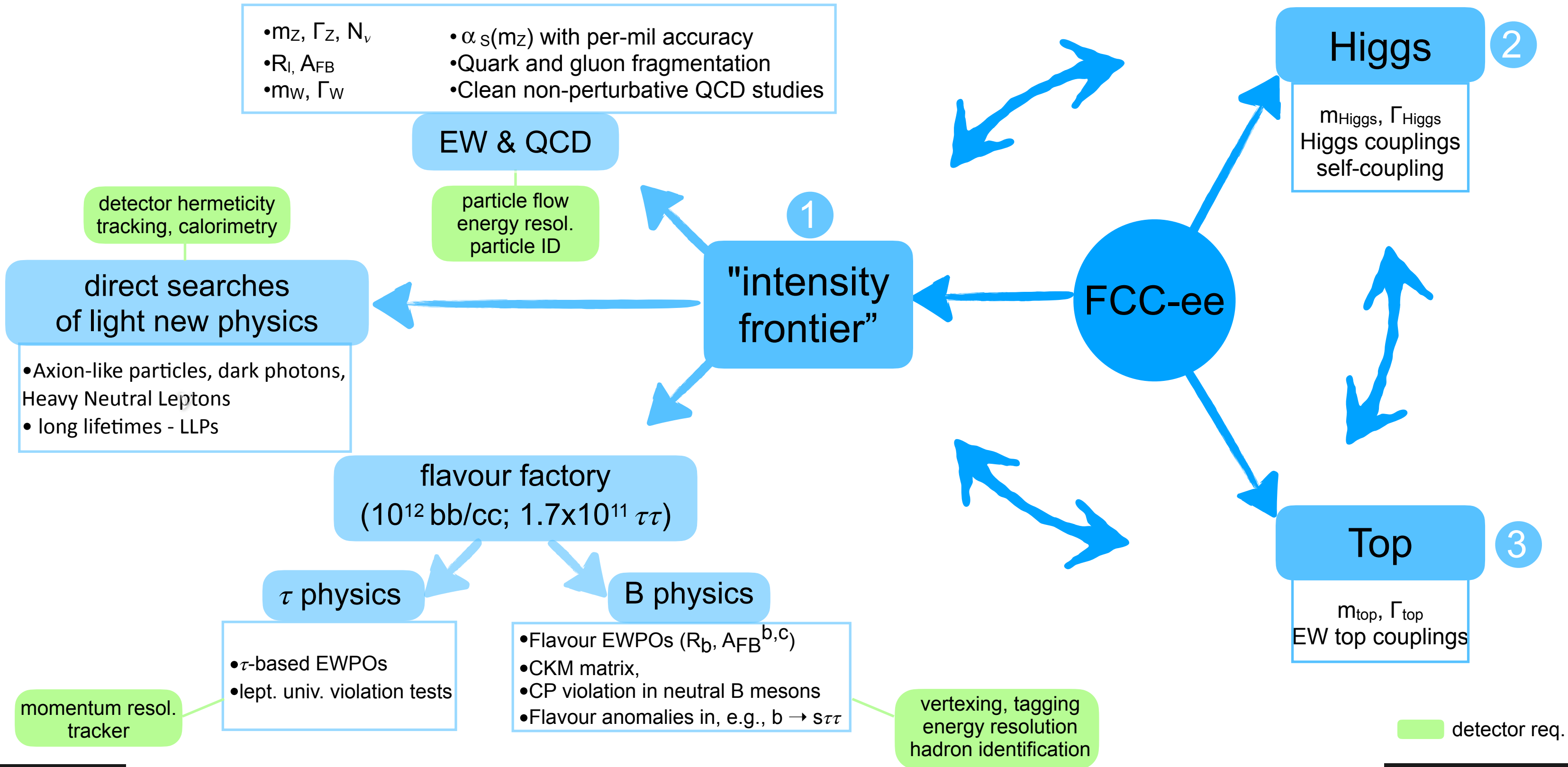
FCC-ee Physics Programme.



FCC-ee Physics Programme.



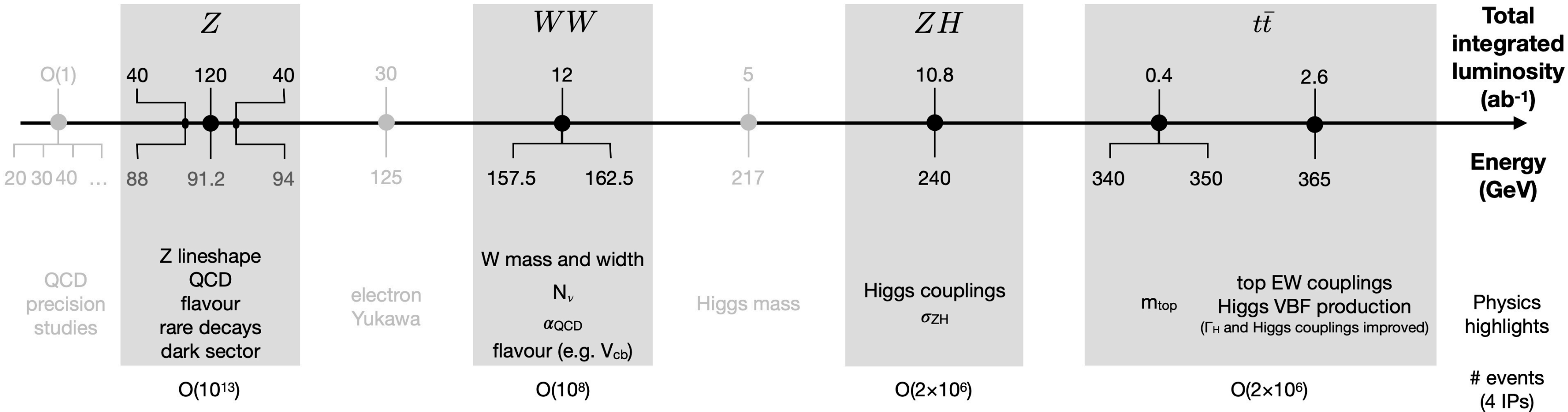
FCC-ee Physics Programme.



Collider Programme (and beyond).

— CDR baseline runs (4IPs)

— Additional opportunities

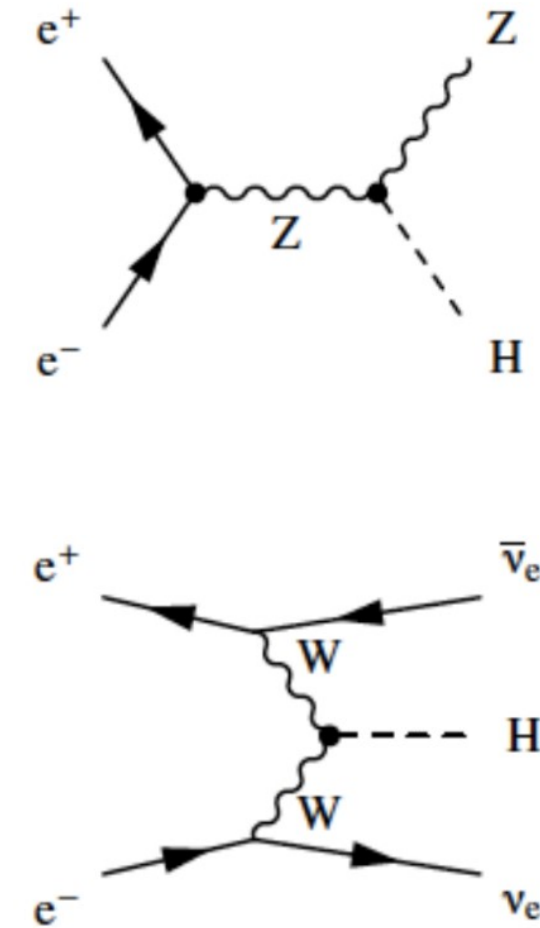
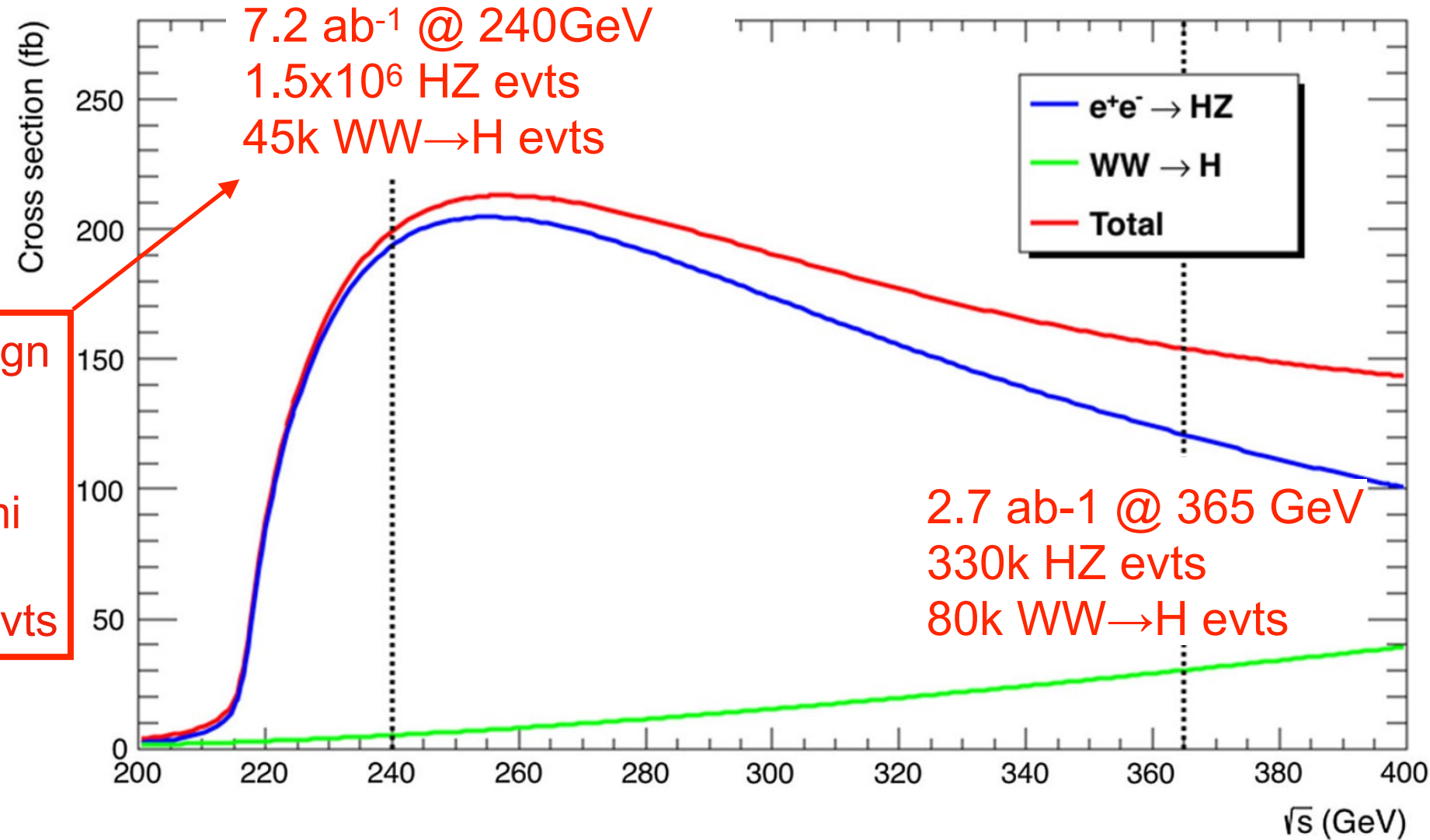


- **Opportunities** beyond the baseline plan (\sqrt{s} below Z, 125GeV, 217GeV; larger integrated lumi...)
- **Opportunities** to exploit FCC facility differently (to be studied more carefully):
 - using the electrons from the injectors for beam-dump experiments,
 - extracting electron beams from the booster,
 - reusing the synchrotron radiation photons.

Higgs Physics

Higgs @ FCC-ee.

Central goal of FCC-ee: model-independent measurement of Higgs width and couplings with (<) % precision. Achieved through operation at two energy points.



new optics design (May 2024) gives 50% more lumi @ 240 GeV ⇒ 2.5x10⁶ HZ evts

Sensitivity to both processes very helpful in improving precision on couplings.

Complementarity with 365 GeV on top of 240 GeV

improvement factor: ∞/3/2/1.5/1.2 on $\kappa_\lambda/\kappa_W/\kappa_b/\kappa_g, \kappa_c/\kappa_\gamma$ (plot in bonus)

Higgs @ FCC-ee.

- Absolute normalisation of couplings (by recoil method). The LHC fit doesn't converge w/o making any assumption.
- Measurement of width (from $ZH \rightarrow ZZZ^*$ and $WW \rightarrow H$)
- $\delta\Gamma_H \sim 1\%$, $\delta m_H \sim 3 \text{ MeV}$ (resp. 25%, 0(20) MeV @ HL-LHC)
- Model-independent coupling determination and improvement factor up to 10 compared to LHC
- (Indirect) sensitivity to new physics up to 70-100 TeV (for maximally strongly coupled models)
($\delta\kappa_X = v^2/f^2$ & $m_{\text{NP}} = g_{\text{NP}} f$)
- Unique access to electron Yukawa

Higgs coupling sensitivity

Coupling	HL-LHC	FCC-ee (240–365 GeV)
		2 IPs / 4 IPs
κ_W [%]	1.5*	0.43 / 0.33
κ_Z [%]	1.3*	0.17 / 0.14
κ_g [%]	2*	0.90 / 0.77
κ_γ [%]	1.6*	1.3 / 1.2
$\kappa_{Z\gamma}$ [%]	10*	10 / 10
κ_c [%]	—	1.3 / 1.1
κ_t [%]	3.2*	3.1 / 3.1
κ_b [%]	2.5*	0.64 / 0.56
κ_μ [%]	4.4*	3.9 / 3.7
κ_τ [%]	1.6*	0.66 / 0.55
BR _{inv} (<%, 95% CL)	1.9*	0.20 / 0.15
BR _{unt} (<%, 95% CL)	4*	1.0 / 0.88

Table from mid-term report

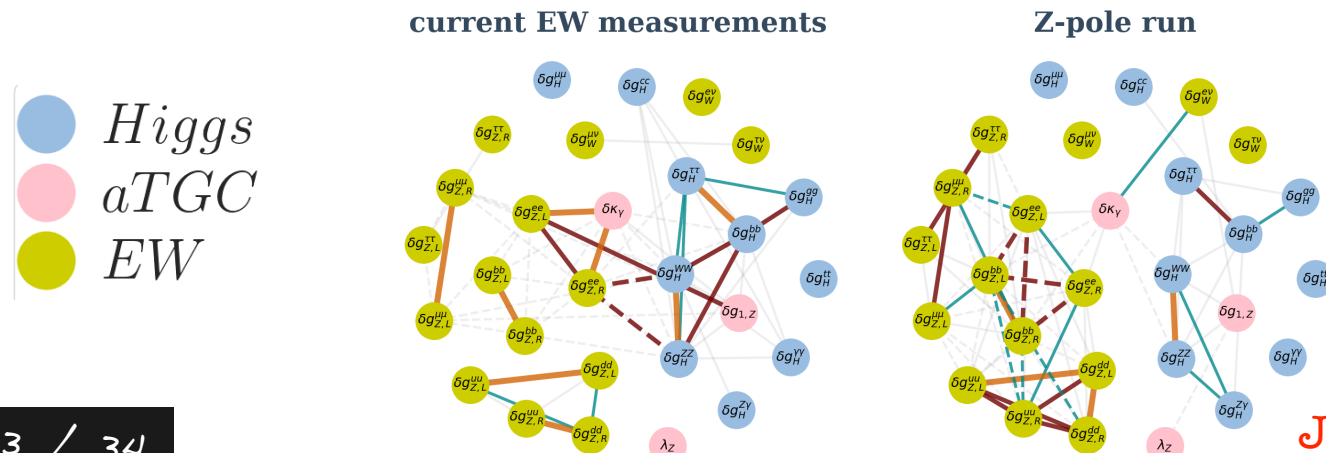
(new luminosity at 240GeV will further improve the coupling reach, e.g. 0.11% for κ_Z)

$$\kappa_X = \frac{g_{hXX}}{g_{hXX}^{\text{SM}}}$$

Higgs @ FCC-ee.

- Absolute normalisation of couplings (by recoil method). The LHC fit doesn't converge w/o making any assumption.
- Measurement of width (from $ZH \rightarrow ZZZ^*$ and $WW \rightarrow H$)
- $\delta\Gamma_H \sim 1\%$, $\delta m_H \sim 3 \text{ MeV}$ (resp. 25%, $O(20) \text{ MeV}$ @ HL-LHC)
- Model-independent coupling determination and improvement factor up to 10 compared to LHC
- (Indirect) sensitivity to new physics up to 70-100 TeV (for maximally strongly coupled models)
($\delta\kappa_X = v^2/f^2$ & $m_{\text{NP}} = g_{\text{NP}} f$)
- Unique access to electron Yukawa

— Higgs programme needs Z-pole —



Higgs coupling sensitivity

Coupling	HL-LHC	FCC-ee (240–365 GeV)
		2 IPs / 4 IPs
κ_W [%]	1.5*	0.43 / 0.33
κ_Z [%]	1.3*	0.17 / 0.14
κ_g [%]	2*	0.90 / 0.77
κ_γ [%]	1.6*	1.3 / 1.2
$\kappa_{Z\gamma}$ [%]	10*	10 / 10
κ_c [%]	—	1.3 / 1.1
κ_t [%]	3.2*	3.1 / 3.1
κ_b [%]	2.5*	0.64 / 0.56
κ_μ [%]	4.4*	3.9 / 3.7
κ_τ [%]	1.6*	0.66 / 0.55
BR_{inv} (<%, 95% CL)	1.9*	0.20 / 0.15
BR_{unt} (<%, 95% CL)	4*	1.0 / 0.88

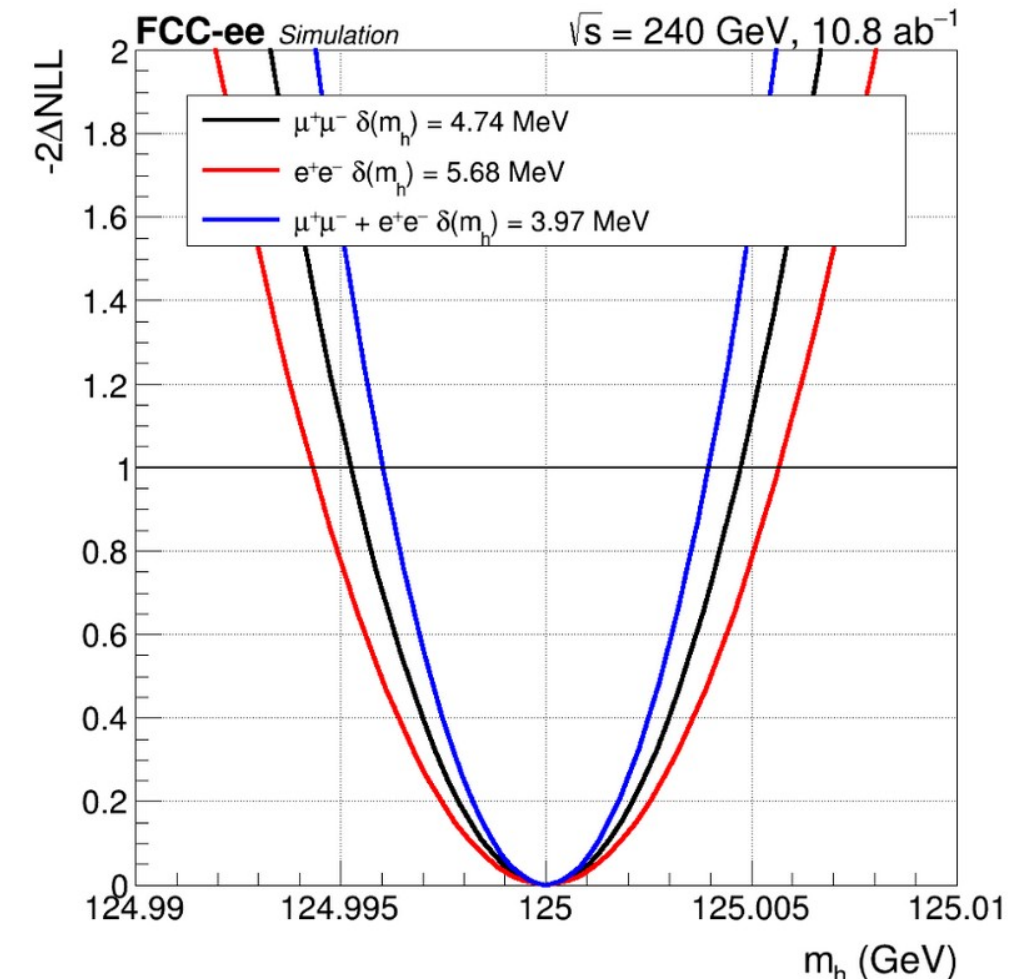
Table from mid-term report

(new luminosity at 240GeV will further improve the coupling reach, e.g. 0.11% for κ_Z)

$$\kappa_X = \frac{g_{hXX}}{g_{hXX}^{\text{SM}}}$$

Higgs Mass

- Recoil mass in Z(l)H events (l=e,μ)
- Thorough study of detector design impact
 - Larger variations from track resolution
 - High field & lighter tracker beneficial



Robust prospects to reach and even go below the natural 4.1 MeV limit set by the SM Higgs width

Nominal configuration

Crystal ECAL to Dual Readout

Nominal 2 T → field 3 T

IDEA drift chamber → CLD Si tracker

Impact of Beam Energy Spread

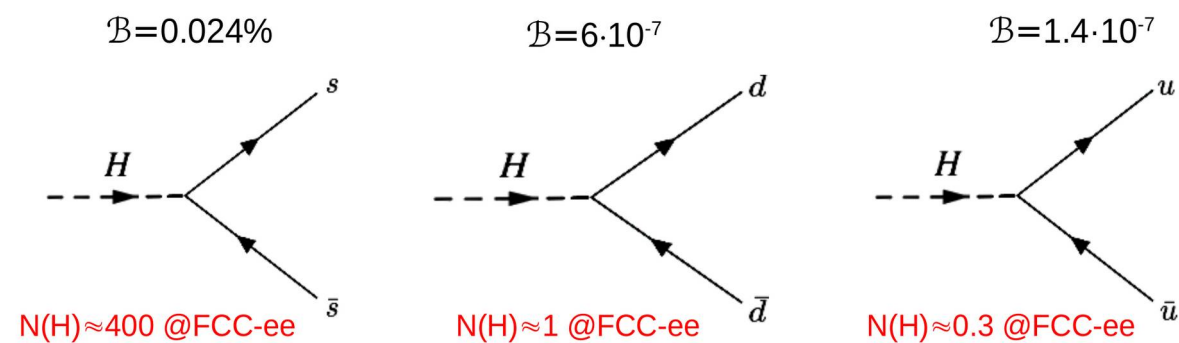
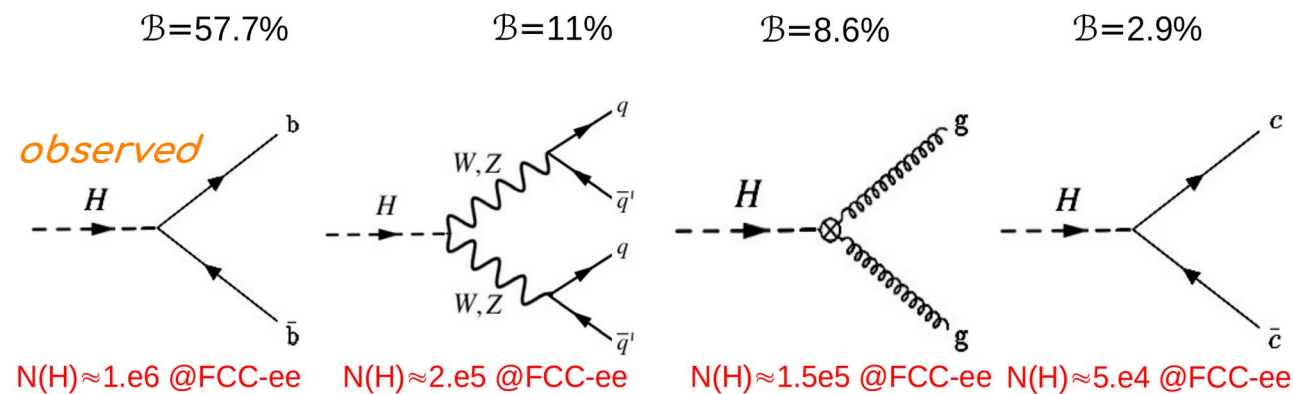
Perfect (=gen-level) momentum resolution

Final state	Muon 240 GeV	Electron 240 GeV	Combination 240 GeV
Nominal	3.92(4.74)	4.95(5.68)	3.07(3.97)
Inclusive	3.92(4.74)	4.95(5.68)	3.10(3.97)
Degradation electron resolution	3.92(4.74)	5.79(6.33)	3.24(4.12)
Magnetic field 3T	3.22(4.14)	4.11(4.83)	2.54(3.52)
Silicon tracker	5.11(5.73)	5.89(6.42)	3.86(4.55)
BES 6% uncertainty	3.92(4.79)	4.95(5.92)	3.07(3.98)
Disable BES	2.11(3.31)	2.93(3.88)	1.71(2.92)
Ideal resolution	3.12(3.95)	3.58(4.52)	2.42(3.40)
Freeze backgrounds	3.91(4.74)	4.95(5.67)	3.07(3.96)
Remove backgrounds	3.08(4.13)	3.51(4.58)	2.31(3.45)

Hadronic Decays

- 80% of the Higgs decays are fully hadronic
- challenging for LHC
- good prospects for FCC-ee thanks to clean environment and optimised tagging algorithms

	$\delta(\sigma \times \text{BR})$ [%]			
	Z(l \bar{l})H	Z($\nu\bar{\nu}$)H	Z(q \bar{q})H	Comb.
H \rightarrow bb	0.7	0.4	0.3	0.22
H \rightarrow cc	4.1	2.2	3.3	1.7
H \rightarrow ss	230	150	440	120
H \rightarrow gg	2.2	1.1	3.1	0.9



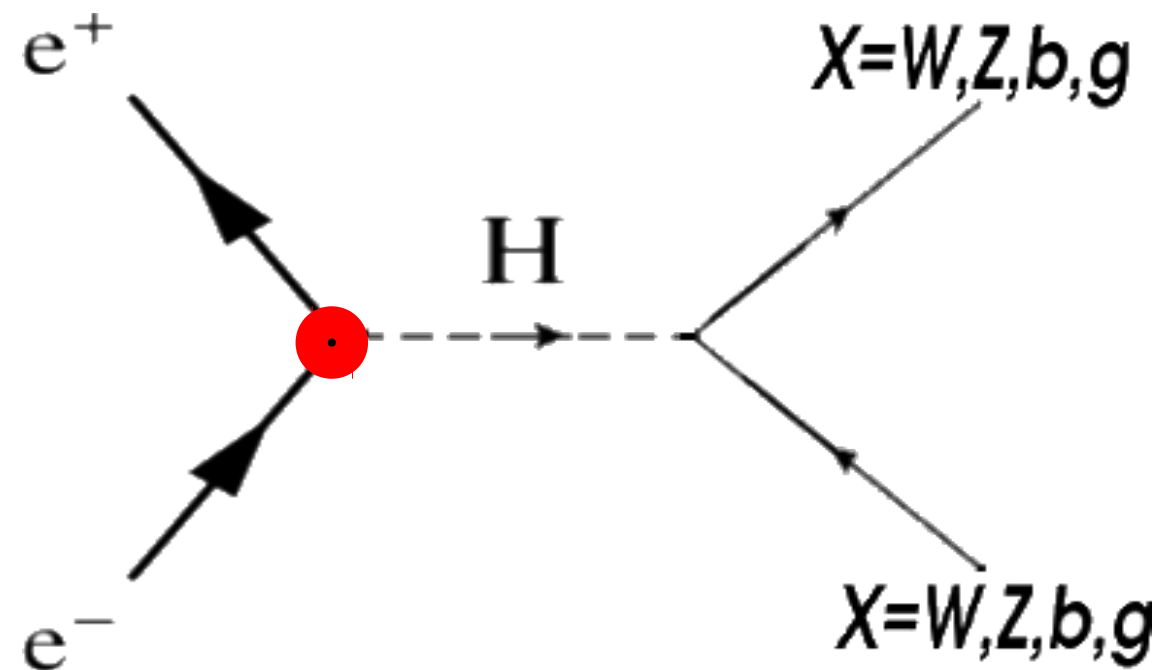
Solid measurements in 2nd generation

Interesting prospects for 1st generation and FCNC decays

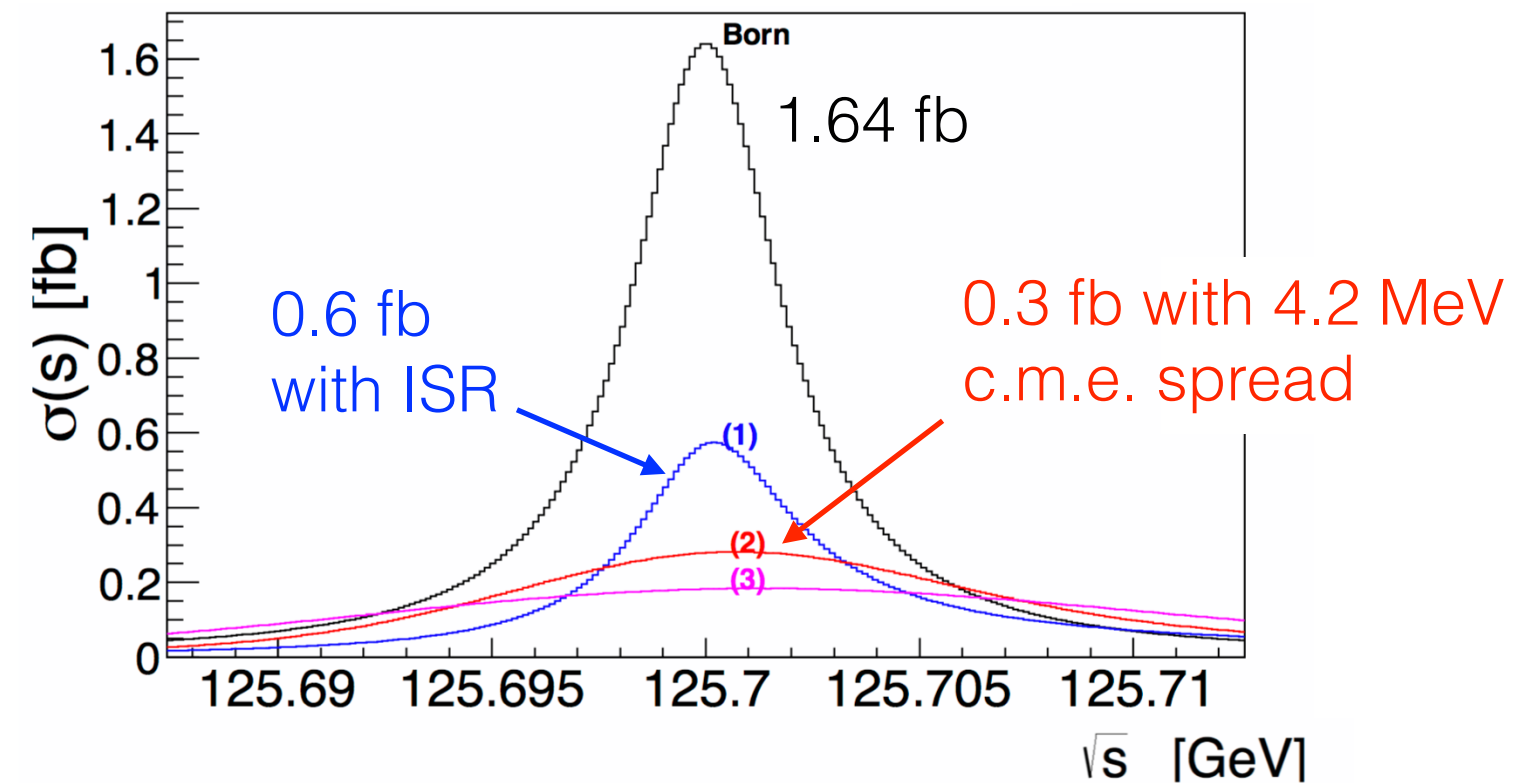
	$\sigma \times \text{BR}$ 95% CL	BR(SM)
H \rightarrow dd	1.4e-03	6e-07
H \rightarrow uu	1.5e-03	1.4e-07
H \rightarrow bs	3.7e-04	e-07
H \rightarrow bd	2.7e-04	e-09
H \rightarrow sd	7.7e-04	e-11
H \rightarrow cu	2.5e-04	e-20

Electron Yukawa

The high luminosity, the precise control of the beam \sqrt{s} , the clean reconstruction of final states make it possible to observe:



Jadach+, arXiv: 1509.02406



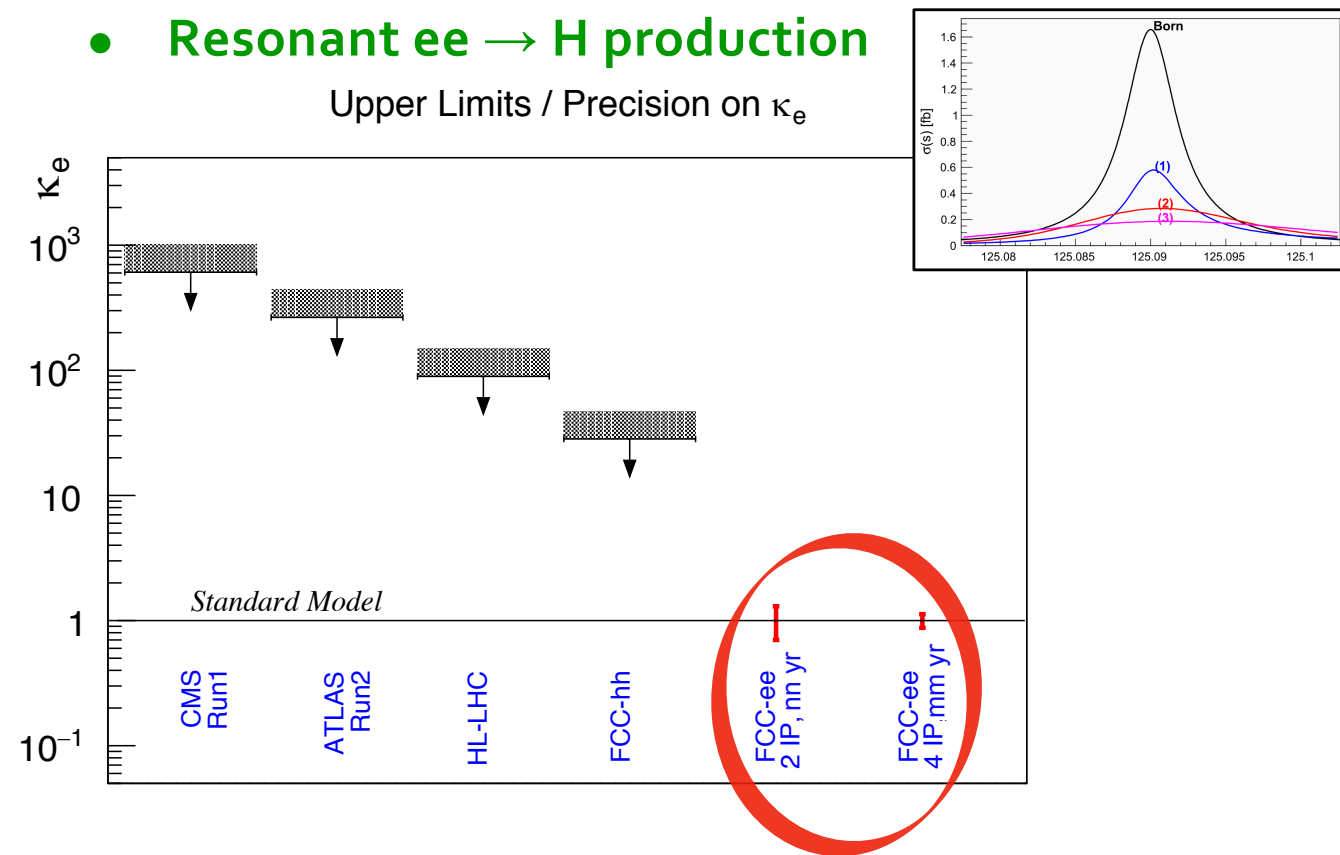
$$\sigma(e^+e^- \rightarrow H) = 1.64 \text{ fb}$$

$$\sigma_{\text{spread+ISR}}(e^+e^- \rightarrow H) = 0.17 \times \sigma(e^+e^- \rightarrow H) = 290 \text{ ab}$$

Electron Yukawa

The high luminosity, the precise control of the beam \sqrt{s} , the clean reconstruction of final states make it possible to observe:

- ◆ $20 \text{ ab}^{-1} / \text{year}$ at $\sqrt{s} = 125 \text{ GeV}$ (not in baseline FCC-ee)
- ◆ Monochromatization $\sigma_{\sqrt{s}} \sim 1\text{-}2 \times \Gamma_H \sim 6 \text{ to } 10 \text{ MeV}$
 - Resonant $ee \rightarrow H$ production



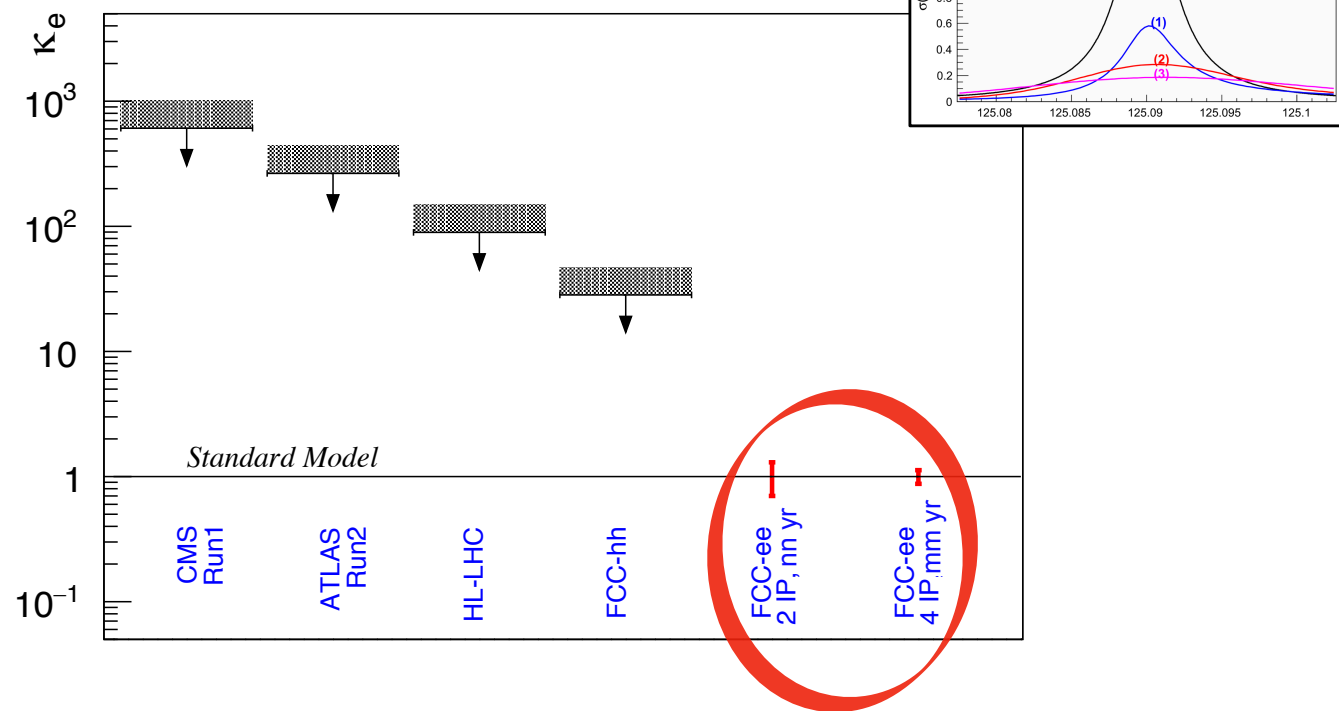
Electron Yukawa

The high luminosity, the precise control of the beam \sqrt{s} , the clean reconstruction of final states make it possible to observe:

- ◆ $20 \text{ ab}^{-1} / \text{year}$ at $\sqrt{s} = 125 \text{ GeV}$ (not in baseline FCC-ee)
- ◆ Monochromatization $\sigma_{\sqrt{s}} \sim 1-2 \times \Gamma_H \sim 6 \text{ to } 10 \text{ MeV}$

- Resonant $ee \rightarrow H$ production

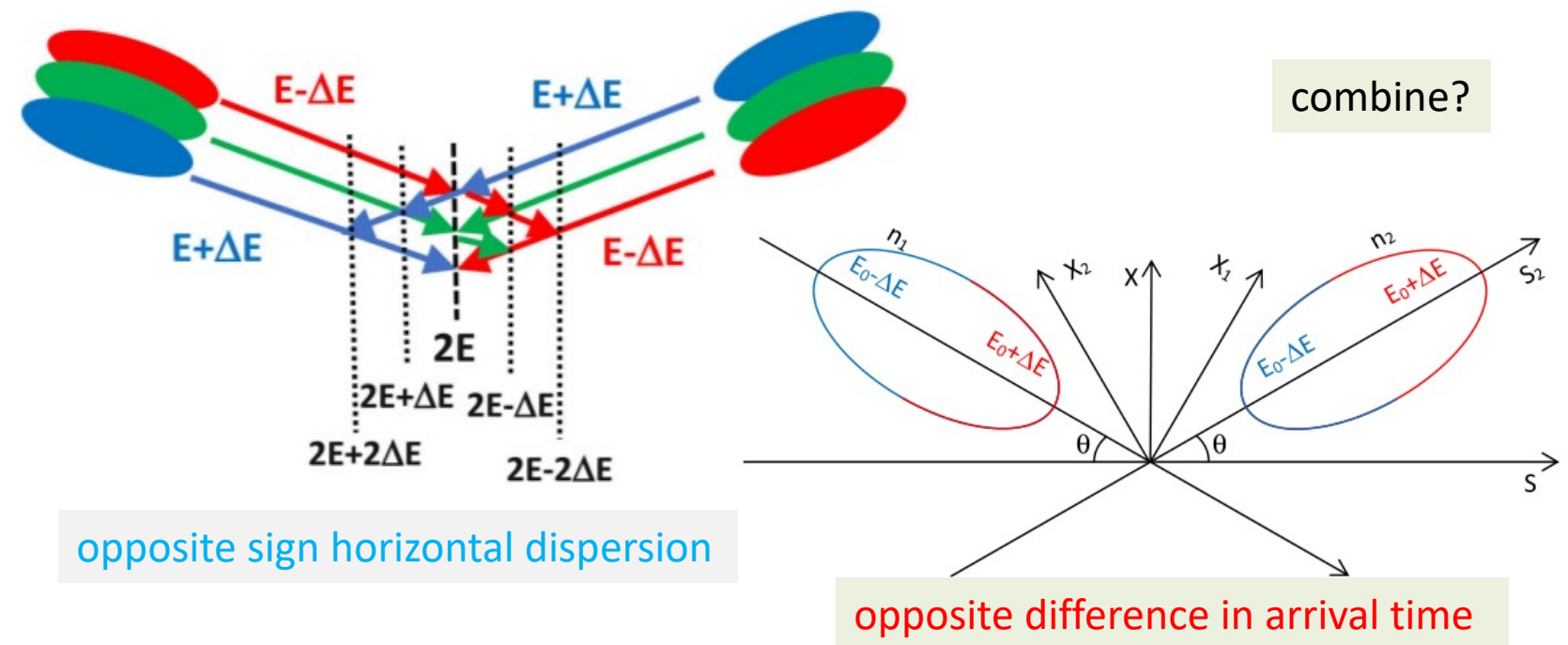
Upper Limits / Precision on κ_e



Monochromatisation

Monochromatization: **UNDER STUDY**

taking advantage of the separate e^+ and e^- rings, one can distribute in opposite way high and low energies in the beam (in x, z time)



combine?

opposite sign horizontal dispersion

opposite difference in arrival time

Electron Yukawa

The high luminosity, the precise control of the beam \sqrt{s} , the clean reconstruction of final states make it possible to observe:

- ◆ 20 ab^{-1} / year at $\sqrt{s} = 125 \text{ GeV}$ (not in baseline FCC-ee)
- ◆ Monochromatization $\sigma_{\sqrt{s}} \sim 1\text{-}2 \times \Gamma_H \sim 6 \text{ to } 10 \text{ MeV}$

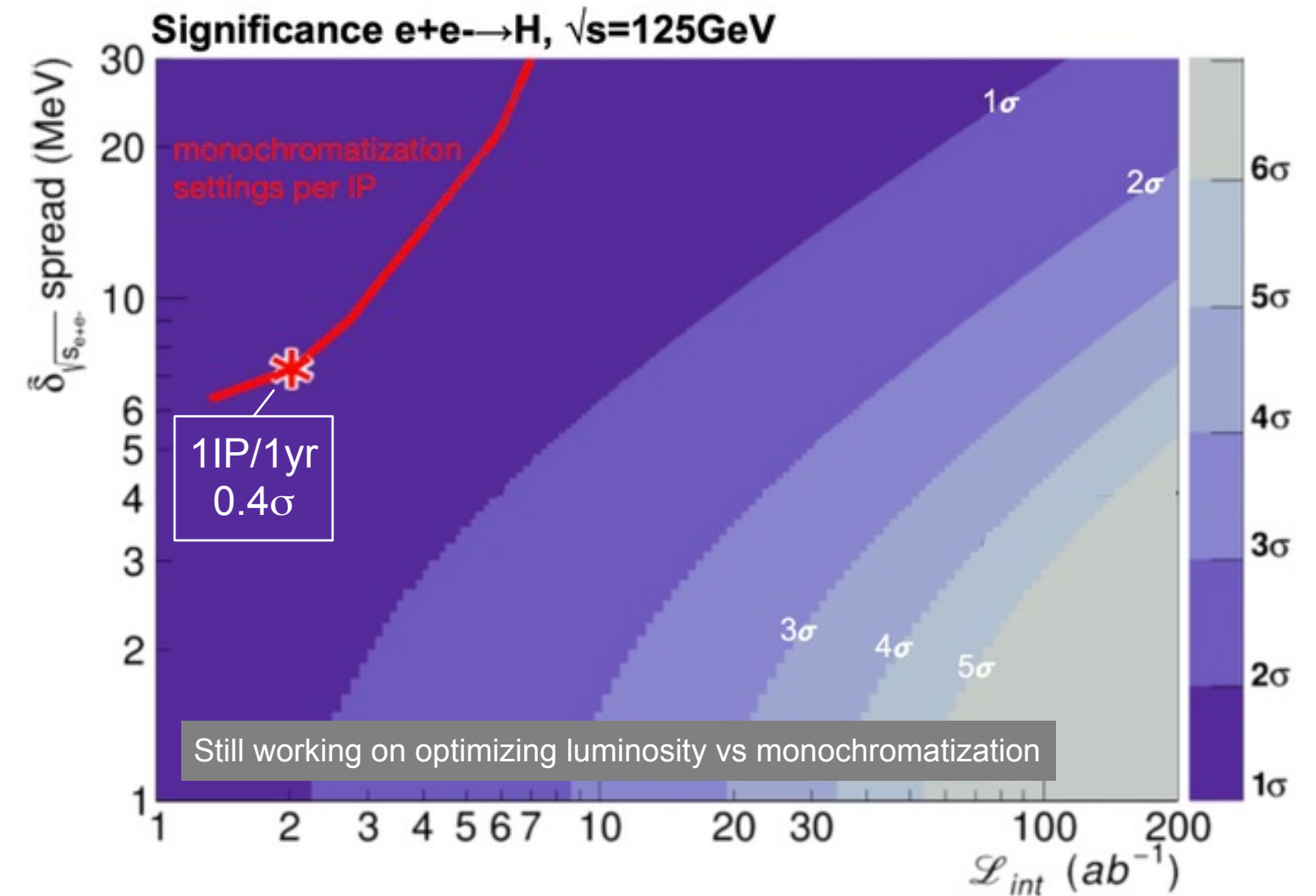
d'Enterria+. arXiv: 2107.02686

Higgs decay channel	\mathcal{B}	$\sigma \times \mathcal{B}$	Irreducible background	σ	S/B
$e^+e^- \rightarrow H \rightarrow b\bar{b}$	58.2%	164 ab	$e^+e^- \rightarrow b\bar{b}$	19 pb	$\mathcal{O}(10^{-5})$
$e^+e^- \rightarrow H \rightarrow gg$	8.2%	23 ab	$e^+e^- \rightarrow q\bar{q}$	61 pb	$\mathcal{O}(10^{-3})$
$e^+e^- \rightarrow H \rightarrow \tau\tau$	6.3%	18 ab	$e^+e^- \rightarrow \tau\tau$	10 pb	$\mathcal{O}(10^{-6})$
$e^+e^- \rightarrow H \rightarrow c\bar{c}$	2.9%	8.2 ab	$e^+e^- \rightarrow c\bar{c}$	22 pb	$\mathcal{O}(10^{-7})$
$e^+e^- \rightarrow H \rightarrow WW^* \rightarrow \ell\nu 2j$	$21.4\% \times 67.6\% \times 32.4\% \times 2$	26.5 ab	$e^+e^- \rightarrow WW^* \rightarrow \ell\nu 2j$	23 fb	$\mathcal{O}(10^{-3})$
$e^+e^- \rightarrow H \rightarrow WW^* \rightarrow 2\ell 2\nu$	$21.4\% \times 32.4\% \times 32.4\%$	6.4 ab	$e^+e^- \rightarrow WW^* \rightarrow 2\ell 2\nu$	5.6 fb	$\mathcal{O}(10^{-3})$
$e^+e^- \rightarrow H \rightarrow WW^* \rightarrow 4j$	$21.4\% \times 67.6\% \times 67.6\%$	27.6 ab	$e^+e^- \rightarrow WW^* \rightarrow 4j$	24 fb	$\mathcal{O}(10^{-3})$
$e^+e^- \rightarrow H \rightarrow ZZ^* \rightarrow 2j 2\nu$	$2.6\% \times 70\% \times 20\% \times 2$	2 ab	$e^+e^- \rightarrow ZZ^* \rightarrow 2j 2\nu$	273 ab	$\mathcal{O}(10^{-2})$
$e^+e^- \rightarrow H \rightarrow ZZ^* \rightarrow 2\ell 2j$	$2.6\% \times 70\% \times 10\% \times 2$	1 ab	$e^+e^- \rightarrow ZZ^* \rightarrow 2\ell 2j$	136 ab	$\mathcal{O}(10^{-2})$
$e^+e^- \rightarrow H \rightarrow ZZ^* \rightarrow 2\ell 2\nu$	$2.6\% \times 20\% \times 10\% \times 2$	0.3 ab	$e^+e^- \rightarrow ZZ^* \rightarrow 2\ell 2\nu$	39 ab	$\mathcal{O}(10^{-2})$
$e^+e^- \rightarrow H \rightarrow \gamma\gamma$	0.23%	0.65 ab	$e^+e^- \rightarrow \gamma\gamma$	79 pb	$\mathcal{O}(10^{-8})$

w. 10/ab

$H \rightarrow gg$	$H \rightarrow WW^* \rightarrow \ell\nu 2j; 2\ell 2\nu; 4j$	$H \rightarrow ZZ^* \rightarrow 2j 2\nu; 2\ell 2j; 2\ell 2\nu$	$H \rightarrow b\bar{b}$	$H \rightarrow \tau_{\text{had}}\tau_{\text{had}}; c\bar{c}; \gamma\gamma$	Combined
1.1 σ	$(0.53 \otimes 0.34 \otimes 0.13)\sigma$	$(0.32 \otimes 0.18 \otimes 0.05)\sigma$	0.13 σ	$< 0.02\sigma$	1.3 σ

w/ 10/ab: S~55, B~2400 \rightarrow 1.1 σ



Electron Yukawa

The high luminosity, the precise control of the beam \sqrt{s} , the clean reconstruction of final states make it possible to observe:

- ◆ 20 ab^{-1} / year at $\sqrt{s} = 125 \text{ GeV}$ (not in baseline FCC-ee)
- ◆ Monochromatization $\sigma_{\sqrt{s}} \sim 1\text{-}2 \times \Gamma_H \sim 6 \text{ to } 10 \text{ MeV}$

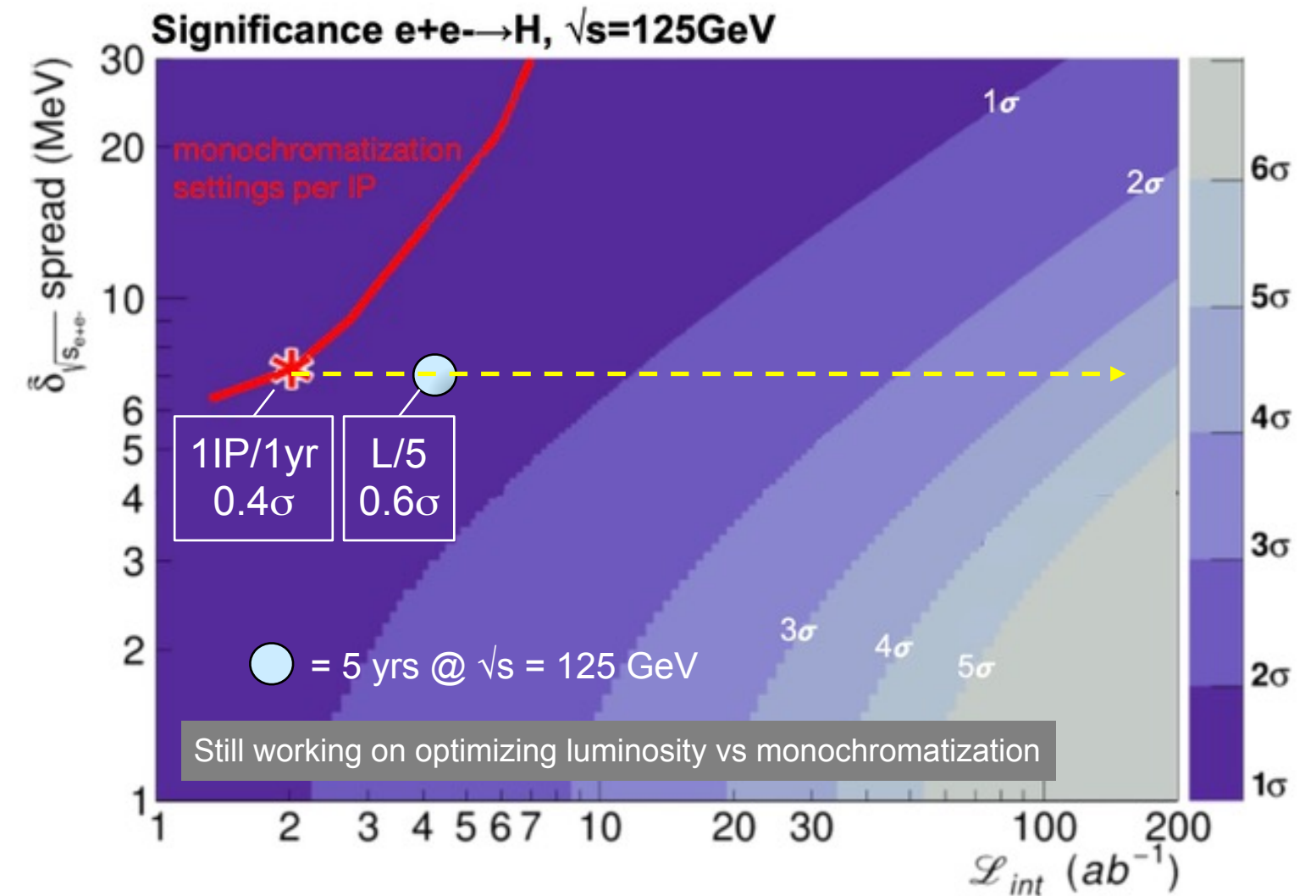
d'Enterria+. arXiv: 2107.02686

Higgs decay channel	\mathcal{B}	$\sigma \times \mathcal{B}$	Irreducible background	σ	S/B
$e^+e^- \rightarrow H \rightarrow b\bar{b}$	58.2%	164 ab	$e^+e^- \rightarrow b\bar{b}$	19 pb	$\mathcal{O}(10^{-5})$
$e^+e^- \rightarrow H \rightarrow gg$	8.2%	23 ab	$e^+e^- \rightarrow q\bar{q}$	61 pb	$\mathcal{O}(10^{-3})$
$e^+e^- \rightarrow H \rightarrow \tau\tau$	6.3%	18 ab	$e^+e^- \rightarrow \tau\tau$	10 pb	$\mathcal{O}(10^{-6})$
$e^+e^- \rightarrow H \rightarrow c\bar{c}$	2.9%	8.2 ab	$e^+e^- \rightarrow c\bar{c}$	22 pb	$\mathcal{O}(10^{-7})$
$e^+e^- \rightarrow H \rightarrow WW^* \rightarrow \ell\nu 2j$	$21.4\% \times 67.6\% \times 32.4\% \times 2$	26.5 ab	$e^+e^- \rightarrow WW^* \rightarrow \ell\nu 2j$	23 fb	$\mathcal{O}(10^{-3})$
$e^+e^- \rightarrow H \rightarrow WW^* \rightarrow 2\ell 2\nu$	$21.4\% \times 32.4\% \times 32.4\%$	6.4 ab	$e^+e^- \rightarrow WW^* \rightarrow 2\ell 2\nu$	5.6 fb	$\mathcal{O}(10^{-3})$
$e^+e^- \rightarrow H \rightarrow WW^* \rightarrow 4j$	$21.4\% \times 67.6\% \times 67.6\%$	27.6 ab	$e^+e^- \rightarrow WW^* \rightarrow 4j$	24 fb	$\mathcal{O}(10^{-3})$
$e^+e^- \rightarrow H \rightarrow ZZ^* \rightarrow 2j 2\nu$	$2.6\% \times 70\% \times 20\% \times 2$	2 ab	$e^+e^- \rightarrow ZZ^* \rightarrow 2j 2\nu$	273 ab	$\mathcal{O}(10^{-2})$
$e^+e^- \rightarrow H \rightarrow ZZ^* \rightarrow 2\ell 2j$	$2.6\% \times 70\% \times 10\% \times 2$	1 ab	$e^+e^- \rightarrow ZZ^* \rightarrow 2\ell 2j$	136 ab	$\mathcal{O}(10^{-2})$
$e^+e^- \rightarrow H \rightarrow ZZ^* \rightarrow 2\ell 2\nu$	$2.6\% \times 20\% \times 10\% \times 2$	0.3 ab	$e^+e^- \rightarrow ZZ^* \rightarrow 2\ell 2\nu$	39 ab	$\mathcal{O}(10^{-2})$
$e^+e^- \rightarrow H \rightarrow \gamma\gamma$	0.23%	0.65 ab	$e^+e^- \rightarrow \gamma\gamma$	79 pb	$\mathcal{O}(10^{-8})$

w. 10/ab

$H \rightarrow gg$	$H \rightarrow WW^* \rightarrow \ell\nu 2j; 2\ell 2\nu; 4j$	$H \rightarrow ZZ^* \rightarrow 2j 2\nu; 2\ell 2j; 2\ell 2\nu$	$H \rightarrow b\bar{b}$	$H \rightarrow \tau_{\text{had}}\tau_{\text{had}}; c\bar{c}; \gamma\gamma$	Combined
1.1 σ	$(0.53 \otimes 0.34 \otimes 0.13)\sigma$	$(0.32 \otimes 0.18 \otimes 0.05)\sigma$	0.13 σ	$< 0.02\sigma$	1.3 σ

w/ 10/ab: S~55, B~2400 \rightarrow 1.1 σ



Electron Yukawa

The high luminosity, the precise control of the beam \sqrt{s} , the clean reconstruction of final states make it possible to observe:

- ◆ 20 ab^{-1} / year at $\sqrt{s} = 125 \text{ GeV}$ (not in baseline FCC-ee)
- ◆ Monochromatization $\sigma_{\sqrt{s}} \sim 1\text{-}2 \times \Gamma_H \sim 6 \text{ to } 10 \text{ MeV}$

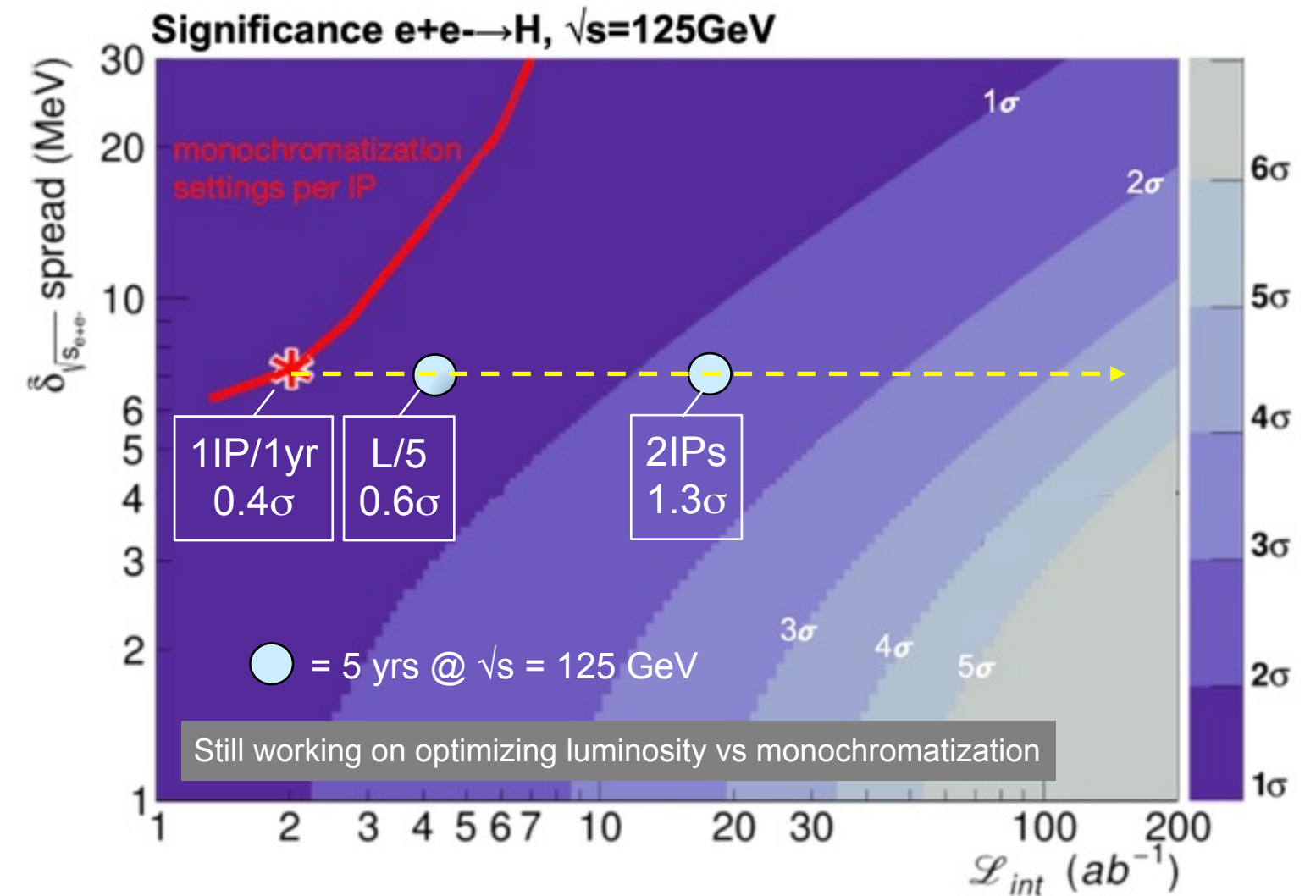
d'Enterria+. arXiv: 2107.02686

Higgs decay channel	\mathcal{B}	$\sigma \times \mathcal{B}$	Irreducible background	σ	S/B
$e^+e^- \rightarrow H \rightarrow b\bar{b}$	58.2%	164 ab	$e^+e^- \rightarrow b\bar{b}$	19 pb	$\mathcal{O}(10^{-5})$
$e^+e^- \rightarrow H \rightarrow gg$	8.2%	23 ab	$e^+e^- \rightarrow q\bar{q}$	61 pb	$\mathcal{O}(10^{-3})$
$e^+e^- \rightarrow H \rightarrow \tau\tau$	6.3%	18 ab	$e^+e^- \rightarrow \tau\tau$	10 pb	$\mathcal{O}(10^{-6})$
$e^+e^- \rightarrow H \rightarrow c\bar{c}$	2.9%	8.2 ab	$e^+e^- \rightarrow c\bar{c}$	22 pb	$\mathcal{O}(10^{-7})$
$e^+e^- \rightarrow H \rightarrow WW^* \rightarrow \ell\nu 2j$	$21.4\% \times 67.6\% \times 32.4\% \times 2$	26.5 ab	$e^+e^- \rightarrow WW^* \rightarrow \ell\nu 2j$	23 fb	$\mathcal{O}(10^{-3})$
$e^+e^- \rightarrow H \rightarrow WW^* \rightarrow 2\ell 2\nu$	$21.4\% \times 32.4\% \times 32.4\%$	6.4 ab	$e^+e^- \rightarrow WW^* \rightarrow 2\ell 2\nu$	5.6 fb	$\mathcal{O}(10^{-3})$
$e^+e^- \rightarrow H \rightarrow WW^* \rightarrow 4j$	$21.4\% \times 67.6\% \times 67.6\%$	27.6 ab	$e^+e^- \rightarrow WW^* \rightarrow 4j$	24 fb	$\mathcal{O}(10^{-3})$
$e^+e^- \rightarrow H \rightarrow ZZ^* \rightarrow 2j 2\nu$	$2.6\% \times 70\% \times 20\% \times 2$	2 ab	$e^+e^- \rightarrow ZZ^* \rightarrow 2j 2\nu$	273 ab	$\mathcal{O}(10^{-2})$
$e^+e^- \rightarrow H \rightarrow ZZ^* \rightarrow 2\ell 2j$	$2.6\% \times 70\% \times 10\% \times 2$	1 ab	$e^+e^- \rightarrow ZZ^* \rightarrow 2\ell 2j$	136 ab	$\mathcal{O}(10^{-2})$
$e^+e^- \rightarrow H \rightarrow ZZ^* \rightarrow 2\ell 2\nu$	$2.6\% \times 20\% \times 10\% \times 2$	0.3 ab	$e^+e^- \rightarrow ZZ^* \rightarrow 2\ell 2\nu$	39 ab	$\mathcal{O}(10^{-2})$
$e^+e^- \rightarrow H \rightarrow \gamma\gamma$	0.23%	0.65 ab	$e^+e^- \rightarrow \gamma\gamma$	79 pb	$\mathcal{O}(10^{-8})$

w. 10/ab

$H \rightarrow gg$	$H \rightarrow WW^* \rightarrow \ell\nu 2j; 2\ell 2\nu; 4j$	$H \rightarrow ZZ^* \rightarrow 2j 2\nu; 2\ell 2j; 2\ell 2\nu$	$H \rightarrow b\bar{b}$	$H \rightarrow \tau_{\text{had}}\tau_{\text{had}}; c\bar{c}; \gamma\gamma$	Combined
1.1 σ	$(0.53 \otimes 0.34 \otimes 0.13)\sigma$	$(0.32 \otimes 0.18 \otimes 0.05)\sigma$	0.13 σ	$< 0.02\sigma$	1.3 σ

w/ 10/ab: S~55, B~2400 \rightarrow 1.1 σ



Electron Yukawa

The high luminosity, the precise control of the beam \sqrt{s} , the clean reconstruction of final states make it possible to observe:

- ◆ 20 ab^{-1} / year at $\sqrt{s} = 125 \text{ GeV}$ (not in baseline FCC-ee)
- ◆ Monochromatization $\sigma_{\sqrt{s}} \sim 1\text{-}2 \times \Gamma_H \sim 6 \text{ to } 10 \text{ MeV}$

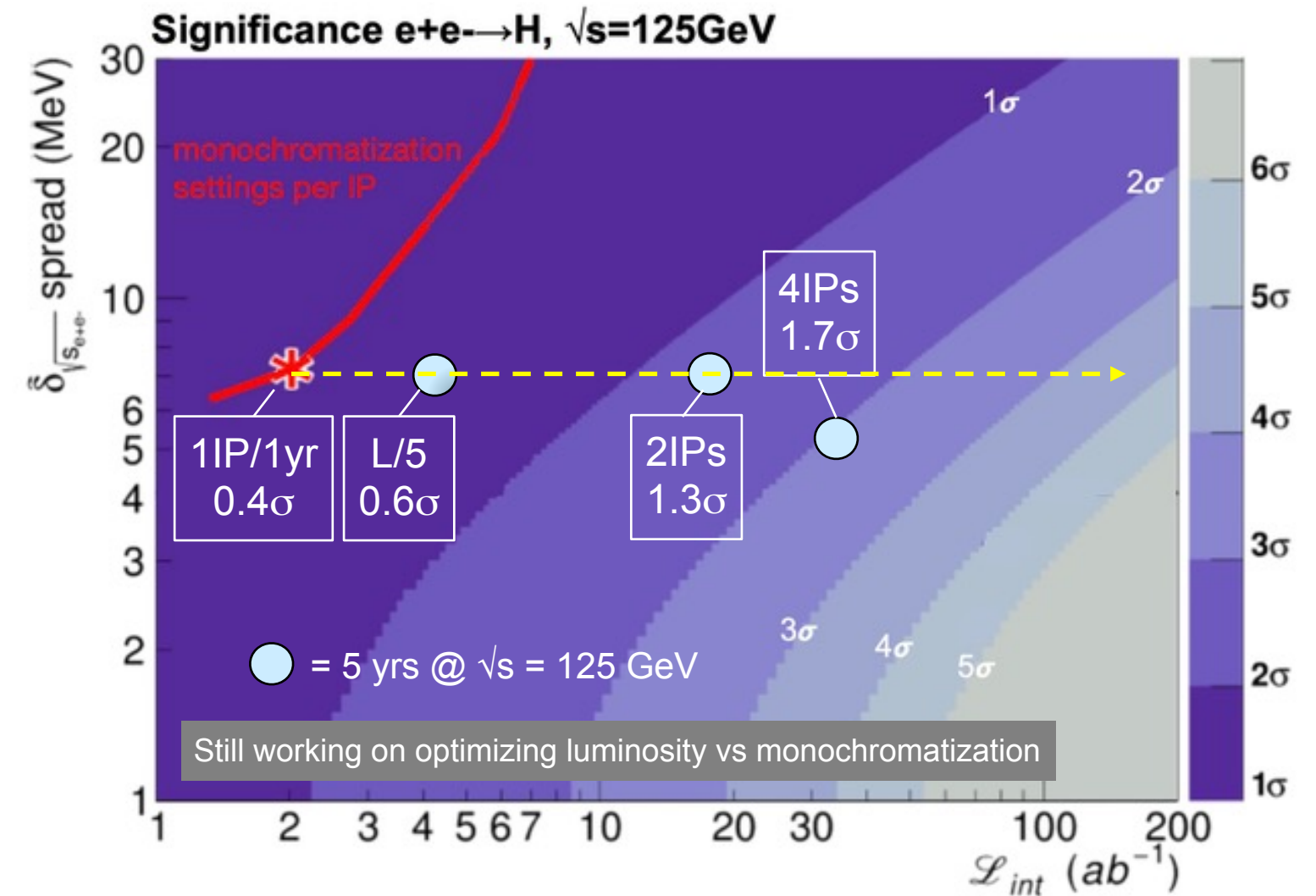
d'Enterria+. arXiv: 2107.02686

Higgs decay channel	\mathcal{B}	$\sigma \times \mathcal{B}$	Irreducible background	σ	S/B
$e^+e^- \rightarrow H \rightarrow b\bar{b}$	58.2%	164 ab	$e^+e^- \rightarrow b\bar{b}$	19 pb	$\mathcal{O}(10^{-5})$
$e^+e^- \rightarrow H \rightarrow gg$	8.2%	23 ab	$e^+e^- \rightarrow q\bar{q}$	61 pb	$\mathcal{O}(10^{-3})$
$e^+e^- \rightarrow H \rightarrow \tau\tau$	6.3%	18 ab	$e^+e^- \rightarrow \tau\tau$	10 pb	$\mathcal{O}(10^{-6})$
$e^+e^- \rightarrow H \rightarrow c\bar{c}$	2.9%	8.2 ab	$e^+e^- \rightarrow c\bar{c}$	22 pb	$\mathcal{O}(10^{-7})$
$e^+e^- \rightarrow H \rightarrow WW^* \rightarrow \ell\nu 2j$	$21.4\% \times 67.6\% \times 32.4\% \times 2$	26.5 ab	$e^+e^- \rightarrow WW^* \rightarrow \ell\nu 2j$	23 fb	$\mathcal{O}(10^{-3})$
$e^+e^- \rightarrow H \rightarrow WW^* \rightarrow 2\ell 2\nu$	$21.4\% \times 32.4\% \times 32.4\%$	6.4 ab	$e^+e^- \rightarrow WW^* \rightarrow 2\ell 2\nu$	5.6 fb	$\mathcal{O}(10^{-3})$
$e^+e^- \rightarrow H \rightarrow WW^* \rightarrow 4j$	$21.4\% \times 67.6\% \times 67.6\%$	27.6 ab	$e^+e^- \rightarrow WW^* \rightarrow 4j$	24 fb	$\mathcal{O}(10^{-3})$
$e^+e^- \rightarrow H \rightarrow ZZ^* \rightarrow 2j 2\nu$	$2.6\% \times 70\% \times 20\% \times 2$	2 ab	$e^+e^- \rightarrow ZZ^* \rightarrow 2j 2\nu$	273 ab	$\mathcal{O}(10^{-2})$
$e^+e^- \rightarrow H \rightarrow ZZ^* \rightarrow 2\ell 2j$	$2.6\% \times 70\% \times 10\% \times 2$	1 ab	$e^+e^- \rightarrow ZZ^* \rightarrow 2\ell 2j$	136 ab	$\mathcal{O}(10^{-2})$
$e^+e^- \rightarrow H \rightarrow ZZ^* \rightarrow 2\ell 2\nu$	$2.6\% \times 20\% \times 10\% \times 2$	0.3 ab	$e^+e^- \rightarrow ZZ^* \rightarrow 2\ell 2\nu$	39 ab	$\mathcal{O}(10^{-2})$
$e^+e^- \rightarrow H \rightarrow \gamma\gamma$	0.23%	0.65 ab	$e^+e^- \rightarrow \gamma\gamma$	79 pb	$\mathcal{O}(10^{-8})$

w. 10/ab

$H \rightarrow gg$	$H \rightarrow WW^* \rightarrow \ell\nu 2j; 2\ell 2\nu; 4j$	$H \rightarrow ZZ^* \rightarrow 2j 2\nu; 2\ell 2j; 2\ell 2\nu$	$H \rightarrow b\bar{b}$	$H \rightarrow \tau_{\text{had}}\tau_{\text{had}}; c\bar{c}; \gamma\gamma$	Combined
1.1 σ	$(0.53 \otimes 0.34 \otimes 0.13)\sigma$	$(0.32 \otimes 0.18 \otimes 0.05)\sigma$	0.13 σ	$< 0.02\sigma$	1.3 σ

w/ 10/ab: S~55, B~2400 \rightarrow 1.1 σ



Electron Yukawa

The high luminosity, the precise control of the beam \sqrt{s} , the clean reconstruction of final states make it possible to observe:

- ◆ 20 ab^{-1} / year at $\sqrt{s} = 125 \text{ GeV}$ (not in baseline FCC-ee)
- ◆ Monochromatization $\sigma_{\sqrt{s}} \sim 1\text{-}2 \times \Gamma_H \sim 6 \text{ to } 10 \text{ MeV}$

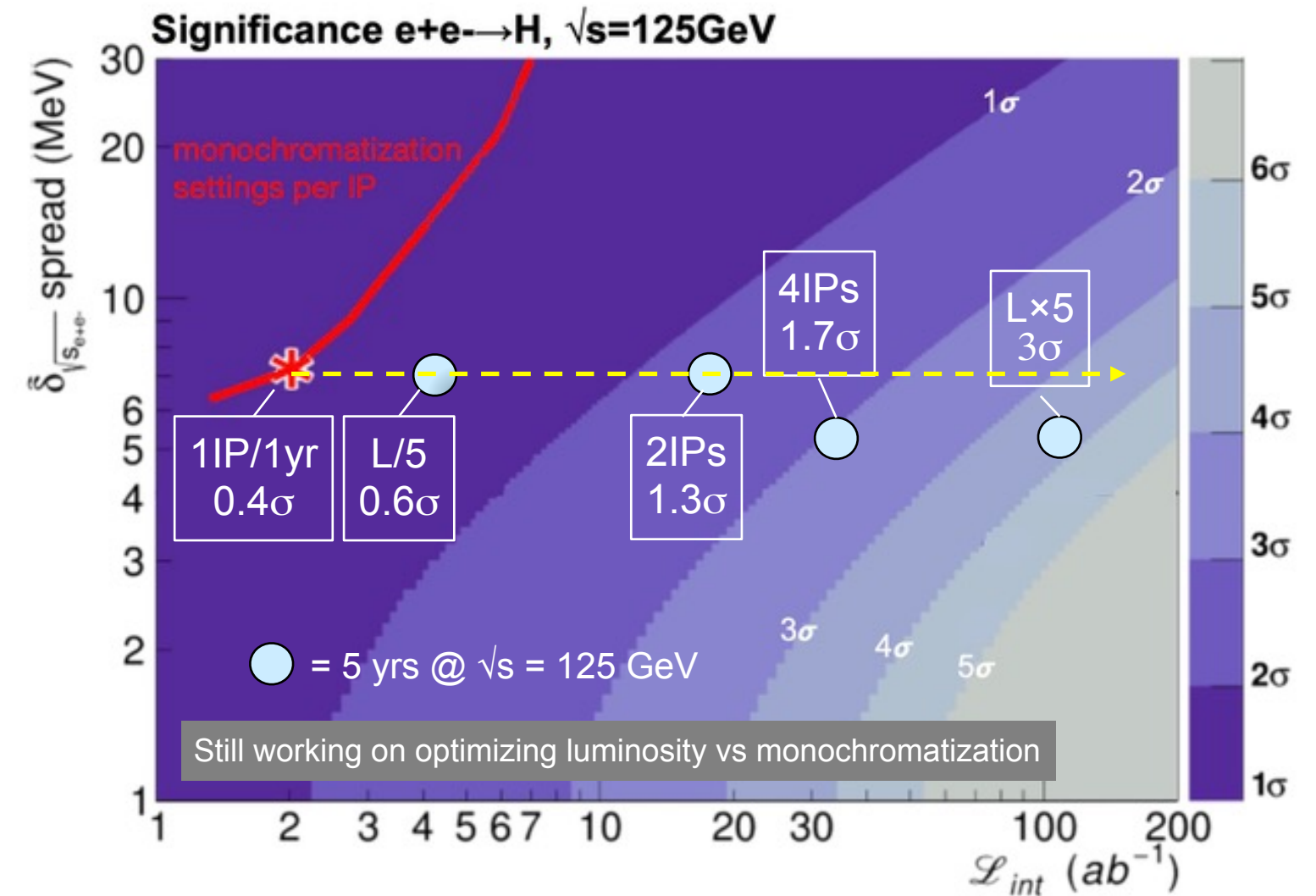
d'Enterria+. arXiv: 2107.02686

Higgs decay channel	\mathcal{B}	$\sigma \times \mathcal{B}$	Irreducible background	σ	S/B
$e^+e^- \rightarrow H \rightarrow b\bar{b}$	58.2%	164 ab	$e^+e^- \rightarrow b\bar{b}$	19 pb	$\mathcal{O}(10^{-5})$
$e^+e^- \rightarrow H \rightarrow gg$	8.2%	23 ab	$e^+e^- \rightarrow q\bar{q}$	61 pb	$\mathcal{O}(10^{-3})$
$e^+e^- \rightarrow H \rightarrow \tau\tau$	6.3%	18 ab	$e^+e^- \rightarrow \tau\tau$	10 pb	$\mathcal{O}(10^{-6})$
$e^+e^- \rightarrow H \rightarrow c\bar{c}$	2.9%	8.2 ab	$e^+e^- \rightarrow c\bar{c}$	22 pb	$\mathcal{O}(10^{-7})$
$e^+e^- \rightarrow H \rightarrow WW^* \rightarrow \ell\nu 2j$	$21.4\% \times 67.6\% \times 32.4\% \times 2$	26.5 ab	$e^+e^- \rightarrow WW^* \rightarrow \ell\nu 2j$	23 fb	$\mathcal{O}(10^{-3})$
$e^+e^- \rightarrow H \rightarrow WW^* \rightarrow 2\ell 2\nu$	$21.4\% \times 32.4\% \times 32.4\%$	6.4 ab	$e^+e^- \rightarrow WW^* \rightarrow 2\ell 2\nu$	5.6 fb	$\mathcal{O}(10^{-3})$
$e^+e^- \rightarrow H \rightarrow WW^* \rightarrow 4j$	$21.4\% \times 67.6\% \times 67.6\%$	27.6 ab	$e^+e^- \rightarrow WW^* \rightarrow 4j$	24 fb	$\mathcal{O}(10^{-3})$
$e^+e^- \rightarrow H \rightarrow ZZ^* \rightarrow 2j 2\nu$	$2.6\% \times 70\% \times 20\% \times 2$	2 ab	$e^+e^- \rightarrow ZZ^* \rightarrow 2j 2\nu$	273 ab	$\mathcal{O}(10^{-2})$
$e^+e^- \rightarrow H \rightarrow ZZ^* \rightarrow 2\ell 2j$	$2.6\% \times 70\% \times 10\% \times 2$	1 ab	$e^+e^- \rightarrow ZZ^* \rightarrow 2\ell 2j$	136 ab	$\mathcal{O}(10^{-2})$
$e^+e^- \rightarrow H \rightarrow ZZ^* \rightarrow 2\ell 2\nu$	$2.6\% \times 20\% \times 10\% \times 2$	0.3 ab	$e^+e^- \rightarrow ZZ^* \rightarrow 2\ell 2\nu$	39 ab	$\mathcal{O}(10^{-2})$
$e^+e^- \rightarrow H \rightarrow \gamma\gamma$	0.23%	0.65 ab	$e^+e^- \rightarrow \gamma\gamma$	79 pb	$\mathcal{O}(10^{-8})$

w. 10/ab

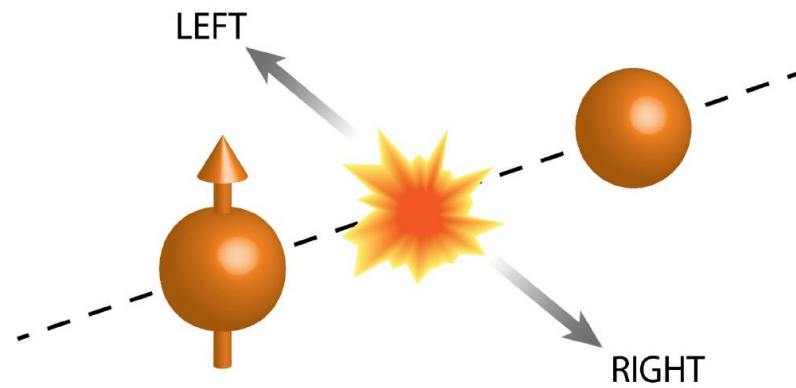
$H \rightarrow gg$	$H \rightarrow WW^* \rightarrow \ell\nu 2j; 2\ell 2\nu; 4j$	$H \rightarrow ZZ^* \rightarrow 2j 2\nu; 2\ell 2j; 2\ell 2\nu$	$H \rightarrow b\bar{b}$	$H \rightarrow \tau_{\text{had}}\tau_{\text{had}}; c\bar{c}; \gamma\gamma$	Combined
1.1 σ	$(0.53 \otimes 0.34 \otimes 0.13)\sigma$	$(0.32 \otimes 0.18 \otimes 0.05)\sigma$	0.13 σ	< 0.02 σ	1.3 σ

w/ 10/ab: S~55, B~2400 \rightarrow 1.1 σ



Electron Yukawa

A recent pheno study ([Boughezal et al 2407.12975](#)) shows that transverse spin asymmetries can increase the sensitivity to the electron Yukawa



$$A = \frac{N}{D}$$

Electron polarized,
positron unpolarized (SP⁰):

$$N = \frac{1}{2}(\sigma^{+0} - \sigma^{-0})$$

$$D = \frac{1}{2}(\sigma^{+0} + \sigma^{-0})$$

Electron transversely
polarized, positron
longitudinally polarized (DP):

$$N = \frac{1}{4}(\sigma^{++} - \sigma^{+-} - \sigma^{-+} + \sigma^{--})$$

$$D = \frac{1}{4}(\sigma^{++} + \sigma^{+-} + \sigma^{-+} + \sigma^{--})$$

Electron transversely
polarized, positron
longitudinally polarized (SP⁺):

$$N = \frac{1}{2}(\sigma^{++} - \sigma^{-+})$$

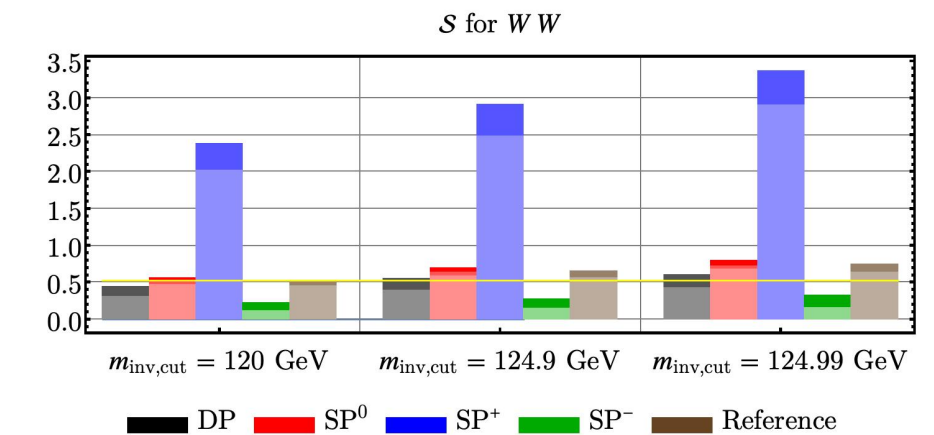
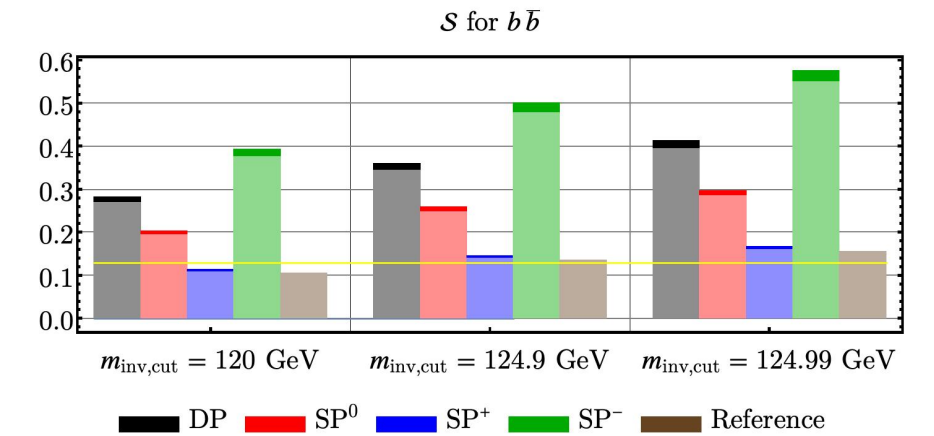
$$D = \frac{1}{2}(\sigma^{++} + \sigma^{-+})$$

Electron transversely
polarized, positron
longitudinally polarized (SP⁻):

$$N = \frac{1}{2}(\sigma^{+-} - \sigma^{--})$$

$$D = \frac{1}{2}(\sigma^{+-} + \sigma^{--})$$

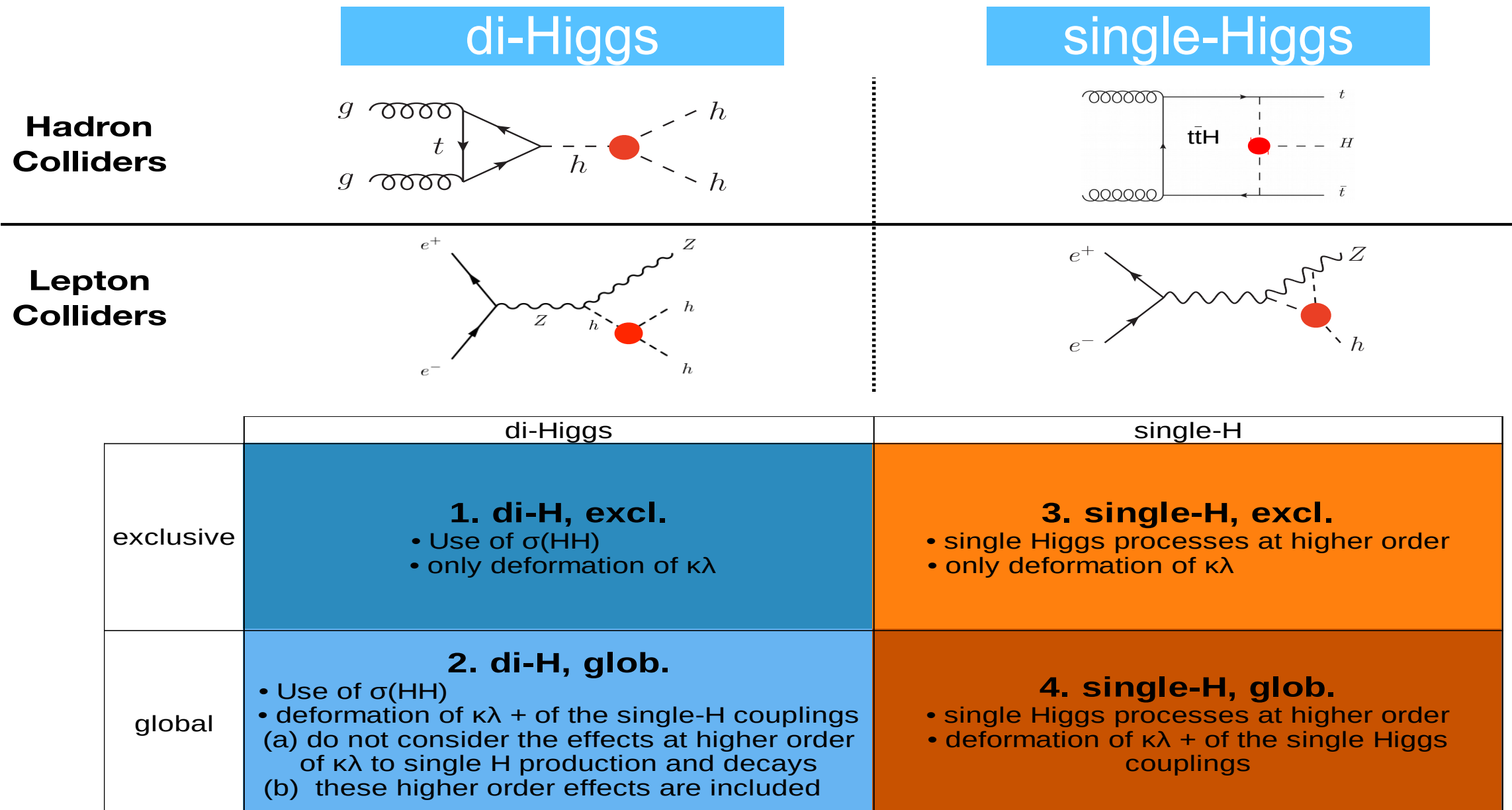
8



Major improvements of up to factors of 6 possible for $b\bar{b}$ and WW (doesn't work for gg)

Higgs Self-Coupling

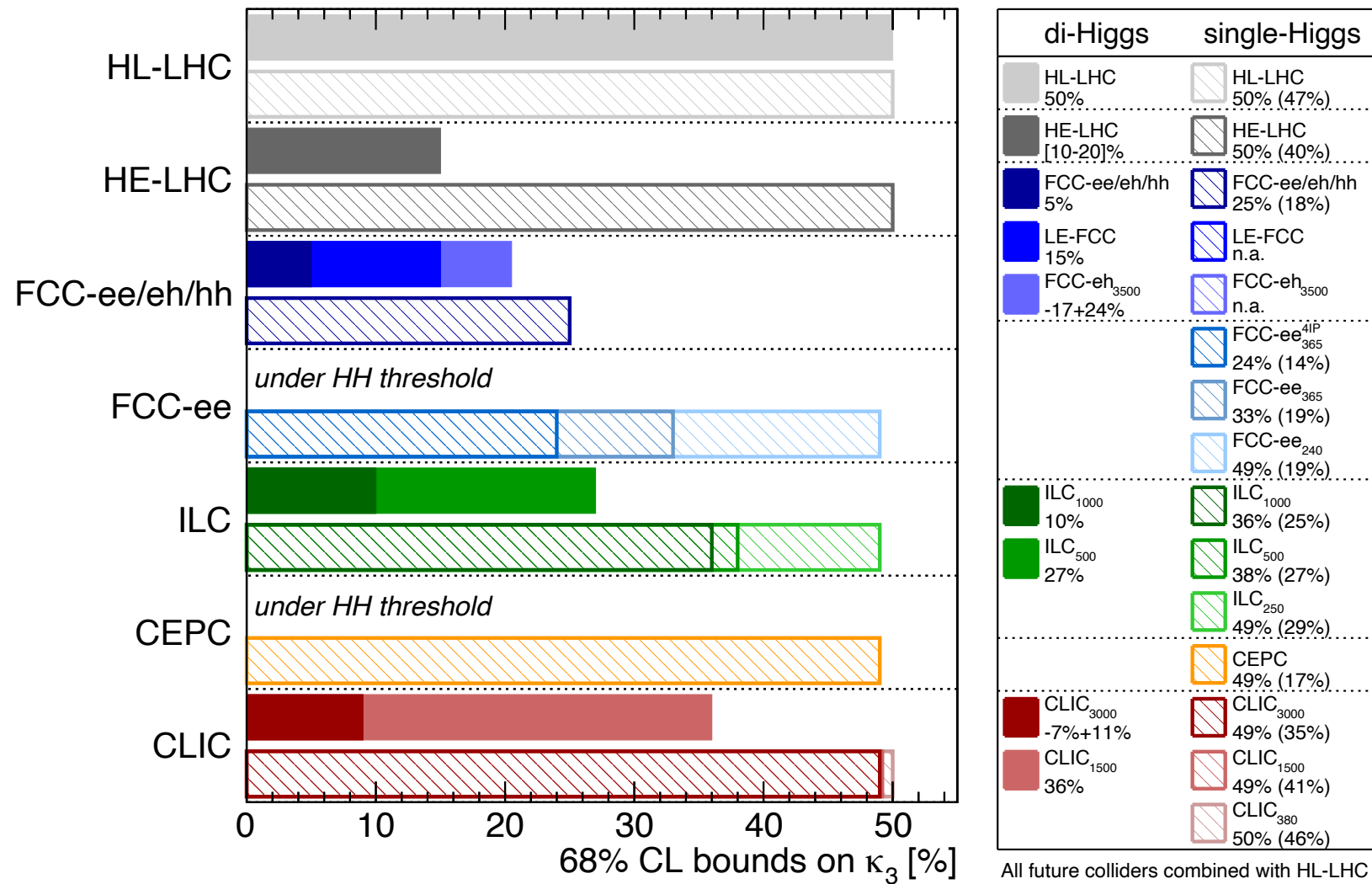
How much can it deviate from SM given the tight constraints on other Higgs couplings?
Do we need to reach HH production threshold to constrain h^3 coupling?



ECFA Higgs study group '19

Higgs Self-Coupling

Higgs@FC WG November 2019



1

Don't need to reach HH threshold to have access to h^3 .
Runs at different energies are essential (e.g. 240 and 365 GeV)

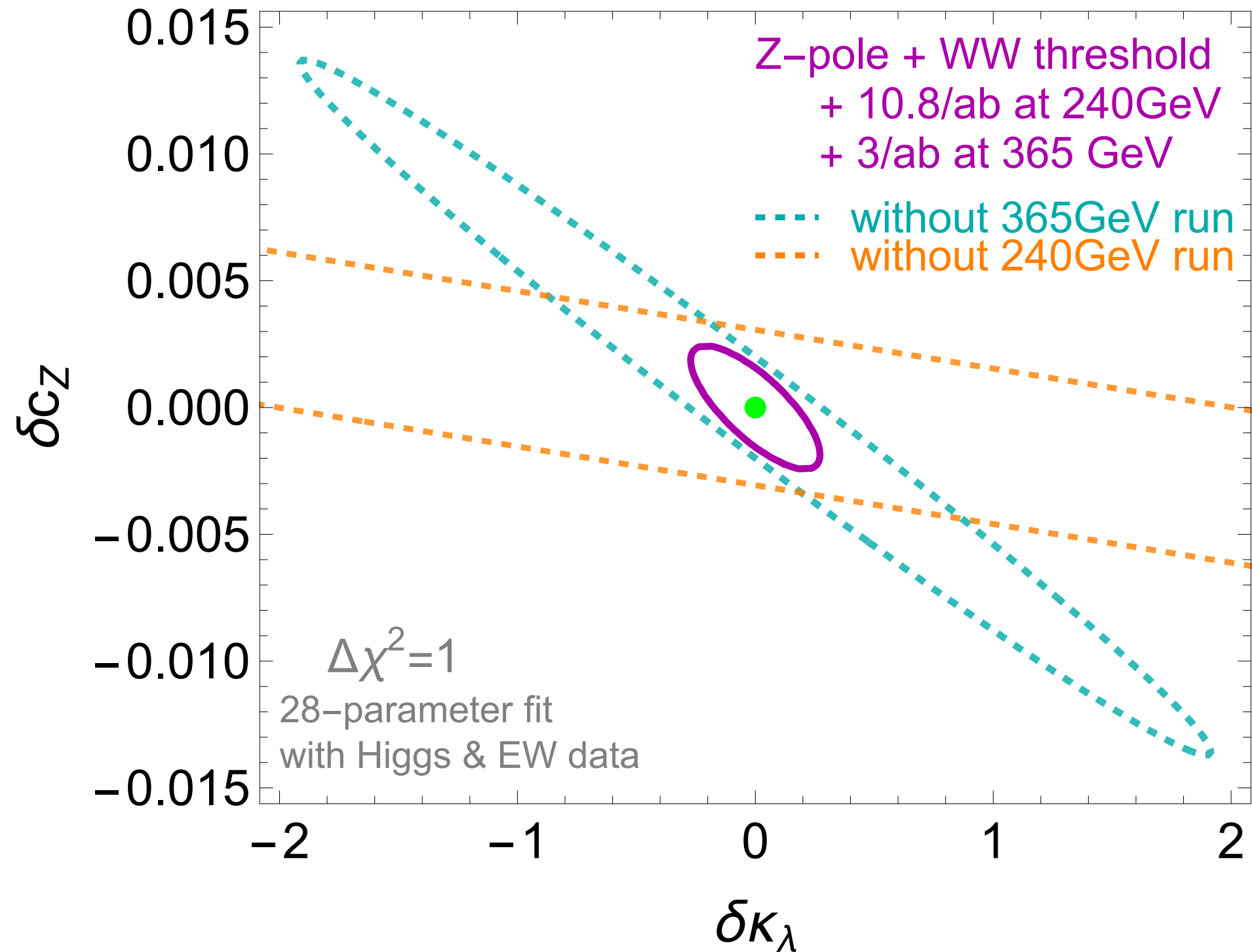
2

The determination of h^3 at FCC-hh relies on HH channel, for which FCC-ee is of little direct help. But the extraction of h^3 requires precise knowledge of y_t .
 $1\% y_t \leftrightarrow 5\% h^3$
Precision measurement of y_t needs FCC-ee.

50% sensitivity: establish that $h^3 \neq 0$ at 95%CL
20% sensitivity: 5σ discovery of the SM h^3 coupling
5% sensitivity: getting sensitive to quantum corrections to Higgs potential

Higgs Self-Coupling

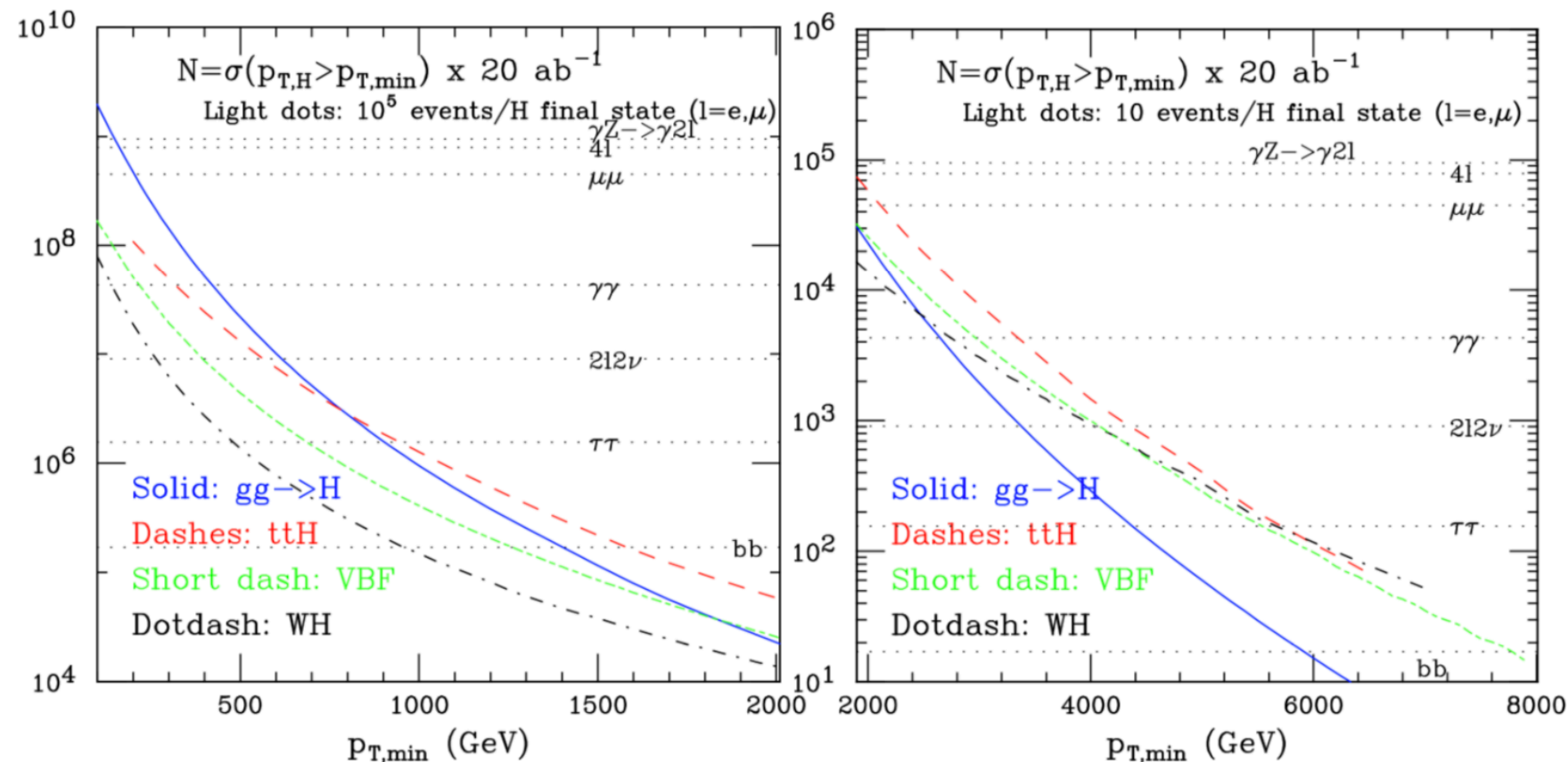
FCC-ee, from SMEFT global fit



Higgs @ FCC-hh.

	ggH (N ³ LO)	VBF (N ² LO)	WH (N ² LO)	ZH (N ² LO)	t \bar{t} H (N ² LO)	HH (NLO)
N100	24×10^9	2.1×10^9	4.6×10^8	3.3×10^8	9.6×10^8	3.6×10^7
N100/N14	180	170	100	110	530	390

(N100 = $\sigma_{100 \text{ TeV}} \times 30 \text{ ab}^{-1}$ & N14 = $\sigma_{14 \text{ TeV}} \times 3 \text{ ab}^{-1}$)



- Large rate ($> 10^{10} \text{ H}$, $> 10^7 \text{ HH}$)
 - unique sensitivity to **rare decays**
 - few % sensitivity to **self-coupling**
- Explore extreme phase space:
 - e.g. 10^6 H w/ $p_T > 1 \text{ TeV}$
 - clean samples with high S/B
 - small systematics

Electroweak Physics

Observable	present value	±	error	FCC-ee Stat.	FCC-ee Syst.	Comment and leading error
m_Z (keV)	91186700	±	2200	4	100	From Z line shape scan Beam energy calibration
Γ_Z (keV)	2495200	±	2300	4	25	From Z line shape scan Beam energy calibration
$\sin^2 \theta_W^{\text{eff}} (\times 10^6)$	231480	±	160	2	2.4	From $A_{\text{FB}}^{\mu\mu}$ at Z peak Beam energy calibration
$1/\alpha_{\text{QED}}(m_Z^2)(\times 10^3)$	128952	±	14	3	small	From $A_{\text{FB}}^{\mu\mu}$ off peak QED&EW errors dominate
$R_\ell^Z (\times 10^3)$	20767	±	25	0.06	0.2-1	Ratio of hadrons to leptons Acceptance for leptons
$\alpha_s(m_Z^2) (\times 10^4)$	1196	±	30	0.1	0.4-1.6	From R_ℓ^Z
$\sigma_{\text{had}}^0 (\times 10^3)$ (nb)	41541	±	37	0.1	4	Peak hadronic cross-section Luminosity measurement
$N_\nu (\times 10^3)$	2996	±	7	0.005	1	Z peak cross-sections Luminosity measurement
$R_b (\times 10^6)$	216290	±	660	0.3	< 60	Ratio of $b\bar{b}$ to hadrons Stat. extrapol. from SLD
$A_{\text{FB},0}^b (\times 10^4)$	992	±	16	0.02	1-3	b-quark asymmetry at Z pole From jet charge
$A_{\text{FB}}^{\text{pol},\tau} (\times 10^4)$	1498	±	49	0.15	< 2	τ polarization asymmetry τ decay physics
τ lifetime (fs)	290.3	±	0.5	0.001	0.04	Radial alignment
τ mass (MeV)	1776.86	±	0.12	0.004	0.04	Momentum scale
τ leptonic ($\mu\nu_\mu\nu_\tau$) B.R. (%)	17.38	±	0.04	0.0001	0.003	e/μ /hadron separation
m_W (MeV)	80350	±	15	0.25	0.3	From WW threshold scan Beam energy calibration
Γ_W (MeV)	2085	±	42	1.2	0.3	From WW threshold scan Beam energy calibration
$\alpha_s(m_W^2)(\times 10^4)$	1010	±	270	3	small	From R_ℓ^W
$N_\nu (\times 10^3)$	2920	±	50	0.8	small	Ratio of invis. to leptonic in radiative Z returns
m_{top} (MeV)	172740	±	500	17	small	From $t\bar{t}$ threshold scan QCD errors dominate
Γ_{top} (MeV)	1410	±	190	45	small	From $t\bar{t}$ threshold scan QCD errors dominate
$\lambda_{\text{top}}/\lambda_{\text{top}}^{\text{SM}}$	1.2	±	0.3	0.10	small	From $t\bar{t}$ threshold scan QCD errors dominate
ttZ couplings		±	30%	0.5 – 1.5 %	small	From $\sqrt{s} = 365$ GeV run

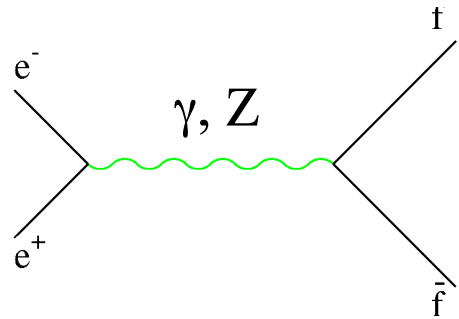
EW Precision Measurements at FCC-ee

Experimental (statistical and systematic) precision of a selection of measurements accessible at FCC-ee, compared with the present world-average precision. FCC-ee syst. scaled down from LEP estimates. Room for improvement with dedicated studies. Note that **syst.** go down also with **stat.** (e.g. beam energy determination from $ee \rightarrow Z/\gamma$ thus the associated uncertainty decreases with luminosity).

Table from mid-term report

Example of EW measurements @ Tera Z

measure $\sigma(e^+e^- \rightarrow \mu^+\mu^-)$ and $A_{FB}^{\mu\mu}$ at (a) judicious \sqrt{s}

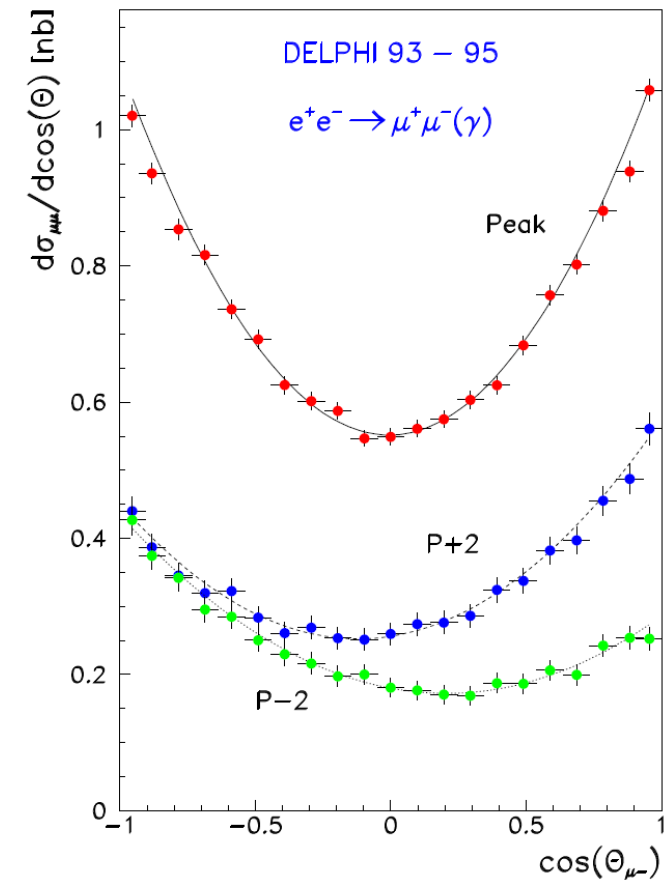


The γ exchange term is proportional to $\alpha_{QED}^2(\sqrt{s})$
 The Z exchange term is proportional to G_F^2 , hence independent of α_{QED}
 The γZ interference is proportional to $\alpha_{QED}(\sqrt{s}) \times G_F$

Excellent experimental control of off-peak di-muon asymmetry motivates campaign to collect 50-80 ab^{-1} off peak to gain highest sensitivity to Z- γ interference

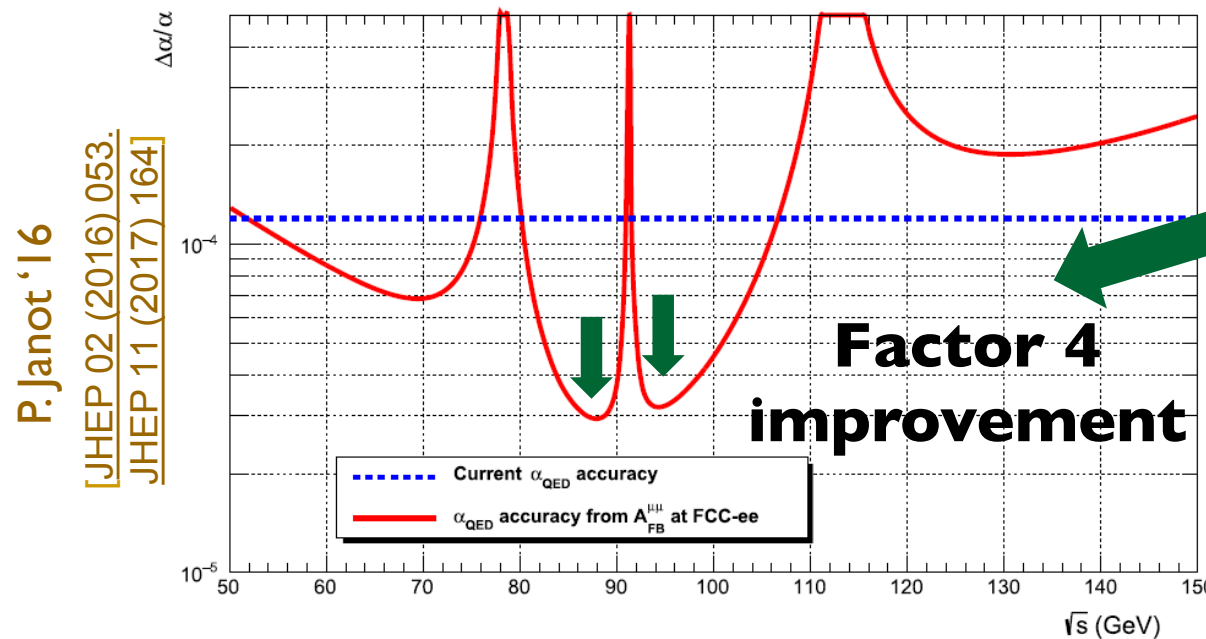
$$A_{FB}^{\mu\mu}(s) \simeq \frac{3}{4} \mathcal{A}_e \mathcal{A}_\mu \times \left[1 + \frac{8\pi\sqrt{2}\alpha_{QED}(s)}{m_Z^2 G_F (1 - 4\sin^2\theta_W^{eff})^2} \frac{s - m_Z^2}{2s} \right]$$

Allows for clean determination of $\alpha_{QED}(m_Z^2)$, which is a *critical* input for m_W closure tests (see later).



strongly depends on \sqrt{s}
direct measurement of $\alpha_{QED}(s)$ at $\sqrt{s} \neq m_Z$
 measure $\sin^2\theta_W$ to high precision (later)

relative α_{QED} uncertainty with 80 ab^{-1}



This dependence, & location of half-integer spin tunes, guides the choice of off-peak energies: 87.8 & 93.9 GeV.

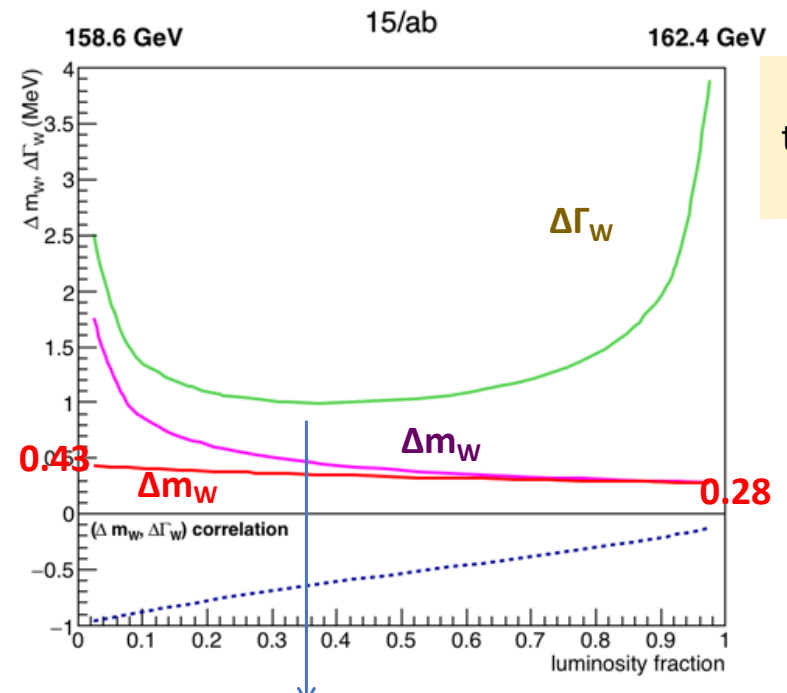
- Measure $\alpha_{QED}(m_Z^2)$ to 3×10^{-5} rel. precision (currently 1.1×10^{-4})
- Stat. dominated; syst. uncertainties $< 10^{-5}$ (dominated by \sqrt{s} calib)
- Theoretical uncertainties $\sim 10^{-4}$, higher order calcs needed

W Mass

Two independent W mass and width measurements @ FCCee :

1. The m_W and Γ_W determinations from the WW threshold cross section lineshape, with 12/ab at $E_{CM} \simeq 157.5-162.5$ GeV $\Delta m_W=0.4$ MeV $\Delta \Gamma_W=1$ MeV
2. Other measurements of m_W and Γ_W from the decay products kinematics at $E_{CM} \simeq 162.5-240-365$ GeV $\Delta m_W, \Delta \Gamma_W= 2-5$ MeV ?

Scans of possible $E_1 E_2$ data taking energies and luminosity fractions f (at the E_2 point)



A - minimum of $\Delta \Gamma_W=0.91$ MeV with $\Delta m_W=0.55$ MeV
 taking data at $E_1=156.6$ GeV $E_2=162.4$ GeV $f=0.25$
 yields $\Delta m_W=0.47$ MeV (as single par)

B- minimum of $\Delta m_W=0.28$ MeV $\Delta \Gamma_W=3.3$ MeV with
 $E_1=155.5$ GeV $E_2=162.4$ GeV $f=0.95$
 yields $\Delta m_W=0.28$ MeV (as single par)

C- minimum of $\Delta \Gamma_W=0.96$ MeV + $\Delta m_W=0.41$ MeV with
 $E_1=157.5$ GeV $E_2=162.4$ GeV $f=0.45$
 yields and $\Delta m_W=0.37$ MeV (as single par)

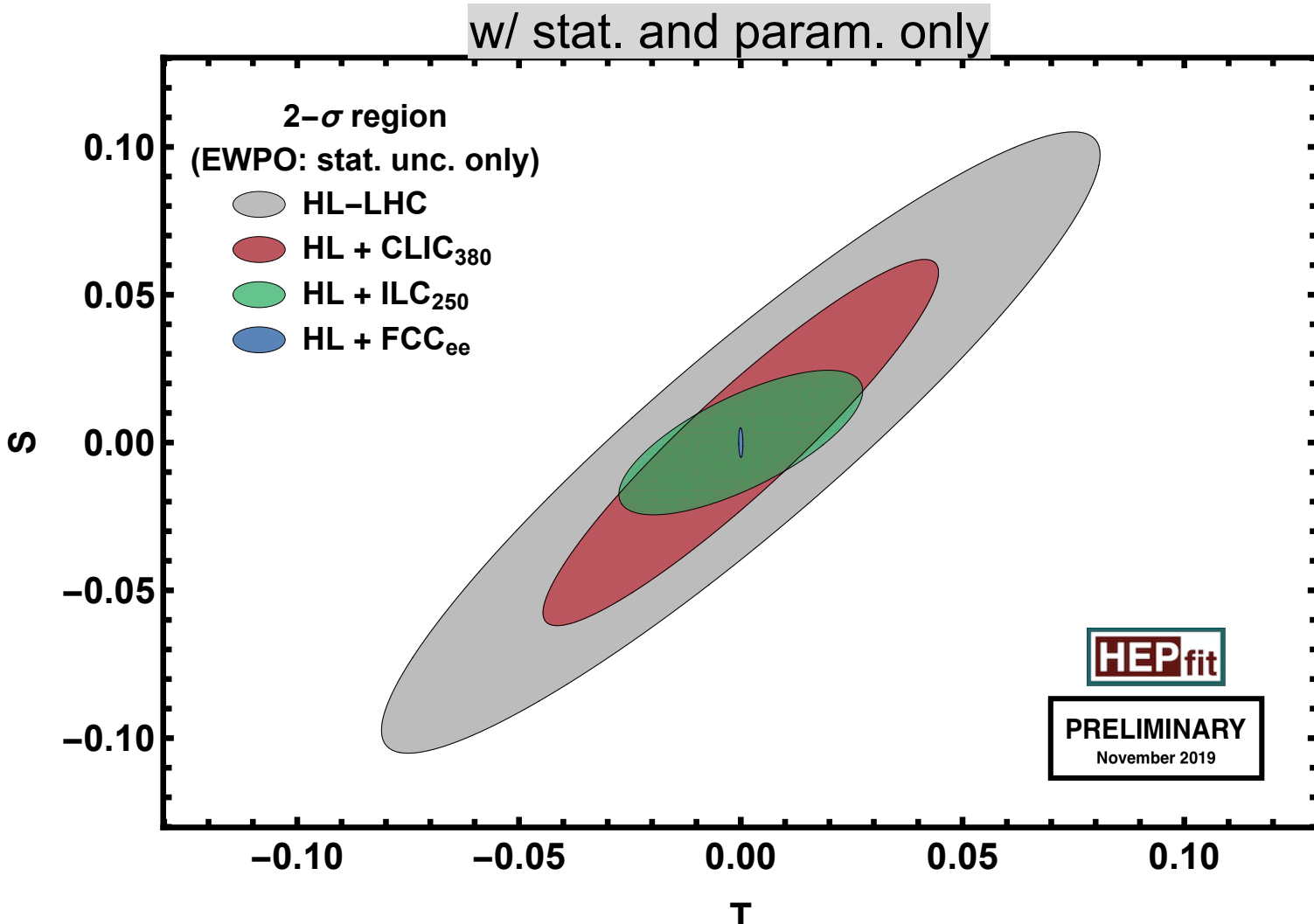
$\Delta m_W=0.45$ MeV, $\Delta \Gamma_W=1$ MeV ($r=-0.6$)
 $\Delta m_W=0.35$ MeV

$\Delta m_W, \Delta \Gamma_W$: error on W mass and width from fitting both
 Δm_W : error on W mass from fitting only m_W

Comparable
 in sensitivity
 with value
 from
 EWPO fit.

Tera-Z EW precision measurements.

- ▶ The target is to reduce syst. uncertainties to the level of stat. uncertainties.
(exploit the large samples and innovative control analyses)
 - ▶ Exquisite \sqrt{s} precision (100keV@Z, 300keV@WW) reduces beam uncertainties (EPOL)
- ➔ ~50 times better precision than LEP/LSD on EW precision observables
(stat. improvement alone is a factor **300-2'000** and innovative analyses/improved detectors can bring syst. down too)

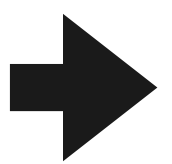


Indirect sensitivity
to 70TeV-scale sector
connected to EW/Higgs

(For the impact of the theory uncertainties on the EW fit, see bonus slides)

Tera-Z EW precision measurements.

- ▶ The target is to reduce syst. uncertainties to the level of stat. uncertainties.
(exploit the large samples and innovative control analyses)
- ▶ Exquisite \sqrt{s} precision (100keV@Z, 300keV@WW) reduces beam uncertainties (EPOL)



~50 times better precision than LEP/LSD on EW precision observables

(stat. improvement alone is a factor **300-2'000** and innovative analyses/improved detectors can bring syst. down too)

Need TH results to fully exploit Tera-Z

Table from mid-term report

Quantity	Current precision	FCC-ee stat. (syst.) precision	Required theory input	Available calc. in 2019	Needed theory improvement [†]
m_Z	2.1 MeV	0.004 (0.1) MeV	non-resonant $e^+e^- \rightarrow f\bar{f}$, initial-state radiation (ISR)	NLO, ISR logarithms up to 6th order	NNLO for $e^+e^- \rightarrow f\bar{f}$
Γ_Z	2.3 MeV	0.004 (0.025) MeV			
$\sin^2 \theta_{\text{eff}}^\ell$	1.6×10^{-4}	$2(2.4) \times 10^{-6}$			
m_W	12 MeV	0.25 (0.3) MeV	lineshape of $e^+e^- \rightarrow WW$ near threshold	NLO ($ee \rightarrow 4f$ or EFT framework)	NNLO for $ee \rightarrow WW$, $W \rightarrow f\bar{f}$ in EFT setup
HZZ coupling	—	0.2%	cross-sect. for $e^+e^- \rightarrow ZH$	NLO + NNLO QCD	NNLO electroweak
m_{top}	100 MeV	17 MeV	threshold scan $e^+e^- \rightarrow t\bar{t}$	N ³ LO QCD, NNLO EW, resummations up to NNLL	Matching fixed orders with resummations, merging with MC, α_s (input)

Indirect sensitivity to 70TeV-scale sector connected to EW/Higgs

[†]The listed needed theory calculations constitute a minimum baseline; additional partial higher-order contributions may also be required.

New Physics Reach @ Z-pole.

There are 48 different types of particles that can have tree-level linear interactions to SM.

de Blas, Criado, Perez-Victoria, Santiago, arXiv: 1711.10391

Name	\mathcal{S}	\mathcal{S}_1	\mathcal{S}_2	φ	Ξ	Ξ_1	Θ_1	Θ_3
Irrep	$(1, 1)_0$	$(1, 1)_1$	$(1, 1)_2$	$(1, 2)_{\frac{1}{2}}$	$(1, 3)_0$	$(1, 3)_1$	$(1, 4)_{\frac{1}{2}}$	$(1, 4)_{\frac{3}{2}}$

Name	ω_1	ω_2	ω_4	Π_1	Π_7	ζ
Irrep	$(3, 1)_{-\frac{1}{3}}$	$(3, 1)_{\frac{2}{3}}$	$(3, 1)_{-\frac{4}{3}}$	$(3, 2)_{\frac{1}{6}}$	$(3, 2)_{\frac{7}{6}}$	$(3, 3)_{-\frac{1}{3}}$

Name	Ω_1	Ω_2	Ω_4	Υ	Φ
Irrep	$(6, 1)_{\frac{1}{3}}$	$(6, 1)_{-\frac{2}{3}}$	$(6, 1)_{\frac{4}{3}}$	$(6, 3)_{\frac{1}{3}}$	$(8, 2)_{\frac{1}{2}}$

Scalars

Name	N	E	Δ_1	Δ_3	Σ	Σ_1
Irrep	$(1, 1)_0$	$(1, 1)_{-1}$	$(1, 2)_{-\frac{1}{2}}$	$(1, 2)_{-\frac{3}{2}}$	$(1, 3)_0$	$(1, 3)_{-1}$

Name	U	D	Q_1	Q_5	Q_7	T_1	T_2
Irrep	$(3, 1)_{\frac{2}{3}}$	$(3, 1)_{-\frac{1}{3}}$	$(3, 2)_{\frac{1}{6}}$	$(3, 2)_{-\frac{5}{6}}$	$(3, 2)_{\frac{7}{6}}$	$(3, 3)_{-\frac{1}{3}}$	$(3, 3)_{\frac{2}{3}}$

Fermions

Name	\mathcal{B}	\mathcal{B}_1	\mathcal{W}	\mathcal{W}_1	\mathcal{G}	\mathcal{G}_1	\mathcal{H}	\mathcal{L}_1
Irrep	$(1, 1)_0$	$(1, 1)_1$	$(1, 3)_0$	$(1, 3)_1$	$(8, 1)_0$	$(8, 1)_1$	$(8, 3)_0$	$(1, 2)_{\frac{1}{2}}$

Name	\mathcal{L}_3	\mathcal{U}_2	\mathcal{U}_5	\mathcal{Q}_1	\mathcal{Q}_5	\mathcal{X}	\mathcal{Y}_1	\mathcal{Y}_5
Irrep	$(1, 2)_{-\frac{3}{2}}$	$(3, 1)_{\frac{2}{3}}$	$(3, 1)_{\frac{5}{3}}$	$(3, 2)_{\frac{1}{6}}$	$(3, 2)_{-\frac{5}{6}}$	$(3, 3)_{\frac{2}{3}}$	$(\bar{6}, 2)_{\frac{1}{6}}$	$(\bar{6}, 2)_{-\frac{5}{6}}$

Vectors

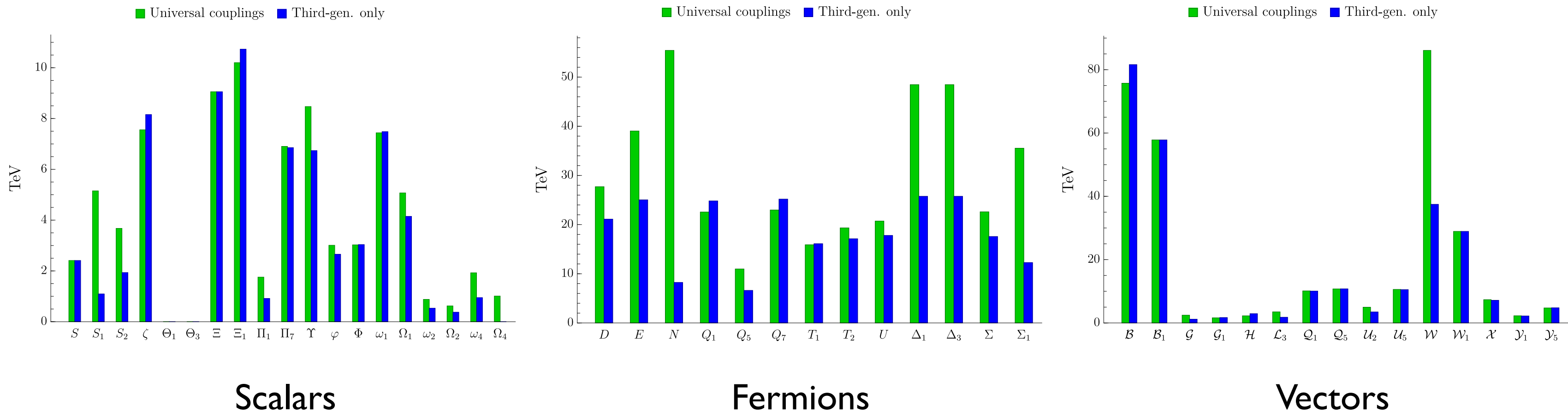
They are not all affecting EW observables at tree-level.

New Physics Reach @ Z-pole.

There are 48 different types of particles that can have tree-level linear interactions to SM.

They are not all affecting EW observables at tree-level.
However, all, but a few, have leading log. running into EW observables.

Allwicher, McCullough, Renner, arXiv: 2408.03992



Tree-level matching and running from 1 TeV to Z mass.
W- and Z-pole observables only (no Higgs, no LEP-2 like observables)

New Physics Reach @ Z-pole.

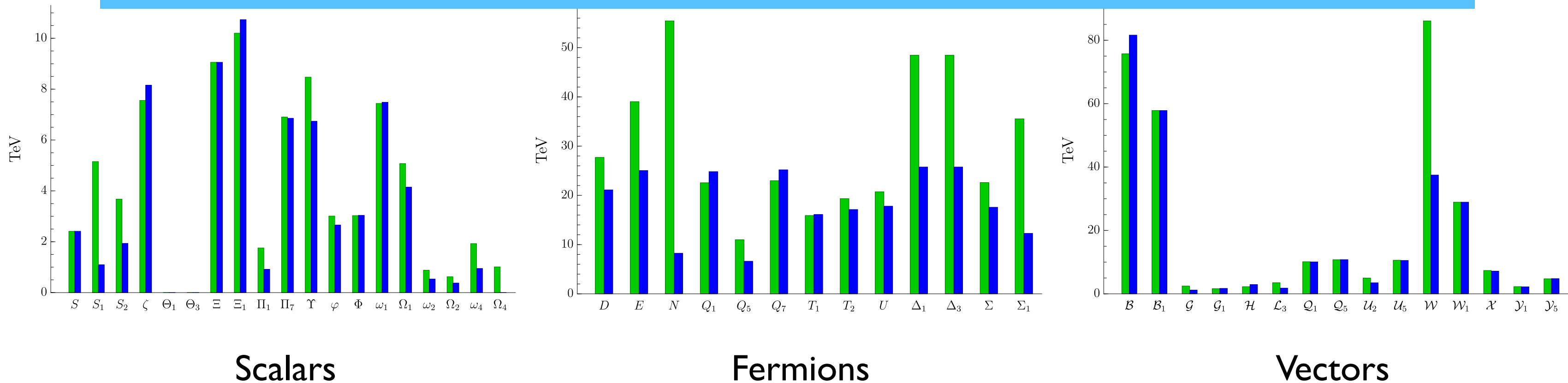
TI

SM.

Tera-Z programme gives comprehensive coverage of new physics coupled to SM.

If a signature shows up elsewhere, it will also show up at Tera-Z.

Tera-Z is not just a high-power LEP exploring the EW sector.



Tree-level matching and running from 1 TeV to Z mass.
W- and Z-pole observables only (no Higgs, no LEP-2 like observables)

Flavour Physics from Z-pole Run

Flavour potential.

At present (Z/h/NewPhysics) FCNCs mostly constrained by low energy observables.
 The large statistics of FCC will open on-shell opportunities.

Particle production (10^9)	B^0 / \bar{B}^0	B^+ / B^-	B_s^0 / \bar{B}_s^0	$\Lambda_b / \bar{\Lambda}_b$	$c\bar{c}$	τ^- / τ^+
Belle II	27.5	27.5	n/a	n/a	65	45
FCC- ee	300	300	80	80	600	150

FCC- ee
 =
 10 x Belle II

Decay mode/Experiment	Belle II (50/ab)	LHCb Run I	LHCb Upgr. (50/fb)	FCC- ee
EW/H penguins				
$B^0 \rightarrow K^*(892)e^+e^-$	~ 2000	~ 150	~ 5000	~ 200000
$\mathcal{B}(B^0 \rightarrow K^*(892)\tau^+\tau^-)$	~ 10	–	–	~ 1000
$B_s \rightarrow \mu^+\mu^-$	n/a	~ 15	~ 500	~ 800
$B^0 \rightarrow \mu^+\mu^-$	~ 5	–	~ 50	~ 100
$\mathcal{B}(B_s \rightarrow \tau^+\tau^-)$				
Leptonic decays				
$B^+ \rightarrow \mu^+\nu_{mu}$	5%	–	–	3%
$B^+ \rightarrow \tau^+\nu_{tau}$	7%	–	–	2%
$B_c^+ \rightarrow \tau^+\nu_{tau}$	n/a	–	–	5%
CP / hadronic decays				
$B^0 \rightarrow J/\Psi K_S (\sigma_{\sin(2\phi_d)})$	$\sim 2 \cdot 10^6 (0.008)$	41500 (0.04)	$\sim 0.8 \cdot 10^6 (0.01)$	$\sim 35 \cdot 10^6 (0.006)$
$B_s \rightarrow D_s^\pm K^\mp$	n/a	6000	~ 200000	$\sim 30 \cdot 10^6$
$B_s(B^0) \rightarrow J/\Psi\phi (\sigma_{\phi_s} \text{ rad})$	n/a	96000 (0.049)	$\sim 2 \cdot 10^6 (0.008)$	$16 \cdot 10^6 (0.003)$

See S. Monteil, Flavour@FCC'22

out of reach
 at LHCb/Belle

boosted b's/ τ 's
 at FCC- ee
 Makes possible
 a topological rec.
 of the decays
 w/ miss. energy

Flavour potential.

At present (Z/h/NewPhysics) FCNCs mostly constrained by low energy observables.
 The large statistics of FCC will open on-shell opportunities.

Particle production (10^9)	B^0 / \bar{B}^0	B^+ / B^-	B_s^0 / \bar{B}_s^0	$\Lambda_b / \bar{\Lambda}_b$	$c\bar{c}$	τ^- / τ^+
Belle II	27.5	27.5	n/a	n/a	65	45
FCC-ee	300	300	80	80	600	150

FCC-ee
 =
 10 x Belle II

Decay mode/Experiment	Belle II (50/ab)	LHCb Run I	LHCb Upgr. (50/fb)	FCC-ee
EW/H penguins				
$B^0 \rightarrow K^*(892)e^+e^-$	~ 2000	~ 150	~ 5000	~ 200000
$\mathcal{B}(B^0 \rightarrow K^*(892)\tau^+\tau^-)$	~ 10			~ 1000

See S. Monteil, Flavour@FCC'22

out of reach at LHCb/Belle

Flavour @ FCC vs Belle/pp

Attribute	$\Upsilon(4S)$	pp	Z^0
All hadron species		✓	✓
High boost		✓	✓
Enormous production cross-section		✓	
Negligible trigger losses	✓		✓
Low backgrounds	✓		✓
Initial energy constraint	✓		(✓)

boosted b's/ τ 's at FCC-ee

Makes possible a topological rec. of the decays w/ miss. energy

Flavour potential.

At present (Z/h/NewPhysics) FCNCs mostly constrained by low energy observables.
 The large statistics of FCC will open on-shell opportunities.

Particle production (10^9)	B^0 / \bar{B}^0	B^+ / B^-	B_s^0 / \bar{B}_s^0	$\Lambda_b / \bar{\Lambda}_b$	$c\bar{c}$	τ^- / τ^+
Belle II	27.5	27.5	n/a	n/a	65	45
FCC-ee	300	300	80	80	600	150

FCC-ee
 =
 10 x Belle II

Decay mode/Experiment	Belle II (50/ab)	LHCb Run I	LHCb Upgr. (50/fb)	FCC-ee
EW/H penguins				
$B^0 \rightarrow K^*(892)e^+e^-$	~ 2000	~ 150	~ 5000	~ 200000
$\mathcal{B}(B^0 \rightarrow K^*(892)\tau^+\tau^-)$	~ 10			~ 1000

See S. Monteil, Flavour@FCC'22

out of reach at LHCb/Belle

Attribute	$\Upsilon(4S)$	pp	Z^0
All hadron species		✓	✓
High boost		✓	✓
Enormous production cross-section		✓	
Negligible trigger losses	✓		✓
Low backgrounds	✓		✓

boosted b's/ τ 's at FCC-ee
 Makes possible a topological rec. of the decays w/ miss. energy

Flavour defines shared (vertexing, tracking, calorimetry) and specific (hadronic PID) detector requirements.

FCC-ee flavour opportunities.

- **CKM element V_{cb}** (critical for normalising the Unitarity Triangle) from WW decays
- **Tau physics** ($>10^{11}$ pairs of tau's produced in Z decays)
 - test of lepton flavour universality: G_F from tau decays @ 10 ppm @ FCC-ee (0.5 ppm from muon decays)
 - lepton flavour violation:
 - $\tau \rightarrow \mu \gamma$: 4×10^{-8} @ Belle2021 $\rightarrow 10^{-9}$ @ FCC-ee
 - $\tau \rightarrow 3\mu$: 2×10^{-8} @ Belle $\rightarrow 3 \times 10^{-10}$ @ BelleII $\rightarrow 10^{-11}$ @ FCC-ee
 - tau lifetime uncertainty:
 - 2000 ppm \rightarrow 10 ppm
 - tau mass uncertainty:
 - 70 ppm \rightarrow 14 ppm
- **Semi-leptonic mixing asymmetries a_{sl}^s and a_{sl}^d**
- ...

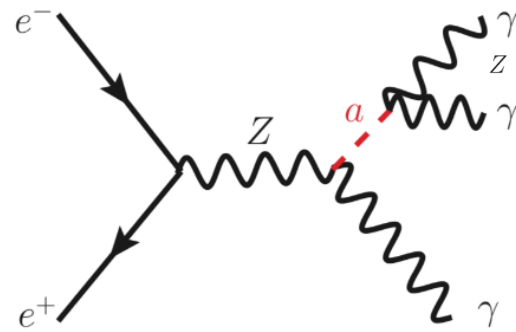
BSM Exploration

FCC-ee: Explore & Discover.

- **ALPs@ colliders**

e.g. $e^+e^- \rightarrow \gamma a$

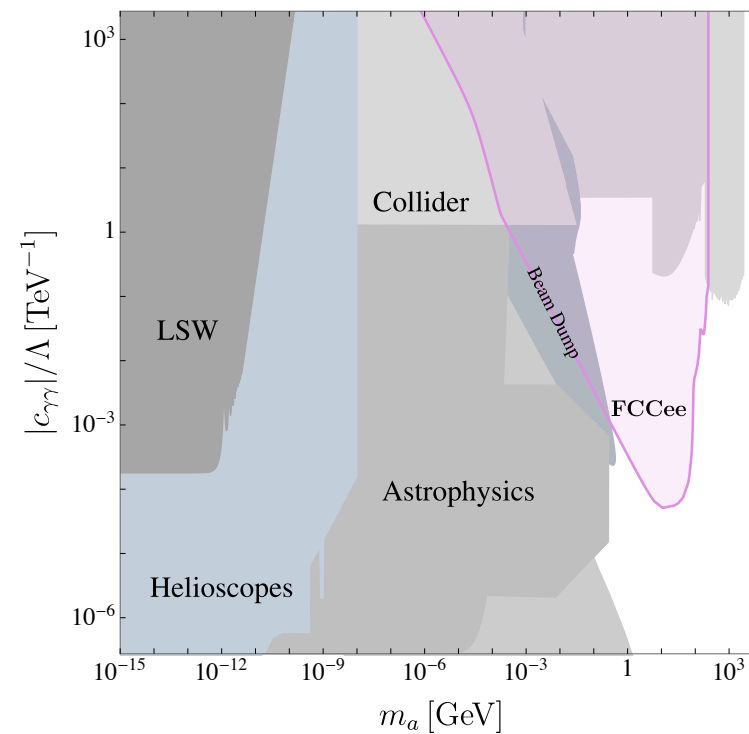
$e^+e^- \rightarrow ha$



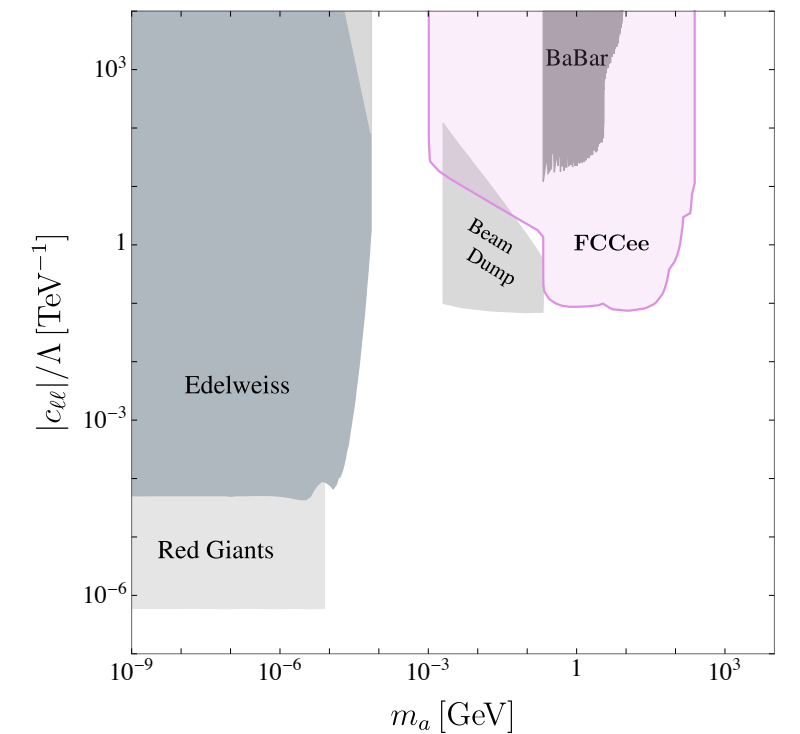
Knapen, Thamm arXiv:2108.08949

Astro/Cosmo → long-lived ALPs
colliders → short-lived ALPs MeV+

ALP coupling to photons

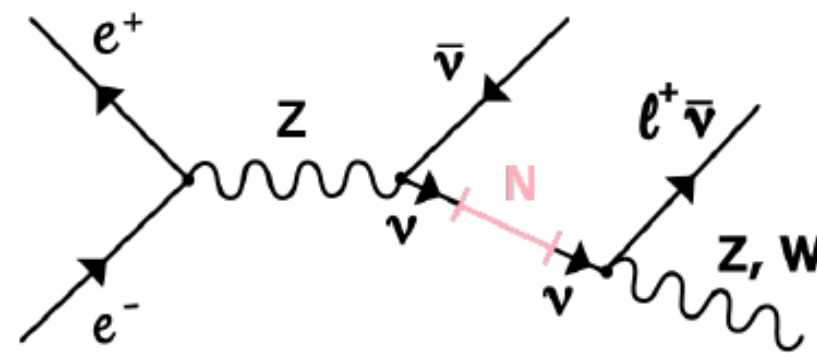


ALP coupling to electrons



- **Search for ν_{RH} .**

Direct observation
in Z decays
from LH-RH mixing



mixing active-sterile neutrinos

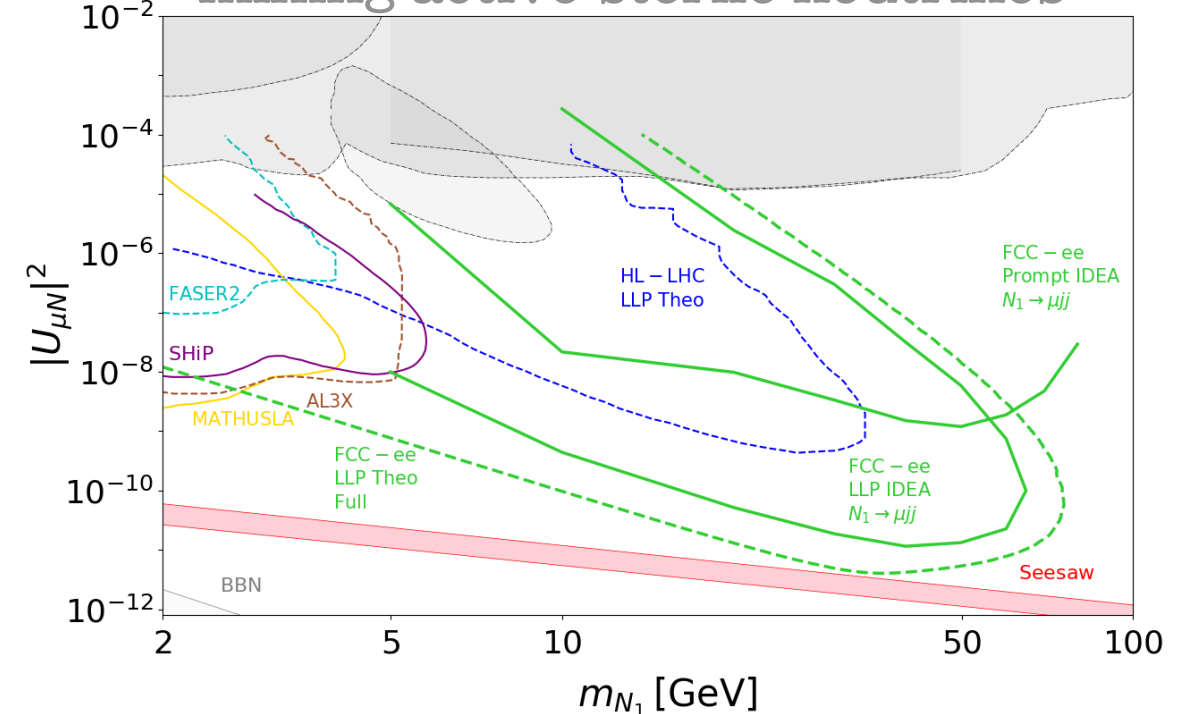


Fig. from mid-term report

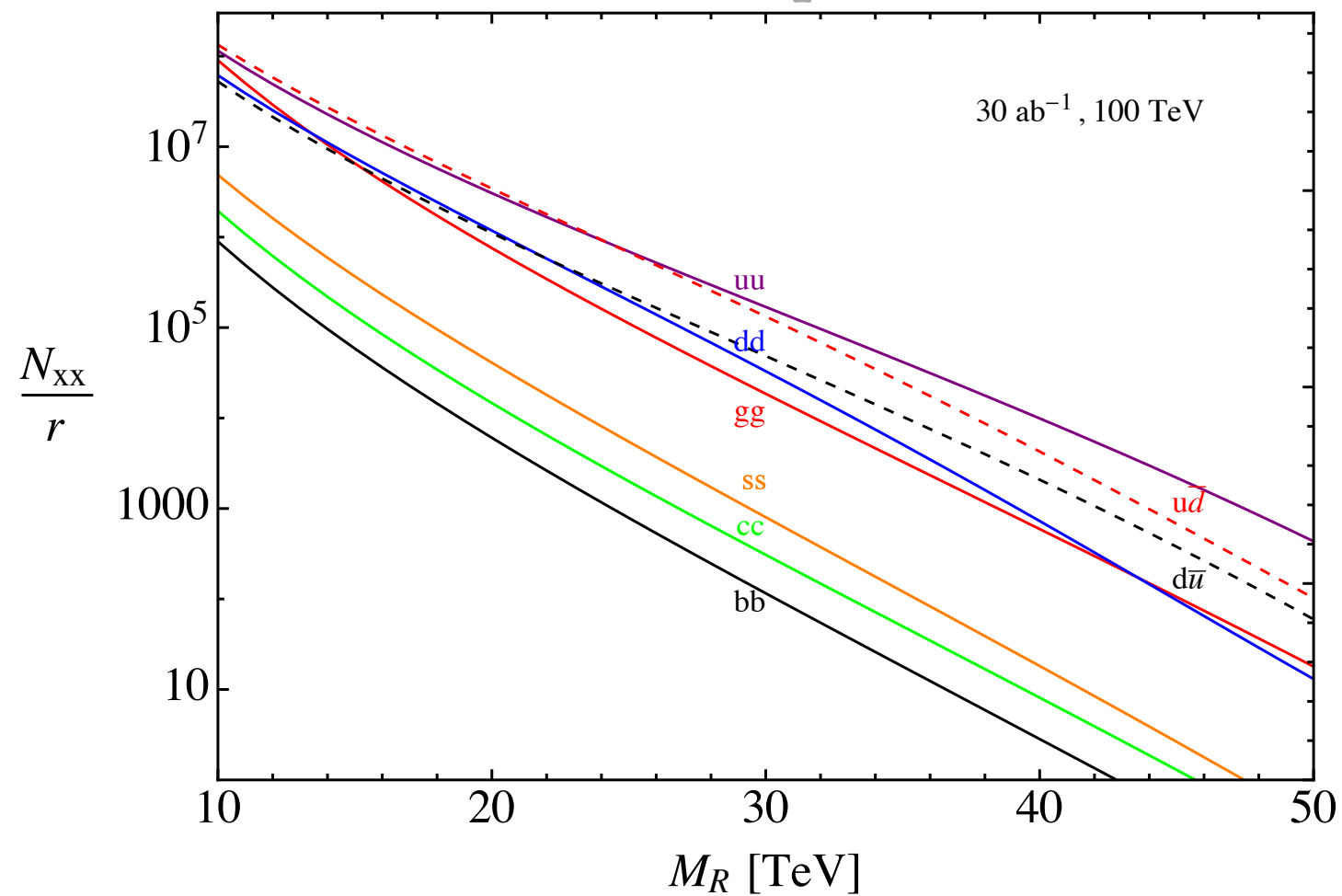
Oct. 4, 2024

Resonance production.

Protons are made of 5 quarks, gluons, photons, W/Z

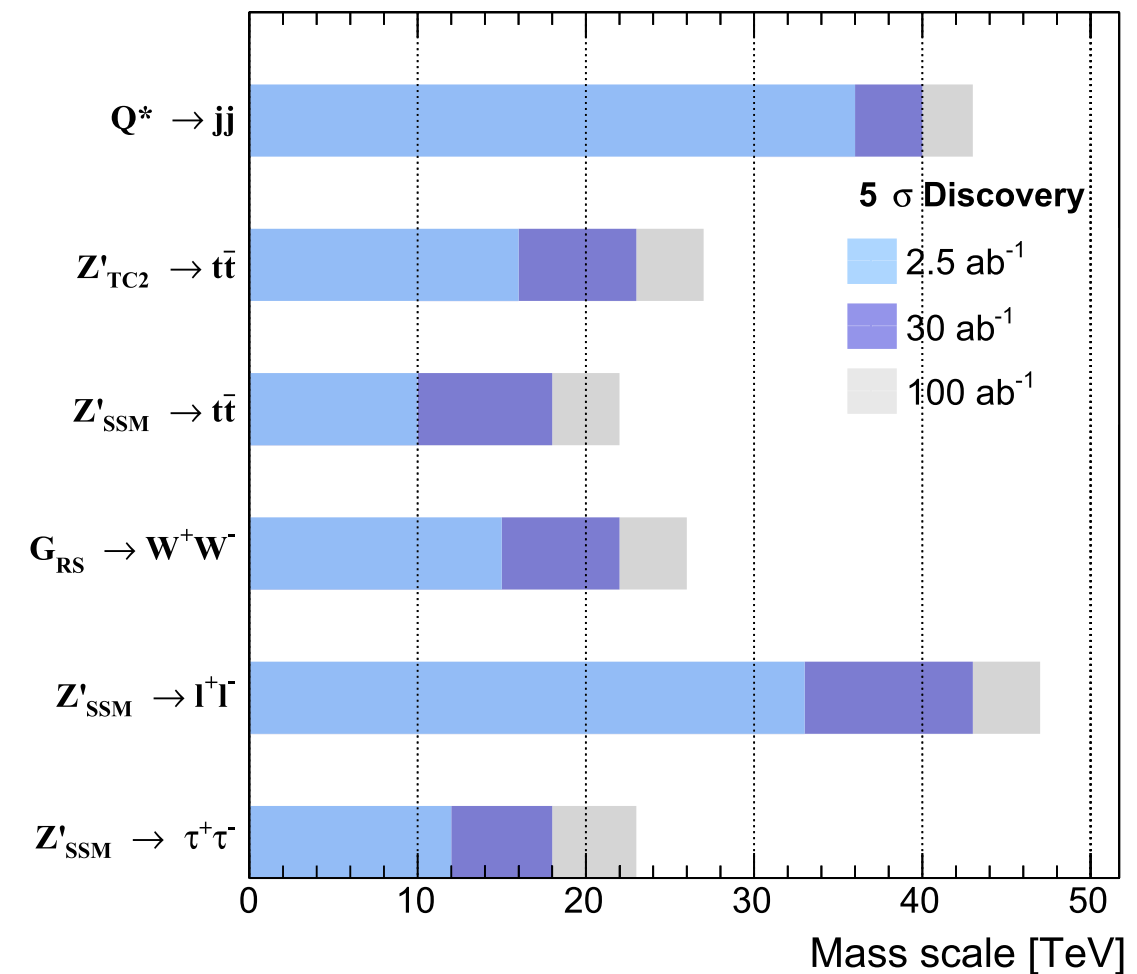
FCC-hh effectively collides 196 different initial states = perfect exploratory machine

resonances produced



Plot from mid-term report

FCC-hh mass reach



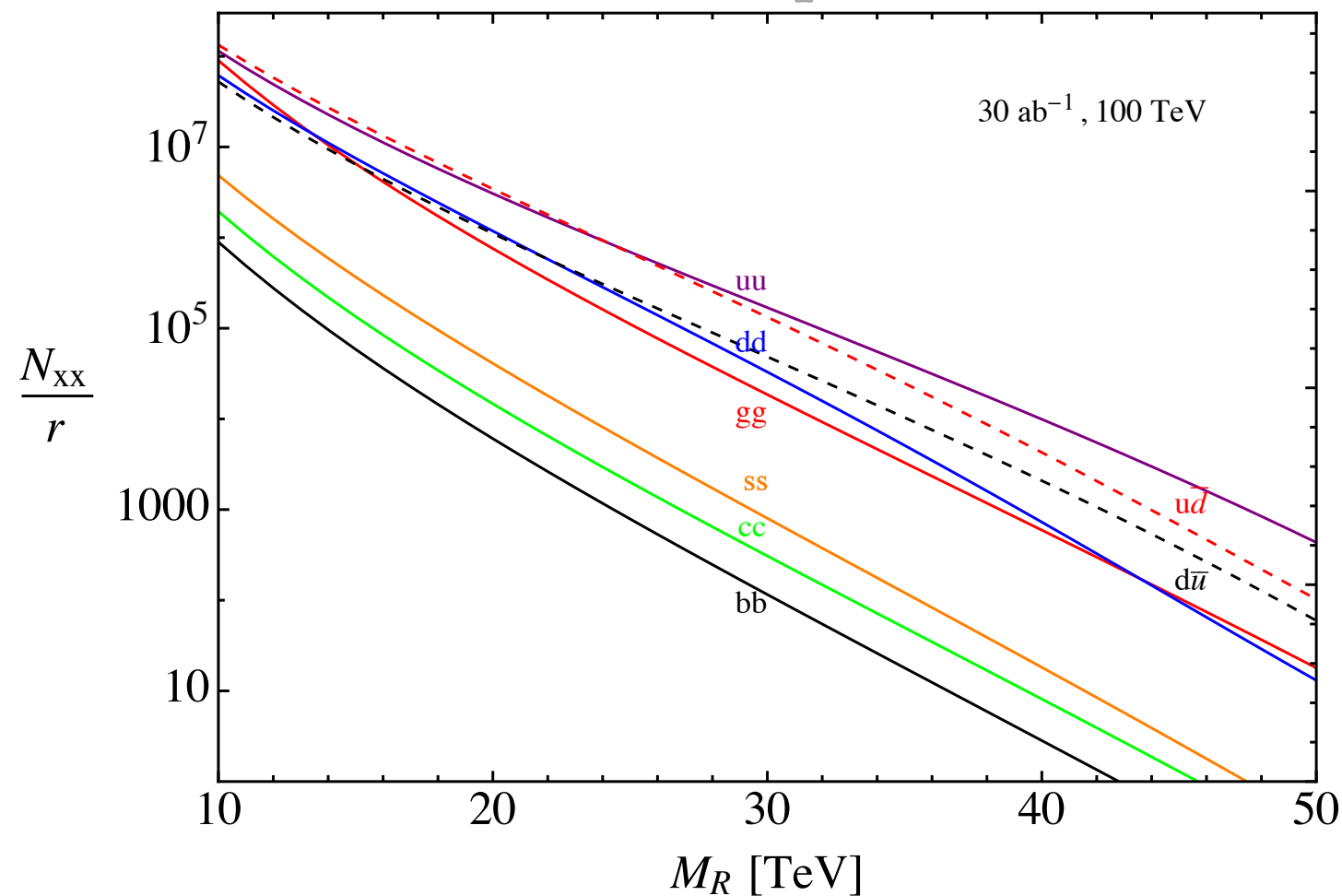
Plot from FCC CDR

Resonance production.

Protons are made of 5 quarks, gluons, photons, W/Z

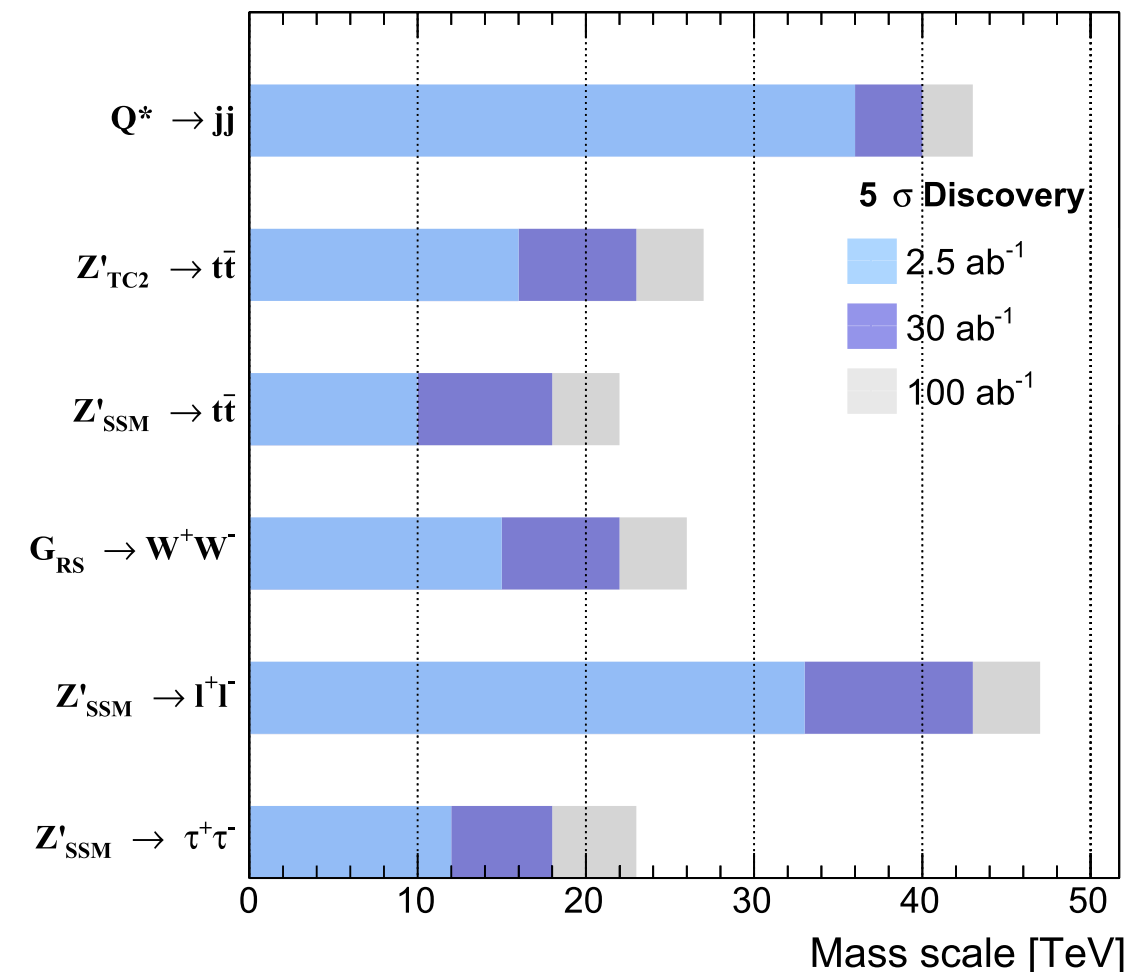
FCC-hh effectively collides 196 different initial states = perfect exploratory machine

resonances produced



Plot from mid-term report

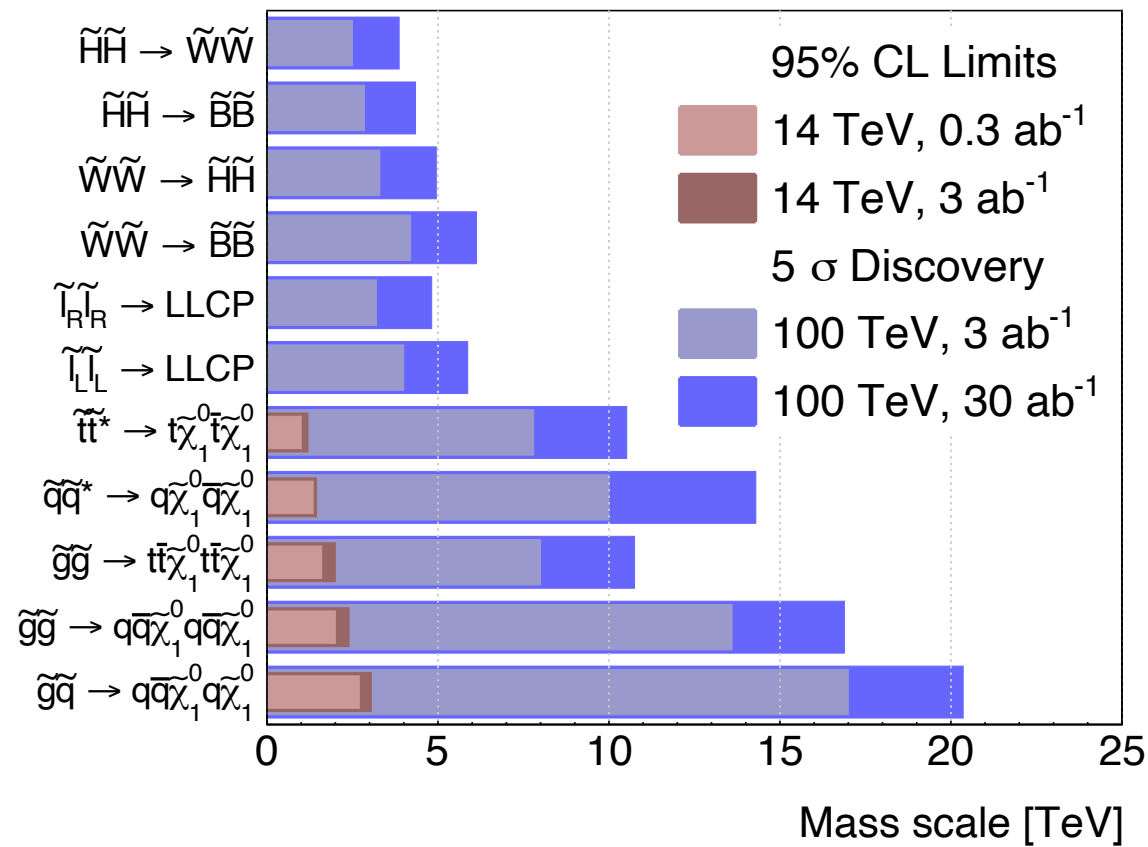
FCC-hh mass reach



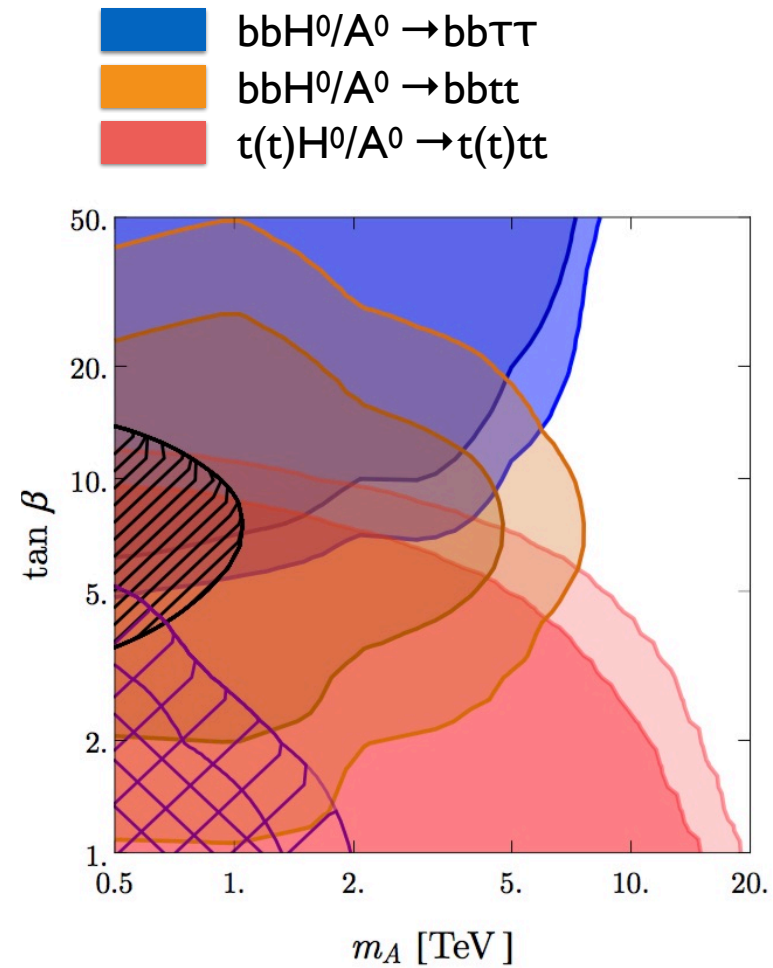
Plot from FCC CDR

FCC-hh allows the direct exploration of new physics at energy scales up to 40 TeV, including any physics that may be indirectly indicated by precision Higgs and EW measurements at FCC-ee.

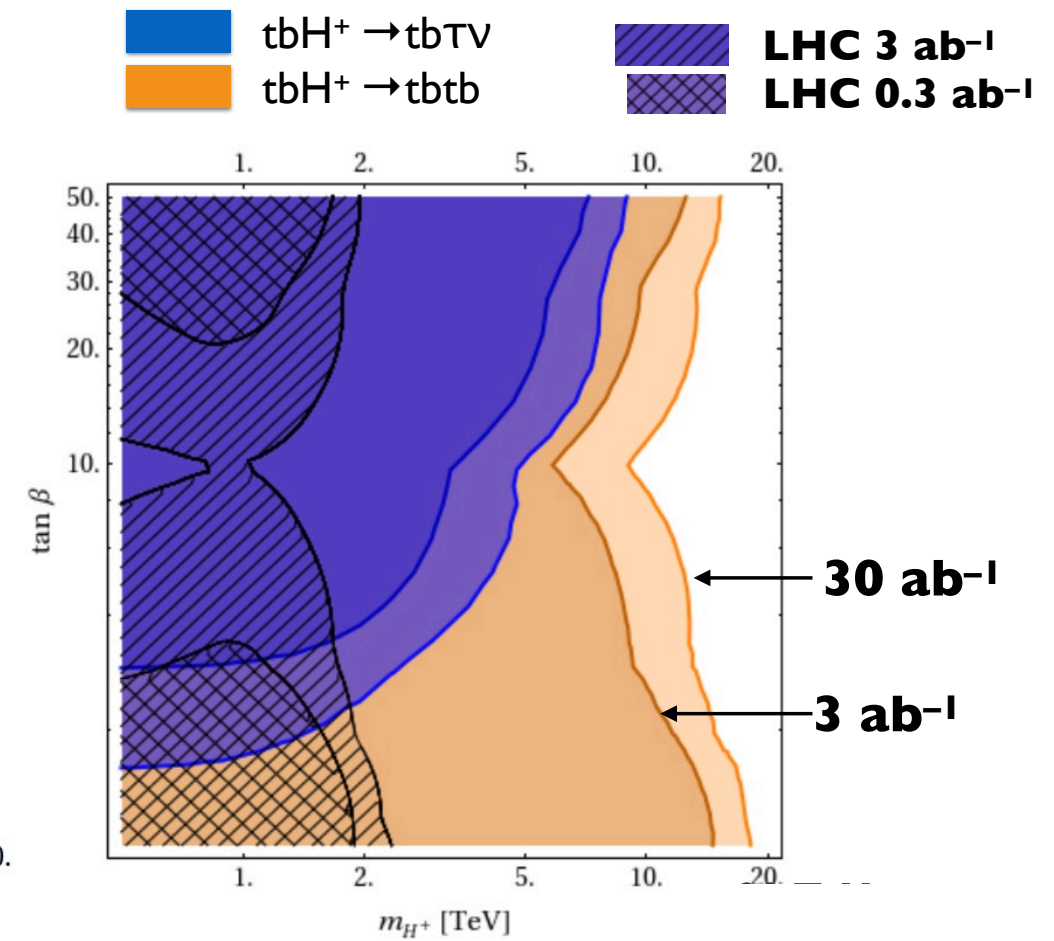
Pushing limits of SUSY.



Plot from [arXiv:1606.00947](https://arxiv.org/abs/1606.00947)



Plot from [arXiv:1605.08744](https://arxiv.org/abs/1605.08744) and [arXiv:1504.07617](https://arxiv.org/abs/1504.07617)



Factor 10 increase on the HL-LHC limits.

15-20TeV squarks/gluinos
 require kinematic threshold 30-40TeV:
 FCC-hh is more than a $\sqrt{s} \sim 10\text{TeV}$ factory

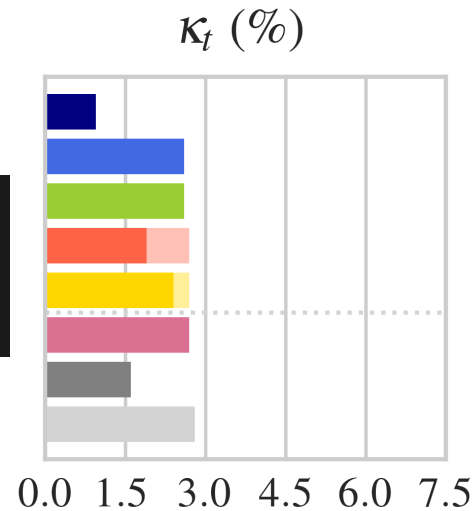
FCC-ee/FCC-hh Interplay

Synergy $ee \leftrightarrow hh$.

1 FCC-hh without ee could bound BR_{inv} but it could say nothing about $BR_{untagged}$ (FCC- ee needed for absolute normalisation of Higgs couplings)

FCC-hh is determining top Yukawa through ratio $t\bar{t}h/t\bar{t}Z$

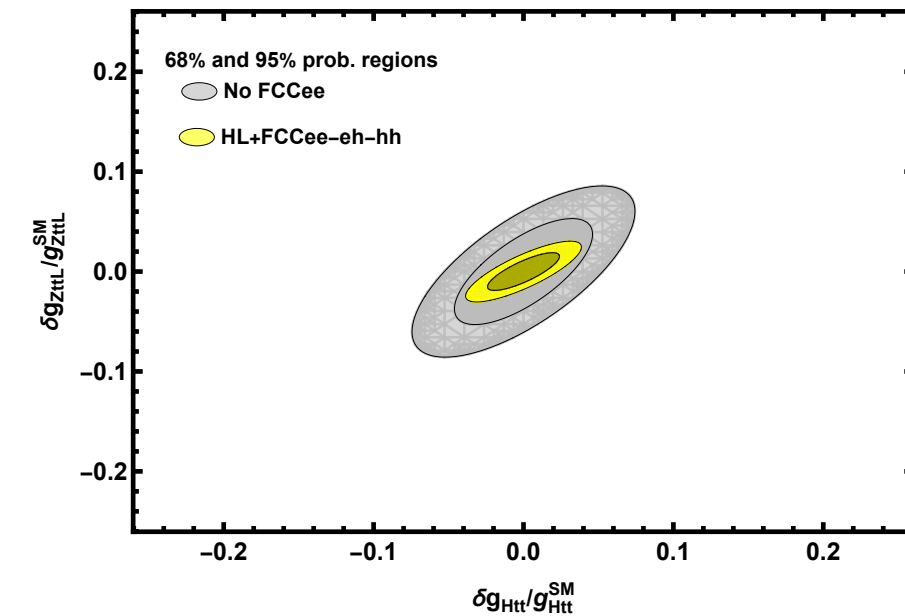
So the extraction of top Yukawa heavily relies on the knowledge of $t\bar{t}Z$ from FCC- ee



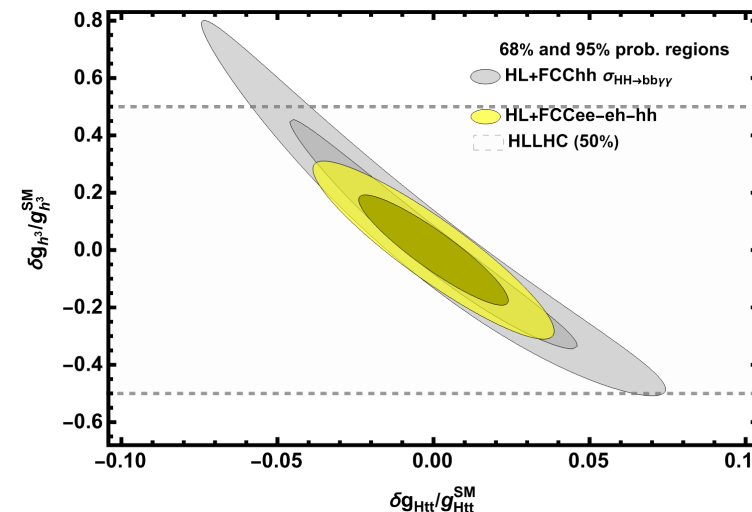
Mangano+ '15

	$\sigma(t\bar{t}H)$ [pb]	$\sigma(t\bar{t}Z)$ [pb]	$\frac{\sigma(t\bar{t}H)}{\sigma(t\bar{t}Z)}$
13 TeV	$0.475^{+5.79\%+3.33\%}_{-9.04\%-3.08\%}$	$0.785^{+9.81\%+3.27\%}_{-11.2\%-3.12\%}$	$0.606^{+2.45\%+0.525\%}_{-3.66\%-0.319\%}$
100 TeV	$33.9^{+7.06\%+2.17\%}_{-8.29\%-2.18\%}$	$57.9^{+8.93\%+2.24\%}_{-9.46\%-2.43\%}$	$0.585^{+1.29\%+0.314\%}_{-2.02\%-0.147\%}$

(uncertainty drops in ratio)



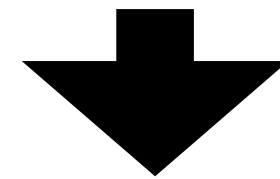
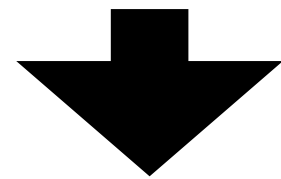
3 Subsequently, the 1% sensitivity on $t\bar{t}h$ is essential to determine h^3 at $O(5\%)$ at FCC-hh



Plots from mid-term report

Conclusions & Outlook

- LHC changed the HEP landscape (Higgs and nothing else - yet?)
- Experimental analyses have been smarter than initially thought:
1 ab⁻¹ in 2024 is more valuable than 1 ab⁻¹ in 2008!



A **circular “Higgs factory”** like FCC-ee has a rich potential:

- ◎ Quantum leap in testing the Standard Model
- ◎ Search directly *and* indirectly for New Physics

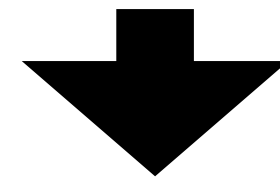
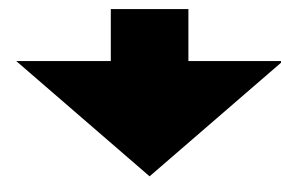
And FCC-ee is an essential part of an integrated programme to probe the energy frontier.

The FCC project perfectly fits the **needs of HEP after LHC**:

▶ guaranteed deliverables & broad exploration potential ◀

Conclusions & Outlook

- LHC changed the HEP landscape (Higgs and nothing else - yet?)
- Experimental analyses have been smarter than initially thought:
1 ab⁻¹ in 2024 is more valuable than 1 ab⁻¹ in 2008!



A **circular “Higgs factory”** like FCC-ee has a rich potential:

- ◎ Quantum leap in testing the Standard Model
- ◎ Search directly *and* indirectly for New Physics

And FCC-ee is an essential part of an integrated programme to probe the energy frontier.

The FCC project perfectly fits the **needs of HEP after LHC**:

▶ guaranteed deliverables & broad exploration potential ◀

**We have profound questions and we need to create opportunities to answer them.
FCC will for sure contribute.**

Acknowledgement.

This project is supported from the European Union's Horizon 2020 research and innovation programme under grant agreement No 951754.



BONUS

Some work ahead of us.

- Development of a common software and the estimate of the computing needs
- Evaluation of the physics performance and requirements for detectors
- Conceptualisation of detectors capable of delivering these requirements
- Mitigation of the interaction region constraints on detectors and vice versa
- Design of methods and tools for centre-of-mass energy calibration, beam polarisation, and monochromatization
- Understanding and optimisation of the physics program
- Exploration of the physics opportunities
- Development of the theoretical tools and observables needed to meet the measurement targets

FCC feasibility study

The launch of the feasibility study.



“An **electron-positron** Higgs factory is the highest-priority next collider. For the longer term, the European particle physics community has the ambition to operate a **proton-proton** collider at the highest achievable energy.”

— CERN council approved the Strategy and CERN management implemented it —
FCC Feasibility Study (FS) started in 2021 and will be completed in 2025.

Mid-term review in 2023.

Objectives of FCC feasibility study.

- Demonstration of the **geological, technical, environmental and administrative feasibility** of the tunnel and surface areas and optimisation of placement and layout of the ring and related infrastructure.
- Pursuit, together with the Host States, of the preparatory **administrative processes** required for a potential project approval to identify and remove any showstopper.
- Optimisation of the design of the **colliders and their injector chains**, supported by R&D to develop the needed key technologies.
- Elaboration of a **sustainable operational model** for the colliders and experiments in terms of human and financial resource needs, as well as environmental aspects and energy efficiency.
- Development of a **consolidated cost estimate**, as well as the **funding and organisational models** needed to enable the project's technical design completion, implementation and operation.
- **Identification of substantial resources** from outside CERN's budget for the implementation of the first stage of a possible future project (**tunnel and FCC-ee**).
- Consolidation of the **physics case and detector concepts** for both colliders.

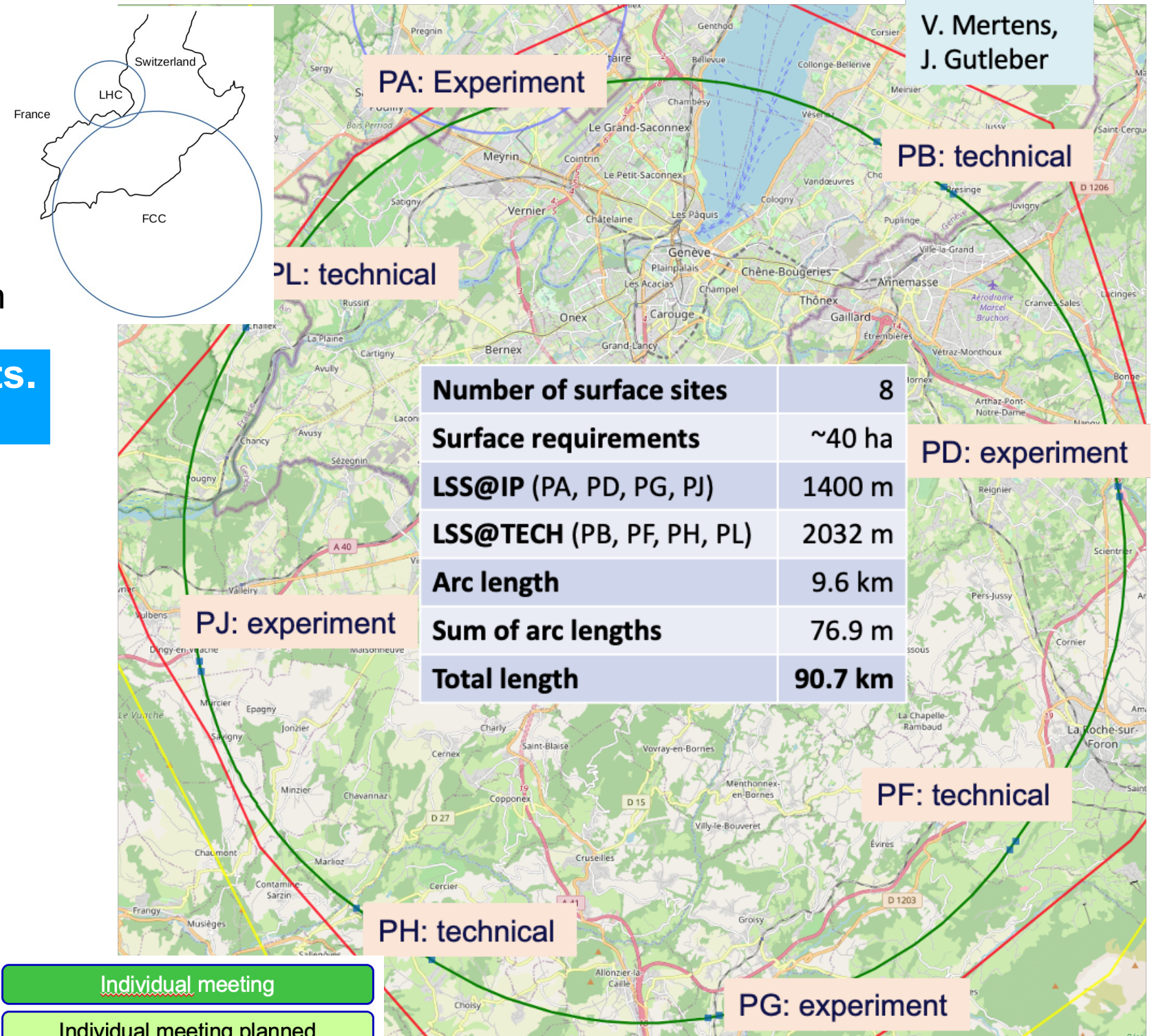
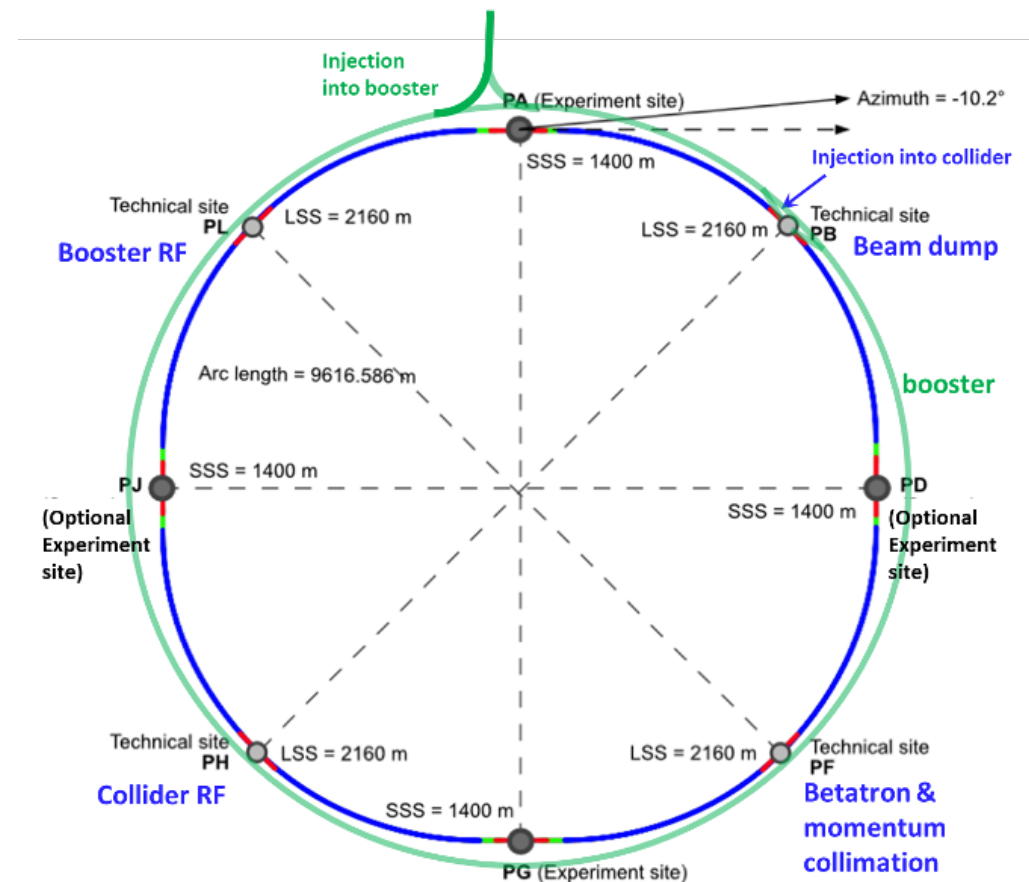
Optimized placement and layout.

M. Benedikt @ CERN 13.02.24

Layout chosen out of ~ 100 initial variants, based on **geology** and **surface constraints** (land availability, access to roads, etc.), **environment**, (protected zones), **infrastructure** (water, electricity, transport), **machine performance** etc.

“Avoid-reduce-compensate” principle of EU and French regulation

Overall lowest-risk baseline: 90.7 km ring, 8 surface points.
Whole project now adapted to this placement



Number of surface sites	8
Surface requirements	~40 ha
LSS@IP (PA, PD, PG, PJ)	1400 m
LSS@TECH (PB, PF, PH, PL)	2032 m
Arc length	9.6 km
Sum of arc lengths	76.9 m
Total length	90.7 km

- Individual meeting
- Individual meeting planned
- Collective meeting

V. Mertens, J. Gutleber

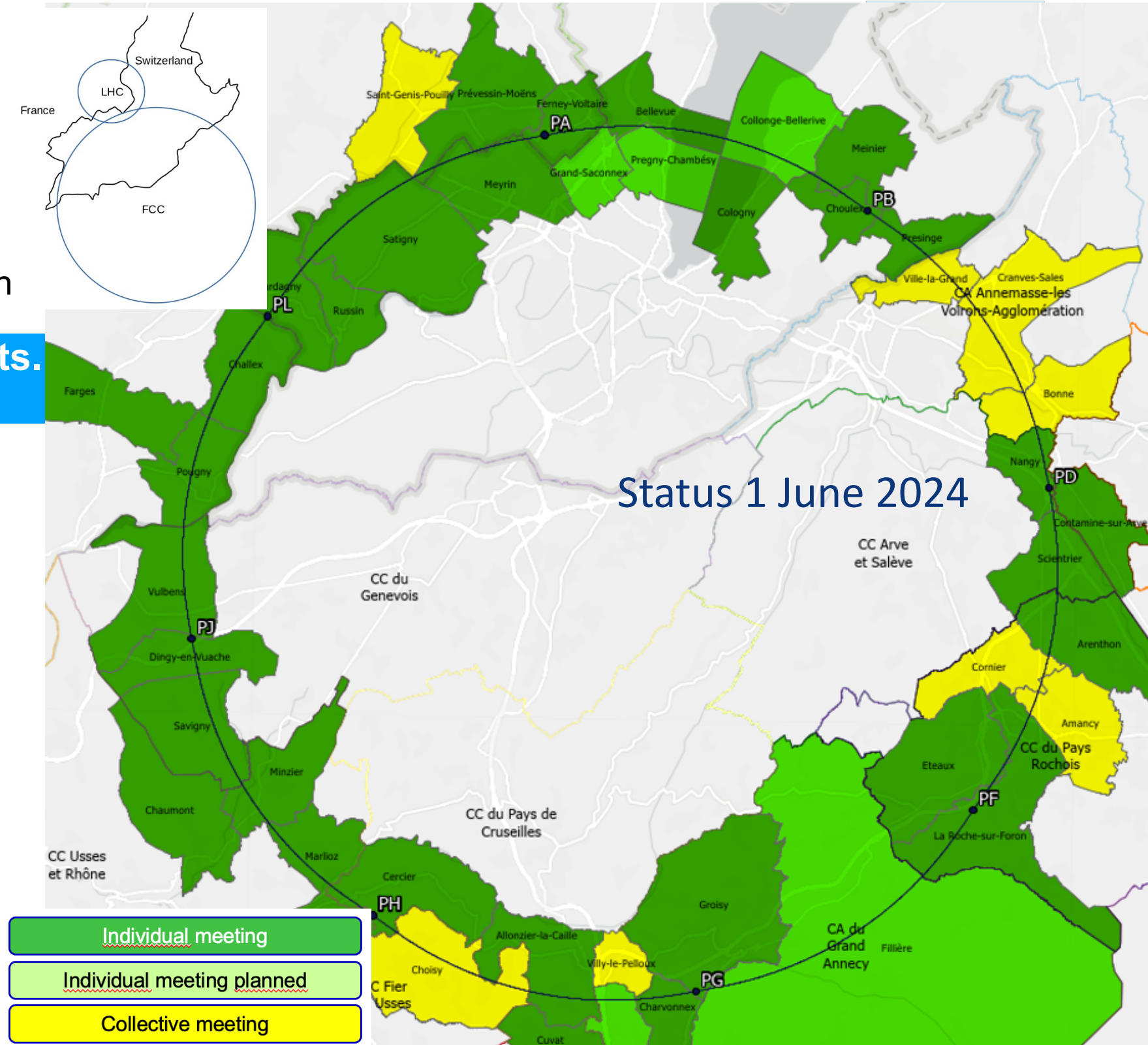
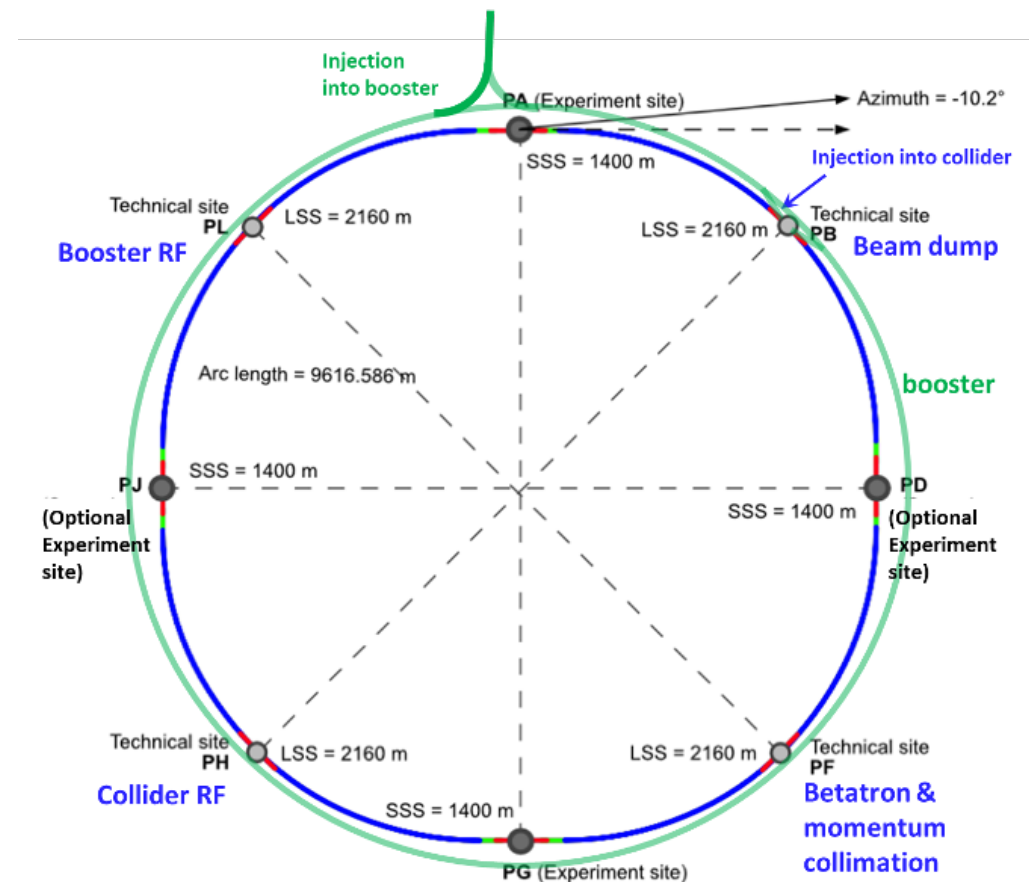
Optimized placement and layout.

M. Benedikt @ CERN 13.02.24

Layout chosen out of ~ 100 initial variants, based on **geology** and **surface constraints** (land availability, access to roads, etc.), **environment**, (protected zones), **infrastructure** (water, electricity, transport), **machine performance** etc.

“Avoid-reduce-compensate” principle of EU and French regulation

Overall lowest-risk baseline: 90.7 km ring, 8 surface points.
Whole project now adapted to this placement



Optimized placement and layout.

M. Benedikt @ CERN 13.02.24

Layout chosen out of ~ 100 initial variants, based on **geology** and **surface constraints** (land availability, access to roads, etc.), **environment**, (protected zones), **infrastructure** (water, electricity, transport), **machine performance** etc.

“Avoid-reduce-compensate” principle of EU and French regulation

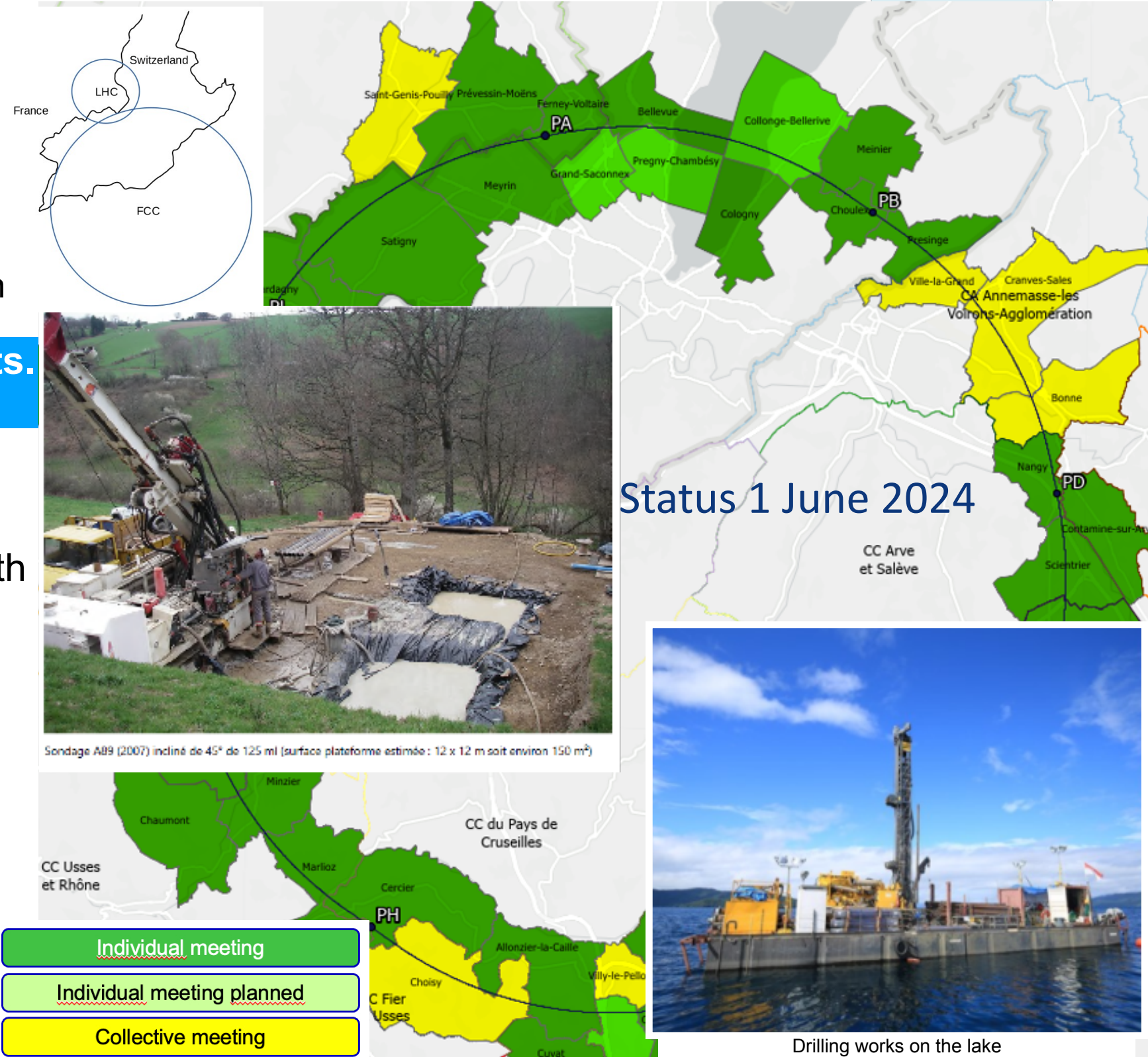
Overall lowest-risk baseline: 90.7 km ring, 8 surface points.
Whole project now adapted to this placement

- **Site investigations in areas with uncertain geological conditions:**

- ▶ Optimisation of localisation of drilling locations ongoing with site visits since end 2022.
- ▶ Alignment with FR and CH on the process for obtaining autorisation procedures. Ongoing for start of drillings in Q2/2024

- **Contracts Status:**

- ▶ Contract for engineering services and role of Engineer during works, active since July 2022
- ▶ Site investigations tendering ongoing towards contract placement in December 2023 and mobilization from January 2024

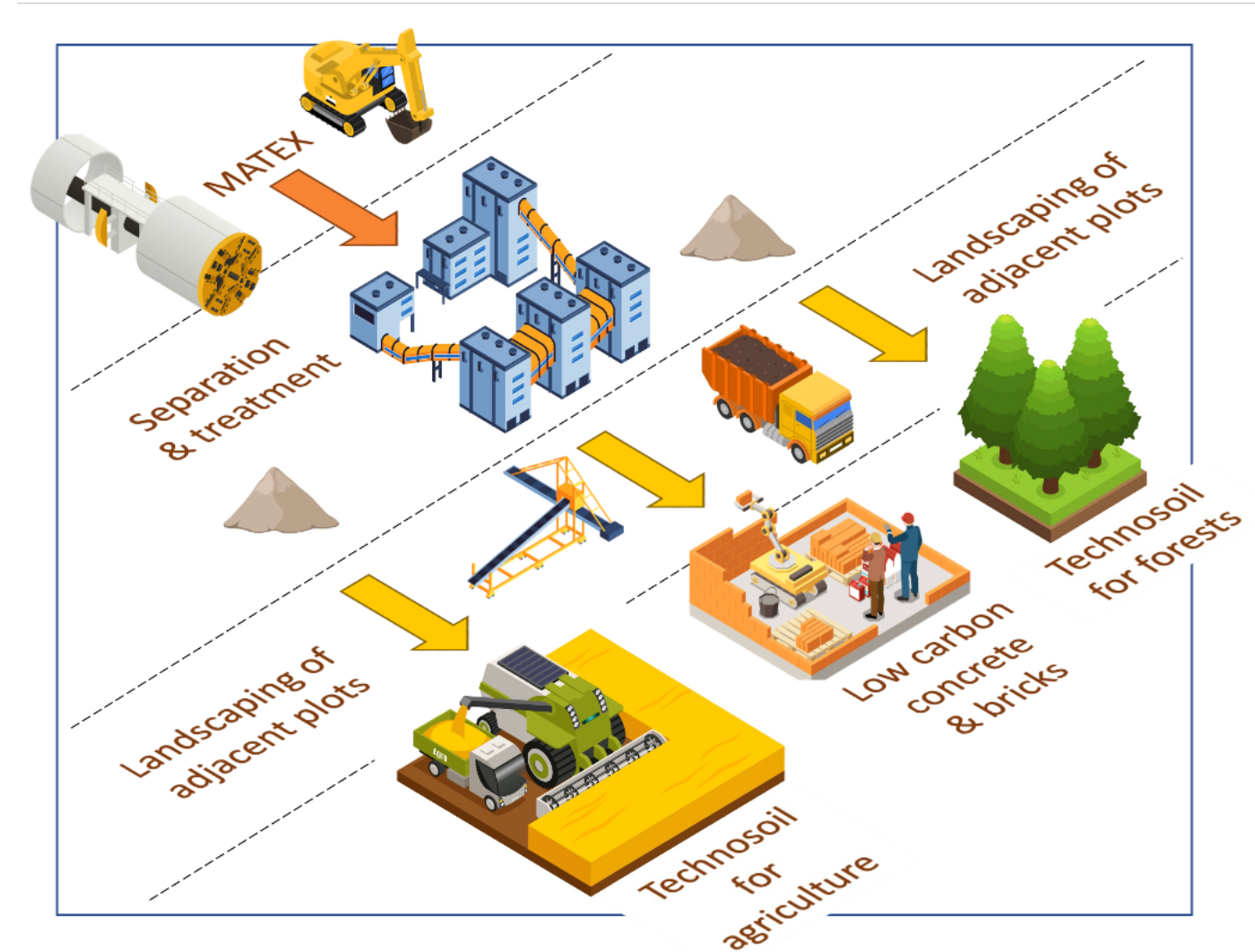


Status 1 June 2024

Environmental considerations.

M. Benedikt @ CERN 13.02.24

- **Excavated material** from FCC subsurface infrastructures: 6.5 Mm³ in situ, 8.4 Mm³ excavated
- **Priority : reuse, minimize disposal**
- 2021-2022: International competition “**Mining the Future**”, launched with the support of the EU Horizon 2020 grant, to find innovative and realistic ideas for the reuse of molasse (96% of excavated materials)
- 2023: “**OpenSky Laboratory**” project: Objective - Develop and test an innovative process to transform sterile “molasse” into fertile soil for agricultural use and afforestation. launched in Jan. 2024: 5500m² near LHC P5 in Cessy (FR). Trial with 5 000t of excavated local molasse → convert it to arable soil (agricultural/forestry)
- **Heat:**
 - heating for local houses
 - cheese factories in Jura and Haute-Savoie expressed special interest

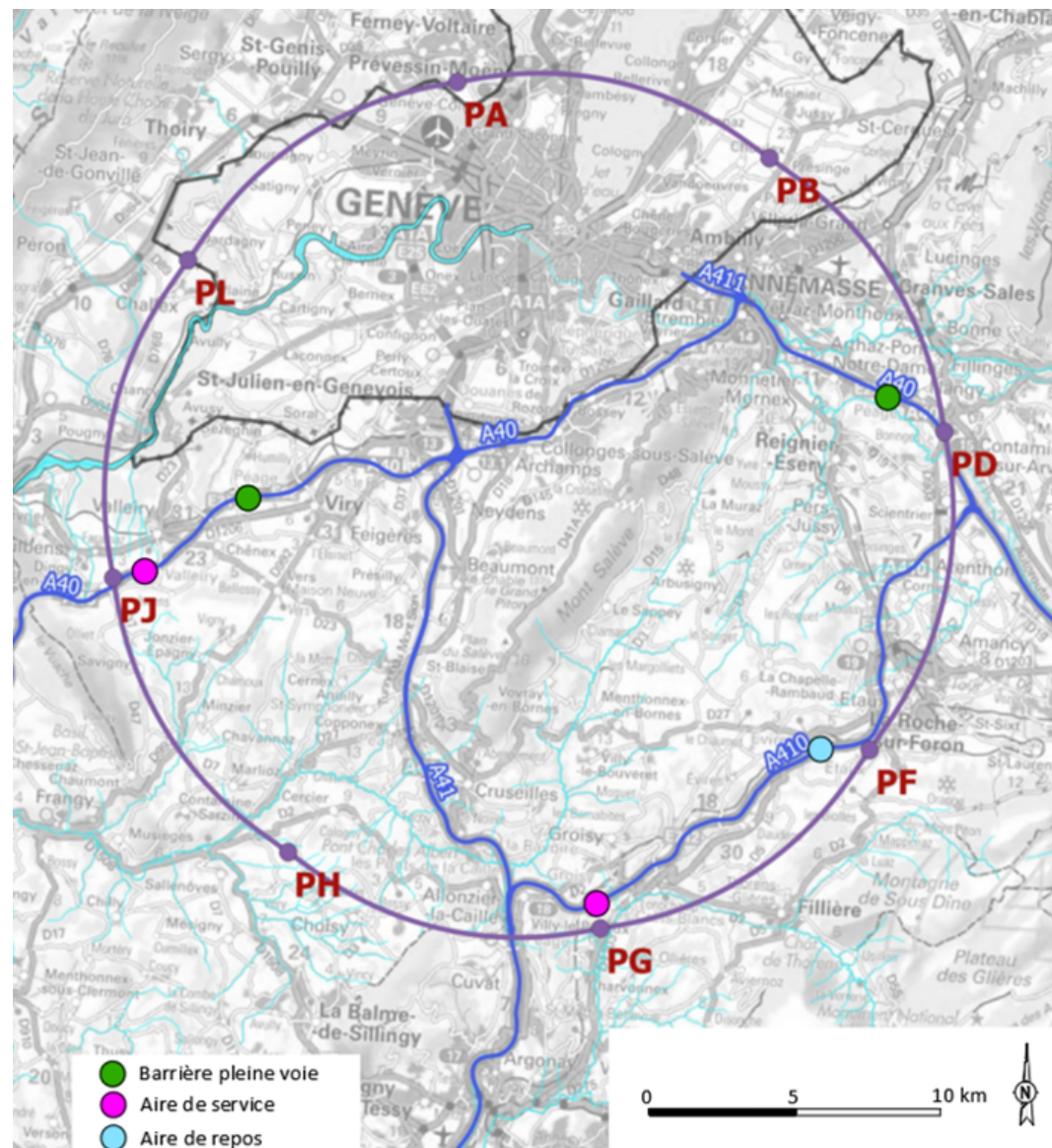


Accelerated soil transformation with funghi

Connections with local infrastructure.

M. Benedikt @ CERN 13.02.24

- **Road accesses** developed for all 8 surface sites
 - ▶ Four possible highway connections defined
 - ▶ Less than 4 km new departmental roads required

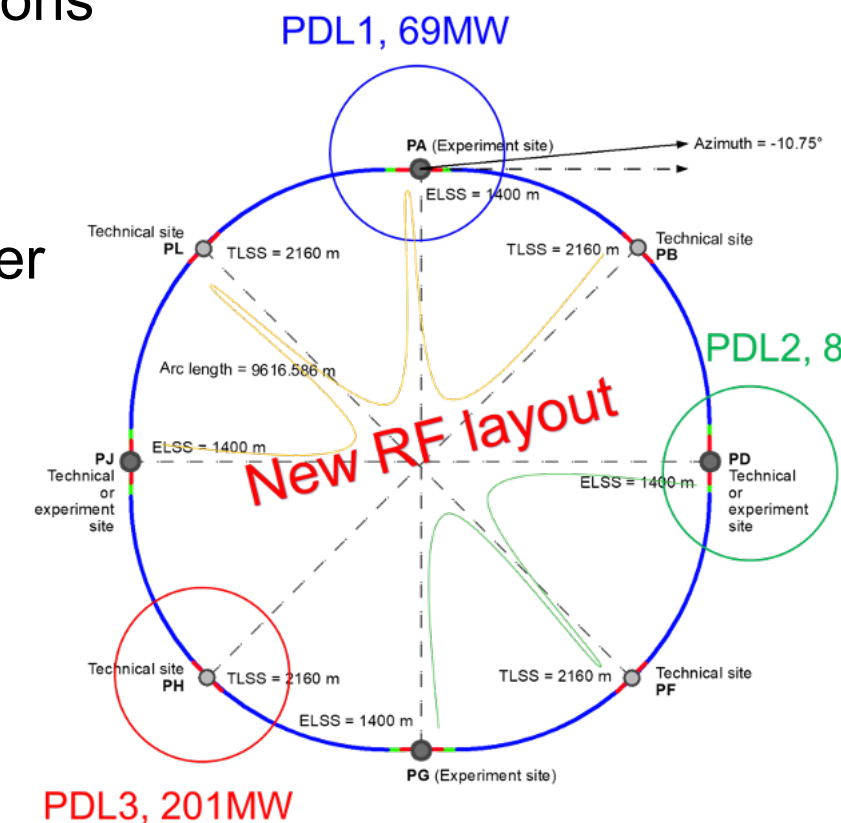


- **Connections to electrical grid**

- ▶ Electrical connection concept studied by RTE (French electrical grid operator) → requested loads have no significant impact on grid

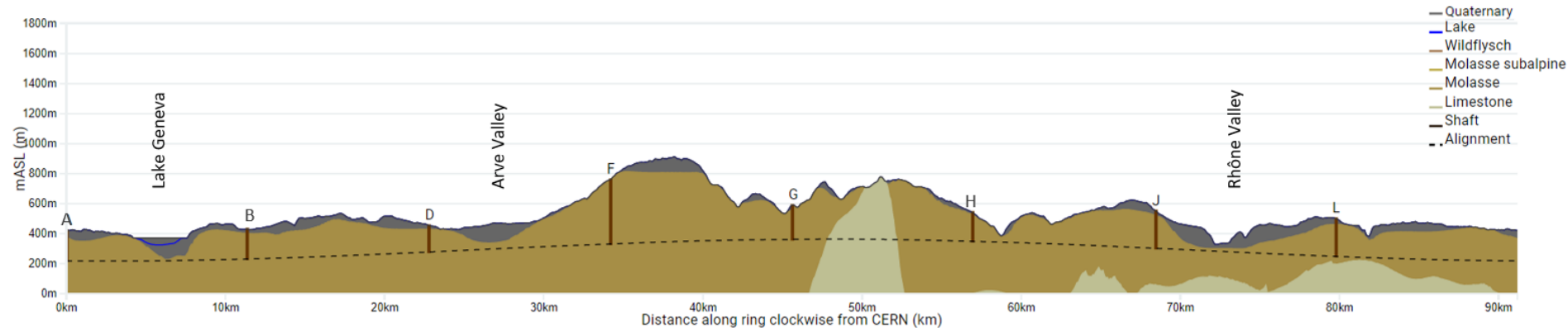
- ▶ Powering concept and power rating of the three sub-stations compatible with FCC-hh

- ▶ R&D efforts aiming at further reduction of the energy consumption of FCC-ee and FCC-hh



Civil engineering

T. Watson @ Anecy FCC Physics '24



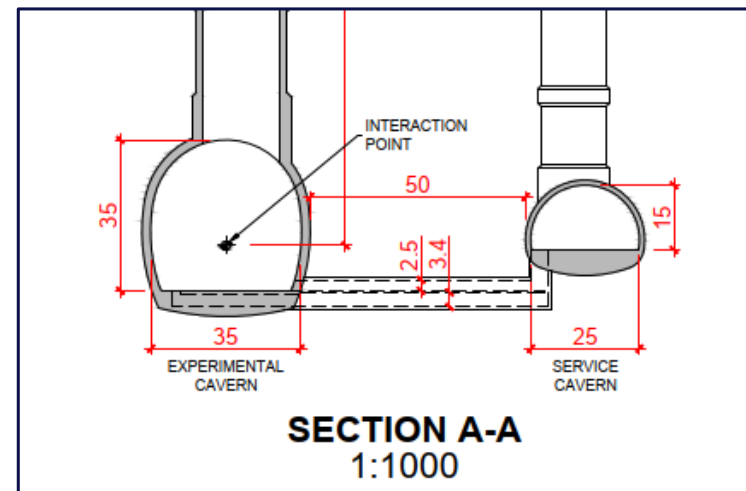
Shaft depths:

A: 201 m B: 201 m D: 181 m F: 400 m G: 226 m H: 235 m J: 253 m L: 250 m

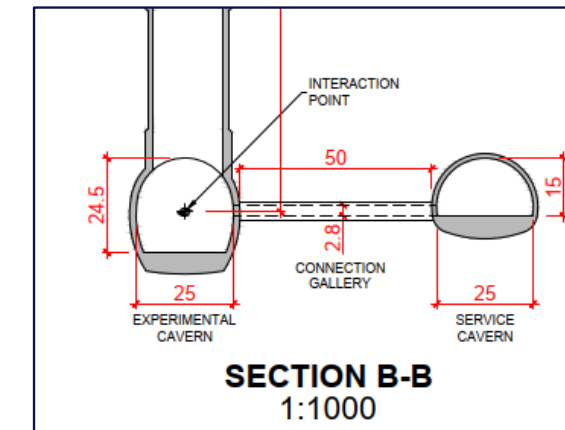


Tunnel Boring Machine (TBM)

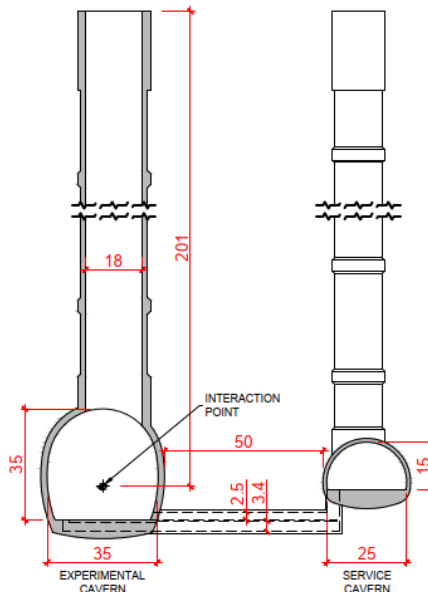
- Tunnel Boring Machines (TBMs) are designed to work almost continuous 24/7 other than periodic maintenance. Rate of 18m/day in the Molasse → 8 years.
- 13 shafts
- 2/2 large/small caverns



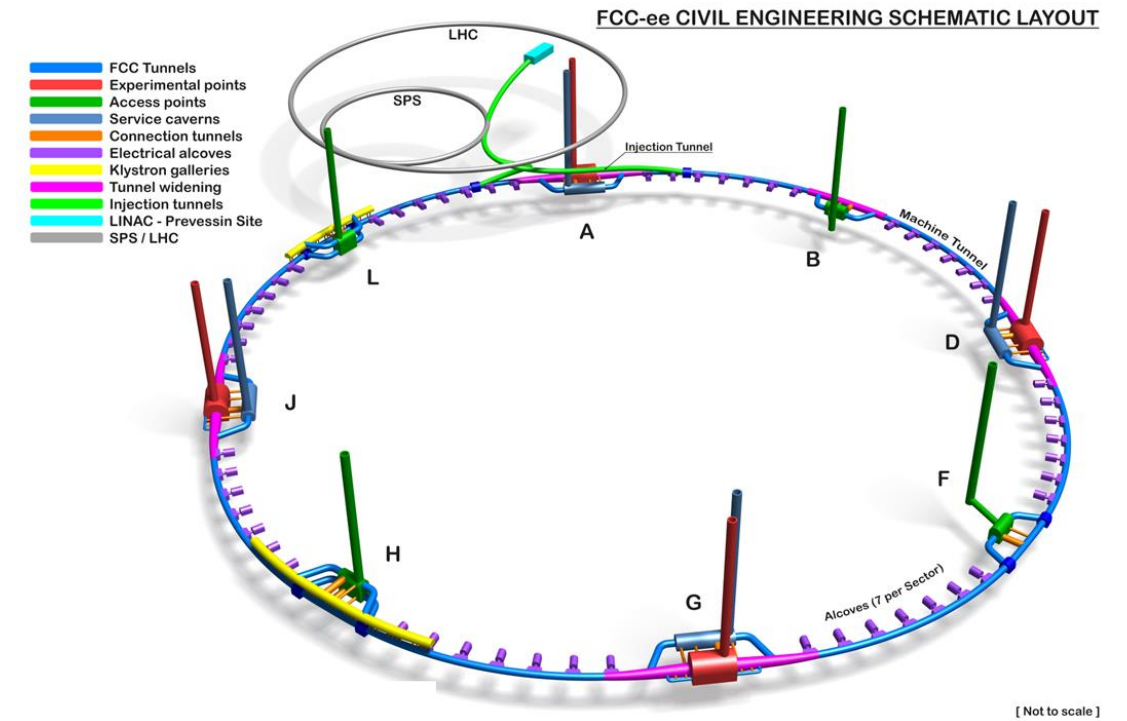
large cavern complex



small cavern complex



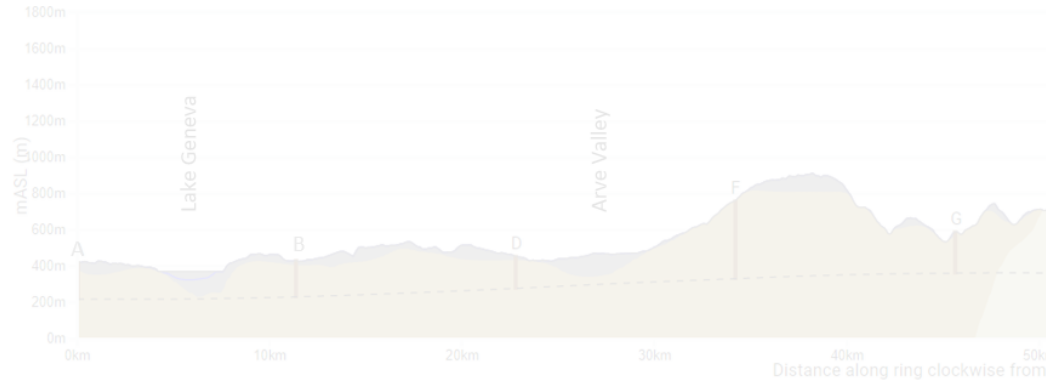
shaft @ exp. site



[Not to scale]

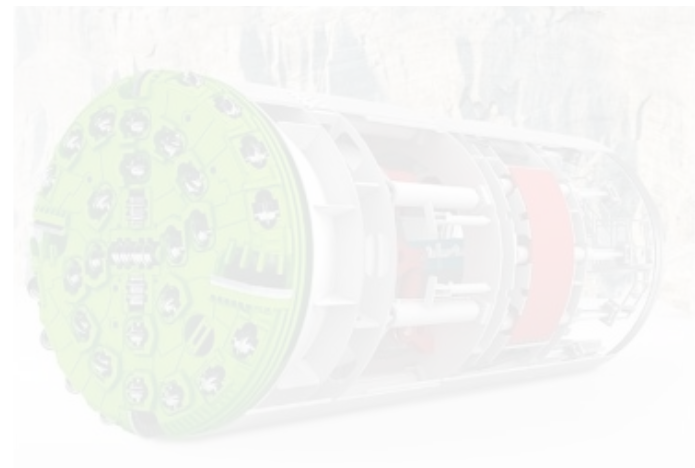
Civil engineering

T. Watson @ Anecy FCC Physics '24



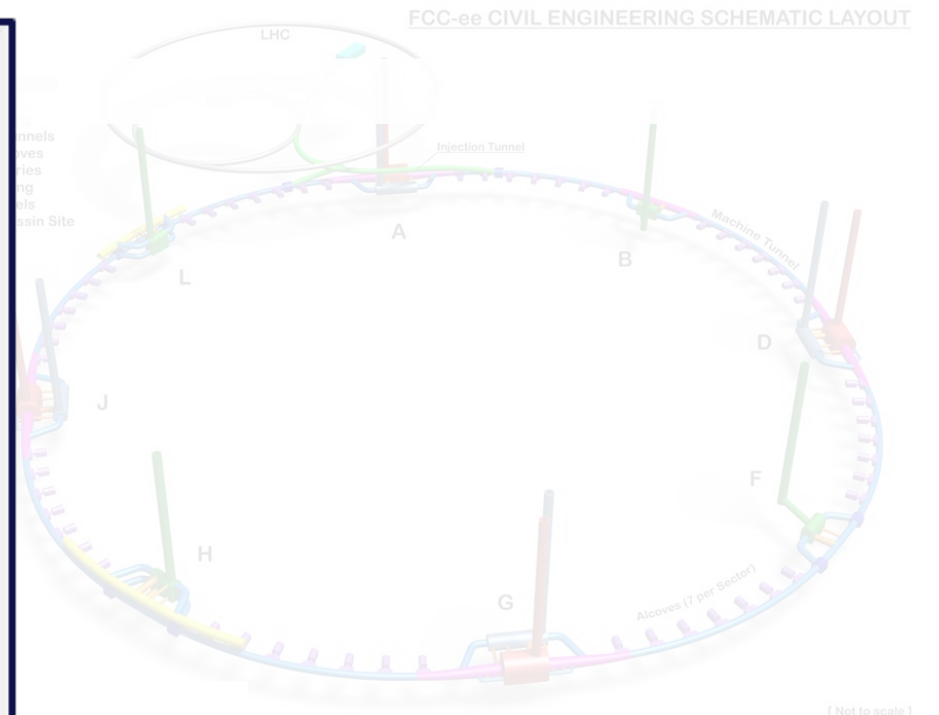
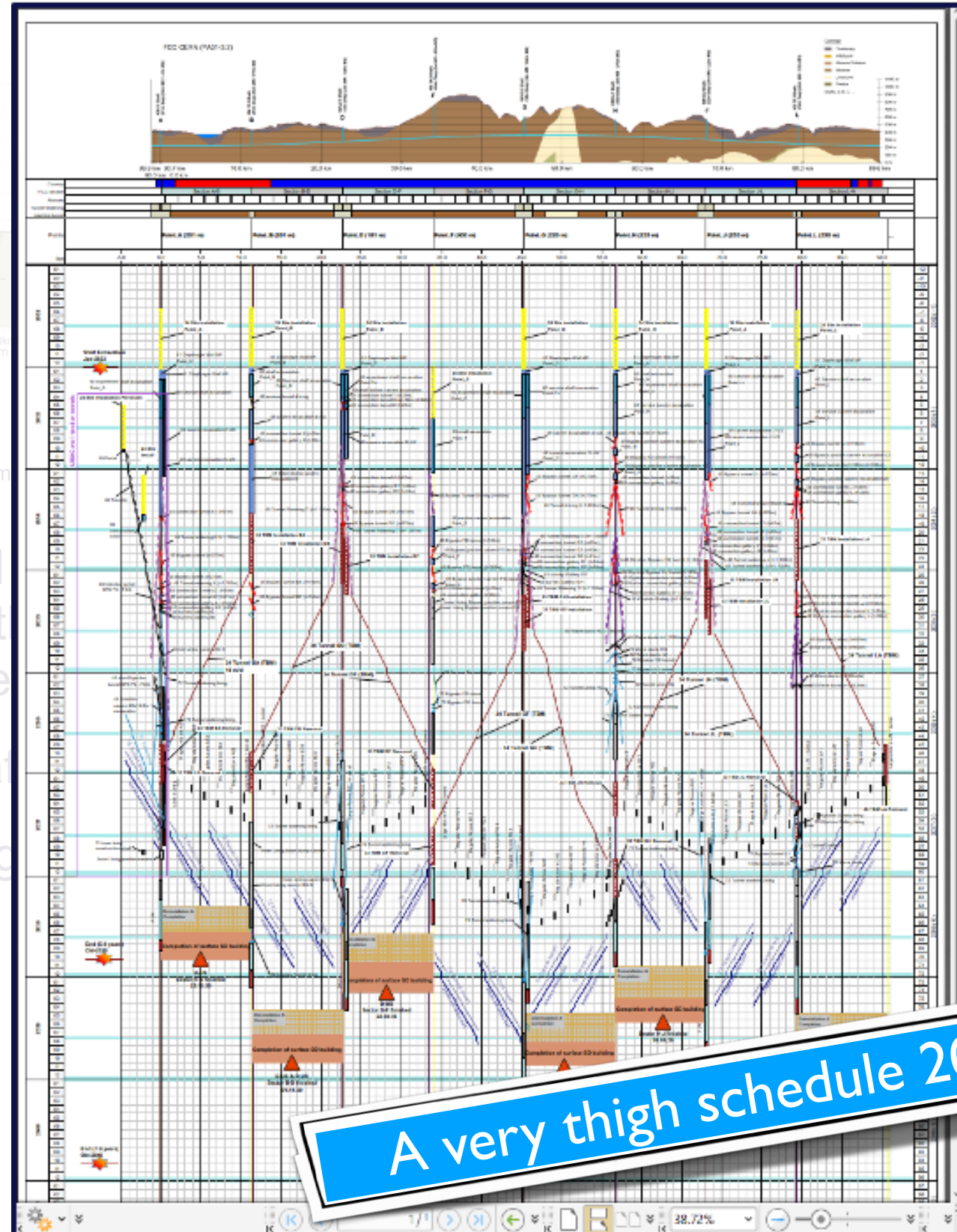
Shaft depths:

A: 201 m B: 201 m D: 181 m F: 400 m G: 226 m

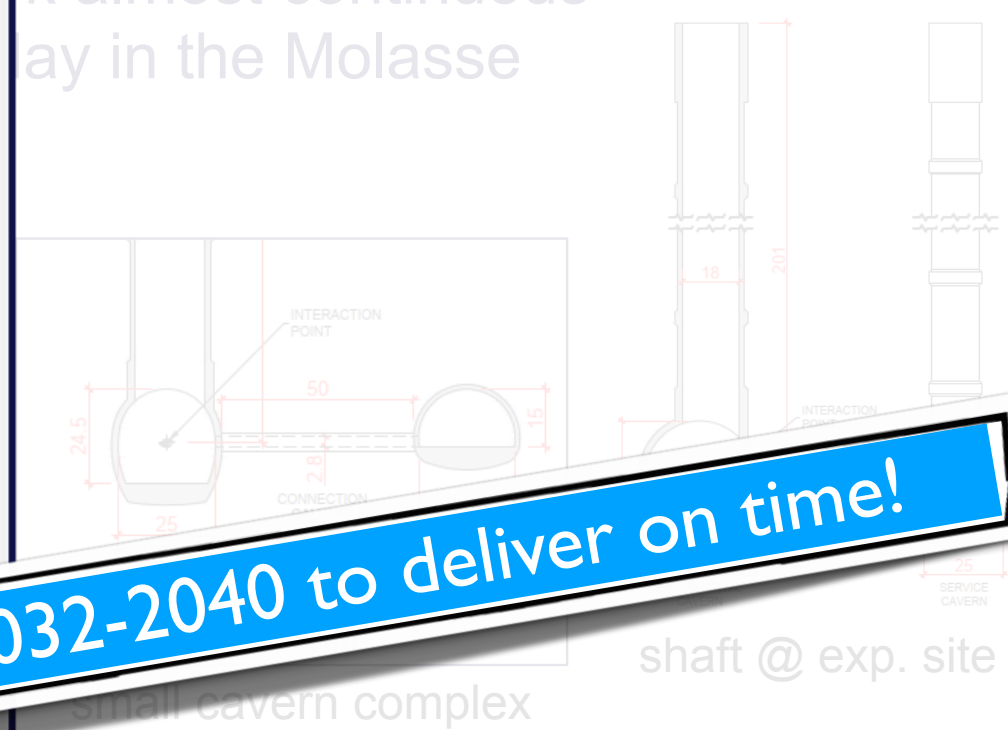


Tunnel Boring Machine (TBM)

- Tunnel 24/7 operation → 8 years
- 13 shafts
- 2/2 large



Work almost continuous day in the Molasse



A very tight schedule 2032-2040 to deliver on time!

FCC feasibility mid-term report.

- **703 pages:** 7 chapters (cost and financial feasibility is a separate document) + refs.

- Placement scenario (75 pages)
- Civil engineering (50 pages)
- Implementation with the host states (45 pages)
- Technical infrastructure (110 pages)
- FCC-ee collider design and performance (170 pages)
- FCC-hh accelerator (60 pages)
- (Cost and financial feasibility)
- Physics and experiments (110 pages)
- References (70 pages)

- **Executive summary:** 44 pages

- Reviewed by

- Scientific Advisory Committee and Cost Review Panel on Oct. 16-18
- Scientific Policy Committee and Financial Committee on Nov. 21-22
- CERN Council Feb. 2

Future Circular Collider Midterm Report

February 2024

528 authors
16 editors

Edited by:

B. Auchmann, W. Bartmann, M. Benedikt, J.P. Burnet, P. Craievich,
M. Giovannozzi, C. Grojean, J. Gutleber, K. Hanke, P. Janot, M. Mangano,
J. Osborne, J. Poole, T. Raubenheimer, T. Watson, F. Zimmermann



This project has received funding under the European Union's
Horizon 2020 research and innovation programme under grant
agreement No 951754.

This document has been produced by the organisations participating in the
FCC feasibility study. The studies and technical concepts presented here
do not represent an agreement or commitment of any of CERN's Member
States or of the European Union for the construction and operation of an
extension to CERN's existing research infrastructures.

The midterm report of the FCC Feasibility Study reflects work in progress
and should therefore not be propagated to people who do not have direct
access to this document.

**confidential documents
(work in progress)
available
to CERN personnel**

Physics, Experiments, Detectors.

- FCC Feasibility Study PED deliverables for mid-term review

8. Physics & Experiments	C. Grojean, P. Janot, M. Mangano	8.1 Overview	} deliverables explicitly requested from SPC & Council
		8.2. Documentation of the specificities of the FCC-ee and FCC-hh physics cases.	
		8.3 Strategic plans for the improved theoretical calculations.	
		8.4 FCC-ee Detector Requirements.	

- Content of the mid-term PED chapter (60 pages were expected → 110 pages delivered)

Physics, Experiments, Detectors.

- FCC Feasibility Study PED deliverables for mid-term review

8. Physics & Experiments	C. Grojean, P. Janot, M. Mangano	8.1 Overview	} deliverables explicitly requested from SPC & Council
		8.2. Documentation of the specificities of the FCC-ee and FCC-hh physics cases.	
		8.3 Strategic plans for the improved theoretical calculations.	
		8.4 FCC-ee Detector Requirements.	

- Content of the mid-term PED chapter (60 pages were expected → 110 pages delivered)

1 Overview	3	4 Detector requirements	54
1.1 FCC-ee: A great Higgs factory, and so much more	4	4.1 Introduction	54
1.2 FCC-hh: The energy-frontier collider with the broadest exploration potential	13	4.2 Machine-detector interface	55
2 Specificities of the FCC physics case	15	4.3 The current detector concepts	56
2.1 Characterisation of the Higgs boson: role of EW measurements and of FCC-hh	16	4.4 Measurement of the tracks of charged particles	58
2.2 Discovery landscape	24	4.5 Requirements on the vertex detector	64
2.3 Flavour advancement	34	4.6 Requirements on charged hadron particle identification	73
2.4 FCC-hh specificities compared to lepton colliders	36	4.7 Requirements on electromagnetic calorimetry	78
3 Theoretical calculations	42	4.8 Requirements on the hadronic calorimeter	88
3.1 Electroweak corrections	44	4.9 Requirements on the muon detector	93
3.2 QCD precision calculations	46	4.10 Precise timing measurements	93
3.3 Monte Carlo event generators	50	5 Outlook and further steps	96
3.4 Organization and support of future activities to improve theoretical precision	53	5.1 Software and Computing	98
		5.2 Physics Performance	99
		5.3 Detector Concepts	101
		5.4 Centre-of-mass energy calibration, polarisation, monochromatisation (EPOL)	103
		5.5 Machine-Detector Interface (MDI)	104
		5.6 Physics Programme	105
		5.7 FCC-hh	106

Physics, Experiments, Detectors.

- FCC Feasibility Study PED deliverables for mid-term review

8. Physics & Experiments	C. Grojean, P. Janot, M. Mangano	8.1 Overview	deliverables explicitly requested from SPC & Council
		8.2. Documentation of the specificities of the FCC-ee and FCC-hh physics cases.	
		8.3 Strategic plans for the improved theoretical calculations.	
		8.4 FCC-ee Detector Requirements.	

- Content of the mid-term PED chapter (90 pages delivered)

1 Overview	54
1.1 FCC-ee	54
1.2 FCC-hh	55
2 Specificities	56
2.1 FCC-ee	58
2.2 FCC-hh	64
2.3 FCC-ee/hh as a Higgs factory	73
2.4 FCC-ee/hh as a flavour factory	78
3 Theoretical potential	88
3.1 Electroweak	93
3.2 QCD	93
3.3 Monochromatisation	96
3.4 Organisation and precision	98
4 Machine-Detector Interface (MDI)	99
5 Physics Programme	101
5.5 Machine-Detector Interface (MDI)	103
5.6 Physics Programme	104
5.7 FCC-hh	105
5.8 FCC-ee ↔ FCC-hh: complementarity and synergy	106

— Main physics topics covered/discussed in the report —

1. FCC-ee: much more than a Higgs factory:
 - precision for discovery
 - tera-Z direct discovery potential
2. FCC-ee/hh as a Higgs/electroweak factory
3. FCC-ee as a flavour factory
4. FCC-hh: the broadest exploration potential at high-energy
5. FCC-ee ↔ FCC-hh: complementarity and synergy

Feedback.

Andy **Parker** (SAC chair), Norbert **Holtkamp** (CRP chair), Hugh **Montgomery** (SPC chair), Laurent **Salzarulo** (FC chair), Eliezer **Rabinovici** (Council president)

“many thanks for the work done, congratulations for the results, impressive quality of the study...”

“Financial Committee underlines the need to make the project attractive from the physics viewpoint and takes the view that it would be unfortunate to sacrifice the attractiveness of the physics for the sake of reducing costs.”



“Si j’ai voulu venir là aujourd’hui c’est pour témoigner ma confiance aux équipes et notre volonté, notre ambition de conserver la première place dans ce domaine.”
[“My visit here bears witness to my trust in CERN personnel and France’s will and ambition to keep the leadership in this domain.”]

E. Macron, CERN 16.11.2023

US Statement of Intent



Deirdre Mulligan

Fabiola Gianotti

“Should the CERN Member States determine the FCC-ee is likely to be CERN’s next world-leading research facility following the high-luminosity Large Hadron Collider, the United States intends to collaborate on its construction and physics exploitation, subject to appropriate domestic approvals.”

White House, April 26, 2024

EU Support

European Union Competitiveness Report

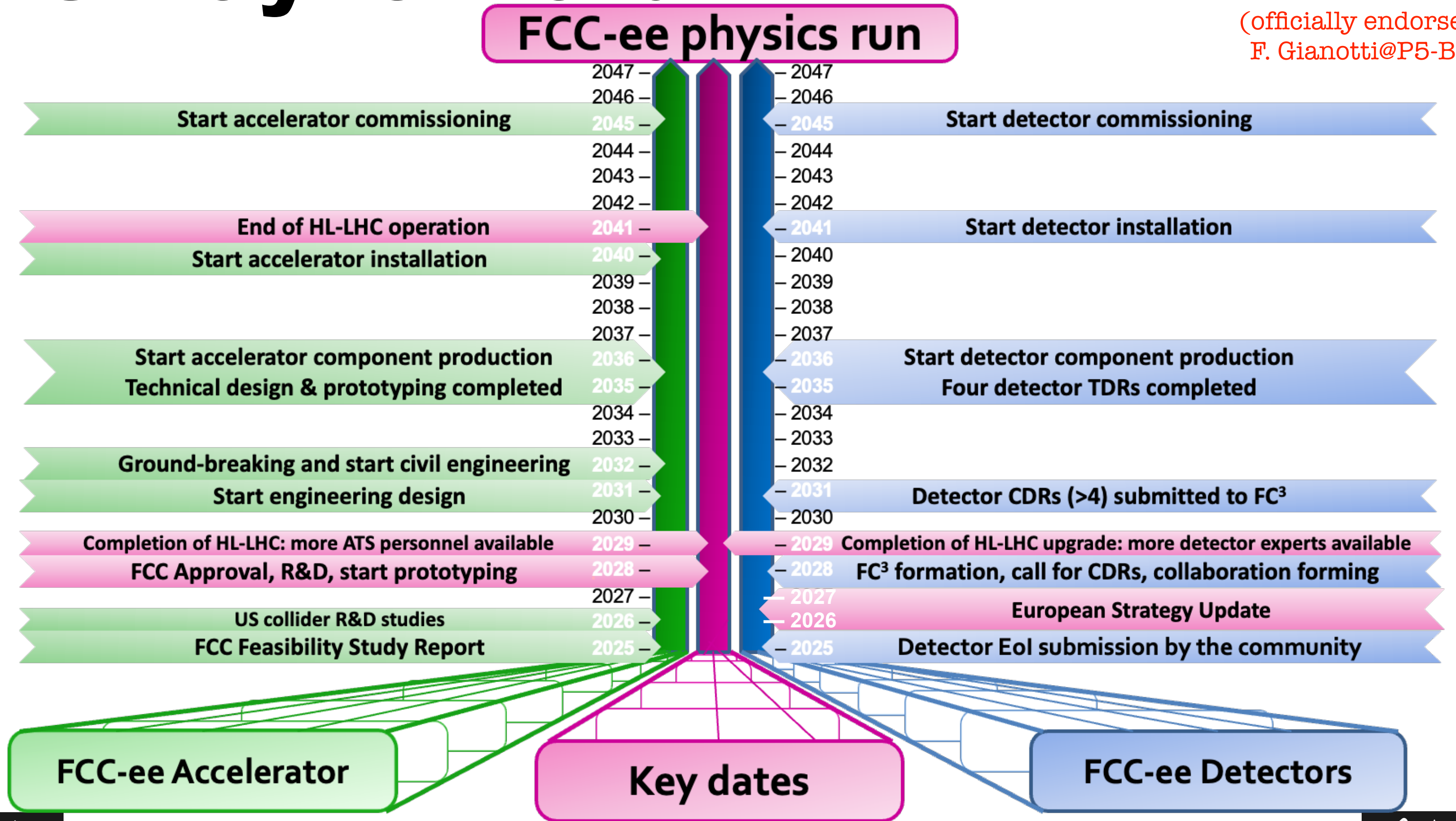
- **400-page report made public by Mario Draghi on Monday 9 September 2024:**
 - Handed to Ursula von der Leyen (European Commission president) for subsequent action
 - Urges the EU to invest 800 billion euros annually [with specific guidance] to close the economic gap between the US and China (consistently seen as a threat throughout the report)
 - CERN mentioned 19 times in the report (and FCC 3 times, unlike any other big scientific facility)!

“Refinancing CERN and ensuring its continued global leadership in frontier research should be regarded as a top EU priority.”

“One of CERN’s most promising current projects, with significant scientific potential, is the construction of the Future Circular Collider (FCC): a 90-km ring designed initially for an electron collider and later for a hadron collider.”

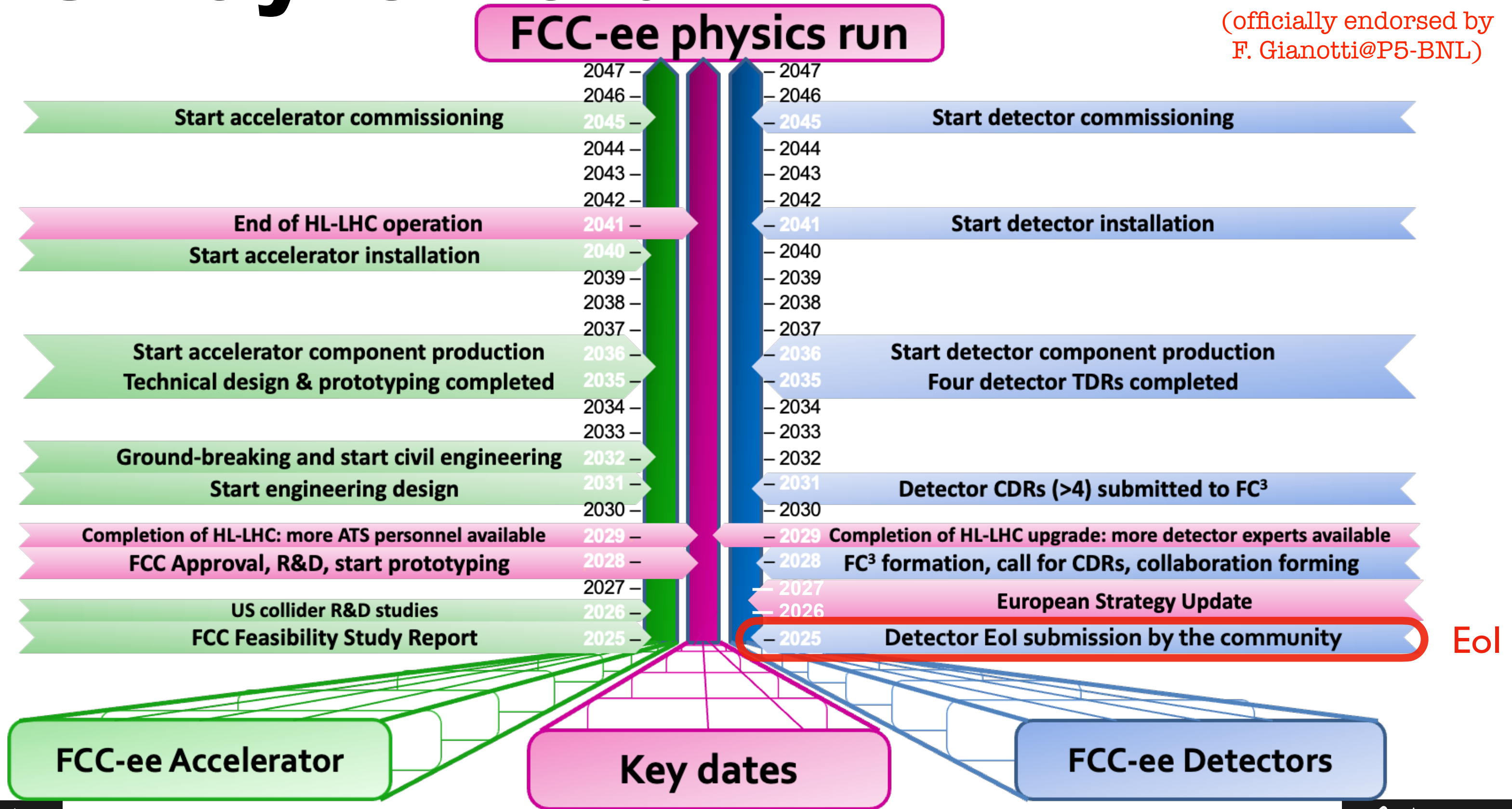
The way forward.

P. Janot
 (officially endorsed by
 F. Gianotti@P5-BNL)



The way forward.

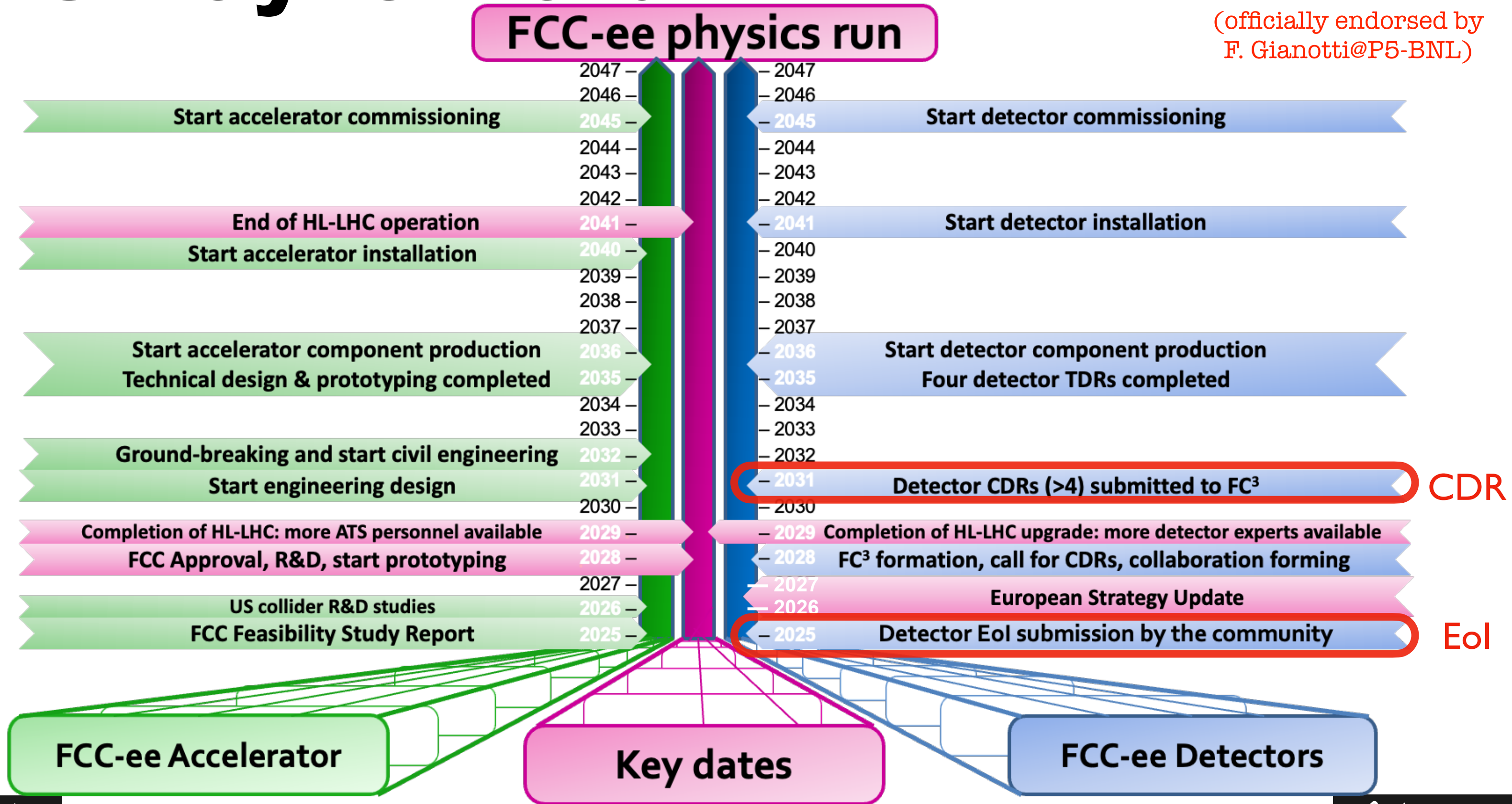
P. Janot
 (officially endorsed by
 F. Gianotti@P5-BNL)



Eol

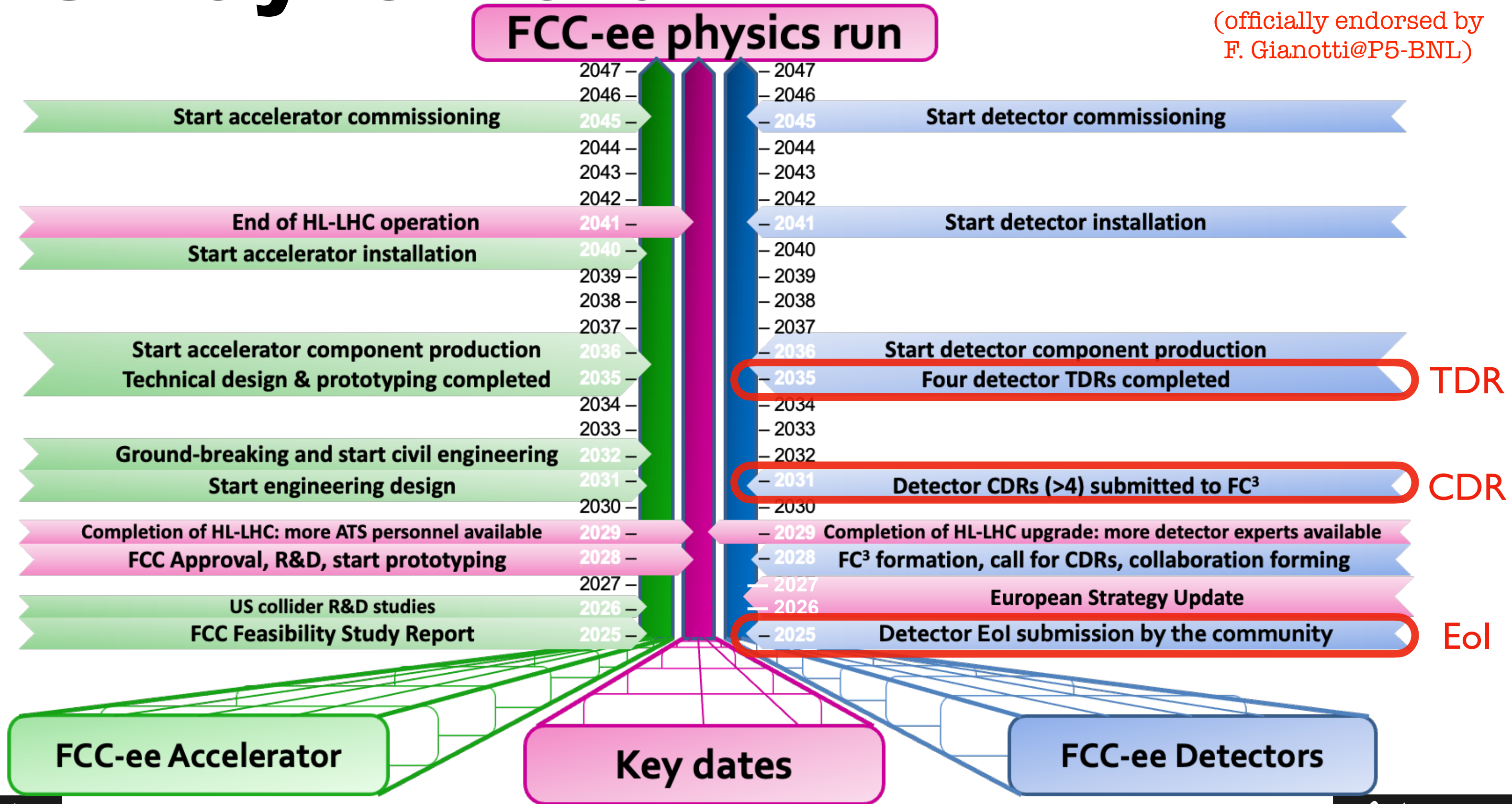
The way forward.

P. Janot
 (officially endorsed by
 F. Gianotti@P5-BNL)



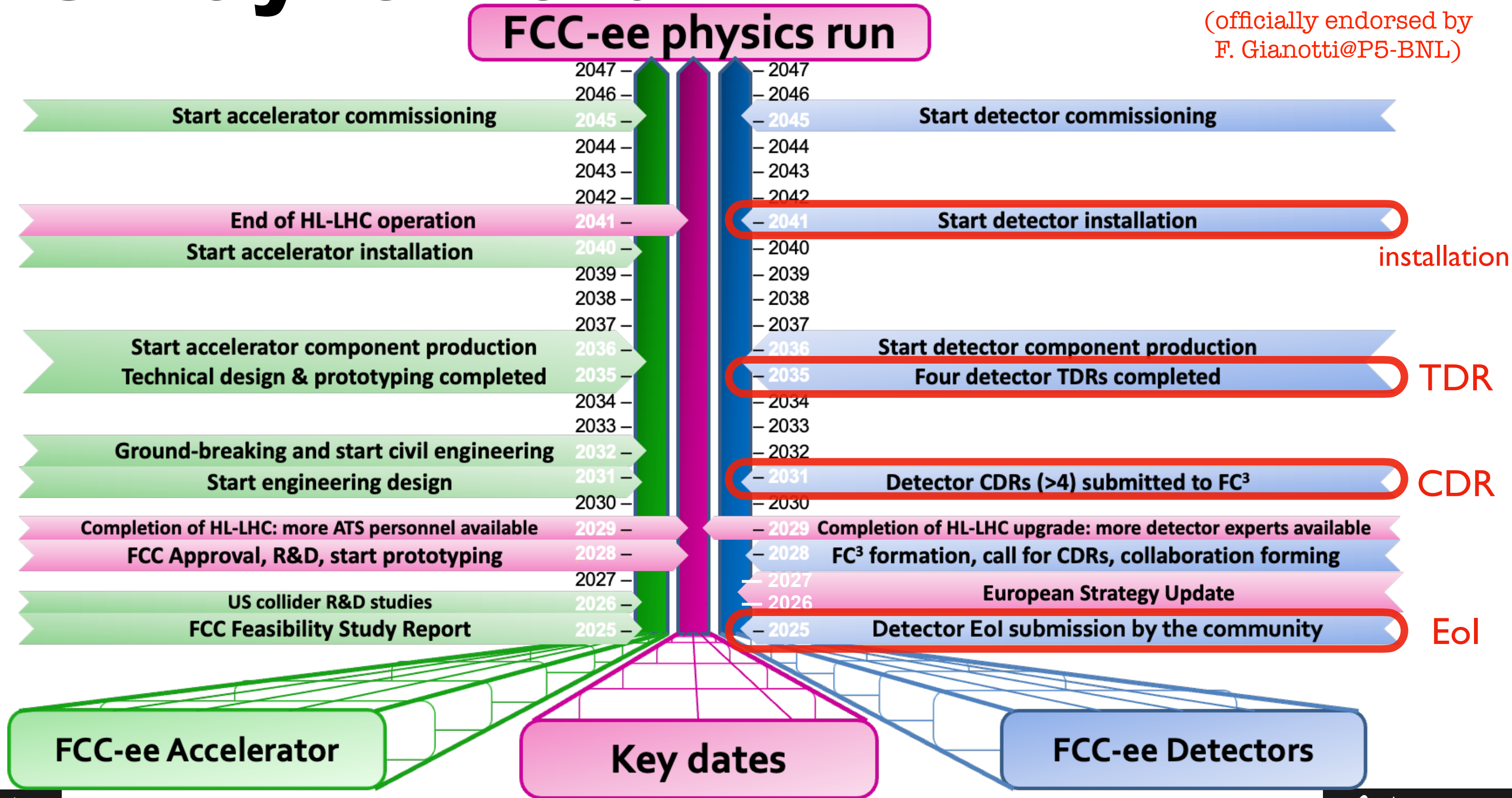
The way forward.

P. Janot
 (officially endorsed by
 F. Gianotti@P5-BNL)



The way forward.

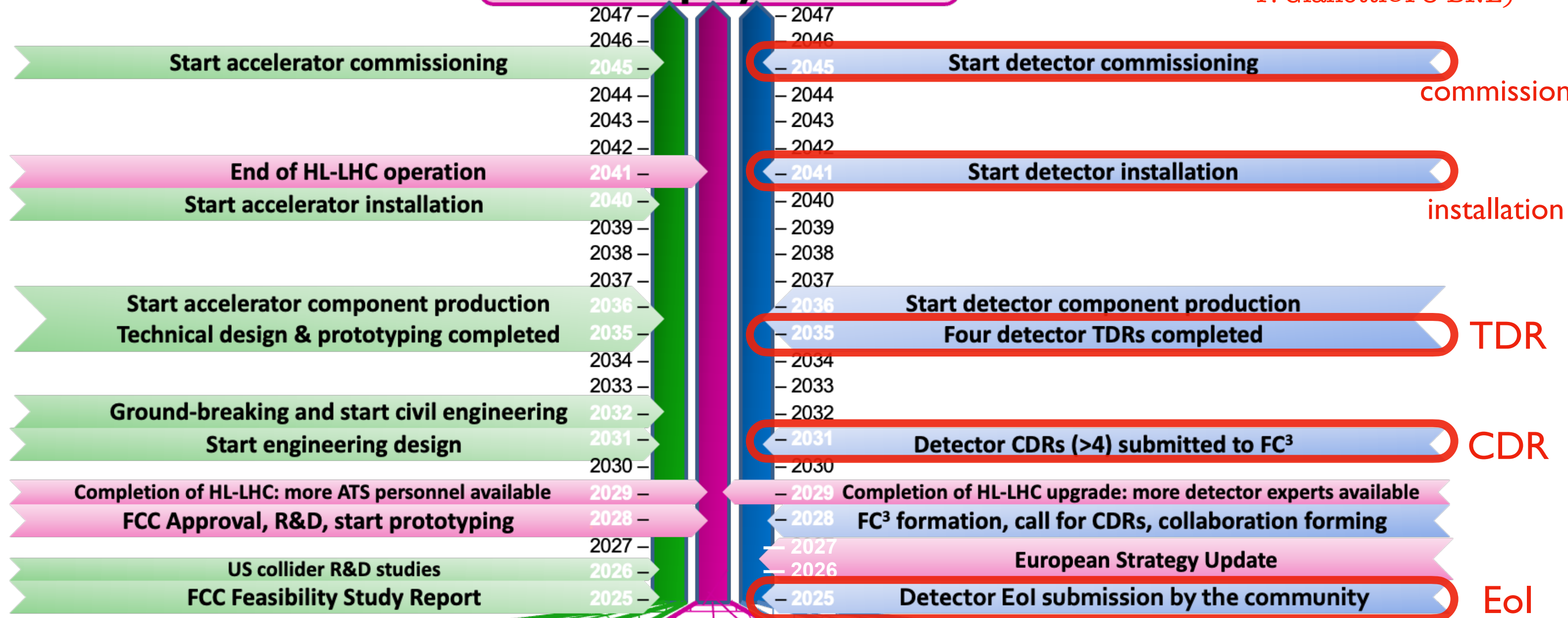
P. Janot
 (officially endorsed by
 F. Gianotti@P5-BNL)



The way forward.

P. Janot
 (officially endorsed by
 F. Gianotti@P5-BNL)

FCC-ee physics run



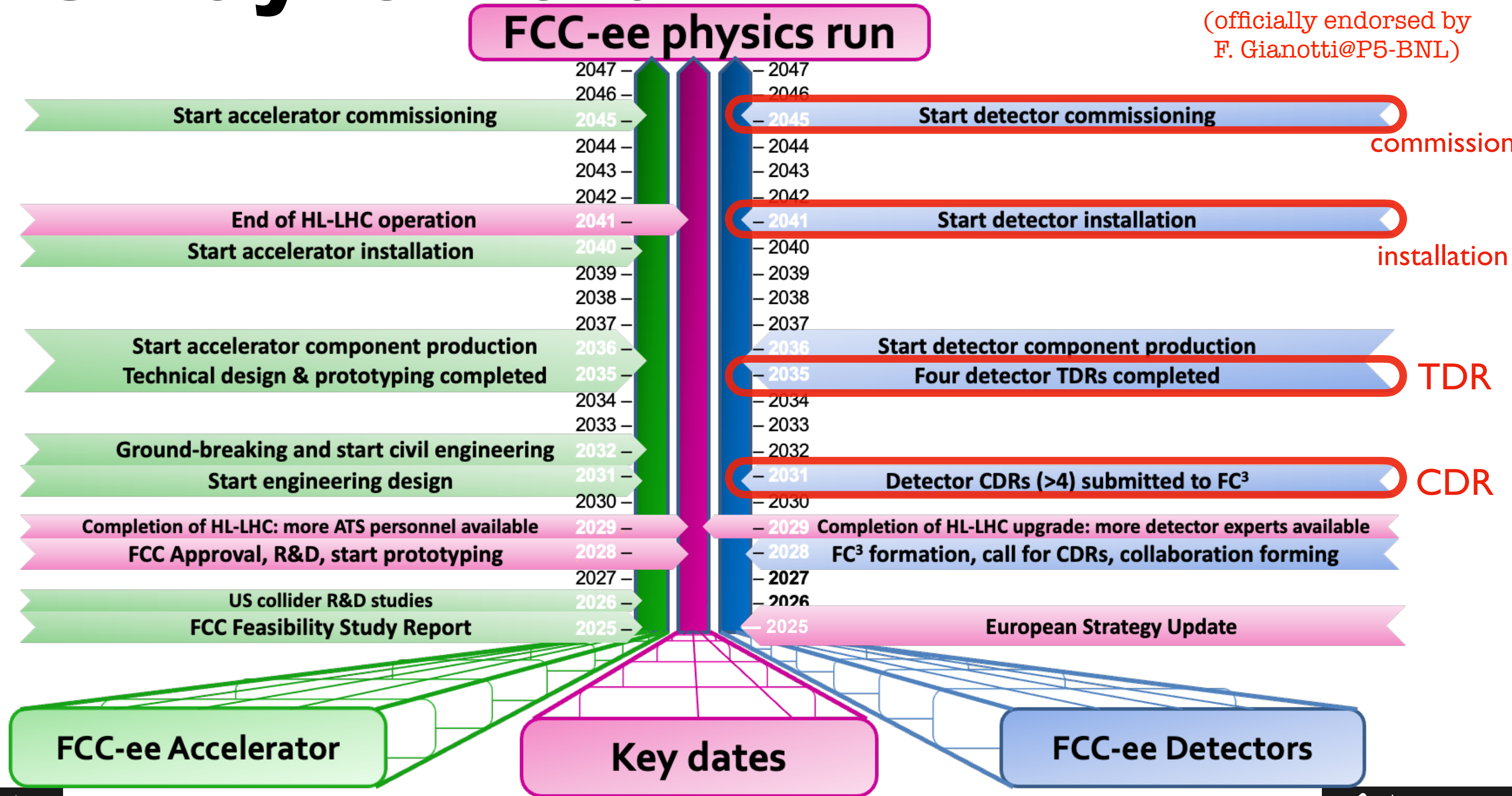
FCC-ee Accelerator

Key dates

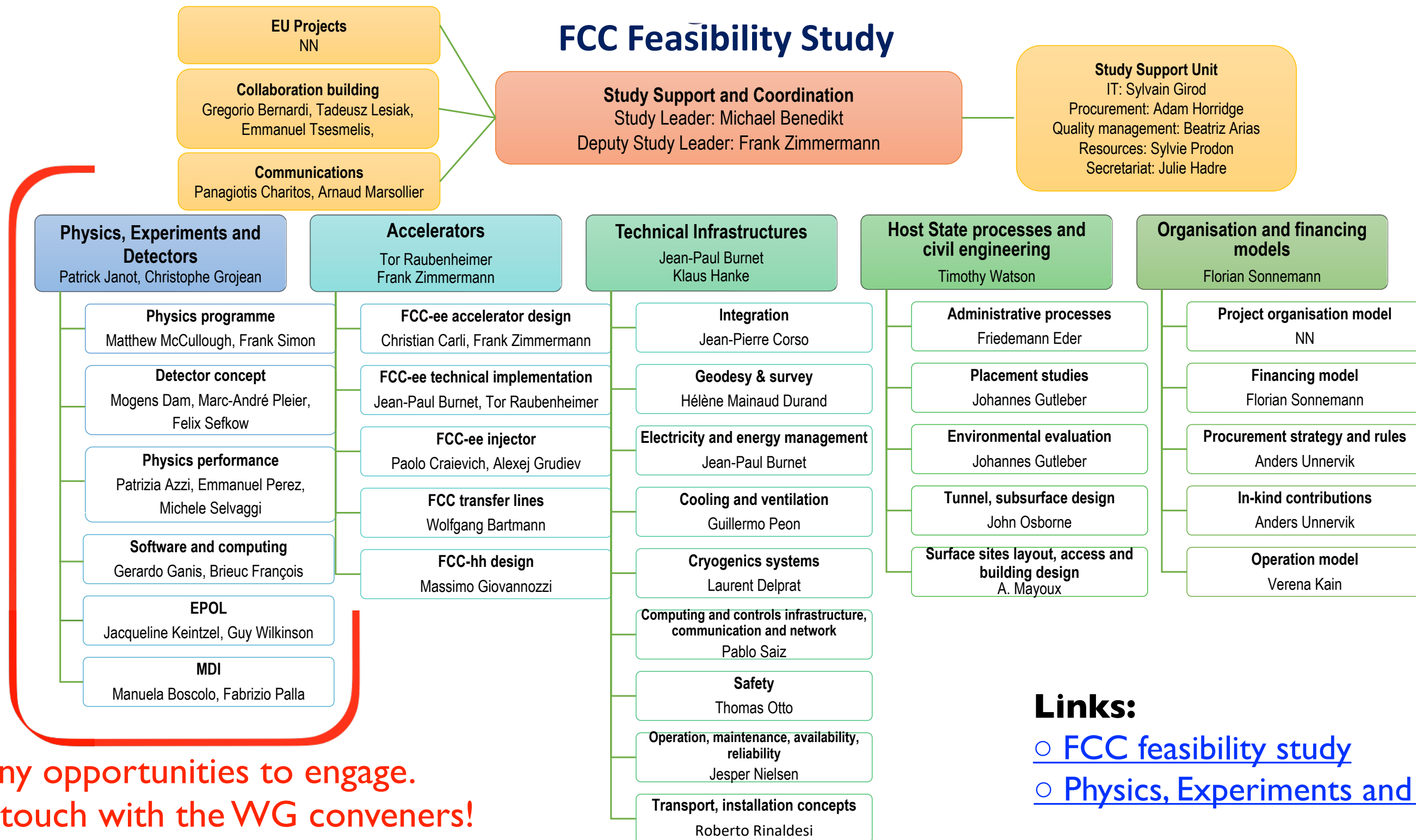
FCC-ee Detectors

The way forward.

P. Janot
(officially endorsed by
F. Gianotti@P5-BNL)



FCC Feasibility Organisation Chart.



Many opportunities to engage.
Get in touch with the WG conveners!

Links:

- [FCC feasibility study](#)
- [Physics, Experiments and Detectors \(PED\)](#)

FCC: an international enterprise.

Increasing international collaboration as a prerequisite for success:

→ links with **science, research & development** and **high-tech industry** will be essential to further advance and prepare the implementation of FCC



FCC Feasibility Study:

Aim is to increase further the collaboration, on all aspects, in particular on Accelerator and Particle/Experiments/Detectors

141
Institutes

32
countries
+
CERN



Construction Cost/Cost of Operation

Construction cost

Main domains for the FCC-ee project :

- Accelerators: 3 847 MCHF
- Injectors & transfer lines: 585 MCHF
- Civil engineering: 5 538 MCHF
- Technical infrastructures: 2 490 MCHF
- Experiments: 150 MCHF
- Territorial development: 191 MCHF

The total cost for FCC-ee, considering two IPs for experiments and the first three stages of operation (Z, W and ZH) is estimated to be **12 801 MCHF**.

- 2 → 4 IPs: + 710 MCHF
- 365 GeV run: +1 465 MCHF

Energy and carbon footprint

- FCC-ee total instantaneous power demand at each centre-of-mass energies

J.-P. Burnet, FCC Week'22

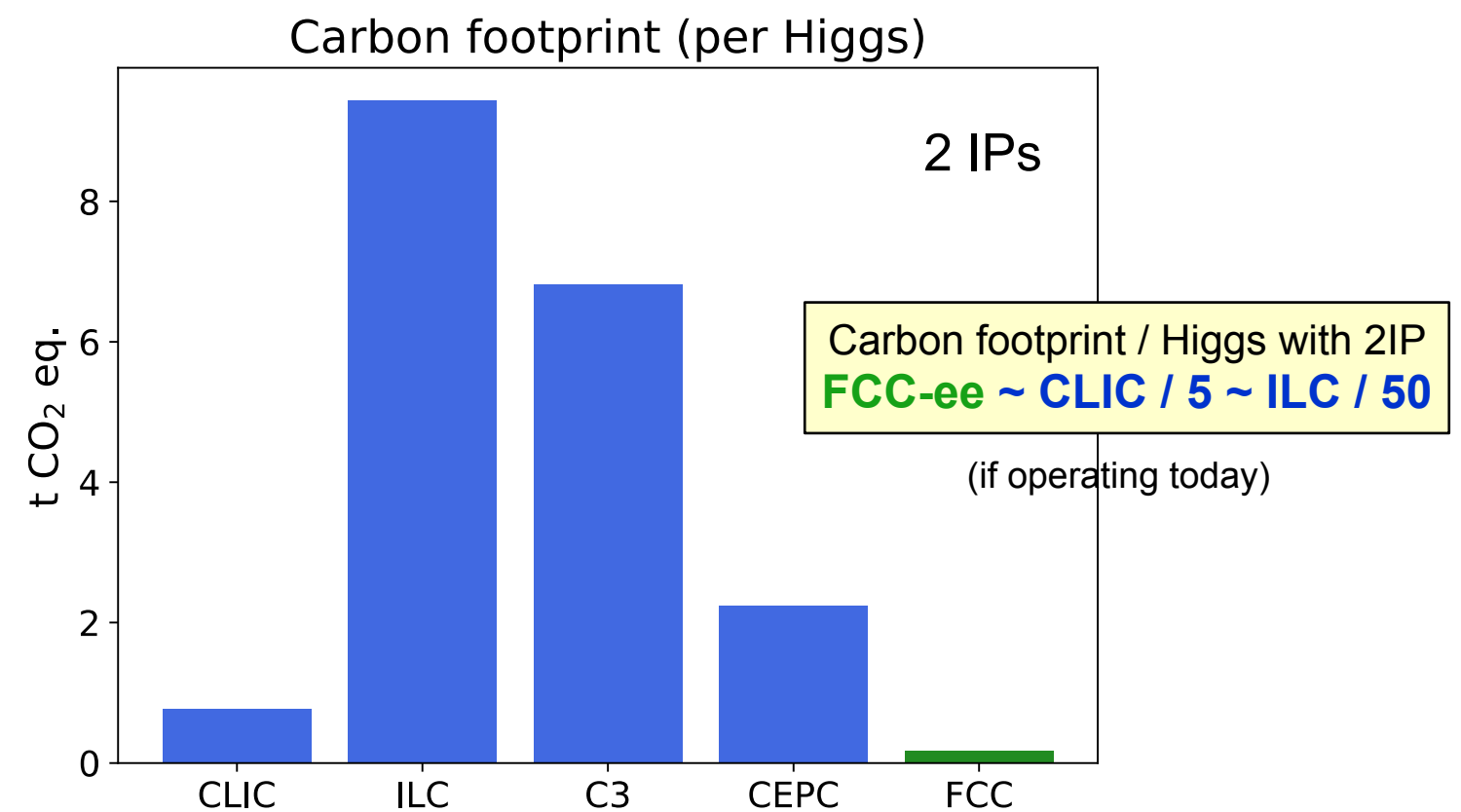
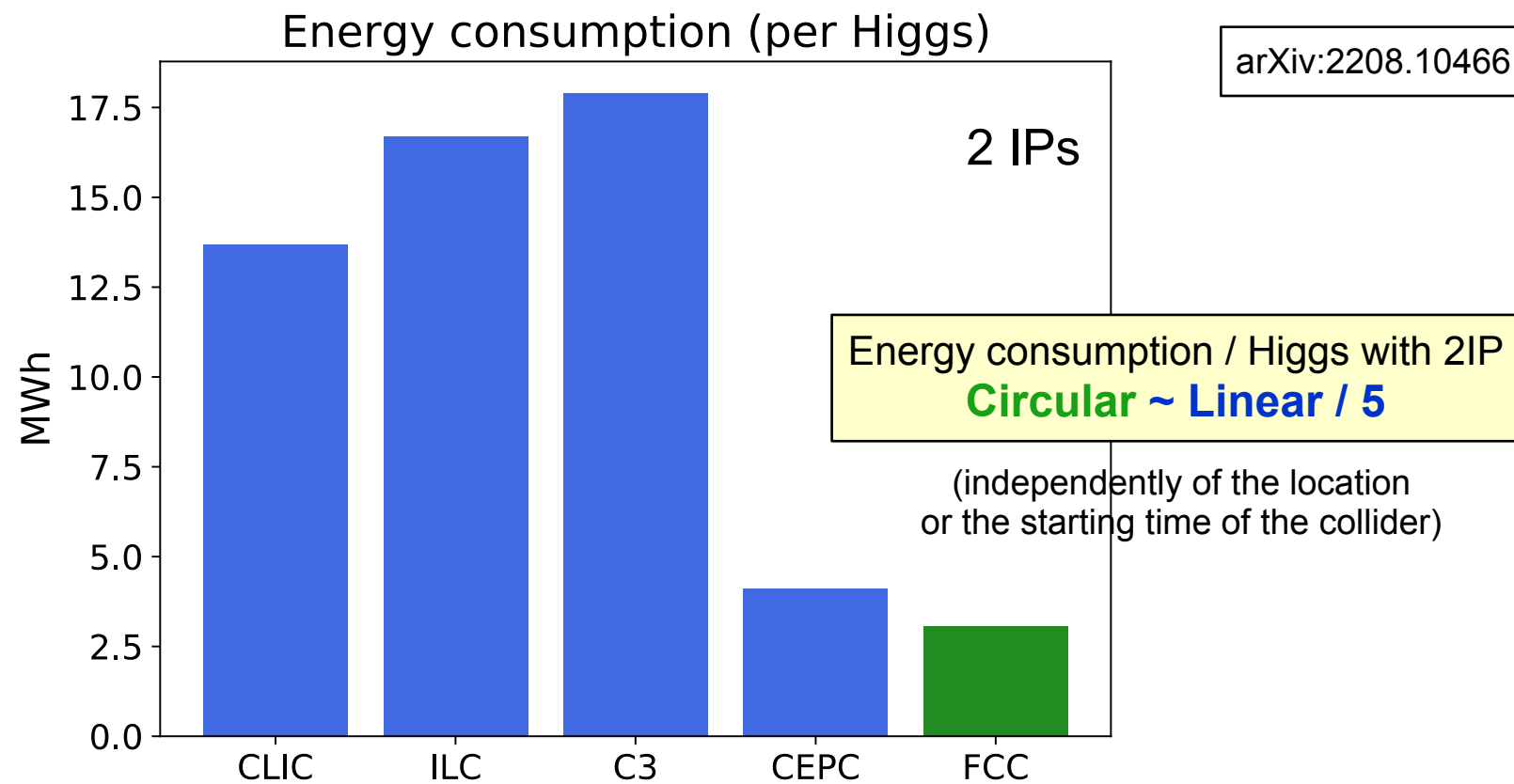
		Z	W	H	TT	
Beam energy (GeV)		45.6	80	120	182.5	
Magnet current		25%	44%	66%	100%	
Power ratio		6%	19%	43%	100%	
PRF EL (MW)	Storage	146	146	146	146	← Ongoing R&D
PRFb EL (MW)	Booster	2	2	2	2	
Pcryo (MW)	all	1,3	12,6	15,8	47,5	← Ongoing R&D
Pcv (MW)	all	33	34	36	40.2	
PEL magnets (MW)	Stroage	6	17	39	89	↓ Potential energy savings
PEL magnets (MW)	Booster	1	3	5	11	
Experiments (MW)	Pt A & G	8	8	8	8	
Data centers (MW)	Pt A & G	4	4	4	4	
General services (MW)		36	36	36	36	
Power during beam operation (MW)		237	262	291	384	

- ◆ At 240 GeV, the instantaneous power of FCC-ee amounts to 291 MW
 - As a comparison, $P(\text{ILC}_{250})=140$ MW, $P(\text{CLIC}_{380})=110$ MW : less power hungry than FCC-ee?
 - Not clear: both produce (2 to 4 times) less Higgs than FCC-ee_{240} , with (3 to 6 times) longer running time

Energy and carbon footprint

- **Our first responsibility (as particle physicists) is to do the maximum of science**
 - ◆ **With the minimal energy consumption and the minimal environmental impact for our planet**
 - Should become one of our top-level decision criteria for design, choice and optimization of a collider

- **All Higgs factories have a “similar” physics outcome (ESU’20 and Snowmass’21)**
 - ◆ **Natural question: what is their energy consumption or carbon footprint for the same physics outcome?**
 - Circular colliders have a much larger instantaneous luminosity and operate several detectors
 - FCC-ee is at CERN, where electricity is already almost carbon-free (and will be even more so in 2048)



Cost of Operation

The total electrical energy consumption over the fourteen years of the FCC-ee research programme is estimated to be around 27 TWh [58], corresponding to an average electricity consumption of 1.9 TWh/year over the entire operation programme, to be compared with the 1.2 TWh/year consumed by CERN today and the expected 1.4 TWh/year for HL-LHC⁹. At the CERN electricity prices from 2014/15, the electricity cost for FCC-ee collider operation would be about 85 MEuro per year. In the HZ running mode, about one million Higgs bosons are expected to be produced in three years, which sets the price of each FCC-ee Higgs boson at 255 Euros. A similar exercise can be done for the first stage of CLIC, expected to consume 0.8 TWh/year over 8 years at 380 GeV to produce about 150,000 Higgs bosons, which sets the price of a CLIC Higgs boson at about 2000 Euros. Finally, with the official ILC operation cost in Japan of 330 MEuro per year [10], its 11.5 to 18.5 years of operation (Section 5), and the 500,000 Higgs bosons produced in total, the price of an ILC Higgs boson is between 7,000 and 12,000 Euros, i.e., between 30 and 50 times more expensive than at FCC-ee. These operation costs are summarized in Table 8.

Table 8: Operation costs of low-energy Higgs factories, expressed in Euros per Higgs boson.

Collider	ILC ₂₅₀	CLIC ₃₈₀	FCC-ee ₂₄₀
Cost (Euros/Higgs)	7,000 to 12,000	2,000	255

Duration of Operation

To get FCC-ee precisions on κ_Z (0.15%), κ_W (0.31%), and κ_H (25%), each collider requires approximately:

FCC-ee (4 expts)	7 years	11 TWh
ILC	~34 years	38 TWh
CLIC	~55 years	33 TWh

To go further and measure κ_H with a 10% precision, the full ILC and CLIC programmes are needed (including phase 3). Alternatively, FCC-hh would reach this precision in 2 to 3 years (and go down to a precision of 3-5% with the full programme). The additional time and energy consumption are as follows:

FCC-hh (100 TeV)	+2 to 3 years	Total 10 years	+6 to 9 (?) TWh	Total 20 TWh
ILC (1 TeV)	+9 years	43 years	+19 TWh	57 TWh
CLIC (1.4 and 3 TeV)	+12 years	67 years	+28 TWh	61 TWh

The integrated FCC project also provide excellent precision to $H\gamma\gamma$, $HZ\gamma$, $H\mu\mu$, and even Hee couplings.

Higgs and EW measurements

Experimental Inputs.

A circular ee Higgs factory starts as a Z/EW factory (**TeraZ**)

A linear ee Higgs factory operating above Z-pole can also preform EW measurements via **Z-radiative** return

A linear ee Higgs factory could also operate on the Z-pole though at lower lumi (**GigaZ**)

	Higgs	aTGC	EWPO	Top EW
FCC-ee	Yes (μ, σ_{ZH}) (Complete with HL-LHC)	Yes (aTGC dom.) <i>Warning</i>	Yes	Yes (365 GeV, Ztt)
ILC	Yes (μ, σ_{ZH}) (Complete with HL-LHC)	Yes (HE limit) <i>Warning</i>	LEP/SLD (Z-pole) + HL-LHC + W (ILC)	Yes (500 GeV, Ztt)
CEPC	Yes (μ, σ_{ZH}) (Complete with HL-LHC)	Yes (aTGC dom) <i>Warning</i>	Yes	No
CLIC	Yes (μ, σ_{ZH})	Yes (Full EFT parameterization)	LEP/SLD (Z-pole) + HL-LHC + W (CLIC)	Yes
HE-LHC	Extrapolated from HL-LHC	N/A → LEP2	LEP/SLD + HL-LHC ($M_W, \sin^2\theta_w$)	-
FCC-hh	Yes ($\mu, BR_i/BR_j$) Used in combination with FCCee/eh	From FCC-ee	From FCC-ee	-
LHeC	Yes (μ)	N/A → LEP2	LEP/SLD + HL-LHC ($M_W, \sin^2\theta_w$)	-
FCC-eh	Yes (μ) Used in combination with FCCee/hh	From FCC-ee	From FCC-ee + Zuu, Zdd	-

Observable	present value	±	error	FCC-ee Stat.	FCC-ee Syst.	Comment and leading error
m_Z (keV)	91186700	±	2200	4	100	From Z line shape scan Beam energy calibration
Γ_Z (keV)	2495200	±	2300	4	25	From Z line shape scan Beam energy calibration
$\sin^2 \theta_W^{\text{eff}} (\times 10^6)$	231480	±	160	2	2.4	From $A_{\text{FB}}^{\mu\mu}$ at Z peak Beam energy calibration
$1/\alpha_{\text{QED}}(m_Z^2)(\times 10^3)$	128952	±	14	3	small	From $A_{\text{FB}}^{\mu\mu}$ off peak QED&EW errors dominate
$R_\ell^Z (\times 10^3)$	20767	±	25	0.06	0.2-1	Ratio of hadrons to leptons Acceptance for leptons
$\alpha_s(m_Z^2) (\times 10^4)$	1196	±	30	0.1	0.4-1.6	From R_ℓ^Z
$\sigma_{\text{had}}^0 (\times 10^3)$ (nb)	41541	±	37	0.1	4	Peak hadronic cross-section Luminosity measurement
$N_\nu (\times 10^3)$	2996	±	7	0.005	1	Z peak cross-sections Luminosity measurement
$R_b (\times 10^6)$	216290	±	660	0.3	< 60	Ratio of $b\bar{b}$ to hadrons Stat. extrapol. from SLD
$A_{\text{FB},0}^b (\times 10^4)$	992	±	16	0.02	1-3	b-quark asymmetry at Z pole From jet charge
$A_{\text{FB}}^{\text{pol},\tau} (\times 10^4)$	1498	±	49	0.15	< 2	τ polarization asymmetry τ decay physics
τ lifetime (fs)	290.3	±	0.5	0.001	0.04	Radial alignment
τ mass (MeV)	1776.86	±	0.12	0.004	0.04	Momentum scale
τ leptonic ($\mu\nu_\mu\nu_\tau$) B.R. (%)	17.38	±	0.04	0.0001	0.003	e/μ /hadron separation
m_W (MeV)	80350	±	15	0.25	0.3	From WW threshold scan Beam energy calibration
Γ_W (MeV)	2085	±	42	1.2	0.3	From WW threshold scan Beam energy calibration
$\alpha_s(m_W^2)(\times 10^4)$	1010	±	270	3	small	From R_ℓ^W
$N_\nu (\times 10^3)$	2920	±	50	0.8	small	Ratio of invis. to leptonic in radiative Z returns
m_{top} (MeV)	172740	±	500	17	small	From $t\bar{t}$ threshold scan QCD errors dominate
Γ_{top} (MeV)	1410	±	190	45	small	From $t\bar{t}$ threshold scan QCD errors dominate
$\lambda_{\text{top}}/\lambda_{\text{top}}^{\text{SM}}$	1.2	±	0.3	0.10	small	From $t\bar{t}$ threshold scan QCD errors dominate
ttZ couplings		±	30%	0.5 – 1.5 %	small	From $\sqrt{s} = 365$ GeV run

EW Precision Measurements at FCC-ee

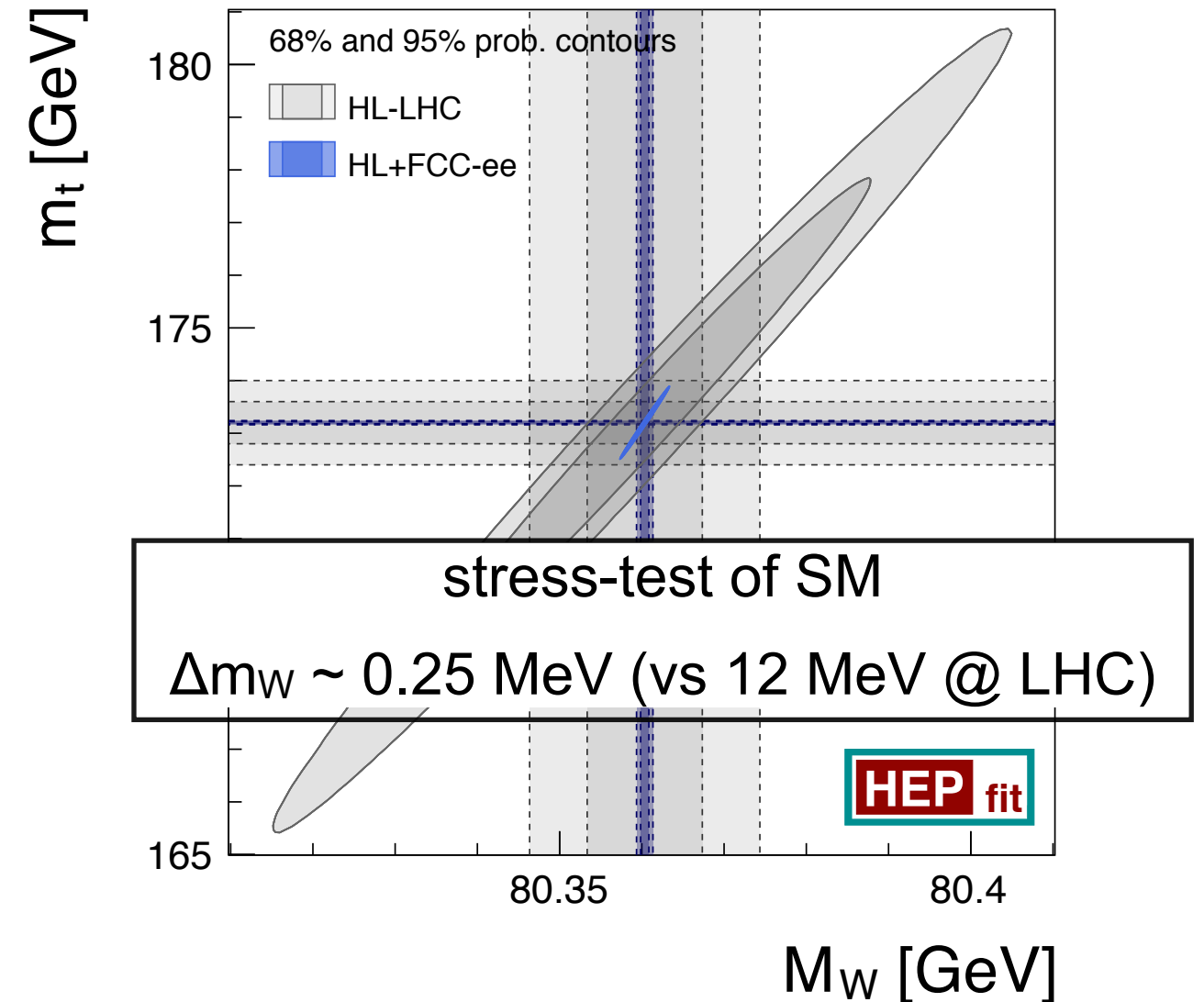
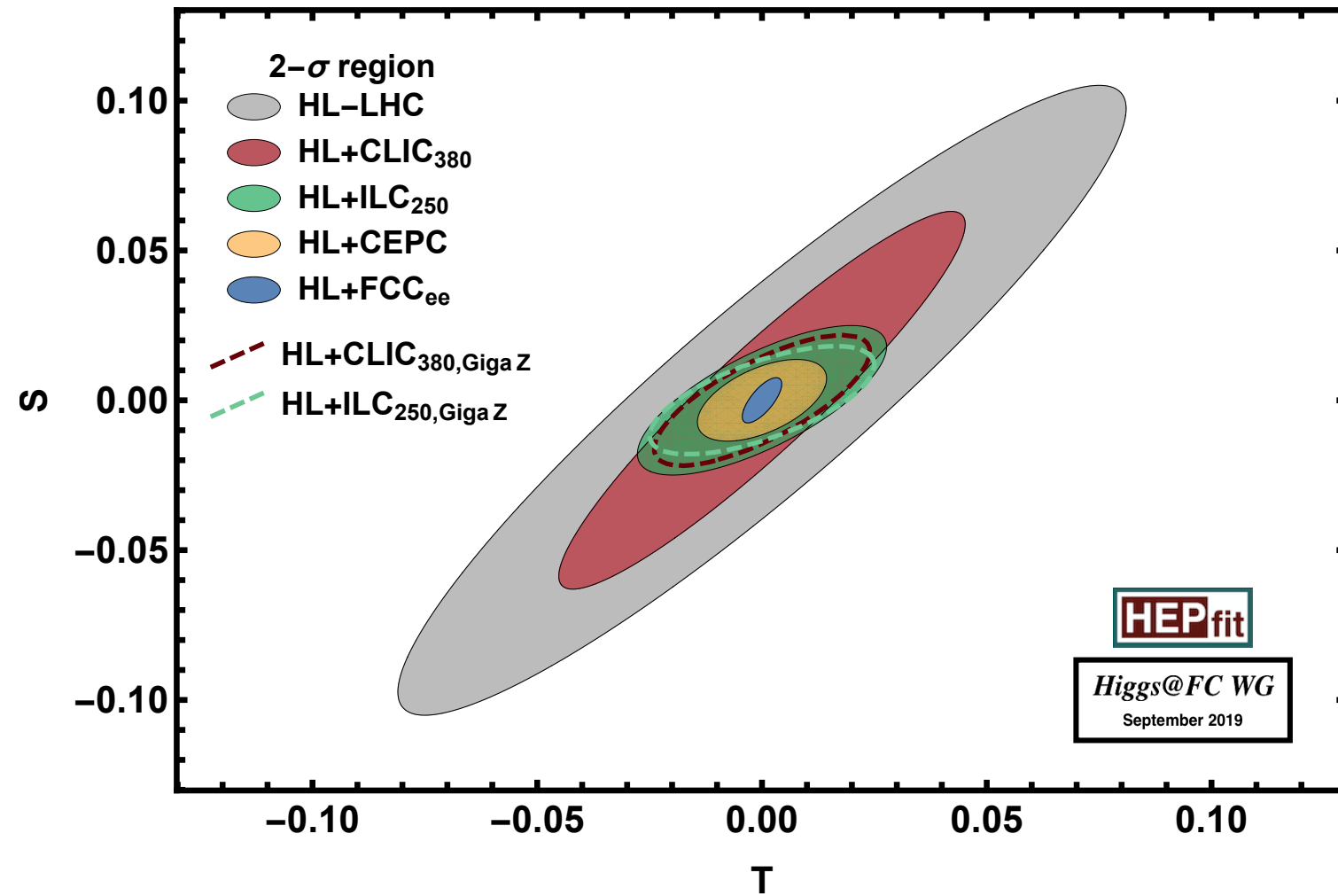
Experimental (statistical and systematic) precision of a selection of measurements accessible at FCC-ee, compared with the present world-average precision. FCC-ee syst. scaled down from LEP estimates. Room for improvement with dedicated studies. Note that syst. go down also with stat. (e.g. beam energy determination from $ee \rightarrow Z/\gamma$ thus goes down with luminosity).

Table from mid-term report

Improvements of EW measurements

Exquisite measurements of m_Z (100 keV), Γ_Z (25 keV), m_W (<500 keV), $\alpha_{\text{QED}}(m_Z)$ ($3 \cdot 10^{-5}$) (all unique to FCC-ee)

w/. stat.+ param. + th-exp syst.



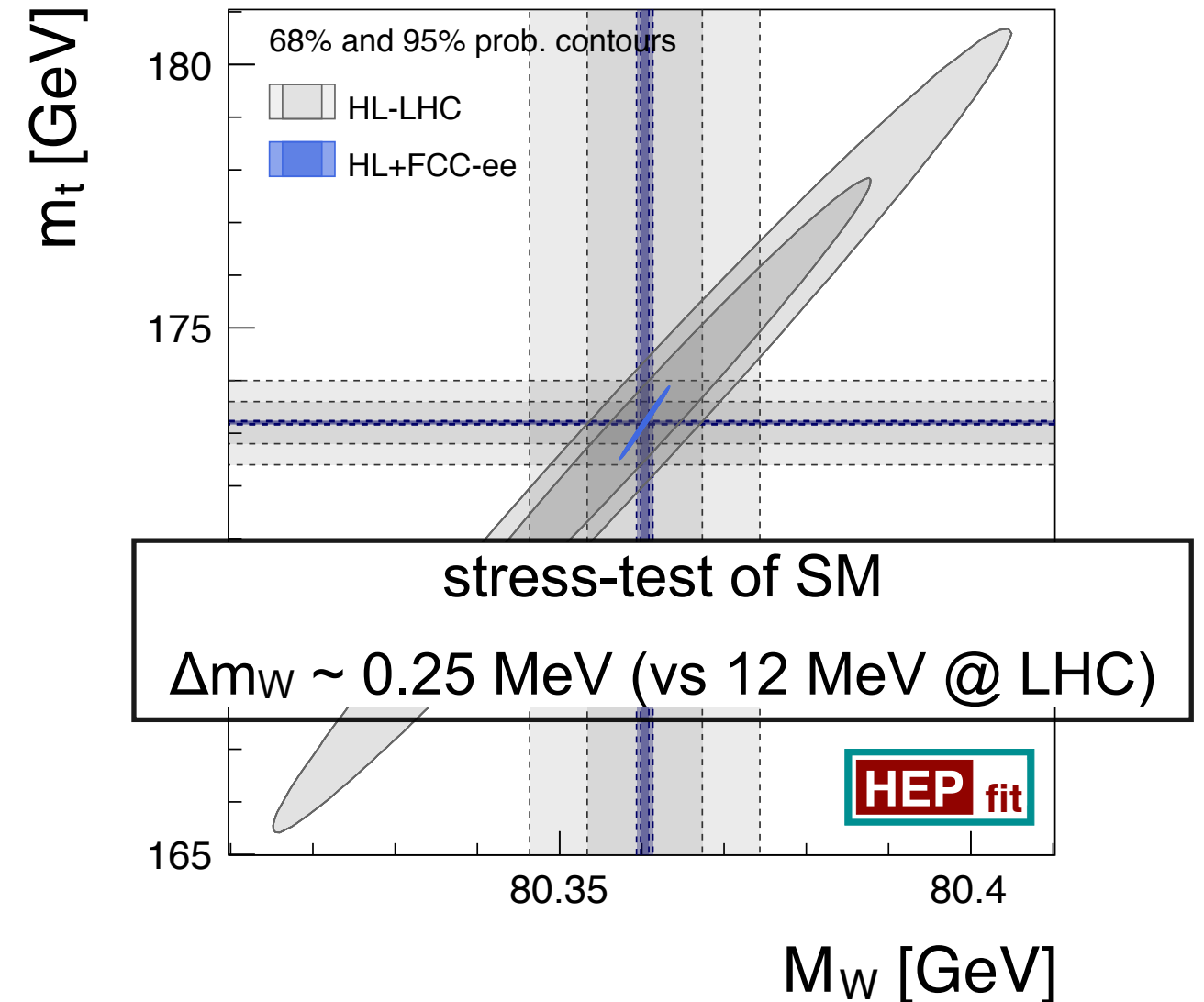
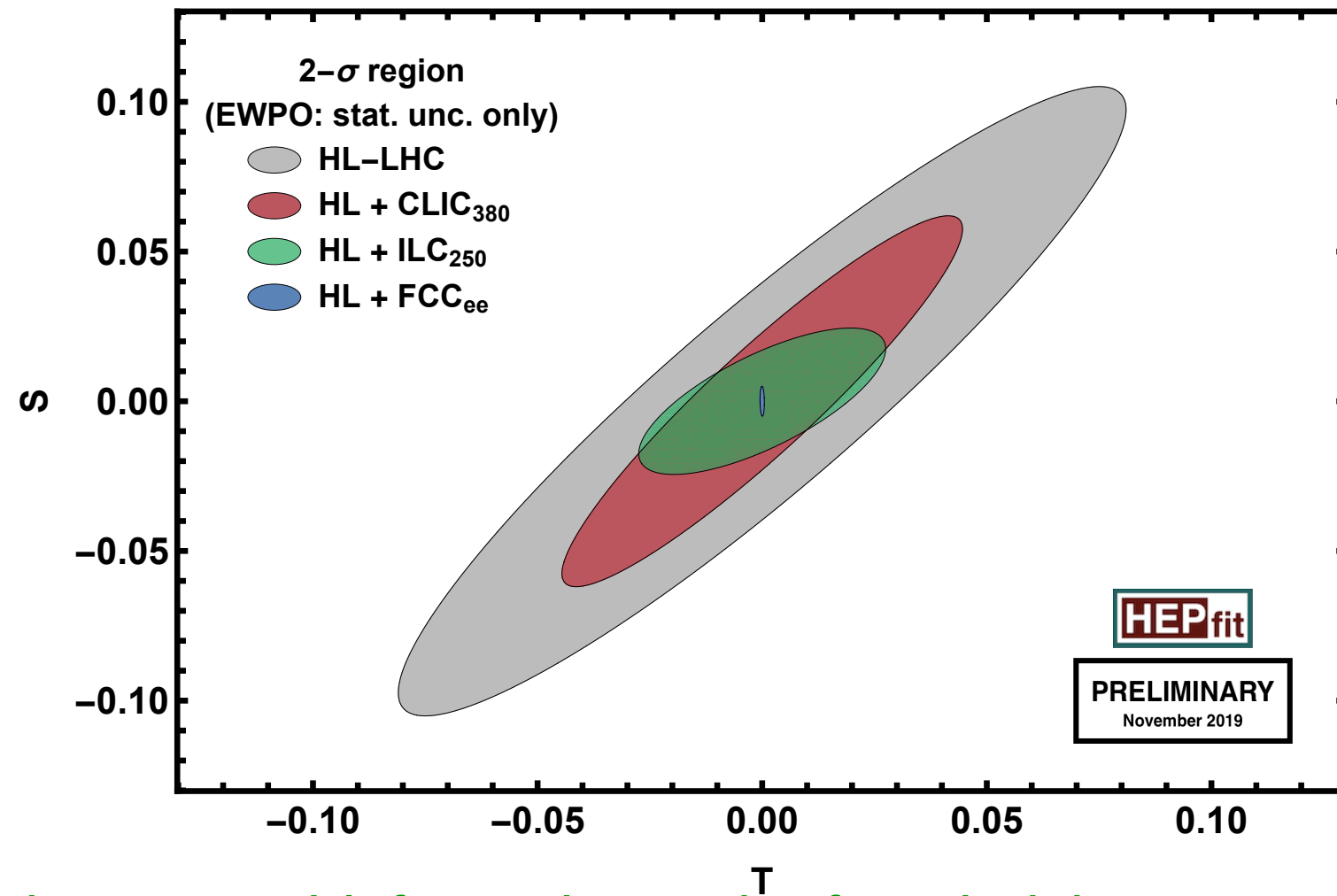
The importance of improved EW measurements is threefold:

- 1) improve mass reach in indirect search for NP ($S \sim 10^{-2} \rightarrow M \sim 70$ TeV)
- 2) reduced parametric uncertainties for other measurements
- 3) reduced degeneracies in a global fit for Higgs couplings

Improvements of EW measurements

Exquisite measurements of m_Z (100 keV), Γ_Z (25 keV), m_W (<500 keV), $\alpha_{\text{QED}}(m_Z)$ ($3 \cdot 10^{-5}$) (all unique to FCC-ee)

w/ stat. and param. only



The importance of improved EW measurements is threefold:

- 1) improve mass reach in indirect search for NP ($S \sim 10^{-2} \rightarrow M \sim 70$ TeV)
- 2) reduced parametric uncertainties for other measurements
- 3) reduced degeneracies in a global fit for Higgs couplings

Systematics vs. Statistics.

PED @ CERN-SPC '2022

- **We often hear that more Z pole statistics is useless, because they are systematics-limited**
 - ◆ This is a passive attitude, which leads to pessimistic expectations and wrong conclusions/planning
 - Experience shows that a careful experimental systematic analysis boils down to a statistical problem
 - If well prepared, theory will go as far as deemed useful : this preparation starts today (and needs SUPPORT)
 - We are working in the spirit of matching systematic errors to expected statistics for all precision measurements
 - ◆ Take the Z lineshape

$\alpha_{\text{QED}}(m_Z)$: Stat. 3×10^{-5}
Obtained at FCC-ee from off-peak asymmetries (87.9 & 94.3 GeV): for the first time, it is a direct measurement of this quantity (game changer)

- Enters as a limiting parametric uncertainties in the new physics interpretation many past and future measurements.
- Is statistics limited and will directly benefit from more luminosity
- No useful impact on $\alpha_{\text{QED}}(m_Z)$ with five times less luminosity

FCC-ee
special

$\sin^2\theta_W^{\text{eff}}$ and Γ_Z (also m_W vs m_Z) : Stat. 2×10^{-6} and 4 keV
Error dominated by point-to-point energy uncertainties.
Based on in-situ comparisons between \sqrt{s} (e.g. with muon pairs), with measurements made every few minutes (100's times per day)
Boils down to

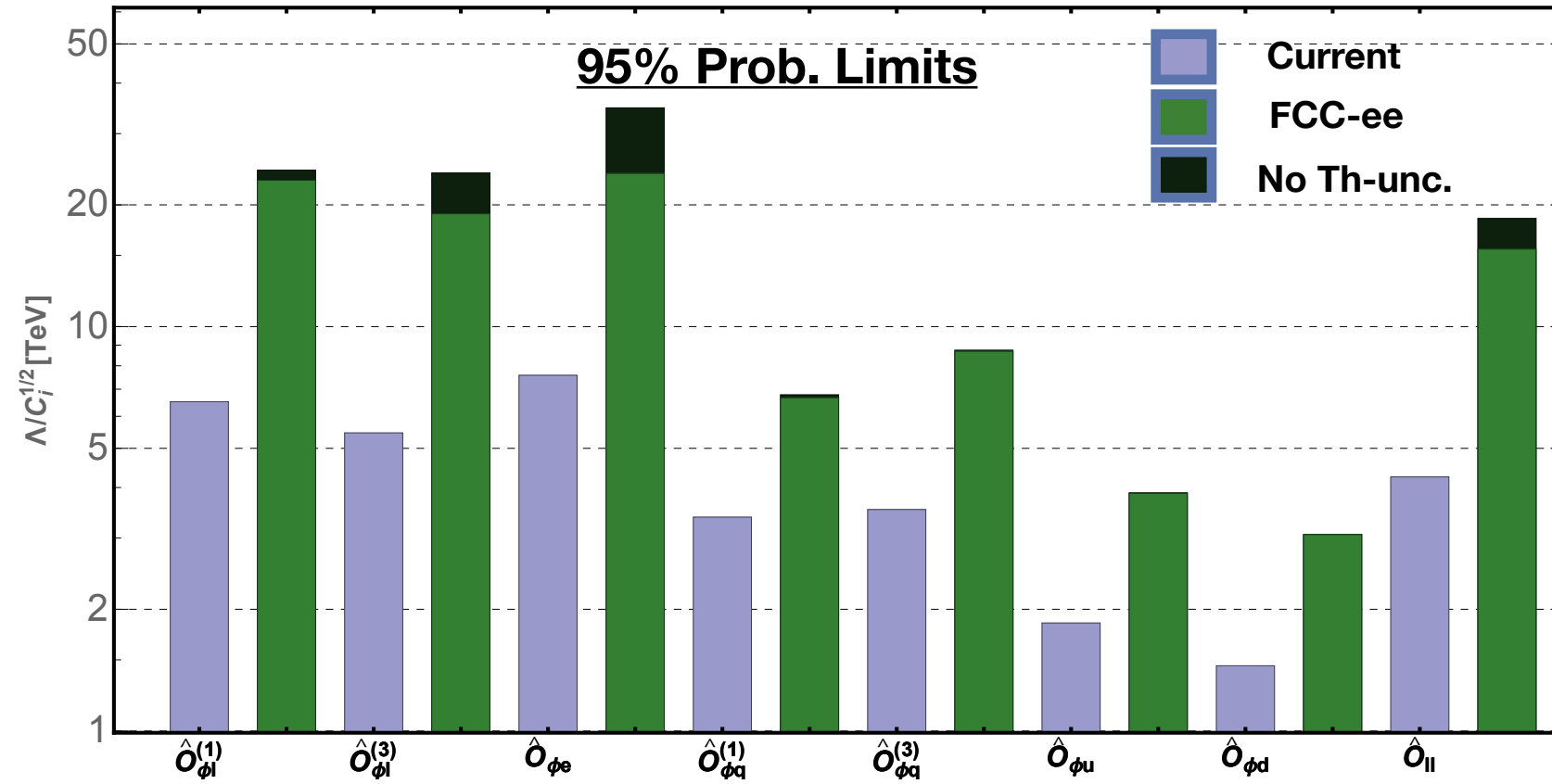
- statistics (the more data the better, scales down as $1/\sqrt{L}$)
- detector systematics (uncorrelated between experiments, scales down a $1/\sqrt{N_{\text{experiments}}}$)

Z (and W) mass: Stat. 4 keV (250 keV)
Error dominated by \sqrt{s} determination with resonant depolarization.
As more understanding is gained, progress are made at a constant pace, and this error approaches regularly the statistical limit

- ◆ Most of the work is (will be) on systematics
 - But huge statistics will turn into better precision
→ A real chance for discovery

Impact of TH uncertainties.

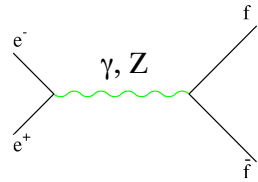
J. de Blas, FCC CDR overview '19



	Current		FCCee		
	Exp.	SM	Exp.	SM (par.)	SM (th.)
δM_W [MeV]	± 15	± 8	± 1	$\pm 0.6/\pm 1$	± 1
$\delta \Gamma_Z$ [MeV]	± 2.3	± 0.73	± 0.1	± 0.1	± 0.2
$\delta \mathcal{A}_\ell$ [$\times 10^{-5}$]	± 210	± 93	± 2.1	$\pm 8/\pm 14$	± 11.8
δR_b^0 [$\times 10^{-5}$]	± 66	± 3	± 6	± 0.3	± 5

Some EW measurements @ Tera

measure $\sigma(e^+e^- \rightarrow \mu^+\mu^-)$ and $A_{FB}^{\mu\mu}$ at (a) judicious \sqrt{s}



The γ exchange term is proportional to $\alpha_{QED}^2(\sqrt{s})$
 The Z exchange term is proportional to G_F^2 , hence independent of α_{QED}
 The γZ interference is proportional to $\alpha_{QED}(\sqrt{s}) \times G_F$

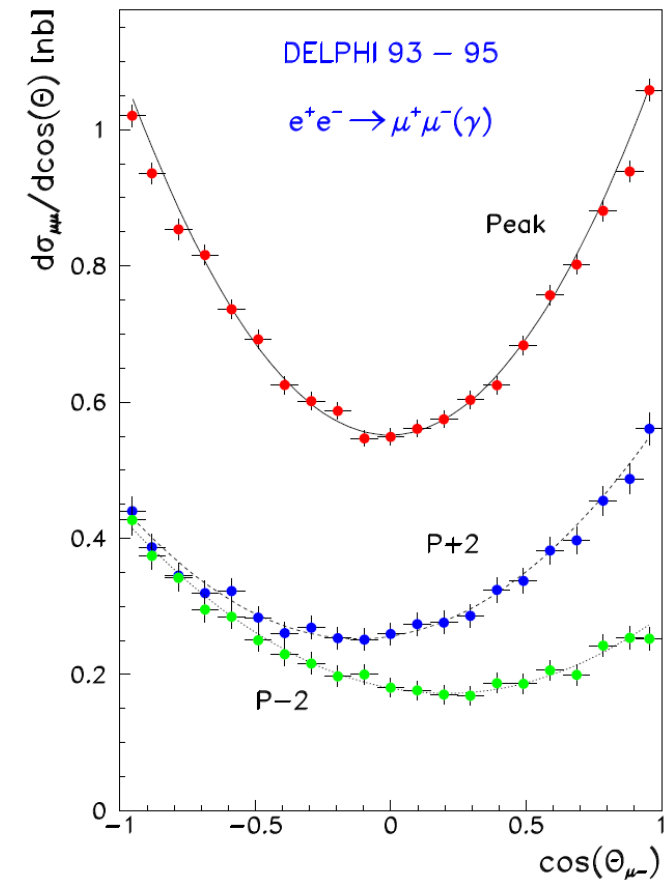
strongly depends on \sqrt{s}

direct measurement of $\alpha_{QED}(s)$ at $\sqrt{s} \neq m_Z$
 measure $\sin^2\theta_W$ to high precision

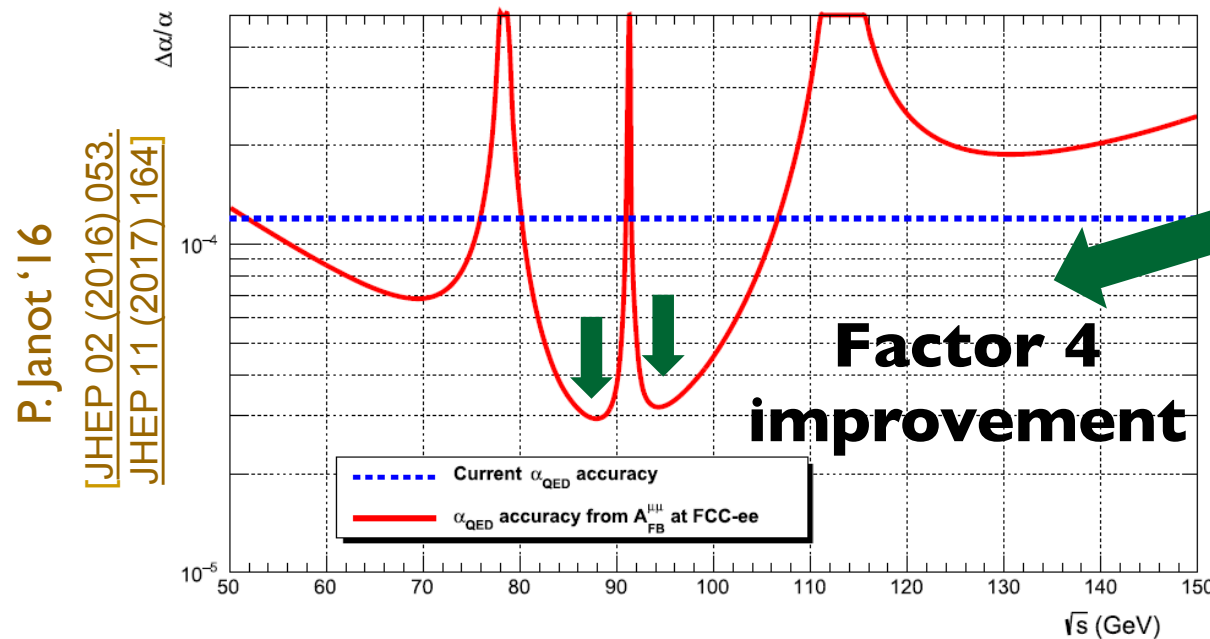
Excellent experimental control of off-peak di-muon asymmetry motivates campaign to collect 50-80 ab^{-1} off peak to gain highest sensitivity to Z- γ interference

$$A_{FB}^{\mu\mu}(s) \simeq \frac{3}{4} \mathcal{A}_e \mathcal{A}_\mu \times \left[1 + \frac{8\pi\sqrt{2}\alpha_{QED}(s)}{m_Z^2 G_F (1 - 4\sin^2\theta_W^{eff})^2} \frac{s - m_Z^2}{2s} \right]$$

Allows for clean determination of $\alpha_{QED}(m_Z^2)$, which is a *critical* input for m_W closure tests (see later).



relative α_{QED} uncertainty with 80 ab^{-1}



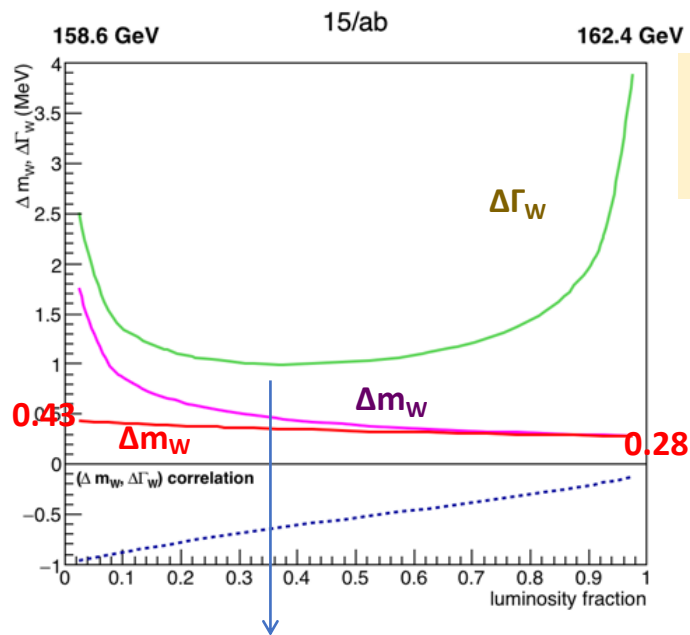
This dependence, & location of half-integer spin tunes, guides the choice of off-peak energies: 87.8 & 93.9 GeV.

- Measure $\alpha_{QED}(m_Z^2)$ to 3×10^{-5} rel. precision (currently 1.1×10^{-4})
- Stat. dominated; syst. uncertainties $< 10^{-5}$ (dominated by \sqrt{s} calib)
- Theoretical uncertainties $\sim 10^{-4}$, higher order calcs needed

M_w.

- Two independent W mass and width measurements @ FCCee :
 - The m_W and Γ_W determinations from the WW threshold cross section lineshape, with 12/ab at $E_{CM} \simeq 157.5-162.5$ GeV** $\Delta m_W=0.4$ MeV $\Delta \Gamma_W=1$ MeV
 - Other measurements of m_W and Γ_W from the decay products kinematics at $E_{CM} \simeq 162.5-240-365$ GeV** $\Delta m_W, \Delta \Gamma_W= 2-5$ MeV ?

Scans of possible $E_1 E_2$ data taking energies and luminosity fractions f (at the E_2 point)



A - minimum of $\Delta \Gamma_W=0.91$ MeV with $\Delta m_W=0.55$ MeV
 taking data at $E_1=156.6$ GeV $E_2=162.4$ GeV $f=0.25$
 yields $\Delta m_W=0.47$ MeV (as single par)

B- minimum of $\Delta m_W=0.28$ MeV $\Delta \Gamma_W=3.3$ MeV with
 $E_1=155.5$ GeV $E_2=162.4$ GeV $f=0.95$
 yields $\Delta m_W=0.28$ MeV (as single par)

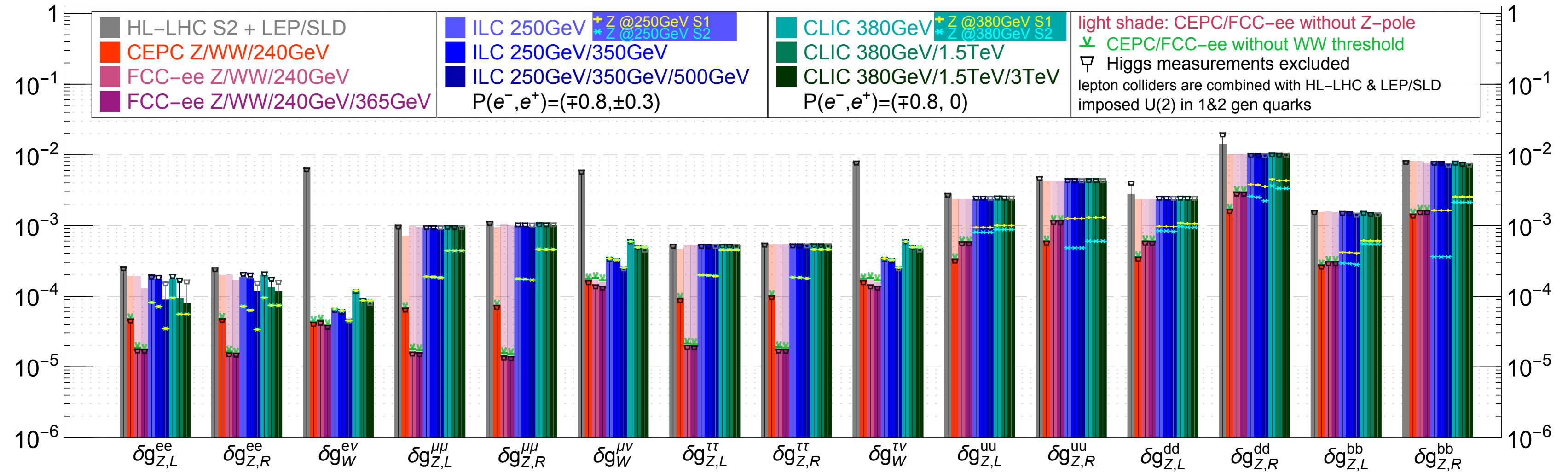
C- minimum of $\Delta \Gamma_W=0.96$ MeV + $\Delta m_W=0.41$ MeV with
 $E_1=157.5$ GeV $E_2=162.4$ GeV $f=0.45$
 yields and $\Delta m_W=0.37$ MeV (as single par)

$\Delta m_W=0.45$ MeV, $\Delta \Gamma_W=1$ MeV ($r=-0.6$)
 $\Delta m_W=0.35$ MeV

$\Delta m_W, \Delta \Gamma_W$: error on W mass and width from fitting both
 Δm_W : error on W mass from fitting only m_W

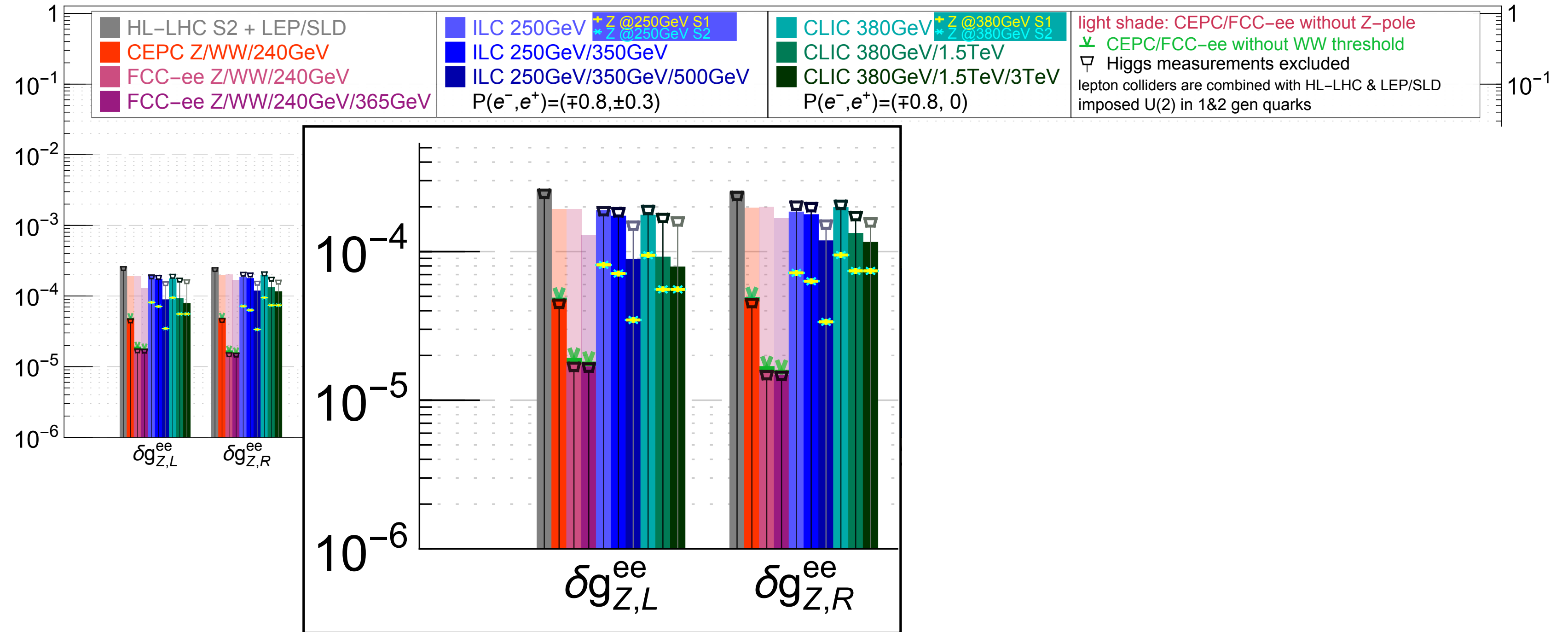
Sensitivity on EW couplings.

J. De Blas, G. Durieux, C. Grojean, J. Gu, A. Paul 1907.04311



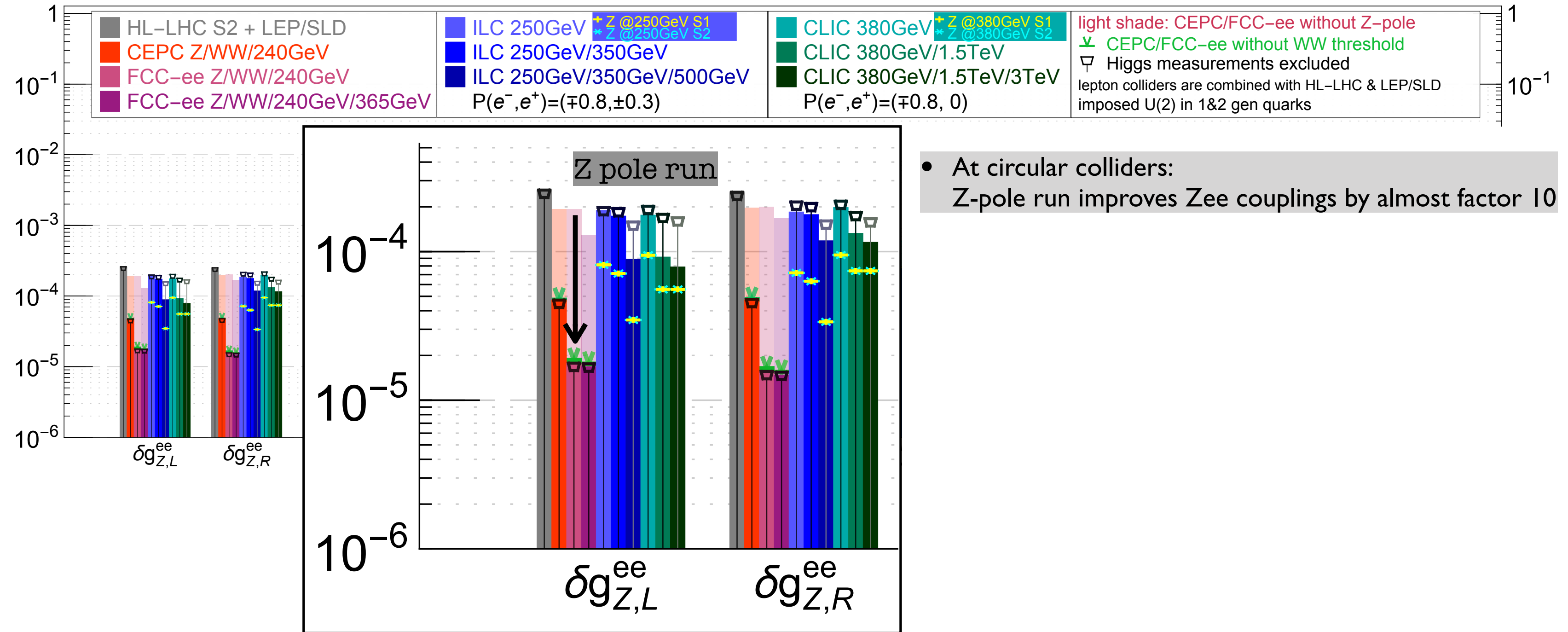
Sensitivity on EW couplings.

J. De Blas, G. Durieux, C. Grojean, J. Gu, A. Paul 1907.04311



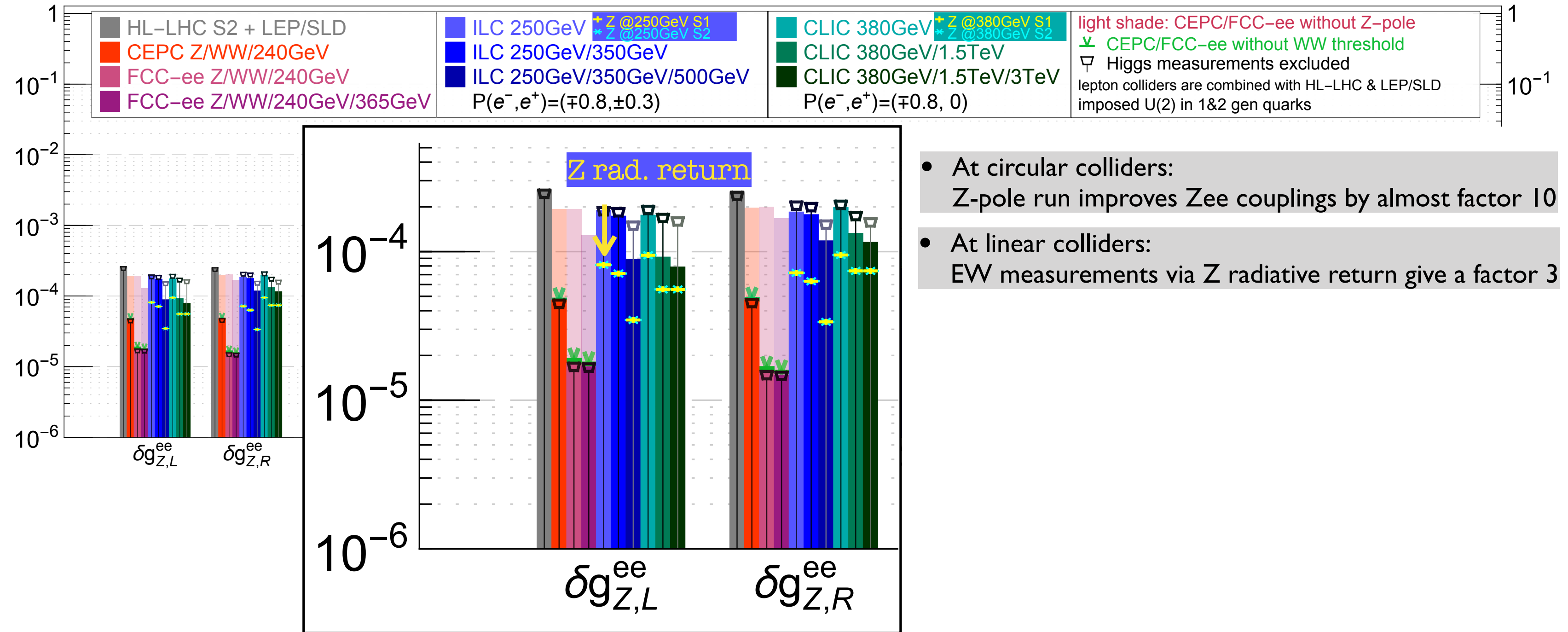
Sensitivity on EW couplings.

J. De Blas, G. Durieux, C. Grojean, J. Gu, A. Paul 1907.04311



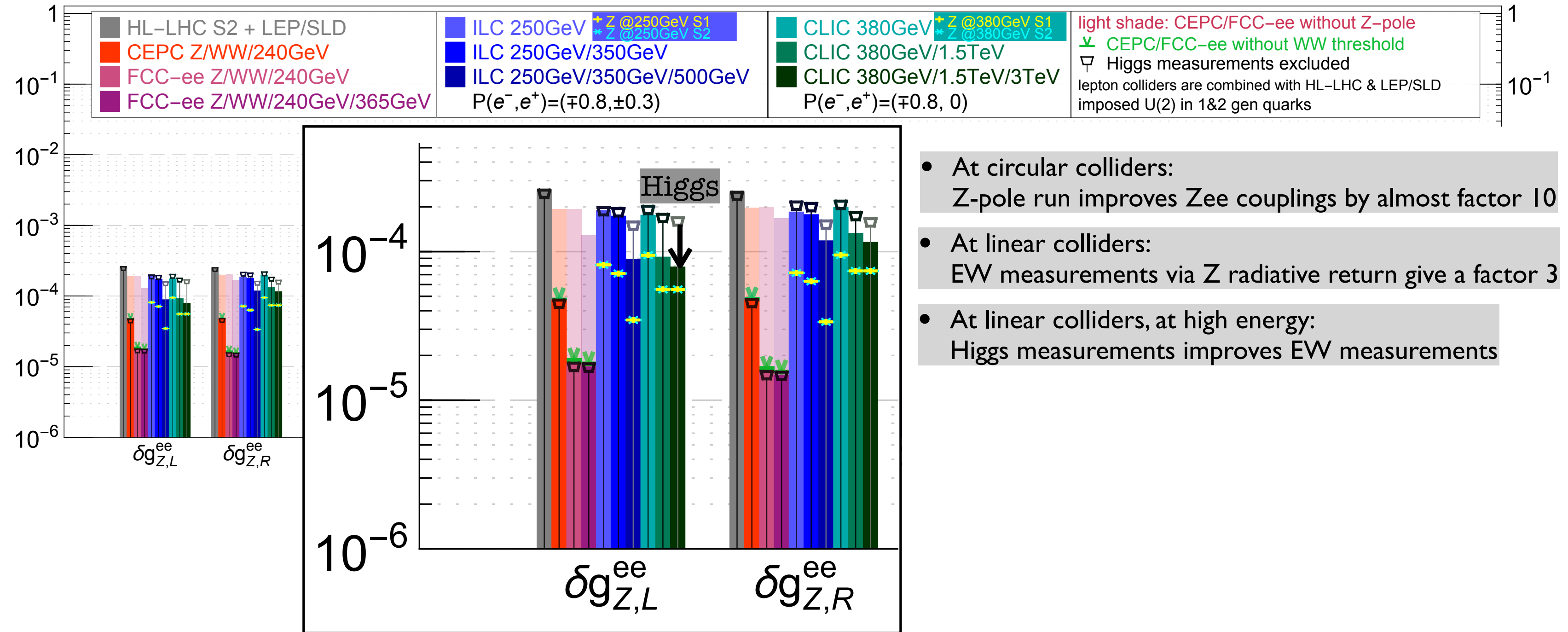
Sensitivity on EW couplings.

J. De Blas, G. Durieux, C. Grojean, J. Gu, A. Paul 1907.04311



Sensitivity on EW couplings.

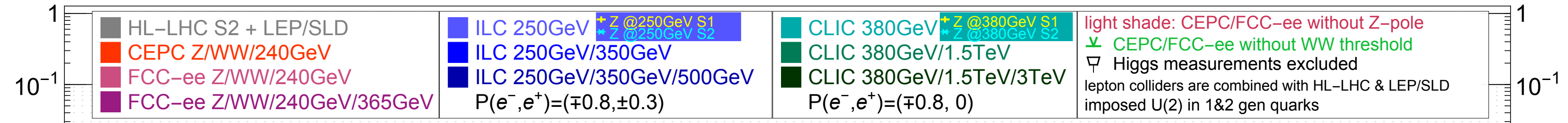
J. De Blas, G. Durieux, C. Grojean, J. Gu, A. Paul 1907.04311



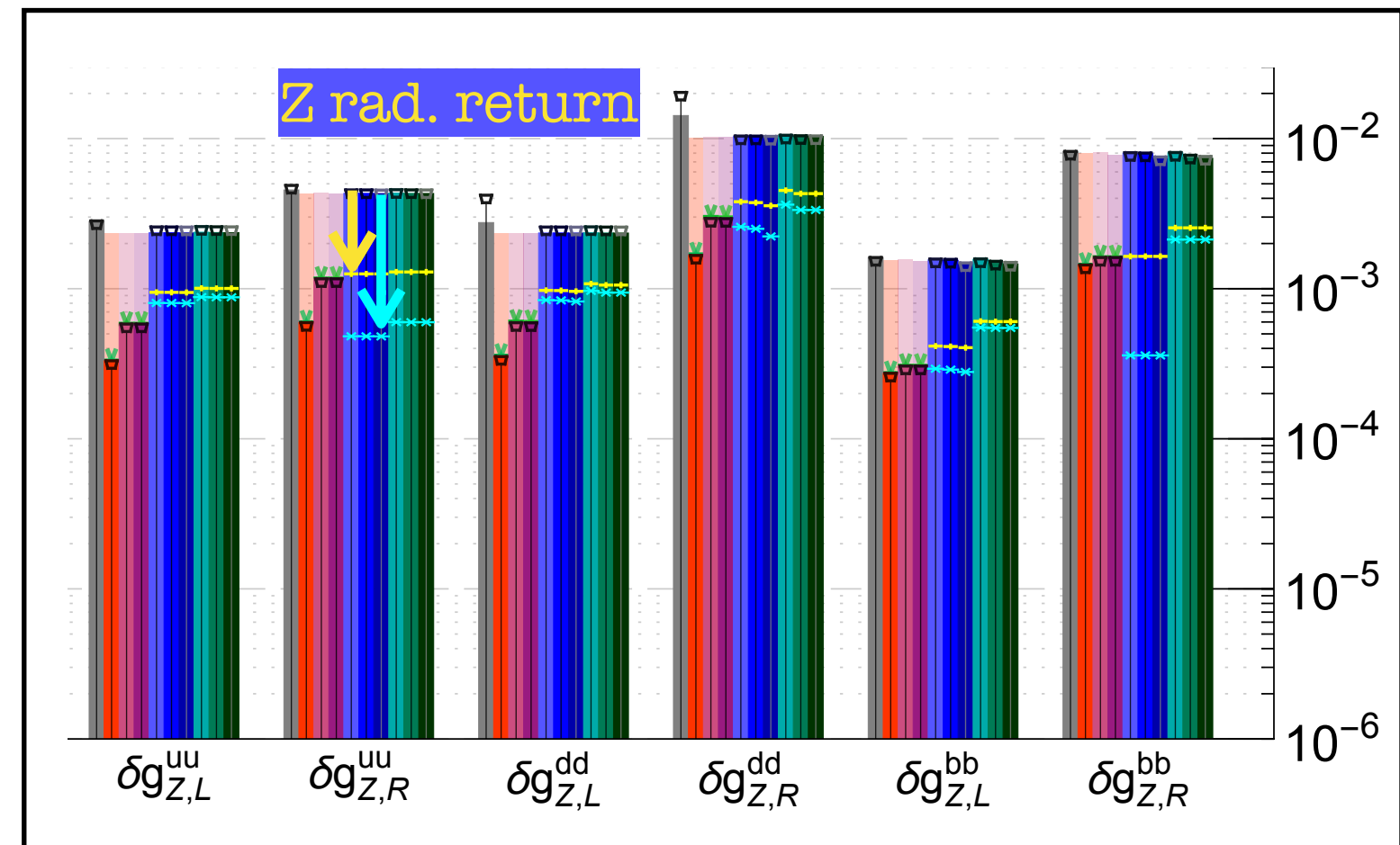
- At circular colliders:
Z-pole run improves Zee couplings by almost factor 10
- At linear colliders:
EW measurements via Z radiative return give a factor 3
- At linear colliders, at high energy:
Higgs measurements improves EW measurements

Sensitivity on EW couplings.

J. De Blas, G. Durieux, C. Grojean, J. Gu, A. Paul 1907.04311

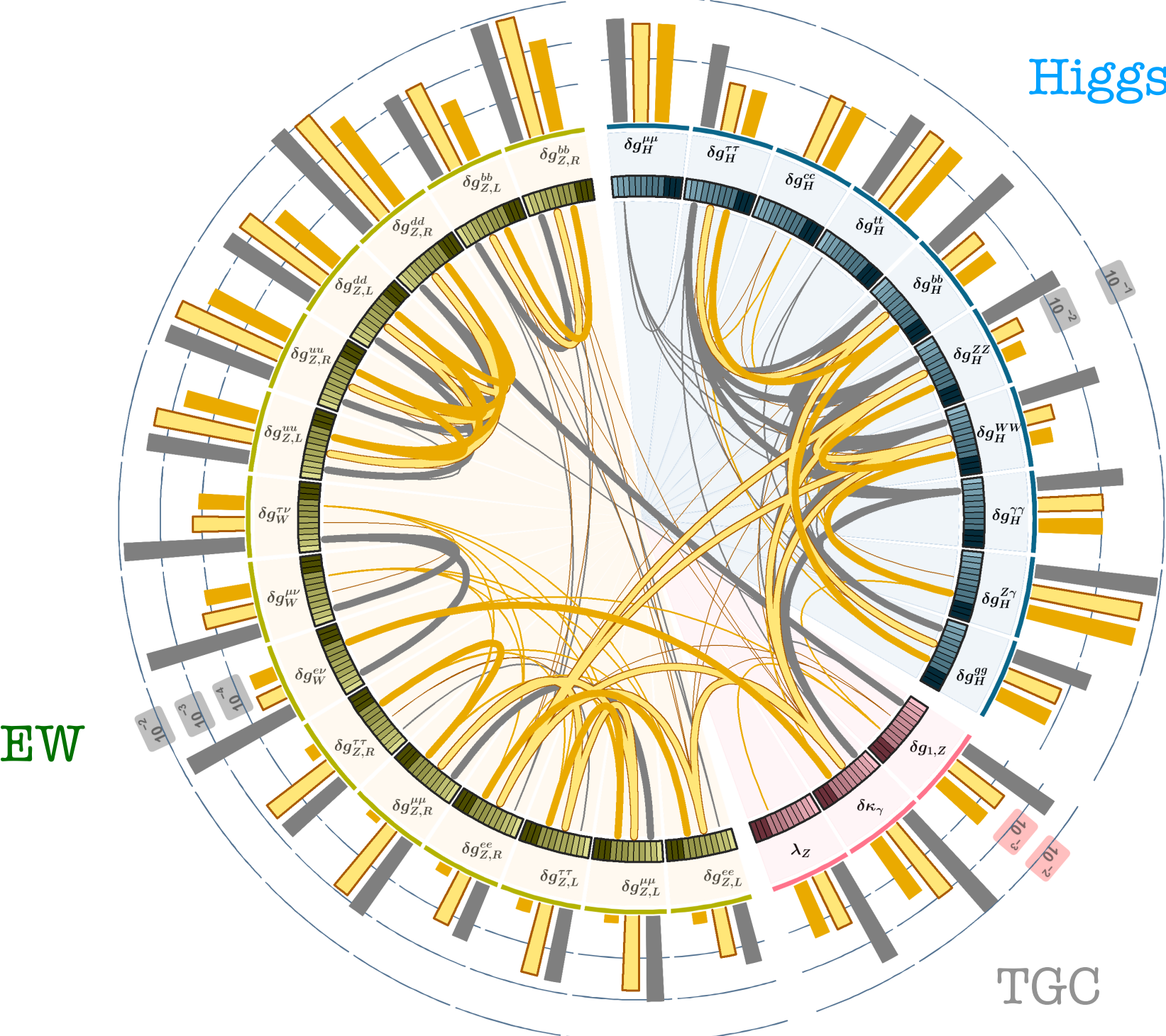


- At linear colliders, at high energy: EW measurements via Z-radiative return has a large impact on $Zq\bar{q}$ couplings
- Improvements depend a lot on hypothesis on systematic uncertainties
 - Yellow: LEP/SLD systematics / 2
 - Blue: small EXP and TH systematics



Why Z-pole for Higgs?

J. De Blas et al. 1907.04311



EW

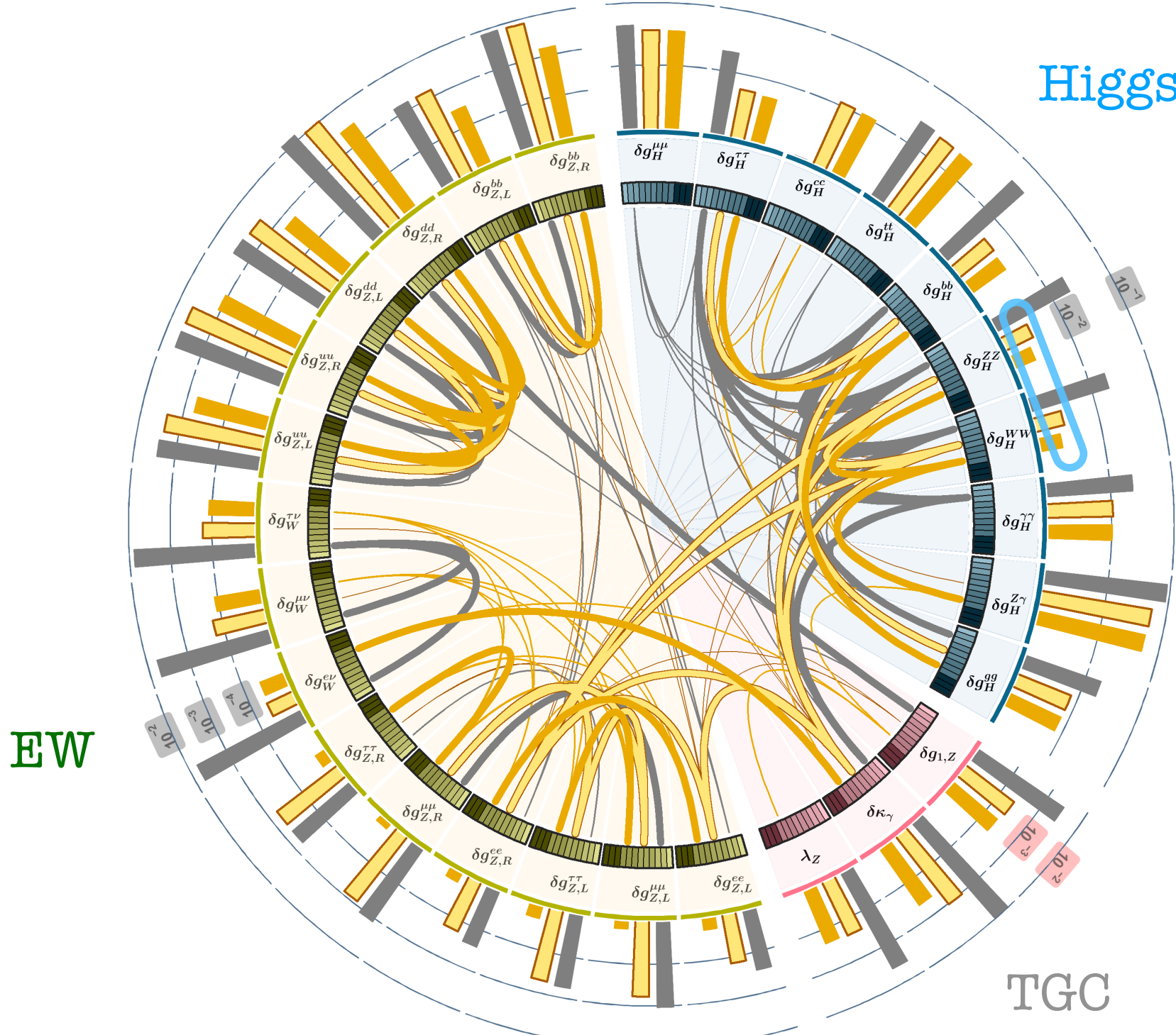
Higgs

TGC

Why Z-pole for Higgs?

J. De Blas et al. 1907.04311

With Z-pole measurements, Higgs coupling determination improves by up to 50%



EW

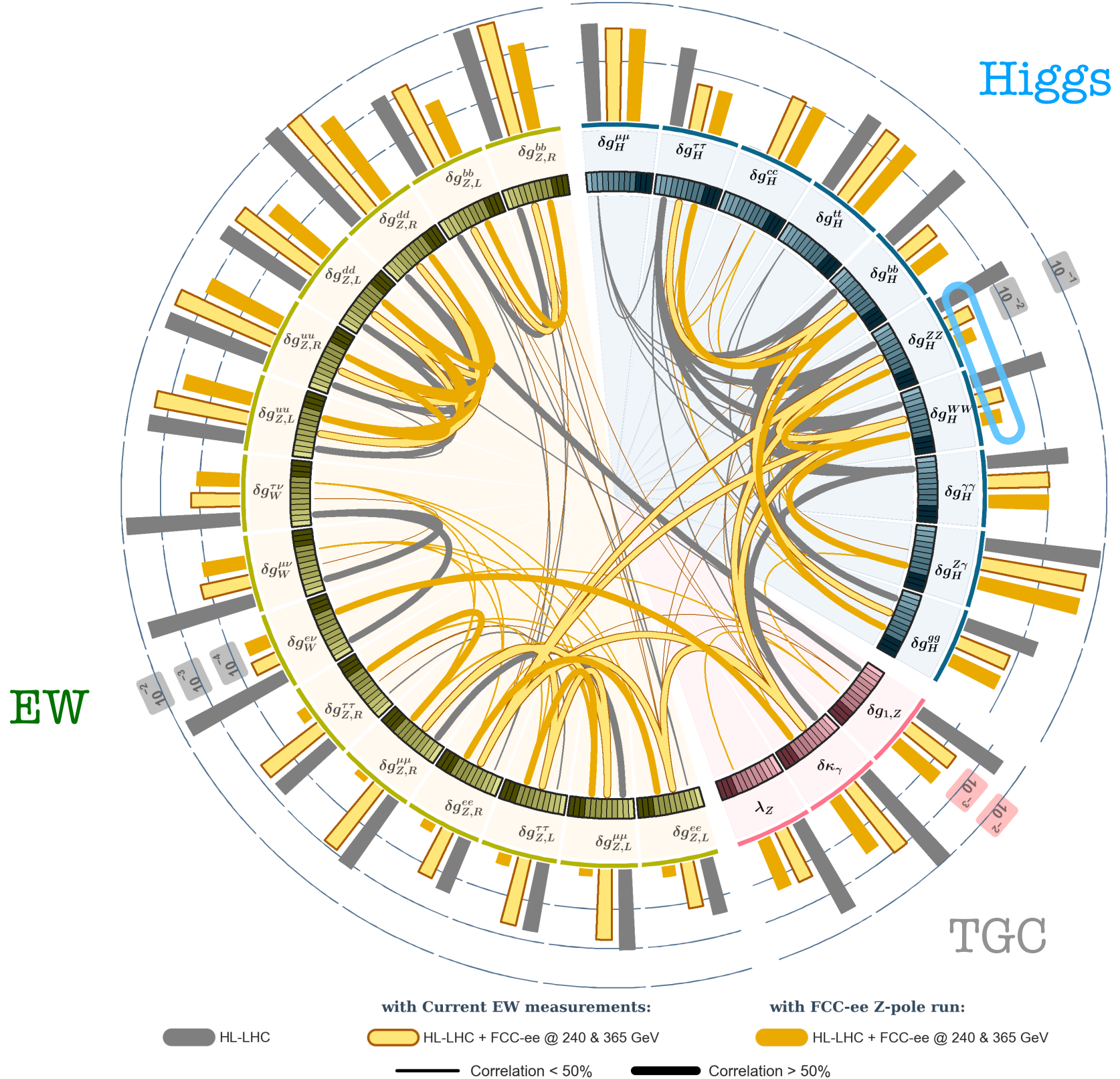
Higgs

TGC

with Current EW measurements: with FCC-ee Z-pole run:
 HL-LHC HL-LHC + FCC-ee @ 240 & 365 GeV HL-LHC + FCC-ee @ 240 & 365 GeV
 Correlation < 50% Correlation > 50%

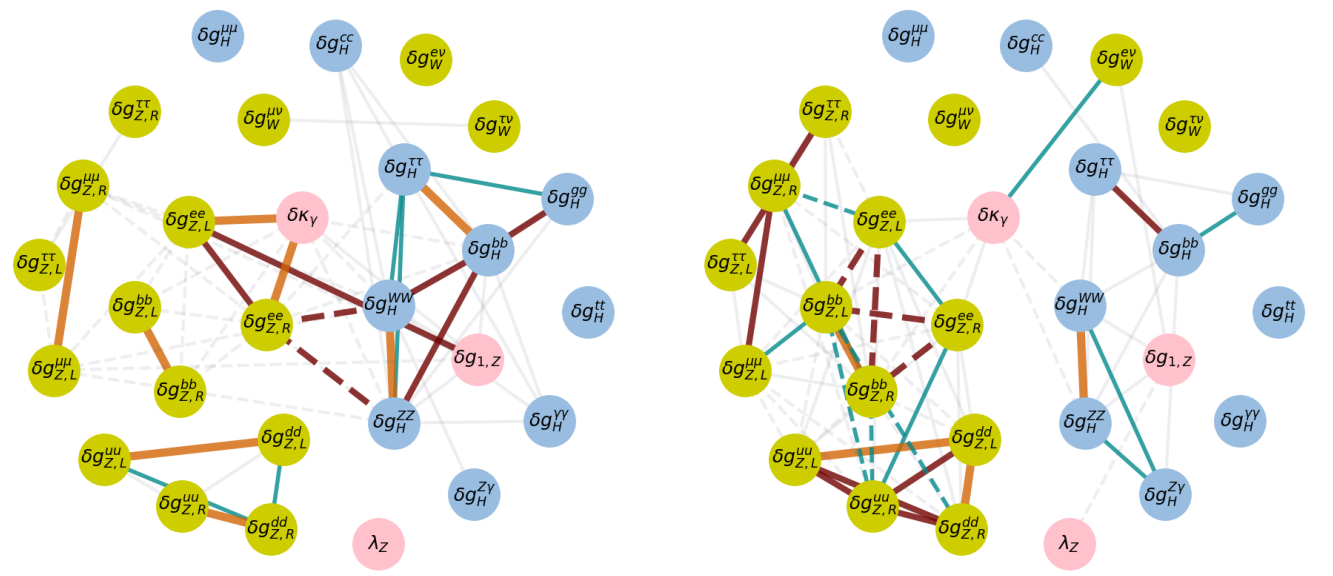
Why Z-pole for Higgs?

J. De Blas et al. 1907.04311



With Z-pole measurements, Higgs coupling determination improves by up to 50%

Z-pole run at circular colliders decorrelates EW and Higgs sectors from each other

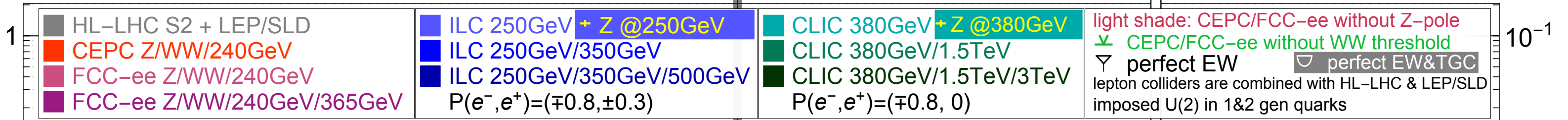


● Higgs
● aTGC
● EW

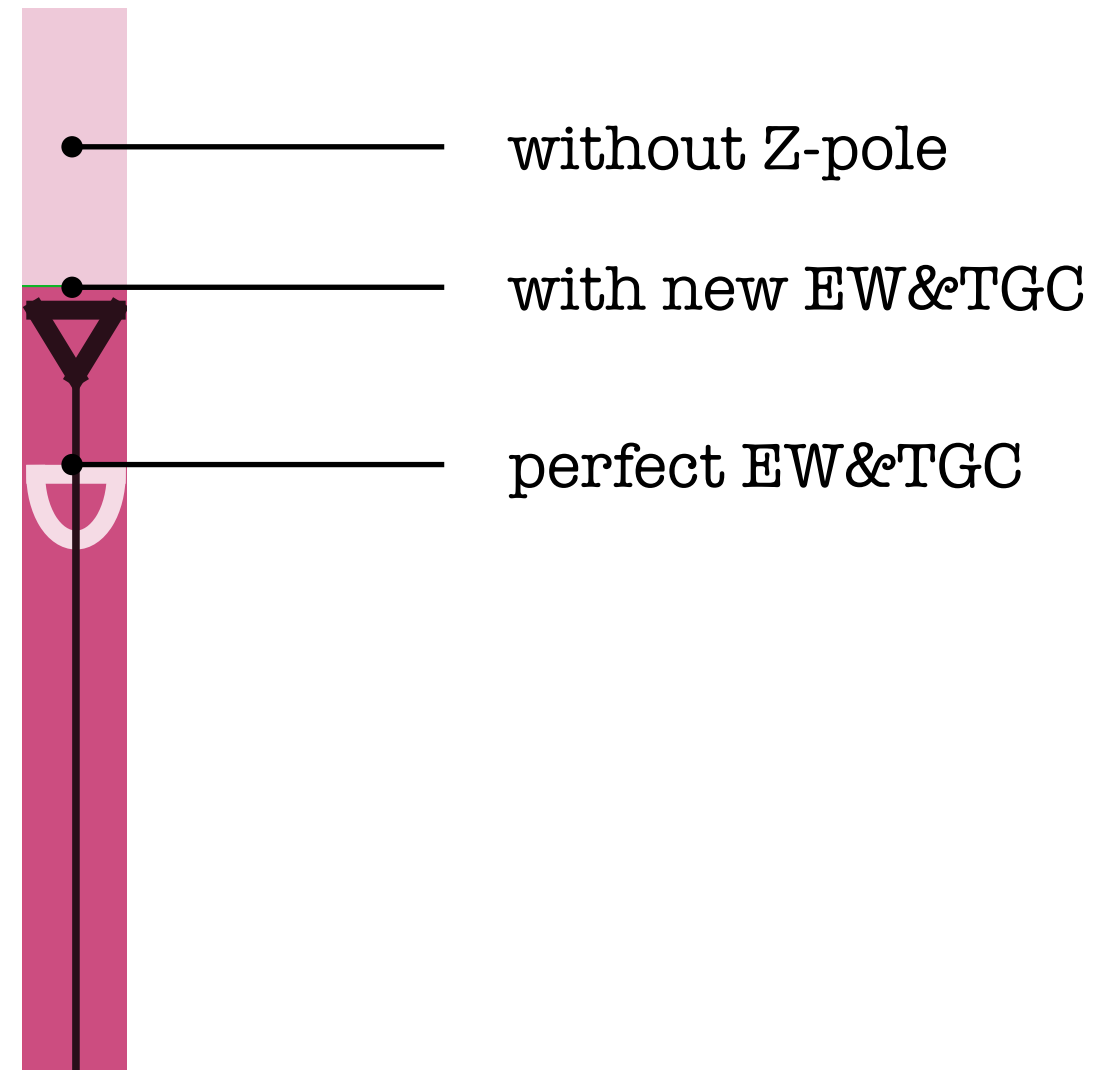
Impact of Z-pole on Higgs.

J. De Blas et al. 1907.04311

Comparing 3 EW scenarios: LEP/SLD, actual EW measurements, perfect EW measurements



Higgs couplings

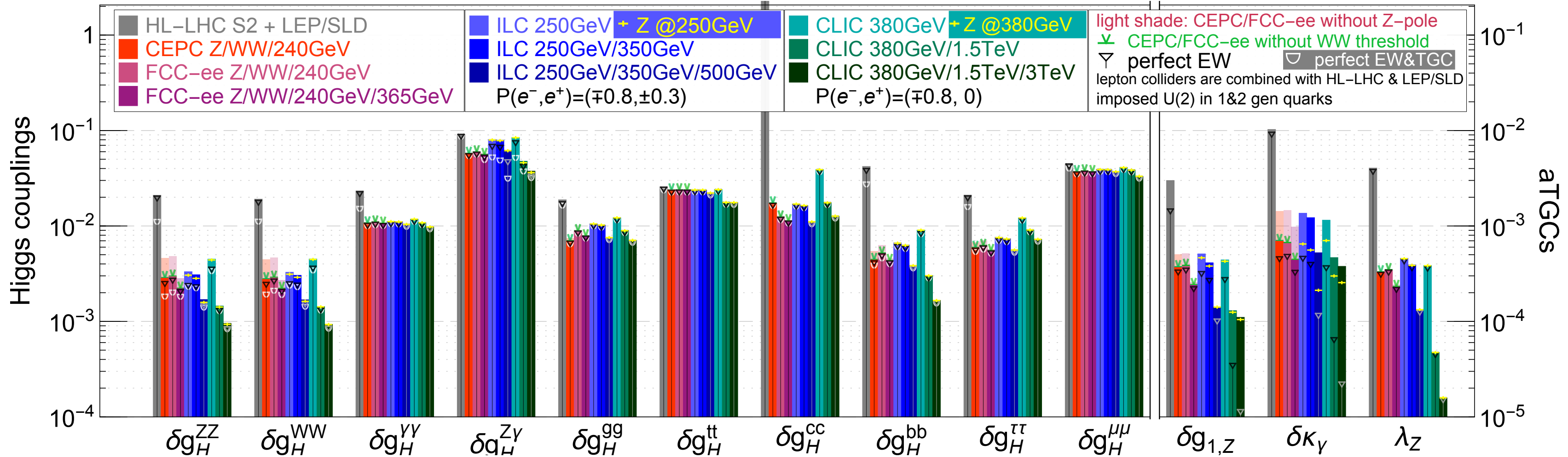


FCC-ee Z/WW/240GeV

Impact of Z-pole on Higgs.

J. De Blas et al. 1907.04311

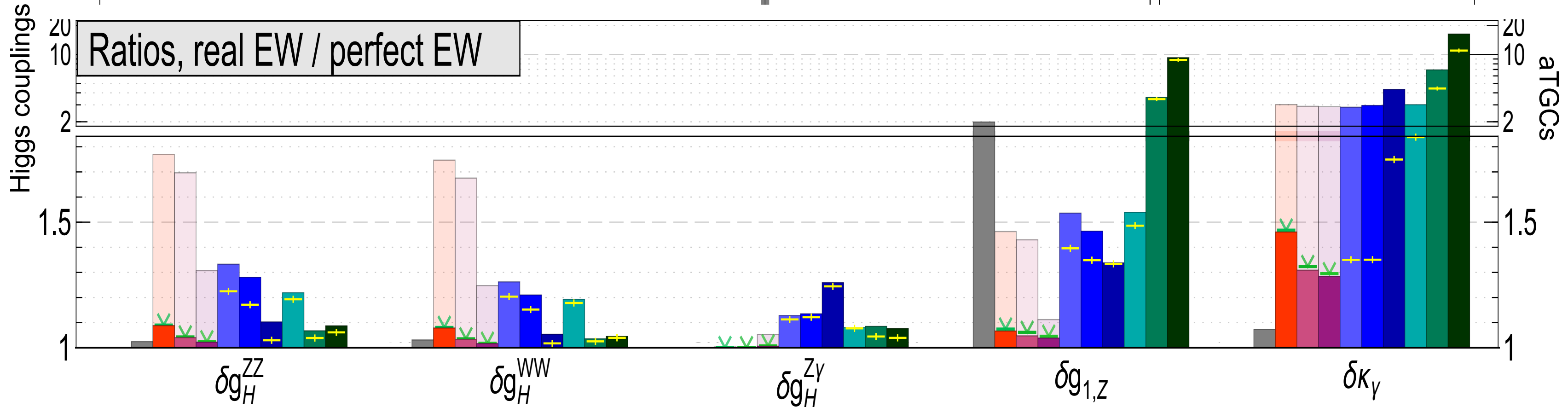
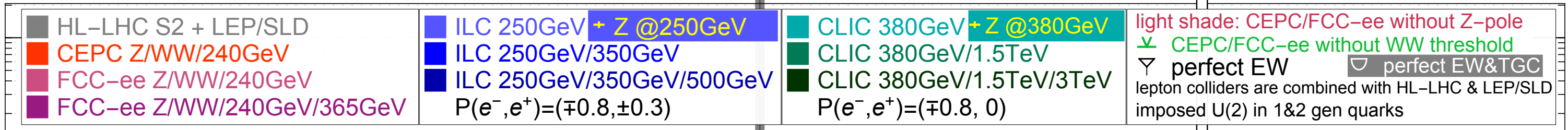
Comparing 3 EW scenarios: LEP/SLD, actual EW measurements, perfect EW measurements



Impact of Z-pole on Higgs.

J. De Blas et al. 1907.04311

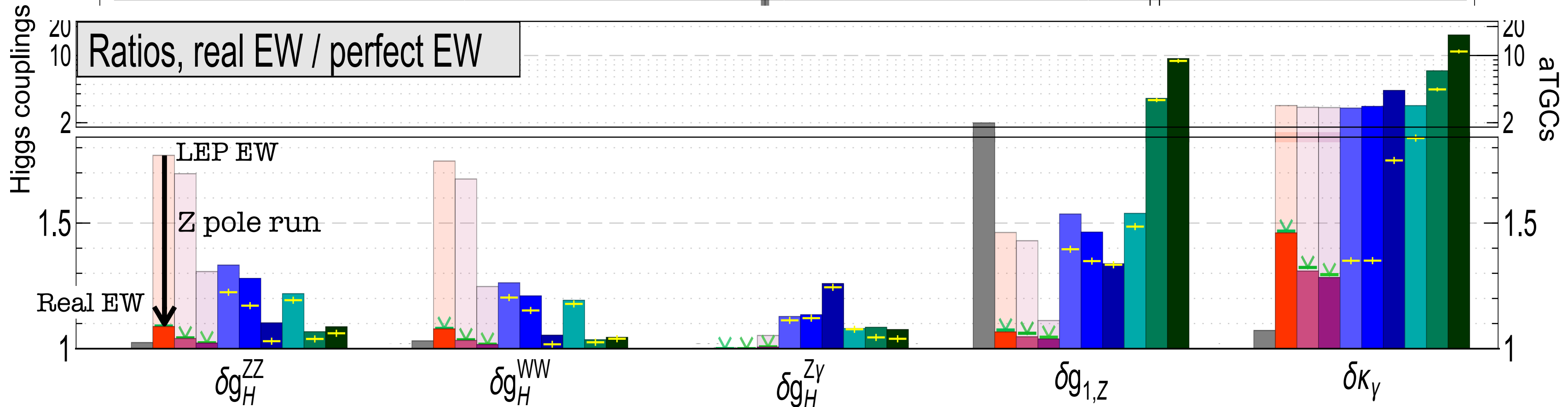
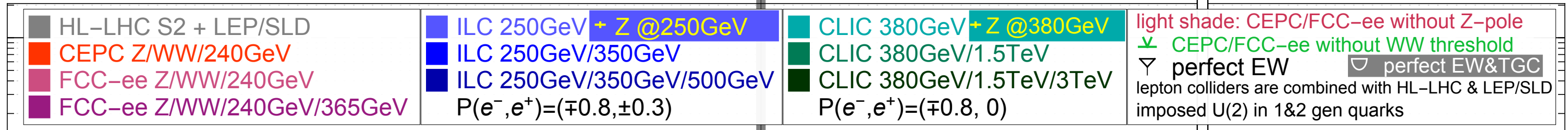
Comparing 3 EW scenarios: LEP/SLD, actual EW measurements, perfect EW measurements



Impact of Z-pole on Higgs.

J. De Blas et al. 1907.04311

Comparing 3 EW scenarios: LEP/SLD, actual EW measurements, perfect EW measurements

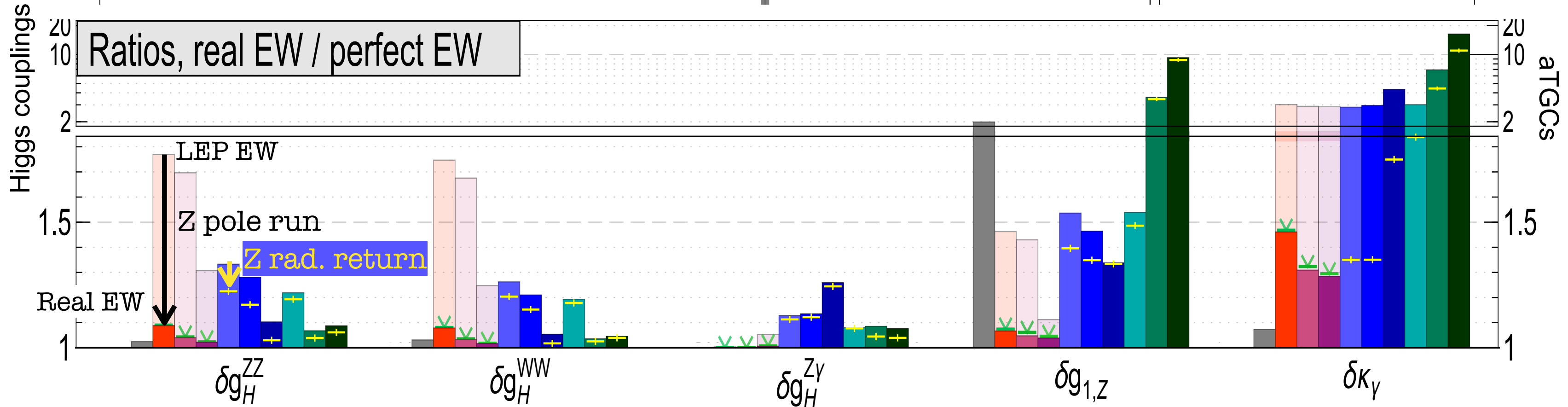
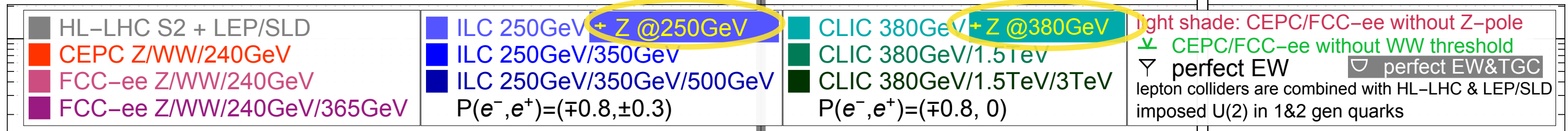


- FCC-ee and CEPC benefit a lot (>50% on HVV) from Z-pole run
- FCC-ee and CEPC EW measurements are almost perfect for what concerns Higgs physics (<10%).

Impact of Z-pole on Higgs.

J. De Blas et al. 1907.04311

Comparing 3 EW scenarios: LEP/SLD, actual EW measurements, perfect EW measurements

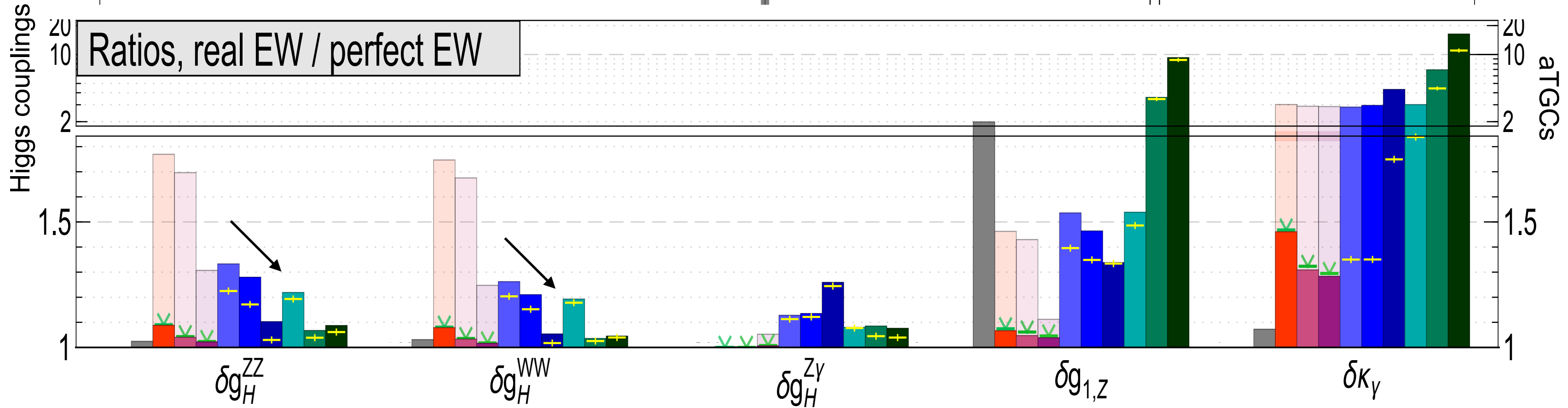
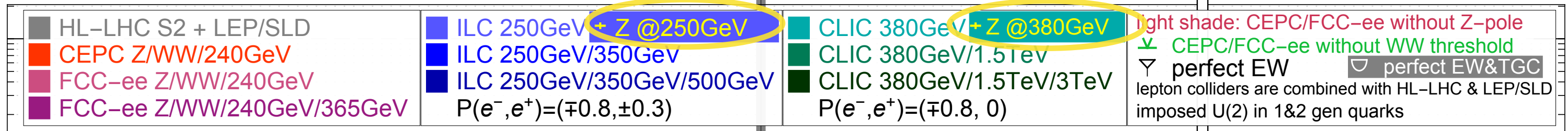


- FCC-ee and CEPC benefit a lot (>50% on HVV) from Z-pole run
- FCC-ee and CEPC EW measurements are almost perfect for what concerns Higgs physics (<10%).
- LEP EW measurements are a limiting factor (~30%) to Higgs precision at ILC, especially for the first runs
But EW measurements at high energy (via Z-radiative return) help mitigating this issue

Impact of Z-pole on Higgs.

J. De Blas et al. 1907.04311

Comparing 3 EW scenarios: LEP/SLD, actual EW measurements, perfect EW measurements



- Higher energy runs reduce the EW contamination in Higgs coupling extraction

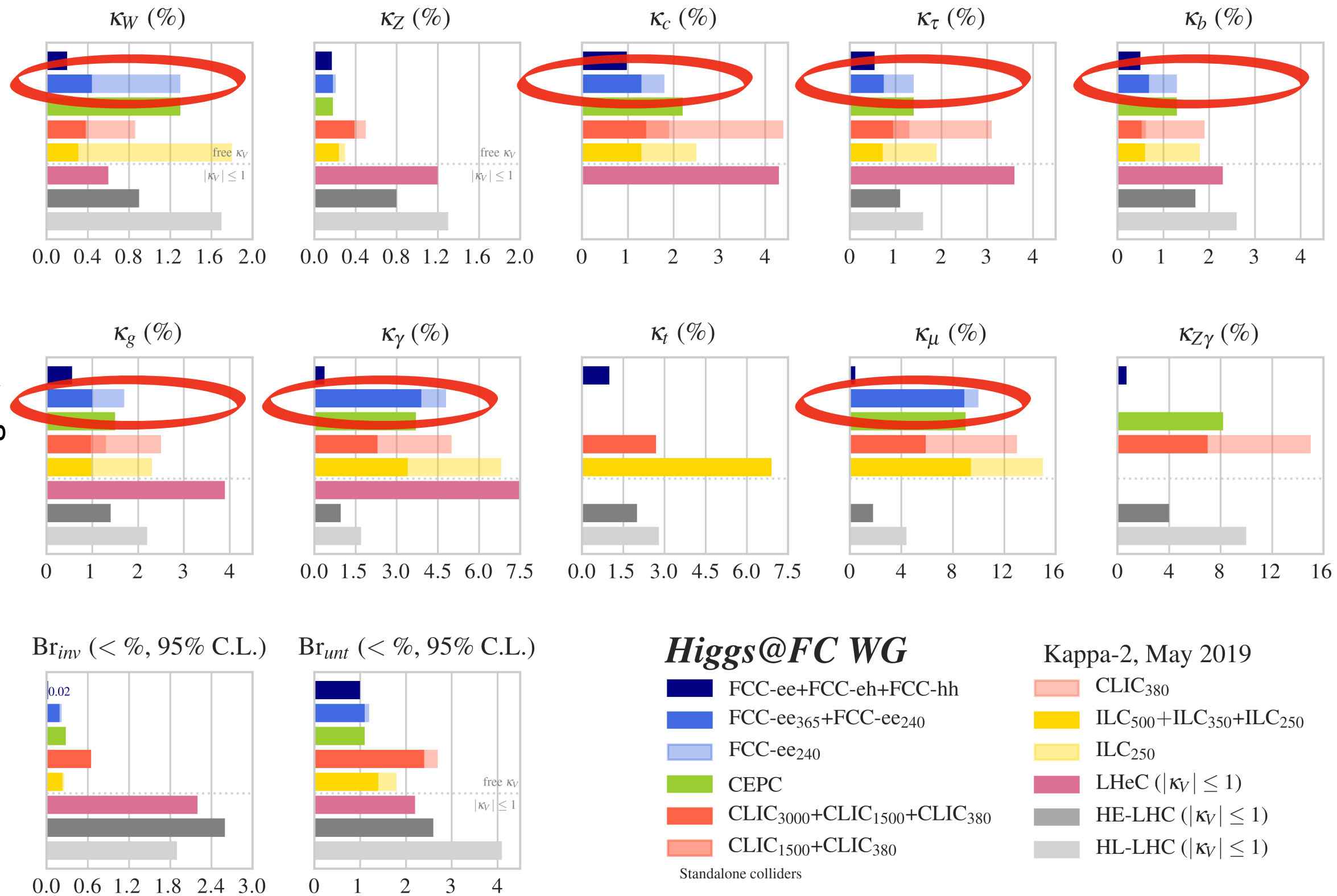
Complementarity 240↔365 GeV.

ECFA Higgs study group '19

Scenario BR_{inv} BR_{unt} include HL-LHC
 kappa-2 measured measured no

hadron collider cannot measure width
 need an assumption to close the fit

e.g. $\kappa_V < 1$



[back to main discussion](#)

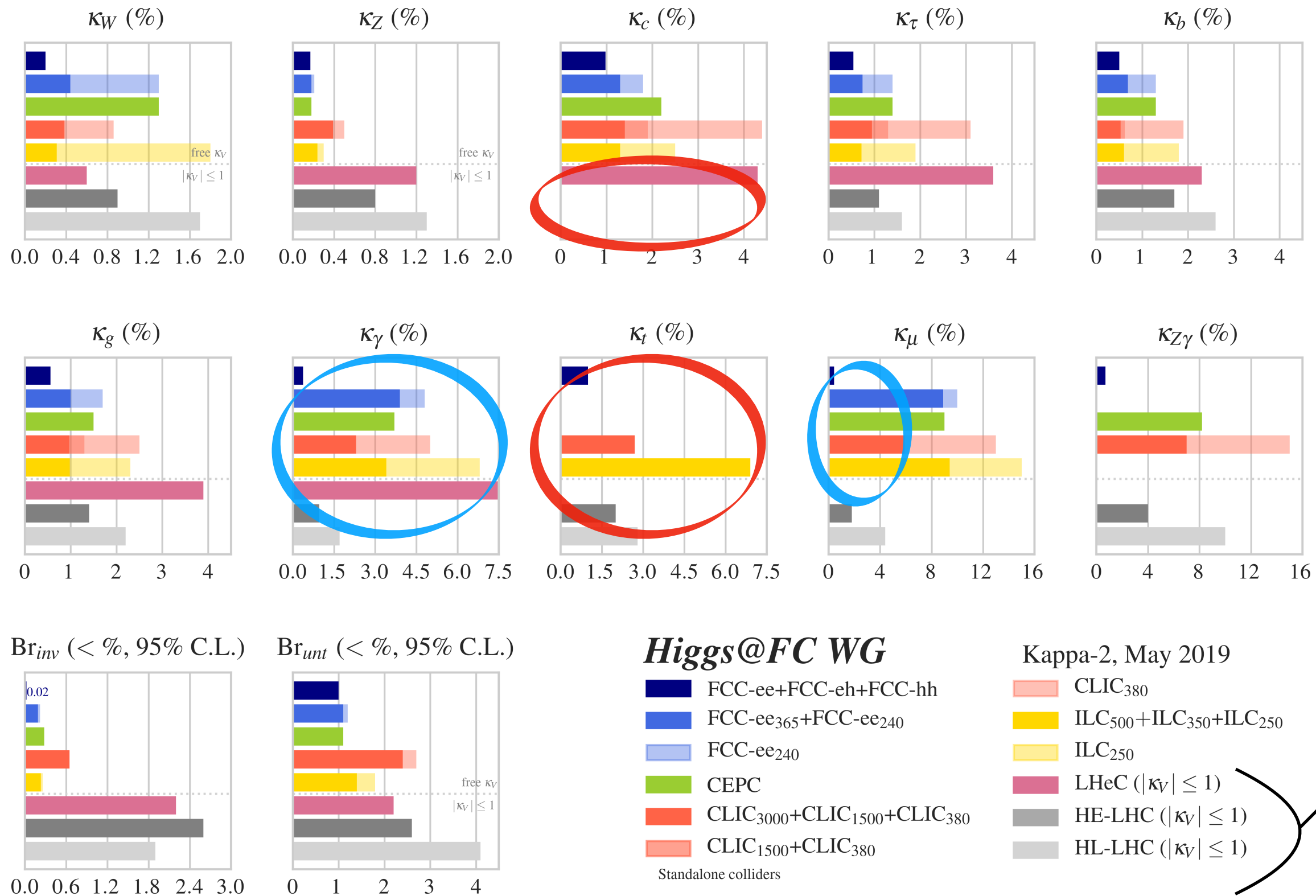
Complementarity FCC-ee ↔ HL-LHC.

ECFA Higgs study group '19

Scenario	BR_{inv}	BR_{unt}	include HL-LHC
kappa-2	measured	measured	no

hadron collider cannot measure width
need an assumption to close the fit

e.g. $\kappa_V < 1$



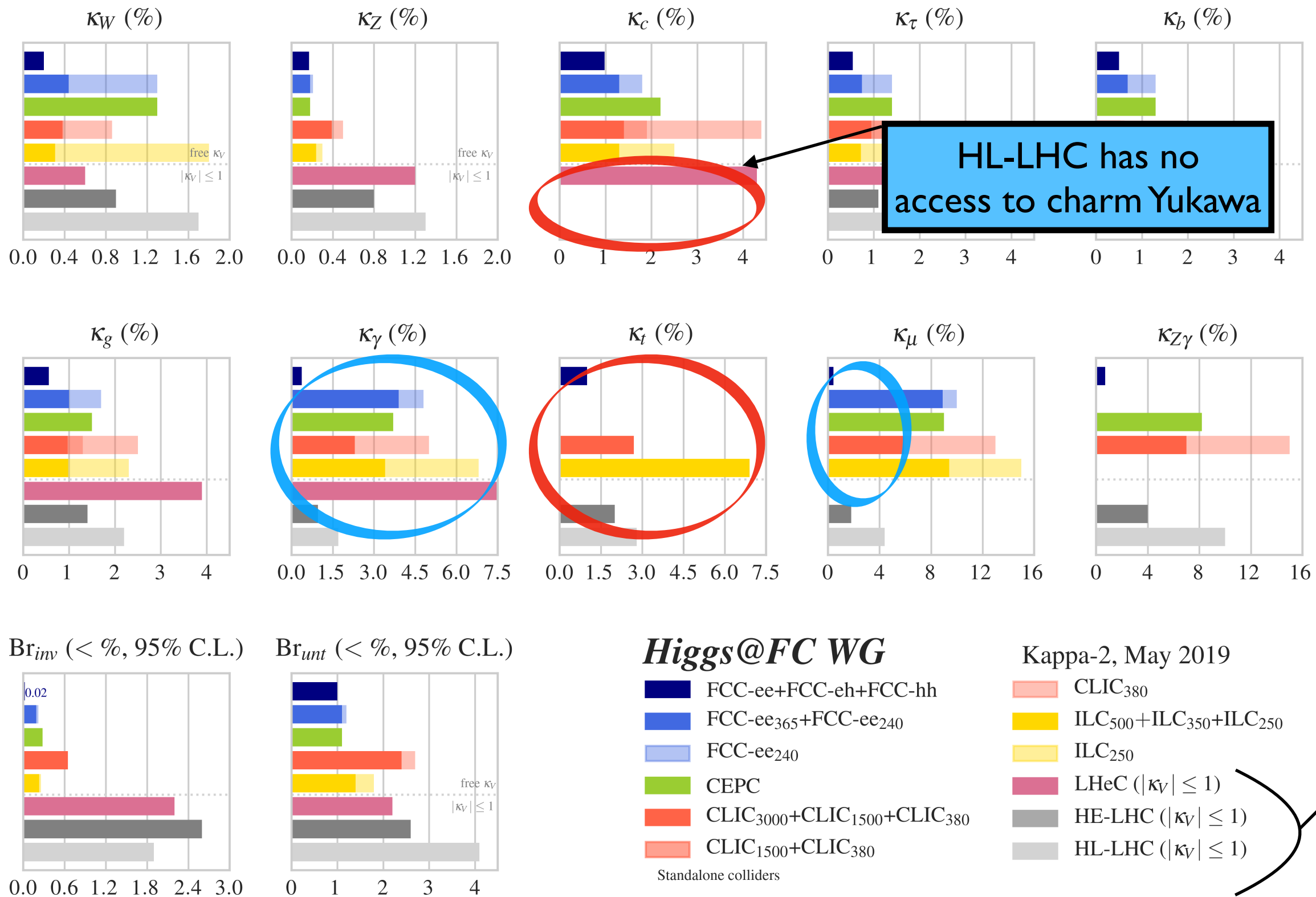
back to main discussion

Complementarity FCC-ee ↔ HL-LHC.

ECFA Higgs study group '19

Scenario	BR_{inv}	BR_{unt}	include HL-LHC
kappa-2	measured	measured	no

hadron collider cannot measure width
need an assumption to close the fit
e.g. $\kappa_V < 1$



HL-LHC has no access to charm Yukawa

[back to main discussion](#)

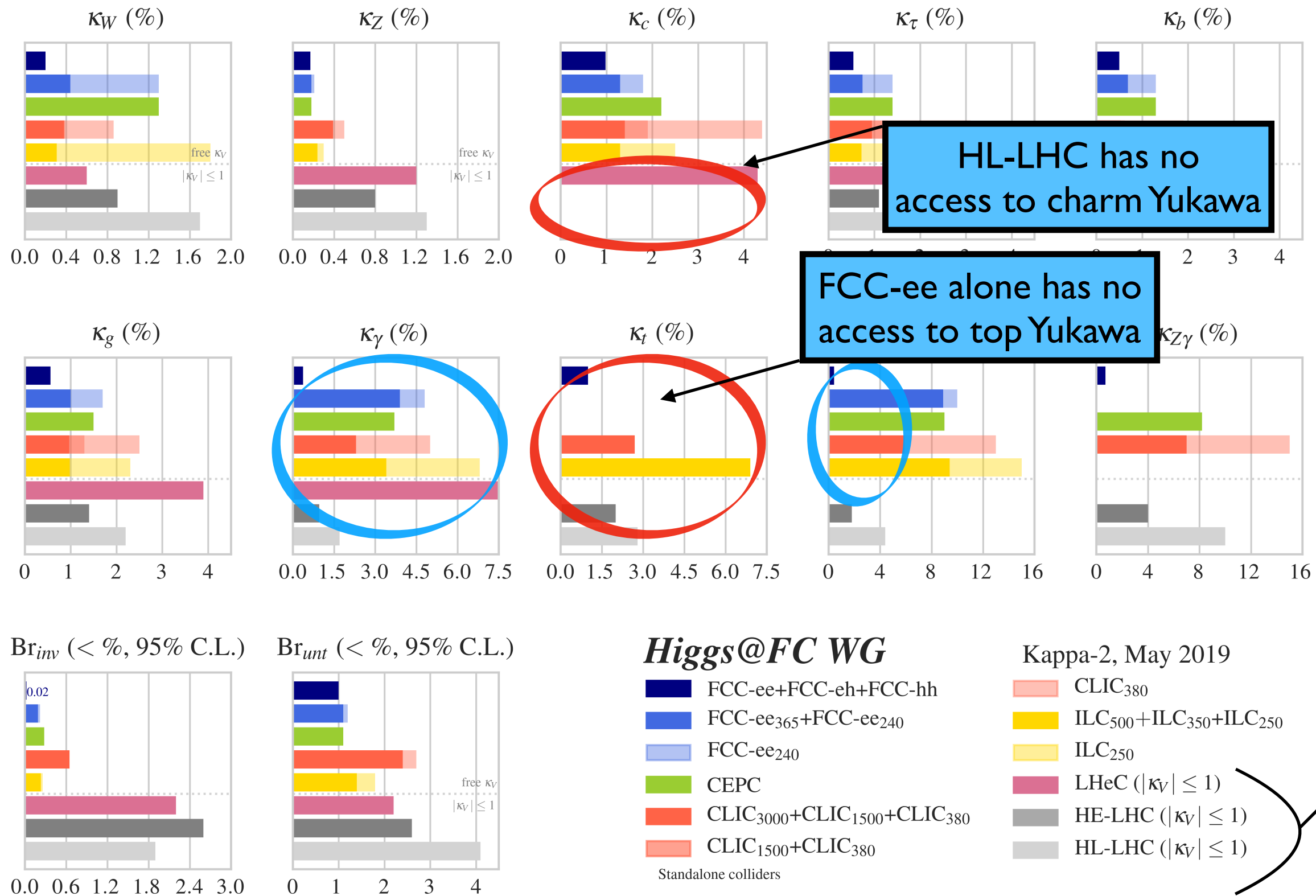
assumption needed for the fit to close at hadron machines

Complementarity FCC-ee ↔ HL-LHC.

ECFA Higgs study group '19

Scenario	BR_{inv}	BR_{unt}	include HL-LHC
kappa-2	measured	measured	no

hadron collider cannot measure width
need an assumption to close the fit
e.g. $\kappa_V < 1$



back to main discussion

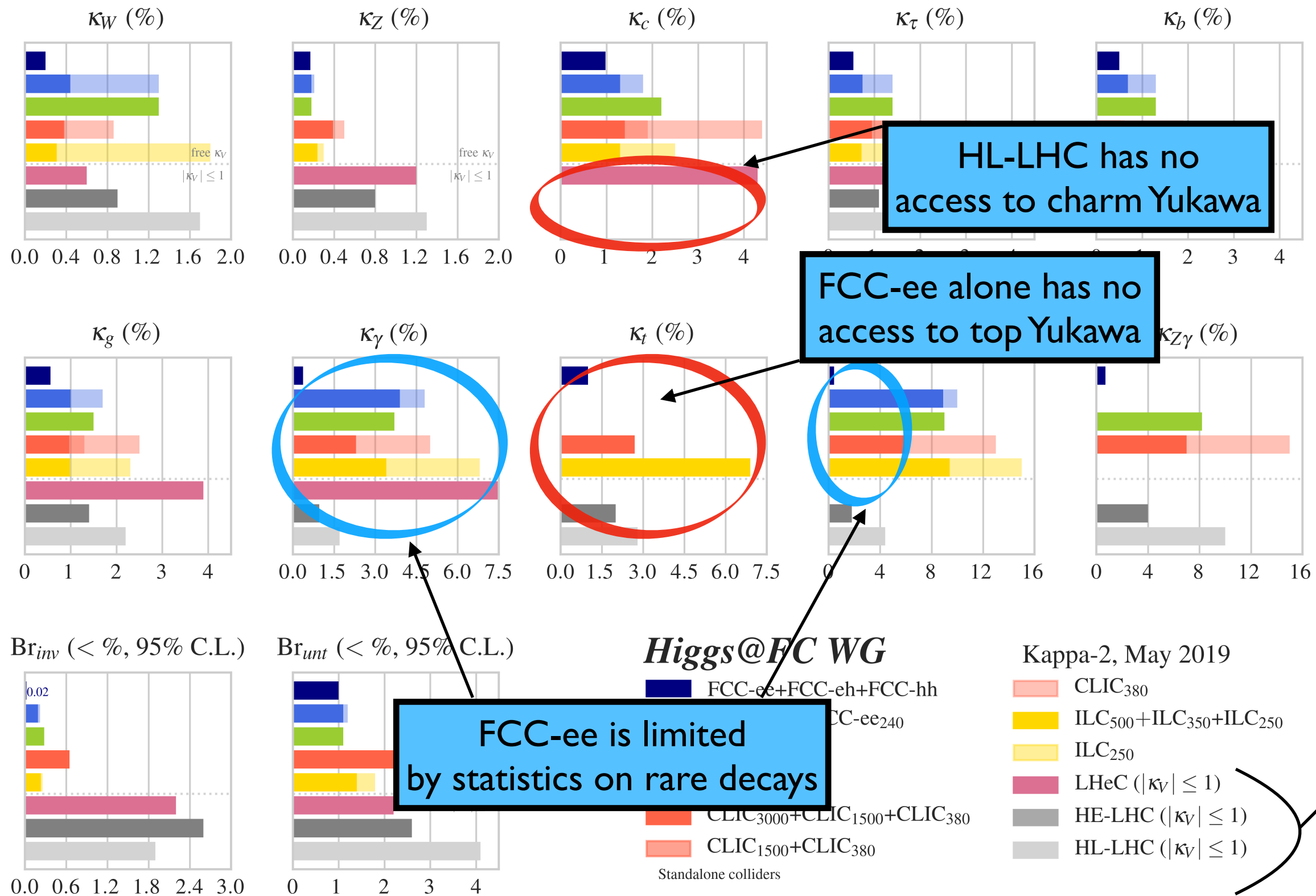
assumption
needed for the fit
to close at hadron
machines

Complementarity FCC-ee ↔ HL-LHC.

ECFA Higgs study group '19

Scenario	BR_{inv}	BR_{unt}	include HL-LHC
kappa-2	measured	measured	no

hadron collider cannot measure width
need an assumption to close the fit
e.g. $\kappa_V < 1$



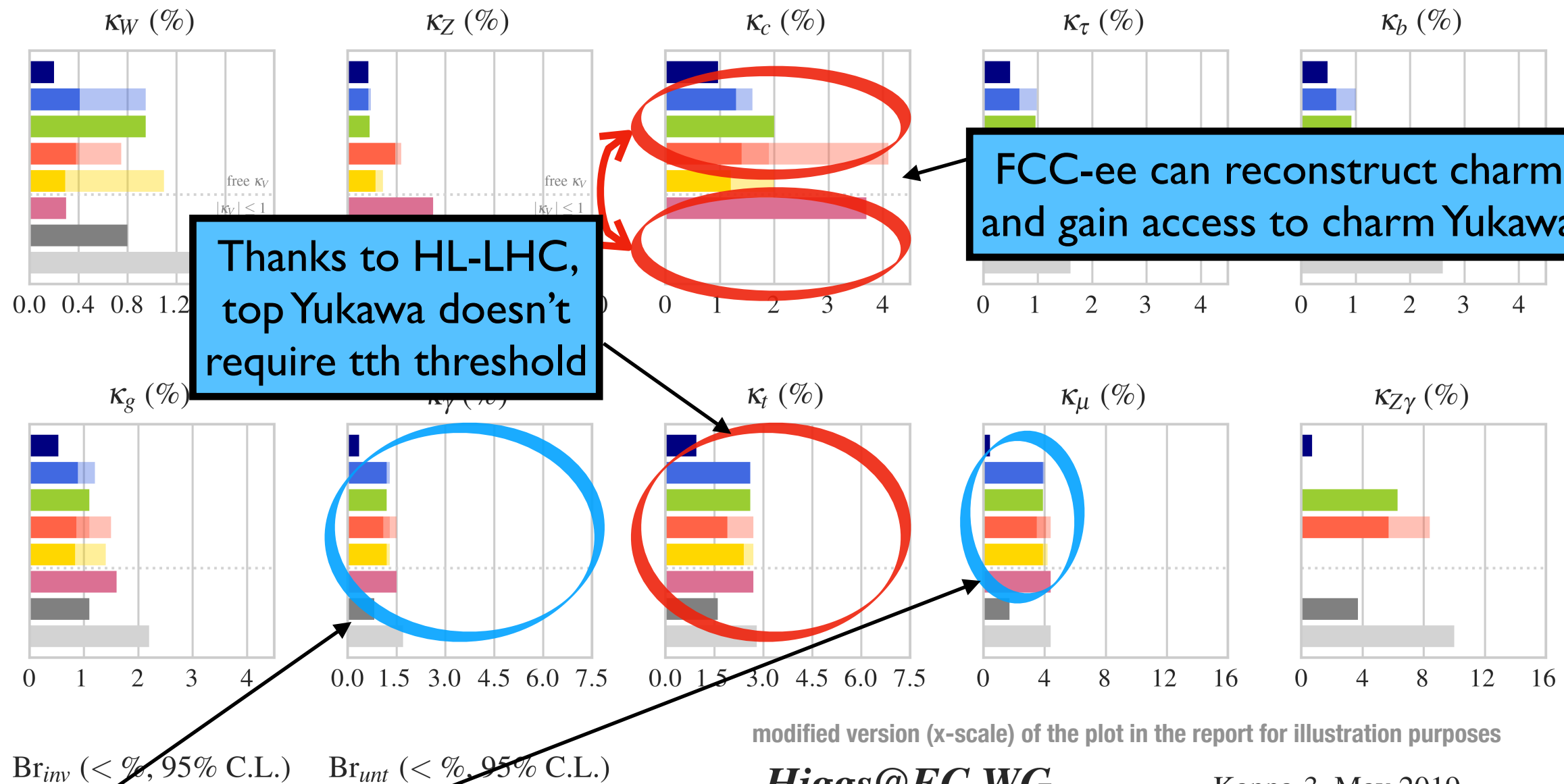
back to main discussion

assumption
needed for the fit
to close at hadron
machines

Complementarity FCC-ee ↔ HL-LHC.

ECFA Higgs study group '19

Scenario	include HL-LHC
kappa-3	yes
BR_{inv}	measured
BR_{unt}	measured



Thanks to HL-LHC, top Yukawa doesn't require tth threshold

FCC-ee can reconstruct charm and gain access to charm Yukawa

LHC brings statistics
FCC-ee adds a bit of sensitivity

Important synergy HL-LHC — low energy lepton colliders

1. Top/Charm Yukawa
2. Statistically limited channels: $\gamma\gamma$, $\mu\mu$

back to main discussion

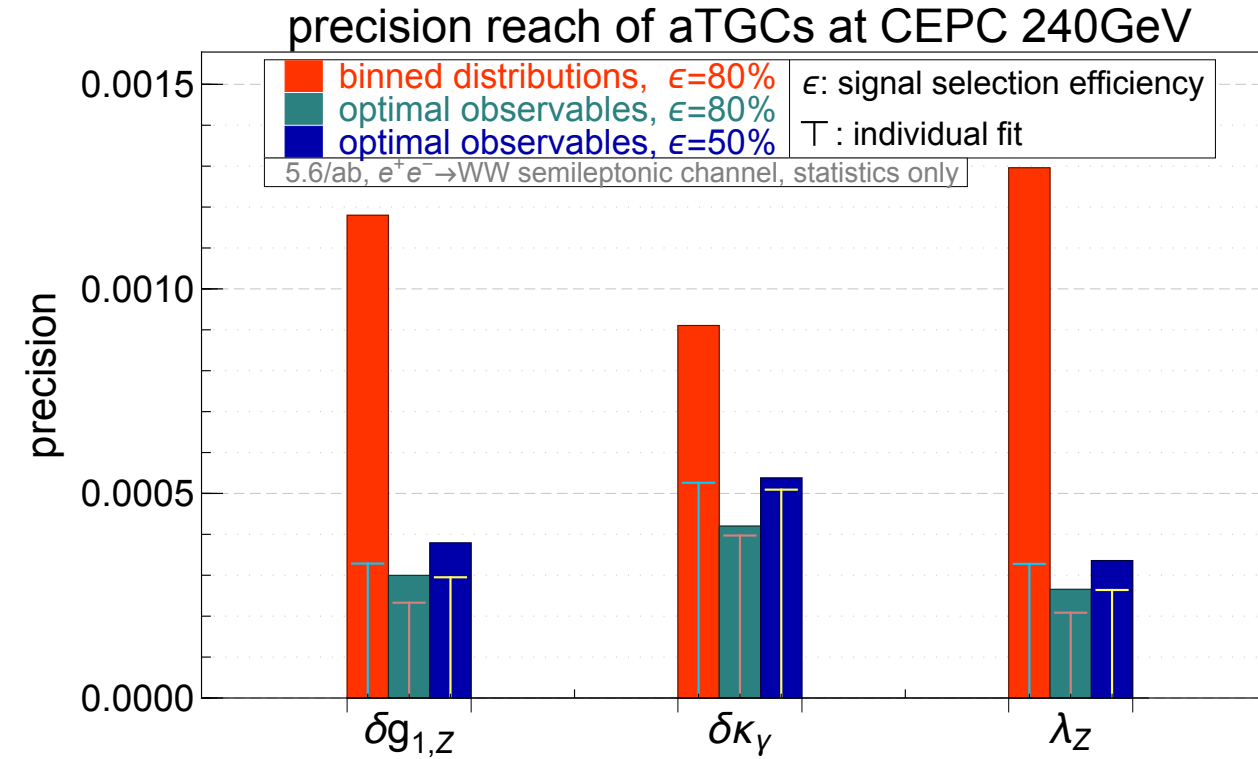
modified version (x-scale) of the plot in the report for illustration purposes

Higgs@FC WG

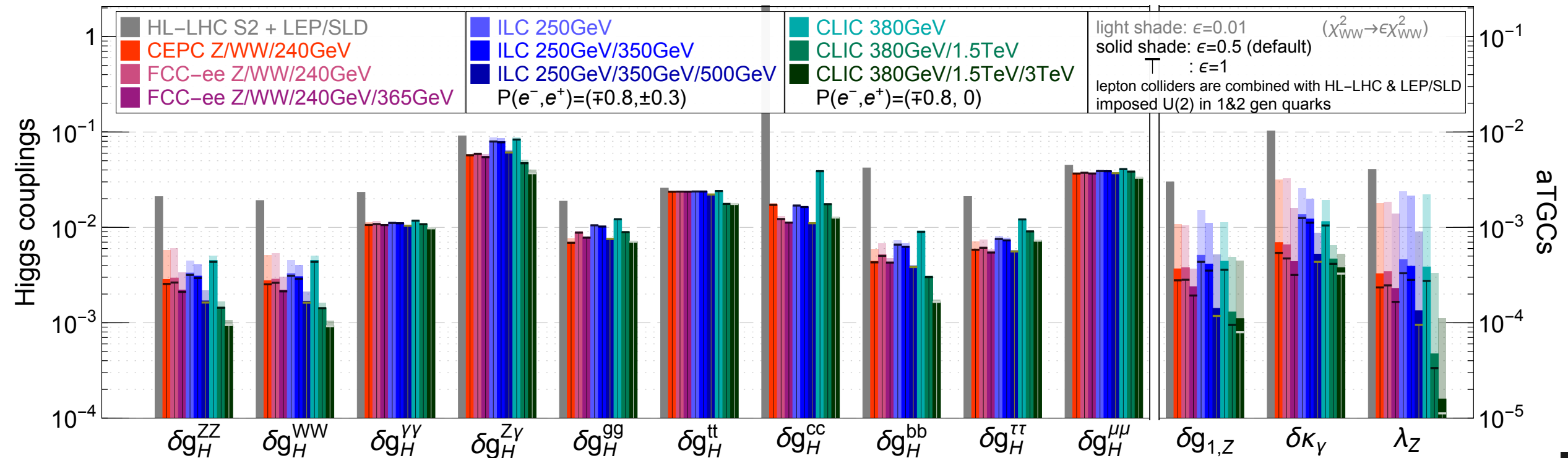
Kappa-3, May 2019

Impact of Diboson Systematics.

J. De Blas, G. Durieux, C. Grojean, J. Gu, A. Paul 1907.04311



precision reach with different assumptions on $e^+e^- \rightarrow WW$ measurements

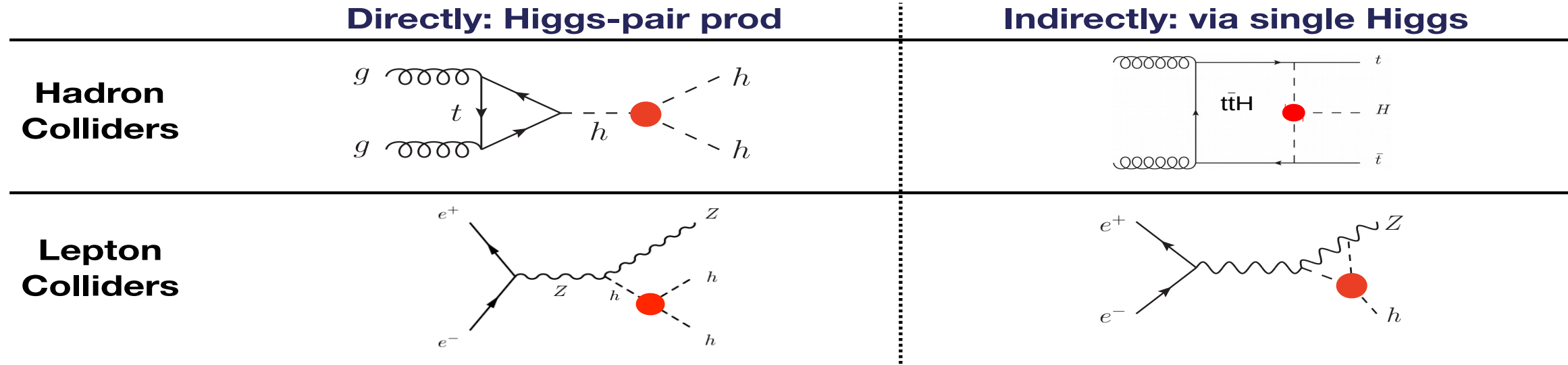


Higgs self-coupling.

Higgs self-couplings is very interesting for a multitude of reasons
(vacuum stability, hierarchy, baryogenesis, GW, EFT probe...).

How much can it deviate from SM given the tight constraints on other Higgs couplings?
Do you need to reach HH production threshold to constrain h^3 coupling?

ECFA Higgs study group '19



	di-Higgs	single-H
exclusive	<p>1. di-H, excl.</p> <ul style="list-style-type: none"> • Use of $\sigma(HH)$ • only deformation of $\kappa\lambda$ 	<p>3. single-H, excl.</p> <ul style="list-style-type: none"> • single Higgs processes at higher order • only deformation of $\kappa\lambda$
global	<p>2. di-H, glob.</p> <ul style="list-style-type: none"> • Use of $\sigma(HH)$ • deformation of $\kappa\lambda$ + of the single-H couplings (a) do not consider the effects at higher order of $\kappa\lambda$ to single H production and decays (b) these higher order effects are included 	<p>4. single-H, glob.</p> <ul style="list-style-type: none"> • single Higgs processes at higher order • deformation of $\kappa\lambda$ + of the single Higgs couplings

Large self-coupling scenarios.

Generically: $\left| \frac{\delta_{h^3}}{\delta_{\text{single } h}} \right| \sim O(1)$ (composite Higgs/susy)

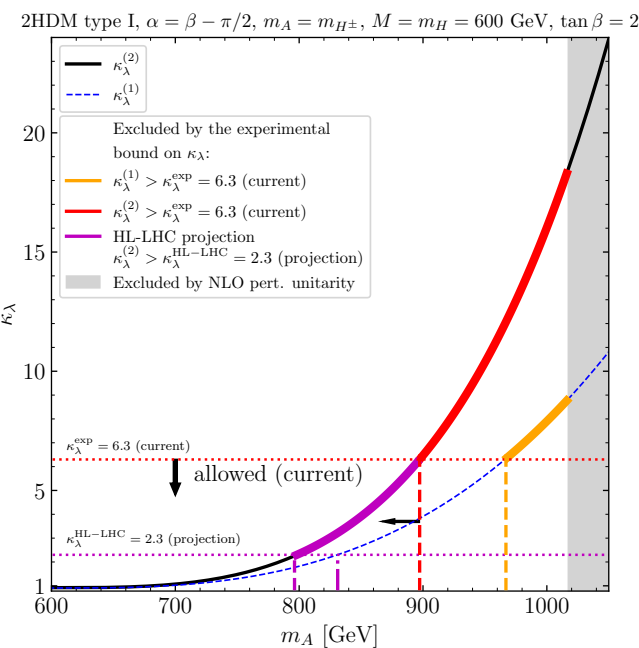
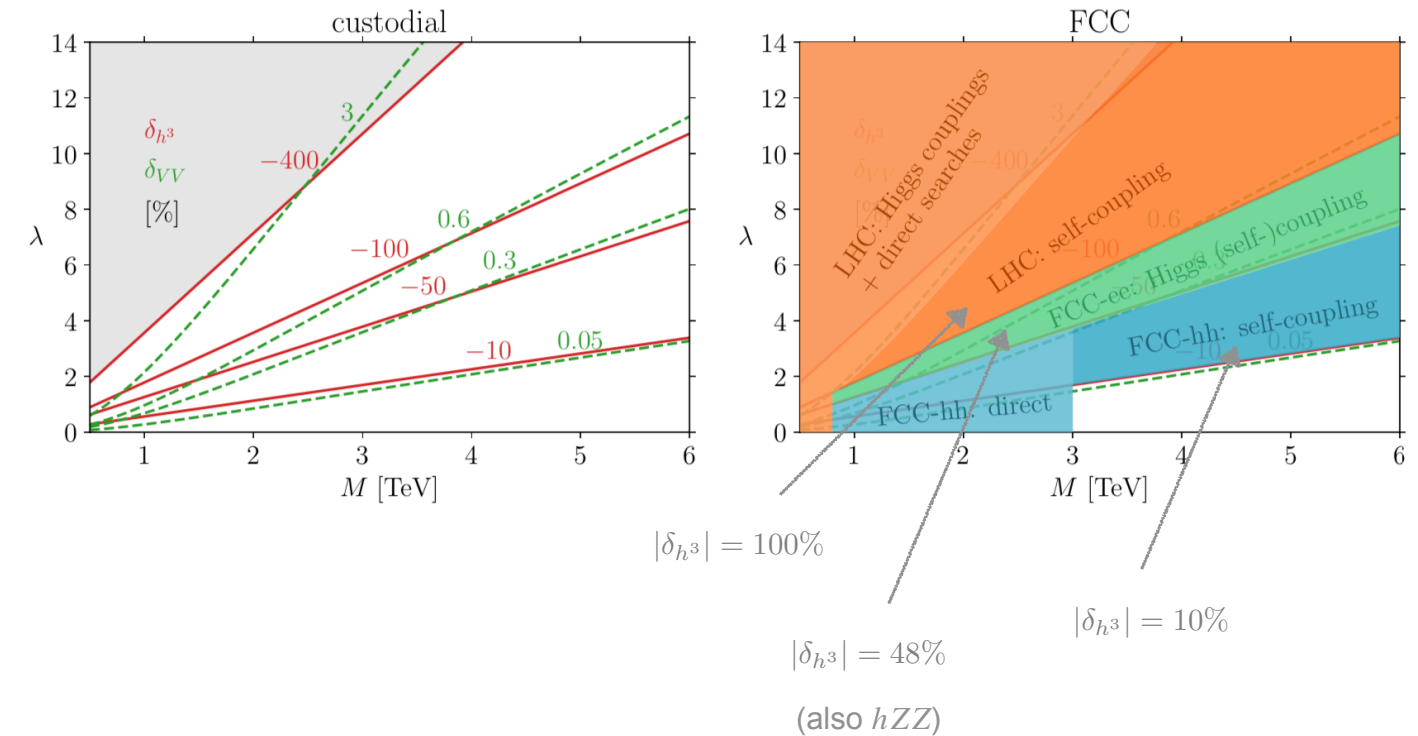
Particular exceptions: Higgs DM-portal models or custodial EW quadruplet

DiVita et al.: 1704.01953

Falkowski, Rattazzi: 1902.05936

Durieux, McCullough, Salvioni: 2209.00666

h^3 generically is not a tool to discover BSM but exceptions exist.

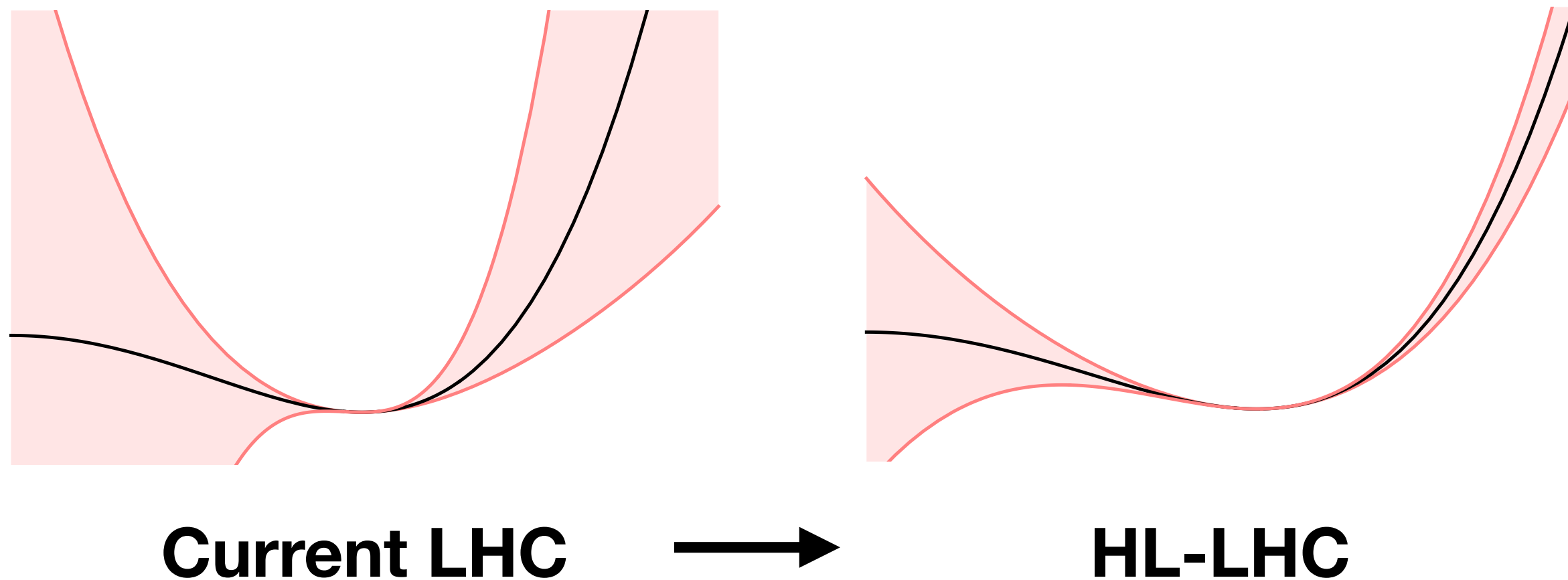


Other exceptions: non-decoupled/fine-tuned spectra

Bahl, Braathen, Weiglein: 2202.03453

Large self-coupling scenarios.

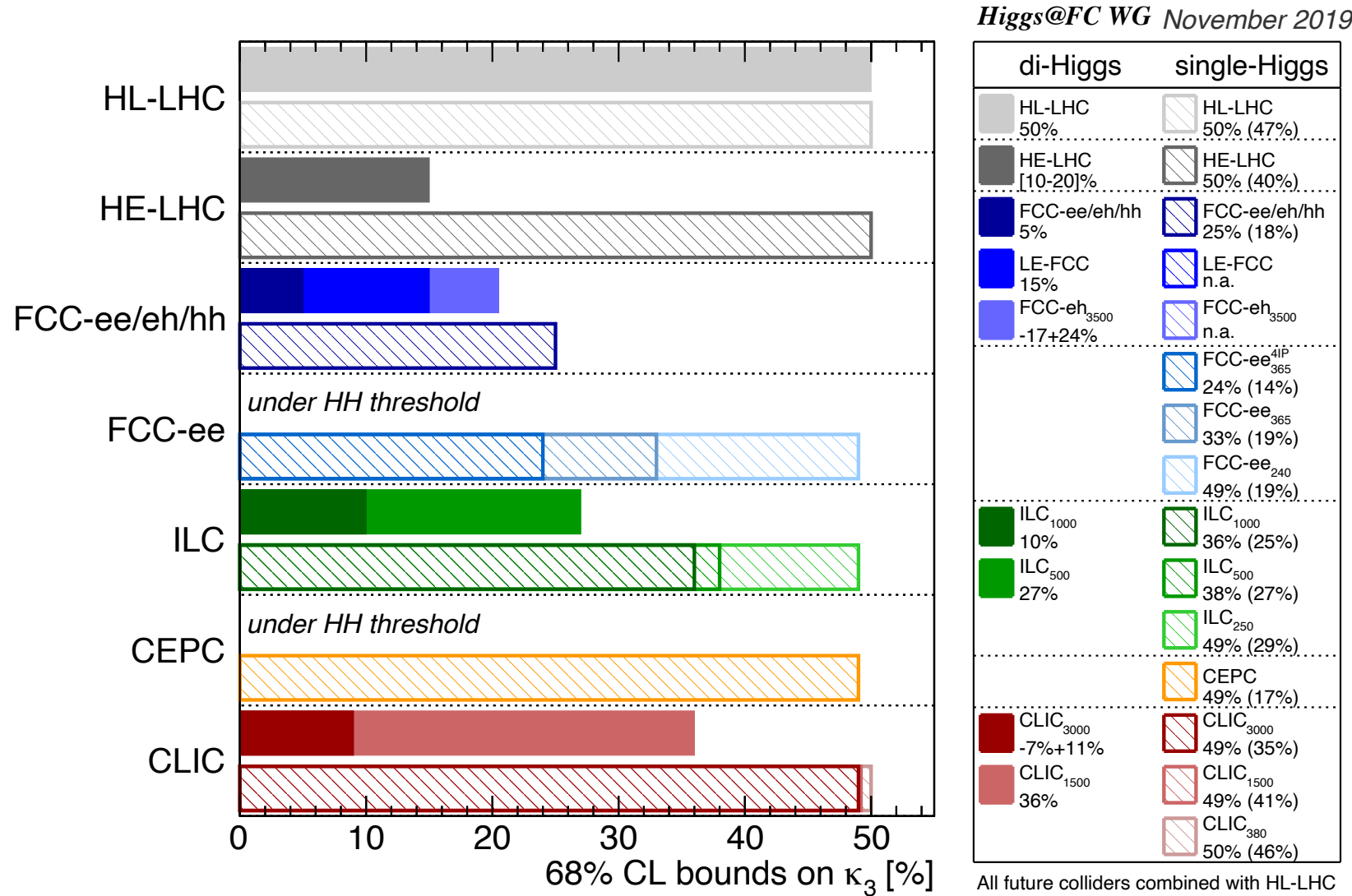
It is true that we haven't "measured" the Higgs potential but there are only peculiar physics scenarios that produce large deviations in the shape of the potential without leaving imprints elsewhere.



R. Petrossian-Byrne/N. Craig @ LCWS'23

Important to understand which dynamics is really probed when embarking into challenging measurements. Actually, double Higgs production is also interesting to probe new physics in its tail rather than near threshold (where the sensitivity to Higgs self-coupling comes from).

Higgs self-coupling.



1

Don't need to reach HH threshold to have access to h^3 .
Z-pole run is very important if the HH threshold cannot be reached

2

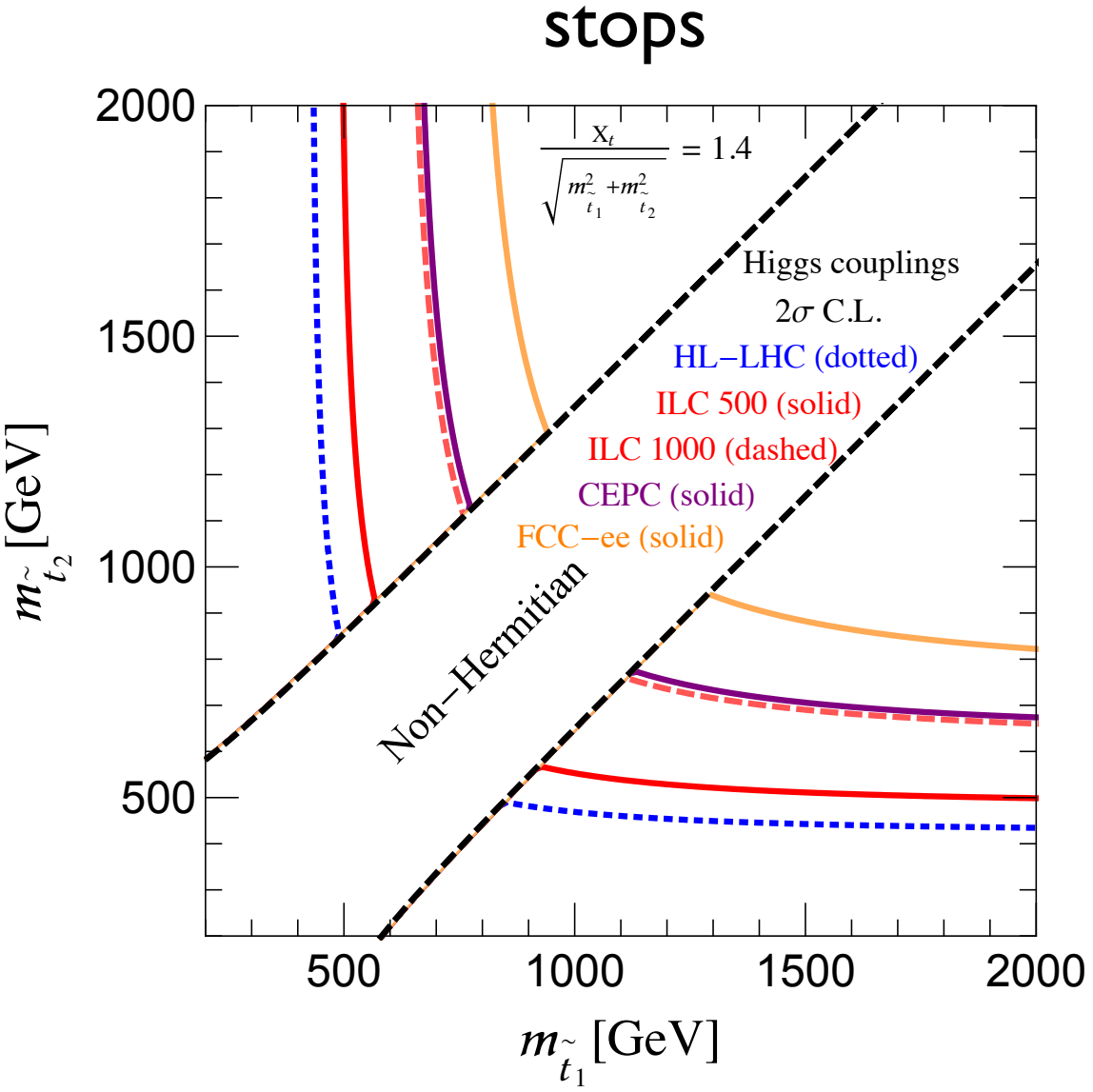
The determination of h^3 at FCC-hh relies on HH channel, for which FCC-ee is of little direct help. But the extraction of h^3 requires precise knowledge of y_t .
 $1\% y_t \leftrightarrow 5\% h^3$
Precision measurement of y_t needs ee

50% sensitivity: establish that $h^3 \neq 0$ at 95%CL
20% sensitivity: 5σ discovery of the SM h^3 coupling
5% sensitivity: getting sensitive to quantum corrections to Higgs potential

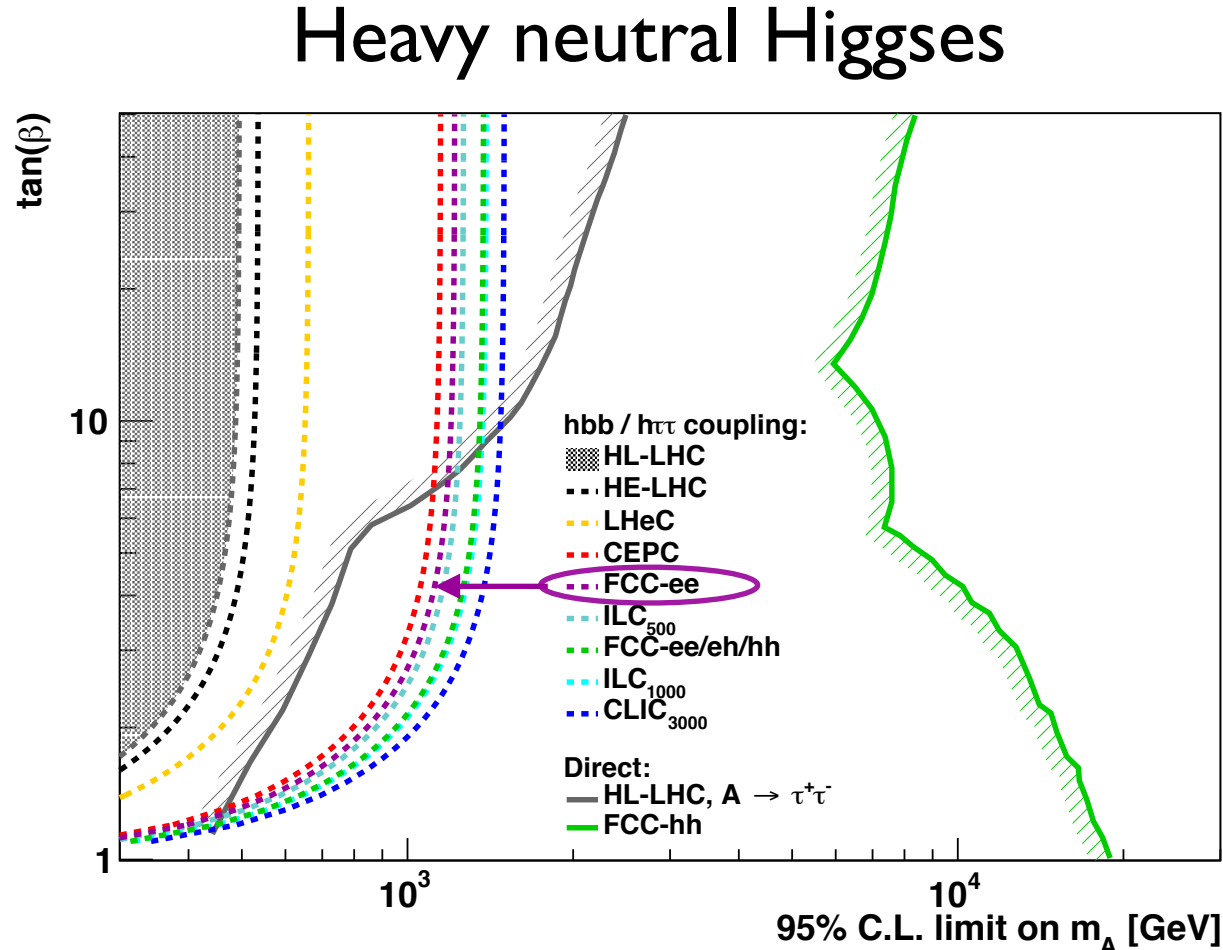
Discovery potential beyond LHC

Discovery Potential Beyond LHC.

Precisely measured EW and Higgs observables are sensitive to heavy New Physics
 Examples of improved sensitivity wrt direct reach @ HL-LHC: SUSY



Fan, Reece, Wang '14



ESU Physics BB '19

Discovery Potential Beyond LHC.

Precisely measured EW and Higgs observables are sensitive to heavy New Physics
 Examples of improved sensitivity wrt direct reach @ HL-LHC: Composite Higgs

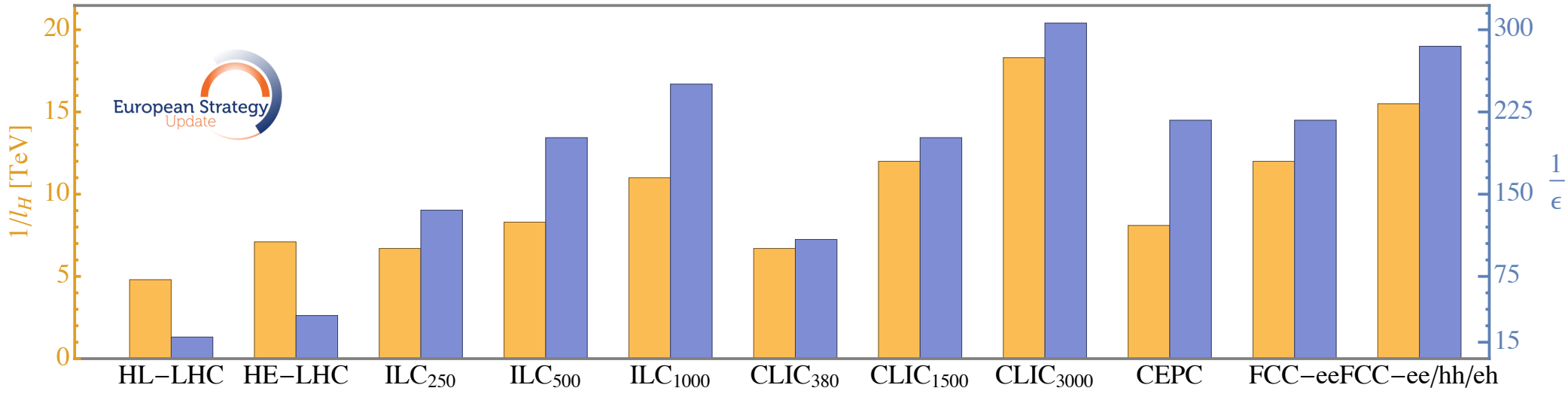
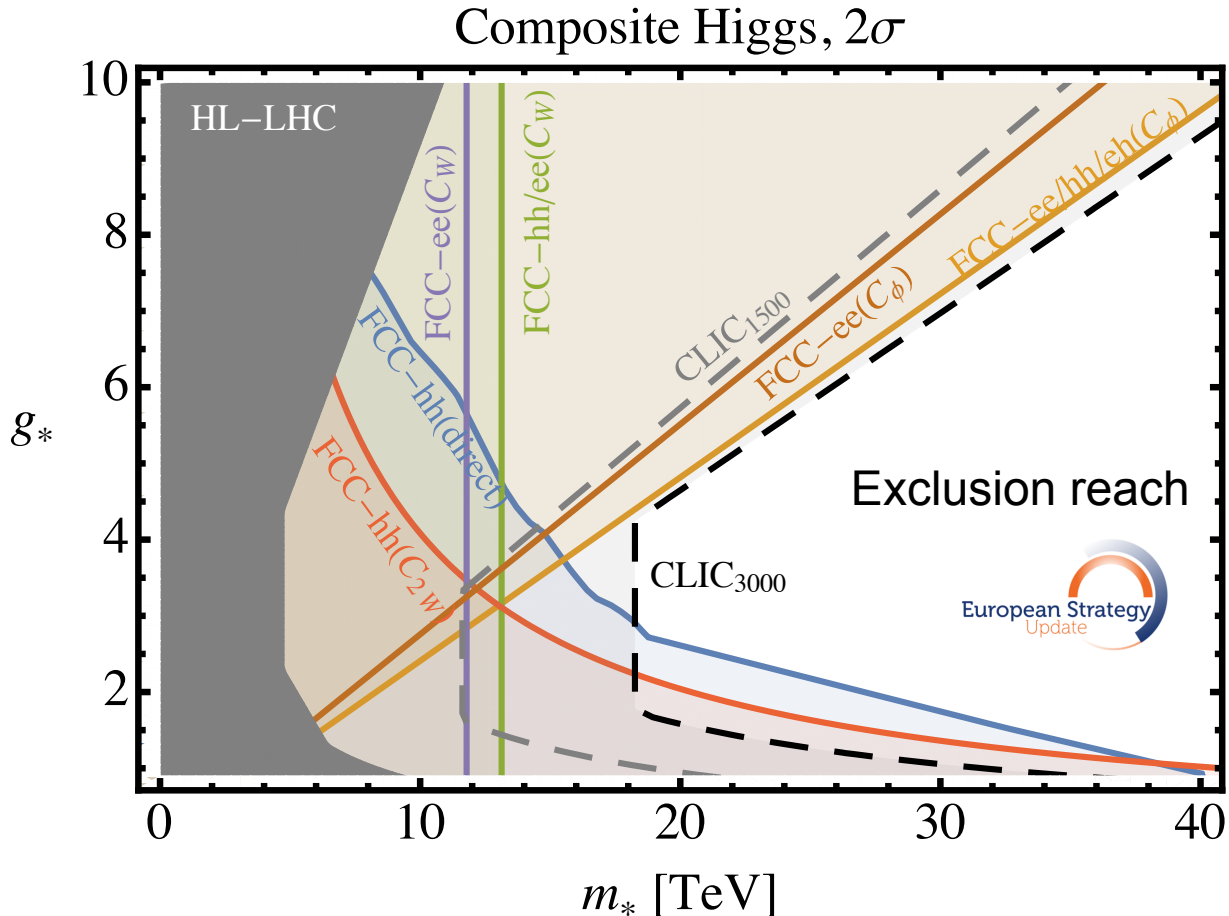


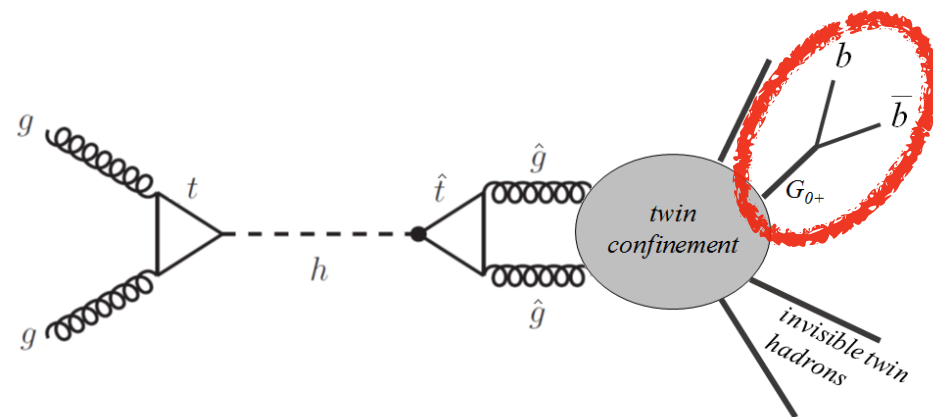
Fig. 8.5: Exclusion reach of different colliders on the inverse Higgs length $1/\ell_H = m_*$ (orange bars, left axis) and the tuning parameter $1/\epsilon$ (blue bars, right axis), obtained by choosing the weakest bound valid for any value of the coupling constant g_* .

Direct Searches for Elusive New Physics

- **LLP searches with displaced vertices**

e.g. in twin Higgs models glueballs that mix with the Higgs and decay back to b-quarks

Craig et al, arXiv:1501.05310



- **Rare decays**

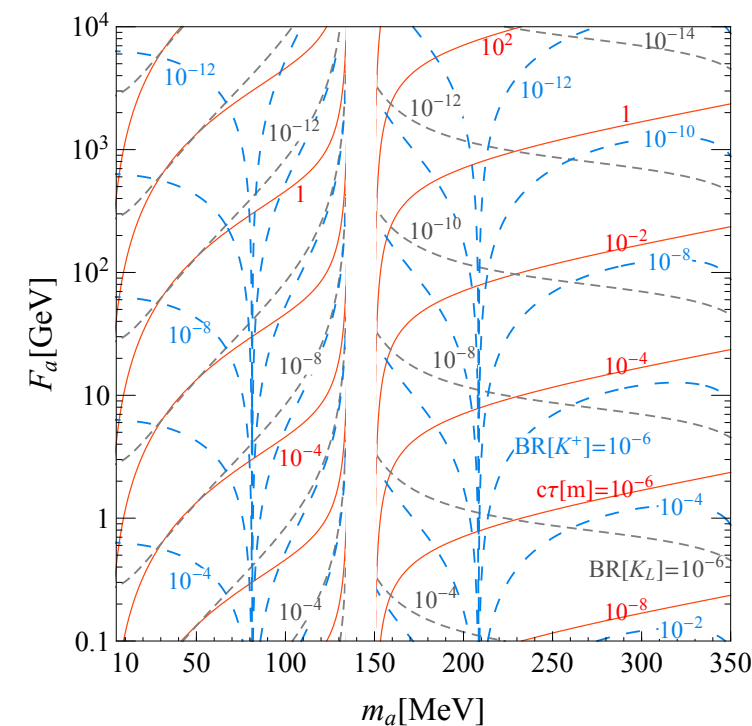
Gori et al arXiv:2005.05170

e.g. ALP mixing w/ SM mesons:

$$K_L \rightarrow \pi^0 a \rightarrow \pi^0 \gamma \gamma \text{ (KOTO)}$$

$$K^+ \rightarrow \pi^+ a \rightarrow \pi^+ \gamma \gamma \text{ (NA62)}$$

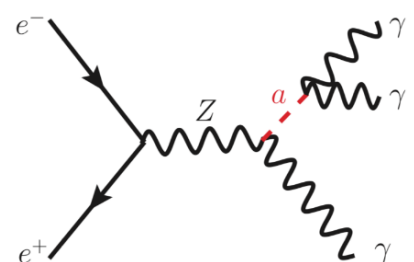
$$\mathcal{L} = \frac{\alpha_s}{8\pi F_a} a G_{\mu\nu} \tilde{G}^{\mu\nu}$$



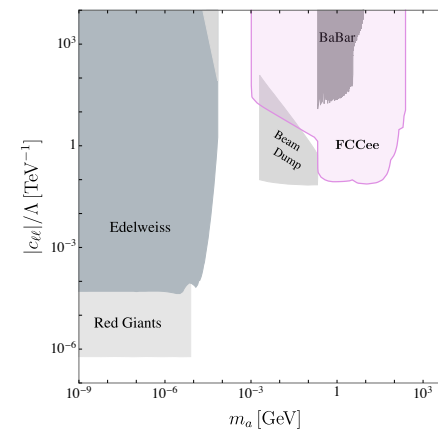
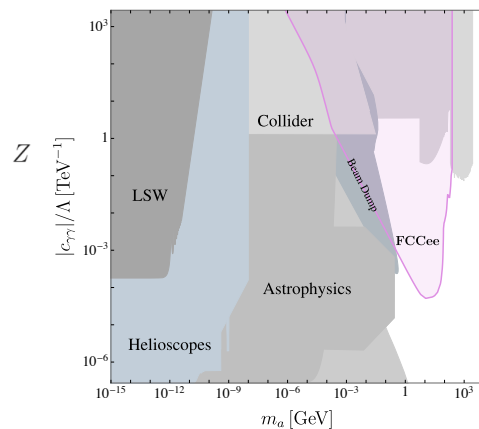
- **ALPs@ colliders**

e.g. $e^+e^- \rightarrow \gamma a$

$e^+e^- \rightarrow ha$



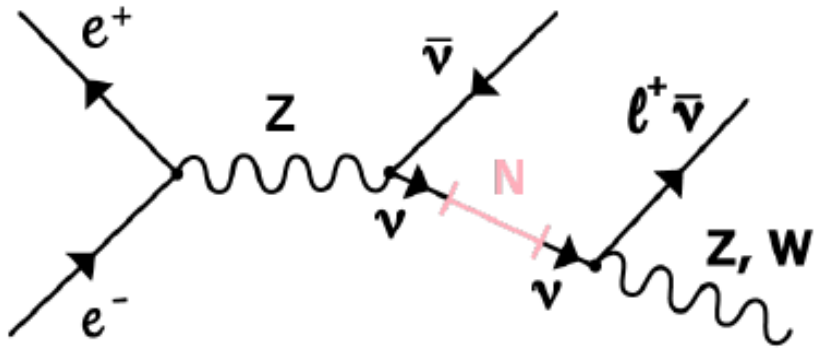
Knapen, Thamm arXiv:2108.08949



Astro/Cosmo \rightarrow long-lived ALPs
colliders \rightarrow short-lived ALPs MeV+

Search for ν_{RH} .

Direct observation
in Z decays
from LH-RH mixing



Important to understand

1. how neutrinos acquired mass
2. if lepton number is conserved
3. if leptogenesis is realised

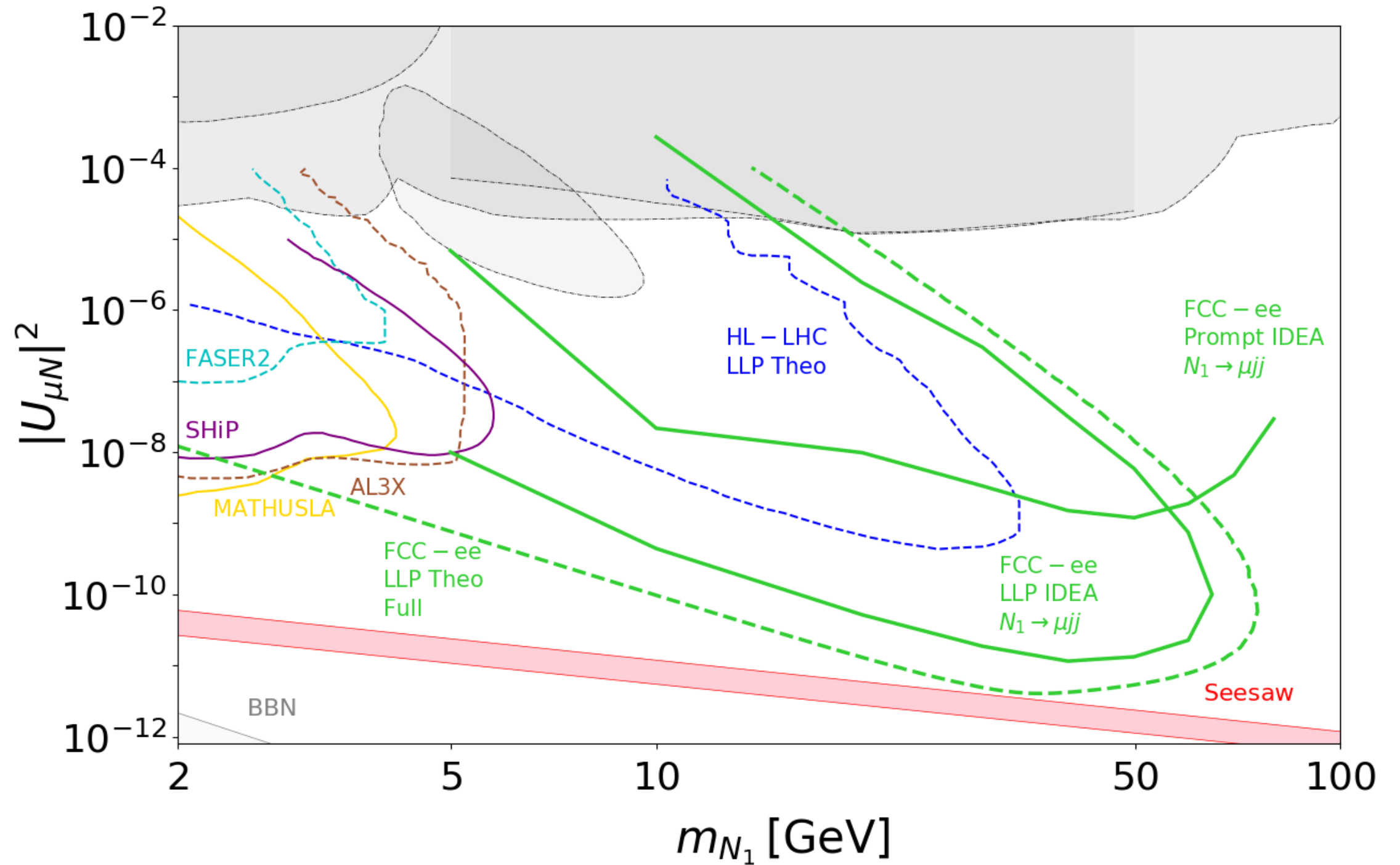
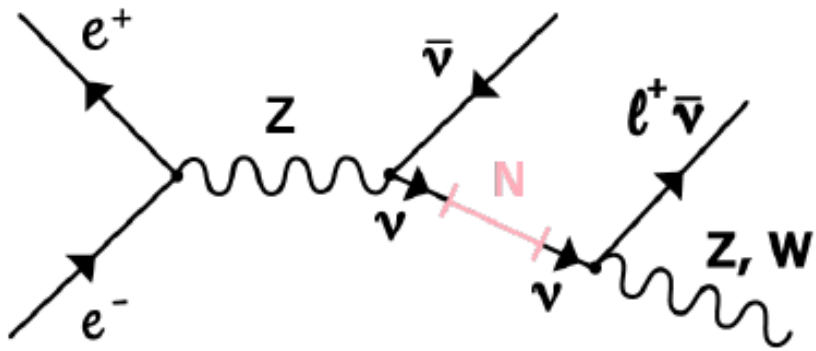


Fig. from mid-term report

Search for ν_{RH} .

Direct observation
in Z decays
from LH-RH mixing



Important to understand

1. how neutrinos acquired mass
2. if lepton number is conserved
3. if leptogenesis is realised

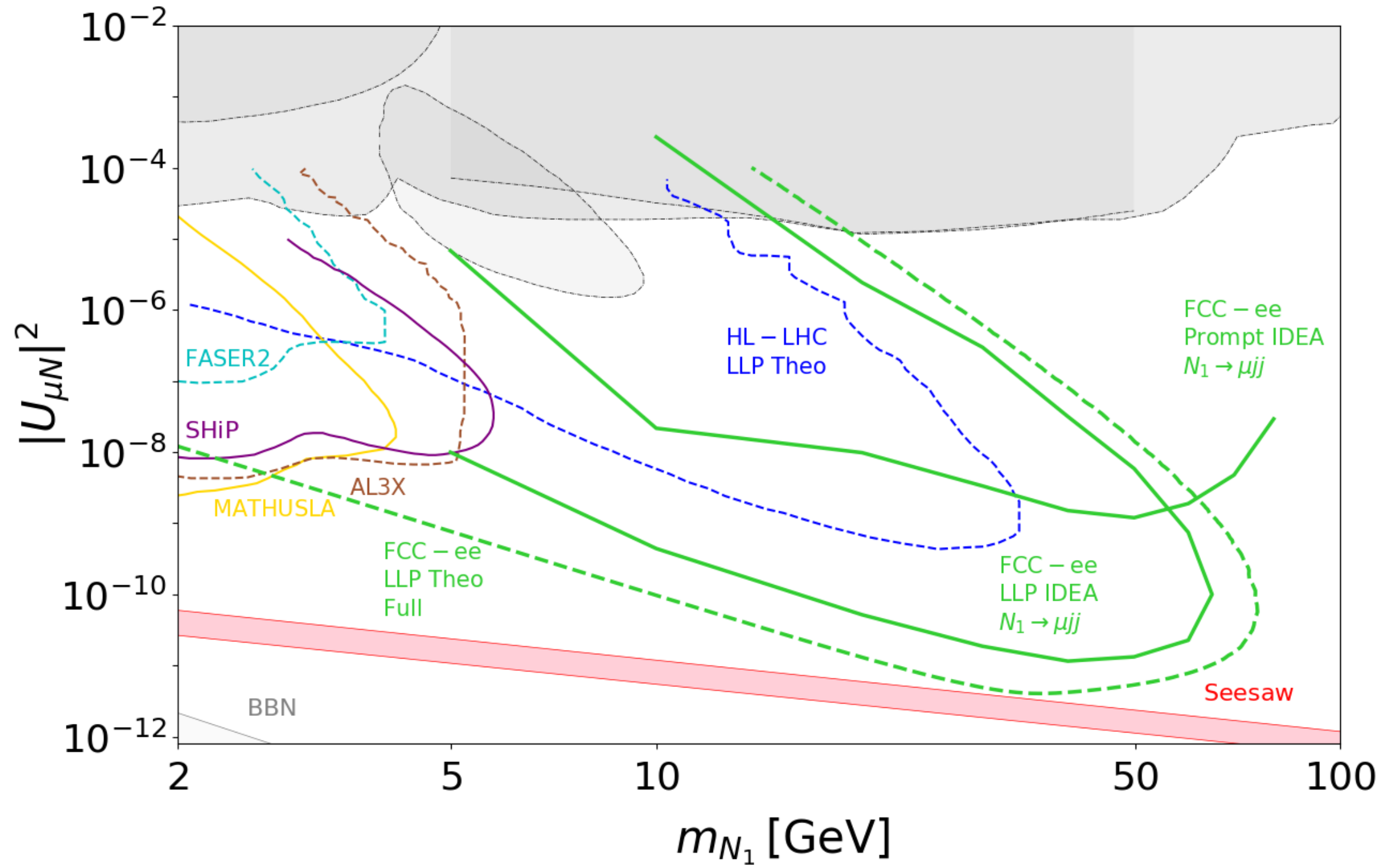
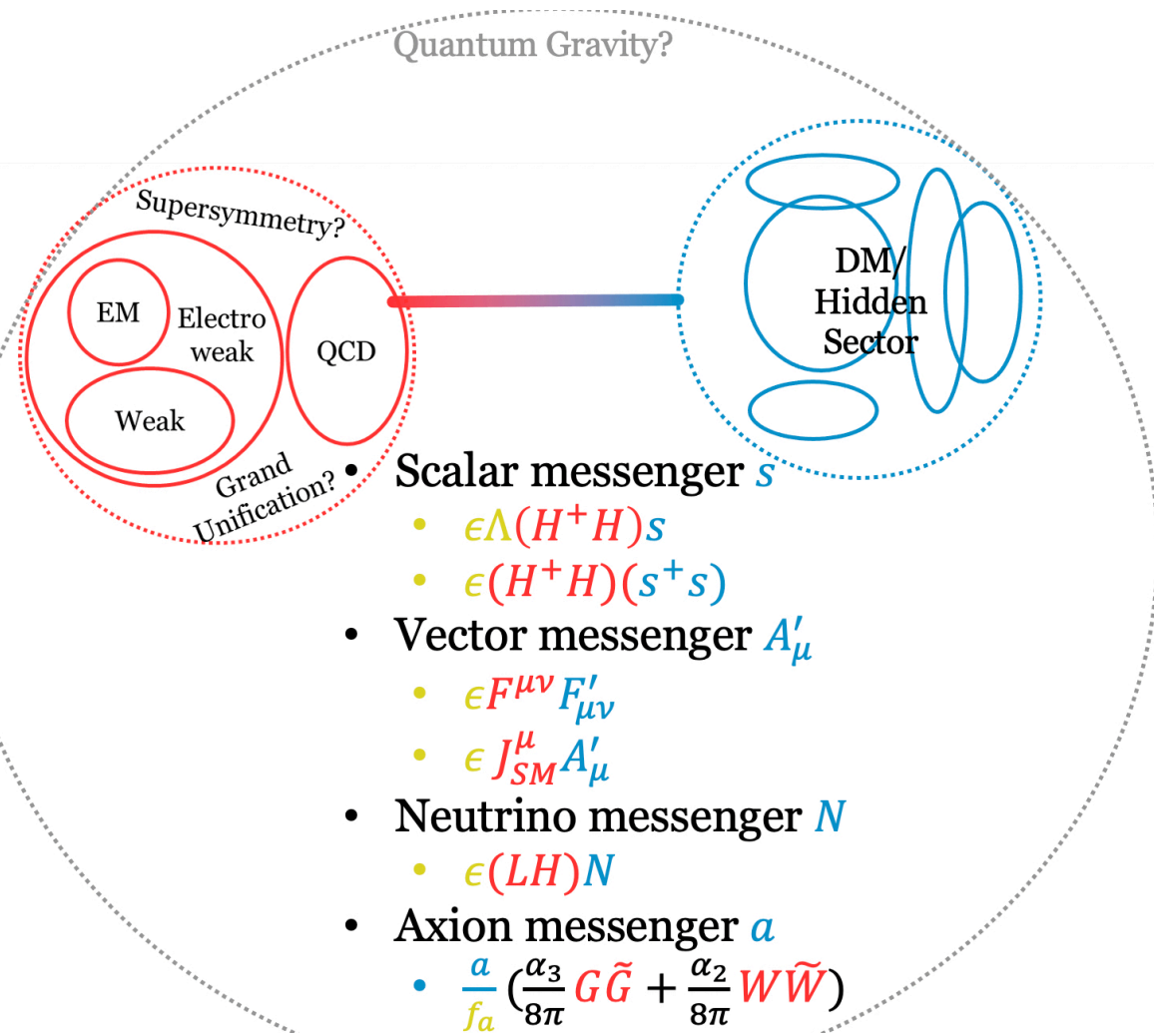


Fig. from mid-term report

Exotics/Long Lived Particles.

Z. Liu @ CEPC 2020

The Higgs could be a good portal to Dark Sector
— rich exotic signatures —



- Scalar messenger s
 - $\epsilon \Lambda (H^+ H) s$
 - $\epsilon (H^+ H) (s^+ s)$
- Vector messenger A'_μ
 - $\epsilon F^{\mu\nu} F'_{\mu\nu}$
 - $\epsilon J_{SM}^\mu A'_\mu$
- Neutrino messenger N
 - $\epsilon (LH) N$
- Axion messenger a
 - $\frac{a}{f_a} \left(\frac{\alpha_3}{8\pi} G\tilde{G} + \frac{\alpha_2}{8\pi} W\tilde{W} \right)$

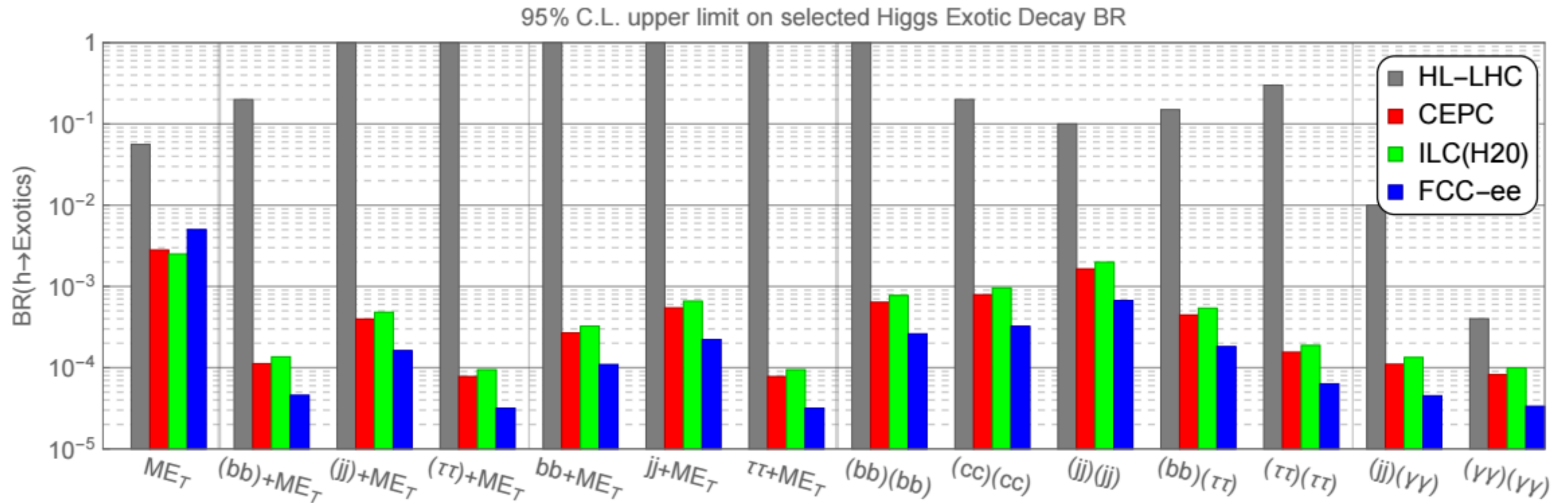
Decay Topologies	Decay mode \mathcal{F}_i	Decay Topologies	Decay mode \mathcal{F}_i
$h \rightarrow 2$	$h \rightarrow \cancel{E}_T$	$h \rightarrow 2 \rightarrow 4$	$h \rightarrow (b\bar{b})(b\bar{b})$
$h \rightarrow 2 \rightarrow 3$	$h \rightarrow \gamma + \cancel{E}_T$		$h \rightarrow (b\bar{b})(\tau^+\tau^-)$
	$h \rightarrow (b\bar{b}) + \cancel{E}_T$		$h \rightarrow (b\bar{b})(\mu^+\mu^-)$
	$h \rightarrow (jj) + \cancel{E}_T$		$h \rightarrow (\tau^+\tau^-)(\tau^+\tau^-)$
	$h \rightarrow (\tau^+\tau^-) + \cancel{E}_T$		$h \rightarrow (\tau^+\tau^-)(\mu^+\mu^-)$
	$h \rightarrow (\gamma\gamma) + \cancel{E}_T$		$h \rightarrow (jj)(jj)$
	$h \rightarrow (\ell^+\ell^-) + \cancel{E}_T$		$h \rightarrow (jj)(\gamma\gamma)$
			$h \rightarrow (jj)(\mu^+\mu^-)$
$h \rightarrow 2 \rightarrow 3 \rightarrow 4$	$h \rightarrow (b\bar{b}) + \cancel{E}_T$		$h \rightarrow (\ell^+\ell^-)(\ell^+\ell^-)$
	$h \rightarrow (jj) + \cancel{E}_T$		$h \rightarrow (\ell^+\ell^-)(\mu^+\mu^-)$
	$h \rightarrow (\tau^+\tau^-) + \cancel{E}_T$		$h \rightarrow (\mu^+\mu^-)(\mu^+\mu^-)$
	$h \rightarrow (\gamma\gamma) + \cancel{E}_T$		$h \rightarrow (\gamma\gamma)(\gamma\gamma)$
	$h \rightarrow (\ell^+\ell^-) + \cancel{E}_T$		$h \rightarrow \gamma\gamma + \cancel{E}_T$
	$h \rightarrow (\mu^+\mu^-) + \cancel{E}_T$		$h \rightarrow (\ell^+\ell^-)(\ell^+\ell^-) + \cancel{E}_T$
$h \rightarrow 2 \rightarrow (1+3)$	$h \rightarrow b\bar{b} + \cancel{E}_T$	$h \rightarrow 2 \rightarrow 4 \rightarrow 6$	$h \rightarrow (\ell^+\ell^-) + \cancel{E}_T + X$
	$h \rightarrow jj + \cancel{E}_T$		$h \rightarrow \ell^+\ell^-\ell^+\ell^- + \cancel{E}_T$
	$h \rightarrow \tau^+\tau^- + \cancel{E}_T$		$h \rightarrow \ell^+\ell^- + \cancel{E}_T + X$
	$h \rightarrow \gamma\gamma + \cancel{E}_T$	$h \rightarrow 2 \rightarrow 6$	
	$h \rightarrow \ell^+\ell^- + \cancel{E}_T$		

LHC's strength
Hard at LHC due to missing energy
Hard at LHC due to hadronic background
Lepton colliders' strength

Exotics/Long Lived Particles.

Z. Liu @ CEPC 2020

The Higgs could be a good portal to Dark Sector
— rich exotic signatures —



How to improve?

> Dedicated detectors, see e.g. talk by [R. Gonzalez Suarez @ FCC week 2021](#)