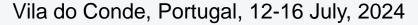


Beyond the Standard Model BrainStorming Meeting:

Particle Physics and Cosmology interface





Axion paradigm with color-mediated neutrino masses

Henrique Brito Câmara

henrique.b.camara@tecnico.ulisboa.pt CFTP/IST, U. Lisbon

In collaboration with: A. Batra, F.R. Joaquim,

R. Srivastava, J.W.F. Valle

arXiv: 2309.06473 [hep-ph]

Phys.Rev.Lett. 132 (2024) 5, 051801













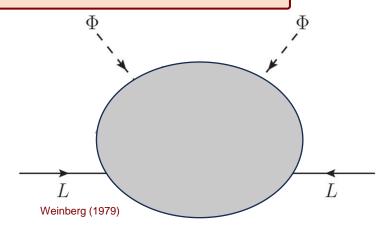
The Standard Model cannot explain:

- Neutrino flavour oscillations which imply massive neutrinos and lepton mixing;
- Observed dark matter abundance;
- Strong CP problem: Lack of a theoretical explanation for the non-observation of the neutron electric dipole moment which indicates that strong interactions preserve CP symmetry.

The Standard Model cannot explain:

- Neutrino flavour oscillations which imply massive neutrinos and lepton mixing;
- Observed dark matter abundance;
- Strong CP problem: Lack of a theoretical explanation for the non-observation of the neutron electric dipole moment which indicates that strong interactions preserve CP symmetry.

Majorana Neutrino masses



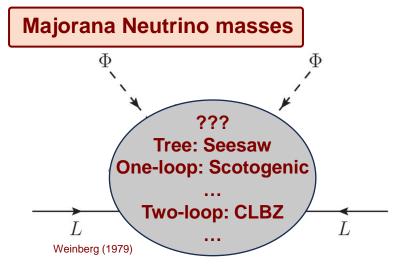
Weinberg (1979)

The Standard Model cannot explain:

- Neutrino flavour oscillations which imply massive neutrinos and lepton mixing;
- Observed dark matter abundance;
- Strong CP problem: Lack of a theoretical explanation for the non-observation of the neutron electric dipole moment which indicates that strong interactions preserve CP symmetry.

The Standard Model cannot explain:

- Neutrino flavour oscillations which imply massive neutrinos and lepton mixing;
- Observed dark matter abundance;
- Strong CP problem: Lack of a theoretical explanation for the non-observation of the neutron electric dipole moment which indicates that strong interactions preserve CP symmetry.





Peccei, Quinn (1977), Weinberg (1978), Wilczek (1978)

$$U(1)_{PQ}$$

$$\sigma = \frac{v_{\sigma} + \rho}{\sqrt{2}} e^{ia_{\sigma}/v_{\sigma}}$$

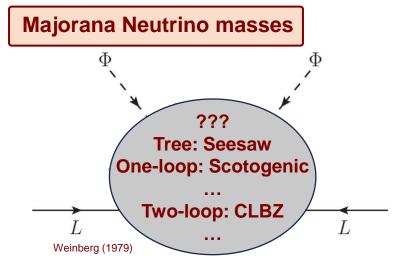
Vector-like quark

KSVZ (1979, 1980)

$$\Psi_{L,R}$$

The Standard Model cannot explain:

- Neutrino flavour oscillations which imply massive neutrinos and lepton mixing;
- Observed dark matter abundance;
- Strong CP problem: Lack of a theoretical explanation for the non-observation of the neutron electric dipole moment which indicates that strong interactions preserve CP symmetry.



Axion

Peccei, Quinn (1977), Weinberg (1978), Wilczek (1978)

$$U(1)_{PQ}$$

$$\sigma = \frac{v_{\sigma} + \rho}{\sqrt{2}} e^{ia_{\sigma}/v_{\sigma}}$$

Vector-like quark

KSVZ (1979, 1980)

$$\Psi_{L,R}$$

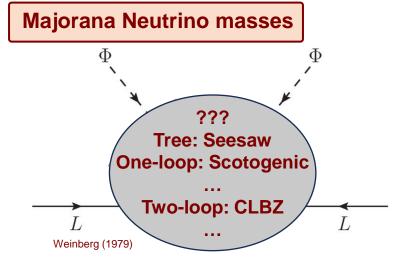
Dynamical solution

$$\left(\frac{a}{f_a} - \bar{\theta}\right) \frac{\alpha_s}{8\pi} G\tilde{G}$$

$$\overline{\theta}_{\text{eff}} = \overline{\theta} + \left\langle \frac{a}{f_a} \right\rangle = 0$$

The Standard Model cannot explain:

- Neutrino flavour oscillations which imply massive neutrinos and lepton mixing;
- Observed dark matter abundance;
- Strong CP problem: Lack of a theoretical explanation for the non-observation of the neutron electric dipole moment which indicates that strong interactions preserve CP symmetry.



Axion

Peccei, Quinn (1977), Weinberg (1978), Wilczek (1978)

$$U(1)_{PQ}$$

$$\sigma = \frac{v_{\sigma} + \rho}{\sqrt{2}} e^{ia_{\sigma}/v_{\sigma}}$$

Vector-like quark

$$\Psi_{L,R}$$

Dynamical solution

$$\left(\frac{a}{f_a} - \bar{\theta}\right) \frac{\alpha_s}{8\pi} G\tilde{G}$$

$$\overline{\theta}_{\text{eff}} = \overline{\theta} + \left\langle \frac{a}{f_a} \right\rangle = 0$$

Our approach:

New class of models where **neutrino masses** are **radiatively generated by colored particles** which **simultaneously** solve through the PQ mechanism the **strong CP problem.** The predicted **axion** particle accounts for **dark matter**.

Fields	$SU(3)_c \otimes SU(2)_L \otimes U(1)_Y$	$U(1)_{PQ}$	Multiplicity
$\overline{\Psi_L}$	$[(p,q),2n\pm 1,0]$	ω	n_{Ψ}
Ψ_R	$[(p,q), 2n \pm 1, 0]$	0	n_{Ψ}
σ	(1, 1, 0)	ω	1
η	[(p,q), 2n, 1/2]	0	n_{η}
χ	$[(p,q),2n\pm 1,0]$	0	n_{χ}

Vector-like q	uarks
---------------	-------

Complex scalar singlet

Colored scalars

Fields	$SU(3)_c \otimes SU(2)_L \otimes U(1)_Y$	$U(1)_{PQ}$	Multiplicity
$\overline{\Psi_L}$	$[(p,q), 2n \pm 1, 0]$	ω	n_{Ψ}
Ψ_R	$[(p,q), 2n \pm 1, 0]$	0	n_{Ψ}
σ	(1, 1, 0)	ω	1
η	[(p,q), 2n, 1/2]	0	n_{η}
χ	$[(p,q),2n\pm 1,0]$	0	n_χ

Vector-like quarks

Complex scalar singlet

Colored scalars

Fields	$SU(3)_c \otimes SU(2)_L \otimes U(1)_Y$	$U(1)_{PQ}$	Multiplicity
$\overline{\Psi_L}$	$[(p,q),2n\pm 1,0]$	ω	n_{Ψ}
Ψ_R	$[(p,q), 2n \pm 1, 0]$	0	n_{Ψ}
σ	(1, 1, 0)	ω	1
η	[(p,q), 2n, 1/2]	0	n_{η}
χ	$[(p,q),2n\pm 1,0]$	0	n_{χ}

Yukawa Lagrangian

$$-\mathcal{L}_{\mathrm{Yuk.}} \supset \mathbf{Y}_{\Psi}\overline{\Psi_{L}}\Psi_{R}\sigma + \frac{1}{2}\mathbf{Y}_{\chi_{j}}\Psi_{R}^{T}C\chi_{j}\Psi_{R} + \mathbf{Y}_{i}\bar{L}\eta_{i}^{*}\Psi_{R} + \mathrm{H.c.}$$

Vector-like qu	ıarks
----------------	-------

Complex scalar singlet

Colored scalars

Fields	$SU(3)_c \otimes SU(2)_L \otimes U(1)_Y$	$U(1)_{PQ}$	Multiplicity
$\left(egin{array}{c} \Psi_L \ \Psi_R \end{array} ight)$	$[(p,q), 2n \pm 1, 0]$	ω	n_{Ψ}
Ψ_R	$[(p,q), 2n \pm 1, 0]$	0	n_{Ψ}
$t \sigma$	(1, 1, 0)	ω	1
η	[(p,q), 2n, 1/2]	0	n_{η}
χ	$[(p,q),2n\pm 1,0]$	0	n_χ

Strong CP problem

Yukawa Lagrangian

$$-\mathcal{L}_{\text{Yuk.}} \supset \mathbf{Y}_{\Psi}\overline{\Psi_{L}}\Psi_{R}\sigma + \frac{1}{2}\mathbf{Y}_{\chi_{j}}\Psi_{R}^{T}C\chi_{j}\Psi_{R} + \mathbf{Y}_{i}\bar{L}\eta_{i}^{*}\Psi_{R} + \text{H.c.}$$

Scalar Potential

$$V \supset \mu_{ijk}\chi_i\chi_j\chi_k + \kappa_{ij}\eta_i^{\dagger}\Phi\chi_j + \lambda_{ijk}\Phi^{\dagger}\eta_i\chi_j\chi_k + \text{H.c.}$$

Vector-like	quarks
--------------------	--------

Complex scalar singlet

Colored scalars

Fields	$SU(3)_c \otimes SU(2)_L \otimes U(1)_Y$	$U(1)_{PQ}$	Multiplicity
$\left(\begin{array}{c} \Psi_L \\ \Psi_R \end{array}\right)$	$[(p,q), 2n \pm 1, 0]$	ω	n_{Ψ}
Ψ_R	$[(p,q), 2n \pm 1, 0]$	0	n_{Ψ}
σ	(1, 1, 0)	ω	1
η	[(p,q), 2n, 1/2]	0	n_{η}
χ	$[(p,q),2n\pm 1,0]$	0	n_χ

Strong CP problem

Yukawa Lagrangian

$$-\mathcal{L}_{\text{Yuk.}} \supset \mathbf{Y}_{\Psi} \overline{\Psi_{L}} \Psi_{R} \sigma + \frac{1}{2} \mathbf{Y}_{\chi_{j}} \Psi_{R}^{T} C \chi_{j} \Psi_{R} + \mathbf{Y}_{i} \bar{L} \eta_{i}^{*} \Psi_{R} + \text{H.c.}$$

Scalar Potential

$$V \supset \mu_{ijk}\chi_i\chi_j\chi_k + \kappa_{ij}\eta_i^{\dagger}\Phi\chi_j + \lambda_{ijk}\Phi^{\dagger}\eta_i\chi_j\chi_k + \text{H.c.}$$

Axion decay constant

$$\sigma = \frac{v_{\sigma} + \rho}{\sqrt{2}} \underbrace{v_{\sigma}/v_{\sigma}}$$

$$f_a = \frac{f_{PQ}}{N} = \frac{v_\sigma}{\sqrt{2}N}$$

Vector-like quarks

Complex scalar singlet

Colored scalars

Fields	$SU(3)_c \otimes SU(2)_L \otimes U(1)_Y$	$U(1)_{PQ}$	Multiplicity
$\left(egin{array}{c} \Psi_L \ \Psi_R \ \sigma \end{array} ight)$	$[(p,q), 2n \pm 1, 0]$	ω	n_{Ψ}
Ψ_R	$[(p,q), 2n \pm 1, 0]$	0	n_{Ψ}
σ	(1, 1, 0)	ω	1
η	[(p,q), 2n, 1/2]	0	n_{η}
χ	$[(p,q),2n\pm 1,0]$	0	n_χ

Strong CP problem

Yukawa Lagrangian

$$-\mathcal{L}_{\mathrm{Yuk.}} \supset \mathbf{Y}_{\Psi} \overline{\Psi_{L}} \Psi_{R} \sigma + \frac{1}{2} \mathbf{Y}_{\chi_{j}} \Psi_{R}^{T} C \chi_{j} \Psi_{R} + \mathbf{Y}_{i} \bar{L} \eta_{i}^{*} \Psi_{R} + \mathrm{H.c.}$$

Scalar Potential

$$V \supset \mu_{ijk}\chi_i\chi_j\chi_k + \kappa_{ij}\eta_i^{\dagger}\Phi\chi_j + \lambda_{ijk}\Phi^{\dagger}\eta_i\chi_j\chi_k + \text{H.c.}$$

Axion decay constant

$$\sigma = \frac{v_{\sigma} + \rho}{\sqrt{2}} e^{ia_{\sigma}/v_{\sigma}}$$

$$f_a = \frac{f_{PQ}}{N} = \frac{v_\sigma}{\sqrt{2}N}$$

Color-anomaly factor

$$N = 2n_{\Psi}\omega(2n \pm 1)T(p,q)$$

vector-line dual na	Vecto	or-like	quarks
---------------------	-------	---------	--------

Complex scalar singlet

Colored scalars

Fields	$SU(3)_c \otimes SU(2)_L \otimes U(1)_Y$	$U(1)_{PQ}$	Multiplicity
$\left(\begin{array}{c} \Psi_L \\ \Psi_R \end{array}\right)$	$[(p,q), 2n \pm 1, 0]$	ω	n_{Ψ}
Ψ_R	$[(p,q), 2n \pm 1, 0]$	0	n_{Ψ}
$t \sigma$	(1, 1, 0)	ω	1
η	[(p,q), 2n, 1/2]	0	n_{η}
χ	$[(p,q),2n\pm 1,0]$	0	n_{χ}

Strong CP problem

Yukawa Lagrangian

$$-\mathcal{L}_{\mathrm{Yuk.}} \supset \mathbf{Y}_{\Psi} \overline{\Psi_{L}} \Psi_{R} \sigma + \frac{1}{2} \mathbf{Y}_{\chi_{j}} \Psi_{R}^{T} C \chi_{j} \Psi_{R} + \mathbf{Y}_{i} \bar{L} \eta_{i}^{*} \Psi_{R} + \mathrm{H.c.}$$

Scalar Potential

$$V \supset \mu_{ijk} \chi_i \chi_j \chi_k + \kappa_{ij} \eta_i^{\dagger} \Phi \chi_j + \lambda_{ijk} \Phi^{\dagger} \eta_i \chi_j \chi_k + \text{H.c.}$$

Axion decay constant

$$\sigma = \frac{v_{\sigma} + \rho}{\sqrt{2}} \underbrace{(ia_{\sigma}/v_{\sigma})}$$

$$f_a = \frac{f_{PQ}}{N} = \frac{v_\sigma}{\sqrt{2}N}$$

Color-anomaly factor

$$N = 2n_{\Psi}\omega(2n \pm 1)T(p,q)$$

QCD axion mass relation

$$m_a = 5.70(7) \left(\frac{10^{12} \; {\rm GeV}}{f_a} \right) {
m \mu eV}$$

Vector-like quarks

Complex scalar singlet

Colored scalars

Fields	$SU(3)_c \otimes SU(2)_L \otimes U(1)_Y$	$U(1)_{PQ}$	Multiplicity
$\overline{\Psi_L}$	$[(p,q),2n\pm 1,0]$	ω	n_{Ψ}
Ψ_R	$[(p,q), 2n \pm 1, 0]$	0	n_{Ψ}
σ	(1, 1, 0)	ω	1
η	[(p,q), 2n, 1/2]	0	n_{η}
χ	$[(p,q),2n\pm 1,0]$	0	n_χ

$$-\mathcal{L}_{\text{Yuk.}} \supset \mathbf{Y}_{\Psi} \overline{\Psi_{L}} \Psi_{R} \sigma + \frac{1}{2} \mathbf{Y}_{\chi_{j}} \Psi_{R}^{T} C \chi_{j} \Psi_{R} + \mathbf{Y}_{i} \bar{L} \eta_{i}^{*} \Psi_{R} + \text{H.c.}$$

$$V \supset \mu_{ijk} \chi_{i} \chi_{j} \chi_{k} + \kappa_{ij} \eta_{i}^{\dagger} \Phi \chi_{j} + \lambda_{ijk} \Phi^{\dagger} \eta_{i} \chi_{j} \chi_{k} + \text{H.c.}$$

Vector-like quarks	Vect	or-like	quarks
--------------------	------	---------	--------

Complex scalar singlet

Colored scalars

Fields	$SU(3)_c \otimes SU(2)_L \otimes U(1)_Y$	$U(1)_{PQ}$	Multiplicity
$\overline{\Psi_L}$	$[(p,q), 2n \pm 1, 0]$	ω	n_{Ψ}
Ψ_R^-	$[(p,q), 2n \pm 1, 0]$	0	n_{Ψ}
σ	(1, 1, 0)	ω	1
η	[(p,q), 2n, 1/2]	0	n_{η}
χ	$[(p,q),2n\pm 1,0]$	0	n_χ

Neutrino mass generation

$$-\mathcal{L}_{\text{Yuk.}} \supset \mathbf{Y}_{\Psi} \overline{\Psi_{L}} \Psi_{R} \sigma + \frac{1}{2} \mathbf{Y}_{\chi_{j}} \Psi_{R}^{T} C \chi_{j} \Psi_{R} + \mathbf{Y}_{i} \bar{L} \eta_{i}^{*} \Psi_{R} + \text{H.c.}$$

$$V \supset \mu_{ijk} \chi_{i} \chi_{j} \chi_{k} + \kappa_{ij} \eta_{i}^{\dagger} \Phi \chi_{j} + \lambda_{ijk} \Phi^{\dagger} \eta_{i} \chi_{j} \chi_{k} + \text{H.c.}$$

Complex scalar singlet

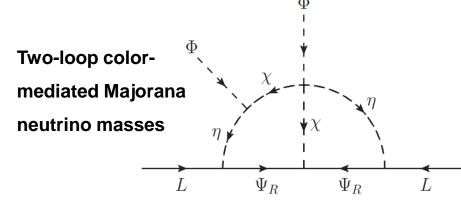
Colored scalars

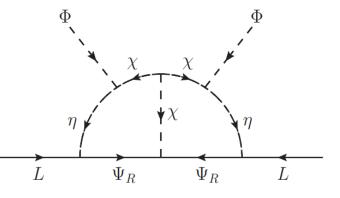
Fields	$SU(3)_c \otimes SU(2)_L \otimes U(1)_Y$	$U(1)_{PQ}$	Multiplicity
$\overline{\Psi_L}$	$[(p,q), 2n \pm 1, 0]$	ω	n_{Ψ}
Ψ_R	$[(p,q), 2n \pm 1, 0]$	0	n_{Ψ}
σ	(1, 1, 0)	ω	1
η	[(p,q), 2n, 1/2]	0	n_{η}
χ	$[(p,q),2n\pm 1,0]$	0	n_{χ}

Neutrino mass generation

$$-\mathcal{L}_{\text{Yuk.}} \supset \mathbf{Y}_{\Psi}\overline{\Psi_{L}}\Psi_{R}\sigma + \frac{1}{2}\mathbf{Y}_{\chi_{j}}\Psi_{R}^{T}C\chi_{j}\Psi_{R} + \mathbf{Y}_{i}\bar{L}\eta_{i}^{*}\Psi_{R} + \text{H.c.}$$

$$V \supset \mu_{ijk}\chi_i\chi_j\chi_k + \kappa_{ij}\eta_i^{\dagger}\Phi\chi_j + \lambda_{ijk}\Phi^{\dagger}\eta_i\chi_j\chi_k + \text{H.c.}$$





Cheng,Li (1980), Zee (1986), Babu (1988)

Vector-like	quarks
--------------------	--------

Complex scalar singlet

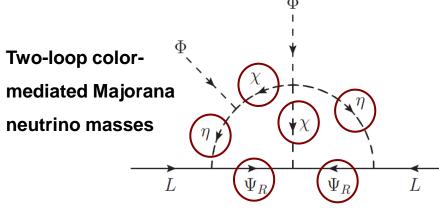
Colored scalars

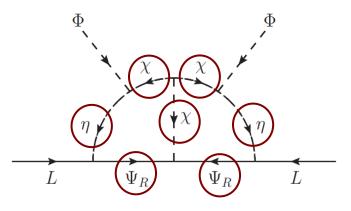
Fields	$SU(3)_c \otimes SU(2)_L \otimes U(1)_Y$	$U(1)_{PQ}$	Multiplicity
$\overline{\Psi_L}$	$[(p,q), 2n \pm 1, 0]$	ω	n_{Ψ}
Ψ_R^-	$[(p,q), 2n \pm 1, 0]$	0	n_{Ψ}
σ	(1, 1, 0)	ω	1
η	[(p,q), 2n, 1/2]	0	n_{η}
χ	$[(p,q),2n\pm 1,0]$	0	n_χ

Neutrino mass generation

$$-\mathcal{L}_{\text{Yuk.}} \supset \mathbf{Y}_{\Psi}\overline{\Psi_{L}}\Psi_{R}\sigma + \frac{1}{2}\mathbf{Y}_{\chi_{j}}\Psi_{R}^{T}C\chi_{j}\Psi_{R} + \mathbf{Y}_{i}\bar{L}\eta_{i}^{*}\Psi_{R} + \text{H.c.}$$

$$V \supset \mu_{ijk}\chi_i\chi_j\chi_k + \kappa_{ij}\eta_i^{\dagger}\Phi\chi_j + \lambda_{ijk}\Phi^{\dagger}\eta_i\chi_j\chi_k + \text{H.c.}$$





Cheng,Li (1980), Zee (1986), Babu (1988)

Vector-like quarks

Complex scalar singlet

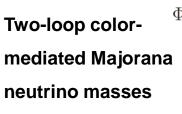
Colored scalars

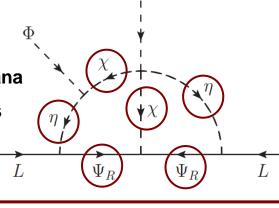
Fields	$SU(3)_c \otimes SU(2)_L \otimes U(1)_Y$	$U(1)_{PQ}$	Multiplicity
$\overline{\Psi_L}$	$[(p,q), 2n \pm 1, 0]$	ω	n_{Ψ}
Ψ_R	$[(p,q), 2n \pm 1, 0]$	0	n_{Ψ}
σ	(1, 1, 0)	ω	1
η	[(p,q), 2n, 1/2]	0	n_{η}
χ	$[(p,q),2n\pm 1,0]$	0	n_χ

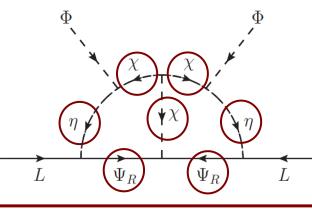
Neutrino mass generation

$$-\mathcal{L}_{\text{Yuk.}} \supset \mathbf{Y}_{\Psi}\overline{\Psi_{L}}\Psi_{R}\sigma + \frac{1}{2}\mathbf{Y}_{\chi_{j}}\Psi_{R}^{T}C\chi_{j}\Psi_{R} + \mathbf{Y}_{i}\bar{L}\eta_{i}^{*}\Psi_{R} + \text{H.c.}$$

$$V \supset \mu_{ijk} \chi_i \chi_j \chi_k + \kappa_{ij} \eta_i^{\dagger} \Phi \chi_j + \lambda_{ijk} \Phi^{\dagger} \eta_i \chi_j \chi_k + \text{H.c.}$$







Cheng,Li (1980), Zee (1986), Babu (1988)

$$(m_{\nu})_{\alpha\beta} \sim 0.1 \text{ eV} \left(\frac{\tilde{Y}_{a\alpha}^{j}(\tilde{Y}_{\chi})_{ab}^{k}\tilde{Y}_{b\beta}^{l}}{10^{-3}}\right) \left(\frac{\tilde{\mu}_{jkl}}{10^{8} \text{ GeV}}\right) \left(\frac{v}{246 \text{ GeV}}\right)^{2} \left(\frac{10^{8} \text{ GeV}}{m_{\zeta}}\right)^{2}$$

Axion-to-photon coupling

$$g_{a\gamma\gamma} = rac{lpha_e}{2\pi f_a} \left[rac{E}{N} - 1.92(4)
ight]$$
 Cortona et al.(2016

Axion-to-photon coupling

$$g_{a\gamma\gamma}=rac{lpha_e}{2\pi f_a}\left[\!\!\left(\!\!rac{E}{N}\!\!\right)\!\!-1.92(4)
ight]$$
 Cortona et al.(2016)

Axion-to-photon coupling

$$g_{a\gamma\gamma} = rac{lpha_e}{2\pi f_a} \left[\left(\!\! rac{E}{N} \!\!
ight) \!\! - 1.92(4)
ight]$$
 Cortona et al.(2016

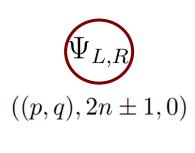
$$\Psi_{L,R}$$

$$((p,q), 2n \pm 1, 0)$$

$$\mathrm{SU}(3)_c \otimes \mathrm{SU}(2)_L \otimes \mathrm{U}(1)_Y$$

Axion-to-photon coupling

$$g_{a\gamma\gamma} = rac{lpha_e}{2\pi f_a} \left[\left(\!\! rac{E}{N} \!\!
ight) \!\! - 1.92(4)
ight]$$
 Cortona et al.(2016)

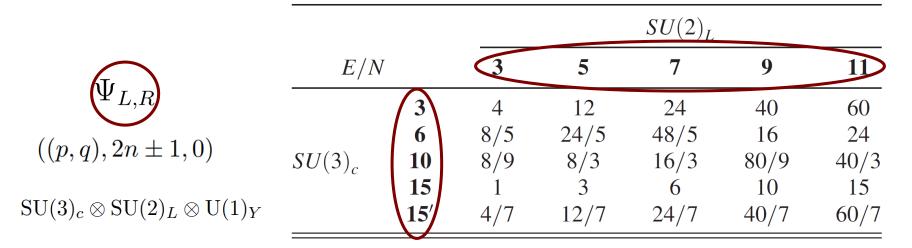


$$\mathrm{SU}(3)_c \otimes \mathrm{SU}(2)_L \otimes \mathrm{U}(1)_Y$$

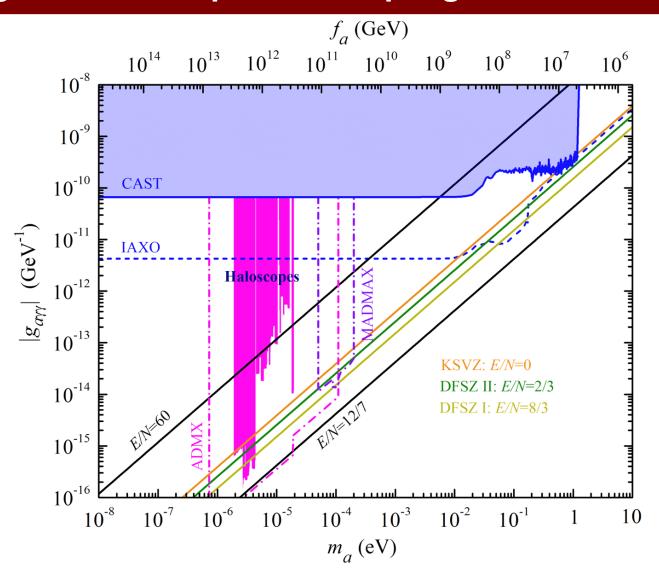
				$SU(2)_L$		
E/N	V	3	5	7	9	11)
	3	4	12	24	40	60
	6	8/5	24/5	48/5	16	24
$SU(3)_c$	10	8/9	8/3	16/3	80/9	40/3
	15	1	3	6	10	15
	15′	4/7	12/7	24/7	40/7	60/7

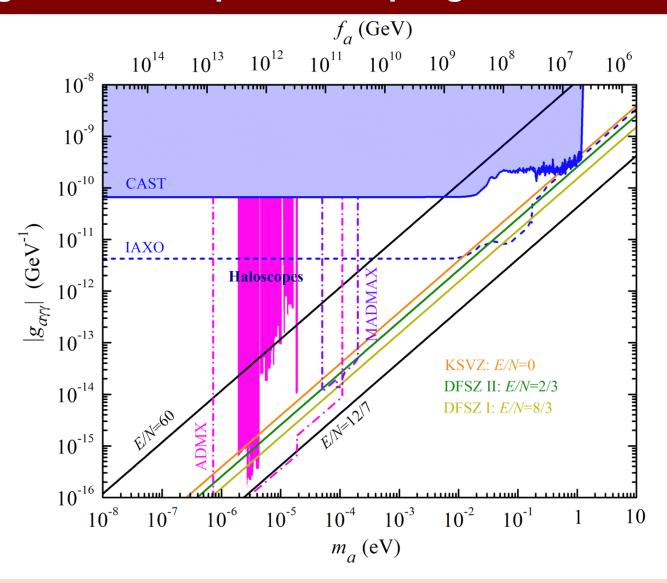
Axion-to-photon coupling

$$g_{a\gamma\gamma} = rac{lpha_e}{2\pi f_a} \left[\!\! \left[\!\! rac{E}{N} \!\!
ight] \!\! - 1.92(4)
ight]$$
 Cortona et al.(2016)



$$\frac{E}{N} = \frac{d(p,q)}{(2n\pm 1)T(p,q)} \sum_{j=0}^{2n\pm 1-1} \left(\frac{2n\pm 1-1}{2} - j\right)^2$$





Axion-to-photon coupling allows to probe the different models at **helioscope** and **haloscope** experiments.

Colored scalars

Vector-like quarks

$$\eta$$
 ((p,q), 2n, 1/2)

$$((p,q), 2n, 1/2)$$
 $\Psi_{L,R}$ $((p,q), 2n \pm 1, 0)$

$$\chi \qquad ((p,q), 2n \pm 1, 0)$$

Colored scalars

$$\eta$$
 ((p,q), 2n, 1/2)

$$\chi \qquad ((p,q), 2n \pm 1, 0)$$

Vector-like quarks

$$((p,q), 2n, 1/2)$$
 $\Psi_{L,R}$ $((p,q), 2n \pm 1, 0)$

Lead to potentially dangerous stable couloured/baryonic and electrically charged relics ...

Colored scalars

Vector-like quarks

$$\eta$$
 $((p,q), 2n$

 $((p,q), 2n \pm 1, 0)$

$$((p,q),2n,1/2)$$
 $\Psi_{L,R}$ $((p,q),2n\pm 1,0)$

Lead to potentially dangerous stable couloured/baryonic and electrically charged relics ...

Axions are naturally light, weakly coupled with ordinary matter, cosmologically stable, and can be nonthermally produced in the early Universe being an excellent DM candidate.

Colored scalars

Vector-like quarks

$$\eta$$

 $((p,q), 2n \pm 1, 0)$

$$((p,q),2n,1/2)$$
 $\Psi_{L,R}$ $((p,q),2n\pm 1,0)$

Lead to potentially dangerous stable couloured/baryonic and electrically charged relics ...

Axions are naturally light, weakly coupled with ordinary matter, cosmologically stable, and can be nonthermally produced in the early Universe being an excellent DM candidate.

Axion dark matter via the misalignment mechanism in pre-inflationary scenario

Callan et al. (1978); Gross et al. (1981); Dimopoulos et al. (2008)

$$\Omega_a h^2 \simeq \Omega_{\rm CDM} h^2 \frac{\theta_0^2}{2.15^2} \left(\frac{f_a}{2 \times 10^{11} \text{ GeV}} \right)^{\frac{7}{6}}$$

Colored scalars

Vector-like quarks

$$\eta = ((p,q), 2n, 1/2)$$

$$((p,q),2n,1/2)$$
 $\Psi_{L,R}$ $((p,q),2n\pm 1,0)$

$$\chi \qquad ((p,q), 2n \pm 1, 0)$$

Lead to potentially dangerous stable couloured/baryonic and electrically charged relics ...

Axions are naturally light, weakly coupled with ordinary matter, cosmologically stable, and can be nonthermally produced in the early Universe being an excellent DM candidate.

Axion dark matter via the misalignment mechanism in pre-inflationary scenario

Callan et al. (1978); Gross et al. (1981); Dimopoulos et al. (2008)

$$\Omega_a h^2 \simeq \Omega_{\text{CDM}} h^2 \left(\frac{\theta_0^2}{2.15^2}\right) \left(\frac{f_a}{2 \times 10^{11} \text{ GeV}}\right)^{\frac{7}{6}}$$

Colored scalars

Vector-like quarks

$$\eta$$
 $((p,q),2n,1/2)$

$$((p,q),2n,1/2)$$
 $\Psi_{L,R}$ $((p,q),2n\pm 1,0)$

$$\chi \qquad ((p,q), 2n \pm 1, 0)$$

Lead to potentially dangerous stable couloured/baryonic and electrically charged relics ...

Axions are naturally light, weakly coupled with ordinary matter, cosmologically stable, and can be nonthermally produced in the early Universe being an excellent DM candidate.

Axion dark matter via the misalignment mechanism in pre-inflationary scenario

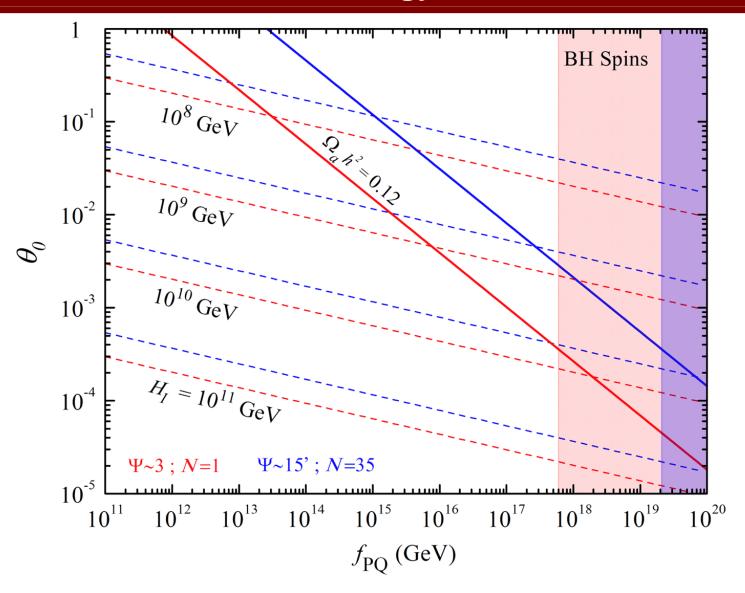
Callan et al. (1978); Gross et al. (1981); Dimopoulos et al. (2008)

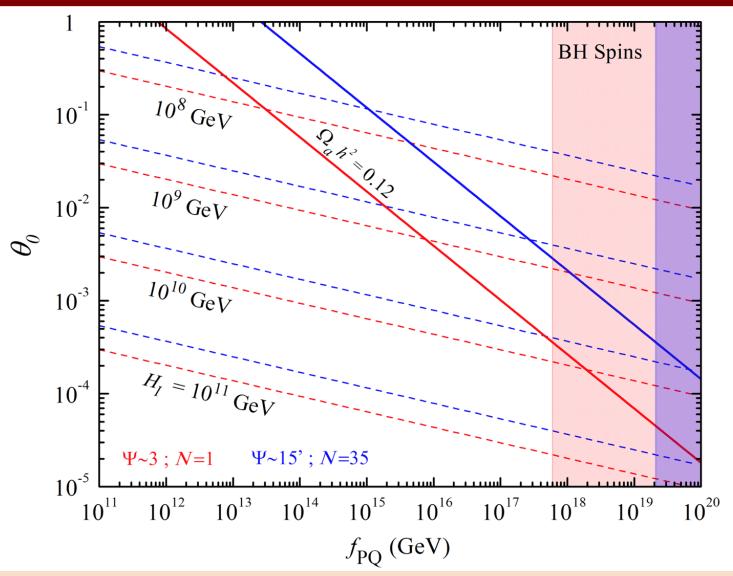
$$\Omega_a h^2 \simeq \Omega_{\rm CDM} h^2 \left(\frac{\theta_0^2}{2.15^2} \right) \left(\frac{f_a}{2 \times 10^{11} \text{ GeV}} \right)^{\frac{7}{6}}$$

Isocurvature fluctuations are constrained by CMB data setting a bound on the inflationary scale

$$H_I \lesssim \frac{0.9 \times 10^7}{\Omega_a h^2 / \Omega_{\rm CDM} h^2} \left(\frac{\theta_0}{\pi} \frac{f_a}{10^{11} \text{ GeV}} \right) \text{ GeV}$$

Di Luzio et al. (2017)





For $\vartheta_0 \sim O(1)$, axions can account for the full CDM budget, provided $f_a \sim 10^{12}$ GeV, a region currently under scrutiny at haloscopes.

 We proposed a connection between two seemingly unrelated facts: small neutrino masses and the strong CP problem. This was achieved within a novel class of KSVZ axion schemes, containing exotic colored fermions and scalars that act as Majorana neutrino mass mediators at the two-loop level.

 We proposed a connection between two seemingly unrelated facts: small neutrino masses and the strong CP problem. This was achieved within a novel class of KSVZ axion schemes, containing exotic colored fermions and scalars that act as Majorana neutrino mass mediators at the two-loop level.

• **Different representation assignments** of the new fields under the SM and PQ symmetries yield **distinct axion-to photon couplings**. This provides a way to differentiate the various realizations of our scheme at future **helioscope and haloscope experiments**.

 We proposed a connection between two seemingly unrelated facts: small neutrino masses and the strong CP problem. This was achieved within a novel class of KSVZ axion schemes, containing exotic colored fermions and scalars that act as Majorana neutrino mass mediators at the two-loop level.

- **Different representation assignments** of the new fields under the SM and PQ symmetries yield **distinct axion-to photon couplings**. This provides a way to differentiate the various realizations of our scheme at future **helioscope and haloscope experiments**.
- Due to potentially dangerous colored relics, we consider **axion DM in the preinflationary scenario**, where the PQ symmetry is broken before inflation.

 We proposed a connection between two seemingly unrelated facts: small neutrino masses and the strong CP problem. This was achieved within a novel class of KSVZ axion schemes, containing exotic colored fermions and scalars that act as Majorana neutrino mass mediators at the two-loop level.

- **Different representation assignments** of the new fields under the SM and PQ symmetries yield **distinct axion-to photon couplings**. This provides a way to differentiate the various realizations of our scheme at future **helioscope and haloscope experiments**.
- Due to potentially dangerous colored relics, we consider **axion DM in the preinflationary scenario**, where the PQ symmetry is broken before inflation.

Thank you!