



Quarkonium polarization measurements

Challenges and opportunities

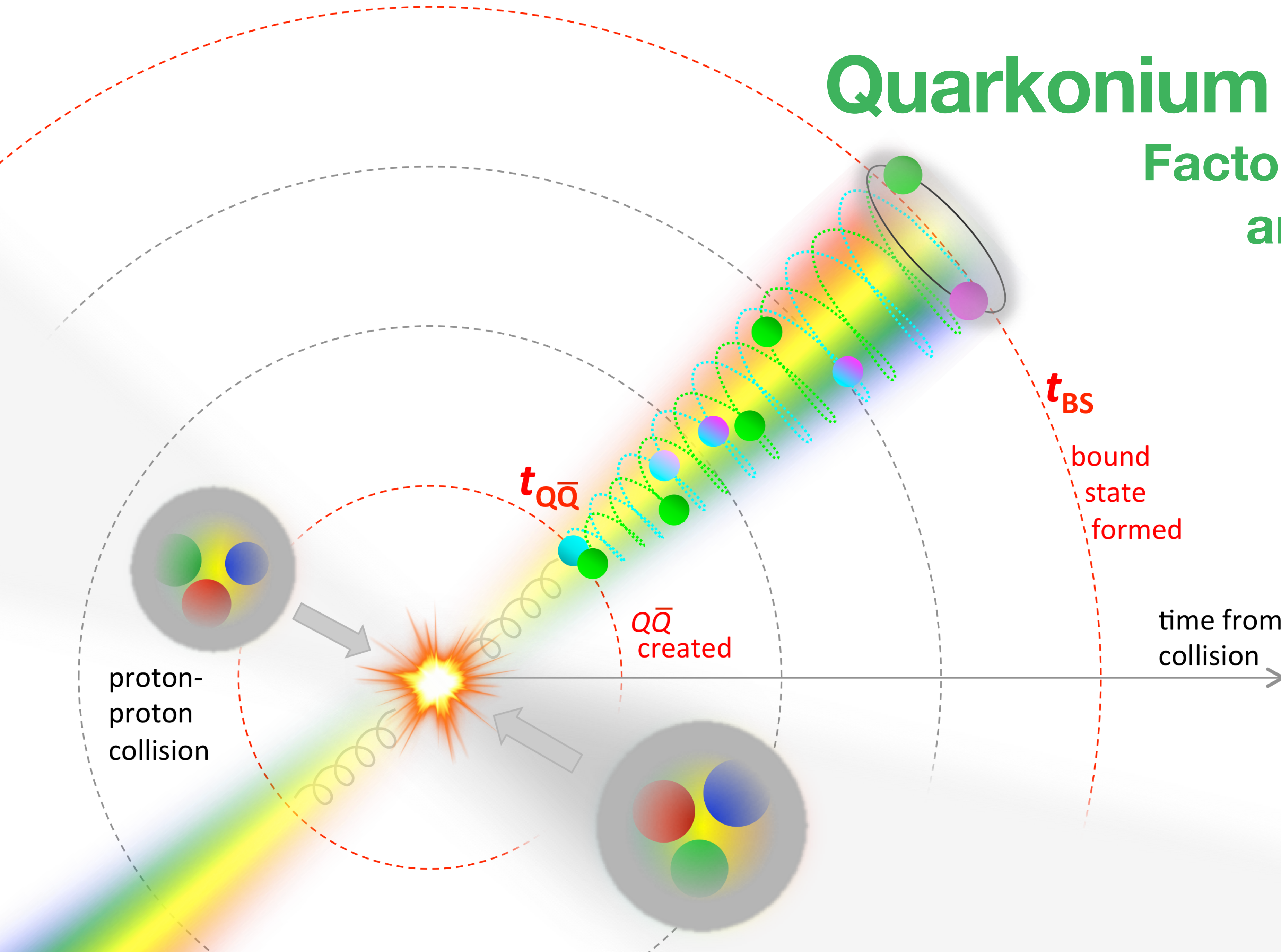
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Quarkonium production

Factorization of short- and long-distance effects



Quarkonium polarization

- Directly reflects mixture of pre-resonance configurations
- Is measured through the angular decay distribution wrt a quantization axis z

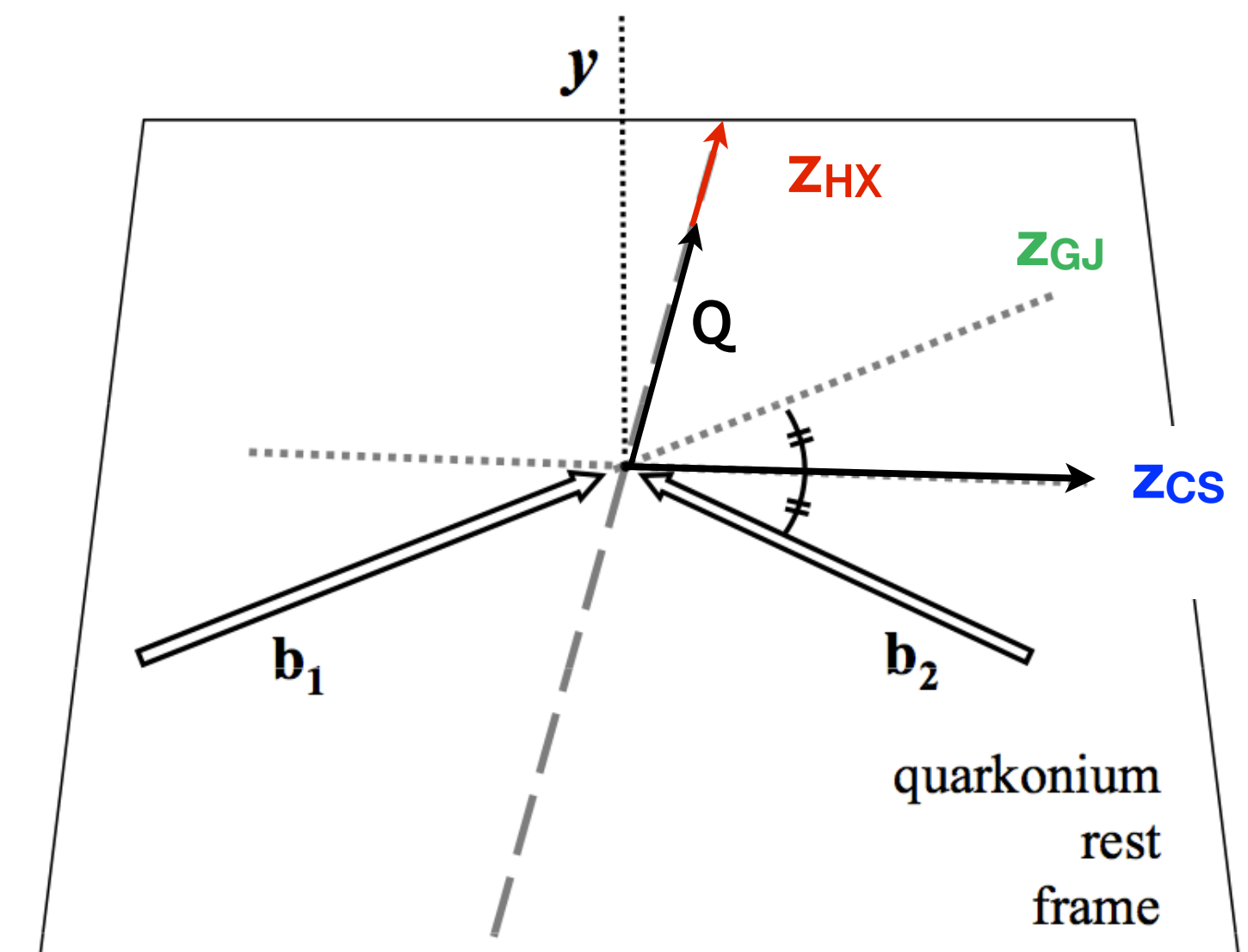
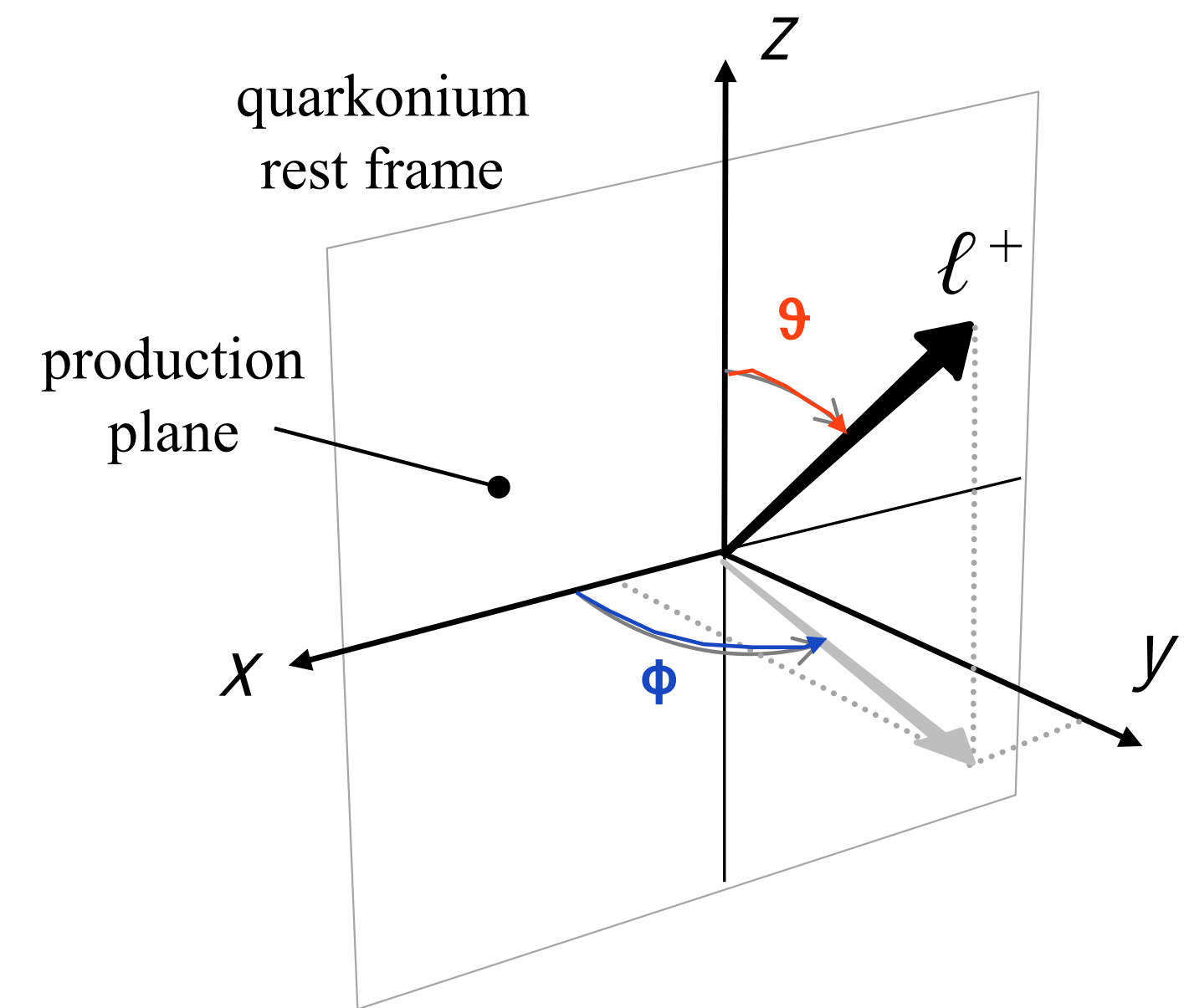
$$\frac{dN}{d \cos \vartheta d\varphi} \propto 1 + \lambda_{\vartheta} \cos^2 \vartheta + \lambda_{\varphi} \sin^2 \vartheta \cos 2\varphi + \lambda_{\vartheta\varphi} \sin 2\vartheta \cos \varphi$$

Reference frames:

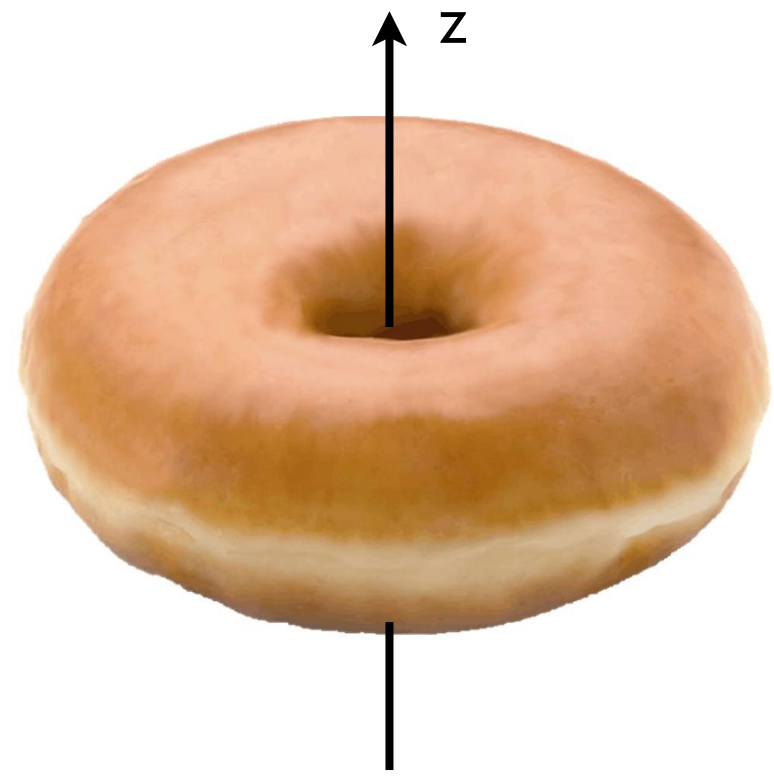
Center-of-mass helicity HX: \mathbf{z}_{HX} \approx direction of quarkonium momentum

Gottfried-Jackson GJ: \mathbf{z}_{GJ} = direction of one beam or the target

Collins-Soper CS: \mathbf{z}_{CS} \approx direction of relative velocity of colliding particles



Angular decay distributions



Longitudinal polarization
 $J_z = 0$

$$\lambda_{\vartheta} = -1$$
$$\lambda_{\varphi} = 0$$
$$\lambda_{\vartheta\varphi} = 0$$



No polarization/
isotropic decay

$$\lambda_{\vartheta} = 0$$
$$\lambda_{\varphi} = 0$$
$$\lambda_{\vartheta\varphi} = 0$$

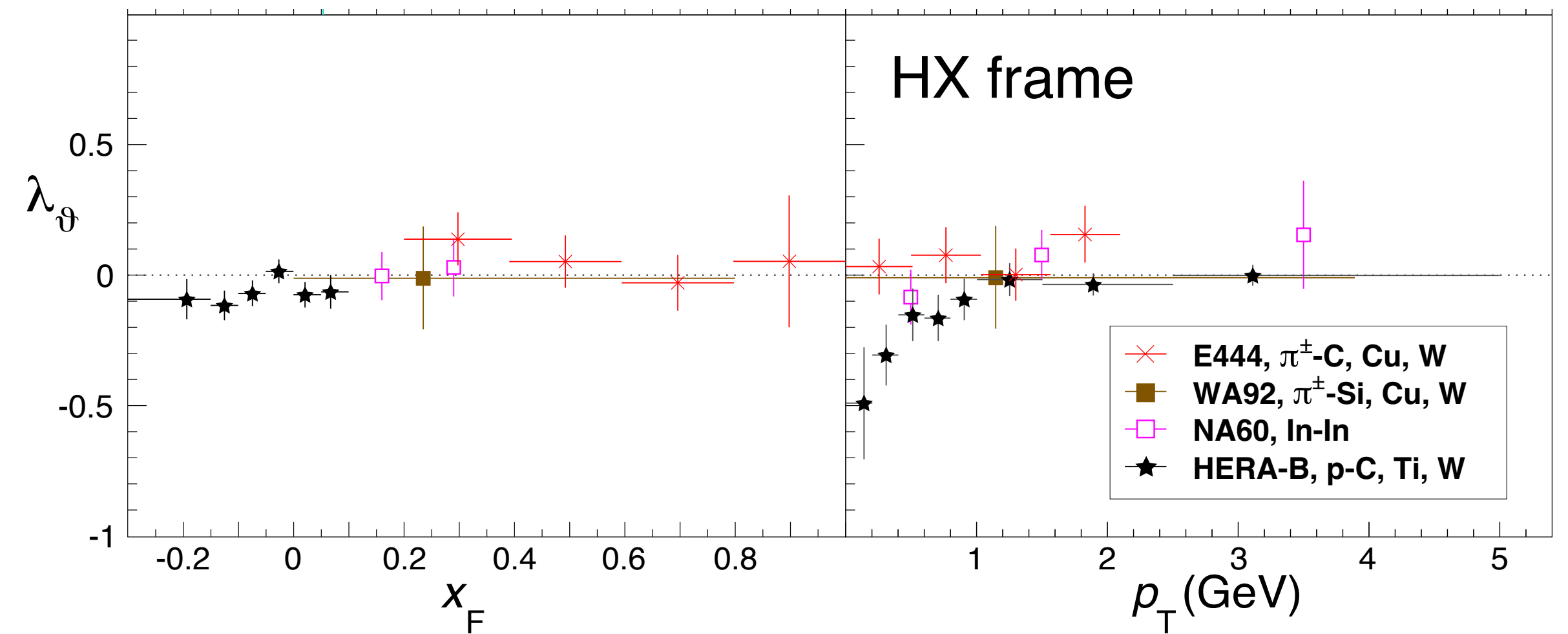
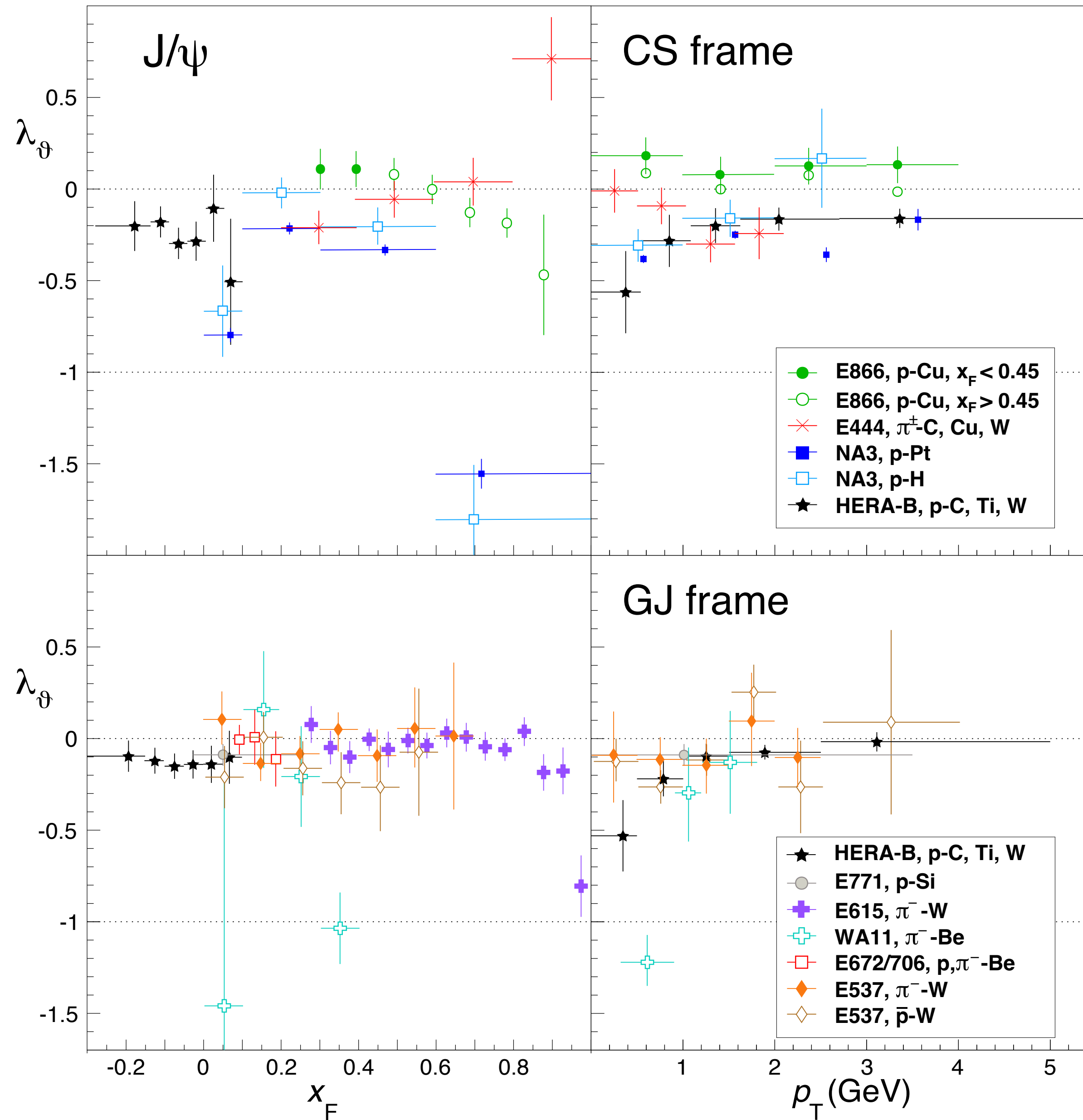


Transverse
polarization
 $J_z = \pm 1$

$$\lambda_{\vartheta} = +1$$
$$\lambda_{\varphi} = 0$$
$$\lambda_{\vartheta\varphi} = 0$$

The shape of the distribution is invariant and can be characterized by the frame-independent quantity $\tilde{\lambda} = (\lambda_{\vartheta} + 3\lambda_{\varphi}) / (1 - \lambda_{\varphi})$

Measurements in fixed-target experiments with several beam particles, target nuclei and collision energies



Indications from existing measurements

1. Hierarchy in λ_θ and λ_ϕ parameters:

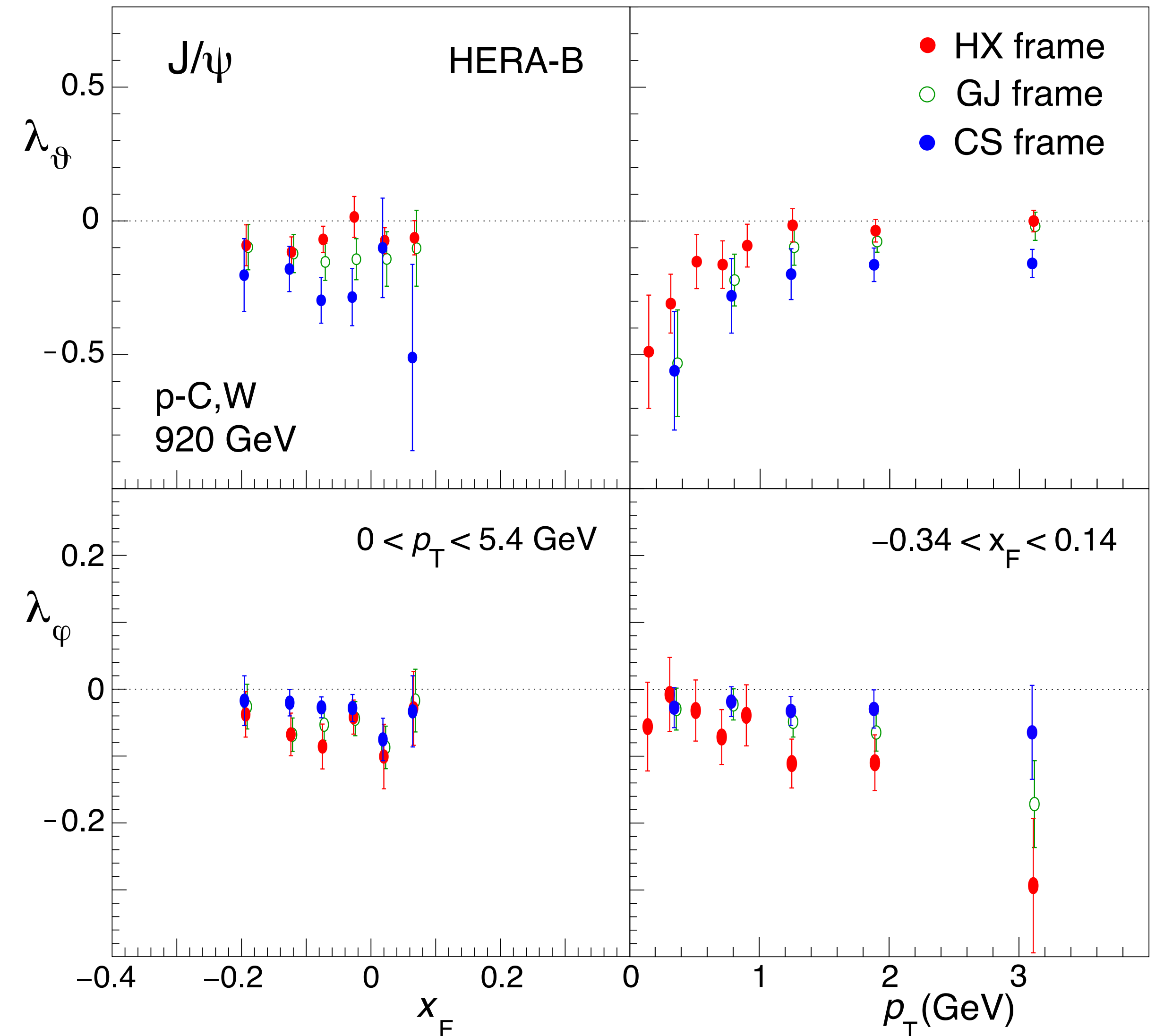
CS - GJ - HX

➔ CS axis more naturally reflects the alignment of the J/ψ angular momentum

➔ Quarkonium production is dominated by 2-to-1 processes, where the produced state is strongly polarized:

- $gg \rightarrow Q$ fully longitudinally polarized
- $q\bar{q} \rightarrow Q$ fully transversely polarized

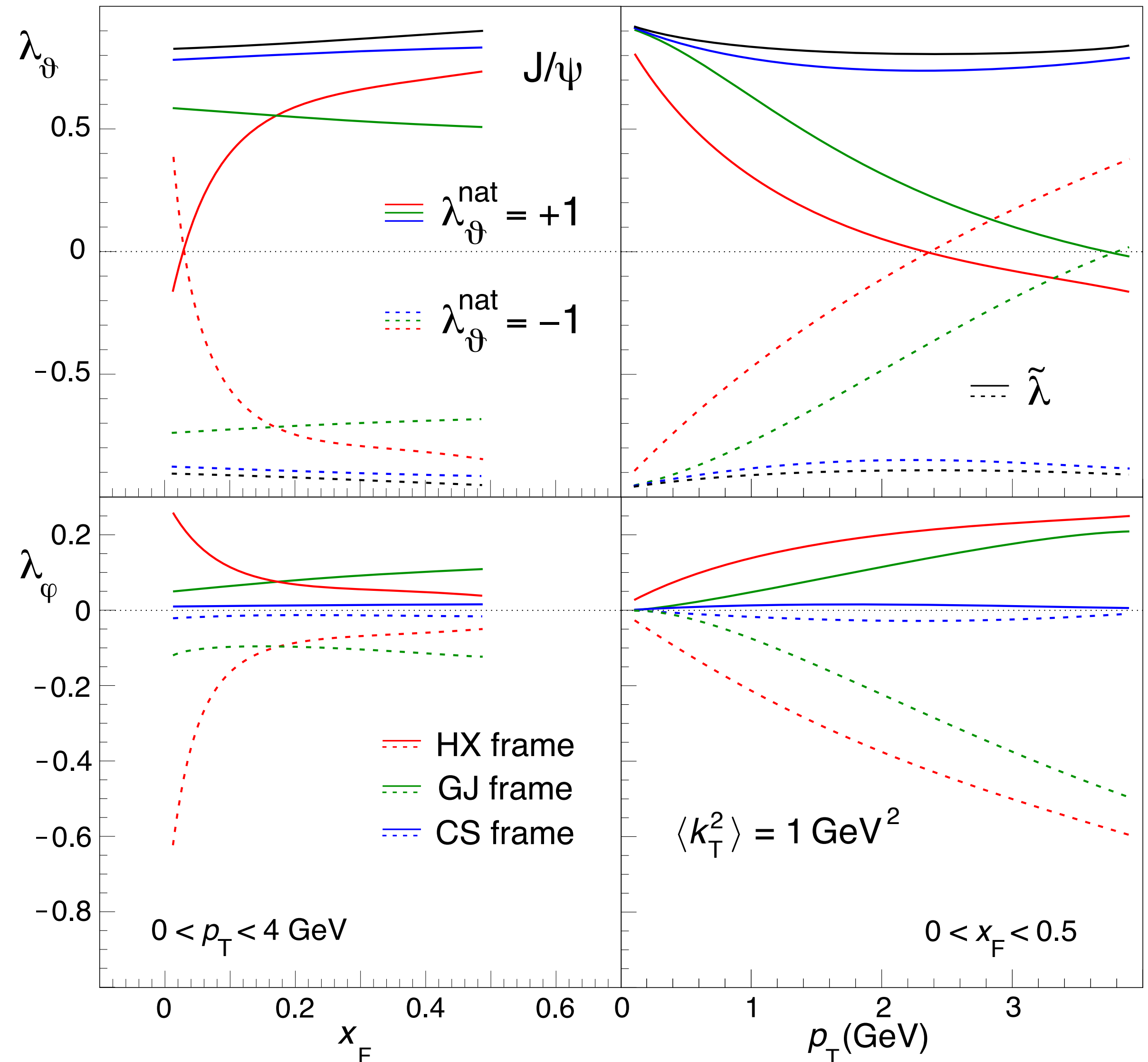
for directly produced quarkonia along the natural polarization axis
(PRD 83 (2011) 056008)



Indications from existing measurements

2. Decreasing magnitude of polarization with increasing p_T

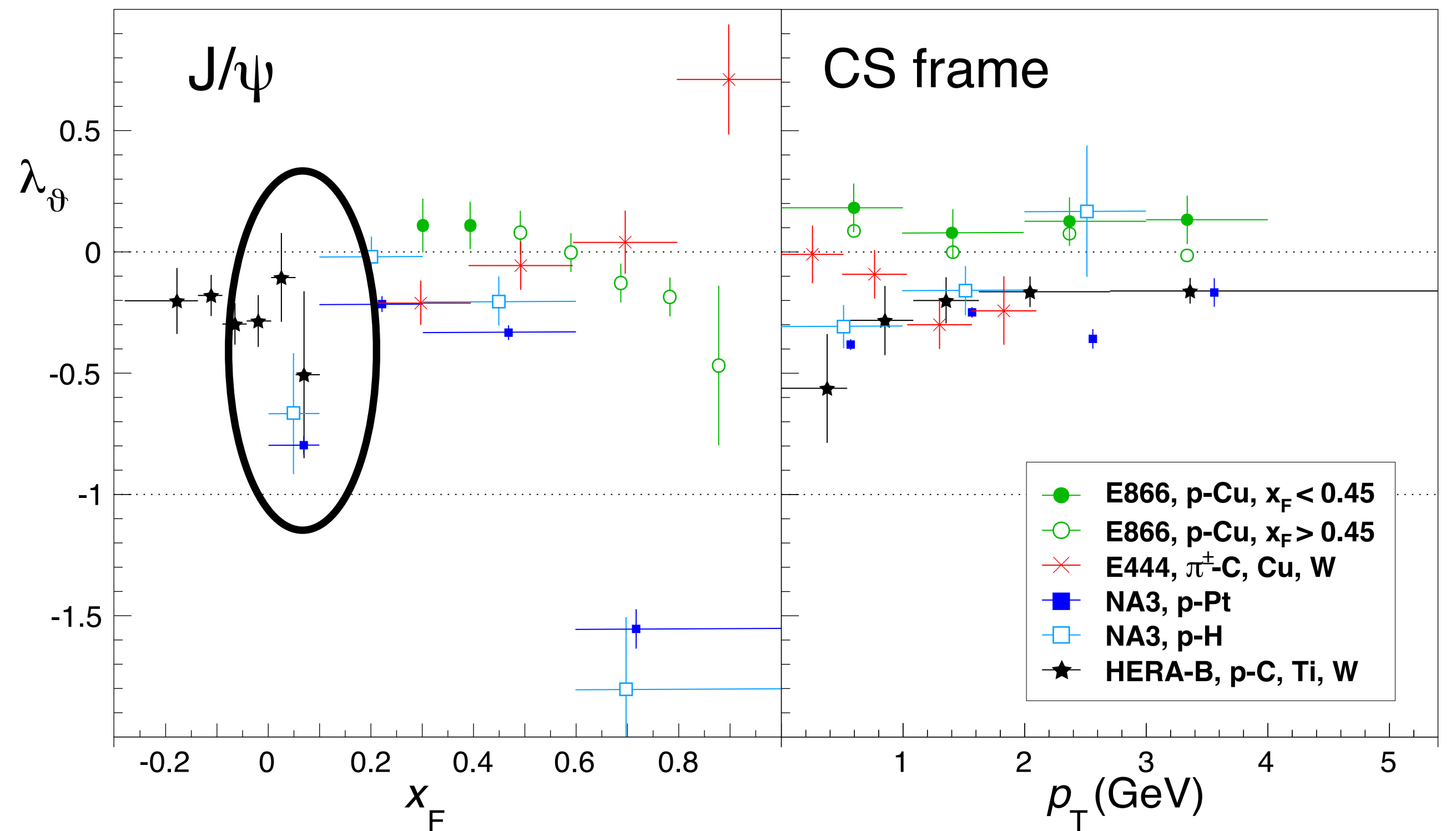
➔ Nonzero transverse momentum distribution of colliding partons has an effect for light quarkonia at low p_T and small x_F



Indications from existing measurements

3. More longitudinal J/ψ polarization at small x_F

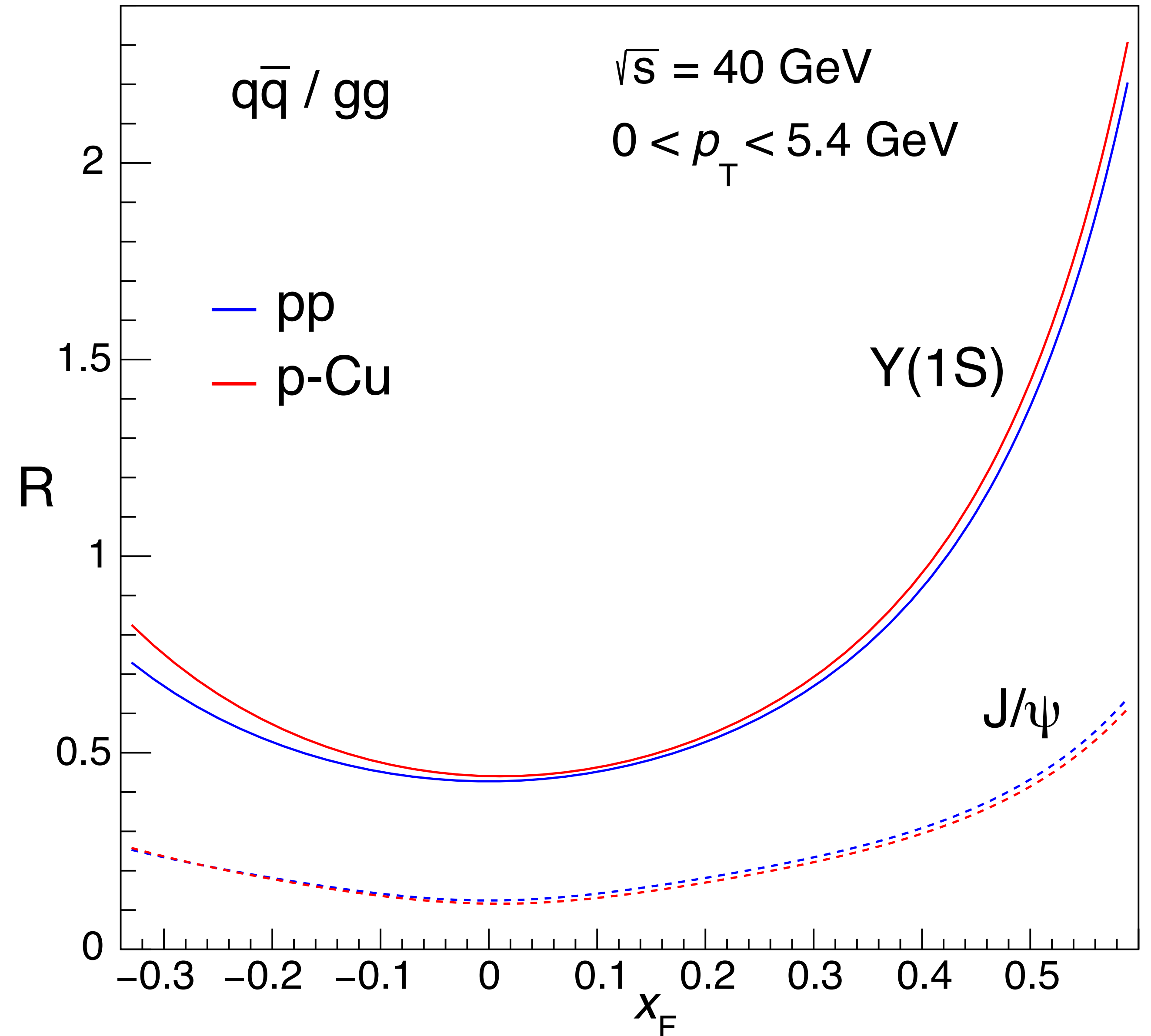
➔ Relative dominance of gg fusion compared to $q\bar{q}$ annihilation at x_F close to 0



$q\bar{q} / gg$ ratios

pp and p-nucleus collisions

- $q\bar{q}$ and gg parton densities computed as the product of the corresponding PDFs using CT14NLO from LHAPDF and EPPS16
- Minimum of ratio around $x_F = 0$
- Nuclear effects are negligible
- $q\bar{q}$ annihilation is more important for heavier quarkonia



Empirical model

Assumptions

- Observed polarization of directly produced quarkonia results from the interplay between $q\bar{q}$ annihilation and gg fusion processes.
- Observable mixture of longitudinal (from gg) and transverse (from $q\bar{q}$) polarizations is fully determined by the product of two ratios:
 1. ratio between $q\bar{q}$ and gg parton densities, R
 2. ratio between $q\bar{q}$ and gg partonic cross sections, r
- Natural polarization parameter in parton-parton CS frame

$$\lambda = \frac{f_{q\bar{q}} \lambda_{\vartheta}^{q\bar{q}} / (3 + \lambda_{\vartheta}^{q\bar{q}}) + f_{gg} \lambda_{\vartheta}^{gg} / (3 + \lambda_{\vartheta}^{gg})}{f_{q\bar{q}} / (3 + \lambda_{\vartheta}^{q\bar{q}}) + f_{gg} / (3 + \lambda_{\vartheta}^{gg})}$$

according to sum rule in
EPJC 69 (2010) 657

$$f_{q\bar{q}} = R \times r / (1 + R \times r) \quad f_{gg} = 1 / (1 + R \times r)$$

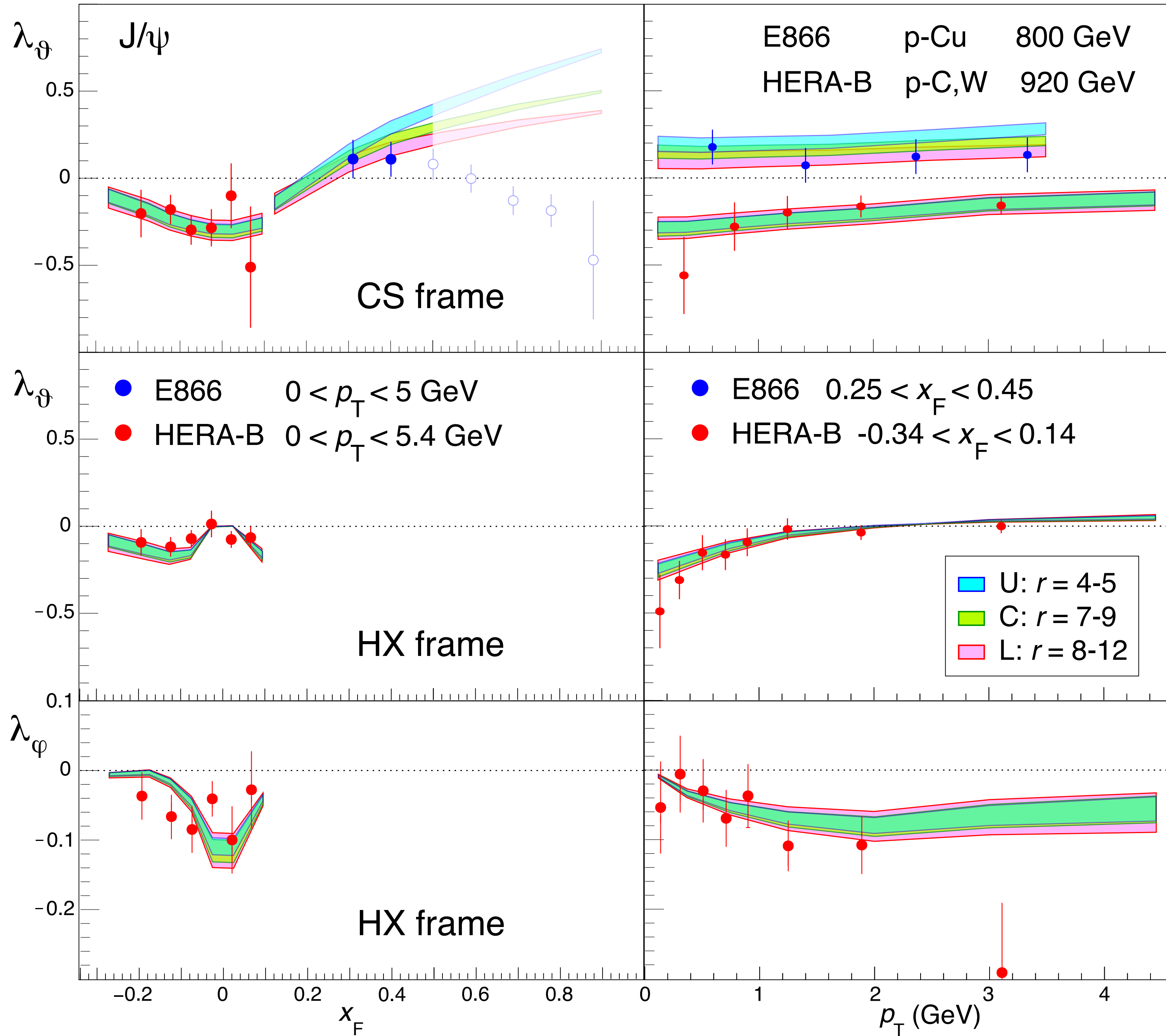
Feed-down from heavier quarkonium states

- S-wave states have the same polarization: $\psi(2S) \rightarrow J/\psi\pi\pi$ (PRD 62 (2000) 032002)
 - P-wave states have different production mechanism due to emission of a transversely polarized gluon (PRD 83 (2011) 096001)
- ➔ Weaker polarization due to mixture from feed-down
- Total feed-down fraction from χ_c : 19% from HERAb (PRD 79 (2009) 012001)

		$\lambda_{\vartheta}^{\chi_1}$	$\lambda_{\vartheta}^{\chi_2}$
central, C	gg	+1	-3/5
	$q\bar{q}$	-1/3	-1/3
lower, L	$gg, q\bar{q}$	-1/3	-3/5
upper, U	$gg, q\bar{q}$	+1	+1

Results

p-nucleus collisions



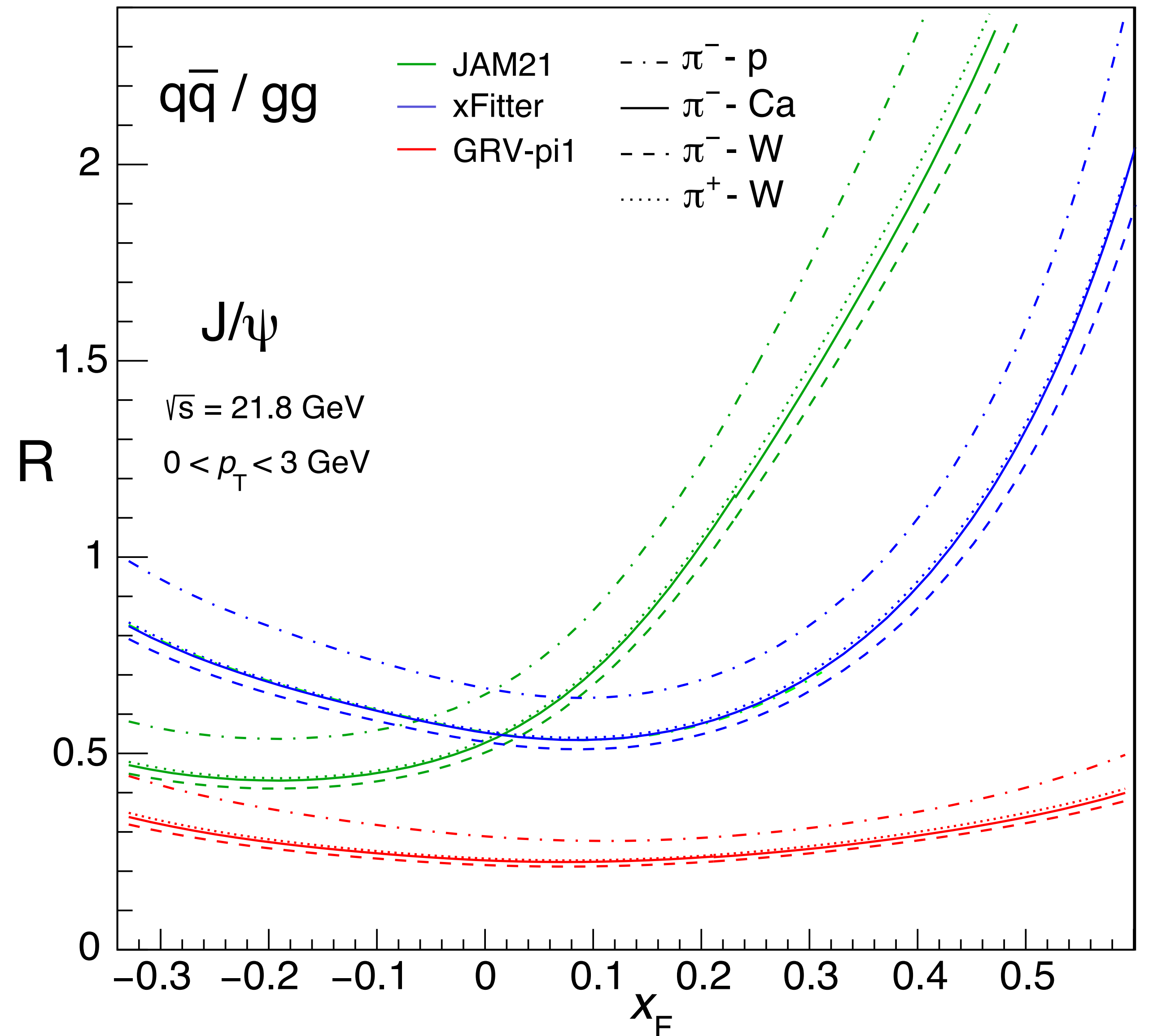
r = ratio between $q\bar{q}$ and gg partonic cross section

$q\bar{q} / gg$ ratio

pion-nucleus collisions

- Significant differences for various pion PDFs because of poorly known gluon densities
- Negligible differences between positive and negative pions
- Nuclear effects have minor impact on x_F dependence

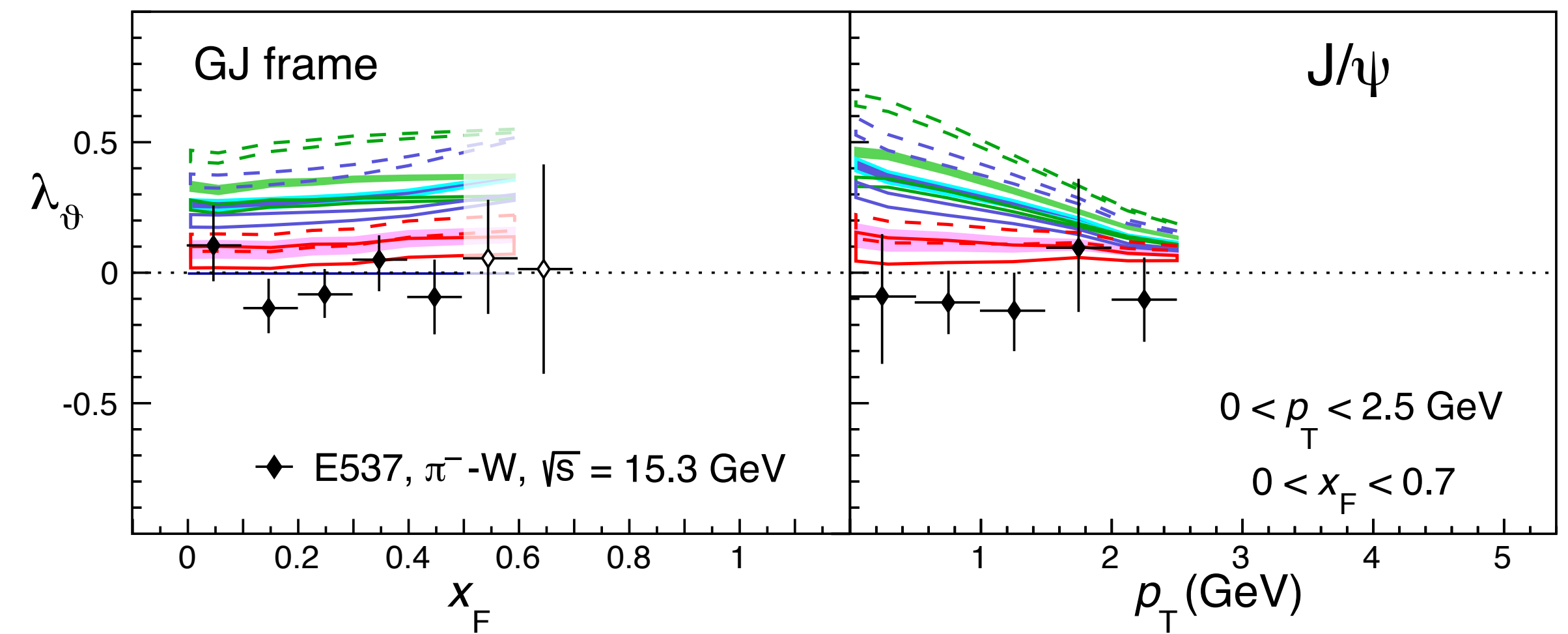
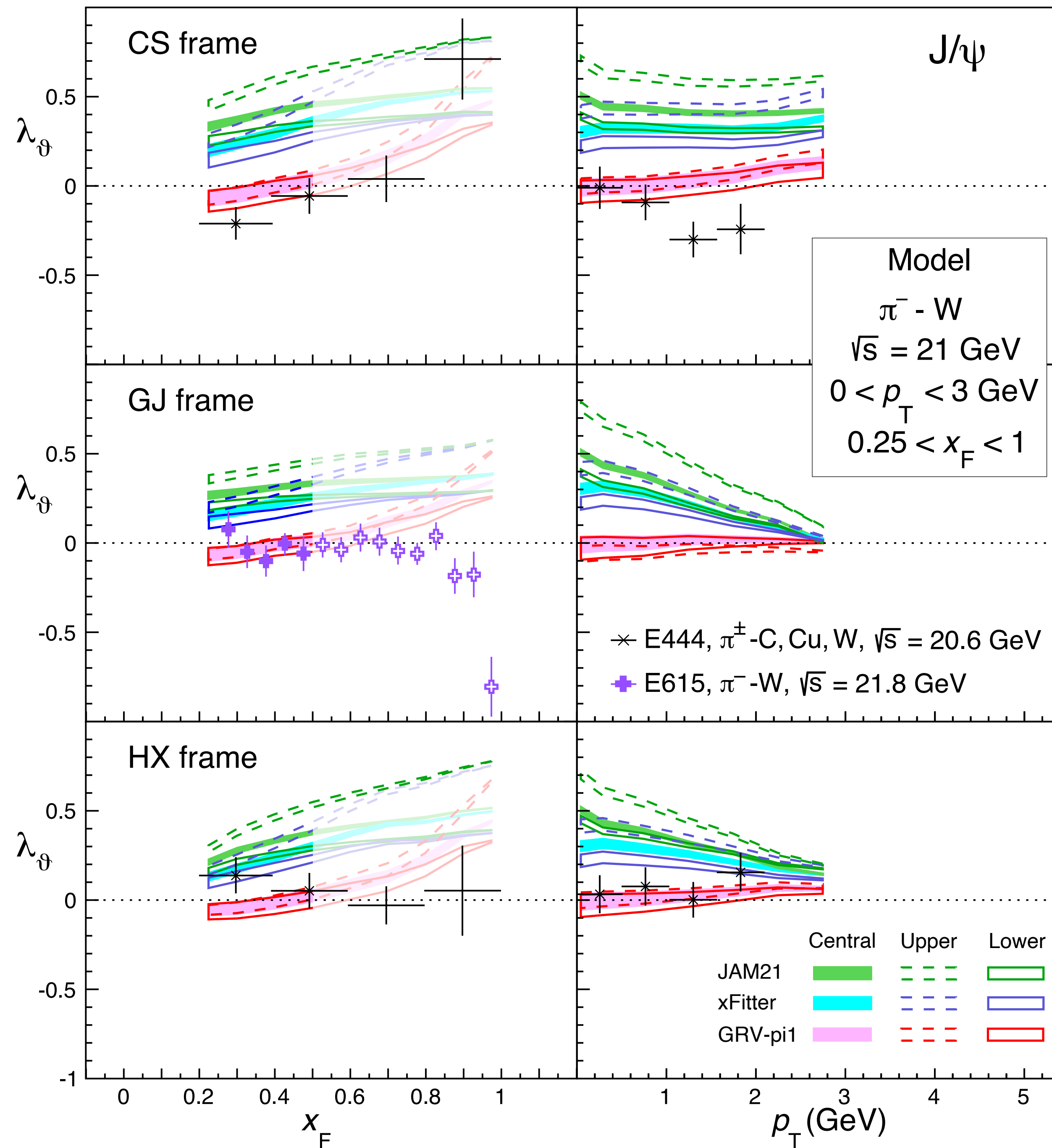
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Results

Pion-nucleus collisions

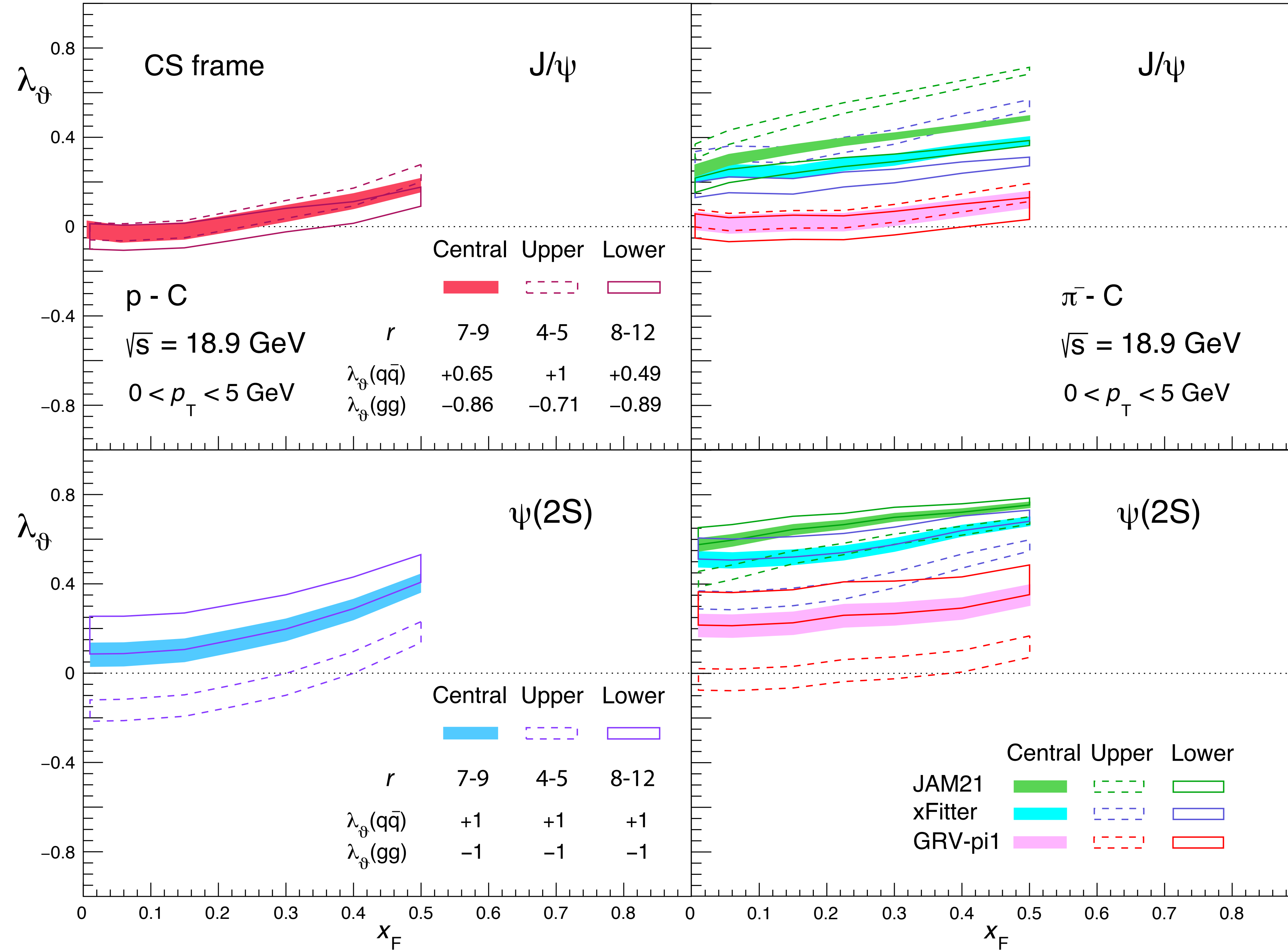
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Predictions



Apparatus for Meson and Baryon
Experimental Research



Conclusions and summary

- Our simple model assumes that observed polarization results from the interplay between $q\bar{q}$ annihilation and gg fusion.
- Future polarization measurements in proton-nucleus collisions can test our model.
- The polarization observable has the potential to provide a strong constraint on the pion PDFs.
- $\psi(2S)$ polarization measurements are particularly interesting since there is no feed-down.

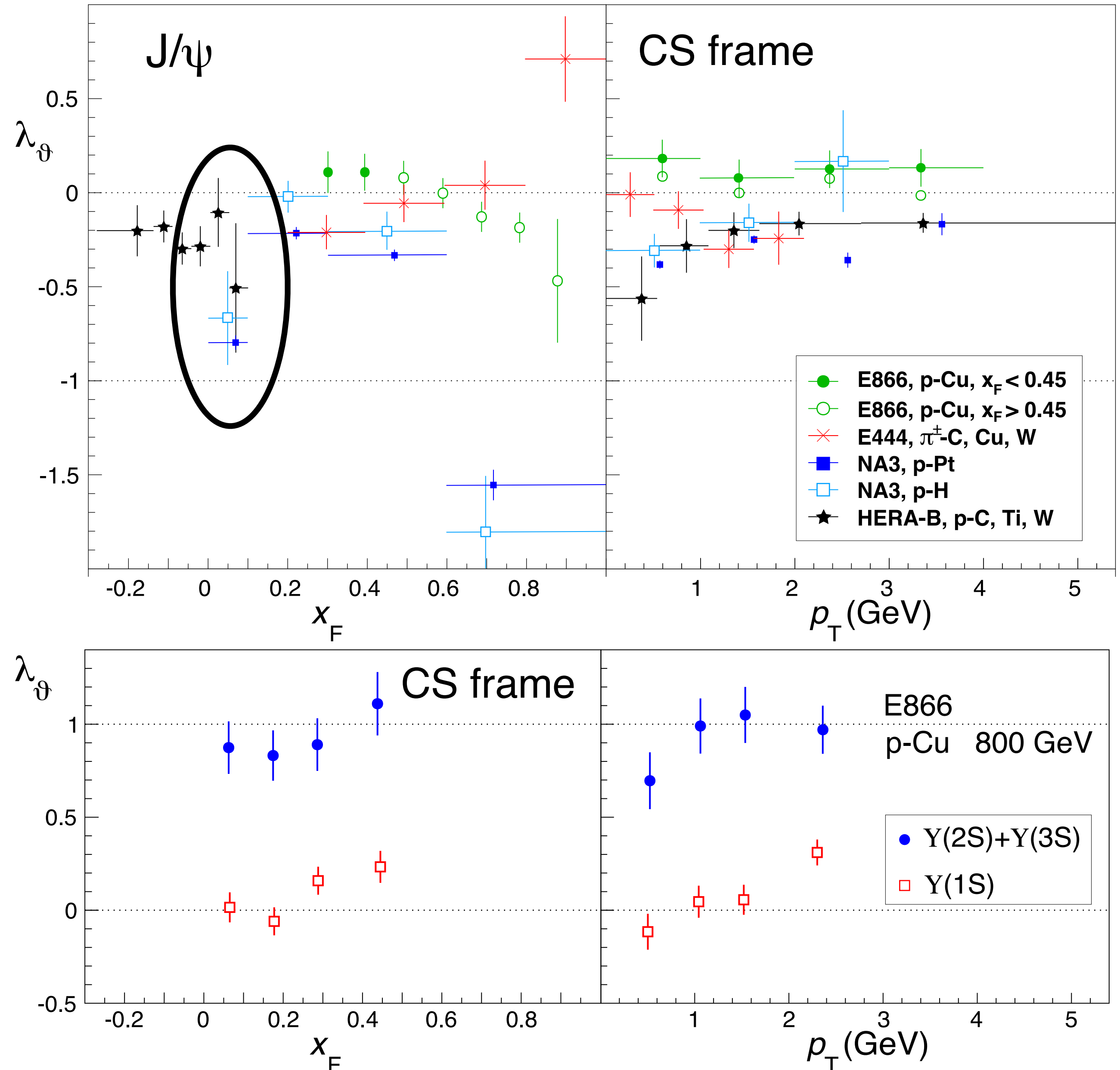
Further reading: “Particle Polarization in High Energy Physics”
<https://link.springer.com/book/10.1007/978-3-031-08876-6>

BACKUP

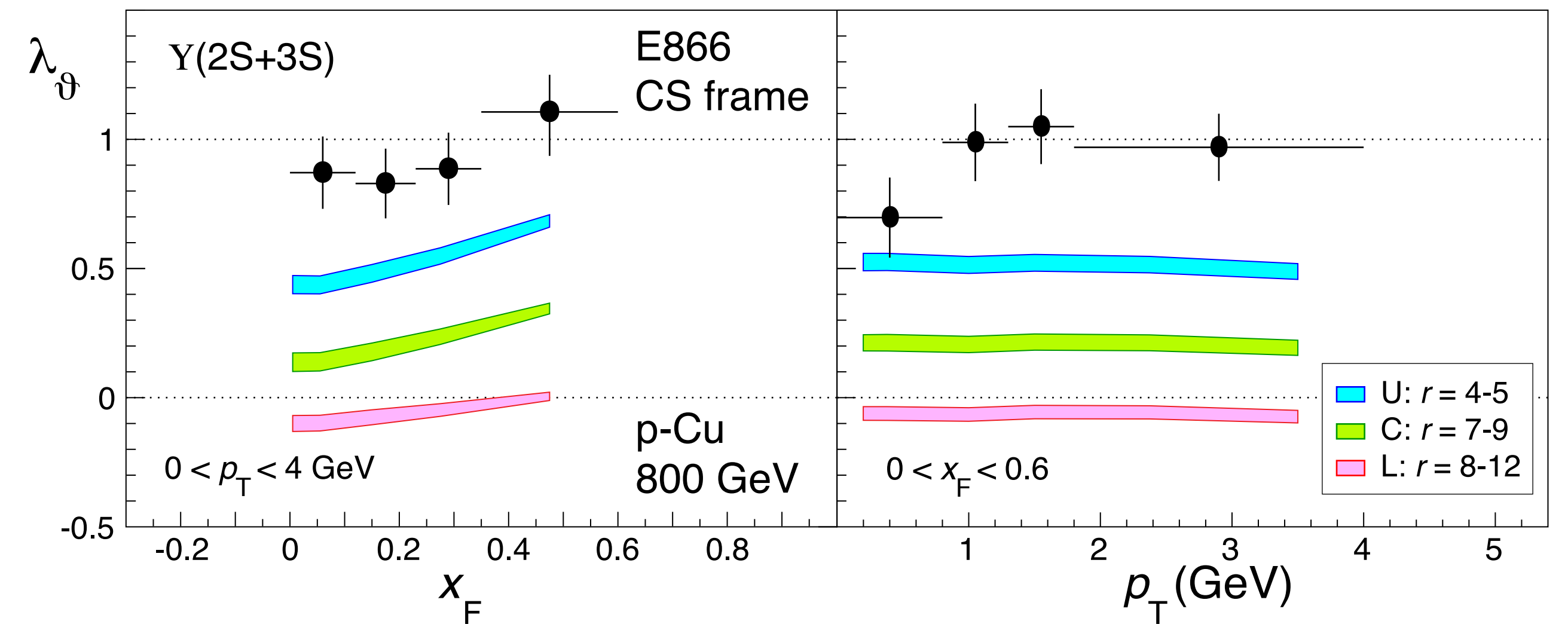
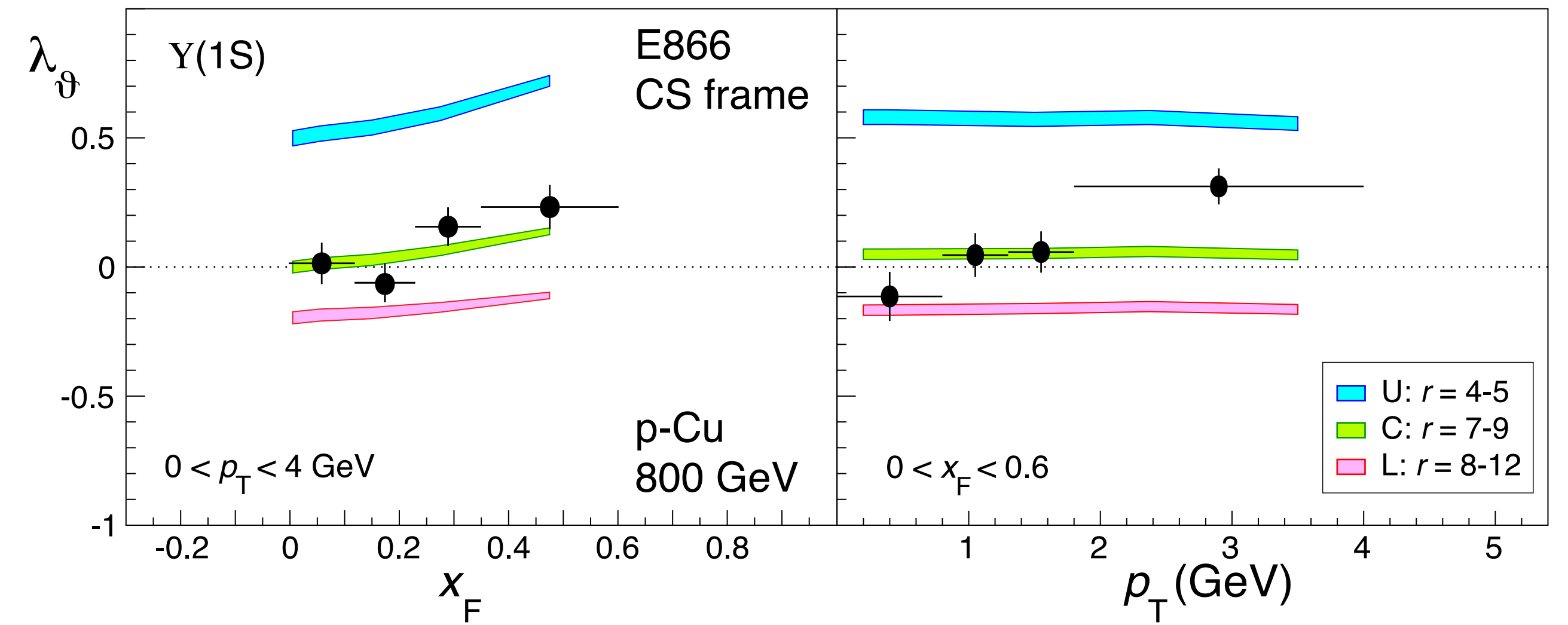
Indications from existing measurements

2. More longitudinal J/ ψ polarization at small x_F

➔ Relative dominance of gluon-gluon fusion at mid-rapidity, qq annihilation is more relevant in more forward region



Results



Constraints

